

Reply to the editor

We would like to thank again both the anonymous reviewers and Editor for reviewing this work.

We have addressed all the technical corrections suggested by Reviewer2 in the modified version.

The main point raised by Reviewer1 is about the compensating error from the LW TOA fluxes. To this point we have now added the biases and rmse values for these TOA fluxes in the Supplementary material for both Southern Ocean and global mean with respect to observational data. As we have now added in the modified Discussion section, the improvement in SW TOA always leads to an increase in the LW component due to thick clouds which is unfortunately unavoidable.

We have also made it clear that even if the same value for capacitance could be applied globally, the only regions where it will make noticeable differences are that of the mixed-phase cloud regions (like SO).

Below are the specific points that we have addressed.

1. Title: Is the Southern Ocean cloud albedo bias reduced compared to observations? This needs to be quantified, see general comment.

We have modified the title to 'Improving the Southern Ocean cloud albedo biases in a general circulation model' to avoid any confusion. The SW CRE bias values have been presented in Table S1 of Supplementary material.

2. P1L1-10: The SW changes are accompanied by LW changes; this needs to be mentioned in the abstract and main text and also be discussed in the main text.

We have modified the Discussion section to stress on LW compensating effect. We have also added the bias values along with the RMSE values (both SO and global) with respect to observational data in tables S2 and S3 of Supplementary material.

3. P2L18: Unless this a named sea surface temperature climatology remove the name here.

Removed

4. P3L1-10: It needs to be discussed here where the new shape parameter is more realistic than the previous one (based on literature), at which latitudes and at which temperature (ranges). In particular it needs to be shown that it's more realistic for the SO region.

Included in the Discussion section

5. P3L18: Do not use contractions in scientific writing.

Corrected

6. P4L22: The shear-dominated unstable layer shows a similar change in LWP in the cap experiment as the stratocumulus boundary layer type clouds. This needs to be discussed.

Addressed in the Supplementary section where the particular plot is shown

7. P4L29-31: The change in the shape parameter will slow down the WBF process but the changes to nucleation temperature will determine if there is ice in the first place, which will also impact the WBF process. So why does the change in the shape parameter affect both LWP and IWP while changes to nucleation temperature mainly impact IWP?

Modified the results section to make it clear.

8. P4L31: LWPs in Fig. 1 are for stratocumulus boundary layer type clouds, why are these dominated by fronts?

Sentence rephrased

9. P5L32-33: Is it cloudy and non-cloudy grid boxes or cloudy and non-cloudy parts of grid boxes?

Rephrased the sentence

10. P6L7: control model, control run, ctrl experiment: choose one expression and use it consistently everywhere

Corrected

11. P6L7-8: Not for the SO region (50S to 70S). What is the net bias over this region?

Same as response 1.

12. P6L8-9: It is stated that SW CRE is reduced over SO, but in the title or abstract the opposite is stated. For which experiments is there an improvement? This needs to be quantified (see general comment).

Same as response 1.

13. P6L9-10: This is in contradiction to the previous sentence.

The SW CRE is calculated by taking the anomaly of TOA SW flux between the clear-sky and all-sky conditions. The SW CRE shows a reduction (as shown in Fig 5) compared to the control run. Whereas Figs. 2b and 3b are only showing the outgoing SW radiation at the TOA which is showing an increase compared to the control run. Basically showing that there is now more outgoing SW at TOA over SO in the experiments w.r.t control run. So an increase in outgoing SW TOA or a reduction in SW CRE over SO would be improvement.

14. P6L11: Do not use contractions in scientific writing

Corrected

15. P6L11-13: Why is the effect of the changed shape parameter in the tropics so much smaller than in the SO region?

Already included in the Results section.

16. P6L12-13: Don't you mean increase in SW CRE here? Fig. 2 suggests an increase in SW CRE in the sensitivity experiments.

Same as response 13.

17. P6L16-17: From Fig. 4b it is not clear that this is actually the case.

Same as response 13.

18. P6L18: Does the model have convective LWP? Rephrase otherwise

Removed the sentence

19. P7L4: It is unclear what “It” refers to.

Modified

20. P7L18: “frontal steady states”: I’m not familiar with this expression

Rephrased the sentence.

21. P7L20-21: As there are undesirable biases in the NH (which need to be discussed more), what is the conclusion for the shape parameter? Should the new shape parameter be used in global climate models or not?

Same as response 4.

22. P7L29-32: Which assumptions about the INP concentration are made?

Rephrased the sentence.

23. P8L9-10: Has the bias compared to observations in the SO region (50S to 70S) indeed been reduced? This needs to be quantified (see specific comments above and general comment).

Same as response 1

24. P15: Fig. 2: The use of upward fluxes for TOA and downward fluxes for the surface should be specified in the title of each panel. For all figures that show fluxes it should be specified in the panel titles if upward or downward fluxes are shown. In the caption of Fig. 2 it is mentioned that upward fluxes are used for TOA and downward fluxes for the surface. This information needs to be added to other figures that show upward or downward fluxes.

Added

Reducing Improving the Southern Ocean cloud albedo biases in a general circulation model

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Abstract. The present generation of global climate models is characterized by insufficient reflection of short-wave radiation over the Southern Ocean due to a misrepresentation of clouds. This is a significant concern as it leads to excessive heating of the ocean surface, sea surface temperature biases, and subsequent problems with atmospheric dynamics. In this study we modify cloud micro-physics in a recent version of the Met Office's Unified Model and show that choosing a more realistic value for the shape parameter of atmospheric ice-crystals, in better agreement with theory and observations, benefits the simulation of short-wave radiation. In the model, for calculating the growth rate of ice crystals through deposition, the default assumption is that all ice particles are spherical in shape. We modify this assumption to effectively allow for oblique shapes or aggregates of ice crystals. Along with modified ice nucleation temperatures, we achieve a reduction in the annual-mean short-wave cloud radiative effect over the Southern Ocean by up to $\sim 4 \text{ Wm}^{-2}$, and seasonally much larger reductions [compared to the control model](#). By slowing the growth of the ice phase, the model simulates substantially more supercooled liquid cloud.

1 Introduction

One of the major known problems in present-day global climate models is an excess in the absorbed short-wave (SW) radiation over the Southern Ocean (SO) (Trenberth and Fasullo, 2010; Ceppi et al., 2012; Hwang and Frierson, 2013; Williams et al., 2013; Hyder et al., 2018). Chapter 9 of the 5th Assessment Report of the Intergovernmental Panel on Climate Change (IPCC AR5) (Flato et al., 2014) points out that most of the 5th Coupled Model Intercomparison Project (CMIP5) models (Taylor et al., 2012) have a positive SW cloud radiative bias of magnitude of up to 20 Wm^{-2} over the SO, suggesting that inadequately simulated clouds allow substantially too much sunlight to reach the ocean surface.

Several studies have focused on the relation between various aspects of cloud representation in the model and radiation biases pronounced over the SO. Bodas-Salcedo et al. (2012, 2014), using cyclone compositing cluster analyses, suggest the need to increase the optical depth of the low-level clouds and improve the simulation of mid-level cloud regime, to help reduce the biases in the model. By modifying the shallow convection detrainment in their global climate model, Kay et al. (2016) showed that the resultant increase in the supercooled liquid clouds (SLC) enable large reductions in long-standing climate model SW radiation biases. Furtado et al. (2016) and Lohmann and Neubauer (2018) point towards the significance of mixed-phase cloud and their representation in the models for better representation of SO. In another study by Furtado and Field

(2017), the importance of ice micro-physics parametrisation in determining the phase composition, and thus the liquid water content of the SO clouds is highlighted.

Discrepancies in the response of clouds to anthropogenic forcings are recognized as a leading reason for a persistent, large spread in the climate sensitivity throughout various generations of climate models (Allen et al., 2014). We thus conjecture that this model problem of cloud albedo bias over the SO contributes to this large spread, and thus solving it would increase confidence in projections of anthropogenic climate change (Tan et al., 2016).

In the present study, we investigate the role of parameters involved in atmospheric ice formation within a global climate model in causing the above mentioned SW radiation bias.

