

# Experimental study of the aerosol impact on fog microphysics

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## Abstract.

Comprehensive field campaigns dedicated to fog life cycle observation were conducted during the winters of 2010-2013 at the SIRTAs observatory in the suburb of Paris. In order to document their properties, in situ microphysical measurements collected during 23 fog events induced by both radiative cooling and stratus lowering are examined here. They reveal large  
5 variability in number, concentration and size of both aerosol background before the fog onset and fog droplets according to the different cases. The objective of this paper is to evaluate the impact of aerosol particles on the fog microphysics. To derive an accurate estimation of the actual activated fog droplet number concentration  $N_{act}$ , we determine the hygroscopicity parameter  $\kappa$ , the dry and the wet critical diameter and the critical supersaturation for each case by using an iterative procedure based on the  $\kappa$ -Köhler theory that combines cloud condensation nuclei (CCN), dry particle and droplet size distribution measurements.

10 Our study reveals low values of the derived critical supersaturation occurring in fog with a median of 0.043 %. Consequently, the median dry and wet activation diameters are 0.39  $\mu\text{m}$  and 3.79  $\mu\text{m}$  respectively, leading to a minor fraction of the aerosol population activated into droplet. The corresponding  $N_{act}$  values are low with median concentrations of 53.5  $\text{cm}^{-3}$  and 111  $\text{cm}^{-3}$  within the percentile 75th. The activated fraction of aerosols exhibit remarkably low correlation with  $\kappa$  values, which reflects the chemical composition of the aerosols. On the contrary the activated fraction exhibits a strong correlation with the  
15 inferred critical diameter throughout the field campaigns. This suggests that the variability in the activated fraction is mostly driven by particle size, while variations in aerosol composition is of secondary importance. Moreover our analysis suggests that the supersaturation reached in fog could be lowered by the aerosol number concentration, which could contribute to the sink term of water vapour during the radiative cooling. Although radiative fogs are usually associated with higher aerosol loading than stratus lowering events, our analysis also reveals that the activated fraction at the beginning of the event is similar for both  
20 types of fog. However the evolution of the droplet concentration during the fog life cycle shows significant differences between both types of fog.

This work demonstrates that an accurate calculation of supersaturation is required to provide a realistic representation of fog microphysical properties in numerical models.

## 1 Introduction

As they reduce visibility, fog events strongly perturb the aviation, marine and land transportation. Gultepe et al. (2009) pointed out the extremely high level of human losses and financial cost related to fog and low visibility events. Moreover present numerical weather prediction models are usually unable to predict the exact location and time evolution of the life cycle of a fog layer (Zhou and Ferrier, 2008; Van der Velde et al., 2010; Bergot, 2013; Boutle et al., 2018).

The occurrence and development of fog result from the nonlinear interaction of competing radiative, thermodynamic, microphysical and dynamical processes (see the review of Gultepe et al. (2007b)). Water droplets are formed by heterogeneous nucleation of aerosol particles when the relative humidity (RH) exceeds 100 %. The Köhler theory (Köhler, 1936) forms the basis of our understanding of cloud droplet formation. The ability of aerosol particles to act as Cloud Condensation Nuclei (CCN) depends largely on their size, composition and phase state. The number size distribution of the droplets activated during the cloud formation phase depends on the supersaturation values reached by the air mass (Pruppacher et al., 1998). Particles having a critical supersaturation below the maximum value are activated and will then further grow by water vapour diffusion as long as the RH remains high enough, while the others particles remain at their equilibrium diameter. At the same time, water droplets are the particles whose diameter are higher than the critical threshold corresponding to that critical supersaturation.

Because of the weak supersaturation inside the fog layers (Hudson, 1980; Pandis et al., 1990; Svenningsson et al., 1992; Hammer et al., 2014; Boutle et al., 2018) the separation between activated fog droplets and non activated particles, so-called hydrated aerosols, is not explicit (Frank et al., 1998). To date the microphysics in Numerical Weather Prediction (NWP) models takes into account activated droplets only. Recently a two moment microphysical cloud scheme, which relies on the prognostic evolution of an aerosol population to predict cloud droplets, has been developed for low supersaturation clouds (Thouren et al., 2012) and needs to be tested against observations. It is thus essential to evaluate precisely total droplet number concentration values and their corresponding size range. Currently their definition mainly deals with the instrumentation limitation that generally begins at  $2 \mu\text{m}$  in diameter to measure droplet size distribution. However Hammer et al. (2014) and Elias et al. (2015) experimentally determined threshold diameter higher than  $2 \mu\text{m}$  to characterize droplets. As smallest hydrated particles are more numerous, considering a smallest threshold diameter could lead to an significant overestimation of the droplet concentration. Moreover as shown by Boutle et al. (2018) an accurate representation of droplet concentration for fog could also impact climate projection.

As a general rule in cloud physics, increase in aerosols causes an increase in the droplet number concentration for a given cloud type (Gulpepe and Isaac, 1996; Ramanathan et al., 2001). Hudson (1980) and Liu et al. (2017) found systematic differences in the fog microstructure between fogs formed in maritime, continental and urban air masses with a simultaneous increase of fog condensation nuclei and aerosol concentration. Recent numerical simulations exhibit a strong positive correlation between aerosol and droplet concentrations (Zhang et al., 2014; Stolaki et al., 2015; Maalick et al., 2016). They also predict fog events with high values of droplet concentrations, typically few hundreds per  $\text{cm}^{-3}$  or more. By contrast, numerical study of Bott (1991) shows that an increase in aerosol particles leads to a decrease of the supersaturation that in turn decreases the activated droplet number.

Recently, Boutle et al. (2018) have obtained impressively low droplet concentrations with LES (Large Eddy Simulation) by numerically solving the condensation equation at every time step and grid point, but they had very few activable aerosol particle.

5 In this study, we examined in situ microphysical measurements collected during the 3 winters campaigns to derive typical fog droplet concentrations in semi-urban condition. From October 2010 to March 2013 a suite of active and passive remote sensing and in situ sensors were deployed at the Instrumented Site for Atmospheric Remote Sensing Research (SIRTA) in the suburb of Paris (Haeffelin et al., 2010). Dupont et al. (2015) analysed 117 fog events to provide typical values of thermodynamical and radiative variables for the fog formation, the mature and dissipation phase. Burnet et al. (2012) have shown that, in term of droplet number concentration and effective diameter, the microphysical properties of fog events sampled at SIRTA present a large variability. Elias et al. (2015) examined the data collected during Nov 2011 and assessed the contribution of hydrated aerosols that led to the extinction of the visible radiation in the mist-fog-mist cycle. Hammer et al. (2014) focused on the activation properties of developed fog events observed during the winter 2012/2013, but they did not provide activated particle concentrations values.

15 Here, we determine the hygroscopicity parameter, the dry and wet critical diameters, and the critical supersaturation by using an iterative procedure based on the  $\kappa$ -Köhler theory that combines CCN measurements, dry particle size distribution and composite wet particle size distribution at ambient humidity. The objectives are to derive accurate estimations of the fog droplet concentrations and to investigate the impact of aerosols on the fog microphysics. This paper provides a comprehensive data set on activation properties and bring explanation on the aerosol impact on activation process.

20 Data and method are described in Sect. 2. and 3., respectively. Results are presented in Sect. 4 with first the statistics of fog activation properties of the 23 fog events analysed in this study. The link between aerosol particles and fog droplets are examined in Sect. 4.2, and the impact of CCN concentration on fog microstructure is discussed in Sect. 5. Conclusions are finally given in Sect. 6.

## 2 The dataset

### 25 2.1 Instrumentation

Data presented here were collected at the SIRTA observatory in the framework of the ParisFog field campaigns (Haeffelin et al., 2010). During the winters 2010 to 2013, specific instrumentation were deployed for the PreViBOSS project (Elias et al., 2012) to provide continuous observation of aerosol and fog microphysics. The experimental set-up was already presented in Burnet et al. (2012), Hammer et al. (2014), Elias et al. (2015) and Dupont et al. (2015). The instruments used in this study are listed in Table 1.

Particle size distribution at ambient humidity is derived from a combination of two optical spectrometers: the WELAS-2000 (Palas GmbH, Karlsruhe, Germany) and the FM-100 (Droplet Measurement Technologies Inn., Boulder, CO, U.S.A.) that do not cover the same size range. The WELAS-2000 (hereafter referred as WELAS) provides particle size spectrum between 0.4

and 40  $\mu\text{m}$  in diameter according to the constructor. However the detection efficiency decreases drastically below  $\sim 1 \mu\text{m}$  (Heim et al., 2008; Elias et al., 2015) resulting in a strong underestimation of the concentration of the submicronic particles. Hammer et al. (2014) choose to consider only data with diameter larger than 1.4  $\mu\text{m}$ . Statistics over the whole data set reveals that the most frequent mode diameter of the WELAS size distribution is 0.96  $\mu\text{m}$ . Thus we choose to use this value as the

5 lowest threshold and only measurements for bin diameter larger than 0.96  $\mu\text{m}$  will be considered in this study. Note that the activated diameter in fog is expected to be larger than 1  $\mu\text{m}$ , thus this instrumental bias will not affect the results presented here, which is next confirmed in this study for SIRTA's fog. The sampling time period was fixed to 5 minutes as a compromise between time resolution and statistical significance of the measurements. Indeed this corresponds to a volume of air of 6.40  $\text{cm}^{-3}$  per sample.

10 The FM-100 provides 1 Hz droplet size distribution from 2 to 50  $\mu\text{m}$  in diameter according to the constructor. Hence providing a large overlapping range with the WELAS measurements. However Burnet et al., 2012 and Elias et al., 2015 show that these distributions only overlap each other at a diameter which fluctuates between 5 to 9  $\mu\text{m}$ . Their comparisons reveal a high discrepancy between both these probes with a large underestimation by the FM-100 for particle less than about 5  $\mu\text{m}$  but conversely a large underestimation by the WELAS for droplets larger than 10  $\mu\text{m}$  (Burnet et al., 2012; Elias et al.,

15 2015). This is illustrated in Fig. 1 which shows the size distributions measured by the WELAS (black) and FM-100 (cyan) for two fog cases with contrasting properties. Figure A.1 shows comparisons of the particle number concentration values of the FM-100 vs the integrated WELAS measurement over [2-10]  $\mu\text{m}$  diameter range that correspond to the four first bins of the FM-100. The default manufacturer's first four bins of the FM-100 are [2-4], [4-6], [6-8] and [8-10]  $\mu\text{m}$  which correspond to 9, 7, 4 and 3 class bins of the WELAS, respectively. Following WELAS sampling time period, FM-100 and WELAS data

20 points are 5 min average that represents 2851 samples for the whole data set. Figure A.1 confirms that the FM-100 strongly underestimates the particle counts in the first size bin and that the WELAS underestimates the concentration of particles larger than 8  $\mu\text{m}$ . The measurements of the two instruments do not match perfectly in the size range of 4 and 8  $\mu\text{m}$ , which reflects the large fluctuations of the crossing diameter over this range. To derive a composite size distribution, Elias et al. (2015) chose a constant value of 7  $\mu\text{m}$ . Because of this high variability, we compute here the composite size distribution by using

25 WELAS data up to 6  $\mu\text{m}$ , FM-100 data above 8  $\mu\text{m}$  and the average of both from 6 to 8  $\mu\text{m}$ . Results are illustrated in Fig. 1 with red segments corresponding to the junction between WELAS and FM-100 distribution. This procedure allows us to take advantage of the finer bin resolution of the WELAS for the smaller particles and also to reduce uncertainties of the FM-100 due to Mie ambiguities for smaller diameter as described by Spiegel et al. (2012) since our FM-100 was using the default manufacturer's bin threshold. Indeed, particularly for smaller diameter, the droplet concentration may be overestimated

30 or underestimated, a few droplets from adjacent classes can be included or some can be counted in an adjacent class (Gonser et al., 2012). Nevertheless because of the large discrepancies between the two sensors on the crossing diameter, next calculation using this crossing diameter also consider a standard deviation using the min and the max size distribution over this size range. Both instruments are located on a scaffolding at about 2.5 m high, close to a PVM-100 from Gerber Scientific Inc. used as a reference for the LWC measurements, which measures on the [2-50] $\mu\text{m}$  diameter. Measurements of visibility and its vertical

35 evolution are given by two Degreanne diffusometers (DF20+ and DF320) located at 4 m and 18 m above ground, respectively.

Confidence is given in FM-100 measurements by comparison of the integrated LWC over its size range with the LWC measured by the PVM-100 and the visibility trend, see Burnet et al, 2012 for 2010-2011 winter. A paper will follow on that issue.

Aerosol particles measurements are performed under dry condition by instruments placed in a shelter. The sampled air mass enters through an aerodynamic size discriminator PM 2.5 inlet and a dryer which reduces the relative humidity to less than 50 %.

5 This value could be higher than the efflorescence point for some particles (Tang et al., 1995). A scanning mobility particle sizer (SMPS) which provides the dry aerosol particle number size spectrum, consists of a differential mobility analyzer (DMA; TSI 3071) which selects particles from 10.6 to 496 nm and of a condensation particle counter (TSI CPC 3022). Hydrophilic filters in the DMA circuit should lower the relative humidity to less than 30% (Denjean et al., 2014) which should be enough to avoid water uptake issues. Another CPC (TSI 3025) measures the total particle number concentration from 2.5 nm to 2.5

10  $\mu\text{m}$ . Finally a continuous flow streamwise thermal gradient CCN chamber (Roberts and Nenes, 2005) is used to measure the CCN number concentration at 5 different supersaturations from 0.1 % to 0.5 %, by step of 0.1 %.

