Scattering and Absorption Cross Sections of Atmospheric Gases in the Ultraviolet-Visible Wavelength Range (307 - 725 nm)

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Abstract

Accurate Rayleigh scattering and absorption cross sections of atmospheric gases are essential for understanding the propagation of electromagnetic radiation in planetary atmospheres. Accurate extinction cross sections are also essential for calibrating high finesse optical cavities and differential optical absorption spectroscopy and for accurate remote sensing. In this study, we measured the scattering and absorption cross sections of carbon dioxide, nitrous oxide, sulfur hexafluoride, oxygen, and methane in the continuous wavelength range of 307–725 nm using Broadband Cavity Enhanced Spectroscopy (BBCES). The experimentally derived Rayleigh scattering cross sections for CO_2, N_2O, SF_6, O_2, and CH_4 agree with refractive index-based calculations, with a difference of (0.37 ±1.24)%, (–0.55 ±1.06)%, (0.91 ±1.35)%, (2.81 ±1.21)%, and (0.89 ±2.18)%, respectively. The O_2-O_2 collision-induced absorption and absorption by methane are obtained with high precision at the 0.8 nm resolution of our BBCES instrument in the 307–725 nm wavelength range. New dispersion relations for N_2O, SF_6, and CH_4 were derived using data in the UV-vis wavelength range. This study provides refractive index dispersion relations, *n*-based Rayleigh scattering cross sections, and absorption cross sections for these gases.

1. Introduction

The dominant interactions of gas-phase molecules with light in Earth's atmosphere can be divided into absorption, where the light energy is converted to internal energy and generally (at atmospheric pressures) transferred to the surrounding environment either as heat or as photoemission, and light scattering where the gases redistribute the light energy in the atmosphere. The knowledge of light extinction (scattering + absorption) by gases is essential for predicting the radiative transfer in the atmospheres of the Earth and other planets. In addition, the light extinction by gases is widely used for determining the effective optical pathlength of high-finesse optical cavities that measure trace gases and aerosols (Washenfelder et al., 2013; Washenfelder et al., 2008; Wilmouth and Sayres, 2019; Jordan et al., 2019) and for Differential Optical Absorption Spectroscopy (DOAS) to infer information about the light extinction properties of aerosols and clouds in the open atmosphere (Baidar et al., 2013; Platt and Stutz, 2008).

The interaction of light with a wavelength much larger than the size of a molecule/particle gives rise to the scattering of light, which is known as Rayleigh scattering(Strutt, 1899). Rayleigh scattering accounts for scattering, local field effects (Lorentz–Lorenz) (Strutt, 1920) as well as depolarization from the non-sphericity of molecule/particles (King correction factor) (King and Eve, 1923; Strutt, 1918). For a gas with known refractive index (n_ν) and King correction factor ($F_k(\nu)$), the wavelength-dependent Rayleigh scattering cross section (σ_ν, cm^2 molecule^{-1}) can be calculated as follows (Sneep and Ubachs, 2005):

$$\sigma_\nu = \frac{24\pi^3\nu^4}{N^2}\left(\frac{n_\nu^2-1}{n_\nu^2+2}\right)^2 F_k(\nu) \quad (1)$$

where N is the number density of the gas (molecules cm^{-3}) and ν is the wavenumber of the light (cm^{-1}). Note that the cross section contains the gas number density but is not in fact dependent on the number density since the refractive index also appears in the expression. This *n*-based method is an advantageous approach for calculating Rayleigh scattering cross sections, but it is vital to note that the accuracy of the calculated cross sections depends on the experimentally-determined refractive indices and the King correction factors. In particular, cautions should be used when applying a dispersion formula derived from measurements in one wavelength region to calculate Rayleigh scattering cross sections in a different wavelength range.

Direct experimental measurement of Rayleigh scattering cross sections is essential given the potential uncertainties in *n*-based calculations. While measurements of the King correction factors and refractive index for gases are well known from the literature (Cuthbertson and Cuthbertson, 1932; Leonard, 1974; Strutt, 1920; Vukovic et al., 1996; Hohm, 1993), there are only a few direct measurements of Rayleigh scattering cross sections (Fuchs et al., 2009; He et al., 2018; Ityaksov et al., 2008a, b; Jordan et al., 2019; Naus and Ubachs, 2000; Sneep and Ubachs, 2005; Thalman and Volkamer, 2013; Thalman et al., 2014; 2017; Wilmouth and Sayres, 2019; 2020), especially measurements with a continuous spectrum from ultraviolet to visible.

Rayleigh scattering cross section measurements were previously performed at a single wavelength (e.g., 458 nm, 532 nm, 632.8 nm) using Nephelometry (Shardanand and Rao, 1977) and cavity-ring down spectroscopy (CRDS) (Ityaksov et al., 2008a, b; Naus and Ubachs, 2000; Sneep and Ubachs, 2005; He et al., 2018). More recently, advanced Broadband Cavity Enhanced Spectroscopy (BBCES) was used to determine the Rayleigh scattering cross sections of gases such as Ar, CO_2, O_2, SF_6, and CH_4. The BBCES technique enables the measurement of Rayleigh scattering cross sections over a broad wavelength range. Thalman et al. (2014) performed measurements over selected wavelength regions between 350 and 660 nm using six BBCES cavities for N_2, Ar, and O_2. The BBCES were calibrated with He and N_2 using Rayleigh scattering cross sections calculated using their refractive index and from cavity-ring down measurements, respectively. They found a good agreement with *n*-based values to within 0.2 ±0.4%. Recent studies using BBCES with 30 nm spectral range were also used for Rayleigh scattering cross section measurement in the UV wavelength region and demonstrated excellent agreement with *n*-based values for Ar and CO_2 (Wilmouth and Sayres, 2020, 2019). Recently, Rayleigh scattering cross sections for CO_2 were measured using BBCES at visible wavelengths between 400 and 650 nm, and agreement with *n*-based values was within 2.4% on average. To the best of our knowledge, there is no direct continuous wavelength measurements of extinction cross sections of gases that covers the ultraviolet across the entire visible range (300–725 nm) as shown in Table 1. Recently, Wilmouth and Sayres (2020) combined refractive index data in the UV region (264-297 nm and 333-363 nm) and at several single wavelengths in the visible, and they derived the dispersion relation of refractive index for SF_6 and CH_4 in the wavelength range of 264-650 nm. However, more data in the visible range are needed in order to further validate these dispersion relations.

In this study, we used a recently-developed BBCES instrument to measure the extinction cross sections of CO_2, N_2O, SF_6, O_2, and CH_4 continuously across the wavelength region 307–725 nm. All of the measurements were done at a single pressure to eliminate effects due to alignment. This requires the use of two gases with different Rayleigh cross sections for the calibration of the BBCES instrument since the reference state is not vacuum. In this study, He and N_2 were used to calibrate the system. By using the *n*-based calculated Rayleigh scattering cross sections of He and N_2 to calibrate the path length of the optical cavity, the other cross sections can be determined relative to the difference between these two gases. We report high accuracy Rayleigh scattering cross sections for all five gases and compared our results with previous *n*-based values. New dispersion relations for N_2O, SF_6, and CH_4 are derived by incorporating data obtained by this study, extinction cross section data in the deep UV, and previously available scattering cross section data in the visible wavelength range.

2. Methods

2.1 Extinction measurement using BBCES

The BBCES systems used in this study are analogous to our previous studies (He et al., 2018; Washenfelder et al., 2016; Bluvshtein et al., 2016). Briefly, our BBCES consists of two channels, one in the UV ($BBCES_{UV}$, 307-350 nm) and one in the UV-vis range ($BBCES_{Vis}$, 338-725 nm). The two channels of the BBCES share a laser-driven Xenon arc lamp source (LDLS EQ−99CAL, Energetiq Technology, Inc., MA, USA) coupled with a high transmission UV-Vis optical fiber from which the light is collimated and focused (BBFIBERX-600-1M, Energetiq Technology, Inc., MA, USA). The light source was purged with high purity N_2 and cooled by an aluminum block (with 15℃ circulating water inside) to maintain stable optical power output. The UV light from the fiber was reflected by a low-pass dichroic mirror and filtered (Schott Glass WG310 and UG11) into the $BBCES_{UV}$ channel, which has a cavity with two 2.5 cm diameter, 1 m radius of curvature mirrors, with manufacturer's reported reflectivity of 0.9995 (per pass loss = 500 parts per million, ppm) at the nominal center wavelength of 330 nm (Advanced Thin Films, Boulder, USA). The transmitted UV-vis light from the beam splitter was reflected and filtered (Schott Glass WG345 and Edmund Optics 15-261) into the $BBCES_{Vis}$ channel consisting of two 2.5 cm, 1 m radius of curvature mirrors (FiveNine Optics, USA) with manufacturer's reported reflectivity above 0.9993

(loss < 700 ppm), see Figure S1. The light emerging through the rear mirror of the cavity was focused using a 0.1 cm F/2 fiber collimator (74−UV, Ocean Optics, Dunedin, FL, USA) into a high transmission UV-vis optical fiber which directs the light into a high-performance spectrometer (QEPro, Ocean Insight, USA). Before gas measurement, the wavelength of the spectrometer was calibrated using an HG-1 mercury argon calibration light source (Ocean Insight, USA) within the wavelength range of 302.15−727.29 nm. During these experiments, a 300 line mm^{-1} grating and a 200 μm entrance slit width were used. The CCD array is a back-illuminated detector with 1024×56 pixels (Hamamatsu S7031-1006, Japan) thermo-electrically cooled to −10 ℃ to reduce thermal noise. Individual spectra at a wavelength resolution of 0.8 nm were acquired with 3.0 s integration time, and a total of 150 spectra were recorded during each measurement.

