# Heini Wernli ETH Zurich, Institute for Atmosphere and Climate Science Universitaetstrasse 16 8092 Zurich Switzerland

Dear Dr. Wernli:

Please find enclosed our revised manuscript, "Vertical profile observations of water vapor deuterium excess in the lower troposphere" for publication as a research article in Atmospheric Chemistry and Physics. We have diligently addressed each comment given by the two reviewers. The main changes include a reorganization of separate results and discussions sections into combined sections for each case study flight day. The reviewers also both questioned our application of isotopic model equations from Stewart and Worden to describe the partial evaporation of liquid droplets and subsequent complete evaporation in a different atmospheric layer. We agree with the reviewers that the assumptions of these models do not hold in our observations. We have replaced this discussion with an improved analysis of the relevant timescales for evaporation within cloud layers.

The changes we have made to the manuscript are consistent with the reviewers' comments, and have greatly improved the paper. The submission includes our responses to each of the reviewers' comments, the manuscript and Supporting Information (SI) with changes indicated in "track changes mode", and the final revised manuscript, SI, and figures. We plan to comply with the journal's data policy by providing a DOI once the manuscript is accepted for publication. In the meantime, the data is accessible at <u>https://vapor-isotope.yale.edu/</u>. We hope you agree that we have completely addressed all the reviewers' concerns, and that the paper is now in publishable form. Thank you for your consideration.

Sincerely,

Olivia Salmon and Lisa Welp Purdue University

# Authors' Responses to Reviewer 1

- Author responses are indicated in blue font.
- Locations of our edits in the "strike-through" and final versions of the manuscript are provided with the following convention: ([strike-through version] pg. #, ln # / [final version] pg. #, ln #).

# 1 Major comments

1. I am not convinced by the way the Rayleigh model is used as an explanation in flights RAY and STC. Even though a Rayleigh model can match much your data, it does not necessarily indicate that the guiding processes are correctly represented. In particular, the presence of a quite homogeneous mixing line of water vapour below the inversion, topped by a pronounced inversion layer, and relative humidity below 10% is in stark contrast to the moist adiabatic profile that would be associated with a Rayleigh model. I recommend to include the possibility of long range advection of air with subsequent mixing into the picture, as well as considering the role of subsidence of air as a contribution to the inversion layer. Can several mixing processes in time and in the vertical combined look like a Rayleigh curve? This analysis should be placed in the context of the critique of Taylor (1984) on Rozanski and Sonntag (1982), and the very valuable study by Gedzelman (1988) to analyze Taylor's (1972) data set. Maybe a conclusion of your study could be a critique rather than a statement of consistency with the Rayleigh model? The Reviewer makes an excellent point that our flight conditions during the RAY (now CLR) flights are not consistent with moist adiabatic conditions that would be required for Rayleigh rainout processes to be actively occurring. However, the Rayleigh prediction is most consistent with our observations, compared to mixing scenarios as we now show in Figure 7. We have changed this discussion to indicate that observed profile is a fingerprint of previous airmass dehydration conditions that is retained by transport and downward mixing of dehydrated higher altitude FT air, consistent with past studies (Taylor, 1984; Gedzelman, 1988). (pg 19, ln 2-28 / pg 8, ln 10-33)

2. The interpretation of the stratocumulus case study as evaporating cloud drops needs further corroboration. The literature cited on the kinetic fractionation of evaporating droplets, such as Stewart (1975) considers rain drops, which are orders of magnitude larger, and have a vertical downward motion. This provides the potential of partial evaporation, leading to a kinetic fractionation signal in the dexcess. Cloud drops, however, will rapidly evaporate completely as they become unstable at smaller radii, leaving no sign of kinetic fractionation in the d-excess. Furthermore, they are suspended in the atmosphere, thus allowing for equilibrium fractionation. Without the consideration of rain evaporation in addition, or other processes, the explanation of the STC case is thus not viable. What is the potential of the occurrence of rain evaporation from these stratocumulus clouds? Can you estimate how much cloud water specific humidity would have to be added to produce the negative d-excess signal, and is this consistent with the saturation specific humidity in the cloud layer? Could there potentially have been a potential for rain evaporation or ice processes in the airmass upstream earlier instead? The Reviewer again brought up some excellent details to consider. Our calculations using the equations in Stewart likely apply to cloud droplets evaporating just like rain droplets. Similarly, we realize that the assumptions of the Worden equations are not valid either because that is for a case of a dehydrating airmass (closed system Rayleigh). However, there are 2 important

timescales to consider when exploring the influence of liquid droplet evaporation in the inversion layer. (1) The timescale at which a liquid droplet isotopically equilibrates with its surrounding vapor. (2) The speed at which droplets move through the inversion layer. We calculated the first timescale using Eqn. B5 is Bolot et al., 2013 which is ~2 seconds for a cloud droplet size of 15 microns. For our observations of vertical wind speeds, the time for a droplet to move from the bottom to the top of the inversion is 19 seconds. These calculations do indicate that cloud droplet evaporation starting in the lower inversion and finishing in the top of the inversion during the flight conditions in unlikely. However, if previous inversion conditions were colder, if the droplets were larger than 50 microns (like drizzle), or if the inversion had faster vertical wind speeds, these values could converge. We have edited the discussion of the STC flight day to reflect these new considerations. (pg 24, ln 19 – pg 25, ln 39 / pg 10, ln 17 – pg 11, ln 2)

3. Section 3.3 (DBL case) describes and interesting case of BL development, but is particularly difficult to follow. Consider reorganising the material in a more logical order, and formulate a clearer take-away in the end.

We have consolidated the DBL case study's results (originally Sect. 3.3) and discussion sections (originally Sect 4.3) so that the case study's key isotope features are identified, and their cause(s) immediately discussed. Furthermore, we have extensively revised Section 3.3 (now Section 4.3), which presents the DBL case study results.

4. The material included in the Supplement only seemed marginally useful. It is also confusing to have both an appendix and a supplement in a manuscript. Consider merging the supplement into the Appendix, keeping only Fig. S1.

This is a good suggestion. We have merged the Appendix and the important components of the Supplemental Information (SI) into the SI in order to further limit the length of the paper.

5. It is difficult to follow the description of the cases without additional context of the flight situations. Maybe Fig. 3 could be split up for each case, and supplemented by a weather chart similar to Fig. S6.

We have split up Fig 3a-c to make three plots for the three case studies (Fig 5, 8, 10) so that they can be placed closer to the text where they are being discussed. Weather charts have been added to the SI for each case studies (Fig. S5.1 - S5.3).

6. The discussion sections 4.3 and 4.4 appear repetitive to what has already been presented earlier and can be deleted.

We have consolidated the Results and Discussion sections so that discussion paragraphs directly follow their respective results sections. This reorganization has reduced repetitive parts of discussion sections 4.3 and 4.4. We have also reorganized the manuscript so that discussion of the campaign-wide observations (previously Sect 4.4, now Sect 3) now precedes discussion of the case study observations.

7. The second half of the first paragraph of the Conclusions section should be moved to the introduction/motivation.

We have made the suggested change.

8. Important studies that should be referred to are Taylor (1972, 1984), and He and Smith (1999) and Ehhalt (1975, 2005) and Tsujimura et al., 2007.

We did not focus on these studies because our primary interest was high-frequency deuteriumexcess observations, but the reviewer is right that they provide great context for our work. This is why we now point to the the detailed overview of airborne water vapor isotope studies reviewed in the introduction of Sodemann et al., 2017 (pg. 3, ln. 37/ pg. 3, ln. 9).

9. I found the structure of the manuscript somewhat confusing, in that first the flights around both Washington DC and Indianapolis are presented, and then almost all analysis focuses on flights from only one site, before returning to a wider overview in section 3.4. It would be more logical for the reader to move Sec. 3.4 ahead of the case studies, and then zoom into individual aspects. Consider adding a more intuitive way to present an overview of all data than a comparison to Rayleigh in 7.

This is a great suggestion. We have followed the Reviewer's suggestion by reorganizing the manuscript so that the campaign-wide results/discussion precede the case studies' results/discussion sections. We have added a new figure (Fig. 3) which shows dD and  $d^{18}$ O vs H<sub>2</sub>O<sub>v</sub> along every VP descent conducted during the campaign. We no longer present our results relative to Rayleigh by removing the right panel of Fig 7 (now Fig 4 post reorganization).

10. Several times in the writing, wording such as "Fig. x shows ..." is used. For conciseness, consider rephrasing this to sentences talking about what is shown in the figure, added by a figure reference "(Fig. x)".

This is a good suggestion, we have made changes where appropriate.

# **2** Detailed comments

P3, L8: "However, relatively few...": As far as I am aware there is only one published dataset of airborne d-excess measurements. Why the citation of Schmidt et al., 2005 here? We have reworded the sentence and deleted the reference to Schmidt et al., 2005, which provides modeled d-excess.

P3, L34: Please state if the inlet or parts of it were heated.

We have added a sentence stating the TWVIA inlet was not heated during the calibrations. (pg 5,  $\ln 27-28 / \text{pg 4}$ ,  $\ln 35-36$ )

P4, L9: How sensitive are the measurements to variations of cavity pressure and temperature during flight?

Once the isotope analyzer has warmed up (i.e. reached the manufacturer recommended pressure and temperature values for operation), cavity pressure and temperature only vary  $(1\sigma)$  by  $\pm 0.02$  Torr and  $\pm 0.08^{\circ}$ C, respectively, over a vertical profile descent on average. This is within the operating specification given by the manufacturer. (pg. 5 ln 13-15 / pg 4, ln 21-22)

P9, L21: Worden et al. 2007 formulate their models for rain evaporation (falling condensate), not cloud evaporation, see their supplementary information.

We have removed all discussion of the Worden et al., 2007 model (see our response to major comment #2 for more explanation).

P10, L10: "Rayleigh-consistent observations": Consider rephrasing in light of major comment 1 This is a great suggestion. We have renamed the RAY case study to CLR, which is now representative of the case study's clear sky conditions rather than a possible interpretation of the day's dominant isotopic process.

Fig.4: Consider adding a panel that shows more mixing lines, e.g. between the bottom and top of the inversion, or the bottom and top of the BL, how do these compare with the complete mixing line?

We have split up Fig. 4 into three figures for the three case studies (Fig. 7 CLR, Fig. 9 STC, and Fig 12 DBL). Each new case study figure has an additional panel showing mixing lines for scenarios relevant to the flight day. The new mixing panels are discussed in each case study's respective sections (Sect 4.1-4.3).

P12, L1: Consider using a standard symbol such as q or m for specific humidity or mixing ratio of water vapour, rather than H2O.

We respectfully choose to retain  $H_2O_v$  to represent water vapor mole fraction, as it is a common convention in the atmospheric chemistry field.

P14, L8: "relative to Rayleigh" rephrase as e.g. "Rayleigh model predictions" The sentence now reflects this suggestion.

P14, L16: "d-excess switches to tracking" rephrase This sentence has been rephrased.

P14, L13-22: hard to follow, rewrite this section for clarity

We thank the Reviewer for identifying this paragraph that could be improved. We have revised this paragraph to include a summarizing, take-away sentence, which highlights the unique d-excess characteristics of the first vertical profile flown on STC (pg 23, ln 17-20 / pg 9, ln 36-40).

P16, L6: "Differences" - difference to what is shown? Are these not absolute values? This sentence has been rephrased.

Fig.6: How are RL and INV defined? You could provide a quantitative comparison of averaged quantities for both layers that supports their distinction.

To address this comment, we have added a sentence that reads, "We define the base of the RL using the same approach described in Section 2.4 ( $d\theta/dz$  and  $|d(H_2O_v)|/dz$  threshold values) for determining the base of the INV ( $z_{INV}$ )."

P17, L14: Not clear what the take-away from this section is. Revise paragraph. We thank the Reviewer for identifying this paragraph as an opportunity for improvement. We have revised this paragraph about our campaign-wide vertical profile observations. (Sect. 3, paragraph 1-2)

P17, L17: "Figure xxx shows": revise according to major comment 11 We have reworded sentences that begin with "Figure # shows..." throughout the manuscript. P18, L12: Not clear what the take-away from this section is. Revise paragraph. We have deleted this paragraph because we decided its inclusion does not benefit the overall story about the campaign-wide observations section.

P18, L16: Why the citation of Lee et al. (2006)? Compare Gedzelman (1988) and the discussion of the Taylor (1972) data (major comment 1).

During our consolidation of the results and discussion sections for each case study, we have removed the indicated sentence. We now discuss the possibility of Rayleigh-consistent condensation occurring upwind of the CLR measurements (see response to major comment #1). We also reference the suggested paper (pg 19, ln 17-28 / pg 8, ln 22-33)

P19, L3: The thermodynamic characteristics of the profiles are not consistent with a Rayleigh processes, even though the isotope composition may be – what is the conclusion from that finding?

We now discuss the possibility of Rayleigh-consistent condensation occurring upwind of the CLR measurements (see response to major comment #1). (pg 19, ln 17-28 / pg 8, ln 22-33)

P19, L15: How would cloud droplets be lofted but not the surrounding vapour? If vapour and droplets move together, how can non-equilibrium fractionation result?

The vapor and droplets do not necessarily have to remain in isotopic equilibrium under changing conditions, including mixing of airmasses. We have revised the indicated sentence so that cloud droplet evaporation is discussed from the framework of timescales, e.g. the time required for condensate and vapor to isotopically equilibrate vs. time required to transport a cloud droplet from the bottom to the top of the inversion. Please see our response to major comment #2 for more detail. (pg 24, ln 19 – pg 25, ln 39 / pg 10, ln 17 – pg 11, ln 2)

P20, L16-20: Consider the possibility of rain evaporation upstream (major comment 2). We now discuss how rain droplet evaporation in dry, cold conditions upwind of the STC flight measurements could possibly explain the minimum in d-excess at the top of the inversion layer. (pg 25, ln 33-36 / pg 10, ln 36-38)

P21, L10: What relative roles could vertical motion and entrainment of dry air from aloft play in the evaporation of the cloud layer?

This is an interesting question that we believe that stable isotope observations can inform about the occurrence of cloud evaporation, but are not sure that it's possible to distinguish the driver of that evaporation (cloud drops moving into the FT versus entrainment of dry air into the cloud) using our observations downwind of clouds. This would require a sophisticated cloud microphysics model, similar to Bolot et al. (2013). Respectfully, it is outside of the scope of this paper, but an exciting future application. We do, however, cite papers that give evidence for stratocumulus cloud droplet evaporation occurring within and above the cloud layer due to differences in entrainment efficiency (pg 24, ln 36- pg 25 ln 2 / pg 10, ln 12-15).

P21, L20-29: Rephrase for clarity. Could ice-phase processes have been relevant further upstream, earlier in time?

We have rephrased the paragraph about the effect condensation in the presence of ice has on vapor d-excess for clarity. We also now include the following sentence, "It is possible, however,

that condensation under ice-supersaturated conditions occurred prior to the STC flight, and that the resulting isotopic imprint was maintained during transport to Indianapolis." (pg 26, ln 28-30 / pg 11, ln 15-17)

P22, L12: What is the Rapid Refresh Model? Could be deleted here. Reference to the model name has been deleted.

P22, L7-34: reorganize and rewrite this section for clarity We have revised the entire results/discussion sections related to the DBL case study to improve their clarity. (now Sect 4.3)

P22, L36: "so we look" rephrase This phrase has been deleted as part of the manuscript reorganization.

Fig. A1: What is the reason for the large scatter between both instruments in the mid-range of humidities? Does longer averaging of the data provide a better agreement? Was the scatter similar for each of the flight? How is the scatter for the downward profiles only? The scatter corresponds to few points relative to most of the points that show good agreement. The mid-range humidities in Fig. A1 (now Fig. S1 due to the Appendix-SI merger) correspond to vertical profiles. These few points that do not show as good agreement result from the TWVIA data being low pass filtered relative to the Picarro due to a longer residence time, as we mention in Sect 2.2.2 (SI pg 5, ln 23 / SI pg 4, ln 31).

# Authors' Responses to Reviewer 2

- Author responses are indicated in blue font.
- Locations of our edits in the "strike-through" and final versions of the manuscript are provided with the following convention: ([strike-through version] pg. #, ln # / [final version] pg. #, ln #).

# **Major Comments**

1. The results could be better distilled. The paper is quite lengthy and detailed. As a result, it is easy for key points to become buried. Usually, I would not list this as a major concern, but in this case, I found much of the critical interpretation for the analysis was lost among the long-winded descriptions of the data. It might help readers if the critical arguments and conclusions were emphasized near the beginning and/or end of each sub-section. One place where this is very much needed is in Section 3.3, which discusses the developing boundary layer case (DBL). Here, more time could be spent on the interpretation of the causes of the fascinating differences in atmospheric structure among the profiles rather than on re-describing Figure 6. To address the Reviewer's comment, we have consolidated the case studies' results (originally Sect. 3.1-3.3) and discussion sections (originally Sect 4.1-4.3) so that each case study's key isotope features are identified, and their possible cause(s) immediately discussed. The consolidated sections are now located in Section 4.1-4.3. We have also extensively revised Section 3.3 (now Section 4.3), which presents the DBL case study results.

2. Section 4.2 is another section that is somewhat difficult to follow, largely because of confusing terminology. How do the "scenarios" relate back to the equations presented in the methods? Also, "partial evaporation" and "near complete evaporation of a semidehydrated drop" could easily be confused as one and the same. As a result, it is not at all clear which "scenario" best represents the data.

We thank the Reviewer for identifying Section 4.2 as an area that could see improvement. We have chosen to remove the cloud evaporation equations (see our response to comment #4), and by association, the paragraphs which discussed the degree of droplet evaporation under two scenarios. In their place, we have added text which evaluate the likelihood of cloud evaporation and its impact on d-excess under meteorologically-relevant time scales. We believe our changes have greatly improved the clarity of Section 4.2. (pg 24, ln 19 – pg 25, ln 39 / pg 10, ln 17 – pg 11, ln 2)

3. I am not convinced by the analysis that Rayleigh distillation is the dominant process determining the vertical isotopic structure of the boundary layer. First and foremost, there are many papers that show the contrary; a few case studies are not sufficient to prove otherwise. Previous papers that have measured water vapor isotope ratio vertical profiles in situ (and have shown profiles consistent with processes other than distillation) include He and Smith 1999, Galewsky et al 2007, Noone et al 2013, Bailey et al 2013, Sodemann et al 2017, and Kelsey et al 2018. These studies contrast, to some degree, with the early work of Ehhalt, which was republished as Ehhalt et al 2005.

We thank the Reviewer for these suggested literature citations. We have revised our interpretation of the RAY (now CLR) case study to express that the Rayleigh-consistent

observations likely reflect the isotopic imprint of prior condensation under saturated conditions, followed by advection of this imprinted signal to the Indianapolis study site. Additionally we note that the lower free troposphere measurements appear more consistent with mixing lines, which is indicative of mixing between subsiding free troposphere air and boundary layer air. We include the some of the suggested references to support this discussion (pg 19, ln 2-28 / pg 8, ln 10-33).

Second, I am not entirely convinced that Rayleigh distillation gives the best physical interpretation for the Indianapolis "Ray" profiles. The paper argues that Rayleigh distillation is a good model when the boundary layer is dry adiabatic (and therefore that no clouds or precipitation exist). This is contrary to expectation: distillation depends on condensation and precipitation. Furthermore, other studies (see above) have shown nearly the opposite: that Rayleigh is appropriate when the boundary layer follows a moist adiabat but not when it follows a dry one. A clear exception, of course, is if the distillation occurs upwind and imprints an isotopic signature that is then advected downwind. One of the earliest papers discussing this phenomenon is Gedzelman 1988. Is it possible that advection is affecting the Indianapolis isotope ratio profiles? If so, this could make for an interesting discussion on whether moist convective processes regionally set the humidity structure of the lower atmosphere locally, which others have argued for tropical/subtropical regions (e.g. Brown et al 2008, Lee et al 2011, Bailey et al 2013).

