



Detection of a climatological short break in the Polar Night Jet in early winter and its relation to cooling over Siberia

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Abstract. The Polar Night Jet (PNJ) is a strong stratospheric westerly circumpolar wind at around 65°N in winter, and the strength of the climatological PNJ is widely recognized to increase monotonically from October through late December. Remarkably, the climatological PNJ temporarily stops increasing during late November. We examined this short break in terms of the atmospheric dynamical balance and found that it results from an increase in the upward propagation of climatological planetary waves from the troposphere to the stratosphere in late November, which coincides with a maximum of the climatological Eliassen-Palm flux convergence in the lower stratosphere. The upward propagation of planetary waves at 100 hPa, which is strongest over Siberia, is related to the climatological strengthening of the tropospheric trough over Siberia. We suggest that longitudinally asymmetric forcing by land–sea heating contrasts caused by their different heat capacities can account for the strengthening of the trough.

1 Introduction

In the Northern Hemisphere (NH) winter, the high-latitude stratosphere is characterized by strong westerly winds around the polar vortex, the so-called Polar Night Jet (PNJ) (e.g., AMS, 2015; Waugh et al., 2017). The PNJ exhibits large interannual and intraseasonal variations dynamically forced by the upward propagation of planetary-scale Rossby waves from the troposphere. On an intraseasonal timescale, the PNJ strength signal propagates downward and poleward from the upper stratosphere to the high-latitude lower stratosphere during winter (e.g., Kuroda and Kodera, 2004; Li et al., 2007). This variation is called the PNJ oscillation. The signal further propagates into the troposphere to produce the Arctic Oscillation (AO; Thompson and Wallace, 1998, 2000) signal at the surface (e.g., Baldwin and Dunkerton, 2001; Deng et al., 2008). The AO, which is the dominant hemispheric seesaw variability in sea level pressure between the polar area and the surrounding mid-latitudes, strongly influences NH



weather patterns and its associated extreme weather events (e.g., Black and McDaniel, 2009; Thompson and Wallace, 2001; Ando et al., 2015; Drouard et al., 2015; Xu et al., 2016; He et al., 2017).

Propagating large-amplitude planetary waves sometimes cause a sudden decrease in the strength of the PNJ accompanied by a sudden increase in polar temperature, a phenomenon known as a sudden stratospheric warming (SSW) event (Matsuno, 1970; Labitzke, 1977). Extreme SSW events occur mostly in mid or late winter; in early winter or early spring, SSWs are weaker and less frequently occur (e.g., Charlton and Polvani, 2007; Maury et al., 2016).

Although the interannual variability of the PNJ has been well studied (e.g., Frauenfeld and Davis, 2003; Kolstad et al., 2010; Woo et al., 2015), the climatological seasonal evolution of the PNJ has been overlooked. Considering the downward propagation of the PNJ strength signal from the lower stratosphere and its effect on tropospheric weather and climate, a detailed understanding of the climatological seasonal evolution of the PNJ is important for weather patterns and extreme weather events. It is generally acknowledged that the climatological PNJ speed increases monotonically from October to December and reaches a maximum in early January. Subsequently, the speed of climatological PNJ decreases monotonically until spring.

We found that in the lower stratosphere the climatological PNJ temporarily stops increasing in late November, and it temporarily stops decreasing in late February (Fig. 1; see Section 3.1 for a detailed explanation). These “short breaks” in the seasonal evolution of the climatological PNJ cannot be detected in monthly averaged data; their detection requires data with a finer temporal resolution. The climatological short break in February is likely due to the fact that SSWs occur less frequently in late February compared with January and early February. The vertical structure and timescale of the short break in late November is different from that of February (see Section 3.1 for a detailed explanation). A detailed understanding of the short break in late November is important in terms of dynamic meteorology of intraseasonal variations in stratosphere. The climatological short break in late November might possibly have been known, but it has not yet been addressed in terms of dynamic meteorology. We examine this climatological short break of the PNJ in late November through a dynamical analysis to infer a possible origin. In Section 2, we briefly describe the data used. Section 3 provides a detailed description of the late November climatological short break. Section 4 discusses a possible cause of the short break, and Section 5 presents our conclusions.

2 Data and methods

We used the 6-hourly Japanese 55-year Reanalysis (JRA-55) dataset with the 1.25° horizontal resolution (Kobayashi et al., 2015; Harada et al., 2016). Because the quality of the stratospheric analysis was improved after the inclusion of satellite data in JRA-55 in 1979, the analysis period was restricted to the period from 1979 through 2016. We therefore defined climatological values as their 38-year average values during 1979–2016.



As our main analysis method, we performed an Eliassen-Palm (EP) flux analysis based on the transformed Eulerian mean (TEM) momentum equation (Equation A3 in Appendix A). This method, which is widely used in dynamic meteorology to diagnose wave and zonal-mean flow interaction, is described in detail in Appendix A.

3 Results

5 3.1 Climatological short break of the Polar Night Jet

First, we outline the seasonal evolution of the PNJ. A latitude–time cross section of the climatological zonal-mean zonal wind at 50 hPa over 50–90°N shows that the strength of the zonal-mean westerlies at 50 hPa ($\bar{U}50$) increases with time from approximately October to late December. Subsequently, $\bar{U}50$ decreases with time from late December through March (Fig. 1a). An examination of the intraseasonal variation of $\bar{U}50$ reveals two short breaks. Between 60° and 80°N, there is a pause in the increasing trend in late November, and there is another pause in the decreasing trend in late February. The short break in late November is statistically significant at the 99% level, and the short break in late February is statistically significant at the 95% level (the two-tailed Student's t test; e.g., Wilks, 2011). The climatological short break in February is likely associated with the less frequent occurrence of SSWs in late February than in January and early February. We note the signal throughout the whole stratosphere. In contrast, the climatological short break in November is restricted to the lower and mid stratosphere (figure not shown).

Here, we defined the zonal-mean zonal wind speed at 65°N and 50 hPa as a PNJ index, following Kodera and Koide (1997). The time series of the climatological 15-day running mean of this PNJ index clearly shows a short break of the PNJ during late November (blue line in Fig. 1b). There are various bumps in the time series of a lower dashed-dotted line in Fig. 1b. This signifies that extreme short breaks occur regardless of the time of the season. The climatological short break of the PNJ, however, occurs only during late November and early February (blue line in Fig. 1b). This suggests that the short breaks in each year more often occur during these two periods than during the other periods. Because it is common that the cause of the February short break corresponds to the occurrence of SSWs, and because the short break in late November more statistically significant than that of February, this paper does not target the February short break. We thus focus on the climatological short break of the PNJ in late November.