During the extinction measurements, the entire 94.0 ± 0.1 cm pathlength between the mirrors was filled with He, N_2, CO_2, N_2O, SF_6, or CH_4. The gases were obtained from several vendors (Airgas, Linde) with the following purities: He, 99.995%; N_2, 99.999%; N_2O, 99.999%, CO_2, 99.999%; SF_6, 99.999%; CH_4, 99.9995%.

The reflectivity of the mirrors (R(λ)) can be determined as a function of wavelength (λ) by taking into account the difference in the extinction due to known literary data of Rayleigh scattering coefficient (α_{Ray}^{gas}) by two different gases such as N_2 ($\alpha_{Ray}^{N_2}(\lambda)$) and He ($\alpha_{Ray}^{He}(\lambda)$) (Washenfelder et al., 2008).

$$R(\lambda) = 1 - \mathrm{d}\,\frac{I_{N_2}(\lambda)\left(\alpha_{Ray}^{N_2}(\lambda)\right) - I_{He}(\lambda)\left(\alpha_{Ray}^{He}(\lambda)\right)}{I_{He}(\lambda) - I_{N_2}(\lambda)} \quad (2)$$

where *d* is the length of the cavity filled by the gas. In this study, the studied gas filled the entire length of the cavity (94.0 ± 0.1 cm) since no purge flows were used. I_{gas} is the light intensity measured by filling the cavity with high purity N_2 ($I_{N_2}(\lambda)$) and He ($I_{He}(\lambda)$). Rayleigh scattering (α_{Ray}^{gas}) is the combined product of Rayleigh scattering cross section (σ) and the gas number density (N) during the measurements. Rayleigh scattering cross sections of N_2 and He were calculated using the data in Table 1. Figure S1 shows typical examples of light intensity when the BBCES cavities are filled with pure N_2. Reflectivity measurements were repeated every three sample measurements to track the stability of the system.

Once the reflectivity is determined, it is possible to calculate the wavelength-dependent extinction cross sections of other gases ($\sigma(\lambda)$) as follows:

$$\sigma(\lambda) = \left[\left(\frac{1-R(\lambda)}{d}\right)\left(\frac{I_{He}(\lambda)-I_{gas}(\lambda)}{I_{gas}(\lambda)}\right) + \left(\frac{I_{He}(\lambda)}{I_{gas}(\lambda)}\right)\left(\alpha_{Ray}^{He}(\lambda)\right)\right]/N \quad (3)$$

Where N is the number density of the gas during the measurements, and I_{gas} (λ) is the light intensity when a target gas fills the cavity. During our experiments, the purge flow of the high reflection mirrors was shut down to ensure that the cavity was filled with target gas completely. To measure the extinction cross sections of CO_2, N_2O, and SF_6, the cavity is filled with pure target gas. Mass flow controller controlled O_2/CH_4 flow was mixed with He in a 2 m Teflon tube (Φ = ¼ inch) to generate a gas mixture with total flow rate of 500 mL min^{-1}. For O_2 experiments, measurements were performed for O_2 + He mixtures by varying the O_2 percentage between 10% and 100% with a 10% step. The CH_4, measurements were performed for CH_4 + He mixtures with CH_4 percentage ranges between 10% to 100% with a 10% step. Additional measurements were also performed for 15%, 25%, 35%, and 45% CH_4.

2.2 Extinction measurements using cavity-ring down spectroscopy (CRDS) at 404 nm and 662 nm.

To obtain independent measurements for the extinction cross sections and to cross-validate of the BBCES technique, we conducted CRD measurements at two fixed wavelengths of 404 nm and 662 nm. CRDS is a highly sensitive technique and uses a different measurement principle than BBCES. The CRDS measured the decay rate of light due to extinction rather than an absolute absorbance (as in the BBCES) and thus was immune to shot-to-shot source light fluctuations. A detailed description of the CRD method for light extinction measurement can be found in Bluvshtein et al. (2016) and He et al. (2018). Briefly, diode lasers (110 mW 404 nm diode laser, iPulse, Toptica Photonics, Munich, Germany; 120 mW 662 nm diode laser, HL6545MG, Thorlabs Inc., NJ, USA) were used as the light source of these CRDS. The 404 nm and 662 nm lasers were modulated at 1383 Hz and 500 Hz with a 50% duty cycle. The diode lasers were optically isolated by quarter waveplates (1/4 λ) and polarizing beam splitters to prevent damage to the laser head by back reflections from the highly reflective CRDS mirror. The back-reflected light beam was directed into a photodiode, which serves as an external trigger source. Light transmitted through the back mirror of the cavity was collected by an optical fiber and detected by a photomultiplier

tube (PMT), which sampled at a rate of 10 to 100 MHz. The time-dependent intensity data was acquired with a 100MHz card (PCI-5122, National Instruments, USA) and processed by data acquisition software in Labview. An exponential curve was fitted to each intensity decay data set (Figure S2). Over 1000 decay time measurements were monitored and averaged on a second basis. The residual of the fit for the averaged intensity decay was obtained and further normalized to the averaged intensity. The derived relative residuals (Figure S2) showed no apparent structure with other time constants, validating the application of CRDS as a good measure of extinction. The resultant 1 Hz decay time was averaged over one measurement duration of five minutes with standard error as the measurement uncertainty.

All of the CRDS measurements were performed under room temperature and pressure downstream from the BBCES instrument. The gas temperature (K-type thermocouple) and cavity pressure (Precision Pressure Transducer, Honeywell International Inc., MN, USA) were recorded between the two cavities for gas number density (N) calculation. During the CRDS measurements, the full cavity was filled with the investigated gases (He, CO_2, N_2O, SF_6, O_2, CH_4, or gas mixtures (O_2 + He and CH_4 + He)). The extinction cross section ($\sigma(\lambda)$) of the studied gas was measured relative to that of He and was calculated by equation (4):

$$\sigma(\lambda) = \frac{L}{clN}\left(\frac{1}{\tau_{gas}} - \frac{1}{\tau_{He}}\right) + \sigma_{He} \quad (4)$$

Where L is the total length of the cavity (l), c is the speed of light, and τ_{gas} and τ_{He} are the ring-down time of the cavity when it is filled by target gas or by the reference gas, He.

2.3 Data processing

For comparison, the scattering cross sections of the gases investigated in this study were also calculated with Equation (1) based on the refractive index and the King correction factors available in the literature that are listed in Table 1. The King correction factors were taken as unity for mono-atomic molecules and spherical molecules (with regards to the depolarization) but deviates for non-spherical molecules. For the 307–725 nm wavelength range of this study, the n-based calculated Rayleigh scattering cross sections from largest to smallest are SF_6 (Sneep and Ubachs, 2005; Wilmouth and Sayres 2020), N_2O (Sneep and Ubachs, 2005), CO_2 (Alms et al. 1975; Bideau-Mehu et al. 1973), CH_4 (Sneep and Ubachs, 2005; Wilmouth and Sayres 2020), N_2 (Bates 1984), O_2 (Bates 1984; Sneep and Ubachs, 2005), and He (Abjean et al., 1970; Cuthbertson and

Cuthbertson, 1932). Additionally, the refractive indices of SF_6, N_2O, and CH_4 were calculated based on Equation (1) using cross section results from this study and the King correction factors listed in Table 1. Our measurements were performed under ~295K and ~1020 hPa. However, the calculated refractive indices were scaled to 288.15K and 1013.25 hPa as in previous studies (Sneep and Ubachs, 2005; Wilmouth and Sayres, 2020).

The extinction of O_2 + He mixtures (α_{O_2+He}) consists of the extinction by O_2 (α_{O_2}) and He (α_{He}), and the O_2–O_2 collision-induced absorption ($\alpha_{O_2-O_2}$). The extinction of O_2 and He is a combined product of extinction cross section (σ_{gas}) and gas number density (N_{gas}). Thus α_{O_2+He} can be described with the following equation:

$$\alpha_{O_2+He} = \sigma_{O_2-O_2} \times N_{O_2}^2 + \sigma_{O_2} \times N_{O_2} + \sigma_{He} \times N_{He} \quad (5)$$

Where N_{O_2} and N_{He} are the number density of the O_2 and He in the cavities. Performing a 2rd order polynomial fit to the extinction obtained by the BBCES with respect to the gas number density thus yields the extinction cross section of O_2 and the O_2-O_2 collision-induced absorption (CIA) cross section.

In addition to the results from 2rd order polynomial fitting, we also used data from pure O_2 measurement to calculate the extinction by O_2 and by CIA of O_2–O_2. The real refractive index of O_2 (n_{O_2}) derived from extinction data measured in the wavelength regions where there is no absorption was fitted using the generalized expression of $(n_{O_2} - 1) \times 10^8 = A + \frac{B}{C - \nu^2}$. Based on the refractive index, the scattering cross sections of O_2 in the wavelength range of 307-725 nm were further calculated. By subtracting the scattering cross section of O_2 from the measured total extinction, we derived the CIA of O_2–O_2. However, the O_2 absorption bands at 580, 630, and 690 nm overlap with those of O_2–O_2 collisions. Additional corrections are thus needed to split the absorption by O_2 and O_2–O_2 collision, which is out of the scope of this study.

Methane has weak vibrational overtone absorption in the UV-vis wavelength range that is comparable to or greater than its Rayleigh scattering. Previous high-resolution spectroscopy studies have identified smooth and unstructured absorption bands across the UV-visible range (Giver, 1978; Smith et al., 1990). The spectral features are substantially broader than 0.8 nm, thus the absorption by CH_4 can be measured by our BBCES. The measured extinction coefficients of

CH_4+He mixtures (α_{CH_4+He}) are linearly correlated with the number concentration of CH_4 (N_{CH_4}) as described by the following equation:

$$\alpha_{CH_4+He} = \sigma_{CH_4} \times N_{CH_4} + \sigma_{He} \times N_{He} \quad (6)$$

A linear fit was used for deriving the extinction cross section of CH_4. The absorption between 300 and 400 nm is negligible as compared to the Rayleigh scattering. Thus extinction data in this UV wavelength range were used to calculate the real part of the refractive index of CH_4 which was further fitted utilizing the expression of $(n_{CH_4} - 1) \times 10^8 = A + \frac{B}{C - \nu^2}$. By applying this dispersion relation, the Rayleigh scattering cross sections in the entire wavelength range of 307–725 nm were derived. Finally, the CH_4 absorption cross sections were calculated by subtraction of the scattering cross section from the extinction cross section.

