<center>**Reply to reviewers**</center>

Reply to Reviewer 1

We thank the reviewer for the careful analysis on our manuscript and for the overall positive opinion on our study. He/she raises very important issues, which have been addressed, largely improving the manuscript, in our opinion. Below, we answer to each of the points raised.

Reviewer says:

On page 2, l. 21–23, the authors say:

> Equivalent analyses focusing on scalar flux cospectra have not been presented as often. Sakai et al. (2001) and Finnigan et al. (2003) used cospectral similarity to conclude that low-frequency contribution could account for missing energy and CO_2 fluxes in their respective budgets, but neither study addressed how the cospectra varied across the canopy.

later, on p. 2, l. 30–33, they say:

> This result indicates that the exchange of scalars between the canopy and the atmosphere at night may occur at longer time scales than those traditionally used in the eddy covariance approach.

and again, on p. 10, l. 10–15:

> Our results support these findings, adding the information that the non-turbulent contribution may dominate the exchange of CO_2 and humidity from the interior of the canopy in very stable nights as well. It is likely that the same process affects other scalars, such as O_3 , whose concentrations are perturbed by intermittent events as shown in Fig. 4b.

However, in the conclusions, they find that low-frequency components are important within the canopy, but that, above the canopy, it is the "turbulent scales" that contribute most of the flux. There seems to be a contradiction between the Introduction (and other parts of the manuscript) and the Conclusions. The introduction should not lead the reader to believe in a situation that will not be supported by the analysis.

We thank the reviewer for raising this issue and for addressing it in great detail. There is certainly a lot of confusion regarding this matter in the original version of the manuscript. Trying to be as simple as possible, what we meant to say is that:

- Above the canopy, turbulent exchange is always important;
- Within the canopy, nonturbulent exchange is always important;
- In very stable conditions, the nonturbulent contribution increases above the canopy as well. In such conditions, both turbulent and non-turbulent exchange become nearly as important.

Now, to correct the confusion, the following portions have been altered in the text:

In the original version (p. 2, l. 29):

> They also found that sensible heat flux cospectra within the canopy peaked at longer time scales, again similar to those of the non-turbulent maxima of horizontal TKE above the canopy. This result indicates that the exchange of scalars between the canopy and the atmosphere at

night may occur at longer time scales than those traditionally used in the eddy covariance approach.

Has been replaced with (changes in red):

They also found that sensible heat flux cospectra within the canopy peaked at longer time scales, again similar to those of the non-turbulent maxima of horizontal TKE above the canopy. Their results indicate that the exchange of scalars within the canopy at night may occur at longer time scales than those traditionally used in the eddy covariance approach. At very stable conditions, such long scales may also contribute to the total exchange between the canopy and the atmosphere.

In the original version (p. 10, l. 9):

Santos et al. (2016) found that the time scales of horizontal turbulent velocity fluctuations within an Amazonian rain forest canopy (at a different site) approach those of the non-turbulent maximum above the forest. They also found that the dominant time scales of the vertical velocity fluctuations and sensible heat flux within the canopy are shifted towards larger values than those above it. Our results support these findings, adding the information that the non-turbulent contribution may dominate the exchange of CO_2 and humidity from the interior of the canopy in very stable nights as well.

Has been replaced with (changes in red):

Santos et al. (2016) found that the dominant time scales of the vertical velocity fluctuations and sensible heat flux within an Amazonian rain forest canopy (at a different site) are shifted towards larger values than those above it. Our results support these findings, adding evidence that the exchanges of CO_2 and humidity within the canopy are also dominated by non-turbulent contribution at very stable nights.

In the original version (p. 17, l. 4, at the conclusion section):

The majority of the fluxes just above the canopy (41-m level, in this case) happens through turbulent exchange. Although no relevant systematic low-frequency contribution to the total fluxes has been found at 41 m for any scalar analyzed (Fig. 7), this result only holds for the average spectra, with appreciable variability among cases, especially in the most stable cases.

Has been replaced with (changes in red, last sentence removed):

Turbulent exchange is always important just above the canopy (41-m level, in this case) but, in very stable nights the non-turbulent contribution has to be also considered.

Reviewer says:

"Bulk" Richardson numbers are used, but these are sensitive to velocity and, most of all, temperature systematic errors between the sensors. Because several analyses are dependent on these Richardson numbers, their reliability must be assessed quantitatively. Have the sensors been intercompared?

In the worst case (no intercomparison, no calibration), a thorough sensitivity analysis must be made of the effects of the temperature (and wind) systematic errors on those Richardson numbers and in the analyses involving them. The reported accuracies for the sensors (assuming optimistically that they did not drift) can be used as a basis for this. The uncertainty introduced by those errors results should then be displayed graphically in all analyses regarding the Richardson number.

Temperature has been measured by a profile of thermohygrometers and by the sonic anemometers, as well. The thermohygrometers have been intercompared, while each sonic temperature has been compared to the closest thermohygrometer. These informations have been added to the manuscript.

Reviewer says:

Turbulent bursts: the criterion for identifying the turbulent bursts and defining the shaded regions in Fig. 2 should be made clear (quantitavely).

It has been assumed that a turbulent burst occurred whenever $\sigma_w > 0.15\,m\,s^{-1}$. This information has been added to the manuscript.

Reviewer says:

Text starting on p. 11, l. 5, says

> Sun et al. (2012) found two regimes of nocturnal turbulence, distinguished by the turbulent kinetic energy (TKE) dependence on the mean wind speed. The fully turbulent regime, typically associated with weakly stable conditions, happens for mean wind speeds larger than a height dependent threshold and is characterized by TKE that steadily increases with wind speed. The other regime, associated with very stable conditions, has reduced turbulence intensities, which are very weakly dependent on the mean wind speed. Dias-Júnior et al. (2017) observed the two regimes above the forest at a site in the southwestern Amazon, finding that each is associated with an independent lognormal frequency distribution of quantities such as the turbulence dissipation rate. For the turbulent night of 15 November 2015 (Fig. 8, crosses), the levels of 41 and 55 m remained in the large wind speed regime for the whole period, while the two different regimes could be observed only at the 80-m level. On the intermittent night, on the other hand (Fig. 8, triangles), both regimes could be observed at all levels. Moreover, the connection intervals, given by shaded areas in Fig. 2, are generally in the large wind speed regime both at 41 m and 55 m (Fig. 8, filled triangles), while the decoupled periods are in most cases in the weak wind regime (Fig. 8, open triangles). **This is an important result, because it indicates that the intermittent bursts of turbulence observed above the canopy are intense enough to cause a regime transition. It means that, during these events, there is likely full vertical coupling over the vertical extent of the SBL (which is, at this time, shallower than 80 m). Therefore, scalars that are emitted from the canopy may be able to escape to higher levels in the boundary layer, as suggested by the episodic mixing of CO_2 and O_3 above 70 m shown in Fig. 4.**

(my emphasis). But high turbulent fluxes above the canopy during the the bursts of turbulent activity are already clearly displayed in Fig 5. The "full coupling" is none other than the relatively high (absolute) values of the fluxes themselves. Given that the fluxes are there, the scalars have already "escaped" the canopy. Therefore, the reasoning in the bold-face text above seems to be rather circular, and nothing new seems to arise from this discussion. Moreover, if the criterion for identifying the bursts was TKE (as I suspect), then it is inevitable that this will be reflected in higher TKE values in Fig. 8. It appears to me that the definition of the bursts and the regime classification in Fig. 8 are one and the same, and that there is nothing to be added here. I strongly suggest deleting this whole passage.

It is possible that an intermittent event is not intense enough to cause regime transition. In this case, they are local, and they lie at the weak wind side of the V vs V_{TKE} diagram (despite the TKE enhancement). However, we understand that this type of intermittence has never been considered in this study, and that its abrupt introduction to the manuscript may be rather confusing to the reader. Therefore, we agree with the reviewer´s suggestion of removing this passage, which we will keep to a specific study on regime transitions.

Reviewer says:

Sections 4 and 5 seem to use all the data from the 15 usable nights. Because the previous section focused strongly on the comparison of the nights of Nov 14 and 15, I had a hard time (at first reading) realizing this. I suggest that both the title and the introduction of each of these sections reinforces the information that, now, data from all 15 nights are being analyzed.

Title of section 4 has been changed to "Mean spectra and cospectra over the 15 nights", while that of section 5 has been changed to "Dependence on stability over the 15 nights". Furthermore first sentence of section 4 now reads "Fig. 7 shows the spectra and cospectra of the turbulent fluctuations and fluxes averaged over the entire period of 15 nights". First sentence of section 5 now reads "The comparison of the fluxes determined with 5-min and 109-min time windows and of their stability dependence for the entire period of 15 nights provides interesting information on the scalar exchange processes within and above the canopy"

Reviewer says:

(p. 12, l. 10): "Figure 9 shows the spectra and cospectra of the turbulent fluctuations and fluxes averaged over the entire period".
Particularly in stable conditions, there is a strong shift of the spectra towards the higher frequencies with increasing stability (Kaimal, 1973). There is no equation describing how the spectra were "averaged", but there should be. The simplest approach (which I suspect is being used here) is to average per frequency. But then, because frequency depends on stability, different stabilities and their spectral densities are being averaged together. The consequences are far from clear to me, and this procedure should not be done without careful justification. Remember, if y = f (x) and f is nonlinear, then $\overline{y} \neq f(\overline{x})$ in general. It is not clear how the fluxes reported in Sect. 5 were calculated. Are they bin averages? Do they come from the integration of the mean spectra? If $F^{(i)}_{wa}$ is the flux from the ith cospectrum, and if $F_{wa,mean}$ is the flux from the mean cospectrum (as depicted in Fig. 9), how do $(1/n)\sum_{i=1}^{n} F^{(i)}_{wa}$ and $F_{wa,mean}$ compare? In this sense, how valuable and correct are the conclusions derived from Fig. 9?

This is an important issue, we thank the reviewer for raising it. First, he/she is correct that we are averaging by frequency. In that sense, Fig. 9 is being affected by this problem, which may cause the spectra to "spread" in the horizontal. We do not think this is a problem, as the paper does not focus on scaling issues but, rather, on reporting physical processes, the contrast between turbulent and nonturbulent exchange. These are clear in Fig. 7, despite some spreading over the time scale axis. Nevertheless, to make it clear, we added an explanation at the caption of Figs. 6, 7, 9, 13 and 14 that they show averages by frequency. ("in all panels, averages are performed over each time scale").
Besides, in figs. 13 and 14, the two classes represent different stabilities, and it can be seen that the scales of the turbulent fluxes vary much less between classes of stability than does the magnitude of the nonturbulent contribution to the fluxes.

Finally, all flux estimates from cospectra have been done for each series separately, so that their average shown are $(1/n)\sum_{i=1}^{n} F_{wa}^{(i)}$ rather than $F_{wa,mean}$ (using the reviewer notation). This has been now properly stated in the manuscript, to avoid confusion from the readers: "Variances and fluxes with a 109-min long time average have been obtained from the integration of the respective multiresolution spectra and cospectra for each series separately, and then averaged, if appropriate."

Reviewer says:

*p. 4, l. 15–16: "Since the different levels of flow structures are analyzed simultaneously, only the data when all levels were available was used." This should be: ". . . Since the different levels of flow structures are analyzed simultaneously, only the data when all levels were available **were** used".*

Done.

