

1 **Light absorption properties and potential sources of particulate brown carbon in**
2 **the Pearl River Delta region of China**

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19 **Abstract:**

20 Brown carbon (BrC) is a special type of organic aerosols (OA), capable of absorbing solar
21 radiation from near-ultraviolet (UV) to visible wavelengths, which may lead to an increased
22 aerosol radiative effect in the atmosphere. While high concentrations of OAs have been
23 observed in the Pearl River Delta (PRD) region of China, the optical properties and
24 corresponding radiative forcing of BrC in the PRD are still not well understood. In this work,
25 we conducted a set of comprehensive measurements of atmospheric particulate matter from 29
26 November 2014 to 2 January 2015 to investigate aerosol compositions, optical properties,
27 source origins and radiative forcing effects at a suburban station in Guangzhou. The particle
28 absorption Ångström exponent (AAE) was deduced and utilized to differentiate light absorption
29 by BrC from that by black carbon (BC). The results showed that the average absorption
30 contributions of BrC were $34.1 \pm 8.0\%$ at 370 nm, $23.7 \pm 7.3\%$ at 470 nm, $16.0 \pm 6.7\%$ at 520 nm,
31 $13.0 \pm 5.4\%$ at 590 nm and $8.7 \pm 4.3\%$ at 660 nm. A sensitivity analysis of the evaluation of the
32 absorption Ångström exponent of BC (AAE_{BC}) was conducted based on the Mie theory
33 calculation assuming that the BC-containing aerosol was mixed with the core-shell and external
34 configurations. The corresponding uncertainty in AAE_{BC} was acquired. We found that
35 variations in the imaginary refractive index (RI) of the BC core can significantly affect the
36 estimation of AAE_{BC} . However, AAE_{BC} was relatively less sensitive to the real part of the RI
37 of the BC core and was least sensitive to the real part of the RI of the nonlight absorbing shell.
38 BrC absorption was closely related to aerosol potassium cation content (K^+), a common tracer
39 of biomass burning emissions, which was most likely associated with straw burning in the rural
40 area of the western PRD. Diurnal variation in BrC absorption revealed that primary organic
41 aerosols had a larger BrC absorption capacity than secondary organic aerosols (SOAs).
42 Radiative transfer simulations showed that BrC absorption may cause $2.3 \pm 1.8 \text{ W m}^{-2}$ radiative
43 forcing at the top of the atmosphere (TOA) and contribute to $15.8 \pm 4.4\%$ of the aerosol warming
44 effect. A chart was constructed to conveniently assess the BrC radiative forcing efficiency in
45 the studied area with reference to certain aerosol single-scattering albedo (SSA) and BrC
46 absorption contributions at various wavelengths. Evidently, the BrC radiative forcing efficiency
47 was higher at shorter wavelengths.

Keywords: Brown carbon, Black carbon, Absorption Ångström exponent, Radiative forcing, Pearl River Delta.

48 **1 Introduction**

49 Black carbon (BC) and organic carbon (OC) are dominant carbonaceous aerosol components
50 that mainly originate from biomass burning in a global scale (Bond et al., 2004) and have
51 attracted great environmental concerns in rapidly developing regions. Carbonaceous aerosols
52 can not only exert adverse impacts on public health, similar to other particulate matters, but
53 also significantly affect the terrestrial radiation balance with enormous uncertainties. In
54 previous studies, BC was often considered to be the only light-absorbing species (Andreae and
55 Gelencser, 2006), and OC was believed to only be able to scatter light, i.e., causing a cooling
56 effect (Bond et al., 2011). Nevertheless, it has been reported that some fraction of organic
57 aerosols (OAs) may also specifically contribute to light absorption from the near-ultraviolet
58 (UV) to visible wavelength range, which is referred to as brown carbon (BrC) (Kirchstetter et
59 al., 2004). BrC optical properties are strongly affected by its chemical composition and physical
60 structure, which are related to different BrC sources. BrC can originate not only from direct
61 emissions, including smoldering, biomass burning or any type of incomplete fuel combustion
62 process (T. C. Bond et al., 1999; Cheng et al., 2011), but also from secondary organic aerosol
63 formation processes, such as aqueous phase reactions in acidic solutions (Desyaterik et al., 2013)
64 or volatile organic compound (VOC) oxidation (Laskin et al., 2015; Sareen et al., 2010). In
65 addition, BrC could have a complicated molecular composition and intermix with other
66 substances, such as BC, non-absorbing OAs and other inorganic materials, making it
67 complicated to investigate BrC optical properties.

68 BC absorption is commonly assumed to be covering the full wavelength-range. However, the
69 light absorption property of BrC is believed to be more wavelength-dependent, which can be
70 represented by distinct absorption Ångström exponent (AAE) values, i.e., the power exponent
71 of the light absorption coefficient. A typical threshold for the AAE of BC (AAE_{BC}) of 1.6 has
72 been recommended to distinguish BrC from BC (Lack and Cappa, 2010), and the AAE of BrC
73 has been reported as having a wider range (2 to 7) (Hoffer et al., 2005). Based on the difference
74 in the wavelength dependence of light absorption between BC and BrC, previous studies have

75 applied the AAE method to differentiate light absorption by BrC through multiwavelength
76 optical measuring apparatus, such as 3-wavelength Photoacoustic Soot Spectrometer (PASS-3)
77 (Lack and Langridge, 2013), multiwavelength Aethalometer (Olson et al., 2015), etc. Based on
78 the AAE method, the BrC absorption contribution has been estimated to be approximately 6 to
79 41% of total aerosol light absorption at short wavelengths, e.g., at 370 nm and 405 nm
80 (Washenfelder et al., 2015). A uniform AAE_{BC} from ~ 300 nm up to ~ 700 nm (Moosmüller et
81 al., 2011) is commonly used when evaluating the BrC absorption contribution using the AAE
82 method. However, it has been reported that the AAE_{BC} can be influenced by the mixing state,
83 BC core size and morphology (Lack and Cappa, 2010). The lensing effect of the coating shell
84 may enhance BC light absorption, the magnitude of which may also depend on wavelength and
85 can alter the value of AAE_{BC} (Liu et al., 2018). Moreover, different values of AAE_{BC} have been
86 found in the near-infrared and UV ranges (Wang et al., 2018). Therefore, using the default value
87 of $AAE_{BC} = 1$ may lead to uncertainty in BrC absorption coefficient estimation.

88 Quantifying BrC optical absorption accurately is essential to interpret aerosol optical depth
89 (AOD), and the corresponding aerosol direct radiative forcing (DRF) on the atmosphere can
90 also be evaluated if the single-scattering albedo (*SSA*) and extinction coefficient of aerosols are
91 known. The estimation of the DRF of BrC has shown a distinct seasonal variation, indicating
92 the influence of different absorption properties of BrC (Arola et al., 2015). A global simulation
93 study indicated that the average warming effect at the TOA caused by BrC absorption can be
94 up to 0.11 W m^{-2} , corresponding to $\sim 25\%$ of that predicted from BC absorption only (Feng et
95 al., 2013).

96 During the last three decades, rapid economic development has led to severe air pollution
97 problems in the PRD region (Chan and Yao, 2008). With rapid increases in the automobile
98 population and factories, high loadings of SOAs have often been observed (Tan et al., 2016b).
99 Biofuel usage may also play a significant role during wintertime air pollution events in the PRD,
100 indicating that the contribution from BrC light absorption cannot be ignored (Wu et al., 2018).
101 Recently, BrC light absorption has been quantified by Qin et al. (2018) using the AAE method
102 in the PRD region. OA chemical composition was simultaneously measured by a high-
103 resolution time-of-flight aerosol mass spectrometer, and it was found that organic aerosols
104 originating from biomass burning possessed the most intense absorption capability and were

105 largely responsible for BrC absorption. Qin et al. (2018) also suggested that correlations
106 between OA chemical compositions and BrC absorption were wavelength-dependent.
107 In this paper, we applied the homologous AAE differentiation method to quantify the fraction
108 of aerosol light absorption by BrC using the measurements from a seven-wavelength
109 Aethalometer. The potential error incurred with this methodology was determined using Mie
110 theory simulations, especially for various complex refractive indexes of the BC core and the
111 coating material. The correlation between BrC light absorption and water-soluble ions, which
112 is used as the source tracer, was employed to identify potential BrC sources. An atmospheric
113 radiative transfer model has also been applied to evaluate the impact of BrC on direct radiative
114 forcing using surface-based aerosol optical properties and satellite-based surface-albedo data.
115 The magnitudes of aerosol radiative forcing at the top of the atmosphere due to BC and BrC
116 were also individually quantified.

117 **2 Methodology**

118 **2.1 Sampling site**

119 Field observations were conducted at the Panyu station (113°21'E, 23°00'N), which is a
120 monitoring site of the Chinese Meteorological Administration (CMA) Atmospheric Watch
121 Network (CAWNET) that is located on the summit of Dazhengang Mountain (approximately
122 150 m above sea level) in Guangzhou, China. Fig. 1 shows the location of the Panyu site, which
123 is situated at the center of the PRD and is separated from residential areas by at least 500 m.
124 Some agricultural fields can be found to the west of the site. Although there were no significant
125 pollution sources nearby, this suburban site was strongly affected by pollutants transported from
126 the urban area of Guangzhou and crop residual fires transported from the rural area of the PRD.
127 The field campaign was conducted from 29 November 2014 to 2 January 2015. During the
128 measurement period, aerosol light scattering and extinction, BC concentration, particle number
129 size distribution (PNSD), OC concentration, and the water-soluble ion concentrations of PM_{2.5}
130 were continuously monitored.

131 **2.2 Measurements and data analysis**

132 All instruments were housed inside the 2nd floor measurement room of a ~5-m tall, 2-story
133 building. The ambient sample was taken on the roof by a 2-m long, 12.7-mm OD stainless steel
134 inlet, and a PM_{2.5} cyclone sampler was also used. The metal tubing was thermally insulated and

135 maintained at a constant temperature of $\sim 25^{\circ}\text{C}$. A diffusion drier was also used in-line to dry
136 the relative humidity (RH) of the air sample below 30% before further analysis.

137 **2.2.1 Measurements of relevant species**

138 A TSI-3936 scanning mobility particle sizer (SMPS) and a TSI-3321 aerodynamic particle sizer
139 (APS) were utilized to measure the 10 to 500 nm mobility diameter and 0.5 to 2.5 μm
140 aerodynamic diameter of the PNSD, respectively. The aerodynamic diameters of the APS data
141 were converted into mobility diameters using a material density of 1.7 g cm^{-3} . A detailed data
142 merging method has been described by Cheng et al. (2006). Furthermore, the pipe diffusion
143 loss of SMPS has been corrected using the empirical formula proposed by Kulkarni et al. (1996).
144 An AE-33 Aethalometer (Magee Scientific Inc.) was utilized for BC mass concentration
145 measurement, which was derived from optical attenuation using a mass absorption cross section
146 (MAC) of $7.77\text{ m}^2\text{ g}^{-1}$ at 880 nm. The sensitivity of AE-33 was approximately $0.03\text{ }\mu\text{g m}^{-3}$ for
147 a 1-min time resolution and a 5.0 liter per minute (LPM) sample flow rate.

