Exploring aerosol cloud interaction using VOCALS-REx aircraft measurements

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Abstract. In situ aircraft measurements obtained during the VAMOS Ocean-Cloud-Atmosphere-Land Study-Regional Experiment (VOCALS-REx) field campaign are analysed to study the aerosol-cloud interactions in the stratocumulus clouds over the southeast Pacific Ocean, with a focus on three understudied topics (separation of aerosol effects from dynamic effects, dispersion effects, and turbulent entrainment mixing processes). Our analysis suggests that an increase in aerosol concentration tends to simultaneously increase both cloud droplet number concentration (N_d) and relative dispersion (ε), while an increase in vertical velocity (w) often increases N_d but decreases ε . After constraining the differences of cloud dynamics, the positive correlation between ε and N_d becomes stronger, implying that perturbations of w could weaken the aerosol influence on ε , and hence result in an underestimation of dispersion effect. A comparative analysis of the difference of cloud microphysical properties between entrainment and non-entrainment zones suggests that the entrainment-mixing mechanism is predominantly extreme inhomogeneous in the stratocumulus that capped by a sharp inversion, whereby the variation of liquid water content (25 %) is similar to that of N_d (29 %) and the droplet size remains approximately constant. In entrainment zone, drier air entrained from the top induces less cloud droplets with respect to total in-cloud particles (0.56 \pm 0.22) than the case in non-entrainment zone (0.73 \pm 0.13) by promoting cloud droplet evaporation. This study is helpful in reducing uncertainties in dispersion effects and entrainment mixing for stratocumulus, and the results of this study may benefit cloud parameterizations in global climate models to more accurately assess aerosol indirect effects.

1 Introduction

Stratocumulus clouds play a key role in the radiative energy budget of the Earth by reflecting incoming shortwave radiation and thus cool the planet surface and offset the warming by greenhouse gases (Hartmann et al., 1992). Stratocumulus clouds are susceptible to aerosols, i.e., aerosol indirect effect (Twomey, 1974; Albrecht, 1989), which currently remains large uncertainties (Lohmann and Feichter, 2005; Chen and Penner, 2005; Carslaw et al., 2013; McCoy et al., 2017).

Globally, marine stratocumulus clouds overlaying the southeast Pacific Ocean (SEP) are the largest and most persistent clouds (Klein and Hartmann, 1993; Bretherton et al., 2004). Sources of anthropogenic aerosols from the Chilean and Peruvian coasts, in contrast with the relatively clean air masses from the Pacific Ocean, make the SEP an ideal region to explore the interaction between aerosols and stratocumulus cloud topped boundary layers. The cloud properties from satellite retrievals exhibit a gradient off the Northern Chile shore. For example, the cloud droplet number concentration decreased from 160 to 40 cm⁻³ (George and Wood, 2010), and the cloud droplet effective radius increased from 8 to 14 µm from the coast to approximately 1000 km offshore (Wood et al., 2006). This gradient is plausibly attributable to anthropogenic aerosols near the coast. Huneeus et al. (2006) found that during easterly wind events, sulfate increased by one order of magnitude over SEP, which resulted in a 1.6 to 2-fold increase in cloud droplet number concentration. Based on observations from satellites and cruises, Wood et al. (2008) suggested that open cellular convection within an overcast stratocumulus is associated with reduced aerosol concentration and air masses not passing through the Chilean coast, which further confirms the impact of aerosols on stratocumulus over the SEP. However, it is difficult to establish the generality of previous studies based on satellite remote sensing due to the absence of in situ observations that provide vertical profiles of cloud and aerosol and detailed in-cloud processes.

The VAMOS (Variability of the American Monsoons) Ocean-Cloud-Atmosphere-Land Study-Regional Experiment (VOCALS-REx), which includes multiple aircraft missions, ship and land-based measurements, took place in the region extending from the near-coast of northern Chile and southern Peru to the remote ocean in the SEP during October–November 2008 (Wood et al., 2011). The data collected during this campaign were examined to investigate the properties of aerosols, clouds, and marine boundary layer over the SEP. For instance, Bretherton et al. (2010) found that the boundary layer was shallow and fairly well mixed near shore but deeper and decoupled offshore. Twohy et al. (2013) found that the clouds near the shore exhibited higher aerosol concentrations, greater droplet concentrations, smaller droplet sizes and smaller liquid water path (*LWP*), and suggested a combination of anthropogenic aerosols and physically thinner clouds near the shore. However, Zheng et al. (2010) found an increase in *LWP* with cloud condensation nuclei (CCN) concentrations under the similar meteorological conditions. Additionally, chemical components and sources of aerosols during VOCALS-REx campaign have been discussed in several studies (Chand et al., 2010; Hawkins et al., 2010; Allen et al., 2011; Twohy et al., 2013; Lee et al., 2014).

Although these studies have improved our understanding of some aspects related to aerosol, cloud and boundary layer properties over SEP, several important factors remain understudied or unexplored. First, the aerosol effect on clouds is often intertwined with the effects of other factors, especially meteorological conditions (Fan et al., 2009; Koren et al., 2010). Currently, the impact of aerosols on the shape of cloud droplet size spectrum (i.e., dispersion effect) is reported to remain large uncertainty. The observed correlations between relative dispersion (ε) and N_d can be positive, negative, or not evident (Table 1),

which could be largely attributable to the coincidentally changing cloud dynamics. Thus, it is necessary to isolate the response of ε to aerosol perturbations from meteorological effects, which to our knowledge, has not received adequate attention in many previous studies. Second, applying different assumptions to the entrainment mixing mechanism can have a significant impact on cloud albedo (Grabowski, 2006; Chosson et al., 2007; Slawinska et al., 2008). Additionally, more recent studies suggested that entrainment mixing may be a possible physical interpretation for the observed anti-Twomey effect (Ma et al., 2018; Jia et al., 2019). However, it remains unclear whether the entrainment-mixing mechanism is predominantly homogeneous, inhomogeneous, or in between (Andrejczuk et al., 2009; Lehmann et al., 2009). By using cloud observations obtained from G-1 aircraft during VOCALS-REx, Yum et al. (2015) found both homogeneous and inhomogeneous mixing in their analysis and attributed it to the uncertainty in the methods they used. Uncertainty in the entrainment mixing mechanism could lead to the inaccurate assessment of aerosol indirect effects. Thus, more attention should be paid to this topic.

