

## Response to Reviewer #2:

Our apologies for missing comments 6 and 7 in our original response. (In converting the pdf to a word file we could copy from these were unintentionally lost by us.) We thank the reviewer for considering the revisions we have already made. Below we provide a response for both comments 6 and 7. Our responses are in **blue** and the initial reviewer comments in **black**.

1. Title should bear the word “laboratory” since the observed results might not be applicable to real world fires. For e.g., something like: “Laboratory-based biomass burning particles from a wide variety of fuels: Part 1: Properties of primary particles”

We understand the reviewers point here. However, we argue that there is nothing in the current title to indicate that these are “real-world” particles. (We assume by “real-world” the reviewer specifically means “ambient.” The fires we sampled from were most assuredly “real.”) In noting that the measurements are for “a wide variety of fuels” we believe it is implied that this is a laboratory experiment. Similarly, this is the case when we note the focus on “primary particles” as ambient studies have had limited success at accessing truly “primary” particles given the logistical challenges of getting close enough to fires and the rapid aging that occurs in the atmosphere. Instead of adjusting the title, we have updated the first sentence of the abstract to include the word “laboratory.”

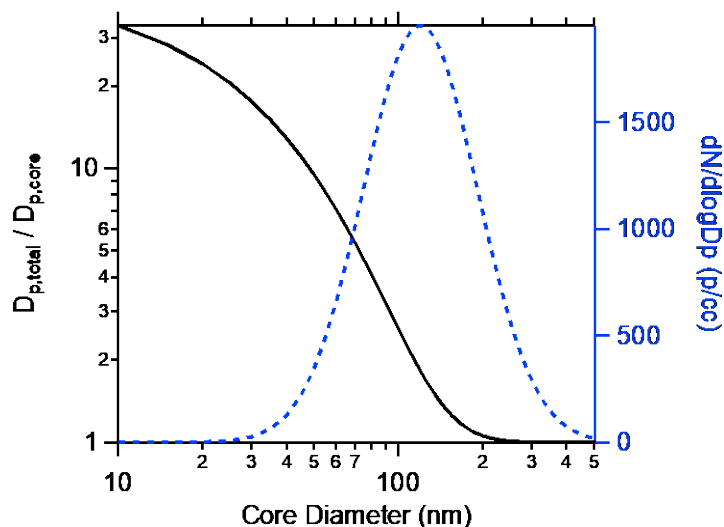
2. The authors provide no explanation (beyond a hand waving argument) to back the statement “The contribution of coating-induced enhancements (i.e. lensing effects) to absorption by black carbon are shown to be negligible for all conditions”. Lensing or focusing of light to the core could also be possible with weakly light-absorbing coating materials such as brown carbon with low imaginary index of refraction. Such a coating would facilitate lensing in addition to itself absorbing. One convincing way to declare that “no lensing” takes place is by looking at the internal field strength of a brown-carbon coated BC aggregate (see methodology in Chakrabarty and Heinson, Phys. Rev. Lett, 2018). I would like to see such a rigorous analysis performed (DDA or T-Matrix and not Mie-based core-shell) by the authors on a few BC aggregates coated with BrC vs non-refractory materials and convincing the reviewer and the community if indeed the “no lensing” claim is valid. If the authors cannot perform such an analysis, then I suggest that they remove all statements from the abstract and the main text regarding “negligible coating-induced enhancements (lensing effects)”. Instead, rephrase or replace the sentences with “brown carbon-coated BC particles yield absorption enhancements of x and y values. . .”

This particular comment follows directly from the Reviewer’s comment #3, questioning whether our conclusion that the “lensing” enhancement is “negligible.” However, here the reviewer (i) asks for calculations including BrC and (ii) indicates that if we cannot do such calculations we should entirely revise the discussion of the absorption enhancement. We address the question of whether the lensing-induced absorption enhancement is “negligible” extensively in our original response to the Reviewer’s comment #3. While we concluded that use of the “negligible” is justified, we have nonetheless revised the manuscript to indicate it as “minor”. (We have additionally gone through now and done a search for “negligible,” replacing it with “minor.”) Here, we address more specifically the utility or appropriateness of inclusion of calculations such as those suggested. In short, we do not think that including results of such

calculations would add substantial value to the manuscript. One reason is that we do not have knowledge of the internal morphology of the BC-containing particles. Given this absence of knowledge, it would be necessary for us to perform systematic calculations over different particle geometries and coating amounts to glean insights. However, such calculations have already been performed by many researchers and results are available in the literature, both for non-absorbing and weakly absorbing coatings (Adachi et al., 2010;Kahnert and Devasthale, 2011;Kahnert et al., 2012;Liu et al., 2015;Kahnert, 2017;Liu et al., 2017a;Zhang et al., 2017;Kanggiesser and Kahnert, 2018;Zhang et al., 2018). Performing new calculations here is therefore, in our opinion, unnecessary. Instead, we believe it sufficient to point to the literature, which we have done by citing the relatively early work by Adachi et al. (2010), where it is demonstrated that non core-shell internal particle morphologies can lead to “lensing” enhancements notably smaller than the core-shell solution.

Further, we reiterate that we measure the mass-weighted mean coating-to-core ratio. The actual population of particles sampled have a diversity of coating-to-core ratios. As has been established previously (Cappa et al., 2012;Fierce et al., 2016;Liu et al., 2017b), when the actual diversity of coating-to-core ratios across the population is accounted for smaller enhancements are obtained compared to that obtained assuming all particles have the mean coating-to-core ratio. For particularly diverse populations of particles, the population-weighted enhancement can be substantially less than that obtained assuming the mean coating-to-core ratio for all particles (Fierce et al., 2016). Thus, a single, or even a “few” calculations for some arbitrary particle morphologies would not, in our opinion, be especially informative. Instead, it would be necessary to perform calculations over a wide range of assumed particle populations and assumed internal morphologies to capture the likely true behavior. Such an exercise is well-beyond the scope of this manuscript, which is predominately observational, and further challenged by the lack of specific information regarding the internal particle morphology or the distribution of coating thicknesses across the particle population. Yet, we can give two examples here to illustrate the general conclusions we are likely to reach should we perform such calculations. Example 1: Assume two 130 nm diameter BC particles. One has a coating-to-core ratio of zero (bare BC) while the other has a coating-to-core ratio of 20. The mean coating to core ratio is 10. The core-shell Mie enhancement for the first particle is 1 while for the second is 2.45. So, the mean enhancement is 1.72, notably lower than an enhancement of 2.1 if the coating-to-core ratio for both particles were 10. While this enhancement is still quite a bit larger than 1, this illustrates the point that consideration of particle diversity generally leads to reductions in the enhancement. Example 2: Assume a log normal distribution of BC cores (mode diameter = 120 nm, width = 1.6). Assume also that the coating amount decreases with core size, consistent with condensational growth, with an (arbitrary) functional form illustrated in the following figure. Set the coating thickness distribution such that the mean coating-to-core ratio is 5, a “thick” value typical of our observations. The resulting coating-induced enhancement is 1.35 for the distribution, compared to a value of 1.79 if all particles had the same coating coating-to-core ratio. Again, a substantial reduction is observed for the more diverse (variable coating-to-core ratio) population compared to the constant coating case. Other functional forms for the coating distribution can be chosen, which would give different—albeit consistent—results, reflecting the particular nature of the assumed distribution; this includes distributions that yield even smaller ensemble-average enhancements for the same coating-to-core ratio. We can include the results of calculations of this sort in the manuscript. However, we do not feel that they are appropriate here given that we do not know the coating distribution. Additionally, we note that we have

assumed a core-shell morphology here, which generally yields larger enhancements compared to more detailed calculations that account for internal particle morphology.



Ultimately, given what the literature says about the importance of considering both internal morphology and of particle populations we believe it is a reasonable inference (and not a “hand waving argument”) to conclude that our observations point to the lensing-induced absorption enhancement being “negligible for all conditions.” We further believe that citation of the existing literature is preferred over the addition of calculations here. Nonetheless, as noted above, we have revised the statements in the abstract and throughout the manuscript as follows to make clearer that (i) we refer to results from our experiments and (ii) that the effect is “minor” rather than “negligible.” Specifically, for the abstract we have revised to state: “The contribution of coating-induced enhancements (i.e., “lensing” effects) to BC absorption are shown minor for all of the burns despite some burns producing particles have large ensemble-average coating-to-core mass ratios.”

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# 1 Biomass-burning derived particles from a wide variety of 2 fuels: Part 1: Properties of primary particles

3 Crystal D. McClure<sup>1</sup>, Christopher Y. Lim<sup>2,\*</sup>, David H. Hagan<sup>2</sup>, Jesse H. Kroll<sup>2</sup>, Christopher D.  
4 Cappa<sup>1,3,\*</sup>

5 <sup>1</sup> Department of Civil and Environmental Engineering, University of California, Davis, CA 95616

6 <sup>2</sup> Department of Civil and Environmental Engineering, Massachusetts Institute of Technology,  
7 Cambridge, MA, USA

8 <sup>3</sup> Atmospheric Sciences Graduate Group, University of California, Davis, CA, USA 95616

9 <sup>\*</sup> Now at Department of Chemistry, University of Toronto, Ontario, Canada

10 <sup>\*</sup> To whom correspondence should be addressed: cdcappa@ucdavis.edu

## 11 ABSTRACT

12 Relationships between various optical, physical, and chemical properties of biomass  
13 combustion derived particles are characterized for particles produced **in the laboratory** from a wide  
14 range of fuels and burn conditions. The modified combustion efficiency (MCE), commonly used  
15 to parameterize biomass particle emissions and properties, is shown to generally have weak  
16 predictive capabilities, especially for more efficient combustion conditions. There is, however, a  
17 strong relationship between many intensive optical properties (e.g. single scatter albedo, Ångstrom  
18 absorption exponent, mass absorption efficiency) and the organic aerosol-to-black carbon  
19 ([OA]/[BC]) mass ratio over a wider range than previously considered (0.3 to 10<sup>5</sup>). The properties  
20 of brown carbon (BrC, i.e. light absorbing organic carbon) also vary with [OA]/[BC]. ~~Coating-~~  
21 induced enhancements (i.e., “lensing” effects) **contribute only a minor amount** to BC absorption  
22 ~~for all of the burns despite some burns producing particles having large ensemble-average coating-~~  
23 ~~to-core mass ratios.~~ The BC-OA mixing state varies strongly with [OA]/[BC]; the fraction of OA  
24 that is internally mixed with BC decreases with [OA]/[BC] while the relative amount of OA coated  
25 on BC increases. In contrast, there is little relationship between many OA bulk chemical properties  
26 and [OA]/[BC], with the O:C and H:C atomic ratios and the relative abundance of a key marker  
27 ion ( $m/z = 60$ , linked to levoglucosan) all showing no dependence on [OA]/[BC]. In contrast, both  
28 the organic nitrate fraction of OA and the OA volatility do depend on the [OA]/[BC]. Neither the  
29 total particle or BC-specific size distributions exhibit any clear dependence on the burn conditions  
30 or [OA]/[BC], although there is perhaps a dependence on fuel type. Overall, our results expand on

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34 existing knowledge to contribute new understanding of the properties of particles emitted from  
35 biomass combustion.

## 36 **1 Introduction**

37 While it is understood that both open and controlled biomass combustion are major sources of  
38 particles to the atmosphere (Andreae and Merlet, 2001), questions remain regarding the properties  
39 of the emitted particles, their relationship with combustion conditions and fuel type, and their  
40 atmospheric evolution. Particles emitted from biomass combustion impact the global radiation  
41 budget and contribute to poor air quality in impacted regions. The emitted primary particles are  
42 primarily composed of organic aerosol (OA) and black carbon (BC), in varying amounts, with  
43 trace inorganic species (Reid et al., 2005;McMeeking et al., 2009;Levin et al., 2010). Particle  
44 intensive properties are often compared against the modified combustion efficiency ( $MCE \sim$   
45  $\Delta[CO_2]/(\Delta[CO]+ \Delta[CO_2])$ ), which provides a measure of the combustion efficiency of a burn. For  
46 example, various particle properties show some relationship with MCE, but often these  
47 relationships are weak, especially for more efficient combustion (higher MCE, corresponding  
48 typically to flaming conditions) (McMeeking et al., 2009;Liu et al., 2013;McMeeking et al., 2014).  
49 Understanding the diversity in the chemical, physical, and optical properties of the emitted  
50 particles is important for establishing the fire- or region-specific emissions and subsequent  
51 impacts.

52 The emitted OA from biomass combustion is somewhat light absorbing (Kirchstetter et al.,  
53 2004). Absorbing OA is commonly referred to as brown carbon (BrC), with properties that appear  
54 to depend on the fuel and combustion conditions (Saleh et al., 2014;Laskin et al., 2018), which  
55 affect particle organic composition (Jen et al., 2019). However, the properties of primary BrC  
56 absorption and, especially, understanding of the relationships between BrC absorption and other  
57 particle properties and burn conditions is only beginning to be unraveled. Additionally, it is  
58 established from theory and laboratory experiments that non-absorbing coatings on black carbon  
59 and other strongly absorbing particles can enhance the absorption (commonly referred to as the  
60 “lensing” effect but more accurately termed here the coating-induced enhancement) (Fuller et al.,  
61 1999;Bond et al., 2006;Lack et al., 2009;Shiraiwa et al., 2010;Cappa et al., 2012). Yet, the extent  
62 to which coating-induced enhancements impact absorption by ambient particles or for mixed-

63 component particles from complex sources, such as biomass burning, remains contentious (Cappa  
64 et al., 2012;Healy et al., 2015;Liu et al., 2015;Peng et al., 2016;Liu et al., 2017).

65 Here, we expand on current understanding of the relationships between various primary  
66 particle properties and burn conditions by analyzing measurements of primary biomass burning  
67 particles produced from combustion of a variety of fuel types, many of particular relevance to the  
68 western U.S.. We demonstrate that various optical properties exhibit a strong relationship with the  
69 [OA]/[BC] mass ratio, much stronger than their relationship with the MCE. We use the  
70 measurements to quantify the individual contributions of BC, BrC and from internal mixing of BC  
71 to the observed light absorption, and examine the variability in the properties of BrC specifically.  
72 We uniquely characterize the mixing state of BC and OA, and how mixing state vary between  
73 individual burns and depend on the mean properties of the emitted particles. We characterize the  
74 variability of OA-specific properties, including OA volatility, bulk chemical composition  
75 (characterized by the O:C and H:C atomic ratio, and the presence of key marker ions), and,  
76 uniquely, the relative abundance of organic nitrate species. We also examine the variability in the  
77 emitted particle size distribution, both for the total particles and for the BC particles specifically.  
78 Some of our analysis serves to support and extend previously determined relationships by  
79 considering a wider range of conditions, while other aspects are unique to this study. These  
80 observations provide a foundation for understanding and interpretation of experiments on the  
81 influence of photochemical aging on biomass particle properties, discussed in a related paper (Lim  
82 et al., 2019).

