

1     **Eastward-propagating planetary wave in the polar**  
2                     **middle atmosphere**

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**Abstract.** We presented the global variations of the eastward propagating wavenumber 1 (E1), 2 (E2), 3 (E3), and 4 (E4) planetary waves (PWs) and their diagnostic results in the polar middle atmosphere, using MERRA-2 temperature and wind datasets in 2019. It is clearly shown that the eastward wave modes exist during winter periods with westward background wind in both hemispheres. The maximum wave amplitudes in the southern hemisphere (SH) are slightly larger and lie lower than those in the northern hemisphere (NH). It is also found that the wave perturbations peak at lower latitudes with smaller amplitude as the wavenumber increases. The period of the E1 mode varies from 3 to 5 days in both hemispheres, while the period of E2 mode is slightly longer in the NH ( $\sim 48$  h) than in the SH ( $\sim 40$  h). The periods of the E3 are  $\sim 30$  h in both SH and NH, and the period of E4 is  $\sim 24$  h. Though the wave periods become shorter as the wavenumber increases, their mean phase speeds are relatively stable, which are  $\sim 53$ ,  $\sim 58$ ,  $\sim 55$ , and  $\sim 52$  m/s at  $70^\circ$  latitudes for W1, W2, W3, and W4, respectively. The eastward PWs occur earlier with increasing zonal wavenumber, which agrees well with the seasonal variations of the critical layers generated by the background wind~~seasonal variations of the background zonal wind through the generation of critical layers~~. Diagnostic analysis also shows that the mean flow instability in the upper stratosphere and upper mesosphere may both contribute to the amplification of the eastward PWs.

## 1 Introduction

Large amplitude planetary waves are dominant in the stratosphere, mesosphere, and lower thermosphere region and their interaction with zonal mean winds is the primary driving force of atmospheric dynamics. In addition, sudden stratospheric warmings (SSWs) and quasi-biennial oscillation (QBO) events can dynamically couple the entire atmosphere from the lower atmosphere to the ionosphere ((Li et al., 2020; Yamazaki et al., 2020; Yadav et al., 2019; Matthias and Ern, 2018; Stray et al., 2015)~~Li et al., 2020; Yamazaki et al., 2020; Yadav et al., 2019; Matthias and Ern, 2018; Stray et al., 2015~~). Westward propagating planetary wave is one of the prominent features during austral and boreal summer ~~periods~~. Westward quasi-2-day waves (Q2DWs) are the most obvious representative waves and one of the most investigated phenomena by planetary wave observations. Most previous studies focused on westward propagating Q2DWs, including zonal wavenumbers of 2 (W2), 3 (W3), and 4 (W4) modes (Lainer et al., 2018; Gu et al., 2018b; Wang et al., 2017; Pancheva et al., 2016; Gu et al., 2016a; Gu et al., 2016b; Lilienthal and Jacobi, 2015; Gu et al., 2013; Limpasuvan and Wu, 2009; Salby, 1981)~~(Lainer et al., 2018; Gu et al., 2018c; Wang et al., 2017; Pancheva et al., 2017; Gu et al., 2016a; Gu et al., 2016b; Lilienthal and Jacobi, 2015; Gu et al., 2013; Limpasuvan and Wu, 2009; Salby, 1981)~~. The ~~s~~Seasonal variations of the occurrence date, peak amplitude, and wave period for

eastward Q2DWs are rarely studied (Gu et al., 2017; Lu et al., 2013; Alexander and Shepherd, 2010; Sandford et al., 2008; Palo et al., 2007; Merzlyakov and Pancheva, 2007; Manney and Randel, 1993; Venne and Stanford, 1979)(~~Gu et al., 2017; Lu et al., 2013; Alexander and Shepherd, 2010; Venne and Stanford, 2010; Merzlyakov and Pancheva, 2007; Palo et al., 2007; Sandford et al., 2007; Manney et al., 1993~~).

Q2DWs usually maximize after the summer solstice in middle latitudes. The largest wave amplitudes are observed near the mesopause in January–February in the Southern Hemisphere (SH), while in the Northern Hemisphere (NH) in July–August (Tunbridge et al., 2011)(~~Tunbridge et al., 2011~~). ~~The~~ W3 and W4 Q2DWs reach ~~a maximum~~ amplitudes during austral and boreal summer ~~periods~~ in the mesosphere and lower thermosphere, respectively. The westward Q2DWs activity has an obvious seasonal variation (Liu et al., 2019; Gu et al., 2018b; Rao et al., 2017)(~~Liu et al., 2019; Gu et al., 2018d; Rao et al., 2016~~). ~~Tunbridge et al. (2011)~~ Tunbridge et al. (2011) have ~~long-term~~ observed long-term Q2DW in the NH and SH, and found that ~~the~~ W3 is generally stronger than ~~those~~ of the other two modes in the SH, reaching ~~an~~ the amplitude of  $\sim 12$ -K, while ~~the~~ W4 is stronger than W3 in the NH, ~~and can reaching~~  $\sim 4$ -K. In addition, W4 is generally longer lived than W3, and W4 is still observed after W3 has ended. The results of Liu et al. (2004)(~~Liu and H. L. (2004)~~) show that ~~the~~ wave source, instability, critical layer, and mean zonal wind



~~structure~~ are the primary reasons for the seasonal variation of Q2DWs. Gu et al. (2019)~~Gu et al. (2018d)~~ have the long-term observation of satellite datasets in the SH and found that the strongest events of W2, W3, and W4 are delayed with increasing the zonal wavenumber and would be ~~indistinguishable~~~~confused~~ during ~~the SSWs period~~. Then wave periods of W4, W3 and W2 vary between ~41-56 h, ~45-52 h, and ~45-48 h, respectively.~~Then the wave period of W3 primarily fluctuated in 45-52 h, while the W4 varies is concentrated between 41-56 h, and the W2 is primarily distributed in 45-48 h.~~ In addition, W2 can be observed in global satellite datasets, showing weaker amplitude than W3 and W4 in the NH and SH (Meek et al., 1996)~~(Meek et al., 1996)~~. The propagation and amplification of Q2DWs are primarily ~~modulated~~dominated by instability, refractive index, and critical layer, while the variation of background wind may cause different zonal wavenumber events (Gu et al., 2016a; Gu et al., 2016b)~~(Gu et al., 2016a; Gu et al., 2016b)~~. ~~Xiong et al. (2018)~~Xiong et al. (2018) studied variations in Q2DWs activity during ~~the SSWs period~~ and found that W1 ~~is~~was generated by the nonlinear interaction between SPW2 and W3. ~~Gu et al. (2018e)~~Gu et al. (2018b) found that the coupling between the NH and SH enhanced the summer easterly during SSWs and promoted the nonlinear interaction between W3 and SPW1 during ~~the SSWs period~~.

Recent studies have found significant eastward planetary waves ~~activity~~ in the polar stratosphere and mesosphere regions, with periods of

~~nearly 2 and 4 days~~with ~~near 2 and 4 days in the wave periods~~ (Gu et al., 2017; Sandford et al., 2008; Merzlyakov and Pancheva, 2007; Coy et al., 2003; Manney and Randel, 1993)~~(Gu et al., 2017; Venne and Stanford, 2010; Merzlyakov and Pancheva, 2007; Sandford et al., 2007; Coy et al., 2003; Manney et al., 1993)~~. Planetary waves with zonal wavenumbers -1 (E1) and -2 (E2) correspond to 4- and 2-day waves, respectively. In addition, planetary waves of 1.2-day with wavenumber -3 (E3) and 0.8-day with wavenumber -4 (E4) have been found to have the same phase speeds as E1 and E2~~In addition, the planetary waves of 1.2-day with wavenumber -3 (E3) and 0.8-day with wavenumber -4 (E4) have been found same phase speeds as the 2-and 4-day waves~~ (Manney and Randel, 1993)~~(Manney et al., 1993)~~. This series of eastward ~~propagating~~ planetary waves ~~has~~ve a significant effect on the thermal and dynamic structure of the polar stratosphere, resulting in significant changes in the wind and temperature ~~variations in~~of the polar stratosphere (Coy et al., 2003; Venne and Stanford, 1979)~~(Venne and Stanford, 2010; Coy et al., 2003)~~. ~~Palo et al. (1999)~~Palo et al. (1999) demonstrated a series of nonlinear interactions between ~~the~~ migrating tides and Q2DWs. Further research, ~~Palo et al. (2007)~~Palo et al. (2007) presented evidence that ~~the eastward~~E2 Q2DW ~~is~~was coupled by ~~the~~ nonlinear planetary wave and ~~the~~ tides in the mesosphere and lower thermosphere.

~~Merzlyakov & Pancheva. (2007)~~Merzlyakov and Pancheva. (2007)

analyzed and studied satellite datasets and ~~they observed~~found ~~that the~~  
~~wave periods of the eastward propagating wave with the zonal wave~~  
~~numbers 1 and 2 in February 2004, while~~ E1 and E2 events ~~is~~ with wave  
~~periods~~ within 1.5-5 days. They found that ~~the direction of~~ EP fluxes ~~for~~  
~~eastward planetary waves~~ is from the upper layer to the lower atmosphere,  
 suggesting that the upper atmosphere has a dynamic influence on the lower  
 atmosphere. Sandford et al. (2008)~~Sandford et al. (2007)~~ reported ~~a~~  
 significant ~~Q2DW~~ fluctuations of E2 Q2DW in the polar mesosphere  
~~with a zonal wavenumber of 2 (E2).~~ They found that changes~~variations of~~  
in mean zonal winds during a major SSW ~~period~~ can influence the  
propagation of polar E2~~affect the propagation of polar E2~~. In addition, they  
 believe that E2 fluctuation is representative in the mesosphere and caused  
 by the instabilities in the polar night jet. Gu et al. (2017)~~Gu et al. (2017)~~  
~~found~~proposed that the amplitude of E2 can reach  $\sim 10\text{K}$ ,  $\sim 20\text{ m/s}$ , and  $\sim 30$   
 m/s in temperature, zonal wind, and meridional wind in the austral winter  
~~period~~, while the amplitude of E2 decreases by near two-thirds in the boreal  
 winter ~~period~~. ~~Lu et al. (2013)~~ Lu et al. (2013) found that ~~the propagation~~  
~~height of~~ eastward planetary wave propagations was limited to ~~the~~ winter  
 high latitudes, which may be because negative refractive indices  
equatorward of  $\sim 45^\circ\text{S}$  result in evanescent wave characteristic~~caused by~~  
~~the negative refractive index of  $45^\circ\text{S}$  at the equator~~, thus preventing the  
propagation of planetary wave ~~propagation to the~~ low latitudes. They

believe that the instability region at  $\sim 50\text{-}60^\circ\text{S}$  may be induced by the stratospheric polar night jet and/or the "double-jet" structure.

The second modern retrospective research and application analysis (MERRA-2) datasets are used to investigate the eastward ~~propagating~~propagation wave characteristics in the polar stratospheric and mesospheric region during 2019~~in the polar stratospheric and mesospheric region~~, including ~~zonal wavenumbers -1 (E1), -2 (E2), -3 (E3) and -4 (E4)~~. Particularly, ~~we our investigate study is to explore~~ the variation of the occurrence date, peak amplitude, and wave period ~~for of~~ eastward waves~~s~~, ~~and as well as~~ the role of instability, background wind structure, and critical layer ~~for in~~ the propagation and amplification of eastward waves ~~propagation and amplification~~. The remaining parts of this paper are organized as follows. In Section 2, the data and methods used in our study are described. In Section 3 analyzes the global latitude-temporal variation structure of eastward waves in the 2019 winter~~of 2019~~. The amplification and propagation features of the eastward planetary waves ~~different wavenumber events for eastward planetary waves~~ in the NH and SH with different wavenumber events are investigated in Section 3.1 and 3.2, respectively. In Section 3.3 ~~compares and analyzes~~~~will compare and analyze~~ eastward waves in the NH and SH. In Section 4 summarizes our research results.

## 2 Data and Analysis

The least-square method is applied to each time window to extract the E1-, E2-, E3-, and E4-wave, with 10-day, 6-day, 4-day, and 4-day, and is used by us to determine the amplitude (Gu et al., 2013)(~~Gu et al., 2013~~). This method has previously been used successfully to identify planetary waves from satellite measurements (Gu et al., 2019; Gu et al., 2018a; Gu et al., 2018b; Gu et al., 2018c; Gu et al., 2013)(~~Gu et al., 2018a; Gu et al., 2018b; Gu et al., 2018c; Gu et al., 2018d~~).

$$y = A \cos[2\pi(\sigma \cdot t + s \cdot \lambda)] + B \sin[2\pi(\sigma \cdot t + s \cdot \lambda)] + C \quad (1)$$

The least-squares method is used to fit the parameters of (  $A, B$ , and  $C$  ). Where (  $\sigma, t, s$ , \_and  $\lambda$  ) are the frequency, UT time, zonal wavenumber, UT time, and longitudes. The amplitude of wave  $R$  ~~of wavenumber~~ could be expressed as  $R = \sqrt{A^2 + B^2}$ .

The second modern retrospective research and application analysis (MERRA-2) data is a set of long-term atmospheric reanalysis datasets started by NASA in 1980, and now using an upgraded version of the Goddard Earth Observing System Model, Version 5 (GEOS-5) data assimilation system. MERRA-2 includes updates to the model (Molod et al., 2014; Molod et al., 2012)(~~Molod et al., 2015; Molod et al., 2012~~) and the Global Statistical Interpolation (GSI) analysis scheme of Wu et al. (2002)(~~Wu et al., (2002)~~). The MERRA-2 data includes various meteorological variables such as net radiation, temperature, relative humidity, wind speed, etc. The MERRA-2 data covers the world, with a

spatial resolution of  $0.5^{\circ} \times 0.625^{\circ}$  and a temporal resolution of 1 hour. This kind of meteorological data is widely used to detect the middle atmosphere such as the planetary wave in the polar atmosphere, global thermal tides, climate variability, and aerosol (Ukhov et al., 2020; Sun et al., 2020; Bali et al., 2019; Lu et al., 2013)(~~Sun et al., 2020; Ukhov et al., 2020; Bali et al., 2019; Lu et al., 2013~~). These studies indicate that MERRA-2 data can be used in our research with high authenticity. The MERRA-2 datasets are used to obtain the variation in background wind, instability, refractive index, and critical layer, and explore the rules of eastward planetary waves propagation and amplification through diagnostic analysis.

The critical layer will absorb or reflect planetary waves during upward propagation from the lower atmosphere. Planetary waves that gain sufficient energy in the instability region will be amplified during reflection.~~The planetary wave will be amplified from the reflection process after gaining sufficient energy in the instability region.~~ This shows that the critical layer plays an important role in regulating the amplification and propagation of planetary waves (Gu et al., 2016a; Gu et al., 2016b; Liu et al., 2004)(~~Gu et al., 2016a; Gu et al., 2016b; Liu and H. L., 2004~~). ~~The baroclinic/barotropic instability in the atmospheric space structure is caused by the simultaneous equalization of the negative latitude gradient and the quasi-geostrophic potential vorticity ( $\overline{q_{\phi}}$ ). Where ( $\Omega$ ) denote the angular speed of the Earth's rotation, and the latitude and zonal mean zonal~~

~~wind are represented by ( $\varphi$  and  $\bar{u}$ ), in the second part, the ( $a$ ) represents the Earth radius, in the last part, ( $\rho$ ,  $f$ , and  $N$ ) denote the background air density, Coriolis parameter, and buoyancy frequency, respectively. The vertical and latitudinal gradients are represented by subscripts ( $z$  and  $\varphi$ ).~~

$$\bar{q}_\varphi = 2\Omega \cos \varphi - \left( \frac{(\bar{u} \cos \varphi)_\varphi}{a \cos \varphi} \right) - \frac{a}{\rho} \left( \frac{f^2}{N^2} \rho \bar{u}_z \right)_z \quad (2)$$

The baroclinic/barotropic instability in the atmospheric space structure is caused by the simultaneous equalization of the negative latitude gradient and the quasi-geostrophic potential vorticity ( $\bar{q}_\varphi$ ). Where ( $\Omega$ ) denote the angular speed of the Earth's rotation, and the latitude and zonal mean zonal wind are represented by ( $\varphi$  and  $\bar{u}$ ), in the second part, the ( $a$ ) represents the Earth radius, in the last part, ( $\rho$ ,  $f$ , and  $N$ ) denote the background air density, Coriolis parameter, and buoyancy frequency, respectively. The vertical and latitudinal gradients are represented by subscripts ( $z$  and  $\varphi$ ).

~~Andrews et al. (1987)~~ Andrews et al. (1987) define the Eliassen-Palm (EP) flux vectors ( $F$ ) to show the properties of planetary wave propagation, calculated as follows:

$$F = \rho a \cos \varphi \begin{bmatrix} \frac{\overline{u'_z v' \theta'}}{\theta'_z} - \overline{v' u'} \\ \left[ f - \frac{(\bar{u} \cos \varphi)_\varphi}{a \cos \varphi} \right] \frac{\overline{v' \theta'}}{\theta'_z} - \overline{w' u'} \end{bmatrix} \quad (3)$$

~~w~~Where  $u'$  and  $v'$  are the planetary wave perturbations in the zonal and

meridional wind,  $\theta'$  and  $w'$  represent potential temperature and vertical wind, respectively. The planetary wave propagation is only favorable where the square of refractive index  $m^2$  was positive:

$$m^2 = \frac{\overline{q_\phi}}{a(\overline{u} - c)} - \frac{s^2}{(a \cos \phi)^2} - \frac{f^2}{4N^2 H^2} \quad (4)$$

Where  $(s)$  denote the zonal wavenumber, the phase speed is represented by  $(c)$ , and the  $(H)$  represents the scale height. The square of the refractive index is taken as the waveguide of planetary waves.

$$c = -v_0 \cos\left(\frac{\phi\pi}{180}\right) / sT \quad (5)$$

where  $(c)$  denote the phase speed, the equatorial linear velocity is represented by  $(v_0)$ , and the  $(\phi)$  represents the latitude,  $s$  and  $T$  represent zonal wavenumber and wave period, respectively.

### 3 Results and Discussion

The global temporal-latitude variation structures of E1, E2, E3, and E4 extracted from the 2019 MERRA-2 temperature datasets using time windows of 10-, 6-, 4- and 4-days respectively are shown in Figure 1. The mean temperature amplitudes of E1, E2, E3, and E4 at 55.4km during periods of 3~5-, 1.5~2.5-, 1~1.5-, and 0.9~1.1-days are shown in Figure 1a, 1b, 1c, and 1d. Eastward waves are characterized by obvious seasonal variations in the SH and NH. ~~The eastward wave modes during winter periods in the SH and NH are characterized by obvious seasonal variations.~~ In addition, ~~the~~ E1, E2 (E3), and E4 reach their ir maximum amplitudes ss at



50-80°(S/N). In the SH, the strongest E1 and E2 events ~~of E1 and E2~~ occur  
 on days 209-218 and 167-172, while ~~those of~~ E3 and E4 events occur on  
days 151-154 and 139-142 ~~days~~. This shows that the ir occurrence date of  
 maximum amplitude occurs earlier with increasing zonal wavenumber. In  
 addition, the maximum amplitudes of E1, E2, E3, and E4 are ~6.0K, ~4.2K,  
~3.6K, and ~2.4K, indicating that the ir peak amplitudes decrease with  
 increasing zonal wavenumber. In the NH, the strongest E1, E2, E3 and E4  
events occur on days 41-50, 69-74, 35-38 and 63-66, and the corresponding  
peak amplitudes are ~5.5K, ~3.8K, ~2.8K and ~1.2K.~~the occurrence dates~~  
~~of the strongest events of E1, E2, E3, and E4 are days 41-50, 69-74, 35-38,~~  
~~and 63-66, and the corresponding peak amplitudes are 5.5K, 3.8K, 2.8K,~~  
~~and 1.2K. The results show~~It is found that the peak amplitude also  
 decreases with increasing zonal wavenumber in the NH, but the occurrence  
date is irregular~~their occurrence dates vary irregularly~~. In addition, ~~the~~ E4  
 is weak in the NH and almost impossible to find, so W4 is out of the  
 discussion in the NH. Figure 2 shows the changes in zonal mean zonal wind  
at 70°S and 70°N in 2019~~that the 2019 zonal mean zonal wind variations~~  
~~at 70°S and 70°N. It is found that the background wind on days 90-240~~  
(70°S) is dominated by westward wind, and reached ~80 m/s at ~50km on  
days 210, while it is dominated by eastward wind in late and early 2019,  
and reaches ~-40 m/s at ~60km. Meanwhile, the background wind is  
primarily westerly wind in late and early 2019 (70°N), and reaches ~90 m/s

at ~ 60km on days 50, while on days 120-240, the background wind is primarily easterly wind, and the amplitude reaches -40 m/s on days 200. Compared with Figure 1, the results show that the eastward wave modes exist during winter periods with westward background wind in both hemispheres. It is clearly shown that the eastward wave modes exist during winter periods with westward background wind in both hemispheres.

### 3.1 In the Southern Hemisphere

Figure 3 shows that maximum amplitude is observed to be ~48.2km and ~70-80°S for E1, and E2 and E3 at ~48.2km and ~60-70°S, while E4 at ~48.2km and ~50-60°S. The E1, E2 (E3), and E4 are distributed in ~70-80°S, ~60-70°S, and ~50-60°S at ~48.2km, respectively. The maximum perturbation observed for E1 occur on days 211-220 with an amplitude of ~8.5K, and the remaining fluctuations occur on days 161-170, 187-196 and 231-240, while the E2 maximizes at days 219-224 with an amplitude of ~7.8K, and also shows three peaks at days 139-144, 173-178, and 187-192. Besides, the strongest E3 occurs on days 151-154 with an amplitude of ~5.2K, and the rest are distributed on days 141-144, 201-204, 209-202. E4 perturbations are distributed on days 127-130, 145-148, 161-164, 213-216, with weak amplitude of ~3K. According to previous studies, the wave period of the eastward wave varies. Therefore, periodic variability of E1, E2, E3, and E4 has also been investigated. The results show that the period corresponding to the maximum perturbation of E1 falls between ~106

(days 187-196) and  $\sim 69$  h (days 211-220), and their wave periods vary greatly. However, the wave period of E2 gradually changes from  $\sim 42$  h (days 139-144) to  $\sim 38$  h (days 219-224), and its stability is stronger than that of E1. In addition, the wave periods of E3 and E4 are about  $\sim 39$  h and  $\sim 24$  h, respectively. Thus, E2, E3, and E4 wave periods are more stable than E1. ~~E1 events occur on days 161-170, 187-196, 211-220, and 231-240, respectively, and the maximum fluctuation on days 211-220 is  $\sim 8.5$ K, as shown in Figure 3a. In addition, the wave period of E1 decreased from the maximum  $\sim 106$  h (days 187-196) to  $\sim 69$  h (days 211-220), indicating the instability of the E1 wave period. E2 events occur on days 139-144, 173-178, 187-192, and 219-224, and the maximum amplitude of  $\sim 7.8$ K occurs on days 219-224, and the wave period of E2 is stable approximate 40 h (Figure 3b). The strongest E3 event occurs on days 151-154, and the rest are distributed on days 141-144, 201-204, 209-202, and shows E3 wave period approximate 29 h, as shown in Figure 3c. From Figure 3d, E4 events are distributed on days 127-130, 145-148, 161-164, 213-216, with a weak amplitude of about  $\sim 3$ K, while the wave period is stable at near 24 h. In addition, we found that the planetary waves E1, E2, E3, and E4 have similar phase speeds.~~

The spectra, spatial (vertical and latitudinal) structures of temperature, zonal and meridional wind, and diagnostic analysis of E1 are extracted from the corresponding two representative events, as shown in Figure 4.

Figure 4a, 4b shows the least-squares fitting spectra for MERRA-2 temperature on days 187–196, 211–220 at  $\sim 48.2\text{km}$  and  $\sim 70\text{--}80^\circ\text{S}$ , when and where the E1 maximizes. An eastward wavenumber -1 signal with the period of  $\sim 106$  h and  $\sim 69$  h clearly dominates the whole spectrum. The temperature spatial structure corresponding to these E1 is shown in Figure 4c, 4d. The temperature spatial structure of E1 presents an obvious amplitude bimodal structure at  $\sim 70\text{--}80^\circ\text{S}$  and  $\sim 50\text{km}$ , and  $\sim 70\text{--}80^\circ\text{S}$  and  $\sim 60\text{km}$ , with the maximum at  $\sim 70\text{--}80^\circ\text{S}$  and  $\sim 50\text{km}$ . The strongest temperature amplitude of E1 occurs at  $\sim 50\text{km}$  and  $\sim 70\text{--}80^\circ\text{S}$  with an amplitude of  $\sim 10\text{K}$  on the days 211–220, and the other peak is  $\sim 9\text{K}$  ( $\sim 70\text{--}80^\circ\text{S}$  and  $\sim 60\text{km}$ ). The temperature amplitude of  $\sim 9\text{K}$  occurs at  $\sim 50\text{km}$  and  $\sim 70\text{--}80^\circ\text{S}$  during days 187–196, and the rest is  $\sim 7\text{K}$  ( $\sim 70\text{--}80^\circ\text{S}$  and  $\sim 60\text{km}$ ). The corresponding spatial structures of zonal wind and meridional wind of these E1 is shown in Figures 4e, 4f, 4g, and 4h. The maximum zonal wind amplitude of E1 occurs at  $\sim 60\text{--}70^\circ\text{S}$  and  $\sim 60\text{km}$  with an amplitude of  $\sim 14$  m/s on days 187–196, and  $\sim 20$  m/s at  $\sim 50\text{--}60^\circ\text{S}$  and  $\sim 60\text{km}$  on days 211–220. The amplitude of E1 meridional wind reaches  $\sim 10$  m/s at  $\sim 70\text{--}80^\circ\text{S}$  and  $\sim 55\text{km}$  (days 187–196) and  $\sim 17$  m/s at  $\sim 70\text{--}80^\circ\text{S}$  and  $\sim 60\text{km}$  (days 211–220), respectively.

The E1 planetary wave has a wave period of  $\sim 106$  h and  $\sim 69$  h on days 187–196 and 211–220 (Figures 4a and 4b). The temperature spatial structure of E1 events presents an obvious dual structure of amplitudes at

~~~50km and ~60km, as displayed in Figures 4c and 4d. The strongest~~  
~~temperature amplitude of E1 occurs at ~50 km and ~70-80°S with an~~  
~~amplitude of ~10K on the days 211-220, and the other peak is ~9K (60km).~~  
~~The temperature amplitude of ~9K occurs on days 187-196, and the rest is~~  
~~~7K (60km). The spatial structures of zonal wind and meridional wind of~~  
~~E1 are shown in Figures 4e, 4f, 4g, and 4h. The maximum amplitudes of~~  
~~zonal and meridional winds occur at ~60°S and ~60km, and ~80°S and~~  
~~~60km. The amplitude of zonal and meridional wind reaches ~14 m/s and~~  
~~~20 m/s, and ~10m/s and ~17m/s on days 187-196 and 211-220,~~  
~~respectively. Figure 4i, 4j shows the diagnostic analysis results for the E1~~  
~~events during days 187-196, 211-220. It is clear that the EP flux vectors is~~  
~~more favorable to propagate in the SH winter and is dramatically amplified~~  
~~by the mean flow instabilities and appropriate background winds at polar~~  
~~and between ~40 and ~80km, with EP flux propagating into the upper~~  
~~atmosphere (Figure 4i). Meanwhile, there is an EP flux at mid-latitudes and~~  
~~~60-80km, which propagates into the lower atmosphere. The wave-mean~~  
~~flow interactions near its critical layer (106 h) of the green curve amplifies~~  
~~E1, and the positive refractive index region surrounded by the yellow curve~~  
~~also enhances E1 propagation. In addition, the strong instability and weak~~  
~~background wind at ~70-80°S and ~40-60km provide sufficient energy for~~  
~~the upward EP flux to propagate and amplify. However, the downward~~  
~~propagating EP flux is amplified by weak instability and strong~~

background wind at  $\sim 50\text{-}60^\circ\text{S}$  and  $\sim 60\text{-}70\text{km}$ . Besides, both upward and  
downward EP flux eventually propagate toward the equator at  $\sim 50\text{km}$ .  
Figure 4j shows that EP flux on days 211-220 propagates downward and  
amplifies after the interaction of the critical layer ( $\sim 69\text{ h}$ ), facilitated by the  
positive refractive index region, which strong instability and weak  
background wind at  $\sim 50\text{-}60^\circ\text{S}$  and  $\sim 60\text{-}70\text{km}$  provide sufficient energy, and  
ultimately point towards the equator at  $\sim 50\text{km}$ . The results show that the  
weak background wind and strong instability in the polar promote the  
upward propagation and amplification of the EP flux. Meanwhile, the  
appropriate background wind and instability in mid-latitude are also  
conductive to the downward propagation and amplification of EP flux.~~From~~  
~~Figure 4i, E1 EP flux presents two directions of propagation. It is clear that~~  
~~the E1 is more favorable to propagate in the winter hemisphere and is~~  
~~dramatically amplified by the mean flow instabilities and appropriate~~  
~~background winds at polar and middle latitudes between  $\sim 40$  and  $\sim 80\text{ km}$ ,~~  
~~where the former propagate to the upper atmosphere and the latter to the~~  
~~lower atmosphere. In addition, E1 is amplified by wave-mean flow~~  
~~interactions near its critical layer ( $106\text{ h}$ ). The strong instability and weak~~  
~~background wind and positive refractive index region provide sufficient~~  
~~energy for the upward EP flux to propagate and amplify. However, the~~  
~~downward EP flux is propagated and amplified by the interaction of the~~  
~~critical layer in the positive refractive index region, where the strong~~

~~background wind and weak instability provide sufficient energy. In addition, both upward and downward EP fluxes eventually propagate toward the equator at ~50km. Figure 4j shows that EP flux propagates downward and amplifies after the interaction of the critical layer (~69 h), which strong instability and strong background wind provide energy, and ultimately point towards the equator. We believe that the weak background wind and strong instability in the polar region promote the upward propagation and amplification of the EP flux. In addition, the strong background wind and weak instability in the middle latitudes are not conducive to the downward propagation and amplification of the EP flux.~~

In other words, instability and appropriate background wind play a dominant role in the propagation and amplification of the E1.

The spectra are observed at ~48.2km and ~60-70°S on days 173-178, 219-224 for E2, when the eastward wavenumber -2 becomes the primary wave mode with the wave period of ~38 h and ~39 h, as shown in Figure 5a, 5b. The temperature spatial structure corresponding to these E2 is shown in Figure 5c, 5d. The temperature spatial structure of E2 shows an obvious amplitude bimodal structure at ~60-70°S and ~50km, and ~60-70°S and ~60km, with the maximum at ~60-70°S and ~50km. The maximum temperature amplitude of E1 occurs at ~50km and ~60-70°S with an amplitude of ~7.5K on the days 173-178, and the other peak is ~6K (~70°S and ~60km). The temperature amplitude of ~10K occurs at ~50km

and  $\sim 60\text{--}70^\circ\text{S}$  during days 219-224, and the rest is  $\sim 6\text{K}$  ( $\sim 70^\circ\text{S}$  and  $\sim 60\text{km}$ ).  
The corresponding spatial structures of zonal wind and meridional wind of these E2 is shown in Figures 5e, 5f, 5g, and 5h. The zonal wind spatial structure of E2 shows an obvious amplitude bimodal structure at  $\sim 50\text{--}60^\circ\text{S}$  and  $\sim 60\text{km}$ , and  $\sim 70\text{--}80^\circ\text{S}$  and  $\sim 60\text{km}$ , with the maximum at  $\sim 50\text{--}60^\circ\text{S}$  and  $\sim 60\text{km}$ . The maximum zonal wind amplitudes of E2 occur at  $\sim 50\text{--}60^\circ\text{S}$  and  $\sim 60\text{km}$  with an amplitude of  $\sim 10\text{ m/s}$  on days 173-178, and the other peak is  $\sim 9\text{ m/s}$  ( $\sim 70\text{--}80^\circ\text{S}$  and  $\sim 60\text{km}$ ). The zonal wind amplitude of  $\sim 20\text{ m/s}$  occurs at  $\sim 50\text{--}60^\circ\text{S}$  and  $\sim 60\text{km}$  on days 219-224, and the rest is  $\sim 15\text{ m/s}$  ( $\sim 70\text{--}80^\circ\text{S}$  and  $\sim 60\text{ km}$ ). The amplitude of E2 meridional wind reaches  $\sim 13\text{ m/s}$  at  $\sim 70\text{--}80^\circ\text{S}$  and  $\sim 60\text{km}$  (days 173-178) and  $\sim 27\text{ m/s}$  at  $\sim 70\text{--}80^\circ\text{S}$  and  $\sim 60\text{km}$  (days 219-224), respectively.

From Figures 5a and 5b, the wave periods of E2 planetary waves are 38 h and 39 h on days 173-178 and 219-224. The temperature spatial structure of E2 events shows an obvious amplitude dual structure (Figures 5c and 5d). The maximum temperature amplitude of the E2 event on days 219-224 and 173-178 is  $\sim 9\text{K}$  ( $\sim 50\text{km}$ ) and  $\sim 7\text{K}$  ( $\sim 50\text{km}$ ). The spatial structures of zonal wind and meridional wind of E2 are shown in Figures 5e, 5f, 5g, and 5h. The maximum amplitudes of zonal and meridional winds occur at  $\sim 60^\circ\text{S}$  and  $\sim 60\text{km}$ , and  $\sim 80^\circ\text{S}$  and  $\sim 60\text{km}$ . The maximum amplitude of zonal wind reaches  $\sim 10\text{ m/s}$  and  $\sim 20\text{ m/s}$ , while the meridional wind is slightly stronger than the zonal wind, which reaches



