### Simultaneous retrievals of biomass burning aerosols and trace gases from the ultraviolet to near-infrared over northern Thailand during the 2019 pre-monsoon season

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**Abstract.** With the advent of spaceborne spectroradiometers in a geostationary constellation, measuring high spectral resolution ultraviolet-visible (UV-VIS) and selected near-/shortwave-infrared (NIR/SWIR) radiances can enable the probing of the life cycle of key atmospheric trace gases and aerosols at higher temporal resolutions over the globe. The UV-VIS measurements are important for retrieving several key trace gases (e.g., O<sub>3</sub>, SO<sub>2</sub>, NO<sub>2</sub>, and HCHO) and particularly for deriving aerosol characteristics (e.g., aerosol absorption and vertical profile). This study examines the merit of simultaneous retrievals of trace gases and aerosols using a ground-based spectroradiometer covering the UV-NIR to monitor their physicochemical processes and to obtain reliable aerosol information for various applications. During the 2019 pre-monsoon season over northern Thailand, we deployed a ground-based SMART-s (Spectral Measurements for Atmospheric Radiative Transferspectroradiometer) instrument, which is an extended-range Pandora with reliable radiometric calibration in the 330-820 nm range, to retrieve remotely sensed chemical and aerosol properties for the first time near biomass burning sources. The high spectral resolution ( $\sim 1.0$  nm full width half maximum with  $\sim 3.7 \times$  oversampling) of sun and sky measurements from SMART-s provides several key trace gases (e.g.,  $O_3$ ,  $NO_2$ , and  $H_2O$ ) and aerosol properties covering the UV where significant light absorption occurs by the carbonaceous particles. During the measurement period, highly correlated total column amounts of NO<sub>2</sub> and aerosol optical thickness ( $\tau_{aer}$ ) retrieved from SMART-s (correlation coefficient, R = 0.74) indicated their common emissions from biomass burning events. The SMART-s retrievals of the spectral single scattering albedo ( $\omega_0$ ) of smoke aerosols showed an abrupt decrease in the UV, which is an important parameter dictating photochemical processes in the atmosphere. The values of  $\omega_0$  and column precipitable water vapor (H<sub>2</sub>O) gradually increase with the mixing of biomass burning smoke particles and higher water vapor concentrations when approaching the monsoon season. The retrieved  $\omega_0$  and weighted mean radius of fine-mode aerosols from SMART-s showed positive correlations with the H<sub>2</sub>O (R = 0.81 for  $\omega_0$  at 330 nm and 0.56 for the volume-weighted mean radius), whereas the real part of the refractive index of fine-mode aerosol  $(n_f)$  showed negative correlations (R = -0.61 at 330 nm), which suggest that aerosol aging processes including hygroscopic growth (e.g., humidification and cloud processing)

can be a major factor affecting the temporal trends of aerosol optical properties. Retrieved  $n_f$  and  $\omega_0$  were closer to those of the water droplet (i.e.,  $n_f$  of about 1.33 and  $\omega_0$  of about 1.0) under lower amounts of NO<sub>2</sub> during the measurement period; considering that the NO<sub>2</sub> amounts in the smoke may indicate the aging of the plume after emission due to its short lifetime, the tendency is also consistent with active hygroscopic processes of the aerosols over this area. Retrieved UV aerosol properties from SMART–s generally support the assumed smoke aerosol models (i.e., the spectral shape of aerosol absorption) used in NASA's current satellite algorithms, and their spectral  $\omega_0$  retrievals from ground and satellites showed good agreements (R = 0.73-0.79). However, temporal and spectral variabilities in the aerosol absorption properties in the UV emphasize the importance of a realistic optical model of aerosols for further improvements in satellite retrievals.

### 1 Introduction

Significant spatiotemporal variabilities in the aerosols in the atmosphere complicate understanding of their scattering and absorption of the solar irradiance, which results in one of

- <sup>5</sup> the largest uncertainties in predicting future climate (IPCC, 2021 [151]; Gliß et al., 2021; Myhre et al., 2013, and references therein). The dominant fraction of the aerosols over the globe cools the atmosphere by reflecting solar irradiance, whereas some species (e.g., black carbon in the smoke plumes) heat
  <sup>10</sup> the air by absorbing sunlight (i.e., direct radiative effects,
- DRE; Chylek and Coakley, 1974; Haywood and Boucher, 2000; Yu et al., 2006). Primary factors of aerosols affecting the DRE are their loading and their absorption properties (e.g., Takemura et al., 2002, and references therein), which
- <sup>15</sup> are often defined as aerosol optical thickness ( $\tau_{aer}$ ; total extinction by aerosols) and single scattering albedo ( $\omega_0$ ; a ratio of the scattering to total extinction by aerosols), respectively. The  $\omega_0$  is calculated from complex refractive indices (n + ik; where n and k are a real and imaginary part, which
- <sup>20</sup> depends on chemical composition) and particle size distribution (PSD), by assuming a spherical (Mie, 1908) or more sophisticated shape (e.g., Mishchenko et al., 2003; Yang et al., 2007). The  $\omega_0$  of non-absorbing aerosols (e.g., sea salt, sulfate, and nitrate particles) is close to 1.0, with a relatively
- <sup>25</sup> flat spectral shape, whereas it decreases down to about 0.7 for absorbing aerosols (e.g., smoke and dust particles) with significant spectral gradients (e.g., Dubovik et al., 2002; Eck et al., 2013; Müller et al., 2011; Sayer et al., 2014). Meteorological condition and aerosol hygroscopicity are also key
- <sup>30</sup> parameters affecting  $\omega_0$ , since increased water content in the particles changes the *n*, *k*, and PSD, which enhances light scattering and results in higher  $\omega_0$  than dry particles (e.g., Jefferson et al., 2017; Li et al., 2019; Tao et al., 2014).

Decades of efforts have led to remote sensing techniques from both ground and satellite providing reliable  $\tau_{aer}$  retrievals over major parts of the globe (e.g., Giles et al., 2019; Hsu et al., 2019; Levy et al., 2013), whereas other aerosol properties retrieved from satellites are limited and relatively more uncertain due to the lower measurement sensitivity and 40 surface contributions (e.g., Jeong et al., 2016; Moosmüller

et al., 2009). However, recent studies using more measure-

ment parameters (e.g., multi-angle polarimetric measurements) showed promising results to provide reliable aerosol properties and constituents from satellites (e.g., Dubovik et al., 2019; Li et al., 2019). Globally networked ground-based instruments have provided reliable optical and physical properties of aerosols (e.g.,  $n, k, \omega_0$ , and PSD), which are less affected by surface reflectance and acquire sufficient information content from multiple observation geometries (e.g., Dubovik and King, 2000; Jeong et al., 2020; Nakajima et al., 50 2020; Sinyuk et al., 2020). Satellite-based retrievals have utilized the aerosol properties from ground-based instruments as key constraints to expand upon their limitations (e.g., Hsu et al., 2019; Levy et al., 2013; Sayer et al., 2014).

Wildfires and prescribed fires have burned about 3.5 % of 55 Earth's ice-free land surface each year from 2001 to 2010 (Randerson et al., 2012) and emit a significant fraction of global aerosols and their precursors into the atmosphere. The biomass burning aerosols (or smoke) consist primarily of carbonaceous aerosols (black and organic carbon), inorganic 60 particles (e.g., potassium, chloride, sulfate, inorganic salts, and trace minerals), and inorganic and organic vapors (Hodshire et al., 2019, and references therein). Particularly primary and secondary organic aerosols, which account for a substantial fraction of fine-mode smoke aerosols, comprise 65 various compounds with enormously different volatility, oxidation, and hygroscopic properties (Xu et al., 2017, and references therein). Due to the reactivity and diversity of smoke particles, the  $\omega_0$  evolves with its environment (i.e., location and season), age, mixing state, and emission source 70 of the plume (e.g., Eck et al., 2013; Haywood et al., 2003; Konovalov et al., 2017). In addition, Petters et al. (2009) reported that a major fraction of the smoke aerosols is already cloud-condensation-nuclei (CCN) active and does not require chemical conversion to be more hygroscopic particles 75 for cloud formation and wet deposition, which adds another complication in understanding the Earth's climate.

Numerous studies have utilized ground-, airborne-, and satellite-based remote sensing techniques to monitor the properties and aging processes of the smoke particles. For <sup>80</sup> example, Haywood et al. (2003) compared aerosol properties (e.g., PSD,  $\tau_{aer}$ , and  $\omega_0$ ) from the collocated AERONET (AErosol RObotic NETwork; Holben et al., 1998) and

airborne in situ measurements at Windhoek, Namibia, in September 2000, which showed excellent agreements. Eck et al. (2013) analyzed the seasonal trend of aerosol properties retrieved from the AERONET and OMI (Ozone Moni-

- <sup>5</sup> toring Instrument) over southern Africa for a 15-year period and reported that the  $\omega_0$  increases significantly as the burning season progresses. Pistone et al. (2019) compared the spectral  $\omega_0$  of smoke aerosols from six independent airborne- and ground-based remote sensing/in situ instruments in Septem-
- <sup>10</sup> ber 2016 at Walvis Bay, Namibia, which showed acceptable agreements within the known uncertainties of each instrument (relative differences of less than about 0.03 in the mid-visible (VIS) range and less than about 0.05 in the near-infrared, depending on the instruments). Over South-
- <sup>15</sup> east Asia, a series of field campaigns, including BASE-ASIA (Biomass-burning Aerosols in South-East Asia: Smoke Impact Assessment) in 2006 and 7-SEAS (Seven SouthEast Asian Studies) from 2008 to the present, aimed to characterize aerosol-meteorological interactions over the region,
- <sup>20</sup> mostly focusing on the smoke plumes. Physicochemical and optical properties of smoke aerosols were analyzed by utilizing intensive ground- and satellite-based instruments during the campaigns (e.g., Lin et al., 2013; Pantina et al., 2016; Reid et al., 2013; Tsay et al., 2013, 2016).
- <sup>25</sup> One of the important characteristics of the carbonaceous aerosols is their significant spectral variabilities in the optical properties in the ultraviolet (UV) wavelengths, which are associated with photolysis processes in the atmosphere, thereby affecting tropospheric photochemistry, hu-
- <sup>30</sup> man health, and agricultural productivity (e.g., George et al., 2015, and references therein). A majority of previous studies utilized direct/diffuse irradiance instruments (e.g., UV MultiFilter Rotating Shadowband Radiometer, UV-MFRSR, and Brewer spectroradiometer) to retrieve  $\omega_0$  in discrete channels
- <sup>35</sup> in the UV. However, as these instruments measure only two observation parameters per channel, the algorithms adopted different sources of measurements/assumptions to complement the insufficient information (e.g., see Table 1 in Corr et al., 2009). For instance, in the absence of additional col-
- <sup>40</sup> located instruments, they assumed fixed asymmetry parameter and surface albedo from previous studies or climatology (e.g., Bais et al., 2005; Petters et al., 2003; Wetzel et al., 2003). Collocated AERONET instruments have provided more realistic constraints of aerosol properties to the UV-
- <sup>45</sup> MFRSR measurements (e.g., PSD and *n* from visible wavelengths; Corr et al., 2009; Krotkov et al., 2005a) for retrieving  $\omega_0$  in the UV. Trace gas absorption (e.g., O<sub>3</sub> and NO<sub>2</sub>) is another source of error for the  $\omega_0$  retrieval using these instruments. To take into account the gas absorptions, Goer-
- <sup>50</sup> ing et al. (2005) simultaneously retrieved total column O<sub>3</sub>, in addition to the  $\tau_{aer}$  and  $\omega_0$ , by using the spectral feature of irradiance. Later, Taylor et al. (2008) added a wavelengthindependent asymmetry parameter to the state vector, where both algorithms are based on the optimal estimation method
- 55 (OEM; Rodgers, 2000). Krotkov et al. (2005a) used the

aerosol-phase function calculated from n at 440 nm, PSD from AERONET, and total column O3 from the Brewer spectroradiometer to retrieve  $\omega_0$  in the UV channels. To account for the NO<sub>2</sub> absorption, which is a significant error source of  $\omega_0$  retrieval for low aerosol loading, they added retrieved 60 NO<sub>2</sub> from the Brewer spectroradiometer for their algorithm (Cede et al., 2006; Krotkov et al., 2005b). SKYNET (SKY radiometer NETwork) instruments are a similar type of sunsky spectroradiometer to the AERONET, which provides  $\omega_0$ at discrete channels in the UV (i.e., 340 and 380 nm). The 65 SKYNET algorithm accounts for the O<sub>3</sub> absorption by using its retrieved total column from its 315 nm channel (Nakajima et al., 2007, 2020). Accuracies of the  $\omega_0$  retrievals from the SKYNET depend on errors in measurement and calibrations for sun and sky scans, surface albedo, cloud con-70 tamination, and the version of the processing software (i.e., SKYRAD.PACK), which showed relatively high biases compared to AERONET (up to 0.07 at longer wavelengths). Recently, Mok et al. (2018) combined AERONET (for n, PSD, and  $\tau_{aer}$ ) and Pandora (for total column O<sub>3</sub> and NO<sub>2</sub>) <sup>75</sup> products to the UV-MFRSR measurements to retrieve spectral  $\omega_0$  in the UV, which showed excellent agreements with SKYNET in the UV (i.e., 340 and 380 nm) but lower correlations in the longer wavelengths (i.e., 673 and 870 nm).