2 Data and experimental set-up

We define a SO region as the latitudinal band between 50°S and 70°S in this study. The control climate model used in this study is the most recent version of the Met Office’s Unified Model, GA7.1 (Walters et al., 2019) with modified micro-physics scheme for riming process and several other scientific changes. Appendix A summarizes the scientific set-up for this model version. The resolution used here is N96L85 (i.e. a horizontal resolution of $1.875^\circ \times 1.25^\circ$ and 85 terrain-following hybrid-height levels extending to 85 km of altitude). It uses the "ENDGAME" dynamical core with a semi-implicit semi-Lagrangian formulation to solve the non-hydrostatic, fully compressible deep-atmosphere equations of motion (Wood et al., 2014)

2.1 Model set-up

The control ~~run~~ [model](#) follows the Atmospheric Model Intercomparison Project (AMIP) climate model development protocol (Gates et al., 1999), using prescribed [Reynolds sea surface temperature climatology covering 1981-2012 \(Reynolds et al., 2007\)](#) [sea surface temperatures](#). Excess atmospheric ice has been a persistent concern in the control version of the model, which is especially pronounced over the SO region. Ice clouds have a significant influence on the global climate through their effects on the Earth’s radiation budget e.g. (Hartmann and Doelling, 1991; Waliser et al., 2009). Hence, sensitivity set-ups in our study are aimed at modifications to the micro-physics scheme such that the ice growth in the model is controlled. We achieve this by modifying those parameters that control the growth of existing ice by vapor deposition and heterogeneous nucleation of new ice. The classical theory of ice crystal growth uses an electrostatic analogy due to the similarity between the equations governing the water vapor distribution around an ice crystal and the electrostatic potential distribution around an electric conductor of the same shape as the ice crystal (Chiruta and Wang, 2003). Thus, the growth rate of ice crystals by diffusion depends on a shape (also known as capacitance) parameter C , which is a function of both ice crystal size and habit. To determine the ice crystal growth rates in models, it is necessary to know the value of C (Chiruta and Wang, 2003; Hobbs, 1976). The standard equation that is used for calculating the growth rate of ice crystals in the model is,

$$\frac{dm}{dt} = 4\pi C (\rho_s - \rho_\alpha) \quad (1)$$

where C is the capacitance, ρ_α and ρ_s are distributions of vapor densities at and away from crystal’s surface (Houghton, 1950).

From eq.1, it is evident that once the value of capacitance C is known, the growth rate of ice crystals can be determined. All other quantities on the right-hand side of eq.1 are independent of the shape (Chiruta and Wang, 2003; Hobbs, 1976). Thus, C in the model effectively defines the shape of ice crystals, which in turn is fed to the ice processes of deposition/sublimation and melting without affecting any other ice processes.

5 Technically the capacitance, C , is defined as $1.0 \times d$ in the model where d is the particle maximum size. However, studies have suggested that this value overestimates the evaporation rate of snowflakes by a factor of 2 (Westbrook et al., 2008). It has been shown in other studies based on theory and observations that by changing the value to $0.5 \times d$, models show a significant improvement compared to the traditional approximations used (Westbrook et al., 2008; Field et al., 2008). Thus, in our sensitivity studies, we modified the value to $0.5 \times d$ (corresponding to any oblate ellipsoid with two unequal axes, thought
10 to be more appropriate for aggregates and plate-like crystals rather than the assumption of spherical crystals alone). The effect of this change in the shape parameter is tested independently as well as in combination with changing the temperatures at which heterogeneous and homogeneous freezing start in the cloud micro-physics scheme. The idea behind modifying the nucleation temperatures is to test the behaviour of capacitance change in a relatively cleaner environment. Basically, the ice nucleation temperature is the temperature at which heterogeneous nucleation of ice first starts to occur in the model. In the
15 control model, this is solely following the temperature dependent function suggested by (Fletcher, 1962). The default value in the model is -10°C for heterogeneous nucleation temperature. However, in much cleaner environments, like the SO, ice nucleation might not start at -10°C as there is a paucity in the ice nucleating particles (INPs). Hence, in reality the nucleation temperatures are much colder for these regions. Since we ~~don't do not~~ have any INP dependency taken into account for ice nucleation temperature in the control model, to test the behaviour of capacitance in such conditions, experiments are
20 conducted by changing the default value of -10°C to -40°C and -20°C . The higher threshold of -40°C has been chosen as it is the maximum temperature at which homogeneous nucleation occurs in the model (i.e. at this temperature all liquid is instantaneously frozen to form ice particles, (Yau and Rogers, 1996)). Along with the nucleation temperature in the micro-physics scheme, convection scheme also impacts the amount of ice produced in the model through its detrainment temperatures. We thus further modified the convection scheme by changing the detrainment temperatures to be ~~very cold from~~ colder than
25 the default values. This thus gives us an overall base to investigate the effect of delaying the heterogeneous ice nucleation in the model. The temperature at which the detraining condensate as ice begins in the model is the start-ice temperature and the temperature at which all condensate is detrained as ice is called the all-ice temperature. Thus, we conducted three sensitivity experiments (henceforth referred to as cap , $c_tnuc=-40$, and $c_tnuc=-20$) to be compared against the control ~~run~~ model. The values used in our numerical simulations are summarised in Table 1.

30 We note that the experiment where nucleation temperature is modified to -40°C (i.e. $c_tnuc=-40$) is applicable mostly for cleaner environments like the SO but not physically realistic for rest of the world, as mentioned in the previous paragraph. However, it is still a much useful sensitivity scenario to study the importance of detrained ice vs. large-scale freezing. All simulations were run for ~~twenty~~ 22 years under steady-state present-day conditions.

2.2 Observational data

We use the National Aeronautics and Space Administration (NASA) Clouds and the Earth's Radiant Energy System - Energy Balanced And Filled (CERES EBAF Ed4.0, Terra-Aqua) for comparing incoming surface long-wave (LW) and short-wave (SW) as well as the top-of-the-atmosphere (TOA) radiation covering the period 2000 to 2018. This data set, in an earlier version, (Loeb et al., 2009), was used in AR5. The overall uncertainty in the monthly all-sky TOA flux for the CERES EBAF Ed4.0 data set is estimated to be 2.5 Wm^{-2} (for for both SW and LW fluxes). For clear-sky TOA, uncertainties in SW and LW fluxes are 5 Wm^{-2} and 4.5 Wm^{-2} respectively (Loeb et al., 2018). Direct observations of ocean surface air-sea fluxes are extremely sparse, particularly over the Southern Ocean. There are large uncertainties in the conventional observational surface heat estimates which can affect the evaluation of simulated surface energy budgets. Hence, for the net surface flux comparison, we are using a combination of both TOA fluxes from satellite and ERA-Interim reanalysis energy divergences (assuming atmospheric column energy conservation). This approach considerably constraints the estimates of net surface flux derived from re-analyses. Further details regarding this approach can be found in (Hyder et al., 2018).

3 Results

Figure 1 represents the anomaly in the annual and DJF mean distributions of ice water path (IWP) and liquid water path (LWP) for stratocumulus boundary layer clouds in the model in various experiments with respect to the control `runmodel`, for the Southern Hemisphere (SH). There are seven boundary layer types that have been identified in the model based on the surface stability and capping cloud (Lock et al., 2000). As our focus is mostly on the stratocumulus boundary layer type clouds in this study, the cloud types considered in this figure are: type 2 = boundary layer with stratocumulus over a stable near-surface layer, type 3 = well-mixed boundary layer and type 4 = unstable boundary layer with a decoupled stratocumulus (DSC) layer not over cumulus. The IWP and LWP are calculated collectively over these types. The other cloud types in the model are those of stable boundary layer (type=1), boundary layer with de-coupled stratocumulus layer over cumulus (type=5), cumulus-capped boundary layer (type=6) and shear-dominated unstable layer (type=7). A similar analysis for the non-stratocumulus cloud types has been provided in the Supplementary material (Figs. S1 and S2) for annual and DJF means.

