



Microphysical investigation of the seeder and feeder region of an Alpine mixed-phase cloud

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Abstract. The seeder-feeder mechanism has been observed to enhance orographic precipitation in previous studies. However, the microphysical processes active in the seeder and feeder region are still being understood. In this paper, we investigate the seeder and feeder region of a mixed-phase cloud passing over the Swiss Alps, focusing on (1) fallstreaks of enhanced radar reflectivity originating from cloud top generating cells (seeder region) and (2) a persistent low-level feeder cloud produced by the boundary layer circulation (feeder region). Observations were obtained from a multi-dimensional set of instruments including ground-based remote sensing instrumentation (Ka-band polarimetric cloud radar, microwave radiometer, wind profiler), in situ instrumentation on a tethered balloon system and ground-based aerosol and precipitation measurements.

The cloud radar observations suggest that ice formation and growth was enhanced within cloud top generating cells, which is consistent with previous observational studies. However, uncertainties exist regarding the dominant ice formation mechanism within these cells. Here we propose different mechanisms that potentially enhance ice nucleation and growth in cloud top generating cells (convective overshooting, radiative cooling, droplet shattering) and attempt to estimate their potential contribution from an ice nucleating particle perspective. Once ice formation and growth within the seeder region exceeded a threshold value, the mixed-phase cloud became fully glaciated.

Local flow effects on the lee side of the mountain barrier induced the formation of a persistent low-level feeder cloud over a small-scale topographic feature in the inner-Alpine valley. In situ measurements within the low-level feeder cloud observed the production of secondary ice particles likely due to the Hallett-Mossop process and ice particle fragmentation upon ice-ice collisions. Therefore, secondary ice production may have been partly responsible for the elevated ice crystal number concentrations that have been previously observed in feeder clouds at mountain-top observatories. Secondary ice production in feeder clouds can potentially enhance orographic precipitation.



20 1 Introduction

Mixed-phase clouds (MPCs), which consist of ice crystals and supercooled cloud droplets, play a crucial role for precipitation formation and are responsible for 30 % to 50 % of the precipitation in the mid-latitudes (Mülmenstädt et al., 2015). Furthermore, MPCs have important implications for the Earth's radiation budget. In particular, the phase partitioning between the liquid and ice phase in MPCs is of major importance, as the radiative properties of ice crystals and cloud droplets differ significantly (Sun and Shine, 1994). Thus, in order to understand the radiative effects and precipitation initiation in MPCs, it is important to understand the microphysical processes that govern MPCs as well as to characterize the vertical distribution of the liquid and ice phase hydrometeors within them.

The coexistence of the ice and liquid phase in MPCs is thermodynamically unstable due to the lower saturation vapor pressure over ice compared to over liquid. Therefore, ice crystals grow rapidly at the expense of the surrounding water droplets if the saturation vapor pressure lies between ice and water saturation. This process is known as the Wegener-Bergeron-Findeisen (WBF) process (Wegener, 1911; Bergeron, 1935; Findeisen, 1938) and can lead to rapid glaciation of the cloud and thus limits the lifetime of MPCs.

In order to sustain mixed-phase regions, two prerequisites need to be fulfilled. Firstly, the environment needs to be supersaturated with respect to liquid water, which can be achieved through sufficiently large updrafts (e.g., Rauber and Tokay, 1991; Harrington et al., 1999). Secondly, the condensate supply rate needs to exceed the diffusional growth rate of the ice crystals. Indeed, persistent MPCs are frequently observed in mountainous regions (e.g., Borys et al., 2003; Lowenthal et al., 2011; Dorsi et al., 2015; Lloyd et al., 2015; Lohmann et al., 2016a; Beck et al., 2017; Lowenthal et al., 2016; Lowenthal et al., 2019) where the local topography produces updrafts capable of providing a continuous source of condensate. In addition, Rauber and Grant (1986) found two further regions where the prerequisites for persistent MPCs are fulfilled: near cloud top and near cloud base. The presence of a supercooled liquid layer at cloud top can increase radiative cooling (e.g., Sun and Shine, 1994; Possner et al., 2017; Eirund et al., 2019). Furthermore, this liquid layer can act as a source region for primary ice nucleation and initial ice growth (i.e., seeder region) and can influence the evolution of the microphysical cloud structure in the lower cloud levels. Meanwhile, the presence of a supercooled liquid layer near cloud base can act as a feeder region on which precipitation particles that formed in the seeder region of the cloud can "feed" on, ultimately enhancing precipitation (e.g., Reinking et al., 2000; Borys et al., 2000; Borys et al., 2003; Lowenthal et al., 2011; Lowenthal et al., 2016; Lowenthal et al., 2019).

Seeder regions were often observed in connection with cloud top generating cells (e.g., Hogan et al., 2002; Crouce et al., 2007; Stark et al., 2013; Kumjian et al., 2014; Rosenow et al., 2014; Plummer et al., 2014; Plummer et al., 2015; Rauber et al., 2015). The term "generating cell" describes a small region of enhanced radar reflectivity at cloud top, which produces an enhanced reflectivity trail, or fallstreak, characteristic of falling hydrometeors. Generating cells have horizontal extents of 1 - 2 km and updraft velocities in the range of 0.75 - 3 m s⁻¹ (Rosenow et al., 2014; Kumjian et al., 2014). Most studies agree that radiative cooling at cloud top is a major driver for the formation and maintenance of generating cells (e.g., Kumjian et al., 2014; Keeler et al., 2016) and that these cells play an important role in primary ice nucleation and growth (e.g., Houze Jr et al., 1981; Hogan et al., 2002; Stark et al., 2013). Moreover, secondary ice production (SIP) processes might be active in generating cells,



55 which can further increase the ice crystal number concentration (ICNC). Indeed, generating cells were found to only account for 10 - 20 % of the total ice growth (e.g., Houze Jr et al., 1981; Plummer et al., 2015), while the majority of the ice growth occurred in the feeder region below.

Ice crystals can grow by various ice processes depending on the ambient conditions and the size distribution of cloud droplets and ice crystals (e.g., Marshall and Langleben, 1954; Fukuta and Takahashi, 1999; Bailey and Hallett, 2009; Connolly et al., 2012). For example, small ice crystals grow initially by diffusion of water vapor and thus their habit is determined by the ambient temperature and supersaturation (Magono and Lee, 1966; Bailey and Hallett, 2009). When ice crystals reach a critical size, they can grow more efficiently by aggregation and riming. Aggregation involves the collision and coalescence between ice particles and is most efficient at temperatures higher than -10°C due to the presence of a thicker quasi-liquid layer, which enhances the stickiness of the ice particles (e.g., Lohmann et al., 2016b). Riming, which involves the collision of an ice particle with a supercooled cloud droplet that freezes upon contact, has often been observed in the feeder regions of clouds (Mitchell et al., 1990; Borys et al., 2000; Borys et al., 2003; Saleeby et al., 2009; Saleeby et al., 2011; Lowenthal et al., 2011; Lowenthal et al., 2019) and has been found to enhance surface precipitation by up to 20 - 50 % (e.g., Mitchell et al., 1990; Borys et al., 2003; Lowenthal et al., 2016). For example, Lowenthal et al. (2016) observed that the precipitation at a mountain-top observatory gained the majority of its mass within 1 km above the mountain-top in the so-called feeder cloud. The efficiency of riming strongly depends on the cloud droplet size distribution (e.g., Borys et al., 2003; Saleeby et al., 2013). Additionally, riming can also produce a large number of ice splinters; e.g., when a cloud droplet of an appropriate size ($> 25\ \mu\text{m}$ in diameter) collides with a rimed ice particle ($> 0.5\ \text{mm}$ in diameter) (Mossop, 1978; Lamb and Verlinde, 2011). This SIP process is called the Hallett-Mossop process (Hallett and Mossop, 1974) and is thought to be active at temperatures between -3°C and -8°C . Other SIP mechanisms include the fragmentation of fragile ice crystals upon collisions with large ice particles (Vardiman, 1978) and the release of small secondary ice particles upon freezing of drizzle-sized droplets (e.g. Langham and Mason, 1958; Mason and Maybank, 1960; Lauber et al., 2018). Indeed, the ICNCs measured in feeder clouds at mountain top research stations frequently exceed the observed ice nucleating particle (INP) concentrations by several orders of magnitude (e.g., Rogers and Vali, 1987; Lloyd et al., 2015; Beck et al., 2018; Lowenthal et al., 2019). Several studies suggested that this discrepancy between the INP concentration and the ICNC can be explained by the influence of surface processes such as blowing snow (Geerts et al., 2015; Beck et al., 2018), hoar frost (Lloyd et al., 2015) or riming on snow-covered surfaces (Rogers and Vali, 1987), which can significantly increase the local ICNC and thereby influence the further evolution of the cloud. So far, it has been difficult to disentangle the contribution of surface processes and SIP mechanisms to the observed ICNC by means of mountain top observations. Therefore, innovative measurement strategies that are located far enough from mountain-tops or mountain slopes are required to reduce the influence of surface processes and to assess the importance of SIP mechanisms and their implications for precipitation formation in feeder clouds.

85 In this study, we investigate the microphysics of a cloud system passing over the Swiss Alps by combining a multi-dimensional set of instruments. A particular emphasis is placed on studying the role of cloud top generating cells and a surface-decoupled feeder cloud for ice growth and precipitation initiation. While most of the studies agree that generating cells have important implications for precipitation formation, less research has focused on the mechanisms that are responsible for the enhanced



ice formation and growth within these cells. We will approach this problem from an INP-cloud perspective, by combining INP
90 and ice crystal measurements. Furthermore, we discuss the role of a low-level feeder cloud for ice growth and SIP processes.
While the lowest part of the boundary layer is usually inaccessible for aircraft in complex terrain or is limited to observations at
mountain-tops, we analyze the microstructure of the low-level feeder cloud by using a tethered balloon system. The presented
case study was observed during the Role of Aerosols and CLOUDs Enhanced by Topography on Snow (RACLETS) campaign,
which took place in the Swiss Alps during February and March 2019. The analysis is based on an extensive set of observa-
95 tions including (1) ground-based remote sensing observations from a cloud radar, microwave radiometer and wind profiler, (2)
balloon-borne in situ observations, (3) INP measurements and (4) surface-based precipitation measurements.

2 Measurement location and instruments

The data presented in this paper was collected during the RACLETS campaign, which took place in the Swiss Alps in the
region around Davos from 8 February 2019 to 28 March 2019 and was designed to observe the pathways of precipitation
100 formation covering the entire aerosol-cloud-precipitation process chain.

Observations of the cloud microphysics were conducted at Wolfgang (1630 m; see Fig. 1) using remote sensing and in situ
instruments. Ground-based remote sensing measurements were obtained with a vertically-pointing Ka-band polarimetric cloud
radar (Mira-36 METEK GmbH, Germany; Görndorf et al., 2015), which provided vertical profiles of radar reflectivity factor,
Doppler velocity, spectral width and linear depolarization ratio (LDR). The radar observations have a vertical resolution of
105 31.17 m and a temporal resolution of 10 s. Moreover, a 14-channel microwave radiometer (HATPRO, Radiometer Physics
GmbH, Germany; Rose et al., 2005) was used to observe vertical profiles of atmospheric temperature and humidity as well as
the column integrated water vapor content (IWV) and liquid water path (LWP). Furthermore, the three-dimensional wind field
was measured at Wolfgang using a radar wind profiler (LAP-3000 Wind profiler, Vaisala, US). The wind profiler data has a
vertical resolution of 200 m and a temporal resolution of 5 min. In situ observations of the vertical low-level cloud structure
110 were measured with the tethered balloon system HoloBalloon (Ramelli et al., 2020). The main component of the measurement
platform is the HOLOGraphic cloud Imager for Microscopic Objects (HOLIMO), which can image an ensemble of cloud
particles in the size range from small cloud droplets (6 μm) to precipitation-sized particles (2 mm) in a three-dimensional
sample volume (Henneberger et al., 2013; Beck et al., 2017; Ramelli et al., 2020). HOLIMO provides information about the
phase-resolved particle size distribution and particle habit. As in Henneberger et al. (2013) and Beck et al. (2017), partitioning
115 between cloud droplets and ice crystals was applied to particles larger than 25 μm , i.e., particles smaller than 25 μm were
classified as cloud droplets. In the present study, cloud droplets were classified using support vector machines, whereas the
ice particles were classified visually (manual classification) in order to reduce the number of misclassified ice particles. Thus,
because of the applied size threshold (25 μm) and the visual classification, the reported ice properties (e.g., ICNC, IWC) can
be considered as a lower estimate. Additionally, the ICNC was derived from the remote sensing observations with the method
120 of Bühl et al. (2019). The particle diameter was estimated from the particle terminal fall velocity and spectral width measured
with the cloud radar. For this case, the particle shapes from Mitchell (1996) were used, assuming 'hexagonal plates' for ice

