## **Reply to Second Review**

NOTE: Reviewer comments are in black below, while the reply to comments are followed in blue.

Second Review of Dodson and Small-Griswold: Turbulent and Boundary Layer Characteristics during VOCALS-REx

Although the manuscript has been improved, my general impressions remain much the same, which can be summed up as: "A large number of pages and figures, not commensurate with the information content." I still see value in the manuscript's broad documentation of data from the campaign.

I had hoped for a clean, sharp, synthesis in the second round but unfortunately that was not the case. In an attempt to be complete, the authors lose sight of the essence of their work. Heavy **Synthesis** is required.

Hopefully you will find that we have made an extensive effort (both from the first round of revisions and this second round) to reduce and simplify the message being presented as much as is realistically possible without losing the "broad documentation" you see value in. Heavy synthesis and broad documentation do not go hand in hand, heavy synthesis removes a lot of the broad documentation, or a lack of heavy synthesis would keep a lot of the broad documentation. We have tried to find a balance between the two. Again, the word "Characteristics" is in the title, this suggests right from the start this paper sets out to create a broad documentation of work.

Through the synthesis:

- Original Figure 2, 12, and 15 have been removed completely, along with Table 5.
- Panels (a), (b) and (d) have been removed in Original Figure 3.
- Original Figures 8 and 9 have been combined into a single Figure, along with removing Panel (c) from the original Figure 8 and Panels (a) and (b) from the original Figure 9. Panels (c) and (d) (i.e., TKE and TKE dissipation) from original Figure 9 have also been combined into a single panel.
- All w'N' data has been removed from the manuscript discussion and associated figures.

Note that roughly 2.5 pages of text has been removed.

We will address more specific concerns/comments below. Please note that all references provided in these comments can be found in the References section of the manuscript.

In addition, grammatical errors persist throughout the paper. The reader encounters sloppiness already in the abstract. (VOMOS?, LFH, "Data was..."). This is inexcusable from native English speakers. I note some grammatical errors below but it's an incomplete list.

We apologize for the typos. We have corrected VOMOS to VAMOS, after checking that it stands for Variability of the American Monsoon Systems. This is the official name of the campaign (VAMOS Ocean-Cloud-Atmosphere-Land Study Regional Experiment (VOCALS-REx)) as named in Wood et. al. (2011). We can find multiple examples (Zheng et al. 2011, Rahn and Garreaud 2010, Bretherton et al. 2010) where they do not explicitly define VAMOS in the abstract.

Secondly, we find the comment "This is inexcusable from native English speakers" to be offensive, inappropriate, and unprofessional coming from a reviewer of a reputable journal. We also note that the second reviewer said that "The article is overall well-written, and easy to read". Regardless, we will make as much changes as realistically possible to improve how the manuscript reads without completely changing our specific writing "voice".

## A comment from my first review:

## 4) Section 2.1: Shortwave absorption doesn't only occur at cloud top.

• You are correct, solar absorption does not occur only at cloud top (although this is where it is primarily confined). We have clarified this statement in lines 157-160 of the new manuscript.

The revised text says "shortwave absorption, which is largest near cloud top due to the scattering of solar radiation limiting absorption lower in the cloud layer (Hignett, 1991).

But Hignett 1991 states on Pg 1474: "In contrast to the longwave cooling, the heating which results from absorption of shortwave solar radiation is distributed more deeply in the cloud layer". Please do some reading on this topic and revise to better reflect the knowledge in the field.

The point of this comment, is that shortwave radiation acts to stabilize the cloud layer (where longwave cooling is the only source at night, leading to enhanced convection). The entire sentence states "...as turbulence typically displays diurnal patterns, with the strongest turbulent mixing occurring during the night when longwave radiational cooling dominates **due to the absence of the stabilizing effect of shortwave absorption**, which is largest near cloud top due to the scattering of solar radiation limiting absorption lower in the cloud layer (Hignett, 1991)."

The Hignett citation has been moved to just after the bold text above. Hignett on pg. 1474 states "The net result of the combined entrainment, shortwave heating and longwave cooling may be that the cloud layer is warmed relative to the subcloud layer. This implies the formation of a stable layer, which serves to limit the vertical extent of mixing from the surface and reduces transport between the upper and lower parts of the boundary layer".

The sentence in questions states "...of shortwave absorption, which is largest near cloud top...". Attached below are two figures from Stephens (1978), which show shortwave heating profiles. These figures show that shortwave radiational heating is largest at cloud top, with a majority of the heating occurring in the upper 20% of the cloud. The statement made in bold is NOT factually incorrect, it is largest near cloud top. To be more precise, it is at a maximum (or is largest) at cloud top. The stratocumulus review from Wood (2012) states on pg. 2392 "solar heating in a cloud layer is largest at the cloud top...." and goes on to cite Stephens (1978). The

Stephens (1978) citation has been added in the location where the original (Hignett 1991) citation was located. Please see lines 147-150 in the new manuscript.

We understand that shortwave absorption depends on a multitude of variables, such as solar zenith angle, cloud optical depth (which is dependent on droplet size). Along with this, roughly 50% of the solar radiation that is absorbed is done so by water vapor, which is a function of cloud thickness and temperature. Seeing as the reviewer already finds the manuscript long and wordy, we do not see the need to go into a detailed discussion of shortwave radiation absorption, especially since this comment is made to simply state that it acts to stabilize (or at the very least limit convection) the boundary layer. Given the references and text provided above, this discussion should be sufficient.

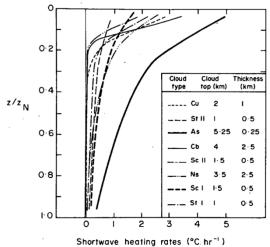


FIG. 5. The shortwave heating profiles in eight cloud models described in Table 3.  $Z_N$  represents the total geometric thickness of the cloud. The surface albedo is 0.3 and normal solar incidence is assumed.

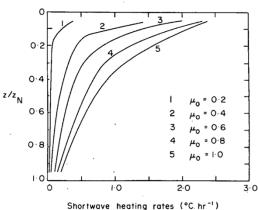


FIG. 6. The shortwave heating rates in a SC I cloud layer for different solar elevation angles. Surface albedo is 0.3 and the cloud is positioned between the 1 and 1.5 km levels in the McClatchey *et al.* (1972) tropical atmosphere.

Below I give examples of problematic text and places that need improved analysis. More than this would require me to sit down and rewrite paragraphs and I'm not going to do that.

We will address all comments (especially in relation to improved analysis or scientific inconsistencies) and reword where we find it improves the clarity of the science being presented.

#### Abstract:

What is "evaporation away from the surface"? Is it evaporation from the surface, or evaporation some distance from the surface? And if the latter, where? You can calculate this from the qc flux profiles. Slight

We have clarified this statement. From Figure 12 in the new manuscript, the orange envelopes bring attention to the layers in which most of the evaporation is occurring. From this, evaporation is occurring between a  $z/z_i$  of 0.4-0.6 (i.e., away from the surface but below cloud base). See lines 14-15 in the new manuscript. This will also be discussed in further detail in specific comments below.

What does "completely offset from one another" mean? Do you mean spatially offset? I waded through the paper and looked at the figures and still don't get it. It's a key result!

Looking at Figure 9 in the new manuscript, Panel (a) shows box plots for in-cloud and sub-cloud TKE measurements. The box represents the 25-75 percent quantiles of the data, with the range of the whiskers representing the range of the data measured. Focusing on Nov. 1<sup>st</sup>, we can see that the boxplots are completely offset from one another with very minimal overlap (i.e., the range of the turbulence measured is completely different for in-cloud and sub-cloud with minimum overlap). This suggest two completely different turbulent environments for in-cloud and sub-cloud.) Simply, the "completely offset from one another" means that the turbulent values for incloud and sub-cloud are completely different.

This is explicitly stated in the manuscript, where lines 473-474 in the original manuscript state "...several cases (including Nov 1<sup>st</sup> and Nov. 2<sup>nd</sup>) where the entire turbulent distribution of the sub-cloud data is shifted to larger values than those of in-cloud data, with minimum overlap". This sentence explicitly states what we mean. We have clarified this in the abstract however. Please see lines 18-19 in the new manuscript.

The inability of the LW cooling to mix through deeper boundary layers is not a new result! Begin your sentence with: "As shown previously ..." and then refer to the papers later (as you do).

This has been added. Please see line 8 in the new manuscript.

## Lines 10/11 require work. Sentence structure, grammar.

"increases" has been changed to "increase". A grammatical check did not find anything wrong, and this appears to be the only possible issue related to the grammar we could find. Grammatical errors are not intentional, if you have specific grammar issues that are not picked up by several standard grammar checkers please let us know what you are thinking. Other than the error corrected above it is unclear what issue you have. Please see line 11 of the new manuscript.

## Introduction:

## Long and unnecessarily wordy "subgride"?

Sentences and unnecessary words (in our opinion at least) have been removed. Roughly a paragraph of text has been removed in the new manuscript from Section 1.1.

"subgride" has been changed to subgrid-scales. Please see line 37 in the new manuscript.

### Line 71: grammar (maintain consistent tense)

Grammar has been corrected to maintain consistent tense. Please see lines 71-73 in the new manuscript.

Line 76: Here and elsewhere: "height throughout the depth of the boundary layer"

"Height throughout the depth of the boundary layer" has been added where necessary. Please see line 77 in the new manuscript.

Line 82: what do you mean by "Convection in the STBL is limited?" By definition, Sc are convective. They may not penetrate very high because of the capping inversion. Please reword.

We were simply stating that convection is limited due to the inversion (which the next two sentence went on to clarify. This phrase has been removed completely in the new manuscript.

Line 91 and throughout: "dryer" is a machine that dries your clothes. "drier".

All three instances where "dryer" is used has been replaced with "drier" Please see lines 91, 515, and 545 of the new manuscript.

## Line 91/92. These parentheses within parentheses are very distracting.

This has been changed. Please see lines 89-92 in the new manuscript.

Line 124: Lewellen et al. 1996

We assume you want us to add the Lewellen et al. (1996) reference. Bretherton and Wyant (1997) state: "A similar conclusion was reached by Lewellen et al. (1996) from numerical simulations of shallow stratocumulus layers subject to prescribed surface fluxes. In this paper, we suggest that surface latent heat fluxes are the most important determinant of decoupling for subtropical CTBLs". Therefore, we will add the Lewellen reference and keep the Bretherton reference. Please see lines 115-116 in the new manuscript.

Line 128: vertical velocity skewness is a strong function of diurnal cycle as shown in VOCALS.

For the sake of simplicity, we will leave out the statement that vertical velocity skewness(w'w'w') is a strong function of diurnal cycle (vertical velocity variance is also a strong function of diurnal cycle, but we do not explicitly state this either). It makes the discussion of w'w'w' even more wordy, and since all but two flights have a 7:00 AM start, this is not necessarily relevant to the results. We have already removed some unnecessary statements in regards to the discussion of w'w'w' to simplify it. We state what the expected w'w'w' is for a well-mixed boundary layer (and consequently a decoupled boundary layer) because this is relevant to the results. Also the vertical profile of w'w'w' for a precipitating boundary layer is relevant to the results, which is why it is briefly discussed.

## **Data and Methods**

Line 153: Although probes failed on some of the days, the key cloud probes are ones that measure LWC. Surely there were other probes measuring bulk LWC (Hot wire, JW, Gerber PVM?) The aerosol data aren't even used in this paper, except as a possible indicator of reduced drizzle. And w'N' analysis is so inconclusive that I see no reason to keep it.

We have removed the w'N' data (where the PDI or Phase Doppler Interferometer is needed for the N' data since this instrument provides a time series of droplet arrival times). Therefore, we have added four more flights for a total of 18 flights (the PDI is what failed on the four flights that were not initially included). Note that there is no cabin data available for the Nov. 5<sup>th</sup> flight. This is also conveyed in Zheng et. al. (2011). w'N' has been removed from the manuscript completely.

A side note that is no longer relevant to the manuscript: Since the PDI provides an accurate count of drops since the PDI provides specific arrival times of each droplet with no instrumentation dead time, we believe the w'N' data is more conclusive than one would believe (as opposed to having only 1-hz CDNC data or some other similar dataset). The lack of discussion on the PDI data may have led to this confusion.

Line 178: "e.g., using Reynolds composition" -- an example of how a little effort can make the sentence more readable.

The use of "e.g.," has been implemented. See line 168 in the new manuscript.

Line 204: "assumption that the flow is isotropic". Why was isotropy never directly calculated (see below)?

The isotropy has been added to Figure 10. See the below discussion on the reviewer comment on the anisotropy ratio

Line 221: structure function (n=2)

The necessary correction has been made. Please see line 210 in the new manuscript.

Line 254: "the error bars would not plot"?? this sentence needs rewriting.

The necessary corrections have been made. Please see lines 242-243 in the revised manuscript.

Line 263: What does "this" refer to? An equation? If so say so and refer to the Eq. number.

Line 263: "this" has been replaced with "Equation 7". Please see line 252 in the revised manuscript.

## Synoptic

## Long and low in information content

Roughly a full-page worth of text has been removed from this section. Also, original sections 3.1 and 3.2 have been combined into a single section. Most of the information removed was from the original section 3.1. We did this to keep the focus on the synoptic variability and boundary layer characteristics, which play a direct role in the turbulence results in various ways.

### Line 286: remove "as was found in"

### This phrase has been removed. Please see line 267 in the new manuscript

### Various places: "minimums" ???

According to Merriam-Webster Dictionary, the plural form of minimum can be either "minimums" or "minima". It is plural since we are talking about the minimum for both the 500 and 700-hPa geopotential heights. This has since been rephrased however. See lines 276-278 in the new manuscript.

#### Line 334: metrics instead of dictators. Line 335, grammar

Metrics has been inserted. Please see line 308 of the new manuscript. The sentence which contained the grammar issues from line 335 has been removed from the new manuscript.

#### Line 342: what does "per measurement" mean?

"per measurement", refers to each 1-hz data point available. If the gradient does not exceed the limit provided in the text between each 1-hz data point, then we are outside the range of the inversion. We have added "1-hz" into the new manuscript. Please see lines 313-314 of the new manuscript.

#### Line 357 grammar

The phrase in question has been removed from the new manuscript.

## Line 365: aircraft typically measure aerosols in their dry state. Please check this was not the case here.

The passive cavity aerosol spectrometer probe (PCASP) is an airborne, optical spectrometer that uses the intensity of scattered laser light to measure an aerosol size distribution from  $0.1-3.0 \,\mu\text{m}$ . The PCASP doesn't specifically dry particles before measuring them unless there is a heater added to the inlet though the sheath air and is dried which can result in the drying of some, but not necessarily all particles We do note that they do not recommend using the PCASP in cloud due to shattering effects. Also, from the blue profiles in Figure 5 of the new manuscript, the N<sub>a</sub> above cloud clearly mirrors the associated profiles of mixing ratio above cloud.

NOTE that the original statement this reviewer comment is based on is no longer in the new manuscript.

https://www.eol.ucar.edu/instruments/passive-cavity-aerosol-spectrometer-probe

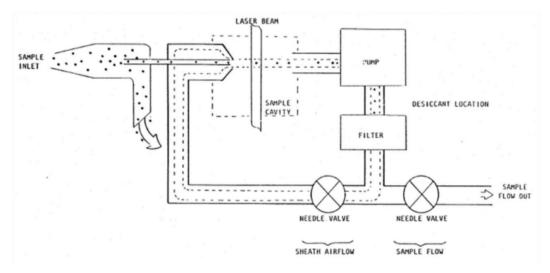


Figure 3: PCASP-100X Airflow Diagram

### Line 376: helped to produce

"Helped to produce" has replaced "in producing" Please see line 335 in the new manuscript.

Line 386 "that"

The phrase in question has been completely reworded. Please see lines 345-346 of the new manuscript.

Line 298: What is mixed layer cloud thickness in this context? Since you're comparing LCL and cloud base, I think you mean mixed layer thickness from surface to cloud base.

As is stated in the manuscript, mixed layer cloud thickness is "the difference between  $z_i$  and the LCL". Per the Jones et al. (2011) (citation provided at the end of the sentence in question) on page 7143: "Decoupling is strongly correlated to the "mixed layer cloud thickness", defined as the difference between the capping inversion height and the LCL".

#### Line 405 grammar

"are" has been replaced with "is" for the sentence in question (we assume this is what you wanted). Please see line 363 in the new manuscript.

#### Line 413: rewrite as "By these metrics 28% of..."

This has been rewritten as suggested. Please see line 372 of the new manuscript.

## Line 420: "relatively consistent" in terms of what? Depth?

In terms of cloud layer thickness. "thickness" has been added to the sentence to clarify. Please see line 378 of the new manuscript.

### Results

Line 435: grammar

This sentence has been slightly rewritten, although we are not sure if the grammar you were referring to has been corrected. Please see lines 392-393 of the new manuscript.

Line 436/437: (Table 3) will suffice.

This has been corrected. Please see line 394 of the new manuscript.

Line 458 "are" not "will be".

"will be" has been changed to "are". Please see lines 408-409 of the new manuscript.

Line 460: "Based on" not "From just analyzing.."

The suggested change has been made. Please see line 412 of the new manuscript.

#### Line 465 grammar

"Coefficient correlation" has been changed to "correlation Coefficient". This is the only thing that we found wrong with this line. Please see line 418 of the new manuscript.

Line 480: driven by

"from" has been replaced with "by", please see line 432 of the new manuscript.

### Change TO11, TO 12 to RF11, RF12 (or vice versa) for consistency

All occurrences of TO have been replaced with RF, both in the main text, figure axis labels and legends, and figure captions.

Line 491: Change "Can be" to "is".

The suggested change has been made. Please see line 449 of the new manuscript.

Line 497: Given that zi decreases...

"acts to" has been removed from "given that  $z_i$  act to decrease...". Please see line 455 of the new manuscript.

Table 4 gives correlations that sometimes don't fit specific cases. The later focus on Nov 1 should be brought in. How well do Nov 1 and 2 follow the broad correlations?

Nov.  $1^{st}$  and  $2^{nd}$  follow the correlations (all correlations were calculated including Nov.  $1^{st}$  and  $2^{nd}$ ). It would not make sense to remove cases that one would not expect to fit a specific correlation. We would consider this manipulating the data to fit a specific narrative. With that said though, we did highlight (see table caption for Table 3) specific cases relating to the decoupling parameter for mixing ratio ( $\alpha_q$ ) where Nov.  $1^{st}$  and  $2^{nd}$  were removed. This is because  $\alpha_q$  had abnormally large values due to the increase in q above the boundary layer.

For an example, though, correlations with and without (respectively) for specific variables are given:

BLH - GPH = -0.49, -0.46 Omega - GPH = -0.71, -0.58 Wind - GPH = 0.37, 0.42 Wind - LHF = 0.56, 0.61Wind - SHF = 0.62, 0.60

We won't go through all the variables listed in Table 4, but as you can see, in general there isn't a large change in the correlations. Since Section 4.3 focuses on the Nov. 1<sup>st</sup> case, we can tie in the broad correlations in that section. See below reply to comments for more detail.

Line 514: "Physically this makes sense" does not make sense. There are many studies showing aerosol causing increases and decreases in LWC.

