



26 separates the attribution from the SCs and the background turbulence, the SCs
27 contribute no more than 10% to the vertical turbulent intensity, but their contributions
28 to the heat flux can be up to 80%. The lake patches produce consistent spatial
29 distributions of wind speed and turbulent stress over the lake–land boundary, and the
30 obvious change of turbulent momentum flux over the boundary of patches can not be
31 neglected. In the entrainment layer, the convective rolls still persist under stronger
32 geostrophic winds of 7–11 m s⁻¹. The increased downdrafts, which mainly occur over
33 the lake patches and carry more warm, dry air down from the free atmosphere. In
34 general, the SCs promote the growth of convective boundary layer, while the
35 background flows inhibit it. The background winds also weaken the patch-induced
36 turbulent intensity, heat flux, and convective intensity.

37 **Key word:** turbulence, heat flux, heterogeneously surface heating, background flows,
38 phase-averaged analysis

39 1. Introduction

40 Turbulence in the planetary boundary layer (PBL), which is derived from surface
41 heating and surface fraction, plays an important role in the exchange of heat,
42 momentum, moisture, and chemical constituents between the surface and free
43 atmosphere (Zhang et al., 2018). Previous studies on the turbulence and turbulent
44 exchange over homogeneous surfaces based on Monin-Obukhov similarity theory
45 were conducted before the 1990s (Sommeria and LeMone, 1978; Moeng, 1984).
46 Turbulence over heterogeneous surfaces was investigated through field campaigns
47 (Wang et al., 2016; Zhao et al., 2018) and numerical simulations (Shao et al., 2013;



48 [Liu et al., 2011](#)) in the past few decades, which has improved our understanding of the
49 transfer and spatial and temporal variability of the turbulence. Thermal surface
50 heterogeneity is a typical issue and leads to the formation of local/secondary
51 circulations. Sea and lake breezes are a well-known example of flows that are
52 generated by heterogeneous surface heating between the land and water ([Crosman and](#)
53 [Horel, 2012](#)). Observations have also revealed the imbalance in the surface energy
54 budget over heterogeneous surfaces ([Foken et al. 2010](#); [Xu et al., 2016](#)). The most
55 widely used eddy covariance (EC) system for a single site has been shown to
56 underestimate the turbulent flux due to the large-eddy transport or secondary
57 circulations not being captured ([Foken et al., 2010](#); [Xu et al., 2016](#)). The simulation
58 studies conducted by [Zhou et al. \(2018\)](#) and [Frederik and Matthias \(2018\)](#) showed that
59 the flux induced by mesoscale or secondary circulations is the main reason for the
60 energy imbalance. Moreover, the PBL parameterization schemes in the
61 weather/climate model over a heterogeneous surface have been continuously
62 improved until now ([Avisar and Pielke, 1989](#); [Shao et al., 2013](#)). Different surface
63 patterns such as mosaic ([Avisar and Schmidt, 1998](#)), chessboard ([Liu et al., 2011](#);
64 [Shen et al., 2016](#)), patchy-like ([Zhou et al., 2018](#)), and strip-like ([Li et al., 2011](#);
65 [Wang et al., 2011](#)) patterns have been utilized to simulate thermodynamic surface
66 heterogeneity. These studies confirmed that the secondary circulation induced by the
67 surface heterogeneity influences the PBL's properties and turbulent characteristics. In
68 addition, several studies have examined the effects of surface heterogeneity on
69 different levels of background winds ([Shen and Leclerc, 1995](#)) and the direction



70 relative to the orientation of the heterogeneity (Wang et al., 2011; Kang and
71 Lenschow, 2014). However, the issues related to the effects of the surface
72 heterogeneity in special areas still need to be explored.

73 The dynamic and thermodynamic influences of the Tibetan Plateau (TP) on the
74 regional and global weather and climate systems are closely related to its PBL, and
75 turbulence plays a significant role in the mass and energy exchange between the TP
76 and the atmosphere (Chen et al., 2013; Chen et al., 2016). Different landscapes make
77 up the heterogeneous land surface over the TP. As the Asian Water Tower, lakes are
78 widely and densely distributed over the TP, which affects the overlying energy and
79 mass transport through the lake-air turbulent heat flux. Biermann et al. (2014) and
80 Wang et al. (2015) discovered that the turbulent flux of Lake Nam co, which is
81 surrounded by wet grasslands, is actually very considerable but was often
82 underestimated in the model. The Source Region of the Yellow River (SRYR) is
83 located in the northeastern part of the TP and is known as the “water tower” of China
84 because it contains 48 lakes. The Ngoring and Gyaring (Sisters) Lakes are two major
85 lakes, and Ngoring Lake is the largest in the SRYR (Li et al., 2015; Wen et al., 2015).
86 In addition to the lakes, the forests, alpine meadow, wetlands, rivers, and glaciers
87 comprise the diverse underlying surfaces in the SRYR, with grassland accounting for
88 about 80% of the area (Mudassar et al., 2018). Consequently, the SRYR is an ideal
89 region for studying the turbulence over a heterogeneous land surface.

90 Observational studies have revealed that water vapor, heat, and energy exchange
91 occur over alpine meadows/wetlands (Zheng et al., 2015; Jia et al., 2017) and lakes



92 ([Li et al., 2015](#); [Wen et al., 2016](#)), and models have been used to simulated the effects
93 of the lakes on the cool and moist regional climate ([Wen et al., 2015](#); [Ao et al., 2018](#)).
94 However, the features of the boundary-layer turbulence over the heterogeneous
95 underlying surfaces and the effects of thermodynamic surface heterogeneity on the
96 turbulent flux in the SRYR remain unclear. Over the last few decades, the lakes have
97 shrunk and the grasslands have degraded in the SRYR due to climate change and the
98 excessive utilization of water resources ([Brierley et al., 2016](#); [Mudassar et al., 2018](#)).
99 It is essential to investigate the variation in the structure of boundary-layer and
100 turbulent heat flux with changes in the surface's thermal properties and the
101 background winds.

102 High resolution field measurements are extremely rare on the TP because of the
103 harsh environmental conditions, so few observational studies on the turbulence
104 characteristics and the turbulent heat flux have been conducted. Large eddy
105 simulation (LES) has the unique advantage of being accurate and able to describe
106 turbulence finely, and thus, it has been widely used to investigate the effects of
107 surface heterogeneity on turbulence ([Hadfield et al., 1991, 1992](#); [Kang and Lenschow,](#)
108 [2014](#)). However, little has been done to improve our understanding of how the surface
109 heterogeneity affects the boundary-layer turbulence, and the contributions of the
110 patch-induced motions to the turbulent flux and energy in the SRYR. Furthermore,
111 modeling the turbulence over the heterogeneous surface in the SRYR can not only lay
112 a basis for the analysis of the local energy and mass transport, but it can also provide a
113 quantitative reference for improving the parameterization schemes over



114 heterogeneous surfaces in weather and climate models.

115 In this study, we used LES to investigate the detailed turbulence characteristics in
116 the SRYR. Our primary focus was the impacts of the surface heat flux anomalies on
117 the turbulent kinetic energy (TKE), turbulent intensity, and turbulent flux. The
118 turbulence characteristics and turbulent fluxes in the surface and the entrainment
119 layers were investigated, too. This paper is arranged as follows Section 2 describes
120 the model and data used in this study. Section 3 discusses the modeled results in detail,
121 and section 4 provides a summary and discussion of our findings.

