Measurement-based climatology of aerosol direct radiative effect, its sensitivities, and uncertainties from a background southeast U.S. site

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Abstract

Aerosol optical properties measured at Appalachian State University's co-located NASA AERONET and NOAA ESRL aerosol network monitoring sites over a nearly four-year period (June 2012 thru Feb 2016) are used, along with satellite-based surface reflectance measurements, to study the seasonal

5 variability of diurnally averaged clear sky aerosol direct radiative effect (DRE) and radiative efficiency (RE) at the top-of-atmosphere (TOA) and at the surface. Aerosol chemistry and loading at the Appalachian State site are likely representative of the background southeast U.S. (SE U.S.), home to high summertime aerosol loading and one of only a few regions not to have warmed during the 20th century. This study is the first multi-year 'ground truth' DRE study in the SE U.S., using aerosol network data

10 products that are often used to validate satellite-based aerosol retrievals. The study is also the first in the SE U.S. to quantify DRE uncertainties and sensitivities to aerosol optical properties and surface reflectance, including their seasonal dependence.

Median DRE for the study period is -2.9 Wm^{-2} at the TOA and -6.1 Wm^{-2} at the surface. Monthly median and monthly mean DRE at the TOA (surface) are -1 to -2 Wm^{-2} (-2 to -3 Wm^{-2}) during winter

15 months and -5 to -6 Wm^{-2} (-10 Wm^{-2}) during summer months. The DRE cycles follow the annual cycle of aerosol optical depth (AOD), which is 9 to 10 times larger in summer than in winter. Aerosol RE is anti-correlated with DRE, with winter values 1.5 to 2 times more negative than summer values. Due to the large seasonal dependence of aerosol DRE and RE, we quantify the sensitivity of DRE to aerosol optical properties and surface reflectance, using a calendar day representative of each season (DEC 21

20 for winter; MAR 21 for spring, JUN 21 for summer, and SEP 21 for fall). We use these sensitivities along with measurement uncertainties of aerosol optical properties and surface reflectance to calculate DRE uncertainties. Aerosol DRE at both the TOA and surface is most sensitive to changes in AOD, followed (in order) by single-scattering albedo (ω_0), scattering asymmetry parameter (g), and surface reflectance (R). One exception is under the high summertime aerosol loading conditions, when sensitivity of TOA

25 DRE to ω_0 is comparable to that of AOD. While DRE sensitivity to AOD varies by only ~25 to 30 % with season, DRE sensitivity to ω_0, g, and R vary by factors of 10 to 20 with season. Since the measurement uncertainties of AOD, ω_0, g, and R are comparable at Appalachian State, their relative contributions to DRE uncertainty are roughly proportional to their (seasonally dependent) DRE sensitivity values, which

suggests that the seasonal dependence of DRE uncertainty must be accounted for. Clear sky aerosol DRE uncertainty at the TOA (surface) ranges from 0.44 Wm^{-2} (0.73 Wm^{-2}) for DEC to 0.90 Wm^{-2} (1.3 Wm^{-2}) for JUN. Expressed as a fraction of DRE computed using monthly median aerosol optical properties and surface reflectance, the DRE uncertainties at TOA (surface) are 16 to 20 % (12 to 20 %) for MAR, JUN, and SEP and 48 % (49 %) for DEC. The relatively low DRE uncertainties are largely due to the low uncertainty in AOD measured by AERONET. Use of satellite-based AOD measurements by MODIS in the DRE calculations increases DRE uncertainties by a factor of 2.5 to 5 and DRE uncertainties are dominated by AOD uncertainty for all seasons.

1 Introduction

Predictions of future climate change resulting from projected increases in carbon dioxide are limited in part by uncertainties in the direct and indirect radiative forcing due to aerosols (Andreae, et.al, 2005). On a global average, the measurement-based estimates of aerosol direct radiative effect (DRE) are 55 to 80 % greater than the model-based estimates. The differences are even larger on regional scales and for the anthropogenic component (Yu, et.al, 2006). Such measurement-model differences are a combination of differences in aerosol amount (optical depth), single-scattering properties, surface albedo, and radiative transfer schemes (Yu et al., 2006). One of the high-priority tasks recommended (Remer, et.al 2009) to reduce the uncertainty in aerosol radiative effects is to "*Maintain, enhance, and expand the surface observation networks measuring aerosol optical properties for satellite retrieval validation, model evaluation, and climate change assessments.*"

The southeast U.S. (SE U.S.) is home to some of the highest warm-season aerosol loading in the U.S. (Goldstein et al., 2009) and is also one of only a few regions not to have not exhibited a warming trend in the 20th century (Menne et al., 2009). Several studies conducted during the past two decades have attempted to quantify aerosol DRE in the SE U.S. Yu et al. (2001) applied 34 days of aerosol optical property measurements near Mount Mitchell, NC from June thru December 1995 to estimate variability in SE U.S. aerosol DRE and atmospheric absorption by aerosols, along with an estimate of 'annually averaged' DRE. Carrico et al. (2003) applied measurements of aerosol optical depth (AOD) and other aerosol optical properties as part of the Atlanta Supersite 1999 study to estimate summer top-of-

atmosphere (TOA) DRE in urban Atlanta, GA. Goldstein et al. (2009) used region and time-averaged AOD near 550 nm, measured by the Multi-angle Imaging Spectrometer (MISR) aboard the polar-orbiting Terra satellite from 2000 thru 2007, as inputs to a first-order radiative transfer calculation (Haywood and Shine, 1995) to show that high summer AOD in the SE US led to more negative aerosol TOA DRE in

5   summer than winter (by 3.9 Wm^{-2}). Goldstein et al. (2009) hypothesized that this summer regional cooling effect was dominated by secondary organic aerosols, resulting from the oxidation of biogenic volatile organic compounds in the presence of anthropogenic NOx and SO$_2$. However, their DRE calculation used assumed values (rather than measured) for surface reflectance and for all aerosol properties except AOD, and they did not consider seasonal variations in these properties. Alston and

10 Sokolik (2016) applied 12 years (2000 thru 2011) of AOD at 550 nm, cloud fraction, and surface albedo measured by the Moderate Resolution Imaging Spectrometer (MODIS) aboard Terra, along with single-scattering albedo near 550 nm from MISR, as inputs to the same TOA DRE equation (Eq. (2) of Haywood and Shine, 1995) used by Goldstein et al. (2009). Their primary objectives were to study TOA DRE seasonal variability and long-term trends in the SE U.S., in the context of changes in AOD, cloud fraction,

15 and surface albedo. They concluded that AOD was a major driver of regional TOA DRE (as compared to surface albedo and cloud fraction) and they also reported a decreasing linear trend in MODIS Terra AOD, which contributed to a small increasing trend (i.e. less negative) in TOA aerosol DRE. However, the sensitivities of DRE to aerosol single-scattering properties and surface reflectance were not explicitly quantified, nor were DRE uncertainties. Estimates of aerosol DRE using MODIS-measured AOD have

20 higher uncertainties than those based on spectral AOD measured at NASA Aerosol Robotic Network (AERONET; Holben et al., 1998) sites, as discussed in Sect. 5.3 of this paper. The MODIS Collection 5.1 AOD also has been found to possess a consistently negative bias (0.02 to 0.03) over the rural SE U.S. Appalachian State AERONET site (Sherman et al., 2016a), which could lead to an under-estimation of aerosol DRE (i.e. less negative).

25   Ground-based sites as part of the NOAA Earth Systems Research Laboratory (NOAA ESRL; Delene and Ogren, 2002), NASA AERONET, and NASA Micro-pulsed Lidar (NASA MPLNET; Welton et al., 2001) federated aerosol monitoring networks possess continuous long-term records of aerosol optical properties used to evaluate aerosol DRE (McComiskey et al., 2008; Michalsky, et al., 2006). Established

in 2009, the Appalachian Atmospheric Interdisciplinary Research Facility at Appalachian State University is home to the only co-located NOAA ESRL, NASA AERONET, and (since 2016) NASA MPLNET sites in the SE U.S. Aerosol chemistry and loading at the semi-rural, high-elevation Appalachian State site (referred to as APP in this paper) are likely representative of the background SE

5 U.S. (Link et al., 2015). As such, APP is well-positioned to improve understanding of aerosol DRE in the SE U.S., including seasonal DRE variability, sensitivities, and uncertainties, as recommended by Remer et al. (2009).

The objective of this paper is to complement previous studies of aerosol DRE in the SE U.S. through a detailed, multi-year study of aerosol DRE seasonal variability, sensitivities, and uncertainties from a

10 single ground-based aerosol network site. Specifically, we

1. Quantify the seasonal variability in diurnally averaged, clear sky aerosol DRE and direct radiative efficiency (RE=DRE per unit AOD) at APP, both at the TOA and at the surface, along with seasonal variability in aerosol and surface properties influencing DRE (Sect. 5.1)

2. Quantify the sensitivity of DRE to key aerosol and surface properties, including any seasonal

15 dependence (Sect. 5.2)

3. Apply sensitivities (2) along with measurement uncertainties to calculate the uncertainty in DRE and the relative contributions of aerosol and surface properties to DRE uncertainty (Sect. 5.3)

Daily averaged aerosol optical properties measured on 418 days between June 2012 and February 2016 are used along with monthly averaged spectral surface reflectance measured by MODIS to study the

20 annual DRE and RE cycles at APP (Sect. 5.1). In Sect. 5.2, we follow a similar approach to that used by McComiskey et al (2008) to quantify DRE sensitivity to AOD, single-scattering albedo (ω_0), scattering asymmetry parameter (g), and surface reflectance (R). Sensitivities of DRE are then used along with measurement uncertainties in AOD, ω_0, g, and R to estimate the resulting uncertainties in DRE at the TOA and at the surface (Sect. 5.3). The use of well-established measurement protocols developed by

25 NOAA ESRL (Delene and Ogren, 2002) and NASA AERONET (Holben et al., 1998), and possessing known uncertainties (Eck et al., 1999; Sherman et al., 2015), facilitates the first study of DRE sensitivities and measurement uncertainty in the SE U.S., with results that are directly comparable with other regions.

This paper differs from the aerosol DRE sensitivity and uncertainty analysis conducted by McComiskey et al. (2008) in that it addresses a different geographic region and in the following additional ways:

1. McComiskey et al (2008) considered generally representative properties of three surface aerosol sites and no seasonal dependence of DRE, while this study uses direct measurements and focuses on the seasonal DRE dependence at a single site

2. The use of measured values for all aerosol properties allows for us to consider their covariances in the DRE uncertainty calculations

3. We compare DRE uncertainties using ground-based AOD measurements made as part of AERONET with those using satellite-based AOD measurements from MODIS

For clarity, it is important to distinguish between aerosol DRE and the often-referenced aerosol direct radiative forcing (DRF), in addition to defining 'clear sky' DRE. Direct radiative effect refers to the difference in net radiative fluxes (Eq. (5)) at a given atmospheric level (often the TOA or surface) with and without the presence of atmospheric aerosols, while DRF refers to the anthropogenic component (Kaufman et al., 2005). Clear sky DRE refers to DRE calculated assuming cloud-free conditions. Most studies neglect cloud effects not only for simplicity but also because satellite-based aerosol retrievals can only be made in the absence of clouds. First-order DRE calculations such as those provided in Haywood and Shine (1995) account for aerosol DRE in the presence of clouds by multiplying the clear sky DRE by the cloud-free sky fraction. For clarity, we also include in Appendix A a table of commonly used acronyms and symbols used in this paper (Table A1).

