



1 **On the Freezing Time of Supercooled Drops in Developing**
2 **Convective Clouds**

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12 **Abstract**

13 In this study, the particle size distributions (PSDs) measured in fresh developing maritime
14 convective clouds sampled during the Ice in Clouds-Tropical (ICE-T) project are shown and
15 compared with the PSDs modeled using a parcel model containing a spectral bin microphysics
16 scheme. The observations suggest that the "first ice" in convective clouds is small. To interpret
17 the observed ice PSDs, the freezing times and temperatures of supercooled drops are analyzed.
18 The results indicate that the freezing time is longer for large drops than it is for small drops. Due



19 to instrumental limitations, freezing drops cannot be identified until they exhibit obvious shape
20 deformation. If the updraft is strong enough, large freezing drops can be carried upwards to a
21 lower temperature than their nucleation temperature before obvious shape deformation occurs. In
22 models, drop freezing is assumed to be instantaneous, which is not realistic; thus, the model
23 yields a broader "first ice" PSD than is observed. This study allows us to interpret the observed
24 ice PSDs in fresh developing convective clouds from the perspective of the freezing time of
25 supercooled drops and notes the deficiency of instantaneous drop freezing in models. To better
26 understand the mechanisms of drop freezing and ice initiation in convective clouds, more
27 laboratory experiments and in situ measurements are needed in the future.

28 **1. Introduction**

29 Ice initiation in convective clouds is still not well understood (Cantrell and Heymsfield, 2005;
30 Lawson et al., 2015; Yang et al., 2016), and it remains one of the main sources of uncertainties in
31 numerical models (Khain et al., 2015). Observations suggest that ice initiation in convective
32 clouds is strongly related to the freezing of supercooled drops (Rangno and Hobbs, 2005;
33 Lawson et al., 2015; Yang et al., 2016; Field et al., 2017). Supercooled drops do not fully freeze
34 instantaneously, and during airborne measurements, freezing drops cannot be observed until they
35 have experienced obvious deformation. The freezing rate of supercooled drops depends on the
36 rate of heat transfer between the drop and ambient air (Pruppacher and Klett, 1997). Typically,
37 the freezing process comprises four stages (Hindmarch et al., 2003): 1) the supercooling stage,
38 during which a drop is supercooled to its nucleation temperature; 2) the recalescence stage,
39 during which rapid kinetic ice nucleation occurs, which results in a rapid drop in temperature that
40 is terminated when the drop temperature reaches 0 °C; 3) the freezing stage, during which the
41 liquid part of a drop continuously freezes and the drop temperature remains at 0 °C; and 4) the



42 cooling stage, during which the frozen drop cools to the ambient temperature.

43 A number of laboratory experiments have been performed to study the freezing of supercooled
44 drops. For example, Johnson and Hallett (1968) showed that the freezing time of supercooled
45 drops decreases with decreasing ambient temperature. In typical air conditions, it takes
46 approximately 400 s for a stationary millimeter-sized drop to completely freeze at $-5\text{ }^{\circ}\text{C}$ under a
47 pressure of 1 atm; this freezing time is reduced to approximately 200 s at $-10\text{ }^{\circ}\text{C}$. They also
48 showed that the freezing rate of supercooled drops is related to the composition of air and that
49 the freezing time of a millimeter-sized drop in helium and hydrogen is only one-fifth of that in
50 air. Hindmarsh et al. (2003) showed that the freezing time increases with increasing drop size. In
51 addition, Hindmarsh et al. (2003) used experimental results to discuss the accuracy of three drop
52 freezing models: the uniform temperature model, the inward freezing model and the outward
53 freezing model. All three of these models have fairly good accuracy in modeling drop
54 temperatures and freezing times, and there are only minor differences between them.

55 In most numerical weather prediction models (NWPMs) and global climate models (GCMs), the
56 freezing of supercooled drops is assumed to be instantaneous, because it is difficult to track the
57 freezing stage of every particle in models and because there are no good observations with which
58 to evaluate the modeled ice microphysics in detail. Phillips et al. (2015) implemented time-
59 dependent freezing for raindrops in a cloud model using spectral bin microphysics (SBM). Their
60 sensitivity tests showed that time-dependent drop freezing delays the formation of hail in
61 convective clouds; however, their model was unable to track the freezing stage of every particle.
62 Using a simplified cloud parcel model and an electromagnetic scattering model, Kumjian et al.
63 (2012) showed that the modeled radar polarimetric variables for convective clouds are more
64 consistent with observations if time-dependent drop freezing is considered. However, drop



65 freezing in fresh developing convective clouds has rarely been discussed. Thus, to better
66 understand ice initiation in convective clouds and to evaluate the modeled microphysics, more
67 observations are needed.

