



Nonlinear resonant interactions of atmospheric tides with annual oscillation based on meteor radar observation and reanalysis data

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Abstract. Nonlinear interactions among gravity waves and among tides and planetary waves (PWs) have extensively been studied, however, resonant interactions between tides and annual (AO) and semiannual (SAO) oscillations were not reported. By using meteor radar observations and reanalysis data for 9 years, we demonstrate that the sum and difference resonant interactions between the diurnal (DT)/semidiurnal (SDT) tides and the AO/SAO do occur in the mesosphere and lower thermosphere (MLT). Both the frequencies and wavenumbers of the secondary waves in the sum (difference) resonant interactions just equal the sum (difference) frequencies and wavenumbers between the DT/SDT and the AO/SAO. Spectral analysis shows that only the DT, SDT, AO, SAO and their secondary waves are the predominant components in both the zonal and meridional winds at 90 km with the spectral amplitudes of 3.5-17.7 ms⁻¹, being much stronger than all the other spectral amplitudes, including the amplitudes (2.1-2.2 ms⁻¹) of the relatively strong terdiurnal tide and 16-day PW. At some altitudes in the MLT, the secondary waves are more intense than the DT/SDT, thus in tidal studies, the magnitude of the secondary waves may be regarded as the tidal one if the observational period is not long enough or their spectral peaks are not distinguished carefully.

31

32 1 Introduction

There are various motions with different temporal and spatial scales in the atmosphere, such as small-scale gravity waves (GWs), large-scale tidal waves and planetary waves (PWs), as well as longer period intraseasonal (ISO), semiannual (SAO) and annual (AO) oscillations. These perturbations play an important role in the circulation and thermal structure of the middle and upper atmosphere due to their ability to transfer energy and momentum among different atmospheric layers and among different latitudinal zones (Andrews et al., 1987). Meanwhile, there are complex coupling and nonlinear interactions among different scale waves and oscillations. For example, GWs and tides modulated by the filtering of PWs or ISOs in the lower atmosphere may drive similar period but out of phase PWs or ISOs in the mesosphere and



lower thermosphere (MLT) as momentum transported by GWs and tides is deposited into the background flow in the MLT due to their dissipation (Eckermann et al., 1997; Jacobi et al., 2006; Cheng et al., 2021). This process is essentially wave-flow coupling through the filtering effect as long as long-period PWs or ISOs are regarded as the temporal varying background field through which short-period GWs and tides propagate. Wave-wave nonlinear interaction is another dynamic process in which new spectral components are excited. In theory, nonlinear interaction comes from the nonlinear advective terms in the dynamic equations. Numerical simulation by Teitelbaum and Vial (1991) clearly showed that two secondary waves with sum and difference frequencies between semidiurnal tide (SDT) and 16-day PW were generated only through nonlinear interactions of SDT and 16-day wave rather than their linear superposition.

Nonlinear resonant interaction in the fluid dynamical context was proposed by Phillips [1960]. For specified wave vectors \vec{k} and frequencies ω , two initial waves with phases $\vec{k}_1 \cdot \vec{r} - \omega_1 t$ and $\vec{k}_2 \cdot \vec{r} - \omega_2 t$ can force a third wave with their sum or difference phase. In other words, if both the wave vector and frequency of a free third mode satisfy the sum or difference matching conditions,

$$\vec{k}_1 \pm \vec{k}_2 = \vec{k}_3 \quad (1)$$

$$\omega_1 \pm \omega_2 = \omega_3 \quad (2)$$

where the subscripts 1, 2 and 3 denote the three interacting waves, resonant interaction occurs and leads to significant energy transfer from the initial waves into the third wave. In the light of the two matching conditions, resonant interactions of small-scale GWs have been extensively investigated (e.g., Yeh and Liu, 1981; Müller et al., 1986; Inhester, 1987; Klostermeyer, 1991; Stenflo, 1998; Yi, 1999; Wüst and Bittner, 2006; Huang et al., 2009; Taklo and Choi, 2020), even in a sheared, dissipative and rotating atmosphere (Fritts et al., 1992; Vanneste, 1995; Axelsson et al., 1996; Huang et al., 2014; Achatz et al., 2017; Chatterjee and Misra, 2021; Voelker et al., 2021), and in the uniform and nonuniform plasmas (Stenflo, 1994; Stenflo and Shukla, 2009). For large-scale tides and PWs, many scientists studied their resonant interactions based



on general circulation models (GCMs) and ground-based and satellite-borne observations by using the matching condition ($n_1 \pm n_2 = n_3$) of zonal wavenumbers n instead of the wave vector matching condition (Kamalabadi et al., 1997; Jacobi et al., 1998, 2001; Pancheva et al., 2000, 2002; Pancheva and Mitchell, 2004; Liu et al., 2007; Palo et al., 2007; Kumar et al., 2008; McCormack et al., 2010; Babu et al., 2011; Chang et al., 2011a, 2011b; Yue et al., 2012; Huang et al., 2013a, 2013b; He and Chau, 2019).