Reviewer says:

p. 4, l. 19–20: "All the time series have been subject to quality control, which caused the removal of those series, which showed multiple spikes or spectra that did not converge to zero at the highest frequencies."

The meaning of this sentence is unclear! What does it mean for a spectrum to "converge to zero" at the highest frequencies? Turbulence spectra decay as $k^{-5/3}$ in the inertial subrange . . . Do you mean spectra displaying noise in the higher frequencies? Not falling off as $k^{-5/3}$, levelling off?

Please explain.

We simply meant "multiresolution spectra that displayed noise at the shortest time scales".

Reviewer says:

p. 4, l. 33 – p. 5, l. 4 There appears to be a conflict of notation between C for the cospectrum and C for the concentration of CO_2.

All occurrences of "C" alone referred to cospectrum.

Reviewer says:

Eq. (1) and (2) How did you calculate θ_{22}, θ_{41} and θ_{80}? From what instrument? Temperature profiles are sensitive to bias in the sensors: were the temperature sensors at these heights intercompared before deployment?

They have been determined from the sonic temperatures, which have been compared to the closest thermohygrometer, as explained in the reply to the second comment from the reviewer, above.

Reviewer says:

p. 5, l. 1–2: "and the standard deviation of the vertical wind component is $\sigma_w = \sum_{\tau} S_w$".

Wrong: the relationship is $\sigma_w^2 = \sum_{\tau} S_w$

Authors: check your calculations carefully to see if this is just a typo, or if you actually calculated (and are reporting) wrong values.

It was just a typo, thank you for noticing it!

Reviewer says:

p. 5, l. 2–3: "Other variables, such as the Richardson number (Ri) and average horizontal wind speed (V) were calculated using the same data series used in the multiresolution decomposition."

Too vague: were the mean velocities from the sonics? Very important (see Main remarks above): from which sensors do the mean temperatures come?

Yes, both came from the sonic data. We used this phrasing to state that the same time intervals were chosen as those used for the multiresolution decomposition. Sentence now reads: "Other variables, such as the Richardson number (Ri) and average horizontal wind speed (V) were calculated using the same time intervals used in the multiresolution decomposition."

Reviewer says:

p. 5, l. 22 Again, how were the θ's measured?

From the sonic data, which have been compared to the thermohygrometers. This information has been added to the manuscript.

Reviewer says:

Section 3: Rename the section to indicate that it is about the comparison of two nights, one fully turbulent and the other intermittently turbulent. Suggestion: Comparison of turbulence characteristics in a fully turbulent with and intermittently turbulent night.

Done.

Reviewer says:

p. 5, l. 25–26 "The nocturnal flow at the site is characterized by the superposition of turbulent and non-turbulent fluctuations. In a fully turbulent night, such as 15 November 2015 (Fig. 1), there is a clear dominant wind direction at all levels."

Figure 1 does not show wind directions at the different levels. It is impossible to infer wind direction at each level from the figure.

Figure direction can be inferred from the signals of the two horizontal wind components. To clarify this matter, wind direction has been associated to the sign change of the components at another sentence (this association has already been made once, at the previous sentence): "In contrast, during the intermittent night of 14 November 2015 (Fig. 2), there is no dominant wind direction at any level above the canopy, as both horizontal components switch sign many times along the night".

Reviewer says:

p. 6, l. 3–8 and Table 1 "The most relevant difference between the two nights regards the magnitude of the turbulent mixing (Table 1). All relevant turbulence statistics are significantly larger on 15 November than on 14 November. The relative difference of the turbulence statistics between nights increases steadily in the vertical. As an example, TKE at 41 m is 3.4 times larger in the turbulent night than in the intermittent case, while at 80 m, TKE is 8.2 times

larger in the turbulent night. Similar increases occur for the corresponding ratios of σ_w and u_* between the two nights."

The authors should reserve the symbol u_* for a single value in each period, which should be the most representative for the friction between the flow above the canopy and the forest. Obviously this would be the value reported at 41 m. The others are "local" values of the kinematic momentum flux, and it would be more appropriate to write them as $\sqrt{\overline{w'u'}}$. Same comment applies for θ_*, etc..

The suggestion has been accepted.

Reviewer says:

Fig. 1-d, Fig. 2-d The title CO_2 is missing from the left vertical axis.

Done.

Reviewer says:

p. 7, l. 10

"All quantities showed much larger variation across the levels in the intermittent night (Fig. 2). Furthermore, sporadic events of coupling occurred during bursts of intermittent turbulence (Fig. 2, shaded areas)."

The authors never explain the exact quantitive criterion for the identification of the shaded areas. It appears to be TKE, but they should give the quantitative criterion in the text.

This issue has already been addressed in the third comment from the reviewer, above.

Reply to Reviewer 2

Despite the rejection suggestion, with which we do not agree, we thank the reviewer for raising the issue of coordinate rotation. This is, as the reviewer points out, a key issue in a study of this kind and it had not been properly clarified in the original version. Below is our answer.

Reviewer says:

The authors analyzed nighttime vertical fluxes of heat, water vapour, carbon dioxide, and ozone within and above a rainforest. They used multiresolution decomposition to determine the scales of atmospheric motions contributing to the vertical fluxes, focusing on low-frequency, non-turbulent fluctuations. I suggest rejection of the current manuscript because the authors failed to address a fundamental issue that no sonic anemometer can be aligned perfectly with the vertical direction perpendicular to the underlying surface. One may use a plumb bob to align sonic anemometer with the vertical direction, but an error of one or two degrees is expected over flat topography, and the error can be larger over slopes. One can use coordinate rotation techniques in the post-processing, but an uncertainty of two degrees still exists (e.g., Forken et al. 2004; Vickers and Mahrt 2006). Aubinet et al. (2003) highlighted that "The 2° offset would induce systematic errors on the vertical velocity up to 0.05 m s^{-1} under typical stable conditions and up to 0.11 m s^{-1} under near-neutral conditions. The resulting error in the vertical advection flux in the presence of a 10 µmol mol^{-1} vertical CO_2 concentration difference may be as high as 5 µmol m^{-2} s^{-1} ". The errors in sonic coordinate system estimates would convert a few percent of variation in horizontal velocity components to variation in the vertical velocity component. On low-frequency, non-turbulent scales, horizontal velocity components are typically two orders of magnitude larger than the vertical velocity component. Consequently, the artificial variation induced by errors in sonic coordinate system estimates is at least comparable to the true variation in vertical fluxes on low-frequency, non-turbulent scales. Using eddy-covariance measurements to draw conclusions about low-frequency, non-turbulent vertical fluxes does not make sense unless the authors can distinguish true variation in vertical fluxes and artificial variation inherited from horizontal velocity components due to errors in the sonic coordinate system estimates. This fundamental issue must be resolved before the manuscript can go to more detailed review.

References

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The analysis of non-turbulent fluxes with time scales from several minutes to a few hours is very sensitive to sonic alignment, as the reviewer points out. However, we do not agree that such sensitivity is so large that it prevents the topic from being studied from field observations. A number of studies have addressed this very important topic of boundary layer research in the past (Monti et al., 2002; Vickers and Mahrt, 2007; Mahrt et al., 2008; Mahrt, 2009; Acevedo and Mahrt, 2010; Hoover et al., 2015; Santos et al., 2016; Cava et al., 2017;

Mortarini et al., 2017) and all of them are, to some degree, subject to the problem raised by the reviewer. Many important results have arisen from these studies, indicating that the problem may, indeed, be studied from field observations. Nevertheless, this is certainly an important matter, to which not enough attention had been given in the original version of the manuscript.

In a previous paper (Acevedo and Mahrt, 2010), vertical profiles of submeso fluxes of sensible heat have been analyzed in detail. It was found that systematic vertical profiles of nonturbulent fluxes of sensible heat exist, and that their inclusion contributes to closing the nocturnal temperature budget near the surface. Not surprisingly, it was also found that those profiles are largely dependent on the type of wind rotation used. Many different wind rotation techniques were tested, and it was found that the "double wind rotation" offered the best results. So large is the importance of coordinate rotation for this matter that an entire subsection of that study is devoted to describing this issue in that paper (section 2.1). Some important results described in that subsection are:

- Double rotation applied to individual time segments leads to heat fluxes that are more coherent with height than using a globally directionally dependent method. This is because "*the measured vertical motion on times scales greater than 5 h may be sufficiently weak and unreliable that the elimination of larger-scale variations of vertical motion through coordinate rotation improves the calculation*" (Acevedo and Mahrt, 2010);

- "*The response of the mesoscale momentum flux to coordinate rotations is more complex than that for the heat flux partly because the rotations affect both the vertical and horizontal velocity components and their correlations. We limit this study to the mesoscale heat flux.*" (Acevedo and Mahrt, 2010)

In the present study, double rotation has been applied to each individual time series, following the suggestions from Acevedo and Mahrt (2010). Besides, only scalar, rather than momentum fluxes are considered. We believe that those results provide strong foundation for us to use this technique in the present study. However, they certainly need to be properly mentioned in the manuscript, and they were not in the original version. Therefore, the following paragraph has been added to section 2:

> Acevedo and Mahrt (2010) used the multiresolution decomposition to analyze vertical profiles of the non-turbulent component of sensible heat fluxes. They found that systematic and organized profiles, whose inclusion contributes for the closure of the nocturnal temperature budget near the surface, are only found when the double wind rotation (Tanner and Thurtell, 1969) is applied to each time series analyzed, separately. They claim that this is more suitable for such analysis than other coordinate rotation procedures, such as globally directionally dependent methods (Lee, 1998; Mahrt et al., 2000; Paw U et al., 2000) because "...*the measured vertical motion on times scales greater than 5 h may be sufficiently weak and unreliable that the elimination of larger-scale variations of vertical motion through coordinate rotation improves the calculation*". For these reasons, the double rotation was applied to each 109-min time series separately.

Finally, we would like to stress that although a very important part of the paper, the contrast between turbulent and nonturbulent modes of the flow is not all that is addressed in the study. Furthermore, not only fluxes ("exchange") are analyzed, as scalar profiles and power spectra are also investigated. For this reason, and to take the focus of the study away from this

matter, we are also changing the paper title to "Nighttime wind and scalar variability within and above an Amazonian canopy".

References:

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~~Turbulent~~ Nighttime wind and ~~non-turbulent exchange of scalars between the forest~~ scalar variability within and ~~the atmosphere at night in Amazonia~~ above an Amazonian canopy

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Abstract. Nocturnal turbulent kinetic energy (TKE) and fluxes of energy, CO_2 and O_3 between the Amazon forest and the atmosphere are evaluated for a 20-day campaign at the Amazon Tall Tower Observatory (ATTO) site. The distinction of these quantities between fully turbulent (weakly stable) and intermittent (very stable) nights is discussed. Spectral analysis indicates that low-frequency, non-turbulent fluctuations are responsible for a large portion of the variability observed on intermittent nights. In these conditions, the low-frequency exchange may dominate over the turbulent transfer. In particular, we show that within the canopy most of the exchange of CO_2 and H_2O happens on temporal scales longer than 100 s. At 80 m, on the other hand, the turbulent fluxes are almost absent in such very stable conditions, suggesting a boundary layer shallower than 80 m. The relationship between TKE and mean winds shows that the stable boundary layer switches from the very stable to the weakly stable regime during intermittent bursts of turbulence. In general, fluxes estimated with long temporal windows that account for the low-frequency effects are more dependent on the stability over a deeper layer above the forest than they are on the stability between the top of the canopy and its interior, suggesting that low-frequency processes are controlled over a deeper layer above the forest.

