

1 **Technical note: Measurement of chemically-resolved volume**  
2 **equivalent diameter and effective density of particles by AAC-**  
3 **SPAMS**

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## 17 **Abstract**

18 Size and effective density ( $\rho_e$ ) are important properties of aerosol particles and are  
19 related to their influences on human health and the global climate. The volume  
20 equivalent diameter ( $D_{ve}$ ) is an intrinsic property that is used to evaluate particle size.  
21  $\rho_e$ , defined as the ratio of particle density to a dynamic shape factor ( $\chi$ ), is used to  
22 characterize the physical property of a particle as an alternative to particle density.  
23 However, it is still challenging to simultaneously characterize the  $D_{ve}$  and  $\rho_e$  of particles.  
24 Here, we present a novel system that classifying particles with their aerodynamic  
25 diameter ( $D_a$ ) by aerodynamic aerosol classifiers (AAC) and determining their vacuum  
26 aerodynamic diameter ( $D_{va}$ ) by single particle aerosol mass spectrometry (SPAMS) to  
27 achieve a measurement of  $D_{ve}$  and  $\rho_e$ . The reliability of the AAC-SPAMS system for  
28 accurately obtaining  $D_{ve}$  and  $\rho_e$  is verified based on the results that the deviation between  
29 the measured values and the theoretical values is less than 6% for the size-resolved  
30 spherical polystyrene latex (PSL). The AAC-SPAMS system is applied to characterize  
31 the  $D_{ve}$  and  $\rho_e$  of  $(\text{NH}_4)_2\text{SO}_4$  and  $\text{NaNO}_3$  particles, suggesting that these particles are  
32 aspherical and their  $\rho_e$  are independent of particle size. Finally, the AAC-SPAMS  
33 system is deployed in a field measurement, showing that it is a powerful technique to  
34 characterize the chemically-resolved  $D_{ve}$  and  $\rho_e$  of particles in real time.

## 35 1. Introduction

36 Size and particle density ( $\rho_p$ ) are critical parameters of aerosol particles in  
37 quantifying the impact of aerosols on air quality, human health and global climate  
38 change (Buseck and Posfai, 1999; Poschl, 2005; Pitz et al., 2003). Effective density ( $\rho_e$ )  
39 has been adopted to characterize the physical property of a particle as an alternative to  
40  $\rho_p$ , since  $\rho_p$  for aspherical aerosol particles is hardly measured (Sumlin et al., 2018;  
41 Katrib et al., 2005). Size and  $\rho_e$  govern the transport properties of a particle both in the  
42 atmosphere and in the human respiratory system (Seinfeld and Pandis, 1998; Liu and  
43 Daum, 2008) and directly/indirectly influence the potential of the particle to absorb or  
44 reflect solar radiation (Tang, 1997; Zhao et al., 2019; Liu and Daum, 2008).  $\rho_e$  can also  
45 provide information concerning particle morphology (Yon et al., 2015) and serve as a  
46 tracer for atmospheric processing (Guo et al., 2014; Yin et al., 2015; Liu et al., 2015).  
47 However, the quantitative relationship between aerosol properties, namely, size and  $\rho_e$ ,  
48 and their effects on air quality, human health and global climate change is not yet well  
49 understood, which is partly because important aerosol properties cannot be measured  
50 by current techniques.

51 **Size.** Size is a fundamental property of particles, which can be parameterized by the  
52 physical quantity of volume equivalent diameter ( $D_{ve}$ ). Defined as the diameter of a  
53 spherical particle with the same volume as the particle (DeCarlo et al., 2004),  $D_{ve}$  is an  
54 intrinsic physical quantity that can be used to evaluate the actual size of the particle.  
55 However, to date, atmospheric science usually describes particle size by other diameter  
56 definitions, such as the electric mobility diameter ( $D_m$ ), aerodynamic equivalent

57 diameter ( $D_a$ ) and vacuum aerodynamic equivalent diameter ( $D_{va}$ ), whose relationships  
 58 with  $D_{ve}$  are shown in Eqs. (1)-(3), respectively:

$$59 \quad \frac{D_m}{C_c(D_m)} = \frac{D_{ve}}{C_c(D_{ve})} \chi_t, \quad (1)$$

$$60 \quad D_a = D_{ve} \sqrt{\frac{\rho_p C_c(D_{ve})}{\chi_t \rho_0 C_c(D_a)}}, \quad (2)$$

$$61 \quad D_{va} = \frac{\rho_p D_{ve}}{\rho_0 \chi_v}, \quad (3)$$

62 where  $C_c(D)$  is the Cunningham slip correction factor,  $\chi_t$  and  $\chi_v$  represent the aerosol  
 63 dynamic shape factor ( $\chi$ ) in the transition regime and in the free-molecule regime,  
 64 respectively, and  $\rho_0$  represents the unit density of 1.0 g/cm<sup>3</sup>. From the definitions, it can  
 65 be seen that  $D_m$ ,  $D_a$ , and  $D_{va}$  are originally derived from  $D_{ve}$ , but in actuality, they do  
 66 not reflect the actual size of the particle. Meanwhile,  $D_{ve}$  cannot be easily obtained,  
 67 which limits its application in the scientific community.

