

Mesospheric semidiurnal tides and near-12-hour waves through jointly analyzing observations of five specular meteor radars from three longitudinal sectors at boreal midlatitudes

Maosheng He¹ and Jorge Luis Chau¹

¹Leibniz-Institute of Atmospheric Physics at the Rostock University, Kühlungsborn, Germany

Correspondence: Maosheng He (he@iap-kborn.de)

Abstract.

In the last decades, mesospheric tides have been intensively investigated with observations from both ground-based radars and satellites. Single-site radar observations provide continuous measurements at fixed locations without horizontal information whereas single spacecraft missions provide typically global coverage however with limited temporal coverage at a given location. In this work, combining eight years (2009-2016) of mesospheric winds collected by five specular meteor radars from three different longitudinal sectors at boreal midlatitudes ($49\pm 8.5^\circ\text{N}$), we develop an approach to investigate the most intense global-scale oscillation, namely, at the period $T=12\pm 0.5\text{hr}$. Resolved are six waves: the semidiurnal westward-traveling tidal modes with zonal wavenumber 1, 2, and 3 (SW1, SW2, SW3), the lunar semidiurnal tide M2, and the upper and lower sidebands (USB and LSB) of the 16-day wave nonlinear modulation on SW2. The temporal variations of the waves are studied statistically with a special focus on their responses to sudden stratospheric warming events (SSWs), and on their climatological seasonal variations. In response to SSWs, USB, LSB, and M2 enhance, while SW2 decreases. However, SW1 and SW3 do not respond noticeably to SSWs, contrary to the broadly-reported enhancements in the literature. The USB, LSB, and SW2 responses could be explained in terms of energy exchange through the nonlinear modulation, while LSB and USB might previously have been misinterpreted as SW1 and SW3, respectively. Besides, we find that LSB and M2 enhancements depend on the SSW classification with respect to the associated split or displacement of the polar vortex. In the case of seasonal variations, our results are qualitatively consistent with previous studies, and show a moderate correlation with an empirical tidal model derived from satellite observations.

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1 Introduction

The availability of observations limits the advance of studies on the mesosphere-lower-thermosphere (MLT). In situ MLT observations are available, e.g., through rockets, only on campaign bases, whereas remote detection allows monitoring MLT

perennially and continuously. Two most common approaches of the remote detection are ground-based radars with all-weather applicability, and satellite-based optical instruments with good mobility.

Both ground- and space-based continuous observations have been used to investigate the global-scale MLT waves. Most of these studies were based on single-point analysis techniques and therefore were subject to inherent spatiotemporal ambiguities (following Paschmann and Daly, 1998, here 'point' refers to a geometric element, either stationary or moving, has no extension in the space). Ground-based observations from single radars could yield high-frequency-resolved spectra of MLT parameters, but can not resolve the global-scale structure (e.g., Azeem et al., 2000). On the other hand, space-based sensors, typically on-board slowly precessing polar satellites (e.g., Oberheide et al., 2002), collect data with global coverage but with limited temporal coverage for given locations. They are capable of determining the horizontal scales, which however cannot distinguish instantaneous temporal variations from spatial variations. The obtained frequency spectra are usually Doppler shifted at limited resolution (e.g., Salby, 1982a, b) under the assumption that the tides are static.

To overcome the spatiotemporal ambiguity, specular meteor radars (SMRs) or medium frequency radars from multi-longitudinal sectors had been combined to resolve the horizontal scale of MLT waves at polar latitudes tentatively. A typical procedure is a least square regression (LSR) fitting of longitudinal harmonic functions with preassigned wavenumber to observations from different longitude sectors (e.g., Murphy, 2003). The LSR procedure was used to decompose the most significant global-scale periodicity, namely, the 12hr tidal oscillation, into the migrating mode SW2 (SW_m represents westward-traveling semidiurnal tidal modes with zonal wavenumber m) and nonmigrating modes SW1 and SW3 (mostly at polar latitudes, e.g., Murphy, 2003; Murphy et al., 2006, 2009; Baumgaertner et al., 2006; Manson et al., 2009). However, as sketched in Figure 1, such decomposition gets complicated by the existence of other waves in the vicinity of 12 hr with wavenumbers identical to those of solar tides. These include the semidiurnal lunar tide (M2), and the lower and upper sidebands (LSB and USB) of the nonlinear modulation of the 16-day planetary wave on SW2. Sharing similar periods and same wavenumbers with the tides, these waves are suspected to have contaminated the interpretation of previous studies. Specifically, LSB and USB might have been detected at low-frequency resolutions and misinterpreted as SW1 and SW3 (cf., He et al., 2018a, b), respectively. Additionally, the M2 estimations might have been contaminated by LSB in spectral studies using single-site observational technique (as explained in He et al., 2017a) or by the power leakage from SW2 in low-frequency-resolved spectral analyses (cf., Section 5.1 in He et al., 2018b).

The main purpose of the current study is to develop an approach to unambiguously separate all six waves sketched in Figure 1 using observations of five SMRs at latitudes near 49 °N between 2009 and 2016. Below, section 2 introduces the six waves and the approach. The results are shown in section 3, and used to investigate the six waves statistically, in particular, their responses to sudden stratospheric warming events (SSWs) and their seasonal variations (Sections 4 and 5, respectively). Note that in the current study, we use the term 'responses to SSWs' to refer to the behaviors associated with SSW, which does not imply causative relations between the behaviors and the phenomenon suggested literally by the term 'SSWs', namely, the sudden increase in the temperature.

2 Data analysis

For the current study, we collect the mesospheric wind observations of SMRs at $49 \pm 8.5^\circ\text{N}$ from three longitudinal sectors, namely, east Asian, Europe, and America. As shown in Figure 2, these SMRs are located at Juliusruh (12°E , 55°N , available since 2007), Collm (13°E , 51°N , since 2004), Beijing (116°E , 40°N , since 2009), Mohe (123°E , 54°N , since 2012), and Tavi-
 5 stock (81°W 43°N , since 2002). The radar system at Tavistock is officially known as Canadian Meteor Orbit Radar (CMOR, e.g., Jones et al., 2005). For details of the radars, e.g., working frequency, power, and configuration of antennas, readers are refer to Liu et al. (2016, 2017); Yu et al. (2013), Singer et al. (2013), Jacobi (2012), and Jones et al. (2005).

The current study uses hourly zonal wind derived at a vertical resolution of 2 km according to the algorithm introduced by Hocking et al. (2001) and Stober et al. (2012). For each SMR, we filter oscillations in the wind at periods $11.6 \pm 0.1\text{hr}$,
 10 $12.0 \pm 0.1\text{hr}$, and $12.4 \pm 0.1\text{hr}$, through high-frequency-resolved wavelet spectral analysis. For each period, we decompose the potential waves with different wavenumber by jointly analyzing the spectral coherency between the SMRs.

2.1 Decomposition approach

A *Morlet* wavelet analysis is applied to the zonal wind at a given altitude for each SMR, resulting in spectra $\tilde{W}_{(f,t)}^n$ where f , t , and $n=1,2,\dots,5$ represent the frequency, time, and an index of SMRs, respectively. $\tilde{W}_{(f,t)}^n$ corresponds to the phasor
 15 representation used (e.g., Murphy, 2002; Baumgaertner et al., 2006). We attribute the coherence among $\tilde{W}_{(f,t)}^n$ to waves traveling in the longitudinal direction with zonal wavenumber m_k , ($k=1,2,\dots, K$) and complex amplitude \tilde{a}_k . At given f and t , we fit $\tilde{\mathbf{a}}_k$ from $\tilde{W}_{(f,t)}^n$ following, e.g., equation 5 in He et al. (2018a),

$$\left(\tilde{W}^1, \tilde{W}^2, \dots, \tilde{W}^5\right)' = \tilde{\mathbf{E}}_{5 \times K} \tilde{\mathbf{a}}_{K \times 1} \quad (1)$$

Here, the k -th entry of $\tilde{\mathbf{a}}_k$ is defined as \tilde{a}_k , and the entry of $\tilde{\mathbf{E}}$ in the k -th row and n -th column is defined as $\tilde{\mathbf{E}}_{n,k} := e^{i2\pi m_k \lambda_n}$,
 20 representing the phase of k -th wave detected by the n -th SMR at longitude λ_n . When $K \leq 5$, Equation 1 allows estimating $\tilde{\mathbf{a}}$ with preassigned m_k , as demonstrated in, e.g., Figure 4 in He et al. (2018a). Since our five SMRs are mainly from three distinct longitudinal sectors, our implementation entails $K \leq 3$. Although two of the five radars provide redundant information as they are in the same longitude sector as other SMRs, we use all five for a broader temporal coverage and higher statistical significance. We assign m following Figure 1 for reasons detailed in Section 2.2.

25 Note that in estimating $\tilde{\mathbf{a}}$, we assume that the meridional variation of all waves is negligible among the SMRs. To test this assumption, we ran the climatological tidal model of the thermosphere (CTMT, Oberheide et al., 2011) derived from TIDI and SABER. Semidiurnal components in the zonal wind at 50°N are highly correlated with those at 40°N : the correlation coefficients associated with SW2 and SW1 are 0.94 and 0.99, respectively (not shown here). For the latitude dependence of the semidiurnal tide and its seasonal variation, readers are referred to, e.g., Yu et al. (2015) and Oberheide et al. (2011).

30 In principle, $\tilde{\mathbf{a}}$ could be estimated through the LSR or a short-time Fourier transform (STFT) within a sliding window (e.g., Murphy, 2003; Baumgaertner et al., 2006). Using a Gaussian window with the proper width, the LSR or STFT might even yield results identical to ours. The width of the window, ΔT , determines proportionally the time resolution $\sigma_t \propto \Delta T$, which

is coupled with the frequency resolution σ_f according to the *Fabor's* uncertainty principle $\sigma_t\sigma_f \geq \frac{1}{4\pi}$. The resolution in our wavelet analysis is determined by the *Morlet* factor as specified in Section 2.2.

