

29

30 **Abstract:**

31 One of the challenges in representing warm rain processes in global climate models
32 (GCM) is related to the representation of the subgrid variability of cloud properties, such as cloud
33 water and cloud droplet number concentration (CDNC), and the effect thereof on individual
34 precipitation processes such as autoconversion. This effect is conventionally treated by
35 multiplying the resolved-scale warm rain process rates by an enhancement factor (E_q) which is
36 derived from integrating over an assumed subgrid cloud water distribution. The assumed subgrid
37 cloud distribution remain highly uncertain. In this study, we derive the subgrid variations of
38 liquid-phase cloud properties over the tropical ocean using the satellite remote sensing products
39 from Moderate Resolution Imaging Spectroradiometer (MODIS) and investigate the
40 corresponding enhancement factors for the GCM parameterization of autoconversion rate. We
41 find that the conventional approach of using only subgrid variability of cloud water is insufficient,
42 and that the subgrid variability of CDNC, as well as the correlation between the two, are also
43 important for the correctly simulating the autoconversion process in GCMs. Using the MODIS
44 data which has the near-global data coverage, we find that E_q shows a strong dependence on
45 cloud regimes, due to the fact that the subgrid variability of cloud water and CDNC is regime-
46 dependent. Our analysis shows a significant increase of E_q from the stratocumulus (Sc) to
47 cumulus (Cu) regions. Furthermore, the enhancement factor E_N due to the subgrid variation of
48 CDNC is derived from satellite observation for the first time, and results reveal several regions
49 downwind of biomass burning aerosols (e.g., Gulf of Guinea, East Coast of South Africa), air
50 pollution (i.e., Eastern China Sea), and active volcanos (e.g., Kilauea Hawaii and Ambae Vanuatu),
51 where the E_N is comparable, or even larger than E_q , suggesting an important role of aerosol in
52 influencing the E_N . MODIS observations suggest that the subgrid variations of cloud liquid water
53 path (LWP) and CDNC are generally positively correlated. As a result, the combined enhancement
54 factor, including the effect of LWP and CDNC correlation, is significantly smaller than the simple
55 product of $E_q \cdot E_N$. Given the importance of warm rain processes in understanding the Earth
56 system dynamics and water cycle, we conclude that more observational studies are needed to
57 provide a better constraint on the warm rain processes in GCMs.

58 1. Introduction

59 Marine boundary layer (MBL) clouds are a strong modulator of Earth’s radiative energy
60 budget (Klein and Hartmann, 1993; Trenberth et al., 2009). They can interact with other
61 components of the climate system, such as aerosols and precipitations, in various ways. The
62 feedback of MBL clouds to climate change remains one of the largest uncertainties in our
63 understanding of the climate sensitivity (Bony and Dufresne, 2005; Soden and Held, 2006).
64 Despite their importance in the climate system, simulating MBL clouds in general circulations
65 models (GCM) has proved to be extremely challenging. A main difficulty is rooted in the fact the
66 typical grid size of GCM (~100km) is much larger than the spatial scale of many cloud
67 microphysical processes, and as a result these subgrid scale processes, as well as the subgrid
68 cloud variations, have to be highly simplified and then parameterized as functions of resolved,
69 grid-level variables.

70 Of particular interest in this study is the warm rain processes in MBL clouds, which have
71 fundamental impacts on the cloud water budget and lifetime. Although in reality it is highly
72 complicated and involves multiple factors, warm rain formation in GCMs is usually parameterized
73 as simple functions of only key cloud parameters. For example, the drizzle in MBL cloud is
74 initialized by the so-called autoconversion process in which the collision-coalescence of cloud
75 droplets gives birth to large drizzle drops (Pruppacher and Klett, 1997). In GCMs, for the sake of
76 efficiency, this process is usually parameterized as a power function of liquid water content (LWC
77 or symbol q_c) and cloud droplet number concentration (CDNC or symbol N_c). One of the most
78 widely used parameterization scheme is developed by Khairoutdinov and Kogan (2000)
79 (“KK2000” hereafter), which has the form

$$\frac{\partial q_r}{\partial t} = C(q_c)^{\beta_q}(N_c)^{\beta_N}, \quad (1)$$

80 where $\frac{\partial q_r}{\partial t}$ is the rain water tendency due to the autoconversion process, q_c has the unit of kg/kg,
81 and N_c of cm^{-3} . The three parameters $C = 1350$, $\beta_q = 2.47$ and $\beta_N = -1.79$ are derived
82 through a simple least-square fitting of the autoconversion rate results from a large-eddy
83 simulation with bin microphysics that can simulate the process-level physics. Even though this is
84 highly simplified, the parametrization scheme still faces a great challenge. The calculation of grid-

85 mean autoconversion efficiency requires the knowledge of subgrid distributions of LWC and
86 CDNC, but in the GCMs only grid-mean quantities $\langle q_c \rangle$ and $\langle N_c \rangle$ are known and available for use
87 in the computation of autoconversion rate. As pointed out by Pincus and Klein (2000), for a
88 process $f(x)$ such as autoconversion that is nonlinearly dependent on subgrid variables, x , the
89 grid-mean value $\langle f(x) \rangle$ is not equal to the value estimated based on the grid-mean $\langle x \rangle$, i.e.,
90 $\langle f(x) \rangle \neq f(\langle x \rangle)$. Mathematically, if $f(x)$ is convex, then $f(\langle x \rangle) < \langle f(x) \rangle$ (Larson and Griffin,
91 2013; Larson et al., 2001). To take this effect into account, a parameter E is often introduced in
92 the GCM as part of the parameterization such that $\langle f(x) \rangle = E \cdot f(\langle x \rangle)$. It is referred to as the
93 “enhancement factor” in many studies and this study too because $E > 1$ for a convex function.
94 Such a nonlinear effect is not just limited to the autoconversion process. Some other examples
95 are the plane-parallel albedo bias (Barker, 1996; Cahalan et al., 1994; Oreopoulos and Davies,
96 1998a), subgrid cloud droplet activation (Morales and Nenes, 2010) and accretion (Boutle et al.,
97 2014; Lebsock et al., 2013).

98 The value of E is determined primarily by two factors: the nonlinearity of $f(x)$ and the
99 subgrid probability density function (PDF) $P(x)$. Given the same subgrid variation of LWC, i.e.,
100 $P(q_c)$, the nonlinear effect impacts the autoconversion parameterization more than it does on
101 the accretion, because the former is a more nonlinear function of q_c than the latter. For the same
102 $f(x)$, a grid box with a narrow and symmetric $P(x)$ would require a smaller E than another grid
103 box with a broader and non-symmetric $P(x)$. Ideally, the value of the enhancement factor E
104 should be diagnosed from the subgrid cloud PDF $P(x)$. Unfortunately, because this is not possible
105 in most conventional GCMs, the value of E is usually assumed to be a constant for the lack of
106 better options. The E for autoconversion due to subgrid LWC variation is assumed to be 3.2 in
107 the two-moment cloud microphysics parameterization schemes by Morrison and Gettelman
108 (2008) (MG scheme hereafter), which is employed in the widely used Community Atmosphere
109 Model (CAM). This choice of $E = 3.2$ is based on an early study by Barker et al. (1996), in which
110 the mesoscale variation of column-integrated optical thickness of the “overcast stratocumulus”,
111 “broken stratocumulus” and “scattered stratocumulus” are studied. The value $E = 3.2$ is derived
112 based on the mesoscale variation of the broken stratocumulus.

113 Clearly, a simple constant E is not adequate. The following is a list of attempts to better
114 understand the subgrid cloud variations and the implications for warm rain simulations in GCMs.
115 Several previous studies have shown that the mesoscale cloud water variation is a strong function
116 of cloud regime—the subgrid cloud water variation of Sc cloud is much different from that of Cu
117 clouds (Barker et al., 1996; Lee et al., 2010; Oreopoulos and Cahalan, 2005; Wood and Hartmann,
118 2006). As the first part of a two-part study, Larson and Griffin (2013) first laid out a systematic
119 theoretical basis for understanding the effects of subgrid cloud property variations on simulating
120 various nonlinear processes in GCM, including not only the autoconversion but also the accretion,
121 condensation, evaporation and sedimentation processes. In the second part, using cloud fields
122 from a large-eddy simulation (LES), Griffin and Larson (2013) showed that inclusion of the
123 enhancement factor indeed leads to more rainwater at surface in single-column simulations and
124 makes them agree better with high-resolution large-eddy simulations. Recently, using the
125 ground-based observations from three Department of Energy (DOE) Atmospheric Radiation
126 Measurement (ARM) sites, Xie and Zhang (2015) developed a scale-aware parameterization
127 scheme for GCMs to account for subgrid cloud water variation. Also using ARM measurement,
128 Ahlgrim and Forbes (2016) analyzed the dependence of cloud water variability on cloud regime.
129 Although these previous studies have shed important light on subgrid cloud variation and the
130 implications for GCM, they lack a global perspective because they are only based on limited data
131 (e.g., LES cases, in situ and ground-based measurement). Currently, satellite remote sensing
132 observation is the only way to achieve a global perspective. Using the observations from the
133 space-borne radar CloudSat, Lebsock et al. (2013) showed that the subgrid cloud water variance
134 is smaller over the Sc region than over the Cu region, and as a result the enhancement factor
135 shows an increasing trend from Sc to Cu region. They also highlighted importance of considering
136 the subgrid co-variability of cloud water and rain water in the computation of the accretion rate.
137 Using a combination of in situ measurement and satellite remote sensing data, Boutle et al.
138 (2014) analyzed the spatial variation of both cloud and rain water, as well as their covariation,
139 and developed a simple parameterization scheme to relate the subgrid cloud water variance to
140 the grid-mean cloud fraction. Later, the study of Boutle et al. (2014) was extended by Hill et al.
141 (2015) who developed a cloud regime dependent and scale-aware parameterization scheme for

142 simulating subgrid cloud water variation. On the modeling side, Guo et al. (2014) investigated the
143 sensitivity of cloud simulation in the Geophysical Fluid Dynamics Laboratory (GFDL) Atmospheric
144 General Circulation Model (AM) to the subgrid cloud water parameterization schemes. A similar
145 study was carried out by Bogenschutz et al. (2013) using the National Center of Atmospheric
146 Research (NCAR) Community Atmospheric Model (CAM). Both studies show that the more
147 sophisticated subgrid parameterization scheme— Cloud Layers Unified by Binormals (CLUBB)
148 (Golaz et al., 2002a; 2002b; Larson et al., 2002)—lead to a better simulation of clouds in the
149 model. However, a more recent study by Song et al. (2018b) reveals that the CLUBB in CAM
150 version 5.3 (CAM5.3) overestimates the enhancement factor in the trade wind cumulus cloud
151 region, which in turn leads to excessive drizzle in the model and “empty clouds” with near-zero
152 cloud water. In addition to CLUBB, the so-called super-parameterization (a.k.a Multiscale
153 Modeling Framework (MMF)), which uses cloud resolving model embedded in the GCM grids to
154 diagnose sub-grid cloud variations (Randall et al., 2003), have also gained increasing popularity.
155 Takahashi et al. (2017) compared the subgrid cloud water variations simulated by a CAM-MMF
156 model with those derived from A-Train observations and found reasonable agreement.

157 Despite these previous studies, many questions remain unanswered. First of all, all the
158 previous studies, as far as we know, have focused on the impact of subgrid cloud water
159 q_c variation. The potential impact of subgrid variation of N_c and the co-variability of N_c with q_c
160 have been overlooked so far. Given the same amount of q_c , a cloud with a smaller N_c would have
161 larger droplets and therefore larger precipitation efficiency than another cloud with a larger N_c .
162 For the same reason, other things equal, a grid with positive correlation of subgrid N_c and q_c
163 would be less efficient in terms of autoconversion than a grid with negative correlation of the
164 two. Secondly, most of previous studies are based on the assumption that the subgrid cloud
165 property variation follows certain well-behaved distributions, usually either Gamma (e.g., Barker,
166 1996; Morrison and Gettelman, 2008; Oreopoulos and Barker, 1999; Oreopoulos and Cahalan,
167 2005) or Lognormal (Boutle et al., 2014; Larson and Griffin, 2013; e.g., Lebsock et al., 2013).
168 However, the validity and performance of the assumed PDF shape are seldom checked.
169 Furthermore, although the study by Lebsock et al. (2013) has depicted a global picture of the

170 enhancement factor for the autoconversion modeling in GCM, the picture is far from clear due
171 to the small sampling rate of CloudSat observations.

172 In this study, we revisit the subgrid variations of liquid-phase cloud properties over the
173 tropical ocean using 10 years of MODIS cloud observations, with the overarching goal to better
174 understand the potential impacts of subgrid cloud variations on the warm rain processes in the
175 conventional GCMs. Similar to previous studies, we will quantify the subgrid cloud water
176 variations based on MODIS observations. Going one step further, we will also attempt to unveil
177 for the first time the subgrid N_c variation, as well as its correlation with cloud water, and
178 investigate the implications for warm rain simulations in GCM. Moreover, we will take advantage
179 of the wide spatial coverage of MODIS data to achieve a more detailed picture of the
180 enhancement factor for the autoconversion simulation. Last but not least, we will evaluate the
181 two widely used distributions, i.e., Lognormal and Gamma, in terms of their performance and
182 limitations for simulating the enhancement factor. We will first explain the theoretical
183 background in Section 2 and introduce the data and methodology in Section 3. The MODIS
184 observations will be presented and discussed in Section 4. The implications for the
185 autoconversion parameterization in the GCMs will be discussed in 5. The main findings will be
186 summarized in Section 6 with an outlook for future studies.

