

Regional impacts of black carbon morphologies on shortwave aerosol-radiation interactions: A comparative study between the US and China

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Abstract. Black carbon (BC) is one of the dominant absorbing aerosol species in the atmosphere. It normally has complex fractal-like structures due to the aggregation process during combustion. A wide range of aerosol-radiation interactions (ARI) of BC has been reported throughout experimental and modeling studies. One reason for the large discrepancies among multiple studies is the application of the over-simplified spherical morphology for BC in ARI estimates. In current climate models, the Mie theory is commonly used to calculate the optical properties of spherical BC aerosols. Here, we employ a regional chemical transport model coupled with a radiative transfer code [¹] that utilizes the non-spherical BC optical simulations to re-evaluate the effects of particles' morphologies on BC [²] shortwave ARI, and the wavelength range of 0.3 – 4.0 μm was considered. Anthropogenic activities and wildfires are two major sources of BC emissions. Therefore, we choose [³] the typical polluted area in eastern China which is dominated by anthropogenic emissions, and [⁴] the fire region in the northwest US [⁵] which is dominated by fire emissions in this study. [⁶] A one-month simulation in eastern China and a seven-days simulation in the fire region in the northwest US [⁷] were performed. The fractal BC model generally presents a larger clear-sky ARI compared to the spherical BC model. Assuming BC particles are externally mixed with other aerosols, the relative differences in the time-averaged clear-sky ARI between the fractal model with a fractal dimension

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⁶removed: Our modeling results show that spherical BC models overestimate the aerosol optical depth (AOD) at 550 nm wavelength up to 0.03 and 0.15 at typical polluted cities in East China and fire regions

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(D_f) of 1.8 [..⁸]and the spherical model are 12.1% – 20.6% and 10.5% – 14.9% for typical polluted urban cities in China and fire sites in the northwest US, respectively[..⁹][..¹⁰][..¹¹][..¹²][..¹³]. Furthermore, the regional-mean clear-sky ARI is also significantly affected by the BC morphology, and relative differences of 17.1% and 38.7% between the fractal model with a D_f of [..¹⁴]1.8 and the [..¹⁵]spherical model were observed in eastern China and the [..¹⁶]northwest US, respectively. However, the existence of clouds would weaken the BC morphological effects. The time-averaged all-sky ARI relative differences between the fractal model with a D_f of 1.8 and the spherical model are 4.9% – 6.4% and 9.0% – 11.3% in typical urban polluted cities and typical fire sites, respectively. Besides, for the regional-mean all-sky ARI, the relative differences between the fractal model and the spherical model are less than 7.3% and 16.8% in the polluted urban area in China and the fire region in the US, respectively. The results imply that current climate modeling may significantly underestimate the BC ARI uncertainties as the morphological effects on BC ARI are ignored in most climate models.

1 Introduction

Black carbon (BC), as the main absorbing aerosol in the atmosphere, exerts a positive radiative forcing, and lofts smoke plumes (Buseck and Buseck, 2000; Streets et al., 2006; Moosmüller et al., 2009). However, there are still large uncertainties in evaluating the BC radiative effects. An important cause of the discrepancy is BC’s complex morphology. BC [..¹⁷]aerosols are assumed as spheres, and the optical properties are calculated using the Mie theory in most climate and atmospheric chemical transport models, such as Community Earth System Model (CESM) (Danabasoglu et al., 2020), the Model for Interdisciplinary Research on Climate (MIROC-SPRINTARS) (Takemura et al., 2005, 2009), Weather Research and Forecasting coupled to Chemistry (WRF-Chem) (Grell et al., 2005; Fast et al., 2006), and GEOS-Chem. However, many studies have shown that BC particles, especially those nascent ones, have fractal-like structures. The spherical assumption for BC can lead to a large deviation from the field measurement data and non-spherical simulated results (Chakrabarty et al., 2007; Luo et al., 2018c; He et al., 2015; Liu and Mishchenko, 2005; Luo et al., 2018d; Mishchenko et al., 2016a; Luo et al., 2021b).

⁸removed: . Besides, spherical BC model underestimates BC aerosol absorption optical depth (AAOD) at 450 nm up to 0.016 and 0.04 at typical polluted cities in East

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35 Based on the sampled BC images, researchers found that the shape of uncoated BC aggregates can be fitted well by a fractal law with monomer number (N_s), mean monomer radius (R), fractal prefactor (k_0), fractal dimension (D_f) and the radius of gyration (R_g) (Mishchenko et al., 2002; Sorensen, 2011; Luo et al., 2021a):

$$N_s = k_0 \left(\frac{R_g}{R} \right)^{D_f} \quad (1)$$

Previous studies have shown that aggregated particle models are more realistic to reproduce the optical measurement results [..¹⁸] (Kahnert, 2010a; Luo et al., 2019, 2018b). Some studies have used the [..¹⁹] fractal BC models to investigate the radiative properties of BC (Wu et al., 2015; Liu et al., 2015a; Liu and Mishchenko, 2005; Yin and Liu, 2010; Teng et al., 2019; Luo et al., 2018a; Kahnert, 2010a). [..²⁰] The direct radiative effect (DRE) was widely used to evaluate the climate effects of aerosol (Bond et al., 2013; Saleh et al., 2015). IPCC (2014) suggested using the new terminology of the aerosol-radiation interactions (ARI) instead of DRE. Thus, in this work, we use the terminology of ARI to replace DRE. However, an extremely limited number of studies have evaluated the ARI of non-spherical BC in regional or global climate models. Kahnert (2010b) has made efforts to simulate the radiative properties of freshly emitted BC using the Multiple-scale Atmospheric Transport and CHemistry (MATCH) model. That study assumed a fixed solar zenith angle (SZA) and restricted the modeling region in Western Europe. [..²¹] Expanding the modeling range to regions with different emission characteristics is important to understand the effects of BC sources on [..²²] ARI.

[..²³] A global mean BC all-sky ARI of +[..²⁴] 0.6 W m⁻² has been reported by [..²⁵] IPCC (2014). However, the BC ARI values estimated based on in-situ optical measurements in some regions are much larger than the rest. BC emissions in China roughly account for one-fourth of its global anthropogenic emission budget (Streets et al., 2001). [..²⁶] eastern China, a typical polluted region, is dominated by anthropogenic emissions (Zhang et al., 2009; Li et al., 2017a). Therefore, BC ARI in [..²⁷] eastern China has gained increasing interest.

Besides the anthropogenic sources, wildfires significantly contribute to the regional BC emissions. The extremely-high BC [..²⁸] concentrations can be found in those fire sites. Recent studies have shown that BC emitted from fire sites can also loft the surrounding atmospheric aerosols to the upper troposphere and lower stratosphere (Yu et al., 2019). Thus, the investigation of BC ARI in these regions is important to understand the detailed plume dynamics and the warming effects of BC. [..²⁹]

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]Northwest US, as one of the [..³⁰]regions where the wildfire is most frequent, has also been investigated in addition to [..³¹]eastern China.

60 In this work, we employed [..³²]WRF-Chem [..³³]to simulate the aerosol mass concentrations. Note here that WRF-Chem assumes aerosols to be spherical. Therefore, the radiative parameters of [..³⁴]fractal BC aggregates were calculated offline using an optical module, Flexible Aerosol Optical Depth (FlexAOD) (Curci, 2012). We calculated the ARI at the top of the atmosphere (TOA) using [..³⁵]a radiative transfer model[..³⁶], libRadtran (Mayer and Kylling, 2005), after the particles' optical properties were calculated. [..³⁷][..³⁸][..³⁹][..⁴⁰][..⁴¹]

65 2 Method

2.1 Aerosol distribution simulation

In this work, WRF-Chem version 4.1.3 was used to simulate the transport of atmospheric species. Two areas were selected to represent the BC sources with different emission characteristics. [..⁴²]Eastern China, a [..⁴³]major polluted region in the world, [..⁴⁴]represents the typical polluted urban region. It consists of 115 east-west grids and 105 south-north grids centered
70 at 112.00°E, 35.00°N with a grid resolution of 18 km. [..⁴⁵]The northwest US, one of the [..⁴⁶]regions where the wildfire is most frequent in the world, was also investigated in this work. The fire region consists of 120 east-west grids and 120 south-north grids centered at 121.48°W, 39.89°N with a grid resolution of 4 km. The schematics of the two [..⁴⁷]regions are shown in Figure S1. Both regions have 33 vertical layers above the ground, with a top pressure of 50 hPa.

We used the Model of Emissions of Gases and Aerosols from Nature [..⁴⁸]version 2.1 (MEGAN2.1) to compute the biogenic
75 emissions over two regions (Guenther et al., 1994, 2006). The anthropogenic [..⁴⁹]inventory for eastern China in the year

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2016 was compiled by the ^[..⁵⁰]Multi-resolution Emission Inventory for China (MEIC) (Li et al., 2014; Liu et al., 2015b)^[..⁵¹]. We used the MIX anthropogenic inventory for the region outside China (Li et al., 2017b). The Regional Acid Deposition Model version 2 (RADM2) atmospheric chemical mechanism (Stockwell et al.) and the Model Aerosol Dynamics for Europe with the Secondary ^[..⁵²]ORGanic Aerosol Model (MADE/SORGAM) were applied in ^[..⁵³]the simulation of eastern China
80 (Seinfeld et al., 2001; Ackermann et al., 1998). Fast-J photolysis scheme (Wild et al., 2000) was used to simulate the photolysis rates. The physical scheme options in WRF-Chem are shown in Table S1. The simulations in eastern China started at 0:00 UTC on November 31st, 2016, and ended at 00:00 UTC on January 1st, 2017. The data from 0:00 UTC on December 1st, 2016 to 00:00 UTC on January 1st, 2017 was used for analysis.

The Emission Database for Global Atmospheric Research for Hemispheric Transport of Air Pollution (EDGAR-HTAP^[..⁵⁴] version 2 emission inventory for 2010 was used in the ^[..⁵⁵]northwest US. MOSAIC aerosol scheme (Zaveri and Peters, 1999; Zaveri et al., 2008) and ^[..⁵⁶]the Carbon Bond Mechanism Z (CBM-Z) photochemical mechanism (Zaveri and Peters, 1999) ^[..⁵⁷]were used in the fire region simulation. The Fire emission was provided by the Fire INventory from NCAR (FINN) (Wiedinmyer et al., 2011). Note here that EDGAR-HTAP anthropogenic inventory and FINN were provided for the MOZART chemical mechanism, so we ^[..⁵⁸]mapped the emission for the ^[..⁵⁹]Model for Ozone and Related chemical
90 Tracers (MOZART) chemical mechanism to the CBM-Z chemical mechanism based on the study of Emmons et al. (2010). ^[..⁶⁰]For both simulations in eastern China and the northwest US, the National Center for Environmental Prediction (NCEP) Global Forecast System ^[..⁶¹](GFS) Final Analysis (FNL) with a horizontal grid spacing of 0.25° and 6 h intervals was used to provide the meteorological initial and boundary conditions. The chemical initial and boundary conditions were obtained from the Model for Ozone and Related ^[..⁶²]chemical Tracers, version 4 (MOZART-4). The simulations in the northwest
95 US started at 0:00 UTC on August 5th, 2016, and ended at 00:00 UTC on August 21st, and the data from 0:00 UTC on August 14th to 00:00 UTC on August 21st was used for analysis.

