Response to reviewer #1: We'd like to thank the reviewer for their critiques. Our responses to individual comments are below. In addition to the proposed changes to the manuscript, we have added a section (Section 3.6, Figures 11-13) in which the parameterizations for D_{pm} and sigma are tested against real smoke plumes observed at the Mount Bachelor Observatory (contribution by J.R. Laing and D.A. Jaffe).

This paper studies the effect of coagulation on particle diameter (Dm) and geometric standard deviation (sigma) in biomass burning plumes using a large-eddy simulation model with an online aerosol microphysical module. The topic is timely and the text is well written; however, there are some issues that need to be clarified before this manuscript can be accepted. In my opinion this paper presents a valuable base case for assessing the role of coagulation and condensation in future biomass burning studies.

Major comments

I have two major concerns regarding this study: the lack of organic aerosol (OA) chemistry and that while interpreting the results little attention is placed on plume dilution and its effect on coagulation.

Running the simulations as "coagulation-only" limits the usability of the results. For instance, new particle formation has been observed in biomass burning plumes (Hennigan et al., 2012; Vakkari et al., 2014), as well as up to a factor of 4 mass increase during the first few hours (Vakkari et al., 2014). Also a recent study by Konovalov et al. (2015) suggests that accounting for OA volatility can improve model performance significantly, although over a much longer time scale than what is considered here. Therefore, without OA evaporation and condensation, can the parameterisations in this paper be a good starting point for global and regional scale models (c.f. line 37)?

We agree that in order to simulate the mass of organic aerosol (or total $PM_{2.5}$) in the plumes properly, the dynamic production and/or loss of organic aerosol must be better understood. Our intention was not to claim to the contrary of this, and we had included the following discussion on line 389-394 of our text:

"One of the limitations of the coagulation-only parameterizations derived in this paper is that they do not include the effects of potential condensation/evaporation of organic aerosol on the aged biomass-burning size distribution. Both condensational growth and evaporative loss of OA has been observed previously in chamber studies and the field due to OA production or evaporation from dilution/chemistry (Cubison et al., 2011; Hecobian et al., 2011; Hennigan et al., 2011; Grieshop et al., 2009; Ortega et al., 2013; Jolleys et al., 2015; Vakkari et al., 2014)."

We agree that it is important to expand this discussion to discuss simulations where aerosol mass is a primary focus and have added the following to the lines above: "Thus, in order to predict biomass-burning aerosol mass, and thus the aerosol size distribution, we must understand how OA evolves in biomass-burning plumes.":

We also have in the text a discussion that a change in diameter through coagulation is different than a change in diameter due to OA condensation/evaporation: “It should be noted, however, that the Dpm growth attributed to OA condensation is not accompanied by a change in particle number (additional OA mass is distributed among existing particles), whereas a similar increase in Dpm growth by coagulation only would have an accompanying decrease in particle number. Thus, the changes to the aerosol size distribution and climatic influence of a size change due to coagulation and condensation are different.”

Our intention was to show that the evolution and variability of the mean/median particle diameter seems to be more dependent on the details of coagulation than condensation/evaporation. For example, if fresh biomass-burning particles have a median diameter of 50 nm, a doubling of OA mass would increase the diameter by 13 nm at most, which cannot explain the sizes of aged biomass-burning particles in the atmosphere. We did not intend to take away the importance of SOA that has a large impact on aerosol direct effects and PM2.5/health and a secondary importance for aerosol size and the indirect effect.

Finally, the Pierce group along with Matt Alvarado at AER have been recently funded by NSF to extend this work to include SOA formation into subgrid plume parameterizations of biomass burning. Adding SOA will take significant effort due to its complex nature and variability between plumes, and we did not feel that we had the time to test SOA properly during Kim Sakamoto’s Masters thesis, which ultimately dictated the scope of this work. We hope to have an updated scheme that includes SOA within the next several years that draws upon recent advances of biomass-burning SOA measurements from the lab and field.

On lines 453-455 it is concluded that SOA formation within the plume has a minor effect compared to coagulation. However, this is based on the assumption that SOA formation does not alter coagulation rate or sigma (lines 388-389). Could you elaborate on the conditions when these assumptions hold? For instance Pierce and Adams (2009) showed that secondary aerosol formation rate is one of the key parameters affecting how large fraction of small particles can grow up to CCN-sizes in new particle formation.

We have added the following sentence to our text at line 393: “For a 25% growth in diameter from SOA, which may be expected from for a factor of 2 increase in OA mass with a small change in sigma, we expect coagulation rates to stay within about 10% (Seinfeld and Pandis, 2006).”

It is correct that SOA can grow new particles to CCN sizes. However, this requires much more than a factor-of-2 increase in OA mass as might be expected in biomass-burning plumes. For example, a factor-of-8 increase in OA mass due to SOA production is required to grow a 40 nm particle to 80 nm. It is certainly plausible that this level of SOA production occurs in some biomass-burning plumes over longer timescales than recent studies; however, this is beyond the scope of this work.