The reviewer make an excellent point that our fight conditions during the RAY (now CLR) flights are not consistent with moist adiabatic conditions. However, the Rayleigh prediction is most consistent with our observations, compared to mixing scenarios. We have changed this discussion to indicate that observed profile is an imprint of previous airmass dehydration conditions, and cite the suggested studies (pg 19, ln 17-28 / pg 8, ln 22-33).

Third, extra care must be taken in matching data to hypothetical Rayleigh curves since these can be designed to fit many data. A good example of this is found in Noone et al 2013. Consequently, it may be difficult to truly distinguish Rayleigh from mixing processes unless the theoretical end-members are well constrained. It is not clear in the manuscript from whence the theoretical end-members for Figure 4 are derived. Some description of the assumptions made would be a valuable addition to the analysis and perhaps make the case that the Ray flights are, in fact, illustrative of (upwind?) Rayleigh distillation in a much more compelling manner. The Reviewer makes a good point. We do note, however, that our Rayleigh curves are calculated using an objective method (Sect. 2.5; Eq (1)). The initial isotopic composition of the Rayleigh air parcel  $(R_0)$  is determined from the average observed boundary layer delta values. The equilibrium fractionation factors are calculated for the lifting condensation level temperature. We also demonstrate in Fig. S4 that varying the equilibrium fractionation factor by observed temperature does not lead to significantly different shaped Rayleigh curves. We now note that mixing endmembers are informed by actual observed delta values in different atmospheric layers (pg 10, ln 26 / pg 6, ln 26). Figure 7 now makes a direct comparison of our observations to Rayleigh theory and mixing scenarios.

4. Some care should be taken in describing the Worden et al 2007 and Stewart 1975 models and applying them to the case of stratiform clouds. Both models were designed to describe freely falling raindrops. In the original presentation of the model, Worden suggests raindrops undergo both an equilibrium fractionation and an effective fractionation, and that "this assumption is

unlikely to be valid when raindrops are small...". I note that Equation 3 substitutes a kinetic fractionation coefficient in place of the effective factor. What impact does this have? How does the model work if the assumptions of large falling drops are violated? The Worden model describes isotopic depletion with a loss of water from an air parcel. How can it be applied to describe a gain of moisture by the atmosphere? I have similar concerns with use of the Stewart model and would like to see more justification for these model choices for the case of stratiform clouds. Also, the equations presented from Stewart are from Equations 2 and 3 of the original paper, and there is an alpha missing in the denominator of the beta equation.

We have opted to remove the Worden et al, 2007 and Stewart, 1975 equations from the paper for the reasons that both reviewers mention. Instead of calculating the impact different degrees of evaporation could have on surrounding vapor, we have reframed this discussion by evaluating the likelihood of cloud droplet evaporation impacts based on transport and equilibration time scales. We now show that a droplet would isotopically equilibrate with surrounding vapor faster than the time required for transport of a droplet from the bottom to the top of the inversion (Equation B5 in Bolot et al., 2013). These calculations suggest cloud droplet evaporation may not be responsible for the observed d-excess minimum during the STC flight. However, we maintain that a negative d-excess signal resulting from cloud or rain droplet evaporation could have been transported from an area upwind of the STC measurements. (pg 24, ln 19 – pg 25, ln 39 / pg 10, ln 17 - pg 11, ln 2)

5. I found it difficult to identify the case studies in the flight figures due to distinct nomenclature. The figures use numbers, the text uses pseudonyms, and only the table provides both these plus dates. I would recommend one naming convention, preferably related to flight number or date. The reason being that a priori, it seems difficult to know whether the "Ray" days will really be Rayleigh-like.

The Reviewer brings up an excellent point. We have deleted any reference to research flight codes (RF#), and now refer to the flights by date. We maintain pseudonyms for the case study flights, but we have made sure that figures include both the flight date and case study pseudonym where appropriate. We have renamed the RAY case study to CLR (for clear skies) so that the pseudonym characterizes the day's meteorology, rather than a possible interpretation of the isotopologue data.

6. The Isotope Theory section suggests there are "three common ways the isotopic composition of the atmospheric H2Ov can change." I would have thought these would be condensation, evaporation, and mixing. Rayleigh distillation is just an example model for condensation processes. Similarly, cloud evaporation is just one type of evaporation that can affect the atmosphere's isotopic composition.

We have reworded this section to indicate that we employ different models to represent condensation and mixing processes. Please refer to our response to comment #4 for our explanation for removing the cloud evaporation equations. (pg 9, ln 24-26 / pg 6, ln 1-4)

7. I have some trouble understanding how partial cloud evaporation can cause a minimum of deuterium excess near the inversion layer. Evaporation tends to favor the diffusion of the D relative to 18O. Why wouldn't partial evaporation result in an enrichment of the surrounding vapor? Perhaps I am missing something, but my hunch is most readers will also have this

impression. It might be worth explaining the physical underpinning behind these conclusions in greater depth, perhaps in Sections 4.2 or 4.4.

The Reviewer is correct that evaporation (from an infinitely large source) typically imparts a positive d-excess signal on the surrounding vapor. This is why we believe cloud evaporation to be responsible for the slight increase in d-excess in the middle of the inversion (Fig. 6b). As a droplet evaporates, its own d-excess signal becomes more negative, so subsequent complete evaporation of the droplet in another region would act to decrease the surrounding vapor d-excess (Aemisegger et al., 2015; Sodemann et al., 2017). We also now include a calculation (Bolot et al., 2013) showing that a droplet would isotopically equilibrate with surrounding vapor faster than the time a droplet would be transported from the bottom of the inversion to the top of the inversion. This calculation indicates cloud evaporation during the STC flight may not be responsible for the observed d-excess minimum. We do, however, maintain that evaporation of cloud/rain droplets upwind might have caused the d-excess minimum. We support this possible explanation with references that show that stratocumulus cloud droplet evaporation occurs at different altitude in and above the cloud layer, and that the inversion layer above stratocumulus clouds are not homogeneous in terms of depth or thermodynamic properties. (pg 24, ln 19 – pg 25, ln 39 / pg 10, ln 17 – pg 11, ln 2)

8. The calibration documentation is quite thorough and comprehensive. I was prompted, however, to ask a few follow up questions regarding the variations in concentration dependence shown for dD. Is it possible one would get a different answer if concentration biases were adjusted first and VSMOW-scaling applied separately? It also appears that there are higher errors in dD at low isotopic values at all water vapor concentrations, not just at the low concentrations. Is it possible that lower precision at low isotope ratios causes the appearance of "irreproducibility" in the concentration dependence?

Our calibration procedure does begin with concentration-dependence corrections (Sect. S2; note that the Appendix has been merged with the SI). We determined that VSMOW scaling of the concentration dependence-corrected delta values was not necessary (Figure S2.3). The Reviewer makes an interesting point about the apparent dD irreproducibility. However, we note that the "irreproducibility" at mid-range humidities (3000-8000 ppmv) only lie on one side of the correction curve (Fig. S2.2b), so it is an offset. On the other hand, variability on either side of the correction curve is observed for drier conditions (550-3000 ppmv; Fig. S2.2b), which indicates the low precision idea could be a possible explanation. We note that the uncertainty is only consequential for very low  $H_2O_v$  mole fractions, where these depleted delta values are observed. We have added a sentence discussing the possible low precision and offset ideas (strike-through SI: pg 5, ln 5 / final SI: pg 5, ln 7).

## **Minor comments**

Page 1 Line 35 – no need for ":" after "include" The ":" has been removed.

Page 2 Line 10 –I had trouble distinguishing the conditions at a moisture source region from "surface H2Ov sources." Perhaps it might be more clear to say "an air parcel's moisture source region, including the geography of the source and its meteorological conditions?" This sentence has been changed to reflect the Reviewer's suggestion.

Page 2 Line 19 – I think "further exchange" is meant instead of "equilibrium?" "Equilibrium" has been replaced by "further exchange".

Page 3 Line 3 – I might remove "point in" before time, since I initially confused "point" with space.

"Point in" has been removed.

Page 3 Line 6 – Perhaps best to say "higher…resolution" since aircraft is not as high resolution as slower-moving platforms. This is a good point, "high" has been changed to "higher".

Page 3 Line 11 – Perhaps best to say that "measurements of vertical profiles" were conducted. We have incorporated this suggestion.

Page 4 Line 28 – Perhaps "produce" for "emit?" Thank you for this suggestion, we have replaced "emit" with "produce".

Page 6 Figure 2 – Could the three analyzed flights be emphasized, perhaps by making the other flight lines dashed?

This is a good suggestion, the case study flight paths are now indicated with solid lines, and all other flight paths are indicated with dashed lines.

Page 7 Table 1 – Table caption/title should provide some explanation of the "codes" and what is meant by "support study"

Per the Reviewer's comment #5, we have modified the flights codes so that the flights are identified by their date.

Page 9 Line 13 – One of many examples of the great care that was taken in the analysis

Page 9 Line 31 – This appears to be the only place where "q" is used instead of "H2Ov." Consistency would help.

Thank you for catching this. As noted in our response to major comment #4, we no longer include the Worden et al., 2007 equation that the Reviewer is referencing.

Page 10 Equations – This appears to be the only place "Rvap" is used instead of "Rv." Again, consistency would help.

Thank you for catching this oversight. As noted in our response to major comment #4, we have removed the Stewart, 1975 equation that the Reviewer is referencing.

Page 11 Figure 4 – All the lines are "solid," thus it might be best to say the "black" line to distinguish it from the "pink" one. Also, I don't fully understand how the average mixing ratio is given by a gray envelope. Shouldn't the average be a point?

Thank you for these suggestions. We have replaced "solid" with "black", and we have clarified that the grey envelope indicates the inversion layer, which is defined by the average  $H_2O_v$  mole fraction observed at  $z_{INV}$  and  $z_{FT}$  during the vertical profiles. (Fig 4 has been split into three figures for each of the case studies, Fig 7- CLR; Fig 9 – STC; Fig 12 – DBL)

Page 12 Line 7+ - Here is an example where it's easy for the reader to become lost in all the number-reporting. This paragraph would greatly benefit from a sentence that provides a bit more of a picture of what is going on physically.

Thank you for this suggestion. We have added text to identify the important meteorological characteristics of each atmospheric layer on this day, and also decreased the amount of number-reporting. (pg 17, ln 6-19 / pg 7, ln 35 - pg 8, ln 9)

Page 12 Line 18 – I might say "predictions" instead of "theory." We have made this substitution.

Page 13 Figure 5 – Caption should explain what the shaded area for the inversion is and what the envelopes around the profiles are.

The caption now indicates that the inversion layer (previously grey bands, now blue bands) lies between  $z_{FT}$  and  $z_{INV}$ , and that the shading around the isotope VP measurements correspond to total measurement uncertainty. (Now Fig 6)

Page 14 Line 6 – How was the average range of the inversion calculated? From how many days or which days of data?

Similar to our changes to the caption of Fig. 4 (now Fig 7, 9 and 12) and Fig. 5 (now Fig. 6), we now specify that the bounds of the inversion layer range, which are unique to each flight, are defined by the average  $H_2O_v$  mole fractions observed at  $z_{FT}$  and  $z_{INV}$  for each flight.

Page 14 Line 16 - I'm not sure I agree that the data start "tracking" the mixing line. There just aren't enough points for me to be convinced of that. Perhaps "approaches" the mixing line? We have made this substitution.

Page 14 Line 26 – This seems like an important argument explaining how is STC different from Ray, and yet it is buried halfway down the page. Perhaps it could be moved up in the subsection.

We have moved this sentence to the second paragraph of Sect 4.2. (pg 21, ln 7 / pg 9, ln 18)

Page 15 Figure 6 – Most axes appear consistent across panels except for theta. Was this purposeful?

Good catch! We have modified the theta range of Fig 6a (now Fig. 11a) for consistency.

Page 17 subtitle – I'm not really sure what "general observations" means. Would "observations from other flights" be more descriptive?

We have renamed this subsection: "Airborne campaign observations of  $H_2O_v$  isotopologues in the lower troposphere". (Now Sect. 3 title)

Page 17 Line 21 – which observations are used here for this argument? We now specify that the "DBL" case study observations are being discussed in this sentence.

Page 19 Line 19 – A sentence or couple words reminding the reader what "Scenario 1" is would be appreciated.

We have revised the STC discussion paragraphs (see our response to major comment #2), and no longer reference cloud evaporation scenarios 1 and 2.

Page 19 Line 22 – I think "drier" is meant. Great catch, we did mean "drier". During our revision of this section, we have deleted "drier".

Page 21 Figure 9 - I would recommend dots (or some symbol) instead of vertical lines to indicate the values of dxs expected. Otherwise, it is not clear that the reader should look for the intersection of the various lines.

We have revised Section 4.2 per major comment #2, and in doing so, we have removed Figure 9 from the manuscript.

Page 22 Line 24 – Excellent synthesis sentence: highlights the key point nicely.

Page 22 Line 36 – I disagree that these are some of the few data of this kind. There are quite a few studies that are not cited in this work. Please see my major comments for ideas. We have reworded this sentence to express that there are few studies that report vertical profiles of d-excess.

Page 23 Line 24 – Kelsey et al 2018 also report dxs profiles. Thank you for this suggestion, we now reference the Kelsey et al, 2018 study here, and in the introduction.

Page 24 Line 4 – Perhaps "investigate" for "interrogate." We have made the suggested replacement.

Page 27 Line 5 – "To check calibration…" against? of? We have rephrased this sentence.

Page 30 Line 14 – The notation seems to change here. Should all isotopologues have subscript "v?"

Thank you for catching this, we have added missing subscripted v's to the isotopologue ratios. (Now in the SI, Section S6)

Page 31 Figure D1 – At first I thought the black square was a symbol in the legend. It might be more clear simply to use the caption to say that the striped region shows the expected range for the observations, or something to that effect.

Thank you for this observation. We have modified the legend to better represent the slanted line region, and we have added a sentence to the caption describing the saturation values that the slanted lines represent. (Now in the SI, Figure S6)

## Vertical profile observations of water vapor deuterium excess in the lower troposphere

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Abstract. We use <u>airborne measurements of water vapor (H<sub>2</sub>O<sub>v</sub>) stable</u> isotopologues-pie vertical profile measurements and complementary meteorological observations to examine how boundary layer <u>dynamics</u>, cloud <u>processing</u>, and <u>atmospheric</u> mixing processes-influence the vertical structure of  $\delta D$ ,  $\delta^{18}O$ , and deuterium-excess (d-excess =  $\delta D - 8 \times \delta^{18}O$ ) in the boundary layer, inversion layer, and lower free troposphere. Airborne measurements of water vapor (H<sub>2</sub>O<sub>v</sub>) stable isotopologuesFlights were conducted around two continental U.S. cities in February – March 2016. Nine research flights were designed to characterize the  $\delta D$ ,  $\delta^{18}O$ , and d excess and included vertical profiles extending from <u>near</u> the surface to  $\leq 2$  km. We examine observations from three unique case study flights in detail. One case study shows <u>observations H<sub>2</sub>O<sub>v</sub> isotopologue vertical profiles</u> that are consistent with Rayleigh isotopic distillation theory coinciding with clear skies, dry adiabatic lapse rates within the boundary layer, and

- 20 relatively constant vertical profiles of wind speed and wind direction. This suggests that the air mass retained the isotopic fingerprint of dehydration during moist adiabatic processes upwind of the study area. Also, observed d-excess values in the free troposphere were sometimes observed-larger than Rayleigh theory predicts, which may indicate mixing of extremely dehydrated air from higher altitudes. The two remaining case studies show that  $H_2O_{\tau}$  isotopic signatures above the boundary layer are sensitive to cloud processes and complex air mass mixing patterns. These two case studies indicate anomalies in the d-excess signature
- 25 relative to Rayleigh theory, such as low d excess values at the interface of the inversion layer and the free troposphere, which is possibly indicative of and indicate cloud evaporation processes and complex boundary layer development. The most notable case study with stratocumulus clouds present had extremely low (negative) d-excess values at the interface of the inversion layer and the free troposphere, which is possibly indicative of cloud or rain droplet evaporation. We discuss how We discuss possible explanations for the observed d excess anomalies, such as cloud evaporation, wind shear, and vertical mixing. *I*in situ H<sub>2</sub>O<sub>v</sub> stable
- 30 isotope measurements, and d-excess in particular, could be useful for improving our understanding of moisture processing and transport mixing occurring between the boundary layer, inversion layer, and free troposphere.

## **1** Introduction

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Water vapor  $(H_2O_v)$  in the lower troposphere modulates processes including cloud formation, precipitation, severe weather development, atmospheric circulation, radiative forcing, and climate feedbacks (Held and Soden, 2000; Kunkel et al., 2012; Tompkins, 2001; Trapp et al., 2007; Trenberth, 2011). Accurately representing these dynamic, mesoscale processes in models can be difficult, and efforts to improve parameterizations are on-going (Gerber et al., 2013; <u>de Lozar and Melledo, 2015;</u> Park et al., 2017; Wood, 2012; Yamaguchi and Feingold, 2013). Some active areas of research include: quantifying the inversion layer entrainment flux (Wood, 2012), refining entrainment-cloud evaporation relationships (Gerber et al., 2013; Yamaguchi and Feingold, 2013), and updating cloud evaporation schemes with new cloud classes (Park et al., 2017).

Free troposphere entrainment and cloud evaporation influence the maintenance of the cloud layer, which in turn influences
radiative forcing (Gerber et al., 2013; Yamaguchi and Feingold, 2013). The nature of H<sub>2</sub>O<sub>v</sub> as a climate feedback agent adds further complexity to our understanding of H<sub>2</sub>O<sub>v</sub>'s role in weather. Anthropogenic greenhouse gas emissions have resulted in increasing global temperatures, enhanced evaporation from soil and the oceans, and higher atmospheric concentrations of H<sub>2</sub>O<sub>v</sub>, the dominant absorber of infrared radiation (Held and Soden, 2006; <u>Hurley and Galewsky, 2010</u>; Wille<u>t</u>t et al., 2007). Warmer temperatures and more humid conditions have caused a shift towards less frequent, but more intense precipitation events, increasing the risk of both
floods and droughts (Roque-Malo and Kumar, 2017; Trenberth, 2011). H<sub>2</sub>O<sub>v</sub> also modulates production of the dominant atmospheric oxidant, the hydroxyl radical (Thompson, 1992). Thus, accurately representing H<sub>2</sub>O<sub>v</sub> in mesoscale processes is of great importance in a warming world.