2.4 Error Propagation for Extinction Measurements

The uncertainty for BBCES measurements can be assessed by the propagation of the errors associated with the measurements. Each parameter (temperature, pressure, light intensity) was measured 150 times for each gas. The standard error of each parameter obtained from the 150 single measurements was used to calculate the uncertainty. The pressure (±0.01%), temperature (±0.1%) and cavity length (94.0 ± 0.1 cm) are combined with the Rayleigh cross section uncertainties for N_2 (±1%) as well as uncertainty in the measurements of the spectral signal by the spectrometer (⩽0.2%) to get an overall relative uncertainty for the effective pathlength curve of ±1.03%. This uncertainty is further propagated to the target gas by consideration of the uncertainties of pressure, temperature, and spectral intensity of the target gas measurements. The overall 1-σ uncertainty of the gas extinction cross section is 1.1%. The precision of the mass flow controllers is 0.5 mL min^{-1}. When the total flow rate is 500 mL min^{-1}, the resulting uncertainty in the gas concentration (10-100%) varies from 0% to 1.0%. Thus, the overall 1-σ uncertainty of extinction coefficients measured for CH_4+He and O_2+He varies from 1.1% to 1.5%. The detailed wavelength-dependent uncertainties were calculated due to the wavelength-dependence of the spectral intensity. The results are shown and discussed in later sections. The uncertainty for the Rayleigh scattering cross section of N_2 is validated up to 468 nm. The uncertainty above this wavelength may be larger than 1%, which is the value used for the calculation in our study. Thus, the uncertainty at wavelengths longer than 468 nm may be underestimated. Moreover, due to the

highly-structured reflectivity curve of the high-reflection mirrors, additional uncertainty could be introduced and this uncertainty can not be quantified in this study.

3 Results and Discussion

3. 1 Performance of the optical system

The reflectivity of the cavity mirrors, measured across the entire range using the difference in Rayleigh scattering of N_2 and He, was very stable throughout the experiments. The measured mirrors reflectivity curves are shown in Figure S1. The mean peak reflectivity of the $BBCES_{UV}$ mirrors was 0.999328 ±0.000006 (672 ±6 ppm) at 330 nm, with a corresponding effective optical pathlength of 1.40 ±0.01 km. The reflectivity curve of the $BBCES_{Vis}$ is much more structured, with reflectivity ranging between 0.999224 ±0.000010 and 0.9999550 ±0.0000006 (776 ±10 ppm > loss > 45 ±0.6 ppm) over a wide wavelength range of 338–725 nm. The reflectivity of the $BBCES_{Vis}$ is much higher than that of our previous system (He et al., 2018) and also covers a much broader wavelength range. Thus the effective pathlength of the BBCESVis varies between 1.3 and 20.4 km, guaranteeing a high sensitivity of the extinction measurement. The mean uncertainty in the effective pathlength across the measured wavelengths as determined from the mirror reflectivity was ±1.03%, which is predominantly due to the uncertainty in the Rayleigh scattering cross section for N_2 derived from *n*-based calculation.

3. 2 Rayleigh scattering cross sections of CO_2, N_2O, SF_6.

Figure 1 shows the extinction cross sections of CO_2, N_2O, and SF_6 measured by the BBCES. The extinction cross sections of these gases monotonically decrease with increasing wavelength, and no absorption (i.e., no structured extinction larger than the smoothly varying Rayleigh curve) is observed in the wavelength range of 307–725 nm, indicating that the measured extinction is due solely to the Rayleigh scattering of these gases. The wavelength-dependent relative standard deviations of the measurements for each gas are shown in Figure 1d. The mean 1-σ uncertainty of the reported cross sections for all three gases across the 307–725 nm wavelength range is 1.04% for CO_2, 1.05% for N_2O, and 1.04% for SF_6. As mentioned above, the derived uncertainty originates predominantly from the uncertainty in the N_2 Rayleigh scattering cross section. Uncertainty in the Rayleigh cross sections of each gas varies with wavelength and generally tracks

the light intensity spectra, which is a combined product of light source spectrum and the mirror reflectivity profile. The uncertainty is much higher when the transmitted light intensity is low (Figure S1).

The BBCES measured Rayleigh scattering cross sections for these three gases agree well with those obtained by our CRDS operating at 404 nm and 662 nm, with deviations smaller than 1.6%. Table 2 listed the Rayleigh scattering cross sections at several wavelengths obtained by the BBCES measurements (Exp) and by the calculations using the refractive index and $F_k(\nu)$ values from Table 1 (*n*-based). The relative differences between these two sets of results are within 1.4%.

Figure 1a–c shows a comparison of the measured Rayleigh scattering cross sections for CO_2, N_2O, and SF_6 with *n*-based calculations and with previous experimental results from the literature. There are a few measurements for the Rayleigh scattering cross sections for CO_2 which cover a wide spectral range (Jordan et al., 2019; Shardanand and Rao, 1977; Sneep and Ubachs, 2005; Wilmouth and Sayres, 2019; He et al., 2018). There are fewer Rayleigh scattering measurements for N_2O and SF_6 in the studied wavelength range. The measured Rayleigh scattering cross sections for CO_2, N_2O, and SF_6 are in excellent agreement with *n*-based calculation. The wavelength-dependent difference of our experimentally derived Rayleigh scattering cross sections with *n*-based calculations are shown in Figure 1e. The mean ratios of our measurements to the n-based values for the entire wavelength range of 307–725 nm are 1.00 ±0.01, 0.99 ±0.01, and 1.01 ±0.01 for CO_2, N_2O, and SF_6, respectively. The relative difference between our measurements and the *n*-based values are (0.37 ±1.24)%, (–0.55 ±1.06)%, (0.91 ±1.35)% (Mean ±SD) for CO_2, N_2O, and SF_6, respectively. Variability of the relative difference is due to structure in the mirror reflectivity that does not fully cancel. The wavelength-dependent Rayleigh scattering cross section is generally described in the form of $\sigma = A \times \lambda^B$ In this study, the measured values and the n based data were both fitted to this function. The relative difference between these two fitted functions is shown in Figure 1(f). That would be a measure of the uncertainty comparing smooth functions to smooth functions. The relative differences were (0.49 ±0.48)%, (-0.41 ±0.30)%, and (0.94 ±0.22)% (Mean ±stdev), for CO_2, N_2O, and SF_6, respectively. The mean values of the relative difference obtained from the fitting function are close to that obtained from the measurements. However, the variabilities are much smaller, which may be related to the cancellation of the influence by the structured mirror reflectivity. Notably, while our results for N_2O agree well with the *n*-based

calculations, previous results obtained by CRDS at 532 nm (Sneep and Ubachs, 2005) and by absorption spectroscopy in the wavelength of 300–315 nm (Bates and Hays, 1967) do not agree well with the n-based calculations. The measurements between 300 and 315 nm were first published by Bates and Hays (1967), who obtained the results from a doctoral thesis. However, the results from our BBCES system are in good agreement with the *n*-based calculations and with experimental results from independent CRDS measurements, thus increasing the confidence in our measured values.

3. 3 Scattering and absorption cross sections of O_2.

The UV-vis spectra of gas-phase molecular oxygen are characterized by discrete structured absorption bands due to the electronic transition ($b^1\sum_g^+(\nu' = 1/2/3) \leftarrow \sum_g^-(\nu'' = 0)$) of O_2 monomer, broader unstructured CIA of O_2–O_2, and structured dimer bands from the bound van der Waals O_2 dimer (Newnham and Ballard, 1998). Under atmospheric conditions, the O_2–O_2 CIA bands are frequently described as "O_4" bands, although absorption by O_2 dimer is thought to be significant only under very low-temperature conditions (Thalman and Volkamer, 2013; Long and Ewing, 1973). Within the wavelength range investigated in this work, the molecular oxygen B band at 688 nm ($b^1\sum_g^+(\nu' = 1) \leftarrow X^3\sum_g^-(\nu'' = 0)$), γ overtone band at 629 nm ($b^1\sum_g^+(\nu' = 2) \leftarrow X^3\sum_g^-(\nu'' = 0)$), and δ overtone band at 580 nm ($b^1\sum_g^+(\nu' = 3) \leftarrow X^3\sum_g^-(\nu'' = 0)$ overlap with O_2–O_2 CIA bands of $^1\sum_g^+(\nu = 1)$, $^1\Delta_g + {}^1\Delta_g\ (\nu = 0)$, and $^1\Delta_g + {}^1\Delta_g\ (\nu = 1)$, respectively. These absorption bands can only be resolved by a high-resolution spectroscopic technique. Absorption cross sections of the B, γ, and δ bands were convoluted from the HITRAN database (Gordon et al., 2017) by considering the temperature, pressure, and wavelength resolution of the instrument. The wings of the oxygen lines also show a quadratic dependence on the pressure due to pressure broadening. However, due to the minimal O_2 absorption contribution below 680 nm and the low instrument wavelength resolution, the extinction cross section of the O_2 monomer can be treated as linearly correlated with the O_2 concentration. Moreover, the O_2–O_2 CIA cross section is correlated with the square of the O_2 concentration. Therefore, these cross sections can be retrieved from measurements at different O_2 concentrations. Due to the discrete structured absorption bands and the wavelength resolution of the instrument, the range of absorption cross sections spans several orders of magnitude within the spectral response of the instrument, limiting the relevance of the absorption cross sections for other researchers. These results are not further

discussed here. However, the data for broader unstructured CIA of O_2–O_2 are still useful for various applications.

