





35 **Abstract**

36

37 A strong downslope wind event under fine weather condition on 13–15 February 2018 was  
38 examined by various observational and high resolution reanalysis datasets during the 2018 Winter  
39 Olympic and Paralympic games in Pyeongchang, Korea. High spatio-temporal resolution of wind  
40 information was obtained by Doppler lidars, automatic weather stations (AWS), wind profiler,  
41 and sounding observations under the International Collaborative Experiments for Pyeongchang  
42 2018 Olympic and Paralympic winter games (ICE-POP 2018). This study aimed to understand  
43 the possible generation mechanisms of localized strong wind event across high mountainous  
44 areas and in the lee side of mountains associated with the underlying large-scale pattern of a low-  
45 pressure system (LPS). The spatial distribution of linear trends for surface wind shows different  
46 patterns, exhibiting increased trend in the lee side and a persistent one in mountainous areas with  
47 the approaching LPS. Surface wind speed was intensified dramatically from  $\sim 3$  to  $\sim 12$   $\text{m s}^{-1}$  (gust  
48 was stronger than  $20$   $\text{m s}^{-1}$  above ground) at a surface station in the lee side (named as GWW).  
49 However, the mountainous station at DGW site appeared to have a persistently strong wind ( $\sim 10$   
50  $\text{m s}^{-1}$ ) during the research period. Budget analysis of horizontal momentum equation and local  
51 reanalysis data suggests that the pressure gradient force (PGF) derived by adiabatic warming  
52 along the downslope and subsequent hydraulic jump in the lee side of mountains was a main  
53 factor in the acceleration of the surface wind at the GWW site. Detailed analysis of the retrieved  
54 3D winds reveals that the PGF also dominate at the DGW site, which causes the persistent strong  
55 wind that is related to the channeling effect across the valley areas in the mountain range. The  
56 observational evidence presented here shows that the different mechanisms in local areas under  
57 the same synoptic condition with LPS are important references in determining the strength and  
58 persistence of the orographic-induced strong winds under fine weather condition.



## 59 **1 Introduction**

60 In mountain regions, wind is an important atmospheric phenomena as enhanced precipitation  
61 may be usually caused by the wind impinging topography (Medina et al., 2007; Yu and Cheng,  
62 2008; Panziera and Germann, 2010; Houze, 2012; Yu and Tsai, 2017; Tsai et al., 2018).  
63 Therefore, topography could significantly affect the behavior of winds to accelerate the wind  
64 speed or to change the wind direction. Such orographically strong wind and mountain waves can  
65 highly induce huge impacts on aviation operations (Clark et al., 2000; Kim and Chun, 2010, 2011;  
66 Kim et al., 2019; Park et al., 2016, 2019), outdoor sport activities, and forest wildfires in a  
67 relatively drier environment under the fine weather conditions (Smith et al., 2018). Downslope  
68 windstorm can produce strong wind in the lee side and plays an essential role in creating and  
69 maintaining the wildfires near the northern California with the easterly winds across Sierra  
70 Nevada and the southern Cascade Mountains (Mass and Ovens, 2019). Lee et al. (2020) also  
71 suggested that the downslope windstorm favored the wildfires along the northeastern coast of  
72 Korea with the westerly winds across the Taebek Mountain Range (TMR). The wind speed was  
73 also usually accelerated locally near the narrow valley or channel in between the mountains like  
74 the “gap wind” occurring along the strait of Juan de Fuca (Reed, 1931; Colle and Mass, 2000) in  
75 Washington, Columbia River Gorge in Oregon (Sharp, 2002), and Jangjeon area in South Korea  
76 (Lee et al., 2020).

77 Large-scale environmental conditions are key factors in determining the location where the  
78 downslope windstorm will be generated. The downslope windstorm usually occurs at the lee side,  
79 and the upstream prevailing wind direction is mostly perpendicular to the orientation of the  
80 mountain range. The elevated inversion-layer and the height of mean-state critical-level are also  
81 important references to evaluate the occurrence of downslope windstorm. The mechanisms of the  
82 downslope windstorm could usually be explained by hydraulic jump, partial reflection, and  
83 critical-level reflection from various numerical and theoretical studies in the past few decades



84 (Long, 1953; Houghton and Kasahara, 1968; Klemp and Lilly, 1975; Smith, 1985; Durran, 1990;  
85 Afanasyev and Peltier, 1998; Epifanio and Qian, 2008; Rögnvaldsson et al., 2011; Cao and  
86 Fovell, 2016). The combination of hydraulic jump and wave breaking can also enhance the  
87 downslope windstorm and increase the wind speed (Shestakova et al., 2018; Tollinger et al.,  
88 2019). The pressure gradient force (PGF) is one of the possible factors that accelerate the wind  
89 speed near the exit of gap between the mountains (Reed, 1931; Finnigan et al., 1994; Colle and  
90 Mass, 2000).

91 A few previous studies provided insightful explanations in numerical aspect about the  
92 development of the strong wind associated with the downslope windstorm along the northeastern  
93 coast of South Korea (in the lee side of the Taebaek Mountain Ranges; TMR). Most strong  
94 downslope wind events can be mainly explained by the three mechanisms in this region: hydraulic  
95 jump, partial reflection, and critical-level reflection (Lee, 2003; Kim and Cheong, 2006; Jang and  
96 Chun, 2008; Lee and In, 2009). The strong wind possibly occurs in any season with appropriate  
97 environmental conditions such as westerly and upstream inversion. Lee et al. (2020) confirms  
98 these conclusions with numerical modeling studies. Furthermore, they also found that PGF is one  
99 of the possible factors to cause the gap wind, and the variability of PGF is highly related to the  
100 localized topographic features. However, it still lacks sufficient observational evidence to  
101 examine the strong downslope wind events in detail because relatively dense wind observations  
102 from ground-based remote sensing techniques could not be easily collected under fine weather  
103 conditions.

104 Pyeongchang hosted the Winter Olympic and Paralympic Games in 2018 (most venues are  
105 located at coastal and higher altitude areas of the TMR). More detailed weather conditions and  
106 accurate prediction for several key parameters such as precipitation, visibility, wind directions,  
107 and wind speed are very important to ensure the safety of all game players and attendees. The  
108 Numerical Modeling Center (NWC), Korea Meteorological Administration (KMA) organized an  
109 intensive field experiment named the International Collaborative Experiments for Pyeongchang



110 2018 Olympic and Paralympic winter games, ICE-POP 2018  
111 ([http://155.230.157.230:8080/Icepop\\_2018/index.jsp](http://155.230.157.230:8080/Icepop_2018/index.jsp)). A very dense observational network was  
112 built to provide a good quality observational dataset in high temporal and spatial resolution either  
113 under precipitation or fine weather conditions. Many kinds of instruments were involved in the  
114 ICE-POP 2018 that allows the investigation of the nature of the strong wind event nearby  
115 mountainous areas in observational aspects.

116 Wider coverage of wind information can be usually obtained by the Doppler radar; however,  
117 regular meteorological Doppler radars cannot detect radial winds in fine weather conditions.  
118 Scanning Doppler lidar can be the best solution in obtaining a more complete wind information  
119 in such conditions even with finer resolutions. Using the Doppler lidar to document the  
120 downslope flow and rotors, Kühnlein et al. (2013) found that the transient internal hydraulic jump  
121 is characterized by turbulence. Menke et al. (2019) identified the recirculation zone over a  
122 complex terrain using six scanning Doppler lidars. The interactions between the winds and terrain  
123 essentially affected the occurrence of flow recirculation. However, only radial winds were used  
124 resulting in incomplete wind observations that can only provide very limited information for the  
125 realistic airflow structures. Complete 3D wind fields could be retrieved from 4DVAR (4D-  
126 Variational Assimilation) using Doppler lidar. The accuracy of wind speed, direction and water  
127 vapor flux will be improved when assimilating the lidar data (Kawabata, 2014). Thus, lidar  
128 observations can indeed provide a good quality of 3D wind information under fine weather  
129 condition.

130 The objective of this study is to use high spatiotemporal resolution of 3D winds and  
131 observational data to investigate the fine-scale structural evolution of strong downslope winds  
132 over the complex terrain in northeastern part of South Korea (i.e., in the Pyeongchang area )  
133 during the period 13–15 February 2018. Multiple Doppler lidars, automatic weather station  
134 (AWS), wind profiler, sounding observations, and local reanalysis (LDAPS: Local Data  
135 Assimilation and Prediction System) datasets were adopted to derive detailed 3D wind fields over



136 TMR and northeastern coastal regions through the WISSDOM (WInd Synthesis System using  
137 Doppler Measurements, Liou and Chang, 2009; Tsai et al., 2018) synthesis. In particular, this  
138 study tries to recognize the mechanisms of the strong wind over the mountainous area and in the  
139 lee side areas of TMR while a LPS passing through the northern side of the Korean peninsula. A  
140 strong downslope wind case was selected for the further analysis in this study not only because  
141 the Olympic games were interrupted due to the strong wind but also because dense observations  
142 are available during the ICE-POP 2018. Interactions between the complex terrain and large-scale  
143 airflow of the LPS could be adequately explored in this study. Furthermore, three scanning  
144 Doppler lidars were established in this area, which provided more sufficient wind information  
145 under fine weather condition.

146

## 147 **2 Data and methodology**

### 148 **2.1 Scanning Doppler lidar**

149 Two different models of scanning Doppler lidars were adopted in this study: (1) “WINDEX-  
150 2000” produced by the manufacturer Laser Systems and (2) the “Stream Line” produced by the  
151 manufacturer HALO Photonics. The scanning Doppler lidar can measure the radial Doppler  
152 velocity by detecting atmospheric aerosols and dusts via laser (class 1M) in a very high spatial  
153 resolution. The radial winds were sufficiently observed by adjustable scanning strategy in three  
154 modes: Plan Position Indicator (PPI), Range Height Indicator (RHI), and Zenith pointing (ZP).  
155 Furthermore, the complete wind information could be constructed using these lidar observations  
156 under fine weather condition via WISSDOM.

157 The WINDEX-2000 lidar operated a full volume scan every ~27 min in which seven PPIs  
158 (the elevation angles of 5°, 7°, 10°, 15°, 30°, 45°, and 80°) and one hemispheric RHI (the azimuth  
159 angle of 0°, that is, started from the north). There are 344 gates along a radar radial direction with



160 360 azimuth angles between  $0^\circ$  and  $360^\circ$ . The gate spacing is 40 m and the maximum observed  
161 radius distance is  $\sim 13$  km. The Stream Line lidar operated a full volume scan every  $\sim 13$  min in  
162 which five PPIs ( $7^\circ$ ,  $15^\circ$ ,  $30^\circ$ ,  $45^\circ$ , and  $80^\circ$  before 10:00 UTC on 14 Feb. 2018 and  $4^\circ$ ,  $8^\circ$ ,  $14^\circ$ ,  $25^\circ$ ,  
163 and  $80^\circ$  after 10:00 UTC) and two hemispheric RHIs (the azimuth angle of  $51^\circ$  and  $330^\circ$ ). There  
164 are 1660 gates along the 360 radar beams with azimuth angles between  $0^\circ$  and  $360^\circ$ . The gate  
165 spacing is 60 m and the maximum observed radius distance can reach to  $\sim 100$  km.

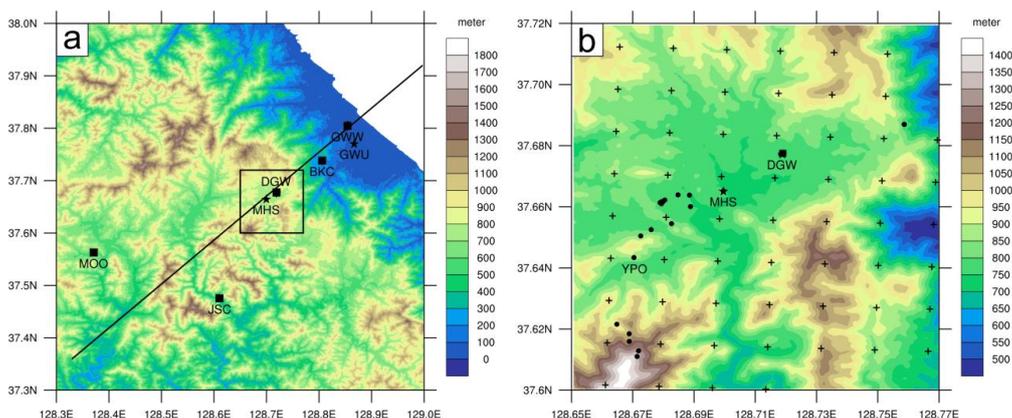
166 Quality control (QC) of the radial winds (in PPI and RHI modes) was done by applying  
167 SNR (signal noise ratio) threshold in advance. To obtain correct and useful measurements, QC is  
168 necessary for each lidar observation, where the non-meteorological echoes are removed as the  
169 threshold of SNR is smaller than 0.04.