2   Site description

The APP site is located at the highest point on the Appalachian State University campus in Boone, NC (Fig.1). Lower tropospheric aerosols are sampled from a 34 m tower as part of NOAA ESRL, from which aerosol optical and microphysical properties (Sect. 3.1.2) are measured by in situ instruments (Sherman et al., 2015). Vertical profiles of aerosols and clouds have also been measured continuously by a micro-pulsed lidar (MPL) as part of NASA MPLNET (Welton et al.,2001) since March 2016. Lidar-measured vertical profiles of normalized relative aerosol backscatter were made periodically from 2011 thru 2014

(prior to joining MPLNET) but have no quality assurance and therefore are not used in this paper, other than qualitative inspection to verify that aerosols are largely confined to the lowest 1 to 2 km of atmosphere above APP. The region surrounding APP is heavily forested and possesses a diversity of elevations (< 300 m to > 2000 m) and a variety of weather regimes (i.e.., winter storms, convective cells,
5   dying tropical cyclones, and stagnant summertime episodes). The region also includes a diversity of anthropogenic and biogenic aerosol sources. Lower tropospheric aerosol light scattering and absorption measured at APP is dominated by particles with diameter less than 1 μm (Sherman et al., 2015) and sub-1 μm aerosol mass consists primarily of organics, with lower levels of sulfates (supplement to Link et al., 2015). Summer AOD in the SE U.S. (including APP) is influenced by isoprene-derived secondary organic
10   aerosol (Goldstein *et al.*, 2009; Link et al., 2015). A biomass burning influence is present in winter aerosol mass concentrations measured at APP (supplement to Link *et al.*, 2015), likely due to residential wood burning in the region. Wood-burning stoves serve as the primary heating source for 6.2 % of occupied housing units in Watauga County (U.S. Census Bureau, 2010) and likely a larger percentage of housing units in the surrounding rural mountain communities.

15   ## 3   Measurements used by the radiative transfer model to calculate aerosol DRE

3.1   Aerosol optical properties

The following aerosol optical properties (including their dependence on wavelength) are standard inputs to radiative transfer models used to calculate aerosol DRE at the TOA and surface: (1) aerosol optical depth (AOD); (2) single-scattering albedo (ω_0); and (3) scattering asymmetry parameter (g). For
20   calculation of broadband, diurnally averaged aerosol DRE, we form daily averages of each optical property and interpolate or extrapolate to their values at 38 equally spaced wavelengths over the 250 to 4000 nm range.

3.1.1   Aerosol optical depth

The CIMEL sunphotometer deployed at APP (known as 'Appalachian_State' within AERONET)
25   measures direct solar radiance at eight wavelengths (λ=340, 380, 440, 500, 675, 870, 940, and 1020 nm) and sky radiance at four of these wavelengths (λ=440, 675, 870, and 1020 nm), using standard AERONET

protocols (Holben, et al., 1998). The direct solar radiance measurements are used to calculate AOD at each of the eight wavelengths except 940 nm, using the Beer-Lambert-Bouguer equation (Holben, et al., 1998). Direct solar radiance measurements are made at optical air mass intervals of 0.25, corresponding to every ~15 minutes near noon and more often near dawn and dusk. Only Level 2 AERONET AOD (cloud-screened, calibrated) is used in this study. The uncertainty for Level 2 AOD is small enough (0.01 to 0.015; Eck et al., 1999) so that AERONET serves as 'ground truth' for comparisons with satellite-derived AOD (Levy et al., 2010; Hyer et al., 2011).

Sky radiance measurements made at AERONET sites are used to derive column-averaged aerosol properties including size distributions and ω_0. Single-scattering albedo can only be reliably retrieved to within ~0.03 for AOD (λ=440 nm) $\geq$ 0.40 (Dubovik et al., 2000). This high loading condition is only satisfied on 2 to 4 days per year at the Appalachian_State site and therefore AERONET ω_0 is not available for use in this study. AOD Ångström exponent (Å_{aod}) in the visible spectral range is typically computed as the slope of a linear fit of log (AOD) versus log (λ) using available wavelengths between 440 and 870 nm. It is used in this paper to wavelength-scale AOD (Fig.2a) using Eq. (1):

$$AOD(\lambda) = AOD(550nm)\left(\frac{550}{\lambda}\right)^{\overset{0}{\text{Å}}_{AOD}}$$ (1)

3.1.2 Single-scattering albedo and scattering asymmetry parameter

The primary aerosol measurements at APP, as part of the NOAA ESRL network, are aerosol light scattering (σ_{sp}), hemispheric backscattering (σ_{bsp}), and absorption (σ_{ap}) coefficients, reported at 450, 550, and 700 nm (Sherman et al., 2015). Each of these parameters is measured for both sub-10 µm particles (PM_{10}) and sub-1 µm particles (PM_1). We use PM_{10} values in this paper, even though σ_{sp}, σ_{bsp}, and σ_{ap} at APP are dominated by sub-1 µm particles (Sherman et al., 2015). Aerosol optical depth measured in the column is representative of aerosol of all sizes and larger particles can contribute greatly to aerosol light scattering. Thus, PM_{10} optical properties measured near the surface will be most comparable to the column AOD measurements.

A three-wavelength integrating nephelometer (Model 3563, TSI Inc., St. Paul, MN) is used for measurement of σ_{sp} (angular range of 7º to 170º) and σ_{bsp} (angular range of 90º to 170º). Aerosol light

absorption coefficients were determined by a three-wavelength Particle Soot Absorption Photometer (PSAP, Radiance Research, Seattle, WA) up until March 2015. A new light absorption instrument developed at NOAA ESRL (Continuous Light Absorption Photometer, CLAP; Ogren et al., 2013) then replaced the PSAP, after a one-year inter-comparison of the PSAP and CLAP instruments at APP. The

5 major difference between the CLAP and PSAP is that the CLAP has eight filter spots (versus one for PSAP) and can thus run nearly eight times longer between filter changes that require human intervention. The CLAP-measured σ_{ap} values are ~5 to 10 % lower than the PSAP (unpublished result). Aerosols entering the instruments are heated as needed to attain relative humidity RH ≤ 40 % to decouple the influences of aerosol amount and RH on σ_{sp} and σ_{bsp}. In-depth discussions of NOAA ESRL aerosol

10 sampling, measurement, and data quality assurance protocols are provided in Sherman et al (2015) and references therein. A scanning humidograph (Sheridan, et al., 2001) is employed at APP to measure the RH dependence of σ_{sp} and σ_{bsp}. We use this and co-located measurements of RH to scale σ_{sp} and σ_{bsp} to ambient RH for the dataset used in this paper.

Single-scattering albedo at each of the 38 wavelengths supplied to the Santa Barbara DISORT

15 Radiative Transfer model (SBDART; Ricchiazzi et al., 1998) is calculated by wavelength-scaling the σ_{sp} and σ_{ap} values at 550 nm, using scattering and absorption Angstrom exponents ($Å_{sp}$, and $Å_{ap}$), which are calculated from the σ_{sp} and σ_{ap} values at 450 and 700 nm (Sherman et al., 2015)

$$\omega_0(\lambda) = \frac{\sigma_{sp}(\lambda)}{\sigma_{sp}(\lambda) + \sigma_{ap}(\lambda)} = \frac{\left(\frac{550}{\lambda}\right)^{A_{sp}^0} \sigma_{sp}(550nm)}{\left(\frac{550}{\lambda}\right)^{A_{sp}^0} \sigma_{sp}(550nm) + \left(\frac{550}{\lambda}\right)^{A_{ap}^0} \sigma_{ap}(550nm)} \qquad (2a)$$

Dividing the numerator and denominator by σ_{sp}(550 nm) + σ_{ap}(550 nm) allows Eq. (2a) to be re-written

20 solely in terms of the intensive aerosol optical properties ω_0 (550 nm), $Å_{sp}$, and $Å_{ap}$.

$$\omega_0(\lambda) = \frac{\left(\dfrac{550}{\lambda}\right)^{A_{sp}^0} \omega_0(550nm)}{\left(\dfrac{550}{\lambda}\right)^{A_{sp}^0} \omega_0(550nm) + \left(\dfrac{550}{\lambda}\right)^{A_{ap}^0}\left(1 - \omega_0(550nm)\right)} \tag{2b}$$

Spectral ω_0 calculated using Eq. (2b) is displayed graphically for each season in Fig.2b. The uncertainty in low-RH aerosol ω_0 is ~0.015 (Sherman et al., 2015). This likely serves as a lower bound to $\Delta\omega_0$, due to unknown uncertainties in the correction of σ_{sp} to ambient RH (Titos et al., 2016).

Scattering asymmetry parameter is calculated at 450 nm, 550 nm, and 700 nm, based on the hemispheric backscatter fraction b= σ_{bsp} / σ_{sp} and the parameterization (Andrews et al., 2007)

$$g = 0.9893 - 3.96b + 7.46b^2 - 7.14b^3 \tag{3}$$

10 The uncertainty Δb is only 2 to 3 % for low-RH aerosol at APP, when covariance effects between σ_{bsp} and σ_{sp} are accounted for (Sherman et al., 2015). By taking the differential of Eq. (3) this leads to an uncertainty estimate Δg <0.01 for a typical value of b~0.15 measured at APP (Sherman et al., 2015). As with $\Delta\omega_0$, the estimate for Δg likely serves as a lower bound, due to unknown uncertainties in the correction of b to ambient RH (Titos et al., 2016). Asymmetry parameter is then wavelength-scaled to the 15 38 wavelengths used by SBDART following McComiskey et al. (2008), their Eq. (8):

$$g(\lambda) = g(550nm)\frac{1 + \left(\dfrac{550}{\lambda_g}\right)^2}{1 + \left(\dfrac{\lambda}{\lambda_g}\right)^2} \tag{4}$$

Spectral g calculated using Eq. (4) is displayed graphically for each season in Fig. 2c. Following McComiskey et al. (2008), we use 5000 nm for λ_g. However, McComiskey et al.(2008) noted that the exact value does not significantly alter the calculated DRE values and sensitivities through spectral

dependence of *g,* but ensures physically reasonable results for very small (larger) particle sizes when reaching the Rayleigh (Mie) limit.

An assumption used in this paper is that ω_0 and *g* measured near the surface are representative of these properties in the column, which is typically valid in a well-mixed boundary layer. Most vertical profiles of aerosol normalized relative backscatter measured by the lidar at APP during part of the study period and afterward (as part of MPLNET) show a qualitatively exponential decay with height and an absence of aerosol layers aloft (unpublished result). In addition, AOD is highly correlated with surface-level aerosol extinction coefficient at APP (r=0.79; Sherman et al., 2016b). These suggest that optical properties may be well represented by measurements made at the surface.