68 **Effective density.** At present, three definitions of  $\rho_e$  are introduced in atmospheric  
 69 science (DeCarlo et al., 2004): the first definition ( $\rho_e^I$ ) is the ratio of the measured  
 70 particle mass ( $m_p$ ) to the particle volume ( $V$ ) calculated assuming a spherical particle  
 71 with a diameter equal to the measured  $D_m$ ; the second definition ( $\rho_e^{II}$ ) is the ratio of  $\rho$   
 72 to  $\chi$  (Hand and Kreidenweis, 2002); and the third definition ( $\rho_e^{III}$ ) is the ratio of  $D_m$  and  
 73  $D_{va}$ , all of which are expressed in Eqs. (4)-(6), respectively.

$$74 \quad \rho_e^I = \frac{6m_p}{\pi D_m^3} \quad (4)$$

$$75 \quad \rho_e^{II} = \frac{\rho_p}{\chi} \quad (5)$$

$$76 \quad \rho_e^{III} = \frac{D_{va}}{D_m} \rho_0 \quad (6)$$

77 The definitions of  $\rho_e^I$  and  $\rho_e^{III}$  can be derived into the final forms, as shown in the Eqs.(7)  
 78 and (8), respectively.

79 
$$\rho_e^I = \frac{\rho}{\chi t^3} \cdot \left( \frac{C_c(D_{ve})}{C_c(D_m)} \right)^3 \quad (7)$$

80 
$$\rho_e^{III} = \rho \cdot \frac{C_c(D_{ve})}{\chi^2 \cdot C_c(D_m)} \quad (8)$$

81 The Eq. (7) is derived from combining the Eq. (1) with Eq. (4) in which  $m_p$  is equal to  
 82  $1/6 \rho \cdot D_{ve}^3$ . The detailed derivation of Eq. (8) was presented in Schneider et al.  
 83 (2006). The  $\rho_e^I$  and  $\rho_e^{III}$  are demonstrated to have the inherent characteristics of  
 84 decreasing with increasing particle size in a separate paper. Therefore, it will introduce  
 85 systemic error when assessing the particle impacts on visibility, human health and  
 86 climate from the physical quantities in  $\rho_e^I$  and  $\rho_e^{III}$ . In contrast,  $\rho_e^{II}$  is independent of  
 87 particle size. Previously,  $\rho_e^{II}$  and the real part in the refractive index ( $n$ ) can be retrieved  
 88 from a fitting procedure that compares the measured light-scattering intensity of  
 89 particles ( $R_{meas}$ ) to the theoretical values ( $R_{theory,test}$ ) calculated by a series of  $n$  and  $\rho_e^{II}$   
 90 values. Moffet and Prather (2005) successfully obtained  $\rho_e^{II}$  for spherical particles by  
 91 single particle mass spectrometry. However, subject to the accuracy of Mie theory for  
 92 the aspherical particles, dry NaCl and calcium-rich dust particles were failed to fit the  
 93  $R_{theory,test}$  well to  $R_{meas}$  (Moffet et al., 2008). Similarly, Zhang et al. (2016a) failed to  
 94 simultaneously retrieve  $\rho_e^{II}$  and  $n$  for  $(NH_4)_2SO_4$  and  $NaNO_3$  particles. To our best  
 95 knowledge, there is no effective technique to achieve the measurement of  $\rho_e^{II}$  for  
 96 aspherical particles. For reference, the symbol  $\rho_e$  in the following text refers to the  
 97 definition of  $\rho_e^{II}$ .

98 The aim of the present work is to develop a method to obtain  $D_{ve}$  and  $\rho_e$ . The  
 99 established system of an aerodynamic aerosol classifier (AAC)-single particle aerosol  
 100 mass spectrometry (SPAMS) is capable of characterizing the  $D_a$  and  $D_{va}$  of particles,

101 which can be applied to theoretically derive  $D_{ve}$  and  $\rho_e$ . To verify the reliability of the  
102 AAC-SPAMS system, we apply it to measure the  $D_{ve}$  and  $\rho_e$  of the spherical particles  
103 of polystyrene latex (PSL). The results are in good agreement with the theoretical values.  
104 Finally, the AAC-SPAMS system is applied to measure the  $D_{ve}$  and  $\rho_e$  for  $(\text{NH}_4)_2\text{SO}_4$   
105 and  $\text{NaNO}_3$  particles and for the chemically-resolved atmospheric particles.

106

## 107 **2. Experimental section**

### 108 **2.1 Measurement system**

109 Figure 1 shows a schematic diagram of the AAC-SPAMS system. The particles are  
110 first dried by a diffusion drying tube (TSI 9302, USA), classified by AAC (Cambustion  
111 Ltd., UK) based on the aerodynamic diameters  $D_a$ , and then transported into SPAMS in  
112 which the  $D_{va}$  and the mass spectra of individual particles are obtained. The working  
113 principle of the AAC is described in detail elsewhere (Tavakoli and Olfert, 2013). AAC  
114 consists of two coaxial cylinders that rotate at the same rotational speed. Polydisperse  
115 particles enter into the space between the cylinders (i.e., classification column) and  
116 experience a centrifugal force that causes them to move toward the outer cylinder. The  
117 particles to be classified can leave the classification column with the particle-free sheath  
118 flow and finally exit the AAC with the sample flow. Thus, the  $D_a$  values of classified  
119 particles can be derived from their relationship with their relaxation time ( $\tau$ ), as shown  
120 in Eq. (9):

$$121 \quad \tau = \frac{c_C(D_a) \cdot \rho_0 \cdot D_a^2}{18\mu} \quad (9)$$

122 where  $\mu$  is the gas dynamic viscosity. Particles with large relaxation times impact and

123 adhere to the outer cylinder, while particles with small relaxation times exit the  
124 classifier with the exhaust flow. The exhaust flow from the AAC was about 0.3 lpm,  
125 and the Size Resolution Parameter ( $R_s$ ) of the AAC was set as 40.

