

# Analyzing the turbulent Planetary Boundary Layer by remote sensing systems: Doppler wind lidar, aerosol elastic lidar and microwave radiometer

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## Abstract

The Planetary Boundary Layer (*PBL*) is the lowermost region of troposphere and endowed with turbulent characteristics, which can have mechanical and/or thermodynamic origins. Such behavior gives to this layer great importance, mainly in studies about pollutant dispersion and weather forecasting. However, the instruments usually applied in studies about turbulence in the *PBL* have limitations in spatial resolution (anemometer towers) or temporal resolution (instrumentation onboard aircraft). Ground-based remote sensing, both active and passive, offers an alternative for studying the *PBL*. In this study we show the capabilities of combining different remote sensing systems (microwave radiometer [*MWR*], Doppler lidar [*DL*] and elastic lidar [*EL*]) for retrieving a detailed picture on the *PBL* turbulent features. The statistical moments of the high frequency distributions of the vertical wind velocity, derived from *DL* and of the backscattered coefficient derived from *EL*, are corrected by two methodologies, namely first lag and -2/3 correction. The corrected profiles, obtained from *DL* data, present small differences when compared against the uncorrected profiles, showing the low influence of noise and the viability of the proposed methodology. Concerning *EL*, in addition to analyze the influence of noise, we explore the use of different wavelengths that usually include *EL* systems operated in extended networks, like EARLINET, LALINET, MPLNET or SKYNET. In this way we want to show the feasibility of extending the capability of existing monitoring networks without strong investments or changes in their measurements protocols. Two case studies were analyzed in detail, one corresponding to a well-defined *PBL* and another one corresponding to a situation with presence of a Saharan dust lofted aerosol layer and clouds. In both cases we discuss results provided by the different instruments showing their complementarity and the cautions to be applied in the data interpretation. Our study shows that the use of *EL* at 532nm requires a careful correction of the signal using the first lag time correction in order to get reliable turbulence information on the *PBL*.

**Keywords:** Turbulence, Planetary Boundary Layer, Doppler lidar, elastic lidar, microwave radiometer, Earlinet.

## 38 **1 Introduction**

39 The Planetary Boundary Layer (*PBL*) is the atmospheric layer directly influenced by the Earth's surface  
40 that responds to its changes within time scales around an hour (Stull, 1988). Such layer is located at the  
41 lowermost region of troposphere, and is mainly characterized by turbulent processes and a daily evolution  
42 cycle. In an ideal situation, some instants after sunrise, the ground surface temperature increases due to the  
43 positive net radiative flux ( $R_n$ ). This process intensifies the convection, where there is an ascension of warm  
44 air masses, causing the downward displacement of colder air masses and consequently originating the  
45 Convective Boundary Layer (*CBL*) or Mixing Layer (*ML*). Such layer has this name due to the mixing  
46 process generated by the ascending air parcels. Slightly before sunset, the gradual reduction of incoming  
47 solar irradiance at the Earth's surface causes the decrease of the positive  $R_n$  and, consequently, its sign  
48 change. In this situation, there is a reduction of the convective processes and a weakening of the turbulence.  
49 In this process the *CBL* leads to the development of two layers, namely a stably stratified boundary layer  
50 called Stable Boundary Layer (*SBL*) close to the surface, and the Residual Layer (*RL*) that contains features  
51 from the previous day's *ML* and is just above the *SBL*.

52 Knowledge of the turbulent processes in the *CBL* is important in diverse studies, mainly for atmospheric  
53 modeling and pollutant dispersion, since turbulent mixing can be considered as the primary process by  
54 which aerosol particles and other scalars are transported vertically in atmosphere. Because turbulent  
55 processes are treated as nondeterministic, they are characterized and described by their statistical properties  
56 (high order statistical moments). When applied to atmospheric studies such analysis provide information  
57 about the field of turbulent fluctuation, as well as, a description of the mixing process in the *PBL* (Pal et  
58 al., 2010).

59 Anemometer towers have been widely applied in studies about turbulence (e.g., Kaimal and Gaynor, 1983;  
60 van Ulden and Wieringa, 1996), however the limited vertical range of these equipment restrict the analysis  
61 to regions close to surface. Aircraft have also been used in atmospheric turbulence studies (e.g., Lenschow  
62 et al., 1980; Williams and Hacker, 1992; Lenschow et al., 1994; Albrecht et al., 1995; Stull et al., 1997;  
63 Andrews et al., 2004; Vogelmann et al., 2012), nevertheless their short time window limits the analysis. In  
64 this scenario, systems with high spatial and temporal resolution and enough range are necessary in order to  
65 provide more detailed results along the day throughout the whole thickness of the *PBL*.

66 In the last decades, lidar systems have been increasingly applied in this kind of study due to their large  
67 vertical range, high data acquisition rate and capability to detect several observed quantities such as vertical  
68 wind velocity [Doppler lidar] (e.g. Lenschow et al., 2000; Lothon et al., 2006; O'Connor et al., 2010), water  
69 vapor [Raman lidar and DIAL] (e.g. Wulfmeyer, 1999; Kiemle et al., 2007; Wulfmeyer et al., 2010; Turner  
70 et al., 2014; Muppa et al., 2015), temperature [rotational Raman lidar] (e.g. Behrendt et al., 2015) and  
71 aerosol [elastic lidar] (e.g. Pal et al., 2010; McNicholas et al., 2015). This allows the observation of a wide  
72 range of atmospheric processes. For example, Pal et al. (2010) demonstrated how the statistical analyses  
73 obtained from high-order moments of elastic lidar can provide information about aerosol plume dynamics  
74 in the *PBL* region. In addition, when different lidar systems operate synergistically, as for example in

75 Engelmann et al. (2008), who combined elastic and Doppler lidar data, it is possible to identify very  
76 complex variables such as vertical particle flux.

77 Different works (Ansmann et al., 2010; O'Connor et al., 2010) have evidenced the feasibility for  
78 characterizing the *PBL* turbulence by *DL*. Pal et al. (2010) have shown the feasibility for retrieving  
79 information on the *PBL* turbulence from high high-order moments of elastic lidar operating at 1064. Such  
80 approaches are even more attractive when considering facilities of networks, e. g. European Aerosol  
81 Research Lidar NETwork (EARLINET) (Pappalardo et al., 2014), Microwave Radiometer Network  
82 (MWRNET) (Rose et al., 2005; Caumont et al., 2016) and ACTRIS CLOUDNET (Illingworth et al., 2007).  
83 For these reasons, and having in mind the wide spread of elastic lidar systems operated at other wavelengths,  
84 like 532 nm or 355 nm, it would be worthy test the feasibility of these other wavelengths in the  
85 characterization of the *PBL* turbulent behavior.

86 The use of simple techniques, applied to the aforementioned remote systems provide robust and similar  
87 information on the *PBL* height (*PBLH*) during the convective period (see for example Moreira et al, 2018),  
88 or a complementary information when the *CBL* is substituted by the presence of the *SBL* and the *RL*  
89 (Moreira et al., in preparation). Thus, the combination of information obtained from the active remote  
90 sensing systems, *DL* and *EL*, acquired with a temporal resolution close to 1 s, and that provided by *MWR*  
91 can provide a detailed understanding about different features of the *PBL*, like structure (*CBL* versus *SBL*  
92 and *RL*), height of the layers, rate of growth of the *PBLH* and turbulence.

93 In this study we show the feasibility of obtaining a clear insight on the *PBL* behavior using a combination  
94 of active and passive remote sensing systems (Elastic Lidar [*EL*], Doppler Lidar [*DL*] and Microwave  
95 Radiometer [*MWR*]) acquired during the SLOPE-I campaign, held at IISTA-CEAMA (Andalusian Institute  
96 for Earth System Research, Granada, Spain) from May to August 2016. One of the goals is to show the  
97 feasibility of using *EL* at 532 nm, considering the widespread use of lidar systems based on laser emission  
98 at this wavelength in different coordinated networks, like as EARLINET (Pappalardo et al., 2014) and  
99 LALINET – Latin American Lidar Network (Guerrero-Rascado et al., 2016). In addition, this study shows  
100 the variety of application that can be done with EARLINET data applying some simple changes in the data  
101 acquisition procedures.

102 This paper is organized as follows. Description of the experimental site and the equipment setup are  
103 presented in Section 2. The methodologies applied are introduced in Section 3. Section 4 presents the results  
104 of the analyses using the different methodologies. Finally, conclusions are summarized in Section 5.

105

## 106 **2 Experimental site and instrumentation**

107 The SLOPE-I (Sierra nevada Lidar aerOsol Profiling Experiment) campaign was performed from May to  
108 September 2016 in South-Eastern Spain in the framework of the European Research Infrastructure for the  
109 observation of Aerosol, Clouds, and Trace gases (ACTRIS). The main objective of this campaign was to  
110 perform a closure study by comparing remote sensing system retrievals of atmospheric aerosol properties,

111 using remote systems operating at the Andalusian Institute of Earth System Research (IISTA-CEAMA)  
112 and in-situ measurements operating at different altitudes in the Northern slope of Sierra Nevada, around 20  
113 km away from IISTA-CEAMA (Bedoya-Velásquez et al., 2018; Román et al., 2018). The IISTA-CEAMA  
114 station is part of EARLINET (Pappalardo et al., 2014) since 2005 and at present is an ACTRIS station  
115 (<http://actris2.nilu.no/>). The research facilities are located at Granada, a medium size city in Southeastern  
116 Spain (Granada, 37.16°N, 3.61°W, 680 m a.s.l.), surrounded by mountains and with Mediterranean-  
117 continental climate conditions that are responsible for cool winters and hot summers. Rain is scarce,  
118 especially from late spring to early autumn. Granada is affected by different kind of aerosol particles locally  
119 originated and medium-long range transported from Europe, Africa and North America (Lyamani et al.,  
120 2006; Guerrero-Rascado et al., 2008, 2009; Titos et al., 2012; Navas-Guzmán et al., 2013; Valenzuela et  
121 al., 2014, Ortiz-Amezcuca et al., 2014, 2017).

122 MULHACÉN is a biaxial ground-based Raman lidar system operated at IISTA-CEAMA in the frame of  
123 EARLINET research network. This system operates with a pulsed Nd:YAG laser, frequency doubled and  
124 tripled by Potassium Dideuterium Phosphate crystals, emitting at wavelengths of 355, 532 and 1064 nm  
125 with output energies per pulse of 60, 65 and 110 mJ, respectively. MULHACÉN operates with three elastic  
126 channels: 355, 532 (parallel and perpendicular polarization) and 1064 nm and three Raman-shifted  
127 channels: 387 (from N<sub>2</sub>), 408 (from H<sub>2</sub>O) and 607 nm (from N<sub>2</sub>). MULHACÉN's overlap is complete at  
128 90% between 520 and 820 m a.g.l. for all the wavelengths, reaching full overlap around 1220 m a.g.l.  
129 (Navas-Guzmán et al., 2011; Guerrero-Rascado et al. 2010). Calibration of the depolarization capabilities  
130 is done following Bravo-Aranda et al. (2013). This system was operated with a temporal and spatial  
131 resolution of 2 s and 7.5 m, respectively. More details can be found at Guerrero-Rascado et al. (2008, 2009).

132 The Doppler lidar (Halo Photonics, model Stream Line XR) is also operated at IISTA-CEAMA. This  
133 system works in continuous and automatic mode from May 2016. It operates at 1.5 μm with pulse energy  
134 and repetition rate of 100 μJ and 15 KHz, respectively. This system records the backscattered signal with a  
135 range resolution of 30 m in 300 range gates with the first range gate starting at 60 m from the instrument.  
136 The telescope focus is set to approximately 800 m. The instrument was operated in vertical stare mode with  
137 a temporal resolution of 2 s.

138 Furthermore, we operated the ground-based passive microwave radiometer (RPG-HATPRO G2,  
139 Radiometer Physics GmbH), which is member of the MWRnet [<http://cetemps.aquila.infn.it/mwrnet/>]. This  
140 system operates in automatic and continuous mode at IISTA-CEAMA since November 2011. The  
141 microwave radiometer (MWR) measures the sky brightness temperature with a radiometric resolution  
142 between 0.3 and 0.4 K root mean square error at 1 s integration time, using direct detection receivers within  
143 two bands: K-band (water vapor – frequencies: 22.24 GHz, 23.04 GHz, 23.84 GHz, 25.44 GHz, 26.24 GHz,  
144 27.84 GHz, 31.4 GHz) and V-band (oxygen – frequencies: 51.26 GHz, 52.28 GHz, 53.86 GHz, 54.94 GHz,  
145 56.66 GHz, 57.3 GHz, 58.0 GHz). From these bands is possible to obtain profiles of water vapor and  
146 temperature, respectively, by inversion algorithms described in Rose et al. (2005). The range resolution of  
147 these profiles vary between 10 and 200 m in the first 2 km and between 200 and 1000 m in the layer between  
148 2 and 10 km (Navas-Guzmán et al., 2014).

149 The meteorological sensor (HMP60, Vaisala) is used to register the air surface temperature and surface  
150 relative humidity, with a temporal resolution of 1 minute. Relative humidity is monitored with an accuracy  
151 of  $\pm 3\%$ , and air surface temperature is acquired with an accuracy and precision of  $0.6^\circ\text{C}$  and  $0.01^\circ\text{C}$ ,  
152 respectively.

153 A CM-11 pyranometer manufactured by Kipp&Zonen (Delft, The Netherlands) is also installed in the  
154 ground-based station. This equipment measures the shortwave (SW) solar global horizontal irradiance data  
155 (305–2800 nm). The CM-11 pyranometer complies with the specifications for the first-class WMO (World  
156 Meteorological Organization) classification of this instrument (resolution better than  $\pm 5\text{ Wm}^{-2}$ ), and the  
157 calibration factor stability has been periodically checked against a reference CM-11 pyranometer (Antón  
158 et. al, 2012).