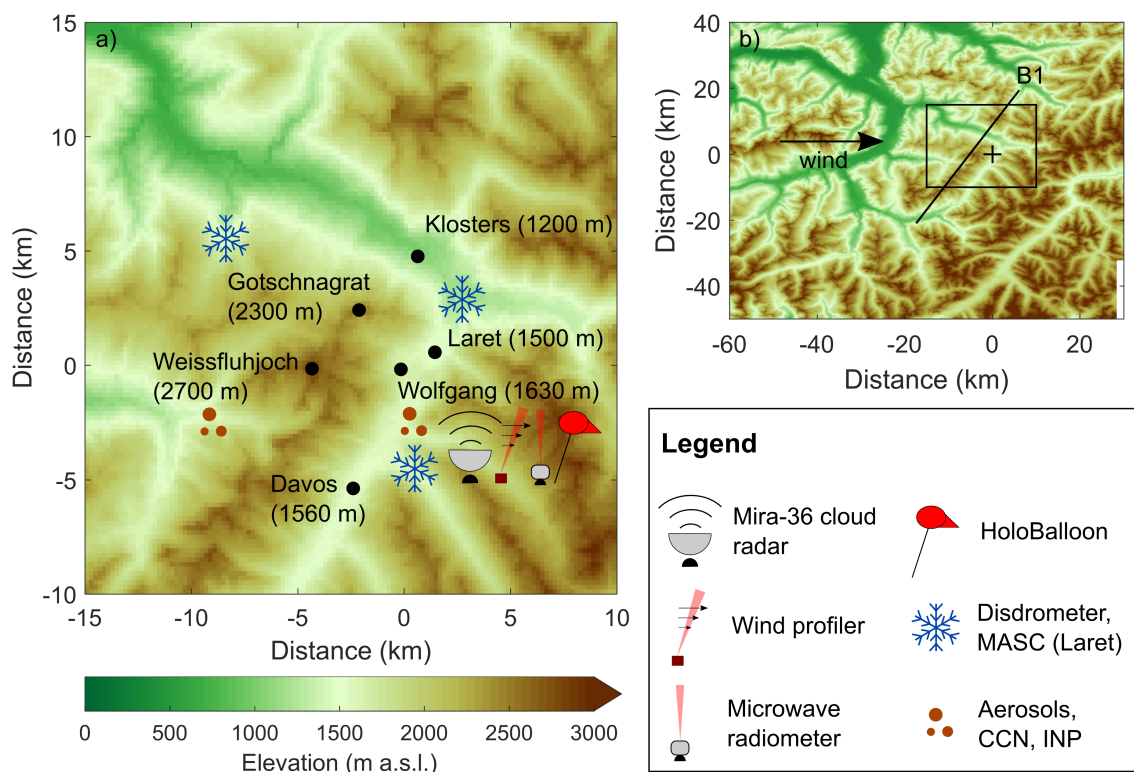


Figure 1. Overview of the measurement locations and the experimental setup (a). The geographical location of Wolfgang (black cross) and the surrounding topography is shown in (b). The large-scale wind direction was from the west as indicated by the black arrow. The most relevant mountain barrier is indicated by B1. An enlarged section of the measurement sites (black rectangle in b) and the instrument setup is shown in panel (a). The elevation data was obtained from the digital height model DHM25 of the Federal Office of Topography swisstopo: https://shop.swisstopo.admin.ch/de/products/height_models/dhm25200, last access: 9 March 2020.

crystals smaller than $600\ \mu\text{m}$ in diameter and ‘aggregates of planar polycrystals in cirrus clouds’ for ice particles larger than $600\ \mu\text{m}$ in diameter. The particle shape was derived from LDR measurements of the cloud radar and the images of ice crystals reconstructed from HOLIMO. The uncertainty of the retrieved ICNC is one order of magnitude.

125 Precipitation was measured at three locations (Wolfgang 1630 m, Laret 1500 m, Gotschnagrät 2300 m; see Fig. 1) using Particle Size Velocity (Parsivel) disdrometers (OTT Parsivel2, OTT HydroMet, Germany; Tokay et al., 2014). Parsivel disdrometers can measure both the size and the fall velocity of hydrometeors that fall through a laser sheet (Löffler-Mang and Joss, 2000). The size of the hydrometeor is estimated from the signal attenuation, whereas the fall velocity of the hydrometeor is obtained from the signal duration. Using the single particle size - fall velocity relationship, the observed particles can be classified
 130 into different hydrometeor classes, by applying different hydrometeor-dependent parameterizations (e.g., Yuter et al., 2006). Precipitation particles in the size range between 0.2 mm and 25 mm are measured. The temporal resolution of the measurements is 30 s. Additionally, a Multi-Angle Snowflake Camera (MASC; Garrett et al., 2012) was installed at Laret (see Fig. 1), which



took photographs of hydrometeors from three different angles and simultaneously measured their fall velocity. The MASC is sensitive to hydrometeors in the size range between 100 μm and 10 cm. Furthermore, a snow drift station was installed at
135 Gotschnagrat, which provided data about wind-driven redistribution of snow (Walter et al., 2020).

Lastly, aerosols and INP properties were measured at the valley station Wolfgang (1630 m) and at the mountain-top station Weissfluhjoch (2700 m) (see Fig. 1). Aerosol instruments were connected to heated inlets for measurements of ambient air at each site. Additionally, ambient aerosols were collected approximately every 1.5 h with a high flow rate impinger (Coriolis μ , Bertin Technologies, France, operation at 300 lpm for 20 mins; Carvalho et al., 2008). The impinger collected aerosol particles
140 larger than 0.5 μm in swirling liquid water and the aqueous solution was analyzed in drop-freezing instruments in order to obtain INP concentration spectra from 0 $^{\circ}\text{C}$ to approximately -20 $^{\circ}\text{C}$. The DRoplet Ice Nuclei Counter Zurich (DRINCZ; David et al., 2019) was operated at Wolfgang and the LED-based Ice Nucleation Detection Apparatus (LINDA; Stopelli et al., 2014) was run at Weissfluhjoch. Both drop-freezing instruments use a digital camera to detect freezing by a change in the light transmission through the aqueous solution. An intercomparison study was conducted between DRINCZ and LINDA. The
145 differences in observations were within the instrumental uncertainty.

The cumulative INP concentration was calculated following eq. (4) in Vali (2019):

$$INPC(T) = -\frac{\ln(1 - FF(T))}{V_a \cdot C}, \quad \text{where } C = \frac{F_{\text{impinger}} \cdot t_{\text{sample}}}{V_{\text{liquid}}} \cdot C_{\text{stdL}} \quad \text{and} \quad C_{\text{stdL}} = \frac{p_{\text{ambient}}}{p_{\text{ref}}} \cdot \frac{T_{\text{ref}}}{T_{\text{ambient}}} \quad (1)$$

using the temperature-dependent frozen fraction $FF(T)$, the volume of an individual aliquot V_a (50 μL at Wolfgang, 100 μL at Weissfluhjoch) and the normalization factor C , which converts the concentration to standard liters of ambient air. C was
150 calculated for each sample by considering the flow rate of the impinger F_{impinger} (300 lpm), the sampling time t_{sample} (usually 20 min), the end volume of the liquid V_{liquid} (approx. 15 mL) and the conversion factor from liters to standard liters C_{stdL} (including the ambient temperature T_{ambient} and pressure p_{ambient} at each site and the reference temperature $T_{\text{ref}} = 273.15\text{K}$ and pressure $p_{\text{ref}} = 1013.25\text{hPa}$). According to the specifications above, the minimal detectable concentration (limit of detection) at Wolfgang was $6.3 \cdot 10^{-4} \text{stdL}^{-1}$ and at Weissfluhjoch $3.5 \cdot 10^{-4} \text{stdL}^{-1}$.

155 3 Description of the case study

The synoptic weather situation over Europe on 8 March 2019 was characterized by a large-scale westerly flow with several low pressure systems (Fig. 2a). This strong westerly flow persisted for several days and brought moist air from the Atlantic towards central Europe. A low-pressure system located over Scandinavia produced a small-scale disturbance on its southern edge, which crossed Switzerland during the day and reached Davos in the afternoon. The presented case study was observed
160 during the passage of this small-scale disturbance, which arrived in Davos at around 15 UTC and lasted until 19 UTC.

During the passage of the cloud system, the temperature at Davos decreased from 3 $^{\circ}\text{C}$ to -2 $^{\circ}\text{C}$ ($\Delta T = -5$ $^{\circ}\text{C}$) and the temperature at Weissfluhjoch decreased from -5 $^{\circ}\text{C}$ to -7.5 $^{\circ}\text{C}$ ($\Delta T = -2.5$ $^{\circ}\text{C}$). The vertical temperature profile of a radiosonde ascent is shown in Figure 2b. The radiosonde was launched from Payerne, which is located around 200 km upstream of Davos. The temperatures measured at Davos, Gotschnagrat and Weissfluhjoch were slightly warmer (1-2 $^{\circ}\text{C}$) than the temperature measured

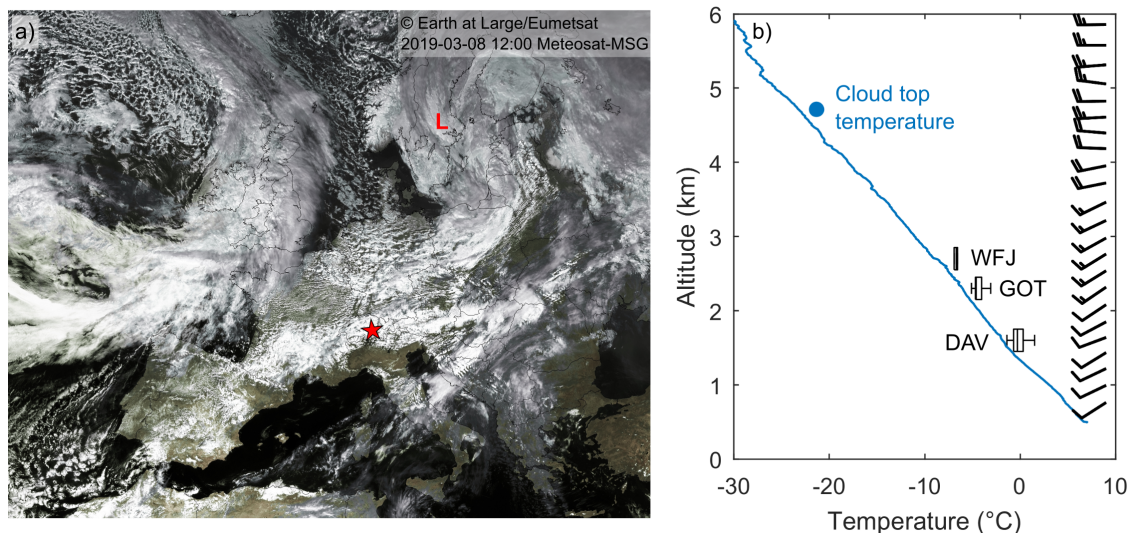


Figure 2. Overview of the synoptic weather situation on 8 March 2019, showing a satellite picture taken over Europe at 12 UTC (a, Eumetsat) and the vertical temperature profile measured by a radiosonde (12 UTC) launched from Payerne (b, MeteoSwiss). The boxplots in (b) indicate the temperature measured at the weather stations Davos (DAV, 1600 m), Gotschnagrat (GOT, 2300 m) and Weissfluhjoch (WFJ, 2700 m) during the passage of the cloud system. The blue dot indicates the cloud top temperature ($-21\text{ }^{\circ}\text{C}$) and cloud top height (4700 m), which was estimated from the cloud radar observations averaged between 16 UTC and 18 UTC. The wind barbs are shown on the right side.

165 by the radiosonde, but the observed lapse rate near Davos was in good agreement with the radiosonde profile measured at Payerne (see boxplots in Fig. 2b). A cloud top temperature of around $-21\text{ }^{\circ}\text{C}$ was estimated from the observed temperature profile, assuming the same temperature deviation as for the ground-based stations and a cloud top height of 4700 m (derived from the cloud radar observations averaged between 16 UTC and 18 UTC).

The horizontal wind fields were measured with a radar wind profiler at Wolfgang (Fig. 3a). In agreement with the Payerne
170 sounding, the wind profiler showed a large-scale wind direction from the west with a mean wind speed in the range of 10 m s^{-1} to 15 m s^{-1} above 3000 m. Below 2400 m, the wind speed was lower ($< 5\text{ m s}^{-1}$) and the flow was coming from the north-east (confined by the Davos valley). This pattern in the low-level wind field can be explained by shielding effects due to the mountain barrier B1 located upstream of Wolfgang (Fig. 1b), resulting in a decoupled low-level flow in the lee of the mountain barrier. A strong decrease in wind speed was observed above 2700 m between 17:45 UTC and 18:30 UTC. In addition, the wind direction
175 veered from 250° to 280° during this time period. This change in the wind pattern coincides with the period of the strongest precipitation event at Wolfgang (Fig. 4e) and could potentially have contributed to the glaciation of the MPC. Furthermore, enhanced wind shear was observed near cloud top ($> 10\text{ m s}^{-1}\text{ km}^{-1}$) with a maximum of $20\text{ m s}^{-1}\text{ km}^{-1}$ corresponding to the most intense precipitation peak (cf. Fig. 3b, Fig. 4e). Another layer of enhanced wind shear was observed between 2500 m and 3000 m, due to the interaction of the large-scale flow with the mountain barrier B1 (Fig. 1).