126 Detailed information about the operation of SPAMS (Hexin Analytical Instrument  
127 Co., Ltd., China) is described elsewhere (Li et al., 2011). Briefly, the particles are  
128 introduced into the vacuum system through a 0.1 mm critical orifice and are gradually  
129 collimated into a beam in the aerodynamic lens. Two continuous diode Nd:YAG laser  
130 beams (532 nm) are used to aerodynamically size the particles, which are subsequently  
131 desorbed/ionized by a pulsed laser (266 nm) that is triggered based on the velocity of a  
132 specific particle. The generated positive and negative ions are recorded with the  
133 corresponding particle size. The  $D_{va}$  of the particle is related to the transit time between  
134 the two laser beams (532 nm) in SPAMS, which can be obtained by using a calibration  
135 curve generated from the measured transit times of a PSL series with predefined sizes  
136 (nominal diameters).

137

## 138 **2.2 Laboratory experiments**

139 Dried spherical PSL (Nanosphere Size Standards, Duke Scientific Corp., Palo Alto)  
140 ( $\rho_p = 1.055 \text{ g/cm}^3$  and  $\chi = 1.0$ ) with  $D_{ve}$  values of  $203.0 \pm 5.0 \text{ nm}$ ,  $310.0 \pm 6.0 \text{ nm}$ ,  $510.0$   
141  $\pm 5.0 \text{ nm}$ , and  $740.0 \pm 6.0 \text{ nm}$  were used in the AAC-SPAMS system, and the  $D_{ve}$  was  
142 verified by Scanning Mobility Particles Sizer (Model 3938, TSI Inc., USA). The PSL  
143 particles were first classified by AAC, and then their  $D_{va}$  values were obtained by  
144 SPAMS. ACC-SPAMS was also applied to the particles of  $(\text{NH}_4)_2\text{SO}_4$  ( $\rho_p = 1.77 \text{ g/cm}^3$ )

145 and  $\text{NaNO}_3$  ( $\rho_p = 2.26 \text{ g/cm}^3$ ) with  $D_a$  values of 250.0 nm, 350.0 nm, 450.0 nm and  
146 550.0 nm.

### 147 **2.3 Ambient sampling**

148 For field observations, the AAC-SPAMS system was placed in science and  
149 technology enterprise accelerator A2 Block, Guangzhou, China, to characterize the  $D_{ve}$ ,  
150  $\rho_e$  and chemical compositions of aerosol particles. The sampling inlet was hung 2.5  
151 meters from the third floor (~12 m above ground level). Ambient aerosol particles were  
152 introduced into the AAC through a 5 m long conductive silicone tube with an inner  
153 diameter of 6 mm and a  $\text{PM}_{2.5}$  cyclone inlet. The overall sampling flow was 3 lpm, and  
154 the residence time was approximately 5 seconds. Sampled particles were classified by  
155 the AAC as one of four  $D_a$ : 250.0 nm, 350.0 nm, 450.0 nm and 550.0 nm. The sampling  
156 time for the particles of each  $D_a$  was approximately 10 minutes. From July 6<sup>th</sup> to 8<sup>th</sup>,  
157 2019, approximately 129,869 ionized particles were obtained from nine rounds of  
158 measurement. The sampling details are shown in Table S1. The number of ionized  
159 particles with the  $D_a$  of 250.0, 350.0, 450.0, and 550.0 nm is about 35,609, 38,374,  
160 31,910, and 23,976, respectively. The sampled ~100,000 particles are first classified by  
161 using an adaptive resonance theory neural network (ART-2a) (Song et al., 1999) with a  
162 vigilance factor of 0.75, a learning rate of 0.05 and 20 iterations.

163

### 164 **2.4 Theoretical derivation of $D_{ve}$ and $\rho_e$ from $D_a$ and $D_{va}$**

165  $D_{ve}$  is the accurate physical quantity of the size of a particle.  $\rho_e$  is an alternative  
166 property for  $\rho_p$ , which is consistent with the property of  $\rho_p$  in terms of being independent



167 of particle size. These two properties cannot yet be measured for unknown particles by  
 168 current techniques. In this study, the calculations of  $D_{ve}$  and  $\rho_e$  for unknown particles  
 169 are theoretically derived from  $D_a$  and  $D_{va}$ . Combining Eqs. (2) and (3), we obtain the  
 170 following Eq. (10):

$$171 \quad C_c(D_a) \frac{D_a^2}{D_{va}} = D_{ve} C_c(D_{ve}) \frac{\chi_v}{\chi_t} \quad (10)$$

172 Based on the approximation between  $\chi_v$  and  $\chi_t$  ( $\chi_v \approx \chi_t = \chi_a$ ) (DeCarlo et al., 2004), Eq.  
 173 (10) becomes Eq. (11):

$$174 \quad C_c(D_a) \frac{D_a^2}{D_{va}} = D_{ve} C_c(D_{ve}) \quad (11)$$

175 The Cunningham Slip Correction Factor is calculated by Eq. (12):

$$176 \quad C_c(D) = 1 + \frac{\lambda}{D} \left( A + B \cdot \exp\left(\frac{C \cdot D}{\lambda}\right) \right), \quad (12)$$

177 where  $\lambda$  is the mean free path of the gas molecules, and  $A$ ,  $B$  and  $C$  are empirically  
 178 determined constants specific to the analysis system. Substituting Eq. (12) into Eq. (11)  
 179 obtains the Eq. (13).

$$180 \quad \frac{D_a^2}{D_{va}} + \frac{D_a \cdot \lambda}{D_{va}} \left( A + B \cdot \exp\left(\frac{C \cdot D_a}{\lambda}\right) \right) = D_{ve} + \lambda \left( A + B \cdot \exp\left(\frac{C \cdot D_{ve}}{\lambda}\right) \right) \quad (13)$$

