1	Five satellite sensor study of the rapid decline of wildfire smoke in the				
2	stratosphere				
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10 Abstract

11

Smoke from western North American wildfires reached the stratosphere in large amounts in 12 August 2017. Limb-oriented satellite-based sensors are commonly used for studies of wildfire 13 aerosol injected into the stratosphere (OMPS-LP (Ozone Mapping and Profiler Suite Limb 14 Profiler) and SAGE III/ISS (Stratospheric Aerosol and Gas Experiment III on the International 15 Space Station)). We find that these methods are inadequate for studies the first 1 - 2 months after 16 such a strong fire event due to event termination ("saturation"). The nadir-viewing lidar CALIOP 17 (Cloud-Aerosol Lidar with Orthogonal Polarization) is less affected due to shorter path in the 18 smoke, and, further, provides means that we could use to develop a method to correct for strong 19 attenuation of the signal. After the initial phase, the aerosol optical depth (AOD) from OMPS-LP 20 and CALOP show very good agreement above the 380 K isentrope, whereas the OMPS-LP tends 21 to produce higher AOD than CALIOP in the lowermost stratosphere (LMS), probably due to 22 reduced sensitivity at altitudes below 17 km. Time series from CALIOP of attenuation-corrected 23 24 stratospheric AOD of wildfire smoke show an exponential decline during the first month after the fire, which coincides with highly significant changes in the wildfire aerosol optical properties. 25 The AOD decline is verified by the evolution of the smoke layer composition, comparing the 26 aerosol scattering ratio (CALIOP) to the water vapor concentration from MLS (Microwave Limb 27 Sounder). Initially the stratospheric wildfire smoke AOD is comparable with the most important 28 29 volcanic eruptions during the last 25 years. Wildfire aerosol declines much faster, 80 - 90% of 30 the AOD is removed with a half-life of approximately 10 days. We hypothesize that this dramatic decline is caused by photolytic loss. This process is rarely observed in the atmosphere. However, 31 32 in the stratosphere this process can be studied with practically no influence from wet deposition, in contrast to the troposphere where this is the main removal path of sub-micron aerosol particles. 33 34 Despite the loss, the aerosol particles from wildfire smoke in the stratosphere are relevant for the 35 climate.

36 37

1. Introduction

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39 Background stratospheric aerosol is composed of sulfuric acid, water, carbonaceous components,

40 and minor extraterrestrial and tropospheric components (Murphy et al., 2007; Kremser et al.,

- 41 2016; Martinsson et al., 2019). Volcanism is a strong source of the stratospheric sulfurous,
- 42 carbonaceous and ash aerosol (Martinsson et al., 2009; Andersson et al., 2013; Friberg et al.,
- 43 2014). Large eruptions, like that of Mt Pinatubo in 1991, affect the stratosphere for several years,
- 44 causing global cooling of several tenths of degrees Kelvin (Kremser et al., 2016). These eruptions
- 45 are scarce, only a few per century (Ammann et al., 2003; Stothers, 2007). Moderate eruptions are
- 46 more frequent contributors to the stratospheric aerosol (Vernier et al., 2011; Andersson et al.,
- 47 2015; Friberg et al., 2018), forming the persistently variable stratospheric background aerosol
- 48 (Solomon et al., 2011).
- 49

50 The stratospheric aerosol is also influenced by pyrocumulonimbus clouds (pyroCb) that form

51 during extreme weather conditions in connection with intense wildfires (Fromm et al., 2010). The

ongoing climate change is projected to increase the frequency of large wildfires (Kasischke et al., 52 2006; Dennison et al., 2014). Interestingly, the two largest events have, in terms of stratospheric 53 impact, occurred during the last few years, in North America 2017 (Peterson et al., 2018) and 54 Australia 2019-2020 (Kablick et al., 2020). Here we investigate the great pyroCbs formed in 55 southern British Columbia, Canada and northern Washington State, USA on August 12-13, 56 2017 (Fromm et al., 2021). Figure 1a shows an example of the strong impact on the stratospheric 57 aerosol of the 2019 Raikoke volcanic eruption, one of the strongest eruptions post Mt Pinatubo in 58 1991. In comparison, Figure 1b demonstrates the formidable early impact of wildfire aerosol. The 59 60 stratospheric impact of that fire has been described in terms of light-backscatter reaching unprecedentedly high values for a non-volcanic aerosol layer (Khaykin et al., 2018), light 61 extinction about 20 times higher than after the Pinatubo volcanic eruption in 1991 (Ansmann et 62 al., 2018), and mass of smoke comparable to that of a moderate sized volcanic eruption (Peterson 63 et al., 2018). The pyroCbs lifted smoke from the fire to the extratropical tropopause region, where 64

- absorption of radiation by black carbon (BC) in the smoke induced additional lift to 23 km
- altitude in 2 months (Yu et al., 2019; Lestrelin et al., 2021).
- 67

68 Smoke particles from wildfires contain a dominating fraction of organic matter by mass

69 (Garofalo et al., 2019). Organic aerosol is susceptible to photochemical loss (Jimenez et al.,

70 2009), and laboratory studies have demonstrated that this phenomenon could be an important

sink of secondary organic aerosol mass (Molina et al., 2004; Sareen et al., 2013). The residence

time of stratospheric air spans months to years depending on its path in the Brewer-Dobson

circulation (Engel et al., 2009; Bönisch et al., 2009). Due to very low probability of clouds, fine

aerosol particles have considerably longer residence times in the stratosphere than in the

r5 troposphere, which further emphasizes the importance of investigating photochemical loss in the

- r6 stratosphere (Martinsson et al., 2019).
- 77

78 The aim of this study is to further understand the stratospheric aerosol sources and its climate 79 impact. We develop methodology to correct for attenuation in dense smoke layers from wildfires 80 to properly deal with intense smoke injections into the stratosphere, with two main questions: 1) 81 does photochemical loss of wildfire smoke occur in the stratosphere, and 2) how does the AOD 82 of smoke from the wildfire studied here compare with volcanic aerosol?

83

84 The first decade of the 21st century was characterized by slower temperature evolution than

anticipated from CMIP5 models (Fyfe et al., 2016). The discrepancy was attributed to inter-

86 decadal Pacific oscillation (Medhaug et al., 2017), variations in solar forcing (Myhre et al., 2013)

and aerosol in the stratosphere from moderate volcanic eruptions (Santer et al., 2014). Should

88 wildfire smoke in the stratosphere be added to this list of phenomena that require more attention

- 89 in climate models?
- 90

91 Our investigation deals with the evolution of the wildfire AOD, and aerosol optical properties

92 obtained from the lidar CALIOP aboard the CALIPSO (Cloud-Aerosol Lidar and Infrared

- 93 Pathfinder Satellite Observation) satellite, OMPS-LP/Suomi and SAGE III/ISS in comparison
- 94 with volcanic injections to the stratosphere. Additionally, the water vapor concentrations of

95 individual smoke layers are investigated by the MLS, the spatial evolution of smoke layers is

96 investigated using OMPS-NM (Ozone Mapping and Profiler Suite Nadir Mapper), and the AODs

and extinction coefficients obtained from CALIOP are compared with that of OMPS-LP andSAGE III/ISS.

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2. Methods

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This study of the dense stratospheric smoke layers from pyro-cumulonimbus formed over
Western North America in August 12 – 13, 2017 is based on five satellite sensors. For four of
them, OMPS-LP, SAGE III/ISS, MLS and OMPS-NM, high level products (Level 2) are used.
The CALIOP data evaluation is based on a Level 1 product. A method to correct for attenuation
of the CALIOP laser beam in the smoke layers is presented. For these reasons CALIOP requires
more space in this section compared to the other methods.

