

# 1 Analyzing the turbulent Planetary Boundary Layer 2 behavior by the synergic use of remote sensing systems: 3 Doppler wind lidar, aerosol elastic lidar and microwave 4 radiometer

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## 15 Abstract

16 The Planetary Boundary Layer (*PBL*) is the lowermost region of troposphere and endowed with turbulent  
17 characteristics, which can have mechanical and/or thermodynamic origins. Such behavior gives to this layer  
18 great importance, mainly in studies about pollutant dispersion and weather forecasting. However, the  
19 instruments usually applied in studies about turbulence in the *PBL* have limitations in spatial resolution  
20 (anemometer towers) or temporal resolution (instrumentation onboard aircraft). In this study we propose  
21 the synergetic use of remote sensing systems (microwave radiometer [*MWR*], Doppler lidar [*DL*] and elastic  
22 lidar [*EL*]) to analyze the turbulent *PBL* behavior. Furthermore, we show how some meteorological  
23 variables such as air temperature, aerosol number density, vertical wind speed, relative humidity and net  
24 radiation might influence the turbulent *PBL* dynamic. The statistical moments of the high frequency  
25 distributions of the vertical wind velocity, derived from *DL* and of the backscattered coefficient derived  
26 from *EL*, are corrected by two methodologies, namely first lag and  $-2/3$  correction. The corrected profiles  
27 present small differences when compared against the uncorrected profiles, showing low influence of noise  
28 and the viability of the proposed methodology. Two case studies were analyzed in detail, one corresponding  
29 to a well-defined *PBL* and another one corresponding to a situation with presence of a Saharan dust lofted  
30 aerosol layer and clouds. In both cases the results provided by the different instruments are complementary,  
31 thus the synergistic use of the different systems allow us performing a detailed monitoring of the turbulent  
32 *PBL* behavior, as well as, a better understanding about how the analyzed variables can interfere in this  
33 process.

34 **Keywords:** Turbulence, Planetary Boundary Layer, Doppler lidar, elastic lidar, microwave radiometer,  
35 Earlinet.

## 36 1 Introduction

37 The Planetary Boundary Layer (*PBL*) is the atmospheric layer directly influenced by the Earth's surface  
38 that responds to its changes within time scales around an hour (Stull, 1988). Such layer is located at the

39 lowermost region of troposphere, and is mainly characterized by turbulent processes and a daily evolution  
40 cycle. In an ideal situation, instants after sunrise, ground surface temperature increase due to the positive  
41 net radiative flux ( $R_n$ ). This process intensifies the convection, thus, the ascending warm air masses heat  
42 the air masses situated in the upper regions of troposphere, originating the Convective Boundary Layer  
43 (*CBL*) or Mixing Layer (*ML*), which has this name due to a mixing process generated by this turbulent  
44 ascending air parcels. Some instants before sunset the gradual reduction of incoming solar irradiance at the  
45 Earth's surface causes the decrease of the positive  $R_n$  and its change in sign. In this situation, there is a  
46 reduction of the convective processes and a weakening of the turbulence. In this process the *CBL* leads to  
47 the development of two layers, namely a stably stratified boundary layer called Stable Boundary Layer  
48 (*SBL*) close to the surface, and the Residual Layer (*RL*) that contains features from the previous day's *ML*  
49 and is just above the *SBL*.

50 Knowledge of the turbulent processes in the *CBL* is important in diverse studies, mainly for atmospheric  
51 modeling and pollutant dispersion, since turbulent mixing can be considered as the primary process by  
52 which aerosol particles and other scalars are transported vertically in atmosphere. Because turbulent  
53 processes are treated as nondeterministic, they are characterized and described by their statistical properties  
54 (high order statistical moments). When applied to atmospheric studies such analysis provide information  
55 about the field of turbulent fluctuation, as well as, a description of the mixing process in the PBL (Pal et  
56 al., 2010).

57 Anemometer towers have been widely applied in studies about turbulence (e.g., Kaimal and Gaynor, 1983;  
58 van Ulden and Wieringa, 1996), however the limited vertical range of these equipment restrict the analysis  
59 to regions close to surface. Aircraft have also been used in atmospheric turbulence studies (e.g., Lenschow  
60 et al., 1980; Williams and Hacker, 1992; Lenschow et al., 1994; Albrecht et al., 1995; Stull et al., 1997;  
61 Andrews et al., 2004; Vogelmann et al., 2012), nevertheless their short time window limits the analysis. In  
62 this scenario, systems with high spatial and temporal resolution and enough range are necessary in order to  
63 provide more detailed results along the day throughout the whole thickness of the *PBL*.

64 In the last decades, lidar systems have been increasingly applied in this kind of study due to their large  
65 vertical range, high data acquisition rate and capability to detect several observed quantities such as vertical  
66 wind velocity [Doppler lidar] (e.g. Lenschow et al., 2000; Lothon et al., 2006; O'Connor et al., 2010), water  
67 vapor [Raman lidar and DIAL] (e.g. Wulfmeyer, 1999; Kiemle et al., 2007; Wulfmeyer et al., 2010; Turner  
68 et al., 2014; Muppa et al., 2015), temperature [rotational Raman lidar] (e.g. Behrendt et al., 2015) and  
69 aerosol [elastic lidar] (e.g. Pal et al., 2010; McNicholas et al., 2015). This allows the observation of a wide  
70 range of atmospheric processes. For example, Pal et al. (2010) demonstrated how the statistical analyses  
71 obtained from high-order moments of elastic lidar can provide information about aerosol plume dynamics  
72 in the PBL region. In addition, when different lidar systems operate synergistically, as for example in  
73 Engelmann et al. (2008), who combined elastic and Doppler lidar data, it is possible to identify very  
74 complex variables such as vertical particle flux. However, this subject requires more exploration, mainly  
75 the synergy among lidar and others remote sensing systems, like microwave radiometer. Thus, the  
76 combination of information obtained from these instruments can provide a more detailed understanding  
77 about the turbulent PBL behavior. Such approach is even more attractive when considering facilities of

78 networks, e. g. European Aerosol Research Lidar NETwork (EARLINET) (Pappalardo et al., 2014),  
79 Microwave Radiometer Network (MWRNET) (Rose et al., 2005; Caumont et al., 2016) and ACTRIS  
80 CLOUDNET (Illingworth et al., 2007).

81 Therefore, considering this scenario, in this study we use synergistically the data of three remote sensing  
82 systems (Elastic Lidar [EL], Doppler Lidar [DL] and Microwave Radiometer [MWR]) acquired during the  
83 SLOPE-I campaign, held at IISTA-CEAMA (Andalusian Institute for Earth System Research, Granada,  
84 Spain) from May to August 2016, in order to analyze the turbulent PBL behavior and to improve our  
85 comprehension about how each analyzed variable influence the PBL dynamics.

86 This paper is organized as follows. Description of the experimental site and the equipment setup are  
87 presented in Section 2. The methodologies applied are introduced in Section 3. Section 4 presents the results  
88 of the analyses using the different methodologies. Finally, conclusions are summarized in Section 5.

89

## 90 **2 Experimental site and instrumentation**

91 The SLOPE-I (Sierra nevada Lidar aerOsol Profiling Experiment) campaign was performed from May to  
92 September 2016 in South-Eastern Spain in the framework of the European Research Infrastructure for the  
93 observation of Aerosol, Clouds, and Trace gases (ACTRIS). The main objective of this campaign was to  
94 perform a closure study by comparing remote sensing system retrievals of atmospheric aerosol properties,  
95 using remote systems operating at the Andalusian Institute of Earth System Research (IISTA-CEAMA)  
96 and in-situ measurements operating at different altitudes in the Northern slope of Sierra Nevada, around 20  
97 km away from IISTA-CEAMA (Bedoya-Velásquez et al., 2018; Román et al., 2018). The IISTA-CEAMA  
98 station is part of EARLINET (Pappalardo et al, 2014) since 2005 and at present is an ACTRIS station  
99 (<http://actris2.nilu.no/>). The research facilities are located at Granada, a medium size city in Southeastern  
100 Spain (Granada, 37.16°N, 3.61°W, 680 m a.s.l.), surrounded by mountains and with Mediterranean-  
101 continental climate conditions that are responsible for cool winters and hot summers. Rain is scarce,  
102 especially from late spring to early autumn. Granada is affected by different kind of aerosol particles locally  
103 originated and medium-long range transported from Europe, Africa and North America (Lyamani et al.,  
104 2006; Guerrero-Rascado et al., 2008, 2009; Titos et al., 2012; Navas-Guzmán et al., 2013; Valenzuela et  
105 al., 2014, Ortiz-Amezcuca et al, 2014, 2017).

106 MULHACEN is a biaxial ground-based Raman lidar system operated at IISTA-CEAMA in the frame of  
107 EARLINET research network. This system operates with a pulsed Nd:YAG laser, frequency doubled and  
108 tripled by Potassium Dideuterium Phosphate crystals, emitting at wavelengths of 355, 532 and 1064 nm  
109 with output energies per pulse of 60, 65 and 110 mJ, respectively. MULHACEN operates with three elastic  
110 channels: 355, 532 (parallel and perpendicular polarization) and 1064 nm and three Raman-shifted  
111 channels: 387 (from N<sub>2</sub>), 408 (from H<sub>2</sub>O) and 607 nm (from N<sub>2</sub>). MULHACEN's overlap is complete at  
112 90% between 520 and 820 m a.g.l. for all the wavelengths, reaching full overlap around 1220 m a.g.l.  
113 (Navas-Guzmán et al ., 2011; Guerrero-Rascado et al. 2010). Calibration of the depolarization capabilities

114 is done following Bravo-Aranda et al. (2013). This system was operated with a temporal and spatial  
115 resolution of 2 s and 7.5 m, respectively. More details can be found at Guerrero-Rascado et al. (2008, 2009).

116 The Doppler lidar (Halo Photonics, model Stream Line XR) is also operated at IISTA-CEAMA. This  
117 system works in continuous and automatic mode from May 2016. It operates at 1.5  $\mu\text{m}$  with pulse energy  
118 and repetition rate of 100  $\mu\text{J}$  and 15 KHz, respectively. This system record the backscattered signal with  
119 300 gates, being the range gate length 30 m, with the first gate at 60 m. The telescope focus is set to  
120 approximately 800 m. For this work the data were collected in stare mode (laser beam is pointed at vertical  
121 with respect to the ground surface) with a time resolution of 2 s.

122 Furthermore, we operated the ground-based passive microwave radiometer (RPG-HATPRO G2,  
123 Radiometer Physics GmbH), which is member of the MWRnet [<http://cetemps.aquila.infn.it/mwrnet/>]. This  
124 system operates in automatic and continuous mode at IISTA-CEAMA since November 2011. The  
125 microwave radiometer (MWR) measures the sky brightness temperature with a radiometric resolution  
126 between 0.3 and 0.4 K root mean square error at 1 s integration time, using direct detection receivers within  
127 two bands: K-band (water vapor – frequencies: 22.24 GHz, 23.04 GHz, 23.84 GHz, 25.44 GHz, 26.24 GHz,  
128 27.84 GHz, 31.4 GHz) and V-band (oxygen – frequencies: 51.26 GHz, 52.28 GHz, 53.86 GHz, 54.94 GHz,  
129 56.66 GHz, 57.3 GHz, 58.0 GHz). From these bands is possible to obtain profiles of water vapor and  
130 temperature, respectively, by inversion algorithms described in Rose et al. (2005). The range resolution of  
131 these profiles vary between 10 and 200 m in the first 2 km and between 200 and 1000 m in the layer between  
132 2 and 10 km (Navas-Guzmán et al., 2014).

133 The meteorological sensor (HMP60, Vaisala) is used to register the air surface temperature and surface  
134 relative humidity, with a temporal resolution of 1 minute. Relative humidity is monitored with an accuracy  
135 of  $\pm 3\%$ , and air surface temperature is acquired with an accuracy and precision of  $0.6^\circ\text{C}$  and  $0.01^\circ\text{C}$ ,  
136 respectively.

137 A CM-11 pyranometer manufactured by Kipp & Zonen (Delft, The Netherlands) is also installed in the  
138 ground-based station. This equipment measures the shortwave (SW) solar global horizontal irradiance data  
139 (305–2800 nm). The CM-11 pyranometer complies with the specifications for the first-class WMO (World  
140 Meteorological Organization) classification of this instrument (resolution better than  $\pm 5\text{ Wm}^{-2}$ ), and the  
141 calibration factor stability has been periodically checked against a reference CM-11 pyranometer (Antón  
142 et. al, 2012).

