

1 Natural Marine Cloud Brightening in the Southern Ocean

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23 Abstract: The number of cloud droplets per unit volume ( $N_d$ ) is a fundamentally  
24 important property of marine boundary layer (MBL) liquid clouds that, at constant liquid  
25 water path, exerts considerable controls on albedo. Past work has shown that regional  
26  $N_d$  has direct correlation to marine primary productivity (PP) because of the role of  
27 seasonally varying biogenically-derived precursor gases in modulating secondary  
28 aerosol properties. These linkages are thought to be observable over the high latitude  
29 oceans where strong seasonal variability in aerosol and meteorology covary in mostly  
30 pristine environments. Here, we examine  $N_d$  variability derived from five years of MODIS  
31 level-2 derived cloud properties in a broad region of the summer Eastern Southern Ocean  
32 and adjacent marginal seas. We demonstrate latitudinal, longitudinal, and temporal  
33 gradients in  $N_d$  that are strongly correlated with the passage of air masses over high PP  
34 waters that are mostly concentrated along the Antarctic Shelf poleward of 60°S. We find  
35 that the albedo of MBL clouds in the latitudes south of 60°S is significantly higher than  
36 similar LWP clouds north of this latitude.

37  
38 Short Summary: The number cloud droplets per unit volume is a significantly  
39 important property of clouds that controls their reflective properties. Computer models of  
40 the Earth's atmosphere and climate have low skill at predicting the reflective properties of  
41 Southern Ocean clouds. Here we investigate the properties of those clouds using satellite  
42 data and find that the cloud droplet number in the Southern Ocean is related to the  
43 oceanic phytoplankton abundance near Antarctica and cause clouds there to be  
44 significantly brighter than clouds further north.

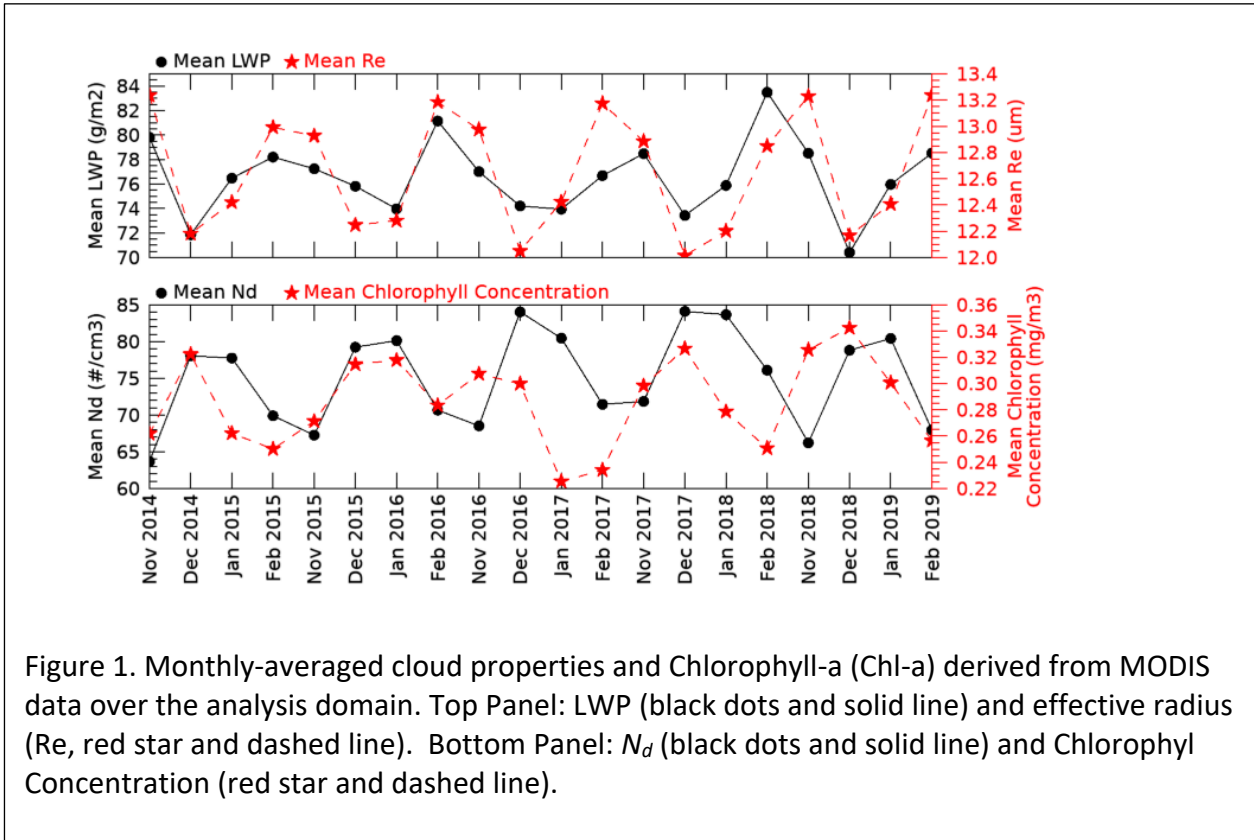
## 45 46 1. Introduction

47  
48 The cloud and precipitation properties of the Southern Ocean (SO) have received  
49 considerable attention since Trenberth and Fasullo (2010) identified a high bias in  
50 surface-absorbed solar energy there (McFarquhar et al., 2020). This bias has been  
51 traced to erroneously small Marine Boundary Layer (MBL) cloud cover in simulations of  
52 the Southern Ocean climate (Bodas-Salcedo, et al., 2016; Naud et al., 2016). The  
53 actual SO cloud climatology and associated albedo are dominated by geometrically thin  
54 MBL clouds (Mace et al., 2010; Mace et al., 2020, 2021). Because the predominant  
55 shallow boundary layer clouds rarely precipitate (Huang et al., 2016), they are sensitive  
56 to cloud condensation nuclei (CCN) concentrations (Twohy and Anderson, 2008;  
57 Painemal et al., 2017).