## 83 **2 Methods**

84 All experiments were conducted during the Fire Influence on Regional to Global Environments  
85 Experiment (FIREX) lab study, which took place at the Missoula Fire Sciences Lab in Missoula,  
86 MT, USA during November, 2016. Numerous types of biomass were combusted in a large  
87 chamber (12 x 12 x 19 m) and the smoke sampled to provide information on the physical, chemical,  
88 and optical properties of the resulting smoke (i.e., particulate and gas emissions). The general fuels  
89 types combusted included (exclusively or in combination): duff, dung, excelsior, straw, litter,  
90 untreated lumber, rotten debris, woody debris, shrub, herbaceous, and canopy biomass. A complete  
91 list of fuels and types is provided in **Table S1**, with further details available on the U.S. National  
92 Oceanic and Atmospheric Administration (NOAA) data archive

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94 (<https://esrl.noaa.gov/csd/projects/firex/>). All data used in this publication are also available on the  
95 NOAA archive, with the processed data summarized in complementary data repository (Cappa et  
96 al., 2019a).

97 Both “room” and “stack” burns were conducted, although here we include results only from  
98 stack burns. During stack burns, the smoke was mixed with background room air and funneled up  
99 a large cylindrical stack (2 m dia. x 15 m height) where it was sampled into a high-flow transfer  
100 line at ca. 0.27 m<sup>3</sup>/s. This flow rate corresponded to sampling approximately 10% of the stack  
101 flow. Smoke was transferred to an adjacent room via the high-flow transfer line (residence time  
102 ca. 2 s) where it was sub-sampled through a PM<sub>2.5</sub> cyclone and injected into a 0.25 m<sup>3</sup> Teflon  
103 photochemical reaction chamber (the mini chamber). Details on the construction and operation of  
104 the mini chamber can be found in (Lim et al., 2019). Here, we focus exclusively on the properties  
105 of particles sampled prior to initiation of photochemical oxidation; results of the photochemical  
106 oxidation experiments are discussed in a series of papers (Coggon et al., 2019; Lim et al., 2019).  
107 In brief, prior to each burn, the chamber was flushed with clean air with a relative humidity (RH)  
108 of approximately 40%. To fill the chamber, smoke was sub-sampled from the high-flow inlet and  
109 injected across the entire burn (typically lasting for 10-20 minutes) or until the chamber  
110 concentration reached a maximum. A suite of instruments sampled from the mini chamber at a  
111 flow rate of approximately 6 lpm. This flow rate varied from burn to burn due to the exact suite of  
112 instruments sampling. Clean makeup air was being injected simultaneously from a zero air  
113 generator to equal the air being sampled out of the chamber. The sampled smoke was diluted by a  
114 factor of ca. seven relative to the air in the high-flow inlet. Subsequent dilution after filling was  
115 characterized by the decay of acetonitrile (ACN). Properties of the primary particles are averaged  
116 over the 5-10 minute period after filling but before the initiation of photochemistry.

117 Particle-phase instrumentation sampled alternately every two minutes through a  
118 thermodenuded or ambient sample line. The thermodenuder was operated at 300 °C with a  
119 residence time of approximately 5 s and volatilized semi-volatile components, including those that  
120 are internally mixed with BC. The ambient line was lined with a charcoal cloth that removed excess  
121 gases (such as VOCs, NO<sub>x</sub>, and O<sub>3</sub>) that could interfere with particle-phase measurements.  
122 Comparison of thermodenuded versus ambient particles allowed for the investigation of coating  
123 amount and volatility. The gas-phase composition in the mini chamber was similar to that sampled  
124 directly from the fire (Koss et al., 2018; Lim et al., 2019). Particle phase instrumentation included:



125 a multi-wavelength cavity-ringdown-photoacoustic absorption spectrometer (CRD-PAS) and a  
126 photoacoustic absorption spectrometer (PASS-3) for characterization of light absorption and  
127 extinction coefficients at 405 nm, 532 nm, and 781 nm; a high resolution aerosol mass  
128 spectrometer (HR-ToF-AMS) for characterization of non-refractory submicron particulate matter  
129 (NR-PM<sub>1</sub>) components (i.e. OA, NO<sub>3</sub>, SO<sub>4</sub>, NH<sub>4</sub>, Cl, K); a soot photometer AMS (SP-AMS) in  
130 laser-only mode for characterization of refractory BC and the NR-components that are internally  
131 mixed with BC; a single particle soot photometer (SP2) for characterization of refractory BC mass  
132 concentrations and size distributions; and a scanning electrical mobility sizer (SEMS) for  
133 measurement of particle mobility size distributions. Further details regarding instrument operation  
134 and calibration are provided in the Supplemental Material and in Lim et al. (2019).

135

### 136 **3 Results and Discussion**

#### 137 **3.1 Bulk optical property relationships**

138 Due to the wide variety of biomass fuels and types used during FIREX, there was a substantial  
139 diversity in the properties of primary particles produced. Previous studies have shown both the  
140 single scatter albedo (SSA) and wavelength-dependence of absorption (the absorption Angstrom  
141 exponent, AAE) depend on the modified combustion efficiency (MCE) (Liu et al.,  
142 2013;McMeeking et al., 2014;Pokhrel et al., 2017). The MCE is defined here as:

$$143 \quad MCE = \frac{[CO_2]}{[CO_2]+[CO]} \quad (1)$$

144 The SSA is defined as:

$$145 \quad SSA = \frac{b_{ext} - b_{abs}}{b_{ext}} \quad (2)$$

146 where  $b_{ext}$  is the wavelength-specific extinction coefficient and  $b_{abs}$  is the wavelength-specific  
147 absorption coefficient. The AAE is defined as:

$$148 \quad AAE = -\log\left(\frac{b_{abs,\lambda_1}}{b_{abs,\lambda_2}}\right) / \log\left(\frac{\lambda_1}{\lambda_2}\right) \quad (3)$$

149 where  $\lambda_1$  and  $\lambda_2$  indicate two different wavelengths, here 405 nm and 532 nm. The MCE  
150 characterizes the overall combustion efficiency, with values closer to unity indicating more

151 complete combustion. In general, higher MCE correspond to more flaming combustion conditions  
152 while smaller MCE correspond to more smoldering conditions. We find a similar relationship  
153 between  $SSA_{405nm}$ , AAE, and  $[OA]/[BC]$  with MCE as previous studies (**Figure 1**) (McMeeking  
154 et al., 2009;Liu et al., 2013;McMeeking et al., 2014;Pokhrel et al., 2017). Specifically, the  
155  $SSA_{405nm}$  is relatively constant and near unity for  $MCE < \sim 0.9$ , but above this value exhibits a  
156 rapid decline, albeit with a substantial amount of scatter (**Figure 1a**). The AAE is also relatively  
157 constant when  $MCE < 0.9$ , with very large values ( $AAE \sim 8$ ). There is a rapid, scattered decrease  
158 in the AAE as MCE increases further (**Figure 1b**). The relationship between  $[OA]/[BC]$  and MCE  
159 is similar, with values generally decreasing as MCE increases but a large amount of scatter (**Figure**  
160 **1d**). There is also a general relationship between the mass absorption coefficient referenced to BC  
161 ( $MAC_{BC}$ ) at 405 nm and the MCE, but with similar scatter as the other properties (**Figure 1c**). The  
162  $MAC_{BC}$  is defined as:

$$163 \quad MAC_{BC} = b_{abs}/[BC] \quad (4)$$

164 The  $MAC_{BC,405nm}$  includes contributions from absorption by BC, BrC, and from coating-induced  
165 enhancement of BC absorption. These results, along with the literature, indicate that MCE can  
166 provide guidance as to the general magnitude of these particle properties, but that the MCE is  
167 ultimately a fairly imprecise metric, especially for the  $SSA_{405nm}$ .

168 However, we find a very strong relationship between the  $SSA_{405nm}$  and the total  $[OA]/[BC]$   
169 ratio (**Figure 1e**). This is consistent with the findings of Pokhrel et al. (2016), who observed  
170 something similar but over a smaller range of  $[OA]/[BC]$ . (Similarly strong relationships are  
171 observed for SSA values at 532 nm and 781 nm (**Figure S1**), or if the  $[NR-PM_1]/[BC]$  are used as  
172 OA averages 95% of the total NR-PM<sub>1</sub> mass.) Smaller  $[OA]/[BC]$  correspond to smaller  $SSA_{405nm}$   
173 values with a sigmoidal relationship observed. (Fit parameters for all fits shown are provided in  
174 **Table S1**.) There is similarly a very strong, sigmoidal relationship between the AAE and  
175  $MAC_{BC,405nm}$  and  $[OA]/[BC]$  (**Figure 1f,g**). The large increase in the  $MAC_{BC,405nm}$  indicates that  
176 BrC contributes substantially to the total absorption. The contributions of coating-induced  
177 enhancements and of BrC are discussed further in Sections 3.4.1 and 3.4.2. The larger range of  
178  $[OA]/[BC]$  and the greater number of individual burns considered here, compared to Pokhrel et al.  
179 (2016), allows for determination of more robust fits. Pokhrel et al. (2017) found that the absorption

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189 enhancement at 405 nm, determined from thermodenuder measurements, increased with  
190 [OA]/[BC] up to [OA]/[BC] ~33 (the largest value reported), consistent with our findings.

191 These observations demonstrate that the optical properties of the primary particles depend on the  
192 relative amount of OA versus BC. This is as expected because OA is generally more scattering,  
193 compared to BC, and light absorbing OA (aka BrC) typically exhibits a much stronger wavelength  
194 dependence than BC. Based on these relationships, we divide the individual burns into different  
195 classes (**Table 1**). We have chosen to classify particles based on the observed  $SSA_{405nm}$  values;  
196 use of [OA]/[BC] for classification yields largely similar results, given the strong relationship  
197 between the two. The dividing lines between classes are selected to yield six classes that span the  
198 entire range of  $SSA_{405nm}$  values, from 0.23 (Class 1) to 0.97 (Class 6), with approximately equal  
199 numbers of individual burns in each class (ca. 8-10). Partitioning the observations into different  
200 particle classes facilitates interpretation of the photochemical evolution of the particles, to be  
201 discussed in future work. In addition, we find that use of the Class average properties versus MCE  
202 generally provides more representative fits to the observations (visually apparent in **Figure 1**, and  
203 supported by the reduced  $\chi^2$  for the fits).

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### 204 3.2 OA composition and volatility

205 Variability in the bulk composition of the OA is characterized by the O:C and H:C atomic  
206 ratios and the fractional abundance ( $f_x$ ) of two marker ions,  $m/z = 44$  and  $m/z = 60$ . The  $f_{44}$  is  
207 complementary to O:C and larger values generally indicate a greater degree of oxygenation and  
208 the presence of carboxylic acids. The  $f_{60}$  is often taken as a marker ion for biomass burning, in  
209 particular a signature of levoglucosan and similar molecules (Schneider et al., 2006;Alfarra et al.,  
210 2007). The high resolution ion  $C_2H_4O_2^+$  contributes to and exhibits similar behavior as  $f_{60}$ ; the  
211 slope for  $f_{C_2H_4O_2^+}$  against  $f_{60}$  is 0.98. While it is known that properties such as  $f_{60}$  vary in different  
212 biomass burning samples (Schneider et al., 2006) or between near-source intercepts of different  
213 ambient plumes (Garofalo et al., 2019), the specific dependence on burn conditions or overall  
214 particle composition (e.g. [OA]/[BC]) has not been systematically explored to our knowledge.

215 The average  $f_{60} = 0.022 \pm 0.01$  ( $1\sigma$ ). The  $f_{60}$  values vary non-monotonically with [OA]/[BC],  
216 exhibiting a slight increase from Class 1 to Class 3 and then a decrease from Class 4 to Class 6  
217 (Figure 2a). This indicates that, while  $f_{60}$  is overall a useful marker ion for biomass burning, it

220 cannot be used to distinguish between different burn conditions. The  $f_{44}$  generally decreases with  
221 [OA]/[BC] (Figure 2b;  $r^2 = 0.33$ .) However, the average  $f_{44}$  values for particle Classes 2-5 differ  
222 negligibly, suggesting that  $f_{44}$  might be useful in discriminating between extreme cases (e.g. Class  
223 1 versus Class 6), but that it is of limited general use in distinguishing between burn conditions  
224 and fuel types. The O:C atomic ratio (average =  $0.37 \pm 0.09$ ) exhibits similar behavior—expected  
225 as  $f_{44}$  is generally related to O:C (Aiken et al., 2008)—with a general decrease as [OA]/[BC]  
226 increases, although a comparably weaker correlation (Figure 2c;  $r^2 = 0.17$ ). The H:C (average =  
227  $1.76 \pm 0.05$ ) exhibits a weak, positive correlation with [OA]/[BC], although the variability is slight  
228 (Figure 2d;  $r^2 = 0.27$ ).

229 The mass fraction of the OA that is composed of nitrated organics ( $f_{\text{ON-OA}} = [\text{ON}]/[\text{OA}]$ ) was  
230 determined using the HR-ToF-AMS measurements and the method of Kiendler-Scharr et al.  
231 (2016) (see the Supplemental Material for further details). The terminology nitrated organics (ON)  
232 includes contributions from both nitro and nitrate functional groups. The fraction of measured  
233 nitrate that was ON ( $f_{\text{ON-N}} = [\text{ON}]/([\text{ON}]+[\text{NO}_3^-])$ ) decreased with [OA]/[BC] and ranged from 0.91  
234 (Class 1) to 0.48 (Class 6) (Figure S2a). The Class-specific average  $f_{\text{ON-OA}}$  also decreased with  
235 [OA]/[BC], although by a much greater extent than the  $f_{\text{ON-N}}$ , ranging from 6.0% (Class 1) to 0.27%  
236 (Class 6) and (Figure 2e). There is a reasonably linear relationship between  $\log(f_{\text{ON-OA}})$  and  
237  $\log([\text{OA}]/[\text{BC}])$  ( $r^2 = 0.47$ ). This indicates that a larger proportion of ON species and  
238 functionalities are produced when particles are, on average, more BC-rich. This does not reflect  
239 differences in fuel nitrogen content as there is no relationship between fuel N and  $f_{\text{ON-OA}}$  (Figure  
240 S2b). Therefore, it seems that the relationship between  $f_{\text{ON-OA}}$  and [OA]/[BC] is related more so to  
241 the burn conditions than the fuel N content, although as with many other properties the relationship  
242 with [OA]/[BC] is clearer than with the MCE (Figure S2c).