### 143 **3 Methodology**

#### 144 **3.1 MWR data analysis**

145 The MWR data are analyzed combining two algorithms, Parcel Method [*PM*] (Holzworth, 1964) and  
146 Temperature Gradient Method [*TGM*] (Coen, 2014), in order to estimate the *PBL* Height ( $PBLH_{MWR}$ ) in  
147 convective and stable situations, respectively. The different situations are discriminated by comparing the  
148 surface potential temperature ( $\theta(z_0)$ ) with the corresponding vertical profile of  $\theta(z)$  up to 5 km. Those

149 cases where all the points in the vertical profile have values larger than  $\theta(z_0)$  are labeled as stable, and  
 150 *TGM* is applied. Otherwise the situation is labeled as unstable and the *PM* is applied. The vertical profile  
 151 of  $\theta(z)$  is obtained from the vertical profile of  $T(z)$  using the following equation (Stull, 2011):

$$152 \quad \theta(z) = T(z) + 0.0098 * z \quad (1)$$

153 where  $T(z)$  is the temperature profile provided by *MWR*,  $z$  is the height above the sea level, and 0.0098  
 154 K/m is the dry adiabatic temperature gradient. A meteorological station co-located with the *MWR* is used  
 155 to detect the surface temperature  $[T(z_0)]$ . In order to reduce the noise,  $\theta(z)$  profiles were averaged  
 156 providing a  $PBLH_{MWR}$  value at 30 minutes intervals. This methodology of *PBLH* detection was selected as  
 157 the reference due to the results obtained during a performed campaign of comparison between *MWR* and  
 158 radiosonde data, where twenty-three radiosondes were launched. High correlations were found between  
 159 *PBLH* retrievals provided by both instruments in stable and unstable cases. Further details are given by  
 160 Moreira et al. (2018a).

### 161 **3.2 Lidar turbulence analysis**

162 Both lidar systems, *DL* and *EL*, gathered data with a temporal resolution of 2 seconds. Then, the data are  
 163 averaged in 1-hour packages, from which the mean value is extracted  $[\bar{q}(z)]$ . Such mean value is subtracted  
 164 from each  $q(z, t)$  profile in order to estimate the vertical profile of the fluctuation for the measured variable  
 165  $[q'(z, t)]$  (i.e. vertical velocity for the *DL*):

$$166 \quad q'(z, t) = q(z, t) - \bar{q}(z) \quad (2)$$

167 Then, from  $q'(z, t)$  is possible to obtain the high-order moments (variance ( $\sigma^2$ ), skewness ( $S$ ) and kurtosis  
 168 ( $K$ )), as well as, the integral time scale ( $\tau$  - which is the time over which the turbulent process are highly  
 169 correlated to itself) as shown in Table 1. These variables can also be obtained from the following  
 170 autocovariance function,  $M_{ij}$ :

$$171 \quad M_{ij} = \int_0^{t_f} [q'(z, t)]^i [q'(z, t + t_f)]^j dt \quad (3)$$

172 where  $t_f$  is the final time,  $i$  and  $j$  indicate the order of autocovariance function.

173 However, it is necessary to considerer that the acquired real data contain instrumental noise,  $\varepsilon(z)$ .  
 174 Therefore, the equation 3 can be rewritten as:

$$175 \quad M_{ij} = \int_0^{\tau} [q(z, t) + \varepsilon(z, t)]^i [q(z, t + \tau) + \varepsilon(z, t + \tau)]^j dt \quad (4)$$

176 The autocovariance function of a time series with zero lag results in the sum of the variances of the  
 177 atmospheric variable and its  $\varepsilon(z)$ . Nevertheless, atmospheric fluctuations are correlated in time, but the  
 178  $\varepsilon(z)$  is random and uncorrelated with the atmospheric signal. Consequently, the noise is only associated  
 179 with lag 0 (Fig. 1). Based on this concept Lenschow et al. (2000) suggested to obtain the corrected

180 autocovariance function,  $M_{11}(\rightarrow 0)$ , from two methods, namely first lag correction or -2/3 law correction.  
 181 In the first method,  $M_{11}(\rightarrow 0)$  is obtained directly by the subtraction of lag 0,  $\Delta M_{11}(0)$ , from the  
 182 autocovariance function,  $M_{11}(0)$ . In the second method  $M_{11}(\rightarrow 0)$  is generated by the extrapolation of  
 183  $M_{11}(0)$  at firsts nonzero lags back to lag zero (-2/3 law correction). The extrapolation can be performed  
 184 using the inertial subrange hypothesis, which is described by the following equation (Monin and Yaglom,  
 185 1979):

$$186 \quad M_{11}(\rightarrow 0) = \overline{q'^2(z, t)} + Ct^{2/3} \quad (5)$$

187 where C represents a parameter of turbulent eddy dissipation rate. The high-order moments and  $\tau$   
 188 corrections and errors are shown in Table 1 (columns 2 and 3, respectively).

189 The same procedure of analysis is applied in studies with *DL* and *EL*, being the main difference the tracer  
 190 used by each system, which are the fluctuation of vertical wind speed ( $w'$ ) for *DL* and aerosol number  
 191 density ( $N'$ ) for *EL*. *DL* provides  $w(z, t)$  directly, and therefore the procedure described in Figure 2 can be  
 192 directly applied. Thus, the two corrections described above are applied separately and finally  $\tau$  and high-  
 193 order moments with and without corrections can be estimated.

194 On the other hand, the *EL* does not provide  $N(z, t)$  directly. Under some restrictions, it is possible to ignore  
 195 the particle hygroscopic growth and to assume that the vertical distribution of aerosol type does not changes  
 196 with time, and to adopt the following relation (Pal et al., 2010):

$$197 \quad \beta_{par}(z, t) \approx N(z, t)Y(z) \Rightarrow \beta'_{par}(z, t) = N'(z, t) \quad (6)$$

198 where  $\beta_{par}$  and  $\beta'_{par}$  represent the particle backscatter coefficient and its fluctuation, respectively, and  $Y$   
 199 does not depends on time.

200 Considering the lidar equation:

$$201 \quad P_{\lambda}(z) = P_0 \frac{ct_d}{2} AO(z) \frac{\beta_{\lambda}(z)}{z^2} e^{-2 \int_0^z \alpha_{\lambda}(z') dz'} \quad (7)$$

202 where  $P_{\lambda}(z)$  is the signal returned from distance  $z$  at time  $t$ ,  $z$  is the distance [m] from the lidar of the  
 203 volume investigated in the atmosphere,  $P_0$  is the power of the emitted laser pulse,  $c$  is the light speed [m/s],  
 204  $t_d$  is the duration of laser pulse [ns],  $A$  is the area [m<sup>2</sup>] of telescope cross section,  $O(z)$  is the overlap  
 205 function,  $\alpha_{\lambda}(z)$  is the total extinction coefficient (due to atmospheric particles and molecules) [(km)<sup>-1</sup>] at  
 206 distance  $z$ ,  $\beta_{\lambda}(z)$  is the total backscatter coefficient (due to atmospheric particles and molecules) [(km·sr)<sup>-1</sup>]  
 207 at distance  $z$  and the subscript  $\lambda$  represents the wavelength. The two path transmittance term related to  
 208  $\alpha(z)$  is considered as nearly negligible at 1064 nm (Pal et al., 2010). Thus, it is possible to affirm that:

$$209 \quad RCS_{1064}(z) = P(z)_{1064} \cdot z^2 \cong G \cdot \beta_{1064}(z) \quad (8)$$

210 and consequently:

$$211 \quad RCS'_{1064}(z, t) \cong \beta'_{1064}(z, t) = \beta'_{par}(z, t) = N'(z, t) \quad (9)$$

212 where  $RCS_{1064}$  and  $RCS'_{1064}$  are the range corrected signal and its fluctuation, respectively,  $G$  is a constant  
213 and the subscripts represent the wavelength.

214 In this way, Pal et al. (2010) have shown the feasibility of using  $EL$  operating at 1064 nm for describing  
215 the atmospheric turbulence. In a recent paper Moreira et al. (2018b), have shown that the use of the  $EL$  at  
216 532 nm, in spite of the larger attenuation expected at this wavelength due to both aerosol and molecules,  
217 provides a description of the turbulence equivalent to that provided by  $EL$  operating at 1064 nm. This result  
218 is interesting having in mind the more extended use of lidar systems based on laser emission at 532 nm in  
219 different coordinated networks. Thus, in EARLINET and LALINET (Latin American Lidar NETwork)  
220 around 76% and 45% of the systems include the wavelength of 1064 nm, while 95% of the EARLINET  
221 systems and 73% of the LALINET systems operate systems that include the wavelength 532 nm (Guerrero-  
222 Rascado et al., 2016). Furthermore, the performance of the lidar systems at 532 nm presents better signal  
223 to noise ratio than that encountered at 1064nm. Thus, in this study we use the  $RCS_{532}$  for analyzing  
224 turbulence using  $EL$ , following the procedure described in Figure 3, which is basically the same  
225 methodology described earlier for  $DL$ .

226 These three methodologies, together with data of net surface radiation (obtained from pyranometer data)  
227 and air temperature (provided by MWR), are used synergistically in order to complement one each other  
228 and consequently generate a detailed picture of how each variable influences the turbulent PBL behavior,  
229 as it will be demonstrated in subsection 4.2.

## 230 **4 Results**

### 231 **4.1 Error Analysis**

232 The influence of random error in noisy observations rapidly grows for higher-order moments (i.e., the  
233 influence of random noise is much larger for the fourth-order moment than for the third-order moment).  
234 Therefore, the first step, in order to ascertain the applied methodology and our data quality, we performed  
235 the error treatment of  $DL$  data as described in Figure 2.

236 Figure 4 illustrates the autocovariance function, generated from  $w'$ , at three different heights. As mentioned  
237 before, the lag 0 is contaminated by noise  $\varepsilon$ , and thus the impact of the noise  $\varepsilon$  increases together with  
238 height, mainly above  $PBLH_{MWR}$  (1100 m a.g.l. in our example).

239 Figure 5-A illustrates the comparison between integral time scale ( $\tau_{w'}$ ) without correction and the two  
240 corrections cited in section 3.2. Except for the first height, under the  $PBLH_{MWR}$  the profiles practically do  
241 not have significant difference, as well as small errors bars. Above  $PBLH_{MWR}$  the first lag correction  
242 presents some differences in relation the other profiles at around 1350 m.

243 Figures 5-B and 5-C show the comparison of variance ( $\sigma_{w'}^2$ ) and skewness ( $\mathcal{S}_{w'}$ ), respectively, with and  
244 without corrections. The profiles corrected by -2/3 law do not present significant differences in comparison  
245 to uncorrected profiles. On the other hand, the profiles corrected by the first lag correction have significant

246 differences under the  $PBLH_{MWR}$ , mainly the  $\sigma_w^2$  ( $\mathbf{S}_w$ , only in the first 50 m), and some slight differences  
247 are evident above  $PBLH_{MWR}$ .

248 For  $EL$  we use the same procedure for the correction and error analysis that we apply to the  $DL$  data. Figure  
249 6 shows the autocovariance function, obtained from  $RCS'$ , at three distinct heights. As expected, the  
250 increase of height produces the increase of  $\varepsilon$ , principally above the  $PBLH_{MWR}$ .

251 Figures 7-A, 7-B, 7-C and 7-D show the vertical profiles of  $\tau_{RCS'}$ ,  $\sigma_{RCS'}^2$ ,  $\mathbf{S}_{RCS'}$  and kurtosis ( $\mathbf{K}_{RCS'}$ ),  
252 respectively, with and without the corrections described in section 3.2. In general, the corrections do not  
253 affect the profiles in a significant way, especially in the region below the  $PBLH_{MWR}$ . Above the  $PBLH_{MWR}$   
254 some small differences are noticed, mainly in the first lag correction. The error bars associated to each  
255 profile also have low values in all cases. When comparing corrected and uncorrected profiles, the largest  
256 differences are observed for the profiles at higher order moments, because of error propagation.  $\mathbf{K}_{RCS'}$   
257 profile is the more affected by corrections, so the kurtosis profile after the first lag correction shows the  
258 largest difference with uncorrected profile.

259 Since the first lag and 2/3 corrections do not have a significant impact within the  $PBL$  region, we adopted  
260 the first lag correction in order to be more careful during the comparison.

## 261 **4.2 Case studies**

262 In this section we present two study cases, in order to show how the synergy of methodologies described  
263 in section 3 can provide a detailed description about the turbulent  $PBL$  behavior. The first case represents  
264 a typical day with a clear sky situation. The second case corresponds to a more complex situation, where  
265 there is presence of clouds and Saharan mineral dust layers.

### 266 **4.2.1 Case study I: clear sky situation**

267 In this case study we use measurements gathered with  $DL$ ,  $MWR$  and pyranometer during 24 hours. The  
268  $EL$  was operated under operator-supervised mode between 08:20 to 18:00 UTC.