In addition, spectral *n* and *k* provide information not only 80 on optical properties but also the chemical composition and physical status. k demonstrates the attenuation of light by particles, which is the key parameter for determining  $\omega_0$ , whereas *n* describes the phase of light scattering by the particles. Numerous studies have focused on measuring/retriev-85 ing n and k by utilizing various techniques to understand the effects of atmospheric particles on climate forcing and tropospheric photochemistry. Table 1 summarizes reported values of n from previous studies and the current study. Kim et al. (2010) retrieved n (at 670 nm) of secondary organic  $_{90}$ aerosols (SOAs) generated by oxidizing  $\alpha$ -pinene,  $\beta$ -pinene, and toluene with  $O_3$ ,  $NO_x$ , and sunlight. The retrieved *n* varied between 1.38 and 1.61, and they suggested that the n of SOAs depends on the aerosol mass concentration, oxidation chemistry, temperature, and aerosol aging. Liu et al. (2013) 95 measured the n and k of SOAs for 220 to 1200 nm, using a variable angle spectroscopic ellipsometer, and reported a rapid increase in n and k in the UV. The n of the three selected SOAs ranged from 1.53-1.58 at 310 nm, 1.49-1.52 at 550 nm, and 1.48–1.50 at 1000 nm. Shepherd et al. (2018) 100 estimated the spectral *n* of urban, remote, and wood smoke aerosols from 460 to 760 nm, based on the optical trapping method, and reported high values of n in the wood smoke aerosols ( $\sim 1.58$ ) compared to the other types (1.47–1.52). They also summarized and compared their values of the spec- 105 tral n to other studies well in their paper. Sumlin et al. (2018) retrieved the spectral n and k (at 375, 405, 532, and 1047 nm) of brown carbon aerosols emitted from controlled fire using burning sources at various geographic origins. They reported that n varies between 1.5 and 1.7 without meaning- 110

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ful dependencies on wavelength, moisture content, source depth, or geographic origin, whereas k increases from 0.003 to 0.014 as wavelengths vary from 532 to 375 nm. Di Biagio et al. (2019) estimated the n and k (at discrete chan-

- <sup>5</sup> nels in 370–950 nm) of 19 mineral dust aerosols from different sources based on Mie calculations combining optical and size measurements. They reported higher k (lower  $\omega_0$ ) of dust particles in the shorter wavelengths, which also depends on the iron content of dust, but the source and wavelength-
- <sup>10</sup> independent values of *n* ranged from 1.48 to 1.55. More recently, Womack et al. (2021) retrieved the *n* and *k* of biomass burning aerosols from 13 controlled fires over a 360-720 nm spectral range using a broadband cavity-enhanced spectrometer combined with PSD measurements. Their algorithm in-
- <sup>15</sup> corporates Mie and Rayleigh–Debye–Gans scattering theories to account for both spherical and non-spherical particles and retrieved *n* at about 1.55–1.60 and *k* to be significantly high ( $\sim 0.25$ ) in the UV.

To be closely in line with and to continue such efforts, we <sup>20</sup> deployed a set of instruments including SMART-s (Spec-

- tral Measurements for Atmospheric Radiative Transferspectroradiometer; Jeong et al., 2018, 2020) and AERONET during the pre-monsoon, yet active biomass burning, season at Fang, Thailand, in 2019. Specifically, we aim to sug-
- <sup>25</sup> gest the benefit of simultaneous retrievals of aerosols and trace gases covering the UV, which may provide useful information on their physicochemical processes. In addition, aerosol properties in the UV are also important for various satellite algorithms for deriving higher-order aerosol param-
- <sup>30</sup> eters (e.g., absorption and vertical distribution) for which reliable measurements remain sparse. The benefits of employing SMART-s, a major instrument we utilized for this study, include the following:

sufficient spectral resolution and coverage for measuring both aerosols and key trace gases (e.g., O<sub>3</sub>, NO<sub>2</sub>, and H<sub>2</sub>O retrievals from direct-sun measurements), and in turn, the high temporal measurements of gaseous absorption help improve the accuracy of ω<sub>0</sub> retrieval;

- instantaneous measurements of the sun/sky spectrum, permitting aerosol spectral properties retrieved from an identical set of volumes;
- reliable radiometric calibration from about 330 to 820 nm by utilizing a NIST-traceable (National Institute of Standards and Technology) uniform spectral radiance source (accuracy of about 1% in the VIS-NIR (near-infrared) and about 2% in the UV wavelengths at an approximate 95% confidence level) to enable accurate retrievals of aerosol column properties (e.g.,  $\tau_{aer}$ , *n*, *k*, and  $\omega_0$ ); and
- stable performance, which is field deployable for a long period. The recent global expansion of the Pandora network operation is based on its reliability at various field

conditions, and SMART–s is nearly identical to the Pandora instrument, except for the spectrometer (extended range from about 280 to 820 nm, with about 1 nm spectral resolution).

As this study is the first attempt to retrieve aerosol properties from the SMART-s near the source region of active and extensive biomass burning, we summarize the experimental design, instrument characteristics, and the radiometric cal- 60 ibration in Sect. 2. In Sect. 3.1, we compare the retrieved aerosol property retrievals (e.g., n, k, and  $\omega_0$ ) from the SMART-s with those collocated from AERONET for a consistency check. Analyses of temporal variations in aerosols and total column trace gases (i.e., NO<sub>2</sub>, H<sub>2</sub>O, and O<sub>3</sub>) re- 65 trieved from the SMART-s are described in Sect. 3.2. We also demonstrated the relationship between aerosol properties and trace gas abundances in this section. Section 3.3 discusses possible applications of the retrieved aerosol parameters for satellite algorithms and preliminary validation/com-70 parison results. The summary and conclusions are given in Sect. 4.

### 2 Measurements and calibrations

#### 2.1 Experimental setup

The ground-based spectroradiometer observations have of-75 fered optimum inversion products of the atmosphere for validating/comparing those from collocated spaceborne sensors; these are less affected by the surface reflectance and can acquire more informative products from their higher resolution of temporal, spectral (including polarization), and angular 80 measurements. In addition, strategically networked groundbased instruments (e.g., Distributed Regional Aerosol Gridded Observation Networks or DRAGON; Holben et al., 2018) can supplement their limited spatial representation. As part of the ongoing 7-SEAS campaign, intensive ob-85 servations were conducted during the pre-monsoon season in April-May 2019 over northern Thailand, specifically the Chiang Mai, Fang, and Doi Ang Khang areas. The international collaborators deployed a sUAS (small Unmanned Aerial System) in a rotary/fixed-wing configuration 90 for  $\sim 130$  flights to measure the boundary layer profiles of thermodynamics and aerosol size/absorption. A mini-lidar, the surface measurements of trace gases, and multiple chemistry samplers are also collocated with the three AERONET instruments and one SMART-s instrument during the cam- 95 paign. As SMART-s is located in the middle of a large source areas of biomass burning during the season, it can provide useful information on carbonaceous aerosols and key trace gases despite its limited spatial coverage.

Figure 1 shows an example of spatial distributions of  $\tau_{aer}$  <sup>100</sup> at 550 nm (Hsu et al., 2019) from the Deep Blue (DB) aerosol algorithm applied to VIIRS (Visible Infrared Imaging Radiometer Suite) aboard the SNPP (Suomi National Polar-

Reference	Aerosol type	Wavelengths	Real part of the refractive index	
Kim et al. (2010)	Secondary organic aerosols	670 nm	1.38–1.61	
Liu et al. (2013)	Secondary organic aerosols	220–1200 nm	1.48–1.58	
Shepherd et al. (2018)	Urban, remote, and wood smoke	460–760 nm	$\sim\!1.58$ for wood smoke aerosols and 1.47–1.52 for urban and remote aerosols	
Sumlin et al. (2018)	Brown carbon aerosols	375, 405, 532, 1047 nm	1.5–1.7	
Di Biagio et al. (2019)	Dust aerosols	370, 470, 520, 590, 660, 880, 950 nm	1.48–1.55	
Womack et al. (2021)	Biomass burning	360–720 nm	1.55–1.6	
This study	Major fraction of biomass burning aerosols mixed with minor fraction of dust particles	330–780 nm	$1.53 \pm 0.03$ for fine mode and $1.51 \pm 0.02$ for coarse mode	

Table 1. Examples of reported real part refractive index of biomass burning and dust aerosols from previous studies and the current study.

orbiting Partnership) satellite and the corresponding truecolor image over Southeast Asia on 30 March 2019, when significant amounts of biomass burning aerosols prevailed ( $\tau_{aer}$  at 550 nm higher than 3.0). The DB aerosol algorithm 5 and its extended family have been applied to various spaceborne spectroradiometers such as AVHRR (Advanced Very

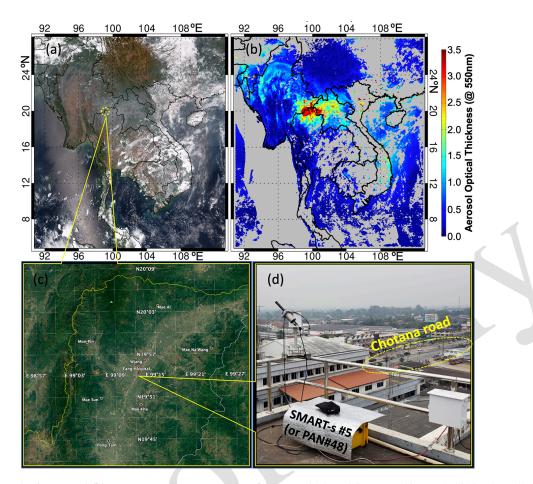
High Resolution Radiometer), SeaWiFS (Sea-viewing Wide Field-of-view Sensor), MODIS (MODerate resolution Imaging Spectroradiometer), VIIRS (Hsu et al., 2019), and current <sup>10</sup> advanced multispectral imagers aboard geostationary satel-

- lites, enabling the construction of long-term aerosol climate data records (CDRs). As previously stated, satellite retrievals provide reliable  $\tau_{aer}$ , as indicated in Fig. 1b; the collocated SMART–s measurement is also presented in the colored cir-
- <sup>15</sup> cle ( $\tau_{aer} = 3.1$ ), which shows an excellent agreement with the DB  $\tau_{aer}$  retrievals nearby (mean  $\tau_{aer} = 2.93$  within 10 km of SMART-s). In an attempt to derive more comprehensive aerosol properties from satellites, previous studies actively utilized UV measurements which are sensitive to the
- <sup>20</sup> aerosol absorption and vertical profile and  $\tau_{aer}$  (e.g., Torres et al., 2013; Lee et al., 2021). Accurate aerosol optical models play a central role in the endeavor. However, due to the lack of a reliable aerosol property database in the UV, they typically made simple assumptions on the spectral features
- <sup>25</sup> of aerosols to extrapolate the properties from longer wavelengths or adopted laboratory measurements. One of the ultimate goals of this study is to contribute to satellite retrievals by providing realistic aerosol optical models over the study domain, particularly in the UV, which will be discussed in <sup>30</sup> Sect. 3.3.

The collocated SMART-s and AERONET instruments are deployed on the rooftop of Fang Hospital in Fang district, as shown in Fig. 1c and d, which is located in a basin in northern Thailand. The population of Fang city was slightly higher <sup>35</sup> than 116 000 in 2010, with a low level of traffic throughout the year. One of the main roads of the city (Chotana Road) is near the building ( $\sim$  50 m). However, we presume that the effects of local emissions from the road to the aerosol and NO<sub>2</sub> amounts are weak, given the low level of local traffic, and that major fractions of the aerosols and trace gases (e.g., 40 NO<sub>2</sub>) during this season are emitted from the biomass burning in this area (Jena et al., 2015; Itahashi et al., 2018; Khodmanee and Amnuaylojaroen, 2021). Figure 1d shows an image of the deployed SMART-s, and Chotana Road is shown behind. Direct-sun measurements of SMART-s started on 45 8 March 2019, and it has acquired additional solar almucantar scans since 19 March after about 10 d of stabilization (e.g., to check the stability under the field conditions and to fine-tune the alignment in tracking). The measurements finished on 2 May 2019. The AERONET instrument is installed on the same rooftop, about 5 m away from SMART-s. In 2019, the surface air temperature at Fang during the premonsoon season reached up to about 42 °C during the daytime, and the relative humidity gradually increased from March ( $\sim 30\%$ ) to early May ( $\sim 50\%$ ). 55

### 2.2 Measurements

The SMART-s instrument was originally developed by the Pandora network group at NASA (National Aeronautics and Space Administration)/GSFC (Goddard Space Flight Center), and the unit (no. 5) used in this study is registered as Pandora no. 48. Most of the components of SMART-s are similar or identical to the standard Pandora instrument, except for the spectrometer. The SMART-s spectrometer is made by the same manufacturer (AvaSpec-ULS2048x64, Avantes; see https://www.avantes.com/, last access: 8 June 2022) as the standard version but covers a wider spectral range (i.e., 280–820 nm) with a lower spectral resolution (~1.0 nm full width at half maxi-



**Figure 1.** (a) True-color image and (b) DB (Deep Blue)  $\tau_{aer}$  product from VIIRS aboard SNPP on 30 March 2019. The yellow circle in panel (a) is the location of Fang, Thailand, and the colored circle in panel (b) shows the collocated  $\tau_{aer}$  retrievals from SMART–s. Panel (c) indicates the location of the measurement site on the rooftop of Fang Hospital, Thailand (19.91° N latitude, 99.21° E longitude, and 480 m above sea level; the map is extracted from https://google.com/maps/, last access: 12 September 2022; © Google Maps 2021). Panel (d) is an image of deployed SMART–s (Pandora no. 48) taken on 8 March 2019. Chotana Road (marked in yellow) is one of the major streets in this area and is about 50 m away from the site.

mum, FWHM, with  $\sim 3.7 \times$  oversampling). As the Pandonia Global Network (PGN; Herman et al., 2009, 2015; see https://www.pandonia-global-network.org, last access: 12 September 2022) is utilizing another type of extended-5 range spectrometer for their dual-detector system, we refer to this modified Pandora as SMART-s in this study. The spectrometer utilizes a 2048 × 64 pixels of back-thinned

- Hamamatsu charge-coupled device (CCD) with a symmetric Czerny–Turner system, and its spectrum covers O<sub>2</sub>, O<sub>3</sub>, <sup>10</sup> NO<sub>2</sub>, SO<sub>2</sub>, HCHO, and H<sub>2</sub>O gas absorption bands (Herman et al., 2015; Jeong et al., 2018). The optical head consists of two rotating filter wheels; one includes neutral density (ND) filters, and the other contains bandpass filters (e.g., U340 and BP300) to block out-of-band (OOB) stray light from
- <sup>15</sup> the near-UV and VIS wavelengths, ground-fused silica diffuser (diffuser hereafter; for NO<sub>2</sub>, H<sub>2</sub>O, and  $\tau_{aer}$  retrieval), and an opaque filter for dark current measurements. By combining a variable exposure time (4–4000 ms) and ND filters,

it can measure radiances with a dynamic range up to an order of 10<sup>7</sup>, which enables the direct-sun and sky scans us- <sup>20</sup> ing a single detector throughout the day. Note that the field of view (FOV) for direct-sun observations using the diffuser of this unit is about 2.8°, which is broadened to evenly distribute light passing through the optical head. Sky observations do not use the diffuser to allow more photons to reach 25 the detector, of which the FOV is about  $1.5^{\circ}$ . The optical head is mounted on a sun/sky scanner and is connected to the spectrometer through a fiber-optic cable of 400 µm in diameter. The spectrometer is thermoelectrically controlled to maintain a near-constant temperature but may vary slightly 30 depending on the ambient temperature (typically less than 1 °C). The spectrometer temperature is recorded with each measurement to monitor data quality.

