



1	Five satellite sensor study of the rapid decline of wildfire smoke in the
2	stratosphere
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4	Bengt G. Martinsson*, Johan Friberg, Oscar S. Sandvik, and Moa K. Sporre
5	Department of Physics, Lund University, Lund, Sweden
6	
7	*Email: bengt.martinsson@nuclear.lu.se
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10 Abstract

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Smoke from Western North American wildfires reached the stratosphere in large amounts in 12 August 2017. Limb-oriented satellite-based sensors are commonly used for studies of wildfire 13 14 aerosol injected into the stratosphere (OMPS-LP (Ozone Mapping and Profiler Suite Limb Profiler) and SAGE III/ISS (Stratospheric Aerosol and Gas Experiment III on the International 15 Space Station)). We find that these methods are inadequate for studies the first 1-2 months after 16 such a strong fire event due to event termination ("saturation"). The nadir-viewing lidar CALIOP 17 (Cloud-Aerosol Lidar with Orthogonal Polarization) is less affected due to shorter path in the 18 smoke, and, further, provides means that we could use to develop a method to correct for strong 19 attenuation of the signal. After the initial phase, the aerosol optical depth (AOD) from OMPS-LP 20 and CALOP show very good agreement above the 380 K isentrope, whereas the OMPS-LP tends 21 to produce higher AOD than CALIOP in the lowermost stratosphere (LMS), probably due to 22 23 reduced sensitivity at altitudes below 17 km. Time series from CALIOP of attenuation-corrected 24 stratospheric AOD of wildfire smoke show an exponential decline during the first month after the fire, which coincides with highly significant changes in the wildfire aerosol optical properties. 25 The AOD decline is verified by the evolution of the smoke layer composition, comparing the 26 aerosol scattering ratio (CALIOP) to the water vapor concentration from MLS (Microwave Limb 27 Sounder). Initially the stratospheric wildfire smoke AOD is comparable with the most important 28 volcanic eruptions during the last 25 years. Wildfire aerosol declines much faster, 80 - 90% of 29 30 the AOD is removed with a half-life of approximately 10 days. We hypothesize that this dramatic 31 decline is caused by photolytic loss. This process is rarely observed in the atmosphere. However, 32 in the stratosphere this process can be studied with practically no influence from wet deposition, in contrast to the troposphere where this is the main removal path of sub-micron aerosol particles. 33 Despite the loss, the aerosol particles from wildfire smoke in the stratosphere are relevant for the 34 climate. 35

1. Introduction

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Background stratospheric aerosol is composed of sulfuric acid, water, carbonaceous components, 39 40 and minor extraterrestrial and tropospheric components (Murphy et al., 2007; Kremser et al., 41 2016; Martinsson et al., 2019). Volcanism is a strong source of the stratospheric sulfurous, 42 carbonaceous and ash aerosol (Martinsson et al., 2009; Andersson et al., 2013; Friberg et al., 2014). Large eruptions, like that of Mt Pinatubo in 1991, affect the stratosphere for several years, 43 causing global cooling of several tenths of degrees Kelvin (Kremser et al., 2016). These eruptions 44 are scarce, only a few per century (Ammann et al., 2003; Stothers, 2007). Moderate eruptions are 45 more frequent contributors to the stratospheric aerosol (Vernier et al., 2011; Andersson et al., 46 2015; Friberg et al., 2018), forming the persistently variable stratospheric background aerosol 47 48 (Solomon et al., 2011).

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50 The stratospheric aerosol is also influenced by pyrocumulonimbus clouds (pyroCb) that form

51 during extreme weather conditions in connection with intense wildfires (Fromm et al., 2010). The





52 ongoing climate change is projected to increase the frequency of large wildfires (Kasischke et al., 2006; Dennison et al., 2014). Interestingly, the two largest events have, in terms of stratospheric 53 impact, occurred during the last few years, in North America 2017 (Peterson et al., 2018) and 54 55 Australia 2019-2020 (Kablick et al., 2020). Here we investigate the great pyroCbs formed in western North America on August 12 - 13, 2017. Figure 1a shows an example on the strong 56 impact on the stratospheric aerosol of the 2019 Raikoke volcanic eruption, one of the strongest 57 eruptions post Mt Pinatubo in 1991. In comparison, Figure 1b demonstrates the formidable early 58 impact of wildfire aerosol. The stratospheric impact of that fire has been described in terms of 59 light-backscatter reaching unprecedentedly high values for a non-volcanic aerosol layer (Khaykin 60 et al., 2018), light extinction about 20 times higher than after the Pinatubo volcanic eruption in 61 62 1991 (Ansmann et al., 2018), and mass of smoke comparable to that of a moderate sized volcanic eruption (Peterson et al., 2018). The pyroCbs lifted smoke from the fire to the extratropical 63 tropopause region, where absorption of radiation by black carbon (BC) in the smoke induced 64 additional lift to 23 km altitude in 2 months (Yu et al., 2019). 65 66 67 Smoke particles from wildfires contain a dominating fraction of organic matter by mass (Garofalo et al., 2019). Organic aerosol is susceptible to photochemical loss (Jimenez et al., 68 2009), and laboratory studies have demonstrated that this phenomenon could be an important 69 sink of secondary aerosol mass (Molina et al., 2004; Sareen et al., 2013). The residence time of 70 71 stratospheric air spans months to years depending on its path in the Brewer-Dobson circulation 72 (Engel et al., 2009; Bönisch et al., 2009). Due to very low probability of clouds, fine aerosol 73 particles have considerably longer residence times in the stratosphere than in the troposphere, 74 which further emphasizes the importance of investigating photochemical loss in the stratosphere 75 (Martinsson et al., 2019). 76 77 The aim of this study is to further understand the stratospheric aerosol sources and its climate impact. We develop methodology to correct for attenuation in dense smoke layers from wildfires 78 79 to properly deal with intense smoke injections into the stratosphere, with two main questions: 1) does photochemical loss of wildfire smoke occur in the stratosphere, and 2) how does the AOD 80 of smoke from the wildfire studied here compare with volcanic aerosol? 81 82 83 The first decade of the 21st century was characterized by slower temperature evolution than 84 anticipated from CMIP5 models (Fyfe et al., 2016). The discrepancy was attributed to interdecadal Pacific oscillation (Medhaug et al., 2017), variations in solar forcing (Myhre et al., 2013) 85 and aerosol in the stratosphere from moderate volcanic eruptions (Santer et al., 2014). Should 86 wildfire smoke in the stratosphere be added to this list of phenomena that require more attention 87 in climate models? 88 89

90 Our investigation deals with the evolution of the wildfire AOD, and aerosol optical properties

91 obtained from the lidar CALIOP aboard the CALIPSO (Cloud-Aerosol Lidar and Infrared

92 Pathfinder Satellite Observation) satellite, OMPS-LP/Suomi and SAGE III/ISS in comparison

93 with volcanic injections to the stratosphere. Additionally, the water vapor concentrations of

94 individual smoke layers are investigated by the MLS, the spatial evolution of smoke layers is





investigated using OMPS-NM (Ozone Mapping and Profiler Suite Nadir Mapper), and the AODs
and extinction coefficients obtained from CALIOP are compared with that of OMPS-LP and
SAGE III/ISS.