269 Figure 8 (A) shows the integral time scale obtained from  $DL$  data ( $\tau_w$ ). The gray areas represents the region  
270 where  $\tau_w$  is lower than the acquisition time of  $DL$  and, therefore, for this region it is not possible to analyze  
271 turbulent processes. However, the gray area is located almost entirely above the  $PBLH_{MWR}$  (white stars).  
272 Thus, the  $DL$  acquisition time allows us to observe the turbulence throughout the whole  $PBL$ . The gray  
273 areas, as well as, the black lines (air temperature), have the same meaning in Figures 8-B and 8-C.

274  $\sigma_w^2$  has low values during the entire period of  $SBL$  (Figure 8-B). Nevertheless, as air temperature begins to  
275 increase (around 07:00 UTC),  $\sigma_w^2$  increases together, as well as,  $PBLH_{MWR}$ .  $\sigma_w^2$  reaches its maximum  
276 values in the middle of the day, when we also observe the maximum values of air temperature and  
277  $PBLH_{MWR}$ . This process is in agreement with the behavior of skewness of  $w'$  ( $\mathbf{S}_w$ ) shown in Figure 8-C.  
278  $\mathbf{S}_w$  is directly associated with the direction of turbulent movements. Thus, positive values correspond with  
279 a surface-heating-driven boundary layer, while negative ones are associated to cloud-top long-wave

280 radiative cooling. If  $S_{w'}$  is positive, both  $\sigma_{w'}^2$  and  $TKE$  (Turbulent Kinetic Energy) are being transported  
281 upwards and consequently, the red regions in Figure 13-C represent positive values of  $S_{w'}$  and the blue  
282 regions refer to negative ones. During the stable period, there is predominance of low values of  $S_{w'}$ .  
283 Nevertheless, as air temperature increases (transition from stable to unstable period),  $S_{w'}$  values begin to  
284 become positive and increase with the ascent of the  $PBLH_{MWR}$  ( $CBL$ ). Air temperature begins to decrease  
285 around 18:00 UTC, causing the reduction of  $S_{w'}$ . In this moment the transition from unstable to stable  
286 period occurs and, therefore, the reduction in  $PBLH_{MWR}$  is due to the  $SBLH$  detection.

287 Figure 8-D shows the values of net surface radiation ( $R_n$ ) that are estimated from solar global irradiance  
288 values using the seasonal model described in Alados et al. (2003). The negative values of  $R_n$  are  
289 concentrated in the stable region.  $R_n$  begins to increase around 06:00 UTC and reaches its maximum in the  
290 middle of the day. Comparing figures 8-C and 8-D, we can observe similarity among the behavior of  $S_{w'}$ ,  
291  $R_n$  and surface air temperature, because these variables increase and decrease together, as expected.

292 The increase of  $R_n$  causes the rise of surface air temperature, which contributes to the positive latent heat  
293 flux from the surface ( $S_{w'}$ ) and, consequently, the growth of the  $PBLH_{MWR}$  ( $CBL$ ).  $R_n$  begins to decrease  
294 certain time before the other variables, but the intense reduction of air temperature and decrease of  $S_{w'}$  and  
295  $SBLH$  detection occurs when  $R_n$  becomes negative again, although there can still be a positive sensible heat  
296 flux, what is characteristic of early evening in urban regions due to the release of the ground heat flux at  
297 that time.

298 Figure 8-E presents the values of surface air temperature and surface relative humidity ( $RH$ ). Air surface  
299 temperature is directly related with  $R_n$  and  $S_{w'}$  values, as aforementioned and expected. On the other hand,  
300  $RH$  is inversely correlated with temperature and, thus, with the rest of variables, due to the relative  
301 constancy of the water vapor mixing ratio characteristic of our site during the study

302 Figure 9 shows the  $RCS_{532}$  profile obtained from 08:00 to 18:00 UTC and the well-defined  $PBLH_{MWR}$   
303 (pink stars). At the beginning of the measurement period (08:20 to 10:00 UTC) it is possible to observe the  
304 presence of a thin residual layer (around 2000 m a.s.l.), and later from 13:00 to 18:00 UTC it is evident a  
305 lofted aerosol layer. The period between 13:00 and 14:00 UTC has been selected to be analyzed. Figure 10-  
306 A presents the profiles of molecular ( $\beta_{Molecular}$ ) and aerosol ( $\beta_{Aerosol}$ ) backscatter coefficients at 532 nm.  
307 Although  $\beta_{532}$  is composed by  $\beta_{Molecular}$  and  $\beta_{Aerosol}$ , it is possible to observe the predominance of  
308  $\beta_{Aerosol}$  in the region below of the  $PBLH_{MWR}$ , as demonstrated in figure 10-B by the  $\beta_{Ratio}$  profile. Similar  
309 results were demonstrated by Moreira et al. (2018b), therefore reinforcing the viability of the use of this  
310 wavelength in studies about turbulence. Figure 11 presents the statistical moments generated from  $RCS'$ ,  
311 which were obtained from 13:00 and 14:00 UTC. The maximum for the variance of RCS can be used as  
312 indicator of  $PBLH$  ( $PBLH_{Elastic}$ ) (Moreira et al., 2015). Thus, the red line in all graphics represent the  
313  $PBLH_{Elastic}$  (2200 m a.s.l.) and the blue one the average value of  $PBLH_{MWR}$  (2250 m a.s.l.), both obtained  
314 between 13 and 14 UTC.

315 Due to well-defined  $PBL$ ,  $PBLH_{Elastic}$  and  $PBLH_{MWR}$  do not have significant differences (50 m).  $\sigma_{RCS'}^2$  has  
316 small values below the  $PBLH$ . Above  $PBLH_{Elastic}$  the values of  $\sigma_{RCS'}^2$  decrease slowly due to location of

317 the lofted aerosol around 2500 m. However, above this aerosol layer the value of  $\sigma_{RCS'}^2$ , is reduced to zero,  
 318 indicating the extreme decreasing in aerosol concentration in the free troposphere. The integral time scale  
 319 obtained from RCS' ( $\tau_{RCS'}$ ) has values higher than *EL* time acquisition throughout the CBL, evidencing  
 320 the feasibility for studying turbulence using this elastic lidar configuration. The skewness values obtained  
 321 from RCS' ( $S_{RCS'}$ ) give us information about aerosol motion. The positive values of  $S_{RCS'}$ , observed in the  
 322 lowest part of profile and above the  $PBLH_{Elastic}$  represents the updrafts aerosol layers. The negative values  
 323 of  $S_{RCS'}$ , indicates the region with low aerosol concentration due to clean air coming from free troposphere  
 324 (*FT*). This movement of ascension of aerosol layers and descent of clean air with zero value of  $S_{RCS'}$ , is  
 325 characteristic of growing PBL and was also detected by Pal et al. (2010) and McNicholas et al. (2014). The  
 326 kurtosis of RCS' ( $K_{RCS'}$ ) determines the level of mixing at different heights. There are values of  $K_{RCS'}$ ,  
 327 larger than 3 in the lowest part of profile and around 2500 m, showing a peaked distribution in this region.  
 328 On other hand, values of  $K_{RCS'}$ , lower than 3 are observed close to the  $PBLH_{Elastic}$ , therefore this region  
 329 has a well-mixed *CBL* regime. Pal et al. (2010) and McNicholas et al. (2014) also detected this feature in  
 330 the region nearby the *PBLH*.

331 The results provided by *DL*, pyranometer and *MWR* data agree with the results observed in Figure 10. In  
 332 the same way, the analysis of high order moments of *RCS'* fully agree with the information in Figure 8.  
 333 Thus, the large values of  $S_{RCS'}$ , and  $K_{RCS'}$ , detected around 2500 m a.s.l, where we can see a lofted aerosol  
 334 layer, suggest the ascent of an aerosol layer and presence of a peaked distribution, respectively.

#### 335 4.2.2 Case study: dusty and cloudy scenario

336 In this case study measurements with *DL*, *MWR* and pyranometer expand during 24 hours, while *EL* data  
 337 are collected from 09:00 to 16:00 UTC.

338 Figure 12-A shows  $\tau_{w'}$ , where the black lines and gray area has the same meaning mentioned earlier.  
 339 Outside the period 13:00 to 17:00 UTC, the grey area is situated completely above the  $PBLH_{MWR}$  (white  
 340 stars), thus *DL* time acquisition is enough to perform studies about turbulence in this case.

341  $\sigma_{w'}^2$ , has values close to zero during all the stable period (Figure 12-B). However, when air temperature and  
 342  $PBLH_{MWR}$  begins to increase (around 06:00 UTC),  $\sigma_{w'}^2$ , also increases and reaches its maximum in the  
 343 middle of the day. In the late afternoon, as air temperature and  $PBLH_{MWR}$  decrease, the values of  $\sigma_{w'}^2$ ,  
 344 decrease gradually, until reach the minimum value associated to the SBL. Figure 12-C shows the profiles  
 345 of  $S_{w'}$ . In the same way of the previous case study, the behavior of  $S_{w'}$  is directly related to the air  
 346 temperature pattern (increasing and decreasing together) and causing the growth and reduction of  
 347  $PBLH_{MWR}$ . The main features of this case are: the low values of  $S_{w'}$ , the slow increase and ascension of  
 348 positive  $S_{w'}$  values and the predominance of negative  $S_{w'}$  values from 12:00 to 13:00 UTC. The first two  
 349 features are likely due to the presence of the intense Saharan dust layer (Figure 13), which reduces the  
 350 transmission of solar irradiance, and consequently the absorption of solar irradiance at the surface,  
 351 generating weak convective process. From Figure 13 we can observe the presence of clouds from 12:00 to  
 352 14:00 UTC. This justifies the intense negative values of  $S_{w'}$ , observed in this period, because, as mentioned

353 before,  $S_w$ , is directly associated with direction of turbulent movements that during this period is associated  
354 to cloud-top long-wave radiative cooling, due to the presence of clouds (Ansmann et al., 2010).

355 The influence of Saharan dust layer can also be evidenced on the  $R_n$  pattern (Figure 12-D), which maintains  
356 negative values until 12:00 UTC and reaches a low maximum value (around 200 W/m<sup>2</sup>). Air surface  
357 temperature and  $RH$  (Figure 12-E) present the same correlation and anti-correlation (respectively) observed  
358 in the earlier case study, where the maximum of air surface temperature and the minimum of  $RH$  are  
359 detected in coincidence with the maximum daily value of  $PBLH_{MWR}$ .

360 As mentioned before, Figure 13 shows the  $RCS$  profile obtained from 09:00 to 16:00 UTC in a complex  
361 situation, with presence of decoupled dust layer (around 3800 m a.s.l.) from 09:00 and 12:00 and clouds  
362 (around 3500 m a.s.l.) from 11:00 to 16:00 UTC. The pink stars represent  $PBLH_{MWR}$ . Figure 14-A presents  
363 the  $\beta_{Molecular}$  and  $\beta_{Aerosol}$  profiles, similarly to Figure 10-A. It is evident the predominance of  $\beta_{Aerosol}$  in  
364 the region below  $PBLH_{MWR}$ , as demonstrated by  $\beta_{Ratio}$  profile in figure 14-B. However due to presence of  
365 dust layer this dominance of  $\beta_{Aerosol}$  is extended to approximately 4500 m a.s.l. Therefore the methodology  
366 proposed by Moreira et al. (2018b), based on considerations of Pal et al. (2010), can be applied.

367 Figure 15 illustrates the statistical moments of  $RCS'$  obtained from 11:00 to 12:00 UTC. The  $\sigma_{RCS'}^2$  profile  
368 presents several peaks due to the presence of distinct aerosol sublayers. The first peak is coincident with  
369 the value of  $PBLH_{MWR}$ . The value of  $PBLH_{elastic}$ , is coincident with the base of the dust layer. This  
370 difficulty to detect the  $PBLH$  in presence of several aerosol layers is inherent to the variance method  
371 (Kovalev and Eichinger, 2004). The values of  $\tau_{RCS'}$  are higher than  $EL$  acquisition time all along the  $PBL$ ,  
372 evidencing the feasibility of  $EL$  time acquisition for studying the turbulence of  $PBL$  in this case. The  $S_{RCS'}$   
373 profile has several positive values, due to the large number of aerosol sublayers that are present. The  
374 characteristic inflection point of  $S_{RCS'}$ , is observed in coincidence with the  $PBLH_{MWR}$ , that confirming the  
375 agreement between this point and the  $PBLH$ .  $K_{RCS'}$  has predominantly values lower than 3 below 2500 m,  
376 thus shown how this region is well mixed as can see in Figure 13. Values of  $K_{RCS'}$ , larger than 3 are observed  
377 in the highest part of profile, where the dust layer is located.

