

# Albedo susceptibility of Northeastern Pacific stratocumulus: the role of covarying meteorological conditions

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**Abstract.** Quantification of the radiative adjustment of marine low-clouds to aerosol perturbations, regionally and globally, remains the largest source of uncertainty in assessing current and future climate. One of the important steps towards quantifying the role of aerosol in modifying cloud radiative properties is to quantify the susceptibility of cloud albedo and liquid water path (LWP) to perturbations in cloud droplet number concentration ( $N_d$ ). We use 10 years of space-borne observations from the polar-orbiting Aqua satellite to quantify the albedo susceptibility of marine low-clouds to  $N_d$  perturbations over the northeast (NE) Pacific stratocumulus (Sc) region. Mutual information analysis reveals a dominating control of cloud state (e.g. LWP and  $N_d$ ) on low-cloud albedo susceptibility, relative to the meteorological states that drive these cloud states. Through a LWP- $N_d$  space decomposition of albedo susceptibilities, we show clear separation among susceptibility regimes (brightening or darkening), consistent with previously established mechanisms through which aerosol modulates cloud properties. These regimes include (i) thin non-precipitating clouds ( $LWP < 55 \text{ g m}^{-2}$ ) that exhibit brightening (occurring 37% of the time), corresponding to the Twomey effect; (ii) thicker non-precipitating clouds, corresponding to entrainment driven negative LWP adjustments that manifest as a darkening regime (36% of the time); and (iii) another brightening regime (22% of the time) consisting of mostly precipitating clouds, corresponding to precipitation-suppression LWP positive adjustments. Overall, we find an annual-mean regional low-cloud brightening potential of  $20.8 \pm 2.68 \text{ W m}^{-2} \ln(N_d)^{-1}$ , despite an overall negative LWP adjustment for non-precipitating marine stratocumulus, owing to the high occurrence of the Twomey brightening regime. Over the NE Pacific, clear seasonal covariabilities among meteorological factors related to the large-scale circulation are found to play an important role in grouping conditions favorable for each susceptibility regime. When considering the covarying meteorological conditions, our results indicate that for the Northeastern Pacific stratocumulus, clouds that exhibit the strongest brightening potential occur most frequently within shallow marine boundary layers over a cool ocean surface with a stable atmosphere and a dry free-troposphere above. Clouds that exhibit a darkening potential associated with negative LWP adjustments occur most frequently within deep marine boundary layers in which the atmospheric instability and the ocean surface are not strong and warm enough to produce frequent precipitation. Cloud brightening associated with warm rain suppression is found to preferably occur either under unstable atmospheric conditions or humid free-tropospheric conditions that co-occur with a warm ocean surface.

Changes in aerosol concentrations in the marine boundary layer, of either natural or anthropogenic origin, can lead to significant changes in the brightness of marine low-level clouds. Examples of aerosol induced changes in cloud reflectivity are observed in aerosol perturbations associated with natural causes, such as volcanic eruptions (e.g. Gassó, 2008; Yuan et al., 2011; Malavelle et al., 2017), and anthropogenic sources across the globe, such as ship emissions, wildfires, and power plants (Toll et al., 2019).

30 Among anthropogenic sources, shiptracks – bright linear cloud features associated with particle emissions (Coakley et al., 1987) – have been used to improve our understanding of cloud responses to aerosol perturbations. The routine and frequent occurrence of global shipping traffic, and constant meteorological conditions in- and out-of-shiptrack make them a ‘natural laboratory’ to improve our understanding of cloud responses to aerosol perturbations. Studies based on satellite observations (e.g. Coakley and Walsh, 2002; Gryspeerd et al., 2019b; Chen et al., 2012; Christensen and Stephens, 2011; Christensen et al., 35 2014) and idealized frameworks such as large-eddy simulations (e.g. Wang et al., 2011; Hill et al., 2009), have been used to improve the quantification of the global aerosol radiative effect (e.g. Diamond et al., 2020). However, to date, our ability to narrow down estimates of climate sensitivity is still limited by uncertainties related to quantifying the radiative adjustment of marine low-clouds to the anthropogenic aerosol (Bellouin et al., 2020).

For non-precipitating warm clouds exhibiting constant liquid water path (LWP), increases in aerosol concentration result 40 in increases in droplet concentration ( $N_d$ ) leading to smaller droplets that make the cloud more reflective (the Twomey effect; Twomey, 1974, 1977). These processes occur at short timescales (order 5 – 10 min, supplementary materials in Glassmeier et al. (2021)). However, LWP is not always constant: LWP adjustments were first suggested to exist in precipitating marine warm 45 clouds: an increase in  $N_d$  leads to smaller cloud droplets that are less likely to grow by collision-coalescence to precipitation-sized raindrops under the same environmental conditions (Albrecht, 1989). The result is a reduction in the loss of cloud water due to precipitation, which then leads to an increase in LWP that enhances cloud brightening associated with the smaller drops.

More recently, negative LWP adjustments in non-precipitating stratocumulus have also been identified: (i) the reduced 50 droplet sizes decrease the sedimentation flux at stratiform cloud-tops, enhancing the evaporative and radiative cooling and thereby the entrainment rate at cloud tops (the sedimentation-entrainment feedback; Ackerman et al., 2004; Bretherton et al., 2007); (ii), smaller cloud droplets evaporate faster, leading to stronger cooling and more turbulent mixing at cloud-tops, which then causes more evaporation, creating a positive feedback loop, known as the evaporation-entrainment feedback (Wang et al., 2003; Xue and Feingold, 2006; Jiang et al., 2006). Both these entrainment-feedbacks reduce cloud LWP in response to the 55 increased concentration of smaller droplets, resulting in less reflective clouds and hence a warming relative to a cloud with constant LWP. A strong offsetting warming effect from the negative LWP adjustment is evident in both observational studies (e.g. Chen et al., 2014; Possner et al., 2020; Gryspeerd et al., 2019a, 2021) as well as large eddy simulation (e.g. Wang et al., Ackerman et al., 2004; Xue et al., 2008). The timescale associated with these negative LWP adjustments is  $t \approx 20$  h (Glassmeier et al., 2021). Because shiptracks exist for only 6 to 7 hours, typically, and are likely to be sampled on average after  $\sim 3$  h, a generalization of shiptrack characterized aerosol-cloud interactions to estimates of anthropogenic aerosol climate

forcing may be substantially overestimated because the shiptrack has not existed for long enough to manifest full negative LWP adjustment (Glassmeier et al., 2021).

60 Moreover, despite routine shipping traffic, ship tracks are only rarely observed over major shipping corridors (only 0.002% of the total ocean-going ship traffic; Campmany et al., 2009), in part due to the narrow range of meteorological conditions required for these bright tracks to form (Durkee et al., 2000). This suggests that the coupled large-scale meteorology and the associated cloud states have a strong impact on the susceptibility of low-clouds to aerosol perturbations. Several studies have tried to constrain the uncertainties in LWP and reflectance adjustments based on cloud states and large-scale meteorological 65 conditions using satellite observations (e.g. Chen et al., 2014; Douglas and L'Ecuyer, 2019; Possner et al., 2020) and found strong meteorological controls on cloud state and cloud albedo susceptibility to aerosol perturbations across the globe: regions with relative dry and unstable conditions tend to be characterized by cloud darkening in response to increased aerosol loading, whereas clouds in stable and moist regions tend to brighten in response to increased aerosol concentrations.

In order to understand and disentangle the impact of individual meteorological drivers, subsampling of data is often applied 70 in these studies to help constrain the degree of freedom of the system within one meteorological variable by limiting that within the other meteorological variables. This results in minimizing and suppressing the influence of the covariability among meteorological drivers, even though large-scale meteorological conditions are spatial-temporally correlated, especially over the eastern subtropical oceans (e.g. Klein and Hartmann, 1993; Eastman et al., 2016). Thus, the “untangling” leads to neglect of 75 important information, i.e. the frequency at which certain environmental conditions co-occur in nature, which profoundly drives the overall radiative impacts of aerosol-cloud interactions. A shift in attention from untangling aerosol and meteorological effects on cloud systems towards embracing and understanding the covariabilities between aerosol and meteorological drivers has been suggested by Mülmenstädt and Feingold (2018). It is the approach adopted here.

In this work, we focus on the potential radiative impact of “intrinsic” cloud adjustments (due to changes in  $N_d$  and LWP). “Extrinsic” cloud adjustment (cloud fraction responses) is not addressed here. We quantify relationships between cloud albedo 80 ( $A_c$ ) and  $N_d$  using satellite-retrieved cloud properties and radiative fluxes (Section 2), following the conceptual framework of using  $N_d$  as an intermediate variable to minimize the influence of confounding meteorology on the causal relationship between aerosol and cloud as in Gryspenert et al. (2016, 2019a). Our target area is the northeast Pacific marine stratocumulus deck, one of the regions contributing most strongly to the overall cooling of the Earth by reflecting incoming solar radiation (Klein and Hartmann, 1993). Cloud albedo susceptibilities are approximated by regressed log-linear relationships between  $N_d$  and  $A_c$  85 within a given satellite snapshot, similar to Painemal (2018), assuming processes are related to the current state of the system captured by the satellite snapshot, with no memory of past states (Section 3). One should note that this Markovian approach of inferring process from composites of satellite snapshots ought to be restricted to informing relationships between cloud properties from a climatological perspective where a sufficient amount of sampling of a time-space varying system creates a robust characterization of the relationships between quantities that describe the system. This contrasts with approaches 90 targeted at the non-Markovian aspect of the system, i.e. quantifying the time derivatives of cloud properties, through tracking properties of the system, either by numerical simulation or temporally-resolved satellite observations (e.g. Glassmeier et al., 2021; Christensen et al., 2020).

The findings of this study (Section 4 and 5) feature two key perspectives: (i) the usage of the LWP- $N_d$  parameter space, supported by mutual information analyses (Section 4.1), helps to show clear separation between albedo susceptibility regimes  
95 that can be linked to physical mechanisms associated with aerosol effects on low clouds (Section 4.2 and 4.3); (ii) distinguished from previous work that minimized the covariability between meteorological drivers (e.g. Douglas and L'Ecuyer, 2019), this study adopts a top-down approach that embraces the covariability among meteorological factors (obtained from ERA5 reanalyses) while identifying conditions under which clouds are more (or less) susceptible to aerosol perturbations and quantifying the frequency of occurrence of these conditions (Section 5).

100 **2 Datasets**

This study focuses on an area of  $10^\circ$  by  $10^\circ$  ( $120\text{--}130^\circ\text{W}$ ,  $20\text{--}30^\circ\text{N}$ ) over the subtropical Northeast (NE) Pacific stratocumulus region, corresponding to an area of regional maximum in annual stratus cloud amount, which is the same region examined in Klein and Hartmann (1993). Marine low-cloud properties and shortwave (SW) radiative measurements are retrieved from the MODerate resolution Imaging Spectroradiometer (MODIS) (Platnick et al., 2003) and the Clouds and the Earth's Radiant  
105 Energy Systems (CERES; Wielicki et al., 1996) sensors onboard the Aqua satellite (overpass  $\sim 1:30$  pm local time), obtained from the CERES Single Scanner Footprint (SSF) product Edition 4 (level 2; Su et al., 2015). Top-of-atmosphere (TOA) SW fluxes, including incoming solar radiation ( $\text{SW}_{TOA_{dn}}$ ) and reflected SW flux ( $\text{SW}_{TOA_{up}}$ ), are derived from the Single Scanner at a CERES footprint resolution of 20 km (Loeb et al., 2005; Su et al., 2015), which are then used to calculate cloud SW albedo ( $A_c$ ) as follows:

$$110 \quad A_c = \frac{(A_{all} - A_{clr}(1 - f_c))}{f_c} \quad (1)$$

where  $A_{all}$  is scene albedo (all-sky albedo), defined as the ratio of  $\text{SW}_{TOA_{up}}$  to  $\text{SW}_{TOA_{dn}}$ ,  $A_{clr}$  is the solar zenith angle (SZA) dependent ocean albedo (clear-sky albedo), derived from the scene albedo under clear sky conditions over the study area, and  $f_c$  is the cloud fraction.

MODIS cloud properties, including cloud optical depth ( $\tau$ ), cloud top effective radius ( $r_e$ ),  $f_c$ , LWP, cloud effective temperature, and cloud top height (CTH) are retrieved using the CERES-MODIS algorithm at MODIS pixels and then aggregated to  
115 the CERES footprint resolution (20 km) and scanning pattern (Minnis et al., 2011b, a). Retrieval of  $r_e$  is based on the  $3.7\text{-}\mu\text{m}$  channel, which has been shown to be less affected by retrieval biases than the  $2.1\text{-}\mu\text{m}$  and  $1.6\text{-}\mu\text{m}$  channels (Grosvenor et al., 2018).  $N_d$  is calculated following Grosvenor et al. (2018) as

$$N_d = \frac{\sqrt{5}}{2\pi k} \left( \frac{f_{ad}c_w(T, P)\tau}{Q_{ext}\rho_w r_e^5} \right)^{1/2} \quad (2)$$

120 where  $k$  is a parameter representing the width of the modified gamma droplet distribution (assumed to be 0.8; Martin et al., 1994),  $f_{ad}$  is the adiabatic fraction (assumed to be 0.8),  $c_w$  is the condensation rate, which is a function of temperature (T) and pressure (P) (Grosvenor and Wood, 2014), calculated using CERES-MODIS cloud effective temperature at a constant pressure of 900 hPa,  $Q_{ext}$  is the extinction efficiency factor, approximated by its asymptotic value of 2 (Grosvenor et al., 2018), and

$\rho_w$  is the density of liquid water. In addition,  $N_d$  is only calculated for CERES footprints with  $f_c > 0.99$  (overcast footprints), 125 cloud effective temperature greater than 273 K (to exclude mixed-phase and ice clouds), CTH less than 3 km,  $\tau > 3$ , and  $r_e > 3 \mu m$ , and SZA  $< 65^\circ$ , to minimize retrieval biases (Painemal et al., 2013; Grosvenor and Wood, 2014; Grosvenor et al., 2018). Furthermore, footprints with a calculated  $N_d$  greater than  $600 \text{ cm}^{-3}$  (outside the 99.9<sup>th</sup> percentile) are discarded to avoid highly unrealistic  $N_d$  retrievals. As discussed further in Section 3, the  $f_c > 0.99$  condition at the CERES 20 km footprint 130 allows for lower  $f_c$  when these 20 km pixels are aggregated to  $1^\circ \times 1^\circ$  scenes. For  $1^\circ \times 1^\circ$  satellite sampled scenes,  $\sim 53\%$  consist of single-layer liquid clouds only, and among these cloudy scenes,  $\sim 41\%$  satisfy the  $N_d$  and  $S_0$  calculation criteria (introduced in Section 3) and are subsequently used in this study.

Meteorological conditions, including sea surface temperature (SST), sea level pressure (SLP), vertical velocity at 700 hPa ( $\omega_{700}$ ), and temperature, humidity, and wind profiles, are obtained from the European Centre for Medium-Range Weather Forecasts (ECMWF) fifth-generation atmospheric reanalysis (ERA5; Hersbach et al., 2020), available every hour at  $0.25^\circ$  135 spatial resolution. Lower-tropospheric-stability (LTS) is calculated as the difference in potential temperature between 700 hPa and 1000 hPa. Free-tropospheric relative humidity ( $\text{RH}_{ft}$ ) is defined as the the mean relative humidity between inversion top and 700 hPa, following Eastman and Wood (2018).

