# Simultaneous retrievals of biomass-burning aerosols and trace gases from the ultraviolet to near-infrared over northern Thailand during the 2019 pre-monsoon season

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**Abstract.** With the advent of spaceborne instruments in a geostationary constellation, measuring high-spectral resolution ultraviolet–visible (UV–VIS) and selected near-/shortwave-infrared (NIR/SWIR) radiances can enable probing the lifecycle of key atmospheric trace gases and aerosols at higher temporal resolutions over the globe. The UV-VIS measurements are important for retrieving several key trace gases (e.g., $O_3$, $SO_2$, $NO_2$, HCHO) and particularly for deriving aerosol characteristics (e.g., aerosol absorption and vertical profile). This study examines the merit of simultaneous retrievals of trace gases and aerosols using a ground-based spectroradiometer covering the UV–NIR to monitor their physicochemical processes, and to obtain reliable aerosol information for various applications. During the 2019 pre-monsoon season over northern Thailand, we deployed a ground-based SMART–s (Spectral Measurements for Atmospheric Radiative Transfer–spectroradiometer) instrument, which is an extended-range Pandora with reliable radiometric calibration in 330–820 nm range, to retrieve remotely sensed chemical and aerosol properties for the first time near biomass-burning sources. The high spectral-resolution (~1.0 nm full-width-half-maximum with ~3.7× oversampling) of Sun and sky measurements from SMART–s provides several key trace gases (e.g., $O_3$, $NO_2$, and $H_2O$) as well as aerosol properties covering the UV where significant light-absorption occurs by the carbonaceous particles. During the measurement period, highly correlated total column amounts of $NO_2$ and aerosol optical thickness ($\tau_{aer}$) retrieved from the SMART–s (correlation coefficient, R = 0.74) indicated their common emissions from biomass-burning events. The SMART-s retrievals of spectral single-scattering albedo ($\omega_0$) of smoke aerosols showed an abrupt decrease in the UV, which is an important parameter dictating photochemical processes in the atmosphere. The values of $\omega_0$ and column precipitable water vapor ($H_2O$) gradually increase with the mixing of biomass-burning smoke particles and higher

35     water vapor when approaching the monsoon season. The retrieved $\omega_0$ and weighted-mean-radius of fine-mode aerosols from the SMART–s showed positive correlations with the $H_2O$ (R = 0.81 for $\omega_0$ at 330 nm and 0.56 for volume-weighted-mean-radius), whereas the real-part of the refractive-index of fine-mode aerosol ($n_f$) showed negative correlations (R = −0.61 at 330 nm), which suggest that aerosol aging processes including hygroscopic growth (e.g., humidification and cloud processing) can be a major factor affecting temporal trends of aerosol optical properties. Retrieved $n_f$ and $\omega_0$ were closer to those of the water

40     droplet (i.e., $n_f$ of about 1.33 and $\omega_0$ of about 1.0) under lower amounts of $NO_2$ during the measurement period; considering that the $NO_2$ amounts in the smoke may indicate aging of the plume after emission due to its short lifetime, the tendency is also consistent with active hygroscopic processes of the aerosols over this area. Retrieved UV aerosol properties from the SMART–s generally support the assumed smoke aerosol models (i.e., the spectral shape of aerosol absorption) used in current NASA's satellite algorithms, and their spectral $\omega_0$ retrievals from ground and satellites showed good agreements (R = 0.73–

45     0.79). However, temporal and spectral variabilities of the aerosol absorption properties in the UV emphasize the importance of a realistic optical model of aerosols for further improvements in satellite retrievals.

## 1 Introduction

        Significant spatiotemporal variabilities of the aerosols in the atmosphere complicate understanding of their scattering and absorption of the solar irradiance, which results in one of the largest uncertainties in predicting future climate (IPCC, 2013;

50     Gliß et al., 2021; Myhre et al., 2013 *and references therein*). The dominant fraction of the aerosols over the globe cools the atmosphere by reflecting solar irradiance, whereas some species (e.g., black carbon in the smoke plumes) heat the air by absorbing sunlight (i.e., direct radiative effects [DRE]: Chylek and Coakley, 1974; Haywood and Boucher, 2000; Yu et al., 2006). Primary factors of aerosols affecting the DRE are their loading and their absorption properties (e.g., Takemura et al., 2002 *and references therein*), which are often defined as aerosol optical thickness ($\tau_{aer}$; total extinction by aerosols) and single-

55     scattering-albedo ($\omega_0$; a ratio of the scattering to total extinction by aerosols), respectively. The $\omega_0$ is calculated from complex refractive indices ($n+ik$; where $n$ and $k$ are a real and imaginary part, which depends on chemical composition) and particle size distribution (PSD), by assuming a spherical (Mie, 1908) or more sophisticated shape (e.g., Mishchenko et al., 2003; Yang et al., 2007). The $\omega_0$ of non-absorbing aerosols (e.g., sea salt, sulfate, and nitrate particles) is close to 1.0 with a relatively flat spectral shape, whereas it decreases down to about 0.7 for absorbing aerosols (e.g., smoke and dust particles) with significant

60     spectral gradients (e.g., Dubovik et al., 2002; Eck et al., 2013; Müller et al., 2011; Sayer et al., 2014). Meteorological condition and aerosol hygroscopicity are also the key parameters affecting $\omega_0$ since increased water content in the particles changes the $n$, $k$, and PSD, which enhances light-scattering and results in higher $\omega_0$ than dry particles (e.g., Jefferson et al., 2017; Li et al., 2019; Tao et al., 2014).

        Decades of efforts have led to remote sensing techniques from both ground and satellite providing reliable $\tau_{aer}$

65     retrievals over major parts of the globe (e.g., Giles et al., 2019; Hsu et al., 2019; Levy et al., 2013), whereas other aerosol properties retrieved from satellites are yet limited and relatively more uncertain due to the lower measurement sensitivity and

surface contributions (e.g., Jeong et al., 2016; Moosmüller et al., 2009). However, recent studies using more measurement parameters (e.g., multi-angle polarimetric measurements) showed promising results to provide reliable aerosol properties and constituents from satellites (e.g., Dubovik et al., 2019; Li et al., 2019). Globally networked ground-based instruments have

70  provided reliable optical and physical properties of aerosols (e.g., $n$, $k$, $\omega_0$, and PSD), which are less affected by surface reflectance and acquire sufficient information content from multiple observation geometries (e.g., Dubovik and King, 2000; Jeong et al., 2020; Nakajima et al., 2020; Sinyuk et al., 2020). Satellite-based retrievals have utilized the aerosol properties from ground-based instruments as key constraints to expand upon their limitations (e.g., Hsu et al., 2019; Levy et al., 2013; Sayer et al., 2014).

75      Wildfires and prescribed fires have burned about 3.5% of Earth's ice-free land surface each year from 2001 to 2010 (Randerson et al., 2012) and emit a significant fraction of global aerosols and their precursors into the atmosphere. The biomass-burning aerosols (or smoke) consist primarily of carbonaceous aerosols (black and organic carbon), inorganic particles (e.g., potassium, chloride, sulfate, inorganic salts, and trace minerals), and inorganic and organic vapors (Hodshire et al., 2019 *and references therein*). Particularly primary and secondary organic aerosols, which account for a substantial fraction

80  of fine-mode smoke aerosols, comprise various compounds with enormously different volatility, oxidation, and hygroscopic properties (Xu et al., 2017 *and references therein*). Due to the reactivity and diversity of smoke particles, the $\omega_0$ evolves with its environment (i.e., location and season), age, mixing state, and emission source of the plume (e.g., Eck et al., 2013; Haywood et al., 2003; Konovalov et al., 2017). In addition, Petters et al. (2009) reported that a major fraction of the smoke aerosols is already cloud-condensation-nuclei (CCN) active, and do not require chemical conversion to be more hygroscopic particles for

85  cloud formation and wet deposition, which adds another complication in understanding the Earth's climate.

      Numerous studies have utilized ground-, airborne-, and satellite-based remote sensing techniques to monitor the properties and aging processes of the smoke particles. For example, Haywood et al. (2003) compared aerosol properties (e.g., PSD, $\tau_{aer}$, and $\omega_0$) from the collocated AERONET (AErosol RObotic NETwork; Holben et al., 1998) and airborne in-situ measurements at Windhoek, Namibia in September 2000, which showed excellent agreements. Eck et al. (2013) analyzed the

90  seasonal trend of aerosol properties retrieved from the AERONET and OMI (Ozone Monitoring Instrument) over southern Africa for a 15-year period and reported that the $\omega_0$ increases significantly as the burning season progresses. Pistone et al. (2019) compared the spectral $\omega_0$ of smoke aerosols from six independent airborne- and ground-based remote-sensing/in-situ instruments in September 2016 out of Walvis Bay, Namibia, which showed acceptable agreements within the known uncertainties of each instrument (relative differences less than about 0.03 in mid-visible and less than about 0.05 in near-

95  infrared, depends on the instruments). Over Southeast Asia, a series of field campaigns including BASE-ASIA (Biomass-burning Aerosols in SouthEast Asia: Smoke Impact Assessment) in 2006 and 7-SEAS (Seven SouthEast Asian Studies) from 2008 to present aimed to characterize aerosol-meteorological interactions over the region, mostly focusing on the smoke plumes. Physicochemical and optical properties of smoke aerosols were analyzed by utilizing intensive ground- and satellite-based instruments during the campaigns (e.g., Lin et al., 2013; Pantina et al., 2016; Reid et al., 2013; Tsay et al., 2013, 2016).

One of the important characteristics of the carbonaceous aerosols is their significant spectral variabilities of optical properties in the ultraviolet (UV) wavelengths, which are associated with photolysis processes in the atmosphere, thereby affecting tropospheric photochemistry, human health and agricultural productivity (e.g., George et al., 2015 *and references therein*). A majority of previous studies utilized direct/diffuse irradiance instruments (e.g., UV–MultiFilter Rotating Shadowband Radiometer [UV–MFRSR], Brewer spectroradiometer) to retrieve $\omega_0$ in discrete channels in the UV. However, as these instruments measure only two observation parameters per channel, the algorithms adopted different sources of measurements/assumptions to complement the insufficient information (e.g., *see Table 1 in* Corr et al., 2009). For instance, in the absence of additional collocated instruments, they assumed fixed asymmetry parameter and surface albedo from previous studies or climatology (e.g., Bais et al., 2005; Petters et al., 2003; Wetzel et al., 2003). Collocated AERONET instruments have provided more realistic constraints of aerosol properties to the UV–MFRSR measurements (e.g., PSD and $n$ from visible [VIS] wavelengths; Corr et al., 2009; Krotkov et al., 2005a) for retrieving $\omega_0$ in the UV. Trace gas absorption (e.g., $O_3$ and $NO_2$) is another source of error for the $\omega_0$ retrieval using these instruments. To take into account the gas absorptions, Goering et al. (2005) simultaneously retrieved total column $O_3$ in addition to the $\tau_{aer}$, and $\omega_0$, by using the spectral feature of irradiance. Later, Taylor et al. (2008) added a wavelength-independent asymmetry parameter to the state vector, where both algorithms are based on the optimal-estimation method (OEM; Rodgers, 2000). Krotkov et al. (2005a) used the aerosol phase function calculated from $n$ at 440 nm and PSD from the AERONET, and total column $O_3$ from the Brewer spectroradiometer, to retrieve $\omega_0$ in the UV channels. To account for the $NO_2$ absorption, which is a significant error source of $\omega_0$ retrieval for low aerosol loading, they added retrieved $NO_2$ from the Brewer spectroradiometer for their algorithm (Cede et al., 2006; Krotkov et al., 2005b). A SKYNET (SKY radiometer NETwork) instruments are a similar type of Sun-sky spectroradiometer to the AERONET, which provides $\omega_0$ at discrete channels in the UV (i.e., 340 and 380 nm). The SKYNET algorithm accounts for the $O_3$ absorption by using its retrieved total column from its 315 nm channel (Nakajima et al., 2007; 2020). Accuracies of the $\omega_0$ retrievals from the SKYNET depend on errors in measurement and calibrations for Sun and sky-scans, surface albedo, cloud contamination, and the version of the processing software (i.e., Skyrad pack), which showed relative high biases compared to the AERONET (up to 0.07 at longer wavelengths). Recently, Mok et al. (2018) combined the AERONET (for $n$, PSD, $\tau_{aer}$) and Pandora (for total column $O_3$ and $NO_2$) products to the UV–MFRSR measurements to retrieve spectral $\omega_0$ in the UV, which showed excellent agreements with SKYNET in the UV (i.e., 340 and 380 nm) but lower correlations in the longer wavelengths (i.e., 673 and 870 nm).

In addition, spectral $n$ and $k$ provide information not only on optical properties but also the chemical composition and physical status. The $k$ demonstrates the attenuation of light by particles, which is the key parameter for determining $\omega_0$, whereas the $n$ describes the phase of light scattering by the particles. Numerous studies have focused on measuring/retrieving the $n$ and $k$ by utilizing various techniques to understand the effects of atmospheric particles on climate forcing and tropospheric photochemistry. Table 1 summarized reported values of the $n$ from the previous and current studies. Kim et al. (2010) retrieved the $n$ (at 670 nm) of secondary organic aerosols (SOA) generated by oxidizing $\alpha$-pinene, $\beta$-pinene, and toluene with $O_3$, $NO_x$, and sunlight. The retrieved $n$ varied between 1.38 and 1.61, and they suggested that the $n$ of SOA depends on aerosol mass

concentration, oxidation chemistry, temperature, and aerosol aging. Liu et al. (2013) measured the $n$ and $k$ of SOA for 220 to 1200 nm using a variable angle spectroscopic ellipsometer, and reported a rapid increase of the $n$ and $k$ in the UV. The $n$ of the three selected SOA ranged from 1.53 to 1.58 at 310 nm, 1.49−1.52 at 550 nm, and 1.48−1.50 at 1000 nm. Shepherd et al. (2018) estimated the spectral $n$ of urban, remote, and wood smoke aerosols from 460 to 760 nm based on the optical trapping method, and reported high values of $n$ of the wood smoke aerosols (~1.58) compared to the other types (1.47–1.52). They also well summarized and compared their values of the spectral $n$ to other studies in their paper. Sumlin et al. (2018) retrieved the spectral $n$ and $k$ (at 375 nm, 405 nm, 532 nm, and 1047 nm) of brown carbon aerosols emitted from controlled fire using burning sources at various geographic origins. They reported that the $n$ varies between 1.5 and 1.7 without meaningful dependencies on wavelength, moisture content, source depth, or geographic origin, whereas the $k$ increases from 0.003 to 0.014 as wavelengths vary from 532 to 375 nm. Biagio et al. (2019) estimated the $n$ and $k$ (at discrete channels in 370–950 nm) of 19 mineral dust aerosols from different sources based on Mie calculations combining optical and size measurements. They reported higher $k$ (lower $\omega_0$) of dust particles in the shorter wavelengths, which also depends on the iron content of dust, but the source and wavelength-independent values of $n$ ranged from 1.48 to 1.55. More recently, Womack et al. (2021) retrieved the $n$ and $k$ of biomass-burning aerosols from 13 controlled fires over a 360–720 nm spectral range using a broadband cavity-enhanced spectrometer combined with PSD measurements. Their algorithm incorporates Mie and Rayleigh–Debye–Gans scattering theories to account for both spherical and non-spherical particles, and retrieved $n$ to be about $1.55 - 1.60$ and $k$ to be significantly high (~0.25) in the UV.

To be closely in line with and continue such efforts, we deployed a set of instruments including the SMART–s (Spectral Measurements for Atmospheric Radiative Transfer–spectroradiometer; Jeong et al., 2018, 2020) and AERONET during the pre-monsoon yet active biomass-burning season at Fang, Thailand in 2019. Specifically, we aim to suggest the benefit of simultaneous retrievals of aerosols and trace gases covering UV, which may provide useful information on their physicochemical processes. In addition, aerosol properties in the UV are also important for various satellite algorithms for deriving higher-order aerosol parameters (e.g., absorption and vertical distribution), for which reliable measurements remain sparse. Benefits of employing SMART–s, a major instrument we utilized for this study, include:

- sufficient spectral resolution and coverage for measuring both aerosols and key trace gases (e.g., $O_3$, $NO_2$, and $H_2O$ retrievals from direct-Sun measurements), in turn, the high-temporal measurements of gaseous absorption help improve the accuracy of $\omega_0$ retrieval;

- instantaneous measurements of the Sun/sky spectrum, permitting aerosol spectral properties retrieved from an identical set of volumes;

- reliable radiometric calibration from about 330 to 820 nm by utilizing a NIST-traceable (National Institute of Standards and Technology) uniform spectral radiance source (accuracy of about 1% in the VIS–NIR [near-infrared] and about 2% in the UV wavelengths at an approximate 95% confidence level) to enable accurate retrievals of aerosol column properties (e.g., $\tau_{aer}$, $n$, $k$, $\omega_0$); and

- stable performance, field-deployable for a long period – the recent expansion globally of Pandora network operation is based on its reliability at various field conditions, and SMART–s is nearly identical to the Pandora instrument except for the spectrometer (extended-range from about 280 to 820 nm, with about 1 nm spectral resolution).

170    As this study is the first attempt to retrieve aerosol properties from the SMART–s near the source region of active and extensive biomass-burning, we summarized the experimental design, instrument characteristics, the radiometric calibration in Section 2. In Section 3.1, we compared the retrieved aerosol property retrievals (e.g., $n$, $k$, $\omega_0$) from the SMART–s with those from collocated AERONET for consistency check. Analyses of temporal variations in aerosols and total column trace gases (i.e., $NO_2$, $H_2O$, $O_3$) retrieved from the SMART–s were described in Section 3.2. We also demonstrated the relationship between

175    aerosol properties and trace gas abundances in this section. Section 3.3 discussed possible applications of the retrieved aerosol parameters for satellite algorithms and preliminary validation/comparison results. Summary and conclusions are given in Section 4.

## 2 Measurements and calibrations

### 2.1 Experimental Setup

180    The ground-based spectroradiometer observations have offered optimum inversion products of the atmosphere for validating/comparing those from collocated space-borne sensors; these are less affected by the surface reflectance and can acquire more informative products from their higher resolution of temporal, spectral (including polarization), and angular measurements. In addition, strategically networked ground-based instruments (e.g., Distributed Regional Aerosol Gridded Observation Networks or DRAGON; Holben et al., 2018) can supplement their limited spatial representation. As a part of the

185    ongoing 7-SEAS, intensive observations were conducted during the pre-monsoon season in April-May 2019 over northern Thailand, specifically the Chiang Mai, Fang, and Doi Angkhang areas. The international collaborators deployed an sUAS (small Unmanned Aerial System) in a rotary-/fixed-wing configuration for ~130 flights to measure boundary-layer profiles of thermodynamics and aerosol size/absorption. A mini-lidar, surface measurements of trace gases, and multiple chemistry samplers are also collocated with the three AERONET and one SMART-s instrument during the campaign. As the SMART-s

190    is located in the middle of large source areas of biomass burning during the season, it can provide useful information on carbonaceous aerosols and key trace gases despite its limited spatial coverage.