3.2 Anomalous upward propagation of the EP flux during late November

In this section, we show that the late November climatological short break of PNJ is caused by anomalous upward propagation of planetary waves. To investigate the dynamical cause of the short break in late November, we compared the time series of the PNJ index (Fig. 2a) with the intraseasonal variation of each term of the TEM equation (Appendix A, Equation (A3)) (Fig. 2b). Here, Term A is the temporal tendency of the



PNJ (i.e., its zonal acceleration); Term B is the Coriolis force acting on the residual mean meridional circulation and the meridional advection of zonal momentum; and Term C is the EP flux divergence at 50 hPa averaged over 60–70°N. The temporal tendency of the zonal wind accords well with the sum of the forcing terms of the TEM momentum diagnostic ($A = B + C$; Appendix A). The EP flux divergence (wave forcing) generally governs the zonal wind tendency, and the short break of the PNJ in November is also caused by wave forcing (Term C in equation (A3)).

The vertical component of the EP flux (F^z) at 100 hPa, averaged over latitudes 50–70°N (Fig. 2c), is used as a measure of planetary-scale Rossby wave propagation into the stratosphere (e.g., Coy et al., 1997; Pawson and Naujokat, 1999; Newman et al., 2001). In late November, the upward EP flux at 100 hPa rapidly increases to its maximum, and this enhanced EP flux is linked to the EP flux convergence at 50 hPa, which brings about the short break in the seasonal evolution of PNJ.

3.3 Calculation of anomalous fields with respect to a linear seasonal evolution in late November

We identified a period between 16 and 30 November for the late-November short break (see Fig. 1). We further defined a climatological meteorological field deviation, \mathcal{A}_{dev} , during the period of the short break as a deviation from the expected linear seasonal evolution of that field (Fig. 3):

$$\mathcal{A}_{dev} = \mathcal{A}_{16-30Nov} - (\mathcal{A}_{1-15Nov} + \mathcal{A}_{1-15Dec})/2, \quad (1)$$

where \mathcal{A} is a climatological meteorological field (e.g., geopotential height) and subscripts indicate the averaging period. $\{(\mathcal{A}_{1-15Nov} + \mathcal{A}_{1-15Dec})/2\}$ is the expected climatological meteorological field during late November given a linear seasonal evolution, and \mathcal{A}_{dev} is the deviation of the actual climatological meteorological field in late November from the expected climatological meteorological field. All anomalous fields during the short break were calculated in this manner (see Figs. 4, A1d, A3d, A4d, A5d, and A6d).

The climatological meridional structures of the EP flux and zonal wind from November to early December are shown in Figs. A1a–c. Deviations of meteorological fields, that is, those that deviate from the expectation of a linear seasonal evolution (see Fig. 3), in late November are also shown in Fig. A1d. An upward EP flux propagation deviation (vectors in Fig. A1d) is seen at 50–80°N from the upper troposphere (300 hPa) through the stratosphere (above 100 hPa) during late November. This flux deviation causes an EP flux convergence deviation in the high-latitude stratosphere (contours in Fig. A1d), which corresponds to the short break in the seasonal evolution of the PNJ. This anomalous upward EP flux originates at mid (40–60°N) and high latitudes (65–80°N). The detailed evolution of the climatological EP flux and zonal wind from November to early December is described in Appendix B1. For reference, other climatological atmospheric fields from November to early December are described in Appendices B2, B3, B4, and B5.



3.4 Links between the anomalous upward propagation of the EP flux and a tropospheric trough over eastern Siberia

This section shows that the anomalous (Term \mathcal{A}_{dev} in Eq. 1) upward propagation of planetary waves coincides with a deepening of the eastern Siberia trough (negative deviation of the geopotential height) in late November. To identify the specific area of the anomalous (Term \mathcal{A}_{dev} in Eq. 1) upward propagation of the EP flux during the period of the short break, we investigated the horizontal distribution of the wave activity flux (WAF; Plumb, 1985). The largest positive deviation of the vertical component of the climatological WAF of the stationary wave component at 100 hPa in late November is centered over Siberia and extends over most of the Eurasian continent (Fig. 4). This distribution implies that the Eurasian area is particularly important for stratosphere–troposphere coupling during late November.

During late November, the Rossby wave deviation propagates upward over central Siberia (60–100°E) in the lower troposphere and around East Siberia in the upper troposphere (Fig. A3d). The WAF divergence deviation is negative (figure not shown), indicating convergence in the stratosphere. We further examined the horizontal structure responsible for the upward WAF at 100 hPa. The vertical component of the WAF is proportional to the meridional eddy heat flux ($v'T'$, where prime denotes the anomaly from the zonal mean). Over Siberia, the area of northerly wind and negative air temperature deviations (Fig. A4d) corresponds to the area of positive WAF deviations (Fig. 4). During late November, the trough over Siberia strengthens with time (see Fig. A3d). These results show that the anomalous (Term \mathcal{A}_{dev} in Eq. 1) upward propagation of planetary waves occurs simultaneously with the deepening of the eastern Siberia trough in late November.

3.5 Geopotential height and air temperature in the middle troposphere

In section 3.4, we showed that the deepening of the trough over Siberia is associated with the strengthening the anomalous (Term \mathcal{A}_{dev} in Eq. 1) vertical propagation of planetary waves and the occurrence of the short break of PNJ in late November. In this section, we show that the deepening of the eastern Siberia trough is associated with geopotential height and air temperature deviations. It is generally known that Rossby waves that propagate into the stratosphere in the high latitudes are planetary-scale waves with wavenumbers 1 to 2 (e.g., Baldwin and Dunkerton, 1999). Here, to identify the source of the deviations, we consider the planetary-scale wave components (i.e., wavenumbers 1 to 2) of geopotential height and air temperature in the troposphere. During late November, deviations of eddy geopotential height at 500 hPa (Z500) are strongly negative over Siberia, whereas they are strongly positive over the Atlantic Ocean (Fig. A5d). This positive–negative contrast means that the climatological trough over Siberia is strengthened and the planetary-scale eddy at Z500 is amplified at high latitudes. Cold deviations of eddy air temperature at 850 hPa (T850) are also seen over Siberia along the Arctic Ocean coast (Fig. A6d), west of the negative geopotential deviation (Fig. A5d). The area of these cold deviation is included in the northerly wind deviation area. Where these areas coincide, the eddy meridional heat flux ($v'T'$) is enhanced. A similar but small enhancement of $v'T'$ is also seen over Greenland, where a positive T deviation is observed (Fig. A6d), and over the North Atlantic Ocean, where a positive geopotential deviation is observed (Fig. A5d).