378 Figure 16 shows the  $RCS'$  high-order moments obtained from 12:00 and 13:00 in presence of cloud cover.  
379 The method based on maximum of  $\sigma_{RCS'}^2$ , locates the  $PBLH_{Elastic}$  at the cloud base, due to the high variance  
380 of  $RCS'$  generated by the clouds.  $\tau_{RCS'}$  presents values larger than  $EL$  time acquisition, therefore this  
381 configuration enable us to study turbulence by  $EL$  analyses.  $S_{RCS'}$  has few peaks, due to the mixing between  
382  $CBL$  and dust layer, generating a more homogenous layer. The highest values of  $S_{RCS'}$  are observed in  
383 regions where there are clouds, and the negative ones (between 3500 and 4000 m) occur due to presence of  
384 air from  $FT$  between the two aerosol layers (Figure 13). The inflection point of  $S_{RCS'}$  profile is observed in  
385  $PBLH_{MWR}$  region.  $K_{RCS'}$  profile has low values in most of the  $PBL$ , demonstrating the high level of mixing  
386 during this period, where dust layer and  $PBL$  are combined. The higher values of  $K_{RCS'}$  are observed in the  
387 region of clouds.

## 388 5 Conclusions

389 In this paper we analyze the turbulent PBL behavior and how each detected variable can influence it. Such  
390 observations were made from the synergy of three different types of remote sensing systems (DL, EL and  
391 MWR) and surface sensors during SLOPE-I campaign. We applied two kind of corrections to the lidar data:  
392 first lag and  $-2/3$  corrections. The corrected DL statistical moments showed little variation with respect to  
393 the uncorrected profiles, denoting a rather low influence of the noise. The statistical moments obtained  
394 from EL also showed a small variation after correction when compared with the uncorrected profiles, except  
395 for  $K_{RCS}$ , that is more affected by noise. The small changes in the profiles after the corrections, specially  
396 inside the PBL, evidence the feasibility of the applied methodology for monitoring the turbulence in the  
397 PBL. Nevertheless, all profiles are corrected by first lag correction, which is more restrictive during the  
398 comparison, in order to be cautious.

399 The case studies present two kind of situations: well-defined PBL and a more complex situation with the  
400 presence of Saharan dust layer and some clouds.  $\sigma_{w'}^2$  and  $S_{w'}$  showed a good agreement with the behavior  
401 of the air temperature,  $R_n$  and  $PBLH_{MWR}$  in both situations, highlighting the feasibility in different  
402 atmospheric conditions.

403 The synergic use of remote sensing systems shows how the results provided by the different instruments  
404 can complement one each other. Thus, it is possible to observe the direct relationship among PBL growth,  
405  $S_{w'}$ ,  $\sigma_{w'}^2$ ,  $\sigma_{RCS}^2$ , and  $R_n$  values. In addition,  $S_{RCS}$  and  $K_{RCS}$  provide a good description about aerosol dynamic.  
406 The combination of these results gives us a detailed description about PBL dynamic and its structure.

407 Therefore, this study shows the feasibility of the described methodology based on remote sensing systems  
408 for studying the turbulence. The feasibility of using the analyses of high order moments of the RCS  
409 collected at 532nm at a temporal resolution of 2 s for the characterization of the atmospheric turbulence in  
410 the PBL offers the possibility for using this procedure in networks such as EARLINET or LALINET with  
411 a reasonable additional effort.

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419 2013".

420 **References**

- 421 Alados, I., Foyo-Moreno, I., Olmo, F. J., Alados-Arboledas, L. Relationship between net radiation and solar  
422 radiation for semi-arid shrub-land. *Agr. Forest Meteorol.*, 116, 221-227, 2003.
- 423 Albrecht, B. A., Bretherton, C. S., Johnson, D., Scubert, W. H., and Frisch, A. S.: The Atlantic  
424 stratocumulus transition experiment—ASTEX, *Bull. Am. Meteorol. Soc.*, 76, 889–904, 1995.
- 425 Andrews, E., Sheridan, P. J. , Ogren, J. A., and Ferrare, R.: In situ aerosol profiles over the Southern Great  
426 Plains cloud and radiation test bed site: 1. Aerosol optical properties, *J. Geophys. Res.*, 109, D06208,  
427 doi:10.1029/2003JD004025, 2004.
- 428 Ansmann, A., Fruntke, J., Engelmann, R. Updraft and downdraft characterization with Doppler lidar: cloud-  
429 free versus cumuli-topped mixed layer. *Atmos. Chem. Phys.*, 10, 7845-7858, 2010.
- 430 Antón, M., Valenzuela, A., Cazorla, A., Gil, J. E., Gálvez-Fernández, J., Lyamani, H., Foyo-Moreno, I.,  
431 Olmo, F. J., Alados-Arboledas, L. Global and diffuse shortwave irradiance during a strong desert dust  
432 episode at Granada (Spain). *Atmos. Res.*, 118, 232 – 239, 2012.
- 433 Bedoya-Velásquez, A. E., Navas-Guzmán, F., Granados-Muñoz, M. J., Titos, G., Román, R., Casquero-  
434 Vera, J. A., Ortiz-Amezcuca, P., Benavent-Oltra, J. A., Moreira, G. de A., Montilla-Rosero, E., Ortiz, C. D.  
435 H., Artiñano, Coz, E., Alados-Arboledas, L., Guerrero-Rascado, J. L. Hygroscopic growth study in the  
436 framework of EARLINET during the SLOPE I campaign: synergy of remote sensing and in-situ  
437 instrumentation. *Atmos. Chem. Phys.*, 18, 7001-7017, 2017.
- 438 Behrendt, A., Wulfmeyer, V., Hammann, E., Muppa, S. K., Pal, S.: Profiles of second- to  
439 fourth-order moments of turbulent temperature fluctuations in the convective boundary layer : first  
440 measurements with rotational Raman lidar. *Atmos. Chem. Phys.*, 15, 5485-5500.  
441 <https://doi.org/10.5194/acp-15-5485-2015>, 2015.
- 442 Bravo-Aranda, J. A., Navas-Guzmán, F., Guerrero-Rascado, J. L., Pérez-Ramírez, D., Granados-Muñoz,  
443 M. J., Alados-Arboledas, L. Analysis of lidar depolarization calibration procedure and application to the  
444 atmospheric aerosol characterization. *Int. J. Remote Sens.*, 34 (9-10), pp. 3543-3560, 2013.
- 445 Caumont, O., Cimini, D., Löhnert, U., Alados-Arboledas, L., Bleisch, R., Buffa, F., Ferrario, M.E., Haeefe,  
446 A., Huet, T., Madonna, F., Pace, G. Assimilation of humidity and temperature observations retrieved from  
447 ground-based microwave radiometers into a convective-scale NWP model. *Q. J. Roy. Meteor. Soc.*, 142  
448 (700), pp. 2692-2704, 2016.
- 449 Engelmann, R.; Wandinger, U.; Ansmann, A.; Müller, D.; Žeromskis, E.;  
450 Althausen, D.; Wehner, B. Lidar Observations of the Vertical Aerosol Flux in the Planetary  
451 Boundary Layer. *J. Atmos. Ocean. Tech.*, v. 25, n. 8, p. 1296-1306, 2008.
- 452 Guerrero-Rascado, J.L., Ruiz, B., Alados-Arboledas, L. Multi-spectral lidar characterization of the vertical  
453 structure of Saharan dust aerosol over Southern Spain. *Atmos. Environ.*, 42, 2668-2681, 2008.

454 Guerrero-Rascado, J.L., Olmo, F.J., Avilés-Rodríguez, I., Navas-Guzmán, F., Pérez-Ramírez, D., Lyamani,  
455 H., Alados-Arboledas, L.: Extreme Saharan dust event over the southern Iberian Peninsula in September  
456 2007: Active and passive remote sensing from surface and satellite. *Atmos. Chem. Phys.*, 9, 21, 8453-8469,  
457 2009.

458 Guerrero-Rascado, J. L., Costa, M. J., Bortoli, D., Silva, A. M., Lyamani, H., Alados-Arboledas, L. Infrared  
459 lidar overlap function: an experimental determination, *Opt. Express*, 18, 20350-20359, 2010

460 Guerrero-Rascado, J. L., Landulfo, E., Antuña, J. C., Barbosa, H. M. J., Barja, B., Bastidas, A. E., Bedoya,  
461 A. E., da Costa, R. F., Estevan, R., Forno, R. N., Gouveia, D. A., Jimenez, C., Larroza, E. G., Lopes, F. J.  
462 S., Montilla-Rosero, E., Moreira, G. A., Nakaema, W. M., Nisperuza, D., Alegria, D., Múnera, M., Otero,  
463 L., Papandrea, S., Pawelko, E., Quel, E. J., Ristori, P., Rodrigues, P. F., Salvador, J., Sánchez, M. F., and  
464 Silva, A.: Latin American Lidar Network (LALINET) for aerosol research: diagnosis on network  
465 instrumentation, *J. Atmos. Sol.-Terr. Phy.*, 138–139, 112–120, 2016.

466 Illingworth, A. J., Hogan, R. J. O' Connor, E. J. Bouniol, D. Brooks, M. E. Delanoe, J. Donovan, D. P.  
467 Eastment, J. D. Gaussiat, N. Goddard, J. W. F. Haeffelin, M. Klein Baltink, H. Krasnov, O. A. Pelon, J.  
468 Piriou, J.-M. Protat, A. Russchenberg, H. W. J. Seifert, A. Tompkins, A. M. Van Zadelhoff, G.-J. Vinit, F.  
469 Willen, U. Wilson, D. R. and Wrench, C. L.: CLOUDNET: Continuous Evaluation of Cloud Profiles in  
470 Seven Operational Models using Ground-Based Observations. *Bull. Am. Meteorol. Soc.*, 88, 883-898,  
471 doi:10.1175/BAMS-88-6-883, 2007.

472 Kaimal, J. C., and Gaynor, J. E.: The Boulder Atmospheric Observatory, *J. Clim. Appl. Meteorol.*, 22, 863–  
473 880, 1983.

474 Kiemle, C., Brewer, W. A., Ehret, G., Hardesty, R. M., Fix, A., Senff, C., Wirth, M., Poberaj, G., and  
475 LeMone, M. A.: Latent heat flux profiles from collocated airborne water vapor and wind lidars during IHOP  
476 2002, *J. Atmos. Ocean. Tech.*, 24, 627–639, 2007.

477 Kovalev, V.A., Eichinger, W.E., *Elastic Lidar*, Wiley 2004.

478 Lenschow, D. H., Wyngaard, J. C., and Pennell, W. T.: Mean-field and second-moment budgets in a  
479 baroclinic convective boundary layer, *J. Atmos. Sci.*, 37, 1313–1326, 1980.

480 Lenschow, D. H., Mann, J., and Kristensen, L.: How long is long enough when measuring fluxes and other  
481 turbulence statistics?, *J. Atmos. Oceanic Technol.*, 11, 661–673, 1994.

482 Lenschow, D. H., Wulfmeyer, V. and Senff, C.: Measuring second- through fourth-order moments in noisy  
483 data, *J. Atmos. Oceanic Technol.*, 17, 1330–1347, 2000.

484 Lenschow, D. H., Lothon, M., Mayor, S. D., Sullivan, P. P., and Canut, G.: A comparison of higher-order  
485 vertical velocity moments in the convective boundary layer from lidar with in situ measurements and large-  
486 eddy simulation, *Bound-Lay. Meteorol.*, 143, 107–123, doi:10.1007/s10546-011-9615-3., 2012.

487 Lothon, M., Lenschow, D. H., and Mayor, S. D.: Coherence and scale of vertical velocity in the convective  
488 boundary layer from a Doppler lidar, *Bound.-Lay. Meteorol.*, 121, 521–536, 2006.

489 Lyamani, H., Olmo, F. J., Alcántara, A., and Alados-Arboledas, L.: Atmospheric aerosols during the 2003  
490 heat wave in southeastern Spain I: Spectral optical depth, *Atmos. Environ.*, 40, 6453–6464, 2006

491 McNicholas, C., Turner, D. D. Characterizing the convective boundary layer turbulence with a High  
492 Spectral Resolution Lidar. *J. Geophys Res-Atmos.*, v. 119, p. 910-927, 2014.

493 Muppa, K.S., Behrendt, A., Späth, F., Wulfmeyer, V., Metzendorf, S., Riede, A.: Turbulent humidity  
494 fluctuations in the convective boundary layer: Cases studies using water vapour differential absorption lidar  
495 measurements. *Bound-Lay. Meteorol.*, 158, 43-66, DOI 10.1007/s10546-015-0078-9, 2014.