## 2.2 Selected fog cases and aerosol properties

During the three winter campaigns of 2010-2013, Dupont et al. (2015) report the occurrence of 117 fog events. However due to

15 instrument failures and technical difficulties in operating the whole set of instrument on a 24/7 mode, 42 events were sampled simultaneously with both the WELAS and the FM-100, and only a subset of 23 cases were also sampled with both the CCNC and the SMPS. They are listed in Table 2 with their classification type, RAD for radiative cooling fog and STL for stratus lowering ones as determined by the scheme of Tardif (2007), and their vertical development based on the comparison of both difusometers : a developed fog produces low visibility condition simultaneously at 4 m and 18 m while a thin fog produces low

20 visibility condition at 4 m only (Elias et al. (2009), Dupont et al. (2015)). About the same proportion (40 %) of RAD and STL fog events occurred at SIRTAsite (Haefelin et al., 2010), Dupont et al. (2015). Among the 23 fog events analysed here, 13 are radiation fog and 10 stratus lowering fog, 19 are developed and 4 thin. Thermodynamics conditions during the field campaigns are described by Dupont et al. (2015).

25 To characterize the aerosol background prior to a fog event, statistics of the total number aerosol particle concentration,  $N_a$ , as measured by the SMPS between [10-496] nm diameter, are computed over the last hour before the beginning of the fog event. This time lapse is chosen to characterize aerosol before fog event because CCNC, which allows to characterize aerosol chemistry and a sampling time of 20 min, that 1 hour time lapse is supposed to provide reliable measurements. Median, 25th and 75th percentiles are reported in Table 2.  $N_a$  values range from  $\approx 2000 \text{ cm}^{-3}$  to  $\approx 20000 \text{ cm}^{-3}$ . Consistently with

30 Haefelin et al. (2010) the smallest values are observed in westerly flow conditions while highest values are associated mostly with an easterly flow when SIRTAsite is exposed to continental conditions. Indeed SIRTAsite is located 25 km south-west of Paris in a semi-urban environment, composed of agricultural fields, wooded areas, housing and industrial developments (Haefelin et al., 2005) and is exposed to air mass charged with pollution originating from regional background according to Crippa et al. (2013).

These various conditions are reflected on the statistics that reveal a large case to case variability. This is illustrated in Fig. 2 that shows the scatter-plot of the droplet number concentration derived by the FM-100,  $N_{FM}$  on [2-50]  $\mu\text{m}$  diameter range as a function of  $Na$ , that represent the aerosol loading, for the 23 fog events (data in table 2). Symbols depend on the fog type : blue and red colors for STL and RAD fogs, respectively, and open and solid symbols for thin and developed fogs, respectively.

5 For STL fogs,  $Na$  ranges from 2000 to 11000  $\text{cm}^{-3}$  with a median value of 4340  $\text{cm}^{-3}$  and 25th and 75th percentiles of 2833 and 6942  $\text{cm}^{-3}$ , respectively, while for RAD cases,  $Na$  are spread over a larger range from 3000 to 20000  $\text{cm}^{-3}$  with a median value of 8822  $\text{cm}^{-3}$  and 25th and 75th percentiles of 5719 and 13094  $\text{cm}^{-3}$ , respectively. Thus over this three year campaign it appears that at SIRTA STL fogs are associated with lower aerosols loading than RAD ones, as already reported by Elias et al. (2015) for the month of November 2011.

10 There is also a significant difference between a STL and RAD fog with respect to  $N_{FM}$ . Despite that the maximum value is reached for a STL fog (147  $\text{cm}^{-3}$  for the case f7)  $N_{FM}$  for a STL fog is generally lower with median values between 38 and 83  $\text{cm}^{-3}$  and 25th-75th percentiles of 31-63 and 53-118  $\text{cm}^{-3}$ , for STL and RAD, respectively. Moreover, for a RAD fog, thin cases exhibit higher values of  $N_{FM}$  with median and 25th-75th percentiles of 97 and 73-131  $\text{cm}^{-3}$  compared to 64 and 49-81  $\text{cm}^{-3}$  for a developed fog. Therefore it appears that a radiative fog is also associated with higher droplet concentrations

15 than stratus lowering fog over the range [2-50]  $\mu\text{m}$  as measured by the FM-100. One can note that these values are rather low for continental condition (Wendisch et al., 1998; Garcia-Garcia et al., 2002; Gultepe and Milbrandt, 2007a; Niu et al., 2012; Price, 2011; Liu et al., 2011; Lu et al., 2013; Zhao et al., 2013; Boutle et al., 2018).

As shown in Fig. 2 the general trend points towards a slight increase in the number of droplets with the aerosol loading but the scatter is very large. In cloud physics the connection between an increase in aerosol particles and an increase in the cloud droplet number concentration (CDNC) has been supported by many in situ observations (Twohy et al., 2005; Lu et al., 2007; Levin and Brenguier, 2009) and even if the discrepancy is large in the compilation of the diverse results (Ramanathan et al., 2001) the general trend is much more pronounced to what is presently observed in Fig. 2. However, CDNC is derived here from FM-100 measurements only between [2-50]  $\mu\text{m}$  diameter. In warm clouds the supersaturation at the cloud base is high enough

25 to activate droplets whose size will increase by water vapour condensation during the ascent of the air parcel. The spectrum then gets narrower because the growth rate of a droplet is inversely proportional to its size, while the remaining interstitial particles keep their own equilibrium diameter at 100 % relative humidity. Higher up in the cloud the spectrum is widened by different processes such as turbulent mixing and collision-coalescence and new CCNs can also be activated. Droplet spectra vary considerably in space and in time but the general trend is that the activated droplet population is clearly separated from the

30 interstitial non-activated particles. See for example measurements into stratocumulus cloud as reported in Martin et al. (1994); Brenguier et al. (2011); Ditas et al. (2012). Measurements by optical counters such as FSSP or CDP over the range [about 2-50  $\mu\text{m}$ ] are considered to provide an accurate estimation of the droplet size distribution.

In contrast, supersaturation in fog is much lower and, as already pointed out by Hudson (1980), in continental air the droplet number size distribution does not exhibit a clear separation. This is clearly illustrated in Fig. 1. To derive an accurate estimate

35 of the CDNC it is thus essential to determine carefully the wet activation diameter in order to integrate the composite size

distribution derived from both WELAS and FM-100 measurements. The method used to discriminate the droplets from the hydrated aerosol particles using kappa-Köhler theory is described in the next section.

### 3 Methods

There are different methods to separate fog droplets from non-activated aerosol particles. A fixed value can be used as a rough estimate as described in previous studies by Noone et al. (1992); Hoag et al. (1999); Elias et al. (2009) among others, with thresholds ranging from 2.5 to 5  $\mu\text{m}$ . To take into account case to case variability, Elias et al. (2015) used WELAS measurements. The volume size distribution is fitted with two log-normal distributions and the transition diameter is defined as the intersection between them. Hammer et al. (2014) investigated the activation properties by measuring the total and interstitial dry particle number size distribution behind two different inlet systems, and with WELAS and CCN measurements. They compared two methods one by fitting surface distributions of SMPS+WELAS measurements similar to Elias et al (2014) and one by retrieving the dry activation diameter from the difference between interstitial and total particle size distribution measurements.

In this study we use a different approach based on an iterative procedure that combines dry particle distribution from SMPS and composite wet particle size distribution derived from both WELAS and FM-100 measurements. The CCN measurements are also used to derive the hygroscopicity parameter ( $\kappa$ ) needed to link the dry and wet activation diameter ( $D_d$  and  $D_w$ ) with the critical supersaturation peak ( $SS_{peak}$ ) using the  $\kappa$ -Köhler theory.

#### 3.1 Kappa-Köhler theory

The Köhler theory (Köhler, 1936) expresses the equilibrium saturation vapor pressure over a solution droplet considering the solute effect (Raoult) and water surface tension effect (Kelvin). Accurate information is needed on the particle dry diameter and about the chemical composition to determine whether or not it can act as a CCN. Recently Petters and Kreidenweis (2007) developed a method, named  $\kappa$ -Köhler, to describe the relationship between particle dry diameter and CCN using a single hygroscopic parameter,  $\kappa$ . This method allows to study the activation process without considering aerosols complex chemical composition (McFiggans et al., 2006). The formulation of the method is expressed in Eq.(1):

$$S(D) = \frac{D^3 - D_d^3}{D^3 - D_d^3(1 - \kappa)} \exp\left(\frac{4\sigma_{s/a}M_w}{RT\rho_w D}\right) \quad (1)$$

where S is the saturation ratio over a solution droplet, D is the droplet diameter,  $D_d$  is the dry droplet diameter,  $\rho_w$  is the density of water,  $M_w$  is the molar mass of water,  $\sigma_{s/a}$  is the surface tension of the solution air/interface (of pure water here), R is the universal gas constant, T is the temperature and  $\kappa$  is the hygroscopicity parameter.  $\kappa$  represents a quantitative measure of aerosol particles water uptake characteristics and CCN activity. The critical supersaturation  $SS_c$  and critical wet activation diameter  $D_w$  correspond to values at maximum of supersaturation. They are linked to a couple ( $D_d, \kappa$ ) through the following relation :

$$\kappa = \frac{4A^3}{27D_d^3 \ln^2 SS_c} \quad A = \frac{4\sigma_{s/a}M_w}{RT\rho_w} \quad (2)$$

$SS_{peak}$  represents the maximum supersaturation that the air mass experienced for a sufficient long time (Hammer et al., 2014), all particles whose  $SS_c$  is less than  $SS_{peak}$  are activated and will further grow by water vapour condensation as long as RH remains high enough while other particles remain stable at their equilibrium diameter at the actual RH value.

### 3.2 Method to determine fog activation properties

- 5 The iterative procedure used to derive fog activation properties is illustrated in Fig. 3. To derive  $D_d$  and  $D_w$ , simultaneously measured CCN concentration ( $N_{ccn}$ ), dry and wet aerosols size distribution were used. From that it follows that all particles become activated to droplets when their dry or wet size is similar to or larger than  $D_d$  and  $D_w$ . Indeed it can exist only one trio  $N_{ccn}$ ,  $D_w$  and  $D_d$  that are linked to  $\kappa$  and  $SS_{peak}$  through Eq.(2). For each fog event the activation properties are determined as follows :  $N_{ccn}$  is set corresponding to a value of  $D_d$ .  $D_w$  and  $SS_{peak}$  are further calculated by using Eq.2
- 10 for  $SS_{peak}$  and by numerically searching for the maximum of Eq. 1 for  $D_w$  for a given value of  $\kappa$ . Then the integral of the droplet distribution from  $D_w$  provides the droplet concentration  $N_d$ . Iterations are made on  $N_{ccn}$  until  $N_{ccn}$  equals  $N_d$  : this is the activated aerosol number concentration for this case hereafter referred as  $N_{act}$ .  $N_{act}$  corresponds to the concentration of activated aerosols at  $SS_{peak}$ . Two sets of iteration are made with the two  $\kappa$  values and the average particle size distribution. To take into account the size distribution variability within the selected time period two other sets of iteration are made by using
- 15 the average size distribution +/- one standard deviation corresponding to the extrema of the WELAS and FM-100 crossing diameter, respectively, and with the  $\kappa$  value maximizing the scatter.

Our method to determine fog activation properties has to assume that the aerosol particles are internally mixed. However atmospheric aerosols are frequently externally mixed with particles of different sizes exhibiting different chemical composition.

- 20 Discussion on this hypothesis is presented in section 3. Moreover, surface tension properties (surface tension) of atmospheric aerosols could be modified by surfactant. This effect has largely been attributed to organics when the organic concentration in the aerosol population is sufficiently large (Facchini et al., 2000; Noziere et al., 2014). The extent to which this will occur is strongly dependent upon the specific molecular properties of the organic molecules. Unfortunately exploring these aspects would need to combine CCN measurements and a detailed analysis of the molecular species contained in organic particles.
- 25 This is out of the scope of this manuscript. In addition, number concentration of aerosol particles larger than 496 nm was assumed to be negligible in our calculation.

- The air mass is sampled through a PM2.5 head and is dried to RH < 50 % before entering the SMPS and the CCN chamber. Hence during a fog event activated aerosols larger than 2.5  $\mu\text{m}$  are missed. Aerosol properties of the air mass must thus be characterized before the occurrence of a fog event. In addition the aerosol particle distribution changes continuously according
- 30 to the boundary layer evolution and aerosol sources (wood burning, road traffic, ...). This is illustrated in Fig. A.2 showing a large variability in the aerosol particle distribution along the diurnal cycle for the fog event f6. There is a large variability along the diurnal cycle. To characterize the air mass just before the fog event, we average data recorded during the last hour before the fog onset. This corresponding time interval is delimited by red segments in Fig. A.2.