108

109 2.1 CALIOP

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111 The evaluation of the CALIOP instrument carried by the CALIPSO satellite is based on version

4-10, level 1B data. CALIOP measures backscattering of laser light at two wavelengths, 532 and

113 1064 nm. For the shorter wavelength, scattered laser light is detected in parallel and

114 perpendicular polarizations relative to the outgoing beam. These almost nadir-viewing aerosol

and cloud measurements result in high resolution vertical profiles. For the altitude ranges <8.2,

116 8.2 - 20.2, 20.2 - 30.1 and 30.1 - 40 km the vertical resolutions are 30, 60, 180, and 300 m,

respectively. CALIPSO orbits between 82° S and 82° N, completing 14 - 15 orbits per day

118 (Winker et al., 2007; Winker et al., 2010).

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120 *2.1.1 AOD*

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Stratospheric AOD was obtained by integrating the backscattering intensity corrected for attenuation (described below) from the tropopause to 35 km altitude. Figure 1b illustrates how attenuation of the laser signal strongly reduced the signal below the dense smoke layer between 11 to 16 km altitude. We use the tropopause height according to MERRA-2 supplied with the

version 4.10 CALIOP data, which is a mixture of a dynamic and a thermal tropopause. The AOD

was averaged in the 20 - 80° N latitude range, where all nighttime swaths available from

128 CALIOP were included. The data were averaged over all longitudes in one-degree latitude bands,

- and these latitude bands were averaged for the $20 80^{\circ}$ N latitude range using area-weighting.
- For dense layers, the lidar ratios estimated for the individual smoke layers were applied(explained below). Apart from the first few days the lidar ratio shows no temporal evolution, it is

(explained below). Apart from the first few days the lidar ratio shows no temporal evolution, it is
 found to have geometrical mean of 48.9 sr with double-sided 95% confidence interval of 47.6 –

50.3 sr (Figure 2a), which is close to the typical background lidar ratio of 50 sr (Jäger and

- 134 Deshler, 2003). For layers that were not dense, the lidar ratio was held at this typical background
- 135 level. The volume depolarization ratio (δ_v) contains information that can be used to classify
- aerosol layers. When δ_v is less than 0.05 the data is considered background and the lidar ratio is
- 137 set to 50 sr (Vernier et al., 2009). Ice-clouds were removed in the lowest 3 km of the stratosphere
- by identifying them in stratospheric layers where the backscattering was high (attenuated

backscattering larger than 0.0025 km⁻¹ sr⁻¹). This limitation is introduced to avoid statistically 139 induced detection of ice clouds from weak signals. Data in these layers were classified as 140 probable ice clouds if their δ_v exceed 0.20, which classifies all the smoke layers in Figure 2c as 141 aerosol since the volume depolarization ratio always is smaller or equal to that of particles for a 142 depolarizing aerosol. The data within each swath were then clustered depending on their location. 143 Noise in the data led to some lone pixels within layers. These were reclassified depending on the 144 surrounding pixels, making sure that no single pixel marked as aerosol occurred within the ice-145 cloud layers. Layers of ice-clouds were then expanded upwards and horizontally to capture faint 146 edges of the clouds (Friberg et al., 2018). Aerosol with δ_v in the range 0.05 to 0.2 were 147 considered to be smoke, and $\delta_v < 0.05$ as background aerosol. In the present work the latter 148 discrimination had little effect because smoke was found to have the same lidar ratio as typical 149 background aerosol. The classification was carried out on data at 8 km resolution along each 150 swath with their highest vertical resolution (30, 60, or 180 m, depending on altitude), after which 151 the tropospheric data were removed. Possible polar stratospheric cloud (PSC) signals north of 152 45°N were excluded by classifying pixels with temperature below 195 K as possible PSC 153

occasions. Underlying pixels were also excluded, to prevent bias from attenuation of the lidar 154 signals or from settling ice-crystals (Friberg et al., 2018). 155

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2.1.2 Attenuation correction and radiative properties of individual smoke layers

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The evolution of the lidar, color and depolarization ratios were investigated using 32 separate 159 smoke layer measurements over the period 3 - 59 days after the fire. CALIOP has a statistical 160 disadvantage compared with lidars at the ground (Baars et al., 2019), because of small solid angle 161 due to long distance to the stratosphere (~700 km) and short measurement time. Optical 162 properties of old and faint individual smoke layers therefore could not be quantified with high 163 164 precision using CALIOP. The faint layers though still affect the AOD determinations described above, where AOD elevation after the fire remains approximately one year. Out of the 32 smoke 165 layers studied, 29 were night-time measurements, whereas the remaining three are defined as 166 167 day-time measurements. These latter ones increased the number of early observations (day 3 - 5)and were taken when the disturbance from solar radiation is small, i.e., shortly before the night. 168

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During the first weeks after the fire the smoke layers could be very dense with layer AODs 170

exceeding 1, causing strong attenuation of the CALIOP signals with two-way transmissions down 171

to below 0.01. Such high AOD values were also observed for this fire by the Deep Space Climate 172

Observatory/Earth Polychromatic Imaging Camera (DSCOVR-EPIC) mapper and the Aerosol 173

174 Robotic Network (AERONET) (Torres et al., 2020). For the 532 nm wavelength the particle lidar

ratio was estimated by aiming the scattering ratio (R; total-to-molecular backscattering ratio) 175

below a smoke layer to a target value. The target value was obtained from the background 176

scattering ratio beside each smoke layer investigated, which on average is R = 1.08, with standard 177

deviation ± 0.05 . To reduce influence from noise, the CALIOP data were averaged along the 178

swath. The averaging range varied between the smoke layers, due to its extension along the 179

swath, the homogeneity of the layer, and avoidance of sub-layer features. 180

The particle lidar ratio of an individual smoke layer was iterated until reaching the target value (R 182 = 1.08) described above from the combined effect of all altitude pixels. Pixels at altitudes outside 183 the smoke layer were set to the background lidar ratio of 50 sr (Jäger and Deshler, 2003). The 184 altitude resolution provided in the CALIOP data was used, where each altitude pixel (j) is 185 corrected for attenuation. The calculation starts at the highest altitude (40 km) and continues 186 downwards in two rounds. In the first round the star-marked quantities of equations 1-3 were 187 computed, correcting for attenuation from overlaying pixels. Before moving to the next altitude, 188 we account for self-attenuation from the pixel itself (equations to the right, without a star): 189 190

191
$$\beta_j^* = \frac{\beta_j'}{\prod_{k=1}^{j-1} T_k^2}; \qquad \beta_j = \frac{\beta_j^*}{\sqrt{T_j^{*2}}}$$
(1)

192

where β' is the attenuated backscattering and T² the two-way transmissions from both particles and molecules. The two-way particle transmission is obtained by first computing the AOD:

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$$AOD_j^* = (\beta_j^* - \beta_{m,j})S_p\Delta z_j; \qquad AOD_j = (\beta_j - \beta_{m,j})S_p\Delta z_j$$
(2)
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where Δz_j is the height of the altitude pixel, $\beta_{m,j}$ is backscattering from air molecules, and S_p the lidar ratio of the aerosol particles. The molecular lidar ratio, for computation of the molecular extinction, was set to 8.70477 sr (Prata et al., 2017). The two-way transmission of altitude pixel j due to the particles present is obtained from:

203
$$T_{p,j}^{*2} = \exp(-2AOD_j^*);$$
 $T_{p,j}^2 = \exp(-2AOD_j)$ (3)
204

These calculations in equations 1 - 3 are carried out until the background layer between altitudes a and b below the smoke layer reaches the target scattering ratio of 1.08 (Figure 3a):

208
$$R = \frac{\sum_{a}^{b} \beta_{j}}{\sum_{a}^{b} \beta_{m,j}}$$
(4)

