

Antarctic atmospheric Richardson number from radiosoundings measurements and AMPS

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Abstract. Monitoring a wide range of atmospheric turbulence over the Antarctic continent is still tricky, while the atmospheric Richardson number (Ri ; a critical/valuable parameter ~~determining which determines~~ the possibility ~~of that~~ turbulence could be triggered) is easier to obtain. The Antarctic atmospheric Ri , calculated using from the potential temperature and wind speed, was investigated using the daily results from the radiosoundings and forecasts of the Antarctic Mesoscale Prediction System (AMPS). Radiosoundings for a year at three sites (McMurdo, South Pole, and Dome C) were used to quantify the reliability of the AMPS forecasts. The AMPS-forecasted Ri (~~inverse of the Richardson number~~) can identify the main spatiotemporal characteristics of atmospheric turbulence over the Antarctic ~~continent in terms of space and time region~~. The correlation coefficients (R_{xy}) of $\log_{10}(Ri)$ at McMurdo, the South Pole, and Dome C are 0.71, 0.6659, and 0.6853, respectively. The Ri , where was generally underestimated by the performance gains during the warm seasons. In addition, a model to improve AMPS-forecasted has been presented. The monthly median at and the AMPS could better capture the trend of $\log_{10}(Ri)$ three sites and the under relatively unstable atmospheric conditions. The seasonal median of $\log_{10}(Ri)$ throughout the along two vertical cross-sections for of the AMPS forecasts are presented. One can observe that the probability of triggering, and it shows some zones where atmospheric turbulence can be highly triggered in Antarctica. The $\log_{10}(Ri)$ is primarily concentrated near the ground. In addition, strong wind shears near escarpment regions have been found in the range of 0.5 km above the ground, thus causing atmospheric instability (or a thick boundary layer). In addition, turbulent atmospheres are likely distributions appear to be reasonably correlated to be triggered some large-scale phenomena or local-scale dynamics (katabatic winds, polar vortices, convection, gravity wave, etc) over the ocean, moving toward the Antarctic Plateau plateau and becoming stable surrounding ocean. Finally, the $\log_{10}(Ri)$ at the planetary boundary layer height ($PBLH$),, has been provided as a reference standard for judging atmospheric stability. The) were calculated and their median value of from the combined data of two vertical cross sections was 0.55, which is 0.316, this median value, in turn, was used to ~~calculate estimate~~ $PBLH$ and agree well with the AMPS-forecasted $PBLH$ ($R_{xy} > 0.69$). Overall, our results suggest that the estimated $\log_{10}(Ri)$ ~~forecasts (> 0.72)~~ by AMPS are reasonable and the turbulence conditions in Antarctica are well revealed.

1 Introduction

The Richardson number (Ri) is a valuable parameter for giving insight into atmospheric stability ~~monitoring~~; it combines both thermodynamic and dynamic profiles, which provides us with valuable insights into turbulent heat fluxes (Town and Walden,

35 2009) and the probability that optical turbulence (Yang et al., 2021; Yang et al., 2022) can be triggered in Antarctica. However,
36 the measurements of atmospheric properties in Antarctica are sparse compared to those in the mid-latitudes and tropics.
37 Atmospheric models have been developed to overcome this limitation (Meso-NH by Lascaux et al., 2009; Polar WRF by
38 Bromwich et al., 2013; MAR by Gallée et al., 2015), allowing researchers to investigate atmospheric variability beyond
39 observational coverage, even for forecasting atmospheric ~~properties~~parameters in the future.

40 The Antarctic Mesoscale Prediction System (AMPS; <https://www2.mmm.ucar.edu/rt/amps/>) runs a real-time atmospheric model
41 and provides numerical forecasts [for Antarctica. The performances of AMPS in forecasting temperature, wind, precipitable water](#)
42 [vapor, cloud, radiation, and heat flux have been examined in previous studies](#) (Monaghan et al., 2005; Seefeldt et al., 2011;
43 Vázquez B and Grejner-Brzezinska, 2012; Wille et al., 2016; Listowski and Lachlan-Cope, 2017; Hines et al., 2019). To our
44 knowledge, using the AMPS to forecast Ri has not been formally validated. Thus, this study will investigate the reliability of
45 the estimated Ri of Antarctica. The atmospheric model employed for AMPS is a using AMPS forecasts. The atmospheric model
46 employed for AMPS is the Polar version of the Weather Research and Forecasting (Polar WRF) model (Powers et al., 2012).
47 Polar WRF (<http://polarmet.osu.edu/PWRF/>) has been modified for use in polar regions, for example, improving the
48 representation of heat transfer through snow and ice (Hines and Bromwich, 2008; Hines et al., 2015). The Polar WRF has been
49 used to simulate the Ri at Dome A in Antarctica, and the simulated Ri basically behaved as expected as the Ri is generally
50 large when the atmosphere is less turbulent (corresponding to the measured astronomical seeing is small; Yang et al., 2021) and
51 performed well in estimating boundary layer height when compared with other methods (Yang et al., 2022). The Polar WRF
52 model was developed and maintained by the Polar Meteorology Group of the Byrd Polar and Climate Research Center (). The
53 simulated by Polar WRF seems to be reasonable in Antarctica when compared to an atmosphere turbulence parameter . In
54 addition, various verification studies of the AMPS have been conducted . However, previous verification studies have
55 concentrated on and around the Ross Ice Shelf (or just a part of Antarctica), which needs to be extended.

56 ~~—And so far~~Presently, monitoring a wide range of atmospheric turbulence over the Antarctic continent is still
57 tremendously ~~tremendously~~ difficult at present, while the, but atmospheric Ri Richardson number (; a critical parameter judging
58 the possibility of the turbulence could be triggered) is easier to obtain, as it can be calculated by~~from~~ the routine meteorological
59 parameters (potential temperature and wind speed). However, ~~a precise evaluation of few studies have evaluated~~
60 models to forecast Ri ~~atmospheric~~ in Antarctica ~~is sketchy~~, because of limited meteorological experiments here. Nevertheless,
61 Geissler and Masciadri (2006) and Hagelin et al. (2008) ~~the analyses from used~~ the European Centre for Medium-Range Weather
62 Forecasts (ECMWF) ~~have been used analyses~~ to calculate atmospheric Ri in Antarctica. The ECMWF analyses were generated
63 from the data assimilation using observations (P. Lönnberg, 1992) ~~(the inverse of the Richardson number) by and~~, and can
64 provide initial states for numerical models (such as Polar WRF). However, their ~~researches have~~ research has some specific
65 ~~shortages~~shortcomings (or problems that need further study): (1) They did not compare Ri A direct comparison of between the
66 model estimations and from measurements has not been conducted. (2) Evaluations of forecast ability have not been done yet and
67 forecasts, while the forecast function is of great significance for practical application. (3) Variability (e.g., astronomical
68 observations, aviation safety, optical communication, etc). (2) How model error of Ri at high horizontal spatial resolution has
69 not been given, as previous studies focus depends on atmospheric conditions has not been analyzed. (3) The correlations between
70 turbulence conditions (indicated by Ri profiles at separate sites) and some large-scale phenomena or local-scale dynamics in
71 Antarctica were not fully investigated. (4) A reasonable reference standard for judging the atmospheric stability by probability of
72 triggering turbulence using the model estimations estimated Ri was not given. To fill these gaps, the scientific goals of this
73 paper are thus as follows:

74 1. To carry out a detailed comparison of ~~the potential~~ temperature and wind speed (on which Ri depends) in the atmospheric
75 column, this study extends the model evaluations ~~performed by above two sites~~ (Hagelin et al., 2008) ~~above two sites~~ to three
76 sites (McMurdo, South Pole, and Dome C) over the Antarctic continent for an entire year. The three sites are considered
77 representative, as the coast (McMurdo), flank (South Pole), and summit (Dome C) of the Antarctic continent will be compared
78 using radiosoundings and AMPS forecasts.

79 2. The radiosonde can measure meteorological parameters, which can estimate Ri . Using the AMPS-forecasted
80 meteorological parameters, ~~we one also can~~ obtain the Ri ~~AMPS forecasted~~. ~~The measured will also be calculated using the~~
81 ~~radiosounding measured meteorological parameters.~~ Then, a ~~direct~~ comparison of Ri between ~~estimated from~~ measurements
82 and forecasts can be achieved, allowing us to evaluate the reliability of AMPS-forecasted Ri ability of the AMPS to forecast in
83 giving insight into the atmospheric stability or instability, turbulence in Antarctica. In addition, we ~~investigate~~ investigated how ~~to~~
84 ~~correct the discrepancies between~~ the AMPS-forecasted ~~models~~ and ~~obtain a result closer to the~~ measurements depend on the
85 atmospheric conditions.

86 -3. ~~We also extended the analysis of done by for three sites to two~~ Two vertical cross-sections for Ri at a high horizontal
87 resolution. Finally, will be given, which may provide a better perspective on the turbulence conditions in both vertical and
88 horizontal dimensions, instead of only focusing on the vertical dimension (or atmospheric column; e.g., Hagelin et al. 2008 and
89 Geissler and Masciadri 2006). This will help to identify regions and periods that are ~~favourable~~ favorable for triggering
90 atmospheric turbulence ~~(or instability) can be identified. Thus, this study provides a better perspective on atmospheric dynamics~~
91 in Antarctica.

92 4. ~~To demonstrate the practicality of the AMPS-forecasted and provide more information about the atmospheric properties~~
93 ~~over the Antarctic continent, the planetary boundary layer (PBL) height, , forecasted by the AMPS, is used~~ Moreover, this will
94 enable us to correlate the Ri compare distribution with the AMPS-forecasted ~~—~~ ansome large-scale phenomena or local-scale
95 dynamics (katabatic winds, polar vortices, convection, gravity wave, etc.) in Antarctica, and the underlying physical processes of
96 Antarctic atmospheric turbulence will be investigated.

97 4. The Planetary Boundary Layer Height ($PBLH$, within which the atmosphere is generally turbulent) can be estimated ~~where~~
98 ~~decreases to using~~ a critical value of Ri , typically 0.25 (Holtslag et al., 1990; Pietroni et al., 2012; Petenko et al., 2019).
99 However, this critical value depends on the vertical resolution of data (Troen and Mahrt, 1986; Holtslag et al., 1990), ~~denoted by~~
100 ~~, and may be different for the AMPS grid resolution.~~ Then, ~~the~~ Ri ~~(i.e. at the AMPS-forecasted PBLH (Ri_{PBLH} PBL height))~~
101 was obtained as a reference standard for judging whether the atmosphere is likely to be laminar flow ($Ri \geq Ri_{PBLH}$ was stable) or
102 turbulent flow ($Ri < Ri_{PBLH}$) or unstable (\rightarrow) when using the AMPS-forecasted Ri .

103 In Sect. 2, we present the experimental data and atmospheric model used in this study, with an explanation of their main
104 characteristics. In Sect. 3, the Richardson number is introduced. In Sect. 4, we compare AMPS forecasts to radiosoundings and
105 ~~analyse~~ analyze the atmospheric ~~propertiesturbulence conditions~~ in Antarctica. Sect. 5 ~~discusses the relationship between the~~
106 ~~distribution and typical atmospheric features in Antaretica. Sect. 6 summarisessummarizes~~ the main ~~aspeetsfindings and primary~~
107 takeaways of this study.

108 2 Data and model

109 2.1 Radiosoundings

110 Daily radiosounding measurements at McMurdo (MM) and South Pole (SP) are available at the Antarctic Meteorological
111 Research Center (AMRC; <ftp://amrc.ssec.wisc.edu/pub>). For Dome C (DC), one can ~~access~~[obtain](#) the measurements at the
112 Antarctic Meteo-Climatological Observatory (<http://www.climantartide.it>). The altitudes of the three sites are 9 m (MM), 2839 m
113 (SP), and 3239 m (DC), where the altitudes correspond to the ~~height~~[heights](#) of the ~~radiosoundings~~[radiosondes](#) at the time of
114 launch. Their locations are shown in Fig. 1. ~~Dome A (DA) is also marked in Fig. 1, which is the highest location (4083 m) on the~~
115 ~~Antarctic plateau, and the atmospheric conditions above it will also be analyzed in this study (Sect. 4.2.3).~~ The
116 ~~radiosounding~~[radiosonde](#)-measured meteorological parameters include pressure, temperature, wind speed, and wind direction.
117 ~~This study used~~; one year ~~of data~~ (from ~~2021~~ March ~~2021~~ to ~~2022~~ February) of these meteorological parameters was used in this
118 study. Generally, the radiosonde was launched once a day at the same hour (sometimes twice a day at MM and SP). In total, 518,
119 508, and 340 profiles were available at MM, SP, and DC from March 2021 to February 2022.

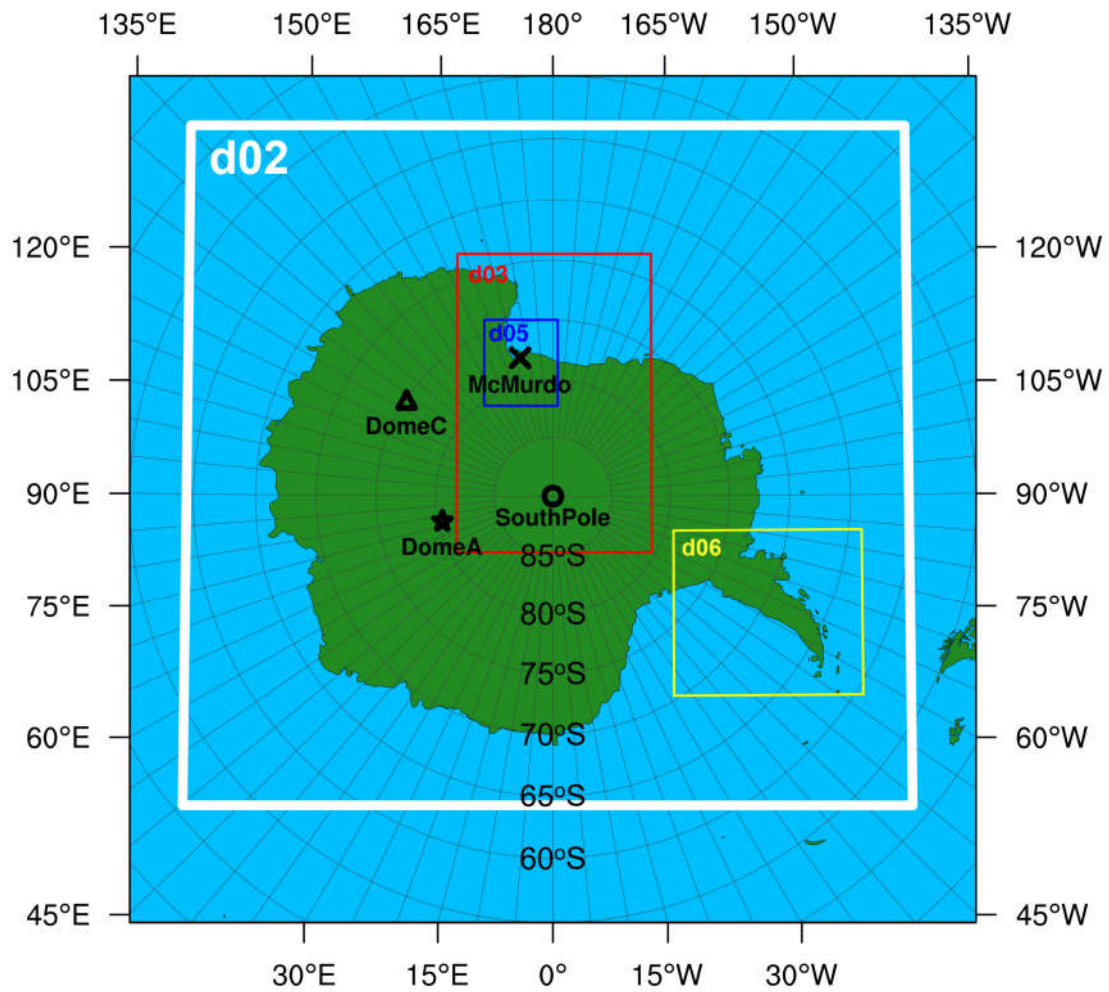
120
121 [Table 1. Main technical specifications of the radiosonde RS41.](#)