### 3 Methods

Despite decades of research addressing the impact of aerosol on cloud radiative effect, the causality problem remains pernicious, 140 in part due to the covarying aerosol and meteorological conditions that make untangling aerosol and meteorological effects extremely hard. In other words, confounding meteorological factors that have influences on both the aerosol and cloud properties (e.g. Mauger and Norris, 2007; Gryspeerdt et al., 2014) often obscure the direct causal relationship between aerosol and cloud properties. As an important step forward, Gryspeerdt et al. (2016, 2019a) show that using  $N_d$  as an intermediary can help reduce the meteorological confounding effect on the causal relationship between aerosol and cloud properties.

145 This work adopts the same logic, that is, it considers  $N_d$  as the independent variable in the cloud system, such that changes in  $N_d$  drive changes in the system, e.g. cloud LWP and albedo (dependent variables), forming a causal relationship. According to the Calculus of Actions (Pearl, 1994), when no confounding effects are present, an observed relationship (seeing) can be used to determine the outcome of an action (doing or causality). In the case of satellite observations, confounding factors 150 can be significantly reduced: for a given satellite snapshot (e.g. covering a  $1^\circ \times 1^\circ$  area), meteorological conditions can be assumed homogenous within a limited space-time frame, enabling one to relate changes in cloud radiative properties to respective changes in  $N_d$  (e.g. Goren and Rosenfeld, 2014; Painemal, 2018). After quantifying the relationship between  $N_d$  and cloud radiative properties in satellite snapshots, we further infer characteristics of the processes governing the cloud system 155 from these relationships with a Markovian methodology, which assumes that processes are related to the observed state of the system with no memory of the past states. One caveat associated with this approach is the difficulty in discerning the causal directions between  $N_d$  and LWP when the system is heavily precipitating and actively removing droplets from the system, as past states of the system cannot be obtained from polar-orbiting satellite snapshots. Because we focus on high cloud fraction

scenes over a marine stratocumulus region, we expect heavily precipitating scenes to be rare in our analyses and assume the observed relationship between  $N_d$  and LWP under precipitating conditions reflects changes in the system if  $N_d$  were perturbed. We leave the validation of this assumption to a future evolution-oriented study that involves the temporal aspect of the cloud system.

To quantify the relationship between  $N_d$  and  $A_c$  in satellite snapshots on a  $1^\circ \times 1^\circ$  grid, we use slopes derived from least squares log-log regressions of  $N_d$  and  $A_c$ , sampled by the MODIS and CERES sensors onboard the polar-orbiting Aqua satellite (1:30 local afternoon overpass, 20-km footprint). We infer this as the cloud albedo susceptibility ( $S_0$ ), represented as follows:

$$S_0 = \frac{d\ln(A_c)}{d\ln(N_d)}. \quad (3)$$

$S_0$  values are only reported if the number of data points is greater than or equal to 5 and the absolute value of the correlation coefficient is greater than 0.2. This provides levels ranging from 25% (minimum required number of samples, 5) to 60% (maximum number of samples within a  $1^\circ$  grid) at which the correlations are statistically significantly according to a Student's t-test. Applying such a threshold on the absolute value of the correlation coefficient between  $A_c$  and  $N_d$  shrinks the sample size of  $S_0$  by  $\sim 19\%$  but does increase the statistical significance of the results by at least 25%. A sensitivity test using  $S_0$  without the correlation coefficient threshold (not shown) indicates no qualitative impacts on the results but a subtle quantitative impact on the occurrence-weighted  $F_0$  (introduced below; from 20.8 to 17.0  $\text{W m}^{-2} \ln(N_d)^{-1}$ ).

Furthermore, the cloud albedo sensitivity to  $N_d$  perturbations is converted to a radiative sensitivity as an intermediate step towards quantifying the radiative forcing, by multiplying the albedo susceptibility by gridbox low-cloud fraction and the incoming solar flux. This is termed radiative susceptibility ( $F_0$ ) hereafter, equivalent to a radiative forcing per  $N_d$  perturbation:

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$$F_0 = \frac{dSW_{\text{TOA}_{\text{up}}}}{d\ln(N_d)} = \frac{dA_c}{d\ln(N_d)} \cdot f_c \cdot SW_{\text{TOA}_{\text{dn}}} [\text{W m}^{-2} \ln(N_d)^{-1}]. \quad (4)$$

Similar forms of this representation of forcing per perturbation have been used in, e.g., Douglas and L'Ecuyer (2019) and Painemal (2018).

Uncertainties embedded in sensors' measuring precision and retrieval techniques/algorithms have been studied and are well understood, and hence minimized in this study by choosing the appropriate sensing channel, and rather strict quality control thresholds for cloud property retrievals (see Section 2 for details). However, uncertainties related to linear regression errors of the slopes ( $\beta_1$ ) of the  $A_c$ – $N_d$  relationship, need to be quantified. A least-squares linear regression takes the form of

$$\hat{y} = \beta_0 + \beta_1 \cdot x \quad (5)$$

where  $\hat{y}$  is the estimated dependent variable of the linear model,  $\beta_0$  is the intercept parameter and  $\beta_1$  is the slope parameter. According to Press et al. (1988), the standard error of the slope parameter ( $S_{\beta_1}$ ) can be expressed as

$$S_{\beta_1} = \sqrt{\frac{SSE/(n-2)}{S_{xx}}} \quad (6)$$

where  $SSE$  is the residual sum of squares, which takes the form of

$$SSE = \Sigma(y_i - \hat{y}_i)^2 = \Sigma(y_i - (\beta_0 + \beta_1 x))^2 \quad (7)$$

$n$  is the number of data points, or the nominal degrees of freedom, in the linear model, and  $S_{xx}$  is the measure of the total amount of variation in the independent variable,  $x$ , which takes the form of:

$$S_{xx} = \Sigma(x_i - \bar{x})^2. \quad (8)$$

To construct confidence intervals around the calculated slope parameter, we use a t-distribution with  $n - 2$  degrees of freedom, implied from the assumptions of the a simple linear regression model (Montgomery and Runger, 2010). As a result, the range of the regressed slopes takes the form of  $\beta_1 \pm t_{\alpha/2, n-2} \cdot S_{\beta_1}$ , where  $100(1 - \alpha)\%$  indicates the confidence interval. We then further

scale the uncertainty associated with the regression slopes by the square root of the ratio of the nominal to effective degree of freedom of  $A_c$  within  $1^\circ \times 1^\circ$  grid boxes to account for the spatiotemporal autocorrelation associated with the regressed field, similar to Myers et al. (2021). We compute the average value of effective degree of freedom using 10 years of CERES data covering the  $10^\circ \times 10^\circ$  study area and the methods of Bretherton et al. (1999). Accordingly, we report the 95% ( $\alpha = 0.05$ ) confidence interval for our regressed slopes that characterize the  $A_c$ - $N_d$  relationship. Note  $t_{0.025, n-2} \approx 2$  for  $n - 2 \geq 6$ .

In order to understand and quantify how cloud albedo susceptibilities vary with changing cloud states (e.g. LWP,  $N_d$ ), meteorological conditions, and aerosol loadings, we aggregate cloud properties, including cloud albedo, and ERA5 meteorological variables (during the Aqua overpass over a 2-hour period) to the same  $1^\circ \times 1^\circ$  grid on which  $S_0$  is calculated. The aggregation method follows a straightforward arithmetic mean of all the pixel-level data points within the grid (0.25° for ERA5 and 20-km for MODIS-CERES), except for cloud properties where we only select overcast footprints for averaging, because  $N_d$  is only retrieved in overcast footprints. Note that requiring overcast conditions for  $N_d$  retrievals at the footprint level does not restrict the  $1^\circ \times 1^\circ$  cloudy scenes analyzed in this study to only overcast scenes, meaning partly cloudy scenes are included in our analyses. In fact, only  $\sim 35\%$  of our  $1^\circ$  cloudy scenes are overcast (see the distribution of  $1^\circ \times 1^\circ$  cloud fraction in Fig. S1). Because the  $1^\circ \times 1^\circ$  cloud fractions of these cloudy scenes analyzed in this work are high (comprising  $\sim 41\%$  of all single-layer liquid cloud scenes over this region), their contribution to the overall cloud radiative effect of the entire cloud population of this region is significant compared to the rest of the (less cloudy) population. Thus, it is important and informative to quantify the response of these high- $f_c$  clouds to aerosol perturbations. That said, it is not the goal of this study to generalize the albedo susceptibility assessment presented here to all marine stratocumulus clouds, especially those with low optical depth, broken or open-cellular structure (high sub-pixel inhomogeneity), conditions under which space-borne  $N_d$  retrievals are highly uncertain (Grosvenor et al., 2018).

## 215 4 Albedo susceptibility in LWP- $N_d$ space

We begin our results section by introducing an informative parameter space, the LWP- $N_d$  space. The choice of these variables is motivated by mutual information analyses that help establish the dominating role of LWP and  $N_d$  in governing albedo susceptibility. Exploring the behavior of these high-level fingerprints (LWP- $N_d$ ) of the system is a pathway to bridge and balance

between the Newtonian and Darwinian approaches that will benefit our understanding of the multi-scale and multidisciplinary

220 nature of the aerosol-cloud system (Mülmenstädt and Feingold, 2018).

#### 4.1 Mutual information analyses reveal primary governance of LWP, $N_d$ and CTH on $S_0$

First, we quantify how much information, treated as entropy (Shannon, 1948), is shared between individual meteorological factors (MFs) and albedo susceptibilities, using a statistical technique called mutual information (MI) analysis (Fig. 1). We follow the methodology in Glenn et al. (2020). Because MI analysis does not require a pre-defined relational function between 225 variables, it handles nonlinear relationships, which is the case for this study (i.e. albedo susceptibility and meteorological factors), just as well as linear relationships. Cloud top heights of marine stratocumulus, marine boundary layer heights, and inversion heights are positively correlated in the setting of the stratocumulus-topped boundary layer (STBL) over the NE Pacific. Therefore, CTH is considered here as a variable indicating one aspect of the cloud state, similar to LWP, while concurrently serving as an indicator of a meteorological condition, namely the depth of the STBL.

230 Although the percentage of shared information between  $S_0$  and meteorological conditions remains very low (less than a percent) for all factors investigated in this study, the MI analysis reveals a leading role of cloud top height ( $\sim 1\%$ ) in terms of covariability with  $S_0$ , whereas the MI of all other factors are comparable to each other (between 0.1% to 0.3%), with boundary layer (BL) meridional winds and  $RH_{ft}$  being the second to highest ( $\sim 0.3\%$ ; Fig. 1a). The leading role of CTH is consistent with the fact that it not only serves as a meteorological index but also often reflects the depth of these marine stratocumulus 235 clouds (a cloud state indicator). The secondary role of BL meridional winds can be explained by the fact that relatively polluted continental flows (northerlies) advect aerosol to our study area ( $120\text{--}130^\circ\text{W}$ ,  $20\text{--}30^\circ\text{N}$ ), whereas southerly flows of a oceanic origin tend to advect cleaner air. The MI between  $S_0$  and the zonal component of the boundary layer wind is half of that with the meridional component (not shown), suggesting meridional winds are more tightly connected to continental/oceanic flows and thereby variations in aerosol loading and  $N_d$  in our study area. This exemplary situation in which meteorology and aerosol 240 conditions covary, points to the importance of considering the covariabilities between aerosol and meteorological drivers.

Next, we examine the unique information contained in individual MFs, if some variable representing a particular cloud state, e.g. LWP,  $N_d$ , or CTH, is known, using the method called conditional MI (CMI) analysis, also following Glenn et al. (2020). When the MI analysis is conditioned on  $N_d$ , LWP, and CTH, the percentage of shared information between  $S_0$  and MFs increases by almost a factor of 10 (Fig. 1b-d), meaning the amount of unique information about  $S_0$  contained in LWP,  $N_d$ , and 245 CTH is almost a factor of 10 greater than that contained in individual MFs. Moreover, we repeat the CMI analysis between the same set of MFs and a randomly permuted  $S_0$  sample space (representing noise; reported as noise-CMI), in order to estimate the baseline signal of these CMIs, by taking the difference between the CMIs and noise-CMIs (Fig. 1, light gray bars). The baseline signal strength suggests that if LWP or  $N_d$  or CTH is known, the unique information remaining in individual MFs that is shared with  $S_0$  is less than a percent different from that which is shared with noise. When one conditions on  $N_d$ , the secondary role of the BL meridional winds is no longer evident, and all MFs beside CTH have almost the same CMI, consistent 250 with the idea of the lower-level wind driving the variability in  $N_d$ . When one conditions on LWP, the leading role of CTH is

much reduced, as CTH correlates with LWP, especially for non-precipitating Sc. Last but not least, when conditioning on CTH, all other MFs have very similar CMIs of about 2%.

From the MI and CMI analyses, we conclude that meteorological conditions affect the albedo susceptibility of low-clouds mainly through governing the states of the clouds, i.e. LWP,  $N_d$  and CTH. If these cloud state indicators are known or pre-defined, e.g. for a given cloud state (LWP,  $N_d$ , CTH), meteorological conditions associated with that state share very little information with the  $S_0$  of those clouds. This is consistent with the concept of "equifinality" (von Bertalanffy, 1950; Mülmenstädt and Feingold, 2018), where multiple, different initial/boundary settings may yield the same realization. In our context, it confirms that many different meteorological conditions can yield the same cloud state (LWP,  $N_d$ , CTH), thereby obscuring unique matchings between meteorological conditions and  $S_0$ , and resulting in overall low MI between MFs and  $S_0$ . These analyses suggest the effective and informative nature of exploring cloud albedo susceptibility in LWP- $N_d$  space.

## 4.2 Mean-state $A_c$ -LWP- $N_d$ relationship

The mean-state  $A_c$ -LWP- $N_d$  relationship of marine low-clouds over the northeast Pacific is shown as an average using equally sized  $N_d$  bins ( $10 \text{ cm}^{-3}$ ). Note that a relationship deduced from equally sized  $N_d$  bins removes the dependence of the relationship on the  $N_d$  distribution, resulting in clearer physical relationships among these properties that are less affected by anthropogenic activities that can cause shifts in the  $N_d$  distribution (Gryspeerdt et al., 2017, 2019a). Moreover, the cloud albedos used in this particular analysis are adjusted to an overhead solar zenith angle ( $\text{SZA} = 0^\circ$ ), in order to obtain a consistent basis for  $A_c$ -LWP- $N_d$  relationships. This is done using the two-stream approximation (Medor and Weaver, 1980), which relates cloud albedo to cloud optical depth and solar zenith angle. Therefore, for a given  $\tau$  we can obtain a theoretical  $A_c$ -SZA relationship using the two-stream approximation. The scattering asymmetry parameter is approximated by a linear function of  $r_e$  following Slingo (1989). We then use the theoretical  $\tau$ -dependent  $A_c$ -SZA relationships to adjust  $A_c$  from measured SZA to overhead SZA.

From a climatological mean-state perspective, precipitating stratocumulus (Sc; approximated by  $r_e > 12 \mu\text{m}$  at cloud top for  $c_w = 2.14 \times 10^6 \text{ kg m}^{-4}$ ) become brighter as  $N_d$  increases (Fig. 2, blue dots). This can be attributed, in part, to the increasing LWP (Fig. 2, black dots), consistent with the cloud lifetime effect (Albrecht, 1989), a macrophysical effect on  $A_c$ . However, the increase in  $A_c$  with increasing  $N_d$  does not stop after the LWP reaches a plateau of  $\sim 120 \text{ g m}^{-2}$  (at  $N_d \approx 20 \text{ cm}^{-3}$ ), suggesting a decrease in cloud effective radius ( $r_e$ ) that contributes to the brightening of the cloud field, a microphysical effect on  $A_c$  (Twomey, 1974, 1977).  $A_c$  reaches a plateau of  $\sim 0.32$  (at  $N_d \approx 100 \text{ cm}^{-3}$ ) when Sc transitions into the non-precipitating regime ( $r_e \leq 12 \mu\text{m}$ ) where negative LWP adjustments to increasing  $N_d$  start to play a dominant role in changes in  $A_c$ .