CCN chamber measurements are used to derive  $\kappa$  by using Eq.2 knowing  $SS_c$  and  $D_d$ . The chamber supplies the concentration of activated aerosol particles  $N_{ccn}$  at five supersaturations from 0.1 % to 0.5 %. Scanning takes about 20 minutes. Therefore 3 scans can be made during the 1 hour interval thus provides meaningful estimation of the aerosol hygroscopicity. However supersaturation values in the chamber are greater than 0.1 % and it is well known that  $SS_{peak}$  in fog is generally lower than 0.1 % (Hammer et al., 2014; Ming and Russell, 2004; Svenningsson et al., 1992; Hudson, 1980). Therefore assumptions must be made to extrapolate  $\kappa$  at lower supersaturation. This is illustrated in Fig. A.3 that shows statistics of  $\kappa$  as a function of supersaturation for the selected time period of Fig. A.2. Data points correspond to calculations with CCN data at each supersaturation and the SMPS aerosol size distribution averaged over the one hour time interval. Solid and dashed lines superimposed to the data points correspond to mean and mean  $\pm$  one standard deviation values, respectively. Two extrema of  $\kappa$  are defined.  $\kappa_{inf}$ , which is computed as the mean value of  $\kappa$  at 0.1 % minus standard deviation, corresponds to the lowest expected value.  $\kappa_{sup}$  is calculated as the linear extrapolation of the mean  $\kappa$  at 0 % plus one standard deviation and corresponds to the highest expected value. Both values are indicated in Fig. A.3 by blue and red dots, respectively. These two extreme values of  $\kappa$  will then provide extreme values of activation properties for a given dry activation diameter.

Once droplets have been activated, their size distribution evolves during the fog life cycle. They can grow by water vapour diffusion or by coalescence with other droplets due to gravitational and turbulent motions and they can also evaporate in case of mixing with clear air or changes in temperature. Following Noone et al. (1992), the mixing with clear air has similar consequences that a external mixture for aerosols, that means that some particles can deactivate and their diameter can be under the critical wet diameter because of evaporation. In convective clouds, the activation occurs mainly within the first tenth of meters above the cloud base and it is thus possible to directly sample the resulting droplet spectra with instrumented aircraft. By contrast, in case of fog, the activation first occurs at the fog onset, the ensuing vertical development of the fog layer depends on many processes among which the radiative cooling at the fog top plays a key role. The fog onset is generally defined by a drop of visibility below the 1 km threshold, but this is still subject to debate (see discussion in Elias et al. (2015)). Time series show indeed various case to case CDNC evolution during the first fog hour of the fog event, with an instantaneous formation, or conversely, a slower increase towards a stable value as illustrated in Fig. A.2. To estimate the most representative fog droplet size distribution of the activated distribution, we average the composite wet particle size distribution derived from both WELAS and FM-100 measurements over a time interval from the beginning of the fog event during which the CDNC reaches a stable value for a sufficient long time. On average, this time period is selected from 30 min to 1 hour after the fog beginning. It is delimited by blue segments in Fig A.2. We therefore also assume that local measurements at 2 m height are representative of the fog layer.

### 3.3 Method validation

Fig. 4 illustrates the impact of  $\kappa$  variability for two fog cases of Fig. 1. The average aerosol dry size distribution measured by the SMPS recalculated at 100 % RH with the  $\kappa_{inf}$  and  $\kappa_{sup}$  values (blue and purple lines, respectively) are superimposed to the composite fog droplet number size distribution. Vertical segments indicate the corresponding  $D_w$  diameters. These values

are 3.37 and 5.45  $\mu\text{m}$  corresponding to  $N_{act}$  of 91 and 61  $\text{cm}^{-3}$ , respectively for the f6 case (Fig. 4-a); and 2.15 and 3.08  $\mu\text{m}$  which correspond to  $N_{act}$  of 116 and 86  $\text{cm}^{-3}$ , respectively for the f20 case (Fig. 4-b). The dashed ones correspond to the two calculations which produce extreme values and thus delimit the range of possible values. If the spectrum remains constant as in Fig. 4-a, these values are rather the identical. But if the shape of the spectrum evolves during the one hour time period, the range of possible values increases as in Fig. 4-b with extreme values of 1.4 and 4.75  $\mu\text{m}$  corresponding to  $N_{act}$  of 266 and 21  $\text{cm}^{-3}$ , respectively, leading to an uncertainty of a factor 10 for  $N_{act}$ . There is a good agreement between both distributions on the overlap area corresponding to hydrated (non-activated) particles. This indicates that the hygroscopic growth of all dry particles measured before the fog event with the derived  $\kappa$  values is consistent with the ambient measurements at the beginning of the fog. This also allows to neglect the mixture issue. Indeed the dry (growth) and wet aerosols particles shape could not fit with our assumption for a significant external mix as there is no reason for an external mix to be homogeneously distributed on the dry aerosol number size distribution. Nevertheless, a significant external mix would have implied less aerosols particles activated per diameter with our method that would lead to an overestimation of  $D_d$ .

Our method provides a satisfactory estimation of  $\kappa$  and the selected time periods are adequate to estimate the activation properties.

## 4 Results

### 4.1 Fog activation properties

Table 3 presents results for the 23 fog events with mean values and uncertainty intervals for each parameter  $\kappa$ ,  $D_d$ ,  $D_w$ ,  $SS_{peak}$  and  $N_{act}$ . The mean value is the average between the two values obtained from the above procedure with the average wet composite size distribution and the two  $\kappa$  values. The uncertainty interval is determined by the extreme values resulting from the average wet composite size distribution minus (plus) one standard deviation and the lowest (highest)  $\kappa$  value, respectively.

Values of  $\kappa_{inf}$  and  $\kappa_{sup}$  are relatively close. Indeed the relative standard deviation to the average value  $\kappa_{mean}$  ranges between 9 and 25 % except for the case f14 which reaches 35 %. It follows that  $\kappa_{mean}$  which corresponds roughly to  $\kappa$  value at  $SS \approx 0.05\%$ , is representative of the  $\kappa$  value at the actual supersaturation. As a result  $\kappa_{mean}$  values range from 0.09 to 0.3 with a median of 0.17. They are in agreement with those determined by Hammer et al. (2014) who found  $\kappa$  between 0.06 and 0.27 centred at 0.14 ( $SS \leq 0.11\%$ ) and Jurányi et al. (2013) that reported  $\kappa$  between 0.08 and 0.24 ( $SS=0.1-1\%$ ) at the same site. They are slightly lower than those observed for continental aerosols : Andreae and Rosenfeld (2008) for example suggested to use  $\kappa = 0.3 \pm 0.1$ . As already pointed out by Hammer et al. (2014) local emissions of road traffic and residential wood burning are likely responsible for these low values of  $\kappa$ . Fig 5 shows the scatterplot of the hygroscopic parameter as a function of the critical supersaturation for each of the 23 fog events with their uncertainty interval. Overall the variability of  $\kappa$  remain very low since 80% of of the  $\kappa$  values range between 0.13 and 0.27. This suggests that the possible variations in aerosol composition occurring during the campaign due to variations in air mass type or source contributions have limited effect on the CCN ability of aerosol in fog conditions. Based on this study, a single  $\kappa$  value of 0.17 could be used for further modelling studies in this area.

The mean dry activation diameter values,  $0.39 \mu\text{m}$  for the median and  $0.35\text{-}0.43 \mu\text{m}$  for the 25th-75th percentiles, is rather high, which indicates that only largest aerosol particle are activated. This is in agreement with previous studies of Noone et al. (1992), Ming and Russell (2004) and Hammer et al. (2014) at the supersaturation occurring in fog. The wet activation diameter  $D_w$  is  $3.79 \mu\text{m}$  in median and range from  $3.03 \mu\text{m}$  and  $4.67 \mu\text{m}$  for the 25th and 75 th percentiles. These values are slightly higher than the median value of  $2.6 \mu\text{m}$  reported by Hammer et al. (2014) but is consistent with the mean value of  $4 \mu\text{m}$  by Elias et al. (2015). The latter determined  $D_{w_{50th}}$  from a November 2011 dataset as the intersection between the two log-normal distributions fitting the particle volume distribution measured by the WELAS. That is also highly consistent with statistics on no-fog event at SIRTAs whereby particle diameters as measured by WELAS do not exceed  $2 \mu\text{m}$ . Except for the f22 event mean values of  $D_w$  are larger than  $2.37 \mu\text{m}$  throughout the field campaign. The differences between hydrated and activated particles can be seen in figures 4-a) and b) and corresponds fairly well to the average wet activation diameter derived for such cases that are equal to  $4.41$  and  $2.62 \mu\text{m}$ , respectively. Moreover it appears that the extreme values indicated by the dashed vertical segments largely maximize the uncertainty of the retrieval.

Critical supersaturation occurring in fog are very lows with median value of  $0.043 \%$  and 25th-75th of  $0.035\text{-}0.051 \%$ . This result was expected for fog conditions, as already shown in experimental and modelling studies (Hudson (1980), Svenningsson et al. (1992), Ming and Russell (2004), Hammer et al. (2014), Shen et al. (2018) and Boutle et al. (2018)).

Finally the corresponding concentration of activated particles  $N_{act}$  are  $53.5 \text{ cm}^{-3}$  for the median and  $28.5\text{-}111 \text{ cm}^{-3}$  for the 25th-75th percentiles. These values are slightly lower than the  $N_{FM}$  ones derived from FM-100 data (reported in Fig. 2) which give a median and 25th-75th percentiles of  $61$  and  $34\text{-}103 \text{ cm}^{-3}$ , respectively. Moreover the median of the average concentration values of particles with diameter in the range of  $[0.96\text{-}50] \mu\text{m}$  at the beginning of the fog events reaches  $389 \text{ cm}^{-3}$  with 25th-75th percentiles of  $260\text{-}660 \text{ cm}^{-3}$ . It follows that only a small fraction of the fog hydrated particle can be activated to droplets due to low critical supersaturations encountered in fog. Obviously because of the uncertainty interval  $N_{act}$  is very broad. For instance this uncertainty is twice the  $N_{act}$  value itself on average but could be as high as a factor 6 for the f1 event. However as mentioned before,  $N_{act_i}$  and  $N_{act_f}$  are the extreme possible values derived by cumulating uncertainties on  $\kappa$  and on the variability of measurements during the one hour time period.

## 4.2 Impact of aerosol particles on fog droplet concentration

Mean values of  $N_{act}$  are shown in Fig. 6 as a function of the other activation parameters  $SS_{peak}$ ,  $D_d$  and  $\kappa$  for the 23 fog events.  $N_{act}$  values increase with  $SS_{peak}$  (Fig. 6-a)) and consequently decreased with  $D_d$  (Fig. 6-b). In contrast to Fig. 6-b), Fig. 6-c) reveals no trend between  $N_{act}$  and  $\kappa$ , suggesting that size matters more than chemical composition for a particle to be activated. Indeed the largest aerosol particles are activated at low SS independently of their  $\kappa$  values. There are some scatter

on Fig. 6-b) for values of  $D_d \leq 0.4\mu\text{m}$  which suggest that some variability of the dry aerosol size number distribution between the different cases occurs mainly below this threshold.

Our results extend to lower supersaturation the conclusion of Fitzgerald (1973) and Andreae and Rosenfeld (2008) on clouds that particles size plays a much greater role than composition in regulating cloud droplet nucleation. The fundamental reason is that soluble mass changes with the third power of particle diameter but only linearly with soluble fraction (Andreae and Rosenfeld (2008)). The color code in Fig. 7-a) shows that, at a given supersaturation, the highest values of the CCN ratio are associated with the highest values of the hygroscopic parameter (i.e. chemical composition), and that at a given CCN ratio, decreases as the supersaturation increases. However more samples are needed before a robust interpretation of that feature to be established.

To remove the influence of the aerosol number concentration we normalize  $N_{act}$  by the number concentration of activable aerosols  $N^*$ , defined as the number concentration of aerosol particles with diameter  $\geq 200$  nm which corresponds to the smallest dry diameter of Table 3. Fig. 7-b shows a good correlation between  $N_{act}$  normalized by  $N^*$  and the inferred critical diameter. We obtained similar correlation when  $N^*$  was calculated for the range of critical diameter inferred during the campaign (150 to 400 nm). This suggest that  $N^*$  using a single critical diameter can be considered as a good proxy of the number of activable particles. This can be understood since the supersaturation occurring in fog has a narrow range of value. CCN at a supersaturation between 0.02 and 0.11 % are mostly composed of accumulation mode particles.

Figure 8 reveals that almost no relationship exists between  $N_{act}$  and  $N^*$ . This surprising behaviour demonstrates that the concentration of fog droplets was roughly independent of the aerosol number concentration. This is most likely due to the low SS occurring in fog leading to very selective criteria on particle diameter for droplet activation. This is also consistent with the very large derived dry diameter. Nevertheless, this figure also shows that, for high concentration of activable particles ( $N^*$ ), lower  $N_{act}$  are obtained. It follows that the number concentration of fog droplets is mainly controlled by the peak supersaturation with low influence from the aerosol loading. This is illustrated in Fig. 9 where  $N_{act}$  are plotted as a function of  $SS_{peak}$  in log scale (black diamond, same data as Fig. 6-a)) superimposed to the statistics of CCN measurements (grey diamond). Compared to the compilation of CCN spectra reported in Andreae and Rosenfeld (2008) (their Fig. 2), our CCN data are spread between SCMS data (Hudson and Yum, 2001) and continental cases. However we observe a strong decrease of  $N_{act}$  for  $SS_{peak} \leq 0.1\%$  similar to the ASTEX data collected in maritime stratocumuli and to Hudson (1980) fog activation spectra measurements. This indicates that only a very small fraction of aerosols are activated for such low values of supersaturation.