### 159 **3 Methodology**

#### 160 **3.1 MWR data analysis**

161 The MWR data are analyzed combining two algorithms, Parcel Method [*PM*] (Holzworth, 1964) and  
162 Temperature Gradient Method [*TGM*] (Coen, 2014), in order to estimate the *PBL* Height ( $PBLH_{MWR}$ ) in  
163 convective and stable situations, respectively. The different situations are discriminated by comparing the  
164 surface potential temperature ( $\theta(z_0)$ ) with the corresponding vertical profile of  $\theta(z)$  up to 5 km. Those  
165 cases where all the points in the vertical profile have values larger than  $\theta(z_0)$  are labeled as stable, and  
166 *TGM* is applied. Otherwise the situation is labeled as unstable and the *PM* is applied. The vertical profile  
167 of  $\theta(z)$  is obtained from the vertical profile of  $T(z)$  using the following equation (Stull, 2011):

$$168 \quad \theta(z) = T(z) + 0.0098 * z \quad (1)$$

169 where  $T(z)$  is the temperature profile provided by *MWR*,  $z$  is the height above the sea level, and 0.0098  
170 K/m is the dry adiabatic temperature gradient. A meteorological station co-located with the *MWR* is used  
171 to detect the surface temperature [ $T(z_0)$ ]. In order to reduce the noise,  $\theta(z)$  profiles were averaged  
172 providing a  $PBLH_{MWR}$  value at 30 minutes intervals. This methodology of *PBLH* detection was selected as  
173 the reference due to the results obtained during a performed intercomparison campaign between *MWR* and  
174 radiosonde data, where twenty-three radiosondes were launched. High correlations were found between  
175 *PBLH* retrievals provided by both instruments in stable and unstable cases. Further details are given by  
176 Moreira et al. (2018a).

#### 177 **3.2 Lidar retrieval of the PBLH.**

178 The simple processing of *DL* and *EL* data allows the estimation of the *CBL* height. Moreira et al. (2018),  
179 have discussed this issue in depth, while Moreira et al. (in preparation) have exploited the complementarity  
180 of the data obtained from distinct remote sensing systems in order to distinguish the sublayers during the

181 period when the *SBL* and *RL* substitute the *CBL*, as well as, in complex situations, like as, presence of dust  
 182 layers.

183 The *PBLH* obtained from *DL* data ( $PBLH_{Doppler}$ ) is estimated from variance threshold method. In this  
 184 method the  $PBLH_{Doppler}$  is attributed to height where the variance of vertical wind speed ( $\sigma_w^2$ ) is lower than  
 185 a determinate threshold, which was adopted as  $0.16 \text{ m}^2/\text{s}^2$  (Moreira et al., 2018). For the  $PBLH_{Doppler}$   
 186 calculations was selected a time interval of 30 minutes. In concerning the *PBLH* obtained from *EL*  
 187 ( $PBLH_{Elastic}$ ), the variance method is applied. Such method assumes the maximum of the variance of  
 188 Range Corrected Signal ( $\sigma_{RCS}^2$ ) as  $PBLH_{Elastic}$  (Moreira et al., 2015). The  $\sigma_{RCS}^2$  is obtained from a time  
 189 interval of 30 minutes.

### 190 3.3 Lidar turbulence analysis

191 Both lidar systems, *DL* and *EL*, gathered data [ $q(z, t)$ ] with a temporal resolution of 2 seconds. Then, the  
 192 data are averaged in 1-hour packages, from which the mean value is extracted [ $\bar{q}(z)$ ]. Such mean value is  
 193 subtracted from each  $q(z, t)$  profile in order to estimate the vertical profile of the fluctuation for the  
 194 measured variable [ $q'(z, t)$ ] (i.e. vertical velocity for the *DL*):

$$195 \quad q'(z, t) = q(z, t) - \bar{q}(z) \quad (2)$$

196 Then, from  $q'(z, t)$  is possible to obtain the high-order moments (variance ( $\sigma^2$ ), skewness ( $S$ ) and kurtosis  
 197 ( $K$ )), as well as, the integral time scale ( $\tau$  - which is the time over which the turbulent process are highly  
 198 correlated to itself) as shown in Table 1. These variables can also be obtained from the following  
 199 autocovariance function,  $M_{ij}$ :

$$200 \quad M_{ij} = \int_0^{t_f} [q'(z, t)]^i [q'(z, t + t_f)]^j dt \quad (3)$$

201 where  $t_f$  is the final time,  $i$  and  $j$  indicate the order of autocovariance function.

202 However, it is necessary to considerer that the acquired real data contain instrumental noise,  $\varepsilon(z)$ .  
 203 Therefore, the equation 3 can be rewritten as:

$$204 \quad M_{ij} = \int_0^{\tau} [q(z, t) + \varepsilon(z, t)]^i [q(z, t + \tau) + \varepsilon(z, t + \tau)]^j dt \quad (4)$$

205 The autocovariance function of a time series with zero lag results in the sum of the variances of the  
 206 atmospheric variable and its  $\varepsilon(z)$ . Nevertheless, atmospheric fluctuations are correlated in time, but the  
 207  $\varepsilon(z)$  is random and uncorrelated with the atmospheric signal. Consequently, the noise is only associated  
 208 with lag 0 (Fig. 1). Based on this concept Lenschow et al. (2000) suggested to obtain the corrected  
 209 autocovariance function,  $M_{11}(\rightarrow 0)$ , from two methods, namely first lag correction or -2/3 law correction.  
 210 In the first method,  $M_{11}(\rightarrow 0)$  is obtained directly by the subtraction of lag 0,  $\Delta M_{11}(0)$ , from the  
 211 autocovariance function,  $M_{11}(0)$ . In the second method  $M_{11}(\rightarrow 0)$  is generated by the extrapolation of

212  $M_{11}(0)$  at firsts nonzero lags back to lag zero (-2/3 law correction). The extrapolation can be performed  
 213 using the inertial subrange hypothesis, which is described by the following equation (Monin and Yaglom,  
 214 1979):

$$215 \quad M_{11}(\rightarrow 0) = \overline{q'^2(z, t)} + Ct^{2/3} \quad (5)$$

216 where  $C$  represents a parameter of turbulent eddy dissipation rate. The high-order moments and  $\tau$   
 217 corrections and errors are shown in Table 1 (columns 2 and 3, respectively).

218 The same procedure of analysis is applied in studies with  $DL$  and  $EL$ , being the main difference the tracer  
 219 used by each system, which are the fluctuation of vertical wind speed ( $w'$ ) for  $DL$  and aerosol number  
 220 density ( $N'$ ) for  $EL$ .  $DL$  provides  $w(z, t)$  directly, and therefore the procedure described in Figure 2 can be  
 221 directly applied. Thus, the two corrections described above are applied separately and finally  $\tau$  and high-  
 222 order moments with and without corrections can be estimated.

223 On the other hand, the  $EL$  does not provide  $N(z, t)$  directly. Under some restrictions, it is possible to ignore  
 224 the particle hygroscopic growth and to assume that the vertical distribution of aerosol type does not changes  
 225 with time, and to adopt the following relation (Pal et al., 2010):

$$226 \quad \beta_{par}(z, t) \approx N(z, t)Y(z) \Rightarrow \beta'_{par}(z, t) = N'(z, t) \quad (6)$$

227 where  $\beta_{par}$  and  $\beta'_{par}$  represent the particle backscatter coefficient and its fluctuation, respectively, and  
 228  $Y(z)$  does not depend on time.

229 Considering the lidar equation:

$$230 \quad P_{\lambda}(z) = P_0 \frac{ct_d}{2} AO(z) \frac{\beta_{\lambda}(z)}{z^2} e^{-2 \int_0^z \alpha_{\lambda}(z' dz')} \quad (7)$$

231 where  $P_{\lambda}(z)$  is the signal returned from distance  $z$  at time  $t$ ,  $z$  is the distance [m] from the lidar of the  
 232 volume investigated in the atmosphere,  $P_0$  is the power of the emitted laser pulse,  $c$  is the light speed [m/s],  
 233  $t_d$  is the duration of laser pulse [ns],  $A$  is the area [m<sup>2</sup>] of telescope cross section,  $O(z)$  is the overlap  
 234 function,  $\alpha_{\lambda}(z)$  is the total extinction coefficient (due to atmospheric particles and molecules) [(km)<sup>-1</sup>] at  
 235 distance  $z$ ,  $\beta_{\lambda}(z)$  is the total backscatter coefficient (due to atmospheric particles and molecules) [(km·sr)<sup>-1</sup>]  
 236 at distance  $z$  and the subscript  $\lambda$  represents the wavelength. The two path transmittance term related to  
 237  $\alpha(z)$  is considered as nearly negligible at 1064 nm (Pal et al., 2010). Thus, it is possible to affirm that:

$$238 \quad RCS_{1064}(z) = P(z)_{1064} \cdot z^2 \cong G \cdot \beta_{1064}(z) \quad (8)$$

239 and consequently:

$$240 \quad RCS'_{1064}(z, t) \cong \beta'_{1064}(z, t) = \beta'_{par}(z, t) = N'(z, t) \quad (9)$$

241 where  $RCS_{1064}$  and  $RCS'_{1064}$  are the range corrected signal and its fluctuation, respectively,  $G$  is a constant  
 242 and the subscripts represent the wavelength.



243 In this way, Pal et al. (2010) have shown the feasibility of using *EL* operating at 1064 nm for describing  
244 the atmospheric turbulence. However, having in mind the more extended use of lidar systems based on  
245 laser emission at 532 nm in different coordinated networks, e.g., in EARLINET and LALINET around 76%  
246 and 45% of the systems include the wavelength of 1064 nm, while 95% of the EARLINET systems and  
247 73% of the LALINET systems operate systems that include the wavelength 532 nm (Guerrero-Rascado et  
248 al., 2016), in this study we evaluate using  $RCS_{532}$  fluctuations to determine turbulence following the  
249 procedure described in Figure 3. This *EL* methodology is very similar to that described earlier for *DL*.

250 we perform the validation of the  $RCS_{532}$  in analyses about turbulence using *EL*, following the procedure  
251 described in Figure 3, which is basically the same methodology described earlier for *DL*.

## 252 4 Results

### 253 4.1 Error Analysis

254 The influence of random error in noisy observations rapidly grows for higher-order moments (i.e., the  
255 influence of random noise is much larger for the fourth-order moment than for the third-order moment).  
256 Therefore, the first step, in order to ascertain the applied methodology and our data quality, we performed  
257 the error treatment of *DL* data as described in Figure 2. For the *DL* analysis we selected the period 08-09  
258 UTC of 19<sup>th</sup> May, the same day that will be presented in Case Study 1. This day is characterized by a well-  
259 defined PBL.

260 Figure 4 illustrates the autocovariance function, generated from  $w'$ , at three different heights. As mentioned  
261 before, the lag 0 is contaminated by noise ( $\epsilon$ ), and thus the impact of the  $\epsilon$  increases together with height,  
262 mainly above  $PBLH_{MWR}$  (1100 m a.g.l. in our example).

263 Figure 5-A illustrates the comparison between integral time scale ( $\tau_{w'}$ ) without correction and the two  
264 corrections cited in section 3.2. Except for the first height-bins, below the  $PBLH_{MWR}$  the profiles have little  
265 differences, as well as small errors bars. Above the  $PBLH_{MWR}$  the first lag correction presents higher  
266 differences in relation to the other profiles at around 1350 m.

267 Figures 5-B and 5-C show the comparison of variance ( $\sigma_{w'}^2$ ) and skewness ( $S_{w'}$ ), respectively, with and  
268 without corrections. The profiles corrected by -2/3 law do not present significant differences in comparison  
269 to uncorrected profiles. On the other hand, the profiles corrected by the first lag correction have slight  
270 differences below the  $PBLH_{MWR}$ , mainly the  $\sigma_{w'}^2$  ( $S_{w'}$  only in the first 50 m). Therefore, considering high  
271 Signal-to-Noise Ratio (*SNR*) conditions, although the presence of  $\epsilon$  can change slightly the value of high  
272 order moments, it is not enough to distort the observed phenomena as shown by the impact of the corrections  
273 applied.

274 For *EL* we use the same procedure for the correction and error analysis that we apply to the *DL* data. The  
275 same day was chosen (19<sup>th</sup> May), however the period selected is between 12 and 13 UTC, due to the  
276 incomplete overlap of MULHACÉN.



277 In this sense, we studied the influence of noise at two wavelengths: 1064 nm, that has been previously  
278 analyzed by Pal et al. (2010) as presented in the section 2 and adopted as reference (considering the rather  
279 low impact of molecular signal and the two ways transmittance shown in 9) and 532 nm, just in order to  
280 check the feasibility of this wavelength for turbulence studies considering its widespread use in observation  
281 networks (Pappalardo et al., 2014; Guerrero-Rascado et al., 2016). Figures 6 and 7 shows the  
282 autocovariance function, obtained from  $RCS'_{1064}$  and  $RCS'_{532}$ , respectively, at three distinct heights. As  
283 expected,  $\varepsilon$  increases with range, principally above the  $PBLH_{MWR}$ . However, the wavelength 532 nm is  
284 more influenced by the noise, what can be verified by the higher peak at lag 0 in figure 7, in comparison  
285 with peaks at same lag in figure 6.

286 Although the level of influence of  $\varepsilon$  in each wavelength depends on the  $SNR$  of them (which is associated  
287 to technical factors such as laser output power, filters, type of detectors), considering the proposed  
288 methodology, to evaluate the composition of each wavelength is also important. The large contribution of  
289  $\beta_{Molecular}^{532}$  to the total  $\beta$  at 532 nm in comparison with the behavior at 1064 nm, can influence the results  
290 obtained from such wavelength, because our methodology is based on the use of  $\beta'_{Aerosol}$ . In addition, the  
291 larger extinction (due to both aerosol particles and molecules) at 532 nm produces a lower two-way  
292 transmittance, resulting in the reduction of the  $SNR$  values at this wavelength. As we used Elastic lidar  
293 technique, we could not calculate aerosol extinction profiles, but an estimation of these transmittances was  
294 done on the basis of Klett method (Klett, 1985). With this method, a constant lidar ratio value was  
295 constrained for each profile using the AOD derived from a collocated AERONET Sun-photometer  
296 (Guerrero-Rascado et al., 2008). Using these constrained lidar ratios, the transmittances were calculated  
297 together with aerosol backscatter profiles, integrated up to 2.5 km. The estimated two-way transmittance  
298 was 0.85 for the case analyzed in this subsection (19<sup>th</sup> May).

299 Figures 8-A, 8-B, 8-C and 8-D show the vertical profiles of  $\tau_{RCS'}$ ,  $\sigma_{RCS'}^2$ ,  $S_{RCS'}$  and kurtosis ( $K_{RCS'}$ ),  
300 respectively, obtained at 1064 nm, with and without the corrections described in section 3.2. In general, the  
301 corrections do not affect the profiles generated from 1064 nm data in a significant way, so that, the higher  
302 influence of corrections is observed in the  $K_{RCS'}$  profile, which is underestimated in some regions. In the  
303 figures 9-A, 9-B, 9-C and 9-D we show same high order moments calculated from 532 nm data. As the  
304 complexity of moments increases, it is possible to observe the larger influence of the corrections, due to  
305 propagation of noise. Nonetheless, the application of the corrections, mainly first lag correction, make these  
306 profiles very similar to those generated from the wavelength 1064 nm, so that the same phenomena can be  
307 observed in both.