   Figure 1 shows an example of spatial distributions of $\tau_{aer}$ at 550 nm (Hsu et al., 2019) from the Deep Blue (DB) aerosol algorithm applied to VIIRS (Visible Infrared Imaging Radiometer Suite) aboard the SNPP (Suomi National Polar-orbiting Partnership) satellite and the corresponding true-color image over Southeast Asia on 30 March 2019, when significant

195    amounts of biomass-burning aerosols prevailed ($\tau_{aer}$ at 550 nm higher than 3.0). The DB aerosol algorithm and its extended family have been applied to various spaceborne spectroradiometers such as AVHRR (Advanced Very High Resolution Radiometer), SeaWiFS (Sea-viewing Wide Field-of-view Sensor), MODIS (MODerate resolution Imaging Spectroradiometer), VIIRS (Hsu et al., 2019), and current advanced multispectral imagers aboard geostationary satellites, enabling the construction

of long-term aerosol climate data records (CDRs). As previously stated, satellite retrievals provide reliable $\tau_{aer}$ as indicated in
Figure 1b; the collocated SMART–s measurement is also presented in the colored circle ($\tau_{aer}$ = 3.1) which shows an excellent
agreement with the DB $\tau_{aer}$ retrievals nearby (mean $\tau_{aer}$ = 2.93 within 10 km of SMART–s). In an attempt to derive more
comprehensive aerosol properties from satellites, previous studies actively utilized UV measurements which are sensitive to
aerosol absorption and vertical profile as well as $\tau_{aer}$ (e.g., Torres et al., 2013; Lee et al., 2021). Accurate aerosol optical models
play a central role in the endeavor. However, due to the lack of a reliable aerosol property database in the UV, they typically
made simple assumptions on the spectral features of aerosols to extrapolate the properties from longer wavelengths or adopted
laboratory measurements. One of the ultimate goals of this study is to contribute to satellite retrievals by providing realistic
aerosol optical models over the study domain, particularly in the UV, which will be discussed in Section 3.3.

The collocated SMART–s and AERONET instruments are deployed on the rooftop of Fang hospital in Fang District,
as shown in Figures 1c and 1d, which is located in a basin of northern Thailand. The population of Fang city is slightly higher
than 116,000 in 2010, with a low level of traffic throughout the year. One of the main roads of the city (Chotana Rd.) is nearby
the building (~50 m). However, we presume the effects of local emissions from the road to the aerosol and $NO_2$ amounts are
weak given the low level of local traffic and that major fractions of the aerosols and trace gases (e.g., $NO_2$) during this season
are emitted from the biomass-burning over this area (Jena et al., 2015; Itahashi et al., 2018; Khodmanee and Amnuaylojaroen,
2021). Figure 1d shows an image of the deployed SMART–s and the Chotana road shown behind. Direct-Sun measurements
of the SMART–s started on 8 March 2019, and acquired additional solar-almucantar scans since 19 March after about 10 days
of stabilization (e.g., for checking stability under the field condition and fine-tuning the alignment in tracking). The
measurements finished on 2 May 2019. The AERONET instrument is installed on the same rooftop, about 5 m away from the
SMART–s. In 2019, surface air temperature at Fang during the pre-monsoon season reached up to about 42°C during the
daytime, and relative humidity gradually increased from March (~30%) to early May (~50%).

## 2.2 Measurements

The SMART–s instrument is originally developed by the Pandora network group at NASA (National Aeronautics
and Space Administration) / GSFC (Goddard Space Flight Center) and the unit (#5) used in this study is registered as Pandora
#48. Most of the components of SMART–s are similar or identical to the standard Pandora instrument except for the
spectrometer. The SMART–s spectrometer is made by the same manufacturer (AvaSpec-ULS2048x64, Avantes, cf.
https://www.avantes.com/ last access on 8 June 2022) as the standard version, but covers a wider spectral range (i.e., 280 –
820 nm) with a lower spectral resolution (~1.0 nm full-width-half-maximum [FWHM] with ~3.7× oversampling). As the
Pandonia Global Network (PGN; Herman et al., 2009, 2015; cf. https://www.pandonia-global-network.org) is utilizing another
type of extended-range spectrometer for their dual-detector system, we refer to this modified Pandora as SMART–s in this
study. The spectrometer utilizes a 2,048 × 64 back-thinned Hamamatsu CCD with a symmetric Czerny-Turner system, and its
spectrum covers $O_2$, $O_3$, $NO_2$, $SO_2$, HCHO, and $H_2O$ gas absorption bands (Herman et al., 2015; Jeong et al., 2018). The optical
head consists of two rotating filter wheels; one includes neutral density (ND) filters, and the other contains bandpass filters

(e.g., U340 and BP300 to block out-of-band (OOB) stray light from the near-UV and VIS wavelengths), ground-fused silica diffuser (diffuser hereafter; for $NO_2$, $H_2O$, and $\tau_{aer}$ retrieval), and an opaque filter for dark current measurements. By combining a variable exposure time (4–4,000 ms) and ND filters, it can measure radiances with a dynamic range up to an order of $10^7$, which enables the direct-Sun and sky-scans using a single detector throughout the day. Note that the field of view (FOV) for direct-Sun observations using the diffuser of this unit is about 2.8°, which is broadened to evenly distribute light passing through the optical head. Sky observations do not use the diffuser to allow more photons to reach the detector, of which the FOV is about 1.5°. The optical head is mounted on a Sun/sky-scanner and is connected to the spectrometer through a fiber-optic cable of 400 μm in diameter. The spectrometer is thermoelectrically controlled to maintain a near-constant temperature but may vary slightly depending on the ambient temperature (typically less than 1°C). The spectrometer temperature is recorded with each measurement to monitor data quality.

The SMART–s algorithm aims to obtain optimal information on aerosols and trace gases with minimum assumptions, which incorporates a series of retrievals from fundamental quantities (i.e., column amounts) to higher-order geophysical parameters (e.g., aerosol physicochemical properties and vertical profiles). Jeong et al. (2018) developed an $\tau_{aer}$ algorithm of the SMART–s based on the spectral Langley method, then compared the retrievals to collocated AERONET measurements at the NASA / GSFC which showed excellent agreements at all overlapping wavelengths (i.e., 330 nm, 380 nm, 440 nm, 500 nm, and 675 nm). Comparisons of the $\tau_{aer}$ from the AERONET and SMART–s during this field deployment are shown in Appendix A (*cf.* Figure A1 and A2). The trace gas algorithm of the SMART–s is designed for a relatively lower spectral resolution (FWHM~1.0 nm) and broader spectral coverage (280–820 nm) compared to the standard Pandora. For retrieving optically thick trace gases such as $O_3$ and $H_2O$, we utilize the spectral Langley method (*see* Appendix B; Jeong et al., 2018), whereas we adopt the spectral-fitting algorithm of Pandora for other optically thin species including $NO_2$ (Herman et al., 2009). Retrieved total columns of $\tau_{aer}$, $O_3$, and $H_2O$ in this study are compared with those from the collocated AERONET (for $\tau_{aer}$ and $H_2O$) and satellite retrievals ($O_3$ from the OMI and $\tau_{aer}$ from the VIIRS) during the measurement period in Section 3.2 and Appendix A. Jeong et al. (2020) developed an OEM-based algorithm using solar-almucantar sky-radiances and total column retrievals (e.g., $\tau_{aer}$, $O_3$, $NO_2$, and $H_2O$) for retrieving spectral $n$, $k$, $\omega_0$ and PSD of aerosols; details of the aerosol-column-property algorithm are in Jeong et al. (2020), key characteristics are included here in Appendix C.

For more than two decades, the AERONET has been supported by NASA to operate a global network of automatic Sun/sky-scanning spectroradiometers for acquiring aerosol information (Holben et al., 1998). The instrument measures discrete channels (i.e., 340, 380, 440, 500, 675, 870, 940, 1,020, and 1,640 nm) of solar irradiance with a 1.2° FOV, which takes about 10 s to scan all spectral filter wheels. The FWHMs of the bandpass filters are 2 nm for 340 nm and 380 nm, 25 nm for 1640 nm, and 10 nm for all other channels, whereas that of the SMART–s is about 1.0 nm for all wavelengths. The estimated uncertainty of $\tau_{aer}$ from the AERONET reference instrument is 0.002, and those from general network instruments are about 0.01 in the VIS–NIR, and are higher (~0.02) in the UV channels (Eck et al., 1999; Giles et al., 2019). Note that the uncertainty of $\tau_{aer}$ from the SMART–s (~0.02 in the VIS-NIR, ~0.03 in the UV) is slightly higher than the AERONET (~0.01 in the VIS–NIR, ~0.02 in the UV) due to the wider FOV (Jeong et al., 2018), which is more susceptible to forward scattering, and

temperature sensitivity of the detector (Kinne et al., 1997). The current AERONET product provides $n$, $k$, $\omega_0$ at 440 nm, 675 nm, 870 nm, and 1020 nm. As the Version 3 algorithm utilizes a vector radiative transfer model (Korkin et al., 2017), it can add a 380 nm channel for the UV-absorbing aerosols (Sinyuk et al., 2020). The recent version of instruments added hybrid sky-scan measurements to allow additional retrievals at solar zenith angle ($\theta_S$) below 50° (Sinyuk et al., 2020). To consider

270 gas absorption, the Version 3 algorithm adopts a monthly climatology (1978–2004) of $O_3$ from the Total Ozone Mapping Spectrometer (TOMS), that of $NO_2$ (2004–2013) from the OMI, and retrieved $H_2O$ using the 940 nm channel measurements (Sinyuk et al., 2020). Further information on the AERONET products is summarized in Giles et al. (2019) and Sinyuk et al. (2020). We utilized the Version 3 and level 2.0 products to compare retrievals from the SMART–s.

## 2.3. A combinative radiometric calibration method for Sun/sky spectroradiometer

275 The standard calibration procedure of the SMART–s includes spectral characterization/registration, linearity and offset correction, radiometric calibration, temperature and flat field correction, and stray light correction (Herman et al., 2015; Jeong et al., 2018; Müller et al., 2020). The PGN also regularly reports updates and standard calibration/validation results on its webpage (https://www.pandonia-global-network.org; Herman et al., 2009, 2015). Utilization of the absolute sky-radiances requires precise radiometric calibrations, which we suggest as a novel and combinative method for Sun/sky spectroradiometers

280 in this study. As this field campaign is the first attempt to deploy a radiometrically calibrated SMART–s using the method, we summarized detailed results of the calibration for the sky scans in this section.

For the initial step of the radiometric calibration, we utilized a uniform spectral radiance light source in the Radiometric Calibration Laboratory (RCL) at NASA / GSFC. The RCL is a class 10,000 cleanroom facility that maintains a number of NIST-traceable integrating sphere sources. The integrating sphere source used in this study is referred to as Grande.

285 Grande is a Spectralon-lined, 101.6 cm diameter integrating-sphere source with a 25.4 cm diameter output aperture, which can generate nine levels of light output. More detailed information and annual calibration reports of Grande are available at https://cf.gsfc.nasa.gov/ or in Gatebe et al. (2007). Figure 2a is an image of the SMART–s mounted in front of the Grande sphere source, and Figure 2b shows Grande's nine levels of spectral output in radiance units. Panel-c presents reported total uncertainty of the Grande sphere's spectral radiance at an approximate 95% confidence level when calibrated using a NIST

290 irradiance standard. Different colors in Figures 2b and 2c indicate the different levels of Grande intensity. Due to relatively low intensity in the UV compared to the VIS-NIR (Figure 2b) in both the NIST-calibrated irradiance standard and the Grande source itself, the calculated uncertainty of the Grande radiance calibration in the UV is higher, and brighter light output provides more accurate intensity as shown in Figure 2c. During the light source calibration, the sensor changed its filter (UV band-pass) to detect the lower intensity of the Grande in the UV which results in the relatively higher uncertainties near 350

295 nm (see Figure 2c). Note that the measurement error covariance matrix of the OEM also accounts for such spectral radiometric uncertainties (*see* Appendix C and Jeong et al., 2020). The SMART–s repeated measurements of Grande ten times for each filter combination (bandpass filters and neutral density filters). The deployment procedure of the SMART–s (or Pandora) includes the organization of the fiber-optic cable and the connection of one end of the cable to the spectrometer (the other end

of the cable is fixed to the optical head). This process can affect light transmittance through the cable. For checking the stability of the fiber-optic cable during deployment, we oriented the cable differently (i.e., re-rolled the cable every time with different diameters or arbitrarily oriented it), then reconnected the ports to the spectrometer at each time of the Grande measurements. Contamination of the front window of the optical head (e.g., raindrops, dew, dust, and insects) is also one of the largest error sources of radiometric measurements. During field deployments, we frequently check the front window and clean it when it is necessary. However, as cleaning the front window also can alter its optical transmittance, we artificially contaminated the front window (by finger, dust, and water) and cleaned it in the same way as we do during field deployments at every time of the Grande observations. In addition, other sources of short- and long-term temporal drift (e.g., spectrometer, filter transmittance) are monitored by pre- and post-mission calibration and the Langley fitting during deployments. Figure 3a shows the results of ten repeat measurements of Grande at its 9-lamps illumination level. The agreement of these repeat measurements indicates good temporal stability in SMART-s responsivity and the Grande output. Figure 3b shows the spectral calibration coefficient calculated from dividing the Grande intensity (Figure 2b) by the average value of the measured voltage count (Figure 3a), while the spectral precision of the Grande calibration (i.e., one standard deviation of ten occurrences of the measurements) is presented in Figure 3c. These results indicate that the precision of the radiometric calibration from the instrument is better than 0.5% in the VIS–NIR channels, and increases at shorter wavelengths to about 0.7% at 330 nm (i.e., the lower limit of the spectral coverage in this study). Radiometric sensitivity of the spectrometer to its temperature ($T_{spec}$) is also tested by controlling the $T_{spec}$, which is less than 0.4% at the entire spectral range for an extreme $T_{spec}$ variability (i.e., $\Delta T_{spec} \sim 3°C$, which is less than 1°C under typical field condition). In general, the uncertainty of the laboratory radiance calibration including the light source and instrument stability is estimated to be better than 2.0% in the VIS–NIR, and 3.0% in the UV at an approximate 95% confidence level.

As typical radiative transfer models (e.g., Spurr 2006; Stamnes et al., 1988) simulate normalized radiances (i.e., radiance divided by solar irradiance as a unit of inverse steradian), reference solar spectrum is also a key parameter of converting raw voltage counts to sky-radiances of the identical unit of forward model calculations. Satellite instruments can directly measure solar irradiance using the same detector with a similar optical path of the earth-reflectance measurements so that a major fraction of calibration uncertainties (e.g., slit function, radiometric coefficient) is canceled out. Although ground-based instruments also measure the solar light using an identical detector as for the sky radiances, they sample the solar irradiance after it passed through the atmosphere. For that reason, their algorithms utilize other sources of the solar spectrum or estimate it from the measurements for the conversion of the sky radiances. The version 3.0 AERONET inversion algorithm utilizes solar irradiance from NOAA's (National Oceanic and Atmospheric Administration) Climate Data Record (Coddington et al., 2016; Sinyuk et al., 2020), and the SKYNET derives a conversion factor of the sky-radiances from direct-Sun measurements based on the solid-view-angle estimation algorithm (Uchiyama et al., 2018a, b). For hyperspectral instruments such as the SMART–s, a combination of the high-resolution solar spectrum and calibrated slit function is a key factor for retrieving the spectral aerosol properties, which is particularly important at wavelengths shorter than 500 nm where spectral variability of the solar irradiance is significant (e.g., see Figure 4a).

In the second step of the radiometric calibration, the SMART–s algorithm estimates the reference solar spectrum by combining direct-Sun measurements, laboratory calibrations, and ancillary solar irradiance data. Raw voltage counts of the sky-scan measurements without bandpass filters can be expressed as follows (e.g., Uchiyama et al., 2018b):

$$V(\lambda) = C(\lambda) \int_{\Omega_{FO}} f(\Omega) I(\lambda, \Omega) d\Omega, \tag{1}$$

where $V$ is the voltage count of a sky-scan measurement, $\lambda$ is the wavelength, $C$ is the sensitivity of the detector to the radiance, $\Omega_{FO}$ is the solid angle of the instrument's field of view (FOV) without bandpass filters, $f$ is the response function of the radiometer's FOV, and $I$ is the sky-radiance. To avoid saturation, the direct-Sun measurements utilize the diffuser with a spectral transmittance of $T_d(\lambda)$ which is measured using the Grande. The measured voltage count of the solar scan can be described as follows:

$$V_S(\lambda) = C(\lambda) T_d(\lambda) \int_{\Omega_S} f(\Omega) I_D(\lambda, \Omega) d\Omega + C(\lambda) T_d(\lambda) \int_{\Omega_{FD}} f(\Omega) I(\lambda, \Omega) d\Omega, \tag{2}$$

where $V_S$ is the voltage count of direct-Sun measurement, $\Omega_S$ and $\Omega_{FD}$ are the solid angle of the Sun and FOV of the SMART–s with the diffuser, respectively. The $I_D$ is direct component of solar measurements. The first term on the right-hand-side of Eq. (2) describes the contribution of direct solar irradiance, and the second term shows the scattered radiance within the FOV. Here we assume that $\Omega_S$ and $\Omega_{FD}$ are not wavelength-dependent within the SMART–s spectral range. We select the Langley calibration dates when $\tau_{aer}$ at 500 nm is less than 0.05 to minimize aerosol impacts and screen cloud-contaminated measurements. In addition, the SMART–s algorithm corrects contribution of Rayleigh scattering of the direct-Sun measurements, which is larger in the shorter wavelength (Jeong et al., 2018). Based on this process and criteria, we assume that the second term of the right-hand-side of Eq. (2) is negligible. However, unscreened thin cirrus cloud may generate diffuse light within the FOV more effectively than the aerosols due to its stronger forward scattering (e.g., Kinne et al., 1997). For plane-parallel solar irradiance measurements ($f = 1$), the $V_S$ can be approximated as:

$$V_S(\lambda) \sim C(\lambda) T_d(\lambda) \int_{\Omega_S} I_D(\lambda, \Omega) d\Omega = C(\lambda) T_d(\lambda) F_{BOA}(\lambda), \tag{3}$$

where $F_{BOA}$ is solar irradiance at the bottom of the atmosphere which can be described as:

$$F_{BOA}(\lambda) = \int_{\Omega_S} I_D(\lambda, \Omega) d\Omega = F_V(\lambda) \left[ \sum_i m_i(\lambda) \tau_i(\lambda) \right]. \tag{4}$$

In Eq. (4), $F_V$ and $\tau$ denote respectively the extraterrestrial solar spectrum and optical thickness of atmospheric constituents, which are derived from the Langley calibration by using the $V_S$ in Eq. (3). The $m$ is the optical air mass of each atmospheric species (e.g., aerosols, clouds, and gases). Then, the last step is normalizing the $F_V$ to a known solar irradiance data for minimizing the remaining systematic calibration error by the following equation:

$$F_{Comb}(\lambda) = \frac{\bar{F}_{Trad}}{\bar{F}_V} F_V(\lambda), \tag{5}$$

where the $\bar{F}_{Trad}$ is the spectral mean value of solar irradiance using the traditional method (i.e., a high-resolution reference solar spectrum convoluted by the instrument's slit function, $F_{Trad}$) and $\bar{F}_V$ is that of $F_V$ at wavelengths between 490–510 nm. The $F_{Comb}$ is the final solar irradiance for the SMART–s algorithm derived by combining the laboratory/Langley calibration,

and reference spectrum (i.e., Coddington et al., 2021 in this study). This spectral window (490–510 nm) is near the middle of the detector, and solar intensity is high with relatively fewer spectral variabilities. In addition, this spectral range avoids strong gas absorptions. By using the spectral shape of the $F_{Comb}$, we expect uncertainties generated from the calibrated slit function are minimized. Despite the Sun/sky measurements undergoing empirical OOB stray light correction (Jeong et al., 2018), the remaining fraction may still be nonnegligible in the shorter wavelengths of UV (particularly wavelengths shorter than about 330 nm; see Figure 7c of Jeong et al., 2018). However, as the $F_{Comb}$ is supposed to be affected by OOB stray light with a comparable degree of the sky-scan measurements due to their similar spectral shape, it can partially cancel out the remaining portion of stray light in the sky-radiances. Figure 4a compares the $F_{Comb}$ to convoluted solar irradiance from Gueymard (2004), Chance and Kurucz (2010), and Coddington et al. (2021), which shows generally good consistency in the entire spectral range of SMART–s. Figure 4b shows an example of spectral radiances using the different solar spectrums in panel-a. The colored symbols in Figure 4b indicate the wavelength node of the aerosol retrieval, which is carefully selected to avoid a strong gas absorption bands and major calibration errors discussed above. Figure 4c depicts relative biases of the sky-radiances using the $F_{Comb}$ to those using Gueymard (2004), Chance and Kurucz (2010), and Coddington et al. (2021) at the wavelength nodes convoluted using the calibrated slit function. The biases are generally smaller than about 2% at wavelengths longer than 500 nm, and are higher in the shorter wavelength up to about 10% at 330 nm. The relatively high discrepancies between the $F_{Comb}$ and the other solar irradiances in the UV are attributable to uncertainties in slit function, remaining OOB stray light, and the $F_V$. Note that uncertainties in the slit function affect $F_{Trad}$ which may be significant in the UV. We estimate the accuracy of $F_V$, which doesn't require spectral convolution, is better than 4% in the UV and 2% in the VIS-NIR based on the accuracy of the spectral $\tau_{aer}$ retrievals. Therefore, the total error of the sky radiance is estimated to be better than 5% in the UV and 3% in the VIS-NIR at an approximate 95% confidence level. The impacts of the different sources of the solar spectrum on aerosol retrievals are discussed in Section 3.