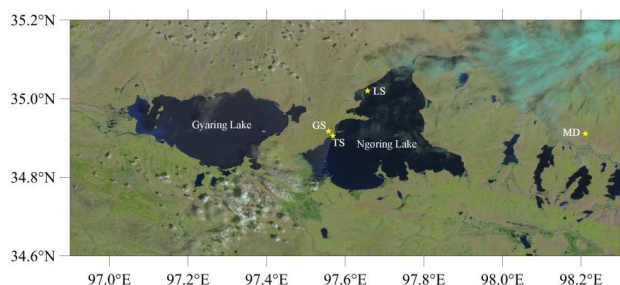
122 **2 data and methods**

123 **2.1 Study area and observations**

124 Ngoring Lake and Gyaring Lake (hereinafter referred to as the two lakes) are
125 located in the SRYR and are surrounded by the alpine meadow. Their mean elevation
126 is 4274 m above sea level. The study area is shown in fig.1. The turbulent flux and
127 standard atmospheric variables were measured over the lake and grassland surface.
128 The GPS radiosonde data from the field campaign on July 29, 2012, 30 m west of
129 Lake Ngoring (near the gradient tower station, TS) and at Madoi station (MD) located
130 30 km the east of the lake (34.918° N, 98.216° E, 4279 m AMSL), as well as the eddy
131 covariance data for Lake Station (LS) above the northwest of the lake (35.026° N,
132 97.652° E) and Grassland Station (GS) (34.913° N, 97.553° E) 1.5 km west of the
133 lake shore were used. For further details on the field campaign and the quality control
134 of the sounding and eddy covariance data, see [Li et al. \(2015\)](#) and [Li et al. \(2017\)](#).
135 The synoptic background near the surface and at 500 hPa and the distribution of the



136 wind components in the vertical and horizontal directions were investigated using the
137 ERA-Interim Reanalysis Data with a $1^\circ \times 1^\circ$ resolution collected at 12:30 LT and
138 18:30 LT (LT: local time, used in the whole study) on July 29, 2012, with a delimiting
139 a range of $32^\circ \text{ N} - 37^\circ \text{ N}$, $95^\circ \text{ E} - 100^\circ \text{ E}$, including the two lakes area ($34.8^\circ \text{ N} - 35^\circ \text{ N}$,
140 $97^\circ \text{ E} - 98^\circ \text{ E}$) and the surrounding grassland.



141 **Fig.1.** Map of the study area obtained using Landsat data, with the location of the observation
142 stations marked by yellow stars. The turbulent fluxes were measured at LS and GS stations. The
143 standard atmospheric variables were observed at the TS. MD is a fixed meteorological observatory
144 of the China Meteorological Administration.

145 2.2 Methods

146 2.2.1 Simulations set-up over the heterogeneous underlying surface

147 The U.K. Met Office large eddy model (LEM) version 2.4 (Gray et al., 2001) was
148 used in this paper. The LEM is a three-dimensional and non-hydrostatic numerical
149 model, which can be used to simulate a wide range of turbulence-scale and
150 cloud-scale problems with a high resolution. The domain size was $135 \text{ km} \times 30 \text{ km} \times$
151 6 km with a horizontal grid-spacing of 200 m. A vertically stretched grid with a
152 minimum spacing of 1.1 m was utilized in the surface layer and a maximum of 64.8 m
153 above 2000 m. Periodic lateral boundary conditions were applied, with a rigid lid at
154 the top of the model domain. To reduce the reflection of the internal gravity waves, a



155 Newtonian damping layer was applied above 3500 m. The surface boundary
156 conditions of the model were derived from the Monin-Obukhov similarity theory
157 using the Businger-Dyer functions. The subgrid model used in the LEM was based on
158 the Smagorinsky-Lilly approach (Brown et al., 1994). The potential temperature, wind
159 (u and v), and relative humidity profiles obtained during the field campaign on July 29,
160 2012, were used to initialize the 3D runs. The LEM was driven by the time-varying
161 sensible heat flux (SHF) and the latent heat flux (LHF) at the surface. The geostrophic
162 wind shear was calculated using the ERA-Interim geostrophic wind at the surface and
163 at 1500 m. The simulation time was 12 hours and the data were output every hour. In
164 this study, twelve 3D runs with different surface heat fluxes under various ambient
165 wind conditions were performed. Two of the runs were horizontally homogeneous
166 with a uniform grass surface under the conditions of wind (HOMW) and no wind
167 (HOM). The other cases were simulated with one (A1L) or two (A2L) lake patches in
168 the middle of the model domain. The surface heat flux anomaly was applied over a 30
169 km wide strip (two strips for A2L) extending the entire 30 km width of the domain in
170 the x-direction. Here the term heat flux refers to both the sensible and latent heat
171 fluxes. This can be viewed as representing Ngoring Lake and Gyaring Lake in the
172 SRYR. It should be noted that in this study the scale of the heterogeneity was large
173 enough to enable the formation of small eddies over the lake patches that could
174 coexist with the large-scale patch-induced circulations (Patton et al., 2005). Four
175 simulations (A1L, A2L, A1LW, and A2LW) were initialized using the surface heat
176 flux over the patch/patches measured at LS and the heat flux outside the patch/patches



177 measured at GS. This means that the heat flux into the modeled atmosphere decreases
178 as the number of patches increases. However, it is helpful to separate the effects of the
179 total increase in heating from the effects of the localization of the heating when
180 considering the consequences of an unresolved spatially changing heat flux for a
181 global model. In order to keep the total heat flux in the modeled atmosphere constant,
182 a “balanced” surface heat flux approach was used. Therefore, if the surface heat flux
183 observed at the GS is denoted as FGS , and the heat fluxes over the patch and outside
184 of the patch are denoted as FL and FG , respectively. FL and FG were calculated using
185 the following equations:

$$186 \quad FL = FGS \times (SL/ST) \quad (1)$$

$$187 \quad FG = FGS \times (SG/ST) \quad (2)$$

188 where ST , SL , and SG are the squares of the model domain, the patch, and the outside
189 of it, respectively. Another four simulations ($A1L_C$, $A2L_C$, $A1LW_C$, and
190 $A2LW_C$) were performed using this balanced surface heat flux approach. The
191 heterogeneous initial conditions were used in the surface heat flux anomaly
192 simulations. The initial profiles over the patch/patches were derived using the data
193 from TS station, and the data observed at the MD station used for the outside
194 patch/patches. Various ambient wind conditions were also used for the surface heat
195 flux anomaly runs. The parameters and the conditions of each run are listed in Table 1
196 for convenience. Sketches of the heterogeneous surface and of the surface heat fluxes
197 over the lake patches and the outside patches for the unbalanced and balanced cases
198 are depicted in fig. 2.



199 The initial potential temperature and special humidity are shown in fig. 2h, and the
200 horizontal components of the wind profiles and the geostrophic wind are shown in fig.
201 2g. A stable layer was found over the grass and a 200 m convective boundary layer
202 (CBL) was found over the lake at 06:30 LT. The special humidity profiles show that
203 the air tends to be moister over the lake (dash lines in fig. 2h). The study area is
204 characterized by a considerable surface heat flux and high wind speeds in the daytime.
205 The stronger initial velocity is from the GS, which recorded wind speeds of up to 10
206 m s^{-1} below 500 m.

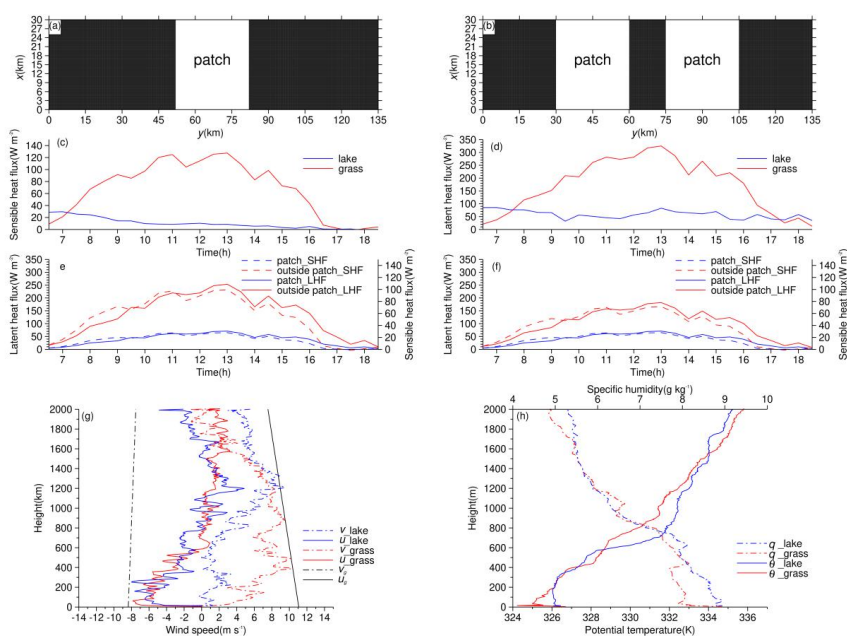
207 **Table 1**

208 Parameters for the 3D simulations.

Name	Wind field	Surface heat flux (SHF and LHF)	Number of lake patches	Size of Lake patch (km)
HOM	without wind	FGS	0	-
HOMW	initial wind + geostrophic wind	FGS	0	-
A1L	without wind	lake patch: FLS (Heat flux that observed at LS); outside patch: FGS	1	30
A2L	without wind	lake patches: FLS (Heat flux that observed at LS); outside patches: FGS	2	30 and 30
A1LW	initial wind + geostrophic wind	Same as A1L	1	30
A2LW	initial wind + geostrophic wind	Same as A2L	2	30 and 30