These results have been fitted using three different equations commonly used in modelling. The typical CCN parameterisation proposed by Twomey (1959) :  $N_{CCN}=CS^k$ , where C represents the CCN number concentration at SS=1 % and the parameter  $k$  that varies significantly (Martins et al., 2009). The formula suggested by Ji et al. (1998) :  $N_{CCN}=N(1-\exp(-BS^k))$  where N is the total number concentration of CCN and B and k are empirical coefficients to be determined. And the more general description of the activation spectra proposed by Cohard et al. (1998) :  $N_{CCN}=CS_{v,w}^k F(\mu, \frac{k}{2}, \frac{k}{2} + 1; -\beta S_{v,w}^2)$  where C is proportional to  $N_{CCN}$  that would be activated when supersaturation tends to infinity and parameters k,  $\mu$  and  $\beta$  are adjustable parameters depending on the aerosol properties. Figure 9 shows that the parameterization of Ji et al. (1998) (green line) with

parameters as indicated by the legend, better reproduces the decrease of  $N_{act}$  for  $SS < 0.1 \%$  compared to the Twomey expression (blue line) but the drop is not as sharp as in the data. In contrast the parameterization of Cohard et al. (1998) (red line) provides the best fit of the data for lower values of  $SS$ .

5 Figure 10 shows the CCN concentration that would be activated at  $SS=0.1 \%$  as a function of the aerosol activable concentration for the 23 fog events. The expected strong correlation between activated particles and aerosol concentration then clearly appears. This confirms that the apparent independence observed in Fig. 8 is simply due to the low supersaturation values experienced in fogs. The symbol color depends on  $\kappa$  as preceding figures. Data are aligned according to their hygroscopic parameter. For a given aerosol particles activable concentration, the number of activated particles at  $SS=0.1 \%$  increases with  
 10  $\kappa$  which is consistent with Fig. 7-a).

One can note in Fig. 8 that the highest values of  $N_{act}$  decrease as  $N^*$  increases. Indeed, while aerosol particles activable concentration increases ( $N^*$ ), variability of  $SS_{peak}$  decreases and tends toward low values as reported in Fig. 11, thus less particles can be activated. This suggests that the supersaturation reached in fog could be limited by the activable aerosol concentration : if numerous aerosol particles are available they are efficient enough to uptake the water vapor excess and  
 15 therefore limit the supersaturation. Hudson (1980) found that aerosol particle concentration values do not have a great effect on fog supersaturation. Our result however supports the conclusions drawn by the numerical study performed by Bott (1991) that the higher the particle concentration, the lower the supersaturation is. Note that no trend appears with  $\kappa$  in Fig. 8, which means that the sensitivity of this process to the hygroscopicity is very weak. As suggested by Bott et al. (1990) this could be due to a lowering of vapour supersaturation with the hydration of the numerous aerosols particles.

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Fig. 6 compared to Fig. 2 shows that there are not obvious differences between STL and RAD cases. As a result, the median values of  $N_{act}$  are 40 and 58  $cm^{-3}$ , and the 25th-75th percentiles are 25-115 and 33-109  $cm^{-3}$ , for STL and RAD cases, respectively. Radiative cases have thus higher values of activated particles but the percentile intervals are rather similar and the difference is less pronounced than the factor of 2 obtained from  $N_{FM}$  derived from the FM-100 measurements only.

### 25 4.3 Impact of CCN concentration on fog microstructure

The mean fog droplet concentration  $Nd_{1h}$  averaged over the one hour time interval at the beginning of the fog event is plotted in Fig. 12-a) as a function of  $N_{act}$ . Obviously they are almost identical apart from small deviations that come from averaging the  $N_{act}$  values derived for both  $\kappa$  values. The color code in Fig. 12-a) shows no differences between radiative

Once droplets have been activated, they compete for the available water vapour: when their number increases the size they  
 30 can reach by water vapour diffusion grow decreases. This inverse relationship between the number and the size of droplets is clearly depicted in Fig. 12-b) that displays the concentration  $Nd_{1h}$  vs. the corresponding mean diameter  $Dm$ . Indeed the highest  $Nd_{1h}$  of 255  $cm^{-3}$  corresponds to the lowest  $Dm$  of 4  $\mu m$ , and for fog events with  $Nd_{1h} > 120 \text{ cm}^{-3}$  the mean diameter can not exceed 7  $\mu m$  while it reaches twice this value for concentration as low as 25  $cm^{-3}$ . However some scatter appears for samples with  $Nd_{1h} < 50 \text{ cm}^{-3}$ . Symbol colour corresponds to the mean LWC values as indicated by the legend.

Lowest values of  $Dm$  are associated with lowest values of LWC which suggests that a lack of available liquid water could also limit the droplet growth, even when there are few of them.

During a fog life cycle several processes contribute to the droplet size distribution. Droplets can grow by water vapour diffusion, or by collision-coalescence with other droplets due to gravitational and turbulent motions. Conversely, they can evaporate if the supersaturation decreases due to heating of the air mass for example or in case of mixing with clear air. Figures 13-a) and b) show the same plots except that  $Nd$ ,  $Dm$  and  $LWC$  correspond to mean values over the complete fog life cycle. The mean droplet concentration values are significantly lower especially for cases with high  $Nact$  values. For instance for cases with  $Nact > 50 \text{ cm}^{-3}$  the ratio  $Nd_{cycle}/Nact$  ranges from 0.25 to 0.92 with an average value of 0.58. Values reported in Brenguier et al. (2011) exhibit a clear separation between cumulus and stratocumulus clouds, with values from 0.32 to 0.56 (0.46 on average) and from 0.72 to 0.96 (0.87 on average), respectively, that were attributed to differences in entrainment-mixing processes in both cloud types. Here it appears that values are spanned over a large range with an intermediate average value. Worthy to be noted Fig. 13-b) reveals that  $Nd_{cycle}$  for a radiative fog can not exceed  $70 \text{ cm}^{-3}$  which suggests that this reduction is more pronounced in case of a radiative fog while STL cases seem less affected. Indeed, for cases with  $Nact > 100 \text{ cm}^{-3}$  the average  $Nd_{cycle}/Nact$  is 0.80 for a STL fog, a value close to stratocumulus clouds, while it drops down to 0.45 for RAD cases. This suggests that the processes involved in the reduction of the CDNC in a radiative fog are as efficient as the entrainment-mixing occurring in cumulus clouds. Note that many studies (Pilié et al., 1975; Choullarton et al., 1981; Gerber, 1991; Bergot, 2013; Maronga and Bosveld, 2017) among others, pointed out the key mechanism of the turbulent mixing in fogs. Recent work of Bergot (2016) has shown, with a LES study, evidence of the key role of dry downdraught at the top of the fog layer on dissipation. These downdraughts would even then reach the surface allowing the sun to warm it. Wærsted et al. (2017) using experimental data, confirms the importance of top processes for the fog evolution. Dilution ratio being as high as in cumulus cloud in radiative fog bring a new element on that point, meaning that turbulent mixing must be of importance.

Figure 13-b) shows that mean diameter values averaged over the fog life cycle are similar to the previous ones. The resulting cluster of data points does not reveal any general trend and the anti-correlation between size and number of droplets is no longer noticeable. Indeed, the evolution of microphysical properties during the fog life cycle is complex and highly varies from case to case and depends on many parameters. As already noticed by Burnet et al. (2012) in their Figure 3 various values of  $D_{eff}$  and  $N_{FM}$  can be observed for a given  $LWC_{FM}$ . A comprehensive analysis on this topic has been performed and will appear in a forthcoming paper.

## 5 Discussion

In this experimental study we derived accurate estimations of the activated droplet concentration at the beginning of a fog event. This was done through a careful estimation of the critical supersaturation and of the wet critical diameter which allowed us to integrate the composite size distribution derived from WELAS and FM-100 measurements. We have shown that the derived parameters are consistent with the parameterization of Cohard et al. (1998). They are also consistent with measurements reported in Hudson (1980) who developed the Isothermal Haze Chamber to extend the CCN measurements to supersaturation

range below 0.1 % and derived fog condensation nuclei (FCN). Indeed the author found effective supersaturation between 0.06 and 0.1 % as well as sharp changes in FCN spectra below 0.1 % with activated concentration at SS=0.4 % ranging from 0.8  $cm^{-3}$  over the sea to 250  $cm^{-3}$  for polluted cases.

To our knowledge, no other study has retrieved experimentally concentration of activated particles (and thus droplets) in fog. Usually measurements of droplet spectra reported in the literature were performed by single particle counters such as FSSP, CDP or FM-100 over the range [2-50]  $\mu m$ . Figure 1 shows typical droplets size distribution observed during the campaign with measures from 2  $\mu m$  up to 50  $\mu m$ .

The corresponding droplet concentrations are in the range of a tenth to several hundred  $cm^{-3}$  (Wendisch et al., 1998; Garcia-Garcia et al., 2002; Gultepe and Milbrandt, 2007a; Niu et al., 2012; Price, 2011; Liu et al., 2011; Burnet et al., 2012; Lu et al., 2013; Zhao et al., 2013; Boutle et al., 2018). In our study droplet concentrations are rather low for semi-urban conditions with maximum value of 150  $cm^{-3}$ . For instance, Liu et al. (2017) reported values in China as high as 1000  $cm^{-3}$ . Moreover our values of critical diameter are large with a median value of 3.8  $\mu m$ . Consequently, the use of FM-100 data over the range [2-50]  $\mu m$  will lead to an overestimation of the droplet concentration, because it includes deliquescent aerosol particles that are not activated. A noticeable exception comes from the f22 case with a maximum value of 264  $cm^{-3}$ . Surprisingly, the median  $N_{FM}$  value for this case is only 36  $cm^{-3}$  which is rather low. This is explained by the fact that this case also corresponds to the lowest critical diameter  $D_w=1.44 \mu m$ , therefore the contribution from WELAS data is the most important and emphasizes the underestimation of the FM-100 in the first classes (see Fig. A.1). Thus for cases with low  $D_w$  WELAS data must be taken into account to avoid an underestimation of the concentration. In a modelling purpose, distinguishing hydrated aerosols particles from droplets allows to calculate an accurate repartition of the vapour deposition on the droplet size distribution. Indeed, behaviour of hydrated particles and droplets is quite different, the second one grow as long as there is vapor in excess while the first one stay at an equilibrium diameter. The second one may grow enough to produce drizzle and have strong interaction with radiation. A clear distinction should improve the representation of processes in numerical weather prediction model and so on the visibilities forecast.

A surprising result is the observed decrease of droplet concentrations with increasing aerosol loading. This is in contrast with recent numerical simulations that exhibit a strong positive correlation between aerosol and droplet concentrations in fog (Zhang et al., 2014; Stolaki et al., 2015; Maalick et al., 2016). This could be due to a limitation of the supersaturation according to Bott (1991). Sensitivity study of Stolaki et al. (2015) indicates that the fog droplet concentration during the mature stage is 2.6 times greater (2.9 times lower) when doubling (halving) the CCN number concentration of the accumulation mode. Indeed at the fog onset, the droplet concentration at the surface reaches 467  $cm^{-3}$  in their simulation with a double CCN concentration. This value is much larger than that of our estimations which is  $\leq 150 \text{ cm}^{-3}$  except in one case where it reaches 264  $cm^{-3}$ . For their reference run also this value is as high as 304  $cm^{-3}$ . The CCN activation scheme used by the model follows the expression of Cohard et al. (1998). Given the CCN activation spectrum derived from their values of CCN size distribution parameters (geometric mean radius and standard deviation of 0.1525  $\mu m$  and 2.33, respectively) and concentration (540  $cm^{-3}$ ) and solubility of aerosols (0.4), this corresponds to an estimation of the maximum supersaturation of 0.15 %. Even for activated concentration of 250  $cm^{-3}$  the corresponding SS is still  $\simeq 0.1 \%$ . We have shown that CCN measurement at

SS=0.1 % exhibits a linear increase of the droplet concentration with the activable aerosol concentration. Consequently with such high supersaturation values it is consistent that their simulations produce a strong positive correlation between aerosol and droplet concentration.

More generally, one shall keep in mind that the use of the CCN activation spectra provides a satisfactory estimation of the activated concentration, providing that the maximum supersaturation is correctly diagnosed. For instance, a current limitation of schemes with adjustment to saturation with parameterized peak supersaturation is that the formulae do not take into account pre-existing liquid water in the model grid box (Thoueron et al., 2012). This leads to a significant overestimation of the supersaturation peak value that in turn will overestimate the activated concentration. Another limitation is the use of minimum values for vertical velocities in typical supersaturation parametrisation as pointed out by Boutle et al. (2018). They are not necessarily consistent with fog dynamic and can lead to an overestimation of the supersaturation peak value. As could do the using of the saturation adjustment assumption instead of a diffusion growth/evaporation equation according to Khain et al. (2015). Moreover, in sensitivity study in which the CCN concentration is increased (or decreased) while keeping other parameters constant, intrinsically introduces a strong dependency between droplet and aerosol particles. Indeed, the pre-existing hygroscopic aerosol particles contribute to the sink term of water vapour during the radiative cooling that will also reduce the supersaturation peak value at the fog onset and then limits the number of activated particles. Thus changing the aerosol properties in models using such a scheme, provides a useful way to modify the droplet concentration and study the impact of this latter on the fog life cycle. However it should not be used to assess the impact of aerosol properties themselves on the fog life cycle.

## 6 Conclusions

In situ microphysical measurements collected during 23 wintertime fog events sampled from October 2010 to March 2013 at the semi-urban site SIRTA near Paris has been examined to document their properties. They reveal a large variability of their characteristic values in terms of number concentration and size of fog droplets depending on the different cases, as well as various evolution of these properties during the fog life cycle. The aerosol background exhibits also a highly variable number concentration of particles  $> 10 \text{ nm}$  before the fog onset. The objective of this paper was to evaluate the impact of aerosol particles on the fog microphysics. As a first step, we focused on the relationship between aerosol and fog droplet number concentrations as we expected that they follow the same trend with more numerous fog droplets with an increase of the aerosol loading.