4 Discussion

Why does the atmospheric trough strengthen over Siberia at this time of the year? We hypothesize that a high-latitude land–sea thermal contrast strengthens the trough. Figure 5 shows the time series of the climatological 15-day running means of Z500 and T850 over Siberia (60–170°E, 50–75°N; inside the brown box in Figs. A5 and A6) and outside of Siberia (170°E–60°W, 50–70°N; inside the blue box in Figs. A5 and A6). The time series of the differences between inside and outside of Siberia (green lines) are also shown. During late November, the rate of increase in the zonal contrast (wave amplitude) of Z500 reaches a maximum (green line in Fig. 5a). Similarly, the rate of increase in the zonal T850 contrast, which roughly corresponds to a high-latitude land–sea thermal contrast, approaches a maximum during late November (green line in Fig. 5b). Siberia is of course a land region whereas the area outside of Siberia is occupied mainly by oceans, in particular, the North Atlantic Ocean. Therefore, we hypothesize that thermal forcing due to the land–sea contrast results in the amplification of the trough over Siberia. It is generally known that there are three main sources of the stationary waves that are responsible for zonally asymmetric circulation in the NH: a land–sea thermal contrast, large-scale orography, and tropical diabatic heating (e.g., Smagorinsky, 1953; Inatsu et al., 2002). Large-scale orography (in the NH, the Himalayas, and Rockies in particular) has been found by many studies to be an important source of planetary waves (e.g., Held et al., 2002; Chang, 2009; Saulière et al., 2012). We demonstrated here that the source of the climatological planetary wave in the troposphere during late November is at higher latitude than the Himalayas (see Figs. 4, A5d, and A6d). Strengthening of the high-latitude land–sea thermal contrast may mainly account for the short break in the climatological PNJ during late November. We did not find any short breaks of the PNJ in the Southern Hemisphere (figure not shown). The absence of a short break in the Southern Hemisphere is logically consistent with our hypothesis, because there are no high-latitude zonal land–sea thermal contrasts there.

The short break of the PNJ is possibly due to the occurrence of SSW events. SSW can occur even in November, although its amplitude is smaller in November than in late winter (Charlton and Polvani, 2007; Maury et al., 2016). The relationship between the short break of the PNJ and SSW in early winter is discussed in detail in Appendix C.

5 Conclusions

We found a short break in the seasonal evolution of climatological PNJ during late November (see Fig. 1). Examination of the atmospheric dynamical balance showed that an increase in upward propagation of climatological planetary waves from the troposphere to the stratosphere in late November is accompanied by convergence of the climatological EP flux in the stratosphere, which brings about this short break in the climatological PNJ (see Fig. 2). The upward propagation of Rossby (planetary) waves over Siberia from the troposphere to the stratosphere is a dominant cause of the late November short break (see Fig. 4). This upward propagation of planetary-scale Rossby waves at high latitudes is associated with amplification of climatological eddy geopotential height and air temperature, that is, with a strengthening of the trough over Siberia.



Further, we inferred that this strengthening of the trough is forced by the high-latitude land–sea thermal contrast around Siberia (see Fig. 5). Influence of the November short break upon tropospheric extreme weather and climate remains to be examined. The short break in the climatological PNJ during late February and its long-term trend are interesting topics for future research. SSW events have occurred almost every year since 2000 (Reichler et al., 2012), and this recent frequent occurrence of SSW events might be associated with Arctic sea ice losses (e.g., Kim et al., 2014; Nakamura et al., 2015; Hoshi et al., 2017). In the future, the typical short break of the PNJ should be examined in a composite analysis as well.

Acknowledgments

We deeply thank Dr. Kunihiro Kodera for very insightful discussions. Students in the Weather and Climate Dynamics Division offered us fruitful advice. This study was supported by the Ministry of Education, Culture, Sports, Science and Technology (MEXT) through a Grant-in-Aid for Scientific Research on Innovative Areas (Grant Number 22106003), the Green Network of Excellence (GRENE) Program Arctic Climate Change Research Project, the Arctic Challenge for Sustainability (ArCS) Project, and Belmont Forum InterDec Project. The work of M. Ogi was supported by the Canada Excellence Research Chairs (CERC) Program. The Grid Analysis and Display System (GrADS) was used to draw the figures.

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Figure 1. (a) Latitude–time cross section of the 15-day running mean of the climatological zonal-mean zonal wind at 50 hPa ($\bar{U}50$; lines and color shading) from 1 October through 31 March. The contour interval is 2.0 m s^{-1} . (b) Time series of the climatological 15-day running mean of the PNJ index (m s^{-1} , blue), defined as zonal-mean zonal wind speed at 65°N and 50 hPa. Dark gray shading indicates the 20th to 80th percentiles, medium gray shading indicates the 10th to 90th percentiles, and light gray shading indicates the 5th to 95th percentiles. The two dashed lines indicated the daily minimum and maximum of the PNJ. The vertical black dotted lines indicate the period of the short break in late November.

Figure 2. Time series of climatological 15-day running means of the (a) PNJ index (m s^{-1} , dark blue) (same as Fig. 1b); (b) the temporal tendency of the PNJ (i.e., zonal acceleration [$\text{m s}^{-1} \text{ day}^{-1}$]; Term A in equation (A3), light blue line), the Coriolis force acting on the residual mean meridional circulation and the meridional advection of zonal momentum ($\text{m s}^{-1} \text{ day}^{-1}$, Term B in equation (A3), purple line) at 50 hPa, and the EP flux divergence ($\text{m s}^{-1} \text{ day}^{-1}$), Term C in equation (A3), red line) at 50 hPa averaged over latitudes $60\text{--}70^\circ\text{N}$; and (c) the vertical component of EP flux (F^z) at 100 hPa ($\text{m}^2 \text{ s}^{-2}$) averaged over latitudes $50\text{--}70^\circ\text{N}$ from 1 October through 31 March. The vertical black dotted lines indicate the period of the late November short break.

Figure 3. Schematic diagram of a late November deviation in the seasonal evolution of a climatological meteorological field. Red stars indicate the values of the meteorological field in early November and early December, the red cross mark indicates its actual value in late November, and the blue cross mark indicates the mean of the early November and early December values. The vertical difference between the actual value (red cross mark) and the expected value (blue cross mark) during late November, which is calculated by equation (1), is the field deviation.