Based on the useful information on the microphysical properties of aerosols and clouds provided by previous studies, in this study, we conduct additional explorations regarding aerosol-cloud interactions over the SEP by employing in situ aircraft data collected by CIRPAS Twin Otter aircraft during VOCALS-REx, which include the following: (a) investigating the controlling factors of cloud droplet formation (e.g., cloud dynamics and aerosols), (b) evaluating the dispersion effect under relatively constant cloud dynamical conditions, and (c) re-examining the entrainment-mixing mechanism by using a different approach from that of Yum et al. (2015).

2 Data and method

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2.1 Aircraft Data

The Twin Otter operated by the Center for Interdisciplinary Remotely Piloted Aircraft Studies (CIRPAS) was aimed to observe aerosol and cloud microphysics, and turbulence near Point Alpha (20° S, 72° W) off the coast of Northern Chile from 16 October to 13 November 2008. A total of 19 flights were carried out, each of which conducting about 3 hours of sampling at Point Alpha and including several soundings and horizontal legs near the ocean surface, below the cloud, near the cloud base, within the cloud, near the cloud top, and above the cloud (Fig. 1). Since all flight tracks are similar, only one track (Oct. 18) is shown in Fig. 1. As cloud and aerosol probe measurements failed during the flight on 5 November and drizzle processes occurred on the flights on 1 November and 2 November, only the observations from other 16 non-drizzling flights are included in this paper.

Both the aerosols below and above clouds and the interstitial aerosols in-cloud were obtained by Passive Cavity Aerosol Spectrometer Probe (PCASP-100), which counted and sized particles from $0.1-2.0 \mu m$ diameter with 20 bins (Zheng et al., 2011; Cai et al., 2013; Twohy et al., 2013). The CCN number concentration was observed by the CCN Spectrometer at a supersaturation of 0.2 % and 0.5% respectively. The cloud data include cloud droplet number concentration (N_d , size range:

2.07–40.2 μ m with 20 bins) from the Cloud, Aerosol and Precipitation probe (CAS), effective radius of cloud droplets (R_e), and liquid water content (LWC) from the PVM-100 probe (Gerber et al.,1994). All data sets used in this study are at a frequency of 1 Hz. Zheng et al. (2011) pointed out that uncertainties of aerosols and clouds measured by these probes are within 15 %. More detailed information about the instruments and measurements on board the CIRPAS Twin Otter aircraft during VOCALS-REx can be found in Zheng et al. (2010) and Wood et al. (2011).

2.2 Data processing

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In this study, the data collected near the land during both take-off and landing are removed to ensure that only the measurements close to Point Alpha (20° S, 72° W) are analysed. Only the clouds with LWC > 0.05 g m⁻³ and $N_d > 15$ cm⁻³ are selected for analysis. We averaged the CCN number concentrations during the legs within 200 m above the cloud top to obtain the average above-cloud CCN, and within 200 m below the cloud base to obtain the mean sub-cloud CCN. During the study period, the CCN Spectrometer constantly measured CCN at a supersaturation of 0.2 % except on the first four flights at a supersaturation of 0.5 %. In order to have a consistent comparison between all flights, we adopted the method by Zheng et al. (2011) to adjust the CCN concentration from supersaturation of 0.5 % to 0.2 % on the first four flights. The effective radius of aerosol particle is calculated from the PCASP-measured aerosol size distributions based on following equation:

$$Ra = \sum n_i r_i^3 / \sum n_i r_i^2 \tag{1}$$

where n_i is the aerosol number concentration in the *i*th bin of PCASP, and r_i represents the arithmetic mean radius of *i*th bin.

To investigate the impact of the entrainment mixing processes on cloud properties and aerosol-cloud interactions, we defined entrainment zone and non-entrainment zone, respectively. Gerber et al. (2005) showed that, in the marine stratocumulus, entrainment occurs when *LWC* begins to decrease from the bottom of the cloud. In this manuscript, entrainment and non-entrainment zone are thus defined as the regions within 20 m above and below the height of maximal *LWC*, respectively. Given that the two zones are both thin layers, there is little difference in the dynamical and thermos-dynamical conditions. It is therefore assumed that the difference of cloud microphysical characteristics between the two zones is only caused by entrainment.

3 Results

3.1 Vertical profiles of aerosol, cloud and meteorological variables

Figure 2 shows the vertical profiles of temperature (a), relative humidity (b), liquid water content (c), cloud droplet effective radius (d), cloud droplet number concentration (e), aerosol effective radius (f), and the ratio of CCN to condensation nuclei (g) during the 16 flights. Note that the vertical altitude is normalized by the inversion height (z_i) defined as the height of the maximum vertical gradient of liquid water potential temperature (θ_L) (Zheng et al., 2011). The normalization minimizes the

effect of the variation of z_i between flights, allowing for better exploration of the average boundary layer (BL) structure during VOCALS-REx.

As shown in Fig. 2a, temperature (T) decreased sharply with height within the BL, which is close to the dry adiabatic lapse rate. A strong inversion occurred at the top of the BL, with an average temperature increase of approximately 10°C. Due to the decrease of T with height and the nearly constant water vapor mixing ratio within the strong mixing BL, the relative humidity (RH) increased rapidly with increasing height (Fig. 2b). T and RH reached the minimum and maximum, respectively, when z/z_i was close to 0.9. Near the top of the BL $(0.9 \le z/z_i \le 1.0)$, the entrainment of the dry and warm air from the free atmosphere aloft resulted in a slight increase in T and a slight decrease in RH. As z/z_i varied from 1 to 1.1, T increased from 11 to approximately 18°C, and RH rapidly decreased to approximately 16 % (Fig. 2a, b). The vertical profiles of T and RH are overall consistent with the observations of other marine stratocumulus clouds (Martinet et al., 1994; Keil and Haywood, 2003). For the cloud properties, an average of all profiles that normalized by z_i only may be insufficient to indicate the vertical variation of clouds due to different cloud base heights of each profile. Thus, the average profiles are not shown in Fig. 2c, d, and e, and the vertical variation of cloud properties can easily be seen from the single profile. Figure 2c shows that the LWC first increased with height from the cloud base, reached the maximum at z/zi = 0.9 and then decreased with further increasing height when $0.9 \le z/zi \le 1.0$. The profile of R_e is similar to that of LWC (Fig. 2d). The profile of N_d remains relatively constant, with a slight increase and decrease near the base and top, respectively (Fig. S1), which is consistent with the results from other VOCALS-REx observations (Painemal and Zuidema, 2011). It is interesting to note that the effective radius of aerosol particles (R_a) below cloud is larger than that above cloud, which is probably attributable to the differences in aerosol sources and aerosol properties (e.g., chemical composition) (Fig. 2f). The profile of CCN/CN is similar to that of R_a (Fig. 2g), suggesting that aerosols with large sizes are more likely to become CCN (Dusek et al., 2006; Zhang et al., 2011). Larger R_a and CCN/CN values are also found in polluted case than clean cases.