For non-precipitating Sc, LWP decreases with increasing  $N_d$ , more markedly when the evaporation-entrainment feedback (EEF; Wang et al., 2003; Xue and Feingold, 2006) becomes more active (right hand side of the EEF isoline on Fig. 2). The strong EEF process that drives a dramatic decrease in LWP ( $\text{dln(LWP)}/\text{dln}(N_d) \sim -0.81$ ) leads to a reduction in  $A_c$  with increasing  $N_d$  until LWP drops below  $\sim 55 \text{ g m}^{-2}$  (Fig. 2, red circular outlines), after which  $A_c$  increases with  $N_d$  despite a continuous reduction in LWP, although more than halved in slope ( $\text{dln(LWP)}/\text{dln}(N_d) \sim -0.38$ ) compared to when LWP is above  $55 \text{ g m}^{-2}$ . This increase in  $A_c$  with increasing  $N_d$  after LWP drops below  $55 \text{ g m}^{-2}$  can be explained by a decrease in

entrainment efficiency as LWP decreases (Hoffmann et al., 2020) and an enhanced Twomey effect for less reflective thin clouds (Platnick and Twomey, 1994). The framework for discussion is the commonly used approximation of cloud albedo response to aerosol perturbations (e.g. Bellouin et al., 2020),

$$S_0 = \frac{d\ln(A_c)}{d\ln(N_d)} = \frac{1 - A_c}{3} \left( 1 + \frac{5}{2} \frac{d\ln(LWP)}{d\ln(N_d)} \right) \quad (9)$$

290 in which  $d\ln(LWP)/d\ln(N_d)$  of -0.4 marks the critical value of the LWP adjustment in the entrainment/non-precipitating regime, as it determines the overall sign of the albedo susceptibility approximation, i.e. a warming (negative) or a cooling (positive) effect (e.g. Glassmeier et al., 2021).

295 The climatological mean-state indicates an overall positive response of  $A_c$  to  $N_d$  perturbations (a cooling effect), despite an overall negative LWP adjustment ( $d\ln(LWP)/d\ln(N_d) \sim -0.58$ ) that would be sufficient to overcome the Twomey effect and lead to warming, for these relatively high  $f_c$  non-precipitating Sc over the NE Pacific region (Fig. 2). The strong and sufficiently negative LWP adjustment derived in this study from long-term satellite observations is in agreement with assessment of Glassmeier et al. (2021) for the same region and regime (a lower bound  $d\ln(LWP)/d\ln(N_d) = -0.64$ ), but based on an ensemble of large-eddy simulations. Such agreement between the results learned from an ensemble of model simulated time-evolving nocturnal stratocumulus systems and results deduced from a large composite of remote satellite sensors captured afternoon 300 stratocumulus properties might suggest a robustness of these characteristics regarding the relationship between  $A_c$ ,  $N_d$  and LWP of marine stratocumulus. The result from this work, in addition, points to the importance and necessity of considering the more strongly entraining regime of thicker clouds ( $LWP > 55 \text{ g m}^{-2}$ ) and the weakly entraining while strongly Twomey-brightening regime of thinner clouds ( $LWP < 55 \text{ g m}^{-2}$ ) separately; the strength of LWP adjustment is more than halved in the latter ( $\sim -0.38$ ) compared to the former regime ( $\sim -0.81$ ), allowing the Twomey effect brightening to prevail.

### 305 4.3 Albedo susceptibility and regimes in the LWP- $N_d$ space

Cloud albedo susceptibility is displayed in LWP- $N_d$  space, with the size of the circles indicating the frequency of occurrence of a particular cloud state (Fig. 3). Precipitating Sc ( $r_e > 12 \mu\text{m}$ ) present an overall cloud brightening potential per  $N_d$  perturbation, indicated by the mostly positive susceptibilities, except for some LWP- $N_d$  states that are in the entrainment-evaporation regime (left of the  $r_e = 12 \mu\text{m}$  isoline and right of the EEF isoline on Fig. 3). An occurrence-weighted mean 310 radiative susceptibility ( $F_0$ ) of  $10.5 \pm 0.91 \text{ W m}^{-2} \ln(N_d)^{-1}$  corresponding to the precipitating Sc with positive  $S_o$ , is consistent with the role of the cloud lifetime effect (Albrecht, 1989, and Fig. 2), such that increases in  $N_d$  suppress the warm rain process, favoring the development of deeper and brighter clouds. This regime is hereafter referred to as *the precipitating-brightening regime*. It occurs  $\sim 22\%$  of the time out of all the high cloud fraction, single-layer liquid clouds we analyzed over the NE Pacific, based on this 10-year satellite-derived climatology.

315 For non-precipitating Sc, two regimes emerge in the LWP- $N_d$  space, indicated by the changing sign of albedo susceptibility at  $LWP \approx 55 \text{ g m}^{-2}$ , with thicker Sc ( $LWP > 55 \text{ g m}^{-2}$ ) showing a cloud darkening potential (negative  $S_0$ ) and thinner Sc ( $LWP < 55 \text{ g m}^{-2}$ ) showing a strong cloud brightening potential (positive  $S_0$ ) per  $N_d$  perturbation (Fig. 3). This is consistent with the “inverted V-shape” dependence of mean-state  $A_c$  as a function of  $N_d$  for non-precipitating Sc shown in Fig. 2 (blue dots),

with the turning point being around  $55 \text{ g m}^{-2}$ . As discussed in Section 4.2, the non-precipitating cloud states with negative  $S_0$  are dominated by the entrainment driven LWP adjustment ( $\sim 0.81$ , Fig. 2 brown fitting line) which is double the critical slope value (-0.4) for entering the warming regime (Glassmeier et al., 2021). This entrainment-evaporation regime cloud state (right of the EEF isoline on Fig. 3) with negative  $S_0$  occurs  $\sim 36\%$  of the time out of the cloudy scenes we analyzed. It produces an occurrence-weighted  $F_0 = -20.2 \pm 1.89 \text{ W m}^{-2} \ln(N_d)^{-1}$ , and is hereafter referred to as *the entrainment-darkening regime* (mostly non-precipitating).

The thinner Sc ( $\text{LWP} < 55 \text{ g m}^{-2}$ ) not only possess strong positive albedo susceptibilities for reasons discussed in Section 4.2, but these cloud states also occur the most frequently ( $\sim 37\%$  of the time; Fig. 3). As a result, a dominating positive occurrence-weighted mean  $F_0$  of  $30.7 \pm 1.60 \text{ W m}^{-2} \ln(N_d)^{-1}$  is associated with these non-precipitating cloud states with positive  $S_0$ , hereafter referred to as *the Twomey-brightening (non-precipitating) regime*. Climatologically, the cloud-state dependent albedo susceptibilities and their corresponding frequency of occurrence together determine that the stratocumulus deck over the NE Pacific presents an overall cloud brightening potential with an occurrence-weighted  $F_0$  of  $20.8 \pm 0.96 \text{ W m}^{-2} \ln(N_d)^{-1}$  (Fig. 3), in agreement with the results shown in Fig. 2.

## 5 Meteorological constraints

One of the main questions we want to address is under what meteorological conditions are marine low-clouds most/least susceptible to aerosol perturbations, or in other words, what is the influence of meteorology on albedo and radiative susceptibilities? Then, by quantifying the frequency of occurrence of susceptible conditions, and the potential radiative effect associated therewith, we have the means to quantify the radiative effect of aerosol-cloud interactions. In this section, we assess meteorological constraints on low-cloud albedo susceptibility from multiple perspectives, with a focus on the covariability among meteorological drivers: where to find susceptible and less susceptible conditions in meteorological factor spaces (5.1), the role of seasonal covariability in meteorological conditions (5.2), and the impact of individual meteorological factors on the occurrence of susceptibility regimes and the overall occurrence-weighted radiative susceptibility (5.3).

### 5.1 Albedo susceptibility in meteorology spaces

We map cloud states in the LWP- $N_d$  space (Fig. 3) directly onto meteorological spaces (Fig. 4), to reveal the association between meteorological conditions and the radiative susceptibility regimes identified in Section 4.2. A clear separation of the entrainment-darkening and Twomey-brightening regimes is evident in all 6 meteorological spaces (Fig. 4, brown and green/blue open circles), more markedly in the direction of cloud top height (Fig. 4a-c). Moreover, these 2 regimes tend to cluster in meteorological spaces: the Twomey-brightening regime clusters at low CTH, highest LTS, relatively low SST, and lowest  $\text{RH}_{ft}$ , and the entrainment-darkening regime clusters at higher CTH, lower LTS, higher SST, and higher  $\text{RH}_{ft}$ , compared to the Twomey-brightening regime (Fig. 4). The clustering of these two regimes in these meteorological spaces is consistent with their states in the LWP- $N_d$  space, as stratocumulus with higher cloud tops usually have higher LWP over the NE Pacific region. Therefore, thicker and deeper clouds are more strongly affected by the cloud-top entrainment feedbacks, leading to decreases

in LWP as  $N_d$  increases, whereas thinner and lower Sc are subject to less effective entrainment processes, maintaining the cloud LWP such that an increase in  $N_d$  can sufficiently decrease  $r_e$  and brighten the clouds. The vertical extent of the subtropical marine stratocumulus or the depth of the stratocumulus-topped boundary layer (STBL) is controlled, to first order, by the LTS at longer time scales (Eastman et al., 2017) and  $RH_{ft}$  at shorter time scales (Eastman et al., 2017; Eastman and Wood, 355 2018), such that enhanced LTS (a stronger buoyancy gradient across the inversion) or higher free-tropospheric humidity (less radiative and evaporative cooling), all else being equal, limits the entrainment of free tropospheric air and thereby suppresses the deepening of marine boundary layers. Hence, the primary occurrence of the Twomey-brightening regime is under the highest LTS conditions, however, perhaps counterintuitively, also under the lowest  $RH_{ft}$  conditions (Fig. 4b, e, and f). This is because large-scale meteorological conditions are strongly correlated over eastern subtropical oceans where the Earth's major 360 marine stratocumulus decks are formed (Wood, 2012), such that LTS and  $RH_{ft}$  are negatively correlated (evident in Fig. 4e and further discussed in 4.3.3), as prevailing free-tropospheric subsidence transports dry upper-level air downward and increases the stability.

In contrast, the precipitating-brightening regime tends to spread out in the meteorological spaces, overlapping with the other two regimes, except in the spaces of  $RH_{ft}$  and LTS (e.g. Fig. 4e). This suggests precipitation-suppression driven cloud 365 brightening tends to occur, first, when LTS is weak (less than 21 K), regardless of  $RH_{ft}$  or SST; second, when the free-troposphere is the moistest (> 45%) co-occurring with the highest SST conditions (> 294.5 K) (Fig. 4f). Despite high SST conditions, the precipitating-brightening branch appears under high  $RH_{ft}$ , indicating a dominant role of the free-tropospheric humidity. Here, enhanced free-tropospheric humidity (a reduced humidity gradient across the cloud top) slows/weakens droplet evaporation, creating favorable conditions for precipitation, which is susceptible to aerosol induced warm-rain suppression 370 process, and thereby cloud brightening. This role of  $RH_{ft}$  is reinforced by the fact that the precipitating-brightening branch is displaced from the non-precipitating branch in Fig. 4f, where  $RH_{ft}$  alone determines which susceptibility regimes the clouds will be in at a constant SST.

The fact that two of the susceptibility regimes cluster while the other spreads out in the meteorological spaces serves to expand our discussion on the concept of “equifinality”. We previously discussed that different meteorological conditions may 375 produce the same cloud state (LWP,  $N_d$ ). Here we see that different meteorological conditions may produce the same  $S_0$ . This ties back to the importance of understanding and quantifying the covariabilities between meteorological factors, as multiple environmental factors may be needed to explain all the variability in cloud states (e.g. Chen et al., 2021) and thereby albedo susceptibility.

## 5.2 The role of seasonal covariability in meteorological conditions

380 Monthly climatologies of ERA5 meteorological factors, including LTS, SST,  $RH_{ft}$ , and 700 hPa subsidence, averaged over the NE Pacific show a strong seasonality and a tight correlation among these factors (Fig. 5a). The annual cycle in SST (blue) and 700 hPa vertical velocity (gray) are correlated and anti-correlated with that of the Northern Hemispheric insolation, respectively (not shown), such that summer time (June–September) SST is the highest whereas free-tropospheric subsidence is the weakest due to a weakened Hadley circulation when insolation is at its annual maximum in the Northern Hemisphere. Moreover, the

385 annual cycle in free-tropospheric humidity (black) is very well anti-correlated with that of the free-tropospheric subsidence, leading to a positive (although lagged) correlation between  $\text{RH}_{ft}$  and SST (also evident in Fig. 4f). As the Hadley circulation starts to strengthen in January, indicated by the enhancing 700 hPa subsidence (January to May), and SST over the subtropical ocean remains cool during boreal spring, LTS (red) increases markedly. SST starts to increase as the Northern Hemisphere enters its summer season, resulting in a weakening of the Hadley circulation and the free-tropospheric subsidence, and leading 390 to a continuous decrease in LTS from June until January. As a result, LTS peaks in June, leading the annual maximum in SST by 3 months (Fig. 5a).

In response to the strengthening LTS during boreal spring, both CTH (black) and cloud LWP (blue) decrease, with cloud LWP reaching its annual minimum in May (Fig. 5b). The thinnest clouds of the year give rise to the annual maximum in the occurrence of the Twomey-brightening regime in May, resulting in an annual maximum of  $F_0$  (Fig. 5c). As LTS decreases and 395 SST continues to warm during boreal summer and fall, cloud LWP and CTH increase until December, when LTS is at its annual minimum and the precipitating-brightening regime is at its annual maximum occurrence, resulting in a secondary peak in the annual cycle of  $F_0$ . During the boreal summer months (June–September), when SST is the highest, the entrainment-darkening regime is at its annual maximum occurrence, resulting in the lowest  $F_0$  throughout the annual cycle. The high summertime  $N_d$  also favors the occurrence of the entrainment-darkening regime through the entrainment feedbacks. This is in agreement with 400 the finding that warmer SST over the northeast (NE) Atlantic leads to mostly darkening clouds (Zhou et al., 2021). Although  $F_0$  responds to SST over the NE Pacific the same way as it does over the NE Atlantic, marine low-clouds over the NE Pacific never enter an overall darkening regime, likely due to the co-occurring high free-tropospheric humidity and high SST conditions and thereby a relatively persistent and high occurrence of the precipitating-brightening regime (July–September), which is rarely the case for the high SST conditions over the NE Atlantic in Zhou et al. (2021).

#### 405 5.3 Meteorology affects the occurrence of albedo susceptibility regimes

As discussed in Section 4.1, meteorological or environmental conditions influence the albedo susceptibility of a cloud field to aerosol perturbations through regulating the state of the clouds, e.g.  $N_d$ , LWP. Because individual meteorological factors tend to co-vary with others, we modify the traditional approach of binning results by individual meteorological factors. Instead we 410 bin by a single meteorological factor and allow all others to co-vary. We present albedo susceptibilities in LWP- $N_d$  space as a function of individual meteorological factors (Figs. 6–9), with a focus on the impact of meteorology on the occurrence and the strength of each albedo susceptibility regime.

