

# Sporadic sodium layer: A possible tracer for the conjunction between the upper and lower atmospheres

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**Abstract.** In this research, we reveal the inter-connection between lightning strokes, reversal of the electric field, ionospheric disturbances, and a sodium layer (N<sub>s</sub>), based on the joint observations by two lidars, an ionosonde, an atmospheric electric mill, a fluxgate magnetometer, and the World Wide Lightning Location Network (WWLLN). Our results suggest that lightning strokes would probably have an influence on the ionosphere and thus affecting the occurrence of N<sub>s</sub>, with the overturning of electric field playing an important role. Statistical results reveal that the sporadic E layers (E<sub>s</sub>) could hardly be formed or maintained when the atmospheric electric field turns upward. A conjunction between the lower and upper atmospheres could be established by these inter-connected phenomena, and the key processes could be suggested as follows: lightning strokes→overturning of electric field→depletion of E<sub>s</sub>/generation of N<sub>s</sub>.

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**Keywords:** sporadic sodium layers, sporadic E layers, atmospheric circuit, lightning stroke, electric field

## 1 Introduction

The upper mesosphere-lower thermosphere (MLT) region, is the junction for momentum and energy exchanges between the Earth's low atmosphere and outer space. However, on account of the limitations of detection methods, this region remains the least known part of our planet's atmosphere (Wang, 2010). Fortunately, the metal layers (especially the sodium layer), which located between about 80 ~ 110 km, could possibly act as a window to detect the MLT parameters by means of fluorescence resonance lidars (Gardner et al., 1986; Gong et al., 2002; Gong et al., 1997). With an active chemical property and high abundance of sodium atoms, the sodium layer has been widely observed and studied all over the world (Marsh et al., 2013; Collins et al., 2002; Plane, 2003; Plane et al., 1999). The sporadic sodium layer (SSL or Na<sub>s</sub>), with the neutral sodium density that could double within several minutes, is the most fantastic phenomenon observed from the sodium layer. Since first reported in 1978 (Clemesha et al., 1978), many mechanisms, involving meteor injection (Clemesha et al., 1980), dust reservoir (von Zahn et al., 1987), recombination of ions and electrons in sporadic E layer (E<sub>s</sub>) (Cox and Plane, 1998), and high temperature theory (Zhou et al., 1993), have all been proposed. Because the Na<sub>s</sub> is suggested to have a connection to so many atmospheric parameters, the metric or phenomenon could be appropriate in acting as a tracer for studying inter-connection between the middle and upper atmospheres. Up to now, a large number of observations report a diversity of the Na<sub>s</sub> features, but the exact mechanism for Na<sub>s</sub> is probably still uncertain (Collins et al., 2002; Cox et al., 1993; Daire et al., 2002; Gardner et al., 1995; Qiu et al., 2015; Zhou and Mathews, 1995; Zhou et al., 1993).

Among all the proposed mechanisms, the E<sub>s</sub> theory is supported by abundant observations and results from numerical simulations (Cox and Plane, 1998; Daire et al., 2002; Dou et al., 2009; Dou et al., 2010; Gardner et al., 1993; Gong et al., 2002; Kane et al., 2001; Kane et al., 1993; Kane et al., 1991; Kirkwood and Nilsson, 2000; Kwon et al., 1988; Mathews et al., 1993; Miyagawa et al., 1999; Nagasawa and Abo, 1995; Nesse et al., 2008; Shibata et al., 2006; Williams et al., 2006). The key process of E<sub>s</sub> theory is the recombination of ions and electrons in the E<sub>s</sub> layer while descending to lower altitudes (Cox and Plane, 1998; Daire et al., 2002). The E<sub>s</sub> layer is mainly influenced by the vertical wind shear (Abdu et al., 2003; Clemesha et al., 1998; Haldoupis et al., 2004; Mathews, 1998; Šauli and Bourdillon, 2008; Wakabayashi and Ono, 2005), the geomagnetic field (Resende et al., 2013; Resende and Denardini, 2012; Zhang et al., 2015; Denardini et al., 2016), and the electric field (Abdu et al., 2003; Damtie et al., 2003; Haldoupis et al., 2004; Kirkwood and Nilsson, 2000; Kirkwood and von Zahn, 1991; Macdougall and Jayachandran, 2005; Matuura et al., 2013; Nygren et al., 2006; Parkinson et al., 1998; Takahashi et al., 2015; Voiculescu et al., 2006; Wakabayashi and Ono, 2005; Wan et al., 2001; Wilkinson et al., 1993). In the Northern Hemisphere, the E<sub>s</sub> layer would descend to a lower altitude during southward electric field (Abdu et al., 2003; Damtie et al., 2003; Haldoupis et al., 2004; Kirkwood and Nilsson, 2000; Kirkwood and von Zahn, 1991; Macdougall and Jayachandran, 2005; Nygrén et al., 2006; Parkinson et al., 1998; Takahashi et al., 2015; Voiculescu et al., 2006; Wakabayashi and Ono, 2005; Wan et al., 2001; Wilkinson et al., 1993), and observations in the polar cap suggest the electric field reversal have an influence on the probability of E<sub>s</sub> occurrences (Macdougall and Jayachandran, 2005).