The SMART-s algorithm aims to obtain optimal information on aerosols and trace gases with minimum assumptions, <sup>35</sup> which incorporates a series of retrievals from fundamental quantities (i.e., column amounts) to higher-order geophysical parameters (e.g., aerosol physicochemical properties and vertical profiles). Jeong et al. (2018) developed a  $\tau_{aer}$  algorithm of the SMART-s based on the spectral Langley method and then compared the retrievals to collocated AERONET

- measurements at the NASA/GSFC, which showed excellent agreements at all overlapping wavelengths (i.e., 330, 380, 440, 500, and 675 nm). Comparisons of the  $\tau_{aer}$  from the AERONET and SMART-s during this field deployment are
- <sup>10</sup> shown in the Appendix (see Figs. A1 and A2). The trace gas algorithm of SMART-s is designed for a relatively lower spectral resolution (FWHM $\sim$  1.0 nm) and broader spectral coverage (280–820 nm) compared to the standard Pandora. For retrieving optically thick trace gases such as O<sub>3</sub> and H<sub>2</sub>O,
- <sup>15</sup> we utilize the spectral Langley method (see Appendix B; Jeong et al., 2018), whereas we adopt the spectral-fitting algorithm of Pandora for other optically thin species, including NO<sub>2</sub> (Herman et al., 2009). Retrieved total columns of  $\tau_{aer}$ , O<sub>3</sub>, and H<sub>2</sub>O in this study are compared with those from
- <sup>20</sup> the collocated AERONET (for  $\tau_{aer}$  and H<sub>2</sub>O) and satellite retrievals (O<sub>3</sub> from the OMI and  $\tau_{aer}$  from the VIIRS) during the measurement period in Sect. 3.2 and Appendix A. Jeong et al. (2020) developed an OEM-based algorithm using solar almucantar sky radiances and total column retrievals (e.g.,
- <sup>25</sup>  $\tau_{aer}$ , O<sub>3</sub>, NO<sub>2</sub>, and H<sub>2</sub>O) for retrieving spectral *n*, *k*,  $\omega_0$ , and PSD of aerosols; details of the aerosol-column-property algorithm are in Jeong et al. (2020) and the key characteristics are included here in Appendix C.

For more than 2 decades, AERONET has been supported 30 by NASA to operate a global network of automatic sun/sky scanning spectroradiometers for acquiring aerosol information (Holben et al., 1998). The instrument measures discrete channels (i.e., 340, 380, 440, 500, 675, 870, 940, 1020, and 1640 nm) of solar irradiance with a 1.2° FOV, which takes 35 about 10 s to scan all spectral filter wheels. The FWHMs of the bandpass filters are 2 nm for 340 and 380 nm, 25 nm for 1640 nm, and 10 nm for all other channels, whereas that of SMART-s is about 1.0 nm for all wavelengths. The estimated uncertainty of  $\tau_{aer}$  from the AERONET reference instrument 40 is 0.002, and those from general network instruments are about 0.01 in the VIS-NIR and are higher ( $\sim 0.02$ ) in the UV channels (Eck et al., 1999; Giles et al., 2019). Note that the uncertainty of  $\tau_{aer}$  from SMART-s (~0.02 in the VIS-NIR;  $\sim 0.03$  in the UV) is slightly higher than AERONET  $_{45}$  (~0.01 in the VIS-NIR; ~0.02 in the UV) due to the wider FOV (Jeong et al., 2018), which is more susceptible to for-

ward scattering and the temperature sensitivity of the detector (Kinne et al., 1997). The current AERONET product provides *n*, *k*, and  $\omega_0$  at 440, 675, 870, and 1020 nm. As the <sup>50</sup> version 3 algorithm utilizes a vector radiative transfer model

(Korkin et al., 2017), it can add a 380 nm channel for the UVabsorbing aerosols (Sinyuk et al., 2020). The recent version of instruments added hybrid sky scanning measurements to allow additional retrievals at the solar zenith angle ( $\theta_S$ ) be-

55 low 50° (Sinyuk et al., 2020). To consider gas absorption,

the version 3 algorithm adopts a monthly climatology (1978–2004) of  $O_3$  from the Total Ozone Mapping Spectrometer (TOMS) and a monthly climatology of NO<sub>2</sub> (2004–2013) from the OMI and retrieved H<sub>2</sub>O using the 940 nm channel measurements (Sinyuk et al., 2020). Further information on <sup>60</sup> AERONET products is summarized in Giles et al. (2019) and Sinyuk et al. (2020). We utilized the version 3 and level 2.0 products to compare retrievals from SMART–s.

### 2.3 A combinative radiometric calibration method for sun/sky spectroradiometer

The standard calibration procedure of SMART-s includes spectral characterization/registration, linearity and offset correction, radiometric calibration, temperature and flat field correction, and stray light correction (Herman et al., 2015; Jeong et al., 2018; Müller et al., 2020). 70 The PGN also regularly reports updates and standard calibration/validation results on its web page (https:// www.pandonia-global-network.org, last access: 12 September 2022; Herman et al., 2009, 2015). Utilization of the absolute sky radiances requires precise radiometric calibrations, 75 which we suggest as a novel and combinative method for sun/sky spectroradiometers in this study. As this field campaign is the first attempt to deploy a radiometrically calibrated SMART-s using the method, we summarize detailed results of the calibration for the sky scans in this section. 80

For the initial step of the radiometric calibration, we utilized a uniform spectral radiance light source in the Radiometric Calibration Laboratory (RCL) at NASA/GSFC. The RCL is a class 10000 clean room facility that maintains a number of NIST-traceable integrating sphere sources. The 85 integrating sphere source used in this study is referred to as Grande. Grande is a Spectralon-lined, 101.6 cm diameter, integrating sphere source with a 25.4 cm diameter output aperture, which can generate nine levels of light output. More detailed information and annual calibration reports of 90 Grande are available at https://cf.gsfc.nasa.gov/ (last access: 12 September 2022) or in Gatebe et al. (2007). Figure 2a is an image of SMART-s mounted in front of the Grande sphere source, and Fig. 2b shows Grande's nine levels of spectral output in radiance units. Figure 2c presents the re- 95 ported total uncertainty of Grande sphere's spectral radiance at an approximate 95 % confidence level when calibrated using a NIST irradiance standard. Different colors in Fig. 2b and c indicate the different levels of Grande intensity. Due to relatively low intensity in the UV compared to the VIS- 100 NIR (Fig. 2b) in both the NIST-calibrated irradiance standard and the Grande source itself, the calculated uncertainty of the Grande radiance calibration in the UV is higher, and a brighter light output provides a more accurate intensity, as shown in Fig. 2c. During the light source calibration, the sen- 105 sor changed its filter (UV bandpass) to detect the lower intensity of Grande in the UV, which results in relatively higher uncertainties near 350 nm (see Fig. 2c). Note that the mea-

surement error covariance matrix of the OEM also accounts for such spectral radiometric uncertainties (see Appendix C and Jeong et al., 2020). SMART–s repeated measurements of Grande 10 times for each filter combination (bandpass fil-

- <sup>5</sup> ters and neutral density filters). The deployment procedure of SMART-s (or Pandora) includes the organization of the fiber-optic cable and the connection of one end of the cable to the spectrometer (the other end of the cable is fixed to the optical head). This process can affect light transmit-
- <sup>10</sup> tance through the cable. To check the stability of the fiberoptic cable during deployment, we oriented the cable differently (i.e., re-rolled the cable every time with different diameters or arbitrarily oriented it) and then reconnected the ports to the spectrometer at each time of the Grande mea-
- <sup>15</sup> surements. Contamination of the front window of the optical head (e.g., raindrops, dew, dust, and insects) is also one of the largest error sources of radiometric measurements. During field deployments, we frequently check the front window and clean it when it is necessary. However, as cleaning the
- <sup>20</sup> front window also can alter its optical transmittance, we artificially contaminated the front window (using a finger, dust, and water) and cleaned it in the same way as we do during field deployments every time there are Grande observations. In addition, other sources of short- and long-term temporal
- <sup>25</sup> drift (e.g., spectrometer and filter transmittance) are monitored by pre- and post-mission calibration and the Langley fitting during deployments. Figure 3a shows the results of 10 repeat measurements of Grande at its nine-lamp illumination level. The agreement of these repeat measurements in-
- <sup>30</sup> dicates good temporal stability in the SMART-s responsivity and the Grande output. Figure 3b shows the spectral calibration coefficient calculated from dividing the Grande intensity (Fig. 2b) by the average value of the measured voltage count (Fig. 3a), while the spectral precision of the Grande calibra-
- <sup>35</sup> tion (i.e., 1 standard deviation of 10 occurrences of the measurements) is presented in Fig. 3c. These results indicate that the precision of the radiometric calibration from the instrument is better than 0.5% in the VIS-NIR channels and increases at shorter wavelengths to about 0.7% at 330 nm (i.e.,
- <sup>40</sup> the lower limit of the spectral coverage in this study). Radiometric sensitivity of the spectrometer to its temperature  $(T_{\text{spec}})$  is also tested by controlling the  $T_{\text{spec}}$ , which is less than 0.4% at the entire spectral range for an extreme  $T_{\text{spec}}$ variability (i.e.,  $\Delta T_{\text{spec}} \sim 3 \,^{\circ}$ C, which is less than 1  $^{\circ}$ C un-<sup>45</sup> der typical field conditions). In general, the uncertainty of
- the laboratory radiance calibration including the light source and instrument stability is estimated to be better than 2.0% in the VIS-NIR and 3.0% in the UV at an approximate 95% confidence level.
- As typical radiative transfer models (e.g., Spurr, 2006; Stamnes et al., 1988) simulate normalized radiances (i.e., radiance divided by solar irradiance as a unit of inverse steradian), the reference solar spectrum is also a key parameter for converting raw voltage counts to sky radiances of the identi-

55 cal unit of forward model calculations. Satellite instruments

can directly measure solar irradiance using the same detector with a similar optical path of the Earth reflectance measurements so that a major fraction of calibration uncertainties (e.g., slit function and radiometric coefficient) is canceled out. Although ground-based instruments also measure the so- 60 lar light using an identical detector to the one used for sky radiances, they sample the solar irradiance after it has passed through the atmosphere. For that reason, their algorithms utilize other sources of the solar spectrum or estimate it from the measurements for the conversion of the sky radiances. 65 The version 3.0 AERONET inversion algorithm utilizes solar irradiance from NOAA's (National Oceanic and Atmospheric Administration) Climate Data Record (Coddington et al., 2016; Sinyuk et al., 2020), and SKYNET derives a conversion factor of the sky radiances from direct-sun measure- 70 ments based on the solid-view-angle estimation algorithm (Uchiyama et al., 2018a, b). For hyperspectral instruments such as SMART-s, a combination of the high-resolution solar spectrum and calibrated slit function is a key factor for retrieving the spectral aerosol properties, which is particu-75 larly important at wavelengths shorter than 500 nm, where spectral variability in the solar irradiance is significant (e.g., see Fig. 4a).

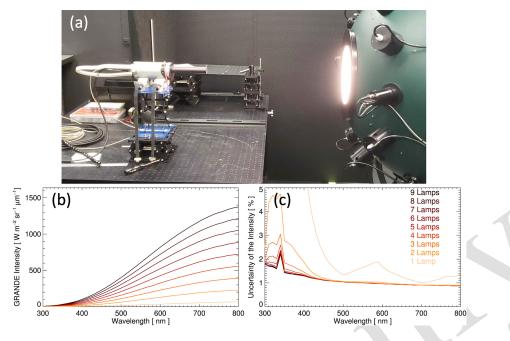
In the second step of the radiometric calibration, the SMART-s algorithm estimates the reference solar spectrum <sup>80</sup> by combining direct-sun measurements, laboratory calibrations, and ancillary solar irradiance data. Raw voltage counts of the sky scan measurements without bandpass filters can be expressed as follows (e.g., Uchiyama et al., 2018b):

$$V(\lambda) = C(\lambda) \int_{\Omega_{\rm FO}} f(\Omega) I(\lambda, \Omega) \,\mathrm{d}\Omega, \qquad (1) \quad \text{se}$$

where *V* is the voltage count of a sky scan measurement,  $\lambda$  is the wavelength, *C* is the sensitivity of the detector to the radiance,  $\Omega_{\text{FO}}$  is the solid angle of the instrument's field of view (FOV) without bandpass filters, *f* is the response function of the radiometer's FOV, and *I* is the sky radiance. To avoid saturation, the direct-sun measurements utilize the diffuser with a spectral transmittance of  $T_d(\lambda)$ , which is measured using Grande. The measured voltage count of the solar scan can be described as follows:

$$V_{\rm S}(\lambda) = C(\lambda) T_{\rm d}(\lambda) \int_{\Omega_{\rm S}} f(\Omega) I_{\rm D}(\lambda, \Omega) d\Omega + C(\lambda) T_{\rm d}(\lambda) \int_{\Omega_{\rm FD}} f(\Omega) I(\lambda, \Omega) d\Omega, \qquad (2) \quad {}_{95}$$

where  $V_{\rm S}$  is the voltage count of direct-sun measurements,  $\Omega_{\rm S}$  and  $\Omega_{\rm FD}$  are the solid angle of the Sun and FOV of SMART-s with the diffuser, respectively.  $I_{\rm D}$  is a direct component of solar measurements. The first term on the righthand side of Eq. (2) describes the contribution of direct solar irradiance, and the second term shows the scattered radiance within the FOV. Here we assume that  $\Omega_{\rm S}$  and  $\Omega_{\rm FD}$ 



**Figure 2.** (a) Image of SMART-s calibration using the NIST-traceable light source (Grande) at the Radiometric Calibration Laboratory, NASA Goddard Space Flight Center. Panel (b) shows the spectral radiance of Grande in 300–800 nm and (c) presents its reported uncertainty. Different colors in panels (b) and (c) indicate nine levels of the Grande radiance. The relatively higher values of uncertainty near 350 nm in panel (c) are due to the filter change of the sensor during the light source calibration.

are not wavelength-dependent within the SMART-s spectral range. We select the Langley calibration dates when  $\tau_{aer}$  at 500 nm is less than 0.05 to minimize aerosol impacts and screen cloud-contaminated measurements. In addition, the

<sup>5</sup> SMART-s algorithm corrects the contribution of Rayleigh scattering of the direct-sun measurements, which is larger in the shorter wavelength (Jeong et al., 2018). Based on this process and these criteria, we assume that the second term on the right-hand side of Eq. (2) is negligible. How-<sup>10</sup> ever, unscreened thin cirrus clouds may generate diffuse light within the FOV mere effectively the success due to

within the FOV more effectively than the aerosols due to their stronger forward scattering (e.g., Kinne et al., 1997). For plane-parallel solar irradiance measurements (f = 1),  $V_S$  can be approximated as follows:

$$V_{\rm S}(\lambda) \sim C(\lambda) T_{\rm d}(\lambda) \int_{\Omega_{\rm S}} I_{\rm D}(\lambda, \Omega) d\Omega$$
$$= C(\lambda) T_{\rm d}(\lambda) F_{\rm BOA}(\lambda), \qquad (3)$$