A1LNG	initial wind	Same as A1L	1	30
A2LNG	initial wind	Same as A2L	2	30 and 30
A1L_C	without wind	lake patch: (SL/ST) ×FGS=(30/135)× FGS; outside patch: (SG/ST)×FGS =(105/135)×FGS	1	30
A2L_C	without wind	lake patches: (SL/ST)×FGS =(30/135)×FGS; outside patches: (SG/ST)×FGS =(75/135)×FGS	2	30 and 30
A1LW_C	initial wind + geostrophic wind	Same as A1L_C	1	30
A2LW_C	initial wind + geostrophic wind	Same as A2L_C	2	30 and 30





209 **Fig. 2.** Sketch of the heterogeneous surface (a and b), (c and d) surface sensible heat flux and
210 latent heat flux over the grassland (red line) and the lake (blue line) from observation. The SHF
211 and the LHF for runs with (e) one and (f) two lake patches and a constant heat flux. Figures 2g
212 and 2h show the initial profiles of the winds (solid lines for u , dash lines for v), potential
213 temperature (solid lines), and special humidity (dash lines) over the lake patches (blue lines) and
214 outside of the patches (red lines). The input geostrophic winds are also shown (black lines).

215 2.2.2 statistical analysis

216 According to turbulence theory, a physical quantity ϕ has two parts, i.e., the
217 horizontal average $\langle\phi\rangle$ and the turbulent fluctuation ϕ' , and

$$218 \quad \phi = \langle\phi\rangle + \phi' \quad (3)$$

219 This equation usually works in cases with a homogeneous surface. The variances of
220 velocity and the potential temperature variances (σ_v^2 , σ_w^2 , σ_θ^2) are calculated from
221 v' , w' , and θ' , respectively. For a heterogeneous surface, phase-averaged analysis
222 helps separate the patch-induced circulations from the random turbulent motions. This
223 method has been applied in studies of the one-dimension and two-dimension
224 heterogeneities (Matthias et al., 2014; Kang and Lenschow, 2014; Shen et al., 2016)
225 and complex and irregular heterogeneities (Maronga and Raasch, 2013). In this study,
226 one-dimensional heterogeneous (in the y direction) simulations were performed for
227 which ϕ can be decomposed into three parts:

$$228 \quad \phi = \langle\phi\rangle + \phi_{hi} + \phi_s \quad (4)$$

229 Where $\langle\phi\rangle$ is the horizontal average; ϕ_{hi} is the heterogeneity-induced part which is
230 the averaged ϕ across the domain in the y direction; and ϕ_s is from the background
231 turbulence. The variances of velocity and the potential temperature induced by the
232 heterogeneity ($[\sigma_v^2]_{hi}$, $[\sigma_w^2]_{hi}$, $[\sigma_\theta^2]_{hi}$) are calculated from v_{hi} , w_{hi} , and θ_{hi} ,



233 respectively.

234 Phase-averaged analysis was also used to obtain the patch-induced component of
235 the turbulent fluxes. We multiplied w and ϕ with both in the forms of Equation (4),
236 and derived the total vertical transport of ϕ :

$$237 \quad \overline{w\phi} = \overline{w}\langle\phi\rangle + \overline{w_{hi}\phi_{hi}} + \overline{w_s\phi_s} \quad (5)$$

238 Since the horizontal average vertical velocity $\langle w \rangle$ is approximately zero in the LES,
239 the turbulent fluxes were divided into two parts: a patch-induced circulation induced
240 part and a background turbulence induced part:

$$241 \quad \overline{w\phi} = \overline{w_{hi}\phi_{hi}} + \overline{w_s\phi_s} \quad (6)$$

242 Moreover, the total kinetic energy e can be written as two parts, e_{hi} and e_s , which
243 represent the contributions from the patch-induced and background turbulence:

$$244 \quad e = e_{hi} + e_s \quad (7)$$

$$245 \quad e_{hi} = \left(\langle u_{hi}^2 \rangle + \langle v_{hi}^2 \rangle + \langle w_{hi}^2 \rangle \right) / 2 \quad (8)$$

$$246 \quad e_s = \left(\langle u_s^2 \rangle + \langle v_s^2 \rangle + \langle w_s^2 \rangle \right) / 2 \quad (9)$$

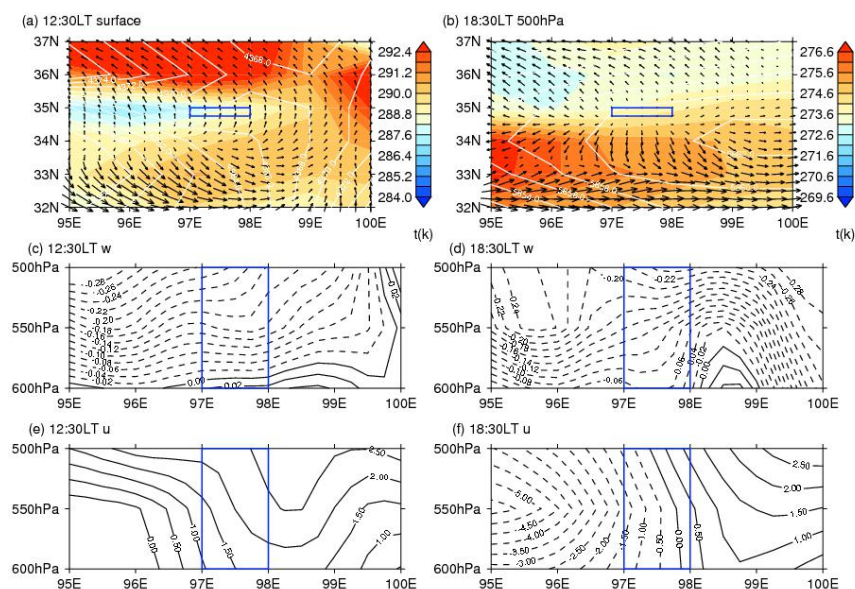
247 **3. Results**

248 **3.1. Synoptic background and wind components' distribution**

249 In order to investigate the existence of a daytime lake breeze (the divergent flows
250 over the lake surface and the downdrafts overlying it) using the ERA-Interim
251 reanalysis data for the two lakes area (34.8° N–35° N, 97° E–98° E; blue box in fig. 3),
252 the wind field, temperature field, and geopotential height field at the surface (600 hPa,
253 ~4200 m) and at 500 hPa (~5500 m) at 12:30 LT and 18:30 LT on July 29, 2012, were
254 analyzed. A cyclone controlled the entire region above the surface at 12:30 LT (Fig. 3a)



255 and divergent flow occurred at 500 hPa at 18:30 LT (Fig. 3b). The vertical sections of
256 the two wind components (u and w) were also depicted to further ascertain the
257 distribution of the wind field in the longitude and latitude directions. It should be
258 noted that downdrafts are dominant below 500 hPa in the two lakes area during the
259 day (Figs. 3c and 3d). As can be seen from fig. 3f, distinct divergent zonal wind (u)
260 flows existed in the two lakes area at 18:30 LT. The wind speed derived from the GPS
261 sounding at 12:30 LT is larger than that at 18:30 LT below 2 km (see fig. S1 in
262 supplement), indicating that the larger background flows covered up the divergent
263 wind flow at 12:30 LT. Evidently, it is difficult to directly observe the lake breeze
264 circulation due to the synoptic background, but the downdrafts and the divergent
265 zonal wind in the two lakes area demonstrate the existence of a lake breeze. In the
266 following sections, the turbulence characteristics over the heterogeneous underlying
267 surface are simulated and the effects of the patch-induced circulation are analyzed.





268 **Fig. 3.** Synoptic background on July 29, 2012. Blue boxes represent the two lakes area. (a) and (b)
269 show the wind field, temperature field, and geopotential height field at the surface (600 hPa,
270 ~4200 m, 10 m wind field, 2 m temperature field; Fig. 3a) at 12:30 LT and at 500 hPa (~5500 m)
271 at 18:30 LT (Fig. 3b). The vertical wind (w , figs. 3c and 3d) and the zonal wind (u , figs. 3e and 3f)
272 below 500 hPa are also shown.