We are looking at the correlations between Na, Nd and drop size. It is general knowledge that enhanced aerosol load leads to larger cloud droplet number concentrations and smaller droplet sizes. We did not calculate the direct correlations between any of these variables and the LWC. However, we have removed "physically this makes sense". We do go on to explain the observed correlations. Please see line 467 in the new manuscript.

Line 516: What are 'enhanced latent heating effects'? Do you mean 'enhanced latent heating'? And if so where?

Yes, we mean the enhanced latent heating (we say effects since latent heating can be a release of energy (condensation) or absorption of energy (evaporation)). We have clarified this statement, and added "in-cloud" to clarify the location. Please see line 469 in the new manuscript.

Fig. 13: Please add a line on your plots for the anisotropy ratio = 2w'w'/(u'u'+v'v'). It goes a long way in showing the source of TKE and its relationship to top-down/bottom-up driven turbulence. Add discussion as necessary.

The mean anisotropy ratio has been added to Panel (e) of Fig. 13. A brief discussion has been added to the text. Please see lines 487-490 of the new manuscript.

Line 581: "downdrafts are smallest". Do you mean smallest areal coverage?

Smallest in spatial coverage. Given we define w'w'w' in the introduction "A positive (negative) vertical velocity skewness indicates that strong, narrow updrafts (downdrafts) are surrounded by larger areas of weaker downdrafts (updrafts)". We shouldn't have to repeat this definition. Especially given the fact that the sentence in question goes on to state "...the downdrafts are smallest, yet strongest at cloud base while updrafts are spatially larger, yet weaker". Regardless, we have added that "...downdrafts are **spatially** smallest, yet strongest at cloud base...". Please see lines 521-522 of the new manuscript.

I'm not convinced that this paper has shown anything interesting about w'N'. The plots are incredibly noisy and no clear results are shown. So why clutter the paper with results like this. Please go back through the paper with a similar view on other plots and remove if they don't yield a clear result. You will be doing yourselves a real service in the long run.

All w'N' data has been removed, along with the included discussion in the text. Our comments at the beginning provide a summary of figures and figure panels that have been removed or changed.

A moot point perhaps but on your figures for w'N', cc-1 is not a standard unit.

We assume you would prefer cm<sup>-3</sup>? Actually, this is a mistake on our part anyways, as N' is the number of drops per second (not a concentration such as Nd).

Figure captions are meant to be cryptic. Please remove all "Note that.." from figure captions.

All "Note that" have been removed from the figure captions.

Figure 16 caption doesn't even say its Nov 1.

Thank you for pointing this out. Nov. 1<sup>st</sup> has been added. Please see the figure caption for Figure 12.

Line 619: Theta is normalized 0 to 1 (min=0; max = 1). But since theta keeps increasing above the BL, this would typically place the unit value at the highest point of your profile. So how did you normalize theta? The plot suggests you did this correctly.

Theta in original Figure 16 is NOT normalized from 0 to 1...please look at the figure axis. Normalizing theta from 0 to 1 was an example to bring attention to the fact that a large amount of entrained air had resulting in increasing theta throughout the boundary layer down to  $z/z_i =$ 0.4. This has been removed to simplify the text and not confuse the reader.

# Line 625: Where is the evaporation occurring? It makes a difference. Why not calculate the divergence of the qc flux?

We do not know what you mean by the qc flux (do you mean the divergence in the cloud droplet mixing ratio flux?).... regardless we do not believe this is necessary. The profile of F $\theta$  shows a clear sink for theta (i.e., cooling) while there is a slight enhancement of Fq (i.e., an enhanced source of vapor). The text explains that this is an indicator that evaporation is occurring. And that it primarily is occurring in the orange envelopes. We do believe that the new Panel (c) with Nd

and drop size proves that precipitation is occurring, and shows that it is primarily confined above  $z/z_i = 0.40$ . Please see lines 575-584 in the new manuscript for the new discussion on this.

## Pgph starting line 629: there's another factor and that is less dilution of cloud water in response to entrainment.

The second factor states "Entrainment of more moist air near cloud top, reducing evaporational cooling...". This is only the case if there is less dilution of cloud water. If there was more dilution of cloud water, there would be no reduction in evaporative cooling. The reduced evaporation from the entrainment of more moist air is in relation to the fact that if the air was not more moist, there would not be reduced evaporation. We have reworded this to read "Entrainment of more moist air near cloud top, reducing evaporational cooling that would otherwise occur through the entrainment of drier air". Hopefully this is sufficient to account for this comment. Please see line 544-546 in the new manuscript.

### Line 632 grammar

We have changed "top and slows" to "top, slowing" hopefully this is what you were referring to. Please see line 546 in the new manuscript.

Line 638 it's more about the cooling profile than the precip rate, as you discuss later. Please shorten/tighten descriptions.

These few sentences have been reworded. Please see lines 551-553 in the new manuscript.

Line 646: What are 'latent heating effects'? Don't you just mean 'latent heating'?

Yes, we mean "latent heating", this has been changed. Please see lines 559-560 in the new manuscript.

Also, note that "latent heating effects" is common terminology. Wood (2012), a paper with over 600 citations, uses the phrase repeatedly. One example on page 2394: "Large buoyancy fluxes are caused mainly by radiative cooling and are additionally enhanced by latent heating effects".

Line 652 and abstract/conclusions, "away from the surface" is highly imprecise language, and since this is a key result please make your arguments crystal clear. The profiles show no subcloud drizzle that would destabilize the sub-cloud layer. Does LWC include drizzle/rain?

Panel (c) in Figure 12 now makes it clear that (1) yes there is sub-cloud drizzle occurring. The text also clarifies where this evaporation is occurring (primarily near  $z/z_i \sim 0.40$  and 0.60, the two orange envelopes from Figure 12). As per the Phase Doppler Interferometer (PDI), and the discussion on lines 575-583 of the new manuscript, there are so few droplets in the sub-cloud layer we would assume this is why the LWC profile does not indicate any liquid water. The PDI is an extremely precise instrument. The droplet arrival time can accurately be measured to < 3.5 µs, resulting in accurately mapping droplets down to 2.1E-4 m (assuming average aircraft speed).

Also, note that in the w'N' profile from the original manuscript, this would have indicated that yes, there was sub-cloud drizzle occurring (as w'N' would be zero when there is no N present, as

was the case for original Figure 17).

As mentioned above, please tie this discussion of Nov 1 back to the broader correlations and point out differences if they exist, or simply note consistency. I couldn't wade through the text one more time to figure it all out.

We will simply point out consistency. Please see lines 603-605 of the new manuscript.

Line 652: "activation" not "activation through condensation". Please stick to accepted terminology.

All discussions relating to w'N' have been removed.

Line 664: I don't believe you have evidence of aerosol activation in updrafts at cloud base. Your measurements are far too noisy to draw this conclusion. Updraft velocities are just not strong enough to raise supersaturation at cloud top. Why is this not simply evaporation in downdrafts? And as I said earlier, since you can't really say much about aerosols, why dilute your message with distracting, uncertain results?

All discussions relating to w'N' have been removed.

Line 655, poorly mixed BLs can be locally well coupled by local, penetrative updrafts.

This is true, but we do not see any evidence of this for the Nov. 1<sup>st</sup> case. Therefor we do not see the need to include such a statement.

Line 671: Can you explain why the consistency in wind direction across the inversion should translate to mixing across the inversion? It's possible that mixing occurs but this isn't a way to address it. The theta jump is about 4K.

This discussion has been removed to simplify the text.

Fig. 13: Why are there 2 identical blue lines for RF12 and 13 when the caption and legend lump them together?

There are two lines, for the two flights (RF11 and RF12)...and they are not identical. Each line represents the profile for each flight (as was the case in original Figure 6 and 12). Perhaps you are looking at a different figure than Fig. 13? We went through all the figures and cannot identify what you are saying in this comment.

We did change the blue color however. Instead of two identical light blue profile lines, we have one dark blue and one light blue profile line in Figures 3, 5, and 10.

General comment on SHF and LHF: these are generally defined at the surface so please clarify. Eqns 2, 3 give the definitions. You should define them as w'q', w'theta'!

This has been implemented throughout the text (i.e.,  $F_{\theta}$ ,  $F_{q}$ , and  $F_{\theta v}$  per the equation definitions) unless we are specifically talking about surface values, in which case we use LHF and SHF.

## Conclusions

Conclusions need to be tightened/clarified.

Clarification has been made where needed (i.e., specific instances outlined below).

Line 703: entrainment acts to increase zi. Zi is a number that can increase or decrease but not deepen.

This has been corrected (throughout the manuscript). Please see lines 622 of the new manuscript.

### Line 710: please add appropriate reference

This statements was also made in Section 4.1. The reference (Bretherton and Wyant 1997) has been inserted at that location. Please see lines 96-98 of the new manuscript. Adding the reference here at the end seems out of place.

#### Line 711: grammar

"As the LHF and SHF increases (decreases), zi increases (decreases)" has been changed to " As the LHF and SHF increase (decrease), zi increases (decreases)". Please see lines 628-631 of the new manuscript.

## Line 712/713: this is a convoluted sentence. Deepening BLs entrain drier air and therefore increase surface LHF

This is incorrect. A deepening boundary layer through the entrainment of drier air does not lead to an increase in surface LHF. The increase in LHF is what leads to the deepening boundary layer. Please see the discussion on lines 457-464 of the new manuscript (back in Section 4.1). This sentence has been restructured to make the point more clear. Please see lines 628-631 of the new manuscript.

Line 715: Again, please tell us what you mean by evaporation away from the surface.

This has been clarified throughout. Please see lines 633-634 of the new manuscript.

### Line 713: again explain: "completely offset one another"

This has been clarified throughout the manuscript.

#### Line 738: grammar

We were unable to determine a grammatical error on or around line 738.

## List of Relevant Changes

- (1) Four flights were added to the analysis, increasing the total from 14 to 18 flights.
- (2) Section 3.1 and Section 3.2 were combined into a single section: Section 3.1 Synoptic Variability at Point Alpha.
- (3) Figure 2, Figure 12, Figure 15, and Table 5 have been removed completely from the original manuscript.
- (4) Panels (a), (b), and (d) have been removed from Figure 3 in the original manuscript.
- (5) Figures 8 and 9 in the original manuscript have been combined into a single Figure (Figure 7 in the new manuscript). Panel (c) from the original Figure 8 and Panels (a) and (b) from the original Figure 9 have been removed. Panels (c) and (d) (i.e., TKE and TKE dissipation) from original Figure 9 have been combined into a single panel.
- (6) All w'N' data have been removed from the manuscript, including Panel (d) in Figure 14 and from Panel (c) in Figures 16 and 17.

## **Turbulent and Boundary Layer Characteristics during VOCALS-REx**

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#### Abstract.

Boundary layer and turbulent characteristics (surface fluxes, turbulent kinetic energy (TKE), turbulent kinetic energy dissipation rate (ε), etc.), along with synoptic scale changes in these properties over time, are examined using data collected from 14–18 research flights made with the CIRPAS Twin Otter Aircraft. Data was were collected during the VOMOS-VAMOS
Ocean-Cloud-Atmosphere-Land Study-Regional Experiment (VOCALS-REx) at Point Alpha (20°S, 72°W) in October and November of 2008 off the coast of South America. The average boundary layer depth is found to be 1175-m1148-m, with 28% of the boundary layer profiles analyzed displaying decoupling. Analysis of correlation coefficients indicate that as atmospheric pressure decreases, the boundary layer height (z<sub>i</sub>) increases. The As has been shown previously, the increase in z<sub>i</sub> is accompanied by a decrease in turbulence within the boundary layer. As z<sub>i</sub> deepensincreases, cooling near cloud top cannot sustain
mixing over the entire depth of the boundary layer, resulting in less turbulence and boundary layer decoupling. As the latent heat flux (LHF) and sensible heat flux (SHF) increasesincrease, z<sub>i</sub> increases, along with the cloud thickness decreasing with increasing LFHLHF. This suggests that an enhanced LHF results in enhanced entrainment which acts to thin the cloud layer

while deepening the boundary layer.

A maximum in TKE on Nov.  $1^{st}$  (both overall average and largest single value measured) is due to <u>sub-cloud</u> precipitation

- 15 acting to destabilize the sub-cloud layer while acting to stabilize the cloud layer (through evaporation occuring away from the surface), while acting to stabilize the cloud layer, primarily confined between a normalized boundary layer height  $(z/z_i)$ of 0.40 to 0.60). Enhanced moisture above cloud top from a passing synoptic system also acts to reduce cloud top cooling, reducing the potential for mixing of the cloud layer. This is observed in both the vertical profiles of the TKE and  $\epsilon$ values, where it is found that the distributions of turbulence for the sub-cloud and in-cloud layer are completely offset from one another  $\frac{1}{2}(i.e.)$
- 20 the range of turbulent values measured have slight or no overlap for the in-cloud and sub-cloud regions), with the TKE in the sub-cloud layer maximizing for the analysis period, while the TKE in the in-cloud layer is below the average in-cloud value for the analysis period. Measures of TKE,  $\epsilon$ , vertical velocity variance, TKE, and the buoyancy flux averaged over all 14-18 flights display a maximum near cloud middle (between normalized in-cloud values of 0.25-0.75). Seven of the fourteen Ten of the 18 flights display two peaks in TKE within the cloud layer, one near cloud base and another near cloud top, signifying evaporative
- 25 and radiational cooling near cloud top and latent heating near cloud base. <del>Overall, it appears that turbulence measured at Point</del> Alpha is weaker than that measured over the open ocean to the west of Point Alpha, and that measured during other scientific

campaignsDecoupled boundary layers tend to have a maximum in turbulence in the sub-cloud layer, with only a single peak in turbulence within the cloud layer.

Copyright statement. TEXT

#### 30 1 Introduction

Stratocumulus (Sc) clouds have a significant impact on climate due to their large spatial extent, covering approximately 20% of Earth's surface (23% over the ocean and 12% over the land) in the annual mean (Randall et al., 1984). According to Wood (2012), the subtropical eastern oceans in particular are marked by extensive regions of Sc sheets (often referred to as semipermanent subtropical marine stratocumulus sheets). Of those, the The largest and most persistent Sc deck in the world,

- 35 the Peruvian Sc deck, lies off the west coast of South America (Bretherton et al., 2004), making its role in climate an essential building block to improved modeling of the overall earth system. A better understanding of Sc decks is therefore necessary to improve our physical understands understanding of mechanisms controlling Sc clouds, and to improve confidence in climate model sensitivity (Zhang et al., 2013), especially considering climate models suffer from order-one uncertainties in Sc cloud representation (Noda and Satoh, 2014; Gesso et al., 2015).
- It is a challenge for models to successfully simulate the Peruvian Sc deck due to the importance of subgride subgrid scales and physical processes which are poorly represented (Wood et al., 2011). Most models continue to struggle with the boundary layer vertical structure (Wyant et al., 2010) which is important for determining Sc cloud properties. One example, as discussed in Akinlabi et al. (2019), is that a robust estimation of the turbulent kinetic energy dissipation rate ( $\epsilon$ ) is needed when creating subgrid models for Lagrangian trajectory analysis of passive scalars (Poggi and Katul, 2006) or large-eddy simulation. Other
- 45 vertical profiles of turbulent fluxes (liquid water, water vapor, energy) determine the mean state of the boundary layer and the resulting properties of the Sc deck (Schubert et al., 1979; Bretherton and Wyant, 1997; Ghate and Cadeddu, 2019).

Although turbulence is critical to atmospheric boundary layer, microphysical, and large scale cloud dynamics, it is difficult to measure, with literature on describing cloud-related turbulence based on in situ data being scarce (Devenish et al., 2012; Shaw, 2003). This study therefore aims to characterize turbulence throughout the vertical profile of the Stratocumulus

- 50 <u>stratocumulus</u> topped marine boundary layer (STBL) over a three-week observation period in October and November of 2008 during the Variability of the American Monsoon Systems (VAMOS) Ocean-Cloud-Atmosphere-Land Study-Regional Experiment (VOCALS-REx). A large in situ dataset was collected throughout the boundary layer <del>, allowing for analysis (on a variety of spatial and temporal scales) in the aims with the goal of improving predictions of the Southeast Pacific coupled oceanatmosphere-land system (Wood et al., 2011). This dataset allows for a classification of turbulent properties not only through</del>
- 55 vertical profiles, but provides an opportunity to analyze how turbulence changes within the boundary layer with varying synoptic conditions.

The main objectives of this paper include a quantification of the amount of turbulence occurring within the boundary layer through the evaluation of turbulent kinetic energy (TKE),  $\epsilon$ , and other turbulent flux measurements. In particular, the main goals include: (1) <u>Analyze analyze</u> day to day variability in turbulent measurements and boundary layer characteristics, relating them

to synoptic changes in meteorological conditions; (2) Determine determine average turbulent values throughout the vertical

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structure of the STBL, classifying the STBL based on different turbulent profiles analyzed.

There has been a plethora of publications stemming from the VOCALS-REx campaign over the last ten years. Papers range from focusing on climatic and synoptic conditions for the VOCALS region (Toniazzo et al., 2011; Rahn and Garreaud, 2010a, b; Rutllant et al., 2013), analyzing cloud-aerosol interactions (Jia et al., 2019; Blot et al., 2013; Painemal and Zuidema,

- 65 2013; Twohy et al., 2013), and analyzing precipitation, boundary layer decoupling, and other boundary layer characteristics (Jones et al., 2011; Bretherton et al., 2010; Terai et al., 2013; Petters et al., 2013; Zheng et al., 2011), to name a few. A total of five aircraft platforms and two ship based platforms were utilized during VOCALS-REx (Wood et al., 2011), with most publications from VOCALS-REx relying and/or focusing on aircraft observations and other data sources outside of those used here (all but Zheng et al. (2011) and Jia et al. (2019) mentioned above). Results found and presented here therefore provide not
- 70 only a collection of in-situ-in situ turbulent measurements, but provide for the opportunity to relate results to other findings at additional measurement locations within the VOCALS domain. An extensive look at turbulent characteristics of the boundary layer during VOCALS-REx does not exist (note that although Zheng et al. (2011) does give a broad analysis of boundary layer characteristics, their focus on turbulence was minimal), which is puzzling given that the Twin Otter aircraft (the data used here, see Section 2.1) was instrumented with an objective to make turbulence measurements.
- 75 Section 1.1 introduces typical boundary layer vertical structure and scientific background. Section 2 provides an overview of the data and methods, followed by synoptic and boundary layer characteristics during VOCALS-REx in Section 3. Section 4 will evaluate and discuss evaluates and discusses the results. Section 5 will provide provides concluding remarks.

#### 1.1 Boundary Layer Vertical Structure

The vertical structure of the boundary layer is strongly tied to the horizontal and vertical structure of Sc clouds (Lilly, 1968; 80 Bretherton et al., 2010). The STBL is characterized by Sc cloud tops located at the base of an inversion, with subsiding air aloft (as part of the descending branch of the Hadley cell circulation) and well mixed conditions and near-constant conserved variables with height throughout the depth of the boundary layer (Wood, 2012). Multiple papers have analyzed typical well mixed STBL vertical structures (i.e., Albrecht et al. (1988); Nicholls (1984)), showing constant potential temperature and mixing ratio with height-throughout the depth of the boundary layer up until the inversion, when the mixing ratio (potential

85 temperature) sharply decreases (increases). Horizontal winds (both direction and velocity) are typically constant with height throughout the throughout the depth of the well mixed boundary layer, with changes in both direction and strength typically present at the top of the STBL, influencing cloud-top entrainment (Mellado et al., 2014; Kopec et al., 2016; Schulz and Mellado, 2018).

