# **Characterising Mass-resolved Mixing State of Black Carbon in Beijing Using a Morphology-Independent Measurement Method**

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#### S1. CPMA-SP2 set up



5 Figure S1: Example for one CPMA-SP2 set up with mono-mass and poly-dispersed scan in sequence

### S2. CPMA-SP2 inversion calculation

The two-variable distribution function  $\left(\frac{\partial^2 N}{\partial \log m_p \partial \log m_{rBC}}\right)$  is calculated by solving a Fredholm integral equation using a Twomey algorithm (Twomey, 1975). The number concentration distribution of rBC-containing particles can be related to the two-variable distribution for a single CPMA setpoint, *i*, using:

$$R_{i} = \frac{\partial^{2} N}{\partial \log m_{\rm rBC}} \bigg|_{i} \cdot \vec{\Gamma} \quad (1)$$

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where  $R_i$  is the number distribution of SP2 response at the *i* th CPMA setpoint, and  $\vec{\Gamma}$  is a kernel function applying a 15 trapezoidal rule approximation for the response of the instrument. The kernel function (Collins et al., 2002) can be described as:

$$\vec{\Gamma} = \overline{\mathsf{e}} \left( m_{\mathrm{p}} \right) \sum_{\phi=1}^{\phi_{max}} \bar{f} \left( d_{\mathrm{p}}, \phi \right) \int_{0}^{\infty} \Omega \left( m_{\mathrm{p}}, Z \right) \mathrm{dlog} m_{\mathrm{p}} \qquad (2)$$

20 where  $\overline{\in}(m_p)$  is a correction factor accounting for particle losses in the CPMA,  $\overline{f}(m_p, \phi)$  is the charge fraction where  $\phi$  represents the number of charges, and  $\Omega(m_p, Z)$  is the CPMA transfer function where Z represents the electrical mobility. The

correction factor  $\overline{\in}$  is determined by the ratio between total number concentration of the SP2 only measurement and total CPMA-SP2 scan number concentration over the same  $m_{rBC}$  range as the two-variable distribution. The correction factor  $\overline{\in}$  can be estimated by comparing the total number concentration measured by the CPMA-SP2 to the SP2-only number concentration, where the total number concentration measured by the CPMA-SP2 is found by iterating the Twomey algorithm. To provide

5 initial guess to the inversion scheme, it was assumed that all the particles were singly charged ( $\phi = 1$ ). Hence, the initial approximation for the two-variable distribution at scan  $i \left( \frac{\partial^2 N_{\text{APPROX}}}{\partial \log m_{\text{p}} \partial \log m_{\text{rBC}}} \right|_{i}$ ) is:

$$\frac{\partial^2 N_{\text{APPROX}}}{\partial \log m_{\text{p}} \partial \log m_{\text{rBC}}}\Big|_{i} = \frac{\frac{dN}{d \log m_{\text{rBC}}}\Big|_{i}}{\bar{f}(m_{\text{p},i},\phi=1)\hat{\beta}}$$
(3)

10 Where  $\frac{dN}{d\log m_{\rm rBC}}\Big|_i$  is the number distribution of  $m_{\rm rBC}$  of response at *i* th CPMA set points, and  $\hat{\beta}$  the is the analytical solution to the integral of a triangular CPMA transfer function (Broda et al., 2018). After the calculation of the initial approximation, the two-variable distribution is corrected by applying the Twomey algorithm iteratively until the total number concentration of the final two-variable solution is close to the CPMA-SP2 original number concentration measurement. Finally, a weighted-average smoothing of the two variable distribution is performed to reduce noise.





#### S3. Original CPMA-SP2 two variable distribution example

Figure S2: Example for one CPMA-SP2 two variable distribution

## S4. Mixing state of rBC-containing particles calculation

The detail of calculation is described in Riemer and West (2013). Consider there are N total rBC-containing particles in the atmosphere, each rBC-containing particle contains 2 types of material: rBC material and non-rBC material. The bins inside the two-variable distribution (see the example shown in Figure 4) represent different rBC-containing particles, and each bin

5 contains a distinct  $m_p$  value and a distinct  $m_{rBC}$  value. The mass fractions of rBC core and non-rBC coating for rBC-containing particle *i* ( $p_i^{rBC}$  and  $p_i^{non-rBC}$ ) are simply derived from the measured mass parameters:

$$p_i^{\text{rBC}} = \frac{m_{\text{rBC}}}{m_{\text{p}}} \qquad (4)$$
$$p_i^{\text{non-rBC}} = \frac{m_{\text{p}} - m_{\text{rBC}}}{m_{\text{p}}} \qquad (5)$$

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For the mass fraction in the population, the mass fraction for total rBC  $(p^{rBC})$  and total non-rBC  $(p^{non-rBC})$  are described as:

$$p^{\text{rBC}} = \frac{M_{\text{rBC}}}{M_{\text{p}}}$$
(6)  
$$p^{\text{non-rBC}} = \frac{M_{\text{p}} - M_{\text{rBC}}}{M_{\text{p}}}$$
(7)

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Where  $M_p$  and  $M_{rBC}$  are the bulk rBC-containing particle mass and the bulk rBC mass which is derived from the CPMA-SP2 two-variable mass distribution results.