148 The $\text{PM}_{2.5}$ mass concentration was measured by an Environment Dust Monitor (Model
149 EDM180, GRIMM Inc.), which monitored the mass concentration of $\text{PM}_{2.5}$ and PM_{10}
150 simultaneously.

151 Water-soluble ions (potassium (K^+), calcium (Ca^{2+}), magnesium (Mg^{2+}), chloride (Cl^-), sulfate
152 (SO_4^{2-}), nitrate (NO_3^-), and ammonium (NH_4^+) were measured with the Monitor for AeRosols
153 and Gases in Air (MARGA) (Model ADI2080, Metrohm Inc.), which is an online analyzer for
154 semi-continuous measurements of gases and water-soluble ions in aerosols (Li et al., 2010).
155 The MARGA was automatically calibrated with standard internal solutions during field
156 measurement. The MARGA utilized its own $\text{PM}_{2.5}$ sampling system provided by the
157 manufacturer.

158 The OC mass concentration was measured by a Sunset online OC/EC analyzer (Model RT-4)
159 with a laser transmittance-based charring correction (Wu et al., 2018). The sample flow rate of
160 the OC/EC analyzer was maintained at 8 LPM. For each measurement cycle (one hour),
161 samples were collected onto a quartz filter within the first 45 min and then thermal-optically
162 analyzed during the remaining 15 min. First, OC was completely volatilized in oxygen-free
163 helium with a stepwise ramped temperature (600°C and 840°C). In the second stage, the
164 temperature was reduced to 550°C , and then EC and pyrolyzed carbon (PC) were combusted

165 in an oxidizing atmosphere (10% oxygen in helium), while the temperature was increased up
 166 to 870 °C step by step. The CO₂ converted from all of the carbon components was then
 167 quantified by a nondispersive infrared absorption CO₂ sensor (Lin et al., 2009). An internal
 168 calibration peak made by 5% methane in helium was applied to quantify OC and EC. To correct
 169 the PC converted from OC to EC, a tunable pulsed diode laser beam was used to monitor the
 170 laser transmittance through the quartz filter throughout the thermal-optical analysis (Bauer et
 171 al., 2012).

172 **2.2.2 Measurements of optical properties**

173 Light extinction by aerosols at 532 nm was detected using a cavity ring-down aerosol extinction
 174 spectrometer (CRDS) (Model XG-1000, Hexin Inc.) by measuring the decay times of laser
 175 intensity through the aerosol-containing sample and the filtered background air sample under
 176 the same conditions. The extinction coefficient (σ_{ext}) was calculated using the procedure
 177 described by Khalizov et al. (2009).

178 Aerosol total scattering (σ_{sp}) was measured by a TSI-3563 integrated nephelometer at three
 179 wavelengths (i.e., 450 nm, 550 nm, and 700 nm) and was calibrated with CO₂ following the
 180 manual instructions. Particle free air was used to check the nephelometer background signal
 181 once every two hours. The scattering coefficients at other wavelengths were extrapolated using
 182 the following equations:

$$183 \quad SAE = - \frac{\ln(\sigma_{scat, \lambda_0}) - \ln(\sigma_{scat, 550nm})}{\ln(\lambda_0) - \ln(550)} \quad (1)$$

$$184 \quad \sigma_{scat}(\lambda) = \sigma_{scat}(550) \cdot \left(\frac{\lambda}{550}\right)^{-SAE} \quad (2)$$

185 where $\lambda_0=450$ nm is for wavelengths less than 550 nm and $\lambda_0=700$ nm is for wavelengths greater
 186 than 550 nm. The corresponding time series of extinction coefficients, scattering coefficients,
 187 and SSA at 532 nm was displayed in Fig. S1.

188 The Aethalometer is also used for multi-wavelength light absorption measurements in this study.
 189 The seven-wavelength aerosol light attenuation coefficients (σ_{ATN}) were converted into aerosol
 190 light absorption coefficients (σ_{abs}) using Eq. (3) (Coen et al., 2010), where k is the parameter
 191 that accounts for the loading effect, ATN is the light attenuation through the filter with sample
 192 loading and C_{ref} is a fixed multiple scattering parameter.

$$193 \quad \sigma_{abs} = \frac{\sigma_{ATN}}{(1 - k \cdot ATN) \cdot C_{ref}} \quad (3)$$

194 The real-time k value was retrieved using the dual-spot loading correction algorithm developed
 195 by Drinovec et al. (2015). The detailed formula of ATN can also be found in Drinovec et al.
 196 (2015). C_{ref} is considered a constant that strongly depends on the filter matrix effect. However,
 197 some studies have suggested that C_{ref} may vary with wavelength (Arnott et al., 2005; Segura et al.,
 198 2014). For internal combustion engines and biomass burning, C_{ref} at 370 nm was expected to be
 199 approximately 12% and 18% less than C_{ref} at 532 nm for the aerosol component, respectively
 200 (Schmid et al., 2006). Different ambient observations also showed that C_{ref} may have regional
 201 specificity, even though it was retrieved by the same methodology (Coen et al., 2010). In this
 202 study, $C_{ref}=3.29$ was used in Eq. (3) at each wavelength, and this value was derived from the
 203 slope of σ_{ATN} measured by the Aethalometer vs. σ_{abs} , which was deduced from the CRDS and
 204 nephelometer measurements. This C_{ref} was also very similar to the C_{ref} of 3.48 determined from
 205 an inter-comparison study between an Aethalometer and a photoacoustic soot spectrometer
 206 during a field campaign conducted in the PRD region in 2004 (Wu et al., 2009).

207 The BC light absorption at certain wavelengths was derived from the absorption coefficient σ_{abs}
 208 according to Beer-Lambert's law, and its variation between different pairs of wavelengths (i.e.,
 209 $\sigma_{abs,BC,\lambda}$) is denoted by the absorption Ångström exponent (AAE) equation developed by
 210 Ångström (1929):

$$211 \quad \sigma_{abs,BC,\lambda} = \sigma_{abs,BC,\lambda_0} \times (\lambda_0 / \lambda)^{-AAE_{BC}} \quad (4)$$

212 It has been suggested that the AAE of BC may vary between short and long wavelength ranges
 213 (Lack and Cappa, 2010); hence, applying a wavelength-independent AAE_{BC} may lead to
 214 uncertainties in the BC absorption calculation from one wavelength to another. In this work,
 215 the light absorptions of BC at various wavelengths were retrieved by a modified wavelength-
 216 dependent AAE differentiation method conducted by Wang et al. (2018):

$$217 \quad \sigma_{abs,BC,\lambda_1} = \sigma_{abs,BC,880nm} \times \left(\frac{880}{\lambda_1}\right)^{AAE_{BC,520-880nm}} \quad (5.1)$$

$$218 \quad \sigma_{abs,BC,\lambda_2} = \sigma_{abs,BC,880nm} \times \left(\frac{880}{520}\right)^{AAE_{BC,520-880nm}} \times \left(\frac{520}{\lambda_2}\right)^{AAE_{BC,370-520nm}} \quad (5.2)$$

219 Here, $\sigma_{abs,BC,\lambda_1}$ represents the absorption coefficient due to only BC greater than 520 nm, and
 220 $\sigma_{abs,BC,\lambda_2}$ represents the absorption coefficient of BC less than 520 nm. $AAE_{BC, \lambda_i-\lambda_{i+1}}$ (i=1, 2
 221 and 3) represents the AAE of BC between a longer and shorter wavelength at $\lambda_i=880, 520$ and
 222 370 nm and was calculated as:

223
$$AAE_{BC,\lambda_i-\lambda_{i+1}} = \frac{\ln(\sigma_{abs,BC,\lambda_i}) - \ln(\sigma_{abs,BC,\lambda_{i+1}})}{\ln(\lambda_i) - \ln(\lambda_{i+1})} \quad (6)$$

224 Accordingly, BrC absorption at a certain wavelength λ ($\sigma_{abs,BrC,\lambda}$) was equal to the value of total
 225 aerosol absorption ($\sigma_{abs,\lambda}$) minus BC absorption ($\sigma_{abs,BC,\lambda}$):

226
$$\sigma_{abs,BrC,\lambda} = \sigma_{abs,\lambda} - \sigma_{abs,BC,\lambda} \quad (7)$$

227 The light absorption data at 880 nm ($\sigma_{abs,880nm}$) were selected to represent BC absorption
 228 ($\sigma_{abs,BC,880nm}$), which shall not be affected by BrC (Drinovec et al., 2015). It has been reported
 229 that the dust-related contributions of PM_{2.5} were normally less than 5% in wintertime in
 230 Guangzhou; therefore, the influence from dust could be negligible in this study (Huang et al.,
 231 2014).

232 2.2.3 Estimation of AAE_{BC}

233 Traditionally, AAE_{BC} was believed to be close to 1.0 (Bodhaine, 1995), which has been
 234 commonly used for BC measurements (Olson et al., 2015). However, studies have
 235 demonstrated that AAE_{BC} can be affected by the refractive index of coating materials, mixing
 236 state, morphology, and BC core size (Liu et al., 2015). Therefore, using the default AAE_{BC} = 1
 237 may lead to uncertainty in BrC absorption estimation. To obtain the correct AAE_{BC}, a series of
 238 Mie theory calculations were conducted using a simplified core-shell model (Bohren and
 239 Huffman, 1983; Wang et al., 2018). We used a modified BHCOAT code and BHMIE code to
 240 calculate the aerosol optical properties of the core-shell and external mixture at different
 241 wavelengths (Cheng et al., 2006). In the Mie theory, a particle is taken as a perfect
 242 homogeneous sphere, and its extinction and scattering efficiencies, $Q_{ext,Mie,\lambda}$ and $Q_{scat,Mie,\lambda}$,
 243 respectively, are expressed as (Mie, 1908; Seinfeld and Pandis, 1998):

244
$$Q_{ext,Mie,\lambda} = \frac{2}{\alpha^2} \sum_{n=1}^{\infty} [(2n+1)Re(a_n + b_n)] \quad (8)$$

245
$$Q_{scat,Mie,\lambda} = \frac{2}{\alpha^2} \sum_{n=1}^{\infty} [(2n+1)(|a_n|^2 + |b_n|^2)] \quad (9)$$

246 where $\alpha = \pi D_p/\lambda$ is the size parameter; a_n and b_n are functions of the complex refractive
 247 index (RI) and α in the Riccati-Bessel form, respectively. Re in Eq. (8) denotes that only the
 248 real part of RI is taken. The absorption efficiency ($Q_{abs,Mie,\lambda}$) is thus the difference between
 249 the extinction and scattering efficiencies:

250
$$Q_{abs,Mie,\lambda} = Q_{ext,Mie,\lambda} - Q_{scat,Mie,\lambda} \quad (10)$$

251 Then, the absorption coefficient $\sigma_{abs,Mie,\lambda}$ was obtained by the following (Bricaud and Morel,
252 1986):

253
$$\sigma_{abs,Mie,\lambda} = \int Q_{abs,Mie,\lambda} \cdot \left(\frac{\pi}{4} D_p^2\right) \cdot N(\log D_p) \cdot d \log D_p \quad (11)$$

254 where $N(\log D_p)$ is the PNSD function. A two-component parameterization of dry particles,
255 i.e., the BC core and the nonlight-absorbing species, was applied to calculate aerosol optical
256 properties here (Wex et al., 2002). \tilde{m}_{core} represents the RI of the BC core, and \tilde{m}_{non} represents
257 the RI of nonlight-absorbing particles.