Measuring element	System resolution	System uncertainty	Data resolution^a
Temperature	0.01 °C	0.15 °C (> 100 hPa) 0.30 °C (<100 hPa)	0.1 °C
Pressure	0.01 hPa	0.5 hPa (> 100 hPa) 0.3 hPa (3-100 hPa)	0.1 hPa
Wind speed	0.1 ms⁻¹	0.15 ms⁻¹	0.1 ms⁻¹ (McMurdo) 0.1 kts (South Pole) 0.1 ms⁻¹ (Dome C)
Wind Direction	0.1 deg	2 deg	0.1 deg (McMurdo) 1 deg (South Pole) 1 deg (Dome C)

122 ^a Resolution in the files that are available for download from the Web (<ftp://amrc.ssec.wisc.edu/pub>, <http://www.climantartide.it> ~~2022~~),
123 ~~).~~

124
125 The radiosonde instrumentation used during this measurement period was the Vaisala RS41 (Technical data:
126 <https://www.vaisala.com/en/products/weather-environmental-sensors/upper-air-radiosondes-rs41>). The accuracy and uncertainty
127 of the ~~radiosonde measurements are listed in Table 1. Vaisala RS41 radiosondes have gradually replaced an older version~~
128 ~~(Vaisala RS92) starting in late 2013. These two radiosondes agree well with global average temperature differences <0.1-0.2 K~~
129 ~~in the lower stratosphere, but RS41 appears to be less sensitive than RS92 to changes in solar elevation angle (Sun et al., 2019).~~
130 ~~Besides, RS41 (1-1.5% dry bias) has better performance than RS92 (3-4% dry bias) relating to the infrared atmospheric sounding~~
131 ~~interferometer as a practical reference (Sun et al., 2021). Near-global radiosonde measurements have been used to calculate the~~
132 ~~Richardson number and derive the boundary layer height, which is positively correlated with the results of four reanalysis~~
133 ~~products (Guo et al., 2021)~~~~radiosounding measurements are listed in Table 1. The balloon scans the atmosphere between the~~
134 ~~ground and an altitude of 10–25 km (low in winter and high in summer); the vertical resolution is approximately 5 m, depending~~
135 ~~on the ascent speed.~~

136 In Antarctica, the radiosondes measure the atmosphere between the ground and an altitude of 10–25 km (low in winter and
137 high in summer) with a typical ascent rate of 5 m s⁻¹, and a logging frequency of 1 Hz; then the vertical resolution is
138 approximately 5 m.
139



140
141 Figure 1. The five two-way interactive horizontal grids (d01, d02, d03, d05, and d06; information online at
142 https://www2.mmm.ucar.edu/rt/amps/information/configuration/maps_2017101012/maps.html

143 Table 1. Main technical specifications of the radiosonde instrumentation.

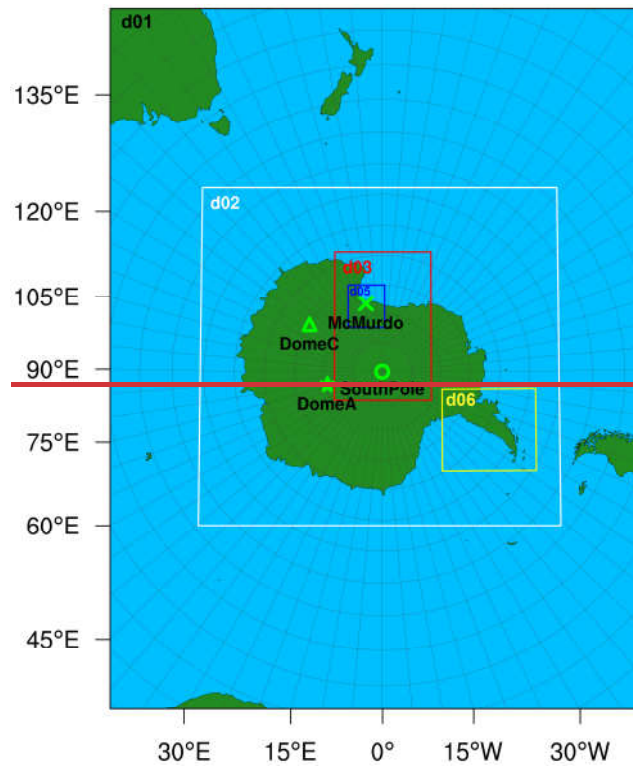
144 ^aResolution in the files that are available for download from the Web.

145) used in the AMPS configuration. The locations of McMurdo (78°S, 167°E), South Pole (90°S, ...°E), Dome C (75°S, 123°E),
146 and Dome A (80°S, 78°E) are shown by the cross, circle, triangle, and star, respectively. This study used grid 2 fields (d02; the
147 white rectangle) that covered the entire Antarctic continent.

148 2.2 AMPS

149 The AMPS can forecast meteorological parameters in four-dimensional space-time in Antarctica, which can be used for
150 comparison with the radiosonde measurements—of radiosoundings. The AMPS grid system consisted of a series of nested
151 domains with 60 vertical levels. This study used grid 2 fields (d02; 8 km horizontal resolution) that covered the entire Antarctic
152 continent (similar to Hines et al., 2019), as shown by the white square in Fig. 1. However, the contributions from the nested grid

153 with higher horizontal resolution (d03: 2.67 km; d05: 0.89 km; d06: 2.67 km) are not entirely lost, as the AMPS used a two-way
 154 nested run and the nest (e.g. d03) feeds its calculation back to the coarser domain (e.g. d02). The original WRF output files for
 155 each AMPS grid were saved in a rolling archive (one can find how to access/download the original WRF output files of AMPS
 156 for each grid at https://www2.mmm.ucar.edu/rt/amps/information/amps_esg_data_info.html). This study used the AMPS outputs
 157 (in original WRF format) for forecast hours 12-21 at three h intervals from the daily AMPS forecasts that began at 12:00 UTC.
 158 Parish and Waight (1987) Thus, our AMPS showed large adjustments to the boundary layer fields had a spin-up time of 12 h.
 159 Finally, the AMPS forecasts for above an ice sheet before the same period (from March 2021 numerical model began to stabilize
 160 after about 10 h. Then, some studies (Hines and Bromwich, 2008; Hines et al., 2019) have discarded the first 12 h forecasts (so-
 161 called 12 h spin-up time). Thus, in this study, only the 12-33-h forecasts from each of the AMPS simulations are combined into a
 162 year-long (2011 March to 2022 February-2022), as the used radiosounding measurements, were downloaded for analysis) output
 163 field at 3-h intervals.



164
 165 Figure 1. The five two way interactive horizontal grids (d01, d02, d03, d05, and d06; information online at) used in the AMPS
 166 configuration. The locations of McMurdo (78°S, 167°E), South Pole (90°S, ...°E), Dome C (75°S, 123°E), and Dome A (80°S,
 167 78°E) are shown by the cross, circle, triangle, and star, respectively.

168 3 Theory of Richardson number

169 The stability of the atmosphere can be estimated using the Richardson number (Ri) is generally defined as (Richardson and
 170 Shaw, 1920; Chan, 2008):

$$171 Ri = \frac{g}{\theta} \frac{\partial \theta / \partial z}{[\partial u / \partial z]^2 + [\partial v / \partial z]^2} \tau \quad (1)$$

172 ~~where~~Where g is the gravitational acceleration (9.8 m s^{-2}), $\theta = T[1000/P]^{0.286}$ is the potential temperature (K), T and P ~~is are~~
173 the temperature (K) and pressure (hPa) of air, respectively. ~~As for wind shear term,~~ u and v are the east-west and north-south
174 components of the wind (m s^{-1}). ~~z is the height (m) above the ground. To calculate Ri (m^{-2}), a centered finite difference~~
175 ~~operation was used to estimate the gradient in Eq. (1).~~

176 The ~~production/development~~ of atmospheric turbulence (~~or unstable atmosphere~~) was shown to be tightly correlated with the
177 Ri . It can, therefore, be an essential indicator of the turbulence characteristics in the atmosphere (Ma et al., 2020; Han et al.,
178 2021; Yang et al., 2021). Atmospheric conditions are ~~favourable~~favorable for the occurrence of turbulence when Ri is less than
179 a critical value (Ri_c ; ~~critical Richardson number~~), and Ri_c is typically chosen as 0.25; ~~however,~~ However, a larger Ri_c should
180 be used in a large-scale model (e.g., 0.5 has been employed by Troen and Mahrt, 1986).

181 ~~In the results of this study, the logarithm of Ri , $\log_{10}(Ri)$, is presented instead of Ri itself, because Ri in this study, the~~
182 ~~inverse of the Richardson number ($-$) was used to provide better evidence of atmospheric stability, as in and . The larger the $-$, the~~
183 ~~higher the probability of triggering turbulence in the atmosphere.~~
184 ~~can vary by two or more orders of magnitude in the atmosphere.~~

185 4 Results and discussion

186 4.1 Temperature/Potential temperature and wind speed

187 The AMPS forecasts are compared to radiosoundings from MM, SP, and DC to investigate the reliability of the AMPS
188 forecasts over the Antarctic continent. ~~The AMPS forecasts and radiosoundings used for this comparison were obtained from~~
189 ~~March 2021 to February 2022.~~ To offer a more convincing result, data corresponding to the altitude at which radiosoundings
190 reached less than five times a ~~month~~season were discarded. In addition, the extracted AMPS forecasts used for comparison were
191 from the nearest grid to the three sites, and the time difference between radiosoundings and AMPS forecasts ~~and radiosoundings~~
192 larger than 1.5 hours was ~~also~~ not used for comparison. Moreover, both radiosoundings and AMPS forecasts ~~and radiosoundings~~
193 were linearly interpolated to the same height series (average annual altitude of the AMPS vertical grid, ~~because, where~~ the
194 altitude of the AMPS grid may vary during the simulation as the AMPS uses the WRF hybrid vertical coordinate, information
195 online at https://www2.mmm.ucar.edu/wrf/users/docs/user_guide_v4/WRFUsersGuide.pdf) for each site. On the other hand, it
196 should be noted that the near-surface radiosonde measurements could be less reliable, as it was just released from the operator's
197 hand (or some machine). Hagelin et al. (2008) conclude that the radiosoundings are $\sim 1 \text{ K}$ colder than the Automatic Weather
198 Station at Dome C and $\sim 2 \text{ K}$ at the South Pole. In this study, the radiosonde measurements in the first $\sim 10 \text{ m}$ above the ground
199 are not used. This is also because the first AMPS grid is $\sim 10 \text{ m}$ above the ground.

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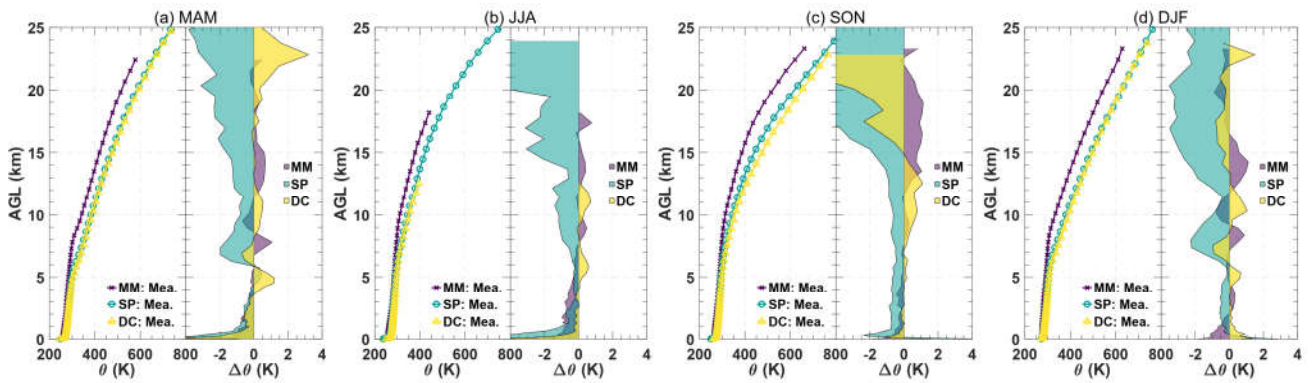


Figure 2. The monthly medians seasonal median of the potential temperature (θ differences (Fig. 2) estimated by the radiosonde measurements (solid lines) and potential temperature difference ($\Delta\theta$ wind speed (Fig. 3) between) calculated by the AMPS forecasts and radiosoundings are presented. The panels in minus the radiosonde measurements, i.e. $\Delta\theta = \theta_{AMPS} - \theta_{Mea.}$ same row of Figs. 2 and 3 correspond to the same season; the first row: Aut. (Autumn (filled areas). Fall: March, April, and May), second row: Win. (Winter (MAM); winter: June, July, and August), third row: Spr. (Spring (JJA); spring: September, October, and November), fourth row: Sum. (Summer (SON); summer: December, January, and February (DJF)).

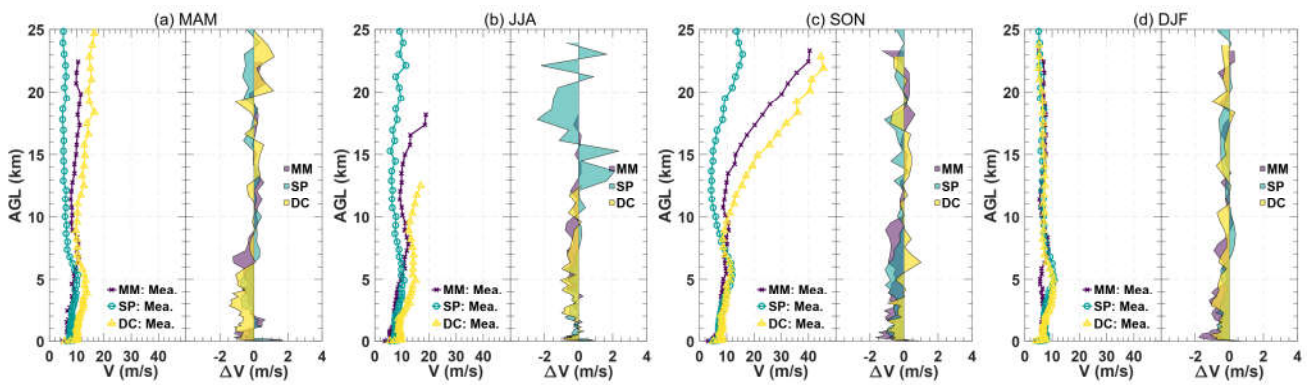
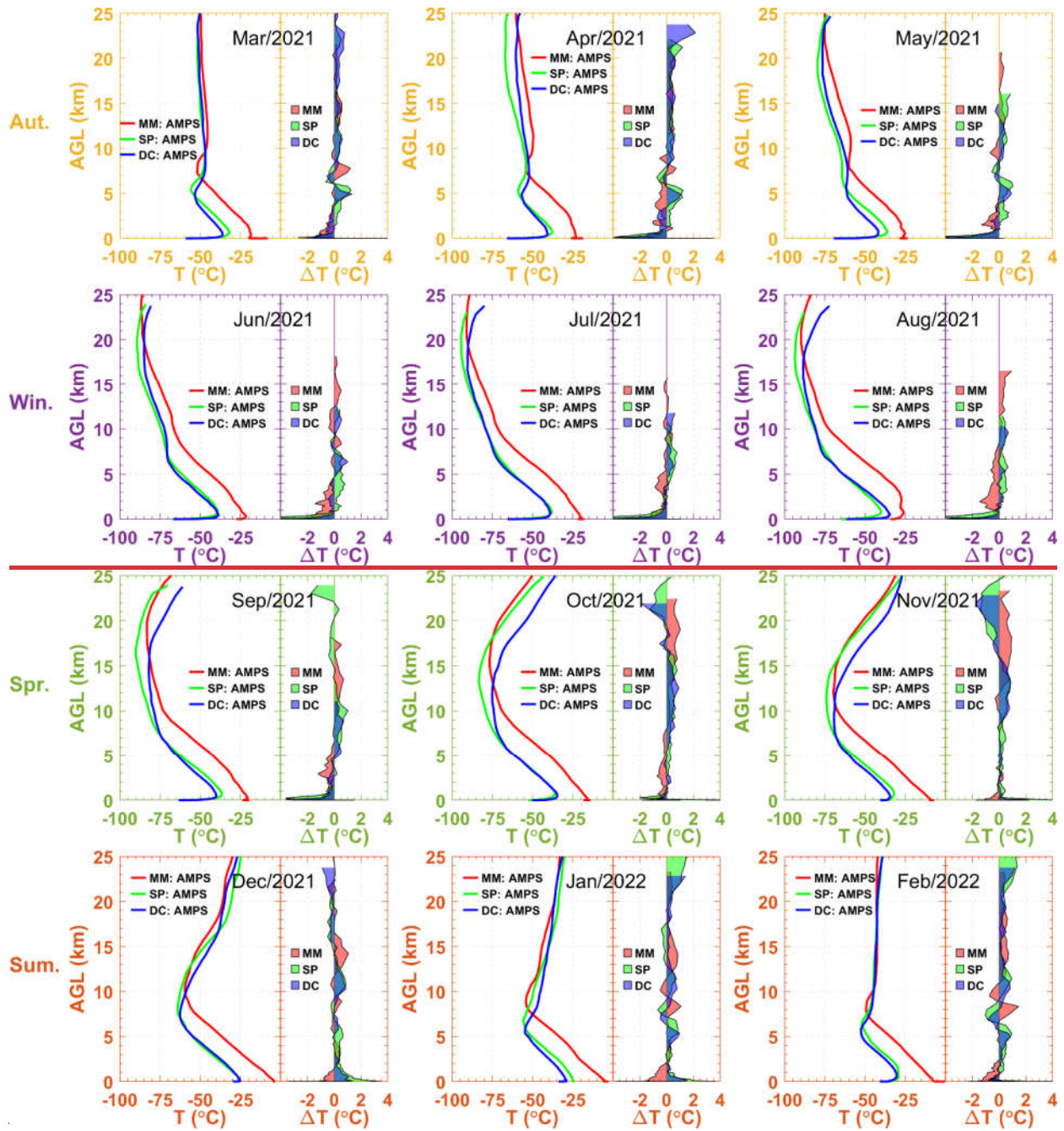