2.2 The morphology of BC

In this work, we only consider externally mixed BC aerosols, which ^[..⁶³]are commonly represented by fractal structures. D_f is a key parameter to describe the compactness of fractal BC (Wang et al., 2017; Yuan et al., 2019)^[..⁶⁴]. D_f increases from

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100 approximately 1.8 to 3 when the BC morphology can vary from a chain-like structure to a spherical structure^[..⁶⁵]. The freshly emitted BC generally exhibits a fluffy structure with a D_f of approximately 1.8 (Heinson et al., 2010, 2017). The laboratory measurements also showed that the freshly emitted BC generally presents a small D_f . Chakrabarty et al. (2006) have shown that D_f of BC emitted from wildland fuels generally exhibits a range of 1.67 – 1.83. A D_f range of 1.6 – 1.9 was observed for BC produced from diesel combustion (Wentzel et al., 2003). China et al. (2013b) indicated that the BC freshly emitted from
 105 wildfire generally exhibits ^[..⁶⁶] D_f with a range of 1.74 – 1.92. However, a more compact structure was commonly observed for BC in the atmosphere with the particle aging (Li et al., 2003; Adachi et al., 2014; Chen et al., 2016; Adachi et al., 2010). A D_f range of 2.2 – 2.4 was observed in the study of Adachi et al. (2010). The fractal structures with a larger D_f are widely used to describe the BC with more compact structures (Adachi et al., 2010). Chakrabarty et al. (2006) further showed that the D_f of aged BC can reach up to 2.6. To represent both fluffy and compact BC, D_f s of 1.8, 2.2, and 2.6 were considered ^[..⁶⁷]
 110]in this study. While the fractal prefactor k_0 was also measured in a wide range in the atmosphere, its impact on the optical properties was relatively small^[..⁶⁸]. We assumed a fixed k_0 of 1.2 in this work. The typical morphologies of fractal BC are shown in Figure 1.

The volume-mean particle radius was commonly used to describe the size of non-spherical BC. ^[..⁶⁹] Previous studies have observed a range of approximately 8 – 57 nm for BC monomer radius (Eggersdorfer and Pratsinis, 2012; Mikhailov et al.,
 115 2006; KÖylü and Faeth, 1992; Lee et al., 2002), while Kahnert and Kanngießer (2020) further showed that the typical range is approximately 10 – 25 nm. In this work, we assumed a fixed monomer radius of 20 nm. We considered an N_s range of 1 – 1000 to represent BC with different sizes. The volume-mean particle radius (r_p) can be calculated using:

$$r_p = RN_s^{1/3} \quad (2)$$

^[..⁷⁰] Note here that BC can also internally mix with other compositions, and the morphology can become more complex
 120 (Wang et al., 2021c, 2017). However, we mainly focus on the freshly emitted BC, and only consider externally mixed BC. Further investigations would be performed for more complex internally mixed BC in the future.

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2.3 The refractive index and size distribution of BC

BC refractive index shows [..⁷¹] spectral dependence (Chang and Charalampopoulos, 1990), while it doesn't vary largely with wavelengths in the short wavelength range [..⁷²] [..⁷³] [..⁷⁴] [..⁷⁵] [..⁷⁶] (Liu et al., 2018; Lack and Cappa, 2010; Bond and Bergstrom, 2006). The suggested BC refractive index values by Bond and Bergstrom (2006) were commonly used. In this work, the median value of $1.85 + 0.71i$ was used, as it was widely used in many regional and global climate models (e.g. WRF-Chem). The size distribution of BC [..⁷⁷] also suffers large uncertainties from different fuels and conditions. The size distribution of BC is commonly fitted by a lognormal size distribution with a geometric mean radius (r_g), and a geometric standard deviation (σ_g) (Schwarz et al., 2008; Mishchenko et al., 2016b):

$$130 \quad n(r_p) = \frac{N_0}{\sqrt{2\pi}r_p \ln(\sigma_g)} \exp \left[- \left(\frac{\ln(r_p) - \ln(r_g)}{\sqrt{2\ln(\sigma_g)}} \right)^2 \right] \quad (3)$$

where $n(r_p)$ is the probability density distribution of particle number; r_p is the [..⁷⁸] volume-mean particle radius, N_0 is the [..⁷⁹] number concentration, which can be calculated by the mass concentration obtained from WRF-Chem by assuming BC mass density, r_g , and σ_g . The details about the calculation of N_0 are shown in Curci (2012). [..⁸⁰]

BC geometric mean radius of $0.05 - 0.06 \mu\text{m}$ is frequently observed by instruments and widely assumed in numerical studies (Alexander et al., 2008; Coz and Leck, 2011; Reddington et al., 2013; Liu et al., 2018). In this work, BC geometric mean radius was assumed to be $0.05 \mu\text{m}$ [..⁸¹], and σ_g was assumed to be [..⁸²] 1.6. We used the volume-equivalent radius of [..⁸³] to characterize the [..⁸⁴] particle size of factual BC. The density of BC was assumed to be 1.8 gm^{-3} based on the suggested values by Bond and Bergstrom (2006).

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2.4 BC radiative properties

140 In this work, BC radiative properties were calculated using the [\[.85\] Multiple Sphere T-matrix](#) [\[.86\] Method](#) (MSTM) (Mackowski and Mishchenko, 2011, 1996). [\[.87\]](#) The MSTM can efficiently calculate the optical properties of spheres without intersecting surfaces. The MSTM has high computational efficiency because it theoretically calculates the optical properties of randomly oriented particles without numerically averaging them over different particle orientations. The MSTM can directly calculate the extinction efficiency (Q_{ext}), scattering efficiency (Q_{sca}), and phase function (P) [\[.88\]](#) with the refractive index, wavelength, and the input shapefile. Then, the extinction cross-section (C_{ext}) and scattering cross-section (C_{sca}) were obtained using:

$$C_{ext} = Q_{ext} \pi r_p^2 \quad (4)$$

$$C_{sca} = Q_{sca} \pi r_p^2 \quad (5)$$

[\[.89\]](#) Extinction coefficient (b_{ext}), Scattering coefficient (b_{sca}) and bulk phase function $\langle P \rangle$ were calculated using the following equations:

$$b_{ext} = \int_{r_{min}}^{r_{max}} C_{ext}(r_p) n(r_p) dr \quad (6)$$

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$$b_{sca} = \int_{r_{min}}^{r_{max}} C_{sca}(r_p) n(r_p) dr \quad (7)$$

$$\langle P(\theta) \rangle = \frac{\int_{r_{min}}^{r_{max}} C_{sca}(r_p) P(\theta, r_p) n(r_p) dr}{b_{sca}} \quad (8)$$

In climate modeling, instead of using the phase function, the Legendre expansion coefficients were commonly used:

$$\langle P(\theta) \rangle = \sum_{S=1}^{S_{max}} \alpha_S P_S(\cos\theta) \quad (9)$$

155 where θ is the scattering angle, P_S are generalized spherical functions, α_S are the Legendre expansion coefficients, and S_{max} is the order of truncation of the [\[.90\] Legendre expansion coefficients](#).

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⁸⁹removed: where r_p represents the volume-equivalent radius of non-spherical BC. In this work, we have calculated the optical properties of BC with a N_s range of 1 – 1000. Bulk extinction cross-section ($\langle C_{ext} \rangle$), scattering cross-section ($\langle C_{sca} \rangle$) and

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In this work, we used the pmom ^[..⁹¹] tool which is available in libRadtran software for calculating the Legendre expansion coefficients. With the inputs of the aerosol bulk phase function and the desired number of Legendre expansion coefficients, the pmom tool can calculate the Legendre expansion coefficients. ^[..⁹²]

160 The radiative properties of fractal BC and BC spheres were calculated at 300 nm – 4000 nm wavelengths. The step size of $\Delta\lambda = 50$ nm was chosen when λ is less than 1000 nm, while $\Delta\lambda = 200$ nm was selected for $1000 \text{ nm} \leq \lambda \leq 2000 \text{ nm}$, and $\Delta\lambda = 400$ nm when $\lambda \geq 2000$ nm. We created look-up tables for ^[..⁹³] b_{ext} , b_{sca} , and the Legendre expansion coefficients of phase functions for each σ_g and r_g . Thus, the optical properties of BC can be obtained by interpolating the look-up tables.

2.5 Flexible Aerosol Optical Depth (FlexAOD)

165 The aerosol mass concentrations from WRF-Chem were inputted into an optical software ^[..⁹⁴] Flexible Aerosol Optical Depth ^[..⁹⁵] (FlexAOD) (Curci, 2012; Curci et al., 2015) to calculate the aerosol radiative properties. FlexAOD is an optical post-processing tool for the atmospheric chemistry-transport model, and it started as a tool for the ^[..⁹⁶] GEOS-Chem model. We have made some modifications ^[..⁹⁷] to FlexAOD to make it accommodate the WRF-Chem outputs. We ^[..⁹⁸] first mapped the aerosols from WRF-Chem into 5 categories: BC, organic carbon(OC), inorganic salt (INS), sea salt (SA), and dust (DST).
170 The mapping details are shown in Tables S2 – S3. After the WRF-Chem species were mapped, the size distribution, refractive indices, and hygroscopic growth factors were then assigned.

^[..⁹⁹] FlexAOD firstly reads the aerosol mass concentrations from WRF-Chem, then ^[..¹⁰⁰] converts them to aerosol volume concentrations based on the assigned mass densities. Based on the assigned normalized size distributions, we can calculate the number concentration (N_0) of each aerosol. FlexAOD pre-calculates the optical properties of each type of aerosol by assuming $N_0 = 1$ with the assumed size distributions. The total scattering/extinction coefficients can be obtained by multiplying the pre-calculated scattering/extinction cross-sections with the number concentrations. The ^[..¹⁰¹] phase function of each type of aerosol is identical to the pre-calculated phase function by assuming $N_0 =$ ^[..¹⁰²] 1, then the total phase function was calculated according to the number concentration of each aerosol. In FlexAOD, aerosol shapes were assumed to be spherical and the corresponding optical properties of each aerosol species were calculated using the Mie code provided by Mishchenko et al. (1999)^[..¹⁰³]. The bulk optical properties were then calculated by combining an assembly of aerosols.

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[..¹⁰⁴] BC optical properties were overwritten using the look-up tables created in Section [..¹⁰⁵] 2.4. As described in Section [..¹⁰⁶] 2.4, we have created look-up tables for non-spherical BC [..¹⁰⁷] with different size distributions. Thus, if the size distribution of BC is assigned, the optical properties of BC with a normalized size distribution can be determined by interpolating the look-up tables. Once the number concentration is calculated, we can determine the total optical properties.

185 Apart from BC, the physical properties of other chemical species were also specified. We used the OC refractive indices suggested by Highwood (2009). The density of OC varies [..¹⁰⁸] under different conditions. The density of the oxidized organic aerosol was reported to be approximately 1.3 g cm^{-3} (Cross et al., 2007), while Nakao et al. (2013) reported that the density of OC with lower oxidation was approximately $1 - 1.2 \text{ g cm}^{-3}$. For freshly formed OC, $0.9 - 1.1 \text{ g cm}^{-3}$ was used by Liu et al. (2017b). In this work, the density of OC was assumed to be 1.2 g m^{-3} . OC size is also commonly fitted by a lognormal size
 190 distribution. In the study of He et al. (2016) and Dentener et al. (2006), $r_g = 0.03 \mu\text{m}$ and $r_g = 0.075 \mu\text{m}$ were assumed for hydrophobic and hydrophilic OC, respectively. In this study, all OC was assumed to be hydrophilic, and we assumed a r_g of $0.075 \mu\text{m}$ and a σ_g of 1.6 for OC. The refractive indices of dust were identical to those used in [..¹⁰⁹] The Goddard Chemistry Aerosol Radiation and Transport (GOCART) Model (Chin et al., 2002). For dust, the gamma distribution was assumed: (Martin et al., 2003; Curci, 2012):

$$195 \quad n(r_p) = N_0 r_p^{(1-3b)/b} \exp\left(-\frac{r_p}{ab}\right) \quad (10)$$

where a and b are two parameters for the distribution, and b is in the range of $0 - 0.5$.