On the other hand, the atmospheric electric circuit is a closed loop (Driscoll et al., 1992; Jánský and Pasko, 2014; Lv et al., 2004; Roble and Hays, 1979; Rycroft and Harrison, 2012; Rycroft et al., 2000; Suparta and Fraser, 2012; Tinsley, 2000), like a capacitor with a positive plate (e.g., the ionosphere) and a negative panel (e.g., the ground), and dielectric medium between them (e.g., the neutral atmosphere). Then the global atmospheric electric circuit formed in the capacitor, with the lightning phenomena generating an upward current (with the atmospheric electric field intensity  $E < 0$ ) and returning a downward current ( $E > 0$ ) under fair weather condition. Nowadays increasing and emerging evidences are pointing to the close link between the upper atmosphere (e.g., the positive plate) and lower atmosphere (e.g., the negative plate) (Harrison et al., 2010; Rycroft, 2006). For example, thunderstorm occurring in the lower atmosphere is suggested to have a direct impact on the  $E_S$  layer based on recent observational results (Bortnik et al., 2006; Christos, 2018; Cummer et al., 2009; Curtius et al., 2006; Davis and Johnson, 2005; Davis and Lo, 2008; England et al., 2006; Fukunishi et al., 1996; Girish and Eapen, 2008; Haldoupis et al., 2012; Immel et al., 2013; Kumar et al., 2009; Kuo and Lee, 2015; Lay et al., 2015; Mangla et al., 2016; Maruyama, 2006; Pasko et al., 2002; Rodger et al., 2001; Rycroft, 2006; Satori et al., 2013; Sentman and Wescott, 1995; Shao et al., 2013; Sharma et al., 2004; Su et al., 2003; Surkov et al., 2006; Yu et al., 2015) or even the sodium layer (Yu et al., 2017). The possible carriers or phenomena connecting the thunderstorm to the upper atmosphere are suggested to be atmospheric tides (England et al., 2006; Haldoupis et al., 2004; Immel et al., 2013), planetary waves (Lv et al., 2004), gravity waves (Davis and Johnson, 2005; Kumar et al., 2009; Lay et al., 2015; Shao et al., 2013), transient luminous event (TLEs) (Cummer et al., 2009; Fukunishi et al., 1996; Haldoupis et al., 2012; Pasko, 2008; Pasko et al., 2002; Sentman and Wescott, 1995; Sharma et al., 2004; Su et al., 2003), the solar activity (Zhang et al., 2020), and also the electric fields (Bortnik et al., 2006; Davis and Johnson, 2005; Davis and Lo, 2008; Immel et al., 2013; Kuo and Lee, 2015; Maruyama, 2006; Rycroft, 2006; Satori et al., 2013; Shao et al., 2013).

In this research, we apply five joint observations for our case studies and statistical works: (1) Two lidars at Hefei (31.8°N, 117.2°E), providing observations of sodium density, mesopause temperature and zonal wind; (2) An ionosonde in Wuhan (30.5°N, 114.6°E), detecting the  $E_S$  and ionospheric echoes in different modes; (3) An atmospheric electric mill (30.5°N, 114.5°E), giving simultaneous electric field variations; (4) A fluxgate magnetometer (30.5°N, 114.5°E), probing the H, D, and Z magnetic field components; and (5) The World Wide Lightning Location Network (WWLLN), observing the location and power of a lightning stroke. The purpose of this study is to examine the possibility of  $Na_S$  acting as a practical robust tracer for the conjunction between the upper and lower atmospheres. Our results suggest that lightning strokes may have an influence on the lower ionosphere leading to the occurrence of  $Na_S$ , with the atmospheric electric field probably playing an important role.

## 2 Observations and results

### 2.1 An Na<sub>s</sub> event during the overturning of electric field

A sporadic sodium layer is detected on the night of June 3<sup>rd</sup>, 2013, by both the USTC sodium lidar (Dou et al., 2009) and the nearby Temperature/Wind (T/W) lidar (Li et al., 2012) (in Fig. 1a and 1b, the same Na<sub>s</sub> event detected with different resolutions). The peak density of Na<sub>s</sub> detected by the wideband lidar is more than 12000 cm<sup>-3</sup>, while the peak density observed by narrowband lidar is about 5000 cm<sup>-3</sup>. The discrepancy of peak densities may come from the wideband lidar. Fig. 1(a) shows the wideband lidar system operated poorly after the midnight. For the sodium lidar, the inversion formula for the sodium number density  $N$  at an altitude of  $z$  is given as follows:

$$N = \frac{\sigma_R n_a(z_0)}{\sigma_{Na}} \cdot \frac{(P(z) - P_B)z^2}{(P(z_0) - P_B)z_0^2}, \quad (1)$$

where  $\sigma_R$  is the Rayleigh backscatter cross section;  $n_a(z_0)$  is the atmosphere density at a reference altitude, given by atmospheric model;  $\sigma_{Na}$  is the effective sodium backscatter cross section;  $P(z)$  is the number of photons detected in the range interval  $(z - \Delta z/2, z + \Delta z/2)$ ;  $P_B$  is the expected photon count per range bin due to background signal and dark counts, calculated through the averaged background signal above 130 km; and,  $P(z_0)$  is the Rayleigh photon count at 30 km altitude, estimated by averaging the measured photon count over a 5-km range interval centered at 30 km (Gardner et al., 1986; Xue, 2007). The averaged photon count  $P(z_0)$  at 30 km is  $P(z_0) = 106$  (smaller than that on normal days), and the expected photon count  $P_B$  at 130 km equals to 18 (larger than that on normal days). Since the error term is inversely proportional to the absolute value of  $P(z_0) - P_B$ , a much smaller  $|P(z_0) - P_B|$  would cause an increase in the deduced sodium number density  $N$ .

This Na<sub>s</sub> event occurs much higher above the centroid height of sodium layer (normally at about 92 km). Both Fig. 1a and 1b show the sodium density begins to increase at about 13:20 UT, while the largest intensity of sodium enhancement occurs from about 14:20 UT, with a peak density on 97.75 km at 14:40 UT. The simultaneous temperature observation by the T/W lidar reveals this Na<sub>s</sub> occurs in a cold region (Fig. 1c), so the high temperature mechanism appears to be inapplicable for this event.

On the other hand, the zonal wind exhibits a suitable wind shear for creating those sporadic E layers, with a westward wind above and an eastward wind below (Fig. 1d). The E<sub>s</sub> layer is predicted to form around the border of the wind shear. Observations by the ionosonde at Wuhan indeed show active sporadic E layers on that day (Fig. 2a and b). The E<sub>s</sub> series keep travelling/propagating downward starting around 6:30 UT, and then disappear at last while the Na<sub>s</sub> occurs coincidentally on about 13:20 UT. So this Na<sub>s</sub> is better explained by the E<sub>s</sub> mechanism, in accord with our previous study which shows that a Na<sub>s</sub> higher than 96 km tends to be controlled by the E<sub>s</sub> mechanism (Qiu et al., 2016). Although the content of sodium ions in E<sub>s</sub> layers seemed to have insufficient concentration (von Zahn et al., 1989), it has also been proposed that the ions could be concentrated by the wind shear effectively (Clemesha et al., 1999; Cox and Plane, 1998; Nesse et al., 2008). On the other hand, laboratory results show that the ligand complexes of Na<sup>+</sup>X would form and thus speed up the recombination of ions in

the mesopause condition (Collins et al., 2002; Cox and Plane, 1998; Daire et al., 2002). The calculated reaction rate suggests the formation of cluster ions is enhanced at lower temperatures, in accordance with the cold region observed in Fig. 1c where the sporadic sodium layer occurs.