Although nonlinear interactions between atmospheric waves have attracted wide attention, interactions of both global-scale tides and annual and semiannual oscillations were not reported yet. As we know, tides have a wave structure, but AO and SAO do not, thus it remains unclear whether they can interact with each other in the realistic atmosphere. Atmospheric tides are the persistent global perturbations and are generated mainly by radiation heating from the Sun (Lindzen and Chapman, 1969; Forbes and Garrett, 1979), in consequence, diurnal tide (DT) has a period of a solar day, and a westward wavenumber $n = -1$ due to the westward migration of the Sun in the sky from the view of the Earth. Besides DT, there are its high-order harmonics in the atmosphere, among which SDT with a period of 12 h and a wavenumber of $n = -2$ is the strongest component (Chapman and Lindzen, 1970; Hagan et al., 1999). Atmospheric AO and SAO are slow periodic changes of the zonal and meridional winds and temperature without phase propagation in the zonal direction. The periods of AO and SAO are 2-3 orders of magnitude larger than those of tides, thus the identification of sum and difference frequencies between tides and AO/SAO from tides is limited due to the requirement of rather long-time observations with high temporal resolution.

In the MLT, atmospheric waves have generally strong amplitudes, thereby, their nonlinear interactions are an important aspect of dynamics in the MLT. In present study, we use the observational and reanalysis data for 9 years to reveal the occurrence of resonant interactions between tides and AO/SAO. In Section 2, the radar observations and reanalysis data are described in brief. In Section 3, nonlinear interactions between DT/SDT and AO/SAO are discussed based on the observational and reanalysis data. A summary is provided



87 in Section 4.

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89 2 Data

90 In this study, we use the meteor radar observation to investigate the nonlinear interactions between the
91 tides and atmospheric oscillations in the MLT. The meteor radar is situated at Wuhan (30.5°N, 114.6°E) in
92 China, established by the Institute of Geology and Geophysics, Chinese Academy of Sciences (IGGCAS).
93 This system is an all-sky interferometric meteor radar (Hocking et al., 2001), operating at 38.5 MHz. Meteor
94 trails are illuminated by a three-element Yagi antenna directed toward the zenith with a transmitted peak
95 power of 7.5 kW and a pulse repetition frequency of 1.98 kHz (Xiong et al., 2004; Jiang et al., 2005; Yu et
96 al., 2013). Echoes reflected by meteor trails are detected by five three-element Yagi antennas oriented along
97 two orthogonal baselines. The receiving antennas are sampled every 13.3 μ s, resulting in a height resolution
98 of 2 km. The hourly horizontal wind profile in the range of 70 to 110 km is deduced from the meteor trail
99 velocities observed by radar, with a typical measurement uncertainty of 2–4 ms^{-1} around 88–90 km (Hocking
100 et al., 2001). The zonal and meridional winds for 9 years from 1 January 2012 to 31 December 2020 are
101 applied to this investigation.

102 The number of meteor counts detected by radar showed a strong dependence on height, with an
103 approximately Gaussian distribution centered at about 88–90 km, thus the data availability is high around
104 88–90 km. As the height goes up and down, the number of detected meteors decreases, and then the
105 acceptable data reduces gradually. In the period of our concern, including several gaps with the maximum
106 missing data of 53 days, the data availability exceeds 80% in the height range of 80–96 km, and then
107 decreases to 69% at 98 km and 67% at 78 km, but is less than 50% above 98 km and below 78 km. In this
108 case, the radar observation in the altitude range of 78–98 km is chosen for our study.

109 Considering that the zonal wavenumber of planetary scale perturbations cannot be deduced from the



110 observations at a single station, we derive the zonal wavenumber from a two-dimensional Fourier transform
 111 with the help of Modern-Era Retrospective analysis for Research and Applications version 2 (MERRA2)
 112 reanalysis data. The upper pressure level at 0.01 hPa in the MERRA2 corresponds to the altitudes of about
 113 78-80 km, which are largely consistent with the lower limit of the radar observation utilized. The MERRA2
 114 data is produced by the Goddard Earth Sciences Data and Information Services Center (GES DISC) of the
 115 National Aeronautics and Space Administration (NASA) with some updates relative to the previous version
 116 (Gelaro et al., 2017), accessible from the online website at <https://gmao.gsfc.nasa.gov/reanalysis/MERRA-2/>.
 117 The reanalysis data is 6-hourly instantaneous analysis fields on a $0.5^\circ \times 0.625^\circ$ latitude-by-longitude grid at
 118 72 pressure levels from the ground up to 0.01 hPa level (Gelaro et al. 2017). The reanalysis zonal and
 119 meridional winds for a same period from 1 January 2012 to 31 December 2020 are used in this study.