3.2 Relationships between aerosol and cloud properties

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The relationships between aerosol and cloud properties are essential to understanding and evaluating aerosol-cloud interactions. Most studies based on satellite data have employed aerosol optical depth or aerosol index as a proxy for CCN number concentration to investigate aerosol-cloud interactions (Koren et al., 2005, 2010; Su et al., 2010; Tang et al., 2014; Ma et al., 2014, 2018; Wang et al., 2014, 2015; Saponaro et al., 2017). However, not all aerosols in the vertical column are actually involved in cloud formation; thus, this assumption is questionable, especially when the cloud layer is decoupled from the aerosol layer. For example, a few studies have shown that aerosols have little effect on cloud properties when aerosol and cloud layers are clearly separated (Costantino and Bréon, 2010, 2013; Liu et al., 2017). To further investigate this issue, the CCN number concentration both below cloud (*sub-CCN*) and above cloud (*abv-CCN*) are examined for their impacts on the cloud properties.

Figure 3 shows the relationships between sub-CCN and cloud properties during all 16 non-drizzling flights. The red dots denote the ten flights with typical well-mixed boundary layer (BL). These flights also shared similar meteorological conditions, such as inversion heights and the jump of potential temperature and total water mixing ratio across the inversion (Zheng et al., 2010), and thus can be used to isolate the response of cloud properties to aerosol perturbations. The blue dots represent the other cases in which the conditions except typical well-mixed BL, such as strong wind shear within the BL, moist layers above clouds, strong decoupled BL and so on, are involved (Table 2). For the cases with typical well-mixed BL, both LWC (Fig. 3a) and N_d (Fig. 3b) exhibited positive correlations with sub-CCN, with correlation coefficients of 0.60 and 0.79, respectively. It is worth highlighting that the similar increases of N_d and LWC led to R_e having no evident correlation with sub-CCN (Fig. 3c), as expected from the conventional first aerosol indirect effect whereby a constant LWC is assumed. For the other cases (blue dots), the sub-CCN impacts on the cloud properties were not evident due to the large differences in the meteorological conditions and the BL structure.

Compared to sub-cloud CCN, the influence of above-cloud CCN on cloud properties is very weak, even for the cases with typical well-mixed BL. The absolute values of the correlation coefficient between the *abv-CCN* and cloud properties are all less than 0.4 (not shown), and none of them passed the significance test ($\alpha = 0.05$). In this study, the above-cloud aerosol number concentration is very low ($129 \pm 60 \text{ cm}^{-3}$), and the inversion capping the cloud top is extremely strong, which weakens the aerosol mixing with cloud layer and hence the aerosol effects on cloud properties. Some previous studies based on aircraft observations for stratocumulus clouds also found that N_d exhibits a significantly positive correlation with *sub-CCN*, but no correlation with *abv-CCN* (Martin et al., 1994; Hudson et al., 2010; Hegg et al., 2012).

Figure 4 contrasts the relationships of N_d (a) and R_e (b) as functions of LWC between the most polluted (Oct. 19) and cleanest (Nov. 09) cases with aerosol concentrations of 647 ± 60 cm⁻³ and 268 ± 35 cm⁻³, respectively. Also shown are the corresponding power-law fits. Although N_d and R_e both increased with first increasing LWC and then levelled off, there were significant detailed differences between the polluted and clean cases. The polluted case exhibits a steeper increase of N_d with increasing LWC than the clean case when LWC is small, whereas the opposite was true for R_e . The low aerosol concentrations under the clean case inhibit the increase of N_d with LWC (Fig. 4a), which hence promotes the rapid increase of R_e with LWC (Fig. 4b). In contrast, there are enough particles that may potentially activated into cloud droplets under the polluted case; thus, N_d increases rapidly with LWC. As a certain amount of water is shared by large amount particles, the increase of R_e is limited. The result is consistent with the study in Beijing by Zhang et al. (2011), but the difference in cloud formations between the clean and polluted conditions is less evident, which is likely attributable to the much smaller difference in aerosol concentration in this study (approximately 400 cm⁻³) than that in Zhang et al. (2011) (approximately 7000 cm⁻³).

3.3 Cloud droplet formation and its controlling factors

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Sub-cloud CCN is considered to be a good proxy for aerosols entering a cloud. However, during the actual flight, it is difficult to simultaneously collect enough samples of sub-cloud CCN and cloud droplets, which may result in statistical analysis uncertainty. This limitation can be remedied by using the total particle concentration, which equals the aerosol concentration outside the clouds and the sum of the droplet concentration and interstitial aerosol concentration inside the clouds. Interstitial aerosols are particles observed inside clouds that either have never activated into cloud droplets or have been deactivated into aerosols after cloud droplet evaporation. Kleinman et al. (2012) pointed out that the number concentration of interstitial aerosols (N_i) can be obtained either directly from the observation of in-cloud aerosols or indirectly from a number balance between sub-cloud and in-cloud particles. In this study, the interstitial aerosol properties are derived from direct measurements in the cloud. By employing aircraft observations over both land and ocean, Gultepe et al. (1996) found that the difference in the number concentration between the total in-cloud particles $(N_d + N_i)$ measured directly and sub-cloud aerosols is very small. Thus, the total in-cloud particles are assumed to characterize the overall level of in-cloud aerosol concentration before activation. Figure 5 shows an example of the Oct. 18 flight to support this assumption. It is shown that the number concentrations of sub-cloud aerosols and total in-cloud particles are very close, with values of 583 ± 55 cm⁻³ and 567 ± 59 cm⁻³, respectively. Similar results are also found in the other flights. The average ratio of $N_d + N_i$ to the sub-cloud aerosol concentration during all flights is 0.94, which is smaller than the value (1.29) found by Kleinman et al. (2012) based on G-1 aircraft measurements during VOCALS-REx. Therefore, the interstitial aerosol observations in this study are unlikely to be significantly interfered with by factors such as cloud droplet shatter and cloud droplet evaporation due to instrument heating, as discussed by Kleinman et al. (2012), which has the potential to create spurious extra aerosols in-cloud.