My second major concern is related to the effect of plume dilution on the coagulation rate. Coagulation depends strongly on aerosol particle number concentration (as stated on lines 145-146). However, the observed changes in Dm and sigma are not discussed in terms of concentration, but only with respect to the input parameters and a rather arbitrary dM/dx (aerosol mass in an infinitesimally thin slice of air perpendicular to wind direction). The effect of dilution on coagulation is mentioned only briefly (e.g. lines 324-326), though Figure 4 shows that in most simulations the Dm and sigma change rapidly near emission, but very slowly later on. Is this decrease in the rate of change in Dm and sigma due to plume dilution and subsequent slowing of coagulation rate?

The aerosol concentration in-plume has two competing reducers: coagulation and plume dilution (both of which we'd expect to contribute to the rate changes in Figure 4). We expect the coagulation rate to slow as coagulation proceeds and number concentration drops (mean particle diameter increases). Plume dilution has a similar effect of reducing number concentration (and slowing coagulation) without a corresponding increase in particle diameter. Which of these factors dominates is primarily dependent on the size of the plume and the stability of the atmosphere. Large-diameter plumes are more resistant to plume dilution than small-diameter plumes. While we do not track dilution rate as a parameter, a lot of the associated variability is captured by the “Fire Area” parameter (large-area fires have larger-diameter plumes). We do not include dilution as an explicit input parameter in our analysis (even though plume dilution does occur in the LES simulations) because dilution is not an inherent property that would be provided by a coarse-grid aerosol model or an emissions inventory (unlike all of the other parameters that we studied), and thus we think it would be generally useless to have “dilution” as an input to a parameterization that will be used in coarse-grid models. We have added the following discussion of the importance of dilution in Line 81-82:

“The coagulational rate is therefore also affected by the rate of plume dilution (through a reduction in N), itself a function of plume size and meteorological conditions.’’

While dM/dx and dM/dxdz are somewhat arbitrary, we wanted a parameter which i) captured a large chunk of both Dpm and sigma variability and ii) was dependent on initial conditions only. These parameters only incorporate total aerosol mass loading, and plume vertical and with-wind extent. When dM/dv (initial concentration) was used in the simple parameterization, it was a slightly worse predictor of Dpm and sigma than dM/dxdz. The is because dM/dxdz is the product of dM/dv and the initial plume width; since wider plumes are less susceptible to dilution than narrower plumes, dM/dxdz captures this plume-width effect while dM/dv does not. This is discussed in Lines 393-396:

“dM/dV was also tested as a parameter within these simplified parameterization, but did not yield better agreements for either D_{pm} or σ than dM/dxdz despite incorporating an additional plume parameter (initial plume y-extent). This is because dM/dxdz is the product of dM/dV and the initial plume width; since wider plumes are less susceptible to dilution than narrower plumes, dM/dxdz captures this plume-width effect while dM/dV does not.”

How does Figure 4 look like if you colour it with concentration instead of dM/dx, or plot

Dm and sigma against concentration? The dM/dx takes into account only dilution along the wind direction, not dilution due to vertical or cross-wind mixing.

Figure 4 does not look qualitatively different if either dM/dxdz, dM/dx, or dM/dv (concentration) are used (see attached). We chose dM/dx as it was the simplest parameter in which the Dm plot shows a qualitative trend. We have updated the figure to dM/dxdz instead of dM/dx to illustrate the stability of this trend when accounting for dilution in the vertical. Initial concentration (dM/dV) was not used for consistency with Figures 9-10 where dM/dV is not an improvement over dM/dxdz in the simple parameterization (for the reasons stated above) despite requiring more fire information.

Can you identify a range (time and space), where coagulation can cause significant changes in the size distribution and after which the plumes become so diluted that coagulation slows down? How would this turning point depend on the initial concentration (emissions) and the meteorological conditions (turbulent mixing) during transport?

Coagulation rates are always slowing down in plumes (proportional to the square of number concentration), so we would need define that it slows down to a subjectively chosen rate. Whether or not a chosen cutoff rate is appropriate depends on the timescale that one cares about, which depends on the modeling application. Since this answer also depends on all of the input factors studied here, the answer is not straightforward and is out of the scope of this paper.

The background aerosol is assumed to be negligible compared to the plume and is set to zero (lines 176-179). However, in ambient air measurements this assumption cannot be made – see e.g. Yokelson et al. (2009). Have you verified that your plumes are so concentrated even after 200km transport that this assumption still holds? When will coagulation rate with the background aerosol become similar to coagulation within the biomass burning mode?

