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## Gravity waves generated by the high graupel/hail loading through buoyancy oscillations in an overshooting hailstorm

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**Abstract.** The convectively generated gravity waves (GWs) have important contributions on the stratospheric and mesospheric momentum and energy budget, and chemical composition, however, large uncertainties still remain about wave source properties and the associated wave-generated mechanisms. The formation mechanism and significant impacts of downward propagating GWs generated by a continental overshooting hailstorm occurred on 19 June 2017 in Beijing in the mid-latitude are reported in this study based on radar observations and simulated results from a three-dimensional cloud model with hail-bin microphysics. It is found that the overshooting storm penetrates the tropopause and enters the lower stratosphere in the mature stage. After the mature stage, the continuous descending process of the upper-level high graupel/hail loading causes the breaking of equilibrium between the buoyancy force and hydrometeor loading established in the mature stage and induces a restoring force of buoyancy, as well as buoyancy oscillations that excite downward propagating GWs. The GWs have a duration of about 20 min and the estimated wavelength of about 3-4 km. The downward propagating GWs not only result in the storm updraft splitting quickly, and significantly change the storm morphology and evolution, but also form the upward propagating GWs through surface reflection process, and induce strong vertical fluctuations in temperature and vertical velocity, and significantly change the dynamic and thermodynamic structure in the lower stratosphere.

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## 33 **1 Introduction**

34 Atmospheric gravity waves (GWs) excited by deep convection have long been  
35 focused (e.g., Pierce and Coroniti, 1966; Stull, 1976; Fovell et al., 1992; Alexander et  
36 al., 1995; Piani et al., 2000; Horinouchi et al., 2002; Snively and Pasko, 2003; Müller  
37 et al., 2018), since convectively generated GWs have been found to have significant  
38 contributions to the momentum and energy budget (Fritts and Alexander, 2003), and  
39 water vapor and tracers transport (Wang et al., 2002; Luderer et al., 2007) in the  
40 troposphere-to-stratosphere transport (TST), and even to the mesospheric chemical  
41 composition (Garcia and Solomon 1985).

42 The momentum flux associated with atmospheric GWs is observed much larger  
43 than those of Kelvin waves and Rossby-gravity waves (Sato and Dunkerton, 1997),  
44 and is an important driving forcing for some climate systems such as the  
45 quasi-biennial oscillation (QBO) (Alexander and Holton, 1997; Piani et al., 2000;  
46 Piani and Durran, 2001; Baldwin et al., 2001; Beres et al., 2002).

47 The properties of atmospheric GWs are described based on the fluctuations of  
48 vertical wind and temperature profiles observed by lidar and radar (Larsen et al., 1982;  
49 Smith et al., 1985; Tsuda et al., 1989, 1994; Sato, 1992), and the generation  
50 mechanisms have been intensively investigated (Weinstock, 1985; Dewan and Good,  
51 1986; Smith, 1987; Hines, 1991; Sato and Yamada, 1994; Warner and McIntyre, 1996;  
52 Nicholls and Pielke, 2000). The wave spectra, number and frequency of GWs have  
53 been also investigated (Lane and Moncrieff, 2008) and summarized (Gardner et al.,  
54 1993). The propagation and breaking of quasi-monochromatic small-scale GWs  
55 induced by thunderstorm activity were found to be closely associated with the  
56 observed airglow at the altitudes through the upper mesosphere and lower  
57 thermosphere (Snively and Pasko, 2003).

58 Atmospheric GWs can be generated by many sources, such as convection, wind  
59 shear, jet streams, frontal systems and topography, as well as pyro-cumulonimbus  
60 clouds (pyroCb) induced by large forest fires (Luderer et al., 2007). Three main  
61 mechanisms for the convectively generated GWs have been proposed. One is referred  
62 to the thermal forcing mechanism, in which, the GWs are generated by convective



63 clouds through the latent heat release (Holton,1973,2002; Salby and Garcia, 1987;  
64 Alexander et al., 1995; Mclandress et al., 2000; Fritts and Alexander, 2003). So that  
65 the latent heating profile in convective storms determines the properties of GWs. Two  
66 other mechanisms are referred to mechanical oscillation, such as the updraft  
67 oscillation (Clark et al.,1986; Fovell et al., 1992; Alexander et al.,1995) and transient  
68 mountain effect. However, the mechanisms for GWs generation by thermal forcing  
69 and updraft oscillation are not easily separated since they are intrinsically coupled in  
70 convective clouds. Lane et al. (2001) modeled GWs in maritime sea-breeze  
71 convection and indicated that the mechanical oscillation mechanism was dominant in  
72 GWs generation, while Song et al. (2003) suggested that the mechanical oscillation  
73 and thermal forcing mechanisms had comparable magnitudes in GWs generation.

74 Numerical models have become an important role in investigation of atmospheric  
75 GWs from single convective cloud models (Alexander et al.,1995; Fovell et al., 1992;  
76 Lane et al.,2001) to General Circulation Models (GCMs) with convection  
77 parameterization and convection-permitting schemes (Liu et al., 2014; Holt et al.,  
78 2016; Müller et al., 2018).

79 Most of previous studies have focused on the convectively generated GWs from  
80 thermal and mechanical oscillations of deep tropical convection and their influences  
81 on the stratospheric atmosphere. The relevant studies on the GWs generated by  
82 continental overshooting convection and their influences on the structure and  
83 evolution of storms, as well as the stratospheric atmosphere remain unclear, merit  
84 further investigation. Since the continent is the main region for human activity,  
85 understanding how GWs generated by continental storms influence the structure and  
86 evolution of storms, and the stratospheric atmosphere could be significant in storm  
87 tracking and forecasting, as well as transport of momentum, energy and pollution  
88 from the low troposphere to the upper atmosphere. The mountain-generated GWs  
89 under certain meteorological conditions have been found to have important roles in  
90 clouds and precipitation, as well as aerosol-cloud-precipitation interactions in  
91 northern China (Guo et al., 2013, 2017). In this study, the properties and generation  
92 mechanism of GWs, as well as the influences on both the storm itself and the



93 stratospheric atmosphere for a continental overshooting hailstorm occurred on 19 June  
94 2017 are reported.

95

## 96 **2 Methods**

### 97 **2.1 Data**

98 The radar data observed by an operational SA-band Doppler radar located in the south  
99 suburban of Beijing city are used to obtain the structure and evolution of the  
100 GWs-generated overshooting storm. The radar data are also used to validate the  
101 modeled storm. The sounding data from Beijing Meteorological station are used to  
102 obtain the environmental conditions for the storm.