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2. Methods

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This study of the dense stratospheric smoke layers from pyro-cumulonimbus formed over
Western North America in August 12 – 13, 2017 is based on five satellite sensors. For four of
them, OMPS-LP, SAGE III/ISS, MLS and OMPS-NM, high level products (Level 2) are used.
The CALIOP data evaluation is based on a Level 1 product. A method to correct for attenuation
of the CALIOP laser beam in the smoke layers is presented. For these reasons CALIOP requires
more space in this section compared to the other methods.

- 107108 2.1 CALIOP
- 109

110 The evaluation of the CALIOP instrument carried by the CALIPSO satellite is based on version 111 4-10, level 1B data. CALIOP measures backscattering of laser light at two wavelengths, 532 and 112 1064 nm. For the shorter wavelength, scattered laser light is detected in parallel and

113 perpendicular polarizations relative to the outgoing beam. These almost nadir-viewing aerosol

and cloud measurements result in high resolution vertical profiles. For the altitude ranges <8.2,

115 8.2 - 20.2, 20.2 - 30.1 and 30.1 - 40 km the vertical resolutions are 30, 60, 180, and 300 m,

respectively. CALIPSO orbits between 82° S and 82° N, completing 14 – 15 orbits per day
(Winker et al., 2007; Winker et al., 2010).

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119 *2.1.1 AOD*

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Stratospheric AOD was obtained by integrating the backscattering intensity corrected for 121 122 attenuation (described below) from the tropopause to 35 km altitude. Figure 1b illustrates how 123 attenuation of the laser signal strongly reduced the signal below the dense smoke layer between 11 to 16 km altitude. We use the tropopause height according to MERRA-2 supplied with the 124 version 4.10 CALIOP data, which is a mixture of a dynamic and a thermal tropopause. The AOD 125 was averaged in the 20 - 80° N latitude range, where all nighttime swaths available from 126 CALIOP were included. The data were averaged over all longitudes in one-degree latitude bands, 127 and these latitude bands were averaged for the 20 - 80° N latitude range using area-weighting. 128 129 For dense layers, the lidar ratios estimated for the individual smoke layers were applied (explained below). Apart from the first few days the lidar ratio shows no temporal evolution, it is 130 found to have geometrical mean of 48.9 sr with double-sided 95% confidence interval of 47.6 – 131 132 50.3 sr (Figure 2a), which is close to the typical background lidar ratio of 50 sr (Jäger and Deshler, 2003). For layers that were not dense, the lidar ratio was held at this typical background 133 level. The volume depolarization ratio (δ_v) contains information that can be used to classify 134 135 aerosol layers. When the depolarization ratio is less than 0.05 the data is considered background and the lidar ratio is set to 50 sr. Ice-clouds were removed in the lowest 3 km of the stratosphere 136 by identifying them in stratospheric layers where the backscattering was high (attenuated 137 backscattering larger than 0.0025 km⁻¹ sr⁻¹). Data in these layers were classified as probable 138 clouds if their δ_v was higher than 0.20, or smoke if δ_v was between 0.05-0.20, after which the data 139





within each swath were clustered depending on their location. Noise in the data led to some lone 140 141 pixels within layers of either ice or smoke. These were reclassified depending on the surrounding pixels, making sure that no single pixel marked as aerosol occurred within the ice-cloud layers. 142 Layers of ice-clouds were then expanded upwards and horizontally to capture faint edges of the 143 clouds (Friberg et al., 2018). The classification was carried out on data at 8 km resolution along 144 each swath with their highest vertical resolution (30, 60, or 180 m, depending on altitude), after 145 146 which the tropospheric data were removed. Possible polar stratospheric cloud (PSC) signals north 147 of 45°N were excluded by classifying pixels with temperature below 195 K as possible PSC occasions. Underlying pixels were also excluded, to prevent bias from attenuation of the lidar 148 signals or from settling ice-crystals (Friberg et al., 2018). 149 150 151 2.1.2 Attenuation correction and radiative properties of individual smoke layers 152 153 The evolution of the lidar, color and depolarization ratios were investigated using 32 separate smoke layer measurements over the period 3 - 59 days after the fire. CALIOP has a statistical 154 disadvantage compared with lidars at the ground (Baars et al., 2019), because of small solid angle 155 due to long distance to the stratosphere (~700 km) and short measurement time. Optical 156 157 properties of old and faint individual smoke layers therefore could not be quantified with high 158 precision using CALIOP. The faint layers though still affect the AOD determinations described 159 above, where AOD elevation after the fire remains approximately one year. Out of the 32 smoke layers studied, 29 were night-time measurements, whereas the remaining three are defined as 160 day-time measurements. These latter ones increased the number of early observations (day 3 - 5)161 and were taken when the disturbance from solar radiation is small, i.e., shortly before the night. 162 163 164 The first weeks after the fire the smoke layers could be very dense with layer AODs exceeding 1, causing strong attenuation of the CALIOP signals with two-way transmissions down to below 165 0.01. For the 532 nm wavelength the particle lidar ratio was estimated by aiming the scattering 166 ratio (R; total-to-molecular backscattering ratio) below a smoke layer to a target value. The target 167 value was obtained from the background scattering ratio beside each smoke layer investigated, 168 169 which on average is R = 1.08, with standard deviation ± 0.05 . To reduce influence from noise, the CALIOP data were averaged along the swath. The averaging range varied between the smoke 170 layers, due to its extension along the swath, the homogeneity of the layer, and avoidance of sub-171 layer features. 172 173 The particle lidar ratio of an individual smoke layer was iterated until reaching the target value (R 174 = 1.08) described above from the combined effect of all altitude pixels. Pixels at altitudes outside 175 the smoke layer were set to the background lidar ratio of 50 sr (Jäger and Deshler, 2003). The 176 altitude resolution provided in the CALIOP data was used, where each altitude pixel (j) is 177 corrected for attenuation. The calculation starts at the highest altitude (40 km) and continues 178 downwards in two rounds. In the first round the star-marked quantities of equations 1-3 were 179 computed, correcting for attenuation from overlaying pixels. Before moving to the next altitude, 180 we account for self-attenuation from the pixel itself (equations to the right, without a star): 181 182





183
$$\beta_j^* = \frac{\beta_j'}{\prod_{k=1}^{j-1} T_k^2}; \qquad \beta_j = \frac{\beta_j^*}{\sqrt{T_j^{*2}}}$$
(1)

where β' is the attenuated backscattering and T² the two-way transmissions from both particles 185 and molecules. The two-way particle transmission is obtained by first computing the AOD: 186 187

