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# **Nonlinear resonant interactions of atmospheric tides with annual oscillation based on meteor radar observation and reanalysis data**

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18 **Abstract.** Nonlinear interactions among gravity waves and among tides and planetary waves (PWs) have  
19 extensively been studied, however, resonant interactions between tides and annual (AO) and semiannual  
20 (SAO) oscillations were not reported. By using meteor radar observations and reanalysis data for 9 years, we  
21 demonstrate that the sum and difference resonant interactions between the diurnal (DT)/semidiurnal (SDT)  
22 tides and the AO/SAO do occur in the mesosphere and lower thermosphere (MLT). Both the frequencies and  
23 wavenumbers of the secondary waves in the sum (difference) resonant interactions just equal the sum  
24 (difference) frequencies and wavenumbers between the DT/SDT and the AO/SAO. Spectral analysis shows  
25 that only the DT, SDT, AO, SAO and their secondary waves are the predominant components in both the  
26 zonal and meridional winds at 90 km with the spectral amplitudes of 3.5-17.7 ms<sup>-1</sup>, being much stronger  
27 than all the other spectral amplitudes, including the amplitudes (2.1-2.2 ms<sup>-1</sup>) of the relatively strong  
28 terdiurnal tide and 16-day PW. At some altitudes in the MLT, the secondary waves are more intense than the  
29 DT/SDT, thus in tidal studies, the magnitude of the secondary waves may be regarded as the tidal one if the  
30 observational period is not long enough or their spectral peaks are not distinguished carefully.

31

## 32 **1 Introduction**

33 There are various motions with different temporal and spatial scales in the atmosphere, such as  
34 small-scale gravity waves (GWs), large-scale tidal waves and planetary waves (PWs), as well as longer  
35 period intraseasonal (ISO), semiannual (SAO) and annual (AO) oscillations. These perturbations play an  
36 important role in the circulation and thermal structure of the middle and upper atmosphere due to their  
37 ability to transfer energy and momentum among different atmospheric layers and among different latitudinal  
38 zones (Andrews et al., 1987). Meanwhile, there are complex coupling and nonlinear interactions among  
39 different scale waves and oscillations. For example, GWs and tides modulated by the filtering of PWs or  
40 ISOs in the lower atmosphere may drive similar period but out of phase PWs or ISOs in the mesosphere and



41 lower thermosphere (MLT) as momentum transported by GWs and tides is deposited into the background  
42 flow in the MLT due to their dissipation (Eckermann et al., 1997; Jacobi et al., 2006; Cheng et al, 2021).  
43 This process is essentially wave-flow coupling through the filtering effect as long as long-period PWs or  
44 ISOs are regarded as the temporal varying background field through which short-period GWs and tides  
45 propagate. Wave-wave nonlinear interaction is another dynamic process in which new spectral components  
46 are excited. In theory, nonlinear interaction comes from the nonlinear advective terms in the dynamic  
47 equations. Numerical simulation by Teitelbaum and Vial (1991) clearly showed that two secondary waves  
48 with sum and difference frequencies between semidiurnal tide (SDT) and 16-day PW were generated only  
49 through nonlinear interactions of SDT and 16-day wave rather than their linear superposition.

50 Nonlinear resonant interaction in the fluid dynamical context was proposed by Phillips [1960]. For  
51 specified wave vectors  $\vec{k}$  and frequencies  $\omega$ , two initial waves with phases  $\vec{k}_1 \cdot \vec{r} - \omega_1 t$  and  $\vec{k}_2 \cdot \vec{r} - \omega_2 t$   
52 can force a third wave with their sum or difference phase. In other words, if both the wave vector and  
53 frequency of a free third mode satisfy the sum or difference matching conditions,

54 
$$\vec{k}_1 \pm \vec{k}_2 = \vec{k}_3 \quad (1)$$

55 
$$\omega_1 \pm \omega_2 = \omega_3 \quad (2)$$

56 where the subscripts 1, 2 and 3 denote the three interacting waves, resonant interaction occurs and leads to  
57 significant energy transfer from the initial waves into the third wave. In the light of the two matching  
58 conditions, resonant interactions of small-scale GWs have been extensively investigated (e.g., Yeh and Liu,  
59 1981; Müller et al., 1986; Inhester, 1987; Klostermeyer, 1991; Stenflo, 1998; Yi, 1999; Wüst and Bittner,  
60 2006; Huang et al., 2009; Taklo and Choi, 2020), even in a sheared, dissipative and rotating atmosphere  
61 (Fritts et al., 1992; Vanneste, 1995; Axelsson et al., 1996; Huang et al., 2014; Achatz et al., 2017; Chatterjee  
62 and Misra, 2021; Voelker et al, 2021), and in the uniform and nonuniform plasmas (Stenflo, 1994; Stenflo  
63 and Shukla, 2009). For large-scale tides and PWs, many scientists studied their resonant interactions based



64 on general circulation models (GCMs) and ground-based and satellite-borne observations by using the  
65 matching condition ( $n_1 \pm n_2 = n_3$ ) of zonal wavenumbers  $n$  instead of the wave vector matching condition  
66 (Kamalabadi et al., 1997; Jacobi et al., 1998, 2001; Pancheva et al., 2000, 2002; Pancheva and Mitchell,  
67 2004; Liu et al., 2007; Palo et al., 2007; Kumar et al., 2008; McCormack et al., 2010; Babu et al., 2011;  
68 Chang et al., 2011a, 2011b; Yue et al., 2012; Huang et al., 2013a, 2013b; He and Chau, 2019).

69 Although nonlinear interactions between atmospheric waves have attracted wide attention, interactions  
70 of both global-scale tides and annual and semiannual oscillations were not reported yet. As we know, tides  
71 have a wave structure, but AO and SAO do not, thus it remains unclear whether they can interact with each  
72 other in the realistic atmosphere. Atmospheric tides are the persistent global perturbations and are generated  
73 mainly by radiation heating from the Sun (Lindzen and Chapman, 1969; Forbes and Garrett, 1979), in  
74 consequence, diurnal tide (DT) has a period of a solar day, and a westward wavenumber  $n = -1$  due to the  
75 westward migration of the Sun in the sky from the view of the Earth. Besides DT, there are its high-order  
76 harmonics in the atmosphere, among which SDT with a period of 12 h and a wavenumber of  $n = -2$  is the  
77 strongest component (Chapman and Lindzen, 1970; Hagan et al., 1999). Atmospheric AO and SAO are slow  
78 periodic changes of the zonal and meridional winds and temperature without phase propagation in the zonal  
79 direction. The periods of AO and SAO are 2-3 orders of magnitude larger than those of tides, thus the  
80 identification of sum and difference frequencies between tides and AO/SAO from tides is limited due to the  
81 requirement of rather long-time observations with high temporal resolution.

82 In the MLT, atmospheric waves have generally strong amplitudes, thereby, their nonlinear interactions  
83 are an important aspect of dynamics in the MLT. In present study, we use the observational and reanalysis  
84 data for 9 years to reveal the occurrence of resonant interactions between tides and AO/SAO. In Section 2,  
85 the radar observations and reanalysis data are described in brief. In Section 3, nonlinear interactions between  
86 DT/SDT and AO/SAO are discussed based on the observational and reanalysis data. A summary is provided



87 in Section 4.