3.2   Spectral surface reflectance

The MODIS spectral surface reflectance product (Justice, et al., 2002) is derived from MODIS bands 1 thru 7. These seven bands (B1 thru B7) are centered near 645, 855, 466, 553, 1243, 1628, and 2113 nm, respectively. We use the MODIS Aqua eight-day surface reflectance product (MYD09A1), downloaded from the Oak Ridge National Laboratory Distributed Active Archive Center (ORNL DAAC), for the DRE studies in this paper. The MYD09A1 product is created by analyzing MODIS spectral observations over eight-day periods and identifying the invariant contributions (i.e., the surface). These products are gridded, reported at 500 m spatial resolution, and have their own quality assurance and error characteristics. Each MYD09A1 pixel contains the best possible observation (with atmospheric correction applied) during an eight-day period as selected by high observation coverage, low view angle, absence of clouds and cloud shadow, and low aerosol loading (Vermote and Kotchenova, 2008). For each eight-day MYD09A1 product, we calculate the mean surface reflectance of all 500 m pixels in a 10 km x 10 km box (corresponding to 20 x 20 pixels) centered at the APP site, for each of the seven MODIS bands. Only eight-day surface reflectance products with at least 50% of pixels in the 10 km x 10 km box passing MODIS quality assurance tests are used in this study. Because surface reflectance varies primarily on seasonal timescales, we form monthly averages at each wavelength from the eight-day products. Uncertainty in the MODIS surface reflectance product under low aerosol loading conditions is 5 % or $5.0*10^{-4}$, whichever is larger (Vermote and Saleous, 2006). Surface reflectance at APP is always large enough so that $\Delta R=0.05*R$. Since R is wavelength dependent for any surface type, the uncertainty ΔR

depends on wavelength. To set an upper bound on ΔR, we note that monthly averaged R at APP is highest in summer (Fig.3b) and for the 855 nm band (B1) and the 1243 nm band (B5), with summer values near 0.40. We use this to estimate ΔR~0.02 for this study. For simplicity, we neglect the wavelength dependence of ΔR, which if considered would result in a smaller ΔR (Figs.3 a-d).

5   The SBDART radiative transfer model used in this study to calculate DRE parameterizes spectral surface reflectance as a linear combination of that due to vegetation, sand, water, and snow. For each month, we calculate spectral surface reflectance in SBDART for a range of relative vegetation, sand, water, and snow contributions and select the combination which minimizes the mean-square difference (weighted by relative solar irradiance) with the monthly averaged MODIS spectral surface reflectance.

10   Note that the relative contributions to R from the different surface types do not need to add up to 1.0. As an example, the mean spectral surface reflectance measured by MODIS above the APP site in June most closely approximates that of 0.75 times the spectral reflectance curve produced by vegetation alone, with no contributions from the sand, water, or snow spectral reflectance curves (Fig.3b). This is due to the darker vegetation from heavy deciduous forest in the region surrounding the site.

4   Methodology

4.1   SBDART radiative transfer model used to calculate DRE

To calculate clear sky, broadband (250 to 4000 nm) aerosol DRE, we run the SBDART model at 5 nm spectral resolution, with clouds and stratospheric aerosols turned off. We configure the model to use four radiation streams (i.e. four zenith and four azimuthal angles), which provides a good combination of

20   computational efficiency and accuracy for calculating fluxes (Ricchiazzi et al. 1998). We use the LowTran-7 atmospheric transmission model, which possesses 20 cm^{-1} resolution. Standard midlatitude summer vertical profiles of pressure, temperature, water vapor density, and ozone density (McClatchey et al., 1972) are used for April thru October and standard mid-latitude winter profiles are used for November thru March. Although some differences from the actual vertical meteorological profiles are

25   expected, the same standard vertical meteorological profiles are used to calculate the radiative fluxes with and without aerosols and thus would not be expected to contribute much to the calculation of DRE, which

is based on the difference in fluxes (Eq. (5)). The aerosol scattering asymmetry parameter (Sect. 3.1.2) supplied to SBDART is used to estimate aerosol scattering phase function, in the Henyey-Greenstein approximation (Henyey and Greenstein, 1941). Vertical distribution of aerosols is believed to be a second-order effect in the calculation of aerosol DRE for primarily scattering aerosols (McComiskey et al., 2008) and we use the SBDART default vertical aerosol density profile in this initial study. The default profile uses an assumed exponential decrease in aerosol density with a scale height inversely proportional to surface-level aerosol light extinction coefficient (Ricchiazzi et al., 1998). The overall curve is scaled by the AOD (Sect.3.1.1). Aerosol density scale heights used by SBDART range from 1.05 to 1.51 km, which qualitatively agree with typical MPL-measured normalized relative backscatter profiles under clear sky conditions at APP (Sect. 2)

4.2   Seasonal variability in aerosol optical properties and DRE

For the study of seasonal DRE variability (Sect. 5.1), we use the SBDART model to calculate diurnally averaged DRE at the TOA and at the surface, for 418 days during the period 14 June 2012 thru 28 June 2016. We then bin the DRE by month (Figs. 4a and 4b). For each of the 418 days, we calculate DRE for each hour to account for the effect of varying solar geometry on aerosol radiative effects, supplying daily-averaged $AOD(\lambda)$, $\omega_0(\lambda)$, and $g(\lambda)$, along with monthly averaged spectral surface reflectance (R), as inputs to the SBDART model. Upwelling and downwelling broadband fluxes for that hour are calculated with average measured aerosol properties and then with no aerosols and their difference is used to calculate DRE using Eq. (5)

$$DRE = (F_{A,\downarrow} - F_{A,\uparrow}) - (F_{NA,\downarrow} - F_{NA,\uparrow}) \tag{5}$$

The process is repeated for all 24 hours and the results averaged to yield diurnally averaged DRE. In addition to DRE, we calculate the direct radiative efficiency (RE) by dividing diurnally averaged DRE by daily averaged AOD at 550 nm. Radiative efficiency is to first-order independent of aerosol amount (i.e. AOD), and dependent on the inherent optical nature of the aerosol, controlled by composition and size. It is a useful quantity for determining whether DRE varies due to changes in aerosol loading or aerosol type. Use of daily averaged DRE in this study integrates over solar angles and the use of daily averaged aerosol optical properties is justified by the small diurnal variability in AOD, ω_0, and g at APP

(Sherman et al., 2015; Sherman et al., 2016a). Monthly DRE and RE statistics are built up using all daily averaged values for that particular month.

4.3   Sensitivity of aerosol DRE to aerosol properties and surface reflectance

To study the sensitivity of clear sky aerosol DRE to changes in AOD, ω_0, g, and R (Sect. 5.2), we follow a similar approach to that used by McComiskey et al (2008). We define the sensitivity S_i of diurnally averaged DRE to parameter ρ_i (where ρ_i stands for either AOD, ω_0, g, or R) as the change in DRE per unit change in ρ_i. Formally, S_i is evaluated as the partial derivative of DRE with respect to ρ_i ($S_i=\partial$ (DRE) / $\partial\rho_i$) evaluated at the 'base case' values for all variables (Table 1). To assess whether these sensitivities S_i are independent of the respective ρ_i values, we plot diurnally averaged DRE versus ρ_i over the largest expected range of each ρ_i at APP (Table 2). For the sensitivity S_R, we scale the entire spectral surface reflectance curve (Figs.3) by values ranging from 0.5 to 2.0, with a value of 1.0 corresponding to the 'base case' curve (i.e. the mean spectral reflectance curve). For example, a value of 2.0 would correspond to twice the mean surface reflectivity (at all wavelengths). Insensitivity of DRE to AOD, ω_0, g, or R is inferred from the degree of linearity of the respective plot (i.e. a constant slope S_i).

Since aerosol optical properties (Figs. 5) and surface reflectance (Figs.3) at APP vary primarily on seasonal scales, we evaluate the S_i separately for each season. It is impractical (and unnecessary) to construct sensitivity curves for each individual month so we choose one representative calendar day to represent each season; December 21 for winter, March 21 for spring, June 21 for summer, and September 21 for fall. We refer to these representative days as DEC, MAR, JUN, and SEP. Inclusion of the equinox days (with equal durations of sunlight and darkness) also provides results which may possibly be indicative of 'annually averaged values'. We use monthly median AOD, ω_0, and g values at 550 nm as base case values but monthly mean values are similar to medians at APP and could also be used with negligible difference in results (Figs.5 a-c). Spectral dependence of each aerosol property is calculated from the values at 550 nm, using the approach of Sect(s). 3.1. The base case R is the monthly mean spectral surface reflectance curve (Figs.3). We then vary the independent variables ρ_i individually about these base case values (Table 1) to generate the 'seasonal' DRE versus ρ_i curves. We evaluate $S_i=\partial$ (DRE) / $\partial\rho_i$, at base case ρ_i value, as the regression slope of the five points on each side of the base case value.

4.4 Estimating uncertainty in aerosol DRE

The DRE sensitivity values $S_i = \partial RE/\partial p_i$ and known (or estimated) measurement uncertainties, ΔAOD, $\Delta\omega_0$, Δg, and ΔR (Table 2) can be used to calculate the corresponding uncertainty in DRE using Eq. (6)

$$\Delta DRE = \sum_{j=1}^{4}\sum_{i=1}^{4} \frac{\partial DRE}{\partial\rho_i}\frac{\partial DRE}{\partial\rho_j}\mathrm{cov}(\rho_i,\rho_j) = \sum_{j=1}^{4}\sum_{i=1}^{4}\frac{\partial DRE}{\partial\rho_i}\frac{\partial DRE}{\partial\rho_j}corr(\rho_i,\rho_j)\Delta\rho_i\Delta\rho_j \qquad (6)$$

Where $\mathrm{cov}(\rho_i,\rho_j)$ is the covariance of ρ_i, and ρ_j, which in turn can be expressed in terms of their linear correlation $corr(\rho_i,\rho_j)$. The summations in Eq. (6) can be explicitly written as

$$\Delta DRE^2 = S_{AOD}^2\Delta AOD^2 + S_{\omega_0}^2\Delta\omega_0^2 + S_g^2\Delta g^2 + S_R^2\Delta R^2 + 2S_{AOD}S_{\omega_0}corr(AOD,\omega_0)\Delta AOD\cdot\Delta\omega_0$$
$$+ 2S_{AOD}S_g corr(AOD,g)\Delta AOD\cdot\Delta g + 2S_{AOD}S_R corr(AOD,R)\Delta AOD\cdot\Delta R + 2S_{\omega_0}S_g corr(\omega_0,g)\Delta\omega_0\cdot\Delta g +$$
$$2S_{\omega_0}S_R corr(\omega_0,R)\Delta\omega_0\cdot\Delta R + 2S_g S_R corr(g,R)\Delta g\cdot\Delta R$$

$$(7)$$

The first four terms of the sum in Eq. (7) facilitate estimates of the contributions to ΔDRE due to the individual sources of uncertainty, neglecting covariance effects (McComiskey et al., 2008). However, correlations amongst some aerosol optical properties are non-negligible at APP during some seasons (Table 3) and must be considered for an improved estimate of ΔDRE. We calculate ΔDRE both with and without the inclusion of the covariance terms to examine their effect on ΔDRE. An equation identical to Eq. (7) can be used to calculate the uncertainty in RE. Results from a similar analysis of the sensitivity of RE to AOD, ω_0, g, and R, along with the associated uncertainties in RE, are provided in Sect(s). S1 and S2 of the Supplement to this paper.