496 Monin, A. S., Yaglom, A. M. *Statistical Fluid Mechanics*, Vol. 2. MIT Press, 874 pp, 1979.

497 Moreira, G. de A., Marques, M. T. A., Nakaema, W., Moreira, A. C. de C. A., Landulfo, E. Planetary  
498 boundary height estimations from Doppler wind lidar measurements, radiosonde and hysplit model  
499 comparisom. *Óptica Pura y Aplicada*, 48, 179-183, 2015

500 Moreira, G. de A., Guerrero-Rascado, J. L., Bravo-Aranda, J. A., Benavent-Oltra, Ortiz-Amezcuca, P.,  
501 Róman, R., Bedoya-Velásquez, A., Landulfo, E., Alados-Arboledas, L. Study of the planetary boundary  
502 layer by microwave radiometer, elastic lidar and Doppler lidar estimations in Southern Iberian Peninsula.  
503 *Atmos. Res.*, 213, 185-195, 2018a.

504 Moreira, G. de A., Lopes, F. J. S., Guerrero-Rascado, J. L., Landulfo, E., Alados-Arboledas, L. Analyzing  
505 turbulence in Planetary Boundary Layer from multiwavelength lidar system: impact of wavelength choice.  
506 *Opt. Express*. Under review, 2018b.

507 Navas Guzmán, F., Guerrero Rascado, J. L., and Alados Arboledas, L.: Retrieval of the lidar overlap  
508 function using Raman signals, *Óptica Pura y Aplicada*, 44, 71–75, 2011.

509 Navas-Guzmán, F., Bravo-Aranda, J.A., Guerrero-Rascado, J.L, Granados-Muñoz, M.J, and Alados-  
510 Arboledas, L.: Statistical analysis of aerosol optical properties retrieved by Raman lidar over Southeastern  
511 Spain. *Tellus B*, 65, 21234, 2013.

512 Navas-Guzmán, F., Fernández-Gálvez, J., Granados-Muñoz, M.J, Guerrero-Rascado, J.L., Bravo-Aranda,  
513 J.A., and Alados-Arboledas, L.: Tropospheric water vapor and relative humidity profiles from lidar and  
514 microwave radiometry. *Atmos. Meas. Tech.*, 7, 1201-1211, 2104.

515 O'Connor, E. J., Illingworth, A. J., Brooks, I. M., Westbrook, C. D., Hogan, R. J., Davies, F., Brooks, B.  
516 J.: A method for estimating the turbulent kinetic energy dissipation rate from a vertically-pointing Doppler  
517 lidar, and independent evaluation from balloon-borne in-situ measurements. *J. Atmos. Ocean. Tech.*, v. 27,  
518 n. 10, 1652-1664, 2010.

519 Ortiz-Amezcuca, P., Guerrero-Rascado, J.L., Granados-Muñoz, M.J., Bravo-Aranda, J. A., Alados-  
520 Arboledas, L. Characterization of atmospheric aerosols for a long range transport of biomass burning  
521 particles from canadian forest fires over the southern iberian peninsula in july 2013. *Optica Pura y Aplicada*,  
522 47 (1), pp. 43-49, 2014.

523 Ortiz-Amezcuca, P., Luis Guerrero-Rascado, J., Granados-Munõz, M.J., Benavent-Oltra, J.A., Böckmann,  
524 C., Samaras, S., Stachlewska, I.S., Janicka, L., Baars, H., Bohlmann, S., Alados-Arboledas, L.  
525 Microphysical characterization of long-range transported biomass burning particles from North America at  
526 three EARLINET stations. *Atmos. Chem. Phys.*, 17 (9), pp. 5931-5946, 2017.

527 Pal, S.; Behrendt, a.; Wulfmeyer, V. Elastic-backscatter-lidar-based characterization of the  
528 convective boundary layer and investigation of related statistics. *Ann. Geophys.*, v. 28, n. 3, p. 825-  
529 847, 2010.

530 Pappalardo, G., Amodeo, A., Apituley, A., Comeron, A., Freudenthaler, V., Linné, H., Ansmann, A.,  
531 Bösenberg, J., D'Amico, G., Mattis, I., Mona, L., Wandinger, U., Amiridis, V., Alados-Arboledas, L.,  
532 Nicolae, D., and Wiegner, M.: EARLINET: towards an advanced sustainable European aerosol lidar  
533 network. *Atmos. Meas. Tech.*, 7, 2389-2409, doi:10.5194/amt-7-2389-2014, 2014.

534 Román, R., Benavent-Oltra, J. A., Casquero-Vera, J. A., Lopatin, A., Cazorla, A., Lyamani, H., Denjean,  
535 C., Fuertes, D., Pérez-Ramirez, D., Torres, B., Toledano, C., Dubovik, O., Cachorro, V. E., Frutos, A. M.,  
536 Olmo, F. J., Alados-Arboledas, L. Retrieval of aerosol profiles combining sunphotometer and ceilometer  
537 measurements in GRASP code. *Atmos. Res.*, 204, 161, 177, 2018.

538 Rose, T., Creewll, S., Löhnert, U., Simmer, C.: A network suitable microwave radiometer for operational  
539 monitoring of cloudy atmosphere. *Atmos. Res.*, 75, 3, 183 – 200, 2005.

540 Stull, R. B., Santoso, E., Berg, L., Hacker, J.: Boundary layer experiment 1996 (BLX96), *Bull. Am.*  
541 *Meteorol. Soc.*, 78, 1149–1158, 1997.

542 Stull, R. B.: An Introduction to Boundary Layer Meteorology, vol. 13, *Kluwer Academic Publishers, the*  
543 *Netherlands*, Dordrecht/Boston/London, 1988.

544 Stull, R. B.: *Meteorology for Scientists and Engineers*, 3<sup>rd</sup> Edition, *Uni. Of British Columbia*, 2011.

545 Titos, G., Foyo-Moreno, I., Lyamani, H., Querol, X., Alastuey, A., and Alados-Arboledas, L.: Optical  
546 properties and chemical composition of aerosol particles at an urban location: An estimation of the aerosol  
547 mass scattering and absorption efficiencies, *J. Geophys.Res.-Atmos.*, 117, D04206,  
548 doi:10.1029/2011JD016671, 2012.

549 Turner, D. D., Ferrare, R. A., Wulfmeyer, V., and Scarino, A. J.: Aircraft evaluation of ground-based  
550 Raman lidar water vapor turbulence profiles in convective mixed layers, *J. Atmos. Oceanic Technol.*, 31,  
551 1078–1088, doi:10.1175/JTECH-D-13-00075-1, 2014

552 van Ulden, A. P., and Wieringa, J. :Atmospheric boundary layer research at Cabauw, Bound-Lay.  
553 *Meteorol.*, 78,39–69, 1996.

554 Valenzuela, A., Olmo, F.J.ab, Lyamani, H.ab, Granados-Muñoz, M.J.ab, Antón, M.c, Guerrero-Rascado,  
555 J.L.ab, Quirantes, A.a, Toledano, C.d, Perez-Ramírez, D.ef, Alados-Arboledas, L.: Aerosol transport over  
556 the western mediterranean basin: Evidence of the contribution of fine particles to desert dust plumes over  
557 alborán island. *J. Geophys Res.*, 119, 24, 14,028-14,044, 2014.

558 Vogelmann, A. M., McFarquhar, G. M., Ogren, J. A., Turner, D. D., Comstock, J. M., Feingold, G., Long,  
559 C. N., Jonsson, H. H., Bucholtz, A., Collins, D. R., Diskin, G. S., Gerber, H., Lawson, R. P., Woods, R. K.,  
560 Andrews, E., Yang, H., Chiu, J. C., Hartsock, D., Hubbe, J. M., Lo, C., Marshak, A., Monroe, J. W.,  
561 Mcfarlane, S. A., Jason, M., and Toto, T.: RACORO extended-term aircraft observations of boundary layer  
562 clouds, *Bull. Am. Meteorol. Soc.*, 93, 861–878, doi:10.1175/BAMS-D-11-00189.1, 2012.

563 Williams, A. G., and Hacker, J. M.: The composite shape and structure of coherent eddies in the convective  
564 boundary layer, *Bound-Lay. Meteorol.*, 61, 213–245, 1992.

565 Wulfmeyer, V.: Investigation of turbulent processes in the lower troposphere with water vapor DIAL and  
566 radar-RASS, *J. Appl. Sci.*, 56, 1055–1076, 1999.

567 Wulfmeyer, V., Pal, S., Turner, D. D., and Wagner, E.: Can water vapour Raman lidar resolve profiles of  
568 turbulent variables in the convective boundary layer?, *Bound-Lay. Meteorol.*, 136, 253–284,  
569 doi:10.1007/s10546-010-9494-z, 2010.

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Table 1 – Variables applied to statistical analysis (Lenschow et al., 2000)

	Without Correction	Correction	Error
<b>Integral Time Scale (<math>\tau</math>)</b>	$\int_0^{\infty} q'(t)dt$	$\frac{1}{q'^2} \int_{t \rightarrow 0}^{\infty} M_{11}(t)dt$	$\tau \cdot \sqrt{\frac{4\Delta M_{11}}{M_{11}(\rightarrow 0)}}$
<b>Variance (<math>\sigma_q^2</math>)</b>	$\frac{1}{T} \sum_{t=1}^T (q(t) - \bar{q})^2$	$M_{11}(\rightarrow 0)$	$q^2 \cdot \sqrt{\frac{4\Delta M_{11}}{M_{11}(\rightarrow 0)}}$
<b>Skewness (<math>S</math>)</b>	$\frac{\bar{q}^3}{\sigma_q^3}$	$\frac{M_{21}(\rightarrow 0)}{M_{11}^{3/2}(\rightarrow 0)}$	$\frac{\Delta M_{21}}{\Delta M_{11}^{3/2}}$
<b>Kurtosis (<math>K</math>)</b>	$\frac{\bar{q}^4}{\sigma_q^4}$	$\frac{3M_{22}(\rightarrow 0) - 2M_{31}(\rightarrow 0) - 3\Delta M_{11}^2}{M_{11}^2(\rightarrow 0)}$	$\frac{4\Delta M_{31} - 3\Delta M_{22} - \Delta M_{11}^2}{\Delta M_{11}^2}$

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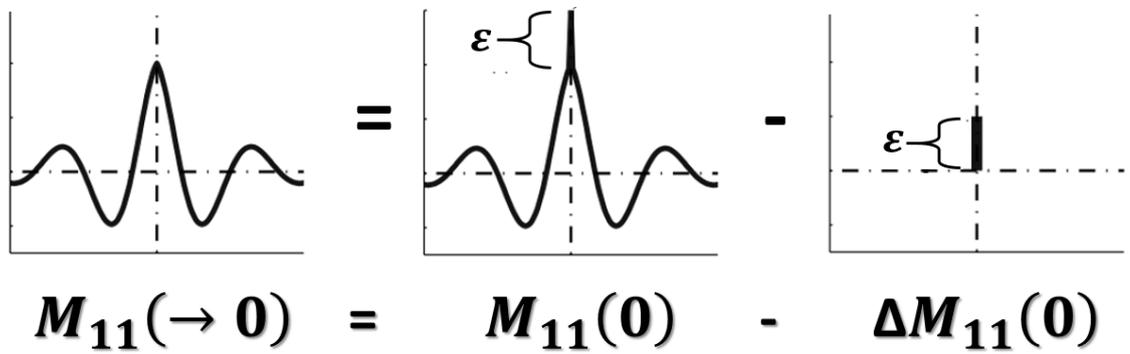


Figure 1 – Procedure to remove the errors of autocovariance functions.  $M_{11}(\rightarrow 0)$  – corrected autocovariance function errors;  $M_{11}(0)$  - autocovariance function without correction;  $\Delta M_{11}(0)$  - error of autocovariance function