More details about the atmospheric parameters are shown in Fig. 3. The time series of sodium density on the peak height display a sharp enhancement from 14:20 UT (marked by the vertical red dashed line in Fig. 3a). The atmospheric electric field detected by the mill exhibits an overturning at around 14:20 UT, alternating from downward direction to upward (Fig. 3b). It can be clearly observed that the enhancement of sodium density occurs coincidentally with the overturning of electric field, as highlighted by the vertical red dashed line in Fig. 3. A nearby fluxgate magnetometer provides the horizontal magnetic field  $H$  (nT) (Fig. 3c), showing disturbances at 14:15 UT. The total magnetic intensity  $B$  could be deduced by the  $H$ ,  $Z$ , and  $D$  components ( $B = \sqrt{H^2 + D^2 + Z^2}$ ) from fluxgate magnetometer observations (the calculated values are plotted in Fig. 3d).

It is worth noting that the overturning of atmospheric electric field discussed here is theoretically rough, since the electric field at the lower ionosphere will be modulated as well (e.g., with a value of several mV/m [Seyler et al., 2004]). Nevertheless, model simulations from the electrodynamics show that the upward electric field in upper atmosphere is proportional to the source current in the troposphere (Driscoll et al., 1992), and that the upward current would continue transmitting to the heights of 100~130 km of the dynamo region where  $E_S$  occurs most frequently (Rycroft et al., 2012). The model, based on rocket observations, shows that the atmospheric electric field has a similar scale and the same polarity from the ground to the altitude of ionosphere (Abdu et al., 2003). Thus, the electric field detected by a ground-based mill could reasonably be a reflection of the actual situation in the lower ionosphere, at least for the trends of variations.

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## 2.2 Statistical results of the $E_S$ and overturning of electric field

A statistical correlation between the overturning of electric field and  $E_S$  variations is summarized in Table 1. The fo $E_S$  values refers to the Critical Frequency, at which the reflection starts while the radio frequency equals the plasma frequency ( $\text{fo}E_S = \omega_{pe} = \left(\frac{n_e e^2}{m_e \epsilon_0}\right)^{\frac{1}{2}} \approx 9\sqrt{10^{-6}n_e}$ , with fo $E_S$  in MHz and  $n_e$  in  $\text{cm}^{-3}$ ) (Bittencourt, 2004). From 2012 to 2014, in the summer season (from May to August, when the  $E_S$  layers occur frequently), there are 242 days with effective observations of  $E_S$  and electric field with the overturning feature. Among all the cases, for about 155/242 days the fo $E_S$  get interrupted during the overturning of electric field, while the fo $E_S$  decreases on 39/242 days. In comparison, fo $E_S$  appears only for 6/242 days, and increases for 26/242 days. Another 28/242 days show no distinguishing features of the  $E_S$  during an overturning of the electric field. These ratios are overlapped, since sometimes there are more than one overturning on a single day. Thus in general, these results suggest that  $E_S$  could hardly be formed for upward electric field situations and that indeed the overturning of electric field causes a depletion of  $E_S$ . This statistical result may also give an implication that the electric field variations detected by ground based mill have a feasible link with the lower ionosphere. It is worth mentioning that there is

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indeed no one-to-one correspondence between the electric field overturning and  $E_s$  depletion, since the  $E_s$  is also influenced by other key parameters and phenomena including wind shears, tides and gravity waves. Sometimes the overturnings recover quickly, without sufficient time for producing an obvious effect on  $E_s$ . Therefore a deeper analysis will be needed in future statistical studies.

## 5 3 Discussions

### 3.1 Possible influences by the atmospheric electric circuit

The atmospheric electric circuit is formed by the ionosphere and ground surface with the dielectric medium (e.g., the neutral atmosphere) sandwiched between them (Driscoll et al., 1992; Harrison, 2020; Jánský and Pasko, 2014; Lv et al., 2004; Roble and Hays, 1979; Rycroft and Harrison, 2012; Rycroft et al., 2000; Rycroft et al., 2012; Rycroft et al., 2007; 10 Suparta and Fraser, 2012; Tinsley, 2000). The lightning phenomena and thunderstorms, acting as the electric generator for the circuit, drive an upward current to the ionosphere. In fair weather regime, the electric field directs downward to the earth surface ( $E > 0$ ), making a closed global electric circuit (see Fig. A1 in Appendix). The electric field could vary through two distinct ways as follows: The first one is the changing magnetic field explained by the Faraday's law (e.g.,  $\nabla \times \vec{E} = -\frac{\partial \vec{B}}{\partial t}$ ). However, observations by the fluxgate magnetometer show that there is just a small disturbance of magnetic field during the 15 overturning of electric field. The other way is the electrostatic induction following the Coulomb's law ( $\vec{E} = \frac{1}{4\pi\epsilon_0} \frac{Q}{r^2} \hat{r}$ ). The connection between the lightning stroke and the overturning of electric field could be explained by a classic thunderstorm charge model through the electric imaging method based on the Coulomb's law (i.e., this model could be supported by a classic electrodynamics textbook written by D.J. Griffiths, 1999). A typical thundercloud (e.g., pairs of ( $Q_1, -Q_2$ ) or ( $-Q_3, Q_4$ ) in Fig. A1), with a dipole of positive charge located above a negative charge part, would produce an upward electric 20 field toward the ionosphere (see more details in Appendix A).

According to the observations from WWLLN, we find two regions (red ovals A and B in Fig. 4) with heavy lightning activities during the period of the  $N_{as}$ . Before the  $N_{as}$  occurrences, there were only a few powerful lightnings detected within about (25.1°N ~ 35.8°N) and (113.8°E ~ 118.1°E) during the period of 12 UT to 13:15 UT (just one strong stroke with a power of 43720.25 kW happening on 12:17 UT, at 25.7229°N and 117.3955°E). The continuous strongest lightnings with a 25 power larger than  $10^4$  kW occur from 13:19 UT to 15:43 UT, mainly concentrating in two areas centered around (35.8°N, 118.1°E) and (25.1°N, 113.8°E). After 15:45 UT, no strong strokes were detected again within this area. Thus, the pairs of ( $Q_1, -Q_2$ ) and ( $-Q_3, Q_4$ ) could be referred to the lightning area of part A and B in Fig. 4. Since thunderstorms could trigger the breakdown process within a rather large area (Leblanc et al., 2008) and influenced the ionosphere with around more than 800 km range horizontally away from the lightning center (Johnson and Davis, 2006; Johnson et al., 1999), the whole area 30 above might undergo a breakdown easily around  $Q_1 \sim Q_4$  (e.g., the whole shadow zone in Fig. 4, involving the two strongest lightning areas and the two observing stations).