Figure 4. Vertical component of the late November deviation of the climatological WAF (Plumb, 1985) at 100 hPa ($10^{-3} \text{ m}^2 \text{ s}^{-2}$) with respect to its linear seasonal evolution, calculated by equation (1) (see Section 3.3). The blue box ($0\text{--}360^\circ\text{E}$, $50\text{--}70^\circ\text{N}$) indicates the averaging area used to calculate the fields shown in Fig. 2c.

Figure 5. Time series of the climatological 15-day running mean (a) Z500 (m) and (b) T850 ($^\circ\text{C}$) over Siberia ($60\text{--}170^\circ\text{E}$, $50\text{--}75^\circ\text{N}$; brown lines), outside Siberia ($170^\circ\text{E}\text{--}60^\circ\text{W}$, $50\text{--}75^\circ\text{N}$; blue lines), and their anomalies within Siberia from their values outside of Siberia (green lines) from 1 October through 31 March. The vertical black dotted lines indicate the period of the late November short break.

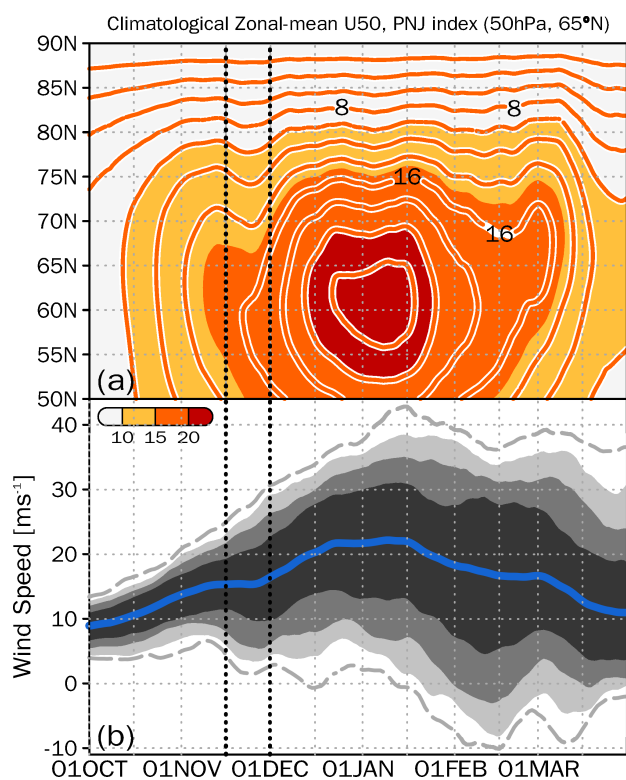


Figure 1. (a) Latitude–time cross section of the 15-day running mean of the climatological zonal-mean zonal wind at 50 hPa (\bar{U}_{50} ; lines and color shading) from 1 October through 31 March. The contour interval is 2.0 m s^{-1} . (b) Time series of the climatological 15-day running mean of the PNJ index (m s^{-1} , blue), defined as zonal-mean zonal wind speed at 65°N and 50 hPa. Dark gray shading indicates the 20th to 80th percentiles, medium gray shading indicates the 10th to 90th percentiles, and light gray shading indicates the 5th to 95th percentiles. The two dashed lines indicated the daily minimum and maximum of the PNJ. The vertical black dotted lines indicate the period of the short break in late November.

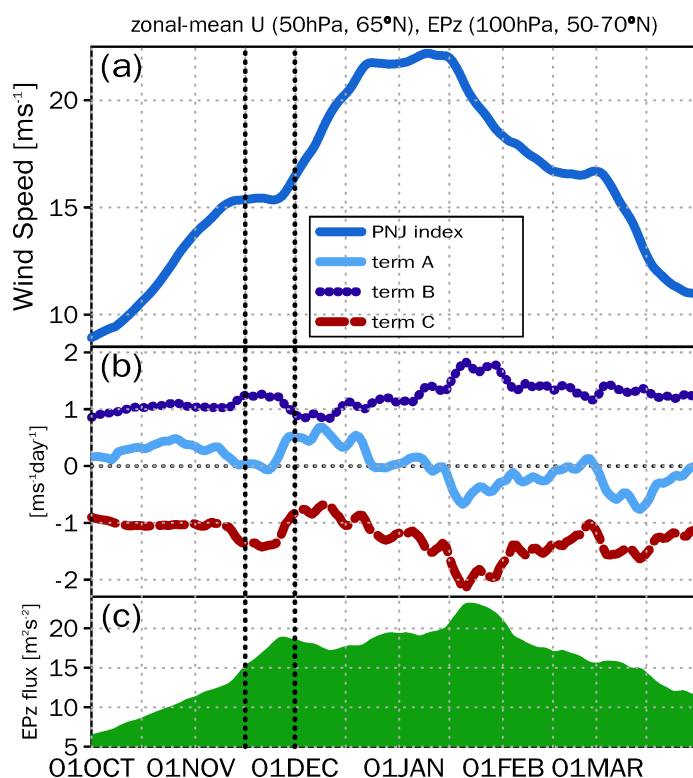


Figure 2. Time series of climatological 15-day running means of the (a) PNJ index (m s^{-1} , dark blue) (same as Fig. 1b); (b) the temporal tendency of the PNJ (i.e., zonal acceleration [$\text{m s}^{-1} \text{ day}^{-1}$]; Term A in equation (A3), light blue line), the Coriolis force acting on the residual mean meridional circulation and the meridional advection of zonal momentum ($\text{m s}^{-1} \text{ day}^{-1}$, Term B in equation (A3), purple line) at 50 hPa, and the EP flux divergence ($\text{m s}^{-1} \text{ day}^{-1}$, Term C in equation (A3), red line) at 50 hPa averaged over latitudes 60–70°N; and (c) the vertical component of EP flux (F^z) at 100 hPa ($\text{m}^2 \text{ s}^{-2}$) averaged over latitudes 50–70°N from 1 October through 31 March. The vertical black dotted lines indicate the period of the late November short break.