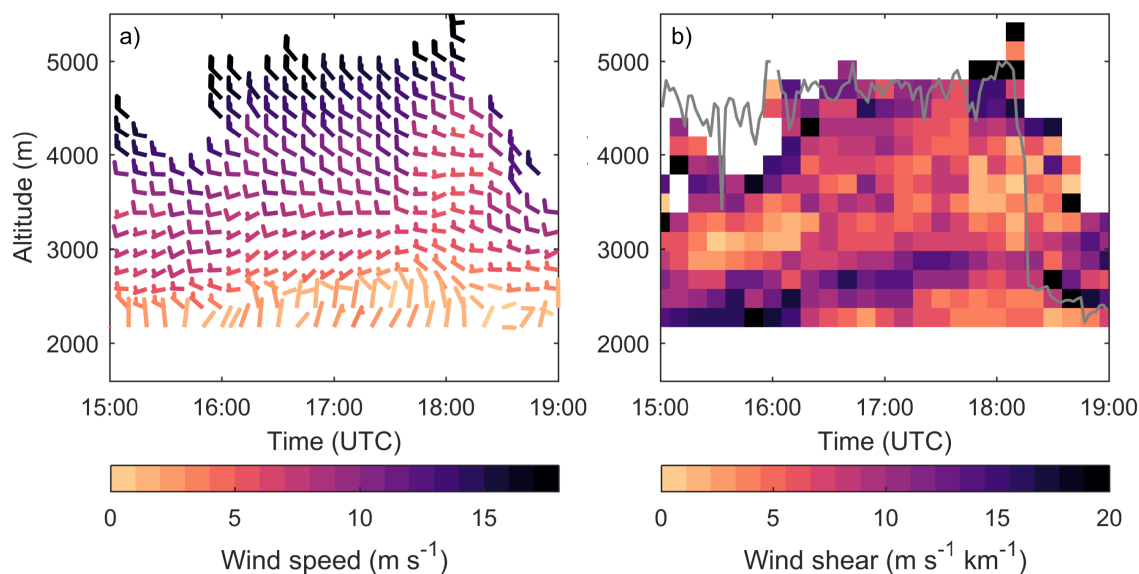


Figure 3. Observations of the wind speed and wind direction (a) and of the wind shear (b) measured by the radar wind profiler located at Wolfgang. The vertical wind shear (s) was calculated from the wind profiler observations, considering changes in the scalar wind speed and direction (u) between two adjacent height levels (z) ($s = \frac{u_2 - u_1}{z_2 - z_1}$). The gray line in (b) shows the cloud radar reflectivity contour of -30 dBZ, which indicates the cloud top height.

180 4 Results and Discussion

4.1 Overview of the microphysical cloud structure

An overview of the observed microphysical cloud structure is shown in Figure 4. The radar reflectivity shows that the precipitation began at 15:10 UTC and was convective in nature (Fig. 4a). At around 17:30 UTC, the reflectivity increased at all altitudes and the highest precipitation rates were observed at the surface (Fig. 4e). The period of high reflectivity (> 10 dBZ) lasted for about one hour. After this period, the cloud top lowered from 5000 m to 2800 m and the precipitation ended shortly after 18:40 UTC. The bulk of the precipitation originated at cloud top as can be seen from the fallstreak pattern of enhanced radar reflectivity (> 10 dBZ, Fig. 4a). The contour frequency by altitude diagram (CFAD, Fig. 5) of the radar reflectivity (Fig. 5a) indicates a rapid increase in the radar reflectivity near cloud top, suggesting that the ice crystals were formed in the layer between 5000 m and 4000 m. The ice crystals rapidly grew to large sizes between 4000 m and 3000 m, before they partly sublimated in the layer between 3000 m and 2000 m, as indicated by the decreasing radar reflectivity (Fig. 4a and 5a) and Doppler velocity (Fig. 4b and 5b) below 3000 m (assuming horizontal homogeneity). The majority of upward motion was observed above 3500 m (Fig. 4b and 5b). It is important to note that the measured vertical Doppler velocity is the sum of the particle fall speed and the air motion. Thus, as the ice particles grow to larger sizes while falling towards the ground, their fall speed increases and therefore mask the updrafts more easily. The Doppler velocity CFAD shows large variations between -4 m s^{-1}

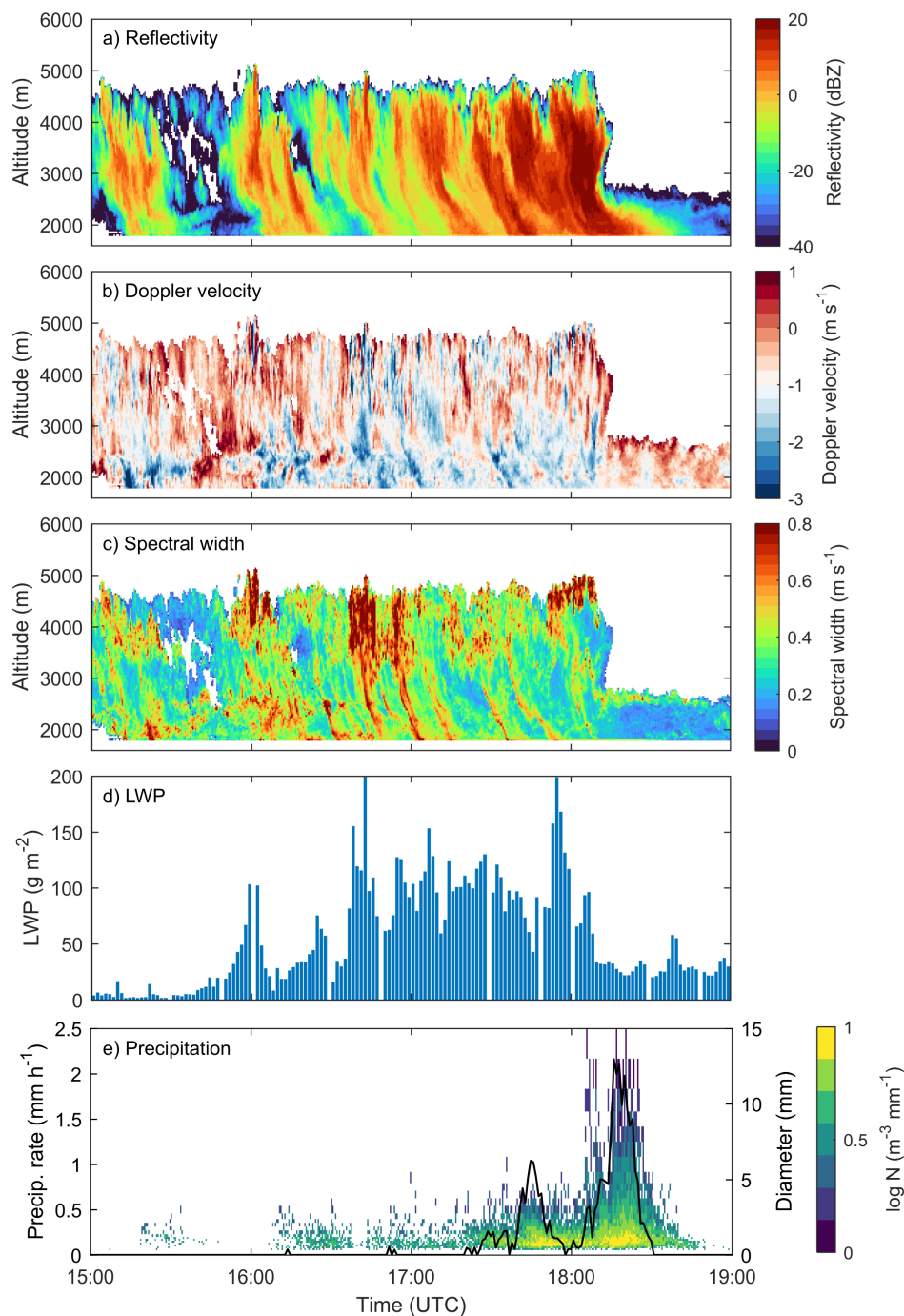


Figure 4. Observations of the cloud structure measured by the cloud radar (a-c) and the microwave radiometer (d) at Wolfgang on 8 March 2019. The cloud radar observations show the radar reflectivity (a), Doppler velocity (b) and spectral width (c). Note that the colorbar in (b) is centered at -1 m s^{-1} to approximately account for the hydrometeor fall speed. The column-integrated LWP measured by the microwave radiometer is shown in (d) and the precipitation measured by the disdrometer at Wolfgang (1630 m) is shown in panel (e).

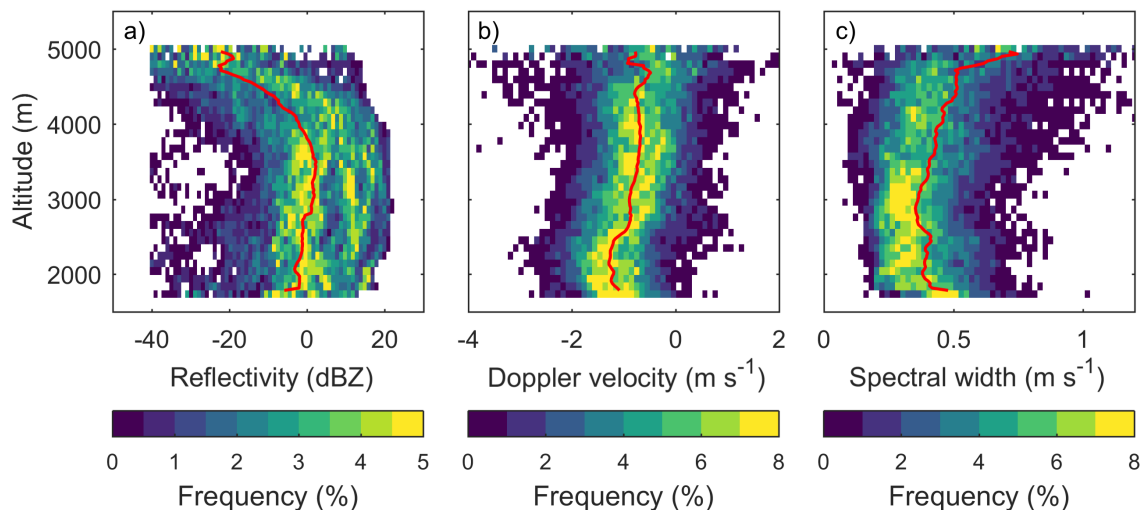


Figure 5. CFADs of the radar reflectivity (a), Doppler velocity (b) and spectral width (c) for the time period between 15:50 UTC and 18:20 UTC. The red line shows the mean vertical profile. The following bin sizes were applied: (1) radar reflectivity from -40 dBZ to 30 dBZ in 1 dBZ intervals, (2) Doppler velocity from -4 m s^{-1} to 3 m s^{-1} in 0.1 m s^{-1} intervals and (3) spectral width from 0 m s^{-1} to 1.2 m s^{-1} in 0.02 m s^{-1} intervals. A height interval of 100 m was used for all radar properties.

195 to 2 m s^{-1} near cloud top (Fig. 5b), indicative of turbulent motions. Indeed, the strong variability in the Doppler velocity was collocated with the enhanced shear layer from the wind profiler (Fig. 3b). Furthermore, the spectral width was also enhanced locally near cloud top (Fig. 4c), which can be attributed to the presence of turbulence (see Fig. 4b) near cloud top. The occurrence of (1) high radar reflectivity fallstreaks (Fig. 4a), (2) positive Doppler velocities (Fig. 4b) and (3) increased spectral width (Fig. 4c) near cloud top, suggest the presence of cloud top generating cells. Cloud top generating cells can enhance ice
200 nucleation and growth and as such have important implications for precipitation formation (e.g., Houze Jr et al., 1981; Hogan et al., 2002; Evans et al., 2005; Ikeda et al., 2007; Crosier et al., 2014; Kumjian et al., 2014; Plummer et al., 2014; Rosenow et al., 2014; Plummer et al., 2015; Rauber et al., 2015) as will be further discussed in Section 4.2.