243 The OA volatility is characterized as the ratio between the OA concentration after  
244 thermodenuding to that without thermodenuding (the mass fraction remaining,  $\text{MFR}_{\text{OA}}$ ). The  
245  $\text{MFR}_{\text{OA}}$  decreases as [OA]/[BC] increases (Figure 2f), indicating that the OA at lower [OA]/[BC]  
246 is less volatile than the OA at higher values. This observation provides support for the proposal by  
247 Saleh et al. (2014) that less volatile, more absorbing species are preferentially formed under  
248 conditions where BC formation is favored, discussed further in Section 3.4.2. The relationship  
249 between  $\text{MFR}_{\text{OA}}$  and [OA]/[BC] is reasonably described by an exponential function.

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### 255 3.3 BC Mixing State

256 As discussed above, the relative amounts of OA and BC vary greatly between fuel types and  
257 combustion conditions. However, the distribution of BC and OA between particles, and how this  
258 varies between very different burn conditions, has not been previously explored in detail to our  
259 knowledge. The bulk average fraction of OA that is internally mixed with BC versus OA that is  
260 externally mixed from BC is determined using the HR-ToF-AMS and SP-AMS measurements.  
261 The HR-ToF-AMS quantifies OA independent of mixing state, whereas the SP-AMS (as operated  
262 here) quantifies only the OA that is internally mixed with BC. The fraction of OA that is internally  
263 mixed with BC ( $f_{OA,int}$ ) is:

$$264 f_{OA,int} = \frac{[OA]_{SP-AMS}}{[OA]_{HR-ToF-AMS}} = \frac{[OA]_{int}}{[OA]_{tot}} \quad (5)$$

265 where the subscript *int* indicates the OA that is internally mixed with BC and the subscript *tot*  
266 indicates the total OA. The  $f_{OA,int}$  should range from 0 to 1. Related, the SP-AMS quantified the  
267 ratio between the OA that is internally mixed with BC and the BC concentration, referred to here  
268 as  $[OA]_{int}/[BC]$ . We find that  $f_{OA,int}$  decreases substantially as  $[OA]/[BC]$  increases, ranging from  
269  $f_{OA,int} = 0.4$  for Class 1 (low SSA) particles to  $f_{OA,int} = 0.01$  for Class 6 (high SSA) particles (**Figure**  
270 **3a**). The data are well-fit by a sigmoidal function. However, the amount of OA coating BC ( $R_{OA-BC}$   
271  $= [OA]_{int}/[BC]$ ) increases with the total  $[OA]/[BC]$ , also with a sigmoidal relationship (**Figure 3b**).  
272 Thus, while a smaller fraction of the total OA is internally mixed with BC for larger total  
273  $[OA]/[BC]$  the amount of OA that coats BC increases. Most likely this behavior reflects that BC  
274 and OA are generated with different efficiencies in different parts of the combusting biomass. BC  
275 is more efficiently generated from flaming combustion while OA is more efficiently generated  
276 from smoldering combustion. These observations demonstrate that the extent to which  
277 atmospheric models can assume that all OA is internally mixed with or externally mixed from BC  
278 at the point of emission will depend on the combustion conditions.

### 279 3.4 Absorption enhancement and brown carbon

#### 280 3.4.1 Coating-induced absorption enhancement

281 Non- or weakly-absorbing coatings on black carbon particles can theoretically increase the  
282 absorption by BC (Fuller et al., 1999; Bond et al., 2006), an effect which has been confirmed by

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285 laboratory experiments (Lack et al., 2009; Shiraiwa et al., 2010; Cappa et al., 2012). The extent to  
286 which coatings on BC actually enhance absorption by BC in the atmosphere remains unclear. Some  
287 studies indicate minor coating-induced enhancements while others indicate substantial  
288 enhancements (Cappa et al., 2012; Healy et al., 2015; Liu et al., 2015; Peng et al., 2016; Zhang et al.,  
289 2016; Liu et al., 2017; Cappa et al., 2019b). Understanding the nature of the coating-induced  
290 enhancement is important for quantifying the radiative impacts of BC (Jacobson, 2001; Bond et al.,  
291 2013). Further, these coating-induced absorption enhancements ( $E_{\text{abs.coat}}$ ) complicate the  
292 determination of brown carbon (BrC) absorption and the two must be separated. Here, we examine  
293 the extent to which coatings on BC for primary biomass burning particles enhance the BC  
294 absorption. Theoretically, the magnitude of  $E_{\text{abs.coat}}$  for an individual particle depends primarily on  
295 the coating thickness and secondarily on the size of the BC core (Bond et al., 2006; Fuller et al.,  
296 1999). Thus, the extent to which coatings enhance BC absorption for a given situation can be  
297 assessed through the relationship between the observed  $MAC_{\text{BC}}$  and the coating-to-core mass ratio  
298 ( $R_{\text{coat-rBC}} = [\text{NR-PM}]_{\text{int}}/[\text{BC}]$ , where *int* indicates that the coating material is internally mixed with  
299 BC). The expectation is that the  $MAC_{\text{BC}}$  increases with  $R_{\text{coat-BC}}$ .

300 However, absorption by BrC can also lead to an apparent increase in the normalized absorption  
301 with  $R_{\text{BC}}$  if the BrC abundance correlates with the total coating amount. Because BrC absorbs more  
302 strongly at shorter wavelengths, the wavelength-dependence of the  $MAC_{\text{BC}}$  to  $R_{\text{BC}}$  relationship can  
303 be used to further separate the influence of coating versus BrC absorption. The  $MAC_{\text{BC}}$  exhibits a  
304 wavelength-dependent relationship with  $R_{\text{coat-rBC}}$  for fresh biomass particles (405 nm, 532 nm and  
305 781 nm) (Figure 4a-c). The  $MAC_{\text{BC}}$  increases notably with  $R_{\text{coat-rBC}}$  at 405 nm and to a lesser extent  
306 at 532 nm. At 781 nm the  $MAC_{\text{BC}}$  is essentially independent of  $R_{\text{coat-rBC}}$  up to  $R_{\text{coat-rBC}}$  values as  
307 large as 10, but does exhibit some increase at  $R_{\text{coat-rBC}} > 10$ . However, this is most likely a result  
308 of absorption by OA at 781 nm and not indicative of an increase in the coating-induced  
309 enhancement, discussed further below. The wavelength dependence provides clear evidence of  
310 BrC absorption at shorter wavelengths.

311 That the  $MAC_{\text{BC}}$  at 781 nm is nearly independent of  $R_{\text{coat-rBC}}$  up to such large  $R_{\text{coat-rBC}}$  values  
312 indicates that there is only a minor coating-induced enhancement for the primary biomass particles,  
313 the magnitude of which is discussed below. Our observations are consistent with McMeeking et  
314 al. (2014), who also investigated the relationship between the  $MAC_{\text{BC}}$  and  $R_{\text{coat-rBC}}$  for a primary  
315 biomass particles from multiple fuel types. Most likely, this lack of a substantial coating-induced

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317 enhancement results from a non-even distribution of non-BC mass across the population of BC  
318 particles (Fierce et al., 2016; Liu et al., 2017) and from the morphology of BC-containing particles  
319 not conforming to an idealized core-shell structure (Adachi et al., 2010). The influence of  
320 photochemical aging on the coating-induced enhancement will be examined in future work.

321 The relationship between  $MAC_{BC}$  and the coating amount ( $R_{coat-rBC}$ ) can be contrasted with the  
322 relationship between  $MAC_{BC}$  and the total  $[OA]/[BC]$ . At all three wavelengths the  $MAC_{BC}$  exhibit  
323 strong, sigmoidal relationships with  $[OA]/[BC]$  (Figure 4d-f). That  $MAC_{BC,781nm}$  exhibits such a  
324 clear relationship with  $[OA]/[BC]$  suggests that even the small apparent coating-induced  
325 enhancement, implied above from the very weak with  $R_{coat-rBC}$ , is largely driven by absorption by  
326 BrC rather than from the impact of coating on BC. Pokhrel et al. (2017) found that the absorption  
327 enhancement, determined from thermogravimetric measurements, increased notably with  $[OA]/[BC]$   
328 up to  $[OA]/[BC] \sim 33$  at 405 nm (the largest value reported by them), but by much less at 660 nm,  
329 consistent with our findings.

330 The observations allow for determination of wavelength-dependent  $MAC_{BC}$  values for pure BC  
331 ( $MAC_{BC,pure}$ ) for each wavelength by extrapolation of the  $MAC_{BC}$  versus  $[OA]/[BC]$  ratio to zero  
332 using sigmoid fits. Since the  $R_{coat-rBC}$  correlates reasonably with  $[OA]/[BC]$  (Figure 3b),  
333 extrapolation against  $[OA]/[BC]$  to zero effectively removes both contributions from BrC and any  
334 coating-induced enhancement. The derived  $MAC_{BC,pure}$  values are  $11.8 \text{ m}^2 \text{ g}^{-1}$  at 405 nm,  $8.8 \text{ m}^2$   
335  $\text{g}^{-1}$  at 532 nm and  $5.5 \text{ m}^2 \text{ g}^{-1}$  at 781 nm, with estimated fit-based uncertainties of  $\sim 10\%$ . The  
336 absolute uncertainties on the  $MAC_{BC,pure}$  are primarily dependent on the uncertainty in the  $b_{abs}$  and  
337  $[rBC]$  measurements, and are  $\sim 35\%$ . The derived  $MAC_{BC}$  values are very similar to those recently  
338 reported by Forestieri et al. (2018) for fresh BC particles:  $MAC_{BC,pure} = 11.9 \text{ m}^2 \text{ g}^{-1}$  at 405 nm and  
339  $8.8 \text{ m}^2 \text{ g}^{-1}$  at 532 nm, with an extrapolated value at 781 nm of  $5.7 \text{ m}^2 \text{ g}^{-1}$ . The value at 532 nm is  
340 somewhat higher than that suggested by Bond and Bergstrom (2006) ( $7.75 \text{ m}^2 \text{ g}^{-1}$  at 532 nm). Our  
341 derived  $MAC_{BC,pure}$  values yield an  $AAE = 1.17$ , determined from a fit to the three wavelengths. An  
342  $AAE$  close to unity indicates absorption is dominated by BC, as expected.

343 Values for the absorption enhancement at 781 nm are calculated as the ratio between the  
344 observed  $MAC_{BC}$  in Figure 4 and the derived  $MAC_{BC,pure}$ . The derived  $E_{abs}$  range from 0.96 to 27.  
345 Values greater than two occur only for the particles having particularly large  $[OA]/[BC]$ ,  $> 400$ .  
346 As  $E_{abs}$  values much greater than two at 781 nm are unlikely to result from mixing-induced

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350 enhancements (Chakrabarty and Heinson, 2018), this again suggests that the OA is somewhat  
 351 absorbing at this wavelength. For the burns where  $[OA]/[BC] < 400$ , the median  $E_{abs} = 1.14$  and  
 352 the arithmetic mean  $E_{abs} = 1.19 \pm 0.14$  ( $1\sigma$ ). Given that some of this enhancement may result from  
 353 BrC absorption at 781, these values can be considered upper-limits on  $E_{abs,coat}$ , and the small  
 354 magnitude is consistent with our conclusion above that, while likely greater than zero, the mixing-  
 355 induced enhancement is generally minor. It is possible that the  $E_{abs,coat}$  values when  $[OA]/[BC] >$   
 356 400 are substantially larger. However, given the general lack of a dependence of the  $MAC_{BC,781nm}$   
 357 for  $R_{BC-coat} < 10$  this seems unlikely.

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### 358 3.4.2 Primary brown carbon absorption

359 The absorption due to brown carbon is determined by difference as:

$$360 \quad b_{abs,BrC} = b_{abs,obs} - MAC_{BC,pure} \cdot [BC] \cdot E_{abs,coat} \quad (6)$$

361 where  $b_{abs,BrC}$  is the absorption due to BrC specifically. Importantly, the use of study-specific  
 362  $MAC_{BC,pure}$  values serves to reduce systematic biases in the  $b_{abs,BrC}$ , compared to direct use of  
 363 literature  $MAC_{BC,pure}$  values. Assuming  $E_{abs,coat} = 1$  provides an upper limit on the BrC absorption,  
 364 which we note is likely most appropriate for the particles sampled here, as discussed in the previous  
 365 section. Therefore, we use the upper-limit values throughout the analysis that follows, unless  
 366 otherwise stated. However, a lower limit for BrC absorption can be determined at 405 nm and 532  
 367 nm assuming that all of the enhancement at 781 nm results from coatings and not from BrC. The  
 368 resulting  $E_{abs,obs}$  ( $= MAC_{BC,obs}/MAC_{BC,pure}$ ) at 781 nm averages 1.19 for  $R_{BC-coat} < 10$ . Using  $E_{abs,coat}$   
 369  $= 1.19$  in Eqn. 7 yields a lower limit for the BrC absorption at the two shorter wavelengths,  
 370 appropriate since  $E_{abs,coat}$  generally has only a small wavelength dependence. A fit to the coating-  
 371 corrected (lower-limit) versus upper-limit  $b_{abs,BrC}$  yields a slope of 0.97 at 405 nm and 0.88 at 532  
 372 nm (**Figure S3**). The smaller difference at 405 nm results from the fractional contribution of BrC  
 373 to the total absorption being larger at this wavelength.

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374 Brown carbon-specific mass absorption coefficients ( $MAC_{BrC}$ ) are determined as the ratio  
 375 between  $b_{abs,BrC}$  and the total OA concentration:

$$376 \quad MAC_{BrC} = \frac{b_{abs,BrC}}{[OA]} \quad (7)$$



379 The  $MAC_{BrC}$  values from Eqn. 7 are bulk-average values, and do not account for different  
380 molecules and classes of molecules likely having different absorptivities. Uncertainties in the  
381  $MAC_{BrC}$  values are determined by error propagation. Similarly, an AAE value for just the brown  
382 carbon ( $AAE_{BrC}$ ) can be calculated using wavelength pairs as:

$$383 \quad AAE_{BrC} = -\log\left(\frac{b_{abs,BrC,\lambda_1}}{b_{abs,BrC,\lambda_2}}\right) / \log\left(\frac{\lambda_1}{\lambda_2}\right); \quad (8)$$

384 The geometric averages of the  $MAC_{BrC}$  values are  $0.76^{+0.65}_{-0.35} \text{ m}^2 \text{ g}^{-1}$ ,  $0.21^{+0.36}_{-0.13} \text{ m}^2 \text{ g}^{-1}$ ,  $0.056^{+0.15}_{-0.04}$   
385  $\text{m}^2 \text{ g}^{-1}$  at 405 nm, 532 nm and 781 nm, with uncertainties the  $1\sigma$  burn-to-burn variability. The  
386  $MAC_{BrC}$  values vary between classes, generally increasing as the [OA]/[BC] ratio decreases at all  
387 wavelengths (shown for 405 nm in **Figure 5a**). For example, the average  $MAC_{405nm} = 2.3 \pm 1 \text{ m}^2$   
388  $\text{g}^{-1}$  for Class 1 and  $0.35 \pm 0.09 \text{ m}^2 \text{ g}^{-1}$  for Class 6. Although the uncertainties on the derived  $MAC_{BrC}$   
389 increase substantially as [OA]/[BC] decreases—because BrC absorption contributes to a smaller  
390 extent at longer wavelengths—the observations nonetheless indicate that the BrC absorptivity  
391 depends on the combustion conditions. The relationship at 405 nm is well-described by a sigmoidal  
392 function in log-log space, with limiting values of  $0.35 \text{ m}^2 \text{ g}^{-1}$  at large [OA]/[BC] and  $11.2 \text{ m}^2 \text{ g}^{-1}$   
393 at small [OA]/[BC]. That the extrapolated zero [OA]/[BC] limit for  $MAC_{BrC}$  is similar to pure BC  
394 suggests an evolution of BrC towards having properties similar to BC when the overall [OA]  
395 content is small. Such behavior is consistent with Saleh et al. (2018), who argue that there is a  
396 continuum of BrC properties that depends on the combustion conditions, as demonstrated in that  
397 study for low-temperature benzene and toluene combustion. The range of the  $MAC_{BrC}$  values  
398 observed here, including that there is notable absorption at 781 nm, encompass many previous  
399 measurements, summarized in **Table S3**. This likely reflects the wide diversity of fuel types and  
400 burn conditions considered here, as exemplified by the very large range of [OA]/[BC].