 $H_2O_v$  stable isotopologue measurements are a potential tool to inform our understanding of the distribution and dynamics of  $H_2O_v$  in the lower troposphere (see review by Galewsky et al., 2016).  $H_2O_v$  stable isotopologue ratios, i.e. the ratio of heavy

- 20 (HDO or  $H_2^{18}O$ ) to light ( $H_2^{16}O$ ) molecules, captures the origin and mixing of moisture sources as well as the condensation and cloud processes that modified that moisture. The  $\delta$ -notation indicates the sample's heavy-to-light isotope ratio reported relative to an international standard ( $\delta = R_{sample}/R_{standard} 1$ ), where  $\delta$  is multiplied by 1000 to report in units of per mil (‰). Both equilibrium fractionation, which separates heavy and light isotopolougues isotopologues based on their unique vapour-pressure differences, and kinetic fractionation, which is controlled by differences in diffusion rates of the isotopologues, lead to different changes to modulate
- 25 the isotopic signatures in the of atmospheric  $H_2O_ve$  (Dansgaard, 1964). Isotopic fractionation processes act to enrich/deplete both HDO and  $H_2^{18}O$  relative to  $H_2^{16}O$  in atmospheric waters, resulting in co-varying  $\delta D$  and  $\delta^{18}O$  signatures. The Rayleigh distillation model describes the equilibrium fractionation between vapour and condensation that dominates the variability of  $H_2O_v$  isotopes in the troposphere. This model can be used to calculate the degree to which condensate is removed from an air parcel as it cools, such as when it is undergoes ascentascends from the surface to higher altitudes. Variability in  $\delta D$  and  $\delta^{18}O$  often co-vary by a factor of
- 30 approximately 8 because the ratio of the  $\delta D$  and  $\delta^{18}O$  equilibrium fractionation factors is approximately 8:1 at typical surface temperatures (Dansgaard, 1964). The second-order isotope parameter deuterium excess (d-excess =  $\delta D - 8 \times \delta^{18}O$ ) can be used to help identify the non-equilibrium kinetic fractionation processes. For example, remoistening of the atmosphere by raindrop evaporation imparts a kinetic isotopic fingerprint on H<sub>2</sub>O<sub>v</sub> d-excess (Field et al. 2010; Lawrence et al. 2004; Risi et al. 2008; Worden et al. 2007; Wright et al. 2009).
- 35 <u>Tropospheric vapor d-excess measurements may provide important information about cloud microphysics, convection</u> processes, precipitation recycling, and free troposphere (FT) entrainment.  $H_2O_v$  stable isotopologue ratios, i.e. the ratio of heavy (HDO or  $H_2$ <sup>48</sup>O) to light ( $H_2$ <sup>46</sup>O) molecules, can contain information about the Local-scale studies have examined the role of moisture source, condensation history, meteorological conditions at an air parcel's moisture source and land surface

evapotranspiration region on d-excess signatures , including the geography of the source and its meteorological conditions surface  $H_2O_x$ -sources, like evapotranspiration, as well as its temperature-dependent phase change history since that point (Benetti et al., 2014; Delattre et al., 2015; Kelsey et al., 2018; Lai and Ehleringer, 2011; Uemura et al., 2008; Welp et al., 2012). The  $\delta$ -notation indicates the sample's heavy to light isotope ratio reported relative to an international standard ( $\delta = R_{sample}/R_{standard} - 1$ ), where  $\delta$  is commonly multiplied by 1000 to report in units of per mil (‰).

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Isotopic fractionation processes act to enrich/deplete both HDO and H<sub>2</sub><sup>18</sup>O relative to H<sub>2</sub><sup>16</sup>O in atmospheric waters, resulting in co-varying δD and δ<sup>18</sup>O signatures. Rayleigh distillation theory can be used to calculate the degree of equilibrium fractionation that occurs when condensate is removed from an air parcel as it cools, such as when it is undergoes ascent from the surface to higher altitudes. Rayleigh theory assumes that when saturation is reached, the condensate is removed immediately from the system via precipitation, thus no equilibrium further exchange\_occurs between the two phases. \_\_The second order isotope parameter deuterium excess (d-excess = δD - 8\*δ<sup>18</sup>O) can be used to identify the type of fractionation occurring, equilibrium or kinetic, given that the ratio of the δD and δ<sup>18</sup>O equilibrium fractionation factors is approximately 8:1 at typical surface level temperatures (Dansgaard, 1964). The ratio of the δD to δ<sup>18</sup>O kinetic fractionation factors is typically less than 8 and decreases with decreasing relative humidity.

- 15 Observations of d excess have been used to deduce meteorological conditions at the evaporation source, assuming it is a conservative tracer not changed by transport and rainout processes (Benetti et al., 2014; Delattre et al., 2015; Steen Larsen et al., 2014; Uemura et al., 2008), but there is also evidence that d excess is not a conserved tracer of evaporative origin if other significant sources of vapor exist, especially outside of the marine environment (<u>Aemisegger et al., 2014;</u> Gorski et al., 2015; Griffis et al., 2016; Fiorella et al., 2018, Parkes et al., 2017; Welp et al., 2012). For example, the unique d excess signature of combustion-derived H<sub>2</sub>O<sub>x</sub>-has been used to quantify the contribution of combustion emissions to boundary layer vapor (Fiorella, et al., 2018;
- Gorski et al., 2015), and several studies have demonstrated the influence of sublimation, vapor deposition, and land surface evapotranspiration on the atmospheric H<sub>2</sub>O<sub>x</sub>-d-excess signature (<u>Aemisegger et al., 2014;</u> Casado et al., 2016; Galewsky, 2015; Griffis et al., 2016; Lai and Ehleringer, 2011; Lowenthal et al., 2016; Moore et al., 2016; Parkes et al., 2017; Schmidt et al., 2005; Samuels Crow et al., 2014; Welp et al., 2012).
- 25 Airborne d-excess measurements may provide information about cloud processes, precipitation recycling, FT entrainment, and more generally, the vertical structure characteristics of d excess over different land cover and in different seasons. Measurements of d-excess have been used to estimate below-cloud precipitation evaporation (Aemisegger et al., 2015; Froehlich et al., 2008; Wang et al., 2016), and mixing between the boundary layer (BL) and <u>free troposphere (FT)</u> from stationary platforms near the surface or at high-altitude mountain sites (Bailey et al., 2015; Benetti et al., 2015; 2018; Froehlich et al., 2008; Galewsky,
- 30 2015; Lowenthal et al., 2016; Samuels-Crow et al., 2014). Kelsey et al. (2018) conducted mobile traverses along leeward and windward slope roads of a mountain in an instrumented vehicle to report vertical profiles of d-excess. While some-high-elevation surface monitoring sites-hasve the advantage of sampling BL and FT air over a diurnal cycle, they-it\_does not provide a complete picture of the H<sub>2</sub>O<sub>v</sub> isotope vertical profile (VP) at a discrete point in-time. Satellite measurements, which can provide discrete VP measurements, currently\_only eurrently\_provide middle troposphere δD profiles (Herman et al., 2014; Worden et al., 2012).
- 35 Airborne platforms are capable of  $\delta D$ ,  $\delta^{18}O$ , and d-excess VP measurements at high<u>er</u> spatiotemporal resolution, and have been conducted since the 1960s extending from the lower troposphere to the stratosphere to investigate a variety of <u>science-scientific</u> questions (overview in Sodemann et al. (2017)). <u>However, relatively fewdDue to either the study</u>'s objectives or limitations of the instrumentation (Dyroff et al., 2015; Herman et al., 2014), however, only one airborne H<sub>2</sub>O<sub>v</sub> isotope study<u>ies, to our knowledge</u>,

has<u>ve</u> reported d-excess measurements in the lower troposphere to provide high vertical resolution snapshots at discrete time points in the boundary later (Sehmidt et al., 2005; Sodemann et al., 2017), due to either the study's objective or limitations of the instrumentation (Dyroff et al., 2015; Herman et al., 2014). Additional airborne d-excess observations may improve our understanding of the role that cloud processes and convection have on determining the moisture distribution of the troposphere

5 (Noone, 2012; Bolot et al., 2013). Our stable H<sub>2</sub>O<sub>3</sub>-isotope measurements over two continental sites is a starting point in filling the field's gap in understanding variability in the lower troposphere d excess profile, and what it reveals about lower troposphere moisture processing on relatively small regional scales.

In this study, we present stable H<sub>2</sub>O<sub>v</sub> isotope measurements over two continental sites. These measurements provide a unique data set for understanding variability in the lower troposphere d-excess profile, and what it reveals about lower troposphere moisture processing on relatively small regional scales. H<sub>2</sub>O<sub>v</sub> stable isotope VP measurements were conducted in the lower troposphere during four flights around the Washington, D.C.-Baltimore, MD area in February 2016 and during five flights around the Indianapolis, IN metropolitan area in March 2016. We compare and contrast observations of the unique vertical structure of δD, δ<sup>18</sup>O, and d-excess from three representative case study days. The case studies provide information about meteorological conditions that produce H<sub>2</sub>O<sub>v</sub> isotopic VP profiles consistent withconsistent with a memory of indicative of past Rayleigh distillationdehydration of the airmass -theory and those where other processes must explain the observations. The case study

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distillation<u>dehydration of the airmass</u> <u>theory</u> and those where other processes must explain the observations. The case study observations reveal d-excess features unique to <u>stratocumulus eloudcloud droplet or rain droplet</u> evaporation and show the influence of synoptic weather patterns and urban versus rural differences on BL development. Interpretations of case study VPs are supported with observations from the remaining flight days in Washington, D.C.-Baltimore and Indianapolis areas.

## 2 Methods

#### 20 **2.1 Study sites**

Flights were conducted around the Washington, D.C.-Baltimore, MD<u>area</u> in February 2016 and around Indianapolis in March 2016. Washington, D.C.-Baltimore- is a metropolitan area of 9.8 million residents that includes the District of Columbia and <u>encompassing encompasses</u> parts of Maryland, Virginia, West Virginia, and Pennsylvania (U.S. Census Bureau, 2018). The Appalachian Mountains lie to the west of Washington, D.C.-Baltimore, and the Chesapeake Bay and the Atlantic Ocean lie to the east side of Washington, D.C.-Baltimore. By contrast, Indianapolis has a population of 2.0 million and is relatively isolated from other metropolitan areas by agricultural fields (U.S. Census Bureau, 2018). The closest large body of water to Indianapolis is Lake Michigan, over 200 km to the north.

### 2.2 Instrumentation

## 2.2.1 Airborne Laboratory for Atmospheric Research (ALAR)

- 30 The Purdue Airborne Laboratory for Atmospheric Research (ALAR) is a modified twin-engine Beacheraft-Beechcraft Duchess aircraft. The ALAR's two rear passenger seats have been removed to make room for scientific instrumentation. Ambient air at the nose of the aircraft is pulled through a forward-facing unheated 5-cm diameter PFA Teflon inlet called the "main manifold" at a flow rate of 1840 L min<sup>-1</sup> using a blower installed at the rear of the aircraft. Residence time in the main manifold is ≤0.1 second. Instruments sample from the main manifold with individual Swagelok "T" connections and Teflon sampling lines.
  35 The Purdue ALAR is equipped with a global positioning and internal navigation system (GPS/INS) for 50 Hz geopositional
- measurements and a Best Air Turbulence (BAT) probe for 50 Hz three-dimensional winds and pressure measurements (Crawford

and Dobosy, 1992; Garman et al., 2006; 2008). Temperature measurements are made with a microbead thermistor installed in the center pressure port of the BAT probe (Garman, 2009). Although not the focus of this study, measurements of carbon dioxide, methane, and  $H_2O_v$  mole fraction were made with a Picarro G2301-m cavity ringdown spectrometer. The Picarro data frequency was 0.5 Hz and the flow rate was 850 sccm. This system provides an independent evaluation of  $H_2O_v$  mole fraction measurements be also be a fully dependent evaluation of  $H_2O_v$  mole fraction measurements in the flow rate was 850 sccm.

5 by the isotope analyser described in the next section. A full description of the ALAR instrumentation suite has been provided by Salmon et al. (2017).

#### 2.2.2 Water vapor mixing ratio and stable isotope measurements

 $H_2O_v$ , δD and δ<sup>18</sup>O measurements (1 Hz) were made with a Los Gatos Research, Inc. (LGR) Triple Water Vapor Isotope Analyzer (TWVIA; model: 911-0034). The TWVIA was configured as a rack-mount, extended-range model, operating with an internal cell pressure of 80 Torr, and is suggested by the manufacturer for isotopic measurements over the  $H_2O_v$  mole fraction range from 4,000 – 60,000 ppmv. The analyser-analyzer can make measurements at  $H_2O_v$  mole fractions below 4,000 ppmv, but the instrument precision worsens (discussed below). The TWVIA sampled ambient air from the main manifold at a flow rate of 500 sccm using the analyzer's internal pump. Cavity pressure and temperature were observed to vary (1σ) by ±0.02 Torr and ±0.08°C, respectively, over a vertical profile descent on average, which is within the operating specification given by the

15 <u>manufacturer</u>. Measurements of  $H_2O_v$ ,  $\delta D$ , and  $\delta^{18}O$  were identically lag adjusted for the sample residence time (average: 8 s) to match geopositional and meteorological measurements. Depending on the ambient air temperature, the cabin of the aircraft was heated to prevent condensation inside tubing and for the comfort of the pilot and mission scientist.

 $H_2O_v$  mole fractions reported by the LGR TWVIA and the Picarro instrument were calibrated on the ground (not in flight) throughout the campaign (on 7 and 17 March 2016) using a LI-COR dewpoint generator (model: LI-610) over the  $H_2O_v$  mole fraction range from 7,000 – 12,000 ppmv. This  $H_2O_v$  mole fraction range corresponds to saturation vapor pressures for temperatures ranging from approximately 3°C – 10°C. The LGR TWVIA (and Picarro)  $H_2O_v$  mole fraction calibration curve slope, y-intercept, and R<sup>2</sup> value are 0.9845 (0.94), -280 ppmv (-200 ppmv), and 0.99978 (0.99895), respectively. Figure A1 shows that t<u>T</u>he calibrated  $H_2O_v$  mole fractions from the Picarro and LGR analysers were consistent in flight (Fig. S1). LGR  $H_2O_v$  measurements are lowpass filtered relative to the Picarro measurements due to the longer LGR residence time.

- 25 The LGR TWVIA isotopic measurements were calibrated in the lab for  $H_2O_v$  concentration dependence before and after the field campaign using an LGR Water Vapor Isotope Standard Source (WVISS; model: 908-0004-9003) with five standards ranging in isotope enrichment from -39.9% to -573.7% in  $\delta D$  and -8.7% to -76.2% in  $\delta^{18}O$  (Table B1S2). Neither the Teflon sampling line between the TWVISS and TWVIA, nor the TWVIA inlet were heated. The range in the standards' ô values brackets the range of  $\delta$  values measured during the campaign. The concentration dependence was characterized over the H<sub>2</sub>O<sub>v</sub> mole fraction range from 550 ppmv – 14,000 ppmv, which corresponds to the lowest  $H_2O_v$  mole fraction the WVISS could consistently emit 30 produce (Appendix BSection S2) and the highest  $H_2O_v$  mole fraction observed during the research flights. The TWVIA's  $H_2O_v$ concentration dependence was monitored between January 2016 and June 2017, with no appreciable instrument drift observed.  $H_2O_v$  concentration-dependence calibration and residual curves are provided in Fig. S2.B1 ( $\delta^{18}O$ ) and B2Fig. -S2.2 ( $\delta D$ ), along with a discussion of the non-linear calibration curve line fitting (Appendix BSection S2). There was no need for an additional correction to normalize to the VSMOW-SLAP (Vienna Standard Mean Ocean Water - Standard Light Antarctic Precipitation) 35 scale (discussed in Appendix BSection S2; Fig. S2.3). Discussion of the instrument precision and calibration uncertainties are provided in Appendix CSection S3. Instrument precision and concentration dependence calibration uncertainties are summed in quadrature to yield the total uncertainty in  $\delta D$ ,  $\delta^{18}O$ , and d excess. Figure 1 shows the total uncertainty in  $\delta D$ ,  $\delta^{18}O$ , and d excess versus H<sub>2</sub>O<sub>v</sub>. Total uUncertainties, the quadrature sums of instrument precision and calibration uncertainties, increase as H<sub>2</sub>O<sub>v</sub>
  - 5

mole fraction decreases below 4,000 ppmv (Fig. 1). Flight measurements of  $\delta D$ ,  $\delta^{18}O$ , and d-excess reported here are smoothed using a 20-second moving average which corresponds to the time required for the TWVIA-reported  $\delta$  values to stabilize after a change in the sample's H<sub>2</sub>O<sub>v</sub> mole fraction or isotopic signature (Appendix-Section S3C).



5 Figure 1: Total uncertainty of  $\delta D$ ,  $\delta^{18}O$ , and d-excess over the range of H<sub>2</sub>O<sub>v</sub> mole fractions observed in flight.

## 2.3 Flight design

Nine daytime research flights were conducted around Washington, D.C.-Baltimore and Indianapolis in February and March of 2016 (Table 1; Fig. 2). Flight paths were designed to maximize the number of vertical profiles (VPs) conducted while
also characterizing upwind/downwind gradients in H<sub>2</sub>O<sub>v</sub> isotopic signature. VPs were sometimes conducted in a spiral pattern to limit the horizontal spatial coverage of the measurements, while other VPs were conducted in a sawtooth pattern ("porpoising"; Gerber et al., 2013) between the BL and FT. The aircraft flew up to ~1600 m above sea level (msl) on average during the VPs. Only data collected on the descents of the VPs, when sampled air transitions from relatively dry to relatively humid, are presented here to minimize the potential influence of memory effects. However, similar features were observed on the ascents and descents.

- 15 The number of VPs (Table 1) conducted on each flight was limited by air traffic and restricted air space (which was worse for the Washington, D.C.-Baltimore study site), cloud cover, and available flight time. The research aircraft typically does not fly through clouds during experimental flights as the BAT probe has electronics exposed to air. Flights included other maneuvers, such as transects conducted upwind, intersecting, and downwind of the urban centers, the interpretation of which is beyond the scope of this paper (Fig. 2).
- 20 Figure 2 shows the flight paths conducted around Washington, D.C. Baltimore and Indianapolis. Specific flight dates and times are provided in Table 1. We focus on three particular daytime experiments conducted around Indianapolis as representative case studies <u>based on their unique features and the consistency and height of the vertical profiles (Table 1).</u> H<sub>2</sub>O<sub>v</sub>-isotope measurements on March 6 (RAY) appear consistent with Rayleigh distillation theory, the observations on March 4 (STC) may reflect moisture processing in a stratocumulus topped-BL, and the March 18 (DBL) observations may reveal differences in urban
- 25 versus rural BL development and the influence of changing synoptic conditions. Conclusions about the processes influencing the

isotopic features observed during the case studies are supported by measurements from the remaining Washington, D.C. and Indianapolis flights (Table 1).



Figure 2: Flight paths conducted around the (a) Washington, D.C.-Baltimore and (b) Indianapolis study sites for the research flights (RF) dates listed in Table 1. Case study flight paths are indicated with solid lines, while all other flight paths are indicated with dotted lines.

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Table 1: Flight log listing flight date, research flight\_(RF) and case study codes used in this manuscript, flight time (local time, LT), number of vertical profiles conducted, and the observed range of potential temperature ( $\theta$ ) and ambient temperature (T) during the flights.

Flight Date	Research	Flight Time	Vertical	θ (°C)	T (°C)
(2016)	Flight Code <u>-</u>	(LT)	Profiles		
	Case Study				
12 February	RF01FEB12	11:45 - 17:30	1	-3.0 - 6.9	-8.9 - 9.3
17 February	FEB17RF02	11:40 - 18:15	1†	6.4 - 12.5	1.6 - 10.4
18 February	FEB18RF03	12:10 - 17:25	1	-0.4 - 17.7	-6.1 - 10.7
19 February	FEB19RF04	11:55 - 17:10	1	0.9 - 14.6	-0.6 - 10.1
4 March	<del>RF05</del> MAR04 -	13:55 - 16:30	5	3.5 - 15.4	-2.8 - 4.2
	<u>STC</u>				
6 March	<u>RF06MAR06 -</u>	12:55 - 15:25	4	9.6 - 21.1	4.4 - 11.6
	CLR				
7 March	RF07MAR07	14:10 - 16:45	6	15.8 - 26.5	10.2 - 18.7
17 March	RF08MAR17	12:15 - 15:00	2	13.6 - 17.5	1.6 – 17.3
18 March	<del>RF09</del> MAR18 -	11:40 - 14:20	$4^{\dagger}$	7.8 - 17.8	0.2 - 10.7
	DBL				

10 \*The supporting research flight days share similarities with the indicated case study, but some caveats exist (Discussion 4.4).