~~~13 m/s and ~27 m/s on days 173-178 and 219-224. The results in Figures~~  
~~5i and 5j show diagnostic analysis during days 173-178, 219-224 for E2. It~~  
~~is clear that E2 is more favorable to propagate in the SH winter and is~~  
~~dramatically amplified by the mean flow instabilities at middle-high~~  
~~latitudes between ~40km and ~80km, with EP flux propagating into the~~  
~~lower atmosphere, and EP flux eventually propagate toward the equator at~~  
~~~50km. Besides, E2 is amplified and propagated by the wave-mean flow~~  
~~interactions near its critical layer (~38 h) of the green curve, and the~~  
~~promoting effect of the positive refractive index region surrounded by the~~  
~~yellow curve. Meanwhile, the weak instability and strong background wind~~  
~~at ~50-60°S and ~50-70km provide energy for the propagation and~~  
~~amplification of EP flux into the lower atmosphere during days173-178~~  
~~(Figure 5i). In the diagnostic analysis of days 219-224, it is found that E2~~  
~~obtains sufficient energy from strong instability and strong background~~  
~~wind at ~50-60°S and ~60-70km, and is amplified and propagated into the~~  
~~lower atmosphere through the critical layer and positive refractive index~~  
~~action, as shown in Figure 5j. The results show that the background wind~~  
~~at ~50-60°S and ~50-70km is weaker on days 173-178 than on days 219-~~  
~~224, and the instability at ~50-60°S and ~60-70km is stronger on days 219-~~  
~~224 than on days 173-178. According to Figure 5a, 5b results shows that~~  
~~E2 has absorbed sufficient energy to be amplified under the background~~  
~~conditions during days 219-224. From Figures 5i and 5j, it is clear that the~~

E2 is more favorable to propagate in the SH winter and is dramatically amplified by the mean flow instabilities at middle high latitudes between  $\sim 40\text{km}$  and  $\sim 80\text{km}$ . Then turns to the equator at  $\sim 50\text{km}$ . We find that the background wind in Figure 5j is weaker than that in Figure 5i, but the instability is stronger. This finding indicates that E2 has absorbed sufficient energy to be amplified under the background conditions during days 219–224 (Figure 5j).

Figures 6a and 6b show the observed spectra of E3 at  $\sim 48.2\text{km}$  and  $\sim 60\text{--}70^\circ\text{S}$  on days 151–154 and 201–204, and the wave period of locked wavenumber -3 is  $\sim 29$  h and  $\sim 29$  h, respectively. The corresponding temperature spatial structures of these E3 is shown in Figure 6c, 6d. The temperature spatial structure of E3 shows an obvious amplitude bimodal structure at  $\sim 60\text{--}70^\circ\text{S}$  and  $\sim 50\text{km}$ , and  $\sim 60\text{--}70^\circ\text{S}$  and  $\sim 60\text{km}$ , with the maximum at  $\sim 60\text{--}70^\circ\text{S}$  and  $\sim 50\text{km}$ . Besides, E3 also has a weak peak at  $\sim 60\text{--}70^\circ\text{S}$  and  $\sim 70\text{km}$ . The strongest temperature amplitude of E3 occurs at  $\sim 50\text{km}$  and  $\sim 60\text{--}70^\circ\text{S}$  with an amplitude of  $\sim 6\text{K}$  on the days 151–154, and the other peak is  $\sim 5\text{K}$  ( $\sim 60\text{--}70^\circ\text{S}$  and  $\sim 60\text{km}$ ). The temperature amplitude of  $\sim 5\text{K}$  occurs at  $\sim 50\text{km}$  ( $\sim 60\text{km}$ ) and  $\sim 60\text{--}70^\circ\text{S}$  during days 201–204. The corresponding spatial structures of zonal wind and meridional wind of these E3 is shown in Figures 6e, 6f, 6g, and 6h. The zonal wind spatial structure of E3 shows an obvious amplitude bimodal structure at  $\sim 70\text{--}80^\circ\text{S}$  and  $\sim 60\text{km}$ , and  $\sim 50\text{--}60^\circ\text{S}$  and  $\sim 60\text{km}$ . The zonal wind amplitudes of E3

466 occur at  $\sim 70\text{-}80^\circ\text{S}$  and  $\sim 60\text{km}$  ( $\sim 50\text{-}60^\circ\text{S}$  and  $\sim 60\text{km}$ ) with an amplitude of  
467  $\sim 9\text{ m/s}$  on days 151-154, and  $\sim 9\text{ m/s}$  at  $\sim 70\text{-}80^\circ\text{S}$  and  $\sim 60\text{km}$  ( $\sim 50\text{-}60^\circ\text{S}$   
468 and  $\sim 60\text{km}$ ) on days 201-204. The amplitude of E3 meridional wind  
469 reaches  $\sim 13\text{ m/s}$  at  $\sim 60\text{-}70^\circ\text{S}$  and  $\sim 55\text{km}$  (days 151-154) and  $\sim 16\text{ m/s}$  at  
470  $\sim 60\text{-}70^\circ\text{S}$  and  $\sim 55\text{km}$  (days 201-204), respectively.

471 EP flux of E3 is similar to that of E2. The instability and appropriate  
472 background wind at mid-high latitudes between  $\sim 50$  and  $\sim 70\text{km}$   
473 dramatically amplify the propagation of E3, which is enhanced by the  
474 interaction near the critical layer ( $\sim 29\text{ h}$ ) and the positive refractive index  
475 region. (Figures 6i and 6j). It is worth noting that the strong instability and  
476 weak background wind at  $\sim 50\text{-}60^\circ\text{S}$  and  $\sim 60\text{-}70\text{km}$  on days 151-154  
477 provide sufficient energy for the propagation and amplification of EP flux  
478 into the lower atmosphere, and ultimately point towards the equator at  
479  $50\text{km}$ . The EP flux propagate to the lower atmosphere during days 201-  
480 204, and is amplified by interaction at the critical layer ( $\sim 29\text{ h}$ ). In addition,  
481 weak instability and weak background wind at  $\sim 50\text{-}60^\circ\text{S}$  and  $\sim 60\text{-}70\text{km}$   
482 provide the energy to amplify E3 propagation. Combine with Figure 6c, 6d,  
483 the stronger the instability at  $\sim 50\text{-}60^\circ\text{S}$  and  $\sim 60\text{-}70\text{km}$ , the stronger the  
484 temperature amplitude of E3. We believe that the background wind and  
485 instability at  $\sim 50\text{-}60^\circ\text{S}$  and  $\sim 60\text{-}70\text{km}$  are the primary reasons for the  
486 propagation and amplification of EP flux into the lower atmosphere.

487 The spectra are observed to be at  $\sim 48.2\text{ km}$  and  $\sim 50\text{-}60^\circ\text{S}$  on days

127-130, 213-216 for E4, when the eastward wavenumber -4 signal with  
 the wave period of ~25 h and ~21 h, as shown in Figure 7a, 7b. The  
 corresponding temperature spatial structures of these E4 is shown in Figure  
 7c, 7d. The temperature spatial structure of E4 shows an obvious amplitude  
 bimodal structure at ~50-60°S and ~50km, and ~50-60°S and ~60km, with  
 the maximum at ~50-60°S and ~50km. The maximum temperature  
 amplitude of E4 occurs at ~50km and ~50-60°S with an amplitude of ~4K  
 on the days 127-130, and the other peak is ~3K (~60-70°S and ~60km).  
 The temperature amplitude of ~3K occurs at ~50km (~60km) during days  
 213-216. The corresponding spatial structures of zonal wind and  
 meridional wind of these E4 is shown in Figures 7e, 7f, 7g, and 7h. The  
 zonal wind spatial structure of E4 shows an obvious amplitude bimodal  
 structure at ~50-60°S and ~55km, and ~60-70°S and ~55km, with the  
 maximum at ~50-60°S and ~55km. The maximum zonal wind amplitude  
 of E4 occurs at ~50-60°S and ~55km with an amplitude of ~9 m/s on days  
 127-130, and the other peak is ~5K (~60-70°S and ~55km). The zonal wind  
 amplitude of ~5m/s occurs at ~50-60°S (~60-70°S) and ~55km on days  
 213-216. The amplitude of E4 meridional wind reaches ~8 m/s at ~60-70°S  
 and ~55km (days 127-130) and ~10 m/s at ~60-70°S and ~55km (days 213-  
 216), respectively. ~~Figures 6a and 6b show that the E3 planetary waves~~  
~~have wave periods of ~29 h on days 151-154 and 201-204. The temperature~~  
~~spatial structure of E3 events presents a dual structure, as shown in Figures~~

6e and 6d. The maximum temperature amplitude of E3 is  $\sim 7\text{K}$  ( $\sim 50\text{km}$ ) on days 151-154, and the other peak is  $\sim 5\text{K}$  ( $\sim 60\text{km}$ ). The temperature amplitude of E3 is  $\sim 5\text{K}$  ( $\sim 50$ ,  $\sim 60\text{km}$ ) on days 201-204. The spatial structures of zonal wind and meridional wind of E3 are shown in Figures 6e, 6f, 6g, and 6h. The maximum amplitudes of zonal and meridional winds occur at  $\sim 50^\circ\text{S}$  and  $\sim 60\text{km}$ , and  $\sim 60^\circ\text{S}$  and  $\sim 60\text{km}$ . The maximum zonal wind is  $\sim 10\text{ m/s}$ , while the meridional wind is slightly stronger than the zonal wind at  $\sim 15\text{ m/s}$ . The EP flux of E3 is similar to that of E2. We find that instability at mid-high latitudes between  $\sim 50$  and  $\sim 70\text{km}$  dramatically amplifies the E3 propagation and that the interaction near the critical layer enhances the process (Figures 6i and 6j). It is worth noting that the region of background wind at  $\sim 50$ - $60^\circ\text{S}$  and  $\sim 60$ - $70\text{km}$  is similar during days 151-154 and 201-204, while the former is strong instability and the latter is weak instability. This finding indicates that the strong instability provides sufficient energy for the amplification for E3 propagation on days 151-154.

Figures 7i and 7j show diagnostic analysis on days 127-130, 213-216 for E4. The results show that E4 is dramatically amplified by the mean flow instabilities at middle-high latitudes between  $\sim 50\text{km}$  and  $\sim 70\text{km}$ , with EP flux propagating into the lower atmosphere, and EP flux eventually propagate toward the equator at  $\sim 50\text{km}$ . E4 is amplified and propagated by the wave-mean flow interaction near the critical layer ( $\sim 25\text{ h}$ ,  $\sim 21\text{ h}$ ), and the positive refractive index region provides the promoting effect. The

strong instability and weak background wind at  $\sim 50\text{-}60^\circ\text{S}$  and  $\sim 60\text{-}70\text{km}$  provide sufficient energy for the propagation and amplification of EP flux into the lower atmosphere during days 127-130. Besides, E4 obtains energy from weak instability and weak background wind at  $\sim 50\text{-}60^\circ\text{S}$  and  $\sim 60\text{-}70\text{km}$  on days 213-216, and is amplified and propagated into the lower atmosphere. The background wind at  $\sim 50\text{-}60^\circ\text{S}$  and  $\sim 60\text{-}70\text{km}$  on days 127-130 is similar to on days 213-216, and the instability at  $\sim 50\text{-}60^\circ\text{S}$  and  $\sim 60\text{-}70\text{km}$  is stronger on days 127-30 than on days 213-216. According to Figure 7a, 7b results shows that E4 has absorbed sufficient energy to be amplified under the background conditions on days 127-130. The temperature amplitude on 127-130 days is stronger.

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The wave period of E4 reaches  $\sim 25\text{ h}$  and  $\sim 21\text{h}$  during days 127-130 and 213-216 in Figures 7a and 7b. The maximum temperature amplitude of E4 occurs on days 127-130, reaching  $\sim 4\text{K}$  ( $\sim 50\text{km}$ ), the other peak is  $\sim 3\text{K}$  ( $\sim 60\text{km}$ ). The temperature amplitude is  $\sim 3\text{K}$  ( $\sim 50, \sim 60\text{ km}$ ) on days 213-216, as displayed in Figure 7d. The spatial structures of zonal wind and meridional wind of E4 are shown in Figures 7e, 7f, 7g, and 7h. The maximum amplitudes of zonal and meridional winds occur at  $\sim 50^\circ\text{S}$  and  $\sim 60\text{km}$ , and  $\sim 60^\circ\text{S}$  and  $\sim 60\text{km}$ . The maximum zonal wind is  $\sim 8\text{ m/s}$ , while the meridional wind is slightly stronger than the zonal wind at  $\sim 10\text{ m/s}$ . We find that the instability in the mid-high latitudes between  $\sim 50$  and  $\sim 70\text{km}$

and the interaction near the critical layer greatly enhance the propagation and amplification of E4 EP flux, as shown in Figures 7i and 7j. The weak background wind and strong instability appear on days 127-130, while the strong background wind and weak instability appear on days 213-216. This finding indicates that E4 is difficult to obtain sufficient energy to be amplified under background conditions during days 213-216. The amplitude on 127-130 days is stronger.

### 3.2 In the Northern Hemisphere

Figure 8 shows that maximum amplitude is observed to be  $\sim 59.2$  km and  $\sim 70-80^\circ\text{N}$  for E1, and the E2 and E3 peaks at  $\sim 59.2$  km and  $\sim 60-70^\circ\text{N}$ . The maximum perturbation of E1 is observed to be at days 41-50 with an amplitude of  $\sim 8\text{K}$ , and the remaining fluctuation occur on days 25-34, 339-348. Besides, the strongest E2 occurs on days 69-74 with an amplitude of  $\sim 7\text{K}$ , and the rest are distributed on days 25-30, 317-322, and 341-346, while the E3 maximizes at days 35-38 with an amplitude of  $\sim 3\text{K}$ , and also shows one peaks at days 53-56. Based on the study of the wave period in the SH for eastward wave, the periodic variability in the NH for E1, E2, and E3 is also investigated. The wave periods of E1 decreased from a maximum of  $\sim 118$  h (days 25-34) to  $\sim 80$  h (days 41-50), indicating that the wave period of E1 is unstable in the NH. The E2 events occur on days 25-30, 69-74, 317-322, and 341-346, of which the corresponding wave periods are  $\sim 36$ ,  $\sim 53$ ,  $\sim 52$ , and  $\sim 48$  h, which are stronger stable than E1. Besides,

the wave period of E3 is relatively stable at  $\sim 29$  h and  $\sim 27$  h. Thus, E2 and E3 wave periods are more stable than E1.

Figure 8 shows that the E1 and E2 (E3) are distributed at  $\sim 70$ – $80^\circ\text{N}$ , and  $\sim 60$ – $70^\circ\text{N}$  at  $\sim 59.2\text{km}$ , respectively. E1 events occur on days 25–34, 41–50, and 339–348 respectively, in which the maximum temperature amplitude of days 41–50 reaches  $\sim 8\text{K}$ , as shown in Figure 8a. In addition, the wave periods of E1 decreased from a maximum of  $\sim 118$  h (days 25–34) to  $\sim 80$  h (days 41–50), indicating that the wave period of E1 is unstable in the NH. Clearly, the E2 events occur on days 25–30, 69–74, 317–322, and 341–346, of which the corresponding wave periods are  $\sim 36$ ,  $\sim 53$ ,  $\sim 52$ , and  $\sim 48$  h, and the maximum temperature amplitude reaches  $\sim 7\text{K}$  on days 69–74, as shown in Figure 8b. Figure 8c shows two E3 events. The strongest temperature amplitude of E3 occurs on days 35–38 and reaches  $\sim 3\text{K}$ , and the other one occurs on days 53–56. The wave period of E3 is relatively stable at  $\sim 29$  h and  $\sim 27$  h. We did not study the E4 event in the NH, because E4 is weak.

The spectra, spatial (vertical and latitudinal) structures of temperature, zonal and meridional wind, and diagnostic analysis of E1 are extracted from the corresponding representative events, as shown in Figure 9. Figures 9a and 9b show the observed spectra of E1 at  $\sim 59.2\text{km}$  and  $\sim 70$ – $80^\circ\text{N}$  on days 25–34 and 41–50, and the wave period of locked wavenumber -1 is  $\sim 118$  h and  $\sim 80$  h, respectively. The corresponding temperature spatial



structures of these E1 is shown in Figure 9c, 9d. The temperature spatial  
 structure of E1 shows an obvious amplitude bimodal structure during days  
 25-34 at  $\sim 60\text{-}70^\circ\text{N}$  and  $\sim 60\text{km}$ , and  $\sim 40\text{-}50^\circ\text{N}$  and  $\sim 70\text{km}$ , with the  
 maximum at  $\sim 60\text{-}70^\circ\text{N}$  and  $\sim 60\text{km}$ . Besides, E1 also has bimodal structure  
 on days 41-50 at  $\sim 60\text{-}70^\circ\text{N}$  and  $\sim 60\text{km}$ , and  $\sim 60\text{-}70^\circ\text{N}$  and  $\sim 70\text{km}$ . The  
 strongest temperature amplitude of E1 occurs at  $\sim 60\text{-}70^\circ\text{N}$  and  $\sim 60\text{km}$  with  
 an amplitude of  $\sim 7\text{K}$  on the days 25-34, and the other peak is  $\sim 4\text{K}$  ( $\sim 40\text{-}$   
 $50^\circ\text{N}$  and  $\sim 70\text{km}$ ). The temperature amplitude of  $\sim 10\text{K}$  occurs at  $\sim 60\text{km}$   
 and  $\sim 60\text{-}70^\circ\text{N}$  during days 41-50, and the rest is  $\sim 8\text{K}$  ( $\sim 60\text{-}70^\circ\text{N}$  and  
 $\sim 70\text{km}$ ). The corresponding spatial structures of zonal wind and meridional  
 wind of these E1 is shown in Figures 9e, 9f, 9g, and 9h. The zonal wind  
 spatial structure of E1 shows an obvious amplitude bimodal structure at  
 $\sim 70\text{-}80^\circ\text{N}$  and  $\sim 70\text{km}$ , and  $\sim 50\text{-}60^\circ\text{N}$  and  $\sim 70\text{km}$ . The zonal wind  
 amplitude of  $\sim 13\text{ m/s}$  occurs at  $\sim 70\text{-}80^\circ\text{N}$  and  $\sim 70\text{km}$  on days 25-34, and  
 the rest is  $\sim 10\text{ m/s}$  ( $\sim 50\text{-}60^\circ\text{N}$  and  $\sim 70\text{ km}$ ). In addition, there is a weak  
 peak of  $9\text{K}$  during days 25-34 ( $\sim 30\text{-}40^\circ\text{N}$  and  $\sim 70\text{ km}$ ). The maximum  
 zonal wind amplitude of E1 occurs at  $\sim 70\text{-}80^\circ\text{N}$  and  $\sim 70\text{km}$  with an  
 amplitude of  $\sim 19\text{ m/s}$  on days 41-50, and the other peak is  $\sim 13\text{K}$  ( $\sim 50\text{-}60^\circ\text{N}$   
 and  $\sim 70\text{km}$ ). The amplitude of E1 meridional wind reaches  $\sim 14\text{ m/s}$  at  $\sim 70\text{-}$   
 $80^\circ\text{N}$  and  $\sim 70\text{km}$  (days 25-34) and  $\sim 22\text{ m/s}$  at  $\sim 70\text{-}80^\circ\text{N}$  and  $\sim 70\text{km}$  (days  
 41-50), respectively. From Figures 9a and 9b, the wave periods of planetary  
 E1 waves on days 25-34 and 41-50 are  $\sim 118\text{ h}$  and  $\sim 80\text{ h}$ , respectively. The

620 ~~temperature spatial structure of the E1 event presents an obvious dual~~  
621 ~~amplitude structure ( $\sim 60\text{km}$ ,  $\sim 70\text{km}$ ). From Figure 9d, the strongest~~  
622 ~~temperature amplitude of E1 appears at  $\sim 60\text{km}$  on days 41-50, with an~~  
623 ~~amplitude of  $\sim 10\text{K}$ , and the rest peak is  $\sim 8\text{K}$  ( $\sim 70\text{km}$ ). The temperature~~  
624 ~~amplitude of  $\sim 7\text{K}$  appears on days 25-34, and the rest is  $\sim 4\text{K}$  ( $\sim 70\text{km}$ ), as~~  
625 ~~shown in Figure 9c. The spatial structures of zonal wind and meridional~~  
626 ~~wind of E1 are shown in Figures 9e, 9f, 9g, and 9h. The maximum~~  
627 ~~amplitudes of zonal and meridional winds occur at  $\sim 50^\circ\text{N}$  and  $\sim 70\text{km}$ , and~~  
628  ~~$\sim 80^\circ\text{N}$  and  $\sim 70\text{km}$ . The maximum zonal wind is  $\sim 14\text{ m/s}$ , while the~~  
629 ~~meridional wind is slightly stronger than the zonal wind at  $\sim 18\text{ m/s}$ .~~

630 The diagnostic analysis results for E1 in Figures 9i and 9j show that  
631 E1 is dramatically amplified by the mean flow instabilities at middle-high  
632 latitudes between  $\sim 50\text{km}$  and  $\sim 70\text{km}$ , with EP flux propagating into the  
633 polar lower atmosphere, and EP flux eventually propagate toward the  
634 equator at  $\sim 50\text{km}$ . The wave-mean flow interaction near the critical layers  
635 ( $\sim 118\text{ h}$ ,  $\sim 80\text{ h}$ ) amplifies and propagates E1, and the promoting effect of  
636 the positive refractive index region amplifies E1. Besides, the weak  
637 instability and strong background wind at  $\sim 40\text{-}50^\circ\text{N}$  and  $\sim 60\text{-}70\text{km}$  provide  
638 energy for the propagation and amplification of EP flux into the polar lower  
639 atmosphere during days 25-34. The E1 obtains sufficient energy from weak  
640 instability and suitable background wind on days 41-50 at  $\sim 40\text{-}50^\circ\text{N}$  and  
641  $\sim 60\text{-}70\text{km}$ , and is amplified and propagated into the polar lower

atmosphere through the critical layer and positive refractive index action. The background wind at  $\sim 40\text{-}50^\circ\text{N}$  and  $\sim 60\text{-}70\text{km}$  is stronger on days 25-34 than on days 41-50, but their instability is similar, indicating that stronger background winds may be unfavorable to E1 propagation and amplification at mid-northern latitudes. According to Figure 9a, 9b results shows that E1 has absorbed sufficient energy to be amplified under the background conditions during days 41-50, showing a stronger temperature amplitude. ~~The results show that E1 is more conducive to the propagation in winter in the NH, and the instability at mid-latitude between  $\sim 60$  and  $\sim 80$  km and the interaction near the critical layer dramatically amplify the propagation of E1, which eventually turns toward the equator in the Arctic (Figures 9i and 9j). Obviously, the instability on days 25-34 and 41-50 is relatively weak, but the former has stronger background winds. This finding indicates that the background condition on days 41-50 is conducive to the propagation and amplification for E1.~~

The spectra are observed to be at  $\sim 59.2$  km and  $\sim 60\text{-}70^\circ\text{N}$  on days 25-30, 69-74 for E2, when the eastward wavenumber -2 signal with the period of  $\sim 36$  h and  $\sim 53$  h, as shown in Figure 10a, 10b. The corresponding temperature spatial structures of these E2 is shown in Figure 10c, 10d. The temperature spatial structure of E2 shows an obvious amplitude bimodal structure at  $\sim 60\text{-}70^\circ\text{N}$  and  $\sim 60\text{km}$ , and  $\sim 60\text{-}70^\circ\text{N}$  and  $\sim 70\text{km}$ , with the maximum at  $\sim 60\text{-}70^\circ\text{N}$  and  $\sim 60\text{km}$ . The maximum temperature amplitude