Figure 2 shows the wavelength-dependent extinction coefficients of O_2+He mixtures. He was used in these experiments to minimize extinction contributions from Rayleigh scattering. Nine absorption peaks centered at 344 nm (CIA), 360 nm (CIA), 380 nm (CIA), 446 nm (CIA), 477 nm (CIA), 532 nm (CIA), 577 nm (δ overtone and CIA), 629 nm (γ overtone and CIA), and 688 nm (B band and CIA) were observed in the wavelength range of 307–725 nm. The absorption coefficients of the central wavelengths for the first eight peaks increase non-linearly with O_2 concentration while that of the 688 nm peak increases in a more linear manner, indicating that the O_2 B band absorption dominates the last absorption peak while the other peaks are mostly associated with CIA of O_2-O_2.

The extinction coefficients obtained by the BBCES correlated well with those measured by the CRD, with slops of 0.990 (R^2=0.9994) and 0.993 (R^2 = 0.9996) at the wavelengths of 404 nm and 662 nm, respectively (Figure 3). This excellent agreement between the instruments further substantiates the BBCES measurements and suggests that the accuracy of the BBCES at these two wavelengths is better than estimated in the error propagation above, where the N_2 refractive index was the largest uncertainty. As explained in the data processing section, the measured extinction coefficients were fitted with a 2rd order polynomial (selected wavelengths at the peaks of the CIA absorption bands are shown in Figure 4). At 476.7, 577.2, and 629.2 nm, the absorption is from the CIA of O_2–O_2. The fit generates positive values matching the absorption cross section of O_2–O_2 CIA. At 687.7 nm where strong B-band absorption appears, the fit yields a small negative coefficient for O_2–O_2 CIA.

Figure 5a shows the extinction cross section measured for 100% O_2. These results agree well with previously reported results by Jordan et al. (2019). For wavelengths where no absorption is detected, the measured extinction cross sections agree well with *n*-based calculations. Figure 5b-c shows the determined extinction cross sections for molecular O_2 and the absorption cross sections of O_2–O_2 CIA. For wavelength ranges without O_2 bands, our extinction cross sections agree well with the *n*-based values with an average deviation of (2.81 ±1.21)%. The absorption cross sections for O_2–O_2 CIA derived in this study mostly agree well with literature data from Thalman and Volkamer (2013). The differences are within 1.1% at 477, 532, 577, and 630 nm but larger

deviations were found at 344 (4.2%), 360 (-29%), 380 (-21%), and 446 (4.2%) nm. These absorption bands are the lowest intensity bands and therefore have the largest relative uncertainties in either measurement. Moreover, the absorptions at 344, 360, 380, and 446 nm contribute a much smaller fraction of the extinction as compared to that of 477, 532, 577, and 630 nm. Thus larger discrepancies were observed during the apportionment of absorption from extinction.

The Rayleigh scattering cross sections of molecular O_2 derived from the 100% O_2 measurement agree well with *n*-based calculations with an average difference of 1.2%. CIA of O_2–O_2 calculated from this single measurement matches the results from the fitting method. Due to strong absorption from O_2 B band and γ overtone band, this method cannot derive the cross sections of CIA of O_2-O_2 at 630 and 688nm.

3. 4 The scattering and absorption cross sections of CH_4.

CH_4 has weak absorption in the UV–vis wavelength range, and these bands dominate the photographic spectra of planets such as Uranus and Neptune (Adel and Slipher, 1934). Figure 6 presents the wavelength-dependent extinction coefficients of CH_4+He mixtures. A total of eleven absorption bands were detected in the wavelength range of 307–725 nm. The extinction coefficients increase as a function of increasing CH_4 concentration. Extinction coefficients obtained by the BBCES correlated well with those measured in parallel by the CRDS, with slopes of 1.002 (R^2=0.9999) and 0.99 (R^2 = 0.999) at the wavelengths of 404 nm and 662 nm (Figure S3). The excellent agreement between these three systems further supports the accuracy of BBCES extinction measurements over a wide working range. The measured extinction coefficients were linearly fit against the CH_4 number concentration. Figure 7 shows the fitted curves at five selected wavelengths. The extinction coefficients have a linear correlation with CH_4 concentration (R^2 > 0.9988) without exception. The calculated slopes represent the extinction cross sections of CH_4 and also indicate a wide dynamic range of our BBCES.

The extinction cross sections for CH_4 retrieved from concentration-dependent measurements are plotted in Figure 8a. BBCES results from this study agree well with results from previous studies using BBCES (Jordan et al., 2019; Wilmouth and Sayres, 2019) and CRDS (Sneep and Ubachs, 2005). Previous studies using a Nephelometer (Shardanand and Rao, 1977) and interferometer (Cuthbertson and Cuthbertson, 1920; Watson et al., 1936) obtained the scattering cross sections and the refractive index of CH_4. The BBCES measures the extinction cross section. For

wavelengths where extinction is dominated by Rayleigh scattering (< 475 nm), our BBCES results agree well with the results from Nephelometer and interferometer measurements. In this study, the refractive index of CH_4 was calculated using the extinction data in the wavelength range of 307-400 nm. The calculated refractive index was fitted to the general expression:

$$(n_{CH_4} - 1) \times 10^8 = 5476 + \frac{4.1579 \times 10^{14}}{1.1568 \times 10^{10} - \nu^2} \quad (7)$$

As shown in Figure 8b, our calculated scattering cross sections are in good agreement with those derived from the newest refractive index developed by Wilmouth and Sayres (2020) (Table 2), with an average difference of (0.89 ±2.18)%. The absorption cross section, which is the difference between the total extinction and the Rayleigh scattering cross section, is shown in Figure 8c. At most spectral ranges, our results are in better agreement with the results from previous studies by Giver, (1978) and Smith et al.(1990). For example, the difference as compared to the results from Giver (1978) at 542, 576.4, 598, 619, 665.7, and 703.6 nm is 4.0% on average. At several wavelength regions (e.g., 520–536nm, 580–605 nm), the results from Fink et al. (1977) differ from all of the other studies. In the wavelength range of 400–725 nm, absorption contributes up to 99.7% of the CH_4 extinction.

3. 5 Dispersion relations for N_2O, SF_6 and CH_4.

SF_6: Wilmouth and Sayres (2020) found that their measured Rayleigh scattering cross sections for SF_6 in the ultraviolet range were lower than those from the *n*-based expression of Sneep and Ubachs (2005). They generated the dispersion formula for SF_6 from the combined fit using refractive index data in the wavelength range of 264-297 nm and 333-363 nm by Wilmouth and Sayres (2020, 2019), and direct refractive index measurement at 632.99 nm (Vukovic et al., 1996) (Figure 9a). In our study, the refractive index of SF_6 in the wavelength range of 307-725 nm was calculated from the measured Rayleigh scattering cross section for 288.15 K and 1013.25 hPa. To better constrain the dispersion formula when extrapolated over a broad wavelength range, we employed an alternative fit of the form $A+B/(C–\nu^2)$ to our data and the data used by Wilmouth and Sayres (2020) for fitting. All sets of data were weighted equally. The resulting dispersion relation for SF_6 in the wavelength range of 264–725 nm is

$$(n_{SF_6} - 1) \times 10^8 = 18997.7 + \frac{8.27663 \times 10^{14}}{1.56833 \times 10^{10} - \nu^2} \quad (8)$$

N_2O: Sneep and Ubachs (2005) derived the refractive index based on polarizability measurements using interferometer at five single wavelengths (457.9, 488, 514.5, 568.2, 647.1 nm) by Alms et al. (1975). In this study, we calculated the refractive index of N_2O from the Rayleigh scattering cross sections in the wavelength range of 307–725 for 288.15K and 1013.25 hPa. Based on this refractive index data set, the dispersion relation (Eq (9)) for N_2O was retrieved for a much broader wavelength range (Figure 9b) compared to that generated by Sneep and Ubachs (2005).

$$(n_{N_2O} - 1) \times 10^8 = 22095 + \frac{1.66291 \times 10^{14}}{6.75226 \times 10^9 - \nu^2} \quad (9)$$

CH_4: The previous study by Wilmouth and Sayres (2019) has shown that their measured Rayleigh scattering cross sections for CH_4 are in substantial disagreement (22%) with those calculated from the refractive index recommended by Sneep and Ubachs (2005). Sneep and Ubachs (2005) formulated the refractive index of CH_4 based on interferometric measurements at wavelengths of 325, 543.5, 594.1, 612, and 633 nm by Hohm (Hohm, 1993). However, the Rayleigh scattering cross sections calculated from their refractive index are much higher than all the measured values listed in Figure 9b. Using equally weighted Rayleigh scattering cross sections data sets in the wavelength range of 264–297 nm, 333–363 nm (Wilmouth and Sayres, 2019, 2020), 307–400 nm from this study, and single wavelength measurements that are not impacted by absorption (Cuthbertson and Cuthbertson, 1920; Watson et al., 1936), we derived the dispersion formula for the refractive index of CH_4 in the combined UV/visible range (Figure 9c) as follows:

$$(n_{CH_4} - 1) \times 10^8 = 3603.09 + \frac{4.40362 \times 10^{14}}{1.1741 \times 10^{10} - \nu^2} \quad (10)$$

The calculated Rayleigh scattering cross sections using the dispersion relations derived in this study were compared with those derived from previously recommended formulations listed in Table 1 (Figure 9). The difference increases significantly towards the longer wavelength in the region of 320–725 nm (Figure S4). The average deviations are 0.1%, 0.9%, and 0.1% for SF_6, N_2O, and CH_4, respectively. Notably, the difference for N_2O is more significant than for the other two gases. This study uses refractive index data in the continuous wavelength ranges of 307–725 nm to derive the dispersion relation, while the formulation for N_2O in Table 1 is derived by Sneep and Ubachs (2005) based on polarizability measurements at five single wavelengths. For the formulation of the refractive index of CH_4, Wilmouth and Sayres (2020) weighted the data sets from Watson and Ramaswamy (1936) and Cuthbertson and Cuthbertson (1920) equally but gave

more weight to their UV measurements when deriving the formulation of the refractive index. In this study, all the CH_4 data set were weighted equally. The derived dispersion relation agrees very well with that from Wilmouth and Sayres (2020), as shown in Figure 9 (c-d).