401 Estimated values of the imaginary component of the refractive index for BrC ( $k_{BrC}$ ) are determined  
402 from Mie theory via optical closure (Zhang et al., 2016), assuming a real part of the refractive  
403 index of 1.5 and a particle diameter of 150 nm, a typical value for these experiments. Imaginary  
404 RI values are of use in atmospheric models for calculation of BrC absorption. There is a linear  
405 relationship between  $MAC_{BrC}$  and  $k_{BrC}$  (**Figure S4a**). Thus, the  $k_{BrC}$  exhibits a similar correlation  
406 with [OA]/[BC] as does the  $MAC_{BrC}$  (**Figure 5a**).

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411 The wavelength-dependence of absorption, i.e. the  $AAE_{405-532}$ , also varies with [OA]/[BC], in this  
412 case with a positive relationship between the two (Figure S**b**). The relationship is reasonably  
413 described by a sigmoidal function. This implies that, while the  $MAC_{BrC}$  varies inversely with  
414 [OA]/[BC] at all wavelengths, the exact variation is wavelength dependent. The  $AAE_{405-532}$   
415 relationship with [OA]/[BC] is well-described by a sigmoidal function (versus  $\log([OA]/[BC])$ ),  
416 with limiting values of 10.4 at large [OA]/[BC] and 1.3 at small [OA]/[BC]. The wavelength-  
417 dependence of the  $k_{BrC}$  ( $w_{BrC}$ ) are also calculated, to facilitate comparison with the literature, as:

$$418 \quad w_{BrC} = -\log\left(\frac{k_{BrC,\lambda 1}}{k_{BrC,\lambda 2}}\right) / \log\left(\frac{\lambda 1}{\lambda 2}\right) \quad (9)$$

419 The  $w_{BrC}$  exhibit a similar dependence on [OA]/[BC] as the  $AAE_{BrC}$ , as the  $w_{BrC}$  and  $AAE_{BrC}$  are  
420 linearly related, albeit with some scatter (Figure S4**b**;  $r^2 = 0.97$ ).

421 Our observations support the results of Saleh et al. (2014), who also found a relationship between  
422 the  $k_{BrC,405nm}$  and [OA]/[BC]. However, our analysis substantially extends the range of [OA]/[BC]  
423 values investigated in that work (they considered [OA]/[BC] from only ca. 2 to 170). In the overlap  
424 region between our two studies the  $k_{BrC,405nm}$  agree reasonably well over the range  $2 < [OA]/[BC]$   
425  $< 50$ , but the  $k_{BrC,405nm}$  from Saleh et al. (2014) are smaller than observed here above [OA]/[BC] =  
426 50. Importantly, our results demonstrate that the linear fit suggested by Saleh et al. (2014) for  
427  $MAC_{BrC}$  is only appropriate over the range of values they considered and that a sigmoidal provides  
428 for a more robust relationship over a wider range of [OA]/[BC]. Related, the wider range of  
429 [OA]/[BC] enables more robust determination of the functional dependence of the wavelength-  
430 dependence of absorption ( $w_{BrC}$ ), with overall larger  $w_{BrC}$  values and a larger plateau at high  
431 [OA]/[BC] compared to the fit by Saleh et al. (2014).

432 The  $MAC_{BrC}$  values also correlate with the nitrated organic fraction of OA, the latter of which, as  
433 noted above, also correlates with the [OA]/[BC] (Figure 6**a**). This observation suggests that  
434 organic nitrate and nitro functionalities may be at least somewhat responsible for the increase in  
435 absorption. Laskin et al. (2018) performed offline molecular level analyses of primary OA  
436 collected during FIREX. They found that nitroaromatics and N-containing polycyclic aromatic  
437 hydrocarbons (PAHs) contribute notably to the total light absorption by BrC, although there are  
438 many non-N-containing species that also contribute to BrC absorption. The variability between  
439 particle Classes is consistent with the results of Lin et al. (2016), which show that the abundance

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443 of N-containing chromophores varies between particles produced from different biomass fuels.  
444 Additionally Mohr et al. (2013) observed a relationship between the concentration of nitrated  
445 phenols and short-wavelength absorption by BrC, although it is possible that for their  
446 measurements these species were produced from chemical processing, as opposed to being directly  
447 emitted. Altogether, our results provide support for the idea that nitrated organic functionalities  
448 are an important contributor to BrC absorption. However, it is very likely that other functional  
449 groups also contribute to the total absorption.

450 The  $MAC_{BrC,405nm}$  exhibits an inverse correlation with the  $f_{60}/f_{44}$  ratio of the OA, although there is  
451 substantial scatter in the  $f_{60}/f_{44}$  ratio for a given particle class (Figure 6b). (The  $f_{44}$  and  $f_{60}$  have no  
452 discernable relationship.) The observed  $MAC_{BrC,405nm}$  relationship with  $f_{60}/f_{44}$  is opposite that  
453 reported by Lack et al. (2013) for ambient measurements of particles a biomass burning plume,  
454 who find a reasonable positive correlation. This difference in behavior results from our sampling  
455 primary particles directly—thereby focusing on the inherent variability in the properties of the  
456 emitted particles—while Lack et al. (2013) sampled ambient particles. For ambient sampling, the  
457 observed relationship will be sensitive to mixing of biomass burning particles with background or  
458 aged biomass particles, which are known to have a smaller  $f_{60}$  (Cubison et al., 2011). Thus, the  
459 relationship observed by Lack et al. (2013) can best be viewed as a mixing line between the fresh  
460 primary particles (having large  $MAC_{BrC,405nm}$  and large  $f_{60}/f_{44}$ ) and background or aged biomass  
461 particles (having small  $MAC_{BrC,405nm}$  and small  $f_{60}/f_{44}$ ), rather than providing information on the  
462 inherent variability in the absorptivity of the fresh particles.

### 463 3.5 Size distributions

464 Total particle mobility size distributions and BC-only size distributions were measured (Figure  
465 7). Primary particle size distributions are important parameters specified in regional and global  
466 models. The number-weighted and volume-weighted size distribution are generally described by  
467 either one or two log-normal modes for individual burns; a two-mode fit provides a more robust  
468 solution across all modes. The mass-weighted BC size distributions are similarly described by one  
469 or two log-normal modes. A fit to the average number-weighted distribution across all particle  
470 classes yields geometric median diameters ( $d_{p,N}$ ) and widths ( $\sigma_g$ ) of 60.3 nm and 1.76, respectively,  
471 for the smaller mode and 153 nm and 1.64 for the larger mode (Figure 8). The amplitude of the  
472 smaller mode is 4.6 times the larger mode. A single mode fit yields  $d_{p,N} = 68$  nm and  $\chi_g = 1.93$ ,

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476 although the fit is poorer. Mann et al. (2014) report  $d_{p,N}$  values used by a variety of global models  
477 for biofuels. The models tend to use either 80 nm or 150 nm, although a few use other values (30  
478 nm, 60 nm, 100 nm). Those using 80 nm typically use  $\sigma_g = 1.80$  while those using 150 nm typically  
479 use  $\sigma_g = 1.59$ , although there are exceptions. Our observations indicate that use of a bimodal  
480 distribution within models would be more representative, but that a single mode can do acceptably.  
481 We find that the volume-weighted distribution calculated from a single-mode fit to the number-  
482 weighted distribution is similar to the observed volume-weighted distribution (Figure 8). Thus,  
483 the use of a single-mode to represent biomass burning size distributions thus appears acceptable,  
484 so long as the appropriate parameters are used. In this context, the widths of the distribution used  
485 by the various global models appear somewhat too small. However, we note that the microphysics  
486 occurring in the fresh smoke sampled here, which will govern the size distributions, may differ  
487 from that in atmospheric plumes.

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488 The average BC-specific mass-weighted size distribution mode is at 148 nm (Figure 8). A  
489 bimodal fit yields values for the mass median diameter ( $d_{p,M}$ ) and  $\sigma_g$  of 137.2 nm and 1.62,  
490 respectively, for the smaller mode and 197.1 nm and 1.24 for the larger mode, with most of the  
491 mass contained in the smaller mode. May et al. (2014) report  $d_{p,M}$  from laboratory biomass  
492 combustion ranging from 140-190 nm, averaging 170 nm. Their average is somewhat larger than  
493 ours, likely reflecting differences in the exact fuels sampled. The mode diameter for the BC-  
494 specific distribution is especially smaller than observed for biomass burning particles from some  
495 ambient observations, which tend to give values closer to 200 nm (Schwarz et al., 2008;Kondo et  
496 al., 2011;Sahu et al., 2012;May et al., 2014;Cappa et al., 2019b). This difference between lab and  
497 field observations was also noted by May et al. (2014). We speculate that the influence of  
498 coagulation may be suppressed in our experiments relative to what occurs in the atmosphere due  
499 to slower overall dilution, leading to smaller BC size distributions. To the extent this is the reason  
500 for the difference, the total particle distributions would also be biased towards too small particles,  
501 compared to the atmosphere. However, there is no relationship between  $d_{p,N}$  and the total particle  
502 number concentration for our experiments. Formation of secondary aerosol in the near-field of a  
503 sampled ambient plume could also contribute to this difference.

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504 There is substantial variability between individual burns within a given particle Class in terms  
505 of the shape of the size distributions (Figure 7). This variability is most evident for Class 1, 2 and

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509 5, but present for all Classes. Nonetheless, the number-weighted mean diameter ( $d_{p,N,mean}$ ) appears  
510 to decrease somewhat with MCE (**Figure 9**;  $r^2 = 0.38$ ). However, the relationship is largely driven  
511 by the Class 6 particles, which generally have lower MCE values, having larger  $d_{p,N,mean}$  values. A  
512 lack of any particularly clear relationship is consistent with Hosseini et al. (2010), who observed  
513 the  $d_{p,N,mean}$  to exhibit a complex relationship with combustion conditions. The  $d_{p,N,mean}$  varies non-  
514 monotonically with [OA]/[BC], with particle size first decreasing slightly as [OA]/[BC] increases  
515 (from Class 1 to Class 3) and then increasing with further increases in [OA]/[BC] (from Class 4 to  
516 Class 6) (**Figure 9**). This is despite the notable burn-to-burn variability. It is important to note that  
517 the mobility-based size is particle shape-dependent; very BC-rich particles are more likely to have  
518 non-spherical shapes and thus have larger mobility diameters. This could explain the minimum in  
519  $d_{p,N}$  around Class 3 particles, for which [OA]/[BC] = 10.

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520 Some of the variability within a class appears related to the presence of different fuel types  
521 within a class. Number-weighted and BC-specific mass-weighted size distributions by fuel type  
522 are shown in **Figure 10**. For the number-weighted distributions, leaf litter and rotten logs exhibit  
523 the greatest variability between different burns, although we note that multiple burns were not  
524 performed for all fuels. The shapes of the leaf litter, peat and “other” fuel types differ most notably  
525 from the other fuel types, with the presence of more than one mode more apparent. (The “other”  
526 category here includes non-traditional biofuels, specifically building materials and excelsior.) For  
527 the BC-specific size distributions, the litter, canopy, and duff exhibited the greatest intra-fuel  
528 variability. For most fuels, the BC-specific distribution peaks around 150 nm, as noted above.  
529 However, for a subset of burns (eight of them) the BC-specific distribution peaks around 100 nm  
530 (**Figure 10**). These small BC-mode distributions occur for the OA-rich particle classes 4, 5 and 6  
531 (**Figure 7**), although there is no clear pattern to their occurrence.

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#### 532 4 Conclusions and Implications

533 Measurements of primary particles produced from combustion of a variety of biomass fuel  
534 types indicate the optical, physical, and chemical properties of the emitted particles exhibit wide  
535 variability. We show that variability in many optical properties (e.g. single scatter albedo,  
536 wavelength dependence of absorption, mass absorptivity of black and brown carbon) is directly  
537 linked to the [OA]/[BC] ratio of the emitted particles; the relationships with [OA]/[BC] are much  
538 stronger than with the commonly used modified combustion efficiency, and mathematical

544 relationships between the various properties are determined. However, the absorption  
545 enhancement due to coating of BC (the so-called “lensing” effect) is shown to be minor and  
546 essentially independent of the amount of coating up to large coating-to-BC mass ratios. The brown  
547 carbon mass absorptivity correlates with the nitrated organic fraction of OA, suggesting that  
548 nitrated organic species contribute to BrC absorption. Many bulk chemical properties (i.e. O:C,  
549 H:C, and the relative concentrations of key marker ions such as  $f_{60}$ ) exhibit limited dependence on  
550 the burn conditions and the [OA]/[BC] ratio. However, both the OA volatility and nitrated organic  
551 fraction of OA decrease with [OA]/[BC]. The fraction of OA that is internally mixed with BC was  
552 shown to decrease strongly with the [OA]/[BC] ratio, from nearly all OA being internally mixed  
553 with BC when the particles are overall BC-rich to only a few percent of OA being mixed with BC  
554 when OA dominates. Yet, the relative amount of OA coating the BC increases with [OA]/[BC];  
555 that is, when more of the OA is externally mixed from BC those particles that do contain BC  
556 nonetheless have thicker OA coatings. The observed total particle size distributions are reasonably  
557 well described by a single log-normal mode, but are better fit using a bimodal distribution. The  
558 BC-specific size distributions are similarly best fit using a bimodal distribution, although a single  
559 mode provides a reasonable representation. The dependence of the geometric median mobility  
560 diameter on the burn conditions or particle state (i.e. the [OA]/[BC]) is complicated by the mobility  
561 diameter being sensitive to variations in particle shape, which depend on the [OA]/[BC] ratio.  
562 Overall, these results expand on previous observations of primary biomass burning particle  
563 properties, considering a wider range of [OA]/[BC] and associated properties. Further, they  
564 provide a foundation for understanding the post-emission evolution of biomass burning smoke due  
565 to photochemical oxidation as discussed in Lim et al. (2019).