of E2 occurs at  $\sim 60\text{-}70^\circ\text{N}$  and  $\sim 60\text{km}$  with an amplitude of  $\sim 5\text{K}$  on days 25-30, and the other peak is  $\sim 4\text{K}$  ( $\sim 60\text{-}70^\circ\text{N}$  and  $\sim 70\text{km}$ ). The temperature amplitude of  $\sim 9\text{K}$  occurs on days 69-74 at  $\sim 60^\circ\text{S}$  and  $\sim 60\text{km}$ , and the other peaks are  $\sim 7\text{K}$  ( $\sim 60\text{-}70^\circ\text{N}$  and  $\sim 70\text{km}$ ),  $\sim 5\text{K}$  ( $\sim 60\text{-}70^\circ\text{N}$  and  $\sim 50\text{km}$ ). The corresponding spatial structures of zonal wind and meridional wind of these E2 is shown in Figures 10e, 10f, 10g, and 10h. The zonal wind spatial structure of E2 shows an obvious amplitude bimodal structure at  $\sim 60\text{-}70^\circ\text{N}$  and  $\sim 60\text{km}$ , and  $\sim 40\text{-}50^\circ\text{N}$  and  $\sim 60\text{km}$ , with the maximum at  $\sim 40\text{-}50^\circ\text{N}$  and  $\sim 60\text{km}$ . The maximum zonal wind amplitude of E2 occurs at  $\sim 60\text{-}70^\circ\text{N}$  and  $\sim 60\text{km}$  ( $\sim 40\text{-}50^\circ\text{N}$  and  $\sim 60\text{km}$ ) with an amplitude of  $\sim 6\text{ m/s}$  on days 25-30. Zonal wind amplitude occurs at  $\sim 40\text{-}50^\circ\text{N}$  and  $\sim 60\text{km}$  with an amplitude of  $\sim 18\text{ m/s}$  on days 41-50, and the other peak is  $\sim 16\text{K}$  ( $\sim 60\text{-}70^\circ\text{N}$  and  $\sim 60\text{km}$ ). The amplitude of E2 meridional wind reaches  $\sim 7\text{ m/s}$  at  $\sim 60\text{-}70^\circ\text{N}$  and  $\sim 70\text{km}$  (days 25-30) and  $\sim 18\text{ m/s}$  at  $\sim 60\text{-}70^\circ\text{N}$  and  $\sim 60\text{km}$  (days 41-50), respectively.

Figures 10i and 10j show diagnostic analysis on days 25-30, 69-74 for E2. It is clear that E2 is dramatically amplified by the mean flow instabilities at middle-high latitudes between  $\sim 40\text{km}$  and  $\sim 70\text{km}$ , with EP flux propagating into the polar lower atmosphere, and EP flux eventually propagate toward the equator at  $\sim 50\text{km}$ . E2 is amplified and propagated by the wave-mean flow interaction near the critical layers ( $\sim 36\text{ h}$ ,  $\sim 53\text{ h}$ ), and the positive refractive index region provides the promoting effect. The

weak instability and strong background wind at  $\sim 50\text{-}60^\circ\text{N}$  and  $\sim 60\text{-}70\text{km}$  provide energy for the propagation and amplification of EP flux into the polar lower atmosphere during days 25-30. Besides, E2 obtains sufficient energy from strong instability and suitable background wind at  $\sim 50\text{-}60^\circ\text{N}$  and  $\sim 60\text{-}70\text{km}$  on days 69-74, and is amplified and propagated into the polar lower atmosphere. The background wind at  $\sim 50\text{-}60^\circ\text{N}$  and  $\sim 60\text{-}70\text{km}$  on days 127-130 is similar to on days 213-216, and the instability at  $\sim 50\text{-}60^\circ\text{S}$  and  $\sim 60\text{-}70\text{km}$  is stronger on days 127-30 than on days 213-216. The background wind at  $\sim 50\text{-}60^\circ\text{N}$  and  $\sim 60\text{-}70\text{km}$  is stronger on days 25-30 than on days 69-74, but the instability at  $\sim 50\text{-}60^\circ\text{N}$  and  $\sim 60\text{-}70\text{km}$  is stronger on days 69-74 than on days 25-30. According to Figure 10a, 10b temperature amplitude results shows that E2 has absorbed sufficient energy to be amplified under the background conditions on days 69-74, with the temperature amplitude on days 69-74 is stronger.

Figures 11a and 11b show the observed spectra of E3 at  $\sim 59.2\text{km}$  and  $\sim 60\text{-}70^\circ\text{N}$  on days 35-38 and 53-56, and the wave period of locked wavenumber -3 is  $\sim 29\text{h}$  and  $\sim 27\text{h}$ , respectively. The corresponding temperature spatial structures of these E3 is shown in Figure 11c, 11d. The temperature spatial structure of E3 shows an obvious amplitude bimodal structure at  $\sim 50\text{-}60^\circ\text{N}$  and  $\sim 60\text{km}$ , and  $\sim 50\text{-}60^\circ\text{N}$  and  $\sim 70\text{km}$ , with the maximum at  $\sim 50\text{-}60^\circ\text{N}$  and  $\sim 60\text{km}$ . The strongest temperature amplitude

of E3 occurs at  $\sim 60\text{km}$  and  $\sim 50\text{-}60^\circ\text{N}$  with an amplitude of  $\sim 6\text{K}$  on the days 35-38, and the other peak is  $\sim 5\text{K}$  ( $\sim 50\text{-}60^\circ\text{N}$  and  $\sim 70\text{km}$ ). The temperature amplitude of  $\sim 4\text{K}$  occurs at  $\sim 60\text{km}$  ( $\sim 70\text{km}$ ) and  $\sim 50\text{-}60^\circ\text{N}$  during days 53-56. The corresponding spatial structures of zonal wind and meridional wind of these E3 is shown in Figures 6e, 6f, 6g, and 6h. The zonal wind spatial structure of E3 shows an obvious amplitude bimodal structure at  $\sim 40\text{-}50^\circ\text{N}$  and  $\sim 70\text{km}$ , and  $\sim 60\text{-}70^\circ\text{N}$  and  $\sim 70\text{km}$ . The zonal wind amplitudes of E3 occur at  $\sim 40\text{-}50^\circ\text{N}$  and  $\sim 70\text{km}$  with an amplitude of  $\sim 15\text{ m/s}$  on days 35-38, and  $\sim 12\text{ m/s}$  at  $\sim 60\text{-}70^\circ\text{N}$  and  $\sim 70\text{km}$ . The maximum zonal wind amplitude of E3 occurs at  $\sim 40\text{-}50^\circ\text{N}$  and  $\sim 70\text{km}$  ( $\sim 60\text{-}70^\circ\text{N}$  and  $\sim 70\text{km}$ ) with an amplitude of  $\sim 7\text{ m/s}$  ( $\sim 6\text{ m/s}$ ) on days 53-56. The amplitude of E3 meridional wind reaches  $\sim 22\text{ m/s}$  at  $\sim 50\text{-}60^\circ\text{N}$  and  $\sim 70\text{km}$  (days 35-38) and  $\sim 12\text{ m/s}$  at  $\sim 60\text{-}70^\circ\text{N}$  and  $\sim 70\text{km}$  (days 53-56), respectively.

Obviously, the instability and appropriate background wind at mid-latitude between  $\sim 50$  and  $\sim 70\text{ km}$  and the interaction near the critical layers ( $\sim 29\text{ h}$ ,  $\sim 27\text{ h}$ ) dramatically amplify the propagation of E3, as shown in Figures 11i and 11j. The background wind is similar on days 35-38 and 53-56, and the former is more unstable. This finding indicates that the E3 in propagation is more likely to get sufficient energy to be amplified on days 35-38. The instability and appropriate background wind at mid-high latitudes between  $\sim 50$  and  $\sim 70\text{km}$  dramatically amplify the propagation of E3, which is enhanced by the interaction near the critical layers ( $\sim 29\text{ h}$ ,  $\sim 27$

h) and the positive refractive index region. (Figure 11i, 11j). It is worth noting that the strong instability and weak background wind at  $\sim 50\text{-}60^\circ\text{N}$  and  $\sim 60\text{-}70\text{km}$  on days 35-38 provide sufficient energy for the propagation and amplification of EP flux into the lower atmosphere, and ultimately point towards the equator at 50km. The EP flux propagate to the lower atmosphere during days 35-38, and is amplified by interactions at the critical layer ( $\sim 29$  h). In addition, weak instability and weak background winds on days 53-56 at  $\sim 50\text{-}60^\circ\text{N}$  and  $\sim 60\text{-}70\text{km}$  provide the energy to amplify E3 propagation. Combine with Figure 11c, 11d, the stronger the instability at  $\sim 50\text{-}60^\circ\text{N}$  and  $\sim 60\text{-}70\text{km}$ , the stronger the temperature amplitude of E3. The results show that the instability on days 35-38 at  $\sim 50\text{-}60^\circ\text{N}$  and  $\sim 60\text{-}70\text{km}$  are the primary reasons for the propagation and amplification of EP flux into the lower atmosphere. From Figures 10a and 10b, the wave period of E2 on days 25-30 and 69-74 reaches  $\sim 36\text{h}$  and  $\sim 53\text{h}$ . The maximum temperature amplitude of E2 appears on days 69-74, reaches  $\sim 10\text{K}$  ( $\sim 60\text{km}$ ), and the other peak is  $\sim 7\text{K}$  ( $\sim 70\text{km}$ ). The maximum temperature amplitude is  $\sim 5\text{K}$  ( $\sim 60\text{km}$ ) during days 25-30, and another peak is  $\sim 4\text{K}$  ( $\sim 70\text{km}$ ). The spatial structures of zonal wind and meridional wind of E2 are shown in Figures 10e, 10f, 10g, and 10h. The maximum amplitudes of zonal and meridional winds occur at  $\sim 50^\circ\text{N}$  and  $\sim 70\text{km}$ , and  $\sim 60^\circ\text{N}$  and  $\sim 70\text{km}$ . The maximum zonal and meridional winds are  $\sim 7$  ( $\sim 18$ ) m/s on days 35-30 (69-74). Clearly, the instability at a mid-high latitude

between  $\sim 40$  and  $\sim 80$  km and the interaction near the critical layer dramatically amplify the propagation of E2, which eventually turns toward the equator, as displayed in Figures 10i and 10j. Obviously, the temperature amplitude is stronger on days 69–74, which indicates that E2 obtains sufficient energy for amplification under the background condition.

The E3 planetary waves have wave periods of  $\sim 29$ h and  $\sim 27$ h on days 35–38 and 53–56, as shown in Figures 11a, and 11b. Clearly, from Figures 11c, and 11d, the maximum temperature amplitude of E3 is  $\sim 6$ K ( $\sim 70$ km) on days 35–38, and another peak is  $\sim 5$ K ( $\sim 60$ km). The peak temperature amplitude is  $\sim 4$ K ( $\sim 60$ ,  $\sim 70$  km) on days 53–56. The spatial structures of zonal wind and meridional wind of E3 are shown in Figures 11e, 11f, 11g, and 11h. The maximum amplitudes of zonal and meridional winds occur at  $\sim 40^\circ$ N and  $\sim 70$ km, and  $\sim 50^\circ$ N and  $\sim 70$ km. The maximum zonal and meridional winds are  $\sim 15$  ( $\sim 7$ ) and  $\sim 22$  ( $\sim 13$ ) m/s on days 35–38 (53–56). Obviously, the instability at mid-latitude between  $\sim 60$  and  $\sim 70$  km and the interaction near the critical layer dramatically amplify the propagation of E3, as shown in Figures 11i and 11j. The background wind is similar on days 35–38 and 53–56, and the former is more unstable. This finding indicates that the E3 in propagation is more likely to get sufficient energy to be amplified on days 35–38.

### 3.3 Comparison between SH and NH

~~We find that t~~The observed latitude and maximum amplitude for



eastward planetary waves (E1, E2, E3, E4) decrease and weaken with increasing zonal wavenumber in the SH, reaching  $\sim 70\text{--}80^\circ\text{S}$ ,  $\sim 60\text{--}70^\circ\text{S}$ ,  $\sim 60\text{--}70^\circ\text{S}$ , and  $\sim 50\text{--}60^\circ\text{S}$ , and  $\sim 10\text{K}$ ,  $\sim 9\text{K}$ ,  $\sim 6\text{K}$ , and  $\sim 3\text{K}$ , respectively. In addition, the occurrence date is earlier with increasing zonal wavenumber. The temperature spatial structure shows a bimodal~~dual~~-peak structure ( $\sim 50$  and  $\sim 60\text{km}$ ), primarily located at  $\sim 50\text{km}$ . The maximum zonal wind amplitudes of E1 and E2, E3 and E4 are almost the equivalents, which are  $\sim 20\text{ m/s}$  and  $\sim 10\text{ m/s}$  respectively. The maximum meridional wind amplitudes of E1, E2, E3 and E4 are  $\sim 17\text{ m/s}$ ,  $\sim 27\text{ m/s}$ ,  $\sim 16\text{ m/s}$ , and  $\sim 11\text{ m/s}$  respectively. The wave period of E1 tends to get shorter from 5 to 3 days, while E2 and E3 are close to  $\sim 40\text{ h}$  and  $\sim 30\text{ h}$ , while E4 remains at  $\sim 24\text{ h}$ . E1, E2, E3, and E4 are more favorable to propagation in the SH winter and are dramatically amplified by the mean flow instabilities at middle latitudes between  $\sim 40$  and  $\sim 70\text{km}$ , with EP flux propagating into the lower atmosphere, and EP flux eventually propagate toward the equator at  $\sim 50\text{km}$ . In addition, the propagation of EP flux for E1 to the upper atmosphere~~E1 upward propagating EP flux~~ may be influenced by the instability and background wind at the Antarctic  $\sim 50\text{km}$ .

The observed latitudes of E1, E2 (E3) decrease with increasing wavenumber in the NH, which are  $\sim 70\text{--}80^\circ\text{N}$ ,  $\sim 60\text{--}70^\circ\text{N}$ , and  $\sim 60\text{--}70^\circ\text{N}$ . The temperature spatial structure of E1, E2, and E3 presents a bimodal~~dual~~-peak structure, primarily located at  $\sim 70\text{km}$ , reaches  $\sim 10\text{K}$ ,  $\sim 9\text{K}$ , and  $\sim 6\text{K}$ .

796 The maximum zonal wind amplitude for E1, E2 and E3 occurs at  $\sim 50-80^{\circ}\text{N}$   
 797 and  $\sim 70\text{km}$ , and their amplitude is almost equal to  $\sim 18\text{ m/s}$ .~~The maximum~~  
 798 ~~zonal wind amplitude appears at  $\sim 50-80^{\circ}\text{N}$  and  $\sim 60\text{km}$ . E1, E2, E3, and E4~~  
 799 ~~are almost the equivalent, which is  $\sim 18\text{ m/s}$  respectively.~~ The maximum  
 800 meridional wind of E1, E2 and E3 occur at  $\sim 50-80^{\circ}\text{N}$  and  $\sim 70\text{km}$ , with  
 801 amplitudes of  $\sim 22\text{ m/s}$ ,  $\sim 18\text{ m/s}$  and  $\sim 22\text{ m/s}$ , respectively.~~The maximum~~  
 802 ~~meridional wind amplitude appears at  $\sim 50-80^{\circ}\text{N}$  and  $\sim 60\text{km}$ . The~~  
 803 ~~maximum amplitudes of E1, E2, and E3 are  $\sim 22\text{ m/s}$ ,  $\sim 18\text{ m/s}$ , and  $\sim 22\text{ m/s}$~~   
 804 ~~respectively.~~ The wave period of E1 tends to be shorter from 5-3 days, and  
 805 E2 and E3 are close to  $\sim 48\text{ h}$  and  $\sim 30\text{ h}$ . In addition, E1, E2, and E3 are  
 806 more favorable to propagation in the NH winter and are dramatically  
 807 amplified by the mean flow instabilities at middle latitudes between  $\sim 40$   
 808 and  $\sim 70\text{-km}$ , with EP flux propagating into the lower atmosphere, and EP  
 809 flux eventually propagate toward the equator at  $\sim 50\text{km}$ .

#### 810 **4 Summary and Conclusions**

811 We present for the first time an extensive study of the global variation  
 812 for eastward planetary wave activity, including zonal wave numbers of -1  
 813 (E1), -2 (E2), -3 (E3), -4 (E4), in the stratosphere and mesosphere using  
 814 the MERRA-2 temperature and wind observations in 2019. The  
 815 temperature and wind amplitudes and wave periods of each event were  
 816 determined by 2-D least-squares fitting. Our study includes the spatial and  
 817 temporal behaviors of the eastward planetary waves in both hemispheres

with a comprehensive diagnostic analysis on their propagation and amplification. The key findings of the study are summarized as follows:

1. The latitudes for the maximum (temperature, zonal and meridional wind) amplitudes of E1, E2, E3, and E4 decrease with increasing wavenumber in the SH and NH. The E1, E2, E3, and E4 events occur earlier with increasing zonal wavenumber in the SH. In addition, eastward wave modes exist during summer periods with westward background wind in both hemispheres.

2. The temperature spatial structures of E1, E2, E3, and E4 present a double-peak structure, which is located at ~50km and ~60km in SH, ~60km, and ~70km in SH. In addition, the lower peak is usually larger than the higher one.

3. The maximum (temperature, zonal and meridional wind) amplitudes of E1, E2, and E3 decrease with increasing zonal wavenumber in the SH and NH. The maximum temperature amplitudes in the SH are slightly larger and lie lower than those in the NH. In addition, the meridional wind amplitudes are slightly larger than the zonal wind in the SH and NH.

4. The wave period of the E1 mode varies from 3 to 5 days in both hemispheres, while the period of E2 mode is slightly longer in the NH ( $\sim$ 48 h) than in the SH ( $\sim$ 40 h). The periods of E3 are  $\sim$ 30 h in both SH and NH, and the period of E4 is  $\sim$ 24 h.

5. The eastward planetary wave is more favorable to propagate in the

winter hemisphere and is dramatically amplified by the mean flow instabilities and appropriate background winds at polar and middle latitudes between  $\sim 40$  and  $\sim 80$ -km. Furthermore, the amplification of planetary waves through wave-mean flow interaction most easily occurs near its critical layer. In addition, the direction of EP flux ultimately points towards the equator.

6. The strong instability and appropriate background wind in the lower layer of the Antarctic region may provide sufficient energy to promote the E1 propagation and amplification to the upper atmosphere~~layer~~.

This study demonstrates how the background zonal wind in the polar middle atmosphere affects the dynamics of eastward planetary waves in the polar middle atmosphere.

853 *Data availability.* MERRA-2 data are available at <http://disc.gsfc.nasa.gov>.

854

855 *Author contributions.* LT carried out the data processing and analysis and  
856 wrote the manuscript. SYG and XKD contributed to reviewing the article.

857

858 *Competing interests.* The authors declare that they have no conflict of  
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860

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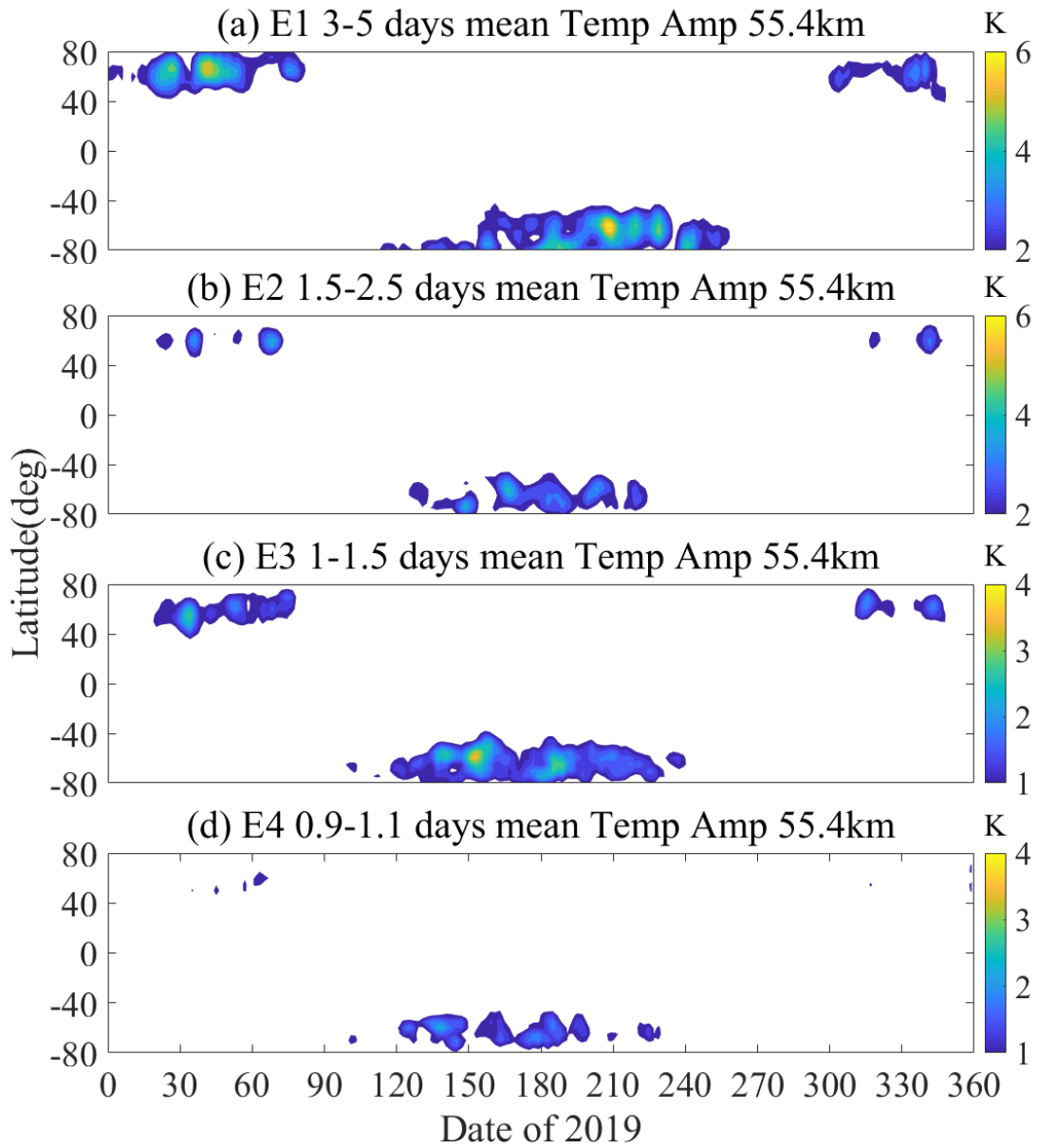
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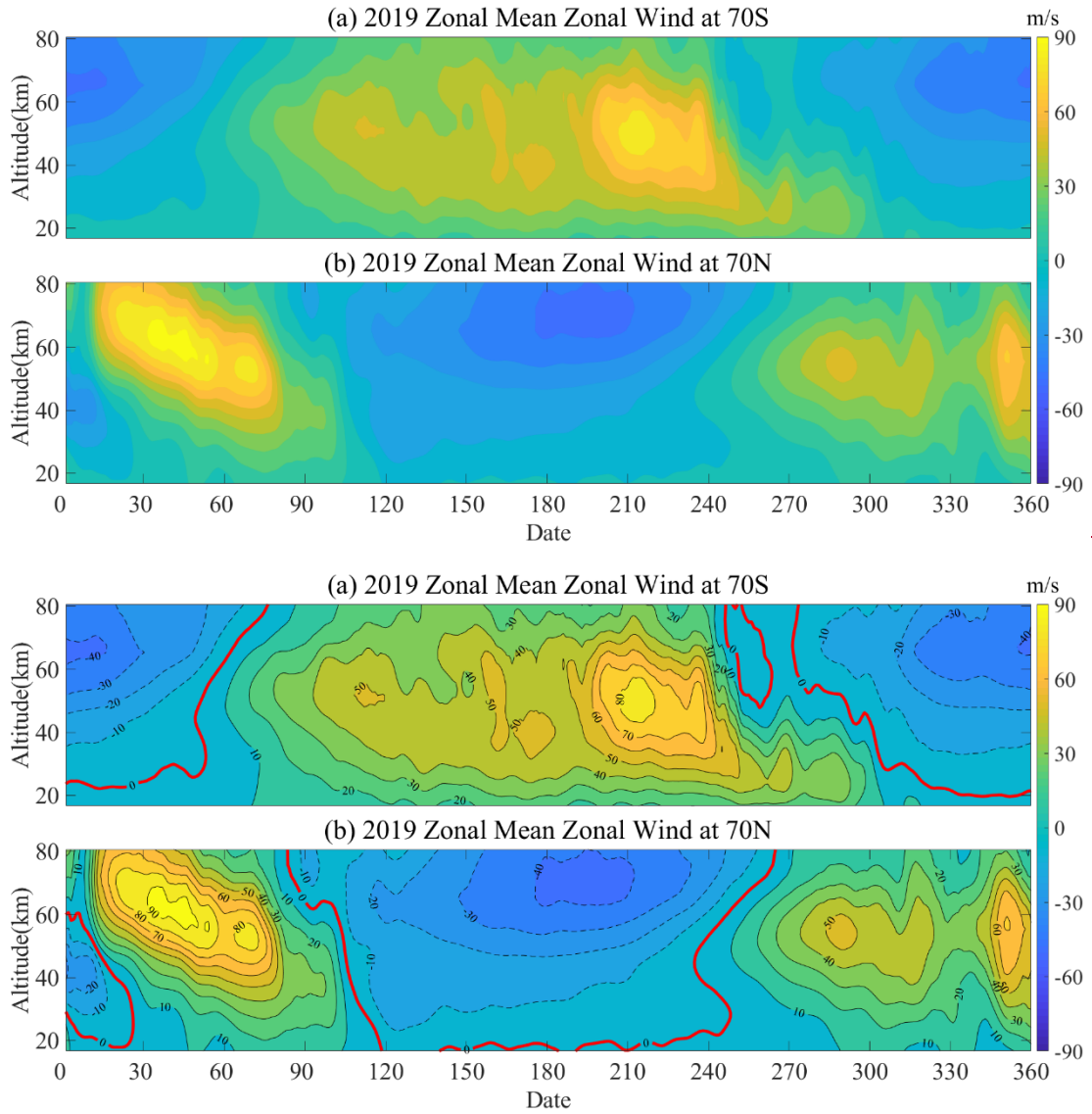
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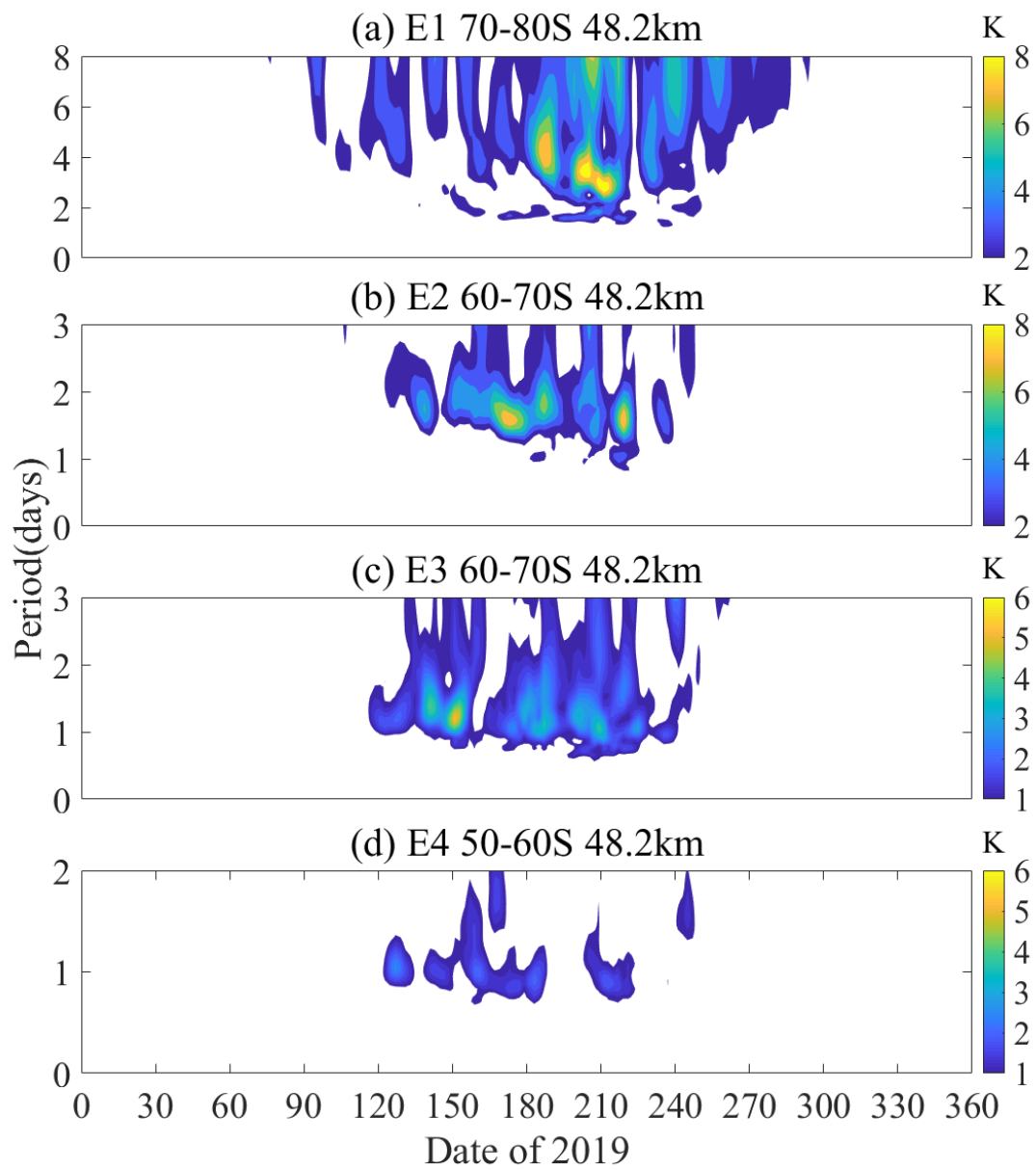
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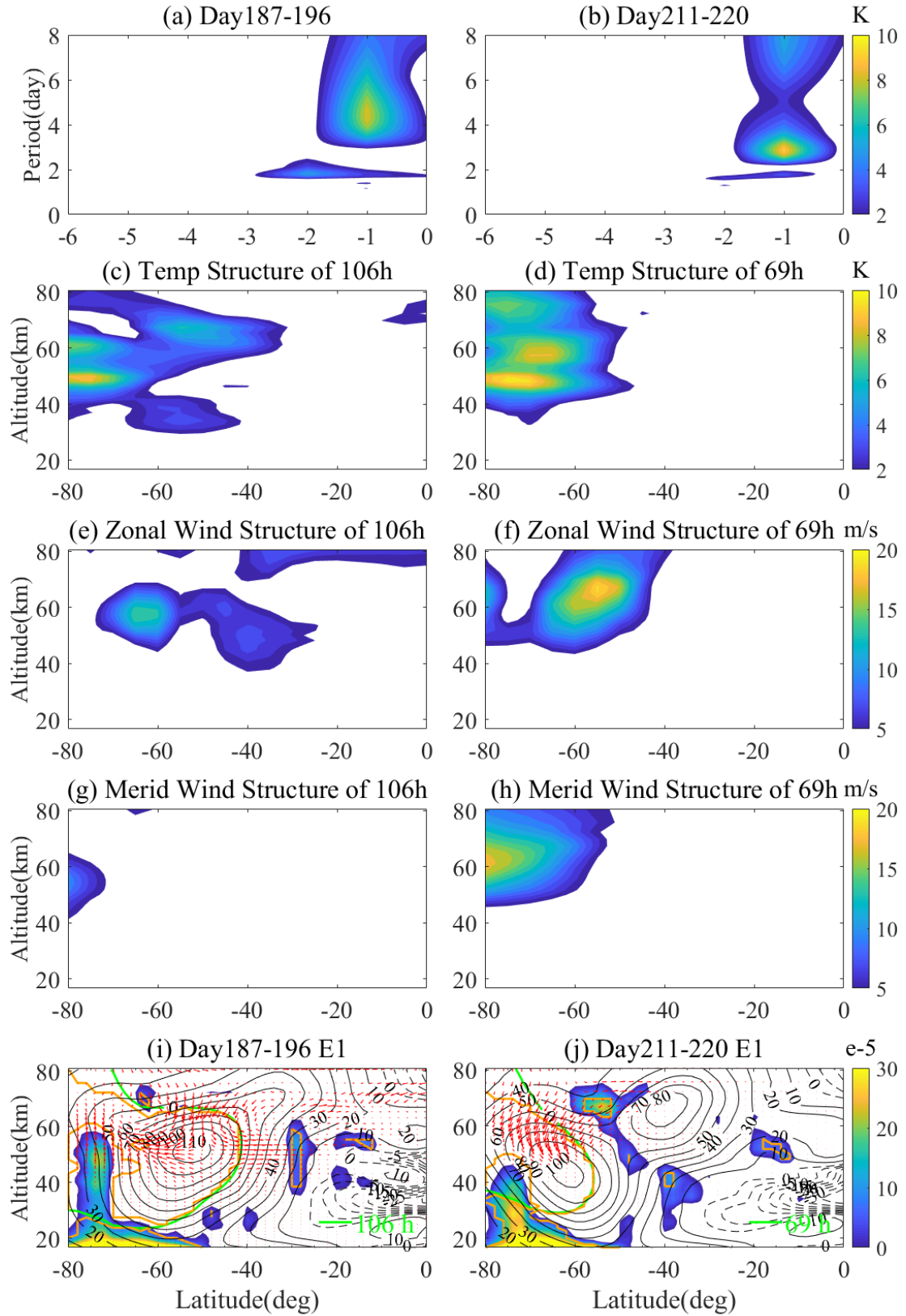
**Figure 1.** The global latitude-temporal variation structures of the (a) E1, (b) E2, (c) E3, (d) E4 planetary waves during 2019. White areas represent small amplitude data (corresponds to the right color bar). The confidence interval is 0.95.



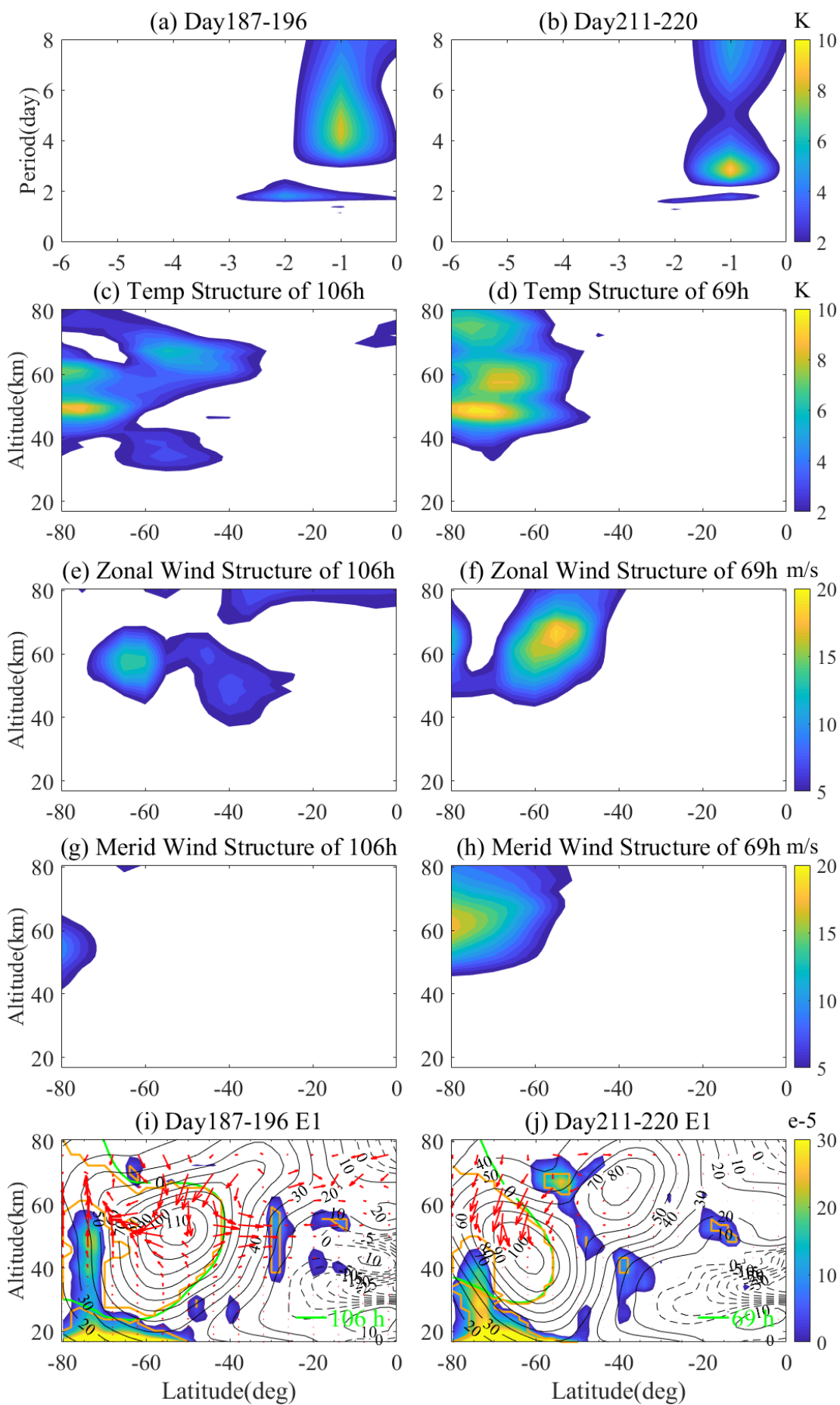
**Figure 2.** The zonal mean zonal wind variations of the (a) the 70S and (b) 70N during 2019. The dotted line represents east wind, the solid line represents west wind, and the red solid line is 0m/s.



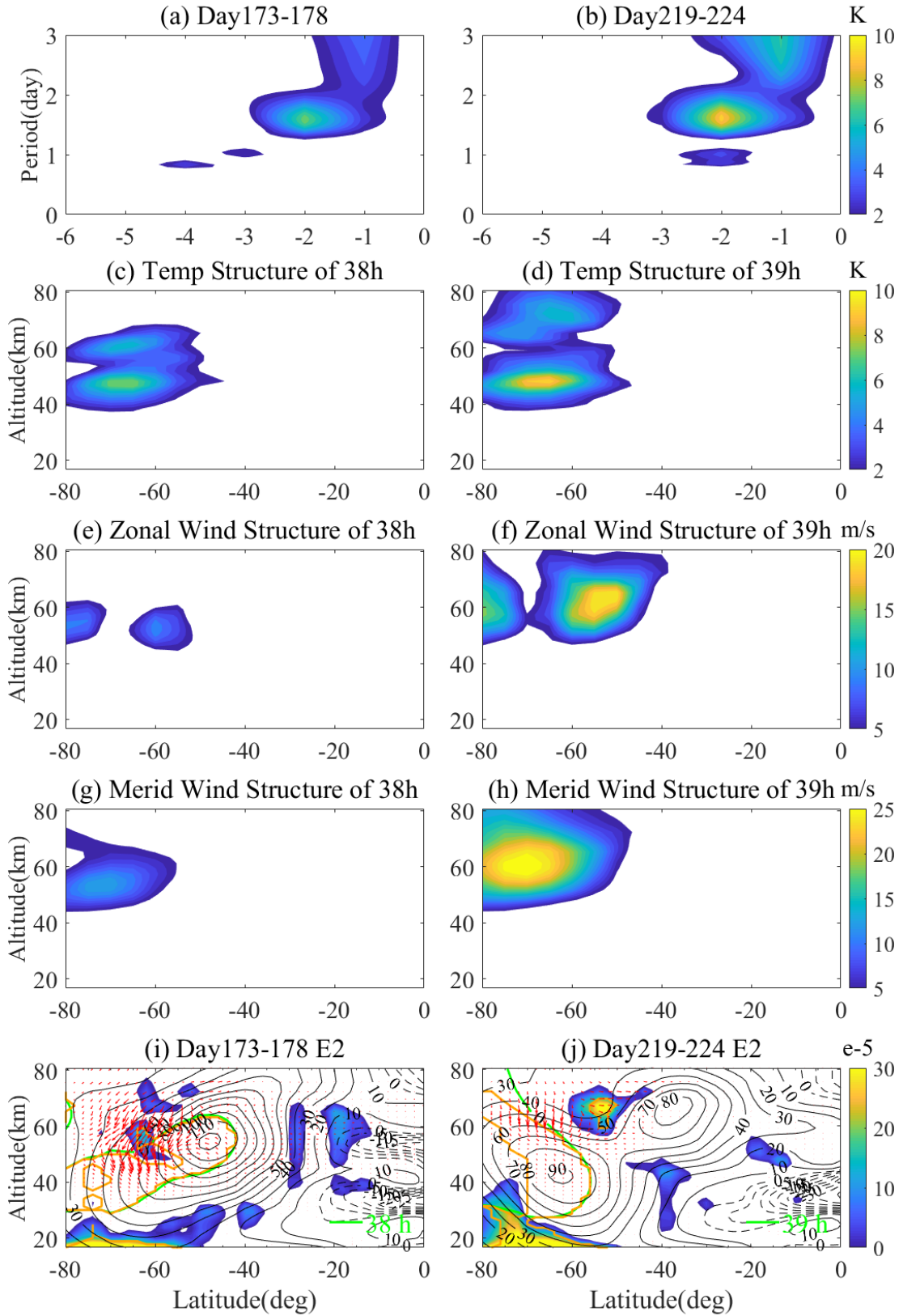
**Figure 3.** The temporal variations of the (a) E1, (b) E2, (c) E3, (d) E4 QTDWs during 2019 austral winter period.

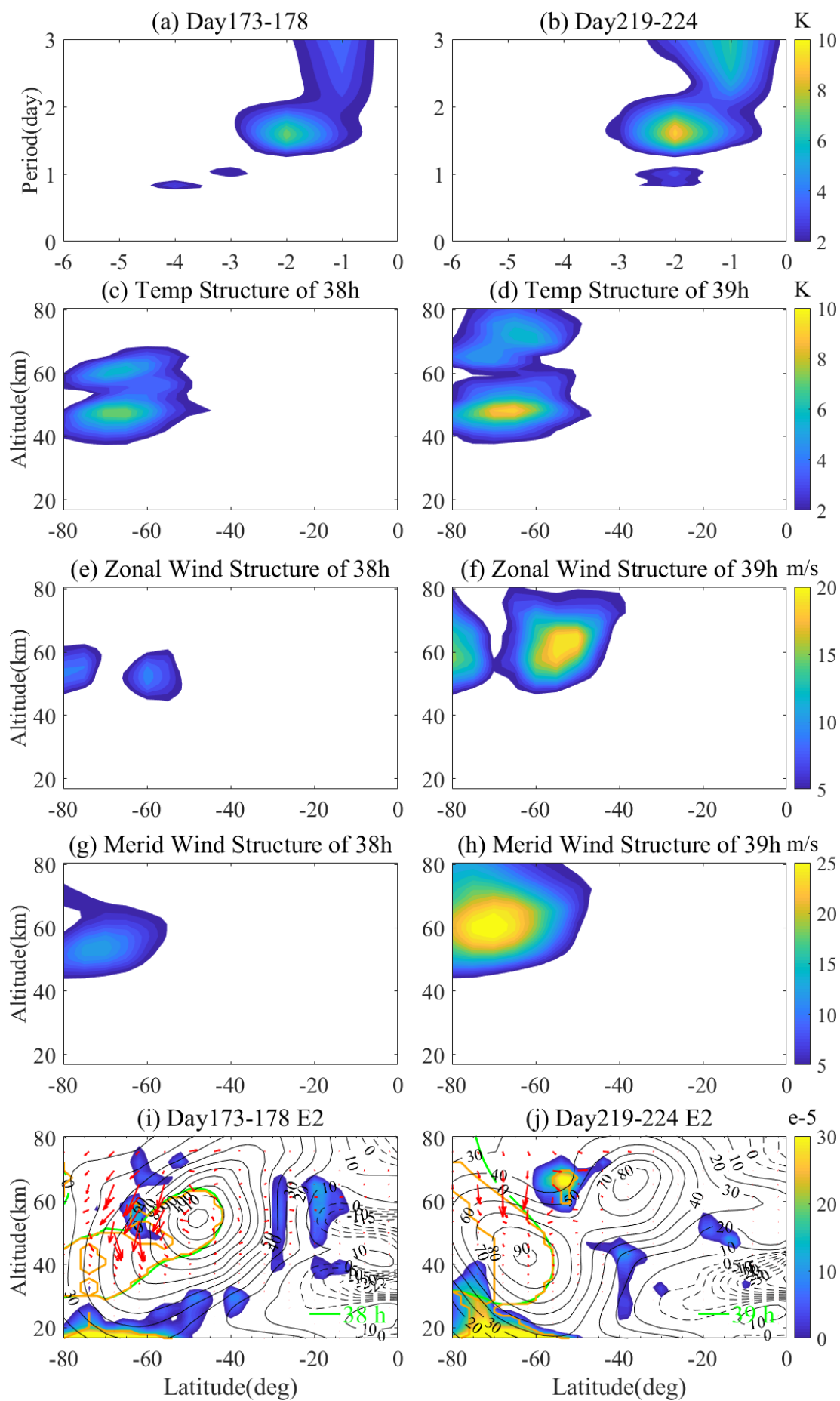




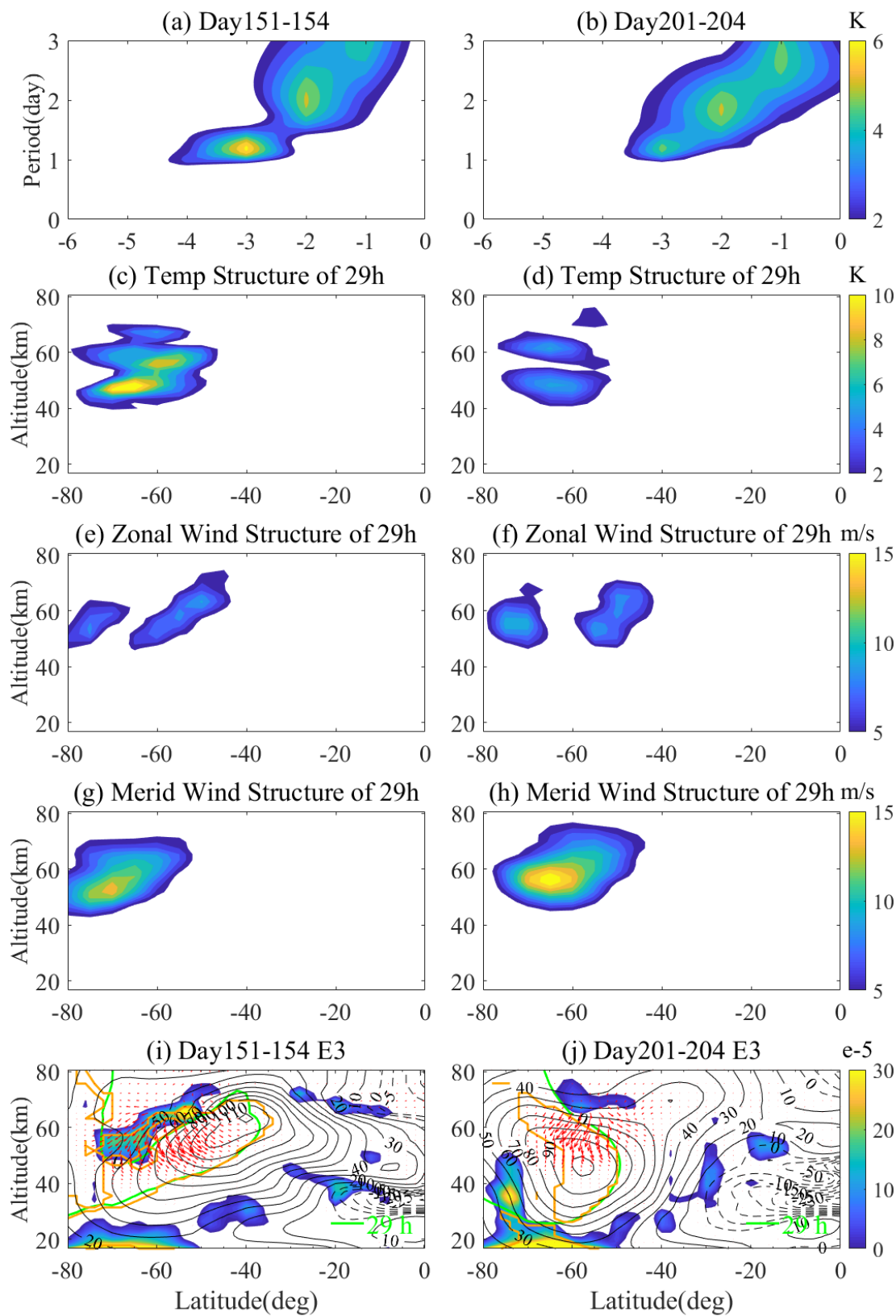


**Figure 4.** The (a, b) spectra, (c, d) temperature spatial structures, (e, f) zonal wind spatial structures, (g, h) meridional wind spatial structures, and (i, j) diagnostic analysis of the E1 typical events during 2019 austral winter period. The MERRA-2 temperature data observations at 48.2km and 70-80°S during days 187–196 (Figure 4a), 211–220 (Figure 4d) are utilized, respectively. The instability (blue shaded region), EP fluxes (red arrow), and critical layers (green line) for E1 typical event. The green line represents critical layers of the E1 with the natural period. Regions enclosed by orange solid lines are characterized by the positive refractive index for the E1.

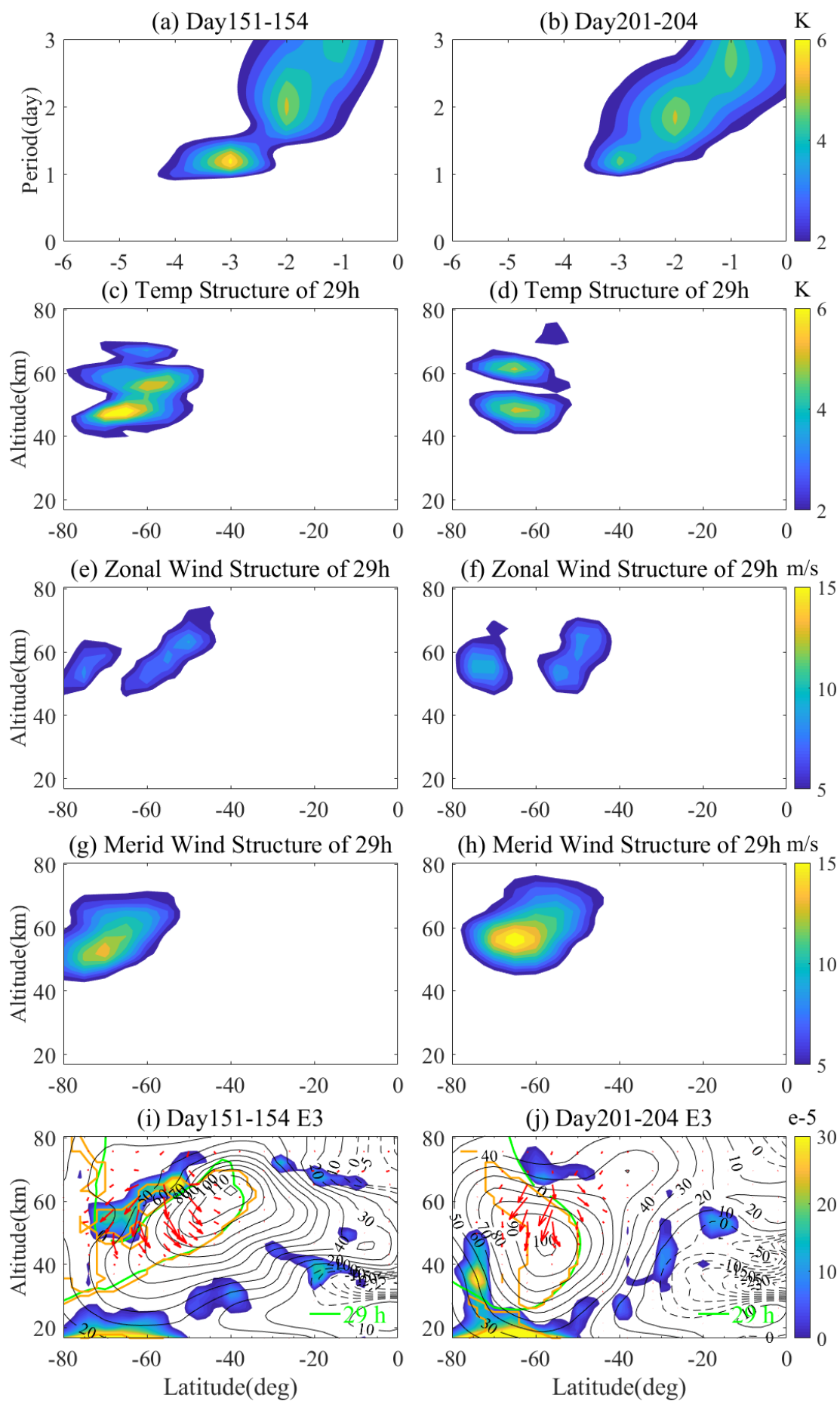




1033 **Figure 5.** The same as Figure 4 but for the E2 during the 2019 austral winter period.

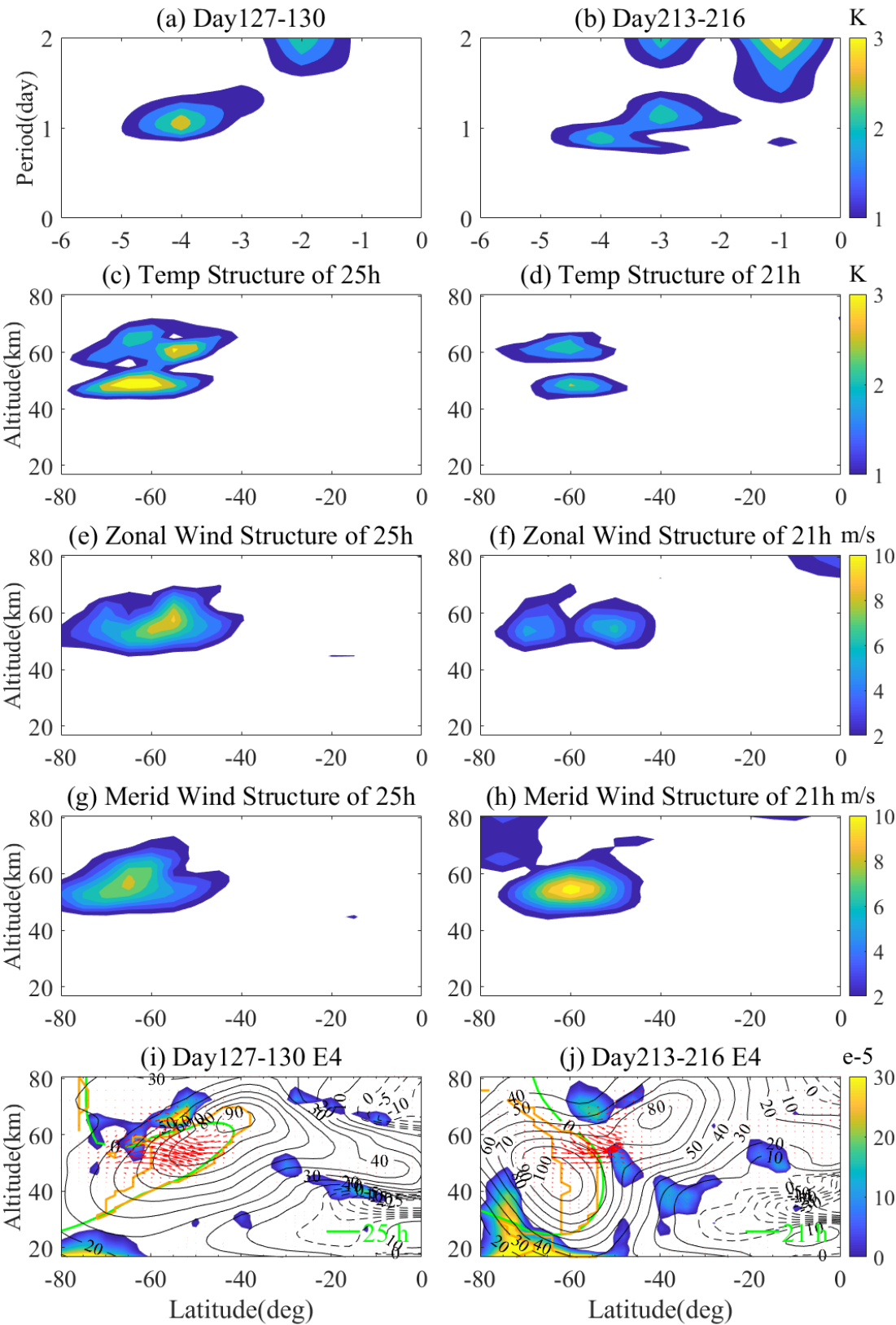


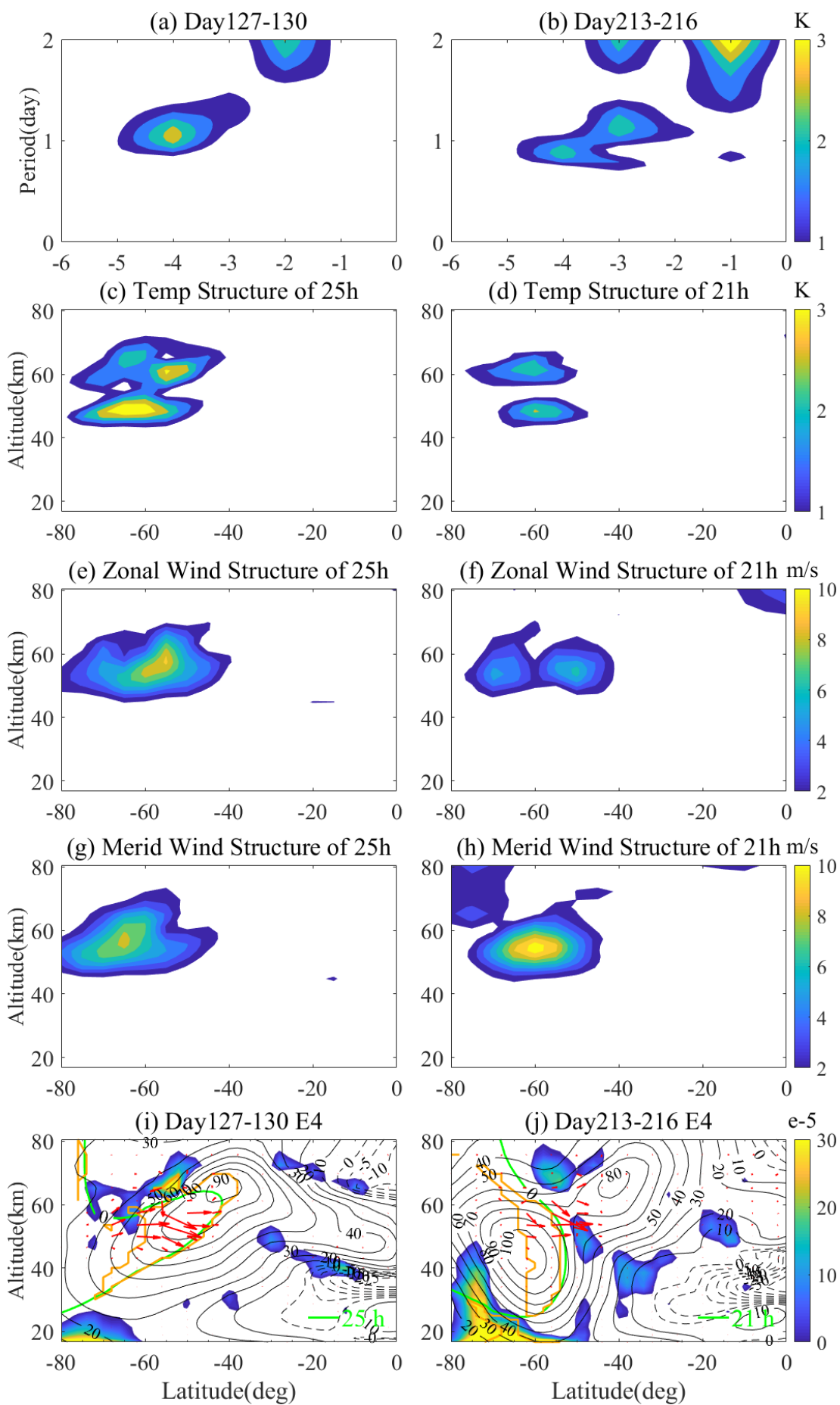
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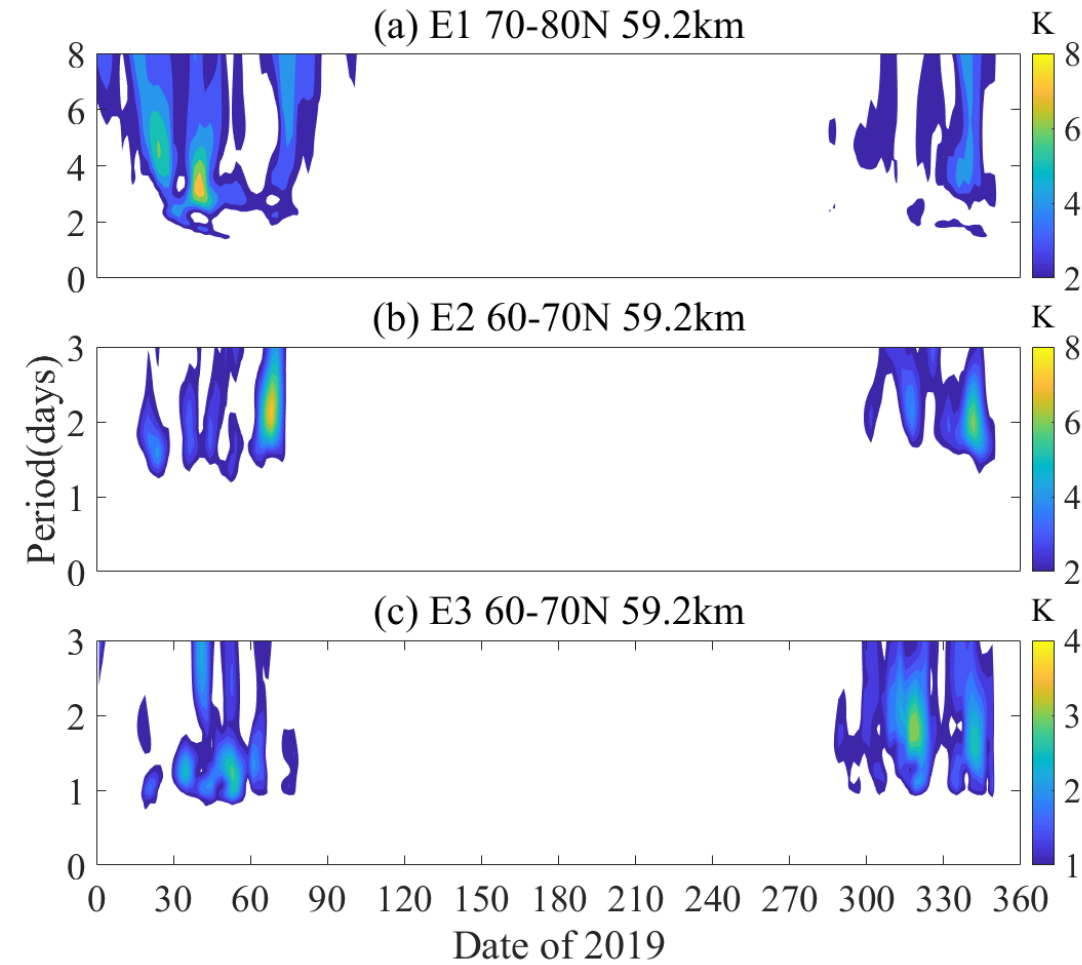
1036 **Figure 6.** The same as Figure 4 but for the E3 during the 2019 austral winter period.



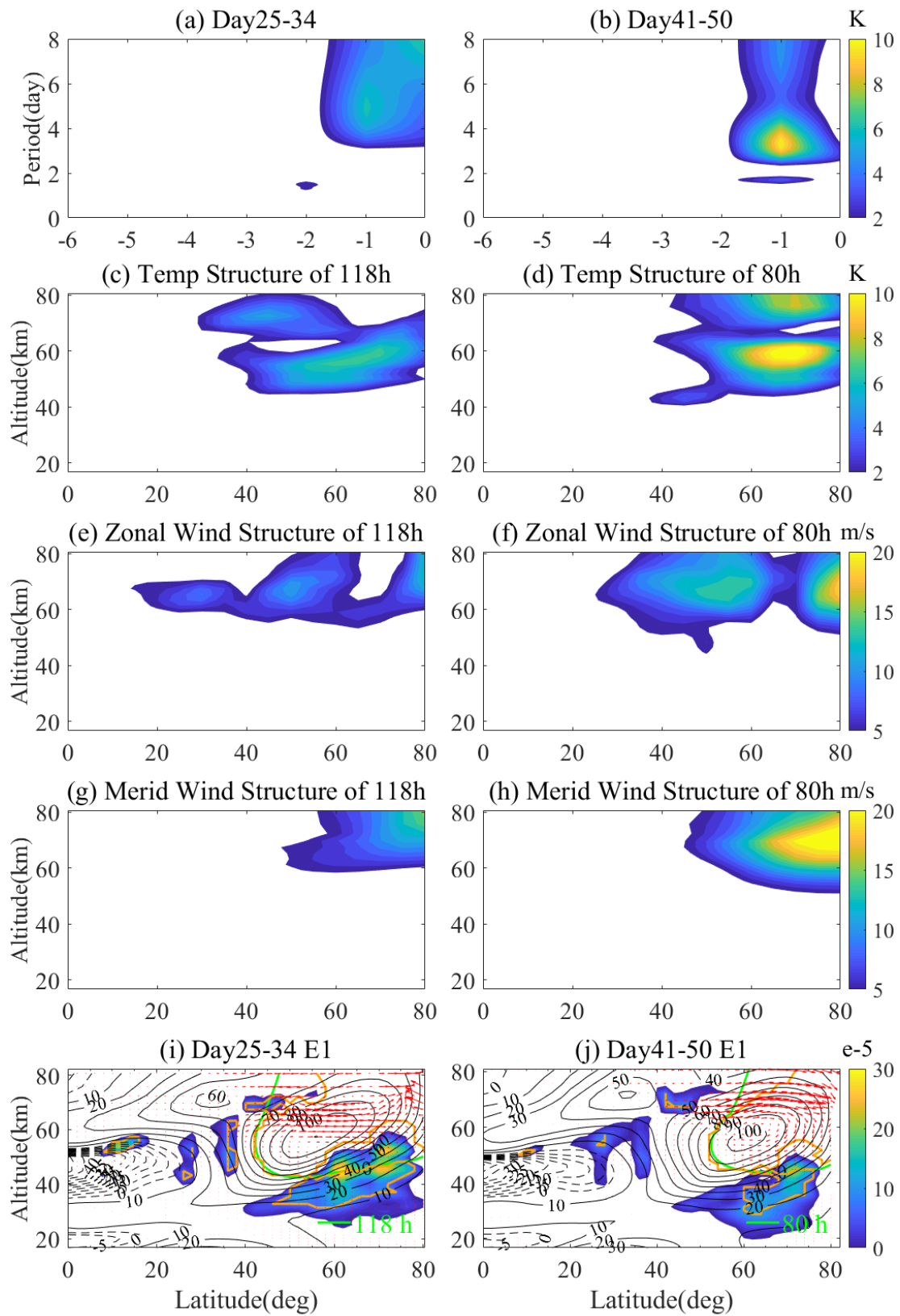


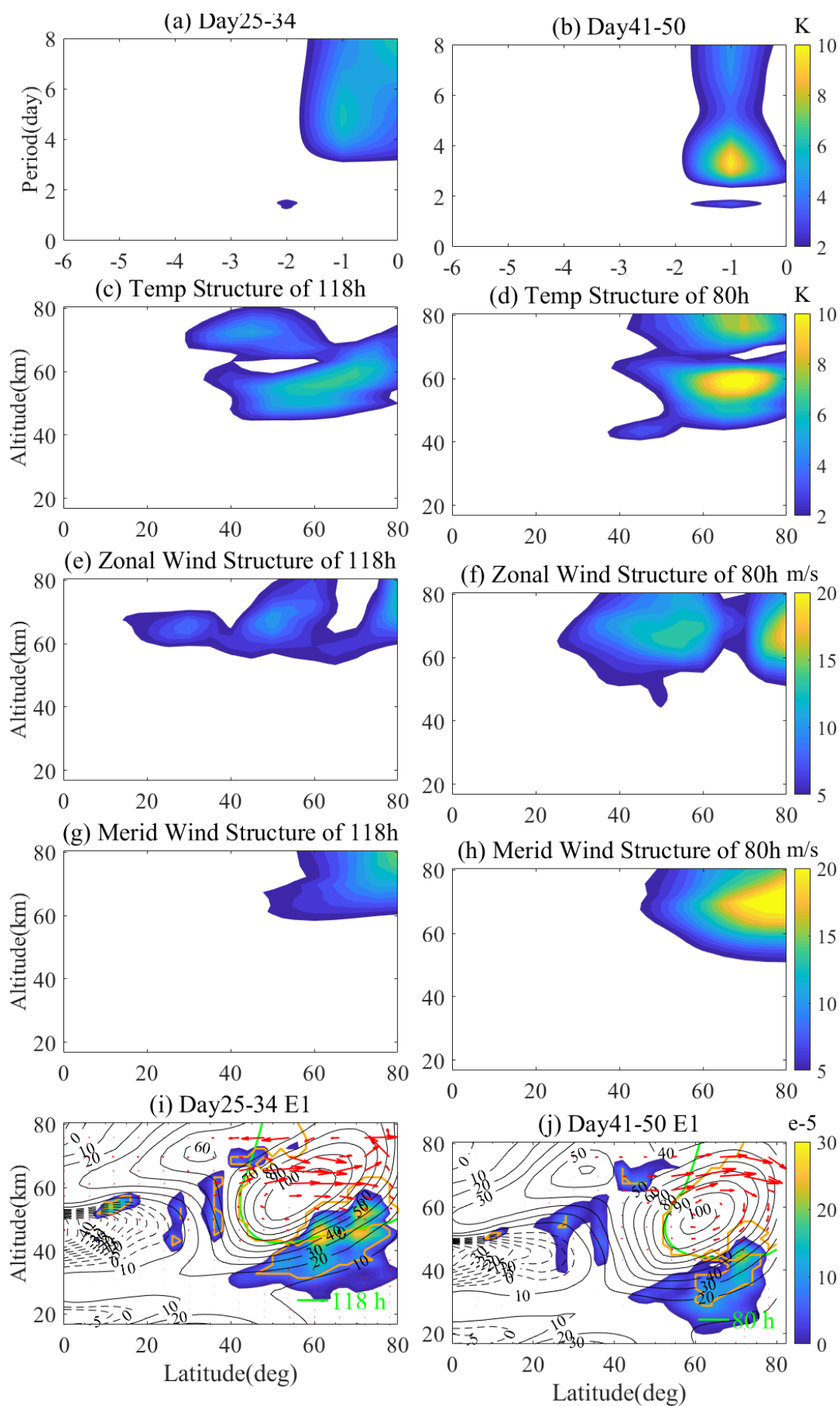


**Figure 7.** The same as Figure 4 but for the E4 during the 2019 austral winter period.

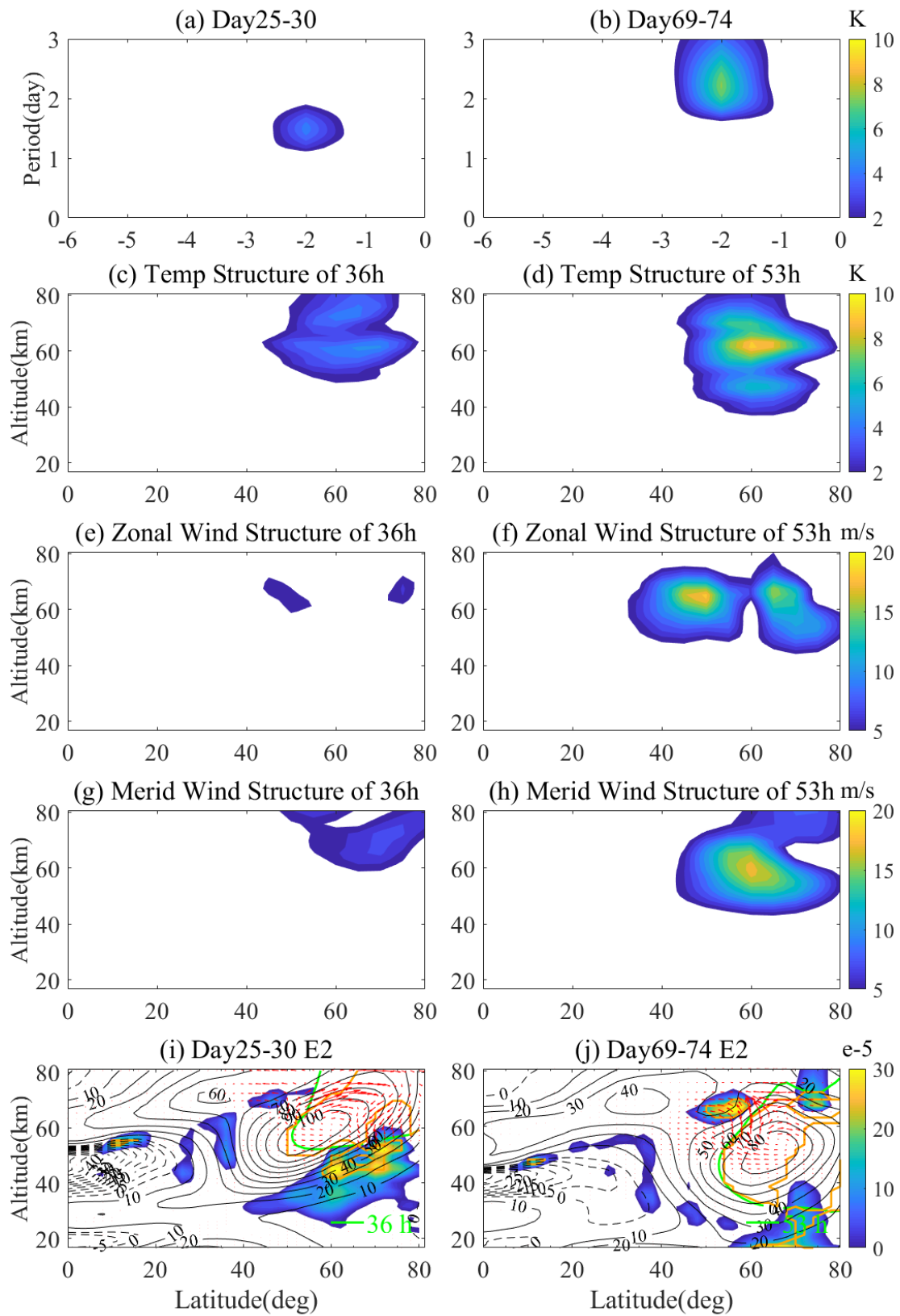


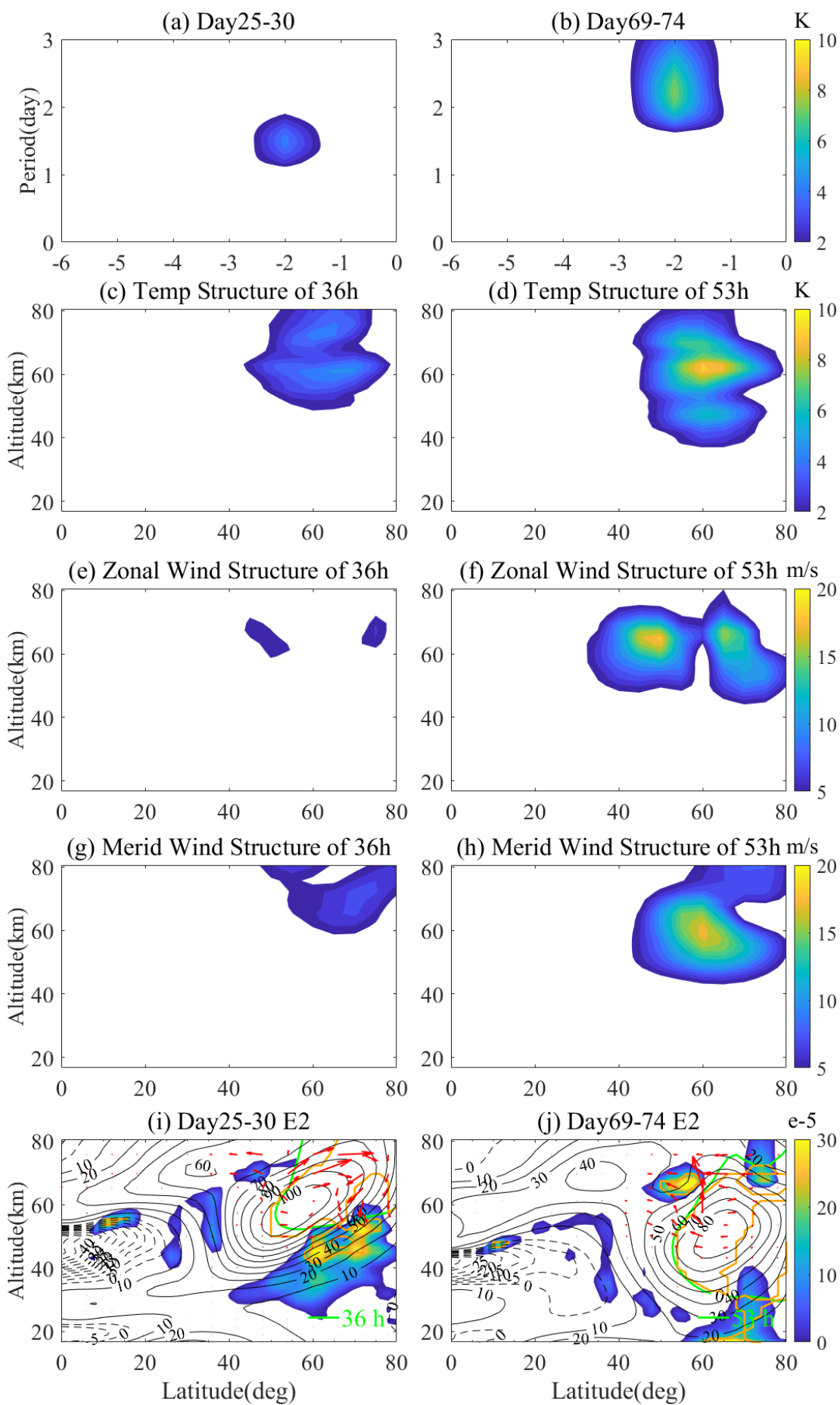
**Figure 8.** The temporal variations of the (a) E1, (b) E2, and (c) E3 QTDWs during the 2019 boreal winter period.



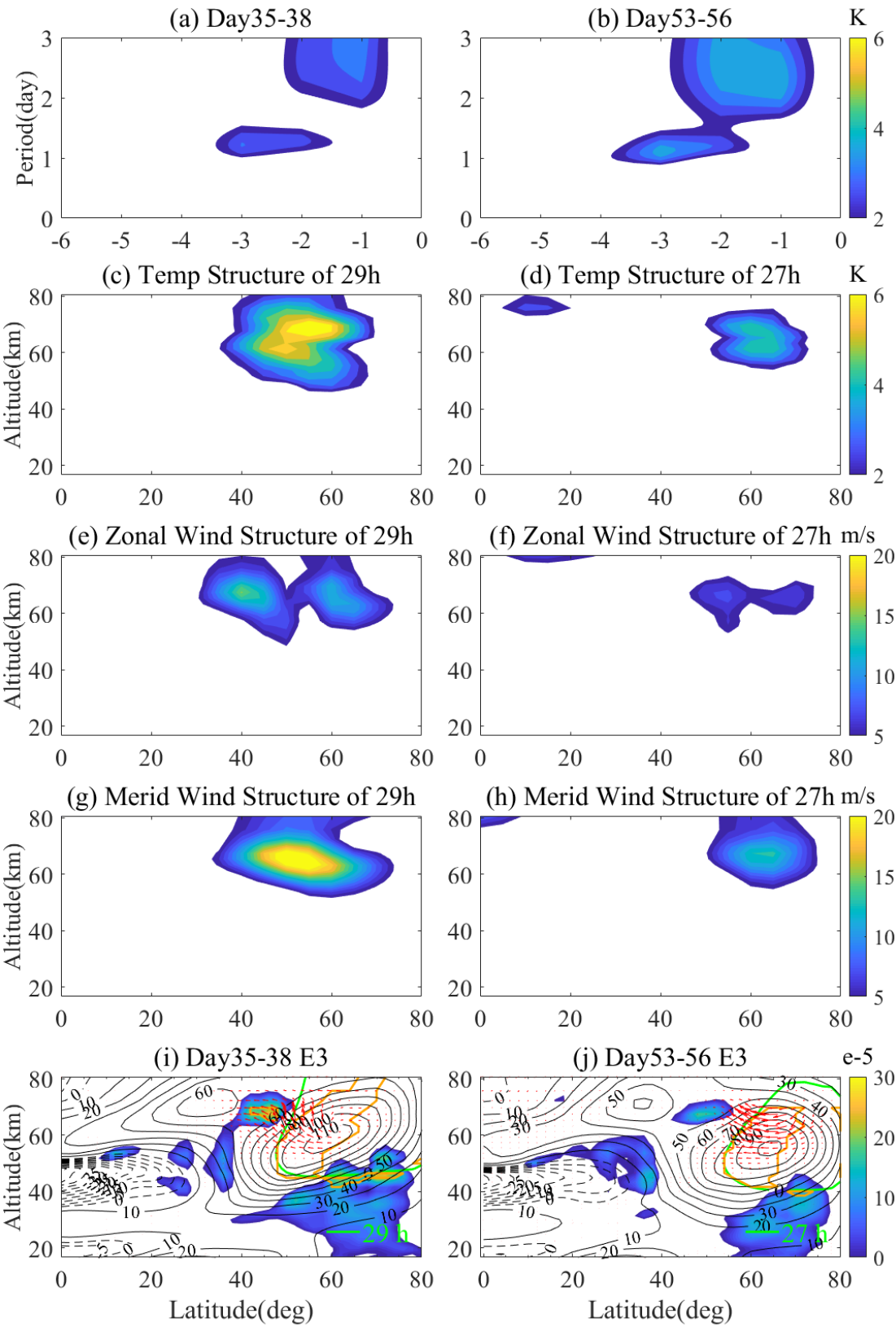


**Figure 9.** he (a, b) spectra, (c, d) temperature spatial structures, (e, f) zonal wind spatial structures, (g, h) meridional wind spatial structures, and (i, j) diagnostic analysis of the E1 typical events during 2019 boreal winter period. The E1 events at 48.2km and 70-80°N were obtained from the MERRA-2 reanalysis.

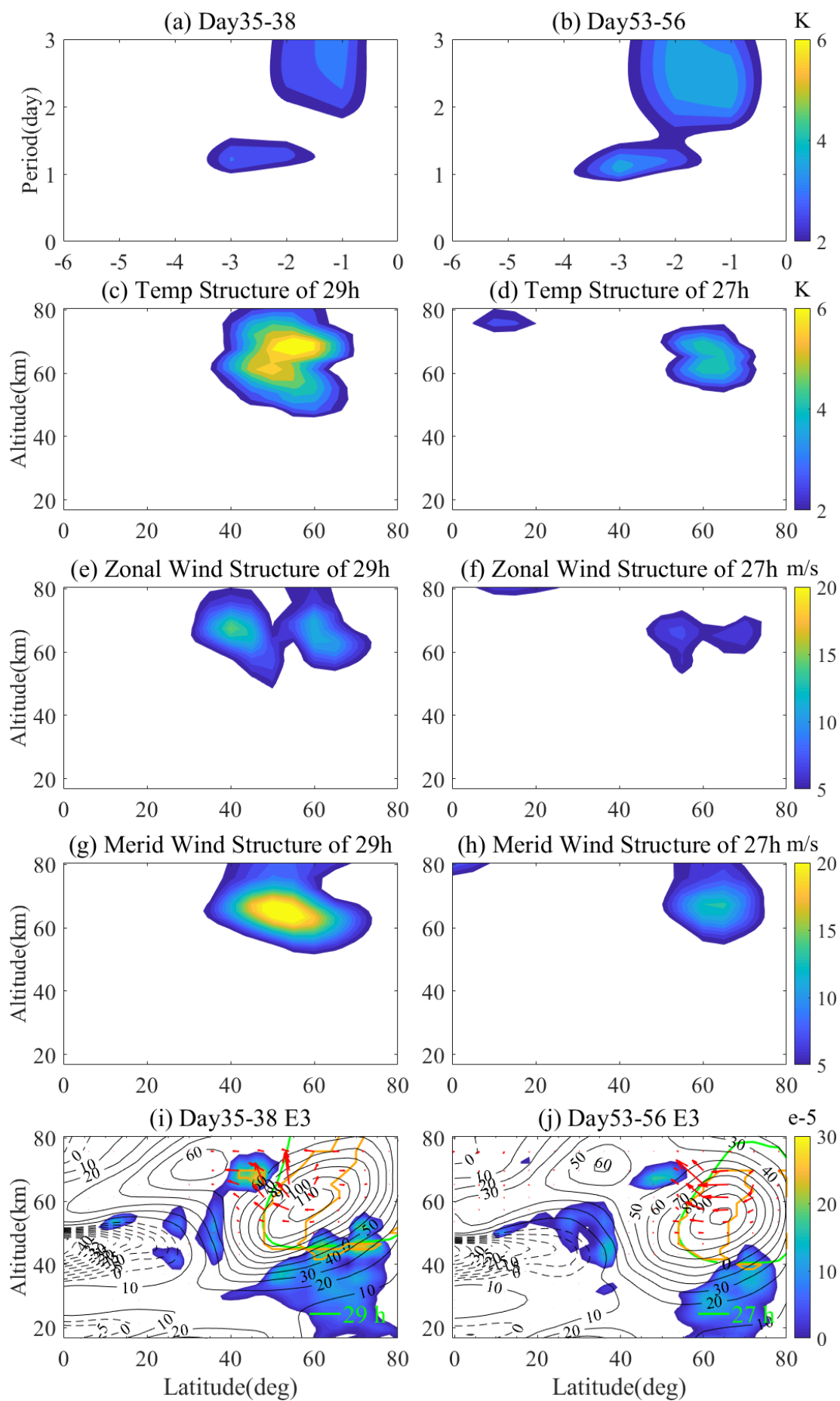




1051 **Figure 10.** The same as Figure 9 but for the E2 during the 2019 boreal winter period.









1054 **Figure 11.** The same as Figure 9 but for the E3 during the 2019 boreal winter period.

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# Eastward-propagating planetary wave in the polar middle atmosphere

**Abstract.** According to MERRA-2 temperature and wind datasets in 2019, this study We presented the global variations of the eastward propagating wavenumber 1 (E1), 2 (E2), 3 (E3), and 4 (E4) planetary waves (PWs) and their diagnostic results in the polar middle atmosphere, ~~using MERRA-2 temperature and wind datasets in 2019. It is clearly shown that.~~ We clearly demonstrated the eastward wave modes exist during winter periods with westward background wind in both hemispheres. The maximum wave amplitudes in the southern hemisphere (SH) are slightly larger and lie lower than those in the northern hemisphere (NH). Moreover, ~~It is also found that~~ the wave perturbations peak at lower latitudes with smaller amplitude as the wavenumber increases. The period of the E1 mode varies ~~from 3 to 3.5~~ 3-5 days in both hemispheres, while the period of E2 mode is slightly longer in the NH (~48 h) than in the SH (~40 h). The periods of the E3 are ~30 h in both SH and NH, and the period of E4 is ~24 h. Despite the shortening of ~~Though the~~ wave periods with the increase of ~~become shorter as the~~ wavenumber ~~increases~~, their mean phase speeds are relatively stable, which are ~53 m/s, ~58 m/s, ~55 m/s, and ~52 m/s at 70° latitudes for W1,

23 W2, W3, and W4, respectively. The eastward PWs occur earlier with  
24 increasing zonal wavenumber, which agrees well with the seasonal  
25 variations of the critical layers generated by the background wind. Our  
26 diagnostic~~Diagnostic~~ analysis also indicated~~shows~~ that the mean flow  
27 instability in the upper stratosphere and upper mesosphere might~~may both~~  
28 contribute to the amplification of the eastward PWs.