In previous studies, lightning strokes in lower atmosphere were reported to cause a reduction of electrons of the ionosphere (Shao et al., 2013), and in reverse an enhancement of sodium density in the metal layer (Yu et al., 2017). Those two scenarios are in accord with our current results presented above, with a depletion of  $E_S$  and a consequential occurrence of  $N_{as}$ . Although such an idea/picture has been proposed long time ago (Griffiths, 1999), this is the first time that one can  
 5 apply the imaging method for observing thunderstorms to explain the link between upper and lower atmospheres through an overturning of upward electric field.

Furthermore, the results from different channels of Wuhan ionosonde exhibit extraordinary echoes in different modes during the lightning period (Fig. 5a to 5l). (a) ~ (c): From 13:15 UT to 13:45 UT, the echoes gradually increase. Note that the powerful lightning period begins on 13:15 UT as well, with the sodium density enhancement and the  $E_S$  depletion occurring  
 10 on about 13:20 UT. (d) ~ (g): Most intense echo signals occur during 14:00 UT to 14:45 UT, while the largest intensity of sodium enhancement begins at 14:20 UT and the sodium density peaks at 14:40 UT. The overturning of electric field also occurs at 14:20 UT. (h) ~ (j): From 15:00 UT to 15:30 UT, the signals weaken gradually; (k) ~ (l): The echoes vanish after 15:45 UT. Afterwards, no strong stroke detected again in the discussed area. Meanwhile, the ionospheric echoes diminish after 15:45 UT, and the overturning of electric field also recovers at about 15:30 UT. Thus, in this case the ionospheric  
 15 echoes and the lightning activities exhibit an obvious synchronous behavior.

### 3.2 Possible mechanism for $N_{as}$

Usually, the mid-latitude  $E_S$  layers would be brought down gradually by tidal fluctuations (Mathews, 1998). The  $E_S$  theory predicts that when a series of  $E_S$  layers descend below 100 km, they will be depleted through the recombination of  
 20 ions and electrons (Cox and Plane, 1998). Since the recombination of  $Na^+ + e^- \rightarrow Na + h\nu$  is inefficient to generate  $N_{as}$ ,  $Na^+$  is believed to first form a ligand  $Na^+ \cdot N_2$  through the recombination reaction:



with a rate coefficient of  $k_1 = 4.8 \times 10^{-30} (T/200)^{-2.2} \text{ cm}^6 \text{ molecule}^{-2} \text{ s}^{-1}$  (Cox and Plane, 1998).  $Na^+ \cdot N_2$  can either switch with  $CO_2$  (which will undergo dissociative electron recombination to form  $Na$ ), or  $O$  (which reforms  $Na^+$ ) (Cox and  
 25 Plane, 1998). Thus the key factor of  $E_S$  mechanism depends on the ratio of  $[O]/[CO_2]$ : Recombination of  $Na^+ \cdot CO_2$  and  $e^-$  will increase rapidly as  $[O]/[CO_2]$  decreases below a value of 100 (Cox and Plane, 1998). Then the sodium atoms could be formed directly from the following chemical reaction:



The chemical reaction rate ( $v$ ) for this second-order reaction could be calculated using the following equation:

$$30 \quad v = k[Na^+ \cdot CO_2]N_e, \quad (4)$$

The reaction rate coefficient  $k_2$  for the chemical reaction (3) is experimentally measured to be:

$$k_2 = 1 \times 10^{-6} \sqrt{\frac{200}{T}} \text{ (cm}^6 \text{ molecule}^{-2} \text{ s}^{-1}) \quad (5)$$

(Collins et al., 2002; Cox and Plane, 1998; Daire et al., 2002), and the electron density  $N_e$  could be calculated using the following equation:

$$N_e = 1.24 \times 10^4 foE_s^2 (cm^{-3}) \quad (6) \text{ (Bittencourt, 2004).}$$

Overall, this  $E_s$  mechanism is most widely accepted, if we neglect  $k_1$  for reaction (2) as being too small with an order of  
 5  $10^{-30}$ . A possible adaptation is to assume a plenty quantity of pre-existing  $Na^+ \cdot N_2$  or  $Na^+ \cdot CO_2$  in the sodium layer, and then the  $E_s$  just needs to provide enough additional electrons. Figure 2(b) shows  $E_s$  descending near 100 km at about 13:20 UT. Then the  $E_s$  depletes, and a moderate enhancement of Na occurs from 13:30 UT to 14:00 UT (shown in Figure 1(a) and 3(a)). This increase in sodium density exhibits no obvious peak, which could probably be in accord with a normally descending  $E_s$  governed by tides. In comparison, the peak profile of the  $Na_s$  shows intense enhancement and sharp peak,  
 10 indicating a distinct mechanism.

On the other hand, a link between the reverse electric field and  $E_s$  variations could be established through the acceleration of electrons. Normally, positive particles will move along the direction of electric field, and negative particles do opposite (Griffith and College, 1999). Since metal ions are much heavier than electrons, the ions would drag electrons to move/drift together, which is called the bipolar diffusion (Griffiths, 1999). In the initial stage, ions and electrons descend  
 15 gradually under the southward electric field. In a partially ionized plasma, the characteristic frequencies for ions and electrons are associated with the collisions of the plasma particles with stationary neutrals (e.g., the electron–neutral collision frequency  $\nu_{en}$  and the ion–neutral collision frequency  $\nu_{in}$ ). The collision frequency  $\nu_{sn}$  for scattering of the plasma species  $s$  by the neutrals is

$$\nu_{sn} = n_n \sigma_s^n V_{Ts}, \quad (7) \text{ (Shukla and Mamun, 2002)}$$

20 where  $n_n$  is the neutral number density,

$\sigma_s^n$  is the scattering cross section (which is typically of the order of  $5 \times 10^{-15} \text{ cm}^2$  and depends weakly on the temperature  $T_s$ ),

and  $V_{Ts} = (k_B T_s / m_s)^{1/2}$  is the thermal speed of the species  $s$ .

