

# The Impact of a Solar Extreme Event on the Middle Atmosphere, a Case Study

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**Abstract.** A possible impact of an extreme solar particle event (ESPE) on the middle atmosphere is studied for the present-day climate and geomagnetic conditions. We consider an ESPE with an occurrence probability of about 1 per millenium. In addition, we assume that the ESPE is followed by an extreme geomagnetic storm (GMS), and we compare the contribution of both extreme events. The strongest known and best documented ESPE of 774/5 CE is taken as a reference example and established estimates of the corresponding ionization rates are applied. The ionization rates due to the energetic particle precipitation (EPP) during an extreme geomagnetic storm are up-scaled from analyzed distributions of electron energy spectra of observed geomagnetic storms. The consecutive buildup of NO<sub>x</sub> and HO<sub>x</sub> by the ionization is modeled in the high top 3D chemistry circulation model KASIMA, using specified dynamics from ERA-Interim analyses up to the stratopause. A specific dynamical situation was chosen which includes an elevated stratosphere event during January and maximizes the vertical coupling between the Northern polar mesosphere-lower thermosphere region and the stratosphere and therefore allows to estimate a maximum possible impact. The particle event initially produces about 65 Gmol of NO<sub>y</sub>, with 25 Gmol of excess NO<sub>y</sub> still after one year. The related ozone loss reaches up to 50% in the upper stratosphere during the first weeks after the event and slowly descends to the mid-stratosphere. After about one year, 20% ozone loss is still observed in the Northern stratosphere. The geomagnetic storm causes strong ozone reduction in the mesosphere but plays only a minor role for the reduction in total ozone. In the Southern hemisphere, the long-lived NO<sub>y</sub> in the polar stratosphere which is produced almost solely by the ESPE, is transported into the Antarctic polar vortex where it experiences strong denitrification into the troposphere. For this special case, we estimate a NO<sub>3</sub> wash-out which could produce a measurable signal in ice cores. The reduction in total ozone causes an increase of the UV erythema dose of less than 5% which maximizes in spring for Northern latitudes of 30° and in summer for Northern latitudes of about 60°.

## 20 1 Introduction

Strong events of solar activity, such as the Carrington flare event (Carrington (1859)), have been observed now for more than 150 years and since the satellite era also their energetic output is now well documented (see for example Aschwanden et al. (2017)). These events and connected strong geomagnetic storms often produce a strong flux of energetic particles which, when precipitating into the Earth's atmosphere, by ionization processes (Usoskin et al. (2011)), cause the formation of the radicals

25 NO<sub>x</sub> and HO<sub>x</sub>. This results in an additional ozone destruction in the middle atmosphere (see for example Sinnhuber et al. (2012)). For solar eruptions observed during the last decades there exist many studies, observational as well as modelling ones, documenting this solar-terrestrial interaction (see for example studies in the context of the WCRP SPARC High Energy Particle Precipitation in the Atmosphere (HEPPA) initiative as Funke et al. (2011a), Funke et al. (2017), Sinnhuber et al. (2022)). In addition, some studies indicate a possible small, but not conclusive surface impact (see for example Seppälä et al. (2009),  
30 Maliniemi et al. (2014), Calisto et al. (2011)). Even stronger events with a probability of  $\sim 1/1000 \text{ y}^{-1}$  have put marks in the concentration of cosmogenic isotopes in natural stratified records (e.g., tree trunks, ice cores) during the last millenia (Usoskin (2017), Cliver et al. (2022)). The impact of such extreme events have been studied mostly in the historical context to simulate the production of isotopes in the lower stratosphere and their transport to the surface to compare with the isotopic records (eg. Sukhodolov et al. (2017)). Here we apply ionization rates estimated for an extreme solar event in our chemical model  
35 of the atmosphere to study its chemical impact and its consequences for the ozone layer under present-day conditions, i.e., if such an event occurred nowadays. The main objective of this study is to estimate the maximum direct impact which such an extraordinary event may have on the chemical state of the atmosphere, and in consequence for the erythemal UV-dose on-ground. As the ozone destroying radicals are generated over a wide range of altitudes in the middle atmosphere, it is also the specific dynamical situation which determines its impact. We focus on the Northern hemisphere (NH) and apply a dynamical  
40 situation where the vertical coupling in the middle atmosphere is particularly efficient.

The paper is structured as follows: First, we construct an extreme solar scenario with properties of an extreme event observed in the past or estimated to have an occurrence rate of about 1/1000 y and estimate the corresponding ionization rate. Then we apply this scenario to the specific atmospheric situation, calculate the additional NO<sub>x</sub> and its impact on ozone. Finally we estimate the consequences for the additional UV dose following the event for NH mid latitudes.

## 45 **2 The extreme scenario setup**

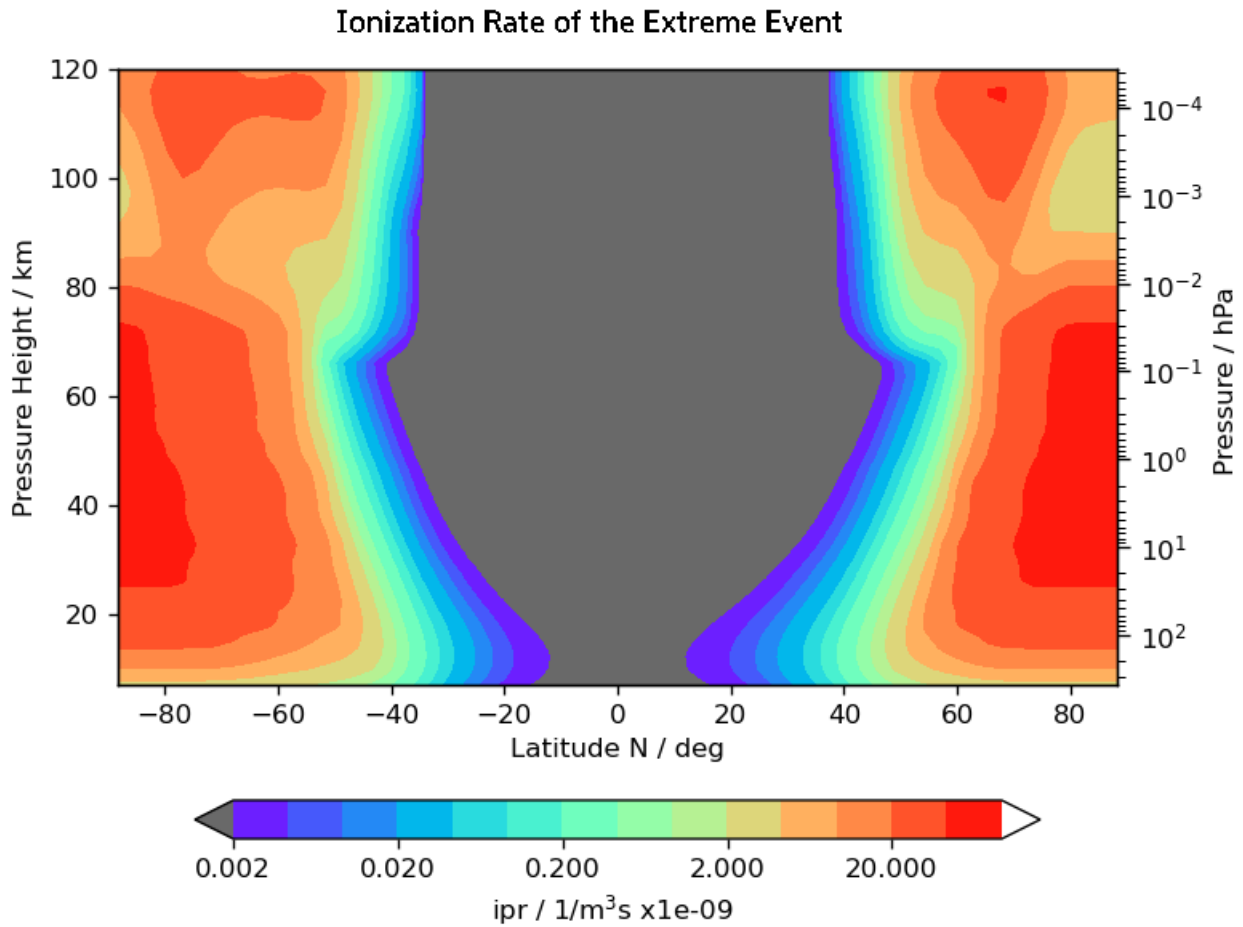
A solar eruption event has several components: the very first signal of such an event detectable on Earth is a flare of electromagnetic radiation emitted in the eruption at the solar surface. Observable from  $\gamma$ -ray energies down to the visual spectrum (white flare), the total impact of the flare on the Earth's middle atmosphere is small and negligible in terms of the chemical impact (Pettit et al. (2018)) and therefore for the possible dynamical coupling. Then, particles accelerated in this explosion and  
50 farther in the corona and interplanetary medium reach the Earth within minutes to hours after the flare, the so called solar proton event (SPE). Here the highest energies of the particles are observed up to several GeV which initiate a nucleonic cascade in the atmosphere. This can reach even the ground and, when colliding with atmosphere's main constituents, cause nuclear reactions which can form, in particular, the cosmogenic isotopes <sup>14</sup>C or <sup>10</sup>Be. As one of the strongest of such events detected in paleo-nuclide records so far, the possible solar eruption dated at 774/5 CE has been studied extensively (see eg. Usoskin  
55 (2017)) including its impact on the atmosphere (Sukhodolov et al. (2017)). Compared to the the strongest directly observed SPE in 1956, Mekhaldi et al. 2015 estimated a factor 40 higher fluence to produce the observed nuclide concentrations. More