120

121 3. Results

122 3.1 Resonant Interactions between DT/SDT and AO/SAO

123 Figure 1 shows the daily mean zonal and meridional winds from 78 to 98 km observed by the meteor
 124 radar at Wuhan for 9 years from 1 January 2012 to 31 December 2020. The different scale perturbations can
 125 be seen in the zonal and meridional wind fields.

126 Lomb-Scargle spectrum analysis is appropriate for decomposing unevenly sampled data into
 127 perturbation components (Scargle, 1982). In order to obtain the dominant periods in the wind fields, a
 128 Lomb-Scargle spectrum analysis, with a 4-times oversampling, is carried out on the zonal and meridional
 129 winds from the meteor radar observation between 1 January 2012 and 31 December 2020. Figures 2a and 2b
 130 present the Lomb-Scargle spectra of the zonal and meridional winds at 90 km, respectively. One can
 131 intuitively see from Figures 2a and 2b that the DT, SDT and low frequency oscillations are the prominent
 132 components, and the terdiurnal tide is also obvious. At Wuhan, the DT has the spectral amplitude of 17.7



133 and 16.9 ms^{-1} in the zonal and meridional winds, which are stronger than those of 9.6 and 12.1 ms^{-1} for the
 134 SDT, respectively. The terdiurnal tide has the magnitude of about 2.2 ms^{-1} in both the zonal and meridional
 135 winds.

136 For the sake of clarity, we zoom in on the several dominant spectral components. Figures 2c and 2d
 137 depict the low-frequency part (period >50 days) of the Lomb-Scargle spectra of the zonal and meridional
 138 winds, respectively, and then the strong AO and SAO can be clearly identified from the low-frequency
 139 components. In the zonal wind, the amplitude (12.3 ms^{-1}) of the AO is more intense than that (6.2 ms^{-1}) of
 140 the SAO, while in the meridional wind, the magnitude (4.6 ms^{-1}) of the SAO is larger than that (3.5 ms^{-1}) of
 141 the AO. Besides the AO, SAO and tides mentioned above, only the 16-day planetary wave (not presented) in
 142 the zonal wind has the spectral magnitude of about 2.1 ms^{-1} , and all the other spectral amplitudes in the
 143 zonal and meridional winds are smaller than 2 ms^{-1} . Hence, from the perspective of 9-year mean results, the
 144 diurnal and semidiurnal tides and annual and semiannual oscillations are the predominant perturbations in
 145 the MLT at 30.5°N .

146 Figures 2e and 2f show the spectra centered at the DT in the zonal and meridional winds, respectively. It
 147 is interesting that the DT spectra in Figures 2a and 2b are divided into multiple spectrum peaks in the
 148 zoomed-in Figures 2e and 2f. We note that these spectral peaks around the DT are right at the sum and
 149 difference frequencies between the DT and the annual and semiannual oscillations. Since the unexpected
 150 components satisfy the sum and different resonant conditions of $\omega_1 \pm \omega_2 = \omega_3$ with the DT and the AO/SAO,
 151 they should be excited through the nonlinear resonant interactions between the DT and the AO/SAO. In the
 152 zonal wind, the secondary waves excited by the tidal difference (sum) resonant interactions with the AO and
 153 SAO have the magnitudes of 10.6 and 6.8 ms^{-1} (6.5 and 8.4 ms^{-1}), respectively. In the meridional wind, the
 154 excited waves in the difference (sum) resonant interactions between the DT and the AO and SAO have the
 155 intensities of 16.9 and 8.6 ms^{-1} (9.9 and 8.5 ms^{-1}), respectively. Hence, these excited waves are much



stronger than the terdiurnal tide with the amplitude of 2.2 ms^{-1} in both the zonal and meridional winds. Moreover, it is surprising that in the tidal difference resonant interaction with the AO, the spectral amplitude (16.9 ms^{-1}) of the secondary wave in the meridional wind is larger than that (14.4 ms^{-1}) of the DT. Therefore, the spectral peak at the DT in Figure 2b is actually the peak of the generated secondary wave rather than the DT. Similarly, Figures 2g and 2h illustrate the occurrence of the resonant interactions between the SDT and the AO and SAO.