The refractive indices of other chemical species were adapted from the [..¹¹⁰] Optical Properties of Aerosols and Clouds (OPAC) package (Hess et al., 1998). The physical properties are displayed in Table S4. Similar to the study of Curci et al. (2019), the hygroscopic growth factors of different aerosols were taken from the OPAC package (Hess et al., 1998). [..¹¹¹]

200 [..¹¹²] Note here that many internally mixed particles exist in the atmosphere, while in this study we mainly aim to study the morphological effects of freshly emitted particles, and more complex particles may be investigated in the future. Effective refractive indices were calculated using the volume mixing method for hydrophilic particles.

The total column [..¹¹³] single-scattering albedo (SSA) and Aerosol Optical Depth (AOD) were calculated by FlexAOD, and [..¹¹⁴] Absorption Aerosol Optical Depth (AAOD) was calculated by:

$$205 \quad \text{AAOD} = \text{AOD}(1 - \text{SSA}) \quad (11)$$

¹⁰⁴removed: FlexAOD assumed that the BC shape was spherical. To consider the non-spherical structures, BC

¹⁰⁵removed: 3.3

¹⁰⁶removed: 3.3

¹⁰⁷removed: for

¹⁰⁸removed: from

¹⁰⁹removed: the GOCARTModel (Chin et al., 2002), and the

¹¹⁰removed: OPAC

¹¹¹removed: For dust, the gamma distribution was assumed: (Martin et al., 2003; Curci, 2012):

¹¹²removed: We must clarify

¹¹³removed: SSA and AOD

¹¹⁴removed: AAOD

2.6 ARI modeling

The optical properties (extinction coefficient, SSA, asymmetric factor (ASY)) calculated using FlexAOD at each WRF-Chem grid were inputted into a radiative transfer model, libRadtran (Mayer and Kylling, 2005), to calculate the radiative fluxes at the top-of-the-atmosphere (TOA). The radiative transfer equation was solved by DIScrete Ordinate Radiative Transfer (DISORT) radiative transfer equation solver (Stamnes et al., 1988; Buras et al., 2011). The libRadtran can select a standard atmosphere background and determine the solar zenith angle (SZA) based on the longitude, latitude, and UTC time. The surface albedo information was obtained from NASA Earth Observations (NEO). The radiative transfer calculations were performed for each hour and then were averaged over one day. In this work, ARI of BC aerosol were calculated using following equations:

$$\text{ARI} = \text{FLUX}_{\text{With BC}} - \text{FLUX}_{\text{Without BC}} \quad (12)$$

$$\text{FLUX} = F^{\downarrow} - F^{\uparrow} \quad (13)$$

where F^{\downarrow} represents downward radiative flux and F^{\uparrow} represents upward radiative flux. In this work, we just considered the ARI at the TOA. Only shortwave ARI was considered, and the wavelength is in the range of 0.3 – 4.0 μm .

In this work, both the clear-sky ARI and all-sky ARI were calculated. The daily-mean cloud optical thickness, cloud effective radius, and cloud cover data from the MODerate Resolution Imaging Spectroradiometer (MODIS) products were used for the all-sky ARI calculations. The regional-mean BC ARI was also calculated. The aerosol optical properties and cloud properties were firstly averaged, then the radiative transfer calculations were performed for only one day (the median day), and similar methods were also used in previous studies (eg. Saleh et al. (2015); Tuccella et al. (2020a)).

3 Results

3.1 Impacts of BC morphology on AOD and AAOD

To verify the modeling performance of the aerosol concentrations, we compared the simulated $\text{PM}_{2.5}$ concentrations with observations at some monitoring sites, and the results are shown in Figure 2. In the figures, the left column represents

¹¹⁵removed: Extinction

¹¹⁶removed: Asymmetric

¹¹⁷removed: DISORT

¹¹⁸removed: Based on the longitude, latitude, and UTC time,

¹¹⁹removed: .. The

¹²⁰removed: EARTH OBSERVATIONS

¹²¹removed: Besides, we filled the missing values with the albedo provided by WRF-Chem. In this work, we just considered the ARI for clear-sky.

¹²²removed: the direct radiative effects (ARI)

¹²³removed: for

¹²⁴removed: $\text{PM}_{2.5}$

the typical cities in eastern China, and the right column represents the sites in ^[..¹²⁵]the northwest US. The calculated PM_{2.5} concentrations in eastern China are generally consistent with the observations. Even though the simulated ^[..¹²⁶]PM_{2.5} concentrations in the fire region are a little higher than the observations, the deviations are not large, and the general trends are consistent. Therefore, it is reasonable to represent the atmospheric aerosol concentrations using WRF-Chem modeling. ^[..¹²⁷]
^[..¹²⁸]In this study, we selected three fire sites to evaluate the morphological effects on the BC ARI. The positions of the selected sites are shown in Table S5, and they represent the fire sites with high aerosol concentrations. As shown in Figure 3, the temporal BC concentrations at fire sites can even exceed approximately 400 μg/m³ ^[..¹²⁹]when the fire occurs, while the BC concentrations are extremely low in other days. As shown in Figure 3, even though the maximum BC concentrations in the polluted urban cities are much smaller than the fire ^[..¹³⁰]sites, the mean BC concentrations ^[..¹³¹]can reach approximately 12 μg/m³. The simulated BC concentrations generally agree with the measurements of Zhang et al. (2012) for the urban region, where BC concentrations were observed to be approximately 4 – 12.7 μg/m³.
^[..¹³²]We further compare the calculated AOD and AAOD ^[..¹³³]with observations from AEROSOL ROBOTIC NETWORK (AERONET). The AERONET data of Beijing is available among all selected cities in Eastern China. As shown in Figure 4, the calculated AAOD and AOD can generally represent the observations.

^[..¹³⁴]
Figure 5 shows the AOD of BC with different morphologies^[..¹³⁵], where BC AOD was calculated by the difference between AOD with BC and that without BC. Zhuang et al. (2019) indicated that the time-averaged BC AOD in Beijing, Hefei, and Taihu ^[..¹³⁶]was approximately 0.05 – 0.06. Our simulations are generally in agreement with their results, and the simulated BC AOD varies from approximately 0.01 to approximately 0.12 in the selected urban cities. ^[..¹³⁷]In Beijing and

¹²⁵removed: North America. As shown in Figure 2, the calculated PM_{2.5} concentrations in

¹²⁶removed: PM_{2.5}

¹²⁷removed: We just compare the calculated AOD and AAOD in Beijing with observations from AEROSOL ROBOTIC NETWORK (AERONET), as the optical observations are limited in our simulation region. As shown in Figure 4, the calculated AAOD and AOD can generally represent the observations.

¹²⁸removed: For the fire region, even though the surrounding BC concentrations are small, BC concentrations at fire sites are extremely high.

¹²⁹removed: , which should have a strong impact on the aerosol radiative effects. In East China, we select Beijing, Shanghai, Tianjin, and Nanjing as the typical sites to represent the polluted urban cities

¹³⁰removed: site

¹³¹removed: are also high, which

¹³²removed: Figure 9 shows the mean BC

¹³³removed: as a percentage of the total AOD and AAOD. In typical China polluted cities, BC AOD accounts for approximately 4.6% – 7% of the total AOD, while BC AOD in fire sites can account for larger than 10% of the total AOD. At 450 nm, in both polluted urban and fire sites, the fractions of BC AAOD are close, which is approximately 30%. This means that the relative proportions of BC and OC in polluted urban sites are close to those of fire sites.

¹³⁴removed: BC morphologies can have significant impacts on the BC AOD and AAOD fractions. As BC morphologies change from from a fractal dimension of 1.8 to a spherical structure, BC AOD fraction can vary in the range of approximately 4.6% – 5.8%, 5.0% – 6.2%, 5.5% – 6.9%, 4.8% – 6.0%, 9.0% – 11.1%, 10.3% – 12.7%, 9.1% – 11.2% in Beijing, Shanghai, Tianjin, Nanjing, Fire Loc1, Fire Loc2, and Fire Loc3, respectively. The variations can be above 25%. For BC AAOD fraction, the values can vary in the range of approximately 28.1% – 31.0%, 25.0% – 27.7%, 30.2% – 33.2%, 26.5% – 29.2% , 25.4% – 28.1%, 27.0% – 30.0%, 25.4% – 28.0% in Beijing, Shanghai, Tianjin, Nanjing, Fire Loc1, Fire Loc2, and Fire Loc3, respectively. The variations are approximately 10%. Therefore, BC morphologies should have significant effects on determining the BC fractions based on the total AOD and AAOD.

¹³⁵removed: . By deriving the BC AOD of three urban cities (

¹³⁶removed:), Zhuang et al. (2019) demonstrated that the mean BC AOD in three urban cities

¹³⁷removed: in

Tianjin, BC AOD can reach approximately 0.12, while in Shanghai and Nanjing, the maximum BC AOD is approximately 0.07 and 0.1, respectively. [..¹³⁸]The maximum BC AOD can reach 0.5 [..¹³⁹]– 0.9 [..¹⁴⁰]in typical fire sites. From Figure 5, we can also see that BC AOD calculated using the [..¹⁴¹]spherical model is relatively higher than those using fractal aggregate models [..¹⁴²], which is consistent with the findings of Liu and Mishchenko (2005). As shown in Figure 6, in the [..¹⁴³]urban area a spherical assumption for BC lead to an [..¹⁴⁴]overestimation of less than 0.03 for AOD compared to the fractal model with a D_f of 1.8, while the [..¹⁴⁵]overestimation can reach approximately 0.15 in fire sites. The [..¹⁴⁶]overestimation accounts for a large proportion of BC AOD, which can exceed 20% of the total BC AOD. [..¹⁴⁷]

Figure 7 shows the calculated BC AAOD using different BC models, where BC AAOD was calculated by the difference between AAOD with BC and that without BC. . Our modeling results show that a more compact structure may lead to a smaller AAOD, and this is consistent with the findings of Liu and Mishchenko (2005) for [..¹⁴⁸]single BC particles. The reason [..¹⁴⁹]is that a more compact structure blocks the light [..¹⁵⁰]from transmitting into the inner part of the particle, and a smaller absorption efficiency was observed (Kahnert and Devasthale, 2011). [..¹⁵¹]Shin et al. (2019a) showed that most BC AAOD [..¹⁵²]at 440 nm in Asia is within the range of 0 – 0.12. Our simulated BC AAOD at 450 nm is generally consistent with their findings. At fire sites, BC AAOD at 450 nm wavelength can reach approximately 0.7. [..¹⁵³]As shown in Figure 8, [..¹⁵⁴]the spherical model underestimates AAOD by approximately 0.016 [..¹⁵⁵]compared to the fractal model with a D_f of 1.8 in typical polluted cities, while the AAOD underestimation using the [..¹⁵⁶]spherical model can reach approximately 0.04 in fire sites. In general, the AAOD underestimation using the [..¹⁵⁷]spherical model is approximately 8% of the total BC AAOD.