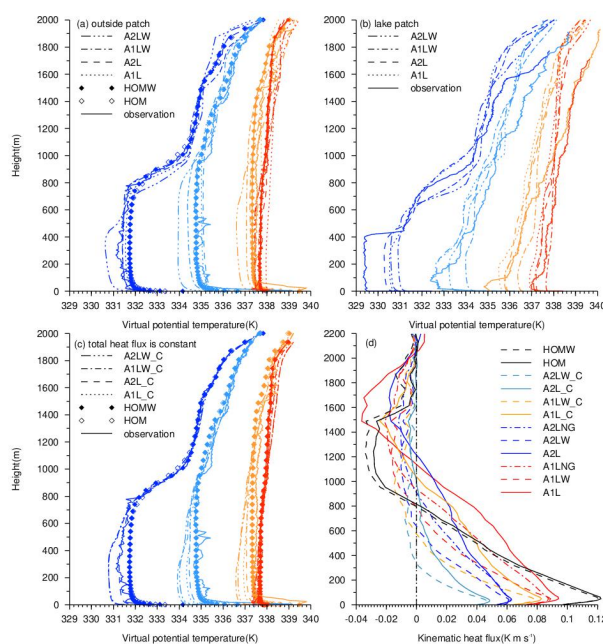
273 **3.2 Effects of the underlying surfaces and background flows on the** 274 **boundary-layer turbulence**

275 **3.2.1 Performance of the LEM and the height of the boundary layer over** 276 **homogeneously heated and heterogeneously heated surfaces**

277 In order to inspect the performance of the LEM over the heterogeneously heated
278 surfaces, the simulated virtual potential temperature (θ_v) over the lake patch/patches
279 and outside were compared with the observations. In addition, by keeping the total
280 surface heat flux into the modeled domain constant, the profiles of the simulated
281 virtual potential temperature over the homogeneous and heterogeneous surfaces were
282 compared in order to investigate the effects of surface heterogeneity on the structure
283 of the boundary layer. The profiles of the kinematic heat flux for all of the runs were
284 used to determine the height of the boundary layer. Figures 4a and 4b compare the
285 simulated profiles of the virtual potential temperature over and outside of the lake
286 patch with the corresponding observations (solid lines) over the grassland and the lake
287 surfaces at different times. In order to account for the effects of the unrepresented
288 large-scale forcing, the simulated horizontally averaged potential temperature, water
289 vapor mixing ratio, and horizontal wind (u and v) were relaxed to those observed
290 using the radiosondes with at a 3 h interval during the simulation (Marshall et al.,



291 2008; Huang et al., 2009). The time series of the kinetic energy (see fig. S2 in
292 supplement) indicates that the equilibration time of the model is approximately 3
293 hours.



294 **Fig. 4.** The profile of the horizontal, averaged virtual potential temperature for the observations
295 and all of the runs over the lake patches (a) and outside the of the (b) lake patches. (c) same as (a)
296 except, for heterogeneous and homogeneous runs with constant surface heat fluxes. Legends for
297 (a), (b), and (c): dark and light blue represent the results at 09:30 and 12:30, respectively, and the
298 orange and red lines show the results at 15:30 and 18:30, respectively. (d) The kinematic heat
299 fluxes for all of the runs.

300 The observation profiles (solid lines in fig. 4a) show that the depth of the
301 convective boundary layer (CBL) over the grassland increases from 700 m at 09:30 to
302 1.1 km at 12:30 to 1.5 km at 15:30 to 1.9 km at 18:30. The inversion layer above the
303 CBL is completely eroded by the turbulence after 12:30. The virtual potential
304 temperature in the well-mixed CBL over the grassland increases approximately 7 K

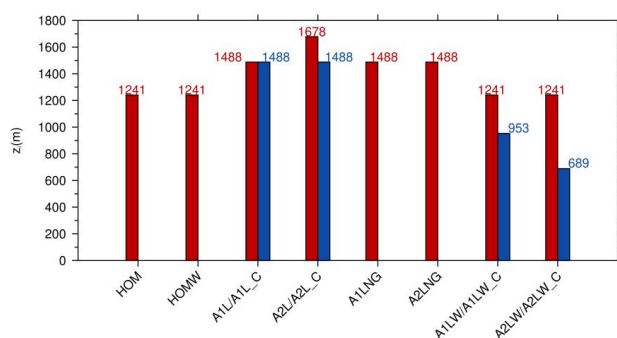


305 from 09:30 to 18:30. The CBL over the heterogeneous surfaces with background wind
306 is cooler and shallower than that over the homogeneous surfaces. This may be
307 because the air blowing from the lake patches cools the CBL of the outside patches
308 that are downwind, which inhibits the development of the CBL. In addition, the model
309 profiles of the virtual potential temperatures over the homogeneously heated and
310 heterogeneously heated surfaces with no background wind have very similar
311 structures and are close to the sounding profiles. This is similar to the modeling
312 results of [Liu et al. \(2011\)](#). The observed virtual potential temperature over the lake
313 surface (solid lines in fig. 4b) shows that the CBL changes to stable stratification as
314 the radiation increases after sunrise, and the modeled θ_v over the patches is about
315 1.0 K warmer than the observed θ_v . As in fig. 4a, fig. 4c also shows that the
316 background wind over the heterogeneous surface inhibits the growth of the CBL.

317 In this study, according to [Sullivan et al. \(1998\)](#), the height of the boundary layer (z_i)
318 was determined using the minimum kinematic heat flux of the simulated results. As
319 can be seen, the maximum surface heat fluxes were relatively large over the
320 homogeneously heated surface, while smaller surface heat fluxes occurred for the case
321 with two lake patches. Compared to the unbalanced cases (A1L, A1LW, A2L, and
322 A2LW), less heat flux was introduced in the balanced cases (A1L_C, A1LW_C,
323 A2L_C, and A2LW_C) and lower CBLs occurred, especially with a wind field (blue
324 bars in fig. 5). The kinematic heat fluxes decreased to zero at higher altitude over the
325 heterogeneously heated surface. When the height continues to rise, the region of
326 negative heat flux is often called the entrainment layer, which is thicker in the cases



327 with background wind. The heights of the CBL indicate (Fig. 5) that the surface heat
328 flux anomaly may contribute to the deepening of the mixed layer, thus increasing the
329 CBL height. However, the shear generated by the background wind strengthens the
330 turbulent exchange between the entrainment layer and the free atmosphere, resulting
331 in an excessively thick entrainment layer, which, however, inhibits the upward
332 development of CBL.



333 **Fig. 5.** Bar chart of the CBL height for each run marked with a concrete value. The red bars from
334 left to right represent runs HOM, HOMW, A1L, A2L, A1LNG, A2LNG, A1LW, and A2LW. The
335 blue bars from left to right represent runs A1L_C, A2L_C, A1LW_C, and A2LW_C.

336 3.2.2 Effects of the surface heat flux anomalies and background winds on the 337 turbulent kinetic energy

338 Local circulations will be induced by differential heating, and the turbulent kinetic
339 energy (TKE) determines the transport of the momentum, heat, and moisture through
340 the boundary layer (Tyagi and Satyanarayana, 2013). Thus, the thermal circulations
341 induced by the lake patches were simulated to investigate the effects of the
342 heterogeneous heating on the spatial distribution of the TKE.