88

## 89 **2 Data**

90 In this study, we use the meteor radar observation to investigate the nonlinear interactions between the  
91 tides and atmospheric oscillations in the MLT. The meteor radar is situated at Wuhan (30.5°N, 114.6°E) in  
92 China, established by the Institute of Geology and Geophysics, Chinese Academy of Sciences (IGGCAS).  
93 This system is an all-sky interferometric meteor radar (Hocking et al., 2001), operating at 38.5 MHz. Meteor  
94 trails are illuminated by a three-element Yagi antenna directed toward the zenith with a transmitted peak  
95 power of 7.5 kW and a pulse repetition frequency of 1.98 kHz (Xiong et al., 2004; Jiang et al., 2005; Yu et  
96 al., 2013). Echoes reflected by meteor trails are detected by five three-element Yagi antennas oriented along  
97 two orthogonal baselines. The receiving antennas are sampled every 13.3  $\mu\text{s}$ , resulting in a height resolution  
98 of 2 km. The hourly horizontal wind profile in the range of 70 to 110 km is deduced from the meteor trail  
99 velocities observed by radar, with a typical measurement uncertainty of 2-4  $\text{ms}^{-1}$  around 88-90 km (Hocking  
100 et al., 2001). The zonal and meridional winds for 9 years from 1 January 2012 to 31 December 2020 are  
101 applied to this investigation.

102 The number of meteor counts detected by radar showed a strong dependence on height, with an  
103 approximately Gaussian distribution centered at about 88-90 km, thus the data availability is high around  
104 88-90 km. As the height goes up and down, the number of detected meteors decreases, and then the  
105 acceptable data reduces gradually. In the period of our concern, including several gaps with the maximum  
106 missing data of 53 days, the data availability exceeds 80% in the height range of 80-96 km, and then  
107 decreases to 69% at 98 km and 67% at 78 km, but is less than 50% above 98 km and below 78 km. In this  
108 case, the radar observation in the altitude range of 78-98 km is chosen for our study.

109 Considering that the zonal wavenumber of planetary scale perturbations cannot be deduced from the



110 observations at a single station, we derive the zonal wavenumber from a two-dimensional Fourier transform  
111 with the help of Modern-Era Retrospective analysis for Research and Applications version 2 (MERRA2)  
112 reanalysis data. The upper pressure level at 0.01 hPa in the MERRA2 corresponds to the altitudes of about  
113 78-80 km, which are largely consistent with the lower limit of the radar observation utilized. The MERRA2  
114 data is produced by the Goddard Earth Sciences Data and Information Services Center (GES DISC) of the  
115 National Aeronautics and Space Administration (NASA) with some updates relative to the previous version  
116 (Gelaro et al., 2017), accessible from the online website at <https://gmao.gsfc.nasa.gov/reanalysis/MERRA-2/>.  
117 The reanalysis data is 6-hourly instantaneous analysis fields on a  $0.5^\circ \times 0.625^\circ$  latitude-by-longitude grid at  
118 72 pressure levels from the ground up to 0.01 hPa level (Gelaro et al. 2017). The reanalysis zonal and  
119 meridional winds for a same period from 1 January 2012 to 31 December 2020 are used in this study.

120

### 121 3. Results

#### 122 3.1 Resonant Interactions between DT/SDT and AO/SAO

123 Figure 1 shows the daily mean zonal and meridional winds from 78 to 98 km observed by the meteor  
124 radar at Wuhan for 9 years from 1 January 2012 to 31 December 2020. The different scale perturbations can  
125 be seen in the zonal and meridional wind fields.

126 Lomb-Scargle spectrum analysis is appropriate for decomposing unevenly sampled data into  
127 perturbation components (Scargle, 1982). In order to obtain the dominant periods in the wind fields, a  
128 Lomb-Scargle spectrum analysis, with a 4-times oversampling, is carried out on the zonal and meridional  
129 winds from the meteor radar observation between 1 January 2012 and 31 December 2020. Figures 2a and 2b  
130 present the Lomb-Scargle spectra of the zonal and meridional winds at 90 km, respectively. One can  
131 intuitively see from Figures 2a and 2b that the DT, SDT and low frequency oscillations are the prominent  
132 components, and the terdiurnal tide is also obvious. At Wuhan, the DT has the spectral amplitude of 17.7



133 and  $16.9 \text{ ms}^{-1}$  in the zonal and meridional winds, which are stronger than those of  $9.6$  and  $12.1 \text{ ms}^{-1}$  for the  
134 SDT, respectively. The terdiurnal tide has the magnitude of about  $2.2 \text{ ms}^{-1}$  in both the zonal and meridional  
135 winds.

136 For the sake of clarity, we zoom in on the several dominant spectral components. Figures 2c and 2d  
137 depict the low-frequency part (period  $>50$  days) of the Lomb-Scargle spectra of the zonal and meridional  
138 winds, respectively, and then the strong AO and SAO can be clearly identified from the low-frequency  
139 components. In the zonal wind, the amplitude ( $12.3 \text{ ms}^{-1}$ ) of the AO is more intense than that ( $6.2 \text{ ms}^{-1}$ ) of  
140 the SAO, while in the meridional wind, the magnitude ( $4.6 \text{ ms}^{-1}$ ) of the SAO is larger than that ( $3.5 \text{ ms}^{-1}$ ) of  
141 the AO. Besides the AO, SAO and tides mentioned above, only the 16-day planetary wave (not presented) in  
142 the zonal wind has the spectral magnitude of about  $2.1 \text{ ms}^{-1}$ , and all the other spectral amplitudes in the  
143 zonal and meridional winds are smaller than  $2 \text{ ms}^{-1}$ . Hence, from the perspective of 9-year mean results, the  
144 diurnal and semidiurnal tides and annual and semiannual oscillations are the predominant perturbations in  
145 the MLT at  $30.5^\circ\text{N}$ .

146 Figures 2e and 2f show the spectra centered at the DT in the zonal and meridional winds, respectively. It  
147 is interesting that the DT spectra in Figures 2a and 2b are divided into multiple spectrum peaks in the  
148 zoomed-in Figures 2e and 2f. We note that these spectral peaks around the DT are right at the sum and  
149 difference frequencies between the DT and the annual and semiannual oscillations. Since the unexpected  
150 components satisfy the sum and different resonant conditions of  $\omega_1 \pm \omega_2 = \omega_3$  with the DT and the AO/SAO,  
151 they should be excited through the nonlinear resonant interactions between the DT and the AO/SAO. In the  
152 zonal wind, the secondary waves excited by the tidal difference (sum) resonant interactions with the AO and  
153 SAO have the magnitudes of  $10.6$  and  $6.8 \text{ ms}^{-1}$  ( $6.5$  and  $8.4 \text{ ms}^{-1}$ ), respectively. In the meridional wind, the  
154 excited waves in the difference (sum) resonant interactions between the DT and the AO and SAO have the  
155 intensities of  $16.9$  and  $8.6 \text{ ms}^{-1}$  ( $9.9$  and  $8.5 \text{ ms}^{-1}$ ), respectively. Hence, these excited waves are much



156 stronger than the terdiurnal tide with the amplitude of  $2.2 \text{ ms}^{-1}$  in both the zonal and meridional winds.  
157 Moreover, it is surprising that in the tidal difference resonant interaction with the AO, the spectral amplitude  
158 ( $16.9 \text{ ms}^{-1}$ ) of the secondary wave in the meridional wind is larger than that ( $14.4 \text{ ms}^{-1}$ ) of the DT. Therefore,  
159 the spectral peak at the DT in Figure 2b is actually the peak of the generated secondary wave rather than the  
160 DT. Similarly, Figures 2g and 2h illustrate the occurrence of the resonant interactions between the SDT and  
161 the AO and SAO.