15

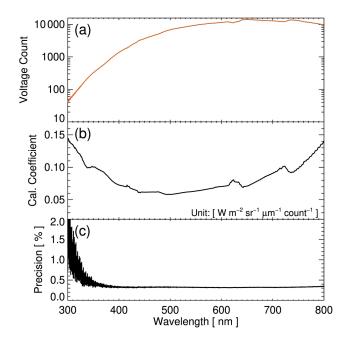
where  $F_{\text{BOA}}$  is the solar irradiance at the bottom of the atmosphere, which can be described as follows:

$$F_{\text{BOA}}(\lambda) = \int_{\Omega_{\text{S}}} I_{\text{D}}(\lambda, \Omega) \, \mathrm{d}\Omega$$
$$\equiv F_{V}(\lambda) \left[ \sum_{i} m_{i}(\lambda) \tau_{i} \left( \sum_{j} m_{j}(\lambda) \tau_{j} \right) \right]$$
(4)

In Eq. (4),  $F_V$  and  $\tau$  denote, respectively, the extraterrestrial solar spectrum and optical thickness of the atmospheric constituents, which are derived from the Langley calibration by using the  $V_S$  in Eq. (3). *m* is the optical air mass of each atmospheric species (e.g., aerosols, clouds, and gases). Then, the last step is normalizing the  $F_V$  to known solar irradiance data for minimizing the remaining systematic calibration error by the following equation:

$$F_{\text{Comb}}(\lambda) = \frac{\overline{F}_{\text{Trad}}}{\overline{F}_{V}} F_{V}(\lambda), \qquad (5)$$

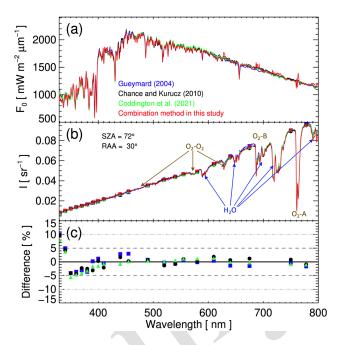
where  $\overline{F}_{\text{Trad}}$  is the spectral mean value of solar irradiance using the traditional method (i.e., a high-resolution reference solar spectrum convoluted by the instrument's slit function; 30  $F_{\text{Trad}}$ ), and  $\overline{F}_V$  is that of  $F_V$  at wavelengths between 490– 510 nm. F<sub>Comb</sub> is the final solar irradiance for the SMART-s algorithm derived by combining the laboratory/Langley calibration and reference spectrum (i.e., Coddington et al., 2021 in this study). This spectral window (490–510 nm) is near the 35 middle of the detector, and the solar intensity is high, with relatively fewer spectral variabilities. In addition, this spectral range avoids strong gas absorptions. By using the spectral shape of the  $F_{\text{Comb}}$ , we expect that uncertainties generated from the calibrated slit function are minimized. Despite 40 sun/sky measurements undergoing empirical OOB stray light correction (Jeong et al., 2018), the remaining fraction may still be non-negligible in the shorter wavelengths of UV (particularly wavelengths shorter than about 330 nm; see Fig. 7c



**Figure 3.** (a) An example of the nine-lamp Grande voltage count measurements from SMART-s (Pandora no. 48) without neutral density or bandpass filters. All 10 colors are plotted as in Fig. 2, indicating the average current measurements over 10 repetitions, which overlap almost completely. Panel (b) is the calibration coefficient, which is calculated from dividing the known Grande intensity by the average value of the measured voltage count. Panel (c) is the precision of the calibration coefficient, which is estimated by calculating 1 standard deviation of the repetitions done 10 times.

of Jeong et al., 2018). However, as  $F_{\text{Comb}}$  is supposed to be affected by OOB stray light with a comparable degree of the sky scan measurements due to their similar spectral shape, it can partially cancel out the remaining portion of stray light

- $_5$  in the sky radiances. Figure 4a compares  $F_{\text{Comb}}$  to convoluted solar irradiance from Gueymard (2004), Chance and Kurucz (2010), and Coddington et al. (2021), which shows generally good consistency in the entire spectral range of SMART-s. Figure 4b shows an example of spectral radi
- <sup>10</sup> ances using the different solar spectrums in Fig. 4a. The colored symbols in Fig. 4b indicate the wavelength node of the aerosol retrieval, which is carefully selected to avoid strong absorption bands of atmospheric gases and major calibration errors discussed above. Figure 4c depicts relative biases
- <sup>15</sup> of the sky radiances, which compares  $F_{\text{Comb}}$  to those using Gueymard (2004), Chance and Kurucz (2010), and Coddington et al. (2021) at the wavelength nodes convoluted from using the calibrated slit function. The biases are generally smaller than about 2 % at wavelengths longer than 500 nm
- <sup>20</sup> and are higher in the shorter wavelength up to about 10% at 330 nm. The relatively high discrepancies between  $F_{\text{Comb}}$  and the other solar irradiances in the UV are attributable to uncertainties in slit function, remaining OOB stray light, and  $F_V$ . Note that uncertainties in the slit function affect  $F_{\text{Trad}}$ ,



**Figure 4. (a)** Spectral solar irradiances ( $F_0$ ) from Gueymard (2004; blue), Chance and Kurucz (2010; black), Coddington et al. (2021; green), and a combination method developed in this study ( $F_{\text{Comb}}$  in red). Panel (**b**) is an example of converted normalized radiance (radiance divided by solar irradiance;  $F_0$ ) measured at Fang, Thailand, on 19 March 2019. Colored lines indicate different sources of  $F_0$  for the conversion (same as panel **a**), and circles, triangles, and rectangles depict selected wavelengths of aerosol inversion in this study. Panel (**c**) presents relative biases of the  $F_{\text{Comb}}$  compared to those of Gueymard (2004; blue rectangle), Chance and Kurucz (2010; black circle), and Coddington et al. (2021; green triangle).

which may be significant in the UV. We estimate that the accuracy of  $F_V$ , which does not require spectral convolution, is better than 4 % in the UV and 2 % in the VIS-NIR, based on the accuracy of the spectral  $\tau_{aer}$  retrievals. Therefore, the total error of the sky radiance is estimated to be better than 5 % in the UV and 3 % in the VIS-NIR at an approximate 95 % 30 confidence level. The impacts of the different sources of the solar spectrum on aerosol retrievals are discussed in Sect. 3.

We applied the aerosol retrieval algorithm to the measurements (both direct sun and solar almucantar scan) with sufficient amounts of photons within the target spectral range (i.e., from 330 to 800 nm), as very high aerosol loading over the area (e.g., Fig. 1b) may result in the low level of voltage counts below the detection limit (e.g., in terms of linearity and noise). Cloud-contaminated direct-sun spectra were screened by using their rapid temporal variability and spectral features (i.e., lower Ångström exponent of clouds), which are described in Jeong et al. (2018). Those for the solar almucantar measurements were removed by checking the horizontal symmetry of the scan (i.e., between clockwise and counter-clockwise half-circle scans) followed by the AERONET strategy (Jeong et al., 2020).

#### 3 Results

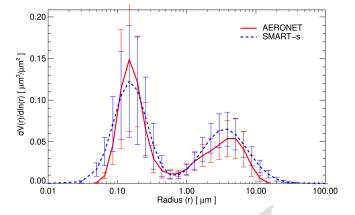
### 3.1 Comparison with the AERONET

- <sup>5</sup> Jeong et al. (2020) applied the SMART–s algorithm to yearlong AERONET sun/sky measurements in 2016 at Kanpur, India, to assess the consistency of the methodology. The retrieved volume size distribution, V(r), from SMART–s showed excellent agreements in the fine mode but mi-
- <sup>10</sup> nor discrepancies in the coarse mode due to the different assumptions and constraints between the two algorithms; AERONET retrieves V(r) at 22 radius nodes over the optically effective range (i.e., from 0.05 to 15.0 µm), constrained by smoothness together with k, whereas SMART-s assumes
- <sup>15</sup> a bimodal, lognormal distribution of the number size distribution N(r) (Dubovik and King, 2000; Jeong et al., 2020). Spectral  $\omega_0$  showed excellent agreements for all wavelengths from 440 to 1020 nm, with *R* ranging from 0.87 to 0.95 and the RMSE (root mean squared error)/MBE (mean bias <sup>20</sup> error) was less than 0.012 during that year-long period. In this section, we performed additional comparisons of aerosol property retrievals from SMART-s to those collocated from

AERONET by utilizing their own measurements.

Figure 5 depicts coincident V(r) values from AERONET <sup>25</sup> and SMART–s, where major fractions of the aerosols are fine-mode smoke particles and fewer but non-negligible portions are coarse particles (e.g., transported dust from the Sahara, Thar Desert, and dry areas of the Indo-Gangetic Plain). Retrieved V(r) values from both instruments show generally

- <sup>30</sup> good agreement, which is consistent with a previous study (Jeong et al., 2020). Regarding the fact that SMART-s V(r)retrievals showed better agreement with AERONET when it is applied to the same measurements (Jeong et al., 2020), another major fraction of the discrepancies in Fig. 5 is likely at-
- st tributable to the different types of measurements (e.g., spectral information and radiometric calibration). Note that the SMART-s measurement is not sensitive to aerosols with a radius greater than the optically effective range ( $\sim 10 \mu m$ ; see Fig. 10 in Jeong et al., 2020), and the SMART-s V(r)
- <sup>40</sup> over this range (see long tails of the blue dashed line) is mostly generated by assuming the lognormal shape of the coarse mode. Further studies to derive optimal information on aerosol size are underway (e.g., additional parameters for size distribution and/or additional modes). According to the-
- <sup>45</sup> oretical error estimates of the OEM-based algorithm in Jeong et al. (2020), a high spectral resolution of SMART-s is beneficial for the fine mode, whereas a broader spectral range of AERONET is advantageous for both modes under the same level of radiometric accuracy. More detailed compar-<sup>50</sup> isons and relevant discussions on the V(r) between the two
- algorithms are summarized in Jeong et al. (2020).



**Figure 5.** Average volume size distribution of aerosols retrieved from SMART-s (blue dashed line) and AERONET (red solid line) from 19 March to 2 May 2019 at Fang, Thailand. The variability in the retrievals (standard deviations at each radius node) during the period are represented as vertical bars in this figure. The SMART-s spectral range is not sensitive to aerosols with a radius greater than about 10  $\mu$ m, and the SMART-s retrievals over this range (see long tails of the blue dashed line) are mostly determined by the lognormal shape assumption, thus with small variability. Note that the vertical bars do not represent retrieval errors at each radius node.

To better understand and assess the two PSD retrievals at each size bin, a precise evaluation through reliable in situ measurements, such as aircraft profiles from DISCOVER-AQ (Deriving Information on Surface Conditions from Col- 55 umn and Vertically Resolved Observations Relevant to Air Quality), KORUS-AQ (KORea-U.S. Air Quality), and a sUAS, are essential. However, to the best of our knowledge, very limited studies compared AERONET V(r) to collocated in situ profile measurements of the PSD. Chauvi- 60 gné et al. (2016) compared the PSD from AERONET (at 410 m a.s.l., above sea level) to in situ measurements at a higher-altitude site in central France (i.e., 1465 m a.s.l.) over a 1-year period, which showed a relative underestimation of AERONET ( $\sim$  40 %). However, the in situ measurement site 65 in Chauvigné et al. (2016) may not fully represent the total column values, which can be associated with the biases. Schafer et al. (2019) compared the AERONET V(r) to in situ aircraft profiles (altitude from about 150 to 5000 m and radius range from 0.03 to 0.5 µm) over Maryland, California, 70 Texas, and Colorado in the United States. They showed that fine-mode PSD parameters derived from AERONET and in situ aircraft measurements generally showed good agreement (average difference of the radius of peak concentration is about 0.011  $\mu$ m and that of the V(r) width is about 0.03  $\mu$ m), 75 whereas differences in the V(r) values depend on the particle radius and location (see Fig. 5 of Schafer et al., 2019). SMARTLabs (Surface-based Mobile Atmospheric Research & Testbed Laboratories; https://smartlabs.gsfc.nasa.gov, last access: 12 September 2022) is developing a sUAS-based 80 aerosol profiling instrument for collocated measurements of

spectral absorption and size distribution, which can provide valuable data for assessments of the PSD and  $\omega_0$  from SMART–s. More precise validation/comparison studies are currently underway.

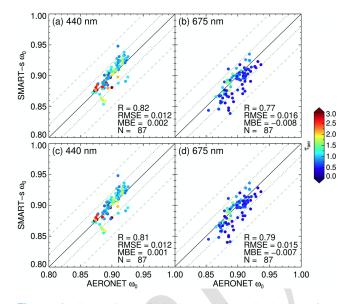
- <sup>5</sup> Figure 6 compares the retrieved  $\omega_0$  from SMART-s and AERONET at overlapping wavelengths (i.e., 440 nm in the left panels and 675 nm in the right panels) during the measurement period. The  $\omega_0$  retrievals from SMART-s in the upper panels of Fig. 6a and b used  $F_{\text{Trad}}$ , whereas those <sup>10</sup> in lower panels (Fig. 6c and d) utilized  $F_{\text{Comb}}$ . In gen-
- eral, all cases showed acceptable agreements, with absolute MBE and RMSE less than 0.02 and *R* ranging from 0.77 to 0.82. The  $\omega_0$  of SMART-s was better correlated with that of AERONET in the shorter wavelength (i.e., 440 nm) <sup>15</sup> due to the higher sensitivity of aerosols from the higher
- $\tau_{aer}$ . In addition,  $F_{Comb}$  generated more consistent  $\omega_0$  retrievals of SMART-s with the AERONET than those using the  $F_{Trad}$ , with a slightly lower RMSE/MBE and higher *R*. Figure 7a presents the mean values of the spectral  $\tau_{aer}$
- <sup>20</sup> from SMART-s and AERONET during the measurement period, which showed excellent agreement over the SMART-s spectral coverage. Figure 7b shows those of the spectral  $\omega_0$ from AERONET and SMART-s using different solar irradiances. While both versions of the  $\omega_0$  values from SMART-s
- <sup>25</sup> exhibit strong absorption in the UV, the one using  $F_{\text{Comb}}$  showed smoother spectral variability, particularly in the UV, which is selected for the SMART–s retrievals. Slight relative low biases of the  $\omega_0$  from SMART–s compared to that from AERONET were also found when the SMART–s algorithm

<sup>30</sup> was applied to the same AERONET measurements at Kanpur, India, which are, however, still within the uncertainty range of the AERONET and SMART–s retrievals (Jeong et al., 2020). Note that AERONET also provides  $\tau_{aer}$  and  $\omega_0$  at longer wavelengths (e.g., 870 and 1020 nm), which are not <sup>35</sup> presented in this figure.