Compared with the previous study by Healy et al. (2014) which calculates the mass fraction parameters based on the 20 morphology and density assumptions or the study by Bondy et al. (2018) which achieves the mass fraction through microscopy estimation, the CPMA-SP2 inversion metrics derives the mass parameters and mass fraction more directly. The information Shannon mixing entropy is calculated via these mass fractions. The connection between the mass fraction of particle *i* and the Shannon mixing entropy within the rBC-containing particle *i* ( $H_i$ ) is given:

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$$H_i = -p_i^{\text{rBC}} \ln p_i^{\text{rBC}} - p_i^{\text{non-rBC}} \ln p_i^{\text{non-rBC}}$$
(8)

The average particle mixing entropy can be expressed as:

$$H_{\alpha} = \sum_{i=1}^{N} p_i H_i \quad (9)$$

The Shannon entropy of the entire bulk population  $(H_{\nu})$  can be expressed as:

$$H_{\gamma} = -p^{\text{rBC}} \ln p^{\text{rBC}} - p^{\text{non-rBC}} \ln p^{\text{non-rBC}}$$
(10)

5 Later, the "diversity parameters" are used to describe the effective number of species in the single particle or the population of particles. The particle diversity for particle *i* is calculated from:

 $D_i = e^{H_i} \tag{11}$ 

10 The average single-particle diversity  $(D_{\alpha})$  and the bulk population diversity  $(D_{\gamma})$  are inferred from:

$$D_{\alpha} = e^{H_{\alpha}}$$
(12)  
$$D_{\gamma} = e^{H_{\gamma}}$$
(13)

15 Finally, the mixing state index ( $\chi$ ) term can be given from the diversity term above:

$$\chi = \frac{D_{\alpha} - 1}{D_{\gamma} - 1} \tag{14}$$

## **S5.** Extrapolation methods

20 Firstly, the number distribution for total rBC-containing particle mass  $m_{\rm p}$  has been fitted the log-normal distribution function:

$$\left. \frac{\mathrm{d}N}{\mathrm{dlog}m_{\mathrm{p}}} \right|_{i} = a_{1} \cdot \exp\left(-\left(\frac{\log\left(m_{\mathrm{p},i}\right) - b_{1}}{c_{1}}\right)\right) + a_{2} \cdot \exp\left(-\left(\frac{\log\left(m_{\mathrm{p},i}\right) - b_{2}}{c_{2}}\right)\right) \quad (15)$$

Where  $a_1$ ,  $b_1$ ,  $c_1$ ,  $a_2$ ,  $b_2$ ,  $c_2$  are all the constants for the distribution function.



Figure S3: Example for the extrapolation for the size distribution of rBC-containing particle number concentration.

As the rBC-containing particle number concentration distribution is derived, the ratio  $(R_i)$  between the rBC-containing particle 5 number concentration at the last measured bin and the number concentration at the extrapolated  $m_p$  bin *i* can be worked out:

$$R_{i} = \frac{\mathrm{d}N/\mathrm{dlog}m_{\mathrm{p}}|_{i}}{\mathrm{d}N/\mathrm{dlog}m_{\mathrm{p}}|_{m_{\mathrm{p}}=15 \mathrm{ fg}}} \quad (16)$$

Where  $dN/d\log m_p|_i$  is the number concentration of rBC-containing particles at the extrapolated  $m_p$  bin *i*, and 10  $dN/d\log m_p|_{m_p=15}$  is the number concentration of rBC-containing particles at the measured largest  $m_p$  bin (the bin with  $m_p = 15$  fg). Then the two-dimensional function distribution  $\frac{\partial^2 N}{\partial \log m_p \partial \log m_{rBC}}$  is extrapolated with two methods; One was based on fitting the two-variable distribution distribution as a function of  $m_{rBC}$  (Fit  $m_{rBC}$  Method), and other based on fitting the two-variable distribution as a function of the mass ratio between the  $m_p$  and  $m_{rBC}$  (Fit Ratio Method):