308 Therefore, in spite of the larger attenuation expected at 532 nm wavelength, which reduces the  $SNR$  of the  
309 profiles in comparison with 1064 nm, the application of the proposed corrections, mainly the first lag,  
310 reduces significantly such influence and enable the observation of the same phenomena detected in the  
311 high-order moments obtained from 1064 nm. Consequently, the wavelength 532 nm will be applied in the  
312 analysis presented in section 4.2. The first lag correction was adopted as default due to it generates more  
313 relevant results in comparison with the -2/3 law correction, providing a more careful analysis.

## 314 4.2 Case studies

315 In this section we present two study cases, in order to show how the products indicated in table 2 can  
316 provide a detailed description about the turbulence in the *PBL*. The first case represents a typical day with  
317 a clear sky situation. The second case corresponds to a more complex situation, where there is presence of  
318 clouds and Saharan mineral dust layers.

### 319 4.2.1 Case study I: clear sky situation

320 In this case study we use measurements gathered with *DL*, *MWR* and pyranometer during 24 hours. The  
321 *EL* was operated under operator-supervised mode between 08:20 to 18:00 UTC.

322 Figure 10 (A) shows the integral time scale obtained from *DL* data ( $\tau_{w'}$ ). The gray area represents the region  
323 where it is not possible to analyze the turbulent process from our *DL* data, either because of the low *SNR*  
324 values, which results in null values of the  $\tau_{w'}$ , or due to regions where the  $\tau_{w'}$  is not null, but it is lower  
325 than the acquisition time of the *DL*. However, the gray area is located almost entirely above the  $PBLH_{MWR}$   
326 (white stars).

327 The  $\sigma_{w'}^2$  has low values during the entire period when the *SBL* is present (Figure 10-B). Nevertheless, as air  
328 temperature begins to increase (around 07:00 UTC), the  $\sigma_{w'}^2$  increases together, as well as, the  $PBLH_{MWR}$ .  
329 The  $\sigma_{w'}^2$  reaches its maximum values in the middle of the day, when we also observe the maximum values  
330 of air temperature and  $PBLH_{MWR}$ . The combination of  $\sigma_{w'}^2$  and  $PBLH_{MWR}$  provides us a better  
331 comprehension about the *PBLH* growth speed, so that, in the moments where high values of  $\sigma_{w'}^2$  are  
332 observed, it means higher values of Turbulent Kinetic Energy (*TKE*), which favor the fast ascension of  
333 *PBLH*.

334 The skewness of  $w'$  ( $S_{w'}$ ) is shown in Figure 11-C. The  $S_{w'}$  describes the distribution of the turbulent  
335 velocities. Thus positive  $S_{w'}$  implies strong but narrow updrafts surrounded by weaker but more widespread  
336 downdrafts, and vice versa for negative  $S_{w'}$ . Consequently, positive values (red regions) correspond with a  
337 surface-heating-driven boundary layer, while negative (blue regions) ones are associated to cloud-top long-  
338 wave radiative cooling. During the stable period, there is predominance of low absolute values of  $S_{w'}$ .  
339 Nevertheless, as air temperature increases (transition from stable to unstable period),  $S_{w'}$  values begin to  
340 become larger. Air temperature begins to decrease around 18:00 UTC, and there is a reduction of  $S_{w'}$ , so  
341 that, the generation rate of convective turbulence decreases. Therefore, the turbulence cannot be maintained  
342 against dissipation, then the *CBL* becomes a *SBL* covered by the *RL*. Thus, the reduction observed in the  
343  $PBLH_{MWR}$  is due to the detection of *SBL* height.

344 Figure 10-D shows the values of net surface radiation ( $R_n$ ) that are estimated from solar global irradiance  
345 values using the seasonal model described in Alados et al. (2003). The negative values of  $R_n$  are  
346 concentrated in the stable region. The  $R_n$  begins to increase around 06:00 UTC and reaches its maximum  
347 in the middle of the day. Comparing figures 8-C and 8-D, we can observe similarity among the behavior of

348  $S_w$ , and  $R_n$ , so that, the joint analysis of these variables reinforce the characterization of this *PBL* as surface-  
349 heating-driven *CBL*.

350 Figure 10-E presents the values of surface air temperature and surface relative humidity (*RH*). Air surface  
351 temperature has a pattern of increase and decrease similar to observed in  $R_n$  and  $S_w$ . On the other hand,  
352 *RH* is inversely correlated with temperature.

353 Figure 11 shows the  $RCS_{532}$  profile obtained from 08:00 to 18:00 UTC. At the beginning of the  
354 measurement period (08:20 to 10:00 UTC) it is possible to observe the presence of a thin residual layer  
355 (around 2000 m a.s.l.), and later from 13:00 to 18:00 UTC it is evident a lofted aerosol layer. In this picture  
356 there are the  $PBLH_{MWR}$  (pink stars), the  $PBLH_{Doppler}$  (blue stars), obtained from the maximum of  $\sigma_w^2$ ,  
357 (Moreira et al., 2018a), and the  $PBLH_{Elastic}$  (black stars), obtained from the maximum of  $\sigma_{RCS}^2$  (Moreira  
358 et al., 2015). In the initial part of measurement, all profiles have similar behavior. However due to distinct  
359 *PBLH* definition and tracer applied by each one, the differences increase as *CBL* becomes more complex,  
360 e.g. the presence of lofted aerosol layer at 14 UTC. The joint observation of the results provided by these  
361 three methods can provide us information about the sublayers in the *PBL*, both in convective and stable  
362 situations. Due to low variability of *PBLH*, the period between 13:00 and 14:00 UTC has been selected to  
363 be analyzed from the high order moments.

364 Figure 12 presents the statistical moments generated from  $RCS'$  of wavelength 532 nm, which were obtained  
365 from 13:00 and 14:00 UTC. The red line in all graphics represent the  $PBLH_{Elastic}$  (2200 m a.s.l.) and the  
366 blue one the average value of  $PBLH_{MWR}$  (2250 m a.s.l.), both obtained between 13 and 14 UTC.

367 Due to presence of a decoupled aerosol layer at 13:30, the average values of  $PBLH_{Elastic}$  and  $PBLH_{MWR}$   
368 have a difference of around 500 m. The  $\sigma_{RCS}^2$  has small and practically constant values between 1000 and  
369 1400m, evidencing the homogeneity of aerosol distribution in this region. From 1400 m the value of  $\sigma_{RCS}^2$ ,  
370 begins to increase, reaching a positive peak at  $PBLH_{MWR}$ , which represents the Entrainment Zone (region  
371 characterized by an intense mixing between air parcels coming from *CBL* and Free Troposphere (*FT*),  
372 causing a high variation in aerosol concentration). The  $PBLH_{Elastic}$  observed at approximately 2900 m  
373 demonstrate an inherent difficulty of variance method to detect the *PBLH* in the presence of several aerosol  
374 layers (Kovalev and Eichinger, 2004). Above  $PBLH_{Elastic}$  the values of  $\sigma_{RCS}^2$  decrease slowly due to  
375 location of the lofted aerosol around 2500 m. However, above this aerosol layer the value of  $\sigma_{RCS}^2$  is  
376 reduced to zero, indicating a large homogeneity in aerosol distribution at this region, what is expected,  
377 because the aerosol concentration at the *FT* is negligible in this case. The integral time scale obtained from  
378  $RCS'$  ( $\tau_{RCS'}$ ) has values higher than *EL* time acquisition throughout the *CBL*, evidencing the feasibility for  
379 studying turbulence using this elastic lidar configuration. The skewness values obtained from  $RCS'$  ( $S_{RCS'}$ )  
380 give us information about aerosol motion. The positive values of  $S_{RCS'}$  observed in the lowest part of profile  
381 and above the  $PBLH_{Elastic}$  represents the updrafts aerosol layers. The negative values of  $S_{RCS'}$  indicates the  
382 region with low aerosol concentration due to clean air coming from *FT*. This movement of ascension of  
383 aerosol layers and descent of clean air with zero value of  $S_{RCS'}$  at *PBLH* (characteristic of the *CBL* growing)  
384 was also detected by Pal et al. (2010) and McNicholas et al. (2014). The kurtosis of  $RCS'$  ( $K_{RCS'}$ ) determines  
385 the level of mixing at different heights. There are values of  $K_{RCS'}$  larger than 3 in the lowest part of profile

386 and around 2500 m, showing a peaked distribution in this region. On other hand, values of  $K_{RCS}$ , lower than  
387 3 are observed close to the  $PBLH_{Elastic}$ , therefore this region has a well-mixed  $CBL$  regime. Pal et al. (2010)  
388 and McNicholas et al. (2014) also detected this feature in the region nearby the  $PBLH$ . In figure 13 are  
389 shown the high-order moments obtained at the same period described above, however from the 1064 nm  
390 data (our reference wavelength). It is possible to observe a similarity between the profiles obtained from  
391 each wavelength, so that, the same phenomena observed in the profiles generated from 532 nm and  
392 described above, also are detected in the profiles obtained from the reference wavelength.

393 The results provided by  $DL$ , pyranometer and  $MWR$  data agree with the results observed in figures 12 and  
394 13. In the same way, the analysis of high order moments of  $RCS'$  fully agree with the information in Figure  
395 10. Thus, the large values of  $S_{RCS'}$  and  $K_{RCS'}$  detected around 2500 m a.s.l, where we can see a lofted aerosol  
396 layer, suggest the ascent of an aerosol layer and presence of a peaked distribution, respectively.

#### 397 4.2.2 Case study: dusty and cloudy scenario

398 In this case study measurements with  $DL$ ,  $MWR$  and pyranometer expand during 24 hours, while  $EL$  data  
399 are collected from 09:00 to 16:00 UTC.

400 Figure 14-A shows  $\tau_{wr}$ . Outside the period 13:00 to 17:00 UTC, the greatest part of grey area is situated  
401 above the  $PBLH_{MWR}$  (white stars), thus  $DL$  time acquisition is enough to perform studies about turbulence  
402 in this case.

403  $\sigma_{wr}^2$  has values close to zero during all the stable period (Figure 14-B). However, when air temperature  
404 begins to increase (around 06:00 UTC), the  $\sigma_{wr}^2$  also increases and reaches its maximum in the middle of  
405 the day. The higher values of  $PBL$  growth speed are observed in the moments where  $\sigma_{wr}^2$  reaches its  
406 maximum values. In the late afternoon, as air temperature decrease, the values of  $\sigma_{wr}^2$  (and consequently  
407 the  $TKE$ ) decrease gradually, until reach the minimum value associated to the  $SBL$ . Figure 14-C shows the  
408 profiles of  $S_{wr}$ . The main features of this case are: the low values of  $S_{wr}$ , the slow increase and ascension  
409 of positive  $S_{wr}$  values and the predominance of negative  $S_{wr}$  values from 12:00 to 13:00 UTC. The first two  
410 features are likely due to the presence of the intense Saharan dust layer (Figure 15), which reduces the  
411 transmission of solar irradiance, and consequently the absorption of solar irradiance at the surface,  
412 generating weak convective process. From figure 16 we can observe the presence of both middle altitude  
413 clouds and very intense dust layers from 12:00 to 15:00 UTC. Such combination contributes to the intense  
414 negative values of  $S_{wr}$  observed in this period until around 2 km, because, as mentioned previously,  $S_{wr}$  is  
415 directly associated with the direction of turbulent movements, which during this period can be characterized  
416 as cloud-top long-wave radiative cooling (Ansmann et al., 2010).

417 The influence of Saharan dust layer can also be evidenced on the  $R_n$  pattern (Figure 14-D), which maintains  
418 negative values until 12:00 UTC and reaches a low maximum value (around 200 W/m<sup>2</sup>). The observation  
419 of  $S_{wr}$  and  $R_n$  between 12:00 and 14:00, as well as, the presence of clouds and geometrically thick dust  
420 layers during this same period, reinforces the hypothesis of a case of the cloud-top long-wave radiative  
421 cooling in the  $CBL$ . Air surface temperature and  $RH$  (Figure 14-E) present the same correlation and anti-

422 correlation (respectively) observed in the earlier case study, where the maximum of air surface temperature  
423 and the minimum of  $RH$  are detected in coincidence with the maximum daily value of  $PBLH_{MWR}$ .

424 As mentioned before, Figure 15 shows the  $RCS$  profile obtained from 09:00 to 16:00 UTC in a complex  
425 situation, with presence of decoupled dust layer (around 3800 m a.s.l.) from 09:00 to 12:00 UTC and the  
426 presence of both middle altitude clouds and very intense dust layers (around 3500 m a.s.l.) from 11:30 to  
427 16:00 UTC. The pink, black and blue stars represent the  $PBLH_{MWR}$ ,  $PBLH_{Doppler}$  and  $PBLH_{Elastic}$   
428 respectively. Due to the presence of dusty layers and clouds, the difference between the methods is more  
429 evident, mainly of the  $PBLH_{Elastic}$ , which uses the aerosol as tracers. This method only produces results  
430 close to the others at 15 UTC, when dust layer is mixed with the  $CBL$ .

431 Figure 16 illustrates the statistical moments of  $RCS'$  of 532 nm wavelength obtained from 11:00 to 12:00  
432 UTC. The  $\sigma_{RCS'}^2$  profile presents several peaks due to the presence of distinct aerosol sublayers. The first  
433 peak is coincident with the value of  $PBLH_{MWR}$ . The value of  $PBLH_{elastic}$ , is coincident with the base of  
434 the dust layer. This difficulty to detect the  $PBLH$  in presence of several aerosol layers is inherent to the  
435 variance method (Kovalev and Eichinger, 2004). However, the joint observation of  $PBLH_{MWR}$  and  
436  $PBLH_{elastic}$ , enable us to characterize and distinguish the several sublayers. The values of  $\tau_{RCS'}$  are higher  
437 than  $EL$  acquisition time all along the  $PBL$ , evidencing the feasibility of  $EL$  time acquisition for studying  
438 the turbulence of  $PBL$  in this case. The  $S_{RCS'}$  profile has several positive values, due to the large number of  
439 aerosol sublayers that are present. The characteristic inflection point of  $S_{RCS'}$  is observed in coincidence  
440 with the  $PBLH_{MWR}$ , that confirming the agreement between this point and the  $PBLH$ . From the analysis of  
441  $S_{RCS'}$  and  $S_w$  is possible to justify this phenomena from the mixing process demonstrated in the earlier case  
442 study. The  $K_{RCS'}$  has predominantly values lower than 3 below 2500 m, thus shown how this region is well  
443 mixed as can see in Figure 16. Values of  $K_{RCS'}$  larger than 3 are observed in the highest part of profile, where  
444 the dust layer is located.