In Fig. 8a, the spectral n of the fine  $(n_f)$  and coarse mode  $(n_c)$  retrieved from SMART-s shows comparable values with the previous studies in Table 1, ranging from about 1.5 to 1.55 with smooth spectral dependencies in the UV; higher values of n of the coarse mode were found in the UV, whereas those of the fine mode were lower. These values were slightly higher than those from the collocated AERONET by about 0.01–0.04. Note that AERONET retrieves a value n for all particle sizes, whereas SMART-s retrieves each size mode 45 (see Appendix C or Jeong et al., 2020), which may result in these differences under the assumption of the lognormal

- size distribution. Regarding the fine-mode-dominated smoke aerosols over the site (see Fig. 5), *n* from AERONET largely represents the contributions of the fine mode, and differ-50 ences between  $n_{\rm f}$  from SMART-s and *n* from AERONET
- are within the known uncertainties of AERONET (Sinyuk et al., 2020) and SMART-s (Jeong et al., 2020). As reported in previous studies and as seen in the spectral  $\omega_0$  retrievals in Fig. 7b, the spectral k from SMART-s increased signif-

55 icantly in the shorter wavelengths, particularly in the fine

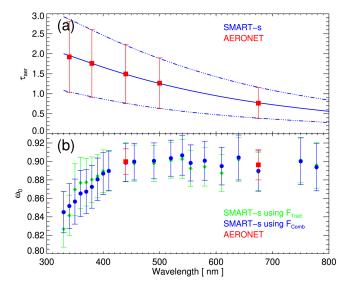


**Figure 6.** Comparison of single scattering albedo ( $\omega_0$ ) from SMART-s and AERONET (version 3, level 2.0) at (**a**) 440 nm and (**b**) 675 nm from 19 March to 2 May 2019. In these two panels, the SMART-s algorithm utilized solar irradiance from Coddington et al. (2021). Panels (**c**) and (**d**) are similar plots to panels (**a**) and (**b**) but for SMART-s retrievals using derived solar irradiance in this study. Colored circles represent values of aerosol optical thickness ( $\tau_{aer}$ ) from the color bar at each wavelength. *R* is the correlation coefficient, RMSE denotes the root mean square error, MBE is the mean bias error, and *N* is the number of samples for the comparison. The dotted and dashed lines represent, respectively, the relative biases of  $\pm 0.02$  and  $\pm 0.05$  from the AERONET product.

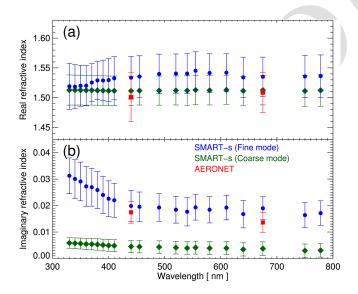
mode (see Fig. 8b). However, as the k retrievals of each mode have cross-correlated measurement sensitivity, separate analysis of each mode may have relatively larger uncertainties than the case of the n retrievals (see Appendix C, Fig. C1). More detailed temporal and spectral analyses of aerosol optical properties from SMART-s are given in the following sections, and scatterplots of the n, k, and PSD between the two instruments are shown in Appendix D.

60

Validation or comparison of the UV aerosol properties from SMART-s remains challenging due to the limited co- 65 incident measurements. As discussed for Fig. 5, in situ profile measurements aboard the aircraft or sUAS platforms may provide reliable sets of data for validation. UV aerosol properties retrieved from other collocated instruments (e.g., Mok et al., 2018; Nakajima et al., 2020) can also offer useful 70 data for checking the consistency and redundancy. The spectral  $\omega_0$  retrievals from SMART-s are mainly determined by the ratio of the sky radiances to the  $\tau_{aer}$ ; highly absorbing aerosols result in the lower level of the sky radiances at a given  $\tau_{aer}$ . Reliable spectral  $\tau_{aer}$  from SMART–s from 330 to 800 nm supports its consistent radiometric performance (e.g., linearity and OOB stray light) over the wavelength range (Jeong et al., 2018). As SMART-s measures sky radiances



**Figure 7.** Mean values of (**a**) aerosol optical thickness ( $\tau_{aer}$ ) from SMART-s (blue line) and AERONET (red square) within 330–800 nm spectral range measured from 19 March to 2 May 2019 at Fang, Thailand. Panel (**b**) shows the spectral single scattering albedo of aerosols ( $\omega_0$ ) from AERONET (red rectangle) and SMART-s, using different solar irradiance; green diamonds used Coddington et al. (2021), and blue circles used the spectrum derived in this study. Variabilities (1 standard deviation) of each value during the deployment period are shown as dash-dotted lines in panel (**a**) and vertical bars in panels (**a**) and (**b**).



**Figure 8.** Mean values of (**a**) the real part of spectral refractive index from AERONET (red square) and SMART-s (blue circle is fine mode; green diamond is coarse mode) within the SMART-s spectral range measured from 19 March to 2 May 2019 at Fang, Thailand. Panel (**b**) shows those of imaginary part. Variabilities (1 standard deviation) in each value during the deployment period are demonstrated as vertical bars in panels (**a**) and (**b**).

using the identical detector, but without a diffuser, we expect comparable radiometric accuracy of the spectral sky measurements with relatively higher uncertainties in the UV, as discussed in Sect. 2.3. Note that our best estimate of the accuracy of the spectral radiance is demonstrated in the measurement error covariance matrix, which is considered in the estimated retrieval error (see Appendix C and Jeong et al., 2020).

### 3.2 Temporal variations and relationship between aerosol properties and trace gases

Total column amounts of trace gases and aerosols are basic and essential quantities not only for understanding their amounts and significance with variations over time but also for providing key constraints for higher-order retrievals (e.g., Jeong et al., 2018, 2020). Direct-sun measurements of 15 SMART-s, AERONET, and Pandora provide very accurate retrievals of these parameters, which thereby have been used to validate/compare various satellite products over the globe. This section presents general characteristics and temporal trends of the basic quantities retrieved from SMART-s (i.e., 20  $\tau_{aer}$  and total column amounts of O<sub>3</sub>, NO<sub>2</sub>, and H<sub>2</sub>O) and then analyzes higher-order retrievals of aerosol properties by comparing them with these quantities.

Figure 9 shows temporal variations in the  $\tau_{aer}$  at 500 nm and total column amounts of NO<sub>2</sub>, H<sub>2</sub>O, and O<sub>3</sub> during the entire deployment period from SMART–s and AERONET. The red circles in Fig. 9a demonstrate aerosol optical thickness at 550 nm from the VIIRS DB product, and those in Fig. 9d depict OMI retrievals based on the TOMS version 8.5 algorithm (Bhartia and Wellemeyer, 2002). In general, the  $\tau_{aer}$ , H<sub>2</sub>O, and O<sub>3</sub> retrievals from SMART–s showed excellent agreement with those from the AERONET, VIIRS, and OMI during the measurement period (see the statistical values in Fig. 9a, c, and d). Scatterplots of these parameters are shown in the Appendix (Fig. A2).

Itahashi et al. (2018) analyzed high values of NO<sub>2</sub> from winter to pre-monsoon seasons over Southeast Asia based on satellite retrievals (i.e., SCanning Imaging Absorption spectroMeter for Atmospheric CHartographY, or SCIAMACHY; Bovensmann et al., 1999) and model calculations (i.e., Com-40 munity Multi-scale Air Quality, or US EPA Office of Research and Development, 2021), and reported that emissions from biomass burning are attributable to the seasonal variation. They also estimated the contributions of biomass burning emissions to the total column NO<sub>2</sub> to be about 28%45 during 2003-2008, which was up to 58 % in March 2004. Khodmanee and Amnuaylojaroen (2021) estimated the contribution of biomass burning to the NO<sub>2</sub> concentration over northern Thailand in March 2014 to be higher than 90%, based on the WRF-Chem (Weather Research and Forecasting 50 model with Chemistry) calculations. Another WRF-Chem study suggested that increase in NO<sub>2</sub> due to the biomass burning emission is up to about 60% over Southeast Asia

**Figure 9.** Temporal variations in total columns of (a) aerosol extinction at 500 nm, (b) nitrogen dioxide, (c) precipitable water vapor, and (d) ozone at Fang, Thailand, in 2019. The blue circles are from SMART-s retrievals, and the gray circles in panels (a) and (c) show those from AERONET. In panel (a), the red circles indicate aerosol optical thickness at 550 nm from VIIRS Deep Blue (DB), while those in panel (d) depict total column ozone retrievals from OMI (TOMS version 8.5). The correlation coefficient (R), root mean squared error (RMSE), and mean bias error (MBE) in panels (a), (c), and (d) are between collocated SMART-s and AERONET/OMI data in black, and those in red in panel (a) are between SMART-s and VIIRS DB retrievals.

from March to May 2005 (Jena et al., 2015). During the measurement period, a large amount of NO<sub>2</sub> from SMART-s is also accompanied by high  $\tau_{aer}$  at 500 nm, with a correlation coefficient (*R*) of about 0.74, which indicates the common 5 emission sources (i.e., biomass burning) of aerosols and NO<sub>2</sub>

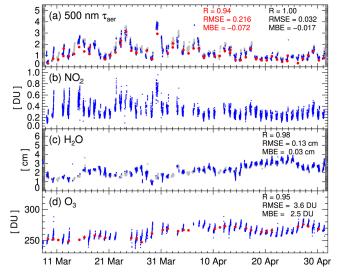
during the events, as shown in Fig. 9a and b.

Interaction between atmospheric  $H_2O$  and aerosols is one of the primary factors in determining aerosol scattering and absorption properties (e.g., Burgos et al., 2019). Particularly

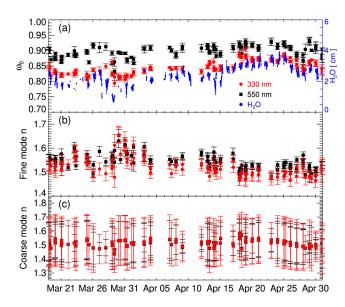
- <sup>10</sup> for organic aerosols, molar mass and water content are dominant parameters for characterizing their phase state (e.g., Koop et al., 2011). The moisture-induced phase transition of organic particles from a glassy to a semisolid state also accelerates the uptake of reactive gases in the atmosphere
- <sup>15</sup> by decreasing viscosity and increasing diffusivity (Shiraiwa et al., 2011). In addition, the condensed water provides a medium for multiphase reactions, thus activating the gas-toparticle conversion of inorganic and organic molecules (Herrmann et al., 2015). The formation of the secondary species
- <sup>20</sup> through heterogeneous reactions generates further feedback to the aerosol-atmosphere system by enhancing the water vapor absorption and hygroscopicity of aerosol particles (e.g., Tang et al., 2016; Wu et al., 2018). Serving as cloud condensation nuclei (CCN), the physicochemical states of aerosols

strongly influence cloud microstructure, thereby affecting 25 the radiative properties of clouds, circulation, and thermodynamics of the atmosphere (e.g., DeMott et al., 2010). During the intermediate period from the dry to monsoon season in 2019, column-precipitable H<sub>2</sub>O gradually increased from about 1 cm in March to over 3 cm in May, as shown in Fig. 9c. 30 Note that these mutually interacting species (i.e.,  $\tau_{aer}$ , NO<sub>2</sub>, and H<sub>2</sub>O) are retrieved simultaneously from the same solar measurements using a ground-fused silica diffuser, which can provide valuable information for further studies. Temporal variation in ozone in Fig. 9d is mostly associated with 35 fractions in the stratosphere, which also gradually increased from March ( $\sim 250$  Dobson units, DU) to May ( $\sim 275$  DU). Although the total column O<sub>3</sub> is not necessarily relevant to the major topic of this study, it is one of the key constraints of the future algorithm for estimating its tropospheric amounts 40 (Jeong et al., 2020). It should also be noted that SMART-s covers the O<sub>3</sub> Chappuis band in the VIS, which can complement its lower spectral resolution than that of the standard Pandora for profile retrieval (e.g., Natraj et al., 2011).

Figure 10 presents the temporal variations in  $\omega_0$  and H<sub>2</sub>O 45 in panel (a) and those of  $n_{\rm f}$  and  $n_c$  in panels (b) and (c), respectively. The red circles and black squares show the retrievals at 330 and 550 nm, respectively, and their error bars indicate the estimated retrieval error ( $\varepsilon_{ret}$ ) based on the optimal estimation method (OEM; e.g., Jeong et al., 2016, 2020; 50 also see Appendix C). Note that  $\varepsilon_{ret}$  is calculated for each retrieval, which is an important merit of the OEM for relevant studies. As shown in Fig. 10a,  $\omega_0$  and H<sub>2</sub>O gradually increased as biomass burning activities decreased approaching the monsoon season (R = 0.65 between  $\omega_0$  and time and 55 0.70 between H<sub>2</sub>O and time). Interestingly, the correlation between the  $\omega_0$  and H<sub>2</sub>O was even higher (R = 0.81 for  $\omega_0$ at 330 nm) than their temporal trends, as shown in this figure. Eck et al. (2013) reported a similar increasing trend of the  $\omega_0$  from long-term AERONET and OMI measurements 60 over southern Africa. They suggested that a trend of decreasing black carbon content (e.g., due to more smoldering rather than flaming combustion) in the aerosol composition during the progress of burning can be a major reason for aerosol growth through aging, coagulation, or hygroscopic swelling, 65 since size-related aerosol parameters such as the Ångström exponent (AE) and volume median radius did not show a meaningful seasonal trend. However, in this study, the retrieved  $n_{\rm f}$  time series in Fig. 10b exhibited a slightly decreasing trend over time with a reliable retrieval accuracy 70  $(\varepsilon_{\rm ret} = 0.031 \pm 0.015)$ ; this is consistent with the effects of the hygroscopic growth of aerosols, for which  $n_{\rm f}$  values decrease (or come closer to the *n* of water) as a particle grows by water vapor uptake (e.g., Michel Flores et al., 2012; Valenzuela et al., 2018). Meanwhile,  $n_c$  retrievals are estimated to be highly uncertain ( $\varepsilon_{ret} = 0.13 \pm 0.04$ ) due to the very limited information on coarse-mode aerosols, as shown by large standard deviations in Fig. 10c; the high values of  $\varepsilon_{ret}$  are attributable to the SMART-s spectral range, which is narrower



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**Figure 10.** Temporal variations in (**a**) total aerosol single scattering albedo ( $\omega_0$ ), the real part of the aerosol refractive index (*n*) of the (**b**) fine mode and (**c**) coarse mode retrieved at Fang, Thailand, in 2019. The red circles and black squares indicate retrievals at 330 and 550 nm, respectively. The vertical bars indicate estimated errors of each retrieval, based on the optimal estimation method. The small blue circles in panel (**a**) represent total column H<sub>2</sub>O retrievals from SMART–s.

than that of the AERONET, and fine-mode-dominated smoke aerosols over this area (see Fig. 5).