343 Fig. 6 shows the vertical distribution of the TKE and the wind vectors over the
344 homogeneous and heterogeneous surfaces with no background wind at 15:30 LT. Over

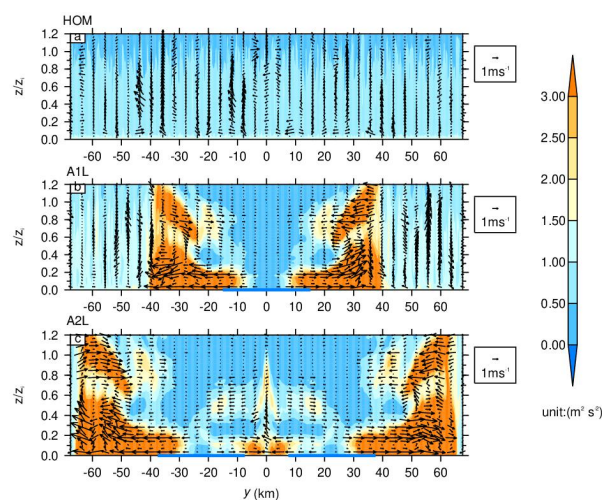


345 the homogeneous surface (Fig. 6a), a relatively uniform TKE with a larger magnitude
346 exists within a much shallower CBL (below $0.1 z_i$), which overlies the scattered and
347 disordered wind vectors throughout the domain. Over the heterogeneous surfaces, the
348 large TKE values are distributed on both sides of the lake patches below $0.5 z_i$, and
349 the divergent winds extend to about 30 km away. In addition, a larger TKE and
350 convergent winds occurred in the upper level of the CBL (Figs. 6b, 6c). Moreover, the
351 air flow between the two lake patches led to a convergent region (updrafts in $y = 0$ km;
352 Fig. 6c). This is consistent with the results of [Avissar and Schmidt \(1998\)](#), who
353 demonstrated that turbulent eddies are randomly distributed over a homogeneous
354 surface, but the TKE exhibits two maxima near the ground surface and the top of the
355 CBL, which is in agreement with the patch-induced circulations. Overall, Figure 6
356 illustrates that the distributions of the TKE and the patch-induced circulations are
357 symmetrical on both sides of the lake patches, while the distribution is random with
358 smaller TKE values over the homogeneous surface.

359 Furthermore, the ratios of the horizontally averaged TKEs in the model domain of
360 the different runs were calculated to examine the effects of the surface anomalies and
361 the ambient winds on the TKE. As is shown in Table 2, the TKEs for the cases with
362 one or two lake patches are about twice that of the TKE of the case without patches
363 (columns 2–3), but the ambient wind leads to a reduction in the impacts of surface
364 flux heterogeneity on the TKE (columns 4–5 and 6–7). This is consistent with the
365 results of [Avissar and Schmidt \(1998\)](#), who reported that a weak background wind of
366 2.5 m/s is strong enough to considerably reduce the impact of the ground-surface



367 heterogeneity on the CBL. For the homogeneous cases, the TKE increases under the
 368 background wind conditions due to the increase in the sheared turbulence. For the
 369 runs with balanced surface heat fluxes (A1L_C, A2L_C, A1LW_C, A2LW_C), the
 370 effects of the heterogeneity on the TKE are less significant, especially for the cases
 371 with more lake patches, but the effects of the background winds on the TKE tend to
 372 be large.



373 **Fig. 6.** The y - z cross sections of the TKE (contour) with superimposed wind vectors composed of
 374 v and w wind over (a) homogeneously heated and (b and c) heterogeneously heated surfaces. The
 375 blue lines on the y -axis represent the lake patches.

376 **Table 2**

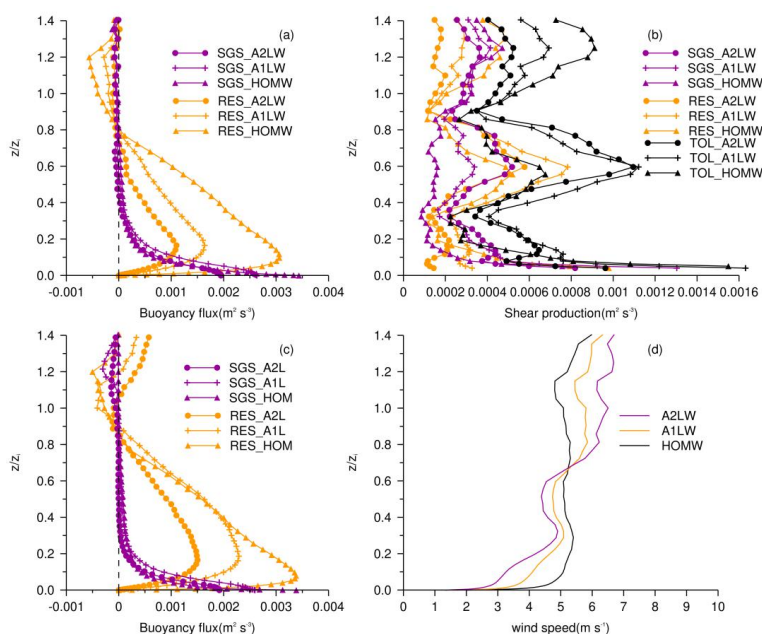
377 The ratio of the TKEs of the different runs. Max, Min, and Mean stand for the maximum,
 378 minimum, and mean ratios of the TKE in the model domain, respectively.

Ratio of TKE	A1L/HOM	A2L/HOM	A1LW/HOMW	A2LW/HOMW	A1LW/A1L	A2LW/A2L	HOMW/HOM
Max	3.31	3.30	1.47	1.42	1.01	0.88	2.09
Min	1.15	1.18	0.80	0.59	0.61	0.40	0.83
Mean	2.00	2.04	1.09	0.95	0.79	0.68	1.41
Ratio of TKE	A1L_C/HOM	A2L_C/HOM	A1LW_C/HOMW	A2LW_C/HOMW	A1LW_C/A1L_C	A2LW_C/A2L_C	HOMW/HOM
Max	2.21	1.44	1.12	0.74	1.84	2.17	2.09
Min	0.43	0.23	0.63	0.31	0.80	0.47	0.83



Mean 1.09 0.75 0.91 0.54 1.22 1.07 1.41

379 To further investigate the effects of the surface heat flux anomalies on the
 380 development of turbulence, it is instructive to examine the vertical profiles of the
 381 buoyancy and shear production terms in the TKE budget equation, which is from the
 382 contributions of the resolved (RES) and subgrid (SGS) eddies (Figs. 7a, 7b, and 7c).



383 **Fig. 7.** Vertical profiles of (a) the buoyancy flux and (b) the shear production term for runs
 384 HOMW, A1LW, and A2LW with background flows, and (c) the profiles of the buoyancy flux for
 385 runs HOM, A1L, and A2L without background flows. (d) The simulated horizontal wind versus
 386 height for runs HOMW, A1LW, and A2LW. The resolved and subgrid results are presented as
 387 yellow and purple lines, respectively. The black lines in (b) are the total (resolved and subgrid
 388 scale) shear production term.

389 Figures 7a and 7c show that the RES buoyancy production decreases as the number
 390 of patches increases and the SGS buoyancy contributions are negligible, except in the
 391 surface layer. Below $0.9 z_i$, the larger RES shear production occurs in the case with
 392 lake patches (Fig. 7b) and the contribution of the SGS shear production is



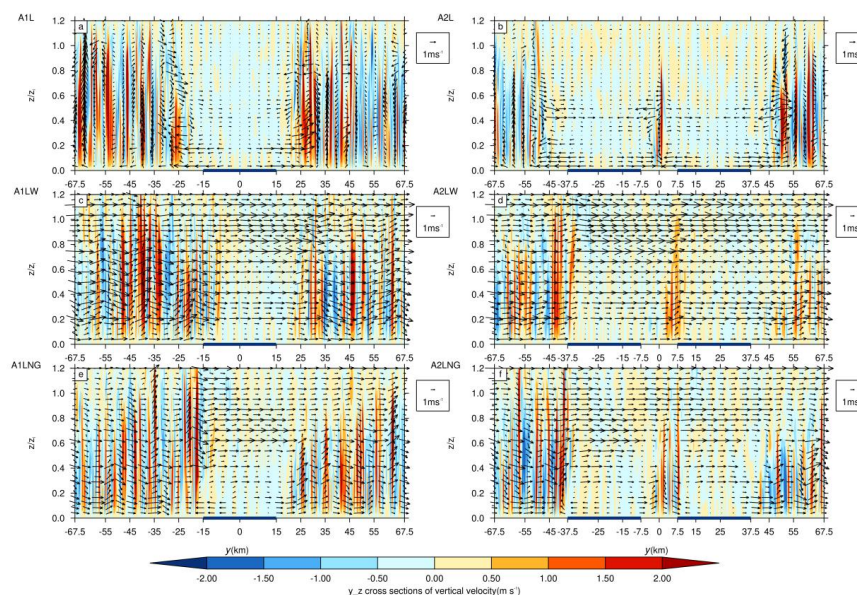
393 considerable (Fig. 7b), which is significant in the CBL for the cases with surface flux
394 anomalies. Thus, the total shear productions (black lines in fig. 7b) of the cases with
395 heterogeneous surfaces are larger. The background winds (Fig. 7d) are weaker for the
396 cases with lake patches below $0.65 z_i$, but the corresponding total shear production
397 term is larger, which shows that the patch-induced circulations are conducive to more
398 shear in the CBL.