445 In order to show the feasibility of 532 nm wavelength, in the figure 17 are presented the high-order moments  
446 obtained between 11-12 UTC from 1064 nm wavelength data. Although the error of  $\sigma_{RCS'}^2$  obtained from  
447 532 nm (pink shadow) is considerably higher than the error of same variable obtained from 1064 nm, all  
448 profiles are very similar, so that, the same phenomena can be observed in both graphics (figure 16 and 17).

449 Figure 18 shows the  $RCS'$  532 nm wavelength high-order moments obtained from 12:00 and 13:00 in  
450 presence of cloud cover. The method based on maximum of  $\sigma_{RCS'}^2$  locates the  $PBLH_{Elastic}$  at the cloud base,  
451 due to the high variance of  $RCS'$  generated by the clouds.  $\tau_{RCS'}$  presents values larger than  $EL$  time  
452 acquisition, therefore this configuration enable us to study turbulence by  $EL$  analyses.  $S_{RCS'}$  has few peaks,  
453 due to the mixing between  $CBL$  and dust layer, generating a more homogenous layer. The highest values  
454 of  $S_{RCS'}$  are observed in regions where there are clouds, and the negative ones (between 3500 and 4000 m)  
455 occur due to presence of air from  $FT$  between the two aerosol layers (Figure 15). The inflection point of  
456  $S_{RCS'}$  profile is observed in  $PBLH_{MWR}$  region.  $K_{RCS'}$  profile has low values in most of the  $PBL$ , demonstrating  
457 the high level of mixing during this period, where dust layer and  $PBL$  are combined. The higher values of  
458  $K_{RCS'}$  are observed in the region of clouds. In the same way of the previous analysis, the high-order moments  
459 of the period mentioned above were calculated for the wavelength of 1064 nm (figure 19). Although there

460 are some differences in the absolute values of some profiles, the high-order moments generated using 1064  
461 and 532 nm have similar profiles, so that, the same phenomena can be observed, demonstrating the viability  
462 of 532 nm wavelength in the proposed methodology.

## 463 **5 Conclusions**

464 In this paper we perform an analysis about the *PBL* turbulent features from three different types of remote  
465 sensing systems (*DL*, *EL* and *MWR*) and surface sensors during SLOPE-I campaign. We applied two kind  
466 of corrections to the lidar data: first lag and  $-2/3$  corrections. The corrected *DL* statistical moments showed  
467 little variation with respect to the uncorrected profiles, denoting a rather low influence of the noise in high  
468 *SNR* conditions. The *EL* high-order moments were obtained from two wavelengths: 1064 nm, adopted as  
469 reference, and 532 nm, in order to verify the viability to use the last one in turbulence analysis. From this  
470 comparison, was possible to observe that the wavelength 532 nm is more affected by noise, in comparison  
471 with 1064 nm, due to the large contribution of the molecular component and the reduced two-way  
472 transmittance at that wavelength. However, the application of proposed corrections, mainly the first lag,  
473 can reduce such influence, so that, the same phenomena can be observed in the high-order moments  
474 provided from both wavelengths

475 The case studies present two kind of situations: well-defined *PBL* and a more complex situation with the  
476 presence of Saharan dust layer and some clouds. In both cases was possible to identify the events describe  
477 in table 2. The-combined use of remote sensing systems shows how the results provided by the different  
478 instruments can complement one each other, providing a detailed observation of some phenomena, mainly  
479 in complex situations.

480 Therefore, this study shows the feasibility of the described methodology based on the combination of  
481 remote sensing systems for retrieving a detailed picture on the *PBL* turbulent features. In addition, the  
482 feasibility of using the analyses of high order moments of the *RCS* collected at 532 nm at a temporal  
483 resolution of 2 s offers the possibility for using the proposed methodology in networks such as EARLINET  
484 or LALINET with a reasonable additional effort.

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493 **References**

- 494 Alados, I., Foyo-Moreno, I., Olmo, F. J., Alados-Arboledas, L. Relationship between net radiation and solar  
495 radiation for semi-arid shrub-land. *Agr. Forest Meteorol.*, 116, 221-227, 2003.
- 496 Albrecht, B. A., Bretherton, C. S., Johnson, D., Scubert, W. H., and Frisch, A. S.: The Atlantic  
497 stratocumulus transition experiment—ASTEX, *Bull. Am. Meteorol. Soc.*, 76, 889–904, 1995.
- 498 Andrews, E., Sheridan, P. J., Ogren, J. A., and Ferrare, R.: In situ aerosol profiles over the Southern Great  
499 Plains cloud and radiation test bed site: 1. Aerosol optical properties, *J. Geophys. Res.*, 109, D06208,  
500 doi:10.1029/2003JD004025, 2004.
- 501 Ansmann, A., Fruntke, J., Engelmann, R. Updraft and downdraft characterization with Doppler lidar: cloud-  
502 free versus cumuli-topped mixed layer. *Atmos. Chem. Phys.*, 10, 7845-7858, 2010.
- 503 Antón, M., Valenzuela, A., Cazorla, A., Gil, J. E., Gálvez-Fernández, J., Lyamani, H., Foyo-Moreno, I.,  
504 Olmo, F. J., Alados-Arboledas, L. Global and diffuse shortwave irradiance during a strong desert dust  
505 episode at Granada (Spain). *Atmos. Res.*, 118, 232 – 239, 2012.
- 506 Bedoya-Velásquez, A. E., Navas-Guzmán, F., Granados-Muñoz, M. J., Titos, G., Román, R., Casquero-  
507 Vera, J. A., Ortiz-Amezcuca, P., Benavent-Oltra, J. A., Moreira, G. de A., Montilla-Rosero, E., Ortiz, C. D.  
508 H., Artiñano, Coz, E., Alados-Arboledas, L., Guerrero-Rascado, J. L. Hygroscopic growth study in the  
509 framework of EARLINET during the SLOPE I campaign: synergy of remote sensing and in-situ  
510 instrumentation. *Atmos. Chem. Phys.*, 18, 7001-7017, 2017.
- 511 Behrendt, A., Wulfmeyer, V., Hammann, E., Muppa, S. K., Pal, S.: Profiles of second- to  
512 fourth-order moments of turbulent temperature fluctuations in the convective boundary layer : first  
513 measurements with rotational Raman lidar. *Atmos. Chem. Phys.*, 15, 5485–5500.  
514 <https://doi.org/10.5194/acp-15-5485-2015>, 2015.
- 515 Bravo-Aranda, J. A., Navas-Guzmán, F., Guerrero-Rascado, J. L., Pérez-Ramírez, D., Granados-Muñoz,  
516 M. J., Alados-Arboledas, L. Analysis of lidar depolarization calibration procedure and application to the  
517 atmospheric aerosol characterization. *Int. J. Remote Sens.*, 34 (9-10), pp. 3543-3560, 2013.
- 518 Caumont, O., Cimini, D., Löhnert, U., Alados-Arboledas, L., Bleisch, R., Buffa, F., Ferrario, M.E., Haeefe,  
519 A., Huet, T., Madonna, F., Pace, G. Assimilation of humidity and temperature observations retrieved from  
520 ground-based microwave radiometers into a convective-scale NWP model. *Q. J. Roy. Meteor. Soc.*, 142  
521 (700), pp. 2692-2704, 2016.
- 522 Engelmann, R.; Wandinger, U.; Ansmann, A.; Müller, D.; Žeromskis, E.;  
523 Althausen, D.; Wehner, B. Lidar Observations of the Vertical Aerosol Flux in the Planetary  
524 Boundary Layer. *J. Atmos. Ocean. Tech.*, v. 25, n. 8, p. 1296–1306, 2008.
- 525 Guerrero-Rascado, J.L., Ruiz, B., Alados-Arboledas, L. Multi-spectral lidar characterization of the vertical  
526 structure of Saharan dust aerosol over Southern Spain. *Atmos. Environ.*, 42, 2668-2681, 2008.



527 Guerrero-Rascado, J.L., Olmo, F.J., Avilés-Rodríguez, I., Navas-Guzmán, F., Pérez-Ramírez, D., Lyamani,  
528 H., Alados-Arboledas, L.: Extreme Saharan dust event over the southern Iberian Peninsula in September  
529 2007: Active and passive remote sensing from surface and satellite. *Atmos. Chem. Phys.*, 9, 21, 8453-8469,  
530 2009.

531 Guerrero-Rascado, J. L., Costa, M. J., Bortoli, D., Silva, A. M., Lyamani, H., Alados-Arboledas, L. Infrared  
532 lidar overlap function: an experimental determination, *Opt. Express*, 18, 20350-20359, 2010

533 Guerrero-Rascado, J. L., Landulfo, E., Antuña, J. C., Barbosa, H. M. J., Barja, B., Bastidas, A. E., Bedoya,  
534 A. E., da Costa, R. F., Estevan, R., Forno, R. N., Gouveia, D. A., Jimenez, C., Larroza, E. G., Lopes, F. J.  
535 S., Montilla-Rosero, E., Moreira, G. A., Nakaema, W. M., Nisperuza, D., Alegria, D., Múnera, M., Otero,  
536 L., Papandrea, S., Pawelko, E., Quel, E. J., Ristori, P., Rodrigues, P. F., Salvador, J., Sánchez, M. F., and  
537 Silva, A.: Latin American Lidar Network (LALINET) for aerosol research: diagnosis on network  
538 instrumentation, *J. Atmos. Sol.-Terr. Phy.*, 138–139, 112–120, 2016.

539 Illingworth, A. J., Hogan, R. J. O' Connor, E. J. Bouniol, D. Brooks, M. E. Delanoe, J. Donovan, D. P.  
540 Eastment, J. D. Gaussiat, N. Goddard, J. W. F. Haeffelin, M. Klein Baltink, H. Krasnov, O. A. Pelon, J.  
541 Piriou, J.-M. Protat, A. Russchenberg, H. W. J. Seifert, A. Tompkins, A. M. Van Zadelhoff, G.-J. Vinit, F.  
542 Willen, U. Wilson, D. R. and Wrench, C. L.: CLOUDNET: Continuous Evaluation of Cloud Profiles in  
543 Seven Operational Models using Ground-Based Observations. *Bull. Am. Meteorol. Soc.*, 88, 883-898,  
544 doi:10.1175/BAMS-88-6-883, 2007.

545 Kaimal, J. C., and Gaynor, J. E.: The Boulder Atmospheric Observatory, *J. Clim. Appl. Meteorol.*, 22, 863–  
546 880, 1983.

547 Kiemle, C., Brewer, W. A., Ehret, G., Hardesty, R. M., Fix, A., Senff, C., Wirth, M., Poberaj, G., and  
548 LeMone, M. A.: Latent heat flux profiles from collocated airborne water vapor and wind lidars during IHOP  
549 2002, *J. Atmos. Ocean. Tech.*, 24, 627–639, 2007.

550 Klett, J. D.: Lidar inversion with variable backscatter/extinction ratios, *Appl. Opt.*, 24, 1638-1643,  
551 <https://doi.org/10.1364/AO.24.001638>, 1985.

552 Kovalev, V.A., Eichinger, W.E., *Elastic Lidar*, Wiley 2004.

553 Lenschow, D. H., Wyngaard, J. C., and Pennell, W. T.: Mean-field and second-moment budgets in a  
554 baroclinic convective boundary layer, *J. Atmos. Sci.*, 37, 1313–1326, 1980.

555 Lenschow, D. H., Mann, J., and Kristensen, L.: How long is long enough when measuring fluxes and other  
556 turbulence statistics?, *J. Atmos. Oceanic Technol.*, 11, 661–673, 1994.

557 Lenschow, D. H., Wulfmeyer, V. and Senff, C.: Measuring second- through fourth-order moments in noisy  
558 data, *J. Atmos. Oceanic Technol.*, 17, 1330–1347, 2000.

559 Lenschow, D. H., Lothon, M., Mayor, S. D., Sullivan, P. P., and Canut, G.: A comparison of higher-order  
560 vertical velocity moments in the convective boundary layer from lidar with in situ measurements and large-  
561 eddy simulation, *Bound-Lay. Meteorol.*, 143, 107–123, doi:10.1007/s10546-011-9615-3., 2012.

562 Lothon, M., Lenschow, D. H., and Mayor, S. D.: Coherence and scale of vertical velocity in the convective  
563 boundary layer from a Doppler lidar, *Bound.-Lay. Meteorol.*, 121, 521–536, 2006.

564 Lyamani, H., Olmo, F. J., Alcántara, A., and Alados-Arboledas, L.: Atmospheric aerosols during the 2003  
565 heat wave in southeastern Spain I: Spectral optical depth, *Atmos. Environ.*, 40, 6453–6464, 2006

566 McNicholas, C., Turner, D. D. Characterizing the convective boundary layer turbulence with a High  
567 Spectral Resolution Lidar. *J. Geophys Res-Atmos.*, v. 119, p. 910–927, 2014.

568 Muppa, K.S., Behrendt, A., Späth, F., Wulfmeyer, V., Metzendorf, S., Riede, A.: Turbulent humidity  
569 fluctuations in the convective boundary layer: Cases studies using water vapour differential absorption lidar  
570 measurements. *Bound-Lay. Meteorol.*, 158, 43-66, DOI 10.1007/s10546-015-0078-9, 2014.

571 Monin, A. S., Yaglom, A. M. *Statistical Fluid Mechanics*, Vol. 2. MIT Press, 874 pp, 1979.

572 Moreira, G. de A., Marques, M. T. A., Nakaema, W., Moreira, A. C. de C. A., Landulfo, E. Planetary  
573 boundary height estimations from Doppler wind lidar measurements, radiosonde and hysplit model  
574 comparisom. *Óptica Pura y Aplicada*, 48, 179-183, 2015

575 Moreira, G. de A., Guerrero-Rascado, J. L., Bravo-Aranda, J. A., Benavent-Oltra, Ortiz-Amezcuca, P.,  
576 Róman, R., Bedoya-Velásquez, A., Landulfo, E., Alados-Arboledas, L. Study of the planetary boundary  
577 layer by microwave radiometer, elastic lidar and Doppler lidar estimations in Southern Iberian Peninsula.  
578 *Atmos. Res.*, 213, 185-195, 2018a.

579 Moreira, G. de A., Lopes, F. J. S., Guerrero-Rascado, J. L., Landulfo, E., Alados-Arboledas, L. Analyzing  
580 turbulence in Planetary Boundary Layer from multiwavelength lidar system: impact of wavelength choice.  
581 *Opt. Express*. Under review, 2018b.

582 Navas Guzmán, F., Guerrero Rascado, J. L., and Alados Arboledas, L.: Retrieval of the lidar overlap  
583 function using Raman signals, *Óptica Pura y Aplicada*, 44, 71–75, 2011.