5 Results and discussion

5.1 Annual cycles of aerosol DRE, RE, and aerosol optical properties

Aerosol DRE and the optical properties influencing DRE demonstrate large seasonal variability at APP. Median diurnally averaged DRE for the nearly four-year study period is -2.9 Wm^{-2} at the TOA and -6.1 Wm^{-2} at the surface (Figs. 4a and 4b). Median DRE at the TOA is ~5 to 6 times larger (i.e. more negative) during summer months (JJA) than during winter months (DJF) and median DRE at the surface is ~4 times larger in summer than in winter. Median DRE at both the TOA and surface is nearly twice as

negative during summer months as the median for the entire period. Variability in DRE is also largest during summer and the surrounding months (Figs. 4a and 4b). While fewer daylight hours during winter obviously contributes to the seasonal differences in diurnally averaged DRE, the annual DRE cycles clearly follow that of AOD (Fig.5a). Median and mean AOD at 550 nm are approximately 9 to 10 times

5   larger in summer than in winter and AOD variability is also largest during summer and the surrounding months.

Further confirmation of the dominant influence of AOD on the annual DRE cycles at APP is seen in the annual RE cycles. Aerosol RE is most sensitive to ω_0 and then g, with a much smaller dependence on AOD and R (Sect.S1 of Supplement to this paper). Based on this, the annual DRE and RE cycles should

10   be qualitatively similar if ω_0 and/or g (rather than AOD) exert the primary influences on the DRE cycle. Instead, median and mean RE are more negative during winter months than during summer months by a factor of ~ 1.5 at the TOA and ~1.5 to 2 at the surface (Figs. 4c and 4d) and RE variability is also largest during winter months. Thus, months with the most negative DRE coincide with the least negative RE. One exception is April, for which median and mean AOD are only half of that during the summer months

15   but surface DRE (and to a lesser extent TOA DRE) is close to that during summer months. April is characterized by a relatively low ω_0 (Fig.5b) and low g (Fig.5c), in addition to a darker surface during spring than summer (Figs.3a and 3b). This leads to a similar RE to that of winter months, coinciding with a high enough AOD to produce surface DRE close to that during summer.

The annual RE cycles at both the TOA and the surface (Figs. 4c and 4d) can be explained using the

20   following information: (1) the sensitivities of RE to increases in both ω_0 and g are the same sign (positive) at the surface but have opposite signs at the TOA (positive for g and negative for ω_0; Table S1 of Supplement); and (2) ω_0 and g are highly correlated for all seasons ($0.78 \leq r \leq 0.85$; Table 3), with both properties exhibiting larger values during summer and the surrounding months and smaller values during most other months (Figs. 5b and 5c). The fact that the RE sensitivities to ω_0 and g are both positive at the

25   surface (i.e. an increase in ω_0 or g drives RE at the surface to less negative values) results in less negative surface RE values during summer (when both ω_0 and g are largest) and more negative RE values during winter and the surrounding months (when both ω_0 and g are smallest). The annual RE cycle at the TOA appears to be influenced more by g than by ω_0, since the decrease in g as one moves away from the

summer months drives the TOA RE more negative than the positive influence on RE due to lower ω_0 (Fig.4c). This is despite greater sensitivity of RE to ω_0 and is likely due to the larger (by a factor of ~2) summer-winter differences in g than in ω_0 (Figs. 5b and 5c).

Alston and Sokolik (2016) reported a mean TOA DRE of approximately -10 Wm^{-2} for mountainous western North Carolina and surrounding areas during summer (their Fig.7), with somewhat more negative values (-12 to -16 Wm^{-2}) for the SE U.S. as a whole. Their DRE values appear to be calculated based on cloud fraction of ~40 % (their Fig.2) while ours are for clear sky conditions (i.e. cloud fraction of zero) so we need to multiply their DRE values by a factor of 1 / (1-0.40) ~1.7 to compare with our clear sky DRE. Our monthly mean clear sky TOA DRE of -5 to -6 Wm^{-2} during summer months (Fig.4a) is approximately three times smaller than Alston and Sokolik's clear sky values (-17 Wm^{-2}). These large DRE differences cannot be explained solely by differences between the two studies in AOD and ω_0, the two parameters to which clear sky TOA DRE is most sensitive (Sect. 5.2). Our summer AOD at 550 nm is ~0.15 to 0.20 (Fig.5a) while theirs was on average ~0.25 to 0.30 (their Fig.1). Using our derived summer DRE sensitivity to AOD at the TOA (Table 3; Sect. 5.2), this difference in AOD gives rise to a difference in TOA DRE of only ~4 Wm^{-2}. Alston and Sokolik (2016) used a similar value of ω_0 (0.95 at 558 nm; their Fig.11) to our summer monthly median values (Fig.5b).

A more likely source for much of the large summer DRE discrepancy is a difference in models used to calculate TOA DRE. Inserting representative summer median values of aerosol optical properties at APP (Table 1) into the first-order DRE equation used by Alston and Sokolik (Eq. (2) of Haywood and Shine, 1995) leads to summer TOA DRE values of ~ -14 Wm^{-2} at APP, which is 2.5 times more negative than monthly mean (and median) TOA DRE calculated using the SBDART radiative transfer model (Fig.4a). The Haywood and Shine (1995) equation used by Alston and Sokolik (2016) is valid for an optically thin atmosphere and uses spectrally weighted aerosol optical properties and surface reflectance as inputs. The degree to which the first assumption holds obviously decreases with the higher AOD of summer months while the degree to which aerosol optical properties at 550 nm are representative of spectrally weighted properties also impacts the resultant TOA DRE. Using an identical procedure to that outlined above for summer, our winter DRE values using the simplified Haywood and Shine (1995) TOA DRE equation are approximately the same as those calculated using the SBDART model. Unlike summer,

a majority of the (~2 to 3 Wm^{-2}) difference between our monthly mean TOA DRF and that reported by Alston and Sokolik (2016) is consistent with differences in AOD and cloud fraction between the studies. Sensitivity of TOA DRF to ω_0, g, and R is so small during winter (Table 3 and Sect. 5.2) that differences in these properties is unlikely to influence TOA DRF agreement.

5.2 Sensitivity of DRE to aerosol optical properties and surface reflectance

The plots of aerosol DRE versus AOD, ω_0, g, and R (Figs. 6 and 7) are for the most part linear, indicating that the sensitivities (i.e. slopes of plots) are independent of the values of aerosol optical properties and surface reflectance, at least over the ranges observed at the APP site. There are a few minor exceptions, namely (1) sensitivity of TOA and surface DRE to AOD for high AOD values in DEC and MAR; and (2) TOA DRE sensitivity to ω_0 in JUN and SEP. The nonlinearity in DRE versus AOD leads to a small dependence of RE on AOD, more so in DEC and MAR (See also Table S1 of Supplement). However, AOD values during the non-summer months at APP are rarely large enough (Fig.5a) to lie on the nonlinear portion of the curves (Figs. 6a and 7a).

The sensitivity S_{AOD} is greatest for MAR (-47 Wm^{-2} AOD^{-1} at TOA and -90 Wm^{-2} AOD^{-1} at surface) but S_{AOD} at the TOA and surface exhibit only modest variation with season (~25 to 30 %). Values of S_{AOD} at the surface are close to twice those at the TOA for all seasons. The sensitivities $S_{\omega o}$, S_g, and S_R vary much more with season than does S_{AOD} (Table 3), with values 10 to 20 times greater in JUN than in DEC. Higher sensitivity of DRE to ω_0, g, and R for higher AOD conditions was also reported by McComiskey et al. (2008), as part of their comparisons of DRE sensitivities at three U.S. Department of Energy-supported Atmospheric Radiation Measurement (DOE ARM) sites. After AOD, aerosol DRE is most sensitive to ω_0 and least sensitive to R, with JUN $S_{\omega o}$ of comparable magnitude to S_{AOD} at the TOA. The fact that $S_{\omega o}$ is negative at the TOA and positive at the surface (Table 3; Figs.6b and 7b) implies that increasing ω_0 leads to a larger cooling effect at TOA and a smaller cooling effect at the surface. The magnitude of $S_{\omega o}$ is ~40 to 60 % greater at the surface than the TOA for all seasons. In contrast, the magnitudes of S_g are nearly identical at TOA and surface for all seasons. The fact that the signs of S_g are positive at TOA and surface is consistent with the fact that larger particles (corresponding to larger g) scatter a greater fraction of light in the forward direction than do smaller particles (smaller g). The

relatively low sensitivity of DRE to surface reflectance at APP during the study period is seen by the low S_R, which range from 3.3 Wm^{-2} (2.7 Wm^{-2}) at the TOA (surface) during JUN to 0.22 (0.20) Wm^{-2} at TOA (surface) during DEC, per unit change in R. This implies that a doubling of R is necessary in to reduce DRE (i.e. make less negative) by 3.3 Wm^{-2} (2.7 Wm^{-2}) at the TOA (surface) during JUN. Changes in R

5 within a given month do not even approach this magnitude at APP, except for extended snow cover during some winter months.

The only similar study of DRE sensitivity and uncertainties to include a continental U.S. site was that of McComiskey et al. (2008), which included the Southern Great Plains DOE ARM site (SGP; Lamont, OK) and the month of SEP. Alston and Sokolik (2016) investigated the effects of changes in AOD, R,

10 and cloud fraction on TOA DRE and concluded that DRE changed most in response to changes in AOD. However, they did not derive numerical values for the sensitivities. McComiskey et al. (2008) used similar aerosol optical properties at 550 nm to our SEP values (Table 1) in their sensitivity study (AOD=0.10; ω_0=0.95; g=0.60). Their spectral surface reflectance for the SGP site was close to twice that of our SEP curve (Fig.3c) because the SGP site is surrounded by wheat, corn, and hay fields (Sherman et al., 2015)

15 rather than forests. McComiskey et al. (2008) also used a scattering Angstrom exponent of 1.0 (for wavelength-scaling ω_0) to reflect size distributions at SGP with a greater influence from super-micron particles (due to local turbidity caused by dust) than that observed at APP (Sherman et al., 2015). The S_{AOD} values reported by McComiskey et al. (2008) were similar (S_{AOD} ~37 Wm^{-2} at TOA and ~67 Wm^{-2} at surface) to our SEP values (Table 3). Their $S_{\omega o}$ and S_g at the TOA and surface were also within ~20 %

20 of our reported SEP values but their S_R were ~4 to 5 time higher than our values. The difference could lie in the methods used. McComiskey used a spectrally flat UV/VIS surface reflectance of 0.10 and a near-IR reflectance of 0.40 for SGP, while the MODIS-measured spectral surface reflectance at APP clearly demonstrates a wavelength dependence in the near-IR during summer and surrounding months (Figs. 3b and 3c), due to the more vegetated landscape near APP than near SGP. We find the more precise treatment

25 of the spectral dependence of R used here and based on input from MODIS measured surface reflectance an improvement in methodology over the McComiskey et al. (2008) study. While good overall agreement in SEP DRE sensitivities is to be expected, based on similar aerosol optical properties input to SBDART,

the results of our study indicate that the seasonal dependence of $S_{\omega o}$, S_g (and to lesser extent, S_R) must be accounted for in DRE uncertainty estimations, at least at APP.