Conclusions and Implications

Rayleigh scattering cross sections between 307 and 725 nm were determined for CO_2, N_2O, SF_6, O_2, and CH_4 by simultaneous BBCES and CRDS measurements. Extinction coefficients obtained by the BBCES show high consistency with those measured by parallel CRDS at 404 and 662 nm (Figure 3 and figure S3), demonstrating that the BBCES measurements provide results with both a wide wavelength range and high accuracy. Comparison of our measurements with the *n*-based calculations for these gases in the entire wavelength range of this study yields excellent agreement with relative differences of (0.37 ±1.24)%, (–0.55 ±1.06)%, (0.91 ±1.35)%, (2.81 ±1.21)%, and (0.89 ±2.18)%, respectively. The O_2-O_2 CIA cross sections obtained from the BBCES measurements are compared with those published by Thalman and Volkamer (2013). The relative differences are within 1.1% at 477, 532, 577, 630 nm. Larger relative differences occur at the weak bands at 344 (4.2%), 360 (–29%), 380 (–21%), and 446 (4.2%) nm. The absorption cross sections of CH_4 in the wavelength range of 400-725 nm agree well with those documented by Giver (1978).

Rayleigh scattering cross sections of CO_2 determined using BBCES and CRDS in this study, and in other studies have shown that the refractive index recommended by Sneep and Ubachs (2005) is suitable for use in the wavelength range of 307–725 nm. By incorporating the refractive index data from previous studies, we developed new dispersion relations for the refractive index of N_2O (307-725 nm), SF_6 (264–725 nm), and CH_4 (264–671 nm). The derived dispersion relations for SF_6 and CH_4 agree well with those provided by Wilmouth and Sayres (2020).

Previous studies measured the Rayleigh scattering and absorption cross sections of CO_2, N_2O, O_2, SF_6, and CH_4 at narrow spectral ranges or single wavelengths. In this study, we used BBCES that covers the broad wavelength range of 307–725 nm to measure total extinction (the sum of absorption and scattering). The measurements validate that refractive index-based methods for calculating Rayleigh extinction cross sections are accurate and provide new fits over more continuous and extended wavelengths range than available in the literature to constrain such methods. The Rayleigh scattering cross sections reported here are useful in several applications.

These include calibration standards based on extinction for optically-based instruments, such as those designed for aerosol optical properties measurements or trace gas concentrations in the field (Jordan et al., 2019; Min et al., 2016; Bluvshtein et al., 2017), especially when high-refractive index gases are used for improved calibration. They will also improve the accuracy of Rayleigh scattering parameterizations for major greenhouse gases in Earth's atmosphere, CO_2, CH_4, and N_2O. Accurate quantitative measurements of Rayleigh scattering coefficients and absorption cross sections of atmospheric gases such as molecular N_2, O_2, CO_2 and the CIA of O_2–O_2 cross sections in the UV-NIR range are of particular importance for the application of Rayleigh LIDAR systems, especially at the Nd:YAG laser harmonics 1064, 532 & 366 nm. These systems analyze the molecular backscattering contributions to the LIDAR's attenuated backscatter signals to retrieve the atmospheric profile of aerosols and clouds in the planetary boundary layer (Tomasi et al., 2005; Herron, 2007). Recent NASA satellite missions have also aimed to measure global carbon dioxide concentrations with high precision (0.25%) (Drouin et al., 2017). These CO_2 global missions use the O_2–O_2 CIA underneath the structured O_2 A-band (760 nm) to evaluate the solar radiation double pathlength in the Earth atmosphere and to determine the atmospheric pressure. The measurements in this study validate the existing literature on the extinction of O_2 collision complexes and molecular oxygen bands, and can be used for calibration purposes in both remote sensing and *in-situ* spectroscopic applications in the atmosphere. In the future, gas extinction measurements at extended wavelengths (near-infrared) and for additional gases (e.g., N_2) will improve the spectroscopic applications in atmospheric studies.

Data availability.

Data are available upon request from the corresponding author (yinon.rudich@weizmann.ac.il).

Author contributions.

Q.H., S.S., and Y.R. designed this study. Q.H., Z.F., and O.S. conducted the experiments. Q.H. prepared the draft and all of the co-authors reviewed it and provided comments.

Competing interests.

The authors declare that they have no conflict of interest.

Acknowledgments

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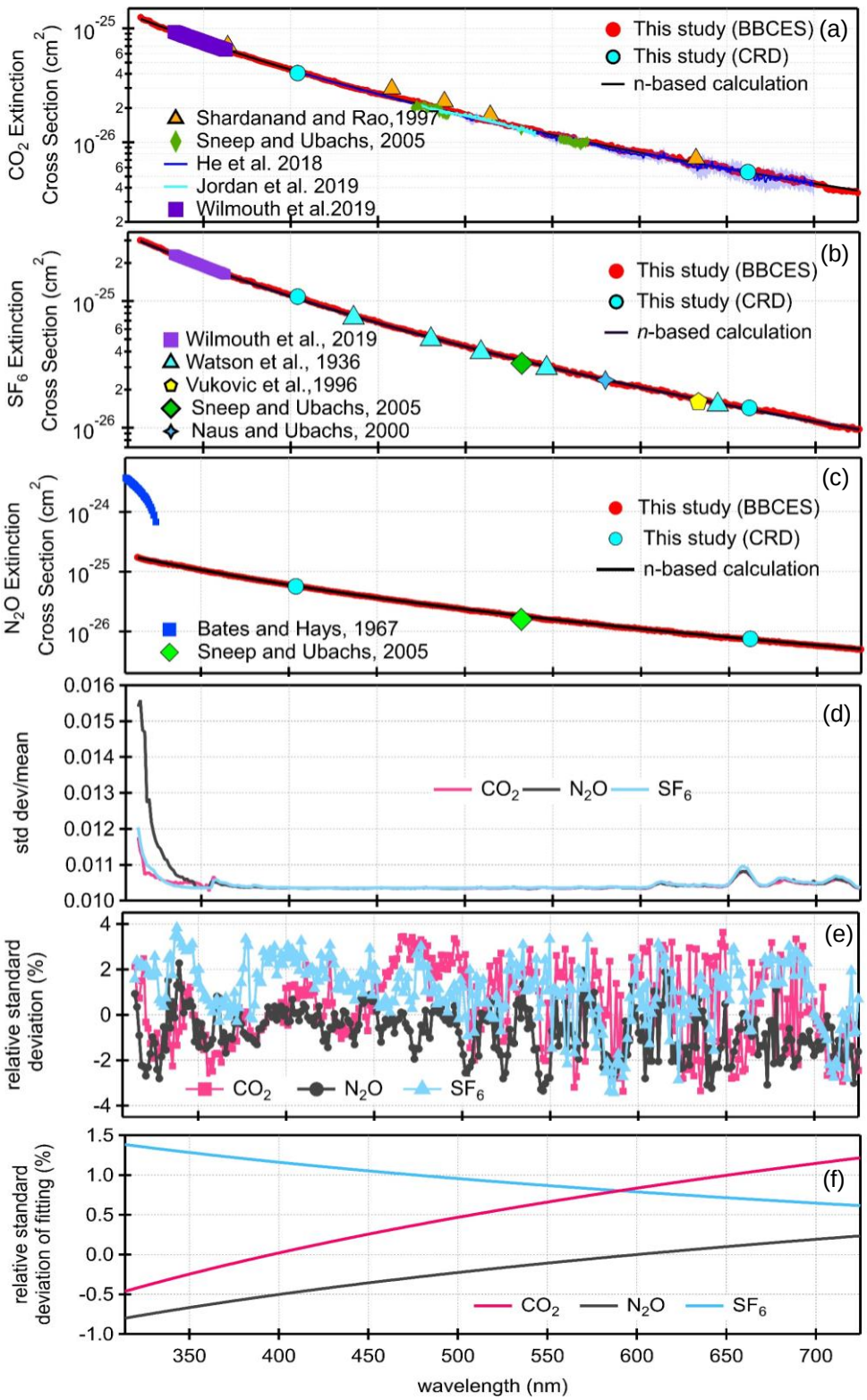


Figure 1. Rayleigh scattering cross sections of CO_2 (a), SF_6 (b), and N_2O (c). Panel (d) shows the relative standard deviations as a function of wavelength for each gas. The relative difference in the cross sections obtained by our measurements and calculations from the refractive index are displayed (e). Panel (f) shows the relative difference after fitting ($\sigma=A\lambda^B$).