Figure 3. As in Fig 2, but for wind speed V ($=\sqrt{u^2 + v^2}$), and $\Delta V = V_{AMPS} - V_{Mea.}$.

The seasonal median difference of potential temperature (see the filled areas in Fig. 2) and wind speed (see the filled areas in Fig. 3) between radiosoundings and the AMPS forecasts are presented. The missing value of the median difference in the upper part of the atmosphere during JJA indicates that the radiosoundings or radiosonde balloon does not reach as high an AGL (Above Ground Level) in winter as they do in summer, probably because the elastic material of the balloons is more fragile in cold seasons and easier to explode (Hagelin et al., 2008). The lack of measurements may also be attributable to some large values of the median difference in the top layer of the profile shown in Figs. 2 and 3, as the AMPS requires the assimilation data from measurements to initialize its numerical model, and the lack of measurements makes it more difficult for AMPS to simulate atmospheric changes that are close to reality.



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Figure 2. The monthly median for temperature forecasted by the AMPS (solid lines) and temperature difference calculated by the AMPS forecasts minus the radiosoundings measurements, i.e. (filled areas).

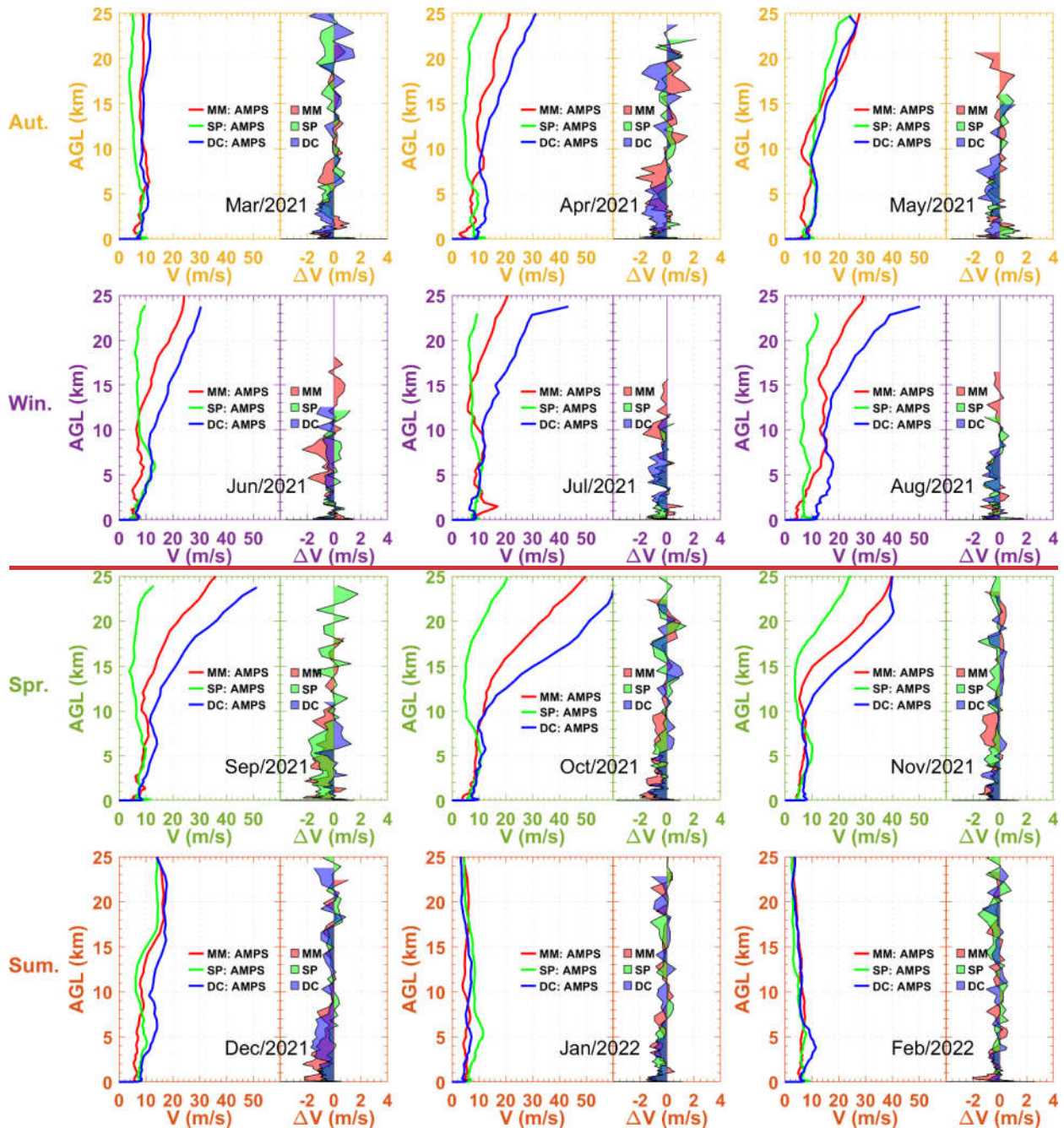


Figure 3. As in Fig. 2, but for wind speed (V), and ΔV .

Fig. 2 shows that the forecasted temperature profiles of SP and DC in the first 5 km are similar, 15–20 K colder than the MM. The median difference in the temperature median difference for θ is of the order of 1 K in the high part of first 5 km (except for the atmosphere (see filled areas in Fig. 2). However, layer in proximity to the ground, the median difference becomes more significant, especially for winter). Above 5 km, the AMPS has obviously underestimated the θ at more than 4 K near MM, while the ground forecasts at SP-

Fig. 3 and DC are closer to measurements. Fig. 3 shows that the forecasted measured wind speed profiles of MM and DC above 10 km at MM and DC are stronger during spring, indicating the occurrence of the Antarctic polar vortex (Boville et al., 1988). However, the change in wind speed above SP is not that obvious, because the Antarctic vortex is roughly pole-centred centered

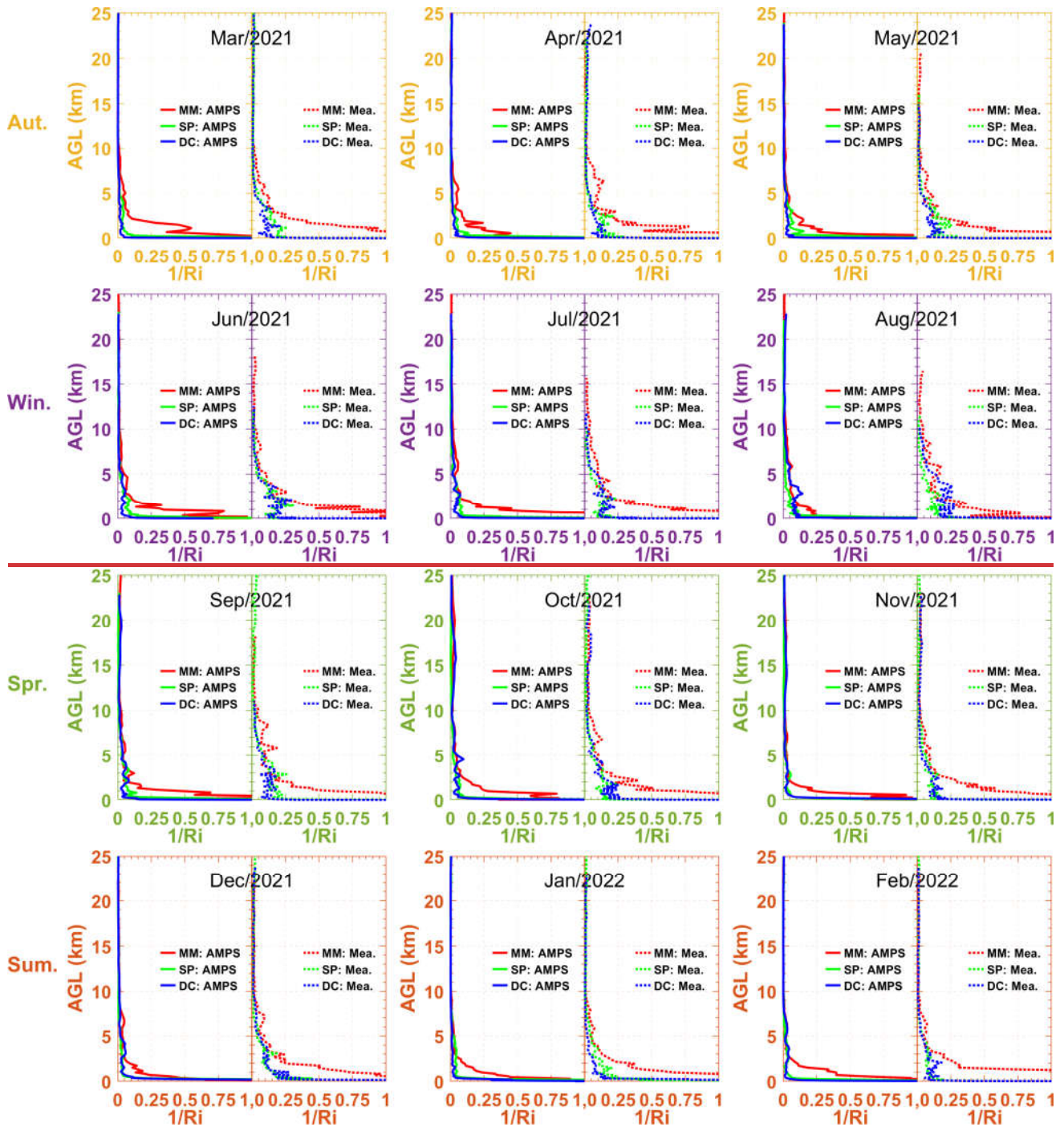
234 (Karpetchko et al., 2005). From the filled areas in Fig. 3, the AMPS forecasts appear consistent with the measurements, as the
235 median difference in wind speed is generally $\sim 1 \text{ m s}^{-1}$ and has barely exceeded 2 m s^{-1} , whether the wind is strong or weak. In the
236 first 10 km, most ΔV at the three sites are less than 0, suggesting that the AMPS underestimated the wind speed. Table 2 shows
237 the statistical evaluations of θ and V forecasted by the AMPS. It seems the AMPS can well capture the trend of θ and V as the
238 correlation coefficient (R_y . In addition, slightly smaller median differences were displayed around 10 km AGL (e.g. September
239 and December, shown in Fig. 3), and the forecasted wind speed also became slightly smaller at that height) are all larger than
240 0.84.

241 ~~In summary, the AMPS almost well forecasted the temperature and wind speed. Nevertheless, it should be noted that both the~~
242 ~~measurement accuracy and position of radiosoundings may be affected by their flight, and the difference may be smaller when~~
243 ~~compared with a fixed measuring instrument, such as the Automatic Weather Station . In other words, the functional~~
244 ~~performance of the AMPS may be better than the results shown in Figs. 2 and 3.~~

245 4.2 Richardson number

246 4.2.1 ~~Monthly statistical~~Statistical analysis

247 To evaluate the performance of ~~the~~ AMPS in forecasting ~~atmospheric stability or instability~~the possibility of triggering
248 turbulence over the Antarctic continent, the Ri ~~forecasted by estimations between~~ AMPS forecasts and ~~measured by~~
249 radiosoundings will be ~~provided~~compared. The calculated value of Ri depends on the vertical resolution of meteorological
250 parameters (Troen and Mahrt, 1986; Holtslag et al., 1990). Thus, the meteorological parameters from the radiosoundings and
251 AMPS forecasts ~~and radiosoundings~~ were interpolated into the same height series (as mentioned in Sect. 4.1) to calculate Ri .
252 Where $\partial\theta/\partial z$ and $(\partial u/\partial z)^2 + (\partial v/\partial z)^2$.



253 Figure 4. The monthly median $1/Ri$ forecasted by the AMPS (solid lines) and measured by radiosoundings
 254 (dashed lines).

255 The NCL (NCAR Command Language) can compute $1/Ri$ from the original WRF outputs, using an NCL function called
 256 “`rigrad_bruntv_atm`” (). The NCL (see Eq. (1)) were both computed using a centered finite difference operation, as we found
 257 that centered difference performed better than forward difference and backward difference (not shown), i.e., better consistency of
 258 $1/Ri$ (not shown) agrees very well with the results obtained using interpolation (i.e. interpolating the hybrid vertical coordinate to
 259 the same height series for each site). between the radiosoundings and AMPS forecasts can be achieved using centered difference.
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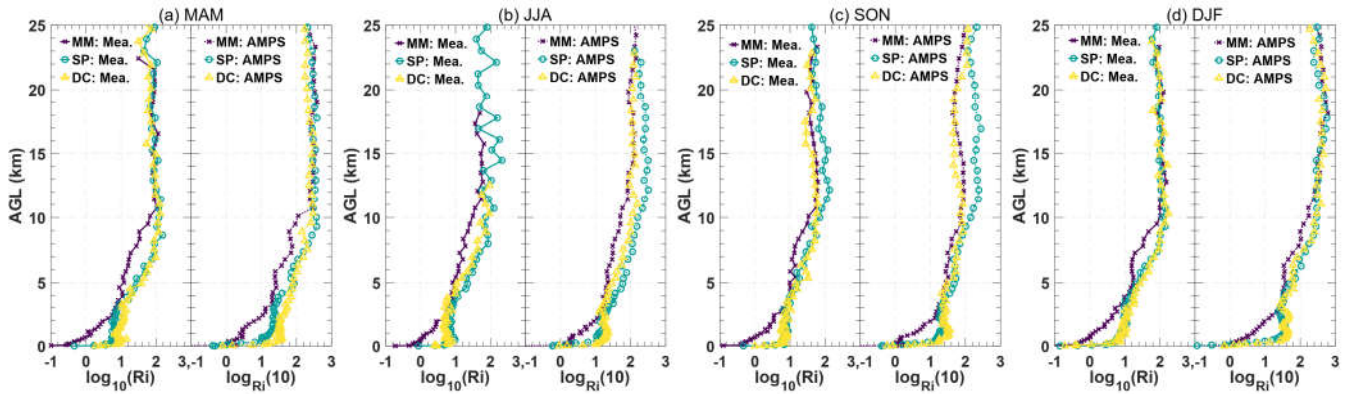


Figure 4. The seasonal median of $\log_{10}(Ri)$ monthly estimated by the radiosonde measurements (solid lines) and the AMPS forecasts (dashed lines).