recent studies hint to a even stronger fluence. Cliver et al. (2022) estimate a 70x particle fluence compared to the 1956 event and we scale the ionization rates of the 1956 event accordingly, viz. 70x.

In addition, a strong eruption on the Sun's surface is often accompanied by a so-called coronal mass ejection which when directed to Earth hits the magnetosphere and causes a major geomagnetic storm (GMS), typically in the range of two days after the flare event. The energies of the particles of the GMS (mostly electrons) are in the range of a few keV to about 1 MeV, and they impact mainly the mesosphere and lower thermosphere. The strength of the disturbance of the geomagnetic field can be expressed by planetary indices like Kp and Ap but their relation to ionization rates for extreme cases could be highly non-linear. Here we take a different approach and assume the results of the study of Meredith et al. (2016) to be applicable for strong geomagnetic events when studying the interaction with the atmosphere. Meredith et al. (2016) conducted an extreme values analysis of the electron flux measured in the MEPED (Medium Energy Proton and Electron Detector) 90°-pitch-angle detector for the three energy intervals provided by the instrument. We assume here that we can apply these results also for particles entering the atmosphere (0° pitch-angle) and scale observed strong GMS events according their distribution. We select extreme events at 30 keV and 300 keV for the values 4.5 and 6 of the parameter  $L_*$  (Meredith et al., 2016, Tbl.3+4) and take the strongest events as representatives of extreme geomagnetic events.  $L_*$  is related to the geomagnetic latitude where the particles enter the atmosphere which is normally the auroral oval ( $L_* = 6$ ), but for disturbed conditions it can be slightly shifted to lower latitudes ( $L_* = 4.5$ ). Note that their list excludes SPEs. The 2003-Nov-20 event is the highest at  $L_* = 4.5$  and third in rank at  $L_* = 6$  for 30 keV, the 2010-Apr-6 09-12UT event is highest at  $L_* = 6$  and the several 2010-Apr-6 events are together highest also at  $L_* = 4.5$  for 300 keV. The electron fluxes for these events have about a 1 in 10 year probability when inspecting their Fig. 6. From this Figure, one can also deduce enhancement factors for 1-in-100 year events which are about 2x for low and mid energy electrons, and 10x for high energy electrons compared to the 1-in-10 year events. Our extreme GMS event is finally constructed in the following way: We take the ionization rates (including protons) calculated with the AIMOS-AISstorm 2.0 model (described in Nesse Tyssøy et al. (2022)) interpolated to our model grid for 2003-Nov-20 and 2010-Apr-06, scale them with factors 2 and 10, respectively, and add them up as one event with a duration of one day. As the AIMOS-AISstorm model includes electrons limited to < 300 keV, the lower boundary of a realistic electron ionization rate is about 70 km. The role of an extreme GMS for the NOy budget has not been studied before. To assess its impact on the ozone layer, high top models are necessary to calculate the NOy in the source region and its subsequent downward transport.

Finally, we assume that SPE and GMS indeed can be combined for this sensitivity study, start with a one-day extreme SPE and after two days with the extreme GMS. Figure 1 shows the three days time and zonal average ionization rate. We note that this approach may be not physically completely consistent but should yield a meaningful upper value for the strength of an extreme event.

The impact of the solar events on ozone is mainly given by the particle induced buildup of NOx which catalytically reacts with ozone. NOy, which includes all nitrogen containing species except N<sub>2</sub>O including the reservoir gases for NOx, has a very long lifetime in the stratosphere where it contributes effectively to ozone destruction. In the middle atmosphere, the impact on ozone is the key parameter for any dynamical atmospheric impact of such events, as short-wave heating rates are dominated by ozone. Changes in ozone will change the temperature distribution the middle atmosphere, and winds and propagating waves



**Figure 1.** The ionization rate of the extreme event (mean of January 23 - 25 2009). One can clearly see the GMS component maximizing in the lower thermosphere following the auroral oval and the polar cap ionization by the SPE.

as a result. In addition, as the ozone density maximizes in the stratosphere, any reduction of ozone will also change the UV intensity at the surface. The SPE induced ionization rate maximizes in the upper stratosphere and lower mesosphere, the GMS has its maximum ionization in the lower thermosphere. In order to estimate a possible maximum impact of a solar event, a dynamical situation is appropriate where the portion of NO<sub>y</sub> built up in the mesosphere/lower thermosphere (MLT) during the event survives photochemical destruction and is transported into the lower stratosphere, where its lifetime is of the order of many years. This sets the date of such an experiment to the winter season.

In satellite observations from instruments like MIPAS on ENVISAT (Fischer et al. (2008)) strong NO<sub>y</sub> intrusions from the lower thermosphere into the NH mesosphere and upper stratosphere have been observed after mid-winter sudden stratospheric warmings (SSWs) when accompanied by so-called elevated stratosphere events. These show a strong downward transport of

air from the MLT (Holt et al. (2013)). For the Northern hemisphere it has been shown that these events also yield the strongest ozone impact (Sinnhuber et al. (2018)). In order to test if an extreme geomagnetic storm can reach a similar impact as the SPE which maximizes lower in the mesosphere, we synchronize the SSW and the geomagnetic storm. As an example for such an dynamical situation, we use the SSW of Jan 21 2009 which has also been studied in the HEPPA-II experiment (Funke et al. (2017)). We note that the situation in the Southern hemisphere may be different. Here the polar middle atmosphere is essentially undisturbed during the whole winter and an early intrusion from the MLT region may reach the mid stratosphere. This has also been observed for the Antarctic winter 2003 (Funke et al. (2005)).

In summary, the extreme event we study is a combination of a SPE and a GMS. The ionization rate is derived from observed events which have been scaled to a strength estimated for an event with a occurrence rate of about 1 in 100 - 1000 years, synchronized with an elevated stratosphere event.