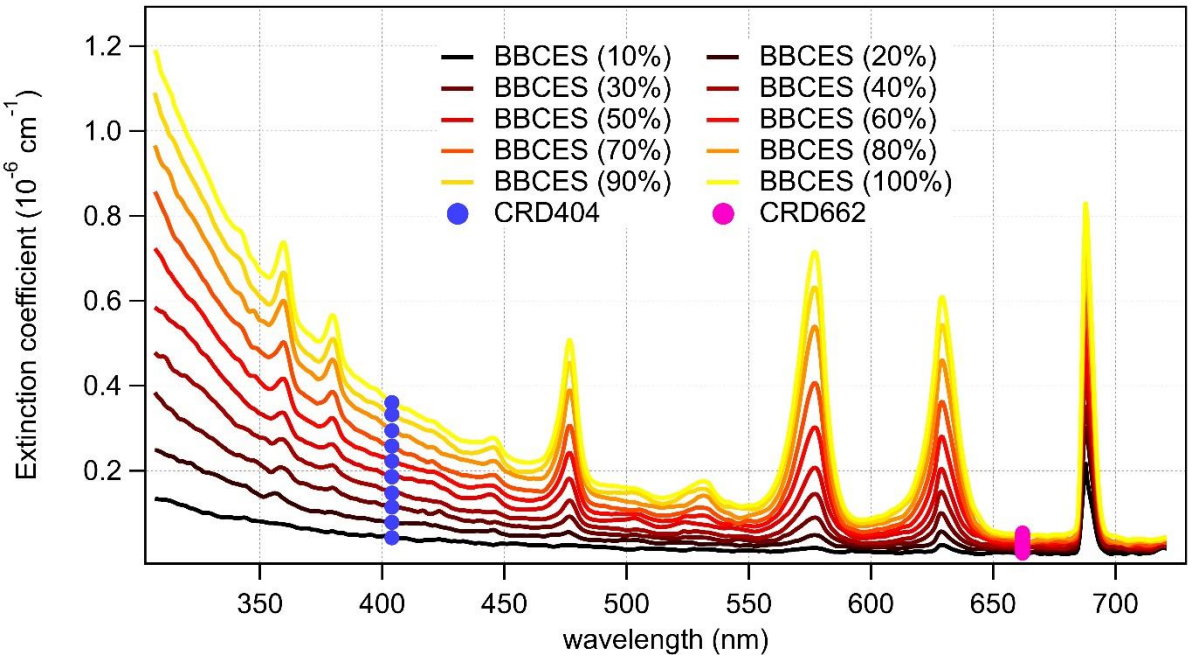
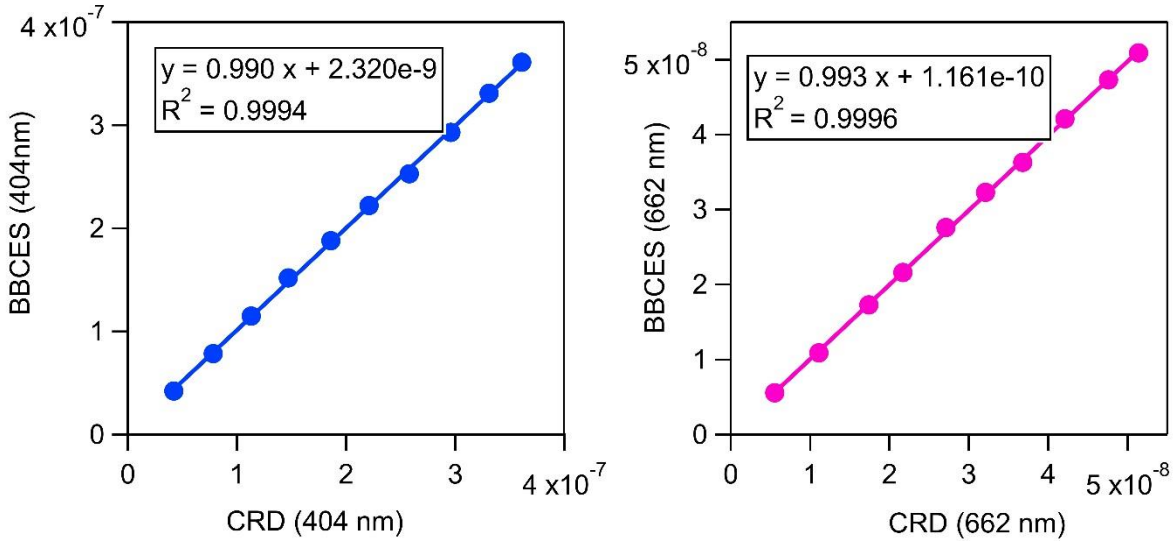


Figure 2. Wavelength-dependent extinction coefficients of O_2 + He mixtures as a function of O_2 concentration. The colored lines represent the extinction coefficients measured by BBCES, and markers represent results from CRDS.



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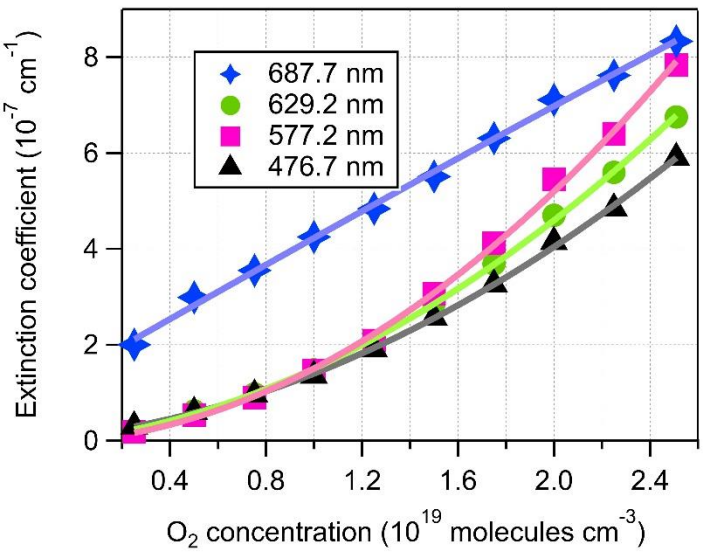
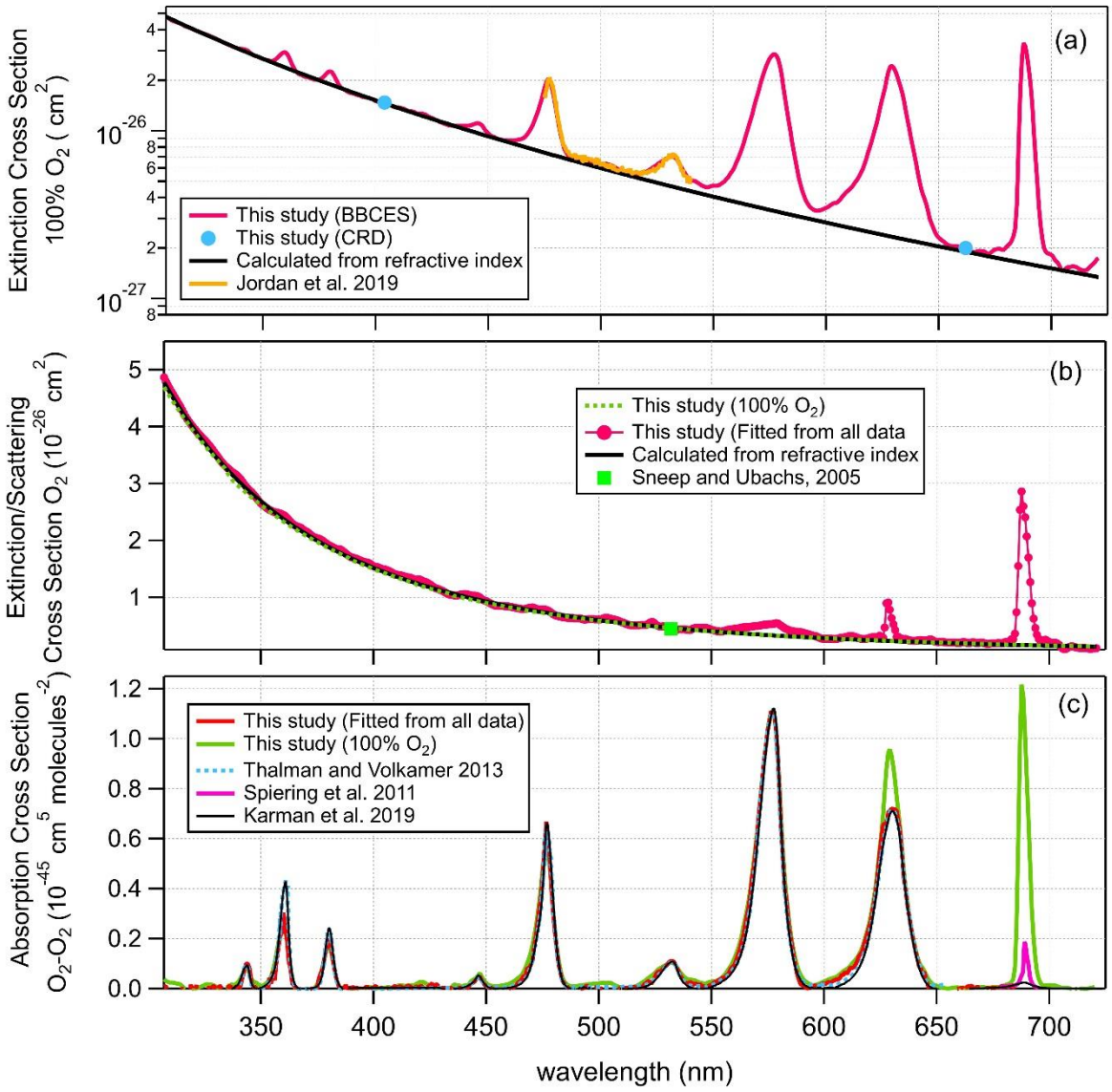
536 Figure 3. Correlations between the extinction coefficients (unit, cm^{-1}) measured by the BBCES
537 and CRDS.



wavelength	a	b	c
687.7 nm	-6.84e-48	2.86e-26	1.34 e-8
629.2 nm	7.29e-46	7.80e-27	-1.09e-10
577.2 nm	1.10e-45	5.27e-27	-1.41e-9
476.7 nm	6.66e-46	7.68e-27	1.53e-9

538

539 Figure 4. 2rd order polynomial fit of extinction coefficients measured by the BBCES. The O_2
540 concentration-dependent extinction coefficients are contributed by the extinction coefficients of
541 O_2 (σ_{O_2}), He (σ_{He}), and the O_2-O_2 CIA cross sections ($\sigma_{O_2\text{-}O_2}$).



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543 Figure 5. Wavelength-dependent extinction cross sections of 100% O_2 (a), extinction cross
544 sections of O_2 (b), and O_2-O_2 CIA cross section (c).