Ice particles that formed within the seeder region interact with other cloud particles while falling through the cloud and thus influence the microphysics of the feeder region below. The low-level cloud structure was observed with the tethered balloon
205 system HoloBalloon (see Fig. 6). The balloon-borne measurements indicate the presence of a low-level liquid layer that was confined to the lowest 300 m of the cloud (see Fig. 6). The cloud droplet number concentration (CDNC) increased from 100 cm^{-3} to 350 cm^{-3} between 16 UTC and 17:45 UTC (Fig. 6a), before the CDNC decreased after 18 UTC. The mean cloud droplet diameter ranged between $8 \mu\text{m}$ and $12 \mu\text{m}$ as shown by the size distribution in Figure 7a. The ICNC was in the range of 1 L^{-1} to 4 L^{-1} between 16 UTC and 18 UTC (Fig. 6b). Ice crystals were especially observed when fallstreaks of enhanced
210 radar reflectivity reached the surface. During the main precipitation event, after 18 UTC, the ICNC increased up to 14 L^{-1} . During the same time period, the ratio between the ice water content (IWC) and total water content (TWC), which is often

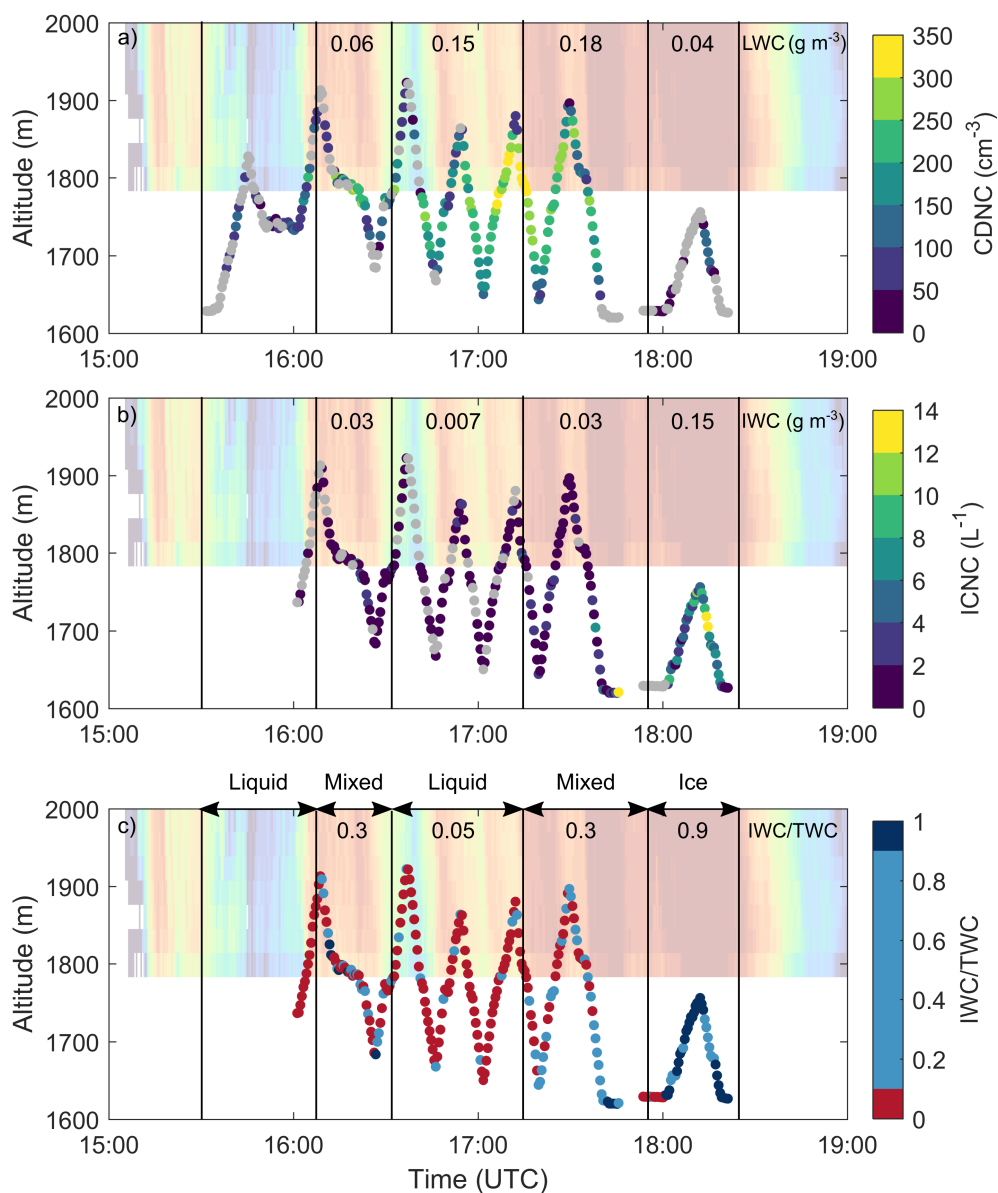


Figure 6. Vertical in situ profiles of the CDNC (a), ICNC (b) and the IWC/TWC ratio (c). The gray dots in (a) and (b) indicate measurement points, which are associated with a liquid water content (LWC) of $< 0.01 \text{ g m}^{-3}$ (for CDNC) or an IWC of 0 L^{-1} (for ICNC). In (c), red colors represents liquid cloud regions ($\text{IWC/TWC} < 0.1$), light blue mixed-phase cloud regions ($0.1 \leq \text{IWC/TWC} < 0.9$) and dark blue indicates ice cloud regions ($\text{IWC/TWC} \geq 0.9$). The cloud radar reflectivity is shown in the background. The numbers in (a), (b) and (c) indicate the mean LWC, IWC and IWC/TWC ratio within the intervals defined by the black vertical lines.

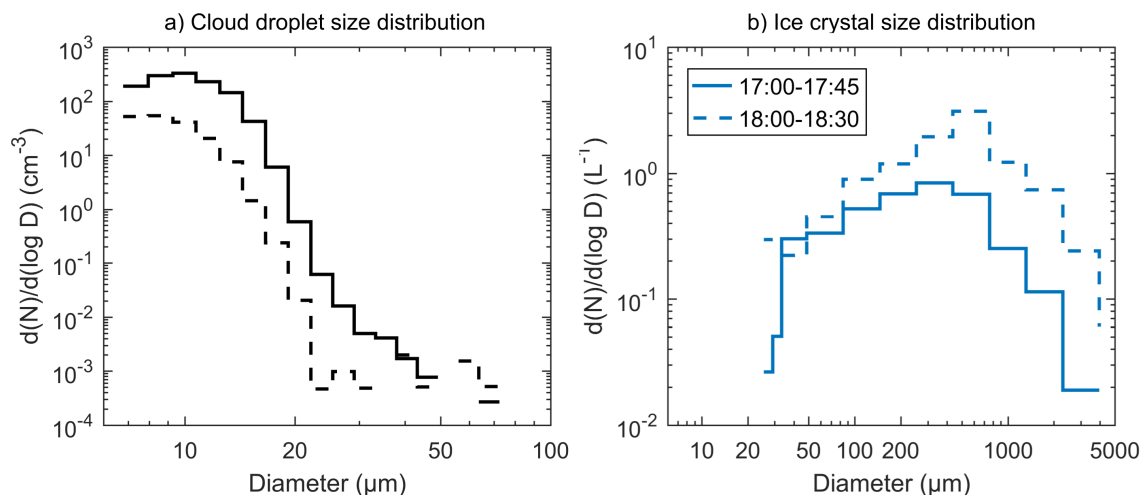


Figure 7. Cloud droplet (a) and ice crystal (b) size distributions observed with the HoloBalloon platform. The size distributions were averaged between 17 UTC and 17:45 UTC (solid line) and between 18 UTC and 18:30 UTC (dashed line).

used to characterize the cloud phase (e.g., Korolev et al., 2003; Lohmann et al., 2016a), increased from 0.05 - 0.3 (liquid to mixed-phase) to 0.9 (ice-phase). Thus, a transition from a mixed-phase low-level cloud (before 18 UTC) to an ice-dominated low-level cloud (after 18 UTC) was observed during the passage of the cloud system (Fig. 6c). The cloud radar and microwave radiometer observations suggest that the entire cloud layer glaciated, as an increase in the radar reflectivity (Fig. 4a) and a decrease in the LWP (Fig. 4d) was observed after 18 UTC. In the absence of sufficiently large updraft velocities for additional cloud droplet activation, the presence of large ice particles or high ICNC can lead to rapid glaciation of the cloud by the WBF process (Korolev and Isaac, 2003).

Even though downward motions were present on the lee side of the mountain barrier (see increased fraction of negative Doppler velocities in Fig. 5b), which contributed to hydrometeor evaporation/sublimation (see decreased reflectivity in Fig. 5a), a persistent low-level liquid layer was observed at Wolfgang. We suggest that this shallow low-level feeder cloud formed due to orographic lifting, as the low-level flow in the lee of the mountain barrier was decoupled from the large-scale flow (Fig. 3a) and was forced to rise from Klosters (1200 m) to Wolfgang (1630 m) over the local topography. It is assumed that this shallow cloud could not generate significant precipitation by itself, due to the limited time available for collision-coalescence of cloud droplets to produce precipitation-sized particles and due to the warm temperatures (> -3 °C), which were limiting the amount of INPs and thus ice formed through primary ice nucleation. However, the hydrometeors that formed in the generating cells can "feed" on the low-level liquid layer and thus enhance precipitation by riming and depositional growth. Additionally, it can provide an environment favorable for the production of secondary ice particles as will be discussed in Section 4.3. In the following section, we investigate the microphysics and dynamics in cloud top generating cells and explore the origin of ice crystals.



4.2 The origin and growth of ice crystals in cloud top generating cells

Observations from the cloud radar, microwave radiometer, HoloBalloon platform and ground-based aerosol measurements were combined to study the dynamics and microphysics within cloud top generating cells. Since no in situ observations within generating cells or near cloud top were available during the RACLETS campaign, the analysis of the microphysics was limited to observations from remote sensing instrumentation and balloon-borne in situ measurements near cloud base. In the first part of this section, the overall dynamical and microphysical structure of generating cells is characterized, whereas in the second part the origin of ice crystals and the microphysical growth processes active within generating cells are investigated from an INP-cloud perspective.

When the strongest generating cells were present, vertical overshooting of up to 500 m was observed at the cloud top (Fig. 8; e.g., at 16 UTC and 16:45 UTC), indicating the presence of strong updrafts. This was also supported by observations of the maximum Doppler velocity (Fig. 8b), which was derived from the Doppler spectra (see Appendix A) and used as a proxy to identify updraft regions. The maximum Doppler velocity suggests that the strongest updrafts were present in the core regions of the cloud top generating cells ($> 3 \text{ m s}^{-1}$), whereas updrafts were weaker outside of the generating cells and at altitudes below 3000 m (Fig. 8b). It is likely that liquid water was produced in these updraft cells, as a positive correlation was found between the vertically-integrated maximum Doppler velocity and the LWP measured by the microwave radiometer (see Fig. A1b). Moreover, anomalies in the cloud top properties and the LWP were observed during the periods with generating cells (Fig. 8d). Coinciding peaks in the anomaly signal were labeled as GC1 (16 UTC), GC2 (16:45 UTC) and GC3 (17:55 UTC). The Spearman's rank correlation coefficients of the anomalies ranged between 0.46 (for reflectivity and spectral width) and 0.73 (for reflectivity and LWP) significant at the 5 % level. Thus, given the significant correlation between updrafts, LWP and radar reflectivity within generating cells, it is likely that the updrafts acted as a major driver for the formation and maintenance of generating cells by providing a continuous source of liquid water and thereby enhancing ice nucleation and growth through immersion freezing, subsequent vapor deposition and riming. Indeed, the comparison of the Doppler spectra inside (GC_{in}) and outside (GC_{out}) of GC2 suggests that the generating cell contributed 10 dBZ to the radar reflectivity (see Appendix B).

To further explore the microphysics within cloud top generating cells, the Doppler spectra along the 17 UTC fallstreak were investigated (Fig. 9). This approach allows to obtain a continuous picture of the evolution of the particle populations along the fallstreak and to draw conclusions regarding the microphysical processes active. Previous studies used the Doppler spectra information for the classification and characterization of ice particle shape and particle populations (e.g., Myagkov et al., 2016; Bühl et al., 2016). The vertical profile of the Doppler spectra shows a broad particle distribution spanning from -5 m s^{-1} to 4 m s^{-1} between 3300 m and 5000 m height, indicative of a turbulent layer. This layer likely marked the extent of the generating cell, where ice crystals were produced and initial growth occurred. The Doppler spectra show a spectral bimodality below 3300 m (Fig. 9; i.e. presence of multiple particle populations with different fall speeds), which extends down to the surface. When analyzing the Doppler spectra of the full period with the peakTree technique (Radenz et al., 2019), multi-peaked situations become evident at the leading edges of the fallstreaks (Fig. 8c). The peakTree algorithm transforms each Doppler spectrum into a full binary tree structure, where the individual sub-peaks are represented as nodes (see Radenz et al.,