### 103 **2.2 The model**

104 A three-dimensional fully compressible nonhydrostatic cloud model with hail-bin  
105 microphysics is employed to investigate the GWs properties, generation mechanism  
106 and the effects in this study (Guo and Huang, 2002). The formation, growth and  
107 conversion processes of cloud water, rainwater, cloud ice, snow and graupel/hail are  
108 included in the model. The Kessler-type scheme is used for the warm microphysical  
109 process (Kessler, 1969). The graupel/hail is categorized into 21 size bins ranging from  
110 100  $\mu\text{m}$  to nearly 7 cm in diameter. The model domain is on a standard spatially  
111 staggered mesh system. The time-splitting integration technique is used to treat  
112 high-frequency acoustic term (Klemp and Wilhelmson, 1978). The large integration  
113 time step is 5 s, while small time step is 0.25 s. The spatial difference terms are of  
114 second-order accuracy except for the advection term that has fourth-order accuracy.  
115 All other derivatives are evaluated with second-order centered differences. The  
116 radiation lateral boundary condition is applied and top boundary is rigid. A Rayleigh  
117 friction zone is used to absorb vertically propagating gravity waves near the top of the  
118 domain. The model uses a first-order closure for subgrid turbulence and a diagnostic  
119 surface boundary layer based on the Monin–Obukhov similarity theory.

120 The single sounding at 20:00 BST (Beijing Standard Time, the same hereafter) on  
121 June 19, 2017 in Beijing is used to initiate the simulation. A thermal bubble located in  
122 the central domain with a horizontal distance of 8 km and vertical distance of 4 km is



123 used for convection initiation in the model, and the maximum temperature  
124 perturbation in the central bubble is  $1.5^{\circ}\text{C}$ . The total integration time is 80 min. The  
125 domain size is 35 km in horizontal with a resolution of 1 km and 18.5 km in vertical  
126 with a resolution of 0.5 km.

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### 128 **3 Results**

#### 129 **3.1 Environmental conditions**

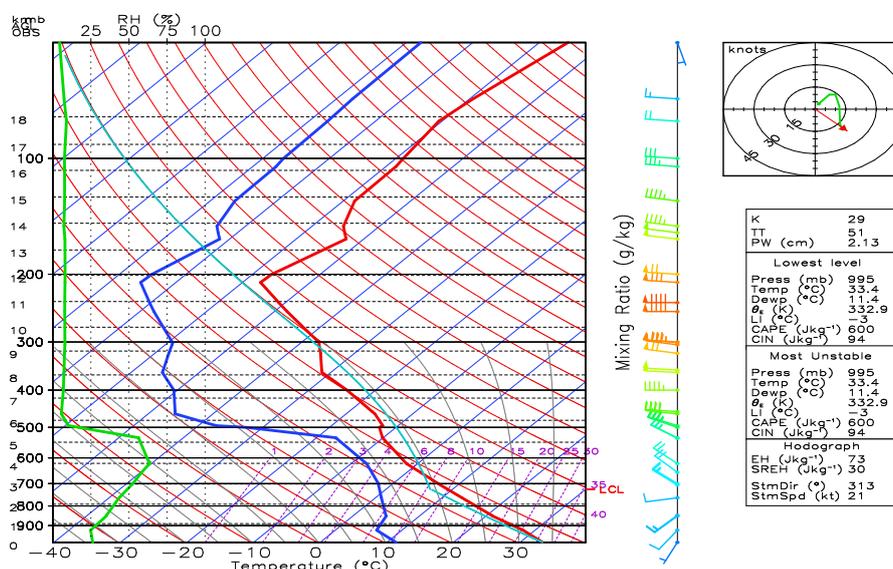
130 The overshooting storm happened in the late afternoon on 19 June 2017 in the  
131 northwestern Beijing city in the mid-latitude when a deep cold trough passed through  
132 the city. The Beijing city was just located in the bottom of the trough with a strong  
133 wind shear. An isolated convection was initially formed in the northeastern mountain  
134 region of Beijing city, and developed as a severe overshooting hailstorm in 30 minutes.  
135 The storm experienced multiple splitting processes, and produced rainfall over 50 mm  
136 and hailstones of about 2.5 cm in diameter. The storm lasted for more than 2 hours.  
137 Radar observations show that the overshooting storm top penetrated the tropopause  
138 ( $\sim 12$  km) and reached up to 16 km in the lower stratosphere, and cloud base was  
139 located around  $20^{\circ}\text{C}$ , indicating that the storm was a severe overshooting storm with a  
140 warm cloud base and favorable for hail formation. The level for zero temperature was  
141 around 3.7 km.

142 Fig.1 is the sounding profiles of temperature, dewpoint temperature and relative  
143 humidity at 20:00 on June 19, 2017 in Beijing meteorological station. It indicates that  
144 the atmospheric layer was relatively dry and the tropopause was located at around 12  
145 km. The Convective Available Potential Energy (CAPE) and Convection Inhibition  
146 (CIN) were 602 J/kg and 94 J/kg, respectively. The wind shear at 0-6 km was 19 m/s  
147 with a southwesterly warm moist advection at the low-level and  
148 northwesterly/westerly cold air advection at the high-level. Therefore, the atmosphere  
149 had a potential unstable condition for convection initiation and development.

150 The hodograph exhibited a clockwise-turning from southwesterly winds near the  
151 surface to northwesterly winds at approximately 5 km, and almost unidirectional  
152 westerly winds at above 5 km. The low-level clockwise-turning hodograph is



153 favorable not only for right-moving storm splitting (Klemp and Wilhelmson, 1978),  
 154 but also for some long-lived, left moving storms (Grasso and Hilgendorf, 2001). The  
 155 Environmental Helicity (EH) and Storm Relative Environmental Helicity (SREH) in  
 156 this study were 73 and 30  $\text{J kg}^{-1}$ , respectively, indicating that there was a relative weak  
 157 “helical” updraft (Johns and Doswell, 1992; Droegemeier et al. 1993).



158  
 159 Fig.1. Profiles of temperature (red), dewpoint temperature (blue) and relative humidity (green) in  
 160 the Skew T-log P diagram at 20:00 on 19 June 2017 at Beijing Meteorological Station (39.8° N,  
 161 116.5° E). The winds, hodograph and environmental conditions are given on the right panel.

162

### 163 3.2 Observed and modeled properties of the storm

164 To understand the properties of the GWs-generated storm, the observed composite  
 165 radar reflectivity and corresponding vertical cross sections for the storm are shown in  
 166 Fig.2. At 19:30, a strong convection had already formed in the northeastern mountain  
 167 region of Beijing (Fig.2a<sub>1</sub>). The corresponding vertical cross section in Fig.2a<sub>2</sub> shows  
 168 that the convection top was located at nearly 14 km and the maximum reflectivity was  
 169 40 dBZ, indicating that the storm had penetrated the tropopause (~12 km) and entered  
 170 the lower stratospheric layer. The storm was considered as a potential hailstorm at  
 171 19:42 since the high reflectivity was forming at the upper levels of the storm.



172 By 19:54, the storm entered the mature stage with the maximum reflectivity more  
173 than 60 dBZ and a pronounced leading stratiform region toward the northeast due to  
174 the influence of the strong southwesterly moist flow at the low- and mid-level  
175 (Fig.2b<sub>1</sub>). Meanwhile, the storm also had an apparent development and extension  
176 toward the southeast. The vertical cross section in x-z in Fig.2b<sub>2</sub> shows that the storm  
177 top was more than 14 km, a large area shows as an overshooting structure although  
178 the storm top is relatively flat. Two high reflectivity cores with more than 50 dBZ  
179 were located at the height of 6-9 km and the reflectivity top with 40 dBZ reached up  
180 to 14 km, indicating that the graupel/hail was forming at the upper levels and the  
181 storm would produce hailfall soon.