The relations between N_d and $N_d + N_i$ during the 16 non-drizzling flights are shown in Fig. 6, where the color represents the in-cloud vertical velocity. All flights exhibited positive correlations between N_d and $N_d + N_i$, representing the aerosol-cloud interaction (IPCC, 2001, 2007, 2013; Hegg et al., 2012). In addition, the effect of dynamical conditions on cloud droplet formation is evident. As shown in Fig. 6, the data are close to the 1:1 line when the vertical velocity is relatively large, namely, the aerosols were almost entirely activated into cloud droplets. However, the data deviate from the 1:1 line when the vertical velocity is small or negative. For example, for all flights, the average ratio of N_d to $N_d + N_i$ with vertical velocity greater than 1 m s⁻¹ is 0.84 ± 0.12 , which is much larger than that with vertical velocity less than -1 m s⁻¹ (0.64 ± 0.14). The regime-dependent behaviour is likely due to the high supersaturation caused by the adiabatic uplift when the vertical velocity is large (Reutter et al., 2009; Chen et al., 2016).

In addition to the dynamical conditions, aerosol microphysical properties such as size distribution and chemical components can also significantly affect activation process (Nenes et al., 2002; Lance et al., 2004; Ervens et al., 2005; Dusek et al., 2006; McFiggans et al., 2006; Zhang et al., 2011; Almeida et al., 2014; Leck and Svensson, 2015). Since part of aerosol

population has activated, it is difficult to obtain information about aerosol size before activation. According to the Köhler theory, larger aerosols have smaller critical supersaturations, and thus, they activate preferentially, suggesting that the effective radius of interstitial aerosols (R_i) is smaller than that of the aerosols before activation. Li et al. (2011) compared the difference in size distribution between interstitial aerosols and aerosols that have been activated to cloud droplets, and found that the peak diameter of the former (0.45 µm) was much smaller than that of the latter (0.8 µm). It can be thus inferred that the size of aerosols activated to cloud droplets, and thus the size of initial aerosols would be larger with an increase in R_i , though the quantitative relationship depends on in-cloud dynamics. Therefore, it is assumed that when compared with the data measured at different sampling locations during flight, the size of the interstitial aerosols can still represent the size of the aerosols before activation to some extent. As indicated in Fig. 7, the larger R_i is, the closer the data are to the 1:1 line, i.e., the higher proportion of cloud droplets in total in-cloud particles ($N_d/(N_d+N_i)$) is. The averaged $N_d/(N_d+N_i)$ for all flights is 0.76 ± 0.13 when R_i is larger than 0.5 µm, but only 0.64 ± 0.23 when R_i is less than 0.25 µm. It is because those aerosols with large sizes are more likely to be activated into cloud droplets. Additionally, as larger aerosol particles form into larger cloud droplets (Twohy et al., 1989, 2013) that are relatively difficult to evaporate, large particles can also inhibit cloud droplet evaporation to a certain extent.

3.4 Dispersion effect

In addition to modulating the cloud droplet number concentration, aerosols also affect the shape of cloud droplet size spectrum (referred to as the "dispersion effect") and thereby affect the cloud albedo (Liu and Daum, 2002). When the dispersion effect is taken into account, the estimated aerosol indirect forcing could be either reduced (Liu and Daum, 2002; Peng and Lohmann, 2003; Kumar et al., 2016; Pandithurai et al., 2012) or enhanced (Ma et al., 2010), i.e., the dispersion effect could act to either offset or enhance the well-known Twomey effect, which mainly depends on the sensitivity of the relative dispersion (ε , the ratio of the standard deviation to the mean radius of the cloud droplet size distribution) to the aerosol number concentration (N_a). However, the dependence of ε on N_a is much less studied and remains even more uncertain than that of N_d . Table 1 summarizes the observed correlations between ε and N_d (or N_a), being positive, negative, or no obvious correlations. The different relationships are indicative of the fact that the effect of aerosol on ε is often intertwined with the effects of other factors, especially cloud dynamical conditions (Pawlowska et al., 2006; Lu et al., 2012). In this section, the relationship between ε and N_d based on in-flight and flight-averaged data are discussed to distinguish the influences of aerosol and cloud dynamics on ε .

Within an individual flight, the aerosol number concentration and chemical components can be assumed to be similar, providing an opportunity to focus on the effect of cloud dynamics to the extent possible. Here, we employ the vertical velocity $(w, \text{ m s}^{-1})$ as a proxy for cloud dynamical conditions. As shown in Fig. 8, the correlations between ε and N_d based on in-flight data are significantly negative during all 16 non-drizzling flights, which is mainly modulated by w, i.e., a larger w corresponds

to a smaller ε but a larger N_d . High supersaturation leads to more cloud droplets to activate and grow to the same size (i.e., narrowing the droplet spectrum) when w is relatively large, but a portion of the cloud droplets may evaporate into smaller sizes and even deactivate into interstitial aerosols when w is small or even negative, resulting in a decrease of N_d and broadening of the droplet spectrum.

It is interesting to see from Table 1 that the correlations between ε and N_d based on in-flight data are generally negative, while the correlations based on the flight-averaged data could be either positive, negative, or even uncorrelated. The latter uncertain relationships may result from variations of the strength of cloud dynamics between flights, which would disrupt or even cancel the real influence of aerosols on relative dispersion (Liu et al., 2006; Peng et al., 2007; Lu et al., 2012). However, many previous studies did not consider the difference in cloud dynamics between flights when correlating ε and N_d , which could result in some degree of overestimation or underestimation of dispersion effect. In this study, the data of all flights were sampled over the same location, i.e., Point Alpha, which can reduce the difference in dynamical conditions caused by variations of horizontal sampling locations. In addition, we also distinguish between the flights of typical mixed BL and others to ensure relatively similar meteorological conditions (see section 3.2). Fig. 9 further shows the probability distribution function of w with mean values and standard deviations for 16 non-drizzling flights. The related statistics are given in Table 2. Except for other cases (crosses; especially Oct. 24, Oct. 29, Nov. 8, and Nov. 13), the difference in the in-cloud dynamics between typical well-mixed BL flights is very small, which confirms the assumption of similar meteorological conditions. As indicated in Fig. 10a, ε and N_d were positively correlated (correlation coefficient of 0.29 and slope of 1.9 \times 10⁻⁴) in the case of the typical well-mixed BL, indicating that increased aerosols concurrently increased ε and N_d . However, the correlation coefficient and slope decrease to 0.11 and 7.7×10^{-5} , respectively, in the all cases (i.e., w is not constrained), implying that the influence of aerosols on the ε - N_d relationship tends to be weaker after intertwining with the effects of cloud dynamics. Although the perturbations of cloud dynamics have been eliminated as much as possible, N_d is still likely determined by both aerosol number concentrations and updraft velocity together. Therefore, a similar statistical analysis is also conducted for sub-cloud CCN. Similar positive correlations between ε and sub-cloud CCN were found, with much improved correlation coefficients (slopes). The correlation coefficients (slopes) were $0.67 (3.1 \times 10^{-4})$ and $0.31 (2.1 \times 10^{-4})$ for the cases with typical well-mixed BL and all cases, respectively (Fig. 10b).