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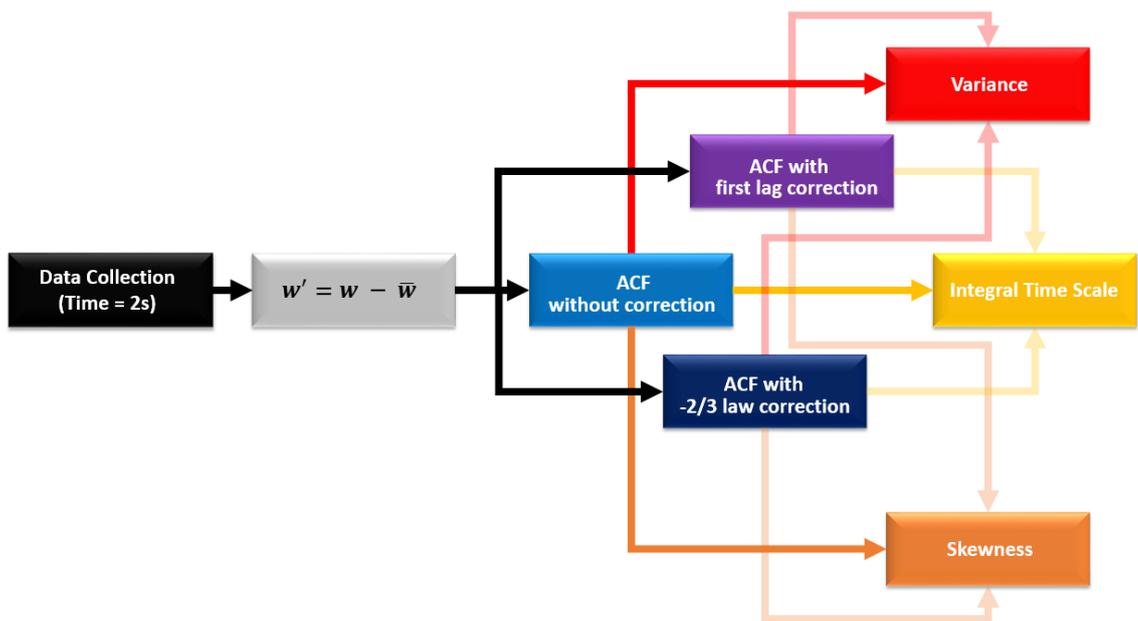


Figure 2 – Flowchart of data analysis methodology applied to the study of turbulence with Doppler lidar

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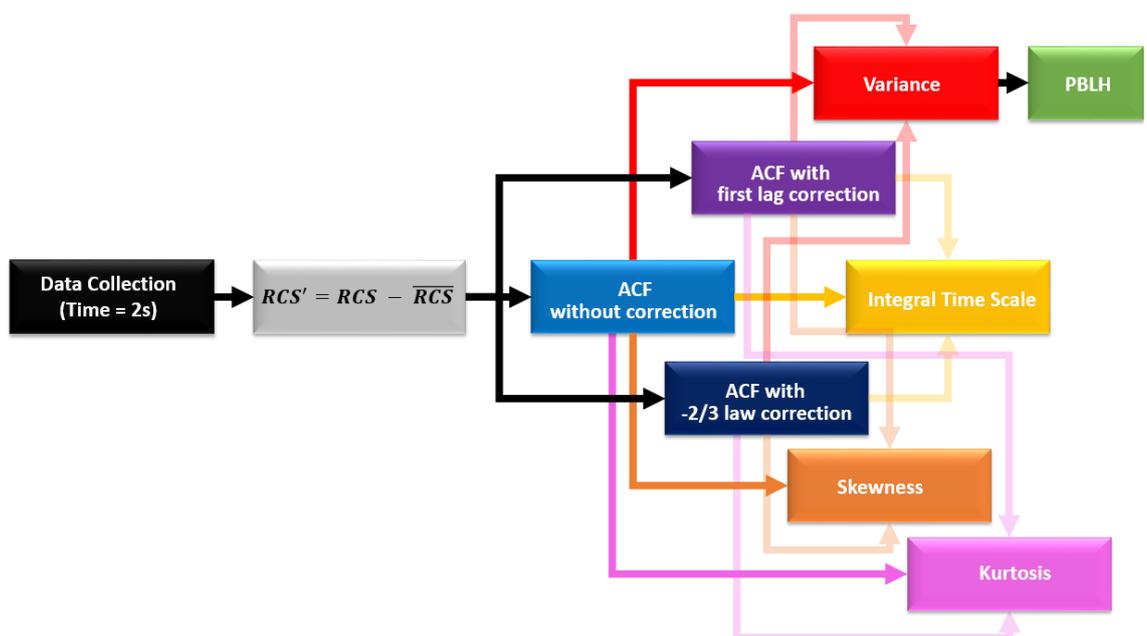


Figure 3 – Flowchart of data analysis methodology applied to the study of turbulence with elastic lidar

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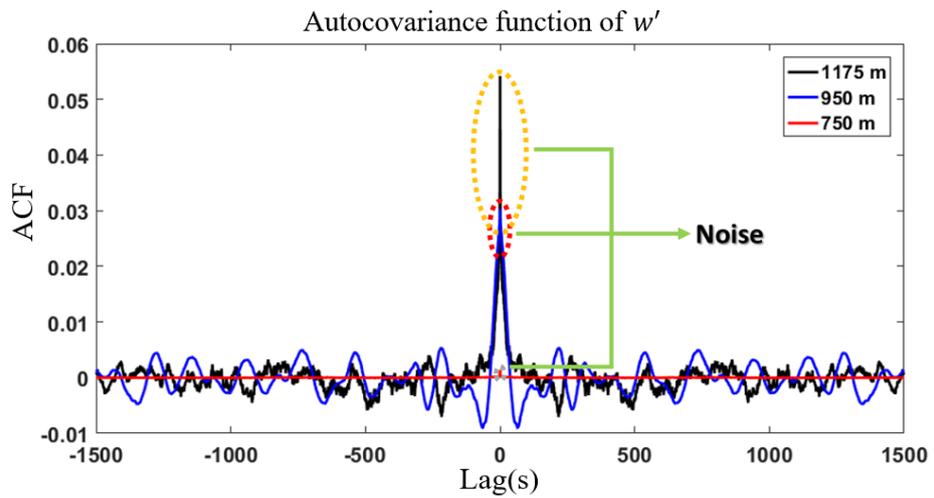


Figure 4 – Autocovariance function (ACF) of  $w'$  at three different heights

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### Profiles obtained from $w'$ - Granada – 19 May 2016 – 08-09 UTC

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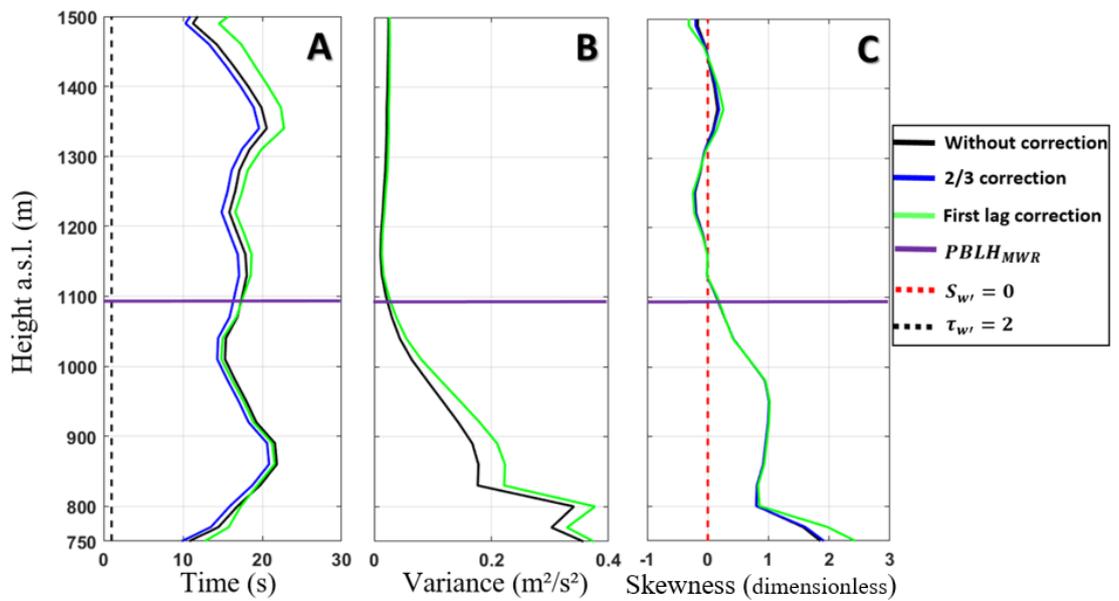


Figure 5 – A - Vertical profile of Integral time scale ( $\tau_{w'}$ ). B - Vertical profile of variance ( $\sigma_{w'}^2$ ). C - Vertical profile of Skewness. ( $S_{w'}$ )

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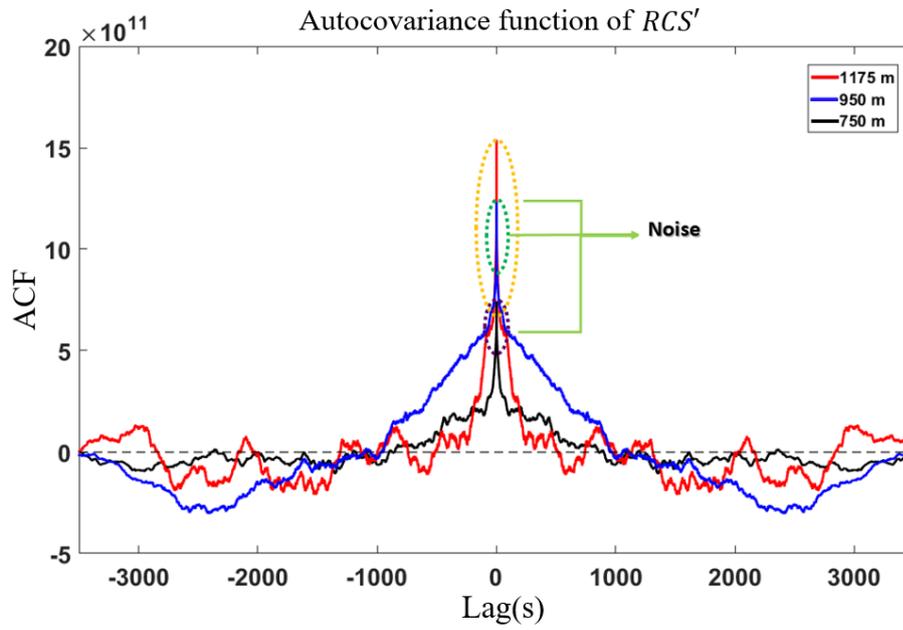


Figure 6 – Autocovariance of  $RCS'$  to three different heights

### Profiles obtained from $RCS'$ - Granada – 19 May 2016 – 09-10 UTC

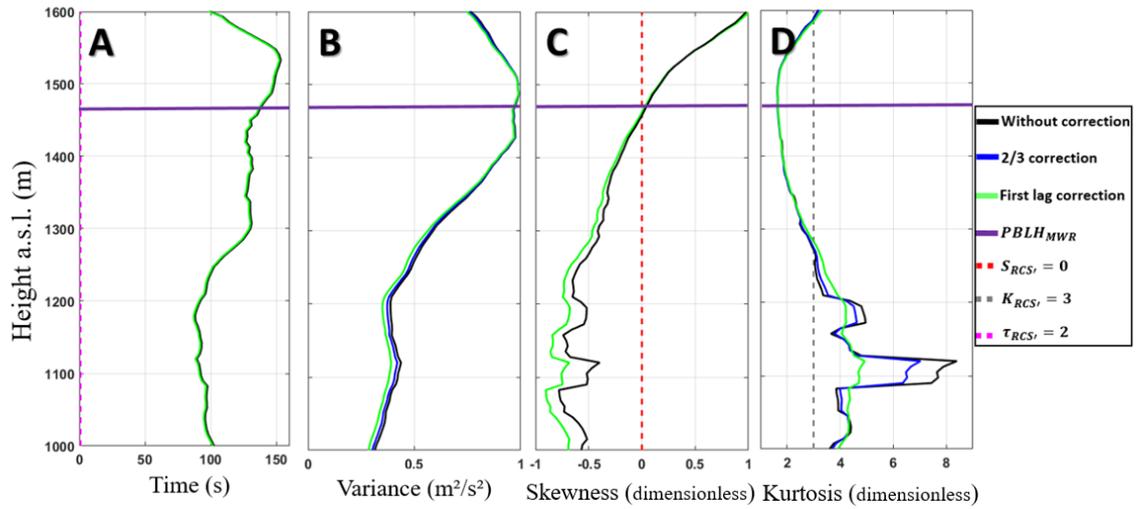


Figure 7 – A- Vertical profile of Integral time scale ( $\tau_{RCS'}$ ). B - Vertical profile of variance ( $\sigma_{RCS'}^2$ ). C - Vertical profile of Skewness ( $S_{RCS'}$ ). D - Vertical profile of Kurtosis ( $K_{RCS'}$ ).