188
$$AOD_j^* = (\beta_j^* - \beta_{m,j})S_p\Delta z_j;$$
 $AOD_j = (\beta_j - \beta_{m,j})S_p\Delta z_j$ (2)
189

where Δz_i is the height of the altitude pixel, $\beta_{m,i}$ is backscattering from air molecules, and S_p the 190 lidar ratio of the aerosol particles. The molecular lidar ratio, for computation of the molecular 191 extinction, was set to 8.70477 sr (Prata et al., 2017). CALIOP measurements are affected by 192 multiple scattering (Wandinger et al., 2010), causing overestimation of the backscattering. The 193 multiple scattering factor (η), the ratio of the apparent to the actual extinction coefficient, is not 194 known. Previous estimates are in the range 0.85 - 0.95 for layers thicker than 500 m (Prata et al., 195 2017). Not correcting for multiple scattering results in determination of the effective lidar ratio, 196 which is lower than the actual lidar ratio by a factor η . In equation 2 thus the backscattering 197 inflated by multiple scattering is multiplied by an underestimated lidar ratio to, at least in part, 198 compensate for the effects of multiple scattering on the AOD. The two-way transmission of 199 200 altitude pixel j due to the particles present is obtained from:

201

202
$$T_p^*$$
203

$$T_{p,j}^{*2} = \exp(-2AOD_j^*);$$
 $T_{p,j}^2 = \exp(-2AOD_j)$ (3)

204 These calculations in equations 1 - 3 are carried out until the background layer between altitudes a and b below the smoke layer reaches the target scattering ratio of 1.08 (Figure 3a): 205 206

207
$$R = \frac{\sum_{a}^{b} \beta_{j}}{\sum_{a}^{b} \beta_{m,j}}$$
(4)

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Error estimates of the effective lidar ratio were obtained by varying the target scattering ratio 209 210 from its average value (R = 1.08) mentioned above, to its ± 0.05 standard deviation range. The 211 fitting uncertainty in these estimates is strongly dependent on the light extinction in the smoke layer. Dense layers result in very small uncertainties in the effective lidar ratio because of the 212 strong impact on R from a slight change in the extinction. Layers with lower extinction 213 progressively increase the uncertainties of the estimate. When the error estimate of the effective 214 215 lidar ratio fit exceeds 25% the result is excluded from the data analysis, which terminates estimates of lidar ratios from day 22 after the fire. 216 217

The color ratio, the ratio between the backscattering at 1064 nm to 532 nm wavelength, is 218

affected by a difference in attenuation of the two wavelengths. This is clearly visible for dense 219

220 smoke layers in the CALIOP browse images by a gradual increase of the color ratio through the

layer because of the weaker attenuation for 1064 nm wavelength than for 532 nm (Figure 1d). 221





Therefore, estimations of the attenuation were undertaken also for the long wavelength. The 222 molecular backscattering is assumed to be 1/16 of that at 532 nm $(1/\lambda^4$ dependence of Rayleigh 223 scattering). Weak molecular scattering at 1064 nm prohibits lidar ratio estimation at that 224 wavelength by CALIOP. Instead, the lidar ratio was assumed to be 60 sr, inducing uncertainties 225 in the color ratio. The volume color ratio is obtained from: 226

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228
$$\chi = \sum_{k=top}^{base} \beta_{1064,k} / \sum_{k=top}^{base} \beta_{532,k}$$
 (5)

To limit influence from attenuation in the color ratio computations, the estimates were based on 230 the upper part of a smoke layer. Starting from the top of the smoke layer, the computations were 231 truncated when the two-way transmission of the 532 nm wavelength fell below 0.7. Varying the 232 1064 nm wavelength lidar ratio in the wide range of 60 ± 20 sr the uncertainty in the color ratio 233 becomes less than $\pm 5\%$ with this constraint applied. From the color ratio we define the particle 234 235 color ratio:

237
$$\chi_p = \sum_{k=top}^{base} (\beta_{1064,k} - \beta_{m,1064,k}) / \sum_{k=top}^{base} (\beta_{532,k} - \beta_{m,532,k}) = \frac{\chi R}{R-1} - \frac{1}{16(R-1)}$$
(6)
238

where we made use of the wavelength dependence of Rayleigh scattering for molecular 239 240 scattering, and the scattering ratio for the 532 nm wavelength was obtained from eqn. 4.

241

We also investigated the depolarization of the scattered laser beam at 532 nm by first forming the 242 volume depolarization ratio: 243

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246

$$\delta_{v} = \sum_{k=ton}^{base} \beta_{k}^{t}$$

$$\delta_{\nu} = \sum_{k=top}^{base} \beta_{532+,k}^{\prime} / \sum_{k=top}^{base} \beta_{532,k}^{\prime}$$
(7)

where symbol \perp indicates scattered light polarized perpendicularly to the incident beam. Having 247 access to the volume depolarization and an estimate of the molecular depolarization ratio $\delta_m \approx$ 248 0.003656 (Prata et al., 2017; Hostetler et al., 2006) the particle depolarization ratio is obtained 249 from: 250

251

252
$$\delta_p = \frac{\delta_v - \delta_m + \delta_v (1 + \delta_m)(R-1)}{\delta_m - \delta_v + (1 + \delta_m)(R-1)}$$
(8)

253

where R is obtained from eqn. 4. 254

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256 2.2 Extinction coefficients and AOD from OMPS-LP

257

The aerosol data from OMPS-LP (Chen et al., 2018; Jaross et al., 2014; Loughman et al., 2018) 258 have lately been used extensively in the literature on volcanic and wildfire impact on the 259 stratospheric aerosol. Several data products are available, here we use the recently released Level 260 261 2 product: Suomi-NPP OMPS LP L2 AER Daily Product, version 2.0 (Taha et al., 2020). The 262 polar-orbiting Suomi satellite completes 14 - 15 laps per day. OMPS-LP is a limb-scattering





263 method that collects data looking backwards along the satellite orbit, and along two other directions separated by 4.25° from the orbit, giving a cross-track separation of approximately 250 264 km at the tangent point. Measurements are undertaken in the wavelength and altitude ranges of 265 290 - 1000 nm and 10 - 80 km, respectively. The vertical resolution of OMPS-LP is 1.5 - 2 km 266 (Rault and Loughman, 2013). The aerosol product used here comprises 6 wavelengths (510, 500, 267 268 675, 745, 869 and 997 nm). The group responsible for the OMPS-LP version 2.0 data (Taha et 269 al., 2020) recommends caution when using data from altitudes below 17 km altitude due to loss of sensitivity. This problem can be reduced by use of the 745 nm and longer wavelengths. Here 270 we will make use of two of wavelengths: 745 nm because of the reduced problem with 271 sensitivity, and 510 nm because it is the wavelength closest to that of CALIOP (532 nm). 272 273 The OMPS-LP aerosol extinction coefficients are provided on a grid with a vertical resolution of 274 275 1 km. To study the smoke from the August 2017 fire we compute the average AOD over all longitudes in the latitude interval $20 - 80^{\circ}$ N for three layers, the LMS (tropopause to 380 K 276 277 isentrope), lower Brewer-Dobson branch (380 - 470 K) and the upper Brewer-Dobson branch 278 (470 K to 35 km altitude). The OMPS-LP version 2 dataset use a cloud detection algorithm (Chen et al., 2016), and comes in two forms: one without filtering out signals from clouds, and the other 279 280 where signals affected by clouds and polar stratospheric clouds are removed. In Figure 4 we show both these varieties for 745 nm wavelength, and, with and without flags regarding data 281 quality including profile retrieval errors (named RetrievalFlags in the OMPS-LP files), high root-282 mean squares (ResidualFlags), and further errors from the South Atlantic anomaly, disturbances 283 284 from the Moon, solar eclipses, planets, and satellite maneuvers (SwathLevelQualityFlags). In the 285 two upper layers (Figures 4a and b) the differences are usually small between the varieties except 286 for some spikes, whereas the LMS data (Figure 4c) show large stochastic variability as well as periods of clear differences between the varieties. Since this data is taken well below 17 km 287 288 altitude, sensitivity issues can be expected (Taha et al., 2020), see above. Days 130 – 190 (during December 2017 to February 2018) several spikes appear in the two higher layers which likely are 289 290 caused by polar stratospheric clouds. The data set filtered for clouds and flagged stands out by comparably small peaks, whereas the differences between the varieties usually are small 291 292 elsewhere. We therefore select the cloud-filtered and flagged data for further analysis in the 293 coming sections.