## 1 Introduction

~~The dominance of large~~ Large amplitude planetary waves ~~are dominant~~ in the stratosphere, mesosphere, and lower thermosphere ~~regions~~ region and their ~~interactions~~ interaction with zonal mean winds ~~is~~ are the primary driving force of atmospheric dynamics. In addition, sudden stratospheric warmings (SSWs) and quasi-biennial oscillation (QBO) events can dynamically couple the entire atmosphere from the lower atmosphere to the ionosphere (Li et al., 2020; Yamazaki et al., 2020; Yadav et al., 2019; Matthias and Ern, 2018; Stray et al., 2015). Westward propagating planetary wave is one of the prominent features during austral and boreal summer. Westward quasi-2-day waves (Q2DWs) are the most obvious representative waves and one of the most investigated phenomena ~~using~~ by planetary wave observations. Most ~~of the~~ previous studies focused on ~~the~~ westward propagating Q2DWs, ~~i.e., including~~ zonal wavenumbers of 2 (W2), 3 (W3), and 4 (W4) modes (Lainer et al., 2018; Gu et al., 2018b; Wang et al., 2017; Pancheva et al., 2016; Gu et al., 2016a; Gu et al., 2016b; Lilienthal and Jacobi, 2015; Gu et al., 2013; Limpasuvan and Wu, 2009; Salby, 1981). ~~However, limited studies were conducted to understand the seasonal~~ Seasonal variations of the occurrence date, peak amplitude, and wave period for ~~the~~ eastward Q2DWs ~~are rarely studied~~ (Gu et al., 2017; Lu et al., 2013; Alexander and Shepherd, 2010; Sandford et al., 2008; Palo et al., 2007; Merzlyakov and

Pancheva, 2007; Manney and Randel, 1993; Venne and Stanford, 1979).

Typically, Q2DWs ~~usually~~ maximize after the summer solstice in the middle latitudes. The largest wave amplitudes ~~are—generally~~ ~~appeared~~observed near the mesopause in January–February in the Southern Hemisphere (SH), while in the Northern Hemisphere (NH) in July–August (Tunbridge et al., 2011). W3 and W4 Q2DWs reach amplitudes during austral and boreal summer in the mesosphere and lower thermosphere, respectively. The seasonal variation of westward Q2DWs activity ~~has an~~is obvious ~~seasonal variation~~ (Liu et al., 2019; Gu et al., 2018b; Rao et al., 2017). By observing the long-term Q2DW in the NH and SH, Tunbridge et al. (2011) ~~have observed long-term Q2DW in the NH and SH, and found~~reported that W3 is generally stronger than the other two modes in the SH, reaching the amplitude of  $\sim 12\text{K}$ , while W4 is stronger than W3 in the NH, reaching  $\sim 4\text{K}$ . Moreover~~In addition~~, W4 ~~is~~ generally lives longer ~~lived~~ than W3, and W4 ~~can~~is still be observed after the ending of W3 ~~has ended~~. A previous study has demonstrated~~The results of Liu et al. (2004) show~~ that wave source, instability, critical layer, and mean zonal wind are the primary reasons for the seasonal variation of Q2DWs (Liu et al., 2004). By studying the long-term satellite datasets in the SH, Gu et al. (2019) ~~have the long-term observation of satellite datasets in the SH and found~~suggested that the strongest events of W2, W3, and W4 could be~~are~~ delayed with increasing the zonal

wavenumber, and [these events](#) would be indistinguishable during SSWs. ~~Then~~The wave periods of W4, W3 and W2 vary ~~between~~[around](#) ~41-56 h, ~45-52 h, and ~45-48 h, respectively. ~~Furthermore~~[In addition](#), W2 can be observed ~~using~~[in](#) global satellite datasets, ~~but it has an~~[showing weaker](#) amplitude [weaker](#) than W3 and W4 in the NH and SH (Meek et al., 1996). The propagation and amplification of Q2DWs are primarily modulated by instability, refractive index, and critical layer, while the variation of background wind may cause different zonal wavenumber events (Gu et al., 2016a; Gu et al., 2016b). [By analyzing the variation of Q2DWs activity during SSWs](#), Xiong et al. (2018) ~~studied variations in Q2DWs activity during SSWs and found~~[noticed](#) that W1 is generated by the nonlinear interaction between SPW2 and W3. [During SSWs](#), ~~Gu et al. (2018b) found that~~ the coupling between the NH and SH [can enhance](#)~~enhanced~~ the summer easterly ~~during SSWs~~ and [promote](#)~~promoted~~ the nonlinear interaction between W3 and SPW1 ~~during SSWs~~[\(Gu et al., 2018b\)](#).

[Some recent](#)~~Recent~~ studies have [discovered](#)~~found~~ significant eastward planetary waves in the polar stratosphere and mesosphere regions, with periods of nearly [2](#)~~two~~ and [4](#)~~four~~ days (Gu et al., 2017; Sandford et al., 2008; Merzlyakov and Pancheva, 2007; Coy et al., 2003; Manney and Randel, 1993). Planetary waves with zonal wavenumbers -1 (E1) and -2 (E2) correspond to 4- and 2-day waves, respectively. ~~In~~

95 ~~addition~~Furthermore, planetary waves of 1.2-day with wavenumber -3  
96 (E3) and 0.8-day with wavenumber -4 (E4) have been reported to  
97 contain~~found to have~~ the same phase speeds as E1 and E2 (Manney and  
98 Randel, 1993). This series of eastward planetary wave can significantly  
99 affect~~has a significant effect on~~ the thermal and dynamic structure of the  
100 polar stratosphere, resulting in profound~~significant~~ changes in the wind  
101 and temperature of the polar stratosphere (Coy et al., 2003; Venne and  
102 Stanford, 1979). Beyond the knowledge about ~~Palo et al. (1999)~~  
103 ~~demonstrated a series of~~ nonlinear interactions between migrating tides  
104 and Q2DWs (Palo et al., 1999), further investigation has confirmed.  
105 ~~Further research, Palo et al. (2007) presented evidence~~ that E2 Q2DW is  
106 coupled by nonlinear planetary wave and tides in the mesosphere and  
107 lower thermosphere (Palo et al., 2007).

108 By studying and analyzing satellite datasets, Merzlyakov and  
109 Pancheva. (2007) indicated~~analyzed and studied satellite datasets and~~  
110 ~~found~~ that the wave periods of E1 and E2 events range~~is within~~ 1.5-5  
111 days. They reported~~found~~ that EP flux travels~~is~~ from the upper ~~layer to~~  
112 the lower atmosphere, meaning~~suggesting~~ that the upper atmosphere has  
113 a dynamic influence on the lower atmosphere. Sandford et al. (2008)  
114 reported about significant fluctuations of E2 Q2DW in the polar  
115 mesosphere. They ~~found~~indicated the influence of~~that~~ changes in mean  
116 zonal winds during a major SSW on~~influenced~~ the propagation of polar

117 E2. In addition, they ~~proposed~~believe that the significance of E2  
 118 fluctuation ~~is representative~~ in the mesosphere ~~and caused~~driven by the  
 119 instabilities in the polar night jet. ~~Gu et al. (2017) found that the~~For E2,  
 120 amplitude of ~~E2 temperature, zonal wind and meridional wind during the~~  
 121 austral winter can reach  $\sim 10$  K,  $\sim 20$  m/s, and  $\sim 30$  m/s, respectively;~~in~~  
 122 ~~temperature, zonal wind, and meridional wind in the austral winter,~~ while  
 123 those during the boreal winter can drop by almost~~the amplitude of E2~~  
 124 ~~decreases by near~~ two-thirds ~~in the boreal winter~~. Lu et al. (2013) found  
 125 that eastward planetary wave propagation ~~is~~was limited to the winter high  
 126 latitudes, ~~which may be probably~~ because the negative refractive indices  
 127 equatorward of  $\sim 45^\circ\text{S}$  result in evanescent wave characteristics, ~~thus~~  
 128 ~~preventing the propagation of planetary wave to low latitudes. That study~~  
 129 suggested that~~They believe that~~ the instability region at  $\sim 50$ - $60^\circ\text{S}$   
 130 might~~may~~ be induced by the stratospheric polar night jet and/or the  
 131 "double-jet" structure.

132 In this study, we use the~~The~~ second modern retrospective research  
 133 and application analysis (MERRA-2) datasets ~~are used~~ to investigate the  
 134 eastward propagating wave characteristics in the polar stratospheric and  
 135 mesospheric region during~~in~~ 2019, including E1, E2, E3 and E4.  
 136 Specifically~~Particularly~~, we investigate the variation of the occurrence  
 137 date, peak amplitude, and wave period of eastward waves, as well as the  
 138 role of instability, background wind structure, and critical layer in the



propagation and amplification of eastward waves. The remaining parts of this paper are organized as follows. ~~In~~ Section 2 ~~describes~~, the data and methods used in ~~this~~our study ~~are described~~. ~~In~~ Section 3 analyzes the global latitude-temporal variation structure of eastward waves ~~in~~during ~~winter~~the ~~in~~ 2019–~~winter~~. The amplification and propagation features of the eastward planetary waves in the NH and SH with different wavenumber events are ~~examined~~investigated in ~~Sections~~Section 3.1 and 3.2, respectively. ~~In~~ Section 3.3 compares and analyzes ~~the~~ eastward waves in the NH and SH. ~~In Section 4 summarizes our~~All research results ~~are summarized in Section 4~~.

## 2 Data and Analysis

The least-square method is applied to each time window to extract the E1-, E2-, E3-, and E4-wave, with 10-day, 6-day, 4-day, and 4-day, and is used by us to determine the amplitude (Gu et al., 2013). This method has ~~been shown to~~previously been used successfully ~~to~~ identify planetary waves from satellite measurements (Gu et al., 2019; Gu et al., 2018a; Gu et al., 2018b; Gu et al., 2018c; Gu et al., 2013).

$$y = A \cos[2\pi(\sigma \cdot t + s \cdot \lambda)] + B \sin[2\pi(\sigma \cdot t + s \cdot \lambda)] + C \quad (1)$$

The least-squares method is used to fit the ~~a set of~~ parameters ~~of~~ (~~A, B, A, B~~ and ~~C~~), ~~Where~~ ~~where~~ ( ~~$\sigma, t, s, \sigma, t, s$~~  and  ~~$\lambda$~~ ) are the frequency, UT time, zonal wavenumber, and longitudes. The amplitude of wave  $R$  ~~can~~ould be expressed as  $-R = \sqrt{A^2 + B^2}$ .

The second modern retrospective research and application analysis (MERRA-2) ~~covers the data is a set of~~ long-term atmospheric reanalysis datasets ~~started~~initiated by NASA in 1980, ~~and now using an~~. It has been upgraded ~~recently using~~version of the Goddard Earth Observing System Model, Version 5 (GEOS-5) data assimilation system. Briefly, MERRA-2 includes some updates to the model (Molod et al., 2014; Molod et al., 2012) and the Global Statistical Interpolation (GSI) analysis scheme of Wu et al. (2002). The MERRA-2 data ~~consist of~~includes various meteorological variables, ~~such as e.g.~~, net radiation, temperature, relative humidity, and wind speed, ~~etc~~. The spatial coverage of MERRA-2 data is ~~the globe covers the world, with a~~ (spatial resolution: ~~of~~  $0.5^{\circ} \times 0.625^{\circ}$ ; ~~and a~~ temporal resolution: ~~of~~ 1 hour). This ~~kind of~~ meteorological data ~~is are~~ widely used to detect the middle atmosphere such as the planetary wave in the polar atmosphere, global thermal tides, climate variability, and aerosol (Ukhov et al., 2020; Sun et al., 2020; Bali et al., 2019; Lu et al., 2013). ~~These~~Many recent studies ~~indicated~~indicate that the feasibility of using MERRA-2 data for the kind of research in present study can be used in our research with high authenticity. ~~Therefore, we apply the~~The MERRA-2 datasets ~~are used~~ to obtain the variation in background wind, instability, refractive index, and critical layer, and explore the rules of eastward planetary waves propagation and amplification through diagnostic analysis.

The critical layer will absorb or reflect planetary waves ~~during upward propagation~~ from the lower atmosphere during upward propagation. Planetary waves that gain sufficient energy in the instability region will be amplified during reflection. In a sense,~~This shows that~~ the critical layer plays an important role in regulating the amplification and propagation of planetary waves (Gu et al., 2016a; Gu et al., 2016b; Liu et al., 2004).

$$\overline{q_\varphi} = 2\Omega \cos \varphi - \left( \frac{(\overline{u} \cos \varphi)_\varphi}{a \cos \varphi} \right) - \frac{a}{\rho} \left( \frac{f^2}{N^2} \rho \overline{u_z} \right)_z \quad (2)$$

The baroclinic/barotropic instability in the atmospheric space structure is caused by the simultaneous equalization of the negative latitude gradient and the quasi-geostrophic potential vorticity ( $\overline{q_\varphi}$ ). In Equation (2), ~~where ( $\Omega$ ) denote~~is the angular speed of the Earth's rotation;  $\varphi$  is the latitude;  ~~$\overline{u}$  is the~~and the latitude and zonal mean zonal wind; ~~are represented by ( $\varphi$  and  $\overline{u}$ ), in the second part, the ( $a$ ) is~~therepresents the Earth radius; ~~in the last part, ( $\rho, f$ , and  $N$ ) denote the~~background  $\rho$  is the air density;  $f$  is the Coriolis parameter;  ~~$N$  and~~is the buoyancy frequency; ~~respectively. subscripts  $z$  and  $\varphi$~~ The are the vertical and latitudinal gradients. ~~are represented by subscripts ( $z$  and  $\varphi$ ).~~

According to Andrews et al. (1987), the properties of planetary wave propagation can be calculated using ~~define~~ the Eliassen-Palm (EP) flux vectors (F), i.e., to show the properties of planetary wave propagation,

calculated as follows:

$$F = \rho a \cos \varphi \left[ \frac{\overline{u_z' \theta'} - \overline{v' u'}}{\overline{\theta_z'}} \left[ f - \frac{(\overline{u \cos \varphi})_\varphi}{a \cos \varphi} \right] \frac{\overline{v' \theta'} - \overline{w' u'}}{\overline{\theta_z'}} \right] \quad (3)$$

where  $u'$  and  $v'$  are the planetary wave perturbations in the zonal and meridional wind, respectively;  $\theta'$  and  $w'$  are represent the potential temperature and vertical wind, respectively. The planetary wave propagation is only favorable where the square of refractive index  $m^2$  is was positive:

$$m^2 = \frac{\overline{q_\varphi}}{a(\overline{u} - c)} - \frac{s^2}{(a \cos \varphi)^2} - \frac{f^2}{4N^2 H^2} \quad (4)$$

where  $(s)$  denote is the zonal wavenumber,  $c$  is the phase speed is represented by  $(c)$ , and the  $(H)$  represents is the scale height. The square of the refractive index is taken as the waveguide of planetary waves, i.e.,

$$c = -v_0 \cos\left(\frac{\varphi\pi}{180}\right) / sT \quad (5)$$

where  $(c)$  denote the phase speed,  $v_0$  is the equatorial linear velocity is represented by  $(v_0)$ , and the  $(\varphi)$  represents the latitude,  $s$  is the zonal wavenumber and  $T$  represent zonal wavenumber and is the wave period, respectively.

### 3 Results and Discussion

Figure 1 shows the The global temporal-latitude variation structures of E1, E2, E3, and E4 extracted from the 2019 MERRA-2 temperature

datasets using time windows of 10-, 6-, 4- and 4-days, respectively ~~are~~  
~~shown in Figure 1~~. The mean temperature ~~amplitude~~ ~~amplitudes~~ of E1, E2,  
 E3, and E4 at 55.4 km during the periods of 3~5-, 1.5~2.5-, 1~1.5-, and  
 0.9~1.1-days are ~~displayed~~ ~~shown~~ in Figure 1a, 1b, 1c, and 1d,  
~~respectively~~. Eastward waves are characterized by obvious seasonal  
 variations in the SH and NH. In addition, E1, E2 (E3), and E4 reach their  
 maximum ~~amplitude~~ ~~amplitudes~~ at 50-80°(S/N). In the SH, the strongest  
 E1 and E2 events occur on days 209-218 and 167-172, while E3 and E4  
 events occur on days 151-154 and 139-142. This ~~means~~ ~~shows~~ that their  
 occurrence date of maximum amplitude ~~get~~ ~~occurs~~ earlier with increasing  
 zonal wavenumber. In addition, the maximum ~~amplitude~~ ~~amplitudes~~ of  
 E1, E2, E3, and E4 are ~6.0\_K, ~4.2\_K, ~3.6\_K, and ~2.4\_K, ~~respectively~~,  
 indicating that their peak ~~amplitude~~ ~~amplitudes~~ ~~drop with rising~~ ~~decrease~~  
~~with increasing~~ zonal wavenumber. In the NH, the strongest E1, E2, E3  
 and E4 events occur on days 41-50, 69-74, 35-38 and 63-66, ~~respectively~~;  
~~and~~ the corresponding peak ~~amplitude~~ ~~amplitudes~~ are ~5.5\_K, ~3.8\_K,  
 ~2.8\_K and ~1.2\_K, ~~respectively~~. ~~Whilst the~~ ~~The~~ results ~~demonstrate~~ ~~show~~  
~~the decline of that~~ the peak amplitude ~~decreases~~ with increasing zonal  
 wavenumber in the NH, ~~but~~ the occurrence date is irregular. ~~Moreover~~ ~~In~~  
~~addition~~, E4 is ~~relatively~~ weak in the NH and ~~difficult~~ ~~almost impossible~~  
 to find, so W4 is ~~insignificant~~ ~~out of the discussion~~ in the NH. Figure 2  
~~presents~~ ~~shows~~ the changes in zonal mean zonal wind at 70°S and 70°N in

2019. It ~~can be seen~~<sup>is found</sup> that the background wind on days 90-240 (70°S) is dominated by westward wind, and ~~reaches~~<sup>reached</sup> ~80 m/s at ~50 km on days 210; ~~while~~ it is dominated by eastward wind in late and early 2019, and reaches ~40 m/s at ~60 km. Meanwhile, the background wind is primarily westerly wind in late and early 2019 (70°N), and reaches ~90 m/s at ~60 km on days 50; while on days 120-240, the background wind is primarily easterly wind, and the amplitude reaches ~40 m/s on days 200. Compared with Figure 1, the results show that the eastward wave modes exist during winter periods with westward background wind in both hemispheres.

### 3.1 In the Southern Hemisphere

Figure 3 shows that ~~observed~~ maximum amplitude is ~~at observed to be~~ ~48.2 km and ~70-80°S for E1; ~~and E2 and E3 at~~ ~48.2 km and ~60-70°S ~~for E2 and E3; while E4 at~~ ~48.2 km and ~50-60°S ~~for E4~~. ~~The~~ For E1, the ~~observed~~ maximum perturbation ~~observed for E1 occur~~ ~~occurs~~ on days 211-220 (with an amplitude of ~8.5 K), and the remaining fluctuations occur on days 161-170, 187-196 and 231-240; ~~while the~~. For E2, the ~~observed maximum perturbation happens~~ ~~maximizes~~ at days 219-224 (with an amplitude of ~7.8 K), and ~~also shows~~ three peaks ~~appear on~~ at days 139-144, 173-178, and 187-192. ~~Regarding E3~~ Besides, the strongest ~~perturbation~~ ~~E3~~ occurs on days 151-154 (with an amplitude of ~5.2 K), ~~while~~ and the rest are distributed on days 141-144, 201-204,

~~and~~ 209-202. E4 perturbations are distributed on days 127-130, 145-148, 161-164, 213-216, with weak amplitude of  $\sim 3$  K. ~~Since~~~~According to~~ ~~earlier~~~~previous~~ studies ~~mentioned that~~, the wave period of the eastward wave ~~can vary~~~~varies~~. ~~Therefore~~, ~~we also investigate the~~ periodic ~~variabilities~~~~variability~~ of E1, E2, E3, and E4 ~~has also been investigated~~.

The results show that the period corresponding to the maximum perturbation of E1 falls between  $\sim 106$  (days 187-196) and  $\sim 69$  h (days 211-220), and their wave periods vary ~~significantly~~~~greatly~~. ~~Nonetheless~~~~However~~, the wave period of E2 gradually changes from  $\sim 42$  h (days 139-144) to  $\sim 38$  h (days 219-224), and its stability is stronger than that of E1. ~~In addition, the~~~~The~~ wave periods of E3 and E4 are about  $\sim 39$  h and  $\sim 24$  h, respectively. ~~These results reflect that~~~~Thus~~, E2, E3, and E4 wave periods are more stable ~~compared to~~~~than~~ E1.

The spectra, spatial (vertical and latitudinal) structures of temperature, zonal and meridional wind, and diagnostic analysis of E1 are extracted from the ~~two~~ corresponding ~~two representative~~ events, (~~refer to as shown in~~ Figure 4). ~~Figures~~~~Figure~~ 4a, 4b ~~shows~~~~shows~~ the least-squares fitting spectra for MERRA-2 temperature on days 187–196, 211-220 at  $\sim 48.2$  km and  $\sim 70$ - $80^\circ$ S, when and where the E1 maximizes. An eastward wavenumber -1 signal with the ~~periods~~~~period~~ of  $\sim 106$  h and  $\sim 69$  h clearly dominates the whole spectrum. The temperature spatial ~~structures~~~~structure~~ corresponding to these E1 (*i.e.*,  $\sim 106$  h and  $\sim 69$  h) are

~~displayed is shown~~ in ~~Figures~~Figure 4c, 4d. The temperature spatial structure of E1 ~~exhibits presents an~~ obvious amplitude bimodal structure at  $\sim 70\text{--}80^\circ\text{S}$  and  $\sim 50$  km, and  $\sim 70\text{--}80^\circ\text{S}$  and  $\sim 60$  km, with the maximum at  $\sim 70\text{--}80^\circ\text{S}$  and  $\sim 50$  km. The strongest temperature amplitude of E1 occurs at  $\sim 50$  km and  $\sim 70\text{--}80^\circ\text{S}$  with an amplitude of  $\sim 10$  K on the days 211-220, and the other peak is  $\sim 9$  K ( $\sim 70\text{--}80^\circ\text{S}$  and  $\sim 60$  km). The temperature amplitude of  $\sim 9$  K occurs at  $\sim 50$  km and  $\sim 70\text{--}80^\circ\text{S}$  during days 187-196, and the rest is  $\sim 7$  K ( $\sim 70\text{--}80^\circ\text{S}$  and  $\sim 60$  km). The corresponding spatial structures of zonal wind and meridional wind of these E1 ~~are is~~ shown in Figures 4e, 4f, 4g, and 4h. The maximum zonal wind amplitude of E1 occurs at  $\sim 60\text{--}70^\circ\text{S}$  and  $\sim 60$  km with an amplitude of  $\sim 14$  m/s on days 187-196, and  $\sim 20$  m/s at  $\sim 50\text{--}60^\circ\text{S}$  and  $\sim 60$  km on days 211-220. The amplitude of E1 meridional wind ~~hits reaches~~  $\sim 10$  m/s at  $\sim 70\text{--}80^\circ\text{S}$  and  $\sim 55$  km (days 187-196) and  $\sim 17$  m/s at  $\sim 70\text{--}80^\circ\text{S}$  and  $\sim 60$  km (days 211-220), respectively.

~~Figures~~Figure 4i, 4j ~~shows~~ shows the diagnostic analysis results for the E1 events during days 187–196, and 211–220, respectively. ~~Apparently, It is clear that~~ the EP flux vectors is more favorable to propagate in the SH winter and is dramatically amplified by the mean flow instabilities and appropriate background winds at polar and between  $\sim 40$  km and  $\sim 80$  km, with EP flux propagating into the upper atmosphere (Figure 4i). Meanwhile, there is an EP flux at the mid-latitudes and



~60-80 km, which propagates into the lower atmosphere. The wave-mean flow interactions near its critical layer (106 h) of the green curve amplifies E1, and the positive refractive index region surrounded by the yellow curve also enhances E1 propagation. In addition, the strong instability and weak background wind at ~70-80°S and ~40-60 km provide sufficient energy for the upward EP flux to propagate and amplify. Nevertheless, the downward propagating EP flux is amplified by weak instability and strong background wind at ~50-60°S and ~60-70 km. Besides, both upward and downward EP fluxes eventually propagate toward the equator at ~50 km. Figure 4j shows that EP flux on days 211-220 propagates downward and amplifies after the interaction of the critical layer (~69 h). Facilitated by the positive refractive index region, which strong instability and weak background wind at ~50-60°S and ~60-70 km provide sufficient energy, and ultimately point towards the equator at ~50 km. The results show that the weak background wind and strong instability in the polar region can promote the upward propagation and amplification of the EP flux. Meanwhile, the appropriate background wind and instability in the mid-latitudes are also conducive to the downward propagation and amplification of EP flux. In other words, instability and appropriate background wind dominate the propagation and amplification of the E1.