58  
59 In the SO, the CCN seasonal cycle (Ayers and Gras, 1991; Vallina et al. 2006; Gras and  
60 Keywood, 2017) is reflected in basin-wide cloud property variations (Krüger and Graßl,  
61 2011). McCoy et al. (2015) and Mace and Avey (2017) also found that MODIS- and A-  
62 Train-derived cloud properties over the SO, demonstrate a similar seasonal cycle in  
63 cloud droplet number concentration ( $N_d$ ) as for CCN. The basin wide variability in CCN  
64 and cloud albedo have been shown to be correlated with marine primary productivity  
65 (PP – defined as the net organic matter, mostly produced by phytoplankton, that is  
66 suspended in the ocean; Vallina et al., 2006; Krüger and Graßl, 2011; McCoy et al.,  
67 2015). McCoy et al. (2020) argue that the SO can be viewed as an analog of the  
68 preindustrial Earth. Given the large natural seasonal variability in CCN and clouds, the

69 SO is a natural laboratory to understand the processes that contribute to simulated  
 70 aerosol-related indirect forcing variability in climate models (Carslaw et al. 2013).

71  
 72 CCN and cloud droplet  $N_d$  in the SO are higher in Summer when significant latitudinal  
 73 gradients have been documented in the SO Australasian sector (Humphries et al.,  
 74 2021). Using time of flight aerosol chemical speciation monitor (ACSM) and ion



75 concentrations from filter samples, Humphries et al., (2021) analyzed the covariance of  
 76 aerosol chemistry, CCN at 0.5% supersaturation, and Condensation Nuclei (CN) larger  
 77 than 10 nm collected aboard Australian research vessels during the 2018 Austral  
 78 Summer (McFarquhar et al., 2021). While sulfates were a major compositional  
 79 component of aerosol at all latitudes during summer these compounds were in higher  
 80 fractional abundance poleward of 65°S where overall CCN numbers were higher by  
 81 ~50%. Chloride derived from sea salt was dominant in the region equatorward of 65°S  
 82 but was mostly absent south of 65°S. The ratio of CCN to CN at 0.5% supersaturation  
 83 increased considerably south of 65°S suggesting unique aerosol chemical processes  
 84 compared to the open ocean. Humphries et al. (2021) also discusses how this  
 85 compositional boundary in aerosol chemistry is often very distinct in the East Antarctic  
 86 waters between 60°S and 65°S. Following Humphries et al. we will refer to this belt as  
 87 the Atmosphere Compositional Front of Antarctica (ACFA). Humphries et al. (2021)  
 88 conclude that aerosol, newly condensed from gas phase sulfur species such as from  
 89 the oxidation of dimethyl sulfide (DMS), are an important component of high latitude  
 90 CCN. These products of phytoplankton physiology are released into the atmosphere

91 from the highly productive waters from ~60°S to the Antarctic – a region well known for  
92 a vast marine food web (Deppler and Davidson, 2017; Behrenfeld et al., 2016).

93  
94 Mace et al. (2021a) derived  $N_d$  and other cloud microphysical properties from non-  
95 precipitating stratocumulus clouds using shipborne remote sensing data. They found  
96 that stratiform clouds poleward of the ACFA had significantly higher  $N_d$  than  
97 equatorward. One particular case took place when the Icebreaker Aurora Australis was  
98 at the Davis Antarctic station just east of Prydz Bay (~77°E) between 1 and 5 January  
99 2018 and featured nearly continuous high  $N_d$  clouds ( $> 150 \text{ cm}^{-3}$ ) occurring in a  
100 southerly flow passing over the ship that had trajectories from the Antarctic Continent.  
101 Similarly, Twohy et al., (2021) report that the highest concentrations of aerosol  
102 composed primarily of non-sea salt sulfates in the free troposphere north of 60°S  
103 observed from research aircraft in Summer 2018 had occurred in airmasses that had  
104 originated recently from over the Antarctic continent. See also Shaw et al. (1988) for an  
105 early examination of the role of biogenic sulfate in modulating summertime aerosol  
106 along coastal Antarctica. Shaw et al. (2007) expands on this idea as does Korhonen et  
107 al., (2008).

## 111 2. Results

113 See Appendix A for methods and definitions. Approximately 40,000 1° latitude by 2°  
 114 longitude MBL cloud scenes per month meet our criteria for liquid phase non  
 115 precipitating clouds in the analysis domain. This number varies by ~25% in a seasonal  
 116 cycle that is due mostly to our solar zenith angle criteria. A seasonal cycle is evident in  
 117 the monthly-averaged cloud properties. LWP and  $r_e$  have seasonal minima in the  
 118 months of December and January. Due to an  $r_e^{-5/2}$  dependence,  $N_d$  is of opposite phase