## 170 **2.2 Surface weather station, sounding, wind profiler and LDAPS**

171 Figure 1 shows an observational domain and domain of mountain clusters during the ICE-  
172 POP 2018. AWS mainly provides the surface wind speed, wind direction, pressure, and  
173 temperature in high spatiotemporal resolution (1-min interval and  $\sim 10$  km distance between each  
174 AWS station). In addition to the regular KMA AWS, additional AWS sites were deployed in the  
175 mountainous area during the ICE-POP 2018. Thus, dense AWS network (black dots in Fig. 1b)  
176 can be utilized to document detailed evolutions of surface parameters and could also be  
177 observational inputs for WISSDOM (the details is Section 2.3). Five soundings are launched  
178 every 3 hours and the sounding sites are all located in the studied domain near the northeastern  
179 part of South Korea (black square in Fig. 1a). Such dense sounding observations provide  
180 environmental winds, temperature profiles, and stability information in a very fine-scale across  
181 the mountainous and coastal areas. The wind profiler is deployed at the GWW site to measure  
182 the winds in case of lacking sounding observations. In addition, the high temporal resolution of  
183 wind profiler measurements (10-min interval) could potentially be reference for the surface and



184 retrieved winds.



185

186 **Figure 1.** (a) Observations used in this study and the topographic features from the DEM (digital elevation model)  
187 in the northeastern Korea (corresponding to the bigger box in the Fig. 3) and (b) WISSDOM synthesis domain  
188 adopted in this study corresponding to the box in Fig. 1a. The locations of the scanning Doppler lidar sites are denoted  
189 by asterisks. The locations of the sounding sites are denoted by squares. Note that the sounding and lidar observations  
190 are both operated at DGW site and a wind profiler is located at GWW site. The locations of the AWS sites and  
191 LDAPS grids are denoted by dots and plus symbols, respectively.

192 The LDAPS is a 3DVAR numerical weather prediction (NWP) product generated from  
193 KMA with a spatial (temporal) resolution of  $\sim 1.5$  km (3 hours) with 70 levels in vertical. High  
194 resolution 3D wind fields from LDAPS will be one of the “observational” inputs in WISSDOM.  
195 This dataset is freely available from the KMA web site (<https://data.kma.go.kr>). Note that a  
196 relatively weaker weighting of the LDAPS input will be set in WISSDOM because this study  
197 prefers to emphasize the contribution from the observations.

### 198 2.3 WInd Synthesis System using DOppler Measurements (WISSDOM)

199 WISSDOM was originally developed by Liou and Chang (2009) and has been applied in the  
200 Pyeongchang area (Tsai et al., 2018). This study adopted a newly improved version, which  
201 includes more observations with more constraints compared with a previous version. In new  
202 version of WISSDOM, the following cost function [eq. (1)] was minimized by using



203 mathematical variational-based method at the retrieval time,

$$204 \quad J = \sum_{M=1}^8 J_M. \quad (1)$$

205 This cost function comprised of eight constraints, and the 3D wind fields were obtained by  
206 variationally adjusting solutions to simultaneously satisfy those constraints at the same time. The  
207 first constraint is the geometric relation between the radial velocity ( $V_r$ ) observations from  
208 multiple lidars and Cartesian winds  $V_t = (u_t, v_t, w_t)$ , control variables, defined as

$$209 \quad J_1 = \sum_{t=1}^2 \sum_{x,y,z} \sum_{i=1}^N \alpha_{1,i} (T_{1,i,t})^2, \quad (2a)$$

$$210 \quad T_{1,i,t} = (V_r)_{i,t} - \frac{(x - P_x^i)}{r_i} u_t - \frac{(y - P_y^i)}{r_i} v_t - \frac{(z - P_z^i)}{r_i} (w_t - W_{T,t}), \text{ and} \quad (2b)$$

$$211 \quad r_i = \sqrt{(x - P_x^i)^2 + (y - P_y^i)^2 + (z - P_z^i)^2}. \quad (2c)$$

212 Any numbers of lidar [subscripts  $i$  in eq. (2a)] can be applied to this constraint at two time levels  
213 (subscripts  $t$ ).  $\alpha_1$  in eq. (2a) is the weighting coefficient corresponding to  $J_1$  (same as the  
214 following equations for  $J_2 - J_8$ ). The subscripts  $i$  and  $t$  in the  $(V_r)_{i,t}$  represent the radial  
215 velocity observed by the  $i$ -th lidar, and  $(u_t, v_t, w_t)$  indicate the 3D wind at location  $(x, y, z)$ ,  
216 and the terminal velocity ( $W_{T,t}$ ) of particles were estimated by radar reflectivity at two time  
217 levels.  $(P_x^i, P_y^i, P_z^i)$  are the coordinates of the  $i$ -th lidar, and the distance between each grid point  
218 and the  $i$ -th lidar were denoted by  $r_i$ . Note that  $W_{T,t}$  will be zero when there is no radar  
219 reflectivity, or the terminal velocity is possibly negligible under fine weather condition.

220 Next constraint is the difference between the  $\mathbf{V}_t$  and the background winds ( $\mathbf{V}_{B,t}$ ) defined  
221 in eq. (3);

$$222 \quad J_2 = \sum_{t=1}^2 \sum_{x,y,z} \alpha_2 (\mathbf{V}_t - \mathbf{V}_{B,t})^2, \quad (3)$$

223 the sounding observations are used to be the background winds in eq. (3). The constraint of the  
224 anelastic continuity equation is



$$J_3 = \sum_{t=1}^2 \sum_{x,y,z} \alpha_3 \left[ \frac{\partial(\rho_0 u_t)}{\partial x} + \frac{\partial(\rho_0 v_t)}{\partial y} + \frac{\partial(\rho_0 w_t)}{\partial z} \right]^2, \quad (4)$$

where  $\rho_0$  is the air density. The fourth constraint was deduced from the vertical vorticity equation given by

$$J_4 = \sum_{x,y,z} \alpha_4 \left\{ \frac{\partial \xi}{\partial t} + \left[ u \frac{\partial \xi}{\partial x} + v \frac{\partial \xi}{\partial y} + w \frac{\partial \xi}{\partial z} + (\xi + f) \left( \frac{\partial u}{\partial x} + \frac{\partial v}{\partial y} \right) + \left( \frac{\partial w}{\partial x} \frac{\partial v}{\partial y} - \frac{\partial w}{\partial y} \frac{\partial u}{\partial z} \right) \right] \right\}^2, \quad (5)$$

where  $f$  indicate the Coriolis parameter and the meaning of overbar in eq. (5) is the temporal average of the two time levels. The constraint about Laplacian smoothing filter:

$$J_5 = \sum_{t=1}^2 \sum_{x,y,z} \alpha_5 [\nabla^2(u_t + v_t + w_t)]^2. \quad (6)$$

The horizontal winds observed by the sounding, AWS and LDAPS, can be interpolated to each given grid in the WISSDOM synthesis domain. The sixth constraint is the difference between the  $\mathbf{V}_t$  and the sounding observations ( $\mathbf{V}_{S,t}$ ), as defined in eq. (7);

$$J_6 = \sum_{t=1}^2 \sum_{x,y,z} \alpha_6 (V_t - V_{S,t})^2. \quad (7)$$

The seventh constraint represents the discrepancy between the surface winds and AWS ( $\mathbf{V}_{A,t}$ ), as expressed in eq. (8);

$$J_7 = \sum_{t=1}^2 \sum_{x,y,z} \alpha_7 (V_t - V_{A,t})^2. \quad (8)$$

Finally, the eighth constraint measures the squared errors between the horizontal winds and the LDAPS ( $\mathbf{V}_{L,t}$ ), as defined in eq. (9);

$$J_8 = \sum_{t=1}^2 \sum_{x,y,z} \alpha_8 (V_t - V_{L,t})^2. \quad (9)$$

Original version of WISSDOM is just only included first five constraints, it had already applied to synthesize high-quality 3D wind field in some of previous studies (Liou and Chang, 2009;

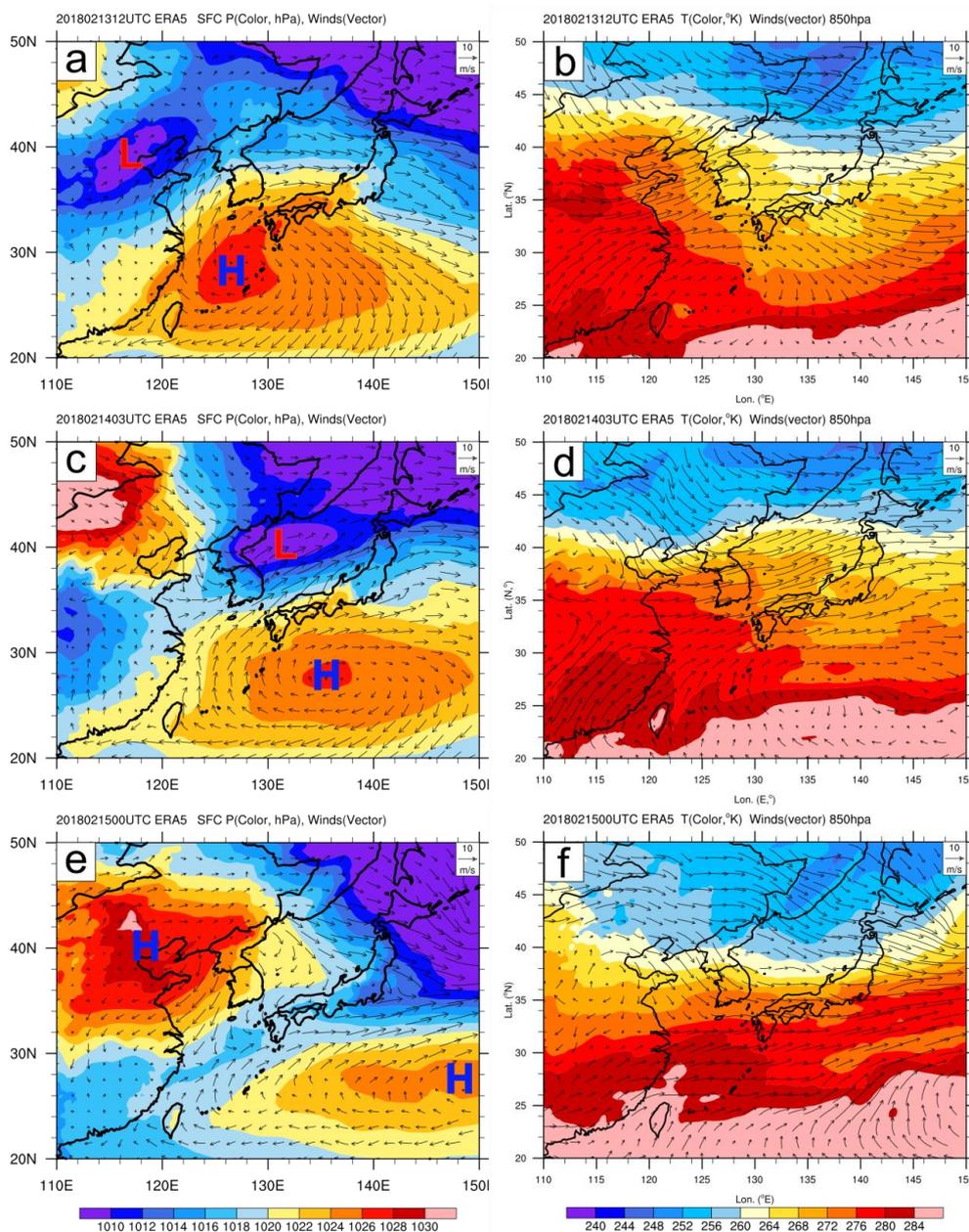


244 Liou et al., 2012, 2013, 2014, 2016; Lee et al., 2017). The primary advantages and additional  
245 details of WISSDOM can be referred to Tsai et al. (2018). The main improvement of the new  
246 version of WISSDOM is that all available wind observations are considered to be one of the  
247 constraints to minimize the cost function. In addition, this new version extends its applicability  
248 by including multiple-lidar observations and thus, realistic wind fields would be retrieved under  
249 fine weather condition.

### 250 **3 Overview of the strong wind case**

#### 251 **3.1 Synoptic condition**

252 Hourly ERA5 dataset was used herein to document the synoptic condition. The ERA5 data is  
253 an atmospheric reanalysis of the globe climate and was generated by the European Centre for  
254 Medium-Range Weather Forecasts (ECMWF, DOI: 10.24381/cds.adbb2d47). In the beginning  
255 of the research period at 12:00 UTC on 13 February 2018, a high-pressure system (HPS) was  
256 located in the southernmost of the Korean peninsula (as Fig. 2a). Surface southwesterly winds  
257 were dominant from the Yellow sea to the western coast of South Korea associated with anti-  
258 cyclonic circulation of the HPS. This stronger southwesterly wind was also related to the cyclonic  
259 circulation of a low-pressure system (LPS) centered at 39°N, 117°E near Beijing, China.  
260 Compared to the western coast, relatively weak wind existed over the land and eastern coast of  
261 Korea. The westerly wind came from China accompanying warm air at a higher layer (850 hPa,  
262 Fig. 2b). This veering wind also indicated that the prevailing southwesterly wind was dominated  
263 by the warm advection. Thus, the temperature gradient existed between the land, western and  
264 eastern coast (exceeding  $\sim 4^{\circ}\text{K}$  difference).