68 Aircraft in situ measurements are necessary to improve our current understanding of ice
69 initiation in convective clouds and to evaluate model simulations. Traditional in situ
70 measurements can rarely identify ice that is smaller than $200 \mu\text{m}$ in diameter. The 3-View Cloud
71 Particle Imager (3V-CPI) is a good tool with which to record small particle images, and it can be
72 used to identify small ice (Field et al., 2017). During the Ice in Clouds-Topical (ICE-T) project,
73 the 3V-CPI that was operated on the SPEC Learjet yielded high-resolution particle images and
74 particle size distributions (PSDs). The 3V-CPI measurements suggest that the observed "first
75 ice" in fresh developing convective clouds are all small ice (Lawson et al., 2015); however, the
76 results of some other studies have suggested that larger supercooled drops may freeze before
77 smaller drops (Bigg, 1953; Heymsfield, 2013). This raises the question: why is the observed
78 "first ice" in convective updrafts small? Understanding the freezing time of supercooled drops is
79 helpful for interpreting the observed ice PSDs in developing convective clouds. In addition,
80 determining the size of "first ice" is important for understanding secondary ice generation
81 process(es). This study aims to analyze the observed PSDs in developing convective clouds
82 using the data collected during the ICE-T project, as well as to interpret these observations
83 through the perspective of the freezing time of supercooled drops. This paper is organized as
84 follows: Section 2 introduces the dataset and the analytical method; Section 3 discusses the
85 results; and a summary is given in Section 4.

86 2. Dataset and Analysis Method



87 **2.1 Calculation of the freezing time of supercooled drops**

88 The calculation of the freezing time and temperature of supercooled drops is governed by a
89 series of heat transfer and phase change equations. These detailed equations have been described
90 in previous studies (e.g., Dye and Hobbs et al., 1968; Hindmarsh et al., 2003). The drop
91 temperature change is balanced by convective heat transfer (i.e., ventilation), radiation and latent
92 heat terms. In this calculation, a supercooled drop is assumed to be carried upward by an updraft,
93 which ascends adiabatically. The terminal velocity of the drop follows that defined by Foote and
94 Du Toit (1969). In this calculation, diffusional growth is included but coalescence is neglected.
95 The initial drop temperature is the same as the ambient air temperature. The temperature inside
96 the drop is assumed to be uniform; this is a reasonable assumption because water and ice have a
97 larger thermal conductivity than air and because of the internal mixing of liquid within the drop
98 (Yao and Schrock, 1976). Hindmarsh et al. (2003) showed that including temperature variations
99 inside the drop has a minor impact on the results. The freezing time is defined as the time period
100 from the start to the end of the drop freezing.

101 **2.2 Aircraft measurements during ICE-T**

102 The ICE-T project was conducted in July 2011 over the Caribbean Sea, near the U.S. Virgin
103 Islands; its goal was to study ice generation in tropical maritime convective clouds. Both the
104 National Center for Atmospheric Research (NCAR) C-130 aircraft and the SPEC Incorporated
105 Learjet were deployed during ICE-T.

106 The SPEC Learjet was equipped with various instruments that were used to study the
107 microphysics in convective clouds during ICE-T. The primary goal of the Learjet was to make
108 repeated penetrations in fresh developing convective updrafts near the cloud top. These



109 instruments included a fast forward-scattering spectrometer probe (FFSSP); a CPI; a two-
110 dimensional stereo (2D-S) probe; a high-volume precipitation spectrometer (HVPS-3), and a
111 Rosemount temperature probe. The measurements obtained using the FFSSP, CPI, 2D-S, and
112 HVPS were combined to generate the PSDs. CPI images were used to identify liquid drops and
113 ice particles that were smaller than $500\ \mu\text{m}$ in diameter, and these percentages of drops and ice
114 particles were applied to the 2D-S PSDs. The 2D-S and HVPS images were used to identify
115 drops and ice particles that were larger than $500\ \mu\text{m}$ in diameter. More information about the
116 processing of the Learjet data can be found in Lawson et al. (2015).

117 The NCAR C-130 was not used to repeatedly penetrate fresh developing convective clouds
118 during ICE-T; instead, it penetrated convective clouds at different stages of their development.
119 Most of these penetrations occurred far below the cloud top, although some were near the cloud
120 top (Heymsfield et al. 2014). The instruments used here included an FFSSP, a two-dimensional
121 cloud (2D-C) probe, a two-dimensional precipitation (2D-P) probe, and a Rosemount
122 temperature probe. The Wyoming Cloud Radar (WCR; Wang et al. 2012) was operated on the C-
123 130 to obtain 2D reflectivity structures, and the Wyoming Cloud Lidar (WCL; Wang et al. 2009)
124 was used to identify liquid-dominated and ice-dominated clouds.

125 **2.3 Parcel model simulation**

126 In this study, we compare the PSDs modeled using a parcel model containing SBM to those
127 observed on the aircraft. The SBM was developed by Hebrew University (Khain et al. 2000) and
128 has been implemented in the Weather Forecast and Research model (WRF; Lynn et al. 2005).
129 Time-dependent drop freezing is not included in this scheme. The purpose of this simulation is
130 not to evaluate the modeled results using observations, but instead to reveal the deficiency of