5.3   DRE uncertainties

Uncertainty in aerosol DRE at both the TOA and surface is largest in JUN and lowest in DEC (Tables 4 and 5), primarily due to the highest base case AOD in summer (JUN) and lowest in winter (DEC) (Table 1). Uncertainties in diurnally averaged DRE at the TOA (surface) are 15 to 20 % (12 to 20 %) of the DRE calculated using base case values for the given month (Table 1), except for DEC, when the DRE uncertainty is close to 50 % at both the TOA and surface. Since the measurement uncertainties of AOD, ω_0, g, and R are all between 0.01 and 0.02 (Table 2), the relative contributions of each parameter to the total DRE uncertainty are approximately proportional to their DRE sensitivity values (Table 3). We note that this is specific to the measurements made at APP (Sect(s). 3) and may not necessarily be true at sites where different measurement protocols are employed.

Uncertainty in DRE at both the TOA and the surface is dominated by AOD uncertainty for the months with lowest base case AOD (DEC and MAR) and covariance effects amongst aerosol optical properties are negligible during these months, so that the cross terms in Eq. (7) do not need to be considered (Tables 4 and 5). In contrast, uncertainties in AOD and ω_0 contribute similarly to DRE uncertainty in JUN and SEP, due to the higher $S_{\omega o}$ during these months. Uncertainties in g and R contribute less to the DRE uncertainty during all seasons. Inclusion of covariance effects in the DRE uncertainty calculations increases the JUN and SEP DRE uncertainty at the TOA (surface) by nearly 0.2 W m^{-2} (~0.25 to 0.3 W m^{-2}), due to the higher correlations amongst AOD, ω_0, and g during the warm-season months at APP (Table 3). It is interesting to note that ω_0 and g at are highly correlated during all seasons, with corr(ω_0,g) between 0.78 and 0.85 (Table 3). This indicates a strong tendency for larger particles at APP to be more reflective. However, the sensitivities $S_{\omega o}$ and S_g are only large enough to lead to non-negligible covariance effects during the higher aerosol loading months (JUN and SEP) at APP (Eq. (7), Tables 4 and 5).

Due to the wide usage of MODIS-measured AOD in aerosol DRE studies, it is instructive to compare DRE uncertainties calculated using AERONET AOD with those using MODIS AOD. In their global inter-comparison of MODIS Collection 5 AOD with AERONET, Levy et al. (2010) estimated the MODIS

AOD error envelope to be ±(0.05+0.15*AOD$_{aeronet}$). For our comparison of DRE uncertainties, we use the lower MODIS AOD uncertainty ΔAOD$_{MODIS}$=0.05 in Eq. (7), in place of the AERONET uncertainty of 0.01 (Table 2). Sherman et al (2016a) reported excellent correlation of MODIS AOD and daily averaged AERONET AOD above APP so we use the same correlation values to calculate the covariance terms involving AOD in Eq. (7) (Table 3). Uncertainty in TOA DRE using MODIS AOD is close to 2 Wm^{-2} for each season (Table 4), which corresponds to 36 % of the base case JUN DRE and 235 % of the DEC base case DRE (Table 4). In terms of absolute DRE uncertainties, those based on MODIS AOD are 2.5 to 5 times the DRE uncertainties using AERONET AOD. Similar fractional DRE uncertainties (35 to 240 %) as that at the TOA are obtained at the surface (Table 5), which correspond to DRE uncertainties ranging from 3.4 Wm^{-2} (SEP) to 4.5 Wm^{-2} (MAR). Due to the higher uncertainty in MODIS AOD (relative to AERONET), the AOD uncertainty is the dominant term in DRE uncertainty for all seasons when MODIS AOD is used. It should be noted that an AOD error envelope for the new MODIS collection (C6) is not yet available but it could be smaller than that of Collection 5, given algorithm improvements (Levy et al., 2013).

6   Summary and conclusions

Daily-averaged aerosol optical properties measured at Appalachian State University's co-located NASA AERONET and NOAA ESRL aerosol monitoring sites over a nearly four-year period are used along with monthly averaged spectral surface reflectance measured by MODIS to study the annual cycles of diurnally averaged clear sky aerosol DRE and RE. This study is the first multi-year 'ground truth' DRE study in the SE U.S., using aerosol network data products that are often used to validate satellite-based aerosol retrievals (Levy et al., 2010; Sherman et al., 2016a). The study is also the first in the SE U.S. to quantify DRE uncertainties and sensitivities to aerosol optical properties and surface reflectance, including their seasonal dependence.

Median diurnally averaged clear sky DRE at APP over the study period is -2.9 Wm^{-2} at the TOA and -6.1 Wm^{-2} at the surface. Monthly median and mean DRE at the TOA (surface) range from -1 to -2 Wm^{-2} (-2 to -3 Wm^{-2}) during winter months to -5 to -6 Wm^{-2} (negative 10 Wm^{-2}) during summer months. While the annual DRE cycle at APP largely follows that of AOD, aerosol RE demonstrates an anti-

correlation with AOD. The most negative RE is observed during November thru April at the TOA and during October thru April at the surface. The least negative RE is observed in June thru September at both the TOA and the surface. Aerosol DRE is most sensitive to AOD, followed (in order) by ω_o, g, and R. One exception is that the sensitivity of TOA DRE with respect to ω_o is comparable to that of AOD during summer (JUN), when the base case AOD at APP is highest. Since the uncertainties in AOD, ω_o, g, and R are of comparable magnitude, the relative contributions of each to the total DRE uncertainty are proportional to the sensitivities. The sensitivities S_{ω_o}, S_g, and S_R vary by factors of ~10 to 20 with season, in contrast to the much more modest seasonal variation (~25 to 30 %) of S_{AOD}. This result supports the assertion that the seasonal dependence of S_{ω_o}, S_g (and to lesser extent, S_R) must be accounted for in DRE uncertainty estimates, at least for sites like APP where there is large seasonality in aerosol loading and in dominant aerosol types.

Using seasonally representative aerosol optical properties and surface reflectance from APP, clear sky aerosol DRE uncertainty at the TOA (surface) ranges from 0.44 Wm^{-2} (0.73 Wm^{-2}) for DEC to 0.90 Wm^{-2} (1.3 Wm^{-2}) for JUN. Expressed as a fraction of DRE calculated using base case aerosol optical properties and surface reflectance (Table 1), the DRE uncertainties at the TOA (surface) are 16 to 20 % (12 to 20 %) for MAR, JUN, and SEP and 48 % (49 %) for DEC. Unlike the McComiskey et al.(2008) study, we include the effect of covariances amongst aerosol optical properties in order to determine their effect on DRE uncertainty. Their impact on DRE uncertainty at APP is negligible for low AOD conditions (DEC and MAR) but does increase ΔDRE by ~0.2 to 0.3 Wm^{-2} under moderate AOD conditions (JUN and SEP). The primary reason for the relatively low DRE uncertainties reported in this study is the small uncertainty in AOD measured by Cimel sunphotometers as part of AERONET (Eck et al., 1999). The DRE uncertainties are dominated by AOD uncertainties for all seasons when MODIS AOD is used as the AOD input to radiative transfer calculations, with uncertainties 2.5 to 5 times larger than those using AERONET AOD.

The results from our study suggest that while satellite-based aerosol measurements provide the necessary global coverage for climate studies, their current levels of uncertainty necessitate complementary ground truth measurements of AOD, ω_0, and g (or some other proxy for scattering phase function) from regionally representative sites to better constrain aerosol DRE in models. Continuous,

long-term aerosol measurements from ground-based aerosol network sites are necessary not only to evaluate satellite-based aerosol retrieval accuracy but also to assess whether AOD at the time of satellite overpass is representative of daily averaged AOD at the site (Sherman et al., 2016a). Our study also illustrates the challenges faced in such 'ground truth' DRE estimates. Aerosol optical depth over much of

5   the non-urban U.S. is too low to retrieve column-averaged ω_0 with an uncertainty less than ~0.03 (at best) at AERONET sites, under most conditions (Dubovik, et al., 2000). Ground-based in situ aerosol networks such as NOAA ESRL can measure ω_0 to ~0.015 (Sherman et al., 2015) even under these low-loading conditions but the dried aerosol properties must be corrected to ambient RH, which requires knowledge of the hygroscopic dependence of aerosol light scattering. Currently such measurements exist at only two

10   NOAA ESRL sites in the U.S. (APP and SGP). In addition, some knowledge of vertical aerosol profiles is necessary to assess whether the near-surface aerosol properties are likely representative of the column-averaged aerosols. Though qualitative inspection of lidar-measured aerosol backscatter profiles at APP indicates that aerosols are generally confined to the first 1 to 2 km, the availability of multi-year, quality-assured profiles of aerosol light extinction (as part of MPLNET) will facilitate expansion of this initial

15   DRE study to include the effects of vertical aerosol structure on DRE at APP.

Data availability: AERONET Level 2 aerosol optical depth data used in this paper is publicly available at the NASA AERONET website (https://aeronet.gsfc.nasa.gov/cgi-bin/webtool_opera_v2_new). The aerosol measurements made as part of NOAA ESRL are publicly available for download at the Global

20   Atmospheric Watch World Data Centre for Aerosols (http://ebas.nilu.no/Default.aspx) and graphical representations of the data are available at the NOAA ESRL website (https://www.esrl.noaa.gov/gmd/aero/net/status_plot.php?sta=app)

Competing interests: The authors declare that they have no conflict of interest.

Appendix A   Table of acronyms and symbols

Table A1. Frequently used acronyms and symbols in this paper

Acronym or symbol	Name or definition
AERONET	NASA Aerosol Robotic Network
AOD	Aerosol optical depth; the vertical integral of aerosol light extinction coefficient
APP	Aerosol monitoring sites at Appalachian State University
DRE	Aerosol direct radiative effect. In this paper, we only consider diurnally averaged DRE
ESRL	National Oceanic and Atmospheric Administration (NOAA) Earth System Research Laboratory
g	Scattering asymmetry parameter
MODIS	Moderate Resolution Imaging Spectrometer, aboard the polar-orbiting NASA Terra and Aqua satellites
MPLNET	NASA Micro-pulsed Lidar Network
R	Surface reflectance
RE	Aerosol radiative efficiency; equal to DRE divided by AOD
RH	Relative humidity
SBDART	Santa Barbara DISORT Radiative Transfer model used to calculate DRE
SE U.S.	Southeastern United States
S_{AOD}	Change in DRE per unit change in AOD[a]
S_g	Change in DRE per unit change in g[a]
S_{ω_0}	Change in DRE per unit change in ω_0[a]
S_R	Change in DRE per unit change in R[a]
TOA	Top-of-atmosphere. For this paper, TOA refers to 100 km above sea level
λ	Wavelength of solar radiation (units: nm)
ω_0	Single-scattering albedo; the fraction of aerosol light extinction due to scattering

[a] Evaluated at base case aerosol properties and surface reflectance (Table 1)

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Table 1. Base case values of aerosol optical properties and vegetation surface cover (SC) coefficients used by SBDART to calculate diurnally averaged DRE for the months DEC, MAR, JUN, and SEP. The vegetation SC coefficients produce the spectral reflectance curves R of Fig.3 and are the linear combinations SC = (snow, water, vegetation, sand) that best match monthly averaged spectral reflectance measured by MODIS. The base case aerosol optical properties are the monthly median values (Fig.5).