So the relaxation times  $\tau = \frac{1}{\nu}$  for ions and electrons are different in a partially ionized plasma, and electrons would respond  
 25 much faster than the heavier sodium ions do (since  $m_i \gg m_e$ ). At the moment when the electric field reverses, electrons will be rapidly accelerated by the northward electric field, and ions would be regarded as essentially remaining northward or unchanged. If we could possibly assume that  $Na^+ \cdot CO_2$  is always excessive (as we see from estimates below it just needs to be a number density of  $100 \text{ cm}^{-3}$ , but regrettably there is no direct measurements and detections up till now), we only need to consider the amount of electrons. When the reactant concentration for R(3) is increased (e.g., when a concentrated electron  
 30 layer accelerated downward below 100 km), this reaction will shift to the right side of R(3). To calculate the chemical reaction rate, we assumed a pre-existing concentration of  $100 \text{ cm}^{-3}$  for  $[Na^+ \cdot CO_2]$ , and used the observed value of 3.1 MHz for fo $E_s$  (means  $N_e = 1.2 \times 10^5 \text{ cm}^{-3}$ ), and 170 K for T. The calculated rate is  $\nu = 13 \text{ cm}^{-3} \text{ s}^{-1}$ , in accord with the required source strength of sodium atoms of 3 sodium atoms  $\text{cm}^{-3} \text{ s}^{-1}$  for the formation of  $Na_s$  (Cox et al., 1993). If we deduct some



influences by eddy diffusion and loss of sodium atoms, this chemical rate can generate a  $\text{Na}_s$  within several minutes. Thus, no matter how many  $\text{Na}^+$  ions are in  $E_s$ , the electrons in  $E_s$  are always sufficient to produce  $\text{Na}_s$ . Perhaps that is the reason why we often observe even a very weak  $E_s$  is accompanied by  $\text{Na}_s$  (Dou et. al., 2010).

Based on the above results, a possible mechanism for  $\text{Na}_s$  could be suggested by the following four steps: (1) Strong lightning strokes produce an upward atmospheric electric field toward the ionosphere; (2) The reverse of electric field would accelerate the electrons in  $E_s$  layer to descend rapidly; (3) The concentrated electrons meet with pre-existing sodium reservoir of  $\text{Na}^+ \cdot \text{CO}_2$  below 100 km, making a faster recombination through the most efficient reaction of  $\text{Na}^+ \cdot \text{CO}_2 + e^- \rightarrow \text{Na} + \text{CO}_2$ ; (4) The depleted  $E_s$  layers generate the formation of  $\text{Na}_s$ . Thus, we propose that there would probably be a connection between the lightning strokes, overturning of the electric field, ionospheric disturbances, and also the  $\text{Na}_s$ . A link between the lower and upper atmospheres could be established by carefully studying and examining these phenomena. However, it is worth noting that the key processes during step (2) and (3) remained still quite uncertain. A more in-depth modelling study concerning both plasma and neutral molecules is needed in the future.

#### 4 Conclusions

In this research, we study the conjunction between the lower and upper atmospheres, through the phenomena and processes of lightning strokes, overturning of the atmospheric electric field, ionospheric disturbance, plasma drift velocity reversal, and the formation and dissipation of sporadic sodium layer. The main findings of our results are summarized as follows:

1. The  $\text{Na}_s$  event discussed in the present case study shows a close relationship with  $E_s$  activities rather than conforming with the prescriptions from high temperature theory.
2. The atmospheric electric field exhibits an overturning, opposite to the fair weather downward field in circuit, in coincident with the depletion of  $E_s$  and the consequent production of  $\text{Na}_s$ .
3. A statistical analysis shows that the  $E_s$  could hardly be formed or maintained when the atmospheric electric field is directed upward.
4. A typical thunderstorm with a positive charge located above a negative charge layer, is shown to produce an upward electric field toward the ionosphere. Two regions with heavy lightning activities nearby are found during the overturning of the atmospheric electric field.
5. Observations by the ionosonde exhibit extraordinary echoes during the lightning period and the temporal property of the echoes behaved synchronously with lightning activities.

Our results support a physical connection between the lightning strokes, overturning of the electric field, ionospheric disturbances, and also possibly the  $\text{Na}_s$  phenomenon. A link between the lower and upper atmospheres could be established by the monitoring of  $\text{Na}_s$  and related phenomena as follows: lightning strokes  $\rightarrow$  overturning of electric field  $\rightarrow$  depletion of  $E_s$ /generation of  $\text{Na}_s$ .

## Appendix A: Calculations for the induced upward electric field in the global electric circuit

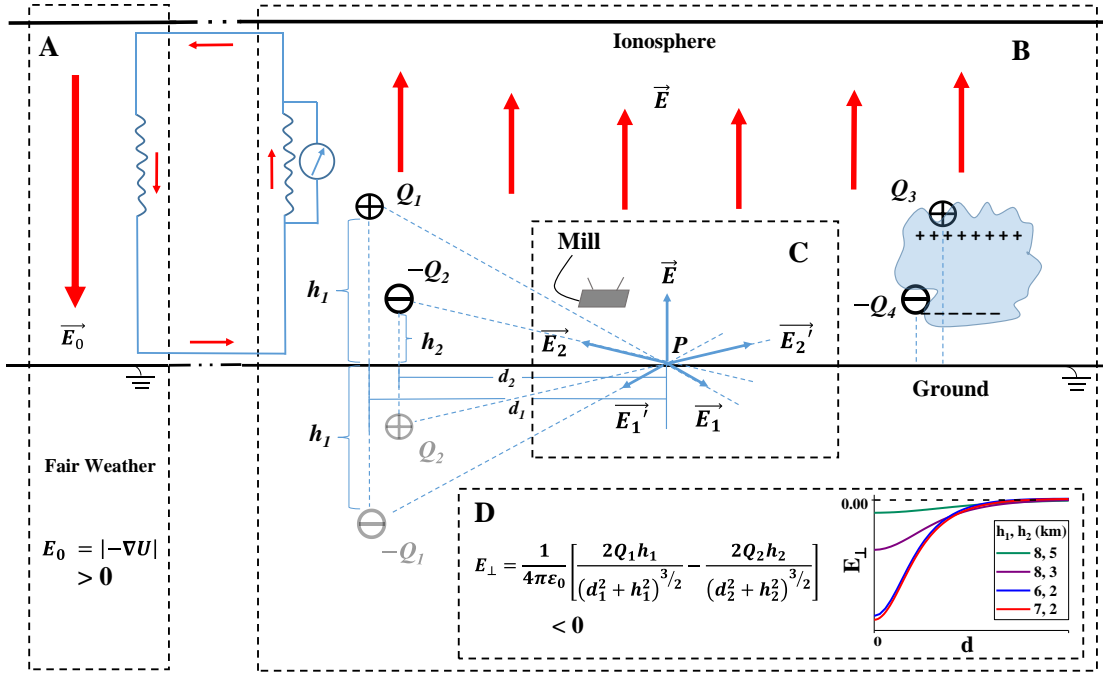


Figure A1: A diagram illustrates the global electric circuit. Part A: The atmospheric electric field under fair weather with a downward field returning from the ionosphere. Part B: The dynamo area, with thunderstorms generating an upward electric field towards the ionosphere. The electric field intensity  $E_{\perp}$  could be deduced through the electric imaging method. Part C: The deduced vertical electric field intensity at any point  $P$  within the thundercloud. Part D: The calculated  $E_{\perp}$  based on the electric imaging method.