162 Now that the wavenumber matching conditions of  $n_1 \pm n_2 = n_3$  should be satisfied in the resonant  
163 interactions, we examine the zonal wavenumber relation with the aid of the MERRA2 reanalysis data. The  
164 reanalysis zonal and meridional winds at 0.01 hPa ( $\sim 80 \text{ km}$ ) at  $30.5^\circ\text{N}$  for the same period of 9 years are  
165 selected to make a two-dimensional Fourier transform. Figure 3 shows the zoomed-in  
166 frequency-wavenumber spectra around the AO and SAO and the DT. In Figure 3, we use the sign  $\delta$  to  
167 denote the frequency of the AO ( $1/365 \text{ day}^{-1}$ ), and the negative wavenumber represents the westward  
168 propagation of wave. One can clearly see from Figure 3 that the zonal wavenumbers in both the zonal and  
169 meridional winds are 0 for the AO and SAO and -1 for the DT, as we all know. The spectral peaks on both  
170 sides of the DT are located not only at the sum and difference frequencies between the DT and the AO/SAO,  
171 similar to the results of the radar observation in Figure 2, but also just at the wavenumber of -1, indicating  
172 that the wavenumber matching conditions of  $n_1 \pm n_2 = n_3$  are satisfied. Hence, this confirms that these new  
173 components are generated through the nonlinear resonant interactions between the DT and the AO/SAO. In  
174 addition, the similar phenomenon arises in the zoomed-in frequency-wavenumber spectra around the SDT  
175 with the zonal wavenumber of -2 (not presented), too.

### 176 3.2 Spectral Amplitude Variation with Altitude

177 Here, we investigate the magnitude variation of these generated waves with height. From the  
178 Lomb-Scargle spectra of the zonal and meridional winds for the period of 9 years observed by the meteor



179 radar, we extract the spectral amplitudes of the DT/SDT, AO/SAO, and secondary waves at altitude range of  
180 78-98 km. Figure 4 shows the amplitude evolution of the four components in the interactions of the DT with  
181 the AO (upper panels) and the SAO (lower panels). In the zonal wind, the amplitude of the AO decreases  
182 sharply from  $21.7 \text{ ms}^{-1}$  at 78 km to only  $1.9 \text{ ms}^{-1}$  at 86 km, and then increases rapidly to  $17.0 \text{ ms}^{-1}$  at 98 km,  
183 while the SAO diminishes gradually from  $14.1 \text{ ms}^{-1}$  at 78 km to  $3.6 \text{ ms}^{-1}$  at 98 km. In the meridional wind,  
184 the AO tends to reduce with height, but the opposite is true for the SAO. As for the DT, its vertical variations  
185 roughly resemble each other in the zonal and meridional winds, with the peak magnitudes of about 21.0 and  
186  $16.9 \text{ ms}^{-1}$  at 94 km, respectively. Interestingly, in the difference resonant interaction between the DT and the  
187 AO, the generated wave has the maximum amplitudes of  $12.5$  and  $18.5 \text{ ms}^{-1}$  in the zonal and meridional  
188 winds at 94 km, and its vertical evolution is similar to that of the DT. In particular, in the meridional wind  
189 above 86 km, the secondary wave is more intense than the DT. In the sum resonant interaction, the  
190 secondary wave in the meridional wind increases to the maximum amplitude of  $10.8 \text{ ms}^{-1}$  at 88 km, and then  
191 decreases monotonously with altitude, while in the zonal wind, after experiencing a decline from 88 km to  
192 92 km, its amplitude grows up with height again. Whereas, in the sum and difference resonant interactions  
193 between the DT and SAO, the excited waves show the trend of first increasing and then decreasing with  
194 height, and have the maximum magnitudes at the heights slightly lower than 94 km.

195 Figure 5 depicts the variation of the four components in the interactions of the SDT with the AO (upper  
196 panels) and the SAO (lower panels). The SDT is strengthened in the zonal wind with the rising height, but is  
197 attenuated in the meridional wind after reaching the peak at 90 km. The amplitudes of these secondary  
198 waves vary with altitude. The secondary wave generated through the sum resonant interaction of the SDT  
199 with the AO is stronger in the meridional wind at 94-98 km than the SDT. On the whole, except at a few  
200 heights, the magnitudes of the SDT and SAO in the MLT at Wuhan are smaller than those of the DT and AO,  
201 respectively, thus the secondary waves in the interactions between the SDT and SAO are generally weaker



202 than those in the interactions between the DT and AO. In addition, since the work due to Reynolds stress in  
203 a sheared flow can lead to significant energy exchanges between wave and flow (Huang et al, 2010) and the  
204 wave dissipation from molecular and eddy diffusions is closely related to the wave scale (Huang et al.,  
205 2014), the magnitudes of the generated waves may also be associated with their dissipation and coupling  
206 with the background flow after they are excited.

207

#### 208 4. Summary

209 In this paper, we use the meteor radar observations and MERRA2 reanalysis data of 9 years to  
210 demonstrate the occurrence of the resonant interactions between the DT/SDT and the AO/SAO in the MLT.

211 The Lomb-Scargle spectra of the zonal and meridional winds at 90 km show that only the AO, SAO, DT  
212 and SDT are the predominant components in both the zonal and meridional winds with the spectral  
213 amplitudes of 3.5-17.7  $\text{ms}^{-1}$ . The terdiurnal tide is a relatively strong component with the magnitude of 2.2  
214  $\text{ms}^{-1}$  in the zonal and meridional winds, and the 16-day PW has the amplitude of 2.1  $\text{ms}^{-1}$  in the zonal wind.  
215 Except these, all the other spectral amplitudes are smaller than 2  $\text{ms}^{-1}$ .

216 When we zoom in on the partial spectra centered the DT, there exist several significant peaks in the  
217 upper and lower sidebands of the DT. These spectral peaks are located right at the sum and difference  
218 frequencies between the DT and the AO and SAO, meaning that these waves are generated through the  
219 nonlinear resonant interactions between the DT and the AO/SAO. The amplitudes of the secondary waves  
220 excited through the difference (sum) resonant interactions between the DT and the AO and SAO attain the  
221 large values of 10.6 and 6.8  $\text{ms}^{-1}$  (6.5 and 8.4  $\text{ms}^{-1}$ ) in the zonal wind, and 16.9 and 8.6  $\text{ms}^{-1}$  (9.9 and 8.5  $\text{ms}^{-1}$ )  
222 in the meridional winds, respectively, which are far stronger than the amplitude (2.2  $\text{ms}^{-1}$ ) of the terdiurnal  
223 tide. In particular, the secondary wave in the difference resonant interaction between the DT and the AO has  
224 the larger amplitude (16.9  $\text{ms}^{-1}$ ) in the meridional wind than the DT (14.4  $\text{ms}^{-1}$ ). The frequency-wavenumber



225 spectra of the reanalysis zonal and meridional winds at 0.01 hPa (~80 km) indicate that the zonal  
226 wavenumber of the DT, AO and secondary waves is -1, 0 and -1, respectively, confirming that the resonant  
227 conditions of wavenumbers in the interactions are met. Similarly, the matching conditions of wavenumbers  
228 and frequencies are also satisfied for the triads of the SDT, AO/SAO and secondary waves. Therefore, the  
229 nonlinear resonant interactions between the DT/SDT and the AO/SAO do occur in the MLT.