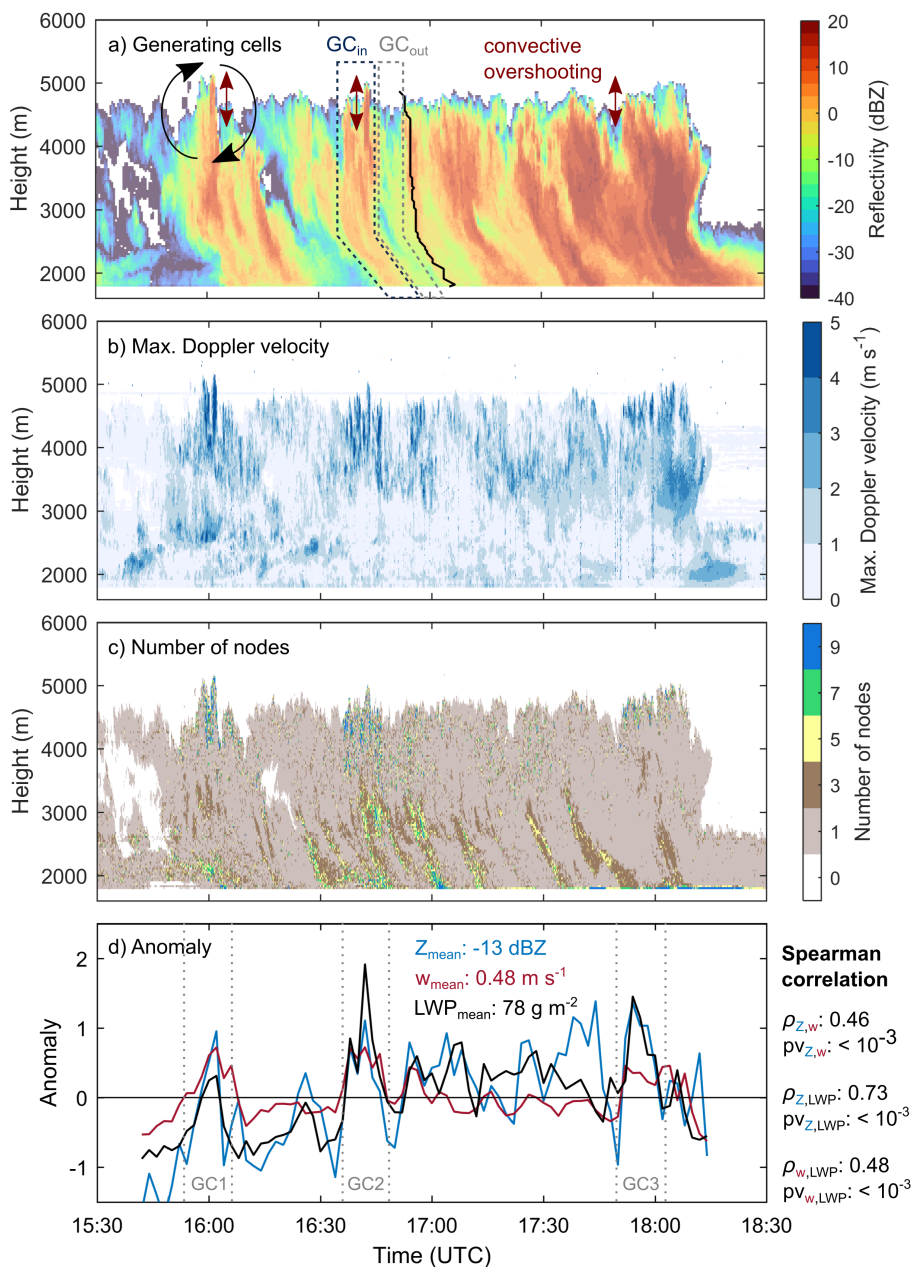


Figure 8. Time series of the radar reflectivity (a), maximum Doppler velocity (b) and number of nodes (c). The black line in panel (a) indicates the 17 UTC fallstreak, which was investigated in Fig. 9. The dashed lines indicate the regions inside (GC_{in}) and outside (GC_{out}) of GC2, which were used for the analysis in Appendix B. The maximum Doppler velocity was derived from the Doppler spectra (see Appendix A). The number of nodes were obtained from the peakTree analysis following the procedure described in Radenz et al. (2019). The evolution of the cloud top anomalies is shown in (d). The radar reflectivity (blue line) and spectral width (red line) were averaged over 600 m from the cloud top. The anomalies were normalized to the mean value, which is indicated in panel (d). The results of the Spearman’s rank correlation are shown to the right of panel (d) with ρ indicating the correlation coefficient and pV the p-value of the Spearman’s rank correlation.

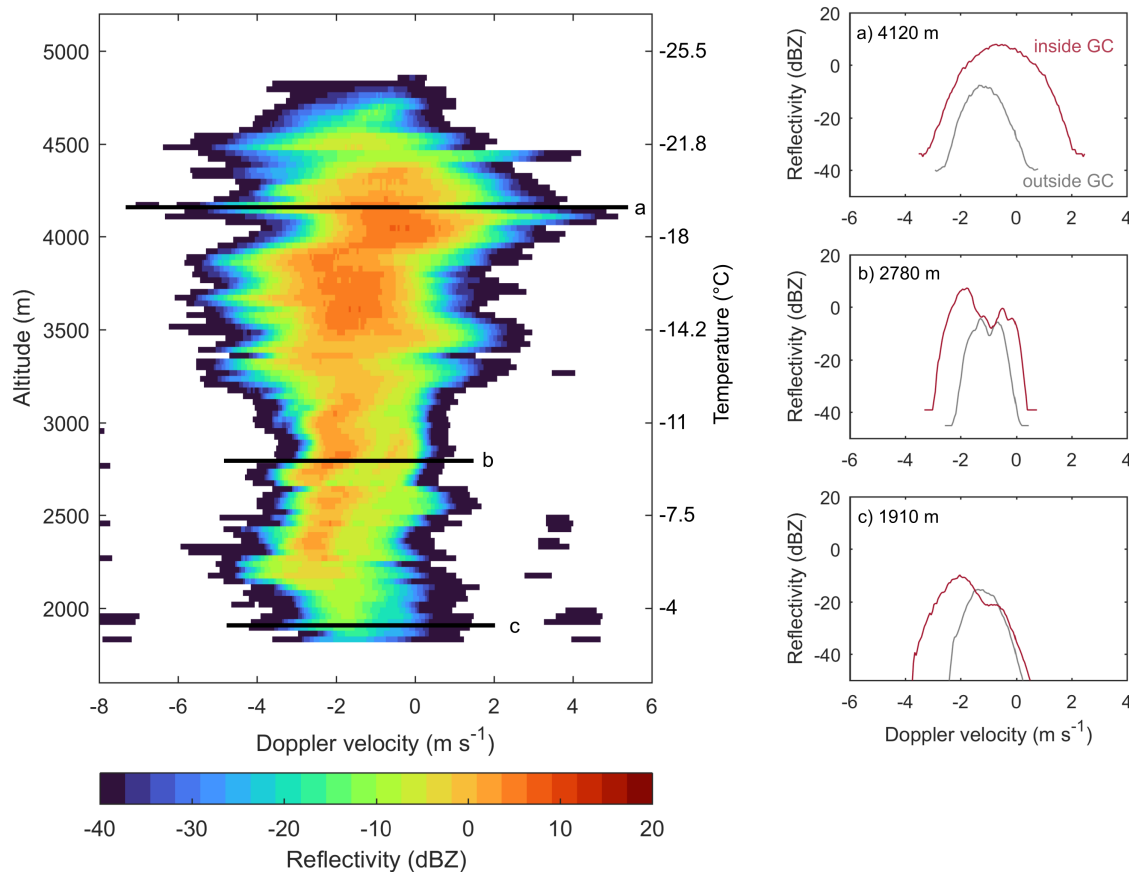


Figure 9. Vertical profile of the Doppler spectra along the 17 UTC fallstreak averaged over 1 min (indicated by black line in Fig. 8a). The Doppler spectra at three selected heights are shown on the right: 4120 m (within turbulent layer), 2780 m (at mountain barrier height), 1910 m (at balloon flight height). The red line indicates the Doppler spectrum inside the 17 UTC fallstreak, whereas the gray spectrum was measured before the fallstreak (inside GC_{out} in Fig. 8a).

265 2019). Thus, three nodes corresponds to two sub-peaks. For example, the Doppler spectrum in Figure 9b (red line) indicates
 the presence of two particle populations, a fast falling one (-2 m s^{-1}) and a slow falling one (-0.5 m s^{-1}). The LDR of the
 slower falling particle population was slightly higher (-25 dB ; not shown) compared to the faster falling population (-28 dB ;
 not shown). These LDR values are characteristic for oblate or plate-like particles (Myagkov et al., 2016). The observed Doppler
 spectra and the ice habits observed near cloud base (Fig. 13) suggest that the faster falling population represents heavily rimed
 270 ice particles/graupel, whereas the slower falling population was associated with stellar dendrites. This is also consistent with
 the observed temperature (dendrite regime; Magono and Lee, 1966; Bailey and Hallett, 2009) and the presence of supercooled
 liquid (riming) within the generating cells. It is likely that these two particle populations were already present above, but only
 separated below the turbulent layer due to the weaker updrafts and their difference in fall speed.

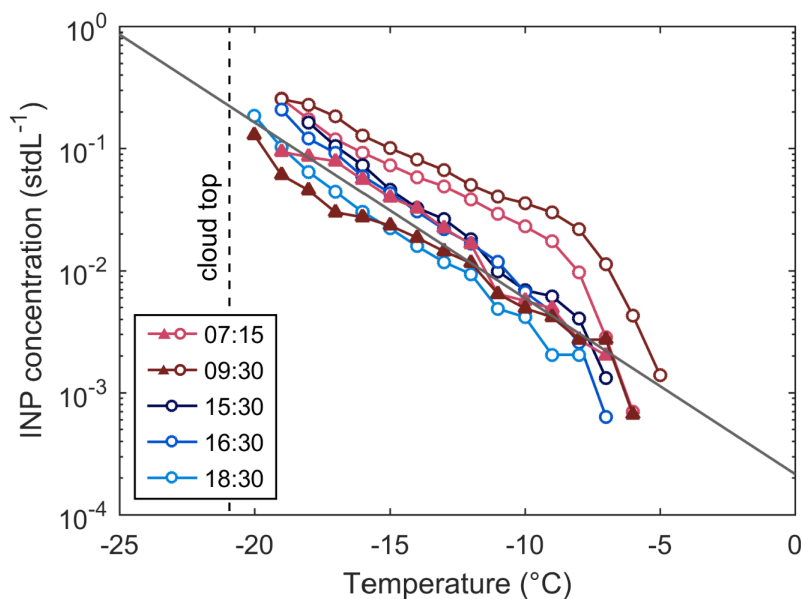


Figure 10. INP concentrations measured at Wolfgang (1630 m, circle) and Weissfluhjoch (2700 m, triangle) for different temperatures and times as indicated in the legend. The cloud top temperature of $-21\text{ }^{\circ}\text{C}$ is shown by the vertical dashed line. The dark gray line is a fit to the INP concentrations measured at Weissfluhjoch at temperatures between $-9\text{ }^{\circ}\text{C}$ and $-17\text{ }^{\circ}\text{C}$ (pre-cloud INP conditions).

In the following, we will further investigate the origin of ice particles that formed within generating cells. Numerous studies
275 have observed enhanced ice formation and growth in these updraft regions (Houze Jr et al., 1981; Hogan et al., 2002; Plummer et al., 2014; Ikeda et al., 2007; Crosier et al., 2014; Kumjian et al., 2014; Rauber et al., 2015). For example, Plummer et al. (2014) found that the ICNC was enhanced by a factor of 2 to 3 within the core region of generating cells compared to the region between the cells. While most of the studies agree that radiative cooling is a major driver for the formation and maintenance of cloud top generating cells, less research has focused on the reason for the enhanced ICNCs that were observed
280 within these cells. Here we provide potential reasons from an INP-cloud perspective and propose possible mechanisms by considering the measured INP concentrations and cloud-base observations of the ICNC and ice particle size. INP concentrations were measured at the valley site Wolfgang (1630 m) and at the mountain-top station Weissfluhjoch (2700 m) (Fig. 10). The observed INP concentrations at a given temperature spanned over one order of magnitude. The INP concentration measured at Wolfgang was a factor of 3 - 10 higher compared to Weissfluhjoch, which was likely a consequence of the decoupled low-level
285 flow (see Fig. 3a) and thus the sampling of different air masses. Based on the INP measurements at Weissfluhjoch, an INP concentration of 0.3 L^{-1} was extrapolated at cloud top (Fig. 10). It is important to note that the cloud top INP concentration was estimated from the Weissfluhjoch measurements in the morning (i.e., representative for pre-cloud INP concentrations), as no INP concentrations were measured at Weissfluhjoch during the passage of the cloud system. Additionally, cloud measurements were conducted by the HoloBalloon platform near cloud base. Since no in situ observations were available within
290 the generating cells, assumptions regarding the upper-level cloud properties were required. We assumed that the largest ice

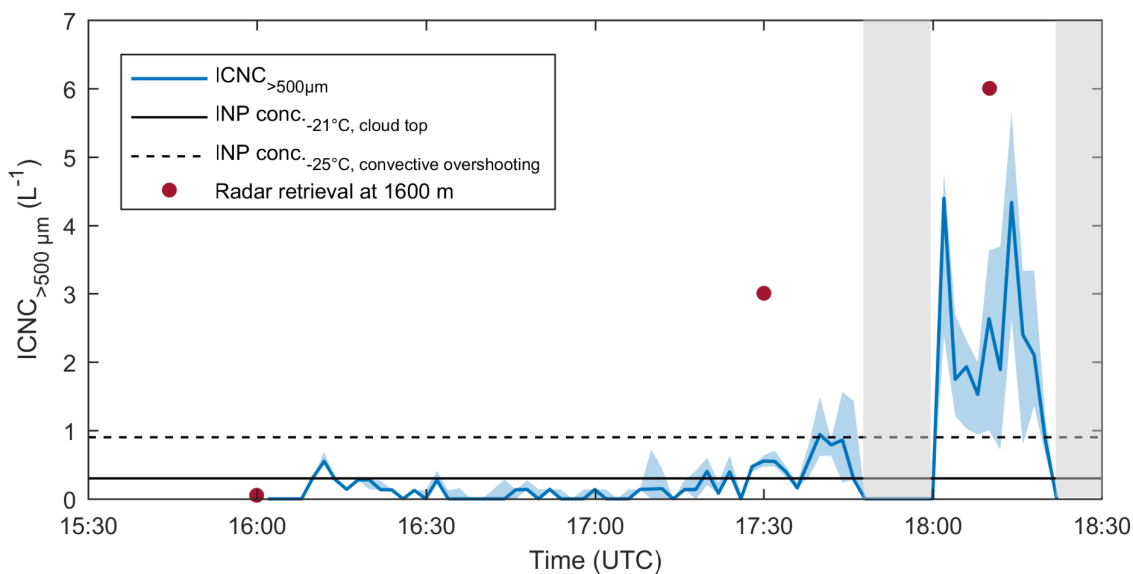


Figure 11. Timeseries of the ICNC of large ice particles ($> 500 \mu\text{m}$, $\text{ICNC}_{>500\mu\text{m}}$, blue line) measured near cloud base by the HoloBalloon platform. Ice particles larger than $500 \mu\text{m}$ in diameter were assumed to have formed near cloud top. The blue shaded area indicates ice particles larger than $400 \mu\text{m}$ (upper bound) and larger than $600 \mu\text{m}$ (lower bound). The INP concentrations extrapolated to -21°C (cloud top, solid line) and to -25°C (convective overshooting, dashed line) are indicated by the black horizontal lines. The ICNC estimated from the radar observations are shown by the red dots using the procedure described in Bühl et al. (2019) (see also Section 2). No microphysical measurements were available within the gray shaded areas.

particles ($> 400 - 600 \mu\text{m}$; derived from particle size distribution in Fig. 7b) formed near cloud top and grew to these large sizes while falling to the surface. This criterion is based on the assumption that the large ice particles did not sublimate completely prior to reaching the surface. The ICNC observed near cloud base was in the same order of magnitude as the ICNC retrieved from the remote sensing observations (red dots in Fig. 11) using the method described in Bühl et al. (2019). The comparison
295 between the observed ICNC of large ice particles ($\text{ICNC}_{>500\mu\text{m}}$) and the INP concentration at cloud top shows a discrepancy between the observed INP concentration and ICNC during certain time periods (Fig. 11), suggesting that the ICNC cannot be solely explained by primary ice nucleation, but that other mechanisms were active.