3.2 [..¹⁵⁸]

¹³⁸removed: In fire sites, BC AOD is much larger, BC AOD can reach approximately

¹³⁹removed: ,

¹⁴⁰removed: , and 0.6 in Loc1, Loc2, and Loc3, respectively

¹⁴¹removed: sphere

¹⁴²removed: . Besides, with the given size distribution, the more compact structure can lead to larger AOD,

¹⁴³removed: polluted urban area ,

¹⁴⁴removed: underestimation

¹⁴⁵removed: underestimation

¹⁴⁶removed: underestimation

¹⁴⁷removed: Therefore, BC AOD and AAOD suffer a non-negligible uncertainty from BC morphologies.

¹⁴⁸removed: a single BC particle

¹⁴⁹removed: may be

¹⁵⁰removed: into inner BC, so the

¹⁵¹removed: By deriving BC AAOD at 440 nm from AERONET,

¹⁵²removed: in Asia are

¹⁵³removed: Our results also show that the sphere model can underestimate BC AAOD.

¹⁵⁴removed: at typical polluted cities, the sphere model can underestimate

¹⁵⁵removed: , while at fire sites,

¹⁵⁶removed: sphere

¹⁵⁷removed: sphere

¹⁵⁸removed: Impacts of BC morphology on AAE and EAE

[..¹⁵⁹] [..¹⁶⁰] [..¹⁶¹] Figure 9 shows the time-averaged BC AOD at 550 nm and AAOD at 450 nm as a percentage of the total AOD and AAOD, respectively. In typical China polluted cities, [..¹⁶²] BC AOD at 550 nm accounts for approximately 4.6% – [..¹⁶³] 7% of the total AOD, while BC AOD in fire sites can account for larger than 10% of the total AOD. At 450 nm, in both polluted urban and fire sites, the fractions of BC AAOD are close, which is approximately [..¹⁶⁴] 30%. This means that the relative proportions of BC and OC in polluted urban sites are close to those of fire sites.

[..¹⁶⁵] BC morphologies also have significant impacts on the BC AOD and AAOD fractions. As BC morphologies change from a fractal dimension of 1.8 to a spherical structure, BC AOD fraction increases from 4.6% – [..¹⁶⁶] 5.5% to 5.8% – [..¹⁶⁷] 6.9% in typical polluted urban cities and from 9.0% – 10.3% to 11.1% – 12.7% in typical fire sites, respectively. The relative differences between the fractal model and the spherical model ((Fractal – Sphere)/Sphere) can be above 25%. For BC AAOD fraction, the values can vary in the range of [..¹⁶⁸] 25.0% – [..¹⁶⁹] 33.2% in typical polluted urban cities and the range of [..¹⁷⁰] 25.4% – [..¹⁷¹] 30.0% in typical fire sites, respectively. The AAOD fraction relative differences between the fractal model with a D_f of 1.8 and the spherical model are approximately 10%.

3.2 Impacts of BC morphology on SSA

280 SSA, as the ratio of scattering to extinction, is widely used to infer aerosol types. Figure 10 shows the comparison of SSA using different BC models. [..¹⁷²] Shin et al. (2019b) showed that the mean SSA at 440 nm was approximately 0.89 – 0.92

¹⁵⁹removed: The absorption

¹⁶⁰removed: Å

¹⁶¹removed: ngström exponent (AAE) calculated using different models is shown in Figure ?? . Obvious deviations between fractal aggregate models and the sphere model are found for AAE estimation. With more compact structures, smaller AAE can be observed, which agrees with the findings of Li et al. (2016) and Liu et al. (2018). At typical

¹⁶²removed: AAE at 450

¹⁶³removed: 850 nm wavelength pair varies from approximately 1.5 to approximately 2.5, while AAE at fire site can reach as large as approximately 3.0. The observed large AAE is caused by large portions of brown carbon, which significantly absorbs radiation in ultra-violet (UV) and short visible wavelengths. However, as this study mainly focuses on the effects of BC morphologies, the AAE uncertainties caused by OC absorption are not considered in this work. Figure ?? shows that the sphere model can underestimate the AAE by approximately 0.17

¹⁶⁴removed: 17% of BC AAE. As AAE is commonly used to determine the aerosol type, BC morphologies may affect the threshold values for distinguishing the aerosol types.

¹⁶⁵removed: EAE, as the spectral-dependence of extinction, is widely used to determine the size information of aerosols. Smaller aerosols normally have larger EAE. Figure ?? shows the EAE calculated using different BC models. By using AERONET observation data, in Beijing, Shaheen et al. (2019) showed that the mean EAE at 440

¹⁶⁶removed: 870 nm in winter was approximately 1.06 ± 0.36 . The Sun-sky radiometer Observation NETwork (SONET) data conducted by Xie et al. (2015) demonstrated that EAE in the winter of Beijing was commonly within the range of approximately 0.5

¹⁶⁷removed: 1.5. Our calculated results generally agree with the observation from previous studies. In most cases, the calculated EAE is commonly within

¹⁶⁸removed: 0.5

¹⁶⁹removed: 1.5, even though some smaller EAE is also observed. In the four typical urban cities, EAE values at 450 – 850 nm are similar, and they commonly vary from approximately 0.3 to approximately 1.5. However, EAE at fire sites is commonly within

¹⁷⁰removed: approximately 1.5

¹⁷¹removed: 1.85, which is much larger than the values in the polluted urban area. The reason may be that dominating aerosols in the fire region is commonly occupied by carbonaceous aerosols, which are commonly in fine mode. Compared to AAE, the effects of BC morphologies on the total EAE is relatively small.

As shown in Figure ??, the EAE differences between fractal BC and spherical BC are within 0.05

¹⁷²removed: By using AERONET data, Shin et al. (2019b) showed

in Beijing using AERONET data, and Shaheen et al. (2019) demonstrated that the mean SSA at 440 nm was approximately 0.897 in winter [..¹⁷³] Beijing. At 450 nm, the calculated SSA [..¹⁷⁴] in typical urban cities is within the range of [..¹⁷⁵] a 0.86 – 0.92, which [..¹⁷⁶] is generally consistent with the observations in previous studies. While in the fire region, SSA varies in a wider range. At 450 nm, SSA at the selected fire sites can vary in the range of [..¹⁷⁷] 0.75 – 0.9, which is a little smaller than that in polluted urban areas due to large portions of carbonaceous aerosols in the fire region.

With more compact structures, SSA presents a larger value, which is consistent with the findings of Kahnert and Devasthale (2011). However, the effects of BC morphologies on total SSA at 450 nm are not obvious due to the small percentage of BC in the atmosphere. As shown in Figure 11, the overestimations of the spherical BC model for SSA at 450 nm are generally within 0.005 at typical polluted cities and within 0.012 at fire sites, which is less than 1% of the total SSA. However, these values may have relatively larger impacts when evaluating the climate effect of BC, as BC commonly presents a relatively small value of 0.2 – 0.4 (Kahnert and Devasthale, 2011).

3.3 Impacts of BC morphology on ARI

BC [..¹⁷⁸] clear-sky ARI varies in different regions and the reported BC ARI varies in previous studies. [..¹⁷⁹] Zhuang et al. (2018) estimated BC clear-sky ARI to be +1.85 Wm^{-2} in East Asia. Much larger BC clear-sky ARI during December (+15 Wm^{-2}) and November (+8 Wm^{-2}) over Ahmedabad and Gurushikhar, respectively, were reported by Rajesh and Ramachandran (2018). [..¹⁸⁰] Zhuang et al. (2019) showed that clear-sky [..¹⁸¹] ARI averaged over East Asia [..¹⁸²] were +0.02 to +1.34 W/m^2 in summer [..¹⁸³] in eastern Asia. Lu et al. (2020) showed daily-mean BC clear-sky ARI were within the range of +[..¹⁸⁴] 1.37 – [..¹⁸⁵] 4.89 Wm^{-2} in [..¹⁸⁶] Beijing. Our calculated daily-mean BC ARI in winter generally agrees with [..¹⁸⁷] those previous studies. The daily-mean clear-sky BC ARI at the TOA in typical sites using different BC models is presented in Figure 12. In winter, BC clear-sky ARI at typically polluted cities [..¹⁸⁸] varies in the range of approximately

¹⁷³removed: in

¹⁷⁴removed: at

¹⁷⁵removed: approximately

¹⁷⁶removed: are

¹⁷⁷removed: approximately

¹⁷⁸removed: ARI at the TOA using different BC models is presented in Figure 12. BC ARI

¹⁷⁹removed: For example, a global mean BC ARI of +1.1 W m^{-2} has been reported by previous studies (Bond et al., 2013), while much larger BC ARI was observed in some regions. Based on situ measurements, Lamb et al. (2018) have estimated the column BC ARI to be +0.48 to +2.01 Wm^{-2} over South Korea.

¹⁸⁰removed: In eastern China,

¹⁸¹removed: TOA

¹⁸²removed: (100–130°E, 20–50°N)

¹⁸³removed: . While in some sites (such as Beijing), the average BC ARI can reach approximately +4.5 Wm^{-2} . Based on different measured BC profiles,

Lu et al. (2020) showed BC ARI was

¹⁸⁴removed: 3.9

¹⁸⁵removed: 5.1

¹⁸⁶removed: the Beijing region

¹⁸⁷removed: their results.

¹⁸⁸removed: commonly

305 $+ [..^{189}] 0.5 - [..^{190}] 5.0 \text{ Wm}^{-2} [..^{191}]$. While in large fire sites, the $[..^{192}]$ daily mean BC clear-sky ARI exceeds $+ [..^{193}] 8.0 \text{ Wm}^{-2} [..^{194}]$. Generally, with a more compact structure, BC presents a smaller clear-sky ARI at the TOA. The reason can be explained from the following aspects. Fractal BC can more efficiently absorb than spherical BC, while the total scattering is not significantly modified. Thus, the fractal BC leads to a larger positive clear-sky ARI.

Table 1. Time-averaged BC Clear-sky ARI at different sites (Wm^{-2}).

Location	$D_f=1.8$	$D_f=2.2$	$D_f=2.6$	Sphere
Beijing	$+ [..^{195}] 1.76$	$+ [..^{196}] 1.68$	$+ [..^{197}] 1.63$	$+ [..^{198}] 1.57$
Shanghai	$+ [..^{199}] 1.52$	$+ [..^{200}] 1.45$	$+ [..^{201}] 1.38$	$+ [..^{202}] 1.26$
Tianjin	$+ [..^{203}] 2.00$	$+ [..^{204}] 1.91$	$+ [..^{205}] 1.86$	$+ [..^{206}] 1.77$
Nanjing	$+ [..^{207}] 1.99$	$+ [..^{208}] 1.90$	$+ [..^{209}] 1.83$	$+ [..^{210}] 1.70$
Fire Loc1	$+ [..^{211}] 5.39$	$+ [..^{212}] 5.16$	$+ [..^{213}] 5.00$	$+ [..^{214}] 4.69$
Fire Loc2	$+ [..^{215}] 8.60$	$+ [..^{216}] 8.27$	$+ [..^{217}] 8.12$	$+ [..^{218}] 7.78$
Fire Loc3	$+ [..^{219}] 5.61$	$+ [..^{220}] 5.38$	$+ [..^{221}] 5.22$	+4.91

$[..^{222}]$

Table 2. Time-averaged BC All-sky ARI at different sites (Wm^{-2}).