Convection in the STBL is limited. Unlike updrafts through convective heating over the ground, updrafts within the STBL 90 do not penetrate the inversion. This is because convection within the STBL is is primarily driven by cooling near cloud top and

not heating at the ocean surface, where cloud top cooling is primarily from a combination of (1) longwave radiational cooling and (2) evaporational cooling from entrainment. The cloud top cooling leads to instability and the convection of warmer, moist air at the surface (Lilly 1968). The cloud cover is greatest when the STBL is shallow  $[0.5 < z_i < 1 \text{ km}]$ , where  $z_i$  is the inversion layer (i.e., boundary layer) height (Wood and Hartmann, 2006).

- 95 It is known that clouds are areas of enhanced turbulence (Pinsky and Khain, 1996). Therefore, Sc sheets are turbulent but in contact with an almost non-turbulent upper atmospheric environment. The boundary layer top is characterized by several strong gradients, including the cloud boundary (gradient in liquid water content), the entrainment zone (gradient in vorticity), where the entrainment zone separates regions of weak and strong mixing between laminar (warmer and dryerdrier) flow above and turbulent (cooler and more moist) flow below), and the capping inversion (gradient in potential temperature). The cloud
- 100 boundary typically lies in the entrainment zone (Albrecht et al., 1985; Kurowski et al., 2009; Malinowski et al., 2013), which in turn lies in the capping inversion, although these layers do not necessarily coincide (Mellado, 2017). Turbulent analysis of these layers in Jen-La Plante et al. (2016) found that turbulence (both TKE and  $\epsilon$ ) decreases moving from cloud top into the free atmosphere above. Through cloud top entrainment, the STBL deepens beyond 1-km and can become decoupled. According to Bretherton and Wyant (1997), due to longwave cooling at the cloud top being unable to maintain mixing of the positively
- 105 buoyant entrained air over the entire depth of the STBL, the upper (cloud containing) layer (turbulence driven primarily by cloud top cooling) becomes decoupled from the surface moisture supply(turbulence driven by surface-fluxes and shear).

The vertical profile of various turbulent fluxes, particularly that of buoyancy buoyancy flux (which is dependent on moisture and heat fluxes which drive buoyancy differences), can tell one a lot about the state of the STBL. For a boundary layer to remain well mixed, the vertical Vertical energy and moisture fluxes must be linear functions of height - According to

- 110 Bretherton and Wyant (1997), the buoyancy flux is not a linear function of height however (unlike that of a dry boundary layer ) for a boundary layer to remain well mixed, which is not the case for a cloud-containing boundary layer (Bretherton and Wyant, 1997). An increase in the buoyancy flux above cloud base is typically proportional to the upward transport of liquid water that is required to sustain the cloud against entrainment drying (i.e., continued mixing of the cloud layer is sustained by surface fluxes). Decoupling of the boundary layer (and the subsequent decrease in cloud cover) can occur when the sub-cloud buoy-
- 115 ancy fluxes become negative, capping convection below cloud base (Albrecht et al., 1988; Ackerman et al., 2009). According to Shaw (2003), one of the main sources of TKE in clouds is evaporative cooling (due to the entrainment of dry air) and condensational heating (due to droplet condensational growth), implying the buoyancy flux is the primary generator of TKE in the STBL (Schubert et al., 1979; Heinze et al., 2015). Given this, the buoyancy flux nearly always has a maximum in the cloud layer (Nicholls and Leighton, 1986; Bretherton and Wyant, 1997), with TKE being generated due to longwave
- 120 and evaporational cooling at cloud top, and condensational heating at cloud base (Moeng et al., 1992). Nicholls (1989) observed through aircraft observations that the largest buoyancy fluxes are close to cloud top, with further observations (Caughey et al., 1982; Nicholls, 1989) suggesting that the descending regions of air originating near cloud top are more a result of radiative cooling rather than evaporative cooling.

The main source of moisture for the STBL is supplied by the surface latent heat flux (LHF), making it an important source of buoyant TKE production (Bretherton and Wyant, 1997), with the surface sensible heat flux (SHF) typically being a much weaker source of turbulence. The SHF and LHF can be compared using the Bowen ratio (the ratio of the sensible to the latent heat flux). A larger LHF (or smaller Bowen ratio) leads to decoupling of the boundary layer due to the LHF concentrating convective energy generation (through condensational and evaporational heating/cooling) within the cloud layer. To state this another way, an An enhanced LHF leads to increased moisture transport to the cloud layer and a thicker cloud, producing

130 more turbulence and enhanced entrainment cooling near cloud top and more turbulence. Enhanced entrainment results in a deepening of the boundary layer, which favors decoupling (Jones et al., 2011). It is argued in (Bretherton and Wyant, 1997). Bretherton and Wyant (1997) and Lewellen et al. (1996) that the surface LHF is the most important determinant of decoupling within the STBL.

Vertical velocity variance typically displays the strongest updrafts and downdrafts in the upper half of the STBL (Hignett,

- 135 1991; Heinze et al., 2015; Mechem et al., 2012), consistent with the largest production of turbulence being contained within the cloud layer. A positive (negative) vertical velocity skewness indicates that strong narrow updrafts (downdrafts) are surrounded by larger areas of weaker downdrafts (updrafts). It has been found that negative vertical velocity skewness is typically contained within most of the cloud layer and below (Nicholls and Leighton, 1986; Nicholls, 1989; Mechem et al., 2012) for well mixed boundary layers, whereas a decoupled boundary layer containing cumulus below stratocumulus may contain positive vertical
- 140 velocity skewness (de Roode and Duynkerke, 1996) due to convection being driven in the surface layer(as compared to cooling near cloud top). The tendency of the vertical velocity skewness to be positive in a strongly precipitating STBL is also well known (Ackerman et al., 2009), with precipitation being a key contributor leading to boundary layer decoupling (Rapp, 2016; Yamaguchi et al., 2017; Feingold et al., 2015).

#### 2 Data and Methods

#### 145 2.1 Data

Data was collected during the Variability of the American Monsoons (VAMOS) Ocean Cloud-Atmosphere-Land Study-Regional Experiment (VOCALS-REx) from the Peruvian Stratocumulus deck off the west Coast of Chili and Peru during October and November of 2008. VOCALS-REx used various platforms, including five aircraft and two research vessels to accumulate an extensive dataset of the boundary layer, lower free troposphere, and cloud deck along 20°S from 70°W to 85°W.
150 Although multiple sampling platforms, locations, and mission types were deployed during the campaign (see Wood et al. (2011)), data collected by the Center for Interdisciplinary Remotely-Piloted Aircraft Studies (CIRPAS) Twin Otter aircraft will be the focus of this paper, which collected data in the vicinity of 20°S, 72°W; from here on termed Point Alpha. The Twin Otter aircraft was operational for 19 flights from October 16<sup>th</sup> to November 13<sup>th</sup>, 2008.

The Twin Otter platform is ideal for a turbulent analysis of the boundary layer due to the aircraft being instrumented to make turbulence and cloud microphysics measurements, with the same location being sampled for each flight. The Twin Otter is also a relatively slow-moving aircraft with a flight speed of roughly  $55 \text{ to } 60 \text{ ms}^{-1}$ , allowing for a higher resolution of spatial sampling as compared to a faster moving aircraft. Each of the Twin Otter flights was carried out using a stacked flight path (Wood et al., 2011), which involved using stacked legs of 50-100 km in length (horizontal flight paths) to sample various levels of the boundary layer and cloud layer, with at least one aircraft vertical sounding (vertical profile) performed for each flight

160 where the aircraft sampled the free upper troposphere and boundary layer in a single ascent or descent. Each flight of five hours originated from Iquique Chile, allowing for roughly three hours of sampling at Point Alpha.

Of the 19 flights performed by the Twin Otter, only 14-18 are used here due to instrumentation failure on five\_one of the flights (Phase Doppler Interferometer and the cloud/aerosol probe Nov.  $5^{th}$ ). Table 1 displays each of the Research Flights (RF) used in this paper. All flights occurred during the day, with all but two flights (RF 8 and RF 17RF8 and RF17) starting around

- 165 7:00 AM local time, with the first vertical profile flown around 8:00 AM local time at Point Alpha. Having each flight sample the same location at roughly the same time is critical, as turbulence typically displays diurnal patterns, with the strongest turbulent mixing occurring during the night when longwave radiational cooling dominates due to the absence of the stabilizing effect of shortwave absorption , which (Hignett, 1991), where solar absorption is largest near cloud top due to the scattering of solar radiation limiting absorption lower in the cloud layer (Hignett, 1991)(Stephens, 1978).
- Meteorological variables were collected at 40-Hz (including u, v, and w wind velocity, water vapor mixing ratio (q) and potential temperature (θ), to name a few) while most cloud and aerosol data were collected at 1-Hz. A 5-port Radome wind gust probe was used with plumbing that effectively trapped liquid water, preventing any liquid water from obstructing the pressure transducer lines. There were zero failures during the campaign, with an accuracy of ± 0.4 ms<sup>-1</sup> for horizontal wind components and ± 0.2 ms<sup>-1</sup> for vertical velocity. The LI-COR 7500 H<sub>2</sub>O/CO<sub>2</sub> gas analyzer was used for all measurements of absolute humidity and q, with an ambient air intake setup that resulted in the LI-COR source and detector window to be
- liquid free, even during prolonged cloud penetrations. The LI-COR accuracy is reported to be within 1% of the actual reading. Further instrumentation information can be found in Zheng et al. (2010) and Wood et al. (2011).

To analyze the synoptic conditions over the study period, data from the National Centers for Environmental Prediction (NCEP) / National Center for Atmospheric Research (NCAR) Reanalysis Project (NNRP, Kistler et al. (2001)) will be used.

180 The data resolution of the NCEP/NCAR reanalysis data is 2.5° x 2.5° x 17 pressure levels, available at six hour intervals. The resolution of this data is suitable for analyzing synoptic scale patterns, but is not ideal for depicting mesoscale variability that may be present on any given day. Boundary layer height is also derived from relative humidity data from the European Centre for Medium-Range Weather Forecasts (ECMWF) Re-Analysis (ERA5), which has a resolution is 0.25° x 0.25° x 37 pressure levels, and is available at an hourly interval (Hersbach et al., 2020).

#### 185 2.2 Turbulent Calculations

The randomness of turbulence makes deterministic description difficult, limiting description to statistics and average values of turbulence, in particular that of e.g. Reynolds decomposition (or averaging). Reynolds decomposition uses a mean value (over some time period, determined by low pass filtering or applying a linear trend) and subtracts it from the actual instantaneous velocity to obtain the turbulent component (or perturbation value). Reynolds decomposition is based on the underlying

190 assumption that the turbulence is isotropic and stationary, conditions that are hardly fulfilled for atmospheric boundary layer flows however, especially when working with data spanning larger timeframes. The problem is defining how to average collected data to best represent the mean and turbulent components for the fluid flow (with shorter subsets of data having more stationary properties in general than that of longer subsets of data). Using the 40-Hz data, a 320-point averaging window is used here for all turbulent analysis, following the methods outlined in Jen-La Plante et al. (2016). A 320-point averaging window

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corresponds to 8 second subsets of data, or a roughly 440-m subset of data in the horizontal spatial scale (assuming average aircraft speed of 55 ms<sup>-1</sup>). Linear regression is then applied to each 320-point averaging window to calculate the mean and determine the perturbation values.

Applying the averaging method discussed above leads to the calculation of the fluctuations of the *u*, *v*, and *w* components of the velocity, along with other parameters used to measure various turbulent fluxes. Variables to be obtained include turbulent 200 kinetic energy, which is given by:

$$TKE = \frac{1}{2} \left( \overline{u'^2} + \overline{v'^2} + \overline{w'^2} \right) \tag{1}$$

where u', v', and w' are the fluctuations of the velocity components. The turbulent sensible heat, latent heat, and buoyancy fluxes will also be obtained, given by:

$$F_{\theta} = C_p \overline{\rho} \overline{w' \theta'} \tag{2}$$

$$F_q = L_v \overline{\rho} \overline{w'q'} \tag{3}$$

$$F_{\theta_v} = C_p \overline{\rho} \overline{w' \theta'_v} \tag{4}$$

respectively. Where  $C_p$  is the specific heat of air (1005 J kg<sup>-1</sup>K<sup>-1</sup>),  $L_v$  is the latent heat of vaporization at 20°C (2.45 · 10<sup>6</sup> J kg<sup>-1</sup>),  $\rho$  is the mean air density, and  $\theta'$ , q', and  $\theta'_v$  are the potential temperature, mixing ratio, and virtual potential temperature perturbations, respectively. Note that  $\theta_v$  (given by  $\theta_v = \theta(1 + 0.61q - q_l)$ ) is commonly used as a proxy for density when calculating the buoyancy. Humid air has a warmer  $\theta_v$  because water vapor is less dense than dry air, while liquid water drops (if falling at terminal velocity) make the air heavier and therefore associates with a colder  $\theta_v$ , where  $q_l$  is the liquid water

mixing ratio.

Just like that of Reynolds decomposition, the calculation of  $\epsilon$  is based on conditions the assumption that the flow is isotropic (i.e., uniformity uniform in all directions), making the measurement of  $\epsilon$  challenging. In particular, classical turbulence theory

- 215 in the inertial subrange from Kolmogorov (1941) is based on assumptions of local isotropy. With that said, there are multiple methods to measure  $\epsilon$ , including the inertial dissipation method, structure functions, and the direct method. Siebert et al. (2006) found that both the inertial dissipation and structure function methods are useful, but the inertial dissipation method sometimes underestimates  $\epsilon$  at low values due to no clear inertial subrange behavior being observed in the power spectral density, which is not the case for the structure function. The structure function method is therefore considered more robust for cases with small
  - 220

and an average dissipation rate will be calculated from the three components.

The calculation of  $\epsilon$  comes from the analysis of the velocity perturbations through the n<sup>th</sup> order structure function (i.e., a statistic to analyze common variation in a time series). The perturbations, as for other turbulent parameters, are determined with respect to an averaging window of 320-points. Each subset of perturbations is then appended to the end of the previous

values of  $\epsilon$ , and will be used here. Due to questions of isotropy,  $\epsilon$  will be evaluated on the u, v, and w components of the wind,

subset to create a single time series of velocity perturbations. The structure function is given by:

$$S_n(l) = \left(\overline{|u(x+l) - u(x)|}\right)^n \tag{5}$$

where *l* is the distance (or in the case of a temporal series, *l* is equivalent to *t* assuming constant flight speed). From Frisch (1995),  $\epsilon$  using the n<sup>th</sup> order structure function can be obtained by using:

$$S_n(l) = C_n |l\epsilon|^{\frac{n}{3}} \tag{6}$$

where  $C_n$  is a constant of the order 1. The second order structure function (n = 2) will be used here(n = 2), where  $C_2 = 2$  for transverse velocity fluctuations and  $C_2 = 2.6$  for longitudinal velocity fluctuations (Chamecki and Dias, 2004), where vertical fluctuations are considered transversal and horizontal fluctuations are considered longitudinal. The structure function follows a 2/3 power law within the inertial subrange, and will only be used to calculate  $\epsilon$  between frequencies of 0.3-5-Hz, neglecting the higher frequency features attributed to interactions with the plane (i.e., vibrations due to the aircraft) and other instrumental artifacts.

Figure 1 Panel (a) provides the power spectral density of vertical velocity and q for three horizontal flight legs within RF3, one in-cloud, one sub-cloud, and one near surface. Note that the power spectral density follows a -5/3 power law fit (red) within the inertial subrange (as opposed to the 2/3 power law fit of the structure function). A spike in energy can be seen at ~10-Hz, which represents the aircraft interactions discussed previously. The power spectral density overlaid in black represents a single

calculation using a 320-point averaging window. The data follows the -5/3 fit well, and the inertial subrange is well resolved for the averaging window used (with the light gray envelope representing the 0.3 to 5-Hz range). A lack of significant flattening within the power spectra at higher frequencies suggests that the random noise level is low (this is more evident in the vertical velocity spectra than that of the *q* spectra).

Analysis of the turbulence as presented here introduces two types of error, including sampling and noise error. This must

- 245 be analyzed to determine the statistical significance when analyzing vertical profiles, especially since error propagation into higher order moments can be significant (McNicholas and Turner, 2014). Sampling errors were estimated using approaches derived and discussed in Lenschow et al. (1994, 2000) and will not be repeated here. Noise error must be considered, as noise within the instrumentation may be significant enough that the atmospheric component of the variance is small compared to the overall measured variance. Noise is measured using the extrapolations of the measured autocovariance functions to lag 0 by the
- 250 structure function. This technique was introduced in Lenschow et al. (2000) to estimate the noise contribution from the second to fourth order moments. Although this technique was traditionally used to estimate lidar noise (Wulfmeyer, 1999; Wulfmeyer et al., 2010, 2016), it has also been extended to in-situ observations (Turner et al., 2014).

Figure 1, Panel (b) provides the autocovariance function of vertical velocity and q for a sub-cloud fight leg in RF3 (black). The fit using the structure function is provided in red (vertical velocity) and green (q). The structure function at lag zero provides the mean variance, while the difference between the autocovariance and structure function at lag zero provides the system noise variance at the corresponding temporal resolution. It is clear that the atmospheric variance and noise can be separated. For example, from Panel (b), looking at the vertical velocity data,  $\overline{w'w'} = 0.20 \text{ m}^2 \text{s}^{-2}$  and the noise variance  $\overline{\delta_w^2} = 0.014 \text{ m}^2 \text{s}^{-2}$ . This results in a noise standard deviation of  $\delta_w = 0.12 \text{ ms}^{-1}$ . Extending this analysis to determine the error propagation within higher order moments, error bars for vertical velocity variance (w'w') and q-variance (q'q'), vertical velocity skewness (w'w'w'), and the kinematic moisture flux (w'q') can be found in Panels (d) through (f), respectively, with noise error bars in red and sampling error bars in black. The noise error is negligible compared to the sampling error, in agreement with results from Turner et al. (2014). Note that some data points do not have noise error bars associated with them. This is due to the fact that the noise was so small, the error bars would not plotwere negligible. The various vertical profiles displayed show that the sampling errors result in a lack of statistical significance between flight legs of different altitudes.

Equations used to determine the noise in the higher order moments from Wulfmeyer et al. (2016) are:

$$\sigma_{w'^2} \cong 2\sqrt{\overline{w'^2}} \sqrt{\frac{\delta^2}{N}} \tag{7}$$

$$\sigma_{w'^3} \cong 3\sqrt{3w'^2} \sqrt{\frac{\delta^2}{N}} \tag{8}$$

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$$\sigma_{w'q'} \cong \sqrt{\overline{q'^2}} \frac{\delta_w^2}{N} + \overline{w'^2} \frac{\delta_q^2}{N} \tag{9}$$

where  $N_N$  is the number of data points. Using this Equation 7, the absolute error for the vertical velocity variance is found to be 0.00068 m<sup>2</sup>s<sup>-2</sup> and the relative error is 0.35% (the relative error for the *q* variance is 1.9%). Both errors are very reasonable, and demonstrate the low noise of the instrumentation.