To further investigate the effects of H<sub>2</sub>O on biomass burning aerosol properties over the experiment regions, we have  $_{5}$  also examined the relationships of variations in  $n_{\rm f}$ ,  $\omega_0$ , and the weighted mean radius of the fine mode and AE with changes in H<sub>2</sub>O amounts. Figure 11a and b compare the retrieved  $n_{\rm f}$  and  $\omega_0$  to the total column amount of H<sub>2</sub>O during the period. As previously discussed,  $n_{\rm f}$  is negatively corto related with the H<sub>2</sub>O (*R* ranging from -0.57 to -0.61), whereas  $\omega_0$  showed a high positive correlation for both the UV and longer wavelengths (R = 0.74-0.81). Figure 11c and d show the relationship between the weighted mean radius of

- the fine mode and AE (y axis) to the column-precipitable <sup>15</sup> H<sub>2</sub>O (x axis) from SMART–s. Figure 11c shows two types of weighted mean radius, namely the area-weighted mean radius ( $r_a$ ; often called the effective radius) and the volumeweighted mean radius ( $r_v$ ).  $r_a$  has a proportional relationship with light extinction by particles, which is utilized by aerosol
- <sup>20</sup> retrievals, including the SMART–s algorithm, whereas  $r_v$  has a linear relationship with the volume growth of aerosols.  $r_a$ and  $r_v$  of the fine mode are calculated by the following equations, where the radius of 0.01 and 0.7 µm are the lower and upper size limit of the fine-mode particles, respectively.

$${}_{25} r_{a} = \frac{\int_{0.01\,\mu\text{m}}^{0.7\,\mu\text{m}} r^{3} N(r) \,\mathrm{d}r}{\int_{0.01\,\mu\text{m}}^{0.7\,\mu\text{m}} r^{2} N(r) \,\mathrm{d}r} \tag{6}$$

$$r_{\rm v} = \frac{\int_{0.01\,\mu\rm{m}}^{0.7\,\mu\rm{m}} r^4 N(r) \,\mathrm{d}r}{\int_{0.01\,\mu\rm{m}}^{0.7\,\mu\rm{m}} r^3 N(r) \,\mathrm{d}r.}$$
(7)

As shown in Fig. 11c, both  $r_a$  and  $r_v$  showed a positive correlation with H<sub>2</sub>O (R = 0.42 and 0.56, respectively), which is higher for  $r_v$ . The short wavelength range AE is another good indicator of the fine-mode particle size (Reid et al., 1999; Eck <sup>30</sup> et al., 2001), which is known to have a negative correlation with the size. The AE and absorbing AE (AAE) are calculated as follows:

$$AE = -\frac{\ln\left[\tau_{aer}(\lambda_1)/\tau_{aer}(\lambda_2)\right]}{\ln(\lambda_1/\lambda_2)},$$
(8)

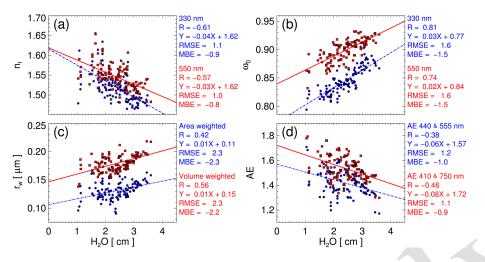
$$AAE = -\frac{\ln[\tau_{abs}(\lambda_1)/\tau_{abs}(\lambda_2)]}{\ln(\lambda_1/\lambda_2)},$$
(9) 35

where  $\tau_{abs}$  is the absorbing aerosol optical thickness (Eq. 10), as follows:

$$\tau_{\rm abs} = (1 - \omega_0) \,\tau_{\rm aer}.\tag{10}$$

The AE calculated using both pairs of wavelengths showed a negative correlation with H<sub>2</sub>O (*R* ranging from -0.46 to  $_{40}$  -0.38).

NO<sub>2</sub> in the smoke plume may contain information about the degree of aging after emission. The lifetime of NO<sub>2</sub> is short (typically less than a few hours), and the brown carbon uptakes NO<sub>2</sub> by photochemical processes; thus, the NO<sub>2</sub> 45 concentration likely decreases through the aging processes of the smoke plumes (e.g., Laskin et al., 2015). Therefore, lower values of  $n_{\rm f}$  (Fig. 12a) and higher values of  $\omega_0$  (Fig. 12b) at lower amounts of NO2 (likely related to the aged plume) support the fact that their temporal variations may also be associ-50 ated with the accelerated aging processes of smoke aerosols by increased H<sub>2</sub>O. Even with the lower correlations, the sizerelated parameters in Fig. 12c and d also indicate that smoke aerosols are likely growing under lower amounts of NO2 during the pre-monsoon period; NO<sub>2</sub> showed a negative correla-55 tion with  $r_v$  and positive correlation with AE. However, note that the correlations of  $n_{\rm f}$ ,  $\omega_0$ ,  $r_{\rm v}$ , and AE with NO<sub>2</sub> also can be attributed to independent trends of emission and photochemical reactions of NO<sub>2</sub>, which are not necessarily associated with the aerosol aging processes. More sophisticated 60 studies combining model simulations and intensive measurements may be able to clarify relationships between NO<sub>2</sub> and aerosol properties. In general, the results in Figs. 10-12 suggest that aerosol aging processes including hygroscopic growth also can be a critical parameter impacting temporal 65 trends of aerosol optical properties over this area, in addition to previously suggested factors over southern Africa (e.g., change in burning sources and conditions in Eck et al., 2013). Overall, such comparisons suggest the potential benefit of simultaneous measures of trace gases and aerosols for under-70 standing atmospheric physicochemical processes.



**Figure 11.** Relationship between aerosol properties and column-precipitable H<sub>2</sub>O retrieved from SMART-s over Fang, Thailand, from 19 March to 2 May 2019. Panels (**a**) and (**b**) compare the real part of the refractive index of the fine mode ( $n_f$ ) and aerosol single scattering albedo ( $\omega_0$ ) to H<sub>2</sub>O, respectively. For the upper panels, blue and red represent  $n_f$  and  $\omega_0$  at 330 and 550 nm, respectively. Panels (**c**) and (**d**) compare the weighted mean radius ( $r_w$ ) and Ångström exponent (AE) to the column precipitable H<sub>2</sub>O, respectively. The blue and red regression in panel (**c**) represent the area- and volume-weighted mean radius, and those in panel (**d**) indicate different wavelength pairs for the AE calculations (blue at 440 and 555 nm and red at 410 and 750 nm).

## 3.3 UV-VIS-NIR aerosol absorption properties for satellite algorithms

The continuous UV-NIR aerosol property information retrieved from SMART-s can be useful not only for validat-5 ing the satellite aerosol products but also for fine-tuning appropriate aerosol models used in the satellite aerosol retrieval algorithm. The operational aerosol algorithm of OMI (OMAERUV; Torres et al., 2013) utilizes radiances at 354 and 388 nm for retrieving the  $\tau_{aer}$  and  $\omega_0$ . They assumed a 10 ratio of k at 354 to 388 nm to be 1.2 for smoke aerosols to account for the spectral absorption effects of organic carbon (Jethva and Torres, 2011; Jeong et al., 2016). They derived a monthly climatology of the aerosol layer height (ALH) from observations by CALIOP (Cloud-Aerosol Lidar with Orthog-15 onal Polarization) as the ancillary data for the OMAERUV algorithm (Torres et al., 2013). Also, recently, the improved ASHE (Aerosol Single-scattering albedo and Height Estimation; Lee et al., 2021) algorithm combines UV measurements from the OMPS-NM (Ozone Mapping and Profiler 20 Suite Nadir Mapper) aboard the SNPP with the VIIRS radiances to provide retrieved  $\omega_0$  and ALH products as part of the VIIRS version 2 DB aerosol CDRs. For the spectral de-

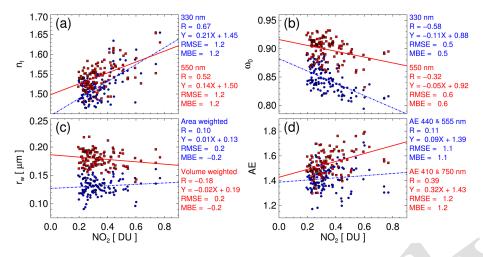
- pendencies of  $\omega_0$  in the UV, they assumed AAE to be 2.0 between 340 and 412 nm. Figure 13 shows an example of  $\omega_0$  retrievals at 340 nm (Fig. 13a), 378 nm (Fig. 13b) and 550 nm (Fig. 13c) from the ASHE algorithm on the same day as Fig. 1 (30 March 2019). Since the ASHE algorithm only performs retrievals when  $\tau_{aer} > 0.5$  and UVAI (UV aerosol index) > 0.7, the spatial coverage of  $\omega_0$  in Fig. 13 is reduced <sup>30</sup> compared to that of  $\tau_{aer}$  shown in Fig. 1. The colored cir-
- cles in this figure depict the collocated  $\omega_0$  retrievals from

SMART–s. In general, the values of  $\omega_0$  retrieved from ASHE nearby the measurement site were comparable to those from SMART–s, with its broader spatial coverage throughout the smoke aerosols of high  $\tau_{aer}$  (see Fig. 1b).

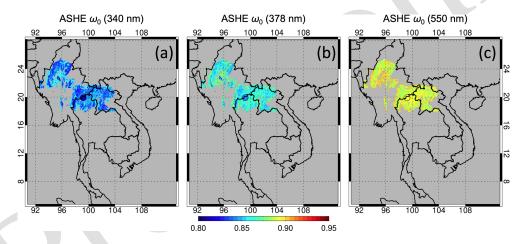
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Temporal variations in  $\omega_0$  from ASHE and SMART-s over the measurement site at overlapping wavelengths demonstrate their reasonable consistency, where both the retrievals indicate increasing trends during the measurement period, with a higher temporal resolution of SMART-s 40 (see Fig. 14a). The ASHE retrievals are available until 18 April 2019, since UVAI over the site decreased lower than the ASHE criteria, likely due to the increased  $\omega_0$  over the period. Figure 14b compares the collocated  $\omega_0$  from ASHE and SMART-s during the measurement period, which 45 showed good agreement between these two at ASHE's retrieval wavelengths. Higher correlations between the ASHE and SMART-s were found in the shorter wavelengths (R =0.79 at 340 nm) due to the higher sensitivity of UV radiances to aerosol absorptions. The MBE and RMSE were lower than 50 0.02 for all wavelengths. Note that the  $\theta_{\rm S}$  at Fang near the overpass time of VIIRS is small, whereas SMART-s measures almucantar radiances  $\theta_{\rm S}$  from 40 to 75°. The average time difference between the ASHE overpass time and the closest SMART-s retrieval during the period was about 3 h, 55 and only samples with a time difference of fewer than 3 h are shown in this figure. However, the difference in time between ASHE and SMART-s still may affect the comparison in Fig. 14, which can be improved by refining the scan strategy (e.g., hybrid scan of AERONET version 3.0) for extending retrieval criteria of the  $\theta_{\rm S}$ .

AERONET sites in regions affected by biomass burning smoke typically had AAE (440-870 nm) ranging from 1.0



**Figure 12.** Comparison of the total column amount of NO<sub>2</sub> to the (**a**) real part of the refractive index of the fine mode ( $n_f$ ) and (**b**) aerosol single scattering albedo ( $\omega_0$ ). The blue and red colors symbolize those retrieved at 330 and 550 nm, respectively. Panels (**c**) and (**d**) compare the weighted mean radius ( $r_w$ ) and Ångström exponent (AE) to the total column NO<sub>2</sub>, respectively. The blue and red symbols in panel (**c**) represent the area- and volume-weighted mean radius, and those in panel (**d**) indicate different wavelength pairs for the AE calculations (blue at 440 and 555 nm and red at 410 and 750 nm).