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CALIOP measurements are affected by multiple scattering (Wandinger et al., 2010), causing 210 211 overestimation of the backscattering described by the multiple scattering factor (η). This factor is not known, previous estimates for fine aerosol are in the range 0.085 - 0.95 for layers thicker 212 than 500 m (Prata et al., 2017). In equation 2 the backscattering inflated by multiple scattering 213 (β) is multiplied by a lidar ratio. The latter is obtained by iterating equations 1-3 until reaching 214 the target scattering ratio, that of the surrounding air, below the cloud. Since the backscattering is 215 inflated by multiple scattering, the lidar ratio obtained will become the product of the actual lidar 216 ratio and the multiple scattering factor, i.e., the effective lidar ratio. Thus, while overestimating 217 218 the backscattering and underestimating the lidar ratio to equal multiplicative degree, the method applied here corrects the AOD for multiple scattering. 219 220

Error estimates of the effective lidar ratio were obtained by varying the target scattering ratio

- from its average value (R = 1.08) mentioned above, to its ± 0.05 standard deviation range. The
- fitting uncertainty in these estimates is strongly dependent on the light extinction in the smoke
- layer. Dense layers result in very small uncertainties in the effective lidar ratio because of the
- strong impact on R from a slight change in the extinction. Layers with lower extinction
- progressively increase the uncertainties of the estimate. When the error estimate of the effective lidar ratio fit exceeds 25% the result is excluded from the data analysis, which terminates
- estimates of lidar ratios from day 22 after the fire.
- 229

The color ratio, the ratio between the backscattering at 1064 nm to 532 nm wavelength, is

affected by a difference in attenuation of the two wavelengths. This is clearly visible for dense

smoke layers in the CALIOP browse images by a gradual increase of the color ratio through the

layer because of the weaker attenuation for 1064 nm wavelength than for 532 nm (Figure 1d).

Therefore, estimations of the attenuation were undertaken also for the long wavelength. The

molecular backscattering is assumed to be 1/16 of that at 532 nm $(1/\lambda^4$ dependence of Rayleigh

scattering). Weak molecular scattering at 1064 nm prohibits lidar ratio estimation at that
wavelength by CALIOP. Instead, the lidar ratio was assumed to be 60 sr, inducing uncertainties

in the color ratio. The volume color ratio is obtained from:

240
$$\chi = \sum_{k=\text{top}}^{\text{base}} \beta_{1064,k} / \sum_{k=\text{top}}^{\text{base}} \beta_{532,k}$$
(5)

241

To limit influence from attenuation in the color ratio computations, the estimates were based on the upper part of a smoke layer. Starting from the top of the smoke layer, the computations were truncated when the two-way transmission of the 532 nm wavelength fell below 0.7. Varying the 1064 nm wavelength lidar ratio in the wide range of 60 ± 20 sr the uncertainty in the color ratio becomes less than $\pm 5\%$ with this constraint applied. From the color ratio we define the particle color ratio:

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249
$$\chi_p = \sum_{k=\text{top}}^{\text{base}} (\beta_{1064,k} - \beta_{m,1064,k}) / \sum_{k=\text{top}}^{\text{base}} (\beta_{532,k} - \beta_{m,532,k}) = \frac{\chi R}{R-1} - \frac{1}{16(R-1)}$$
 (6)
250

where we made use of the wavelength dependence of Rayleigh scattering for molecular scattering, and the scattering ratio for the 532 nm wavelength was obtained from eqn. 4.

We also investigated the depolarization of the scattered laser beam at 532 nm by first forming the volume depolarization ratio:

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257
$$\delta_{\nu} = \sum_{k=\text{top}}^{\text{base}} \beta'_{532 \perp, k} / \sum_{k=\text{top}}^{\text{base}} \beta'_{532, k}$$
(7)

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where symbol \perp indicates scattered light polarized perpendicularly to the incident beam. Having access to the volume depolarization and an estimate of the molecular depolarization ratio $\delta_m \approx$ 0.003656 (Prata et al., 2017; Hostetler et al., 2006) the particle depolarization ratio is obtained from:

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$$\delta_p = \frac{\delta_v - \delta_m + \delta_v (1 + \delta_m)(R - 1)}{\delta_m - \delta_v + (1 + \delta_m)(R - 1)} \tag{8}$$

where R is obtained from eqn. 4.

268 2.2 Extinction coefficients and AOD from OMPS-LP

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270 The aerosol data from OMPS-LP (Chen et al., 2018; Jaross et al., 2014; Loughman et al., 2018) have lately been used extensively in the literature on volcanic and wildfire impact on the 271 stratospheric aerosol. Several data products are available, here we use the recently released Level 272 273 2 product: Suomi-NPP OMPS LP L2 AER Daily Product, version 2.0 (Taha et al., 2020). The 274 polar-orbiting Suomi satellite completes between 14 and 15 orbits per day. OMPS-LP is a limb-275 scattering method that collects data looking backwards along the satellite orbit, and along two 276 other directions separated by 4.25° from the orbit, giving a cross-track separation of approximately 250 km at the tangent point. Measurements are undertaken in the wavelength and 277 278 altitude ranges of 290 - 1000 nm and 10 - 80 km, respectively. The vertical resolution of OMPS-279 LP is 1.5 - 2 km (Rault and Loughman, 2013). The measurements are evaluated by the Gauss-Seidel limb scattering (GSLS) radiative transfer model. By improving calculations of the multiple 280 scattering source function, the total radiance error has become 1 - 3% (Loughman et al., 2015). 281 282 The aerosol product used here comprises 6 wavelengths (510, 500, 675, 745, 869 and 997 nm). The group responsible for the OMPS-LP version 2.0 data (Taha et al., 2020) recommends caution 283 when using data from altitudes below 17 km altitude due to loss of sensitivity. This problem can 284 285 be reduced by use of the 745 nm and longer wavelengths. Here we will make use of two of wavelengths: 745 nm because of the reduced problem with sensitivity, and 510 nm because it is 286 the wavelength closest to that of CALIOP (532 nm). 287 288 The OMPS-LP aerosol extinction coefficients are provided on a grid with a vertical resolution of

- The OMPS-LP aerosol extinction coefficients are provided on a grid with a vertical resolution of 1 km. To study the smoke from the August 2017 fire we compute the average AOD over all longitudes in the latitude interval $20 - 80^{\circ}$ N for three layers, the LMS (tropopause to 380 K
- isentrope), lower Brewer-Dobson branch (380 470 K) and the upper Brewer-Dobson branch
 (470 K to 35 km altitude). The OMPS-LP version 2 dataset use a cloud detection algorithm (Chen
- et al., 2016), and comes in two forms: one without filtering out signals from clouds, and the other where signals affected by clouds and polar stratospheric clouds are removed. In Figure 4 we
- show both these varieties for 745 nm wavelength, and, with and without flags regarding data
- quality including profile retrieval errors (named RetrievalFlags in the OMPS-LP files), high root-
- mean squares (ResidualFlags), and further errors from the South Atlantic anomaly, disturbances
 from the Moon, solar eclipses, planets, and satellite maneuvers (SwathLevelQualityFlags). In the
- two upper layers (Figures 4a and b) the differences are usually small between the varieties except
- for some spikes, whereas the LMS data (Figure 4c) show large stochastic variability as well as
- 302 periods of clear differences between the varieties. Since this data is taken well below 17 km
- altitude, sensitivity issues can be expected (Taha et al., 2020), see above. Days 130 190 (during
- 304 December 2017 to February 2018) several spikes appear in the two higher layers which probably

are caused by polar stratospheric clouds. The data set filtered for clouds and flagged stands out by
 comparably small peaks, whereas the differences between the varieties usually are small
 elsewhere. We therefore select the cloud-filtered and flagged data for further analysis in the
 coming sections.