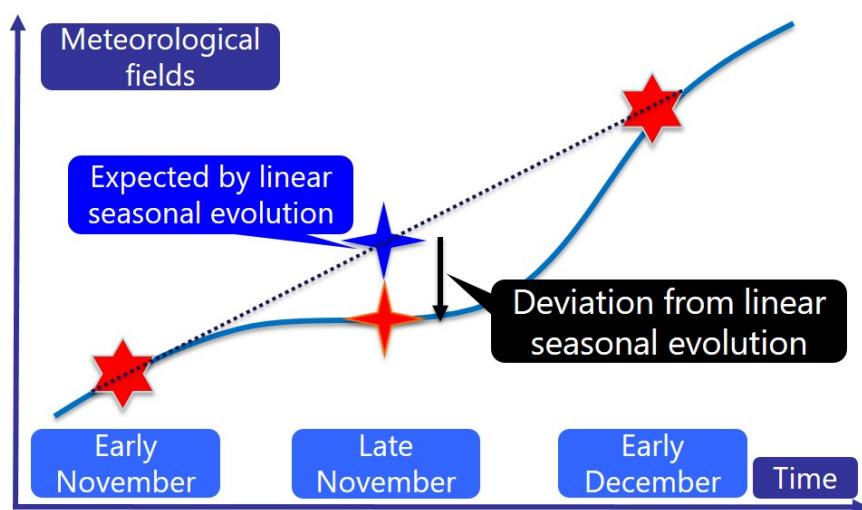


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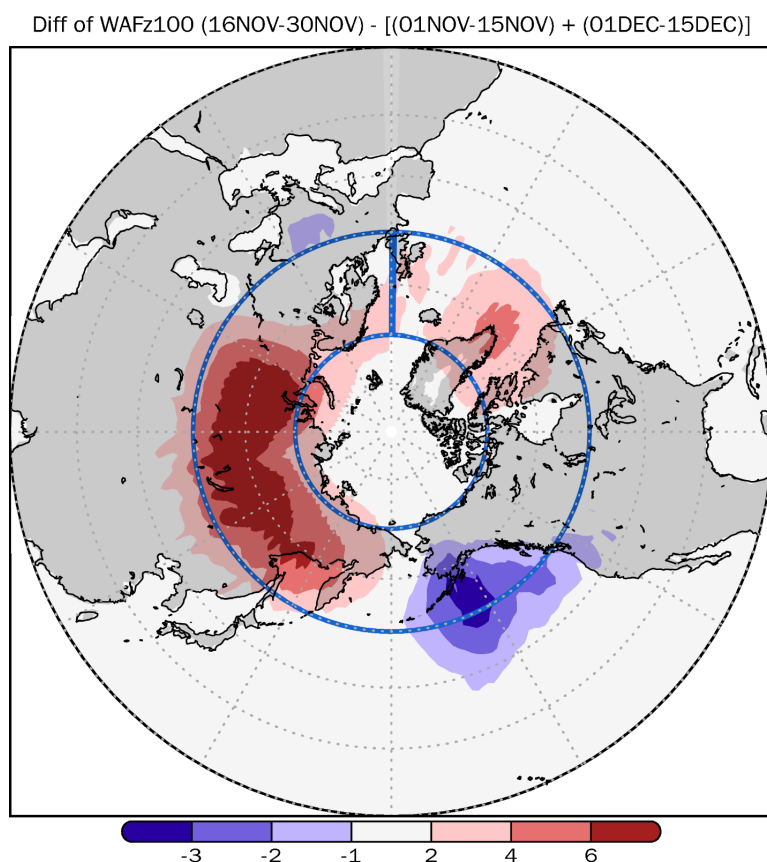


Figure 4. Vertical component of the late November deviation of the climatological WAF (Plumb, 1985) at 100 hPa ($10^{-3} \text{ m}^2 \text{ s}^{-2}$) with respect to its linear seasonal evolution, calculated by equation (1) (see Section 3.3). The blue box ($0\text{--}360^\circ\text{E}$, $50\text{--}70^\circ\text{N}$) indicates the averaging area used to calculate the fields shown in Fig. 2c.

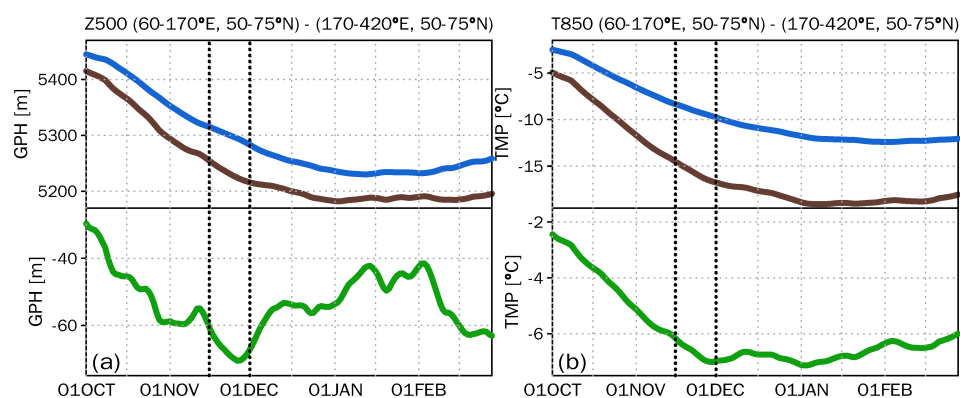


Figure 5. Time series of the climatological 15-day running mean (a) Z500 (m) and (b) T850 (°C) over Siberia (60–170°E, 50–75°N; brown lines), outside Siberia (170°E–60°W, 50–75°N; blue lines), and their anomalies within Siberia from their values outside of Siberia (green lines) from 1 October through 31 March. The vertical black dotted lines indicate the period of the late November short break.



Appendices

Appendix A Transformed Eulerian Mean (TEM) Diagnostics

Eliassen-Palm (EP) flux analysis is widely used in dynamic meteorology to diagnose wave and zonal-mean flow interactions. The EP flux shows the propagation of Rossby (planetary) waves (Andrews and McIntyre, 1976). The meridional (F^ϕ) and vertical (F^z) components of the EP flux (\mathbf{F}) are defined as follows:

$$F^\phi \equiv \rho_0 a \cos \phi \left[(\partial \bar{u} / \partial z) \overline{v' \theta'} / \bar{\theta}_z - \overline{u' v'} \right] \quad (\text{A1})$$

$$F^z \equiv \rho_0 a \cos \phi \left\{ [f - (a \cos \phi)^{-1} \partial (\bar{u} \cos \phi) / \partial \phi] \overline{v' \theta'} / \bar{\theta}_z - \overline{w' u'} \right\}, \quad (\text{A2})$$

where a is the radius of the Earth, f is the Coriolis parameter, ϕ is latitude, θ is potential temperature, u is zonal wind, and v is meridional wind. Overbars denote zonal means, primes denote anomaly from the zonal mean, z is a log-pressure coordinate, and ρ_0 is air density. $\bar{\theta}_z = \partial \bar{\theta} / \partial z$ is computed from the zonal mean of the potential temperature in log-pressure coordinates. The eddy-flux terms $u'v'$ and $v'\theta'$ are computed from the zonal anomalies in the 6-hourly data, and the product is zonally averaged and then time averaged to obtain 15-day means.