We applied the aerosol retrieval algorithm to the measurements (both direct Sun and solar-almucantar scan) with sufficient amounts of photons within the target spectral range (i.e., from 330 to 800 nm) as very high aerosol loading over the area (e.g., Figure 1b) may result in the low level of voltage counts below the detection limit (e.g., in terms of linearity and noise). Cloud-contaminated direct-Sun spectrums were screened by using their rapid temporal variability and spectral features (i.e., lower Ångström exponent of clouds), which are described in Jeong et al., (2018). Those for the solar-almucantar measurements were removed by checking the horizontal symmetry of the scan (i.e., between clockwise and counter-clockwise half-circle scans) followed by the AERONET strategy (Jeong et al., 2020).

# 3 Results

## 3.1 Comparison with the AERONET

Jeong et al. (2020) applied the SMART–s algorithm to a year-long AERONET Sun/sky measurements in 2016 at Kanpur, India to assess the consistency of methodology. Retrieved volume-size-distribution, V(r), from the SMART–s showed excellent agreements in the fine-mode but minor discrepancies in the coarse-mode due to the different assumptions and constraints between the two algorithms; the AERONET retrieves V(r) at 22 radius nodes over the optically effective range (i.e., from 0.05 μm to 15.0 μm) constrained by smoothness together with the $k$, whereas the SMART–s assumes a bi-modal, lognormal distribution of the number-size-distribution N(r) (Dubovik and King, 2000; Jeong et al., 2020). Spectral $\omega_0$ showed excellent agreements for all wavelengths from 440 nm to 1020 nm, with R ranging from 0.87–0.95 and RMSE (Root-Mean-Squared Error) / MBE (Mean-Bias Error) less than 0.012 during that year-long period. In this section, we performed additional comparisons of aerosol property retrievals from SMART–s to those from collocated AERONET by utilizing their own measurements.

Figure 5 depicts coincident V(r)s from the AERONET and SMART–s, where major fractions of the aerosols are fine-mode smoke particles, and less but nonnegligible portions are coarse particles (e.g., transported dust from Saharan and Thar desert, dry areas of the Indo-Gangetic plain). Retrieved V(r)s from both instruments show generally good agreement, which is consistent with the previous study (Jeong et al., 2020). Regarding that the SMART–s V(r) retrievals showed better agreements to the AERONET when it is applied to the same measurements (Jeong et al., 2020), another major fraction of the discrepancies in Figure 5 is likely attributable to the different types of measurements (e.g., spectral information and radiometric calibration). Note that the SMART–s measurement is not sensitive to aerosols with a radius greater than the optically effective range (~10 μm; *see Figure 10 in* Jeong et al., 2020), and the SMART–s V(r) over this range (see long tails of the blue dashed-line) is mostly generated by assuming the lognormal shape of coarse-mode. Further studies to derive optimal information on aerosol size are underway (e.g., additional parameters for size distribution and/or additional modes). According to theoretical error estimates of the OEM-based algorithm in Jeong et al. (2020), high-spectral resolution of the SMART–s is beneficial for the fine-mode whereas a broader spectral-range of the AERONET is advantageous for both modes under the same level of radiometric accuracy. More detailed comparisons and relevant discussions on the V(r) between the two algorithms are summarized in Jeong et al. (2020).

To better understand and assess the two PSD retrievals at each size-bin, precise evaluation through reliable in-situ measurements, such as aircraft profiles from DISCOVER-AQ (Deriving Information on Surface Conditions from Column and Vertically Resolved Observations Relevant to Air Quality), KORUS-AQ (KORea U.S.-Air Quality) and sUAS, are essential. However, to the best of our knowledge, very limited studies compared the AERONET V(r) to collocated in-situ profile measurements of the PSD. Chauvigné et al. (2016) compared the PSD from AERONET (at 410 m a.s.l. altitude) to in-situ measurements at a higher altitude site in central France (i.e., 1465 m a.s.l.) over a one-year period, which showed relative

425  underestimation of AERONET (~40%). However, the in-situ measurement site in Chauvigné et al. (2016) may not fully represent the total column values, which can be associated with the biases. Schafer et al. (2019) compared the AERONET $V(r)$ to in-situ aircraft profiles (altitude from about 150 to 5000 m, and radius range from 0.03 to 0.5 µm) over Maryland, California, Texas, and Colorado in the United States. They showed that fine-mode PSD parameters derived from AERONET and in-situ aircraft measurements showed generally good agreement (average difference of radius of peak concentration about 0.011 µm,

430  and that of $V(r)$ width about 0.03 µm), whereas differences of the $V(r)$s depend on particle radius and location (*see Figure 5 of* Schafer et al., 2019). The SMARTLabs (Surface-based Mobile Atmospheric Research & Testbed Laboratories; https://smartlabs.gsfc.nasa.gov) is developing a sUAS-based aerosol profiling instrument for collocated measurements of spectral absorption and size distribution, which can provide valuable data for assessments of the PSD and $\omega_0$ from SMART–s. More precise validation/comparison studies are currently underway.

435       Figure 6 compares the retrieved $\omega_0$ from SMART–s and AERONET at overlapping wavelengths (i.e., 440 nm in left panels and 675 nm in right panels) during the measurement period. The $\omega_0$ retrievals from SMART–s in the upper panels of Figure 6 (a and b) used the $F_{\text{Trad}}$, whereas those in lower panels (c and d) utilized the $F_{\text{Comb}}$. In general, all cases showed acceptable agreements with absolute mean-bias errors (MBE) and root-mean-square errors (RMSE) less than 0.02, and R ranging from 0.77 to 0.82. The $\omega_0$ of SMART–s was better correlated with that of AERONET in the shorter wavelength (i.e.,

440  440 nm) due to the higher sensitivity of aerosols from the higher $\tau_{\text{aer}}$. In addition, the $F_{\text{Comb}}$ generated more consistent $\omega_0$ retrievals of SMART–s with the AERONET than those using the $F_{\text{Trad}}$, with the slightly lower RMSE / MBE and higher R. Figure 7a presents mean values of the spectral $\tau_{\text{aer}}$ from SMART–s and AERONET during the measurement period, which showed excellent agreement over the SMART–s spectral coverage. Panel-b of Figure 7 shows those of the spectral $\omega_0$ from the AERONET and SMART–s using different solar irradiances. While both versions of the $\omega_0$ values from SMART–s exhibit

445  strong absorption in the UV, the one using $F_{\text{Comb}}$ showed smoother spectral variability, particularly in the UV, which is selected for the SMART–s retrievals. Slight relative low biases of the $\omega_0$ from SMART–s compared to that from AERONET were also found when the SMART–s algorithm was applied to the same AERONET measurements at Kanpur, India, which are however still within the uncertainty range of the AERONET and SMART–s retrievals (Jeong et al., 2020). Note that the AERONET also provides $\tau_{\text{aer}}$ and $\omega_0$ at longer wavelengths (e.g., 870 nm and 1020 nm), which are not presented in this figure.

450       In Figure 8a, the spectral $n$ of fine- ($n_{\text{f}}$) and coarse-mode ($n_{\text{c}}$) retrieved from the SMART–s shows comparable values with the previous studies in Table 1 ranging from about 1.5 to 1.55 with smooth spectral dependencies in the UV; higher values of $n$ of the coarse-mode were found in the UV whereas those of the fine-mode were lower. These values were slightly higher than those from the collocated AERONET by about 0.01–0.04. Note that the AERONET retrieves a value $n$ for all particle sizes, whereas the SMART–s retrieves each size mode (*see* Appendix C or Jeong et al., 2020), which may result in these

455  differences under the assumption of lognormal size distribution. Regarding the fine-mode dominated smoke aerosols over the site (*cf*. Figure 5), $n$ from the AERONET largely represents the contributions of the fine-mode, and differences between the $n_{\text{f}}$ from SMART–s and $n$ from AERONET are within known uncertainties of the AERONET (Sinyuk et al., 2020) and SMART–s (Jeong et al., 2020). As reported in previous studies and spectral $\omega_0$ retrievals in Figure 7b, the spectral $k$ from SMART–s

increased significantly in the shorter wavelengths, particularly in the fine-mode (*cf.* Figure 8b). However, as the *k* retrievals of each mode have cross-correlated measurement sensitivity, separate analysis of each mode may have relatively larger uncertainties than the case of the *n* retrievals (*cf.* Appendix C and Figure A5). More detailed temporal and spectral analyses of aerosol optical properties from the SMART–s are given in the following sections and scatter plots of the *n*, *k*, and PSD between the two instruments are shown in Appendix D.

Validation or comparison of the UV aerosol properties from the SMART–s remains challenging due to the limited coincident measurements. As discussed for Figure 5, in-situ profile measurements onboard the aircraft or sUAS platforms may provide reliable sets of data for validation. UV aerosol properties retrieved from other collocated instruments (e.g., Mok et al., 2018; Nakajima et al., 2020) can also offer useful data for checking consistency and redundancy. The spectral $\omega_0$ retrievals from SMART–s are mainly determined by the ratio of the sky radiances to the $\tau_{aer}$; highly absorbing aerosols result in the lower level of the sky radiances at a given $\tau_{aer}$. Reliable spectral $\tau_{aer}$ from the SMART–s from 330 nm to 800 nm supports its consistent radiometric performance (e.g., linearity, OOB stray light) over the wavelength range (Jeong et al., 2018). As the SMART–s measures sky radiances using the identical detector but without a diffuser, we expect comparable radiometric accuracy of the spectral sky measurements with relatively higher uncertainties in the UV as discussed in Section 2.3. Note that our best estimate of the accuracy of the spectral radiance is demonstrated in the measurement error covariance matrix, thereby considered in the estimated retrieval error (*see* Appendix C and Jeong et al., 2020).

## 3.2 Temporal variations and relationship between aerosol properties and trace gases

Total column amounts of trace gases and aerosols are basic and essential quantities not only for understanding their amounts and significance with variations over time but also for providing key constraints for higher-order retrievals (e.g., Jeong et al., 2018; 2020). Direct-Sun measurements of the SMART–s, AERONET, and Pandora provide very accurate retrievals of these parameters, which thereby have been used to validate/compare various satellite products over the globe. This section presents general characteristics and temporal trends of the basic quantities retrieved from the SMART–s (i.e., $\tau_{aer}$, total column amounts of $O_3$, $NO_2$, and $H_2O$), then analyzes higher-order retrievals of aerosol properties by comparing them with these quantities.

Figure 9 shows temporal variations of the $\tau_{aer}$ at 500 nm and total column amounts of $NO_2$, $H_2O$, and $O_3$ during the entire deployment period from the SMART–s and AERONET. The red circles in Figure 9a demonstrate aerosol optical thickness at 550 nm from the VIIRS DB product, and those in Figure 9d depict OMI retrievals based on the TOMS (Total Mapping Spectrometer) Version 8.5 algorithm (Bhartia and Wellemeyer, 2002). In general, the $\tau_{aer}$, $H_2O$, and $O_3$ retrievals from SMART–s showed excellent agreements with those from the AERONET, VIIRS, and OMI during the measurement period (see statistical values in Figures 9a, 9c, and 9d). Scatter plots of these parameters are shown in Appendix A (Figure A2).

Itahashi et al. (2018) analyzed high values of $NO_2$ from winter to pre-monsoon seasons over Southeast Asia based on satellite retrievals (i.e., SCanning Imaging Absorption spectroMeter for Atmospheric CHartographY, or SCIAMACHY;

Bovensmann et al., 1999) and model calculations (i.e., Community Multi-scale Air Quality, or CMAQ), and reported that emissions from biomass-burning are attributable to the seasonal variation. They also estimated contributions of biomass-burning emissions to the total column $NO_2$ to be about 28% during 2003–2008 which was up to 58% in March 2004. Khodmanee and Amnuaylojaroen (2021) estimated the contribution of biomass-burning to the $NO_2$ concentration over northern Thailand in March 2014 to be higher than 90% based on the WRF-Chem (Weather Research and Forecasting model with Chemistry) calculations. Another WRF-Chem study suggested that increase of $NO_2$ due to the biomass-burning emission is up to about 60% over Southeast Asia from March to May 2005 (Jena et al., 2015). During the measurement period, a large amount of $NO_2$ from the SMART–s is also accompanied by the high $\tau_{aer}$ at 500 nm with a correlation coefficient (R) of about 0.74, which indicates the common emission sources (i.e., biomass-burning) of aerosols and $NO_2$ during the events as shown in Figure 9a and 9b.

Interaction between atmospheric $H_2O$ and aerosols is one of the primary factors in determining aerosol scattering and absorption properties (e.g., Burgos et al., 2019). Particularly for organic aerosols, molar mass and water content are dominant parameters for characterizing their phase state (e.g., Koop et al., 2011). The moisture-induced phase transition of organic particles from a glassy to a semisolid state also accelerates the uptake of reactive gases in the atmosphere by decreasing viscosity and increasing diffusivity (Shiraiwa et al., 2011). In addition, the condensed water provides a medium for multiphase reactions, thus activating the gas-to-particle conversion of inorganic and organic molecules (Herrmann et al., 2015). The formation of the secondary species through the heterogeneous reactions generates further feedback to the aerosol-atmosphere system by enhancing water vapor absorption and hygroscopicity of aerosol particles (e.g., Tang et al., 2016; Wu et al., 2018). Serving as cloud-condensation-nuclei (CCN), the physicochemical states of aerosols strongly influence cloud microstructure, thereby affecting the radiative properties of clouds, circulation and thermodynamics of the atmosphere (e.g., DeMott et al., 2010). During the intermediate period from dry to monsoon season in 2019, column precipitable $H_2O$ gradually increases from about 1 cm in March to over 3 cm in May as shown in Figure 9c. Note that these mutually interacting species (i.e., $\tau_{aer}$, $NO_2$, $H_2O$) are retrieved simultaneously from the same solar measurements using the ground-fused silica diffuser, which can provide valuable information for further studies. Temporal variation of ozone in Figure 9d is mostly associated with fractions in the stratosphere, which also gradually increased from March (~250 DU) to May (~275 DU). Although the total column $O_3$ is not necessarily relevant to the major topic of this study, it is one of the key constraints of the future algorithm for estimating its tropospheric amounts (Jeong et al., 2020). It should also be noted that the SMART–s covers the $O_3$ Chappuis band in the VIS, which can complement its lower spectral resolution than the standard Pandora for profile retrieval (e.g., Natraj et al., 2011).

Figure 10 presents temporal variations of the $\omega_0$ and $H_2O$ in panel-a, those of the $n_f$ and $n_c$ in panels (b) and (c), respectively. The red circles and black squares show the retrievals at 330 nm and 550 nm, respectively, and their error bars indicate estimated retrieval error ($\varepsilon_{ret}$) based on the optimal-estimation-method (OEM; e.g., Jeong et al., 2016; 2020, *and also see* Appendix C). Note that the $\varepsilon_{ret}$ is calculated for each retrieval, which is an important merit of the OEM for relevant studies. As shown in Figure 10a, the $\omega_0$ and $H_2O$ gradually increased as biomass-burning activities decreased approaching the monsoon season (R = 0.65 between $\omega_0$ and time, and 0.70 between $H_2O$ and time). Interestingly, the correlation between the $\omega_0$ and $H_2O$

was even higher (R = 0.81 for $\omega_0$ at 330 nm) than their temporal trends as shown in this figure. Eck et al. (2013) reported a similar increasing trend of the $\omega_0$ from long-term AERONET and OMI measurements over Southern Africa. They suggested that a trend of decreasing black carbon content (e.g., due to more smouldering rather than flaming combustion) in the aerosol composition during the progress of burning can be a major reason than aerosol growth through aging, coagulation, or

530    hygroscopic swelling, since size-related aerosol parameters such as Ångström exponent (AE) and volume-median-radius did not show a meaningful seasonal trend. However in this study, the retrieved $n_f$ time series in Figure 10b exhibited a slightly decreasing trend over time with reliable retrieval accuracy ($\varepsilon_{ret} = 0.031 \pm 0.015$); this is consistent with the effects of the hygroscopic growth of aerosols, for which $n_f$ values decrease (or get closer to the $n$ of water) as a particle grows by water vapor uptake (e.g., Flores et al., 2012; Valenzuela et al., 2018). Meanwhile, the $n_c$ retrievals are estimated to be highly uncertain ($\varepsilon_{ret}$

535    $= 0.13 \pm 0.04$) due to the very limited information on coarse-mode aerosols as shown by large standard deviations in Figure 10c; the high values of $\varepsilon_{ret}$ are attributable to the SMART–s spectral range, which is narrower than that of the AERONET, and fine-mode dominated smoke aerosols over this area (*cf.* Figure 5).