Location	$D_f=1.8$	$D_f=2.2$	$D_f=2.6$	Sphere
Beijing	$+ [..^{223}] 1.77$	$+ [..^{224}] 1.70$	$+ [..^{225}] 1.68$	$+ [..^{226}] 1.67$
Shanghai	+2.26	+2.18	+2.15	+2.15
Tianjin	+2.00	+1.93	+1.90	+1.88
Nanjing	+2.79	+2.69	+2.66	+2.66
Fire Loc1	+5.14	+4.93	+4.81	+4.62
Fire Loc2	+7.38	+7.11	+6.99	+6.77
Fire Loc3	+5.22	+5.02	+4.89	+4.69

Table 1 shows the time-averaged BC clear-sky ARI (averaging BC clear-sky ARI over the simulation period) in typical polluted urban sites and fire sites. Using the spherical model, the time-averaged clear-sky ARI is $+1.26 - +1.77 \text{ Wm}^{-2}$

¹⁸⁹removed: 2

¹⁹⁰removed: 9.0

¹⁹¹removed: , and the mean BC ARI is approximately $+3.5 - 5 \text{ Wm}^{-2}$

¹⁹²removed: mean BC ARI can reach approximately

¹⁹³removed: 7.0

¹⁹⁴removed: , and the maximum BC ARI can exceed $+11 \text{ Wm}^{-2}$.

²²²removed: Generally, with a more compact structure, BC presents smaller ARI at the TOA. Figure ?? shows the difference of BC ARI at the TOA between fractal aggregate models and the sphere model. In the polluted urban area, the spherical assumption can underestimate BC ARI by approximately 1.0 Wm^{-2} , and the underestimation at fire sites can reach approximately 0.6 Wm^{-2} . As shown in Table 1, using the sphere model, mean BC ARI are

in [..²²⁷] typical polluted urban cities, while it increases to $+1.52 - +2.00 \text{ Wm}^{-2}$ using a fractal aggregate model with a D_f of 1.8 [..²²⁸]. The relative [..²²⁹] differences between the fractal model with a D_f of 1.8 and spherical model ((Fractal – Sphere)/Sphere) can reach approximately 12.1% – 20.6% in typical urban cities. In fire sites, [..²³⁰] when modifying BC structure from a sphere to a fractal aggregate with a D_f of 1.8, [..²³¹] the time-averaged BC clear-sky ARI increases from [..²³²] $4.69 - +[..²³³] 7.78 \text{ Wm}^{-2}$, [..²³⁴] to $+ [..²³⁵] 5.39 - + [..²³⁶] 8.60 \text{ Wm}^{-2}$ [..²³⁷], [..²³⁸] and the relative differences between the two models are 10.5% – 14.9%. Lu et al. (2020) showed that BC shapes can introduce approximately 5% relative uncertainties in [..²³⁹] eastern China using different measured BC profiles. However, our results show that much larger uncertainties can be introduced from BC morphologies. The reason is that [..²⁴⁰] $D_f = 2.8$ was assumed for BC aggregates [..²⁴¹] in the study of Lu et al. (2020), which are close to spherical shape. This D_f value is [..²⁴²] larger than the observed D_f . Besides, due to different [..²⁴³] solar zenith angles, our results show the ARI uncertainties caused by BC morphologies may vary in different regions. Therefore, the BC morphological effects on the BC ARI should be carefully considered in different regions.

4 [..²⁴⁴]

[..²⁴⁵] IPCC (2014) has reported a global mean BC all-sky ARI of $+0.6 \text{ W m}^{-2}$, while Wang et al. (2014) estimated a smaller all-sky ARI of $+0.13 \text{ W m}^{-2}$ based on the constraints from the mass and absorption observations. Tuccella et al. (2020b) further showed that the global mean BC all-sky ARI is in the range of $+0.13$ and $+0.25 \text{ W m}^{-2}$. However, the BC all-sky ARI in some specific regions is relatively large. Based on in-situ measurements, Lamb et al. (2018) have estimated the mean column all-sky BC ARI to be $+0.48$ to $+2.01 \text{ Wm}^{-2}$ over South Korea. In Beijing, Sun et al. (2022) found that the mean BC all-sky ARI decreased from $+3.36 \text{ W m}^{-2}$ in [..²⁴⁶] 2012 to $+1.09 \text{ W m}^{-2}$ in 2020. The daily-mean all-sky ARI in typical urban polluted cities in eastern China and fire sites in the northwest US are shown in Figure 13. The daily-mean

²²⁷removed: Beijing, Shanghai, Tianjin, and Nanjing, respectively, while

²²⁸removed: , BC ARI can reach $+3.83, +4.29, +4.92, +4.03 \text{ Wm}^{-2}$, respectively

²²⁹removed: deviations are approximately 10.4%, 15.3%, 10.6%, 13.8% in Beijing, Shanghai, Tianjin, and Nanjing, respectively

²³⁰removed: the BC ARI differences using different models are relatively smaller. When

²³¹removed: mean BC

²³²removed: 5.52 to

²³³removed: 5.90

²³⁴removed: $+6.64$

²³⁵removed: 7.05 Wm^{-2} ,

²³⁶removed: 4.91 to $+5.23$

²³⁷removed: at Loc1, Loc2, and Loc3, respectively, and the relative uncertainties are approximately 6.9%

²³⁸removed: 6.2%, 6.5%, respectively. By using different measured BC profiles, Lu et al. (2020) showed BC shape

²³⁹removed: East China

²⁴⁰removed: in the study of Lu et al. (2020),

²⁴¹removed: , which is

²⁴²removed: maybe a little

²⁴³removed: mean solar

²⁴⁴removed: Summary and Conclusions

²⁴⁵removed: In this study, we numerically investigated the effects of BC morphologies on ARI at four cities

²⁴⁶removed: East China and three locations

all-sky ARI estimated in this study generally varies in the range of approximately $+0.2 - +4.5 \text{ Wm}^{-2}$ in typical polluted
330 cities. Table 2 shows the time-averaged BC all-sky ARI (averaging BC all-sky ARI over the simulation period) in different
sites. The time-averaged all-sky ARI are $+1.67 - +2.79 \text{ Wm}^{-2}$ in urban cities, and it is relatively high compared to the
regional-mean all-sky ARI due to the high BC emissions. Our estimated BC all-sky ARI is generally in the range reported
by Sun et al. (2022).

The differences of all-sky ARI between the fractal model and spherical model are smaller than those of clear-sky ARI
335 in typical polluted cities in eastern China. The relative differences in the time-averaged all-sky ARI between fractal model
and the spherical model are below 6.5% in typical polluted cities. The all-sky ARI at fire sites in the northwest US is
smaller than the clear-sky ARI. The time-averaged all-sky ARI are $+4.62 - +7.38 \text{ Wm}^{-2}$ in typical fire sites. The relative
differences in the time-averaged all-sky ARI between the fractal model with a D_f of 1.8 and the spherical model are 9.0%
– 11.3% in typical fire sites, which is relatively smaller than the differences for clear-sky ARI.

340 The regional-mean ARI are shown in Figures 14 – 15. As shown in Figure 14, the BC clear-sky ARI of exceeding
 $+3.0 \text{ W/m}^2$ is observed in eastern China. With a spherical BC model, the regional-mean clear-sky ARI in eastern China
is estimated as $+1.35 \text{ Wm}^{-2}$, and it agrees well with the reported regional-mean clear-sky ARI of $+1.34 \text{ Wm}^{-2}$ in East
Asia by Zhuang et al. (2019). BC morphologies also have a non-ignorable impact on the BC ARI. The regional-mean ARI
deviations between the fractal model and spherical model are $+0.23$, $+0.15$, and $+0.1 \text{ Wm}^{-2}$ when D_f is 1.8, 2.2, and
345 2.6, respectively, and the relative differences are 17.1%, 11.1%, and 7.4%, respectively.

The BC morphologies have a relatively small impact on the all-sky ARI. As shown in the lower panels of Figure 14, the
all-sky ARI is generally larger than the clear-sky ARI, while the deviations between the fractal model and spherical model
are smaller compared to those for clear-sky. The BC all-sky ARI can exceed $+4.0 \text{ W/m}^2$ in eastern China. A relatively
larger regional-mean all-sky ARI is observed than clear-sky ARI, which is $+1.79 \text{ Wm}^{-2}$ in eastern China when using
350 the spherical model. The regional-mean all-sky ARI differences between the fractal model and spherical model are 0.13,
0.06, 0.03 Wm^{-2} when D_f is 1.8, 2.2, and 2.6, respectively. The relative differences for all-sky ARI between the two
models are less than 7.3%.

The regional-mean clear-sky ARI in the fire region in the northwest US is shown in Figure 15. The clear-sky ARI in
the fire sites is obviously larger than the other sites. The BC clear-sky ARI of exceeding $+5.0 \text{ W/m}^2$ is observed, and the
355 regional-mean clear-sky ARI is $+0.93 \text{ W/m}^2$. The differences of regional-mean clear-sky ARI between the fractal model
and the spherical model in the fire region in the northwest US are more substantial than those in eastern China, which
reach approximately $+0.36$, $+0.29$, and $+0.19 \text{ W/m}^2$ when D_f is 1.8, 2.2, and 2.6, respectively. The relative differences
between the fractal model and spherical model are 38.7%, 31.2%, and 20.4% when D_f is 1.8, 2.2, and 2.6, respectively.
Smaller all-sky ARI than clear-sky ARI are observed at the sites where fire occurs, while the regional-mean all-sky ARI in
360 the fire region in the northwest US is generally larger than clear-sky ARI, which is $+1.67 \text{ W/m}^2$. The differences between
the two models are $+0.28$, $+0.19$, and $+0.11 \text{ W/m}^2$ when D_f is 1.8, 2.2, and 2.6, respectively, and the relative differences
are 16.8%, 11.4%, and 6.6%.

4 Discussion

In current climate models, such as CESM, MIROC-SPRINTARS, and WRF-Chem, the Mie theory was commonly used to calculate the optical properties of BC aerosols. However, fractal-like BC aerosols were often observed in the atmosphere. In this work, we found that the effects of BC morphology are spatially-dependent. Compared to the spherical BC model, the fractal BC model generally presents a larger clear-sky ARI, which may lead to the underestimations of BC ARI in the climate models. The relative differences in the time-averaged clear-sky ARI are 12.1% – 20.6% and 10.5% – 14.9% in typical polluted urban cities and fire sites, respectively. Furthermore, the regional-mean clear-sky ARI is also significantly affected by the BC morphology, and relative differences of 17.1% and 38.7% between the fractal model were observed in eastern China and in the northwest US^[..²⁴⁷], respectively, while the existence of cloud would weaken the BC morphological effects. The results imply that current climate modeling may significantly underestimate the BC ARI uncertainties as the morphological effects on BC ARI are ignored in most climate models.

However, this work is by no means exhaustive. This work assumed that BC aerosols are externally mixed with other chemical components, while BC aerosols are often internally mixed with other components, such as organic aerosols, sulfate, etc (China et al., 2013a; Adachi et al., 2010; Wang et al., 2021b). BC absorption can be significantly enhanced by the "lensing Effect" even if BC aerosols are internally mixed with non-absorbing materials, which may lead to larger BC ARI (Chung et al., 2012; Liu et al., 2017a). Previous studies have shown that the morphologies of internally mixed BC would significantly affect its absorption enhancement (Luo et al., 2019; Wang et al., 2021a; Luo et al., 2021c), so lead to larger uncertainties in the estimation of BC ARI. Thus, the sensitivities of BC morphologies on the ARI estimated in this work may be smaller than those in real cases.

Futhermore, we found that the spherical assumption generally unerestimates the clear-sky ARI for externally mixed BC, while oppsite phenomenon may be found for internally mixed BC. A core-shell spherical morphology was widely used to represent the internally mixed BC. However, many partially coated BC aerosols exist in the atmosphere, while the core-shell spherical BC model commonly assumes the BC is fully embedded in a coating shell. The core-shell morphology may overestimate the absorption of partially coated BC (Wang et al., 2021a; Zhang et al., 2018), so overestimate the ARI. Thus, the ARI of internally mixed BC with complex morphologies should be further investigated in the future.