We used the primitive form of the Transformed Eulerian Mean (TEM) momentum equation to examine the diagnostics of the zonal-mean momentum (e.g., Andrews et al., 1987; Holton and Hakim, 2012; Vallis, 2017):

$$\underbrace{\partial \bar{u} / \partial t}_{(\text{A})} = \underbrace{\bar{v}^* \{ f - (a \cos \phi)^{-1} \partial (\bar{u} \cos \phi) / \partial \phi \}}_{(\text{B})} - \underbrace{\bar{w}^* \partial \bar{u} / \partial z + (\rho_0 a \cos \phi)^{-1} \nabla \cdot \mathbf{F} + \bar{X}}_{(\text{C})}, \quad (\text{A3})$$

where \bar{v}^* and \bar{w}^* are the meridional and vertical components of the residual mean meridional circulation, \bar{X} is a residual Term that includes internal diffusion and surface friction as well as sub-grid scale forcing such as gravity wave drag. Term A in equation (A3) is the temporal tendency of the zonal-mean zonal wind, Term B is the Coriolis force acting on the residual mean meridional circulation and the meridional advection of zonal momentum, and Term C is the divergence of the EP flux vector, i.e., wave forcing.

The vertical component of the wave activity flux (WAF; Plumb, 1985) at 100 hPa provides a useful diagnostic for identifying the source region of vertically propagating stationary planetary waves. The vertical component of the WAF is proportional to the vertical component of the EP flux. The eddy terms are computed from the zonal deviations relative to each 15-day mean.



Wavenumber decomposition was carried out by applying Fourier analysis to the geopotential height and air temperature fields.

Appendix B Climatological fields from early November to early December and their late November deviations

Appendix B1 Zonal mean zonal wind, EP flux, and EP flux divergence

Figures A1a–c show the climatological zonal mean zonal wind, EP flux, and EP flux divergence in the NH during (a) early November, (b) late November, and (c) early December. The subtropical jet is commonly centered in the upper troposphere at 35°N, 200 hPa, and the PNJ is centered in the stratosphere at 65°N. These two westerly maxima gradually strengthen with time. The EP flux propagates upward from the lower troposphere to the mid- and upper troposphere in the low latitudes, and it propagates into the stratosphere in the high latitudes. The EP flux also gradually propagates upward with time. Figure A1d shows the departures of the fields shown in Fig. A1b from the average of the fields shown Figs. A1a and A1c (calculated by equation (1)). Thus, Fig. A1d shows the late November field deviations from a linear seasonal evolution.

10 Appendix B2 Vertical component of the wave activity flux of the stationary wave component at 100 hPa

Figures A2a–c show the vertical component of the climatological wave activity flux (WAF) of the stationary wave component at 100 hPa during early November, late November, and early December, respectively. During all three periods, a strong positive signature is centered in the Russian far east and extends from eastern Europe to the east coast of Asia.

Appendix B3 Eddy component of geopotential height and zonal and vertical components of the WAF averaged over 50–70°N

15 Figures A3a–c show the eddies (anomalies from the zonal mean) of climatological geopotential height and the zonal and vertical components of the climatological WAF distribution, averaged over 50–70°N (inside the blue box in Fig. A2) during early November, late November, and early December, respectively. Over East Siberia (100–120°E), an area of strong negative eddies (i.e., a geopotential height trough) extends from the middle troposphere to the stratosphere with a westward-upward tilt, and an area of positive eddies (i.e., a ridge) occurs near the surface over East Siberia (i.e., the area of Siberian High). Over 180°E–120°W, there is an area of strong positive anomalies in the stratosphere (i.e., the Aleutian High). Rossby waves propagate upward over East Siberia from the lower troposphere to the upper troposphere. Figure A3d shows the late November anomalies. Note that the WAF was calculated with equation (1), not from the zonal anomalies of climatological geopotential height shown in Fig. A3d.



Appendix B4 Zonal anomalies of meridional wind and air temperature at 100 hPa

Figures A4a–c show the zonal anomalies of climatological meridional wind and air temperature at 100 hPa during early November, late November, and early December, respectively. During all three periods, northerly winds and negative air temperatures occur over Siberia and southerly winds and positive air temperatures occur over the northwest Pacific Ocean. This collocation corresponds to the area of positive anomalies of the WAF over Siberia (see Figs. A2a–c).

Appendix B5 Geopotential height and air temperature in middle troposphere

Figures A5a–c show the climatological eddy geopotential height at 500 hPa (Z500) during early November, late November, and early December, respectively. Negative anomalies (trough) are seen from East Siberia to East Asia, whereas positive anomalies (ridge) are over the North Atlantic Ocean to Europe. Figure A5d shows the planetary-scale eddy geopotential height deviation at 500 hPa. Figure A6 is the same as Fig. A5, but for air temperature at 850 hPa (T850). Negative anomalies (cold air) are seen over East Siberia to East Asia, and positive anomalies (warm air) are apparent over the North Atlantic Ocean to Europe (Figs. A6a–c).

Appendix C. Relationship between the early winter (November) short break of the PNJ and SSW events

As illustrated in Fig. 1b, the interannual spread of the PNJ is larger during January and February than from October to November. The larger spread in winter can be attributed to the fact that most SSW events occur during January and February, and SSW event magnitudes are also large during this period (Charlton and Polvani, 2007; Maury et al., 2016). The upper percentile wind speed values do not show a short break of the PNJ during November, whereas they do show the late February short break. This result is consistent with the fact that SSW events are rare and their amplitudes are smaller in November than in late winter.

Manney et al. (2001) indicated that minor early winter SSWs that occurred in November 2000, which they called Canadian Warmings (CWs; Labitzke, 1977, 1982), may have had a profound impact on the development of a vortex and a low-temperature region in the lower stratosphere. Waugh and Randel (1999) presented an overview of climatological stratospheric polar vortices, including during the early winter period, and examined them by an elliptical diagnostics analysis. Elliptical diagnostics define the area, center, elongation, and orientation of each vortex and are used to quantify their structure and evolution. They found that the PNJ in the NH becomes more distorted and its position shifts away from the pole from October through December. They also recognized a climatological southward shift of the center of the polar vortex in the NH in late November (Fig. 4d in Waugh and Randel, 1999). This finding is consistent with our study results.