258 In a realistic atmosphere, some nonlight-absorbing particles may exist independently without
259 BC (Liu et al., 2013; Cheung et al., 2016). In this work, the portion of nonlight-absorbing
260 particles at a certain size (D_p) was determined by our previous measurements at the same site
261 using a Volatility Tandem Differential Mobility Analyzer (V-TDMA), during which completely
262 vaporized (CV) particles at 300°C were referred to as nonlight-absorbing particles that
263 externally mixed with other BC-containing particles. Thus, the PNSD of CV particles
264 ($N(\log D_p)_{CV}$) and BC-containing particles ($N(\log D_p)_{BC}$) can be given by the following
265 equations:

266
$$N(\log D_p)_{CV} = N(\log D_p)_{measure} \cdot \Phi(D_p)_{N,CV} \quad (12)$$

267
$$N(\log D_p)_{BC} = N(\log D_p)_{measure} \cdot (1 - \Phi(D_p)_{N,CV}) \quad (13)$$

268 where $N(\log D_p)_{measure}$ is the PNSD of the measured particles from SMPS and APS.
269 $\Phi(D_p)_{N,CV}$ was the number fraction of CV particles in different size bin.

270 A previous study applied three kinds of BC mixture models to calculate the aerosol optical
271 properties, including external, homogenously internal and core-shell mixtures (Bohren and
272 Huffman, 2007; Seinfeld and Pandis, 1998). To quantify the mixing state of BC, r_{ext} was defined
273 as the mass fraction of externally mixed BC (M_{ext}) in total BC (M_{BC}):

274
$$r_{ext} = \frac{M_{ext}}{M_{BC}} \quad (14)$$

275 Tan et al. (2016) suggested that two extreme conditions of external and core-shell mixtures
276 comprised the actual mixing state of BC in the PRD. Hence, we simply divided the PNSD of
277 BC into the PNSD from an external mixture of BC and a core-shell mixture of BC. The PNSDs

278 of externally mixed BC particles and core-shell mixed BC particles were referred to by the
 279 following equations with a given r_{ext} .

$$280 \quad N(\log D_p)_{ext} = N(\log D_p)_{BC} \cdot f_{BC} \cdot r_{ext} \quad (15)$$

$$281 \quad N(\log D_p)_{core-shell} = N(\log D_p)_{BC} \cdot (1 - f_{BC} \cdot r_{ext}) \quad (16)$$

282 f_{BC} was defined as the BC volume fraction in the BC-containing particle volume, which can be
 283 converted from the BC mass concentration:

$$284 \quad f_{BC} = \frac{M_{BC}}{\rho_{BC} \cdot \sum_{D_p} N(\log D_p)_{BC} \cdot \left(\frac{\pi}{6} \cdot D_p^3\right)} \quad (17)$$

285 where ρ_{BC} is the density of BC and is assumed to be 1.5 g cm^{-3} (Ma et al., 2012); M_{BC} is the BC
 286 mass concentration derived from the multi-angle absorption photometer (MAAP), which was
 287 obtained by an empirical formula from the Aethalometer that measured the BC concentration
 288 ($M_{BC,AE}$), as proposed by Wu et al. (2009):

$$289 \quad M_{BC} = 0.897 \cdot M_{BC,AE} - 0.062 \quad (18)$$

290 The PNSDs of externally mixed nonlight-absorbing particles and externally mixed BC particles
 291 were input into the BHMIE code, and the PNSD of the core-shell mixed particles was imported
 292 into the BICOAT code. Another critical parameter for the core-shell model was the diameter
 293 of the BC core. For the simplified core-shell model we applied, the visualization was that a
 294 homogeneous BC core sphere was encapsulated in a shell of non-absorbing coating (Bohren
 295 and Huffman, 2007). Without size-resolved coating thickness measurements, core-shell mixed
 296 particles simply assumed that cores with the same diameter had the same coating thickness.
 297 Therefore, the diameter of the BC core was calculated as follows:

$$298 \quad D_{core} = D_p \cdot \left(\frac{f_{BC} - f_{BC} \cdot r_{ext}}{1 - f_{BC} \cdot r_{ext}} \right)^{\frac{1}{3}} \quad (19)$$

299 D_{core} and D_p are inputted as parameters into a_n and b_n , respectively, which was described by
 300 Bohren and Huffman (2007). The corresponding time series of size distribution of the derived
 301 external BC and internal BC core were illustrated in Fig. S2. Thus, the $\sigma_{abs,BC,Mie, \lambda_i}$ values of
 302 all six wavelengths were calculated through the Mie model, and then the AAE_{BC} values of those
 303 five wavelengths were obtained using Eq. (6). The performance of this empirically determined
 304 calculation method has been compared with other possible BC mixing schemes in details (see
 305 Table 1).

306 **2.2.4 Atmospheric radiative transfer model**

307 In this work, the Santa Barbara DISORT Atmospheric Radiative Transfer (SBDART) model
 308 was employed to estimate the DRF of BrC absorption, i.e., its effects on the downward and
 309 upward fluxes (F in W m^{-2}) of solar radiation at the TOA. SBDART is a software tool that can
 310 be used to compute plane-parallel radiative transfer under both clear and cloudy conditions
 311 within the atmosphere. More details about this model have been described by Ricchiazzi et al.
 312 (1998). Both ground measurements and remote sensing data were used in the simulation. The
 313 surface albedo was derived from a 500 m resolution MODIS BRDF/albedo model parameter
 314 product (MCD43A3, daily). The MCD43A3 products are the total shortwave broadband black-
 315 sky albedo (α_{BSA}) and white-sky albedo (α_{WSA}), while the actual surface albedo (α) was
 316 computed from a linear combination of α_{WSA} and α_{BSA} , which were weighted by the diffuse ratio
 317 (r_d) and direct ratio ($1-r_d$), respectively:

$$318 \quad \alpha = (1 - r_d) \cdot \alpha_{BSA} + r_d \cdot \alpha_{WSA} \quad (20)$$

319 r_d was obtained from an exponential fit of Eq. (21) based on empirical observations (Stokes and
 320 Schwartz, 1994; Roesch, 2004):

$$321 \quad r_d = 0.122 + 0.85e^{-4.8\mu_0} \quad (21)$$

322 where μ_0 is the cosine of the zenith angle, which is calculated by the model for any specified
 323 date, time, and latitude and longitude of the site. The surface-based aerosol optical properties,
 324 including the aerosol light absorption coefficients of both BC and BrC, i.e., differentiated from
 325 each other under the assumption of uniform AAE_{BC} , along with the nephelometer-measured
 326 aerosol scattering coefficients, were used to calculate the SSA at different wavelengths
 327 according to Eq. (22),

$$328 \quad SSA(\lambda) = \frac{\sigma_{scat, \lambda}}{\sigma_{abs, BrC, \lambda} + \sigma_{abs, BC, \lambda} + \sigma_{scat, \lambda}} \quad (22)$$

329 which was then used in the model calculation. Finally, the AOD and asymmetry factor (ASY)
 330 at 440, 675 and 870 nm were derived from the Aerosol Robotic Network (AERONET)
 331 measurements at the Hong Kong Polytechnic University site (Holben et al., 1998), which is
 332 approximately 115 km to the southeast of the Panyu site. The tropical atmospheric profile was
 333 used in the SBDART model based on the prevailing weather conditions in the PRD. The aerosol
 334 DRF (ΔF) was calculated as the difference between the downward and upward radiation fluxes:

$$335 \quad \Delta F = F \downarrow - F \uparrow \quad (23)$$

336 **3 Results and discussion**

337 **3.1 Aerosol light absorption**

338 The AAE_{BC} is widely defined as the uniform representation of the wavelength dependence of a
339 BC particle (Olson et al., 2015). In reality, AAE_{BC} may vary significantly with BC containing
340 aerosols of different sizes, mixing states, and morphologies (Scarnato et al., 2013; Lack and
341 Langridge, 2013). In fact, some studies showed that the AAE of a large-size, pure BC core may
342 be less than 1.0 (Liu et al., 2018) and that the AAE of BC coated with a non-absorbing shell
343 may be larger than that under uniformity (Lack and Cappa, 2010).

344 It has been suggested that a significant fraction of smaller size particles is non BC-containing
345 (Ma et al., 2017; Cheung et al., 2016). BC and non-BC materials can also be externally or
346 internally mixed. Although size resolved BC measurements were not available during this work,
347 we have conducted size resolved Volatility Tandem Differential Mobility Analyzer (V-TDMA)
348 measurements at 300°C for 40, 80, 110, 150, 200 and 300 nm, respectively, during an earlier
349 field campaign (February 2014) at the same site as in this work. At 300°C, all non-BC particle
350 will be completely vaporized (CV) and thus the portion of non BC particles at such size, denoted
351 as $\Phi_{N,CV}$, can be determined. The average $\Phi_{N,CV}$ values were 0.384, 0.181, 0.180, 0.158, 0.143
352 and 0.137, corresponding to 40, 80, 110, 150, 200 and 300 nm (see Fig. S3), respectively
353 (Cheung et al., 2016; Tan et al., 2016a). The size-independent $\Phi_{N,CV}$ was interpolated linearly
354 with these six diameters. For particle size larger than 300 nm and less than 40 nm, $\Phi_{N,CV}$ values
355 were set to 0.137 and 0.384, respectively. For particle size larger than 300 nm and less than 40
356 nm, $\Phi_{N,CV}$ values were set to 0.137 and 0.384, respectively. Accordingly, the complete
357 distribution of $\Phi_{N,CV}$ for the whole PNSD was obtained. The mixing states of BC particles were
358 also estimated here, i.e., the mass portion of externally mixed BC with respect to total BC,
359 denoted as r_{ext} . The value of r_{ext} was taken as 0.58, which was obtained using an optical closure
360 method during a previous field experiment at this site (Tan et al., 2016a). During the following
361 Mie theory calculation, a fixed refractive index ($\tilde{m}_{core}=1.80-0.54i$, $\tilde{m}_{non}=1.55-10^{-7}i$) was
362 adopted for the whole size range. Accordingly, the calculated BC absorption at 880 nm (Abs_{880})
363 was 21.869 Mm^{-1} , which is reasonably close to the measured mean value of 21.199 Mm^{-1} . To
364 further validate our calculation scheme (Base Case), we have considered several extreme cases.
365 Case 1: BC is completely externally mixed with non-BC particles, i.e., $\Phi_{N,CV} = 0$ and $r_{ext} = 1$;

366 Case 2: BC is present in every size bin and BC is completely internally mixed with non-BC
367 material, i.e., $\Phi_{N,CV} = 0$ and $r_{ext} = 0$; Case 3: BC is both internally and externally mixed but there
368 is no non BC-containing particles, i.e., $\Phi_{N,CV} = 0$ and $r_{ext} = 0.58$; Case 4: BC is internally mixed
369 with non-BC material and there is non-BC particles present, i.e., $\Phi_{N,CV}$ ranges from 0.384 to
370 0.137 and $r_{ext} = 0$; Case 5: the same as case 4 except assuming a fixed non-BC to BC ratio of
371 0.5, i.e., $\Phi_{N,CV} = 0.5$, $r_{ext} = 0$; Case 6: the same as case 5 except that some externally mixed BC
372 is also present, i.e., $\Phi_{N,CV} = 0.5$, $r_{ext} = 0.58$. The calculation results are listed in Table 1. Evidently,
373 case 1 (complete externally mixed) will significantly underestimate the measured Abs_{880} ,
374 indicating that most BC particles were not likely externally mixed at the Panyu site. Complete
375 internal mixing state (case 2, 4, and 5), on the contrary, would substantially overestimate the
376 BC absorption regardless the form of BC core distribution function. However, when the r_{ext}
377 were considered (case base, 3, and 6), the calculated Abs_{880} values were all very close to the
378 measured value.