182 At 20:06, the high reflectivity in the storm had an apparent development and  
183 expansion toward the east (Fig.2c<sub>1</sub>). The southeastern extension of high reflectivity  
184 was also obvious. The striking phenomenon is that the upper-level reflectivity had an  
185 apparent V-shaped splitting structure, which should be closely related to the high  
186 reflectivity descending process (Fig.2c<sub>2</sub>). Meanwhile, the storm top experienced an  
187 explosive growth for about 2 km from 14 to 16 km and the overshooting structure  
188 became pronounced. The simulated results in the next section will show that the  
189 downward propagating GWs are generated at this stage.

190 At 20:12, the development and extension of high reflectivity toward the east  
191 became more obvious than that toward the southeast (Fig.2d<sub>1</sub>), suggesting that the  
192 storm splitting in the west-east direction was faster than that in the south-north  
193 direction, although the developments toward the both directions were initiated almost  
194 at the same time. The V-shaped reflectivity splitting structure became more obvious  
195 on the eastern flank of the storm due to the further descending of the upper-level high  
196 reflectivity (Fig.2d<sub>2</sub>). Corresponding to the high-reflectivity descending process, the  
197 cloud top was decreased to 14 km.

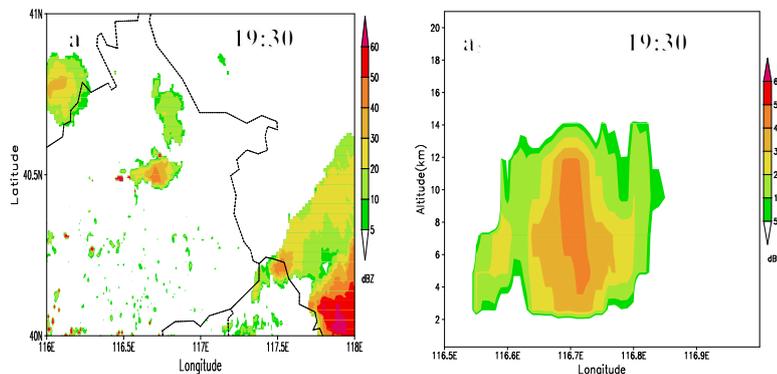
198 By 20:18, the mid- and upper-level high reflectivity in the storm had already split  
199 in the west-east direction (Fig.2e<sub>1</sub>, e<sub>2</sub>), indicating the main updraft of the storm had  
200 split into two independent updrafts. After 20:18, the storm development and  
201 expansion toward the southeast tended to enhance and became more pronounced. The



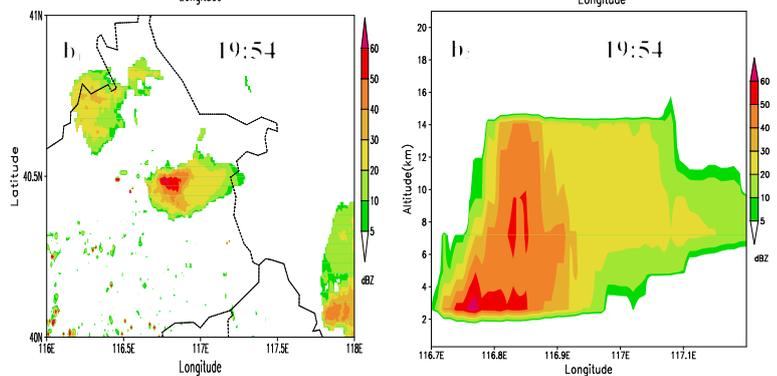
202 splitting in the south-north was quite similar to that in the west-east direction, but the  
203 reflectivity splitting occurred in the central storm. Since the paper mainly focuses on  
204 the properties of the storm and associated generation mechanism of GWs, the splitting  
205 mechanism is out of the scope of this study.

206 As stated above, the observed storm had two pronounced features, one was that the  
207 storm top penetrated the tropopause and entered the lower stratosphere in the mature.  
208 The other was the V-shaped reflectivity splitting structure associated with the  
209 descending of the upper-level high reflectivity, and the accompanied explosive growth  
210 of storm top and the overshooting structure after the mature stage.

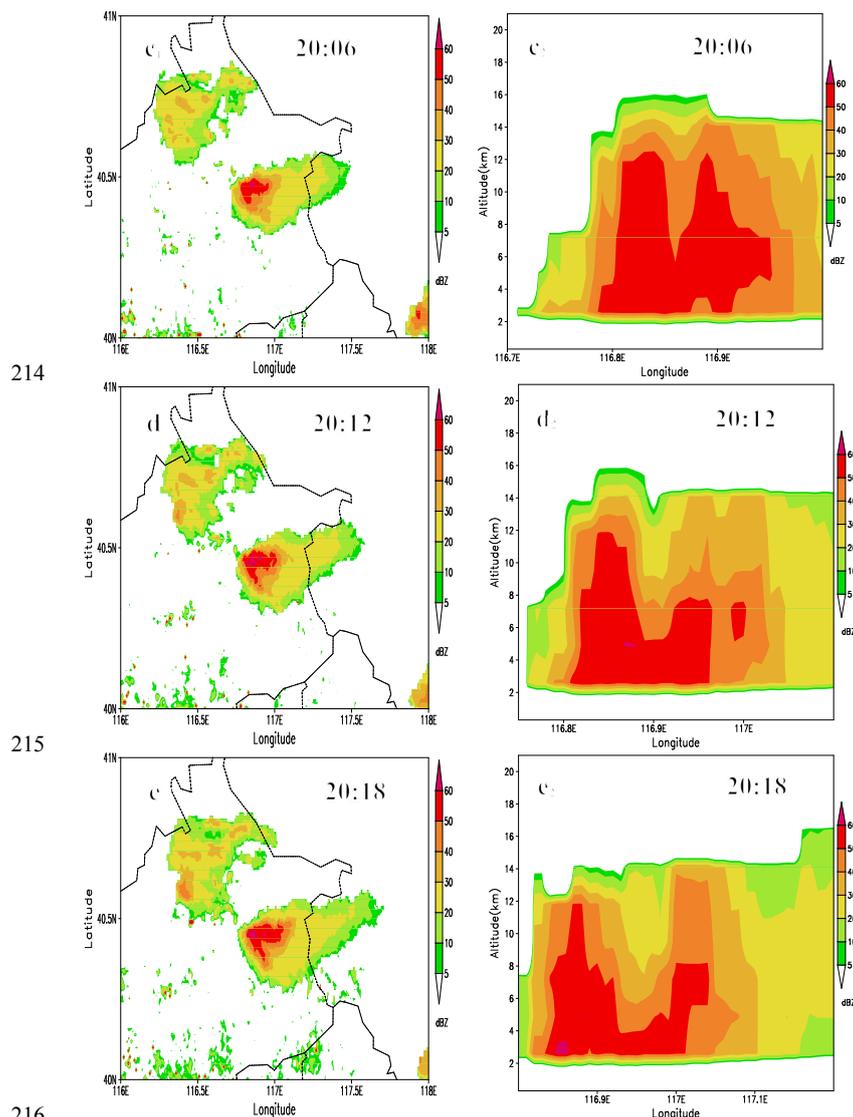
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216  
217 Fig.2. Observed composite radar reflectivity (a<sub>1</sub>-e<sub>1</sub>) and corresponding vertical cross sections  
218 along the high echo cores in the west-east direction (a<sub>2</sub>-e<sub>2</sub>) at (a) 19:30, (b) 19:54, (c) 20:06, (d)  
219 20:12, and (e) 20:18 on 19 June 2017 in Beijing.