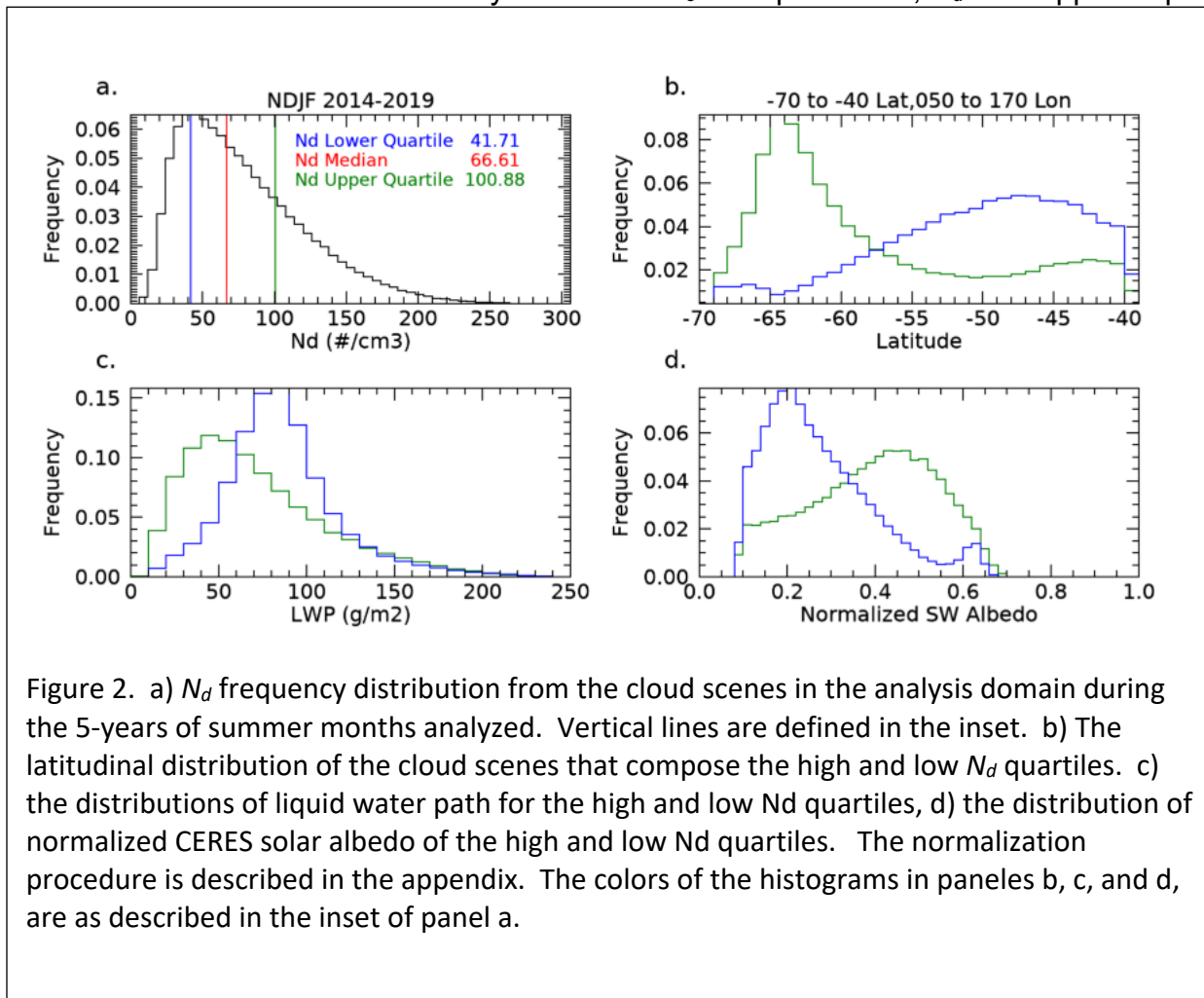


Figure 2. a)  $N_d$  frequency distribution from the cloud scenes in the analysis domain during the 5-years of summer months analyzed. Vertical lines are defined in the inset. b) The latitudinal distribution of the cloud scenes that compose the high and low  $N_d$  quartiles. c) the distributions of liquid water path for the high and low  $N_d$  quartiles, d) the distribution of normalized CERES solar albedo of the high and low  $N_d$  quartiles. The normalization procedure is described in the appendix. The colors of the histograms in panels b, c, and d, are as described in the inset of panel a.

119 with  $r_e$  and correlated with it at -0.93. The seasonal variability in LWP ( $r_e$ ) is on the  
 120 order of 7% (4%) and is small in comparison to  $N_d$  (~25%).  $\tau$  and  $r_e$  are derived from  
 121 the visible and near infrared reflectances with the MODIS level 2 retrieval algorithm  
 122 (Nakajima and King, 1990). LWP is, then, derived from

$$\tau = \frac{3}{2\rho_w} \frac{LWP}{r_e} \quad \text{Equation 1.}$$

125 that is derived in Stephens (1978). It is reasonable to consider whether seasonal  
 126 variations in  $N_d$ , perhaps linked to CCN, might be associated with variability in LWP.  
 127 We find that LWP decreases as  $N_d$  increases with a correlation coefficient in the  
 128 monthly means of -0.60 in the monthly means.