#### 275 3 Synoptic and Boundary Layer Characteristics

#### 3.1 Mean Synoptic Conditions Variability at Point Alpha

The Southeast Pacific Ocean is found on the eastern edge of the south-Pacific semipermanent subtropical anticyclone, characterized by large scale upper tropospheric subsidence leading to a strong temperature inversion with a well-mixed boundary layer below. The surface pressure therefore is controlled in part by the location of the south-Pacific subtropical anticyclone. This

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0 anticyclone is routinely interrupted (especially between fall and spring) by periods of relative relatively low pressure which is associated with localized troughing or the passage of midlatitude cyclones to the south. Several papers (Toniazzo et al., 2011; Rahn and Garreaud, 2010a) have analyzed the synoptic characteristics during VOCALS-REx. These papershowever, however, tend to focus on the VOCALS-REx region as a whole and not specifically on Point Alpha, which is done in this section.

Figure ?? shows the mean of large-scale meteorological conditions, including Spatial maps of the mean sea level pressure

285 , omega ( $\omega$ , representing vertical velocity in pressure coordinates) and 700-hPa geopotential height from NCEP reanalysis data over the study region between Oct. 19<sup>th</sup> to Nov. 12<sup>th</sup>. The mean sea level pressure (panel (a )) displays (not shown here, see Zheng et al. (2011) or Toniazzo et al. (2011) for a visual) display the anticyclone near its climatological position of 30°S, 100°W(Toniazzo et al., 2011). The overlaid sea level pressure standard deviation (only displayed up to 20 hPa) shows variability increasing southward, indicating enhanced midlatitude storm tracks. Enhanced variability that is in line with

290 midlatitude troughing from panel (b) also decreases toward the coast, suggesting more variation in the synoptic pattern over the open ocean as compared to the near-coastal region. This is as expected, as Barret et al. (2009) found that synoptic systems tend to weaken as they move towards the coast of South America.

The mean 700-hPa geopotential height is displayed in panel (b), overlaid with  $\omega$  data. Subsidence (green shading) dominated the VOCALS-REx region, with Point Alpha having an average value of 57 hPa day<sup>-1</sup> at the 700-hPa level. While enhanced

295 storm. While enhanced strom tracks were primarily contained within the mid-latitudes, the standard deviation in the 700hPa geopotential height displays midlatitude troughing extending map display midlatitude troughing that extended between Point Alpha and the subtropical high (as was found in (Zheng et al., 2011))(Zheng et al., 2011), suggesting that meteorological conditions at Point Alpha were influenced by both midlatitude synoptic systems and the subtropical anticyclone.

The sea-level pressure was also measured using both reanalysis data and aircraft 30-m level horizontal flight legs. Figure 2, panel (a) shows that the reanalysis data at Point Alpha tended to be on average 1.5-hPa greater than the aircraft measured sea level pressure. The pressure decreased by roughly 3-hPa from Oct. 19<sup>th</sup> to Nov. 12<sup>th</sup>, however, this decrease cannot be considered a seasonal signal because it is within synoptic scale variation. The sea surface temperature (SST) and atmospheric surface temperature (both measured during 30-m horizontal flight legs) increased steadily throughout the observation period, increasing by 2.79 and 2.28 °C, respectively.

#### 305 3.2 Synoptic Variability at Point Alpha

Synoptic variability at Point Alpha is summarized in Figure 2 by time series of geopotential height at various levels 500 and 700-hPa. Higher geopotential heights are associated with ridging aloft while decreases in geopotential heights are associated with synoptic disturbances or troughs. The 500-hPa geopotential height (see Figure 2) varied between 5840 and 5900-m, with an increase of 9-m a decrease of 27-m between Oct.  $1916^{th}$  and Nov.  $1213^{th}$ . Figure 2 also displays enhanced synoptic scale variation during October, with several disturbances effecting Point Alpha. The 500, 700 (panel (c)), 850, and 1000 (panel (d))

310 variation during October, with several disturbances effecting Point Alpha. The 500 , 700 (panel (c)), 850, and 1000 (panel (d)) hPa-and 700-hPa geopotential heights alternate between areas of high and low height through Nov. 2<sup>nd</sup>. After Nov. 2<sup>nd</sup>, the 500-hPa geopotential height is more consistent, with height increasing over Point Alpha until Nov. 10<sup>th</sup>, at which point the height begins to decrease.

Besides minor disturbances in October, there are two main disturbances that stand out. The first disturbance occurs on Nov.

315 1<sup>st</sup> and 2<sup>nd</sup> (green shading in Figure 2), where both the 500 and 700-hPa heights have minimums a minimum (5842 and 3134 m, respectively) due to the influence of a synoptic system. The 850 and 1000-hPa heights also have secondary minimums. The second disturbance was the formation of a costal low, which can be seen by decreasing geopotential heights on Nov. 12<sup>th</sup>. Both the 850 and 1000-hPa geopotential heights reached minimums on Nov. 12<sup>th</sup> (1498 and 104 m, respectively)after Nov. 10<sup>th</sup>. This costal low reached a minimum (the coastal low was strongest) after the analysis period , on Nov. 15<sup>th</sup> (Rahn and 10<sup>th</sup>).

Garreaud, 2010a). The ridging which formed after Nov.  $2^{nd}$  leads to the formation of the coastal low through the warming of the lower and middle troposphere (Garreaud and Rutllant, 2003).

The 700-hPa geopotential height map (not shown here) displayed a midlatitude trough developing and extending past Point Alpha from Oct.  $29^{th}$  through Nov.  $3^{rd}$ . A deep midlatitude trough forms off the west coast of South America by Oct.  $30^{th}$ . extending past 15°S. The trough axis begins to move over Point Alpha by  $\frac{\text{October-Oct. } 31^{st}}{\text{Oct. } 31^{st}}$ , with the main impacts of the trough on Point Alpha (in terms of lowest geopotential height) being observed on Nov. 1<sup>st</sup> and 2<sup>nd</sup>. The 500-hPa geopotential height map (not shown here) shows the ridge axis directly over Point Alpha on Nov.  $1^{st}$ .

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Figure 3 (panels-Panels (a) through (c)) show atmospheric wind direction and velocity using data collected from horizontal flight legs. Panel Panels (d) and (e) displays display wind direction and wind speed using data collected from aircraft vertical soundings, repsectively. Atmospheric winds near the surface (measured during 30-m horizontal flight legs) at Point Alpha were mostly southerly (150 to  $\frac{180230^{\circ}}{1000}$ ) with a mean of  $\frac{176179^{\circ}}{1700}$ . Strong wind shear was present near the inversion, with winds 330 above the marine boundary layer (measured during horizontal flight legs above the inversion) having a mostly northwesterly component (mean of  $\frac{273276^{\circ}}{1000}$ ) while having more variability in direction than that of the boundary layer (300 to 360°). Although on most flight days the wind speed and direction were mostly constant with height in the throughout the depth of the boundary layer (see panel-Panels (d) and (e)), on Nov.  $1^{st}$  and  $4^{th}$  (blue lines) the wind direction shifted sharply within the boundary layer from southerly to northeasterly, along with varying wind speed. On Nov.  $2^{nd}$  (green line), the wind direction had its strongest a westerly component (214°). Shear within the boundary layer is not common. Zheng et al. (2011) suggest that this shear is linked to coastal processes such as the propagation of the upsidence wave. It should also be noted however

- 335 that the wind shear within the boundary layer is present on the same day (November Nov.  $1^{st}$ ) that the trough axis is located over point Point Alpha. On the proceeding day, the surface winds experience their second most westerly component -(where
- Oct. 24<sup>th</sup> has the most westerly component). According to Rahn and Garreaud (2010a), as troughs approach the coast of South 340 America, southeast winds are typically replaced by southwest winds, Between Oct.  $29^{th}$  and Nov.  $2^{nd}$ , wind direction within the boundary layer shows its most variation, gradually shifting from  $153^{\circ}$  (most easterly component measured) to  $\frac{213}{214^{\circ}}$ (second most westerly componentmeasured), respectively. While the trough approaches the coast of Chile, southeast winds are replaced by southwest winds, as is typical of synoptic scale disturbances in the region (Rahn and Garreaud, 2010a). Note that
- the most westerly component of the boundary layer wind measured on Oct. 24<sup>th</sup> coincides with a dip in the geopotential height 345 (with winds quickly shifting back to easterly), suggesting a weak disturbance on Oct.  $24^{th}$ .

#### 3.2 **Boundary Layer Characteristics**

Boundary layer height is perhaps the most important feature of the marine boundary layer (MBL), with  $z_i$  being one of the main dictators metrics for boundary layer characteristics such as decoupling and cloud cover (Albrecht et al., 1995). Findings 350 from Rahn and Garreaud (2010a) at a separate observation point within the VOCALS-REx region suggests that  $z_1$  tended to be either low (600-m) or high (1500-m) with periods of high or low depth interrupted by rapid transitions between the two states over 12 to 36 hour periods due to synoptic variability. Figure 4 shows the thickness of the Sc cloud layer, the thickness of the inversion layer, and subsequently the MBL height for each flight. The expected lifted condensation level (LCL) for a well mixed boundary layer is also provided, using  $z_{LCL} = 123(T - T_d)$ , where  $T_d$  is dew point temperature.  $z_i$  is also 355 provided from extrapolating relative humidity data from ECMWF reanalysis (Engeln and Teixeira, 2013). The cloud layer was

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identified using a liquid water content (LWC) greater than or equal to 0.01 g m<sup>-3</sup>, while the inversion layer was identified by the region of greatest change in q (absolute change > 0.10 g kg<sup>-1</sup> per 1-Hz measurement) and  $\theta$  (absolute change > 0.20 K per 1-Hz measurement) within the vertical profiles. This results in the bottom of the inversion layer characterized by the profiles beginning to lose the boundary layer features, while the top of the inversion layer had lost all boundary layer features.

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The average  $z_i$  was  $\frac{1175 \text{-m}}{1148 \text{-m}}$  (see Table 2 for boundary layer characteristics), with the average cloud layer and inversion thickness being  $\frac{239 \text{ and } 59 \text{ m} 229}{239 \text{ and } 55 \text{ m}}$ , respectively. Figure 4 shows that  $z_i$  varied between 996 and 1450-m, with mostly gradual changes in height from flight day to flight day (note that the mean difference between  $z_i$  and ECMWF- $z_i$ ) was  $\frac{43.44}{26}$  +  $\frac{26}{26}$  +  $\frac{26}{26}$  +  $\frac{26}{2}$  +  $\frac{26}{$ of a rate of change above 100 m day<sup>-1</sup>. After Oct. 27<sup>th</sup> is when For cloud thickness, the most significant changes took place

- to the cloud thickness and  $z_7$ . Between Oct. after October  $27^{th}$  and  $29^{th}$ ,  $z_7$  increased from 995 to 1300-m (152 m day<sup>-1</sup>, the 365 second largest rate of change), where the ECMWF-z, shows that the increase was mostly confined from Oct. 27<sup>th</sup> to Oct. 28<sup>th</sup>. The next four flight days recorded the thickest cloud layers, peaking on Nov.  $1^{st}$  and  $2^{nd}$  with thicknesses of 382 and 472-m, respectively. It should also be noted that between Oct.  $29^{th}$  and  $30^{th}$ ,  $z_t$  decreased from 1300 to 1177-m (124 m day<sup>-1</sup>, the third largest rate of change, although this is not conveyed in the ECMWF-z, data). After Nov, 2<sup>nd</sup>, the cloud layer thinned
- and z; increased from 1136-m to 1450-m between November  $4^{th}$  and Nov.  $8^{th}$ . Although this is a rate of 79 m day<sup>-1</sup>, there 370 is no in situ data in-between November  $4^{th}$  and  $8^{th}$ . The ECMWF- $z_i$  provides a mean rate of change for this period of 93 m  $day^{-1}$ , with the largest change of 160 m  $day^{-1}$  between November 7<sup>th</sup> and 8<sup>th</sup>, suggesting a rapid rise in z<sub>i</sub>, in concurrence from findings in Rahn and Garreaud (2010a). After  $z_i$  peaks on November  $8^{th}$ ,  $z_i$  falls rapidly over the next two days, showing decreases of 174 m day<sup>-1</sup> and 102 m day<sup>-1</sup> from Nov. 8<sup>th</sup> to Nov. 10<sup>th</sup>, respectively.
- Although the Although the time series of cloud droplet number concentration  $(N_d)$  is not shown here, it showed we observed 375 a notable dip to a minimum on Nov.  $1^{st}$  of 81 cm<sup>-3</sup> (where the average is  $\frac{292-280}{292-280}$  cm<sup>-3</sup>), corresponding with minimums-a minimum in both boundary layer cloud condensation nuclei (CCN) and aerosol number concentration . Above boundary layer aerosol number concentration had a maximum on Nov. 1st. However, this can most likely be attributed to enhanced moisture (see Figure 5) above the boundary layer due to the passing synoptic system, where enhanced moisture can increase the sizeof
- hygroscopic acrosols that would otherwise be too small to be measured under dryer conditions ( $N_a$ ), along with a maximum in 380 average drop size.

Figure 5 shows vertical profiles (where the height (z) is normalized with by the inversion height to give a non-dimensional vertical coordinate of  $z/z_i$ ) of  $\theta$ , q, LWC, and the aerosol number concentration N<sub>q</sub>. Individual flight profiles are in gray, with the red profile representing the mean and the blue profiles representing the flights conducted on November  $1^{st}$  (RF11, light blue) and Nov. 2<sup>nd</sup> (RF12, dark blue). Mean profiles show that on average the MBL is well mixed up to the inversion, which

385

then prevents mixing into the free atmosphere above (as evident by the decrease in aerosol number concentration between the boundary layer and free atmosphere).

The largest deviations from the mean in the profiles occur during the passage of the synoptic system on Nov.  $1^{st}$  and  $2^{nd}$ . At this time, both RF11 and RF12 measured (1) The the thickest Sc cloud layer, with Nov.  $1^{st}$  having the largest average cloud droplet size (20.8  $\mu$ m) and in-cloud drizzle rates, while November 2<sup>nd</sup> had the lowest recorded cloud base and largest recorded 390

LWC; (2) A-a larger mixing ratio above the boundary layer. This suggests the presence of a moist layer aloft which may have helped in producing to produce the thickest cloud layers observed; (3) The the smallest differences in both  $\theta$  and q from the bottom to the top of the inversion layer. During the passage of strong events as described by Rahn and Garreaud (2010a), the inversion defining the MBL erodes, making it hard to define  $z_i$ . This process is partially displayed by the small differences in temperature and moisture across the inversion layer during the passage of the synoptic disturbance.

The differences in q and  $\theta$  can be better visualized in Figure 6, which shows the differences between below and above inversion values in panel (a).  $z/z_i$  values between 0.85 and 0.95 were used for the averages below the inversion, while data between  $z/z_i$  values of 1.10 and 1.20 were used for the averages above the inversion. Besides Nov. 1<sup>st</sup>, 2<sup>nd</sup>, and to a lesser degree Nov. 4<sup>th</sup>, the average difference in  $\theta$  across the inversion was 17-K, while the average difference in q was -6.2 g kg<sup>-1</sup>.

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400 On Nov.  $1^{st}$  when both reached a minimum difference, the difference between q and  $\theta$  across the inversion was 1.9 g kg<sup>-1</sup> and 14-K, respectively, where a weaker inversion allows for more entrainment mixing near cloud top (Galewsky, 2018).

There are multiple methods which can be used to To analyze whether the boundary layer is well mixed or decoupled. Methods usedhere include well-mixed or decoupled, two methods are used: (1) decoupling parameters and (2) analysis of the expected LCL for a well-mixed layer in relation to actual cloud base. Decoupling parameters  $\alpha_{\theta}$  and  $\alpha_{q}$  depend on the profiles of  $\theta$  and q, respectively (Wood and Bretherton, 2004). The decoupling parameters measure the relative difference in q and  $\theta$  between the bottom (near the surface) and top (near the inversion) portions of the boundary layer, and are given by:

$$\alpha_{\theta} = \frac{\theta(z_i^-) - \theta(0)}{\theta(z_i^+) - \theta(0)} \tag{10}$$

$$\alpha_q = \frac{(z_i^-) - q(0)}{q(z_i^+) - q(0)},\tag{11}$$

where  $z_i^+$  ( $z_i^-$ ) is the level ~25 m above (below)  $z_i$ , and  $\theta(0)$  and q(0) are the potential temperature and mixing ratio at the surface. Here,  $z_i^+$  is calculated using data between  $z/z_i$  values of 1.03 to 1.05, while  $z_i^-$  is calculated using data between  $z/z_i$ values of 0.95 to 0.97 (this is roughly 25 m above and below  $z_i$ , respectively). The closer to zero the decoupling parameters are, the more well-mixed the boundary layer is. Previous observations suggest that if the parameters exceed ~ 0.30, the boundary layer is decoupled (Albrecht et al., 1995).

Mixed layer cloud thickness represents the difference between z<sub>i</sub> and the LCL (Δz<sub>m</sub>), and was found to be strongly correlated to decoupling in Jones et al. (2011). The difference between cloud base (z<sub>b</sub>) and the LCL represents another decoupling index (Δz<sub>b</sub>) related to the LCL presented in Jones et al. (2011). Decoupling of the boundary layer occurs when the boundary layer deepens, resulting in a larger difference between the inversion and the LCL as the LCL diverges from cloud base. A well-mixed boundary layer would have z<sub>b</sub> and LCL measurements which are in close agreement, while a decoupled boundary layer would have a divergence in the similarities between the two values. Previous observations within the VOCALS-REx
domain from Jones et al. (2011) found that the boundary layer tended to be decoupled if Δz<sub>b</sub> > 150-m and if Δz<sub>m</sub> > 500-m.

Figure 6 shows the decoupling parameters in panel (b). The average value of  $\alpha_{\theta}$  ( $\alpha_{q}$ ) are 0.15 (0.07 is 0.14 (0.08), both of which are within the regime of well mixed. During RF11 and RF12, q increases above the inversion leading to large values for  $\alpha_{q}$ , while  $\Delta\theta$  is relatively small as compared to other flights, with  $\alpha_{\theta}$  being above 0.30 during November 1<sup>st</sup>, (where Zheng

et al. (2011) suggest drizzle processes act to stabilize the boundary layer, leading to decoupling). Panel (c) provides values for

425  $\Delta z_b$  and  $\Delta z_m$  for each flight, with average values of  $\frac{125 \text{ and } 363 \text{-m}111 \text{ and } 340 \text{-m}}{125 \text{ m} 111 \text{ and } 340 \text{-m}}$ , respectively. Again, both values are within the regime of well mixed regime.