220  
221 To compare with the observed storm, the temporal evolution of the simulated  
222 mixing ratio of total hydrometeors for the modeled storm in the x-z (west-east) cross  
223 section is displayed in Fig.3. At 10 min, which is roughly corresponding to the  
224 observed time 20:10 since the model simulation is initiated with the sounding data at  
225 20:00 (the starting time for operational sounding is 19:15). A vigorous convection is



226 formed with the cloud-top height of 10 km, and the maximum mixing ratio of total  
227 hydrometeors reaches more than 15 g/kg (Fig.3a). By 12 min (Fig.3b), the modeled  
228 storm has the cloud-top height of 14 km and the maximum mixing ratio of 20 g/kg,  
229 indicating that the storm enters the mature stage with a very high loading of  
230 hydrometeors (graupel/hail particles, see Fig.4a<sub>1</sub>-a<sub>3</sub>) at the upper levels. The  
231 overshooting storm top penetrates the tropopause (~12 km) and enters the height of 14  
232 km in the lower stratosphere, which is well consistent with radar observations. The  
233 modeled storm has the maximum updraft of about 60 m/s in the mature stage and  
234 downdraft of about -35 m/s.

235 At 14 min (Fig.3c), the upper-level high total hydrometeor mixing ratio  
236 significantly decreases from 20 to 15 g/kg due to the strong descending process of  
237 upper-level graupel/hail on the eastern flank of the storm. Meanwhile, the modeled  
238 storm top increases from 14 to about 16 km. The overshooting storm structure become  
239 more pronounced. All modeled features are well consistent with those observed by  
240 radar (Fig.2b<sub>2</sub>, c<sub>2</sub>).

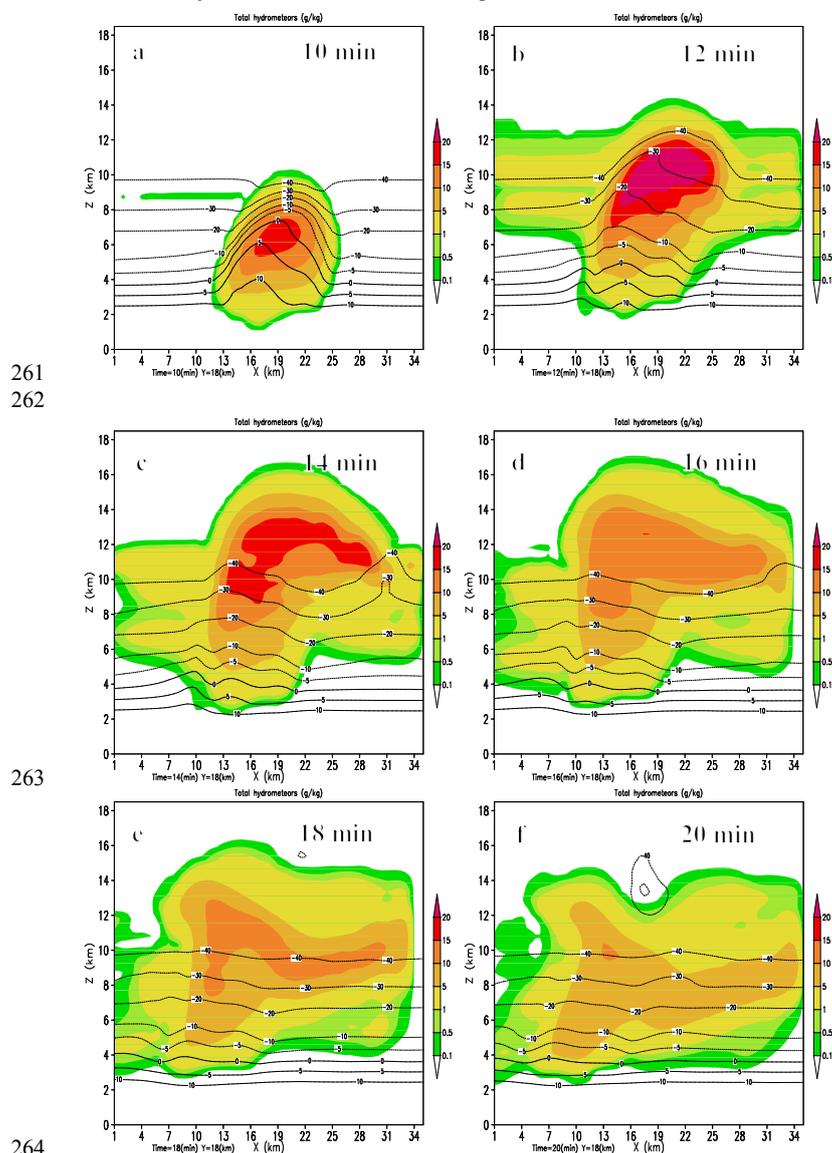
241 By 16 min (Fig.3d), the total hydrometeor mixing ratio decreases from 15 to 10  
242 g/kg and the continuous descending of upper-level graupel/hail further strengthens the  
243 cloud-top height. At 18 min (Fig.3e), the total hydrometeor distribution tends to have  
244 an obvious V-shaped splitting structure, which is also consistent with the V-shaped  
245 reflectivity splitting structure observed by radar, indicating that the V-shaped splitting  
246 structure is closely associated with the descending of precipitating hydrometeors. The  
247 area with the mixing ratio of 10 g/kg decrease significantly, indicating that the  
248 cloud-top height tends to decrease in the region with apparent descending  
249 hydrometeors.

250 By 20 min (Fig.3f), the further descending of precipitating hydrometeors cause the  
251 maximum mixing ratio to decrease from 10 to 5 g/kg in the almost whole storm, and  
252 the V-shaped splitting structure of hydrometeor mixing ratio descends to the lower  
253 levels. The cloud-top decreases to about 14 km.

254 As described above, the properties of the modeled storm and descending processes  
255 of upper-level precipitating hydrometeors are generally consistent with radar



256 observations. It is shown that the storm penetrates the tropopause and enter the lower  
257 stratosphere in the mature stage. The strong descending of precipitating hydrometers  
258 causes more pronounced overshooting structure. The continuous descending of  
259 precipitating hydrometeors can induce an apparent V-shaped splitting structure as that  
260 observed by radar after the mature stage.



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266 Fig.3. Temporal evolution of the simulated mixing ratio of total hydrometeors in the x-z



267 (west-east) cross section for the modeled storm from 10 to 20 min on 19 June, 2017. The  
268 horizontal solid and dashed lines in figures are environmental positive and negative temperatures,  
269 respectively.

270

### 271 **3.3 The generation mechanism and relevant properties for the GWs**

272 To investigate the generation mechanism and relevant properties of the GWs, the  
273 temporal evolution of graupel/hail mixing ratio, pressure perturbation, temperature  
274 perturbation and vertical velocity from 14 to 18 min in the x-z cross section are shown  
275 in Fig.4.