At the point where the plumes are diluted to ambient-air concentrations, subgrid processing is no longer different from standard grid-size processing and model schemes are sufficient (no parameterization is needed). We have added an explanation of this at Line 178:

“In cases where the plume dilutes to similar concentrations to the ambient background, subgrid-plume coagulation schemes are no longer necessary, and grid-resolved coagulation will properly account for coagulation.”

How is turbulent mixing handled in the simulations? Table 1 (page 23) lists “Mixing depth of aerosol layer” as an input parameter, yet on line 214 “mixing depth” is calculated from the simulated vertical profile of aerosol mass. Is this related to mixed layer height (e.g. height of convective planetary boundary layer)?

SAM calculates turbulence explicitly in the LES using thermodynamic profiles and surface heat/momentum fluxes. There is no assigned PBL in SAM, though it can be calculated from the

resulting profiles. We have used “mixing depth” and “injection height” to refer to the aerosol layer and *not* the convective planetary boundary layer in general (we do not explicitly calculate it anywhere).

Injection height is the height at which the biomass-burning aerosol is emitted into the SAM-LES gridboxes after initialization, after which it is subject to turbulent mixing within the model. Depending on meteorological factors, this initialized aerosol layer can remain stable above the PBL or be mixed down through the PBL to the ground (see revised Figure 3). *Mixing depth* is the term used to describe the height extent of the aerosol layer after it is mixed in the LES (at a given point in the LES simulation) to the final time/distance of the simulation, which may vary significantly from the injection height and/or PBL height.

We have changed Line 153 to emphasize our definition of mixing depth:

“...and aerosol mixing depth (hereafter referred to as *mixing depth;* the vertical extent of the aerosol plume)”.

Table 1 has been changed to “Mixing Depth” only for consistency with this definition.

If Figure 3 is a representative sample of the simulations it seems as if majority of the plumes are not in the mixed layer but above it, as they do not reach the surface. Again, I would expect the turbulent mixing in convective PBL (or the lack of convective mixing in free troposphere or the residual layer) to have a significant effect on plume dilution and therefore the coagulation rate. Is it so?

Line 824, Figure 3 Please provide a legend for the lines (indicating input and meteorological parameters).

We chose our 6 profiles at random for the Figure 3. For 5 of the 6 these cases, the initial height at which we put the plume in the SAM-TOMAS simulation was not in the mixing boundary layer in the model (which is calculated independently by SAM-TOMAS using the meteorological and surface conditions that were randomly chosen for each case). Some simulations were at night where the mixing layer was very small <100m and other cases the injection heights were higher than the daytime mixing layer. We do not expect this to have a great impact on our results as the results should be most sensitive to the depth of the layer the plume mixes to and not whether the plume touches the ground or not.

SAM-TOMAS incorporates a multitude of meteorological parameters (e.g. atmospheric stability) for each simulation in addition to the 7-inputs we varied. Instead of citing specific simulations, we simplified Figure 3 to two representative SAM-TOMAS vertical profiles - one showing the case where the aerosol plume mixes down to the ground through the mixing layer, and one where the plume remains suspended at the injection height throughout the simulation.

The following caption now accompanies Figure 3:

“Final vertical profiles for two representative SAM-TOMAS simulations after four hours, normalized to individual aerosol load and averaged horizontally across the domain. The black profile shows a simulation where the aerosol fully mixed through the boundary layer to the ground, while the red profile shows a simulation where the aerosol plume still stable at the emission injection layer.”

Minor comments:
Line 69-71 There are some more recent studies which you might want to look up. For instance Akagi et al. (2012), Hennigan et al. (2012), Ortega et al. (2013), Vakkari et al. (2014), Jolleys et al. (2015), Konovalov et al. (2015) and May et al. (2015) come tomy mind.

These references have been added to the appropriate lines with the exception of Ortega et al. (2013),which was already cited at these lines. Konovalov et al. (2015) was added to Section 3.5:

“Konovalov et al. (2015) has emphasized the importance of OA simulation in modeling in long-range (>1000 km) plume evolution.”

Line 382-385 Also for this statement some more recent references could be considered.

More recent studies have been appended.

Line 71 “This SOA condenses onto existing particles causing growth of the aerosol size distribution.” Please reconsider this statement as there are observations of new particle formation in biomass burning plumes (Hennigan et al., 2012; Vakkari et al., 2014).

We have added the following discussion of new-particle formation in biomass-burning plumes in lines 71-74:
“This SOA can condense onto existing particles causing growth of the aerosol size distribution. It can also spur new-particle formation in biomass-burning plumes as has been observed in lab studies (Hennigan et al., 2012) and field campaign analyses (Vakkari et al., 2014).”

And in Lines 233-235:

“We do not address new-particle formation in biomass-burning plumes in this work. In plumes where new-particle formation in biomass-burning plumes occurs, our parameterizations will underestimate the number of particles and overestimate the mean diameter of the plume particles.”