Now that the wavenumber matching conditions of $n_1 \pm n_2 = n_3$ should be satisfied in the resonant interactions, we examine the zonal wavenumber relation with the aid of the MERRA2 reanalysis data. The reanalysis zonal and meridional winds at 0.01 hPa ($\sim 80 \text{ km}$) at 30.5°N for the same period of 9 years are selected to make a two-dimensional Fourier transform. Figure 3 shows the zoomed-in frequency-wavenumber spectra around the AO and SAO and the DT. In Figure 3, we use the sign δ to denote the frequency of the AO ($1/365 \text{ day}^{-1}$), and the negative wavenumber represents the westward propagation of wave. One can clearly see from Figure 3 that the zonal wavenumbers in both the zonal and meridional winds are 0 for the AO and SAO and -1 for the DT, as we all know. The spectral peaks on both sides of the DT are located not only at the sum and difference frequencies between the DT and the AO/SAO, similar to the results of the radar observation in Figure 2, but also just at the wavenumber of -1, indicating that the wavenumber matching conditions of $n_1 \pm n_2 = n_3$ are satisfied. Hence, this confirms that these new components are generated through the nonlinear resonant interactions between the DT and the AO/SAO. In addition, the similar phenomenon arises in the zoomed-in frequency-wavenumber spectra around the SDT with the zonal wavenumber of -2 (not presented), too.

3.2 Spectral Amplitude Variation with Altitude

Here, we investigate the magnitude variation of these generated waves with height. From the Lomb-Scargle spectra of the zonal and meridional winds for the period of 9 years observed by the meteor



179 radar, we extract the spectral amplitudes of the DT/SDT, AO/SAO, and secondary waves at altitude range of
 180 78-98 km. Figure 4 shows the amplitude evolution of the four components in the interactions of the DT with
 181 the AO (upper panels) and the SAO (lower panels). In the zonal wind, the amplitude of the AO decreases
 182 sharply from 21.7 ms^{-1} at 78 km to only 1.9 ms^{-1} at 86 km, and then increases rapidly to 17.0 ms^{-1} at 98 km,
 183 while the SAO diminishes gradually from 14.1 ms^{-1} at 78 km to 3.6 ms^{-1} at 98 km. In the meridional wind,
 184 the AO tends to reduce with height, but the opposite is true for the SAO. As for the DT, its vertical variations
 185 roughly resemble each other in the zonal and meridional winds, with the peak magnitudes of about 21.0 and
 186 16.9 ms^{-1} at 94 km, respectively. Interestingly, in the difference resonant interaction between the DT and the
 187 AO, the generated wave has the maximum amplitudes of 12.5 and 18.5 ms^{-1} in the zonal and meridional
 188 winds at 94 km, and its vertical evolution is similar to that of the DT. In particular, in the meridional wind
 189 above 86 km, the secondary wave is more intense than the DT. In the sum resonant interaction, the
 190 secondary wave in the meridional wind increases to the maximum amplitude of 10.8 ms^{-1} at 88 km, and then
 191 decreases monotonously with altitude, while in the zonal wind, after experiencing a decline from 88 km to
 192 92 km, its amplitude grows up with height again. Whereas, in the sum and difference resonant interactions
 193 between the DT and SAO, the excited waves show the trend of first increasing and then decreasing with
 194 height, and have the maximum magnitudes at the heights slightly lower than 94 km.

195 Figure 5 depicts the variation of the four components in the interactions of the SDT with the AO (upper
 196 panels) and the SAO (lower panels). The SDT is strengthened in the zonal wind with the rising height, but is
 197 attenuated in the meridional wind after reaching the peak at 90 km. The amplitudes of these secondary
 198 waves vary with altitude. The secondary wave generated through the sum resonant interaction of the SDT
 199 with the AO is stronger in the meridional wind at 94-98 km than the SDT. On the whole, except at a few
 200 heights, the magnitudes of the SDT and SAO in the MLT at Wuhan are smaller than those of the DT and AO,
 201 respectively, thus the secondary waves in the interactions between the SDT and SAO are generally weaker



than those in the interactions between the DT and AO. In addition, since the work due to Reynolds stress in a sheared flow can lead to significant energy exchanges between wave and flow (Huang et al, 2010) and the wave dissipation from molecular and eddy diffusions is closely related to the wave scale (Huang et al., 2014), the magnitudes of the generated waves may also be associated with their dissipation and coupling with the background flow after they are excited.

4. Summary

In this paper, we use the meteor radar observations and MERRA2 reanalysis data of 9 years to demonstrate the occurrence of the resonant interactions between the DT/SDT and the AO/SAO in the MLT.

The Lomb-Scargle spectra of the zonal and meridional winds at 90 km show that only the AO, SAO, DT and SDT are the predominant components in both the zonal and meridional winds with the spectral amplitudes of 3.5-17.7 ms^{-1} . The terdiurnal tide is a relatively strong component with the magnitude of 2.2 ms^{-1} in the zonal and meridional winds, and the 16-day PW has the amplitude of 2.1 ms^{-1} in the zonal wind. Except these, all the other spectral amplitudes are smaller than 2 ms^{-1} .