187 **2. Theoretical Background**

188 **2.1. Theoretical Distributions to describe subgrid cloud property variations**

189 In previous studies, the spatial variations of cloud properties, such as cloud optical thickness
190 (COT), cloud liquid water path (LWP) and cloud liquid water content (LWC), are often described
191 using either of two theoretical distributions—the Gamma and Lognormal distribution. The
192 probability density function (PDF) from a Gamma distribution is a two-parameter function as
193 follows (Barker, 1996; Oreopoulos and Davies, 1998b):

$$P_G(x) = \frac{1}{\Gamma(\nu)} \alpha^\nu x^{\nu-1} \exp(-\alpha x), \quad (2)$$

194 where Γ is the Gamma function, ν is the so-called inverse relative variance, and α the so-called
195 rate parameter. If x follows the Gamma distribution, its mean value is given by

$$\langle x \rangle = \int_0^{\infty} x P_G(x) dx = \frac{v}{\alpha}, \quad (3)$$

196 and variance given by

$$Var(x) = \int_0^{\infty} (x - \langle x \rangle)^2 P_G(x) dx = \frac{v}{\alpha^2}. \quad (4)$$

197 It follows from Eq. (3) and (4) that the so-called inverse relative variance is

$$v = \frac{1}{\eta} = \frac{\langle x \rangle^2}{Var(x)}, \quad (5)$$

198 where $\eta = \frac{Var(x)}{\langle x \rangle^2}$ is the relative variance. If x follows the Gamma distribution, for a physical

199 process $M(x)$ that is a power function of x ,

$$M(x) = Kx^\beta, \quad (6)$$

200 then the expected value $\langle M(x) \rangle$ is given by

$$\langle M(x) \rangle_G = K \int_0^{\infty} x^\beta P_G(x) dx = \frac{\Gamma(v+\beta)}{\Gamma(v)v^\beta} K \langle x \rangle^\beta, \quad \beta > -v. \quad (7)$$

201 As explained in the introduction, for a nonlinear process $M(x)$, $\langle M(x) \rangle \neq M(\langle x \rangle)$. The ratio
202 between the two E is by definition the enhancement factor:

$$E(P_G, v, \beta) = \frac{\langle Kx^\beta \rangle}{K \langle x \rangle^\beta} = \frac{1}{\langle x \rangle^\beta} \int_0^{\infty} x^\beta P_G(x) dx = \frac{\Gamma(v+\beta)}{\Gamma(v)v^\beta}, \quad (8)$$

203 The PDF of a Lognormal distribution is given as follows (Larson and Griffin, 2013;
204 Lebsack et al., 2013):

$$P_L(x) = \frac{1}{\sqrt{2\pi}x\sigma} \exp\left(-\frac{(\ln x - \mu)^2}{2\sigma^2}\right), \quad (9)$$

205 where $\mu = \langle \ln x \rangle$ and $\sigma^2 = Var(\ln x)$ correspond to the mean and variance of $\ln x$, respectively.

206 The mean value of the Lognormal distribution is given by

$$\langle x \rangle = \int_0^{\infty} x P_L(x) dx = e^{\mu + \frac{\sigma^2}{2}}, \quad (10)$$

207 and the variance by

$$Var(x) = \int_0^{\infty} (x - \langle x \rangle)^2 P_L(x) dx = e^{2\mu + \sigma^2} (e^{\sigma^2} - 1). \quad (11)$$

208 It follows from Eq. (10) and (11) that the inverse relative variance can be derived from the
209 following equation

$$e^{\sigma^2} = 1 + \frac{Var(x)}{\langle x \rangle^2} = 1 + \frac{1}{v}. \quad (12)$$

210 If x follows the Lognormal distribution, the expected value of $\langle M(x) \rangle$ is

$$\langle M(x) \rangle_L = K \int_0^\infty x^\beta P_L(x) dx = \left(1 + \frac{1}{v}\right)^{\frac{\beta^2 - \beta}{2}} K \langle x \rangle^\beta. \quad (13)$$

211 Evidently, the corresponding enhancement factor is given by

$$E(P_L, v, \beta) = \frac{\langle Kx^\beta \rangle}{K \langle x \rangle^\beta} = \left(1 + \frac{1}{v}\right)^{\frac{\beta^2 - \beta}{2}}. \quad (14)$$

212 Note that Eq. (7) and (8) are only valid when $\beta > -v$ because Gamma function $\Gamma(v + \beta)$ can
 213 run into singular values when $v + \beta < 0$. In contrast, Eq. (13) and (14) are valid for any real value
 214 β . This is one advantage of the Lognormal distribution over the Gamma distribution.

215 An example of the Gamma and Lognormal distributions for q_c is shown in Figure 1a. In
 216 this example, both distributions have the same mean $\langle q_c \rangle = 0.5g/kg$ and also the same inverse
 217 relative variance $v_q = 3$. Although the general shapes of the two PDFs are similar, they differ
 218 significantly at the two ends: the Gamma PDF is larger than Lognormal PDF over the small values
 219 of q_c , and the opposite is true over the large values of q_c . The Gamma and Lognormal
 220 distributions can also be used to describe the spatial variation of N_c (Gultepe and Isaac, 2004).
 221 An example is given in Figure 1c, in which q_c is a constant of $0.5g/kg$, $\langle N_c \rangle = 50 cm^{-3}$, and $v_N =$
 222 5.0 . Figure 1 b shows the autoconversion rate based on the KK2000 parameterization scheme for
 223 the Gamma $P_G(q_c)$ and Lognormal $P_L(q_c)$ that are shown in Figure 1a. Interestingly, although
 224 the cumulative autoconversion rates based on the two types of PDFs are almost identical, the
 225 contribution to the total autoconversion rate from the different LWC bins are quite different. As
 226 show in Figure 1a, the $P_L(q_c)$ has a longer tail than the $P_G(q_c)$, i.e., the occurrence probability of
 227 large q_c (e.g., $q_c > 2.0g/kg$) is much higher in the Lognormal than in Gamma PDF. This
 228 difference is further amplified in the autoconversion rate computation in Figure 1b because the
 229 autoconversion rate is proportional to $q_c^{2.47}$.

230 The enhancement factors based on the Gamma (i.e., $E(P_G, \beta)$ in Eq. (8)) and Lognormal
 231 (i.e., $E(P_L, \beta)$ in Eq. (14)) PDF for $\beta_q = 2.47$ are plotted as a function of the inverse relative
 232 variance v in Figure 2. When subgrid clouds are more homogenous i.e., $v > 1$, the enhancement
 233 factor based on the two PDFs are similar. However, for more inhomogeneous grids with i.e., $v <$

234 1, the $E(P_L, \beta)$ is significantly larger than that $E(P_G, \beta)$, which is probably because of the longer
 235 tail of $P_L(q_c)$ as shown in Figure 1 a and b.

236

237 **2.2. Impacts of subgrid cloud variations on warm rain parameterization in GCM**

238 The warm rain process in MBL clouds involves many interacting microphysical processes. In
 239 this study, we only focus only on the simulation of autoconversion in GCM. Other nonlinear
 240 processes, such as accretion and evaporation have been investigated in previous studies (Boutle
 241 et al., 2014; Lebsock et al., 2013).

242 Ideally, if the subgrid variations of q_c and N_c are known, then the grid-mean in-cloud
 243 autoconversion rate should be derived from the following integral

$$\langle \frac{\partial q_r}{\partial t} \rangle = \int_0^\infty \int_0^\infty C(q_c)^{\beta_q} (N_c)^{\beta_N} P(q_c, N_c) dq_c dN_c, \quad (15)$$

244 where $P(q_c, N_c)$ is the joint PDF of q_c and N_c . Unfortunately, most conventional GCMs lack the
 245 capability of predicting the subgrid variations of cloud properties, with only a couple of
 246 exceptions (Thayer-Calder et al., 2015). What is known from the GCM is usually the in-cloud grid-
 247 mean values $\langle q_c \rangle$ and $\langle N_c \rangle$. As a result, instead of using Eq. (15), the autoconversion rate in GCMs
 248 is usually computed from the following equation

$$\langle \frac{\partial q_r}{\partial t} \rangle = E \cdot C(\langle q_c \rangle)^{\beta_q} (\langle N_c \rangle)^{\beta_N}, \quad (16)$$

249 where E is the enhancement factor defined as:

$$E = \frac{\int_0^\infty \int_0^\infty (q_c)^{\beta_q} (N_c)^{\beta_N} P(q_c, N_c) dq_c dN_c}{(\langle q_c \rangle)^{\beta_q} (\langle N_c \rangle)^{\beta_N}}. \quad (17)$$

250 The value of the enhancement factor depends on the subgrid variations of q_c and N_c . If clouds
 251 are homogenous on the subgrid scale, then $E \sim 1$. The more inhomogeneous the clouds are, the
 252 larger the E is. In the special case where q_c and N_c are independent, then the joint PDF $P(q_c, N_c)$
 253 becomes $P(q_c, N_c) = P(q_c)P(N_c)$, where $P(q_c)$ and $P(N_c)$ are the PDF of the subgrid q_c and
 254 N_c . Consequently, Eq. (15) reduces to

$$\langle \frac{\partial q_r}{\partial t} \rangle = C \int_0^\infty (q_c)^{\beta_q} P(q_c) dq_c \int_0^\infty (N_c)^{\beta_N} P(N_c) dN_c, \quad (18)$$

255 and Eq.(17) to

$$E = E_q \cdot E_N, \quad (19)$$

256 where E_q is the enhancement factor due to the subgrid variation of cloud water which has the
257 form,

$$E_q = \frac{\int_0^\infty (q_c)^{\beta_q} P(q_c) dq_c}{(\langle q_c \rangle)^{\beta_q}}, \quad (20)$$

258 and the E_N is the enhancement factor due to the subgrid variation of CDNC which has the form,

$$E_N = \frac{\int_0^\infty (N_c)^{\beta_N} P(N_c) dN_c}{(\langle N_c \rangle)^{\beta_N}}. \quad (21)$$

259 Obviously, if $P(q_c)$ and $P(N_c)$ follow either Gamma or Lognormal distribution, then the above
260 equations reduce to Eq. (8) or (14), respectively.

261 If q_c and N_c both have significant subgrid variations and they are not independent, the
262 enhancement factor should ideally be diagnosed from Eq. (17). However, the joint PDF $P(q_c, N_c)$
263 may not be known and the integration can be time-consuming. Some previous studies proposed
264 to approximate the $P(q_c, N_c)$ as a bivariate lognormal distribution as follows:

$$P(q_c, N_c) = \frac{1}{2\pi q_c N_c \sigma_q \sigma_N \sqrt{1 - \rho^2}} \exp\left(-\frac{\zeta}{2}\right) \quad (22)$$

$$\zeta = \frac{1}{1 - \rho^2} \left[\left(\frac{\ln q_c - \mu_q}{\sigma_q} \right)^2 - 2\rho \left(\frac{\ln q_c - \mu_q}{\sigma_q} \right) \left(\frac{\ln N_c - \mu_N}{\sigma_N} \right) + \left(\frac{\ln N_c - \mu_N}{\sigma_N} \right)^2 \right],$$

265 where ρ is the correlation coefficient between q_c and N_c (Larson and Griffin, 2013; Lebsock et
266 al., 2013). As such, both q_c and N_c follow a marginal lognormal distribution in Eq. (9). Substituting
267 Eq. (22) into Eq. (17), we obtain the enhancement factor for the bivariate lognormal distribution
268 that consists of three terms

$$E = E_q(P_L, v_q, \beta_q) \cdot E_N(P_L, v_N, \beta_N) \cdot E_{COV}(\rho, \beta_q, \beta_N, v_q, v_N), \quad (23)$$

269 where $E_q(P_L, v_q, \beta_q) = \left(1 + \frac{1}{v_q}\right)^{\frac{\beta_q^2 - \beta_q}{2}}$ and $E_N(P_L, v_N, \beta_N) = \left(1 + \frac{1}{v_N}\right)^{\frac{\beta_N^2 - \beta_N}{2}}$ correspond to the
270 impacts of subgrid q_c and N_c variance, respectively (i.e., Eq. (14)), and the third term

$$E_{COV}(\rho, \beta_q, \beta_N, v_q, v_N) = \exp(\rho \beta_q \beta_N \sigma_q \sigma_N), \quad (24)$$

271 corresponds to the impact of the co-variation of q_c and N_c on the enhancement factor. Obviously,
 272 Eq. (23) reduces to Eq. (19) when q_c and N_c are uncorrelated (i.e., $\rho = 0, E_{COV} = 1$). If q_c and N_c
 273 are negatively correlated (i.e., $\rho < 0$ and $E_{COV} > 1$), clouds with larger q_c would tend to have
 274 smaller N_c . The autoconversion rate in such a case would be larger than that in the case where
 275 q_c and N_c are positively correlated (i.e., i.e., $\rho > 0$ and $E_{COV} < 1$). A positive correlation would
 276 exist, for instance, if all droplets in cloud were the same size, but some parcels had more droplets
 277 than other parcels.

278 Most current GCMs do not have the capability to simulate the subgrid cloud property
 279 variations. They usually have to use pre-defined subgrid cloud variations in the computation of
 280 grid-mean autoconversion rate instead of using prognostic values. For example, in the MG
 281 scheme for the CAM5.3, the subgrid q_c is assumed to follow the Gamma distribution in Eq. (2)
 282 with a fixed $v_q = 1$ and as a result constant $E_q = 3.2$. Lately, advanced subgrid parameterization
 283 schemes, such as CLUBB, have been implemented in several GCMs, including CAM6 and GFDL
 284 AM model (Bogenschutz et al., 2017; Guo et al., 2015; 2014), which provides information on the
 285 subgrid q_c variation to the host model. The information can then be used to dynamically diagnose
 286 the enhancement factor E_q , which will help the model simulate the cloud regime dependence of
 287 E_q (Guo et al., 2010; 2014).

288 However, as explained above, not only the subgrid variation of q_c but the subgrid
 289 variation of N_c can also influence the enhancement factor. Unfortunately, this aspect has been
 290 ignored by almost all GCMs, even the latest CAM6 with CLUBB. Physically, provided the same q_c ,
 291 a cloud with smaller N_c would have larger droplet size and therefore larger precipitation
 292 efficiency than the cloud with larger N_c . Because the autoconversion rate depends nonlinearly
 293 on N_c , the grid-mean autoconversion rate computed based on a skewed PDF of N_c (i.e.,
 294 $\int_0^\infty (N_c)^{\beta_N} P(N_c) dN_c$) would be different from that computed based on the mean of N_c (i.e.,
 295 $(\langle N_c \rangle)^{\beta_N}$). The autoconversion enhancement factor based on the Lognormal PDF $E(P_L, \beta)$ for
 296 $\beta_N = -1.79$ is given in Figure 2. Interestingly, at the same inverse relative variance v , the
 297 enhancement factor based on the same Lognormal PDF $E(P_L, \beta)$ for $\beta_N = -1.79$ is actually
 298 larger than that for $\beta_q = 2.47$ because of the formula of the exponent in Eq. (14) (i.e., $\frac{\beta^2 - \beta}{2}$).
 299 Moreover, the correlation between N_c and q_c can also be important. Going back to Eq.(23),

300 evidently, $E > E_q$ if and only if $E_N \cdot E_{COV} > 1$. After some manipulation, we can show that if
301 $\beta_N < 0$ and $\sigma_N > 0$, then

$$E_N \cdot E_{COV} > 1, \text{ if } \rho < \frac{\sigma_N}{\sigma_q} \cdot \frac{(1-\beta_N)}{2\beta_q}. \quad (25)$$

302 This equation reveals that when q_c and N_c are weakly or negatively correlated ($\rho \leq 0$),
303 considering only E_q would tend to underestimate E . On the other hand, however, if q_c and N_c
304 are highly positively correlated ($\rho \sim 1$) then considering E_q only would tend to overestimate E .