<sup>†</sup>Measurements of meteorological variables are completely or partially unavailable during one of the vertical profiles due to temporary failure of winds measurement system.

- Nine daytime research flights were conducted around Washington, D.C. Baltimore and Indianapolis in February and March of 2016 (Table 1; Fig. 2). Flight paths were designed to maximize the number of VPs conducted while also characterizing upwind/downwind gradients in H<sub>2</sub>O<sub>x</sub>-isotopic signature. VPs were sometimes conducted in a spiral pattern to limit the horizontal spatial coverage of the measurements, while other VPs were conducted in a sawtooth pattern ("porpoising"; Gerber et al., 2013) between the BL and FT (b and c panels of Figures #, #, and #) when the research aircraft travelled between the West Lafayette, IN, Purdue airport and the Indianapolis study site. Figure 3 shows examples of these two types of VPs conducted during the case study
- 20 flights. The aircraft flew up to ~1600 m above sea level (msl) on average during the VPs. Only data collected on the descents of the VPs, when sampled air transitions from relatively dry to relatively humid, are presented here to minimize the potential influence of memory effects. However, similar features were observed on the ascents and descents. The number of VPs (Table 1) conducted on each flight was limited by air traffic and restricted air space (which was worse for the Washington, D.C. Baltimore study site), cloud cover, and available flight time. The research aircraft typically does not fly through clouds during experimental flights.
- 25 Flights included other maneuvers, such as transects conducted upwind, intersecting, and downwind of the urban centers, the interpretation of which is beyond the scope of this paper (Fig. 32).



Figure 3: The (a) RAY, (b) STC, and (c) DBL flight path and altitude time series with vertical profiles (VP) highlighted.
 Flight paths overlay the study site's cloud cover captured at approximately 12:30 local time by the Terra MODIS satellite (<u>https://worldview.earthdata.nasa.gov/</u>). The teal outline indicates the Indianapolis city boundaries. The West Lafayette, IN, Purdue airport is indicated by the airplane marker.

#### 2.4 Atmospheric layer identification

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We classify regions of the atmosphere into the boundary layer (BL), inversion layer (INV), and free troposphere (FT), to compare and contrast features observed in  $\delta D$ ,  $\delta^{18}O$ , and d excess isotopic signatures features during the research flights. The altitude at the base of the INV ( $z_{INV}$ ) is defined as the lowest altitude at which the change in potential temperature ( $\theta$ ) with altitude exceeds 0.5 K for a 10 m change in altitude ( $d\theta/dz > 0.05$  K m<sup>-1</sup>). Rates of  $d\theta/dz > 0.05$  K m<sup>-1</sup> were commonly observed within the INV during the research flights. The altitude at the base of the FT ( $z_{FT}$ ) is defined as the altitude above  $z_{INV}$  at which  $d\theta/dz$  transitions to  $z_{O} 0.5$  K m<sup>-1</sup> at the base of the FT ( $z_{FT}$ ) is defined as the altitude above  $z_{INV}$  at which  $d\theta/dz$  transitions to  $z_{O} 0.5$  K m<sup>-1</sup>.

- to <0.05 K m<sup>-1</sup>. A recent evaluation of methods for determining boundary layer height from aircraft measurements indicate the potential temperature gradient approach is most reliable (Dai et al., 2014). However, if layers are not definable using the  $d\theta/dz >$  or < 0.05 K m<sup>-1</sup> criterion, the secondary criterion of  $|d(H_2O_v)|/dz > 20$  ppmv m<sup>-1</sup> and  $|d(H_2O_v)|/dz < 20$  ppmv m<sup>-1</sup> is used to define  $z_{INV}$  and  $z_{FT}$ , respectively. These threshold values are appropriate for our wintertime, mid-latitude observations, but may not be universally appropriate in different locations or seasons. If neither criterion is met, the profiles of  $\theta$ ,  $d\theta/dz$ ,  $H_2O_v$ , and  $d(H_2O_v)/dz$
- 20 are collectively considered, and  $z_{INV}$  is visually defined as the point at which  $H_2O_v$  and  $\theta$  begin decreasing and increasing, respectively. Similarly,  $z_{FT}$  is visually determined as the altitude at which the rate of change of  $H_2O_v$  and  $\theta$  with altitude begins to decrease.

## 2.5 Isotope theory

There are three common ways many processes that can influence the isotopic composition of atmospheric H<sub>2</sub>O<sub>v</sub>; here we 25 examine two common ways can change: condensation and mixing. We use the Rayleigh distillation model to represent condensation as of an ascending air parcel and we use a simple two-member mixing model to represent atmospheric mixing. is dehydrated as it cools with altitude, mixing of different air masses, and We calculate the influence of these processes using R notation, but present the results by converting to delta notation ( $\delta = R_{sample}/R_{standard} - 1$ ) using VSMOW as the standard.-cloud formation and evaporation. We show here how each of these processes is <u>are</u> expected to change the isotopic signatures of atmospheric H<sub>2</sub>O<sub>x</sub> as H<sub>2</sub>O<sub>x</sub>-mixing ratios change.

As air is dehydrated, for example during ascent, the heavier isotopologues are preferentially condensed first. The Rayleigh 5 distillation model describes the effects of equilibrium fractionation on the isotopic composition of a dehydrating air parcel. for example during ascent (Dansgaard, 1964). Condensate that is formed as an ascending air parcel expands and cools is isotopically enriched in the heavier isotope relative to the vapor and in the open form of the Rayleigh model is assumed to be immediately removed from the system. The isotopic composition of the parcel as predicted by Rayleigh theory is given by the (eq. (1)).

$$R_{Ray} = R_o \left(\frac{H_2 O_v}{H_2 O_{v_o}}\right)^{\alpha_e - 1} \tag{1}$$

Here R<sub>o</sub> and R<sub>Ray</sub> are the heavy to light isotopologue ratios (<sup>HDO</sup>/<sub>H2O</sub> or <sup>H218</sup>O<sub>4PO</sub>/<sub>H2O</sub>) of the parcel prior to the ascent and at any point throughout the ascent, respectively. The remaining fraction of H<sub>2</sub>O<sub>v</sub> left in the ascending parcel relative to initial conditions is given by <sup>H2Ov</sup>/<sub>H2Ovo</sub>. We determined the initial R<sub>o</sub> and H<sub>2</sub>O<sub>vo</sub> input values for each day from the average BL values measured along the VP descents. The temperature-dependent equilibrium fractionation factor, α<sub>e</sub>, for each isotopologue is calculated for the temperature corresponding toof the air parcel's lifting condensation level (LCL) altitude using Horita and Wesolowski (1994) for
LCL temperatures greater than 0°C and Ellehøj et al. (2013) for LCL temperatures less than 0°C. The LCL is the height at which an air parcel would become saturated if lifted adiabaticallyalong a dry adiabatic ascent and is often used as an estimate of cloud base height (Romps, 2017). The VP observations show that ambient temperatures vary with altitude along the vertical profiles. However, Rayleigh distillation curves calculated with α<sub>e</sub> values defined by the varying ambient temperatures measured along the vertical profiles are nearly identical to Rayleigh curves calculated with a single LCL-defined α<sub>e</sub> value (Fig.ure S41).

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The mixing of two air parcels (*A* and *B*) results in a heavy-to-light isotopologue ratio of an air parcel,  $R_{mix}$ , given by eq. (2) using HDO<sub>v</sub> and H<sub>2</sub>O<sub>v</sub>-(H<sub>2</sub><sup>16</sup>O<sub>v</sub>)-as an example,  $\frac{HDO}{H_2O_{mix}}$ . In eq. (2),  $R_{mix}$  is the ratio of the weighted average of the heavy isotopologue to the weighted-average of the light isotopologue. The fraction of air parcel *A*,  $f_A$ , and air parcel *B*,  $f_B$ , sum to unity. The mixture's H<sub>2</sub>O<sub>v</sub> mole fraction is simply the weighted average of each parcels' individual H<sub>2</sub>O<sub>v</sub> mole fraction. H<sub>2</sub><sup>18</sup>O<sub>v</sub> can replace HDO<sub>v</sub> in eq. (2).

25 
$$R_{mix} = \left(\frac{HDO_v}{H_2O_v}\right)_{mix} = \frac{f_A[HDO_v]_A + f_B[HDO_v]_B}{f_A[H_2O_v]_A + f_B[H_2O_v]_B}$$
(2)

We consider mixing processes in our observations by choosing two regions of the VPs as potential end-members, for example, theBL and the FT, and use the observed  $H_2O_v$  and isotopic ratios to define end-member values.The isotopic influence ofcloud evaporation on the surrounding water vapour is complicated and depends on the mass of water in the vapour and liquidphases (Noone, 2012). Here, we compare two simplified approaches, described in detail below, but outlined here. In the first

- 30 approach, we use the model from Worden et al. (2007) to describe cloud evaporation into a completely dry atmosphere. In the second approach using the model from Stewart (1975), we assume that cloud evaporation happens in two distinct regions of the inversion layer. First, cloud liquid is formed in equilibrium with atmospheric vapour at the LCL temperature. Next, that liquid is partially evaporated in the lower portion of the INV, changing its isotopic composition. Finally, that partially evaporated cloud droplet is moved to the upper portion of the INV where it evaporates completely.
- 35 The Worden et al. (2007) model describes the isotopic signature of an air parcel that is influenced by the evaporation of cloud droplets using a modified Rayleigh model shown in eq. (3).

<del>ð8</del>	$\frac{1}{a} \left[ a \left( \frac{1 - f_{evap}/a_k}{a_k} \right) \right]$	_ 1]	(	3)
∂H <sub>2</sub> 0 <sub>₽</sub>	$H_2 O_{\psi} \begin{bmatrix} u_e \\ 1 - f_{evap} \end{bmatrix}$	1]		57

Here  $\frac{\partial \delta}{\partial q}$  represents the change in the air parcel's  $\delta$  signature with change in  $H_2O_{\psi}$  concentration. The fraction of the cloud droplet that has evaporated is given by  $f_{evap}$ . The kinetic fractionation coefficient is given by  $a_k$ , and is calculated according to Merlivat and Jouzel (1979). This is a simplified model that assumes the relative humidity at the surface of the evaporating cloud droplet is 0% (Worden et al., 2007; Noone, 2012).

The isotopic signature of a cloud droplet that undergoes partial evaporation and kinetic isotope fractionation within the INV is calculated according to eqs. (4) through (6), from Stewart (1975).

$$R_{cloud} = \gamma R_{vap} + (R_{cloud,o} - \gamma R_{vap}) f_{cloud}^{\beta}$$

$$\gamma = \frac{\alpha_e h}{1 - \alpha_e \left(\frac{p}{D_t}\right)^n (1-h)}$$
(5)

$$10 \quad \beta = \frac{1 - \alpha_{e} \left(\frac{D}{D_{t}}\right)^{n} (1-h)}{\left(\frac{D}{D_{t}}\right)^{n} (1-h)} \tag{6}$$

Here  $R_{cloud}$  is the isotopic ratio of the remaining cloud droplet,  $R_{cloud,o}$  is the initial isotopic ratio of the cloud droplet,  $R_{vap}$  is the isotopic ratio of the atmospheric vapor,  $f_{cloud}$  is the fraction of the cloud droplet remaining, and h is relative humidity. The ratio of the diffusivity of light water to heavy water,  $\frac{D}{D_t}$ , is 1.02512 for H<sub>2</sub><sup>46</sup>O:HDO and 1.02849 for H<sub>2</sub><sup>46</sup>O:H<sub>2</sub><sup>48</sup>O (Merlivat, 1978). The sealing constant, n, is 0.58 and determines the magnitude of kinetic isotope fractionation (Stewart, 1975).

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## 3 Airborne campaign observations of H<sub>2</sub>O<sub>v</sub> isotopologuesd-excess in different layers of the lower troposphere

<u>The campaign-wide observations show that values of  $\delta D$  and  $\delta^{18}O$  decrease as  $H_2O_v$  mole fractions approach zero (Fig. 3). This relationship results from preferential removal of the heavier isotopologues during condensation processes, which becomes more pronounced at colder temperatures, and is consistent with previous airborne and high-altitude measurements of</u>

20 <u>H<sub>2</sub>O<sub>v</sub> stable isotopologues (Bailey et al 2013; Galewsky et al 2007; He and Smith, 1999; Noone et al 2013; Samuels-Crow et al.,</u> 2014; Sodemann et al 2017). Later we will return to examine individual profiles of  $\delta D$  and  $\delta^{18}O$  on case study days, but here we focus on identifying common patterns in d-excess signatures observed in the BL, INV, and FT that can be used to understand processes controlling moisture in the lower troposphere.



Figure 3: Measurements of (a)  $\delta D$  and (b)  $\delta^{18}O$  made during vertical profile (VP) descents on the nine research flights (Table 1).

- The case study days presented above were chosen for their distinct isotopic features and because several VPs were
  5 conducted each day. However, they only represent 30% of the research flight days (Table 1). Figure 7a c shows tThe
  Washington, D.C.-Baltimore and Indianapolis VP absolute d-excess observations within the FT, INV, and BL, are presented as a function of H<sub>2</sub>O<sub>v</sub> mole fraction in Fig. 4a-c, respectively. Figures 7d f show the measured d excess relative to Rayleigh distillation theory, i.e. Rayleigh predicted d excess has been subtracted from the observations. Overall, BL d excess
  observations at the Indianapolis and Washington, D.C. Baltimore study sites are relatively consistent with Rayleigh theory when
- 10 <u>using observations on each day as the initial conditions (Fig. 7c).</u> Generally, the air became drier and the d-excess signature exhibited greater variability with increasing altitude (Fig. 4). The greatest departures from Rayleigh theoryvariability -wasere most commonly-observed in the INV (Fig. 47b), where the d-excess signature deviated both positively and negatively from the global average precipitation d-excess value of 10‰ (which is provided in Fig. 4 for reference only). Observations of d excess in the INV at the Washington, D.C. Baltimore study site in particular were significantly more negative than Rayleigh predictions,
- 15 <u>by up to 80 % (Fig. 7b)</u>. The FT showed low  $H_2O_v$  mole fractions as well as large positive d-excess values (Fig. 4a), as which are predicted by Rayleigh distillation theory for such very low H2O<sub>v</sub> mole fractions (Bony et al., 2008).

<u>The Very lowest  $H_2O_{v}$ -mole fractions of the campaign were commonly observed in the FT, as were large, positive d excess</u> values (Fig. 4a). Large positive d excess values are predicted by Rayleigh distillation theory at low  $H_2O_{v}$ -mole fractions. However, observed FT d excess values are even more positive than the Rayleigh predictions in Fig. 7a at the lowest  $H_2O_{v}$ -mole

20 <u>fractions.</u> High FT d-excess signatures have been reported by other studies, which hypothesize that veryextremely dry, depleted <u>FT-air masses in the mid-to-upper troposphere earrying</u>with large positive d-excess signatures, mix downward towards flightlevel altitudes in the lower FT (Bony et al., 2008; Samuels-Crow et al., 2014; Sodemann et al., 2017). These FT air masses likely originated from another source region and possibly underwent multiple condensation cycles to achieve such isotopic depletion prior to mixing with more humid air across the INV. Thus, FT d-excess values likely act as a record of the condensation and isotopic depletion history of a transported air parcel.

- The 7 March 2016 flight-day in Indianapolis (MARRF07) revealsstands out as an unusual set of meteorological
  conditions because, unlike all other flight days of the campaign, the H<sub>2</sub>O<sub>v</sub> mole fraction increased with altitude in the lower troposphere-increased in H<sub>2</sub>O<sub>v</sub>-mole fraction with altitude (Fig. 47a-c; gold trace), and it was the most humid day of the entire campaign (Fig. 3; gold trace). A warm, southerly front moved into the Indianapolis study area on this day, and rain preceded the flight measurementsobservations. -The relatively high H<sub>2</sub>O<sub>v</sub> mole fractions in the INV and FT likely reflect residual humidity from the storm. -Overall, the MARRF07 VP observations do not exhibit distinctive isotopic features, and d-excess features,
- 10 generally aligns with the global average d excess value of precipitation varies around Rayleigh predictions for the BL, i.e. -not deviating much from 10 permil-‰ in the INV<sub>7</sub> and FT (Fig. 74). These unique data likely reflect the imprint of the precipitation's isotopic signature on the surrounding vapor via isotopic equilibrium.



Figure 47: Observations of Rayleigh-subtracted ddD-excesss observations in the (a) free troposphere (FT), (b) inversion (INV), and (c) boundary layer (BL) during all the VP descents. Colors indicate each research flight. BL, INV, and FT observations are defined by the rate of change of atmospheric variables, and not defined by H<sub>2</sub>O<sub>\*</sub> mole fraction. The dashed lines correspond to the global average precipitation d-excess value of 10% for reference only.

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 Most measurements of H<sub>2</sub>O<sub>v</sub>-mole fraction in the FT were below 4000 ppmv; the corresponding FT measurements of <sup>548</sup>O, δD, and d excess are plotted as a function of H<sub>2</sub>O<sub>v</sub>-in Fig. S5. While each Washington, D.C. Baltimore VP (Table 1) extended into the FT, H<sub>2</sub>O<sub>v</sub>-mole fractions were often lower than the humidity range over which the TWVIA was calibrated
 (lower limit: 550 ppmv; Appendix B). Contributing to the scarcity of Washington, D.C. Baltimore VP measurements is the fact that only one VP was conducted per flight, mainly as a result of congested and restricted air spaces near the capitol. Values of δD range from 30 to 60‰, δ<sup>18</sup>O ranges from 200 to 400‰, and d excess can be close to 10‰, but also increase up to 100‰ below 1000 ppmv.

### 4 DiscussionCase Studies

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We focus on three particular Indianapolis flights as representative case studies based on their distinct features and vertical profiles over large altitude ranges (Table 1).  $H_2O_v$  isotope measurements conducted within and above the cloudless, well-mixed boundary layer on March 6 (CLR) represent the meteorologically-simplest observations of the entire campaign, as no clouds,

5 precipitation, or shifting synoptic conditions were observed. Isotopic observations on March 4 (STC) may reflect moisture processing in a stratocumulus topped-BL, and the March 18 (DBL) observations may reveal differences in urban versus rural BL development and the influence of changing synoptic conditions.

#### 34.1 Rayleigh-consistent observations Clear sky observations of a well-mixed boundary layer (RAYCLR)

Four VPs consistent with Rayleigh distillation theory were conducted on 6 March 2016 ("CLR" case study) aroundin Indianapolis during clear sky conditions. (Fig. 5a; weather map is presented in Fig. S5.1). Cloud-top height estimated from the Terra MODIS satellite retrievals (https://worldview.earthdata.nasa.gov/) indicate that the sparse cloud cover shown in Fig. 5a corresponds to higher altitude (>4800 m) clouds. The CLR measurements were made below 1400 m above sea level (Fig. 5b), and as a result, were likely not impacted by higher cloud processes. In terms of meteorology, the CLR case study is the simplest flight
 day of the airborne campaign. It is a useful case study to examine isotopic signatures across the BL, INV, and FT without the



Figure 5. (a) CLR flight path overlaying the study site's cloud cover (https://worldview.earthdata.nasa.gov/) captured at 12:30 local time. The teal outline indicates the Indianapolis city boundaries. The West Lafayette, IN, Purdue airport is indicated by the airplane marker. (b) CLR case study altitude time series.