To further investigate the effects of $H_2O$ on biomass-burning aerosol properties over the experiment regions, we have also examined the relationships of variations in $n_f$, $\omega_0$, weighted-mean-radius of fine-mode and AE with changes in $H_2O$

540    amounts. The upper panels of Figure 11 (panel-a and -b) compare the retrieved $n_f$ and $\omega_0$ to the total column amount of $H_2O$ during the period. As previously discussed, the $n_f$ is negatively correlated with the $H_2O$ (R ranging from −0.57 to −0.61), whereas the $\omega_0$ showed a high positive correlation for both the UV and longer wavelengths (R = 0.74 – 0.81). Lower panels of Figure 11 (panel-c and -d) show the relationship between the weighted-mean-radius of fine-mode and AE (*y*-axis) to column precipitable $H_2O$ (*x*-axis) from the SMART–s. Figure 11c shows two types of weighted mean radius; area-weighted-mean-

545    radius ($r_a$; or often called effective radius) and volume-weighted-mean-radius ($r_v$). The $r_a$ has a proportional relationship with light extinction by particles, which is utilized by aerosol retrievals including the SMART–s algorithm, whereas the $r_v$ has a linear relationship with volume-growth of aerosols. The $r_a$ and $r_v$ of fine-mode are calculated by the following equations, where the radius of 0.01 $\mu m$ and 0.7$\mu m$ are the lower and upper size-limit of the fine-mode particles, respectively:

$$r_a = \frac{\int_{0.01\mu m}^{0.7\mu m} r^3 N(r)dr}{\int_{0.01\mu m}^{0.7\mu m} r^2 N(r)dr} \tag{6}$$

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$$r_v = \frac{\int_{0.01\mu m}^{0.7\mu m} r^4 N(r)dr}{\int_{0.01\mu m}^{0.7\mu m} r^3 N(r)dr} \tag{7}$$

As shown in Figure 11c, both the $r_a$ and $r_v$ showed a positive correlation with $H_2O$ (R = 0.42 and 0.56, respectively), which is higher for the $r_v$. The short wavelength range AE is another good indicator of the fine-mode particle size (Reid et al., 1999; Eck et al., 2001), which is known to have a negative correlation with the size. The AE and absorbing AE (AAE) are calculated as follows:

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$$AE = -\frac{\ln[\tau_{aer}(\lambda_1)/\tau_{aer}(\lambda_2)]}{\ln(\lambda_1/\lambda_2)}, \tag{8}$$

$$AAE = -\frac{\ln[\tau_{abs}(\lambda_1)/\tau_{abs}(\lambda_2)]}{\ln(\lambda_1/\lambda_2)}, \tag{9}$$

where $\tau_{abs}$ is the absorbing aerosol optical thickness (Equation 10):

$$\tau_{abs} = (1 - \omega_0)\tau_{aer}. \tag{10}$$

The AE calculated using both pairs of wavelengths showed a negative correlation with the $H_2O$ (R ranging from −0.46 to −0.38).

NO$_2$ in the smoke plume may contain information on the degree of aging after emission; the lifetime of $NO_2$ is short (typically less than a few hours) and the brown carbon uptakes $NO_2$ by photochemical processes, thus $NO_2$ concentration likely decreases through the aging processes of the smoke plumes (e.g., Laskin et al., 2015). Therefore, lower values of the $n_f$ (Figure 12a) and higher values of the $\omega_0$ (Figure 12b) at lower amounts of $NO_2$ (likely related to the aged plume) support that their temporal variations may also be associated with accelerated aging processes of smoke aerosols by the increased $H_2O$. Even with the lower correlations, the size-related parameters in Figures 12c and 12d also indicate that smoke aerosols are likely growing under lower amounts of $NO_2$ during the pre-monsoon period; $NO_2$ showed a negative correlation with $r_v$ and positive correlation with AE. However, note that the correlations of $n_f$, $\omega_0$, $r_v$ and AE with $NO_2$ also can be attributed to independent trends of emission and photochemical reactions of $NO_2$, which are not necessarily associated with the aerosol aging processes. More sophisticated studies combining model simulations and intensive measurements may be able to clarify relationships between $NO_2$ and aerosol properties. In general, the results in Figures 10–12 suggest that aerosol aging processes including hygroscopic growth also can be a critical parameter impacting temporal trends of aerosol optical properties over this area, in addition to previously suggested factors over South Africa (e.g., change of burning sources and conditions in Eck et al., 2013). Overall, such comparisons suggest the potential benefit of simultaneous measures of trace gases and aerosols for understanding atmospheric physicochemical processes.

**3.3 UV-VIS-NIR aerosol absorption properties for satellite algorithms**

The continuous UV-NIR aerosol property information retrieved from SMART-s can be useful not only for validating the satellite aerosol products, but also for fine-tuning appropriate aerosol models used in the satellite aerosol retrieval algorithm. The operational aerosol algorithm of OMI (OMAERUV; Torres et al., 2013) utilizes radiances at 354 nm and 388 nm for retrieving the $\tau_{aer}$ and $\omega_0$. They assumed a ratio of $k$ at 354 to that at 388 nm as 1.2 for smoke aerosols to account for the spectral absorption effects of organic carbon (Jethva and Torres, 2011; Jeong et al., 2016). They derived a monthly climatology of aerosol layer height (ALH) from observations by CALIOP (Cloud-Aerosol Lidar with Orthogonal Polarization) as ancillary data for the OMAERUV algorithm (Torres et al., 2013). Also, recently the improved ASHE (Aerosol Single-scattering albedo and Height Estimation; Lee et al., 2021) algorithm combines UV measurements from the OMPS-NM (Ozone Mapping and Profiler Suite Nadir Mapper) onboard the SNPP with the VIIRS radiances to provide retrieved $\omega_0$ and ALH products as part of the VIIRS version 2 DB aerosol CDRs. For spectral dependences of the $\omega_0$ in the UV, they assumed AAE as 2.0 between 340 nm and 412 nm. Figure 13 shows an example of $\omega_0$ retrievals at 340 nm (panel-a), 378 nm (panel-b) and 550 nm (panel-c) from the ASHE algorithm on the same day as Figure 1 (30 March 2019). Since the ASHE algorithm only performs retrievals when $\tau_{aer} > 0.5$ and UVAI (UV Aerosol Index) > 0.7, the spatial coverage of $\omega_0$ in Figure 13 is reduced

590    compared to that of $\tau_{aer}$ shown in Figure 1. The colored circles in this figure depict the collocated $\omega_0$ retrievals from SMART–s. In general, the values of $\omega_0$ retrieved from ASHE nearby the measurement site was comparable to that from the SMART–s, with its broader spatial coverage throughout the smoke aerosols of high $\tau_{aer}$ (*cf.* Figure 1b).

Temporal variations in $\omega_0$ from the ASHE and SMART–s over the measurement site at overlapping wavelengths demonstrate their reasonable consistency, where both the retrievals indicate increasing trends during the measurement period

595    with a higher temporal resolution of the SMART–s (*cf.* Figure 14a). The ASHE retrievals are available until 18 April 2019, since UVAI over the site decreased lower than the ASHE criteria likely due to the increased $\omega_0$ over the period. Figure 14b compares the collocated $\omega_0$ from ASHE and SMART–s during the measurement period, which showed good agreement between these two at ASHE's retrieval wavelengths. Higher correlations between the ASHE and SMART–s were found in the shorter wavelengths (R = 0.79 at 340 nm) due to the higher sensitivity of UV radiances to aerosol absorptions. The MBE and

600    RMSE were lower than 0.02 for all wavelengths. Note that the $\theta_S$ at Fang near the overpass time of VIIRS is small, whereas the SMART–s measures almucantar radiances $\theta_S$ from 40 to 75°. The average time difference between the ASHE overpass time and the closest SMART–s retrieval during the period was about 3 hours, and only samples with a time difference of fewer than 3 hours are shown in this figure. However, the difference of time between the ASHE and SMART–s still may affect the comparison in Figure 14, which can be improved by refinements of the scan strategy (e.g., hybrid scan of the AERONET

605    version 3.0) for extending retrieval criteria of the $\theta_S$.

AERONET sites in regions affected by biomass-burning smoke typically had AAE (440–870 nm) ranging from 1.0 to 2.0 with values closer to 1.0 indicating a greater contribution of black carbon and AAE near 2.0 indicating increased organic-to-black carbon ratios (Giles et al., 2012). Figure 15a presents temporal variations in the AAE using different wavelength pairs calculated from the SMART–s $\tau_{aer}$ and $\omega_0$. The AAEs calculated from UV wavelength and 550 nm pairs were comparable to

610    the assumed value of the ASHE algorithm (2.04 ± 0.27 for 340–550 nm pair, 1.94 ± 0.33 for 378 nm and 550 nm pair) with non-negligible temporal variabilities ranging from about 1.3 to about 2.6. The AAE calculated using the 340–412 nm wavelengths pair, which is the actual pair for the ASHE inversion, showed much higher values (2.69 ± 0.35). Such discrepancies between the SMART–s retrievals and assumptions of aerosol properties in the ASHE algorithm may propagate to the differences in Figures 13 and 14. However, the retrieval errors are contextual, meaning that other error sources (such as

615    uncertainties in the retrieved $\tau_{aer}$, assumed size distribution in the aerosol optical models, etc.) can also contribute to the discrepancies, which makes it difficult to quantify the contribution of the AAE uncertainties to the retrieval errors. A longer-term data record in the UV is therefore highly desired. Figure 15b shows temporal variation in the ratio of $k_f$ at 354 nm to 388 nm (1.17 ± 0.05), which is in good agreement with the assumed value of the OMAERUV algorithm (i.e., 1.2 for smoke aerosols), but with significant temporal variabilities. Both panels in Figure 15 suggest that current assumptions of the UV

620    aerosol properties in OMAERUV and ASHE algorithms are generally good approximations. However, temporal and spectral variabilities of the aerosol optical properties, which are presented throughout this section, also emphasize the importance of realistic aerosol models in the UV for further improvements of satellite algorithms. High spectral-resolution of the aerosol optical properties covering the UV can also benefit recently launched or upcoming hyperspectral satellite sensors targeting

atmospheric composition (e.g., Chance et al., 2019; Ingmann et al., 2012; Kim et al., 2020). The National Institute of Environmental Research of South Korea recently started to deploy standard Pandoras and a few SMART–s units over Asia to validate GEMS (Geostationary Environment Monitoring Spectrometer; Kim et al., 2020) aerosol and trace gas products as well as to improve the satellite algorithms. Thus, deployments of networked SMART–s can contribute to comparing/validating spatiotemporal variations of aerosol $\omega_0$, which is a key parameter for understanding their aging processes and interaction with other environmental conditions (e.g., terrain and meteorology; *cf*. Figure 13 for an example of the spatial variability). These ground-based measurements will provide important long-term records of UV aerosol properties at multiple strategic sites over Asia.

## 4 Summary and Conclusions

The SMART–s was deployed during the pre-monsoon season in northern Thailand to perform direct-Sun and sky-radiance measurements near biomass-burning sources. In this study, we summarized the detailed radiometric calibration procedures and results. To optimize solar irradiance for the radiometric conversion, we combined the Langley and NIST-traceable integrating sphere calibration data with the high-resolution reference spectrum from Coddington et al. (2021). We estimate the total uncertainties in the sky-radiance measurements are about 5% in the UV and better than 3% in the VIS–NIR wavelengths. The total column amount of the $\tau_{aer}$ and $H_2O$ from SMART–s showed excellent agreements with those from collocated AERONET measurements (R = 1.0 and 0.98, respectively). Total column $O_3$ retrievals from the OMI showed good consistency with those from the SMART–s (R = 0.95, RMSE and MBE less than 3.6 DU). During the measurement period from mid-March to early-May in 2019, the $\tau_{aer}$ was mostly large (frequently exceeding 2.0 at 500 nm) and strongly correlated with total column $NO_2$ (correlation coefficient, R = 0.74), likely due to the high emissions of biomass-burning smoke. The $\omega_0$ from SMART–s and AERONET at overlapping wavelengths (i.e., 440 nm and 675 nm) showed acceptable agreements within uncertainties of these instruments (R = 0.79 – 0.81 with RMSE and MBE less than 0.015). The SMART–s retrievals showed good agreements of fine-mode V($r$) with those from the AERONET, which is dominated by the smoke aerosols during the period. The spectral $\omega_0$ of smoke aerosols showed an abrupt decrease in the UV consistent with the understanding of absorption by carbonaceous aerosols.

One of the major merits of the simultaneous retrieval of trace gases and aerosols from SMART–s is that it allows informative analysis of physicochemical interactions in the atmosphere. Our analyses comparing the trace gases (i.e., $H_2O$ and $NO_2$) and aerosol properties (e.g., $\omega_0$, $n$, $r_a$, and $r_v$) suggest that aerosol aging processes including hygroscopic growth can be a critical factor affecting temporal trends of aerosol optical properties during the pre-monsoon period over northern Thailand. Firstly, the $\omega_0$ and column precipitable $H_2O$ gradually increased together as it approached the monsoon season, and the correlation between $\omega_0$ and $H_2O$ was generally higher (R = 0.74 – 0.81) than their temporal trends (R = 0.65 for $\omega_0$ and 0.70 for $H_2O$). Secondly, the area/volume-weighted radius of fine-mode also showed a positive correlation with the $H_2O$ (R = 0.42 and 0.56, respectively). The third result supporting the conclusion is that the $n_f$ from SMART–s showed a negative correlation with the total column $H_2O$ (R = −0.61 for 330 nm and –0.57 for 550 nm), which is supposed to decrease (or get close to the $n$

of water about 1.33) as the particle uptake water vapor. However, in this study, measurements are performed at only one location during a limited period, while characteristics of smoke aerosols can vary significantly by region and time due to different fuel types, combustion efficiency, and aging processes. A longer period of measurements from multiple sites may help to clarify/understand such relationships.

The UV radiances are useful for satellite algorithms to retrieve higher-order aerosol parameters (e.g., single-scattering albedo and aerosol layer height), as these are sensitive to aerosol absorption and vertical profile. However, due to the lack of information contents and reliable aerosol model covering the UV, the algorithms typically assumed the spectral optical properties of aerosols or adopted from laboratory databases (e.g., Jethva and Torres, 2011; Lee et al., 2021). The retrieved UV aerosol properties from the SMART–s showed generally good agreements with the current assumptions of the ASHE and OMAERUV algorithms, thereby reasonable consistency of the $\omega_0$ from between the SMART–s and ASHE retrievals (R = 0.73 – 0.79, MBE and RMSE less than 0.02). However, temporal and spectral variabilities of aerosol absorption properties (e.g., $\omega_0$ and $k$) in the UV emphasize the importance of a realistic aerosol model for further improvements of satellite retrievals. Recently launched (e.g., GEMS) and upcoming hyperspectral sensors on the geostationary orbit (e.g., Tropospheric Emissions: Monitoring of Pollution; Chance et al., 2019, and Sentinel–4; Ingmann et al., 2012) aim to derive diurnal variations of trace gases and aerosols. The SMART–s can provide key parameters of trace gases and aerosols for constraining and validating satellite algorithms from its higher temporal resolution retrievals. Further improvements of algorithms and instruments for acquiring better information content and more reliable products (e.g., adding spectral polarization measurements, and refinement of scanning strategy such as adding a hybrid scan of the AERONET) are currently underway.

**Appendix A. Scatter plots between the SMART–s direct Sun retrievals and AERONET/satellite observations.**

Direct-Sun retrievals from the SMART–s are analyzed and compared with other sources of retrievals in Section 3. However, this section additionally shows scatter plots of $\tau_{aer}$, total column $H_2O$ and $O_3$ from the SMART–s with those from AERONET and satellite measurements during the campaign period for validation and checking consistency. As reported by Jeong et al. (2018), spectral $\tau_{aer}$ from the SMART–s showed excellent agreements with the AERONET at all overlapping wavelengths as shown in Figures A1 and A2-a, which is also in good agreement with the VIIRS DB product (Figure A2-b). The total precipitable water vapor ($H_2O$) and total column ozone from the SMART–s also showed excellent agreements with the AERONET and OMI as shown in Figures A2-c and A2-d.

**Appendix B. SMART–s ozone and water vapor retrieval algorithm**

The spectral Langley method (Jeong et al., 2018) retrieves spectral $\tau_{aer}$ by subtracting gas optical thickness ($\tau_{gas}$) from the total optical thickness ($\tau_{tot}$); the SMART–s observation is beneficial for this procedure as it measures spectral features of

the $\tau_{gas}$. The SMART–s retrieves optically thick trace gases (i.e., $O_3$ and $H_2O$) by using a similar method but with narrower fitting windows of the trace absorption bands (i.e., 315–335 nm for $O_3$ and 550–680 nm for $H_2O$). Examples of the fitting results for the $O_3$ and $H_2O$ are shown in Figure A3 and A4, respectively, and intercomparison results with other data are in

690 Appendix A during the campaign. The fitting model for $O_3$ retrieval includes linear polynomials, $O_3$ and $SO_2$ cross-sections, and Raman spectrum, whereas that for the $H_2O$ utilizes linear polynomials, $H_2O$, $O_3$, and $O_2$ cross-sections. More detailed design and calibration procedures for the spectral Langley method are demonstrated in Jeong et al. (2018) and measurements for long-term validation/comparison results are currently underway, which will be reported in the following paper (Jeong et al., 2022; *to be submitted*).

695 **Appendix C. SMART–s aerosol inversion algorithm**

Spectral bands for the aerosol property retrieval are carefully selected to avoid strong absorption by the gases and to efficiently obtain maximum information on aerosols. For example, we avoided major $O_2$-A, $O_2$-B, $O_2$-$O_2$, and $H_2O$ bands where each corresponding $F_{Comb}$ is highly uncertain (see Figures 4a and 4b). Nodes of relative azimuth angle ($\phi_r$) for the solar-almucantar scan and those of wavelengths for the retrieval are summarized in Table A1 with other parameters. VLIDORT

700 (linearized pseudo-spherical vector Discrete Ordinate Radiative Transfer) code generates the full Stokes' parameters and analytic weighting function of atmospheric and surface variables including aerosol properties (Spurr, 2006; Spurr et al., 2012, Spurr and Christi, 2014). The state vector ($\mathbf{x}$; a vector with elements of retrieval parameters and control variables to fit the measurements using a forward model) consists of fine and coarse mode $n$ and $k$, five parameters of the PSD, two aerosol layer-height parameters, and spectral surface albedo. The algorithm assumes aerosol number-PSD as a bi-modal lognormal shape as

705 follows:

$$N(r) = \frac{F_{num}}{\sqrt{2\pi}\ln\sigma_f}\frac{1}{r_f}\exp\left[-\frac{1}{2}\left(\frac{\ln r - \ln r_f}{\ln\sigma_f}\right)^2\right] + \frac{(1-F_{num})}{\sqrt{2\pi}\ln\sigma_c}\frac{1}{r_c}\exp\left[-\frac{1}{2}\left(\frac{\ln r - \ln r_c}{\ln\sigma_c}\right)^2\right], \qquad (A1)$$

where $N(r)$ is the number-size distribution, and the $r_f$ and $r_c$ are fine- and coarse-mode mean radius, and the $\sigma_f$ and $\sigma_c$ are geometric standard deviation of each mode. The $F_{num}$ is number fraction of fine mode. However, the assumed bi-modal lognormal shape may not adequately represent the true fine- or coarse-mode distribution in some cases such as volcanic

710 eruptions, aerosol aging, and cloud processing (Eck et al., 2010; 2013; 2018). The aerosol extinction profile is assumed to be a Gaussian shape as follows:

$$ALH(z) = W\frac{e^{-h(z-z_p)}}{\left[1+e^{-h(z-z_p)}\right]^2}, \qquad (A2)$$

where $ALH(z)$ stands for the aerosol layer height (i.e., vertical profile of aerosol extinction), and the $W$ is normalization factor. $z_p$ is the peak height and $h$ is the vertical dispersion parameter of the Gaussian profile shape. *A priori* information of the $z_p$ and

715 $h$ is extracted from climatology of reanalysis data (e.g., Modern-Era Retrospective analysis for Research and Applications, Version 2; Gelaro et al., 2017), which is however not sensitive to the solar almucantar measurements. In this study, we assumed

the surface reflectance ($\rho$) is Lambertian, of which *a priori* data are obtained from merged satellite measurements. More details of the parameters and design of the algorithm can be found at Jeong et al. (2020).