131 instantaneous drop freezing in SBM and its inability to capture the observed rapid ice generation.
132 The modeled parcel has a depth of 500 m. The observed drop size distribution at $-6\text{ }^{\circ}\text{C}$ is used as
133 an input. The vertical air velocity is 10 m/s, which is a typical mean updraft strength in the
134 convective clouds sampled during ICE-T. The hydrometeor types include cloud drop/rain,
135 ice/snow, and graupel; the PSD of each hydrometeor type has 33 mass bins. The ice nucleation
136 mechanisms include immersion freezing using Bigg's parameterization (1953),
137 deposition/condensation nucleation (Meyer et al., 1992), contact nucleation (Meyer et al., 1992),
138 and the Hallett-Mossop process (Hallett and Mossop, 1974). Other ice microphysics processes
139 include riming, coalescence and diffusional growth. During every time step, 1% of the aerosols
140 in the ambient air are assumed to become entrained into the cloud parcel. The ambient aerosol
141 size distribution is observed using a high-flow dual-channel differential mobility analyzer
142 (HDDMA; DeMott et al., 2016) and a Passive Cavity Aerosol Spectrometer Probe (PCASP;
143 Baumgardner et al., 2011) operated on the C-130.

144 **3. Results and Discussion**

145 **3.1 Comparison of observed and modeled particle size distributions**

146 Fig. 1 shows the size distributions measured by the Learjet and those modeled using a parcel
147 model with SBM. The simulation data on the left panels include all of the ice physics
148 implemented in the SBM, while liquid-ice collision is turned off for the right panels. The Learjet
149 measurements suggest that the ice particles observed in fresh developing convective clouds are
150 relatively small ($20\text{-}300\text{ }\mu\text{m}$ in diameter) between $-7\text{ }^{\circ}\text{C}$ and $-10\text{ }^{\circ}\text{C}$ and that the PSD of ice
151 broadens as the temperature decreases. The modeled ice PSD is much broader than that observed
152 between $-7\text{ }^{\circ}\text{C}$ and $-10\text{ }^{\circ}\text{C}$. The deposition/condensation nucleation exhibits the largest



153 contribution to the modeled ice PSDs (Fig. 1d). Immersion freezing, contact freezing and the
154 Hallett-Mossop process contribute less to the modeled ice PSDs. Small ice particles are mostly
155 formed by deposition/condensation nucleation, whereas large ice is produced by immersion
156 freezing and drop-ice collision (Fig. 1d and h).

157 An obvious difference between the observed and modeled ice PSDs is that large ice is not
158 observed between $-7\text{ }^{\circ}\text{C}$ and $-10\text{ }^{\circ}\text{C}$ but is found in the modeled results (Fig. 1d). There are three
159 possible explanations for this: first, large freezing (or frozen) drops cannot be identified from the
160 images taken by the probes, or the sampling volume of the probes is too small; second, the
161 modeled results are not realistic; third, there could be a combination of the first and second
162 possibilities. Previous studies have suggested that during immersion freezing, large drops have a
163 higher probability of freezing than small drops at the same temperature (Bigg, 1953). In addition,
164 small ice that is generated by other mechanisms (e.g., deposition/condensation nucleation,
165 secondary ice) can be quickly collected by large drops in convective clouds, which results in the
166 freezing of large drops. There is no evidence that large drops do not freeze between $-7\text{ }^{\circ}\text{C}$ and -
167 $10\text{ }^{\circ}\text{C}$. In the observations, only non-spherical particles are regarded as ice, but freezing drops
168 exhibit no (obvious) shape deformation during the early stage of freezing (Johnson and Hallett,
169 1968; Hindmarsh et al., 2003). Due to the limitations of the instruments, freezing drops that do
170 not exhibit obvious shape deformation cannot be identified; thus, the first possibility may apply.
171 On the other hand, in the model simulations, drop freezing is assumed to be instantaneous, which
172 could result in a broad ice PSD at warm temperatures; because this is not true in natural clouds,
173 the second possibility may also apply. Therefore, the large difference between the measured and
174 simulated ice PSDs is probably both observation- and model-related.

175 Examples of particle images collected by the 2D-C on the C-130 and the CPI on the Learjet are



176 shown in Fig. 2. Both the 2D-C and CPI images were measured near the cloud top in the updraft
177 cores of developing convective clouds. As noted in the figure, the observed ice particles mostly
178 comprise small frozen drops between $-8\text{ }^{\circ}\text{C}$ and $-10\text{ }^{\circ}\text{C}$ (Fig. 2c). Some particles may have just
179 begun freezing because they exhibit slight shape deformation, as shown by the particle images in
180 the red box in Fig. 2c; however, we have no other evidence with which to confirm this. Between
181 $-10\text{ }^{\circ}\text{C}$ and $-13\text{ }^{\circ}\text{C}$, we observe more ice particles, including both large and small frozen drops, as
182 well as rimed graupels (Fig. 2a and b). Columns and plates were also observed. Considering the
183 time that is needed for columns and plates to grow, they were probably generated at a warmer
184 temperature than is observed. Due to the relatively low resolution of the 2D-S, 2D-C, HVPS and
185 2D-P images, large freezing (or frozen) drops that exhibit no obvious shape deformation cannot
186 be identified, and they are thus regarded as drops. In some spherical CPI particle images, it is
187 also difficult to determine whether the particles have begun freezing or not, because freezing
188 drops exhibit no (or no obvious) shape deformation during the early stages of freezing (e.g.,
189 Johnson and Hallett, 1968; Hindmarsh et al., 2003).