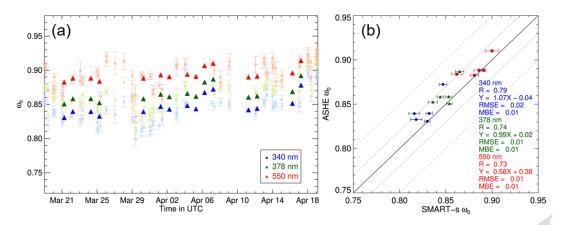


**Figure 13.** Aerosol single scattering albedo ( $\omega_0$ ) retrieved at (**a**) 340 nm, (**b**) 378 nm, and (**c**) 550 nm from ASHE (Aerosol Single-scattering albedo and Height Estimation; Lee et al., 2021) algorithm on 30 March 2019. Colored circles present values of collocated SMART–s retrievals which are indicated by the color bar at each wavelength.

to 2.0, with values closer to 1.0 indicating a greater contribution of black carbon and AAE near 2.0 indicating increased organic-to-black carbon ratios (Giles et al., 2012). Figure 15a presents the temporal variations in the AAE us-

- <sup>5</sup> ing different wavelength pairs calculated from SMART–s  $\tau_{aer}$  and  $\omega_0$ . The AAEs calculated from the UV wavelength and 550 nm pairs were comparable to the assumed value of the ASHE algorithm (2.04±0.27 for the 340–550 nm pair and 1.94±0.33 for the 378 and 550 nm pairs), with non-
- <sup>10</sup> negligible temporal variabilities ranging from about 1.3 to about 2.6. The AAE calculated using the 340–412 nm wavelengths pair, which is the actual pair for the ASHE inversion, showed much higher values ( $2.69 \pm 0.35$ ). Such discrepancies between the SMART–s retrievals and assumptions

of aerosol properties in the ASHE algorithm may propagate <sup>15</sup> to the differences in Figs. 13 and 14. However, the retrieval errors are contextual, meaning that other error sources (uncertainties in the retrieved  $\tau_{aer}$ , assumed size distribution in the aerosol optical models, etc.) can also contribute to the discrepancies, which makes it difficult to quantify the contribution of the AAE uncertainties to the retrieval errors. A longer-term data record in the UV is therefore highly desired. Figure 15b shows temporal variation in the ratio of  $k_f$  at 354 to 388 nm (1.17 ± 0.05), which is in good agreement with the assumed value of the OMAERUV algorithm (i.e., 1.2 for smoke aerosols) but with significant temporal variabilities. Both panels in Fig. 15 suggest that current assumptions of the UV aerosol properties in OMAERUV and ASHE algorithms



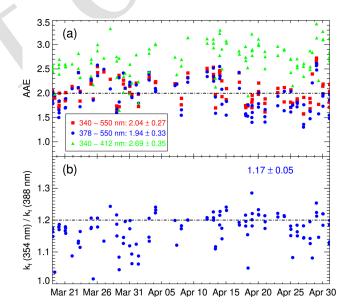
**Figure 14. (a)** Temporal variations in aerosol spectral single scattering albedo ( $\omega_0$ ; 340 nm in blue, 378 nm in green, and 550 nm in red) retrieved from ASHE (dark colors; Lee et al., 2021) and SMART-s (faint colors) over Fang, Thailand, in 2019. ASHE retrievals are not available after 18 April 2019, since UVAI over the site decreased lower than the ASHE criteria, likely due to the increased  $\omega_0$  over the period. Panel (**b**) compares retrieved  $\omega_0$  from ASHE and SMART-s at each wavelength during the period. The dotted and dashed lines represent, respectively, relative biases of  $\pm 0.02$  and  $\pm 0.05$  from the SMART-s product. The collocated samples are limited to have a time difference of less than 3 h. The vertical and horizontal bars in panels (**a**) and (**b**), respectively, indicate estimated errors of each retrieval based on the optimal estimation method.

are generally good approximations. However, temporal and spectral variabilities of the aerosol optical properties, which are presented throughout this section, also emphasize the importance of realistic aerosol models in the UV for further im-

- <sup>5</sup> provements of satellite algorithms. High spectral resolution of the aerosol optical properties covering the UV can also benefit recently launched or upcoming hyperspectral satellite sensors targeting atmospheric composition (e.g., Chance et al., 2019; Ingmann et al., 2012; Kim et al., 2020). The Na-
- <sup>10</sup> tional Institute of Environmental Research of South Korea recently started to deploy standard Pandora devices and a few SMART-s units over Asia to validate GEMS (Geostationary Environment Monitoring Spectrometer; Kim et al., 2020) aerosol and trace gas products and to improve the satellite
- <sup>15</sup> algorithms. Thus, deployments of networked SMART–s can contribute to comparing/validating spatiotemporal variations in aerosol  $\omega_0$ , which is a key parameter for understanding their aging processes and interaction with other environmental conditions (e.g., terrain and meteorology; see Fig. 13 for <sup>20</sup> an example of the spatial variability). These ground-based
- measurements will provide important long-term records of UV aerosol properties at multiple strategic sites over Asia.

### 4 Summary and conclusions

SMART-s was deployed during the pre-monsoon season in <sup>25</sup> northern Thailand in 2019 to perform direct sun and sky radiance measurements near biomass burning sources. In this study, we summarized the detailed radiometric calibration procedures and results. To optimize solar irradiance for the radiometric conversion, we combined the Langley <sup>20</sup> and NIST-traceable integrating sphere calibration data with the high-resolution reference spectrum from Coddington et



**Figure 15.** Temporal variations in the (a) AAE of aerosols and (b) ratio of fine-mode imaginary refractive indices  $(k_f)$  of aerosols at 354 and 388 nm retrieved at Fang, Thailand, in 2019. The color-coded symbols in panel (a) represent different wavelength pairs, with red for 340–550 nm, blue for 378–550 nm, and green for 340–412 nm. The dash-dotted lines in panels (a) and (b) are the assumed values of the ASHE and OMAERUV algorithms for smoke aerosols, respectively.

al. (2021). We estimate that the total uncertainties in the sky radiance measurements are about 5% in the UV and better than 3% in the VIS-NIR wavelengths. The total column amount of  $\tau_{aer}$  and H<sub>2</sub>O from SMART–s showed excellent  $\tau_{aer}$  agreements with those from collocated AERONET measure-

- ments (R = 1.0 and 0.98, respectively). Total column O<sub>3</sub> retrievals from the OMI showed good consistency with those from SMART-s (R = 0.95; RMSE and MBE of less than 3.6 DU). During the measurement period from mid-March
- <sup>10</sup> to early May in 2019,  $\tau_{aer}$  was mostly large (frequently exceeding 2.0 at 500 nm) and strongly correlated with total column NO<sub>2</sub> (R = 0.74), likely due to the high emissions of biomass burning smoke.  $\omega_0$  from SMART–s and AERONET at overlapping wavelengths (i.e., 440 and 675 nm) showed
- <sup>15</sup> acceptable agreements within uncertainties of these instruments (R = 0.79-0.81, with RMSE and MBE of less than 0.015). SMART-s retrievals showed good agreement of finemode V(r) with those from the AERONET, which is dominated by the smoke aerosols during this period. The spectral
- $_{20}$   $\omega_0$  of smoke aerosols showed an abrupt decrease in the UV, consistent with the understanding of absorption by carbonaceous aerosols.

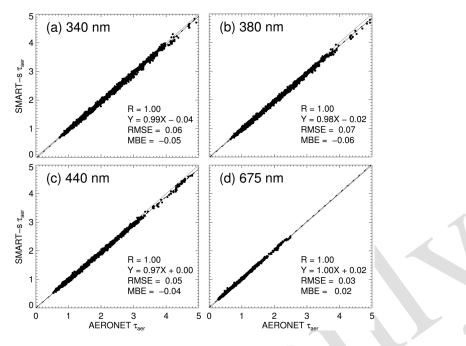
One of the major merits of the simultaneous retrieval of trace gases and aerosols from SMART-s is that it allows an

- <sup>25</sup> informative analysis of physicochemical interactions in the atmosphere. Our analyses comparing the trace gases (i.e., H<sub>2</sub>O and NO<sub>2</sub>) and aerosol properties (e.g.,  $\omega_0$ , *n*, *r*<sub>a</sub>, and *r*<sub>v</sub>) suggest that aerosol aging processes, including hygroscopic growth, can be a critical factor affecting temporal
- <sup>30</sup> trends of aerosol optical properties during the pre-monsoon period over northern Thailand. First, the  $\omega_0$  and column precipitable H<sub>2</sub>O gradually increased together as it approached the monsoon season, and the correlation between  $\omega_0$  and H<sub>2</sub>O was generally higher (R = 0.74-0.81) than their tempo-
- <sup>35</sup> ral trends (R = 0.65 for  $\omega_0$  and 0.70 for H<sub>2</sub>O). Second, the area-/volume-weighted radius of the fine mode also showed a positive correlation with the H<sub>2</sub>O (R = 0.42 and 0.56, respectively). The third result supporting the conclusion is that  $n_f$  from SMART-s showed a negative correlation with the
- <sup>40</sup> total column H<sub>2</sub>O (R = -0.61 for 330 nm and -0.57 for 550 nm), which is supposed to decrease (or come close to the *n* of water about 1.33) as the particle uptake water vapor. However, in this study, measurements are performed at only one location during a limited period, while characteristics of
- <sup>45</sup> smoke aerosols can vary significantly by region and time due to different fuel types, combustion efficiency, and aging processes. A longer period of measurements from multiple sites may help to clarify/provide understanding of such relationships.

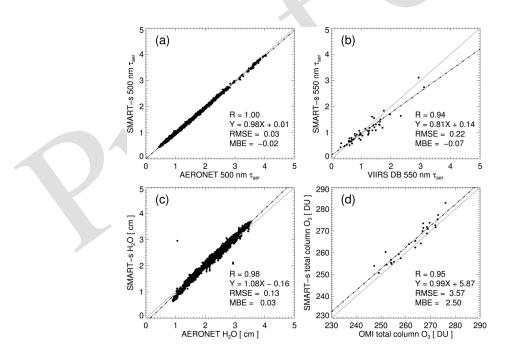
The UV radiances are useful for satellite algorithms to re- 50 trieve higher-order aerosol parameters (e.g., single scattering albedo and aerosol layer height), as these are sensitive to aerosol absorption and the vertical profile. However, due to the lack of information and reliable aerosol models covering the UV, the algorithms typically assumed the spectral 55 optical properties of aerosols or adopted them from laboratory databases (e.g., Jethva and Torres, 2011; Lee et al., 2021). The retrieved UV aerosol properties from SMART-s showed generally good agreements with the current assumptions of the ASHE and OMAERUV algorithms; thereby, 60 there is a reasonable consistency in the  $\omega_0$  from SMART-s and ASHE retrievals (R = 0.73-0.79 and MBE and RMSE of less than 0.02). However, temporal and spectral variabilities of aerosol absorption properties (e.g.,  $\omega_0$  and k) in the UV emphasize the importance of a realistic aerosol model for fur-65 ther improvements of satellite retrievals. Recently launched (e.g., GEMS) and upcoming hyperspectral sensors on the geostationary orbit (e.g., Tropospheric Emissions: Monitoring of Pollution; Chance et al., 2019; Sentinel-4; Ingmann et al., 2012) aim to derive diurnal variations in trace gases 70 and aerosols. SMART-s can provide key parameters of trace gases and aerosols for constraining and validating satellite algorithms from its higher temporal resolution retrievals. Further improvements of algorithms and instruments for acquiring better information content and more reliable products 75 (e.g., adding spectral polarization measurements and refinement of scanning strategy, such as adding a hybrid scan of the AERONET) are currently underway.

# Appendix A: Scatterplots between SMART-s direct-sun retrievals and AERONET/satellite observations

Direct-sun retrievals from SMART–s are analyzed and compared with other sources of retrievals in Sect. 3. However, this section additionally shows scatterplots of  $\tau_{aer}$  and total column H<sub>2</sub>O and O<sub>3</sub> from the SMART–s with those from <sup>85</sup> AERONET and satellite measurements during the campaign period for validation and checking consistency. As reported by Jeong et al. (2018), spectral  $\tau_{aer}$  from SMART–s showed excellent agreements with AERONET at all overlapping wavelengths, as shown in Figs. A1 and A2a, which is also <sup>90</sup> in good agreement with the VIIRS DB product (Fig. A2b). The total precipitable water vapor (H<sub>2</sub>O) and total column ozone from SMART–s also showed excellent agreement with AERONET and OMI, as shown in Fig. A2c and d.



**Figure A1.** Comparison of aerosol optical thickness ( $\tau_{aer}$ ) from SMART-s and AERONET (version 3, level 2.0) at (**a**) 340 nm, (**b**) 380 nm, (**c**) 440 nm, and (**d**) 675 nm measured at Fang, Thailand, from 8 March to 2 May 2019. *R* is the correlation coefficient, RMSE denotes the root mean square error, and MBE is the mean bias error. The black dot-dashed line and the gray solid line represent regression and one-to-one lines, respectively.

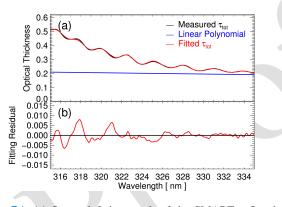


**Figure A2.** Panel (a) compares aerosol optical thickness ( $\tau_{aer}$ ) at 500 nm from SMART-s and AERONET (version 3, level 2.0), and panel (b) compares that at 550 nm from VIIRS DB and SMART-s. Total precipitable water vapor (H<sub>2</sub>O) products from SMART-s and AERONET are compared in panel (c), and total column ozone retrievals from OMI and SMART-s are compared in panel (d). The retrievals are obtained at Fang, Thailand, from 8 March to 2 May 2019. *R* is the correlation coefficient, RMSE denotes root mean square error, and MBE is the mean bias error. The black dot-dashed line and the gray solid line represent regression and one-to-one line, respectively.

### Appendix B: SMART-s ozone and water vapor retrieval algorithm

The spectral Langley method (Jeong et al., 2018) retrieves spectral  $\tau_{aer}$  by subtracting gas optical thickness ( $\tau_{gas}$ ) from  $\tau_{gas}$  the total optical thickness ( $\tau_{tot}$ ); the SMART-s observation is beneficial for this procedure as it measures spectral features of the  $\tau_{gas}$ . SMART-s retrieves optically thick trace gases (i.e., O<sub>3</sub> and H<sub>2</sub>O) by using a similar method but with narrower fitting windows of the trace absorption bands

- <sup>10</sup> (i.e., 315-335 nm for O<sub>3</sub> and 550-680 nm for H<sub>2</sub>O). Examples of the fitting results for the O<sub>3</sub> and H<sub>2</sub>O are shown in Figs. B1 and B2, respectively, and intercomparison results during the campaign, together with other data, are shown in Appendix A. The fitting model for O<sub>3</sub> retrieval includes lin-
- <sup>15</sup> ear polynomials,  $O_3$  and  $SO_2$  cross-sections, and the Raman spectrum, whereas that for the H<sub>2</sub>O utilizes linear polynomials and H<sub>2</sub>O, O<sub>3</sub>, and O<sub>2</sub> cross sections. More detailed design and calibration procedures for the spectral Langley method are demonstrated in Jeong et al. (2018). Measure-
- <sup>20</sup> ments for long-term validation/comparison results are currently underway and will be reported in a follow-up paper (Jeong et al., 2022).

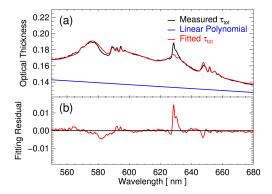


**Figure B1. (a)** Spectral fitting result of the SMART-s  $O_3$  algorithm at 16:50 UTC on 14 October 2019 measured on the rooftop of the NASA Goddard Space Flight Center Building 33 (38.99° N latitude, 76.84° W longitude). Panel (**b**) presents the fitting residual (measured optical thickness subtracted by model optical thickness).