Line 153 “Mixing depth had a range of 150-2500 m” but Table 1 (page 23) gives “Mixing depth” limits as 120 m and 2500 m. Which one is it?

The mixing-depth limit is 150m, and we have updated Table 1.

Line 169-170 "The algorithm simulated the size distribution across 15 logarithmically-spaced size bins spanning 3 nm-10 μm." This leaves quite few bins for the size range of interest. Can coarse size resolution become an issue for the coagulation calculation?

There are 10 bins between 10 nm and 1 μm capture the area of biomass-burning aerosol emission and growth. There are only two bins above 1 μm for coarse aerosol. As we track two independent moments (number and mass) within each bin, the model fidelity is much higher than single-moment schemes with similar resolution. The TOMAS microphysics algorithm has been evaluated in Lee and Adams, (2012) and it generally captures the processes similarly to higher-resolution versions of the model. We now clarify the bin structure: "The algorithm simulated the size distribution across 135 logarithmically -spaced size bins spanning 3 nm-1 μm with 2 additional bins spanning 1-10 μm."

The coarse size resolution can have an impact on the sigma calculation of the mode near smaller modal sizes (~1.32) where the distribution is less than a single bin-width. This is discussed in lines 263-265.

Line 191-192 "We ran 100 SAM-TOMAS simulations at 500 m x 500 m horizontal resolution (total horizontal extent = 100 km)," but Figure 4 x-axis extends to > 200 km. I assume these are the same data because on lines 244-245 it is stated that "Figure 4 shows the Dpm (panels a and c) and σ (panels b and d) as a function of distance for each of the 100 SAM-TOMAS simulations used to train the emulator (Sect. 3.2)." What was the horizontal extent?

By "total horizontal extent" we were referring to the cross-wind (y-direction) extent, not the with-wind (x-direction) extent as plotted in Figure 4. Lines 190-196 have been clarified to:

"We ran 100 SAM-TOMAS simulations at 500 m x 500 m horizontal resolution (total cross-wind (y-direction) horizontal extent = 100 km), and constant 40 m vertical resolution (total vertical extent = 4 km)... The output from each SAM-TOMAS simulation was recorded at four different times (400 total time slices across 100 simulations) as the plume progressed along the with-wind (x-direction) axis."

Line 871, Figure 4 There are so many overlying lines that it is getting difficult to read. Please consider if you can clarify it.

We like that Figure 4 shows the evolution of diameter and sigma across all of our simulations so that the reader can see what variability is captured by our simulations. We dedicate the rest of the paper to generalizing the results, so we prefer to keep Figure similar (though we have updated to use dM/dxdz rather than dM/dx).

Response to reviewer #2: We'd like to thank the reviewer for their critiques. Our responses to individual comments are below. In addition to the proposed changes to the manuscript, we have added a section (Section 3.6, Figures 11-13) in which the parameterizations for Dpm and sigma are tested against real smoke plumes observed at the Mount Bachelor Observatory (contribution by J.R. Laing and D.A. Jaffe).

This paper investigates the influence of coagulation on the particle number size distribution, notably on the mean diameter (Dm) and geometric standard deviation (Sigma) of a single particle mode, in biomass burning plumes. The work is based on a large number of sophisticated model simulations. The authors investigate how Dm and Sigma evolve with time in biomass burning plumes, and how their evolution is related to several parameters associated with primary particle emissions, fire conditions and atmospheric conditions. The authors compare briefly the influence of coagulation to that caused by organic aerosol formation/loss in a plume. The authors finally parameterize their results to a form that is applicable in large-scale modeling frameworks.

The is scientifically sound and original. The text is well organized and easy to read (with a couple of minor exceptions mentioned below). The authors are able to explain very well the numerous results obtained from model simulations. I have only a few minor suggestions for revisions.

Scientific comments:
I have a hard time of understanding Figure 3, even after reading the text on lines 210-213. I recommend that the authors work a bit more to make their message here clearer to readers.

In response to this comment and a comment by reviewer 1, we have simplified Figure 3 to show only two representative profiles of SAM-TOMAS simulations: one which shows the aerosol plume mixing through the boundary layer to the ground, and a second which shows the plume still suspended at the emission injection height at the end of the simulated run. These are the two “types” of mixing depths possible (to ground through the PBL and suspended).

The Figure 3 caption was updated to reflect the changed figure:

“Final vertical profiles for two representative SAM-TOMAS simulations after four hours, normalized to individual aerosol load and averaged horizontally across the domain. The black profile shows a simulation where the aerosol fully mixed through the boundary layer to the ground with some aerosol trapped in a stable layer above the boundary layer, while the red profile shows a simulation where the aerosol plume still stable at the emission injection layer.”

Lines 254-255. The authors state that the initial mode mean diameter have little effect on Dm. I do not get this point when looking at Figures 4a and c: if Dm is initially large, it seems to typically lead to higher values of Dm at later plume times compared to cases where Dm is initially small. Could the authors specify what they mean here?