1 Introduction

The turbulence structure above forested canopies has been an important subject of research over the past decades. Such knowledge is essential to answer relevant scientific questions such as the quantification of the exchange of scalars between forested ecosystems and the atmosphere. Some of the precursor studies in this field have been performed in the Amazon region during projects such as ABLE 2A and 2B (Fitzjarrald et al., 1988; Garstang et al., 1990; Fan et al., 1990). Subsequent projects in this region that kept the focus on this subject include ABRACOS (Grace et al., 1995; Malhi et al., 1998; Kruijt et al., 2000), LBA

(Araújo et al., 2002; Saleska et al., 2003; Miller et al., 2004), GO-Amazon (Fuentes et al., 2016; Santos et al., 2016) and, most recently, ATTO (Andreae et al., 2015; Zahn et al., 2016).

Ultimately, one of the most relevant questions that these projects aimed to answer is the role of the Amazon rainforest as either a net sink or source of CO_2 to the atmosphere. Results diverge largely among the studies, from a net sink of 5.9 T C ha^{-1} yr^{-1} found by Grace et al. (1995) to a net source of 1.3 T C ha^{-1} yr^{-1} found by Saleska et al. (2003). Although some of this variability may be accepted as genuine, caused by site or interannual differences, it is now well established that methodological problems had affected the estimates that found enhanced carbon uptake. Most of these issues regard the treatment of nocturnal data, as a consequence of the complex nature of the atmospheric flow during night under stable conditions. In particular, during very stable nights, turbulent mixing is reduced and constrained to small temporal scales (Vickers and Mahrt, 2006; Acevedo et al., 2014). The exchange of properties such as CO_2 from the forest to the atmosphere may occur mostly through non-turbulent motion, such as drainage flows (Staebler and Fitzjarrald, 2004; Aubinet et al., 2003; Feigenwinter et al., 2004; Tóta et al., 2008) or by transport on temporal scales longer than those that characterize the turbulent flow (Santos et al., 2016).

The motion with temporal fluctuations longer than turbulence but smaller than those produced by mesoscale systems has been referred to as "submeso" by Mahrt (2009) and it has become an important subject of micrometeorological research since then. Typically, these non-turbulent fluctuations may be larger in magnitude than their turbulent counterpart, and they may introduce fluxes that are larger as well. On the other hand, these fluxes are not driven by local gradients, so that they are also much more variable than the turbulent fluxes, and of either sign, in such a way that their overall contribution frequently averages out over longer periods (Vickers and Mahrt, 2003). Nevertheless, their contribution may be important for closing the budgets over smaller time periods (Acevedo and Mahrt, 2010; Kidston et al., 2010).

Many studies on turbulence above and within forested canopies have presented an analysis of the spectral distribution of the turbulence velocity components and of their vertical variation with respect to the canopy top (Baldocchi and Meyers, 1988; Blanken et al., 1998; Dupont and Patton, 2012). In general, these studies focused on the time scale of the turbulent exchange and how it depends on factors such as distance from canopy top and atmospheric stability. Equivalent analyses focusing on scalar flux cospectra have not been presented as often. Sakai et al. (2001) and Finnigan et al. (2003) used cospectral similarity to conclude that low-frequency contribution could account for missing energy and CO_2 fluxes in their respective budgets, but neither study addressed how the cospectra varied across the canopy. Other studies (Campos et al., 2009; Fares et al., 2014) looked at scalar flux cospectra with the specific purpose of identifying the proper temporal scale for turbulent flux determination. Santos et al. (2016) found that horizontal turbulent kinetic energy (TKE) spectra are bimodal above an Amazonian rain forest canopy, with the peak at short time scales being related to turbulence and the peak at longer time scales being associated with non-turbulent, submeso fluctuations. Within the canopy, on the other hand, only the peak at longer time scale is preserved, indicating that non-turbulent fluctuations above the canopy propagate downward more efficiently than the turbulent ones. They also found that sensible heat flux cospectra within the canopy peaked at longer time scales, again similar to those of the non-turbulent maxima of horizontal TKE above the canopy. This result indicates that the exchange of scalars between the canopy and the atmosphere within the canopy at night may occur at longer time scales than those traditionally

used in the eddy covariance approach. At very stable conditions, such long scales may also contribute to the total exchange between the canopy and the atmosphere. Their study, however, did not include the analysis of CO_2 or latent heat fluxes and reactive trace gases like O_3, so that the question whether these quantities are affected by similar processes remains open.

A comparison of scalar flux cospectra within and above a forested canopy, aimed specifically at addressing the contribution of non-turbulent flow to the total fluxes at the different heights, has not been presented previously. The present study aims at addressing this issue and to evaluate how these exchange processes affect the scalar profiles that are routinely measured at ATTO.

2 Data and Methods

2.1 Experimental site

The dataset was collected during the ATTO (Amazon Tall Tower Observatory)-IOP1 campaign in October/November 2015 at Reserva de Desenvolvimento Sustentável Uatumã (Uatumã Sustainable Development Reserve – USDR), in the Amazon region. The site is located on a plateau at 120 m above sea level, approximately 150 km northeast of Manaus and 12 km northeast of the Uatumã River. The average height of the highest trees near the tower is 37 m. Further information regarding terrain, soil and vegetation can be found at Andreae et al. (2015).

Micrometeorological observations were carried out on an 80-m walk-up tower with rectangular cross section at five different levels: 14, 22, 41, 55, and 80 m above the ground, the first two levels being within the forest canopy, while the three others are above it. Fast-response wind measurements were performed at all levels (CSAT3, Campbell Scientific Inc., at 14, 41 and 55 m, IRGASON, Campbell Scientific Inc., at 22 m, and Windmaster, Gill Instruments Limited, at 80 m). Temperature has been measured by a profile of thermohygrometers and by the sonic anemometers, as well. The thermohygrometers have been intercompared, while each sonic temperature has been compared to the closest thermohygrometer. Scalar concentrations of CO_2 and water vapor were measured at 22 m (IRGASON, Campbell Scientific Inc.), 41 and 80 m (LI-7500A, LI-COR Inc.). The diurnal cycle of the H_2O mixing ratios at 41 m was erroneous for unidentified reasons. The short-term (up to 20 min) variations were correct, but the longer trend did neither agree with the other open path instruments at 22 m, and 80 m, a nearby psychrometer (Frankenberger type, Theodor Friedrichs GmbH, Germany), or with the profile measurements (see below). Therefore, the water vapor mixing ratios at that level haven been corrected by separating the short term fluctuations from the trend by applying a running mean with a window size of 5 min and adding this high frequency contribution to the running mean (window size 5 min) of the nearby psychrometer. Scalar concentrations of ozone were measured at 41 m with chemiluminescence O_3 sondes (Enviscope, Germany). In front of the fast O_3 instrument there was a 5 m long 3/4 inch (7.52 mm inner diameter) Teflon tubing with a Teflon inlet filter (47 mm diameter, 5 μm pore size). The flows were varying due to filter clogging. After a filter change the flows were $21\,\mathrm{l\,min^{-1}}$ and $23.5\,\mathrm{l\,min^{-1}}$, respectively, whereas before the filter change they were $16\,\mathrm{l\,min^{-1}}$ and $14\,\mathrm{l\,min^{-1}}$, respectively. The resulting lag times were 0.6-0.95 sec and the Reynolds numbers in the tubing were 2400 to 4000 at 35 °C. On the days considered for the case studies (14 and 15 November), the flow was about $16\,\mathrm{l\,min^{-1}}$ and the residence time was therefore 0.8 sec. All the data were collected at 10-Hz rate. As the signal of the

fast O_3 sondes has a considerable drift, it was calibrated to a slow O_3 analyzer (TEI 49i, Thermo Scientific) as described by Zhu et al. (2015). The CO_2 profiles were measured sequentially by CO_2/H_2O analyzers (LI-7000, LI-COR Inc.) connected to heated inlets at 8 heights (0.05 m, 0.5 m, 4.0 m, 12.0 m, 24.0 m, 38.3 m, 53.0 m and 79.3 m). During the case study nights (November 14 and 15) only the LI-7000 after the Nafion®dryer was running and therefore the water vapor values could not be

5 used. The O_3 profiles were also measured using the same inlet system with an Ozone Analyzer (TEI 49i, Thermo Scientific). Ambient air was continuously drawn from the inlets through non-transparent PTFE-tubing (3/8") to a valve block, which switched between the different inlet levels, so that one intake height was purged by the sample pump (PTFE coated) while all the others were purged by the bypass pump. A time interval of 1 min was necessary for getting a constant and reliable signal for each concentration level: a complete cycle took 8 x 2 = 16 min, providing 2 measurements per level. Three 16-min

10 measurement cycles plus one shorter 12-min cycle were performed every hour. During that last cycle, a small compromise was made to fit 4 cycles into the hour, and valve switches occurred every 90 s, thereby allowing for only one concentration measurement at each level. The ambient air inlets mounted on the tower were protected from rain entering the inlet line by polyethylene funnels and from insects by polyethylene nets. A PTFE-filter (5 μm) was installed right after the inlet. The tubing was insulated with Styrofoam and heated. The internal temperature and pressure corrections of the LI-7000 were used, but to

15 further minimize pressure effects, the air drawn from the inlets for analysis was sampled at the exit of the Teflon pump, so that the measurements were made close to ambient pressure for all measurement levels. The entire setup was comparable to the profile system employed by Mayer et al. (2011).

2.2   Data Analysis

In the present study, 20 days of nocturnal data were analyzed, from 1 to 20 November 2015. To avoid sampling intense events

20 associated with the transitional characteristics between daytime and nocturnal boundary layers, the evening period between the sunset and 20:00 local time (LT) was not considered for this analysis. For this reason, nocturnal periods were restricted from 20:00 to 05:00 (LT). Since the different levels of flow structures are analyzed simultaneously, only the data when all levels were available ~~was~~ were used. The 14-m level frequently presented gaps and was not considered for this study. Therefore, the levels included are: one inside the canopy (22 m), one just above canopy top (41 m), and two levels well above the canopy (55

25 m and 80 m).

All the time series have been subject to quality control, which caused the removal of those series ~~,~~ which showed multiple spikes or ~~spectra that did not converge to zero at the highest frequencies~~ multiresolution spectra that displayed noise at the shortest time scales. For any case where a given series was discarded for a given variable, it has not been used for any of the variables. With these restrictions, 15 nights were kept for the final analysis. Ozone measurements started on 11 November

30 2015, so that only 9 nights of ozone flux data were available.

The data were analyzed using two different time windows: 5 and 109 min. The multiresolution decomposition (Howell and Mahrt, 1997; Vickers and Mahrt, 2003; Voronovich and Kiely, 2007) was applied to 109 min, which corresponds to groups of 2^{16} data points. In contrast to the Fourier transform, which determines periodicity, this technique mainly extracts the width of the dominant turbulent events by locally decomposing the variances. For this reason, the multiresolution spectrum (S)

and cospectrum (C) have the property that the integration up to a given time scale t is equal to the variance and covariance, respectively, for a t-long time series. Consequently, the multiresolution value for a given time scale captures the physical processes (and the flux) whose duration is smaller than that time scale.