276 As stated above, the overshooting severe storm is formed when the simulated storm  
277 is in the mature stage at 12 min. In this stage, the storm reaches its maximum updraft  
278 of 60 m/s and an equilibrium between the buoyancy force and hydrometeor loading is  
279 established. As long as the updraft cannot further sustain the upper-level high  
280 hydrometeor loading, the high precipitating hydrometeors may break the equilibrium  
281 and descend after the mature stage.

282 At 14 min, the upper levels are dominated by high graupel/hail loading with the  
283 maximum mixing ratio of more than 15 g/kg (Fig.4a<sub>1</sub>), indicating the upper-level high  
284 reflectivity observed by radar and hydrometeors simulated by the model are due to the  
285 graupel/hail particles. The corresponding pressure perturbation distribution in Fig.4b<sub>1</sub>  
286 shows that there is a strong positive pressure perturbation of more than +3 hPa at the  
287 upper levels of 8-15 km on the western flank of the storm, which is related to the  
288 strong updraft and latent heating on the flank (Fig.4d<sub>1</sub>). The high graupel/hail loading  
289 is just located in the region of strong positive pressure perturbations, so that the  
290 fluctuation of the high graupel/hail loading may significantly change the pressure  
291 perturbation. A tilting and relatively uniform positive pressure perturbation with +1  
292 hPa penetrates the middle and lower levels and corresponds an obvious wavelike  
293 temperature perturbation (Fig.4c<sub>1</sub>), indicating that downward propagating GWs have  
294 already occurred at this stage since the cloud-top height has an apparent upward  
295 extension. The upper-level high negative temperature perturbation region is closely  
296 associated with the strong outflow at the cloud top. A small negative temperature



297 perturbation region located just below the high negative temperature region should be  
298 caused by the downward propagating GWs. The positive temperature area located on  
299 the eastern flank is due to the adiabatic warming of downdraft. The propagating GWs  
300 cannot be seen clearly in the pressure perturbation due primarily to that strong  
301 background pressure perturbations offset the effect induced by the GWs. Since there  
302 is no apparent change in pressure perturbation, the vertical velocity in the area is still  
303 dominated by updraft (Fig.4d<sub>1</sub>).

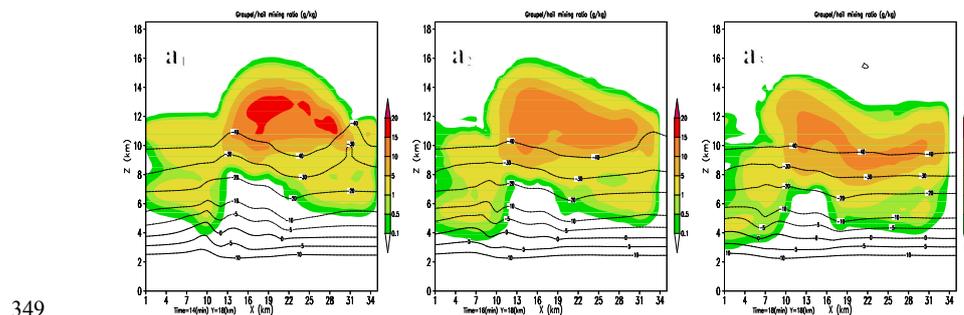
304 By 16 min, the maximum graupel/hail mixing ratio decreases to be lower than 15  
305 g/kg due to the apparent descending of graupel/hail particles at the upper levels on the  
306 eastern flank of the storm (Fig.4a<sub>2</sub>). In response to the significant decrease of the  
307 upper-level graupel/hail loading, the equilibrium between the buoyance force and  
308 hydrometeor loading is destructed and a strong restoring force of buoyancy is  
309 produced in the stratosphere. The formation of the restoring force of buoyancy causes  
310 the overshooting structure to be more prominent. The buoyancy oscillations induced  
311 by the continuous descending of the upper-level graupel/hail in the overshooting  
312 storm induce a pronounced downward propagating GWs, which is can be clearly seen  
313 in pressure perturbation (Fig.3b<sub>2</sub>). In the pressure perturbation distribution, the  
314 positive pressure perturbation is generally not as obvious as the negative pressure  
315 perturbation, this is because that when the downward propagating GWs penetrate the  
316 high-pressure region dominated by the updraft, the positive pressure perturbation  
317 induced by the GWs should be much smaller than that induced by the updraft of the  
318 storm.

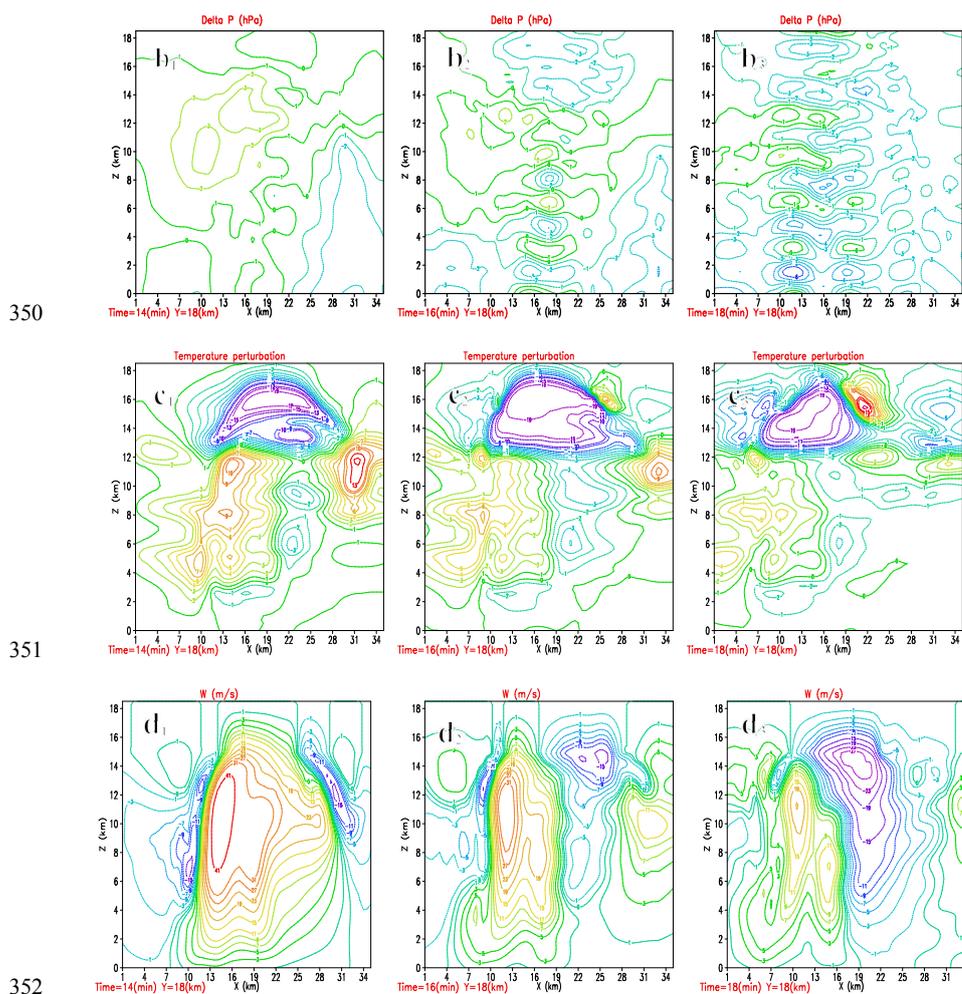
319 The estimated wavelength of the GWs is around 3-4 km. Accompanying with the  
320 strengthening downward propagating GWs, the temperature perturbation is further  
321 enhanced (Fig.4c<sub>2</sub>). As a result, the main updraft in the storm is split into two  
322 independent updrafts as shown in Fig.4d<sub>2</sub>, indicating that the rapid updraft splitting is  
323 closely associated with the downward propagating GWs generated by buoyancy  
324 oscillations induced by the descending of the upper-level high graupel/hail. The weak  
325 downdrafts of -1~-3 m/s are distributed along the path of downward propagating GWs,  
326 indicating the downward momentum transport associated with the downward