305 **3. Data and Methodology**

306 To derive the above-mentioned enhancement factors, we will use 10 years (2007 ~ 2016) of
307 the latest collection 6 (C6) *daily mean* level-3 cloud retrieval product from the Aqua-MODIS
308 instrument (product name "MYD08_D3"), which contains the gridded statistics of cloud
309 properties computed from pixel-level (i.e., level-2) retrievals. As summarized in Platnick et al.
310 (2003; 2017), the operational level-2 MODIS cloud product provides cloud masking (Ackerman et
311 al., 1998), cloud top height (Menzel et al., 1983), cloud top thermodynamic phase determination
312 (Menzel et al., 2006), and COT, cloud effective radius (CER) and LWP retrievals based on the bi-
313 spectral solar reflectance method (Nakajima and King, 1990). All MODIS level-2 atmosphere
314 products, including the cloud, aerosol and water vapor products, are aggregated to $1^\circ \times 1^\circ$ spatial
315 resolution on a daily, eight-day, and monthly basis. Aggregations include a variety of scalar
316 statistical information, including mean, standard deviation, max/min occurrences, as well as
317 histograms including both marginal and joint histograms. For COT, CER and LWP, the MODIS
318 level-3 product provides both their "in-cloud" grid-mean values ($\langle x \rangle$) and subgrid standard
319 deviations (σ_x). The inverse relative variance v can then be derived from Eq. (5), i.e., $v =$
320 $\langle x \rangle^2 / \sigma_x^2$. Note that the operational MODIS product provides two CER retrievals, one based on
321 the observation from the band 7 centered around $2.1 \mu\text{m}$ and the other from band 20 at $3.7 \mu\text{m}$.
322 As discussed in several previous studies (Cho et al., 2015; Zhang and Platnick, 2011; Zhang et al.,
323 2012; 2016), the $3.7 \mu\text{m}$ band CER retrieval is more resilient to the 3-D effects and retrieval failure
324 than the $2.1 \mu\text{m}$ band retrievals. For these reasons, it is used as the observational reference in
325 this study.

326 Given the COT and CER retrieval, the operational MODIS product estimates the LWP of cloud

327 using

$$LWP = \frac{2}{3} \rho_w COT \cdot CER, \quad (26)$$

328 where ρ_w is the density of water. Several studies have argued that a smaller coefficient of 5/9,
329 instead of 2/3, should be used in estimation of LWP (Lebsock et al., 2011; Seethala and Horváth,
330 2010; Wood and Hartmann, 2006). The choice of coefficient does not matter in this study because
331 it is a common factor in the calculation of ν . The choice of the coefficient has no impact on our
332 study, because we are interested in the relative inverse variance $\nu = \langle x \rangle^2 / \sigma_x^2$. We note here that
333 it is the LWC q_C , instead of the LWP, that is used in the KK2000 scheme. So, the spatial variability
334 of LWC is what is most relevant. However, the remote sensing of cloud water vertical profile from
335 satellite sensor for liquid-phase clouds is extremely challenging even with active sensors. It is why
336 most previous studies using the satellite observations analyzed the spatial variation of LWP,
337 rather than LWC. In fact, even Lebsock et al. (2013), who used the level-2 CloudSat observations,
338 had to use the vertical averaged LWC in their analysis. Airborne in situ measurement faces similar
339 challenge. For example, Boutle et al. (2014) use the LWC observation along “horizontal flight
340 tracks” to study the spatial variability of cloud water, which only samples the LWC at certain levels
341 of MBL clouds. Ground-based observations are much better than satellite and airborne
342 observation in this regard. Recently, Xie and Zhang (2015) analyzed the cloud water profiles
343 retrieved using ground-based radars from the three ARM sites and found no obvious in-cloud
344 vertical dependence of the spatial variability of LWC. Following these previous studies, we
345 assume that the horizontal subgrid variation of LWC is *not* strongly dependent on height and its
346 value can be inferred from the spatial variability of the vertical integrated quantity LWP. The
347 uncertainty caused by this assumption will be assessed in future studies.

348 The current MODIS level-3 cloud product does *not* provide CDNC retrievals. Following
349 previous studies (Bennartz, 2007; Bennartz and Rausch, 2017; Grosvenor and Wood, 2014;
350 McCoy et al., 2017a), we estimate N_c of liquid-phase clouds from the MODIS retrieved COT (τ)
351 and CER (r_e) based on the classic adiabatic cloud model

$$N_c(\tau, r_e) = \frac{\sqrt{5}}{2\pi k} \frac{\sqrt{f_{ad}\Gamma_w}}{\sqrt{\rho_w Q_e}} \tau^{\frac{1}{2}} r_e^{-\frac{5}{2}} = \frac{\sqrt{15}}{2\pi k} \frac{\sqrt{f_{ad}\Gamma_w}}{\rho_w \sqrt{2Q_e}} LWP^{\frac{1}{2}} r_e^{-3}, \quad (27)$$

352 where ρ_w is the density of water; $Q_e \approx 2$ is the extinction efficiency of cloud droplets; k is the
 353 ratio of r_e to mean volume-equivalent radius; f_{ad} is the adiabaticity of the cloud; Γ_w is the LWC
 354 lapse rate. Following previous studies, we assume $k = 0.8$ and $f_{ad} = 1.0$ to be constant and
 355 compute Γ_w from the grid mean liquid cloud top temperature and pressure. The theoretical
 356 basis and main uncertainty sources of the CDNC estimation based on the adiabatic cloud model
 357 from MODIS-like passive cloud retrievals are nicely reviewed by Grosvenor et al. (2018).

358 Ideally, the values of LWP and CDNC should be estimated on pixel-by-pixel basis from the
 359 level-2 MODIS product. However, pixel-by-pixel estimation is highly time consuming, which
 360 makes it difficult to achieve a global perspective. Using an alternative method, many previous
 361 studies estimate the grid-level CDNC statistics from the joint histogram of COT vs. CER provided
 362 in the level-3 MODIS cloud products (Bennartz, 2007; McCoy et al., 2017a; 2017b). For a given
 363 $1^\circ \times 1^\circ$ grid-box, the liquid-phase COT-CER joint histogram provides the counts of successful cloud
 364 property retrievals with respect to 108 joint COT-CER bins that are bounded by 13 COT bin
 365 boundaries, ranging from 0 to 150, and 10 CER bin boundaries, ranging from $4 \mu\text{m}$ to $30 \mu\text{m}$. With
 366 the joint histogram, which is essentially the joint PDF of COT and CER $P(\tau, r_e)$, we can estimate
 367 the grid mean and variance of CDNC from the following equations

$$\langle x \rangle = \int \int x(\tau, r_e) P(\tau, r_e) d\tau dr_e, \quad (28)$$

$$Var(x) = \int \int (x(\tau, r_e) - \langle N_c \rangle)^2 P(\tau, r_e) d\tau dr_e, \quad (29)$$

368 where x can be either LWP or CDNC. Figure 3a shows the LWP in Eq. (26) as a function of the 13
 369 COT bins and 10 CER bins from the MODIS level-3 product. As expected, the largest LWP values
 370 are found when both COT and CER are large. Figure 3b shows the CDNC in Eq. (27) as a function
 371 of the COT and CER bins. As expected, the largest CDNC values are found when both COT is large
 372 and CER is small. Figure 3c shows an example of the COT-CER joint histogram from the Aqua-
 373 MODIS daily level-3 product "MYD08_D3" on January 09th, 2007 at the grid box 1°S and 1°W . In
 374 this particular grid box, a combination of $\sim 2\text{-}4$ COT and $\sim 10\text{-}12 \mu\text{m}$ CER is the most frequently

375 observed cloud value. Using the joint histogram in Figure 3c, we can derive the mean and variance
376 of both LWP and COT using the Eqs. (28) and (29).

377 The efficiency of using the level-3 MODIS product is accompanied by three important
378 limitations. First of all, as mentioned earlier MODIS provides only LWP retrievals while LWC is
379 needed in the KK2000 scheme. Second, the current level-3 MODIS cloud product has a fixed $1^\circ \times 1^\circ$
380 spatial resolution. Although this resolution is highly relevant to the current generation of GCMs,
381 i.e., Coupled Model Intercomparison Project Phase 6 (CMIP5) (Eyring et al., 2016), future GCMs
382 may have significantly finer resolution. Third, it is difficult to sub-sample the pixels with the best
383 retrieval quality. These limitations will have to be addressed in future studies.

384 **4. Grid-mean and subgrid variations of liquid-phase cloud properties**

385 In this study, we limit our analysis to tropical oceans only where warm rain is frequent and
386 MODIS cloud retrievals have a relatively better quality than over land or over high latitude. The
387 annual mean total cloud fraction (f_{tot}), liquid-phase cloud fraction (f_{liq}), in-cloud COT, CER from
388 the $3.7 \mu\text{m}$ band, LWP and estimated CDNC over the tropical oceans based on 10 years Aqua-
389 MODIS retrievals are shown in Figure 4. The highest f_{liq} in the tropics is usually found in the
390 stratocumulus (Sc) decks over the Eastern boundary of the ocean, e.g., SE Pacific off coast of Peru,
391 NE Pacific off the coast of California and SE Atlantic off the coast of Namibia. The liquid-cloud
392 fraction reduces significantly toward the open ocean trade wind regions, where the dominant
393 cloud types are broken cumulus (Cu). Close to the continents, the Sc decks are susceptible to the
394 influence of continental air mass with higher loading of aerosols in comparison with pristine
395 ocean environment, which is probably the reason the SC decks have smaller CER and higher CDNC
396 than the open-ocean trade cumulus (Figure 4 d and f). The in-cloud COT (Figure 4 c) and LWP
397 (Figure 4 e) generally increase from the Sc decks to the open-ocean Cu regime, although less
398 dramatically than the transition of cloud fraction. The Sc decks and the Sc-to-Cu transition are
399 the most prominent features of liquid-phase clouds in the tropics. However, as mentioned in the
400 introduction, simulating these features in the GCMs proves to be an extremely challenging task,
401 and most GCMs suffer from some common problems, such as the “too few too bright” problem
402 and the abrupt Sc-to-Cu transition problem (Kubar et al., 2014; Nam et al., 2012; Song et al.,
403 2018a).

404 Switching the focus now from grid-mean values to subgrid variability, we will show the
 405 grid-level inverse relative variances $v = \langle x \rangle^2 / Var(x)$ for several key cloud properties. Here, we
 406 first derive the daily mean v and then aggregate the result to monthly mean values. Therefore,
 407 for each grid box we have 120 samples (i.e., 10 years x 12 months) of monthly mean v for analysis
 408 and visualization. Because the value of v can be ill-behaved when $Var(x)$ approaches zero,
 409 instead of the mean value, we plot the median value of \tilde{v} based on 120 months of MODIS
 410 observations in Figure 5. There are several interesting and important features in Figure 5. First of
 411 all, the \tilde{v} of all four sets of cloud properties (i.e., COT, CER, LWP and CDNC) all exhibits a clear
 412 and similar Sc-to-Cu transition, with larger values in the Sc region and smaller value in the broken
 413 Cu regions. This indicates that cloud properties, including both optical and microphysical
 414 properties, are more homogenous, in terms of spatial distribution within the grid, in the Sc region
 415 than in the Cu region. Secondly, the value of \tilde{v} of CER (i.e., 10~100 in Figure 5b) is larger than
 416 that of the other properties (i.e., 1~10) by almost an order of magnitude, indicating that the
 417 subgrid variability of CER is very small. On the other hand, however, it is important to note that
 418 the \tilde{v} of CDNC (Figure 5d) is comparable with that of COT (Figure 5a) and LWP (Figure 5c). The
 419 reason is probably in part because the highly nonlinear relationship between CDNC and CER (i.e.,
 420 $N_c \sim r_e^{-\frac{5}{2}}$) leads to a stronger variability of CDNC than CER, and also in part because the variability
 421 of CDNC is also contributed by the subgrid variation of COT. In some regions, the Gulf of Guinea,
 422 East and South China Sea, and Bay of Bengal for example, the \tilde{v} of CDNC is close to unity,
 423 indicating the subgrid standard deviation of CDNC is comparable to the grid-mean values in these
 424 regions. As discussed in the next section, the significant subgrid variability of CDNC in these
 425 regions should be taken into account when modeling the nonlinear processes, such as the
 426 autoconversion, in GCM to avoid systematic biases due to the nonlinearity effect.

427 The values of \tilde{v} in Figure 5 from this study are in reasonable agreement with previous
 428 studies. Barker (1996) selected a few dozens of cloud scenes, each about 100 ~ 200 km in size,
 429 from the Landsat observation and analyzed their spatial variability of COT. It is found that the
 430 typical value of v for “overcast stratocumulus”, “broken stratocumulus” and “scattered cumulus”
 431 is 7.9, 1.2, and 0.7, respectively (see their Table 3), which is consistent with the Sc-to-Cu transition
 432 pattern seen in Figure 5. Oreopoulos and Cahalan (2005) derived the subgrid inhomogeneity of

433 COT on a global scale from the level-3 Terra-MODIS retrievals. Although using a different metric
434 (i.e., their inhomogeneity parameter is defined as $\chi = \exp(\ln\langle\tau\rangle) / \langle\tau\rangle$), they also found
435 systematic increase of inhomogeneity (decreasing value of χ) from the Sc region to cu region.
436 Also using the MODIS cloud property retrievals, Wood and Hartmann(2006) investigated the
437 meso-scale spatial variability of LWP in the NE Pacific and SE Pacific region. The ν of LWP is found
438 to increase systematically with meso-scale cloud fraction and the relationship between the two
439 can be reasonably explained by a simple PDF cloud thickness model in Considine et al. (1997).
440 See also Kawai and Teixeira (2010).