584 Navas-Guzmán, F., Bravo-Aranda, J.A., Guerrero-Rascado, J.L, Granados-Muñoz, M.J, and Alados-  
585 Arboledas, L.: Statistical analysis of aerosol optical properties retrieved by Raman lidar over Southeastern  
586 Spain. *Tellus B*, 65, 21234, 2013.

587 Navas-Guzmán, F., Fernández-Gálvez, J., Granados-Muñoz, M.J, Guerrero-Rascado, J.L., Bravo-Aranda,  
588 J.A., and Alados-Arboledas, L.: Tropospheric water vapor and relative humidity profiles from lidar and  
589 microwave radiometry. *Atmos. Meas. Tech.*, 7, 1201-1211, 2104.

590 O'Connor, E. J., Illingworth, A. J., Brooks, I. M., Westbrook, C. D., Hogan, R. J., Davies, F., Brooks, B.  
591 J.: A method for estimating the turbulent kinetic energy dissipation rate from a vertically-pointing Doppler  
592 lidar, and independent evaluation from balloon-borne in-situ measurements. *J. Atmos. Ocean. Tech.*, v. 27,  
593 n. 10, 1652-1664, 2010.

594 Ortiz-Amezcuca, P., Guerrero-Rascado, J.L., Granados-Muñoz, M.J., Bravo-Aranda, J. A., Alados-  
595 Arboledas, L. Characterization of atmospheric aerosols for a long range transport of biomass burning

596 particles from canadian forest fires over the southern iberian peninsula in july 2013. *Optica Pura y Aplicada*,  
597 47 (1), pp. 43-49, 2014.

598 Ortiz-Amezcuca, P., Luis Guerrero-Rascado, J., Granados-Munõz, M.J., Benavent-Oltra, J.A., Böckmann,  
599 C., Samaras, S., Stachlewska, I.S., Janicka, L., Baars, H., Bohlmann, S., Alados-Arboledas, L.  
600 Microphysical characterization of long-range transported biomass burning particles from North America at  
601 three EARLINET stations. *Atmos. Chem. Phys.*, 17 (9), pp. 5931-5946, 2017.

602 Pal, S.; Behrendt, a.; Wulfmeyer, V. Elastic-backscatter-lidar-based characterization of the  
603 convective boundary layer and investigation of related statistics. *Ann. Geophys.*, v. 28, n. 3, p. 825–  
604 847, 2010.

605 Pappalardo, G., Amodeo, A., Apituley, A., Comeron, A., Freudenthaler, V., Linné, H., Ansmann, A.,  
606 Bösenberg, J., D'Amico, G., Mattis, I., Mona, L., Wandinger, U., Amiridis, V., Alados-Arboledas, L.,  
607 Nicolae, D., and Wiegner, M.: EARLINET: towards an advanced sustainable European aerosol lidar  
608 network. *Atmos. Meas. Tech.*, 7, 2389-2409, doi:10.5194/amt-7-2389-2014, 2014.

609 Román, R., Benavent-Oltra, J. A., Casquero-Vera, J. A., Lopatin, A., Cazorla, A., Lyamani, H., Denjean,  
610 C., Fuertes, D., Pérez-Ramirez, D., Torres, B., Toledano, C., Dubovik, O., Cachorro, V. E., Frutos, A. M.,  
611 Olmo, F. J., Alados-Arboledas, L. Retrieval of aerosol profiles combining sunphotometer and ceilometer  
612 measurements in GRASP code. *Atmos. Res.*, 204, 161, 177, 2018.

613 Rose, T., Creewll, S., Löhnert, U., Simmer, C.: A network suitable microwave radiometer for operational  
614 monitoring of cloudy atmosphere. *Atmos. Res.*, 75, 3, 183 – 200, 2005.

615 Stull, R. B., Santoso, E., Berg, L., Hacker, J.: Boundary layer experiment 1996 (BLX96), *Bull. Am.*  
616 *Meteorol. Soc.*, 78, 1149–1158, 1997.

617 Stull, R. B.: An Introduction to Boundary Layer Meteorology, vol. 13, *Kluwer Academic Publishers, the*  
618 *Netherlands*, Dordrecht/Boston/London, 1988.

619 Stull, R. B.: *Meteorology for Scientists and Engineers*, 3<sup>rd</sup> Edition, *Uni. Of British Columbia*, 2011.

620 Titos, G., Foyo-Moreno, I., Lyamani, H., Querol, X., Alastuey, A., and Alados-Arboledas, L.: Optical  
621 properties and chemical composition of aerosol particles at an urban location: An estimation of the aerosol  
622 mass scattering and absorption efficiencies, *J. Geophys.Res.-Atmos.*, 117, D04206,  
623 doi:10.1029/2011JD016671, 2012.

624 Turner, D. D., Ferrare, R. A., Wulfmeyer, V., and Scarino, A. J.: Aircraft evaluation of ground-based  
625 Raman lidar water vapor turbulence profiles in convective mixed layers, *J. Atmos. Oceanic Technol.*, 31,  
626 1078–1088, doi:10.1175/JTECH-D-13-00075-1, 2014

627 vanUlden, A. P., and Wieringa, J. :Atmospheric boundary layer research at Cabauw, Bound-Lay. *Meteorol.*,  
628 78,39–69, 1996.

629 Valenzuela, A., Olmo, F.J.ab, Lyamani, H.ab, Granados-Muñoz, M.J.ab, Antón, M.c, Guerrero-Rascado,  
630 J.L.ab, Quirantes, A.a, Toledano, C.d, Perez-Ramírez, D.ef, Alados-Arboledas, L.: Aerosol transport over

631 the western mediterranean basin: Evidence of the contribution of fine particles to desert dust plumes over  
632 alborán island. *J. Geophys Res.*, 119, 24, 14,028-14,044, 2014.

633 Vogelmann, A. M., McFarquhar, G. M., Ogren, J. A., Turner, D. D., Comstock, J. M., Feingold, G., Long,  
634 C. N., Jonsson, H. H., Bucholtz, A., Collins, D. R., Diskin, G. S., Gerber, H., Lawson, R. P., Woods, R. K.,  
635 Andrews, E., Yang, H., Chiu, J. C., Hartsock, D., Hubbe, J. M., Lo, C., Marshak, A., Monroe, J. W.,  
636 Mcfarlane, S. A., Jason, M., and Toto, T.: RACORO extended-term aircraft observations of boundary layer  
637 clouds, *Bull. Am. Meteorol. Soc.*, 93, 861–878, doi:10.1175/BAMS-D-11-00189.1, 2012.

638 Williams, A. G., and Hacker, J. M.: The composite shape and structure of coherent eddies in the convective  
639 boundary layer, *Bound-Lay. Meteorol.*, 61, 213–245, 1992.

640 Wulfmeyer, V.: Investigation of turbulent processes in the lower troposphere with water vapor DIAL and  
641 radar-RASS, *J. Appl. Sci.*, 56, 1055–1076, 1999.

642 Wulfmeyer, V., Pal, S., Turner, D. D., and Wagner, E.: Can water vapour Raman lidar resolve profiles of  
643 turbulent variables in the convective boundary layer?, *Bound-Lay. Meteorol.*, 136, 253–284,  
644 doi:10.1007/s10546-010-9494-z, 2010.

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Table 1 – Variables applied to statistical analysis (Lenschow et al., 2000)

	Without Correction	Correction	Error
<b>Integral Time Scale (<math>\tau</math>)</b>	$\int_0^{\infty} q'(t) dt$	$\frac{1}{\overline{q'^2}} \int_{t \rightarrow 0}^{\infty} M_{11}(t) dt$	$\tau \cdot \sqrt{\frac{4\Delta M_{11}}{M_{11}(\rightarrow 0)}}$
<b>Variance (<math>\sigma_q^2</math>)</b>	$\frac{1}{T} \sum_{t=1}^T (q(t) - \bar{q})^2$	$M_{11}(\rightarrow 0)$	$q^2 \cdot \sqrt{\frac{4\Delta M_{11}}{M_{11}(\rightarrow 0)}}$
<b>Skewness (<math>S</math>)</b>	$\frac{\overline{q^3}}{\sigma_q^3}$	$\frac{M_{21}(\rightarrow 0)}{M_{11}^{3/2}(\rightarrow 0)}$	$\frac{\Delta M_{21}}{\Delta M_{11}^{3/2}}$
<b>Kurtosis (<math>K</math>)</b>	$\frac{\overline{q^4}}{\sigma_q^4}$	$\frac{3M_{22}(\rightarrow 0) - 2M_{31}(\rightarrow 0) - 3\Delta M_{11}^2}{M_{11}^2(\rightarrow 0)}$	$\frac{4\Delta M_{31} - 3\Delta M_{22} - \Delta M_{11}^2}{\Delta M_{11}^2}$

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Table 2 – Products and their respective meaning, provided by each system

Product	System	Meaning
$\tau_{w'}(z)$	Doppler lidar	Measurement in time of length of turbulent eddies
$\sigma_{w'}^2(z)$	Doppler lidar	Turbulent Kinetic Energy
$S_{w'}(z)$	Doppler lidar	Direction of turbulent movements
$PBLH_{Doppler}$	Doppler lidar	Top of CBL obtained from variance threshold method
$\tau_{RCS'}(z)$	Elastic lidar	Measurement in time of length of turbulent eddies
$\sigma_{RCS'}^2(z)$	Elastic lidar	Homogeneity of aerosol distribution
$S_{RCS'}(z)$	Elastic lidar	Aerosol motion ( $S < 0 \rightarrow$ Downdrafts, $S > 0 \rightarrow$ Updrafts)
$K_{RCS'}(z)$	Elastic lidar	Level of aerosol mixing ( $K < 3 \rightarrow$ Well-Mixed, $K > 3 \rightarrow$ Low Mixing)
$PBLH_{Elastic}$	Elastic lidar	Top of aerosol layer obtained from variance method
$PBLH_{MWR}$	MWR	Top of CBL/SBL layer obtained from Potential Temperature

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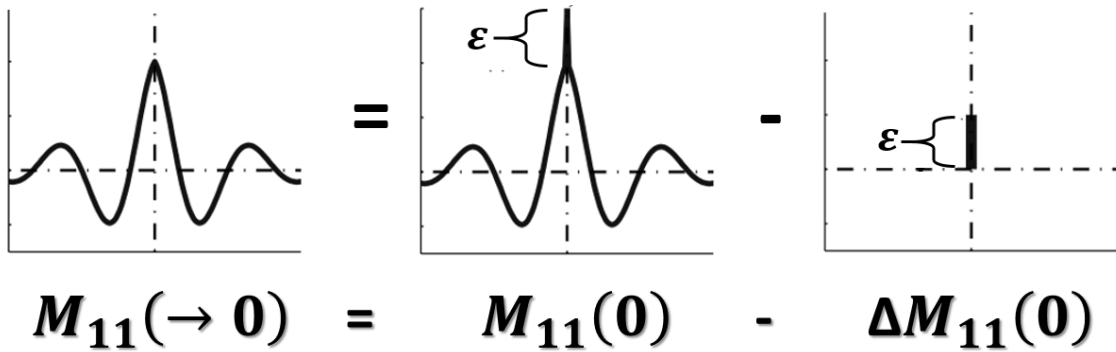


Figure 1 – Procedure to remove the errors of autocovariance functions.  $M_{11}(\rightarrow 0)$  – corrected autocovariance function errors;  $M_{11}(0)$  - autocovariance function without correction;  $\Delta M_{11}(0)$  - error of autocovariance function.

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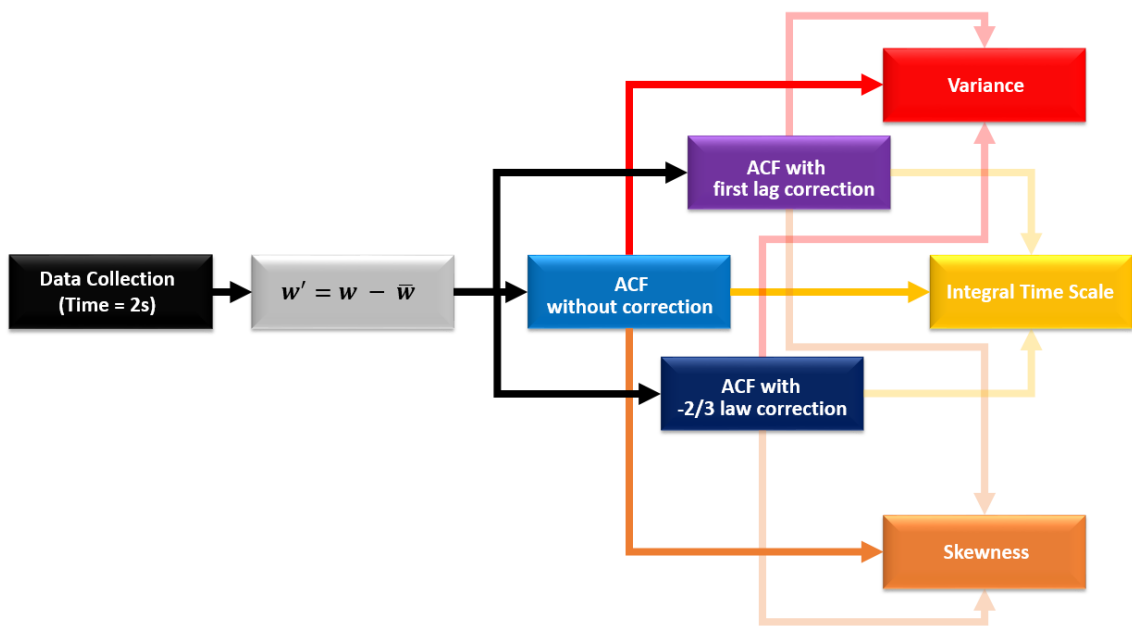


Figure 2 – Flowchart of data analysis methodology applied to the study of turbulence with Doppler lidar