190 **3.2 Freezing time of supercooled drops**

191 To better understand the observed PSDs, we analyze the freezing time of supercooled drops in
192 this section. Fig. 3 shows the changes in drop temperature and ice mass fraction with changes in
193 time and ambient temperature. The updraft velocity is assumed to be 10 m/s. Drops and air
194 parcels ascend from $-6\text{ }^{\circ}\text{C}$ ($\sim 520\text{ mb}$, $\sim 5600\text{ m}$). The nucleation temperature, which is the
195 temperature at which drops begin to freeze, is assumed to be $-8\text{ }^{\circ}\text{C}$. The figure demonstrates that
196 a drop with a radius of $100\text{ }\mu\text{m}$ cools from $-6\text{ }^{\circ}\text{C}$ to $-8\text{ }^{\circ}\text{C}$ and begins to freeze at approximately
197 23 s. The latent heat released due to freezing leads to a sudden drop in temperature from $-8\text{ }^{\circ}\text{C}$ to
198 $0\text{ }^{\circ}\text{C}$ (Fig. 3a), and the ice mass fraction increases from 0 to 0.1 (Fig. 3b). It takes approximately



199 4 seconds for the drop to fully freeze; during freezing, the drop temperature remains at 0 °C (Fig.
200 3a), and the ice mass fraction continuously increases (Fig. 3b). After completely freezing, the
201 frozen drop rapidly cools due to the large difference between the ambient temperature and the
202 drop surface temperature. The cooling rate slows down when the frozen drop temperature
203 approaches the ambient temperature. According to its equations, the cooling rate for a drop in the
204 updraft is largely controlled by convective heat transfer, rather than radiation or diffusional
205 growth. If significant riming occurs on the freezing (frozen) drop surface, the cooling rate could
206 be slower, and the freezing time could thus be longer due to the latent heat release that occurs
207 during riming (Phillips et al., 2015). The drop temperature changes in a similar way for larger
208 drops as it does for small drops. However, due to their higher terminal velocity, it takes longer
209 for larger drops to reach their nucleation temperature (-8 °C). Drops with radii of 250 μm and
210 500 μm begin to freeze at 28 s and 43 s, respectively (Fig. 3a), and their ambient temperatures
211 are approximately -8.1 °C and -8.15 °C (Fig. 3c), respectively. In addition, it takes longer for
212 larger drops to completely freeze. Drops with radii of 250 μm and 500 μm require approximately
213 15 s and 35 s, respectively, to fully freeze (Fig. 3a); these frozen drops are found at temperatures
214 of -9.2 °C and -9.95 °C, respectively (Fig. 3c).

215 Fig. 4 shows the freezing time and frozen temperature as functions of the drop radius for
216 different vertical air velocities and nucleation temperatures. The freezing time represents the
217 time period from the start of drop freezing to the end of drop freezing. The figure shows that the
218 freezing time increases as the radius increases. For the same nucleation temperature, drops freeze
219 faster in stronger updrafts than they do in weaker ones (Fig. 4a); however, their frozen
220 temperatures are colder in stronger updrafts (Fig. 4b). In addition, for the same updraft strength,
221 a drop freezes faster when its nucleation temperature is lower, and it fully freezes at colder



222 temperatures. Moreover, for the same drop radius, the effect of the updraft strength on the
223 freezing time is smaller if a drop nucleates at a lower temperature, as is indicated by the smaller
224 differences between the solid, dashed and dotted lines for colder nucleation temperatures (Fig.
225 4a); however, its impact on frozen temperature does not vary substantially with different
226 nucleation temperatures (Fig. 4b).

227 Large drops may begin to freeze at warmer temperatures than small drops (Bigg, 1953). Fig. 5
228 shows the nucleation temperature and frozen temperature as functions of the drop radius. The
229 nucleation temperature is the temperature at which drops have a $10^{-4}\%$ probability of freezing, as
230 determined based on Bigg's parametrization for immersion freezing. This probability is low
231 because of the low concentration of immersion ice nuclei that are present at warm temperatures.
232 The figure shows that large drops may begin to freeze at warmer temperatures than small drops;
233 however, due to their longer freezing times, large drops may fully freeze at colder temperatures
234 than small drops if the updraft is strong enough. Immersion freezing is not the only ice
235 nucleation mechanism. In convective clouds, small ice can be generated at warmer temperatures
236 by other mechanisms (e.g., condensation/deposition nucleation). The ice PSD measured by the
237 Learjet indicates that large frozen drops were observed at colder temperatures than small ice, but
238 it is not known whether these large drops started to freeze before or after the small droplets, and
239 the mechanisms that lead to drop freezing are not well understood.