181 Thus, if the  $D_a$  and  $D_{va}$  of an unknown particle can be measured, its  $D_{ve}$  will be  
 182 calculated according to Eq. (13). Finally, the  $\rho_e$  value of the particles is calculated by  
 183 the  $D_{va}$  and  $D_{ve}$  values according to Eq. (14), which is obtained by combining Eq.(3)  
 184 and Eq.(5):

$$185 \quad \rho_e = \frac{\rho_p}{\chi_a} = \frac{D_{va}}{\rho_0 \cdot D_{ve}} \quad (14)$$

186 Thus, we can obtain both the  $D_{ve}$  and  $\rho_e$  values of unknown particles based on the  $D_a$   
 187 and  $D_{va}$  values. Because the AAC and SPAMS instruments have the ability to determine  
 188  $D_a$  and  $D_{va}$ , the AAC-SPAMS system, which is developed in this study, can be used to

189 obtain the  $D_{ve}$  and  $\rho_e$  values for unknown particles.

190

### 191 3. Results and discussion

#### 192 3.1 Verification of the AAC-SPAMS system to obtain $D_{ve}$ and $\rho_e$

193 The  $D_{va}$  distribution of PSL particles with predefined  $D_{ve}$  values after screening by  
194 the AAC is shown in Figure S1. We used Gaussian fitting to obtain the peak  $D_{va}$  for  
195 each size PSL with an R-squared fitting coefficient ( $R^2$ ) over 0.98. Each fitting has a  
196 full width at half maximum (FWHM) of 6.6%, 4.4%, 2.3% and 2.2%, and the  
197 corresponding peaks are 215.8 nm, 319.0 nm, 532.1 nm and 803.5 nm, respectively.

198 Substituting the  $D_a$  and  $D_{va}$  values of PSL into Eq. (11), the measured  $D_{ve}$  ( $D_{ve,me}$ ) of  
199 PSL from AAC-SPAMS system is 203.6 nm, 309.7 nm, 511.6 nm and 737.2 nm,  
200 respectively (Figure 2a). Thus, the deviations between the theoretical  $D_{ve}$  ( $D_{ve,th}$ ) and  
201  $D_{ve,me}$  values are 0.3%, -0.1%, 0.3% and -0.4%, respectively. On the other hand, the  
202 measured  $\rho_e$  ( $\rho_{e,me}$ ) values of the particles are calculated from the  $D_{va}$  and  $D_{ve,me}$  values  
203 with Eq. (14), and the  $\rho_{e,me}$  values are 1.1 g/cm<sup>3</sup>, 1.0 g/cm<sup>3</sup>, 1.0 g/cm<sup>3</sup>, and 1.1 g/cm<sup>3</sup>  
204 (Figure 2b). The deviations of  $\rho_{e,me}$  are determined to be 4.3%, -5.2%, -5.2%, and 4.3%,  
205 respectively, by comparing to the theoretical  $\rho_e$  ( $\rho_{e,th}$ ) that is equals to the  $\rho_p$  for the  
206 spherical particles (i.e. 1.055 g/cm<sup>3</sup> of PSL particles). That is, the deviations of  $D_{ve,me}$   
207 and  $\rho_{e,me}$  characterized by the AAC-SPAMS system are within 1% and 6%, respectively.  
208 We therefore conclude that the AAC-SPAMS system is highly accurate for obtaining  
209 aerosol  $D_{ve}$  and  $\rho_e$ .

210

211 **3.2 Application of the AAC-SPAMS system for obtaining  $D_{ve}$  and  $\rho_e$  of  $(\text{NH}_4)_2\text{SO}_4$**   
212 **and  $\text{NaNO}_3$**

213 Figure S2 shows the  $D_{va}$  distributions of  $(\text{NH}_4)_2\text{SO}_4$  and  $\text{NaNO}_3$  particles, which have  
214  $D_a$  values of 250.0, 350.0, 450.0, and 550.0 nm, as screened by the AAC. The  $D_{va}$  peaks  
215 are obtained by Gaussian fitting, with  $R^2$  values over 0.93 and FWHM values ranging  
216 from 7.6% to 10.6%. The  $(\text{NH}_4)_2\text{SO}_4$  particles have  $D_{va}$  values of 300.0, 418.0, 551.1,  
217 and 695.1 nm (Figure S2), which correspond to particles possessing  $D_{ve,me}$  values of  
218 177.3, 254.4, 331.8, and 409.3 nm, respectively, according to Eq. (11). Substituting the  
219 values of  $D_{va}$  and  $D_{ve,me}$  into Eq. (12), the  $\rho_{e,me}$  values are 1.7, 1.6, 1.6, and 1.7  $\text{g}/\text{cm}^3$   
220 (Figure 3a), respectively. Similarly, the selected  $\text{NaNO}_3$  particles are determined to  
221 have  $D_{va}$  values of 321.0, 454.9, 599.8, and 755.3 nm (Figure S2), corresponding to  
222  $D_{ve,me}$  values of 150.1, 218.2, 287.0, and 355.9 nm, respectively. The  $\rho_{e,me}$  values of the  
223  $\text{NaNO}_3$  particles are 2.2, 2.0, 2.0, and 2.1  $\text{g}/\text{cm}^3$  for the four particle sizes (Figure 3b),  
224 respectively. Figure 3 also shows that the  $\rho_{e,me}$  values of the  $\text{NaNO}_3$  and  $(\text{NH}_4)_2\text{SO}_4$   
225 particles at four size deviate from their average values with the maximum of 5.9 % and  
226 4.8%, respectively, which are identical with the deviation phenomenon for the  $\rho_{e,me}$  of  
227 PSL particles. These deviations may be derived from the calibration of particle  $D_{va}$   
228 from the SPAMS. While the R-square of size calibration curve is 0.999, the curve of  
229 exponential function is found to slightly deviate from the data points measured by  
230 SPAMS. For example, size calibration function has the deviation of -4.4% and 3.1%  
231 from the data points of 310 and 740 nm, respectively.