### 3 Description of the model and experiment

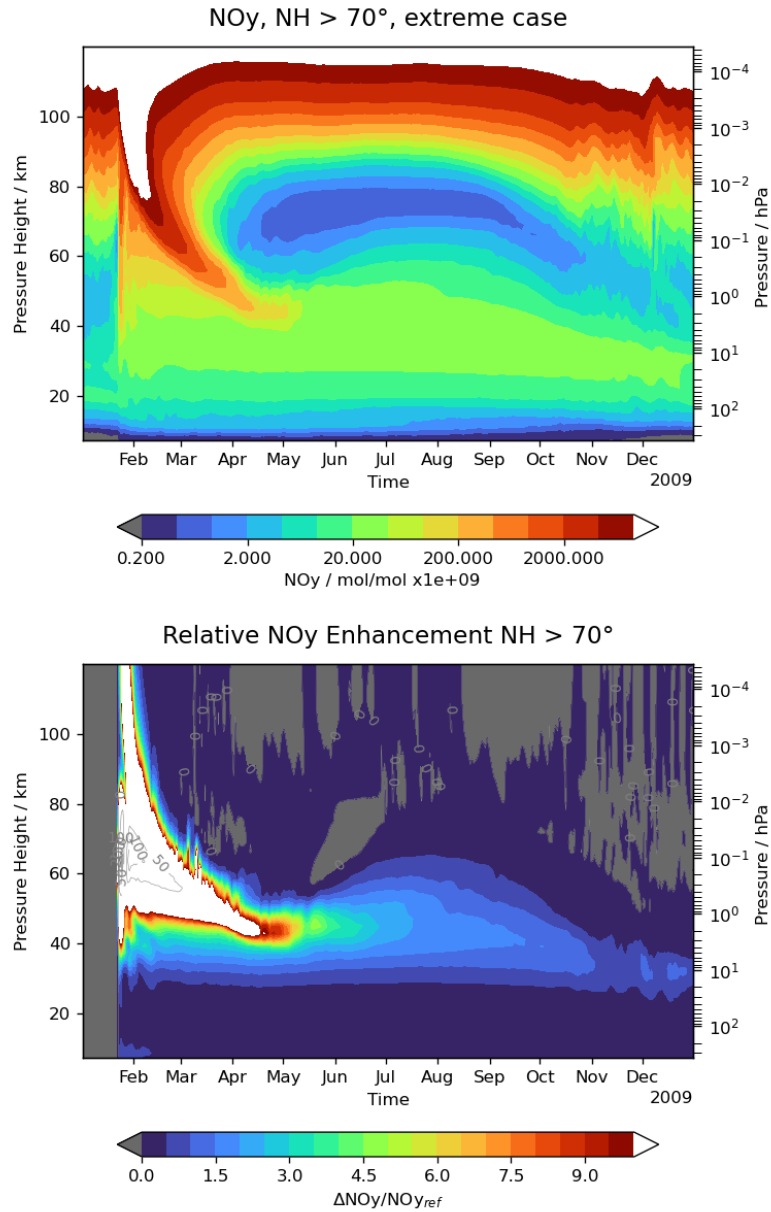
For the purpose of this sensitivity study we use the KARlsruhe SIMulation Model of the middle Atmopshere (Kouker et al. (1999)) in the version described in Sinnhuber et al. (2022). The model solves the meteorological basic equations in spectral form in the altitude range between 300 hPa and  $3.6 \times 10^{-5}$  hPa with the pressure height  $z = H \log(p/p_0)$  ( $H = 7$  km and  $p_0 = 1013.25$  hPa) as vertical coordinate. It uses radiative forcing terms for UV-Vis and IR, and a gravity wave drag scheme. In order to yield a realistic meteorology, the model is relaxed (nudged) to ERA-Interim meteorological analyses (Dee et al. (2011)) between the lower boundary of the model and 1 hPa. A full stratospheric chemistry including heterogeneous processes is adapted to include source terms related to particle ionization. For the production of HOx the parameterization of Solomon et al. (1981) is used. For the production of NOx, 0.7 NO molecules are produced per ion pair and 0.55 N atoms in ground state. The model also includes a HNO<sub>3</sub> production from proton hydrates based on the parameterization of de Zafra and Smyshlyaev (2001) which has been modified to be dependent on actual ionization rates. The model participated in all the three HEPPA inter-comparisons (Funke et al. (2011b), Funke et al. (2017), Sinnhuber et al. (2022)). The model has proven to simulate a realistic chemistry in the middle atmosphere for NOy intrusions which is sufficient to study the direct impact.

The following simulations are performed: one with the extreme event (EXT), where the ionization rates from the extreme event are applied from January-21 to January-23., and one with a background ionization only (REF). The background ionization rate is a mean from AIMOS ionization rates (Wissing and Kallenrode (2009)) for minimum Ap indices. In addition, a simulation which only includes the solar proton event is used to compare the contributions of the two components.