RF11,13, and 15 are shown to be decoupled, with both  $\Delta z_b$  and  $\Delta z_m$  at or above the 150 and 500-m threshold values, respectively. RF12 is decoupled according to  $\Delta z_m$  only and RF6 and 16 are decoupled according to  $\Delta z_m$  only. Looking at raw profiles of q and  $\theta$  (not shown here), RF11, RF6, 11, 12, 13, and 15 appear to be decoupled due to distinct humidity changes

430 within the sub-cloud profiles, including the presence of a cumulus layer below the Sc deck that is visible from analyzing the LWC profiles (not displayed here) during RF11 (Nov. 8<sup>th</sup>). This results in By these metrics 28% of profiles analyzed being the profiles analyzed are decoupled.

The comparison between Panels (b) and (c) demonstrate that determining decoupling using  $\Delta z_b$  and  $\Delta z_m$  appears to be more accurate than the decoupling parameters when comparing the results to the raw vertical profiles. A more accurate value for determining decoupling using  $\alpha_{\theta}$  and  $\alpha_q$  for the data presented here is 0.20, as compared to the 0.30 stated in Albrecht et al. (1995). A value of 0.20 would lead to better agreement between the two methods. Note that the correlation between  $\Delta z_b$  and  $\Delta z_m$  is 0.79-0.76 (i.e., when the mixed layer cloud thickness increases, the difference between the LCL and cloud base increases). This suggests that when the boundary layer deepens, the cloud layer thickness remains relatively consistent, in agreement with findings from Jones et al. (2011).

#### 440 **4 Results**

Here, we will quantify the amount of turbulence occurring within the boundary layer. In particular, analysis includes: (1)
Analyze analyze day to day variability in turbulent measurements and boundary layer characteristics, relating them to synoptic changes in meteorological conditions; (2) Determine determine average turbulent values throughout the vertical structure of the STBL, classifying the STBL based on different turbulent profiles analyzed. For each flight analyzed here, the Sc deck lies directly below a strong inversion. This It should be noted that this extreme vertical gradient can cause instrument response issues with the measurement of both the dry bulb and dew point temperature for some distance beneath cloud top (Nicholls and Leighton, 1986). Therefore, data collected during both vertical profiles and horizontal legs will be used and compared.

#### 4.1 Synoptic Variability of Turbulence

- Figure 7 shows the mean surface (30-m horizontal flight leg) LHF (panel Panel (a)), SHF (panel Panel (b)), and Bowen ratio (panel TKE and  $\epsilon$  (Panel (c)) for each flight day with the standard deviation represented by the gray envelopeshaded envelopes. Note that for days with two or more mean values, there were two or more 30-m horizontal flight legs, with good agreement between mean leg values within the same flight. The LHF peaks on Oct.  $26^{th}$  with a value of 53.3-50.7 W m<sup>-2</sup>, and from that point decreases steadily to its minimum values of 19.7 and 18.5-19.2 and 18.4 W m<sup>-2</sup> just as and after the minimum in geopotential height on Nov.  $2^{nd}$  and  $4^{th}$ , respectively. The SHF has a sharp increase to its maximum value of 17.1-17.0 W m<sup>-2</sup>
- 455 on Nov.  $1^{st}$  and decreases to its secondary minimum a below average value of 5.4 W m<sup>-2</sup> on Nov.  $2^{nd}$  (note that mean values

of surface fluxes can be found in Table 3). The Bowen ratio is typically small (less than 0.20), especially for the first half of the eampaign. The Bowen ratio has a sharp increase on Nov. 1<sup>st</sup> to match the increase in the SHF (and remains above 0.20 for the remainder of the analysis period), suggesting that the liquid water flux in the cloud layer should not be taken to be proportional to the upward LHF after Nov. 1<sup>st</sup>. Note that the average surface values of the LHF and SHF are generally in agreement with

460 those found in Zheng et al. (2011), who found values of 48.5 and 7.1 W m<sup>-2</sup>, respectively. The differences most likely arise due to different averaging techniques.

Figure ?? gives the surface friction velocity (vertical transport of horizontal momentum), vertical velocity variance, TKE, Surface TKE and  $\epsilon$  in Panels (a) through (d), respectively. One commonality between each parameter is that the maximum value is reached both reach a maximum on Nov. 1<sup>st</sup> followed by the minimum value both reaching a minimum on Nov. 2<sup>nd</sup>

- 465 (see Table 3 for the mean and range of the values). For all four variables, there is Both TKE and  $\epsilon$  show very little variation between measurements, except for between Oct.  $30^{th}$  and Nov.  $2^{nd}$ , where turbulence shows a large increase in turbulence is observed before followed by a rapid decrease. Overall, there is good agreement between mean values for within the same flight, with the exception of Nov.  $12^{th}$ , which contains the largest difference between mean values for each variable in discussion here. This large difference was not observed however for the surface LHF and for TKE and  $\epsilon$  and Nov.  $13^{th}$  for the SHF.
- 470 Shifting focus to the entire depth of the boundary layer, Figure 8 shows boxplots (made up of leg mean values) of subcloud (white) and in-cloud (blue) values of LHF  $\underline{F}_{q_e}$  (Panel (a)) and buoyancy flux  $\underline{F}_{\theta_{u_e}}$  (Panel ((c)). Panels (b) and (d) display histograms of the LHF and buoyancy flux  $\underline{F}_{q}$  and  $\underline{F}_{\theta_{u_e}}$  data with normal distribution fits for reference, respectively. The overall LHF was 11.03  $\underline{F}_{q}$  was 10.63  $\pm$  12.97 3.66 Wm<sup>-2</sup>, with the a sub-cloud mean being 15.74 of 16.43  $\pm$  16.4 6.84 Wm<sup>-2</sup> and the an in-cloud mean being 6.01 of 5.0  $\pm$  3.75 3.05 Wm<sup>-2</sup>. The sub-cloud LHF  $\underline{F}_{q_e}$  is clearly offset to larger values, owing
- 475 to surface evaporation and subsequent transport of moisture. The red dots in Panel (a) represent the surface values, which are always the largest within the entirety of the vertical layer. The lowest mean values occurred on the same days as the minimum in geopotential height, Nov. 1<sup>st</sup> and 2<sup>nd</sup>, with values of 5.51 and 4.67 5.60 and 4.68 Wm<sup>-2</sup>, respectively. Although Statistically speaking these two data sets are visually different, statistically speaking they are similar, with a p-value of 0.22 (sub-cloud vs. in-cloud) are statistically different, with a p-value of 2.7E-6 (note that all statistical significance testing will be are carried out
- 480 using the Wilcoxon-Sum-Rank-Test). Removing the surface 30-m horizontal flight leg data however results in the these two data sets being statistically similr, with a *p*-value of 0.21.

The buoyancy flux As demonstrated in Panel (c)displays that, the overall mean buoyancy flux was  $4.89 \pm 4.86$  F<sub> $\theta_{w}$ </sub> was  $5.84 \pm 2.86$  Wm<sup>-2</sup>, with the a sub-cloud mean being 4.64 value of  $5.10 \pm 3.94 \pm 2.99$  Wm<sup>-2</sup> and the an in-cloud being 5.12 value of  $5.99 \pm 5.64 \pm 0.03$  Wm<sup>-2</sup>. From just analyzing the mean values of flight legs Based on the mean flight leg values,

- there does not appear to be a large difference in the buoyancy flux  $\underline{F}_{\theta_{u}}$  between the sub-cloud and in-cloud sections of the boundary layer, which is not as expected. In-cloud buoyancy in general is enhanced due to latent heating and cooling effects. There is no statistical significance statistically significant difference between the in-cloud and sub-cloud data, with a p-value of 0.39p-value of 0.43. While the medians in the data populations are similar, the buoyancy flux  $\underline{F}_{\theta_u}$  in-cloud has a clear increase in variance and a much larger range , suggesting (-13.2 - 38.1 Wm<sup>-2</sup>), suggesting more variation and isolated occurrences of
- 490 extremely large buoyancy fluxes or small  $F_{\theta_u}$  within the cloud as compared to sub-cloud (range of -1.7 19.5). Connecting

back to concepts discussed in the introduction, the coefficient correlation coefficient between the surface LHF and the in-cloud buoyancy is 0.40, suggesting  $F_{\theta_{u}}$  is 0.34, providing some evidence that a larger surface LHF leads to a larger in-cloud buoyancy flux, as suggested by Bretherton and Wyant (1997) and Lewellen et al. (1996).

- Figure 9 provides the same format as that of Figure 8, except for TKE (Panel (a)) and  $\epsilon$  (Panel(c)). The total mean TKE was  $0.132-0.126 \pm 0.03 \text{ m}^2\text{s}^{-2}$ , with a sub-cloud mean of  $0.133-0.127 \pm 0.05-0.04 \text{ m}^2\text{s}^{-2}$  and an in-cloud mean of  $0.132 - 0.124 \pm 0.04 - 0.035 \text{ m}^2\text{s}^{-2}$ . The total mean  $\epsilon$  was  $3.97-3.74 \pm 1.28-1.34 \text{ cm}^2\text{s}^{-3}$ , with a sub-cloud mean of  $4.14-3.87 \pm 2.45 - 1.58 \text{ cm}^2\text{s}^{-3}$  and an in-cloud mean of  $3.80-3.50 \pm 1.81-1.40 \text{ cm}^2\text{s}^{-3}$ . Overall, very consistent values (when looking at the means) between sub-cloud and in-cloud existance observed, resulting in statistical similarity between the data populations for both TKE and  $\epsilon$ . However, in looking at the boxplots, one can see that there are several cases (including Oct.  $24^{th}$ , Nov.  $1^{st}$  and
- Nov.  $2^{nd}$ ) where the entire turbulent distribution of the sub-cloud data is shifted to larger values than those of in-cloud data, with minimal overlap. This implies that the two layers have limited mixing between them, perhaps due to a more turbulent decoupled lower boundary layer. This will be explored in further detail in Section 4.2. Along with having different turbulent distributions between in-cloud and sub-cloud, both the TKE and the  $\epsilon$  had maximum average values on Nov  $1^{st}$  (0.163 m<sup>2</sup>s<sup>-2</sup> and 6.13 cm<sup>2</sup>s<sup>-3</sup>, respectively) and minimum average values on Nov.  $2^{nd}$  (0.065 m<sup>2</sup>s<sup>-2</sup> and 1.30 cm<sup>2</sup>s<sup>-3</sup>, respectively).
- The analysis to this point clearly shows a maximum in turbulent properties on Nov. 1<sup>st</sup> and a minimum on Nov. 2<sup>nd</sup>. This maximum is driven from by turbulence below the cloud however, with the in-cloud TKE (0.128-0.113 m<sup>2</sup>s<sup>-2</sup>) and  $\epsilon$  (2.78-2.66 cm<sup>2</sup>s<sup>-3</sup>) being below normal for in-cloud values, where the normal is 0.129-0.124 m<sup>2</sup>s<sup>-2</sup> and 3.68-3.50 cm<sup>2</sup>s<sup>-3</sup>, respectively. Panel (d) shows the total  $\epsilon$  distribution for in-cloud and sub-cloud. An increase in in-cloud frequency for  $\epsilon$  is clear for the lowest values (first two histogram bars). Eight of the 15 measurements from the first two histogram bars came from RF11
- 510 and RF12 (From Section 3.2, it is known that each of the three cases outlined above (Oct. 24<sup>th</sup>, Nov. 1<sup>st</sup> and Nov. 2<sup>nd</sup>), which includes all are decoupled. Although Nov. 4<sup>th</sup> and 8<sup>th</sup> are also decoupled, these boxplot profiles do not show the same shifted turbulent distributions between the sub-cloud and in-cloud values for those flights. The other seven measurements were all sampled above a normalized boundary layer height of 0.90, suggesting entrainment mixing of more laminar flow near the top of the entrainment layer into the upper cloud layer, reducing the turbulent energylayers. However, Nov. 4<sup>th</sup> and 8<sup>th</sup> do
- 515 have lower than average turbulence (all five of the decoupled boundary layers have lower than average turbulence, except for Nov.  $1^{st}$ , which will be looked at closer in Section 4.2). There is a strong negative correlation between the variables used to determine decoupling in Section 3.2 and the in-cloud turbulence (bottom portion of Table 4), displaying that more decoupled boundary layers correspond to less in-cloud turbulence. The decoupling variables also have a strong positive correlation with  $z_i$ , indicating an increase in  $z_i$  leads to an increased chance of boundary layer decoupling. Negative correlations between
- 520 in-cloud turbulence and  $z_i$  reinforce this idea. The increase in  $z_i$  is accompanied by a decrease in turbulence within the cloud layer. As  $z_i$  increases, cooling near cloud top cannot sustain mixing over the entire depth of the boundary layer, resulting in less turbulence and boundary layer decoupling (Bretherton and Wyant, 1997).

It Along with analyzing correlations between  $z_i$ , turbulence, and decoupling, it is important to analyze other turbulent fluxes of energy, momentum, and moisture as they act to determine boundary layer structure and characteristics, along with analyzing how these variables are related to synoptic scale properties such as geopotential height. The correlation coefficients between boundary layer characteristics and synoptic scale properties can be found in the top portion of Table 4. The 700hPa geopotential height (i.e., pressure) is fairly strongly correlated with  $z_i$ , although this correlation is negative with a value of -0.37-0.49, suggesting that as the pressure increases,  $z_i$  decreases. The rate of change in  $z_i$  can be is governed by the entrainment rate ( $\omega_e$ ) and <u>omega ( $\omega$ -, i.e., vertical velocity in pressure coordinates</u>). If the rate of subsidence increases to the

530 point that it is larger than  $\omega_e$ , then  $z_i$  will decrease with time.  $\omega$  depends primarily on synoptic scale patterns, in particular that of geopotential height. Pressure and  $\omega$  have a correlation of -0.89-0.71, suggesting that as pressure increases, the subsidence increases (or at the very least, upward vertical motion is diminished). Entrainment on the other hand, can depend on multiple variables including the inversion layer thickness, wind shear, and surface fluxes. Increases in  $\omega_e$  result in a higher LCL for the entrained air and a resulting increase in boundary layer height as a result. Given that  $z_i$  acts to decrease as the pressure increases, 535 this suggests that the subsidence becomes the dominating component that governs  $z_i$  over that of rather than entrainment.

this suggests that the subsidence becomes the dominating component that governs  $z_i$  over that of rather than entrainment. The surface LHF provides the main source of moisture in the STBL, which in turn is an important source of buoyant TKE production. An enhanced (reduced) LHF will generate thicker (thinner) clouds with larger (smaller) LWC values, resulting

- in enhanced (reduced) evaporative cooling near cloud top leading to enhanced (reduced) buoyancy driven entrainment, and a subsequent deepening (thinning) of the boundary layer. This process is demonstrated well when analyzing the correlation coefficients. Both the LHF and SHF are positively correlated with  $z_i$  (correlation coefficients of 0.36 and 0.440.31 and 0.43, respectively) while the LHF is negatively correlated with the Sc cloud thickness (correlation coefficients of -0.50-0.34). Therefore, a larger LHF tends to result in a thinner Sc cloud layer but a larger  $z_i$ , suggesting enhanced entrainment which acts to thin the cloud layer while deepening the boundary layer. It should also be noted that the correlation between the SHF and wind speed is significant, as anticipated, since the SHF is expected to increase-increase linearly with wind speed (Palm et al., 1999).
- 545 Both TKE and  $\epsilon As N_d$  and  $N_a$  increase in-cloud with respect to pressure (correlation coefficient of 0.23 and 0.24, respectively) and decreases with respect to  $z_i$  (-0.32 and -0.34, respectively). The observed decrease in boundary layer turbulence with increasing  $z_i$  is due to decoupling and an inability for the entire boundary layer to be mixed (leading to a subsequent decrease in turbulence), while a shallow boundary layer can be easily mixed through cooling near cloud top.
- As the cloud droplet number concentration and aerosol number concentration increase (accompanied by a decrease in average droplet size), the TKE and  $\epsilon$  increase (with correlation coefficients of 0.35, 0.42, and -0.32 0.54, and -0.21 in relation to TKE, respectively). Physically this makes sense. As precipitation is suppressed due to larger number concentrations and smaller droplet sizes, a reduced moisture loss from the STBL can result, leading to thicker clouds, a larger buoyancy flux, and a larger TKE. Smaller droplets will also evaporate more readily, leading to enhanced latent heating effects in-cloud latent heating (i.e., absorption of energy through evaporation) and a resultant increase in turbulence through the bouyancy flux.

#### 555 4.2 Vertical Profiles

It has been shown through the boundary layer vertical structure in Figure 5 that the boundary layer is, on average, well mixed when considering thermodynamic variables. Figure ?? represents vertical profiles of the buoyancy flux (Panel (a)), LHF (Panel (b))Data collected during aircraft soundings (as opposed to mean values of horizontal flight legs), *u*-variance (u'u'), *v*-variance (y'v'), vertical velocity variance (Panel(c)), and TKE w'w'), and the TKE are displayed in Figure 10 Panels (a) through (Panel

- (d)), where each dot represents a leg mean value, with in-cloud values in red and values measured during, respectively, with the red line representing the mean profile and each gray line representing individual flight profiles. The blue lines represent flight profiles for Nov. 1<sup>st</sup> (light blue) and Nov. 2<sup>nd</sup> in blue. The buoyancy flux in the sub-cloud layer (on average) varied between -2 and 20 Wm<sup>-2</sup> and decreased with height until increasing within the cloud layer with values ranging between -5 and 43 Wm<sup>-2</sup>. The standard deviation (in orange) was produced using data from vertical flight profiles as opposed to the horizontal legs due
- to data uniformity throughout the boundary layer depth. The buoyancy flux has a clear increase in variance within the cloud layer. The LHF peaks near the surface, ranging between -1 and 55 Wm<sup>-2</sup> below the cloud layer and generally decreases with height. The variance peak of 33 Wm<sup>-2</sup> occurs (dark blue). Panel (e) displays the mean values from each of Panels (a) through (d), along with the anisotropy ratio  $\left(\frac{2w'w'}{w'a'+w'w'}\right)$ . The profile of each variable in question shows a near constant value below cloud base, with an increase in-cloud before beginning to decrease near cloud top. Both w'w' and TKE reach their peak values
- 570 at  $z/z_i = 0.99$ , signifying the large gradient in q near  $z_i$  and the variation in evaporative cooling due to entrainment mixing at cloud top between flight days.

Vertical velocity variance (from here on w'w')ranged from 0.008 to 0.20 m<sup>2</sup>s<sup>-2</sup>. The observed 0.89 (or a normalized incloud location of 0.43). TKE values plummet above the inversion due to the dominance of clear, stable, and subsiding air aloft. The overall maximum in TKE measured (for all 18 flights) is found near z/zi = 0.60 (looking at the light blue profile line in Figure 10, Panel (d)) during RF11 (Nov. 1<sup>st</sup>). This will be discussed in more detail in Section 4.3.