232 Taking the systematic error into account, the slight difference of the  $\rho_{e,me}$  values for

233 the four sizes suggests that the  $\rho_e$  of  $(\text{NH}_4)_2\text{SO}_4$  and  $\text{NaNO}_3$  particles is independent of  
234 particle size from 250.0 nm to 550.0 nm. It is determined by the definition of effective  
235 density used in this study, which keeps constant as long as the  $\chi_a$  of the particles does  
236 not change with particle size for pure compound. The average  $\rho_{e,me}$  values of  $(\text{NH}_4)_2\text{SO}_4$   
237 and  $\text{NaNO}_3$  particles are calculated to be  $1.7 \pm 0.1$  and  $2.1 \pm 0.1$   $\text{g}/\text{cm}^3$ , respectively.  
238 The average  $\rho_{e,me}$  values are lower than that the  $\rho_p$  of  $(\text{NH}_4)_2\text{SO}_4$  ( $1.77$   $\text{g}/\text{cm}^3$ ) and  
239  $\text{NaNO}_3$  ( $2.27$   $\text{g}/\text{cm}^3$ ), which is caused that the  $\rho_{e,me}$  is determined by both of  $\rho_p$  and  $\chi_a$ .  
240 According to Eq. (14), the  $\chi_a$  of  $(\text{NH}_4)_2\text{SO}_4$  and  $\text{NaNO}_3$  particles with different  $D_a$  are  
241 calculated to be 1.04, 1.11, 1.11, and 1.04 and to be 1.03, 1.14, 1.14, and 1.08,  
242 respectively. Thus, the average  $\chi_a$  values of the  $(\text{NH}_4)_2\text{SO}_4$  and  $\text{NaNO}_3$  particles are  
243 determined to be  $1.07 \pm 0.04$  and  $1.10 \pm 0.05$ , respectively, which can be used to  
244 parameterize their morphology.

245 The average  $\chi_a$  values of the  $(\text{NH}_4)_2\text{SO}_4$  and  $\text{NaNO}_3$  particles indicate that these  
246 particles are aspherical. The asphericity of  $(\text{NH}_4)_2\text{SO}_4$  determined by AAC-SPAMS  
247 system is consistent with the previous studies reporting that the  $\chi_a$  of  $(\text{NH}_4)_2\text{SO}_4$  were  
248 larger than the value of 1.03 (Zelenyuk et al., 2006; Beranek et al., 2012; Zhang et al.,  
249 2016a). However, previous studies found that the  $\text{NaNO}_3$  particles had different  
250 morphology. Zhang et al. (2016a) observed that  $\text{NaNO}_3$  had the  $\chi_a$  of 1.09-1.13,  
251 indicating its asphericity, while Hoffman et al. (2004) found that  $\text{NaNO}_3$  particle had a  
252 round droplet-like shape even at 15% RH, which was supported by the consistence  
253 between the measured value of “anhydrous” droplet density and the calculated value of  
254 “anhydrous” solution droplet (Zelenyuk et al., 2005). Eclectically, Tang and

255 Munkelwitz (1994) studied that most of the  $\text{NaNO}_3$  particles crystallized between 20%  
256 and 30% RH but some persisted down to 10% RH to form solution droplets. Notably,  
257 the spherical  $\text{NaNO}_3$  particles at low RH observed by Hoffman et al. (2004) were dried  
258 in the sticky carbon tape which might affect the phase transition of droplet-like  $\text{NaNO}_3$   
259 particles. In this study, most  $\text{NaNO}_3$  particles were crystallized because the RH of the  
260 aerosol flow carrying the  $\text{NaNO}_3$  particles was reduced to below 20% through the  
261 diffusion drying tube. Besides, the result that the crystallized  $\text{NaNO}_3$  particles are  
262 aspherical is supported by their FWHM values of the  $D_{va}$  distributions which are  
263 consistent with that of aspherical  $(\text{NH}_4)_2\text{SO}_4$  but wider than spherical PSL (Figures S1  
264 and S2).

265

### 266 **3.3 Application of the AAC-SPAMS system for measuring the chemically-resolved** 267 **$D_{ve}$ and $\rho_e$**

268 SPAMS can obtain information on the chemical composition of individual particles,  
269 implying that the AAC-SPAMS system has the ability to simultaneously characterize  
270  $D_{ve}$ ,  $\rho_e$  and the chemical compositions of particles in real time. It is worth noting that  
271 the particles with the largest  $\chi$  in the actual atmosphere should be freshly emitted soot,  
272 which  $\chi$  is 2.5 (Peng et al., 2016). This largest  $\chi$  fitly meets the upper limit for the  
273 approximation between the  $\chi_t$  and  $\chi_v$  (DeCarlo et al., 2004). Therefore, the AAC-SPAMS  
274 system can obtain the chemically-resolved  $D_{ve}$  and  $\rho_e$  values for unknown aerosol  
275 particles in the field observation.