5 Summary and Conclusions

The current climate modeling commonly assumes a spherical morphology for BC, while the fractal structure is more realistic than the spherical morphology for externally mixed BC. In this work, we used the fractal model to re-evaluate the BC ARI in a typical polluted urban area in eastern China and a fire region in the northwest US. We found that BC morphologies have non-ignorable impacts on the aerosol optical properties. At 550 nm wavelength, the spherical BC

²⁴⁷removed: .

model can overestimate the AOD up to 0.03 and 0.15 [..²⁴⁸] in typical polluted cities in [..²⁴⁹] China and fire [..²⁵⁰] sites in US, respectively. The overestimations roughly account for 20% of the total BC AOD. Besides, the spherical BC model [..²⁵¹]
395]underestimates BC AAOD at 450 nm up to 0.016 and 0.04 at typical polluted cities in [..²⁵²] China and fire sites in the US, respectively, [..²⁵³] compared to the fractal model with a D_f of 1.8. The underestimations account for approximately 8% of the total BC AAOD. [..²⁵⁴]

[..²⁵⁵] Both the morphological effects on clear-sky and all-sky ARI are evaluated. With a spherical BC model, the [..²⁵⁶]
]estimated time-averaged clear-sky ARI is generally in the range of +[..²⁵⁷] 1.26 – +[..²⁵⁸] 1.77 Wm^{-2} in typical urban
400 polluted cities in eastern China, while this range increases to approximately +[..²⁵⁹] 1.52 – +[..²⁶⁰] 2.00 Wm^{-2} when using the fractal model with a D_f of 1.8. The [..²⁶¹] clear-sky ARI relative differences between the two models are approximately 12.1% – 20.6% in typical urban polluted cities. In fire sites, [..²⁶²] when modifying BC structure from a sphere to a fractal aggregate, the time-averaged BC clear-sky ARI increases from +[..²⁶³] 4.69 – +[..²⁶⁴] 7.78 Wm^{-2} [..²⁶⁵] to +[..²⁶⁶] 5.39 – +8.60 Wm^{-2} [..²⁶⁷] in typical fire sites. The relative differences between the two models are approximately 10.5% –
405 14.9% in typical fire sites. The existence of clouds weaken the effects of BC morphologies on the ARI. The all-sky ARI relative differences between the fractal model and the spherical model are approximately 4.9% – 6.4% and 9.0% – 11.3% in typical urban polluted cities and typical fire sites, respectively, which is relatively smaller than those of clear-sky ARI.

The impacts of BC morphologies on the regional-mean ARI were also evaluated. The regional-mean clear-sky ARI was estimated as +[..²⁶⁸] 1.35 Wm^{-2} and +[..²⁶⁹] 0.93 Wm^{-2} [..²⁷⁰] in the polluted urban area and the fire region, respectively,
410 using the spherical BC model. The regional-mean clear-sky ARI differences between the fractal model and the spherical

²⁴⁸removed: at

²⁴⁹removed: East

²⁵⁰removed: regions in the northwest

²⁵¹removed: can underestimate

²⁵²removed: East

²⁵³removed: when modifying BC structure from spheres to fractal aggregates

²⁵⁴removed: Even though BC morphologies have small impacts on total SSA and EAE, their effects on the AAE are rather significant. At 450 – 850 nm wavelength pair, AAE is approximately 1.5 – 2.5 and 2.0 – 3.0 in typical polluted urban area and fire region, respectively. The calculated AAE difference between the spherical BC model and fractal aggregate model can reach approximately 0.17.

²⁵⁵removed: BC morphologies have non-negligible impacts on the BC ARI. Using the

²⁵⁶removed: calculated mean ARI is approximately

²⁵⁷removed: 3.47 Wm^{-2} ,

²⁵⁸removed: 3.72 Wm^{-2} , +4.45 Wm^{-2} , +3.54 Wm^{-2} in Beijing, Shanghai, Tianjin, and Nanjing, respectively, while the values increase

²⁵⁹removed: 3.83 Wm^{-2} ,

²⁶⁰removed: 4.29 Wm^{-2} , +4.92 Wm^{-2} , and +4.03 Wm^{-2} , respectively, when using a fractal aggregate

²⁶¹removed: relative variations are 10.4%, 15.3%, 10.6%, and 13.8%, respectively. At

²⁶²removed: the uncertainties are smaller but still non-negligible. When modifying BC from spheres to fractal aggregates with a D_f of 1.8, the mean BC

²⁶³removed: 5.52 to

²⁶⁴removed: 5.90

²⁶⁵removed: , +6.64

²⁶⁶removed: 7.05

²⁶⁷removed: ,

²⁶⁸removed: 4.91 to

²⁶⁹removed: 5.23

²⁷⁰removed: at Loc1, Loc2,

model are approximately +0.23 and $[..^{271}] + 0.36 \text{ Wm}^{-2}$ in these two regions, respectively, and the relative $[..^{272}]$ differences between the two models are approximately 17.1% and 38.7%, respectively. The all-sky ARI differences between the fractal model and the spherical model are relatively smaller. The relative differences in the regional-mean all-sky ARI between the fractal model and the spherical model are less than 7.3% and 16.8% in these two regions, $[..^{273}]$ respectively. $[..^{274}]$

415]Thus, the effects of BC morphologies on the $[..^{275}]$ ARI should be carefully $[..^{276}]$ considered in different regions.

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420

²⁷¹removed: Loc3

²⁷²removed: variations are approximately 6.9%

²⁷³removed: 6.2%, 6.5% ,

²⁷⁴removed: Therefore

²⁷⁵removed: regional radiative effects

²⁷⁶removed: evaluated

²⁷⁷removed: The look-up tables calculated in this work can be obtained from https://figshare.com/articles/dataset/Look_up_tables_zip/13096241. The PM2.5 data in China was obtained from <https://www.aqistudy.cn/historydata/>, and PM2.5 data in North America can be found from <https://www.epa.gov/outdoor-air-quality-data/download-daily-data>. FlexAOD can be obtained on request (<http://pumpkin.aquila.infn.it/flexaod/>)

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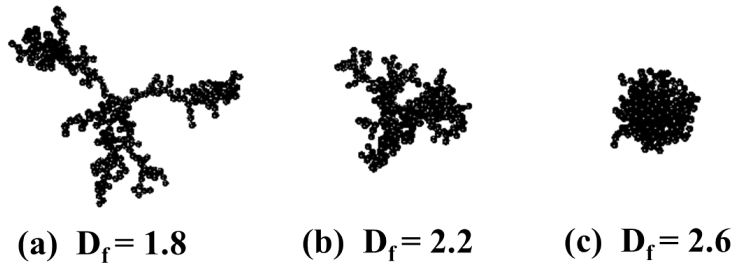


Figure 1. Typical morphologies of fractal BC.

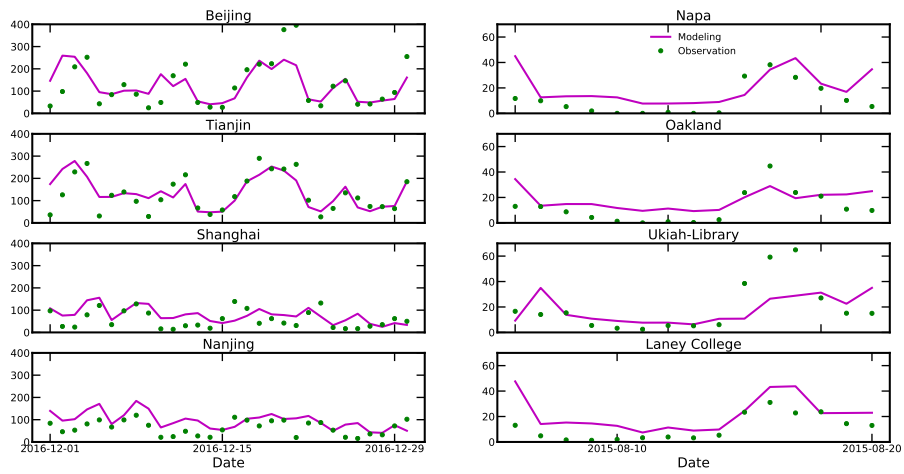


Figure 2. Comparison of measured and calculated PM_{2.5} concentrations.

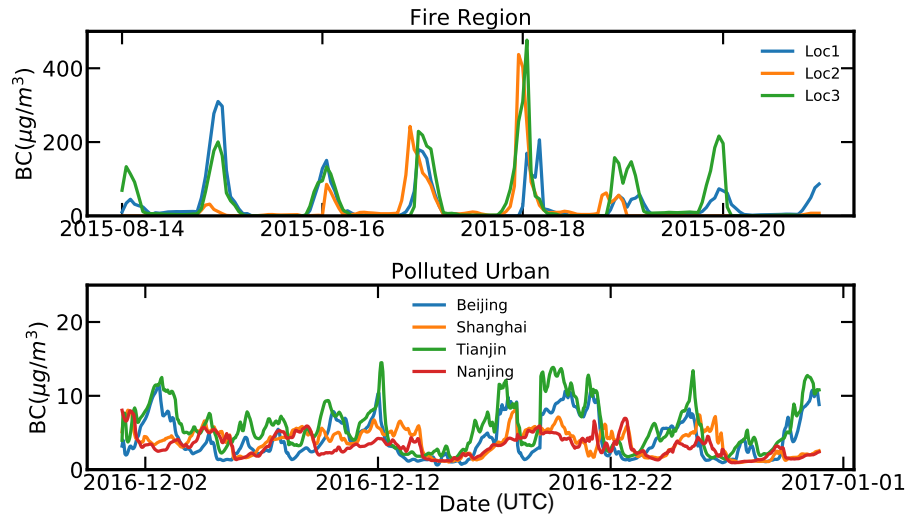


Figure 3. The time series of BC concentrations in typical northwest US and eastern China.

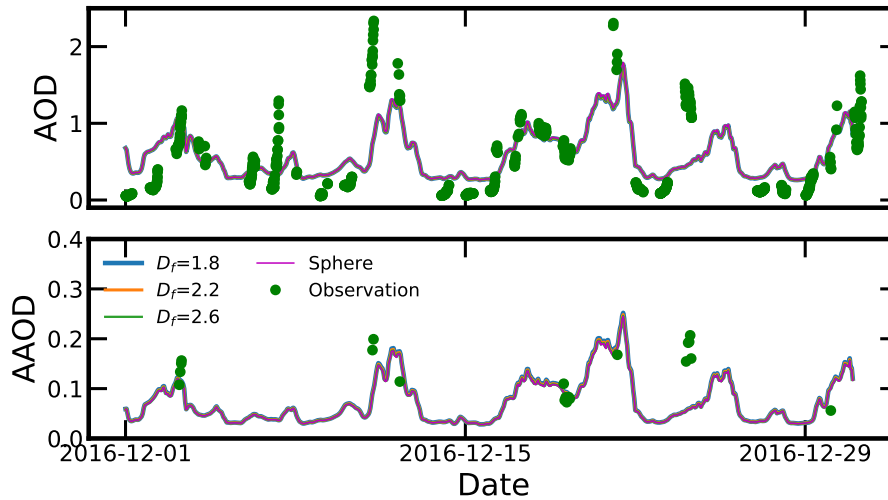


Figure 4. Comparison of measured and calculated AOD and AAOD in Beijing, where $\lambda=500$ nm for AOD. $\lambda=440$ nm and 450 nm for measured and calculated AAOD, respectively, and the observations were obtained from AERONET.