3.5 Entrainment in stratocumulus

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Entrainment is a key process that occurs in the clouds and plays an important role in the formation and evolution of clouds and the change of droplet spectrum, as well as the aerosol indirect effect (Chen et al., 2014, 2015; Andersen and Cermak, 2015). The nature of entrainment is related to the cloud type. Entrainment in cumulus is primarily lateral with strong dilution of the cloud, which induces *LWC* to decrease rapidly to approximately 20% of its adiabatic value (Warner, 1955). Entrainment in stratocumulus is mainly determined by the strength of the gradients in buoyancy and horizontal winds (Wang and Albrecht

1994; Gerber et al. 2005; de Roode and Wang 2007; Wood, 2012), and proceeds from the top and mostly affects a thin layer (Gerber et al., 2005), whose dilution effect is much weaker than that in cumulus (Warner, 1955, 1969a, 1969b; Blyth et al., 1988; Gerber et al., 2008; Burnet and Brenguier, 2007; Haman et al., 2007). Aircraft observations of marine stratocumulus clouds showed that the vertical profile of *LWC* is essentially the same as the adiabatic profile, i.e., the cloud is almost adiabatic (Keil and Haywood, 2003). Furthermore, it remains unclear whether the subsequent entrainment-mixing mechanism is predominantly homogeneous, inhomogeneous, or in between (Andrejczuk et al., 2009; Lehmann et al., 2009). Some previous studies have shown that stratocumulus is generally dominated by the inhomogeneous mechanism (Pawlowska et al., 2000; Burnet and Brenguier, 2007; Haman et al., 2007; Lu et al., 2011; Yum et al., 2015). By employing a different vertical description in characterizing the region near cloud top (Malinowski et al., 2013), Gerber et al. (2016) noted that both extreme inhomogeneous mixing and homogenous mixing play a role in unbroken stratocumulus, but the reduction in cloud droplet effective radius appears to be secondary in comparison to the dilution process that preserves the relative shape of the droplet spectrum.

To explore entrainment in stratocumulus during VOCALS-REx, we first compare the differences in cloud microphysics between the entrainment and non-entrainment zones near the cloud top. Here, the entrainment and non-entrainment zones are defined as the regions within 20 m above and below the height of the maximal LWC, respectively. As anticipated, the adiabatic fraction (AF, the ratio of the measured LWC to its adiabatic value) in the entrainment zone (AF_{ent}) is generally lower than that in the non-entrainment zone ($AF_{non-ent}$), with mean values for all flights of 0.64 and 0.77, respectively (Table 2), which further confirms the rationality of dividing the two zones. Compared with the non-entrainment zone, the peak radius of cloud droplets in the entrainment zone has little changes (Fig. 11), and the effective radius of cloud droplets (R_e) increases by only 2 % (Table 2). However, N_d and LWC decrease significantly on average, by 29 % and 25 %, respectively (Table 2), especially during the flights on Oct. 18, Nov. 04, Nov. 09 and Nov. 13, where N_d decreases by 60 %, 56 %, 56 % and 59 %, and LWC decreases by 56 %, 62 %, 56 % and 59 %, respectively (Table 2). It is suggested that dry and warm air entrained from cloud top dilutes N_d and LWC by similar amounts, while the size of droplets is relatively unaffected, which is thought as extreme inhomogeneous entrainment-mixing process. Additionally, both P_{LWC} and P_{Nd} are negatively correlated with $AF_{ent}/AF_{non-ent}$, with correlation coefficients of -0.60 and -0.47, respectively, indicating the dependence of the LWC and N_d changes on the adiabatic fraction changes (Fig. S2), where P_{LWC} and P_{Nd} are the reduction percentages in LWC and N_d within the entrainment zone relative to the non-entrainment zone.

The flight on Oct. 18 with strong entrainment is chosen to investigate the difference between the entrainment and non-entrainment zones. As shown in Fig. 12b, dry and warm air entrained from the top reduced the RH in the entrainment zone by 9 % on average and hence acted to accelerate cloud droplet evaporation. Consequently, $N_d/(N_d + N_i)$ in the entrainment zone (0.56 ± 0.22) is much lower than that in non-entrainment zone (0.73 ± 0.13) (Fig. 12c). Additionally, the relative dispersion in

the entrainment zone is generally larger than that in the non-entrainment zone (Fig. 12d), implying that drier air entrained from the top could broaden the cloud droplet spectrum by promoting cloud droplet evaporation. Some previous observations also showed that ε with a low AF tends to be larger than that with high AF, and attributed it to the effect of entrainment mixing (Pawlowska et al., 2006; Lu et al., 2009). It is noted that the occurrence frequency of R_i in the entrainment zone is significantly higher than that in the non-entrainment zone when $R_i < 0.38$ µm, but the opposite is true when $R_i > 0.5$ µm (Fig. 12a). This result suggests that in addition to dry and warm air, small particles are also entrained into clouds from the top (Fig. 2f) and large particles are detrained out of the clouds simultaneously. However, the inversion capping a typical stratocumulus is usually too strong to excite significant updrafts near cloud top (Stevens, 2002; Wood, 2012; Malinowski et al., 2013). Ghate et al. (2010) found that vertical velocities near the top of stratocumulus tend towards zero overall with only approximately 4% of updrafts being stronger than 0.5 m s⁻¹. Therefore, although smaller aerosols are entrained into the entrainment zone, these aerosols seem unlikely to influence droplet formation. The effect of entrainment mixing on stratocumulus is mainly governed by the entrained dry air rather than small aerosols.

4 Summary

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By using in situ aircraft data collected by CIRPAS Twin Otter aircraft at Point Alpha during VOCALS-REx from 16 October to 13 November 2008, aerosol-cloud interactions are investigated with a focus on understudied factors including separation of aerosol effects from dynamic effects, dispersion effects, and turbulent entrainment mixing processes.