When we zoom in on the partial spectra centered the DT, there exist several significant peaks in the upper and lower sidebands of the DT. These spectral peaks are located right at the sum and difference frequencies between the DT and the AO and SAO, meaning that these waves are generated through the nonlinear resonant interactions between the DT and the AO/SAO. The amplitudes of the secondary waves excited through the difference (sum) resonant interactions between the DT and the AO and SAO attain the large values of 10.6 and 6.8 ms^{-1} (6.5 and 8.4 ms^{-1}) in the zonal wind, and 16.9 and 8.6 ms^{-1} (9.9 and 8.5 ms^{-1}) in the meridional winds, respectively, which are far stronger than the amplitude (2.2 ms^{-1}) of the terdiurnal tide. In particular, the secondary wave in the difference resonant interaction between the DT and the AO has the larger amplitude (16.9 ms^{-1}) in the meridional wind than the DT (14.4 ms^{-1}). The frequency-wavenumber



spectra of the reanalysis zonal and meridional winds at 0.01 hPa (~80 km) indicate that the zonal wavenumber of the DT, AO and secondary waves is -1, 0 and -1, respectively, confirming that the resonant conditions of wavenumbers in the interactions are met. Similarly, the matching conditions of wavenumbers and frequencies are also satisfied for the triads of the SDT, AO/SAO and secondary waves. Therefore, the nonlinear resonant interactions between the DT/SDT and the AO/SAO do occur in the MLT.

In the height range of 78-98 km, the secondary waves in the interactions between the SDT and SAO are weaker than those in the interactions between the DT and AO because the amplitude of the SDT and SAO are generally smaller than those of the DT and AO, respectively. At some altitudes, the secondary waves have the larger amplitude than the DT/SDT, thus in the tidal studies, the magnitude of the secondary waves may be treated as the intensity of the tides when the time series of data is not long enough or the spectral components are not distinguished carefully.

Data availability. The meteor radar observations are provided by the IGGCAS and the Chinese Meridian Project at the website of <http://wdc.geophys.ac.cn/>, and the MERRA2 reanalysis data is from the NASA GES DISC at <https://gmao.gsfc.nasa.gov/reanalysis/MERRA-2/> (Gelaro et al., 2017).

Author contributions. KH proposed the scientific ideas. XH and KH completed the analysis and the manuscript. HC, SZ, CH, and YG discussed the results in the manuscript.

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- 402



403 **Caption:**

404 **Figure 1.** Daily mean (a) zonal and (b) meridional winds at 78-98 km between 1 January 2012 and 31
 405 December 2020 from meteor radar observation at Wuhan.

406 **Figure 2.** Lomb-Scargle spectra of (a) zonal and (b) meridional winds at 90 km from radar observation of 9
 407 years and their zoomed-in views around frequencies corresponding to (c, d) AO and SAO, (e, f) DT and (g,
 408 h) SDT. The left and right columns denote the spectra of the zonal and meridional winds, respectively. The
 409 red (blue) dashed vertical line is marked at the frequency of the AO (SAO) in Panels (c, d), and the sum and
 410 difference frequencies between the DT and the AO (SAO) in Panels (e, f) and between the SDT and the AO
 411 (SAO) in Panels (g, h).