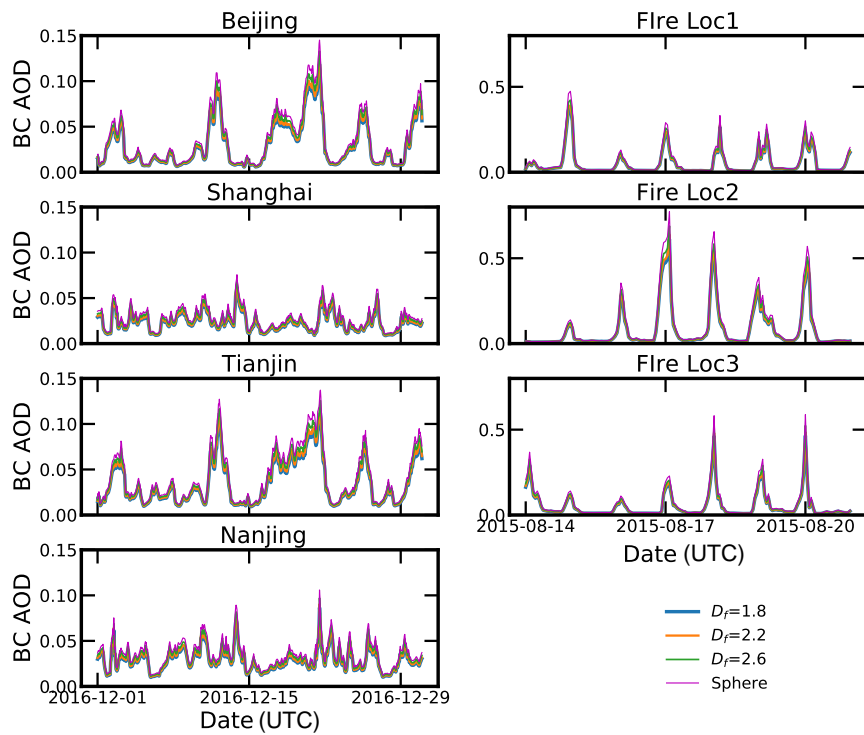


Figure 5. The comparison of BC AOD for different BC morphologies, $\lambda=550$ nm.

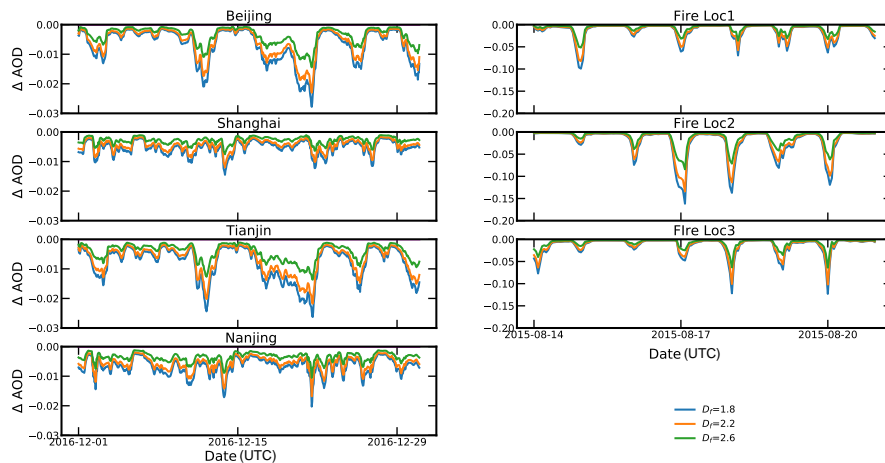


Figure 6. The AOD difference between fractal aggregate models and the spherical model, $\lambda=550$ nm.

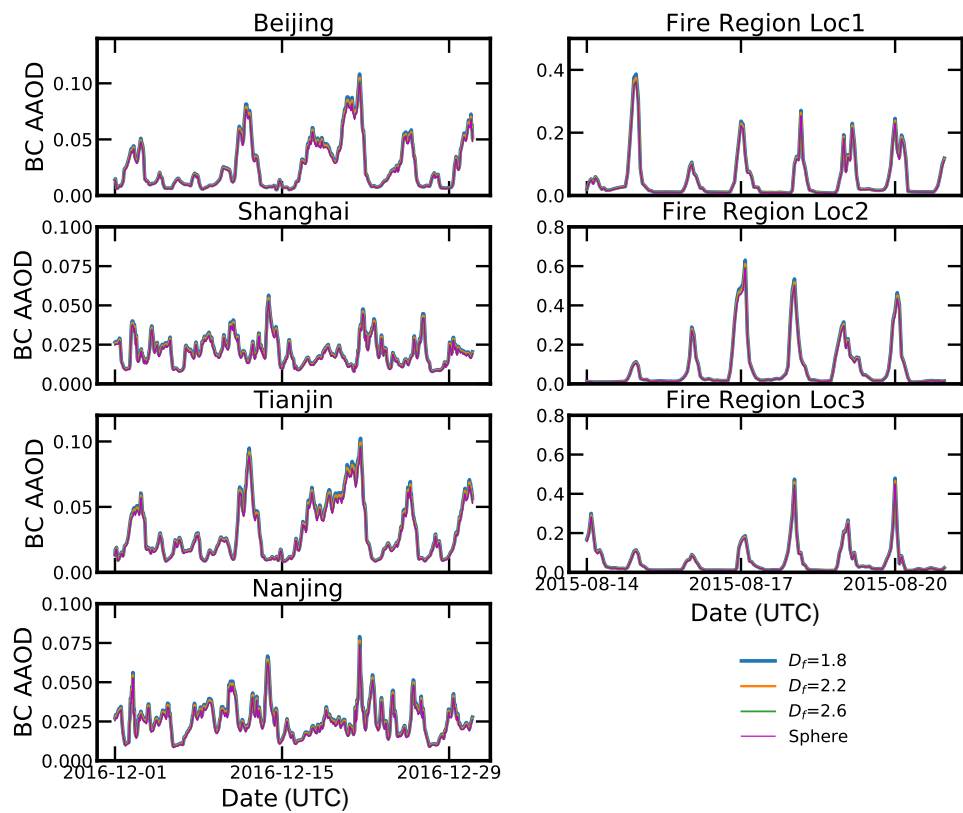


Figure 7. The comparison of BC AAOD for different BC morphologies, $\lambda=450$ nm.

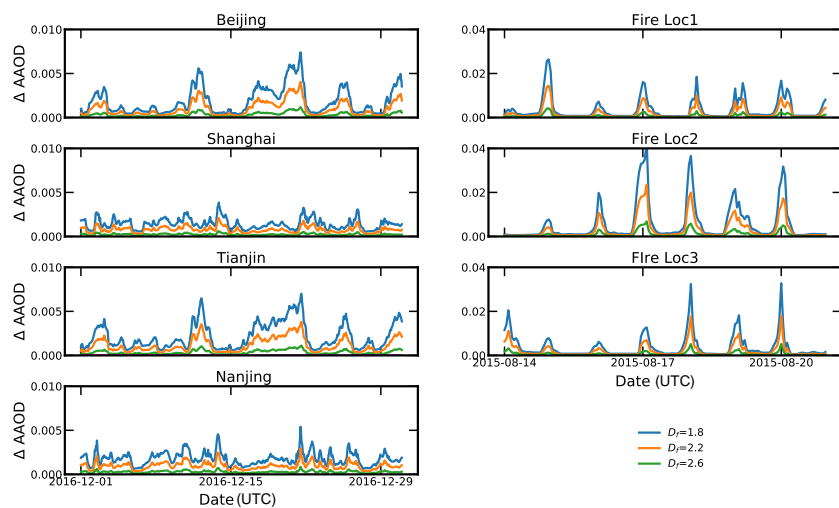


Figure 8. The AAOD difference between fractal aggregate models and the spherical model, $\lambda = 450$ nm.

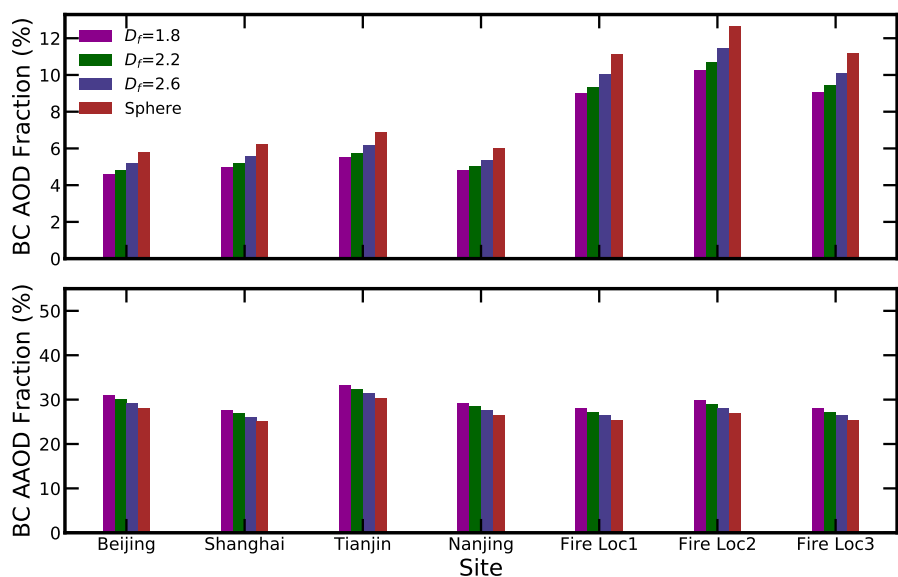


Figure 9. BC AOD (550 nm) and AAOD(450 nm) as a percentage of total AOD and AAOD.

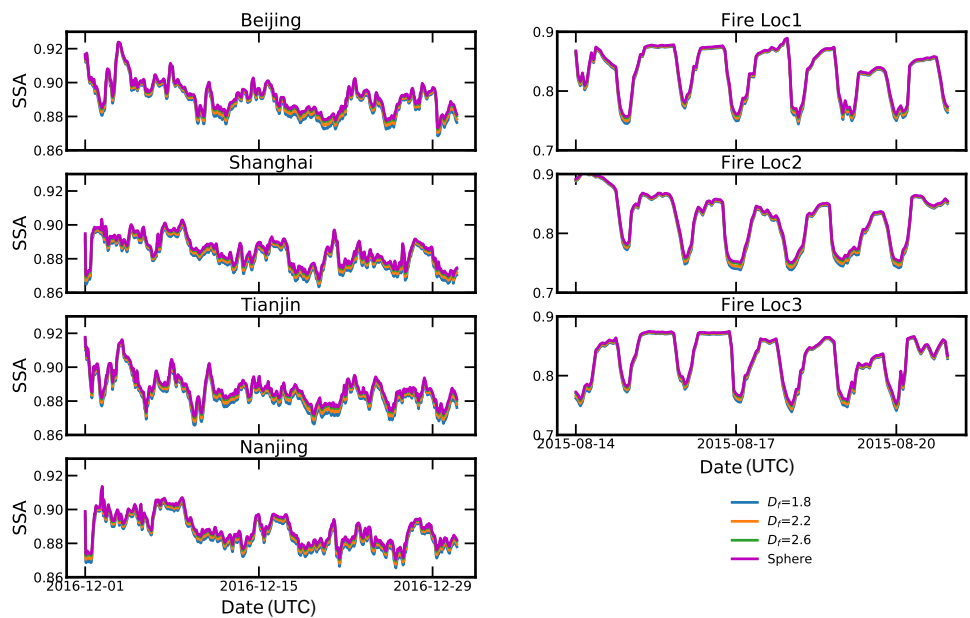


Figure 10. The comparison of SSA using different BC models, $\lambda=450$ nm.