Observations of  $\delta D$ ,  $\delta^{18}O$ , and d excess measured along the <u>CLR</u> VP descents are plotted as a function of H<sub>2</sub>O<sub>v</sub>-mole fraction in Fig. 4a c, respectively, along with predictions from Rayleigh distillation theory. The mixing lines in Fig. 4 show an air parcel's isotopic signature if varying proportions of BL and FT air are mixed (Methods 2.5). For the most part, the <u>shape of the</u> Rayleigh predicated  $\delta^{18}O$ ,  $\delta D$ , and d excess values <u>curves</u>along the <u>are similar to the observations</u> four VPs are consistent with the observations from the BL up to the top of the INV (Fig. 4a c). However, <u>observed  $\delta D$ ,  $\delta^{18}O$ , and d excess observations values</u> <u>near</u> at the interface of the INV and FT exhibit a hyperbolic shape (Fig. 4<u>a</u>c), which is associated with mixing between distinct air parcels (Noone, 2012).

5



STC. Shading around the  $\delta^{18}$ O,  $\delta$ D, and d-excess VP measurements define measurement uncertainty. Measurements in the boundary layer (BL), inversion layer (INV), and free troposphere (FT) are indicated for reference. The inversion layers,

- VP2 observations are presented as a representative example of RAY because VP2 was conducted approximately midway through the flight, it covers the largest vertical range relative to the remaining VPs, and it was conducted in a spiral formation to minimize the horizontal spatial extent over which the measurements were made (Fig. 3a). The VP measurements made during CLR indicate a well-mixed BL and FT, and wind speed and wind direction were relatively constant from the BL to FT, indicating that BL, INV, and FT air parcels shared recent advection histories.<sup>2</sup> VP2 measurements in the BL, from 380 m – 780 m above ground level, of δ<sup>18</sup>O, δD, and d-excess are relatively constant with altitude within the BL (surface to z<sub>INV</sub>), varying by 1.2‰, 15.3‰, and
- 10 10.9‰, respectively (Fig. 6a). VP2 observations are presented as a representative example of RAYCLR because VP2 was conducted approximately midway through the flight (Fig. 5b), it covers the largest vertical range-relative to the remaining VPs, and it was conducted in a spiral formation to minimize the horizontal spatial extent over which the measurements were made (Fig. 35aa). The ambient temperature profile approximately follows the dry adiabatic lapse rate to the top of the BL (Fig. 5a6a), above which, temperature increases within the INV layer. VP2-H<sub>2</sub>O<sub>v</sub> mole fraction decreases by 5095 ppmvrapidly in the INV between
- 15  $z_{INV}$  and  $z_{FT}$  before becoming relatively stable constant in the FT. The VP2 INV  $\delta^{18}O$  and  $\delta D$  values track the H<sub>2</sub>O<sub>v</sub> profile within the INV, decreasing by 30.8‰ and 193.2‰, respectively, as the heavier isotopologues preferentially condense. Observed d-excess values in the INV initially decrease slightly with altitude, and then increase, varying overall by 66.6‰. Just above the INV H<sub>2</sub>O<sub>v</sub> mole fractions near a minimum, and  $\delta D$  increases while  $\delta^{48}O$  decreases, causing FT d excess values to increase rapidly. Above ~1100 m in the FT, the VP2-H<sub>2</sub>O<sub>v</sub>,  $\delta^{18}O$ ,  $\delta D$ , and d-excess signatures are relatively constant with altitude. Wind speed along the
- 20 RAY VP2 ranged from 4.3 m s<sup>-1</sup> to 10.3 m s<sup>-1</sup>, and wind direction only varied by 60° were relatively constant from the BL to FT<sub>2</sub> indicating that BL, INV, and FT air parcels shared recent advection histories. In terms of meteorology, the CLR case study is the simplest flight day of the airborne campaign. It is a useful case study to examine isotopic signature across the BL, INV, and FT without the influence of complex atmospheric circulations or vapor condensate interactions from clouds or precipitation. Cloud top height estimated from the Terra MODIS satellite retrievals (<u>https://worldview.earthdata.nasa.gov/</u>) indicate that
- 25 the sparse cloud cover shown in Fig. 3a corresponds to higher altitude (>4800 m) clouds. The RAY measurements were made below 1400 m above ground level (Fig. 3a), and as a result, were likely not impacted by cloud processes from the sparse, higher altitude clouds.

Observations on 17 March 2016 (RF08) were also consistent with Rayleigh theory (Table 1). Like RF06 (RAY), skies were clear of clouds and wind speed and wind direction was relatively constant from the BL to the FT and a nearly dry adiabatic
 lapse rate was present from the surface up to ~3 km on RF08 (Fig. S2). RF06 was chosen for the Rayleigh case study over RF08 because RF08 H<sub>2</sub>O<sub>v</sub> mole fractions covered a smaller range and only two VPs were conducted that day.



Figure 7: Comparison of vertical profile  $\delta^{18}$ O (top panel),  $\delta$ D (middle panel), and d-excess (bottom panel) measurements to Rayleigh theory (left panel) and mixing (right panel) curves for CLR (MAR06). Individual VP descents are indicated by the different-colored points. The bounds of the inversion layer (INV), indicated by grey shading, are defined by the average H<sub>2</sub>O<sub>v</sub> mole fractions observed at z<sub>INV</sub> and z<sub>FT</sub>.

<u>CLRRAY is the simplest case study day where  $\delta$  values and d-excess observations generally track the Rayleigh predictions</u> in Fig. 4a c. Observations of  $\delta D$ ,  $\delta^{18}O$ , and d-excess measured along the CLR VP descents are plotted as a function of H<sub>2</sub>O<sub>v</sub> mole fraction in Fig. 7, along with predictions from Rayleigh distillation theory (Fig. 7a-c) and different BL-FT mixing scenarios (Fig.

- 5 7d-f, Methods 2.5). The measured δ values more closely match Rayleigh curves than the mixing lines from the BL through the INV (Fig. 7). However, -positive deviations relative to from Rayleigh d-excess exist in the upper INV toand lower FT (Fig. #7c)<sub>x</sub>. and-Wwe hypothesize that dry, isotopically-depleted FT air parcels carrying large, positive d-excess values can mix downward into more humid air parcels of smaller d-excess values near the top of the INV (Sodemann et al., 2017). As the H<sub>2</sub>O<sub>v</sub> mole fraction approaches zero, Rayleigh-predicted d-excess approaches 7000‰ (Bony et al., 2008). Thus, dehydrated lower-altitude FT air
- 10 masses can carry a more positive d-excess signature for a given H<sub>2</sub>O<sub>v</sub> mole fraction due to vertical mixing, than is predicted by Rayleigh theory due to vertical mixing in the FT-(Fig. S5c). The d-excess signature in the upper INV and lower FT more closely follows the FT-INVBL mixing lines -(see VP1- and VP4-defined endmember mixing scenarios in Fig. 47fe), which supports this hypothesis.- Our results are consistent with Dyroff et al. (2015), who report lower troposphere δD observations over the Atlantic Ocean, and explain the vertical structure of δD at lower altitudes using Rayleigh theory, while mixing scenarios dictate the δD
- 15 profile based on -higher altitude observations-indicate mixing scenarios dictate the  $\delta D$  profile.

Despite apparent similarities between observed and Rayleigh d-excess in the BL and INV, the meteorological conditions within the BL are contrary to assumptions of the Rayleigh distillation model. Rayleigh distillation theory describes a two-phase system at saturation, however, no clouds (i.e. condensate) were observed at the flight altitudes where the CLR measurements were made. Furthermore, the BL was dry adiabatic (Fig. 6a). The agreement between the CLR observations and the Rayleigh-predicted d-excess, however, could be described by upwind condensation consistent with Rayleigh distillation theory, followed by subsequent advection to the Indianapolis study site. Given that wind direction and wind speed were observed to be relatively constant throughout the lower troposphere (Fig. 6a), air parcels in the BL, INV, and FT likely share similar trajectories. Thus, the Rayleigh-consistent CLR VP observations at lower flight altitudes likely result from upwind condensation that imprinted an isotopic signature on the air parcels that was maintained during transport, while lower FT observations point to mixing between subsiding FT and INV air. Past studies have also reported on the usefulness of H<sub>2</sub>O<sub>y</sub> stable isotopes as "imprints" of condensation, mixing, and atmospheric transport (Bailey et al 2013; Brown et al., 2008; Galewsky et al 2007; Gedzelman, 1988; He and Smith

1999; Samuel-Crow et al., 2014; Taylor, 1984).

From our measurements, we can comment on what types of atmospheric conditions are likely to produce profiles 30 consistent with Rayleigh theory, i.e. equilibrium fractionation and no vapor condensate (cloud droplet) equilibration. The RAY meteorological and isotopic observations are very similar to those made on 17 March 2018 (RF08; Table 1). These two flight days' observations provide possible criteria for when it is appropriate to use Rayleigh theory to predict isotopic signature throughout the BL and INV. These criteria include that the study area is free of clouds/precipitation (supported by the presence of a dry adiabatic lapse rate throughout the BL) and wind speed and direction is relatively constant throughout the BL, INV, and FT. The STC and

35 <u>DBL case studies, which do not follow Rayleigh theory, violate one or more these criteria.</u>

#### 34.2 Stratocumulus-topped boundary layer observations (STC)

Figure 3b shows that the <u>The</u> center and eastern portions of the Indianapolis study area were covered by stratocumulus clouds during part of the 4 March 2016 flight (RF05, "STC" case study), on which five VPs were conducted (Fig. 8). The cloud

- 5 cover map in Fig. 8a is provided to show the cloud type and extent during the afternoon of STC, however, it does not necessarily represent the cloud cover conditions throughout the 2.5 h flight. Satellite images (not shown) captured during the early afternoon of STC show that a thick cloud cover was sustained from the beginning into the middle of the flight, particularly over the city of Indianapolis, but eventually transitioned to scattered cover throughout the afternoon. This is consistent with visual observations made by the pilot and mission scientist that thick clouds persisted from the beginning of the flight until approximately 15:30 local
- 10 time. Vertical profile temperature measurements collected on this day tThesupport the presence of an earlier cloud layer-on STC is supported by the vertical profile temperature measurements collected on this day (Fig. 6b). BL air-becomes was nearly saturated at 788 m ( $z_{CB}$  for "cloud base") along VP2 on STC (Fig. 6b). The ambient temperature lapse rate is 8.8 K km<sup>-1</sup> (close to the dry adiabatic lapse rate of 9.8 K km<sup>-1</sup>) near the surface until an altitude of  $z_{CB}$ , where the lapse rate transitions to 2.8 K km<sup>-1</sup>. These observations are indicative of a stratocumulus cloud layer, which sits directly below the INV, and sustains the temperature inversion
- 15 via radiative cooling (Wood, 2012). Indeed, a sharp decrease in  $\theta$  is localized at  $z_{INV}$  (Fig. 6b). Evidence of a stratocumulus cloud layer wasere apparent on VP1 and VP2, but a clear change in lapse rate below the INV was not observed on the remaining later STC VPs, indicating the air was not saturated below the INV during the later portion of the flight, consistent with visual observations of the thinning cloud layer. The STC flight is a useful case study to investigate how a stratocumulus cloud layer can influence the vertical structure of H<sub>2</sub>O<sub>v</sub> isotopologues in the lower troposphere.



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Figure 8. (a) STC flight path overlaying the study site's cloud cover (https://worldview.earthdata.nasa.gov/) captured at 12:30 local time. The teal outline indicates the Indianapolis city boundaries. The West Lafayette, IN, Purdue airport is indicated by the airplane marker. (b) STC case study altitude time series.

<u>Measurements of isotopic and meteorological variables made along VP2 are shown in Fig. 6b. VP2 data is presented</u> because it was conducted approximately mid-flight and it was flown in a spiral formation. <del>minimizing the horizontal spatial extent</del> over which the measurements were made. Measurements of  $\delta^{18}$ O,  $\delta$ D, and d-excess within the BL varied by 3.3‰, 27.4‰, and

- 5 <u>19.1‰</u>, respectively. This is approximately double the variability in  $\delta D$ ,  $\delta^{18}O$ , and d-excess observed within the BL along the CLR VP2 (Fig. 6a). Within the INV, H<sub>2</sub>O<sub>v</sub>,  $\delta^{18}O$ , and  $\delta D$  values decrease by 1930 ppmv, 17.6‰, and 159.7‰, respectively from BL values. Similar to CLR, d-excess steadily increases in the FT on STC as H<sub>2</sub>O<sub>v</sub> mole fractions decrease. Despite general similarities in the vertical structure of the  $\delta$  values on the CLR and STC flight days, there are notable differences in the d-excess structure near the INV on the two days (Fig. 6). Contrary to the CLR INV d-excess profile (Fig. 6a), d-excess first increases with altitude within
- 10 the lower INV before decreasing to a minimum at  $z_{FT}$  on STC (Fig. 6b). Obvious anomalies in the STC d-excess signature relative to CLR are also apparent when plotted as a function of  $H_2O_v$  mole fraction (Fig. 9c and Fig. 7c, respectively).



Figure 9: Comparison of vertical profile  $\delta^{18}$ O (top panel),  $\delta$ D (middle panel), and d-excess (bottom panel) measurements to Rayleigh theory (left panel) and mixing (right panel) curves for STC (MAR04). Individual VP descents are indicated by the different-colored points. The bounds of the inversion layer (INV), indicated by grey shading, are defined by the average H<sub>2</sub>O<sub>v</sub> mole fractions observed at z<sub>INV</sub> and z<sub>FT</sub>.

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Measurements on STC reveal unique d excess features within the INV that may reflect stratocumulus cloud evaporation. Five VPs were conducted on this day.<u>Here we compare the STC observations to Rayleigh curves using the assumptions described</u> in Section 2.5. As mentioned above, the Rayleigh distillation model assumes condensation in a two phase system at saturation (Section 4.1). While the atmosphere on STC was unsaturated at most flight altitudes (except near the cloud base; Fig. 6b), we deem the comparison to open-system Rayleigh curves is a useful exercise -as we showed above in the discussion of the CLR case study. Mixing processes produce δ<sup>18</sup>O and δD value that plot above the open system Rayleigh curve and processes involving nonequilibrium liquid vapor interactions or closed system Rayleigh processes can plot below the curve. STC VP δ<sup>18</sup>O and δD

- 5 observations are relatively consistent with Rayleigh predictions (Fig. 4d), with the exception of deviateding but are more negatively from the Rayleigh curve in drier portions of the INV and FT and (Fig. 4b). The grey shading in Fig. 4 is the average range of observed INV H<sub>2</sub>O<sub>x</sub> mole fractions and does not represent the INV location for every VP on a single day. <u>STC deviations from Rayleigh theory wereare more pronounced for δD and d excess than δ<sup>18</sup>O (Fig. 49ad be f). The INV ranges (grey shading in Fig. 4) are defined by each day's average H<sub>2</sub>O<sub>x</sub> mole fractions observed at z<sub>INV</sub> and z<sub>FT2</sub> STC deviations from Rayleigh theory are</u>
- 10 more pronounced for δD and d excess than δ<sup>18</sup>O (Fig. 4e f). With the exception of VP1, the rest-Most of the VPs' δD values are more negative relative to Rayleigh in the INV and plateau in the FT which is unusual in water vapour isotope observations (Fig. 4e9b). Mixing processes can produce δ<sup>18</sup>O and δD values that plot very nearly on or above the open-system Rayleigh curve depending on the mixing endmembers (e.g. Fig. 9), whereas processes involving non-equilibrium liquid-vapor interactions or closed-system Rayleigh processes can plot below the Rayleigh curve. While the mixing scenarios presented in Fig 9 (right panel)
- 15 show remarkable agreement with the VP d-excess observations on STC, the mixing scenarios do not provide an explanation for the minimum at z<sub>FT</sub>.

The d-excess measurements along VP2 through VP5 (Fig. 9c) reveal two anomalies that yield some insight: (Fig. 4f), (1) the slight increase in d-excess in the middle of the inversion layer (particularly for VP2 and VP5) and the (2) the d-excess minimum at the INV-FT interface ( $z_{FT}$ ). From this minimum at  $z_{FT}$ , the FT d-excess signature becomes more positive with increasing altitude

20 (as the air becomes drier), and eventually transitions to being more positive than the Rayleigh curve (Fig.  $4f_{9c}$ ). <u>Unlike RAY, decreasing to a minimum at  $z_{FT_2}$ </u>

<u>The FT-d-excess values observed during VP1 represent the only FTVP measurements on STC that elosely track open</u> system Rayleigh curve and the BL FT mixing line, and do not exhibit a minimum at the INV-FT interface (Fig. 49cf). <u>VP1 was</u> conducted immediately after take-off from the Purdue airport. Sky conditions in the vicinity of the airport were clear, but a layer

- 25 of stratocumulus clouds was observed over Indianapolis. Only one VP (VP1) was conducted before the research aircraft encountered the <u>thick stratocumulus</u> cloud layer <u>over Indianapolis</u> (Fig. <u>8a3b</u>). Unlike VP2 – 5, VP1 d-excess tracks the Rayleigh <u>curveline</u> at the INV-FT interface. Slightly above  $z_{FT}$ , the VP1 H<sub>2</sub>O<sub>v</sub> mole fraction <u>began increasing increases</u> and <u>the d-excess</u> <u>switches to tracking approaches</u> the mixing line (Fig. 9c,f). We believe this VP represents conditions prior to cloud process <u>influences</u>. The FT d excess values observed during VP1 represent the only FT measurements on STC that closely track the BL-
- 30 <u>FT mixing line, and do not exhibit a minimum at the INV FT interface (Fig. 4f).</u> The differences between VP1 and the other VPs are described in more detail later.

Figure 5b shows vertical profiles of  $\theta$ , H<sub>2</sub>O<sub>v</sub>,  $\delta^{18}$ O,  $\delta$ D, RH, d excess, wind direction, wind speed, vertical wind variance (W  $\sigma^2$ ), and ambient temperature measured along VP2 conducted on STC. VP2 data is presented because it was conducted approximately mid flight and it was conducted in a spiral formation minimizing the horizontal spatial extent over which the

35 measurements were made. Measurements of δ<sup>18</sup>O, δD, and d excess within the BL varied by 3.3‰, 27.4‰, and 19.1‰, respectively. This is approximately double the variability in δD, δ<sup>18</sup>O, and d excess observed within the BL along the RAY (RF06; 6 March 2016) VP2 (Fig. 5a). Within the INV, H<sub>2</sub>O<sub>v</sub>, δ<sup>18</sup>O, and δD values decrease by 1930 ppmv, 17.6‰, and 159.7‰, respectively from BL values. Unlike RAY, d excess first increases with altitude within the lower INV before decreasing to a minimum at z<sub>FT</sub>. Similar to RAY, d excess steadily increases in the FT on STC as H<sub>2</sub>O<sub>v</sub>-mole fractions decrease.