240 **3.3 Discussion**

241 The above analysis indicates that large frozen drops are observed at relatively colder
242 temperatures than small ice in strong updrafts of convective clouds but that they may begin to
243 freeze at warmer temperatures. If the vertical air velocity is not strong enough, large drops may



244 descend or remain at the same level for long periods of time, and they may freeze if their
245 temperature reaches the nucleation temperature. An example of this is shown in Fig. 6. In this
246 case, penetration occurred approximately 500 m below the cloud top, as is indicated by the WCR
247 reflectivity (Fig. 6a). The WCL power (Fig. 6c) quickly attenuated and the WCL depolarization
248 ratio (Fig. 6d) is relatively low, which indicates that this cloud was dominated by liquid drops.
249 At the flight level, the temperature (Fig. 6e) ranges from $-4\text{ }^{\circ}\text{C}$ to $-4.5\text{ }^{\circ}\text{C}$ in the updraft and is
250 approximately $-5\text{ }^{\circ}\text{C}$ near the cloud edge. The maximum updraft velocity is 7 m/s, and the mean
251 updraft velocity is approximately 3 m/s. The Doppler velocity (Fig. 6b) is negative in most areas
252 of the clouds, and its maximum value is approximately 4 m/s. The 2D-C images clearly show the
253 existence of ice (Fig. 6f). Most of the ice particles are frozen drops and graupel, and some are
254 needles and columns. The graupel may fall from above; thus, they may start freezing at a colder
255 temperature than the flight level temperature. Considering the time that is needed for the drops to
256 freeze and for the needles and columns to grow through vapor diffusion, this ice may have
257 nucleated when the cloud top was lower than observed.

258 The freezing of supercooled drops may be associated with some corresponding processes. For
259 example, drops may break up or shatter during freezing, which can produce multiple ice
260 fragments and splinters. Mason and Maybank (1960) showed that the freezing of a millimeter-
261 sized drop may produce more than a hundred splinters. These ice splinters can enhance ice
262 initiation in convective clouds. In addition, the change in drop temperature during freezing may
263 exert impacts on the Hallett-Mossop process. Heymsfield and Mossop (1984) showed that the
264 Hallett-Mossop process is not only related to the ambient temperature but is also related to the
265 graupel surface temperature. In the SBM used in this study, the Hallett-Mossop process is only
266 parameterized for ambient temperatures between $-3\text{ }^{\circ}\text{C}$ and $-8\text{ }^{\circ}\text{C}$. However, the Hallett-Mossop



267 process may occur at colder ambient temperatures if the frozen drop (or graupel) surface
268 temperature is appropriate (Heymsfield and Mossop, 1984). Fig. 3 shows that the drop
269 temperature cools from 0 °C to its ambient temperature after being fully frozen and that the
270 cooling rate may be even slower if there is significant riming on the surface of the particle
271 (Phillips et al. 2015). During this process, if the drop surface temperature and other ambient
272 conditions are suitable, the Hallett-Mossop process may occur at an air temperature that is colder
273 than -8 °C, which could also enhance the initiation of ice in developing convective clouds. For
274 example, a millimeter-sized frozen drop can collect approximately 600 droplets in five seconds,
275 assuming that the droplet concentration is 50 cm⁻³ and its diameter is 20 μm. Thus, two or three
276 ice splinters may be produced if the ambient conditions are suitable. Moreover, time-dependent
277 freezing can have an impact on the dynamics in developing clouds. The instantaneous freezing of
278 a supercooled drop results in the sudden release of a large amount of latent heat, which may lead
279 to an overestimation of the vertical velocity in modeled convective clouds. In contrast, time-
280 dependent drop freezing can affect the cloud dynamics in a different way because its latent heat
281 is gradually released. Future studies are needed to explore these drop freezing-related processes.

282 This study reveals the importance of understanding drop freezing in convective clouds and
283 allows us to interpret the observed ice PSDs; however, it also raises some specific questions
284 about ice initiation. For example, it is not known why the observed "primary ice" concentration
285 is much higher than the ice nuclei concentration (DeMott et al., 2016) and the modeled ice
286 concentration (Fig. 1). There are several possibilities for this, including the production of ice
287 fragments and splinters during drop freezing or the Hallett-Mossop process; droplet collisional
288 freezing (Alkezweeny, 1969); or the electrofreezing of drops (Pruppacher, 1973). In addition, it
289 is not known whether large drops begin freezing before or after small droplets. Answering these



290 questions requires a better understanding of the primary drop freezing mechanisms in convective
291 clouds, which in turn requires more laboratory experiments to be performed and more in situ
292 measurements to be obtained in the future.

293 **4. Summary**

294 In this study, the PSDs measured in fresh developing maritime convective clouds sampled during
295 ICE-T are shown and the deficiency of instantaneous drop freezing in models is discussed. The
296 observations presented here suggest that the "first ice" that is observed is small. To interpret the
297 observed ice PSDs, the freezing times and temperatures of supercooled drops are calculated. This
298 analysis indicates that the freezing time is longer for large drops than it is for small drops. Due to
299 the limitations of airborne instruments, freezing drops cannot be identified until they exhibit
300 obvious shape deformation. If the updraft is strong enough, large freezing drops may be brought
301 up to a colder temperature than their nucleation temperature before they begin to exhibit obvious
302 shape deformation. This study allows us to interpret the observed ice PSDs in fresh developing
303 convective clouds from the perspective of drop freezing. However, the mechanisms of drop
304 freezing and ice initiation are still not well known. Future studies are required to evaluate model
305 simulations using time-dependent drop freezing, to understand the impact of time-dependent
306 drop freezing on the microphysics and dynamics of convective clouds, and to further explore the
307 mechanisms of drop freezing and ice initiation.