From Figs. (1a) and (1b), it is evident that there is noticeable decrease in the IWP in the stratocumulus boundary layer clouds as a result of modified micro-physics, which is captured in all sensitivity experiments in both annual and DJF means. The experiment, `c_tnuc=-40` (solid black line in Figs. 1a and 1b) shows the maximum response and the experiment `cap` (solid red line Figs. 1a and 1b) has the minimum decrease in IWP with respect to the control `runmodel`. Conforming to the decrease in the IWP, there is a corresponding increase in the LWP as well, over the SO region (Figs. 1c and 1d). However, the response of LWP is more or less similar in all the 3 experiments with respect to the control `runmodel`. It is because modification to capacitance value affects both the liquid and ice water contents while changes to nucleation temperature will have an impact predominantly on IWP. ~~The zonal mean LWPs as shown in Fig.-(1) are mostly dominated by the fronts. So, even if the changes to nucleation temperature shows some sensitivity, they are mostly away from the frontal systems and hence will be mostly~~ Even any small sensitivity on LWP due to the changes in nucleation temperatures will not have much impact on the zonally

averaged stratocumulus types, as the effects are mostly restricted to the shallow boundary ~~layer types thus not impacting the~~
~~LWP much types and hence not evident for Fig. 1.~~ Thus, experiments where both capacitance and nucleation temperature are
modified will have an added impact on the IWP.

Figure 2 shows the zonal-mean changes in the annual-mean distributions of various radiative fluxes in the model for the
5 SH. Upward fluxes are used for all TOA figures and downward fluxes for surface figures. In all the model experiments there
is a general decrease in the outgoing LW flux at the TOA in the SO region (solid lines in Fig. 2a) with respect to the control
~~run~~model. This is accompanied by a corresponding increase in the outgoing SW flux at the TOA (dotted lines in Fig. 2b),
indicating that in all model experiments the planetary albedo has increased versus the control. Except in *cap*, the decrease in
LW radiation at the TOA, in absolute terms, is larger than the increase in SW TOA over the SO. This is visible in the distribution
10 of net radiation at the TOA (i.e. LW plus SW at TOA) (dashed red line in Fig. 2c). For *cap* experiment, there is an increase in
the net outgoing TOA radiation whereas for *c_tnuc=-40* and *c_tnuc=-20* experiments, it shows a decrease over the SO region.

The surface distributions of the radiative fluxes are represented in Figs. 2d to 2f. In all the experiments, the net downward
LW radiation at the surface shows an increase over the SO (solid lines in Fig. 2d) with respect to the control ~~run~~model. The
corresponding SW component shows a decrease over SO (dotted lines in Fig. 2e). The distribution of anomaly in the radiative
15 heat fluxes with respect to the control ~~run~~model is represented by dashed lines in Fig 2f. It primarily represents the difference
between total net downward surface radiation and total heat flux at the surface i.e (incoming LW + SW at surface) - (sensible
heat flux + latent heat flux). Although there is an improvement (i.e. reduction) in the downward SW component (dotted lines
in Fig 2e), due to the compensating increase in the LW component (solid lines in Fig 2d), there is a net increase in the heat
flux into the surface over the SO (dashed lines in Fig 2f) in almost all experiments. However, the net radiative heat flux shows
20 a tendency of slight decrease over SO for *cap* experiment with respect to the control ~~run~~model (dashed green line in fig. 2f).

Figure 3 shows the distributions of various radiative fluxes in the model for the SH for the DJF season. The radiative fluxes
show a more pronounced response during the austral summer season. The net radiative flux at TOA (dashed lines in Fig. 3c)
~~is showing shows~~ an increase over the SO for all the experiments unlike the annual-mean distribution where only the *cap*
experiment ~~showed shows~~ an increase. Similarly, the net radiative heat flux at the surface (dashed lines in Fig. 3f) shows a
25 general decrease over the SO in all the experiments with respect to the control ~~run~~model for the DJF season. Whereas in
annual-mean, only *cap* experiment showed a decrease in net surface flux.

Figure 4 represents the difference between the model experiments and observational data for annual-mean. It is mainly
intended to provide a reference for the model behaviour in terms of radiative fluxes. By comparing Figs. 2 and 4, it can be seen
that it is the LW component that shows similar signs of response (solid lines in Figs. 2a, 2d and 4a, 4d) than SW for SO region.
30 But the net fluxes show mostly similar signs of response in both TOA and surface in Figs. 2c, 4c and 2f, 4f except for the TOA
flux in *cap* experiment (dashed red line in Fig. 2c).

Figure 5 shows the zonally averaged annual and DJF mean distributions of the anomaly in the SW cloud radiative effect (SW
CRE) between different model experiments with respect to the control ~~run~~model. The SW CRE shows the impact of cloud on
TOA SW flux and is calculated by taking the anomaly of TOA SW flux between the ~~cloudy and non-cloudy grid boxes~~clear-sky

and all-sky conditions. The reduction in the SW radiative flux over SO is more pronounced in the DJF season (Fig. 5b) with $c_tnuc=-40$ experiment showing the minimum reduction in SW CRE compared to the control ~~run~~model.

Figure 6 shows the annual and DJF mean spatial distributions of the anomaly in SW CRE between different model experiments with respect to the control ~~run~~model. It is evident from Fig. 6 that there is a general improvement (i.e. a reduction) in the SW CRE over the SO region in all three experiments compared to the control ~~run~~model. The reduction is more pronounced in the *cap* experiment (Figs. 6a and 6d). The response in $c_tnuc=-40$ is the minimum (Figs. 6b and 6e). As expected, the response is more robust for the DJF season in all experiments with respect to the control ~~run~~model (Figs. 6d to Fig. 6f).

Figure 7 shows the spatial distribution of TOA SW CRE anomaly between various model simulations with respect to the CERES TOA data. As evident from Fig. 7a, the control model does have an excess in the absorbed SW radiation over the SH high latitudes like many other global models. Comparing Fig. 7a with the Figs. 7b to 7d, it can be seen that the changes in micro-physics parametrisation have improved (~~i.e. reduced~~) ~~the SW CRE~~ the SW biases over the SO. In general there is an increase in the reflected SW radiation over the SO compared to the control model.

The values for annual-mean SW CRE biases in the model experiments (in comparison to both control model as well as observational data) have been summarised in Table S1 of the supplementary material.

It is to be noted that the tropical regions ~~don't do not~~ show much reduction in the SW CRE in response to the capacitance changes. While the SO does show significant reduction in SW CRE with respect to the control ~~run~~model, especially for the *cap* experiment (Figs. 6a and 6d), the tropical regions still show an increase. In an earlier study by (Furtado et al., 2016), using a similar version of the control model (Numerical Weather Prediction NWP configuration), it has been shown that for tropics and subtropics there is a general tendency by the model to overpredict the LWP in response to ~~microphysies~~ micro-physics modifications. Increasing the stratiform cloud LWP will cause more SW radiation to be reflected back to space. But over the SO, this effect is beneficial because the control model already has a large negative bias in outgoing SW radiation in that region. ~~Some potential reasons could be the flaws in parametrizations or uncertainties in the estimation of LWP in the convection scheme.~~ Basically, in ~~a frontal steady state~~ tropics or quasi steady state features like that of the fronts, the capacitance does not have much of an impact compared to more dynamic sites like that of supercooled liquid clouds (Furtado et al., 2016, 2015).

Generally, the impact of ice nucleation temperature experiments ($c_tnuc=-40$ and $c_tnuc=-20$) on fluxes is more mixed than using capacitance change alone (*cap*). The detrimental effects due to changes in the ice nucleation temperature could be mostly attributed to the changes in the vertical distribution of clouds affecting not just the low clouds but also the high clouds. By changing the nucleation temperature, essentially the level at which freezing occurs is modified. When the freezing of water in lower levels is delayed it results in cirrus clouds at higher atmospheric levels. This can change the high cloud characteristics thus affecting both LW and SW.