##### a. Cloud top height (CTH)

415 As cloud top heights of marine Sc increase or as the Sc-topped boundary layers deepen, clouds are more likely to develop higher LWPs and are more likely to precipitate. A pronounced decrease in occurrence-weighted radiative susceptibility with increasing CTH, from  $60.5 \text{ W m}^{-2} \ln(N_d)^{-1}$  to  $-40.3 \text{ W m}^{-2} \ln(N_d)^{-1}$ , is noted (Fig. 6a). The  $F_0$  uncertainties in this section are reported in Tables S1–S4 of the supplementary material. The remarkable decrease in  $F_0$  can be reasoned through two

contributing mechanisms, i) changes in the magnitude of  $S_0$  and ii) a shift in the frequency of occurrence of cloud states (LWP, 420  $N_d$ ), as cloud top elevates. First, regarding changes in the magnitude of  $S_0$ , a clear enhancement in the negative susceptibilities in the entrainment-darkening regime, by -0.16, is evident as CTH increases (Fig. 6a and dashed curves in 6c), consistent with an increasing influence of the entrainment feedbacks as cloud deepens. For the precipitating-brightening Sc,  $S_0$  decreases slightly with increasing CTH, by -0.05, leading to a steady decrease in regime-mean  $F_0$ , by  $-10.1 \text{ W m}^{-2} \ln(N_d)^{-1}$  (Fig. 6b), given little change in the occurrence of the regime. This could reflect two possible balancing mechanisms: i) a balance between warm 425 rain suppression and the increasing precipitation (droplet removal) efficiency with deeper/higher clouds; ii) a balance between warm rain suppression and the strengthening entrainment drying with higher cloud tops.

Second, a pronounced shift in the occurrence of the albedo susceptibility regimes (Fig. 6a and solid curves in 6c) is perhaps more evident, such that the marine Sc over the NE Pacific are more likely to be found in the entrainment-darkening regime (55%) rather than the Twomey-brightening regime (11%) in the highest CTH quartile. This is in contrast to the lowest CTH 430 quartile, where the Twomey-brightening regime (55%) is much more likely to occur than the entrainment-darkening regime (14%). This shift in regime occurrence (and the MFs that define them) as CTH increases is the primary driver of the significant changes in the overall occurrence-weighted  $F_0$ , in which the contribution from the Twomey-brightening regime shrinks by  $44.6 \text{ W m}^{-2} \ln(N_d)^{-1}$ , and the contribution from the entrainment-darkening regime increases by  $55.5 \text{ W m}^{-2} \ln(N_d)^{-1}$  (Fig. 6b). In a nut shell, stronger entrainment and more entrainment drying are expected for clouds with higher cloud tops.

435

### **b. Lower-tropospheric stability (LTS)**

When lower-troposphere stability is low (unstable conditions, leftmost panel on Fig. 7a), clouds are most frequently observed 440 in high-LWP states, consisting of the most frequently occurring precipitating-brightening regime (41% of the time) whose radiative susceptibility contribution is almost entirely offset by that of the less frequently occurring entrainment-darkening regime (33% of the time). This is consistent with the governing role of LTS on stratocumulus-topped marine boundary layer characteristics, such that weaker LTS allows stronger entrainment of free-tropospheric air into the boundary layer, resulting in on-average deeper boundary layers and thicker clouds. As LTS increases, the precipitating-brightening regime occurs less and less frequently (from 41% to 11%), whereas the occurrence of the Twomey-brightening regime increases from 22% to 51% 445 (Fig. 7a), as expected from the suppressing effect of high atmospheric stability on the deepening of STBL.

The impact of LTS on cloud-top entrainment can be directly seen by comparing the strength (darkness of the color) and the weighted  $F_0$ -contribution (labelled) of the entrainment-darkening regime across LTS quartiles (Fig. 7a). However, one should be mindful of the obscuring effect of the covarying  $\text{RH}_{ft}$ : when LTS is low,  $\text{RH}_{ft}$  is high (discussed in Section 5.1 and shown in Fig. 4e), which suppresses the enhanced entrainment-drying that would have occurred if the free-troposphere above were 450 dry. Nevertheless, the entrainment-darkening regime is weakest under the highest LTS condition (rightmost panel on Fig. 7a), compared to the other LTS quartiles. However, the lowest LTS quartile does not exhibit the strongest entrainment-darkening regime, owing to the co-occurring high  $\text{RH}_{ft}$  conditions.

In a nutshell, increasing LTS mostly affects the occurrence of the precipitating-brightening regime (by -25%) and the Twomey-brightening regime (by +24%) (Fig. 7c), leading to changes in the occurrence-weighted- $F_0$  of -11.4 and +30.6 W 455  $m^{-2} \ln(N_d)^{-1}$ , respectively (Fig. 7b).

### c. Free-tropospheric humidity ( $RH_{ft}$ )

The effect of  $RH_{ft}$  on radiative susceptibility has two aspects. First, moister air above cloud top reduces the humidity gradient across the cloud-top inversion thereby weakening the evaporation-entrainment feedback. This is evident in comparisons 460 between the leftmost and rightmost panels on Fig. 8a, where fewer cloud states in LWP- $N_d$  space are represented by a weakened entrainment-darkening regime under higher  $RH_{ft}$  conditions, consistent with the findings in Ackerman et al. (2004) and Chen et al. (2014). Second, as conditions in the free-troposphere become more humid with decreasing atmospheric stability 465 (negatively correlated LTS and  $RH_{ft}$ ), marine low-level clouds are more likely to possess higher LWP and reside in a more favorable environment for precipitation, indicated by the high occurrence of the precipitating-brightening regime (39%) in the highest  $RH_{ft}$  quartile (also consistent with Ackerman et al. (2004)). ERA5 humidity profiles also indicate a positive correlation 470 between  $RH_{ft}$  and the RH within the boundary layer (not shown), further supporting higher LWP. The increase in LWP with increasing  $RH_{ft}$  leads to a shift in cloud state away from the Twomey-brightening regime, towards the other two regimes, but mostly towards the precipitating-brightening regime (Fig. 8a). Worth noting is that the magnitude of these two effects of  $RH_{ft}$  on albedo 475 susceptibility and their occurrence amplify as  $RH_{ft}$  increases: note the steep changes at the highest 20 percentile of  $RH_{ft}$  in Fig. 8b-c.

The Twomey-brightening and the precipitating-brightening regimes are more sensitive (manifested more in their frequency-of-occurrence rather than their strength) to variations in  $RH_{ft}$  and LTS (Fig. 7-8), whereas the sensitivity of the entrainment-darkening regime to these two factors is largely suppressed by the negative correlation between large-scale  $RH_{ft}$  and LTS 475 conditions over this region. This again points to the important role of covarying meteorological conditions in affecting albedo susceptibility.

### d. Sea surface temperature (SST)

480 As sea surface temperature increases over the NE Pacific, radiative susceptibility decreases from 39.9 to  $6.1 \text{ W m}^{-2} \ln(N_d)^{-1}$  (Fig. 9a). First, SST changes are the driver of changes in many other meteorological factors, e.g. surface fluxes, MBL height, LTS, and humidity. Here, we do not attempt to separate out the role of SST on radiative susceptibilities while controlling for other MFs, but rather explore the radiative susceptibility as a function of SST, with all the inherent covariability between SST 485 and other MFs. In general, cloud states shift towards higher LWP and lower  $N_d$ , an indication of thicker clouds with larger SST conditions (more circles to the left of the  $12 \mu\text{m}$  isoline on Fig. 9a rightmost panel). This is consistent with an increase in SST leading to an increase in surface fluxes and a weaker LTS in a well-mixed marine boundary layer, both supporting

the development of deeper Sc with higher LWPs (similar to the response of trade-wind cumulus to warming in Vogel et al. (2016)). Another effect associated with thicker clouds is the creation of favorable conditions for the entrainment feedbacks, 490 which is shown as a strengthening of the entrainment-darkening  $S_0$  (Fig. 9a, brown circles getting darker). As a result, as SST increases, the increasing occurrence of the strengthening entrainment-darkening regime and the decreasing occurrence of the Twomey-brightening regime (Fig. 9c) lead to the overall decrease in  $F_0$ , by  $\sim 34 \text{ W m}^{-2} \ln(N_d)^{-1}$  (Fig. 9a, leftmost vs rightmost).

In the current climate, the free-tropospheric humidity over the NE Pacific correlates well with SST through the seasonality 495 in large-scale circulation (i.e. the free-tropospheric subsidence related to the Hadley circulation), such that higher SST is associated with enhanced above-cloud humidity, favoring the occurrence of the precipitating-brightening regime (Fig. 9c, the “U” shaped occurrence variation of the precipitating-brightening regime). The rebounding of the precipitating-brightening regime at high SST (Fig. 9b and c) partially offsets the darkening potential that would otherwise dominate the overall radiative 500 susceptibility, leading to a warming effect, in the absence of the enhanced free-tropospheric humidity (similar to over the NE Atlantic; Zhou et al., 2021). However, if SST continues to rise in the coming decades, and assuming the same trend observed in Fig. 9, we might expect the NE Pacific stratocumulus region to exhibit an overall darkening potential to aerosol perturbations.

In the assessment of the role of individual MFs, it is important to emphasize that a change in one MF is usually associated 505 with changes in other MFs (the seasonal covariability in meteorological conditions as an example). Our goal in this section has been to retain this covariability between MFs in our analyses, with the aim of quantifying influences of meteorology on radiative susceptibility in the manner in which nature is observed. In selecting one variable for stratification, and allowing all others to co-vary, we come closer to reality than traditional investigations of individual MFs in which all others are held constant. The latter approach only represents a small portion of the natural variability, and the role of covariabilities between MFs is missed.

## 6 Concluding remarks

510 This study quantifies the albedo susceptibility and radiative susceptibility to  $N_d$  perturbations of high  $f_c$  single layer, marine low-clouds over the NE Pacific stratocumulus region, using 10 years of MODIS-retrieved daytime cloud properties and CERES-measured radiative fluxes at the top-of-atmosphere. A novel aspect of this study is the assessment of susceptibility across a LWP- $N_d$  space, such that albedo susceptibility associated with individual cloud states (LWP,  $N_d$ ) and, more importantly, their frequencies of occurrence are quantified. Moreover, the effects of ERA5 meteorological factors and their covariability, 515 on the albedo susceptibility are explored. This allows us to quantify conditions under which low-clouds are most/least susceptible to aerosol perturbations, and how frequently these conditions occur. Robust establishment of three albedo susceptibility regimes is found regardless of meteorological states or environmental conditions, however, the occurrence and strength of these regimes are clearly modified by meteorological conditions. Key findings are:

1. Based on mutual information analysis, LWP,  $N_d$  and CTH are shown to be the governing factors of low-cloud albedo 520 susceptibility. Individual meteorological factors add very little (less than a percent) shared information with  $S_0$  if the

aforementioned three variables are known (Fig. 1). That said, meteorological factors are shown to affect the overall radiative susceptibility of marine Sc but mainly through governing the frequency of occurrence of cloud states, i.e. LWP and  $N_d$ , and thereby the occurrence of each of the susceptibility regimes (Figs. 6-9). This led us to use LWP- $N_d$  as our parameter space, in which we further explore albedo susceptibilities and the brightening versus darkening regimes.

525 2. From a climatological mean-state perspective, LWP and  $N_d$  are negatively correlated for non-precipitating Sc (Fig. 2), consistent with previous polar-orbiting satellite based studies (e.g. Gryspeerdt et al., 2019a; Possner et al., 2020). Results from the current study, however, indicate that despite the negative LWP adjustment, cloud albedo increases with increasing  $N_d$  for non-precipitating Sc overall, pointing to the importance of considering the high-LWP cloud states separately from the low-LWP cloud states, as the negative LWP adjustments are clearly different for thicker versus 530 thinner Sc (Fig. 2).

535 3. When cloud albedo susceptibility is mapped onto the LWP- $N_d$  state space, three susceptibility regimes emerge: i) the Twomey-brightening regime (occurring 37%, contributing  $30.7 \pm 1.55 \text{ W m}^{-2} \ln(N_d)^{-1}$ ), consisting of non-precipitating thinner clouds ( $\text{LWP} < \sim 55 \text{ g m}^{-2}$ ) and consistent with a dominating Twomey effect for clouds of relatively low albedo and weaker entrainment; ii) the entrainment-darkening regime (occurring 36%, contributing  $-20.2 \pm 1.86 \text{ W m}^{-2} \ln(N_d)^{-1}$ ), comprising mostly non-precipitating thicker clouds ( $\text{LWP} > \sim 55 \text{ g m}^{-2}$ ) and consistent with entrainment feedbacks that drive a decrease in LWP with increasing  $N_d$ ; iii) the precipitating-brightening regime (occurring 540 22%, contributing  $10.5 \pm 1.45 \text{ W m}^{-2} \ln(N_d)^{-1}$ ), comprising precipitating clouds with effective radii mostly greater than  $12 \mu\text{m}$  and consistent with the cloud lifetime effect due to a suppressed warm rain process (Fig. 3). An overall cloud brightening potential of  $20.8 \pm 0.96 \text{ W m}^{-2} \ln(N_d)^{-1}$  is found for the marine low-clouds over the NE Pacific stratocumulus region, after the frequency of occurrence of each regime is accounted for.

545 4. When cloud states, along with their associated radiative susceptibilities, are mapped to meteorological spaces of LTS, SST, CTH, and  $\text{RH}_{ft}$ , the entrainment-darkening regime and the Twomey-brightening regime are clearly associated with distinct meteorological conditions. The Twomey-brightening regime occurs most frequently under low CTH, highest LTS, low SST, and lowest  $\text{RH}_{ft}$  conditions. Such a combination of these meteorological factors occurs in May as a result of the seasonally covarying meteorological conditions related to the large-scale circulation over the NE Pacific. The entrainment-darkening regime occurs most frequently under relatively high CTH and intermediate LTS, SST,  $\text{RH}_{ft}$  550 conditions, which prevail during the boreal summer months (July–September). The precipitating-brightening regime mostly manifests in unstable conditions (low LTS), occurring during winter months (November–January), but a very moist free troposphere (co-occurring with high SST in August) also promotes the occurrence of this regime (Figs. 4–5).

555 5. As cloud-top height or marine boundary layer height increases, cloud states shift towards larger LWP, resulting in a pronounced decrease in the Twomey-brightening regime occurrence and a marked increase in the occurrence of the entrainment-darkening regime. This is accompanied by an enhanced entrainment-darkening susceptibility strength and

555 a reduced precipitating-brightening susceptibility strength. As a result,  $F_0$  decreases substantially with increasing CTH, from 60 to  $-40 \text{ W m}^{-2} \ln(N_d)^{-1}$  (Fig. 6).

560 6. The influence of LTS on  $F_0$  is mainly exerted via the frequency of occurrence of each susceptibility regime, rather than its mean  $S_0$ . Strong stability (high LTS) leads to shallower Sc that mostly occur in the Twomey-brightening regime, whereas unstable conditions (low LTS) allow clouds to grow deeper and become more prone to precipitation, leading to high occurrence of the precipitating-brightening regime (Fig. 7).

565 7. A moist free-troposphere has two major impacts on the radiative susceptibility, i) a reduced humidity gradient across the cloud-top inversion weakens the evaporation-entrainment process, leading to a less negative LWP adjustment for thicker non-precipitating clouds; ii) a moist free-troposphere, co-occurring with low LTS, gives rise to a higher occurrence of thicker and deeper clouds, driving a major shift of cloud states away from the Twomey-brightening regime, mostly into the precipitating-brightening regime (Fig. 8).

570 8. The negative correlation between large-scale LTS and  $RH_{ft}$  conditions obscures their individual role in affecting the cloud-top entrainment-evaporation process (Figs. 7-8).