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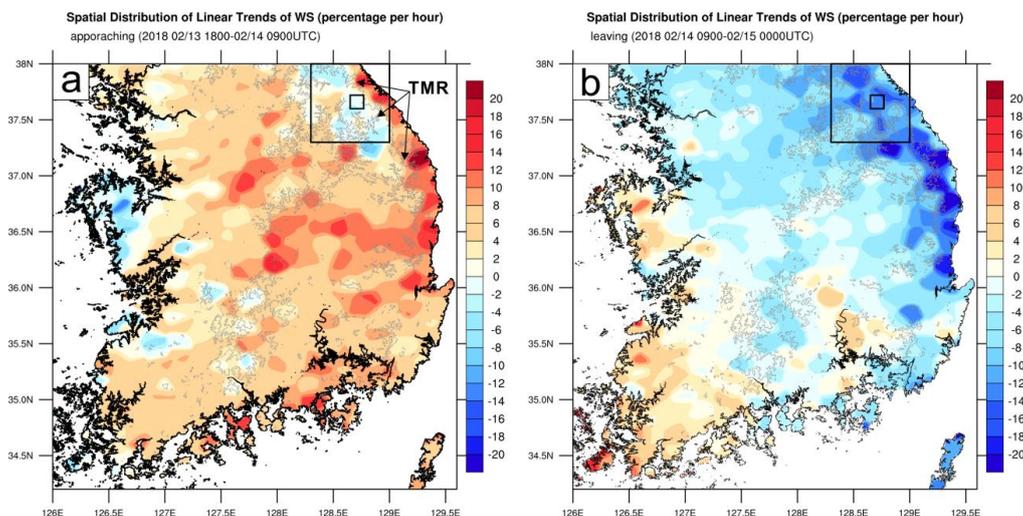
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**Figure 2.** (Left panels) Horizontal winds (vectors) and surface pressure (hPa, color shading) at the surface level obtained from the ERA5 reanalysis dataset at (a) 12:00 UTC on 13 Feb., (c) 03:00 UTC on 14 Feb. and (e) 00:00 UTC on 15 Feb. 2018. (Right panels) Horizontal winds (vectors) and temperature (°K, color shading) at 850 hPa level from the same reanalysis at (b) 12:00 UTC on 13 Feb., (d) 03:00 UTC on 14 Feb. and (f) 00:00 UTC on 15 Feb. 2018. The locations of low (high) pressure systems are marked by “L” (“H”).



271           Consequently, the LPS and HPS were both moving eastward. The surface wind became  
272 stronger and turned to westerly over the Korean peninsula associated with the confluences  
273 between these two systems (Fig. 2c). Horizontal pressure gradient was intensified along the  
274 northeastern coast of Korea as the LPS moved to the East Sea. A relatively lower temperature  
275 was detected over the mountain area (i.e., near northeastern coast of South Korea) even when the  
276 warm advection was approaching Korea (Fig. 2d). The other HPS came from China at around  
277 00:00 UTC on 15 February 2018 and the environmental winds surrounding Korea were switched  
278 to relatively weak northerly winds. Relatively weaker pressure gradient and smaller temperature  
279 differences between the western and eastern coasts were shown in Figs. 2e and 2f. No  
280 precipitation existed along the northeastern coast of South Korea according to the AWS  
281 observations during the research period (not show).

282           The research period can be separated into two stages based on the synoptic conditions as the  
283 LPS was approaching (between 18:00 UTC on 13 February and 09:00 UTC on 14 February 2018)  
284 and leaving (09:00 UTC on 14 February and 00:00 UTC on 15 February 2018) from Korea.  
285 Spatial distribution of linear trends (percentage  $\text{h}^{-1}$ , Chou et al., 2013) of surface wind speed for  
286 these two stages were shown in Fig. 3. Generally, the surface wind speed was increased by ~2–8  
287 percentage  $\text{h}^{-1}$  while the LPS was approaching Korea and higher increasing trend (over than 10  
288 percentage  $\text{h}^{-1}$ ) was related to the upstream of mountain ranges and the east coast (lee side of the  
289 TMR) in Korea (Fig. 3a). In addition, the strong positive spatial gradient appeared from the TMR  
290 to the east coast along with negative trends over the mountain ranges as described in the Fig. 3a.  
291 That is, the surface wind speed slightly decreased or nearly constant ( $\sim 10 \text{ m s}^{-1}$ ) over the TMR  
292 while the LPS was approaching. The surface wind speed dramatically increased with the  
293 approaching LPS. When the LPS was moving away from the east coast, the negative trend is  
294 dramatic, in particular, over the TMR and east coastal regions (Fig. 3b).



295  
296 **Figure 3.** (a) Spatial distribution of linear trends of surface wind speed (unit: percentage  $h^{-1}$ , color shading)  
297 calculated from the AWS (automatic weather station) observations during the low-pressure system (LPS)  
298 approaching from 18:00 UTC on 13 to 09:00 UTC on 14 Feb. 2018. (b) as in (a) but for the time period when the  
299 LPS was leaving from 09:00 UTC on 14 to 00:00 UTC on 15 Feb. 2018. The topographic height thresholds of 600  
300 m MSL (mean sea level) is indicated by gray lines and the arrows marked the location of the TMR.

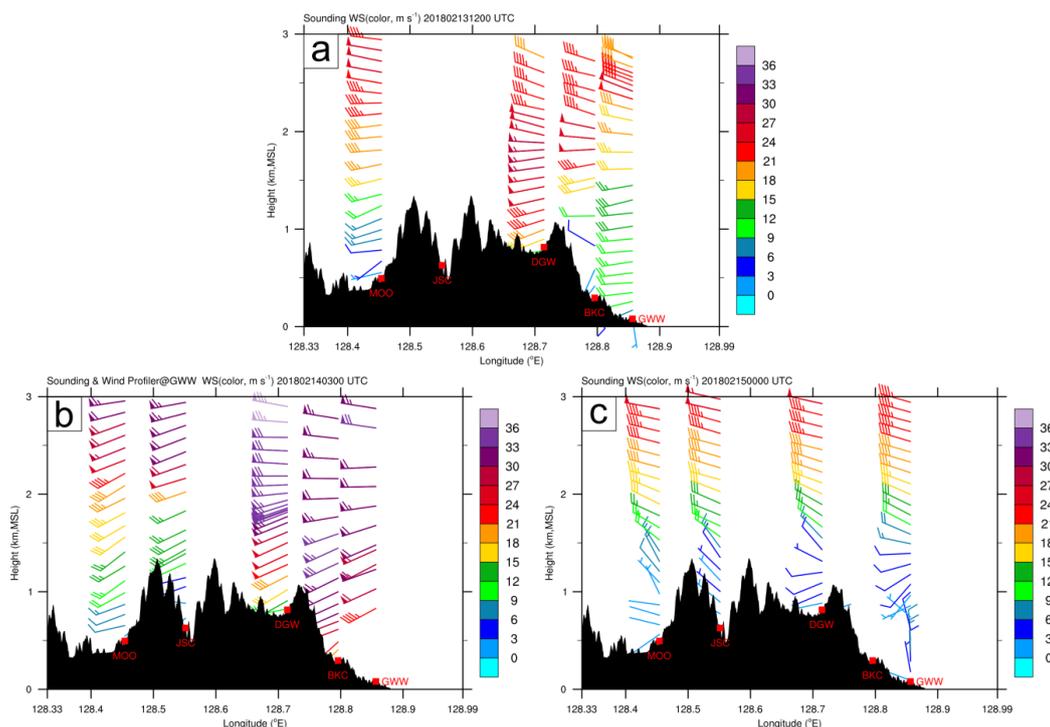
### 301 3.2. Environmental conditions in the local area near northeastern Korea

302 Because the spatial distribution of linear trends of surface wind speed revealed obvious  
303 gradient just near the mountainous area and in the lee side of the TMR, two domains were selected  
304 in this study that are shown as boxes in Fig. 3. All available intensive observations are also  
305 marked in Figs. 1a and 1b, corresponding to the bigger and smaller boxes in Fig. 3. Three  
306 scanning Doppler lidars were deployed at DGW, GWU (Stream Line), and MHS (WINDEX-  
307 2000), indicated by asterisks in Fig. 1a. Five sounding stations aligned from the mountainous  
308 area to the coastal area (i.e., perpendicular with the orientation of TMR). The sounding stations  
309 MOD and JSC were located at the southwestern TMR with gentle slope and the station DGW  
310 was the closest site to most outdoor venues of the Olympic games near the crest of TMR. Station  
311 BKC and GWW were located at the northeastern slope of TMR and the coast area, respectively  
312 (as Fig. 1a). In addition, a wind profiler is located at the lee side (GWW). The WISSDOM



313 synthesis domain was set over the mountainous area with a horizontal spatial coverage of  $12 \times$   
314  $12 \text{ km}^2$  as shown in Fig. 1b. The horizontal and vertical grid size were both set to 50 m, and  
315 vertical extension is from 0- km to 3- km height MSL (mean sea level). Additional AWS stations  
316 were deployed around the venues (black dots in Fig. 1b) during the ICE-POP 2018. Furthermore,  
317 the location of the LDAPS reanalysis dataset was also marked by the plus symbols.

318 Figure. 4 shows the variations in the environmental winds observed by the soundings and/or  
319 wind profiler along the crossline (black line in Fig. 1a) from the mountainous area to the lee side.  
320 Instead of lacking sounding observations at GWW site, the wind profiler observation is used to  
321 provide the wind information near the coastal area when the LPS was passing Korea. In the  
322 beginning of the research period, prevailing westerly winds are dominant at all sounding sites  
323 (Fig. 4a). However, a relatively stronger wind was only measured in heights below  $\sim 1.5 \text{ km}$   
324 at the DGW site near the crest of the TMR ( $\sim 25 \text{ m s}^{-1}$ ), and weaker wind ( $< 15 \text{ m s}^{-1}$ ) were  
325 observed at other sites (MOO, BKC, and GWW) in lower layers on both windward slope and lee  
326 side. The wind direction still revealed westerly at 03:00 UTC on 14 February 2018 (as Fig. 4b).  
327 However, strong winds were detected at the TMR (DGW) and in the downslope with a wind  
328 speed larger than  $20 \text{ m s}^{-1}$  above the BKC and GWW sites. Although the wind speed becomes  
329 stronger above 1.5 km MSL over the DGW site, it did not exhibit a significant change near the  
330 surface. The results demonstrated that persistent strong winds existed over the mountainous area  
331 (i.e., near the DGW site) while the LPS was approaching. After the LPS passed Korea, all  
332 sounding sites exhibited the wind speed weakened near the surface (Fig. 4c).



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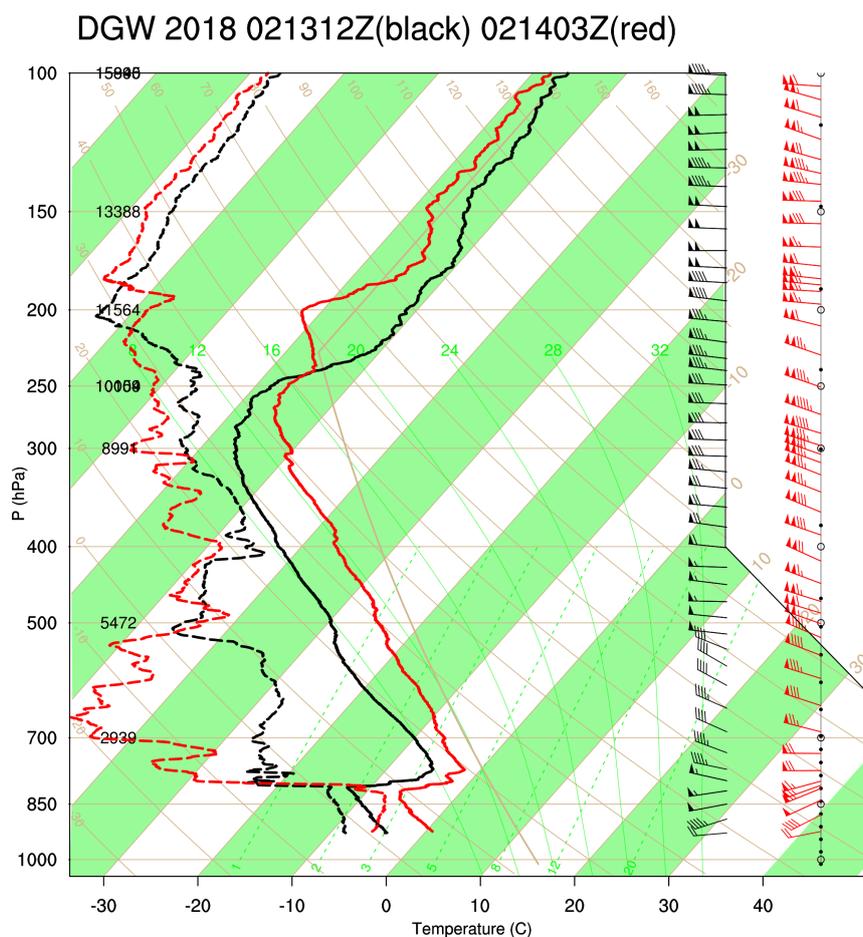
334 **Figure 4.** Horizontal winds observed by sounding and wind profiler along the cross line corresponding to Fig. 1a at  
335 (a) 12:00 UTC on 13 Feb. 2018, (b) 03:00 UTC on 14 Feb., and (c) 00:00 UTC on 15 Feb. 2018. A full wind barb  
336 corresponds to  $5 \text{ m s}^{-1}$ ; a half barb corresponds to  $2.5 \text{ m s}^{-1}$ . The color indicates the wind speed correspond to bars.  
337 The black shading in the lower portion indicates the averaged topography along the line in Fig. 1a.