The multiresolution decomposition was applied sequentially to the time series, starting at 20:00 LT, with an overlap of 30 min between the subsequent series, totaling 14 decompositions for each night. A total of 200 series was used in the study, considering all nights. Acevedo and Mahrt (2010) used the multiresolution decomposition to analyze vertical profiles of the non-turbulent component of sensible heat fluxes. They found that systematic and organized profiles, whose inclusion contributes for the closure of the nocturnal temperature budget near the surface, are only found when the double wind rotation (Tanner and Thurtell, 1969) is applied to each time series analyzed, separately. They claim that this is more suitable for such analysis than other coordinate rotation procedures, such as globally directionally dependent methods (Lee, 1998; Mahrt et al., 2000; Paw U "...the measured vertical motion on times scales greater than 5 h may be sufficiently weak and unreliable that the elimination of larger-scale variations of vertical motion through coordinate rotation improves the calculation". For these reasons, the double rotation was applied to each 109-min time series separately.

Variances and fluxes with a 109-min long time average were obtained from the integration of the respective multiresolution spectra and cospectra for each series separately, and then averaged, if appropriate. Therefore, sensible and latent heat, CO_2 and ozone fluxes are given by $F_H = \sum_\tau C_{\overline{w'\theta'}}$, $F_q = \sum_\tau C_{\overline{w'q'}}$, $F_C = \sum_\tau C_{\overline{w'C'}}$, and $F_{O_3} = \sum_\tau C_{\overline{w'O'}}$, turbulent kinetic energy is $TKE = 0.5 \sum_\tau (S_u + S_v + S_w)$ and the standard deviation of the vertical wind component is $\sigma_w = (\sum_\tau S_w)^{1/2}$. Other variables, such as the Richardson number (Ri) and average horizontal wind speed (V) were calculated using the same ~~data series~~ time interval used in the multiresolution decomposition.

At night, it is expected that the temporal scales of turbulent transport are smaller. Campos et al. (2009) showed at another site in the Amazon forest that the contribution of turbulence to the nocturnal fluxes above the canopy occurs at temporal scales smaller than 200 s. The use of a 109-min long time window is necessary to determine the contribution of turbulent and non-turbulent motions to the fluxes. However, in order to attempt to reduce any contribution from non-turbulent transport, statistical moments, fluxes and other variables were also calculated using a 5-min time window, as used by Dupont and Patton (2012). Quantities such as sensible heat $\left(F_H = \overline{w'\theta'}\right)$, latent heat $\left(F_q = \overline{w'q'}\right)$, CO_2 $\left(F_{CO_2} = \overline{w'C'}\right)$, and ozone $\left(F_{O_3} = \overline{w'O'}\right)$ fluxes, turbulent kinetic energy (TKE), the average horizontal wind speed (V), Richardson number (Ri), and σ_w were determined for both 5 and 109 min. The turbulent velocity scale, defined as $V_{TKE} = \sqrt{TKE} = [0.5\,(\sigma_u + \sigma_v + \sigma_w)]^{1/2}$, was calculated for 5-min time windows only. This study comprises of a total of 1,577 5 min windows.

The bulk Richardson number was used to quantify atmospheric stability. The choice of the bulk instead of the flux Richardson number for the analysis has two reasons: to avoid self-correlation (Hicks, 1978; Klipp and Mahrt, 2004; Baas et al., 2006) and to quantify better the stability in very stable conditions, when fluxes are expected to approach zero. Similarly as used by

Bosveld et al. (1999); Mammarella et al. (2007); Oliveira et al. (2013), a "within-canopy Richardson number" (Ri_{can}) and an "above-canopy Richardson number" (Ri_{top}) (Santos et al., 2016) were defined as

$$Ri_{can} = \frac{g}{\Theta} \Delta z \frac{\theta_{41m} - \theta_{22m}}{(V_{41m} - V_{22m})^2} \tag{1}$$

and

$$Ri_{top} = \frac{g}{\Theta} \Delta z \frac{\theta_{80m} - \theta_{41m}}{(V_{80m} - V_{41m})^2} \tag{2}$$

where g is the gravitational acceleration, Θ is the average potential temperature in the layer, Δz is the height difference between the two levels, and θ and V are the mean potential temperature and average horizontal wind speed at each level, respectively.

3 ~~Case studies~~ Comparison of turbulence characteristics in a fully turbulent with and intermittently turbulent night.

The nocturnal flow at the site is characterized by the superposition of turbulent and non-turbulent fluctuations. In a fully turbulent night, such as 15 November 2015 (Fig. 1), there is a clear dominant wind direction at all levels. In this case, it is very rare that the horizontal wind components switch sign above the canopy. In contrast, during the intermittent night of 14 November 2015 (Fig. 2), there is no dominant wind direction at any level above the canopy, as both horizontal components switch sign many times along the night. Low-frequency fluctuations are superposed on the turbulent fluctuations, causing the mean wind direction to change quadrants frequently throughout the night. Such fluctuations have been recently attributed to submeso flow (Mahrt, 2009), while in the pollutant dispersion community similar phenomena are often referred to as meandering (Oettl et al., 2005). The most relevant difference between the two nights regards the magnitude of the turbulent mixing (Table 1). All relevant turbulence statistics are significantly larger on 15 November than on 14 November. The relative difference of the turbulence statistics between nights increases steadily in the vertical. As an example, TKE at 41 m is 3.4 times larger in the turbulent night than in the intermittent case, while at 80 m, TKE is 8.2 times larger in the turbulent night. Similar increases occur for the corresponding ratios of σ_w and ~~u_*~~ $(\overline{u'w'}^2 + \overline{v'w'}^2)^{1/4}$ between the two nights.

Another interesting characteristic that indicates a contrast between the two nights shown in Fig. 1 and Fig. 2 regards the degree of vertical coupling across the levels, a phenomenon that has been observed by van Gorsel et al. (2011); Oliveira et al. (2013); Jocher et al. (2017). In the turbulent night, temperatures were always similar between the levels of 41 and 55 m, while at 80 m it was slightly warmer, but with the same cooling tendency throughout the period. CO_2 was correspondingly similar between 22 and 41 m, with the same tendencies and slightly lower values at 80 m. Although the mean trend is similar at 22 m and 41 m, substantial short time deviations towards higher CO_2 values were observed at 22 m (Fig. 1d). This is in line with the higher variability of CO_2 values in the lower canopy, as can be seen from the profiles (Fig. 3a). This higher variability and stronger gradient (in both CO_2 and O_3) in the lower canopy point to a decoupling of the sub canopy even in the turbulent night.
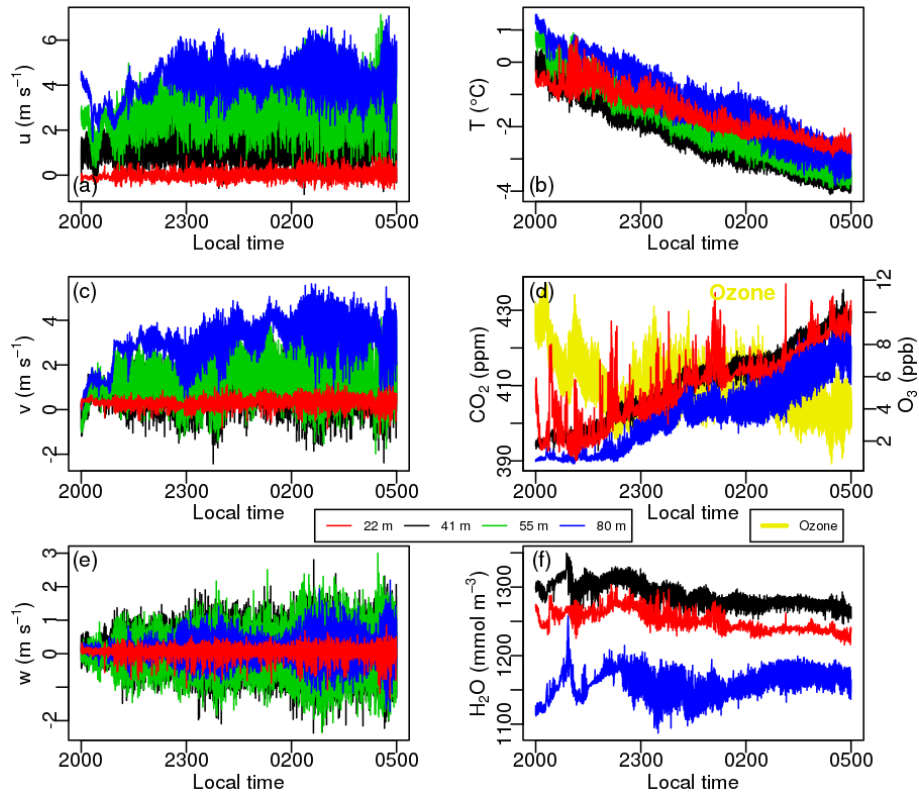
Figure 1. Time series of horizontal (a and c) and vertical (e) wind components, temperature perturbation from the 20:00 LT value at 41 m (b), CO_2 and O_3 (d) and water vapor (f) concentrations for the turbulent night.

As the 22 m level is within the maximum of the LAI (~24 m), which separates the upper canopy from the lower canopy, it will be influenced by both regimes. The gradients between 24 m and 38 m are always positive for O_3 and negative for CO_2. This can be related to the reactivity of O_3 as it reacts with compounds emitted from the soil (mainly NO) and plants (alkenes) and is not only taken up by stomata, but is also deposited to leaf surfaces in considerable amounts, especially under humid conditions
5 (Fuentes and Gillespie, 1992; Rummel et al., 2007). At night, CO_2 is emitted by soils and plants due to respiration, causing a negative gradient.