9. Increases in SST lead to a deeper marine boundary layer, lower LTS and thicker clouds. As a result,  $F_0$  decreases with increasing SST, owing to a higher occurrence of deeper clouds (meaning less occurrence of the Twomey-brightening regime) and a stronger entrainment-darkening regime associated with the weakened stability. In contrast to the NE Atlantic (Zhou et al., 2021), moist free-tropospheric conditions, co-occurring with high SSTs, during summertime over the NE Pacific, hamper the role of the strengthening entrainment-darkening regime, by shifting clouds towards the precipitating-brightening regime (Fig. 9).

575 By focusing on this marine stratocumulus dominated region/ regime over the NE Pacific, we were able to separate cloud states into three clearly defined susceptibility regimes in the LWP- $N_d$  space and link the responses to existing understanding of marine stratocumulus. Future work quantifying the occurrence and strength of these three regimes at various oceanic locations, associated with different meteorological regimes/conditions, will enable an extended satellite-based assessment of the radiative susceptibility of global marine low-clouds. Moreover, if aerosol perturbations, natural or anthropogenic, are estimated in some form, the characterization and quantification of radiative susceptibility regimes provided in this study can be used to provide a global estimate of radiative forcing or radiative effect, due to aerosol-marine low-cloud interactions. Such assessments are planned for a follow-on study.

580 *Data availability.* The CERES SSF data are publicly available from NASA's Langley Research Center (<https://satcorps.larc.nasa.gov/>). The fifth-generation ECMWF (ERA5) atmospheric reanalyses of the global climate data are available through the Copernicus Climate Change Service (C3S, <https://cds.climate.copernicus.eu/>).

*Author contributions.* JZ carried out the analysis and wrote the manuscript. XZ, TG and GF contributed to developing the basic ideas, discussing the results, and editing the paper.

585 *Competing interests.* Graham Feingold is a co-editor of ACP. Other than this, the authors declare that they have no conflict of interests

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## References

595 Ackerman, A. S., Kirkpatrick, M. P., Stevens, D. E., and Toon, O. B.: The impact of humidity above stratiform clouds on indirect aerosol climate forcing, *Nature*, 432, 1014–1017, 2004.

Albrecht, B. A.: Aerosols, Cloud Microphysics, and Fractional Cloudiness, *Science*, 245, 1227–1230, <https://doi.org/10.1126/science.245.4923.1227>, 1989.

Bellouin, N., Quaas, J., Gryspenert, E., Kinne, S., Stier, P., Watson-Parris, D., Boucher, O., Carslaw, K., Christensen, M., Daniell, A.-L., Dufresne, J.-L., Feingold, G., Fiedler, S., Forster, P., Gettelman, A., Haywood, J., Lohmann, U., Malavelle, F., Mauritzen, T., and Stevens, B.: Bounding global aerosol radiative forcing of climate change, *Rev. of Geophys.*, 58, e2019RG000660, <https://doi.org/10.1029/2019RG000660>, 2020.

Bretherton, C. S., Widmann, M., Dymnikov, V. P., Wallace, J. M., and Bladé, I.: The Effective Number of Spatial Degrees of Freedom of a Time-Varying Field, *J. Climate*, 12, 1990–2009, [https://doi.org/10.1175/1520-0442\(1999\)012<1990:TENOSD>2.0.CO;2](https://doi.org/10.1175/1520-0442(1999)012<1990:TENOSD>2.0.CO;2), 1999.

600 Bretherton, C. S., Blossey, P. N., and Uchida, J.: Cloud droplet sedimentation, entrainment efficiency, and subtropical stratocumulus albedo, *Geophys. Res. Lett.*, 34, L03 813, [https://doi.org/https://doi.org/10.1029/2006GL027648](https://doi.org/10.1029/2006GL027648), 2007.

Campmany, E., Grainger, R. G., Dean, S. M., and Sayer, A. M.: Automatic detection of ship tracks in ATSR-2 satellite imagery, *Atmos. Chem. Phys.*, 9, 1899–1905, <https://doi.org/10.5194/acp-9-1899-2009>, 2009.

Chen, Y.-C., Christensen, M. W., Xue, L., Sorooshian, A., Stephens, G. L., Rasmussen, R. M., and Seinfeld, J. H.: Occurrence of lower cloud albedo in ship tracks, *Atmos. Chem. Phys.*, 12, 8223–8235, <https://doi.org/10.5194/acp-12-8223-2012>, 2012.

610 Chen, Y.-C., Christensen, M., Stephens, G. L., and Seinfeld, J. H.: Satellite-based estimate of global aerosol–cloud radiative forcing by marine warm clouds, *Nature Geosci.*, 7, 643–646, <https://doi.org/10.1038/ngeo2214>, 2014.

Chen, Y.-S., Yamaguchi, T., Bogenschutz, P. A., and Feingold, G.: Model evaluation and intercomparison of marine warm low cloud fractions with neural network ensembles, *J. Adv. Model. Earth Syst.*, 13, e2021MS002 625, <https://doi.org/https://doi.org/10.1029/2021MS002625>, 2021.

615 Christensen, M. W. and Stephens, G. L.: Microphysical and macrophysical responses of marine stratocumulus polluted by underlying ships: Evidence of cloud deepening, *J. Geophys. Res.-Atmos.*, 116, <https://doi.org/https://doi.org/10.1029/2010JD014638>, 2011.

Christensen, M. W., Suzuki, K., Zambri, B., and Stephens, G. L.: Ship track observations of a reduced shortwave aerosol indirect effect in mixed-phase clouds, *Geophys. Res. Lett.*, 41, 6970–6977, [https://doi.org/https://doi.org/10.1002/2014GL061320](https://doi.org/10.1002/2014GL061320), 2014.

620 Christensen, M. W., Jones, W. K., and Stier, P.: Aerosols enhance cloud lifetime and brightness along the stratus-to-cumulus transition, *P. Natl. Acad. Sci. USA*, 117, 17 591–17 598, <https://doi.org/10.1073/pnas.1921231117>, 2020.

Coakley, J. A., Bernstein, R. L., and Durkee, P. A.: Effect of Ship-Stack Effluents on Cloud Reflectivity, *Science*, 237, 1020–1022, <https://doi.org/10.1126/science.237.4818.1020>, 1987.

625 Coakley, J. A. J. and Walsh, C. D.: Limits to the Aerosol Indirect Radiative Effect Derived from Observations of Ship Tracks, *J. Atmos. Sci.*, 59, 668–680, [https://doi.org/10.1175/1520-0469\(2002\)059<0668:LTTAIR>2.0.CO;2](https://doi.org/10.1175/1520-0469(2002)059<0668:LTTAIR>2.0.CO;2), 2002.

Diamond, M. S., Director, H. M., Eastman, R., Possner, A., and Wood, R.: Substantial Cloud Brightening From Shipping in Subtropical Low Clouds, *AGU Advances*, 1, e2019AV000 111, <https://doi.org/https://doi.org/10.1029/2019AV000111>, 2020.

Douglas, A. and L'Ecuyer, T.: Quantifying variations in shortwave aerosol–cloud–radiation interactions using local meteorology and cloud state constraints, *Atmos. Chem. Phys.*, 19, 6251–6268, <https://doi.org/10.5194/acp-19-6251-2019>, 2019.

630 Durkee, P. A., Noone, K. J., and Bluth, R. T.: The Monterey Area Ship Track Experiment, *J. Atmos. Sci.*, 57, 2523–2541, [https://doi.org/10.1175/1520-0469\(2000\)057<2523:TMASTE>2.0.CO;2](https://doi.org/10.1175/1520-0469(2000)057<2523:TMASTE>2.0.CO;2), 2000.

Eastman, R. and Wood, R.: The Competing Effects of Stability and Humidity on Subtropical Stratocumulus Entrainment and Cloud Evolution from a Lagrangian Perspective, *J. Atmos. Sci.*, 75, 2563–2578, <https://doi.org/10.1175/JAS-D-18-0030.1>, 2018.

635 Eastman, R., Wood, R., and Bretherton, C. S.: Time Scales of Clouds and Cloud-Controlling Variables in Subtropical Stratocumulus from a Lagrangian Perspective, *J. Atmos. Sci.*, 73, 3079–3091, <https://doi.org/10.1175/JAS-D-16-0050.1>, 2016.

Eastman, R., Wood, R., and O. K. T.: The Subtropical Stratocumulus-Topped Planetary Boundary Layer: A Climatology and the Lagrangian Evolution, *J. Atmos. Sci.*, 74, 2633–2656, <https://doi.org/10.1175/JAS-D-16-0336.1>, 2017.

Gassó, S.: Satellite observations of the impact of weak volcanic activity on marine clouds, *J. Geophys. Res.-Atmos.*, 113, <https://doi.org/https://doi.org/10.1029/2007JD009106>, 2008.

640 Glassmeier, F., Hoffmann, F., Johnson, J. S., Yamaguchi, T., Carslaw, K. S., and Feingold, G.: Aerosol-cloud-climate cooling overestimated by ship-track data, *Science*, 371, 485–489, <https://doi.org/10.1126/science.abd3980>, 2021.

Glenn, I. B., Feingold, G., Gristey, J. J., and Yamaguchi, T.: Quantification of the Radiative Effect of Aerosol Cloud Interactions in Shallow Continental Cumulus Clouds, *J. Atmos. Sci.*, 77, 2905–2920, <https://doi.org/10.1175/JAS-D-19-0269.1>, 2020.

645 Goren, T. and Rosenfeld, D.: Decomposing aerosol cloud radiative effects into cloud cover, liquid water path and Twomey components in marine stratocumulus, *Atmos. Res.*, 138, 378–393, <https://doi.org/https://doi.org/10.1016/j.atmosres.2013.12.008>, 2014.

Grosvenor, D. P. and Wood, R.: The effect of solar zenith angle on MODIS cloud optical and microphysical retrievals within marine liquid water clouds, *Atmos. Chem. Phys.*, 14, 7291–7321, <https://doi.org/10.5194/acp-14-7291-2014>, 2014.

650 Grosvenor, D. P., Sourdeval, O., Zuidema, P., Ackerman, A., Alexandrov, M. D., Bennartz, R., Boers, R., Cairns, B., Chiu, J. C., Christensen, M., Deneke, H., Diamond, M., Feingold, G., Fridlind, A., HÄEnerbein, A., Knist, C., Kollias, P., Marshak, A., McCoy, D., Merk, D., Painemal, D., Rausch, J., Rosenfeld, D., Russchenberg, H., Seifert, P., Sinclair, K., Stier, P., van Diedenhoven, B., Wendisch, M., Werner, F., Wood, R., Zhang, Z., and Quaas, J.: Remote Sensing of Droplet Number Concentration in Warm Clouds: A Review of the Current State of Knowledge and Perspectives, *Rev. Geophys.*, 56, 409–453, <https://doi.org/10.1029/2017RG000593>, 2018.

Gryspeerdt, E., Stier, P., and Partridge, D. G.: Satellite observations of cloud regime development: the role of aerosol processes, *Atmos. Chem. Phys.*, 14, 1141–1158, <https://doi.org/10.5194/acp-14-1141-2014>, 2014.

655 Gryspeerdt, E., Quaas, J., and Bellouin, N.: Constraining the aerosol influence on cloud fraction, *J. Geophys. Res.-Atmos.*, 121, 3566–3583, <https://doi.org/10.1002/2015JD023744>, 2016.

Gryspeerdt, E., Quaas, J., Ferrachat, S., Gettelman, A., Ghan, S., Lohmann, U., Morrison, H., Neubauer, D., Partridge, D. G., Stier, P., Takemura, T., Wang, H., Wang, M., and Zhang, K.: Constraining the instantaneous aerosol influence on cloud albedo, *Proc. Natl. Acad. Sci.*, 114, 4899–4904, <https://doi.org/10.1073/pnas.1617765114>, 2017.

660 Gryspeerdt, E., Goren, T., Sourdeval, O., Quaas, J., Mülmenstädt, J., Dipu, S., Unglaub, C., Gettelman, A., and Christensen, M.: Constraining the aerosol influence on cloud liquid water path, *Atmos. Chem. Phys.*, 19, 5331–5347, <https://doi.org/10.5194/acp-19-5331-2019>, 2019a.

Gryspeerdt, E., Smith, T. W. P., O’Keeffe, E., Christensen, M. W., and Goldsworth, F. W.: The Impact of Ship Emission Controls Recorded by Cloud Properties, *Geophys. Res. Lett.*, 46, 12 547–12 555, <https://doi.org/https://doi.org/10.1029/2019GL084700>, 2019b.

Gryspeerdt, E., Goren, T., and Smith, T. W. P.: Observing the timescales of aerosol–cloud interactions in snapshot satellite images, *Atmos. Chem. Phys.*, 21, 6093–6109, <https://doi.org/10.5194/acp-21-6093-2021>, 2021.

Hersbach, H., Bell, B., Berrisford, P., Hirahara, S., Horányi, A., Muñoz-Sabater, J., Nicolas, J., Peubey, C., Radu, R., Schepers, D., Simmons, A., Soci, C., Abdalla, S., Abellan, X., Balsamo, G., Bechtold, P., Biavati, G., Bidlot, J., Bonavita, M., De Chiara, G., Dahlgren, P., Dee,

D., Diamantakis, M., Dragani, R., Flemming, J., Forbes, R., Fuentes, M., Geer, A., Haimberger, L., Healy, S., Hogan, R. J., Hólm, E., Janisková, M., Keeley, S., Laloyaux, P., Lopez, P., Lupu, C., Radnoti, G., de Rosnay, P., Rozum, I., Vamborg, F., Villaume, S., and Thépaut, 670 J.-N.: The ERA5 global reanalysis, *Q. J. Roy. Meteor. Soc.*, 146, 1999–2049, <https://doi.org/10.1002/qj.3803>, 2020.

Hill, A. A., Feingold, G., and Jiang, H.: The Influence of Entrainment and Mixing Assumption on Aerosol–Cloud Interactions in Marine Stratocumulus, *J. Atmos. Sci.*, 66, 1450–1464, <https://doi.org/10.1175/2008JAS2909.1>, 2009.

Hoffmann, F., Glassmeier, F., Yamaguchi, T., and Feingold, G.: Liquid Water Path Steady States in Stratocumulus: Insights from Process-Level Emulation and Mixed-Layer Theory, *J. Atmos. Sci.*, 77, 2203–2215, <https://doi.org/10.1175/JAS-D-19-0241.1>, 2020.

675 Jiang, H., Xue, H., Teller, A., Feingold, G., and Levin, Z.: Aerosol effects on the lifetime of shallow cumulus, *Geophys. Res. Lett.*, 33, L14 806, [https://doi.org/https://doi.org/10.1029/2006GL026024](https://doi.org/10.1029/2006GL026024), 2006.

Klein, S. A. and Hartmann, D. L.: The Seasonal Cycle of Low Stratiform Clouds, *J. Climate*, 6, 1587–1606, [https://doi.org/10.1175/1520-0442\(1993\)006<1587:TSCOLS>2.0.CO;2](https://doi.org/10.1175/1520-0442(1993)006<1587:TSCOLS>2.0.CO;2), 1993.

680 Loeb, N. G., Kato, S., Loukachev, K., and Manalo-Smith, N.: Angular Distribution Models for Top-of-Atmosphere Radiative Flux Estimation from the Clouds and the Earth's Radiant Energy System Instrument on the Terra Satellite. Part I: Methodology, *J. Atmos. Ocean. Technol.*, 22, 338–351, <https://doi.org/10.1175/JTECH1712.1>, 2005.