4 Discussion

Errors in the representation of ice clouds is one of the major shortcomings in many of the present-day global climate models and this can have a significant influence on the global climate through their effects on the Earth's radiation budget (Hartmann and Doelling, 1991; Waliser et al., 2009). ~~It is~~ These errors are mostly coupled to an underestimation of ~~SLCs~~. ~~This problem~~
5 ~~is supercooled liquid clouds and are~~ of particular importance in the SO ~~region~~ regions that are characterized by abundant ~~SLCs~~ supercooled liquid clouds (Kay et al., 2016; Bodas-Salcedo et al., 2016; Huang et al., 2012; Hu et al., 2010). When ice and supercooled liquid coexist, the ice grows at the expense of the liquid by the Wegener-Bergeron-Findeisen (WBF) mechanism (Wegener, 1911; Bergeron, 1935; Findeisen, 1938). Acknowledging the complexities in representing the many possible background micro-physical processes that are responsible for this in a global climate model, the primary idea of
10 modifying the shape parameter of ice-crystals is to reduce the rate of depositional growth of ice particles. This reduction essentially slows down the deposition growth of ice crystals, which leaves more water vapor to be available for condensation into liquid phase particles. By lowering the value of capacitance to $0.5 \times d$, we model a decrease in the IWP and an associated increase in the LWP over the SO region (Fig. 1). As a result of this increase in LWP, the outgoing SW fluxes are increased (dotted lines in Figs. 2b and 3b), i.e. ~~an increased LWP corresponds to brighter clouds reflecting more sunlight~~. This results in
15 a decrease of the downwelling short-wave radiation reaching the surface.

Our choice of $0.5 \times d$ for capacitance is based on theory and observational studies (Field et al., 2008; Westbrook et al., 2008). The control model has improved in terms of SO cloud albedo bias with this value than with the default value of $1.0 \times d$. However, as seen in Fig.7, even though the SO has shown signs of improvement, biases still remain especially in the tropics ~~and other frontal steady states~~. ~~Previous~~. As already mentioned before, previous studies using a similar NWP configuration of the control model have suggested that the value of depositional capacitance is important mostly for non-equilibrium situations like mixed-phase whereas for quasi-steady frontal states and tropics, capacitance might not be very significant (Furtado et al., 2016, 2015). Therefore, even if the same value for capacitance could be applied globally, the only regions where it will make noticeable differences are that of the mixed-phase cloud regions (like SO).
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Earlier studies have also suggested that some cloud micro-physics parameterisations produce unrealistically bright clouds, especially over the Northern Hemisphere (NH) (Furtado et al., 2016). In our control model, the SW biases over the NH were smaller than those over the SO and further brightening of modelled NH clouds is undesirable (Furtado et al., 2016). While the changes to nucleation temperature has significant impact on the tropics as well, the capacitance changes are more localized to the high latitudes.
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Our study shows that while modifying the capacitance is clearly benefiting the SO region SW radiative fluxes, there is some ~~detrimental behaviour of fluxes due to the impact of ice nucleation temperature modifications~~. compensating effect from the LW TOA fluxes. As already discussed, the SW anomaly is thought to be mostly due to a lack of liquid clouds in the atmosphere. By targeting the capacitance, the loss of liquid in mixed-phase conditions is reduced which results in an improvement in the outgoing SW TOA. However, this improvement in SW TOA always leads to an increase in the LW component due to thick clouds which is unfortunately unavoidable. The annual-mean biases in TOA radiative fluxes along with RMSE values in
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comparison with observational data have been provided in Tables S2 and S3 of the Supplementary material. The SO RMSE values for SW TOA tend to be lower than global values as expected. Although the bias in SW TOA over SO has improved in general compared to the observational data, that improvement does not reflect much in the RMSE values. While an in-depth ~~analyses~~ analysis of compensating errors from LW TOA is beyond the scope of this particular study, it is a significant aspect and provides an interesting outlook for future studies.

Also, several recent studies point towards the significance of INP for cloud phase (Kanji et al., 2017; Vergara-Temprado et al., 2018). A further development of the research outlined here could be to make glaciation explicitly dependent on INP concentration to attend to the persisting model biases in other regions of the world as well. At present, in most global climate models, cloud phase is determined only by a threshold temperature like in our control model. Vergara-Temprado et al. (2018), using by combining a high-resolution ~~numerical weather prediction NWP model and making assumptions on the~~ NWP model with estimates of INP concentration over the SO, simulate clouds that are far more reflective than those in current global climate models, in better agreement with satellite observations.

5 Conclusions

In this study we reduce the annual-mean SW radiation biases over SO in a recent version of the UK Met Office's Unified Model by up to $\sim 4 \text{ Wm}^{-2}$, and seasonally up to $\sim 8 \text{ Wm}^{-2}$ compared to the control model. This and other contemporary climate models are characterized by excess cloud ice causing biases in SW radiation which are especially pronounced over the SO. Here, we modify the capacitance or shape parameter which represents ice crystal shape and habit. In our sensitivity studies, we reduce this parameter from $1.0 \times d$ to $0.5 \times d$ (corresponding to any oblate sphere shape in general, where the horizontal axes are longer than the vertical axis and more representative of an aggregate or flat ice crystal) and thus delaying the depositional growth of ice particles. This leads to an increase in liquid water in stratiform clouds and consequently increases the outgoing SW over the SO. We also examine the impact of changing other temperature thresholds in the cloud micro-physics scheme for the onset of heterogeneous ice production. Our analysis shows that the SW radiation bias has significantly reduced over the SO after the modification of these parameters. However, disparities still exist in other regions. INPs that are currently not represented in the cloud micro-physics scheme might be a factor in this model behavior. The fact that nucleation temperature changes currently is associated with the same effects globally is undesirable, it further motivates the future work to couple the nucleation temperature to a prognostic or the least a regionally specified INP concentration.

Data availability. Model data is available at: <http://doi.org/10.5281/zenodo.3775170>.

Observational data is available at: https://ceres.larc.nasa.gov/order_data.php and

Appendix A

Several changes were introduced in the control model used in this study relative to its predecessor GA7.1 (Walters et al., 2019; Brown et al., 2012). These changes range from minor bug fixes and optimisation techniques to major science changes in the convection, large-scale-precipitation, boundary layer and radiation schemes. As far as our study is concerned, the main modification to GA7.1 is the inclusion of the modified micro-physics scheme which includes a shape dependence of riming rates using the parameterization by Heymsfield and Miloshevich (2003), as a measure to prevent small liquid droplets from riming (Furtado and Field, 2017).

Some of the modifications in the convection scheme include that of a prognostic based convective entrainment rate, implementation of a new melting scheme to remove larger spikes in convective heating in the mid-troposphere, a revised forced detrainment calculation and a corrected evaporation of convective precipitation to remove existing errors.

The modified boundary layer scheme also includes changes to reduce vertical resolution sensitivity and an improved turbulent kinetic energy diagnostic and how it is used for aerosol activation.

The change in the radiation scheme is the implementation of spectral dispersion suggested by Liu et al. (2008) to improve the simulation of the first aerosol indirect effect.

A brief overview of the science changes is available in the Unified Model Newsletter Dec 2017 edition (Research and Model Development News, pages 10 -12. This document is included along with the model data at: <http://doi.org/10.5281/zenodo.3775170>).

Author contributions. VV carried out the model runs, performed analysis, created figures, wrote the manuscript and was also involved in the design and conceptualization of the study. OM was involved with obtaining the project grant, supervised the study and analyses of results. PF, KF and PH provided guidance in designing the model runs and analyses of results. JW provided technical support in setting up global climate model in super-computer environment. All authors have read and approved the final paper.

Competing interests. The authors declare that they have no conflict of interest.