327 propagating GWs can damage the updraft of the storm and induce the storm splitting  
328 rapidly. The strong compensating subsidence with the magnitude of -15 m/s in the  
329 stratosphere should be primarily induced by the downward propagating GWs  
330 (Bretherton and Smolarkiewicz, 1989), and the descending of graupel/hail particles.

331 At 18 min, the continuous descending of graupel/hail on the eastern flank of the  
332 storm significantly decreases the upper-level graupel/hail loading (Fig.4a<sub>3</sub>).  
333 Meanwhile, the descending of graupel/hail particles tends to shift toward the west and  
334 induce a westward shifting of downward propagating GWs (Fig.4b<sub>3</sub>). An interesting  
335 phenomenon at this stage is that the GWs also have an apparent upward propagation  
336 to the lower stratospheric layer. Since the downward propagating GWs are generated  
337 through buoyancy oscillations induced by the descending of the upper-level  
338 graupel/hail in the stratospheric layer as shown in Fig.4b<sub>2</sub>, the upward propagating  
339 GWs should be induced by the surface-reflected GWs when the downward  
340 propagating GWs reach to the surface as proposed by Kim et al. (2012). It will be  
341 seen in the following section that the upward propagating GWs can generate a strong  
342 horizontally propagating GWs in the lower stratospheric layer. In fact, the temperature  
343 perturbation has already shown an obvious horizontal wavelike distribution in the  
344 layers above the tropopause (~12 km) (Fig.3c<sub>3</sub>). The temperature perturbation below  
345 the tropopause tends to weaken due to the effect of the upward propagating GWs.  
346 With the descending of the upper-level graupel/hail and the formation of strong  
347 downdraft, the storm tends to weaken significantly (Fig.4d<sub>3</sub>). Note that the vertical  
348 velocity in the stratospheric layer also tends to have a wavelike perturbation.





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353 Fig.4. Temporal evolution of (a<sub>1</sub>-a<sub>3</sub>) graupel/hail mixing ratio (g/kg), (b<sub>1</sub>-b<sub>3</sub>) pressure perturbation  
354 (hPa), (c<sub>1</sub>-c<sub>3</sub>) temperature perturbation, and (d<sub>1</sub>-d<sub>3</sub>) vertical velocity (m/s) in x-z cross section at  
355 y=18 km from 14 to 18 min. The horizontal solid and dashed lines in (a) are environmental  
356 positive and negative temperatures, respectively. The updraft is solid lines and downdraft is  
357 dashed lines in (c).

358

359 As described above, the storm penetrates the tropopause and enters the lower  
360 stratospheric layer and forms an overshooting severe storm in the mature stage. The  
361 descending of the upper-level high graupel/hail loading in the overshooting storm  
362 breaks the equilibrium between the buoyancy force and hydrometeor loading



363 established in the mature and induce a strong restoring force of buoyancy. The  
364 continuous descending processes produce buoyancy oscillations that excite the  
365 downward propagating GWs. The downward propagating GWs can produce apparent  
366 downward perturbations in pressure, temperature and vertical velocity, resulting in a  
367 rapid storm splitting. The initially induced downdraft by the GWs is around -1~ -3  
368 m/s. The GWs generated by the overshooting severe storm have a duration of about  
369 20 minutes and an estimated wavelength of about 3-4 km. Although the generation  
370 mechanism of the GWs in this study is different from those proposed by previous  
371 studies, the relevant features of the GWs are generally similar to convectively  
372 generated GWs through other mechanisms (e.g., Larsen et al.1982; Nolan and Zhang,  
373 2017; Jewtoukoff et al. 2013). Larsen et al. (1982) observed GWs generated by  
374 afternoon thunderstorms with a vertically-pointing 430 MHz radar and found that  
375 when the cloud-top height reached the tropopause, gravity-wave oscillations in the  
376 vertical velocity above the tropopause would develop, with an amplitude of 2 m/s, and  
377 period of close to 6 min. The aircraft measurements by Nolan and Zhang (2017)  
378 indicated that the GWs have radial wavelengths of 2–10 km and vertical velocity  
379 magnitudes from 0.1 to 1.0 m/s. Jewtoukoff et al. (2013) reported the GWs near a  
380 tropical cyclone with wavelengths of around 1 km observed by a balloon at 19 km  
381 altitude. In addition, the upward propagating GWs induced by the surface reflection  
382 are also obvious in the simulation and will be further discussed in the next section.

383

### 384 **3.4 The influences of the surface-reflected GWs on the stratosphere**

385 As shown above, both the downward and upward propagating GWs are formed  
386 through buoyancy oscillations and surface reflections in the continental overshooting  
387 hailstorm. It is shown that the downward momentum transport associated with the  
388 downward propagating GWs can induce the rapid updraft splitting and change the  
389 storm morphology and evolution. One of important issues is that whether the  
390 surface-reflected upward propagating GWs can also affect the stratospheric  
391 atmosphere through the upward momentum and energy transport as proposed by  
392 previous studies (e.g., Alexander and Holton, 1997; Piani et al., 2000; Baldwin et al.,



393 2001; Beres et al., 2002; Fritts and Alexander,2003). To investigate this issue, the  
394 subsequent evolution of cloud total hydrometeor, pressure and temperature  
395 perturbations, and vertical velocity from 20 to 50 min for the simulated storm is  
396 displayed in Fig. 5.

397 At 20 min, the upper-level mixing ratio of graupel/hail decreases to be less than  
398 10 g/kg (Fig.5a<sub>1</sub>). The Fig.5b<sub>1</sub> shows that the wavelike pressure perturbation induced  
399 by the downward propagating GWs continues to shift toward the west due to the  
400 westward shifting of descending process of graupel/hail. The pressure perturbation  
401 tends to weaken due to the weakening of graupel/hail loading. It can be clearly seen  
402 that the surface-reflected upward propagating GWs induce an obvious positive  
403 pressure perturbation in the lower stratosphere (Fig.5b<sub>1</sub>). This phenomenon is more  
404 prominent in the temperature perturbation (Fig.5c<sub>1</sub>). In the lower stratosphere, the  
405 pronounced wavelike fluctuations in temperature perturbation can be clearly seen,  
406 indicating that the momentum and energy transport associated with the upward  
407 propagating GWs can enter the lower stratospheric layer and generate strong  
408 horizontally propagating GWs as observed by aircraft (Nolan and Zhang (2017)).