379 When the AAE_{BC} was assumed to be uniform, the campaign-averaged σ_{BrC} values were
380 $17.6 \pm 13.7 \text{ Mm}^{-1}$ at 370 nm, $9.7 \pm 7.9 \text{ Mm}^{-1}$ at 470 nm, $5.8 \pm 5.1 \text{ Mm}^{-1}$ at 520 nm, $4.0 \pm 3.5 \text{ Mm}^{-1}$ at
381 590 nm and $2.3 \pm 2.1 \text{ Mm}^{-1}$ at 660 nm. At the corresponding wavelengths, BrC absorption
382 contributed $26.2 \pm 8.5\%$, $20.0 \pm 7.3\%$, $14.3 \pm 6.5\%$, $11.7 \pm 5.3\%$, and $7.8 \pm 4.1\%$ to the total aerosol
383 absorption, respectively. When the AAE_{BC} was applied as the result of the Mie model
384 calculation, the corrected campaign-averaged $\sigma_{abs,BrC}$ values were $23.5 \pm 17.7 \text{ Mm}^{-1}$ at 370 nm,
385 $11.8 \pm 9.5 \text{ Mm}^{-1}$ at 470 nm, $6.7 \pm 5.7 \text{ Mm}^{-1}$ at 520 nm, $4.6 \pm 3.9 \text{ Mm}^{-1}$ at 590 nm and $2.6 \pm 2.3 \text{ Mm}^{-1}$
386 at 660 nm. At the corresponding wavelengths, BrC absorption contributed $34.1 \pm 8.0\%$,
387 $23.7 \pm 7.3\%$, $16.0 \pm 6.7\%$, $13.0 \pm 5.4\%$, and $8.7 \pm 4.3\%$ to the total aerosol absorption (see Fig. 2),
388 respectively. Evidently, aerosol light absorption was predominantly due to BC; however, BrC
389 also played a significant role, especially at shorter wavelengths. Table 2 shows the
390 intercomparison of BrC light absorption in the near UV range between this work and other
391 studies in the East Asian region. Clearly, the reported values vary substantially, and our result
392 is toward the lower end of values. Figure S4 displayed the time series of particle AAE measured
393 by the Aethalometer and AAE_{BC} was derived from Mie model calculation. The AAE_{BC} was
394 almost always lower than AAE, indicating appreciable BrC light absorption at the Panyu site.

395 **3.2 Uncertainty in BC and BrC optical differentiation**

396 Theoretically, the magnitude of BC absorptions can be affected by both parts of the complex
397 refractive indexes (RIs); thus, AAE_{BC} may also vary with the RIs of both the BC core and
398 coating shell. In fact, RI was also one of the least known properties of BC and other coating
399 materials with negligible absorbing capabilities. The refractive index of the BC core (\tilde{m}_{core})
400 displays a wide range of variations (Liu et al., 2018). Typically, the real and imaginary parts of
401 the RI can vary from 1.5 to 2.0 and 0.5 to 1.1, respectively. In addition, the shell was assumed
402 to consist of non-absorbing material in the core-shell model, i.e., its imaginary RI was set to be
403 close to zero (10^{-7}). The real part of the non-absorbing material RI (\tilde{m}_{non}) may vary from 1.35
404 to 1.6 due to the presence of OA (Zhang et al., 2018; Redmond and Thompson, 2011) and
405 inorganic salts (Erlick et al., 2011). Hence, it is necessary to investigate the uncertainties
406 associated with the variations in AAE_{BC} by varying the RIs of both the BC core and the non-
407 absorbing materials.

408 Figure 3 shows the impacts of RI on the evaluations of AAE_{BC} based on core-shell and external
409 configuration, where the RI of the BC core was set to be constant, i.e., $\tilde{m}_{core}=1.80-0.54i$, and
410 the real part of \tilde{m}_{non} varied from 1.35 to 1.6 at an interval of 0.05, with the imaginary part of
411 \tilde{m}_{non} set at 10^{-7} . As shown in Fig. 3a, the calculated AAE_{BC} for the core-shell model was higher
412 than 1.0 at longer wavelengths (520 to 880 nm) and lower than 1.0 at shorter wavelengths (370
413 to 520 nm) (the red line in Fig. 3 denotes $AAE_{BC}=1$). The averaged $AAE_{BC,370-520nm}$ ranged from
414 0.84 to 0.87, and the $AAE_{BC,520-880nm}$ ranged from 1.07 to 1.15, indicating that the $AAE_{BC,520-}$
415 $880nm$ appeared to be more sensitive to the shell's real part than $AAE_{BC,370-520nm}$. Even if the shell
416 material was assumed to be non-absorbing, the variation in the real RI of the shell, which was
417 referred to as the real part of \tilde{m}_{non} , still led to changes in the shell's refractivity and
418 correspondingly altered its lensing effect, causing a change in AAE_{BC} . Meanwhile, $AAE_{BC,370-}$
419 $520nm$ and $AAE_{BC,520-880nm}$ generally increased with an increasing real part of the shell. In Fig. 3b,
420 under the externally mixed conditions, $AAE_{BC,370-520nm}$ and $AAE_{BC,520-880nm}$ were both less than
421 1.0. The average $AAE_{BC,370-520nm}$ was 0.33, and the average $AAE_{BC,520-880nm}$ was 0.63. These
422 values were far less than the values under core-shell mixture conditions. In the external mixture
423 model, the BC core and nonlight-absorbing materials were assumed to exist dependently, and
424 then the optical properties of these two components were considered separately. Therefore,
425 altering the real part of the externally mixed non-absorbing material would not affect the light

426 absorption property of the BC core or AAE_{BC} .

427 The impacts of the BC core on AAE_{BC} are shown in Fig. 4, where the refractive index of
428 nonlight-absorbing materials was assumed to be $\tilde{m}_{non}=1.55-10^{-7}i$ and \tilde{m}_{non} was wavelength-
429 independent. Fig. 4 was obtained with a core-shell mixture model (Fig. 4a and 4b) and an
430 external mixture model (Fig. 4c and 4d) by varying the real part of \tilde{m}_{core} from 1.5 to 2.0 with
431 a step of 0.05 and varying the imaginary part of the \tilde{m}_{core} from 0.4 to 1.0 with a step of 0.05,
432 respectively. As shown in Figs. 4a and 4b, for the core-shell mixture, the averaged $AAE_{BC,370-}$
433 $520nm$ ranged from 0.55 to 0.99, and the averaged $AAE_{BC,520-880nm}$ ranged from 0.84 to 1.27. The
434 AAE_{BC} at a certain wavelength generally increased when increasing the real part of \tilde{m}_{core} but
435 decreased when increasing the imaginary part of \tilde{m}_{core} . The AAE_{BC} appeared to be more
436 sensitive to the imaginary part of \tilde{m}_{core} than the real part of \tilde{m}_{core} because the imaginary part
437 of \tilde{m}_{core} was directly related to the light-absorbing properties of particles. In Fig. 4c and 4d,
438 for the external mixture, the averaged $AAE_{BC,370-520nm}$ ranged from 0.04 to 0.45 and the averaged
439 $AAE_{BC,520-880nm}$ ranged from 0.28 to 0.79, while the averaged $AAE_{BC,370-520nm}$ and $AAE_{BC,520-}$
440 $880nm$ were both less than 1.0. Similar to the core-shell mixture, the $AAE_{BC,520-880nm}$ increased
441 when increasing the real part of \tilde{m}_{core} but decreased when increasing the imaginary part of
442 \tilde{m}_{core} . However, the variation patterns of $AAE_{BC,370-520nm}$ were different from those of $AAE_{BC,520-}$
443 $880nm$. The $AAE_{BC,370-520nm}$ values were not changed by altering the real part of \tilde{m}_{core} within the
444 low imaginary part of \tilde{m}_{core} , whereas the $AAE_{BC,370-520nm}$ values still increased when increasing
445 the real part of \tilde{m}_{core} within the high imaginary part of \tilde{m}_{core} . A possible explanation was that
446 the externally mixed BC core had weak light absorption within the low imaginary part of \tilde{m}_{core} ,
447 causing the $AAE_{BC,370-520nm}$ values to be insensitive to the real part of \tilde{m}_{core} . The $AAE_{BC,520-880nm}$
448 values were higher than the $AAE_{BC,370-520nm}$ values regardless of whether in they were for the
449 core-shell mixture or external mixture. In addition, the AAE_{BC} values conducted by the core-
450 shell mixture were higher than those conducted by the external mixture.

451 Figure 4 demonstrates that the variation in the imaginary RI of the BC core has the most
452 significant impact on the estimated AAE_{BC} , indicating that the chemical component of BC
453 emitted from different sources leads to a large uncertainty in AAE_{BC} estimation. At the same
454 time, the influence arising from varying the real RI of the BC core was relatively moderate.
455 Nevertheless, Fig. 3 demonstrated that change in the real RI of the non-absorbing materials

456 caused the least/no impact compared to that caused by the variations in the complex RI of the
457 BC core.

458 It should be pointed out that most BC-containing particles are often observed as fractal rather
459 than spherical in shape (Katrinak et al., 1993). Because the Mie model assumes that all particles
460 are spherical, it may lead to potential uncertainty for the estimation of AAE_{BC} and BrC
461 absorption contributions. Moreover, the externally mixed soot aggregates were “chain-like” or
462 “puff-like” in the PRD dry season (Feng et al., 2010), in which the fractal dimension (D_f) was
463 between 1.5 and 2.0. Coating soot aggregates were likely sphere (D_f approaches 3) from the
464 high-resolution transmission electron microscopy (TEM) measurements taken in Hongkong
465 ((Zhou et al., 2014)). A soot aggregate sensitivity study with the superposition T-matrix method
466 indicated that using the assumption of volume-equivalent spheres for the soot aggregates may
467 result in an overestimation of approximately up to 15% and an underestimation of
468 approximately up to 50% in the predicted 870 nm light absorption when the D_f is between 1.5
469 and 3.0 (Liu et al., 2008). However, it should be recognized that the complex shapes or positions
470 of the BC core inside the particle make it impractical to be numerically simulated in the exact
471 details. By far the Mie model with a core-shell configuration would be the most practical and
472 effective simulation scheme for BC particle optical property simulation.