2.2 Targeting waves and assignment of zonal wavenumber

Tides are characterized by oscillations at periods which are integral fractions of a solar or lunar day. In the atmosphere, the solar tides are primarily forced by daily variation in the absorption of sunlight (Chapman and Lindzen, 1970). At 12hr period, the migrating component SW2 is known to be the dominant tide (e.g., Pancheva and Mukhtarov, 2012), while the nonmigrating components SW1 and SW3 are also frequently reported (e.g., Angelats I Coll and Forbes, 2002; Manson et al., 2009). At the latitude for our study (49°N), SW1 and SW3 are expected to be more intensive than other semidiurnal nonmigrating tides on climatological averages (not shown here) according to the tidal model (cf., Oberheide et al., 2011). These solar tides, according to the classic tidal theory (Chapman and Lindzen, 1970), have amplitudes ~ 20 times larger than those of lunar-gravitationally-forced tides. Despite these theoretical predictions, oscillations at 12.4hr have been clearly detected in the upper atmosphere and ionosphere and explained as the lunar tide M2, particularly around SSWs (e.g., Stening, 2011; Fejer et al., 2011; Chau et al., 2015). The occurrence of M2 was also confirmed by a wavenumber identification using a dual-SMR network ($m=2$ at 12.4hr during SSW 2013, cf., He et al., 2018b). The significant M2 tide was attributed to the lunar forcing resonance due to a shift of a local maximum (namely the so-called *Pekeris* peak) in the atmospheric frequency response, which is supported by a comparison in a numerical experiment using GSWM driven by two specifications of a climatological-mean background atmosphere and that during SSWs (Forbes and Zhang, 2012).

In addition to the M2, also oscillating at the period 12.4 hr is a westward-traveling structure with zonal wavenumber $m=1$, namely, the lower sideband (LSB) of the modulation of the 16-day planetary wave (PW) on SW2 tide (as explicitly detected and explained in He et al., 2018a). LSB's m and f are determined by their parent waves according to the nonlinear interaction resonance conditions $\tilde{\Psi}_{LSB} = \tilde{\Psi}_{SW2}\tilde{\Psi}_{PW}^*$. Here, $\tilde{\Psi}_{\bullet} := e^{i2\pi(f\bullet t + m\bullet\lambda)}$ represents the phase of a wave \bullet (e.g., He et al., 2017a). The 16-day PW is a normal wave, and its intrinsic period of 12.5d is determined by the resonant properties of the atmosphere (e.g., Ahlquist, 1982; Longuet-Higgins, 1968; Madden, 2007; Salby, 1984). Doppler-shifted by the prevailing eastward wind during winter, the PW is observed at a period up to 20 days, with an average of 16 days (for the climatology of the 16-day PW, cf., Luo et al., 2002; Day and Mitchell, 2010). The corresponding LSB occurs in the frequency range of $f^{LSB} = (2 - 1/12.5, 2 - 1/20) d^{-1}$, associated with LSB at $T^{LSB} = 12.4 \pm 0.1$ hr. Similar to LSB, an upper sideband (USB), at $T^{USB} = 11.6 \pm 0.1$ hr and $m=3$, might also be excited by the modulation, following the resonance conditions $\tilde{\Psi}_{USB} = \tilde{\Psi}_{SW2}\tilde{\Psi}_{PW}$ (as explicitly detected in He et al., 2018a). To encompass the most potential LSB and USB periods, in our wavelet analysis (cf., Grossmann et al., 1990), we set the *Morlet* factor to 128 so that the passed frequency band corresponds to 12.4 ± 0.1 hr, 11.6 ± 0.1 hr, and 12.0 ± 0.1 hr. These period bands are narrow enough to prevent power leakage or aliasing between each other.

As sketched in Figure 1, above mentioned six waves occupy three near-12hr periods associated with three zonal wavenumbers. To implement Equation 1 for quantifying the waves, we assume that in comparison with the mentioned waves, other potential waves are negligible at each of the periods. Specifically, we assume that at $T=12.0 \pm 0.1$ hr the most important waves are tides SW1, SW2 and SW3 ($m_1=1$, $m_2=2$ and $m_3=3$), at $T=12.4 \pm 0.1$ hr M2 and LSB are dominant ($m_1=1$ and $m_2=2$), and at

$T=11.6\pm 0.1$ hr exist only the USB ($m_1=3$). With these assignments of m and according to Equation 1, we repeat the estimation of \bar{a} on the grids of date t and altitude h at each of the three periods, resulting in the amplitudes for all six waves, $\bar{a}^\bullet(h, t)$, where \bullet represents LSB, M2, the USB, SW1, SW2, or SW3. The corresponding amplitude $|\bar{a}^\bullet(h, t)|$ is displayed in Figure 3.

3 Results

5 In Figure 3, the decomposition is based on observations from five SMRs between 2012 and 2016, whereas before 2012 only four SMRs are available (Mohe SMR started operation in 2012). The different SMR combinations are designated by the yellow and cyan lines at the bottom of Figure 3f. Using the four SMRs, we also produced the results between 2012 and 2016, which are highly consistent with the results from the five SMRs: the corresponding correlation coefficients are 0.92, 0.96, 0.93, 0.95, 0.98, and 0.95 for the six components, respectively. In Figures 3a-c, the horizontal yellow line around January 2013 shows that
10 the amplitudes are quantitatively consistent with the recent estimation using only the two SMRs at Juliusruh and Mohe: the components $m=1, 2, \text{ and } 3$ maximize at roughly 4, 8, and 8m/s in both Figures 3a-c here and Figure 4 in He et al. (2018b). The correlation and consistency suggest that the decomposition is overly not sensitive to the absence of one SMR between 2009 and 2011.

The temporal variations in Figures 3a-f share some similarities. First, in Figures 3a-c LSB, USB, and M2 are often enhanced
15 noticeably in the month following the vertical magenta dashed lines which indicate the polar vortex weakening (PVW, cf., Zhang and Forbes, 2014) as a reference of SSWs in the current study. Second, as shown in Figures 3d-f, SW1, SW2, and SW3 are characterized by repeating annual patterns, as separated by the calendar new year indicated by the solid white lines. For a statistical study on the SSW responses and the seasonal variations, we average the amplitudes of the six components with respect to the time since the PVW epoch and the calendar new year, following the composite analysis approach (CA, e.g., Chau
20 et al., 2015). CA is also known as superposed epoch analysis, SEA, in geophysics and solar physics (e.g., Chree, 1914). The PVW and calendar results are shown in Figure 4 and discussed in Sections 4 and 5, respectively.

4 Responses to SSWs

As the most radical manifestation of stratosphere-troposphere coupling, SSWs impact the upper atmosphere in broad altitude and latitude ranges (e.g., Goncharenko and Zhang, 2008; Goncharenko et al., 2013). One type of impact is the broadly-reported
25 enhancements of various waves at periods near 12 hr, including M2, SW1, SW3 and LSB and USB (e.g., Chau et al., 2015; Angelats I Coll and Forbes, 2002; Liu et al., 2010, and references therein). Recently, He et al. (2017a) argued that there might not be SW1 and SW3 enhancements during SSWs and instead suggested that the reported enhancements are just misinterpreted signatures of LSB and USB at low-frequency resolution, respectively. These arguments about SW1 and SW3 were supported observationally by two case studies (He et al., 2018a, b, , respectively). Here, we extend this earlier interpretation statistically
30 in Section 4.1 and investigate their year-to-year variability in Section 4.2.

The current section links observationally the secondary waves, LSB and USB, with SSWs through the interaction between SW2 and the 16-day PW. Such a link entails two more associations, one among SW2, the PW and the secondary waves and the other between PW and SSWs. Both associations were established through single radar analysis approaches. Triple co-occurrence and triple coherence among the PW, SW2, and the secondary waves during SSWs were reported in case studies (e.g., He et al., 2017a), and the PW amplifications during SSWs were also reported, e.g., by Pancheva et al. (2008). While the current work investigates only the near-12hr waves using multi-radar analysis approaches, in future work we will investigate the associations using the same approach.

4.1 Multi-year average

The SSW CA results in Figures 4a-f suggest that among the six components, only three, namely, LSB, M2, and USB, exhibit a sharp maximum immediately following PVW, whereas the others, namely, SW1, SW2 and SW3, do not: their intensities are largely decreasing from 40 days before PVW to 50 days after. The enhancements of LSB, USB, and M2 around SSWs are consistent with existing studies, both statistical studies with single radar approaches (e.g., Chau et al., 2015) and case studies (e.g., He et al., 2017a). However, our finding that SW1 and SW3 do not show enhanced intensity during SSWs are at variance with most existing studies (e.g., Liu et al., 2010; Pedatella and Forbes, 2010; Pedatella et al., 2012; Pedatella and Liu, 2013; Wu and Nozawa, 2015). LSB and USB enhancements associated with non-enhancing SW1 and SW3 support the hypothesis that LSB and USB were detected at low frequency-resolution and misinterpreted as SW1 and SW3, respectively (He et al., 2017a). In a case study on SSW 2009, evidence for the SW1 misinterpretation was extracted with an intercontinental-scale dual-SMR network extending along 80°N (He et al., 2018a), while in another case study on SSW 2013, similar evidence was identified for the SW3 misinterpretation with a similar network at 54°N (He et al., 2018b). Here, we report the first multi-year statistical evidence. Besides the responses of LSB and USB, also supporting the hypothesis is the decreasing SW2 decrease at PVW (note that the color is scaled for SW2 amplitude in a range broader than those of others). The declining SW2 feeds the LSB and USB enhancements: SW2 provides 100% and 97% of the energy of the LSB and USB, respectively, according to the *Manley-Rowe* relations detailed in He et al. (2017a).

4.2 Year-to-year variability during SSW

Although LSB, USB, and M2 composite behaviors look similar to each other in Figures 4a-c, their patterns show remarkably different year-to-year variability as shown in Figures 3a-c. To investigate the year-to-year variability, we conduct a CA similar to $\langle |\tilde{a}| \rangle$ displayed in Figures 4a-f, but for the complex amplitude $\langle \tilde{a} \rangle$. Different from the $\langle |\tilde{a}| \rangle$ in Figures 4a-c where all three components maximize during SSW, in $\langle \tilde{a} \rangle$ (not shown here) only M2 maximizes whereas LSB and USB do not. Determined by the phases of both SW2 and the PW at SSWs, $\langle \tilde{a} \rangle$ of LSB and USB exhibit more randomness than that of M2 whose phase is determined only by the M2 phase at SSW. The consistency between $|\langle \tilde{a} \rangle|$ and $\langle |\tilde{a}| \rangle$ of M2 might be attributed either to a potential association between SSW and a particular lunar phase (as suggested by, e.g., Fejer et al., 2010) or simply to the limited sampling number of M2 enhancement events during SSWs (see Figure 3b).

To explore possible relationships between the enhancements of different waves, we search, in Figures 3a-c, the maximum amplitude in a 30d-wide window following each PVW, as a measure of the intensity of the corresponding enhancements. The maxima are marked as magenta plus symbols in Figure 3. A clear association is found between LSB and M2 enhancements. As illustrated in Figure 5, the seven events are clustered mainly into three groups. In the case of other combinations, i.e., USB vs. M2, or LSB vs. USB, we have not found any noticeable relationship. This result suggests that LSB and USB are independent of each other during SSWs. The lack of coupling between the sidebands has been discussed in detail in Section 4.5 in He et al. (2017a).