230 In the height range of 78-98 km, the secondary waves in the interactions between the SDT and SAO are  
231 weaker than those in the interactions between the DT and AO because the amplitude of the SDT and SAO  
232 are generally smaller than those of the DT and AO, respectively. At some altitudes, the secondary waves  
233 have the larger amplitude than the DT/SDT, thus in the tidal studies, the magnitude of the secondary waves  
234 may be treated as the intensity of the tides when the time series of data is not long enough or the spectral  
235 components are not distinguished carefully.

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238 **Data availability.** The meteor radar observations are provided by the IGGCAS and the Chinese Meridian  
239 Project at the website of <http://wdc.geophys.ac.cn/>, and the MERRA2 reanalysis data is from the NASA  
240 GES DISC at <https://gmao.gsfc.nasa.gov/reanalysis/MERRA-2/> (Gelaro et al., 2017).

241

242 **Author contributions.** KH proposed the scientific ideas. XH and KH completed the analysis and the  
243 manuscript. HC, SZ, CH, and YG discussed the results in the manuscript.

244

245 **Competing interests.** The contact author has declared that neither they nor their co-authors have any  
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247



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253

#### 254 **References**

255 Achatz, U., Ribstein, B., Senf, F., and Klein, R.: The interaction between synoptic-scale balanced flow and a  
256 finite-amplitude mesoscale wave field throughout all atmospheric layers: weak and moderately strong  
257 stratification, *Q. J. Roy. Meteor. Soc.*, 143, 342-361, <https://doi.org/10.1002/qj.2926>, 2017.

258 Andrews, D. G., Holton, J. R., and Leovy, C. B.: Middle atmosphere dynamics, 1st edn., vol. 40, edited by:  
259 Dmowska, R., and Holton, J. R., San Diego, Calif., Academic Press Inc, p. 489, ISBN: 9780120585762,  
260 1987.

261 Axelsson, P., Larsson, J., and Stenflo, L.: Nonlinear interaction between acoustic gravity waves in a rotating  
262 atmosphere, *Nonlinear Proc. Geoph.*, 3, 216-220, <https://doi.org/10.5194/npg-3-216-1996>, 1996.

263 Babu, V. S., Kumar, K. K., John, S. R., Subrahmanyam, K. V., and Ramkumar, G.: Meteor radar  
264 observations of short-term variability of quasi 2 day waves and their interaction with tides and planetary  
265 waves in the mesosphere-lower thermosphere region over Thumba (8.5°N, 77°E), *J. Geophys. Res.*, 116,  
266 D16121, <https://doi.org/10.1029/2010JD015390>, 2011.

267 Chang, L. C., Liu, J.-Y., and Palo, S. E.: Propagating planetary wave coupling in SABER MLT temperatures  
268 and GPS TEC during the 2005/2006 austral summer, *J. Geophys. Res.*, 116, A10324,  
269 <https://doi.org/10.1029/2011JA016687>, 2011a.

270 Chang, L. C., Palo, S. E., and Liu, H.-L.: Short-term variability in the migrating diurnal tide caused by  
271 interactions with the quasi 2 day wave, *J. Geophys. Res.*, 116, D12112,  
272 <https://doi.org/10.1029/2010JD014996>, 2011b.

273 Chapman, S. and Lindzen, R. S.: Atmospheric tides, 1st edn., Springer Dordrecht, Dordrecht, Holland,  
274 <https://doi.org/10.1007/978-94-010-3399-2>, 1970.



- 275 Chatterjee, D. and Misra, A. P.: Effects of Coriolis force on the nonlinear interactions of acoustic-gravity  
276 waves in the atmosphere, *J. Atmos. Sol.-Terr. Phy.*, 222, 105722,  
277 <https://doi.org/10.1016/j.jastp.2021.105722>, 2021.
- 278 Cheng, H., Huang, K. M., Liu, A. Z., Zhang, S. D., Huang, C. M., and Gong, Y.: A quasi-27-day oscillation  
279 activity from the troposphere to the mesosphere and lower thermosphere at low latitudes, *Earth Planets  
280 Space*, 73, 183, <https://doi.org/10.1186/s40623-021-01521-1>, 2021.
- 281 Eckermann, S. D., Rajopadhyaya, D. K., and Vincent, R. A.: Intraseasonal wind variability in the equatorial  
282 mesosphere and lower thermosphere: Long-term observations from central Pacific, *J. Atmos. Sol.-Terr.  
283 Phy.*, 59, 603-627, [https://doi.org/10.1016/S1364-6826\(96\)00143-5](https://doi.org/10.1016/S1364-6826(96)00143-5), 1997.
- 284 Fritts, D. C., Sun, S. J., and Wang, D. Y.: Wave-wave interactions in compressible atmosphere: 1. A general  
285 formulation including rotation and wind shear, *J. Geophys. Res.*, 97, 9975-9988,  
286 <https://doi.org/10.1029/92JD00347>, 1992.
- 287 Forbes, J. M. and Garrett, H. B.: Theoretical studies of atmospheric tides, *Rev. Geophys.*, 17, 1951-1981,  
288 <https://doi.org/10.1029/RG017i008p01951>, 1979.
- 289 Gelaro, R., McCarty, W., Suárez, M. J., Todling, R., Molod, A., Takacs, L., Randles, C. A., Darmenov, A.,  
290 Bosilovich, M. G., Reichle, R., Wargan, K., Coy, L., Cullather, R., Draper, C., Akella, S., Buchard, V.,  
291 Conaty, A., da Silva, A. M., Gu, W., Kim, G.-K., Koster, R., Lucchesi, R., Merkova, D., Nielsen, J. E.,  
292 Partyka, G., Pawson, S., Putman, W., Rienecker, M., Schubert, S. D., Sienkiewicz, M., and Zhao, B.: The  
293 modern-era retrospective analysis for research and applications, version 2 (MERRA-2), *J. Climate*, 30,  
294 5419-5454, <https://doi.org/10.1175/JCLI-D-16-0758.1>, 2017.
- 295 Hagan, M. E., Burrage, M. D., Forbes, J. M., Hackney, J., Randel, W. J., and Zhang, X.: GSWM-98: results  
296 for migrating solar tides, *J. Geophys. Res.*, 104, 6813-6827, <https://doi.org/10.1029/1998JA900125>,  
297 1999.
- 298 He, M. and Chau, J. L.: Mesospheric semidiurnal tides and near-12 h waves through jointly analyzing  
299 observations of five specular meteor radars from three longitudinal sectors at boreal midlatitudes, *Atmos.  
300 Chem. Phys.*, 19, 5993-6006, <https://doi.org/10.5194/acp-19-5993-2019>, 2019.
- 301 Hocking, W. K., Fuller, B., and Vandeppeer, B.: Real-time determination of meteor-related parameters