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0 2.3 Extinction coefficients from SAGE III/ISS

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SAGE III/ISS is a limb-viewing instrument based on solar occultation. Here we make use of 312 313 Level 2 aerosol extinction coefficients (SAGE III/ISS User's Guide, 2018), version 5.10, supplied with a vertical resolution of 0.5 km. The upper limit of the slant path optical depth is 314 about 8, translating to a vertical optical depth of approximately 0.02 (SAGE III/ISS User's Guide, 315 2018). The orbiting of ISS differs markedly from the polar orbiting satellites CALIPSO 316 (CALIOP) and Soumi (OMPS-LP). This causes sporadic coverage by ISS of the latitudes of 317 interest here, resulting in that no average AODs over the $20 - 80^{\circ}$ N latitude range could be 318 formed with adequate time resolution. However, daily maximum extinction coefficients from 319 SAGE III/ISS could, when available, be included in a comparison with CALIOP and OMPS-LP. 320

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322 2.4 Water vapor measurements from MLS

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Water vapor concentrations (mixing ratio) in individual smoke layers was obtained from the MLS 324 325 instrument aboard the Aura satellite (Waters et al., 2006) in 12 vertical steps per decade of 326 pressure (version 5.0-1.0a, level 2). In nighttime measurements from days 6 – 59 after the fire, the smoke layers studied by CALIOP were also investigated with MLS in almost simultaneous 327 measurements, both instruments being on satellites that are members of the A-train (L'Ecuyer 328 329 and Jiang, 2010). Data in the 10 - 316 hPa atmospheric pressure range were used, with vertical 330 resolution 1.3 – 3.2 km (Lambert et al., 2020; Livesey et al., 2020). Limited vertical resolution 331 induces problems to obtain well defined observation of H₂O concentration of smoke layers close to the strong H₂O concentration gradient across the tropopause. H₂O from MLS for this fire have 332 previously been reported by Pumphrey et al. (2021). Close to the tropopause, but in the 333 stratosphere, no H₂O peak from a smoke layer can be detected. As the distance to the tropopause 334 increases, an H₂O peak from the smoke layer becomes discernible. Further up from the 335 tropopause, when the peak H₂O concentration is well above the extratropical tropopause at 336 atmospheric pressure of less than 110 hPa, a deep minimum appears between the tropopause 337 gradient and the peak from the smoke layer. All H₂O peaks were fitted with a Gaussian 338 distribution operating on logarithmic pressure and H₂O concentration to obtain estimates of the 339 peak concentration and the corresponding atmospheric pressure. To investigate a time 340 dependence in the smoke layer composition the peak H₂O concentration (C_{H2O}) was compared 341 with the attenuation-corrected aerosol scattering ratio (R) from CALIOP, the optical equivalent of 342 the mixing ratio, where the latter was obtained by forming the geometrical mean over 900 m 343 around the peak scattering ratio. The ratio of the of the two quantities (R/C_{H2O}) was formed, and 344 its dependence on time from the fire was studied. Out of the 13 smoke layers available with peak 345 water vapor concentrations above the altitude of 110 hPa atmospheric pressure, one was flagged 346 347 as low quality in the MLS data set, leaving 12 observations for the study of the R/C_{H2O} evolution.

349 2.5 UV aerosol index from OMPS-NM

350 The UV aerosol index of OMPS-NM based on measurements at two wavelengths, 340 and 378.5 351 nm, is the official NASA aerosol index product according to OMPS-NM (NMMIEAI-L2 V2.1.1) 352 release notes (Torres, 2019). For strongly UV absorbing aerosols, like black carbon from 353 wildfires, the UV aerosol indexes strongly increases with altitude (Herman et al., 1997). UV 354 aerosol index can be used to quantify AOD when layer altitude is available (Torres et al., 2020). 355 356 However, here the OMPS-NM UV aerosol index was used to map the geographical evolution of the smoke layers, that according to CALIOP measurements were distributed in both the 357 troposphere and the stratosphere. 358

359

3. Results

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Here we use an approach based on five satellite sensors to study the influence on the stratosphere 362 of the great North American fire in August 2017. We start by briefly describing results from the 363 method to correct CALIOP data for attenuation of the backscattered laser light. Then follows a 364 comparison of AODs obtained from OMPS-LP and CALIOP. Absorption aerosol index from 365 OMPS-NM is used to describe the dispersion of the wildfire aerosol in the stratosphere. To 366 explain differences in AOD between OMPS-LP and CALIOP, a comparison of extinction 367 coefficients follows, where results from SAGE III/ISS also are included in the comparison. The 368 369 evolution of the optical properties of the wildfire aerosol is then described, before the North American wildfire aerosol is compared with volcanic influence on the stratospheric AOD. 370 Finally, the fifth data set, water vapor from the MLS, is introduced in the discussion section, 371 372 where the evolution of the wildfire aerosol in the stratosphere is analyzed.

373

374 *3.1 Correction for attenuation*

375

The smoke layers usually were 1 - 3 km thick and could extend several degrees in longitude and 376 latitude. Measurements with the CALIOP lidar provide, in addition to short, nadir-viewing 377 measurement path in dense layers, the advantage that the signal is retrieved as a function of 378 position along the laser path with high resolution, which can be used to correct for attenuation of 379 the signal. Figure 3a shows the attenuated scattering ratio (R'; the measured backscattering 380 divided by the calculated molecular backscattering) from an example-smoke-layer measured on 381 August 16, 2017. The scattering ratio should be close to 1 in air layers with low aerosol 382 concentration, whereas values below 1 is caused by attenuation from particles. As can be seen in 383 Figure 3a, the attenuated scattering ratio first increases (starting from above the layer). Then the 384 signal decreases and reaches well below unity from 11 km altitude and downwards, i.e., well 385 below the scattering ratio of particle-free air. By techniques described in the Methods section we 386 correct for attenuation and fit the lidar ratio (the ratio of extinction to backscattering) (Figure 2a) 387 to obtain an estimate of the backscattering without attenuation, as illustrated by the scattering 388 ratio (R) in Figure 3a. 389 390

391 The evolution of wildfire aerosol from day 3 to 59 after the North American PyroCbs on August

- 12, 2017, is first investigated by comparing 32 smoke layers from individual CALIOP swaths.
- 393 The influence from attenuation is shown in Figure 3b. Clear deviation from the 1:1 line appears
- already at layer attenuated (uncorrected) AODs (AOD_{att}) of 0.12, and 50% reduction of the signal
- appears at layer AOD_{att} of approximately 0.25. Reduction by more than 50% appears until day 10 $\frac{1}{20}$
- after the fire, whereas those measurements close to the 1:1 line were taken after day 30. The
 AOD, i.e., the AOD corrected for attenuation, exceeds the AOD_{att} by more than a factor of 5 in
- 398 the densest layers of this study (Figure 3b).
- 399

400 *3.2 Comparison of CALIOP and OMPS-LP*

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To study the evolution of the stratospheric AOD, we form a 3-dimensional box in the

stratosphere extending over all longitudes in the $20 - 80^{\circ}$ N latitude range. In this box we use all

daily profiles, 14 - 15 CALIOP and 42 - 45 OMPS-LP, to form the average AOD. We apply the

405 method to correct CALIOP data for attenuation, as described in the Methods section. AODs are

406 computed for three layers, the LMS, the lower Brewer-Dobson branch, and the upper Brewer-407 Dobson branch, as shown in Figure 5.

408

409 When comparing AODs, the measurement wavelengths should be as close as possible, due to the

410 wavelength dependence of scattering. CALIOP AODs are shown for 532 nm wavelength, and the

411 OMPS-LP data are shown for the close wavelength of 510 nm. In addition, the 745 nm AODs

from OMPS-LP is shown. The response to the 2017 North American fire is weak in the upper

Brewer-Dobson branch (Figure 5a), whereas the two lower layers (Figures 5b-c) show clear
increase of the AOD. Comparing the two methods, they agree well in the upper Brewer-Dobson

- 414 Increase of the AOD. Comparing the two methods, they agree wen in the upper Brewer-Dooson 415 branch. In the lower Brewer-Dobson branch we see good agreement between the two methods,
- except for the first 1 2 months after the fire where much higher AODs are recorded by CALIOP
- 417 (Figure 5b). The latter is also true for the LMS, whereas the general agreement between the two

417 (Figure 50). The latter is also true for the LMS, whereas the general agreement between the two 418 methods is poor (Figure 5c). The OMPS-LP documentation advise against using data from below

- 419 approximately 17 km altitude, approximately the upper limit of the LMS, due to loss of
- sensitivity (Taha et al., 2020). We therefore do not perform any further comparisons in the LMS.
- 421 The stratosphere above the LMS (above the 380 K isentrope) shows good agreement between the

422 two methods, except for the first 1 - 2 months after the fire (Figure 5d).