The shift recognized by Waugh and Randel (1999) may be related to the occurrence of wavenumber 1-type minor SSW events (CWs) in late November (Labitzke and Naujokat, 2000; Manney et al., 2001), and these CWs might affect the short break in the seasonal evolution of



climatological PNJ during late November. Moreover, Maury et al. (2016) defined an SSW event in terms of a comprehensive description of stratospheric warming events without a priori distinctions between major and minor events. Small-amplitude warmings occur during late November but not in early November or in December. Therefore, the late November climatological short break of the PNJ may be related to early winter SSW events. Our study results are consistent with these previous studies.

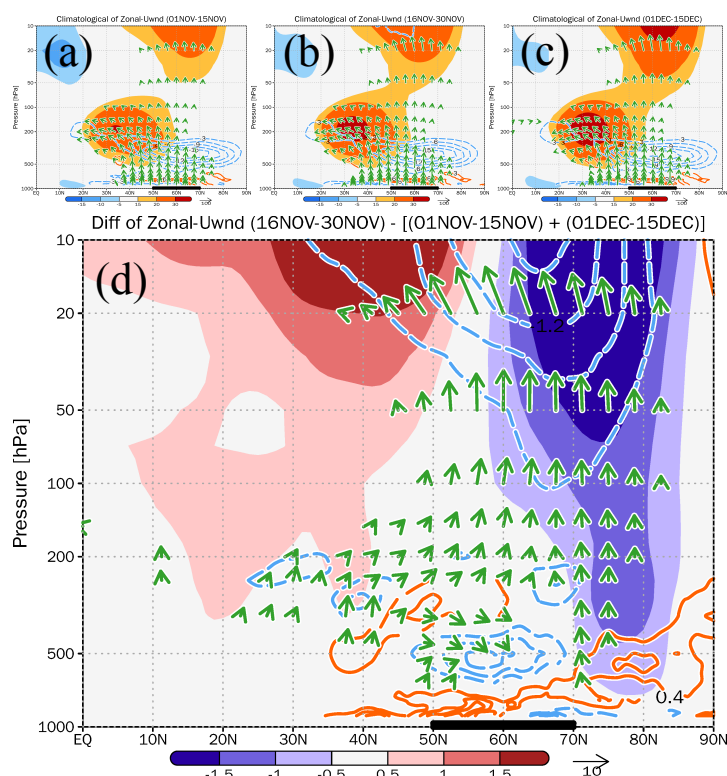


Figure A1. Climatological zonal-mean zonal wind speed (m s^{-1} , color shading), EP flux ($\text{m}^2 \text{s}^{-2}$, vectors), and the flux divergence ($\text{m s}^{-1} \text{day}^{-1}$, contours) during (a) early November (1–15 November), (b) late November (16–30 November), and (c) early December (1–15 December). (d) Late November deviations (late November deviations from the mean of the early November and early December values calculated with equation (1); see Section 3.3). The EP flux is standardized by density (1.225 kg m^{-3}) and the radius of the Earth ($6.37 \times 10^6 \text{ m}$). The vertical component of the vectors is multiplied by a factor of 250. The bold black line indicates the longitudinal range for Siberia (50–70°N).

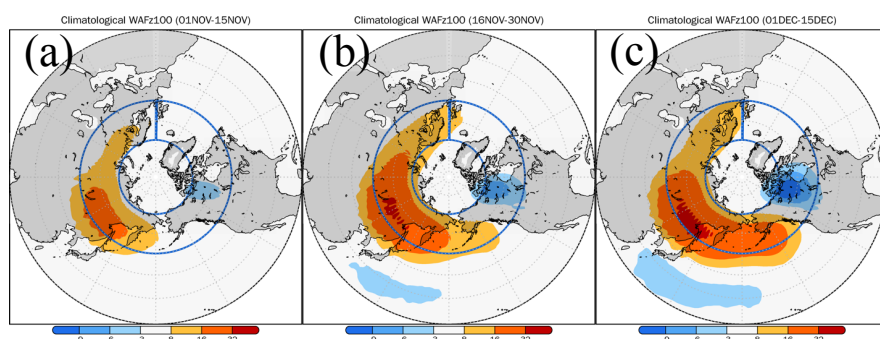


Figure A2. Vertical component of the climatological wave activity flux (Plumb, 1985) at 100 hPa ($10^{-3} \text{ m}^2 \text{ s}^{-2}$) during (a) early November, (b) late November, and (c) early December. The box outlined in blue (0–360°E, 50–70°N) indicates the averaging area used for calculating the fields shown in Fig. A3.

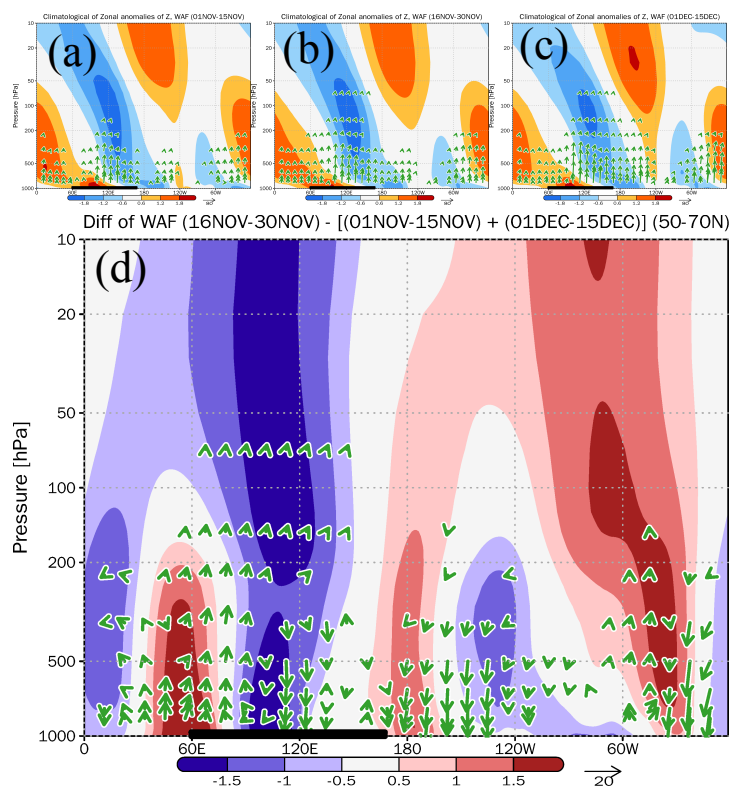


Figure A3. Zonal anomalies of climatological geopotential height (m, color shading) and zonal and vertical components of WAF ($10^{-3} \text{ m}^2 \text{ s}^{-2}$, vectors), averaged over latitude $50\text{--}70^\circ\text{N}$ (inside the blue box in Fig. A2) during (a) early November, (b) late November, and (c) early December. (d) Late November field deviations calculated by equation (1) (see Section 3.3). The geopotential height is normalized by the standard deviation at each height. The WAF magnitude is standardized by pressure (p/p_s , p_s is a standard sea-level pressure) and the square of the radius of the Earth ($6.37 \times 10^6 \text{ m}$). The vertical components of the vectors are multiplied by a factor of 500. The black line indicates the latitudinal range for Siberia ($60\text{--}170^\circ\text{E}$).