## 566 **5 Data Availability**

567 All data are available from the NOAA FIREX-AQ data repository  
568 (<https://esrl.noaa.gov/csd/projects/firex/firelab/>). This includes a summary of the fuel types used  
569 for each burn and the measurement time-series for each burn. The primary particle averages used  
570 in this work are additionally collected in the UC DASH data repository (Cappa et al., 2019a).

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572 **6 Author Contributions**

573 CDC and JHK designed the experiments. CDC, CYL, and DHH carried out the measurements  
574 and data processing. CDC, CDM, and CYL analyzed data. CDC and CDM wrote the manuscript,  
575 with contributions from all co-authors.

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918 **9 Tables**

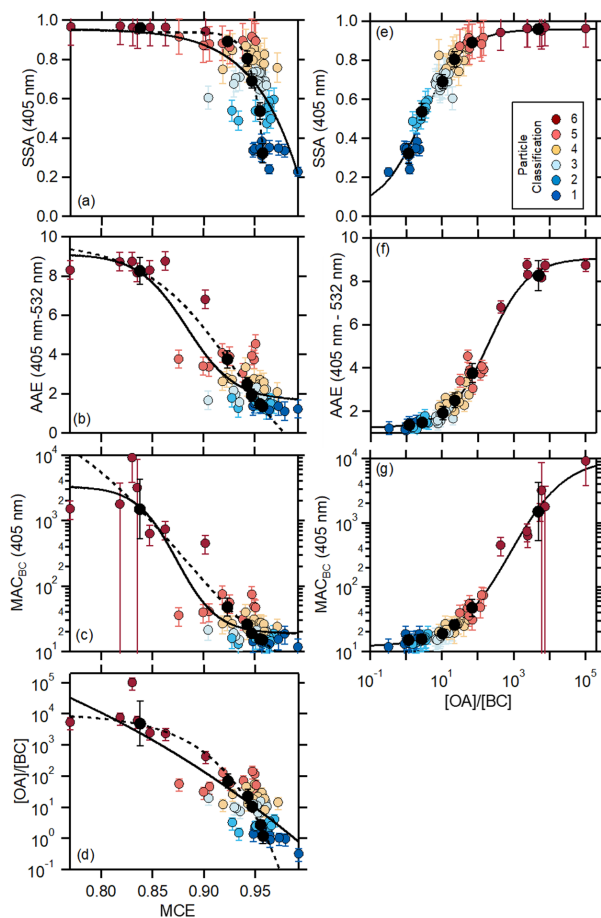
919 **Table 1.** Fuels by particle Class.

<b>Class</b>	<b>Fuel</b>	<b>SSA range</b>	<b>[OA]/[BC] range</b>
Class 1	Chaparral, canopy, litter (pine), building materials, excelsior	0.23-0.43	0.3-2.4
Class 2	Manzanita, Sage, litter (fir)	0.43-0.60	1.5-4.1
Type 3	Pine, fir, litter, canopy, juniper	0.60-0.74	6.6-20
Class 4	Pine, fir, canopy, rotten log, ceonothos	0.74-0.87	8.3-55
Class 5	Canopy (pine), rice, bear grass, duff	0.87-0.93	31-143
Class 6	Rotten log, duff, peat, dung	0.93-1.00	431-10 <sup>5</sup>

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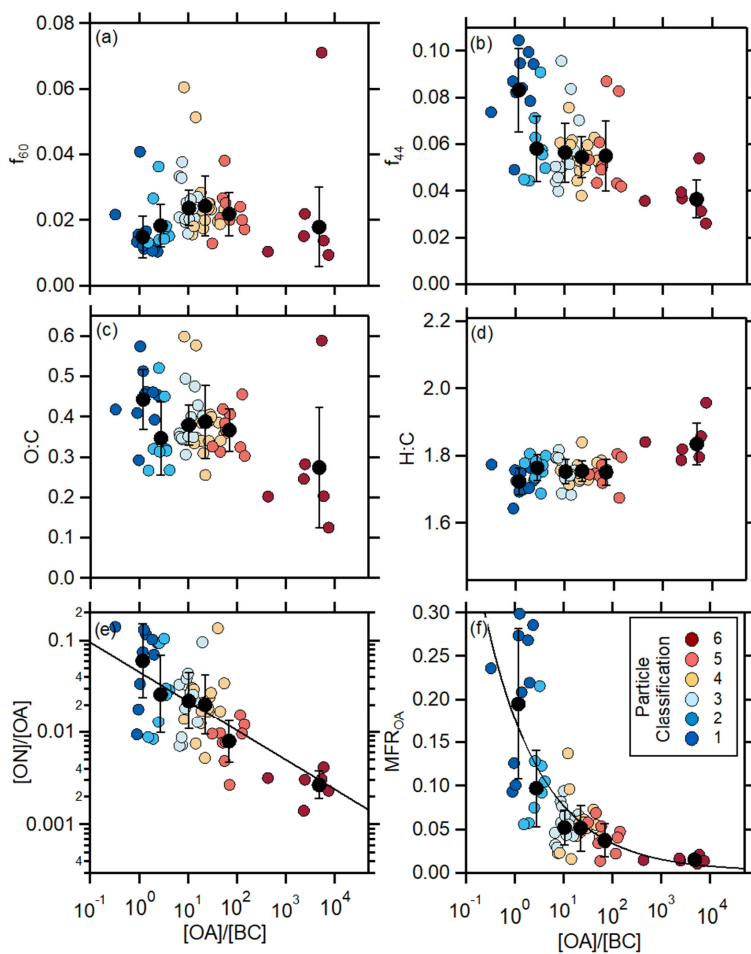
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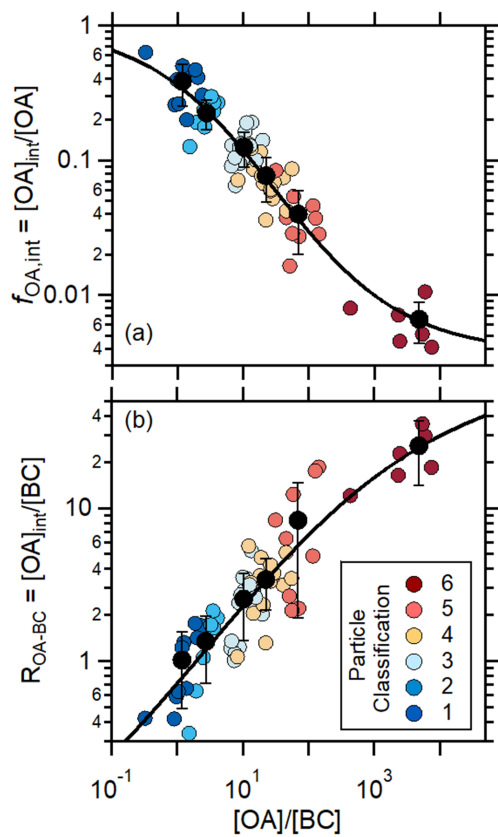
924 **Figure 1.** (left panels) Relationship between (a) the  $SSA_{405nm}$ , (b) the  $AAE_{405-532}$ , (c) the  $MAC_{BC}$ ,  
 925 and (d) the  $[OA]/[BC]$  mass ratio and the modified combustion efficiency, MCE. Results for  
 926 individual burns are shown as points colored by the particle Class, and Class average values are  
 927 shown as black circles. Uncertainties on the Class averages are  $1\sigma$  based on measurement  
 928 variability and uncertainties on for the individual burns are from error propagation of measurement  
 929 uncertainties. The solid black lines are fits to the individual burns (colored points) while the dashed  
 930 black lines are fits to the Class averages (Table S2). (right panels) Relationship between (e) the  
 931  $SSA_{405nm}$ , (f) the  $AAE_{405-532}$ , and (g) the  $MAC_{BC}$  on the  $[OA]/[BC]$  mass ratio. The solid black  
 932 lines here are sigmoidal fits to the individual burns. Fits to the Class averages are similar.



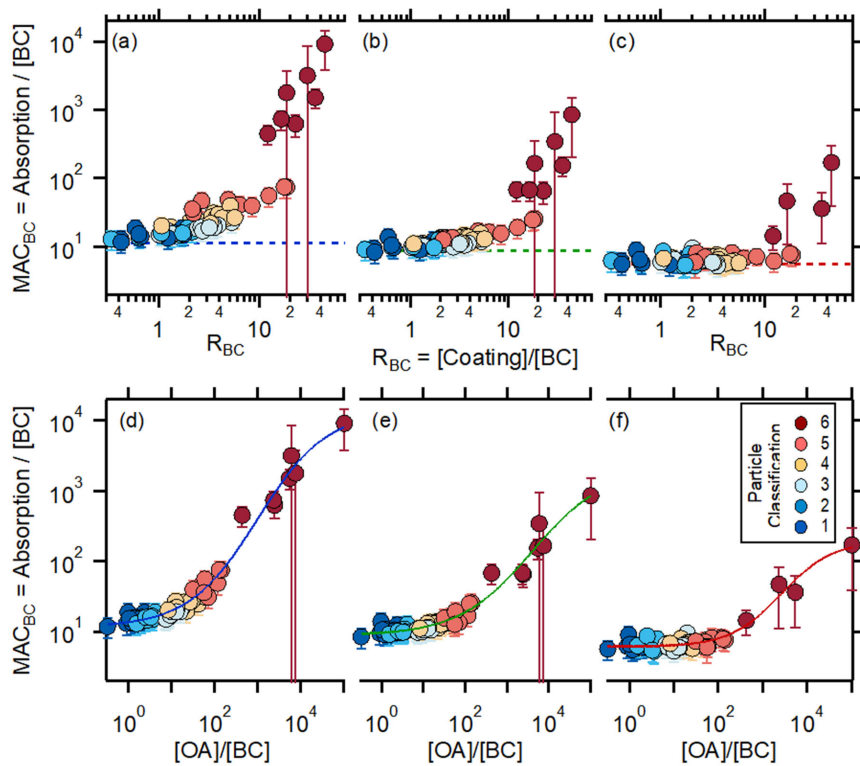
934

935 **Figure 2.** Dependence of (a)  $f_{60}$ , (b)  $f_{44}$ , (c) O:C, (d) H:C, (e) the nitrated organic fraction of OA,  
 936  $f_{\text{ON-OA}}$ , and (f) the OA volatility, characterized as the mass fraction remaining after heating. Results  
 937 for individual burns are shown as points colored by the particle Class, and Class average values  
 938 are shown as black circles. Uncertainties on the Class averages are  $1\sigma$  based on measurement  
 939 variability. For  $f_{\text{ON-OA}}$  and  $MFR_{\text{OA}}$ , fits to the observations are shown (see text).

940



941  
 942 **Figure 3.** Relationship between (a) the fraction of OA that is internally mixed with BC,  $f_{\text{OA,int}}$  and  
 943 (b) the OA-to-BC mass ratio for only the internally mixed OA, and the total  $[\text{OA}]/[\text{BC}]$  mass ratio.  
 944 Results for individual burns are shown as points colored by the particle Class, and Class average  
 945 values are shown as black circles. Uncertainties on the Class averages are  $1\sigma$  based on  
 946 measurement variability. Black lines are sigmoidal fits to the data, in log-log space.  
 947

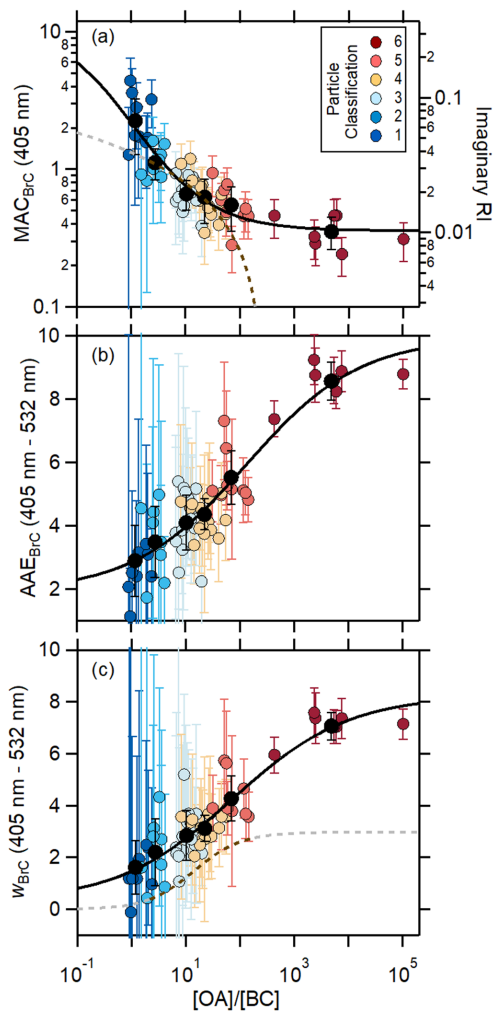


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949 **Figure 4.** (Top Panels) The relationship between the wavelength-specific  $MAC_{BC}$  and the coating-  
 950 to-BC mass ratio for (a) 405 nm, (b) 532 nm and (c) 781 nm. The horizontal dashed lines show the  
 951 derived  $MAC_{BC,pure}$  values. (Bottom Panels) The relationship between the wavelength-dependent  
 952  $MAC_{BC}$  and the total  $[OA]/[BC]$  mass ratio for (d) 405 nm, (e) 532 nm and (f) 781 nm. The lines  
 953 are sigmoidal fits. Uncertainties for the individual burns are determined from error propagation.  
 954 Graphs of the wavelength-specific  $MAC_{BC}$  versus  $[OA]/[BC]$  with each shown using independent  
 955 y-axis scales are provided for comparison in [Figure S1](#).