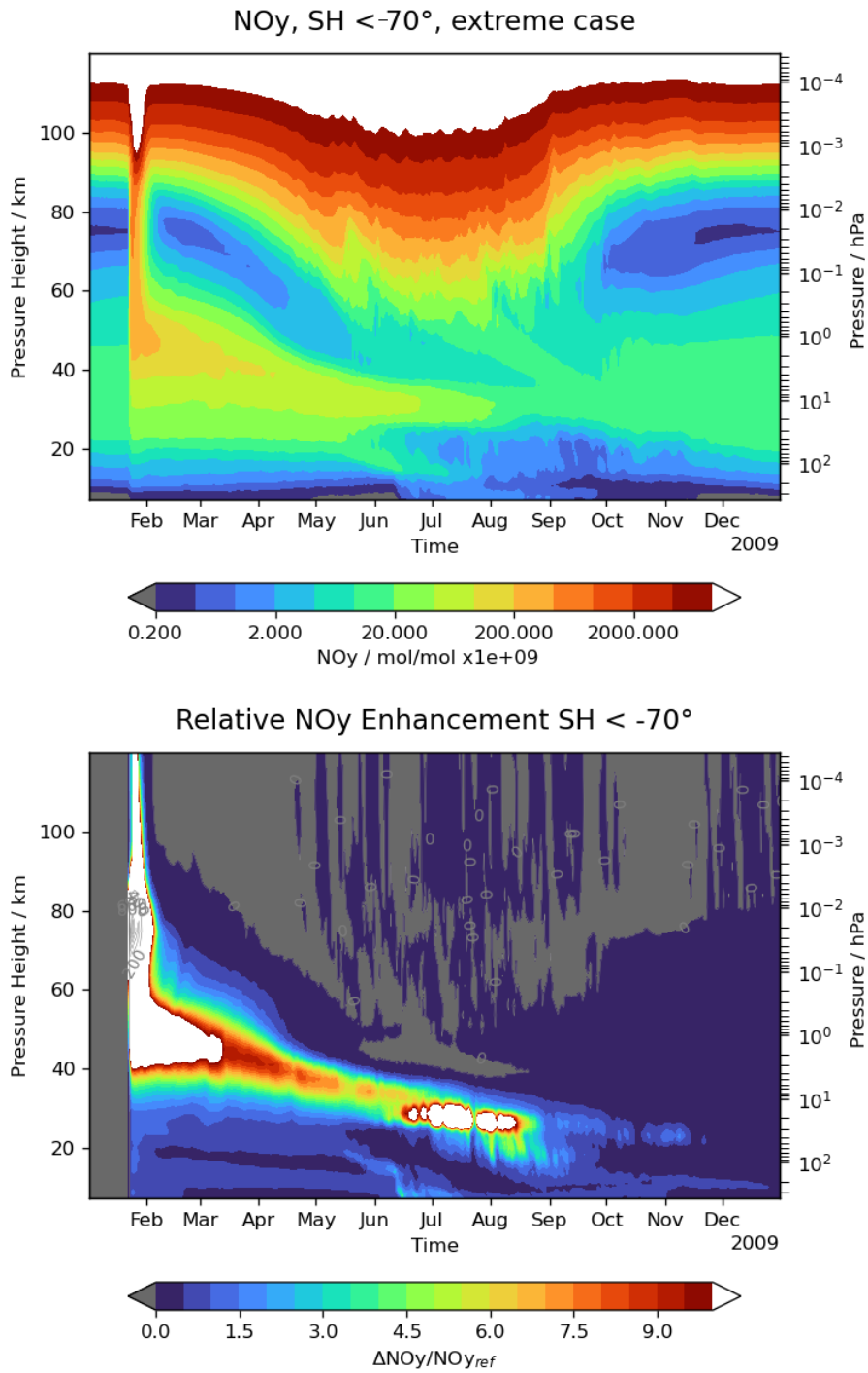
## 4 Results

### 4.1 The NOy intrusion

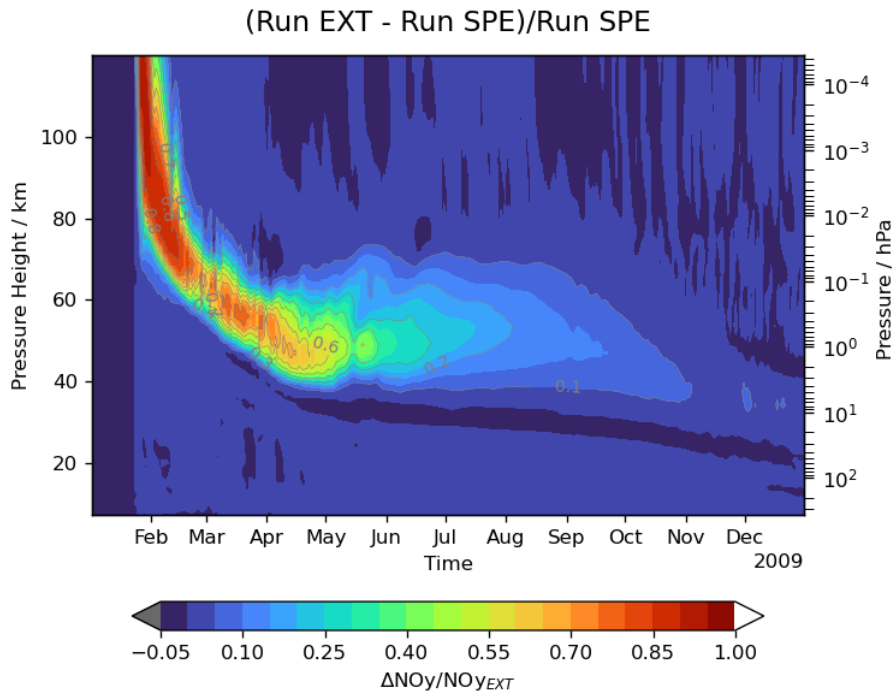
Figure 2 shows the mean of NOy ( $= \text{NO}_x + \text{HNO}_3 + \text{ClONO}_2 + \text{BrONO}_2 + \text{HONO} + \text{N}_2\text{O}_5 + \text{BrONO}$ ) (gas phase only) where  $\text{NO}_x = \text{N} + \text{NO} + \text{NO}_2$ , over the Northern polar cap (latitude  $> 70^\circ\text{N}$ ) for a year with the dynamics and chemical boundary



**Figure 2.** Top: NO<sub>y</sub> mixing ratio in the course of one year with specified dynamics of the year 2009 (ERA-Interim) and applying the extreme scenario case as described in the text. For the polar cap the intrusion connected to the MLT part dominates around the stratopause, below NO<sub>y</sub> stems from the SPE. Bottom: Relative enhancement of NO<sub>y</sub> compared to reference run.



**Figure 3.** As in 2, but for 70S

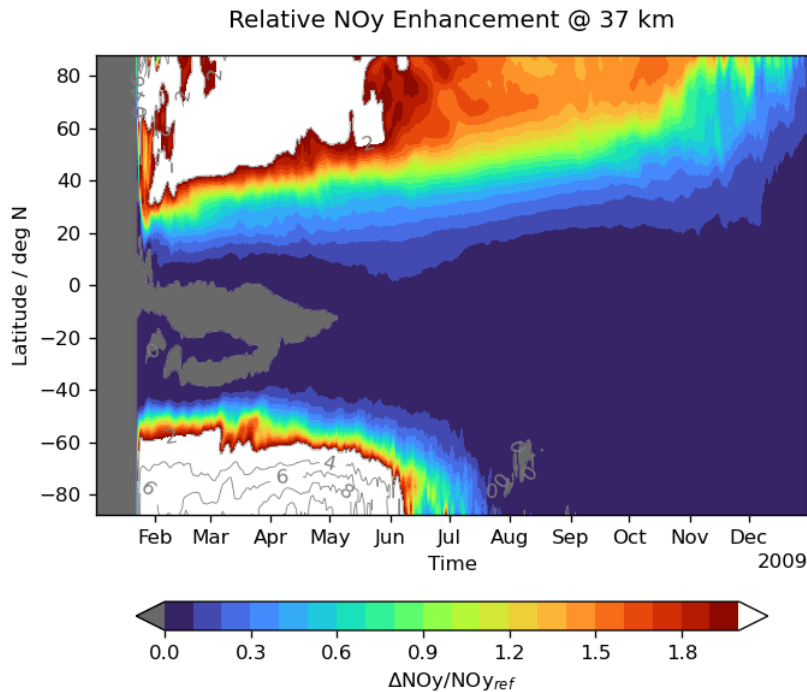


**Figure 4.** The relative contribution of the GMS component to the NO<sub>y</sub> intrusion (latitude > 70N), calculated as the difference between the EXT run and a run with the SPE component only, relative to the full event.

conditions of 2009, but with the extreme solar event described in Section 2 included. The two components of the event (SPE and GMS) can clearly be seen as an instantaneous increase over all altitudes and a tongue of increased NO<sub>y</sub> values transported from the MLT within the elevated stratopause event, respectively. In the specific dynamical situation, the downward transport of the GMS component from the MLT is fast enough to reach the photolytic 'safe' stratosphere at the end of the NH winter which then causes a significant NO<sub>y</sub> enhancement in the upper stratosphere. It stays there and is diluted during the summer, and is then transported further down in the beginning of the winter season. The enhancement related to the SPE is seen down to the tropopause and stays there over the whole year with relative enhancement of more than 20%. NO<sub>x</sub> is lost through transport into the troposphere where the main reservoir in the lowest stratosphere HNO<sub>3</sub> is washed out, or into the mesosphere where it is photolytically destroyed.

In the Southern hemisphere (shown in Figure 3), the NO<sub>y</sub> enhancement is essentially from the SPE, as the MLT component is photolytically destroyed during the summer season. We also observe a higher relative enhancement here, as the excess NO<sub>y</sub> in the upper stratosphere is transported downward in the beginning winter, compared to the upward draft in the NH summer. In addition, de- and renitrification by sedimentation of condensed HNO<sub>3</sub> can be observed, resulting in an effective removal of nitrate from the stratosphere into the troposphere. The relative maxima in the south winter middle stratosphere are related to the higher supply of HNO<sub>3</sub> where in the reference run this air mass already was depleted.