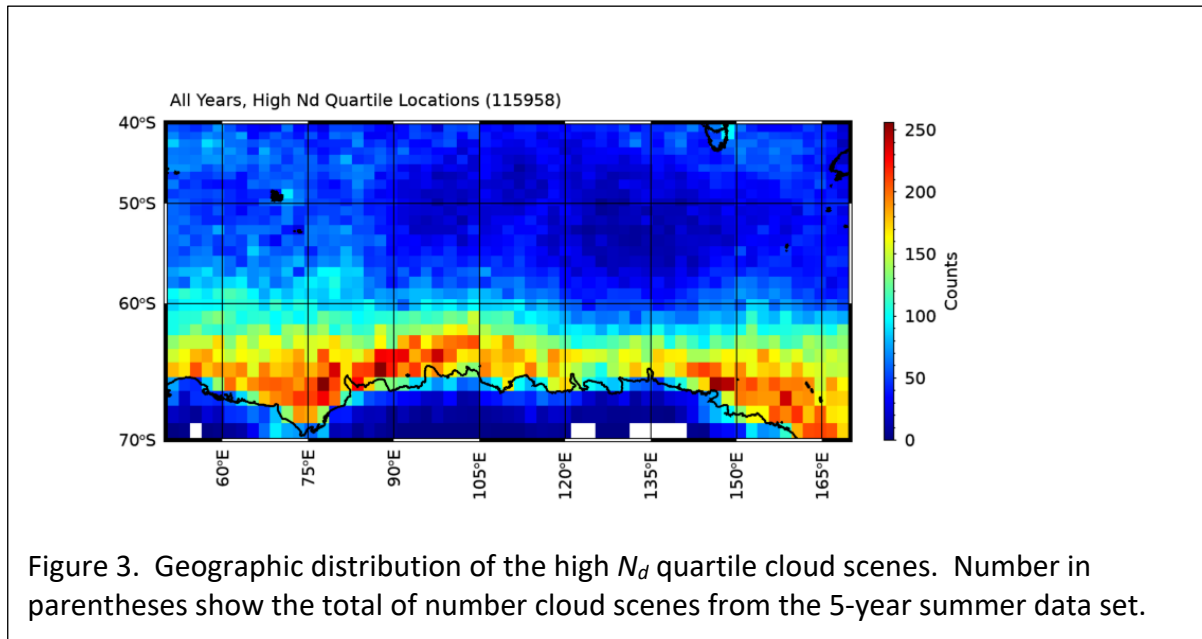
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131 In four of the five years, we see by inspection of Figure 1 that Chl-a leads changes in  $N_d$   
132 by approximately 1 month. The correlation coefficient of  $N_d$  and Chl-a increases from  
133 0.27 to 0.60 when  $N_d$  is lagged from 0 to 1 month in the Figure 1 time series although  
134 this result should be interpreted with caution given the break between February and  
135 November in the time series. These results are broadly like those presented by McCoy  
136 et al., (2015) and Mace and Avey (2017). McCoy et al. (2015) link  $N_d$  variations to PP  
137 using regression analysis of MODIS derived  $N_d$  against a biogeochemical  
138 parameterization of biogenic sulfate and organic mass fraction (See also Lana et al.,  
139 2012).

140  
141 We find a broad distribution of scene-averaged  $N_d$  (Figure 2a) with median, lower and  
142 upper quartile values of  $66 \text{ cm}^{-3}$ ,  $42 \text{ cm}^{-3}$  and  $101 \text{ cm}^{-3}$  respectively. Henceforth, we  
143 focus our analysis on the groups of scenes that are less than and greater than the  
144 upper and lower quartiles. The high and low  $N_d$  scenes have distinct latitudinal  
145 occurrence distributions (Figure 2b) with low  $N_d$  scenes peaking broadly at  $48^\circ\text{S}$  while  
146 the high  $N_d$  scenes demonstrate a modal occurrence near  $64^\circ\text{S}$ . Overall, the  $N_d$  gradient  
147 implied by Figure 2 is correlated with the latitudinal distribution of imager-derived Chl-a  
148 (i.e., Deppler and Davidson, 2017). The seasonally averaged  $N_d$  gradient is also  
149 discussed in McCoy et al., (2020). Differentiating seasonally varying properties north  
150 and south of the ACFA (not shown), we find a clear differentiation in  $r_e$  and  $N_d$  with  
151 smaller  $r_e$  south of the ACFA (mean  $r_e \sim 11 \mu\text{m}$ ,  $N_d \sim 100$ ) compared to north (mean  
152  $r_e \sim 13 \mu\text{m}$ ,  $N_d \sim 67 \text{ cm}^{-3}$ ). LWP is slightly larger by  $\sim 7\%$  south of the ACFA. Both  
153 regions have a distinct seasonal cycle in cloud properties shown in Figure 1 although  
154 the southern latitudes have larger interannual variability likely owing to variations in  
155 annual sea ice extent and melt. The LWP distribution of the high  $N_d$  quartile is  
156 significantly shifted to lower values compared to the low  $N_d$  quartile LWP distribution  
157 (Figure 2c). This finding is in accordance with the observational and theoretical work  
158 presented in Glassmeier et al., (2021) who argue that closed cell stratocumulus that  
159 dominate the clouds examined here have increased entrainment drying under higher  $N_d$   
160 conditions. Figure 2c and 2d illustrate that even though the high  $N_d$  quartile scenes tend  
161 to have lower LWP, their solar albedo ( $A$ ) tends to be significantly higher than the low  
162  $N_d$  quartile scenes illustrating the influence of cloud microphysics on the radiative  
163 forcing of these clouds.

164  
165 The high  $N_d$  scenes occur predominantly poleward of the ACFA (Figure 3). Interestingly  
166 we find that the latitudinal gradient weakens slightly west of  $90^\circ\text{E}$  with a broad region of  
167 higher  $N_d$  occurrence in the vicinity of the Kerguelen Rise where PP is higher (Cavagna  
168 et al., 2015). Establishing causality between regions of high PP and cloud properties is  
169 challenging (i.e., Meskhidze and Nenes, 2006; Miller and Yuter, 2008). While we find  
170 seasonal associations over broad regions here, the chain of causality between  
171 phytoplankton and clouds is not immediate or even necessarily direct because the  
172 chemical processes take time to evolve and can move along chemical pathways that  
173 have divergent outcomes (Woodhouse et al., 2013). To increase cloud  $N_d$ , new CCN  
174 must be formed. Formation of new CCN can occur when sulfur compounds emitted  
175 from the ocean surface nucleate after oxidation in the presence of sunlight. This  
176 process of new particle formation occurs in the absence of other aerosol and often