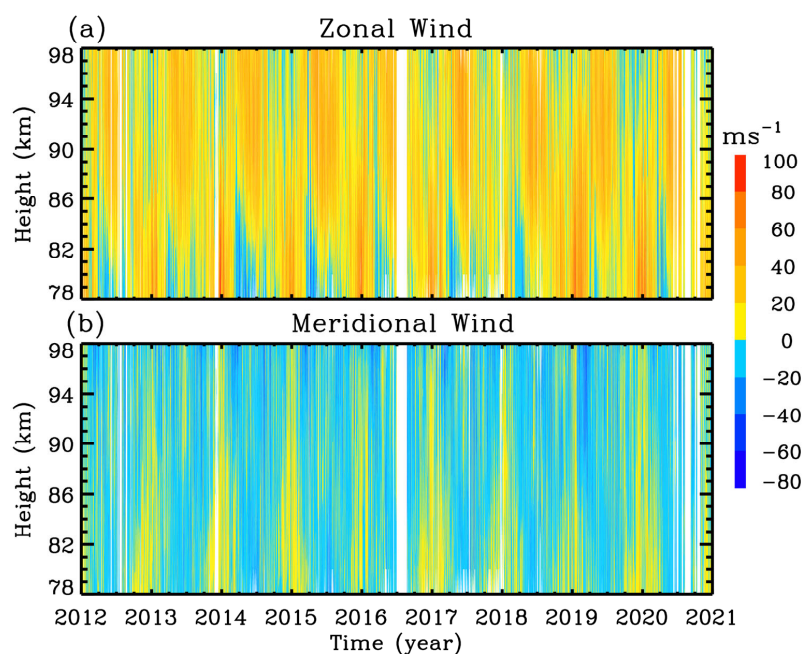
412 **Figure 3.** Zoomed-in views of frequency-wavenumber spectra of zonal and meridional winds at 0.01 hPa at
 413 30.5°N around frequencies corresponding to (a, b) AO and SAO and (c, d) DT derived from MERRA2
 414 reanalysis data of 9 years. The sign δ is the frequency of AO ($1/365 \text{ day}^{-1}$), and the left and right columns
 415 denote the spectra of the zonal and meridional winds, respectively. The negative wavenumber represents the
 416 westward phase progression. The dotted horizontal lines are marked at the frequencies of the AO and SAO
 417 in Panels (a, b), and at the frequency of the DT and the sum and difference frequencies between the DT and
 418 the AO/SAO in Panels (c, d). The dotted vertical lines are marked at the wavenumber of 0 in Panels (a, b)
 419 and -1 in Panels (c, d).

420 **Figure 4.** Spectral amplitudes of (blue) DT, (red) AO, (yellow) SAO, and (purple solid) sum and (purple
 421 dashed) difference frequency secondary waves generated through resonant interactions between DT and
 422 AO/SAO in (left) zonal and (right) meridional winds at 78-90 km. Panels (a, b) show the amplitudes of the
 423 DT, AO and secondary waves, and Panels (c, d) show the amplitudes of the DT, SAO and secondary waves.

424 **Figure 5.** Spectral amplitudes of (green) SDT, (red) AO, (yellow) SAO, and (purple solid) sum and (purple
 425 dashed) difference frequency secondary waves generated through resonant interactions between SDT and



426 AO/SAO in (left) zonal and (right) meridional winds at 78-90 km. Panels (a, b) show the amplitudes of the
427 SDT, AO and secondary waves, and Panels (c, d) show the amplitudes of the SDT, SAO and secondary
428 waves.
429



430

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432 December 2020 from meteor radar observation at Wuhan.