For E2, the spectra are observed at ~48.2 km and ~60-70°S on

days 173-178, and 219-224 for E2, when the eastward wavenumber -2 becomes the primary wave mode with the wave ~~periods~~period of ~38 h and ~39 h, respectively (as shown in ~~Figures~~Figure 5a, 5b). The temperature spatial ~~structures~~structure corresponding to these E2 (~~~38 h and ~39 h~~) ~~are~~is ~~presented~~shown in ~~Figures~~Figure 5c, 5d. The temperature spatial structure of E2 shows an obvious amplitude bimodal structure at ~60-70°S and ~50 km, and ~60-70°S and ~60 km, with the maximum at ~60-70°S and ~50 km. The maximum temperature amplitude of E1 occurs at ~50 km and ~60-70°S with an amplitude of ~7.5 K on the days 173-178, and the other peak is ~6 K (~70°S and ~60 km). The temperature amplitude of ~10 K ~~happensee~~occurs at ~50 km and ~60-70°S during days 219-224, and the rest is ~6 K (~70°S and ~60 km). The corresponding spatial structures of zonal wind and meridional wind of these E2 ~~are~~is ~~illustrated~~shown in Figures 5e, 5f, 5g, and 5h. The zonal wind spatial structure of E2 shows an obvious amplitude bimodal structure at ~50-60°S and ~60 km, and ~70-80°S and ~60 km, with the maximum at ~50-60°S and ~60 km. The maximum zonal wind ~~amplitude~~amplitudes of E2 ~~appearsee~~appear at ~50-60°S and ~60 km with an amplitude of ~10 m/s on days 173-178, and the other peak is ~9 m/s (~70-80°S and ~60 km). The zonal wind amplitude of ~20 m/s occurs at ~50-60°S and ~60 km on days 219-224, and the rest is ~15 m/s (~70-80°S and ~60 km). The amplitude of E2 meridional wind reaches ~13 m/s at

~70-80°S and ~60 km (days 173-178) and ~27 m/s at ~70-80°S and ~60 km (days 219-224), respectively.

~~The results in~~ Figures 5i and 5j ~~illustrateshow the~~ diagnostic analysis during days 173-178, ~~and~~ 219-224 for E2, ~~respectively~~. ~~Obviously, it is clear that~~ E2 is more ~~likelyfavorable~~ to propagate in the SH winter and is dramatically amplified by the mean flow instabilities at ~~the~~ middle-high latitudes between ~40 km and ~80 km, ~~with~~ ~~With~~ EP flux propagating into the lower atmosphere, ~~and EP flux it~~ eventually ~~propagatespropagate~~ toward the equator at ~50 km. Besides, E2 is amplified and propagated by the wave-mean flow interactions near its critical layer (~38 h) of the green curve, and the promoting effect of the positive refractive index region surrounded by the yellow curve. Meanwhile, the weak instability and strong background wind at ~50-60°S and ~50-70 km provide ~~the~~ energy for the propagation and amplification of EP flux into the lower atmosphere during days 173-178 (Figure 5i). ~~According to~~ ~~the~~ diagnostic analysis of days 219-224, ~~it is found that~~ E2 obtains sufficient energy from strong instability and strong background wind at ~50-60°S and ~60-70 km, ~~and~~. ~~It~~ is amplified and propagated into the lower atmosphere through the critical layer and positive refractive index action, (as shown in Figure 5j). The results show that the background wind at ~50-60°S and ~50-70 km is weaker on days 173-178 than on days 219-224, ~~;~~ and the instability at ~50-60°S and ~60-70 km is stronger on

days 219-224 than on days 173-178. ~~According to Figure 5a, 5b~~ Our results ~~show~~ shows that E2 has absorbed sufficient energy to be amplified under the background conditions during days 219-224 (~~Figures 5a, 5b~~).

Figures 6a and 6b show the observed spectra of E3 at  $\sim 48.2$  km and  $\sim 60$ - $70^\circ$ S on days 151-154 and 201-204, and the wave ~~periods~~ period of locked wavenumber -3 ~~are~~ is  $\sim 29$  h and  $\sim 29$  h, respectively. The corresponding temperature spatial structures of these E3 (~~i.e.,  $\sim 29$  h and  $\sim 29$  h~~) ~~are displayed~~ ~~is shown~~ in ~~Figures~~ Figure 6c, 6d. The temperature spatial structure of E3 shows an obvious amplitude bimodal structure at  $\sim 60$ - $70^\circ$ S and  $\sim 50$  km, and  $\sim 60$ - $70^\circ$ S and  $\sim 60$  km, with the maximum at  $\sim 60$ - $70^\circ$ S and  $\sim 50$  km. Besides, E3 also has a weak peak at  $\sim 60$ - $70^\circ$ S and  $\sim 70$  km. The strongest temperature amplitude of E3 occurs at  $\sim 50$  km and  $\sim 60$ - $70^\circ$ S with an amplitude of  $\sim 6$  K on the days 151-154, and the other peak is  $\sim 5$  K ( $\sim 60$ - $70^\circ$ S and  $\sim 60$  km). The temperature amplitude of  $\sim 5$  K ~~happense~~ occurs at  $\sim 50$  km ( $\sim 60$  km) and  $\sim 60$ - $70^\circ$ S during days 201-204. The corresponding spatial structures of zonal wind and meridional wind of these E3 ~~is~~ are shown in Figures 6e, 6f, 6g, and 6h. The zonal wind spatial structure of E3 shows an obvious amplitude bimodal structure at  $\sim 70$ - $80^\circ$ S and  $\sim 60$  km, and  $\sim 50$ - $60^\circ$ S and  $\sim 60$  km. The zonal wind amplitudes of E3 occur at  $\sim 70$ - $80^\circ$ S and  $\sim 60$  km ( $\sim 50$ - $60^\circ$ S and  $\sim 60$  km) with an amplitude of  $\sim 9$  m/s on days 151-154, and  $\sim 9$  m/s at  $\sim 70$ - $80^\circ$ S and  $\sim 60$  km ( $\sim 50$ - $60^\circ$ S and  $\sim 60$  km) on days 201-204. The amplitude of

E3 meridional wind ~~hitsreaches~~ ~13 m/s at ~60-70°S and ~55 km (days 151-154) and ~16 m/s at ~60-70°S and ~55 km (days 201-204), respectively.

EP flux of E3 is similar to that of E2. The instability and appropriate background wind at ~~the~~ mid-high latitudes between ~50 km and ~70 km dramatically amplify the propagation of E3, which is enhanced by the interaction near the critical layer (~29 h) and the positive refractive index region- (Figures 6i and 6j). ~~Notably, It is worth noting that~~ the strong instability and weak background wind at ~50-60°S and ~60-70 km on days 151-154 provide sufficient energy for the propagation and amplification of EP flux into the lower atmosphere, and ultimately point ~~towardtowards~~ the equator at 50 km. ~~During days 201-204, the~~The EP flux ~~propagatespropagate teinto~~ the lower atmosphere ~~during days 201-204,~~ and ~~isgets~~ amplified by interaction at the critical layer (~29 h). ~~BesidesIn addition,~~ weak instability and weak background wind at ~50-60°S and ~60-70 km provide the energy to amplify ~~the~~ E3 propagation. ~~Combine with Figure~~Figures 6c, 6d ~~indicate that,~~ the stronger the instability at ~50-60°S and ~60-70 km, the stronger the temperature amplitude of E3. We believe that the background wind and instability at ~50-60°S and ~60-70 km are the ~~main~~primary reasons for the propagation and amplification of EP flux into the lower atmosphere.

~~The~~For E4, the spectra ~~appearare observed to be~~ at ~48.2 km and

~50-60°S on days 127-130, ~~and 213-216 for E4~~, when the eastward  
 wavenumber -4 signal with the wave period of ~25 h and ~21 h, ~~as shown~~  
~~in Figure~~ (see Figures 7a, 7b). The corresponding temperature spatial  
 structures of these E4 (i.e., ~25 h and ~21 h) ~~are~~ is shown in  
 Figures ~~Figure~~ 7c, 7d. The temperature spatial structure of E4 shows an  
 obvious amplitude bimodal structure at ~50-60°S and ~50 km, and  
 ~50-60°S and ~60 km, with the maximum at ~50-60°S and ~50 km. The  
 maximum temperature amplitude of E4 occurs at ~50 km and ~50-60°S  
 with an amplitude of ~4 K on the days 127-130, and the other peak is ~3  
 K (~60-70°S and ~60 km). The temperature amplitude of ~3 K occurs at  
 ~50 km (~60 km) during days 213-216. The corresponding spatial  
 structures of zonal wind and meridional wind of these E4 ~~are presented~~  
~~shown~~ in Figures 7e, 7f, 7g, and 7h. The zonal wind spatial structure of  
 E4 shows an obvious amplitude bimodal structure at ~50-60°S and ~55  
 km, and ~60-70°S and ~55 km, with the maximum at ~50-60°S and ~55  
 km. The maximum zonal wind amplitude of E4 ~~happense~~ occurs at  
 ~50-60°S and ~55 km with an amplitude of ~9 m/s on days 127-130, and  
 the other peak is ~5 K (~60-70°S and ~55 km). The zonal wind amplitude  
 of ~5 m/s occurs at ~50-60°S (~60-70°S) and ~55 km on days 213-216.  
 The amplitude of E4 meridional wind reaches ~8 m/s at ~60-70°S and  
 ~55 km (days 127-130) and ~10 m/s at ~60-70°S and ~55 km (days  
 213-216), respectively.

Diagnostic analysis for E4 on days 127-130, and 213-216 are shown in Figures 7i and 7j, respectively. The results show that E4 is dramatically amplified by the mean flow instabilities at the middle-high latitudes between ~50 km and ~70 km, with. With EP flux propagating into the lower atmosphere, and EP flux it finally propagates eventually propagate toward the equator at ~50 km. E4 is amplified and propagated by the wave-mean flow interaction near the critical layer (~25 h, ~21 h), and the positive refractive index region generates provides the promoting effect. The strong instability and weak background wind at ~50-60°S and ~60-70 km provide sufficient energy for the propagation and amplification of EP flux into the lower atmosphere during days 127-130. Besides, E4 obtains energy from weak instability and weak background wind at ~50-60°S and ~60-70 km on days 213-216, and it is amplified and propagated into the lower atmosphere. The background wind at ~50-60°S and ~60-70 km on days 127-130 is similar to on days 213-216, and the instability at ~50-60°S and ~60-70 km is stronger on days 127-30 than on days 213-216. According to Figure 7a, and 7b, results shows that E4 absorbs has absorbed sufficient energy to be amplified under the background conditions on days 127-130, and. The the temperature amplitude on 127-130 days is stronger.

### 3.2 In the Northern Hemisphere

Figure 8 shows that ~~the observed~~ maximum amplitude ~~appears at is~~ ~~observed to be~~ ~59.2 km and ~70-80°N for E1, and the E2 and E3 peaks at ~59.2 km and ~60-70°N. The maximum perturbation of E1 ~~occurs is~~ ~~observed to be at on~~ days 41-50 (with an amplitude of ~8K), ~~while and~~ the remaining ~~fluctuations occur fluctuation occur~~ on days 25-34, ~~and~~ 339-348. Besides, the strongest E2 occurs on days 69-74 (with an amplitude of ~7 K), and the rest are distributed on days 25-30, 317-322, and 341-346, ~~while~~. ~~By contrast~~, the E3 maximizes ~~on at~~ days 35-38 (with an amplitude of ~3K), and also shows ~~a peak one peaks at on~~ days 53-56. Based on the study of the wave period in the SH for eastward wave, the periodic ~~variabilities of E1, E2 and E3 variability~~ in the NH ~~for E1, E2, and E3 is are~~ also ~~examined investigated~~. The wave ~~period periods~~ of E1 ~~decreases decreased~~ from a maximum of ~118 h (days 25-34) to ~80 h (days 41-50), indicating ~~the instability of that~~ the wave period of E1 ~~is unstable~~ in the NH. The E2 events occur on days 25-30, 69-74, 317-322, and 341-346, of which the corresponding wave periods are ~36, ~53, ~52, and ~48 h, which are stronger ~~and more~~ stable than E1. Besides, the wave period of E3 is relatively stable at ~29 h and ~27 h. Thus, E2 and E3 wave periods are more stable than E1.

The spectra, spatial (vertical and latitudinal) structures of temperature, zonal and meridional wind, and diagnostic analysis of E1 are extracted from the corresponding representative events ~~(,~~ as shown in



Figure 9). Figures 9a and 9b show the observed spectra of E1 at  $\sim 59.2^\circ\text{N}$  and  $\sim 70\text{--}80^\circ\text{N}$  on days 25-34 and 41-50, and the wave period of locked wavenumber -1 is  $\sim 118$  h and  $\sim 80$  h, respectively. The corresponding temperature spatial structures of these E1 ( $\sim 118$  h and  $\sim 80$  h) are shown in Figure 9c, 9d. The temperature spatial structure of E1 shows an obvious amplitude bimodal structure during days 25-34 at  $\sim 60\text{--}70^\circ\text{N}$  and  $\sim 60$  km, and  $\sim 40\text{--}50^\circ\text{N}$  and  $\sim 70$  km, with the maximum at  $\sim 60\text{--}70^\circ\text{N}$  and  $\sim 60$  km. Besides, E1 also has bimodal structure on days 41-50 at  $\sim 60\text{--}70^\circ\text{N}$  and  $\sim 60$  km, and  $\sim 60\text{--}70^\circ\text{N}$  and  $\sim 70$  km. The strongest temperature amplitude of E1 occurs at  $\sim 60\text{--}70^\circ\text{N}$  and  $\sim 60$  km with an amplitude of  $\sim 7$  K on the days 25-34, and the other peak is  $\sim 4$  K ( $\sim 40\text{--}50^\circ\text{N}$  and  $\sim 70$  km). The temperature amplitude of  $\sim 10$  K occurs at  $\sim 60$  km and  $\sim 60\text{--}70^\circ\text{N}$  during days 41-50, and the rest is  $\sim 8$  K ( $\sim 60\text{--}70^\circ\text{N}$  and  $\sim 70$  km). The corresponding spatial structures of zonal wind and meridional wind of these E1 are illustrated in Figures 9e, 9f, 9g, and 9h. The zonal wind spatial structure of E1 presents an obvious amplitude bimodal structure at  $\sim 70\text{--}80^\circ\text{N}$  and  $\sim 70$  km, and  $\sim 50\text{--}60^\circ\text{N}$  and  $\sim 70$  km. The zonal wind amplitude of  $\sim 13$  m/s occurs at  $\sim 70\text{--}80^\circ\text{N}$  and  $\sim 70$  km on days 25-34, and the rest is  $\sim 10$  m/s ( $\sim 50\text{--}60^\circ\text{N}$  and  $\sim 70$  km). In addition, there is a weak peak of 9K during days 25-34 ( $\sim 30\text{--}40^\circ\text{N}$  and  $\sim 70$  km). The maximum zonal wind amplitude of E1 occurs at  $\sim 70\text{--}80^\circ\text{N}$  and  $\sim 70$  km with an amplitude of  $\sim 19$  m/s on days

41-50, and the other peak is  $\sim 13\text{K}$  ( $\sim 50\text{-}60^\circ\text{N}$  and  $\sim 70\text{ km}$ ). The amplitude of E1 meridional wind ~~hits~~reaches  $\sim 14\text{ m/s}$  at  $\sim 70\text{-}80^\circ\text{N}$  and  $\sim 70\text{ km}$  (days 25-34) and  $\sim 22\text{ m/s}$  at  $\sim 70\text{-}80^\circ\text{N}$  and  $\sim 70\text{ km}$  (days 41-50), respectively.

The diagnostic analysis results for E1 (in Figures 9i and 9j) ~~suggests~~show the dramatic amplification of that E1 ~~is dramatically amplified~~ by the mean flow instabilities at the middle-high latitudes between  $\sim 50\text{ km}$  and  $\sim 70\text{ km}$ , ~~with~~. With the propagation of EP flux ~~propagating~~ into the polar lower atmosphere, ~~and EP flux it~~ eventually ~~propagates~~propagate toward the equator at  $\sim 50\text{ km}$ . The wave-mean flow interaction near the critical layers ( $\sim 118\text{ h}$ ,  $\sim 80\text{ h}$ ) amplifies and propagates E1, and the promoting effect of the positive refractive index region amplifies E1. ~~Furthermore~~Besides, the weak instability and strong background wind at  $\sim 40\text{-}50^\circ\text{N}$  and  $\sim 60\text{-}70\text{ km}$  ~~generate~~the provide energy for the propagation and amplification of EP flux into the polar lower atmosphere during days 25-34. The E1 obtains sufficient energy from weak instability and suitable background wind on days 41-50 at  $\sim 40\text{-}50^\circ\text{N}$  and  $\sim 60\text{-}70\text{ km}$ , and is amplified and propagated into the polar lower atmosphere through the critical layer and positive refractive index action. The background wind at  $\sim 40\text{-}50^\circ\text{N}$  and  $\sim 60\text{-}70\text{ km}$  is stronger on days 25-34 than on days 41-50, but their instability is similar, indicating that stronger background winds ~~might weaken~~may be unfavorable to E1

propagation and amplification at the mid-northern latitudes. ~~According to~~  
~~Figure 9a, 9b~~ Our results ~~show~~ shows that E1 ~~absorbs~~ has absorbed  
~~sufficient~~ adequate energy to be amplified under the background  
conditions during days 41-50, ~~reflecting~~ showing a stronger temperature  
amplitude (see Figures 9a, 9b).

For E2, the The spectra are ~~observed to be~~ at  $\sim 59.2$  km and  $\sim 60$ - $70^\circ\text{N}$   
on days 25-30, and 69-74 ~~for E2~~, when the eastward wavenumber -2  
signal with the period ~~of~~  $\sim 36$  h and  $\sim 53$  h, (as shown in ~~Figures~~ Figure 10a,  
10b). The corresponding temperature spatial structures of these E2 (i.e.,  
 $\sim 36$  h and  $\sim 53$  h) ~~are presented~~ is shown in ~~Figures~~ Figure 10c, 10d. The  
temperature spatial structure of E2 ~~demonstrates~~ shows an obvious  
amplitude bimodal structure at  $\sim 60$ - $70^\circ\text{N}$  and  $\sim 60$  km, and  $\sim 60$ - $70^\circ\text{N}$  and  
 $\sim 70$  km, with the maximum at  $\sim 60$ - $70^\circ\text{N}$  and  $\sim 60$  km. The maximum  
temperature amplitude of E2 occurs at  $\sim 60$ - $70^\circ\text{N}$  and  $\sim 60$  km with an  
amplitude of  $\sim 5$  K on days 25-30, and the other peak is  $\sim 4$  K ( $\sim 60$ - $70^\circ\text{N}$   
and  $\sim 70$  km). The temperature amplitude of  $\sim 9$  K occurs on days 69-74 at  
 $\sim 60^\circ\text{S}$  and  $\sim 60$  km, and the other peaks are  $\sim 7$  K ( $\sim 60$ - $70^\circ\text{N}$  and  $\sim 70$  km),  
 $\sim 5$  K ( $\sim 60$ - $70^\circ\text{N}$  and  $\sim 50$  km). The corresponding spatial structures of  
zonal wind and meridional wind of these E2 ~~is~~ are shown in Figures 10e,  
10f, 10g, and 10h. The zonal wind spatial structure of E2 shows an  
obvious amplitude bimodal structure at  $\sim 60$ - $70^\circ\text{N}$  and  $\sim 60$  km, and  
 $\sim 40$ - $50^\circ\text{N}$  and  $\sim 60$  km, with the maximum at  $\sim 40$ - $50^\circ\text{N}$  and  $\sim 60$  km. The

maximum zonal wind amplitude of E2 ~~appears~~ occurs at  $\sim 60\text{-}70^\circ\text{N}$  and  $\sim 60$  km ( $\sim 40\text{-}50^\circ\text{N}$  and  $\sim 60$  km) with an amplitude of  $\sim 6$  m/s on days 25-30. Zonal wind amplitude occurs at  $\sim 40\text{-}50^\circ\text{N}$  and  $\sim 60$  km with an amplitude of  $\sim 18$  m/s on days 41-50, and the other peak is  $\sim 16$  K ( $\sim 60\text{-}70^\circ\text{N}$  and  $\sim 60$  km). The amplitude of E2 meridional wind reaches  $\sim 7$  m/s at  $\sim 60\text{-}70^\circ\text{N}$  and  $\sim 70$  km (days 25-30) and  $\sim 18$  m/s at  $\sim 60\text{-}70^\circ\text{N}$  and  $\sim 60$  km (days 41-50), respectively.

~~Figures 10i and 10j show~~ The diagnostic analysis of E2 on days 25-30, and 69-74 ~~are shown in Figures 10i and 10j, respectively for E2. It is clear that~~ Apparently, E2 is ~~significantly~~ dramatically amplified by the mean flow instabilities at the middle-high latitudes between  $\sim 40$  km and  $\sim 70$  km, with EP flux propagating into the polar lower atmosphere, and EP flux eventually propagate toward the equator at  $\sim 50$  km. E2 is amplified and propagated by the wave-mean flow interaction near the critical layers ( $\sim 36$  h,  $\sim 53$  h), and the positive refractive index region provides the promoting effect. The weak instability and strong background wind at  $\sim 50\text{-}60^\circ\text{N}$  and  $\sim 60\text{-}70$  km provide the energy for the propagation and amplification of EP flux into the polar lower atmosphere during days 25-30. ~~Moreover~~ Besides, E2 obtains sufficient energy from strong instability and suitable background wind at  $\sim 50\text{-}60^\circ\text{N}$  and  $\sim 60\text{-}70$  km on days 69-74, and it is amplified and propagated into the polar lower atmosphere. The background wind at  $\sim 50\text{-}60^\circ\text{N}$  and  $\sim 60\text{-}70$  km on days

127-130 is similar to on days 213-216, and the instability at  $\sim 50\text{-}60^\circ\text{S}$  and  $\sim 60\text{-}70$  km is stronger on days 127-30 than on days 213-216. Although the background wind at  $\sim 50\text{-}60^\circ\text{N}$  and  $\sim 60\text{-}70$  km is stronger on days 25-30 than on days 69-74, but the instability at  $\sim 50\text{-}60^\circ\text{N}$  and  $\sim 60\text{-}70$  km is stronger on days 69-74 than on days 25-30. According to Figure 10a, 10b, the temperature amplitude results indicates that E2 absorbs sufficient energy to be amplified under the background conditions on days 69-74, with stronger temperature amplitude on days 69-74 is stronger (Figures 10a, 10b).

Figures 11a and 11b show the observed spectra of E3 at  $\sim 59.2$  km and  $\sim 60\text{-}70^\circ\text{N}$  on days 35-38 and 53-56, and the wave period of locked wavenumber -3 is  $\sim 29$  h and  $\sim 27$  h, respectively. The corresponding temperature spatial structures of these E3 (i.e.,  $\sim 29$  h and  $\sim 27$  h) are shown in Figure 11c, 11d. The temperature spatial structure of E3 shows an obvious amplitude bimodal structure at  $\sim 50\text{-}60^\circ\text{N}$  and  $\sim 60$  km, and  $\sim 50\text{-}60^\circ\text{N}$  and  $\sim 70$  km, with the maximum at  $\sim 50\text{-}60^\circ\text{N}$  and  $\sim 60$  km. The strongest temperature amplitude of E3 occurs at  $\sim 60$  km and  $\sim 50\text{-}60^\circ\text{N}$  with an amplitude of  $\sim 6$  K on the days 35-38, and the other peak is  $\sim 5$  K ( $\sim 50\text{-}60^\circ\text{N}$  and  $\sim 70$  km). The temperature amplitude of  $\sim 4$  K occurs at  $\sim 60$  km ( $\sim 70$  km) and  $\sim 50\text{-}60^\circ\text{N}$  during days 53-56. The corresponding spatial structures of zonal wind and meridional