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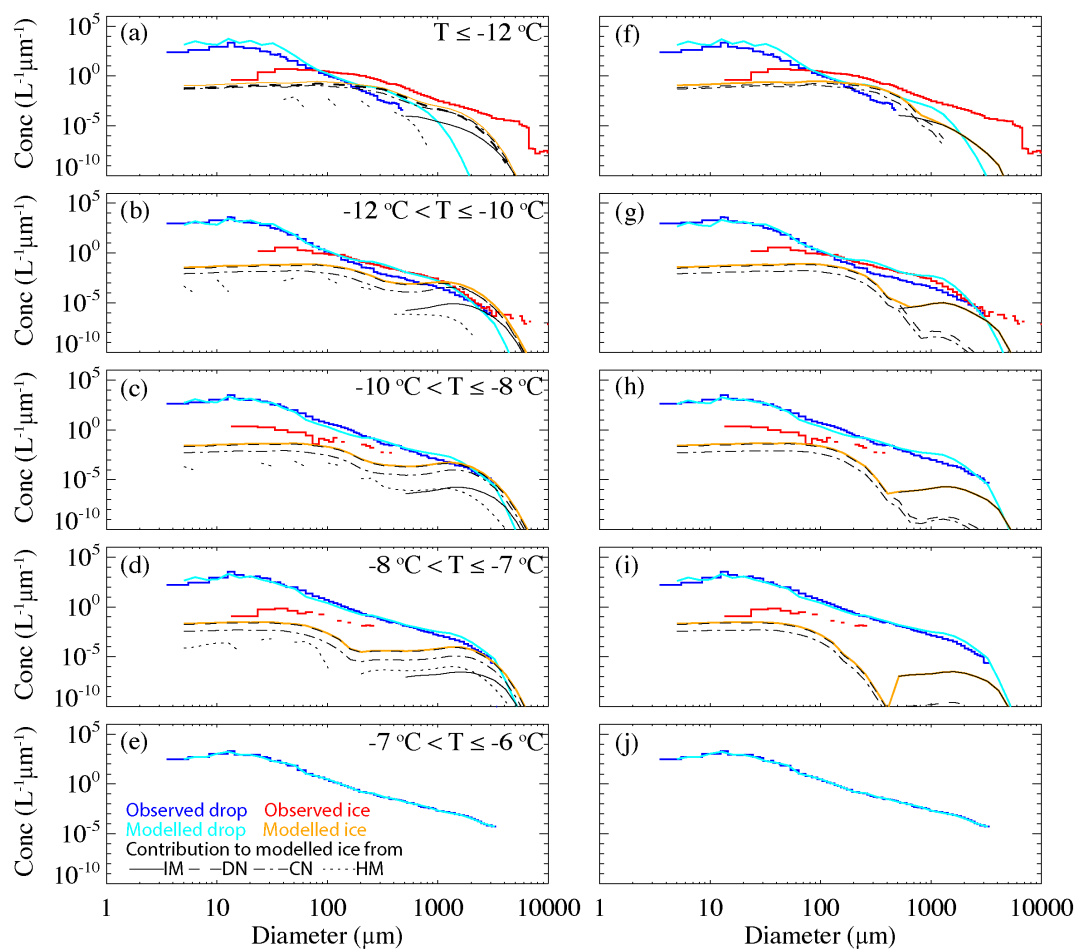


Figure 1. Particle size distributions in fresh developing convective clouds observed by the Learjet during ICE-T and those modeled using a parcel model with SBM. In the left panels, all of the ice physics implemented in the SBM are included; in the right panels, liquid-ice collision is excluded. The black solid, dashed, dashed-dotted, and dotted lines represent the contributions from immersion freezing (IM), deposition/condensation nucleation (DN), contact nucleation (CN), and the Hallett-Mossop process (HM), respectively, to the modeled ice size distributions.