409 It should be noted here that the temperature distribution pattern with a warm  
410 center surrounded a U-shaped or V-shaped cold region in the lower stratospheric layer  
411 over the storm top is quite similar to those found in the pyroCb<sub>s</sub> (Luderer et al.,2007)  
412 and intense thunderstorms (Wang et al., 2002) induced by GWs. However, the GWs in  
413 the lower stratosphere in this study are generated by the surface reflection of  
414 downward propagating GWs rather than that directly produced on the storm top. The  
415 formation of the wavelike distribution in vertical velocity can be also seen in the  
416 lower stratosphere, although the vertical velocity distribution is still dominated by  
417 main updraft and the compensating subsidence.

418 By 30 min, the graupel/hail descends to the lower levels and some of them melt as  
419 rainwater, so that high mixing ratio of hydrometeor presents at the surface (Fig.5a<sub>2</sub>).  
420 The cooling caused by both the melting and evaporating processes causes the pressure  
421 at the near-surface to decrease significantly (Fig.5b<sub>2</sub>). The small surface positive  
422 pressure should be related to the cold downdraft. The surface-reflected upward



423 propagating GWs induce a new temperature fluctuation at the lower levels of the  
424 stratosphere (Fig.5c<sub>2</sub>). A strong cold pool with the thickness of 4 km is formed at the  
425 near-surface layer with the minimum temperature of -15 °C. The downdraft is  
426 dominated in the cold pool (Fig.5d<sub>2</sub>). The surface-reflected upward propagating GWs  
427 also induce apparent fluctuations in vertical velocity in the stratosphere.

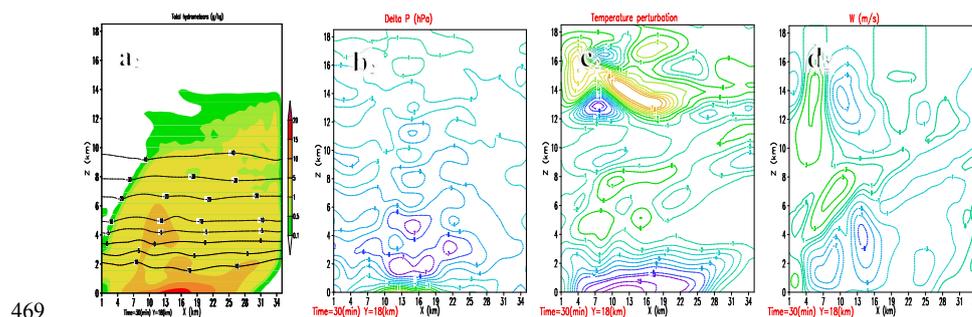
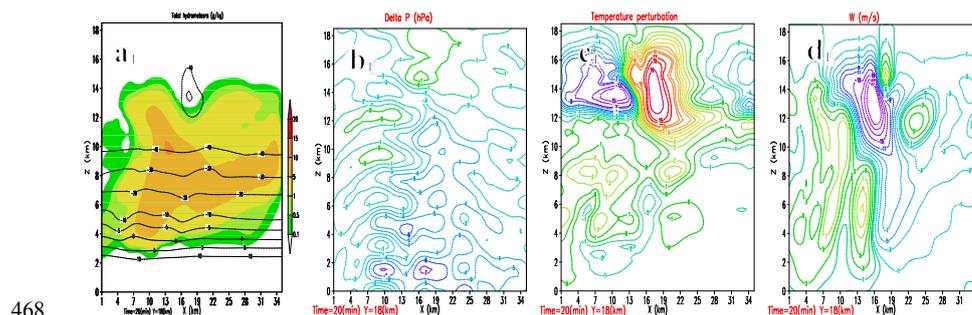
428 At 40 min, the cloud-top descends to the height below the tropopause and the  
429 graupel/hail descending process has weakened significantly (Fig.5a<sub>3</sub>). Since the  
430 restoring force and buoyancy oscillations cannot be formed in the troposphere, the  
431 downward propagating GWs cannot be also generated (Fig.5b<sub>3</sub>), instead, the low-level  
432 evaporative cooling produces a strong negative pressure perturbation in the  
433 near-surface layers. The strong cold pool spreading cause the surrounding air to lift  
434 and condense (Fig.5c<sub>3</sub>). Meanwhile, it seems that the strong cold pool spreading also  
435 generates the weak upward propagating GWs in the stable low layers, and induces  
436 relatively weak horizontal temperature fluctuations in the low stratosphere. The  
437 vertical velocity distribution in Fig.5d<sub>3</sub> shows that within the cold pool there is  
438 downdraft while above the cold pool there is a weak updraft due to the lifting of  
439 spreading outflow. In the stratosphere, there are horizontally propagating weak  
440 fluctuations in vertical velocity with an amplitude of around 1-3 m/s. Therefore, the  
441 strong cold pool spreading at the low levels could also generate upward propagating  
442 GWs and induce the momentum and energy transport from the low tropospheric  
443 levels to the upper stratospheric layers.

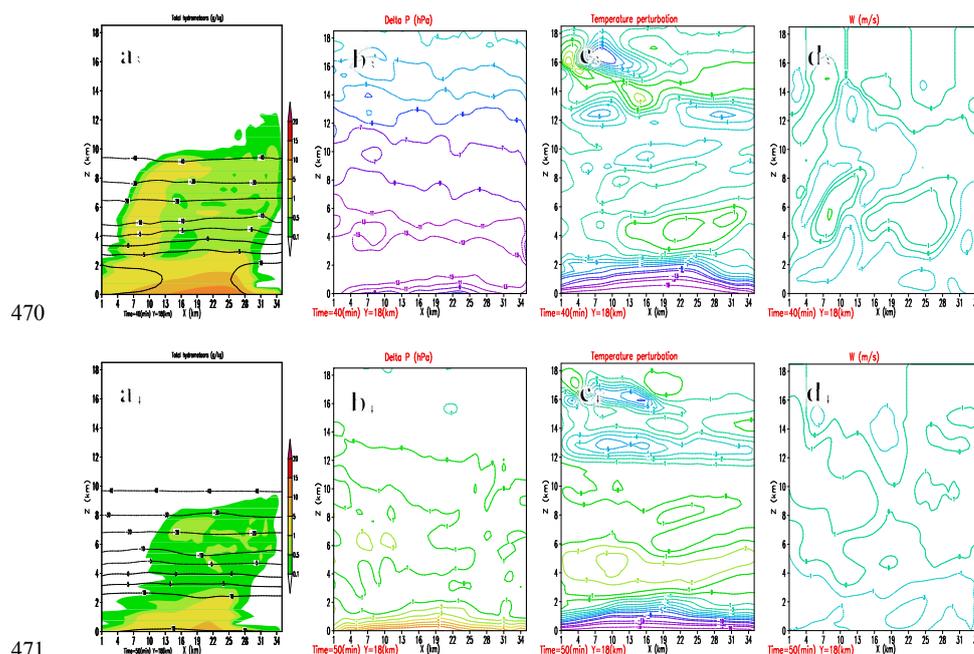
444 By 50 min, the convective cloud has evolved as a stratiform cloud (Fig.5a<sub>4</sub>). Since  
445 precipitation and the associated melting and evaporative cooling weaken significantly  
446 and the downward cold airflow become dominant (Fig.5d<sub>4</sub>), the near-surface layer is  
447 dominated by positive pressure perturbation (Fig.5b<sub>4</sub>). The spreading outflow induced  
448 by the cold pool continuously lifts and condensates the air above it and form a weak  
449 positive temperature perturbation (Fig.5c<sub>4</sub>). The vertical velocity distribution in  
450 Fig.5d<sub>4</sub> shows that there is a downdraft in the cold pool and weak uplifting velocity  
451 above the cold pool. The horizontal GWs are no longer to propagate in the  
452 stratosphere. Therefore, comparing with the GWs generated by buoyancy oscillations