423

424 3.3 Early evolution of the smoke layers

425

The daily AOD averages from CALIOP show large variability during the first days after the fire

because the lidar measures narrow curtains through the atmosphere, Figure 5e. The variability
remains until the smoke layers become sufficiently dispersed, allowing several daily

- 428 remains until the shoke layers become sufficiently dispersed, anowing several daily 429 measurements of the smoke layers. The nadir-viewing OMPS-NM provides UV (ultraviolet)
- 429 absorbing aerosol index, where strong signal for strongly UV light absorbing aerosol is obtained
- 431 in the upper troposphere and the stratosphere. Figure 6 shows the geographic evolution of the
- 431 In the upper troposphere and the stratosphere. Figure 6 shows the geographic evolution of the432 smoke layers from August 14 to 22, 2017 together with the orbits followed by the CALIOP
- 433 measurements. Up to August 16 the smoke is found in a rather confined area. From August 17 the

- 434 smoke layers are stretched in eastward direction, and after that the smoke spreads rapidly to the
- east. The dispersion gradually increases the number of daily CALIOP observations of the smoke.
- 436 This can also be seen in Figure 5e, where the variability in the daily AOD data becomes
- 437 successively smaller. From day 10 (August 22) we see a clear pattern of decline of the AOD.
- 438
- Figure 5e shows the total stratospheric AOD according to CALIOP from the tropopause to 35 km altitude. We see a strong decline of the stratospheric AOD the first 1.5 months after the fire, and a fitted exponential function has a half-life of 6.5 ± 0.9 days. Such a decline cannot be found in the
- 442 OMPS-LP AODs, which instead are increasing during the first month.
- 443

444 To further investigate this clear difference between the two methods, individual smoke layers are 445 investigated with respect to extinction coefficients. Figure 7a-d show the extinction coefficient of strong smoke layers from four days in August and September 2017. From CALIOP we show the 446 attenuated extinction coefficients as well as the profiles corrected for attenuation. Together with 447 the CALIOP data the OMPS-LP data closest by are shown. It is obvious that OMPS-LP shows 448 very much smaller reaction to the smoke layers than CALIOP. However, we cannot be sure that 449 the two instruments viewed the same airmasses in these four examples, because the two 450 instruments do not belong to the same satellite constellation. To remove that obstacle, the daily 451 maximum stratospheric extinction coefficient from OMPS-LP was extracted and compared with 452 32 selected profiles' peak extinction coefficients from CALIOP. SAGE III/ISS was also included 453 454 in the comparison from day 19 after the fire. Unfortunately, the orbiting of ISS did not permit 455 measurements of the fire studied here before that day. The very strong signals from CALIOP are not reflected in the OMPS-LP or SAGE III/ISS measurements, see Figure 7e. In part, this can be 456 explained by difference in vertical resolution, but as shown in Figures 7a-d, these high extinction 457 458 coefficients extend to broad vertical ranges that should allow detection of strong signals also by 459 OMPS-LP and SAGE III/ISS.

460

461 There is one principal difference between CALIOP on one hand and OMPS-LP and SAGE III/ISS on the other hand: whereas the former is nadir-viewing (vertical) the latter two methods 462 operate in limb orientation (horizontal). This is important, because the horizontal extension of 463 smoke layers is much larger, e.g., the smoke layer in Figure 1b has a vertical extension of 464 approximately 2 km, whereas the horizontal extension is approximately 700 km. The vertical, 465 two-way transmission to the CALIOP sensor through this layer is approximately 0.01, which we 466 correct for. The horizontal path through this layer is 350 times longer, implying that the one-way 467 limb transmission becomes 10^{-350} for the same wavelength. Even if the horizontal extension 468 would be just one tenth the transmission is still as low as 10^{-35} . Obviously, the radiation used for 469 detection in OMPS-LP and SAGE III/ISS is rapidly eliminated in such smoke layers. Therefore, 470 these two methods are inadequate for studies of dense aerosol layers. The upper limit in terms of 471 vertical AOD is estimated to 0.02 (SAGE III/ISS Users Guide, 2018), corresponding to the 472 extinction coefficient of 0.02 km⁻¹ for a 1 km thick layer. This problem is also acknowledged for 473 OMPS-LP (Chen et al., 2018; DeLand, 2019), and has been pointed out for other limb-oriented 474 satellite-based instruments (Fromm et al., 2014). Failure to properly handle this methodological 475 476 shortcoming could seriously affect attempts to verify results by modeling (Lurton et al., 2018).

477 Despite the clear limitation of OMPS-LP and SAGE III/ISS in this respect, the large body of

- information on wildfires is based on these methods, e.g., Bourassa et al., (2019), Das et al.,
- 479 (2021), Khaykin et al., (2020), Kloss et al., (2019), Torres et al., (2020) and Yu et al., (2019). By
- 480 comparing with CALIOP we here show that the limb-oriented techniques miss the dramatic
- events during the first 1 2 months after the fire. The rapid decline of the wildfire smoke will be further analyzed below.
- 483
- 484 3.4 Aerosol optical properties
- 485

486 To further investigate the unusual evolution of the AOD, we turn to the optical properties of the wildfire aerosol. The particle color and depolarization ratios are shown in Figure 2b and c. The 487 former is the ratio of backscattering at 1064 nm wavelength to that at 532 nm, where a smaller 488 color ratio indicates smaller particles, and the latter is the ratio of perpendicularly polarized to 489 total scattering at 532 nm, where a low ratio indicates particle shape close to spherical. To test the 490 significance in the evolution the data were temporally divided into two equal halves by number of 491 data points, and geometric averages were formed (black lines in Figure 2). The particle color ratio 492 shows a highly significant decrease comparing the first to the last half of the data points, whereas 493 the particle depolarization ratio increases with high significance. The change in the optical 494 properties takes place up to 15 - 30 days after the fire. This coincides with the decline of the 495 AOD, thus connecting a change of the aerosol properties to the AOD decline. 496

- 497
- 498 3.5 Stratospheric AOD variability caused by volcanism and wildfires
- 499

500 The stratospheric AOD varies considerably over time mainly due to influence from explosive volcanic eruptions as demonstrated in Figure 8, showing the period 2008 - 2018. In this time 501 502 span, nine volcanic eruptions clearly, but to varying degree, affected the stratospheric AOD. We also identify two cases of influence from wildfires, the Victoria fire (Australia, 2009) and the fire 503 studied here (Western North America, 2017). The residence time in the stratosphere varies from 504 several years for tropical injections into the upper layer representing the upper branch in the 505 Brewer-Dobson circulation (BD) (Figure 8a), the order of a year in the shallow branch of the BD 506 circulation (Figure 8b), to months in the LMS (Figure 8c) (Friberg et al., 2018). The sum of the 507 three layers is shown in Figure 8d. The volcanic eruptions in these 11 years mainly affected the 508 two lower stratospheric layers, only the Kelut eruption (2014) clearly reached to the deep BD 509 branch. Fire aerosol contains black carbon, which absorbs radiation, heats surrounding air and 510 induces lifting, as observed after the fire studied here (Khaykin et al., 2018; Yu et al., 2019). 511 512 After both fires, we see weak AOD elevation in the deep BD branch (Figure 8a), but for the fire studied here the two lower layers dominate the AOD, like most of the volcanic eruptions in the 513 eleven-year period. 514

515

516 Comparing the evolution of the AOD of the North American wildfire with the evolution of the

- aerosol from two of the most important volcanic eruptions during the last 25 years (Figure 9), we
- 518 find that the maximum stratospheric AOD after the fire is similar to that after the 2011 Nabro and
- 519 2009 Sarychev eruptions. During the first couple of months after volcanic events the AOD grows

520 due to formation of condensable sulfuric acid from the emitted volcanic gas sulfur dioxide. In

521 contrast, the wildfire aerosol displays a rapid decline during the first few weeks, before the AOD

stabilizes (Figure 9). This is followed by a period of rather stable AOD of more than 6 months,

523 before the AOD evolution turns to a slower decline towards background conditions, with similar 524 seasonality as the aerosol from the volcanic eruptions discussed (Figure 9). This latter decline is

525 mainly caused by springtime transport out from the stratosphere at mid and high latitudes

- 526
- 527 528

4. Discussion

(Bönisch et al., 2009; Martinsson et al., 2017).