685 Malavelle, F. F., Haywood, J. M., Jones, A., Gettelman, A., Clarisse, L., Bauduin, S., Allan, R. P., Karset, I. H. H., Kristjánsson, J. E., Oreopoulos, L., Cho, N., Lee, D., Bellouin, N., Boucher, O., Grosvenor, D. P., Carslaw, K. S., Dhomse, S., Mann, G. W., Schmidt, A., Coe, H., Hartley, M. E., Dalvi, M., Hill, A. A., Johnson, B. T., Johnson, C. E., Knight, J. R., O'Connor, F. M., Partridge, D. G., Stier, P., Myhre, G., Platnick, S., Stephens, G. L., Takahashi, H., and Thordarson, T.: Strong constraints on aerosol–cloud interactions from volcanic eruptions, *Nature*, 546, 485–491, <https://doi.org/10.1038/nature22974>, 2017.

Martin, G. M., Johnson, D. W., and Spice, A.: The Measurement and Parameterization of Effective Radius of Droplets in Warm Stratocumulus Clouds, *J. Atmos. Sci.*, 51, 1823–1842, [https://doi.org/10.1175/1520-0469\(1994\)051<1823:TMAPOE>2.0.CO;2](https://doi.org/10.1175/1520-0469(1994)051<1823:TMAPOE>2.0.CO;2), 1994.

690 Mauger, G. S. and Norris, J. R.: Meteorological bias in satellite estimates of aerosol–cloud relationships, *Geophys. Res. Lett.*, 34, <https://doi.org/https://doi.org/10.1029/2007GL029952>, 2007.

Meador, W. E. and Weaver, W. R.: Two-Stream Approximations to Radiative Transfer in Planetary Atmospheres: A Unified Description of Existing Methods and a New Improvement, *J. Atmos. Sci.*, 37, 630–643, [https://doi.org/10.1175/1520-0469\(1980\)037<0630:TSATRT>2.0.CO;2](https://doi.org/10.1175/1520-0469(1980)037<0630:TSATRT>2.0.CO;2), 1980.

695 Minnis, P., Sun-Mack, S., Chen, Y., Khaiyer, M. M., Yi, Y., Ayers, J. K., Brown, R. R., Dong, X., Gibson, S. C., Heck, P. W., Lin, B., Nordeen, M. L., Nguyen, L., Palikonda, R., Smith, W. L., Spangenberg, D. A., Trepte, Q. Z., and Xi, B.: CERES edition-2 cloud property retrievals using TRMM VIRS and Terra and Aqua MODIS data—Part II: Examples of average results and comparisons with other data, *IEEE Trans. Geosci. Remote Sens.*, 49, 4401–4430, <https://doi.org/10.1109/TGRS.2011.2144602>, 2011a.

700 Minnis, P., Sun-Mack, S., Young, D. F., Heck, P. W., Garber, D. P., Chen, Y., Spangenberg, D. A., Arduini, R. F., Trepte, Q. Z., Smith, W. L., Ayers, J. K., Gibson, S. C., Miller, W. F., Hong, G., Chakrapani, V., Takano, Y., Liou, K., Xie, Y., and Yang, P.: CERES edition-2 cloud property retrievals using TRMM VIRS and Terra and Aqua MODIS data—Part I: Algorithms, *IEEE Trans. Geosci. Remote Sens.*, 49, 4374–4400, <https://doi.org/10.1109/TGRS.2011.2144601>, 2011b.

Montgomery, D. C. and Rung, G. C.: *Applied Statistics and Probability for Engineers*, 5th Edition, Wiley, 2010.

705 Myers, T. A., Scott, R. C., Zelinka, M. D., Klein, S. A., Norris, J. R., and Caldwell, P. M.: Observational constraints on low cloud feedback reduce uncertainty of climate sensitivity, *Nat. Clim. Chang.*, 11, 501–507, <https://doi.org/https://doi.org/10.1038/s41558-021-01039-0>, 2021.

Mülmenstädt, J. and Feingold, G.: The Radiative Forcing of Aerosol–Cloud Interactions in Liquid Clouds: Wrestling and Embracing Uncertainty, *Curr. Clim. Change Rep.*, 4, 23–40, <https://doi.org/10.1007/s40641-018-0089-y>, 2018.

Painemal, D.: Global Estimates of Changes in Shortwave Low-Cloud Albedo and Fluxes Due to Variations in Cloud Droplet Number Concentration Derived From CERES-MODIS Satellite Sensors, *Geophys. Res. Lett.*, 45, 9288–9296, [https://doi.org/https://doi.org/10.1029/2018GL078880](https://doi.org/10.1029/2018GL078880), 2018.

710 Painemal, D., Minnis, P., and Sun-Mack, S.: The impact of horizontal heterogeneities, cloud fraction, and liquid water path on warm cloud effective radii from CERES-like Aqua MODIS retrievals, *Atmos. Chem. Phys.*, 13, 9997–10 003, <https://doi.org/10.5194/acp-13-9997-2013>, 2013.

Pearl, J.: A Probabilistic Calculus of Actions, in: *Uncertainty Proceedings 1994*, edited by de Mantaras, R. L. and Poole, D., pp. 454–462, 715 Morgan Kaufmann, San Francisco, CA, <https://doi.org/https://doi.org/10.1016/B978-1-55860-332-5.50062-6>, 1994.

Platnick, S. and Twomey, S.: Determining the susceptibility of cloud albedo to changes in droplet concentration with the advanced very high resolution radiometer, *J. Appl. Meteorol.*, 33, 334–347, [https://doi.org/10.1175/1520-0450\(1994\)033<0334:DTSOCA>2.0.CO;2](https://doi.org/10.1175/1520-0450(1994)033<0334:DTSOCA>2.0.CO;2), 1994.

Platnick, S., King, M. D., Ackerman, S. A., Menzel, W. P., Baum, B. A., Riedi, J. C., and Frey, R. A.: The MODIS cloud products: algorithms and examples from Terra, *IEEE Trans. Geos. Remote Sens.*, 41, 459–473, <https://doi.org/10.1109/TGRS.2002.808301>, 2003.

720 Possner, A., Eastman, R., Bender, F., and Glassmeier, F.: Deconvolution of boundary layer depth and aerosol constraints on cloud water path in subtropical stratocumulus decks, *Atmos. Chem. Phys.*, 20, 3609–3621, <https://doi.org/10.5194/acp-20-3609-2020>, 2020.

Press, W. H., Flannery, B. P., Teukolsky, S. A., and Vetterling, W. T.: *Numerical Recipes in C: The Art of Scientific Computing*, Cambridge University Press, USA, 1988.

725 Shannon, C. E.: A mathematical theory of communication, *Bell Syst. Tech. J.*, 27, 379–423, <https://doi.org/10.1002/j.1538-7305.1948.tb01338.x>, 1948.

Slingo, A.: A GCM Parameterization for the Shortwave Radiative Properties of Water Clouds, *J. Atmos. Sci.*, 46, 1419–1427, [https://doi.org/10.1175/1520-0469\(1989\)046<1419:AGPFTS>2.0.CO;2](https://doi.org/10.1175/1520-0469(1989)046<1419:AGPFTS>2.0.CO;2), 1989.

730 Su, W., Corbett, J., Eitzen, Z., and Liang, L.: Next-generation angular distribution models for top-of-atmosphere radiative flux calculation from CERES instruments: methodology, *Atmos. Meas. Tech.*, 8, 611–632, <https://doi.org/10.5194/amt-8-611-2015>, 2015.

Toll, V., Christensen, M., Quaas, J., and Bellouin, N.: Weak average liquid-cloud-water response to anthropogenic aerosols, *Nature*, 572, 51–55, <https://doi.org/10.1038/s41586-019-1423-9>, 2019.

Twomey, S.: Pollution and the planetary albedo, *Atmospheric Environment*, 8, 1251–1256, [https://doi.org/10.1016/0004-6981\(74\)90004-3](https://doi.org/10.1016/0004-6981(74)90004-3), 1974.

735 Twomey, S.: The Influence of Pollution on the Shortwave Albedo of Clouds, *J. Atmos. Sci.*, 34, 1149–1152, [https://doi.org/10.1175/1520-0469\(1977\)034<1149:TIOPOT>2.0.CO;2](https://doi.org/10.1175/1520-0469(1977)034<1149:TIOPOT>2.0.CO;2), 1977.

Vogel, R., Nuijens, L., and Stevens, B.: The role of precipitation and spatial organization in the response of trade-wind clouds to warming, *J. Adv. Model. Earth Syst.*, 8, 843–862, <https://doi.org/https://doi.org/10.1002/2015MS000568>, 2016.

von Bertalanffy, L.: The Theory of Open Systems in Physics and Biology, *Science*, 111, 23–29, <https://doi.org/10.1126/science.111.2872.23>, 1950.

740 Wang, H., Rasch, P. J., and Feingold, G.: Manipulating marine stratocumulus cloud amount and albedo: a process-modelling study of aerosol-cloud-precipitation interactions in response to injection of cloud condensation nuclei, *Atmos. Chem. Phys.*, 11, 4237–4249, <https://doi.org/10.5194/acp-11-4237-2011>, 2011.

Wang, S., Wang, Q., and Feingold, G.: Turbulence, Condensation, and Liquid Water Transport in Numerically Simulated Nonprecipitating Stratocumulus Clouds, *J. Atmos. Sci.*, 60, 262–278, [https://doi.org/10.1175/1520-0469\(2003\)060<0262:TCALWT>2.0.CO;2](https://doi.org/10.1175/1520-0469(2003)060<0262:TCALWT>2.0.CO;2), 2003.

745 Wielicki, B. A., Barkstrom, B. R., Harrison, E. F., Lee, R. B., Smith, G. L., and Cooper, J. E.: Clouds and the Earth's Radiant Energy System (CERES): An Earth Observing System Experiment, *B. Am. Meteor. Soc.*, 77, 853–868, [https://doi.org/10.1175/1520-0477\(1996\)077<0853:CATERE>2.0.CO;2](https://doi.org/10.1175/1520-0477(1996)077<0853:CATERE>2.0.CO;2), 1996.

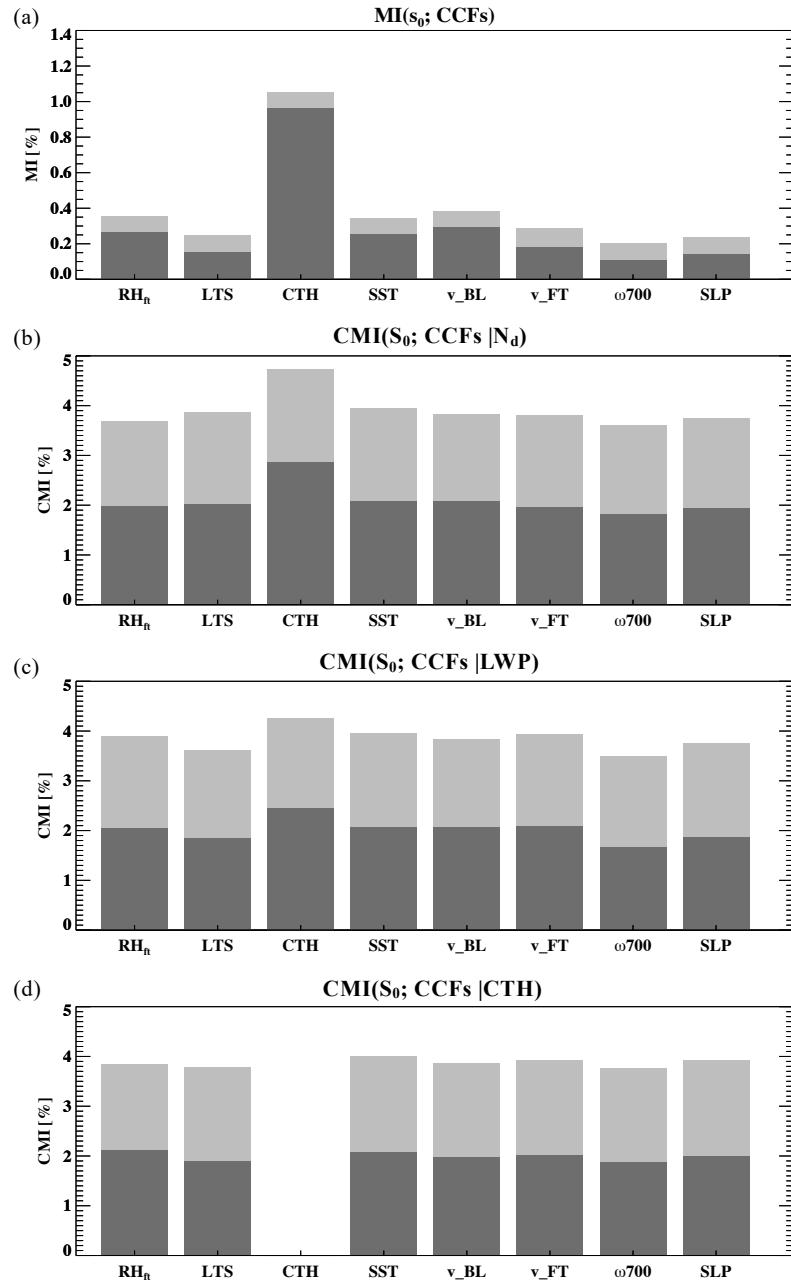
Wood, R.: Stratocumulus Clouds, *Mon. Wea. Rev.*, 140, 2373–2423, <https://doi.org/10.1175/MWR-D-11-00121.1>, 2012.

Xue, H. and Feingold, G.: Large-Eddy Simulations of Trade Wind Cumuli: Investigation of Aerosol Indirect Effects, *J. Atmos. Sci.*, 63, 750 1605–1622, <https://doi.org/10.1175/JAS3706.1>, 2006.

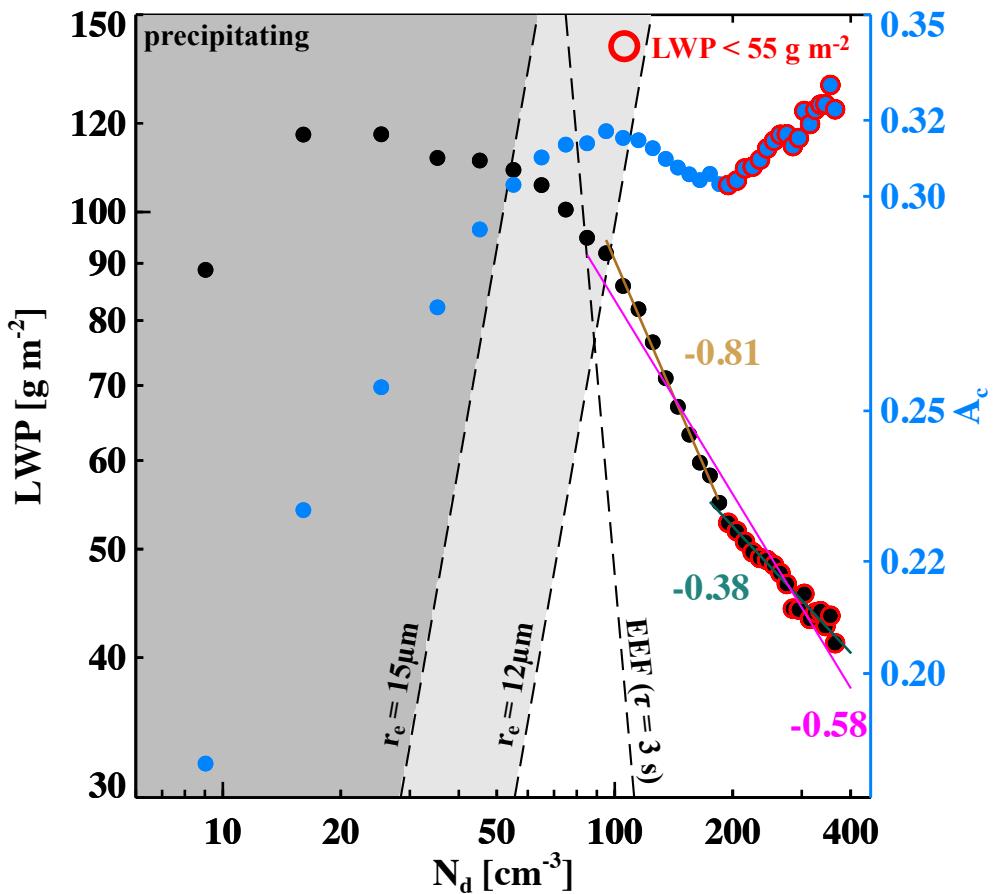
Xue, H., Feingold, G., and Stevens, B.: Aerosol Effects on Clouds, Precipitation, and the Organization of Shallow Cumulus Convection, *J. Atmos. Sci.*, 65, 392–406, <https://doi.org/10.1175/2007JAS2428.1>, 2008.