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References

- Allen, M. R., Barros, V. R., Broome, J., Cramer, W., Christ, R., Church, J. A., Clarke, L., Dahe, Q., Dasgupta, P., Dubash, N. K., et al.: IPCC fifth assessment synthesis report-climate change 2014 synthesis report, 2014.
- Bergeron, T.: On the Physics of Cloud and Precipitation, Procès Verbaux de la Séance de VU, GGI à Lisbonne 19S3, Paris, 1935.
- 5 Bodas-Salcedo, A., Williams, K., Field, P., and Lock, A.: The surface downwelling solar radiation surplus over the Southern Ocean in the Met Office model: The role of midlatitude cyclone clouds, *Journal of Climate*, 25, 7467–7486, 2012.
- Bodas-Salcedo, A., Williams, K. D., Ringer, M. A., Beau, I., Cole, J. N., Dufresne, J.-L., Koshiro, T., Stevens, B., Wang, Z., and Yokohata, T.: Origins of the solar radiation biases over the Southern Ocean in CFMIP2 models, *Journal of Climate*, 27, 41–56, 2014.
- Bodas-Salcedo, A., Hill, P., Furtado, K., Williams, K., Field, P., Manners, J., Hyder, P., and Kato, S.: Large contribution of supercooled liquid clouds to the solar radiation budget of the Southern Ocean, *Journal of Climate*, 29, 4213–4228, 2016.
- 10 Brown, A., Milton, S., Cullen, M., Golding, B., Mitchell, J., and Shelly, A.: Unified modeling and prediction of weather and climate: A 25-year journey, *Bulletin of the American Meteorological Society*, 93, 1865–1877, 2012.
- Ceppi, P., Hwang, Y.-T., Frierson, D. M., and Hartmann, D. L.: Southern Hemisphere jet latitude biases in CMIP5 models linked to shortwave cloud forcing, *Geophysical Research Letters*, 39, 2012.
- 15 Chiruta, M. and Wang, P. K.: The capacitance of rosette ice crystals, *Journal of the atmospheric sciences*, 60, 836–846, 2003.
- Field, P. R., Heymsfield, J., Bansemer, A., and Twohy, C. H.: Determination of the combined ventilation factor and capacitance for ice crystal aggregates from airborne observations in a tropical anvil cloud, *Journal of the Atmospheric Sciences*, 65, 376–391, 2008.
- Findeisen, W.: Kolloid-meteorologische Vorgänge bei Neiderschlags-bildung, *Meteor. Z.*, 55, 121–133, 1938.
- Flato, G., Marotzke, J., Abiodun, B., Braconnot, P., Chou, S. C., Collins, W., Cox, P., Driouech, F., Emori, S., Eyring, V., et al.: Evaluation of climate models, in: *Climate change 2013: the physical science basis. Contribution of Working Group I to the Fifth Assessment Report of the Intergovernmental Panel on Climate Change*, pp. 741–866, Cambridge University Press, 2014.
- 20 Fletcher, N.: *NH The Physics of Rain Clouds*, 1962.
- Furtado, K. and Field, P.: The role of ice microphysics parametrizations in determining the prevalence of supercooled liquid water in high-resolution simulations of a Southern Ocean midlatitude cyclone, *Journal of the Atmospheric Sciences*, 74, 2001–2021, 2017.
- 25 Furtado, K., Field, P., Cotton, R., and Baran, A.: The sensitivity of simulated high clouds to ice crystal fall speed, shape and size distribution, *Quarterly Journal of the Royal Meteorological Society*, 141, 1546–1559, 2015.
- Furtado, K., Field, P., Boutle, I., Morcrette, C., and Wilkinson, J.: A physically based subgrid parameterization for the production and maintenance of mixed-phase clouds in a general circulation model, *Journal of the Atmospheric Sciences*, 73, 279–291, 2016.
- Gates, W. L., Boyle, J. S., Covey, C., Dease, C. G., Doutriaux, C. M., Drach, R. S., Fiorino, M., Gleckler, P. J., Hnilo, J. J., Marlais, S. M., et al.: An overview of the results of the Atmospheric Model Intercomparison Project (AMIP I), *Bulletin of the American Meteorological Society*, 80, 29–56, 1999.
- 30 Hartmann, D. L. and Doelling, D.: On the net radiative effectiveness of clouds, *Journal of Geophysical Research: Atmospheres*, 96, 869–891, 1991.
- Heymsfield, A. J. and Miloshevich, L. M.: Parameterizations for the cross-sectional area and extinction of cirrus and stratiform ice cloud particles, *Journal of the atmospheric sciences*, 60, 936–956, 2003.
- 35 Hobbs, P. V.: *Ice Physics*, Oxford University Press, p. 837, 1976.
- Houghton, H. G.: A preliminary quantitative analysis of precipitation mechanisms, *Journal of Meteorology*, 7, 363–369, 1950.

- Hu, Y., Rodier, S., Xu, K.-m., Sun, W., Huang, J., Lin, B., Zhai, P., and Josset, D.: Occurrence, liquid water content, and fraction of super-cooled water clouds from combined CALIOP/IIR/MODIS measurements, *Journal of Geophysical Research: Atmospheres*, 115, 2010.
- Huang, Y., Siems, S. T., Manton, M. J., Protat, A., and Delanoë, J.: A study on the low-altitude clouds over the Southern Ocean using the DARDAR-MASK, *Journal of Geophysical Research: Atmospheres*, 117, 2012.
- 5 Hwang, Y.-T. and Frierson, D. M.: Link between the double-Intertropical Convergence Zone problem and cloud biases over the Southern Ocean, *Proceedings of the National Academy of Sciences*, 110, 4935–4940, 2013.
- Hyder, P., Edwards, J. M., Allan, R. P., Hewitt, H. T., Bracegirdle, T. J., Gregory, J. M., Wood, R. A., Meijers, A. J., Mulcahy, J., Field, P., et al.: Critical Southern Ocean climate model biases traced to atmospheric model cloud errors, *Nature communications*, 9, 2018.
- Kanji, Z. A., Ladino, L. A., Wex, H., Boose, Y., Burkert-Kohn, M., Cziczo, D. J., and Krämer, M.: Overview of ice nucleating particles, *Meteorological Monographs*, 58, 1–1, 2017.
- Kay, J. E., Wall, C., Yettella, V., Medeiros, B., Hannay, C., Caldwell, P., and Bitz, C.: Global climate impacts of fixing the Southern Ocean shortwave radiation bias in the Community Earth System Model (CESM), *Journal of Climate*, 29, 4617–4636, 2016.
- Liu, Y., Daum, P. H., Guo, H., and Peng, Y.: Dispersion bias, dispersion effect, and the aerosol–cloud conundrum, *Environmental Research Letters*, 3, 045 021, 2008.
- 15 Lock, A., Brown, A., Bush, M., Martin, G., and Smith, R.: A new boundary layer mixing scheme. Part I: Scheme description and single-column model tests, *Monthly weather review*, 128, 3187–3199, 2000.
- Loeb, N. G., Wielicki, B. A., Doelling, D. R., Smith, G. L., Keyes, D. F., Kato, S., Manalo-Smith, N., and Wong, T.: Toward optimal closure of the Earth’s top-of-atmosphere radiation budget, *Journal of Climate*, 22, 748–766, 2009.
- Loeb, N. G., Doelling, D. R., Wang, H., Su, W., Nguyen, C., Corbett, J. G., Liang, L., Mitrescu, C., Rose, F. G., and Kato, S.: Clouds and the
20 earth’s radiant energy system (CERES) energy balanced and filled (EBAF) top-of-atmosphere (TOA) edition-4.0 data product, *Journal of Climate*, 31, 895–918, 2018.
- Lohmann, U. and Neubauer, D.: The importance of mixed-phase and ice clouds for climate sensitivity in the global aerosol-climate model ECHAM6-HAM2, *Atmospheric Chemistry and Physics*, 18, 8807–8828, 2018.
- Reynolds, R. W., Smith, T. M., Liu, C., Chelton, D. B., Casey, K. S., and Schlax, M. G.: Daily high-resolution-blended analyses for sea
25 surface temperature, *Journal of Climate*, 20, 5473–5496, 2007.
- Tan, I., Storelvmo, T., and Zelinka, M. D.: Observational constraints on mixed-phase clouds imply higher climate sensitivity, *Science*, 352, 224–227, 2016.
- Taylor, K. E., Stouffer, R. J., and Meehl, G. A.: An overview of CMIP5 and the experiment design, *Bulletin of the American Meteorological Society*, 93, 485–498, 2012.
- 30 Trenberth, K. E. and Fasullo, J. T.: Simulation of present-day and twenty-first-century energy budgets of the southern oceans, *Journal of Climate*, 23, 440–454, 2010.
- Vergara-Temprado, J., Miltenberger, A. K., Furtado, K., Grosvenor, D. P., Shipway, B. J., Hill, A. A., Wilkinson, J. M., Field, P. R., Murray, B. J., and Carslaw, K. S.: Strong control of Southern Ocean cloud reflectivity by ice-nucleating particles, *Proceedings of the National Academy of Sciences*, 115, 2687–2692, 2018.
- 35 Waliser, D. E., Li, J.-L. F., Woods, C. P., Austin, R. T., Bacmeister, J., Chern, J., Del Genio, A., Jiang, J. H., Kuang, Z., Meng, H., et al.: Cloud ice: A climate model challenge with signs and expectations of progress, *Journal of Geophysical Research: Atmospheres*, 114, 2009.
- Walters, D., Baran, A., Boutle, I., Brooks, M., Furtado, K., Hill, P., Lock, A., Manners, J., Morcrette, C., Mulchay, J., et al.: The Met Office Unified Model Global Atmosphere 7.0/7.1 and JULES Global Land 7.0 configurations, *Geosci. Model Dev*, 2019.