We mean to emphasize that the variability in Dpm can mainly be attributed to factors outside of initial Dpm. While those simulations with higher initial diameters do climb to higher final diameters than others, the final Dpm variability is not principally driven by starting diameter. As the dominant Dpm variability factors form a discussion in Section 3.3, we have removed this line to alleviate confusion.

Line 390: Is this correct? Condensation of a non-volatile vapor into a single mode tend to narrow this mode, not widen it, as stated here.

Your are correct, we wrote the opposite of what we meant. Line 390 is now: “These assumptions are imperfect as irreversible condensation (evaporation) decreases (increases) σ…”.

The authors analyze shortly the influence OA production/loss on their results (section 3.5), and discuss also the potential effects of cloud processing (lines 463-469). This is clearly sufficient for these two processes in this paper. However, the authors do not mention at all new particle formation (NPF) that has been estimated to be a frequent process in biomass burning plumes. NPF might have notable effects on aerosol size distribution, and thereby on both Dm and Sigma, in evolving biomass burning plume. The authors should spend at least a few lines on discussing the relevance of this process in biomass burning plumes and on the potential effects of NPF on their results.

We have changed the following discussion of SOA pathways in Section 1.1 (Lines 77-79):

“This SOA can condense onto existing particles causing growth of the aerosol size distribution. It can also spur new-particle formation in biomass-burning plumes as has been observed in lab studies (Hennigan et al., 2012) and field campaign analyses (Vakkari et al., 2014).”

We have added the following lines to our methods:

“We do not address new-particle formation in biomass-burning plumes in this work. In plumes where new-particle formation in biomass-burning plumes occurs, our parameterizations will underestimate the number of particles and overestimate the mean diameter of the plume particles.”

Technical issues:

Line 265: Figures 5 shows Sigma versus Dm rather than Dm versus Sigma.

Line 382: "..OA has been. . ."

These have been corrected in text.

The evolution of biomass-burning aerosol size distributions due to coagulation: dependence on fire and meteorological details and parameterization

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Abstract

Biomass-burning aerosols have a significant effect on global and regional aerosol climate forcings. To model the magnitude of these effects accurately requires knowledge of the size distribution of the emitted and evolving aerosol particles. Current biomass-burning inventories do not include size distributions, and global and regional models generally assume a fixed size distribution from all biomass-burning emissions. However, biomass-burning size distributions evolve in the plume due to coagulation and net organic aerosol (OA) evaporation or formation, and the plume processes occur on spacial scales smaller than global/regional-model grid boxes. The extent of this size-distribution evolution is dependent on a variety of factors relating to the emission source and atmospheric conditions. Therefore, to account for biomass-burning aerosol size in global models accurately requires an *effective* aerosol size distribution that accounts for this sub-grid evolution and can be derived from available emissions-inventory and meteorological parameters.

In this paper, we perform a detailed investigation of the effects of coagulation on the aerosol size distribution in biomass-burning plumes. We compare the effect of coagulation to that of OA evaporation and formation. We develop coagulation-only parameterizations for effective biomass-

burning size distributions using the SAM-TOMAS large-eddy simulation plume model. For the most-sophisticated parameterization, we use the Gaussian Emulation Machine for Sensitivity Analysis (GEM-SA) to build a parameterization of the aged size distribution based on the SAM-TOMAS output and seven inputs: emission median dry diameter, emission distribution modal width, mass emissions flux, fire area, mean boundary-layer wind speed, plume mixing depth, and time/distance since emission. This parameterization was tested against an independent set of SAM-TOMAS simulations, and yields R^2 values of 0.83 and 0.89 for D_{pm} and modal width, respectively. The size distribution is particularly sensitive to the mass emissions flux, fire area, wind speed, and time, and we provide simplified fits of the aged size distribution to just these input variables. The simplified fits were tested against eleven aged biomass-burning size distributions observed at the Mt. Bachelor Observatory in August 2015. The simple fits captured over half of the variability in observed D_{pm} and modal width even though the freshly emitted D_{pm} and modal widths were unknown. These fits may be used in global and regional aerosol models. Finally, we show that coagulation generally leads to greater changes in the particle size distribution than does OA evaporation/formation using estimates of OA production/loss from the literature.