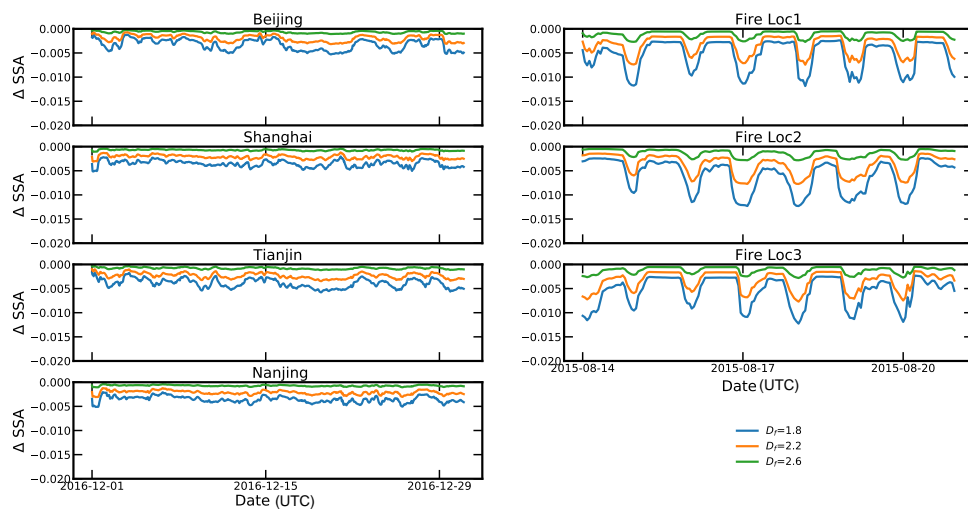


Figure 11. The SSA differences between fractal aggregate models and the spherical model, $\lambda=450$ nm.

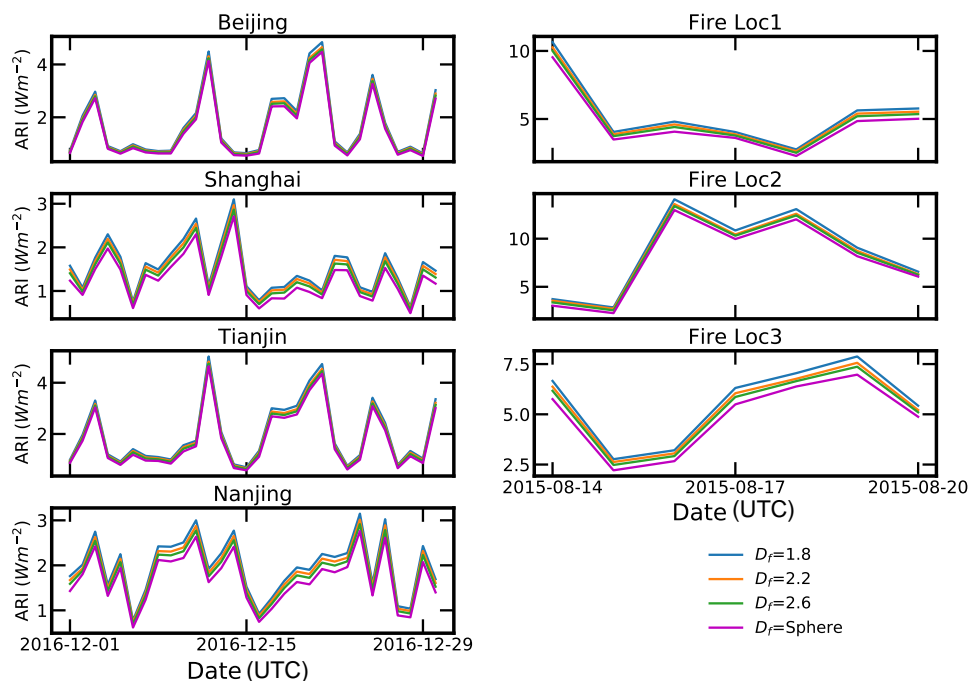


Figure 12. The clear-sky BC ARI at the typical urban polluted cities in eastern China and fire sites in the northwest US calculated using different BC models.

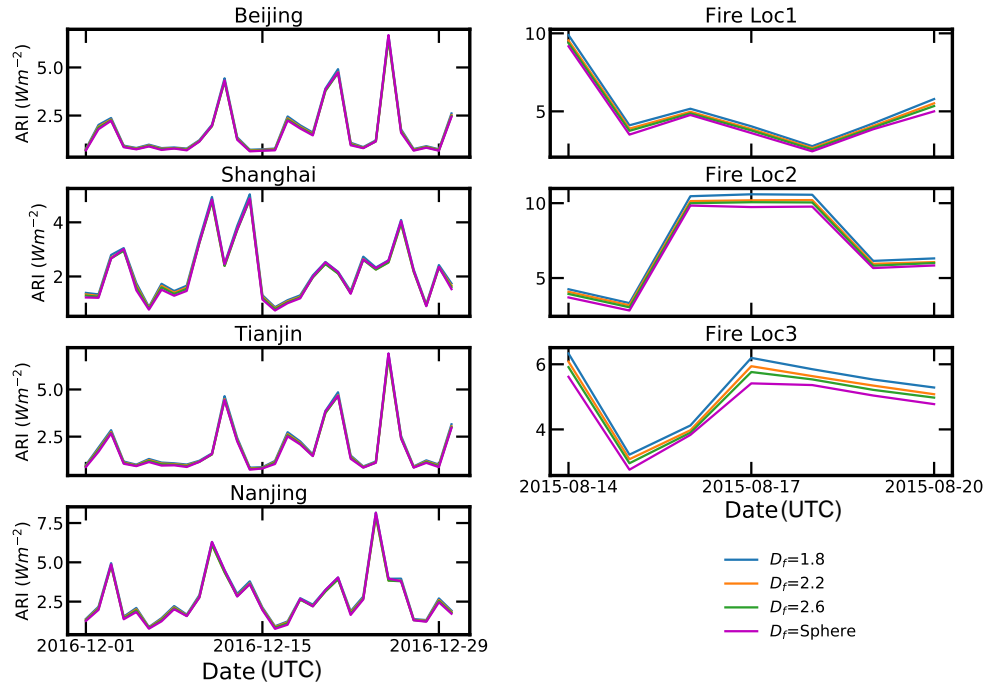


Figure 13. The all-sky BC ARI at the typical urban polluted cities in eastern China and fire sites in the northwest US calculated using different BC models..

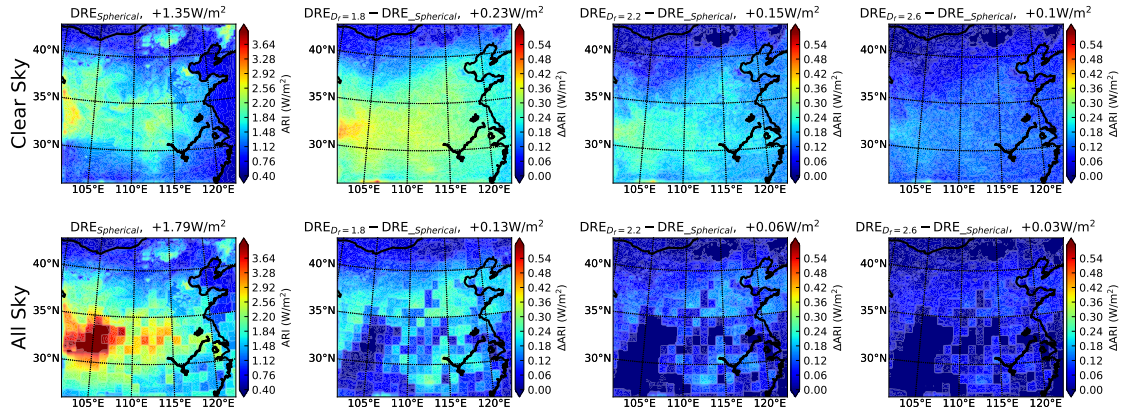


Figure 14. The regional-mean BC ARI at the TOA in eastern China for different particle models.

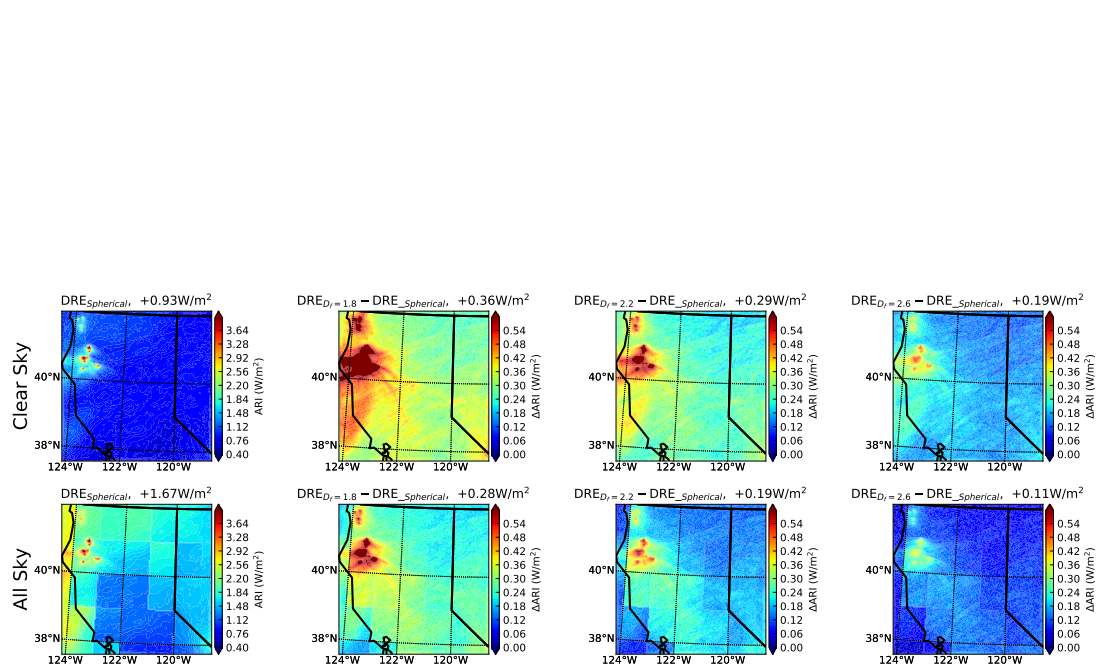


Figure 15. Similar to Figure 14, but for northwest US.

705 *Code availability.* The WRF-Chem is publicly available from <https://ruc.noaa.gov/wrf/wrf-chem/>; MSTM can be download from <https://www.eng.auburn.edu/~ckwski/scatcodes/>; The libRadtran is available from <http://www.libradtran.org/doku.php>; The FlexAOD can be requested from Prof. Curci (<http://pumpkin.aquila.infn.it/flexaod/>).

Data availability. The look-up tables calculated in this work can be obtained from https://figshare.com/articles/dataset/Look_up_tables_zip/13096241. The PM2.5 data in China was obtained from <https://www.aqistudy.cn/historydata/>, and PM2.5 data in the northwest US can be found from <https://www.epa.gov/outdoor-air-quality-data/download-daily-data>.
710

Author contributions. JL and QXZ conceived the presented idea. JL developed the models, performed the computations, and wrote the paper. ZQL, CZ, YMZ, RKC, YZ verified the simulation methods and results. QXZ revised the paper and supervised the findings of this work. GC developed the FlexAOD model. All authors discussed the results and contributed to the final paper.

Competing interests. The authors declare that they have no conflict of interest.