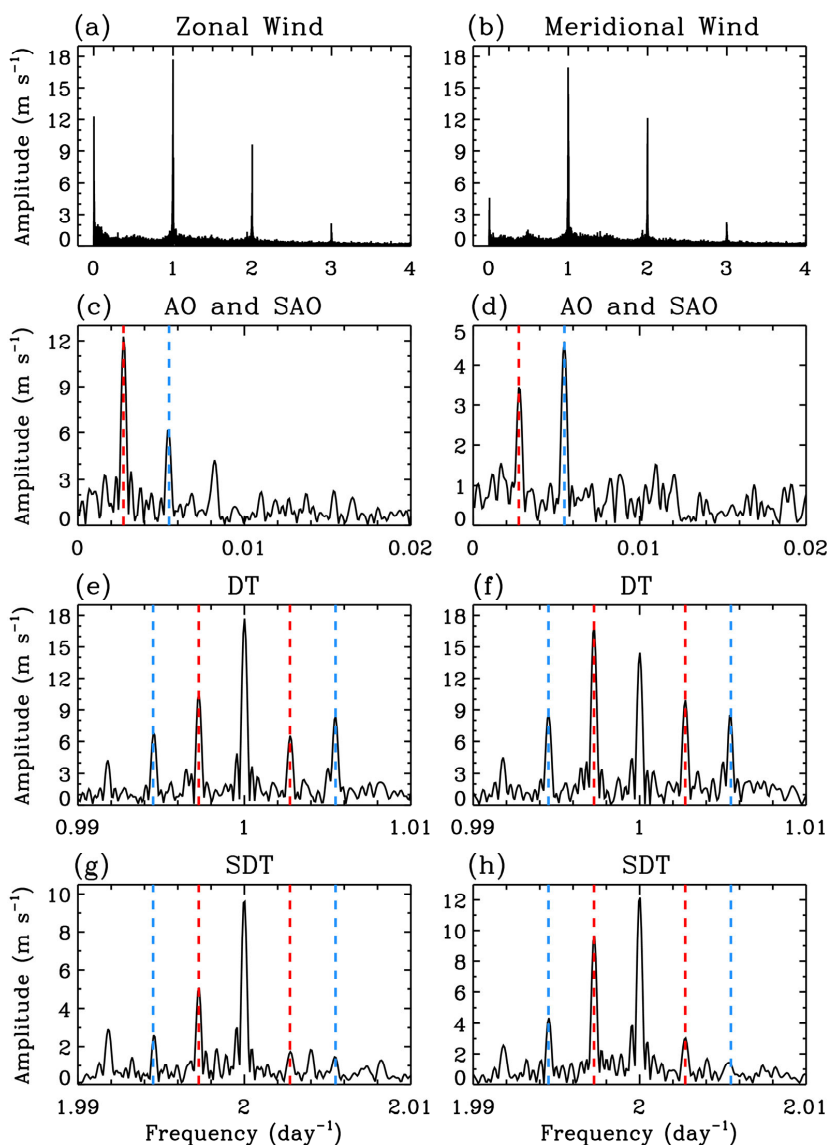


Figure 2. Lomb-Scargle spectra of (a) zonal and (b) meridional winds at 90 km from radar observation of 9 years and their zoomed-in views around frequencies corresponding to (c, d) AO and SAO, (e, f) DT and (g, h) SDT. The left and right columns denote the spectra of the zonal and meridional winds, respectively. The red (blue) dashed vertical line is marked at the frequency of the AO (SAO) in Panels (c, d), and the sum and difference frequencies between the DT and the AO (SAO) in Panels (e, f) and between the SDT and the AO (SAO) in Panels (g, h).