To derive accurate estimations of the actual activated fog droplet number concentration ( $N_{act}$ ), we determined the hygroscopicity parameter, the dry and wet critical diameters, and the critical supersaturation for the 23 events by using an iterative procedure based on the  $\kappa$ -Köhler theory which combined CCN measurements, dry particle distribution from an SMPS and composite wet particle size distribution at ambient humidity derived from WELAS and FM-100 measurements. These data are averaged over a one hour time period before and during the fog onset to characterize the air mass and the fog properties, respectively.



Values of the hygroscopicity parameter  $\kappa$  extrapolated at  $SS \simeq 0.05\%$  were found to range from 0.09 to 0.3 which are characteristics for urban aerosol. They are rather similar from case to case thus we recommended to use  $\kappa = 0.17 \pm 0.05$  for fog modelling studies in this urban area. Our study reveals low values of the derived critical supersaturations with median of 0.043 % and 25th-75th percentiles of 0.035-0.051 %. Consequently wet and dry activation diameters are high and the corresponding  $N_{act}$  are low for continental conditions with a median concentration of  $53.5 \text{ cm}^{-3}$  and 25th-75th percentiles of 28.5-111  $\text{cm}^{-3}$ . Our results depict a sharp drop of  $N_{act}$  as critical supersaturations decrease which is best fitted by the parameterization of Cohard et al. (1998).

Concentration of activated particles using a single critical diameter can be considered as a good proxy of the number of activable particles. No detectable trend between concentration of aerosol particles with diameter  $> 200 \text{ nm}$  and  $N_{act}$  was observed. But  $N_{act}$  could be modulated for high concentration of aerosols particles. Hydration of numerous aerosols particles could indeed lower the supersaturation. Our analysis corroborates modelling study by Bott (1991) suggesting that an increase of the aerosol concentration limits the SS values. In contrast the CCN data at 0.1 % supersaturation exhibits a strong correlation with these aerosol concentrations. We therefore conclude that the droplet number concentration is roughly independent of the aerosol one because the actual supersaturations reached in these fog events are too low. Given the very diverse origin of atmospheric particles, the complexity of their composition has long been seen as a major obstacle to modelling aerosol effects on fog properties. Understanding that particle's ability to act as CCN at the supersaturation reached in fog is largely controlled by aerosol size rather than composition greatly facilitates the treatment of aerosol impacts on fog microphysics in NWP models. Nevertheless, some reservation must be issued on that point. Aerosol size distribution impacts fog microphysics more than chemical composition on the measured range on the hygroscopic parameter during the three winters of observations. Higher or lower values of this parameter could bias this conclusion. Other measurements of aerosol and fog microphysics in other environmental conditions would be needed to pursue this issue.

Despite that a stratus lowering fog appears to be associated with lower aerosol loading than a radiation fog event, no significant differences were observed with respect to the droplet number concentration at the fog formation when calculations are performed by taking into account the wet critical diameter. In contrast, statistics over the complete life cycle indicate that a radiation fog is subject to a pronounced decrease in the droplet concentration while stratus lowering cases seem unaffected. In a radiation fog this decrease, which occurs mainly for events with  $N_{act} > 50 \text{ cm}^{-3}$ , is significant with a ratio of average droplet concentration to  $N_{act}$  ranging from 0.25 to 0.67. This reduction is similar to the dilution which results from entrainment-mixing in cumulus clouds meaning that mixing with clear air may be of importance for fog microphysic evolution. The expected inverse relationship between the number and the size of the droplets at the formation phase is clearly depicted except for some cases with low liquid water content values. However this trend is less pronounced on average over the fog life cycle.

Nenes and Seinfeld (2003) have shown that kinetic limitations in droplet activation may retard the growth of CCN sufficiently to limit the number of activated droplets formed. The extent to which potential kinetic limitations to droplet formation exist in fog conditions will be explored in a subsequent publication focusing on the temporal evolution of droplets spectra and visibility during fog life cycle.

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## ACRONYMS

**CCN** : Cloud Condensation Nuclei

**N<sub>d</sub>** : Concentration of particles above the wet critical diameter (derived from WELAS+FM-100)

**CN** : Cloud Nuclei

5 **N<sub>FM</sub>** : Concentration of particles measured by the FM-100 between [2-50] micron diameter

**N<sub>a</sub>** : Concentration of dry particles measured by the SMPS between [10-495] nm diameter

**N<sub>act</sub>** : Concentration of activated particles

**N\*** : Concentration of dry particles between [200-496] nm diameter (derived from SMPS)

**N<sub>ccn</sub>** : Concentration of activated aerosol measured by the CCN at different supersaturation

10 **D<sub>d</sub>** : Dry critical diameter

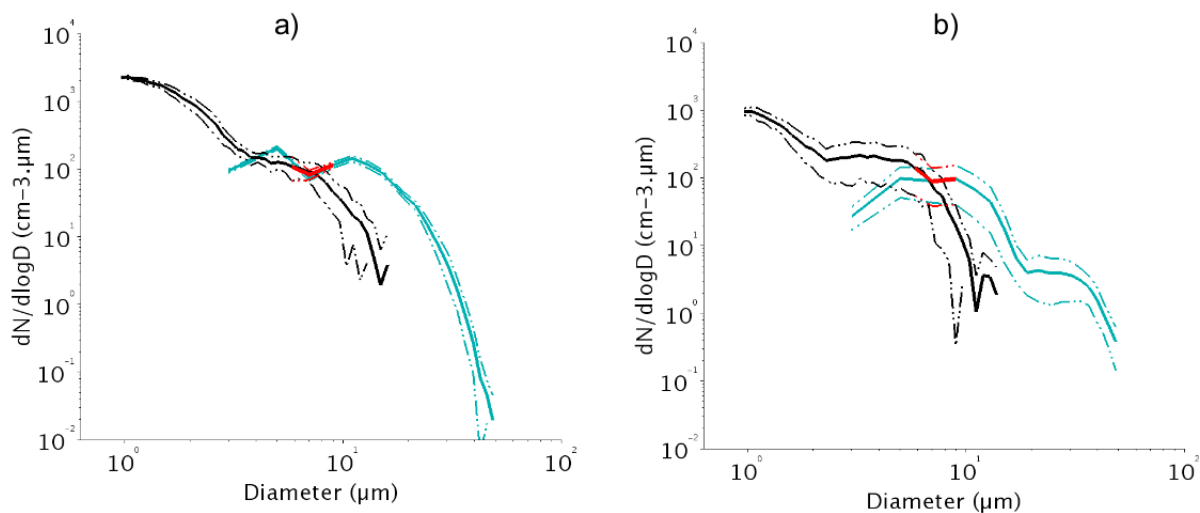
**D<sub>w</sub>** : Wet critical diameter

**D<sub>m</sub>** : Median diameter

**SS<sub>peak</sub>** : Critical supersaturation peak

**SS** : Supersaturation

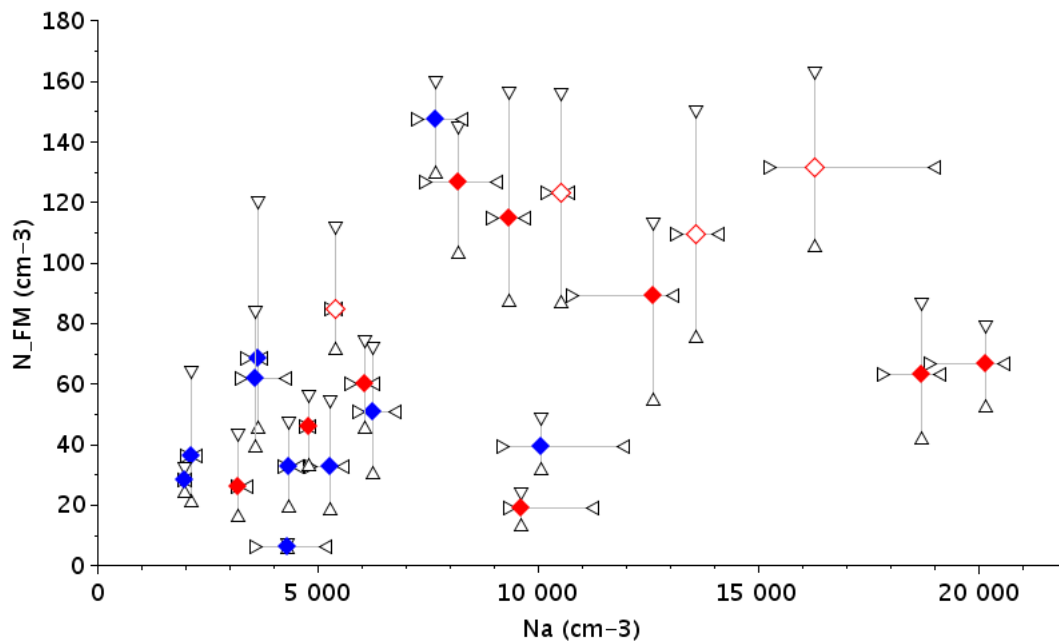




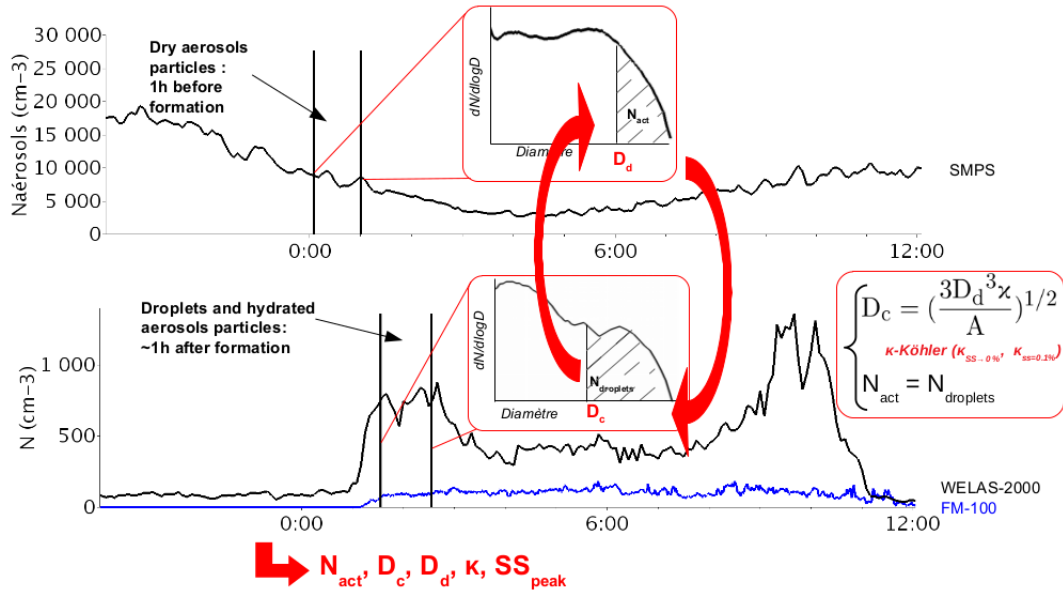
**Figure 1.** One hour average composite number size distributions at ambient humidity derived from WELAS (black) and FM-100 (cyan) at the beginning of the fog event for cases a) f6 and b) f20. The solid line represents mean value and the dotted lines mean  $\pm$  one standard deviation. The red segments correspond to the result of merging both distributions between 6 and 8  $\mu\text{m}$ .

**Table 1.** Instrumentation deployed during the campaigns that are used for this study.

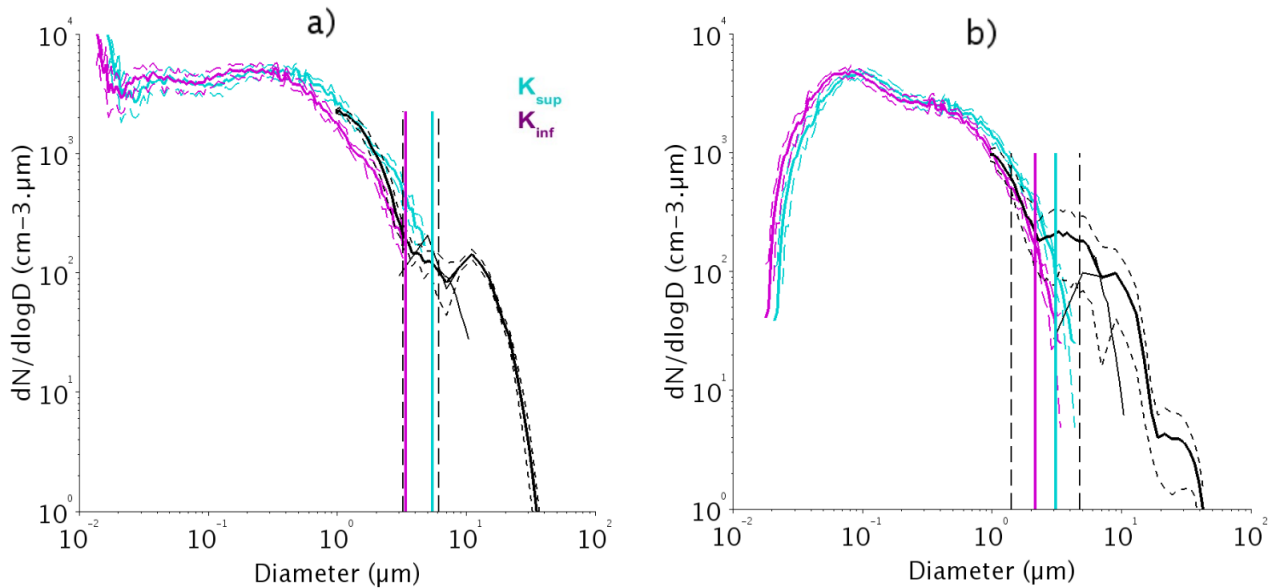
Instruments	Measured parameters	Time resolution
PALAS WELAS - 2000	Hydrated and activated aerosol particles number size distribution $D=[0.39-42] \mu\text{m}$	5 min
DMT Fog Monitor (FM-100)	Droplet number size distribution $D=[2-50] \mu\text{m}$	1 min
PVM Gerber	Liquid water content $D=[3-50] \mu\text{m}$	1 min
Degreanne DF20+	Horizontal visibility	1 min
TSI SMPS	Dry aerosol particles distribution $D=[10.6-493] \text{nm}$	5 min
CCNC DMT	CCN concentration $SS=[0.1-0.5] \%$	20 min for a cycle
CPC (TSI 3025)	Aerosol particles concentration $D=[2.5-2500] \text{nm}$	5 min