The seasonal median profiles of $\log_{10}(Ri)$ from the AMPS and radiosoundings and AMPS forecasts are shown in Fig. 4. However, the median differences are not similar to those of presented like the θ previous temperature and V wind speed. This is because, the Ri value of can oscillate vary massively (by two or more orders of magnitude) in the atmosphere, and a precise quantification seems less plausible. Considering this, we initially intended to examine whether AMPS can reconstruct an accurate shape of $\log_{10}(Ri)$. Nevertheless, the AMPS-forecasted profile (while median difference is not suitable for this purpose), and the results from radiosoundings and AMPS forecasts are both presented. Nevertheless, the model bias is by all means of great significance, and it will be discussed later (see Table 2, and Fig. 6). In Fig. 4, one can see that the AMPS-forecasted Ri can identify that the atmosphere above MM tends to be more unstable (or turbulent, Ri is larger) than SP and DC, per the measurements from radiosoundings. In addition, in the vertical height direction, the AMPS forecasts can roughly capture the height that can easily trigger turbulence. On the one hand, for example, one can observe that the Ri from radiosoundings and AMPS forecasts and radiosoundings both show large values very close to the ground at Dome C and the South Pole, which is per the fact that strong atmospheric turbulence is concentrated within the surface layer above the high plateau (Marks et al., 1999; Agabi et al., 2006). A very calm atmosphere (Ri is small) at high altitudes was also consistent with the results given by Travouillon et al. (2003), Aristidi et al. (2005), Trinet et al. (2008) and Vernin et al. (2009). On the other hand, the AMPS can reconstruct the near-ground “convex-concave-convex” (hereafter “C-C-C”) shaped $\log_{10}(Ri)$ profiles indicated by ~~and~~. On the other hand, in a more detailed comparison, for example, in March 2021, both forecasts and the radiosonde measurements (see more details in Fig. 5). In terms of time, the AMPS can forecast that the free-atmosphere Ri show that profiles are inclined to the y-axis at ~ 10 km AGL, and in December 2021, the AMPS forecasts are able to capture a small bump of that occurred at around 7 km AGL, decreased during spring (SON), this decrease is obvious for MM and DC (where the wind speed are significantly stronger during SON, as in Fig. 3).

Some researchers are afraid to conduct quantitative analysis for the model-estimated Ri from the numerical models was generally missed as it always varies dramatically (e.g., Hagelin et al., 2008, who focused on the qualitative analysis). Nevertheless, quantitative analysis has been tried in this study since that can give a precise evaluation of the forecast ability of AMPS. Then, the R_{xy} correlation coefficient (ρ), mean bias ($Bias$; AMPS-radiosoundings/radiosonde), and root mean square error ($RMSE$) are calculated. Here, using the combined data of all profiles for each season. Where the time difference between

radiosoundings and AMPS forecasts and radiosoundings, with lesser was limited to less than 1.5 hours, is used. This is the same as the meteorological parameters mentioned in Sect. 4.1. Besides, all profile data meeting the time constraints are also limited to a value range of (0, 1), which is the same as the x-axis in Fig. 4. Finally, the monthlyseasonal values of the three statistical operators are obtainedcalculated, as shownlisted in Fig. 5Table 2. However, we want to emphasiseemphasize that one should focus on the value of R_{xy} that reflects the tendency, instead of *Bias* and *RMSE*, as a precise quantification remains in doubt (Hagelin et al., 2008). The mean values of R_{xy} , the combined 12 monthly in Fig. 5a for MM, SP, and DC, are 0.71, 0.66, and 0.68, respectively. This suggests that the AMPS forecasted can identify the main characteristics of atmospheric turbulence over the Antarctic continent in terms of space and time. In terms of time, one can observe larger during warm seasons (e.g. Autumn from March 2021 to May 2021 and summer from December 2021 to February 2022). In sum, the AMPS seems to perform better when more measurements can be obtained. As we know, more field experiments could be conducted at the coast (MM) in the space dimension (one can reach the coast easier than the internal Antarctic Plateau) and during warm seasons (e.g., summer) in the time dimension (the cold is much harder to endure in winter than it is in summer).

Fig. 5b shows that is underestimated every season, and the mean values of the 12 monthly for MM, SP, and DC over four seasons are 0.71, 0.59, and 0.53, respectively. The highest R_{xy} is at MM for DJF (0.77) and the lowest is at DC for JJA (0.45). We found that these two cases correspond to the most unstable and stable atmospheric conditions, their median $[\theta_{1000m} - \theta_{0m}]/[1000\text{ m} - 0\text{ m}]$ are -0.097, -0.031, and -0.071 equal to 0.0038 and 0.0721, respectively. This suggests that the AMPS can better capture the trend of $\log_{10}(Ri)$ This may be because the model results were generally smoother than the measurements. The atmosphere is favourable for the occurrence of turbulence (is large) under rapidly changing meteorological parameters. Fig. 5c shows that at MM is the largest, which may be due to the at a relatively unstable atmosphere above it and could fluctuate massively.

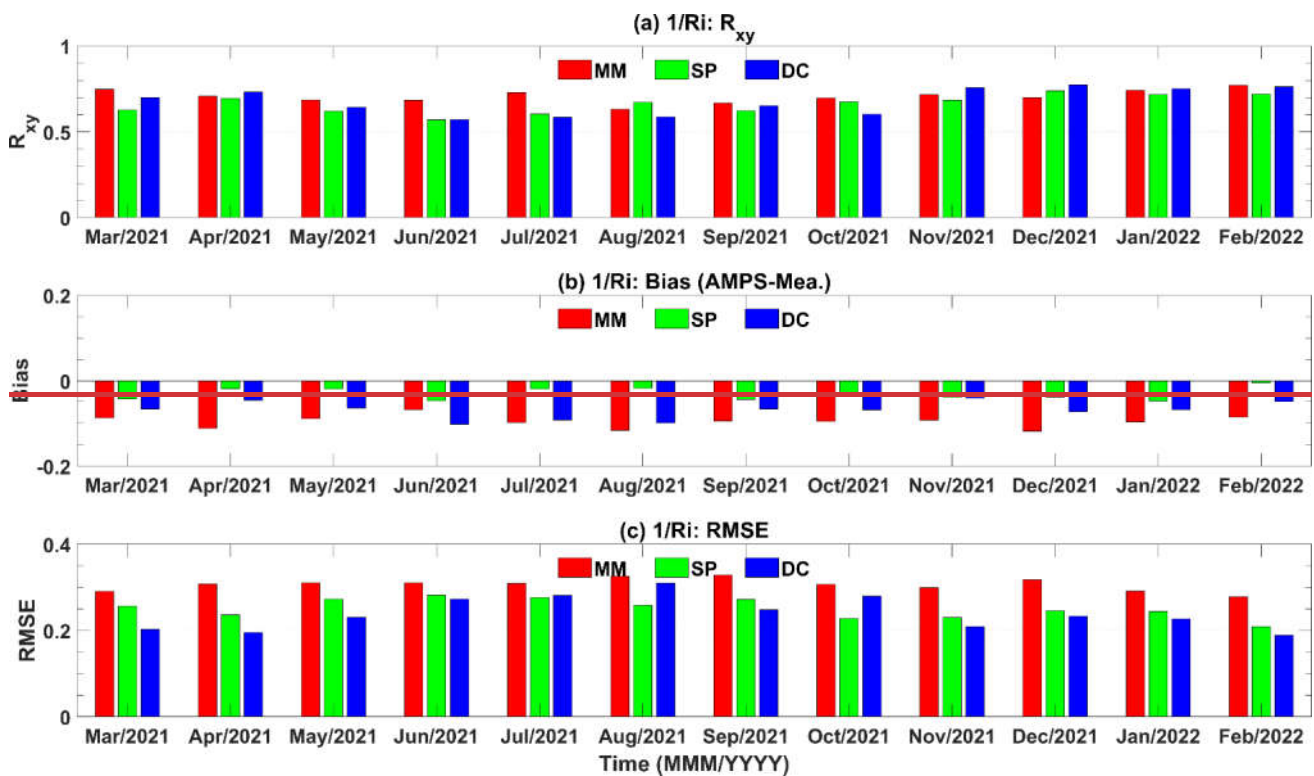


Figure 5. The monthly statistics for three sites.

4.2.2 Vertical distribution

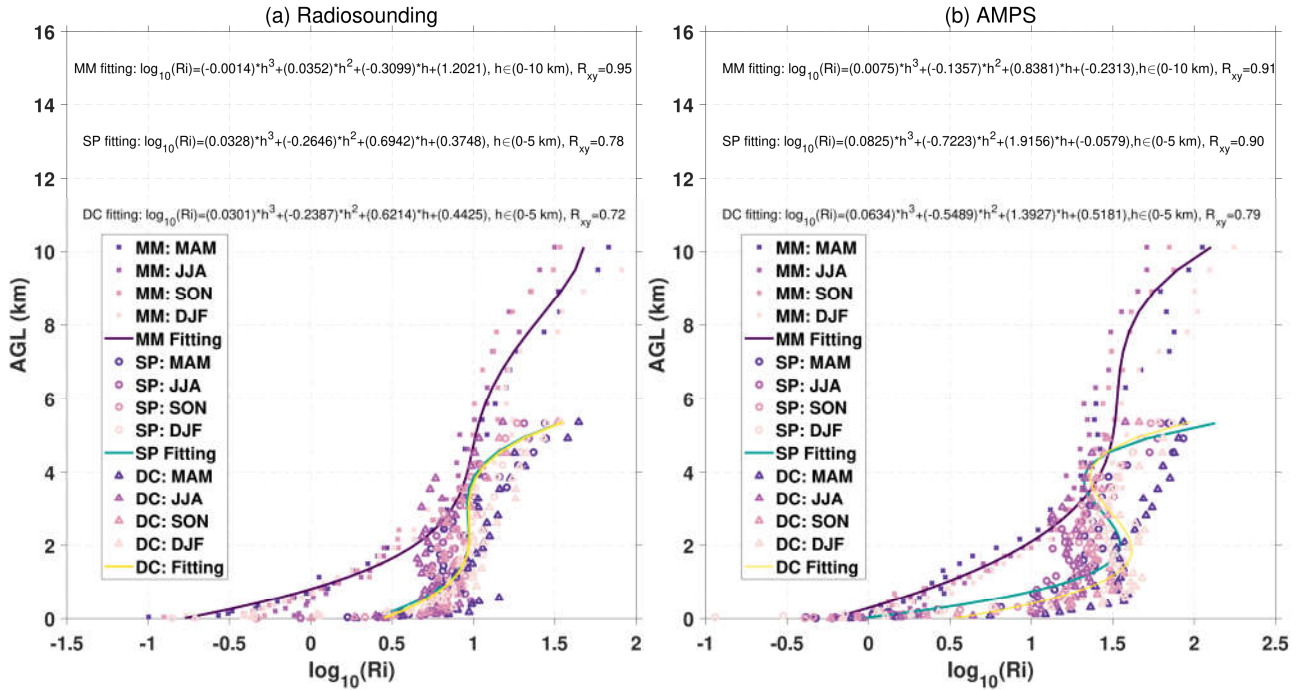
To further investigate the difference between the radiosounding and model results, the vertical distributions of the whole year. However, the *Bias* at the three sites are shown in Fig. 6. The reason is the largest (0.47) in the most unstable case. This is because the AMPS overestimated the potential temperature gradient under an unstable atmosphere (see Fig. 6a, which will be discussed later). For the stable atmosphere, the lowest R_{xy} for $\log_{10}(Ri)$ using (not, as presented in Sect. 4.2.1) is seems to be consistent with the fact that model errors increase with increasing stability (Nigro et al., 2017) we find that the distribution will be more linear at the same height. Fig. 6a shows that the data points are distributed mainly over the diagonal line between the measurements and the forecasts, this suggests the AMPS can make a better forecast at McMurdo than at the South Pole (Fig. 6b) and Dome C (Fig. 6c).

Table 2 also shows an interesting result: the R_{xy} of $\log_{10}(Ri)$ is higher when the *RMSE* of θ and V are smaller. Moreover, Hines et al. (2019) showed that using the Morrison microphysics scheme in the numerical model resulted in a smaller *RMSE* for temperature and wind, than the default scheme (WSM5C) in AMPS. Therefore, we may conclude that replacing WSM5C with Morrison could improve the AMPS-forecasted $\log_{10}(Ri)$. In other words, using Morrison may lead to higher R_{xy} for $\log_{10}(Ri)$, as it simulates dynamic stability with less variability (the *RMSE* for temperature and wind could be smaller). On the other hand, larger *RMSE* for θ and V are mainly found during cold months (JJA, SON), indicating that winter dynamic stability is more variable (similar to Bromwich et al., 2013).

Table 2 summarises that the $\log_{10}(Ri)$ was overestimated by the AMPS at each site for every season (all *Bias* are positive). This may be due to some local-scale dynamics not being represented properly (see Fig. 6, which will be discussed later). From another perspective, the model results were generally smoother than the measurements, and the atmosphere is less favorable for the occurrence of turbulence under slowly changing meteorological parameters, then the AMPS-forecasted Ri could be larger.

Table 2. Statistical evaluations of the potential temperature (θ), wind speed (V), and logarithmic Richardson number ($\log_{10}(Ri)$) forecasted by the AMPS when compared with the results from radiosonde measurements.

Season	McMurdo				South Pole				Dome C			
	MAM	JJA	SON	DJF	MAM	JJA	SON	DJF	MAM	JJA	SON	DJF
θ : R_{xy}	0.99	0.99	0.99	0.99	0.99	0.99	0.99	0.99	0.99	0.99	0.99	0.99
θ : <i>Bias</i>	-0.32	-0.61	-0.10	-0.23	-1.30	-1.41	-1.45	-0.94	-0.74	-0.58	-0.25	0.19
θ : <i>RMSE</i>	1.82	1.91	1.78	1.56	2.65	2.88	3.48	2.43	4.28	2.22	2.38	1.76
V : R_{xy}	0.85	0.89	0.95	0.90	0.86	0.84	0.89	0.90	0.92	0.92	0.97	0.95
V : <i>Bias</i>	-0.25	-0.25	-0.59	-0.67	-0.16	-0.29	-0.64	-0.40	-0.62	-0.43	-0.24	-0.50
V : <i>RMSE</i>	3.16	3.63	3.23	2.81	2.52	2.90	2.69	2.37	2.50	2.94	2.69	1.89
$\log_{10}(Ri)$: R_{xy}	0.75	0.65	0.68	0.77	0.61	0.50	0.56	0.70	0.51	0.45	0.50	0.66
$\log_{10}(Ri)$: <i>Bias</i>	0.41	0.32	0.33	0.47	0.36	0.23	0.29	0.35	0.45	0.39	0.30	0.46
$\log_{10}(Ri)$: <i>RMSE</i>	0.90	0.86	0.88	0.93	0.90	0.84	0.85	0.90	0.93	0.86	0.81	0.91



339

340 Figure 5. The polynomial curve fitting of near-ground median profiles of $\log_{10}(Ri)$ from Fig. 4; (a) the $\log_{10}(Ri)$ was
 341 estimated by the radiosonde measurements, (b) the $\log_{10}(Ri)$ was estimated by the AMPS forecasts.