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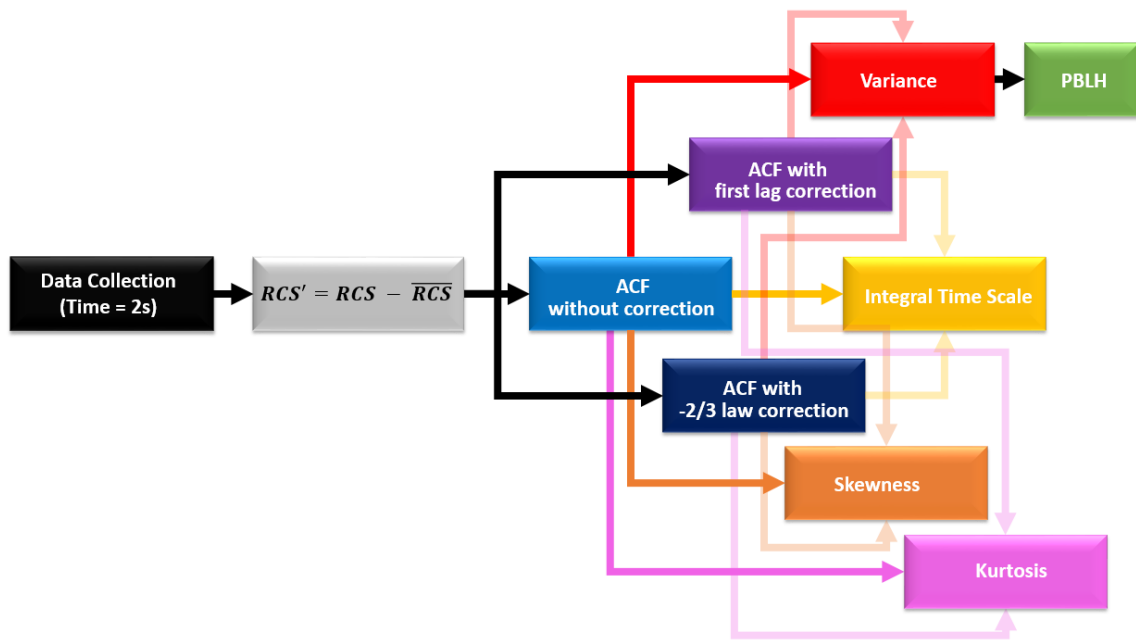


Figure 3 – Flowchart of data analysis methodology applied to the study of turbulence with elastic lidar

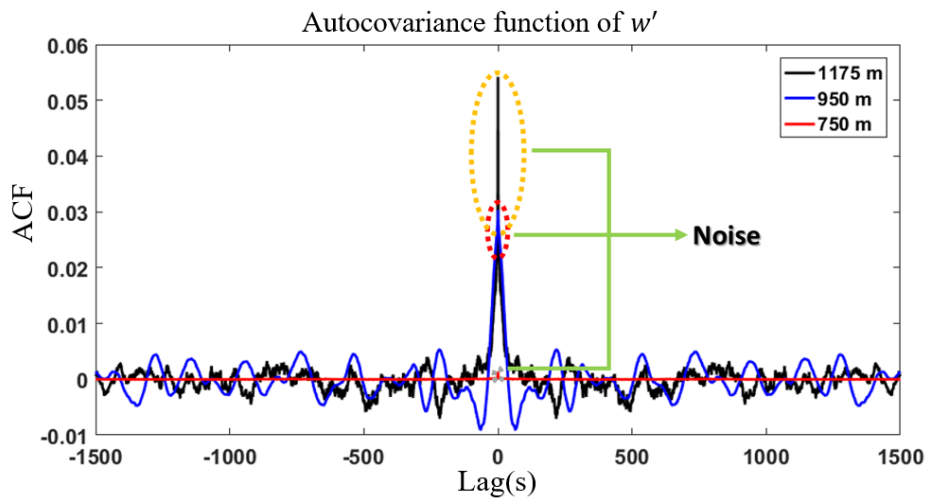


Figure 4 – Autocovariance function (ACF) of  $w'$ , obtained from Doppler lidar at three different heights on 19th May 2016 at 08-09 UTC in Granada.

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### Profiles obtained from $w'$ - Granada – 19 May 2016 – 08-09 UTC

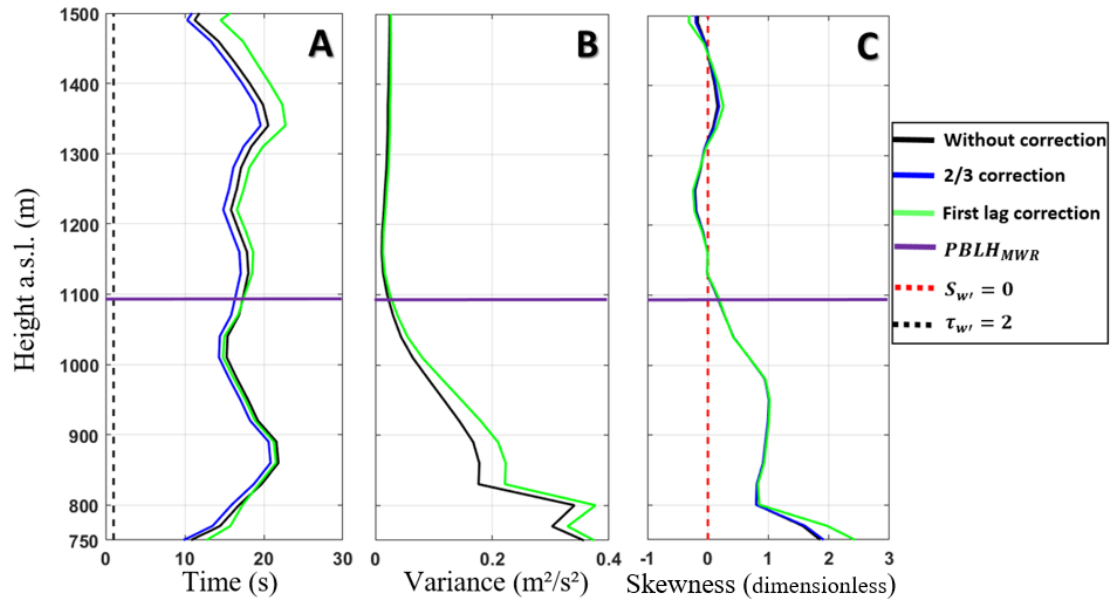


Figure 5 – A - Vertical profile of Integral time scale ( $\tau_{w'}$ ). B - Vertical profile of variance ( $\sigma_{w'}^2$ ). C - Vertical profile of Skewness ( $S_{w'}$ ). All profiles were obtained from Doppler lidar data on 19th May 2016 at 08-09 UTC in Granada.

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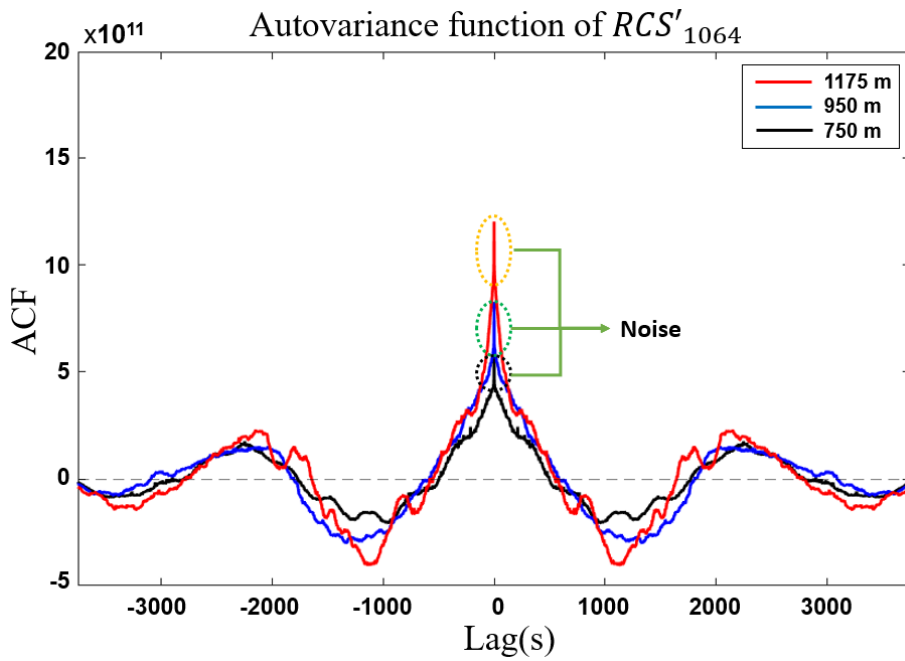


Figure 6 – Autocovariance of  $RCS'_{1064}$  obtained from MULHACÉN elastic lidar data to three different heights on 19th May 2016 at 12-13 UTC in Granada.

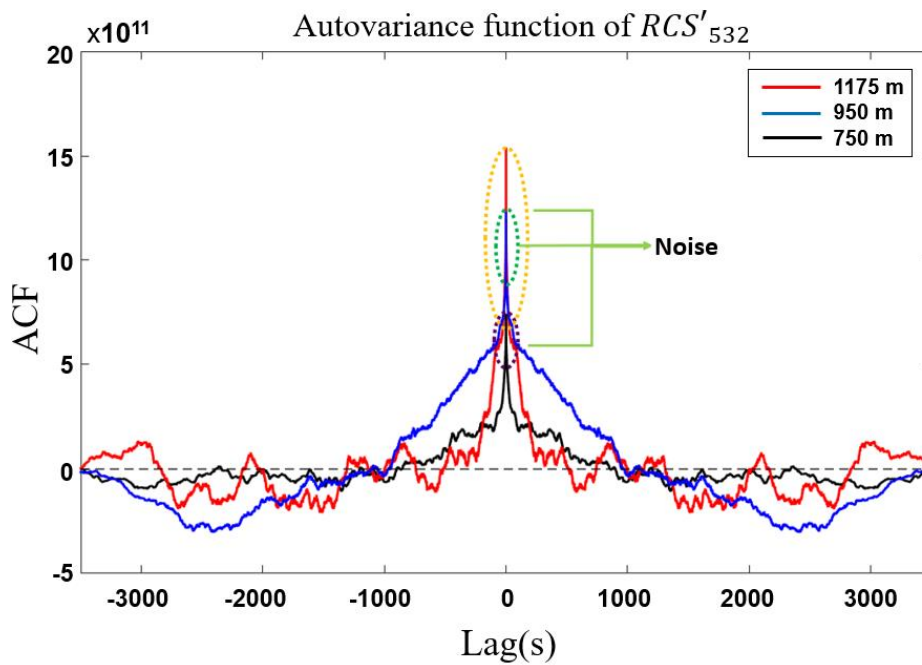


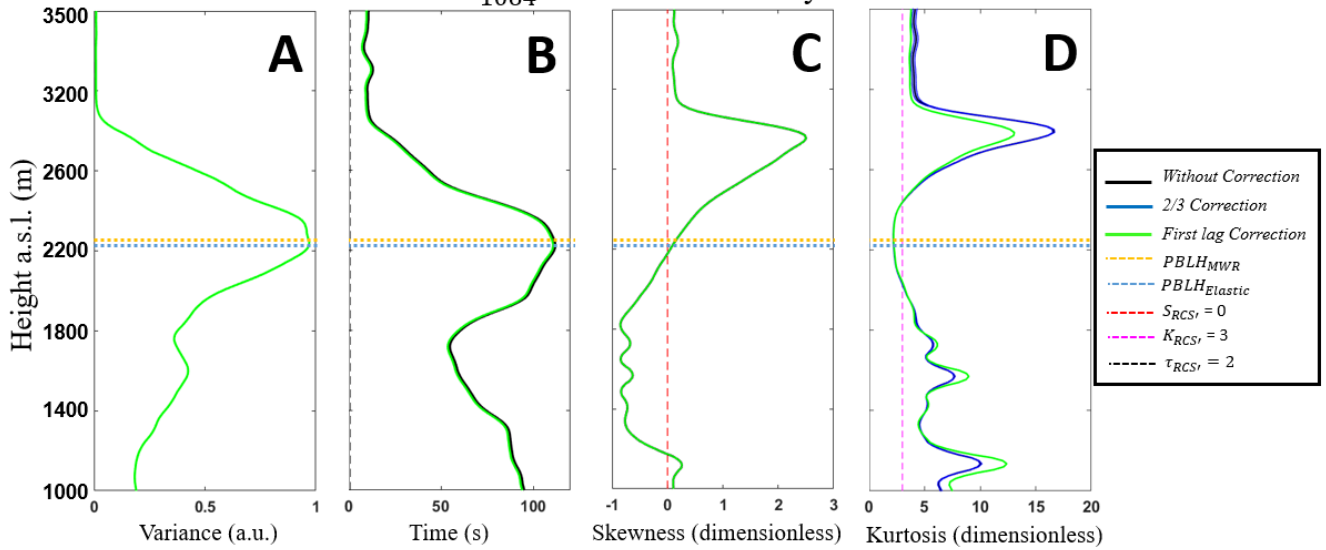
Figure 7 – Autocovariance of  $RCS'_{532}$  obtained from MULHACÉN elastic lidar data to three different heights on 19th May 2016 at 12-13 UTC in Granada.

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Profiles obtained from  $RCS'_{1064}$  – Granada – 19 May 2016 – 12-13 UTC



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Figure 8 – A- Vertical profile of Integral time scale ( $\tau_{RCS'}$ ). B - Vertical profile of variance ( $\sigma_{RCS'}^2$ ). C - Vertical profile of Skewness ( $S_{RCS'}$ ). D - Vertical profile of Kurtosis ( $K_{RCS'}$ ). All profiles were obtained from MULHACÉN elastic lidar data on 19th May2016 in Granada between 12-13 UTC.

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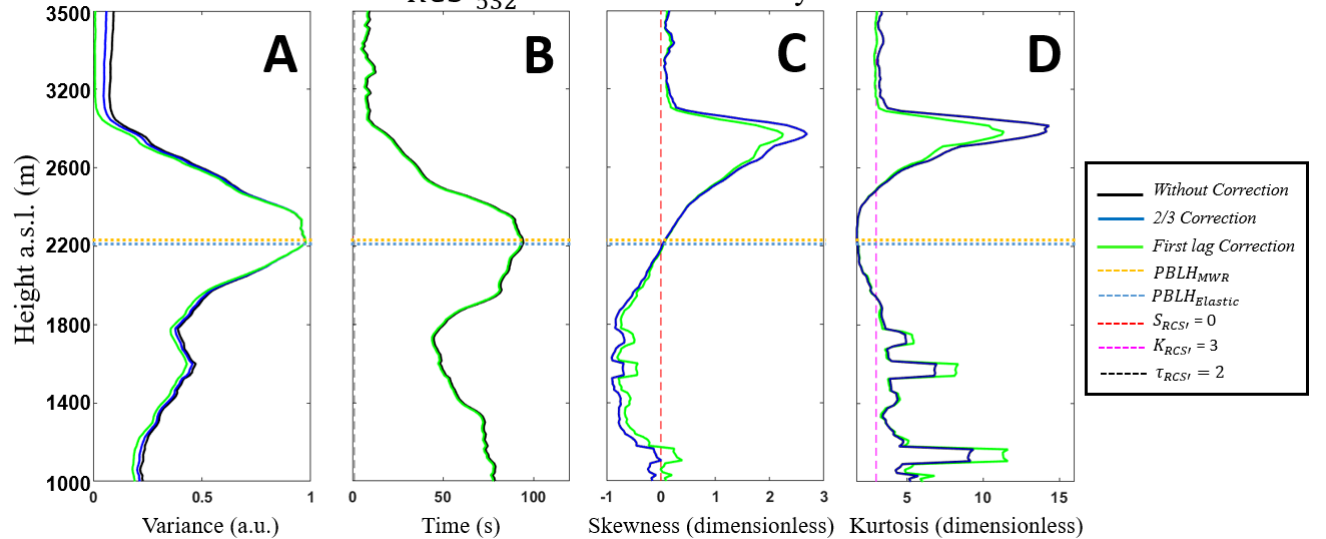
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Profiles obtained from  $RCS'_{532}$  – Granada – 19 May 2016 – 12-13 UTC



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Figure 9 – A- Vertical profile of Integral time scale ( $\tau_{RCS'}$ ). B - Vertical profile of variance ( $\sigma_{RCS'}^2$ ). C - Vertical profile of Skewness ( $S_{RCS'}$ ). D - Vertical profile of Kurtosis ( $K_{RCS'}$ ). All profiles were obtained from MULHACÉN elastic lidar data on 19th May2016 in Granada between 12-13 UTC.