We further investigate three clusters in Figure 5 according to a classification of associated major SSWs (cf., Seviour et al., 2016; Esler and Matthewman, 2011): vortex-split or displacement marked by black solid and open circles, respectively. Clearly, three clusters circled in the magenta lines in Figure 5 are associated with the SSW classification: (a) the strongest LSB and intermediate M2 occur in vortex-displacement events, (b) intermediate LSB and strongest M2 occur in vortex-split events, and (c) weakest M2 and weakest LSB occur mostly in non-major-SSW events. Here, non-major-SSW events refer to the minor and final SSWs (cf., Butler et al., 2015, 2017; Limpasuvan et al., 2005). The only exception in this classification is the 2015 event.

The association between the SSW classification and M2 strength is consistent with the conclusion drawn from more SSW events using equatorial magnetic field observations (e.g., Siddiqui et al., 2018). Here, our multi-SMR-jointed analysis allows us to separate LSB and M2 components that share the same period. Our results imply that LSB has contaminated previous M2 estimations based on single-site observations, particularly during vortex-displacement SSWs. The association between the vortex-displacement SSW and M2 implies an alternative explanation for the LSB signatures at $T=12.4\text{hr}$ with $m=1$. Although has never been proposed in existing literature, the LSB signature might be, according to the resonant condition, a secondary wave of the nonlinear interaction between stationary PW with zonal wavenumber 1 structure (sPW1) and M2. Although our analysis is *NOT* sufficient to exclude the possibility of the sPW1-M2 interaction, evidence from case studies was reported only supporting the PW-SW2 interaction, including the triple co-occurrence and triple coherency of the three involved waves, and the accompany or occurrence of the USB (e.g., He et al., 2017b). On the contrary, against the sPW1-M2 interaction is the fact that the LSB signature was observed without co-occurrence of significant M2 (e.g., He et al., 2018b). Accordingly, throughout the current work, we explain the LSB signature as a secondary wave of the PW-SW2 interaction.

5 Climatological seasonal variations of the solar tides

In the current section, we change our focus to the seasonal climatology of the identified six waves. Similar to Figures 4a-f showing the SSW CA with respect to PVW, Figures 4g-l display the CA results with respect to the start of the calendar year. Figures 4g-l exhibit similarities with Figures 4a-f, e.g., similar vertical and temporal extensions of the primary peaks. The similarities are not surprising since the time epochs are close to each other: PVWs always occurred in winter near the start of the new year. In comparison with the SSW CA results, in the calendar CA the primary peaks of LSB, USB, and M2 (Figures 4g-i) are slightly smeared out. In contrast, the peaks of the solar tides (SW1, SW2, and SW3 in Figures 4j, 4k, and 4l, respectively) have not been smeared out in the calendar CA, the peaks of SW2 and SW3 are even sharper and stronger. These

results suggest that the temporal variations of LSB, USB, and M2 characterized more by their responses to SSWs than by their seasonal variations, whereas those of the solar tides are characterized more by the seasonal variations.

5.1 Comparison to previous studies

In the amplitude plots shown in Figures 4j-l and 3d-f the vertical variations are characterized by larger amplitudes at higher altitudes. MLT waves are often excited in and propagated from the stratosphere or troposphere. The upward propagating waves amplify exponentially with increasing altitude as the air density decreases, and eventually dissipate. Such a simple vertical structure is associated with the fact that in Figure 3 the temporal variations, enhance and fading, of waves extent typically in a broad altitude range rather than at limited altitude levels. We average vertically the amplitudes shown in Figures 3d-f for an one-dimensional representation, and display the average as a function of the day of year (DoY) in Figure 6a. The averaged components are shown as scatter plot in Figures 6b, 6c, and 6d, against each other, i.e., for SW1 vs SW2, SW3 vs SW2, and SW3 vs SW1, respectively. The most salient feature of the scatters is that SW2 is almost always the dominant component, except in late October when SW3 is comparable to SW2. The scatters are further averaged seasonally displayed as the red, green and blue solid lines, summarizing the main seasonal variations of SW1, SW2, and SW3, respectively. SW2 is characterized by two comparable peaks in September and in December, and steep decreases in September-October and March-April (DoY 250-300 and 0-80), which is consistent with the seasonal variation of the 12.0hr harmonic amplitude (S2) observed from single radar analyses (as used in, e.g., Conte et al., 2018) although the SW1 and SW3 are not negligible in comparison with SW2. SW1 is characterized by a single peak appearing in winter and a minimum in summer, which is largely consistent with thermospheric seasonal variation of SW1 at 50°N according to CHAMP observations (Figure 12, in Oberheide et al., 2011). SW3 is characterized by two peaks in earlier May and October (at DoY 130 and 280), respectively. Similar annual dual peaks of SW3 were observed from SABER measurements (Figure 2.7 in Hartwell, 1994), and also obtained at 50°N at 88km altitude from the three years (2006-2008) of simulated data from the Canadian Middle Atmosphere Model Data Assimilation System (Figures 6b and 10c in Xu et al., 2012). Interestingly, in Figure 6a, d, the relative importance of SW1 and SW3 switches around early April and November (DoY 90 and 310) : in summer SW3 is stronger than SW1, but SW1 is stronger in winter. These seasonal variations might be associated with the climatology of the background mean wind (e.g., Laskar et al., 2016; Conte et al., 2017).

In comparison with some previous studies, the amplitudes in Figure 3 appear to be weaker (e.g., Jacobi, 2012), for at least three potential reasons. First, based on single-site observations, most existing studies did not separate waves with different wavenumbers but had to explain the total oscillations at 12hr or 12.4hr as approximations of SW2 or M2 (e.g., Chau et al., 2015). Second, most existing studies used windowing functions much narrower than ours, resulting in broader passbands and capturing more energy (e.g., Chau et al., 2015; Forbes and Zhang, 2012). Third, some studies present the amplitude of total wind including both zonal and meridional (e.g., Chau et al., 2015; Conte et al., 2017) while here we focus only on the zonal component. For a quantitative comparison, in the next section, we present a comparison with an independent empirical tidal model CTMT.

5.2 A comparison with an empirical model

Figures 7a-f present composite analysis in the same manner as in Figures 4j-l but for the amplitude of complex average, $|\langle \tilde{a} \rangle|$. The similarities between Figures 4j-l and 7a-c, indicate that the phases of solar tides are consistent from year to year.

For an independent quantitative comparison, we present the seasonal variation of the solar tides according to CTMT (Oberheide et al., 2011), in Figures 7g-l. The CTMT results exhibit some consistency with our results, especially on SW2. SW2 in Figure 7h maximizes during August-September and December-January, in between these periods there is a minimum. The vertical gradient is steeper during the August-September maximum than that during December-January. These features are similar to those in Figure 7b. However, in the CTMT results, the maximums or the minimum can hardly be observed (Figure 7h). These morphological discrepancies might arise from the low temporal resolution of the model. The effective resolution is about two months during which the satellite observations used for the model cover the whole local time once. In contrast, the September maximum and the minimum are narrower than two months; therefore they might be smeared out. In the case of phase, our result in Figure 7e also exhibit similarities as the CTMT results in Figure 7k.

SW3 is compared in Figures 7c, 7f, 7i and 7l, from which similarities in both amplitude and phase occur mainly in fall. Although SW3 in the CTMT results also exhibit a second maximum, it occurs during February-March, up to two months before the second annual peak in early May from our results (see Figures 7c and 6a). This difference might be associated with the seasonally uneven sampling of observations used for the model: the satellite takes two months to cover all local time sectors. In the case of SW1, major discrepancies are found in both amplitudes and phases. For instance, the December maximum in our results (Figures 7a) could not be found in the CTMT (Figures 7g).

These qualitative findings are supported quantitatively by Figures 8a, 8b, and 8c where the in-phase and quadrature components of our estimations versus model are shown as scatter plots for SW1, SW2, and SW3, respectively. The highest correlation coefficient is observed in SW2. Since the temporal and vertical resolutions of our results are higher than the CTMT, we smear our results down to the resolutions of CTMT and calculate the correlation coefficients again, yielding correlations coefficients slightly higher than those in Figure 8 by up to 0.01 (not shown). The correlations are overall not high, reflecting mainly the different assumptions used by the two approaches. Our approaches assume the meridian tidal variation among our SMRs are negligible whereas CTMT, as well as any other tidal analyses using single-satellite approaches, assumes the tides are static in the data-binning window. Evaluating these assumptions comparatively entails an independent model with high resolutions in both time and space. Besides, our results might be contaminated by the neglected tides, which is quantified in the next section, whereas the CTMT tidal components might be contaminated by aliasing from waves with similar periods, e.g., the secondary waves and M2.

30 5.3 Bias of our estimation due to the existence of neglected solar tidal components

For estimating the solar tides and as explained in Section 2.2, in our approach we assume that our targeting components, i.e., SW1, SW2, and SW3, are the dominantly components at 12.0hr. This assumption might be too strong given that other

neglected semidiurnal tidal components have also been reported (e.g., Oberheide et al., 2011; He et al., 2011). The current section quantifies the bias due to the existence of other neglected components according to CTMT.