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2.3 Extinction coefficients from SAGE III/ISS 295

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SAGE III/ISS is a limb-viewing instrument based on solar occultation. Here we make use of 297 Level 2 aerosol extinction coefficients (SAGE III/ISS User's Guide, 2018), version 5.10, 298 supplied with a vertical resolution of 0.5 km. The upper limit of the slant path optical depth is 299 about 8, translating to a vertical optical depth of approximately 0.02 (SAGE III/ISS User's Guide, 300 2018). The orbiting of ISS differs markedly from the polar orbiting satellites CALIPSO 301 302 (CALIOP) and Soumi (OMPS-LP). This causes sporadic coverage by ISS of the latitudes of

interest here, resulting in that no average AODs over the $20 - 80^{\circ}$ N latitude range could be 303

304 formed with adequate time resolution. However, daily maximum extinction coefficients from

SAGE III/ISS could, when available, be included in a comparison with CALIOP and OMPS-LP. 305





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2.4 Water vapor measurements from MLS 307

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Water vapor concentrations (mixing ratio) in individual smoke layers was obtained from the MLS 309 instrument aboard the Aura satellite (Waters et al., 2006) in 12 vertical steps per decade of 310 311 pressure (version 5.0-1.0a, level 2). In nighttime measurements from days 6-59 after the fire, 312 the smoke layers studied by CALIOP were also investigated with MLS in almost simultaneous measurements, both instruments being on satellites that are members of the A-train. Data in the 313 10-316 hPa atmospheric pressure range were used, with vertical resolution 1.3-3.2 km 314 (Livesley et al., 2020). Limited vertical resolution induces problems to obtain well defined 315 observation of H₂O concentration of smoke layers close to the strong H₂O concentration gradient 316 across the tropopause. Close to the tropopause, but in the stratosphere, no H₂O peak from a 317 318 smoke layer can be detected. As the distance to the tropopause increases, an H₂O peak from the smoke layer becomes discernible. Further up from the tropopause, when the peak H₂O 319 concentration is well above the extratropical troppause at atmospheric pressure of less than 110 320 321 hPa, a deep minimum appears between the tropopause gradient and the peak from the smoke layer. All H₂O peaks were fitted with a Gaussian distribution operating on logarithmic pressure 322 323 and H₂O concentration to obtain estimates of the peak concentration and the corresponding atmospheric pressure. To investigate a time dependence in the smoke layer composition the peak 324 H₂O concentration (C_{H2O}) was compared with the attenuation-corrected aerosol scattering ratio 325 (R) from CALIOP, the optical equivalent of the mixing ratio, where the latter was obtained by 326 327 forming the geometrical mean over 900 m around the peak scattering ratio. The ratio of the of the 328 two quantities (R/C_{H2O}) was formed, and its dependence on time from the fire was studied. Out of the 13 smoke layers available with peak water vapor concentrations above the altitude of 110 hPa 329 atmospheric pressure, one was flagged as low quality in the MLS data set, leaving 12 330 observations for the study of the R/C_{H2O} evolution. 331

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333 2.5 UV aerosol index from OMPS-NM

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335 The UV aerosol index of OMPS-NM based on measurements at two wavelengths, 340 and 378.5

nm, is the official NASA aerosol index product according to OMPS-NM (NMMIEAI-L2 V2.1.1) 336 337 release notes (Torres, 2019). For strongly UV absorbing aerosols, like black carbon from wildfires, the UV aerosol indexes strongly increases with altitude (Herman et al., 1997). Here the 338 339 OMPS-NM UV aerosol index was used to map the geographical evolution of the smoke layers, 340 that according to CALIOP measurements were distributed in both the troposphere and the 341 stratosphere.

342

3. Results

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345 Here we use an approach based on five satellite sensors to study the influence on the stratosphere of the great North American fire in August 2017. We start by briefly describing results from the 346 347 method to correct CALIOP data for attenuation of the backscattered laser light. Then follows a comparison of AODs obtained from OMPS-LP and CALIOP. Absorption aerosol index from 348





349 OMPS-NM is used to describe the dispersion of the wildfire aerosol in the stratosphere. To

explain differences in AOD between OMPS-LP and CALIOP, a comparison of extinction
 coefficients follows, where results from SAGE III/ISS also are included in the comparison. The

evolution of the optical properties of the wildfire aerosol is then described, before the North

American wildfire aerosol is compared with volcanic influence on the stratospheric AOD.

Finally, the fifth data set, water vapor from the MLS, is introduced in the discussion section,

355 where the evolution of the wildfire aerosol in the stratosphere is analyzed.

356

357 3.1 Correction for attenuation

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The smoke layers usually were 1 - 3 km thick and could extend several degrees in longitude and 359 latitude. Measurements with the CALIOP lidar provide, in addition to short, nadir-viewing 360 measurement path in dense layers, the advantage that the signal is retrieved as a function of 361 position along the laser path with high resolution, which can be used to correct for attenuation of 362 the signal. Figure 3a shows the attenuated scattering ratio (R'; the measured backscattering 363 364 divided by the calculated molecular backscattering) from an example-smoke-layer measured on August 16, 2017. The scattering ratio should be close to 1 in air layers with low aerosol 365 366 concentration, whereas values below 1 is caused by attenuation from particles. As can be seen in Figure 3a, the attenuated scattering ratio first increases (starting from above the layer). Then the 367 signal decreases and reaches well below unity from 11 km altitude and downwards, i.e., well 368 below the scattering ratio of particle-free air. By techniques described in the Methods section we 369 370 correct for attenuation and fit the lidar ratio (Figure 2a) to obtain an estimate of the

371 backscattering without attenuation, as illustrated by the scattering ratio (R) in Figure 3a.

372

The evolution of wildfire aerosol from day 3 to 59 after the North American PyroCbs on August

12, 2017, is first investigated by comparing 32 smoke layers from individual CALIOP swaths.