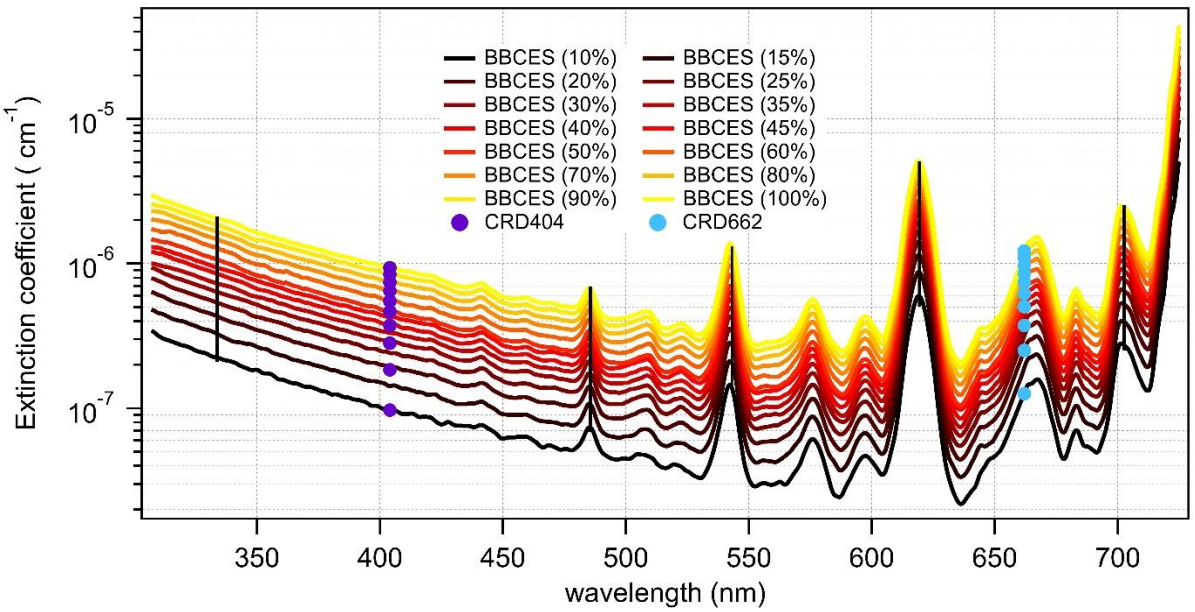
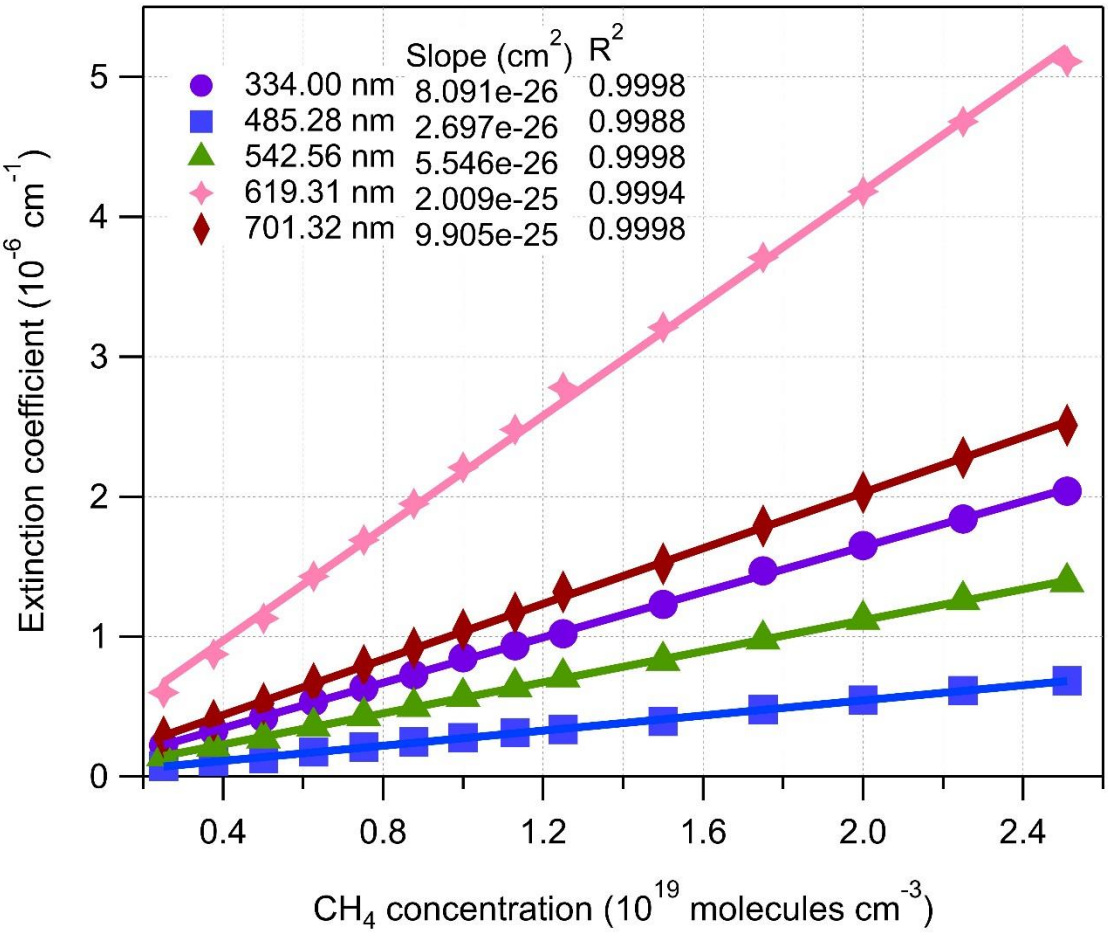


Figure 6. Wavelength-dependent extinction coefficients of CH_4 + He mixtures as a function of CH_4 mixing ratio. The colored lines represent extinction coefficients obtained from BBCES and markers represent results from CRDS. Measurements were performed with CH_4 percentage within 10% and 100% with a 10% step. Moreover, BBCES measurements were also performed for 15%, 25%, 35%, and 45% CH_4. The number concentration of 100% methane was 2.50143×10^{19} molecules cm^{-3}. Data at selected wavelengths (vertical lines) are shown in Figure 7.



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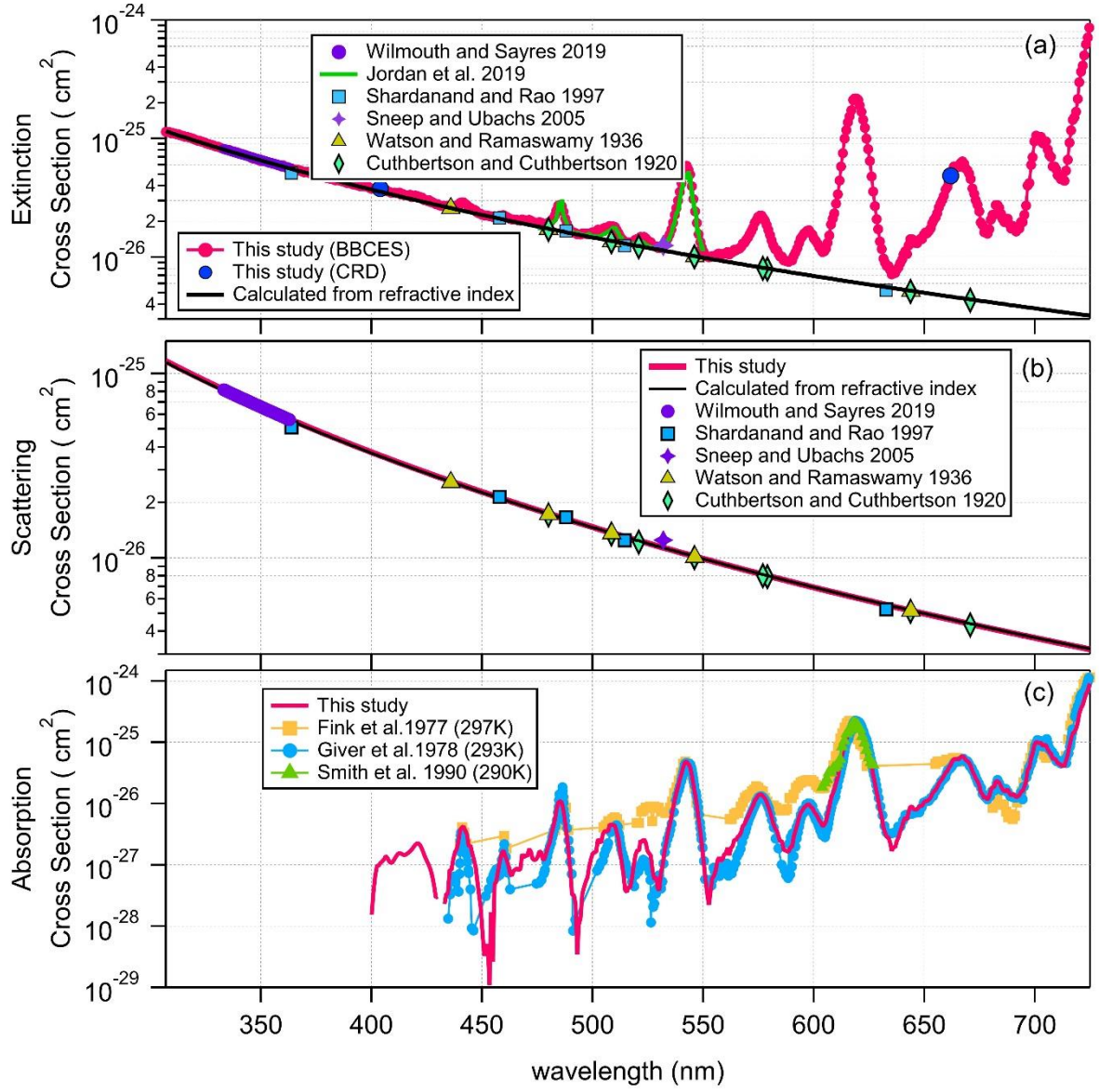
553 Figure 7. The relationship between BBCES measured extinction coefficients of CH_4+He mixtures
554 and CH_4 concentration. The selected wavelengths were located in Figure 6 by vertical lines.