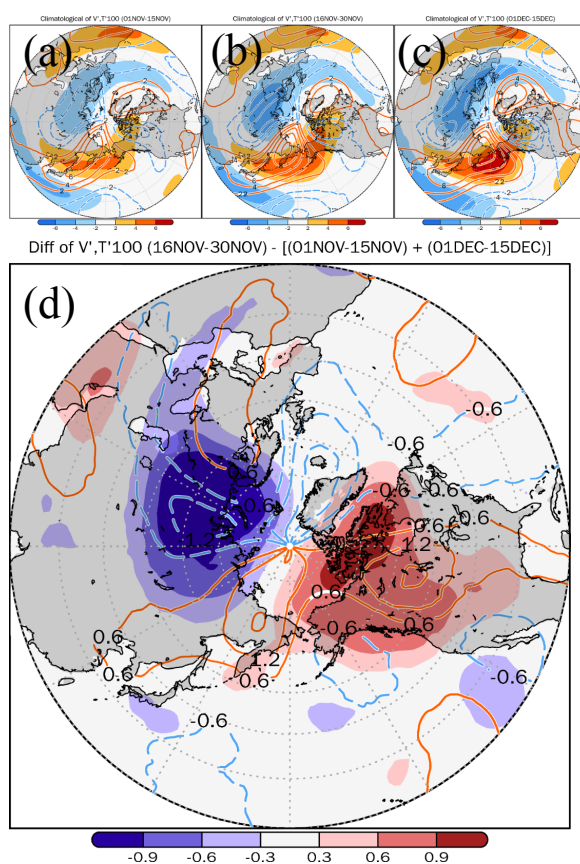


Figure A4. Zonal anomalies of climatological meridional wind (m s^{-1} , contours) and air temperature ($^{\circ}\text{C}$, color shading) at 100 hPa during (a) early November, (b) late November, and (c) early December. (d) Late November deviations calculated by equation (1) (see Section 3.3).

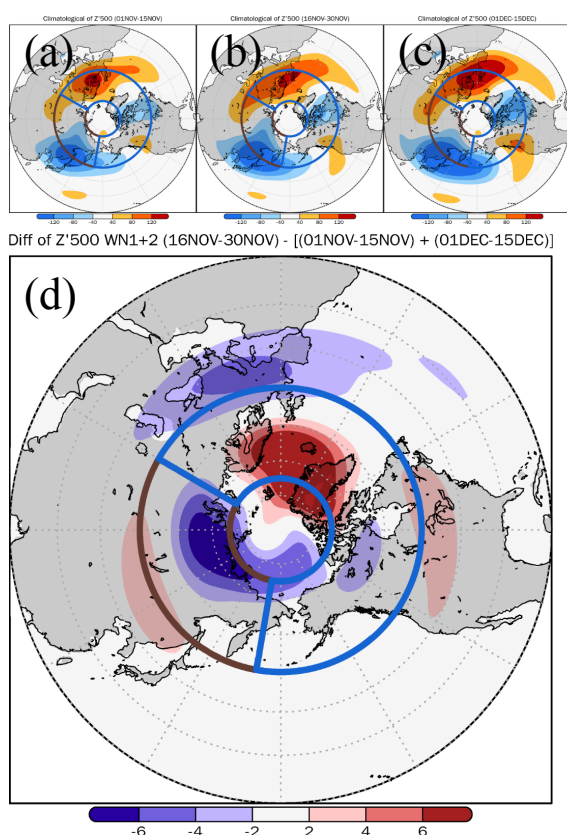


Figure A5. Zonal anomalies of climatological geopotential height at 500 hPa (m) during (a) early November, (b) late November, and (c) early December. (d) Late November deviations calculated by equation (1) (see Section 3.3) with wavenumber decomposition; only planetary-scale components, wavenumbers 1 to 2, were used. The brown (60–170°E, 50–75°N) and blue (170°E–60°W, 50–75°N) boxes indicate the averaging areas used for calculating the fields shown in Fig. 5.

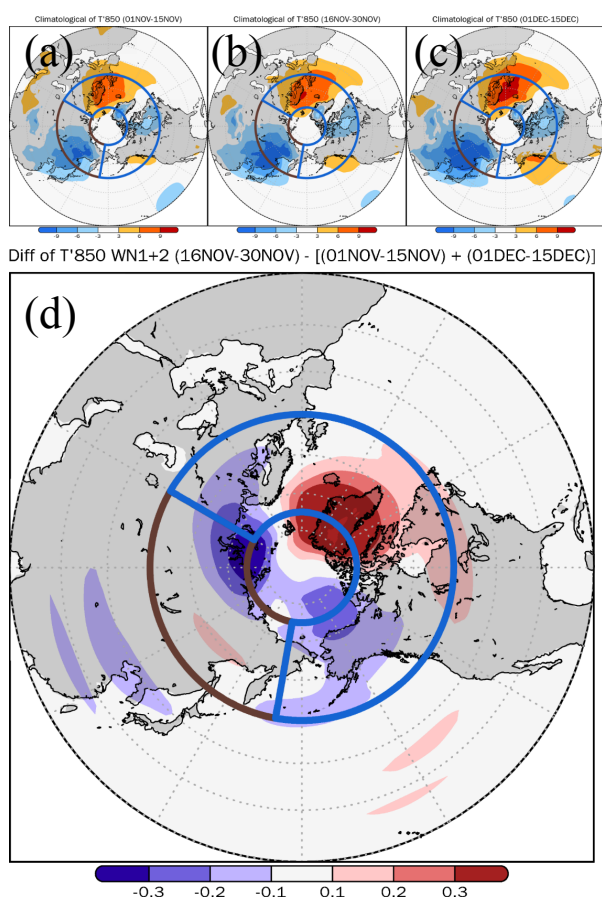


Figure A6. Same as Fig. A5, but for air temperature at 850 hPa (°C).