Static instability driven by cloud top radiative cooling can produce strong updrafts (Fig. 8b) and lead to convective overshooting of cloud top generating cells (see red arrows in Fig. 8a). This convective overshooting can decrease the cloud top temperature and therefore increase the ICNC formed by primary ice nucleation. For example, the cloud top height during GC1 increased by
300 500 m from 4500 m to 5000 m. Considering the observed temperature profile in Figure 2b, the cloud top temperature decreased by 3.6°C from -21°C (at the average cloud top height) to -24.6°C (at 5000 m) upon convective overshooting. Consequently, the INP concentration increased by a factor of 3.3 from 0.3 L^{-1} to 1 L^{-1} (Fig. 10) due to the colder cloud top temperature. The $\text{ICNC}_{>500\mu\text{m}}$ measured near cloud base lied below the extrapolated INP concentration at -25°C before 18 UTC (Fig. 11),



305 suggesting that the observed $\text{ICNC}_{>500\mu\text{m}}$ near cloud base can be solely explained by primary ice nucleation and convective overshooting. After 18 UTC, the $\text{ICNC}_{>500\mu\text{m}}$ measured near cloud base lied above the convective overshooting line (Fig. 11), suggesting that other processes were occurring.

For example, the positive feedback between supercooled liquid water, radiative cooling and turbulence that has been observed near cloud tops (e.g., Morrison et al., 2012) might have contributed to enhanced ice formation. The presence of supercooled
310 liquid can lead to strong longwave radiative cooling (e.g., Possner et al., 2017). This radiative cooling decreases the stability near cloud top, which causes turbulent motions, which in turn can produce further supercooled liquid water. The magnitude of the longwave radiative cooling strongly depends on the cloud phase, the liquid water content and particle size distribution, among other factors (e.g., Turner et al., 2018). Indeed, the LWP, as measured by the microwave radiometer, was enhanced within generating cells (see Fig. 8d) and thus likely increased the longwave radiative cooling at cloud top. The question is by
315 how much the radiative cooling was enhanced within generating cells due the increased cloud liquid water compared to their surrounding regions. Previous studies observed longwave radiative cooling rates in the range of $1 - 5 \text{ K h}^{-1}$ near cloud top (e.g., Chen and Cotton, 1987; Pinto, 1998; Jiang et al., 2000; Rasmussen et al., 2002; Morrison et al., 2011; Morrison et al., 2012; Possner et al., 2017; Turner et al., 2018; Eirund et al., 2019). Additionally, Turner et al. (2018) computed radiative heating rate (RHR) profiles in the atmosphere as a function of cloud type and LWP by using an observational data set. According to Turner
320 et al. (2018), an increase in the LWP from 50 g m^{-2} to 150 g m^{-2} (e.g., GC2 in Fig. 8) can cause an increase in the longwave radiative cooling rate from around 1.7 K h^{-1} to 2.9 K h^{-1} ($\Delta\text{RHR} = 1.2 \text{ K h}^{-1}$). This could potentially cool the cloud top temperature by 0.3 K , if a lifetime of around 15 min is assumed for cloud top generating cells (i.e., $1.2 \text{ K h}^{-1} \times 15 \text{ min} = 0.3 \text{ K}$), and increase the INP concentration from 0.3 L^{-1} to 0.35 L^{-1} (factor 1.2, see Fig. 10). Thus, longwave radiative cooling only plays a minor role in enhancing primary ice nucleation. Nevertheless, longwave radiative cooling is of major importance for
325 the production of radiatively driven turbulence near cloud top and thus for maintaining generating cells.

Other mechanisms must be active to explain the increased ICNCs after 18 UTC. For instance, the enhanced updrafts in generating cells allow all hydrometeors to grow to larger sizes. It is unlikely that the larger cloud droplet size would significantly increase primary ice nucleation by immersion freezing, which is the dominant ice nucleation mechanism in MPCs (e.g., Ansmann et al., 2008; De Boer et al., 2011; Westbrook and Illingworth, 2011), but it can play an important role for SIP. For
330 example, the freezing of drizzle-sized droplets can release a large number of small secondary ice particles (e.g., Langham and Mason, 1958; Mason and Maybank, 1960; Lauber et al., 2018; Korolev and Leisner, 2020). This process is known as droplet shattering and has been observed to be strongly dependent on the cloud droplet size and to be potentially effective over a large temperature range (Lauber et al., 2018). Previous field studies have observed the presence of drizzle-sized droplets in the size range of $100 \mu\text{m} - 300 \mu\text{m}$ in regions of strong vertical updrafts (e.g., Hauf and Schröder, 2006; Ikeda et al., 2007). A single
335 shattering event has the potential to produce hundreds of ice crystals. Thus, droplet shattering could increase the ICNC in generating cells by several orders of magnitude if supercooled drizzle drops are present in the updraft regions. As the discrepancy between the INP concentration at cloud top and $\text{ICNC}_{>500\mu\text{m}}$ was up to a factor 8 after 18 UTC, droplet shattering might have contributed to enhanced ice formation and growth and to the glaciation of the MPC. However, in situ observations within generating cells would be necessary to further investigate this hypothesis.

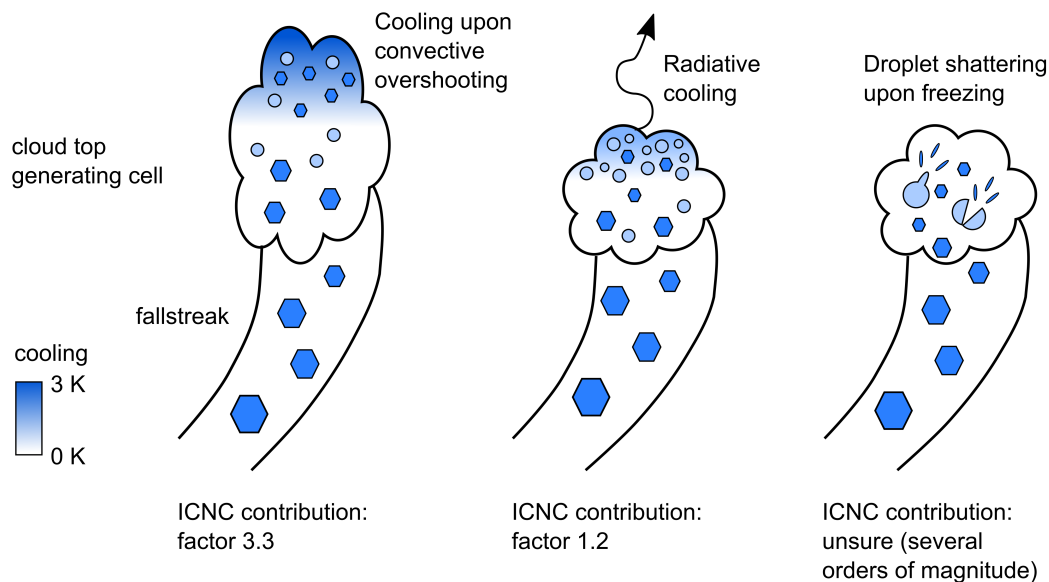


Figure 12. Potential mechanisms that could enhance the ICNC in cloud top generating cells: convective overshooting (left), radiative cooling (center) and droplet shattering upon freezing (right). Their potential contributions to the ICNC are estimated at the bottom based on the present case study (see text for more details).

340 In summary, the increased ICNC in generating cells can be the result of different mechanisms or a combination of several mechanisms. Three possible mechanisms have been proposed in this study and their potential contributions are summarized in Figure 12. Firstly, primary ice nucleation in generating cells can be increased due to convective overshooting or radiative cooling. The ICNC observed before 18 UTC can likely be explained by these two mechanisms, since the INP concentration and the $ICNC_{>500\mu m}$ measured near cloud base agreed within the same order of magnitude (Fig. 11). We found that the contribution
345 from convective overshooting (factor 3.3) was larger than that of radiative cooling (factor 1.2) in the present study. On the other hand, the ICNC of large particles measured at the surface after 18 UTC exceeded the INP concentration by almost one order of magnitude, suggesting that SIP processes such as droplet shattering might be active within generating cells. However, more targeted studies are necessary to understand which mechanisms are responsible for enhanced ice formation and growth within
350 generating cells. In particular, in situ measurements of the cloud properties within generating cells and their environmental conditions (e.g., temperature, updrafts, INP conditions) are of major importance to address these questions.

4.3 Secondary ice production processes in feeder cloud

Ice crystals that formed in the seeder region can grow by microphysical interactions with other cloud particles while falling through the cloud layer and thus influence the microphysics of the entire cloud. For example, if large ice particles fall through a supercooled liquid layer, they can initiate the glaciation of the cloud layer through the WBF process and/or grow by riming.
355 Furthermore, the total number of nodes (Fig. 8c) shows multi-peaked situations below 3300 m, indicating the presence of



multiple particle populations with different fall speed. This suggest that secondary ice particles might be produced in the feeder region of the cloud. In the following, we investigate the importance of ice growth and SIP in the feeder region by analyzing the phase-resolved cloud properties measured in situ with the HoloBalloon platform. In particular, the analysis of the ice crystal habit and size can provide important information about the formation and growth history of ice particles.

360 Figure 13 shows a representative set of ice particle images observed by HOLIMO as a function of height and time. It can be seen that ice crystal habits varied greatly during the passage of the cloud system. For example, the images indicate the presence of numerous columns between 17:00 UTC and 17:20 UTC at altitudes above 1780 m (yellow boxes), which are known to grow at temperatures between -3°C and -10°C (Magono and Lee, 1966; Bailey and Hallett, 2009). Furthermore, irregular shaped particles were abundant (green boxes), consistent with previous studies (e.g., Korolev et al., 1999; Stoelinga et al., 2007). A

365 large fraction of graupel and rimed particles was observed between 17 UTC and 17:40 UTC (red boxes). After 18 UTC, the ice crystals became more aggregated (blue boxes) and less rimed (see also MASC data in Fig. 14d), suggesting a decrease in the amount of liquid water available for riming. Furthermore, small pristine ice crystals (plates and columns) were present over the entire period (see Fig. 15c and purple boxes in Fig. 13).

The large variability in ice crystal habit and size suggests that the ice crystals have formed and grown in different regions. As

370 discussed in Section 4.2, it is likely that the heavily rimed ice particles and large dendrites (Fig. 13) were produced within the seeder region of the cloud and gained mass by riming and deposition while falling through the cloud. On the other hand, the small pristine ice crystals were likely formed within the feeder region of the cloud. Previous studies have found that small pristine ice crystals ($<100\mu\text{m}$) were spatially correlated with their environment of origin (e.g., Korolev et al., 2020). For example, it is possible that the observed columns originated within the multi-peaked structures (Fig. 8c), as the temperature

375 below 3000 m was in the temperature regime of columnar growth (Bailey and Hallett, 2009). Pristine plates likely grew in the lowest part of the cloud, where the prevailing temperature was above -3°C . These small ice crystals ($<100\mu\text{m}$) could have formed either by primary ice nucleation or by SIP processes within the feeder cloud and rapidly grown by diffusion to larger sizes (e.g., Korolev et al., 2020). The contribution of primary ice nucleation to the observed ICNC can be estimated from the measured INP concentration (Fig. 10), which was below the minimal detectable concentration ($6.3 \cdot 10^{-4} \text{stdL}^{-1}$; see Section

380 2) at a temperature of -3°C . Thus, the minimal detectable concentration of $6.3 \cdot 10^{-4} \text{stdL}^{-1}$ represents an upper limit for the INP concentration within the feeder region. The ICNC of particles smaller than $100\mu\text{m}$ in diameter observed in the feeder cloud ($1 - 2 \text{L}^{-1}$; Fig. 14a) exceeded the INP concentration by three orders of magnitude, suggesting that primary ice nucleation alone cannot explain the small ice crystals observed.