473 Furthermore, we have performed Monte Carlo simulations to evaluate the uncertainties of the
474 Mie calculation performed during this work. In the simulation, a sequence of random numbers
475 or errors were applied to the input parameters, and then the corresponding uncertainties of
476 particle light absorption and AAE_{BC} by were computed using the Mie model. Five hundred of
477 reiteration were conducted during the simulation such that the random errors will be normally
478 distributed. The standard deviations (σ) of all input parameters are listed in Table S1. In order
479 to cover the effect of extreme value, we used a range of $\pm 3\sigma$, or a confidence level of 99%, in
480 the Monte Carlo simulation. Table S2 listed the Monte Carlo simulation results, i.e., the average
481 relative standard deviations (σ_{Mie}) of the calculated BC light absorption at 880 nm (Abs_{880}),
482 $AAE_{BC,370-520}$, and $AAE_{BC,520-880}$. The uncertainties of the calculated Abs_{880} , $AAE_{BC,370-520}$, and
483 $AAE_{BC,520-880}$ at 2 times of σ_{Mie} , i.e., at a confidence coefficient of 95%, were approximately
484 $\pm 31\%$, $\pm 16\%$, and $\pm 13\%$, respectively. Figure S5a showed the time series of the uncertainties
485 of Abs_{880} , $AAE_{BC,370-520}$ and $AAE_{BC,520-880}$ from Monte Carlo simulation for the campaign period.

486 These uncertainties will certainly be propagated into the calculated BrC absorption
487 contributions, too. Hence, we also estimated the corresponding uncertainties in the BrC
488 absorption contribution results, as shown in Fig. S5b. Accordingly, the averaged lower limits
489 of BrC absorption contributions were $26.8\% \pm 9.1\%$ at 370 nm, $17.5\% \pm 8.1\%$ at 470 nm,
490 $10.1\% \pm 7.3\%$ at 520 nm, $8.5\% \pm 5.8\%$ at 590 nm and $5.3\% \pm 4.5\%$ at 660 nm, respectively, and
491 the averaged upper limits of BrC absorption contribution ratios were $40.7\% \pm 7.2\%$ at 370 nm,
492 $29.5\% \pm 6.7\%$ at 470 nm, $21.1\% \pm 6.2\%$ at 520 nm, $17.3\% \pm 5.2\%$ at 590 nm and $12.0\% \pm 4.1\%$ at
493 660 nm, respectively.

494 **3.3 Characteristics of BrC light absorption, water-soluble ions and OC concentrations**

495 Globally, BrC has been observed to be highly correlated with biomass and biofuel burning
496 emissions (Laskin et al., 2015). Since large quantities of sylvite are present in biomass burning
497 particles, the K^+ abundance has often been used as a biomass burning tracer (Levine, 1991).
498 Figure 5 presents the time series of the OC mass concentration, K^+ concentration, and BrC
499 absorption from 29 November 2014 to 2 January 2015 at the Panyu site. The range of the OC
500 concentration obtained from the OC/EC online analyzer was from 1.5 to $65.2 \mu\text{g cm}^{-3}$, and the
501 campaign average was $12.5 \pm 7.3 \mu\text{g cm}^{-3}$. The BrC absorption hourly mean data were between
502 0.2 and 123.2 Mm^{-1} , and the campaign average was $23.5 \pm 17.7 \text{ Mm}^{-1}$. On the other hand, the
503 average K^+ concentration was $1.0 \pm 0.7 \mu\text{g cm}^{-3}$ (ranging from 0 to $5.4 \mu\text{g cm}^{-3}$). Clearly, similar
504 trends among OC, K^+ , and BrC absorption can be seen during this field campaign (Fig. 5).

505 To investigate the origins of these observed OC, K^+ , and BrC, wind rose plots (as shown in Fig.
506 6) were generated for OC, K^+ , and BrC absorption, respectively. All three panels of Fig. 6
507 consistently show that the three substances were associated with the same wind pattern. For the
508 entire campaign period, the highest values of OC, K^+ , and $\sigma_{abs,BrC,370nm}$ were mostly associated
509 with southwesterly winds with a relatively low wind speed ($\sim 2 \text{ m s}^{-1}$). The relatively higher OC
510 and K^+ concentrations were highly related to the seasonal straw burning in the countryside of
511 the PRD located to the west of the Panyu station. In contrast, OC and K^+ concentrations during
512 periods with easterly winds were substantially lower than those during periods with westerly
513 winds. The wind rose plot of $\sigma_{abs,BrC,370nm}$ is shown in Fig. 6c. Similar to OC and K^+ , $\sigma_{abs,BrC,370nm}$
514 showed higher values under weak ($< 2 \text{ m s}^{-1}$) westerly winds and lower values from the north
515 and south, indicating that BrC absorption was likely attributed to local sources and was

516 accumulated under calm wind conditions. Figure S6 showed the 3-day backward trajectory and
517 the fire counts for 5 to 7 (Fig. S6a), 12 to 14 (Fig. S6b) and 24 to 26 (Fig. S6c) in November
518 2014, representing low loading, moderate loading and high loading period. Clearly, the high
519 loading period concurred with stagnant air movement and higher fire counts, indicating the
520 contribution from open fire burning sources. However, there was a detectable difference among
521 the three rose plots of Fig. 6 in the maximum concentration direction. A possible explanation
522 was that although biomass burning emissions were believed to be the dominant and primary
523 source of OC, K^+ , and BrC, their emission ratios were highly variable and may change with the
524 type of biofuel and burning condition and may even vary during different stages of burning
525 (Burling et al., 2012). Although biomass burning emissions contain substantial light-absorbing
526 BrC, further atmospheric aging processes may significantly reduce its light-absorbing
527 capability (Satish et al., 2017). Moreover, secondary formation may also lead to BrC formation
528 inside these primary aerosols, such as humic-like substances formed through aqueous-phase
529 reactions, which have been suggested to be an important component of BrC (Andreae and
530 Gelencser, 2006).

531 To further explore the possible sources of BrC optical absorption, the diurnal variations in OC,
532 K^+ , $\sigma_{abs,BrC,370nm}$, and $\sigma_{abs,BrC,370nm}/OC$ values are plotted in Fig. 7. The diurnal variation in OC
533 at the Panyu site appeared to be dominated by the development of the planetary boundary layer
534 (PBL) height, i.e., primary emissions accumulated at night and were swiftly diluted by vertical
535 mixing in the morning. The slight increase in OC in the afternoon indicated that photochemistry
536 may have still weakly contributed to SOA formation. Fig. 7b shows the diurnal variation in K^+ .
537 Unlike OC, K^+ shows a small peak at approximately 6 AM, which was consistent with breakfast
538 time and was very likely due to cooking activities using biofuel. No lunch and dinner time K^+
539 peaks were observed. The most likely explanation is that the boundary layer height is much
540 higher during lunch and dinner time than in the early morning, providing a much better
541 atmospheric diffusion condition for air pollutants. It is still a common practice to collect straw
542 as biofuel in local rural areas, which can be visually spotted but is not heavily utilized in the
543 region. However, the diurnal profile of $\sigma_{abs,BrC,370nm}$ (see Fig. 7c) shows the combined features
544 of OC and K^+ since both primary and secondary processes affect its intensity. The nighttime
545 increasing trend was most likely attributed to straw burning activities in early winter in nearby

546 rural areas that continued to accumulate within the shallow PBL (Jiang et al., 2013).
547 $\sigma_{abs,BrC,370nm}/OC$, i.e., the mass absorption coefficient of BrC (MAC_{BrC}) (Fig. 7d), showed a
548 relatively flat pattern, with a pronounced dip in the afternoon and higher values at nighttime,
549 which was likely due to enhanced primary emissions and stable stratification at nighttime.
550 Declining trends during the late morning and afternoon hours indicated that the aging process
551 and photochemical production may reduce the light-absorbing capacity of BrC (Qin et al.,
552 2018).

553 Furthermore, Fig. 8 shows the linear regression analysis results used to evaluate the correlations
554 of $\sigma_{abs,BrC,370nm}$ with the OC, K^+ , Ca^{2+} , Mg^{2+} , Cl^- , SO_4^{2-} , NO_3^- , and NH_4^+ concentrations. The best
555 correlations can be found between $\sigma_{abs,BrC,370nm}$ and K^+ ($R^2=0.6148$), followed by those between
556 $\sigma_{abs,BrC,370nm}$ and OC ($R^2=0.4514$), NO_3^- ($R^2=0.4224$) and NH_4^+ ($R^2=0.4656$). Source
557 apportionment analysis of OA and BrC absorption in Beijing and Guangzhou illustrated that
558 biomass burning organic aerosols (BBOAs) correlated well with BrC light absorption (Xie et
559 al., 2018; Qin et al., 2018). Thus, the significant correlation between BrC absorption and K^+
560 reaffirmed that biomass burning was the crucial emission source of BrC observed in this work.
561 Although the geographic location of the observation site was situated in a coastal area and K^+
562 could also be found in sea salt (Pio et al., 2008), it should be noted that the prevailing wind
563 direction during winter was from the north (see Fig. 3), which drives maritime air parcels away
564 from the site. Hence, the effect of sea salt and crustal materials to K^+ was slight, which was
565 demonstrated in the supplementary information as shown in Fig. S7. Other earlier studies also
566 suggested that the sea salt contribution to the K^+ concentrations of $PM_{2.5}$ was trivial in the PRD
567 region during the winter (Lai et al., 2007). Another possible K^+ source was coal combustion.
568 The coal consumption in the PRD region was dominated by coal-fired power plants. The
569 emission from power plants was usually very steady and was less likely to affect the diurnal
570 correlation between K^+ and BrC absorption. As shown in Fig. S8, the ratios of $K^+/PM_{2.5}$ varies
571 approximately from 0.015 and 0.020 and the diurnal profile of $K^+/PM_{2.5}$ shows very little
572 variation. Yu et al. (2018) have suggested that K^+ usually accounted for 2.34-5.49% of $PM_{2.5}$ in
573 the laboratory biomass burning study. However, K^+ was normally lower than 1% of coal
574 combustion $PM_{2.5}$. Therefore, the ratio range of K^+ to $PM_{2.5}$ observed in this work likely
575 indicated aged biomass burning particles. Both nitrogen oxides (NO_x) and ammonia (NH_3) can

576 be found in biomass burning plumes (Andreae and Merlet, 2001). For NO_3^- and NH_4^+ , nitrate
577 can be converted from NO_x through atmospheric reactions, and ammonium may originate from
578 NH_3 . However, similar to the diurnal variation in $\sigma_{abs,BrC,370nm}$, diurnal variations in NH_4^+ and
579 NO_3^- also increased in the afternoon and appeared at nighttime in Fig S8. However, $\text{NO}_3^-/\text{PM}_{2.5}$
580 and $\text{NH}_4^+/\text{PM}_{2.5}$ reached their peaks at noon, indicating that ammonium nitrate formed from the
581 secondary reaction at this time. Along with the reduced boundary layer height and ambient
582 temperature, NO_3^- was accumulated until the photochemical reaction stopped at night. The
583 diurnal variation in NH_4^+ was similar to that in NO_3^- due to the acid/base neutralization reaction.
584 The overlapping of the $\sigma_{abs,BrC,370nm}$, NH_4^+ and NO_3^- diurnal variations would lead to a significant
585 correlation between BrC absorption and NO_3^- or NH_4^+ . High concentrations of Ca^{2+} and Mg^{2+}
586 are often found in dust-related aerosols (Lee et al., 1999). $\sigma_{abs,BrC,370nm}$ showed poor correlations
587 with both Ca^{2+} and Mg^{2+} , indicating that dust-related aerosol components contribute
588 insignificantly to the total aerosol mass loading and, thus, dust may not affect the AAE
589 differentiation method used in this work. Although sulfur dioxide (SO_2) may also be emitted by
590 biomass burning, SO_4^{2-} is often believed to be secondary in nature, and the presence of other
591 intense SO_2 sources (e.g., automobile and industrial emissions) further reduces the correlation
592 between BrC and SO_4^{2-} . Sources of Cl^- include both combustion and sea salt spray (Waldman
593 et al., 1991). Although the prevailing wintertime wind direction was from the north, sea salt
594 can still be carried to the site by a weak sea breeze, and thus, Cl^- may not show considerable
595 correlation with BrC.