One difference between STC and RAY is the presence of a stratocumulus cloud layer for STC (Fig. 3). Figure 5b shows that STC VP2 air becomes nearly saturated at 788 m ( $z_{CB}$  for "cloud base"). The ambient temperature lapse rate is 8.8 K km<sup>-1</sup> (close to the dry adiabatic lapse rate of 9.8 K km<sup>-1</sup>) near the surface until an altitude of  $z_{CB}$ , where the lapse rate transitions to 2.8 K km<sup>-1</sup>. These observations could indicate a stratocumulus cloud layer, which sits directly below the INV, and sustains the

5 temperature inversion via radiative cooling (Wood, 2012). Indeed, 0 decreases sharply at z<sub>INV</sub> (Fig. 5b). Indications of a stratocumulus cloud layer were apparent on VP1 and VP2, but a clear change in lapse rate below the INV was not observed on the remaining STC VPs, indicating the air was not saturated below the INV during the later portion of the flight.

Figure 3b shows that stratocumulus clouds covered most of the study area at approximately 12:30 local time on STC. The cloud cover map in Fig. 3b is provided to show the cloud type and extent during the afternoon of STC, however, it does not

- 10 necessarily represent the cloud cover conditions throughout the 2.5 h flight. Figure S3 is a GIF of satellite cloud cover images captured during the afternoon of STC that shows the cloud cover evolution over the course of the flight. Thick cloud cover was sustained into the middle of the flight, particularly over the city of Indianapolis. As a result, VPs were not conducted within the eity limits of Indianapolis. Toward the latter half of the flight the cloud cover transitioned to scattered. Four additional VPs were conducted at the end of the flight west of Indianapolis (Fig. 3b).
- 15

In contrast to RAY (Fig. 5a), STC wind speed is highly variable from the BL to the FT (Fig. 5b). Wind speed values ranged from a minimum of 0.4 m s<sup>-1</sup> in the BL to a maximum in the FT of 9.8 m s<sup>-1</sup>. A distinctive wind shear is obvious at z<sub>FT</sub>. Highly variable wind direction in the BL is a result of low BL wind speeds.

## **4.2 Stratocumulus Cloud Evaporation**

- Given that Tthe presence of the stratocumulus cloud layer is a defining meteorological characteristic of the STC case study day, we first therefore, we evaluate the potential for Ccloud and rain processes may to cause the observed d-excess anomalies from Rayleigh theory in the observed during the STC case study day. As mentioned above, two d-excess anomalies were observed within the INV, (1) a slight increase in d excess near the middle of the INV, and (2) the minimum in d-excess at the INV FT interface on VPs 2 5 (Fig. 4f). Sodemann et al. (2017) also describe VP observations of a minimum in d-excess (negative values) at the BL top, and hypothesize the negative d-excess signal results from rain droplet evaporation directly below cumulus clouds at
- 25 the top of the boundary layer. Rehydration processes, like cloud and rain droplet evaporation, have been proposed as mechanisms that could produce negative d-excess anomalies (Bolot et al., 2013-and; -Sodemann et al., 2017). HDO molecules preferentially evaporate relative to  $H_2^{18}O$  molecules (Dansgaard, 1964).-. The result is a relatively positive vapor d-excess, while the d-excess signature of the residual droplet becomes progressively more negative as it evaporates (Aemisegger et al., 2015). Therefore, if liquid droplet evaporation <del>occuring</del>occurr<del>ing</del>ed in separate atmospheric layers from start to finish, such as from the bottom of the
- 30 INV to the top of the INV, could produce a positive d-excess anomaly could be transferred to the surrounding vapor as the droplet begins evaporating starting in the lower INV. As the liquid droplet is subsequently <u>-followed by</u>-transported to the upper INV where, a negative d-excess anomaly iscould be transferred to the surrounding vapor-in the upper INV as the evaporation of the liquid droplet nears completion.-. When evaluating the effect that cloud evaporation could have on vertical profiles VPs of water vapourH<sub>2</sub>O<sub>v</sub> isotopologuesic variability, we must consider the altitudes at which clouds form and evaporate. The main difference
- 35 between the cloud and rain droplet evaporation processes is the size of the liquid droplets: <50 micron for clouds versus >50 micron for rain droplets. Evidence exists that stratocumulus cloud droplets can evaporate at different altitudes, specifically within the cloud layer and within the inversion layer (Kollias and Albrecht, 2000; de Lozar and Melledo, 2015). Furthermore, inversion layer depths are not homogeneous above a stratocumulus cloud layer. Observations show inversion layers to be thicker above downdrafts and thinner near updrafts (Kollias and Albrecht, 2000). Differing INV depths, and associated thermodynamic properties, above the

stratocumulus cloud layer on STC may in part explain why unique d-excess anomalies were observed near the middle of the INV and at  $z_{FT}$  (Fig. 6b).

When evaluating the effect that cloud evaporation could have on VPs of  $H_2O_v$  isotopologues, we must consider the altitudes at which clouds form and evaporate and the speed of the relevant processes. In the case of evaporation, this may not

- 5 <u>occur in a single layer, but throughout a range of altitudes or atmospheric layers.</u> Stratocumulus cloud tops are typically present directly below z<sub>INV</sub> (Wood, 2012). The top of the INV (z<sub>FT</sub>) is approximately the upper limit of BL mixing (Wood, 2012). Lofting of cloud droplets into the INV would cause droplet evaporation, as the INV was <u>under not</u> saturated (Fig 56b). <u>In the cloud</u> evaporation scenarios discussed below, it is assumed that the cloud droplets form at the top of the BL, and the droplets undergo evaporation within the INV. As mentioned above, two d-excess anomalies were observed within the INV, (1) a slight increase in
- 10 <u>d excess near the middle of the INV, and (2) the minimum in d excess at the INV FT interface on VP2 VP5 (Fig. 4f). Disregarding for the moment the potential isotopic enrichment or depletion cloud droplet evaporation would have on the surrounding vapor, we first This hypothesis requires that a negative d-excess anomaly is retained by an partitially evaporated ing droplet while it is transported to the upper INV. This requires us to consider two important timescales. First, t(1) The timescale at which a liquid droplet isotopically equilibrates with its surrounding vapor vapour, and second, -(2) tThe speed at which droplets move through</u>
- 15 the INV by way of vertical winds. Using the method described by Bolot et al., 2013, which is -based off of on prior work by Jouzel, 1986 (Chapter 2), we estimate that the e-folding time required for a cloud droplet with a 15 μm radius to isotopically equilibrate with the surrounding vapor is approximately two seconds, under the conditions observed in the middle of the INV, or about five seconds at the top of the INV. For our observations of vertical wind speeds, Tthe time for a droplet to move from the bottom to the top of the INV is 19 seconds, based on vertical wind speed observations. These calculations indicate that a cloud droplet located
- 20 near the bottom of the INV would isotopically equilibrate approach isotopic equilibrium with surrounding vapor before reaching the top of the INV under the observed flight conditions. Thus it is unlikely that these calculations do not strongly support the hypothesis that evaporation of stratocumulus cloud droplets observed during the flight had any relationship produced to-the observed minimum in d-excess at z<sub>FT</sub>. However, cloud droplet evaporation may explain the local maximum in d excess at the middle of the INV (Fig. 4f; Fig. 6b).
- 25 <u>HDO molecules preferentially evaporate relative to  $H_2^{18}O$  molecules (Dansgaard, 1964). For the case of a rain droplet, evaporation would result in a relative increase in d excess of the surrounding vapor and a decrease in the droplet's d excess signature (Aemisegger et al., 2015). Given the time scale of isotopic equilibrium and vertical wind speeds, it is possible that stratocumulus cloud droplet evaporation could be responsible for the slight increase in d excess within the middle of the INV (Fig. <u>4f. Fig. 6b).</u></u>
- 30 The question of what process or processes are responsible for the minimum in d-excess at the INV-FT interface  $(z_{FT})$ remains. If previous INV conditions on STC were colder, had faster vertical wind speeds, or<del>and</del> the liquid droplet radii were larger than 50 microns (e.g. rain drizzle), the timescales associated with isotopic equilibrium and transport across the INV could converge. The main difference between the cloud and rain droplet evaporation processes is the size of the liquid droplets: <50 micron for cloud versus >50 micron for rain droplets (Kollias and Albrecht, 2000). It is possible that such conditions were present prior to
- 35 when the STC measurements were made, and that the minimum in d-excess is the isotopic imprint of rain droplet evaporation under cold, unsaturated, turbulent conditions. We find discussion of these potential d-excess anomalies in the literature in reference to rain droplets evaporating below the cloud layer (Aemisegger et al., 2015; Gat, 1996), but we believe the same isotopic fingerprint would occur on vapor could occur as liquid cloud or rain drizzle droplets <u>-that are</u>lofted in drier environments- finish evaporating would occur for the vaporat the top of the INV, or at the top of the BL, as hypothesized by Sodemann et al. (2017).

The slight easterly movement of the cloud layer as shown in Fig. S3 could be interpreted as the stratocumulus cloud layer being advected out of the Indianapolis study area by westerly winds rather than evaporating away. However, Fig. 5b shows that wind speeds at the cloud layer were relatively low (<6 m s<sup>-1</sup>), and the clouds in Figure S3 transition from opaque to semi-transparent over the flight period. This information taken together suggests there was evaporation of the stratocumulus cloud layer during the

5 <u>flight.</u>

We believe our observations provide evidence that the process of cloud evaporation may be spread throughout the INV layer with the beginning and end of evaporation separated in space. Within the INV, turbulent updrafts and downdrafts are warmer and cooler than the surrounding air, respectively (Betts, 1985). Updrafts potentially carrying partially evaporated cloud droplets to the top of the INV may facilitate complete evaporation of the droplets due to the low RH and possibly wind shear-promoted

- 10 mechanical turbulence at z<sub>FT</sub> (Fig. 5b). Due to higher RH values within the middle of the INV relative to z<sub>FT</sub>, it is possible cloud droplet evaporation did not proceed to completion at these INV altitudes. Partial evaporation would impart a positive d excess signal on atmospheric vapor and act to increase RH (Aemsigger et al., 2015; Sodemann et al., 2017), both of which exhibit a local maximum at ~985 m agl (Fig. 5b).
- Other studies have reported The negative water vapor-H<sub>2</sub>O<sub>v</sub> d-excess observations at z<sub>FT</sub>and hypothesize that the could also have resulted se are from kinetic fractionation of vapor during deposition on ice crystals or snow (i.e. in ice supersaturated conditions) (Bolot et al., 2013; Casado et al., 2016; Galewsky, 2015; Lowenthal et al., 2016; Moore et al., 2016; Samuels-Crow et al., 2014; Schmidt et al., 2005).

Low d-excess values (relative to Rayleigh atcurves assuming RH = 100%) were sometimes observed within the INV layers during the Washington, D.C.-Baltimore flights (Fig. 4b). Ambient temperatures observed in flight in Washington, D.C.-Baltimore were

- 20 sometimes less than 0°C (Table 1), thus vapor deposition on ice crystals could be possible for those scenarios. It is unlikely that ice supersaturationvapor deposition on ice occurred during the -is responsible for the minimum in vapor d excess observed at  $z_{FT}$ on STC case study flight because temperatures were greater than 0°C (Table 1). However, as an example, Figure <del>DS61</del> shows that the theoretical d-excess values of STC vapor under ice supersaturated conditions (Appendix DSection S6). We reiterate that icesupersaturation is an unlikely explanation for the STC  $z_{FT}$  d-excess minimum because flight altitudes were less than 2 km, and ice
- 25 (cirrus) clouds are typically present at ~6 km. It is unlikely that ice hydrometeors falling from higher altitudes could be sustained at the top of the inversion and contribute to the low d-excess signal observed on VP2–VP5 through vapor deposition given the >0°C temperatures. There was, however, a region of high-humidity upwind (northwest) of the study site at altitudes between 3-5.5 km where ice or mixed-phase condensate could have been present (Fig. S5.2). It is possible that condensation under ice-supersaturated conditions occurred prior to the STC flight, and that the resulting isotopic imprint was maintained during transport
- 30 to Indianapolis and subsequently mixed downward via subsiding FT air (Fig. 9f). Both of these explanations-we explore for the minimum in d-excess at z<sub>FT</sub> on STC require an advected signal of a prior process (complete cloud droplet evaporation or vapor deposition on ice). This process must have happened relatively close in time to the STC flight, since the minimum in d-excess at z<sub>FT</sub> was not observed on VP1, but the anomaly was observed two hours later when VP2-VP5 were conducted. While vapor deposition in ice supersaturated conditions might not be relevant for the STC flight, low d excess (relative to Rayleigh at RH =
- 35 <u>100 %) was sometimes observed in the INV layers during the Washington, D.C. Baltimore flights (Fig. 7b). Ambient temperatures observed in flight in Washington, D.C. Baltimore were sometimes less than 0°C (Table 1), thus vapor deposition on ice crystals would be more likely for those scenarios.</u>

## 34.3 Observations of a developingDeveloping boundary layer observations (DBL)

The <u>third-final airborne</u> case study<u>, measurements wasere</u> conducted on 18 March 2016 in Indianapolis (Fig. 10), ("DBL" for developing boundary layer, RF09). <u>This case study</u> is referred to as DBL for "developing boundary layer" because <u>Mm</u>easurements on this day reveal considerable spatiotemporal variability in the vertical structure of the observed

- 5 meteorological and isotopic variables. The boundary layer height increased over the course of the flight and may reflect a combination of a residual layer from the previous day, urban vs. rural differences in BL development, and the effects of a frontal pattern moving across the Indianapolis study area. Figure 3c shows that Indianapolis was cloud free at approximately 12:30 LT and clear skies continued throughout the flight. Scattered clouds developed over Indianapolis late in the afternoon after the research flight (17:00 local time; Fig. S4). Within the BL, wind direction and wind speed were relatively constant. Wind shearing is apparent at the z<sub>FT</sub> (Fig. 6). Wind speeds increase from ~5 m s<sup>-1</sup> to ~18 m s<sup>-1</sup> between the BL and FT, which is a larger gradient
- 10 is apparent at the z<sub>FT</sub> (Fig. 6). Wind speeds increase from ~5 m s<sup>+</sup> to ~18 m s<sup>+</sup> between the BL and FT, which is a larger gradient than was observed on STC. Wind speed stabilizes within the FT (Fig. 6). Wind direction varies only by ~30<sup>-e</sup> during each of the four VPs.



Figure 10. (a) DBL flight path overlaying the study site's cloud cover (https://worldview.earthdata.nasa.gov/) captured at approximately 12:30 local time. The teal outline indicates the Indianapolis city boundaries. The West Lafayette, IN, Purdue airport is indicated by the airplane marker. (b) DBL case study altitude time series.



Figure <u>116</u>: DBL vertical profile (VP) measurements from <u>18 March 2016</u>, in the boundary layer (BL), inversion layer (INV<u>; blue</u>), previous day's residual layer (RL<u>; yellow</u>), and free troposphere (FT). Observations corresponding to VP1-4 are shown in (a-d), respectively.

- 5 Differences in the vertical structure of  $\delta^{18}$ O,  $\delta$ D, and d excess along the four DBL VPs are shown in Figure 6. <u>A</u> defining characteristic of the DBL case study is the variability in both the meteorological and isotopic variables between each of the four VPs (Fig. 11). Observations along VP1 show an INV layer, marked by a characteristic increase in  $\theta$  and a corresponding decrease in H<sub>2</sub>O<sub>v</sub> (Fig. 11a), separating the BL and FT. However, T there appears to be two distinct atmospheric layers betweenseparating the BL and FT in VP2 (Fig. <del>6b</del><u>11b</u>). The layer directly below the FT in VP2 is <u>consistent with athe</u> residual
- 10 layer (RL) from the previous day's boundary layer (Fig. 611b). We define the base of the RL using the same approach described in Section 2.4 (d0/dz and |d(H<sub>2</sub>O<sub>v</sub>)|/dz threshold values) for determining the base of the INV ( $z_{INV}$ ). Indications for the presence of a RL are discussed in Section 4.3. Both the RL and the INV (directly below the RL) show characteristic decreases in H<sub>2</sub>O<sub>v</sub>,  $\delta^{18}$ O, and  $\delta$ D values. The presence of multiple layers is supported by the increase in the variance of the vertical wind speed (W  $\sigma^2$ ), indicating wind shear, at the interface of atmospheric layers (Fig. 6b11b-d). The temporal factors influencing the presence of
- 15 the RL are discussed below.

The  $\delta^{48}$ O and  $\delta$ D values were relatively constant with altitude along VP3 (Figure 6c), however the vertical structure of H<sub>2</sub>O<sub>v</sub> and d excess is similar in shape to VP2 (Fig. 6b). Similarities in the vertical structure of H<sub>2</sub>O<sub>v</sub> and d-excess features between VP2 and VP3 give indications that the RL persisted for an additional hour after VP2 was conducted (Fig. 11c) and INV were present during VP3 as well as VP2. Despite The winds measurement system on the aircraft had a temporary failure of the aircraft's

- 20 <u>winds measurement system</u>-halfway through the VP3 descent, <u>but</u>-available measurements show an increase in wind speed and  $\frac{1}{7}$  as well as an increase in-vertical wind variance (W  $\sigma^2$ ), as well as and sharp small-temperature changes inversions at the base and top of the <u>RLINV</u>, supporting VP3's d-excess indications of the persisting RL (Fig. <u>116</u>c). <u>The DBL case study is an example of how d-excess can support meteorological measurements in identifying distinct atmospheric layers.</u> The VP4 vertical structure of  $\delta^{48}$ O and  $\delta$ D near the surface and higher in the FT (>1400 m) are constant, and are of similar enrichment to VP3  $\delta$  values.
- 25 However, VP4  $\delta^{48}$ O and  $\delta$ D values decrease to a minimum within the INV before increasing and then plateauing in the FT.

<u>A unique feature of DBL is the presence of a residual layer (RL), which was incorporated into the BL throughout the</u> <u>duration of the flight.</u> This residual layer hypothesis is supported by Fig. S2, which shows that on the day prior to DBL, 17 March <u>2016, Rapid Refresh Model</u> ambient temperature profiles for the Indianapolis International Airport (KIND) from the day prior to DBL, 17 March 2016, which follow a nearly dry adiabatic lapse rate all the way to an altitude of 3 km (Fig. S7). This was a

- 30 relatively warm, turbulent day. and H<sub>2</sub>O<sub>v</sub> isotopologue observations were consistent with Rayleigh theory (Table 1, RF08). A cold front moved into the Indianapolis study area on DBL (18 March 2016). The ambient temperature profile on DBL shows the previous day's residual layer persisted into the early afternoon, between approximately 1 3 km (Fig. \$7\$2), before being incorporated into the BL. The distinction between the RL and BL blurs as surface heating progresses throughout the day and the RL is incorporated into the BL. It would also be expected that the downwind edge of the Indianapolis city boundaries (winds were
- 35 from the northwest; Fig. 11a3e) would have a more well mixed BL due to stronger turbulent mixing from the urban heat island and increased surface roughness (Grimmond et al., 2010; Stull, 1988). -Support for this is given by measurements made alongDBL VP3, which was conducted on the downwind edge of the Indianapolis city boundaries-observations (Fig. 6c), which- and reveal a considerably more homogenous structure in  $\delta D$  and  $\delta^{18}O$  relative to VP1 and VP2 (Fig. 11c). The ambient temperatures measured along VP3 in the FT and RL are warmer relative to ambient temperatures along the other three VP's (Fig. 611),

indicatingdemonstrating the influence of the urban heat island. Although the VP3  $H_2O_v$ ,  $\delta D$ , and  $\delta^{18}O$  values are relatively more homogenous in the vertical dimension, the d-excess signatures still maintain indications of the RL, as the vertical structure of dexcess is similar to VP2 d-excess observations. This is an example of how d-excess can provide more clues about atmospheric eirculation than can  $\delta D$  or  $\delta^{18}O$  alone.