wind of these E3 ~~are illustratedis shown~~ in Figures 6e, 6f, 6g, and 6h. The zonal wind spatial structure of E3 shows an obvious amplitude bimodal structure at  $\sim 40\text{-}50^\circ\text{N}$  and  $\sim 70$  km, and  $\sim 60\text{-}70^\circ\text{N}$  and  $\sim 70$  km. The zonal wind amplitudes of E3 occur at  $\sim 40\text{-}50^\circ\text{N}$  and  $\sim 70$  km with an amplitude of  $\sim 15$  m/s on days 35-38, and  $\sim 12$  m/s at  $\sim 60\text{-}70^\circ\text{N}$  and  $\sim 70$  km. The maximum zonal wind amplitude of E3 ~~appearsoccurs~~ at  $\sim 40\text{-}50^\circ\text{N}$  and  $\sim 70$  km ( $\sim 60\text{-}70^\circ\text{N}$  and  $\sim 70$  km) with an amplitude of  $\sim 7$  m/s ( $\sim 6$  m/s) on days 53-56. The amplitude of E3 meridional wind reaches  $\sim 22$  m/s at  $\sim 50\text{-}60^\circ\text{N}$  and  $\sim 70$  km (days 35-38) and  $\sim 12$  m/s at  $\sim 60\text{-}70^\circ\text{N}$  and  $\sim 70$  km (days 53-56), respectively.

Obviously, the instability and appropriate background wind at ~~the mid-latitudesmid-latitude~~ between  $\sim 50$  km and  $\sim 70$  km and the interaction near the critical layers ( $\sim 29$  h,  $\sim 27$  h) dramatically amplify the propagation of E3, ~~as shown in~~ (see Figures 11i and 11j). The background wind is similar on days 35-38 and 53-56, and the former is ~~relativelymore~~ unstable. This finding indicates that the E3 in propagation is more likely to ~~gatherget~~ sufficient energy to be amplified on days 35-38. The instability and appropriate background wind at ~~the~~ mid-high latitudes between  $\sim 50$  km and  $\sim 70$  km ~~drasticallydramatically~~ amplify the propagation of E3, which is enhanced by the interaction near the critical layers ( $\sim 29$  h,  $\sim 27$  h) and the positive refractive index region. ~~(FiguresFigure~~ 11i, 11j). ~~In particular,It is worth noting that~~ the strong

instability and weak background wind at  $\sim 50\text{-}60^\circ\text{N}$  and  $\sim 60\text{-}70$  km on days 35-38 ~~generate~~provide sufficient energy for the propagation and amplification of EP flux into the lower atmosphere, and ultimately point ~~toward~~towards the equator at 50 km. The EP flux ~~propagates~~propagate to the lower atmosphere during days 35-38, and ~~it~~ is amplified by interactions at the critical layer ( $\sim 29$  h). In addition, weak instability and weak background winds on days 53-56 at  $\sim 50\text{-}60^\circ\text{N}$  and  $\sim 60\text{-}70$  km provide the energy to amplify E3 propagation. Combine with ~~Figures~~Figure 11c, 11d, the stronger the instability at  $\sim 50\text{-}60^\circ\text{N}$  and  $\sim 60\text{-}70$  km, the stronger the temperature amplitude of E3. The results show that the instability on days 35-38 at  $\sim 50\text{-}60^\circ\text{N}$  and  $\sim 60\text{-}70$  km are the primary reasons for the propagation and amplification of EP flux into the lower atmosphere.

### 3.3 Comparison between SH and NH

The observed latitude and maximum amplitude for eastward planetary waves (~~i.e.~~, E1, E2, E3, E4) decrease and weaken with increasing zonal wavenumber in the SH, reaching  $\sim 70\text{-}80^\circ\text{S}$ ,  $\sim 60\text{-}70^\circ\text{S}$ ,  $\sim 60\text{-}70^\circ\text{S}$ , and  $\sim 50\text{-}60^\circ\text{S}$ , and  $\sim 10$  K,  $\sim 9$  K,  $\sim 6$  K, and  $\sim 3$  K, respectively. In addition, the occurrence date ~~gets~~is earlier with increasing zonal wavenumber. The temperature spatial structure ~~demonstrates~~shows a bimodal-peak structure ( $\sim 50$  and  $\sim 60$  km), ~~mainly~~primarily located at  $\sim 50$  km. The maximum zonal wind amplitudes of E1 and E2, E3 and E4 are

almost the ~~same~~equivalents, which are  $\sim 20$  m/s and  $\sim 10$  m/s, respectively. The maximum meridional wind amplitudes of E1, E2, E3 and E4 are  $\sim 17$  m/s,  $\sim 27$  m/s,  $\sim 16$  m/s, and  $\sim 11$  m/s, respectively. The wave period of E1 tends to get shorter from 5 to 3 days, while E2 and E3 are close to  $\sim 40$  h and  $\sim 30$  h, ~~and~~while E4 remains at  $\sim 24$  h. E1, E2, E3, and E4 are more favorable to propagation in the SH winter and are ~~abruptly~~dramatically amplified by the mean flow instabilities at ~~the~~ middle latitudes between  $\sim 40$  km and  $\sim 70$  km, ~~with~~. ~~With the propagation of~~ EP flux ~~propagating~~ into the lower atmosphere, and ~~it finally propagates~~ EP flux ~~eventually propagate~~ toward the equator at  $\sim 50$  km. In addition, the propagation of EP flux for E1 to the upper atmosphere ~~might~~may be influenced by the instability and background wind at the Antarctic  $\sim 50$  km.

The observed latitudes of E1, E2 (E3) decrease with increasing wavenumber in the NH, which are  $\sim 70$ - $80^\circ$ N,  $\sim 60$ - $70^\circ$ N, and  $\sim 60$ - $70^\circ$ N. ~~With bimodal-peak structure located at  $\sim 70$  km, the~~ The temperature spatial ~~structures~~structure of E1, E2, and E3 ~~presents a bimodal-peak structure, primarily located at  $\sim 70$  km, reaches~~reach  $\sim 10$  K,  $\sim 9$  K, and  $\sim 6$  K, respectively. The maximum zonal wind amplitude for E1, E2 and E3 ~~occure~~occurs at  $\sim 50$ - $80^\circ$ N and  $\sim 70$  km, and their amplitude ~~are~~is almost equal to  $\sim 18$  m/s. The maximum meridional ~~winds~~wind of E1, E2 and E3 occur at  $\sim 50$ - $80^\circ$ N and  $\sim 70$  km, with amplitudes of  $\sim 22$  m/s,  $\sim 18$  m/s and  $\sim 22$  m/s, respectively. The wave period of E1 tends to be shorter from 5-3



days, and E2 and E3 are close to  $\sim 48$  h and  $\sim 30$  h. In addition, E1, E2, and E3 are more favorable to propagation in the NH winter and are dramatically amplified by the mean flow instabilities at the middle latitudes between  $\sim 40^\circ$  and  $\sim 70^\circ$  km, with the propagation of EP flux propagating into the lower atmosphere and then, and EP flux eventually propagate toward the equator at  $\sim 50^\circ$  km.

#### 4 Summary and Conclusions

Based on the MERRA-2 temperature and wind observations in 2019, we present for the first time an extensive study of the global variation for eastward planetary wave activity, including zonal wave numbers of -1 (E1), -2 (E2), -3 (E3), -4 (E4), in the stratosphere and mesosphere, using the MERRA-2 temperature and wind observations in 2019. The temperature and wind amplitudes and wave periods of each event were determined using 2-D least-squares fitting. Our study covered the spatial and temporal patterns of the eastward planetary waves in both hemispheres with a comprehensive diagnostic analysis on their propagation and amplification. The key findings of this study are summarized below as follows:

1. The latitudes for the maximum (temperature, zonal and meridional wind) amplitudes of E1, E2, E3, and E4 decrease with increasing wavenumber in the SH and NH. The E1, E2, E3, and E4 events occur earlier with increasing zonal wavenumber in the SH. In addition,

eastward wave modes exist during summer periods with westward background wind in both hemispheres.

2. The temperature spatial structures of E1, E2, E3, and E4 present a double-peak structure, which is located at ~50 km and ~60 km in SH, ~60 km, and ~70 km in NH. Furthermore, the lower peak is usually larger than the higher one.

3. The maximum (temperature, zonal and meridional wind) amplitudes of E1, E2, and E3 decrease with increasing zonal wavenumber in the SH and NH. The maximum temperature amplitudes in the SH are slightly larger and lie lower than those in the NH. In addition, the meridional wind amplitudes are slightly larger than the zonal wind in the SH and NH.

4. The wave period of the E1 mode ranges from 3 to 5 days in both hemispheres, while the period of E2 mode is slightly longer in the NH (~48 h) than in the SH (~40 h). The periods of E3 are ~30 h in both SH and NH, while the period of E4 is ~24 h.

5. The eastward planetary wave is more favorable to propagate in the winter hemisphere and is dramatically amplified by the mean flow instabilities and appropriate background winds at polar region and the middle latitudes between ~40 km and ~80 km. Furthermore, the amplification of planetary waves through wave-mean flow interaction

~~occurs~~<sup>most</sup> easily ~~occurs near~~<sup>close to</sup> its critical layer. In addition, the direction of EP flux ultimately points towards the equator.

6. The strong instability and appropriate background wind in the lower layer of the Antarctic region ~~might generate adequate~~<sup>may provide sufficient</sup> energy to promote the E1 propagation and amplification to the upper atmosphere.

~~Overall, this study demonstrated~~<sup>This study demonstrates</sup> how the background zonal wind in the polar middle atmosphere affects the dynamics of eastward planetary waves in the polar middle atmosphere.

717

718 *Data availability.* MERRA-2 data are available at  
719 <http://disc.gsfc.nasa.gov>.

720

721 *Code availability.* Code is available at  
722 <http://hdl.pid21.cn/21.86116.7/04.99.01293>.

723

724 *Author contributions.* LT carried out the data processing and analysis and  
725 wrote the manuscript. SYG and XKD contributed to reviewing the article.

726

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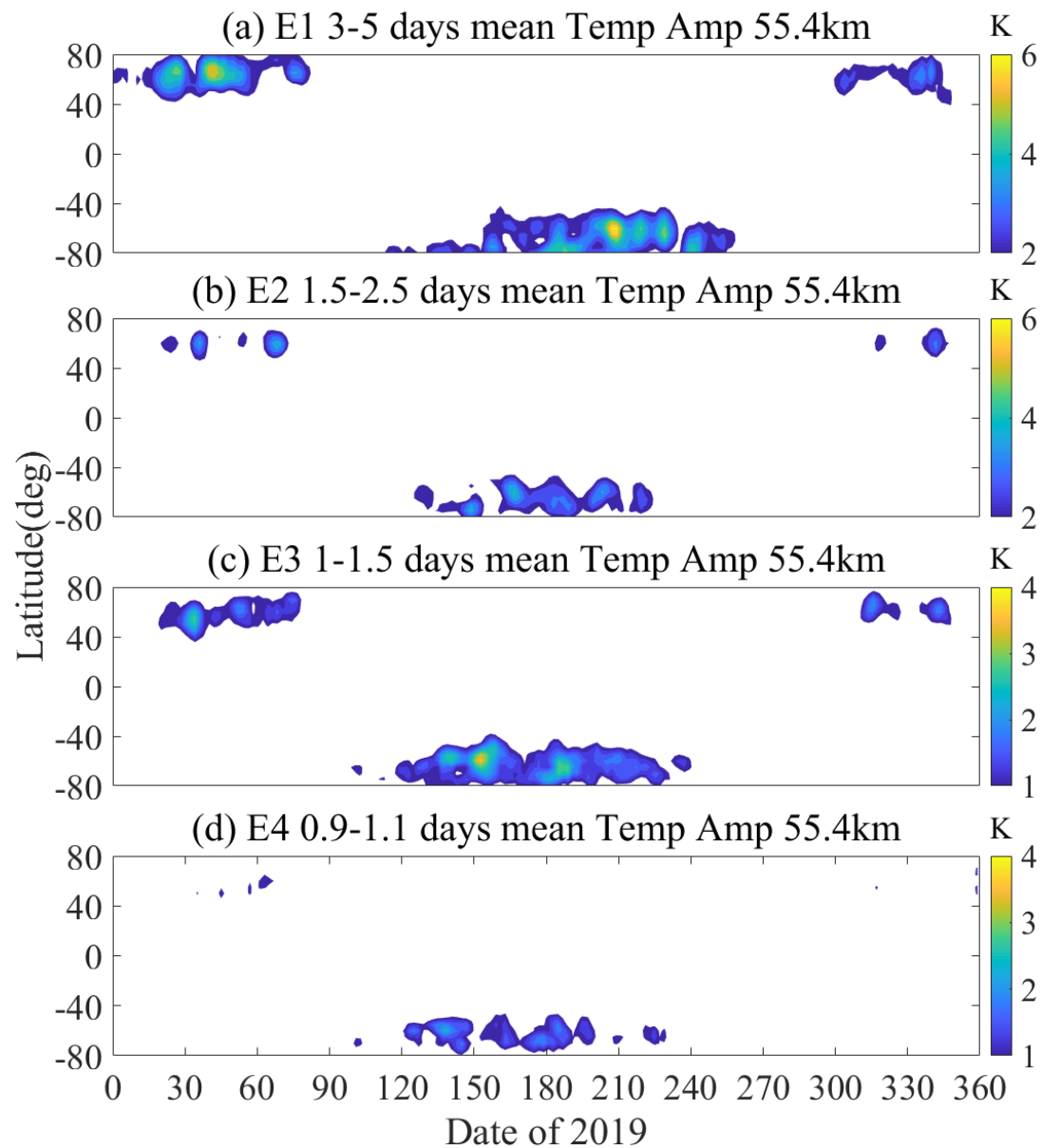
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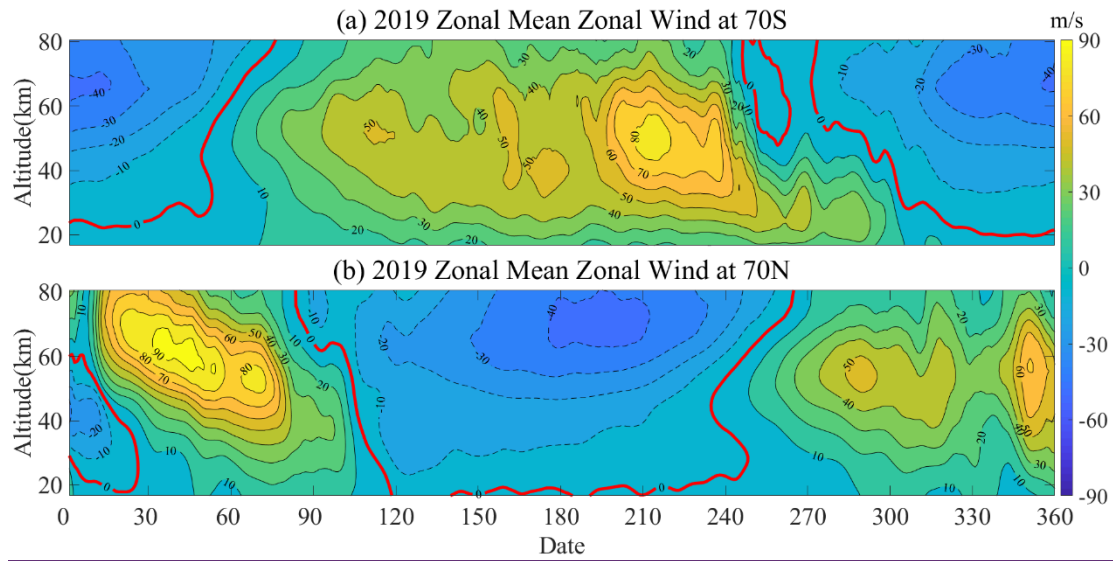


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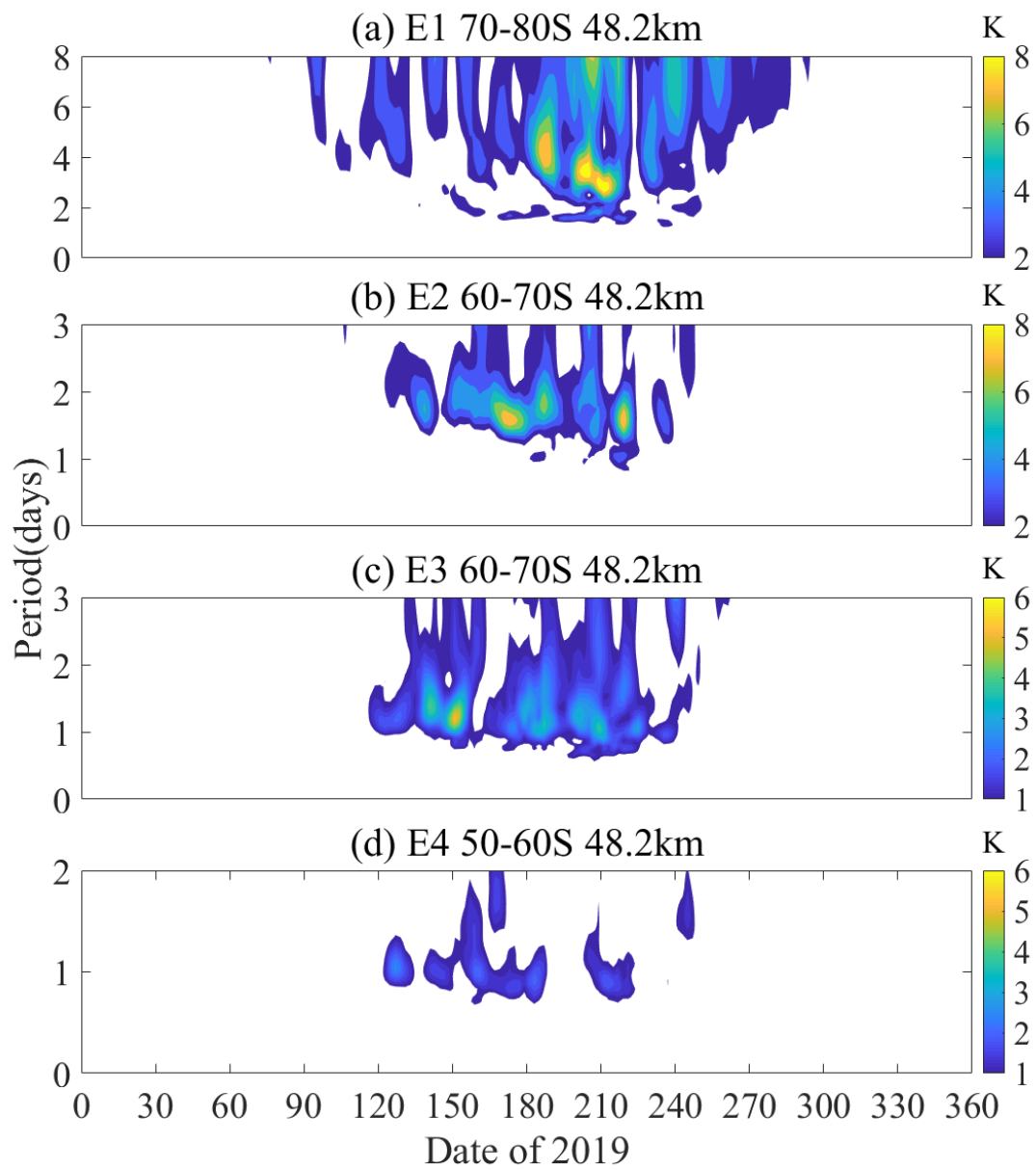
880 **Figure 1.** The global latitude-temporal variation structures of the (a) E1, (b) E2, (c)

881 E3, and (d) E4 planetary waves during 2019. White areas represent small amplitude

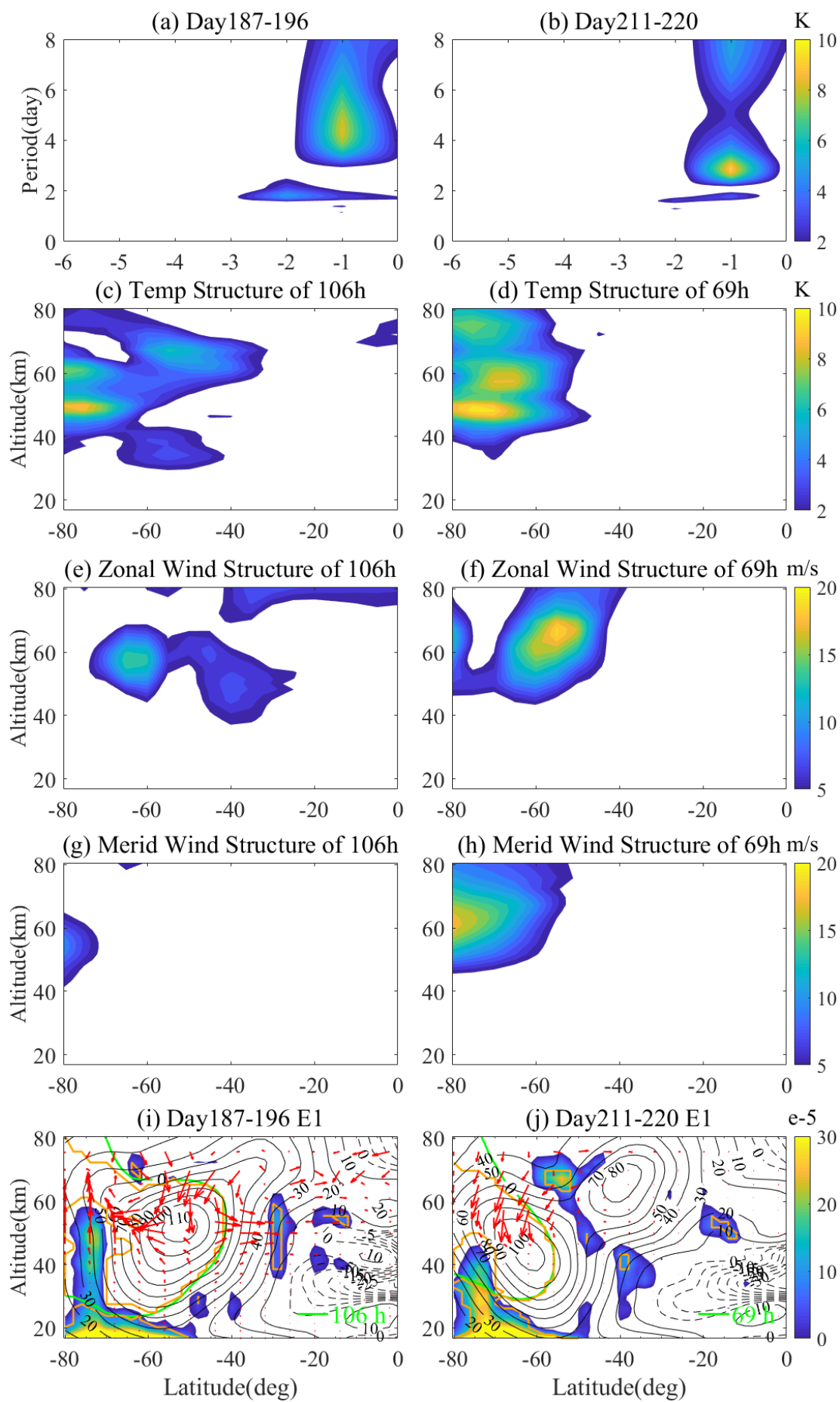
882 data (corresponds to the right color bar). The confidence interval is 0.95.



**Figure 2.** The zonal mean zonal wind variations of ~~the~~ (a) ~~the~~ 70°S and (b) 70°N during 2019. The dotted line represents east wind, the solid line represents west wind, and the red solid line is 0 m/s.

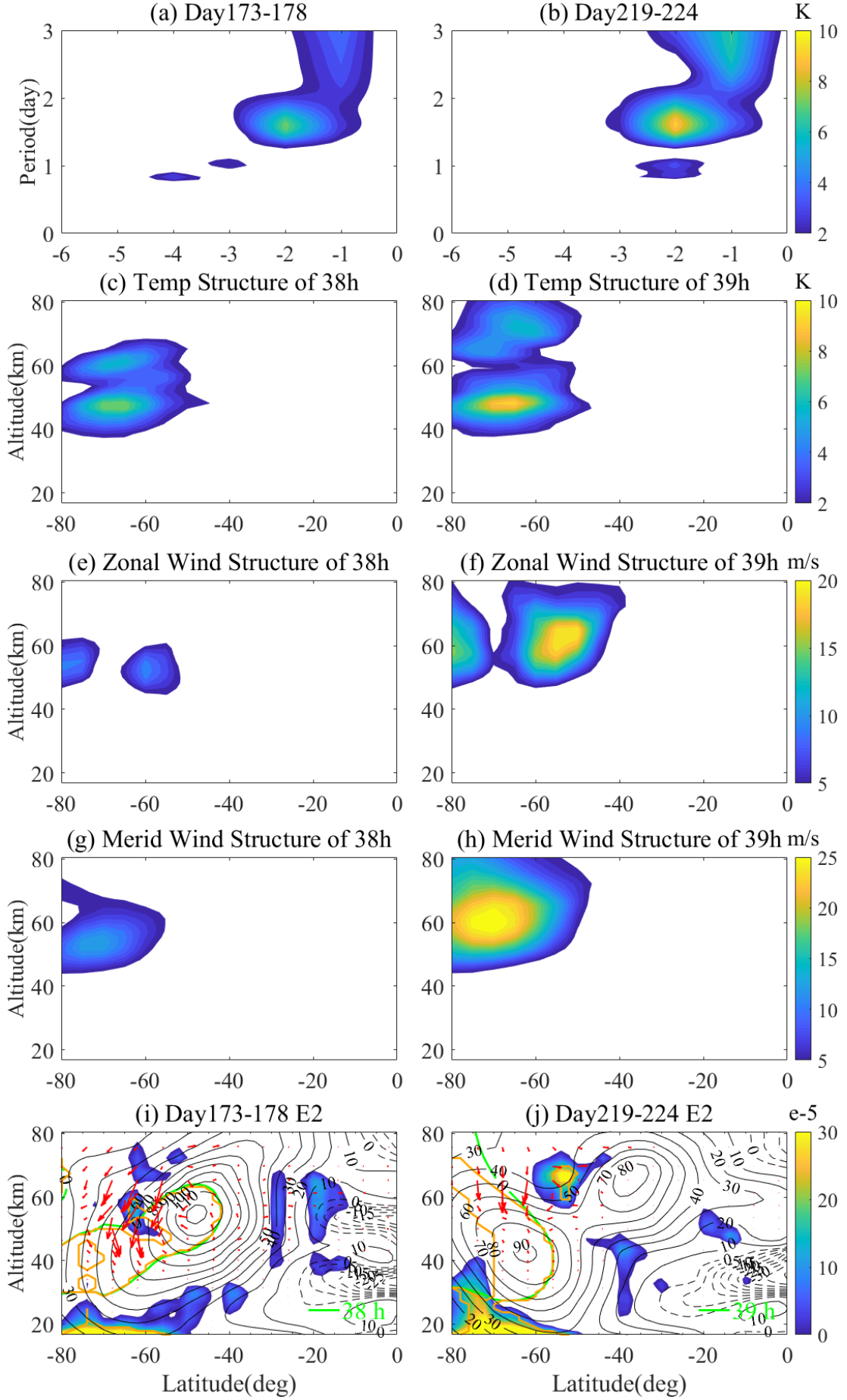


**Figure 3.** The temporal variations of the (a) E1, (b) E2, (c) E3, and (d) E4 QTDWs during 2019 austral winter period.

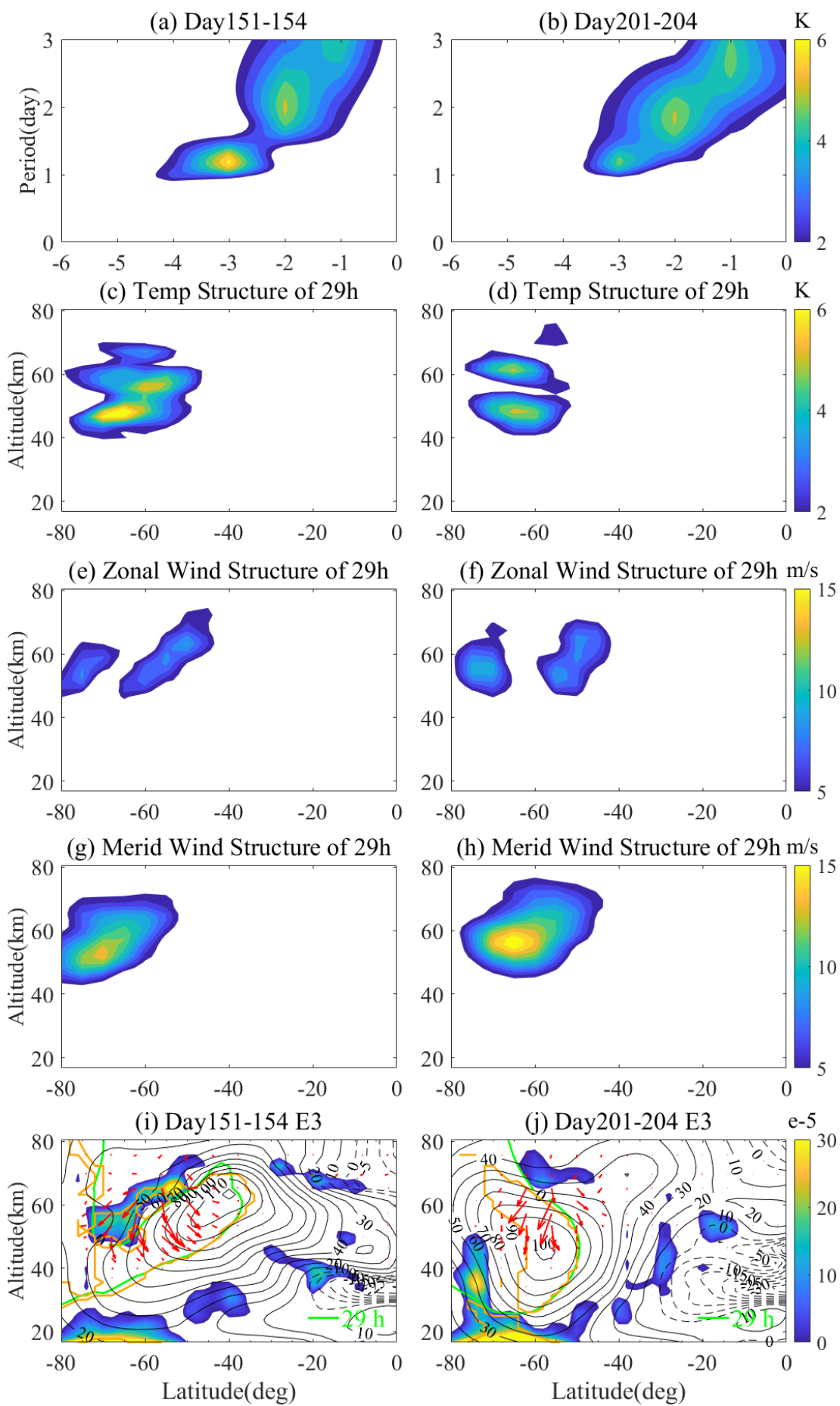




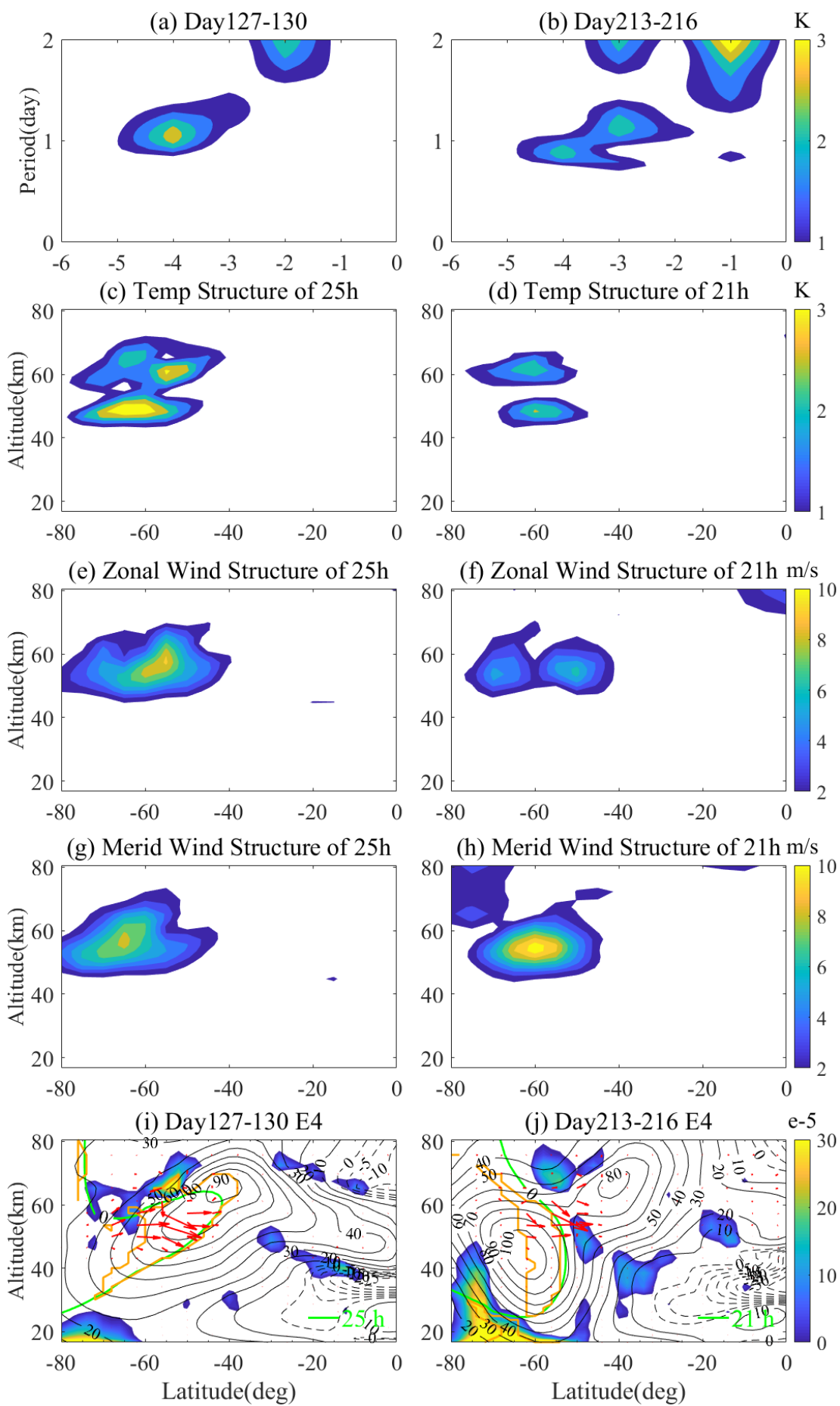
**Figure 4.** The (a, b) spectra, (c, d) temperature spatial structures, (e, f) zonal wind spatial structures, (g, h) meridional wind spatial structures, and (i, j) diagnostic analysis of the E1 typical events during 2019 austral winter period. The MERRA-2 temperature data observations at 48.2 km and 70-80°S during days 187–196 (Figure 4a), 211–220 (Figure 4d) are utilized, respectively. The instability (blue shaded region), EP fluxes (red arrow), and critical layers (green line) for E1 typical event. The green line represents critical layers of the E1 with the natural period. Regions enclosed by orange solid lines are characterized by the positive refractive index for the E1.



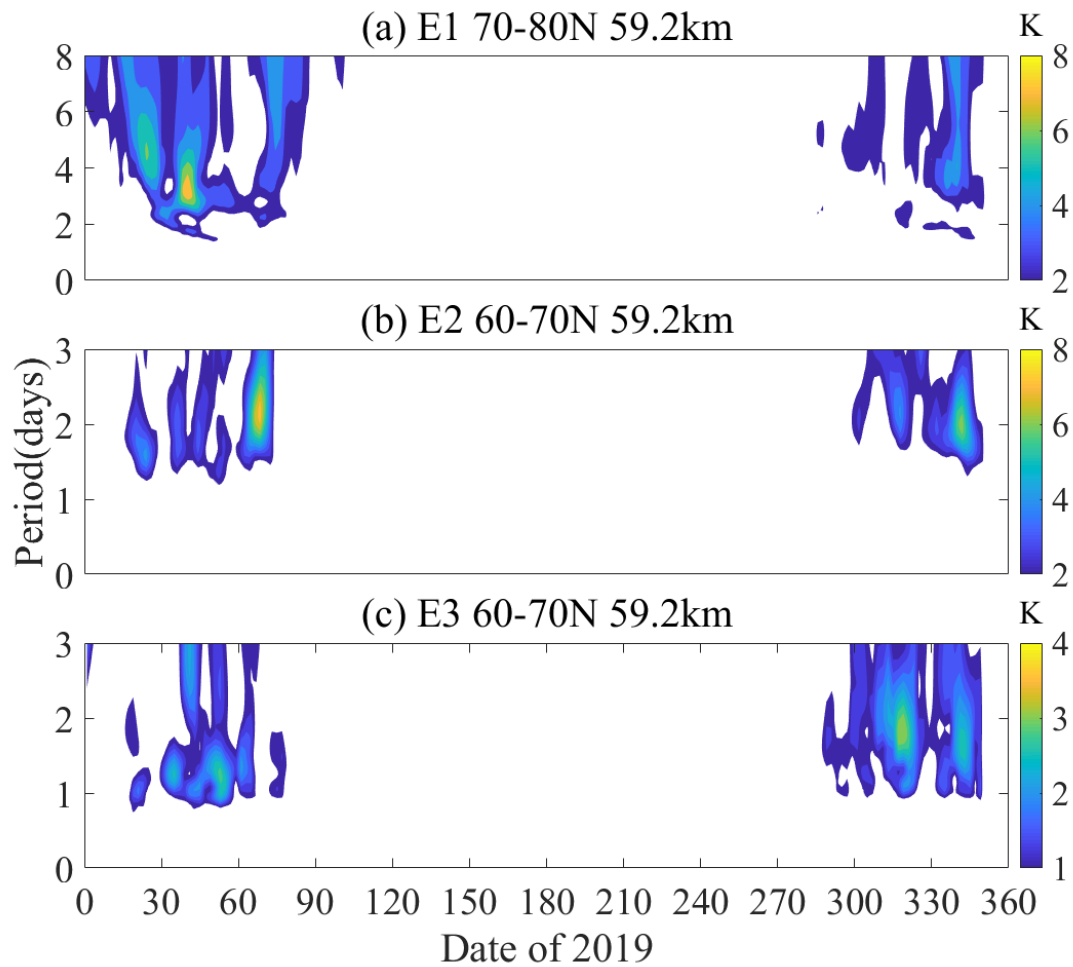
901 **Figure 5.** ~~The same~~Same as Figure 4 but for ~~the~~-E2 during the 2019 austral winter  
902 period.



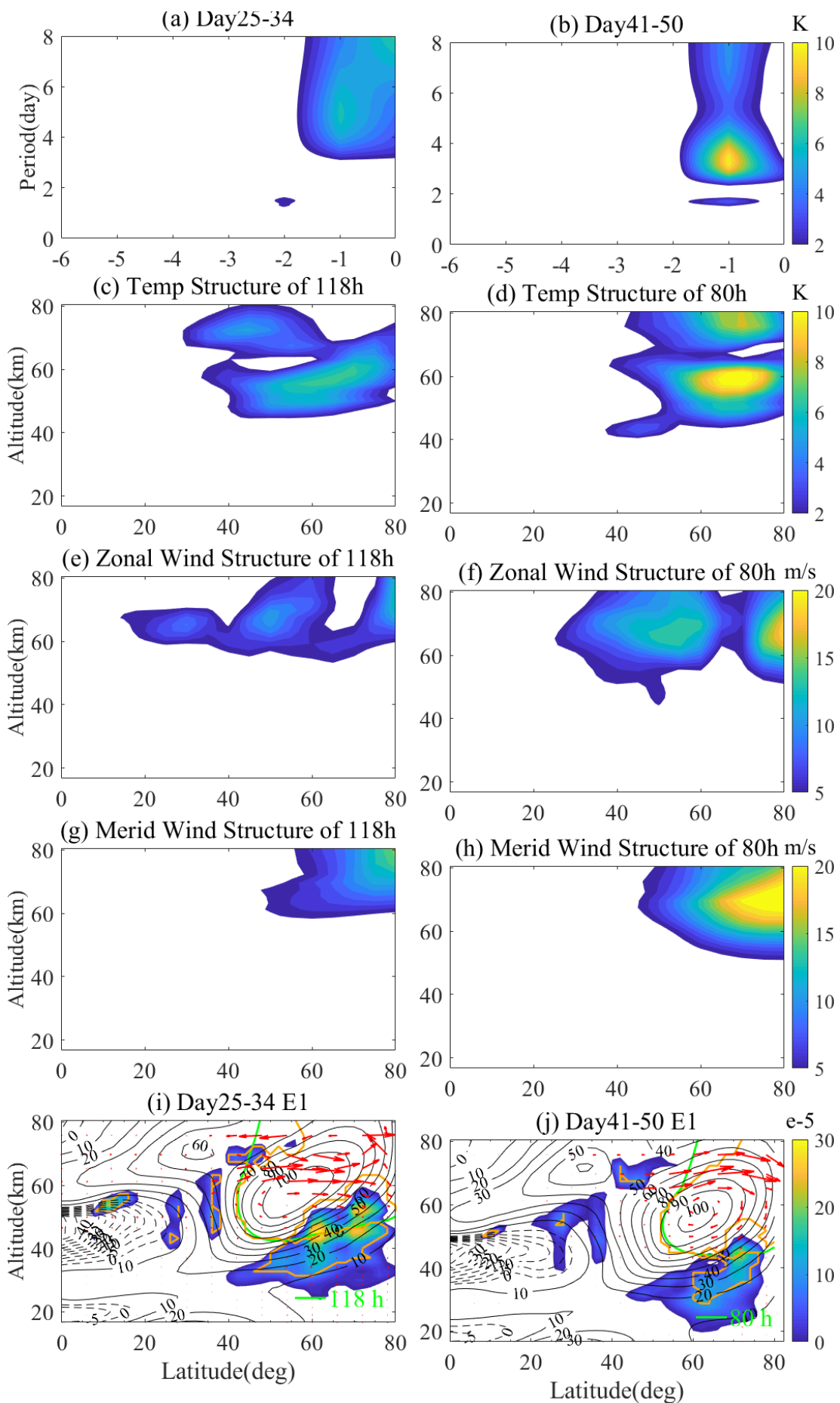
904 **Figure 6.** ~~The same~~Same as Figure 4 but for ~~the~~-E3 during the 2019 austral winter  
905 period.



**Figure 7.** ~~The same~~ Same as Figure 4 but for ~~the~~ E4 during the 2019 austral winter period.

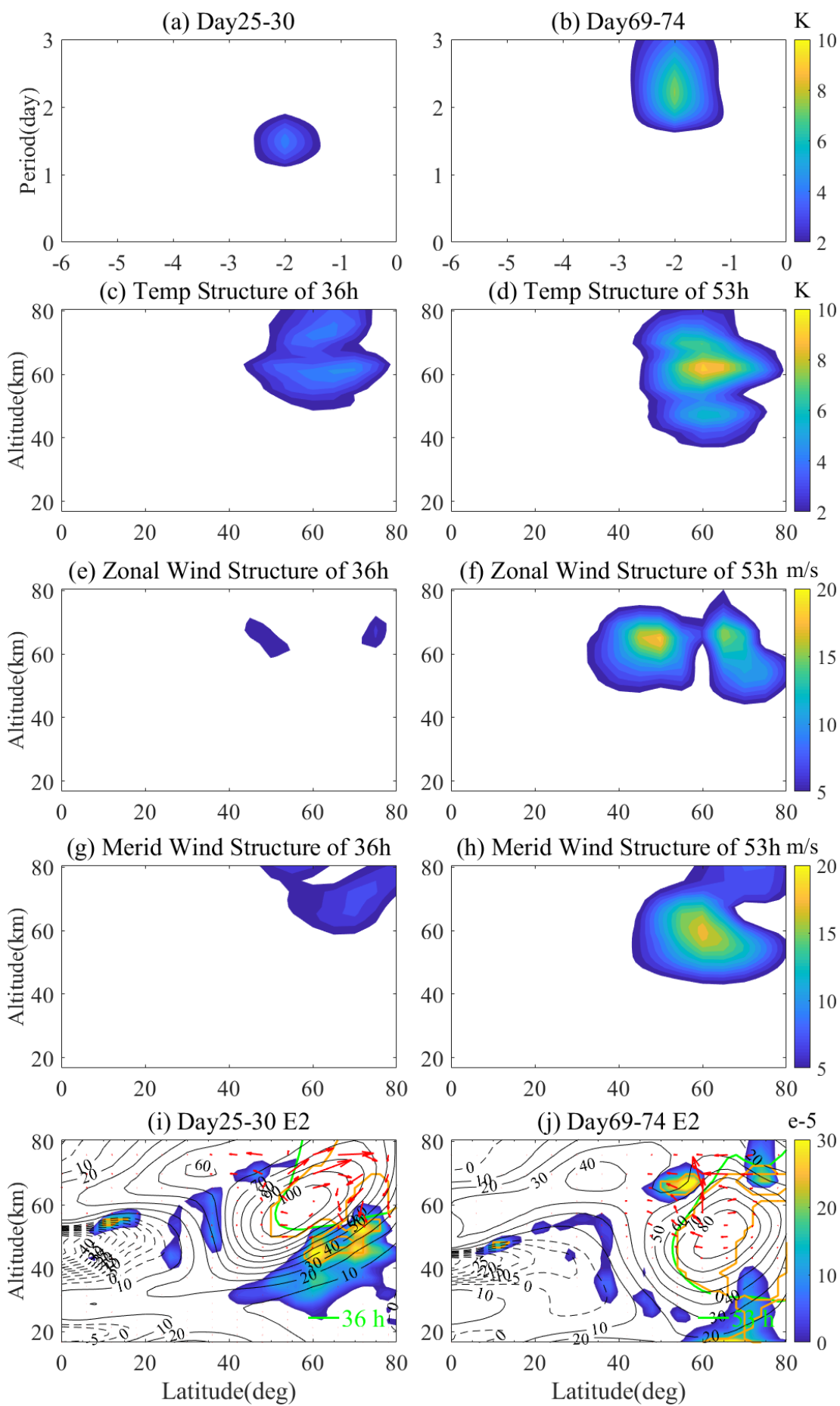


**Figure 8.** The temporal variations of ~~the~~ (a) E1, (b) E2, and (c) E3 QTDWs during the 2019 boreal winter period.

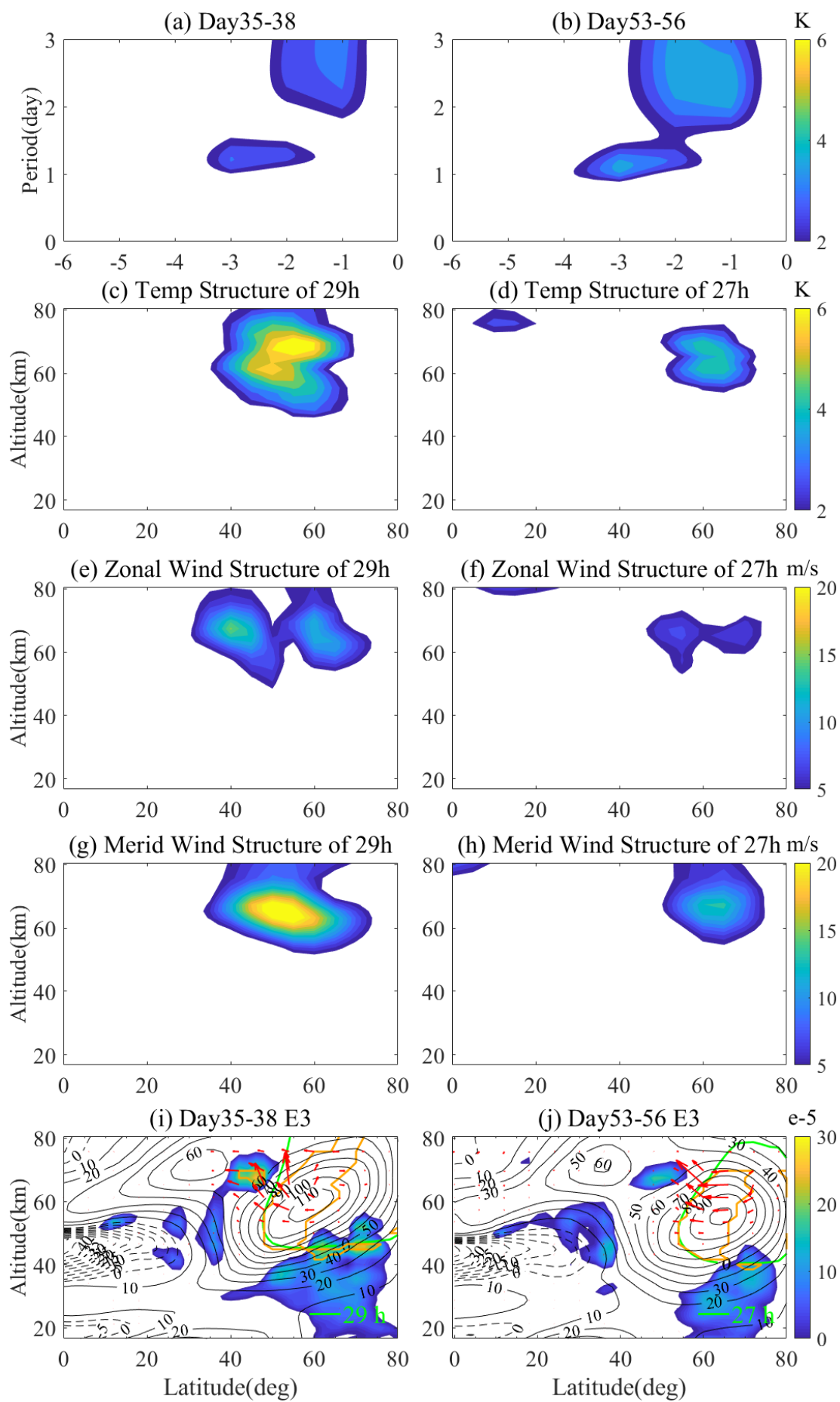




**Figure 9.** [heThe](#) (a, b) spectra, (c, d) temperature spatial structures, (e, f) zonal wind spatial structures, (g, h) meridional wind spatial structures, and (i, j) diagnostic analysis of the E1 typical events during 2019 boreal winter period. The E1 events at 48.2 km and 70-80°N were obtained from the MERRA-2 reanalysis.



918 **Figure 10.** ~~The same~~Same as Figure 9 but for ~~the~~ E2 during the 2019 boreal winter  
919 period.



921 **Figure 11.** ~~The same~~Same as Figure 9 but for ~~the~~ E3 during the 2019 boreal winter  
922 period.  
923

1     **Eastward-propagating planetary wave in the polar**  
2                     **middle atmosphere**

3  
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**Abstract.** According to MERRA-2 temperature and wind datasets in 2019, this study presented the global variations of the eastward propagating wavenumber 1 (E1), 2 (E2), 3 (E3) and 4 (E4) planetary waves (PWs) and their diagnostic results in the polar middle atmosphere. We clearly demonstrated the eastward wave modes exist during winter periods with westward background wind in both hemispheres. The maximum wave amplitudes in the southern hemisphere (SH) are slightly larger and lie lower than those in the northern hemisphere (NH). Moreover, the wave perturbations peak at lower latitudes with smaller amplitude as the wavenumber increases. The period of the E1 mode varies 3-5 days in both hemispheres, while the period of E2 mode is slightly longer in the NH ( $\sim 48$  h) than in the SH ( $\sim 40$  h). The periods of the E3 are  $\sim 30$  h in both SH and NH, and the period of E4 is  $\sim 24$  h. Despite the shortening of wave periods with the increase of wavenumber, their mean phase speeds are relatively stable, which are  $\sim 53$  m/s,  $\sim 58$  m/s,  $\sim 55$  m/s and  $\sim 52$  m/s at  $70^\circ$  latitudes for W1, W2, W3 and W4, respectively. The eastward PWs occur earlier with increasing zonal wavenumber, which agrees well with the seasonal variations of the critical layers generated by the background wind. Our diagnostic analysis also indicated that the mean flow instability in the upper stratosphere and upper mesosphere might contribute to the amplification of the eastward PWs.

## 1 Introduction

The dominance of large amplitude planetary waves in the stratosphere, mesosphere and lower thermosphere regions and their interactions with zonal mean winds are the primary driving force of atmospheric dynamics. In addition, sudden stratospheric warmings (SSWs) and quasi-biennial oscillation (QBO) events can dynamically couple the entire atmosphere from the lower atmosphere to the ionosphere (Li et al., 2020; Yamazaki et al., 2020; Yadav et al., 2019; Matthias and Ern, 2018; Stray et al., 2015). Westward propagating planetary wave is one of the prominent features during austral and boreal summer. Westward quasi-2-day waves (Q2DWs) are the most obvious representative waves and one of the most investigated phenomena using planetary wave observations. Most of the previous studies focused on the westward propagating Q2DWs, i.e., zonal wavenumbers of 2 (W2), 3 (W3) and 4 (W4) modes (Lainer et al., 2018; Gu et al., 2018b; Wang et al., 2017; Pancheva et al., 2016; Gu et al., 2016a; Gu et al., 2016b; Lilienthal and Jacobi, 2015; Gu et al., 2013; Limpasuvan and Wu, 2009; Salby, 1981). However, limited studies were conducted to understand the seasonal variations of the occurrence date, peak amplitude and wave period for the eastward Q2DWs (Gu et al., 2017; Lu et al., 2013; Alexander and Shepherd, 2010; Sandford et al., 2008; Palo et al., 2007; Merzlyakov and Pancheva, 2007; Manney and Randel, 1993; Venne and



Stanford, 1979).

Typically, Q2DWs maximize after the summer solstice in the middle latitudes. The largest wave amplitudes generally appear near the mesopause in January–February in the Southern Hemisphere (SH), while in the Northern Hemisphere (NH) in July–August (Tunbridge et al., 2011). W3 and W4 Q2DWs reach amplitudes during austral and boreal summer in the mesosphere and lower thermosphere, respectively. The seasonal variation of westward Q2DWs activity is obvious (Liu et al., 2019; Gu et al., 2018b; Rao et al., 2017). By observing the long-term Q2DW in the NH and SH, Tunbridge et al. (2011) reported that W3 is generally stronger than the other two modes in the SH, reaching the amplitude of  $\sim 12\text{K}$ ; while W4 is stronger than W3 in the NH, reaching  $\sim 4\text{K}$ . Moreover, W4 generally lives longer than W3, and W4 can still be observed after the ending of W3. A previous study has demonstrated that wave source, instability, critical layer and mean zonal wind are the primary reasons for the seasonal variation of Q2DWs (Liu et al., 2004). By studying the long-term satellite datasets in the SH, Gu et al. (2019) suggested that the strongest events of W2, W3 and W4 could be delayed with increasing the zonal wavenumber, and these events would be indistinguishable during SSWs. The wave periods of W4, W3 and W2 vary around  $\sim 41\text{-}56\text{ h}$ ,  $\sim 45\text{-}52\text{ h}$ , and  $\sim 45\text{-}48\text{ h}$ , respectively. Furthermore, W2 can be observed using global satellite datasets, but it has an amplitude weaker than W3

and W4 in the NH and SH (Meek et al., 1996). The propagation and amplification of Q2DWs are primarily modulated by instability, refractive index and critical layer, while the variation of background wind may cause different zonal wavenumber events (Gu et al., 2016a; Gu et al., 2016b). By analyzing the variation of Q2DWs activity during SSWs, Xiong et al. (2018) noticed that W1 is generated by the nonlinear interaction between SPW2 and W3. During SSWs, the coupling between the NH and SH can enhance the summer easterly and promote the nonlinear interaction between W3 and SPW1 (Gu et al., 2018b).

Some recent studies have discovered significant eastward planetary waves in the polar stratosphere and mesosphere regions, with periods of nearly two and four days (Gu et al., 2017; Sandford et al., 2008; Merzlyakov and Pancheva, 2007; Coy et al., 2003; Manney and Randel, 1993). Planetary waves with zonal wavenumbers -1 (E1) and -2 (E2) correspond to 4- and 2-day waves, respectively. Furthermore, planetary waves of 1.2-day with wavenumber -3 (E3) and 0.8-day with wavenumber -4 (E4) have been reported to contain the same phase speeds as E1 and E2 (Manney and Randel, 1993). This series of eastward planetary wave can significantly affect the thermal and dynamic structure of the polar stratosphere, resulting in profound changes in the wind and temperature of the polar stratosphere (Coy et al., 2003; Venne and Stanford, 1979). Beyond the knowledge about nonlinear interactions

97 between migrating tides and Q2DWs (Palo et al., 1999), further  
98 investigation has confirmed that E2 Q2DW ~~is—ecouldoupled—~~ be  
99 generated by the nonlinear interaction between planetary wave and tides  
100 in the mesosphere and lower thermosphere (MLT) (Palo et al., 2007). We  
101 should note that the E2 Q2DW generated in the MLT region is different  
102 from that in the polar stratosphere, which is discussed in this paper.

103 By studying and analyzing satellite datasets, Merzlyakov and  
104 Pancheva. (2007) indicated that the wave periods of E1 and E2 events  
105 range 1.5-5 days. They reported that EP flux travels from the upper to the  
106 lower atmosphere, meaning that the upper atmosphere has a dynamic  
107 influence on the lower atmosphere. Sandford et al. (2008) reported about  
108 significant fluctuations of E2 Q2DW in the polar mesosphere. They  
109 indicated the influence of changes in mean zonal winds during a major  
110 SSW on the propagation of polar E2. In addition, they proposed the  
111 significance of E2 fluctuation in the mesosphere driven by the  
112 instabilities in the polar night jet. For E2, amplitude of temperature, zonal  
113 wind and meridional wind during the austral winter can reach ~10 K, ~20  
114 m/s and ~30 m/s, respectively; while those during the boreal winter can  
115 drop by almost two-thirds. Lu et al. (2013) found that eastward planetary  
116 wave propagation is limited to the winter high latitudes probably because  
117 the negative refractive indices equatorward of ~45°S result in evanescent  
118 wave characteristics. That study suggested that the instability region at

~50-60°S might be induced by the stratospheric polar night jet and/or the "double-jet" structure.

In this study, we use the second modern retrospective research and application analysis (MERRA-2) datasets to investigate the eastward propagating wave characteristics of the stratosphere and mesosphere in polar region in 2019, including E1, E2, E3 and E4. Specifically, we investigate the variation of the occurrence date, peak amplitude and wave period of eastward waves; as well as the role of instability, background wind structure and critical layer in the propagation and amplification of eastward waves. The remaining parts of this paper are organized as follows. Section 2 describes the data and methods used in this study. Section 3 analyzes the global latitude-temporal variation structure of eastward waves during winter in 2019. The amplification and propagation features of the eastward planetary waves in the NH and SH with different wavenumber events are examined in Sections 3.1 and 3.2, respectively. Section 3.3 compares and analyzes the eastward waves in the NH and SH. All research results are summarized in Section 4.

## **2 Data and Analysis**

To extract the E1-, E2-, E3- and E4-wave, we apply the least-square method to each time window (e.g., 10-day, 6-day, 4-day and 4-day), and then use time window to determine the amplitude. (Gu et al., 2013). This method has been shown to successfully identify planetary waves from

satellite measurements (Gu et al., 2019; Gu et al., 2018a; Gu et al., 2018b; Gu et al., 2018c; Gu et al., 2013).

$$y = A \cos[2\pi(\sigma \cdot t + s \cdot \lambda)] + B \sin[2\pi(\sigma \cdot t + s \cdot \lambda)] + C \quad (1)$$

The least-squares method is used to fit the a set of parameters ( $A$ ,  $B$  and  $C$ ), where  $\sigma$ ,  $t$ ,  $s$  and  $\lambda$  are the frequency, UT time, zonal wavenumber and longitudes. The amplitude of wave  $R$  can be expressed as  $R = \sqrt{A^2 + B^2}$ .