Vertical profiles of aerosol, cloud and meteorological variables indicated that the BL was capped by a sharp inversion during 16 non-drizzling flights. Cloud properties, such as LWC and N_d , are positively correlated with sub-cloud CCN number concentration, with correlation coefficients of 0.60 and 0.79, respectively. No evident correlation was found between cloud properties and above-cloud CCN number concentrations. This is mainly due to the low aerosol number concentrations above-cloud (129 \pm 60 cm⁻³) and the extremely strong inversion capping the cloud top, which inhibits the mixing of the above-cloud aerosols with the cloud layer. Therefore, the influence of the above-cloud CCN on cloud properties is weaker than the sub-cloud CCN.

The results showed that both dynamical conditions and aerosol microphysical properties have significant effects on cloud droplet formation. In the case of large vertical velocity and aerosol size, the cloud droplet proportion of total in-cloud particles is relatively high (e.g., 0.84 ± 0.12 and 0.76 ± 0.13 , respectively), i.e., cloud droplets form more easily. Although aerosol chemical components are also critical to cloud droplet formation (Nenes et al., 2002; Lance et al., 2004; Ervens et al., 2005; McFiggans et al., 2006; Wang et al., 2008; Almeida et al., 2014), this topic was not discussed in this study due to the unavailability of measurements.

The correlations between ε and N_d based on the in-flight data, representing the w-induced correlation, are significantly negative, while the correlations derived from flight-averaged data (i.e., aerosol-induced correlation) are positive. This finding implies that an increase in aerosol concentration tends to concurrently increase ε and N_d , while an increase in w often increases N_d but decreases ε , which agrees with the theoretical analysis (Liu et al., 2006). After constraining the differences in cloud dynamics between flights, positive ε - N_d correlations become stronger, indicating that perturbations of w could weaken the influence of aerosols on ε , and hence may result in an underestimation of aerosol dispersion effect. Thus, this finding highlights the necessity of isolating the relative dispersion response to aerosol perturbations from dynamic effects when investigating aerosol dispersion effect and estimating aerosol indirect forcing.

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Overall, the entrainment in stratocumulus is quite weak and close to being adiabatic in some cases. In this study, the difference in cloud microphysics between the entrainment and non-entrainment zones indicated that the entrainment in stratocumulus is mostly dominated by extreme inhomogeneous entrainment-mixing mechanism. On average, the entrainment reduced N_d and LWC by 29 % and 25 %, respectively, while had little effect on R_e (only increases by 1.8 %). During the flights on Oct. 18, Nov. 04, Nov. 09 and Nov. 13, the entrainment was relatively strong and diluted N_d and LWC by about 50 %. In the entrainment zone, the drier air entrained from the top resulted in a smaller $N_d/(N_d+N_i)$ (0.56 ± 0.22) than that in the non-entrainment zone (0.73 ± 0.13) . This implies that entrainment may significantly influence cloud droplet formation and therefore influence the cloud properties near the top by promoting cloud droplet evaporation. Furthermore, we found that the relative dispersion in the entrainment zone is larger than that in the non-entrainment zone. In addition to the dry and warm air, aerosols with smaller sizes are also entrained into the entrainment zone, but these aerosols seem unlikely to influence cloud droplet formation due to the negligible droplet nucleation near the stratocumulus top. That is, the effect of entrainment mixing on stratocumulus is mainly determined by the entrained dry air instead of the aerosols with properties that are different from those near the cloud base. These results seem at odds with some studies on cumulus clouds. Slawinska et al. (2012) found that in a shallow cumulus, a significant part (40 %) of aerosols is activated above cloud base (secondary activation), which is dominated by entrained aerosols. Using large-eddy simulations (LES), Hoffmann et al. (2015) suggested that sub-cloud aerosols and laterally entrained aerosols contribute to all activated aerosols inside the cloud by fractions of 70 % and 30 %, respectively. Evidently, the topics of how and to what extent entrained aerosols with properties that are different from sub-cloud aerosols can affect the formation and evolution of clouds merits further exploration.

Data availability. The aircraft measurements data during VOCALS-REx was obtained from the public ftp at http://data.eol.ucar.edu/master_list/?project=VOCALS.

Author contributions. HJ and XM designed the study and the statistical analysis. HJ processed the data and drafted the manuscript. XM and YL validated and debugged the results. All authors contributed to revising the manuscript.

Competing interests. The authors declare that they have no conflict of interest.

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Table 1. Correlations between ε and N_d (N_d) from observation studies.

Observations	Observation type	Location	Data for correlation analysis	Correlation		
Liu and Daum, 2002	Aircraft	Ocean & coast	Flight-averaged	Positive		
Peng and Lohmann, 2003	Aircraft	Coast	Flight-averaged	Positive		
Pawlowska et al., 2006	Aircraft	Ocean	In-flight Flight-averaged	Negative Positive		
Zhao et al., 2006	Aircraft	Land, ocean, and coast	In-flight	ε converges to a small range of values with increasing N_d		
Lu et al., 2007	Aircraft	Ocean	In-flight Flight-averaged	Negative None for N_a ; Positive for N_a		
Lu et al., 2012	Aircraft	Land	In-flight Flight-averaged	Negative Negative		
Hudson et al., 2012	Aircraft	Ocean	Flight-averaged	Negative		
Ma et al., 2012	Aircraft	Land	Flight-averaged	Negative		
Pandithurai et al., 2012	Aircraft	Land	Flight-averaged	Positive		
Kumar et al., 2016	ground- based	Land	_	Positive		

Table 2. Flight information and parameters that represent the properties of entrainment during all 16 non-drizzling flights.

Flight number	RF01	RF02	RF03	RF04	RF05	RF06	RF07	RF08	RF09
Date	10.16	10.18	10.19	10.21	10.22	10.24	10.26	10.27	10.29
	Typical	Typical	Typical	Typical	Typical	Other	Typical	Typical	Other
BL type	J.F.	J.F.	7 F	J.F	J.F	Wind shear	J.F	J.F	Decoupled
w ave ^a	0.09	0.08	0.11	0.08	0.08	-0.06	0.06	0.08	-0.13
w std ^b	0.42	0.55	0.58	0.51	0.51	0.30	0.56	0.41	0.61
w skew ^c	-0.38	-0.16	-0.27	-0.21	-0.27	0.00	-0.23	0.08	-0.27
$P_{LWC}^{\mathbf{d}}$	26	56	33	25	25	29	-3	11	3
$P_{Nd}^{\mathbf{e}}$	32	60	30	39	28	34	5	20	6
$P_{Re}{}^{ m f}$	-2	-6	1	-7	-2	0	-4	-2	-2
AF_{ent}^{g}	0.77	0.52	0.58	0.85	0.49	0.52	0.51	0.76	0.81
$AF_{non-ent}^{\mathrm{h}}$	0.95	0.84	0.82	0.77	0.74	0.78	0.73	0.82	0.80
Flight number	RF10	RF11	RF12	RF13	RF14	RF15	RF16	Total	
Date	10.30	11.04	11.08	11.09	11.10	11.12	11.13		
	Other	Other	Other	Typical	Typical	Typical	0.1		
BL type	Wind shear	Wind shear,	Decoupled				Other		
		Decoupled					Wind Shear		
w ave	0.02	0.07	0.02	0.08	0.09	0.09	-0.02		
w std	0.45	0.41	0.42	0.47	0.49	0.51	0.41		
w skew	-0.13	-0.48	-0.03	-0.48	-0.26	-0.27	-0.42		
P_{LWC}	10	62	3	56	3	-2	59	25	
P_{Nd}	8	56	24	56	-2	8	59	29	
P_{Re}	0	4	-8	-2	3	-3	-1	-2	
AF_{ent}	0.73	0.66	0.84	0.28	0.70	0.67	0.56	0.64	
AFnon-ent	0.82	0.97	0.77	0.50	0.79	0.60	0.64	0.77	