596 **3.4 BrC radiative forcing efficiency**

597 The radiative effects of aerosol scattering, BrC absorption, and BC absorption were investigated
598 by the SBDART model. For each investigated variable under cloud-free conditions, we run the
599 model twice to calculate the DRF at the TOA with and without the investigated variable.
600 Accordingly, the difference of ΔF between the two simulations was considered as the radiative
601 effect of the investigated variable. The results showed that the average radiative forcings at the
602 TOA by scattering, BrC absorption, and BC absorption were $-21.4 \pm 5.5 \text{ W m}^{-2}$, $2.3 \pm 1.8 \text{ W m}^{-2}$,
603 and $10.9 \pm 5.1 \text{ W m}^{-2}$, respectively. Furthermore, BrC absorption was attributed to $15.8 \pm 4.4\%$
604 of the warming effect caused by aerosol light absorption, demonstrating the nonnegligible role
605 of BrC in radiative forcing evaluation.

606 We also calculated the BrC radiative forcing efficiency (RFE) under various *SSA* (ranging from
607 0.7 to 0.99) at three wavelengths, i.e., 440 nm, 675 nm, and 870 nm. The RFE was denoted as
608 the radiative forcing normalized by the AOD. The average AOD and ASY at the three
609 wavelengths were 0.365 and 0.691 at 440 nm, 0.212 and 0.632 at 675 nm, and 0.154 and 0.619
610 at 870 nm, respectively. A solar zenith angle of 55° and an average shortwave broadband surface
611 albedo (0.119) were used in the calculation. The results were plotted as a set of RFE lookup
612 charts as a function of the surface BrC absorption contribution (see Fig. 9).

613 In general, for any wavelength, the RFE increased with increasing BrC absorption contribution
614 for a certain *SSA*, indicating that BrC was a more efficient radiative forcing agent due to the
615 preferential absorbance of BrC in a shorter wavelength range. However, for a certain BrC
616 absorption contribution, the RFE increased with decreasing *SSA*, i.e., a higher portion of light-
617 absorbing aerosol components can lead to more efficient radiative forcing. The trend among
618 panels (a), (b), and (c) in Fig. 9 demonstrated that the effect of BrC absorption contribution on
619 RFE was wavelength-dependent, i.e., BrC was a weaker radiative forcing agent at longer
620 wavelengths, which is also consistent with the wavelength-dependent light-absorbing property
621 of BrC. The red stars in Fig. 9 denote the average *SSA* and BrC absorption contribution
622 conditions during this campaign, i.e., 0.029 W m⁻² per unit AOD at 440 nm (Fig. 9a), 0.007 W
623 m⁻² per unit AOD at 675 nm (Fig. 9b), and 0.0002 W m⁻² per unit AOD at 870 nm (Fig. 9c).
624 These results suggested that the average value of RFE decreased distinctly from 440 nm to 870
625 nm not only because of the lower BrC absorption contribution but also because of the
626 wavelength-dependence of the BrC RFE. It should also be noted that the simulations were based
627 on *SSA* measured under dry conditions. Under the typical ambient conditions of the PRD, the
628 *SSA* might be markedly enhanced by aerosol water uptake (Jung et al., 2009), and then, the BrC
629 radiative forcing efficiency might be less. Moreover, Fig. 9 also serves as a lookup table to
630 conveniently assess the BrC radiative forcing efficiency at different wavelengths with different
631 BrC absorption contributions for a certain *SSA*.

632 **4 Conclusion**

633 In this work, light absorption due to BrC in the PRD region of China was quantitatively deduced
634 during the winter season of 2014. The AAE of ambient particles and BC core were derived
635 from the measurements. For ambient particles, AAE_{370-520nm} and AAE_{520-880nm} ranged from 0.81

636 to 2.31 and 0.91 to 2.13, respectively. In the case of BC, $AAE_{BC,370-520nm}$ and $AAE_{BC,520-880nm}$
637 ranged from 0.59 to 0.98 and 0.82 to 1.15, respectively. Using the absorption coefficients of
638 BC calculated according to the Mie theory and the observed total aerosol absorption coefficients,
639 we estimated the AAE_{BC} and hence the BrC absorption contribution for the optically equivalent
640 mixture configuration. The average BrC light absorption contribution ranged from $8.7 \pm 4.3\%$ at
641 660 nm up to $34.1 \pm 8.0\%$ at 370 nm when AAE_{BC} was set as uniform. The sensitivity of AAE_{BC}
642 estimation associated with different RI and mixing state assumptions was further investigated.
643 The results showed that variations in the real RI of the non-absorbing material (1.35 to 1.6) may
644 decrease $AAE_{BC,370-520nm}$ from 0.84 to 0.87 and $AAE_{BC,520-880nm}$ from 1.07 to 1.15 for core-shell
645 mixtures, with an $AAE_{BC,370-520nm}$ of 0.33 and $AAE_{BC,520-880nm}$ of 0.63 for external mixtures.
646 Variations in the core's real RI (1.5 to 2.0) and imaginary RI (0.4 to 1.0) may reduce $AAE_{BC,370-}$
647 $520nm}$ from 0.55 to 0.99 and $AAE_{BC,520-880nm}$ from 0.84 to 1.27 for the core-shell mixture and
648 reduce $AAE_{BC,370-520nm}$ from 0.04 to 0.45 and $AAE_{BC,520-880nm}$ from 0.28 to 0.79 for the external
649 mixture. These results indicate that the optical properties of the BC core and non-absorbing
650 material can significantly affect the accuracy of AAE_{BC} and corresponding BrC absorption
651 contribution estimations. Compared to the values of BrC light absorption coefficient and BrC
652 light absorption contribution from other similar studies conducted in the East Asia region, the
653 BrC measured in this work showed relatively lower values of light absorption coefficient but
654 was found responsible for relatively higher portion of light absorption. It should be noted that
655 the calculated BrC light absorption may vary exponentially with the value of AAE_{BC} . According
656 to Monte Carlo simulations under 95% confidence level, we found that BrC light absorption
657 contribution ratios in this work can range roughly from 18% to 48% at 370 nm, 10% to 37% at
658 470 nm, 3% to 27% at 520 nm, 3% to 22% at 590 nm, and 1% to 16% at 660 nm, respectively.
659 Therefore, proper values of AAE_{BC} have to be carefully obtained for a particular study area,
660 especially needed to be constrained by the BC mass concentration, size-distribution, and mixing
661 state measurements.

662 Additionally, the measurements of major water-soluble inorganic ions (including K^+ , NO_3^- , and
663 NH_4^+) and particulate OC showed consistent features with those of $\sigma_{abs,BrC,370nm}$, implying that
664 BrC was associated with biomass-burning emissions from nearby rural areas. Moreover, the
665 diurnal trend in $\sigma_{abs,BrC,370nm}/OC$ indicated that primary biomass burning emissions can produce

666 intense light-absorbing BrC, while the photochemical aging process may weaken the light-
667 absorbing capability of BrC.

668 Using a radiative transfer model (i.e., SBDART), we estimated the BrC effects on aerosol
669 radiative forcing. The average shortwave aerosol direct radiative forcings due to scattering, BrC
670 absorption, and BC absorption at the TOA were evaluated to be $-21.4 \pm 5.5 \text{ W m}^{-2}$, $2.3 \pm 1.8 \text{ W}$
671 m^{-2} , and $10.9 \pm 5.1 \text{ W m}^{-2}$, respectively. BrC absorption accounted for $15.8 \pm 4.4\%$ of the total
672 shortwave solar absorption warming effect at the TOA, indicating that BrC might be an
673 important climate forcing agent, which is largely neglected in current climate models. To
674 facilitate the estimation of the climate effects of BrC, a set of look-up charts was constructed
675 for the investigated area based on the default tropical atmosphere profile, average surface
676 albedo, average asymmetry factor, and surface-measured aerosol properties (i.e., BrC
677 absorption contribution, *SSA*, and AOD). Therefore, the role of the BrC radiative forcing
678 efficiency at three wavelengths can be conveniently assessed for certain *SSA* and BrC
679 absorption contributions.

680

681 *Author contributions*

682 HT, JZ, YM, and CC designed the experiments, and ZL, LL, YQ, NW, FL, YL, and MC carried
683 out the field measurements and data analysis. ZL and YQ performed Mie theory simulation.
684 ZL, JZ, and HT prepared the manuscript with comments from all coauthors.

685

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





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941 **Table 1.** Inter-comparison of the performance of various Mie-calculation schemes. The base
 942 case is based on the empirical distribution function and mixing states of BC particles obtained
 943 from previous field measurements at the same site. $\Phi_{N,CV}$ denotes the portion of non-BC
 944 particles and r_{ext} is the mass portion of externally mixed BC with respect to total BC. AAE_{BC} is
 945 the absorption Ångström exponent of BC, and the subscript represents the wavelength range.
 946 $Abs_{BrC,370\%}$ and $Abs_{BrC,520\%}$ are the BrC absorption contributions at 370 nm and 520 nm,
 947 respectively. $Calcabs_{880}$ is the calculated absorption at 880 nm using Mie model. $Measabs_{880}$ is
 948 the measured absorption by an Aethalometer at 880 nm. b is the intercept of the regression
 949 analysis between $Measabs_{880}$ and $Calcabs_{880}$, i.e., $Calcabs_{880}=b * Measabs_{880}$. R^2 is the
 950 correlation coefficient of the equation. The refractive index of BC core (\tilde{m}_{core}) and nonlight-
 951 absorbing particles (\tilde{m}_{non}) is set to be 1.80-0.54i and 1.55-10⁻⁷i, respectively (Tan et al., 2016a),

| Case # | Scheme | $\Phi_{N,CV}$ | r_{ext} | AAE | AAE | Abs | Abs | Calc | Meas | b | R^2 |
|--------|---|-------------------|-----------|------------|------------|-----------|-----------|--------------------|--------------------|-------|-------|
| | | | | BC,370-520 | BC,520-880 | BrC,370 % | BrC,520 % | abs ₈₈₀ | abs ₈₈₀ | | |
| Base |  | 0.384 to 0.137 | 0.58 | 0.723 | 0.962 | 34.13% | 15.96% | 21.869 | | 1.019 | 0.979 |
| 1 |  | 0 | 1 | 0.331 | 0.626 | 51.64% | 29.57% | 15.832 | | 0.747 | 0.968 |
| 2 |  | 0 | 0 | 0.856 | 1.128 | 24.76% | 8.28% | 27.827 | | 1.295 | 0.976 |
| 3 |  | 0 | 0.58 | 0.745 | 0.974 | 33.22% | 15.46% | 21.936 | 21.199 | 1.029 | 0.979 |
| 4 |  | 0.384 to 0.137 | 0 | 0.835 | 1.111 | 26.01% | 9.14% | 27.302 | | 1.269 | 0.975 |
| 5 |  | 0.5 | 0 | 0.778 | 1.043 | 29.96% | 12.30% | 24.921 | | 1.150 | 0.968 |
| 6 |  | 0.5 | 0.58 | 0.674 | 0.928 | 36.39% | 17.49% | 20.897 | | 0.977 | 0.975 |