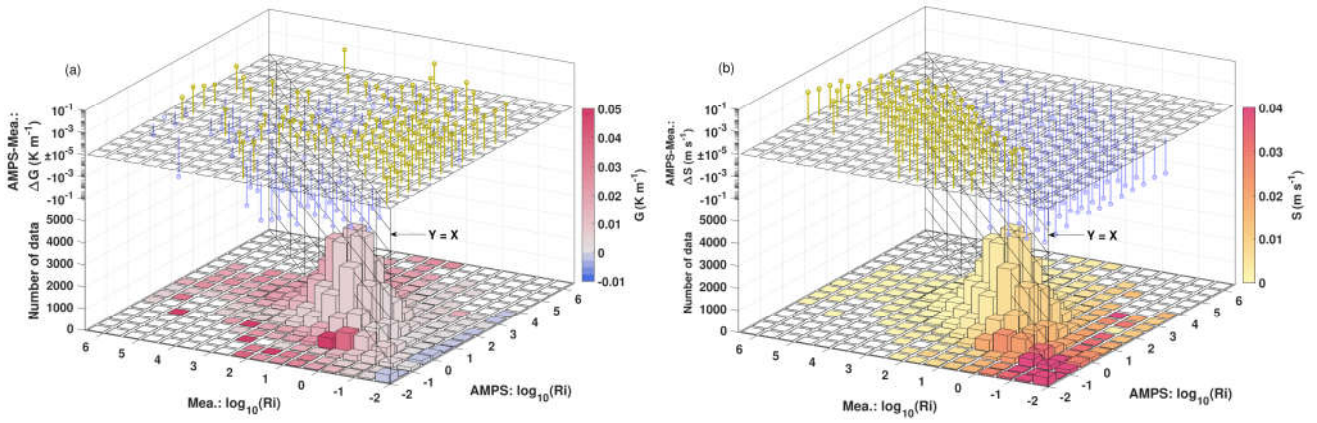
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343 The near-ground atmosphere in Antarctica is an important turbulence source, then an analytical function for $\log_{10}(Ri)$ profiles
 344 near the ground was fitted to better contextualize the results (as shown in Fig. 5). Fig. 5 shows that the near-ground $\log_{10}(Ri)$
 345 profiles are the “C-C-C” shape. The “concave” structure in the “C-C-C” shape could be attributed by the near-ground jet stream
 346 (Mihalikova et al., 2012). A cubic polynomial function was used (see the upper part of the plots in Fig. 5) instead of a
 347 logarithmic function, because the “C-C-C” shape seems hard for logarithmic function fitting. Moreover, each fitted curve used
 348 all four-seasons data points in Fig. 4, as the seasonal variation are not too significant. Nevertheless, one can see more details
 349 about the temporal variation of $\log_{10}(Ri)$ near the ground in Sect. 4.2.3.

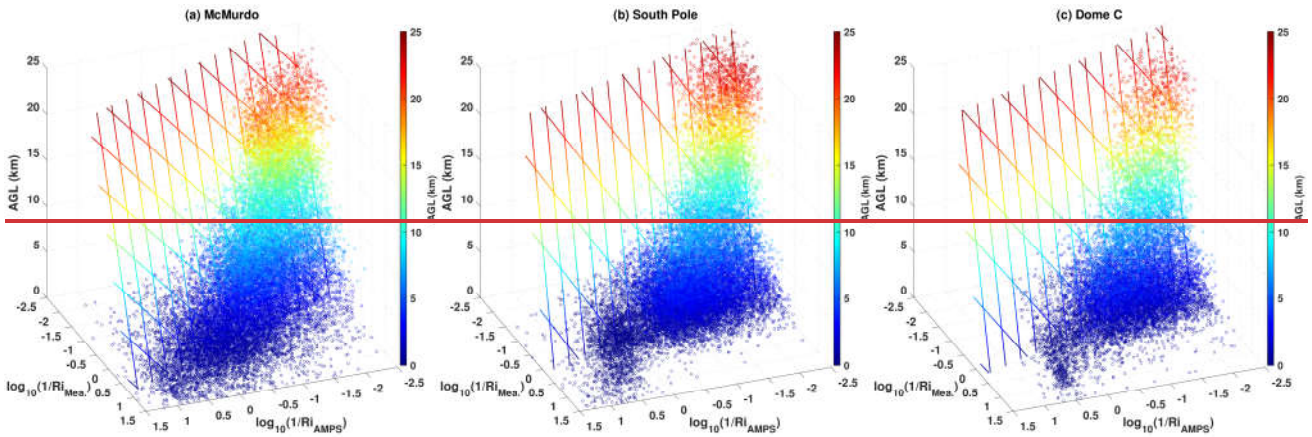
350 Fig. 6 shows the AMPS performance under different potential temperature gradient ($G = \partial\theta/\partial z$) and wind shear ($S = [(\partial u/\partial z)^2 + (\partial v/\partial z)^2]^{1/2}$).

351 The statistical results presented in Fig. 6 were counted based on all the collected data points at the
 352 three sites (MM, SP, and DC) for an entire year. One can see that the Ri was overestimated by the AMPS at an unstable
 353 atmosphere (see light blue bin in Fig. 6a), where the AMPS has overestimated the potential temperature gradient (i.e. $\Delta G > 0$).
 354 But for strong temperature inversion (see dark red bin in Fig. 6a), the AMPS has underestimated the G and Ri . As for strong
 355 wind shear conditions (see dark red bin in Fig. 6b), when the Ri is small (basically corresponding to a near-surface layer with a
 356 high probability of triggering strong turbulence, as in Fig. 4), the AMPS has underestimated the intensity of wind shear ($\Delta S < 0$
 357). This may be caused by the AMPS has underestimated the wind speed near the ground (as in Fig. 3). In sum, if the model aims
 358 for a more accurate forecast of Ri , the biases under these atmospheric conditions need to be corrected.

359



360
 361 Figure 6. Performance of the AMPS under different atmospheric conditions. (a) and (b) are respectively the case of potential
 362 temperature gradient ($G = \partial\theta/\partial z$) and wind shear ($S = [(\partial u/\partial z)^2 + (\partial v/\partial z)^2]^{1/2}$). G (color of the bin in (a)), S (color of the bin
 363 in (b)), ΔG (stem above the bin in (a)) and ΔS (stem above the bin in (b)) are presented using the median value for each
 364 0.5×0.5 bin of $\log_{10}(Ri)$



365
 366 Figure 6. the vertical distributions of whole year from the AMPS and the measurements at McMurdo (a), the South Pole (b), and
 367 Dome C (c). The coloured grid lines represent the fitted planes using the presented data points.

368
 369 We assume that the difference between the measurements and model results was correlated with height (or AGL, here
 370 indicated by), and a simple model to improve the AMPS forecasted by considering (km) is given by the following linear
 371 equation:

372
 373 The three undetermined coefficients (, , and) can be determined using a linear fitting by the measurements, (i.e. will be
 374 replaced with for fitting).

375
 376 Table 2. The fitted model for improving the AMPS forecasted.

Site	Model: $\log_{10}(1/Ri_{\text{ImpAMPS}}) = C_1 + C_2 \log_{10}(1/Ri_{\text{AMPS}}) + C_3 H$	R_{xy} (Before) ^a	R_{xy} (After) ^b
McMurdo	$\log_{10}(1/Ri_{\text{ImpAMPS}}) = 0.1222 + 0.4826 \log_{10}(1/Ri_{\text{AMPS}}) - 0.0660 H$	0.7215	0.7691

South Pole	$\log_{10}(1/Ri_{\text{ImpAMPS}}) = 0.4414 + 0.3827 \log_{10}(1/Ri_{\text{AMPS}}) - 0.0387 H$	0.6009	0.6391
Dome C	$\log_{10}(1/Ri_{\text{ImpAMPS}}) = 0.3596 + 0.3977 \log_{10}(1/Ri_{\text{AMPS}}) - 0.0457 H$	0.4730	0.5235

^aCorrelation coefficient between $\log_{10}(1/Ri_{\text{ImpAMPS}})$ and $\log_{10}(1/Ri_{\text{AMPS}})$ (i.e. before using the fitted model).

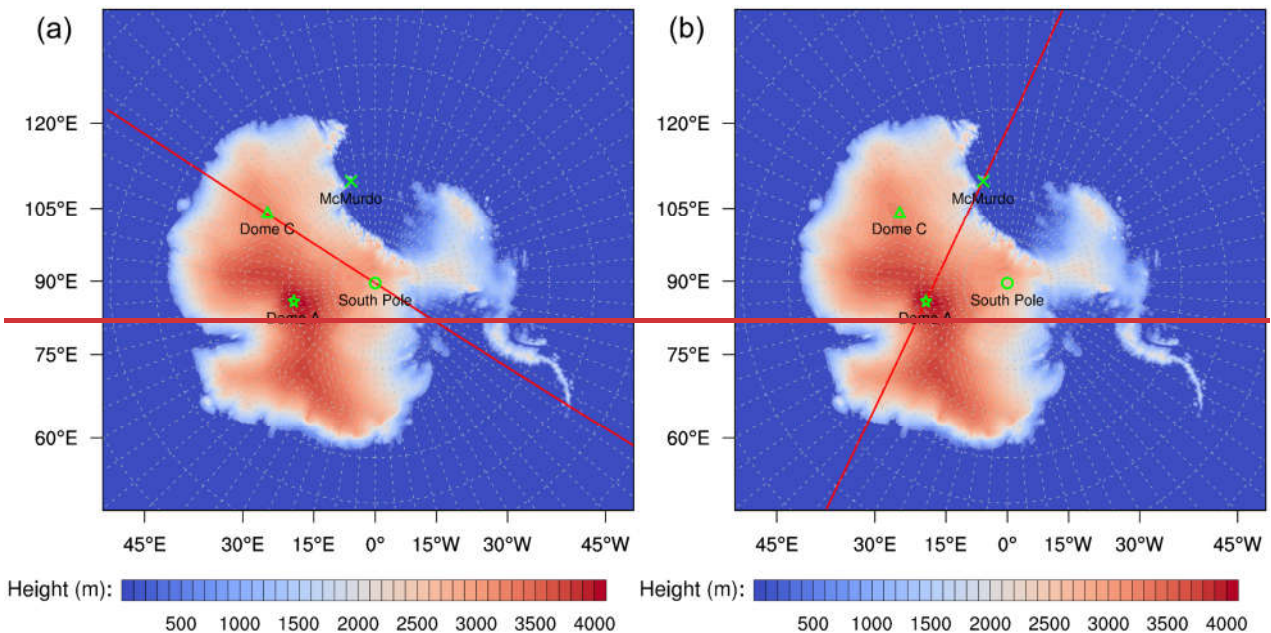
^bCorrelation coefficient between $\log_{10}(1/Ri_{\text{ImpAMPS}})$ and $\log_{10}(1/Ri_{\text{AMPS}})$ (i.e. after using the fitted model).

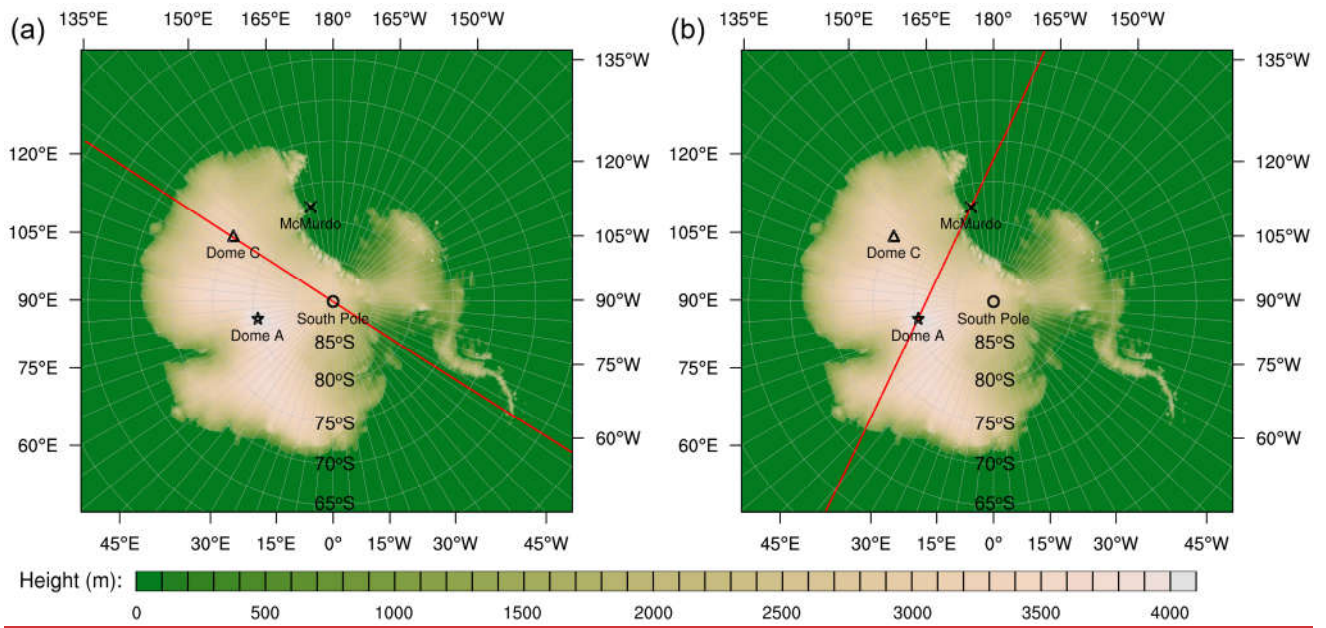
Table 2 lists the fitted model and its corresponding improvements. The correlation coefficient at all three sites increased after using the fitted model, which means that the fitted model can be used to modify the AMPS forecasts (Ri) and obtain an improved result (Ri_{Imp}). Through analysis of the fitted model, one can see the undetermined coefficient is less than 1, we find this is due to the AMPS underestimated when $Ri < 0$; the undetermined coefficient is negative, indicating the negative regulating action for Ri with increasing height.

4.2.3.

4.2.2 Vertical cross-section

The results are given in Sect. 4.2.1 that shows show the AMPS can forecast the main tendency of $\log_{10}(Ri)$. Then, we consider that it is worth a try to use the AMPS-forecasted $\log_{10}(Ri)$ to comprehend the characteristics of atmospheric turbulence in Antarctica. The results of the AMPS-forecasted $\log_{10}(Ri)$ were presented through interpolation of the AMPS grid 2 field with two vertical cross-sections, which provides us with a broader perspective on the probability of turbulence triggered in four-dimensional space-time. One vertical cross-section is interpolated through two specified points (the South Pole SP and Dome C), DC, and another is through Dome ADA and McMurdo (MM, as shown in Fig. 7). The corresponding AMPS forecasts are shown in Figs. 8 and 9, respectively.