441 As explained in section 2, the correlation between cloud water and CDNC can also
442 influence the computation of enhancement factor and thereby the grid-mean autoconversion
443 rate. Figure 5e shows the median value of the LWP and CDNC correlation coefficient $\tilde{\rho}$. Similar to
444 the derivation of median $\tilde{\nu}$, we first compute the monthly mean ρ from daily MODIS observations
445 and then derive the median value of $\tilde{\rho}$ for each grid from the 120 months of observation. As
446 shown in Figure 5e, at the subgrid level, the LWP and CDNC tend to be positively correlated
447 almost over all tropical oceans. Mathematically, this is not surprising because as shown in Figure
448 5b and c, the subgrid variability of r_e is order of magnitude smaller than that of LWP. Since CDNC
449 is proportional to $LWP^{\frac{1}{2}}r_e^{-3}$ according to Eq. (27), the subgrid variability of CDNC is mainly
450 determined by the variability of LWP, leading to the positive correlation. Physically, the
451 correlation can be explained by several mechanisms. For example, Wood et al. (2018) and O et
452 al. (2018) found that a large amount of low-level water clouds over the stratocumulus to cumulus
453 transition are “optically thin veil clouds”. These clouds are usually associated with low LWP and
454 low CDNC (therefore positive correlation) and probably caused by the strong precipitation
455 scavenging process in the active cumulus. Note that our definition of ρ is the subgrid spatial
456 correlation of LWP and CDNC. It may be different from the definition used in many aerosol
457 indirect effect studies where the temporal correlation of monthly mean LWP and CDNC is more
458 interested.
459

460 5. Implications for warm-rain simulations in GCM

461 5.1. Influence of subgrid variation of cloud water

462 As discussed in Section 2.2, most current GCMs only considers the impact of subgrid cloud
463 water variation on autoconversion rate but ignore the impact of subgrid CDNC variation. To make
464 our analysis relevant to the current GCMs, we first analyze E_q in Eq. (20) based on observation.
465 The impacts of subgrid CDNC variation (i.e., E_N) and its correlation with cloud water (i.e., E_{COV})
466 will be analyzed in the next section.

467 We derive E_q using two approaches. First, we derive it from the observed LWP PDF based on
468 Eq. (20). As such, we do not have to make any assumption about the shape of LWP PDF although
469 solving the integration in Eq. (20) is time-consuming. In the second approach, we first derive the
470 relative inverse relative variance ν of LWP and then derive the enhancement factor by assuming
471 the subgrid PDF to be either Gamma or Lognormal. This approach is more efficient, but it may be
472 subject to error if the true PDF deviates from the assumed PDF shape. Figure 6a shows the annual
473 mean enhancement factor E_q in the tropical region derived based on Eq. (20) (i.e., the first
474 approach) from 10 years of MODIS observation. Figure 6 b and c show the annual mean
475 enhancement factor E_q derived by assuming the subgrid cloud water follows the Lognormal (i.e.,
476 Eq. (14)) and Gamma distribution (i.e., Eq. (8)), respectively. There are a couple of interesting and
477 important points to note. First of all, similar to the grid-mean quantities in Figure 4, the
478 enhancement factor E_q also shows a clear Sc-to-Cu transition. Over the Sc decks, because clouds
479 are more homogeneous ($\tilde{\nu} > 5$), the enhancement factor E_q is only around 1 ~ 2.5, while over
480 the Cu regions, the more inhomogeneous clouds with $\tilde{\nu} < 1$ leads to a larger enhancement
481 factor E_q around 3~5. As aforementioned, in the current CAM5.3, E_q is assumed to be a constant
482 of 3.2. While this value is within the observational range, it obviously cannot capture the Sc-to-
483 Cu transition. In fact, the constant value 3.2 overestimates the E_q over the Sc region and
484 underestimates the E_q over the Cu region, which could lead to unrealistic drizzle production in
485 both regions and to consequential impacts on cloud water budget, radiation and even aerosol
486 indirect effects on the model. The second point to note is that the E_q based on the Lognormal
487 PDF assumption in Figure 6 b agrees well with the results in Figure 6 a derived directly from the
488 observation. In contrast, the E_q based on the Gamma PDF assumption in Figure 6 c tends to be

489 smaller, especially in the Cu regions. This result seems to suggest that the Lognormal distribution
490 provides a better fit to the observed subgrid cloud water variation than the Gamma distribution,
491 which has rarely been noted and reported in the previous studies.

492 A flexible, cloud-regime dependent E_q could help improve the simulation of Sc-to-Cu
493 transition in the GCM. If a GCM employs an advanced cloud parameterization scheme, such as
494 CLUBB, that is able to provide regime-dependent information on subgrid cloud variation, i.e., v ,
495 then the enhancement factor E_q could be diagnosed from v . However, most traditional cloud
496 parameterization schemes do not provide information on subgrid cloud variation. In such case, if
497 one does not wish to use a constant E_q , but a varying regime-dependent scheme, then either v
498 or E_q need to be parameterized as a function of some grid-mean cloud properties resolved by
499 the GCM. In fact, several attempts have been made along this line. Based on the combination air-
500 borne in situ measurement and satellite remote sensing product, Boutle et al. (2014)
501 parameterized the “fractional standard deviation” (which is equivalent to $1/\sqrt{v}$ in our definition)
502 of liquid-phase cloud as a function of grid-mean cloud fraction. This scheme was later updated
503 and tested in a host GCM in Hill et al. (2015), and was found to reduce the shortwave cloud
504 radiative forcing biases in the model. In a recent study, Xie and Zhang (2015) derived the subgrid
505 cloud variations from the ground-based observations from three Department of Energy (DOE)
506 Atmospheric Radiation Measurement (ARM) sites, and then parameterize the inverse relative
507 variance v as a function of the atmospheric stability.

508 Figure 7a shows the variation of inverse relative variance v as a function of the grid-mean
509 liquid-phase cloud fraction f_{liq} . In general, the value of v increases with the increasing f_{liq} , which
510 is expected from the Sc-to-Cu increase of f_{liq} in Figure 4b and the Sc-to-Cu decrease of v in Figure
511 5c. The $v(f_{liq})$ pattern in Figure 7a is also consistent with the results reported in Wood and
512 Hartmann (2006) and Lebsock et al. (2013). In the hope of obtaining a simple parameterization
513 scheme for $v(f_{liq})$ that can be used in GCMs, we fit the median value of v as a simple 3rd order
514 polynomial of f_{liq} as follows:

$$v(f_{liq}) = 2.38 - 4.95f_{liq} + 8.74f_{liq}^2 - 0.49f_{liq}^3, \quad f_{liq} \in [0,1]. \quad (30)$$

515 For comparison, we also plotted the $v(f_{liq})$ parameterization developed by Boutle et al. (2014)
 516 for grid size of 100 km (dashed line in the figure). Apparently, comparing Eq. (30), it generates
 517 slightly smaller v for a given the f_{liq} . This difference is probably because our results are based on
 518 different data.

519 To test the performance of this simple parameterization, we first substitute the f_{liq} from
 520 MODIS daily mean level-3 product into the above equation and then use the resultant v to
 521 compute the enhancement factor E_q . Unfortunately, the enhancement factor E_q computed
 522 based on the parameterized $v(f_{liq})$ as shown in Figure 8a substantially underestimate the
 523 observation-based results in Figure 6, especially over the Cu regions. The deviation is probably
 524 because the relationship between E_q and v is highly nonlinear (e.g., Eq. (8) and (14)) and
 525 therefore the above parameterization scheme that only fits the = value of v is not able to capture
 526 the variability of E_q . Based on this consideration, we tried an alternative approach. Instead of
 527 parameterization of v , we directly parameterize the enhancement factor E_q as a function of f_{liq} .
 528 Figure 7b shows the variation of E_q as a function of f_{liq} . As expected, E_q generally decreases with
 529 increasing f_{liq} . The median value of E_q is fitted with the following 3rd order polynomial of f_{liq}

$$E_q(f_{liq}) = 2.72 + 7.33f_{liq} - 19.17f_{liq}^2 + 10.69f_{liq}^3, \quad f_{liq} \in [0,1]. \quad (31)$$

530 As shown in Figure 8b, the value of E_q based on the above equation clearly agrees with the
 531 observation-based values in Figure 6 better than that based on the parameterization of $v(f_{liq})$.
 532 The elimination of the middle step indeed improves the parameterization results. While this is
 533 encouraging, it should be kept in mind that the Eq. (31) has very limited application, i.e., it is only
 534 useful for the autoconversion rate computation for a particular value of the autoconversion
 535 exponent beta, i.e., $\beta_q = 2.47$. A good parameterization of v could be useful for not only
 536 autoconversion, but also for accretion and radiation computations. Another caution is that, if
 537 applied to a GCM, the performance of the $E_q(f_{liq})$ parameterization in Eq. (31) will be dependent
 538 on the simulated accuracy of f_{liq} in the model.

539

5.2. Influence of subgrid variance of CDNC

540
541 Now we will investigate the impacts of subgrid CDNC variation on the autoconversion rate
542 simulation. For the moment, we will consider E_N only. The impact of CDNC and cloud water
543 correlation will be discussed in the next section. Similar to E_q we first derive E_N from the CDNC
544 PDF based on Eq. (21). The annual mean result based on 10 years of MODIS observations is shown
545 in Figure 9a. There are several intriguing points to note. First of all, the value of E_N is actually
546 larger than E_q in Figure 9 such that we even have to use a different color scale for this plot.
547 Secondly, E_N the regions with escalated E_N seem to coincide with the downwind regions of
548 biomass burning aerosols (e.g., Gulf of Guinea, East Coast of South Africa), air pollution (i.e.,
549 Eastern China Sea), and, most interestingly, active volcanos (e.g., Kilauea Hawaii and Ambae
550 Vanuatu). We have also checked the seasonal variation of the E_N and the results also support
551 this observation. Another interesting feature to note is that, although the dust outflow regions
552 such as Tropical East Atlantic and Arabian Sea, have heavy aerosol loading, the value of E_N there
553 is only moderate. Figure 9b shows the value of E_N computed based on Eq. (14) from the inverse
554 relative variance of v , assuming that the subgrid CDNC follows a Lognormal PDF. Although the
555 overall pattern is consistent with Figure 9a, the assumption of Lognormal PDF seems to
556 underestimate E_N . A closer examination indicates that the Lognormal PDF tend to underestimate
557 the population of clouds with small CDNC, and therefore underestimate the variance of CDNC as
558 well as E_N . We did not compute the E_N based on the Gamma distribution because of the singular
559 value problem aforementioned in Section 2.1.

560 We could not find any previous observation-based study on the global pattern of the
561 subgrid variation of CDNC and the corresponding E_N . So, it is difficult for us to corroborate our
562 results. On one hand, the magnitude of E_N is surprisingly large. As explained in Section 3, the
563 CDNC is estimated based on Eq. (27) from the MODIS retrieval of COT and CER. Several previous
564 studies have shown that the sub-pixel level surface contamination, subpixel cloud
565 inhomogeneity, and three-dimensional radiative transfer effects, can cause significant errors in
566 the MODIS CER retrievals especially over broken cloud regions (Zhang and Platnick, 2011; Zhang
567 et al., 2012; 2016). Given the fact that the CDNC retrieval is highly sensitive to CER error as a
568 result of $N_d \sim r_e^{-\frac{5}{2}}$, the influence of retrieval uncertainty on subgrid CDNC variation cannot be

569 ruled out. On the other hand, the pattern of E_N in Figure 9a seems to suggest that there are some
570 underlying physical mechanisms controlling the subgrid variation of CDNC, in which aerosols
571 seem to play an important role. To achieve a better understanding, we analyzed the dependence
572 of E_N on liquid cloud fraction and grid-mean CDNC in Figure 10, which reveals that E_N has a
573 stronger dependence on CDNC than cloud fraction. This result seems to indicate that the pattern
574 of E_N in Figure 9 is largely determined by physical mechanisms rather than retrieval
575 uncertainties. Interestingly, the largest E_N is usually found when liquid cloud fraction is small and
576 CDNC is large and decreases with decreasing CDNC and increasing cloud fraction. This pattern
577 leads us to the following hypothesis: In the regions where aerosol is limited, even weak updraft
578 can activate most cloud condensation nuclei (CCN). As a result, even if there is significant subgrid
579 variation of turbulence at cloud base, the subgrid variation of CDNC remains small. In contrast,
580 in regions where aerosol is abundant, the subgrid variation of turbulence becomes important.
581 The subgrid variation of updraft leads to subgrid variation CDNC and thereby large E_N .

582 As far as we know, the results in Figure 9 and Figure 10 mark the first attempt based on
583 satellite observations to unveil the global pattern of the subgrid variations of CDNC and
584 investigate the consequential impacts on warm rain simulations in GCMs. Although obscured by
585 satellite retrieval uncertainties, the results still provide valuable insights. First of all, the
586 enhancement factor E_N due to the subgrid variations of CDNC is nonnegligible, even comparable
587 the effect of subgrid cloud water variation (i.e., E_q). Second, the global pattern of E_N in Figure 9
588 provides a valuable map for future studies.

589 **5.3. The combined effect of subgrid variations of cloud water and CDNC**

590 Finally, in this section we examine the combined effect of subgrid variations of cloud
591 water and CDNC, as well as their correlation, on the autoconversion rate simulation. The annual
592 mean combined enhancement factor E derived based on Eq. (17) from 10 years of MODIS COT
593 and CER observation is shown in Figure 11a. Comparing to the E_q in Figure 6 and E_N in Figure 9,
594 the combined enhancement factor is generally larger. It is easy to see that the in some regions
595 (e.g., Gulf of Guinea, East Coast of South Africa and Eastern China Sea) the combined
596 enhancement factor E resembles the E_N while in other regions (i.e., trade wind cumulus regions
597 over open ocean) it resembles more of E_q . Interestingly, because both E_q and E_N are small over

598 the Sc decks, those regions have the smallest combined enhancement factor E . As discussed in
599 Section 2.2, only when the subgrid variation of cloud water is uncorrelated with the subgrid
600 variation of CDNC can the combined enhancement factor E be decomposed into the simple
601 product of E_q and E_N (i.e., Eq. (19)). Figure 11b shows the annual mean value of the simple
602 product $E_q \cdot E_N$, without considering the correlation between cloud water and CDNC. Evidently,
603 the simple product substantially overestimates the combined enhancement factor derived from
604 the joint PDF of LWP and CDNC. This result can be explained by the mostly positive subgrid
605 correlation between LWP and CDNC in Figure 5e. As explained in section 2.2, the positive
606 correlation means that clouds with more water also tend to have more CDNC. The
607 autoconversion rate of such configuration is lower than that when LWP and CDNC have no
608 correlation.

609 Together, the E_q in Figure 6, E_N in Figure 9 and the combined enhancement factor in
610 Figure 11 lead us to the following important conclusion. It is not sufficient to consider only the
611 impact of subgrid variation of cloud water (i.e., E_q) on the autoconversion rate simulation. The
612 influences of subgrid CDNC variation, as well as the correlation between cloud water and CDNC,
613 must also be taken into account to avoid significant error.

614 Finally, the combined enhancement factor derived based on Eq. (23) assuming that the
615 LWP and CDNC follow the bi-variate lognormal distribution is shown in Figure 11c. Despite the
616 tendency of overestimation, the result agrees reasonably well with that based on observed joint
617 PDF in Figure 11a, clearly better than the simple product $E_q \cdot E_N$. This is encouraging as it
618 suggests that the bi-variate lognormal distribution can be used in the future to model the
619 combined effect of cloud water and CDNC on autoconversion rate simulation in GCMs.