177 requires mixing of the gaseous compounds from the boundary layer into the low-aerosol  
178 free-troposphere where the newly formed aerosol can be transported widely (Shaw,  
179 2007; Korhonen et al., 2008). Other pathways are possible such as deposition of  
180 sulfate compounds onto primary sea salt particles that modify the chemical properties of  
181 existing CCN rather than nucleating new CCN (Fossum et al., 2020) or even removal of  
182 sulfur compounds from the gas phase via aqueous phase oxidation in clouds  
183 (Woodhouse et al., 2013).  
184



185 Given the foregoing discussion, it seems reasonable that an air mass that is producing  
186 clouds with certain features could be interacting with an aerosol population that has  
187 evolved over periods of days (Brechtel et al., 1998). In addition, natural cloud  
188 processes such as collision and coalescence of drops tend to cause  $N_d$  to decrease  
189 while precipitation efficiently scavenges CCN, thereby lowering CCN concentration and  
190 even modifying their composition and size through aqueous processing (Hoppel et al.,  
191 1986). With larger  $r_e$  north of the ACFA, the collision-coalescence process is likely more  
192 active (Freud and Rosenfeld, 2012) and could explain the latitudinal difference in  
193 adiabaticity (see methods) found in in situ data. For instance, Kang et al. (2022)  
194 analyzed data collected from Macquarie Island (54.6°S, 158.9°E) and found that, not  
195 only were most clouds drizzling, but that precipitation as light as 0.01 mm hr<sup>-1</sup> could  
196 reduce  $N_d$  by ~50%. Therefore, a cloud field should be considered as the product of  
197 both local dynamics and thermodynamics primarily with modulation by a local  
198 population of CCN. To examine the role of air mass history, we calculate the 5-day back  
199 trajectories using the Hybrid Single-Particle Lagrangian Integrated Trajectory (HYSPLIT;  
200 Stein et al., 2015) model using the Global Data Assimilation System (GDAS; Kamitsu,  
201 1989) as input. The parcel's endpoint is the central latitude and longitude of the cloud  
202 scene, and the location and model output are stored hourly.  
203

204 South of the ACFA, the histories of the populations tend to be statistically different  
205 (Figure 4). The low  $N_d$  clouds are more likely to be observed in airmasses that have  
206 trajectories that originated in the open ocean region to the north of the ACFA. High  $N_d$   
207 scenes rarely evolve in airmasses that originate in the open ocean to the north of the  
208 ACFA. The likelihood is that an airmass that has produced a high  $N_d$  cloud scene south  
209 of the ACFA latitude has spent most of the previous 5 days over latitudes south of the  
210 ACFA. North of the ACFA, the latitude distributions during the months of November and  
211 February (not shown) are essentially identical for the high and low  $N_d$  quartiles.  
212 However, for December and January, we find that the high  $N_d$  clouds observed north of  
213 the ACFA have an increased likelihood of trajectories emanating from south of the  
214 ACFA during the 5-days prior to the MODIS observation.

215

### 216 3. Discussion and Conclusions

217

218 Using MODIS level 2 cloud property retrievals and the technique developed in  
219 Grosvenor et al. (2018; hereafter G18) to estimate  $N_d$ , we examine the latitudinal and  
220 seasonal cycles of non-precipitating liquid-phase clouds in the Australasian sector of the  
221 Summertime Southern Ocean. The  $r_e$  and  $N_d$  have distinctive differences north and  
222 south of the ACFA but demonstrate similar seasonal cycles. We infer that the spatial  
223 and temporal variability in cloud  $N_d$ , and  $r_e$  are at least partially a function of the  
224 geographic and temporal variability in CCN that, in turn, is related to the seasonality of  
225 primary sources such as sea salt and the latitudinal variability in marine PP. The  
226 highest  $N_d$  clouds tend to be overwhelmingly found along the East Antarctic coastal  
227 waters south of the ACFA.