1. Introduction

1.1 Biomass-burning aerosols

Biomass burning (including wildfires, prescribed fires, and agricultural fires) releases significant amounts of gas- and particle-phase species to the atmosphere (Andreae and Merlet, 2001; Reid et al., 2005). The particle-phase emissions are composed primarily of a mixture of organic aerosol (OA) and black carbon (BC) with some inorganic species (e.g. potassium), and the ratios of these species depend on the source fire conditions (Capes et al., 2008; Carrico et al., 2010; Cubison et al., 2011; Hecobian et al., 2011; Hennigan et al., 2011; Reid et al., 2005). These aerosols affect the global radiation budget through the indirect and direct aerosol effects (Boucher et al., 2013). The smoke particles themselves are able to act as cloud condensation nuclei (CCN) and increase cloud albedo and lifetime (indirect aerosol effect; Lee et al., 2013; Pierce et al., 2007; Spracklen et al., 2011) as well as scattering/absorbing incoming solar-radiation directly (direct aerosol effect; Alonso-Blanco et al., 2014; Boucher et al., 2013; Haywood and Boucher, 2000; Jacobson, 2001).

Particle size has a significant effect on the magnitude of both the direct and indirect aerosol

effects (Lee et al., 2013; Seinfeld and Pandis, 2006; Spracklen et al., 2011). The composition and diameter of the particles affect their absorption/scattering efficiencies, which dictate the amount of solar radiation absorbed/scattered per emitted mass of particles (Seinfeld and Pandis, 2006). Particle diameter and hygroscopicity determine the particles' ability to act as a CCN and influence cloud processes, and the total number of emitted particles increases with decreased particle size if total mass emissions are fixed. Spracklen et al., (2011) found that a reduction by a factor of two in particle size for all carbonaceous aerosols (for a fixed total aerosol mass) resulted in a ~300% increase in the cloud albedo indirect effect globally, as more particles were available to act as CCN. Lee et al., (2013) determined that CCN concentrations in the GLOMAP model were very sensitive to uncertainties in biomass-burning emission diameter on both the regional and global scale (its attributable CCN uncertainty ranked third of 28 factors tested globally). Therefore, to ascertain the role of biomass-burning aerosols in climate forcings accurately, biomass-burning size distributions must be well represented in aerosol-climate models.

Size distributions are subject to physical and chemical processing in the plume. The formation of secondary organic aerosol (SOA) has been observed in lab studies of biomass-burning aerosol (Cubison et al., 2011; Grieshop et al., 2009; Hennigan et al., 2011; Heringa et al., 2011; Ortega et al., 2013) and in field campaigns (DeCarlo et al., 2010; Lee et al., 2008; Reid et al., 1998; Yokelson et al., 2009). This SOA can condenses onto existing particles causing growth of the aerosol size distribution. It can also spur new--particle formation in biomass-burning plumes as has been observed in lab studies (Hennigan et al., 2012) and field campaign analyses (Vakkari et al., 2014).- Conversely, recent lab and field studies have characterized primary organic aerosol (POA) as semi-volatile, with plume dilution allowing the evaporation of organic aerosol from particles (Huffman et al., 2009; Cubison et al., 2011; May et al., 2013, 2015; Jolleys et al., 2015,–). The cumulative net effects of OA production/loss within biomass-burning plumes has been found to be highly variable from fire to fire (Akagi et al., 2012; Hennigan et. al, 2011).

_____Coagulation is also important for size-distribution evolution as it reduces particle number and shifts the distribution to larger sizes. Coagulation rates are proportional to the square of the particle number concentration (all else remaining fixed), so the high number concentrations in biomass-burning plumes relative to background can lead to rapid coagulational growth of the size distribution. The coagulation rate is therefore also affected by the rate of plume dilution (through a reduction in N), itself a function of plume size and meteorological conditions. The rate and magnitude of the aerosol growth

caused by these combined processes is a function of aging time, emission source characteristics, aerosol properties at emission, and atmospheric conditions.

These condensation/evaporation and coagulation aging processes affect both the composition and size of the aerosol size distribution – both properties that influence the extent to which smoke particles affect climate. While fresh smoke is generally composed of fine particles between 20-60 nm in diameter (Levin et al., 2010), condensation and coagulation cause rapid aerosol growth to larger sizes (over 100 nm) on timescales of often less than 24 hrs (Janhäll et al., 2010). However, Janhäll et al., (2010) found the observed geometric mean diameter of aged biomass-burning particles varied between 170-300 nm, with geometric standard deviations (hereafter referred to as "modal width") between 1.3-1.7 with significant dependence on fuel type and modified combustion efficiency. It is currently unclear to what extent these factors and others drive the variability in aged size distributions.

As stated earlier, an accurate representation of aged biomass-burning aerosol size is necessary for predictions of aerosol climate effects in regional and global models (Lee et al., 2013). Current wildfire inventories are mass-based (neglecting aerosol size data), and thus regional and global models used for aerosol-climate effects generally specify fixed, "aged" size distributions that do not account for sub-grid processing of the emitted particles (Reid et al., 2009; van der Werf et al., 2010; Wiedinmyer et al., 2011). Any variability in the biomass-burning size distribution due to fire or emissions characteristics and meteorology are not accounted for, nor is it clear what the best "aged" size distribution to use is in these models.