- 302 utilizing modern digital technology, *J. Atmos. Sol.-Terr. Phys.*, 63, 155-169,  
303 [https://doi.org/10.1016/S1364-6826\(00\)00138-3](https://doi.org/10.1016/S1364-6826(00)00138-3), 2001.
- 304 Huang, K. M., Zhang, S. D., and Yi, F.: Gravity wave excitation through resonant interaction in a  
305 compressible atmosphere, *Geophys. Res. Lett.*, 36, L01803, <https://doi.org/10.1029/2008GL035575>,  
306 2009.
- 307 Huang, K. M., Zhang, S. D., and Yi, F.: Reflection and transmission of atmospheric gravity waves in a stably  
308 sheared horizontal wind field, *J. Geophys. Res.*, 115, D16103, <https://doi.org/10.1029/2009JD012687>,  
309 2010.
- 310 Huang, K. M., Liu, A. Z., Lu, X., Li, Z., Gan, Q., Gong, Y., Huang, C. M., Yi, F., and Zhang, S. D.:  
311 Nonlinear coupling between quasi 2 day wave and tides based on meteor radar observations at Maui, *J.*  
312 *Geophys. Res.-Atmos.*, 118, 10936-10943, <https://doi.org/10.1002/jgrd.50872>, 2013a.
- 313 Huang, K. M., Liu, A. Z., Zhang, S. D., Yi, F., Huang, C. M., Gan, Q., Gong, Y., and Zhang, Y. H.: A  
314 nonlinear interaction event between a 16-day wave and a diurnal tide from meteor radar observations,  
315 *Ann. Geophys.*, 31, 2039-2048, <https://doi.org/10.5194/angeo-31-2039-2013>, 2013b.
- 316 Huang, K. M., Zhang, S. D., Yi, F., Huang, C. M., Gan, Q., Gong, Y., and Zhang, Y. H.: Nonlinear  
317 interaction of gravity waves in a nonisothermal and dissipative atmosphere, *Ann. Geophys.*, 32, 263-275,  
318 <https://doi.org/10.5194/angeo-32-263-2014>, 2014.
- 319 Inhester, B.: Numerical modeling of hydromagnetic wave coupling in the magnetosphere, *J. Geophys.*  
320 *Res.-Space*, 92, 4751-4756, <https://doi.org/10.1029/JA092iA05p04751>, 1987.
- 321 Jacobi, Ch., Schindler, R., and Kürschner, D.: Non-linear interaction of the quasi 2-day wave and long-term  
322 oscillations in the summer midlatitude mesopause region as seen from LF D1 wind measurements over  
323 Central Europe (Collm, 52°N, 15°E), *J. Atmos. Sol.-Terr. Phys.*, 60, 1175-1191,  
324 [https://doi.org/10.1016/S1364-6826\(98\)00076-5](https://doi.org/10.1016/S1364-6826(98)00076-5), 1998.
- 325 Jacobi, Ch., Portnyagin, Yu. I., Merzlyakov, E. G., Kashcheyev, B. L., Oleynikov, A. N., Kürschner, D.,  
326 Mitchell, N. J., Middleton, H. R., Muller, H. G., and Comley, V. E.: Mesosphere/lower thermosphere  
327 wind measurements over Europe in summer 1998, *J. Atmos. Sol.-Terr. Phys.*, 63, 1017-1031,  
328 [https://doi.org/10.1016/S1364-6826\(01\)00012-8](https://doi.org/10.1016/S1364-6826(01)00012-8), 2001.
- 329 Jacobi, Ch., Fröhlich, K., and Pogoreltsev, A.: Quasi two-day-wave modulation of gravity wave flux and



- 330 consequences for the planetary wave propagation in a simple circulation model, *J. Atmos. Sol.-Terr. Phys.*,  
331 68, 283-292, <https://doi.org/10.1016/j.jastp.2005.01.017>, 2006.
- 332 Jiang, G., Xiong, J.-G., Wan, W.-X., Ning, B.-Q., Liu, L.-B., Vincent, R. A., and Reid, I.: The 16-day waves  
333 in the mesosphere and lower thermosphere over Wuhan (30.6°N, 114.5°E) and Adelaide (35°S, 138°E),  
334 *Adv. Space Res.*, 35, 2005-2010, <https://doi.org/10.1016/j.asr.2005.03.011>, 2005.
- 335 Kamalabadi, F., Forbes, J. M., Makarov, N. M., and Portnyagin, Yu. I.: Evidence for nonlinear coupling of  
336 planetary waves and tides in the Antarctic mesopause, *J. Geophys. Res.*, 102, 4437-4446,  
337 <https://doi.org/10.1029/96JD01996>, 1997.
- 338 Klostermeyer, J.: Two- and three-dimensional parametric instabilities in finite-amplitude internal gravity  
339 waves, *Geophys. Astro. Fluid*, 61, 1-25, <https://doi.org/10.1080/03091929108229035>, 1991.
- 340 Kumar, K. K., Deepa, V., Antonita, T. M., and Ramkumar, G.: Meteor radar observations of short-term tidal  
341 variabilities in the low-latitude mesosphere-lower thermosphere: Evidence for nonlinear wave-wave  
342 interactions, *J. Geophys. Res.*, 113, D16108, <https://doi.org/10.1029/2007JD009610>, 2008.
- 343 Lindzen, R. S. and Chapman, S.: Atmospheric tides, *Space Sci. Rev.*, 10,  
344 3-188, <https://doi.org/10.1007/BF00171584>, 1969.
- 345 Liu, H.-L., Li, T., She, C.-Y., Oberheide, J., Wu, Q., Hagan, M. E., Xu, J., Roble, R. G., Mlynczak, M. G.,  
346 and Russell III, J. M.: Comparative study of short-term diurnal tidal variability, *J. Geophys. Res.*, 112,  
347 D18108, <https://doi.org/10.1029/2007JD008542>, 2007.
- 348 McCormack, J. P., Eckermann, S. D., Hoppel, K. W., and Vincent, R. A.: Amplification of the quasi-two day  
349 wave through nonlinear interaction with the migrating diurnal tide, *Geophys. Res. Lett.*, 37, L16810,  
350 <https://doi.org/10.1029/2010GL043906>, 2010.
- 351 Müller, P., Holloway, G., Henyey, F., and Pomphrey, N.: Nonlinear interactions among internal gravity  
352 waves, *Rev. Geophys.*, 24, 493-536, <https://doi.org/10.1029/RG024i003p00493>, 1986.
- 353 Palo, S. E., Forbes, J. M., Zhang, X., Russell III, J. M., and Mlynczak, M. G.: An eastward propagating  
354 two-day wave: Evidence for nonlinear planetary wave and tidal coupling in the mesosphere and lower  
355 thermosphere, *Geophys. Res. Lett.*, 34, L07807, <https://doi.org/10.1029/2006GL027728>, 2007.
- 356 Pancheva, D. V. and Mitchell, N. J.: Planetary waves and variability of the semidiurnal tide in the  
357 mesosphere and lower thermosphere over Esrange (68°N, 21°E) during winter, *J. Geophys. Res.*, 109,