228



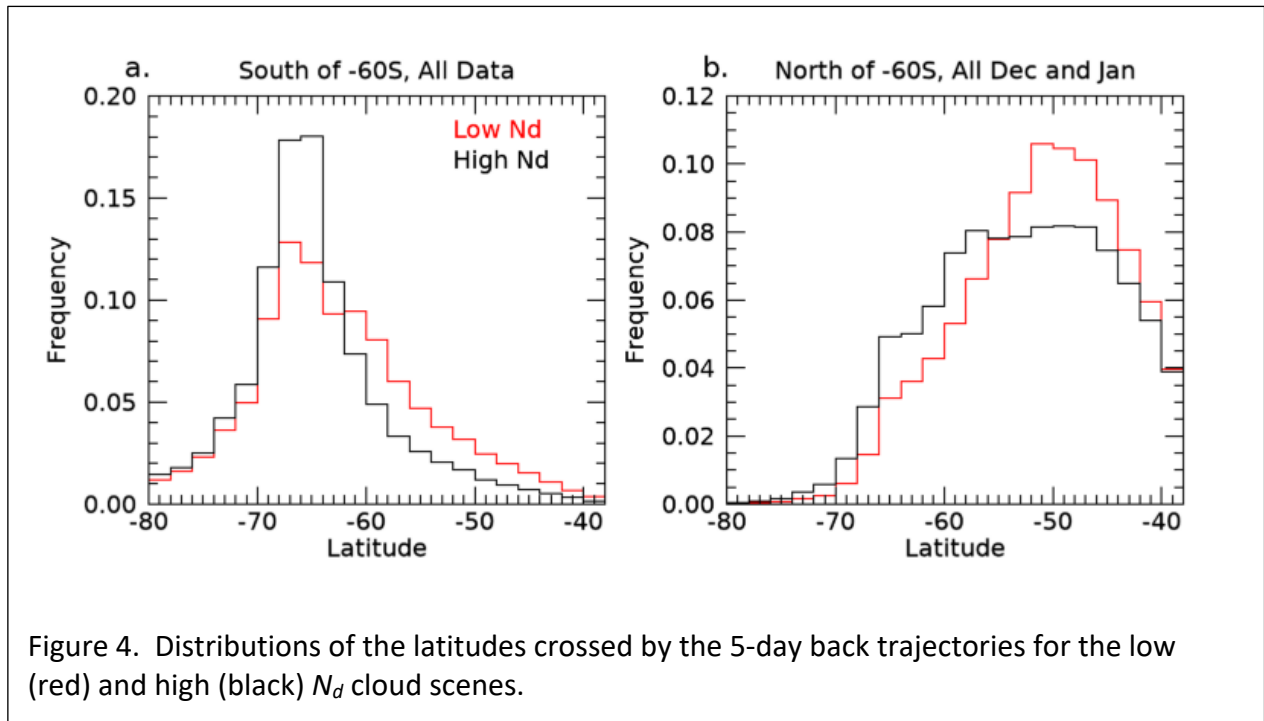


Figure 4. Distributions of the latitudes crossed by the 5-day back trajectories for the low (red) and high (black)  $N_d$  cloud scenes.

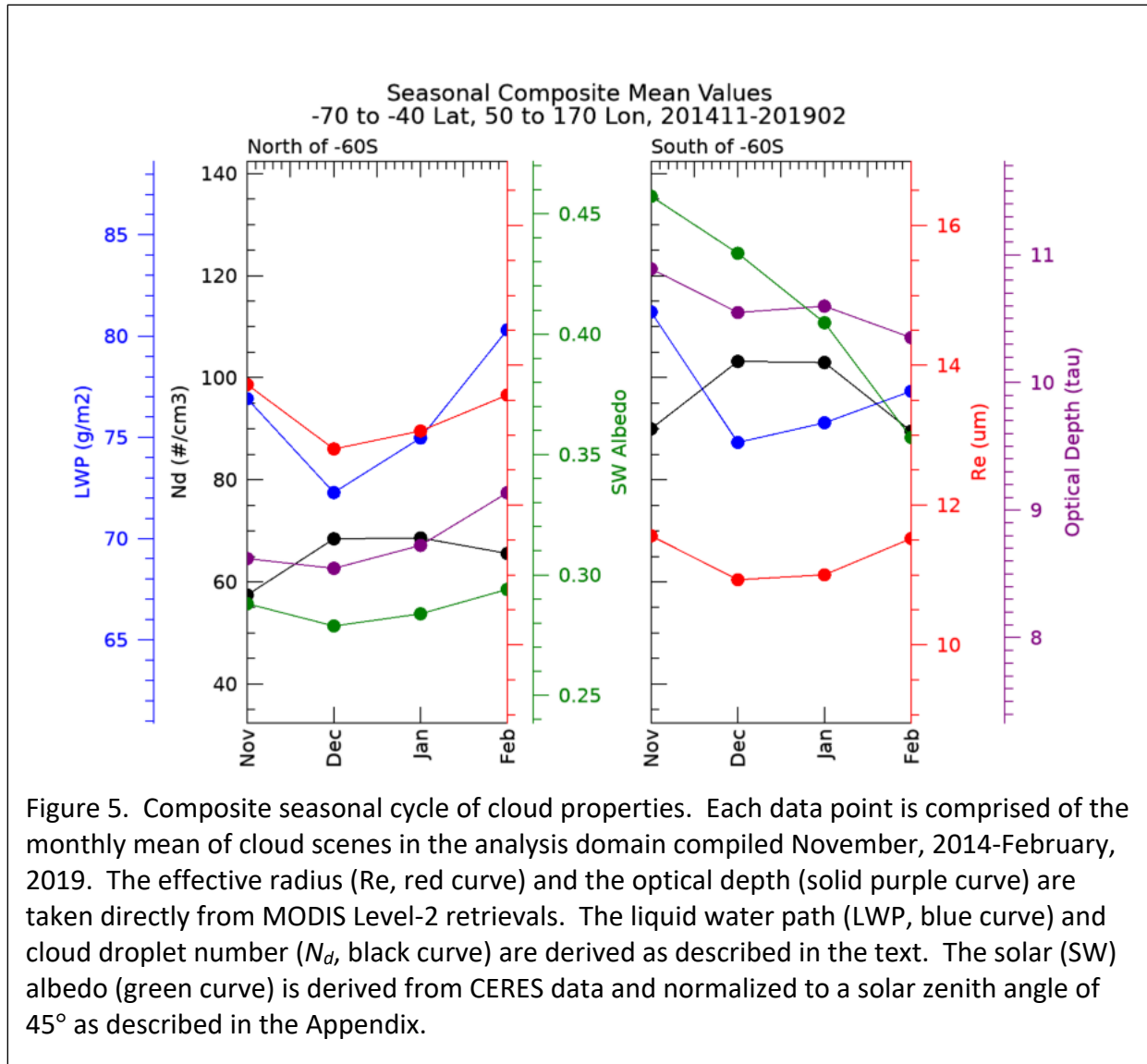
229 Because aerosol precursor gasses like DMS often require trajectories through the free  
 230 troposphere to nucleate new particles that then take time to reach CCN sizes  
 231 (Korhonen et al., 2008; Shaw et al. 2007), we examine the back trajectories of the  
 232 airmasses observed with high and low  $N_d$  south of the ACFA and find significant  
 233 differences. Low  $N_d$  cloud scenes are more likely to have arrived south of the ACFA  
 234 from northerly trajectories that would have transported low CCN air dominated by sea  
 235 salt. The high  $N_d$  cloud scenes are more likely to have trajectories that have remained  
 236 adjacent to or had passed over the Antarctic continent. North of the ACFA, while the  
 237 trajectory statistics for the high and low  $N_d$  quartiles in November and February are  
 238 nearly identical, during December and January the high  $N_d$  clouds scenes tend to have  
 239 an increased likelihood of arriving north of the ACFA from southerly trajectories,  
 240 suggesting that high CCN airmasses are being transported northward especially during  
 241 December and January.