In this paper, we perform a detailed investigation of coagulation in biomass-burning plumes and compare to the effects of OA evaporation and formation. We investigate the factors that influence coagulational growth of the particles in the plume. These factors include fire area, particle-emissions mass flux, particle-emissions size, and meteorological conditions. We create parameterizations of varying degrees of complexity for median dry diameter (D_{pm}) and lognormal modal width (σ) of the aged biomass-burning size distributions as a function of these input parameters, based on detailed numerical simulations using a large-eddy model with embedded aerosol microphysics (SAM-TOMAS). Finally, we compare the effect of coagulation on the aerosol size distribution to that of OA production/evaporation.

We describe the parameterization building process, including the use of a Gaussian emulator, in Sect. 2. A discussion of input and output ranges, processing, and constraints of the parameters we have

chosen is provided in Sect. 2.1. We discuss the SAM-TOMAS model and the emulation process in Sect. 2.2-2.3. Sections 3.1-3.2 contain the results of the SAM-TOMAS model and the emulator. We discuss emulator sensitivities to the inputs in Sect. 3.3 and present a series of simplified fit equations for the effective size distributions in Sect. 3.4. We discuss the effects of potential OA production/loss on our size distribution estimates in Sect. 3.5. The simplified--fit equations are tested against biomass-burning plumes observed at the Mt. Bachelor Observatory in Sect. 3.6. Finally, we conclude in Sect. 4, including future plans for testing the parameterization and known existing limitations.

2. Methods

Figure 1 provides an overview of our methods that will be described in detail in the subsections below. In short, we used a Large-Eddy Simulation model, the System for Atmospheric Modelling (SAM; Khairoutdinov and Randall, 2003), with the online aerosol microphysics module, TwO Moment Aerosol Sectional (TOMAS, Adams and Seinfeld, 2002; Stevens et al., 2012) to simulate the evolution of the biomass-burning aerosol size distribution by coagulation across a wide range of emission and meteorological conditions. We used the SAM-TOMAS size distributions to build parameterizations to predict aged D_{pm} and σ using: (1) a statistical emulator of the SAM-TOMAS model itself and (2) simplified fits to the SAM-TOMAS output data. The statistical emulator was built by the Gaussian Emulation Machine for Sensitivity Analysis (GEM-SA), and we used the emulator and SAM-TOMAS data to determine the relative importance of various inputs to shaping the aged size distribution.

2.1 Investigated factors that may lead to variability in aged size distributions

We investigated seven parameters that may affect the aging of the biomass-burning aerosol size distribution. These can be divided into those representing the initial lognormal-mode size parameters (D_{pm0}, σ_0), fire conditions (mass flux, fire area), atmospheric conditions (wind speed, plume mixing depth), and time. Each of these parameters is generally available in large-scale aerosol models, which means a parameterization for aged biomass-burning size distributions based on these parameters may be used in these models. Table 1 lists these input parameters and the ranges of values tested in this work.

We assumed that the initial size distributions were a single lognormal mode (described by dry median diameter, D_{pm}, and modal width, σ), which is sufficient when representing both fresh and aged observed biomass-burning size distributions (Capes et al., 2008; Janhäll et al., 2010; Levin et al., 2010; Sakamoto et al., 2015). The initial size-distribution parameters specify the median dry diameter (D_{pm0}) and modal width (σ_0) of the freshly emitted aerosol distribution. We varied these parameters between 20-100 nm for D_{pm0} and 1.2-2.4 for σ_0. The large ranges are due to variability in combustion efficiency and fuel-type factors as seen in lab and observational studies (Janhäll et al., 2010; Levin et al., 2010).

Fire area, mass flux, wind speed and aerosol mixing depth (hereafter referred to as *mixing depth;* **the** vertical extent of the aerosol plume) all affect the aerosol number concentration (N) within the plume, which in turn affects the coagulation rate (proportional to N^2). In our simulations, we constrained mass flux to 2×10^{-8} - 5×10^{-6} kg m^{-2} s^{-1} using approximate maximum and minimum values of summed black carbon and organic carbon flux (BC+OC) found in the Global Fire Emissions Database ver. 3 (GFED3; van der Werf et al., 2010; available from http://www.globalfiredata.org). Fire area ranged from 1 - 49 km^2 (simulated as a square), which was found to represent the range of fire sizes in GFED3. Boundary layer wind speed varied between 2 m s^{-1} and 20 m s^{-1} and was based on ranges in the National Center for Environmental Prediction (NCEP) North American Regional Reanalysis (NARR) meteorology (Mesinger et al., 2006) during the fire season (specifically, July, 2010). Mixing depth had a range of 150-2500 m (based on SAM-TOMAS output; see Sect. 2.2).

The aging time was the final input parameter, and we used 5 hr (300 min) as an upper time bound due to this being a typical timescale for transport across large global model gridboxes.