Secondary ice production processes are necessary to explain the observed ICNC in the low-level liquid layer. As the cloud

385 droplets in the low-level feeder cloud were small ($<50\mu\text{m}$ in diameter, Fig. 7a), droplet shattering was likely not the responsible mechanism. However, as the temperature at 1900 m was around -3°C and large rimed particles (Fig. 14a) and cloud droplets larger than $25\mu\text{m}$ in diameter (Fig. 14b) were observed in the low-level liquid layer, the Hallett-Mossop process may have been active (Hallett and Mossop, 1974; Mossop, 1978). Another mechanism that could have led to the production of secondary ice particles in the low-level feeder cloud is ice particle fragmentation upon ice-ice collisions (e.g., Vardiman, 1978; Takahashi

390 et al., 1995). As the low-level liquid layer contained small pristine and large rimed ice particles (Fig. 13), which have different

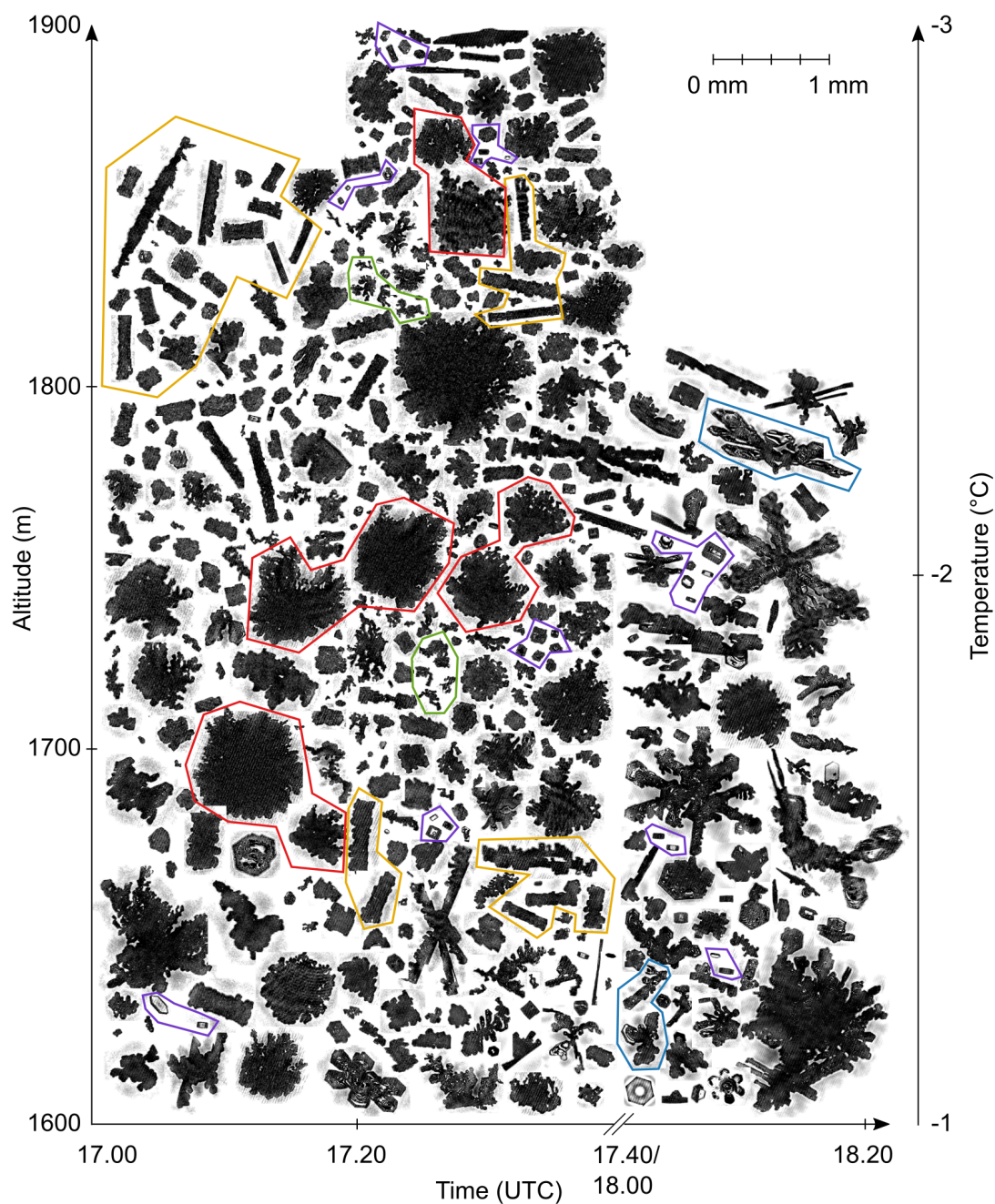


Figure 13. Example images of the ice crystals observed with HOLIMO as a function of height and time. The height-corresponding temperature is shown on the y-axis on the right side. The boxes indicate columns (yellow), pristine ice particles (purple), large rimed particles (red), irregular particles (green) and aggregates (blue).

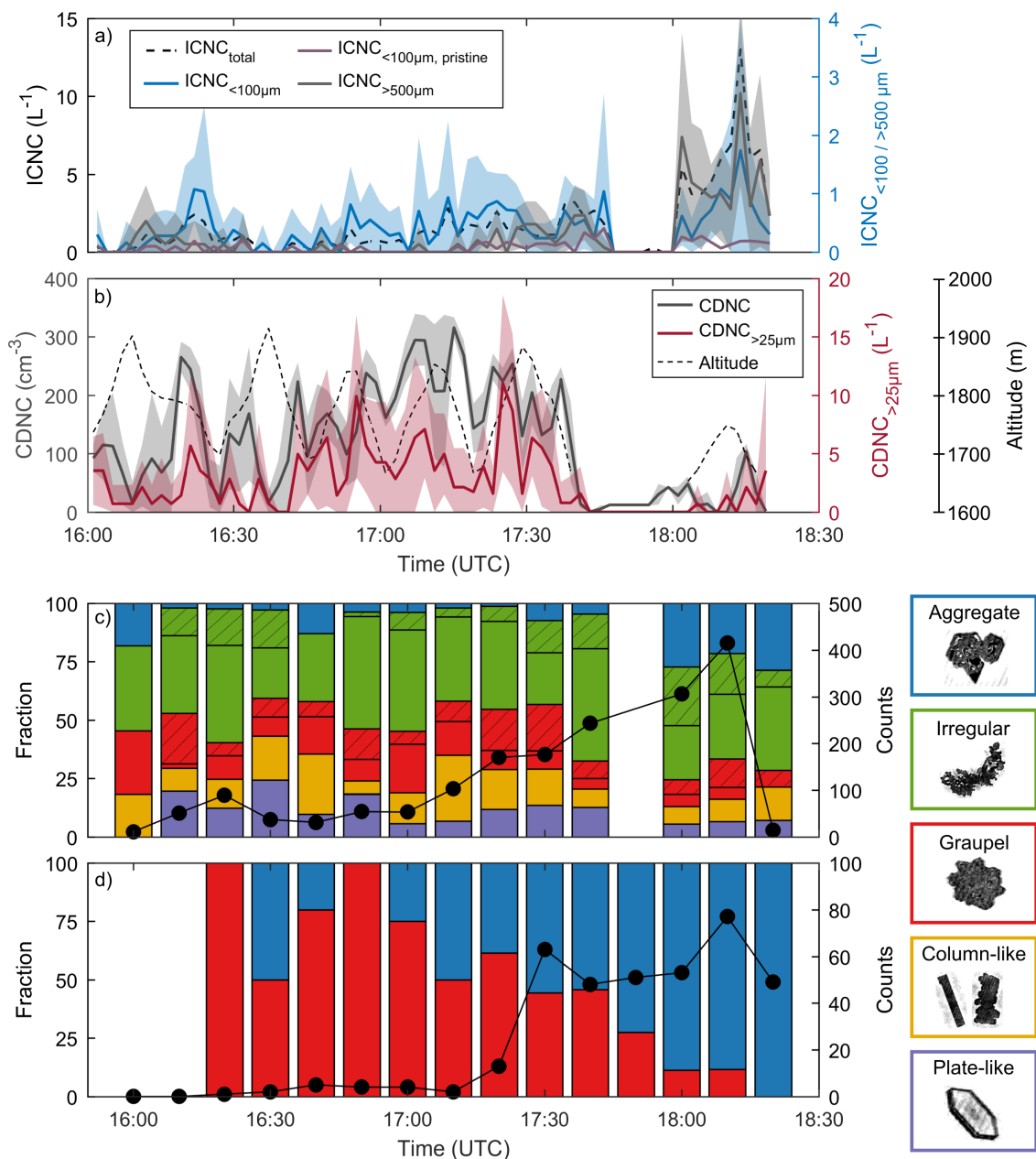


Figure 14. Timeseries of the ice (a) and liquid (b) cloud properties measured by the HoloBalloon platform. The shaded areas indicate the standard deviations. The dashed line in (b) shows the altitude of the balloon. The temporal evolution of the ice habit fraction is shown in (c, HOLIMO) and (d, MASC). The total counts during the 10 min-interval are indicated by the black dots. The ice particles observed with HOLIMO were classified into 5 categories: (1) plate-like, (2) column-like, (3) graupel (roundish particles), (4) irregular (unidentified habit) and (5) aggregates. Example ice particles are shown on the right. Only particles larger than $50\mu m$ were considered in the habit analysis. Shaded areas in (c) indicate particles with a higher degree of riming. The hydrometeors observed by the MASC were classified into graupel and aggregates/snowflakes.

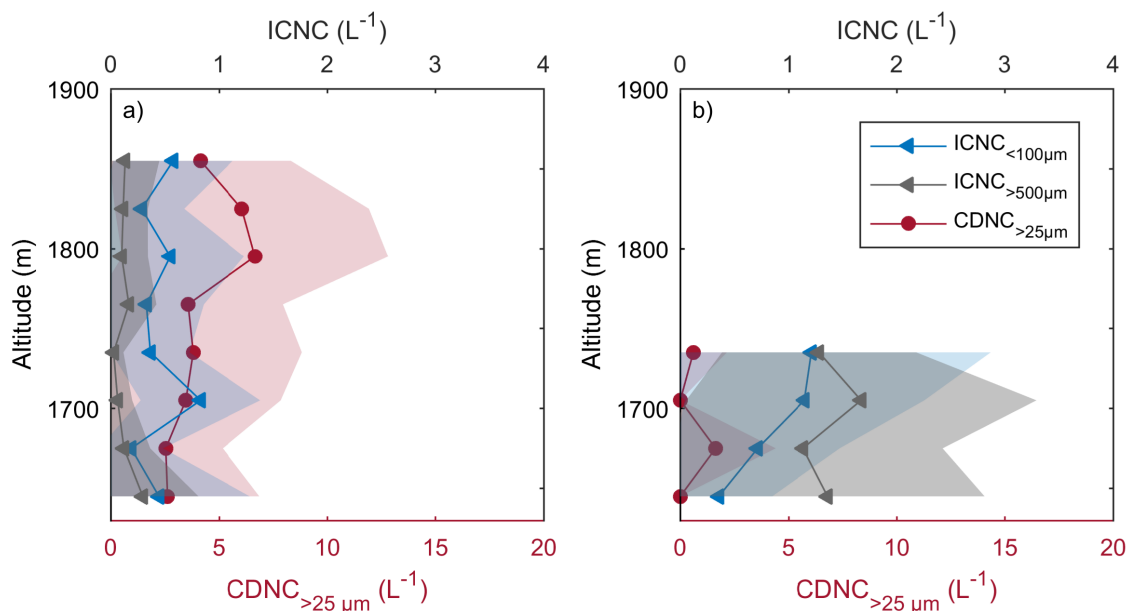


Figure 15. Mean vertical profiles of the liquid and ice properties measured between 16.30 - 17.30 UTC (a) and between 17:50 - 18:20 UTC (b) with the HoloBalloon platform. The shaded areas indicate the standard deviations.

terminal fall velocities and therefore enhanced collision efficiencies, this suggests that collisional ice fragmentation may have been occurring. Indeed, the ice crystal images in Figure 13 indicate the presence of ice fragments. Therefore, based on the temporal evolution of the cloud properties (Fig. 14a,b), we suggest that ice particle fragmentation upon collision was the dominant SIP process after 18 UTC, since the CDNC and in particular the number of large cloud droplets decreased after 395 18 UTC (Fig. 14b). In contrast, the presence of large cloud droplets ($>25\ \mu\text{m}$) before 18 UTC suggests that both the Hallett-Mossop process and collisional ice fragmentation contributed to the observed ICNC.