^{a, b, c} w ave, w std, and w skew are the average, standard deviation, and skewness of in-cloud vertical velocities, respectively..

d, e, f P_{LWC} , P_{Nd} , and P_{Re} are the percentages of reduction in LWC, N_d and R_e within entrainment zone relative to non-entrainment zone.(unit: %)

 $^{^{}g, h}AF_{ent}$ and AF_{non-en} are adiabatic fraction in entrainment zone and non-entrainment zone, respectively. Here, adiabatic fraction is defined as the ratio of the measured to its adiabatic LWC that is calculated using pressure and temperature near cloud base.

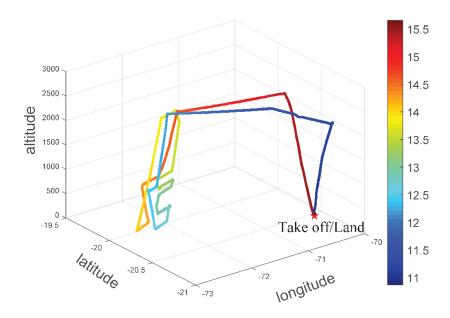


Fig. 1. The flight track in Oct. 18, and the colors represent flight time in hour (UTC).

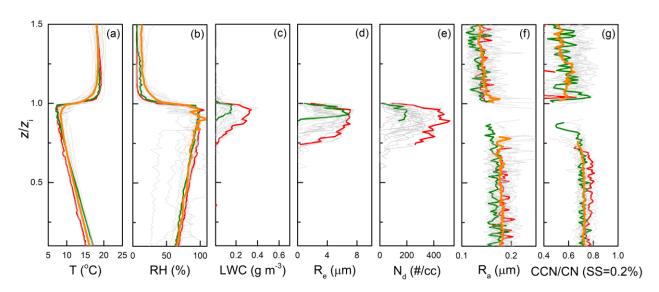


Fig. 2. Vertical profiles scaled by the inversion height. (a) temperature (K); (b) relative humidity (%); (c) liquid water content (g m ⁻³); (d) cloud droplet effective radius (μm); (e) cloud droplet number concentration (cm ⁻³); (f) aerosol effective radius (μm), and (g) the number concentration ratio of CCN to aerosols for all 16 non-drizzling flights. The gray lines show all individual flights, and the orange lines indicate the average profiles. The red and green lines represent the polluted (Oct. 18) and clean (Nov. 9) cases, respectively.

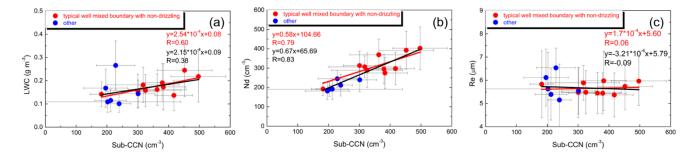


Fig. 3. (a) LWC (g cm⁻³); (b) N_d (cm⁻³); (c) R_e (μ m) as a function of sub-cloud CCN concentrations (SS=0.2%) for all 16 non-drizzling flights. The error bars through these symbols indicate the standard deviation. Red symbols are the cases with typical well-mixed BL discussed in Zheng et al. (2011), and blue symbols for others. Red (black) texts are the correlation coefficient for typical well-mixed cases (all cases).

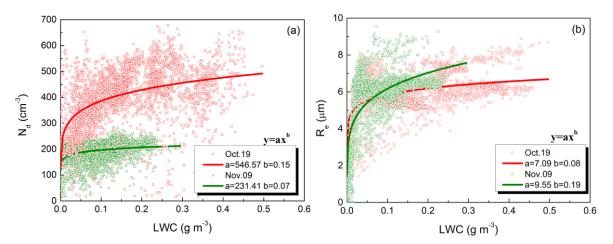


Fig. 4. Correlations between (a) N_d (cm⁻³), (b) R_e (μ m) and LWC (g m⁻³) for clean (green) and polluted (red) cases, respectively.

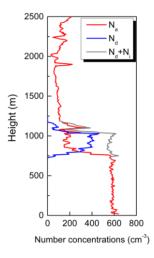


Fig. 5. Vertical profiles of number concentrations of aerosols (N_a) , cloud droplets (N_d) and total in-cloud particles $(N_d + N_i)$ during the flight on Oct. 18.

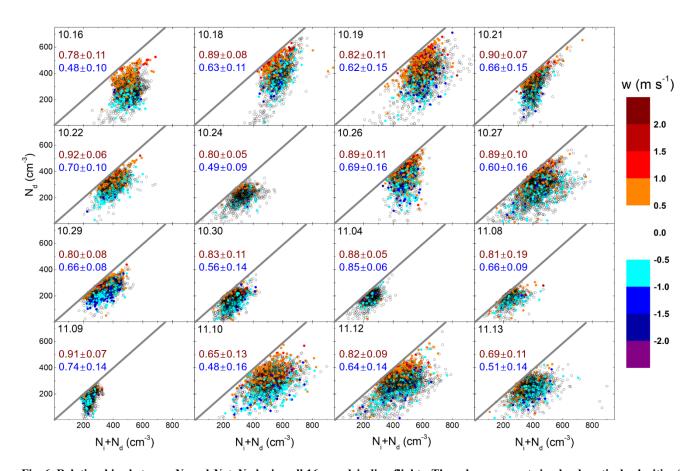


Fig. 6. Relationships between N_d and $N_i + N_d$ during all 16 non-drizzling flights. The color represents in-cloud vertical velocities (m s⁻¹), and gray line is 1:1 line. The mean and standard deviation of $N_d/(N_d+N_i)$ for vertical velocity greater than 1 m s⁻¹ (red) and less than -1 m s⁻¹ (blue) are shown.