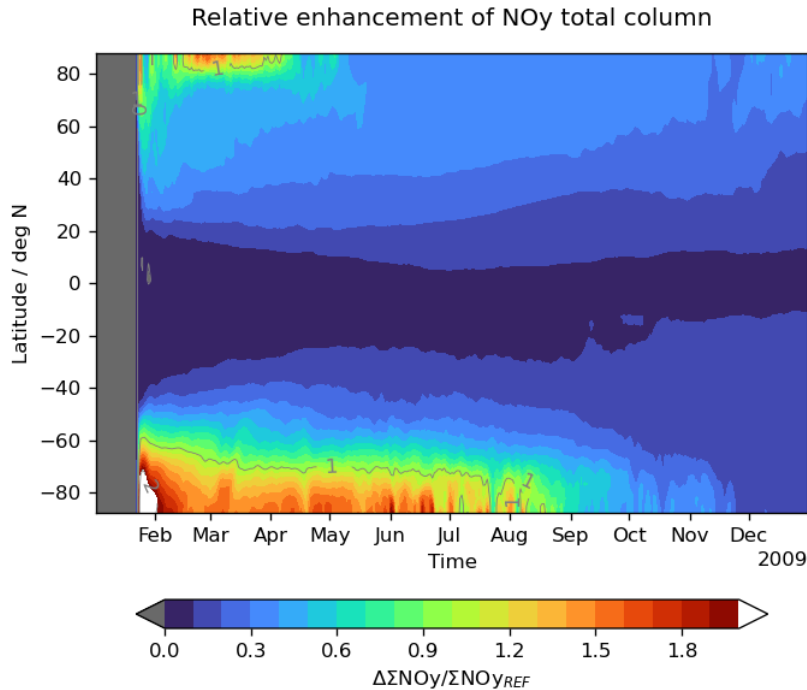
**Figure 5.** Relative enhancement of NOy by the the extreme event at 37 km pressure height.

The contribution of the GMS to the NOy at high Northern latitudes ( $>70^{\circ}\text{N}$ ) in the intrusion is shown in Figure 4. It is limited in altitude to above about 40 km and is diluted to values below 0.1 in autumn.

The distribution of NOy over latitudes at 37 km altitude is shown in Figure 5. At this altitude the contribution of NOy to ozone depletion maximizes (see for example Brasseur and Solomon (2005), their Figure 6.1). A closer inspection shows that at this altitude also NOy from the NH is transported to the SH and significantly contributes to the NOy after some months, as part of the diabatic circulation from the summer to the winter pole.

#### 4.2 The Impact on the Global NOy Budget

To assess the impact of the NOy intrusions on the global NOy budget, we first calculated the total column of NOy with and without the event, as shown in Figure 6. The two hemispheres show remarkable differences, with a higher total relative NOy enhancement at high Southern latitudes in the first six months but a fast decay after, and a longer lasting NOy enhancement in mid and high latitudes in the Northern hemisphere (see also 7). From Figure 3 which shows mixing ratios, one can deduce that the strong relative enhancement is affected by differences in the depletion of gaseous NOy by denitrification in the reference and extreme run. This gets clearer when showing the absolute change of the total columnar content. The change of total amount NOy (in units of mol) for different latitude bands by the event is shown in Figure 7. In the Southern hemisphere, the fast decay shortly after the event connected to the photolytical destruction in the polar mesosphere, then some stabilisation in the winter



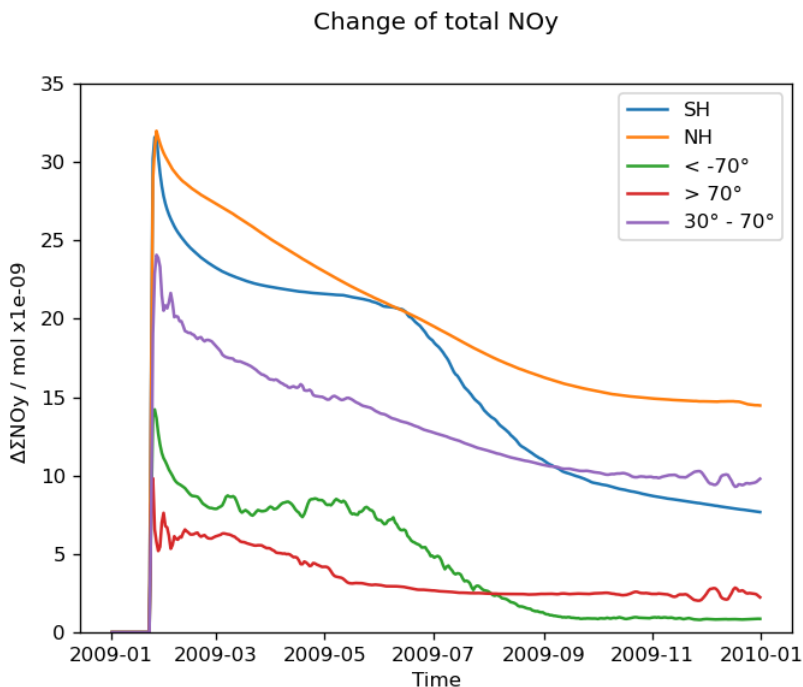
**Figure 6.** Enhancement of the total NO<sub>y</sub> column (> 7 km pressure height) relative to the reference run.

followed by the denitrification in the cold polar stratosphere can clearly be observed. The Northern hemisphere is characterized by a more continuous decay, stabilizing in the next winter season.

### 4.3 Impact on ozone

165 The NO<sub>y</sub> enhancement causes additional ozone depletion via the catalytic reaction of NO<sub>x</sub> with O<sub>3</sub>. This is shown in Figure 8 as the relative change of ozone. In the MLT, there is also a short lived ozone enhancement which is connected to the dissociation of O<sub>2</sub> by atomic nitrogen. In addition, the ozone decrease in the stratosphere causes less absorption of solar UV and a small build-up of ozone below the depleted zone. Mirroring the different amount of NO<sub>y</sub> transported into the stratosphere, the corresponding ozone loss is higher in the Southern hemisphere, whereas the ozone loss in the Northern hemisphere is lasting  
170 for a longer time. This just reflects the loss of NO<sub>y</sub> in the Southern hemisphere in the course of the polar winter.

The change of the ozone column relative to the reference run is shown in Figure 9. It shows the already described asymmetry in the two hemispheres, with exceptional high changes of total ozone loss in the Antarctic vortex. Total ozone recovers to nearly normal values at the end of the simulation year in the Southern hemisphere. On the other hand, in the Northern Hemisphere the maximum loss in midlatitudes is less than 5% but a deficit of total ozone of a few percent stays permanently throughout the  
175 total year. The contribution of the GMS to the loss in total ozone is negligible (not shown).