Averaging kernel (**A**) of the OEM is a useful matrix for understanding the information content of a set of measurements and inversion method, of which elements show the sensitivity of retrievals to the true state (e.g., Rodgers, 2000; Jeong et al., 2020). The **A** is defined as follows:

$$\mathbf{A} = \mathbf{GK} = \frac{\partial \hat{\mathbf{x}}}{\partial \mathbf{x}}, \tag{A3}$$

where **G** is the gain matrix for representing the sensitivity of retrievals to the measurements and **K** is the weighting function matrix of which the elements are partial derivatives of each measurement with respect to the state vector. In the Equation (A3), each element of the **A** characterizes how the retrieval ($\hat{\mathbf{x}}$) responses to the true state (**x**). Diagonal elements of **A** ($D_A$) indicate the sensitivity of each retrieval parameter using a set of measurements and an inversion method, whereas off-diagonal elements of an $i^{th}$ row ($R_A$) demonstrate retrieval errors of $\mathbf{x}_i$ by cross-correlation with other parameters or by insufficient information contained in the measurements. Therefore, for an ideal inversion with an observing system, its **A** is close to an identity matrix (Rodgers, 1990). More discussion of **A** for the SMART–s is summarized in Jeong et al. (2020). Figure A5 shows an example of **A** from the SMART–s retrieval at Fang on 10 April 2019, when fine- and coarse-mode volume fractions were comparable and $\omega_0$ was about 0.87 with $\tau_{aer}$ about 1.06 at 440 nm. Figure A5a presents the whole **A**, and the panel-b zooms to the PSD (indices from 1 to 5) and ALH parameters (indices of 6 and 7) as indicated as a green square in panel-a. The $D_A$ of PSD parameters (diagonal elements of 1–5; order of $r_f$, $\sigma_f$, $r_c$, $\sigma_c$, and $F_{num}$) are close to one, which shows their sufficient retrieval sensitivity from the measurement. Particularly, $R_A$ of the $r_f$ and the $F_{num}$ (see off-diagonal elements of the $1^{st}$ and $5^{th}$ rows in Figure A5b) have small absolute values, whereas those of $\sigma_f$, $r_c$ and $\sigma_c$ (rows from 2 to 4) are relatively high. As solar-almucantar measurements are not sensitive to the vertical profile of aerosols, $D_A$ of the ALH parameters ($z_p$ and $h$ of indices 6 and 7) are low with their relatively higher values of the $R_A$ for $k_f$ and $k_c$. $D_A$ of the $n_f$ and $n_c$ show their sufficient retrieval sensitivity at all wavelengths. In addition, the $R_A$ of $n_f$ for $n_c$ and that of $n_c$ for $n_f$ were negligible which suggests that their retrieval sensitivity for each mode is independent and can be retrieved separately. However, the $R_A$s of $n_f$ and $n_c$ demonstrate that they are also affected by the $k_f$ and $k_c$. On the contrary, $R_A$ of the $k_f$ and $k_c$ are low at most of the other parameters (i.e., $n$, PSD, ALH, and $\rho$), whereas the $R_A$ of $k_f$ at state vector $k_c$ (or *vice versa*) shows high values of diagonal elements indicating retrieval sensitivity across the fine- and coarse-mode; therefore, we analyzed the $\omega_0$ for both modes in this study. As well-recognized, retrieval sensitivity of the $\rho$ is negligible which results in low values of the whole rows of $\rho$ in Figure A5. We also limit the retrieval range of solar zenith angle ($\theta_S$) up to 75° to minimize the effects of surface reflectance at high $\theta_S$ and viewing zenith angle ($\theta v$). Note that the **A** varies significantly depending on cases, and retrieval errors due to the interferences between different parameters and/or lack of information content are considered by the error estimation method of the OEM (Rodgers, 2000; Jeong et al., 2020).

One of the important merits of the OEM is its theoretical formulations of retrieval errors, which are classified into four categories (Rodgers, 1990; 2000; Jeong et al., 2016; 2020): the smoothing error ($\varepsilon_s$), retrieval noise ($\varepsilon_m$), forward model

750 error ($\varepsilon_r$), and model parameter error ($\varepsilon_f$). As the **x** contains the most dominant parameters of aerosols and quantifying uncertainties in the radiative transfer model is challenging, we neglect the ($\varepsilon_r$) and ($\varepsilon_f$). We defined retrieval error ($\varepsilon_{ret}$) as the square root of the sum of squared $\varepsilon_s$ and $\varepsilon_m$, which represents the minimum uncertainty of the SMART-s aerosol inversions. Detailed descriptions and discussions of the $\varepsilon_{ret}$ are summarized by Jeong et al. (2020).

**Appendix D. Comparison of aerosol optical properties from the SMART–s and AERONET**

755 Coincident retrievals of the $\omega_0$, $n$, $k$, and V(r) are compared and discussed in Figures 5 to 10. The remaining scatter plots of these inversion parameters of the SMART–s and AERONET are shown in this section. Figure A6 compares aerosol inversion products from the SMART–s and AERONET (Version 3, Level 2.0). The upper panels compare the real part of the refractive index at (a) 440 nm and (b) 675 nm, and the middle panels are for the imaginary part. The lower panels compare (e) area- and (f) volume-weighted-mean-radii. Larger discrepancies between the SMART–s and AERONET were found for the $n$

760 due to the lack of information content and different assumptions of the algorithms, whereas higher consistency appeared for the k as discussed for the $\omega_0$. High agreements of the $r_s$ and $r_v$ were found between the two instruments as shown in Figures A6e and A6f (R = 0.86 with MBE and RMSE less than 0.016 $\mu$m).

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**Code/Data availability**

The SMART–s data are also described at https://earth.gsfc.nasa.gov/climate/instruments/smartlabs, and available
775 from the authors Jeong and Tsay. The AERONET data are available at the https://aeronet.gsfc.nasa.gov website.

**Author contribution**

The first and corresponding author (Ukkyo Jeong) led overall algorithm development, instrument calibration and manuscript writing. Si-Chee Tsay supervised the overall activities for this study and is also in charge of the campaign for this study. Christina Hsu and Jaehwa Lee provided the satellite retrievals for comparing the aerosol optical properties analyzed in this study. David Giles and Brent Holben provided critical revision for aerosol retrievals as well as the overall manuscript. John Cooper and James Butler provided the NIST-traceable light source for radiometric calibration, and also reviewed instrument calibrations. Robert Swap was in charge of the Pandora network operation for the campaign and also supported instrument maintenance. Sheng-Hsiang Wang and Comporn Chantara supported the local operation of the instruments throughout the campaign periods and also reviewed the final manuscript. Hyungkee Hong, Donghee Kim and Jhoon Kim are in charge of the Asian Network of the Pandora and the Geostationary Environment Monitoring Spectrometer, which provided critical insight for this study. They also provided critical revisions for the manuscript.

**Competing interests**

The authors do not have any competing interests.

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**Table 1: Examples of reported real part refractive index of biomass burning and dust aerosols from previous and current study.**

| Reference | Aerosol type | Wavelengths | Real part of the refractive index |
|---|---|---|---|
| Kim et al. (2010) | Secondary organic aerosols | 670 nm | 1.38 – 1.61 |
| Liu et al. (2013) | Secondary organic aerosols | 220 – 1200 nm | 1.48 – 1.58 |
| Sheperd et al. (2018) | Urban, remote, wood smoke | 460 – 760 nm | ~1.58 for wood smoke aerosols<br>1.47 – 1.52 for urban and remote aerosols |
| Sumlin et al. (2018) | Brown carbon aerosols | 375 nm, 405 nm, 532 nm, 1047 nm | 1.5 – 1.7 |
| Biagio et al. (2019) | Dust aerosols | 370 nm, 470 nm, 520 nm, 590 nm, 660 nm, 880 nm, 950 nm | 1.48 – 1.55 |
| Womack et al. (2021) | Biomass burning | 360 – 720 nm | 1.55 – 1.6 |
| This study | Major fraction of biomass-burning aerosols mixed with minor fraction of dust particles | 330 – 780 nm | $1.53 \pm 0.03$ for fine mode<br>$1.51 \pm 0.02$ for coarse mode |

1255  **Table A1: Relative azimuth angles, wavelength node, sources of measurement error covariance matrix, and parameters of state vector of SMART–s algorithm.**

| Algorithm parameter | Description |
|---|---|
| Relative azimuth angles | 3.0°, 3.5°, 4.0°, 5.0°, 6.0°, 7.0°, 8.0°, 10.0°, 12.0°, 14.0°, 16.0°, 18.0°, 20.0°, 25.0°, 30.0°, 35.0°, 40.0°, 45.0°, 50.0°, 60.0°, 70.0°, 80.0°, 90.0°, 100.0°, 120.0°, 140.0°, 160.0°, and 180.0° |
| Wavelength node | 330 nm, 340 nm, 350 nm, 360 nm, 370 nm, 380 nm, 390 nm, 400 nm, 410 nm, 440 nm, 455 nm, 490 nm, 520 nm, 540 nm, 555 nm, 580 nm, 610 nm, 640 nm, 675 nm, 750 nm, and 778 nm |
| Sources of measurement error covariance matrix | Estimated from Langley and laboratory calibration results |
| State vector | Lognormal parameters of aerosol number-size distribution for fine- and coarse-mode, number-fine-mode fraction, two parameters of gaussian vertical profile shape of aerosols, spectral complex refractive indices of fine- and coarse-mode, spectral surface reflectance |

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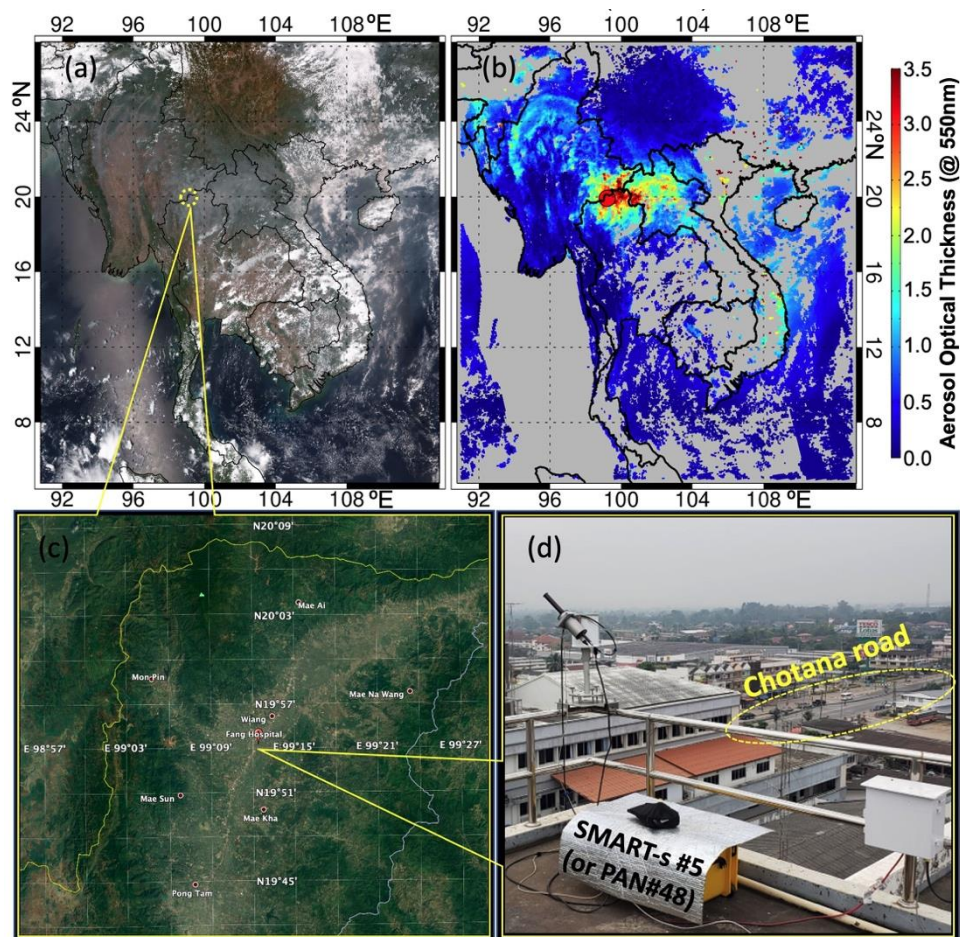
**Figure 1: (a) True color image and (b) DB (Deep Blue) $\tau_{aer}$ product from VIIRS (Visible Infrared Imaging Radiometer Suite) onboard SNPP (Suomi National Polar-orbiting Partnership) on 30 March 2019. Yellow circle in panel (a) is the location of Fang, Thailand, and colored circle of panel (b) shows collocated $\tau_{aer}$ retrievals from the SMART–s. Panel (c) indicates the location of measurement site at the rooftop of Fang hospital, Thailand (19.91°N latitude and 99.21°E longitude, 480 m above sea level; the map is extracted from https://google.com/maps/; © Google Maps 2021). Panel (d) is an image of deployed SMART–s (Pandora#48) taken on 8 March 2019. The Chotana road (marked in yellow) is one of the major streets at this area and is about 50 m away from the site.**

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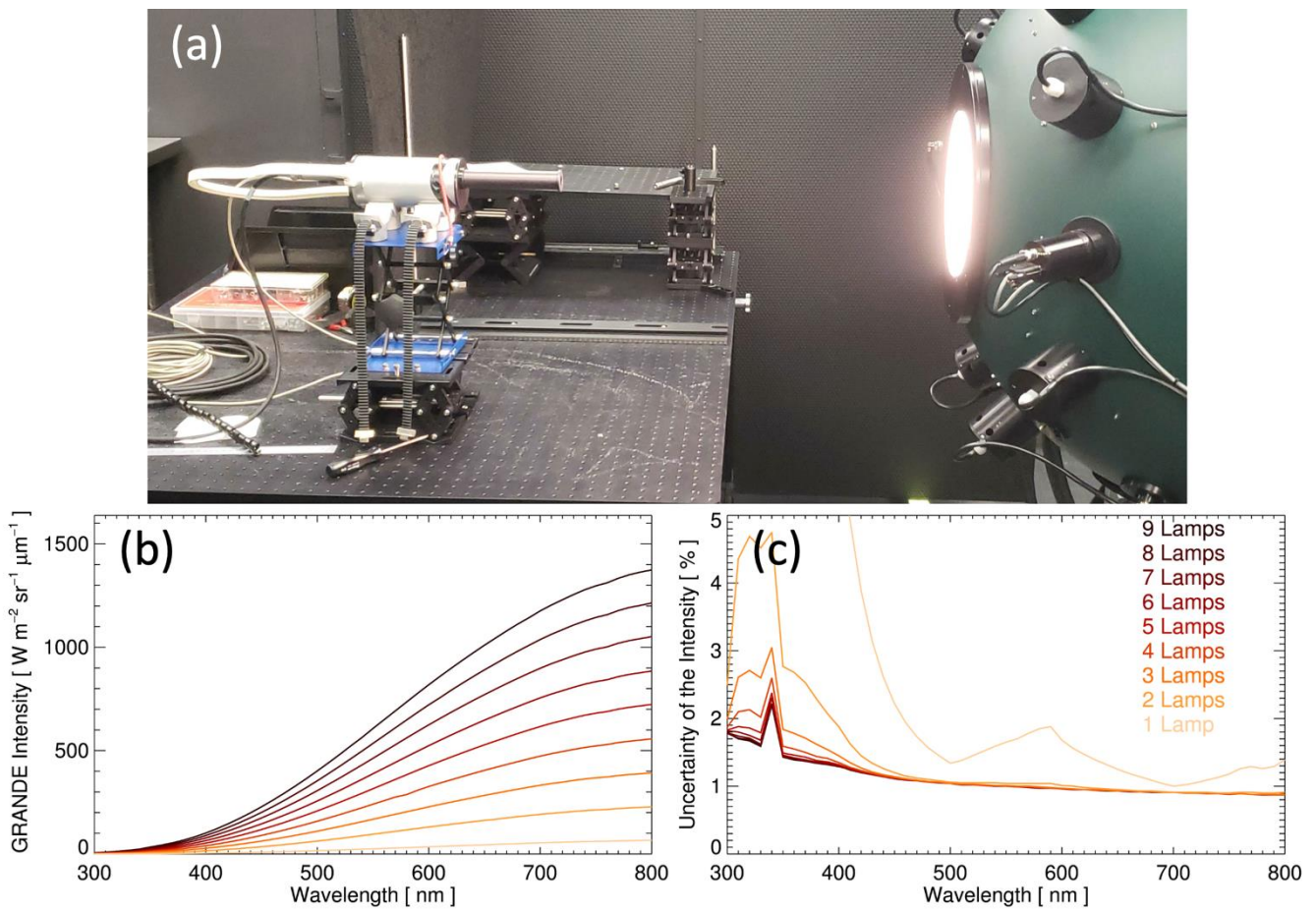
**Figure 2: (a) Image of SMART–s calibration using the NIST-traceable light source (Grande) at the Radiometric Calibration Laboratory, NASA Goddard Space Flight Center. Panel (b) shows spectral radiance of the Grande in 300 – 800 nm, and panel (c) presents its reported uncertainty. Different colors in (b) and (c) indicate nine levels of the Grande radiance. The relatively higher values of uncertainty near 350 nm in panel (c) are due to filter change of sensor during the light-source calibration.**

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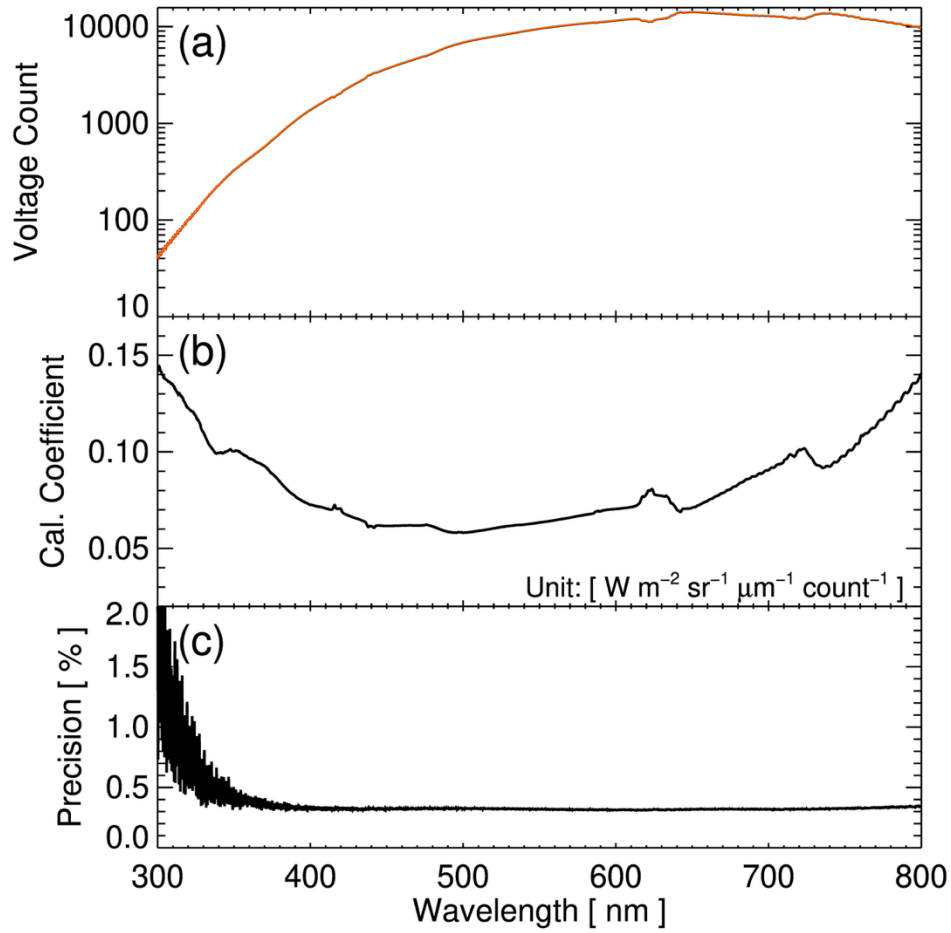
**Figure 3: (a) An example of nine-lamps Grande voltage count measurements from SMART–s (Pandora#48) without neutral density or band-pass filters. Different colors indicate ten times of repetitions, which overlap almost on top of each other. Panel (b) is calibration coefficient, which is calculated from dividing known Grande intensity by average value of the measured voltage count. (c) is precision of the calibration coefficient, which is estimated by calculating one standard deviation of the ten times of repetitions.**