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621 **6. Summary and Outlook**

622 One of the difficulties in GCM simulation of the warm rain parameterization is how to
623 account for the impact of subgrid variations of cloud properties, such as cloud water and CDCN,
624 on nonlinear precipitation processes such as autoconversion. In practice, this impact is often
625 treated by adding the enhancement factor term to the parameterization scheme. In this study,
626 we derived the subgrid variations of liquid-phase cloud properties over the tropical ocean using

627 the satellite remote sensing products from MODIS and investigated the corresponding
628 enhancement factors for parameterizations of autoconversion rate. In comparison with previous
629 work, our study is able to shed some new light on this problem in the following regards:

- 630 1. A theoretical framework is presented to explain the importance of the subgrid
631 variation of CDNC and its correlation with cloud water on the autoconversion rate
632 simulation in GCMs.
- 633 2. The wide spatial coverage of the Level-3 MODIS product enables us to depict a
634 detailed quantitative picture of the enhancement factor E_q , which shows a clear
635 cloud regime dependence, i.e., a Sc-to-Cu increase. The constant $E_q = 3.2$ used in
636 the current CAM5.3 model overestimates and estimates the observed E_q in the Sc
637 and Cu regions, respectively.
- 638 3. The E_q based on the Lognormal PDF assumption performs significantly better than
639 that based on the Gamma PDF assumption. A simple parameterization scheme is
640 provided to relate E_q to the grid-mean liquid cloud fraction, which can be readily
641 used in GCMs.
- 642 4. For the first time, the enhancement factor E_N due to the subgrid variation of CDNC
643 is derived from satellite observation, and the results reveal several regions
644 downwind of biomass burning aerosols (e.g., Gulf of Guinea, East Coast of South
645 Africa), air pollution (i.e., Eastern China Sea), and active volcanos (e.g., Kilauea
646 Hawaii and Ambae Vanuatu). The largest E_N is usually found where CDNC is large
647 and liquid cloud fraction is small and decreases with decreasing CDNC and
648 increasing cloud fraction.
- 649 5. MODIS observations suggest that the subgrid LWP and CDNC are mostly positively
650 correlated. As a result, the combined enhancement factor is significantly smaller
651 than the simple product of $E_q \cdot E_N$ (i.e., assuming no correlation). The combined
652 enhancement factor derived assuming LWP and CDNC to follow the bi-variate
653 lognormal distribution agree with the observation-based results reasonably well.

654 As noted in the previous sections, this study has several important limitations, most of
655 which are a result of using the level-3 MODIS observations. The fixed $1^\circ \times 1^\circ$ spatial resolution of

656 MODIS level-3 product makes it impossible for us to investigate the scale-dependence of subgrid
657 cloud variation. Similar to previous studies, we have to make several assumptions when
658 estimating the CDNC from level-3 MODIS product. Furthermore, the retrieval uncertainties
659 associated with the optically thin clouds in MODIS product pose a challenging obstacle for the
660 quantification of subgrid cloud property variations and the corresponding enhancement factors.
661 These limitations have to be addressed using additional independent observations from, for
662 example, ground based remote sensing product and/or in situ measurement from air-borne field
663 campaigns. Recently, a few novel methods have been developed to provide certain information
664 on the subgrid cloud property variations to the host GCM. Most noticeable examples are the
665 super-parameterization method (a.k.a. multi-scale modeling framework) (Wang et al., 2015) and
666 the PDF-based higher-order turbulence closure methods (e.g., Cloud Layer Unified By Binormals,
667 CLUBB (Golaz et al., 2002a; Guo et al., 2015; Larson et al., 2002) and Eddy-Diffusivity Mass-Flux
668 (EDMF) (Sušelj et al., 2013)). The subgrid cloud property variations derived in this study provide
669 the valuable observational basis for the evaluation and improvement of these schemes.
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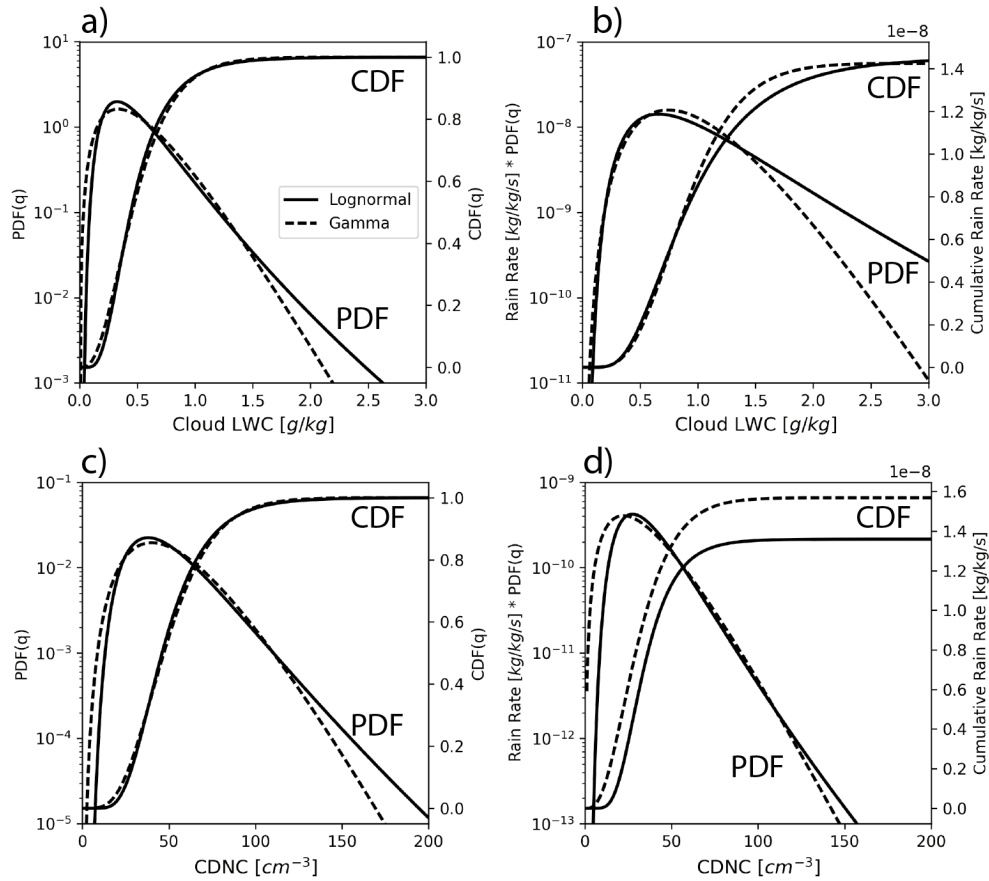
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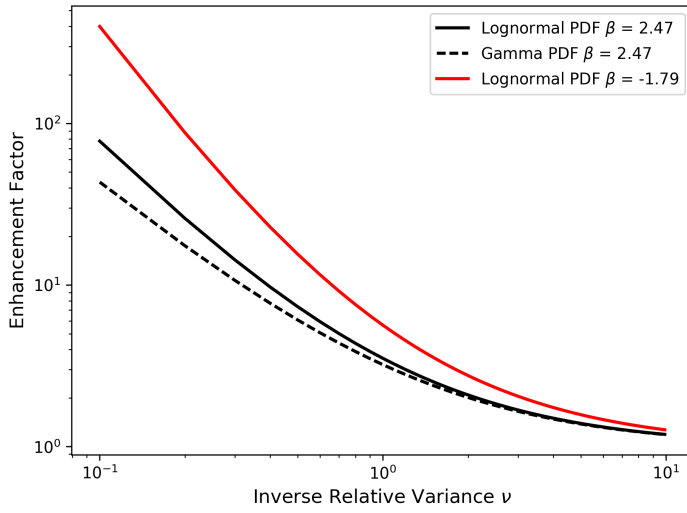
691 Figures:



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693 *Figure 1* a) The probability density function (PDF) and cumulative distribution function (CDF) of
694 cloud LWC (q_c) that follow the Gamma (dashed) and Lognormal (solid) distribution. For the both
695 distributions, $\langle q_c \rangle = 0.5g/kg$ and $v_q = 3.0$. b) The PDF and CDF of autoconversion rate
696 computed based on the KK2000 scheme in Eq.(15) and the PDF of q_c . In the computation, the N_c
697 is kept at a constant of $50 cm^{-1}$. c) The PDF and CDF of N_c that follow the Gamma (dashed) and
698 Lognormal (solid) distribution. For the both distributions, $\langle N_c \rangle = 50cm^{-3}$ and $v_N = 5.0$. d) the
699 PDF and CDF of the autoconversion rate computed based on the KK2000 scheme in Eq. (15) and
700 the PDF of N_c . The q_c is kept at $0.5g/kg$ in the computation.

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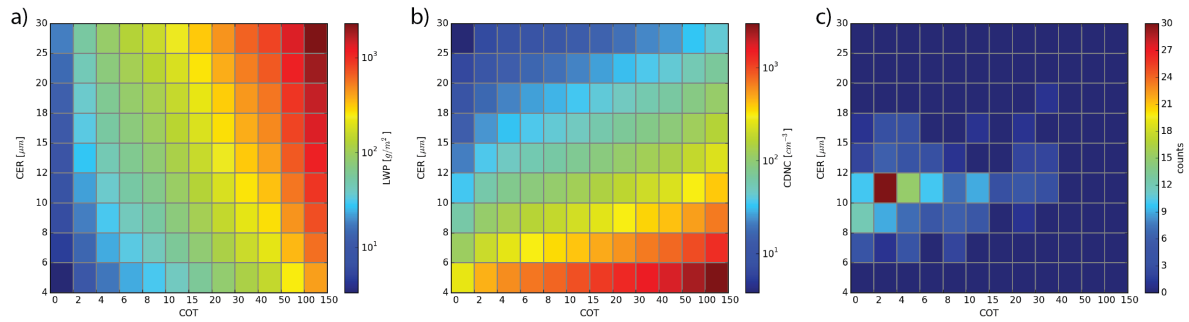
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703 *Figure 2 Enhancement factors based on Lognormal $E(P_L, \beta)$ and Gamma $E(P_G, \beta)$ subgrid PDF*

704 *for different β as a function of the inverse relative variance ν .*

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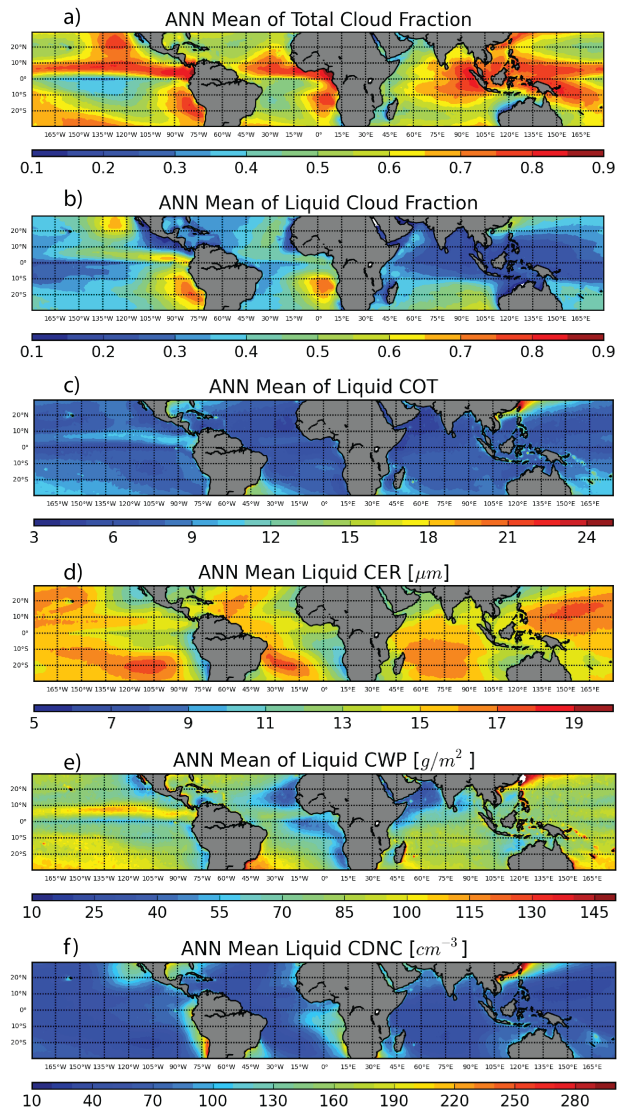
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Figure 3 The (a) LWP and (b) CDNC as a function of COT and CER. (c) An example of the COT-CER joint histogram observed by Aqua-MODIS on Jan. 09th, 2007 at 1°S and 1°W.

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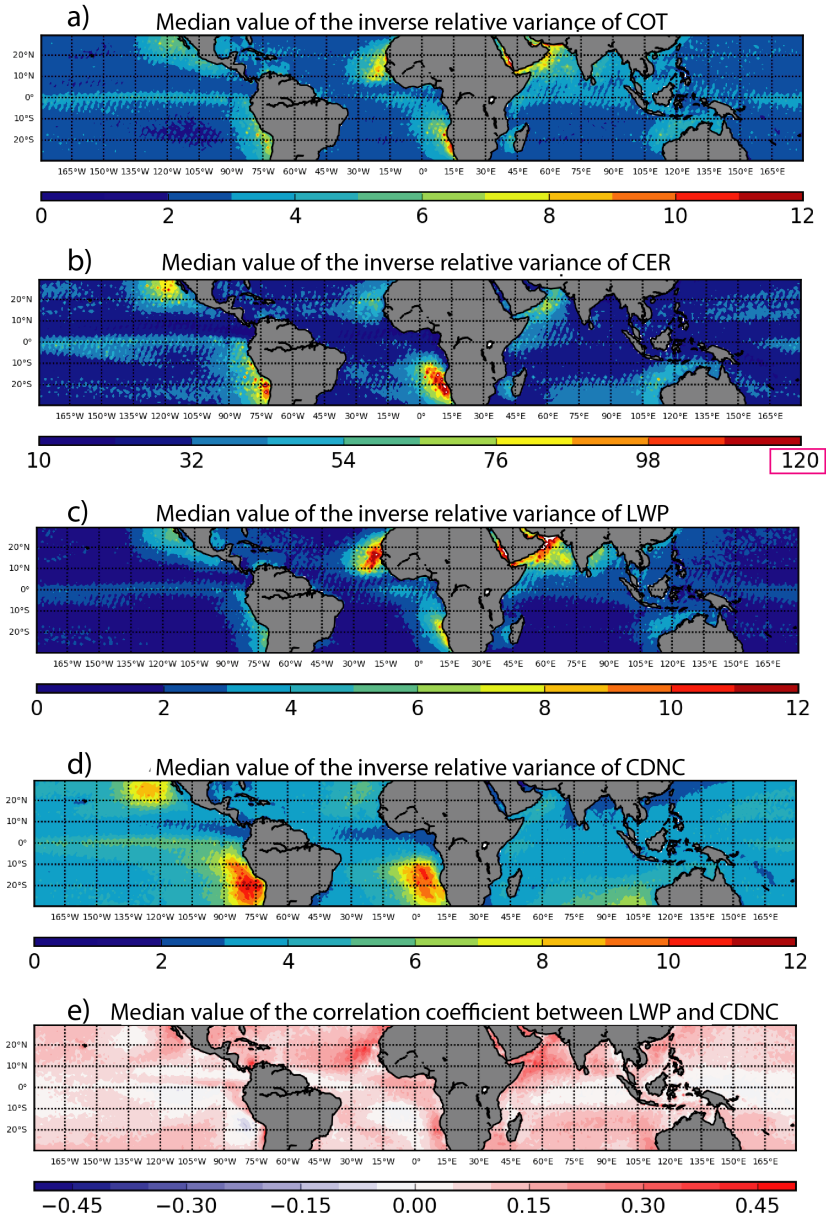
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714 *Figure 4* 10-year (2007~2016) averaged annual mean a) total cloud fraction, b) liquid cloud
715 fraction, c) cloud optical thickness, d) cloud effective radius retrieved from the 3.7 μm band, e)
716 cloud water path and f) cloud droplet concentration retrievals from Aqua-MODIS over the
717 tropical (30° S-30° N) oceans. All quantities are “in-cloud” mean that are averaged over the
718 cloudy-part of the grid only.