338 Detailed environmental conditions in the upstream of the lee site (or the TMR area) was  
339 investigated by the sounding observations at the DGW site (Fig. 5). An inversion layer existed  
340 around 800 hPa height mainly due to the warm advection accompanied by the southwesterly at  
341 850 hPa ahead of the LPS (Figs. 2a and 2b at 12:00 UTC on 13 Feb. 2018) until it was passing  
342 through the Korea peninsula (at 03:00 UTC on 14 Feb. 2018). The temperature was increasing  
343 near the surface and became drier above the inversion layer between the two time steps. The wind  
344 direction comprised of westerly at all levels while the LPS was passing through. The wind speed  
345 became stronger above the inversion layer, but it exhibited no clear changes below  $\sim 800 \text{ hPa}$ . It  
346 is worth to mention that the inversion layer was mainly developed by two contributions: (1) large-  
347 scale warm advection and (2) stable boundary layer. Their separate contributions would require



348 a modeling study for this case in the future. The sounding observations showed a good condition  
349 for generating hydraulic jump and downslope windstorm in the lee side (Lee et al. 2020).

350 In summary, the environmental winds were mostly westerly in this case. However, the wind  
351 speeds revealed different characteristics at mountainous area and lee side. Strong winds ( $\sim 10$  m  
352  $s^{-1}$ ) persist near the surface on the mountainous area and the wind speed dramatically increased  
353 in the lee side.



354  
355 **Figure 5.** Profiles of temperature (solid lines), dew point (dashed lines), and horizontal winds observed by sounding  
356 at the DGW site at 12:00 UTC on 13 Feb. (black lines) and 03:00 UTC on 14 Feb. (red lines) 2018. A full wind barb  
357 corresponds to  $5 \text{ m s}^{-1}$  and a half barb corresponds to  $2.5 \text{ m s}^{-1}$ .



## 358 **4 Stronger winds in the lee side**

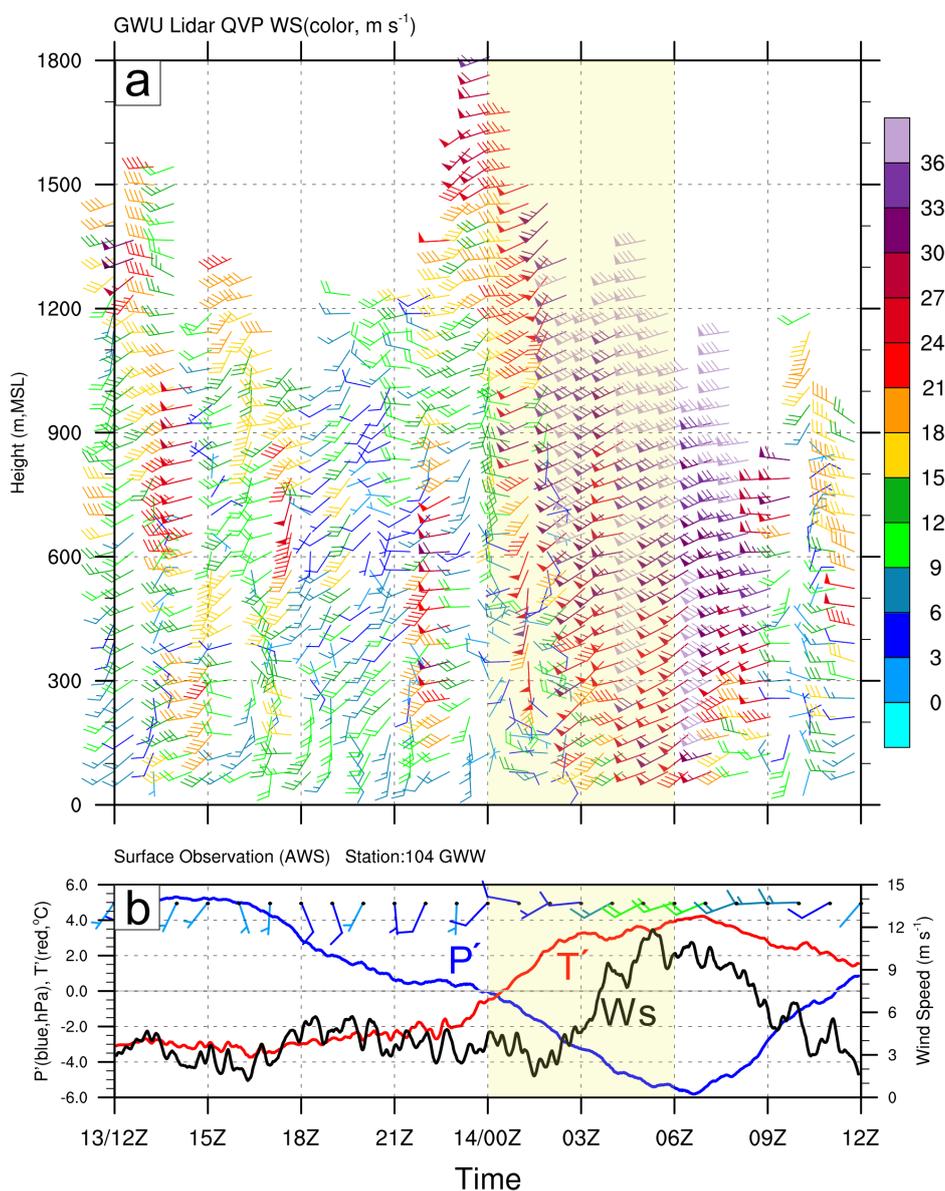
### 359 **4.1 The lidar and AWS observations**

360 Although the environmental wind direction likely westerly, the wind speed had a dramatic  
361 increasing trend in the lee side of the TMR. The detailed wind speed and surface fluctuations  
362 were documented by lidar quasi-vertical profile (QVP, Ryzhkov et al., 2016) at GWU site (upper  
363 panel) and the AWS observations at GWW site (lower panel) as shown in Fig. 6. The wind speed  
364 was relatively weak around 6- 9 m s<sup>-1</sup> at the lowest layer in the beginning of the research period.  
365 Strong winds were then measured by the lidar QVP reaching ~36 m s<sup>-1</sup> up to ~1.5 km MSL after  
366 00:00 UTC on 14 February 2018 (as Fig. 6a). It was clear that strong upper winds propagate  
367 toward the lower layer, intensify the wind speed near surface between 03:00 UTC and 09:00  
368 UTC. Finally, the wind speed became weak after 09:00 UTC on 14 February. Winds observed  
369 from the sounding and wind profiler were consistent with these QVP winds (cf. Fig. 4).

370 Fluctuations in wind speed, direction, perturbation pressure and perturbation temperature at  
371 GWW site were shown in Fig. 6b. The changes of wind speed were similar to the lowest layer of  
372 lidar observations (cf. Fig. 6a). Relatively weak winds were measured at the early stage of the  
373 period and the surface wind speed was intensified dramatically exceeding ~12 m s<sup>-1</sup> between  
374 00:00 and 06:00 UTC on 14 February (named as speed-up stage and highlighted by shaded area  
375 in the Fig. 6). The surface wind direction also showed similar patterns to the lidar observations  
376 as it had minor changes from more southerly to westerly. Although these two stations were at  
377 different locations along the northeastern coast of Korea, they revealed consistent changes on  
378 wind fields. The results also implied that the wind fields along the coast and in the lee side of the  
379 TMR have almost the same characteristics, which could be verified by the linear trend distribution  
380 of the surface wind speed (cf. Fig. 3a). Negative perturbation temperature was measured within  
381 the first 12 hours in the beginning of the period and increased after 00:00 UTC on 14 February



382 from  $\sim -3$  °C to 4 °C. The fluctuation of perturbation pressure showed an opposite phase with the  
383 perturbation temperature with the changing magnitude from around 5 to  $-6$  hPa. The wind speed  
384 was increasing just after the perturbation temperature (pressure) was rising (dropping). That is, a  
385 significant lag between  $T'$  ( $P'$ ) and wind speed is evident. Their specific relationships and  
386 mechanisms will be clarified through a more detailed analysis in Section 4.2.



387



388 **Figure 6.** (a) Time series of QVP (quasi-vertical profile) from lidar observations at GWU site during 12:00 UTC on  
389 13 Feb. to 12:00 UTC on 14 Feb. 2018. A full wind barb corresponds to  $5 \text{ m s}^{-1}$ ; a half barb corresponds to  $2.5 \text{ m}$   
390  $\text{s}^{-1}$  and the color indicates the wind speed ( $\text{m s}^{-1}$ ) correspond to color scale. (b) Time series of horizontal winds (wind  
391 barbs), wind speed ( $\text{m s}^{-1}$ , black line), perturbation pressure (hPa, P', blue line) and perturbation temperature ( $^{\circ}\text{C}$ , T',  
392 red line) observed from the AWS at GWW site. The time period with accelerating wind speed is also highlighted by  
393 light yellow shading (i.e., speed-up stage).

#### 394 4.2 Possible mechanism of strong winds in the lee side

395 The winds can usually be accelerated by the PGF between the two sites as the stronger wind  
396 usually occurred at the site where lower pressure was located. Therefore, the DGW site was  
397 selected as the upstream from the GWW site, and the differences between the surface temperature  
398 and sea level pressure will be analyzed. A relatively warm environment was presented in the lee  
399 side of the TMR and the temperature between the DGW and GWW sites suddenly increases from  
400  $\sim 7^{\circ}\text{C}$  in the beginning of the research period to  $\sim 8.5^{\circ}\text{C}$  after 00:00 UTC on 14 February (Fig.  
401 7a). The expected temperature difference between the two sites is about  $6.9^{\circ}\text{C}$  (adiabatic cooling  
402 rate for 0.7 km height difference) when the adiabatic heating assumed. The sea level pressure  
403 also decreased from  $\sim -1 \text{ hPa}$  to  $-4 \text{ hPa}$  at the time when the temperature was raising. The  
404 observed wind speed at the GWW site showed no obvious changes in the beginning. However,  
405 the wind speed has significantly increased just  $\sim 1$  hour after the sea level pressure (temperature)  
406 was decreasing (raising). This result revealed that the changes of wind speed are possibly related  
407 to the fluctuation of temperature and pressure. To clarify the role of pressure gradient on the wind  
408 speed at the DGW site, the local accelerations between the two sites could be approximated based  
409 on the horizontal momentum equation expressed as