The observed average in-cloud w'w' at Point Alpha was  $\frac{0.105 \cdot 0.127 \pm 0.051}{10.027 \pm 0.051} \text{ m}^2 \text{s}^{-2}$  with values fluctuating considerably more than those in the sub-cloud layer ( $\frac{0.091 \pm 0.025 \text{ m}^2 \text{s}^{-2}}{10.025 \text{ m}^2 \text{s}^{-2}}$ ), in agreement with findings from Bretherton et al. (2010), who measured a larger standard deviation in vertical velocity in-cloud vs. sub-cloud). The average in-cloud value of w'w' found here is significantly lower than what was found over more remote ocean areas ( $80^{\circ}\text{W} - 85^{\circ}\text{W}$ ,  $20^{\circ}\text{S}$ ) of  $0.36 \text{ m}^2\text{s}^{-2}$ 

- (Bretherton et al., 2010). Nocturnal measurements of the Californian Sc deck during DYCOMS-II also revealed a stronger turbulent structure than that measured at Point Alpha, with observations showing in-cloud w'w' larger than 0.4 m<sup>2</sup>s<sup>-2</sup> with a maximum of 0.5 m<sup>2</sup>s<sup>-2</sup> near the base of the Sc deck (Stevens et al., 2005). As discussed in Wood (2012), w'w' is typically more vigorous at night due to the buoyancy production being larger from the lack of shortwave radiation absorption, which acts to stabilize the layer. As is As is found here, Hignett (1991); Nicholls (1984); Ghate et al. (2014) also found that w'w' is peaked in the upper half of the STBL away from any boundaries such as cloud top. Note that the TKE mirrors that of w'w' in
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terms of vertical spatial tendencies.

Considering data collected during aircraft soundings (as opposed to mean values of horizontal flight legs)Overall, the flow is not isotropic (anisotropy ratio is equal to one for isotropic flow, where vertical turbulence dominates for values greater than one and horizontal turbulence dominates for values less than one). Vertical turbulence has its largest component near the

590 surface (z/z) = 0.11, u-variance (u'u'), v-variance (v'v'), w'w', and the TKE are displayed in Figure 10 Panels (a)through (d), respectively, with the red line representing the mean profile and each gray line representing individual flight profiles. The blue lines represent flight profiles for Nov. 1<sup>st</sup> and Nov. 2<sup>nd</sup>. Panel (e) displays the mean values from each of Panels (a) through (d). The profile of each variable in question shows a near constant value below cloud base, with an increase in-cloud before beginning to decrease near cloud top. Both w'w' and TKE reach their peak values while having a secondary peak in-cloud in 595 accordance with the peak in TKE at  $z/z_i = 0.88$  (or a normalized in-cloud location of 0.40). 0.89.

Simulations and observations from Pasquier and Jonas (1998) of in-cloud TKE showed that the maximum TKE occurred in two locations, near cloud top and near cloud base, suggesting that turbulence is being generated through two processes: (1) Cooling cooling at or near cloud top (through evaporation or longwave cooling), resulting in cool, dry downdrafts; (2) Warming warming near cloud base from the release of latent heat through condensation, resulting in positively buoyant up-

- 600 drafts. However, no conclusions can be made here on whether or not there are two sources of TKE due to the low vertical resolution of the mean values (i.e., averaging over 14 flight profiles). TKE values plummet above the inversion due to the dominance of clear, stable, and subsiding air aloft. The overall maximum in TKE measured (for all 14 flights) is found near z/zi = 0.60 (looking at the blue profile line in Figure 10, Panel (d)) during RF11 (Nov. 1<sup>st</sup>). This will be discussed in more detail in Section 4.3.
- 605 Looking at individual profiles of TKE, (not shown here), only six of the fourteen eight of the eighteen flights have a maximum TKE within the cloud layer. Modeling and observations of boundary layer profiles of turbulence from Pasquier and Jonas (1998) showed that mixing and overturning of the boundary layer profile due to buoyancy effects leads to a maximum in turbulence commonly being reached in the sub-cloud layer. Seven of the fourteen Ten of the eighteen flights display two peaks in TKE within the cloud layer, one near cloud base and another near cloud top, signifying evaporative cooling near cloud top and latent
- 610 heating near cloud base. Of the six eight flights that have a maximum TKE within the cloud layer, all six eight display two peaks in the TKE within the cloud layer, one near cloud base and one near cloud top. Having the maximum in TKE in the sub-cloud layer can signify decoupling (Durand and Bourcy, 2001). A slight decoupling can lead to less moisture transport into the Sc layer, resulting in less latent heat release due to condensation. This could be why only one flight has two two flights have two peaks in TKE within the cloud when the turbulence maximum is reached below cloud, due to latent heat release at
- 615 cloud base being suppressed. All decoupled flights identified in Section 3.2 (with the exception of Nov. 1<sup>st</sup> and Nov. 2<sup>nd</sup>) have a single peak in TKE in the cloud layer, with the maximum TKE value being reached within the sub-cloud layer.

Figure 11 provides the same format as Figure 10, except for values of  $\frac{1}{2}$  (Panel (a)),  $\frac{1}{2}$  (Panel (b)),  $\frac{1}{2}$  (Panel (b)),  $\frac{1}{2}$  (Panel (c)), and the cloud droplet number flux (Panel (d)). Note that Figure 11 displays the range of data in the gray envelope, as opposed to showing each individual profile with a single gray line. Both the buoyancy flux

- 620 and the droplet number concentration flux (from here on w'N') have maximum values  $F_{\theta_m}$  has a maximum value at  $z/z_i =$  0.93 (normalized in-cloud height of 0.59). The peak near cloud middle is due to a combination of the warm/moist updrafts and cool/dry downdrafts meeting, formed by evaporative cooling at cloud top and latent heating near cloud base. The same concept can be extended to w'N', where droplets are activating near cloud base while evaporating near cloud top, suggesting that in the lower cloud the cloud droplet number concentration increases with updrafts (condensation), while the cloud droplet
- 625 number concentration decreases with downdrafts (evaporation) in the upper cloud region. According to Pasquier and Jonas (1998), the buoyancy flux  $F_{g_{u}}$  should reach a minimum near cloud top from the entrainment of warm, dry air down into the cloud layer. Although the mean profile does not show a decrease at cloud top, the raw data (i.e., unsmoothed) does show a negative buoyancy flux at cloud top. For individual flights, only RF11 (Nov 1<sup>st</sup>) had a maximum in the buoyancy flux  $F_{g_u}$  in

the sub-cloud layer. The LHF-F<sub>q</sub> peaks at the surface, but also sees a secondary maximum at  $z/z_i = 0.99$ . The maximum at cloud top can be attributed to the strong q gradient and to entrainment of drier air down into the cloud (i.e., also a positive flux

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since both w' and q' are negative).

Well-mixed STBLs tend to show characteristics of downdrafts that are spatially smaller, but stronger, than updrafts. This results in a negative vertical velocity skewness (from here on w'w'w') through most of the cloud and sub-cloud layer (Nicholls, 1989; Hogan et al., 2009; Ghate et al., 2014). Panel (c) displays indicates that w'w'w' on average is negative throughout the cloud layer and through most of the sub-cloud layer, having a maximum value near the surface. The minimum values value in w'w'w' occurs at cloud base (normalized in-cloud value of 0.04), suggesting that overall, the downdrafts are spatially smallest,

yet strongest at cloud base while updrafts are spatially larger, yet weaker. largest, but weakest.

Figure ?? shows the average buoyancy flux (Panel (a)), LHF (Panel (b)), TKE (Panel (c)), and c (Panel (d)) averaged over all flights for sub-cloud and for different layers within the cloud. Each black dot represents the average value for individual
flights using horizontal leg averages. The blue dots represent mean values using horizontal flight legs, while the red dots represent mean values using flight vertical profile data. Values for in-cloud are calculated for layers between normalized in-cloud height values of 0-0.25 (cloud base), 0.25-0.50 (bottom-middle), 0.50-0.75 (top-middle), and 0.75-1.0 (cloud top). Table ?? summarizes the mean values for each layer. A clear difference in values and trends can be seen between sampling methods. For example, looking at the LHF in Panel (b), we see that the horizontal leg sampling correctly captures the larger

- 645 LHF at the surface due to evaporation from the ocean surface, whereas the profile samples do not capture this increase at the surface (the profile data is terminated at the start of the 30-m horizontal flight legs, meaning there is limited samples near the surface for the profile method). Conversely, the profile method observed a large increase in TKE at cloud top from evaporative cooling due to entrainment mixing, which is not observed in the horizontal leg method. Another example is the buoyancy flux, which is seen to have a large increase in-cloud as compared to sub-cloud using the profile method. The horizontal leg method
- 650 displays a maximum in the top-middle region of the cloud, but the overall buoyancy flux increase in-cloud vs. sub-cloud is compressed as compared to the profile method. In analyzing Table ??, it is clear that the average turbulence (both TKE and *e*) peaks either in the bottom-middle or top-middle of the cloud (i.e., between a normalized in-cloud height of 0.25-0.75). This is also the two layers in which the buoyancy flux is at a maximum. TKE production near cloud base from latent heat release moves up through the cloud layer, while TKE production near cloud top from evaporative cooling moves down through the cloud layer, resulting in a maximum within the middle of the cloud.

Panels (e) and (f) represent the u, v, and w components of the TKE and  $\epsilon$ , respectively. The anisotropic conditions present within the turbulent boundary flow can clearly be seen due to the differing values in each component. Although the u and v components are similar for most layers, differences are evident in the w-component. If the flow was perfectly isotropic, one would expect the same values for each component of the TKE and  $\epsilon$ .

#### 660 4.3 RF 11 (November 1st)

Turbulent and boundary layer characteristics have been shown to be abnormal on Nov.  $1^{st}$ , with a minimum in 500-hPa geopotential height, aerosol number concentration, and cloud droplet number concentration. November  $N_{a}$ , and  $N_{d}$ , Nov.  $1^{st}$ 

also had overall mean maximum values of TKE and  $\epsilon$  within the sub-cloud layer, along with maximum values in the surface SHF and in-cloud drizzle rate. The average drizzle rate in-cloud on November-Nov. 1<sup>st</sup> was the largest recorded (a mean

- 665 in-cloud drizzle water content of 0.025 gm<sup>-3</sup> measured by the CIP probe) and roughly 4.5 times that of the second largest in-cloud average recorded on November Nov.  $2^{nd}$  (0.0055 gm<sup>-3</sup>), where the average for all other flights was 0.0014 gm<sup>-3</sup>. A moist layer is present above the boundary layer from looking at profiles of q in Figure 5, leading to the secondary maximum in LWC and cloud thickness (November Nov.  $2^{nd}$  had the largest cloud thickness and LWC). Also, visible in Figure 3 is the presence of wind shear near  $z/z_i = 0.60$ .
- 670 In order to explore this case further, Figure 12 shows profiles of multiple thermodynamic and turbulent variables as a function of  $z/z_i$ . Panel (a) shows profiles of  $\theta$  (blue), LWC (black), and q (red). The gray envelope represents the cloud layer, while the orange envelopes represent areas in the sub-cloud layer where the SHIF  $\underline{F}_{\theta}$  is negative and TKE and  $\epsilon$  are enhanced. The potential temperature at the base of the lowest orange envelope begins to deviate from its surface value, decreasing significantly. Normalizing  $\theta$  from 0 to 1 (where the surface is 0 (the minimum temperature) and the top is 1 (the maximum temperature), we
- 675 find that the value of  $\theta$  is 0.32 at cloud top and 0.10 at cloud base, increasing steadily up to the inversion, inferring significant entrainment of the warmer, less buoyant air aloft down to  $z/z_i \sim 0.40$ . However, q within the boundary layer stays relatively constant. This is due to the fact that the entrainment of the warmer air aloft has a larger q than that near the surface of the boundary layer. Significant decoupling is occurring in the sub-cloud layer, near  $z/z_i = 0.60$  (where the largest TKE and  $\epsilon$ are located) and 0.40 (secondary maximum in the TKE and  $\epsilon$ ). It is suggested here that precipitation from the Sc deck acts
- 680 to decouple the boundary layer and enhance sub-cloud turbulence due to evaporative cooling of precipitation from the Se deck above. Zheng et al. (2011) states that the cloud liquid water path reached a maximum on Nov. 1<sup>st</sup> and Nov. 2<sup>nd</sup> due to the total-water specific humidity above the inversion being larger than that within the boundary layer. The inversion strength became significantly weaker on these two days (as evident from Figure 6) and the boundary layer was decoupled due to drizzle. occuring primarily in the regions outlined by the orange envelopes.
- Several variables must be considered here. First, the moist layer above the Sc deck can have two effects, including (1) changing the radiative balance at cloud top through increased downwelling longwave radiation (Christensen et al., 2013) and (2) Entrainment entrainment of more moist air near cloud top, reducing evaporational cooling that would otherwise occur through the entrainment of drier air (Eastman et al., 2017). Both effects act to reduce cooling (both evaporational and radiational) near cloud topand slows, slowing the rate of boundary layer deepening through decreases in entrainment. Eastman and Wood
- 690 (2018) found that high humidity above the Sc deck acts to slow boundary layer deepening while the entrainment of increased water vapor into the boundary layer results in enhanced cloud cover.

Second, drizzle can have multiple effects on boundary layer structure, including (1) precipitation removes liquid water from the Sc deck, resulting in cloud thinning if the <u>surface</u> LHF is not large enough to maintain the Sc deck (Austin et al., 1995); (2) Warming warming of the drizzle producing cloud layer occurs through latent heating, acting to stabilize the cloud layer; (3)

695 Changing the stability of the sub-cloud layer depending on the rate of precipitation profile of sub-cloud evaporation. Significant proportions of precipitation are known to evaporate before reaching the surface (Comstock et al., 2004; Wood et al., 2015; Zhou et al., 2015). The profile of sub-cloud evaporation, and where this evaporation occurs determines whether the layer will become more or less unstable. When precipitation is heavier and in the form of large drops it tends to stabilize the boundary layer from evaporational cooling spread over the depth of the sub-cloud layer, with substantial evaporation near the

500 surface stabilizing the boundary layer. When precipitation is lighter and in the form of small drops, cooling persists in the uppermost part of the sub-cloud region, resulting in destabilization of the sub-cloud layer (Feingold et al., 1996; Wood, 2005; Mechem et al., 2012; Rapp, 2016; Ghate and Cadeddu, 2019; Wood, 2012).

Here, precipitation promotes STBL decoupling by reducing the diabatic cooling in the cloud layer through in-cloud latent heating effects resulting in a stabilization of the cloud layer (where the average in-cloud turbulence is the 4<sup>th</sup> lowest measured

- on Nov. 1<sup>st</sup> and lowest measured on is the 5<sup>th</sup> lowest measured while Nov. 2<sup>nd</sup> is the lowest measured, see Figure 9). The sub-cloud evaporation leads to cooling below cloud and a resultant local minimum in the buoyancy flux is created (Bretherton and Wyant, 1997). It is known from Wood (2005) that evaporative cooling shows cooler and more moist characteristics than that of non-precipitating regions. The SHF  $F_{q}$  is observed to be negative from  $z/z_i \sim 0.4$  up to cloud base, with the minimum and local minimum outlined in the orange envelopes. The LHF  $F_{q}$  is also shown to be slightly enhanced within these regions
- 710 -(i.e., an enhanced source of vapor from evaporation). This suggests that evaporational cooling is occurring away from the surfacein these regions, resulting in the largest average turbulence being measured in the sub-cloud layer on this day due to sub-cloud destabilization. From earlier, it was mentioned that Zheng et al. (2011) suggested drizzle processes acted to stabilize the boundary layer, leading to decoupling. This is partially true, as the precipitation does lead to decoupling, however, the precipitation actually destabilized the sub-cloud layer while stabilizing the cloud layer.
- Normally, this process the process just explained will result in the cloud layer being decoupled form the surface moisture source, leading to a thinning cloud layer. However, the Sc deck is receiving moisture from the upper atmosphere (as seen in the negative LHF  $\mathbf{F}_{q}$  above cloud, where w' is negative but q' is positive). This process acts to moisten the boundary layer, which will lower the LCL, and assuming that  $z_i$  does not change, this will thicken the cloud (Randall, 1984). Note that the cloud layer on Nov. 2<sup>nd</sup> is thicker than that on Nov. 1<sup>st</sup> by roughly 100 m, while  $z_i$  is roughly 50-m lower.
- T20 Looking at w'N', an increase occurs near cloud base up to the middle region of the cloud, before decreasing to negative values in the upper half of the cloud. The positive values near cloud base occur due to droplet activation through condensation, while the negative values occur in the upper half of the cloud from upward vertical velocity perturbations having less droplets than that of negative vertical velocity perturbations, suggesting that droplet activation may be occurring near cloud top as well. Panel (c), N<sub>d</sub> is provided by the Phase Doppler Interferometer (see Chuang et al. (2008) for more information), which
- provides a time series of droplet arrival times with no instrumentation dead time. The average drop size is also provided. The average drop size in-cloud (sub-cloud) is 20.7  $\mu$ m (6.49  $\mu$ m). The average drop size below cloud base and above the top orange envelope is 0.26  $\mu$ m (the maximum value for the profile). The average N<sub>d</sub> in-cloud (sub-cloud) is 81.7 cm<sup>-3</sup> (15.23 cm<sup>-3</sup>). N<sub>d</sub> below the bottom orange envelope is 7.15 cm<sup>-3</sup> whereas N<sub>d</sub> above the bottom orange envelope to cloud base is 25.33 cm<sup>-3</sup>. Panel (c) infers two conclusions: (1) the rapid evaporation of large drops within the first orange envelope ( $z/z_i \sim 0.60$ ); (2)
- 730 The evaporation of a majority of the remaining smaller droplets within the second orange envelope  $(z/z_i \sim 0.40)$ , reinforcing the fact that evaporation away from the boundary layer surface results in decoupling while enhancing sub-cloud turbulence.

Note that w'w'w' also varies between positive and negative values within the sub-cloud layer, providing more evidence that decoupling is occurring.

To summarize, it appears that the sub-cloud layer is decoupled from the Sc deck due to the evaporative cooling of pre-

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- cipitation. This increases turbulence within the sub-cloud layer while reducing turbulence in the cloud layer. However, the cloud layer is still supplied with moisture through the entrainment of the more moist air aloft, driving cloud deepening and sustaining the Sc deck. The wind direction shifts from the south in the lower portion of the boundary layer to from the north northerly near  $z/z_i = 0.60$ . The fact that the free atmosphere wind direction extends into the sub-cloud layer indicates that significant entrainment mixing has occurred, resulting in the upper 40% of the boundary layer to share characteristics with the
- 740 free atmosphere (whereas Zheng et al. (2011) attribute this to an upsidence wave). Note that the maximum value in TKE that is measured on Nov.  $1^{st}$  at  $z/z_i = 0.60$  (see the light blue profile line in Figure ??10) matches the location at which the wind shear is occurring. However, this spike in TKE cannot be attributed to the wind shear alone, as wind shear that occurs at the inversion for each flight day and within the boundary layer on Nov.  $4^{th}$  do does not result in large increases in turbulence. The increase in turbulence seen on Nov. 1<sup>st</sup> is related to latent heating affects and the resulting changes in the buoyancy fluxes.
- Although not displayed here, profiles for Nov.  $2^{nd}$  (the day with the lowest average turbulence, both in-cloud and sub-745 cloud) shows show a very consistent turbulent profile (no large spikes within or below the cloud layer). It is suggested here that between Nov.  $1^{st}$  and  $2^{nd}$  one of two things occurred, either (1) Precipitation precipitation stopped (i.e., the source of instability in the sub-cloud layer) and enhanced turbulent mixing of the sub-cloud layer ceased (while the cloud layer continued to deepen from the entrainment of more moist air reducing the LCL) or the more likely candidate (2) Precipitation precipitation continued
- 750 to occur, leading to evaporation near the surface and a stabilization of the entire boundary layer - Note that although (there is evidence that a small number of droplets are already reaching the surface from Figure 12 Panel (c)). Although in-cloud drizzle is occurring (stabilizing the cloud layer through latent heating) on Nov.  $2^{nd}$ , there is no evidence of sub-cloud evaporation -There or drizzle. Therefore, there are limited sources of turbulent production until dryer air moves in and enhanced entrainment cooling near cloud top can resume mixing of the boundary layer, or if precipitation restarts and acts to destabilize the sub-cloud
- laver. It should be noted that although attention has been brought to Nov. 1<sup>st</sup> and Nov. 2<sup>nd</sup> throughout the manuscript (due to 755 the passing synoptic system leading to unique characteristics that warranted further investigation), these two flight days do fit the overall correlations that were presented in Table 4.