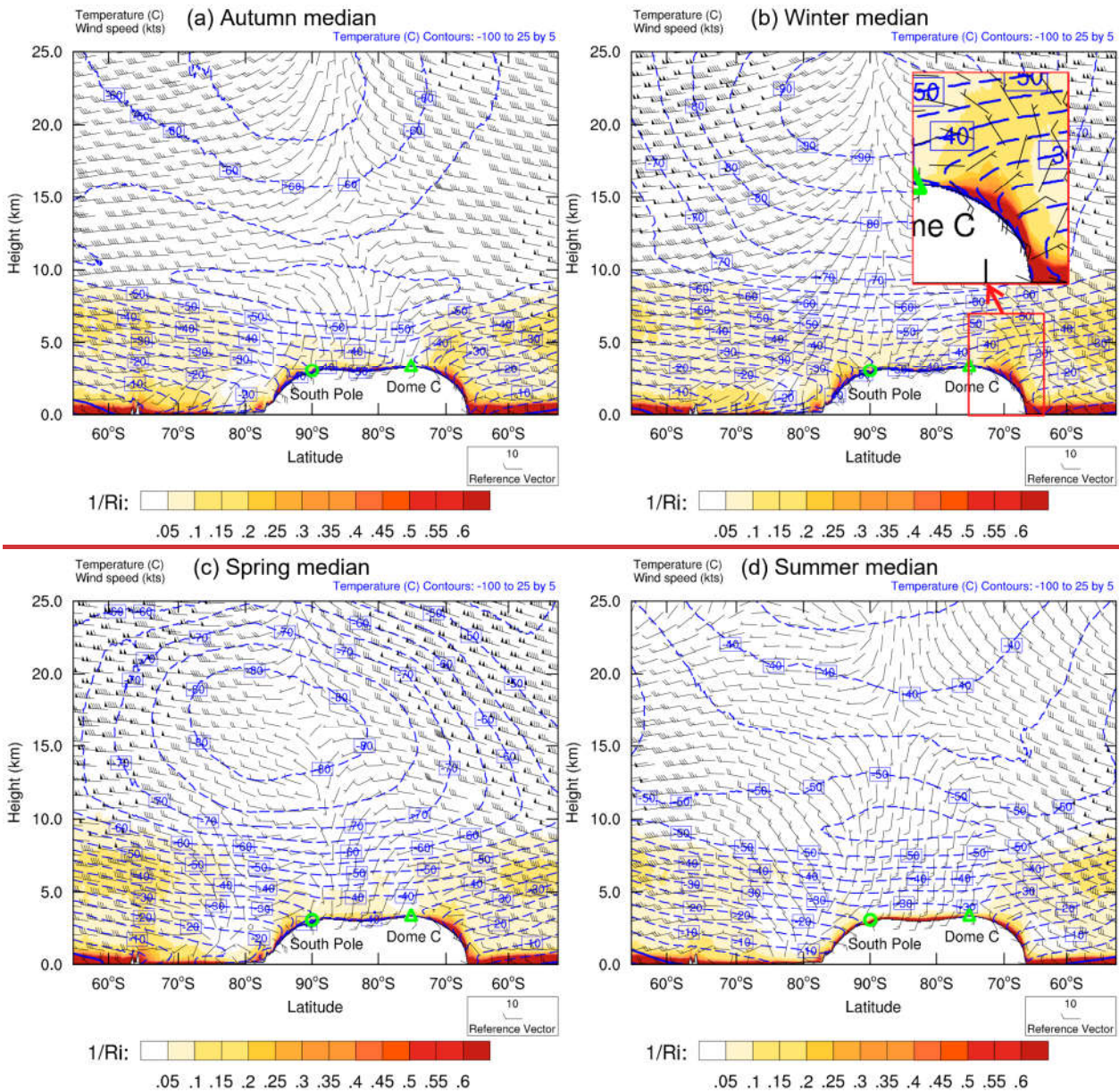
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Figure 7. Two lines (marked by red lines) are used to create vertical cross-sections. (a) a line through Dome C and the South Pole, (b) a line through Dome A and McMurdo. The colour scale indicates the terrain height (m).



400
 401 **Figure 8.** Median vertical cross section of temperature, wind speed, and $1/Ri$ along the red line through the South Pole and Dome C
 402 (shown in Fig. 7a). The height (km) on the y axis represents the elevation above sea level. The seasonal median for autumn:
 403 March 2021 to May 2021 (a), winter: June 2021 to August 2021 (b), spring: September 2021 to November 2021 (c), and summer:
 404 December 2021 to February 2022 (d), where terrain fields are given. The reference vector is 10 kts ($\sim 5.2 \text{ m s}^{-1}$) for the wind
 405 barb. Winds are depicted as blowing generated from the direction the flags are facing (earth coordinates), e.g., the reference
 406 vector represents wind blowing from the west. RAMP2 data set (<https://nsidc.org/data/nsidc-0082/versions/2>).

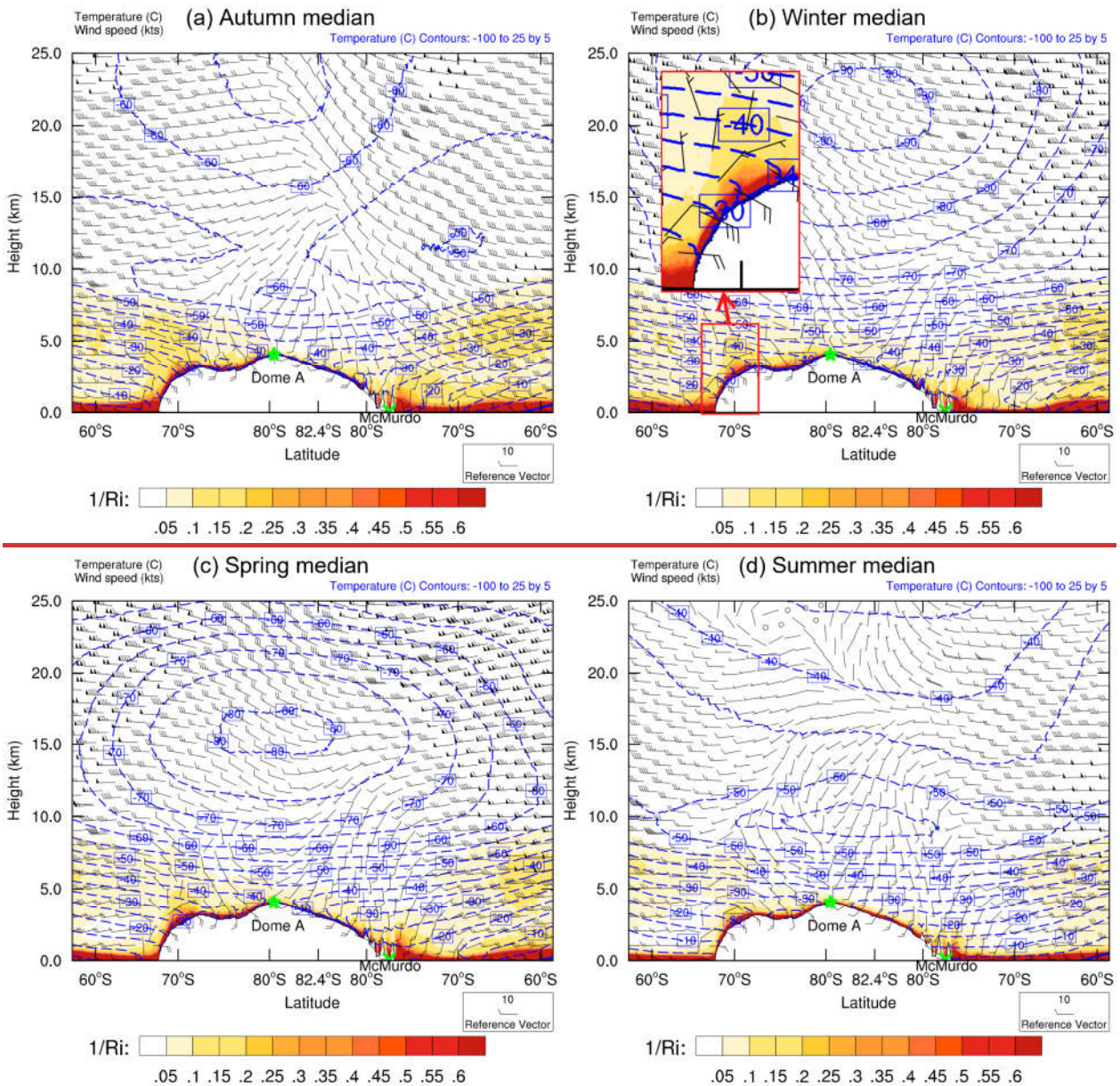


Figure 9. As in Fig. 8, but for the vertical section along the red line through Dome A and McMurdo (shown in Fig. 7b).

Figs. 8 and 9 show the seasonal median of the AMPS forecasts. The temperature and wind speed were lower above the Antarctic Plateau than over the ocean. In the polar winter, (JJA), the temperature contours are dense near the ground above the interior Plateau, representing a strong surface-layer temperature inversion (such inversion has been observed by Yagüe et al., 2001; Argentini et al., 2013; Hu et al., 2019). The surface-layer wind speeds increase from the summit to the escarpment region (caused by the well-known katabatic wind over the surface slope area in Antarctica) and then decrease toward the coast, which is consistent with previous measurements (Ma et al., 2010; Rinke et al., 2012). The Ri is obviously larger above the summits (e.g. DA and DC), suggesting the PBLH could be thin, this agrees with the results from Swain and Gallée (2006); Bonner et al. (2010); and Aristidi et al. (2015). The depths of large Ri near the ground are thinner near the summits (e.g. Dome A and Dome C) of the Antarctic Plateau. Notably, a thick depth of large Ri (which could be considered thick and will be discussed later) occurred near the escarpment region from the side of the interior plateau, as clearly shown during winter (see the enlarged drawings in

Figs. 8 and 9). Owing to the forecast ability of AMPS in vertical space, one can see that this high is caused by the wind shear, where the wind direction is southeast within the surface layer and changes to the northwest at ~5 km AGL.

The results of Ri distribution from the AMPS outputs provided us with valuable insights into the atmospheric turbulence in the Antarctic region (while using the radiosonde measurements is hard to do so). Here, we attempt to relate the features of atmospheric turbulence to some large-scale phenomena or local-scale dynamics over the Antarctic plateau and the ocean surrounding it: the shear-induced turbulence (katabatic winds, polar vortices), convection (cloud cooling, boundary layer convection), temperature inversion, and the wave-induced turbulence (orographic gravity waves, trapped lee waves, inertia-gravity waves). Table 3 lists their possible functional areas that are marked in Figs. 8 and 9. This is dedicated to qualitatively evaluating the AMPS outputs and investigating the underlying physical processes of triggering atmospheric turbulence.

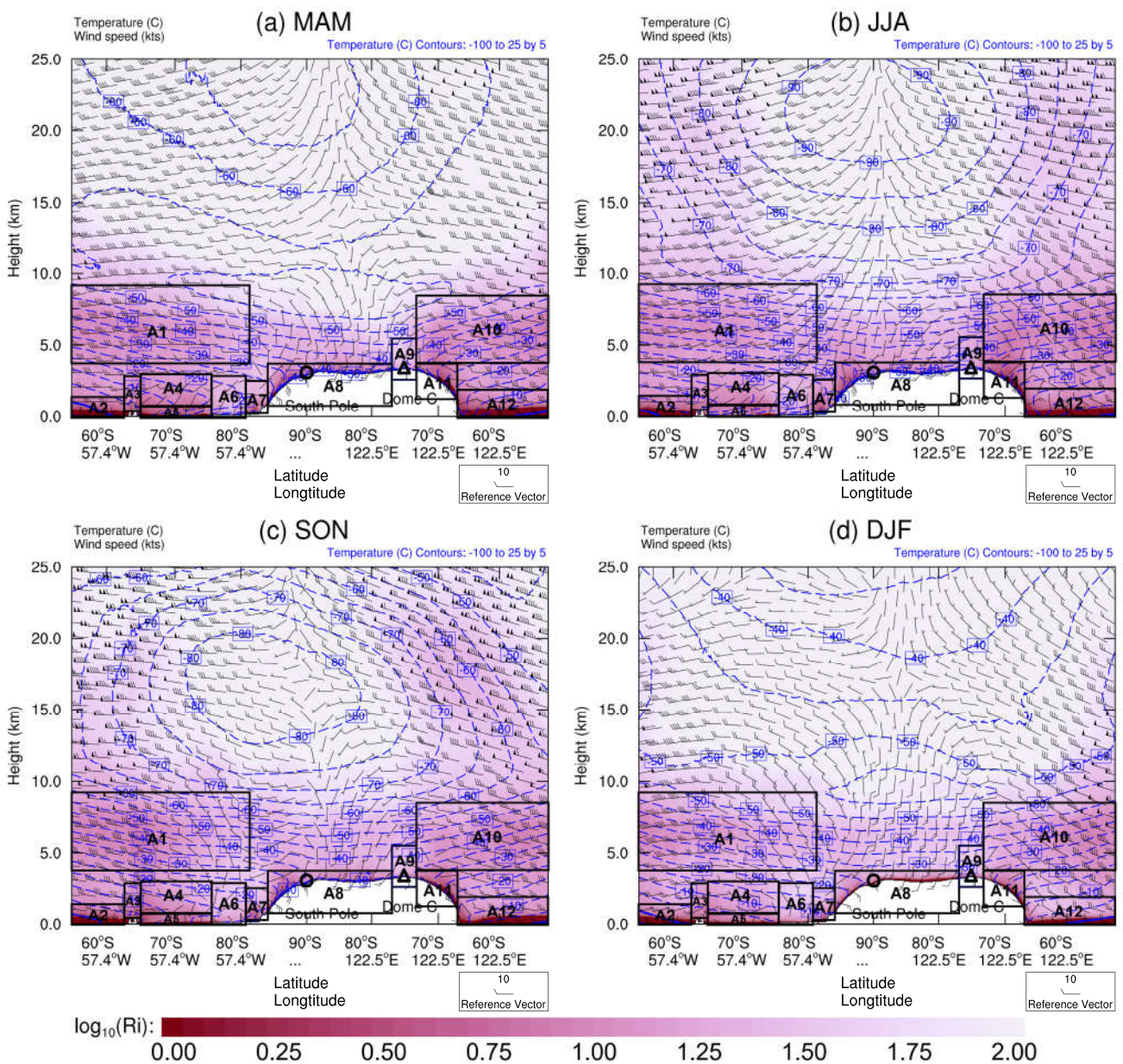


Figure 8. The seasonal median of temperature (instead of θ), wind speed, and $\log_{10}(Ri)$ along the vertical cross-section through the South Pole (black circle) and Dome C (black triangle), as shown by the red line in Fig. 7a. The height (km) on the y-axis

represents the elevation above sea level. The AX in each plot are used to mark the possible functional areas of some atmospheric activities (as listed in Table 3).

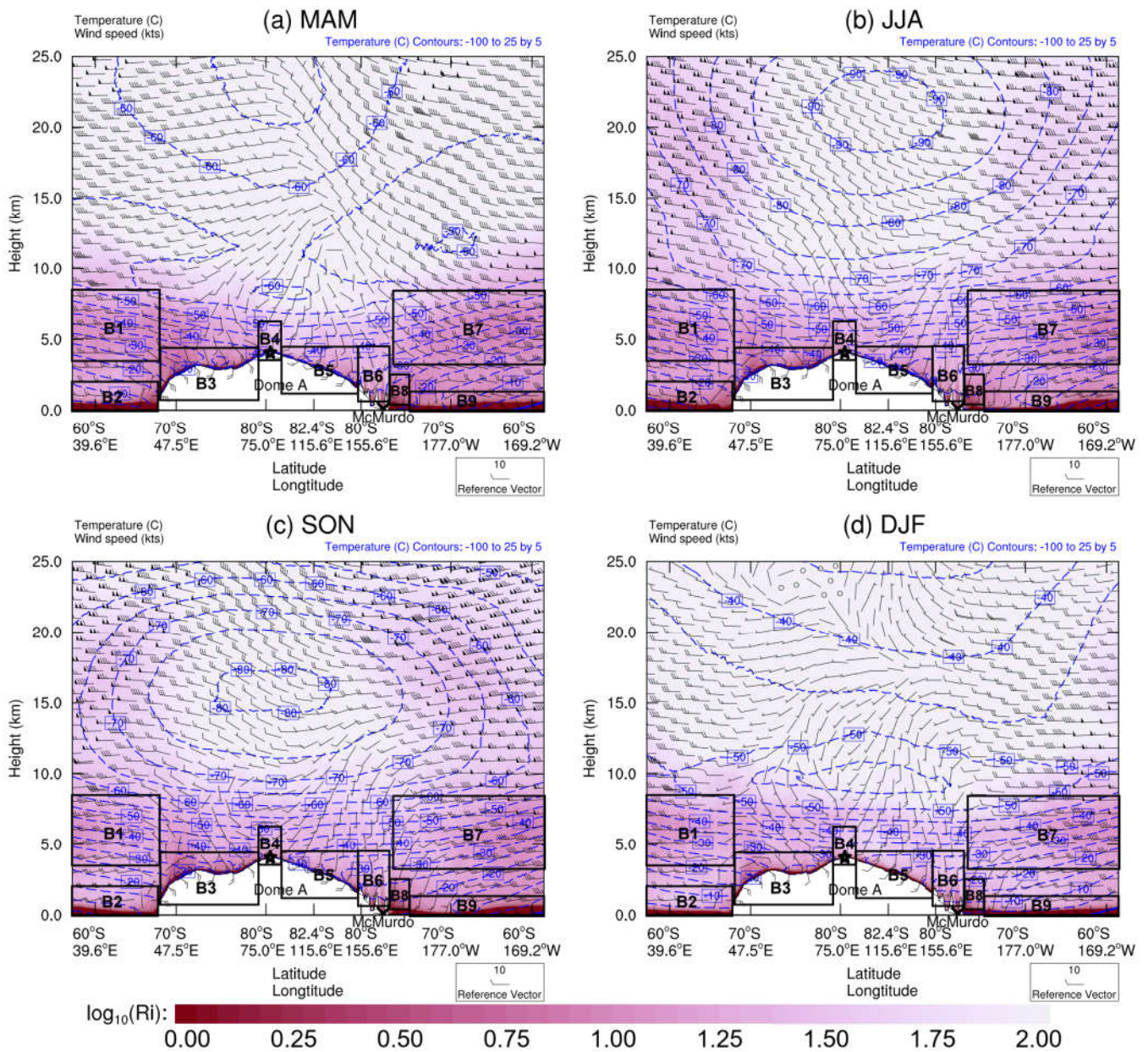


Figure 9. As in Fig. 8, but for the vertical cross-section through Dome A (black star) and McMurdo (black cross), as shown by the red line in Fig. 7b.