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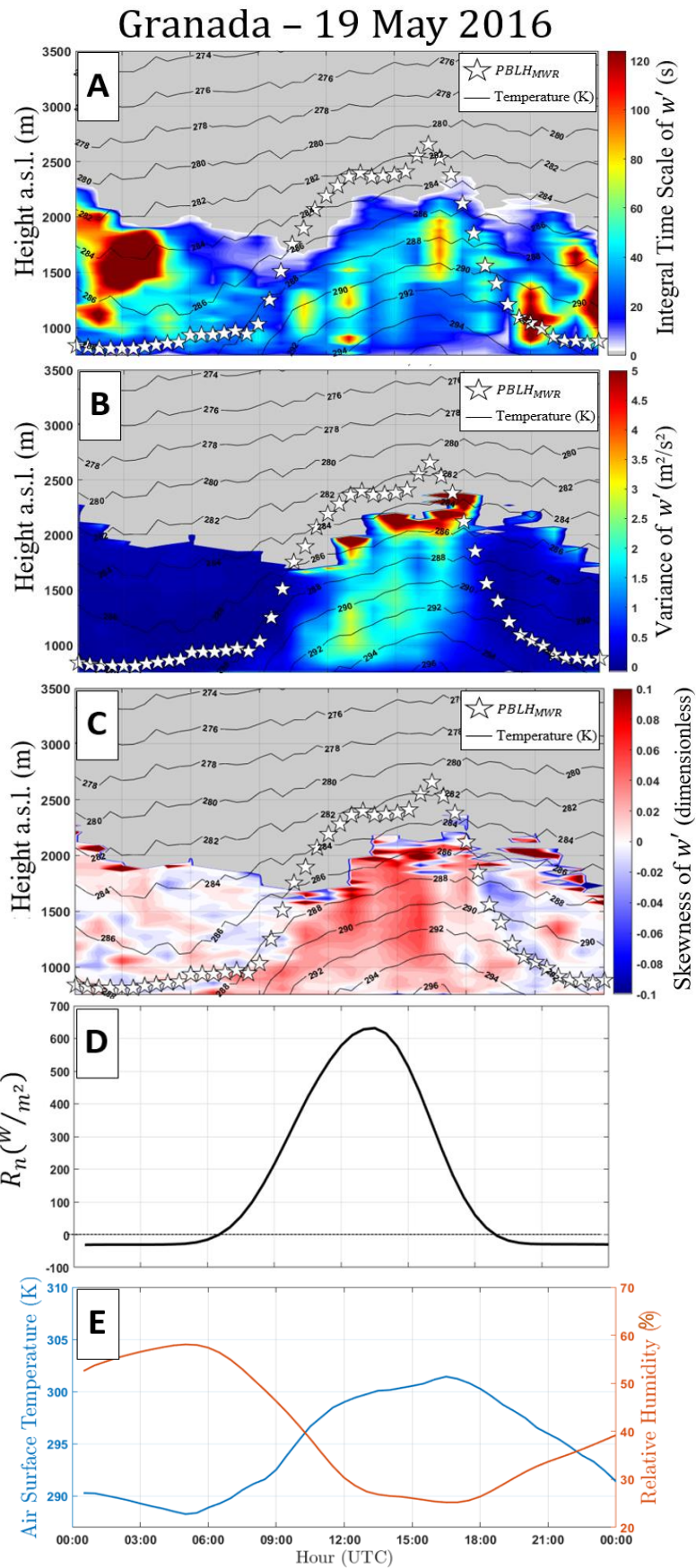


Figure 10 – A – integral time scale obtained from Doppler lidar data [ $\tau_{w'}$ ], B – variance obtained from Doppler lidar data [ $\sigma_{w'}^2$ ], C – skewness obtained from Doppler lidar data [ $S_{w'}$ ], D – net radiation obtained from pyranometer data [ $R_n$ ], E – Air surface temperature [blue line] and surface relative humidity [RH - orange line] both were obtained from surface sensors. All profiles were acquired on 19th May 2016 in Granada. In A, B and C black lines and white stars represent air temperature and  $PBLH_{MWR}$ , respectively.

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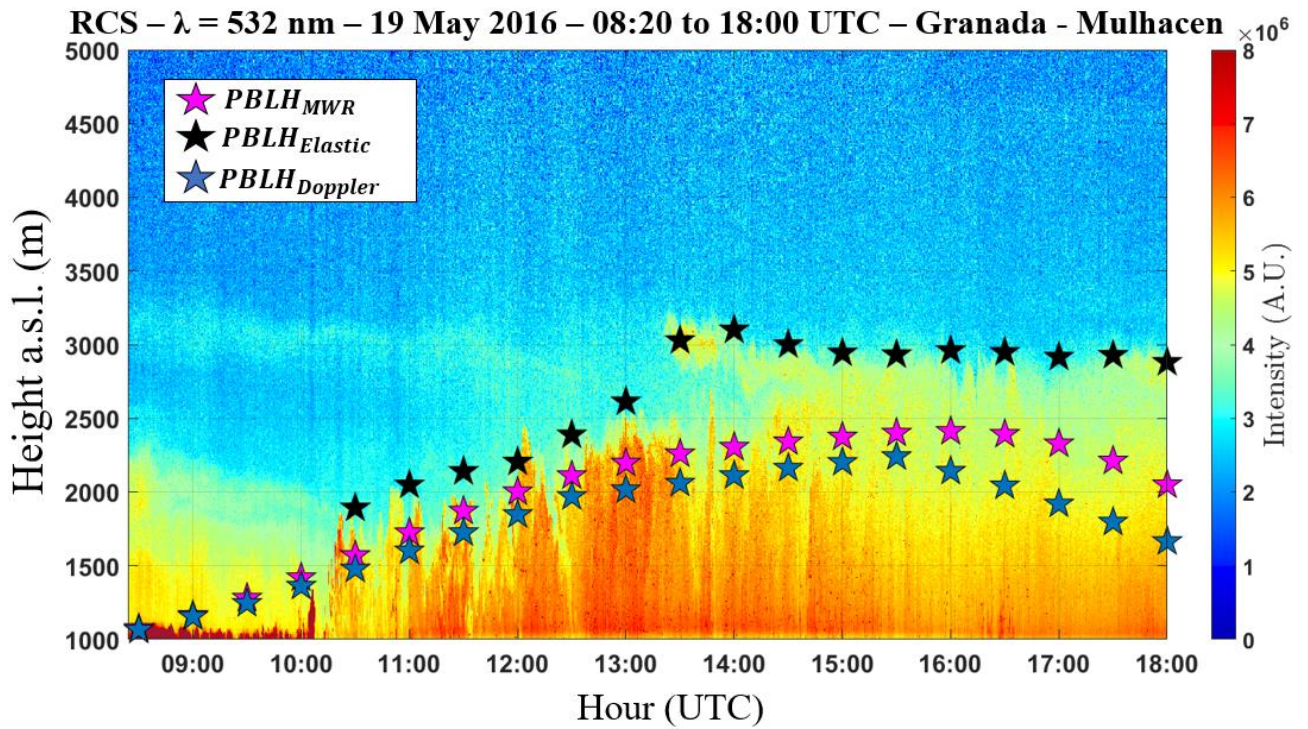


Figure 11 – Time-Height plot of RCS obtained on 19 May 2016 in Granada. Pink stars represent the  $PBLH_{MWR}$ , black stars represent the  $PBLH_{Elastic}$  and blue stars represent the  $PBLH_{Doppler}$ .

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## Granada – 19 May – 13-14 UTC

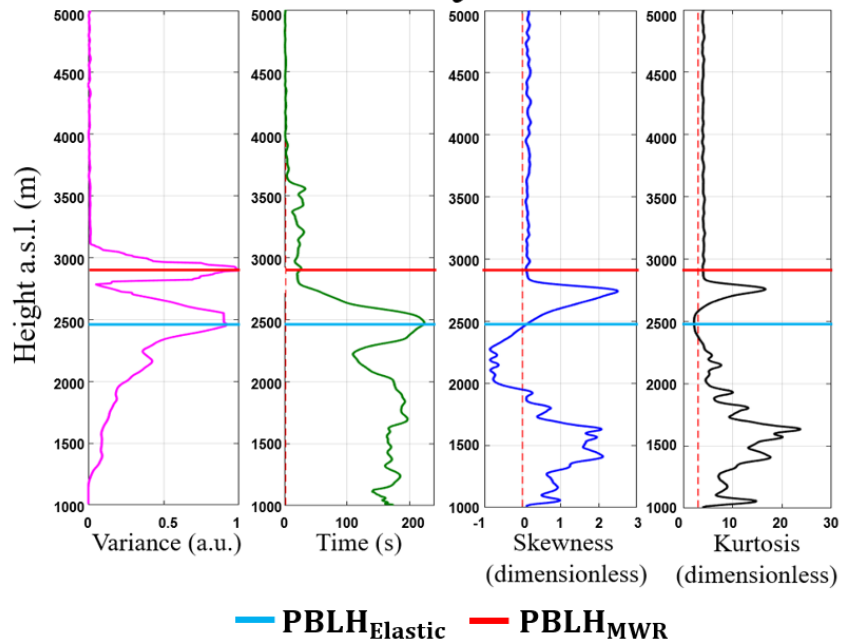


Figure 12 – Statistical moments obtained from 532 nm wavelength data of elastic lidar (MULHACÉN) in Granada at 13 to 14 UTC - 19 May 2016. From left to right: variance [ $\sigma_{RCS'}^2$ ], integral time scale [ $\tau_{RCS'}$ ], skewness [ $S_{RCS'}$ ] and kurtosis [ $K_{RCS'}$ ].



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## Granada – 19 May – 13-14 UTC

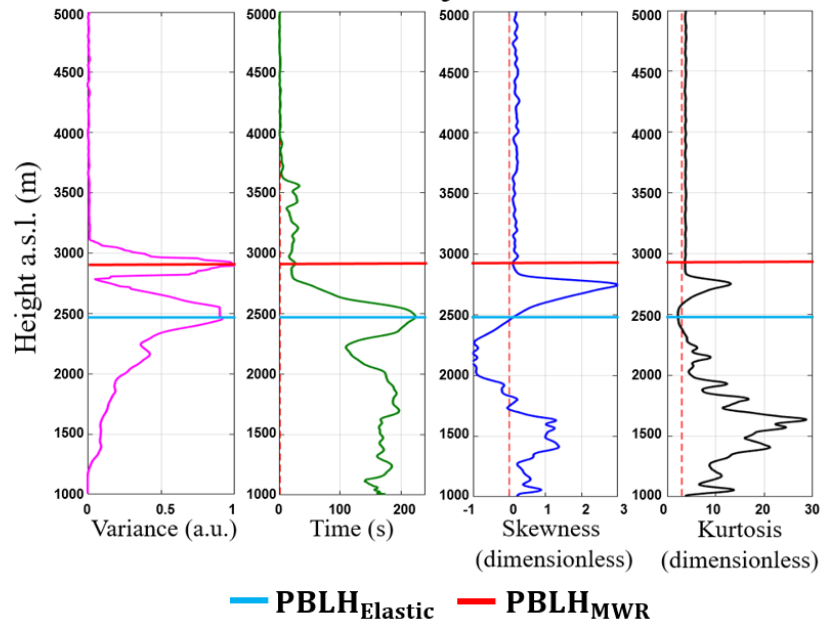


Figure 13 – Statistical moments obtained from 1064 nm wavelength data of elastic lidar(MULHACÉN) in Granada at 13 to 14 UTC - 19 May 2016. From left to right: variance [ $\sigma_{RCS'}^2$ ], integral time scale [ $\tau_{RCS'}$ ], skewness [ $S_{RCS'}$ ] and kurtosis [ $K_{RCS'}$ ].