The influence from attenuation is shown in Figure 3b. Clear deviation from the 1:1 line appears

already at layer attenuated (uncorrected) AODs (AOD_{att}) of 0.12, and 50% reduction of the signal

appears at layer AOD_{att} of approximately 0.25. Reduction by more than 50% appears until day 10

after the fire, whereas those measurements close to the 1:1 line were taken after day 30. The

AOD, i.e., the AOD corrected for attenuation, exceeds the AOD_{att} by more than a factor of 5 in
the densest layers of this study (Figure 3b).

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382 *3.2 Comparison of CALIOP and OMPS-LP*

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384 To study the evolution of the stratospheric AOD, we form a 3-dimensional box in the

stratosphere extending over all longitudes in the $20 - 80^{\circ}$ N latitude range. In this box we use all

daily profiles ,14 - 15 CALIOP and 42 - 45 OMPS-LP, to form the average AOD. We apply the method to correct CALIOP data for attenuation, as described in the Methods section. AODs are

method to correct CALIOP data for attenuation, as described in the Methods section. AODs are
 computed for three layers, the LMS, the lower Brewer-Dobson branch, and the upper Brewer-

389 Dobson branch, as shown in Figure 5.





391 When comparing AODs, the measurement wavelengths should be as close as possible, due to the wavelength dependence of scattering. CALIOP AODs are shown for 532 nm wavelength, and the 392 OMPS-LP data are shown for the close wavelength of 510 nm. In addition, the 745 nm AODs 393 from OMPS-LP is shown. The response to the 2017 North American fire is weak in the upper 394 Brewer-Dobson branch (Figure 5a), whereas the two lower layers (Figures 5b-c) show clear 395 396 increase of the AOD. Comparing the two methods, they agree well in the upper Brewer-Dobson 397 branch. In the lower Brewer-Dobson branch we see good agreement between the two methods, except for the first 1-2 months after the fire where much higher AODs are recorded by CALIOP 398 (Figure 5b). The latter is also true for the LMS, whereas the general agreement between the two 399 methods is poor (Figure 5c). The OMPS-LP documentation advise against using data from below 400 approximately 17 km altitude, approximately the upper limit of the LMS, due to loss of 401 sensitivity (Taha et al., 2020). We therefore do not perform any further comparisons in the LMS. 402 The stratosphere above the LMS (above the 380 K isentrope) shows good agreement between the 403 two methods, except for the first 1 - 2 months after the fire (Figure 5d). 404

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406 *3.3 Early evolution of the smoke layers*

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408 The daily AOD averages show large variability the first days after the fire because the lidar measures narrow curtains through the atmosphere, Figure 5e. The variability remains until the 409 smoke layers become sufficiently dispersed, allowing several daily measurements of the smoke 410 layers. The nadir-viewing OMPS-NM provides UV (ultraviolet) absorbing aerosol index, where 411 412 strong signal for strongly UV light absorbing aerosol is obtained in the upper troposphere and the 413 stratosphere. Figure 6 shows the geographic evolution of the smoke layers from August 14 to 22, 414 2017 together with the orbits followed by the CALIOP measurements. Up to August 16 the smoke is found in a rather confined area. From August 17 the smoke layers are stretched in 415 416 Easterly direction, and after that the smoke spreads rapidly to the East. The dispersion gradually increases the number of daily CALIOP observations of the smoke. This can also be seen in 417 418 Figure 5e, where the variability in the daily AOD data becomes successively smaller. From day 10 (August 22) we see a clear pattern of decline of the AOD. 419 420 421

Figure 5e shows the total stratospheric AOD according to CALIOP from the tropopause to 35 km altitude. We see a strong decline of the stratospheric AOD the first 1.5 months after the fire, and a fitted exponential function has a half-life of 6.5 ± 0.9 days. Such a decline cannot be found in the OMPS-LP AODs, which instead are increasing during the first month.

425

To further investigate this clear difference between the two methods, individual smoke layers areinvestigated with respect to extinction coefficients. Figure 7a-d show the extinction coefficient of

strong smoke layers from four days in August and September 2017. From CALIOP we show the

429 attenuated extinction coefficients as well as the profiles corrected for attenuation. Together with

- the CALIOP data the OMPS-LP data closest by are shown. It is obvious that OMPS-LP shows
- 431 very much smaller reaction to the smoke layers than CALIOP. However, we cannot be sure that
- the two instruments viewed the same airmasses in these four examples, because the two
- 433 instruments do not belong to the same satellite constellation. To remove that obstacle, the daily





434 maximum stratospheric extinction coefficient from OMPS-LP was extracted and compared with 32 selected profiles' peak extinction coefficients from CALIOP. SAGE III/ISS was also included 435 in the comparison from day 19 after the fire. Unfortunately, the orbiting of ISS did not permit 436 measurements of the fire studied here before that day. The very strong signals from CALIOP are 437 not reflected in the OMPS-LP or SAGE III/ISS measurements, see Figure 7e. In part, this can be 438 439 explained by difference in vertical resolution, but as shown in Figures 7a-d, these high extinction coefficients extend to broad vertical ranges that should allow detection of strong signals also by 440 OMPS-LP and SAGE III/ISS. 441

442

443 There is one principal difference between CALIOP on one hand and OMPS-LP and SAGE

444 III/ISS on the other hand: whereas the former is nadir-viewing (vertical) the latter two methods 445 operate in limb orientation (horizontal). This is important, because the horizontal extension of

smoke layers is much larger, e.g., the smoke layer in Figure 1b has a vertical extension of

447 approximately 2 km, whereas the horizontal extension is approximately 700 km. The vertical,

two-way transmission to the CALIOP sensor through this layer is approximately 0.01, which we correct for. The horizontal path through this layer is 350 times longer, implying that the one-way

449 correct for. The horizontal path through this layer is 350 times longer, implying that the one-way 450 limb transmission becomes 10^{-350} for the same wavelength. Even if the horizontal extension

451 would be just one tenth the transmission is still as low as 10^{-35} . Obviously, the radiation used for

452 detection in OMPS-LP and SAGE III/ISS is rapidly eliminated in such smoke layers. Therefore,

- 453 these two methods are inadequate for studies of dense aerosol layers. The upper limit in terms of
- vertical AOD is estimated to 0.02 (SAGE III/ISS Users Guide, 2018), corresponding to the

extinction coefficient of 0.02 km⁻¹ for a 1 km thick layer. This problem is also acknowledged for
OMPS-LP (Chen et al., 2018; DeLand, 2019). Despite the clear limitation of OMPS-LP and

456 SAGE III/ISS in this respect, the large body of information on wildfires is based on these

458 methods, e.g., Bourassa et al., (2019), Das et al., (2021), Khaykin et al., (2020), Kloss et al.,

459 (2019), Torres et al., (2020) and Yu et al., (2019). By comparing with CALIOP we here show

that the limb-oriented techniques miss the dramatic events during the first 1 - 2 months after the fire. The rapid decline of the wildfire smoke will be further analyzed below.