As a result of katabatic winds (Rinke et al., 2012), the near-surface wind speeds increase from the interior plateau to the steep slope (Figs. 8 and 9), which is driven by gravity. Strong winds can lead to strong wind shear and increased levels of mechanical turbulence (e.g., Huang et al., 2021; Solanki et al., 2022), as one can see the surface layer with small Ri at the escarpment region (see A8 and A11 areas in Fig. 8, B3 and B5 areas in Fig. 9). Where the regions between the SP and DC (in A8 area) are also located on the slope (see Fig. 7a) and show a relatively small Ri near the ground.

A strong polar vortex implies that the zonal winds are intense, and atmospheric turbulence is more prone to occur. The Antarctic polar vortex reaches its maximum intensity in the winter-spring season (Zuev and Savelieva, 2019a), which corresponds to the relatively turbulent free atmosphere with low Ri values over the ocean during JJA and SON (see Figs. 8 and 9). Moreover, the strongest zonal winds are located over the ocean (Zuev and Savelieva, 2019b), the interaction between the zonal wind and the ocean surface may generate wind shear and facilitate the development of turbulence (see areas A2 and A12 in Fig. 8, B2 and B9 in Fig. 9).

Cloud cooling refers to two kinds of cooling-induced turbulence in this study: Cloud Top Cooling (CTC) and Below Cloud-base Turbulence (BCT). The CTC is contributed by radiative cooling, which could be one of the driving mechanisms of the mixed-layer turbulence (Deardorff, 1976). The BCT usually occurs below the bases of midlevel clouds accompanied by precipitation that does not reach the ground, cooling by evaporation or sublimation seems to contribute to the turbulence (Kudo, 2013; Kantha et al., 2019). In sum, regions with clouds may advance the development of turbulence. The cloud fraction observed by satellite lidar is higher above the ocean than the Antarctic plateau (Spinhirne et al., 2005; Saunders et al., 2009). Thus, cloud may benefit small $\log_{10}(Ri)$ above the ocean (A1 and A10 areas in Fig. 8, plus the B1 and B7 areas in Fig. 9).

Boundary layer convection is generated by forcing from the ground, solar heating of the ground during sunny days causes thermals of warmer air to rise and convection will form (He et al., 2020), then the turbulence could be developed forced by buoyancy (Verma et al., 2017). The albedo of fresh snow over sea ice is very high, while that for open water is relatively small (Hines et al., 2015). Thus, solar heating will be much more stand out over open water and lead to the thermal convection boom. This can be used to reasonably explain the results in Fig. 8, that the Ri over ocean (A2 area) is smaller than over ice shelf (A6 area), and the A5 area can be regarded as a “transition region” (sea ice and open water could both exist) between them with an intermediate value of Ri .

Table 3. The possible functional areas of some typical large-scale phenomena or local-scale dynamics over the Antarctic plateau and the ocean surrounding it.

<u>Atmospheric activities</u>	<u>Possible functional areas</u>		<u>Contribution for triggering turbulence</u>
	<u>Marked areas in Fig. 8</u>	<u>Marked areas in Fig. 9</u>	
<u>Katabatic winds</u>	<u>A8, A11</u>	<u>B3, B5</u>	<u>Positive</u>
<u>Polar vortices</u>	<u>A1, A2, A5, A10, A12</u>	<u>B1, B2, B7, B9</u>	<u>Positive</u>
<u>Cloud cooling</u>	<u>A1, A10</u>	<u>B1, B7</u>	<u>Positive</u>
<u>Boundary layer convection</u>	<u>A2, A5, A12</u>	<u>B2, B9</u>	<u>Positive</u>
<u>Temperature inversion</u>	<u>A6, A9</u>	<u>B4</u>	<u>Negative</u>
<u>Orographic gravity waves</u>	<u>A3</u>	<u>B6</u>	<u>Positive</u>
<u>Trapped lee waves</u>	<u>A4, A7</u>	<u>B8</u>	<u>Positive</u>
<u>Inertia-gravity waves</u>	<u>A1, A10</u>	<u>B1, B7</u>	<u>Positive</u>

The strength of the near-ground temperature inversion forecasted by the AMPS increases from the coast to the high interior, and its strength weakens during polar summer, such a phenomenon has also been observed in previous studies (Hudson and Brandt, 2005; Ma et al., 2010). The general increase in temperature-inversion strength was considered to correspond to a less turbulent atmosphere (when the boundary layer is shallower), owing to large stability suppressing turbulence. This corresponds to the larger Ri in the summit area where a stronger temperature inversion occurred (see A9 area in Fig. 8 and B4 area in Fig.

9). There is a similar phenomenon occurred over the Ronne ice shelf (A6 area in Fig. 8), especially for JJA (when the temperature inversion is more obvious). Importantly, it should be noted that it is the range of turbulence (or $PBLH$) that would be suppressed by the temperature inversion and the turbulence intensity could be strong within the inversion layer (Petenko et al., 2019). For example, the turbulence above Dome C is mainly concentrated in the first tens of meters above the ground (Aristidi et al., 2015).

The development of Orographic Gravity Wave (OGW) is the interaction between near-surface wind and a mountain barrier (Lv et al., 2021; Zhang et al., 2022a; Zhang et al., 2022b). The OGW breaking could be a source of turbulence. Obviously, OGW can be triggered above the Antarctic Peninsula (A3 area in Fig. 8) and Transantarctic Mountains (B6 area in Fig. 9). But the atmosphere just above the top of the mountain seems to be laminar (e.g., see the larger value of Ri in B6 area), this may be due to that the breaking of the OGW may not happen immediately after being generated above the mountains.

Trapped Lee Waves (TLW) belongs to OGW. Specially, TLW, as its name implies, tends to form on the lee side of mountains and turbulence may be developed in the downstream (Xue et al., 2022). Thus, the small Ri in A4 area in Fig. 8 can be attributed by the TLW forced by the Antarctic Peninsula (see its position in Fig. 7a). It is the same case for B8 area in Fig. 9 (but forced by the Transantarctic Mountains). The katabatic winds could be linked to TLW and result in enhanced turbulence. This could explain the A7 area (Fig. 8) have small Ri on the lee side of the mountain.

Inertia-Gravity Waves (IGW) are influenced by the Coriolis effect (increasing with wind speed), and the frequency of IGW is close to inertial frequency. IGW and Kelvin-Helmholtz instability (which can be characterized by the Richardson number) are generally presumed to be closely linked. At high latitudes, the IGW energy density's maxima occur at around 5 km AGL (Zhang et al., 2022a). This may suggest that the IGW can also be a contributor to the small Ri above the ocean (A1 and A10 areas in Fig. 8, plus the B1 and B7 areas in Fig. 9).

In addition, one can see the temporal evolution of Ri vertical cross-sections for a year from the video supplement (vertical cross-section through the red line shown in Fig. 7a: <https://doi.org/10.5446/60761>; and Fig. 7b: <https://doi.org/10.5446/60760>). It shows that the atmospheric conditions are variable, and a significant transition between laminar flow and turbulent flow could occur at any time. Some activities in Antarctica require a non-turbulent atmosphere, such as astronomical observations (Burton, 2010) and aviation safety (Gultepe and Feltz, 2019) atmosphere exhibited strong daily variability. Thus, Therefore, real-time forecasting is necessary if one wishes to avoid a turbulent atmosphere in Antarctica, instead of counting only on the displayed of the Richardson number is important and helpful, rather than relying solely on the statistical results. Moreover, unstable atmospheres are presented in this study. Furthermore, the video shows that atmospheric turbulence is likely to be triggered over the ocean, moving toward the Antarctic Plateau and becoming stable/weakening. This was likely may be due to the obstruction of the high plateau, and the which creates a calm atmosphere above it behaved calmly.

4.2.43 Richardson number at the planetary boundary layer height

The Richardson number is used to determine the boundary layer height using a critical value, typically 0.25 (Troen and Mahrt, 1986; Holtslag et al., 1990; Pietroni et al., 2012). The, Thus, the critical value (or the value of Ri at the $PBLH$, Ri_{PBLH}) is worth studying, which is a significant application for the Richardson number. In addition, previous studies have suggested that the Ri_{PBLH} Richardson number depends on the vertical resolution of the model data (Troen and Mahrt, 1986; Holtslag et al., 1990). As for the resolution of the AMPS grid, it is difficult necessary to recalculate Ri_{PBLH} define what strong turbulence layer (or atmosphere instability) may appear, whenever a reference standard based on the AMPS outputs, since the value of Ri_{PBLH} is

not given. would thus be a helpful reference standard for judging whether an atmospheric layer turbulence is stable likely to be suppressed ($Ri \lesssim Ri_{PBLH}$) or unstable developed ($Ri \gtrsim Ri_{PBLH}$).

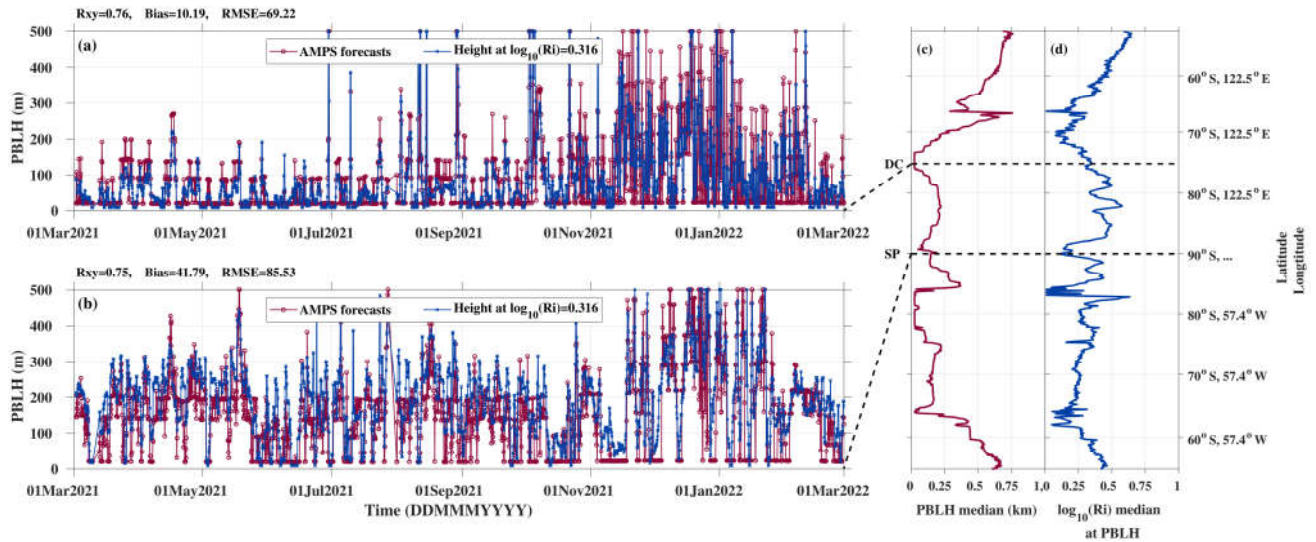


Figure 10. Temporal evolution of $PBLH$ directly forecasted by the AMPS (red circles) and estimated by the height corresponding to $\log_{10}(Ri) = 0.316$ (blue crosses) at DC (a) and SP (b). Median annual $PBLH$ (c) and $\log_{10}(Ri)$ at the $PBLH$ (d) along the red line through DC and SP shown in Fig. 7a.

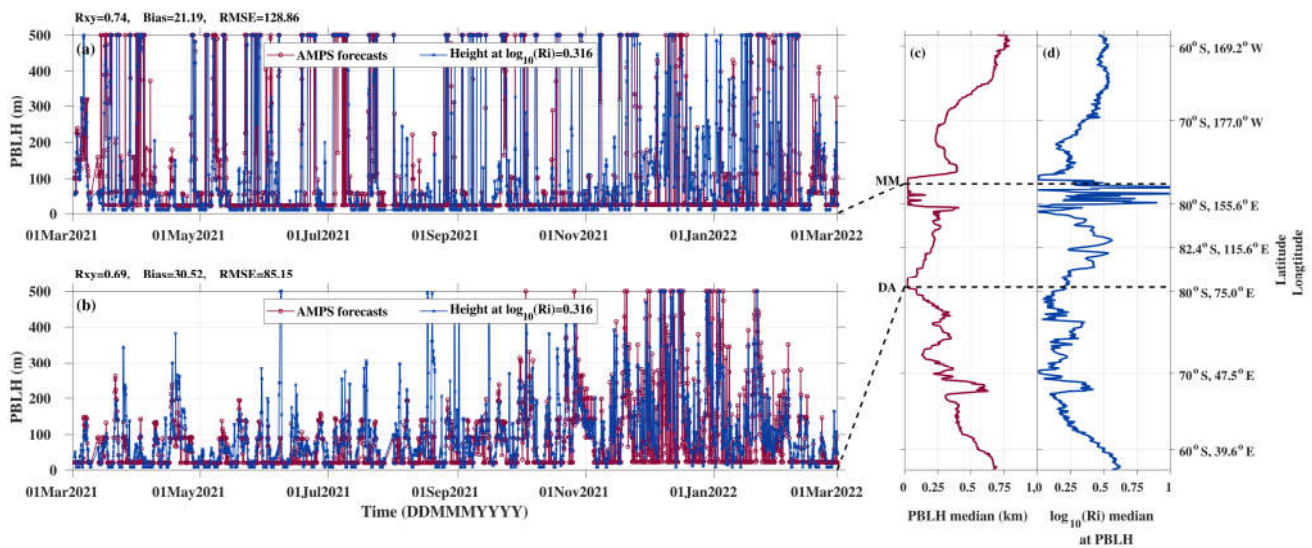


Figure 11. As in Fig. 10, but for data above MM and DA, plus the red line through MM and DA shown in Fig. 7b.

The PBL -planetary boundary layer scheme of Polar WRF in the AMPS used was the Mellor-Yamada-Janjić (MYJ; Janjić, 1994) scheme, and the ability of the AMPS to model the Antarctic PBL -was-boundary layer has been examined by Wille et al. (2017). The MYJ scheme defines the $PBLH$ where turbulent kinetic energy decreases to a prescribed value of $0.1 \text{ m}^2 \text{ s}^{-2}$ (Xie et al., 2012), and the AMPS forecasts include the values of $PBLH$. Figs. 10a (DC) and 11b (DA) show that the $PBLH$ directly forecasted by the AMPS (red lines) was mostly less than 100 m in the summit (DC and DA) during the polar winter, which such variation range is consistent with the SODAR observations (DC: Petenko et al., 2014; DA: Bonner et al.,

2010); Fig. 10b also displays a result being in accordance with the SODAR observations at the SP, as the most PBLH was shown to be within 100-300 m (Travouillon et al., 2003). Thus, the AMPS-forecasted PBLH is considered believable-to-be realistic.