## Granada - 08 Jul 2016

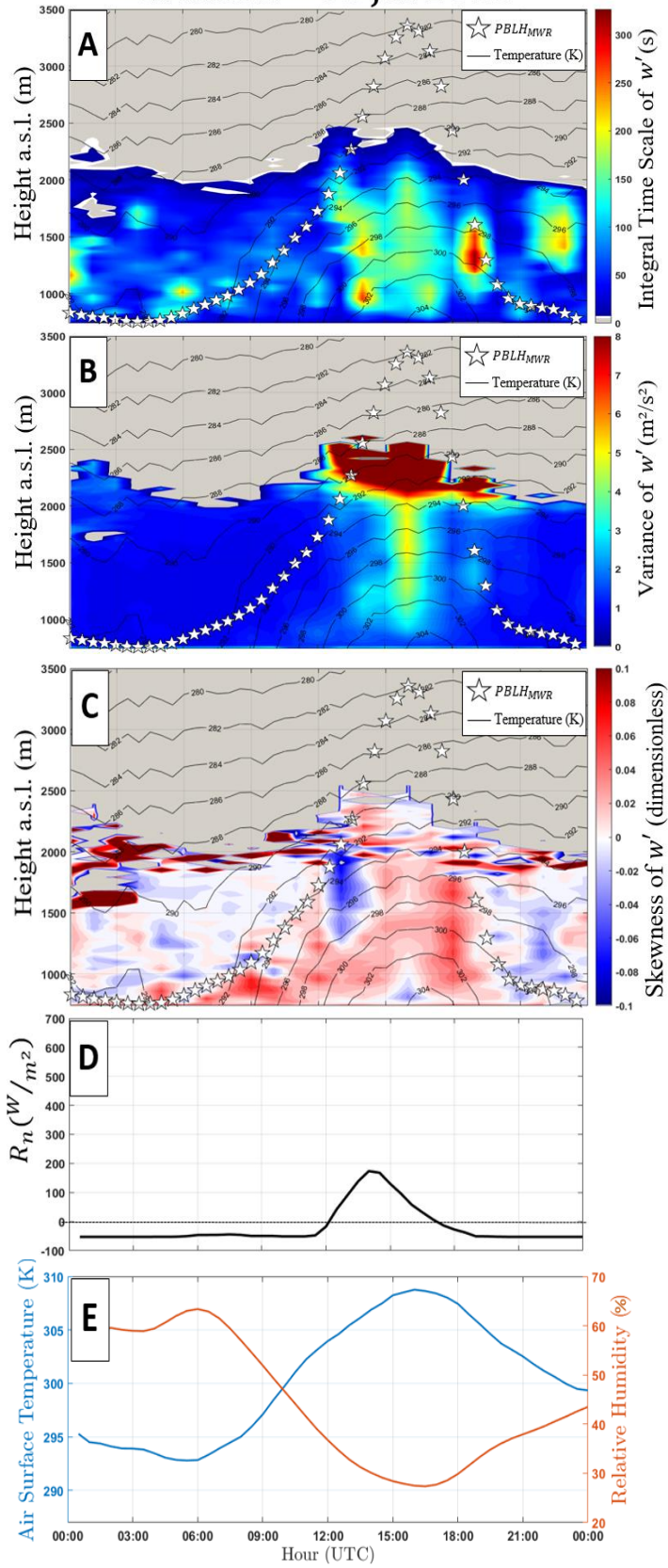


Figure 14 - A – integral time scale from Doppler lidar data [ $\tau_{w'}$ ], B – variance from Doppler lidar data [ $\sigma_{w'}^2$ ], C – skewness from Doppler lidar data [ $S_{w'}$ ], D – net radiation from pyranometer data [ $R_n$ ], E – Air surface temperature [blue line] and surface relative humidity [RH – orange line] from surface sensor data. All profiles were obtained in Granada on 08 July 2016. In A, B and C black lines and white stars represent air temperature and  $PBLH_{MWR}$ , respectively.



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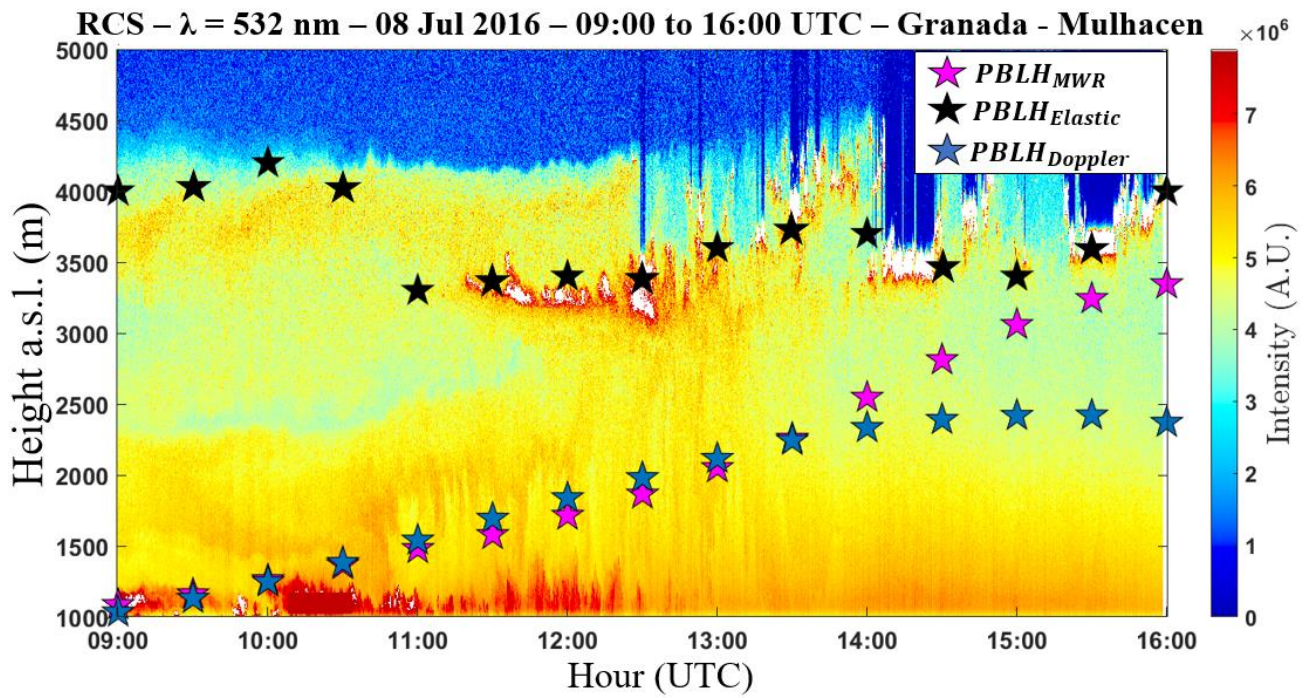


Figure 15 – Time-Height plot of RCS obtained from MULHACÉN elastic lidar data on 08 July 2016 in Granada. Pink stars represent the  $PBLH_{MWR}$ , black stars represent the  $PBLH_{Elastic}$  and blue stars represent the  $PBLH_{Doppler}$ .

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## Granada – 08 Jul – 11-12 UTC

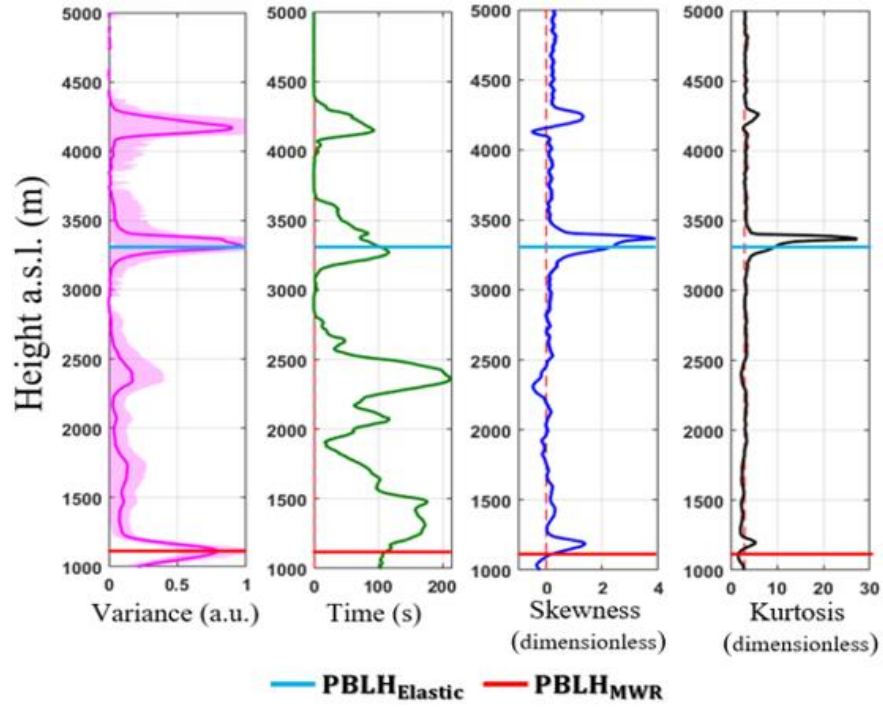


Figure 16 - Statistical moments obtained from 532 nm wavelength data of elastic lidar(MULHACÉN) in Granada between 11-12 UTC on 08th July 2016. From left to right: variance [ $\sigma_{RCS'}^2$ ], integral time scale [ $\tau_{RCS'}$ ], skewness [ $S_{RCS'}$ ] and kurtosis [ $K_{RCS'}$ ].

## Granada – 08 Jul – 11-12 UTC

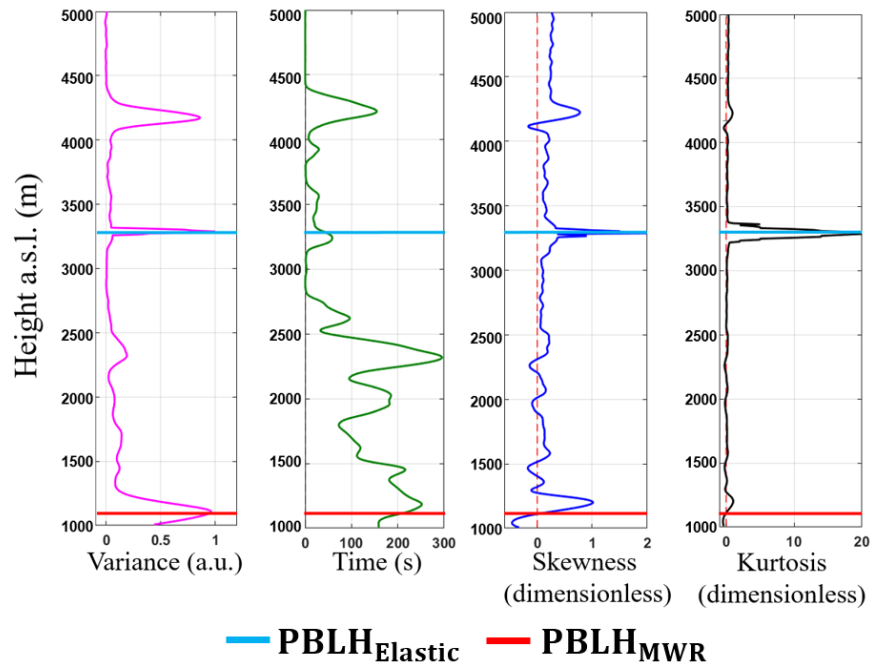


Figure 17 - Statistical moments obtained from 1064 nm wavelength data of elastic lidar(MULHACÉN) in Granada between 11-12 UTC on 08th July 2016. From left to right: variance [ $\sigma_{RCS'}^2$ ], integral time scale [ $\tau_{RCS'}$ ], skewness [ $S_{RCS'}$ ] and kurtosis [ $K_{RCS'}$ ].

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## Granada – 08 Jul – 12-13 UTC

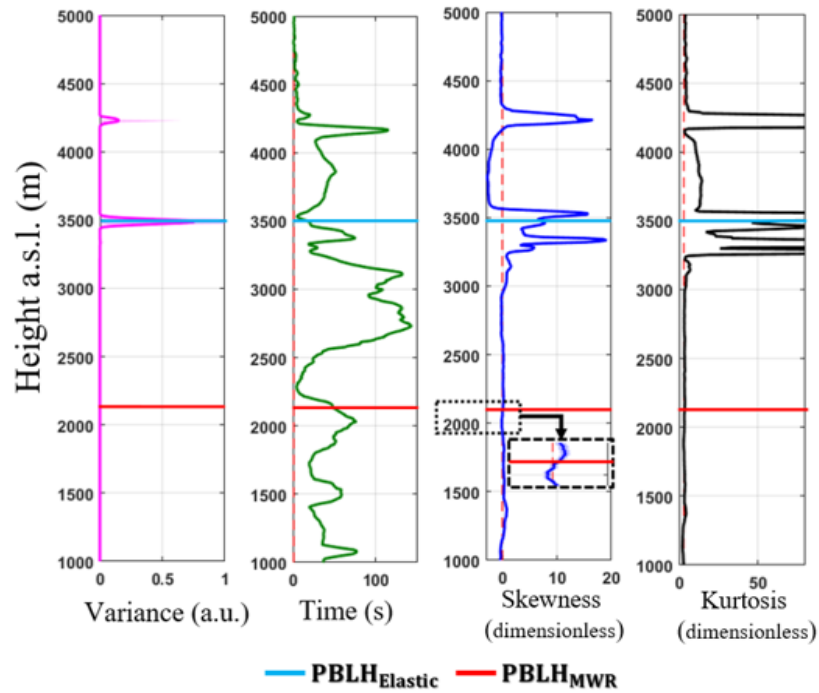


Figure 18 - Statistical moments obtained from 532 nm wavelength data of elastic lidar (MULHACÉN) in Granada between 12 -13 UTC on 08 July 2016. From left to right: variance [ $\sigma_{RCS'}^2$ ], integral time scale [ $\tau_{RCS'}$ ], skewness [ $S_{RCS'}$ ] and kurtosis [ $K_{RCS'}$ ].

## Granada – 08 Jul – 12-13 UTC

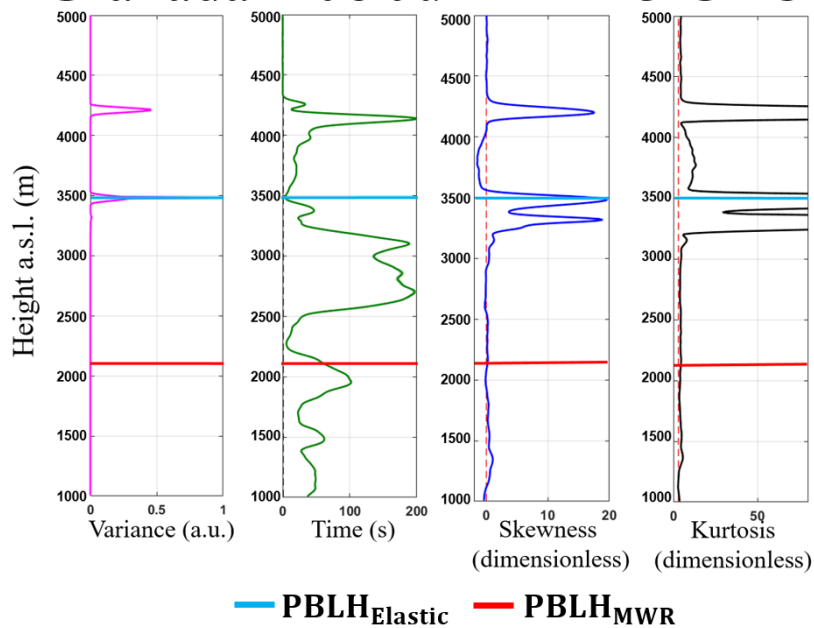


Figure 19 - Statistical moments obtained from 1064 nm wavelength data of elastic lidar (MULHACÉN) in Granada between 12 -13 UTC on 08 July 2016. From left to right: variance [ $\sigma_{RCS'}^2$ ], integral time scale [ $\tau_{RCS'}$ ], skewness [ $S_{RCS'}$ ] and kurtosis [ $K_{RCS'}$ ].