529 530 The smoke aerosol is distributed both in the LMS and in the lower BD branch like aerosol from several volcanic eruptions (Figure 8). The rapid decline of the smoke aerosol during the first 531 month after the fire thus cannot be explained by transport out of the stratosphere. Measurements 532 with Raman lidars at three wavelengths indicate that the smoke from this North American fire 533 contain an accumulation mode but no coarse mode (Haarig et al., 2018; Hu et al., 2019). To leave 534 the extratropical stratosphere particles must pass through the LMS. The influence from 535 sedimentation on submicron diameter particles is small at that level of the stratosphere, e.g., for 536 0.6 µm diameter sulfuric acid/water particles the sedimentation velocity is 0.15 km/month, which 537 is slow compared to the large-scale transport down to the troposphere from the LMS (Martinsson 538 et al., 2005). Moreover, the change in the particle depolarization ratio (Figure 2c) indicates 539 change of the aerosol particle properties, and the particle color ratio decrease after the fire (Figure 540 541 2b) is the expected outcome for reduced particle sizes. Based on these arguments we turn the attention to loss of material from the aerosol particles to the gas phase to explain the rapid 542 543 decrease in AOD seen in Figure 5e.

544

545 Smoke layers contain water vapor that could induce hygroscopic growth/shrinkage. Water vapor 546 profiles for individual smoke layers from days 6-60 after the fire were obtained from the MLS. 547 Measurements close to the troppause (Figure 10a) are affected by a steep gradient in H_2O concentration. The profiles well above the gradient peaking at atmospheric pressure of less than 548 549 110 hPa are shown in Figure 10b. For the latter category the peak H₂O concentration is in the range 7 - 14 ppmv, implying a maximum H₂O vapor pressure of 0.16 Pa. For typical conditions 550 in the extratropics that vapor pressure corresponds to a relative humidity of a few percent or less 551 (Murphy and Koop, 2005). 552

553

To further investigate the smoke layers, the temporal evolution of the composition is studied by 554 forming the ratio of the mixing ratios of two components: aerosol backscattering and H₂O at the 555 peak of respective vertical distribution. As pointed out above, the strong H₂O gradient around the 556 tropopause affects the MLS measurements. But for the smoke layers higher up, peaking above 557 110 hPa, we find a rapid decrease in the aerosol scattering ratio compared with the H₂O 558 concentration (Figure 10c). Fitting an exponential function $\left(\frac{R}{C_{H2O}} = a + be^{-t/\tau}\right)$, the half-life 559 560 becomes 9.7±3.2 days, which is somewhat longer than that computed from the AOD (half-life 6.5±0.9 days). The rapid AOD decline (Figure 5e) is thus verified by relative concentrations of 561

aerosol and H₂O under well-controlled humidity conditions, whereas the low relative humidity
 rules out hygroscopic growth and influence from clouds as the explanation of the AOD decline.

564

The near-field wildfire aerosol contains, besides black carbon (Bond et al., 2013; Ditas et al., 565 2018), approximately 90% organic material (Garofalo et al., 2019). After emission, secondary 566 organic aerosol (SOA) is formed by oxidation of gas phase compounds (Shrivastava et al., 2017). 567 Knowledge of processes controlling formation and removal in the atmosphere is limited (Hodzic 568 et al., 2016). Global aerosol models usually remove SOA mainly by wet (90%) and, to a smaller 569 extent, by dry deposition (Tsigaridis et al., 2014). In contrast to the species dominating the 570 stratospheric aerosol and its precursor compounds during background conditions and volcanic 571 influence (sulfuric acid and sulfur dioxide), organic species are not the ultimate 572 573 thermodynamically stable compounds (Hallquist et al., 2009). Organic aerosol is an intermediate state on routes, with little known rates, from emitted compositions to the highly oxidized gaseous 574 products CO and CO₂ (Jimenez et al., 2009). Modeling and numerous laboratory studies find 575 evidence for photolytic removal rates of organic aerosol similar to that of wet deposition in the 576 troposphere (Hodzic et al., 2016; Zawadowics et al., 2020). Recently, photolytic removal of 577 particulate SOA was included in the Whole Atmosphere Community Climate Model (WACCM6) 578 (Gettelman et al., 2019). Hodzic et al. (2015) estimate the photolytic loss over a 10-day period to 579 50% for most organic species at mid tropospheric conditions. 580

581

These high rates are disputed by Yu et al. (2019), claiming a lifetime of 150 days (halflife 104 582 583 days) of organic aerosol from the fire studied here, whereas Das et al. (2021) explain a similar half-life of the same fire by large-scale circulation and particle sedimentation using OMPS-LP 584 and modeling. The experimental data used here cannot differentiate these two explanations, 585 although the slow part of the smoke decline is similar in seasonality to that of volcanic aerosol 586 587 (Figure 9) where photochemical loss is less important. The modeling study by Yu et al. (2019) was based on mimicking the extinction according to SAGE III/ISS at 1020 nm wavelength at 18 588 km altitude. For three reasons their study misses the strong decline of the AOD during the first 589 month. Firstly, because the orbiting of ISS prohibits studies of the wildfire smoke the first 19 590 days after the fire, secondly because of the time required to transport the wildfire aerosol to 18 591 km altitude is approximately one month (Yu et al., 2019) and thirdly because problems with 592 event termination ("saturation"), see Figure 7e. We therefore conclude that Yu et al. (2019) could 593 not observe the main decline of the aerosol taking place during the first 1 - 2 months after the 594 595 fire, see section 3.3 for further details.

596

597 Submicron aerosol particles have much longer residence time in the stratosphere than in the 598 troposphere due to sparsity of clouds, thus inhibiting the sink that traditionally is considered the 599 most important in the troposphere, i.e., wet deposition. This provides unique possibilities to study 600 photolytic loss without competition from other aerosol sinks. Interpreting the body of evidence 601 on the strong and rapid decline of the stratospheric AOD during the first month after the fire, we

- find that photolytic loss of organic aerosol is a highly likely explanation. The rate of photolytic
- 103 loss is likely better described by the evolution of R/C_{H2O} than by the AOD, because the latter
- 604 could to some degree be affected by transport across the tropopause. Our strong experimental

- evidence leads us to the hypothesis that the rapid decline of the wildfire aerosol in the
- stratosphere with a half-life of 10 days is caused by photochemical loss of organic material. This
- should be further investigated by modeling, but that is outside the scope of the present study.
- 608

To further put the strong early decline of wildfire aerosol into context, we compare the AOD