Arrange Equation 1 into two parts, namely, the targeting components with amplitudes $\tilde{\mathbf{a}}_{3 \times 1}^{tar}$ and the neglected components with $\tilde{\mathbf{a}}_{(K-3) \times 1}^{neg}$,

$$5 \quad \left(\tilde{W}^1, \tilde{W}^2, \dots, \tilde{W}^5 \right)' = \tilde{\mathbf{E}}_{5 \times K} \tilde{\mathbf{a}}_{K \times 1} := \tilde{\mathbf{E}}_{5 \times 3}^{tar} \tilde{\mathbf{a}}_{3 \times 1}^{tar} + \tilde{\mathbf{E}}_{5 \times (K-3)}^{neg} \tilde{\mathbf{a}}_{(K-3) \times 1}^{neg} \quad (2)$$

Multiply $(\tilde{\mathbf{E}}_{5 \times 3}^{tar})^{-1} := \left((\tilde{\mathbf{E}}_{5 \times 3}^{tar})^T \tilde{\mathbf{E}}_{5 \times 3}^{tar} \right)^{-1} (\tilde{\mathbf{E}}_{5 \times 3}^{tar})^T$, resulting in,

$$(\tilde{\mathbf{E}}_{5 \times 3}^{tar})^{-1} \left(\tilde{W}^1, \tilde{W}^2, \dots, \tilde{W}^5 \right)' = \tilde{\mathbf{a}}_{3 \times 1}^{tar} + (\tilde{\mathbf{E}}_{5 \times 3}^{tar})^{-1} \tilde{\mathbf{E}}_{5 \times (K-3)}^{neg} \tilde{\mathbf{a}}_{(K-3) \times 1}^{neg} \quad (3)$$

Here, the term on the left side is the estimated amplitude of the targeting components, while the first term on the right is the corresponding real amplitude. Therefore, their difference, i.e., the second term on the right, corresponds to the bias due to the neglected tidal components. According to CTMT, we estimate the bias and display its absolute value in Figure 9. Contributing to the bias are semidiurnal components SE3, SE2, SE1, S0, and SW4. The bias is overall below 2 m/s, except above 90km in summer, which suggests our main conclusions in previous sections are not affected by our assumption that SW1, SW2, and SW3 are the dominant components. Actually, determined by the configuration of the SMRs, the matrices of both $(\tilde{\mathbf{E}}_{5 \times 3}^{tar})^{-1}$ and $(\tilde{\mathbf{E}}_{5 \times 3}^{tar})^{-1} \tilde{\mathbf{E}}_{5 \times (K-3)}^{neg}$ are very well-conditioned (with condition number of 1.6 and 2.4, respectively), suggesting our estimations are not sensitively affected by the errors in both the wavelet spectra and the neglected semidiurnal components.

Our comparison from the previous section has stressed the additional information that our results bring, specifically, those on SW1 and SW3. In future efforts, we plan to add more ground-based observations and try to combine them with satellite-based wind and temperature observations (cf, Zhou et al., 2018), to improve our understanding of mesospheric tides. Although the current work focuses on the near 12 hr waves at midlatitudes, our joint-dataset analysis approach could be extended to other periods, e.g., diurnal or terdiurnal tides.

6 Conclusions

Combining mesospheric zonal wind observations collected by five midlatitude SMRs from three longitudinal sectors, we develop an approach to investigate statistically six waves at periods close to 12 hr, namely, three solar tides (SW1, SW2, and SW3), two sidebands of nonlinear modulation of 16-day wave on SW2 (LSB and USB), and a lunar tide (M2). We first filter the observation from each SMR into three narrow frequency bands through a high-frequency-resolved wavelet analysis. Then, in each of the three bands, wavelet spectra from all SMRs are combined to fit the potential waves. The results suggest that the temporal variations of the waves are characterized by responses to SSWs (enhancements of LSB, USB, and M2, and a decrease in SW2) and climatological seasonal variations of the solar tides. Our main results are:

(1) Contrary to most extensive previous literature, our results suggest that SW1 and SW3 do not statistically enhance during SSWs. The LSB and USB enhancements have been misinterpreted as SW1 and SW3 signatures, respectively. Meanwhile, the

enhancements are associated with a decrease in SW2, which could be explained in terms of the energy exchange through the nonlinear interaction.

(2) Both enhancements of LSB and M2 depend on the SSW classification with respect to the polar vortex split or displacement. M2 enhancement is stronger during vortex-split SSWs than that during the vortex-displacement, whereas LSB is the other way around. Overall, M2 is stronger than LSB, except during the vortex-displacement SSW when they are comparable, 5 implicating that LSB might contaminate the existing M2 estimations based on single-site observations.

(3) The seasonal variations of solar tides are in reasonable agreements with existing observational studies: SW2 is the dominant component which maximizes around September and December followed by two minimums; SW1 maximizes in winter, and SW3 maximizes in fall and spring. In October when SW3 is at its annual maximum and SW2 is at a minimum, 10 when their strengths are comparable to each other. These results suggest that the 12.0hr harmonic amplitude from single radar analyses is dominated by SW2 for most of seasons except in October.

Code availability. TEXT

Data availability. TEXT

Our main results, namely, the complex amplitudes of the six waves as a function of time and altitude, are shared at 15 <ftp://ftp.iap-kborn.de/data-in-publications/HeACP2018/>.

Code and data availability. TEXT

Sample availability. TEXT

Author contributions. TEXT

Conceptualization, M.H.and J.L.C.; Methodology, M.H.; Software, M.H.; Writing– original draft, M.H.; Writing– review&editing, 20 M.H.and J.L.C.; Funding acquisition, J.L.C..

Competing interests. TEXT

The authors declare having no conflict of interest.

Disclaimer. TEXT

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References

- Ahlquist, J. E.: Normal-mode global Rossby waves.pdf, [http://journals.ametsoc.org/doi/abs/10.1175/1520-0469\(1982\)039{%}3C0193:NMGRWT{%}3E2.0.CO;2](http://journals.ametsoc.org/doi/abs/10.1175/1520-0469(1982)039{%}3C0193:NMGRWT{%}3E2.0.CO;2), 1982.
- Angelats I Coll, M. and Forbes, J. M.: Nonlinear interactions in the upper atmosphere: The $s = 1$ and $s = 3$ nonmigrating semidiurnal tides, *J. Geophys. Res. Sp. Phys.*, 107, 1–18, <https://doi.org/10.1029/2001JA900179>, 2002.
- Azeem, S. M., Killeen, T. L., Johnson, R. M., Wu, Q., and Gell, D. A.: Space-time analysis of TIMED Doppler Interferometer (TIDI) measurements, *Geophys. Res. Lett.*, 27, 3297–3300, <https://doi.org/10.1029/1999GL011289>, <https://doi.org/10.1029/1999gl011289>, 2000.
- Baumgaertner, A. J., Jarvis, M. J., McDonald, A. J., and Fraser, G. J.: Observations of the wavenumber 1 and 2 components of the semi-diurnal tide over Antarctica, *J. Atmos. Solar-Terrestrial Phys.*, 68, 1195–1214, <https://doi.org/10.1016/j.jastp.2006.03.001>, <http://www.sciencedirect.com/science/article/pii/S1364682606000691>[https://www.sciencedirect.com/science/article/pii/S1364682606000691?](https://www.sciencedirect.com/science/article/pii/S1364682606000691?via{%}3Dihub) via{%}3Dihub, 2006.
- Butler, A. H., Seidel, D. J., Hardiman, S. C., Butchart, N., Birner, T., and Match, A.: Defining Sudden Stratospheric Warmings, *Bulletin of the American Meteorological Society*, 96, 1913–1928, <https://doi.org/10.1175/BAMS-D-13-00173.1>, <https://doi.org/10.1175/BAMS-D-13-00173.1>, 2015.
- 15 Butler, A. H., Sjöberg, J. P., Seidel, D. J., and Rosenlof, K. H.: A sudden stratospheric warming compendium, *Earth System Science Data*, 9, 63–76, <https://doi.org/10.5194/essd-9-63-2017>, <https://www.earth-syst-sci-data.net/9/63/2017/>, 2017.
- Chapman, S. and Lindzen, R. S.: Atmospheric Tides: Thermal and Gravitational: Nomenclature, Notation and New Results, [https://doi.org/10.1175/1520-0469\(1970\)027<0707:ATTAGN>2.0.CO;2](https://doi.org/10.1175/1520-0469(1970)027<0707:ATTAGN>2.0.CO;2), 1970.
- Chau, J. L., Hoffmann, P., Pedatella, N. M., Matthias, V., and Stober, G.: Upper mesospheric lunar tides over middle and high latitudes during sudden stratospheric warming events, *J. Geophys. Res. Sp. Phys.*, 120, 3084–3096, <https://doi.org/10.1002/2015JA020998>, <https://agupubs.onlinelibrary.wiley.com/doi/abs/10.1002/2015JA020998>, 2015.
- 20 Chree, C.: Some Phenomena of Sunspots and of Terrestrial Magnetism . Part II, *Philos. Trans. R. Soc. London. Ser. A, Contain. Pap. a Math. Phys. Character*, 213, 245–277, <https://doi.org/10.1098/rsta.1913.0003>, <https://doi.org/10.1098/rsta.1913.0003>, 1914.
- Conte, J. F., Chau, J. L., Stober, G., Pedatella, N., Maute, A., Hoffmann, P., Janches, D., Fritts, D., and Murphy, D. J.: Climatology of semidiurnal lunar and solar tides at middle and high latitudes: Interhemispheric comparison, *J. Geophys. Res. Sp. Phys.*, 122, 7750–7760, <https://doi.org/10.1002/2017JA024396>, <http://dx.doi.org/10.1002/2017JA024396><http://onlinelibrary.wiley.com/store/10.1002/2017JA024396/asset/jgra53687.pdf?v=1{%}&t=jansn8up{%}&s=72eb34b0c842fc65e21955e9b8b7bfd8028227ee>, 2017.
- Conte, J. F., Chau, J. L., Laskar, F. I., Stober, G., Schmidt, H., and Brown, P.: Semidiurnal solar tide differences between fall and spring transition times in the Northern Hemisphere, *Ann. Geophys.*, 36, 999–1008, <https://doi.org/10.5194/angeo-36-999-2018>, <https://www.ann-geophys.net/36/999/2018/>, 2018.
- 30 Day, K. A. and Mitchell, N. J.: The 16-day wave in the Arctic and Antarctic mesosphere and lower thermosphere, *Atmos. Chem. Phys.*, 10, 1461–1472, <https://doi.org/10.5194/acp-10-1461-2010>, 2010.
- Esler, J. G. and Matthewman, N. J.: Stratospheric Sudden Warmings as Self-Tuning Resonances. Part II: Vortex Displacement Events, *J. Atmos. Sci.*, 68, 2505–2523, <https://doi.org/10.1175/JAS-D-11-08.1>, <http://journals.ametsoc.org/doi/abs/10.1175/JAS-D-11-08.1>, 2011.
- 35 Fejer, B. G., Olson, M. E., Chau, J. L., Stolle, C., Luehr, H., Goncharenko, L. P., Yumoto, K., and Nagatsuma, T.: Lunar-dependent equatorial ionospheric electrodynamic effects during sudden stratospheric warmings, *J. Geophys. Res. Sp. Phys.*, 115, 1–9, <https://doi.org/10.1029/2010JA015273>, 2010.