276 As an example, we applied the AAC-SPAMS system to illustrate how the measured

277  $D_a$ ,  $D_{va}$ , and chemical composition of an individual particle can be used to calculate the  
278  $D_{ve}$  and  $\rho_e$  for unknown particles. The sampled  $\sim 100,000$  particles are classified into  
279 eight major particle types with distinct chemical patterns of K-rich, EC-S, K-Na, Amine,  
280 EC-N-S, OC-N-S and OC-EC-N-S and Metal-rich, representing 97% of the detected  
281 particle population. [Details of the chemical composition and number fraction of the](#)  
282 [eight types of particles are presented in the Figure S3 and Figure S4, respectively, which](#)  
283 [are discussed in the Supporting Information.](#)

284 We used Gaussian fitting to obtain the  $D_{va}$  peaks for each particle type with  $D_a$  values  
285 of 250.0 nm, 350.0 nm, 450.0 nm, and 550.0 nm. Then, we calculated the  $D_{ve}$  values of  
286 the atmospheric particles with Eq. (11). Table 1 presents the average  $D_{ve}$  values of the  
287 eight particle types, for which the standard deviation is calculated based on nine  
288 samples. The average  $D_{ve}$  at  $D_a$  values of 250.0 nm, 350.0 nm, 450.0 nm, and 550.0 nm  
289 has the following wide ranges: from 188.5 nm to 200.8 nm, 271.9 nm to 295.7 nm,  
290 342.5 nm to 428.9 nm, and 397.3 nm to 570.9 nm, respectively, which are caused by  
291 the chemical composition differences. The result indicates that particles with  
292 significantly different  $D_{ve}$  might possess the same  $D_a$ . Furthermore, the large standard  
293 deviation of  $D_{ve}$ , such as 21.9 nm for K-Na at 250.0 nm, 32.3 nm for OC-EC-N-S at  
294 350.0 nm, and 44.3 nm for OC-N-S at 450.0 nm, indicates that the  $D_{ve}$  of particles is  
295 remarkably different even for particles with the same type and same  $D_a$ .

296 According to  $D_{ve}$  and  $D_{va}$ , we calculated the  $\rho_e$  of each particle type with Eq. (12).  
297 Figure 4 shows the variations of the  $\rho_e$  with  $D_{ve}$  for nine particle samples. For pure  
298 compounds, such as  $(\text{NH}_4)_2\text{SO}_4$  and  $\text{NaNO}_3$  particle,  $\rho_e$  theoretically does not change

299 with particle size. However, the sampled particles have experienced complex  
300 atmospheric processes. Therefore,  $\rho_e$  has a very wide distribution for each type of  
301 particle with a similar  $D_{ve}$ . Specifically, the  $\rho_e$  of K-Na increases with  $D_{ve}$ , while the  $\rho_e$   
302 of OC-N-S and OC-EC-N-S decreases with  $D_{ve}$ , which may be influenced by the  
303 particle shape or the material density. Additionally, the average  $\rho_e$  of each type of  
304 particle is in the order from small to large:  $1.2 \pm 0.2 \text{ g/cm}^3$  for OC-EC-N-S,  $1.3 \pm 0.2$   
305  $\text{g/cm}^3$  for OC-N-S,  $1.4 \pm 0.1 \text{ g/cm}^3$  for K-rich,  $1.4 \text{ g/cm}^3$  for Amine,  $1.5 \text{ g/cm}^3$  for EC-  
306 N-S,  $1.5 \text{ g/cm}^3$  for EC-S,  $1.6 \pm 0.1 \text{ g/cm}^3$  for K-Na and  $1.6 \pm 0.1 \text{ g/cm}^3$  for Metal-rich.  
307 It is reasonable to find that the average  $\rho_e$  of internally mixed particles distributes in the  
308 range of their material densities ( $\rho_m$ ). For instance, mainly comprised of internally  
309 mixed sulfate and organics, the OC-EC-N-S, OC-N-S, K-rich, and Amine particles have  
310 the average  $\rho_e$  between that of sulfate with  $\rho_m$  of  $1.77 \text{ g/cm}^3$  and organic aerosols with  
311  $\rho_m$  of  $1.2 \text{ g/cm}^3$  (Cross et al., 2007).

312

#### 313 **4. Conclusion**

314 We first develop an AAC-SPAMS system to achieve the measurement of the  $D_{ve}$  and  
315  $\rho_e$  of the particles through characterizing their  $D_a$  and  $D_{va}$ . The reliability of the AAC-  
316 SPAMS system is verified by accurately measuring the  $D_{ve}$  and  $\rho_e$  of PSL. Applying the  
317 AAC-SPAMS system to determine the  $D_{ve}$  and  $\rho_e$  of  $(\text{NH}_4)_2\text{SO}_4$  and  $\text{NaNO}_3$  particles  
318 shows that these particles are aspherical and their  $\rho_e$  are independent of particle size.  
319 Coupled with the ability of SPAMS to characterize the chemical composition of  
320 individual particles, we conducted a sample proof of the AAC-SPAMS equipment in

321 Guangzhou to first characterize the  $D_{ve}$ ,  $\rho_e$  and chemical compositions of atmospheric  
322 particles, showing the potential application of this system in field observations. The  
323 approach achieves the measurement of chemically-resolved  $D_{ve}$  and  $\rho_e$ , and provides  
324 the possibility to determine their quantitative relationship with other particle properties,  
325 which would be benefit for further reduction of the uncertainty associated with the  
326 effects of particles on air quality, human health and radiative forcing.

327

328 **Data availability.** Data in this study is available at [https://github.com/longer1217/All-](https://github.com/longer1217/All-figures-data)  
329 [figures-data](https://github.com/longer1217/All-figures-data).