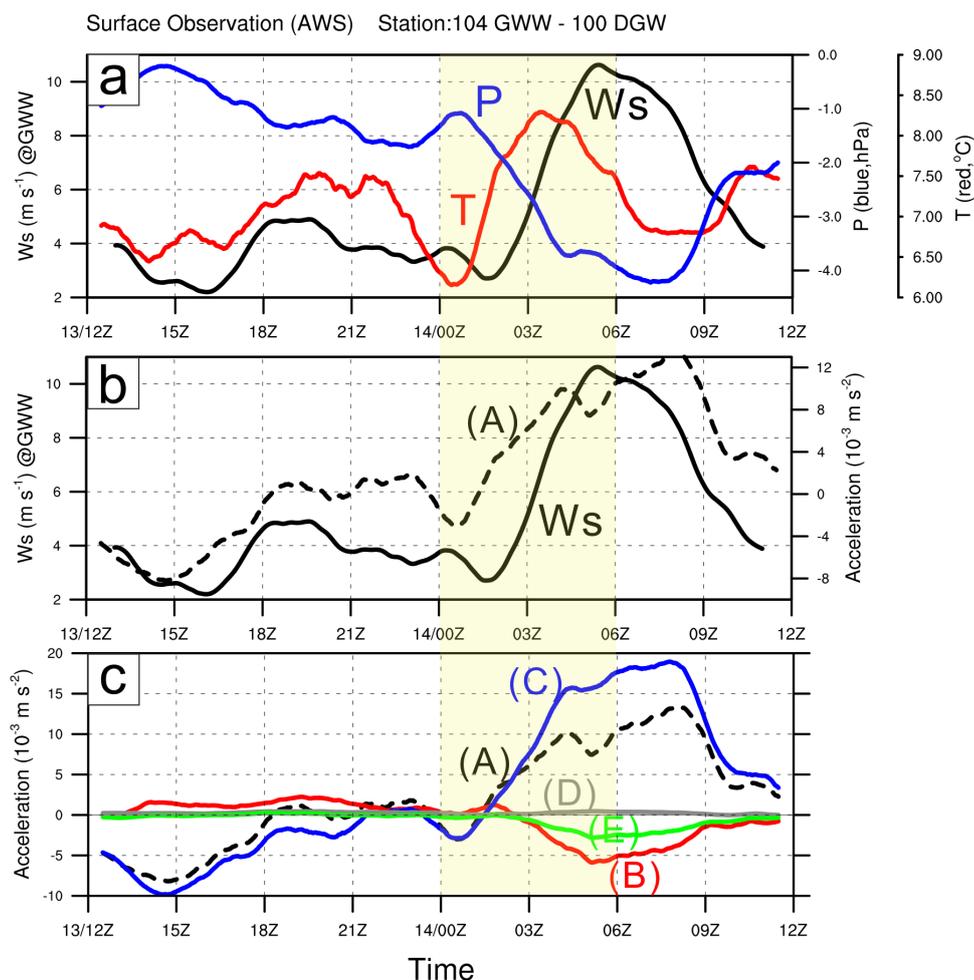
$$410 \quad \underbrace{\frac{\partial u}{\partial t}}_A = \underbrace{-u \frac{\partial u}{\partial x}}_B - \underbrace{\frac{1}{\rho} \frac{\partial P}{\partial x}}_C + \underbrace{fv}_D + \underbrace{\frac{C_d W_s u}{H}}_E. \quad (10)$$

411 In equation (10), Term A is the changes of  $u$  component with the time, which also  
412 corresponds to the wind accelerations along the west-east direction and Term B is the advective



413 acceleration amount to the distance ( $x$ ) between these two selected sites. Only the  $u$  component  
414 was considered in this study since the  $v$  and  $w$  components could be neglected because the  
415 environmental winds mostly comprised of westerlies (Yu et al., 2020). The PGF was indicated  
416 by Term C, where the  $\rho$  is air density and the  $P$  is sea level pressure. Coriolis acceleration and  
417 friction were indicated by Term D and Term E, respectively, where the  $C_d$ ,  $W_s$ , and  $H$  in Term  
418 E are the drag coefficient, wind speed and boundary layer height. The value of drag coefficient  
419 would most likely be a unitless constant based on Stull (1988) and was set as  $\sim 3.9 \times 10^{-3}$  in this  
420 study. The height of  $H$  used in this study was  $\sim 150$  (1500) m MSL according to the GWW  
421 (DGW) sounding observations at 00:00 UTC on 14 February 2018 (not shown).

422 Basically, the wind acceleration (i.e., Term A) that is derived from the equation (10) by  
423 adding terms from B to E is in good agreement with the fluctuation of wind speed at GWW site  
424 (Fig. 7b). A relatively weak wind speed occurred in the beginning and is coincident with negative  
425 and weak accelerations. Consequently, the wind speed jumped at GWW site in the speed-up stage  
426 (i.e., shading area in Fig. 7) associated with the increased and positive accelerations (i.e., Term  
427 A). Furthermore, the contribution of Terms B~E into Term A could also be evaluated individually  
428 by calculating each Term. PGF (Term C) dominated the changes of Term A with almost the same  
429 magnitudes during the entire research period as shown in Fig. 7c. In the beginning, the advective  
430 acceleration (Term B) can provide a little bit of positive contribution to Term A while PGF term  
431 is negative. However, both Term B and friction (Term E) gave negative feedback to Term A in  
432 the speed-up stage. Coriolis acceleration (Term D) always exhibited an almost zero acceleration  
433 to Term A due to sub-synoptic scale feature. The results suggested that the PGF would be the  
434 main factor to dominate the changes of wind speed at the GWW site in the lee side.



435  
 436 **Figure 7.** (a) Time series of wind speed ( $\text{m s}^{-1}$ , black line) observed from the AWS at GWW site, and the differences  
 437 of sea level pressure (hPa, blue line), temperature ( $^{\circ}\text{C}$ , red line) between the GWW and DGW sites from 12:00 UTC  
 438 on 13 Feb. to 12:00 UTC on 14 Feb. 2018. (b) Time series of the u component acceleration ( $10^{-3} \text{ m s}^{-2}$ , Term A,  
 439 black dashed line) estimated from the horizontal momentum equation [eq. (10)] between GWW and DGW sites. (c)  
 440 Time series of the u component acceleration (Term A, black dashed line), the advective acceleration (Term B, red  
 441 line), the PGF (Term C, blue line), Coriolis acceleration (Term D, gray line), and friction (Term E, green line). The  
 442 time period with accelerating wind speed is also highlighted by light yellow shading (i.e., speed-up stage).

443 Since the gusty wind can be explained by PGF, this result is also consistent with the  
 444 fluctuations in the perturbation pressure from the AWS observations at GWW site (cf. Fig. 6b)  
 445 as the observed perturbation pressure always comprised of negative values with the maximum of  
 446  $\sim -6$  hPa during the speed-up stage. To understand the possible causes of the relatively lower



447 pressure occurring in the lee side area, more detailed analysis is needed. The surface perturbation  
448 pressure can be estimated by solving the integration of the vertical momentum equation (Yu and  
449 Tsai, 2010),

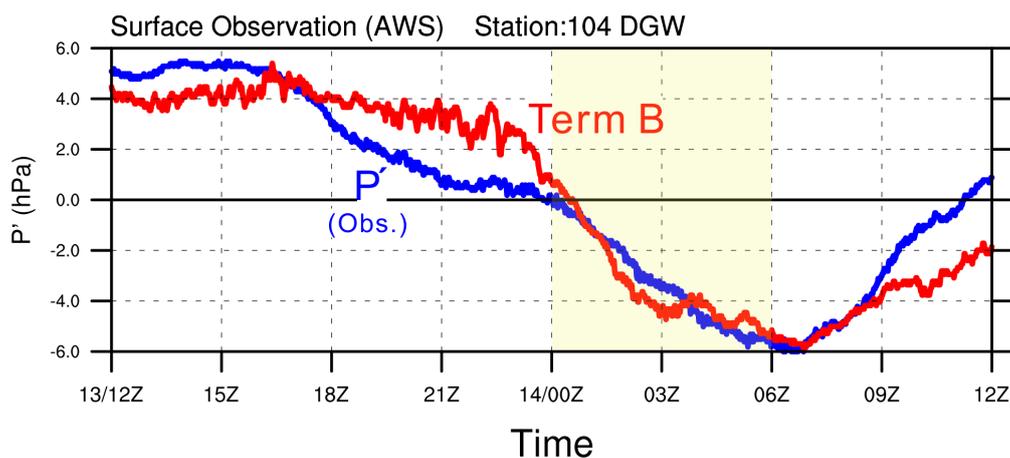
$$450 \quad P'_S = \underbrace{\int_0^{Z_T} \rho_0 \left( \frac{\partial w}{\partial t} + u \frac{\partial w}{\partial x} + v \frac{\partial w}{\partial y} + w \frac{\partial w}{\partial z} \right)}_A + \underbrace{\int_0^{Z_{LCL}} -\rho_0 g \frac{T'_v}{T_{v0}} dZ}_B$$
$$451 \quad + \underbrace{\int_{Z_{LCL}}^{Z_T} -\rho_0 g \frac{T'_v}{T_{v0}} dZ}_C + \underbrace{\int_0^{Z_T} \rho_0 g q_r dZ}_D, \quad (11)$$

452 where  $P'_S$  is surface perturbation pressure,  $\rho_0$  is the reference air density function of height, and  
453  $T_{v0}$  and  $T'_v$  are the reference and perturbation virtual temperature obtained from surface  
454 observations. The  $g$  and  $q_r$  in equation (11) correspond to the gravity acceleration and  
455 rainwater mixing ratio, respectively.  $Z_{LCL}$  and  $Z_T$  indicate the height of the lifting condensation  
456 level (LCL) and the cloud top, respectively. The LCL (~1900 m) can be calculated from the  
457 averaged sounding observations at GWW during the period. The terms on the right-hand side in  
458 equation (11) are different contributions to the estimated surface pressure ( $P'_S$ ) including  
459 nonhydrostatic effect (Term A), sub-cloud and in-cloud warming/cooling (Terms B and C), and  
460 water loading (Term D).

461 The Terms C and D would be possibly neglected since the studied case just occurred under  
462 fine weather condition. Remaining contributions for the estimated perturbation pressure may  
463 come from the nonhydrostatic effect (Term A) and in-cloud warming/cooling (Term B) because  
464 this study ignored the in-cloud warming/cooling and water loading. The estimated values of Term  
465 B have ~1 (~2) hPa underestimated (overestimated) from the observed perturbations pressure at  
466 ~15:00 (21:00) UTC on 13 February in advance of the speed-up stage and there is about 2 hPa ~  
467 3 hPa difference after 09:00 UTC on 14 February. Nevertheless, the estimated values of Term B  
468 are in good agreement with the observed perturbations pressure without significant differences  
469 (Fig. 8) in speed-up stage. The results suggested that the origin of the lower pressure in the lee  
470 side of TMR was deduced by sub-cloud warming. Such PGF could accelerate the wind speed in



471 the speed-up stage in the lee side of the mountains.

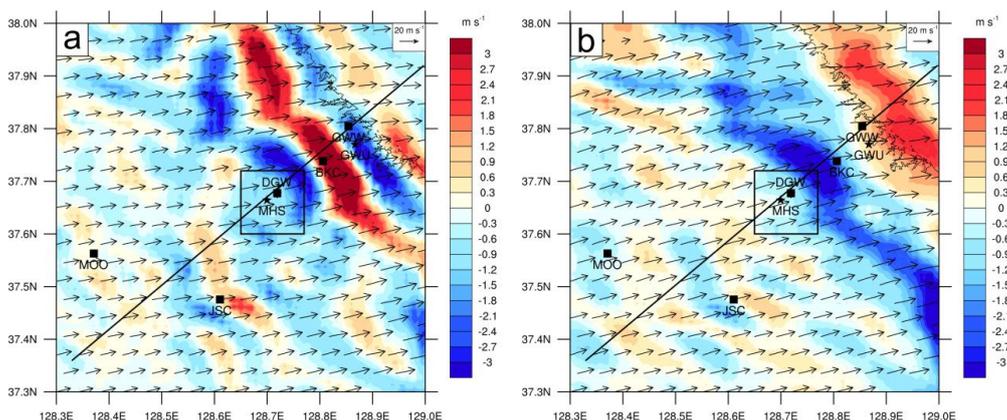


472

473 **Figure 8.** Time series of perturbation pressure (unit: hPa) observed by the AWS (blue line) and estimated Term B  
474 (sub-cloud cooling/warming, red line) in the vertical momentum equation [eq. (11)] at the GWW site from 12:00  
475 UTC on 13 Feb. to 12:00 UTC on 14 Feb. 2018. The time period with accelerating wind speed is also highlighted  
476 by light yellow shading (i.e., speed-up stage).

### 477 4.3 Mountain wave, hydraulic jump, and downslope windstorm

478 The mountain wave feature was detected in the local reanalysis data of the LDAPS (Fig. 9).  
479 Alternative downdraft and updraft were presented near the crest (near DGW site) and the lee side  
480 of the TMR at 21:00 UTC on 13 February 2018 (at 3 hours prior to the speed-up stage, Fig. 9a).  
481 The mountain wave propagated toward the northeastern direction (parallel to The TMR)  
482 associated with the interactions between the prevailing west-southwesterly winds and topography  
483 (lee wave in Fig. 9a). Stronger downdraft and updraft were characterized by positive and negative  
484 phases stronger than  $3 \text{ m s}^{-1}$  at DGW, BKC and GWW sites and the phase lines were parallel  
485 with the orientation of the TMR. Subsequently, in the speed-up stage, the mountain wave  
486 structures significantly changed at 03:00 UTC on 14 February 2018. The wavelength became  
487 longer but the wave still parallel to the TMR and the northeastern coast (Fig. 9b).