Previous studies have observed large discrepancies between the INP concentration and ICNC in the feeder region of clouds (e.g., Rogers and Vali, 1987; Lloyd et al., 2015; Beck et al., 2018; Lowenthal et al., 2019). These observations were frequently conducted at mountain-top research stations or near mountain slopes, where ICNCs of several hundreds to thousands per liter 400 have been reported (e.g., Rogers and Vali, 1987; Lloyd et al., 2015; Lowenthal et al., 2019). These large ICNCs were attributed to the influence of surface processes such as blowing snow (Rogers and Vali, 1987, Geerts et al., 2015), hoar frost (Lloyd et al., 2015), riming on snow-covered surfaces (Rogers and Vali, 1987) or ice crystal enhancement through turbulence and convergence (Beck et al., 2018), whereas the contribution of SIP processes has been suggested to be minor or has been difficult to assess (Lloyd et al., 2015, Beck et al., 2018). By performing balloon-borne measurements in a mountain valley, we measured 405 ICNC two order of magnitude lower than previous observations at mountain-tops ($1 - 10\ \text{L}^{-1}$ instead of $100 - 1000\ \text{L}^{-1}$) and thus were able to significantly reduce the impact of surface processes. Based on the observed INP concentration (Fig. 10) and ICNC (Fig. 14a), we suggest that SIP processes contributed up to $1-2\ \text{L}^{-1}$ to the observed ICNC and thus accounted for up to



50 % of the total ICNC before 18 UTC. On the other hand, the increase in the ICNC from 3 L^{-1} up to 15 L^{-1} after 18 UTC (Fig. 14a) cannot be solely explained by SIP within the feeder cloud, since the observed increase was primarily due to large ice particles ($> 300 \mu\text{m}$, see Fig. 7b). This increase in ICNC can likely be attributed to a change in the microphysics within the seeder region, which initiated the glaciation of the MPC.

If only a small concentration of secondary ice particles is captured by updrafts or turbulence within the feeder region and lifted aloft, they can initiate further ice formation and growth at temperatures well above typical INP activation temperatures and have a significant impact on the development of the cloud (e.g., cloud properties, glaciation, lifetime). While the CDNC decreased above 1850 m, the vertical profiles of the ICNC showed no height dependence over the 200 m height interval (Fig. 15a). This suggests that SIP was active over the entire low-level liquid layer. However, due to the limited vertical extent of the profiles, we cannot make a final statement regarding the impact of SIP within the feeder region on the cloud microphysics aloft. Further observations in 'surface-decoupled' environments (i.e., reduced influence of surface processes) with a larger vertical extent are required to assess the role of SIP in feeder clouds. This is important, as it can lead to the formation of precipitation in clouds which otherwise may not have produced significant precipitation.

5 Conclusions

In this paper, we investigated the microphysical evolution of a mixed-phase cloud passing over the Swiss Alps using a multi-dimensional set of observations and instruments including (1) ground-based remote sensing, (2) in situ microphysical observations on a tethered balloon system, (3) INP measurements and (4) surface precipitation measurements. A particular emphasis was placed on studying the microphysics within cloud top generating cells and a persistent low-level feeder cloud from an aerosol-cloud-precipitation perspective. The key findings are summarized as follows:

- The microphysical structure of the MPC was observed with a vertically-pointing Ka-band polarimetric cloud radar and with a tethered balloon system. The phase transition from a liquid to an ice cloud was observed during the passage of the cloud system. It is likely that the Wegener-Bergeron-Findeisen process contributed to the glaciation of the MPC. Regarding the vertical cloud structure, generating cells with enhanced radar reflectivity were observed near the cloud top, which acted as a seeder region and produced fallstreaks of enhanced radar reflectivity. Furthermore, the decoupled boundary layer circulation in the lee of the mountain barrier produced local updrafts and turbulence, which led to the formation of a persistent low-level feeder cloud.
- The cloud radar and microwave radiometer observations suggest that ice formation and growth as well as liquid water production was enhanced within cloud top generating cells. While numerous studies have observed enhanced ICNCs within generating cells, uncertainties exist regarding their ice formation mechanism. Here we proposed different processes and discussed their potential contribution. Cooling associated with convective overshooting was suggested to increase the INP concentration by a factor of 3.3, whereas radiative cooling was estimated to increase the ICNC formed by primary ice nucleation only by a factor of 1.2. In addition, secondary ice production through droplet shattering was



440 proposed to potentially increase the ICNC by several orders of magnitude and might have contributed to the glaciation of the MPC.

– The co-existence of small pristine ice crystals and large rimed ice particles was observed in the low-level feeder cloud, suggesting the occurrence of secondary ice production. By using a tethered balloon to observe the feeder cloud in the mountain valley, we were able to significantly reduce the influence of surface processes compared to previous obser-
445 vations at mountain tops and to investigate the contribution of secondary ice production in the feeder region of clouds. The ICNC of small ice crystals ($< 100 \mu\text{m}$) measured near cloud base exceeded the INP concentration by three orders of magnitude. Conditions favorable for the Hallett-Mossop process and ice particle fragmentation upon ice-ice collisions were found. We suggest that secondary ice production in the feeder cloud increased the ICNC by a factor of 2.

Overall, this study observed the temporal and spatial evolution of the microphysics within the seeder and feeder region of a
450 MPC passing over the Swiss Alps. We found that a significant increase in ice formation and growth within the seeder region can induce the glaciation of the MPC. In addition, we found that secondary ice production mechanisms were active in the feeder cloud, which initiated ice formation at temperatures where no INP were detectable. This case study demonstrates that secondary ice production can occur in different cloud regions and have important implications for precipitation initiation and the lifetime of MPCs in general. Further studies are required to understand the role of secondary ice production both in the
455 seeder and feeder regions of clouds. These studies should include vertically-resolved in situ observations of the microphysical properties, aerosol properties (e.g., INP) and environmental conditions (e.g., temperature, vertical updraft velocity) over the entire cloud depth and should be performed in a 'surface-decoupled' environment (i.e., reduced influence of surface processes).

Appendix A: The use of the maximum Doppler velocity as a proxy for regions with updrafts and liquid water

In the framework of the present study, the maximum Doppler velocity was used as a proxy to identify regions with updrafts
460 and liquid water. The maximum Doppler velocity v_{max} was derived from the Doppler spectra as shown in Figure A1a. In order to be more robust regarding the presence of extreme values, v_{max} was defined as follows:

$$v_{max} = \text{maximum Doppler velocity where } Z \geq (Z_{min} + 0.1 \cdot (Z_{max} - Z_{min})) \quad (\text{A1})$$

where Z_{min} and Z_{max} are the minimum and maximum radar reflectivity. To validate whether v_{max} can also be used to identify regions with liquid water, it was compared to the LWP measured by the microwave radiometer. Since the LWP is integrated over
465 the whole vertical column, the vertically-integrated v_{max} is shown in Figure A1b. A positive correlation was found between v_{max} and the LWP with a Spearman rank correlation coefficient of 0.5 significant at the 5% level. Based on this result, we assume that v_{max} can be used as a proxy for updraft regions and regions with liquid water.

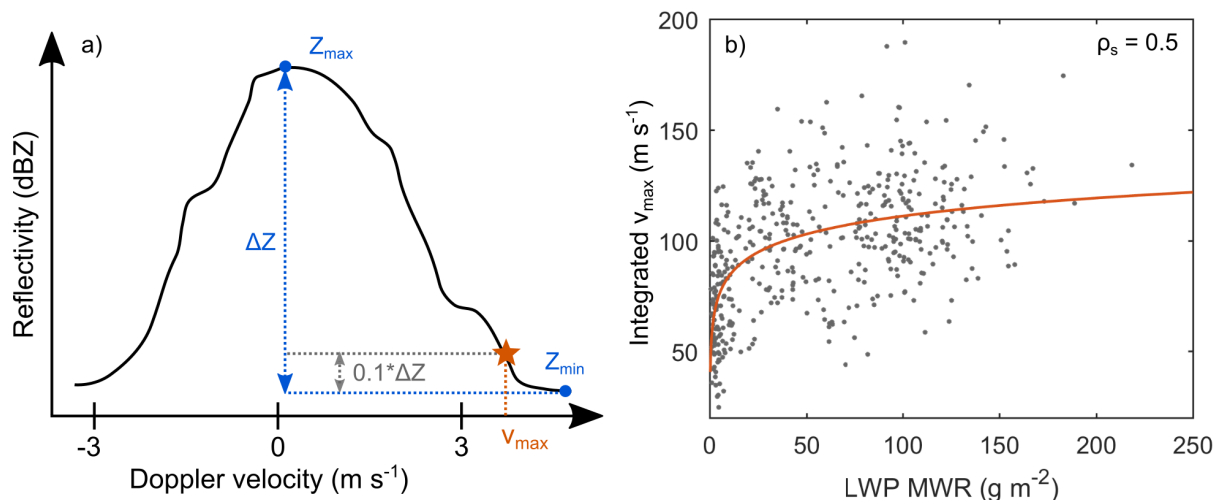


Figure A1. An example Doppler spectrum is shown in (a) to demonstrate the derivation of the maximum Doppler velocity v_{max} (orange star), where Z_{min} and Z_{max} are the minimum and maximum radar reflectivity (see text for more details). The relationship between the vertically-integrated v_{max} and the LWP measured by the microwave radiometer for the time period 15 - 18 UTC is shown in panel (b). The orange line is a logarithmic fit through the data points. ρ_s indicates the Spearman's rank correlation coefficient.

Appendix B: Cloud properties inside and outside of generating cells

To estimate the contribution of generating cells to the observed radar reflectivity signal, the cloud properties inside (a) and
470 outside (b) of GC2 were compared (Fig. B1). The region inside the generating cell was determined by tracking the fallstreak of
enhanced radar reflectivity, whereas the region outside of the generating cell was defined by the region between two adjacent
cells that was characterized by a lower radar reflectivity (see Fig. 8a). The difference between the Doppler spectra inside and
outside of generating cells (Fig. B1c) suggests that the generating cell at 16:45 UTC contributed 10 dBZ to the radar reflectivity
and thus enhanced ice and precipitation formation. For a quantitative analysis, a large number of generating cells needs to be
475 analyzed, which was beyond the scope of the present study.

Author contributions. FR analyzed the observational data and prepared the figures of the manuscript. FR, JH, JP and AL performed the HoloBalloon measurements. JB computed the ICNC retrievals from the remote sensing observations. MR performed the peakTree analysis. PS processed the remote sensing data and Doppler spectra of MIRA-36. JW and CM collected and processed the INP data. RE operated the OCEANET container during the RACLETS campaign. MH operated the radar wind profiler and processed the wind profiler data. JB,
480 MR, PS helped in interpreting the remote sensing data. FR, JH, RD and UL analyzed and interpreted the observational data. FR prepared the manuscript with contributions from all authors.

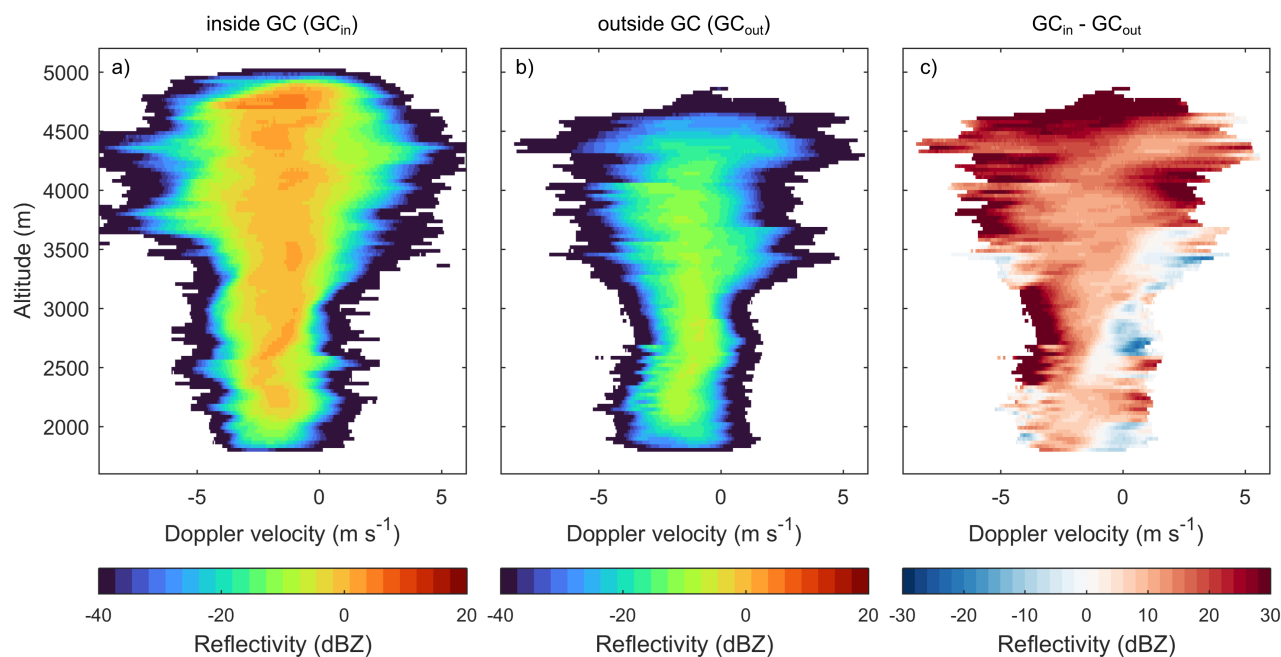


Figure B1. Vertical profile of the Doppler spectra inside (a; averaged over the region GC_{in}) and outside (b; averaged over the region GC_{out}) of the generating cell at 16:45 UTC (GC_{in} and GC_{out} are specified in Fig. 8a). The difference in the Doppler spectra between GC_{in} and GC_{out} is shown in panel (c).

Competing interests. The authors declare that they have no conflict of interest.

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