- 358 A08307, <https://doi.org/10.1029/2004JA010433>, 2004.
- 359 Pancheva, D. V., Beard, A. G., and Mitchell, N. J.: Nonlinear interactions between planetary waves in the  
360 mesosphere/lower-thermosphere region, *J. Geophys. Res.*, 105, 157-170,  
361 <https://doi.org/10.1029/1999JA900332>, 2000.
- 362 Pancheva, D., Merzlyakov, E., Mitchell, N. J., Portnyagin, Y., Manson, A. H., Jacobi, Ch., Meek, C. E., Luo,  
363 Y., Clark, R. R., Hocking, W. K., MacDougall, J., Muller, H. G., Kürschner, D., Jones, G. O. L., Vincent,  
364 R. A., Reid, I. M., Singer, W., Igarashi, K., Fraser, G. I., Fahrutdinova, A. N., Stepanov, A. M., Poole, L.  
365 M. G., Mailga, S. B., Kashcheyev, B. L., and Oleynikov, A. N.: Global-scale tidal variability during the  
366 PSMOS campaign of June-August 1999: Interaction with planetary waves, *J. Atmos. Sol.-Terr. Phys.*, 64,  
367 1865-1896, [https://doi.org/10.1016/S1364-6826\(02\)00199-2](https://doi.org/10.1016/S1364-6826(02)00199-2), 2002.
- 368 Phillips, O. M.: On the dynamics of unsteady gravity waves of finite amplitude Part 1. The elementary  
369 interactions, *J. Fluid Mech.*, 9, 193-217, <https://doi.org/10.1017/S0022112060001043>, 1960.
- 370 Scargle, J. D.: Studies in astronomical time series analysis: II. Statistical aspects of spectral analysis of  
371 unevenly spaced data, *Astrophys. J.*, 263, 835-853, <https://doi.org/10.1086/160554>, 1982.
- 372 Stenflo, L.: Resonant three-wave interactions in plasmas, *Phys. Scripta*, T50, 15-19,  
373 <https://doi.org/10.1088/0031-8949/1994/T50/002>, 1994.
- 374 Stenflo, L.: Nonlinear equations for acoustic gravity waves, *Phys. Scripta*, T75, 306,  
375 <http://dx.doi.org/10.1238/physica.topical.075a00306>, 1998.
- 376 Stenflo, L. and Shukla, P. K.: Nonlinear acoustic-gravity waves, *J. Plasma Phys.*, 75, 841-847,  
377 <https://doi.org/10.1017/S0022377809007892>, 2009.
- 378 Taklo, T. M. A. and Choi, W.: Group resonant interactions between surface and internal gravity waves in a  
379 two-layer system, *J. Fluid Mech.*, 892, A14, <https://doi.org/10.1017/jfm.2020.180>, 2020.
- 380 Teitelbaum, H., and Vial, F.: On tidal variability induced by nonlinear interaction with planetary waves, *J.*  
381 *Geophys. Res.*, 96, 14169-14178, <https://doi.org/10.1029/91JA01019>, 1991.
- 382 Vanneste, J.: The instability of internal gravity waves to localized disturbances, *Ann. Geophys.*, 13, 196-210,  
383 <https://doi.org/10.1007/s00585-995-0196-7>, 1995.
- 384 Voelker, G. S., Akylas, T. R., and Achatz, U.: An application of WKB theory for triad interactions of  
385 internal gravity waves in varying background flows, *Q. J. Roy. Meteor. Soc.*, 147, 1112-1134,



- 386 <https://doi.org/10.1002/qj.3962>, 2021.
- 387 Wüst, S. and Bittner, M.: Non-linear resonant wave-wave interaction (triad): Case studies based on rocket  
388 data and first application to satellite data, *J. Atmos. Sol.-Terr. Phy.*, 68, 959-976,  
389 <https://doi.org/10.1016/j.jastp.2005.11.011>, 2006.
- 390 Xiong, J.-G., Wan, W., Ning, B., and Liu, L.: First results of the tidal structure in the MLT revealed by  
391 Wuhan meteor radar (30°40'N, 114°30'E), *J. Atmos. Sol.-Terr. Phy.*, 66, 675-682,  
392 <https://doi.org/10.1016/j.jastp.2004.01.018>, 2004.
- 393 Yeh, K. C., and Liu, C. H.: The instability of atmospheric gravity waves through wave-wave interactions, *J.*  
394 *Geophys. Res.*, 86, 9722-9728, <https://doi.org/10.1029/JC086iC10p09722>, 1981.
- 395 Yi, F.: Resonant interactions between propagating gravity wave packets, *J. Atmos. Sol.-Terr. Phy.*, 61,  
396 675-691, [https://doi.org/10.1016/S1364-6826\(99\)00026-7](https://doi.org/10.1016/S1364-6826(99)00026-7), 1999.
- 397 Yu, Y., Wan, W., Ning, B., Liu, L., Wang, Z., Hu, L., and Ren, Z.: Tidal wind mapping from observations of  
398 a meteor radar chain in December 2011, *J. Geophys. Res.*, 118, 2321-2332,  
399 <https://doi.org/10.1029/2012JA017976>, 2013.
- 400 Yue, J., Liu, H.-L., and Chang, L. C.: Numerical investigation of the quasi 2 day wave in the mesosphere  
401 and lower thermosphere, *J. Geophys. Res.*, 117, D05111, <https://doi.org/10.1029/2011JD016574>, 2012.
- 402



403 **Caption:**

404 **Figure 1.** Daily mean (a) zonal and (b) meridional winds at 78-98 km between 1 January 2012 and 31  
405 December 2020 from meteor radar observation at Wuhan.

406 **Figure 2.** Lomb-Scargle spectra of (a) zonal and (b) meridional winds at 90 km from radar observation of 9  
407 years and their zoomed-in views around frequencies corresponding to (c, d) AO and SAO, (e, f) DT and (g,  
408 h) SDT. The left and right columns denote the spectra of the zonal and meridional winds, respectively. The  
409 red (blue) dashed vertical line is marked at the frequency of the AO (SAO) in Panels (c, d), and the sum and  
410 difference frequencies between the DT and the AO (SAO) in Panels (e, f) and between the SDT and the AO  
411 (SAO) in Panels (g, h).