Suppose there is positive charge  $Q_1$  at the top of a thunderstorm, with a distance of  $d_1$  above the ground; and a negative charge  $-Q_2$  at the bottom with a distance of  $d_2$ . Since the ground surface could be regarded as an infinite conducting plane, it would generate an induced charge. The boundary conditions here is:

$$U = 0 \quad \text{at } z=0$$

$$U \rightarrow 0 \quad \text{at infinity}$$

Under the uniqueness theorem, we can remove the ground surface if we put the postulated image charges of  $-Q_1$  and  $Q_2$  to the corresponding mirror points. Then for an arbitrary point  $P$  near the boundary, the vertical electric field equals to the following expression according to the Coulomb's law:

$$E_{\perp} = \frac{1}{4\pi\epsilon_0} \left[ \frac{2Q_1 h_1}{(d_1^2 + h_1^2)^{3/2}} - \frac{2Q_2 h_2}{(d_2^2 + h_2^2)^{3/2}} \right]. \quad (\text{A1}) \quad (\text{Griffiths, 1999})$$

In the simplest case, when  $Q_1$  equals to  $Q_2$  and  $d_1=d_2=d$ ,  $E_{\perp}$  varies with the distance  $d$ . If  $Q_2$  is larger than  $Q_1$  ( $Q_2 > Q_1 > 0$ ), and the negative charge  $-Q_2$  is more closed to the observing point  $P$  ( $d_2 < d_1$ ),  $E_{\perp}$  would acquire negative values (e.g., with the upward direction). A brief simulation result is shown by part D, exhibiting a persistent negative values for  $E_{\perp}$ .

## Data availability

The data sets of sodium fluorescence lidar at Hefei and three kinds of instruments at Wuhan (the ionosonde, electric mill and the fluxgate magnetometer) are publicly available from the Chinese Meridian Project database at <http://data.meridianproject.ac.cn/>. The access to the sodium density and temperature data by the USTC T/W lidar is referred to National Space Science Data Center, National Science & Technology Infrastructure of China (<http://www.nssdc.ac.cn>). The lightning location and power data can be downloaded from the World- Wide Lightning Location Network (<http://wwlln.net/>).

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### **Contributions**

Shican Qiu conceived this study and wrote this manuscript. She also prepared Fig. 2~5 in the main text and Fig. A1 in the  
5 Appendix.

Ning Wang performed data analysis.

Willie Soon was in charge of the organization and English polishing of the whole manuscript.

Gaopeng Lu added some materials about thunderstorms and lightning strokes.

Mingjiao Jia prepared Fig. 1 and gave some useful comments on the content.

10 Xianghui Xue wrote the response to reviewers and added some materials in the discussion.

Tao Li helped with the response to reviewers.

Xiankang Dou conceived this study and provided data from the Chinese Meridian Project.

### **Competing interests**

15 The authors declare no conflict of interest.

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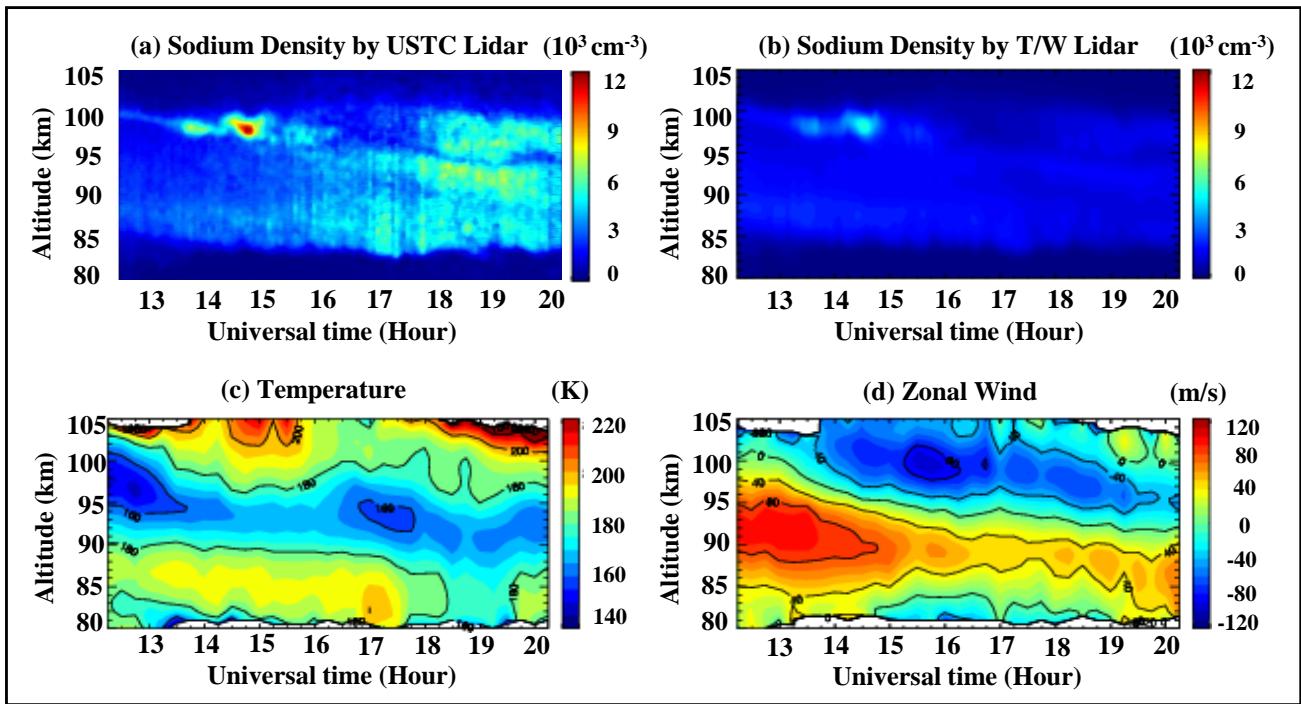


Figure 1: Observations on June 3<sup>rd</sup>, 2013, by the USTC sodium lidar and T/W lidar. (a) The sodium density detected by the USTC lidar. A moderate increase of sodium density appears at about 13:20 UT, while the largest intensity of sodium enhancement begins at about 14:20 UT. The sodium density peaks at 14:40 UT around 97.75 km. (b) The sodium density profile by T/W lidar, with a different resolution. (c) Temperature profile observed by the nearby T/W lidar, showing a cold region where the Nas occurs. (d) The zonal wind detected by the T/W lidar, exhibiting a suitable wind shear for the creation or formation of Es.