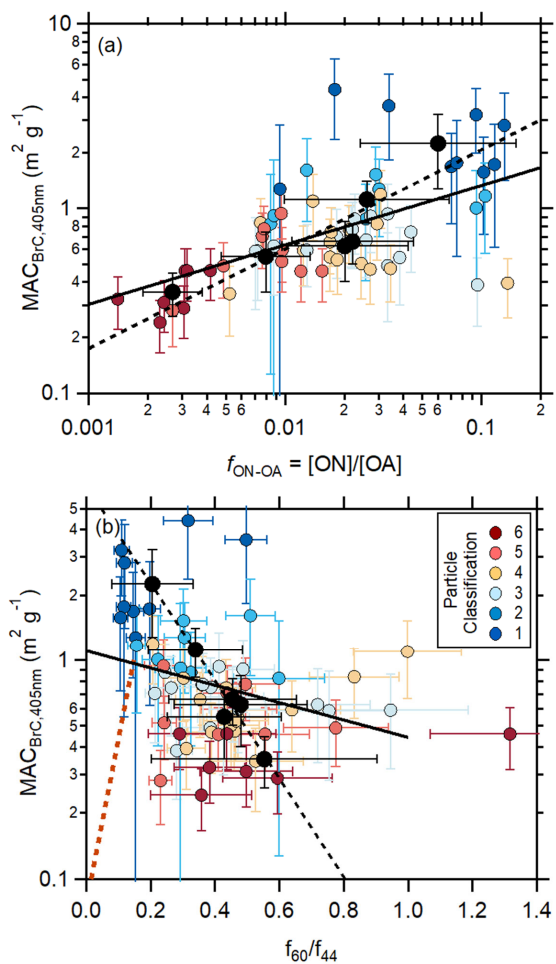
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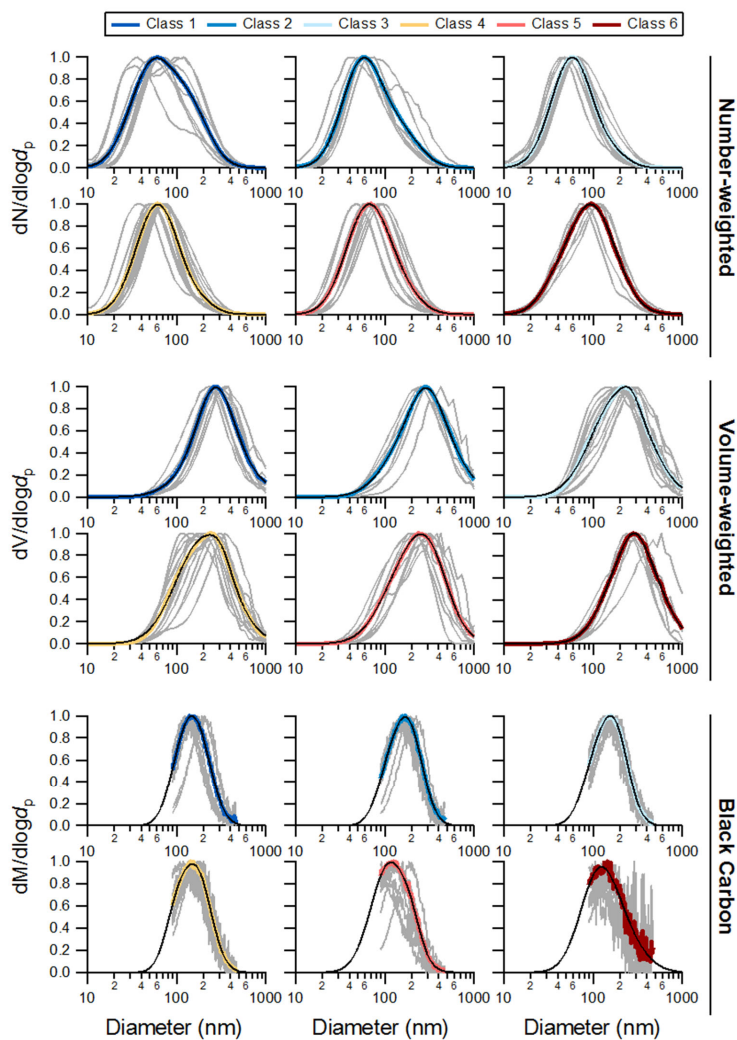


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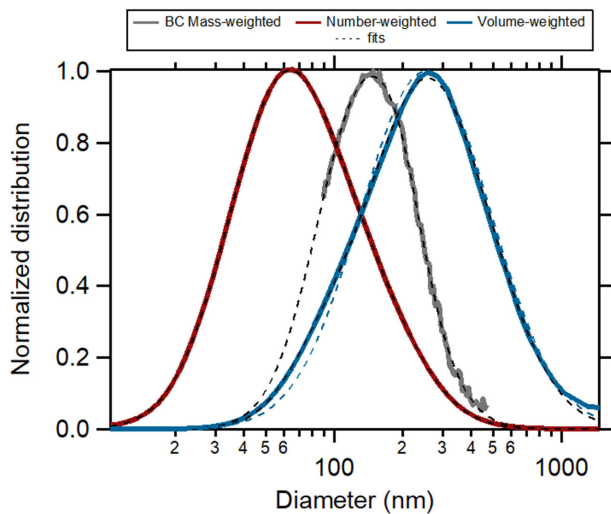
959 **Figure 5.** Relationship between (a)  $MAC_{BrC,405nm}$ , (b)  $AAE_{BrC,405-532}$ , and (c)  $w_{BrC,405-532}$  and  
 960  $[OA]/[BC]$  mass ratio. The solid lines are sigmoidal fits to the observations, against  
 961  $\log([OA]/[BC])$ . The dashed lines are based on the parameterization of Saleh et al. (2014), with  
 962 the brown color indicating the measuring range in that study and the gray color extrapolated.  
 963 Results for individual burns are shown as points colored by the particle Class, and Class average  
 964 values are shown as black circles. Uncertainties on the Class averages are  $1\sigma$  based on  
 965 measurement variability. Uncertainties for the individual burns are determined from error  
 966 propagation.



969 **Figure 6.** Relationship between the  $MAC_{BRC,405nm}$  and (a) the nitrated organic fraction of total  
 970 organic aerosol,  $f_{ON-OA}$ , and (b) the  $f_{60}/f_{44}$  ion ratio for organic aerosol. Results for individual burns  
 971 are shown as points colored by the particle Class, and Class average values are shown as black  
 972 circles. Uncertainties on the Class averages are  $1\sigma$  based on measurement variability. Uncertainties  
 973 for the individual burns are determined from error propagation. Solid black lines are fits to all  
 974 burns and dashed black lines are fits to the Class averages. The dashed brown line in panel (b) is  
 975 the relationship reported by Lack et al. (2013) for ambient particles in a biomass burning plume.



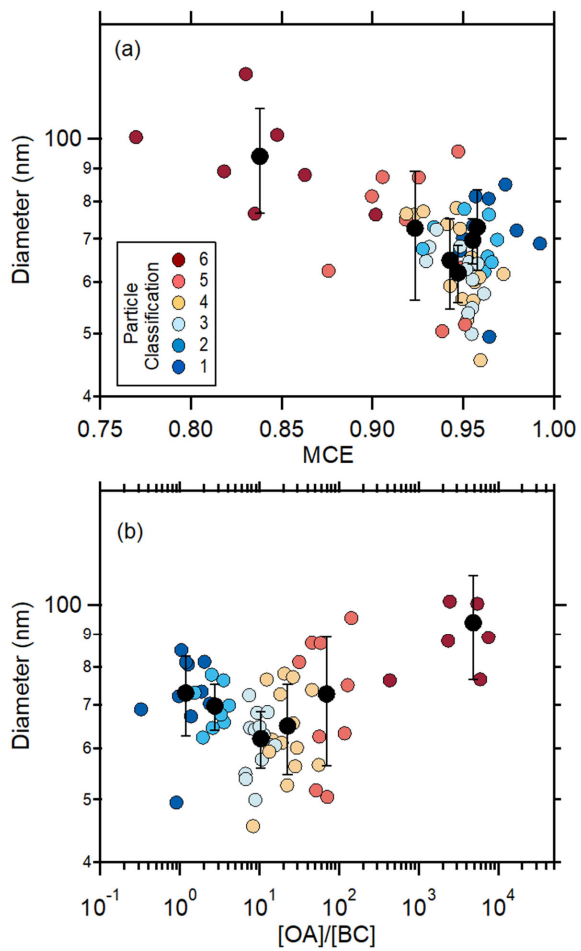
977  
 978 **Figure 7.** Class-specific total particle number-weighted (top) and volume-weighted (middle)  
 979 mobility size distributions, and the BC-only mass-weighted (bottom) size distribution. Individual  
 980 burns are shown in gray and class averages are shown as colors. Bimodal log-normal fits are thin  
 981 black lines. Note that the number-weighted and volume-weighted distributions are graphed versus  
 982 mobility diameter and the BC mass-weighted distribution against the BC volume equivalent  
 983 diameter.



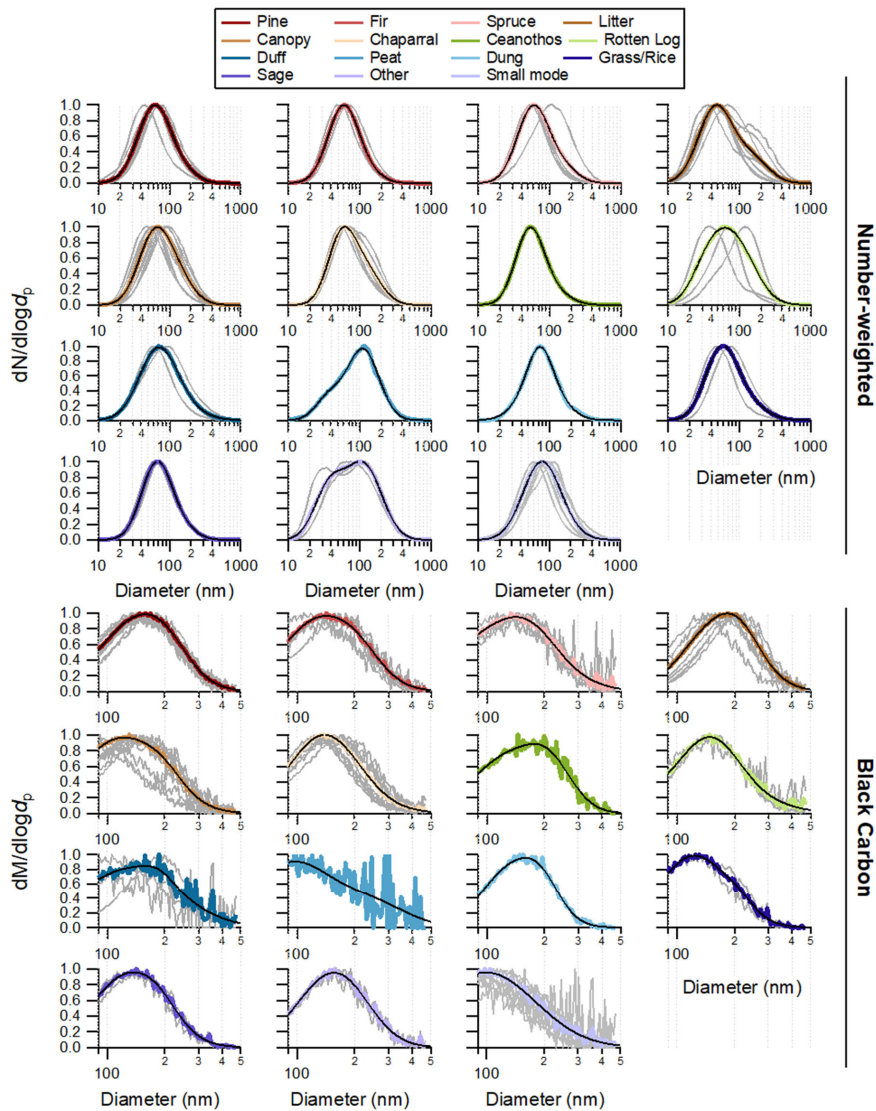
984  
 985 **Figure 8.** Average total particle number-weighted (red) and volume-weighted (blue) size  
 986 distributions and the BC-specific mass-weighted size distributions. Black dashed lines are bimodal  
 987 log-normal fits. The dashed blue line is the total particle volume-weighted distribution calculated  
 988 from a single-mode fit to the number-weighted distribution. Note that the number-weighted and  
 989 volume-weighted distributions are graphed versus mobility diameter and the BC mass-weighted  
 990 distribution against the BC volume equivalent diameter.

991





992  
 993 **Figure 9.** Relationship between number-weighted particle median diameter and (a) the MCE and  
 994 (b) the [OA]/[BC] ratio. Colored circles are for individual burns and black circles for particle class  
 995 averages.



996  
 997 **Figure 10.** Normalized total particle number-weighted (top) and the BC-only mass-weighted  
 998 (bottom) size distributions shown by fuel type (see legend). Individual burns are gray and averages  
 999 for a fuel type colors. For some fuels there is only one size distribution. Bimodal log-normal fits  
 1000 are the black lines. The “other” category includes non-traditional biofuels, specifically building  
 1001 materials and excelsior.

1        **Supplementary Material for:**  
2        **Biomass-burning derived particles from a wide variety of**  
3        **fuels: Part 1: Properties of primary particles**

4        Crystal D. McClure<sup>1</sup>, Christopher Y. Lim<sup>2,%</sup>, David H. Hagan<sup>2</sup>, Jesse H. Kroll<sup>2</sup>, Christopher D.  
5        Cappa<sup>1,3,\*</sup>

6        <sup>1</sup> Department of Civil and Environmental Engineering, University of California, Davis, CA 95616

7        <sup>2</sup> Department of Civil and Environmental Engineering, Massachusetts Institute of Technology,  
8        Cambridge, MA, USA

9        <sup>3</sup> Atmospheric Sciences Graduate Group, University of California, Davis, CA, USA 95616

10        <sup>%</sup> Now at Department of Chemistry, University of Toronto, Ontario, Canada

11        \* To whom correspondence should be addressed: cdcappa@ucdavis.edu

12

13        **10.1 Description of Instrumentation**

14        **Particle Optical Property Measurements**

15        Particle optical properties for PM<sub>1</sub> were measured at 405 nm and 532 nm using the UC Davis  
16        Cavity Ringdown-Photoacoustic Spectrometer (CRD-PAS). In the UC Davis CRD-PAS, Light  
17        absorption coefficients ( $b_{\text{abs}}$ ; units = Mm<sup>-1</sup>) for dry particles are determined using photoacoustic  
18        spectroscopy (Lack et al., 2012b). Light extinction coefficients ( $b_{\text{ext}}$ ; units = Mm<sup>-1</sup>) for dry (<20%  
19        relative humidity) particles are measured at 405 nm and 532 nm via cavity ringdown spectroscopy  
20        (Langridge et al., 2011). Humidified light extinction measurements (RH ~85%) are also measured  
21        at 532 nm by cavity ringdown spectroscopy. The absorption measurements from the PAS were  
22        calibrated relative to the extinction measurement from the CRD using gas-phase O<sub>3</sub> and NO<sub>2</sub> with  
23        an estimated accuracy of 5% at 532 nm and 8% at 405 nm. Light absorption and scattering  
24        coefficients were also measured at 781 nm using a commercial PASS-3 photoacoustic  
25        spectrometer (DMT, Inc.). In the PASS-3, light absorption coefficients are measured by  
26        photoacoustic spectroscopy. Light scattering coefficients ( $b_{\text{sca}}$ ; units = Mm<sup>-1</sup>) are determined for  
27        dry particles with the PASS-3 using reciprocal nephelometry. The absorption measured by the

28 PASS-3 was calibrated relative to the UC Davis PAS using polydisperse fullerene soot and  
29 assuming that the absorption Ångstrom exponent was 1.4 (Metcalf et al., 2013). The estimated  
30 uncertainty in  $b_{\text{abs}}$  at 781 nm is 10%.