488

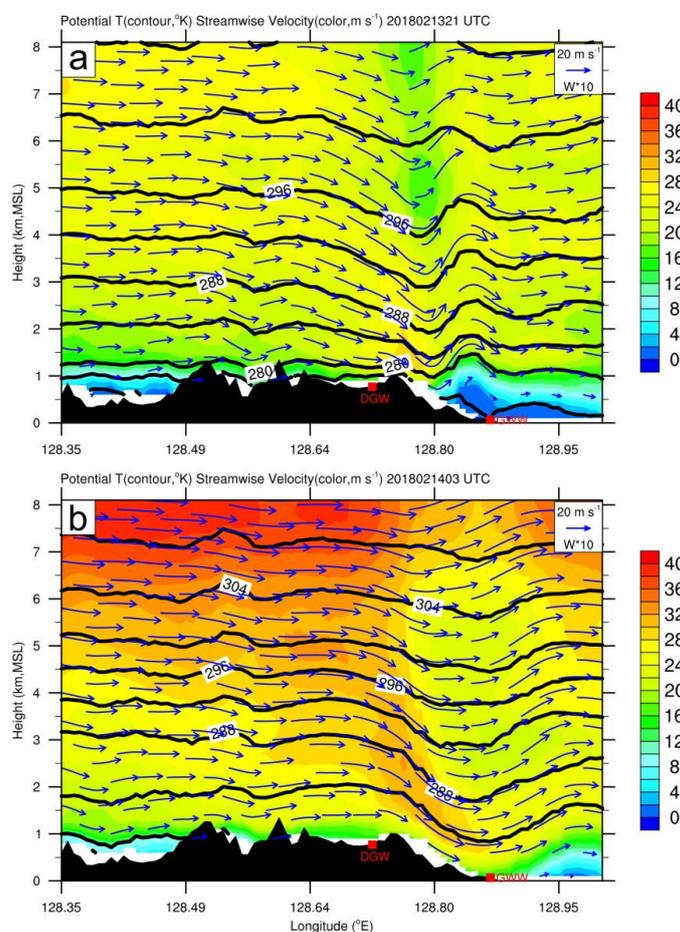
489 **Figure 9.** Horizontal distribution of the vertical velocity ( $\text{m s}^{-1}$ , color shading) and horizontal winds (vectors) at 2  
490 km MSL from the LDAPS in the domain corresponding to the Fig. 1a at (a) 21:00 UTC on 13 Feb. 2018 and (b)  
491 03:00 UTC on 14 Feb. 2018. The locations of the scanning Doppler lidars (soundings) sites are denoted by asterisks  
492 (squares).

493

494 The cross section of the potential temperature (thick solid line in Fig. 10) and streamwise  
495 velocity (colors in Fig. 10) perpendicular with the orientation of the TMR demonstrated the  
496 mountain wave characteristics in the lee side (between the DGW and GWW sites) at 21:00 UTC  
497 on 13 February 2018 (Fig. 10a). In this time period, a relatively stronger streamwise velocity only  
498 occurred near the downslope of the TMR ( $\sim 128.78^\circ\text{E}$ ,  $\sim 1$  km MSL) and coincides with the  
499 stronger downdraft. Weaker streamwise velocity ( $< 4 \text{ m s}^{-1}$ ) appeared near the GWW site in the  
500 coastal area. However, the potential temperature pattern in the speed-up stage was characterized  
501 by a longer wavelength with higher amplitude of the mountain wave (Fig. 10b), which is  
502 consistent with the vertical velocity field (Fig. 9b). A stronger wind exceeding  $30 \text{ m s}^{-1}$  (shaded  
503 orange colors in Fig. 10b) stemmed from higher altitudes as jet stream was approaching the  
504 Korean peninsula at this time. The strong wind propagated down to 0.5 km toward the GWW  
505 site. Note that the maximum values of Froude number related to the environmental winds at the  
506 DGW sounding site were estimated around 0.55–0.89. These Froude number were calculated  
507 from the sensitivity test using dry and saturated Brunt-Vaisala frequency and increasing the  
topographic height between 1000 and 2000 MSL. These structures were similar to the hydraulic



508 jump reported from several previously numerical studies in the northeastern coast of Korea (Kim  
509 and Cheong, 2006; Jang and Chun, 2008; Lee and In, 2009; Lee et al., 2020). The stronger  
510 streamwise velocity extended from the upper layers to the downslope of the TMR (exceeding  
511  $\sim 36 \text{ m s}^{-1}$ ) coincident with downdraft, which is similar to the downslope windstorm that  
512 contributed to the increased temperature via adiabatic warming on the slope. Along with this,  
513 surface velocity was intensified exceeding  $\sim 12 \text{ m s}^{-1}$  near the surface at the GWW site associated  
514 with the downslope wind. Note that the magnitudes of streamwise velocity are consistent with  
515 the fluctuation of the surface wind speed observed from the AWS (wind speed in Fig. 6). The  
516 result from this LDAPS analysis suggested that the downslope wind dominated the acceleration  
517 of wind speed in the lee side of the TMR.



518



519 **Figure 10.** Vertical cross section of the LDAPS potential temperature ( $^{\circ}\text{K}$ , contours), streamwise velocity ( $\text{m s}^{-1}$ ,  
520 color shading), and airflow (vectors) along the black lines in Fig. 9a at (a) 21:00 UTC on 13 Feb. 2018 and (b) 03:00  
521 UTC on 14 Feb. 2018. The black shading in the lower portion indicates the topography along the black line in Fig.  
522 9.

## 523 **5 Stronger winds in mountainous area**

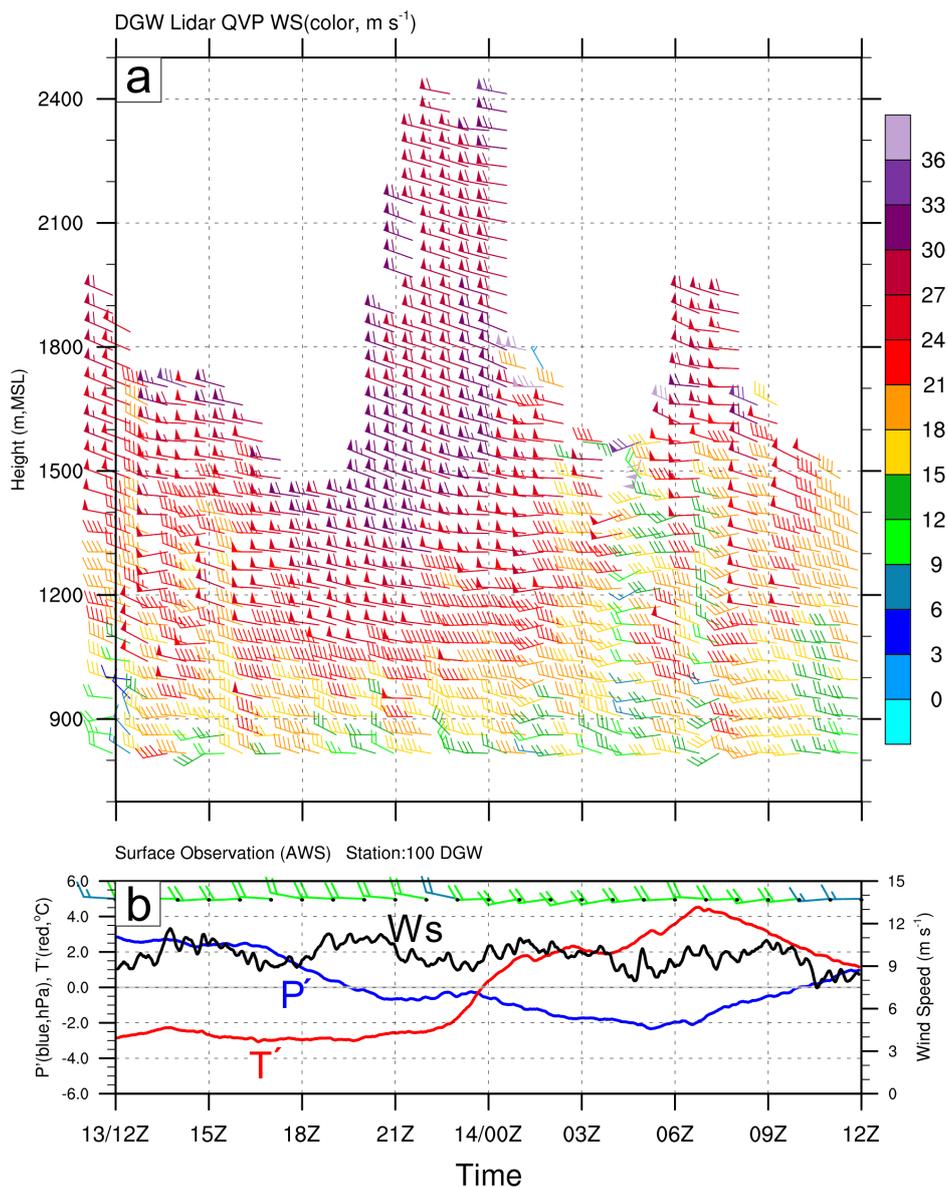
### 524 **5.1 The lidar and AWS observations**

525 The combination of the LPS and HPS provided a large-scale environmental wind favorable  
526 for westerly over mountainous area. According to the DGW QVP from observations (Fig. 11a),  
527 the wind speed ranged from  $\sim 12$  to  $36 \text{ m s}^{-1}$  at the low-level layers ( $\sim 900$  to  $1800 \text{ m MSL}$ ) during  
528 12:00 UTC on 13 February to 12:00 UTC on 14 February 2018. After this time period, the wind  
529 was decaying so quickly becoming nearly calm associated with the approaching HPS (Fig. 2e).  
530 The surface wind fluctuates in the range of  $7 \text{ m s}^{-1}$  to  $12 \text{ m s}^{-1}$  with the periodicity of 6h at the  
531 DGW site, similar to the pattern in the lidar QVP (Fig. 11b). These characteristics were quite  
532 different from the AWS and lidar observations in the lee side of mountains (for example GWW  
533 site). Unlike the coastal site, the strong wind can sustain for a day in the mountainous area. In  
534 particular, there was persistent westerly at the entire altitude in the mountainous area (i.e., DGW  
535 site). However, the wind direction was quite variable from southerly to westerly at the lee side of  
536 the mountainous area (GWU or GWW site). Significant strong winds were measured at the DGW  
537 site above  $1000 \text{ m MSL}$  on 13 February (Fig. 11), wind was weak at the GWU site (Fig. 6).  
538 Although the wind was getting stronger at the GWU and GWW sites on  $\sim 02:00$  UTC on 14  
539 February, the low-level or surface winds are getting slightly weaker at the DGW site. This trend  
540 is consistent with the wind fields from the sounding observations at the DGW site (Figs. 4a and  
541 4b).

542 The AWS observations at the DGW site demonstrated sustained strong westerlies ( $\sim 10 \text{ m s}^{-1}$ )



543 with periodic fluctuation from 12:00 UTC on 13 February to 12:00 UTC on 14 February 2018  
544 (Fig. 11b). Although the wind speed fluctuated periodically, no periodicity was shown in the  
545 perturbation pressure and perturbation temperature. Instead, the perturbation pressure  
546 monotonically dropped from 3 hPa at 12:00 UTC on 13 February to -2 hPa at 05:00 UTC on 14  
547 February and then increased to 1 hPa at 12:00 UTC on 14 February 2018. The perturbation  
548 temperature showed nearly opposite trend of the perturbation pressure. The perturbation  
549 temperature was nearly steady until 22:00 UTC and then increased to 4.5 °C at 07:00UTC on 14  
550 February 2018. Although the movement of LPS affected the changes of the surface perturbation  
551 pressure and temperature at the DGW site, the changes of the wind speed had no clear relation  
552 with the surface pressure and temperature.



553

554 **Figure 11.** (a) Time series of QVP (Quasi-Vertical profile) from lidar observations at DGW site from 12:00 UTC  
555 on 13 Feb. to 12:00 UTC on 14 Feb. 2018. A full wind barb corresponds to  $5 \text{ m s}^{-1}$ ; a half barb corresponds to  $2.5$   
556  $\text{m s}^{-1}$  and the color indicates the wind speed correspond to bar. (b) Time series of the horizontal winds (wind bars),  
557 wind speed ( $\text{m s}^{-1}$ , black line), perturbation pressure (hPa,  $P'$ , blue line), and perturbation temperature ( $^{\circ}\text{C}$ ,  $T'$ , red  
558 line) observed from the AWS at DGW site.