   All quantities showed much larger variation across the levels in the intermittent night (Fig. 2). Furthermore, sporadic events of coupling occurred during bursts of intermittent turbulence. (Fig. 2, shaded areas). It has been assumed that the turbulent bursts occurred whenever $\sigma_w > 0.15$ ms^{-1}. During these events of coupling, the gradients of temperature and CO_2 concen-
10 tration became sporadically smaller across the vertical, except for the 80-m level, indicating that the coupling induced by the events extended over a layer shallower than 80 m. In general, the temporal evolution of all scalars show a monotonic increase (of CO_2) or decrease (of temperature and O_3) throughout the turbulent night at all levels (Fig. 1). In the intermittent night, on the other hand, large increases and decreases of all scalars occur in small periods of time at all levels, except at 80 m. As
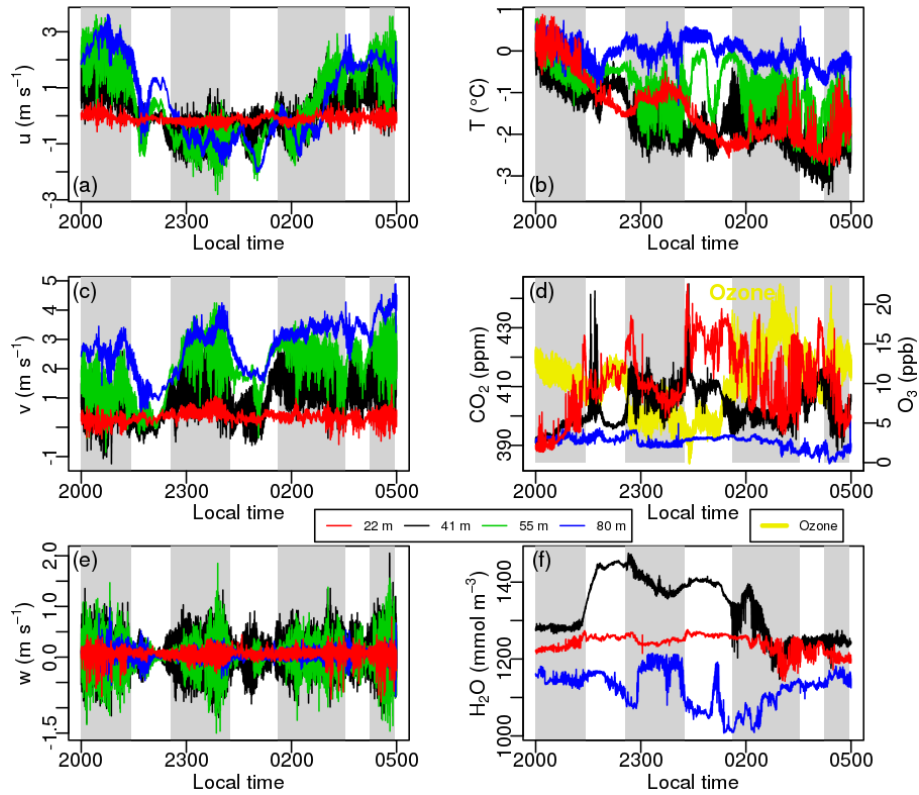
Figure 2. Same as in Fig. 1, but for the intermittent night. Shaded areas indicate intermittent turbulence bursts.
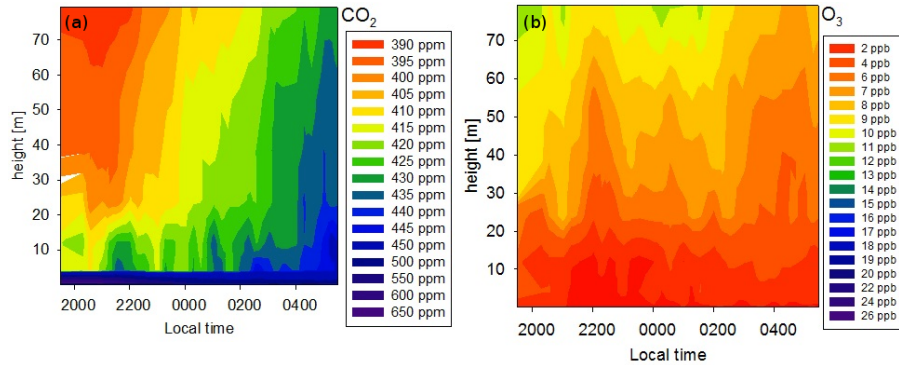

Figure 3. Concentrations of CO_2 (a) and O_3 (b) as a function of time and height for the turbulent night.

CO_2 has a clear source at the ground and O_3 has clear sink at the ground one can identified from the profiles if air is coming from aloft or from below (Fig. 4). Air from above is rich in O_3 and lower in CO_2, whereas air from below is rich in CO_2 but depleted in O_3. From this perspective, in the first event air is mixed down from aloft, while in the second event air is mixed

both upward and downward from the canopy top. In the third event, air is first mixed down and finally there is a burst of air going upwards from the canopy. At 80 m, temperature (Fig. 2b) and CO_2 (Fig. 2d) show much smaller fluctuations than at the other levels. This is further evidence that the stable boundary layer (SBL) thickness is shallower during the intermittent night, such that the canopy exchange fluxes do not affect the state of the atmosphere at 80 m. This fact contrasts strongly with the
5 steady trends of both scalars at 80 m during the turbulent night (Fig. 1b and Fig. 1d), which indicates that in this case, this level is fully coupled through turbulence to the canopy top. While in the turbulent nighttime, scalar fluxes did not vary substantially throughout the period (Fig. 5a, c, e, g), the most intense turbulent fluxes of sensible heat (Fig. 5b), CO_2 (Fig. 5d), O_3 (Fig. 5f), and latent heat (Fig. 5h) during the intermittent night occurred during these coupling periods (Table 2).
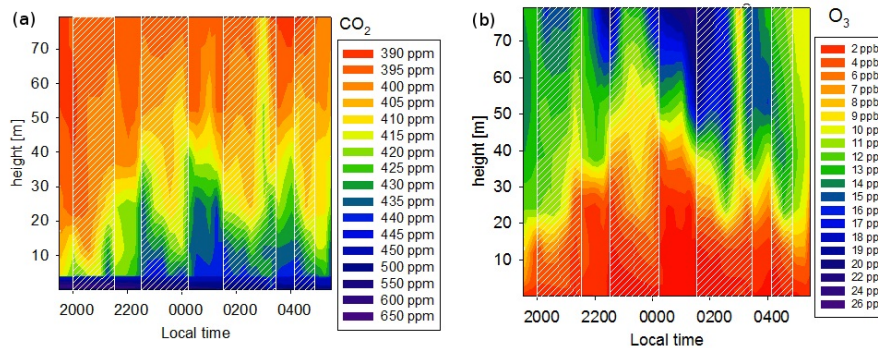
Figure 4. Same as in Fig. 3, but for the intermittent night.

Previous studies have reported that non-turbulent modes of the flow only become relevant when turbulence is weak (Acevedo
10 et al., 2014), a likely consequence of the diffusive nature of turbulence destroying the non-turbulent temporal and spatial variability of the atmospheric variables. This relationship between the turbulent and non-turbulent modes of the flow is illustrated by the TKE spectrum during both nights (Fig. 6). In the turbulent case, most of the energy is associated with turbulence, so that the most energetic time scale is near 10 s at all levels, except for 80 m (Fig. 6a). At this level, the longest time scales are the most energetic, but the 10-s turbulent maximum and a cospectral gap (near 100 s) are still evident. In contrast, in the
15 intermittent night of November 14, at all levels most energy prevails at the longest timescales provided by the decomposition method. This energy is associated with the low-frequency fluctuations responsible for the variability of the wind direction visible in Fig. 2a and Fig. 2c. These spectra confirm that when fully turbulent conditions prevail, the energy of the non-turbulent, low frequency modes of the flow is reduced considerably. The cospectra of the fluxes of sensible heat (Fig. 7a, d, f, i), CO_2 (Fig. 7b, g, j), O_3 (Fig. 7e), and latent heat (Fig. 7c, h, k) confirm the enhanced turbulent exchange of all quantities in the
20 turbulent night compared to the intermittent one. They also show that, consistently to what occurs with TKE, the non-turbulent exchange of these scalars is enhanced in the intermittent case. In particular, a significant low-frequency flux of CO_2 occurs at 22 m in the intermittent night, such that the total flux at this level is larger during the intermittent night (4.0 µmol m^2 s^{-1}, Table 3) than during the turbulent one (1.1 µmol m^2 s^{-1}), when all scales of the motion that are captured by the decomposition window are considered. This is in line with the idea that non-turbulent motions better penetrate the canopy. The same occurs
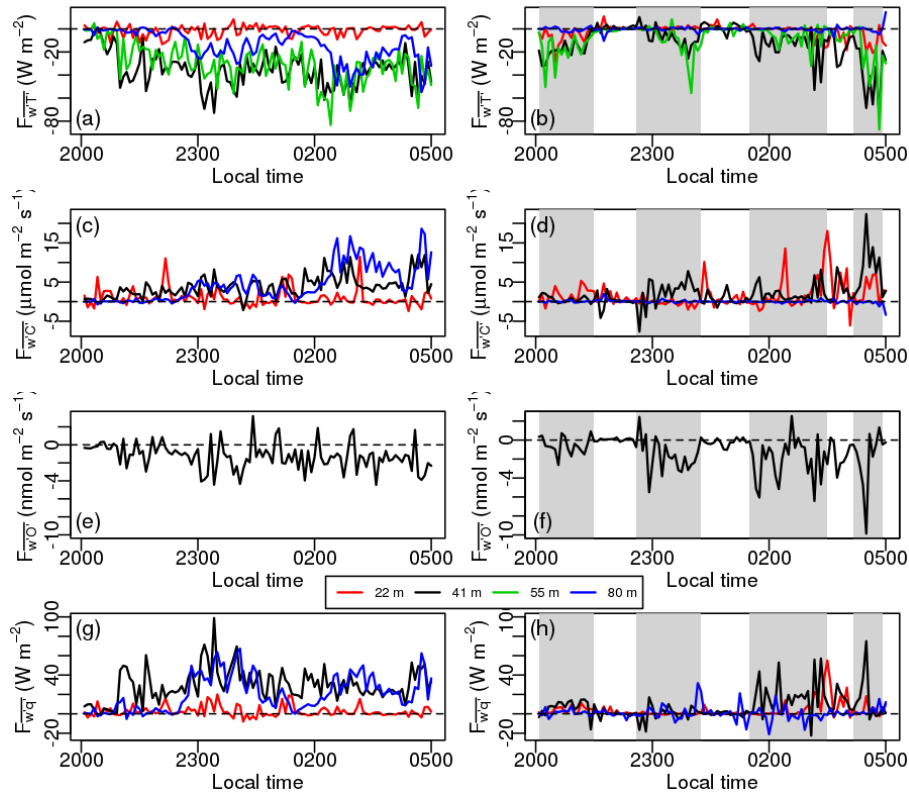
Figure 5. Fluxes of sensible heat (a and b), CO_2 (c and d), O_3 (e and f), and latent heat (g and h) for the turbulent night (left panels) and for the intermittent night (right panels). Shaded areas indicate intermittent turbulence bursts as shown in Fig. 2.

for latent heat, which shows a mean flux of 8.8 W m^{-2} in the intermittent night and of 2.8 W m^{-2} in the turbulent night, when all scales are considered. Even for 5-min fluxes, the larger fluxes occur still in the intermittent night (1.7 µmol m^2 s^{-1} versus 1.0 µmol m^2 s^{-1} in the turbulent night for CO_2, and 3.8 W m^{-2} versus 2.5 W m^{-2} for latent heat), but the differences are smaller. These numbers show that the low-frequency contributions dominate the exchange of CO_2 and moisture in the interior

5 of the forest in the intermittent night. Santos et al. (2016) found that the ~~dominant~~ time scales of ~~horizontal turbulent velocity fluctuations~~ the vertical velocity fluctuations and sensible heat flux within an Amazonian rain forest canopy (at a different site) ~~approach those of the non-turbulent maximum above the forest. They also found that the dominant time scales of the vertical velocity fluctuations and sensible heat flux within the canopy~~ are shifted towards larger values than those above it. Our results support these findings, adding the information that the ~~non-turbulent contribution may dominate the exchange~~ exchanges of

10 CO_2 and humidity ~~from the interior of the canopy in~~ within the canopy are also dominated by non-turbulent contribution at very stable nights~~as well~~. It is likely that the same process affects other scalars, such as O_3, whose concentrations are perturbed by intermittent events as shown in Fig. 4b.
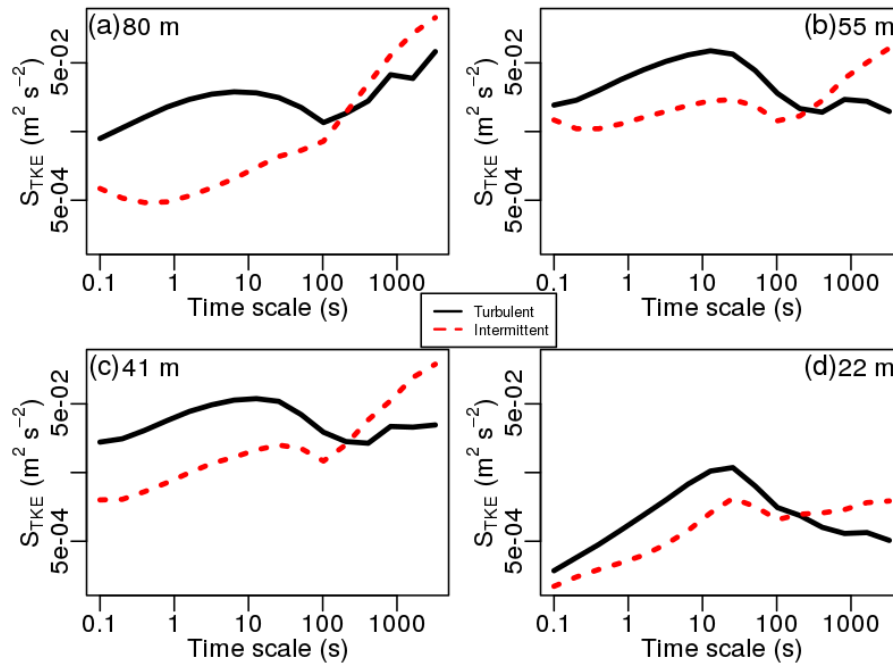
Figure 6. Average TKE spectra for the turbulent night (black solid lines) and the intermittent night (red dashed lines) for all levels, as indicated in each panel. In all panels, averages are performed over each time scale.