Property p_i	MAR	JUN	SEP	DEC
AOD (550 nm)	0.05	0.15	0.10	0.02
ω_0 (550 nm)	0.91	0.95	0.95	0.91
g (550 nm)	0.57	0.62	0.67	0.61
$\mathring{A}_{sp}$	2.1	2.0	2.2	2.1
$\mathring{A}_{ap}$	1.2	0.6	0.8	1.2
$\mathring{A}_{AOD}$	1.3	1.7	1.8	1.4
Vegetation SC coefficients	(0,0,0.30,0.20)	(0,0,0,0.75)	(0,0,0.05,0.55)	(0,0,0.40,0.15)

Table 2. Range over which aerosol optical properties and relative spectral surface reflectance used to calculate DRE are varied, along with the measurement uncertainties used in the DRE uncertainty calculations. Relative surface reflectance of 1.0 corresponds to the mean spectral reflectance curve R used for a given month (Fig.3) and the entire curve is scaled over the range of scale factors shown in the table.

Parameter	Range varied in sensitivity curves	Measurement Uncertainty (Source)	
AOD at 550nm	0.0-0.3	0.01	(Eck et al., 1999)
ω_0 at 550nm	0.75-1.0	0.015[a]	(Sherman et al., 2015)
g at 550nm	0.50-0.75	0.01[a]	(Sherman et al., 2015)
Relative surface reflectance	0.5-2.0	0.05*R[b]	(Vermote and Saleous, 2006)

5 [a] Value is for aerosols at low relative humidity (RH<40%);

[b] For uncertainty in R, we use a wavelength-independent ΔR~0.020, which corresponds to reflectance R of 40% and thus represents an upper bound for APP (Fig.3).

Table 3. Sensitivity of top-of-atmosphere (TOA) and surface DRE to AOD, ω_0, g, and R. Sensitivities S_i are calculated as the slope of DRE versus ρ_i, evaluated at base case values (Table 1). All sensitivities are in units of W m^{-2} per unit change in the parameter ρ_i. The correlations between aerosol optical properties are used along with uncertainties (Table 2) to calculate the covariances used in the DRE uncertainty calculations
5 (Tables 4 and 5).

Property p_i	MAR	JUN	SEP	DEC
TOA S_{AOD}	-47	-35	-34	-43
TOA $S_{\omega 0}$	-9.1	-39	-18	-2.6
TOA S_g	5.9	19	12	1.8
TOA S_R	0.92	3.3	1.4	0.22
Surface S_{AOD}	-90	-69	-61	-72
Surface $S_{\omega 0}$	16	54	30	4.3
Surface S_g	6.1	19	12	1.9
Surface S_R	0.77	2.7	1.4	0.20
Corr(AOD,ω_0)	-0.02	0.25	0.57	0.10
Corr(AOD,g)	-0.08	0.30	0.56	0.15
Corr(ω_0,g)	0.78	0.79	0.85	0.84

Table 4. Calculated uncertainties in DRE at the TOA, using the sensitivities and correlations given in Table 3 and measurement uncertainties given in Table 2 as inputs to Eq. (7). Units of ΔDRE are W m^{-2}. The uncertainties associated with aerosol optical depth are calculated twice; once using AERONET AOD uncertainties and once using the lower bound for MODIS AOD uncertainty (shown in parentheses).

	MAR	JUN	SEP	DEC
ΔDRE$_{AOD}$	0.47 (2.3)	0.35 (1.8)	0.34 (1.7)	0.43 (2.1)
ΔDRE$_{\omega0}$	0.14	0.58	0.27	0.039
ΔDRE$_g$	0.059	0.18	0.12	0.018
ΔDRE$_R$	0.023	0.084	0.034	0.0050
Sum of covariance terms	0.0060 (-0.022)	0.31 (0.87)	0.20 (0.80)	0.0070 (0.030)
ΔDRE (covariance terms not included)	0.49 (2.3)	0.71 (1.9)	0.45 (1.7)	0.43 (2.1)
ΔDRE (covariance terms included)	0.50 (2.3)	0.90 (2.1)	0.64 (2.0)	0.44 (2.1)
DRE[a] (Base case)	-2.4	-5.7	-3.6	-0.91
ΔDRE[a] / DRE (Base Case)	0.20 (0.96)	0.16 (0.36)	0.18 (0.54)	0.48 (2.4)

[a] Uncertainty includes covariance terms

Table 5. Calculated uncertainties in DRE at the surface, using the sensitivities and correlations given in Table 2 and measurement uncertainties given in Table 1 as inputs to Eq. (7). Units of ΔDRE are W m^{-2}. The uncertainties associated with aerosol optical depth are calculated twice; once using AERONET AOD uncertainties and once using the lower bound for MODIS AOD uncertainties (shown in parentheses).

	MAR	JUN	SEP	DEC
ΔDRE_{AOD}	0.90 (4.5)	0.69 (3.5)	0.61 (3.0)	0.79 (3.6)
$\Delta DRE_{\omega 0}$	0.23	0.80	0.45	0.064
ΔDRE_g	0.061	0.19	0.12	0.019
ΔDRE_R	0.019	0.069	0.036	0.0050
Sum of covariance terms	0.0050 (-0.063)	0.60 (2.0)	0.49 (2.1)	0.015 (0.068)
ΔDRE (covariance terms not included)	0.93 (4.5)	1.1 (3.6)	0.77 (3.1)	0.72 (3.6)
ΔDRE (covariance terms included)	0.93 (4.5)	1.3 (3.8)	1.0 (3.4)	0.73 (3.6)
DRE (Base case)	-4.6	-11	-6.3	-1.5
ΔDRE[1] / DRE (Base Case)	0.20 (0.98)	0.12 (0.36)	0.16 (0.54)	0.49 (2.4)