555

556 Figure 8. Extinction (a), scattering (b), and absorption (c) cross sections of CH_4.

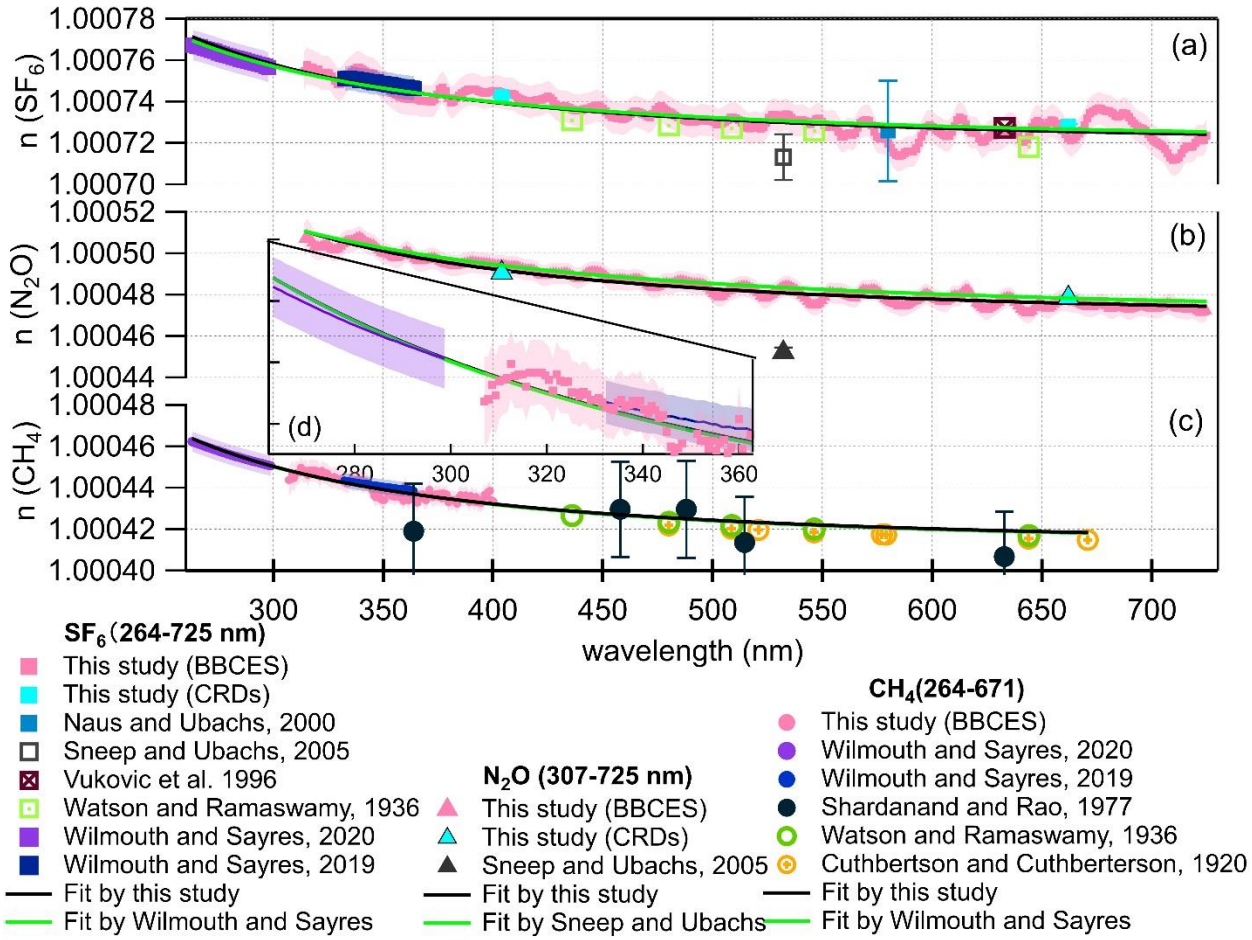


Figure 9. Real refractive index (n) for SF_6 (a), N_2O (b), and CH_4 (c). Comparison of Refractive index from this work with previous studies (Cuthbertson and Cuthbertson, 1920; Naus and Ubachs, 2000; Shardanand and Rao, 1977; Sneep and Ubachs, 2005; Vukovic et al., 1996; Watson et al., 1936; Wilmouth and Sayres, 2019, 2020) over the wavelength range of 264–725 nm. The green line represents the dispersion relation given in Table 1. The black line represents the dispersion relation given in Eq. (8–10) derived from a fit to our data and references results. The shading represents 1-σ uncertainty of the n. The n values for Shardanand and Rao (1977), Sneep and Ubachs (2005), Naus and Ubachs (2000) were calculated from their reported Rayleigh scattering cross sections. Refractive index data from Sneep and Ubachs (2005) are not used in the fitting since these results are away from others. Data from Shardanand and Rao (1977) are not used due to large uncertainties. All of the data sets are equally weighted during fitting. Panel (d) is a close-up view of panel (c) in the wavelength range of 264–363 nm.

Table 1. Refractive index and King correction factors for calculating Rayleigh scattering cross-sections and available measurements in the wavelength range of 300–725 nm. Measurements for He and N_2 are not summarized in this table.

Gas	Refractive index and King correction factors				Measurements	
	(n-1) $\times 10^8$	$F_k(\nu)$	ν (cm^{-1})	References	λ (nm)	References
He	$2283 + \frac{18102 \times 10^{13}}{1.5342 \times 10^{10} - \nu^2}$	1.0	14285-33333	Abjean, 1970; ***Leonard, 1974;*** Cuthbertson, 1932		
N_2	$5677.465 + \frac{318.81874 \times 10^{12}}{1.44 \times 10^{10} - \nu^2}$	$1.034 + 3.17 \times 10^{-12}\ \nu^2$	21360-39370	***Bates 1984*** Sneep, 2005; Naus, 2000		
CO_2	$\times \left(\frac{1.1427 \times 10^{11}}{5799.25 - \nu^2} + \frac{120.05}{(89223.8)^2 - \nu^2} + \frac{5.3334}{(75037.5)^2 - \nu^2} + \frac{4.3244}{(67837.7)^2 - \nu^2} + \frac{1.218145 \times 10^{-5}}{(2418.136)^2 - \nu^2}\right)$	$1.1364 + 2.53\ \times 10^{-11}\ \nu^2$	39417-55340	*Alms, 1975*; **Bideau-Mehu, 1973**; Sneep, 2005	333-725	Jordan, 2019; Shardanand, 1977; Sneep, 2005; Wilmouth, 2019; He, 2018
CH_4	$4869.8 + \frac{4.1023 \times 10^{14}}{1.133 \times 10^{10} - \nu^2}$	1.0	15385-40000	*Sneep, 2005*; **Wilmouth, 2020**	264-297 333-363, 434-725	Cuthbertson 1920; Jordan, 2019; Shardanand, 1977; Sneep, 2005; Watson, 1936; Wilmouth, 2019;2020; Smith, 1990; Giver, 1978; Fink, 1977
N_2O	$46890 + 4.12 \times 10^{-6}\ \nu^2$	$\frac{3.3462 + 70.8 \times 10^{-12}\ \nu^2}{2.7692 - 47.2 \times 10^{-12}\ \nu^2}$	15453-21838	Alms, 1975; ***Sneep, 2005***	300-320, 532	Johnston, 1975; Sneep, 2005
SF_6	$18611.4 + \frac{8.9566 \times 10^{14}}{1.680 \times 10^{10} - \nu^2}$	1.0	15385-40000	*Sneep, 2005*; Vukovic, 1996; **Wilmouth, 2020**	264-297 333-363, 532, 633	Sneep, 2005; Vukovic, 1996; Wilmouth, 2019,2020
O_2 [a]	$20564.8 + \frac{2.480899 \times 10^{13}}{4.09 \times 10^{9} - \nu^2}$	$1.09 + 1.385 \times 10^{-11}\ \nu^2 + 1.448 \times 10^{-20}\ \nu^4$	18315-34722	**Bates 1984**; Hohm, 1993; *Sneep, 2005*	328-667	Thalman, 2013; Jordan, 2019; Hermans, 1999; Greenblatt, 1990; Spiering, 2011

Unless noted, the refractive index is scaled to 288.15 K and 1013.25 hPa. N = 2.546899 $\times 10^{19}$ molecules cm^{-3}.

Due to limited space, only the first name of each reference is shown in the table.

The references in bold and italics describe the formulation of refractive index and King correction factor for *n*-based calculation, respectively.

[a] The refractive index was obtained at 273.15 K and 1013.25 hPa, N = 2.68678 $\times 10^{19}$ molecules cm^{-3} is used in Eq. (1)

580 Table 2. The Rayleigh scattering cross sections (10^{-27} cm^2) calculated from the refractive index
581 (*n*-based) and obtained from BBCES (Exp) of selected wavelengths.

λ(nm)	CO_2		SF_6		N_2O		O_2		CH_4	
	n-based	Exp	*n*-based	Exp	*n*-based	Exp	*n*-based	Exp	*n*-based	Exp
330	98.22	96.8	241.5	239.4	137.9	136.7	34.71	35.1	84.12	85.3
404	41.67	41.6	104.5	105.7	57.71	57.9	14.57	14.8	35.57	35.9
532	13.32	13.3	33.92	34.1	18.19	18.3	4.642	4.55	11.34	11.3
660	5.516	5.52	14.16	14.2	7.483	7.47	1.924	1.95	4.693	4.68
710	4.101	4.08	10.55	10.4	5.551	5.48	1.430	1.41	3.487	3.47

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