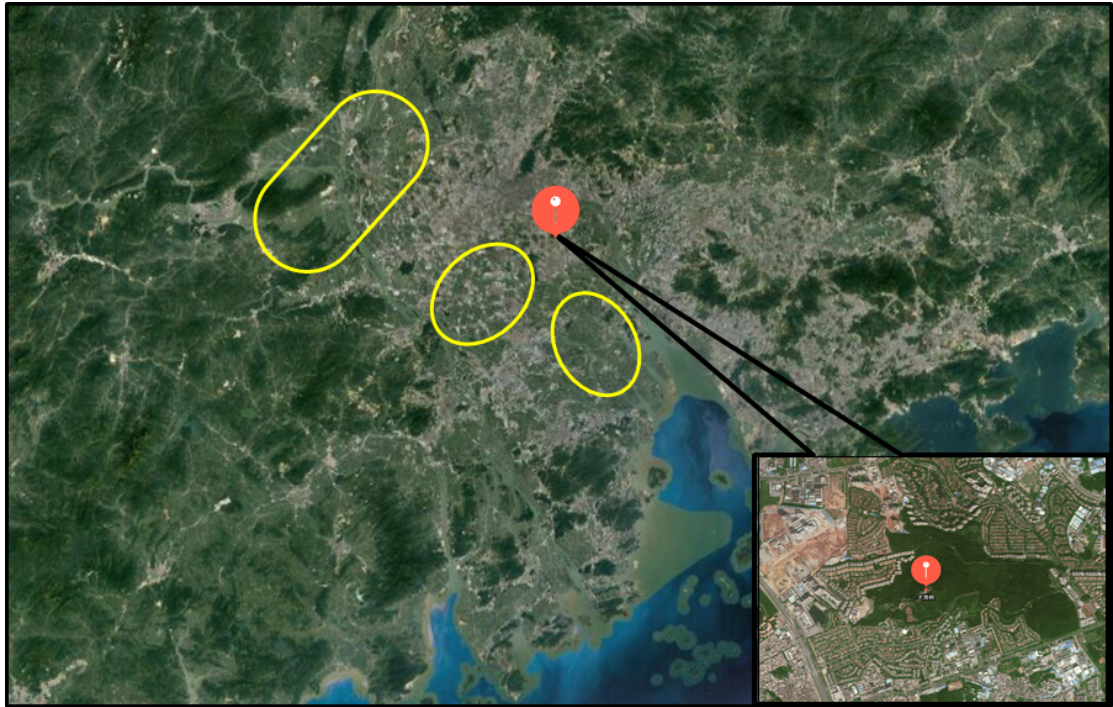
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955 **Table 2.** Observational studies of the BrC light absorption coefficient and contribution in the
 956 near-ultraviolet wavelength range in East Asia.

| Periods | Location | λ (nm) | Mean BrC absorption coefficient | Mean BrC absorption contribution | Instrumentation | Reference |
|---|-------------------|----------------|---|---|----------------------------|--------------------------|
| Nov. 2014 – Jan. 2015 | Guangzhou (China) | 370 | 17.6 Mm ⁻¹ 23.5 Mm ⁻¹ | 26.2% (AAE _{BC} =1) 34.1% (corrected) | Aethalometer r AE-33 | This study |
| Jan. 2014 – Feb. 2014; Sep. 2014 – Oct. 2014 | Shenzhen (China) | 405 | 3.0 Mm ⁻¹ 1.4 Mm ⁻¹ | 11.7% (winter) 6.3% (fall) | PASS-3 | Yuan et al. (2016a) |
| Nov. 2014 | Heshan (China) | 405 | 3.9 Mm ⁻¹ | 12.1% | PASS-3 | Yuan et al. (2016a) |
| Nov. 2016 – Dec. 2016 | Beijing (China) | 370 | 106.4 Mm ⁻¹ 93.8 Mm ⁻¹ | 46% (at the ground) 48% (at 260 m) | Aethalometer r AE-33 | Xie et al. (2018) |
| Jun. 2013 – May. 2016 | Nanjing (China) | 370 | 35.8 Mm ⁻¹ | 16.7% | Aethalometer r AE-31 | Wang et al. (2018) |
| Jan. 2012 | Nagoya (Japan) | 405 | Not detected | 11% (300°C) 17% (400°C) | Thermometer r PASS-3 | (Nakayama et al. (2015)) |

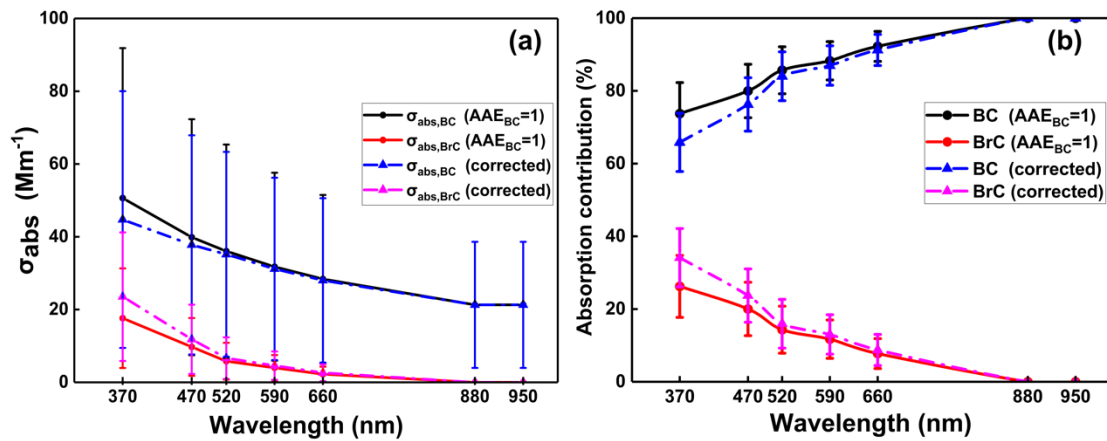


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959 Figure 1. The location of Panyu station (CAWNET) in the PRD region (indicated by the red

960 dot). The plain areas within the yellow circles are the main rural areas of western PRD.

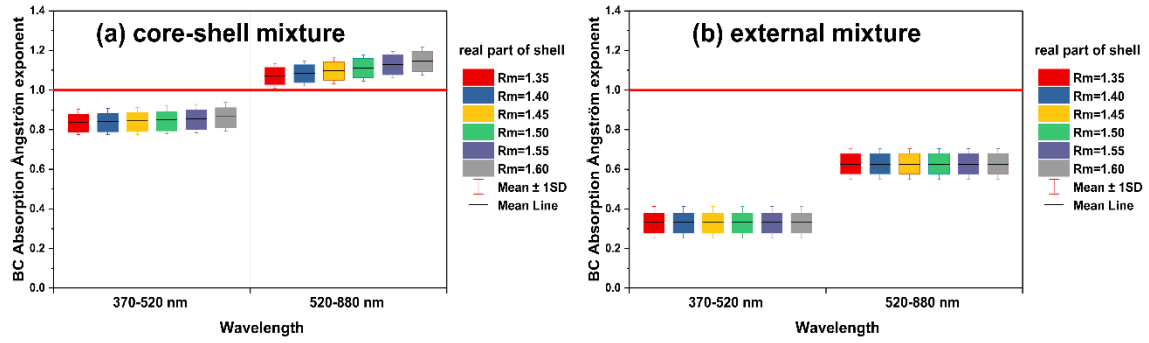
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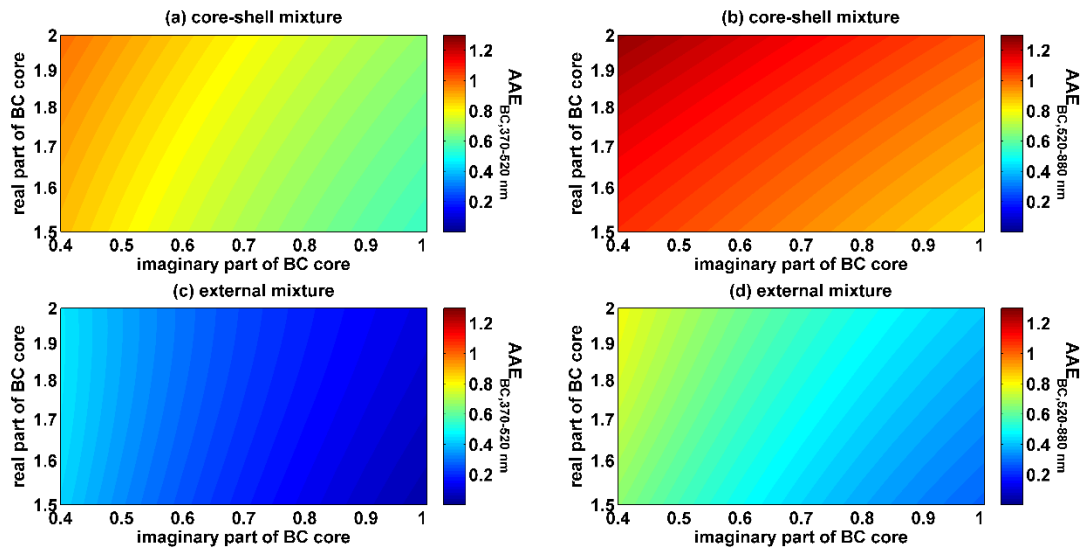
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964 Figure 2. (a) BC and BrC particle average light absorption coefficients at different
 965 wavelengths under different AAE_{BC} assumptions; the whiskers represent an error of
 966 one standard deviation. (b) Contributions of BC and BrC to the total light absorption
 967 coefficient at different wavelengths under different AAE_{BC} assumptions; the whiskers
 968 represent an error of one standard deviation.