242  
 243 Given that the main difference between the source regions north and south of the ACFA  
 244 is the magnitude of the marine PP, and given previous analyses of CCN compositional  
 245 sensitivity to marine biological factors (e.g. Humphries et al., 2021; Vallina et al., 2006;  
 246 Lana et al., 2012; McCoy et al, 2015), we conclude that the biological source of sulfate  
 247 precursor gasses and the slackening of surface winds with latitude during Summer  
 248 plays a dominating role in controlling the latitudinal gradients in the properties of weakly  
 249 precipitating MBL cloud fields over the Southern Ocean. Figure 5 summarizes our  
 250 findings by presenting composite seasonal cycles of MBL cloud scenes north and south  
 251 of 60°S. The LWP in both latitudinal bands go through a weak seasonal cycle. The  
 252 significant contrast in optical depth between the northern and southern bands is, we  
 253 infer, mostly caused by the latitudinal contrast in  $N_d$ . Based on available evidence, we  
 254 conclude that the differences in  $r_e$  in MODIS retrievals are causally linked to oceanic PP

255 gradients that drive CCN, and thereby  $N_d$ , to be higher over the southern region. This  
256 sensitivity, in turn, plays a significant role in modulating the regional albedo ( $A$ ) and,  
257 thereby, influences the input of sunlight to the surface ocean. We note that the  
258 seasonal cycle in  $A$  is different between the northern and southern latitude domains (a  
259 topic for future work), however, always  $A$  of the southern domain is higher than that of  
260 the northern domain. However, we should be careful not to overstate this case. Cloud  
261 processes that consume  $N_d$  and modify CCN (i.e. precipitation and cloud processing)  
262 also play a role in modulating cloud  $N_d$  and therefore regional  $A$  (Kang et al., 2022;  
263 McCoy et al., 2020). The airmass history and source region, while apparently  
264 important, are among many factors involved.

265  
266 Since the magnitude of PP is significantly lower north of the ACFA throughout the  
267 summer season, a similar seasonal cycle in  $N_d$  and  $r_e$  suggests that CCN derived from  
268 DMS oxidation of precursor gasses emitted primarily from Antarctic coastal waters  
269 perhaps seeds much of the rest of the Southern Ocean with biogenic sulfate aerosol as  
270 observed in recent airborne observations (Twohy et al., 2021). The northerly transport  
271 of these high sulfate airmasses out of the Antarctic coastal waters (Figure 4b) and  
272 southerly transport of low sulfate air masses into the Antarctic coastal region near the  
273 surface (Figure 4a) have been reported by Humphries et al. (2016, 2021) and Shaw  
274 (1988) and observed in the free troposphere with recent research aircraft  
275 measurements (Twohy et al. 2021).

276  
277 Our ability to identify natural marine cloud brightening (Latham et al., 2008) due to  
278 aerosol-cloud coupling is a direct result of the absence of other anthropogenic and  
279 continental influences in the pristine SO. As argued by McCoy et al. (2020), it seems  
280 clear that in several important ways, the Southern Ocean is the last vestige of the  
281 preindustrial atmosphere allowing us to constrain processes that remain important to  
282 our understanding of the global climate (Carslaw et al., 2013).



## 284 Appendix. Methods

285

286 We use MODIS imager-derived Level-2 retrievals (Platnick et al., 2015) of effective  
 287 radius ( $r_e$ ) and optical depth ( $\tau$ ) from five summer periods (2014-2019) collected  
 288 between the latitudes of  $45^\circ\text{S}$  and  $76^\circ\text{S}$  and longitudes of  $40^\circ\text{E}$  and  $170^\circ\text{E}$  to focus  
 289 roughly on where the ships and aircraft sampled in Summer 2017-18. We calculate  $N_d$   
 290 using the method derived and evaluated in G18:

291

$$292 \quad N_d = \frac{\sqrt{5}}{2\pi\kappa} \left( \frac{f_{ad}c_w\tau}{Q_{ext}\rho_w r_e^5} \right)^{1/2} \quad (\text{A1})$$

293

294 where  $\rho_w$  is the density of liquid water ( $1 \text{ g cm}^{-3}$ ),  $f_{ad}$  is an adiabaticity assumption,  $c_w$  is  
 295 the vertical derivative of the adiabatic liquid water content,  $Q_{ext}$  is the extinction efficiency  
 296 that is typically assumed to be 2 for cloud droplets, and  $\kappa$  is the cubed ratio of  $r_e$  to  $r_v$ . As

297 noted by G18,  $N_d$  depends on  $r_e^{-5/2}$ , which implies that the sensitivity or the rate of change  
298 of  $N_d$  to retrieved  $r_e$  goes as the  $-7/2$  exponent. Any biases in  $r_e$ , then would significantly  
299 bias  $N_d$ . G18 provide a thorough evaluation of the sources of uncertainty in  $N_d$  due to  
300 assumptions of adiabaticity, scene heterogeneity, etc., and conclude that  $N_d$  derived  
301 using equation 1 applied to MODIS cloud retrievals has an overall uncertainty of  $\sim 80\%$ .