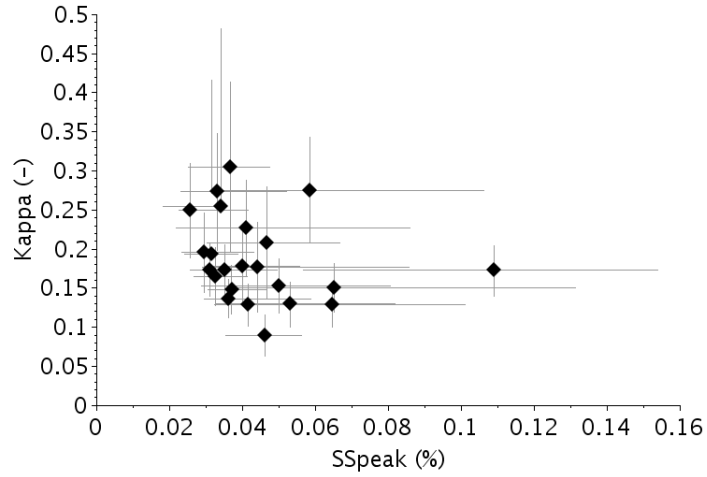
**Figure 2.** Droplet number concentration as derived from the FM-100 collected during the whole fog event as a function of the aerosol number concentration as derived from the SMPS measurements collected over the last hour before the fog beginning. Median values (diamond) are indicated for each case of table 2 with the 25th and 75th percentiles (triangle) to represent the variability. Symbols depend on the fog type STL (blue) and RAD (red) with open and solid symbols for thin and developed fogs, respectively.



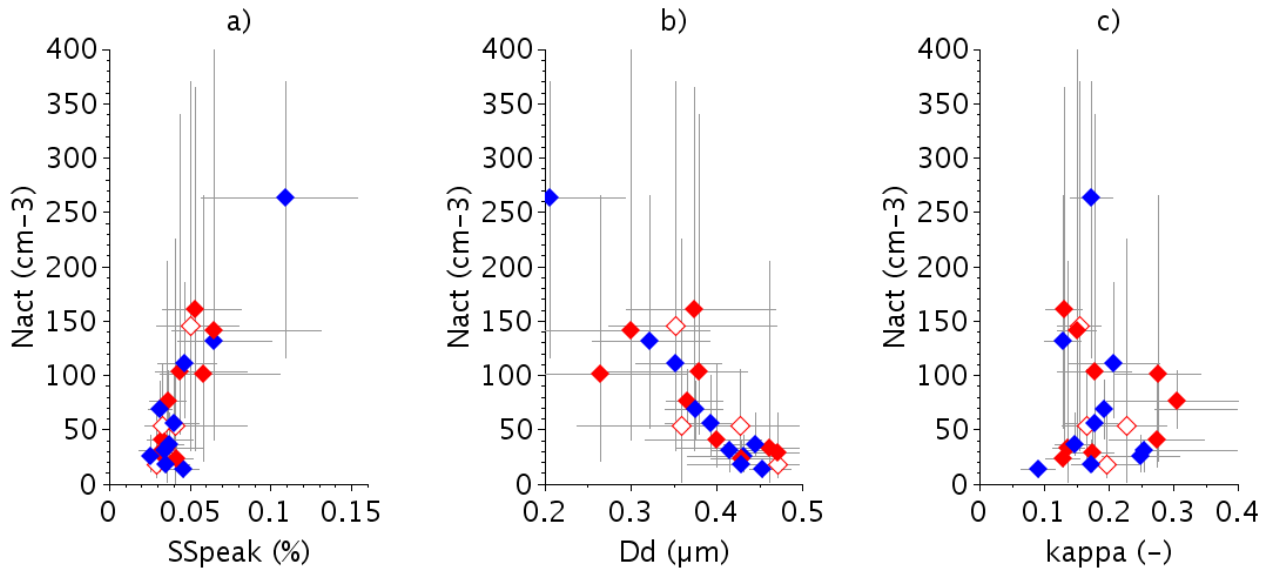
**Figure 3.** Schematic of the iterative method used to retrieve fog activation properties from dry and ambient humidity particle size distributions.



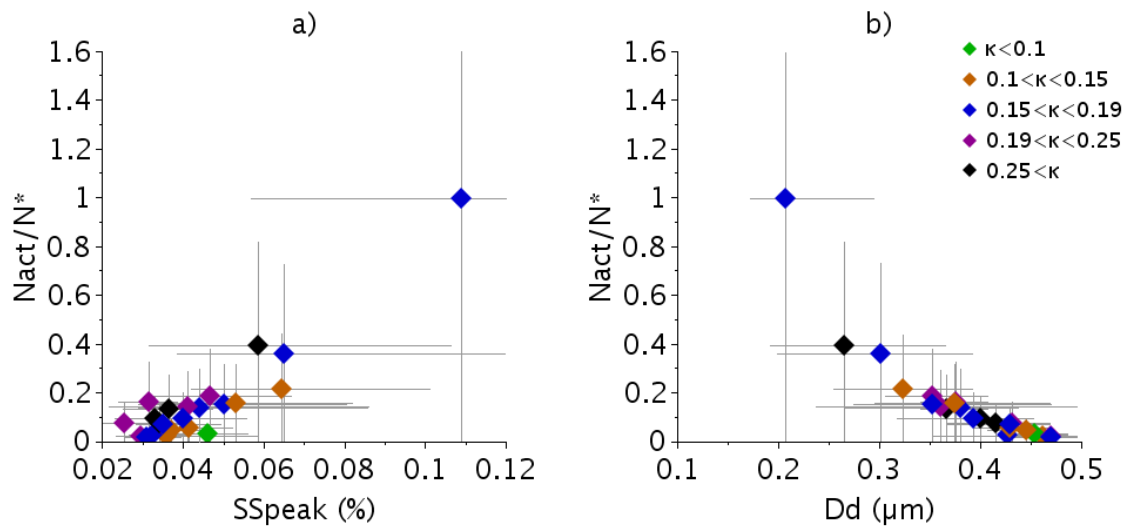
**Figure 4.** One hour wet number size distribution before the fog event resulting from the hygroscopic growth at RH=100% of the dry distribution measured by the SMPS calculated with  $\kappa_{inf}$  (purple) and  $\kappa_{sup}$  (cyan) for case a) f6 and b) f20. The in-fog composite size distribution of Fig. 1 is superimposed (black). The solid lines represent mean values and the dashed lines mean  $\pm$  one standard deviation. The vertical segments indicate the mean values of  $D_w$  corresponding to each  $\kappa$  (solid color lines) and the extreme values (dashed lines).



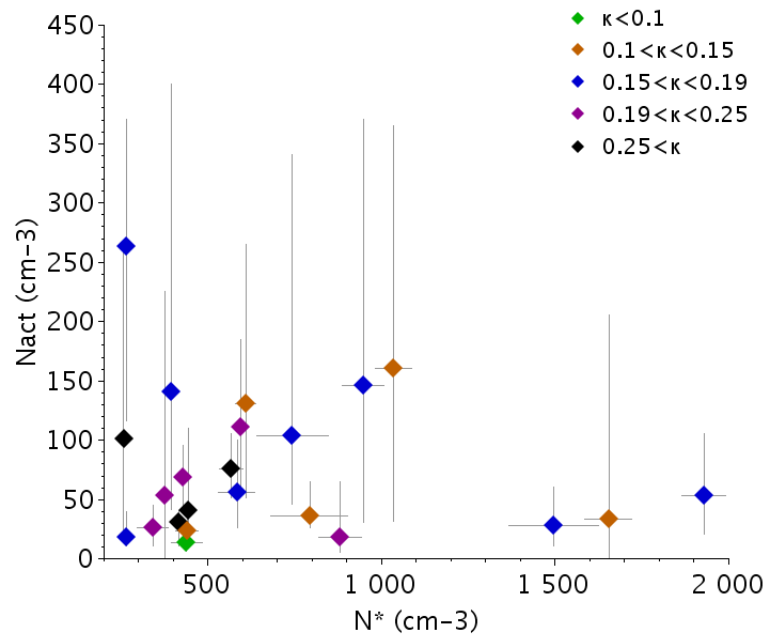
**Figure 5.** Hygroscopicity parameter  $\kappa$  as a function of the critical supersaturation for the 23 fog events. Diamonds correspond to mean values and error bars to uncertainty intervals.



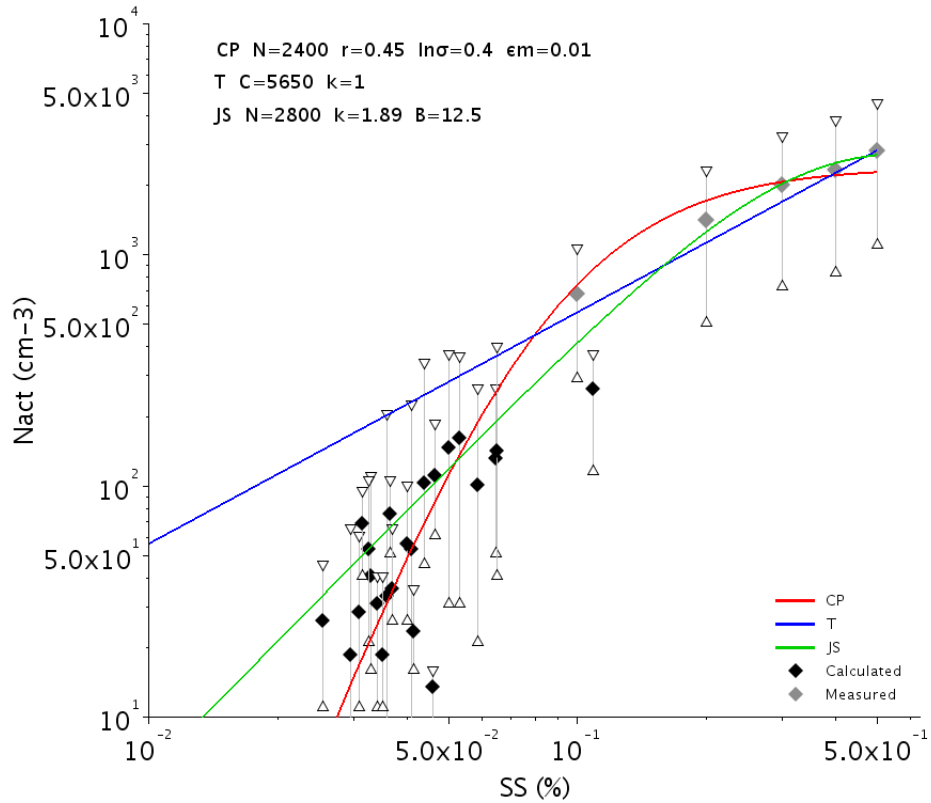
**Figure 6.** Number concentration of activated particles  $N_{act}$  as a function of a) the critical supersaturation, b) the dry critical diameter and c)  $\kappa$ , for the 23 fog cases. Diamonds correspond to mean values and error bars to uncertainty intervals. Same symbols as in Fig. 2.



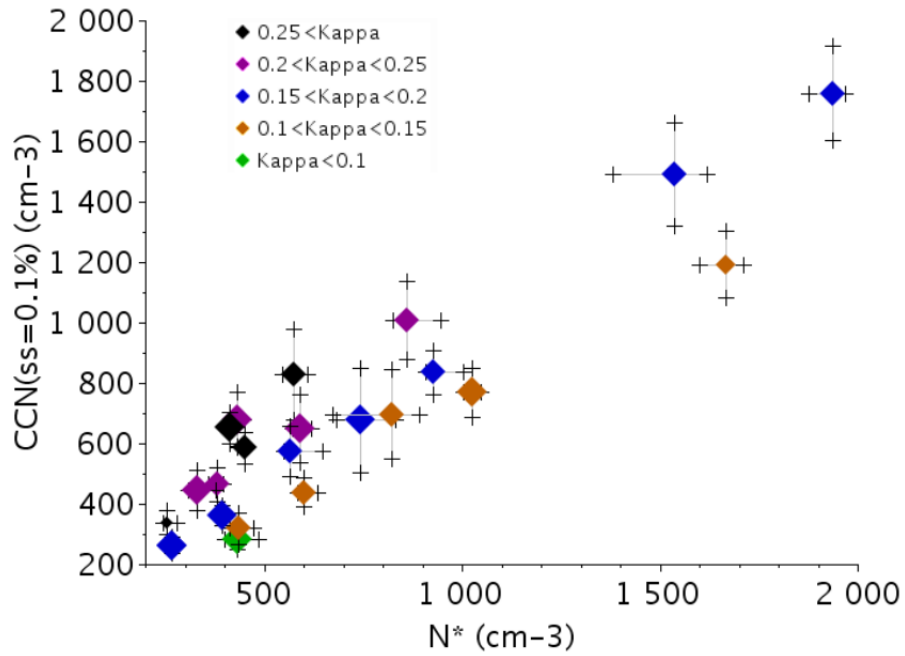
**Figure 7.** Ratio of activated particles to concentration of aerosol particles with diameter > 200 nm (activated fraction) for the 23 fog cases function of a) the critical supersaturation and b) the dry critical diameter. Diamonds correspond to mean values and error bars to uncertainty intervals. Color as a function of  $\kappa$  as indicated in the legend.