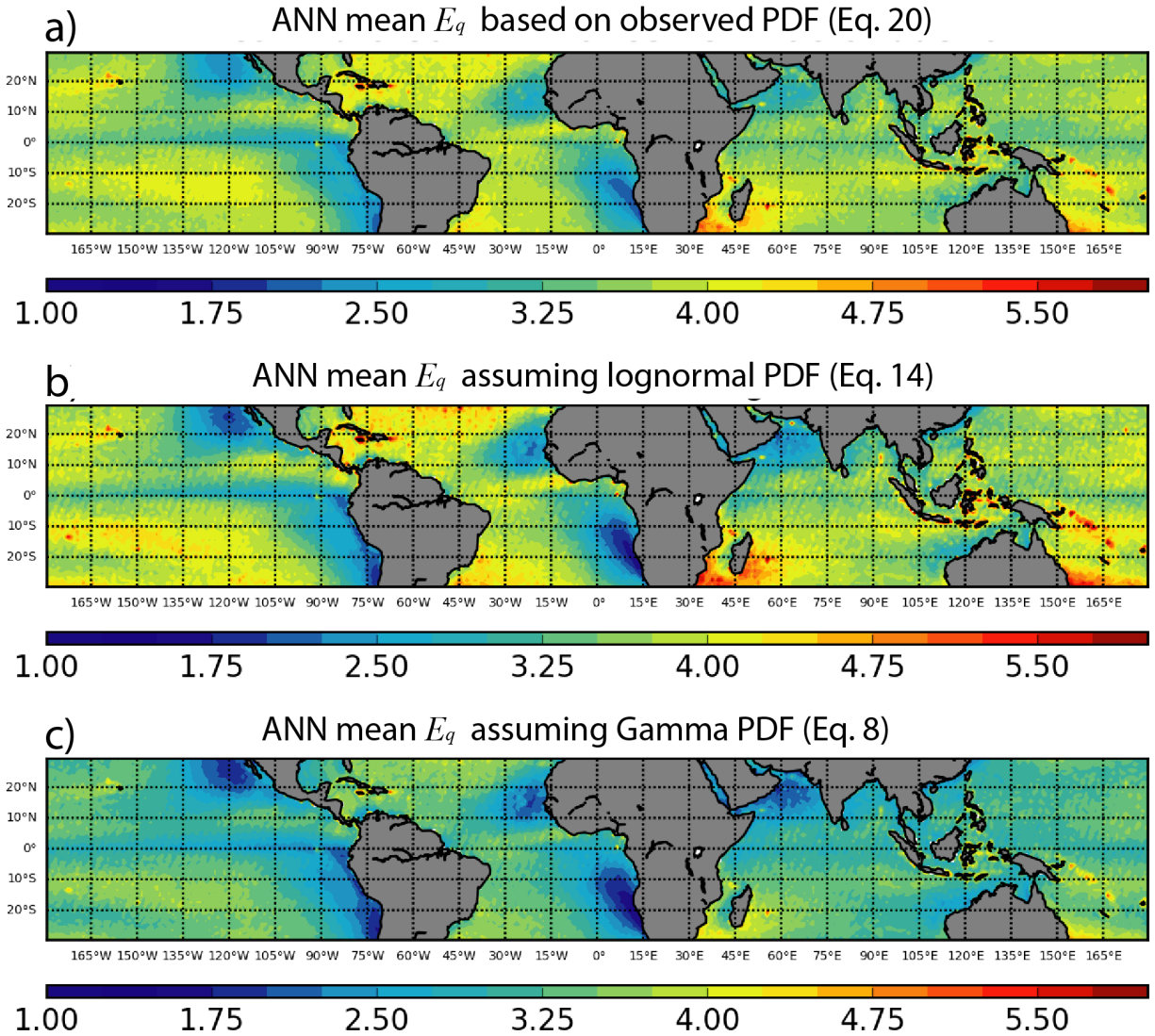
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721 *Figure 5* Median value of the inverse relative variance (i.e., $v = \langle x \rangle^2 / Var(x)$) for a) COT, b)
 722 CER, c) LWP and d) CDNC, and e) median value of the correlation coefficient between LWP and
 723 CDNC derived from 10 years of MODIS observations. Note that the color scale of CER is
 724 different from others’.

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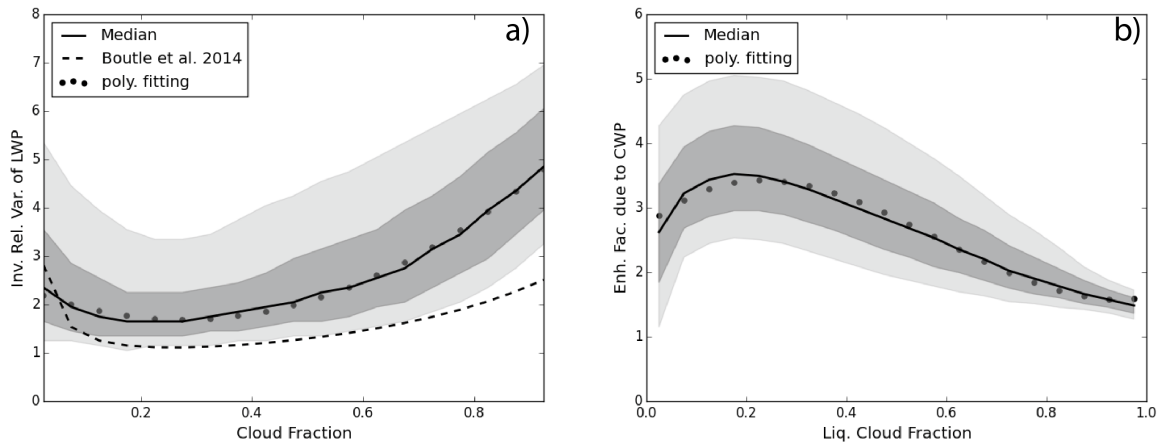
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Figure 6 The annual mean factor for the KK2000 scheme due to subgrid variation of LWP computed a) directly from observation, i.e., E_q in Eq.(20), b) from relative variance assuming Lognormal PDF of LWP, i.e., E_q in Eq.(14) and c) from relative variance assuming the Gamma PDF of LWP i.e., E_q in Eq.(8).

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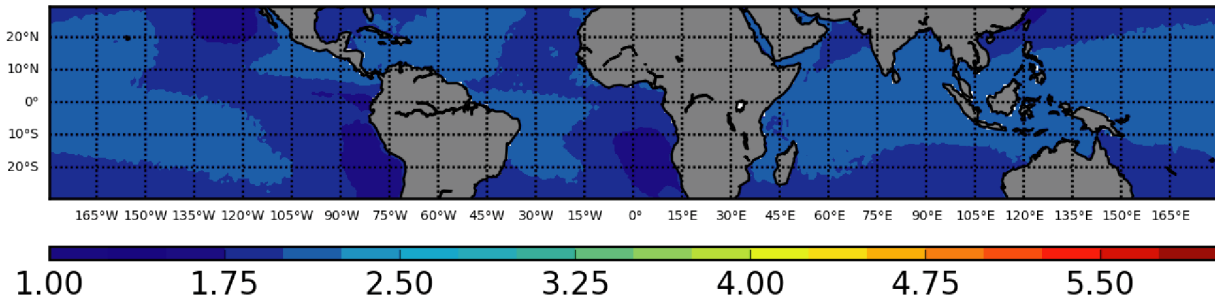


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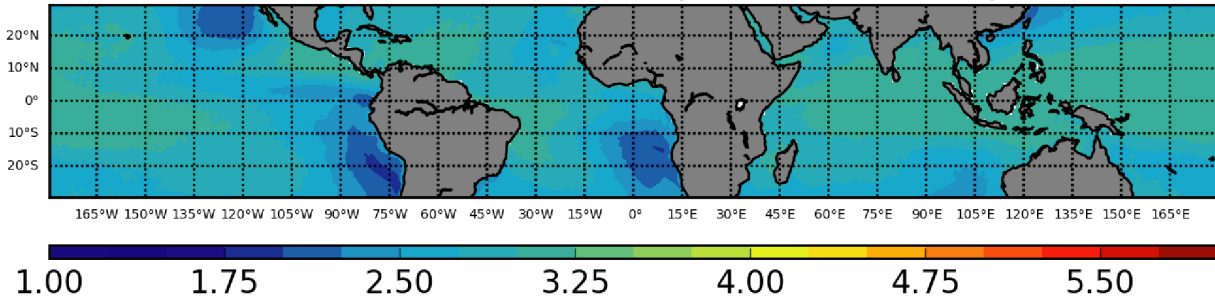
736 *Figure 7* a) The inverse relative variance v and b) autoconversion enhancement factor due to
737 LWP subgrid variability assuming Log-normal PDF as a function of grid-mean liquid cloud
738 fraction, where the solid line, dark shaded area, and light shaded area correspond to the
739 median value, 25%~75% percentiles, and 10~90% percentiles, respectively. The dotted lines
740 correspond to simple 3-rd order polynomial fitting. For reference, the $v_q(f_{liq})$ parameterization
741 scheme based on Boutle et al. (2014) is also plotted in a).

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a) ANN mean E_q derived from parameterized v_q (Eq. 30)



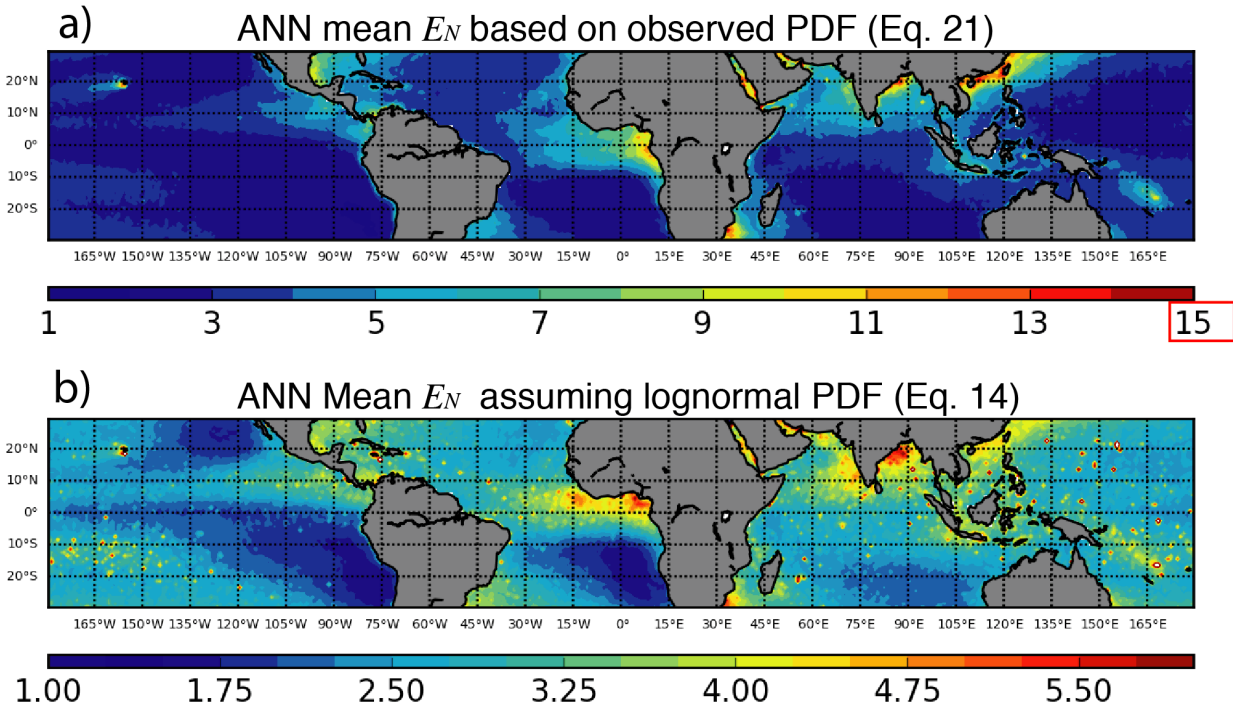
b) ANN mean E_q derived from E_q parameterization (Eq. 31)



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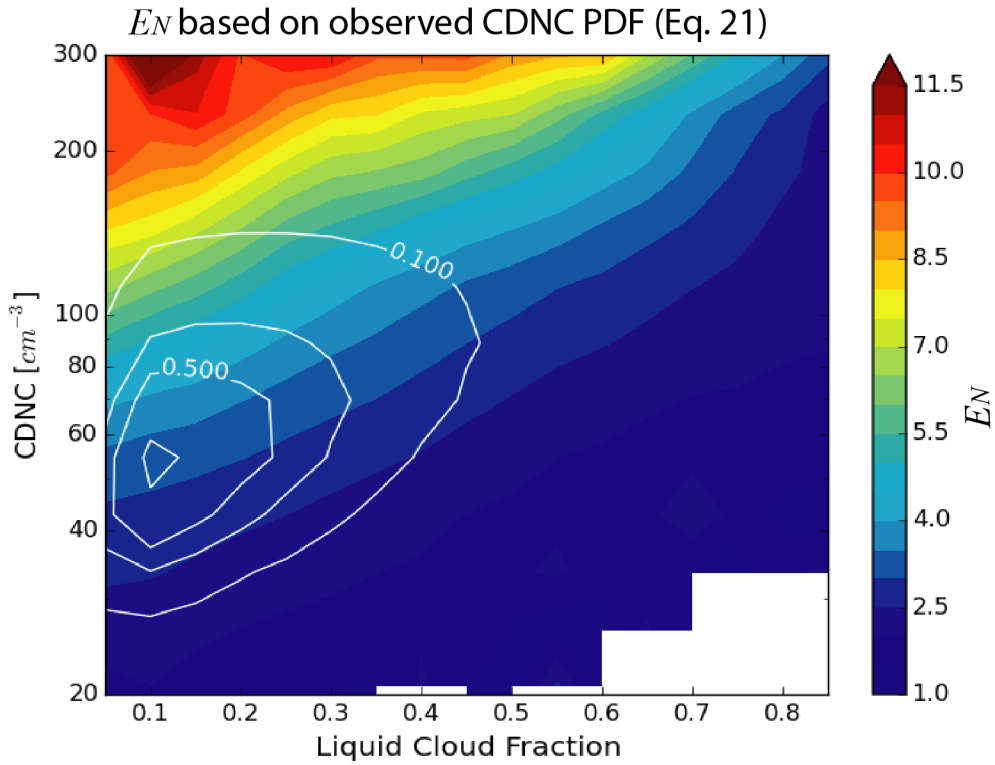
Figure 8 Annual mean value of the enhancement factor E_N computed based on the a) $v(f_{liq}) = 2.38 - 4.95f_{liq} + 8.74f_{liq}^2 - 0.49f_{liq}^3$ parameterization scheme in Eq. (30) and b) $E_q(f_{liq}) = 2.72 + 7.33f_{liq} - 19.17f_{liq}^2 + 10.69f_{liq}^3$ parameterization scheme in Eq. (31).