- Wegener, A.: *Thermodynamik der atmosphäre*, JA Barth, 1911.
- Westbrook, C. D., Hogan, R. J., and Illingworth, A. J.: The capacitance of pristine ice crystals and aggregate snowflakes, *Journal of the atmospheric sciences*, 65, 206–219, 2008.
- Williams, K. D., Bodas-Salcedo, A., Déqué, M., Fermepin, S., Medeiros, B., Watanabe, M., Jakob, C., Klein, S. A., Senior, C. A., and
5 Williamson, D. L.: The Transpose-AMIP II experiment and its application to the understanding of Southern Ocean cloud biases in climate models, *Journal of Climate*, 26, 3258–3274, 2013.
- Wood, N., Staniforth, A., White, A., Allen, T., Diamantakis, M., Gross, M., Melvin, T., Smith, C., Vosper, S., Zerroukat, M., et al.: An inherently mass-conserving semi-implicit semi-Lagrangian discretization of the deep-atmosphere global non-hydrostatic equations, *Quarterly Journal of the Royal Meteorological Society*, 140, 1505–1520, 2014.
- 10 Yau, M. K. and Rogers, R. R.: *A short course in cloud physics*, Elsevier, 1996.

Table 1. Values used in the model runs

	Capacitance	Ice nucleation temperature (°C)	Start-ice temperature (°C)	All-ice temperature (°C)
control	1.0	-10	-10	-20
cap	0.5	-10	-10	-20
c_tnuc=-40	0.5	-40	-40	-41
c_tnuc=-20	0.5	-20	-40	-41

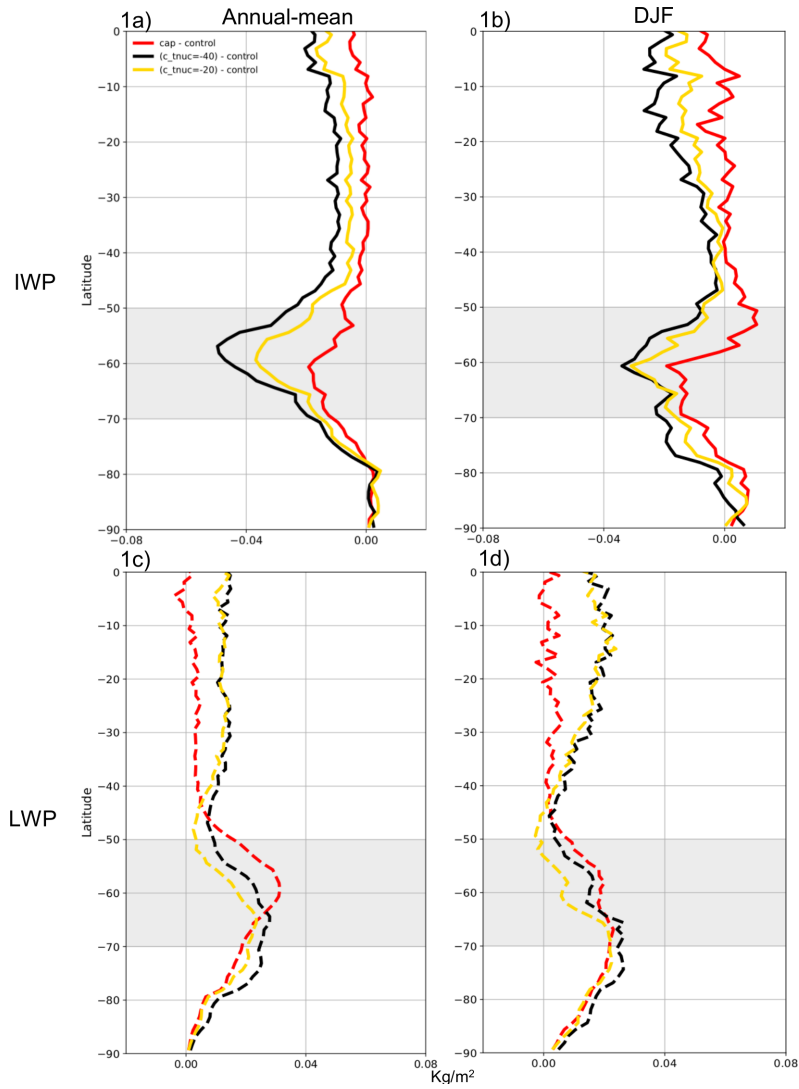


Figure 1. Distribution of zonally averaged anomalies in IWP (solid lines in 1a and 1b) and LWP (dashed lines in 1c and 1d) over the stratocumulus boundary layer type clouds in the model for the SH. (1a) and (1c) represent annual-mean; (1b) and (1d) represent DJF mean. The colour codes are as follows: red = anomaly of *cap* with respect to control, black = anomaly of *c_tnuc*=-40 with respect to control, yellow = anomaly of *c_tnuc*=-20 with respect to control. Values are calculated from 12 hourly instantaneous model output over ~~20-22~~ years. The SO region identified in this study is highlighted in gray.

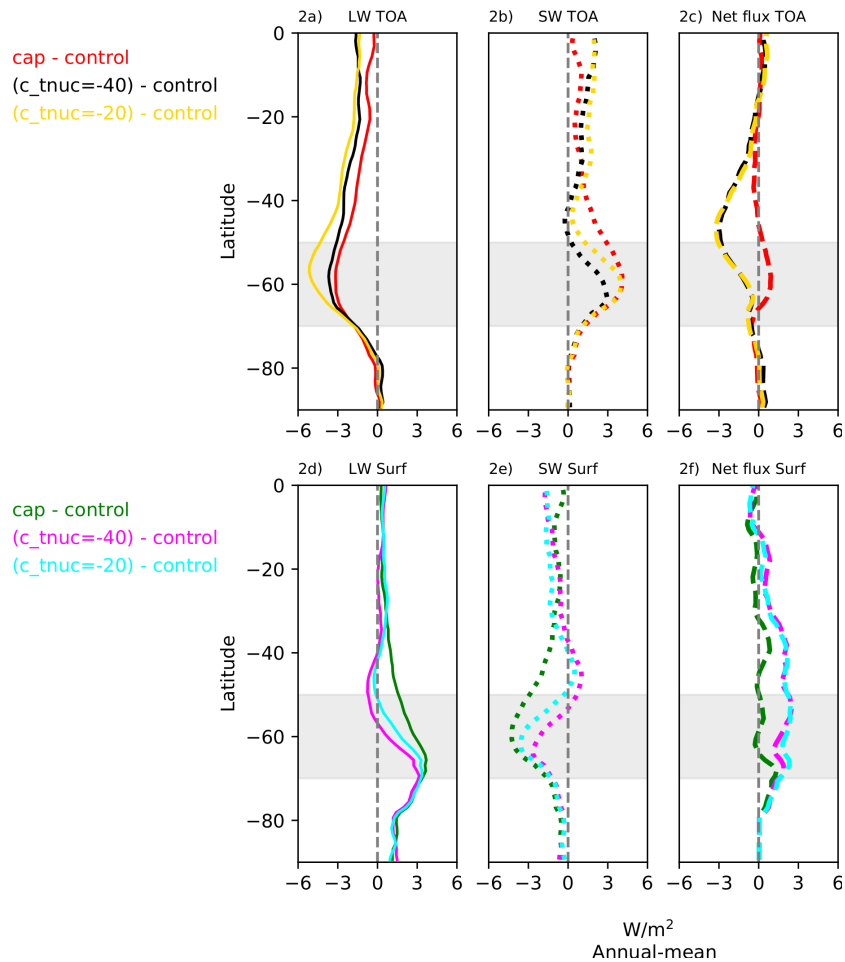


Figure 2. Distribution of zonally averaged annual-mean radiative flux anomalies in various experiments with respect to the model control ~~run~~ model for the SH. (2a) to (2c) represent the TOA (upward) radiative flux anomalies and (2d) to (2f) represent the surface (downward) flux anomalies. The solid lines in (2a) = LW at TOA, dotted lines in (2b) = SW at TOA and dashed lines in (2c) = net flux at TOA (i.e. LW + SW at TOA). Color codes for TOA fluxes are: red = *cap* - control ; black = (*c_tnuc*=-40) - control and yellow = (*c_tnuc*=-20) - control. Similarly, the solid lines in (2d) = LW at surface, dotted lines in (2e) = SW at surface and dashed lines in (2f) = net radiative heat flux at surf (i.e. incoming LW + SW at surface - (sensible heat + latent heat)). Color codes for surface fluxes are: green = *cap* - control ; magenta = (*c_tnuc*=-40) - control and cyan = (*c_tnuc*=-20) - control. Annual-mean values for model are calculated from daily-mean output over ~~20-22~~ years. The SO region identified in this study is highlighted in gray.