Yuan, T., Remer, L. A., and Yu, H.: Microphysical, macrophysical and radiative signatures of volcanic aerosols in trade wind cumulus observed by the A-Train, *Atmos. Chem. Phys.*, 11, 7119–7132, <https://doi.org/10.5194/acp-11-7119-2011>, 2011.

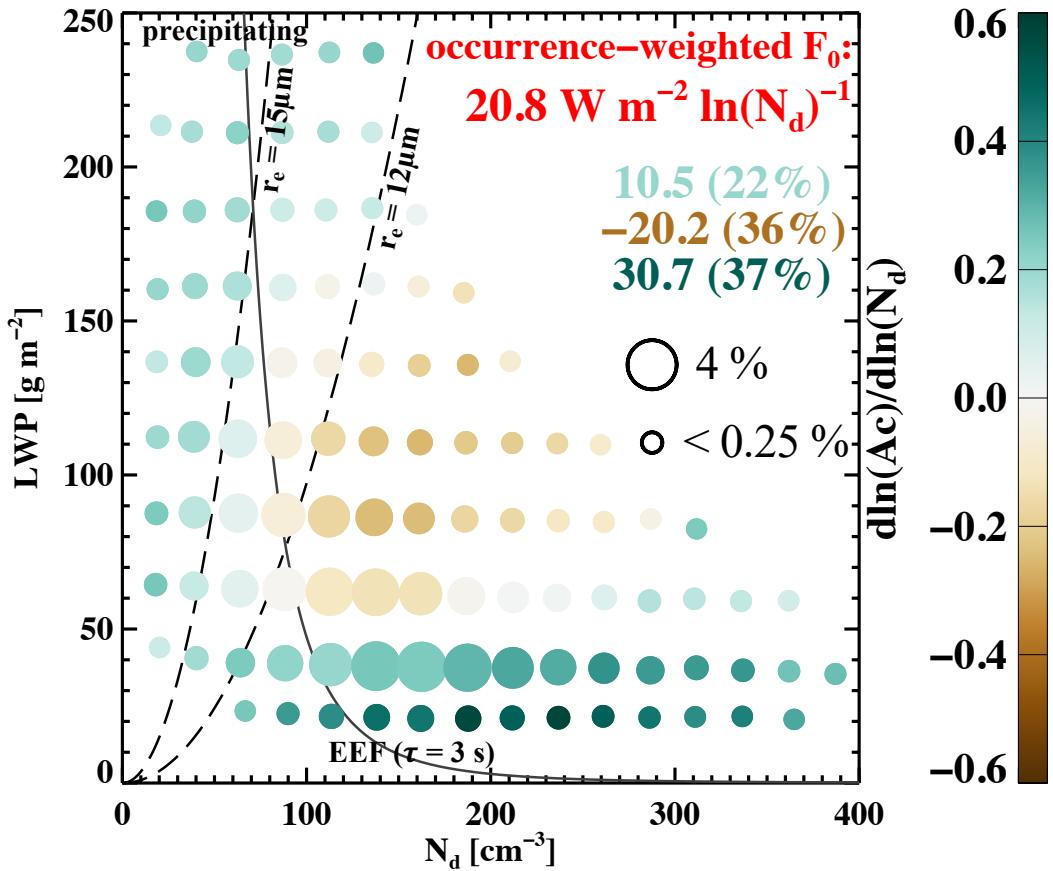
755 Zhou, X., Zhang, J., and Feingold, G.: Sea surface temperature control on the aerosol-induced brightness of marine clouds over the North Atlantic Ocean, *Geophys. Res. Lett.*, manuscript, 2021.



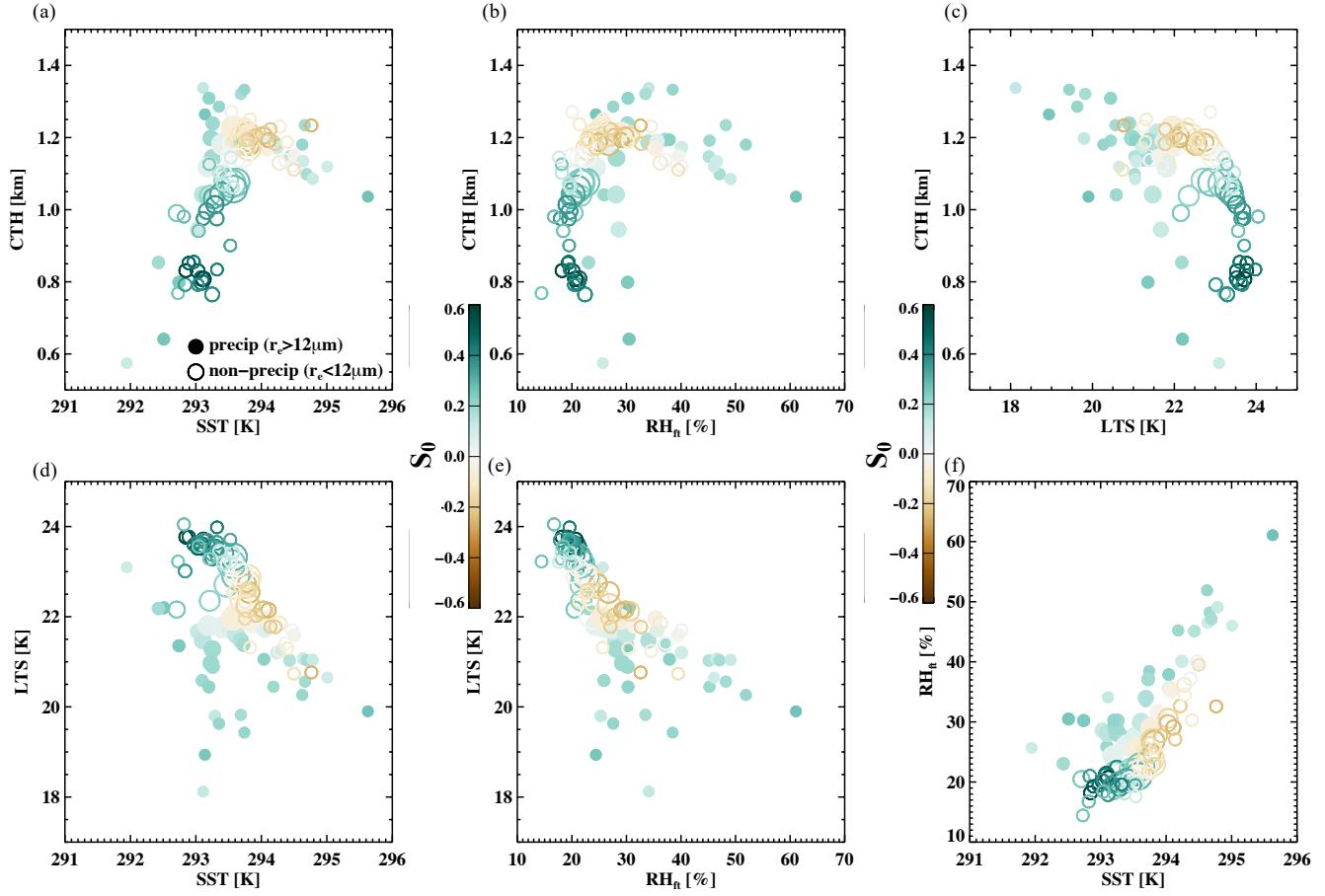
**Figure 1.** **a)** Mutual information (MI; dark gray) for  $s_0$  and 8 meteorological factors (MFs), including  $RH_{ft}$ , LTS, CTH, SST, BL and FT winds, 700 hPa vertical velocity ( $\omega 700$ ), and sea level pressure (SLP). **b-d)** Conditional MI (CMI; dark gray) for  $s_0$  and the 8 MFs, conditioned by  $N_d$ , LWP, and CTH, respectively. Noise-CMI (light gray) is represented by the (conditional) mutual information between MFs and randomly permuted  $s_0$  sample space (effectively noise).



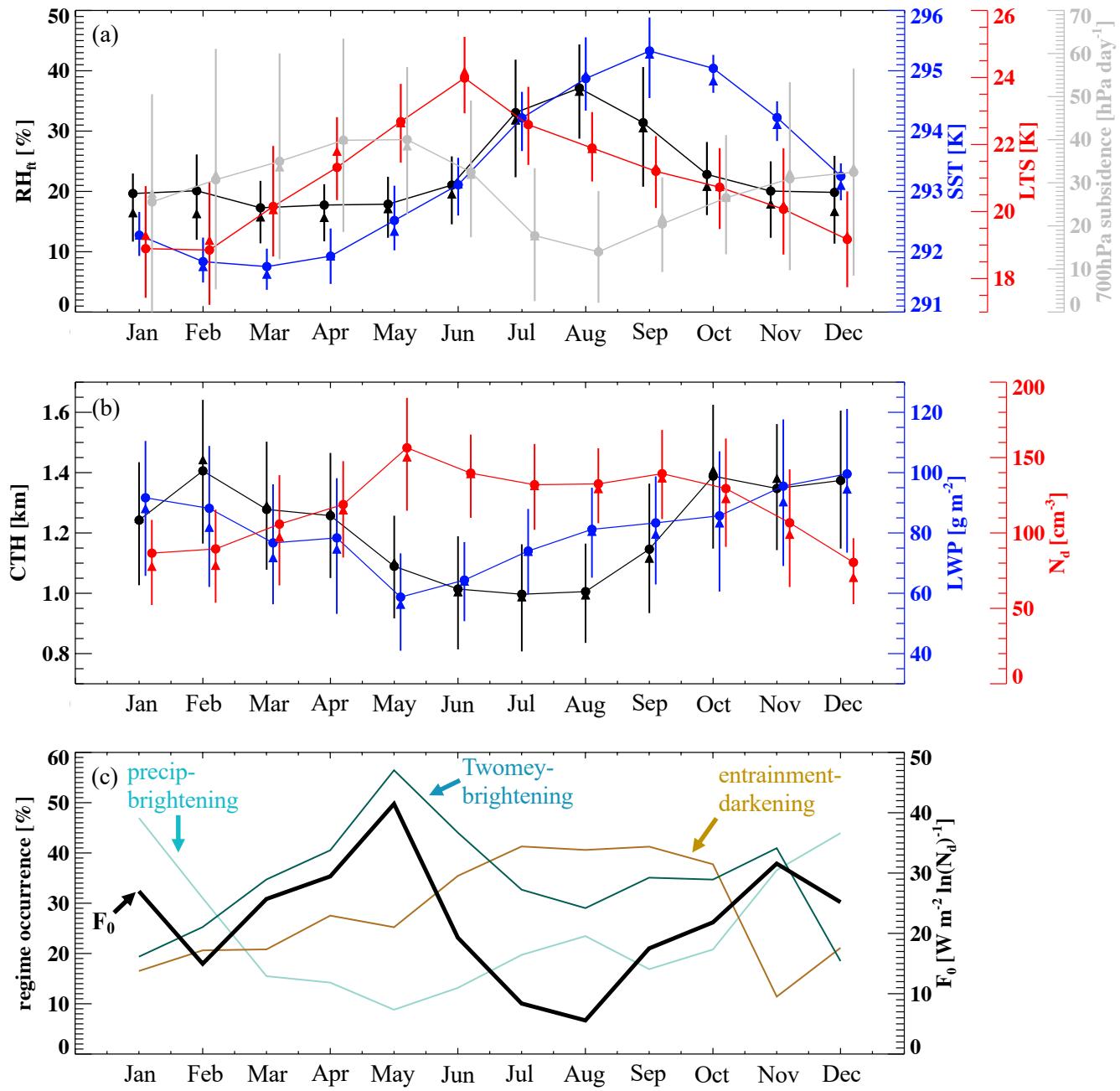
**Figure 2.** Mean liquid water path (LWP; black dots) and cloud albedo ( $A_c$ ; blue dots) of each cloud droplet number concentration ( $N_d$ ) bin (bin-size of  $10 \text{ cm}^{-3}$ ). Values are shown on logarithm scales. Isolines of evaporation-entrainment feedback (EEF; phase relaxation timescale of 3 seconds), effective radius ( $r_e$ ) of  $12 \mu\text{m}$  and  $15 \mu\text{m}$  (commonly used measures of precipitation) based on an adiabatic condensation rate of  $2.14 \times 10^6 \text{ kg m}^{-4}$ , shades of grey background colors represent a general indicator of likelihood of precipitation, and bin-mean LWP less than  $55 \text{ g m}^{-2}$  are highlighted by red circular outlines. The linear regressed slopes of  $\ln(\text{LWP}) - \ln(N_d)$  for all non-precipitating clouds (magenta), non-precipitating clouds with  $\text{LWP} > 55 \text{ g m}^{-2}$  (brown), and non-precipitating clouds with  $\text{LWP} < 55 \text{ g m}^{-2}$  (green) are also indicated.