330

331 **Author contributions.** The idea for the study was conceived by LP and GHZ. All  
332 experiments were performed by LP with the assistance of LL. LP wrote the paper which  
333 was reviewed by GHZ and XHB. All co-authors discussed the results and commented  
334 on the manuscript.

335

336 **Competing interests.** The authors declare they have no conflict of interest.

337

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345

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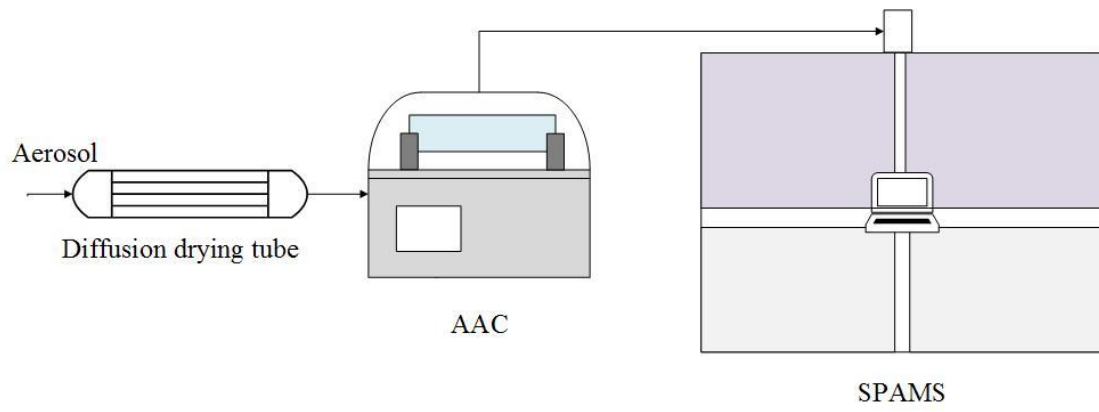
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445

446 **Table 1.**  $D_{ve}$  and its standard deviation for the eight particle types at  $D_a$  values of 250.0 nm, 350.0  
 447 nm, 450.0 nm, and 550.0 nm from nine round measurement.  
 448

$D_a$ (nm)	K-rich	EC-S	K-Na	Amine
<b>250.0</b>	193.1 ± 8.2	192.2 ± 8.1	193.8 ± 21.9	190.6 ± 4.6
<b>350.0</b>	284.0 ± 28.4	280.8 ± 9.3	271.9 ± 18.0	284.8 ± 18.2
<b>450.0</b>	364.7 ± 21.1	357.8 ± 6.9	342.5 ± 7.3	367.9 ± 9.7
<b>550.0</b>	416.6 ± 28.3	439.5 ± 5.4	397.3 ± 29.7	442.5 ± 7.4
$D_a$ (nm)	EC-N-S	OC-N-S	OC-EC-N-S	Metal-rich
<b>250.0</b>	188.5 ± 5.9	200.8 ± 17.9	195.4 ± 8.9	189.0 ± 6.7
<b>350.0</b>	281.3 ± 9.3	295.7 ± 29.8	294.0 ± 32.3	277.0 ± 9.1
<b>450.0</b>	358.0 ± 5.8	398.3 ± 44.3	428.9 ± 24.0	342.9 ± 10.0
<b>550.0</b>	453.2 ± 16.4	547.4 ± 14.7	570.9	407.4 ± 14.5

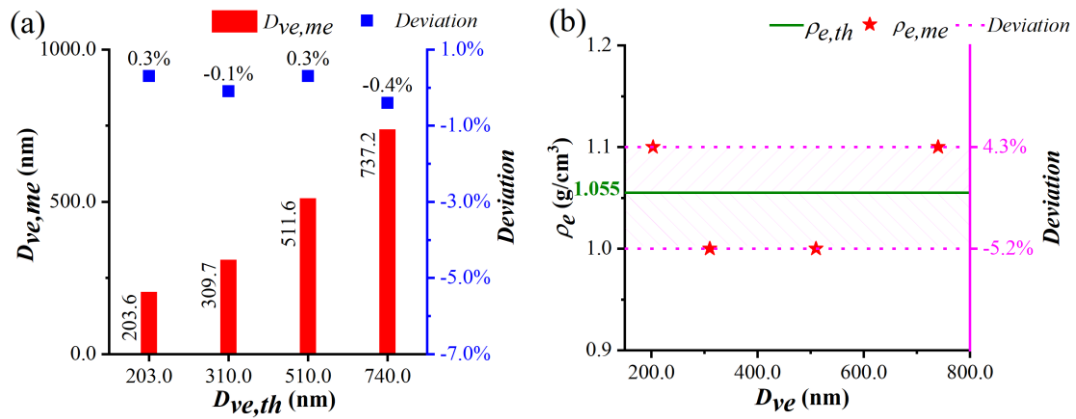
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450

451 **Figure 1.** Schematic diagram of the AAC-SPAMS system (0.3 lpm). The diffusion drying tube is

452 filled with orange silica gel, which reduces the RH to 5-15%.