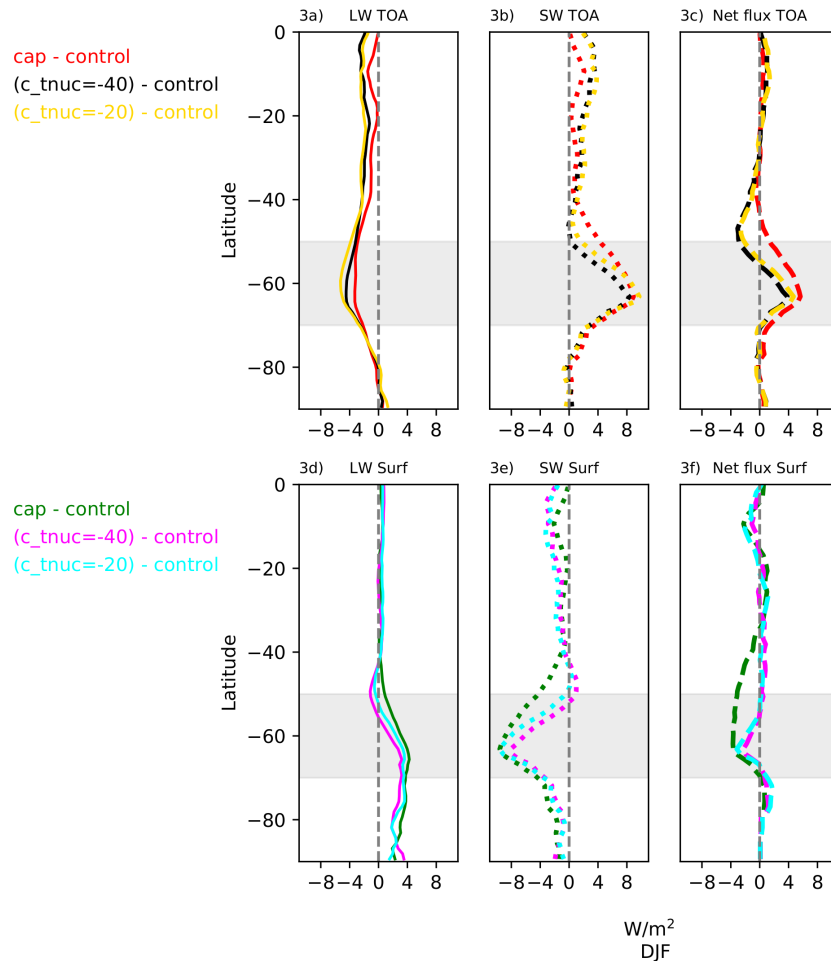


Figure 3. Similar to Fig 2. but for DJF season

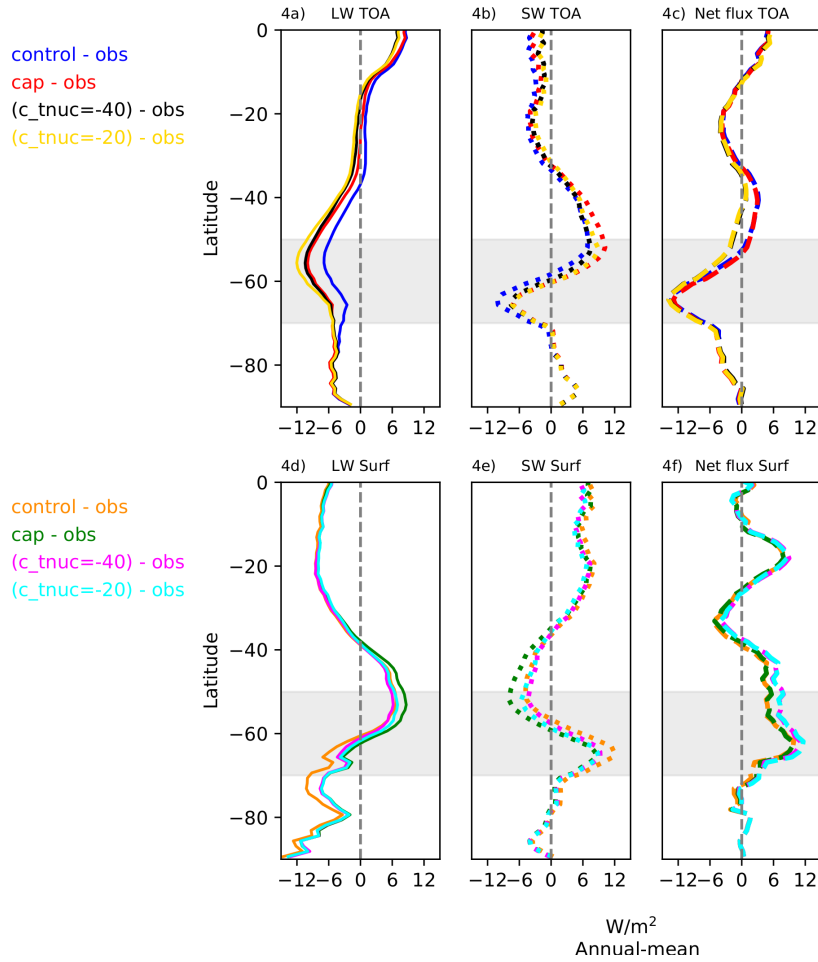


Figure 4. Distribution of zonally averaged annual-mean radiative flux anomalies in various model experiments with respect to observations for the SH. (4a) to (4c) represent the TOA (upward) radiative flux anomalies and (4d) to (4f) represent the surface (downward) flux anomalies. The solid lines in (4a) = LW at TOA, dotted lines in (4b) = SW at TOA and dashed lines in (4c) = net flux at TOA (i.e. LW + SW at TOA). Color codes for TOA fluxes are: blue = control - obs; red = *cap* - obs; black = (*c_tnuc*=-40) - obs and yellow = (*c_tnuc*=-20) - obs. Similarly, the solid lines in (4d) = LW at surface, dotted lines in (4e) = SW at surface and dashed lines in (4f) = net radiative heat flux at surf (i.e. incoming LW + SW at surface - (sensible heat + latent heat)). Color codes for surface fluxes are: orange = control - obs; green = *cap* - obs; magenta = (*c_tnuc*=-40) - obs and cyan = (*c_tnuc*=-20) - obs. Annual-mean values for model are calculated from daily-mean output over ~~20-22~~ years. The net surface flux observational data is a combination of both TOA fluxes from satellite and ERA-Interim reanalysis energy divergences (details in Section 2.2). The SO region identified in this study is highlighted in gray.