462

463 3.4 Aerosol optical properties

464

465 To further investigate the unusual evolution of the AOD, we turn to the optical properties of the wildfire aerosol. The particle color and depolarization ratios are shown in Figure 2b and c. To test 466 467 the significance in the evolution the data were temporally divided into two equal halves by 468 number of data points, and geometric averages were formed (black lines in Figure 2). The particle color ratio shows a highly significant decrease comparing the first to the last half of the data 469 points, whereas the particle depolarization ratio increases with high significance. The change in 470 the optical properties takes place up to 15 - 30 days after the fire. This coincides with the decline 471 of the AOD, thus connecting a change of the aerosol properties to the AOD decline. 472 473

474 *3.5 Stratospheric AOD variability caused by volcanism and wildfires*





476 The stratospheric AOD varies considerably over time mainly due to influence from explosive volcanic eruptions as demonstrated in Figure 8, showing the period 2008 - 2018. In this time 477 span, nine volcanic eruptions clearly, but to varying degree, affected the stratospheric AOD. We 478 479 also identify two cases of influence from wildfires, the Victoria fire (Australia, 2009) and the fire studied here (Western North America, 2017). The residence time in the stratosphere varies from 480 481 several years for tropical injections into the upper layer representing the upper branch in the Brewer-Dobson circulation (BD) (Figure 8a), the order of a year in the shallow branch of the BD 482 circulation (Figure 8b), to months in the LMS (Figure 8c) (Friberg et al., 2018). The sum of the 483 three layers is shown in Figure 8d. The volcanic eruptions in these 11 years mainly affected the 484 two lower stratospheric layers, only the Kelut eruption (2014) clearly reached to the deep BD 485 branch. Fire aerosol contains black carbon, which absorbs radiation, heats surrounding air and 486 induces lifting, as observed after the fire studied here (Khaykin et al., 2018; Yu et al., 2019). 487 After both fires, we see weak AOD elevation in the deep BD branch (Figure 8a), but for the fire 488 studied here the two lower layers dominate the AOD, like most of the volcanic eruptions in the 489 490 eleven-year period.

491

Comparing the evolution of the AOD of the North American wildfire with the evolution of the 492 493 aerosol from two of the most important volcanic eruptions during the last 25 years (Figure 9), we find that the maximum stratospheric AOD after the fire is similar to that after the 2011 Nabro and 494 2009 Sarychev eruptions. During the first couple of months after volcanic events the AOD grows 495 due to formation of condensable sulfuric acid from the emitted volcanic gas sulfur dioxide. In 496 497 contrast, the wildfire aerosol displays a rapid decline during the first few weeks, before the AOD 498 stabilizes (Figure 9). This is followed by a period of rather stable AOD of more than 6 months, 499 before the AOD evolution turns to a slower decline towards background conditions, with similar seasonality as the aerosol from the volcanic eruptions discussed (Figure 9). This latter decline is 500 mainly caused by springtime transport out from the stratosphere at mid and high latitudes 501 (Bönisch et al., 2009; Martinsson et al., 2017). 502

4. Discussion

504 505

503

The smoke aerosol is distributed both in the LMS and in the lower BD branch like aerosol from 506 507 several volcanic eruptions (Figure 8). The rapid decline of the smoke aerosol during the first 508 month after the fire thus cannot be explained by transport out of the stratosphere. Measurements 509 with Raman lidars at three wavelengths indicate that the smoke from this North American fire 510 contain an accumulation mode but no coarse mode (Haarig et al., 2018; Hu et al., 2019). The influence from sedimentation on submicron diameter particles is small (Martinsson et al., 2005). 511 Moreover, the change in the particle depolarization ratio (Figure 2c) indicates change of the 512 aerosol particle properties, and the particle color ratio decrease after the fire (Figure 2b) is the 513 expected outcome for reduced particle sizes. Based on these arguments we turn the attention to 514 515 loss of material from the aerosol particles to the gas phase to explain the rapid decrease in AOD seen in Figure 5e. 516





- 518 Smoke layers contain water vapor that could induce hygroscopic growth/shrinkage. Water vapor
- 519 profiles for individual smoke layers from days 6 60 after the fire were obtained from the MLS.
- 520 Measurements close to the tropopause (Figure 10a) are affected by a steep gradient in H_2O
- 521 concentration. The profiles well above the gradient peaking at atmospheric pressure of less than
- 522 110 hPa are shown in Figure 10b. For the latter category the peak H_2O concentration is in the 523 range 7 – 14 ppmv, implying a maximum H_2O vapor pressure of 0.16 Pa. For typical conditions
- in the extratropics that vapor pressure corresponds to a relative humidity of a few percent or less
- 525 (Murphy and Koop, 2005).
- 526

527 To further investigate the smoke layers, the temporal evolution of the composition is studied by

528 forming the ratio of the mixing ratios of two components: aerosol backscattering and H₂O at the

529 peak of respective vertical distribution. As pointed out above, the strong H_2O gradient around the

tropopause affects the MLS measurements. But for the smoke layers higher up, peaking above

531 110 hPa, we find a rapid decrease in the aerosol scattering ratio compared with the H_2O

532 concentration (Figure 10c). Fitting an exponential function $\left(\frac{R}{C_{H2O}} = a + be^{-t/\tau}\right)$, the half-life

becomes 9.7±3.2 days, which is somewhat longer than that computed from the AOD (half-life

534 6.5±0.9 days). The rapid AOD decline (Figure 5e) is thus verified by relative concentrations of

aerosol and H₂O under well-controlled humidity conditions, whereas the low relative humidity
 rules out hygroscopic growth and influence from clouds as the explanation of the AOD decline.

537

538 The near-field wildfire aerosol contains, besides black carbon (Bond et al., 2013; Ditas et al., 539 2018), approximately 90% organic material (Garofalo et al., 2019). After emission, secondary 540 organic aerosol (SOA) is formed by oxidation of gas phase compounds (Shrivastava et al., 2017). Knowledge of processes controlling formation and removal in the atmosphere is limited (Hodzic 541 et al., 2016). Global aerosol models usually remove SOA mainly by wet (90%) and, to a smaller 542 543 extent, by dry deposition (Tsigaridis et al., 2014). In contrast to the species dominating the stratospheric aerosol and its precursor compounds during background conditions and volcanic 544 545 influence (sulfuric acid and sulfur dioxide), organic species are not the ultimate

thermodynamically stable compounds (Hallquist et al., 2009). Organic aerosol is an intermediate
state on routes, with little known rates, from emitted compositions to the highly oxidized gaseous
products CO and CO₂ (Jimenez et al., 2009). Modeling and numerous laboratory studies find

evidence for photolytic removal rates of organic aerosol similar to that of wet deposition in the

troposphere (Hodzic et al., 2016; Zawadowics et al., 2020). Recently, photolytic removal of

551 particulate SOA was included in the Whole Atmosphere Community Climate Model (WACCM6)

(Gettelman et al., 2019). Hodzic et al. (2015) estimate the photolytic loss over a 10-day period to

- 553 50% for most organic species at mid tropospheric conditions.
- 554

These high rates are disputed by Yu et al. (2019), claiming a lifetime of 150 days (halflife 104 days) of organic aerosol from the fire studied here, whereas Das et al. (2021) explain a similar

half-life of the same fire by large-scale circulation and particle sedimentation using OMPS-LP