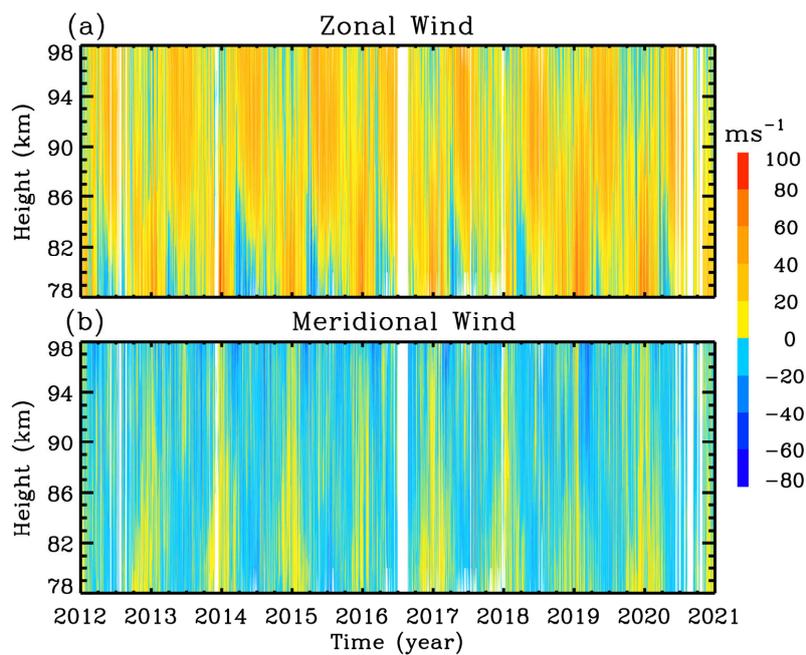
412 **Figure 3.** Zoomed-in views of frequency-wavenumber spectra of zonal and meridional winds at 0.01 hPa at  
413 30.5°N around frequencies corresponding to (a, b) AO and SAO and (c, d) DT derived from MERRA2  
414 reanalysis data of 9 years. The sign  $\delta$  is the frequency of AO ( $1/365 \text{ day}^{-1}$ ), and the left and right columns  
415 denote the spectra of the zonal and meridional winds, respectively. The negative wavenumber represents the  
416 westward phase progression. The dotted horizontal lines are marked at the frequencies of the AO and SAO  
417 in Panels (a, b), and at the frequency of the DT and the sum and difference frequencies between the DT and  
418 the AO/SAO in Panels (c, d). The dotted vertical lines are marked at the wavenumber of 0 in Panels (a, b)  
419 and -1 in Panels (c, d).

420 **Figure 4.** Spectral amplitudes of (blue) DT, (red) AO, (yellow) SAO, and (purple solid) sum and (purple  
421 dashed) difference frequency secondary waves generated through resonant interactions between DT and  
422 AO/SAO in (left) zonal and (right) meridional winds at 78-90 km. Panels (a, b) show the amplitudes of the  
423 DT, AO and secondary waves, and Panels (c, d) show the amplitudes of the DT, SAO and secondary waves.

424 **Figure 5.** Spectral amplitudes of (green) SDT, (red) AO, (yellow) SAO, and (purple solid) sum and (purple  
425 dashed) difference frequency secondary waves generated through resonant interactions between SDT and



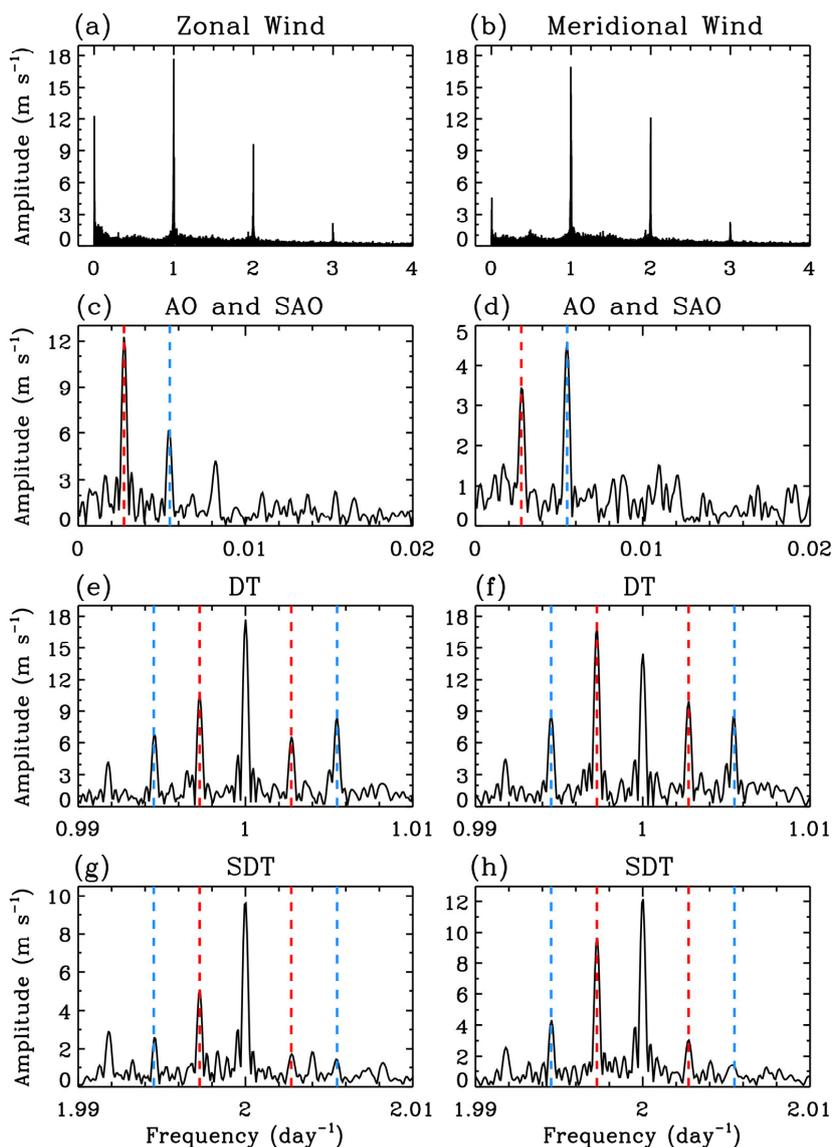
426 AO/SAO in (left) zonal and (right) meridional winds at 78-90 km. Panels (a, b) show the amplitudes of the  
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429



430

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433

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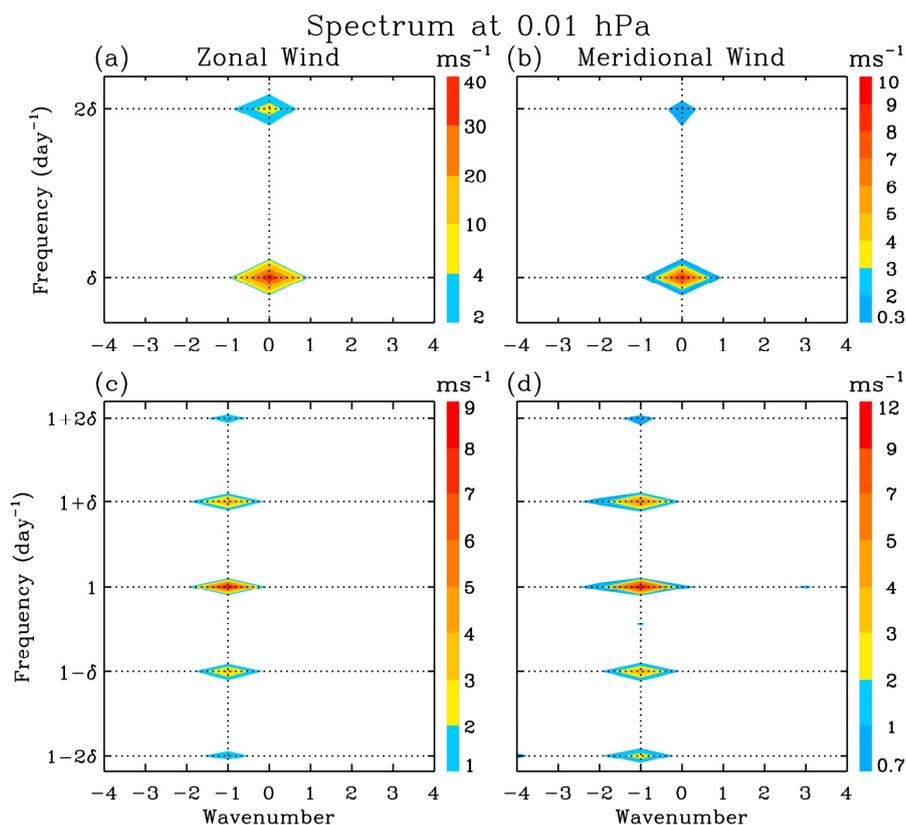
435 years and their zoomed-in views around frequencies corresponding to (c, d) AO and SAO, (e, f) DT and (g,