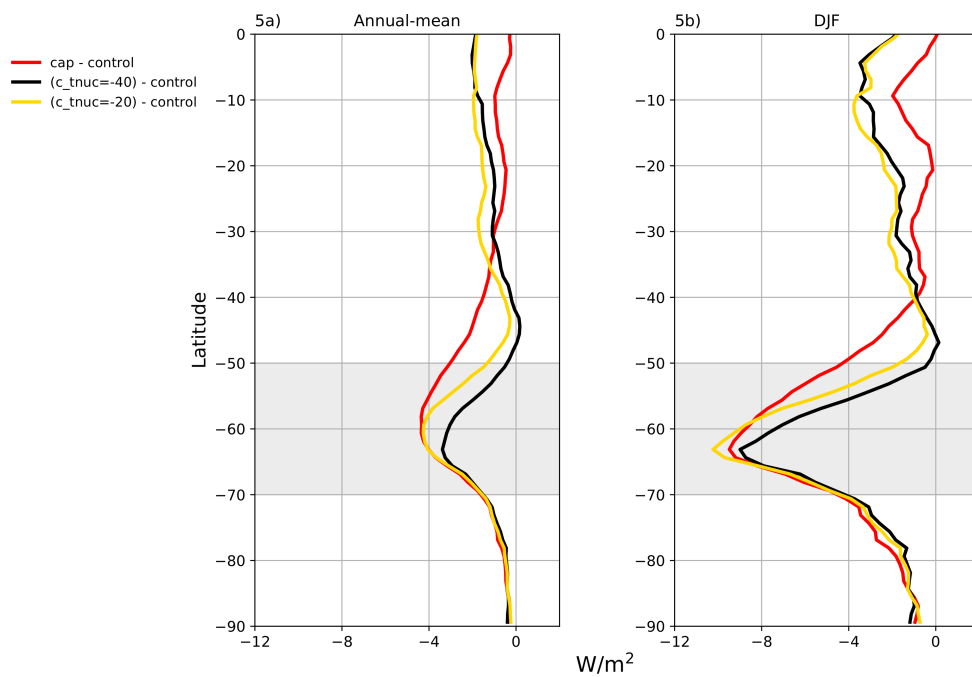


Figure 5. Distribution of zonally averaged SW CRE anomalies over SH in various experiments with respect to the model control [run-model](#) for (5a) annual-mean and (5b) DJF mean. The colour codes are as follows: red = *cap* - control, black = (*c_tnuc*=-40) - control and yellow = (*c_tnuc*=-20) - control. Values are calculated from daily-mean output over 20-22 years. The SO region identified in this study is highlighted in gray.

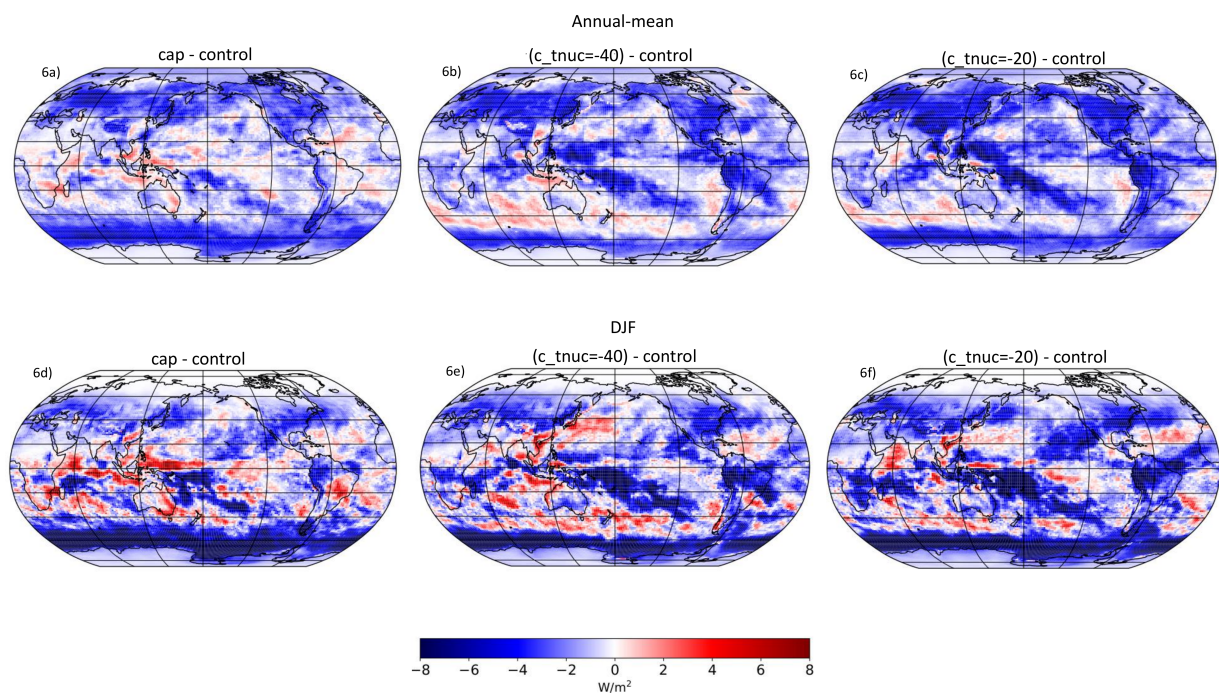


Figure 6. Spatial distributions of annual-mean (6a) to (6c) and DJF mean (6d) to (6f) SW CRE anomaly at TOA in different experiments with respect to the control [FHMmodel](#). (6a) and (6d) = *cap* - control, (6b) and (6e) = (*c_tnuc*=-40) - control, (6c) and (6f) = (*c_tnuc*=-20) - control. Values calculated from daily-mean output over [20-22](#) years.

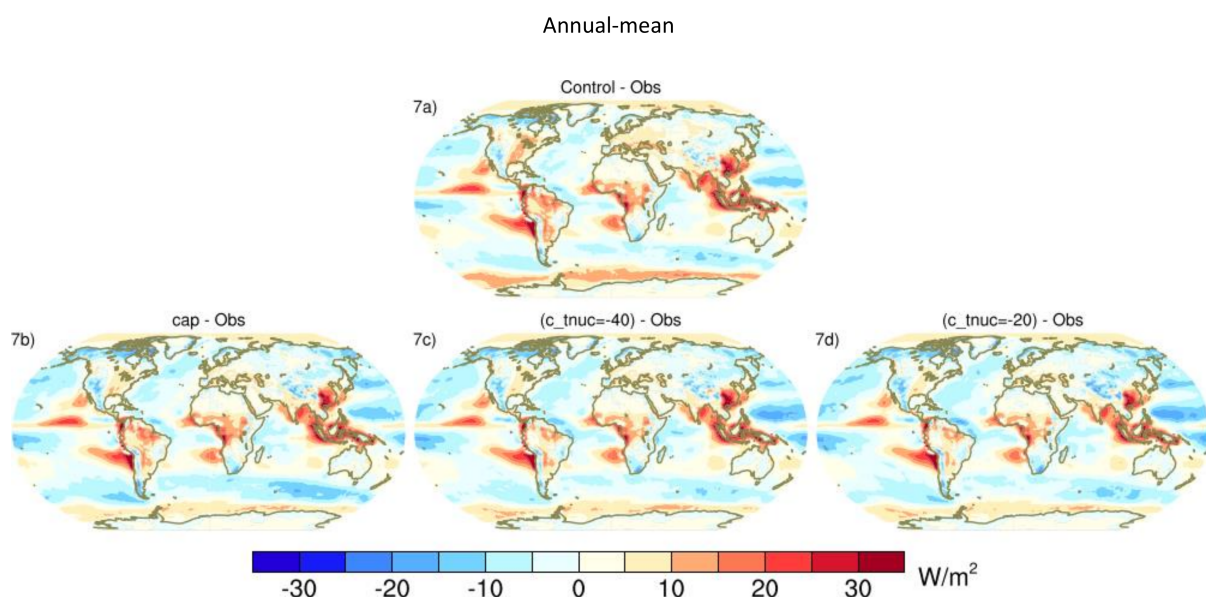


Figure 7. Spatial distributions of annual mean SW CRE anomaly at TOA in different model simulations with respect to the CERES TOA data. (7a) = control - Obs ; (7b) = *cap* - Obs, (7c) = (*c_tnuc*=-40) - Obs; (7d) = (*c_tnuc*=-20) - Obs. Model data calculated from daily-mean output over 20-22 years. Observational (CERES EBAF TOA) data consist of monthly mean values covering the period 2000-2018