Comparing RF11 to a well-mixed boundary layer, Figure 13 provides the In Figure 13 (same format as that of Figure 12, except for Figure 12), a well-mixed boundary layer in RF03 (Oct. 19<sup>th</sup>) -is analyzed. Also, note that Panel (c) provides the

- TKE flux as opposed to N<sub>d</sub> and drop size (since there is no sub-cloud precipitation to analyze). Both  $\theta$  and q appear to be 760 well-mixed throughout the boundary layer, with a slight decrease in  $\theta$  throughout the cloud layer. TKE,  $\epsilon$ , the LHF, and the SHF  $F_{\alpha}$ , and  $F_{\theta}$  all have two peaks, one near cloud base and one near cloud top, suggesting latent heating near cloud base and evaporative cooling near cloud top. The SHF- $F_{\theta}$  also has a negative value above cloud top due to the entrainment of warm, dry air down into the cloud. The droplet number concentration flux increases near cloud base owing to droplet activation, and sees
- 765 a sharp decrease near a normalized in-cloud height of 0.50, suggesting most of the activation is occurring in the bottom half of the cloud layer. The vertical velocity skewness has a maximum negative value near cloud base, and never has an increase

to positive values. The negative TKE flux within the cloud layer suggest that upward moving air is transporting less TKE than that of downward moving air -(i.e., the main source of turbulence is from entrainment mixing near cloud top resulting in evaporative cooling).

## 770 5 Conclusion

Variations in turbulent and meteorological properties within the boundary layer on a flight by flight-flight-by-flight basis (synoptic variation) have been examined. It has been shown that the influence of a synoptic system on Nov. 1<sup>st</sup> and Nov. 2<sup>nd</sup> leads to a deepening of the cloud layer during passage from due to a moist layer directly above the boundary layer. A large increase in z<sub>i</sub> is observed after passage. Although the pressure is increasing (and subsidence becomes stronger) after the passage of the synoptic system, it is proposed that the moist layer above the boundary layer limits z<sub>i</sub>-boundary layer deepening due to reduced evaporational and radiational cooling near cloud top, limiting entrainment (counteracting the fact that subsidence is weaker). As the synoptic system passes and the upper atmosphere dries, cloud top cooling is enhanced and entrainment acts to deepen increase z<sub>i</sub>, counteracting the fact that subsidence is increasing. Turbulence is shown to be rather weak as compared to other observational studies of Se decks. TKE is shown to vary around 0.13 m<sup>2</sup>s<sup>-2</sup>, except on the days leading up to and following the synoptic system passage, where the TKE increases rapidly to a maximum on Nov. 1<sup>st</sup> due to precipitation leading to enhanced turbulence in the sub-cloud layer and then decreases significantly to a minimum on Nov. 2<sup>nd</sup>. Analysis over the observation period indicates:

 As the pressure decreases (increases), z<sub>i</sub> increases (decreases), accompanied by a decrease (increase) in turbulencewithin the boundary layer of in-cloud turbulence. As z<sub>i</sub> deepens, increases, the chance of boundary layer decoupling increases due to cooling near cloud top cannot being unable to sustain mixing over the entire depth of the boundary layer, resulting in less turbulence and decoupling as compared to a shallow, well mixed boundary layer.

- As Correlation coefficients indicate that as the LHF and SHF increases (decreases) increase,  $z_i$  increases(decreases). When the LHF increases however, the cloud thickness decreases (increases) tends to decrease. A larger LHF tends to produce thinner Sc clouds but a larger  $z_i$ , suggesting enhanced entrainment which at cloud top generated from the larger LHF (through more moisture being available for evaporation) acts to thin the cloud layer while deepening the boundary layer.

- A maximum in TKE on Nov. 1<sup>st</sup> (both overall average and largest single value measured) is due to precipitation acting to destabilize the sub-cloud layer (through evaporation occuring away from the surface, primarily near z/z<sub>i</sub> ~ 0.4 and z/z<sub>i</sub> ~ 0.6), while acting to stabilize the cloud layer. This is observed in both the vertical profiles of RF11 and the TKE and ε values in Figure 9, where it is shown that the distributions of turbulence the turbulent data for the sub-cloud and eloud layer-in-cloud layers are completely offset from one another, with the TKE in the sub-cloud layer maximizing for the analysis period, while the TKE in the eloud-in-cloud layer is below the average value for the analysis period. Nov. 2<sup>nd</sup> has the lowest average turbulence measured (both in-cloud and sub-cloud), and is believed to be a result of (1) lack of cooling near cloud top due to the enhanced moist layer above and (2) Heavy-heavy precipitation from the

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previous day (or sometime prior to the measurements being made) leading to evaporation through the entire sub-cloud laver, stabilizing it.

- Six of the fourteen Eight of the 18 flights have a maximum TKE within the cloud layer. Seven of the fourteen Ten of the 18 flights display two peaks in TKE within the cloud layer, one near cloud base and another near cloud top, signifying evaporative cooling near cloud top and latent heating near cloud base. Of the six eight flights that have a maximum TKE within the cloud layer, all six eight display two peaks in the TKE within the cloud layer, one near cloud base and one near 805 cloud top. This suggests that enhanced turbulence below the cloud can act to reduce latent heating and cooling effects within the cloud layer which generate turbulence near cloud top and bottom. Perhaps, enhanced Enhanced sub-cloud turbulence (as compared to in-cloud) could be an initial indicator that the process of boundary layer decoupling has begun, but has not developed to the point that classical measurement techniques of decoupling (like those discussed in Section 3.3.2) can measure decoupling yet, the decoupling. All five of the decoupled flights, with the exception of Nov.  $1^{st}$  and Nov.  $2^{nd}$ , have a single peak in TKE in the cloud layer, with the maximum TKE value being reached within the 810

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sub-cloud layer.

- Analyzing different layers of turbulence over the 14-eighteen flights shows that TKE,  $\epsilon$ , the vertical velocity variance, TKE, and the buoyancy flux, on average, all reach maximum values near cloud middle (between normalized in-cloud values of 0.25- 0.75).
- 815 The results presented here represent a snapshot of data through 14-18 aircraft flights, with at least a day between any two flights. Therefore, the results presented represent boundary layer conditions that were present at the time of measurement, limiting any analysis of continuously evolving boundary layer and turbulent conditions. For example, being able to analyze the changing thermodynamic and dynamic conditions that resulted in large turbulent changes between Nov. 1<sup>st</sup> and Nov. 2<sup>nd</sup> would be ideal, especially since multiple papers have called for observational studies to assess the impact of drizzle evaporation
- induced cooling on boundary layer turbulence (Wood et al., 2016; Zheng et al., 2016, 2017). It has also been displayed that 820 how turbulence is analyzed is important to understanding the true extent of how turbulence varies within the boundary layer. Taking large scale averages of turbulent parameters (such as over entire horizontal flight legs) may lead to important smaller resolution variations being averaged out. For example, the vertical profiles presented in Figures 12 and 13 show much more detail in the vertical trends as compared to the averaged results of horizontal leg means displayed in Figure ??.
- Data availability. All cabin data from different aircraft platforms can be found on the VOCALS-REx website at https://archive.eol.ucar.edu/ 825 projects/vocals/rex.html (last access: 27 October 2020). All NCEP/NCAR reanalysis data can be found from NOAA at https://www.esrl.noaa .gov/psd/data/gridded/data.ncep.reanalysis.html.

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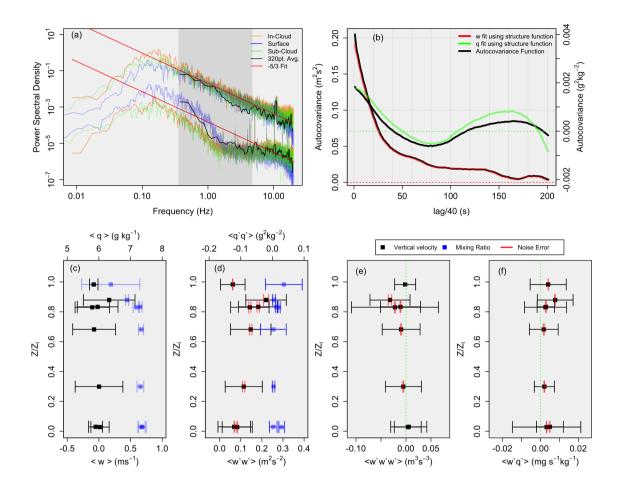
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**Figure 1.** All data presented is from RF3. Panel (a): Power spectral density for three different horizontal flight legs, including in-cloud (orange), sub-cloud (green), and near surface (blue) fit with a -5/3 power law (red), where the upper fit is vertical velocity  $(m^2 s^{-1})$ , and the lower fit is  $q (g^2 kg^{-2})$ . The black specra overlaid represents a single specra using a 320-point average. The light gray envelope represents the 0.3 to 5-Hz range; Panel (b): Autocovariance Autocovariance functions of vertical velocity and q (black) with the fit structure function (green for q and red for vertical velocity); Panel (c): leg mean vertical velocity (black) and q (blue), where the error bars represent the square root of the total variance; Panel (d): As in Panel (c), except for the variance. Note that red Red error bars represent the noise error, while the remaining error bars represent the sampling error; Panel (e): As in Panel (d), except for vertical velocity skewness; Panel (f); As in Panel (d), except for the flux w'q'.

Panel (a): Mean sea level pressure (hPa) between Oct. 19<sup>th</sup> and Nov. 12<sup>th</sup> with the standard deviation overlaid. Note that
 the contours are every 3-hPa; Panel(b): Mean 700-hPa geopotential height with mean omega (ω, hPa/day) overlaid at the same level. Contours are every 10-m. The red dot in both panels represents the location of Point Alpha.

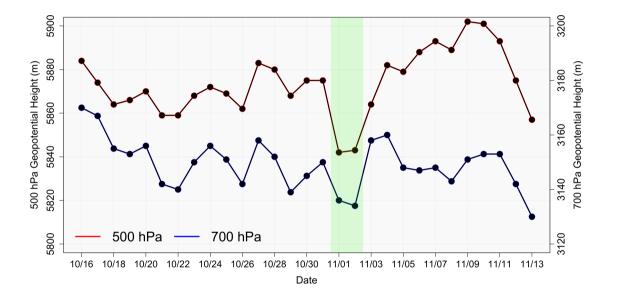
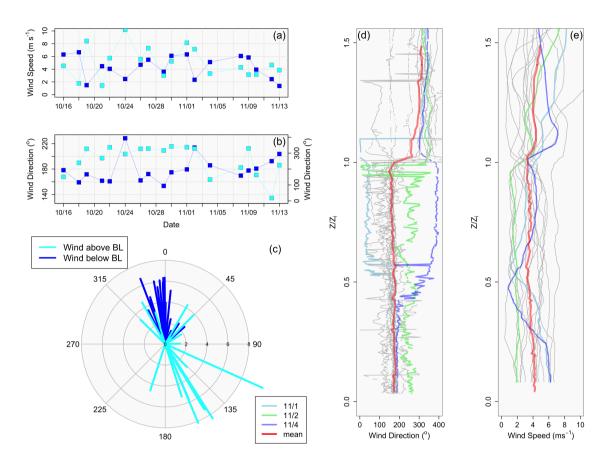
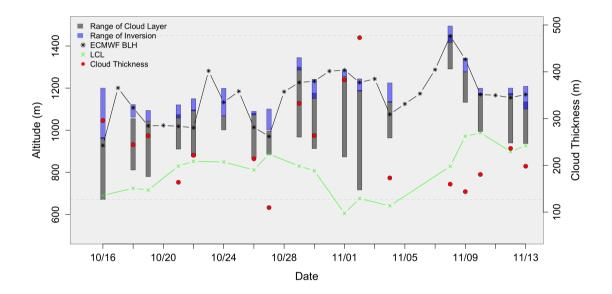


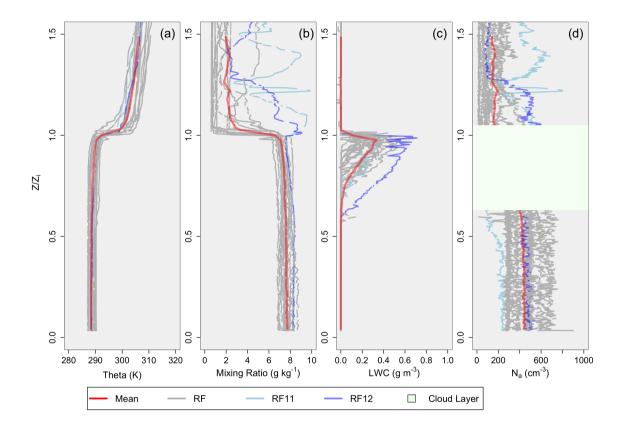
Figure 2. Panel (a): NCEP/NCAR reanalysis (open squares) and flight data (solid squares); Panel (b): Sea surface temperature and atmospheric surface temperature collected from flight data; Panel (c): 500-hPa (red) and 700-hPa (blue) geopotential height data from NCEP/NCAR reanalysis data; Panel (d): As in panel (c), except for 850-hPa (red) and 1000-hPa (blue).



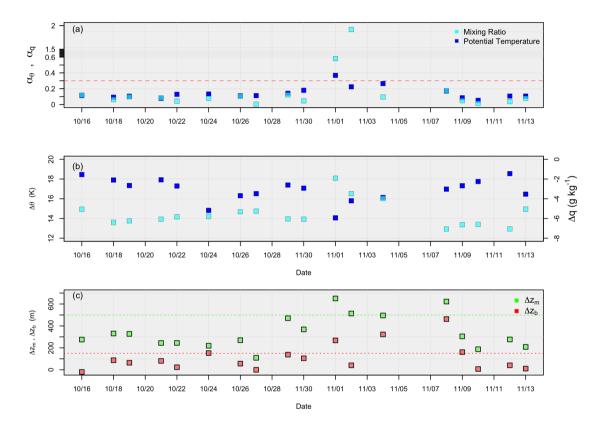
**Figure 3.** Panel (a): Wind speed  $(ms^{-1})$  at the surface collected during 30-m horizontal flight legs (dark blue) and above the inversion collected during horizontal flight legs above the boundary layer (light blue); Panel (b): As in Panel (a), except for wind direction (degree); Panel (c): Vectors showing wind direction from panel (b); Panel (d): Vertical profiles (collected during aircraft soundings) of wind direction for each flight day plotted vs. normalized boundary layer height. Nov.  $1^{st}$  and  $4^{th}$  are displayed in light blue and dark blue, respectively. Nov.  $2^{nd}$  is green, and the mean wind direction is represented by red; Panel (e): As in Panel (d), except for wind speed  $(ms^{-1})$ .



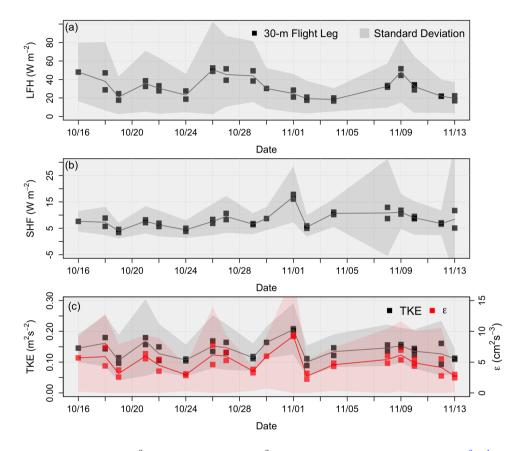
**Figure 4.** Shows the The range of the cloud (gray) and inversion (blue) layer as a function of altitude for each RF. The top of each gray profile represents cloud top and  $z_i$ . The bottom of each gray profile represents cloud base. Cloud thickness (represented as a single value) is represented by each red dot (right y-axis). The LCL and ECMWF- $z_i$  are provided with the black star and green x, repsectively.



**Figure 5.** Profiles scaled by  $z_i$ . (a)  $\theta$  (K); (b) q (g kg<sup>-1</sup>); (c) LWC (g m<sup>-3</sup>); (d) Aerosol number concentration  $N_d$  (cm<sup>-3</sup>). The red profile represents the mean value, and the two blue profiles represent RF 11–RF11 (light blue) and RF 12. RF12 (dark blue). The green layer represents the relative cloud layer for panel (d), as aerosol data cannot be collected in the cloud layer.

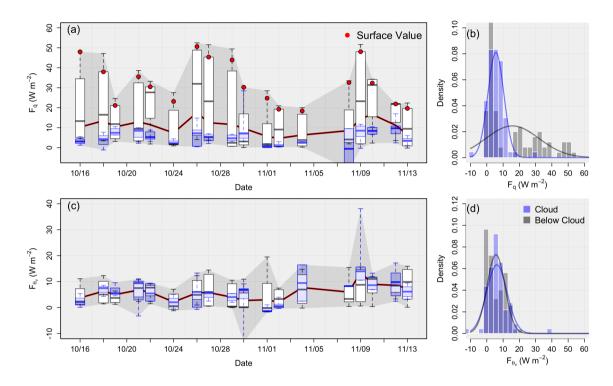


**Figure 6.** (a)  $\theta$  (left y-axis, dark blue) and q (right y-axis, light blue) differences across the inversion, for all 14-18 flights; (b) The decoupling parameters for mixing ratio (light blue) and potential temperature (dark blue), where the red dashed line represents the 0.30 value; (c) Mixed layer cloud thickness (green) and the difference between cloud base and the LCL (red), where the red dashed line represents the 500 value and the green dashed line represents the 150 value.

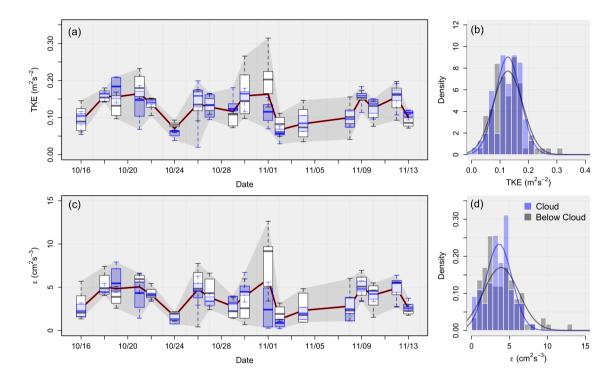


**Figure 7.** Values of (a) surface LHF (Wm<sup>-2</sup>); (b) surface SHF (Wm<sup>-2</sup>); and (c) the Bowen ratio surface TKE (m<sup>2</sup>s<sup>-1</sup>) in black and surface  $\epsilon$  (cm<sup>2</sup>s<sup>-3</sup>) in red, for each flight day. Note that each black Each square is a mean of a 30-m horizontal flight leg, while the gray each envelope represents the standard deviation.