302  
303 The most uncertain quantity in the assumptions used in Equation A1 is  $f_{ad}$  since the cloud  
304 vertical structure is not constrained by MODIS measurements. Using cloud thickness from  
305 ship-based cloud radar and lidar along with retrieved LWP from collocated microwave  
306 radiometer (Mace et al., 2021a), we estimate the value of  $f_{ad}$  in nonprecipitating  
307 stratocumulus observed during the summer of 2018 (McFarquhar et al., 2021). We find  
308 that the mean and standard deviation of  $f_{ad}$  north of the ACFA is 0.66 and 0.48,  
309 respectively. South of the ACFA, the mean and standard deviation of  $f_{ad}$  is 0.93 and 0.60,  
310 respectively. For the calculations of  $N_d$  in equation A1, we use a constant value for  $f_{ad}$  of  
311 0.8.  $N_d$  is proportional to the square root of  $f_{ad}$ , therefore,  $\frac{\partial \ln N_d}{\partial \ln f_{ad}} = \frac{1}{2}$  and a fractional  
312 variation in  $f_{ad}$  on the order of 0.5 would imply an uncertainty in  $N_d$  of 0.25. Furthermore,  
313 we expect in regions with  $f_{ad}$  higher (lower) than 0.8 the  $N_d$  would be biased low (high).  
314 As we show, the regions with higher  $N_d$  tend to be in the south and lower  $N_d$  in the north  
315 counter to these expected biases. Additionally in this study, we will be examining  
316 differences in spatially averaged  $N_d$  that are greater than a factor of 2. These results  
317 imply that bias and random error due to uncertainty in  $f_{ad}$  is unlikely to significantly  
318 influence the qualitative findings of this study.

319  
320 Another source of systematic bias could be from the quantity  $\kappa$  that can be shown to be  
321 a function of the variance of the droplet size distribution and is assumed to be a constant  
322 at 0.7. G18 discusses this issue in some detail and concludes that there may be  
323 systematic biases on the order of 12% that could be a function of  $N_d$  in pristine conditions.  
324 While this quantity can be investigated with data collected in situ, no such data exists in  
325 stratocumulus clouds south of the ACFA. Therefore, we recognize a potential source of  
326 bias due to  $\kappa$  that is likely much smaller than the systematic latitudinal differences we find.

327  
328 Given the uncertainties in  $N_d$  at the pixel level, we implement a filtering and averaging  
329 scheme to focus on liquid phase, weakly precipitating cloud scenes. We define a scene  
330 as a  $1^\circ$  latitude by  $2^\circ$  longitude domain where pixels are reported in the MODIS L2 data  
331 to be of liquid-phase. We assume that clouds are weakly precipitating clouds if the cloud  
332 liquid water path (LWP)  $< 300 \text{ g m}^{-2}$ . We require that the sensor and solar zenith angles  
333 ( $\theta$ ) at that pixel are less than  $30^\circ$  and  $60^\circ$ , respectively. The maximum  $\theta$  requirement is  
334 motivated by the findings of Grosvenor and Wood (2014) who find that systematic errors  
335 in MODIS retrievals increase significantly for  $\theta > 60^\circ$ . The  $\theta$  requirement causes us to  
336 focus on the months from November through February. We require at least 1000 1-km  
337 resolution pixels with these characteristics to exist within a scene (typical number  
338  $> 10000$ ). In addition, we require that no more than 10% of the pixels have a cloud top  
339 temperature less than  $-20^\circ\text{C}$  to ensure the absence of ice phase hydrometeors. Cloud  
340 properties within a scene are averaged.

341

342 Collocated cloud albedos ( $A$ ) of the cloud scenes are analyzed.  $A$  is derived from the  
343 Clouds and the Earth's Radiant Energy System (CERES) Energy Balanced and Filled  
344 (EBAF) version 4.0 (Loeb et al, 2018) data collected using instruments on board Aqua  
345 and Terra. The albedo is derived by dividing the upwelling shortwave flux at the top of  
346 the atmosphere (TOA) by the downwelling shortwave flux at TOA. Because  $A$  has a  
347 solar zenith angle dependence, (Minnis et al. 1998), we normalize all albedo values to  
348  $\theta=45^\circ$  (approximately the mean value of  $\theta$  for the analysis domain and months  
349 analyzed) with an empirical method using theoretically calculated  $A$  ( $\hat{A}$ ) as a function of  
350 latitude presented in Minnis et al. (1998 – their figure 7). The normalization is  
351 implemented by first approximating the latitudinal dependence of  $A$  for various cloud  
352 optical depths ( $\tau$ ) using the following regression equation:  $\hat{A} = 0.51 - 0.43\mu_0^{1/2} +$   
353  $0.17\ln \tau$  where  $\mu_0 = \cos \theta$ .  $\hat{A}$  approximates the variation of  $A$  with latitude within ~15%  
354 at  $\tau=8$ . The fit decreases in accuracy at higher and lower  $\tau$  increasing to an uncertainty  
355 of ~30% for  $\tau=2$  and  $\tau=32$  (these values of  $\tau$  (2, 8, 32) are those presented in Minnis et  
356 al., 1998, Figure 7). The averaged  $\tau$  of the MBL cloud scenes in our analysis is  
357 approximately between 9 and 11 (Figure 5) so we expect that  $\hat{A}$  is typically a reasonable  
358 approximation of  $A$ . The normalization of all  $A$  to  $\theta = 45^\circ$  is accomplished by  
359 multiplying the CERES  $A$  by the ratio  $\frac{\hat{A}(\mu_0(\theta=45),\tau)}{\hat{A}(\mu_0,\tau)}$  where  $\tau$  is from the MODIS cloud  
360 scene. The magnitude of the ratio applied to the data ranges from 0.85 at higher  
361 latitudes to 1.2 at lower latitudes with an average near 1.

362

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364 interpretation of the results. SB was responsible for implementing data analysis code  
365 and generation of figures. RH provided background on aerosol and provided insight  
366 regrading various aspects of the study. MPG and ES assisted GM in the study design  
367 and implementation.

368

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370

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376

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