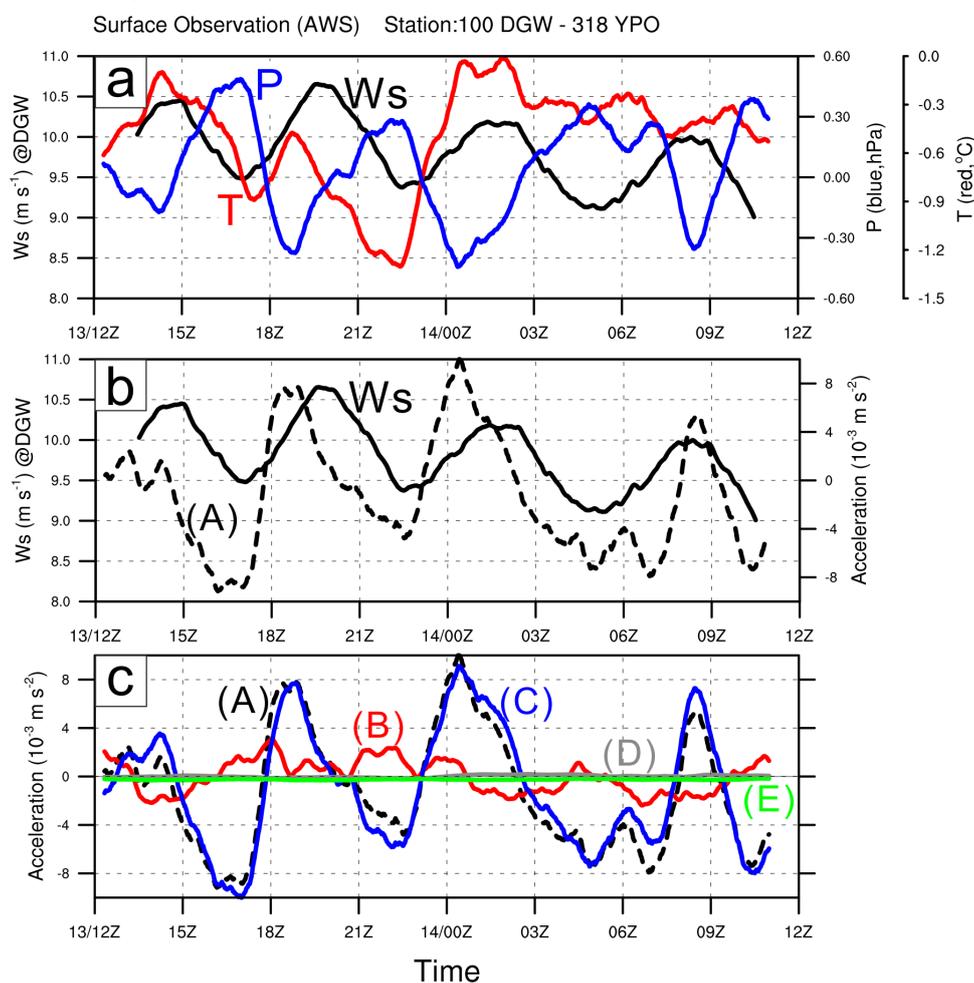
## 559 5.2 Possible mechanism of strong wind in mountainous area

560 To document the possible mechanisms about sustained strong wind occurring at DGW site  
561 over the mountainous area, difference of temperature and pressure were analyzed in detail.  
562 Similar to the DGW and GWW sites in Fig. 7, an upstream surface station (YPO site in Fig. 1b)  
563 was selected to calculate the temperature and pressure differences with the DGW site. Figure 12a  
564 reveals that the fluctuation of pressure differences (blue line in Fig. 12a) had an almost out of  
565 phase with the fluctuations in wind speed (black line in Fig. 12a) at the DGW site. Furthermore,  
566 the wind speed was gently decreasing with periodicity (wavelength of about 6h). This result  
567 provided a clue that the pressure gradient likely dominated the wind speed in this local area.  
568 Compared to the lee side of the mountains at the GWW site (Fig. 7), negative values of the  
569 temperature differences (minimum  $-1.3\text{ }^{\circ}\text{C}$ ) were calculated in the mountainous area and even  
570 became smaller ( $-0.5\text{ }^{\circ}\text{C}$ ) after 12:00 UTC on 14 February. Thus, the differences of pressure  
571 seemed to affect the wind speed patterns, and the fluctuation of wind speed did less relate to  
572 differences of temperature between these two sites. Note that the estimated values of Term B in  
573 eq.(11) are revealed significant different from observed perturbation pressure at the DGW site  
574 (between  $\sim 1\text{ hPa}$  and  $2\text{ hPa}$ ) during entire period, this result also indicated that the fluctuation of  
575 pressure is more related to nonhydrostatic effects [Term A in eq. (11)] but sub-cloud  
576 warming/cooling (not show). The periodic characteristic of the surface wind may have linked to  
577 nonlinear dynamic such as gap flow and gravity wave mechanisms (Shun et al., 2003).

578 The acceleration of wind speed at the DGW site can also be estimated by equation (10). Most  
579 of the estimated Term A and wind speed were also in a good agreement except for the short time  
580 period (Fig. 12b). Basically, the wind speed was increasing (decreasing) when estimated  
581 acceleration (Term A) is positive (negative). To understand the main contributor that dominates  
582 such local strong winds in this area, a detailed budget analysis of the momentum equation was  
583 performed (Fig. 12c). The PGF (Term C) was the most important factor for the estimated



584 acceleration, which means that the PGF could possibly determine the changes of the wind speed  
 585 at the DGW site. The advective acceleration was relatively small. The Coriolis force and friction  
 586 had no clear impacts on the acceleration (Term A).

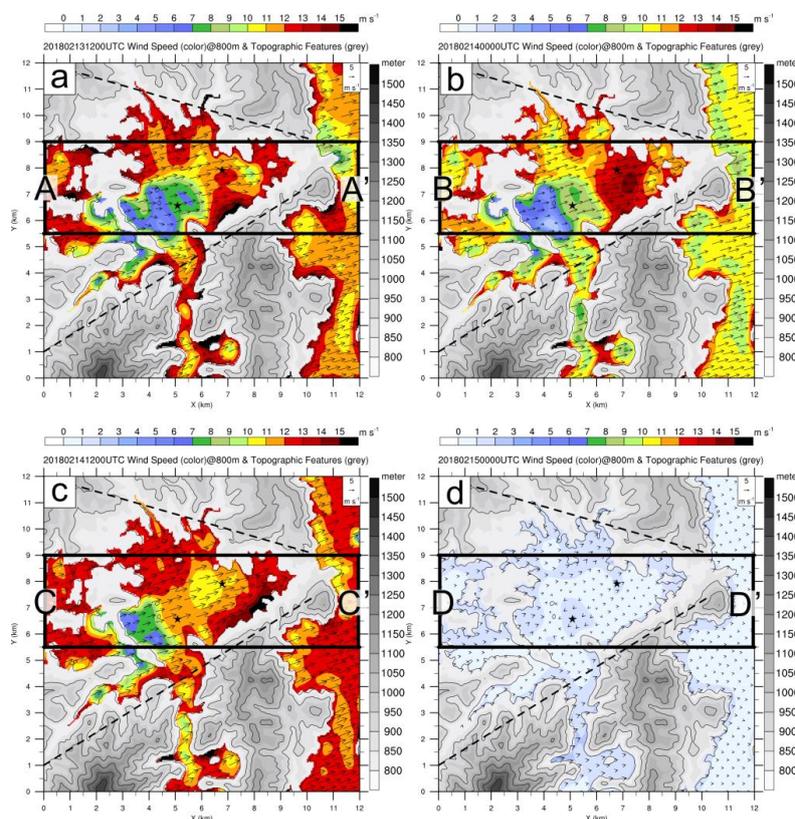


587  
 588 **Figure 12.** (a) Time series of wind speed ( $\text{m s}^{-1}$ , black line) observed from the AWS at DGW site, and the differences  
 589 of sea level pressure (hPa, blue line), temperature ( $^{\circ}\text{C}$ , red line) between the DGW and YPO sites from 12:00 UTC  
 590 on 13 Feb. to 12:00 UTC on 14 Feb. 2018. (b) Time series of the u component acceleration ( $10^{-3} \text{ m s}^{-2}$ , Term A,  
 591 black dashed line) estimated from the horizontal momentum equation [eq. (10)] between YPO and DGW sites. (c)  
 592 Time series of the u component acceleration (Term A, black dashed line), the advective acceleration (Term B, red  
 593 line), the PGF (Term C, blue line), Coriolis acceleration (Term D, gray line), and friction (Term E, green line).

594 The above results show that PGF is the main factor to accelerate wind speed, but temperature  
 595 is not a critical factor to change the PGF over the mountainous area. To determine the possible



596 factor that contributes into PGF, more detailed analysis of horizontal winds was performed with  
597 WISSDOM synthesis. Figure 13 demonstrates the fine-scale wind fields at 800 m MSL (near  
598 surface in the studied domain). In this height, a unique topographic feature can be explored as it  
599 is composed of a relatively wide (narrow) area in the western (eastern) side along the valley. This  
600 channel like feature could be emphasized and was marked by the area between two thin dashed  
601 lines in Fig. 13. Four periods (00:00 UTC on 13 February 2018, 00:00 UTC on 14 February,  
602 12:00 UTC on 14 February, and 00:00 UTC on 15 February 2018) were selected to investigate  
603 the changes of wind patterns in this channel along the valley. The prevailing wind was westerly  
604 with a slight deflection near the center of the domain and the eastern side of the valley while the  
605 LPS approached Korea (Figs 13a, 13b, and 13c). Nevertheless, a relatively weak wind ( $\sim 6 \text{ m s}^{-1}$ )  
606 always existed in the center of the domain nearby the MHS lidar site (wide segment of the valley)  
607 and a stronger wind ( $14 \text{ m s}^{-1}$ ) was observed in the downstream nearby the DGW site (narrow  
608 segment of the valley). Wind speed decreased and nearly became calm after the LPS moved away  
609 from Korea (Fig. 13d).



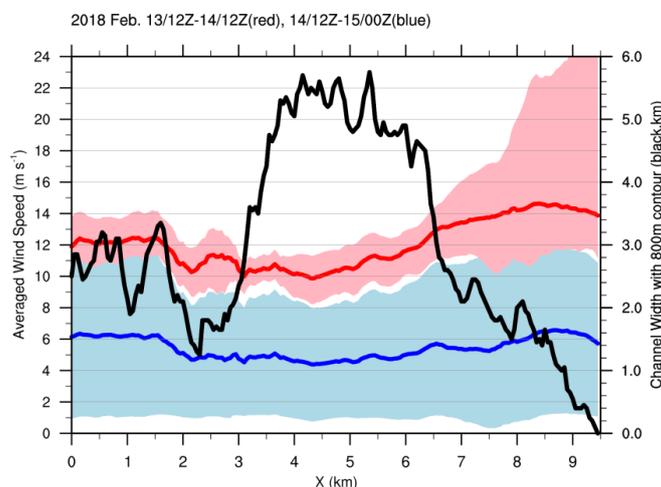
610  
611 **Figure 13.** Horizontal distribution of the wind speed ( $\text{m s}^{-1}$ , color shading) at 800 m (MSL) retrieved in the  
612 WISSDOM domain at (a) 00:00 UTC on 13 Feb. 2018, (b), 00:00 UTC on 14 Feb., (c) 12:00 UTC on 14 Feb., and  
613 (d) 00:00 UTC on 15 Feb. 2018. The black dashed lines mark the area of the channel to calculate the averaged wind  
614 speed and channel width as shown in Fig. 14. The inserted box indicates the averaged area in the vertical cross  
615 sections along the valley (A-A'). Topographic features indicated by the gray shading and contours. Locations of the  
616 scanning Doppler lidar sites are denoted by asterisks.

617 To understand the relations between the topography and winds, the average wind speed  
618 (colored thick lines in Fig. 14) and the channel width (thick black line in Fig. 14) along the valley  
619 at 800 m MSL were calculated in two time periods when the LPS approaches (before 12:00 UTC  
620 on 14 February 2018) and leaves (after 12:00 UTC on 14 February 2018). The channel width was  
621 around 2 km at  $x = 0$  km to 3 km (western side) and became wider  $\sim 5.5$  km at  $x = 3$  km to 6.5  
622 km. The channel width then decreased significantly to nearly 0 km at  $x = 6.5$  km to 9.5 km.

623 When LPS was approaching (average wind speed in red line and range of minimum and



624 maximum wind speed in shading in Fig. 14), the average wind speed increased from  $\sim 10 \text{ m s}^{-1}$   
625 to  $\sim 14 \text{ m s}^{-1}$ , which was coincident with the changing of channel width from  $\sim 5.5 \text{ km}$  to  $0 \text{ km}$   
626 along the valley. The maximum wind speed was larger than  $24 \text{ m s}^{-1}$  near the narrowest segment  
627 of the valley. When LPS was leaving (blue line and shading), the averaged wind speed had a  
628 nearly constant strength ( $\sim 5 \text{ m s}^{-1}$ ) without obvious changes and no clear relations with the width  
629 of the valley. This analysis reveals that the channeling effect may play an important role to  
630 dominate the spatial distribution of wind speed with the valley. Furthermore, downstream  
631 acceleration in the narrow segment of valley could be more significant with strong upstream  
632 winds.