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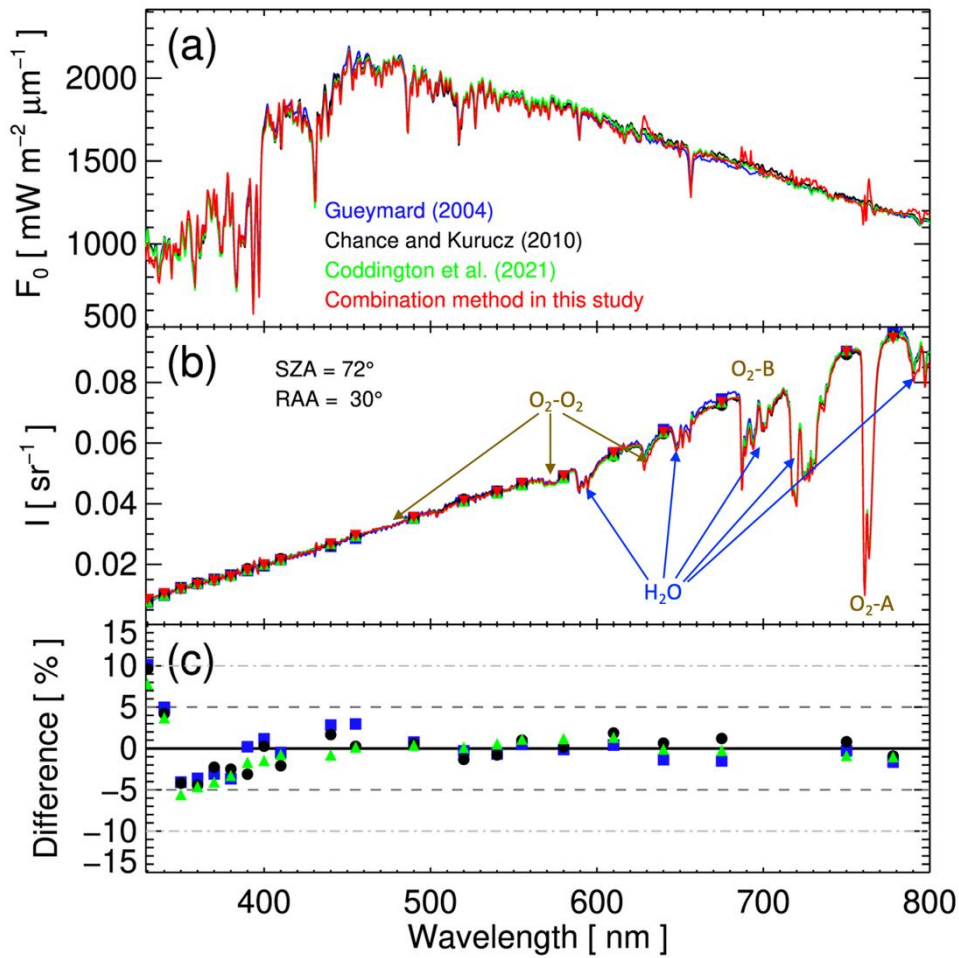
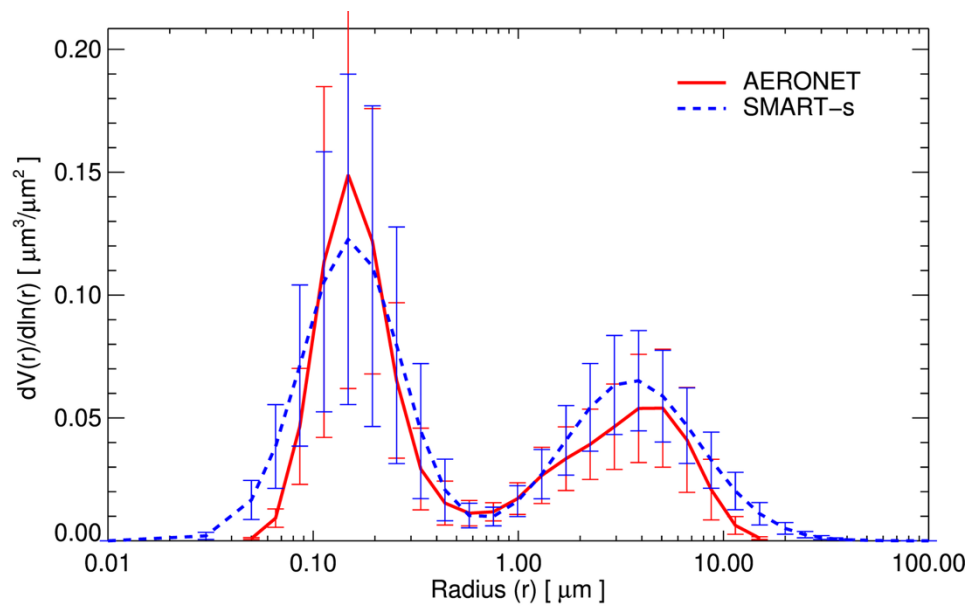


**Figure 4: (a) Spectral solar irradiances ($F_0$) from Gueymard, 2004 (blue), Chance and Kurucz, 2010 (black), Coddington et al., 2021 (green), and combination method developed in this study ($F_{Comb}$ in red). Panel (b) is an example of converted normalized radiance (radiance divided by solar irradiance, $F_0$) measured at Fang, Thailand on 19 March 2019. Colored lines indicate different sources of $F_0$ for the conversion (same as in panel-a), and circles, triangles, and rectangles depict selected wavelengths of aerosol inversion in this study. Panel (c) presents relative biases of the $F_{Comb}$ compared to those of Gueymard, 2004 (blue rectangle), Chance and Kurucz, 2010 (black circle), and Coddington et al.,2021 (green triangle).**

1290 **Figure 5: Average volume-size-distribution of aerosols retrieved from SMART–s (blue dashed line) and AERONET (red solid line) from 19 March to 2 May 2019 at Fang, Thailand. Standard deviations at each radius node during the period are represented as vertical bars in this figure. The SMART–s spectral range is not sensitive to aerosols with radius greater than about 10 μm, and the SMART–s retrievals over this range (see long tails of the blue dashed-line) are mostly determined by the lognormal-shape assumption.**

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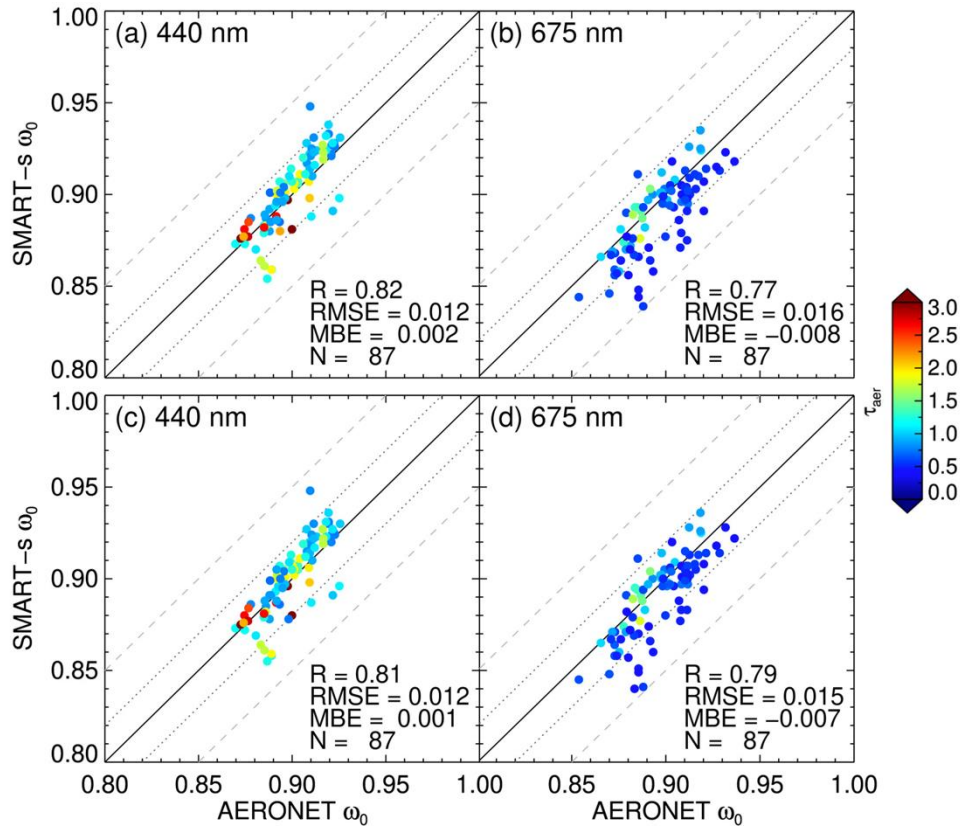


**Figure 6: Comparison of single-scattering albedo ($\omega_0$) from SMART–s and AERONET (Version 3, Level 2.0) at (a) 440 nm and (b) 675 nm from 19 March to 2 May 2019. In these two panels, SMART–s algorithm utilized solar irradiance from Coddington et al. (2021). Panels (c) and (d) are similar plots to (a) and (b) but for SMART–s retrievals using derived solar irradiance in this study. Colored circles represent values of aerosol optical thickness ($\tau_{aer}$) from the colorbar at each wavelength. The R is the correlation coefficient, RMSE denotes root-mean-square error, and MBE is the mean-bias error, and N is the number of samples for the comparison. The dotted and dashed lines represent respectively relative biases of ±0.02 and ± 0.05 from the AERONET product.**

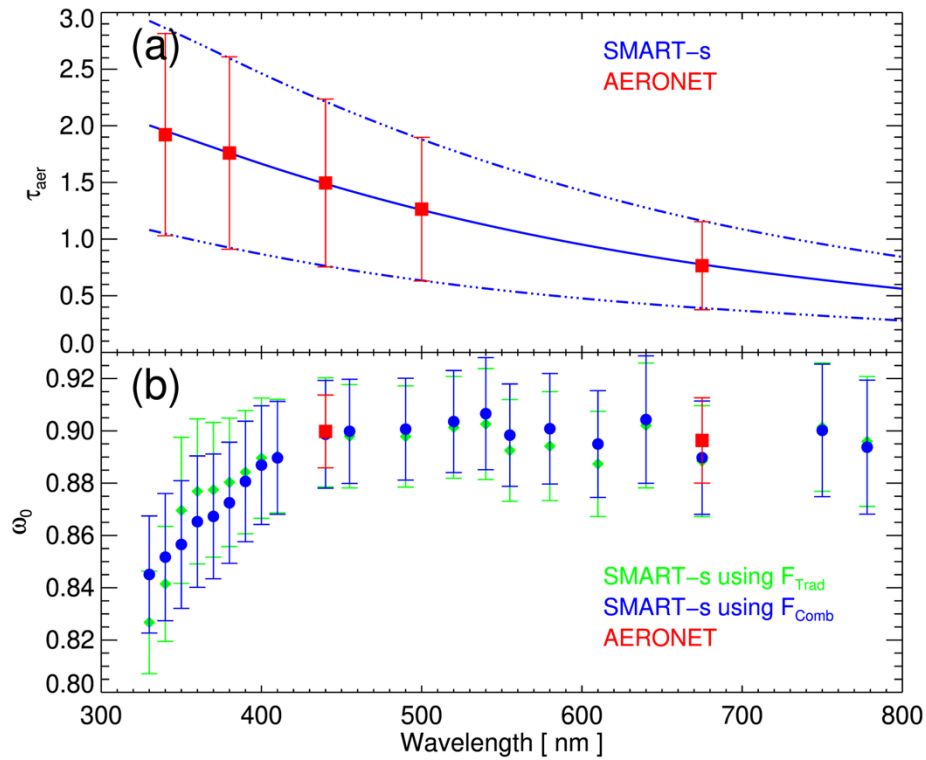



**Figure 7: Mean values of (a) aerosol optical thickness ($\tau_{aer}$) from SMART–s (blue line) and AERONET (red square) within 330–800 nm spectral range measured from 19 March to 2 May 2019 at Fang, Thailand. Panel (b) shows those of spectral single-scattering albedo of aerosols ($\omega_0$) from AERONET (red rectangle) and SMART–s using different solar irradiance; green diamonds used that of Coddington et al. (2021) and blue circles used spectrum derived in this study. Variabilities (one standard deviations) of each value during the deployment period are shown as dashed-dotted lines in panel (a) and vertical bars in panel (a) and (b).**

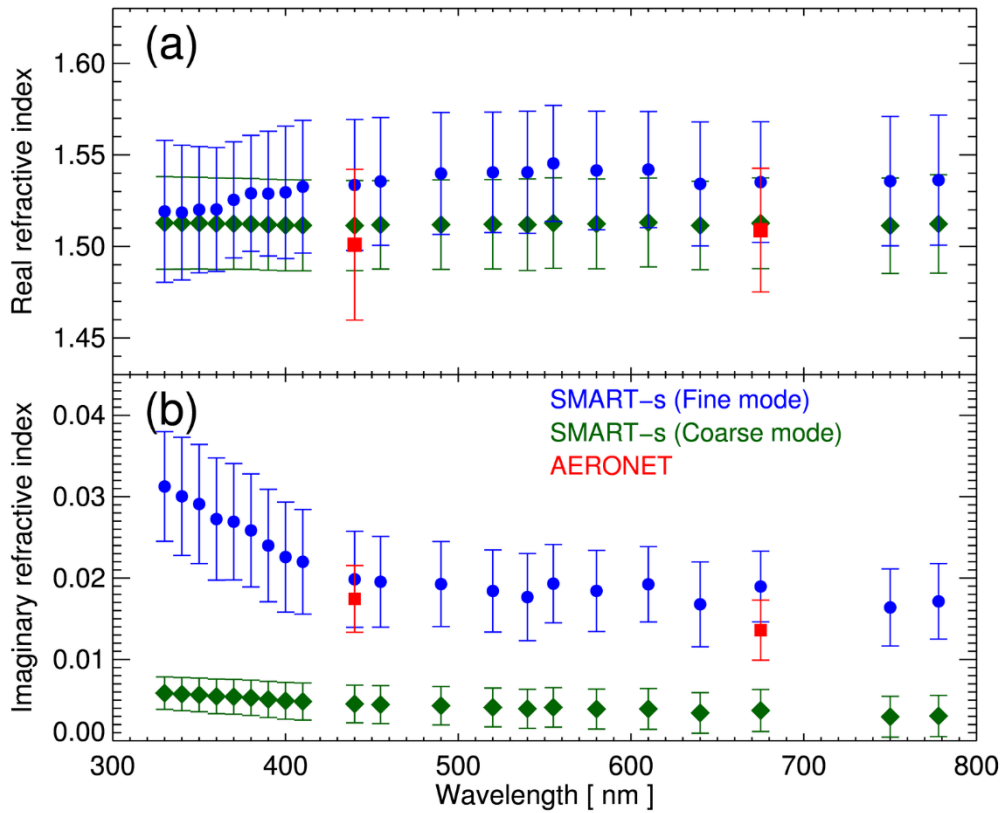



**Figure 8: Mean values of (a) real part of spectral refractive index from AERONET (red square) and SMART–s (blue circle: fine-mode, green diamond: coarse-mode) within SMART–s spectral range measured from 19 March to 2 May 2019 at Fang, Thailand. Panel (b) shows those of Imaginary part. Variabilities (one standard deviations) of each value during the deployment period are demonstrated as vertical bars in panel (a) and (b).**