- Fejer, B. G., Tracy, B. D., Olson, M. E., and Chau, J. L.: Enhanced lunar semidiurnal equatorial vertical plasma drifts during sudden stratospheric warmings, *Geophys. Res. Lett.*, 38, 7271, <https://doi.org/10.1029/2011GL049788>, 2011.
- Forbes, J. M. and Zhang, X.: Lunar tide amplification during the January 2009 stratosphere warming event: Observations and theory, *J. Geophys. Res. Sp. Phys.*, 117, 1–13, <https://doi.org/10.1029/2012JA017963>, 2012.
- 5 Goncharenko, L. and Zhang, S. R.: Ionospheric signatures of sudden stratospheric warming: Ion temperature at middle latitude, *Geophys. Res. Lett.*, 35, 4–7, <https://doi.org/10.1029/2008GL035684>, 2008.
- Goncharenko, L., Chau, J. L., Condor, P., Coster, A., and Benkevitch, L.: Ionospheric effects of sudden stratospheric warming during moderate-to-high solar activity: Case study of January 2013, *Geophys. Res. Lett.*, 40, 4982–4986, <https://doi.org/10.1002/grl.50980>, 2013.
- Grossmann, A., Kronland-Martinet, R., and Morlet, J.: Reading and Understanding Continuous Wavelet Transforms, in: *Wavelets*, edited by 10 Combes, J.-M., Grossmann, A., and Tchamitchian, P., pp. 2–20, Springer Berlin Heidelberg, Berlin, Heidelberg, 1990.
- Hartwell, F. P.: Wiring methods for patient care areas, vol. 93, <https://doi.org/10.1007/978-94-007-0326-1>, <http://www.springerlink.com/index/10.1007/978-94-007-0326-1> <http://link.springer.com/10.1007/978-94-007-0326-1>, 1994.
- He, M., Liu, L., Wan, W., and Wei, Y.: Strong evidence for couplings between the ionospheric wave-4 structure and atmospheric tides, 15 *Geophys. Res. Lett.*, 38, <https://doi.org/10.1029/2011GL047855>, 2011.
- He, M., Chau, J. L., Stober, G., Hall, C. M., Tsutsumi, M., and Hoffmann, P.: Application of Manley-Rowe Relation in Analyzing Nonlinear Interactions Between Planetary Waves and the Solar Semidiurnal Tide During 2009 Sudden Stratospheric Warming Event, *J. Geophys. Res. Sp. Phys.*, 122, 10, 783–10,795, <https://doi.org/10.1002/2017JA024630>, 2017a.
- He, M., Vogt, J., Heyner, D., and Zhong, J.: Solar wind controls on Mercury’s magnetospheric cusp, *J. Geophys. Res. Sp. Phys.*, 122, 20 6150–6164, <https://doi.org/10.1002/2016JA023687>, 2017b.
- He, M., Chau, J. L., Hall, C., Tsutsumi, M., Meek, C., and Hoffmann, P.: The 16-day planetary wave triggers the SW1-tidal-like signatures during 2009 sudden stratospheric warming, *Geophys. Res. Lett.*, <https://doi.org/10.1029/2018GL079798>, 2018a.
- He, M., Chau, J. L., Stober, G., Li, G., Ning, B., and Hoffmann, P.: Relations Between Semidiurnal Tidal Variants Through Diagnosing the Zonal Wavenumber Using a Phase Differencing Technique Based on Two Ground-Based Detectors, *J. Geophys. Res. Atmos.*, 25 123, 4015–4026, <https://doi.org/10.1002/2018JD028400>, <https://agupubs.onlinelibrary.wiley.com/doi/pdf/10.1002/2018JD028400><http://doi.wiley.com/10.1002/2018JD028400>, 2018b.
- Hocking, W., Fuller, B., and Vandeeper, B.: Real-time determination of meteor-related parameters utilizing modern digital technology, *Journal of Atmospheric and Solar-Terrestrial Physics*, 63, 155 – 169, [https://doi.org/https://doi.org/10.1016/S1364-6826\(00\)00138-3](https://doi.org/https://doi.org/10.1016/S1364-6826(00)00138-3), <http://www.sciencedirect.com/science/article/pii/S1364682600001383>, radar applications for atmosphere and ionosphere research - PIERS 30 1999, 2001.
- Jacobi, C.: 6 year mean prevailing winds and tides measured by VHF meteor radar over Collm (51.3°N, 13.0°E), *Journal of Atmospheric and Solar-Terrestrial Physics*, 78-79, 8 – 18, <https://doi.org/https://doi.org/10.1016/j.jastp.2011.04.010>, <http://www.sciencedirect.com/science/article/pii/S1364682611001210>, structure and Dynamics of Mesosphere and Lower Thermosphere, 2012.
- Jones, J., Brown, P., Ellis, K., Webster, A., Campbell-Brown, M., Krzeminski, Z., and Weryk, R.: The Canadian Meteor Orbit Radar: system 35 overview and preliminary results, *Planetary and Space Science*, 53, 413 – 421, <https://doi.org/https://doi.org/10.1016/j.pss.2004.11.002>, <http://www.sciencedirect.com/science/article/pii/S0032063304002302>, 2005.

- Laskar, F. I., Chau, J. L., Stober, G., Hoffmann, P., Hall, C. M., and Tsutsumi, M.: Quasi-biennial oscillation modulation of the middle- and high-latitude mesospheric semidiurnal tides during August–September, *J. Geophys. Res. A Sp. Phys.*, 121, 4869–4879, <https://doi.org/10.1002/2015JA022065>, 2016.
- Limpasuvan, V., Hartmann, D. L., Thompson, D. W., Jeev, K., and Yung, Y. L.: Stratosphere-troposphere evolution during polar vortex intensification, *J. Geophys. Res. Atmos.*, 110, 1–15, <https://doi.org/10.1029/2005JD006302>, 2005.
- Liu, H. L., Wang, W., Richmond, A. D., and Roble, R. G.: Ionospheric variability due to planetary waves and tides for solar minimum conditions, *J. Geophys. Res. Sp. Phys.*, 115, <https://doi.org/10.1029/2009JA015188>, <https://agupubs.onlinelibrary.wiley.com/doi/abs/10.1029/2009JA015188>, 2010.
- Liu, L., Liu, H., Chen, Y., Le, H., Sun, Y.-Y., Ning, B., Hu, L., and Wan, W.: Variations of the meteor echo heights at Beijing and Mohe, China, *Journal of Geophysical Research: Space Physics*, 122, 1117–1127, <https://doi.org/10.1002/2016JA023448>, <https://agupubs.onlinelibrary.wiley.com/doi/abs/10.1002/2016JA023448>, 2016.
- Liu, L., Liu, H., Le, H., Chen, Y., Sun, Y. Y., Ning, B., Hu, L., Wan, W., Li, N., and Xiong, J.: Mesospheric temperatures estimated from the meteor radar observations at Mohe, China, *J. Geophys. Res. Sp. Phys.*, 122, 2249–2259, <https://doi.org/10.1002/2016JA023776>, 2017.
- Longuet-Higgins, M. S.: The Eigenfunctions of Laplace’s Tidal Equations over a Sphere, *Philos. Trans. R. Soc. A Math. Phys. Eng. Sci.*, 262, 511–607, <https://doi.org/10.1098/rsta.1968.0003>, <http://rsta.royalsocietypublishing.org/cgi/doi/10.1098/rsta.1968.0003>, 1968.
- Luo, Y., Manson, A. H., Meek, C. E., Meyer, C. K., Burrage, M. D., Fritts, D. C., Hall, C. M., Hocking, W. K., MacDougall, J., Riggan, D. M., and Vincent, R. A.: The 16-day planetary waves: multi-MF radar observations from the arctic to equator and comparisons with the HRDI measurements and the GSWM modelling results, *Ann. Geophys.*, 20, 691–709, <https://doi.org/10.5194/angeo-20-691-2002>, <http://www.ann-geophys.net/20/691/2002/>, 2002.
- Madden, R. A.: Large-scale, free Rossby waves in the atmosphere - An update, *Tellus, Ser. A Dyn. Meteorol. Oceanogr.*, 59, 571–590, <https://doi.org/10.1111/j.1600-0870.2007.00257.x>, <http://dx.doi.org/10.1111/j.1600-0870.2007.00257.xhttps://www.tandfonline.com/doi/pdf/10.1111/j.1600-0870.2007.00257.x?needAccess=true>, 2007.
- Manson, A. H., Meek, C. E., Chshyolkova, T., Xu, X., Aso, T., Drummond, J. R., Hall, C. M., Hocking, W. K., Jacobi, C., Tsutsumi, M., and Ward, W. E.: Arctic tidal characteristics at Eureka (80° N, 86° W) and Svalbard (78° N, 16° E) for 2006/07: Seasonal and longitudinal variations, migrating and non-migrating tides, *Ann. Geophys.*, 27, 1153–1173, <https://doi.org/10.5194/angeo-27-1153-2009>, 2009.
- Murphy, D. J.: Variations in the phase of the semidiurnal tide over Davis, Antarctica, *J. Atmos. Solar-Terrestrial Phys.*, 64, 1069–1081, [https://doi.org/10.1016/S1364-6826\(02\)00058-5](https://doi.org/10.1016/S1364-6826(02)00058-5), 2002.
- Murphy, D. J.: Observations of a nonmigrating component of the semidiurnal tide over Antarctica, *J. Geophys. Res.*, 108, 4241, <https://doi.org/10.1029/2002JD003077>, <http://doi.wiley.com/10.1029/2002JD003077>, 2003.
- Murphy, D. J., Forbes, J. M., Walterscheid, R. L., Hagan, M. E., Avery, S. K., Aso, T., Fraser, G. J., Fritts, D. C., Jarvis, M. J., McDonald, A. J., Riggan, D. M., Tsutsumi, M., and Vincent, R. A.: A climatology of tides in the antarctic mesosphere and lower thermosphere, *J. Geophys. Res. Atmos.*, 111, 1–17, <https://doi.org/10.1029/2005JD006803>, 2006.
- Murphy, D. J., Aso, T., Fritts, D. C., Hibbins, R. E., McDonald, A. J., Riggan, D. M., Tsutsumi, M., and Vincent, R. A.: Source regions for antarctic MLT non-migrating semidiurnal tides, *Geophys. Res. Lett.*, 36, 1–5, <https://doi.org/10.1029/2008GL037064>, 2009.
- Oberheide, J., Hagan, M. E., and Roble, R. G.: Tidal signatures and aliasing in temperature data from slowly precessing satellites, *Journal of Geophysical Research: Space Physics*, 108, <https://doi.org/10.1029/2002JA009585>, <https://agupubs.onlinelibrary.wiley.com/doi/abs/10.1029/2002JA009585>, 2002.