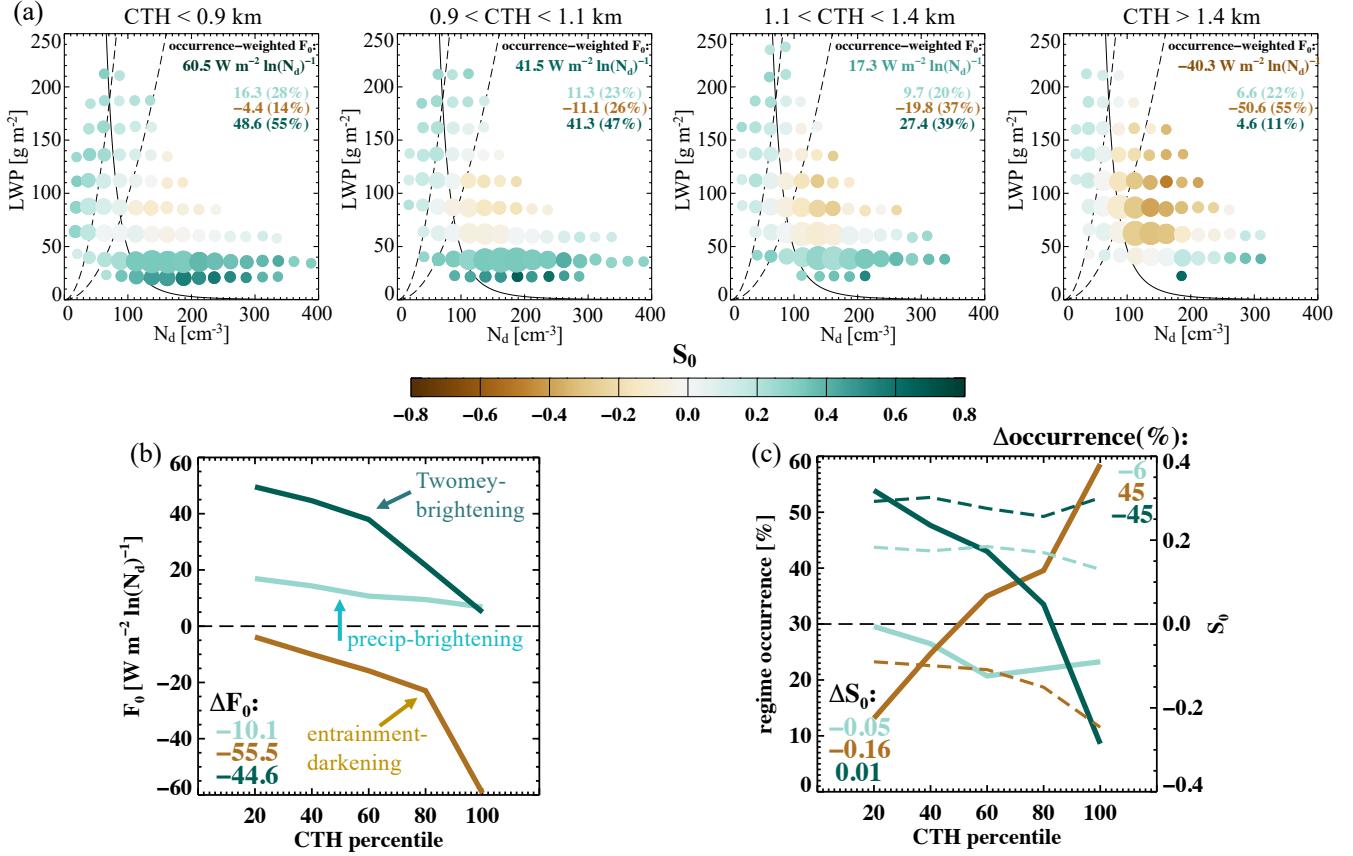
**Figure 3.** Cloud albedo susceptibility ( $S_0$ , colored filled circles) in LWP- $N_d$  space, as bin means (bin-size of  $25 \text{ g m}^{-2}$  and  $25 \text{ cm}^{-3}$ ). Isolines of  $r_e$  of  $12 \mu\text{m}$  and  $15 \mu\text{m}$  (black dashed) and EEF (as in Fig. 2) are indicated. Size of the filled circles in each panel indicates the relative frequency of occurrence of each bin (reference circle sizes with corresponding occurrence are indicated). Mean radiative susceptibility ( $F_0$ ) weighted by the frequency of occurrence of each LWP- $N_d$  bin is printed in red (named “occurrence-weighted  $F_0$ ”), under which is a decomposition of  $F_0$  into precipitating-brightening (light green; positive susceptibility states with effective radii greater than  $12 \mu\text{m}$ ), entrainment-darkening (brown; negative susceptibility states and right-hand side of the EEF isoline), and Twomey-brightening (dark green; non-precipitating states with positive susceptibilities) regimes, with the occurrence of each regime in parentheses.



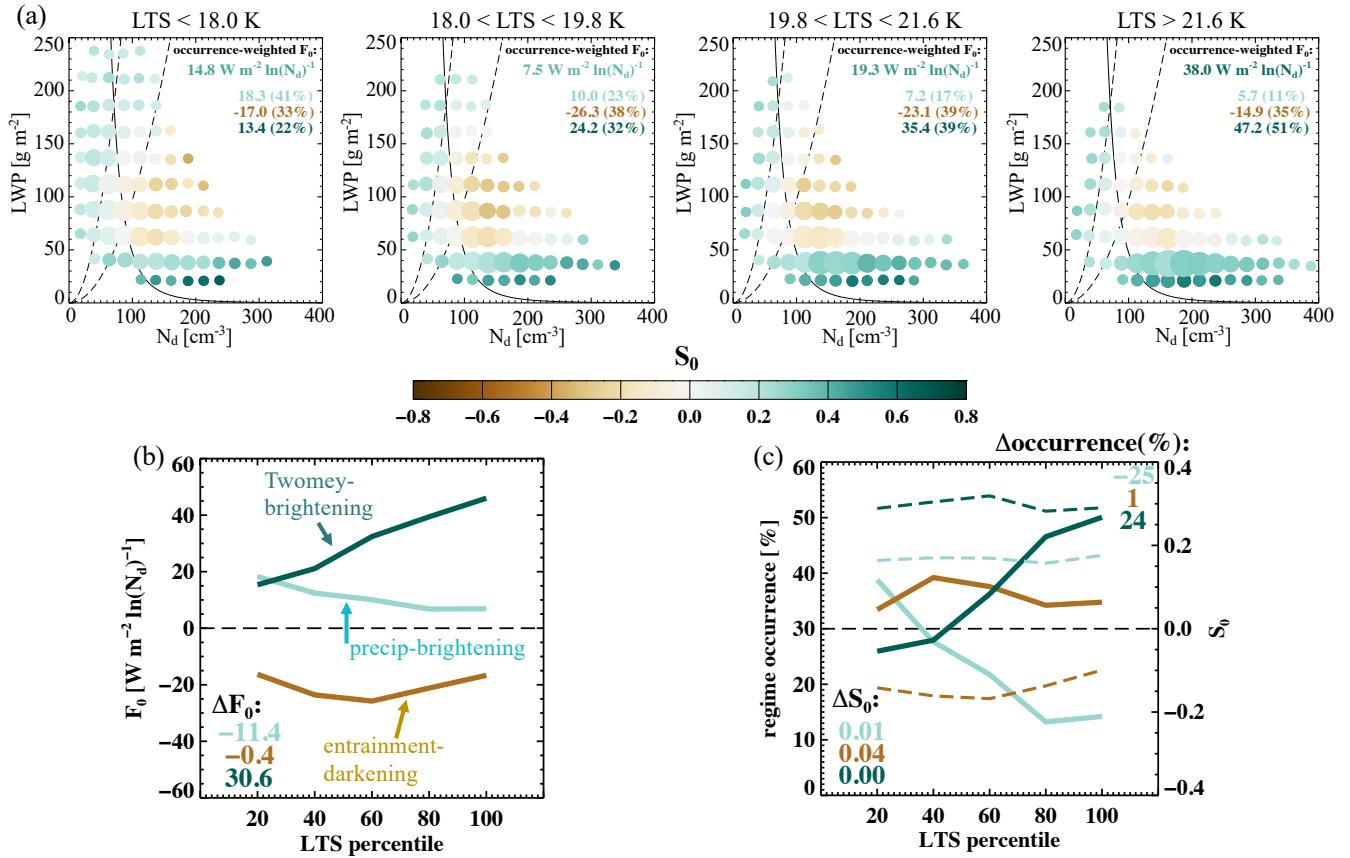
**Figure 4.** Mean meteorological/cloud state conditions associated with each LWP- $N_d$  bin in Fig. 3, in the space of **a)** CTH-SST, **b)** CTH- $RH_{ft}$ , **c)** CTH-LTS, **d)** LTS-SST, **e)** LTS- $RH_{ft}$ , and **f)**  $RH_{ft}$ -SST. Size and color of the circles represent the frequency of occurrence and the mean  $S_0$  of that LWP- $N_d$  bin, respectively, as shown in Fig. 3. Precipitating clouds (based on a  $r_e$  threshold of  $12 \mu m$ ) and non-precipitating clouds are indicated by filled and open circles, respectively.



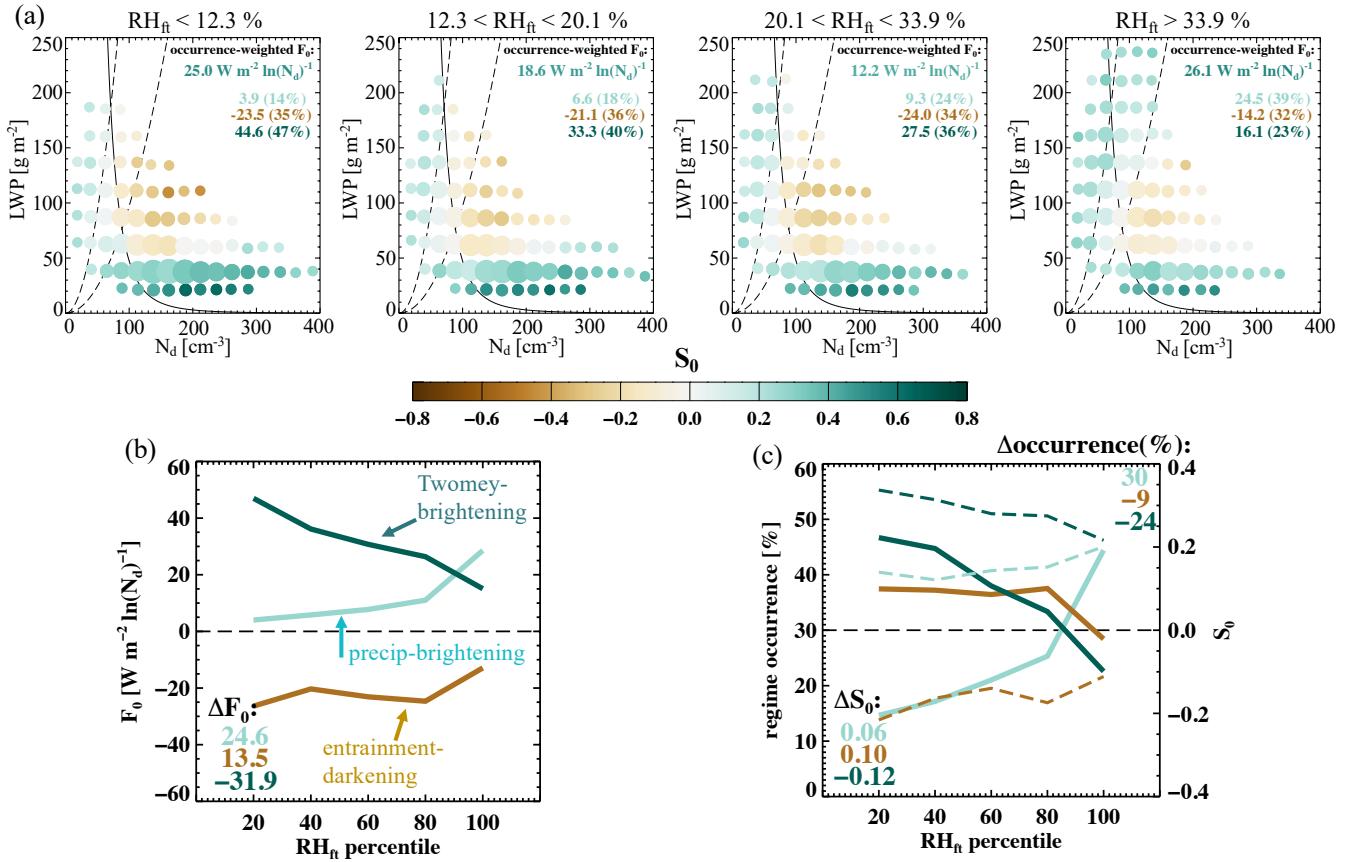
**Figure 5.** Annual cycle of **a)** ERA5 RH<sub>ft</sub> (black), SST (blue), LTS (red), 700 hPa subsidence (gray), **b)** MODIS CTH (black), LWP (blue), and N<sub>d</sub> (red), as monthly means (filled circles), medians (filled triangles), and interquartile ranges (vertical bars). **c)** Annual cycle of occurrence-weighted F<sub>0</sub> (black) and the occurrence of each albedo susceptibility regime (colored).



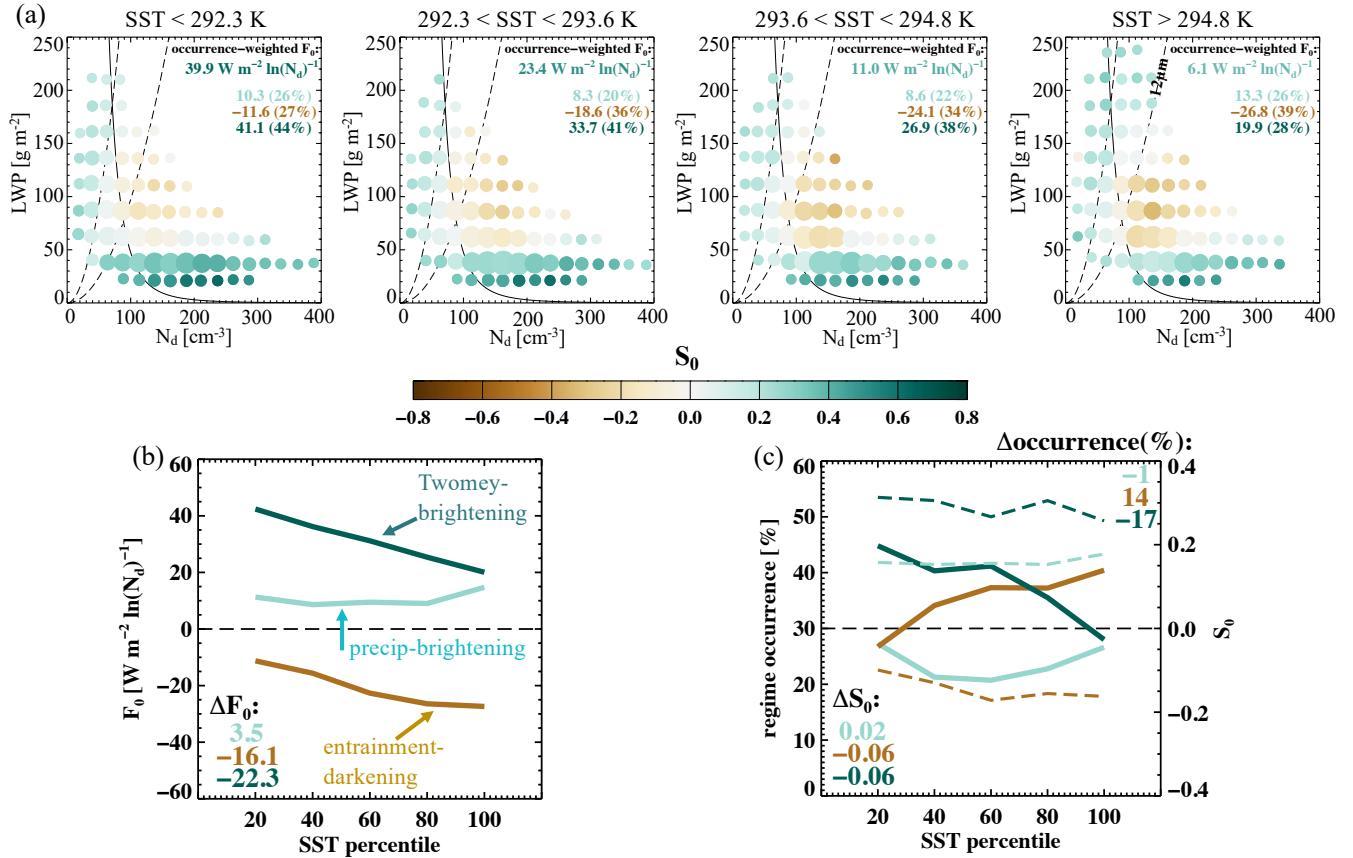
**Figure 6.** **a)** as in Fig. 3, but conditioned on cloud top height (CTH) quartiles. Occurrence-weighted mean **b)**  $F_0$  and **c)** regime-mean  $S_0$  (dashed curves) and regime-occurrence (solid curves) of the 3 albedo susceptibility regimes (defined in Fig. 3) as a function of CTH, increment of 20 percentile.



**Figure 7.** As in Fig. 6, but for lower-tropospheric stability (LTS).



**Figure 8.** As in Fig. 6, but for free-tropospheric relative humidity (RH<sub>ft</sub>).



**Figure 9.** As in Fig. 6, but for sea surface temperature (SST).