- and modeling. The experimental data used here cannot differentiate these two explanations,
- although the slow part of the smoke decline is similar in seasonality to that of volcanic aerosol
- 560 (Figure 9) where photochemical loss is less important. The modeling study by Yu et al. (2019)





- 561 was based on mimicking the extinction according to SAGE III/ISS at 1020 nm wavelength at 18 km altitude. For three reasons their study misses the strong decline of the AOD during the first 562 month. Firstly, because the orbiting of ISS prohibits studies of the wildfire smoke the first 19 563 days after the fire, secondly because of the time required to transport the wildfire aerosol to 18 564 km altitude is approximately one month (Yu et al., 2019) and thirdly because problems with 565 event termination ("saturation"), see Figure 7e. We therefore conclude that that Yu et al. (2019) 566 could not observe the main decline of the aerosol taking place during the first 1-2 months after 567 the fire, see section 3.3 for further details. 568
- 569

570 Submicron aerosol particles have much longer residence time in the stratosphere than in the 571 troposphere due to sparsity of clouds, thus inhibiting the sink that traditionally is considered the 572 most important in the troposphere, i.e., wet deposition. This provides unique possibilities to study 573 photolytic loss without competition from other aerosol sinks. Interpreting the body of evidence 574 on the strong and rapid decline of the stratospheric AOD during the first month after the fire, we

575 find that photolytic loss of organic aerosol is a highly likely explanation. The rate of photolytic

577 could to some degree be affected by transport across the tropopause. Our strong experimental

evidence leads us to the hypothesis that the rapid decline of the wildfire aerosol in the
stratosphere with a half-life of 10 days is caused by photochemical loss of organic material. This

- 580 should be further investigated by modeling, but that is outside the scope of the present study.
- 581

582 To further put the strong early decline of wildfire aerosol into context, we compare the AOD 583 during background conditions (years 2013 and 2014) with the year of the fire. When the 584 contribution of the exponential term is very small of the wildfire aerosol (after 7 half-lives), the background is approximately 2/3 of the wildfire AOD (Figure 9). Taking the background into 585 account, the excess stratospheric aerosol due to the wildfire declines by 83% from the R/C_{H2O} 586 value day 10 after the fire. The process starts before day 10, indicating that almost all the organic 587 588 aerosol constituting approximately 90% of the near-field wildfire aerosol mass (Garofalo et al., 2019) could be lost by photolysis. Residual wildfire aerosol particles, likely stripped off by a 589 590 large fraction of its original organic content, remain in the stratosphere up to approximately one 591 year (Figure 9).

592

Finally, we investigate the stratospheric aerosol load from the wildfire by comparing with the more studied volcanic impact (Table 1). The AOD growth, the average AOD over one year from the fire/eruption subtracted by the average background AOD (2013 - 2014), is approximately 1/4 and 1/3 of that of two of the most important volcanic eruptions for the stratospheric aerosol in the last 25 years (Sarychev 2009, Nabro 2011). The average excess aerosol during the year following the fire corresponds to a radiative forcing of -0.06 W m⁻² in the region 20 - 80° N, using standard conversion as an approximation (Solomon et al., 2011).

600 601 Co

601 **Conclusions** 602





603 In this study we investigate massive injections of smoke into the stratosphere from the August 2017 North American wildfires using five satellite sensors. Methodology was developed to 604 correct CALIOP data for attenuation of the laser signal. The CALIOP AOD and extinction 605 coefficients were compared with OMPS-LP and SAGE III/ISS. From 1 - 2 months after the fire 606 we find that OMPS-LP and CALIOP AOD agree very well at altitudes above the 380 K 607 608 isentrope, where the former demonstrates high sensitivity with small statistical fluctuations. The methods differ dramatically during the first 1-2 months after the fire when the smoke layers are 609 dense, because the long optical path through the smoke of the limb-oriented instruments OMPS-610 LP and SAGE III/ISS cause event termination ("saturation"). This is clearly demonstrated by the 611 low daily maximum extinction coefficients of the two instruments, being orders of magnitude 612 lower than the peak extinction coefficients of CALIOP. The nadir viewing CALIOP experiences 613 a much shorter optical path, because the vertical extension of smoke layers usually are orders of 614 magnitude shorter than for limb orientation. We find that CALIOP is an indispensable tool for 615 studies of dense smoke layers entering the stratosphere after intense wildfires, providing signal 616 along the laser path that can be used to correct for attenuation. Once the smoke layers are 617 618 sufficiently thin, the limb technique OMPS-LP provide sensitive measurements of the AOD that can be used together with CALIOP. 619 620

The AOD from the wildfire declines exponentially with a half-life of 6.5 days. This decline is 621 further studied by the evolution of the ratio of the aerosol and water vapor mixing ratios of the 622 smoke layers, resulting in a massive decline of 80 - 90% of the wildfire aerosol with a half-life of 623 624 approximately 10 days. We find transport out of the stratosphere, sedimentation, influence from 625 clouds or hygroscopic growth/shrinkage to be highly unlikely explanations for the rapid decline of wildfire aerosol in the stratosphere. Based on strong experimental evidence we hypothesize 626 that photochemical loss of organic aerosol causes the rapid decline, which would mean that 627 almost the entire organic fraction of the wildfire aerosol would be lost in the exponential decline. 628 The half-life according to this study agrees well with results from laboratory studies and global 629 630 modeling. Our unique result could be obtained because of the long residence time of aerosol particles in the stratosphere, whereas tropospheric studies of photochemical loss are extremely 631 632 difficult because it is masked by SOA formation and wet deposition due to short residence time. The residual aerosol leaves the stratosphere within a year in the Brewer-Dobson circulation. 633 634 Despite the initial loss, the long-term effects of wildfire smoke on the stratospheric AOD and radiative forcing are considerable. The ongoing climate change is projected to increase the 635 636 frequency of wildfires, prompting the need for inclusion of wildfire impact on the stratospheric aerosol load in the climate models. 637

638

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647 648	Author Contributions
649	
650	B.G.M. designed the study, designed methodology, undertook part of the data analysis, and wrote
651	most of the paper. J.F. contributed to the design of the study, designed methodology, did part of
652	the data analysis, and wrote parts of the text. O.S.S. contributed to the data analysis and M.K.S.
653	contributed to the design of methodology. In addition, all authors participated in discussions and
654	commented on the manuscript.
655	-
656	Data availability
657	·
658	CALIOP V4.10 lidar data (https://search.earthdata.nasa.gov/search?fp=CALIPSO) are publicly
659	available.
660	OMPS-NM UV aerosol index was obtained from the publicly available site
661	https://worldview.earthdata.nasa.gov/.
662	OMPS-LP stratospheric aerosol optical depths were obtained from
663	https://disc.gsfc.nasa.gov/datasets/OMPS_NPP_LP_L2_AER_DAILY_2/summary
664	MLS water vapor concentrations were obtained from
665	https://disc.gsfc.nasa.gov/datasets?page=1&keywords=ML2H2O_005
666	SAGE III/ISS aerosol data were obtained from
667	https://asdc.larc.nasa.gov/project/SAGE%20III-ISS/g3bssp_51.
668	
669	Competing Interest
670	
671	The authors declare no competing interests.
672	
673	Additional Information
674	
675	Correspondence and requests for materials should be addressed to B.G.M.
676	
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959 Tables

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Table 1. Maximum and yearly average stratospheric AOD during background conditions and

962 during one year after the fire and after the two volcanic eruptions in Figure 9.