399 **3.2.3 Effects of the background flows on the circulations**

400 In order to investigate the effects of the background winds on the patch-induced
401 circulations, the vertical distributions of the vertical velocity and wind fields for the
402 runs with and without background winds were compared. In fig. 8, the patch-induced
403 circulations are not easy to distinguish in the cases with background winds (about
404 13.9 m/s above a height of 1.2 km) due to the cancellation of the local pressure
405 gradient by the synoptic pressure gradient, which is consistent with the results of
406 [Crosman and Horel \(2010\)](#). This also indicates that the boundary-layer convection
407 tends to weaken as the number of lake patches increases (the maximum updrafts are
408 4.8 m s^{-1} , 4.2 m s^{-1} , 3.5 m s^{-1} , and 3.3 m s^{-1} for runs A1L, A2L, A1LW, and A2LW,
409 respectively). Moreover, the wind fields for the cases without geostrophic winds
410 exhibit divergent flows over the lake patches. As in the study of [Kang and Lenschow](#)
411 [\(2014\)](#), our study also confirms that the symmetrical patch-induced circulations and
412 the intensity of the convection become indistinguishable and weak under the
413 background flow conditions. However, the divergent flows in the lower level are still
414 visible when the geostrophic wind is removed (A1LNG and A2LNG in figs. 8e and



415 8f).



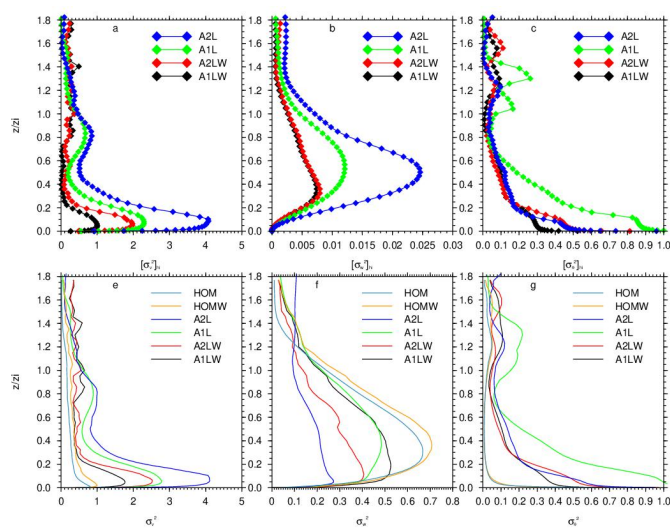
416 **Fig. 8.** Instantaneous y - z cross sections of the vertical velocity (m s^{-1}) and wind vectors above the
417 heterogeneous surfaces for runs (a and b) without and (c and d) with background winds, and (e
418 and f) with the geostrophic wind removed. The blue lines represent the lake patches.

419 **3.3 Effects of patch-induced circulation on the turbulent intensity** 420 **and heat flux**

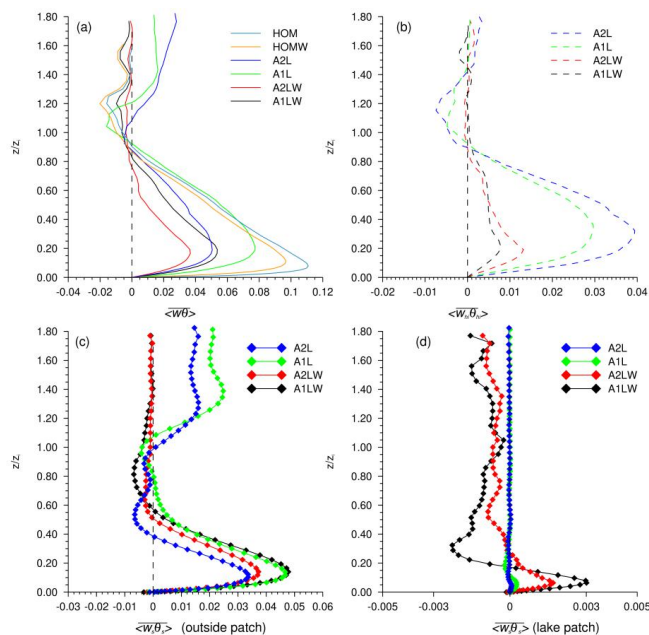
421 We used the phase-averaged method to decompose the contributions of the
422 turbulent intensity and the heat flux from the patch-induced circulations and the
423 background turbulence and to quantitatively analyze the heterogeneity-induced
424 contribution to the turbulent intensity. For the variance of the velocity, the horizontal
425 (Fig. 9a) and vertical (Fig. 9b) variances induced by the heterogeneity increase as the
426 number of lake patches increases, and the horizontal variance is larger than the
427 vertical variance. However, the background flows tend to decrease both the



428 patch-induced (Figs. 9a, 9b) and total (Figs 9e, 9f) turbulent intensity.



429 **Fig. 9.** (a, b, c) Heterogeneity-induced and (e, f, g) total dimensionless turbulence statistics for
 430 runs HOM, HOMW, A1L, A1LW, A2L, and A2LW. Shown are the profiles of the (a, e) v variance,
 431 (b, f) w variance, and (c, g) θ variance.



432 **Fig. 10.** (a) Area-averaged total turbulent heat flux (solid lines) and (b) heterogeneity-induced



433 turbulent heat flux (dash lines). The background turbulence (lines with diamond symbols) of heat
434 flux over the (c) grassland and (d) lake patches.

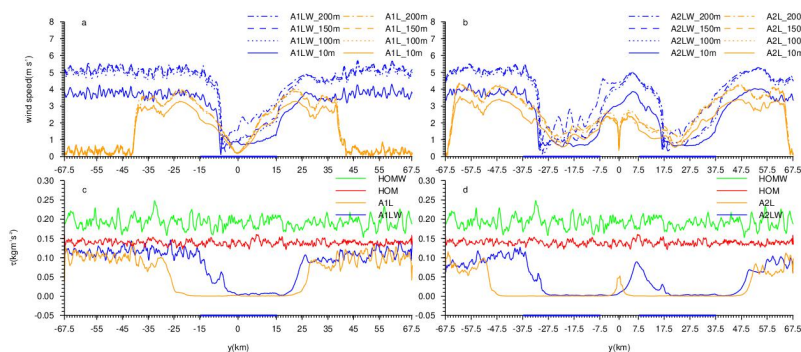
435 This also shows that a larger difference in the variances of the horizontal velocity
436 occurs in the surface layer and gradually decreases with height (Figs. 9a and 9e),
437 which means that the effects of the surface properties on the horizontal turbulence
438 diminish with height in the CBL. In this respect, our results are similar to those found
439 by Wang et al. (2011), Shao et al. (2013), and Frederik and Matthias (2018). The total
440 horizontal turbulent intensity is mainly from the contribution of the patch-induced
441 circulations and is larger than that of the homogeneous cases (Fig. 9e), which tends to
442 become stronger as the number of patches increases but becomes weaker as the total
443 vertical turbulent intensity increases (Fig. 9f, same as in the cases with balanced
444 surface fluxes). It should be noted that the contribution of the patch-induced
445 circulations to the vertical velocity variance is no more than 10% (Fig. 9b and 9f),
446 which implies that the background turbulence contributes more to the fluctuations in
447 the vertical velocity than to those in the horizontal velocity. Figures 9c and 9g show
448 that the patch-induced motions make the largest contribution to the variances of the
449 potential temperature. However, the background winds decrease the variances of the
450 potential temperature and decrease the impact of the surface heterogeneity on the
451 variances of the potential temperature.