436 h) SDT. The left and right columns denote the spectra of the zonal and meridional winds, respectively. The

437 red (blue) dashed vertical line is marked at the frequency of the AO (SAO) in Panels (c, d), and the sum and

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439 (SAO) in Panels (g, h).



440

441 **Figure 3.** Zoomed-in views of frequency-wavenumber spectra of zonal and meridional winds at 0.01 hPa at

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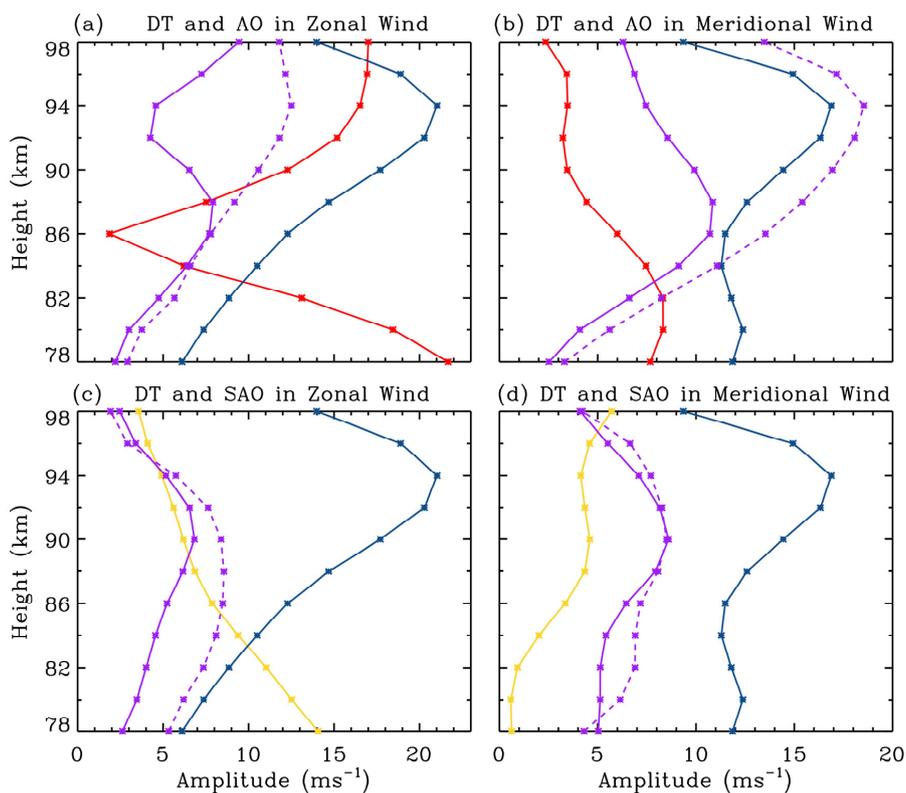
445 westward phase progression. The dotted horizontal lines are marked at the frequencies of the AO and SAO

446 in Panels (a, b), and at the frequency of the DT and the sum and difference frequencies between the DT and

447 the AO/SAO in Panels (c, d). The dotted vertical lines are marked at the wavenumber of 0 in Panels (a, b)

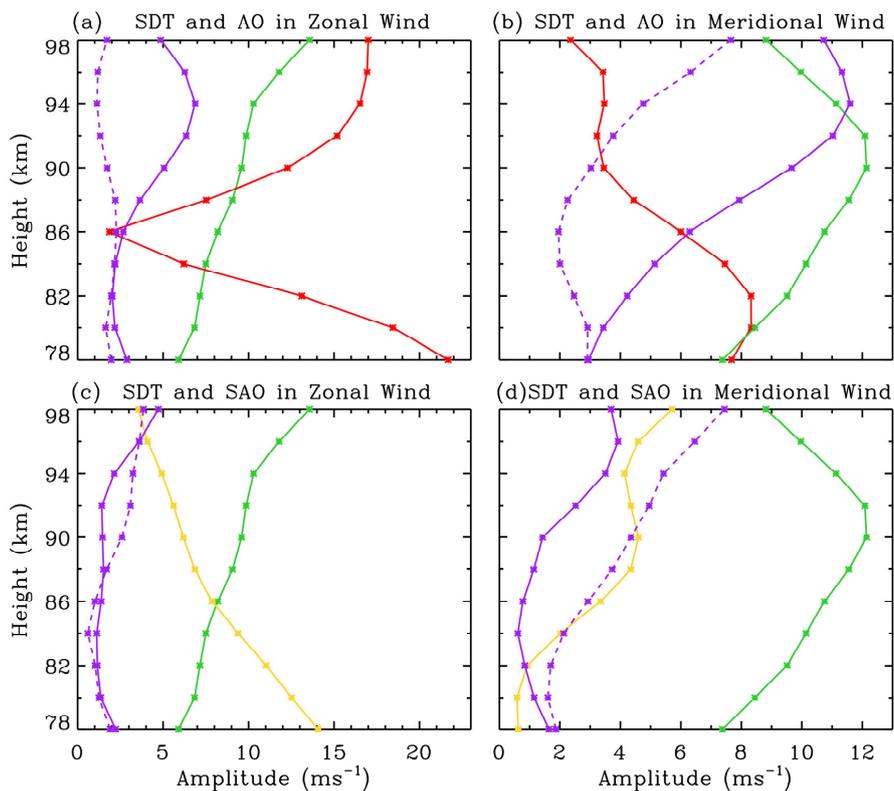
448 and -1 in Panels (c, d).

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451 **Figure 4.** Spectral amplitudes of (blue) DT, (red) AO, (yellow) SAO, and (purple solid) sum and (purple  
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453 AO/SAO in (left) zonal and (right) meridional winds at 78-90 km. Panels (a, b) show the amplitudes of the  
454 DT, AO and secondary waves, and Panels (c, d) show the amplitudes of the DT, SAO and secondary waves.



455

456 **Figure 5.** Spectral amplitudes of (green) SDT, (red) AO, (yellow) SAO, and (purple solid) sum and (purple  
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458 AO/SAO in (left) zonal and (right) meridional winds at 78-90 km. Panels (a, b) show the amplitudes of the  
459 SDT, AO and secondary waves, and Panels (c, d) show the amplitudes of the SDT, SAO and secondary  
460 waves.