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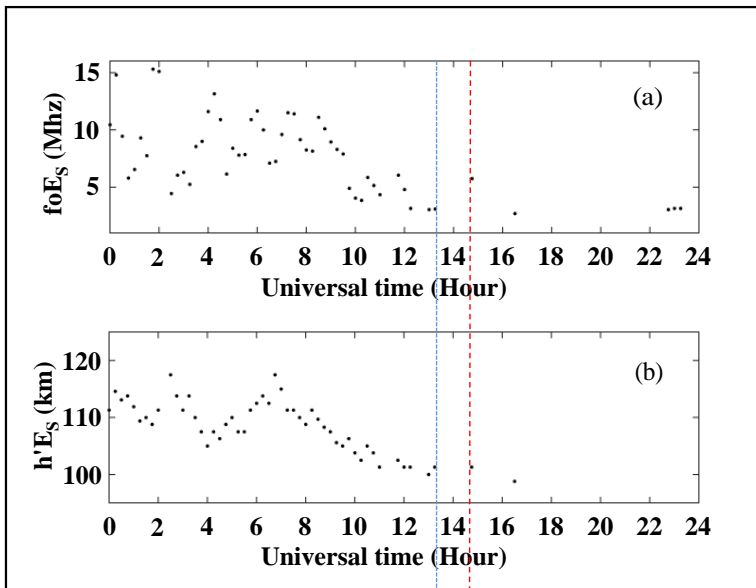


Figure 2: Sporadic E layers observed by the ionosonde at Wuhan (30.5°N, 114.6°E). (a) The time series of the critical frequency for Es (foEs). The Es layers travel/propagate downward starting around 6:30 UT, and deplete altogether at about 13:20 UT. (b) The visual height of Es (h'Es). The vertical blue dotted line annotates the beginning of the Nas around 13:20 UT, and the vertical red dashed line points out the time when the most intense sodium enhancement starts on 14:20 UT.

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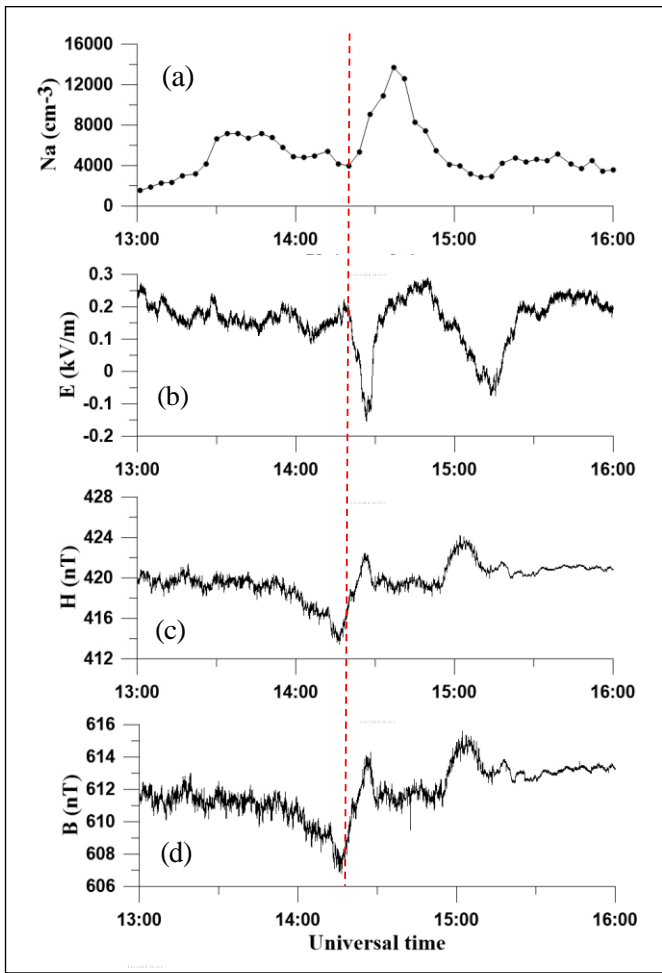


Figure 3: Observations of some atmospheric parameters and some deduced results. (a) Time series of sodium density variations at peak height 97.75 km. The sodium density begins to increase at about 14:20 UT. (b) Atmospheric electric field variations, exhibiting a synchronous overturning from 14:20 UT with the enhancement of sodium density (also pointed out by the vertical red dashed line). Note that there is another overturning peaking at 15:15 UT, without another  $N_{as}$  being produced, which could be explained by a depletion of ions in the Es. The electric field recovers at about 15:30 UT. (c) Horizontal magnetic field observed by the fluxgate magnetometer. (d) The deduced magnetic induction intensity from observations of ( $H$ ,  $D$ ,  $Z$ ) components by the fluxgate magnetometer ( $B = \sqrt{H^2 + D^2 + Z^2}$ ).