Another interesting contrast between the turbulent and intermittent nights can be seen for the sensible heat (Fig. 7a), CO_2 (Fig. 7b), and latent heat (Fig. 7c) cospectra at 80 m, which show an almost total absence of turbulent fluxes (time scale smaller than 100 s) of these quantities at this height in the intermittent case. This result explains the reduced fluctuations of temperature (Fig. 2b) and CO_2 (Fig. 2d) at 80 m in the intermittent night, adding evidence to the suggestion that the SBL may be rather shallow in that night, possibly such that the 80-m level is near its top. The large fluctuations of water vapor at 80 m (Fig. 2f) may be related to the enhanced low frequency flux of this quantity during this night (Fig. 7c).

Sun et al. (2012) found two regimes of nocturnal turbulence, distinguished by the turbulent kinetic energy (TKE) dependence on the mean wind speed. The fully turbulent regime, typically associated with weakly stable conditions, happens for mean wind speeds larger than a height dependent threshold and is characterized by TKE that steadily increases with wind speed. The other regime, associated with very stable conditions, has reduced turbulence intensities, which are very weakly dependent on the mean wind speed. Dias-Júnior et al. (2017) observed the two regimes above the forest at a site in the southwestern Amazon, finding that each is associated with an independent lognormal frequency distribution of quantities such as the turbulence dissipation rate. For the turbulent night of 15 November 2015 (Fig. **??**, crosses), the levels of 41 and 55 m remained in the large wind speed regime for the whole period, while the two different regimes could be observed only at the 80-m level. On the intermittent night, on the other hand (Fig. **??**, triangles), both regimes could be observed at all levels. Moreover, the connection intervals, given by shaded areas in Fig. 2, are generally in the large wind speed regime both at 41 m and 55 m (Fig. **??**, filled
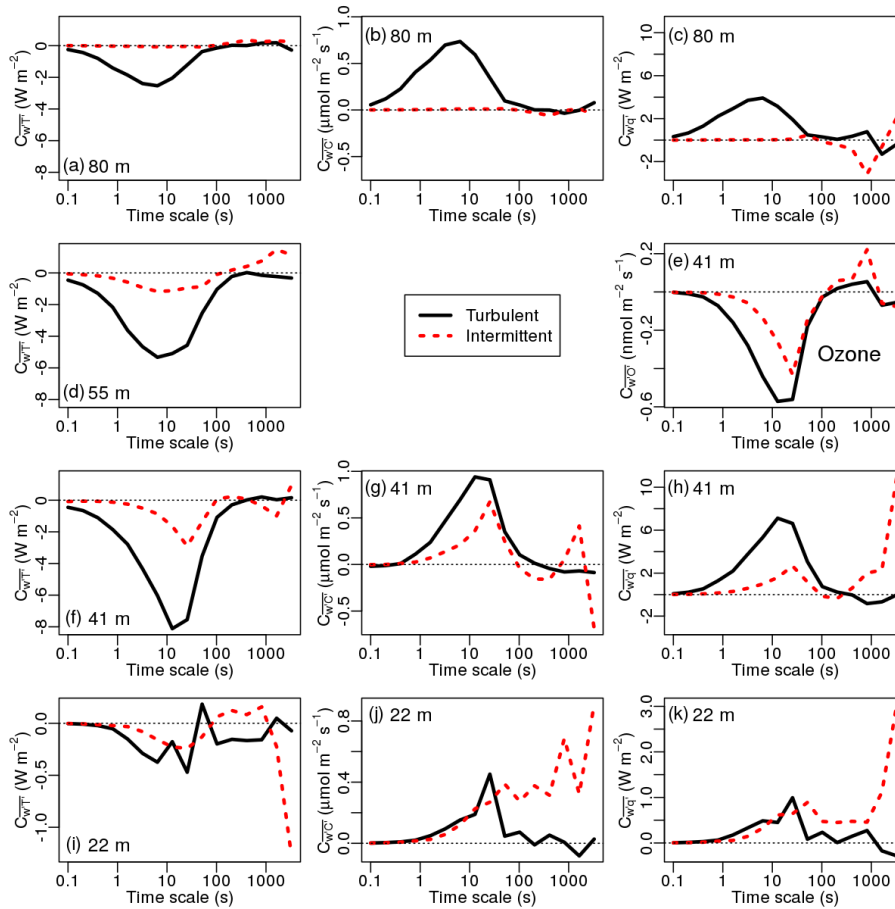
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Figure 7. Average cospectra of sensible heat (left panels), CO_2 (central panels), O_3 (e), and latent heat fluxes (c, h and k) for the turbulent night (black solid lines) and the intermittent night (red dashed lines) for all levels, as indicated in each panel. In all panels, averages are performed over each time scale.

triangles), while the decoupled periods are in most cases in the weak wind regime (Fig. **??**, open triangles). This is an important result, because it indicates that the intermittent bursts of turbulence observed above the canopy are intense enough to cause a regime transition. It means that, during these events, there is likely full vertical coupling over the vertical extent of the SBL (which is, at this time, shallower than 80 m). Therefore, scalars that are emitted from the canopy may be able to escape to higher levels in the boundary layer, as suggested by the episodic mixing of and above 70 m shown in Fig. 4.

Turbulent velocity scale (V_{TKE}) dependence on mean wind speed (V) for the turbulent night (crosses), shaded areas in Fig. 2 (filled triangles) and non-shaded areas in Fig. 2 (open triangles).

4 Mean spectra and cospectra over the 15 nights

Figure 8 shows the spectra and cospectra of the turbulent fluctuations and fluxes averaged over the entire period of 15 nights. The distinction between turbulent and non-turbulent fluctuations across the vertical is evident in the averaged TKE spectra (Fig. 8a). On average, a time scale around 100 s separates the two types of fluctuations. This can be seen by the fact that S_{TKE}
5 decreases with height above the canopy, following what typically happens to turbulent fluctuations in the stable boundary layer (Kruijt et al., 2000; Santos et al., 2016) only for time scales smaller than 100 s. For longer time scales, S_{TKE} increases with height. This result has been previously shown by Andreae et al. (2015) at the same tower, but for a different period. The average cospectra of all scalars show the largest turbulent fluxes at 41 m, with a sharp maximum at a timescale around 30 s. Within the canopy, systematic positive fluxes of CO_2 and latent heat happen at all time scales. The magnitude increases with increasing
10 time scale. Therefore, this is an indication that the CO_2 and humidity exchanges in the interior of the forest may be, to a large extent, caused by non-turbulent motion. This process will be further investigated later in this paper, when the scale dependence of the fluxes is compared to the stability within and above the canopy. The average sensible heat cospectrum at 22 m is negative for time scales smaller than 50 s and positive for larger time scales. Santos et al. (2016) showed that sensible heat fluxes near the forest floor are positive for all time scales, while they are negative for all time scales at the canopy top. At intermediate
15 heights, they tend to be like those shown in Fig. 8c. These authors also showed that the height within the canopy where the total sensible heat flux switches sign from upward at lower levels to downward at higher levels is stability dependent, increasing as it becomes more stable.

5 Dependence on stability over the 15 nights

The comparison of the fluxes determined with 5-min and 109-min time windows and their stability dependence for the entire
20 period of 15 nights provides interesting information on the scalar exchange processes within and above the canopy. Sensible heat flux at 22 m (Fig. 9i) switches sign at intermediate stability for both 5-min and 109-min averaging periods. This is in agreement with the result obtained by Santos et al. (2016), that the height where the sensible heat flux switches sign is stability dependent. In the present case, the critical value at which the 5-min sensible heat flux at 22 m switches sign is 0.25. At the same height, CO_2 (Fig. 9j) and latent heat (Fig. 9k) fluxes are similar with both averaging times at near-neutral conditions,
25 but become appreciably larger with 109-min windows than with 5-min windows as conditions become more stable. This result confirms the idea that the CO_2 and humidity exchange within the canopy at very stable conditions are dominated by processes with long time scales. At 41 m (Fig. 9e, f, g, h), 55 m (Fig. 9d), and 80 m (Fig. 9a, b, c) the 109-min fluxes of scalars are more erratic, sometimes with no clear dependence on $\mathrm{Ri_{can}}$. The 5-min fluxes, on the other hand show a tendency to decrease in magnitude with stability at all heights above the canopy. The 5-min TKE decreases with stability at all levels (Fig. 10), but at
30 80 m it becomes near zero for $\mathrm{Ri_{can}} > 0.2$. This is the same range of stability for which the 5-min fluxes of heat and CO_2 are virtually suppressed, indicating that in these very stable conditions the stable boundary layer thickness may be close to 80 m or shallower. When the same quantities are compared to $\mathrm{Ri_{top}}$ (Fig. 11), the most significant difference is that the 109-min fluxes at 41 m (Fig. 11e, f, g, h, black lines) and 55 m (Fig. 11d, black line) show a more systematic dependence on stability. This
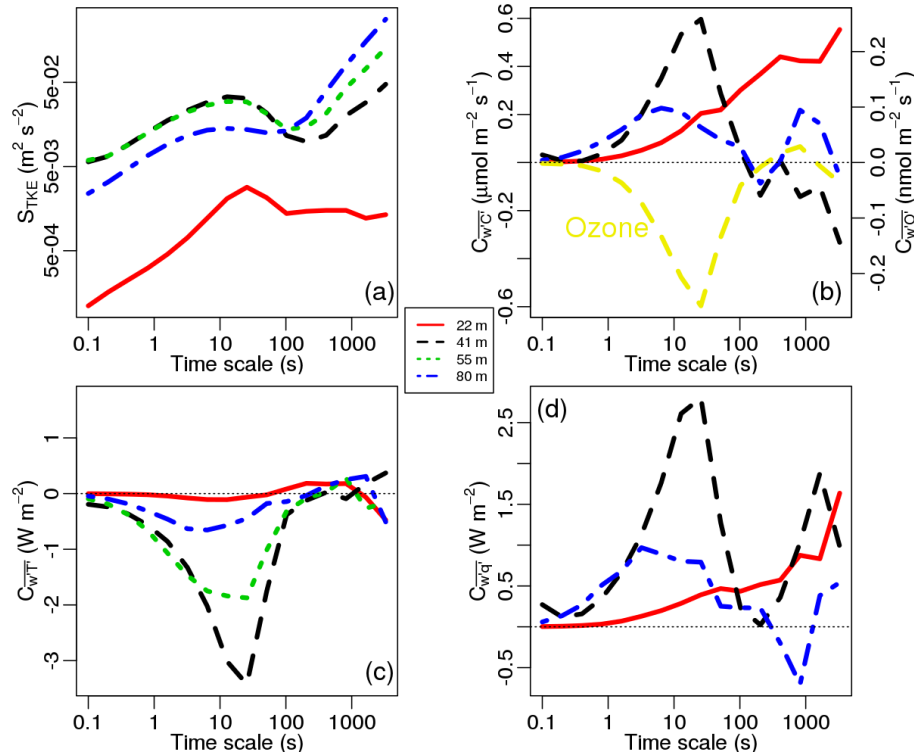
Figure 8. Average spectra of TKE (a) and cospectra of CO_2 and O_3 (b), sensible heat (c), and latent heat (d) fluxes for the entire dataset. In all panels, averages are performed over each time scale.

result indicates that the low-frequency exchange at the canopy top and above is controlled by the stability over a large distance above the forest.