453 induced by the descending of the upper-level high graupel/hail, the GWs excited by  
454 the strong spreading outflow of the cold pool are relatively weak and have less impact  
455 on the stratosphere. However, the cold pool has an important role in maintaining the  
456 subsequent clouds and precipitation through the lifting process.

457 Therefore, the surface-reflected upward propagating GWs have significant impacts  
458 on the temperature and vertical velocity distributions in the stratosphere, indicating  
459 that the GWs generated by the overshooting severe hailstorm not only influence the  
460 storm morphology and evolution through downward propagating process, but also  
461 significantly affect the stratospheric atmosphere through the surface-reflected upward  
462 propagating process. The GWs excited by the strong spreading outflow of the cold  
463 pool are relatively weak. But the spreading outflow of the cold pool has an important  
464 role in maintaining the subsequent development of clouds and precipitation through  
465 the lifting process. When the environmental air is under the unstable condition, the  
466 lifting could induce convection and a longer duration of the storm. This property is  
467 apparent in the splitting in the south-north direction (not shown).





470  
471  
472 Fig.5. Vertical cross sections of (a<sub>1</sub>-a<sub>4</sub>) graupel/hail mixing ratio (g/kg), (b<sub>1</sub>-b<sub>4</sub>) pressure  
473 perturbation (hPa), (c<sub>1</sub>-c<sub>4</sub>) temperature perturbation, and (d<sub>1</sub>-d<sub>4</sub>) vertical velocity (m/s) in x-z at  
474 y=18 km from 20 to 50 min. The horizontal solid and dashed lines in (a) are environmental  
475 positive and negative temperatures, respectively. The cloud boundary is superimposed as a thick  
476 solid curve for 0.1 g/kg in (b). The updraft downdraft is solid lines and is dashed lines in (d).

477  
478 **4 Conclusions and discussion**

479 The GWs generated by a continental overshooting hailstorm occurred on 19 June  
480 2017 in Beijing in the mid-latitude are first reported in this study based on radar  
481 observations and modeled results. The main conclusions are summarized as follows.

482 The GWs-generated overshooting hailstorm has the maximum cloud-top height of  
483 over 16 km, updraft of 60 m/s and graupel/hail mixing ratio of over 20 g/kg in the  
484 mature stage. The storm penetrates the tropopause and enters the lower stratosphere  
485 and forms a typical overshooting storm. After the mature stage, the descending of the  
486 upper-level high graupel/hail loading causes the breaking of equilibrium between the  
487 buoyancy force and hydrometeor loading established in the mature stage, and induce a



488 strong restoring force of buoyancy. The continuous descending processes of the  
489 upper-level high graupel/hail loading produces buoyancy oscillations that excite  
490 downward propagating GWs. The GWs have the estimated wavelength of about 3-4  
491 km and duration of about 20 min.

492 The momentum flux associated with the downward propagating GWs produces  
493 downdraft, and cause the main updraft splitting quickly in the storm, and significantly  
494 change the storm structure and evolution. The downdraft magnitude induced by the  
495 GWs is about -1~3 m/s in the initial stage. The upward propagating GWs can be also  
496 formed through the surface reflection of the downward propagating GWs. The upward  
497 propagating GWs are trapped in the lower stratosphere and induce the large  
498 fluctuations in temperature and vertical velocity, causing significant influences on the  
499 dynamic and thermodynamic structure in the low stratosphere.

500 The generation mechanism of the GWs reported in this study is different from the  
501 convectively generated GWs mechanisms through mechanical, thermal and mountain  
502 forcing proposed by previous studies, since the convectively generated GWs through  
503 mechanical and thermal forcing mechanisms are closely associated with latent heating  
504 release and updraft fluctuation, and generally propagate upward with a restoring force  
505 of gravity, so that the GWs have significant contributions to the stratospheric  
506 momentum and energy budget (Fritts and Alexander, 2003), while the GWs reported  
507 in this study are excited by buoyancy oscillations caused by the continuous  
508 descending processes of graupel/hail in an overshooting hailstorm. The restoring force  
509 is buoyancy. The GWs propagate downward and have important impacts on the storm  
510 morphology and evolution, as well as lower stratosphere through the surface  
511 reflection process.

512 The properties of the GWs generated by the overshooting hailstorm in this study  
513 are generally consistent with radar and aircraft observations (Larsen et al.,1982;  
514 Jewtoukoff et al., 2013; Nolan and Zhang, 2017). The temperature distribution pattern  
515 with a warm center surrounded a U-shaped or V-shaped cold region in the lower  
516 stratospheric layer over the storm top is quite similar to that found in the pyroCb  
517 (Luderer et al.,2007) and intense thunderstorms (Wang et al., 2002). Luderer et



518 al.(2007) proposed that small-scale mixing processes are strongly enhanced by the  
519 formation and breaking of a stationary gravity wave induced by the overshoot.  
520 However, the GWs in the lower stratosphere in this study are generated by the surface  
521 reflection process of downward propagating GWs rather than that directly induced by  
522 the storm overshoot. In addition, it should be noted that the overshooting hailstorm in  
523 this study has the maximum updraft of 60 m/s, cloud-top height up to 16 km and  
524 graupel/hail mixing ratio up to 20 g/kg in the mature stage, so that the strong  
525 downward propagating GWs can be generated through buoyancy oscillations induced  
526 by the continuous descending processes of the upper-level high graupel/hail loading.  
527 Whether the GWs can be generated in general continental severe hailstorms through  
528 these processes remain uncertain and needs further study in the future.

529

530 Data availability. Radar and sounding data used in this study are available from the National  
531 Meteorological Information Center (NMIC), China Meteorological Administration (CMA),  
532 website: <http://www.nmic.cn/>.

533

534 Author contributions. XLG conceptualized and designed the study.  
535 XG, DF and XLG performed data analysis and numerical simulations.  
536 XG and DF conducted formal analysis and wrote the manuscript.  
537 All authors read and approved the final manuscript.

538

539 Competing interests. The authors declare that they have no conflict of interest.

540

541 *Acknowledgements.* This work was supported by the National Natural Science Foundation of  
542 China (42105173) and the Second Tibetan Plateau Comprehensive Scientific Expedition  
543 (2019QZKK0104).

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