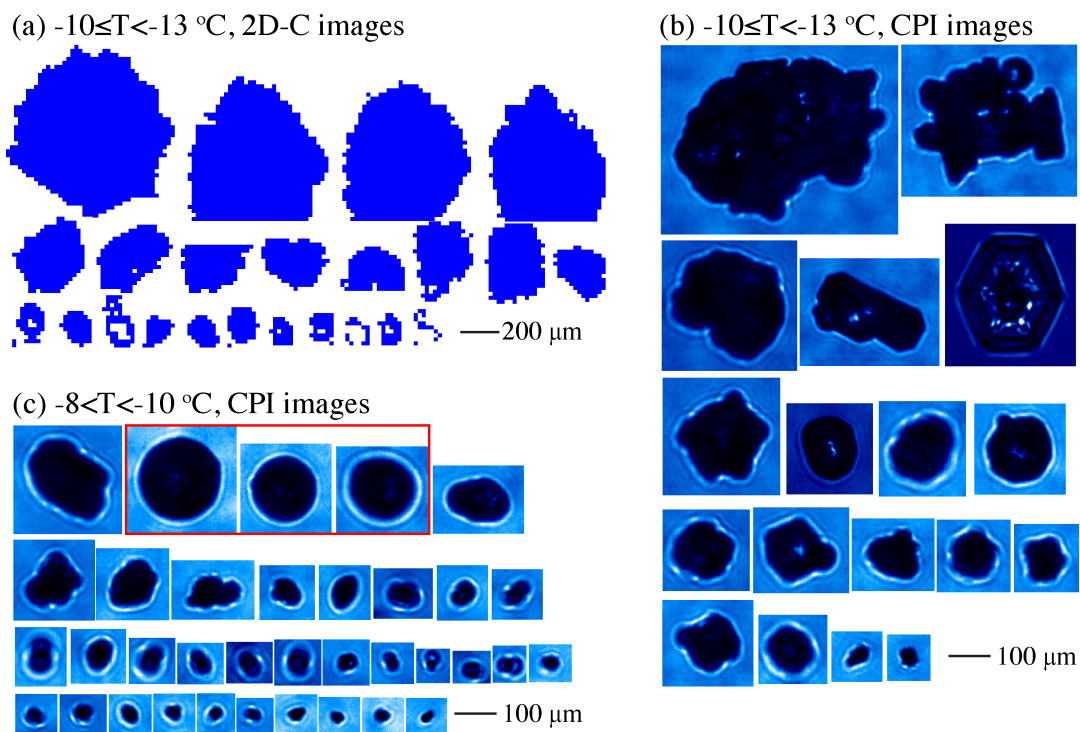


Figure 2. Examples of the 2D-C and CPI images measured in the developing convective clouds sampled during the ICE-T project.

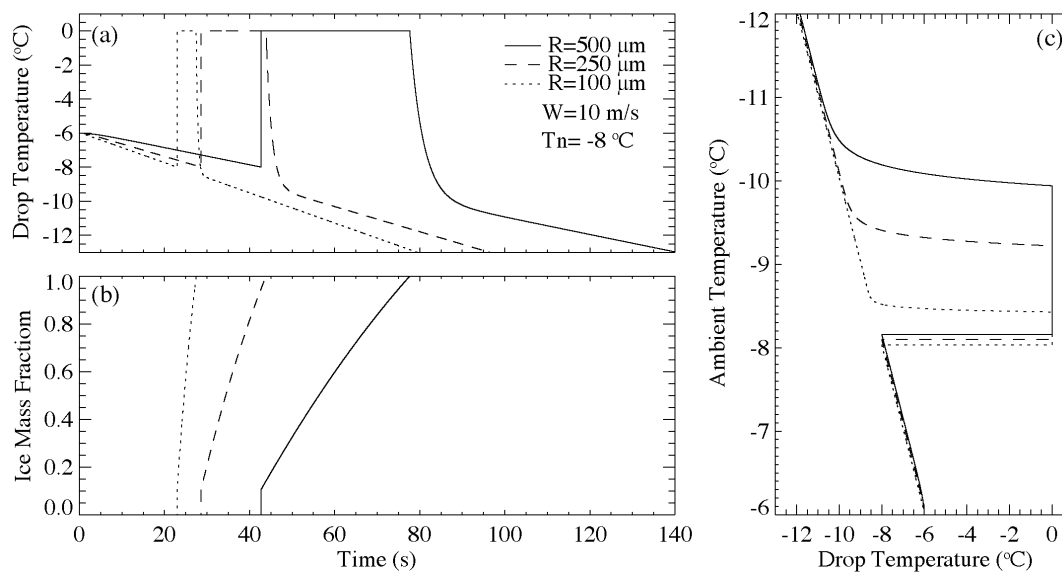


Figure 3. (a) Changes in drop temperature over time for drops with different radii. Vertical air velocity (W) is assumed to be 10 m/s and nucleation temperature (T_n) is -8 C; (b) same as (a) but for ice mass fraction; (c) ambient temperature versus drop temperature for drops with different radii.

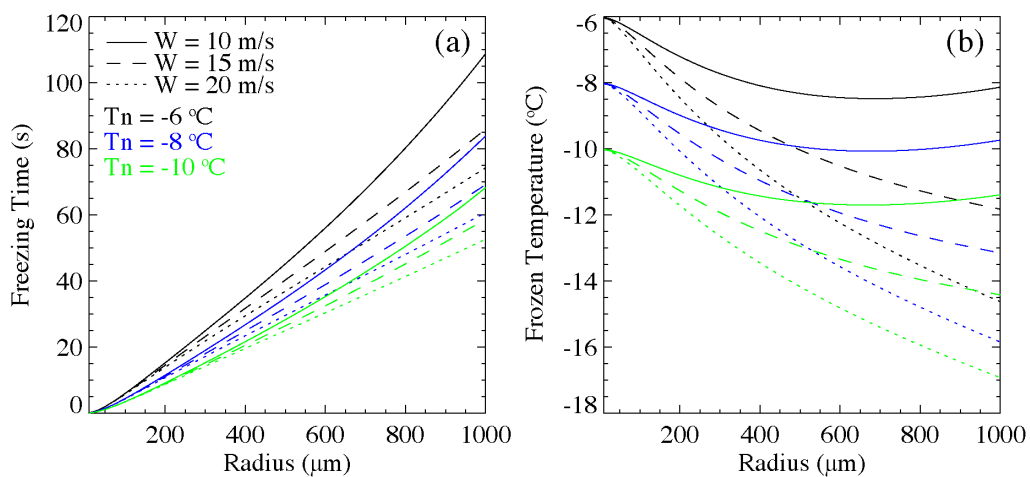


Figure 4. (a) Freezing time and (b) frozen temperature as functions of drop radius for different values of vertical air velocity (W) and nucleation temperature (T_n).