## 31 **Particle Composition Measurements**

### 32 **Refractory black carbon measurement**

33 Refractory black carbon (rBC) concentrations and BC-specific particle size distributions were  
34 measured using a single particle soot photometer (SP2). The SP2 measures the concentration of  
35 rBC within individual rBC-containing particles. Sampled particles pass through a 1064 nm  
36 intracavity laser. Absorption of this light by rBC leads to rapid heating of the particles. If heating  
37 outweighs conductive cooling the particles will reach a sufficiently high temperature (i.e. their  
38 boiling point) that they will incandesce. The intensity of this incandescent light is proportional to  
39 the rBC mass of that particle (usually on the order of 0.1 – 10 fg per particle). Size distributions of  
40 only the rBC (exclusive of any other internally mixed material) are generated by converting the  
41 per particle mass to a volume equivalent diameter ( $d_{\text{p,VED}}$  here, assuming  $\rho_{\text{rBC}} = 1.8 \text{ g cm}^{-3}$ ) and  
42 binning the particles by size. The SP2 was calibrated using size-selected fullerene particles (Lot  
43 L20W054, Alfa Aesar, Ward Hill, MA, USA).

44 When the number concentration of rBC-containing or non-rBC-containing particles is large,  
45 the SP2 may suffer from negative biases in the concentration measurement. This can happen when  
46 the SP2 detectors are triggered by one particle and a second passes through the viewing volume  
47 during the detection window (typically  $\sim 50 \mu\text{s}$ ). Such particle coincidence effects can be  
48 minimized by decreasing the sample flowrate into the SP2 to decrease the likelihood that two  
49 particles are simultaneously in the viewing volume. Here, the SP2 sample flowrate was varied  
50 from 5 sccm to 120 sccm in a step-wise manner over the course of an experiment to deal with the  
51 very large dynamic range of concentrations in the mini chamber. The flow rate was increased to  
52 maintain an approximately constant particle count rate in the instrument while minimizing the  
53 influence of particle coincidence. Inspection of individual particle detection events indicates that  
54 particle coincidence was generally avoided.

55 The SP2 data were processed using the SP2 Toolkit from the Paul Scherer Institute (PSI),  
56 developed by Martin Gysel. The SP2 size-dependent counting efficiency was determined by

57 simultaneously measuring the concentration of the calibration particles with a mixing condensation  
58 particle counter (BMI Model 2002). The particle counting efficiency was found to be unity for  
59 particles with  $d_{p,VED} > 100$  nm. The SP2 used in this study measured particles over the size range  
60  $90 \leq d_{p,VED} \leq 822$ . Below the lower size limit, the detection efficiency falls off rapidly due, in part,  
61 to the large surface area-to-volume (SA-to-V) ratio of these particles. When the SA-to-V ratio is  
62 sufficiently large conductive cooling competes effectively with the radiative heating from the laser  
63 and the particles do not emit enough incandescent light at short enough wavelengths to trigger  
64 detection. Above the upper size limit, the incandescence level is sufficient to saturate the detector,  
65 leading to an underestimate in particle mass. All SP2 mass concentration measurements were  
66 corrected for the missing mass contained in particles below the lower and upper size limit, using a  
67 multi-mode fitting approach.

68 The observed campaign average distribution mode peak is around 150 nm. The observed  
69 distributions (1 min averages) were fit to a four-mode log-normal distribution to estimate and  
70 correct for the rBC outside of the SP2 detection window, i.e. for rBC “missing mass”. The average  
71 ratio between the observed rBC concentration and the total estimated from fitting was  $0.83 \pm 0.06$   
72 ( $1\sigma$ ). There was some experiment-to-experiment and time-dependent variability in the missing  
73 mass fraction that is accounted for by fitting the observations at 1 min time resolution. This  
74 approach follows that of Zhang et al. (2016). While a single mode fit provides a reasonably  
75 representation of the overall campaign average distribution, inspection of the individual  
76 distributions across the experiments indicates that a multi-mode fitting approach provides a  
77 substantially more robust description of the observed size distribution, especially as particle aging  
78 proceeds.

### 79 Composition and concentration of NR-PM

80 The concentration of non-refractory particulate matter (NR-PM) species in  $PM_{10}$  were measured  
81 using a high-resolution time-of-flight aerosol mass spectrometer (HR-ToF-AMS, henceforth HR-  
82 AMS) (Canagaratna et al., 2007) during both the Fresno and Fontana studies, as discussed in detail  
83 by (Lim et al., 2019). The NR-PM components are functionally defined as those materials that  
84 evaporate rapidly after impaction onto a heated surface *in vacuo* at  $\sim 600$  °C. The NR-PM  
85 components characterized include particulate sulfate, nitrate, ammonium, chloride and organic  
86 matter. The data were processed using the PIKA toolkit in IGOR (Wavemetrics, Inc.). The

87 collection efficiency (CE) of the HR-AMS was determined by comparison with size distributions  
88 measured using the scanning electrical mobility spectrometer (SEMS). The collection efficiency  
89 differed between primary and secondary and secondary particles and was found to co-vary with  
90 the volatility of the organic aerosol. The variation in the CE was empirically accounted for, as  
91 discussed in (Lim et al., 2019). The estimated uncertainty for the HR-AMS measurements is  $\pm 30\%$ ,  
92 although the precision is much better than this.

### 93 Particulate nitrated organics characterization

94 The concentration of nitrated organic functional groups ( $ON_f$ ) is determined from the HR-AMS  
95 measurements. Kiendler-Scharr et al. (2016) showed that the fraction of total nitrate measured by  
96 the HR-AMS that derives from organic nitrate functional groups ( $f_{ON-N}$ ) relates to the measured  
97  $[NO_2^+]/[NO^+]$  ratio ( $R_{meas}$ ):

$$98 \quad f_{ON-N} = \frac{(1+R_{ON})(R_{meas}-R_{calib})}{(1+R_{meas})(R_{ON}-R_{calib})} \quad (S1)$$

99 where  $R_{ON} = 0.1$  and  $R_{calib}$  is an instrument-specific factor determined from calibration with  
100  $NH_4NO_3$  and here equaling 0.45. The Kiendler-Scharr et al. (2016) approach focused on the  
101 behavior of organic nitrates. We assume here that nitro-organics behave similarly and thus that  
102  $ON_f$  encompasses contributions from both nitrate and nitro functional groups. Equation S1 is  
103 thought reliable when the  $f_{ON-N} > 0.15$  (Bruns et al., 2010). The average  $f_{ON-N}$  for the FIREX  
104 measurements is  $0.74 \pm 0.24$  ( $1\sigma$ ). The concentration of particulate ON functional groups is then  
105  $[ON_f] = f_{ON-N}[NO_3^-]$ . Note that this includes only the mass of the functional group; the total mass  
106 concentration of the ON species (including the carbon backbone) can be estimated by multiplying  
107 the  $ON_f$  concentration by the ratio between an assumed MW for the ON species and that for the  
108 nitrate functional group. We assume that ON species have a MW = 200 amu, and thus  $[ON] =$   
109  $3.22[ON_f]$ . If the ON signal is dominated by nitro functional groups, rather than nitrate, then the  
110 estimated  $[ON]$  is a lower limit.

### 111 Composition and concentration of BC-containing particles

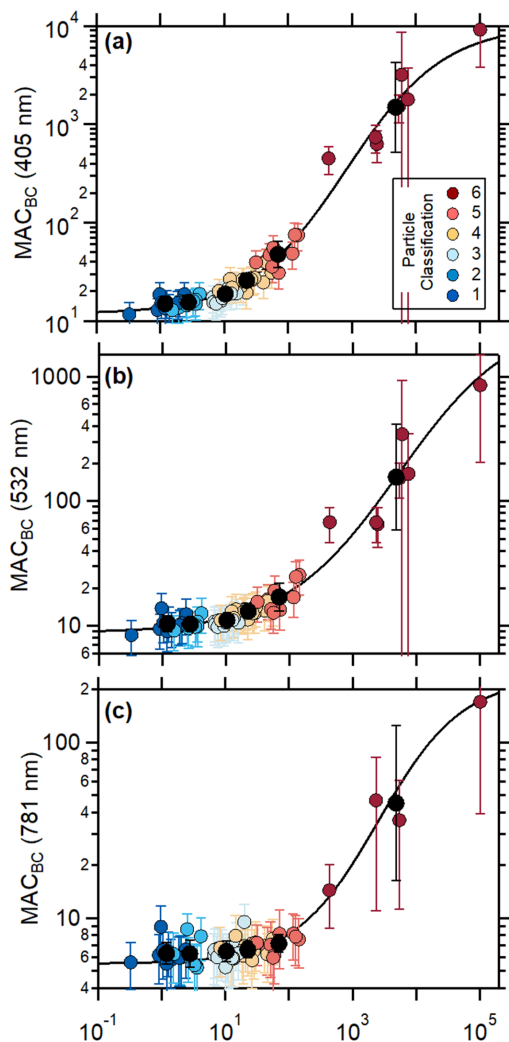
112 The concentrations and composition of only BC-containing particles were determined using a  
113 soot particle aerosol mass spectrometer (SP-AMS) (Onasch et al., 2012). In the SP-AMS, a focused  
114 particle beam is intersected with an intra-cavity Nd:YAG laser operating at 1064 nm. Particles

115 containing BC are rapidly heated by the laser, leading to evaporation of both the NR-PM materials  
116 and the refractory BC. In these studies, the standard HR-AMS tungsten vaporizer was removed so  
117 that particles that do not contain BC are not vaporized and are therefore not detected. Thus, the  
118 SP-AMS is specific to BC-containing particles, as operated here. In addition to BC, the SP-AMS  
119 measures the internally mixed particulate inorganic (sulfate, nitrate, ammonium, and chloride) and  
120 organic mass loading. The NR-PM species that are associated with BC will be distinguished from  
121 the bulk average NR-PM species (from the HR-AMS) using the subscript BC (i.e. NR-PM<sub>rBC</sub>).  
122 The SP-AMS particle detection efficiency is determined in large part by the extent of overlap  
123 between the particle and laser beam. Particles were sampled through a PM<sub>1</sub> aerodynamic lens, with  
124 particles measured down to ~40 nm vacuum aerodynamic diameter. The SP-AMS detection  
125 efficiency was determined by referencing the rBC concentration measured by the SP-AMS to that  
126 measured by the SP2, as in (Collier et al., 2018). The SP-AMS/SP2 ratio depended on the ratio  
127 between the NR-PM<sub>rBC</sub> and BC, with the NR-PM/rBC ratio decreasing as the SP-AMS/SP2 ratio  
128 increases. However, throughout this work we use only the [NR-PM<sub>rBC</sub>]/[rBC] or [OA<sub>rBC</sub>]/[rBC]  
129 ratios, which are not dependent on the absolute instrument calibration, but only the relative  
130 detection efficiency of these species. The coating-to-core mass ratio for both campaigns is  
131 calculated directly from the SP-AMS measurements as  $R_{BC} = [\text{NR-PM}]_{BC}/[\text{BC}]$ .

### 132 **Gas Composition Measurements**

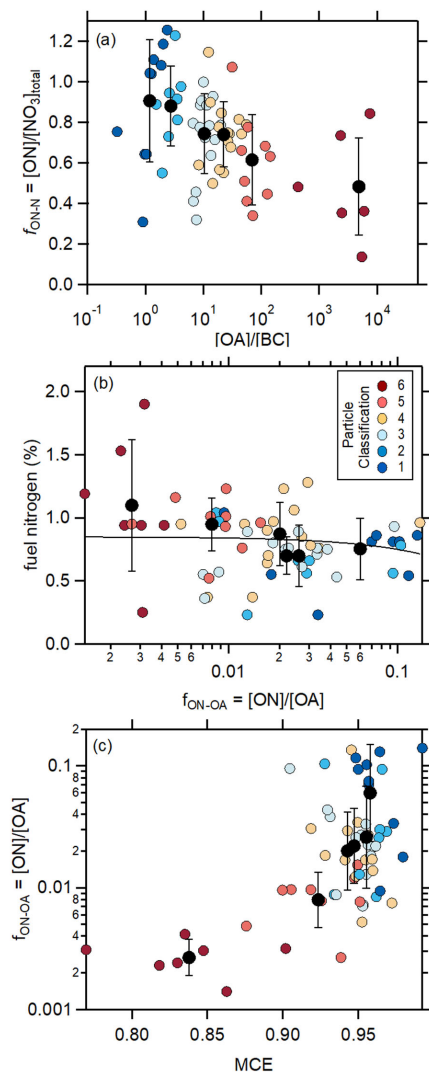
133 The concentrations of select gas-phase non-methane organic gases (NMOG) and some  
134 inorganic species (e.g. HONO) were measured using H<sub>3</sub>O<sup>+</sup> and I<sup>-</sup> chemical ionization mass  
135 spectrometers (CIMS), that included high-resolution time-of-flight mass spectrometers. Only the  
136 measurements from the PTR-TOF-MS, operated by the National Oceanic and Atmospheric  
137 Administration, are used here. The PTR-TOF-MS measurements are described in detail in (Koss  
138 et al., 2018) and (Sekimoto et al., 2018). In addition to the NMOG measurements, other inorganic  
139 gases (O<sub>3</sub>, CO, CO<sub>2</sub>, SO<sub>2</sub>) were measured using commercial instrumentation.