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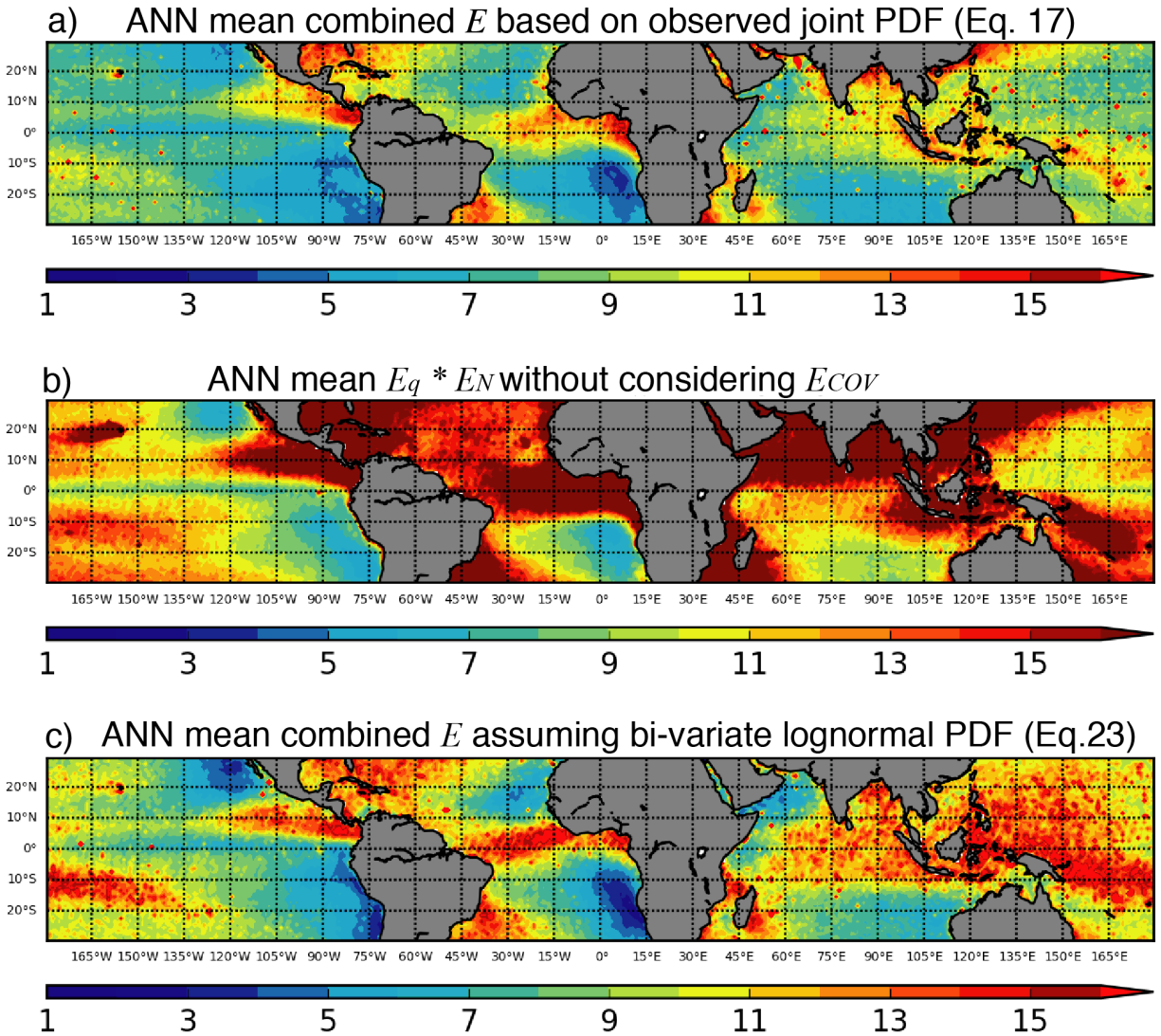
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Figure 9 Annual mean value of the enhancement factor E_N derived from a) observation based on Eq. (21) and b) from Eq. (14) assuming Lognormal subgrid CDNC distribution.



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 756 *Figure 10 Dependence of E_N on f_{liq} and N_d . The color map corresponds to the mean value of E_N*
 757 *for a given N_d and f_{liq} bin. The white contour lines correspond to the relative sampling*
 758 *frequency of N_d and f_{liq} bins (i.e., the most frequently observed combination is $N_d \sim 50\text{cm}^{-3}$*
 759 *and $f_{liq} \sim 0.1$).*

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764 *Figure 11 The combined enhancement factor derived a) based on Eq. (17) from the observed*
 765 *joint PDF of LWP and CDNC, b) assuming that subgrid variations of LWP and CDNC are*
 766 *uncorrelated, i.e., $E_q \cdot E_N$ only and c) based on Eq. (23) assuming that the subgrid LWP and*
 767 *CDNC following the bi-variate lognormal distribution.*

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772 **Reference:**

773 Ackerman, S., Strabala, K., Menzel, W., Frey, R., Moeller, C. and Gumley, L.: Discriminating clear
774 sky from clouds with MODIS, *Journal of Geophysical Research*, 103(D24), 32,141–32,157, 1998.

775 Ahlgrimm, M. and Forbes, R. M.: Regime dependence of cloud condensate variability observed
776 at the Atmospheric Radiation Measurement Sites, *Quarterly Journal of the Royal
777 Meteorological Society*, 142(697), 1605–1617, doi:10.1002/qj.2783, 2016.

778 Barker, H. W.: A Parameterization for Computing Grid-Averaged Solar Fluxes for
779 Inhomogeneous Marine Boundary Layer Clouds. Part I: Methodology and Homogeneous Biases,
780 *J. Atmos. Sci.*, 53(16), 2289–2303, doi:10.1175/1520-0469(1996)053<2289:APFCGA>2.0.CO;2,
781 1996.

782 Barker, H. W., Wiellicki, B. A. and Parker, L.: A Parameterization for Computing Grid-Averaged
783 Solar Fluxes for Inhomogeneous Marine Boundary Layer Clouds. Part II: Validation Using
784 Satellite Data, [http://dx.doi.org/10.1175/1520-0469\(1996\)053<2304:APFCGA>2.0.CO;2](http://dx.doi.org/10.1175/1520-0469(1996)053<2304:APFCGA>2.0.CO;2), 53(16),
785 2304–2316 [online] Available from: [http://journals.ametsoc.org/doi/pdf/10.1175/1520-
786 0469%281996%29053%3C2304%3AAPFCGA%3E2.0.CO%3B2](http://journals.ametsoc.org/doi/pdf/10.1175/1520-0469%281996%29053%3C2304%3AAPFCGA%3E2.0.CO%3B2), 1996.

787 Bennartz, R.: Global assessment of marine boundary layer cloud droplet number concentration
788 from satellite, *Journal of Geophysical Research-Atmospheres*, 2007.

789 Bennartz, R. and Rausch, J.: Global and regional estimates of warm cloud droplet number
790 concentration based on 13 years of AQUA-MODIS observations, *Atmospheric Chemistry and
791 Physics*, 1–32, doi:10.5194/acp-2016-1130, 2017.

792 Bogenschutz, P. A., Gettelman, A., Hannay, C., Larson, V. E., Neale, R. B., Craig, C. and Chen, C.-
793 C.: The Path to CAM6: Coupled Simulations with CAM5.4 and CAM5.5, *Geosci. Model Dev.*, 1–
794 38, doi:10.5194/gmd-2017-129, 2017.

795 Bogenschutz, P. A., Gettelman, A., MORRISON, H., Larson, V. E., Craig, C. and Schanen, D. P.:
796 Higher-Order Turbulence Closure and Its Impact on Climate Simulations in the Community
797 Atmosphere Model, *J. Climate*, 26(23), 9655–9676, doi:10.1175/JCLI-D-13-00075.1, 2013.

798 Bony, S. and Dufresne, J.-L.: Marine boundary layer clouds at the heart of tropical cloud
799 feedback uncertainties in climate models, *Geophysical Research Letters*, 32(20), L20806,
800 doi:10.1029/2005GL023851, 2005.

801 Boutle, I. A., Abel, S. J., Hill, P. G. and Morcrette, C. J.: Spatial variability of liquid cloud and rain:
802 observations and microphysical effects, *Quarterly Journal of the Royal Meteorological Society*,
803 140(679), 583–594, doi:10.1002/qj.2140, 2014.

804 Cahalan, R. F., Ridgway, W., Wiscombe, W. J., Bell, T. L. and Snider, J. B.: The Albedo of Fractal
805 Stratocumulus Clouds, *J. Atmos. Sci.*, 51(16), 2434–2455, doi:10.1175/1520-
806 0469(1994)051<2434:TAOFSC>2.0.CO;2, 1994.

807 Cho, H. M., Zhang, Z., Meyer, K., Lebsock, M., Platnick, S., Ackerman, A. S., Di Girolamo, L., C
808 Labonnote, L., Cornet, C., Riedi, J. and Holz, R. E.: Frequency and causes of failed MODIS cloud
809 property retrievals for liquid phase clouds over global oceans, *Journal of Geophysical Research-*
810 *Atmospheres*, 120(9), 2015JD023161–n/a, doi:10.1002/2015JD023161, 2015.

811 Considine, G., Curry, J. A. and Wielicki, B.: Modeling cloud fraction and horizontal variability in
812 marine boundary layer clouds, *J. Geophys. Res.*, 102(D12), 13517–13525, 1997.

813 Eyring, V., Bony, S., Meehl, G. A., Senior, C. A., Stevens, B., Stouffer, R. J. and Taylor, K. E.:
814 Overview of the Coupled Model Intercomparison Project Phase 6 (CMIP6) experimental design
815 and organization, *Geosci. Model Dev.*, 9(5), 1937–1958, doi:10.5194/gmd-9-1937-2016, 2016.

816 Golaz, J.-C., Larson, V. E. and Cotton, W. R.: A PDF-Based Model for Boundary Layer Clouds. Part
817 I: Method and Model Description, *JAS*, 59(24), 3540–3551, doi:10.1175/1520-
818 0469(2002)059<3540:APBMFB>2.0.CO;2, 2002a.

819 Golaz, J.-C., Larson, V. E. and Cotton, W. R.: A PDF-Based Model for Boundary Layer Clouds. Part
820 II: Model Results, [http://dx.doi.org/10.1175/1520-0469\(2002\)059<3552:APBMFB>2.0.CO;2](http://dx.doi.org/10.1175/1520-0469(2002)059<3552:APBMFB>2.0.CO;2),
821 59(24), 3552–3571, doi:10.1175/1520-0469(2002)059<3552:APBMFB>2.0.CO;2, 2002b.

822 Griffin, B. M. and Larson, V. E.: Analytic upscaling of a local microphysics scheme. Part II:
823 Simulations, *Quarterly Journal of the Royal Meteorological Society*, 139(670), 58–69,
824 doi:10.1002/qj.1966, 2013.

825 Grosvenor, D. P. and Wood, R.: The effect of solar zenith angle on MODIS cloud optical and
826 microphysical retrievals within marine liquid water clouds, *Atmospheric Chemistry and Physics*,
827 14(14), 7291–7321, doi:10.5194/acpd-14-303-2014, 2014.

828 Grosvenor, D. P., Sourdeval, O., Zuidema, P., Ackerman, A., Alexandrov, M. D., Bennartz, R.,
829 Cairns, B., Chiu, C., Christensen, M., Diamond, M., Feingold, G., Fridlind, A., Hunerbein, A., Knist,
830 C., Kollias, P., Marshak, A., McCoy, D., Merk, D., Painemal, D., Rausch, J., Rosenfeld, D.,
831 Russchenberg, H., Seifert, P., Sinclair, K., Stier, P., vanDiedenhoven, B., Wendisch, M., Werner,
832 F., Wood, R., Zhang, Z. and Quaas, J.: **Remote sensing of droplet number concentration in**
833 **warm clouds: A review of the current state of knowledge and perspectives**, *Reviews of*
834 *Geophysics*, (in review), 2018.

835 Gultepe, I. and Isaac, G. A.: Aircraft observations of cloud droplet number concentration:
836 Implications for climate studies, *Quarterly Journal of the Royal Meteorological Society*,
837 130(602), 2377–2390, doi:10.1256/qj.03.120, 2004.

838 Guo, H., Golaz, J. C., Donner, L. J., Larson, V. E., Schanen, D. P. and Griffin, B. M.: Multi-variate
839 probability density functions with dynamics for cloud droplet activation in large-scale models:
840 single column tests, *Geosci. Model Dev.*, 3(2), 475–486, doi:10.5194/gmd-3-475-2010, 2010.

841 Guo, H., Golaz, J. C., Donner, L. J., Wyman, B., Zhao, M. and Ginoux, P.: CLUBB as a unified cloud
842 parameterization: Opportunities and challenges, *Geophysical Research Letters*, 42(11), 4540–
843 4547, doi:10.1002/2015GL063672, 2015.