### Appendix C: SMART-s aerosol inversion algorithm

Spectral bands for the aerosol property retrieval are care-<sup>25</sup> fully selected to avoid strong absorption by the gases and to efficiently obtain maximum information on aerosols. For example, we avoided major O<sub>2</sub>-A, O<sub>2</sub>-B, O<sub>2</sub>-O<sub>2</sub>, and H<sub>2</sub>O bands where each corresponding  $F_{\text{Comb}}$  is highly uncertain (see Fig. 4a and b). Nodes of the relative azimuth angle ( $\phi_r$ )

<sup>30</sup> for the solar almucantar scan and those of wavelengths for the retrieval are summarized in Table C1 with other parameters. The VLIDORT (linearized pseudo-spherical vector discrete ordinate radiative transfer) code generates the full Stokes'



**Figure B2.** Similar to Fig. B1 but for  $H_2O$ . (a) Spectral fitting result of the SMART–s  $H_2O$  algorithm at 16:55 UTC on 14 October 2019 measured at the same location as in Fig. B1. Panel (b) presents the fitting residual (measured optical thickness subtracted by model optical thickness).

parameters and analytic weighting function of atmospheric and surface variables, including the aerosol properties (Spurr,  $_{35}$ 2006; Spurr et al., 2012, Spurr and Christi, 2014). The state vector (x, which is a vector with elements of retrieval parameters and control variables to fit the measurements using a forward model) consists of fine- and coarse-mode n and k, five parameters of the PSD, two aerosol layer height parameters, and spectral surface albedo. The algorithm assumes the aerosol number PSD as a bimodal lognormal shape, as follows:

$$N(r) = \frac{F_{\text{num}}}{\sqrt{2\pi} \ln \sigma_{\text{f}}} \frac{1}{r_{\text{f}}} \exp\left[-\frac{1}{2} \left(\frac{\ln r - \ln r_{\text{f}}}{\ln \sigma_{\text{f}}}\right)^2\right] + \frac{(1 - F_{\text{num}})}{\sqrt{2\pi} \ln \sigma_{\text{c}}} \frac{1}{r_{\text{c}}} \exp\left[-\frac{1}{2} \left(\frac{\ln r - \ln r_{\text{c}}}{\ln \sigma_{\text{c}}}\right)^2\right], \quad (C1)$$

where N(r) is the number size distribution, and the  $r_{\rm f}$  and  $_{45}$  $r_{\rm c}$  are fine- and coarse-mode mean radius.  $\sigma_{\rm f}$  and  $\sigma_{\rm c}$  are the geometric standard deviation of each mode.  $F_{\rm num}$  is number fraction of the fine mode. However, the assumed bimodal lognormal shape may not adequately represent the true fineor coarse-mode distribution in some cases, such as volcanic eruptions, aerosol aging, and cloud processing (Eck et al., 2010, 2013, 2018). The aerosol extinction profile is assumed to be a Gaussian shape, as follows:

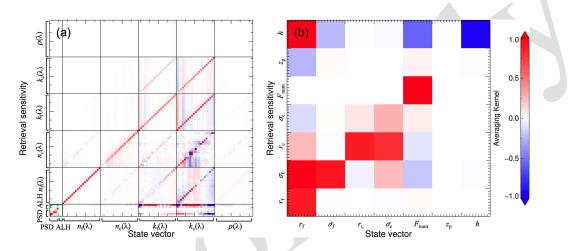
ALH(z) = 
$$W \frac{e^{-h(z-z_p)}}{\left[1 + e^{-h(z-z_p)}\right]^2}$$
, (C2)

where ALH(*z*) stands for the aerosol layer height (i.e., vertical profile of aerosol extinction), and the *W* is the normalization factor.  $z_p$  is the peak height, and *h* is the vertical dispersion parameter of the Gaussian profile shape. A priori information of  $z_p$  and *h* are extracted from the climatology of reanalysis data (e.g., Modern-Era Retrospective analysis for Research and Applications, Version 2; Gelaro et al.,

Table C1. Relative azimuth angle	es, wavelength node	, sources of	measurement	error covariance	matrix, and	d parameters of	state vector of
SMART-s algorithm.							

Algorithm parameter	Description	
Relative azimuth angles	3.0, 3.5, 4.0, 5.0, 6.0, 7.0, 8.0, 10.0, 12.0, 14.0, 16.0, 18.0, 20.0, 25.0, 30.0, 35.0, 40.0, 45.0, 50.0, 60.0, 70.0, 80.0, 90.0, 100.0, 120.0, 140.0, 160.0, and $180.0^{\circ}$	
Wavelength node	330, 340, 350, 360, 370, 380, 390, 400, 410, 440, 455, 490, 520, 540, 555, 580, 610, 640, 675, 750, and 778 nm	
Sources of measurement error covariance matrix	r Estimated from Langley and laboratory calibration results	
State vector	Lognormal parameters of aerosol number-size distribution for fine and coarse mode, number	

fine-mode fraction, two parameters of Gaussian vertical profile shape of aerosols, spectral complex refractive indices of fine and coarse mode, and spectral surface reflectance



**Figure C1.** Averaging kernel matrix (**A**) of SMART–s retrievals at 08:31 UTC on 10 April 2019 at Fang, Thailand, when fine-/coarse-mode volume fractions were comparable and aerosol single scattering albedo was about 0.87, with the aerosol optical thickness of about 1.06 at 440 nm. Indices from 1 to 5 correspond to fine-mode mean radius ( $r_f$ ) and geometric standard deviation ( $\sigma_f$ ), and those pair for the coarse mode ( $r_f$  and  $\sigma_f$ ) and number fine-mode fraction ( $F_{num}$ ). Indices 6 and 7 are the peak height ( $z_p$ ) and dispersion parameter (h) of the assumed aerosol extinction profile. Each element of the real (n) and imaginary (k) part of the refractive index for fine (subscript f) and coarse mode (subscript c) indicates its retrieval sensitivity at each wavelength. Panel (**a**) shows the whole of **A**, and panel (**b**) magnifies to the particle size distribution (PSD) and aerosol layer height (ALH) parameters (indices from 1 to 7), as indicated by a green square at the bottom left of panel (**a**).

2017), which is, however, not sensitive to the solar almucantar measurements. In this study, we assumed that the surface reflectance ( $\rho$ ) is Lambertian, of which a priori data are obtained from merged satellite measurements. More details of 5 the parameters and design of the algorithm can be found at

Jeong et al. (2020).

The averaging kernel (**A**) of the OEM is a useful matrix for understanding the information content of a set of measurements and the inversion method, of which elements show the <sup>10</sup> sensitivity of retrievals to the true state (e.g., Rodgers, 2000; Jeong et al., 2020). **A** is defined as follows:

$$\mathbf{A} = \mathbf{G}\mathbf{K} = \frac{\partial \hat{x}}{\partial x},\tag{C3}$$

where G is the gain matrix for representing the sensitivity of retrievals to the measurements, and K is the weighting

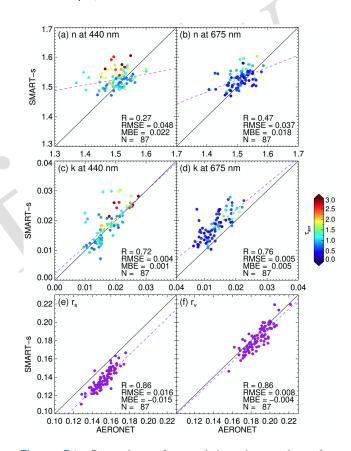
function matrix of which the elements are partial derivatives of each measurement with respect to the state vector. In Eq. (C3), each element of A characterizes how the retrieval  $(\hat{x})$  responses to the true state (x). Diagonal elements of A  $(D_{\rm A})$  indicate the sensitivity of each retrieval parameter, using a set of measurements and an inversion method, whereas 20 the off-diagonal elements of an *i*th row  $(R_A)$  demonstrate retrieval errors of  $x_i$  by cross-correlation with other parameters or by insufficient information contained in the measurements. Therefore, for an ideal inversion with an observing system, its A is close to an identity matrix (Rodgers, 1990). 25 More discussion about A for SMART-s is summarized in Jeong et al. (2020). Figure C1 shows an example of A from the SMART-s retrieval at Fang on 10 April 2019, when fineand coarse-mode volume fractions were comparable, and  $\omega_0$ 

was about 0.87, with  $\tau_{aer}$  about 1.06 at 440 nm. Figure C1a presents the whole **A**, and Fig. C1b magnifies to the PSD (indices from 1 to 5) and ALH parameters (indices of 6 and 7), as indicated with a green square in Fig. C1a. The  $D_A$  of PSD

- <sup>5</sup> parameters (diagonal elements of 1–5; order of  $r_{\rm f}$ ,  $\sigma_{\rm f}$ ,  $r_{\rm c}$ ,  $\sigma_{\rm c}$ , and  $F_{\rm num}$ ) are close to one, which shows their sufficient retrieval sensitivity from the measurement. Particularly,  $R_{\rm A}$  of  $r_{\rm f}$  and  $F_{\rm num}$  (see the off-diagonal elements of the first and fifth rows in Fig. C1b) have small absolute values, whereas
- <sup>10</sup> those of  $\sigma_{\rm f}$ ,  $r_{\rm c}$  and  $\sigma_{\rm c}$  (rows from 2 to 4) are relatively high. As solar almucantar measurements are not sensitive to the vertical profile of aerosols,  $D_{\rm A}$  of the ALH parameters ( $z_{\rm p}$  and *h* of indices 6 and 7) are low, with their relatively higher values of the  $R_{\rm A}$  for  $k_{\rm f}$  and  $k_{\rm c}$ .  $D_{\rm A}$  of the  $n_{\rm f}$  and  $n_{\rm c}$  show their
- <sup>15</sup> sufficient retrieval sensitivity at all wavelengths. In addition, the  $R_A$  of  $n_f$  for  $n_c$  and that of  $n_c$  for  $n_f$  were negligible, which suggests that their retrieval sensitivity for each mode is independent and can be retrieved separately. However, the  $R_A$ s of  $n_f$  and  $n_c$  demonstrate that they are also affected by
- <sup>20</sup>  $k_{\rm f}$  and  $k_{\rm c}$ . On the contrary,  $R_{\rm A}$  of  $k_{\rm f}$  and  $k_{\rm c}$  are low at most of the other parameters (i.e., n, PSD, ALH, and  $\rho$ ), whereas the  $R_{\rm A}$  of  $k_{\rm f}$  at state vector  $k_{\rm c}$  (or vice versa) shows high values of diagonal elements, indicating retrieval sensitivity across the fine and coarse mode; therefore, we analyzed  $\omega_0$  for both
- <sup>25</sup> modes in this study. It is well recognized that the retrieval sensitivity of  $\rho$  is negligible, which results in low values of the whole rows of  $\rho$  in Fig. C1. We also limit the retrieval range of the solar zenith angle ( $\theta_S$ ) up to 75° to minimize the effects of surface reflectance at a high  $\theta_S$  and zenith view-
- <sup>30</sup> ing angle ( $\theta_v$ ). Note that **A** varies significantly, depending on the case, and retrieval errors due to the interferences between different parameters and/or the lack of information are considered by the error estimation method of the OEM (Rodgers, 2000; Jeong et al., 2020).
- One of the important merits of the OEM is its theoretical formulations of retrieval errors, which are classified into four categories (Rodgers, 1990, 2000; Jeong et al., 2016, 2020), i.e., the smoothing error ( $\varepsilon_s$ ), retrieval noise ( $\varepsilon_m$ ), forward model error ( $\varepsilon_r$ ), and model parameter error ( $\varepsilon_f$ ). As *x* con-
- <sup>40</sup> tains the most dominant parameters of aerosols and quantifying uncertainties in the radiative transfer model is challenging, we neglect ( $\varepsilon_r$ ) and ( $\varepsilon_f$ ). We defined the retrieval error ( $\varepsilon_{ret}$ ) as the square root of the sum of squared  $\varepsilon_s$  and  $\varepsilon_m$ , which represents the minimum uncertainty of the SMART-s
- <sup>45</sup> aerosol inversions. Detailed descriptions and discussions of the  $\varepsilon_{ret}$  are summarized by Jeong et al. (2020).

### Appendix D: Comparison of aerosol optical properties from the SMART-s and AERONET

Coincident retrievals of  $\omega_0$ , *n*, *k*, and *V*(*r*) are compared <sup>50</sup> and discussed in Figs. 5 to 10. The remaining scatterplots of these inversion parameters of SMART–s and AERONET are shown in this section. Figure D1 compares the aerosol inversion products from SMART-s and AERONET (version 3, level 2.0). The upper panels compare the real part of the refractive index at (Fig. D1a) 440 nm and e (Fig. D1b) 675 nm, and the middle panels are for the imaginary part. The lower panels compare (Fig. D1e) area- and (Fig. D1f) volume-weighted mean radii. Larger discrepancies between SMART-s and AERONET were found for *n* due to the lack of information content and different assumptions of the algorithms, whereas a higher consistency appeared for *k*, as discussed for  $\omega_0$ . High agreements of  $r_s$  and  $r_v$  were found between the two instruments, as shown in Fig. D1e and f (R = 0.86, with MBE and RMSE of less than 0.016 µm).



**Figure D1.** Comparison of aerosol inversion products from SMART-s and AERONET (version 3, level 2.0) from 19 March to 2 May 2019. Upper panels compare the real part of the refractive index at (a) 440 nm and (b) 675 nm, and middle panels compare the imaginary part at these wavelengths. Colored circles in panels (a)–(d) represent the values of aerosol optical thickness ( $\tau_{aer}$ ) from the color bar at each wavelength. Lower panels compare (e) area- and (f) volume-weighted mean radii ( $r_s$  and  $r_v$ , respectively). *R* is the correlation coefficient, RMSE denotes the root mean square error, MBE is the mean bias error, and *N* is the number of samples for the comparison. The black solid line and purple dashed line depict one-to-one and regressions, respectively.

**Data availability.** SMART-s data are also described at https://earth.gsfc.nasa.gov/climate/instruments/smartlabs (last access: 12 September 2022) and are available on request from the authors (Ukkyo Jeong and Si-Chee Tsay). AERONET data are s available at https://aeronet.gsfc.nasa.gov (Giles et al., 2019; Sinyuk et al., 2020).

Author contributions. The first and corresponding author (UJ) led the overall algorithm development and instrument calibration and wrote the paper. SCT supervised the overall activities for this <sup>10</sup> study and was also in charge of the campaign for this study. NCH and JL provided the satellite retrievals to compare the aerosol optical properties analyzed in this study. DMG and BNH provided critical revision for the aerosol retrievals and the overall paper. JWC and JJB provided the NIST-traceable light source for the radiometric <sup>15</sup> calibration and also reviewed the instrument calibrations. RJS was

- in charge of the Pandora network operation for the campaign and also supported instrument maintenance. SHW and SC supported the local operation of the instruments throughout the campaign periods and also reviewed the final paper. HH, DK, and JK are in charge of
- <sup>20</sup> the Asian Network of the Pandora and the Geostationary Environment Monitoring Spectrometer, which provided critical insight for this study. They also provided critical revisions for the paper.

**Competing interests.** The contact author has declared that none of the authors has any competing interests.

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