- during background conditions (years 2013 and 2014) with the year of the fire. When the
- 611 contribution of the exponential term is very small of the wildfire aerosol (after 7 half-lives), the
- background is approximately 2/3 of the wildfire AOD (Figure 9). Taking the background into
- account, the excess stratospheric aerosol due to the wildfire declines by 83% from the R/C_{H2O}
- value day 10 after the fire. The process starts before day 10, indicating that almost all the organicaerosol constituting approximately 90% of the near-field wildfire aerosol mass (Garofalo et al.,
- 616 2019) could be lost by photolysis. Residual wildfire aerosol particles, likely stripped off by a
- 617 large fraction of its original organic content, remain in the stratosphere up to approximately one
- 618 year (Figure 9).
- 619

Finally, we investigate the stratospheric aerosol load from the wildfire by comparing with the more studied volcanic impact (Table 1). The AOD growth, the average AOD over one year from the fire/eruption subtracted by the average background AOD (2013 - 2014), is approximately 1/4 and 1/3 of that of two of the most important volcanic eruptions for the stratospheric aerosol in the

- last 25 years (Sarychev 2009, Nabro 2011). The average excess aerosol during the year following the fire corresponds to a radiative forcing of -0.06 W m⁻² in the region 20 - 80° N, using standard
- 626 conversion as an approximation (Solomon et al., 2011).
- 627

628 Conclusions

629

630 In this study we investigate massive injections of smoke into the stratosphere from the August 631 2017 North American wildfires using five satellite sensors. Methodology was developed to correct CALIOP data for attenuation of the laser signal. The CALIOP AOD and extinction 632 coefficients were compared with OMPS-LP and SAGE III/ISS. From 1 - 2 months after the fire 633 we find that OMPS-LP and CALIOP AOD agree very well at altitudes above the 380 K 634 isentrope, where the former demonstrates high sensitivity with small statistical fluctuations. The 635 methods differ dramatically during the first 1-2 months after the fire when the smoke layers are 636 dense, because the long optical path through the smoke of the limb-oriented instruments OMPS-637 LP and SAGE III/ISS cause event termination ("saturation"). This is clearly demonstrated by the 638 low daily maximum extinction coefficients of the two instruments, being orders of magnitude 639 640 lower than the peak extinction coefficients of CALIOP. The nadir viewing CALIOP experiences a much shorter optical path, because the vertical extension of smoke layers usually are orders of 641 magnitude shorter than for limb orientation. We find that CALIOP is an indispensable tool for 642 studies of dense smoke layers entering the stratosphere after intense wildfires, providing signal 643 along the laser path that can be used to correct for attenuation. Once the smoke layers are 644 sufficiently thin, the limb technique OMPS-LP provide sensitive measurements of the AOD that 645 can be used together with CALIOP. 646

The AOD from the wildfire declines exponentially with a half-life of 6.5 days. This decline is 648 649 further studied by the evolution of the ratio of the aerosol and water vapor mixing ratios of the smoke layers, resulting in a massive decline of 80 - 90% of the wildfire aerosol with a half-life of 650 approximately 10 days. We find transport out of the stratosphere, sedimentation, influence from 651 clouds or hygroscopic growth/shrinkage to be highly unlikely explanations for the rapid decline 652 of wildfire aerosol in the stratosphere. Based on strong experimental evidence we hypothesize 653 that photochemical loss of organic aerosol causes the rapid decline, which would mean that 654 almost the entire organic fraction of the wildfire aerosol would be lost in the exponential decline. 655 656 The half-life according to this study agrees well with results from laboratory studies and global modeling. Our unique result could be obtained because of the long residence time of aerosol 657 particles in the stratosphere, whereas tropospheric studies of photochemical loss are extremely 658 difficult because it is masked by SOA formation and wet deposition due to short residence time. 659 The residual aerosol leaves the stratosphere within a year in the Brewer-Dobson circulation. 660 Despite the initial loss, the long-term effects of wildfire smoke on the stratospheric AOD and 661 radiative forcing are considerable. The ongoing climate change is projected to increase the 662 frequency of wildfires, prompting the need for inclusion of wildfire impact on the stratospheric 663 aerosol load in the climate models. 664

665

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- 674

675 Author Contributions

676

B.G.M. designed the study, designed methodology, undertook part of the data analysis, and wrote
most of the paper. J.F. contributed to the design of the study, designed methodology, did part of
the data analysis, and wrote parts of the text. O.S.S. contributed to the data analysis and M.K.S.
contributed to the design of methodology. In addition, all authors participated in discussions and
commented on the manuscript.

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683 Data availability

684

685 CALIOP V4.10 lidar data (https://search.earthdata.nasa.gov/search?fp=CALIPSO) are publicly
 686 available.

- 687 OMPS-NM UV aerosol index was obtained from the publicly available site
- 688 https://worldview.earthdata.nasa.gov/.
- 689 OMPS-LP stratospheric aerosol optical depths were obtained from
- 690 https://disc.gsfc.nasa.gov/datasets/OMPS_NPP_LP_L2_AER_DAILY_2/summary

691 MLS water vapor concentrations were obtained from 692 https://disc.gsfc.nasa.gov/datasets?page=1&keywords=ML2H2O 005 SAGE III/ISS aerosol data were obtained from 693 https://asdc.larc.nasa.gov/project/SAGE%20III-ISS/g3bssp 51. 694 695 696 **Competing Interest** 697 The authors declare no competing interests. 698 699 700 **Additional Information** 701 702 Correspondence and requests for materials should be addressed to B.G.M. 703 704 References 705 706 Ammann C.M., Meehl G.A., Washington W.M., and Zender C.S., A monthly and latitudinally varying volcanic forcing dataset in simulations of 20th century climate. Geophys Res. Lett., 707 30, 1567-1661, doi:10.1029/2003GL016875, 2003. 708 709 Andersson, S. M., Martinsson, B. G., Friberg, J., Brenninkmeijer, C. A. M., Rauthe-Schöch, A., 710 Hermann, M., van Velthoven, P. F. J., and Zahn, A., Composition and evolution of volcanic aerosol from eruptions of Kasatochi, Sarychev and Eyjafjallajökull in 2008 - 2010 based on 711 CARIBIC observations. Atmos. Chem. Phys. 13, 1781-1796, doi:10.5194/acp-13-1781-2013, 712 2013. 713 714 Andersson, S. M., Martinsson, B. G., Vernier, J. P., Friberg, J., Brenninkmeijer, C. A. M., Hermann, M., Van Velthoven, P.F. J., and Zahn, A, Significant radiative impact of volcanic 715 716 aerosol in the lowermost stratosphere, Nat. Commun. 6:7692 doi:10.1038/ncomms8692, 2015. 717 718 Ansmann, A., Baars, H., Chudnovsky, A., Mattis, I., Veselovskii, I., Haarig, M., Seifert, P., Engelmann, R., and Wandinger, U., Extreme levels of Canadian wildfire smoke in the 719 stratosphere over central Europe on 21-22 August 2017. Atmos. Chem. Phys. 18, 11831-720 11845, doi.org/105194/acp-18-11831-2018, 2018. 721 722 Baars, H., Ansmann, A., Ohneiser, K., Haarig, M., Engelmann, R., Althausen, D., Hanssen, I., 723 Gausa, M., Pietruczuk, A., Szkop, A., Stachlewska, I. S., Wang, D., Reichardt, J., Skupin, A., 724 Mattis, I., Trickl, T., Vogelmann, H., Navas-Guzmán, F., Haefele, A., Acheson, K., Ruth, A. A., Tatarov, B., Müller, D., Hu, Q., Podvin, T., Goloub, P., Veselovskii, I., Pietras, C., 725 Haeffelin, M., Fréville, P., Sicard, M., Comerón, A., Fernández García, A. J., Molero 726 Menéndez, F., Córdoba-Jabonero, C., Guerrero-Rascado, J. L., Alados-Arboledas, L., Bortoli, 727 D., Costa, M. J., Dionisi, D., Liberti, G. L., Wang, X., Sannino, A., Papagiannopoulos, N., 728 729 Boselli, A., Mona, L., D'Amico, G., Romano, S., Perrone, M. R., Belegante, L., Nicolae, D., Grigorov, I., Gialitaki, A., Amiridis, V., Soupiona, O., Papayannis, A., Mamouri, R.-E., 730 Nisantzi, A., Heese, B., Hofer, J., Schechner, Y. Y., Wandinger, U., and Pappalardo, G., The 731 unprecedented 2017-2018 stratospheric smoke event: Decay phase and aerosol properties 732 observed with the EARLINET. Atmos. Chem. Phys. 19, 15183-15198, doi.org/10.5194/acp-733 19-15183-2019, 2019. 734

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1015 Tables

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Table 1. Maximum and yearly average stratospheric AOD during background conditions andduring one year after the fire and after the two volcanic eruptions in Figure 9.