[1]Uncertainty includes covariance terms
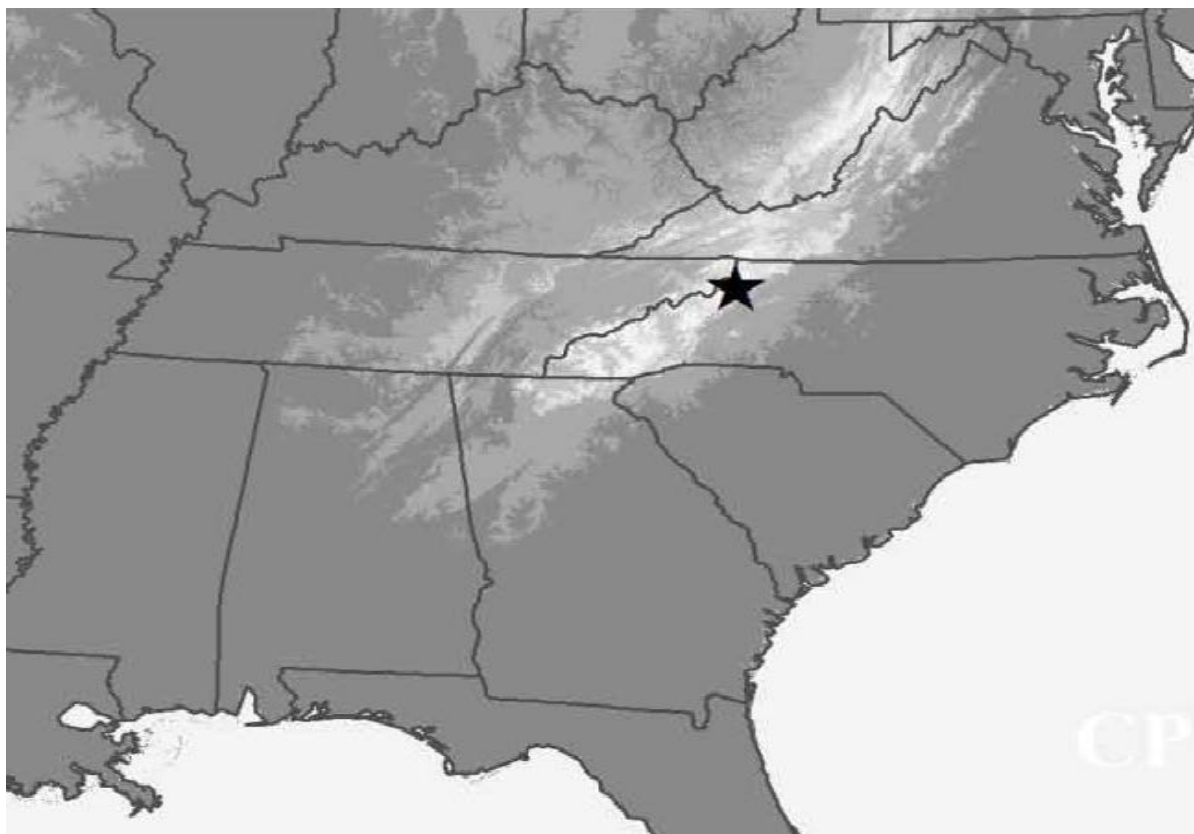
Figure 1. Map of the SE U.S. showing the location of Appalachian State University (APP) in Boone, NC
$(36.21^0$ N, 81.69^0 N, 1080 m above sea level). White shades denote mountain elevations.
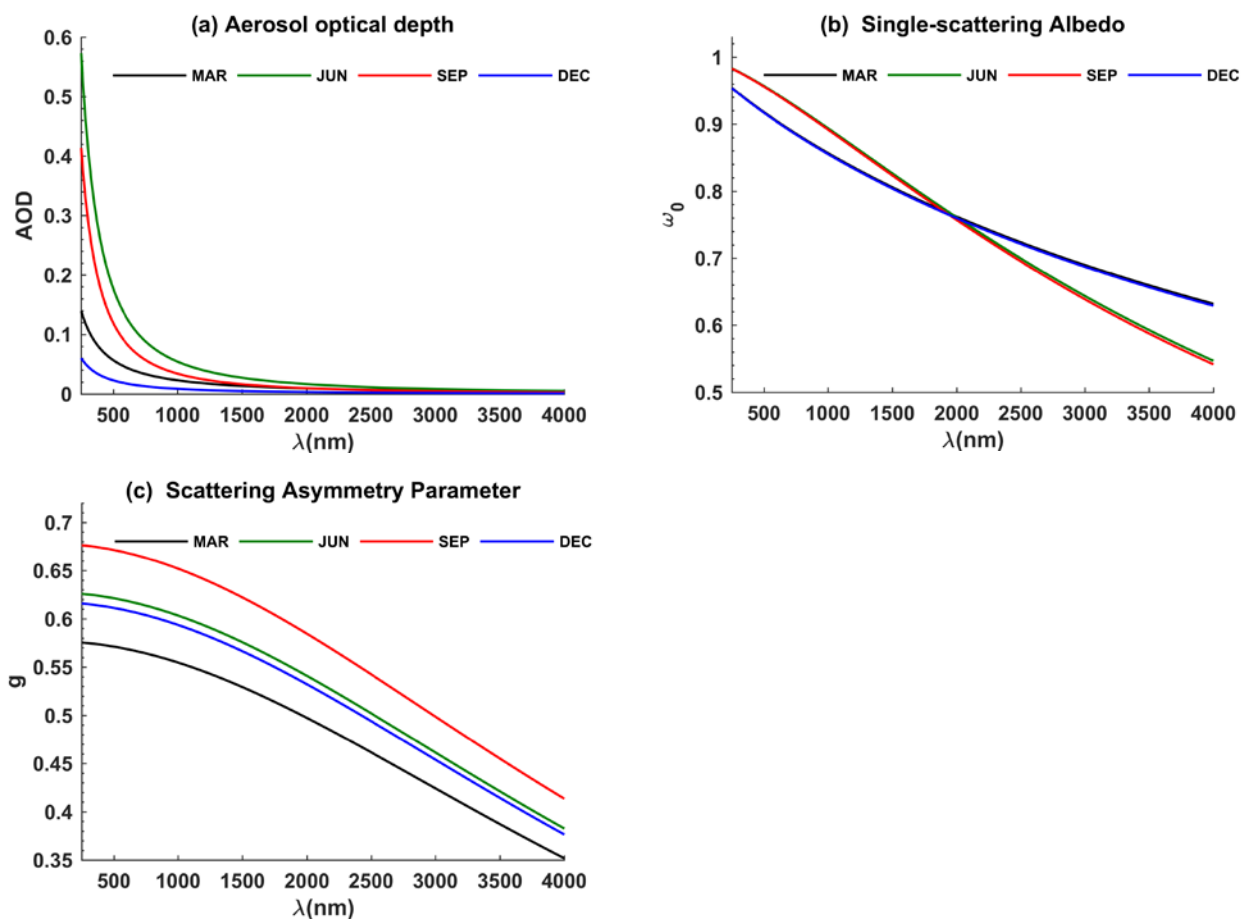
Figure 2. Wavelength dependence of AOD, ω_0, and g, calculated using Eq(s). 1 thru 3 and base case aerosol optical properties for each season (Table 1). The MAR and DEC curves are nearly identical in (b), as are the JUN and SEP curves in (b), due to the use of nearly identical base case aerosol optical properties in generating the curves.
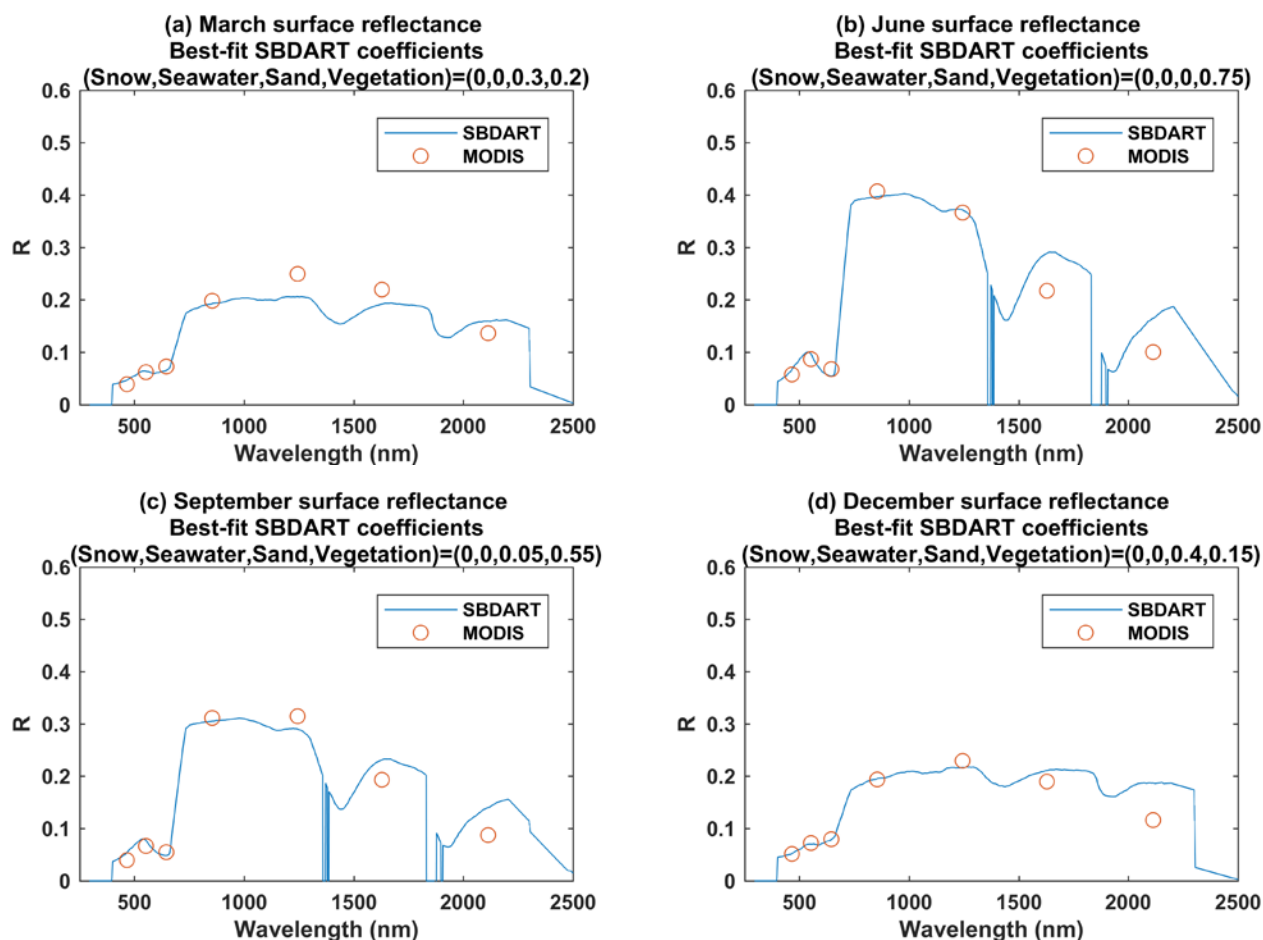
Figure 3. Spectral surface reflectance over the wavelength range 250 to 2500 nm for (a) March; (b) June; (c) September; and (d) December. The plot for each month contains average of MODIS Aqua spectral surface reflectance values for that month, over the study period. The SBDART curve is based on spectral reflectance produced by the vegetation types whose combination provided best fit with MODIS-measured spectral reflectance.
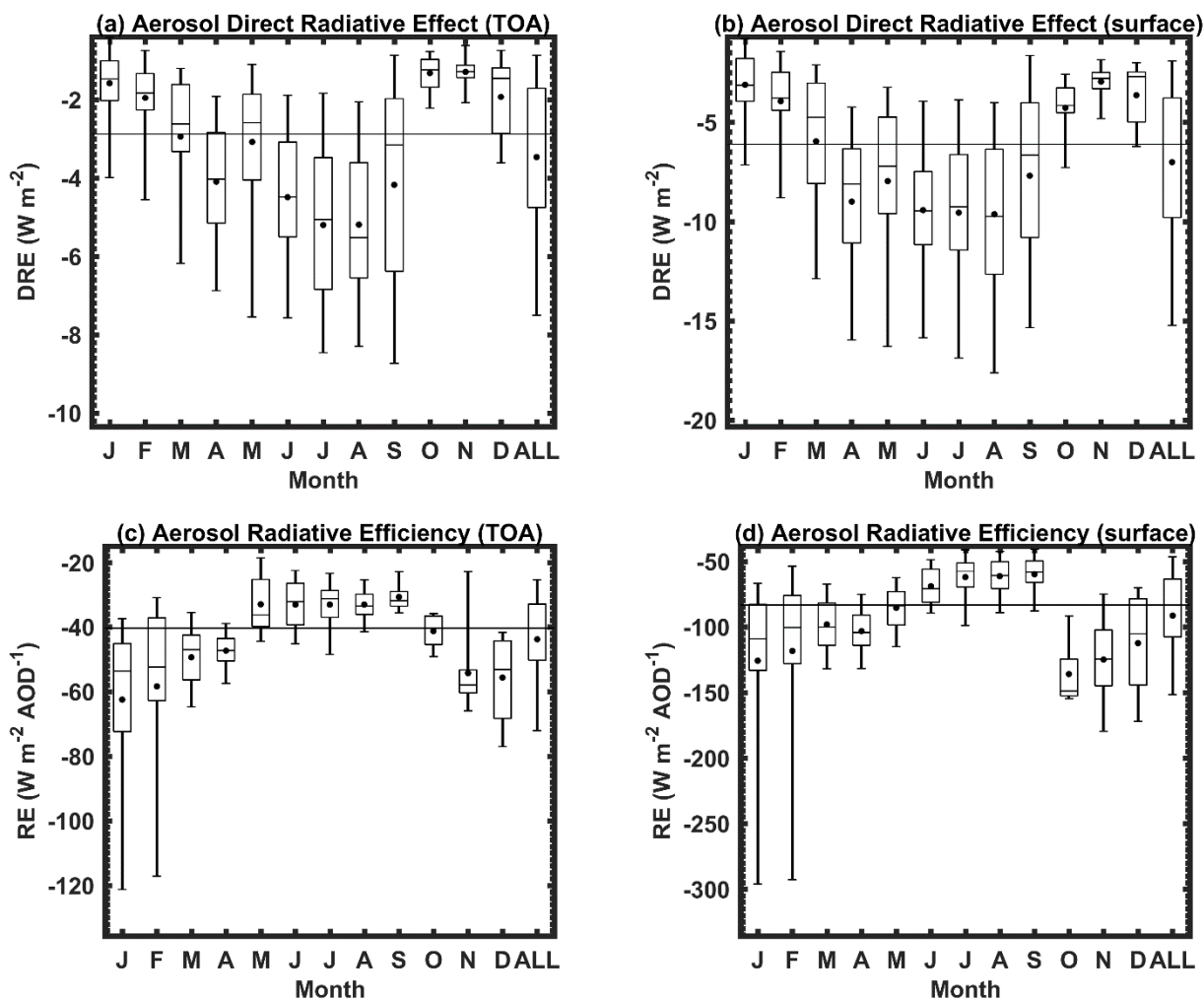
Figure 4. Boxplots of calculated monthly binned aerosol DRE and RE at the top of atmosphere and the surface. The 'ALL' box provides the statistics for all days in June 2012 thru February 2016 period. The mean for each month is denoted by the dot while the horizontal bar represents the median. The top