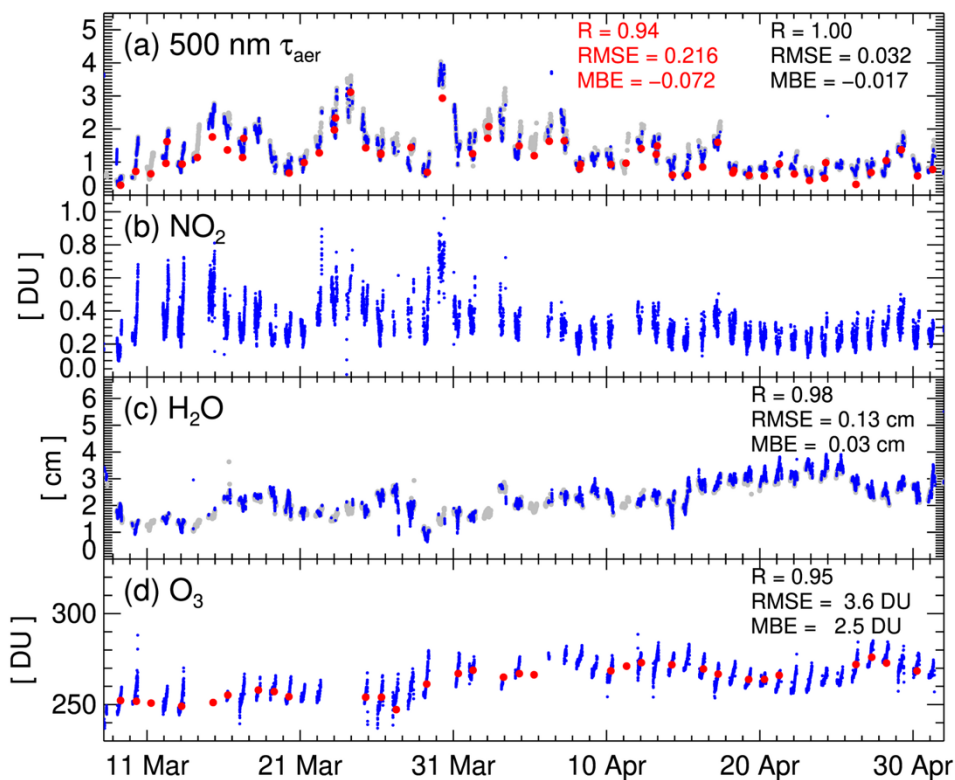
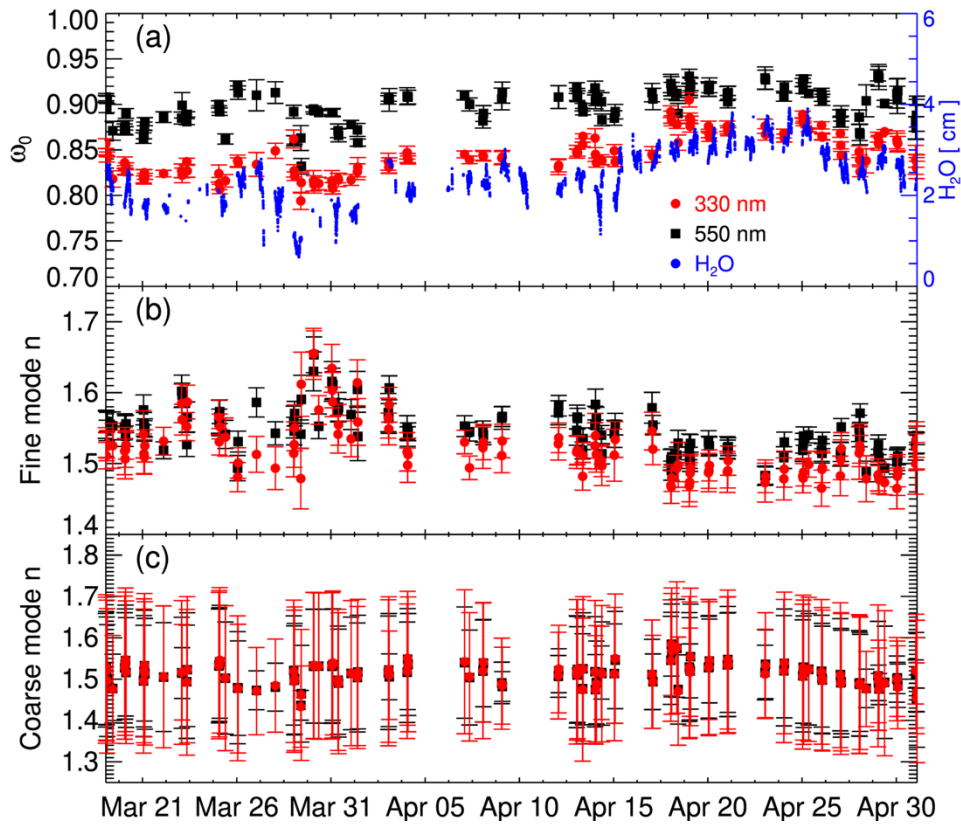



Figure 9: Temporal variations of total columns of (a) aerosol extinction at 500 nm, (b) nitrogen dioxide, (c) precipitable water vapor, and (d) ozone at Fang, Thailand in 2019. The blue circles are from SMART–s retrievals, and the grey circles in panels (a) and (c) show those from the AERONET. In panel (a), the red circles indicate aerosol optical thickness at 550 nm from VIIRS Deep Blue (DB), while those in panel (d) depict total column ozone retrievals from OMI (TOMS Version 8.5). The correlation coefficient (R), root-mean-squared-error (RMSE), and mean-bias-error (MBE) at panels (a), (c), and (d) are between collocated SMART–s and AERONET/OMI data in black, and those of red color in panel (a) are between SMART–s and VIIRS DB retrievals.

1330    **Figure 10: Temporal variations of (a) total aerosol single-scattering albedo ($\omega_0$), real-part of aerosol refractive index ($n$) of (b) fine-mode and (c) coarse-mode retrieved at Fang, Thailand in 2019. The red circles black squares indicate retrievals at 330 nm and 550 nm, respectively. The vertical bars indicate estimated errors of each retrieval based on the optimal-estimation method. The small blue circles in panel (a) represent total column $H_2O$ retrievals from SMART–s.**

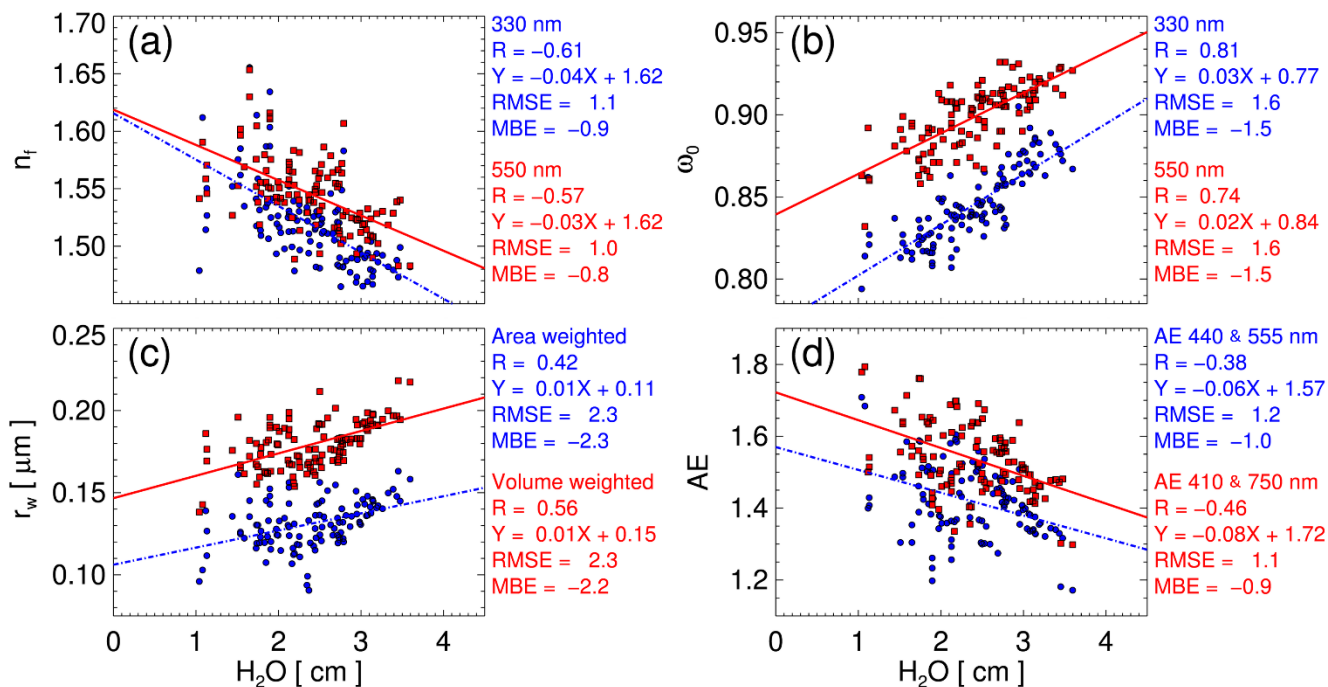



Figure 11: Relationship between aerosol properties and column precipitable $H_2O$ retrieved from SMART−s over Fang, Thailand from 19 March to 2 May 2019. Panels (a) and (b) compare real-part of the refractive index of fine-mode ($n_f$), and aerosol single-scattering albedo ($\omega_0$) to $H_2O$, respectively. For the upper panels, blue and red color represent $n_f$ and $\omega_0$ at 330 nm and 550 nm, respectively. Panels (c) and (d) compare weighted-mean-radius ($r_w$) and Ångström exponent (AE) to the column precipitable $H_2O$, respectively. The blue and red symbols in panel (c) represent area- and volume-weighted mean radius, and those in panel (d) indicate different wavelength pairs for the AE calculations (blue using 440 nm and 555 nm, and red using 410 nm and 750 nm).

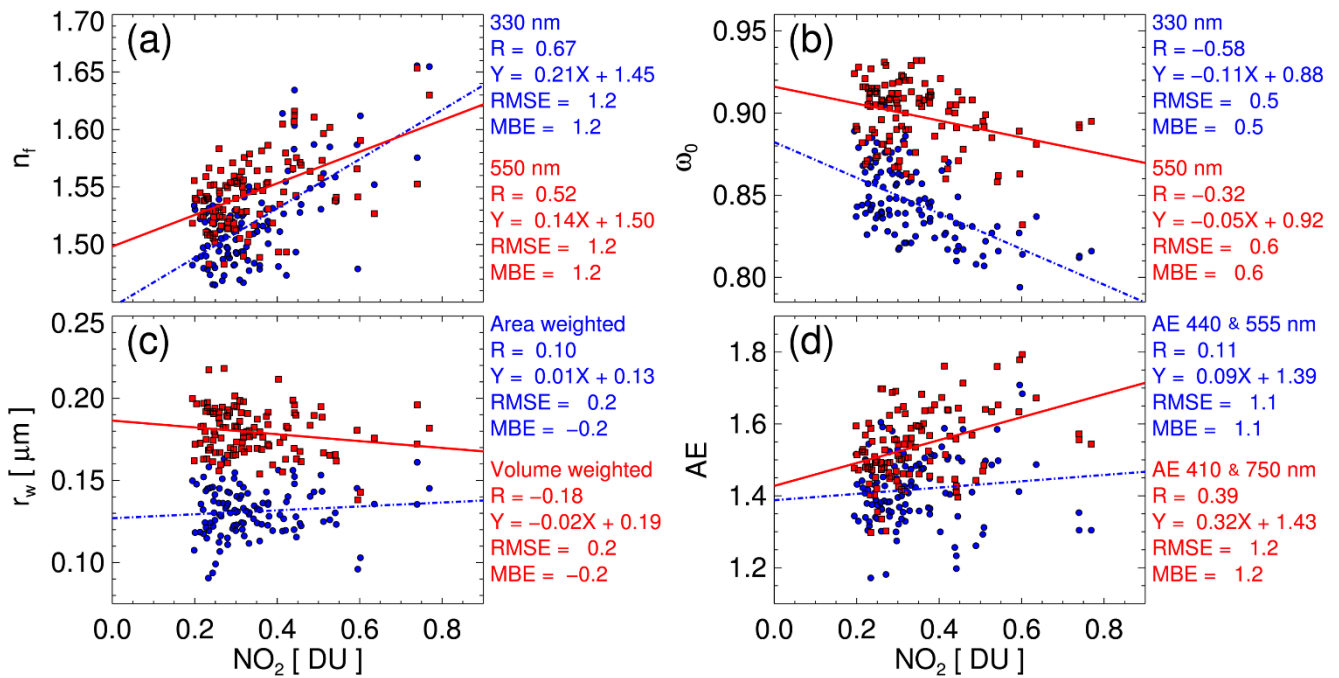



**Figure 12: Comparison of total column amount of NO₂ to (a) real-part of the refractive index of fine-mode ($n_f$) and (b) aerosol single-scattering albedo ($\omega_0$). The blue and red colors symbolize those retrieved at 330 nm and 550 nm, respectively. Panels (c) and (d) compare weighted-mean-radius ($r_w$) and Ångström exponent (AE) to the total column NO₂, respectively. The blue and red symbols in panel (c) represent area- and volume-weighted mean radius, and those in panel (d) indicate different wavelength pairs for the AE calculations (blue using 440 nm and 555 nm, and red using 410 nm and 750 nm).**

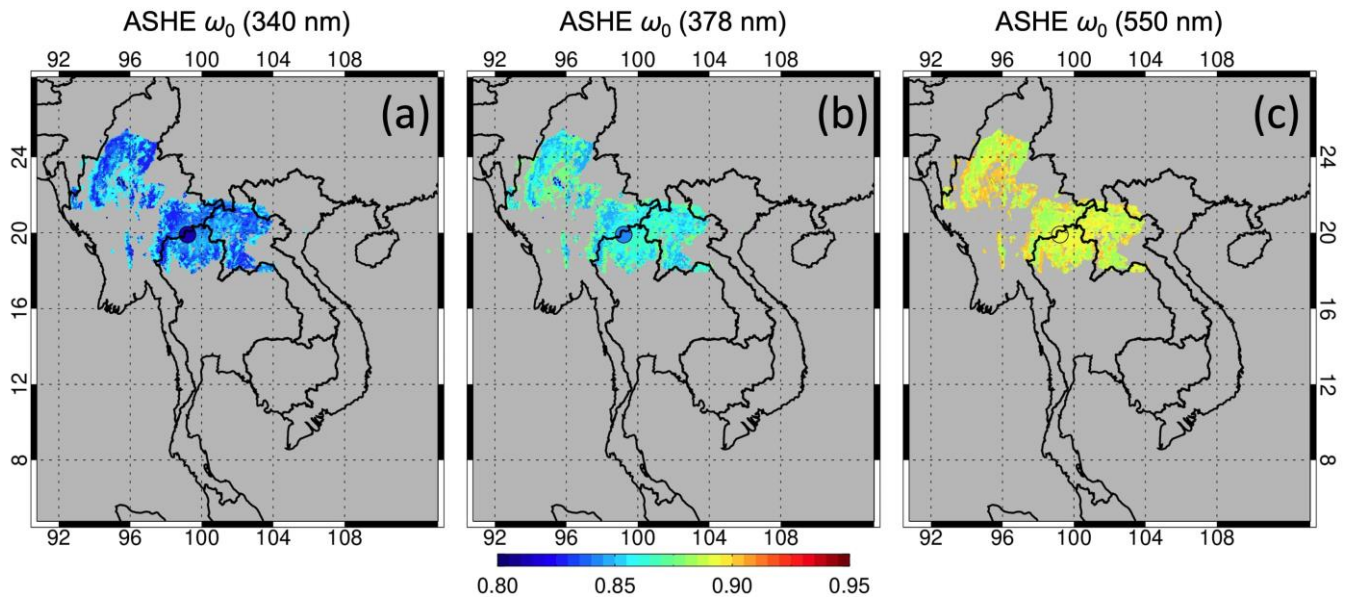



**Figure 13:** Aerosol single-scattering albedo ($\omega_0$) retrieved at (a) 340 nm, (b) 378 nm and (c) 550 nm from ASHE (Aerosol Single-scattering albedo and Height Estimation; Lee et al., 2021) algorithm on 30 March 2019. Colored circles present values of collocated SMART–s retrievals indicated by the colorbar at each wavelength.

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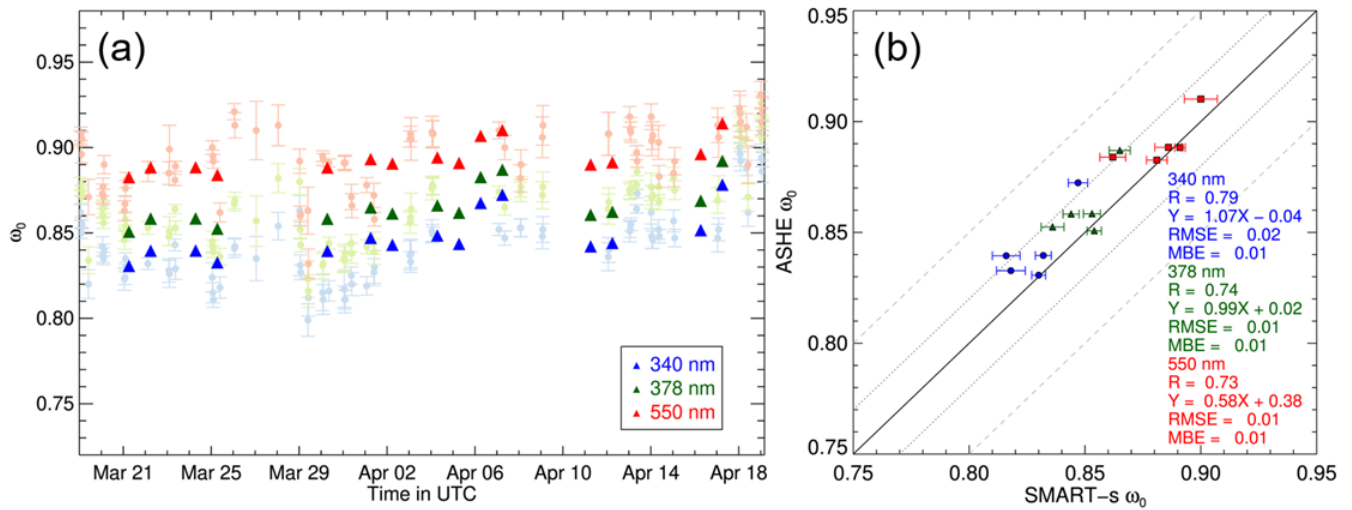
**Figure 14: (a) Temporal variations of aerosol spectral single-scattering albedo ($\omega_0$; 340 nm in blue, 378 nm in green, and 550 nm in red) retrieved from ASHE (dark colors; Aerosol Single-scattering albedo and Height Estimation; Lee et al., 2021) and SMART–s (faint colors) over Fang, Thailand in 2019. The ASHE retrievals are not available after 18 April 2019, since UVAI over the site decreased lower than the ASHE criteria likely due to the increased $\omega_0$ over the period. Panel (b) compares retrieved $\omega_0$ from ASHE and SMART–s at each wavelength during the period. The dotted and dashed lines represent respectively relative biases of ±0.02 and ± 0.05 from the SMART–s product. The collocated samples are limited to have a time difference less than 3 hours. The vertical and horizontal bars in panels (a) and (b), respectively, indicate estimated errors of each retrieval based on the optimal-estimation method.**

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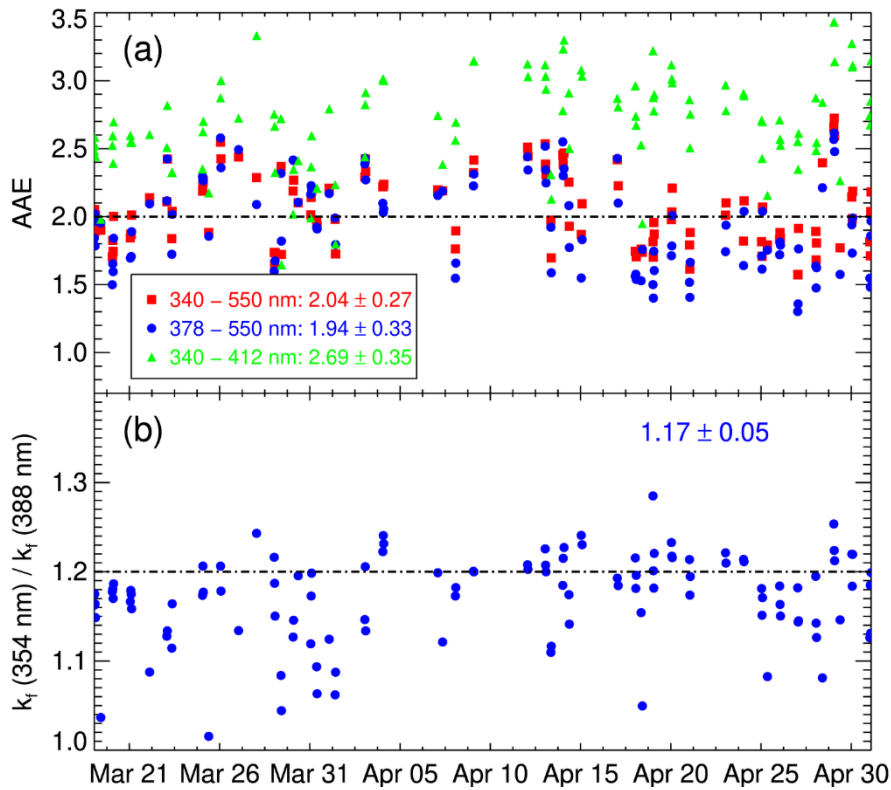
**Figure 15: Temporal variations of (a) absorbing Ångström exponent (AAE) of aerosols, and (b) ratio of fine-mode imaginary refractive indices ($k_f$) of aerosols at 354 nm and 388 nm retrieved at Fang, Thailand in 2019. The color-coded symbols in panel (a) represent different wavelength pairs: red for 340 – 550 nm, blue for 378 – 550 nm, and green for 340 – 412 nm. The dash-dot lines in panel (a) and (b) are assumed values of ASHE and OMAERUV algorithms for smoke aerosols, respectively.**