452 Using the same method, we analyzed the contributions of the patch-induced and
453 background turbulence to the heat flux. Figure 10a shows that as the number of lake
454 patches increases, the area-averaged total heat flux decreases in both the mixed and



455 entrainment layers, and the balanced surface heat flux cases exhibit similar variations
456 (see fig.03 in supplement). The patch-induced transport of the heat flux increases as
457 the number of lake patches increases (Fig. 10b). The patch-induced motions
458 contribute up to 80% of the heat flux in run A2L, which has unbalanced surface fluxes
459 (Fig. 10b), and 61% in run A2L_C which has balanced surface fluxes (see fig. S3 in
460 supplement). It should also be noted that the background winds tend to decrease the
461 heat flux transport over the heterogeneous surfaces. As is shown in figs. 10c and 10e,
462 the contribution of the background turbulence to the local heat flux is larger over the
463 region outside of the lake patches than over the patches. We hope that the results of
464 our analysis of the contributions of the heterogeneity-induced circulation and
465 background turbulence to the turbulence intensity and the heat flux over a
466 heterogeneous surface will provide a basis for further studies of the local energy and
467 mass transport in the SRYR over the TP.

468 3.4 Turbulence in the surface layer

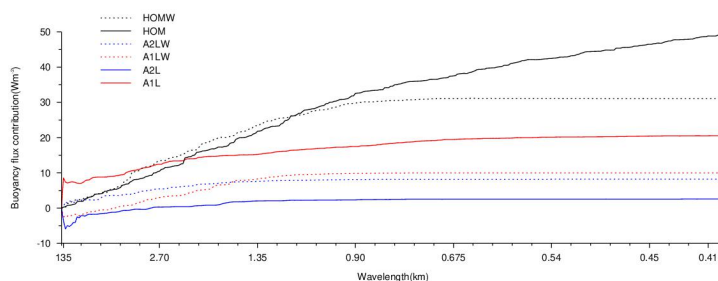


469 **Fig. 11.** Variations in the (a and b) wind speed and (c and d) Reynolds stress in the horizontal
470 direction below 200 m for the cases with (blue lines) and without (yellow lines) background flows.

471 The frictional velocity (u^*) is a critical parameter in the turbulence exchange near



472 the surface, and it plays an important role in the transport of momentum in the
473 boundary layer. Patil et al. (2016) reported that the frictional velocity increases with
474 increasing wind speed under lower wind speed conditions in the surface layer. Thus,
475 we focused on the variations in the wind speed (Figs. 11a and 11b) and Reynolds
476 stress (Figs. 11c and 11d) in the horizontal direction below 200 m in order to
477 investigate the effects of the patch-induced motions on the momentum flux in the
478 surface layer for various background winds. It was found that the inland extension of
479 the patch-induced divergent flows reached about 25 km with no background winds
480 (yellow curves in figs. 11a and 11b). The speed of the divergent winds increases from
481 the lake patches to the outside and increases with height below 200 m with and
482 without background winds. The wind speeds decrease rapidly (4.0 m s^{-1}) within 10 km
483 along with the wind blowing from west of the lake patches to east of the lake patch,
484 and then, the wind speeds increase steadily (blue curves in figs. 11a and 11b). The
485 changes in the surface winds and surface properties have significant effects on the
486 turbulent momentum flux. Figures 11c and 11d show that the transport of the
487 momentum flux is smaller over the heterogeneous surface. The consistent variations
488 in the wind speeds and the turbulent stresses illustrate that the lake patches alter the
489 spatial distribution of the turbulent stress, which would further affect the surface wind
490 speeds, especially over the land-lake boundary regions.



491 **Fig. 12.** The cumulative contribution of the buoyancy fluxes of all wavelengths (km) at a height of
492 50 m for runs HOM, HOMW, A1L, A2L, A1LW, and A2LW.

493 In order to quantify the contributions of the buoyancy fluxes, due to the different
494 scales of the eddies, we calculated the ogives, which are the running integrals of the
495 cospectral densities (Friehe et al., 1991), and used these values to show the
496 cumulative contribution to the fluxes of all of the wavelengths (Brooks and Rogers,
497 2000). The ogive curves (Fig. 12) show that the small eddies make a significant
498 contribution to the buoyancy fluxes over the homogeneous surfaces with no
499 background winds (solid black line). The background wind increases the buoyancy
500 flux for wavelengths larger than about 1.1 km and decreases it for smaller
501 wavelengths based on a comparison of cases HOM (solid black line) and HOMW
502 (black dotted line). The above results confirms that the heat transport is enhanced by
503 the large eddies but is weakened by the small eddies, especially under the control of
504 the background wind. The buoyancy flux for a wavelength larger than about 2.2 km
505 makes a greater contribution in the case with one lake patch without background wind
506 (solid red line). The buoyancy fluxes for wavelengths of greater than 2.7 km are
507 transported downward for the case with two lake patches. For the case with one lake
508 patch, the background flows tend to decrease the transport of the buoyancy flux for



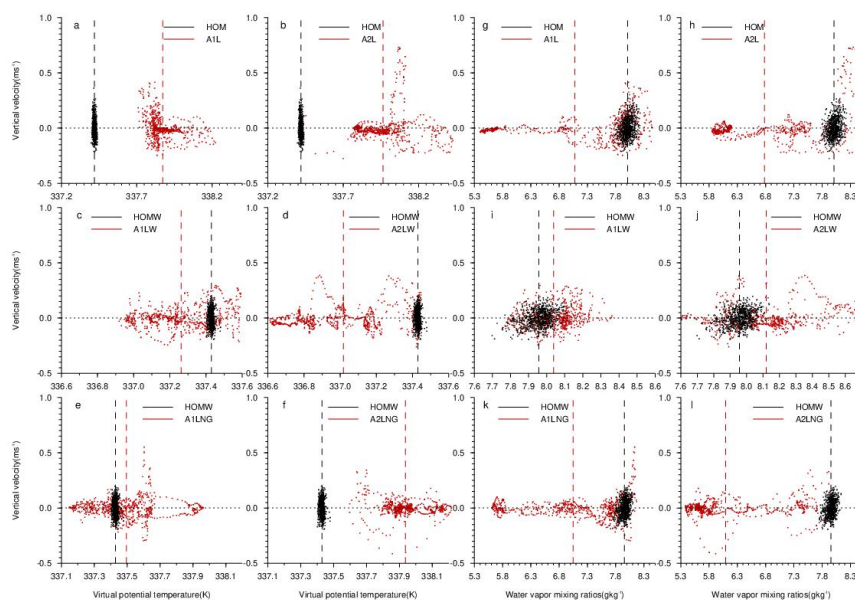
509 larger wavelengths near the surface (red dotted line) due to the stronger horizontal
510 wind (Fig. 7d), and they help to transport the buoyancy fluxes downward for
511 wavelengths larger than 3.3 km. However, for the case with two lake patches, the
512 background wind causes the large eddies to transport the buoyancy fluxes upward.
513 Thus, increasing the number of lake patches leads to more patch-induced motions, but
514 this does not tend to enhance the ability of the wind to transport heat. It is concluded
515 that slightly more of the buoyancy flux of the case with one lake patch is transported
516 by the small eddies with wavelengths of less than 1.5 km compared with the case with
517 two lake patches and background wind conditions (red and blue dotted lines).

518 **3.5 The characteristics of the boundary-layer turbulence in the** 519 **entrainment layer**

520 The LES study conducted by [Matthias et al. \(2014\)](#) found that there is increased
521 entrainment from the more strongly heated surface patch cases compared to the
522 homogeneous cases, and the impact of the heterogeneity on entrainment vanishes due
523 to horizontal mixing if the mean flow is aligned perpendicular to the border between
524 the differentially heated patches. To investigate the effects of the thermal properties of
525 the heterogeneous surface and the background flows on the turbulence in the
526 entrainment layer, the characteristics of the heat flux in the entrainment layer were
527 analyzed. Our simulated results show that the downward transport of the heat flux
528 decreases as the number of lake patches increases in the entrainment layer for both the
529 wind and no wind cases (Figs. 4d and 10a), which is also true in the balanced heat
530 flux runs.



531 By comparing the maximum and minimum vertical velocities at the top of the
532 boundary layer (Table 3), we found that the convective intensity of the entrainment
533 layer in the case with two lake patches and no wind fields is stronger, but it is
534 weakened by the background flows. Whereas, it decreases as the number of lake
535 patches increases in the balanced heat flux cases (A1L_C and A2L_C), corresponding
536 to a smaller TKE (Table 2) and total turbulent intensity (Fig. 9f). [Huang et al. \(2007\)](#)
537 pointed out that an appropriate surface heat flux and background flows maintain the
538 convective roll, and our simulations demonstrate this roll-like convection (see fig. S4
539 in supplement), which is mainly induced by the persistence of the background
540 turbulence with stronger geostrophic winds of 7–11 m s⁻¹ (black lines in fig. 2g).
541 However, [Maronga and Raasch \(2013\)](#) found that a higher wind speed of 6 m s⁻¹
542 generates convective rolls derived from the secondary circulation over a complex
543 heterogeneous surface.