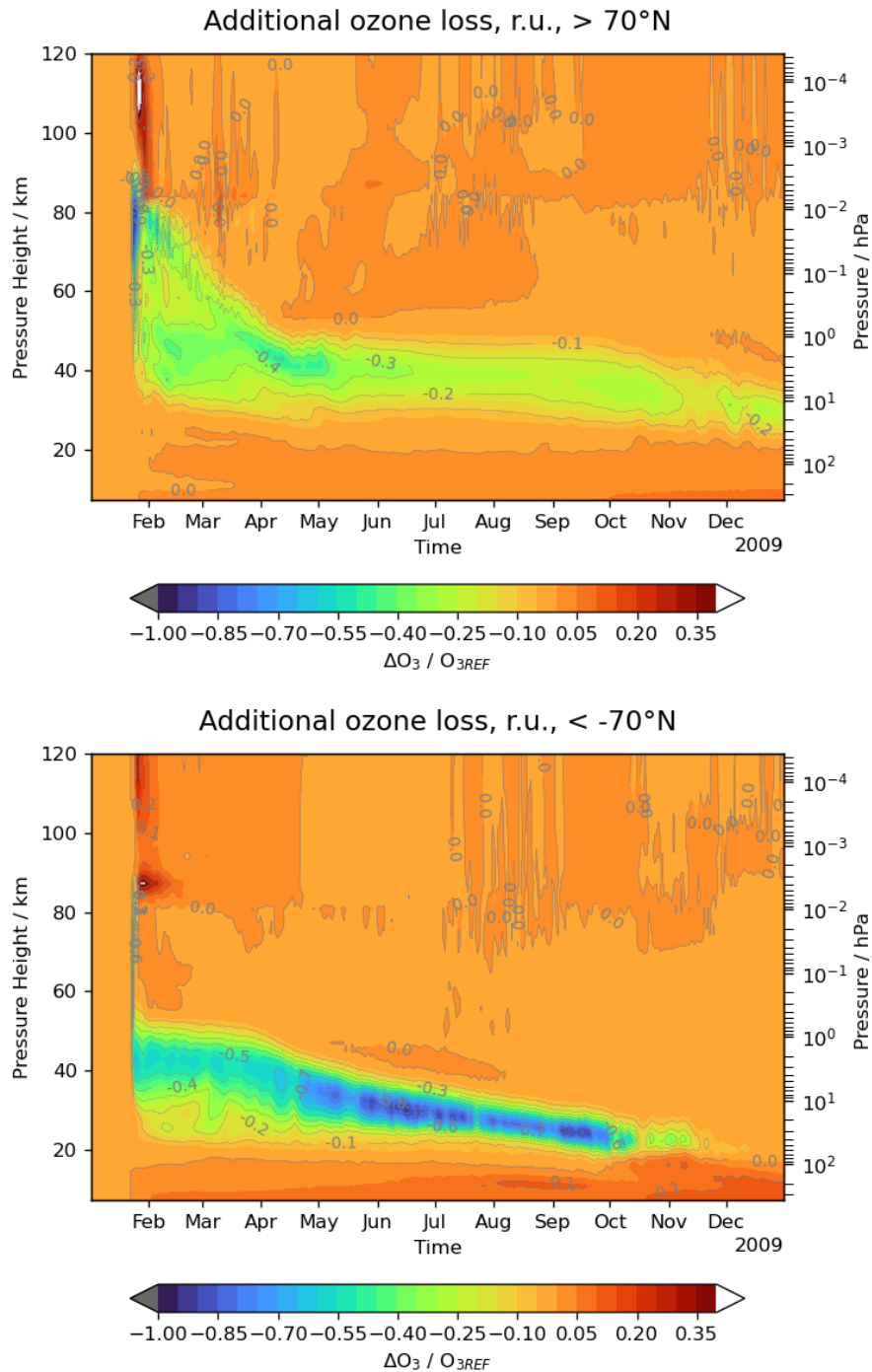


**Figure 7.** The NO<sub>y</sub> signal from the extreme event (EXT- REF) for different latitude bands as the total amount of NO<sub>y</sub> in mol.

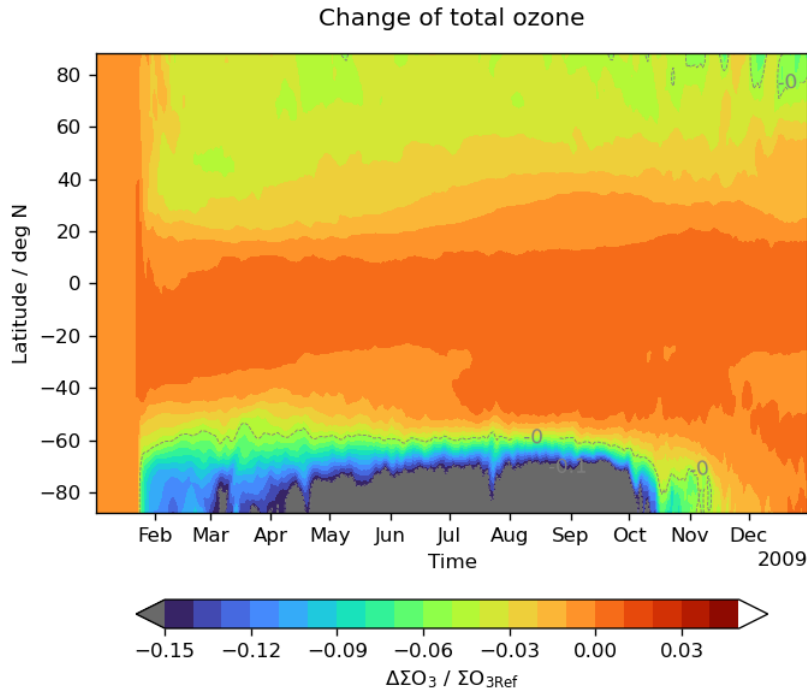
## 5 Discussion

### 5.1 The specific dynamical situation and the total NO<sub>y</sub> input

Inspecting Fig. 1, we note that the ionization rate of the GMS shows its maximum in the thermosphere and at the auroral oval latitudes, whereas the ESPE is concentrated to the poles. This has some consequences for the respective impact of the two components. As discussed in Section 2, the specific dynamical situation maximizes the impact of a solar extreme event in the Northern hemisphere by a long-lasting downward transport of NO<sub>y</sub> in the winter polar middle atmosphere. Fig. 4 shows that the contribution of the GMS to the additional NO<sub>y</sub> is of minor importance compared to the impact of the ESPE. The higher altitude of the source region together with its lower latitude makes the photochemical destruction of NO<sub>y</sub> at the end of the downward transport regime more probable. Accordingly, an additional experiment with a 10xGMS event, but without the ESPE component, was carried out. In this event, a NO<sub>y</sub> enhancement in the NH of about 2-3 GMol in late summer is observed. We therefore conclude that an extreme GMS in the range discussed will not pose a strong and long lasting disturbance in the NH middle atmosphere. Still, the mesospheric part of the ESPE profits from this dynamical situation in the NH: NO<sub>y</sub> which has not reached the mid stratosphere would have been destroyed photochemically when transported by an instable vortex or minor SSWs to sunlit mid-latitudes.



**Figure 8.** Ozone change in the extreme run relative to the reference run ((REF-EXT)/REF) related to the NO<sub>y</sub> intrusion through the extreme event (EXT - REF) for the Northern and Southern polar cap (latitude  $|\phi| \geq 70^\circ$ ).



**Figure 9.** The change of total ozone caused by the solar extreme event relative to the reference run.

190 For comparison, we tested also other dynamical situations, from early to late winter (only in the winter season NO<sub>y</sub> is stable in the mesosphere): in winter 2009/2010 but the event put on doy 300 in 2009, a rather early and strong mid-winter warming (December 1998), and the SSW in mid-January 2004, which from the date is quite similar to the set-up studied. For the EXT scenario, we estimate the initial total NO<sub>y</sub> input to about 65 Gmol which decreases to about 25 Gmol after one year (see Figure 7). The undisturbed total NO<sub>y</sub> amount in the atmosphere prior to the event above about 7 km altitude is simulated to 135 Gmol.

195 The experiments for 2004 and 1998 yield the same initial amount of the intrusion, but decay to about 10 Gmol and 25 Gmol after one year, respectively. The event set to 2009 doy 300 yielded with 53 Gmol slightly less initial NO<sub>y</sub> and showed a rather fast decay to only a few Gmol after one year.

Some insight might be gained also from a comparison with Sukhodolov et al. (2017) in this respect (albeit the results of Sukhodolov et al. (2017) are deduced from an ensemble mean and for pre-industrial conditions): their Fig. 4 shows a decay of the NO<sub>x</sub> enhancement to 100% within about 140 days in the NH. In the EXT experiment after 140 days the enhancements still exceeds 300% and reaches 100% after about 270 days.

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From these comparisons, the scenario selected seems to be an example of a strong and long enduring impact of the event. To estimate the range of the impact of strong solar events under different dynamical situations needs more simulations, also including feedback effects on the dynamics. This is beyond this study and will be addressed in a follow-up paper.

205 The total NO<sub>y</sub> mass produced in experiment EXT of about 65 Gmol corresponding to about 40% of the total atmospheric  
NO<sub>y</sub> content is substantially greater than what has been found in observation during the last decades. For example, Vitt and  
Jackman (1996) estimate a few percent increase of NO<sub>y</sub> after strong SPEs in the Northern polar stratosphere, Funke et al. (2014)  
estimated the total NO<sub>y</sub> input from satellite observations for the period 2002 to 2012 to a few Gmol per year, Reddmann et al.  
(2010) deduced the total NO<sub>y</sub> input from a model simulation for the period 2003 to 2004 to about 2 Gmol. In terms of NO<sub>y</sub>  
210 input, the NO<sub>y</sub> buildup by the solar event studied is about 10 to 30 times higher than what is estimated for solar particle events  
observed in the satellite era which is of the order of the scaling factor we apply to the ionization rates of the SPE. In this sense,  
the solar extreme event is also an atmospheric extreme.