- Oberheide, J., Forbes, J. M., Zhang, X., and Bruinsma, S. L.: Climatology of upward propagating diurnal and semidiurnal tides in the thermosphere, *J. Geophys. Res. Sp. Phys.*, 116, A11 306, <https://doi.org/10.1029/2011JA016784>, <http://dx.doi.org/10.1029/2011JA016784http://www.agu.org/journals/ja/ja1111/2011JA016784/2011JA016784.pdf>, 2011.
- Pancheva, D. and Mukhtarov, P.: Global response of the ionosphere to atmospheric tides forced from below: Recent progress based on satellite measurements : *Espace of the ionosphere*, vol. 168, <https://doi.org/10.1007/s11214-011-9837-1>, 2012.
- Pancheva, D., Mukhtarov, P., Mitchell, N. J., Merzlyakov, E., Smith, A. K., Andonov, B., Singer, W., Hocking, W., Meek, C., Manson, A., and Murayama, Y.: Planetary waves in coupling the stratosphere and mesosphere during the major stratospheric warming in 2003/2004, *J. Geophys. Res. Atmos.*, 113, 1–22, <https://doi.org/10.1029/2007JD009011>, 2008.
- Paschmann, G. and Daly, P. W.: *Analysis methods for multi-spacecraft data*, Published for the International Space Science Institute by ESA Publications Division, 1998.
- Pedatella, N. M. and Forbes, J. M.: Evidence for stratosphere sudden warming-ionosphere coupling due to vertically propagating tides, *Geophys. Res. Lett.*, 37, n/a—n/a, <https://doi.org/10.1029/2010GL043560>, <http://doi.wiley.com/10.1029/2010GL043560>, 2010.
- Pedatella, N. M. and Liu, H. L.: The influence of atmospheric tide and planetary wave variability during sudden stratosphere warmings on the low latitude ionosphere, *J. Geophys. Res. Sp. Phys.*, 118, 5333–5347, <https://doi.org/10.1002/jgra.50492>, <http://doi.wiley.com/10.1002/jgra.50492>, 2013.
- Pedatella, N. M., Liu, H. L., Richmond, A. D., Maute, A., and Fang, T. W.: Simulations of solar and lunar tidal variability in the mesosphere and lower thermosphere during sudden stratosphere warmings and their influence on the low-latitude ionosphere, *J. Geophys. Res. Sp. Phys.*, 117, n/a—n/a, <https://doi.org/10.1029/2012JA017858>, <http://doi.wiley.com/10.1029/2012JA017858>, 2012.
- Salby, M. L.: *Sampling Theory for Asynoptic Satellite Observations. Part I: Space-Time Spectra, Resolution, and Aliasing*, *J. Atmos. Sci.*, 39, 2577–2600, [https://doi.org/10.1175/1520-0469\(1982\)039<2577:STFASO>2.0.CO;2](https://doi.org/10.1175/1520-0469(1982)039<2577:STFASO>2.0.CO;2), [https://doi.org/10.1175/1520-0469\(1982\)039%3C2577:STFASO%3E2.0.CO;2](https://doi.org/10.1175/1520-0469(1982)039%3C2577:STFASO%3E2.0.CO;2), 1982a.
- Salby, M. L.: *Sampling Theory for Asynoptic Satellite Observations. Part I: Space-Time Spectra, Resolution, and Aliasing*, *J. Atmos. Sci.*, 39, 2577–2600, [https://doi.org/10.1175/1520-0469\(1982\)039<2577:STFASO>2.0.CO;2](https://doi.org/10.1175/1520-0469(1982)039<2577:STFASO>2.0.CO;2), [https://doi.org/10.1175/1520-0469\(1982\)039%3C2577:STFASO%3E2.0.CO;2](https://doi.org/10.1175/1520-0469(1982)039%3C2577:STFASO%3E2.0.CO;2), 1982b.
- Salby, M. L.: Transient disturbances in the stratosphere: implications for theory and observing systems, *J. Atmos. Terr. Phys.*, 46, 1009–1047, [https://doi.org/10.1016/0021-9169\(84\)90007-2](https://doi.org/10.1016/0021-9169(84)90007-2), <http://linkinghub.elsevier.com/retrieve/pii/0021916984900072>, 1984.
- Seviour, W. J., Gray, L. J., and Mitchell, D. M.: Stratospheric polar vortex splits and displacements in the high-top CMIP5 climate models, *J. Geophys. Res.*, 121, 1400–1413, <https://doi.org/10.1002/2015JD024178>, <http://doi.wiley.com/10.1002/2015JD024178>, 2016.
- Siddiqui, T. A., Yamazaki, Y., Stolle, C., Lühr, H., Matzka, J., Maute, A., and Pedatella, N.: Dependence of Lunar Tide of the Equatorial Electrojet on the Wintertime Polar Vortex, Solar Flux, and QBO, *Geophys. Res. Lett.*, 45, 3801–3810, <https://doi.org/10.1029/2018GL077510>, <https://agupubs.onlinelibrary.wiley.com/doi/abs/10.1029/2018GL077510>, 2018.
- Singer, W., Hoffmann, P., Kishore Kumar, G., Mitchell, N. J., and Matthias, V.: *Atmospheric Coupling by Gravity Waves: Climatology of Gravity Wave Activity, Mesospheric Turbulence and Their Relations to Solar Activity*, pp. 409–427, Springer Netherlands, Dordrecht, https://doi.org/10.1007/978-94-007-4348-9_22, https://doi.org/10.1007/978-94-007-4348-9_22, 2013.
- Stening, R. J.: Lunar tide in the equatorial electrojet in relation to stratospheric warmings, *J. Geophys. Res. Sp. Phys.*, 116, n/a—n/a, <https://doi.org/10.1029/2011JA017047>, <http://doi.wiley.com/10.1029/2011JA017047>, 2011.

- Stober, G., Jacobi, C., Matthias, V., Hoffmann, P., and Gerding, M.: Neutral air density variations during strong planetary wave activity in the mesopause region derived from meteor radar observations, *Journal of Atmospheric and Solar-Terrestrial Physics*, 74, 55 – 63, <https://doi.org/https://doi.org/10.1016/j.jastp.2011.10.007>, <http://www.sciencedirect.com/science/article/pii/S136468261100280X>, 2012.
- Wu, Q. and Nozawa, S.: Mesospheric and thermospheric observations of the January 2010 stratospheric warming event, *J. Atmos. Solar-Terrestrial Phys.*, 123, 22–38, <https://doi.org/10.1016/j.jastp.2014.11.006>, <http://www.sciencedirect.com/science/article/pii/S1364682614002703>, 2015.
- Xu, X., Manson, A. H., Meek, C. E., Riggin, D. M., Jacobi, C., and Drummond, J. R.: Mesospheric wind diurnal tides within the Canadian Middle Atmosphere Model Data Assimilation System, *J. Atmos. Solar-Terrestrial Phys.*, 74, 24–43, <https://doi.org/10.1016/j.jastp.2011.09.003>, 2012.
- 10 Yu, Y., Wan, W., Ning, B., Liu, L., Wang, Z., Hu, L., and Ren, Z.: Tidal wind mapping from observations of a meteor radar chain in December 2011, *J. Geophys. Res. Sp. Phys.*, 118, 2321–2332, <https://doi.org/10.1029/2012JA017976>, <https://agupubs.onlinelibrary.wiley.com/doi/abs/10.1029/2012JA017976>, 2013.
- Yu, Y., Wan, W., Ren, Z., Xiong, B., Zhang, Y., Hu, L., Ning, B., and Liu, L.: Seasonal variations of MLT tides revealed by a meteor radar chain based on Hough mode decomposition, *J. Geophys. Res. A Sp. Phys.*, 120, 7030–7048, <https://doi.org/10.1002/2015JA021276>, 2015.
- 15 Zhang, J. T., Forbes, J. M., Zhang, C. H., Doornbos, E., and Bruinsma, S. L.: Lunar tide contribution to thermosphere weather, *Sp. Weather*, 12, 538–551, <https://doi.org/10.1002/2014SW001079>, <http://doi.wiley.com/10.1002/2014SW001079>, 2014.
- Zhang, X. and Forbes, J. M.: Lunar tide in the thermosphere and weakening of the northern polar vortex, *Geophys. Res. Lett.*, 41, 8201–8207, <https://doi.org/10.1002/2014GL062103>, <http://doi.wiley.com/10.1002/2014GL062103>, 2014.
- Zhou, X., Wan, W., Yu, Y., Ning, B., Hu, L., and Yue, X.: New Approach to Estimate Tidal Climatology From Ground- and Space-Based
20 Observations, *J. Geophys. Res. Sp. Phys.*, 123, 5087–5101, <https://doi.org/10.1029/2017JA024967>, 2018.

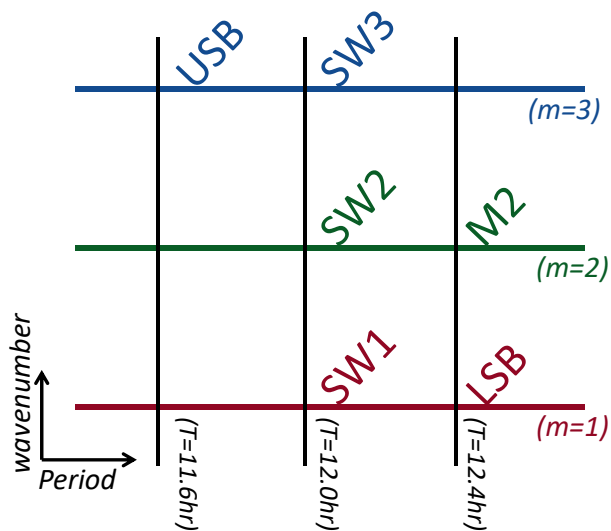


Figure 1. Distribution of near-12hr waves in the frequency and zonal wavenumber space (adapted from He et al., 2018a). In the current study, the colors red, green, and blue represent waves with zonal wavenumber $m=1, 2,$ and $3,$ respectively.

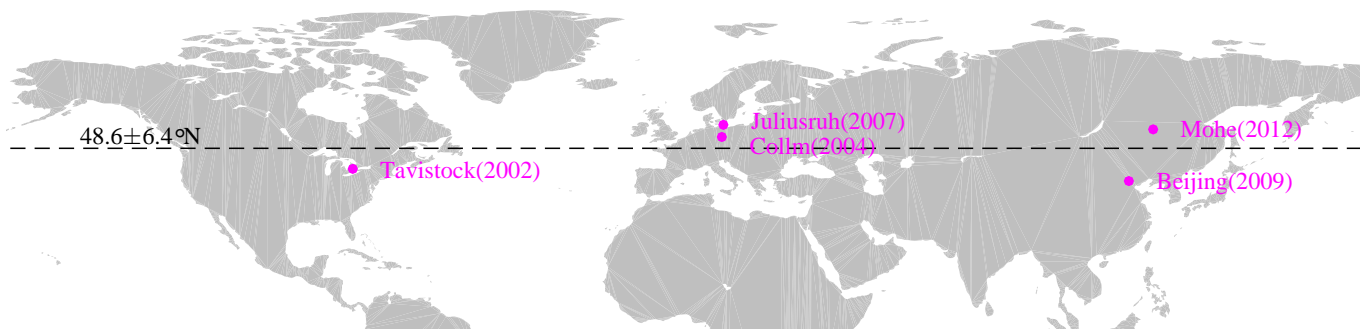


Figure 2. Distribution of five SMRs used in the current study. The numbers following the location names present the earliest available years of the corresponding observations.