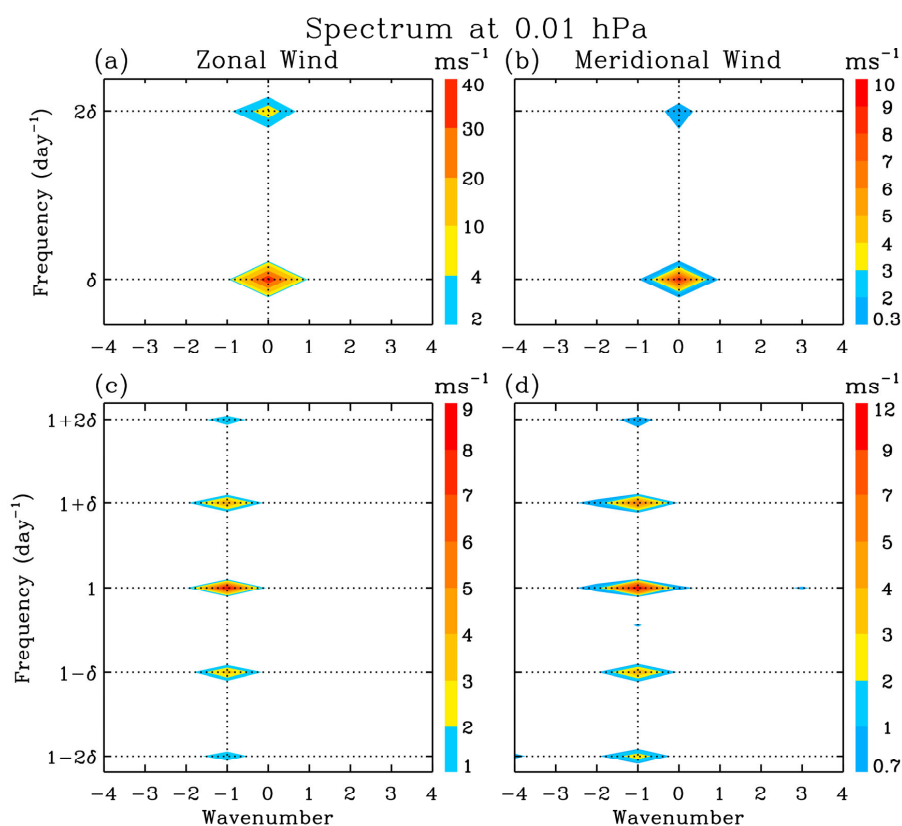
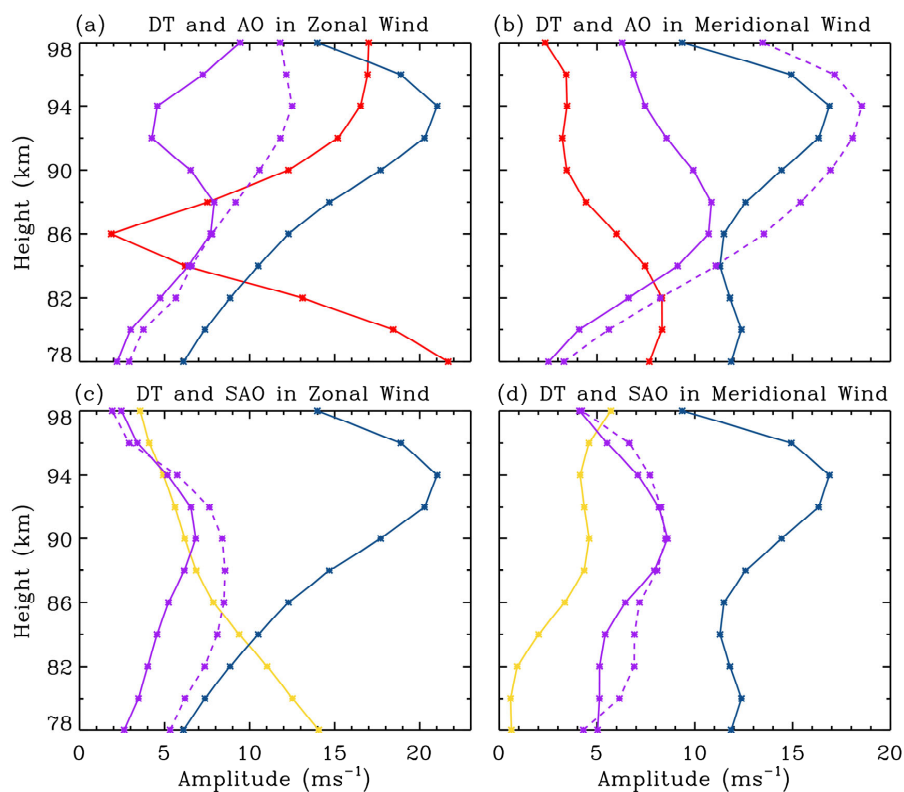
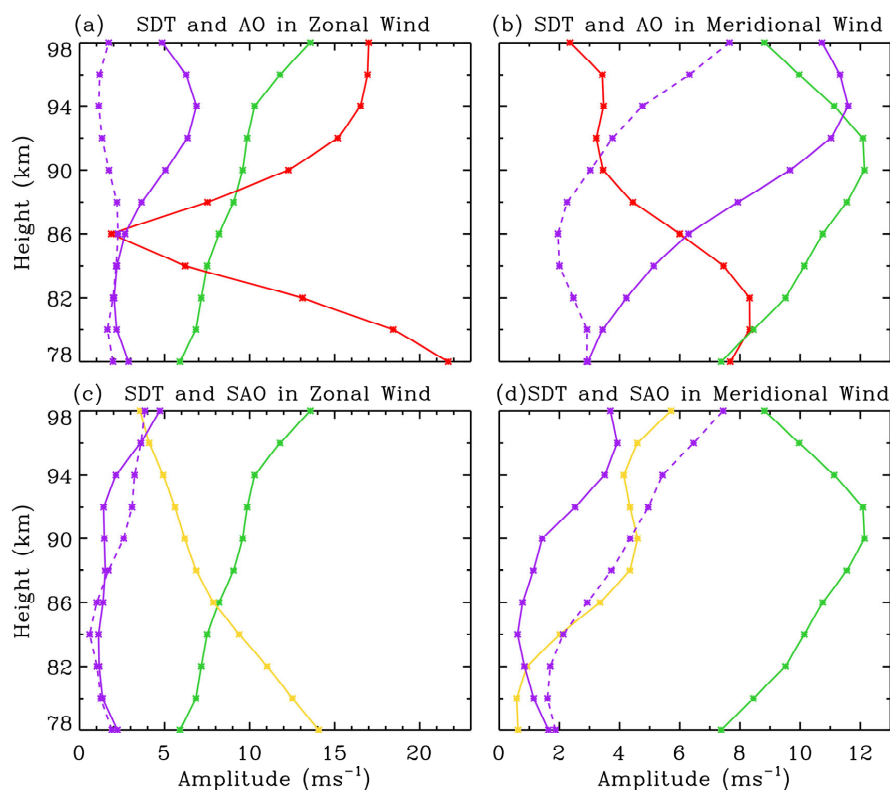


Figure 3. Zoomed-in views of frequency-wavenumber spectra of zonal and meridional winds at 0.01 hPa at 30.5°N around frequencies corresponding to (a, b) AO and SAO and (c, d) DT derived from MERRA2 reanalysis data of 9 years. The sign δ is the frequency of AO ($1/365 \text{ day}^{-1}$), and the left and right columns denote the spectra of the zonal and meridional winds, respectively. The negative wavenumber represents the westward phase progression. The dotted horizontal lines are marked at the frequencies of the AO and SAO in Panels (a, b), and at the frequency of the DT and the sum and difference frequencies between the DT and the AO/SAO in Panels (c, d). The dotted vertical lines are marked at the wavenumber of 0 in Panels (a, b) and -1 in Panels (c, d).



450

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455

456 **Figure 5.** Spectral amplitudes of (green) SDT, (red) AO, (yellow) SAO, and (purple solid) sum and (purple
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