140

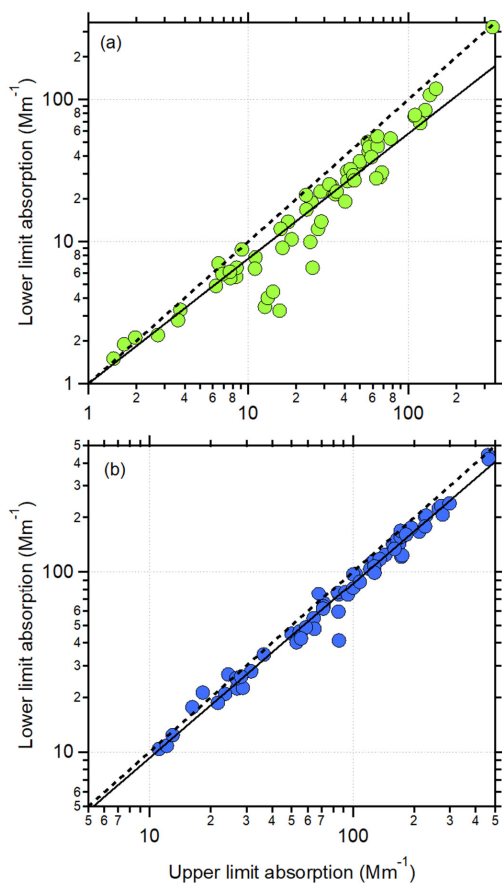


142  
 143 **Figure S1.** Relationship between the observed ambient particle  $MAC_{BC}$  and the total particle  
 144  $[OA]/[rBC]$  at (a) 405 nm, (b) 532 nm, and (c) 781 nm. Individual points colored by Class (see  
 145 text) and class averages as black circles.





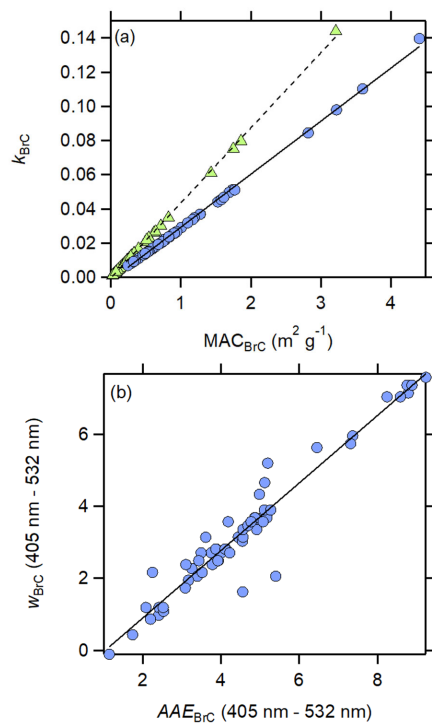
146  
 147 **Figure S2.** (a) Relationship between fuel nitrogen and the fraction of OA that is organic nitrate,  
 148  $f_{\text{ON-N}}$ . There is no correlation between the two. (b) Relationship between  $f_{\text{ON-OA}}$  and the modified  
 149 combustion efficiency, MCE. Results for individual burns are shown as points colored by the  
 150 particle Class, and Class average values are shown as black circles. Uncertainties on the Class  
 151 averages are  $1\sigma$  based on measurement variability.



152

153 **Figure S3.** The derived lower limit brown carbon absorption versus the upper limit brown carbon  
 154 absorption at (a) 532 nm and (b) 405 nm. The lower limit estimate for BrC absorption accounts for  
 155 the potential influence of coating-induced enhancements. The dashed line is the one-to-one line  
 156 and the solid line is a linear fit with slopes equaling 0.88 at 532 nm and 0.97 at 405 nm.

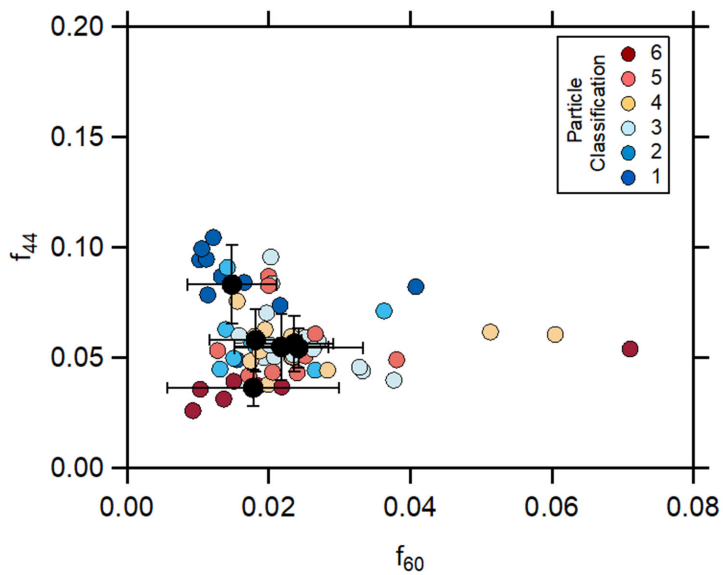
157



158

159 **Figure S4.** (a) Relationship between the imaginary refractive index for BrC,  $k_{\text{BrC}}$ , at 405 nm (blue  
 160 circles) and the observed  $\text{MAC}_{\text{BrC}}$  at 405 nm or at 532 nm (green triangles). Lines are linear fits to  
 161 the observations. (b) Relationship between the wavelength dependence of  $k_{\text{BrC}}$ ,  $w_{\text{BrC}}$ , determined  
 162 for the 405 nm – 532 nm pair, and the  $\text{AAE}_{\text{BrC}}$  for the same wavelengths.  
 163

164



165

166 **Figure S5.** The relationship between the fractional abundance of the  $m/z = 44$  ( $f_{44}$ ) and  $m/z = 60$   
167 ( $f_{60}$ ) ions from organic aerosol. Points are colored by particle class for individual burns, and the  
168 class averages shown in black.

169

170 **Table S1.** Fuels combusted. Further details regarding fuel properties are available at the NOAA  
171 data repository, in particular in the summary spreadsheet  
172 ([https://esrl.noaa.gov/csd/groups/csd7/measurements/2016firex/FireLab/DataDownload/FIREX\\_](https://esrl.noaa.gov/csd/groups/csd7/measurements/2016firex/FireLab/DataDownload/FIREX_BurnListComplete_V5.xlsx)  
173 [BurnListComplete\\_V5.xlsx](https://esrl.noaa.gov/csd/groups/csd7/measurements/2016firex/FireLab/DataDownload/FIREX_BurnListComplete_V5.xlsx); access date 04 February 2019)

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<b>Fuel Type</b>
Bear Grass
Building Material - Untreated Wood
Ceanothos
Chapparral (canopy)
Chamise
Manzanita
Douglas Fir (litter, canopy, mixture, rotten log)
Dung
Engelmann spruce (canopy, mixture, duff)
Excelsior (wood wool)
Jeffrey Pine (duff)
Juniper (canopy)
Loblolly pine (litter)
Lodgepole (canopy, litter, mixture)
Peat
Ponderosa pine (litter, canopy, mixture, rotten log)
Rice Straw, Arkansas
Sage
Sage Brush
Subalpine fir (canopy, litter, mix, duff)

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174

175

176 **Table S2.** Fit coefficients for the various fits performed, organized by fit type (e.g. sigmoidal,  
 177 power law, linear, exponential). Note: continues on second page.

y	x	c <sub>1</sub>	c <sub>2</sub>	c <sub>3</sub>	c <sub>4</sub>	r <sup>2</sup>
$y = c_1 + \frac{c_2}{1 + \frac{\exp(c_3 - x)}{c_4}}$						
SSA <sub>405nm</sub>	log([OA]/[BC])	0.03	0.93	0.444	0.579	
SSA <sub>532nm</sub>	log([OA]/[BC])	0.085	0.91	0.623	0.520	
SSA <sub>781nm</sub>	log([OA]/[BC])	0.10	0.90	0.700	0.538	
AAE <sub>405-532</sub>	log([OA]/[BC])	1.25	7.81	2.298	0.554	
log(MAC <sub>BC,405nm</sub> )	log([OA]/[BC])	1.072	2.94	2.914	0.765	
log(MAC <sub>BC,532nm</sub> )	log([OA]/[BC])	0.94	2.56	3.721	0.900	
log(MAC <sub>BC,781nm</sub> )	log([OA]/[BC])	0.74	1.62	3.411	0.655	
log( <i>f</i> <sub>OA,int</sub> )	log([OA]/[BC])	0	-2.43	1.477	0.987	
log( <i>R</i> <sub>OA,BC</sub> )	log([OA]/[BC])	-1.76	3.70	0.462	1.823	
log(MAC <sub>BrC,405nm</sub> )	log([OA]/[BC])	1.072	-1.519	0.053	0.732	
AAE <sub>405nm</sub>	MCE	9.124	-7.476	0.884	0.0236	
AAE <sub>405nm</sub>	MCE*	9.723	-13.142	0.932	0.0452	
log(MAC <sub>BC,405nm</sub> )	MCE	3.523	-2.251	0.874	0.0198	
log(MAC <sub>BC,405nm</sub> )	MCE*	4.949	-5.073	0.882	0.0705	
log([OA]/[BC])	MCE*	4.030	-39.57	1.072	0.0500	
$y = c_1 + c_2 \cdot x^{c_3}$						
SSA <sub>405nm</sub>	MCE	0.954	-0.880	22.76		
SSA <sub>405nm</sub>	MCE*	0.939	-26.34	88.29		
log([OA]/[BC])	MCE	7.952	-8.272	3.351		
$y = c_1 \cdot x + c_2$						
<i>f</i> <sub>44</sub>	log([OA]/[BC])	-0.0097	0.0686			0.33
O:C	log([OA]/[BC])	-0.0345	0.407			0.17
H:C	log([OA]/[BC])	0.0228	1.737			0.27
log(MAC <sub>BrC,405nm</sub> )	log( <i>f</i> <sub>ON-OA</sub> )	0.322	0.446			0.33
log(MAC <sub>BrC,405nm</sub> )	log( <i>f</i> <sub>ON-OA</sub> )*	0.856	0.538			0.81
log(MAC <sub>BrC,405nm</sub> )	<i>f</i> <sub>60</sub> / <i>f</i> <sub>44</sub>	-0.396	0.043			0.11
log(MAC <sub>BrC,405nm</sub> )	<i>f</i> <sub>60</sub> / <i>f</i> <sub>44</sub> *	-2.242	0.803			0.96
log([ON]/[OA])	log([OA]/[BC])	-1.342	-0.320			0.47
log([ON]/[OA])	log([OA]/[BC])*	-1.342	-0.320			0.47
<i>k</i> <sub>BrC,405nm</sub>	MAC <sub>BrC,405nm</sub>	0.03104	-0.00177			0.99
<i>k</i> <sub>BrC,532nm</sub>	MAC <sub>BrC,532nm</sub>	0.0440	-0.00048			0.99

<b>y</b>	<b>x</b>	<b>c<sub>1</sub></b>	<b>c<sub>2</sub></b>	<b>c<sub>3</sub></b>	<b>c<sub>4</sub></b>	<b>r<sup>2</sup></b>
W <sub>BrC,405-532</sub>	AAE <sub>BrC,405-532</sub>	0.938	0.976			0.96
	$y = c_1 + c_2 \cdot \exp(-c_3 \cdot x)$					
MFR <sub>OA</sub>	log([OA]/[BC])	0.00175	0.1760	0.8520		0.157

\* Fits were performed to the Class averages, rather than to the individual burns.

178

179

180 **Table S3.** Literature imaginary refractive index and *MAC* values for biomass burning derived  
 181 brown carbon.

	$\lambda$ , nm	$k_{\text{BBOA}}$	$MAC_{\text{BBOA}}$ $\text{m}^2 \text{g}^{-1}$	Optical Measurement	Aerosol type sampled	Sampling Location or note	Literature
Laboratory	550	0.02-0.06		Aethalometer	Oak burning POA	-	(Saleh et al., 2013)
	550	0.015-0.04		Aethalometer	Pocosin Pine burning POA	-	(Saleh et al., 2013)
	550	0.0055- 0.022		Aethalometer	Galberry burning POA	-	(Saleh et al., 2013)
	400	0.038	1.1	UV/Vis (filter methanol extracts)	Pine/Oak wood burning	-	(Chen and Bond, 2010)
	405	0.015		Photo-Acoustic Spectrometer	Ponderosa Pine Duff burning	-	(Chakrabarty et al., 2010)
	405	0.0076		Photo-Acoustic Spectrometer	Tar balls from Alaskan Duff burning	-	(Chakrabarty et al., 2010)
	550		0.8-3.2	CLAP	Tar balls from liquid tar (turkey oak)	-	(Hoffer et al., 2016)
	405	0.01	0.35	Photo-Acoustic Spectrometer	Alaskan Peat	-	(Sumlin et al., 2017)
	355, 405, 532, 1064	0.012, 0.0065, 0.0024, 0.0023		Photo-Acoustic Spectrometer	Alaskan & Indonesian Peat	central values reported here	(Sumlin et al., 2018)
	600/400 ratio		0.04	Water soluble organic carbon	Florida peat	Ratio between wavelengths reported	(Sengupta et al., 2018)
Ambient	404	0.01	1.0-1.1	Photo-Acoustic Spectrometer	Wild fire, near- source emission	Four Mile Canyon, Colorado	(Lack et al., 2012a)
	470		1.01	Aethalometer	Biomass burning influenced	Beijing, China	(Yang et al., 2009)
	400	0.112	2.9	Filter transmission	Wood burning and biomass smoke aerosols	Savanna	(Kirchstetter et al., 2004)
	532	0.0016- 0.0019	0.029- 0.031	Photo-Acoustic Spectrometer	HULIS from biomass burning aerosols	Amazon basin	(Hoffer et al., 2006)
	Broadband	0.05-0.07		Airborne lidar	Upwind of forest fires	Northern Canada	(Wandinger et al., 2002)
	Broadband	0.07±0.03/ 0.04±0.01		White light optical particle counter	Open fire/ Smoldering phase	Urban Rehovot, Israel	(Adler et al., 2011)
	405	0.037	0.79 or 1.22	Photo-Acoustic Spectrometer	Residential biomass burning influenced	Fresno, CA	(Zhang et al., 2016)
	405		0.84	Photo-Acoustic Spectrometer	Residential biomass burning influenced	Fresno, CA	(Cappa et al., 2019b)
	405		2.3	Aethalometer	Biomass burning influenced	Guangzhou, China	(Qin et al., 2018)
	365		0.32	Water soluble organic carbon	Plume intercept – closest point to fire	Western US	(Forrister et al., 2015)
365		1.35	Water soluble organic carbon	Regional biomass burning	SE US	(Washenfelder et al., 2015)	



405	0.037		Water soluble organic carbon	Regional	Kanpur, India	(Shamjad et al., 2016)
405		0.7-1.3	Water soluble organic carbon	Bonfire festival	Rehovot, Israel	(Bluvstein et al., 2017)
405		0.6	Methanol soluble organic carbon	Prescribed burn	NW US	(Xie et al., 2017)
400, 600, 800	0.31, 0.26, 0.22		Electron loss	Asian outflow	Downwind of Asia	(Alexander et al., 2008)
800/400 ratio		0.26	Methanol soluble organic carbon	Ambient particles (ratio between wavelengths reported)	Athens, Georgia	(Phillips and Smith, 2017)
400, 550, 700	0.112, 0.030, 0.001		Acetone treatment + attenuation	African biomass burning	Southern Africa	(Kirchstetter et al., 2004)