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# Granada - 19 May 2016

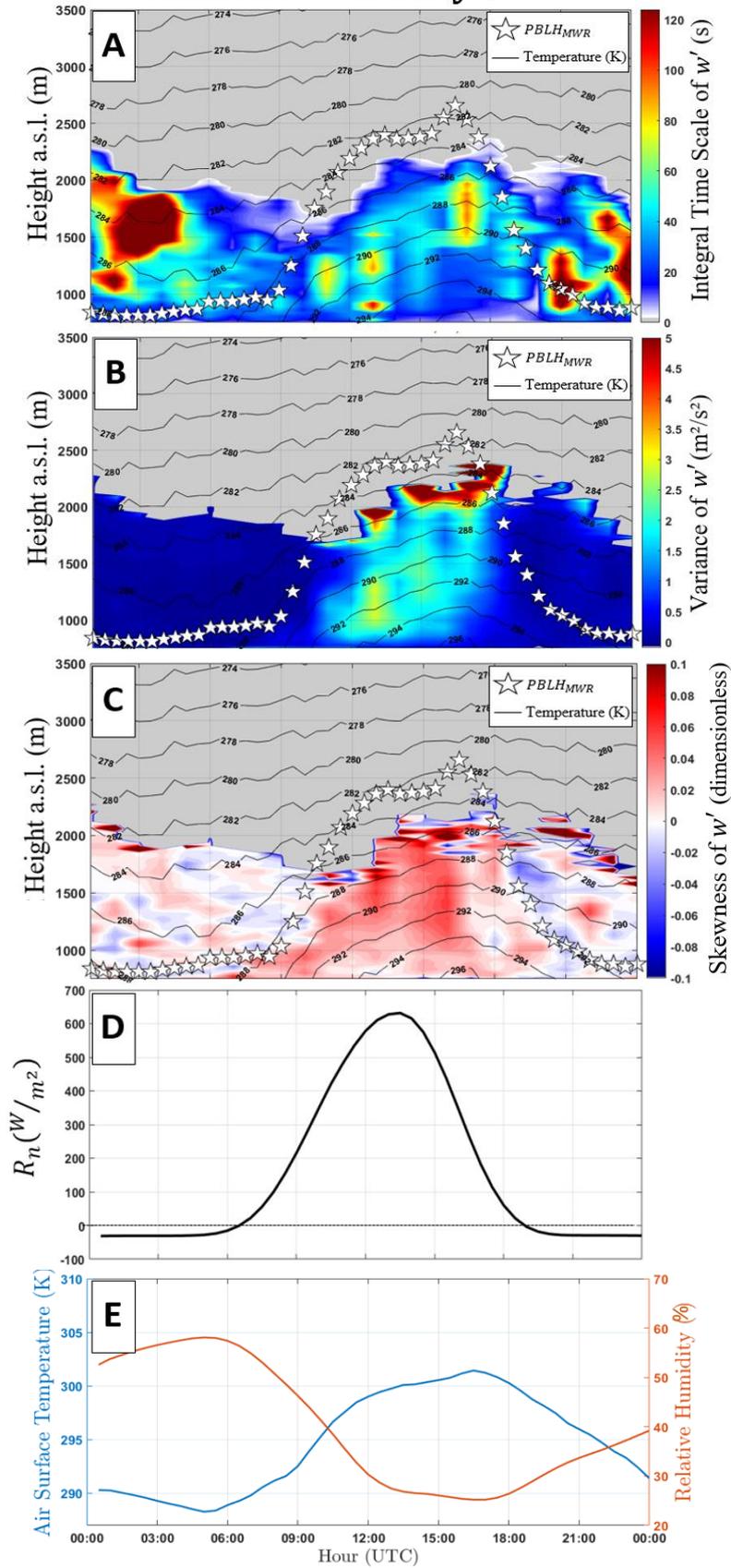


Figure 8 – A – integral time scale [ $\tau_{w'}$ ], B – variance [ $\sigma_{w'}^2$ ], C – skewness [ $S_{w'}$ ], D – net radiation [ $R_n$ ], E – Air surface temperature [blue line] and surface relative humidity [RH – orange line]. In A, B and C black lines and white stars represent air temperature and  $PBLH_{MWR}$ , respectively.

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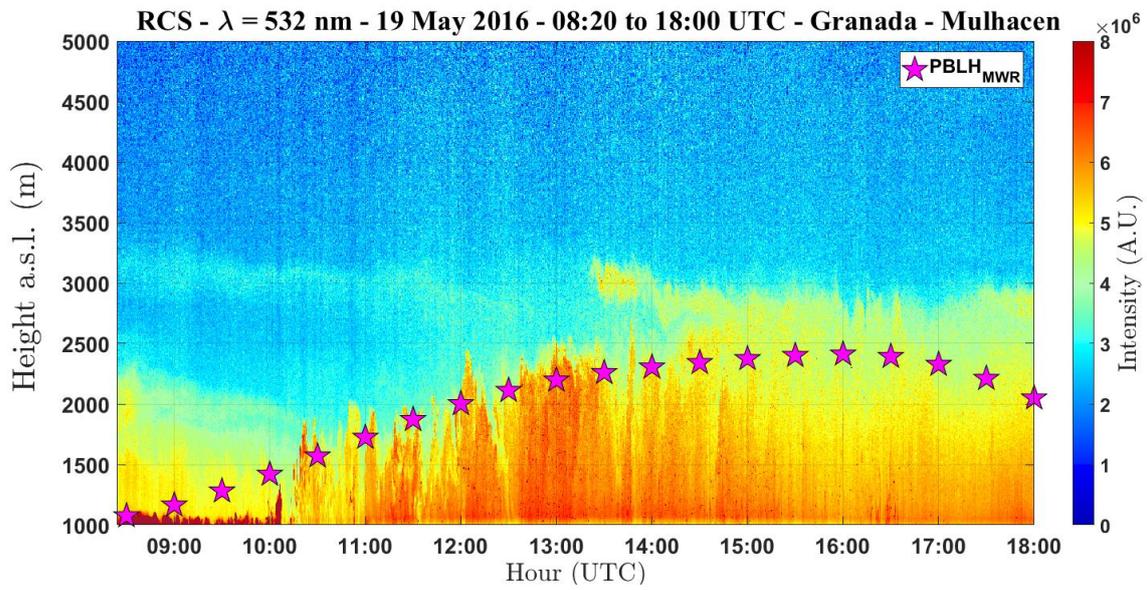


Figure 9 – Time-Height plot of RCS - 19 May 2016. Pink stars represent  $PBLH_{MWR}$

### Granada – 19 May – 13-14 UTC

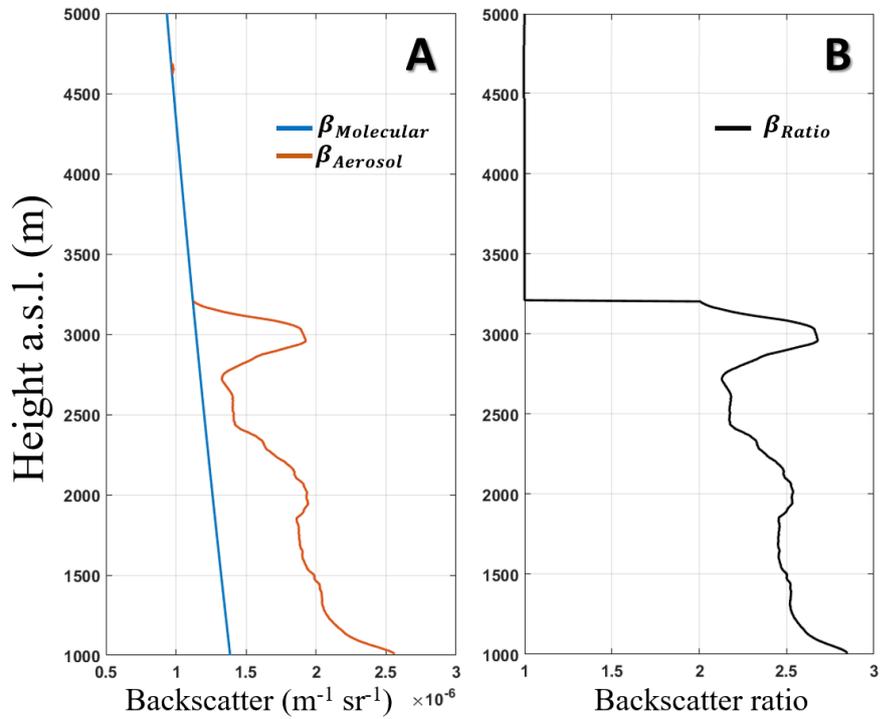


Figure 10 – (A)  $\beta_{Molecular}$  (blue line) and  $\beta_{Aerosol}$  (orange line). (B)  $\beta_{Ratio}$  (black line). All profiles were obtained from the 532 nm lidar signal

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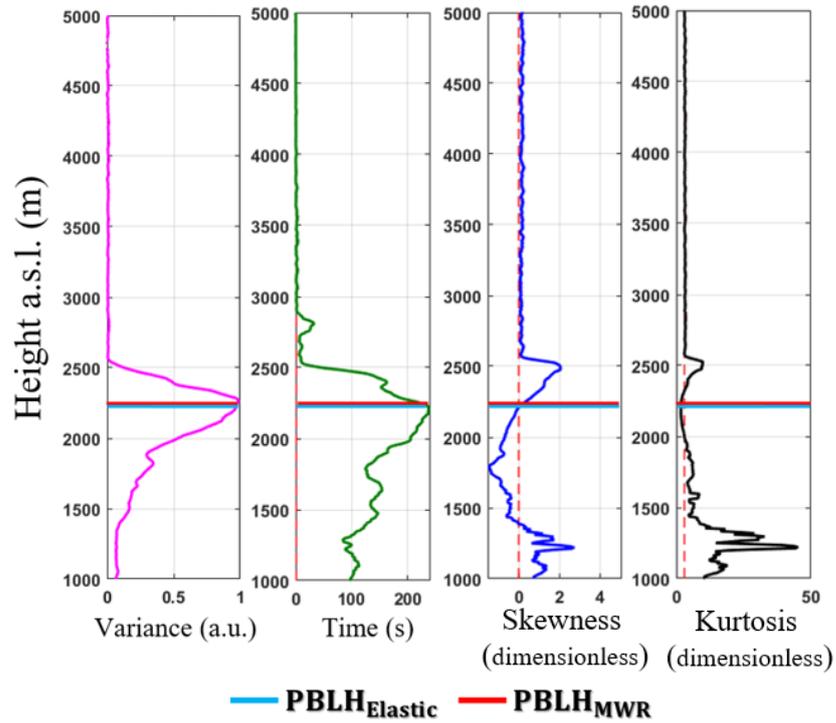


Figure 11 – Statistical moments obtained from elastic lidar data at 13 to 14 UTC - 19 May 2016. From left to right: variance [ $\sigma_{RCS'}^2$ ], integral time scale [ $\tau_{RCS'}$ ], skewness [ $S_{RCS'}$ ] and kurtosis [ $K_{RCS'}$ ].