The 0.25-Ri_{can} threshold, over which the 22-m turbulent sensible heat flux switches sign, is used to classify each series as weakly stable ($Ri_{can} \leq 0.25$) or very stable ($Ri_{can} > 0.25$). The average spectra for each of these classes are shown in Fig. 12.

5 At all levels, the low-frequency TKE is almost independent of stability, as for time scales larger than 100 s the TKE spectra of the very stable series approach those of the weakly stable ones. For smaller time scales, on the other hand, a significant distinction prevails, with TKE being many times larger under weakly stable conditions than during very stable conditions. The turbulent portions of the scalar fluxes above the canopy (41 m and above) respond accordingly, always being much larger in magnitude in the weakly stable cases (Fig. 13).

10 At 22 m, the criterion for classifying the time series ensures that the sensible heat flux at small time scales must be negative under the weakly stable conditions and positive under the very stable ones. This is shown in Fig. 13i, but it is remarkable that even in the weakly stable cases, there is a range of time scales for which the sensible heat flux within the canopy is upward. This range is also observed (however broader) in the very stable class, although in this case the negative maximum associated with turbulent exchange is absent. This result shows that low-frequency exchange of heat within the canopy is consistently
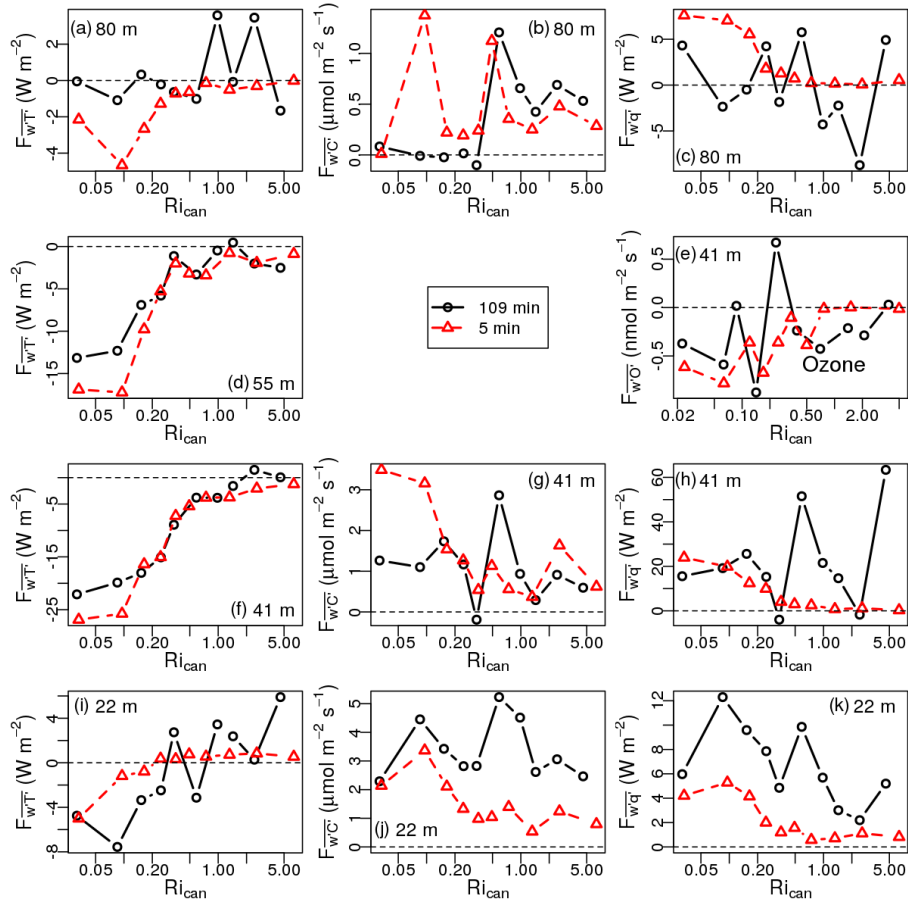
14

Figure 9. The dependence of sensible heat (left panels), CO_2 (central panels), O_3 (e), and latent heat (c, h and k) fluxes on canopy Richardson number (Ri_{can}) for all levels, as indicated in each panel. Fluxes have been determined with 5-min (red lines, triangles) and 109-min (black lines, circles) time windows.

upward, regardless of stability. Therefore, the main control exerted by stability regards the depth within the canopy where the downward turbulent heat flux penetrates (Santos et al., 2016). The total 22-m fluxes of CO_2 (Fig. 13j) are larger in the very stable class than in the weakly stable one, which is mainly caused by the contribution of time scales larger than 100 s, corroborating the results from the case studies (Sect. 3). For latent heat (Fig. 13k), the total flux is slightly larger in the weakly stable cases, and this result is mainly caused by the contributions of fluxes with time scales smaller than 100 s.
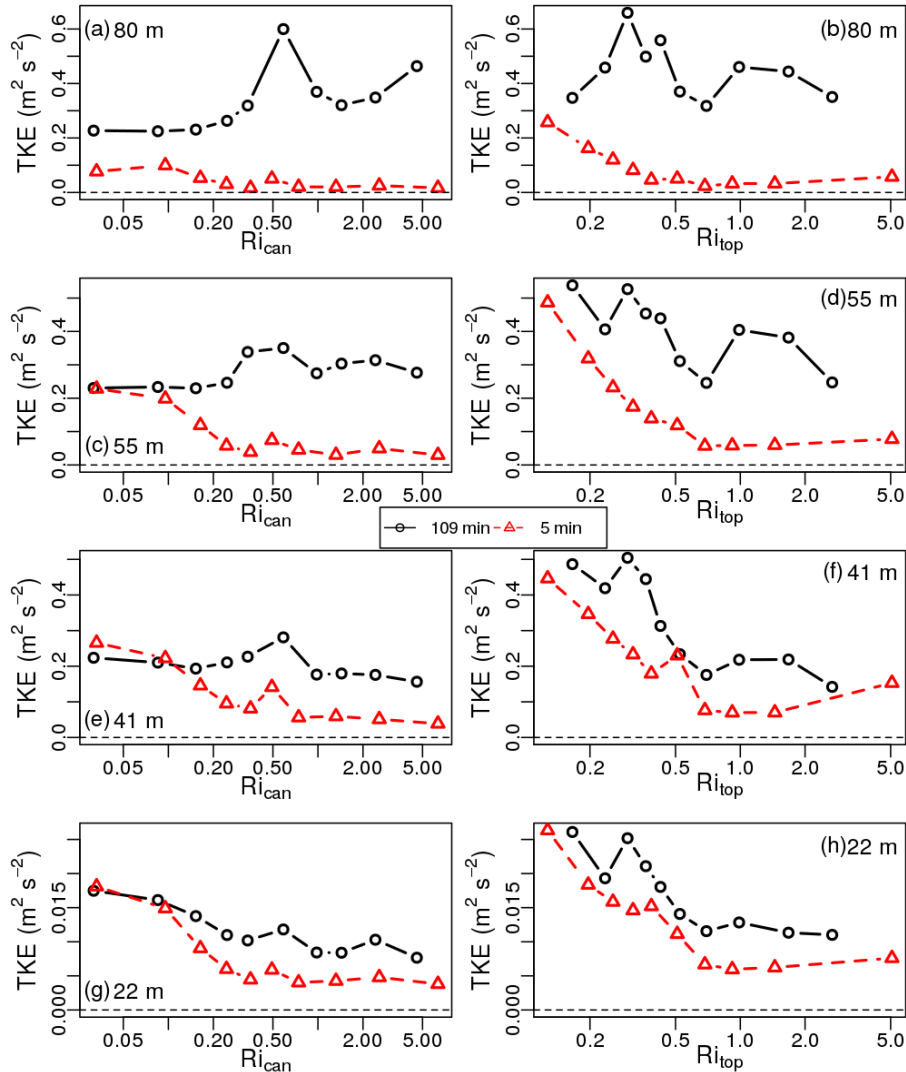
Figure 10. The dependence of TKE on canopy Richardson number (Ri_{can}, left panels) and top Richardson number (Ri_{top}, right panels) for all levels, as indicated in each panel, using 5-min (red lines, triangles) and 109-min (black lines, circles) time windows.

6   Conclusions

The main novelty of the present study has been a detailed analysis of different scalar fluxes and their time scales within and above a rain forest canopy at night. The data was collected at the Amazon Tall Tower Observatory (ATTO) and included fluxes of CO_2, O_3, latent and sensible heat. The most relevant findings include:
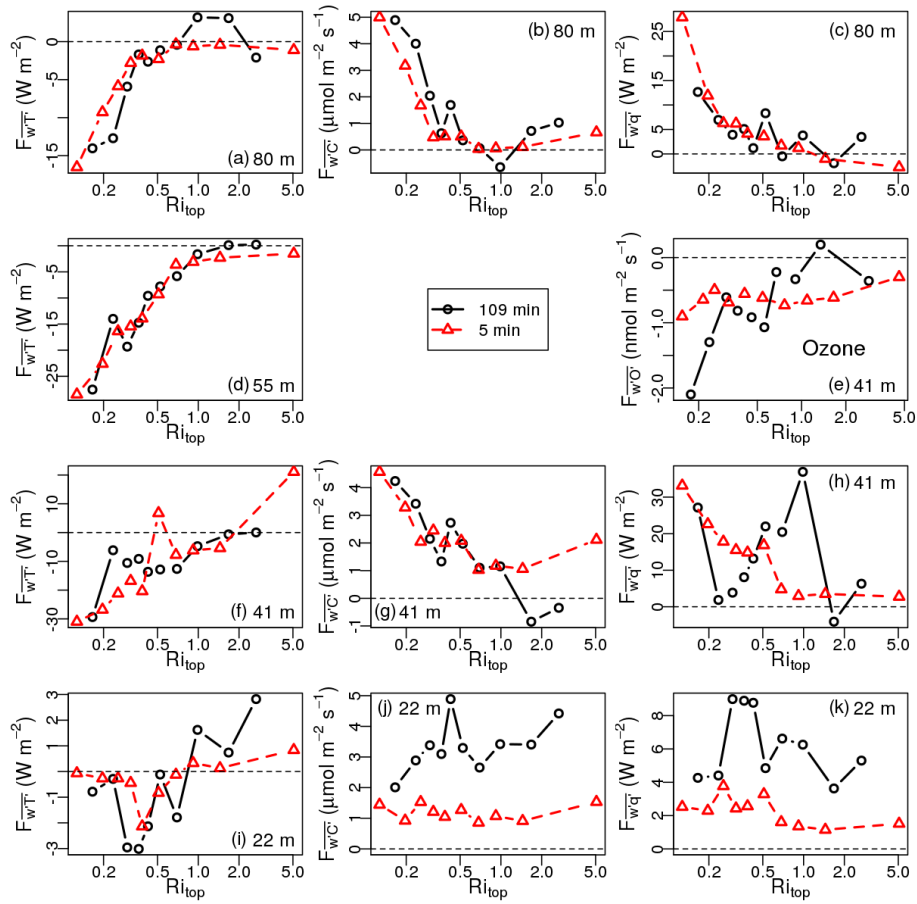
Figure 11. Same as in Fig. 9, but for dependence on top Richardson number (Ri_{top}).