	Background	Background	Wildfire	Sarychev	Nabro
Year	2013	2014	2017	2009	2011
AOD max	0.009	0.009	0.020	0.028	0.017
AOD	0.0075	0.0074	0.0097	0.0169	0.0138
AOD growth ^a	-	-	0.0023	0.0095	0.0064
RF ^b	-	-	-0.06	-0.24	-0.16

aGrowth of AOD due to influence from wildfire/volcanism obtained by subtracting the average of 2013
and 2014 AOD.

^bRadiative forcing (W m⁻²) of the background-subtracted AOD.

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974 aerosol layers in the stratosphere three days after the 2019 Raikoke eruption and b) a

stratospheric smoke layer from the August 12, 2017, North American wildfire. c) Volume 975

976 depolarization ratio at 532 nm and d) attenuated color ratio (1064 to 532 nm) for the curtain in b).

The white lines in a) and b) show the position of the tropopause. 977







Figure 2. Particle optical properties during the first 60 days after the fire. Black error bars show 980 981 standard error and the double-sided 95% probability range of the geometric means. a) Particle lidar ratios for 532 nm where data points with fitting error exceeding 25% are discarded. The 982 black line shows the geometric mean after day 4, and the full and dotted blue lines show the 983 standard deviation and the double-sided 95% probability range of the distribution. b) Particle 984 color ratio (1064 nm divided by 532 nm wavelength backscattering) with exponential fit ($R^2 =$ 985 0.48, $P < 10^{-10}$), and c) particle depolarization ratio with exponential fit ($R^2 = 0.76$, $P < 10^{-10}$). 986 The color and depolarization ratios were divided in two equal groups by number of observations 987 to illustrate the highly significant changes with time of the optical properties, where the long and 988 short error bars are the standard error and the double-sided 95% probability range of the 989 geometric means. 990





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Figure 3. Illustration of methodology and its effect. a) The attenuated and corrected scattering
ratios as a function of altitude. Example of methodology for one smoke layer, where the
scattering ratio between 7.5 - 10 km altitude, below the smoke layer at 10.5 - 14 km, is targeted
to a value of 1.08 (explained in the method section) by iteratively fitting the lidar ratio for 532 nm
wavelength. b) The attenuated layer AOD (AOD_{att}) related to the layer AOD corrected for
attenuation. The 1:1 relation is shown by the full line.







Figure 4. OMPS-LP layer AODs averaged over 20 to 80° North for 745 nm wavelength using
data filtered and not filtered from clouds and polar stratospheric clouds, and with and without
data flagged for data quality. Layer AOD for a) the upper Brewer-Dobson branch (470 K
isentrope – 35 km), b) the lower brewer-Dobson branch (380 – 470 K) and c) the LMS
(tropopause – 380 K) are shown.







Figure 5. AOD evolution of the stratospheric AOD (daily average) from 75 days before to 310 1007 days after the 2017 western North American fires. Comparisons of AOD from CALIOP (532 nm) 1008 with OMPS-LP (510 and 745 nm) with cloud filtering and flags activated for a) the upper 1009 1010 Brewer-Dobson branch (470 K isentrope – 35 km, b) the lower Brewer-Dobson branch (380 – 470 K) c) the LMS (tropopause – 380 K), d) from 380 K to 35 km (sum of layers in a and b) and 1011 e) the stratosphere of CALIOP from the tropopause to 35 km (sum of layers in a, b and c). The 1012 black, full line is an exponential fit ($R^2 = 0.79$, $P < 10^{-10}$) to the AOD over days 10 - 115 after the 1013 fire. The total stratospheric AOD half-life of the fit is 6.5±0.9 days. 1014







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Figure 6. Daily OMPS-NM aerosol absorbing index (UV) August 14 – 22, 2017 over all
longitudes and latitudes 20 - 80° N. This index is sensitive to UV absorbing aerosol particles in
the upper troposphere and the stratosphere, where signals from tropospheric aerosol declines
faster than from stratospheric due to short residence time. The yellow lines indicate nighttime
swaths of the CALIPSO satellite, and the faint lines show CALIPSO daytime swaths.







Figure 7. Extinction coefficients according to CALIOP, OMPS-LP and SAGE III/ISS in the 20 80° North latitude range during the first 60 days following the North American fire. a – d)
selected profiles (attenuated and corrected for attenuation) from CALIOP compared with closest
profiles according to OMPS-LP. e) Peak extinction coefficient from selected CALIOP profiles
compared with daily maximum extinction coefficients from OMPS-LP and SAGE III/ISS. Note
that SAGE III/ISS data are missing the first 19 days because of irregular coverage of the latitude
range of interest.







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Figure 8. Zonally and eight-day moving average aerosol optical depth (AOD) of the stratosphere. 1031 1032 a - c) AOD in three layers obtained from CALIOP data (level 1B): a) 470 K potential temperature 1033 to 35 km (deep Brewer-Dobson branch), b) 380 – 470 K (shallow Brewer-Dobson branch), c) the 1034 tropopause to 380 K (LMS). d) The total AOD from the tropopause to 35 km altitude. Volcanic eruptions marked by white triangles: Kasatochi (Ka), Sarychev (Sa), Merapi (Me), Grimsvötn 1035 (Gr), Puyehue-Cordón Caulle (Pu), Nabro (Na), Kelut (Ke), Calbuco (Ca), and Ambae (Am), and 1036 1037 wildfires marked by orange circles: Victoria fire (Vi) and Western North American fires (Wn) at time and latitude of eruption/fire. The AODs are corrected for attenuation. 1038







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1040 Figure 9. Evolution of the AOD in the 20 - 80° N interval (8-day moving average) over two

1041 years: close to background conditions in the latitude interval studied (2013 - 2014), the year and 1042 the following year of the August 12, 2017, fire (2017 - 2018), and the same for two volcanic 1043 eruptions, the June 12, 2009, Sarychev (2009 - 2010) and June 12, 2011, Nabro (2011 - 2012)

1044 eruptions.







Figure 10. Water vapor in the smoke layer. Microwave Limb Sounder (MLS) measurements of water vapor concentrations (ppmv) Vs. atmospheric pressure for smoke layers a) close to the tropopause and b) well above the tropopause (atmospheric pressure < 110 hPa at the H₂O peak) for individual smoke layers available days 6 - 60 after the fire. c) The peak scattering ratio (R) according to CALIOP divided by the peak water vapor concentration (C_{H2O}) from MLS. The full line is an exponential fit (R² = 0.88, P < 3×10⁻¹⁰) to smoke layers peaking in water vapor concentration at a pressure less than 110 hPa. The half-life of the fit is 9.7±3.2 days.