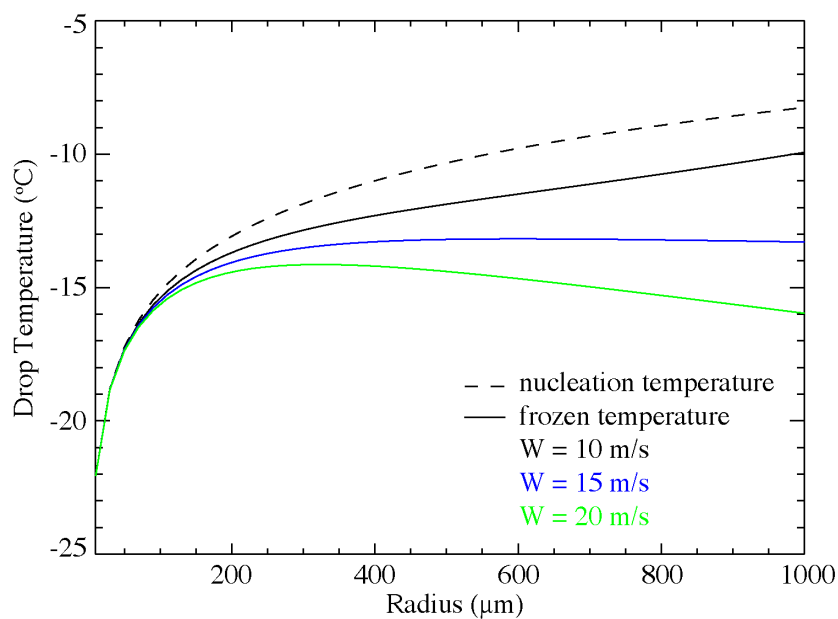


Figure 5. Drop temperature as a function of drop radius for different vertical air velocity (W) values. The nucleation temperature is the temperature at which drops have a $10^{-4}\%$ probability of freezing, as determined based on Bigg's parameterization for immersion freezing.

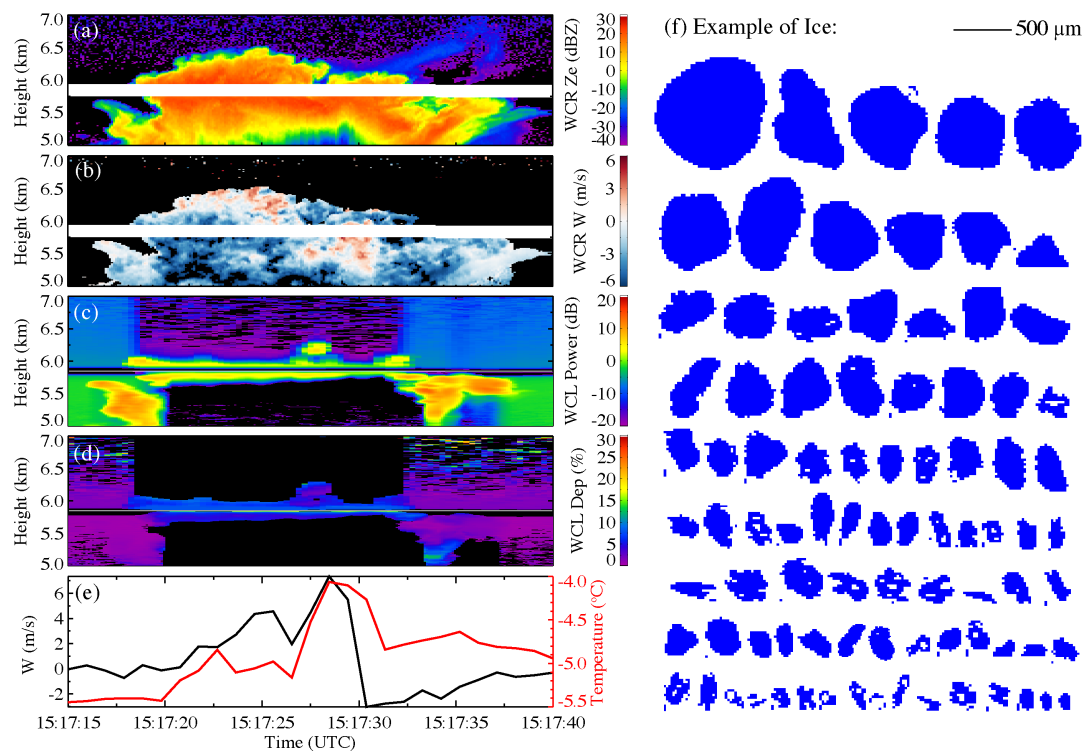


Figure 6. An example of the penetration of the C-130 in a developing cloud sampled on 23 July 2011: (a) WCR reflectivity; (b) WCR Doppler velocity; (c) WCL power; (d) WCL depolarization ratio; (e) ambient temperature and in situ vertical air velocity; and (f) examples of ice particles measured using 2D-C.