544 **Fig. 13.** (a, b, c, d, e, f) The joint vertical velocities and virtual potential temperatures and (g, h, i, j,
545 k, l) water vapor mixing ratios at the top of the CBL for the homogeneous and heterogeneous runs.
546 The black dotted lines represent the mean vertical velocity, and the black and red dashed lines
547 show the mean virtual potential temperatures and water vapor mixing ratios, respectively.

548 In addition, the boundary layer variables (including the vertical velocities, virtual
549 potential temperatures, and water vapor mixing ratios) in the entrainment layer are
550 also subject to the effects of the surface heterogeneity. Figure 13 shows the joint
551 distribution of the vertical velocities and the virtual potential temperatures, as well as
552 the vertical velocities and water vapor mixing ratios. Comparing to the
553 homogeneously heated cases, the increased downdrafts mainly occur over the lake
554 patches, and they carry more warm, dry air down from the free atmosphere (Figs. 13a,
555 13b, 13g, and 13h), which is due to the convergent airflow caused by the
556 patch-induced circulations at the top of the CBL. This effect is much more evident in
557 the case with two lake patches, but it is weakened by the gradually strengthening
558 background flows (except figs. 13a, 13b, 13g, and 13h). We obtained the same results
559 for the balanced cases. In particular, colder and moister air exists in the entrainment
560 layer in the cases with ambient winds (A1LW_C and A2LW_C).

561 **Table 3**

562 The maximum and minimum vertical velocities at the top of the boundary layer in cases A1LW,
563 A2LW, A1LW_C, A2LW_C, A1LNG, A2LNG, A1L, A2L and HOM, HOMW.

Case	A1LW	A2LW	A1L	A2L	HOM	HOMW
W(max)	4.01	3.54	5.42	5.55	4.37	4.48
W(min)	-2.29	-1.98	-2.40	-3.12	-2.15	-2.50
Case	A1LW_C	A2LW_C	A1L_C	A2L_C	A1LNG	A2LNG
W(max)	2.97	2.06	4.26	2.46	3.90	4.46
W(min)	-2.03	-1.25	-1.91	-1.27	-2.23	-1.97



564 **4 Summary and discussion**

565 The downdrafts and divergent zonal wind in the two lakes area obtained from the
566 ERA-Interim reanalysis data indicate the existence of a lake breeze in the SRYR. Ten
567 runs of the 1D strip-like distribution of the surface heat flux and two homogeneously
568 heated runs based on the observations made during the summer of 2012 in the
569 Ngoring Lake Basin were conducted in order to investigate the effects of the
570 patch-induced circulations on the boundary-layer turbulence and its energy transport
571 at the lake-air and grass-air interfaces, and the influence of the background flows also
572 be considered.

573 The thermodynamic heterogeneity of the surface is conducive to deepening the
574 mixed layer, thus increasing the CBL height and enhancing the TKE when there are
575 no background flows. The background flows bring shear, resulting in an excessively
576 thick entrainment layer, which inhibits the growth of CBL and reduce the effects of
577 the heterogeneously heated surface on the TKE. The distribution of the TKE over the
578 heterogeneously heated surface is consistent with the patch-induced circulations
579 described by [Avissar and Schmidt \(1998\)](#). In addition, the surface heat anomaly and
580 background winds have similar effects on the CBL in the cases with a balanced
581 surface heat flux, but the enhanced effects on the TKE are far lower in the cases with
582 an unbalanced surface heat flux. Thus, it is more beneficial to consider the ambient
583 winds. By analyzing the buoyancy and shear production terms in the TKE budget
584 equation and separating the contribution of the resolvable-scale (RES) and
585 subgrid-scale (SGS) eddies, we found that the contributions of the wind shear to the



586 TKE from the SGS eddies are considerable in the CBL (below $0.9z_i$) over a
587 heterogeneously heated surface. The total shear production term is larger below $0.65z_i$
588 in the heterogeneously heated cases with weaker background winds, demonstrating
589 that the patch-induced circulations are conducive to producing more shear in the CBL.
590 We obtained the same conclusion as [Kang and Lenschow \(2014\)](#), that is, the
591 patch-induced circulations become indistinguishable under background flows
592 conditions, and the ambient winds also weaken the convective intensity.

593 Then, we conducted a phase-averaged analysis to separate the contributions of the
594 turbulent intensity and the transport of the total heat flux from those of the
595 patch-induced circulations and the background turbulence field. The patch-induced
596 turbulent intensity increases with increasing lake patches. It mainly contributes to the
597 horizontal turbulent intensity and the potential temperature variance, while it
598 contributes no more than 10% to the vertical turbulent intensity, of which the
599 background turbulence contributes the most. The ambient winds weaken the
600 patch-induced and horizontal turbulent intensities but strengthen the vertical turbulent
601 intensity. The contribution of the patch-induced heat flux was up to 80% in the
602 unbalanced cases and 60% in the balanced cases. The background turbulence made a
603 larger contribution to the heat flux over the area outside of the patches, which have a
604 stronger surface heat flux than that over the lake patches. The background flows also
605 inhibit the transport of the heat flux.

606 To understand the effects of the patch-induced motions on the momentum flux in
607 the surface layer under various background wind conditions, we focused on the



608 variations in the wind speed and the Reynolds stress in the horizontal direction below
609 200 m. Without ambient winds, the inland extent of the patch-induced flows was
610 about 25 km. When the background winds flowed into the lake patches, they
611 decreased by 4.0 m s^{-1} within about 10 km and increased steadily when flowing out of
612 the patches. The synchronized variations in the wind speed and momentum flux in the
613 horizontal direction illustrate that the lake patches alter the spatial distribution of the
614 turbulent stress, which further affects the surface wind speeds, especially over the
615 land-lake boundary regions. We also analyzed the cumulative contribution of eddies
616 with different scales to the buoyancy flux near the surface. It was found that without
617 background flows, the buoyancy flux is transmitted upward by the eddies with larger
618 wavelengths for the case with one lake patch; while there is a negative buoyancy flux
619 in the case with two lake patches. Thus, increasing the number of lake patches leads
620 to more patch-induced motions, which do not tend to enhance the heat transport
621 ability. The background flows promote the opposite results.

622 In the entrainment layer, in contrast to [Matthias et al. \(2014\)](#) who found that the
623 entrainment increased for the stronger heated surface patch cases compared to the
624 homogeneous case, we found that the entrainment flux decrease as the number of lake
625 patches increases. For the unbalanced cases, the convective intensity increases as the
626 number of lake patches increases, but the background flows weaken it. For the
627 balanced cases, the convective intensity weakens as the number of lake patches
628 increases, corresponding to a smaller TKE and total turbulent intensity. In this study,
629 whether the convective rolls persist mainly depended on the background turbulence



630 field with a higher geostrophic wind of 7–11 m s⁻¹, while [Maronga and Raasch \(2013\)](#)
631 reported a higher wind speed of 6 m s⁻¹. As the number of lake patches increases, the
632 increased downdrafts are mainly located over the lake patches, and they carry more
633 warm, dry air down from the free atmosphere in both the balanced and unbalanced
634 cases. The background winds weaken this effect even when there is cooler, moister air
635 in the entrainment layer in the balanced cases.

636 Our study provides ideal simulations of the boundary-layer turbulence over the
637 heterogeneously heated surface in the SRYR. It mainly focused on the influences of
638 the heterogeneous distribution of the surface heat flux and the background winds. In
639 the future, we plan to conduct further research that will take into consideration the
640 topography and additional physical processes to provide a reference for the study of
641 the energy and water exchange processes over the complex surface of the SRYR.

642 **Acknowledgments**

643 This research was supported by the second Tibetan Plateau Scientific Expedition
644 and Research Program (STEP, 2019QZKK0604), the National Natural Science
645 Foundation of China (NSFC) (91837208, 41775013). Thanks to the Zoige Plateau
646 Wetland Ecosystem Research Station, Chinese Academy of Sciences for the field
647 observation data, it supported by the Science and Technology Plan of Gansu Province
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