Despite differences in δ value features along the four DBL VPs, the relationship between d excess and H<sub>2</sub>O<sub>4</sub> mole
fraction appears relatively consistent throughout the day (Fig. 4i). A positive d excess anomaly relative to Rayleigh is present at the INV FT interface, in contrast to the negative d excess anomaly in STC (Fig. 4i). Inspection of the δ<sup>18</sup>O (Fig. 4g) and δD values (Fig. 4h) show that δ values along VP1 and VP2 are consistent with Rayleigh predictions in the BL and in the more humid portion of the INV. As H<sub>2</sub>O<sub>4</sub> mole fractions decrease in the INV, the δ<sup>18</sup>O and δD signatures observed on VP1 and VP2 transition to more negative values than the Rayleigh prediction, but approximately retain their original slope, suggesting agreement with

- 10 the Rayleigh model if more isotopically depleted initial conditions were considered. DBL VP3Measurements of  $-\delta^{18}$ O and  $\delta$ D values along DBL VP3 are the only observations during our airborne campaign measurements of are enriched, and relatively constant,  $\delta$  values relative to VP1 and VP2, and are that extend from-relatively constant between the BL through to the and FT (Fig. 11c)recall the grey shading in Fig. 4 represents the average range of INV H<sub>2</sub>O<sub>v</sub>-mole fraction). Despite VP3 extending into the FT<sub>7</sub> H<sub>2</sub>O<sub>v</sub> mole fractions of 1700 ppmv and less
- 15 were observed in the FT of VP1 and, VP2, and VP4 (Fig. 4g i11a-c and Fig. 12). Relatively humid conditions were also observed at the highest altitudes flown in the FT along VP4 (Fig. 11d). Enriched, vertically unvarying VP4 δ<sup>18</sup>O and δD values are unique to DBL VP3. VP4 δ<sup>18</sup>O and δD values decrease across the INV (Fig. 11d). are and track a δ-H<sub>2</sub>O<sub>v</sub> path similar -to VP1 and VP2 in-from the BL upthrough to the lower INV-FT interface, which are similar in shape to both Rayleigh predictions or BL-FT mixing scenarios (Fig. 12). However, the trend in the VP4 H<sub>2</sub>O<sub>v</sub>-mole fraction reverses in the FT, at approximately 1370
- 20 m, and begins increasing. Interestingly, VP4 FT δ<sup>18</sup>O and δD values become enriched at higher altitudes, correspondsing to an increase in H<sub>2</sub>O<sub>v</sub> mole fractions (Fig. 11d), and instead appear to track a mixing line with the VP3 δ values (red trace in Fig. 12d-e). Despite differences in δ value features along the four DBL VPs, the relationship between d-excess and H<sub>2</sub>O<sub>v</sub> mole fraction appears remarkablylatively consistent throughout the day (Fig. 412ci). A positive d excess anomaly relative to Rayleigh is present at the INV FT interface, in contrast to the negative d excess anomaly in STC (Fig. 4i).
- 25 The relatively humid, isotopically-enriched FT air observed along VP3 and VP4 could have been caused by Athe shortwave trough in the mid-troposphere (3-5 km) which carried moist air into the Indianapolis study air in the late afternoon on this day-and likely influenced FT measurements (Fig.ure S5.3S6).- The dewpoint profile in Fig. S7S2 shows this relatively moist mid-tropospheric air descendingsubsiding over the course of the afternoon, reaching flight altitudes by the time VP3 and VP4 wereas flownconducted. The vertically unvarying, isotopically enriched VP3 observations likely reflect a combination of enhanced
   30 turbulence from the urban area and humid air from the shortwave trough mixing downward into the lower FT. In contrast, VP4 was conducted over a rural area north of Indianapolis-at the end of the flight, and likely did not experience the same degree of vertical mixing as the atmosphere downwind of Indianapolis (where VP3 was flown; Fig. 10a)-and a RL is not obvious (Fig. 6d). However, unlike VP1-3Potentially due to weaker turbulence in the rural area, a sharp decrease in δD and δ<sup>18</sup>O was observed at z<sub>FT</sub>
- 35 The DBL case study shows the  $\delta^{18}$ O,  $\delta$ D and d excesshow isotopic water vapor measurements can be used as tracers to track the development of different atmospheric structures and circulations, including residual layers, urban heat island impacts, and passing fronts.  $\delta$ D,-  $\delta^{18}$ O, and particularly d-excess can support meteorological measurements in identifying distinct atmospheric layers difficult to identify solely based on meteorological data.

on VP4 before increasing with altitude, <u>until</u> reaching enriched  $\delta$  values observed in the VP4 BL and throughout VP3.

#### 3.4 General observations of H<sub>2</sub>O<sub>+</sub> isotopologues in the lower troposphere

The case study days presented above were chosen for their distinct isotopic features and because several VPs were conducted each day. However, they only represent 30% of the research flight days (Table 1). Figure 7a c shows the Washington, D.C. Baltimore and Indianapolis VP absolute d excess observations in the FT, INV, and BL, respectively. Figures 7d f show the

- 5 measured d excess relative to Rayleigh distillation theory, i.e. Rayleigh predicted d excess has been subtracted from the observations. Overall, BL d excess observations at the Indianapolis and Washington, D.C. Baltimore study sites are relatively consistent with Rayleigh theory when using observations on each day as the initial conditions (Fig. 7c). The greatest departures from Rayleigh theory were most commonly observed in the INV (Fig. 7b). Observations of d excess in the INV at the Washington, D.C. Baltimore study site in particular were significantly more negative than Rayleigh predictions, by up to 80 %
- 10 (Fig. 7b). Very low  $H_2O_{\tau}$  mole fractions were commonly observed in the FT, as were large, positive d excess values. Large positive d excess values are predicted by Rayleigh distillation theory at low  $H_2O_{\tau}$ -mole fractions. However, observed FT d-excess values are even more positive than the Rayleigh predictions in Fig. 7a at the lowest  $H_2O_{\tau}$ -mole fractions.

The 7 March 2016 flight day in Indianapolis (RF07) stands out as an unusual set of conditions because the lower troposphere increased in H<sub>2</sub>O<sub>x</sub>-mole fraction with altitude (Fig. 7a c; gold trace). A warm, southerly front moved into the

15 Indianapolis study area on this day, and rain preceded the flight observations. The relatively high H<sub>2</sub>O<sub>x</sub>-mole fractions in the INV and FT likely reflect residual humidity from the storm. Overall, the RF07 VP observations do not exhibit distinctive isotopic features, and d excess generally varies around Rayleigh predictions for the BL, INV, and FT (Fig. 7).





Most measurements of  $H_2O_x$  mole fraction in the FT were below 4000 ppmv; the corresponding FT measurements of  $\delta^{48}O$ ,  $\delta D$ , and d-excess are plotted as a function of  $H_2O_x$  in Fig. S5. While each Washington, D.C.-Baltimore VP (Table 1) extended into the FT,  $H_2O_x$  mole fractions were often lower than the humidity range over which the TWVIA was calibrated (lower limit: 550 ppmv; Appendix B). Contributing to the scarcity of Washington, D.C. Baltimore VP measurements is the fact that only one VP was conducted per flight, mainly as a result of congested and restricted air spaces near the capitol. Values of  $\delta D$ 

range from 30 to 60‰, 8<sup>48</sup>O ranges from 200 to 400‰, and d excess can be close to 10‰, but also increase up to 100‰ below 1000 ppmv.

#### **4 Discussion**

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### 4.1 Rayleigh-Consistent Conditions

- 10 Previous observations have shown that Rayleigh distillation theory successfully explains most of the variability in H<sub>2</sub>O<sub>\*</sub> isotopic composition observed in the BL (Lee et al., 2006). RAY is the simplest case study day where δ values and d excess observations generally track the Rayleigh predictions in Fig. 4a c. However, positive deviations from Rayleigh d excess exist in the FT, and we hypothesize that dry FT air parcels carrying large, positive d excess values can mix downward into more humid air parcels of smaller d excess values near the top of the INV (Sodemann et al., 2017). As the H<sub>2</sub>O<sub>\*</sub> mole fraction approaches zero, Rayleigh predicted d excess approaches 7000‰ (Bony et al., 2008). Thus, lower altitude FT air masses can carry a more positive d excess signature for a given H<sub>2</sub>O<sub>\*</sub> mole fraction due to vertical mixing, than is predicted by Rayleigh theory (Fig. S5c). The d-
- excess signature in the FT more closely follows the FT BL mixing line (Fig. 4c), which supports this hypothesis. Our results are consistent with Dyroff et al. (2015) who report lower troposphere δD observations over the Atlantic Ocean, and explain the vertical structure of δD at lower altitudes using Rayleigh theory, while higher altitude observations indicate mixing scenarios dictate the 20 δD profile.

From our measurements, we can comment on what types of atmospheric conditions are likely to produce profiles consistent with Rayleigh theory, i.e. equilibrium fractionation and no vapor condensate (cloud droplet) equilibration. The RAY meteorological and isotopic observations are very similar to those made on 17 March 2018 (RF08; Table 1). These two flight days' observations provide possible criteria for when it is appropriate to use Rayleigh theory to predict isotopic signature throughout the BL and INV. These criteria include that the study area is free of clouds/precipitation (supported by the presence of a dry adiabatic lapse rate throughout the BL) and wind speed and direction is relatively constant throughout the BL, INV, and FT. The STC and

DBL case studies, which do not follow Rayleigh theory, violate one or more these criteria.

#### **4.2 Stratocumulus Cloud Evaporation**

Cloud processes may cause the d excess anomalies from Rayleigh theory in the STC case study day. When evaluating the effect that cloud evaporation could have on vertical profiles of water vapour isotopic variability, we must consider the altitudes at which clouds form and evaporate. In the case of evaporation, this may not occur in a single layer, but throughout a range of altitudes or atmospheric layers. Stratocumulus cloud tops are typically present directly below z<sub>INV</sub> (Wood, 2012). The top of the INV (z<sub>FT</sub>) is approximately the upper limit of BL mixing (Wood, 2012). Lofting of cloud droplets into the INV was under saturated (Fig 5b). In the cloud evaporation scenarios discussed below, it is assumed that the cloud droplets form at the top of the BL, and the droplets undergo evaporation within the INV.

Figure 8a-b shows H<sub>2</sub>O<sub>+</sub>  $\delta$ D and d-excess predicted for several cloud evaporation scenarios. We choose to not show  $\delta^{18}$ O under these cloud droplet evaporation scenarios because the results are similar to those for  $\delta$ D. Scenario 1 in Figure 8a shows that

the  $\delta D$  isotopic signature of atmospheric vapor that is influenced by 35% cloud droplet evaporation air (eq. 3 in Section 2.5) tracks the STC observations for VP3 and VP4 within the INV. While cloud evaporation scenario 1 tracks with observed d excess in more humid portions of the INV (Figure 8b), it predicts larger positive d excess in dryer portions of the INV (i.e. <2500 ppmv). This is the opposite of what is observed in d excess at the upper (drier) part of the INV for VP2 through VP5. Scenario 1 does describe the steep slope in  $\delta D$  (and  $\delta^{18}O$ ) values in the INV but does not explain the minimum in d-excess (Figure 8a).



**Figure 8:** STC vertical profile (VP) (a)  $\delta D$  and (b) d-excesss observations plotted with different cloud droplet evaporation scenarios. The two scenarios consider partial cloud droplet evaporation (Scenario 1) or the complete evaporation of previously dehydrated cloud droplets followed by mixing (Scenario 2).

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Cloud evaporation scenario 2 considers the effect of complete evaporation of a previously dehydrated cloud droplet (from eqs. 4 6 in Section 2.5). Figure 9 shows the d excess signature of an evaporating cloud droplet that was formed at the top of the BL (consistent with stratocumulus cloud formation) evaporating into air that has an isotopic composition consistent with that observed at the top of the INV on STC. Figure 9 shows the d excess signature of cloud droplets as they near complete evaporation,

- 15 highlighting 65%, 80%, and 95% evaporation values. Figure 8 shows the effect of complete evaporation of these semi-dehydrated (65%, 80%, and 95%) cloud droplets at the top of the INV, followed by subsequent mixing with INV air (eq. 2 in Section 2.5). As can be seen from Figure 8b, scenario 2 describes the minimum d-excess signature at the top of INV. However, Figure 8a shows that scenario 2 does not agree with VP2 4 & D observations within the INV. It is possible that partial evaporation (scenario 1) occurs in the lower half of the INV, followed by complete evaporation of previously dehydrated cloud droplets at the top of the INV.
- 20 (scenario 2). We believe the minimum d excess anomaly at z<sub>FT</sub> for VP2 5 may result from complete evaporation of a partially dehydrated cloud droplet (Fig. 8b). We find discussion of these potential d excess anomalies in the literature in reference to

raindrops evaporating below the cloud layer, but we believe the same isotopic fractionation would occur as liquid cloud droplets evaporate in unsaturated environments, such as at the top of the INV (Aemisegger et al., 2015; Gat, 1996; Sodemann et al., 2017).



5 Figure 9. Cloud droplet (condensate) d-excess as a function of the fractional amount of the droplet lost to evaporation (calculated using eqs. 4-6). As cloud droplets dehydrate, d-excess values can become extremely negative. The d-excess signature of cloud droplets that have evaporated by 65%, 80%, and 95% are highlighted. The impact on d-excess of complete evaporation of these previously dehydrated cloud droplets, followed by mixing, is shown in Figure 8.

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The slight easterly movement of the cloud layer as shown in Fig. S3 could be interpreted as the stratocumulus cloud layer being advected out of the Indianapolis study area by westerly winds rather than evaporating away. However, Fig. 5b shows that wind speeds at the cloud layer were relatively low (<6 m s<sup>-1</sup>), and the clouds in Figure S3 transition from opaque to semi-transparent over the flight period. This information taken together suggests there was evaporation of the stratocumulus cloud layer during the flight.

- 15 We believe our observations provide evidence that the process of cloud evaporation may be spread throughout the INV layer with the beginning and end of evaporation separated in space. Within the INV, turbulent updrafts and downdrafts are warmer and cooler than the surrounding air, respectively (Betts, 1985). Updrafts potentially carrying partially evaporated cloud droplets to the top of the INV may facilitate complete evaporation of the droplets due to the low RH and possibly wind shear promoted mechanical turbulence at z<sub>FT</sub> (Fig. 5b). Due to higher RH values within the middle of the INV relative to z<sub>FT</sub>, it is possible cloud
- 20 droplet evaporation did not proceed to completion at these INV altitudes. Partial evaporation would impart a positive d excess signal on atmospheric vapor and act to increase RH (Aemsigger et al., 2015; Sodemann et al., 2017), both of which exhibit a local maximum at ~985 m agl (Fig. 5b).

Other studies have reported negative water vapor d excess and hypothesize that these are from kinetic fractionation of vapor during deposition on ice crystals or snow (i.e. in ice supersaturated conditions) (Casado et al., 2016; Galewsky, 2015;

25 Lowenthal et al., 2016; Moore et al., 2016; Samuels Crow et al., 2014; Schmidt et al., 2005). It is unlikely that ice supersaturation is responsible for the minimum in vapor d excess observed at z<sub>FT</sub> on STC because temperatures were greater than 0°C (Table 1). However, as an example, Figure D1 shows the theoretical d excess values of STC vapor under ice supersaturated conditions (Appendix D). We reiterate that ice supersaturation is an unlikely explanation for the STC z<sub>FT</sub> d excess minimum because flight

altitudes were less than 2 km, and ice (cirrus) clouds are typically present at ~6 km. It is unlikely that ice hydrometeors falling from higher altitudes could be sustained at the top of the inversion and contribute to the low d-excess signal observed on VP2– VP5 through vapor deposition given the >0°C temperatures. While vapor deposition in ice supersaturated conditions might not be relevant for the STC flight, low d-excess (relative to Rayleigh at RH = 100 %) was sometimes observed in the INV layers during

5 the Washington, D.C.-Baltimore flights (Fig. 7b). Ambient temperatures observed in flight in Washington, D.C.-Baltimore were sometimes less than 0°C (Table 1), thus vapor deposition on ice crystals would be more likely for those scenarios.

#### 4.3 Developing Boundary Layer

Unlike RAY and STC, the vertical structures of H<sub>2</sub>O<sub>v</sub>, δD, δ<sup>18</sup>O, and d excess observed in DBL are highly variable in
 both space and time (Fig. 6). Despite variability in the vertical structure, a fairly consistent d excess H<sub>2</sub>O<sub>v</sub> relationship is maintained throughout the flight (Fig. 4i). Here we attempt to interpret and explain the heterogeneity in the DBL VPs.

A unique feature of DBL is the presence of a residual layer (RL), which was incorporated into the BL throughout the duration of the flight. This residual layer hypothesis is supported by Fig. S2, which shows that on the day prior to DBL, 17 March 2016, Rapid Refresh Model ambient temperature profiles follow a nearly dry adiabatic lapse rate all the way to an altitude of 3 km.

- 15 This was a relatively warm, turbulent day and H<sub>2</sub>O<sub>v</sub> isotopologue observations were consistent with Rayleigh theory (Table 1, RF08). A cold front moved into the Indianapolis study area on DBL (18 March 2016). The ambient temperature profile on DBL shows the previous day's residual layer persisted into the early afternoon, between approximately 1 3 km (Fig. S2), before being incorporated into the BL. The distinction between the RL and BL blurs as surface heating progresses throughout the day and the RL is incorporated into the BL. It would also be expected that the downwind edge of the Indianapolis city boundaries (winds were
- 20 from the northwest; Fig. 3c) would have more well mixed BL due to stronger turbulent mixing from the urban heat island and increased surface roughness (Grimmond et al., 2010; Stull, 1988). Support for this is given by DBL VP3 observations (Fig. 6c), which reveal a considerably more homogenous structure in δD and δ<sup>18</sup>O relative to VP1 and VP2. The ambient temperatures measured along VP3 in the FT and RL are warmer relative to ambient temperatures along the other three VP's (Fig. 6), indicating the influence of the urban heat island. Although the VP3 H<sub>2</sub>O<sub>v7</sub> δD, and δ<sup>18</sup>O values are relatively more homogenous in the vertical
- 25 dimension, the d excess signatures still maintain indications of the RL, as the vertical structure of d excess is similar to VP2 d excess observations. This is an example of how d excess can provide more clues about atmospheric circulation than can  $\delta D$  or  $\delta^{18}O$  alone.

A shortwave trough in the mid troposphere (3-5 km) carried moist air into the Indianapolis study air in the late afternoon on this day and likely influenced FT measurements (Figure S6). The dewpoint profile in Fig. S2 shows this relatively moist mid-

- 30 tropospheric air descending over the course of the afternoon, reaching flight altitudes by the time VP4 was conducted. VP4 was conducted over a rural area north of Indianapolis at the end of the flight and a RL is not obvious (Fig. 6d). However, unlike VP1-3, a sharp decrease in δD and δ<sup>48</sup>O was observed at z<sub>FT</sub> on VP4 before increasing with altitude until reaching δ values observed in the VP4 BL and throughout VP3. The DBL case study shows the δ<sup>48</sup>O, δD and d excess measurements can be used as tracers to track the development of different atmospheric structures and circulations, including residual layers, urban heat island impacts,
- 35 and passing fronts.

#### 4.4 Features of the lower troposphere d-excess vertical profile

The data presented here are some of the few lower troposphere water vapour isotope observations published, so we look for common patterns that can be used to predict values in other areas. Nearly all BL observations of d excess at the Indianapolis and Washington, D.C. Baltimore study sites agree well with Rayleigh theory (Fig. 7c). This is consistent with previous

5 observations showing that Rayleigh distillation theory successfully explains most of the variability in  $H_2O_v$  isotopic composition observed in the BL (Lee et al., 2006).