The second modern retrospective research and application analysis (MERRA-2) covers the long-term atmospheric reanalysis datasets initiated by NASA in 1980. It has been upgraded recently using the Goddard Earth Observing System Model, Version 5 (GEOS-5) data assimilation system. Briefly, MERRA-2 includes some updates to the model (Molod et al., 2014; Molod et al., 2012) and the Global Statistical Interpolation (GSI) analysis scheme of Wu et al. (2002). The MERRA-2 data consist of various meteorological variables, e.g., net radiation, temperature, relative humidity and wind speed. The spatial coverage of MERRA-2 data is the globe (spatial resolution:  $0.5^\circ \times 0.625^\circ$ ; temporal resolution: 1 h). This meteorological data are widely used to detect the middle atmosphere such as the planetary wave in the polar atmosphere, global thermal tides, climate variability and aerosol (Ukhov et al., 2020; Sun et al., 2020; Bali et al., 2019; Lu et al., 2013). Many recent studies indicated the feasibility of using MERRA-2 data for the kind of research

in present study. Therefore, we apply the MERRA-2 datasets to obtain the variation in background wind, instability, refractive index and critical layer; and explore the patterns of eastward planetary waves propagation and amplification through diagnostic analysis.

The critical layer will absorb or reflect planetary waves from the lower atmosphere during upward propagation. Planetary waves that gain sufficient energy in the unstable region will be amplified during reflection. In a sense, the critical layer plays an important role in regulating the amplification and propagation of planetary waves (Gu et al., 2016a; Gu et al., 2016b; Liu et al., 2004).

$$\overline{q_\varphi} = 2\Omega \cos \varphi - \left( \frac{(\overline{u} \cos \varphi)_\varphi}{a \cos \varphi} \right)_\varphi - \frac{a}{\rho} \left( \frac{f^2}{N^2} \rho \overline{u}_z \right)_z \quad (2)$$

The baroclinic/barotropic instability in the atmospheric space structure is caused by the simultaneous equalization of the negative latitude gradient and the quasi-geostrophic potential vorticity ( $\overline{q_\varphi}$ ). In Equation (2),  $\Omega$  is the angular speed of the Earth's rotation;  $\varphi$  is the latitude;  $\overline{u}$  is the zonal mean zonal wind;  $a$  is the Earth radius;  $\rho$  is the air density;  $f$  is the Coriolis parameter;  $N$  is the buoyancy frequency; subscripts  $z$  and  $\varphi$  are the vertical and latitudinal gradients.

According to Andrews et al. (1987), the properties of planetary wave propagation can be calculated using the Eliassen-Palm (EP) flux vectors (F), i.e.,

$$F = \rho a \cos \varphi \left[ \frac{\overline{u_z v' \theta'}}{\overline{\theta_z}} - \overline{v' u'} \right] \left[ f - \frac{(\overline{u \cos \varphi})_\varphi}{a \cos \varphi} \right] \frac{\overline{v' \theta'}}{\overline{\theta_z}} - \overline{w' u'} \quad (3)$$

where  $u'$  and  $v'$  are the planetary wave perturbations in the zonal and meridional wind, respectively;  $\theta'$  and  $w'$  are the potential temperature and vertical wind, respectively. The planetary wave propagation is only favorable where the square of refractive index  $m^2$  is positive:

$$m^2 = \frac{\overline{q_\varphi}}{a(\overline{u} - c)} - \frac{s^2}{(a \cos \varphi)^2} - \frac{f^2}{4N^2 H^2} \quad (4)$$

where  $s$  is the zonal wavenumber,  $c$  is the phase speed and  $H$  is the scale height. The square of the refractive index is taken as the waveguide of planetary waves, i.e.,

$$c = -v_0 \cos\left(\frac{\varphi\pi}{180}\right) / sT \quad (5)$$

where  $v_0$  is the equatorial linear velocity,  $s$  is the zonal wavenumber and  $T$  is the wave period.

### 3 Results and Discussion

Figure 1 shows the global temporal-latitude variation structures of E1, E2, E3 and E4 extracted from the 2019 MERRA-2 temperature datasets using time windows 10-, 6-, 4- and 4-days, respectively. The mean temperature amplitude of E1, E2, E3 and E4 at 55.4 km during the periods 3~5-, 1.5~2.5-, 1~1.5- and 0.9~1.1-day are displayed in Figure 1a, 1b, 1c and 1d, respectively. Eastward waves are characterized by obvious

203 seasonal variations in the SH and NH. In addition, E1, E2 (E3) and E4  
 204 reach their maximum amplitude at 50-80°(S/N). In the SH, the strongest  
 205 E1 and E2 events occur on days 209-218 and 167-172; while E3 and E4  
 206 events occur on days 151-154 and 139-142. This means that their  
 207 occurrence date of maximum amplitude gets earlier with increasing zonal  
 208 wavenumber. In addition, the maximum amplitude of E1, E2, E3 and E4  
 209 are ~6.0 K, ~4.2 K, ~3.6 K and ~2.4 K, respectively, indicating that their  
 210 peak amplitude drop with rising zonal wavenumber. In the NH, the  
 211 strongest E1, E2, E3 and E4 events occur on days 41-50, 69-74, 35-38  
 212 and 63-66, respectively; the corresponding peak amplitude are ~5.5 K,  
 213 ~3.8 K, ~2.8 K and ~1.2 K, respectively. Whilst the results demonstrate  
 214 the decline of the peak amplitude with increasing zonal wavenumber in  
 215 the NH, the occurrence date is irregular. Moreover, E4 is relatively weak  
 216 in the NH and difficult to find, so W4 is insignificant in the NH. Figure 2  
 217 presents the changes in zonal mean zonal wind at 70°S and 70°N in 2019.  
 218 It can be seen that the background wind on days 90-240 (70°S) is  
 219 dominated by westward wind, and reaches ~80 m/s at ~50 km on days  
 220 210; it is dominated by eastward wind in late and early 2019, and reaches  
 221 ~-40 m/s at ~60 km. Meanwhile, the background wind is primarily  
 222 westerly wind in late and early 2019 (70°N), and reaches ~90 m/s at ~60  
 223 km on days 50; while on days 120-240, the background wind is primarily  
 224 easterly wind, and the amplitude reaches -40 m/s on days 200. Compared



with Figure 1, the results show that the eastward wave modes exist during winter periods with westward background wind in both hemispheres.

### 3.1 In the Southern Hemisphere

Figure 3 shows that observed maximum temperature amplitude is at  $\sim 48.2$  km and  $\sim 70$ - $80^\circ$ S for E1;  $\sim 48.2$  km and  $\sim 60$ - $70^\circ$ S for E2 and E3;  $\sim 48.2$  km and  $\sim 50$ - $60^\circ$ S for E4. For E1, the observed maximum perturbation occurs on days 211-220 (with an amplitude of  $\sim 8.5$  K), and the remaining fluctuations occur on days 161-170, 187-196 and 231-240. For E2, the observed maximum perturbation happens at days 219-224 (with an amplitude of  $\sim 7.8$  K), and three peaks appear on days 139-144, 173-178 and 187-192. Regarding E3, the strongest perturbation occurs on days 151-154 (with an amplitude of  $\sim 5.2$  K), while the rest are distributed on days 141-144, 201-204 and 209-202. E4 perturbations are distributed on days 127-130, 145-148, 161-164, 213-216, with weak amplitude of  $\sim 3$  K. Since earlier studies mentioned that the wave period of the eastward wave can vary, we also investigate the periodic variabilities of E1, E2, E3 and E4. The results show that the period corresponding to the maximum perturbation of E1 falls between  $\sim 106$  (days 187-196) and  $\sim 69$  h (days 211-220), and their wave periods vary significantly. Nonetheless, the wave period of E2 gradually changes from  $\sim 42$  h (days 139-144) to  $\sim 38$  h (days 219-224), and its stability is stronger than that of E1. The wave periods of E3 and E4 are about  $\sim 39$  h and  $\sim 24$  h, respectively. These

results reflect that E2, E3 and E4 wave periods are more stable compared to E1.

The spectra, spatial (vertical and latitudinal) structures of temperature, zonal and meridional wind, and diagnostic analysis of E1 are extracted from the two corresponding events (refer to Figure 4). Figures 4a, 4b show the least-squares fitting spectra for MERRA-2 temperature on days 187–196, 211–220 at  $\sim 48.2$  km and  $\sim 70$ – $80^\circ\text{S}$ , when and where the E1 maximizes. An eastward wavenumber -1 signal with the periods of  $\sim 106$  h and  $\sim 69$  h clearly dominates the whole spectrum. The temperature spatial structures corresponding to these E1 (i.e.,  $\sim 106$  h and  $\sim 69$  h) are displayed in Figures 4c, 4d. The temperature spatial structure of E1 exhibits obvious amplitude bimodal structure at  $\sim 70$ – $80^\circ\text{S}$  and  $\sim 50$  km, and  $\sim 70$ – $80^\circ\text{S}$  and  $\sim 60$  km, with the maximum at  $\sim 70$ – $80^\circ\text{S}$  and  $\sim 50$  km. The strongest temperature amplitude of E1 occurs at  $\sim 50$  km and  $\sim 70$ – $80^\circ\text{S}$  with an amplitude of  $\sim 10$  K on the days 211–220, and the other peak is  $\sim 9$  K ( $\sim 70$ – $80^\circ\text{S}$  and  $\sim 60$  km). The temperature amplitude of  $\sim 9$  K occurs at  $\sim 50$  km and  $\sim 70$ – $80^\circ\text{S}$  during days 187–196, and the rest is  $\sim 7$  K ( $\sim 70$ – $80^\circ\text{S}$  and  $\sim 60$  km). The corresponding spatial structures of zonal wind and meridional wind of these E1 are shown in Figures 4e, 4f, 4g and 4h. The maximum zonal wind amplitude of E1 occurs at  $\sim 60$ – $70^\circ\text{S}$  and  $\sim 60$  km with an amplitude of  $\sim 14$  m/s on days 187–196, and  $\sim 20$  m/s at  $\sim 50$ – $60^\circ\text{S}$  and  $\sim 60$  km on days 211–220. The amplitude of E1 meridional

wind hits  $\sim 10$  m/s at  $\sim 70$ - $80^\circ$ S and  $\sim 55$  km (days 187-196) and  $\sim 17$  m/s at  $\sim 70$ - $80^\circ$ S and  $\sim 60$  km (days 211-220), respectively.

Figures 4i, 4j show the diagnostic analysis results for the E1 events during days 187–196 and 211–220, respectively. Apparently, the EP flux vectors is more favorable to propagate in the SH winter and is dramatically amplified by the mean flow instabilities and appropriate background winds at polar region and between  $\sim 40$  km and  $\sim 80$  km, with EP flux propagating into the upper atmosphere (Figure 4i). Meanwhile, there is an EP flux at the mid-latitudes and  $\sim 60$ - $80$  km, which propagates into the lower atmosphere. The wave-mean flow interaction near its critical layer (106 h) of the green curve amplifies E1, and the positive refractive index region surrounded by the yellow curve also enhances E1 propagation. In addition, the strong instability and weak background wind at  $\sim 70$ - $80^\circ$ S and  $\sim 40$ - $60$  km provide sufficient energy for the upward EP flux to propagate and amplify. Nevertheless, the downward propagating EP flux is amplified by weak instability and strong background wind at  $\sim 50$ - $60^\circ$ S and  $\sim 60$ - $70$  km. Besides, both upward and downward EP fluxes eventually propagate toward the equator at  $\sim 50$  km. Figure 4j shows that EP flux on days 211-220 propagates downward and amplifies after the interaction of the critical layer ( $\sim 69$  h). The positive refractive index region, strong instability and weak background wind at  $\sim 50$ - $60^\circ$ S and  $\sim 60$ - $70$  km provide sufficient energy for E1 amplification and

propagation, and ultimately point towards the equator at  $\sim 50$  km. The results show that the weak background wind and strong instability in the polar region can promote the upward propagation and amplification of EP flux. Meanwhile, the appropriate background wind and instability in the mid-latitudes are also conducive to the downward propagation and amplification of EP flux. In other words, instability and appropriate background wind dominate the propagation and amplification of E1.

For E2, the spectra are observed at  $\sim 48.2$  km and  $\sim 60-70^\circ\text{S}$  on days 173-178 and 219-224 when the eastward wavenumber -2 becomes the primary wave mode with the wave periods  $\sim 38$  h and  $\sim 39$  h, respectively (as shown in Figures 5a, 5b). The temperature spatial structures corresponding to these E2 ( $\sim 38$  h and  $\sim 39$  h) are presented in Figures 5c, 5d. The temperature spatial structure of E2 shows an obvious amplitude bimodal structure at  $\sim 60-70^\circ\text{S}$  and  $\sim 50$  km, and  $\sim 60-70^\circ\text{S}$  and  $\sim 60$  km, with the maximum at  $\sim 60-70^\circ\text{S}$  and  $\sim 50$  km. The maximum temperature amplitude of E1 occurs at  $\sim 50$  km and  $\sim 60-70^\circ\text{S}$  with an amplitude of  $\sim 7.5$  K on the days 173-178, and the other peak is  $\sim 6$  K ( $\sim 70^\circ\text{S}$  and  $\sim 60$  km). The temperature amplitude of  $\sim 10$  K happens at  $\sim 50$  km and  $\sim 60-70^\circ\text{S}$  during days 219-224, and the rest is  $\sim 6$  K ( $\sim 70^\circ\text{S}$  and  $\sim 60$  km). The corresponding spatial structures of zonal wind and meridional wind of these E2 are illustrated in Figures 5e, 5f, 5g and 5h. The zonal wind spatial structure of E2 shows an obvious amplitude bimodal structure at

~50-60°S and ~60 km, and ~70-80°S and ~60 km, with the maximum at ~50-60°S and ~60 km. The maximum zonal wind amplitude of E2 appear at ~50-60°S and ~60 km with an amplitude of ~10 m/s on days 173-178, and the other peak is ~9 m/s (~70-80°S and ~60 km). The zonal wind amplitude of ~20 m/s occurs at ~50-60°S and ~60 km on days 219-224, and the rest is ~15 m/s (~70-80°S and ~60 km). The amplitude of E2 meridional wind reaches ~13 m/s at ~70-80°S and ~60 km (days 173-178) and ~27 m/s at ~70-80°S and ~60 km (days 219-224), respectively.

Figures 5i and 5j illustrate the diagnostic analysis during days 173-178 and 219-224 for E2, respectively. Obviously, E2 is more likely to propagate in the SH winter and is dramatically amplified by the mean flow instabilities at the middle-high latitudes between ~40 km and ~80 km. With EP flux propagating into the lower atmosphere, it eventually propagates toward the equator at ~50 km. Besides, E2 is amplified and propagated by the wave-mean flow interactions near its critical layer (~38 h) of the green curve, and the promoting effect of the positive refractive index region surrounded by the yellow curve. Meanwhile, the weak instability and strong background wind at ~50-60°S and ~50-70 km provide the energy for the propagation and amplification of EP flux into the lower atmosphere during days 173-178 (Figure 5i). According to the diagnostic analysis of days 219-224, E2 obtains sufficient energy from strong instability and strong background wind at ~50-60°S and ~60-70

km. It is amplified and propagated into the lower atmosphere through the critical layer and positive refractive index action (as shown in Figure 5j). The results show that the background wind at  $\sim 50\text{-}60^\circ\text{S}$  and  $\sim 50\text{-}70$  km is weaker on days 173-178 than on days 219-224; and the instability at  $\sim 50\text{-}60^\circ\text{S}$  and  $\sim 60\text{-}70$  km is stronger on days 219-224 than on days 173-178. Our results show that E2 has absorbed sufficient energy to be amplified under the background conditions during days 219-224 (Figures 5a, 5b).

Figures 6a and 6b show the observed spectra of E3 at  $\sim 48.2$  km and  $\sim 60\text{-}70^\circ\text{S}$  on days 151-154 and 201-204, and the wave periods of locked wavenumber -3 are  $\sim 29$  h and  $\sim 29$  h, respectively. The corresponding temperature spatial structures of these E3 (i.e.,  $\sim 29$  h and  $\sim 29$  h) are displayed in Figures 6c, 6d. The temperature spatial structure of E3 shows an obvious amplitude bimodal structure at  $\sim 60\text{-}70^\circ\text{S}$  and  $\sim 50$  km, and  $\sim 60\text{-}70^\circ\text{S}$  and  $\sim 60$  km, with the maximum at  $\sim 60\text{-}70^\circ\text{S}$  and  $\sim 50$  km. Besides, E3 also has a weak peak at  $\sim 60\text{-}70^\circ\text{S}$  and  $\sim 70$  km. The strongest temperature amplitude of E3 occurs at  $\sim 50$  km and  $\sim 60\text{-}70^\circ\text{S}$  with an amplitude of  $\sim 6\text{K}$  on the days 151-154, and the other peak is  $\sim 5$  K ( $\sim 60\text{-}70^\circ\text{S}$  and  $\sim 60$  km). The temperature amplitude of  $\sim 5$  K happens at  $\sim 50$  km ( $\sim 60$  km) and  $\sim 60\text{-}70^\circ\text{S}$  during days 201-204. The corresponding spatial structures of zonal wind and meridional wind of these E3 are shown in Figures 6e, 6f, 6g and 6h. The zonal wind spatial structure of E3

shows an obvious amplitude bimodal structure at  $\sim 70\text{-}80^\circ\text{S}$  and  $\sim 60$  km, and  $\sim 50\text{-}60^\circ\text{S}$  and  $\sim 60$  km. The zonal wind amplitudes of E3 occur at  $\sim 70\text{-}80^\circ\text{S}$  and  $\sim 60$  km ( $\sim 50\text{-}60^\circ\text{S}$  and  $\sim 60$  km) with an amplitude of  $\sim 9$  m/s on days 151-154, and  $\sim 9$  m/s at  $\sim 70\text{-}80^\circ\text{S}$  and  $\sim 60$  km ( $\sim 50\text{-}60^\circ\text{S}$  and  $\sim 60$  km) on days 201-204. The amplitude of E3 meridional wind hits  $\sim 13$  m/s at  $\sim 60\text{-}70^\circ\text{S}$  and  $\sim 55$  km (days 151-154) and  $\sim 16$  m/s at  $\sim 60\text{-}70^\circ\text{S}$  and  $\sim 55$  km (days 201-204), respectively.

EP flux of E3 is similar to that of E2. The instability and appropriate background wind at the mid-high latitudes between  $\sim 50$  km and  $\sim 70$  km dramatically amplify the propagation of E3, which is enhanced by the interaction near the critical layer ( $\sim 29$  h) and the positive refractive index region (Figures 6i and 6j). Notably, the strong instability and weak background wind at  $\sim 50\text{-}60^\circ\text{S}$  and  $\sim 60\text{-}70$  km on days 151-154 provide sufficient energy for the propagation and amplification of EP flux into the lower atmosphere, and ultimately point toward the equator at  $50$  km. During days 201-204, the EP flux propagates into the lower atmosphere and gets amplified by interaction at the critical layer ( $\sim 29$  h). Besides, weak instability and weak background wind at  $\sim 50\text{-}60^\circ\text{S}$  and  $\sim 60\text{-}70$  km provide the energy to amplify the E3 propagation. Figures 6c 6d indicate that the stronger the instability at  $\sim 50\text{-}60^\circ\text{S}$  and  $\sim 60\text{-}70$  km, the stronger the temperature amplitude of E3. We believe that the background wind and instability at  $\sim 50\text{-}60^\circ\text{S}$  and  $\sim 60\text{-}70$  km are the main reasons for the

propagation and amplification of EP flux into the lower atmosphere.

For E4, the spectra appear at  $\sim 48.2$  km and  $\sim 50$ - $60^\circ\text{S}$  on days 127-130 and 213-216 when the eastward wavenumber -4 signal with the wave period of  $\sim 25$  h and  $\sim 21$  h (see Figures 7a, 7b). The corresponding temperature spatial structures of these E4 (i.e.,  $\sim 25$  h and  $\sim 21$  h) are shown in Figures 7c, 7d. The temperature spatial structure of E4 shows an obvious amplitude bimodal structure at  $\sim 50$ - $60^\circ\text{S}$  and  $\sim 50$  km, and  $\sim 50$ - $60^\circ\text{S}$  and  $\sim 60$  km, with the maximum at  $\sim 50$ - $60^\circ\text{S}$  and  $\sim 50$  km. The maximum temperature amplitude of E4 occurs at  $\sim 50$  km and  $\sim 50$ - $60^\circ\text{S}$  with an amplitude of  $\sim 4$  K on the days 127-130, and the other peak is  $\sim 3$  K ( $\sim 60$ - $70^\circ\text{S}$  and  $\sim 60$  km). The temperature amplitude of  $\sim 3$  K occurs at  $\sim 50$  km ( $\sim 60$  km) during days 213-216. The corresponding spatial structures of zonal wind and meridional wind of these E4 are presented in Figures 7e, 7f, 7g and 7h. The zonal wind spatial structure of E4 shows an obvious amplitude bimodal structure at  $\sim 50$ - $60^\circ\text{S}$  and  $\sim 55$  km, and  $\sim 60$ - $70^\circ\text{S}$  and  $\sim 55$  km, with the maximum at  $\sim 50$ - $60^\circ\text{S}$  and  $\sim 55$  km. The maximum zonal wind amplitude of E4 happens at  $\sim 50$ - $60^\circ\text{S}$  and  $\sim 55$  km with an amplitude of  $\sim 9$  m/s on days 127-130, and the other peak is  $\sim 5$  K ( $\sim 60$ - $70^\circ\text{S}$  and  $\sim 55$  km). The zonal wind amplitude of  $\sim 5$  m/s occurs at  $\sim 50$ - $60^\circ\text{S}$  ( $\sim 60$ - $70^\circ\text{S}$ ) and  $\sim 55$  km on days 213-216. The amplitude of E4 meridional wind reaches  $\sim 8$  m/s at  $\sim 60$ - $70^\circ\text{S}$  and  $\sim 55$  km (days 127-130) and  $\sim 10$  m/s at  $\sim 60$ - $70^\circ\text{S}$  and  $\sim 55$  km (days 213-216), respectively.



Diagnostic analysis for E4 on days 127-130 and 213-216 are shown in Figures 7i and 7j, respectively. The results demonstrate that E4 is dramatically amplified by the mean flow instabilities at the middle-high latitudes between  $\sim 50$  km and  $\sim 70$  km. With EP flux propagating into the lower atmosphere, it finally propagates toward the equator at  $\sim 50$  km. E4 is amplified and propagated by the wave-mean flow interaction near the critical layer ( $\sim 25$  h,  $\sim 21$  h), and the positive refractive index region generates the promoting effect. The strong instability and weak background wind at  $\sim 50$ - $60^\circ$ S and  $\sim 60$ - $70$  km provide sufficient energy for the propagation and amplification of EP flux into the lower atmosphere during days 127-130. Besides, E4 obtains energy from weak instability and weak background wind at  $\sim 50$ - $60^\circ$ S and  $\sim 60$ - $70$  km on days 213-216, and it is amplified and propagated into the lower atmosphere. The background wind at  $\sim 50$ - $60^\circ$ S and  $\sim 60$ - $70$  km on days 127-130 is similar to on days 213-216, and the instability at  $\sim 50$ - $60^\circ$ S and  $\sim 60$ - $70$  km is stronger on days 127-30 than on days 213-216. According to Figures 7a and 7b, E4 absorbs sufficient energy to be amplified under the background conditions on days 127-130, and the temperature amplitude on 127-130 days is stronger.