453

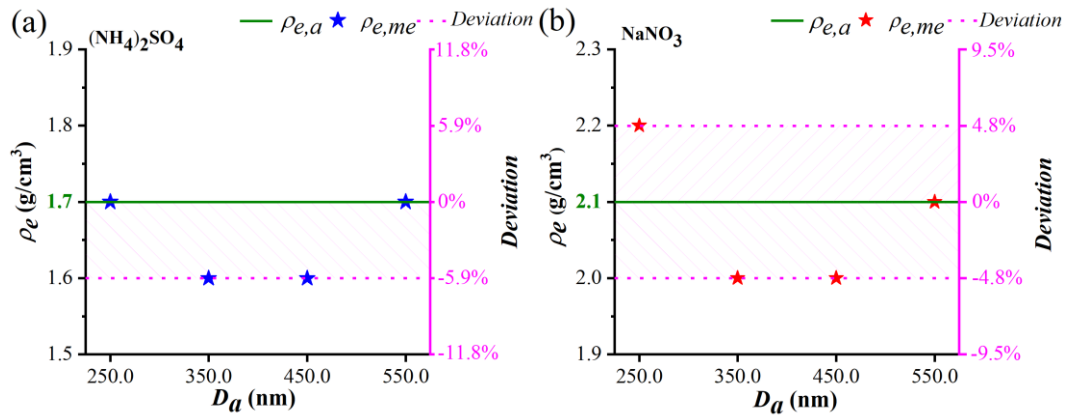
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455 **Figure 2.** (a) Comparison between the measured  $D_{ve}$  ( $D_{ve,me}$ ) and the theoretical  $D_{ve}$  ( $D_{ve,th}$ ) of the

456 PSL particles. (b) Comparison between the measured  $\rho_e$  ( $\rho_{e,me}$ ) and the theoretical  $\rho_e$  ( $\rho_{e,th}$ ) of the

457 PSL particles.





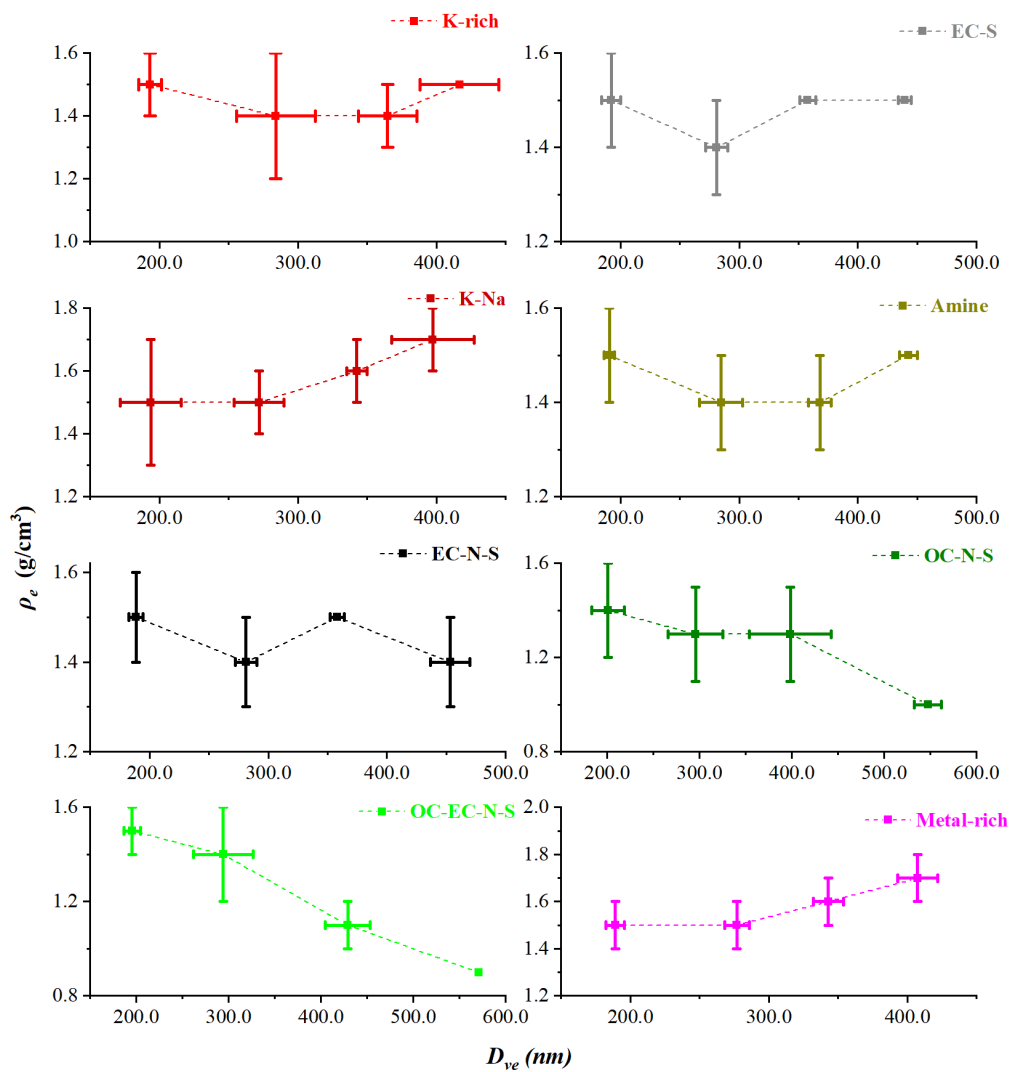
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459

460 **Figure 3.** (a) Comparison between the measured  $\rho_e$  ( $\rho_{e,me}$ ) and average  $\rho_e$  ( $\rho_{e,a}$ ) values of the

461  $(\text{NH}_4)_2\text{SO}_4$  particles. (b) Comparison between the measured  $\rho_e$  ( $\rho_{e,me}$ ) and average  $\rho_e$  ( $\rho_{e,a}$ ) values of

462 the  $\text{NaNO}_3$  particles.



463

464

465 **Figure 4.** Variation in  $\rho_e$  of the eight particle types with  $D_{ve}$ . The solid lines represent the rang of

466 the  $\rho_e$  and  $D_{ve}$  measured from nine rounds, and the data points stand for the average values.