The budget of NO<sub>y</sub> in the atmosphere on the longer term depends on the source strength of the event and the loss terms by  
photochemistry and exchange processes to the troposphere where NO<sub>y</sub> is removed by precipitation.

215 For the source strength we use ionization rates estimated for the event in 774 CE. There are several steps to estimate  
ionization rates for such events which all have factors of uncertainty. This is extensively discussed in Usoskin et al. (2011) and  
also in Sukhodolov et al. (2017). One of the main uncertainties of the SPE part is the energy spectrum of the protons which  
has been scaled from an event in 1956 with a rather hard spectrum, compared to other events in the satellite era Usoskin et al.  
(2020). A softer spectrum would shift the maximum NO<sub>y</sub> production to higher altitudes where the lifetime of NO<sub>y</sub> is reduced.  
220 Therefore, the initial NO<sub>y</sub> production by the ionization we estimate is probably on the upper limit of what can be expected.

The GMS contribution has been estimated from events observed in the past, but data at high flux levels are rare and have  
therefore substantial uncertainty. On the other hand, theoretical considerations (Vasyliunas (2011)) set an upper limit to the  
strength of a GMS in terms of Dst to about 2500 nT, which is probably outside the estimations we used. In addition, the  
geographical distribution of the auroral oval depends on the strengths of the event. Here we used the distributions from the  
225 AISstorm model for observed events and therefore miss the shift of the auroral oval to lower latitudes when the GMS strength  
is much higher. The scaling factors taken from Meredith et al. (2016) have been derived for a particle population that is partly  
trapped in the magnetosphere. Applying them to particle precipitation includes another source of uncertainty. But even a factor  
of two higher contribution from the GMS part would not bring this component above the SPE contribution.

For the efficiency of the NO production we use in the model the results of Porter et al. (1976). Some studies indicate a  
230 dependency of these factors on the ionization rate itself (Nieder et al. (2014)), with a kind of saturation effect, especially in the  
thermosphere. This can be a major uncertainty for the production term and needs detailed ion chemistry studies.

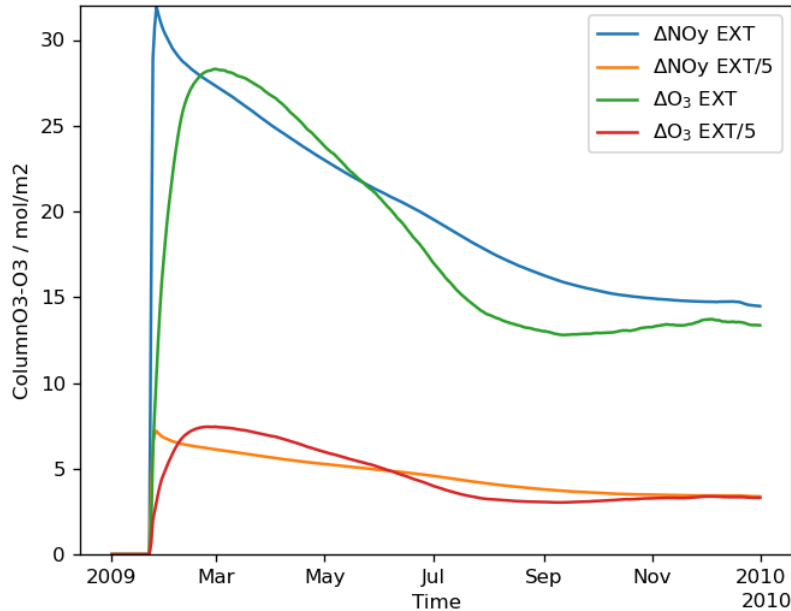
Besides the source strength, the downward transport from the source region in the mesosphere to the stratosphere is essential.  
From the studies within the SPARC HEPPA initiative (Funke et al. (2011a), Funke et al. (2017)) we know that many models do  
not correctly simulate this transport in the polar winter. The KASIMA model we use here in combination with ERA-Interim  
235 analyses yield a downward transport from the MLT which is general in good agreement with observations, also for a situation  
with an elevated stratosphere as we use here. Finally, the atmospheric lifetime of NO<sub>y</sub> is determined also by the correct pattern  
and strength of the Brewer-Dobson circulation. It determines the upward transport into the region with photolytic loss, and the  
exchange to the troposphere. Mean age of air simulations with KASIMA yield here also a good agreement with observations  
(Stiller et al. (2008), Haenel et al. (2015)).

240 In conclusion, the main uncertainties for the strength of the NO<sub>y</sub> intrusion for a selected dynamical scenario as in our simulation is the energy spectrum of the particle flux and the initial ion chemistry production terms.

## 5.2 The NO<sub>y</sub> loss through the Antarctic winter

The loss of NO<sub>y</sub> in the Southern polar winter atmosphere deserves special attention. The KASIMA model includes heterogeneous processes in the polar winter stratosphere causing denitrification and renitrification of the cold air masses in the Antarctic polar vortex through the nitric acid trihydrate (NAT) condensation of HNO<sub>3</sub> on water ice and the adsorption on sulfate aerosols, followed by sedimentation of the larger particles into lower layers as part of the annual cycle. In the KASIMA model, this process seems to be saturated: in the core of the Antarctic polar vortex denitrification is complete. Additional NO<sub>y</sub> brought into the lower stratosphere by the intrusion does not change that but by the changed saturation pressure, more volume will take part in the denitrification. The change of total NO<sub>y</sub> for the Southern polar cap in Figure 7 supports that finding: at the end of the winter, most of the additional NO<sub>y</sub> of about 6 Gmol in the Southern polar cap is lost to the troposphere.

Inspection of Figure 3 shows that the major NO<sub>y</sub> renitrification in the Southern polar lower stratosphere takes place between end of July to mid of August. NO<sub>y</sub> here is mainly in the form of HNO<sub>3</sub>. Once transported to the troposphere, NO<sub>y</sub> is deposited onto the surface by rain or snow within a few days. If we assume that the NO<sub>y</sub> is deposited in the form of NO<sub>3</sub> on the polar cap, we deduce an average deposition value of  $2.4 \cdot 10^{-6} \text{g/cm}^2$ . Using the approximation that the column integrated particle produced NO<sub>y</sub> is deposited, Melott et al. (2016) estimate a deposition of about  $1.4 \cdot 10^{-7} \text{g/cm}^2$  for the 1956 event, corresponding to  $1 \cdot 10^{-5} \text{g/cm}^2$  for the extreme event just by applying the scaling factor. As in our model loss processes as photolysis and transport to lower latitudes is included, our lower value is in line with their estimation. With an average precipitation of 166 mm/year for Antarctica and assuming that the nitrate is deposited within one month we estimate a mass mixing ratio for nitrate of 190ng/g in a monthly layer of (evenly distributed) deposited snow. A typical detection limit for nitrate events measured in ice cores with monthly to sub-annual resolution is of the order of 100ng/g (Smart et al. (2014)), so we would expect that a nitrate signal could be observable for such an event in Antarctica. Note, that our model does not include tropospheric meteorology, so an analysis of regional precipitation patterns is outside the scope of this paper. Sukhodolov et al. (2017) estimate from their analysis that a nitrate signal would not be detectable for a Greenland ice core. This is essentially in agreement with our findings, which show a much smaller change of total NO<sub>y</sub> in the Northern polar cap in early spring. Here, the higher polar temperatures and less stable polar vortex do not allow strong and fast denitrification, the primary cause of the high NO<sub>y</sub> depletion in the Southern polar lower stratosphere. In addition, we use analyzed winds and temperatures near to a realistic meteorology, whereas free running models may suffer in several aspects to correctly describe the polar dynamical and chemical processes in the lower stratosphere. For the case of the 774/5 event, ice core data for Antarctica do not show nitrate enhancements (Sukhodolov et al. (2017); Jong et al. (2022)). From Fig. 3 one may note that a strong denitrification signal is only expected when the NO<sub>y</sub> tongue reaches the lower stratosphere when there are also low temperatures. A mismatch, for example an event occurring in Antarctic spring, would presumably result in a much broader, and therefore undetectable nitrate peak.