The FT is generally drier than the BL (Fig. 7a), and FT d excess values observed at very low H<sub>2</sub>O<sub>v</sub> mole fractions are often more positive than predicted by Rayleigh theory (Fig. 7a). This has been explained by very dry, depleted FT air masses, which carry large positive d excess signatures, mixing downward towards flight level altitudes (Bony et al., 2008; Sodemann et

10 al., 2017). These FT air masses likely would have originated from another source region and possibly underwent multiple condensation cycles to achieve such isotopic depletion prior to mixing with more humid air across the INV. Thus, we do not necessarily expect FT air to have a d-excess signature consistent with Rayleigh theory of an ascending BL air parcel (Dyroff et al., 2015).

We observed large departures from Rayleigh theory in the INV. Figures 7b and 7e show that d excess observations in
 the Washington, D.C. Baltimore INV can deviate negatively by ~80‰ relative to Rayleigh predictions. We believe the minimum in STC d excess at z<sub>FT</sub>-is a result of stratocumulus cloud evaporation. Partly cloudy or overcast conditions were also present over the Washington, D.C. Baltimore study site on all four Washington, D.C. Baltimore flights (https://worldview.earthdata.nasa.gov/; https://www.ncdc.noaa.gov). It is possible that the very negative Washington, D.C. Baltimore INV d excess measurements were a result of complete evaporation of semi-evaporated cloud droplets within the

20 inversion. We note that the most negative Washington, D.C. Baltimore d excess values correspond to the driest INV observations (Fig. 7b), where we would expect cloud top evaporation resulting from free troposphere entrainment to be the most prevalent.

#### **5** Conclusions

The aim of this study is to observe and interpret the vertical structure of H<sub>2</sub>O<sub>v</sub> stable isotopic composition, specifically dexcess, in the continental lower troposphere. Current literature regardingPrevious studies presenting d-excess observations areis heavily focused on ocean evaporation at coastal or island surface sites (Benetti et al., 2014; 2015; 2018; Delattre et al., 2015; Steen-Larsen et al., 2014; Uemura et al., 2008). Few reported observations of d-excess in the INV and FT exist (Galewsky et al., 2015; Lowenthal et al., 2016; Kelsey et al., 2018; Samuels-Crow et al., 2014; Schmidt et al., 2005; Sodemann et al., 2017), and, to our knowledge, only one study has used airborne measurements to provide high vertical resolution snapshots of the lower troposphere d-excess profile at discrete time points (Sodemann et al., 2017). Our stable H<sub>2</sub>O<sub>v</sub> isotope measurements over two continental sites

is a starting point in filling the field's gap in understanding variability in the lower troposphere d-excess profile, and what it reveals about lower troposphere moisture processing on relatively small regional scales.

Our observations reaffirm the dominant role that Rayleigh distillation processes have on  $\underline{H_2O_v}$  water vapour isotopic variabilitysignature, and that isotopic signatures can be retained by air masses as they are advected from previous points of

35 <u>condensation. Our measurements indicate that  $H_2O_v$  isotopologues, and d-excess in particular, can act as fingerprints of earlier processes. especially in the boundary layer.</u> This process is true for the CLR case study, in which Rayleigh distillation curves well-representeen predict vertical profile observations from near the surface up through the atmospheric inversion layer, despite that the

temperature profile within the BL followeding the dry, rather than the (Rayleigh-assumed) moist<sub>7</sub> adiabat. -in clear sky conditions with reasonably constant wind conditions with height. Similarly, the evidence of upwind processes was also retained in the  $H_2O_v$ isotopic signature on STC. The STC measurements give clues that the observed air mass experienced prior evaporation of cloud or rain droplets. Measurements made during the rapidly changing atmospheric conditions on the DBL case study also show that

- 5  $\underline{H_2O_v}$  isotopic signature can reflect near-instantaneous changes in the atmosphere. These new results highlight the potential for  $\underline{H_2O_v}$  water vapour i isotope ratios, especially d-excess, to diagnose-identify complex processes across the atmospheric inversion layer including cloud condensation, evaporation, and mixing or entrainment of free tropospheric air into the boundary layer. These types of measurements may become increasingly valuable as we seek to understand the physical processes that sustain cloud layers and spatial-otemporally variable boundary layer mixing depths.
- Our interpretation of the d-excess VPs could be further evaluated by isotope-enabled circulation and weather models (Aemisegger et al., 2015; Pfahl et al., 2012; Schmidt et al., 2005). However, the simulation of convective boundary layer-BL processes with isotope-enabled models is complex (Bolot et al., 2013; Benetti et al., 2018). The measurements reported here could help further develop current and forthcoming isotope enabled models, particularly for simulating wintertime, continental lower troposphere processes or stratocumulus evaporation. Our observations of the d-excess profile in a stratocumulus cloud-topped BL boundary layer and the d-excess observations reported by Sodemann et al. (2017) near marine cumulus clouds seem to indicate represent an opportunity to investigate cloud-class specific RH and d-excess relationships. Future studies could interrogate investigate the sensitivity of the d-excess signature to different classes of clouds and their associated unique cloud processes.

## Data availability.

GeospositionalGeopositional, meteorological, greenhouse gas, and water vapor isotope measurements are available for the
 Washington, D.C.-Baltimore and Indianapolis flight days (Table 1) are available by request and will be archived upon publication in the stable water vapour isotope database: <u>https://vapor-isotope.yale.edu/</u>. The authors request that they be notified if the data is to be used in publication.

#### Appendices

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#### Appendix A. Comparison of LGR and Picarro H<sub>2</sub>O<sub>+</sub> mole fraction

Figure A1 compares calibrated H<sub>2</sub>O<sub>4</sub> mole fractions from the Los Gatos Research (LGR) Triple Water Vapor Isotope Analyzer (TWVIA) and the Picarro cavity ringdown spectrometer during the entire flight conducted on 6 March 2016 (RAY).

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Figure A1: Comparison of calibrated LGR and Picarro H<sub>2</sub>O<sub>4</sub>-mole fraction measurements from the entire research flight conducted on 6 March 2016 (RAY).

## 5 Appendix B. Water vapor concentration-dependence calibration

A Los Gatos Research (LGR) Water Vapor Isotope Standard Source (WVISS; model: 908 0004 9003) equipped with a secondary dry air mixing chamber for extended range operation was used to characterize the LGR Triple Water Vapor Isotope Analyzer's (TWVIA; model: 911-0034) non-linear response to water vapor (H2O,) concentration (Rambo et al., 2011). The WVISS samples liquid water with a known isotopic composition from a reservoir. The standard sample is then nebulized using zero (dry) 10 air into a heated chamber (75°C), where it evaporates completely and is further diluted with zero (dry) air with programmable flow rates to output a range of H<sub>2</sub>O<sub>x</sub> fractions with the same isotopic signature as the liquid standard. Different combinations of nebulizer sizes (flow rates) and standard versus extended range operation were required to span a large range of H<sub>2</sub>O<sub>x</sub> values. The TWVIA's H<sub>2</sub>O<sub>2</sub>-dependence (while operating in extended range mode, ~80 Torr) was evaluated over the range from 550 ppmv 14,000 ppmv, consistent with range of H<sub>2</sub>O<sub>\*</sub> mole fractions observed during the research flights (Table 1). Free troposphere H<sub>2</sub>O<sub>\*</sub> mole fractions were sometimes less than 550 ppmv, but the lowest H<sub>2</sub>O<sub>v</sub>-mole fraction the WVISS can emit is 500 ppmv and we found 15 stable flows of  $H_2O_x$  mole fractions lower than 550 ppmv were difficult to achieve with the WVISS. We opt not to report  $\delta D$  and  $\delta^{18}$ O values for instances when H<sub>2</sub>O<sub>2</sub> mole fraction is less than 550 ppmy. The  $\delta$ D and  $\delta^{18}$ O values of the H<sub>2</sub>O<sub>2</sub> isotope standards, which bracket the ranges observed during the research flights (Table 1), are listed in Table B1. The WVISS was programmed to sample each H<sub>2</sub>O<sub>4</sub> mole fraction for  $\geq 20$  min. The  $\delta D$  and  $\delta^{+8}O$  H<sub>2</sub>O<sub>4</sub> dependence calibration curves were constructed from the average  $\delta D$  and  $\delta^{18}O$  values reported during the last 200 s of each calibration period in order to remove any influence of transition 20

- instability caused by water moving onto and off of the walls of the system during the calibration  $H_2O_x$  step changes. The  $\delta^{18}O$  and  $\delta D H_2O_x$  dependence curves shown in Fig. B1 and B2, respectively, were fit using the locally weighted polynomial regression "locpoly" function from R's "locfit" package (Bailey et al., 2015). A 100 ppmv sliding window was used for the local polynomial regression fitting over the range from 550 ppmv 14,000 ppmv  $H_2O_x$ -
- 25 Table B1: Calibration standards

Standard*	<del>ðD (‰)</del>	<del>ð<sup>18</sup>O (‰)</del>	<del>d-excess (‰)</del>
Purdue tap water	<del>-39.9</del>	-8.7	<del>29.7</del>

Boulder tap water	<del>-117.3</del>	<del>-15.4</del>	<del>5.9</del>
USGS 46	<del>-235.8</del>	<del>-29.8</del>	<del>2.6</del>
South Pole Glacier Water	- <u>434.5</u>	<del>-54.3</del>	<del>-0.1</del>
Custom Light Blend <sup>‡</sup>	<del>-573.7</del>	<del>-76.2</del>	<del>36.1</del>

\*Standard values are reported relative to the VSMOW SLAP scale

<sup>+</sup>The Custom Light Blend is a mixture of Purdue tap water, Sigma Aldrich deuterium depleted water ( $\leq 1$  ppm HDO), and Isotee 95% H<sub>2</sub><sup>18</sup>O (<sup>18</sup>O enriched) to achieve a depleted isotopic signature that brackets the most depleted research flight observations of  $\delta D$  and  $\delta^{18}O$  that also has a realistic d excess signature. Because the Custom Light Blend is isotopically more depleted than our standards, known amounts of the Custom Light Blend and Purdue tap water were combined to make three mixtures, which were analysed using an LGR liquid water isotope analyser (T-LWIA-45-EP; model: 912-0050-0001) to determine the Custom Light Blend's isotopic signature.

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The TWVIA's H<sub>2</sub>O<sub>\*</sub> dependence curve was reproducible over all  $\delta^{48}$ O isotope standard signatures considered (Table B1). The  $\delta D H_2O_*$  dependence curve was reproducible for the three relatively enriched isotope standards, more enriched than 235.8‰, in Table B1 but was not always reproducible using the most depleted standards (South Pole Glacier and Custom Light Blend) over the H<sub>2</sub>O<sub>\*</sub> range of ~3000 ppmv to ~ 8000 ppmv as is shown in the Figure B1a and B2a. At H<sub>2</sub>O<sub>\*</sub> mole fractions above and below that range, the calibration curve remained reproducible. The cause of the 3000 ppmv – 8000 ppmv irreproducibility of the  $\delta D$ -

- 15 H<sub>2</sub>O<sub>v</sub>-dependence curve associated with very depleted δD values remains unknown, perhaps small leaks in the experimental setup or uncertainty associated with curve fitting. To our knowledge this behaviour has not been described in the literature. However, δD values consistent with the two most depleted standards (Table B1) were only observed in the free troposphere and correspond to low H<sub>2</sub>O<sub>v</sub> mole fractions (<1000 ppmv) and were outside of the irreproducible window of H<sub>2</sub>O<sub>v</sub> values. Therefore, it was not consequential to actual flight observations in this experiment. We note that there also appears to be large variability in the TWVIA-
- 20 reported δD values <1000 ppmv H<sub>2</sub>O<sub>v</sub> for the two depleted standards, but there is also relatively larger variability in this H<sub>2</sub>O<sub>v</sub> range for the enriched standards as well. To avoid biases resulting from the depleted δD irreproducibility, the δD water vapor dependence curve is defined using calibration data from the three relatively enriched standards. However, δD calibration data from each of the five standards is used to define uncertainties (see below in Appendix C).

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**Figure B1:**  $\delta^{18}$ O-H<sub>2</sub>O<sub>\*</sub> dependence (a) calibration curve and (b) residuals. The true  $\delta^{18}$ O signature of each standard (Table B1) has been subtracted from the TWVIA measurements to give the "adjusted"  $\delta^{18}$ O signature in (a).



**Figure B2: \deltaD-H**<sub>2</sub>**O**<sub>\*</sub> **dependence** (a) **calibration curve and** (b) **residuals. The true \deltaD signature of each standard (Table B1) has been subtracted from the measurements to give the "adjusted" \deltaD signature in (a).** 

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To check calibration the VSMOW SLAP scale, Figure B3 shows that the linear regressions of the isotope standards' H<sub>2</sub>O<sub>\*</sub> concentration dependence corrected  $\delta$  values versus true gas phase isotopic signature for (a)  $\delta^{48}$ O and (b)  $\delta$ D have slopes near unity and intercepts near zero.  $\delta^{48}$ O had a slope of 1.009(±0.001), a y intercept of 0.08(±0.03), and an R<sup>2</sup> of 0.997254. The  $\delta$ D ordinary least squares regression line had a slope of 0.9954(±0.0005), a y intercept of  $-0.5(\pm0.09)$ , and an R<sup>2</sup> of 0.99958. A VSMOW SLAP correction was not applied because it would be negligible compared to the uncertainty associated with the concentration dependence correction and the instrument precision (Appendix C).



Figure B3. VSMOW-SLAP calibration curves for (a)  $\delta^{18}$ O and (b)  $\delta$ D. H<sub>2</sub>O<sub>4</sub>-concentration-corrected isotopic signatures are plotted against the standard's true isotopic signature. Linear regression fit slopes and intercepts are included in the figure insets.

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## Appendix C. Water vapor $\delta D$ , $\delta^{18}O$ , and d-excess error propagation

#### **Instrument precision:**

10 The TWVIA instrument precision was calculated as the  $1\sigma$  standard deviation for the last 20 seconds of every calibration period (Appendix B). The interval used to smooth the  $\delta D$ ,  $\delta^{18}O$ , and d excess values reported in this paper is 20 s, which corresponds to the time required for the TWVIA signal to stabilize after a change in the sample's H<sub>2</sub>O<sub>\*</sub> mole fraction or isotopic signature. The  $\delta D$  and  $\delta^{18}O$  precision values are calculated as a function of H<sub>2</sub>O<sub>\*</sub> mole fraction using power functions (Fig. C1).



Figure C1: TWVIA δ<sup>18</sup>O and δD 20-s instrument precision (1σ) as a function of water vapor (H<sub>2</sub>O<sub>\*</sub>) mole fraction.

## H<sub>2</sub>O<sub>\*</sub>-dependence calibration uncertainty:

5 The uncertainty associated with the TWVIA δD and δ<sup>14</sup>O H<sub>2</sub>O<sub>v</sub> dependence corrections is determined from the calibration residuals shown in Fig. B1b and Fig. B2b. We note that the calibration residuals do include a small instrument precision component, as the calibration values are the average of 200 s sampling periods. The absolute value of the δD and δ<sup>48</sup>O residuals from all five reference waters tested were filtered into bins defined by 100 ppmv H<sub>2</sub>O<sub>v</sub>-mole fraction increments. Averages of the absolute δD and δ<sup>48</sup>O residuals were calculated for each bin. For relatively dry conditions (i.e. below 3500 ppmv H<sub>2</sub>O<sub>v</sub>), the bin-averaged calibration residuals increase as H<sub>2</sub>O<sub>v</sub>-mole fractions decrease. A best fit linear regression was determined for the bin-averaged residuals as a function of H<sub>2</sub>O<sub>v</sub>-mole fraction (from 550 – 3500 ppmv for δD and 550 – 3700 ppmv for δ<sup>48</sup>O). Bin-averaged residuals were relatively constant for H<sub>2</sub>O<sub>v</sub>-mole fractions greater than 3500 ppmv for δD and 3700 ppmv for δ<sup>48</sup>O. Average H<sub>2</sub>O<sub>v</sub>-dependence calibration uncertainties of 1.8‰ for δD and 0.9‰ for δ<sup>48</sup>O were calculated from the bin averaged residuals from 3500 – 14000 ppmv for δD and 3700 – 14000 ppmv for δ<sup>48</sup>O. Higher uncertainties in the δ values at low H<sub>2</sub>O<sub>v</sub>-mole fractions is not surprising, as the manufacturer suggests the TWVIA be used for sampling air ranging from 4,000 – 60,000 ppmv H<sub>3</sub>O<sub>v</sub>-

#### **Total uncertainty:**

Total  $\delta D$  and  $\delta^{18}O$  uncertainty is calculated by propagating the error resulting from instrument precision,  $S_{precision}$ , and 20 from the calibration,  $S_{calibration}$ , as in eq. (C1):

The total d excess uncertainty is determined according to eq. (C2):

(C1)

where  $S_{total,\delta D}$  and  $S_{total,\delta^{10}D}$  are the total  $\delta D$  and  $\delta^{18}O$  uncertainties (given be eq. (C1)). The total uncertainty for  $\delta D$ ,  $\delta^{18}O$ , and d-excess as function of H<sub>2</sub>O<sub>+</sub> mole fraction is presented in Fig. 1.

## 5 Appendix D. Fractionation of water vapor in ice supersaturated conditions

H<sub>2</sub>O<sub>v</sub>-undergoing deposition in ice supersaturated conditions is impacted by equilibrium and kinetic fractionation. The kinetic fractionation factor is calculated via Galewsky (2015) eq. (D1):

where is saturation with respect to ice, expressed as a fraction. The equilibrium fractionation factor calculated for the temperature
 at the lifting condensation level (LCL) and is discussed in Methods 2.5. The ratio of the molecular diffusivity of the light to heavy isotopologue, is 1.02849 for <sup>18</sup>O and 1.02512 for D (Merlivat, 1978).

The isotopic signature of an air parcel in ice supersaturated conditions () can be calculated according to eq. (D2):

-is the heavy to light isotopologue ratio ( or ) of the parcel prior to the ascent. The remaining fraction of  $H_2O_{v}$  left in the ascending parcel relative to initial conditions is given by .

Figure D1 shows the STC VP d excess observations along with Raleigh vapor calculated from RH = 100 % (Methods 2.5) and vapor in ice supersaturated conditions (eq. (D2)). To match the most negative d excess value observed at the top of the INV on STC, a supersaturation () of 1.17 ( $RH_i = 117\%$  in Fig. D1) was used but does not necessarily reflect reality for the temperature or altitude of the observations. The <u>Rayleigh</u> curve is presented for reference.

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**Figure D1: STC VP d-excess observations, Rayleigh vapor d-excess, and calculated d-excess of vapor in ice supersaturated** (**RHi**) conditions up to **RHi = 117 %. Ice supersaturated conditions were chosen merely to match the INV-FT interface dexcess observations, and do not necessarily reflect a realistic RH**<sub>i</sub> for the STC flight day.

(C2)

(D1)

(D2)

Supplement link. Figures Sections S1-6-S7 are provided in the accompanying Supporting Information. Figure S3 and Figure S4 are GIFs available in separate files.

## Competing interests.

The authors declare no competing interests.

## 5 Author contribution.

OS, LW, and PS designed the experiments. OS collected the airborne data, with the help of KH. OS analysed the data. OS, LW, MB, and PS interpreted the results. OS prepared the manuscript with contributions from all co-authors. BS maintained the experimental aircraft.

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