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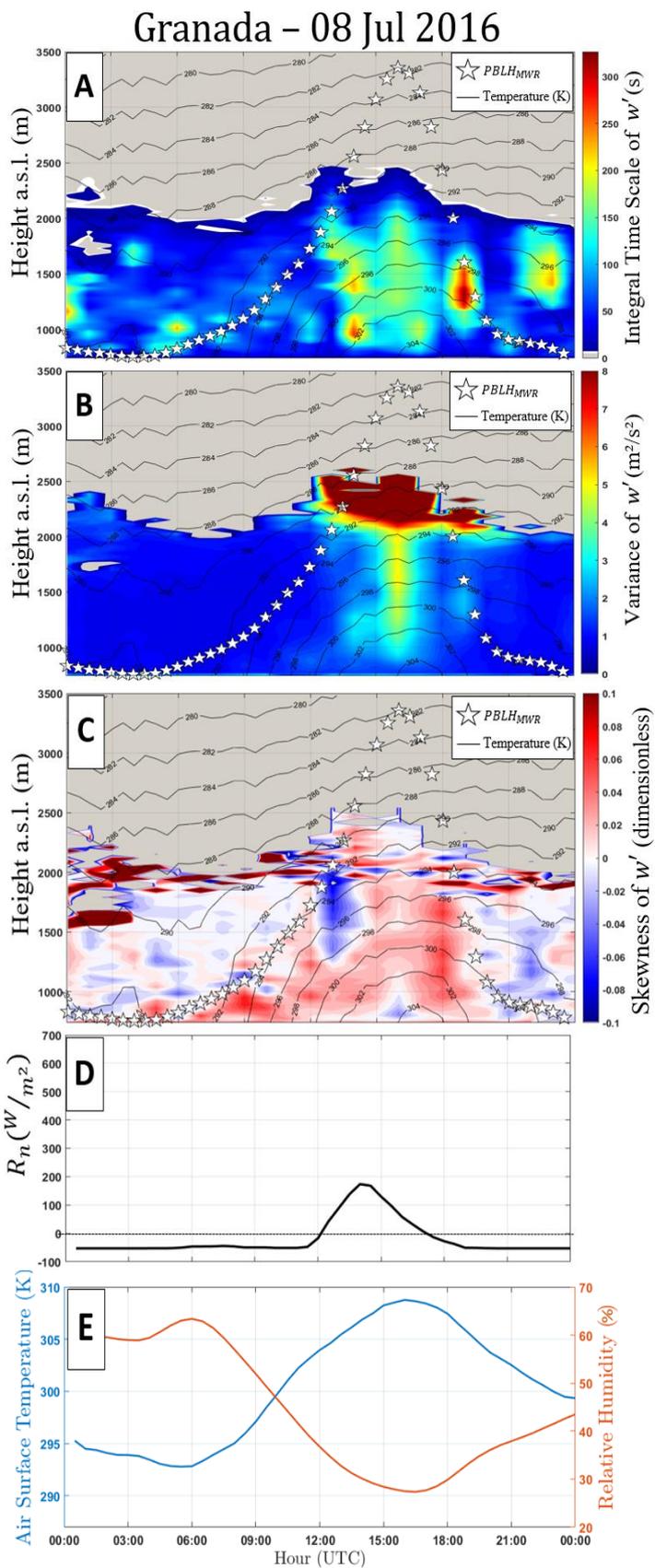
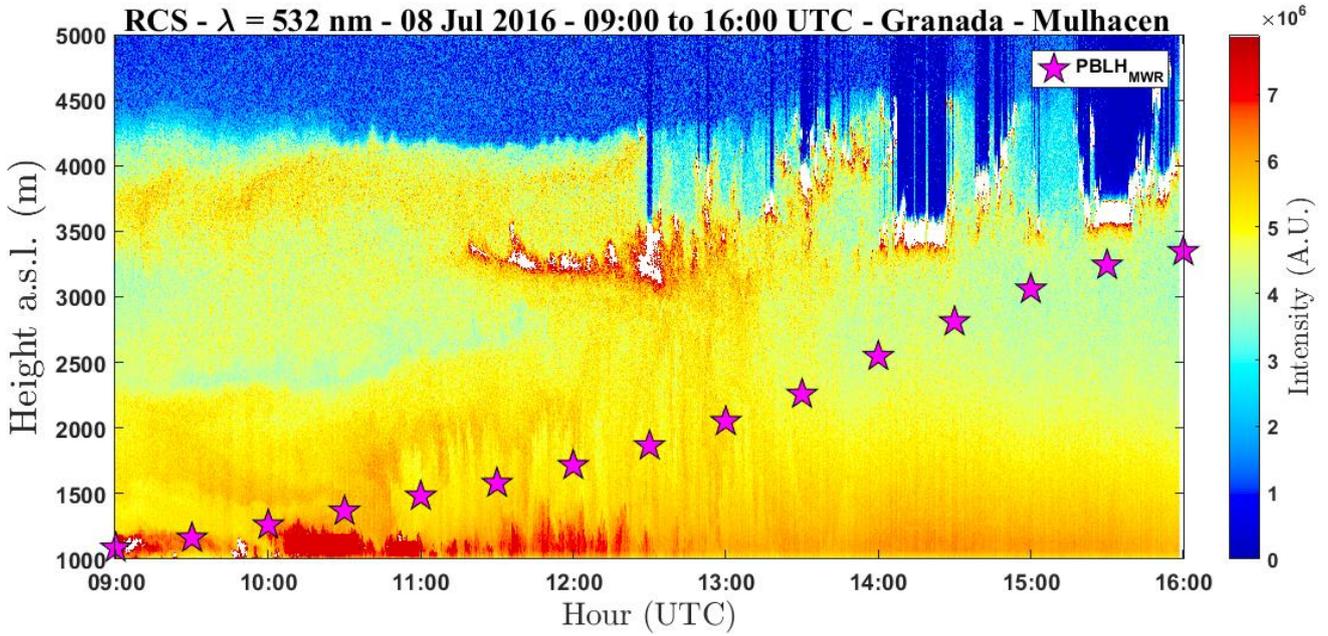


Figure 12 - A – integral time scale [ $\tau_{w'}$ ], B – variance [ $\sigma_{w'}^2$ ], C – skewness [ $S_{w'}$ ], D – net radiation [ $R_n$ ], E – Air surface temperature [blue line] and surface relative humidity [RH – orange line]. In A, B and C black lines and white stars represent air temperature and  $PBLH_{MWR}$ , respectively.

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Figure 13 – Time-Height plot of RCS - 08 July 2016. Pink stars represent  $PBLH_{MWR}$ .

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### Granada – 08 Jul – 12-13 UTC

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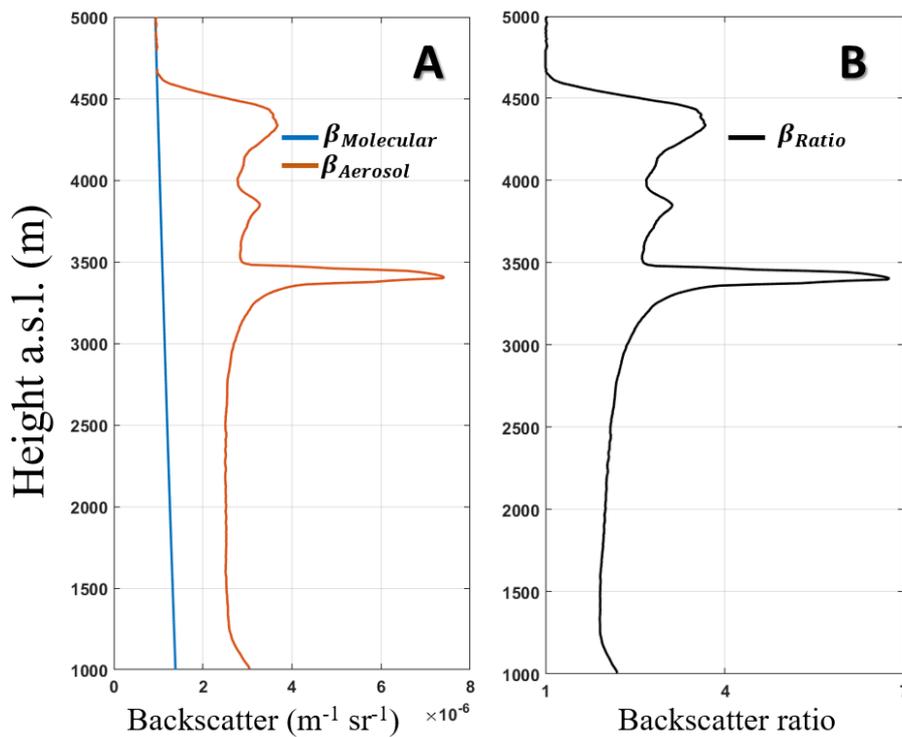
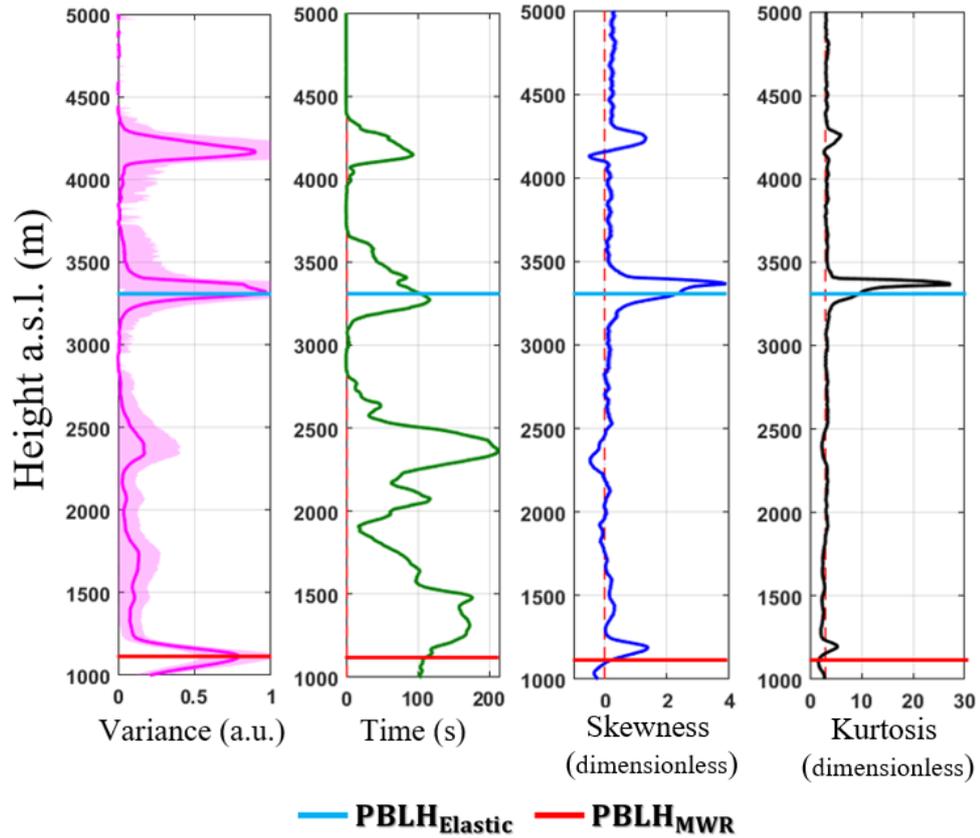


Figure 14 – (A)  $\beta_{Molecular}$  (blue line) and  $\beta_{Aerosol}$  (orange line). (B)  $\beta_{Ratio}$  (black line). All profiles were obtained from the 532 nm lidar signal

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Figure 15 - Statistical moments obtained from elastic lidar data at 11 to 12 UTC - 08 July 2016. From left to right: variance [ $\sigma_{RCS'}^2$ ], integral time scale [ $\tau_{RCS'}$ ], skewness [ $S_{RCS'}$ ] and kurtosis [ $K_{RCS'}$ ].

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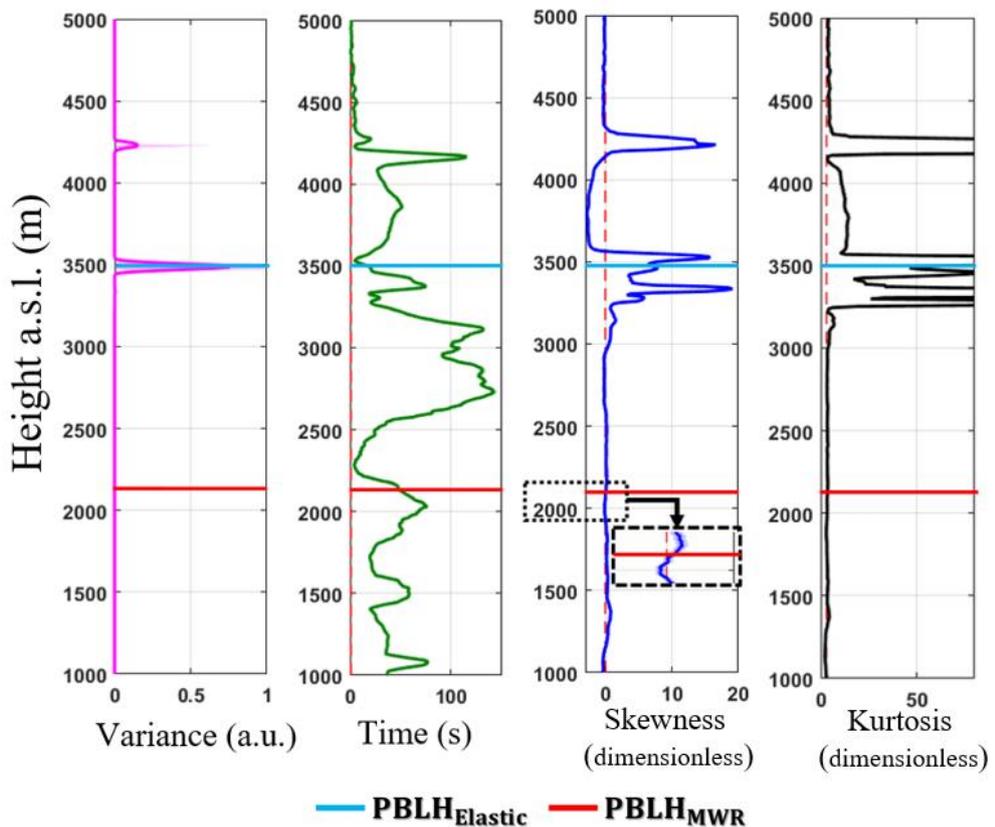


Figure 16 - Statistical moments obtained from elastic lidar data at 12 to 13 UTC - 08 July 2016. From left to right: variance [ $\sigma_{RCS'}^2$ ], integral time scale [ $\tau_{RCS'}$ ], skewness [ $S_{RCS'}$ ] and kurtosis [ $K_{RCS'}$ ].