As in Figure 7, except for (a) Friction velocity  $(m^2s^{-2})$ ; (b) Vertical velocity variance  $(m^2s^{-2})$ ; (c) TKE  $(m^2s^{-2})$ ; (d)  $\epsilon$   $(em^2s^{-3})$ .



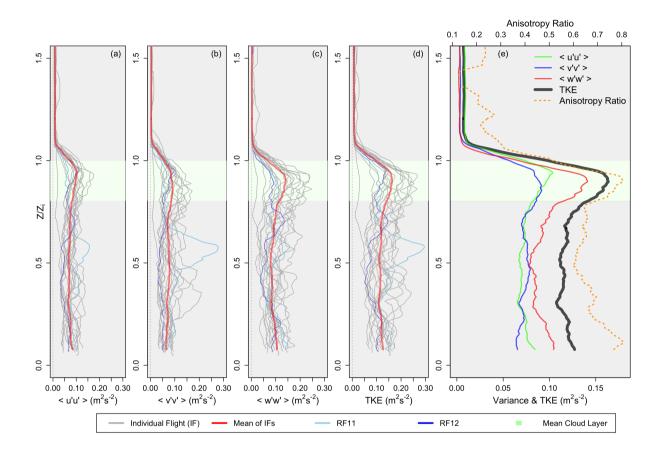
**Figure 8.** Boxplots of in-cloud (blue) and sub-cloud (white) data using mean values of horizontal flight legs for (a) the LHF  $F_q$  (Wm<sup>-2</sup>) and (c) the Buoyancy flux  $F_{\theta_u}$  (Wm<sup>-2</sup>). Note that the The gray envelope represents the range of the data, while the red line represents the mean value for each flightmean values. Panels (b) and (d) shows provide the distributions distribution of the data populations (with normal distributions overlaid for reference) for the LHF  $F_q$  and buoyancy flux  $F_{\theta_u}$ , respectively.



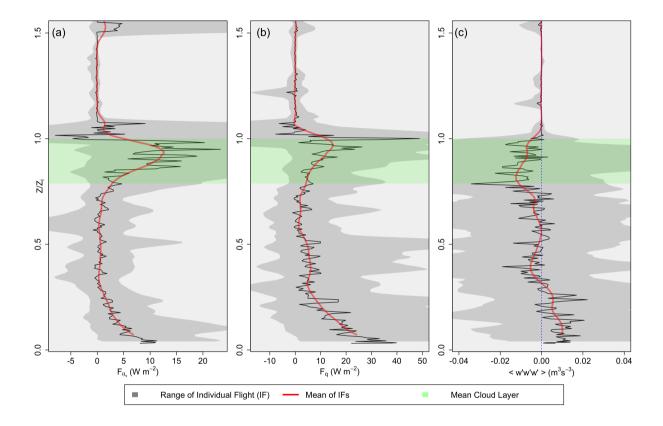
**Figure 9.** As in Figure 8, except for TKE  $(m^2s^{-2})$  in (a) and (b) and  $\epsilon$   $(cm^{-2}s^{-3})$  in (c) and (d)

Mean values of horizontal flight legs plotted as a function of normalized boundary layer height. In-cloud data is red while

1030 data collected during Nov.  $1^{st}$  and  $2^{nd}$  is blue. The standard deviation is represented in orange from vertical sounding data. (a) Buoyancy flux (Wm<sup>-2</sup>); (b) LHF (Wm<sup>-2</sup>); (c) Vertical velocity variance (m<sup>2</sup>s<sup>-2</sup>); (d) TKE (m<sup>2</sup>s<sup>-2</sup>).

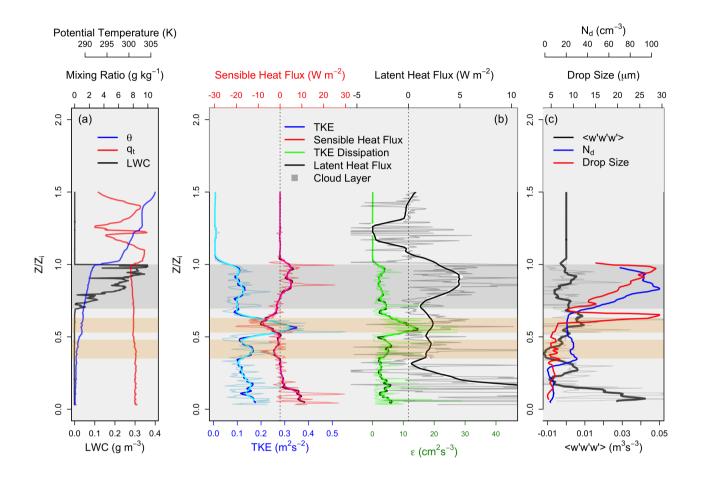


**Figure 10.** Vertical profiles (from data collected during flight profiles) of (a) <u>u-variance</u> <u>u-variance</u>  $(m^2 s^{-2})$ ; (b) <u>v-variance</u> <u>v-variance</u>  $(m^2 s^{-2})$ ; (c) <u>w-variance</u>  $(m^2 s^{-2})$ ; (d) TKE  $(m^2 s^{-2})$ . Individual flights are displayed in gray, the mean value is displayed in red, with RF11 and RF12 shown in <u>light blue and dark blue</u>, respectively. Panel (e) shows the mean Mean values from each of panels (a) through (d), along with the anisotropy ratio in orange.

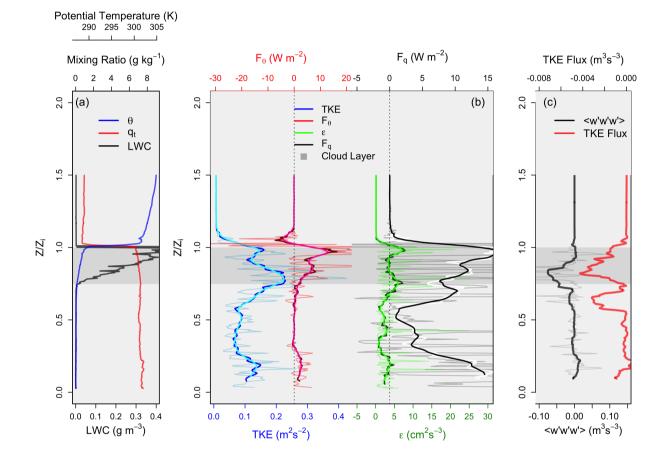


**Figure 11.** Vertical profiles (from data collected during flight profiles) of (a) Buoyancy flux  $F_{\theta_{ac}}$  (Wm<sup>-2</sup>); (b) LHF  $F_{q}$  (Wm<sup>-2</sup>); (c) Vertical velocity skewness (m<sup>3</sup>s<sup>-3</sup>); (d) cloud droplet flux (ms<sup>-1</sup>ce<sup>-1</sup>). Note that the The red line is the smoothed average of the raw data (black), while the gray envelope represents the range of values encountered.

(a) Buoyancy flux (Wm<sup>-2</sup>); (b) LHF (Wm<sup>-2</sup>); (c) TKE (m<sup>2</sup>s<sup>-2</sup>); (d)  $\epsilon$  (cm<sup>2</sup>s<sup>-3</sup>). Data is divided into layers, including below-cloud (BEL-C), bottom of cloud (BC), bottom-middle of cloud (BMC), top-middle cloud (TMC), and top of cloud (TC). Red represents mean values for each layer using data collected during flight vertical profiles while blue represents mean values for each layer using data collected during flight legs. Black dots represent mean values for each flight using horizontal flight leg data. Note that the black dots in Panel (c) represent individual leg mean values as opposed to mean flight values. Panels (c) and (f) represent *u* (blue), *v* (green), and *w* (red) components of the TKE and  $\epsilon$ , respectively.



**Figure 12.** Vertical profiles as functions of normalized boundary layer height for RF11 (Nov. 1<sup>st</sup>) of (a)  $\theta$  (K) in blue, q (g kg<sup>-1</sup>) in red, and LWC (gm<sup>-3</sup>) in black; (b) TKE (m<sup>2</sup>s<sup>-2</sup>) in blue, SHF F<sub> $\theta$ </sub> (Wm<sup>-2</sup>) in red,  $\epsilon$  (cm<sup>2</sup>s<sup>-3</sup>) in green, and the LHF F<sub>g</sub> (Wm<sup>-2</sup>) in black. Note that the The thin light colored lines represent raw values, while the dark thick lines represent smoothed averages; (c)  $\langle w'w'w' \rangle$  (m<sup>3</sup>s<sup>-3</sup>) in black, the TKE Flux  $N_d$  (m<sup>3</sup>scm<sup>-3</sup>) in redblue, and the corresponding average droplet number concentration flux size ( $\mu$ ms<sup>-1</sup>ee<sup>-1</sup>) in bluered. Note that the The gray envelope represents the cloud layer, and the orange envelopes represent areas of interest (main location of decoupling locations and evaporation).



**Figure 13.** As in Figure 12, except for the well-mixed boundary layer case of  $\frac{RF03}{RF3}$  (Oct 19<sup>th</sup>). Panel (c) provides a profile of TKE flux (m<sup>3</sup>s<sup>-3</sup>) as opposed to N and average drop size.

Flight	Date	Time (UTC)	BL Conditions
RF_1	<u>Oct 16, 2008</u>	<u>15:10 - 17:50</u>	Well Mixed
$\frac{RF2}{2}$	<u>Oct 18, 2008</u>	12:15 - 14:40	Well Mixed
RF 3	Oct 19, 2008	12:05 - 14:40	Well Mixed
RF 4	Oct 21, 2008	12:10 - 14:50	Well Mixed
RF 5	Oct 22, 2008	12:00 - 14:40	Well Mixed
RF <u>6</u>	<u>Oct 24, 2008</u>	<u>12:15 - 15:00</u>	Decoupled
<u>RF</u> 7	Oct 26, 2008	12:00 - 15:00	Well Mixed
RF 8	Oct 27, 2008	15:55 - 19:00	Well Mixed
RF 9	Oct 29, 2008	11:50 - 15:00	Well Mixed
RF 10	Oct 30, 2008	11:50 - 15:00	Well Mixed
RF 11	Nov 01, 2008	12:05 - 15:05	Wind Shear / Moisture Above
RF 12	Nov 02, 2008	11:55 - 15:00	Moisture Above
RF 13	Nov 04, 2008	11:50 - 14:40	Wind Shear
RF 15	Nov 08, 2008	11:50 - 15:00	Decoupled
RF 16	Nov 09, 2008	11:50 - 15:05	Well Mixed
RF 17	Nov 10, 2008	14:45 - 18:00	Well Mixed
RF 18	Nov 12, 2008	11:50 - 15:15	Well Mixed
<u>RF 19</u>	<u>Nov 13, 2008</u>	12:00 - 14:50	Well Mixed

**Table 1.** Column (1): Research Flight (RF) identification; (2) The corresponding date; (3) Flight start and end times at Point Alpha. Note that local time: UTC – 4; (4) Boundary Layer conditions for each flight.

**Table 2.** Mean, standard deviation, and range of values for select variables over the 14-18 flights analyzed, with standard deviation values in parenthesis.

	Mean	Range
$z_i$ (m)	<del>1173-1148 (</del> 119)	<del>993-<u>996</u> -</del> 1450
Cloud Base (m)	<del>936 (141</del> 918 (142)	<del>716<u>670</u> -</del> 1291
Cloud Thickness (m)	<del>237 (101229 (98</del> )	<del>107-67</del> - <del>475</del> 473
Boundary Layer $\theta$ (K)	289 ( <del>1.12</del> 1.06)	287 - 291
Boundary Layer $q$ (g kg <sup>-1</sup> )	<del>7.51 (0.48</del> 7.53 (0.43)	<del>6.83</del> - <u>6.82</u> - <del>8.32</del> 8.34
$\Delta \theta$ (K)	<del>16.8 (1.10<u>16.9 (1.18</u>)</del>	13.89 - <del>18.42<u>18.54</u></del>
$\Delta q  (\mathrm{g \ kg^{-1}})$	<del>-5.53 (1.50</del> -5.55 (1.32)	-7.10 - 1.46
Boundary Layer PCASP (cc <sup>-1</sup> $N_a$ (cm <sup>-3</sup> )	410 (127418 (124)	<del>642 -</del> 230 <u>- 673</u>
$\frac{\text{CDNC}\left(\text{cc}^{-1}\text{N}_{d}\left(\text{cm}^{-3}\right)\right)}{\text{CDNC}\left(\text{cc}^{-1}\text{N}_{d}\left(\text{cm}^{-3}\right)\right)}$	<del>262 (110280 (111</del> )	80.5 - 423
Drop Size (µm)	<del>12.33 (2.83</del> 12.35 (2.81)	9.6 - 20.5
Boundary Layer Wind Speed (m $s^{-1}$ )	<del>4.42 (1.444.38 (1.76)</del>	<del>2.03-</del> 1.37 - <del>6.31</del> 6.66
Boundary Layer Wind Direction (°)	<del>170 (46179 (29</del> )	NA
Free Atmosphere Wind Speed (m $s^{-1}$ )	<del>5.16 (1.89</del> 5.33 (3.19)	2.83 - <del>8.34<u>15.14</u></del>
Free Atmosphere Wind Direction ( $^{\circ}$ )	<del>280 (115</del> 276 (93)	NA
$lpha_ heta$	0.15 0.14 (0.08)	<del>0.05 0.052</del> - 0.37
$lpha_q$	<del>0.071 (0.049</del> 0.075 (0.044)	0.002 - 1.94
$\Delta z_m$ (m)	<del>363 (164</del> 340 (151)	<del>108-<u>109</u> - <u>653651</u></del>
$\Delta z_b$ (m)	<del>125 (136</del> 111 (127)	<del>1.8</del> -20 - 463

**Table 3.** Mean and range of values for select surface variables over the 14-18 flights analyzed, with standard deviation and the research flight number in parentheses for column mean and range, respectively.

Latent heat flux $(Wm^{-2})$
Sensible heat flux $(Wm^{-2})$
Friction velocity (ms <sup>-1</sup> )0.17 (0.023)0.13 (RF 11) - 0.22 (RF12) Vertical velocity variance (ms <sup>-1</sup> )0.097 (0.017)0.073 (RF 11) - 0.114 (RF 12) TKE (m <sup>2</sup> )
TKE dissipation rate ( $cm^2s^{-3}$ )

**Table 4.** Correlation coefficient values in the right-panel , with the and variables in the left panel. Note that , with  $\leftrightarrow$  divides dividing the variables being compared. GPH is geopotential height (i.e., a proxy for pressure).  $N_a$ -Values with \* were calculated excluding RF 11 and  $N_D$  represent RF 12 due to the aerosol number concentration and the cloud droplet number concentrationincrease in mixing ratio above  $z_i$ , respectively resulting in abnormally large  $\alpha_g$  values.

$z_i \leftrightarrow  ext{GPH}$
$\omega\leftrightarrow \mathrm{GPH}$
Wind speed $\leftrightarrow$ GPH
$SHF \leftrightarrow wind speed$
$LHF \leftrightarrow wind speed$
$z_i \leftrightarrow \text{wind speed}$
$z_i \leftrightarrow \mathrm{SHF}$
$z_i \leftrightarrow \mathrm{LHF}$
$LHF \leftrightarrow cloud thickness$
$SHF \leftrightarrow cloud thickness$
$\operatorname{in-cloud} \epsilon \leftrightarrow z_i$
$\frac{\text{sub-cloud }\epsilon \leftrightarrow z_i - 0.13}{\text{in-cloud TKE}}$ in-cloud TKE $\leftrightarrow$
$\frac{\text{sub-cloud TKE} \leftrightarrow z_i \text{-}0.20}{N_a} \leftrightarrow \text{in-cloud}$
$N_D \leftrightarrow \underline{\text{in-cloud}} \operatorname{TKE}$
drop size $(\mu m) \leftrightarrow \underline{\text{in-cloud}}$ TKE
$N_a \leftrightarrow \operatorname{in-cloud} \epsilon$
$N_D \leftrightarrow \operatorname{in-cloud} \epsilon$
drop size $(\mu m) \leftrightarrow \underset{\sim}{\text{in-cloud}} \epsilon$
Mean values for each layer discussed in Figure ??, where the top rows represent data calculated using flight vertical profiles, while the bottom rows represent data calculated using flight vertical profiles, while the bottom rows represent data calculated using flight vertical profiles, while the bottom rows represent data calculated using flight vertical profiles, while the bottom rows represent data calculated using flight vertical profiles, while the bottom rows represent data calculated using flight vertical profiles, while the bottom rows represent data calculated using flight vertical profiles, while the bottom rows represent data calculated using flight vertical profiles, while the bottom rows represent data calculated using flight vertical profiles, while the bottom rows represent data calculated using flight vertical profiles, while the bottom rows represent data calculated using flight vertical profiles, while the bottom rows represent data calculated using flight vertical profiles, while the bottom rows represent data calculated using flight vertical profiles, while the bottom rows represent data calculated using flight vertical profiles, while the bottom rows represent data calculated using flight vertical profiles, while the bottom rows represent data calculated using flight vertical profiles.
$\frac{\text{Below Cloud } z_i \leftrightarrow \Delta z_b}{z_i \leftrightarrow \Delta z_b}$
$\frac{\text{Cloud Base}}{z_i \leftrightarrow \alpha_{\theta}}$
$\frac{\text{Bottom Middle}_{z_i}\leftrightarrow \alpha_q}{\text{Middle}_{z_i}\leftrightarrow \alpha_q}$
$\frac{\text{Top Middle} \Delta z_{m} \leftrightarrow \text{in-cloud TKE}}{\Delta z_{m} \leftrightarrow \text{in-cloud TKE}}$
$\frac{\text{Cloud Top} \Delta z_{\mathcal{R}} \leftrightarrow \text{in-cloud } \epsilon}{\epsilon}$
$\frac{\text{Horizontal Leg Data} \Delta z_b \leftrightarrow \text{in-cloud } ]}{\text{Horizontal Leg Data} \Delta z_b \leftrightarrow \text{in-cloud } ]}$

 $\begin{array}{l} \textbf{Below Cloud} \Delta z_k \leftrightarrow \text{in-cloud} \ \epsilon \\ \hline \textbf{Cloud Base} \alpha_{\theta} \leftrightarrow \text{in-cloud TKE} \\ \hline \textbf{Bottom Middle} \alpha_{\theta} \leftrightarrow \text{in-cloud} \ \epsilon \\ \hline \textbf{Top Middle} \alpha_g \leftrightarrow \text{in-cloud TKE} \\ \hline \textbf{Cloud Top} \alpha_g \leftrightarrow \text{in-cloud} \ \epsilon \\ \hline \end{array}$