**Figure 8.** Number concentration of activated particles  $N_{act}$  as a function of  $N^*$  the concentration of aerosol particles with diameter  $> 200$  nm for the 23 fog events. For  $N^*$  diamonds correspond to median values and error bars to 25th-75th percentile intervals. Colour as a function of  $\kappa$  as indicated in the legend.

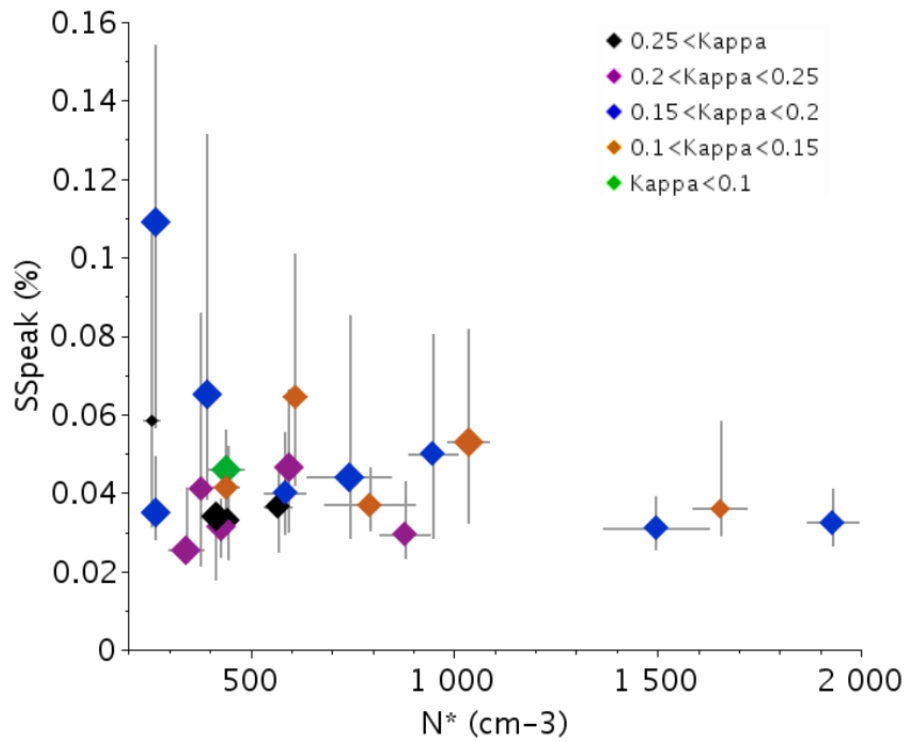


**Figure 9.** Number concentration of activated particles  $N_{act}$  as a function of the critical supersaturation for the 23 fog events (black) superimposed for the statistics of CCN chamber measurements (grey). For CCN measurements diamonds correspond to median values and error bars to 25th-75th percentile intervals. Colour lines correspond to fitted results of  $N_{CCN}=CS_{v,w}^k F(\mu, \frac{k}{2}, \frac{k}{2} + 1; -\beta S_{v,w}^2)$  (red),  $N_{CCN}=CS^k$  (blue) and  $N_{CCN}=N (1-\exp(-BS^k))$  (green), with parameters as indicated on the legend.

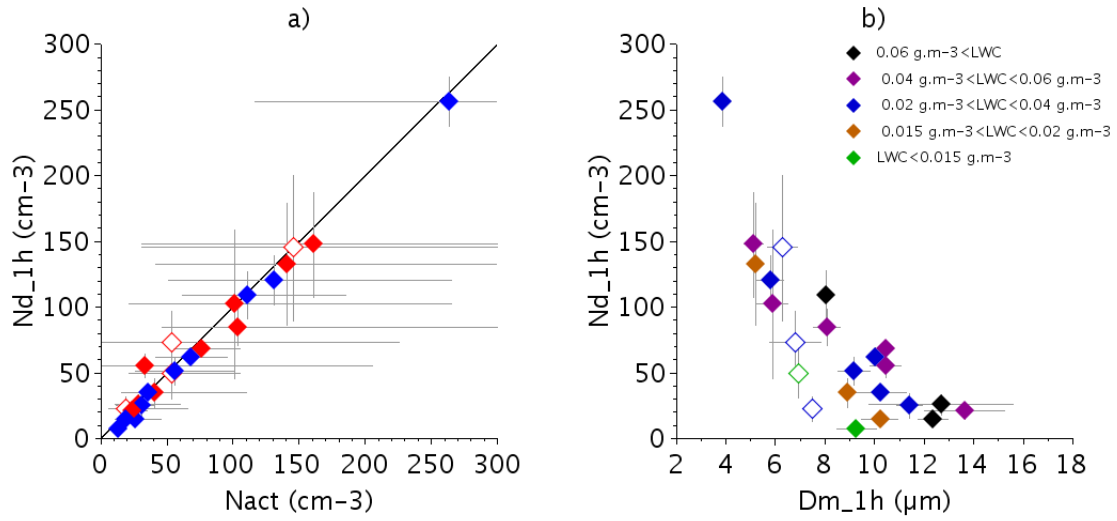


**Figure 10.** Median values for the 23 fog events of the number concentration measured by the CCN chamber at  $SS=0.1\%$  as a function of the concentration of aerosol particles with diameter  $> 200$  nm. Colour as a function of  $\kappa$  as in Fig. 8. Symbol size proportional to the mean diameter.

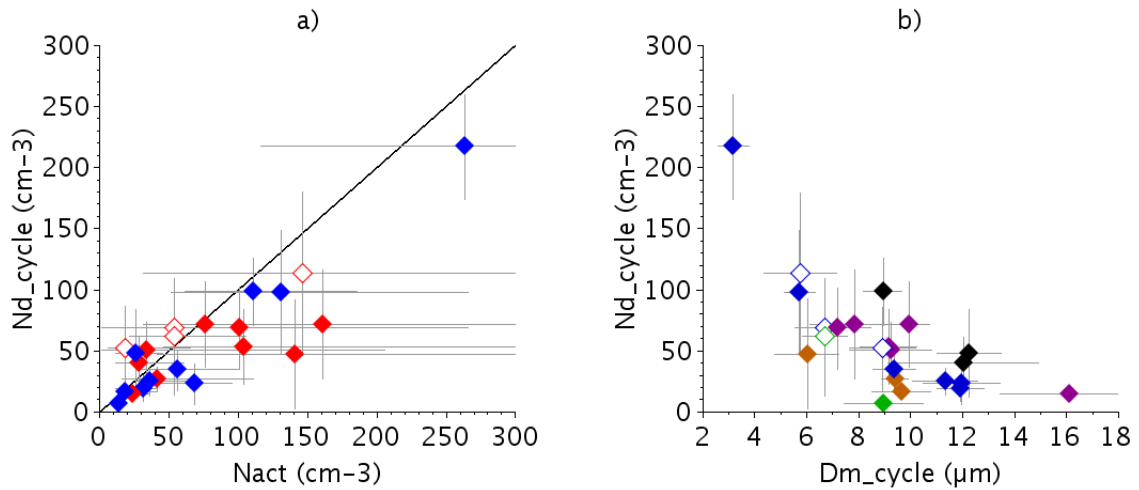




**Figure 11.** Critical supersaturation as a function of the concentration of aerosol particles with diameter > 200 nm for the 23 fog cases. Diamonds correspond to mean values and error bars to uncertainty intervals. Colour and size of symbols as in Fig. 10.



**Figure 12.** Droplet number concentration at the beginning of the fog for the 23 cases as a function of a) the activated particle concentration  $N_{act}$  and b) the mean fog diameter. Diamonds correspond to median values and error bars to 25th-75th percentile intervals. Symbol colors as a function of a) fog type as in Fig. 2 and b) LWC values as indicated on the legend.



**Figure 13.** Same as Fig. 12 except that statistics of  $Nd$ ,  $Dm$  and LWC are computed over the complete fog life cycle.

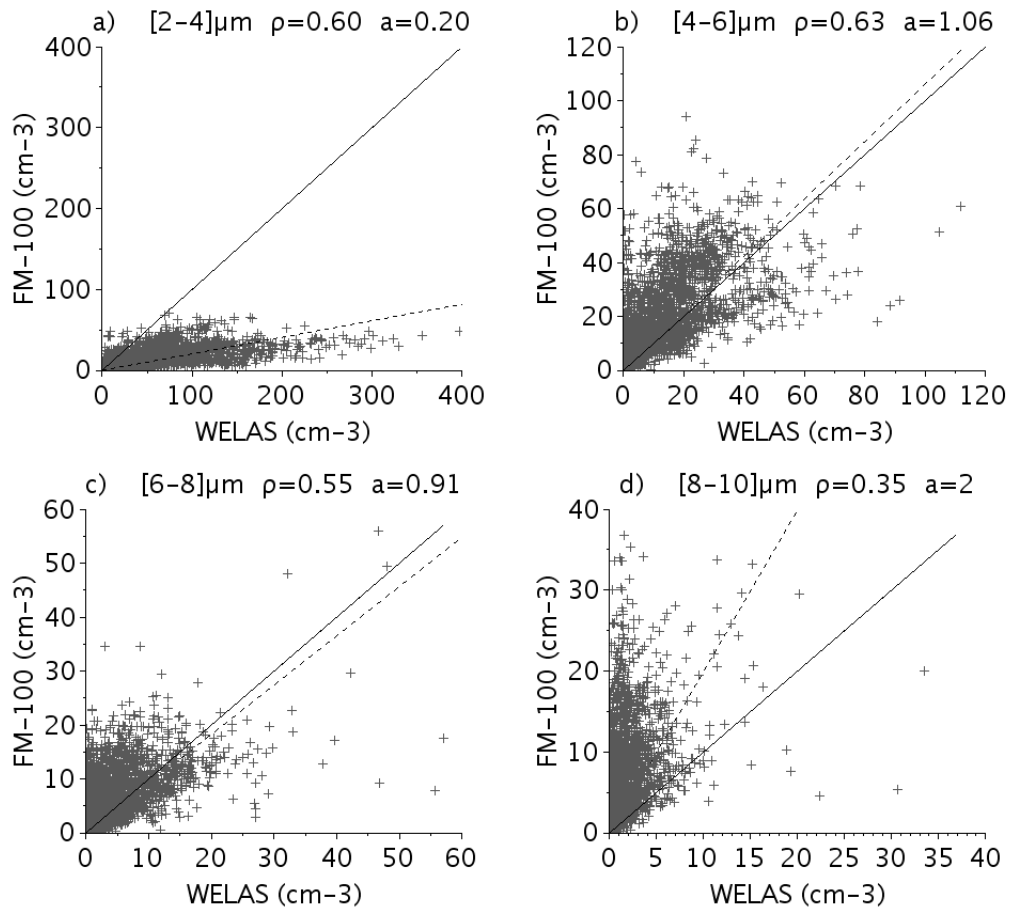
**Table 2.** List of fog events analysed here. Type RAD corresponds to radiation fog and STL to stratus lowering, "thick" to fog developed on the vertical, and "thin" to fog layer with top altitude lower than 18 m.  $N_a$  is the number concentration of aerosol particles derived from SMPS data. 25th, 50th and 75th percentiles are computed from the distribution of 5 minutes samples over the last hour before the fog beginning.  $N_{FM}$  is the droplet number concentration as measured by the FM-100 over the range [2-50]  $\mu\text{m}$ . Statistics of  $N_{FM}$  are computed with minute average data of samples with  $\text{LWC} > 0.005 \text{ g.m}^{-3}$  collected during the whole fog event. Ultimate column provides the minimum of visibility observed during the event.