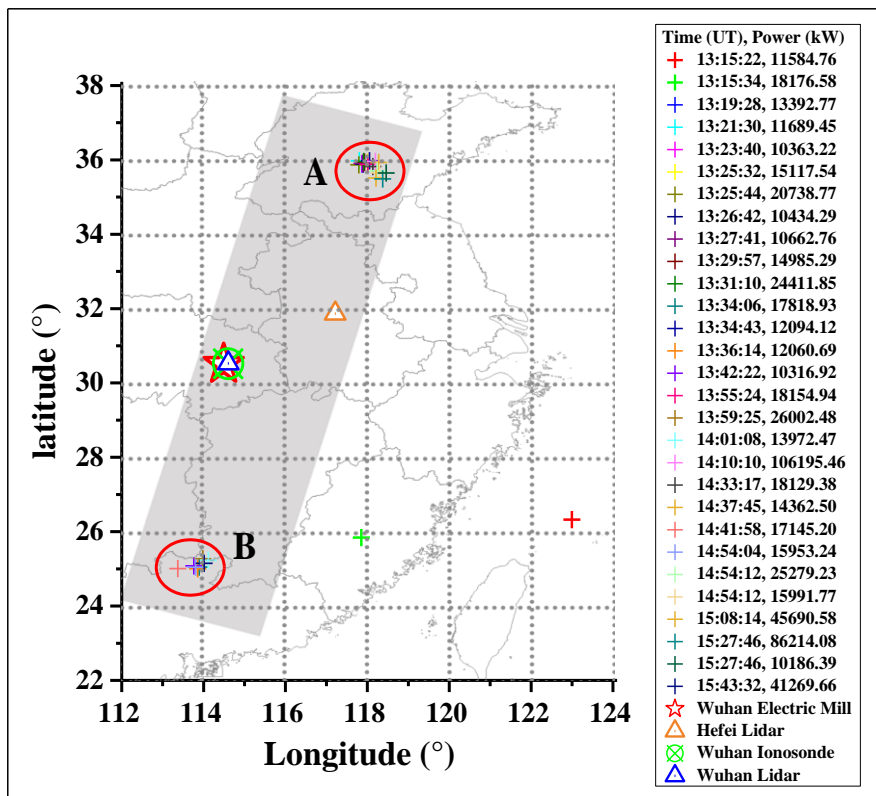


Figure 4: The lightning strokes are detected by WWLLN. The continuous strongest lightnings with a power larger than  $10^4$  kW occur from 13:19 UT to 15:43 UT, mainly concentrating around areas within (35.8°N, 118.1°E) and (25.1°N, 113.8°E), respectively.

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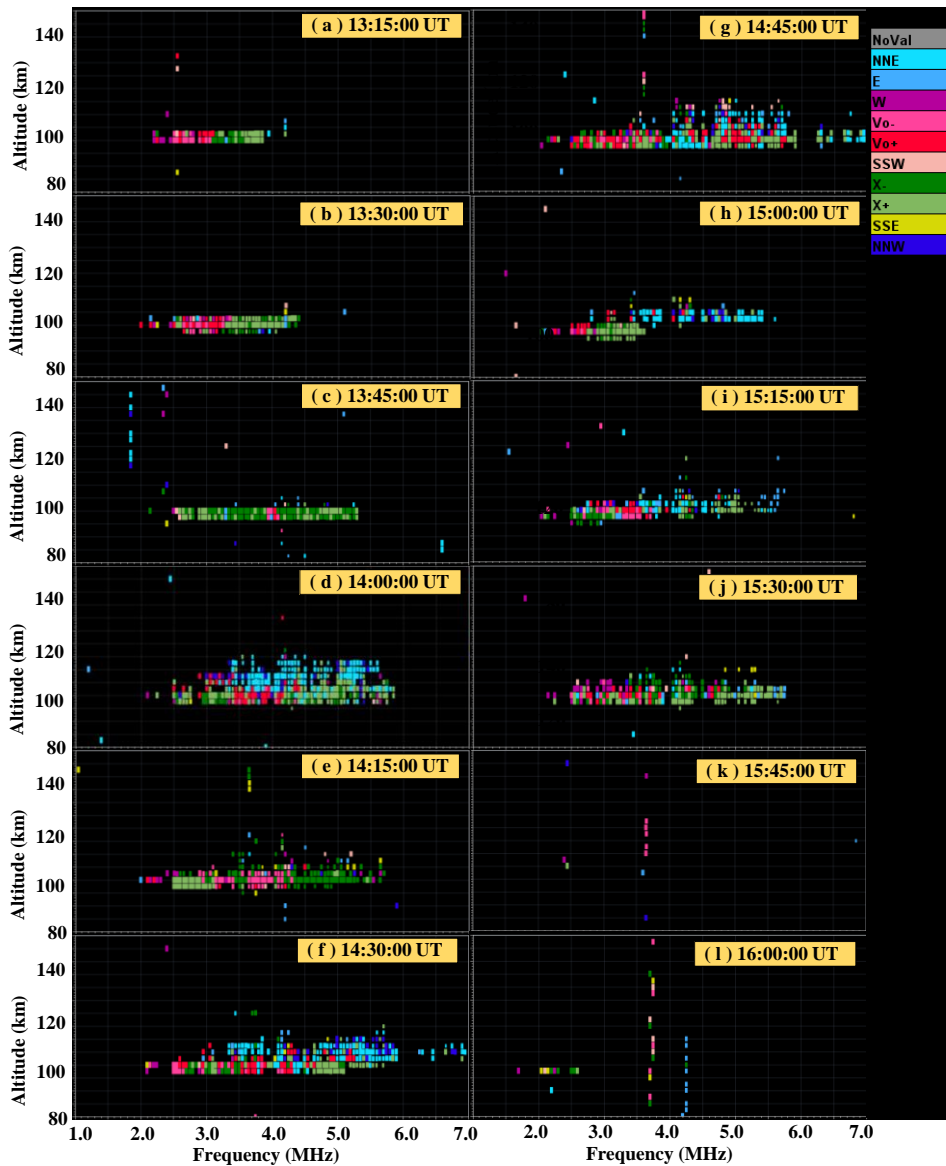


Figure 5: Extraordinary echoes by Wuhan ionosonde in different modes: (a) ~ (c): From 13:15 UT to 13:45 UT, the echoes gradually increase. Note that the powerful lightning period begins on 13:15 UT as well, with the sodium density enhancement and the Es depletion occurring on about 13:20 UT. (d) ~ (g): Most intense echo signals occur during 14:00 UT to 14:45 UT, while the largest intensity of sodium enhancement begins at 14:20 UT and the sodium density peaks at 14:40 UT. The overturning of electric field also occurs at 14:20 UT. (h) ~ (j): From 15:00 UT to 15:30 UT, the signals weaken gradually. (k) ~ (l): The echoes vanish after 15:45 UT. Afterwards, no strong stroke detected again within the discussed area. Meanwhile, the ionospheric echoes diminish after 15:45 UT, and the overturning of electric field also recovers at about 15:30 UT.



**Table 1 : The nature of foEs variations under the situations of the overturning of electric field**

<b>foEs Variations</b>	<b>Termination</b>	<b>Decrease</b>	<b>Generation</b>	<b>Increase</b>	<b>Others</b>
<b>Ratios</b> (Case days/Total days)	<b>155/242</b>	<b>39/242</b>	<b>6/242</b>	<b>26/242</b>	<b>28/242</b>