and bottom of the box represent 75th and 25th percentiles while the top and bottom whisker extend to the 95th and 5th percentiles, respectively. The horizontal line drawn through all boxes of each plot represents the median value over the entire period (all months).
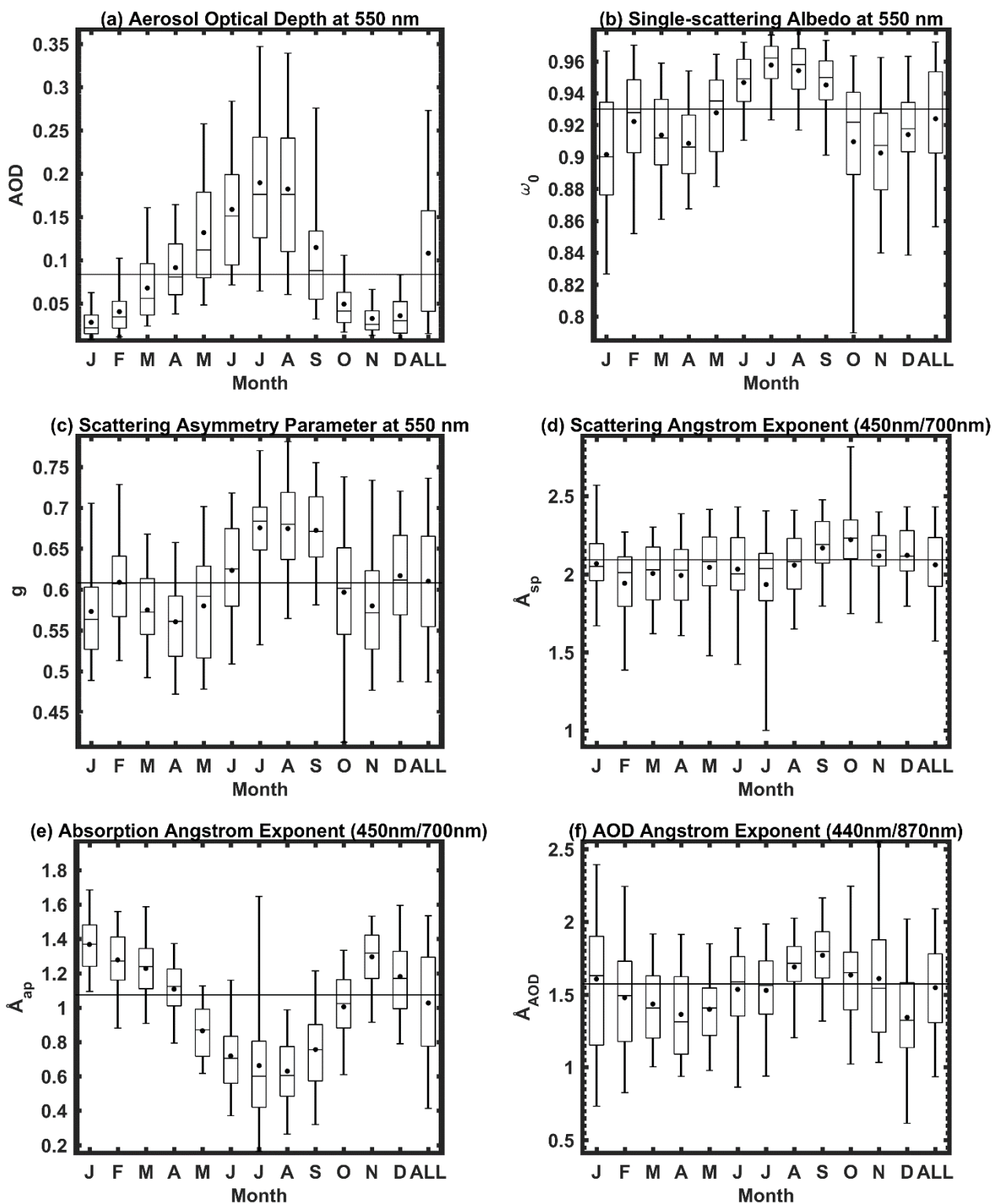
Figure 5. Boxplots of monthly binned aerosol optical properties at APP. The 'ALL' box provides the statistics for all days in June 2012 thru February 2016 period. The mean for each month is denoted by the dot while the horizontal bar represents the median. The top and bottom of the box represent 75th and 25th percentiles while the top and bottom whisker extend to the 95th and 5th percentiles, respectively. The horizontal line drawn through all boxes of each plot represents the median value over the entire period (all months).
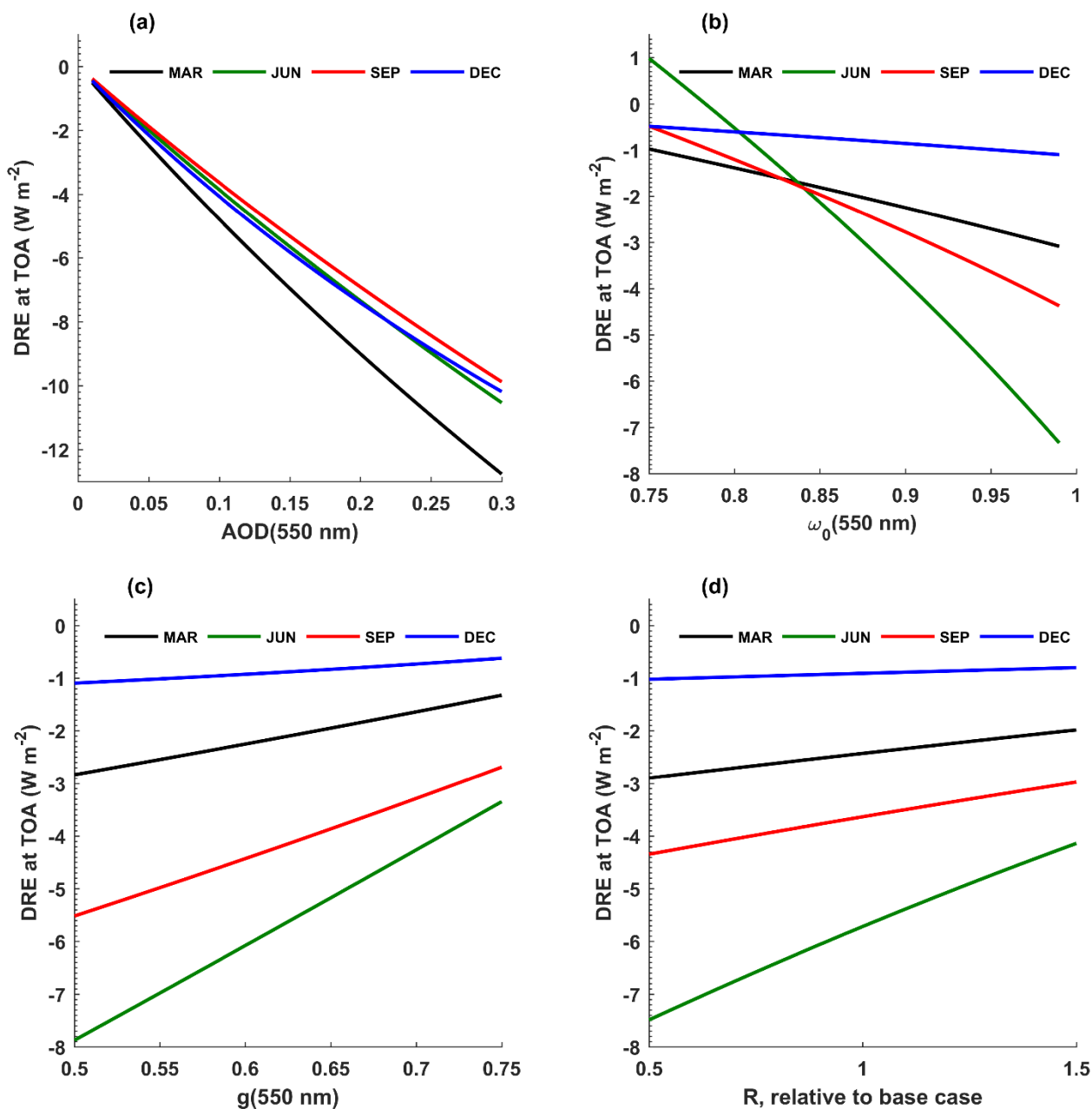
Figure 6. Seasonal dependence of the sensitivity of top-of-atmosphere (TOA) aerosol DRE to (a) AOD;
(b) ω_0; (c) g; and (d) spectral surface reflectance R, relative to the base case curves shown in Fig.3. For
example, a relative surface reflectance of 1.5 corresponds to scaling of the base case spectral reflectance
curves (Fig.3) by a factor of 1.5
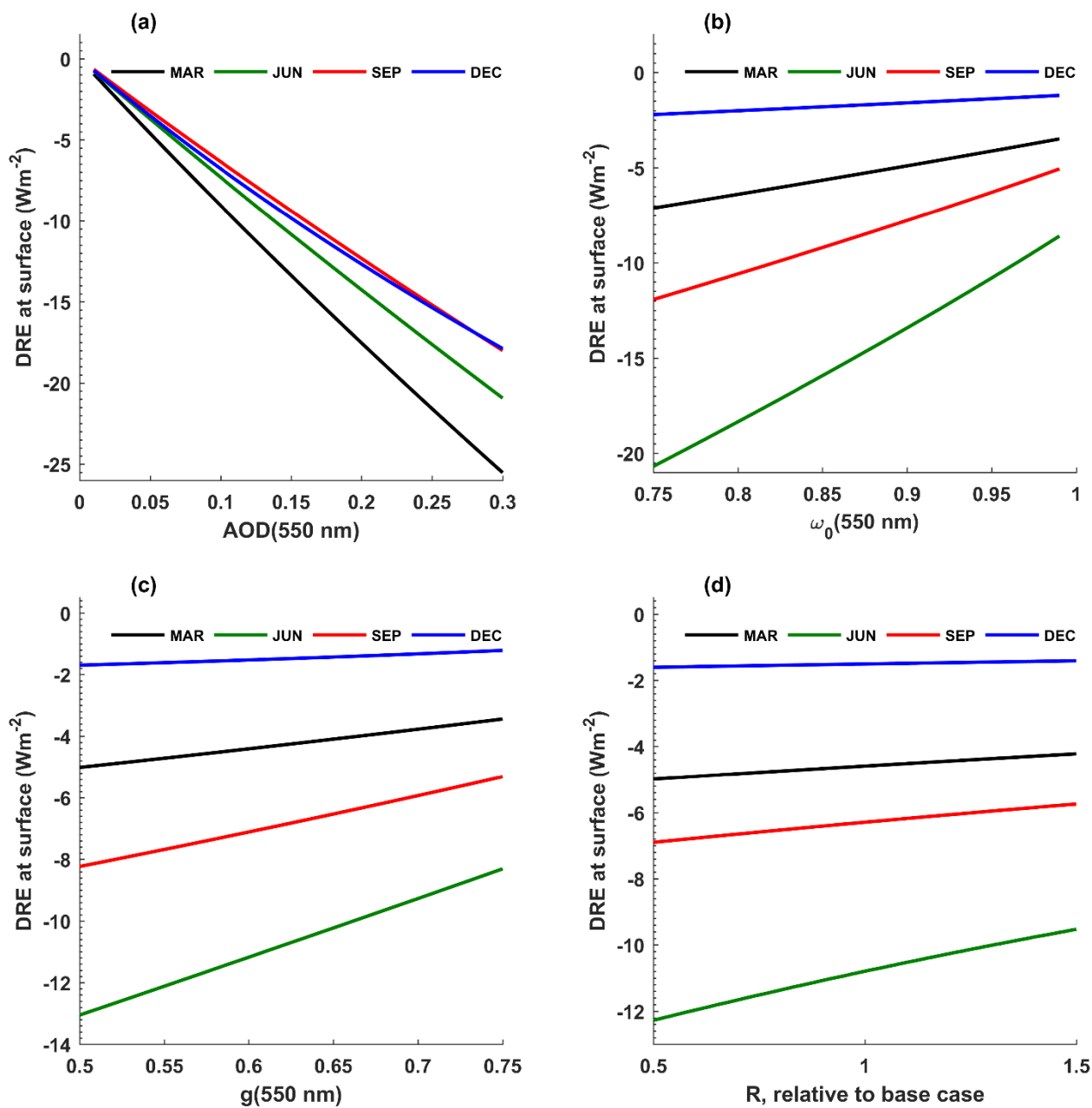
Figure 7. Seasonal dependence of the sensitivity of aerosol DRE at the surface to (a) AOD; (b) ω_0; (c) g; and (d) spectral surface reflectance R, relative to the base case curves shown in Fig.3.