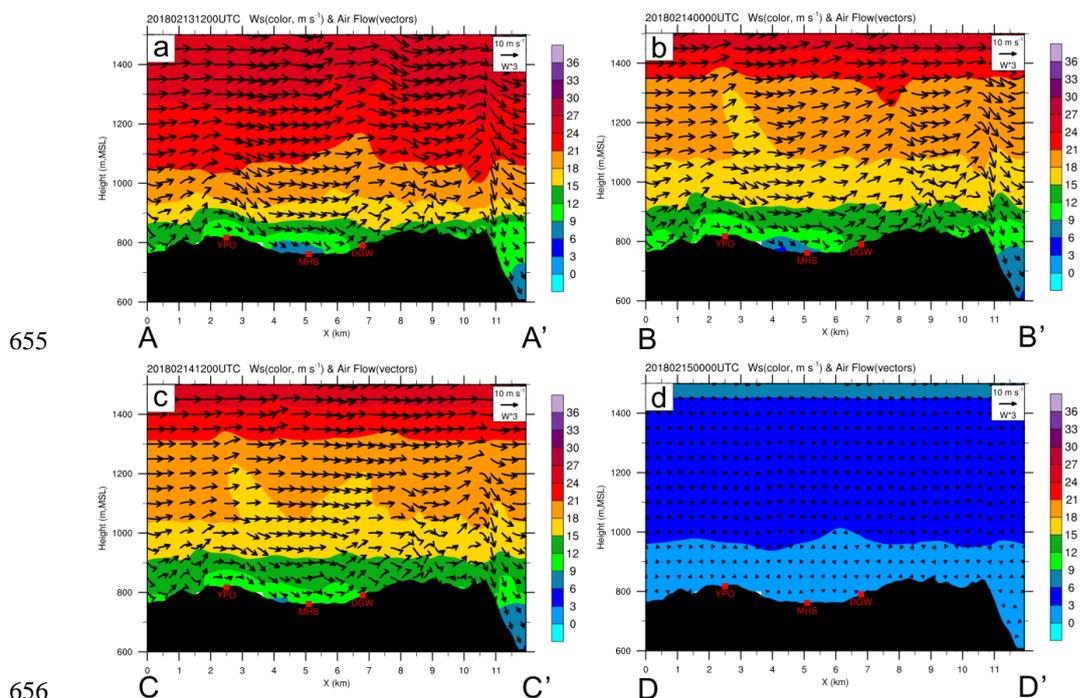


633  
634 **Figure 14.** Averaged wind speed and its range along the valley corresponding to the area indicated by the dashed  
635 lines in Fig. 13 at two times: 12:00 UTC on 13 Feb. to 12:00 UTC on 14 Feb. (red line and shading) and 12:00 UTC  
636 on 14 Feb. to 00:00 UTC on 15 Feb. 2018 (blue line and shading). The red and blue shading show the maximum and  
637 minimum values along the valley for the two times. Averaged channel width along the valley was plotted by a thick  
638 black line.

639 Figure 15 shows the mean vertical structures of wind speed, airflow, and topographic features  
640 from each cross section along the boxes in Fig. 13. The boxes were set on our main focus area  
641 from wider to narrow segments along the valley and parallel with the environmental wind  
642 direction (westerly). These analyses allow us to investigate detailed airflow features from near



643 the surface to higher altitudes and their interactions with topography. The four time periods were  
644 12:00 UTC on 13 February 2018, 00:00 UTC and 12:00 UTC on 14 February, and 00:00 UTC  
645 on 15 February 2018. The mean vertical structures in the first three periods (when the LPS  
646 approaching) revealed similar characteristics that the uniform and stronger westerly winds (larger  
647 than  $\sim 18 \text{ m s}^{-1}$ ) were above the layers higher than 1 km MSL. In contrast, the airflow had more  
648 significant variances near the surface layers. In the layers below 1 km MSL, the westerly winds  
649 were lifted on the upslope and become downdraft behind the mountain crests. The wind speed  
650 was quite weak (strong) near the MHS (DGW) site, which is coincident with the relatively wider  
651 (narrow) segment in the valley on the three time periods. In particular, the high wind speed area  
652 was only presented between  $x = \sim 6.5 \text{ km}$  and  $9.5 \text{ km}$  (i.e., the narrowest segment of the valley).  
653 The winds became more uniform and weaker in the upper layers and near the surface when the  
654 LPS moved away from Korea at 00:00 UTC on 15 February 2018 (Fig. 15d).

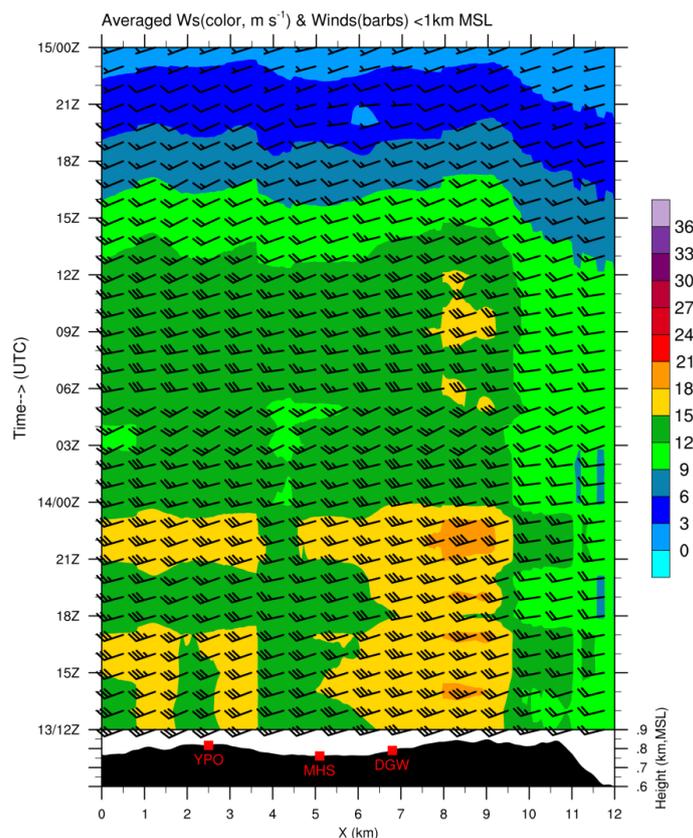


656  
657 **Figure 15.** Averaged vertical cross section of the WISSDOM-derived wind speed ( $\text{m s}^{-1}$ , color shading) and wind  
658 vectors (combined cross barrier flow and threefold vertical velocity) at four time periods (a) 12:00 UTC on 13 Feb.  
659 2018, (b) 00:00 UTC on 14 Feb. 2018, (c) 12:00 UTC on 14 Feb. 2018, and (d) 00:00 UTC on 15 Feb. 2018. The



660 area of the cross section is shown in the black box in Fig. 13. The black shading in the lower portion indicates the  
661 topography along the box.

662 Because the winds manifested clear variations only near the surface layers, the mean vertical  
663 structures of wind speed and directions could be further averaged just below 1 km MSL. Fig. 16  
664 shows the continuous time series of averaged wind field during the entire period with the same x  
665 axis in Fig. 15. Results demonstrate that the winds near the surface layers were accelerated in the  
666 narrow segment between  $x = \sim 6.5$  km and 9.5 km for strong enough upstream winds (before  
667 00:00 UTC on 14 February). This characteristic is similar to the gap wind or channeling effect  
668 from the previous simulation and observational studies (Overland and Walter, 1981; Neiman et  
669 al., 2006; Heinemann, 2018). Consequently, a relatively weak channeling effect induced weaker  
670 winds in the narrow segment of valley during 00:00–15:00 UTC on 14 February 2018 because  
671 the environmental winds from upstream became weaker. Finally, the channeling effect was no  
672 longer existing when the upstream winds became calm after 15:00 UTC. The wind would possibly  
673 accelerate when it blows from wider to narrow segments of the valley due to PGF as it is related  
674 to the Bernoulli's Law, i.e., the pressure would reduce when the flow speed is increased and vice  
675 versa. Observational analysis reveals a relatively low pressure in the narrow segment of valley  
676 and thus, PGF would locally dominate the airflow accelerating over the mountainous area.



677

678 **Figure 16.** Temporal variation of the averaged wind speed ( $\text{m s}^{-1}$ , color shading) and the horizontal winds (wind  
679 barbs) from WISSDOM derived in the valley from 00:00 UTC 13 Feb. on 13 to 00:00 UTC on 15 Feb. 2018. The  
680 low-level winds (below 1 km MSL) within the black boxes in Fig. 15 were averaged in a direction normal to the  
681 orientation of the boxes. The black shading in the lower portion indicates the averaged topography along the boxes.

## 682 6. Conclusion

683 This study uses Doppler lidars, wind profiler, soundings, and surface observations to  
684 examine a strong downslope wind event during the ICE-POP 2018. Detailed characteristics of  
685 wind fields and possible mechanisms were explored during the passage of a low-pressure system  
686 (LPS) over the northern part of the Korea Peninsula on 13–15 February 2018. Although the wind  
687 speed is generally increased in South Korea when the LPS is approaching, it comprised of more  
688 significant increasing trend along the downslope and in the lee side of the TMR (Taebek



689 Mountain Range). The wind speed has no obvious changes but are persistently strong over the  
690 TMR. Conspicuous gradient of linear trend of the wind speed only existed between the  
691 mountainous areas and in the lee side areas. Moreover, the wind speed shows a decreasing trend  
692 synchronously after the LPS moves away from Korea.

693 From the sounding observations, low-level environmental winds revealed high variability  
694 from the mountainous area to the lee side of the mountains. The wind direction is comprised of  
695 most westerly associated with the LPS and the wind speed are sustainedly strong ( $\sim 10 \text{ m s}^{-1}$ ) at  
696 the DGW site (i.e., mountainous area) during the research period. However, the wind speed at  
697 the lee side (GWW) clearly changed from being relatively weak to a stronger one. The winds  
698 then become nearly calm both in the mountainous or lee side areas after the LPS moved away  
699 from the Korean peninsula. In addition, upstream inversion layers (at  $\sim 850 \text{ hPa}$  level) were also  
700 detected by sounding observations at the DGW site while the strong wind occurred in the lee side  
701 of the mountains.

702 In the lee side of the mountains, the surface wind speed has dramatically increased (from  $\sim 3$   
703 to  $12 \text{ m s}^{-1}$ ) at the GWW site during the research period. The surface temperature (pressure)  
704 perturbation is also changed to positive (negative) values and showed significant time lag with  
705 wind speed change. The sea level pressure and temperature differences between the mountainous  
706 station at the DGW and the lee side station at the GWW demonstrate that the wind speed is  
707 suddenly raised with increasing temperature (exceeded  $\sim 8.5^\circ\text{C}$ ) and decreasing pressure (from  $-1$   
708 hPa to  $-4 \text{ hPa}$ ). The estimated wind accelerations [Term A in eq. (10)] are in good agreement  
709 with the observed wind speed, which are mainly contributed by the PGF [Term C in eq. (10)].  
710 The negative surface perturbation pressure at the GWW site has nearly the same magnitude with  
711 the estimated surface perturbation pressure from the integration of the vertical momentum  
712 equation. The fluctuation of surface perturbation pressure is mainly dominated by the sub-cloud  
713 warming/cooling [i.e., Term B in eq. (11)]. Results indicate that the adiabatic warming plays an  
714 important role to reduce the surface pressure and the winds are accelerated by PGF in the lee side



715 of the TMR. Furthermore, the downslope winds were also dominated by the stronger wind  
716 occurring along the lee side based on the LDAPS analysis. The development of the strong  
717 downslope wind is highly related to the mountain wave and hydraulic jump.

718 In the mountainous area, persistent surface strong wind (with the fluctuation of  $4.5 \text{ m s}^{-1}$ )  
719 was observed at the DGW site when the LPS was approaching (leaving). The surface wind has  
720 no clear relationship with the surface perturbation pressure and perturbation temperature.  
721 However, the sea level pressure differences between the upstream station YPO and DGW show  
722 similar amplitudes (out of phase) in the fluctuation of surface wind speed. In contrast, the  
723 temperature differences are small (between  $-0.5^\circ\text{C}$  and  $1.2^\circ\text{C}$ ) with no clear relations with the  
724 fluctuation of surface wind at the DGW site. Although the temperature has no clear relation with  
725 the strong wind, estimated wind accelerations [Term A in eq. (10)] results are in good agreement  
726 with the observed surface wind speed. It means that the PGF is still the main contributor for the  
727 wind acceleration at the DGW site. The 3D winds derived from WISSDOM synthesis also reveal  
728 that the wind speed at the DGW site (narrow segment in the valley) is always stronger than the  
729 YPO site (wider segment in the valley) when the LPS was approaching. Thus, the channeling  
730 effect is possible mechanism to dominate the wind acceleration in the mountainous area.

731 In this study, the observational evidence shows that the different mechanisms are important  
732 references to determine the strength and persistence of the orographically strong winds in the  
733 same underlying LPS under fine weather condition. In the future, high-resolution numerical  
734 modeling analysis will be performed for all strong wind events during the ICE-POP 2018,  
735 because the detailed thermodynamic information was desired to give more complete descriptions  
736 about the distribution of potential temperature across the mountainous area. The kinematic and  
737 thermodynamic information from the simulations will be important indicators to further  
738 investigate the existing of mountain wave, included the hydraulic jump, wave breaking, and  
739 partial reflection for the generation of the downslope windstorm. More cases will be included to  
740 provide comprehensive explanations of the strong downslope wind in the northeastern



741 mountainous part of South Korea. More importantly, we aim to extend our understanding on the  
742 variability of winds around the terrain in a very fine-scale even at different seasons.

743

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751

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753

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774



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