	Background	Background	Wildfire	Sarychev	Nabro
Year	2013	2014	2017	2009	2011
AOD max	0.009	0.009	0.020	0.028	0.017
AOD	0.0075	0.0074	0.0097	0.0169	0.0138
AOD growth ^a	-	-	0.0023	0.0095	0.0064
RF ^b	-	-	-0.06	-0.24	-0.16

aGrowth of AOD due to influence from wildfire/volcanism obtained by subtracting the average of 2013
and 2014 AOD.

^bRadiative forcing (W m⁻²) of the background-subtracted AOD.

1022



1029 Figure 1. CALIOP curtains of total attenuated backscatter (km⁻¹ sr⁻¹) at 532 nm from a) volcanic

aerosol layers in the stratosphere three days after the 2019 Raikoke eruption and b) a

1031 stratospheric smoke layer from the August 12, 2017, North American wildfire, where "aE-b" in

the scale refers to $a10^{-b}$. c) Volume depolarization ratio at 532 nm and d) attenuated color ratio

1033 (1064 to 532 nm) for the curtain in b). The white lines in a) and b) show the position of the

tropopause.



1037 Figure 2. Particle optical properties during the first 60 days after the fire. Black error bars show standard error and the double-sided 95% probability range of the geometric means. a) Particle 1038 lidar ratios for 532 nm where data points with fitting error exceeding 25% are discarded. The 1039 black line shows the geometric mean after day 4, and the full and dotted blue lines show the 1040 standard deviation and the double-sided 95% probability range of the distribution. b) Particle 1041 color ratio (1064 nm divided by 532 nm wavelength backscattering) with exponential fit ($R^2 =$ 1042 0.48, $P < 10^{-10}$), and c) particle depolarization ratio with exponential fit ($R^2 = 0.76$, $P < 10^{-10}$). 1043 The color and depolarization ratios were divided in two equal groups by number of observations 1044 to illustrate the highly significant changes with time of the optical properties, where the long and 1045 short error bars are the standard error and the double-sided 95% probability range of the 1046 1047 geometric means.

Figure 3. Illustration of methodology and its effect. a) The attenuated and corrected scattering ratios as a function of altitude. Example of methodology for one smoke layer, where the scattering ratio between 7.5 - 10 km altitude, below the smoke layer at 10.5 - 14 km, is targeted to a value of 1.08 (explained in the method section) by iteratively fitting the lidar ratio for 532 nm wavelength. b) The attenuated layer AOD (AOD_{att}) related to the layer AOD corrected for attenuation. The 1:1 relation is shown by the full line.

Figure 4. OMPS-LP layer AODs averaged over 20 to 80° North for 745 nm wavelength using
data filtered and not filtered from clouds and polar stratospheric clouds, and with and without
data flagged for data quality. Layer AOD for a) the upper Brewer-Dobson branch (470 K
isentrope – 35 km), b) the lower brewer-Dobson branch (380 – 470 K) and c) the LMS
(tropopause – 380 K) are shown.

Figure 5. AOD evolution of the stratospheric AOD (daily average) from 75 days before to 310 1064 days after the 2017 western North American fires. Comparisons of AOD from CALIOP (532 nm) 1065 with OMPS-LP (510 and 745 nm) with cloud filtering and flags activated for a) the upper 1066 Brewer-Dobson branch (470 K isentrope – 35 km, b) the lower Brewer-Dobson branch (380 – 1067 470 K) c) the LMS (tropopause – 380 K), d) from 380 K to 35 km (sum of layers in a and b) and 1068 e) the stratosphere of CALIOP from the tropopause to 35 km (sum of layers in a, b and c). The 1069 black, full line is an exponential fit ($R^2 = 0.79$, $P < 10^{-10}$) to the AOD over days 10 - 115 after the 1070 fire. The total stratospheric AOD half-life of the fit is 6.5 ± 0.9 days. 1071

Figure 6. Daily OMPS-NM aerosol absorbing index (UV) August 14 – 22, 2017 over all
longitudes and latitudes 20 - 80° N. This index is sensitive to UV absorbing aerosol particles in
the upper troposphere and the stratosphere, where signals from tropospheric aerosol declines
faster than from stratospheric due to short residence time. The yellow lines indicate nighttime
swaths of the CALIPSO satellite, and the faint lines show CALIPSO daytime swaths.

Figure 7. Extinction coefficients according to CALIOP, OMPS-LP and SAGE III/ISS in the 20 80° North latitude range during the first 60 days following the North American fire. a – d)
selected profiles (attenuated and corrected for attenuation) from CALIOP compared with closest
profiles according to OMPS-LP. e) Peak extinction coefficient from selected CALIOP profiles
compared with daily maximum extinction coefficients from OMPS-LP and SAGE III/ISS. Note
that SAGE III/ISS data are missing the first 19 days because of irregular coverage of the latitude

1085 range of interest.

Julo8 Jan09 Julo9 Jan10 Jul10 Jan11 Jul11 Jan12 Jul12 Jan13 Jul13 Jan14 Jul14 Jan15 Jul15 Jan16 Jul16 Jan17 Jul17 Jan18 Jul18

Figure 8. Zonally and eight-day moving average aerosol optical depth (AOD) of the stratosphere.
a - c) AOD in three layers obtained from CALIOP data (level 1B): a) 470 K potential temperature
to 35 km (deep Brewer-Dobson branch), b) 380 – 470 K (shallow Brewer-Dobson branch), c) the

1090 tropopause to 380 K (LMS). d) The total AOD from the tropopause to 35 km altitude. Volcanic

1091 eruptions marked by white triangles: Kasatochi (Ka), Sarychev (Sa), Merapi (Me), Grimsvötn

1092 (Gr), Puyehue-Cordón Caulle (Pu), Nabro (Na), Kelut (Ke), Calbuco (Ca), and Ambae (Am), and

1093 wildfires marked by orange circles: Victoria fire (Vi) and Western North American fires (Wn) at

time and latitude of eruption/fire. The AODs are corrected for attenuation.

Figure 9. Evolution of the AOD in the 20 - 80° N interval (8-day moving average) over two years: close to background conditions in the latitude interval studied (2013 - 2014), the year and the following year of the August 12, 2017, fire (2017 - 2018), and the same for two volcanic eruptions, the June 12, 2009, Sarychev (2009 - 2010) and June 12, 2011, Nabro (2011 - 2012) eruptions.

Figure 10. Water vapor in the smoke layer. Microwave Limb Sounder (MLS) measurements of water vapor concentrations (ppmv) Vs. atmospheric pressure for smoke layers a) close to the tropopause and b) well above the tropopause (atmospheric pressure < 110 hPa at the H₂O peak) for individual smoke layers available days 6 - 60 after the fire. c) The peak scattering ratio (R) according to CALIOP divided by the peak water vapor concentration (C_{H2O}) from MLS. The full line is an exponential fit (R² = 0.88, P < 3×10⁻¹⁰) to smoke layers peaking in water vapor concentration at a pressure less than 110 hPa. The half-life of the fit is 9.7±3.2 days.