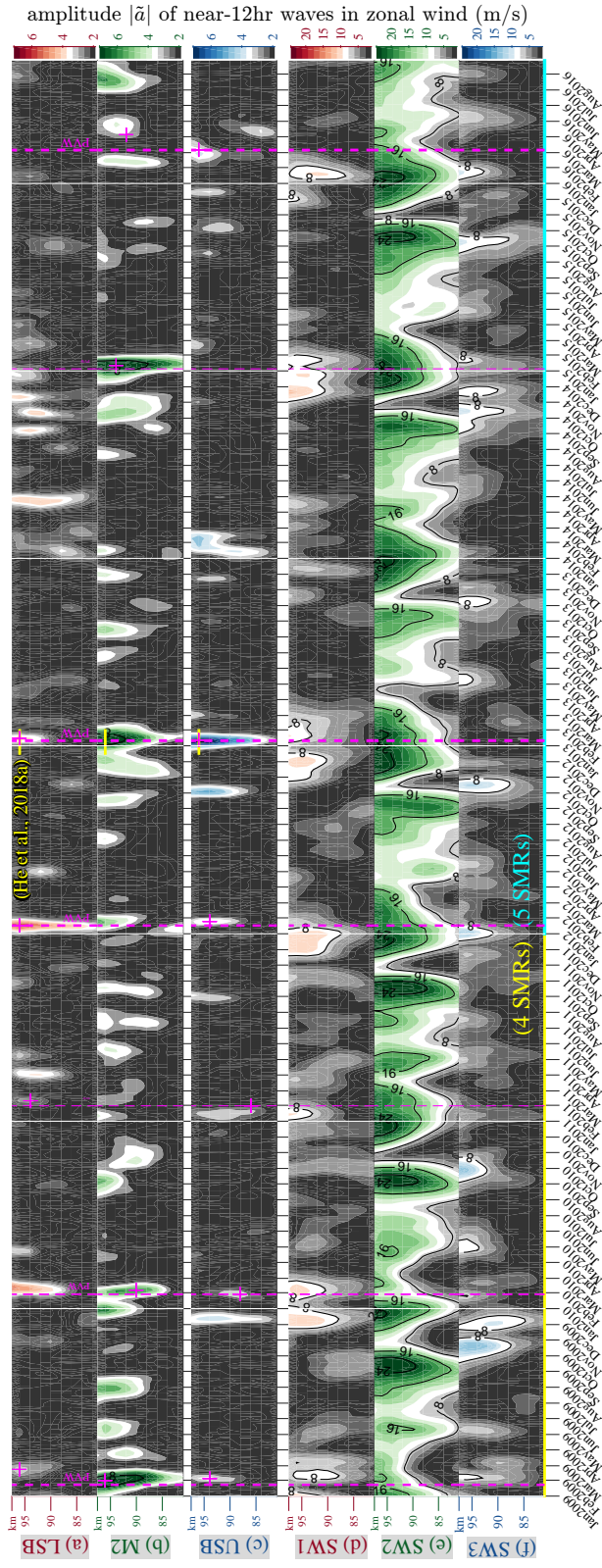


Figure 3. (a) The amplitude of the lower sideband (LSB) of the nonlinear modulation of the 16-day wave on the semidiurnal tide SW2 as a function of time and altitude. (b-f) The same plots as (a) but for the lunar tide M2, the upper sideband (USB), and the solar tides SW1, SW2, and SW3, respectively. (g) The altitude averages of Panels a, b, and c (LSB, M2, and USB), and (h) those of Panels a, b, and c (SW1, SW2, and SW3). In each Panel, the solid white vertical lines display the first day of each year, and the dashed magenta lines display PVWs. In (f), the cyan line on the bottom illustrates the interval from 2012 to 2016 in which all decomposition is based on five SMRs whereas the yellow line represents that MSR observations are not available at Mohe. In (a-c), the magenta plus symbols illustrate the maximum amplitude in each 30-day window following each PVW.

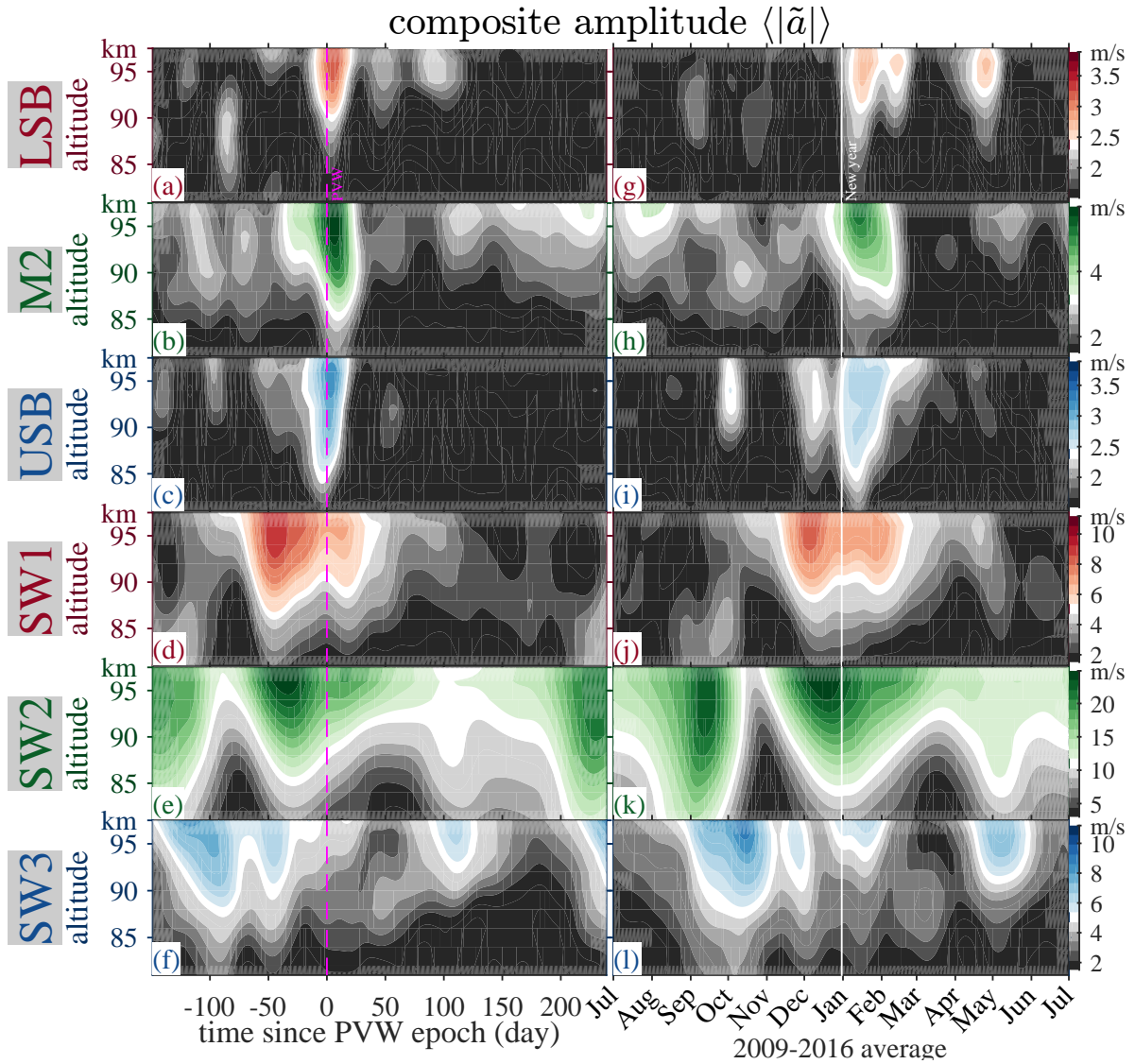


Figure 4. (a) Composite analysis of LSB from Figure 3a with respect to the occurrence of PVWs, namely, the magenta dashed lines in Figure 3a. (b-f) Same plots as (a) but for M2, USB, SW1, SW2, and SW3 from Figures 3b-f, respectively. (g-l) Same plots as (a-f) but with respect to the start of the calendar years, namely, the white lines in Figure 3.

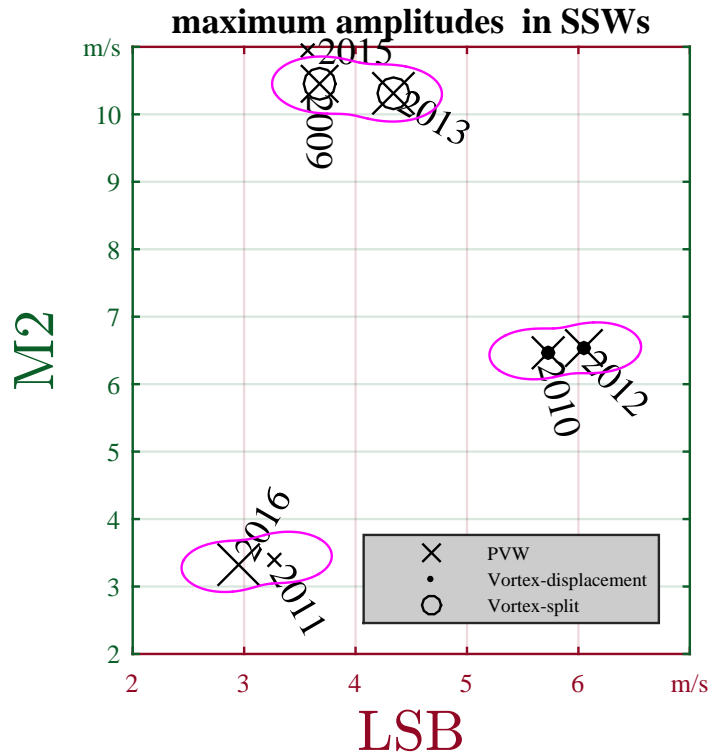


Figure 5. Scatter plots of the maximum amplitudes of LSB and M2 during SSWs, read from Figure 3. The size of the cross is proportional to the PVW strength defined by Zhang et al. (2014). The magenta circles cluster the PVWs into three main groups according to the SSW classification according to the associated polar vortex split or displacement.