### **3.2 In the Northern Hemisphere**

Figure 8 shows that the observed maximum temperature amplitude appears at  $\sim 59.2$  km and  $\sim 70$ - $80^\circ$ N for E1, and the E2 and E3 peaks at

~59.2 km and ~60-70°N. The maximum perturbation of E1 occurs on days 41-50 (with an amplitude of ~8K), while the remaining fluctuations occur on days 25-34 and 339-348. Besides, the strongest E2 occurs on days 69-74 (with an amplitude of ~7 K), and the rest are distributed on days 25-30, 317-322 and 341-346. By contrast, the E3 maximizes on days 35-38 (with an amplitude of ~3K), and also shows a peak on days 53-56. Based on the study of the wave period in the SH for eastward wave, the periodic variabilities of E1, E2 and E3 in the NH are also examined. The wave period of E1 decreases from a maximum of ~118 h (days 25-34) to ~80 h (days 41-50), indicating the instability of the wave period of E1 in the NH. The E2 events occur on days 25-30, 69-74, 317-322 and 341-346, of which the corresponding wave periods are ~36, ~53, ~52 and ~48 h, which are stronger and more stable than E1. Besides, the wave period of E3 is relatively stable at ~29 h and ~27 h. Thus, E2 and E3 wave periods are more stable than E1.

The spectra, spatial (vertical and latitudinal) structures of temperature, zonal and meridional wind, and diagnostic analysis of E1 are extracted from the corresponding representative events (as shown in Figure 9). Figures 9a and 9b show the observed spectra of E1 at ~59.2 km and ~70-80°N on days 25-34 and 41-50, and the wave periods of locked wavenumber -1 are ~118 h and ~80 h, respectively. The corresponding temperature spatial structures of these E1 (~118 h and ~80 h) are shown

in Figures 9c, 9d. The temperature spatial structure of E1 shows an  
 obvious amplitude bimodal structure during days 25-34 at  $\sim 60-70^\circ\text{N}$  and  
 $\sim 60$  km, and  $\sim 40-50^\circ\text{N}$  and  $\sim 70$  km, with the maximum at  $\sim 60-70^\circ\text{N}$  and  
 $\sim 60$  km. On top of that, E1 also has bimodal structure on days 41-50 at  
 $\sim 60-70^\circ\text{N}$  and  $\sim 60$  km, and  $\sim 60-70^\circ\text{N}$  and  $\sim 70$  km. The strongest  
 temperature amplitude of E1 occurs at  $\sim 60-70^\circ\text{N}$  and  $\sim 60$  km with an  
 amplitude of  $\sim 7$  K on the days 25-34, and the other peak is  $\sim 4$  K  
 ( $\sim 40-50^\circ\text{N}$  and  $\sim 70$  km). The temperature amplitude of  $\sim 10$  K occurs at  
 $\sim 60$  km and  $\sim 60-70^\circ\text{N}$  during days 41-50, and the rest is  $\sim 8$  K ( $\sim 60-70^\circ\text{N}$   
 and  $\sim 70$  km). The corresponding spatial structures of zonal wind and  
 meridional wind of these E1 are illustrated in Figures 9e, 9f, 9g and 9h.  
 The zonal wind spatial structure of E1 presents an obvious amplitude  
 bimodal structure at  $\sim 70-80^\circ\text{N}$  and  $\sim 70$  km, and  $\sim 50-60^\circ\text{N}$  and  $\sim 70$  km.  
 The zonal wind amplitude of  $\sim 13$  m/s occurs at  $\sim 70-80^\circ\text{N}$  and  $\sim 70$  km on  
 days 25-34, and the rest is  $\sim 10$  m/s ( $\sim 50-60^\circ\text{N}$  and  $\sim 70$  km). In addition,  
 there is a weak peak of 9K during days 25-34 ( $\sim 30-40^\circ\text{N}$  and  $\sim 70$  km).  
 The maximum zonal wind amplitude of E1 occurs at  $\sim 70-80^\circ\text{N}$  and  $\sim 70$   
 km with an amplitude of  $\sim 19$  m/s on days 41-50, and the other peak is  
 $\sim 13$  K ( $\sim 50-60^\circ\text{N}$  and  $\sim 70$  km). The amplitude of E1 meridional wind hits  
 $\sim 14$  m/s at  $\sim 70-80^\circ\text{N}$  and  $\sim 70$  km (days 25-34) and  $\sim 22$  m/s at  $\sim 70-80^\circ\text{N}$   
 and  $\sim 70$  km (days 41-50), respectively.

The diagnostic analysis results for E1 (in Figures 9i and 9j) suggest

the dramatic amplification of E1 by the mean flow instabilities at the middle-high latitudes between  $\sim 50$  km and  $\sim 70$  km. With the propagation of EP flux into the polar lower atmosphere, it eventually propagates toward the equator at  $\sim 50$  km. The wave-mean flow interaction near the critical layers ( $\sim 118$  h,  $\sim 80$  h) amplifies and propagates E1, and the promoting effect of the positive refractive index region amplifies E1. Furthermore, the weak instability and strong background wind at  $\sim 40$ - $50^\circ$ N and  $\sim 60$ - $70$  km generate the energy for the propagation and amplification of EP flux into the polar lower atmosphere during days 25-34. The E1 obtains sufficient energy from weak instability and suitable background wind on days 41-50 at  $\sim 40$ - $50^\circ$ N and  $\sim 60$ - $70$  km, and is amplified and propagated into the polar lower atmosphere through the critical layer and positive refractive index action. The background wind at  $\sim 40$ - $50^\circ$ N and  $\sim 60$ - $70$  km is stronger on days 25-34 than on days 41-50, but their instability is similar, indicating that stronger background winds might weaken E1 propagation and amplification at the mid-northern latitudes. Our results show that E1 absorbs adequate energy to be amplified under the background conditions during days 41-50, reflecting stronger temperature amplitude (see Figures 9a, 9b).

For E2, the spectra are at  $\sim 59.2$  km and  $\sim 60$ - $70^\circ$ N on days 25-30 and 69-74 when the eastward wavenumber -2 signal with the period  $\sim 36$  h and  $\sim 53$  h (as shown in Figures 10a, 10b). The corresponding temperature

spatial structures of these E2 (i.e.,  $\sim 36$  h and  $\sim 53$  h) are presented in  
 Figures 10c, 10d. The temperature spatial structure of E2 demonstrates an  
 obvious amplitude bimodal structure at  $\sim 60$ - $70^\circ\text{N}$  and  $\sim 60$  km, and  
 $\sim 60$ - $70^\circ\text{N}$  and  $\sim 70$  km, with the maximum at  $\sim 60$ - $70^\circ\text{N}$  and  $\sim 60$  km. The  
 maximum temperature amplitude of E2 occurs at  $\sim 60$ - $70^\circ\text{N}$  and  $\sim 60$  km  
 with an amplitude of  $\sim 5$  K on days 25-30, and the other peak is  $\sim 4$  K  
 ( $\sim 60$ - $70^\circ\text{N}$  and  $\sim 70$  km). The temperature amplitude of  $\sim 9$  K occurs on  
 days 69-74 at  $\sim 60^\circ\text{S}$  and  $\sim 60$  km, and the other peaks are  $\sim 7$  K ( $\sim 60$ - $70^\circ\text{N}$   
 and  $\sim 70$  km),  $\sim 5$  K ( $\sim 60$ - $70^\circ\text{N}$  and  $\sim 50$  km). The corresponding spatial  
 structures of zonal wind and meridional wind of these E2 are shown in  
 Figures 10e, 10f, 10g and 10h. The zonal wind spatial structure of E2  
 shows an obvious amplitude bimodal structure at  $\sim 60$ - $70^\circ\text{N}$  and  $\sim 60$  km,  
 and  $\sim 40$ - $50^\circ\text{N}$  and  $\sim 60$  km, with the maximum at  $\sim 40$ - $50^\circ\text{N}$  and  $\sim 60$  km.  
 The maximum zonal wind amplitude of E2 appears at  $\sim 60$ - $70^\circ\text{N}$  and  $\sim 60$   
 km ( $\sim 40$ - $50^\circ\text{N}$  and  $\sim 60$  km) with an amplitude of  $\sim 6$  m/s on days 25-30.  
 Zonal wind amplitude occurs at  $\sim 40$ - $50^\circ\text{N}$  and  $\sim 60$  km with an amplitude  
 of  $\sim 18$  m/s on days 41-50, and the other peak is  $\sim 16$  K ( $\sim 60$ - $70^\circ\text{N}$  and  
 $\sim 60$  km). The amplitude of E2 meridional wind reaches  $\sim 7$  m/s at  
 $\sim 60$ - $70^\circ\text{N}$  and  $\sim 70$  km (days 25-30) and  $\sim 18$  m/s at  $\sim 60$ - $70^\circ\text{N}$  and  $\sim 60$  km  
 (days 41-50), respectively.

The diagnostic analysis of E2 on days 25-30 and 69-74 are shown in  
 Figures 10i and 10j, respectively. Apparently, E2 is significantly

amplified by the mean flow instabilities at the middle-high latitudes between  $\sim 40$  km and  $\sim 70$  km, with EP flux propagating into the polar lower atmosphere, and EP flux eventually propagate toward the equator at  $\sim 50$  km. E2 is amplified and propagated by the wave-mean flow interaction near the critical layers ( $\sim 36$  h,  $\sim 53$  h), and the positive refractive index region provides the promoting effect. The weak instability and strong background wind at  $\sim 50$ - $60^\circ$ N and  $\sim 60$ - $70$  km provide the energy for the propagation and amplification of EP flux into the polar lower atmosphere during days 25-30. Moreover, E2 obtains sufficient energy from strong instability and suitable background wind at  $\sim 50$ - $60^\circ$ N and  $\sim 60$ - $70$  km on days 69-74, and it is amplified and propagated into the polar lower atmosphere. The background wind at  $\sim 50$ - $60^\circ$ N and  $\sim 60$ - $70$  km on days 127-130 is similar to on days 213-216, and the instability at  $\sim 50$ - $60^\circ$ S and  $\sim 60$ - $70$  km is stronger on days 127-30 than on days 213-216. Although the background wind at  $\sim 50$ - $60^\circ$ N and  $\sim 60$ - $70$  km is stronger on days 25-30 than on days 69-74, the instability at  $\sim 50$ - $60^\circ$ N and  $\sim 60$ - $70$  km is stronger on days 69-74 than on days 25-30. The temperature amplitude results indicate that E2 absorbs sufficient energy to be amplified under the background conditions on days 69-74, with stronger temperature amplitude on days 69-74 (Figures 10a, 10b).

Figures 11a and 11b show the observed spectra of E3 at  $\sim 59.2$ km and  $\sim 60$ - $70^\circ$ N on days 35-38 and 53-56, and the wave periods of locked

wavenumber -3 are  $\sim 29$  h and  $\sim 27$  h, respectively. The corresponding temperature spatial structures of these E3 (i.e.,  $\sim 29$  h and  $\sim 27$  h) are shown in Figures 11c, 11d. The temperature spatial structure of E3 shows an obvious amplitude bimodal structure at  $\sim 50$ - $60^\circ\text{N}$  and  $\sim 60$  km, and  $\sim 50$ - $60^\circ\text{N}$  and  $\sim 70$  km, with the maximum at  $\sim 50$ - $60^\circ\text{N}$  and  $\sim 60$  km. The strongest temperature amplitude of E3 occurs at  $\sim 60$  km and  $\sim 50$ - $60^\circ\text{N}$  with an amplitude of  $\sim 6$  K on the days 35-38, and the other peak is  $\sim 5$  K ( $\sim 50$ - $60^\circ\text{N}$  and  $\sim 70$  km). The temperature amplitude of  $\sim 4$  K occurs at  $\sim 60$  km ( $\sim 70$  km) and  $\sim 50$ - $60^\circ\text{N}$  during days 53-56. The corresponding spatial structures of zonal wind and meridional wind of these E3 are illustrated in Figures 6e, 6f, 6g and 6h. The zonal wind spatial structure of E3 shows an obvious amplitude bimodal structure at  $\sim 40$ - $50^\circ\text{N}$  and  $\sim 70$  km, and  $\sim 60$ - $70^\circ\text{N}$  and  $\sim 70$  km. The zonal wind amplitudes of E3 occur at  $\sim 40$ - $50^\circ\text{N}$  and  $\sim 70$  km with an amplitude of  $\sim 15$  m/s on days 35-38, and  $\sim 12$  m/s at  $\sim 60$ - $70^\circ\text{N}$  and  $\sim 70$  km. The maximum zonal wind amplitude of E3 appears at  $\sim 40$ - $50^\circ\text{N}$  and  $\sim 70$  km ( $\sim 60$ - $70^\circ\text{N}$  and  $\sim 70$  km) with an amplitude of  $\sim 7$  m/s ( $\sim 6$  m/s) on days 53-56. The amplitude of E3 meridional wind reaches  $\sim 22$  m/s at  $\sim 50$ - $60^\circ\text{N}$  and  $\sim 70$  km (days 35-38) and  $\sim 12$  m/s at  $\sim 60$ - $70^\circ\text{N}$  and  $\sim 70$  km (days 53-56), respectively.

Obviously, the instability and appropriate background wind at the mid-latitudes between  $\sim 50$  km and  $\sim 70$  km and the interaction near the critical layers ( $\sim 29$  h,  $\sim 27$  h) dramatically amplify the propagation of E3

(see Figures 11i and 11j). The background wind is similar on days 35-38 and 53-56, and the former is relatively unstable. This finding indicates that the E3 in propagation is more likely to gather sufficient energy to be amplified on days 35-38. The instability and appropriate background wind at the mid-high latitudes between  $\sim 50$  km and  $\sim 70$  km drastically amplify the propagation of E3, which is enhanced by the interaction near the critical layers ( $\sim 29$  h,  $\sim 27$  h) and the positive refractive index region (Figures 11i, 11j). In particular, the strong instability and weak background wind at  $\sim 50$ - $60^\circ$ N and  $\sim 60$ - $70$  km on days 35-38 generate sufficient energy for the propagation and amplification of EP flux into the lower atmosphere, and ultimately point toward the equator at 50 km. The EP flux propagates to the lower atmosphere during days 35-38, and it is amplified by interactions at the critical layer ( $\sim 29$  h). In addition, weak instability and weak background winds on days 53-56 at  $\sim 50$ - $60^\circ$ N and  $\sim 60$ - $70$  km provide the energy to amplify E3 propagation. Combine with Figures 11c, 11d, the stronger the instability at  $\sim 50$ - $60^\circ$ N and  $\sim 60$ - $70$  km, the stronger the temperature amplitude of E3. The results show that the instability on days 35-38 at  $\sim 50$ - $60^\circ$ N and  $\sim 60$ - $70$  km are the primary reasons for the propagation and amplification of EP flux into the lower atmosphere.

### **3.3 Comparison between SH and NH**

The observed latitude and maximum temperature amplitude for



eastward planetary waves (i.e., E1, E2, E3, E4) decrease and weaken with increasing zonal wavenumber in the SH, reaching  $\sim 70\text{--}80^\circ\text{S}$ ,  $\sim 60\text{--}70^\circ\text{S}$ ,  $\sim 60\text{--}70^\circ\text{S}$  and  $\sim 50\text{--}60^\circ\text{S}$ , and  $\sim 10\text{ K}$ ,  $\sim 9\text{ K}$ ,  $\sim 6\text{ K}$ , and  $\sim 3\text{ K}$ , respectively. In addition, the occurrence date gets earlier with increasing zonal wavenumber. The temperature spatial structure demonstrates a bimodal-peak structure ( $\sim 50$  and  $\sim 60\text{ km}$ ), mainly located at  $\sim 50\text{ km}$ . The maximum zonal wind amplitudes of E1 and E2, E3 and E4 are almost the same, which are  $\sim 20\text{ m/s}$  and  $\sim 10\text{ m/s}$ , respectively. The maximum meridional wind amplitudes of E1, E2, E3 and E4 are  $\sim 17\text{ m/s}$ ,  $\sim 27\text{ m/s}$ ,  $\sim 16\text{ m/s}$  and  $\sim 11\text{ m/s}$ , respectively. The wave period of E1 tends to get shorter from 5 to 3 days; while E2 and E3 are close to  $\sim 40\text{ h}$  and  $\sim 30\text{ h}$ ; and E4 remains at  $\sim 24\text{ h}$ . E1, E2, E3 and E4 are more favorable to propagation in the SH winter and are abruptly amplified by the mean flow instabilities at the middle latitudes between  $\sim 40\text{ km}$  and  $\sim 70\text{ km}$ . With the propagation of EP flux into the lower atmosphere, and it finally propagates toward the equator at  $\sim 50\text{ km}$ . In addition, the propagation of EP flux for E1 to the upper atmosphere might be influenced by the instability and background wind at the Antarctic  $\sim 50\text{ km}$ .

The observed latitudes of E1, E2 (E3) decrease with increasing wavenumber in the NH, which are  $\sim 70\text{--}80^\circ\text{N}$ ,  $\sim 60\text{--}70^\circ\text{N}$  and  $\sim 60\text{--}70^\circ\text{N}$ . With bimodal-peak structure located at  $\sim 70\text{ km}$ , the temperature spatial structures of E1, E2 and E3 reach  $\sim 10\text{ K}$ ,  $\sim 9\text{ K}$  and  $\sim 6\text{ K}$ , respectively.

The maximum zonal wind amplitude for E1, E2 and E3 occur at  $\sim 50\text{-}80^\circ\text{N}$  and  $\sim 70$  km, and their amplitude are almost equal to  $\sim 18$  m/s. The maximum meridional winds of E1, E2 and E3 occur at  $\sim 50\text{-}80^\circ\text{N}$  and  $\sim 70$  km with amplitudes of  $\sim 22$  m/s,  $\sim 18$  m/s and  $\sim 22$  m/s, respectively. The wave period of E1 tends to be shorter from 5-3 days; and E2 and E3 are close to  $\sim 48$  h and  $\sim 30$  h. In addition, E1, E2 and E3 are more favorable to propagation in the NH winter and are dramatically amplified by the mean flow instabilities at the middle latitudes between  $\sim 40$  km and  $\sim 70$  km, with the propagation of EP flux into the lower atmosphere and then toward the equator at  $\sim 50$  km.

#### **4 Summary and Conclusions**

Based on the MERRA-2 temperature and wind observations in 2019, we present for the first time an extensive study of the global variation for eastward planetary wave activity, including zonal wave numbers of -1 (E1), -2 (E2), -3 (E3), -4 (E4) in the stratosphere and mesosphere. The temperature and wind amplitude and wave periods of each event were determined using 2-D least-squares fitting. Our study covered the spatial and temporal patterns of the eastward planetary waves in both hemispheres with a comprehensive diagnostic analysis on their propagation and amplification. The key findings of this study are summarized below:

1. The latitudes for the maximum (temperature, zonal and meridional

wind) amplitudes of E1, E2, E3 and E4 decrease with increasing wavenumber in the SH and NH. The E1, E2, E3 and E4 events occur earlier with increasing zonal wavenumber in the SH. In addition, eastward wave modes exist during summer periods with westward background wind in both hemispheres.

2. The temperature spatial structures of E1, E2, E3 and E4 present a double-peak structure, which is located at ~50 km and ~60 km in SH, ~60 km and ~70 km in SH. Furthermore, the lower peak is usually larger than the higher one.

3. The maximum (temperature, zonal and meridional wind) amplitude of E1, E2 and E3 decline with rising zonal wavenumber in the SH and NH. The maximum temperature amplitude in the SH are slightly larger and lie lower than those in the NH. In addition, the meridional wind amplitude are slightly larger than the zonal wind in the SH and NH.

4. The wave period of the E1 mode ranges 3-5 days in both hemispheres, while the period of E2 mode is slightly longer in the NH (~48 h) than in the SH (~40 h). The periods of E3 in both SH and NH are ~30 h, while the period of E4 is ~24 h.

5. The eastward planetary wave is more favorable to propagate in the winter hemisphere and is drastically amplified by the mean flow instabilities and appropriate background winds at polar region and the middle latitudes between ~40 km and ~80 km. Furthermore, the

643 amplification of planetary waves through wave-mean flow interaction  
644 occurs easily close to its critical layer. In addition, the direction of EP flux  
645 ultimately points towards the equator.

646 6. The strong instability and appropriate background wind in the lower  
647 layer of the Antarctic region might generate adequate energy to promote  
648 the E1 propagation and amplification to the upper atmosphere.

649 Overall, this study demonstrated how the background zonal wind in  
650 the polar middle atmosphere affects the dynamics of eastward planetary  
651 waves in the polar middle atmosphere.

*Data availability.* MERRA-2 data are available at  
<http://disc.gsfc.nasa.gov>.

*Code availability.* Code is available at  
<http://hdl.pid21.cn/21.86116.7/04.99.01293>.

*Author contributions.* LT carried out the data processing and analysis and wrote the manuscript. SYG and XKD contributed to reviewing the article.

*Competing interests.* The authors declare that they have no conflict of interest.

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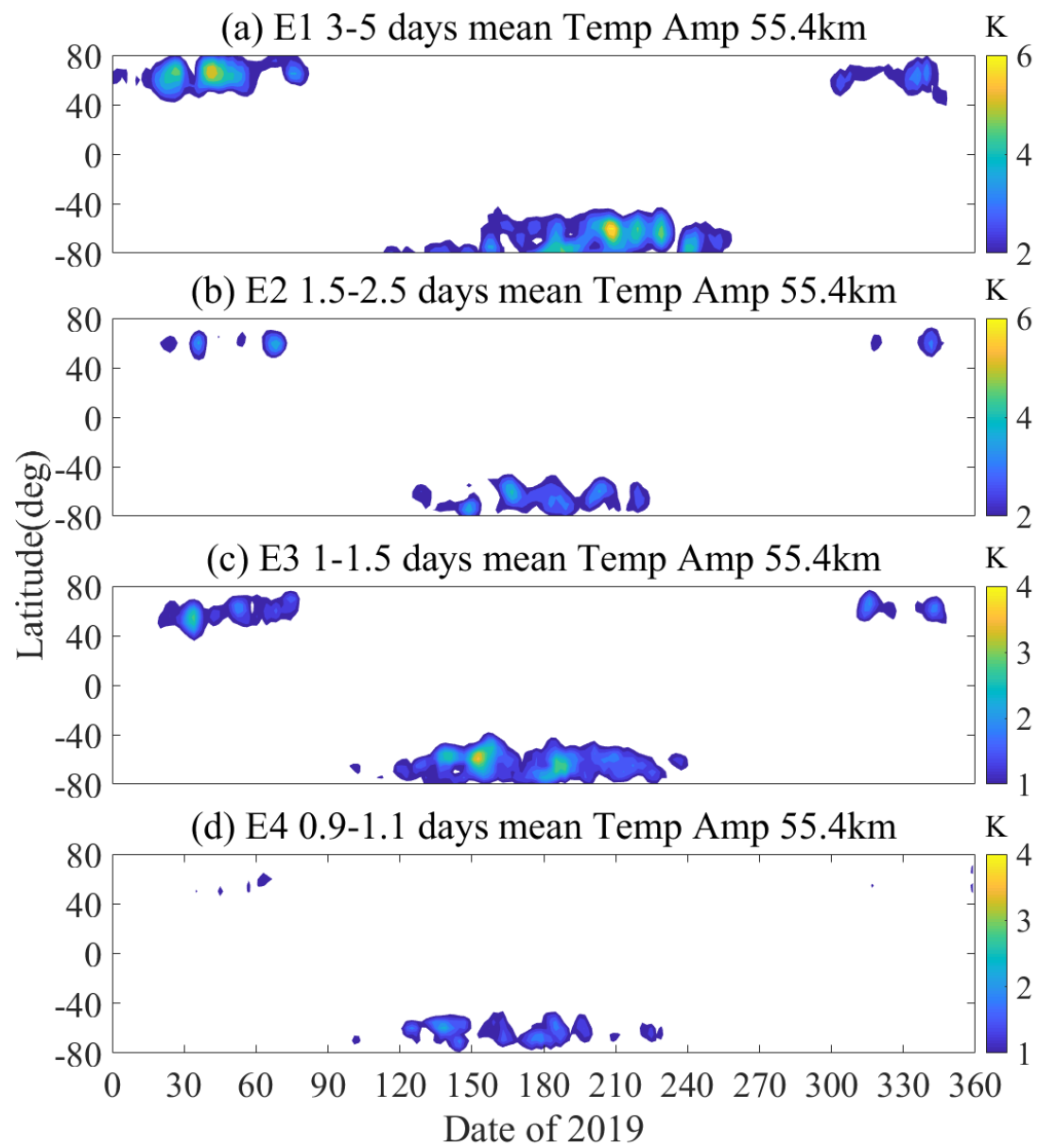
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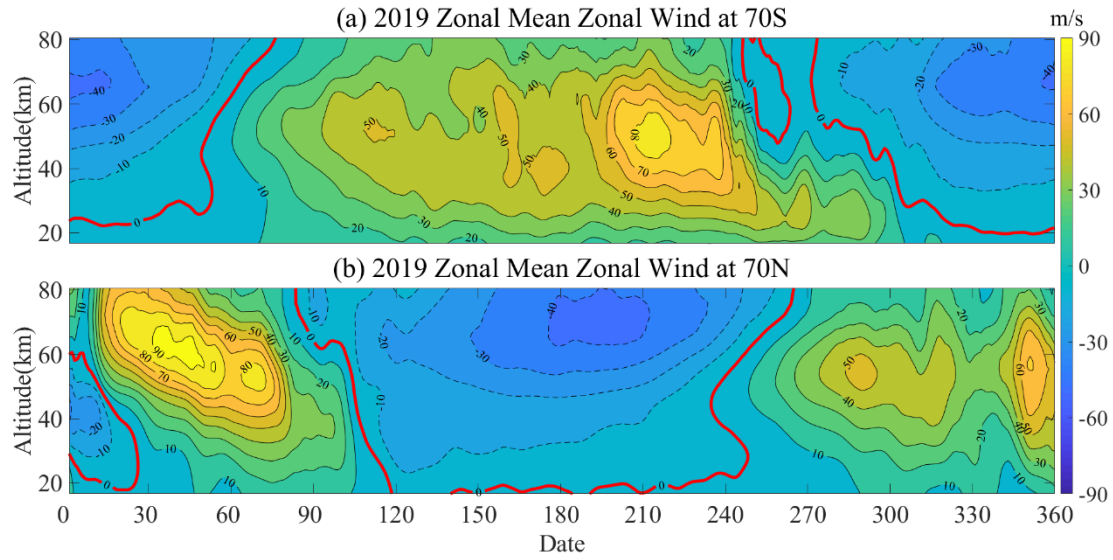
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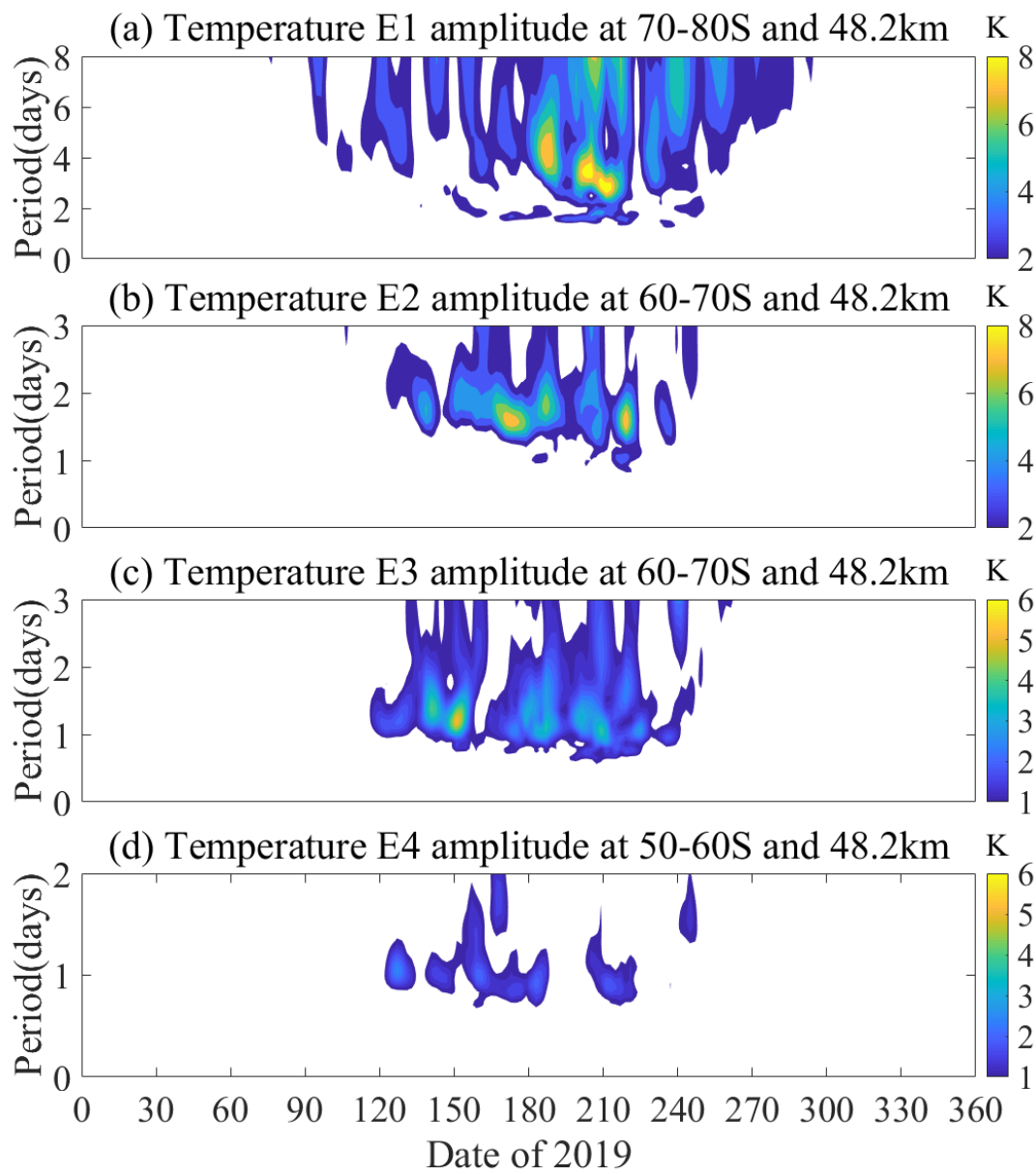


**Figure 1.** The global latitude-temporal variation structures of the (a) E1, (b) E2, (c) E3 and (d) E4 planetary waves during 2019. White areas represent small amplitude data (corresponds to the right color bar).

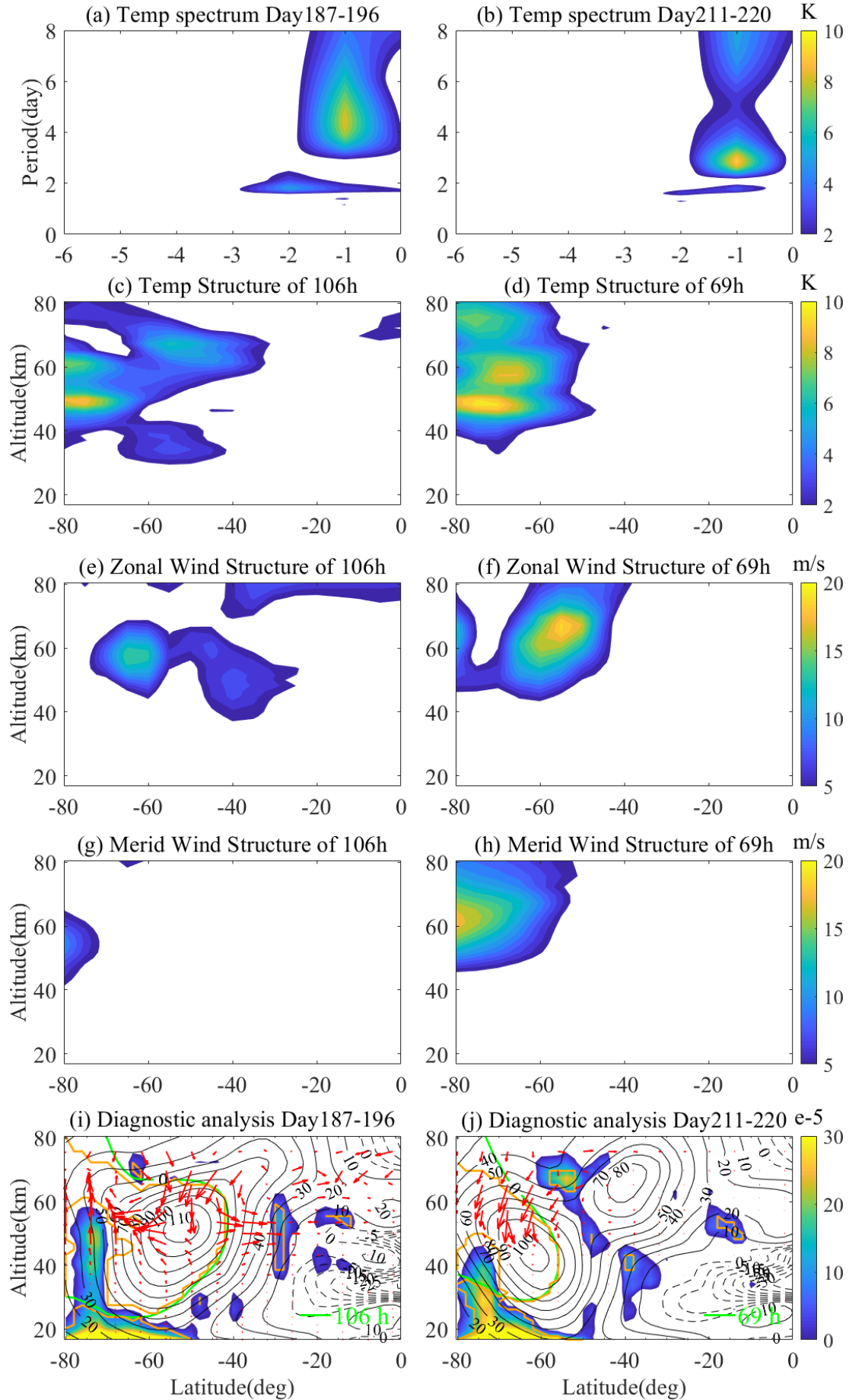


**Figure 2.** The zonal mean zonal wind variations of (a) 70°S and (b) 70°N during 2019.

The dotted line represents eastward wind, the solid line represents westward wind, and the red solid line is 0 m/s.

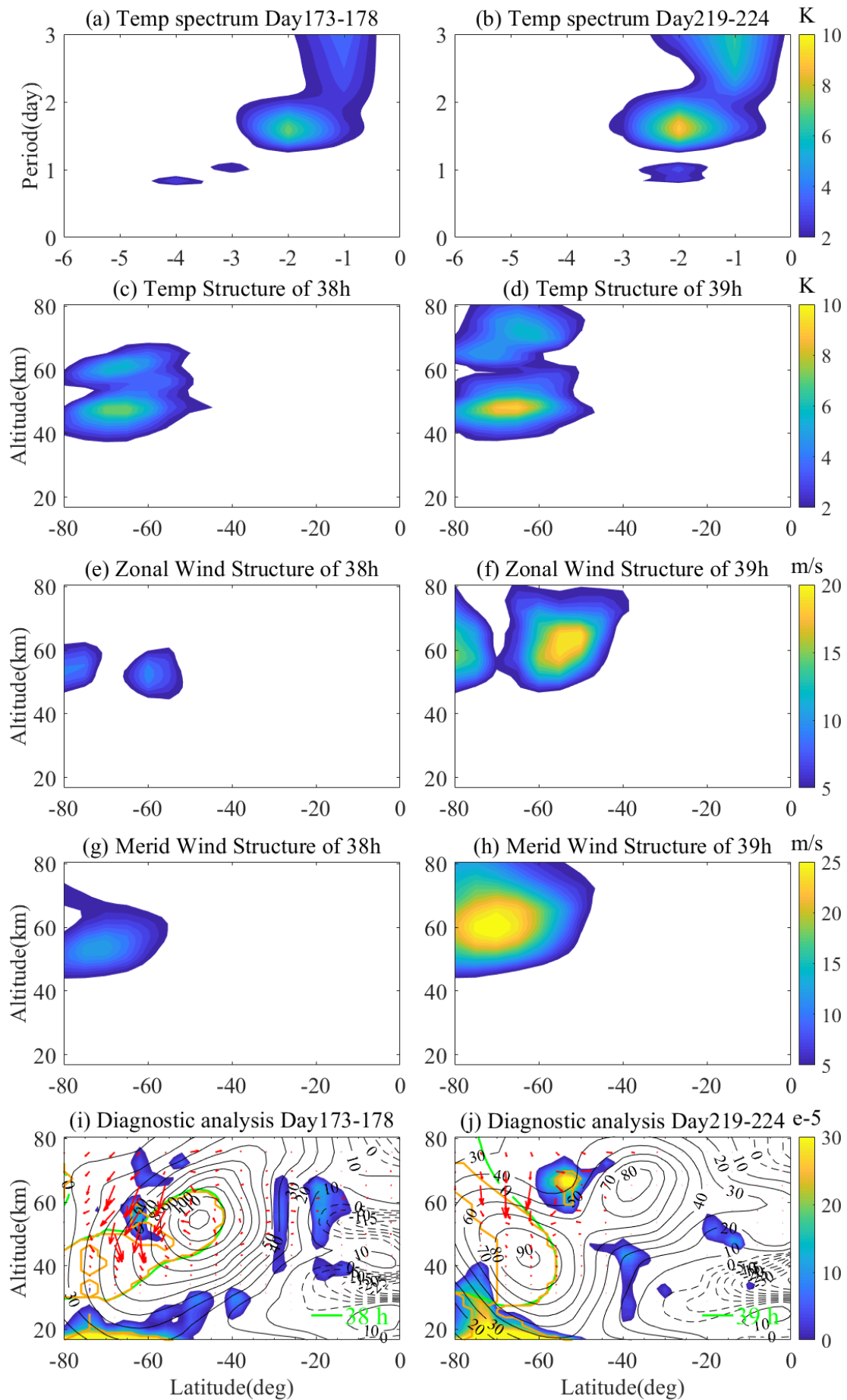


**Figure 3.** The temporal variations of (a) E1, (b) E2, (c) E3 and (d) E4 QTDWs during 2019 austral winter period.

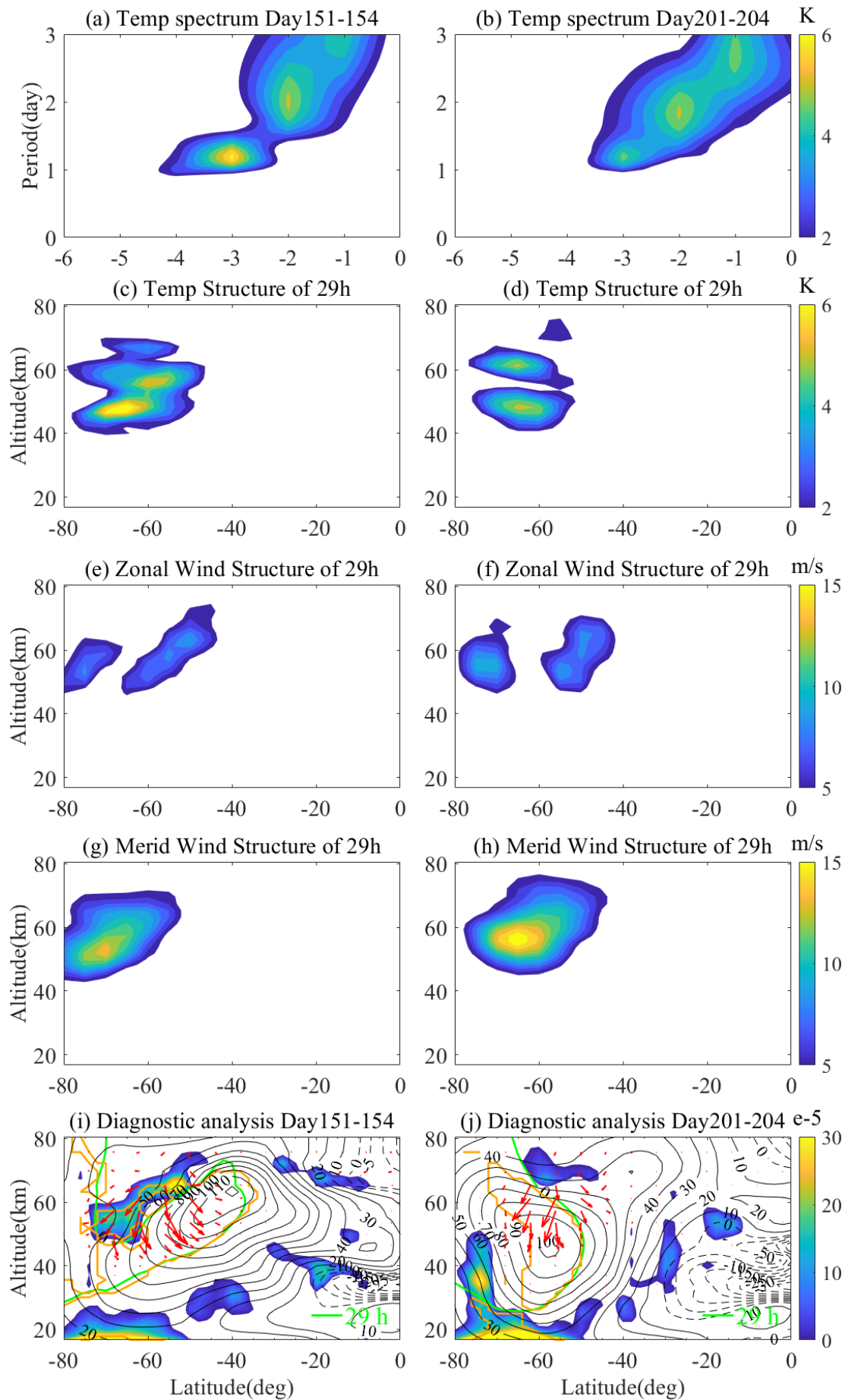




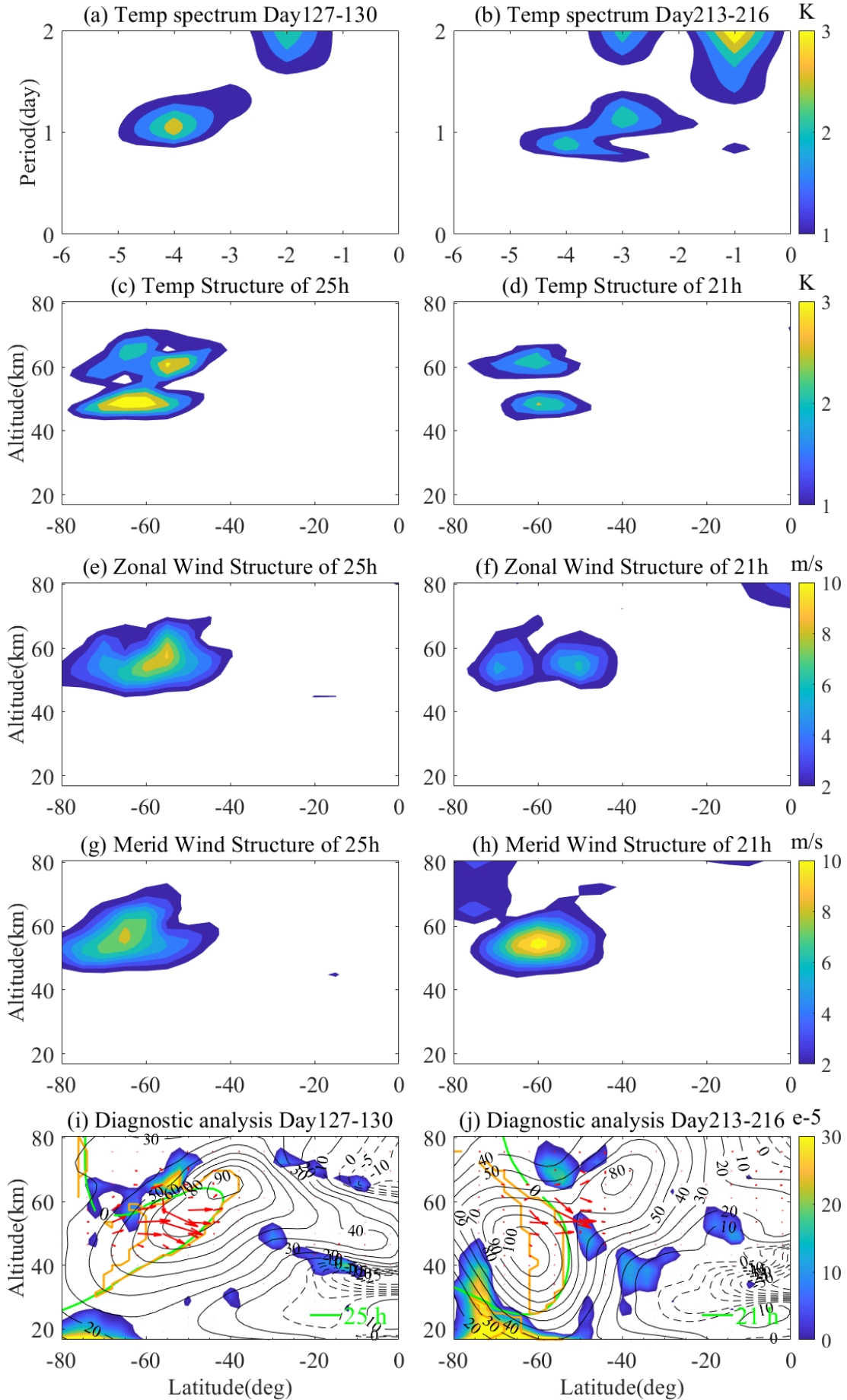
**Figure 4.** The (a, b) spectra, (c, d) temperature spatial structures, (e, f) zonal wind spatial structures, (g, h) meridional wind spatial structures, and (i, j) diagnostic analysis of the E1 typical events during 2019 austral winter period. The MERRA-2 temperature data observations at 48.2 km and 70-80°S during days 187–196 (Figure 4a), 211–220 (Figure 4d) are utilized, respectively. In the diagnostic analysis of E1 events, the blue shaded region is instability, the red arrow is EP flux, and the green line is critical layer. The green line represents critical layers of E1 with the natural period. Regions enclosed by orange solid lines are characterized by the positive refractive index for the E1.



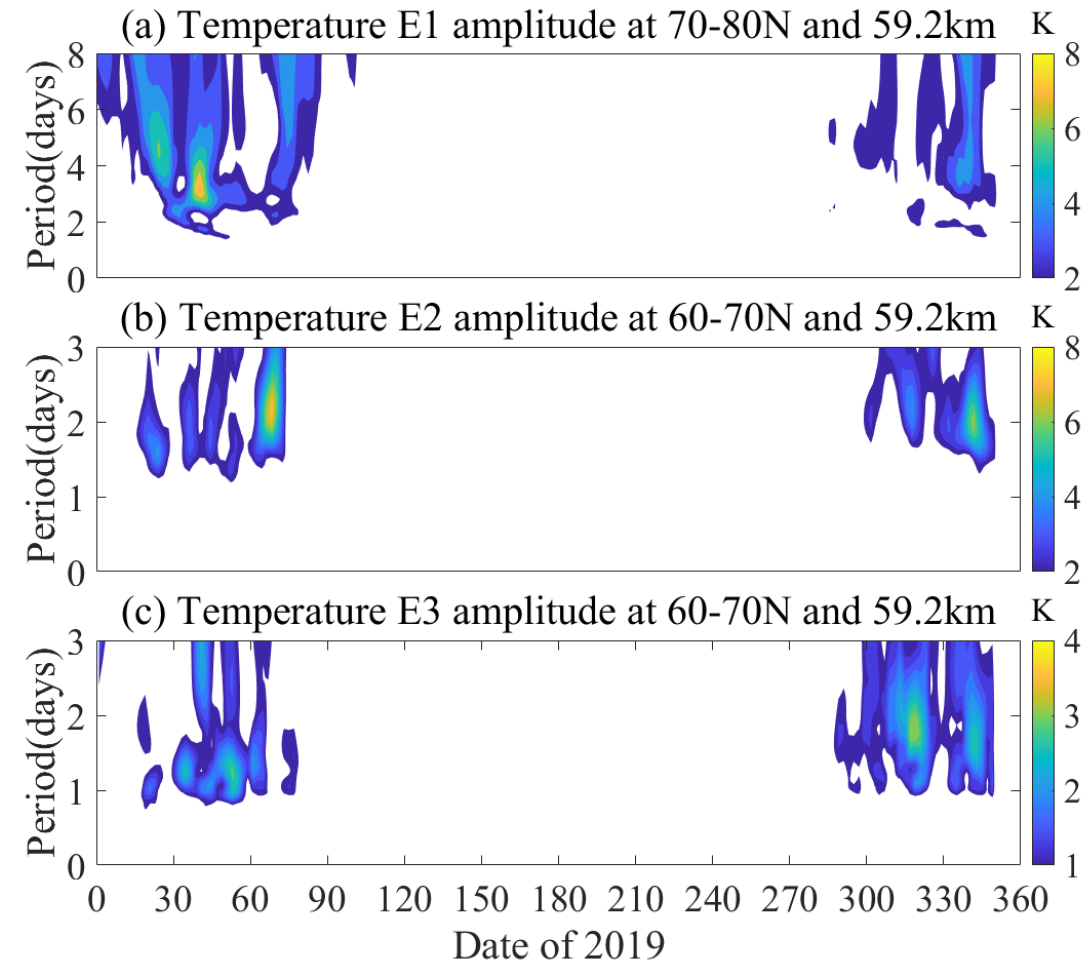
834 **Figure 5.** Same as Figure 4 but for E2 during the 2019 austral winter period.



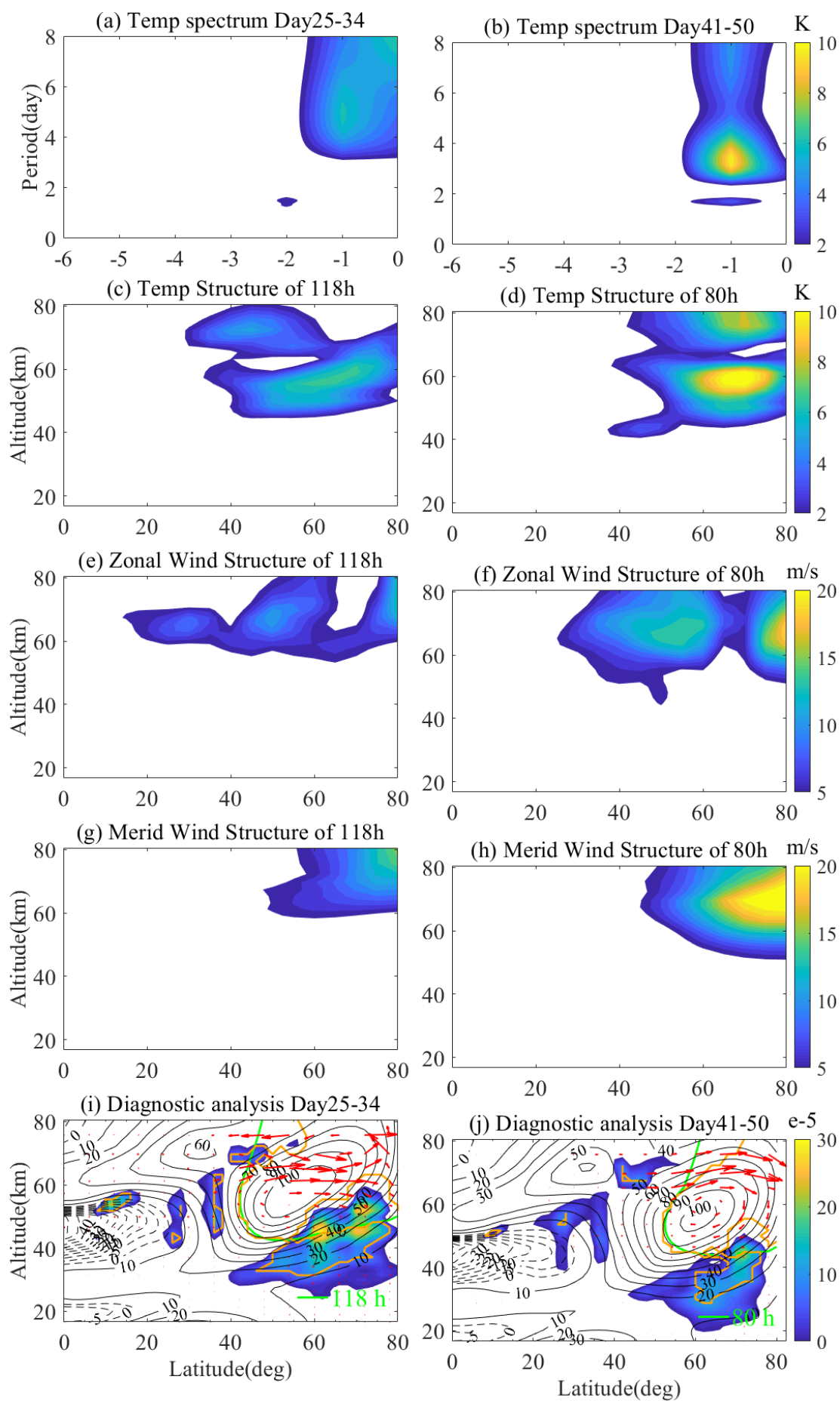
836 **Figure 6.** Same as Figure 4 but for E3 during the 2019 austral winter period.



**Figure 7.** Same as Figure 4 but for E4 during the 2019 austral winter period.

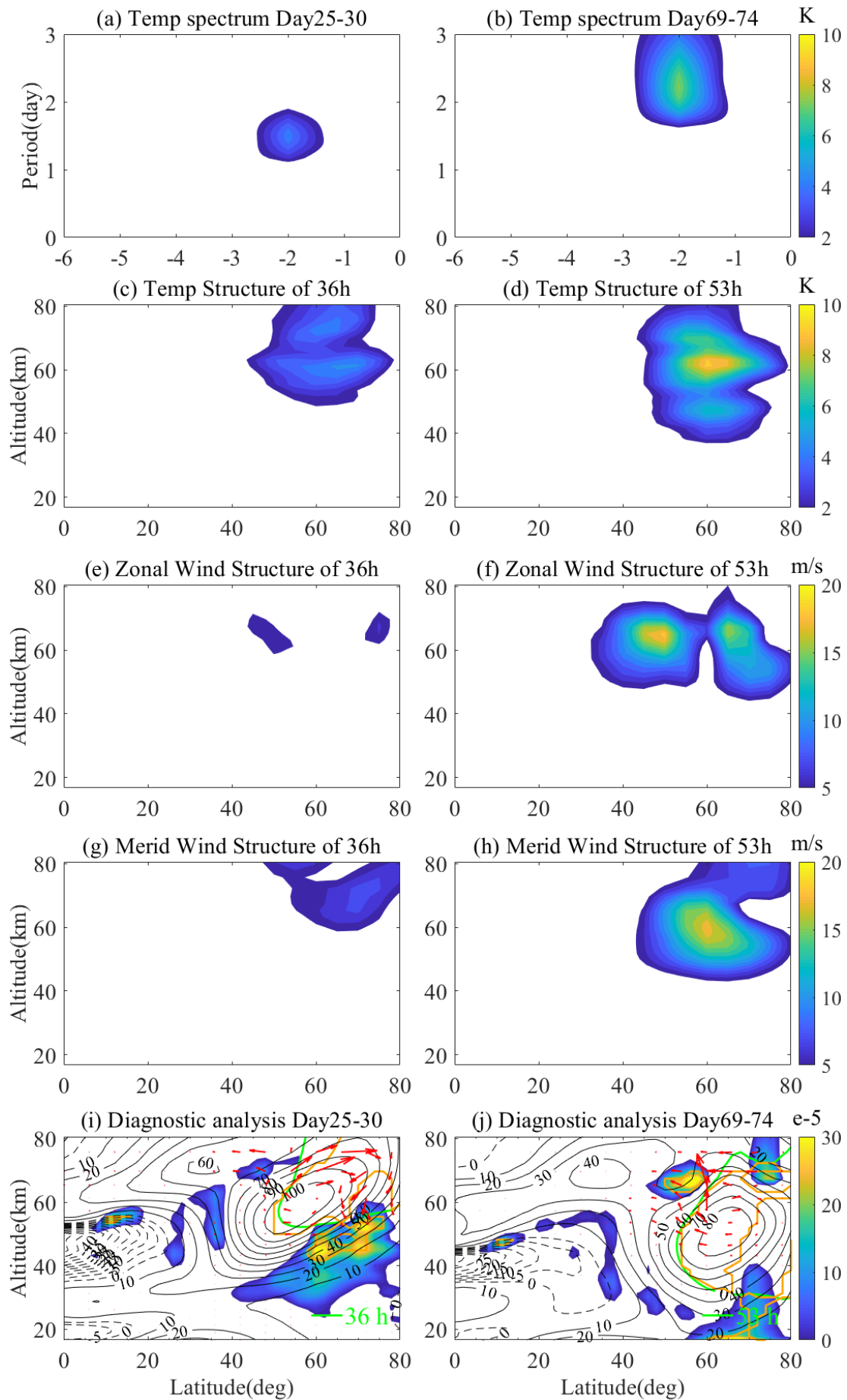


**Figure 8.** The temporal variations of (a) E1, (b) E2 and (c) E3 QTDWs during the 2019 boreal winter period.

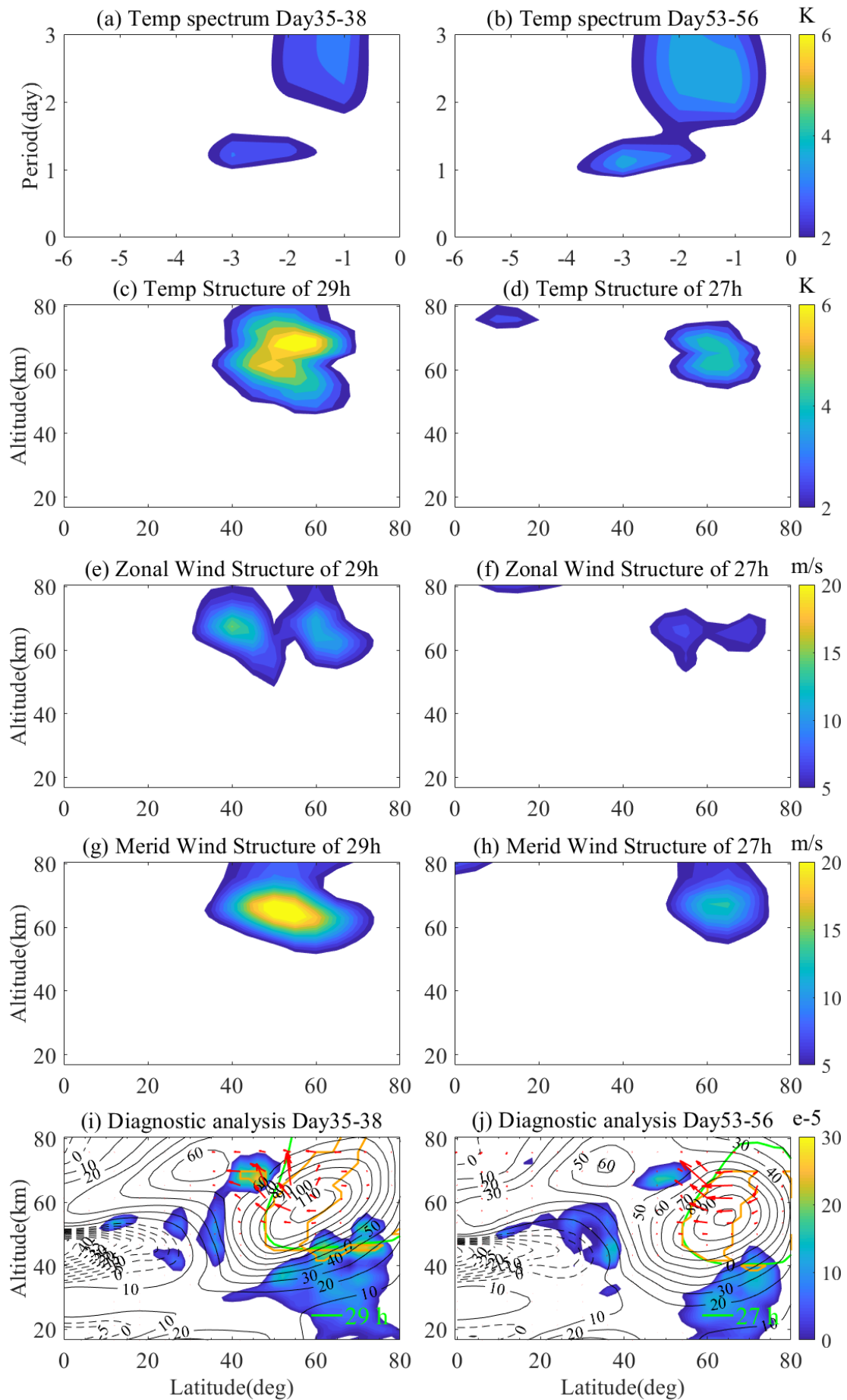




843 **Figure 9.** The (a, b) spectra, (c, d) temperature spatial structures, (e, f) zonal wind  
844 spatial structures, (g, h) meridional wind spatial structures and (i, j) diagnostic  
845 analysis of the E1 typical events during 2019 boreal winter period. The E1 events at  
846 48.2 km and 70-80°N were obtained from the MERRA-2 reanalysis.



848 **Figure 10.** Same as Figure 9 but for E2 during the 2019 boreal winter period.



850 **Figure 11.** Same as Figure 9 but for E3 during the 2019 boreal winter period.

851