Figs. 10c and 11c show the median annual PBLH forecasted by AMPS along the red lines in Figs. 7a and 7b, respectively. One can observe that the AMPS forecasted is a thin PBLH over the plateau can be observed, especially forat the summits Domes (e.g., Dome ADA and Dome CDC), which agreesis consistent with previous studies (Swain and Gallée, 2006). WhileIn contrast, a thick PBLH areis shown near the escarpment region, as seen by the dump in 67 (e.g., ~68°S (up), 122.5°E in Fig. 10c and ~68°S (down) in Fig. 11e. This; this corresponds to the relatively low Ri thick depth of near the large shown in enlarged drawings (Figs. ground in A11 area in Fig. 8 and 9).

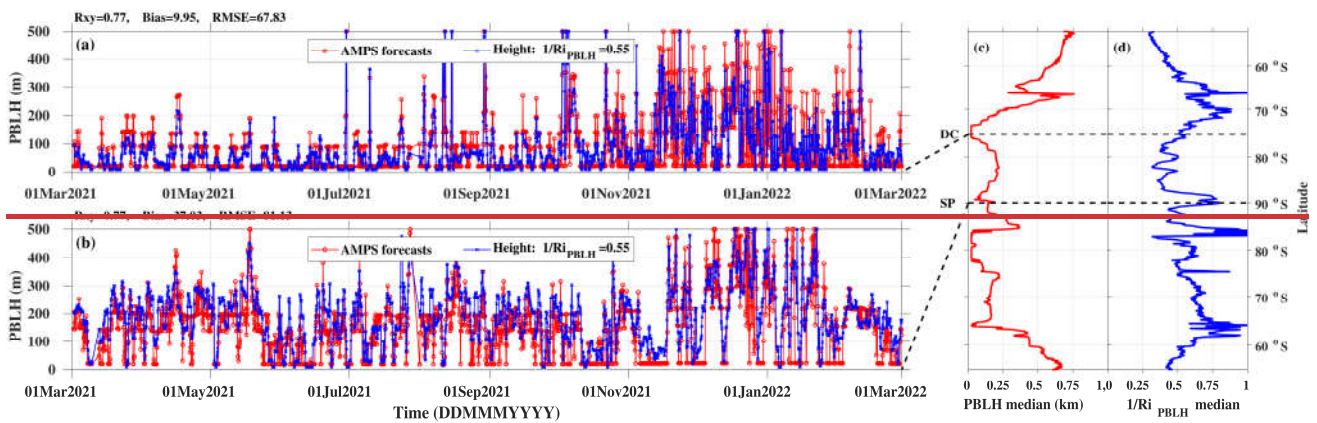


Figure 10. Temporal evolution of directly forecasted by the AMPS (red circles) and estimated by the height corresponding to $=0.55$ (blue crosses) at DC (a) and SP (b). Median annual (c) and (d) along the red line through DC and SP shown in Fig. 7a.

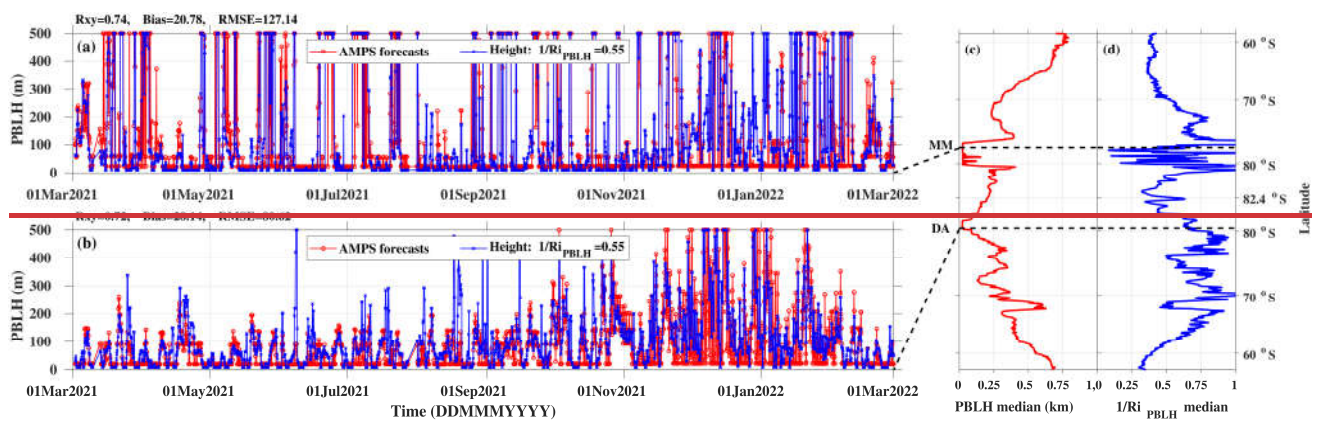


Figure 11. As in Fig. 10, but for data along the red line through MM and DA showed in Fig. 7b.

Then, The $\log_{10}(Ri_{PBLH})$ was computed using linear interpolation between the AMPS grid with height equals to the AMPS-forecasted PBLH. The median value of $\log_{10}(Ri_{PBLH})$ from the combined data of two vertical cross-sections (i.e. Figs. 10d and 11d) was calculated as 0.316, i.e., median $\log_{10}(Ri_{PBLH})$ equals to 0.316 (or Ri_{PBLH} 55 to give only one typical value.=2.07). However, some researchers have proposedemployed $Ri_{PBLH}=0.25$ when radiosoundingsthe radiosonde measurements with the high vertical resolution arewas used (e.g. at Dome C; Pietroni et al., 2012). Here, the smaller critical value larger Ri_{PBLH} for the

AMPS forecasts may be caused by the coarse vertical grid resolution (as implied from Troen and Mahrt, 1986) of the AMPS and its data smoothness (as mentioned in Sect. 4.2.1; the AMPS forecasts show smaller/larger Ri than the radiosoundings that even though they have already been interpolated to the same vertical resolution as the AMPS) grid).

To test the credibility of the critical value (i.e. $\log_{10}(Ri_{PBLH})=0.55316$), $PBLH$ was also derived as the height where the AMPS-forecasted $\log_{10}(Ri)$ decreases to 0.316. The R_{xy} , Bias, and RMSE (blue lines in Figs. 10a-b and 11a-b). Even if larger critical values are used, no substantial variations would occur in the estimate because of the large gradient of near the ground (Fig. 4), also mentioned in . In the internal Antarctic Plateau (DC, SP, and DA), the shows a slight change in magnitude during the cold seasons (e.g. from March to October 2021). Such seasonal variation is not significant along the coast of Antarctica (MM), which fluctuates every season. During the polar summer (e.g. January 2022) at all four sites, estimated by $=0.55$ and read directly from the AMPS forecasts fluctuate considerably. The , , and of $PBLH$ between the estimations using the critical value ($=0.55$) and AMPS forecasts, (blue lines in Figs. 10a-b and 11a-b) and the direct forecasts of AMPS (red lines in Figs. 10a-b and 11a-b) are depicted in the top left of the plot, where Bias indicates the former minus the latter. It appears that the values of R_{xy} (\rightarrow all larger than 0.7269) are almost satisfactory, and then we may conclude that $\log_{10}(Ri_{PBLH})=0.55316$ is a reliable critical value for judging the behaviour/behavior of atmospheric turbulence. The atmosphere layer could be considered turbulent when for $\log_{10}(Ri_{PBLH}) \geq 0.55$ (this 316 (when the turbulence intensity could be comparable to that within the boundary layer). However, this critical value may only be valid for using the AMPS forecasts). $\rightarrow 9$ indicates the results calculated by this critical value overestimate the . And >67 suggests that one should be careful when estimating a precise value for with this critical value (0.55).

5 Discussion

The above results of the distribution provided us with valuable insights into the atmospheric turbulence in the Antarctic region. Here, we attempt to relate the features of atmospheric turbulence to some typical atmospheric conditions in Antarctica: temperature inversion, katabatic winds, and polar vortices.

The strength of the near ground temperature inversion increases from the coast to the interior of Antarctica, and its strength weakens during summer (Fig. 2), such a phenomenon has also been observed in previous studies. Combined with the analysis of Figs. 11 and 12, the general increase in temperature inversion strength was found to correspond to a less turbulent atmosphere with smaller (when the boundary layer is thinner), owing to large stability suppressing turbulence.

As a result of typical katabatic winds, the wind speeds increase from the interior plateau to the steep slope (Figs. 8 and 9). Strong winds can lead to increased levels of mechanical turbulence, as one can see a thick depth of large layer at the escarpment region.

A strong polar vortex means that zonal winds are strong, and atmospheric turbulence is more likely to be triggered. The Antarctic polar vortex reaches its peak intensity in the winter spring and weakens during summer. One can also see that the free atmosphere in Antarctica becomes relatively calm during summer, as shown by the smaller in Figs. 8d and 9d. On the other hand, the strongest zonal winds exist over the ocean, which suggests that the atmosphere above the ocean can be an important source of turbulence (Figs. 8 and 9 show larger over the ocean).

6 Conclusion

We have examined the ability of AMPS to forecast the ~~inverse of the~~ Richardson number (~~the larger the, the higher the probability of triggering turbulence in the atmosphere~~) in the Antarctic atmosphere. This includes ~~quantifying~~evaluating the accuracy of meteorological parameters (~~Temperature~~ θ and ~~Wind speed~~, V , on which the Ri depends), and comparing the $\log_{10}(Ri)$ ~~from the estimations between radiosoundings and~~ AMPS forecasts ~~and radiosoundings~~. In addition, the analysis of atmospheric $\log_{10}(Ri)$ over the entire Antarctic continent and the ocean surrounding it, was presented on an annual time scale. Finally, the $\log_{10}(Ri)$ at the Planetary Boundary Layer Height (PBLH) Further, the forecasted by the AMPS was employed to understand how to evaluate atmospheric stability or instability using the value of θ and V has been calculated.

From the analysis presented above, we deduce the following:

1. Comparisons of ~~the~~ AMPS forecasts with ~~the~~ radiosoundings from three representative sites (coast: McMurdo, flank: ~~the~~ South Pole, summit: Dome C) show that the forecasts can accurately describe the trend of atmospheric meteorological parameters above the Antarctic continent, as the R_{xy} for θ reached as high as 0.99 and the R_{xy} for V in almost the whole range from the ground to 25 km AGL. The monthly median difference in the air temperature is of the order of 1 K in a large part of the atmosphere, where the cold temperature bias occurs near the ground during winter. As for the wind speed, the median difference was -1 m s^{-1} and rarely exceeded 2 m s^{-1} , are all larger than 0.85 (Table 2).

2. We proved that the AMPS forecasts can identify the main characteristics of atmospheric turbulence over the Antarctic continent in terms of both space and time. The R_{xy} of $\log_{10}(Ri)$ at MM, SP, and DC are 0.71, 0.59, and 0.53, respectively. And the AMPS can reconstruct the near-ground “convex-concave-convex” shaped $\log_{10}(Ri)$ profiles indicated by the radiosonde measurements (Fig. 5). We also find that the R_{xy} of $\log_{10}(Ri)$ would be higher when the RMSE of θ and V are smaller (Table 2). Besides, the AMPS can better capture the trend of $\log_{10}(Ri)$ (R_{xy} would be larger) at a relatively unstable atmosphere (weaker temperature inversion). Moreover, the values of $\log_{10}(Ri)$ 66, and 0.68, respectively, and larger was obtained during warm seasons. Moreover, the value of $\log_{10}(Ri)$ are all were generally overestimated at the three sites; this is partly the result of the potential temperature gradients at the unstable atmosphere being overestimated by the AMPS, and the AMPS has generally underestimated at these three sites. This may be because the model results are generally smoother than the measurements. Furthermore, a model depending on height to improve the AMPS forecasted has been proposed the wind shear when it was strong.

3. The seasonal medians of the AMPS forecasts from two vertical cross-sections were presented, (Figs. 8 and ~~the AMPS forecasts~~), which provides us with a broader perspective on when and where atmospheric turbulence could be highly triggered in the Antarctic region. The AMPS-forecasted $\log_{10}(Ri)$ were qualitatively verified again. it can be observed that, as its statistical distribution behaved as the AMPS can forecast the main expected atmospheric properties as expected: strong temperature inversion near the ground, attributed by some typical large-scale phenomena or local-scale dynamics (katabatic winds, polar vortices, convection, gravity wave, etc.) over the Antarctic plateau and the ocean surrounding it. For example, a very calm laminar atmosphere (is small) at high altitudes, and a thin boundary layer above the Antarctic Plateau and a shallow boundary layer in the Domes area are illustrated by the AMPS forecasts.

4. The $\log_{10}(Ri)$, especially in the summit area. Moreover, we also obtained a new result, that strong wind shears near the escarpment regions have occurred from at the PBLH ground to 5 km AGL, causing a thick turbulent atmosphere with large

622 values of (or thick boundary layer), which is more evident during the polar winter. In addition, according to the temporal
623 evolution of vertical cross-sections for a year from the video supplement, we find that unstable atmospheres are likely to be
624 triggered over the ocean, move toward the Antarctic Plateau, and become stable.

625 4. The AMPS can forecast the boundary layer height (z_{BL}). And the z_{BL} at the (z_{BL}) was were calculated in this study, which could be a
626 helpful reference standard for judging whether the atmospheric layer is stable (\leftarrow) or unstable (\rightarrow) when using AMPS forecasted.
627 The and their median value of z_{BL} from the combined data of two vertical cross-sections is 0.316. $\log_{10}(Ri) = 0.31655$, which has
628 been in turn was used for estimating the z_{BL} and agrees agree well with the AMPS-forecasted $PBLH$ ($R_{xy} \geq$
629 0.69). The atmosphere layer could be considered turbulent at $\log_{10}(Ri) \leftarrow 0.72$. <0.316 (when the turbulence intensity could be
630 comparable to that within the boundary layer).

631 The overall results show that the AMPS can forecast the a realistic behaviour of Ri and could be applied to, and the turbulence
632 conditions in Antarctica are well revealed; furthermore, some practical operations that want to avoid a turbulent atmosphere —
633 such as astronomical observations (Burton 2010), aviation safety (Gultepe and Feltz 2019) and free space optical communication
634 (Jianjun Yin et al. 2017) — can apply the AMPS-forecasted Ri astronomy in Antaretica, which is interested in the impacts of
635 atmospheric turbulence..

636 Data availability

637 The meteorological parameters measured by radiosoundings the radiosondes at McMurdo, South Pole that support the findings of
638 this study are available at the Antarctic Meteorological Research Center (<ftp://amrc.ssec.wisc.edu/pub>), while the meteorological
639 parameters at Dome C are available at the Antarctic Meteo-Climatological Observatory (<http://www.climantartide.it>). The
640 original WRF output files of AMPS used in this study can be found at
641 https://www2.mmm.ucar.edu/rt/amps/information/amps_esg_data_info.html.

642 Video supplement

643 The video supplement related to the vertical cross-section through the South Pole and Dome C in this article (Fig. 7a) is available
644 online at <https://doi.org/10.5446/60761>. Another vertical cross-section through Dome A and McMurdo (Fig. 7b) is
645 <https://doi.org/10.5446/60760>.

646 Author contributions

647 QY and XW planned the investigation; QY, XH, XW, and ZW analyzed the data; QY and YG wrote the manuscript draft; QY
648 finished the visualization; QY, XW, XH, XQ, and ZW performed the valuation; XW, CQ, TL, XQ, and PW reviewed and edited
649 the manuscript.

650 Competing interests

651 The authors declare that they have no conflict of interest.

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