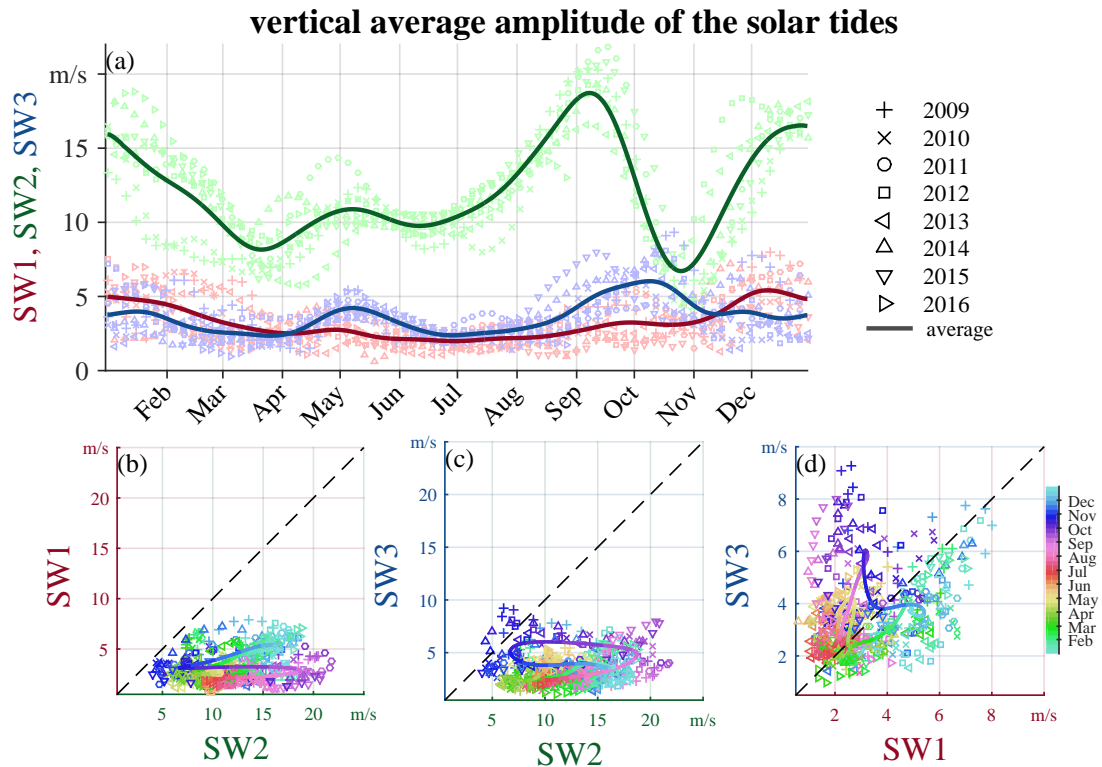


Figure 6. (a) Vertical average amplitude of SW1, SW2, and SW3 scattered as a function of date. (b, c, d) Scatter plot between the vertical average amplitudes of SW2 vs SW1, SW2 vs SW3, and SW1 vs SW3, respectively. In each panel, each point corresponds to a five day interval in Figure 3; and the solid colored line represents the multi-year average.

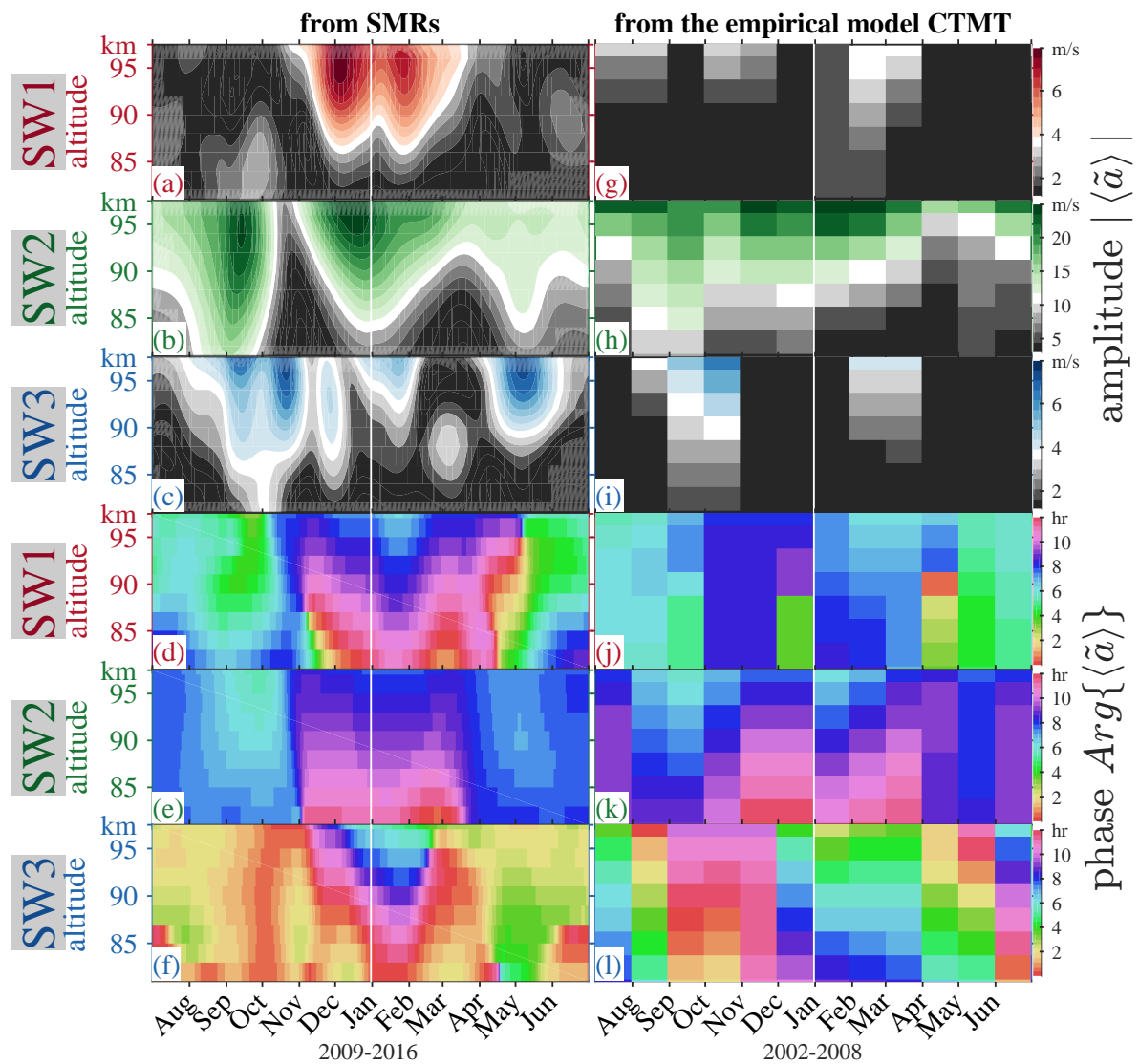


Figure 7. (a, b, c) Same plots as Figures 4j, 4k, 4l but for complex amplitude $|\langle \tilde{a} \rangle|$ of SW1, SW2, and SW3, with their phase shown in (d, e, f), respectively. (g-l) Similar plots as (a-f) but according to the climatological tidal model of the thermosphere (CTMT) derived from SABER and TIDI observations (Oberheide et al., 2011).

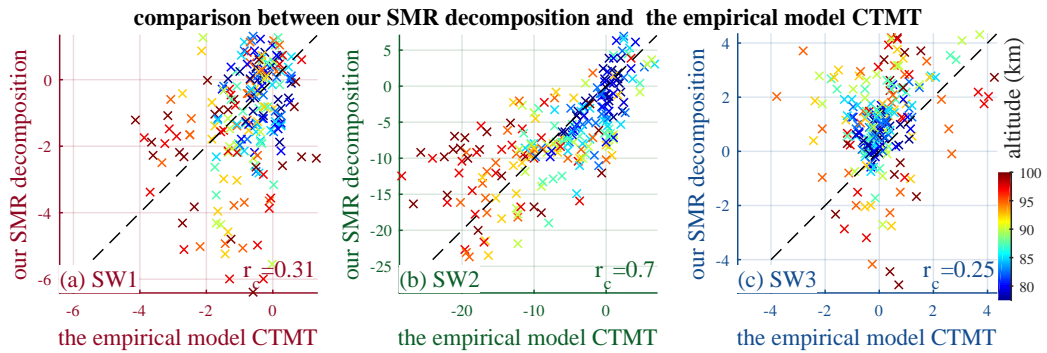


Figure 8. Scatter plots of SW1, SW2 and SW3 shown in Figures 4a-f, against those shown in Figures 4g-l. Each cross in Panel (a,b,c) represents the real or imaginary part of one pixel in Figures 7g, 7h, 7i, respectively.

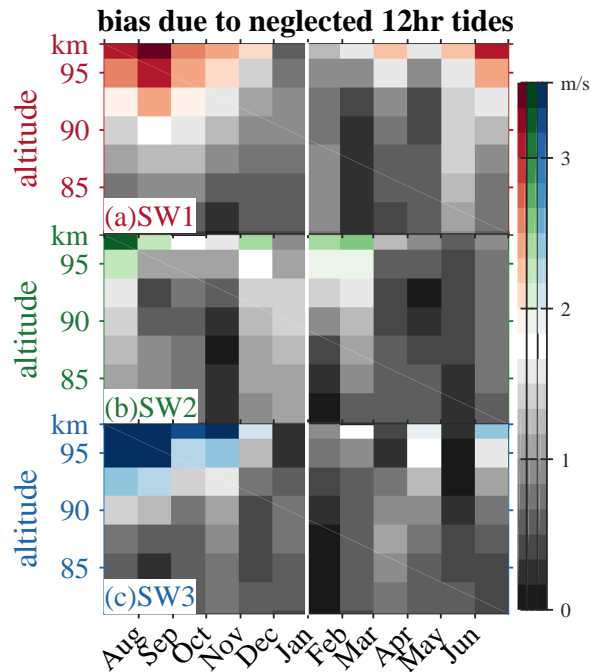


Figure 9. The bias of our tidal estimation due to the existence of neglected 12hr tides (including SE3, SE2, SE1, S0, and SW4) predicted by CTMT (Oberheide et al., 2011).