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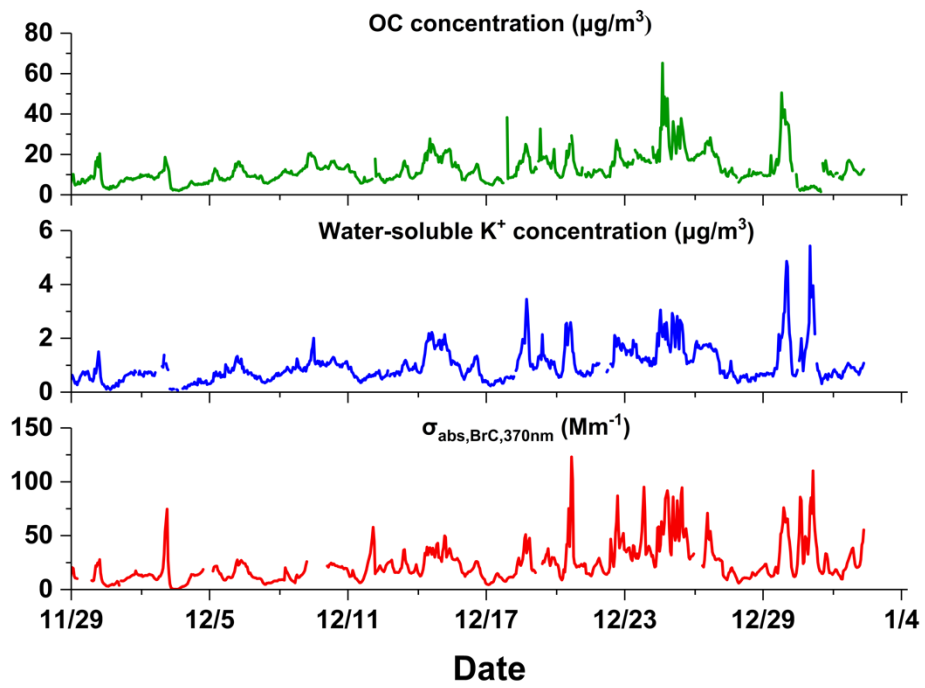
970 Figure 3. Influence of the wavelength-independent refractive index of the non-
 971 absorbing materials on the (a) AAEs of the core-shell mixture and (b) AAEs of the
 972 external mixture with a constant BC core refractive index ($\tilde{m}_{core}=1.80-0.54i$). The
 973 imaginary part of the non(less)-absorbing shell was set to 10^{-7} , while the real part varied
 974 from 1.35 to 1.60. In each panel, the boundaries of the box represent the 75th and 25th
 975 percentiles; the whiskers above and below each box indicate an error of one standard
 976 deviation; the black lines in the boxes denote the average values. In panels a and b, the
 977 red line indicates where $AAE_{BC}=1$.
 978



980

981 Figure 4. Influence of the wavelength-independent refractive index of the BC core on
 982 AAEs with a constant shell refractive index ($\tilde{m}_{shell}=1.55-10^{-7}i$). A core-shell mixture
 983 was used for panels a and b, and an external mixture was used for panels c and d. The
 984 real part of \tilde{m}_{core} varied from 1.5 to 2.0, with a step of 0.05, and the imaginary part of
 985 \tilde{m}_{core} varied from 0.4 to 1.0, with a step of 0.05.

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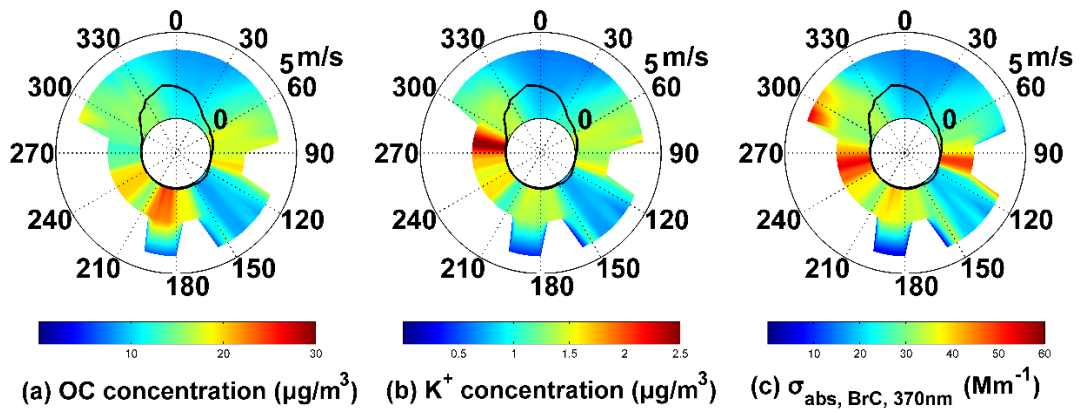
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988 Figure 5. Time series of the OC aerosol mass concentration (green line), water-soluble K^+ mass

989 concentration (blue line), and BrC light absorption (red line).

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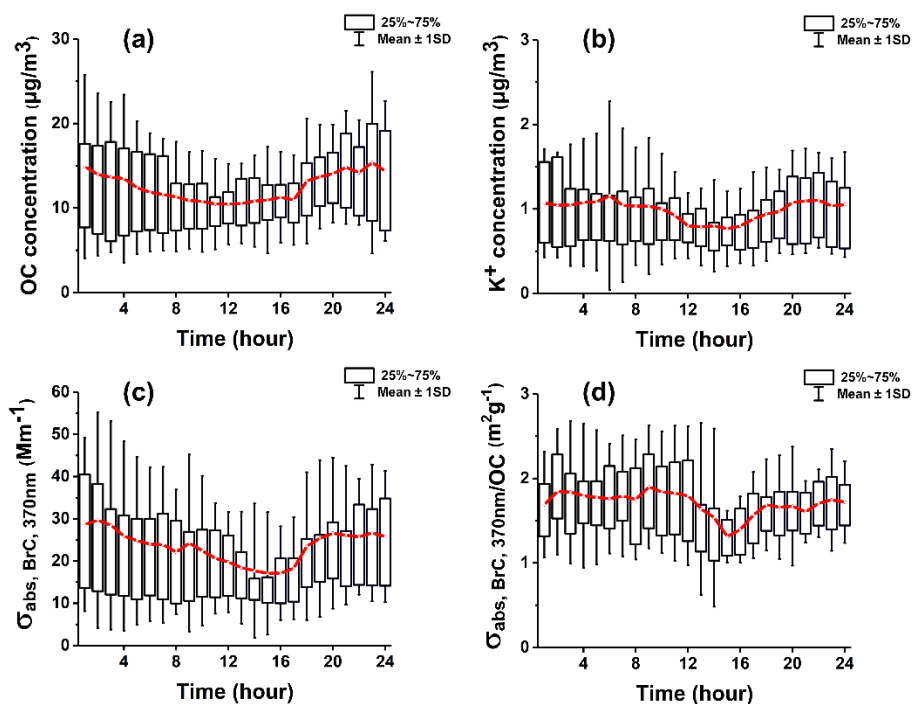
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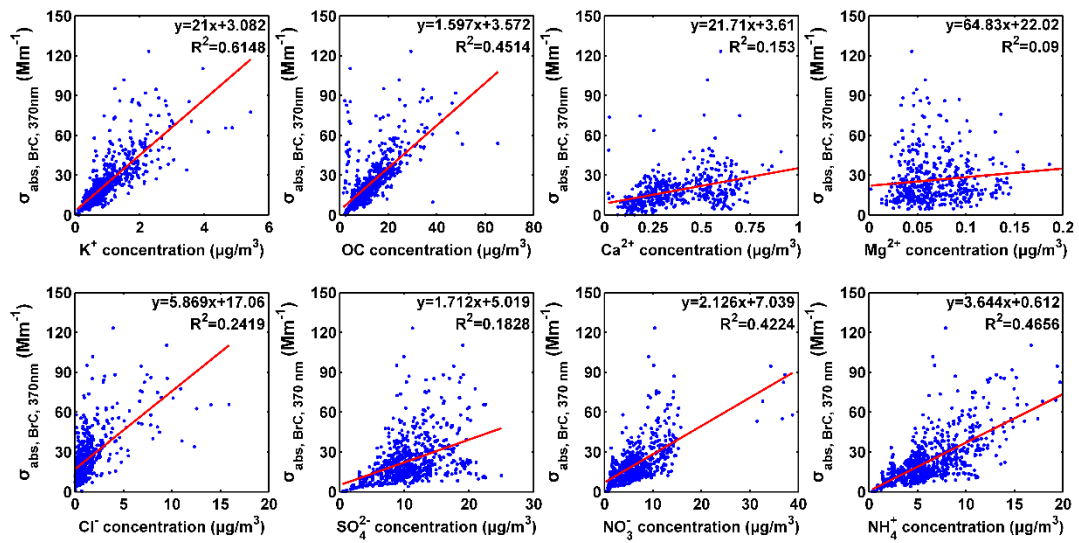
993 Figure 6. Wind rose plots of OC (a), K^+ (b), and $\sigma_{\text{abs, BrC, 370nm}}$ (c). In each panel, the black solid
994 lines denote the frequency of the wind direction. The shaded contour represents the average
995 values of the corresponding species for that wind speed (radial length) and wind direction
996 (transverse direction) in polar coordinates.

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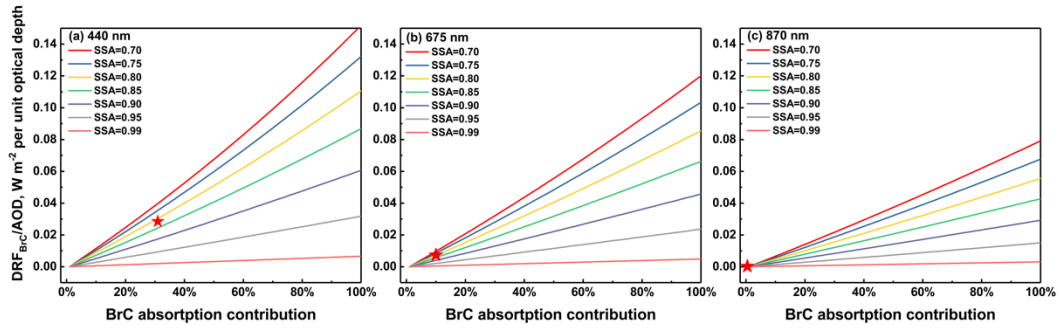
1000 Figure 7. Box-whisker plots of diurnal trends in the OC concentration (a), water-soluble K^+
 1001 concentration (b), $\sigma_{\text{abs,BrC,370nm}}$ (c), and $\sigma_{\text{abs,BrC,370nm}} / \text{OC}$ (d). The red traces represent the
 1002 variation in the average value. The upper and lower boundaries of the box represent the 75th
 1003 and 25th percentiles, respectively; the whiskers above and below each box represent an error
 1004 of one standard deviation.
 1005



1006

1007 Figure 8. Correlations of the BrC absorption coefficient at 370 nm with OC, water-soluble K^+ ,

1008 Ca^{2+} , Mg^{2+} , Cl^- , SO_4^{2-} , NO_3^- , and NH_4^+ aerosol concentrations.



1009

1010 Figure 9. BrC radiative forcing efficiencies, which are defined as the BrC TOA direct radiative
 1011 forcing divided by the AOD, as a function of the BrC to BC absorption ratio and *SSA* measured
 1012 at the surface. The average AOD of the three wavelengths, the average ASY of the three
 1013 wavelengths, a solar zenith angle of 55° , and average shortwave broadband surface albedo were
 1014 used in the calculation. The red star corresponds to the average *SSA* and BrC absorption
 1015 contributions determined from this campaign.