844 Guo, Z., Wang, M., Qian, Y., Larson, V. E., Ghan, S., Ovchinnikov, M., Bogenschutz, P. A., Zhao,
845 C., Lin, G. and Zhou, T.: A sensitivity analysis of cloud properties to CLUBB parameters in the
846 single-column Community Atmosphere Model (SCAM5), *J. Adv. Model. Earth Syst.*, 6(3), 829–
847 858, doi:10.1002/2014MS000315, 2014.

848 Hill, P. G., Morcrette, C. J. and Boutle, I. A.: A regime-dependent parametrization of subgrid-
849 scale cloud water content variability, *Quarterly Journal of the Royal Meteorological Society*,
850 141(691), 1975–1986, doi:10.1002/qj.2506, 2015.

851 Kawai, H. and Teixeira, J.: Probability Density Functions of Liquid Water Path and Cloud Amount
852 of Marine Boundary Layer Clouds: Geographical and Seasonal Variations and Controlling
853 Meteorological Factors, <http://dx.doi.org/10.1175/2009JCLI3070.1>, 23(8), 2079–2092,
854 doi:10.1175/2009JCLI3070.1, 2010.

855 Khairoutdinov, M. and Kogan, Y.: A New Cloud Physics Parameterization in a Large-Eddy
856 Simulation Model of Marine Stratocumulus, *Mon. Wea. Rev.*, 128(1), 229–243 [online] Available
857 from: [http://journals.ametsoc.org/doi/abs/10.1175/1520-](http://journals.ametsoc.org/doi/abs/10.1175/1520-0493(2000)128%3C0229%3AANCPPI%3E2.0.CO%3B2)
858 [0493\(2000\)128%3C0229%3AANCPPI%3E2.0.CO%3B2](http://journals.ametsoc.org/doi/abs/10.1175/1520-0493(2000)128%3C0229%3AANCPPI%3E2.0.CO%3B2), 2000.

859 Klein, S. and Hartmann, D.: The seasonal cycle of low stratiform clouds, *Journal of Climate*, 6(8),
860 1587–1606, 1993.

861 Kubar, T. L., Stephens, G. L., Lebsock, M., Larson, V. E. and Bogenschutz, P. A.: Regional
862 Assessments of Low Clouds against Large-Scale Stability in CAM5 and CAM-CLUBB Using MODIS
863 and ERA-Interim Reanalysis Data, *J. Climate*, 28(4), 1685–1706, doi:10.1175/JCLI-D-14-00184.1,
864 2014.

865 Larson, V. E. and Griffin, B. M.: Analytic upscaling of a local microphysics scheme. Part I:
866 Derivation, *Quarterly Journal of the Royal Meteorological Society*, 139(670), 46–57,
867 doi:10.1002/qj.1967, 2013.

868 Larson, V. E., Golaz, J.-C. and Cotton, W. R.: Small-Scale and Mesoscale Variability in Cloudy
869 Boundary Layers: Joint Probability Density Functions, *J. Atmos. Sci.*, 59(24), 3519–3539,
870 doi:10.1175/1520-0469(2002)059<3519:SSAMVI>2.0.CO;2, 2002.

871 Larson, V. E., Wood, R., Field, P. R., Golaz, J.-C., Vonder Haar, T. H. and Cotton, W. R.: Systematic
872 Biases in the Microphysics and Thermodynamics of Numerical Models That Ignore Subgrid-Scale

873 Variability, *J. Atmos. Sci.*, 58(9), 1117–1128, doi:10.1175/1520-
874 0469(2001)058<1117:SBITMA>2.0.CO;2, 2001.

875 Lebsock, M. D., L'Ecuyer, T. S. and Stephens, G. L.: Detecting the Ratio of Rain and Cloud Water
876 in Low-Latitude Shallow Marine Clouds, *Journal of Applied Meteorology and Climatology*, 50(2),
877 419–432, doi:10.1175/2010JAMC2494.1, 2011.

878 Lebsock, M., MORRISON, H. and Gettelman, A.: Microphysical implications of cloud-
879 precipitation covariance derived from satellite remote sensing, *Journal of Geophysical*
880 *Research-Atmospheres*, 118(12), 6521–6533, doi:10.1002/jgrd.50347, 2013.

881 Lee, S., Kahn, B. H. and Teixeira, J.: Characterization of cloud liquid water content distributions
882 from CloudSat, *J. Geophys. Res.*, 115(D20), D00A23, doi:10.1029/2009JD013272, 2010.

883 McCoy, D. T., Bender, F. A. M., Grosvenor, D. P., Mohrmann, J. K., Hartmann, D. L., Wood, R.
884 and Field, P. R.: Predicting decadal trends in cloud droplet number concentration using
885 reanalysis and satellite data, *Atmospheric Chemistry and Physics*, 1–21, 2017a.

886 McCoy, D. T., Bender, F. A. M., Mohrmann, J. K. C., Hartmann, D. L., Wood, R. and Grosvenor, D.
887 P.: The global aerosol-cloud first indirect effect estimated using MODIS, MERRA, and AeroCom,
888 *Journal of Geophysical Research-Atmospheres*, 122(3), 1779–1796, doi:10.1002/2016JD026141,
889 2017b.

890 Menzel, P., Frey, R., Baum, B. and Zhang, H.: Cloud Top Properties and Cloud Phase Algorithm
891 Theoretical Basis Document. 2006.

892 Menzel, W., Smith, W. and Stewart, T.: Improved Cloud Motion Wind Vector and Altitude
893 Assignment Using VAS, *Journal of Applied Meteorology*, 22(3), 377–384, 1983.

894 Morales, R. and Nenes, A.: Characteristic updrafts for computing distribution-averaged cloud
895 droplet number and stratocumulus cloud properties, *J. Geophys. Res.*, 115(D18), 1227, 2010.

896 Morrison, H. and Gettelman, A.: A new two-moment bulk stratiform cloud microphysics scheme
897 in the Community Atmosphere Model, version 3 (CAM3). Part I: Description and numerical
898 tests, *Journal of Climate*, 2008.

899 Nakajima, T. and King, M. D.: Determination of the Optical Thickness and Effective Particle
900 Radius of Clouds from Reflected Solar Radiation Measurements. Part I: Theory, *J. Atmos. Sci.*,
901 47(15), 1878–1893, doi:10.1175/1520-0469(1990)047<1878:DOTOTA>2.0.CO;2, 1990.

902 Nam, C., Bony, S., Dufresne, J. L. and Chepfer, H.: The “too few, too bright” tropical low-cloud
903 problem in CMIP5 models, *Geophysical Research ...*, doi:10.1029/2012GL053421, 2012.

904 O, K. T., Wood, R. and Tseng, H. H.: Deeper, Precipitating PBLs Associated With Optically Thin
905 Veil Clouds in the Sc-Cu Transition, *Geophysical Research Letters*, 27(1), 1, 2018.

- 906 Oreopoulos, L. and Barker, H. W.: Accounting for subgrid-scale cloud variability in a multi-layer
907 1d solar radiative transfer algorithm, *Quarterly Journal of the Royal Meteorological Society*,
908 125(553), 301–330, doi:10.1002/qj.49712555316, 1999.
- 909 Oreopoulos, L. and Cahalan, R. F.: Cloud Inhomogeneity from MODIS, *Journal of Climate*,
910 18(23), 5110–5124, doi:10.1175/JCLI3591.1, 2005.
- 911 Oreopoulos, L. and Davies, R.: Plane Parallel Albedo Biases from Satellite Observations. Part I:
912 Dependence on Resolution and Other Factors, *Journal of Climate*, 11(5), 919–932, 1998a.
- 913 Oreopoulos, L. and Davies, R.: Plane Parallel Albedo Biases from Satellite Observations. Part II:
914 Parameterizations for Bias Removal, *J. Climate*, 11(5), 933–944, 1998b.
- 915 Pincus, R. and Klein, S. A.: Unresolved spatial variability and microphysical process rates in
916 large-scale models, *J. Geophys. Res.*, 105(D22), 27059–27065, doi:10.1029/2000JD900504,
917 2000.
- 918 Platnick, S., King, M. D., Ackerman, S. A., Menzel, W. P., Baum, B. A., Riédi, J. C. and Frey, R. A.:
919 The MODIS cloud products: algorithms and examples from Terra, *IEEE TRANSACTIONS ON*
920 *GEOSCIENCE AND REMOTE SENSING*, 41(2), 459–473, doi:10.1109/TGRS.2002.808301, 2003.
- 921 Platnick, S., Meyer, K. G., King, M. D., Wind, G., Amarasinghe, N., Marchant, B., Arnold, G. T.,
922 Zhang, Z., Hubanks, P. A., Holz, R. E., Yang, P., Ridgway, W. L. and Riedi, J.: The MODIS Cloud
923 Optical and Microphysical Products: Collection 6 Updates and Examples From Terra and Aqua,
924 *IEEE TRANSACTIONS ON GEOSCIENCE AND REMOTE SENSING*, 55(1), 502–525,
925 doi:10.1109/TGRS.2016.2610522, 2017.
- 926 Pruppacher, H. R. and Klett, J. D.: *Microphysics of Clouds and Precipitation: With an*
927 *Introduction to Cloud Chemistry and Cloud Electricity*, 954 pp. 1997.
- 928 Randall, D., Khairoutdinov, M., Arakawa, A. and Grabowski, W.: Breaking the Cloud
929 Parameterization Deadlock, *Bulletin of the American Meteorological Society*, 84(11), 1547–
930 1564, doi:10.1175/BAMS-84-11-1547, 2003.
- 931 Seethala, C. and Horváth, Á.: Global assessment of AMSR-E and MODIS cloud liquid water path
932 retrievals in warm oceanic clouds, *J Geophys Res*, 115(D13), D13202, 2010.
- 933 Soden, B. and Held, I.: An assessment of climate feedbacks in coupled ocean–atmosphere
934 models, *Journal of Climate*, 2006.
- 935 Song, H., Song, H., Zhang, Z., Ma, P.-L., Ghan, S. J. and Wang, M.: An Evaluation of Marine
936 Boundary Layer Cloud Property Simulations in the Community Atmosphere Model Using
937 Satellite Observations: Conventional Subgrid Parameterization versus CLUBB, *Journal of*
938 *Climate*, 31(6), 2299–2320, doi:10.1175/JCLI-D-17-0277.1, 2018a.

939 Song, H., Zhang, Z., Ma, P.-L., Ghan, S. J. and Wang, M.: An Evaluation of Marine Boundary Layer
940 Cloud Property Simulations in the Community Atmosphere Model Using Satellite Observations:
941 Conventional Subgrid Parameterization versus CLUBB, *Journal of Climate*, 31(6), 2299–2320,
942 doi:10.1175/JCLI-D-17-0277.1, 2018b.

943 Sušelj, K., Teixeira, J. and Chung, D.: A Unified Model for Moist Convective Boundary Layers
944 Based on a Stochastic Eddy-Diffusivity/Mass-Flux Parameterization, *J. Atmos. Sci.*, 70(7), 1929–
945 1953, doi:10.1175/JAS-D-12-0106.1, 2013.

946 Takahashi, H., Lebsock, M., Suzuki, K., Stephens, G. and Wang, M.: An investigation of
947 microphysics and subgrid-scale variability in warm-rain clouds using the A-Train observations
948 and a multiscale modeling framework, *Journal of Geophysical Research-Atmospheres*, 138(669),
949 2151, 2017.

950 Thayer-Calder, K., Gettelman, A., Craig, C., Goldhaber, S., Bogenschutz, P. A., Chen, C. C.,
951 Morrison, H., Höft, J., Raut, E., Griffin, B. M., Weber, J. K., Larson, V. E., Wyant, M. C., Wang, M.,
952 Guo, Z. and Ghan, S. J.: A unified parameterization of clouds and turbulence using CLUBB and
953 subcolumns in the Community Atmosphere Model, *Geosci. Model Dev.*, 8(12), 3801–3821,
954 doi:10.5194/gmd-8-3801-2015, 2015.

955 Trenberth, K. E., Fasullo, J. T. and Kiehl, J.: Earth's Global Energy Budget, *Bull. Amer. Meteor.*
956 *Soc.*, 90(3), 311–323, doi:10.1175/2008BAMS2634.1, 2009.

957 Wang, M., Larson, V. E., Ghan, S., Ovchinnikov, M., Schanen, D. P., Xiao, H., Liu, X., Rasch, P. and
958 Guo, Z.: A multiscale modeling framework model (superparameterized CAM5) with a higher-
959 order turbulence closure: Model description and low-cloud simulations, *J. Adv. Model. Earth*
960 *Syst.*, n/a–n/a, doi:10.1002/2014MS000375, 2015.

961 Wood, R. and Hartmann, D.: Spatial variability of liquid water path in marine low cloud: The
962 importance of mesoscale cellular convection, *Journal of Climate*, 2006.

963 Wood, R., O, K.-T., Bretherton, C. S., Mohrmann, J., ALBRECHT, B. A., Zuidema, P., Ghate, V.,
964 Schwartz, C., Eloranta, E., Glienke, S., Shaw, R., Fugal, J. and Minnis, P.: Ultraclean layers and
965 optically thin clouds in the stratocumulus to cumulus transition: part I. Observations, *J. Atmos.*
966 *Sci.*, JAS–D–17–0213.1, doi:10.1175/JAS-D-17-0213.1, 2018.

967 Xie, X. and Zhang, M.: Scale-aware parameterization of liquid cloud inhomogeneity and its
968 impact on simulated climate in CESM, *Journal of Geophysical Research-Atmospheres*, 120(16),
969 8359–8371, doi:10.1002/2015JD023565, 2015.

970 Zhang, Z. and Platnick, S.: An assessment of differences between cloud effective particle radius
971 retrievals for marine water clouds from three MODIS spectral bands, *J Geophys Res*, 116(D20),
972 D20215, doi:10.1029/2011JD016216, 2011.

973 Zhang, Z., Ackerman, A. S., Feingold, G., Platnick, S., Pincus, R. and Xue, H.: Effects of cloud
974 horizontal inhomogeneity and drizzle on remote sensing of cloud droplet effective radius: Case

975 studies based on large-eddy simulations, J Geophys Res, 117(D19), D19208–,
976 doi:10.1029/2012JD017655, 2012.

977 Zhang, Z., Werner, F., Cho, H. M., Wind, G., Platnick, S., Ackerman, A. S., Di Girolamo, L.,
978 Marshak, A. and Meyer, K.: A framework based on 2-D Taylor expansion for quantifying the
979 impacts of sub-pixel reflectance variance and covariance on cloud optical thickness and
980 effective radius retrievals based on the bi-spectral method, Journal of Geophysical Research-
981 Atmospheres, 2016JD024837, doi:10.1002/2016JD024837, 2016.

982