- Within the canopy, fluxes of CO_2 and latent heat are dominated by processes with long time scales. Given that such low-frequency exchange tends to be enhanced in very stable conditions, the total scalar flux within the canopy may be larger in very stable nights than in weakly stable ones;

- In very stable nights, turbulent fluxes are effectively suppressed at 80 m, indicating that a very shallow stable boundary layer (SBL) may exist in these situations;

- Intermittent turbulence may produce very large fluxes and affect concentrations of CO_2 and O_3 from near the SBL top down to the middle of the canopy.

Although low-frequency contributions to the fluxes are enhanced during very stable nights, their inclusion into the scalar budgets is not enough to bring the nocturnal fluxes in the very stable nights close to those observed during fully turbulent conditions. Processes such as drainage flows or local storage may account for the differences.
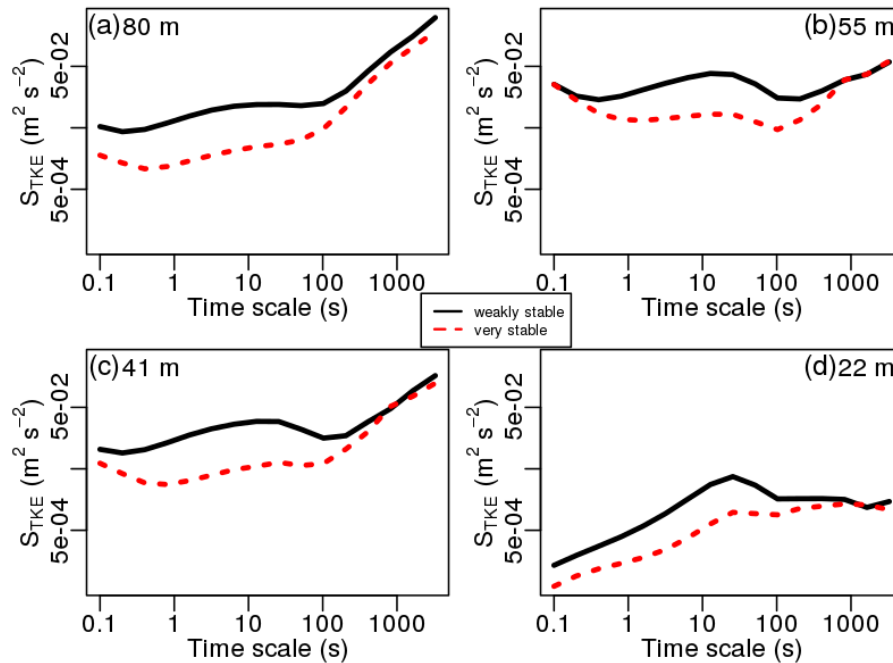
17

Figure 12. Average TKE spectra for the weakly stable (black solid lines) and the very stable (red dashed lines) cases for all levels, as indicated in each panel. In all panels, averages are performed over each time scale.

~~The majority of the flux~~ Turbulent exchange is always important just above the canopy (41-m level, in this case) ~~happens through turbulent exchange. Although no relevant systematic low-frequency contribution to the total fluxes has been found at 41 m for any scalar analyzed (Fig. 8), this result only holds for the average spectra, with appreciable variability among cases, especially in the most stable cases~~ but, in very stable nights, the non-turbulent contribution has to be also considered. Campos et al. (2009) found the low-frequency contribution of CO_2 fluxes above a similar Amazonian canopy to be seasonally dependent, a result that could not be examined with the present data set.

A fully instrumented 320-m tower is expected to start operating continually in the near future at the ATTO site. It will allow addressing questions such as the seasonality of the exchange at different scales, as well as the thickness of the SBL and the nature of the scalar exchange within and above the canopy in much more detail. In this sense, the results of the present study will provide important guidelines for the future investigations.
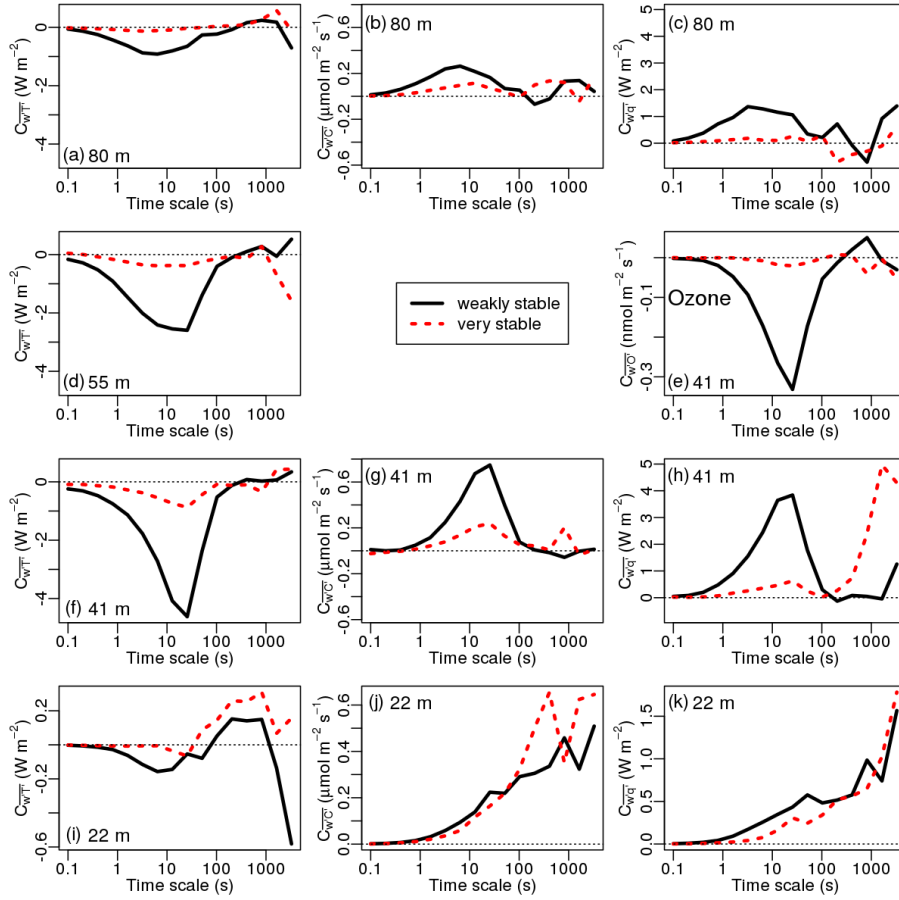
Figure 13. Average cospectra of sensible heat (left panels), CO_2 (central panels), O_3 (e), and latent heat fluxes (c, h and k) for the weakly stable (black solid lines) and the very stable (red dashed lines) cases for all levels, as indicated in each panel. In all panels, averages are performed over each time scale.

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Table 1. 5-min turbulence statistics averaged for each night analyzed in Sect. 3.

level	σ_w	$u_* \left(\overline{u'w'}^2 + \overline{v'w'}^2 \right)^{1/4}$	TKE	σ_w	$u_* \left(\overline{u'w'}^2 + \overline{v'w'}^2 \right)^{1/4}$	TKE
		14 November 2015 Intermittent night			15 November 2015 Turbulent night	
(m)	$(\mathrm{m\,s^{-1}})$	$(\mathrm{m\,s^{-1}})$	$(\mathrm{m^2\,s^{-2}})$	$(\mathrm{m\,s^{-1}})$	$(\mathrm{m\,s^{-1}})$	$(\mathrm{m^2\,s^{-2}})$
22	0.07	0.04	0.01	0.11	0.07	0.03
41	0.19	0.14	0.13	0.39	0.30	0.44
55	0.15	0.10	0.10	0.37	0.27	0.41
80	0.06	0.04	0.02	0.18	0.13	0.16

Table 2. 5-min TKE and fluxes averaged for shaded and non-shaded areas in the intermittent night (see Fig. 2).

	level (m)	F_C $\mu mol\, m^2\, s^{-1}$	F_H $W\, m^{-2}$	F_q $W\, m^{-2}$	TKE $(m^2\, s^{-2})$	F_O $nmol\, m^2\, s^{-1}$
Shaded	22	1.6	-2.9	3.7	0.01	
	41	2.6	-15.9	9.6	0.16	-1.5
	55	-	-12.5	-	0.13	
	80	0.0	-0.4	0.9	0.02	
Non-shaded	22	1.6	-0.4	3.7	0.01	
	41	1.0	-3.2	2.9	0.06	-0.4
	55	-	-2.4	-	0.03	
	80	0.0	-0.1	0.4	0.02	

The column span header: 14 November 2015 – Intermittent night

Table 3. TKE and fluxes averaged for each night analyzed in Sect. 3 using a time window of 5 and 109 min.

	level (m)	F_C μmol m^2 s^{-1}	F_H W m^{-2}	F_q W m^{-2}	TKE (m^2 s^{-2})	F_O nmol m^2 s^{-1}	F_C μmol m^2 s^{-1}	F_H W m^{-2}	F_q W m^{-2}	TKE (m^2 s^{-2})	F_O nmol m^2 s^{-1}
			14 November 2015 Intermittent night					15 November 2015 Turbulent night			
109 min	22	4.0	-2.0	8.8	0.02		1.1	-2.1	2.8	0.03	
	41	1.2	-7.9	22.9	0.29	-0.9	3.5	-37.4	29.7	0.49	-2.3
	55	-	-2.4	-	0.48		-	-32.4	-	0.49	
	80	-0.1	1.0	-2.4	0.47		3.9	-13.4	20.4	0.29	
5 min	22	1.6	-2.1	3.7	0.01		1.0	-1.3	2.6	0.03	
	41	2.1	-11.9	7.5	0.13	-1.2	3.8	-35.6	29.4	0.44	-1.3
	55	-	-9.3	-	0.10		-	-30.2	-	0.41	
	80	0.0	-0.3	0.7	0.02		4.3	-14.1	20.8	0.16	