$$15 \quad \frac{\partial^2 N}{\partial \log m_{\rm p} \partial \log m_{\rm rBC}} \Big|_{\rm Fit \ m_{\rm rBC},(i,j)} = R_i \cdot \left(a_1 \cdot \exp\left(-\left(\frac{\log \left(m_{\rm rBC,j}\right) - b_1}{c_1}\right)\right) + a_2 \cdot \exp\left(-\left(\frac{\log \left(m_{\rm rBC,j}\right) - b_2}{c_2}\right)\right)\right) \quad (17)$$

$$\frac{\partial^2 N}{\partial \log m_{\rm p} \partial \log m_{\rm rBC}} \Big|_{\rm Fit \ Ratio,(i,j)} = R_i \cdot \left(a_1 \cdot \exp\left(-\left(\frac{\log \left(m_{\rm p,i}/m_{\rm rBC,j}\right) - b_1}{c_1}\right)\right) + a_2 \cdot \exp\left(-\left(\frac{\log \left(m_{\rm p,i}/m_{\rm rBC,j}\right) - b_2}{c_2}\right)\right)\right) \quad (18)$$

While  $\frac{\partial^2 N}{\partial \log m_p \partial \log m_{rBC}}\Big|_{(i,j)}$  represents the number concentration of two variable distribution at  $m_p$  bin *i* and  $m_{rBC}$  bin *j*;  $a_1$ ,

 $b_1$ ,  $c_1$ ,  $a_2$ ,  $b_2$ ,  $c_2$  are all the constants for the distribution function.

The examples for the extrapolation for the two periods are shown in Figure S4, the original measured part is on the left of the 5 black line while the extrapolated part is on the right. Particles with larger rBC core are preferred for the Fit Ratio method.



Figure S4: Examples for the two different extrapolation methods

# 10 S6. Bulk MR and NR-PM1/rBC mass concentration

Figure S5 illustrates the bulk MR value increased together with the NR-PM1 mass concentration and rBC mass concentration during the winter time. Compare to the winter campaign period, the average bulk MR varied less significantly in summer and the NR-PM1 and rBC concentration is lower.



Figure S5: CPMA-SP2 bulk MR results as a function of NR-PM1 Mass Concentration and rBC Mass Concentration





Figure S6: Frequency of Mixing State Index.

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periods in winter. The investigation of haze period show that large fraction of rBC-containing particles is well internal mixed. For the rBC-containing particles with their bulk MR around 1, the  $\chi$  varies from around 57% to around 75%. This widely variation range is the result of well internal mixture in summer and the poorly internal mixture in winter light polluted periods. Due to most of the rBC-containing particles with bulk MR~1 are found in the summer period, the  $\chi$  distributed for MR~1 is

5 similar to the distributed for the summer period.

## Reference

Bondy, A. L., Bonanno, D., Moffet, R. C., Wang, B., Laskin, A., and Ault, A. P.: The diverse chemical mixing state of aerosol particles in the southeastern United States, Atmospheric Chemistry and Physics, 18, 12595-12612, 10.5194/acp-18-12595-2018, 2018.

- Broda, K. N., Olfert, J. S., Irwin, M., Schill, G. P., McMeeking, G. R., Schnitzler, E. G., and Jäger, W.: A novel inversion method to determine the mass distribution of non-refractory coatings on refractory black carbon using a centrifugal particle mass analyzer and single particle soot photometer, Aerosol Science and Technology, 1-12, 10.1080/02786826.2018.1433812, 2018.
- 15 Collins, D. R., Flagan, R. C., and Seinfeld, J. H.: Improved Inversion of Scanning DMA Data, Aerosol Science and Technology, 36, 1-9, 10.1080/027868202753339032, 2002. Healy, R. M., Riemer, N., Wenger, J. C., Murphy, M., West, M., Poulain, L., Wiedensohler, A., O'Connor, I. P., McGillicuddy, E., Sodeau, J. R., and Evans, G. J.: Single particle diversity and mixing state measurements, Atmospheric Chemistry and Physics, 14, 6289-6299, 10.5194/acp-14-6289-2014, 2014.
- Riemer, N., and West, M.: Quantifying aerosol mixing state with entropy and diversity measures, Atmospheric Chemistry and Physics, 13, 11423-11439, 10.5194/acp-13-11423-2013, 2013.
   Twomey, S.: Comparison of constrained linear inversion and an iterative nonlinear algorithm applied to the indirect estimation of particle size distributions, Journal of Computational Physics, 18, 188-200, <u>https://doi.org/10.1016/0021-9991(75)90028-5</u>, 1975.
- 25