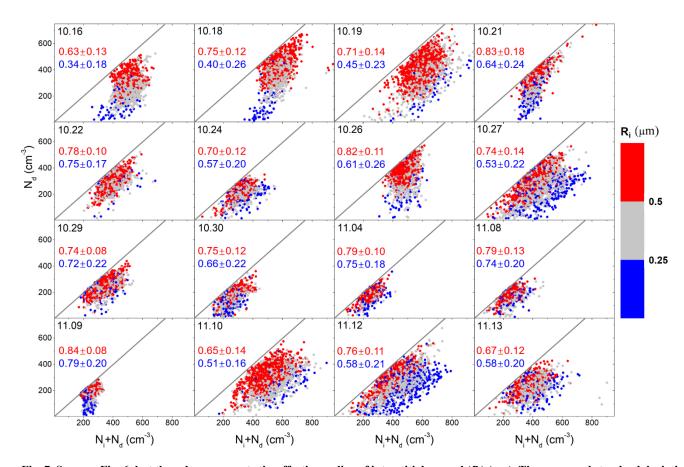


Fig. 7. Same as Fig. 6, but the color represents the effective radius of interstitial aerosol (R_i) (μ m). The mean and standard deviation of $N_d/(N_d+N_i)$ for R_i greater than 0.5 μ m (red) and less than 0.25 μ m (blue) are shown.

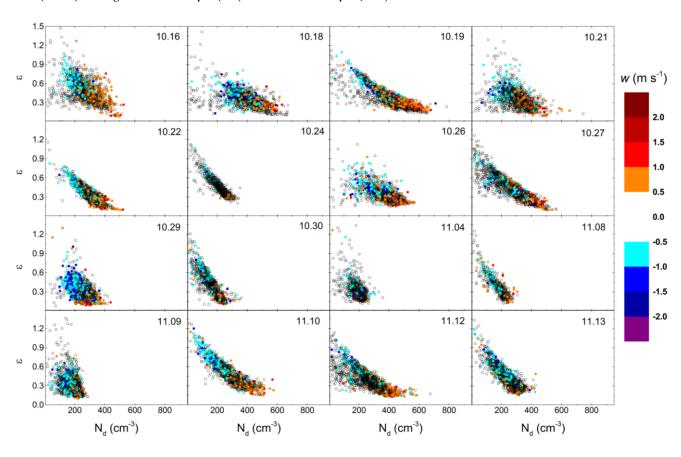


Fig. 8. Relationships between relative dispersion (ϵ) and N_d during all 16 non-drizzling flights, in which the color represents in-cloud

vertical velocities (m s⁻¹).

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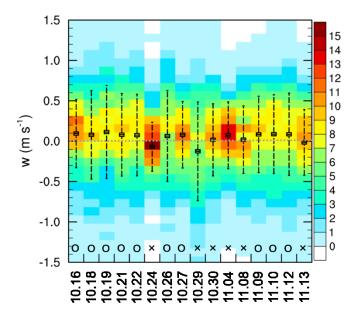


Fig. 9. Probability distribution function (units: %) of vertical velocity (w) for 16 non-drizzling flights. Black symbols are mean values of w, and error bars through these symbols indicate the standard deviation. Circles are the cases with typical well-mixed BL, and crosses represents the other cases.

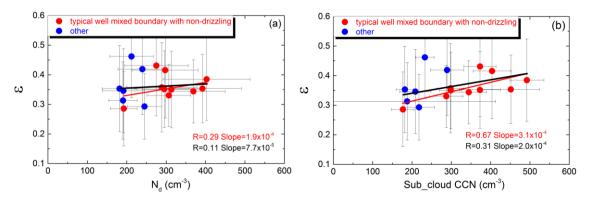


Fig. 10. Relative dispersion (ε) as a function of (a) N_d and (b) sub-cloud CCN concentrations (SS=0.2%) for all flights. The error bars through these symbols indicate the standard deviation. Red symbols are the cases with typical well-mixed BL, and blue symbols for others. Red (black) texts are the correlation coefficient and slope for typical well-mixed cases (all cases).



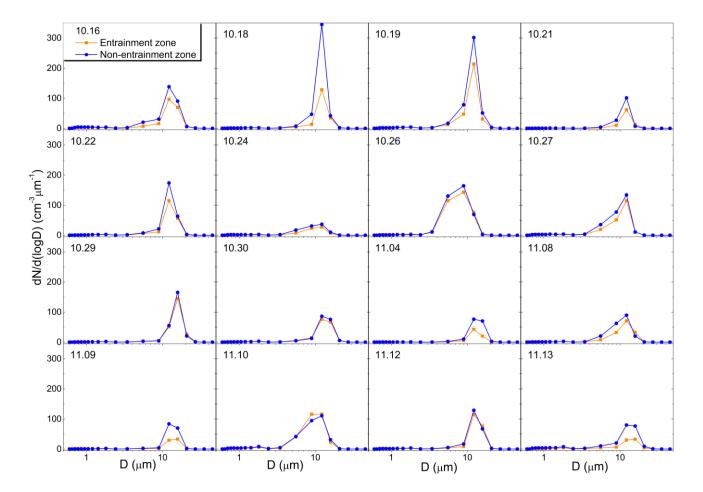


Fig. 11. Number size distributions of cloud droplets in the entrainment (yellow) and non-entrainment zones (blue) during all 16 non-drizzling flights.

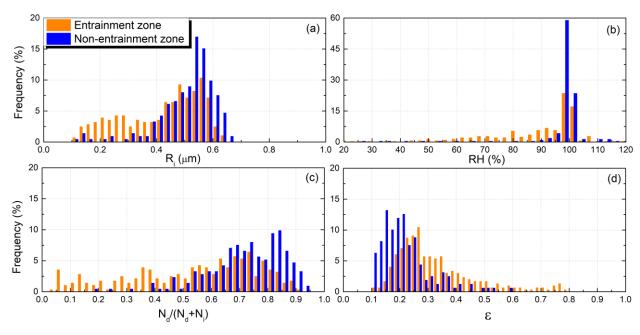


Fig. 12. Probability density functions of (a) R_i (μ m), (b) RH (%), (c) $N_d/(N_d + N_i)$, and (d) ε in the entrainment (yellow) and non-entrainment zones (blue) during the flight on Oct. 18.