**Figure 10.** Total NO<sub>y</sub> input by the EXT event and an event with a strength of 1/5 in Gmol into the Northern hemisphere. Shown is also the corresponding decrease of the ozone in Gmol, scaled by a factor of -0.03.

### 5.3 The ozone loss caused by the event

Inspecting Fig. 8, one may note the nearly complete ozone loss immediately after the event in the polar mesosphere. This rises  
 275 the question, to what extent the ozone loss scales with the NO<sub>y</sub> enhancement, and if the loss is near saturation, at least during  
 the first weeks after the event. For that reason we performed an additional sensitivity experiment where we scaled the strength  
 of the ionization rates of the EXT experiment by a factor of 0.2. Fig 10 gives the result of that experiment for the Northern  
 hemisphere in terms of the total amount of the enhancement of NO<sub>y</sub> and the corresponding ozone decrease. First, we can see  
 280 that the total NO<sub>y</sub> input into the atmosphere scales nearly perfectly with the strength of the forcing. Surprisingly, also the ozone  
 decrease scales with the ionization rate with about the same factor, with a only slightly relative higher ozone decrease in the  
 reduced case. In addition, the ozone loss is strongly correlated with the additional NO<sub>y</sub>.

The reason for this behaviour lies in the fact that the ESPE reaches deep into the middle atmosphere where the relative ozone  
 reduction related to enhanced NO<sub>y</sub> is small, again see 2. For total ozone, which essentially stems from the lower stratosphere,  
 this results in quasi-linear scaling of  $\Delta\text{NO}_y$  and  $\Delta\text{O}_3$ .

285 Considering radiative feedback effects, the situation may be different: following Gray et al. (2010), an efficient top down  
 mechanism for radiative feedback involves temperature gradients induced by ozone changes at mid-latitudes around the



stratopause. For our setup, the induced decrease of the ozone amount above 40 km for latitudes  $> 40^\circ\text{N}$  shows a more saturated scaling (3x instead of 5x) and vanishes after about one year, whereas the total ozone loss stays at some constant level in the second half of the year, stemming mainly from the lower stratosphere. In our setup we would therefore expect only a smaller feedback on the circulation compared to the setup of Sukhodolov et al. (2017).

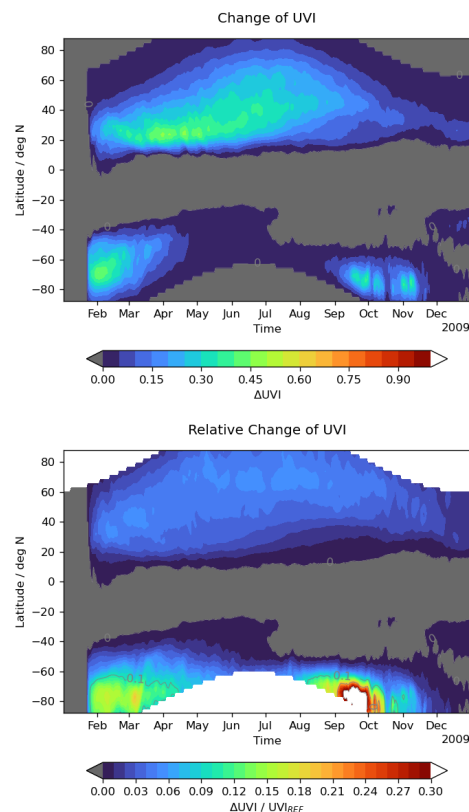
#### 5.4 Increase of UV-Index caused by the ozone loss

In order to estimate any health related impact by the reduced total ozone column, we calculate the corresponding change in the UV index. The UV-Index (UVI) is calculated from the spectrally weighted and integrated solar irradiance where the weight function takes the erythemal action spectrum into account (CIE International Commission on Illumination (2019)). Here we use the simplified formula of Allaart et al. (2004), their equation (8) together with equation (6), to estimate the increase of UVI connected to the decrease of the total ozone column ( $\Sigma\text{O}_3$ ) after a solar extreme event. In this formula, besides the dependence on the extraterrestrial flux and the solar zenith angle, for cloud-free situations UVI depends on  $\Sigma\text{O}_3$  only. For the column down to 7 km, the lower boundary of the model, we get an increase of the UVI for noon as shown in Fig. 11. The absolute highest change of about 0.5 UVI occurs in the months March and April at latitudes of  $25^\circ$  North which include the region on Earth with the highest population density. The relative change of UVI reaches about 4 - 5% in the northern midlatitudes where also the summer months are affected. In order to set this change in a context, we remark that the ozone depletion by CFCs reached a maximum reduction of total ozone of about 3% in northern midlatitudes in 1995 (Weber et al. (2022)), quite similar to the values we found for the extreme event (see Fig. 9), but lasting for more than a decade.

There is also a significant increase of UVI at high Southern latitudes, but not at latitudes which have a high population density.

## 6 Conclusions

The purpose of this study was to analyze direct effects of an extreme solar eruption on the chemical state of the middle atmosphere. We considered an extreme solar event with a strong SPE connected also with a strong GMS. The setup of the numerical experiment was chosen to yield a maximum chemical impact in the Northern Hemisphere. This was achieved by setting the event to Northern Hemisphere winter in combination with an elevated stratosphere SSW event. The main findings are that the GMS component from energetic electrons even in this dynamical situation leads only to a minor ozone reduction. The ozone reduction in our extreme scenario stems essentially from the ESPE and survives through the following boreal summer. The additional  $\text{NO}_y$  amount in the middle atmosphere scales with the strength of the event, as does the ozone loss. The studied event yields a significant but limited enhancement of the UV index in populated latitudes in the first months following the event of a few percent. The ESPE causes a strong impact in terms of additional  $\text{NO}_y$  and ozone loss in southern high latitudes, despite the event had been put to Northern Hemisphere winter to minimize photochemical loss of  $\text{NO}_y$ . This is related to the deep penetration of the energetic particles at altitudes where the loss of  $\text{NO}_y$  is negligible, and the early on-set of the downward transport with the beginning of the Antarctic winter in the middle atmosphere. Besides the related ozone loss, the denitrification



**Figure 11.** Change of the UV index by the extreme event for noon and the correspondent relative change of the UV index.

of NO<sub>y</sub> in the lower Antarctic stratosphere and the following washout in the troposphere can yield potentially measurable NO<sub>3</sub> signals in ice cores. This, however, has not been observed so far, perhaps because the signal strength strongly depends on the timing of the event. Therefore, nitrate signals in ice cores are not reliable indicators for solar activity.

*Author contributions.* TR prepared the experiments and wrote the paper, IU, JMW and OY provided ionization rates, MS critically discussed and improved the manuscript

*Competing interests.* The authors declare that no competing interests are present.

*Acknowledgements.* This work was supported by the German Federal Ministry of Education and Research within the research program ROMIC-2 and by the Academy of Finland (projects ESPERA, No. 321882). The AIMOS model is funded by the German Science Foundation

(DFG project WI4417/2-1). J.M. Wissing also thanks to support from the German Aerospace Center (DLR). The authors acknowledge the NOAA National Centers for Environmental Information (<https://ngdc.noaa.gov/stp/satellite/poes/dataaccess.html>) for the POES and Metop particle data used in this study.

330 Model data used for this analysis are accessible through <https://dx.doi.org/10.35097/917>.

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