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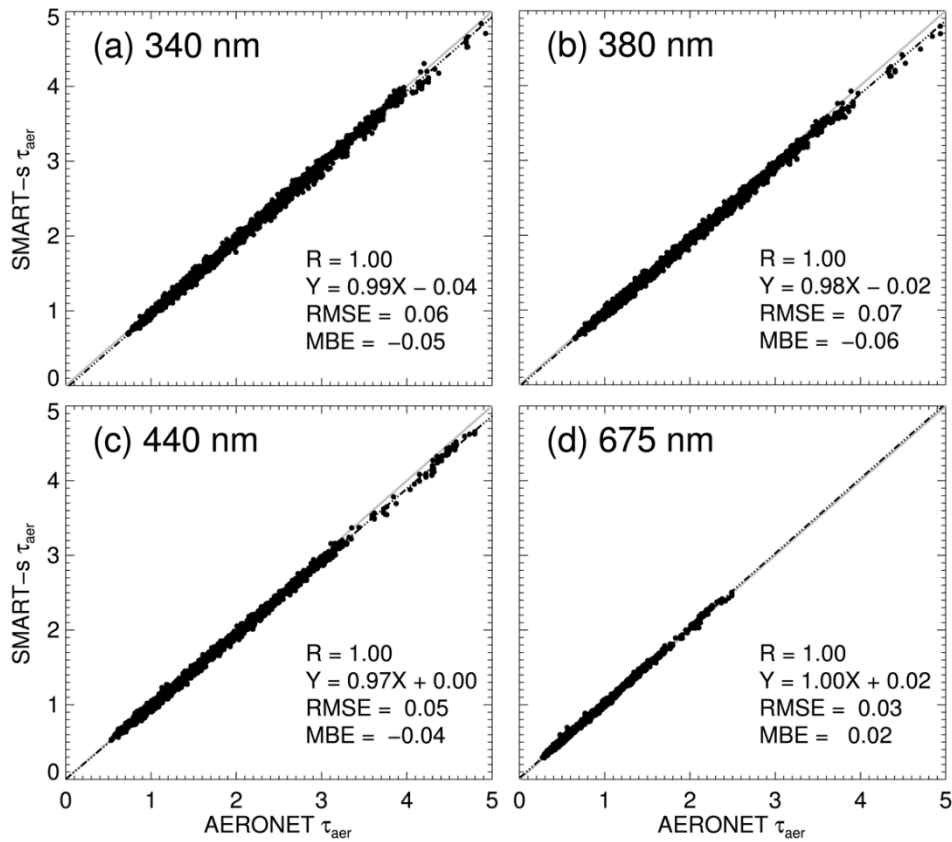
**Figure A1:** Comparison of aerosol optical thickness ($\tau_{aer}$) from SMART–s and AERONET (Version 3, Level 2.0) at (a) 340 nm, (b) 380 nm, (c) 440 nm, and (d) 675 nm measured at Fang, Thailand from 8 March to 2 May in 2019. The R is the correlation coefficient, RMSE denotes root-mean-square error, and MBE is the mean-bias error. The black dot-dashed line and the grey solid line represent regression and one-to-one lines, respectively.
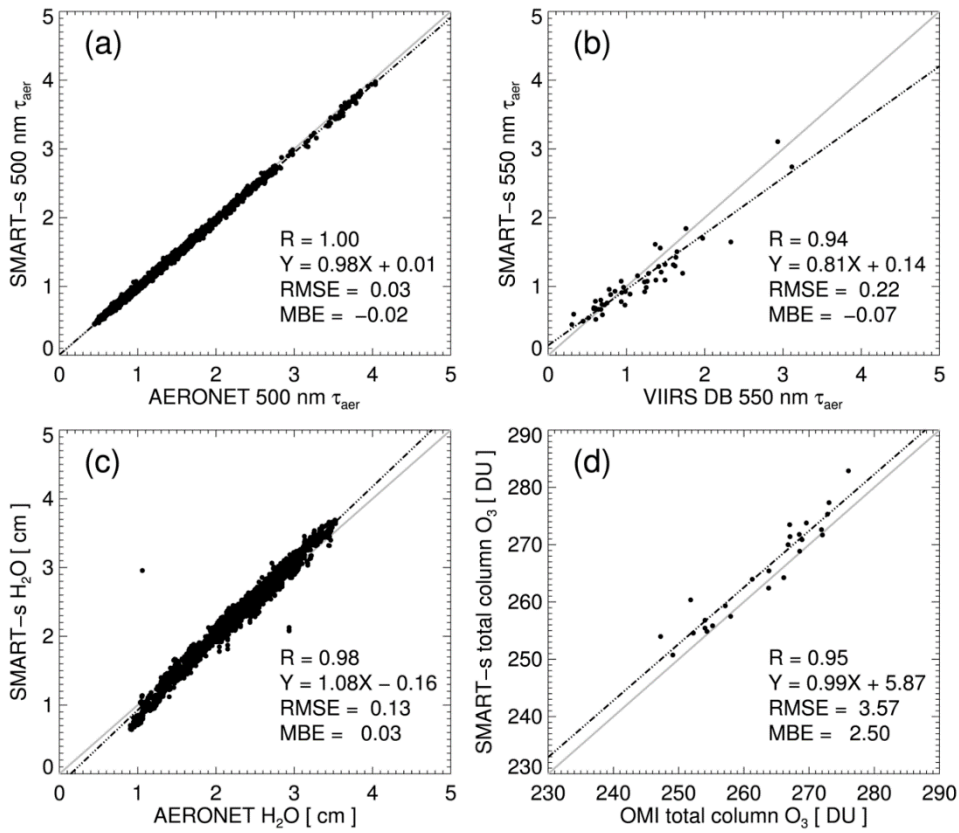
**Figure A2: Panel (a) compares aerosol optical thickness ($\tau_{aer}$) at 500 nm from SMART–s and AERONET (Version 3, Level 2.0), and (b) compares that at 550 nm from VIIRS DB and SMART–s. Total precipitable water vapor (H₂O) products from SMART–s and AERONET are compared in panel (c), and total column ozone retrievals from OMI and SMART–s are compared in panel (d). The retrievals are obtained at Fang, Thailand from 8 March to 2 May in 2019. The R is the correlation coefficient, RMSE denotes root-mean-square error, and MBE is the mean-bias error. The black dot-dashed line and the grey solid line represent regression and one-to-one line, respectively.**
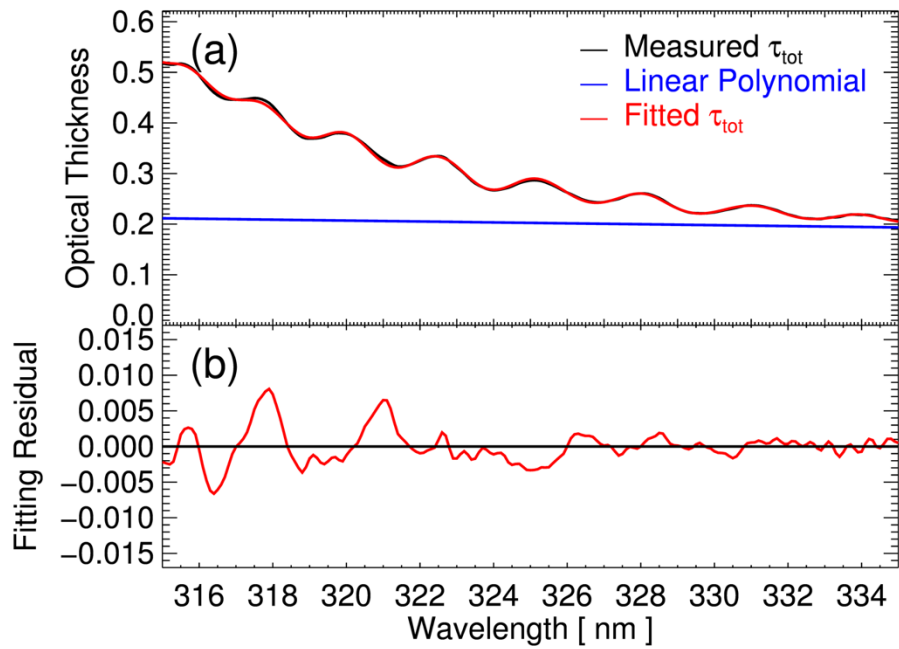
**Figure A3: (a) Spectral fitting result of the SMART–s O₃ algorithm at 16:50 UTC on 14 October 2019 measured on the rooftop of the NASA Goddard Space Flight Center Building 33 (38.99°N latitude, 76.84°W longitude). Panel (b) presents fitting residual (measured optical thickness subtracted by model optical thickness).**

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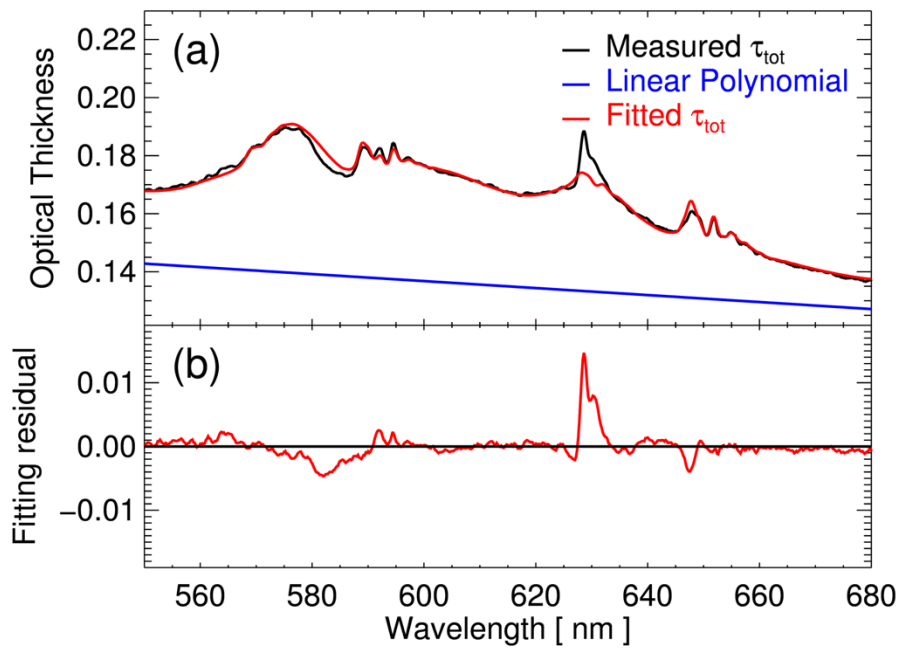
**Figure A4: Similar figure of Figure A3 but for H₂O; (a) Spectral fitting result of the SMART–s H₂O algorithm at 16:55 UTC on 14 October 2019 measured at the same location. Panel (b) presents fitting residual (measured optical thickness subtracted by model optical thickness).**

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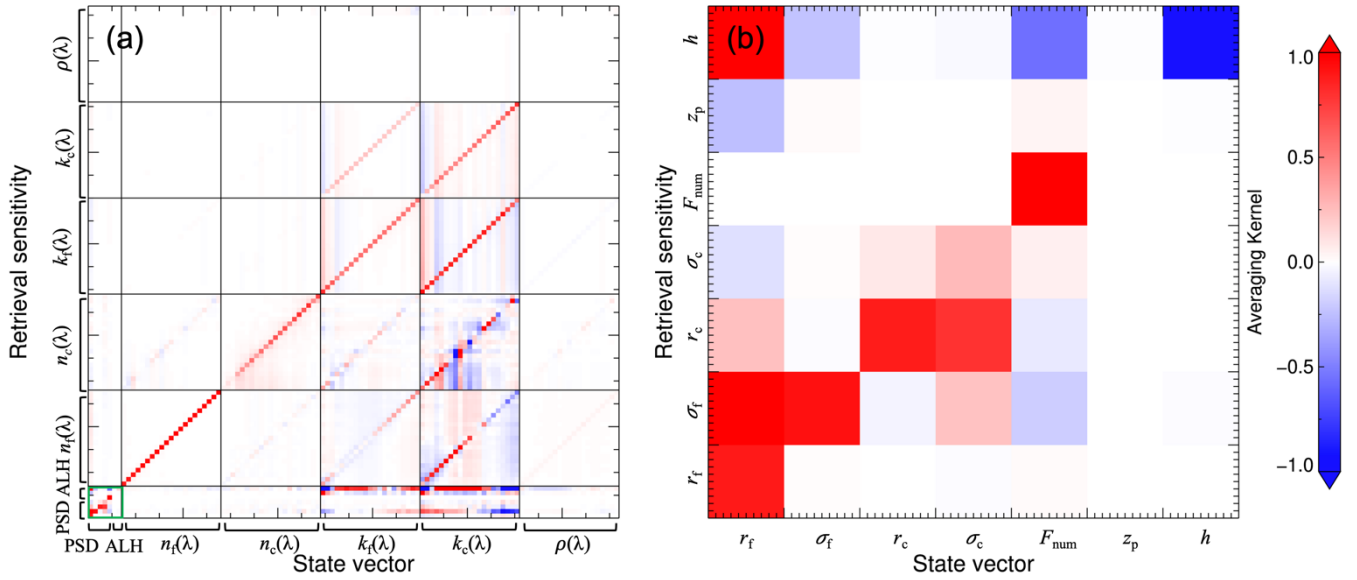
**Figure A5:** Averaging kernel matrix (A) of SMART–s retrievals at 08:31 UTC on 10 April 2019 at Fang, Thailand, when fine-/coarse-mode volume fractions were comparable, and aerosol single-scattering albedo was about 0.87 with aerosol optical thickness about 1.06 at 440 nm. Indices from 1 to 5 correspond to fine-mode mean radius ($r_f$) and geometric standard deviation ($\sigma_f$), those pair for coarse mode ($r_c$ and $\sigma_c$), and number fine-mode fraction ($F_{num}$). Indices 6 and 7 are peak height ($z_p$) and dispersion parameter ($h$) of the assumed aerosol extinction profile. Each element of real ($n$) and imaginary ($k$) part of the refractive index for fine-(subscript f) and coarse-mode (subscript c) indicates its retrieval sensitivity at each wavelength. Panel (a) shows the whole A, and panel (b) zooms to particle-size-distribution (PSD) and aerosol-layer-height (ALH) parameters (indices from 1 to 7) as indicated as a green square in panel (a).
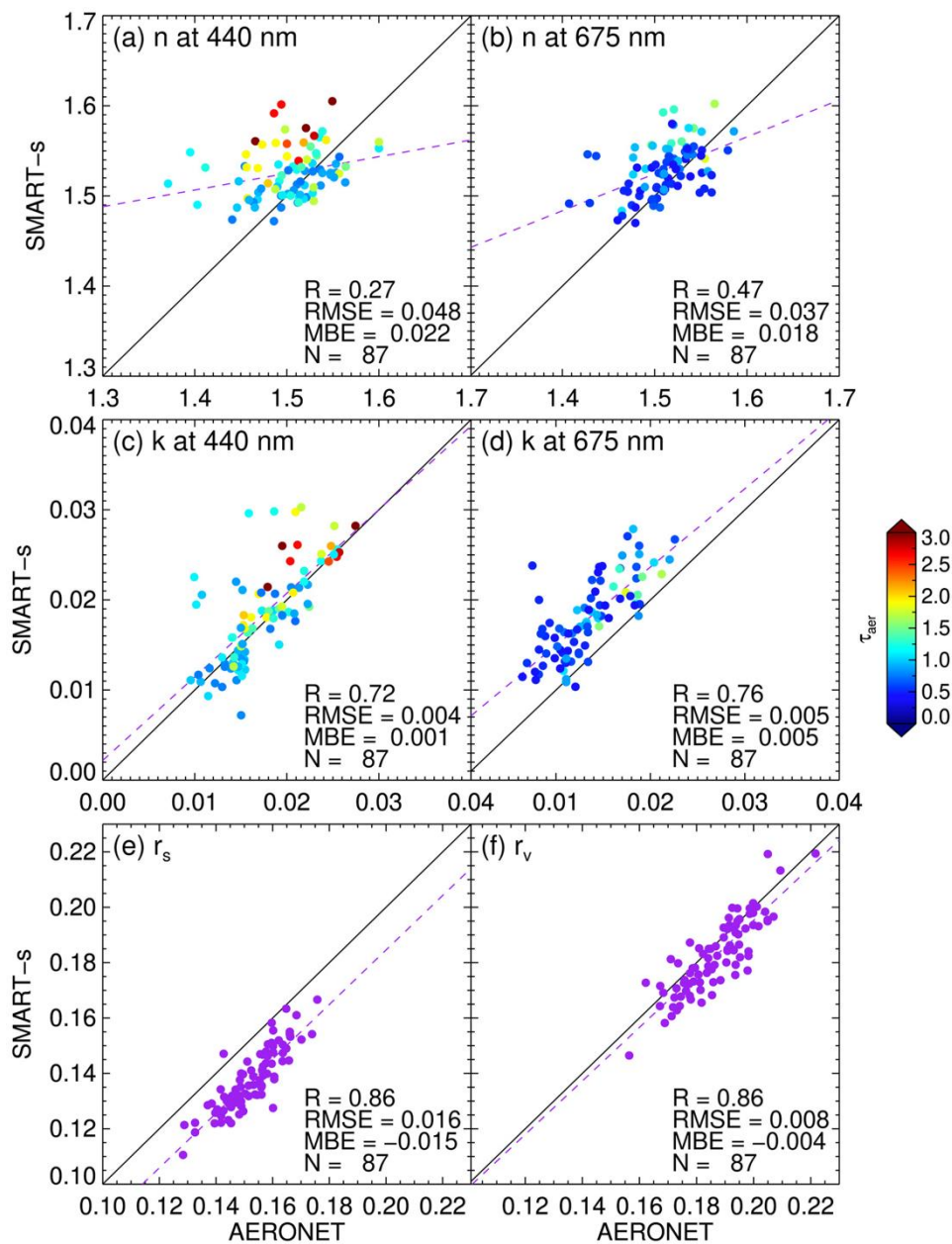
**Figure A6: Comparison of aerosol inversion products from SMART–s and AERONET (Version 3, Level 2.0) from 19 March to 2 May 2019. Upper panels compare the real part of the refractive index at (a) 440 nm and (b) 675 nm, and middle panels compare the imaginary part at these wavelengths. Colored circles in panel (a)-(d) represent values of aerosol optical thickness ($\tau_{aer}$) from the colorbar at each wavelength. Lower panels compare (e) area- and (f) volume-weighted-mean-radii ($r_s$ and $r_v$, respectively). The R is the correlation coefficient, RMSE denotes root-mean-square error, and MBE is the mean-bias error, and N is the number of samples for the comparison. Black solid-line and purple dashed-line depict one-to-one and regressions, respectively.**