Fog ID	Start time [UTC]	End time [UTC]	Type		$N_a$ ( $\text{cm}^{-3}$ )			$N_{FM}$ ( $\text{cm}^{-3}$ )			VISI <sub>min</sub> (m)
			(1)	(2)	25th	50th	75th	25th	50th	75th	
f1	16/11/10 2200	17/11/10 0540	RAD	thick	18868	20175	20625	52	67	80	53
f2	19/11/10 0540	19/11/10 1010	STL	thick	5877	6235	6783	30	51	73	108
f3	19/11/10 1540	19/11/10 1750	STL	thick	3540	4296	5213	5.3	6.3	7.4	389
f4	10/11/11 1800	11/11/11 1730	RAD	thick	10723	12602	13086	54	89	114	122
f5	15/11/11 0230	15/11/11 0940	RAD	thick	5703	6043	6285	45	60	75	235
f6	16/11/11 0110	16/11/11 1330	RAD	thick	7391	8188	9087	103	127	145	128
f7	16/11/11 1600	17/11/11 0010	STL	thick	7213	7648	8310	129	147	160	112
f8	18/11/11 0130	18/11/11 0410	RAD	thin	5245	5396	5447	71	85	112	167
f9	19/11/11 2200	20/11/11 0830	RAD	thin	13094	13587	14135	75	110	150	90
f10	21/11/11 2350	22/11/11 0810	RAD	thin	10160	10526	10737	87	123	156	87
f11	22/11/11 2050	22/11/11 2230	RAD	thin	15209	16290	19031	105	132	164	116
f12	23/11/11 0325	23/11/11 1005	RAD	thick	8920	9320	9723	87	115	157	115
f13	24/11/11 0620	24/11/11 1400	STL	thick	4770	5258	5586	18	33	55	193
f14	24/11/11 1610	24/11/11 1815	STL	thick	4180	4340	4566	19	33	48	267
f15	25/11/11 2140	26/11/11 1030	STL	thick	3328	3618	3762	45	68	121	106
f16	28/11/11 0630	28/11/11 1040	RAD	thick	9308	9618	11258	13	19	24	195
f17	16/11/12 2045	17/11/12 0920	STL	thick	9145	10051	11959	31	39	49	151
f18	20/11/12 0300	20/11/12 0910	RAD	thick	3108	3168	3410	16	26	44	177
f19	20/11/12 2015	20/11/12 2250	STL	thick	3198	3567	4282	39	62	84	139
f20	22/11/12 0315	22/11/12 0910	RAD	thick	4673	4780	4827	33	46	56	144
f21	30/11/12 1900	01/12/12 0245	RAD	thick	17791	18712	19151	42	63	87	102
f22	10/01/13 0215	10/01/13 0336	STL	thick	1975	2098	2273	21	36	64	233
f23	10/01/13 0500	10/01/13 0640	STL	thick	1904	1955	2009	24	29	33	351

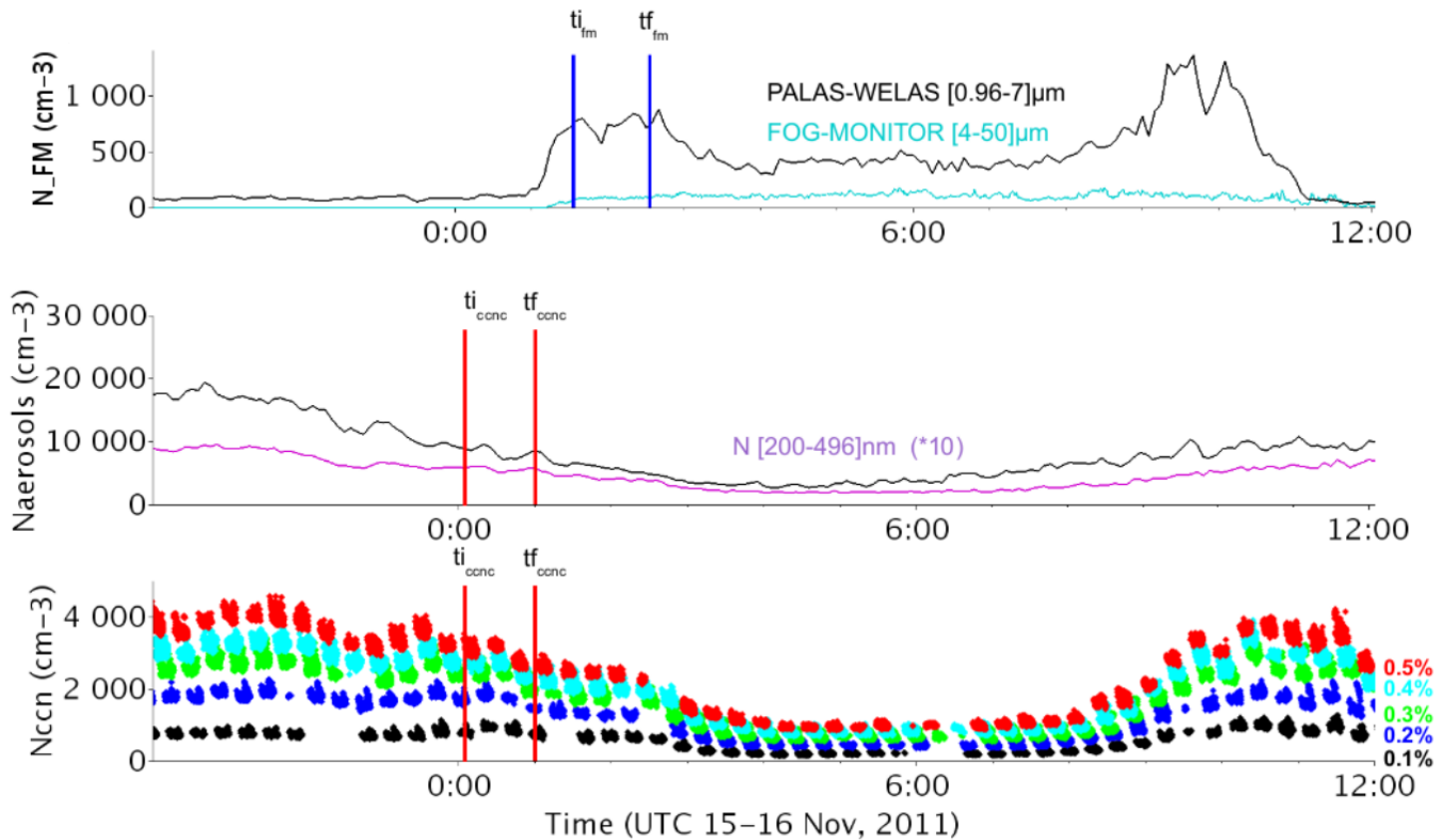
**Table 3.** Activation properties of the 23 fog events : hygroscopicity parameters  $\kappa_{in,f}$  and  $\kappa_{sup}$ , mean values and uncertainty intervals of the dry diameter  $D_d$ , the wet diameter  $D_w$ , the supersaturation  $SS_{peak}$  and the number concentration of activated particles  $N_{act}$ . 5th, 25th, 50th, 75th and 95th percentiles of the distribution of the 23 cases are indicated on the last lines for each parameter.

Fog ID	$\kappa_{in,f}$	$\kappa_{sup}$	$D_d$ $\mu\text{m}$	$D_{di}$ $\mu\text{m}$	$D_{ds}$ $\mu\text{m}$	$D_w$ $\mu\text{m}$	$D_{wi}$ $\mu\text{m}$	$D_{ws}$ $\mu\text{m}$	SS %	$SS_i$ %	$SS_s$ %	$N_{act}$ $\text{cm}^{-3}$	$N_{act,i}$ $\text{cm}^{-3}$	$N_{act,s}$ $\text{cm}^{-3}$
f1	0.11	0.16	0.46	0.35	0.50	4.24	2.56	5.14	0.04	0.03	0.06	33.5	1	206
f2	0.10	0.16	0.32	0.25	0.39	2.42	1.49	3.57	0.06	0.04	0.10	131	51	266
f3	0.06	0.12	0.45	0.44	0.49	3.31	2.66	4.25	0.05	0.04	0.06	13.5	6	16
f4	0.12	0.24	0.38	0.26	0.44	3.65	1.73	5.15	0.04	0.03	0.09	104	46	341
f5	0.20	0.35	0.4	0.32	0.45	4.83	2.89	6.56	0.03	0.02	0.05	41	16	111
f6	0.20	0.41	0.37	0.34	0.41	4.41	3.19	6.08	0.04	0.02	0.05	76	51	106
f7	0.14	0.28	0.35	0.31	0.41	3.46	2.27	5	0.05	0.03	0.07	111	61	186
f8	0.17	0.29	0.36	0.24	0.50	3.79	1.73	6.9	0.04	0.02	0.09	53.5	1	226
f9	0.14	0.25	0.47	0.39	0.50	5.22	3.43	6.37	0.03	0.02	0.04	18.5	6	66
f10	0.12	0.19	0.35	0.27	0.47	3.03	1.83	5.15	0.05	0.03	0.08	146	31	371
f11	0.13	0.20	0.43	0.42	0.49	4.67	3.59	5.63	0.03	0.03	0.04	53.5	21	106
f12	0.14	0.21	0.47	0.42	0.50	4.9	3.76	5.84	0.03	0.03	0.04	28.5	11	61
f13	0.17	0.22	0.34	0.24	0.38	3.25	1.73	4.04	0.05	0.04	0.09	106	71	281
f14	0.13	0.38	0.41	0.34	0.47	4.91	2.66	7.25	0.03	0.02	0.06	31	11	81
f15	0.19	0.31	0.43	0.37	0.47	5.15	3.54	6.6	0.03	0.02	0.04	21	11	46
f16	0.10	0.16	0.43	0.39	0.47	3.69	2.87	4.65	0.04	0.03	0.05	23.5	16	36
f17	0.12	0.18	0.45	0.41	0.47	4.15	3.23	4.97	0.04	0.03	0.05	36	26	66
f18	0.12	0.18	0.30	0.20	0.39	2.37	1.13	3.86	0.07	0.04	0.13	141	41	401
f19	0.13	0.22	0.39	0.34	0.44	3.82	2.66	5.02	0.04	0.03	0.06	56	26	101
f20	0.21	0.34	0.26	0.19	0.37	2.62	1.4	4.75	0.06	0.03	0.11	101	21	266
f21	0.10	0.16	0.37	0.29	0.47	3.03	1.85	4.67	0.05	0.03	0.08	161	31	366
f22	0.14	0.21	0.21	0.17	0.29	1.44	0.98	2.65	0.11	0.06	0.15	264	116	371
f23	0.14	0.21	0.43	0.37	0.47	4.27	3.03	5.33	0.04	0.03	0.05	18.5	11	41
5th	0.12	0.18	0.35	0.25	0.41	3.03	1.73	4.65	0.03	0.02	0.05	18.5	11	41
25th	0.13	0.21	0.39	0.34	0.47	3.79	2.66	5.14	0.04	0.03	0.05	28.5	11	66
50th	0.14	0.21	0.43	0.37	0.47	4.27	3.03	5.33	0.04	0.03	0.06	53.5	21	111
75th	0.17	0.29	0.43	0.39	0.49	4.67	3.19	6.08	0.05	0.03	0.09	111	46	281
95th	0.21	0.41	0.47	0.44	0.50	5.22	3.76	7.25	0.11	0.06	0.15	264	116	401

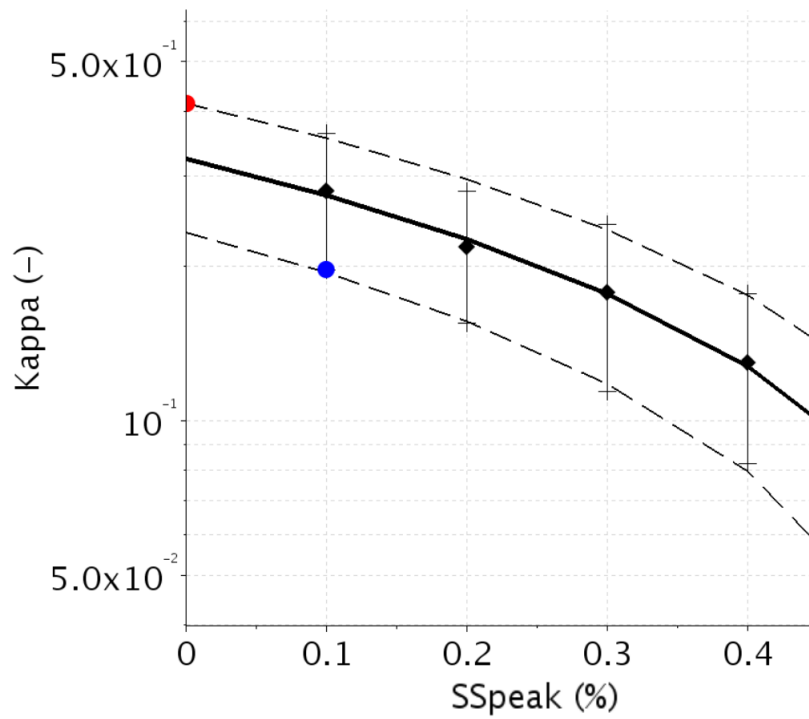
## Appendix A: Supplementary material



**Figure A.1.** Scatterplot of the five minutes average particle number concentration values as measured by the FM-100 for the four first bins vs the integrated WELAS measurement over the corresponding diameter range a) [2-4]  $\mu\text{m}$ , b) [4-6]  $\mu\text{m}$ , c) [6-8]  $\mu\text{m}$  and d) [8-10]  $\mu\text{m}$ . Solid line corresponds to 1:1 line and dashed line corresponds to best fit line with correlation coefficient  $\rho$  and slope  $a$ , as indicated in the legend.



**Figure A.2.** Time series of measurements for the f6 case : (a) particle number concentration at ambient humidity from FM-100 (cyan) and WELAS (black); (b) dry particle number concentration from SMPS : total (black) and particle with diameter > 200 nm (purple); and (c)  $N_{ccn}$  from the CCN chamber at each supersaturation with color as indicated on the label. Vertical segments indicate the selected time periods before (red) and during (blue) the beginning of the fog event.



**Figure A.3.** Derived  $\kappa$  values as a function of CCN chamber supersaturations from measurements over the 1 hour time period before the fog event for the f6 case. Mean (diamond, solid line) and mean  $\pm$  one standard deviation (dashed lines). Extrapolated  $\kappa_{inf}$  (blue dot) and  $\kappa_{sup}$  (red dot) are also indicated.