2.2 The SAM-TOMAS model

We used the SAM-TOMAS model to simulate the evolution of biomass-burning aerosol size distributions due to coagulation across the range of input parameters described above. SAM (Khairoutdinov and Randall, 2003) is a dynamical large-eddy simulation (LES) model, which has previously been used to model emissions plumes (Lonsdale et al., 2012; Stevens et al., 2012; Stevens and Pierce, 2013). We ran the model in Lagrangian 2D mode (Stevens and Pierce, 2013), in which a wall oriented normal to the mean boundary layer wind moves at the mean boundary-layer wind speed. This moving wall tracks the radial dispersion of a plume as it travels downwind (Fig. 2). This 2D mode is computationally efficient compared to the full 3D model with minor differences due to axial plume

symmetry (Stevens and Pierce, 2013).

The size distributions of the aerosol particles in SAM were simulated using the TwO Moment Aerosol Sectional (TOMAS; Adams and Seinfeld, 2002) microphysical scheme embedded into SAM. The algorithm simulated the size distribution across 135 logarithmically -spaced size bins spanning 3 nm-10 μm with 2 additional bins spanning 1-10 μm. The aerosol size distribution was tracked via two independent moments for each bin of the size distribution (mass and number). TOMAS calculated coagulation explicitly in each grid cell assuming a Brownian diffusion kernel (Seinfeld and Pandis, 2006). Our SAM-TOMAS simulations included only coagulation, and particles were assumed to be a single species (no differentiating between BC and OA). The SAM-TOMAS model had previously been tested against observations in Stevens et al. (2012) and Lonsdale et al. (2012) for power plant plumes.

We set background aerosol concentrations to zero as the biomass-burning aerosol concentrations emitted into SAM-TOMAS were orders of magnitude larger than those present in a remote background location, and as such the lack of background aerosol would have had an insignificant effect on the rate of in-plume coagulational processing. In cases where the plume dilutes to similar concentrations to the ambient background, subgrid--plume coagulation schemes are no longer necessary, and grid-resolved coagulation will properly account for coagulation. The biomass-burning aerosol was assumed to have a constant density of 1400 kg m^{-3} as primarily a mix of organic compounds, thus we do not consider how changes in BC/OA composition may affect density and coagulation rates. The hygroscopicity of the aerosol particles was set to zero, allowing no water uptake. This assumption is not true of real world biomass-burning aerosol and has been characterized in other works finding hygroscopicities of fresh (κ=0.02-0.8; Petters et al., 2009) and aged smoke (κ=0.1-0.3; Engelhart et al., 2012) with a strong dependence on fuel type. In terms of their effect on the size distribution, a constant κ across all particle sizes has the simple effect of increasing the effective diameter of the particles via water uptake by a scalar factor. This initial increase should only have a relatively minor effect on the final dry D_{pm} or σ of the plume after coagulational processing as the mean coagulation rates are relatively insensitive to the size shifting of a particle population (Seinfeld and Pandis, 2006; Stuart et al., 2013).

We ran 100 SAM-TOMAS simulations at 500 m x 500 m horizontal resolution (total cross-wind (y-direction) horizontal extent = 100 km), and constant 40 m vertical resolution (total vertical extent = 4 km). This resolution accommodated the chosen plume parameters (see Sect. 2.1). The model was run with a master timestep of 2 seconds (varied internally for accuracy in the coagulation calculation) for a

duration of 5 model hours (300 minutes). The output from each SAM-TOMAS simulation was recorded at four different times (400 total time slices across 100 simulations) as the plume progressed along the with-wind (x-direction) axis.

The seven inputs to the SAM-TOMAS model were constrained to capture a range of biomass-burning characteristics in realistic scenarios and are summarized in Table 2. The ranges of values used for D_{pm0}, σ_0, fire area and mass flux are the same as those listed in Table 1. The meteorological fields were supplied by NCEP reanalysis meteorology from over North America (land only, lat: 30º - 70º N, lon: 70º -135º W) during the July 2010 fire season. The SAM-TOMAS wall speed was set equal to the mean boundary layer wind speed from NCEP. We filtered these inputs by requiring wind speed > 2 m s^{-1} to eliminate stagnation situations over the source. The injection height (lower bound) and injection depth of the aerosol were specified at between 50-1500 m and 500-2000 m respectively. No emission injection parameterization (e.g. Freitas et al., 2007) was used as we were only trying to capture a range of mixing depths for our aging calculation, and the absolute height was relatively unimportant. All the SAM-TOMAS simulation inputs were chosen using semi-random Latin hypercube sampling across the ranges listed above (Lee et al., 2012). The results of the full SAM-TOMAS simulation set are summarized in Sect. 3.1.

We calculated the time-dependent mixing depth of the plume from vertical profiles averaged horizontally across the entire simulation wall at each time slice. Figure 3 shows a sample of ~~six~~two vertical profiles from different SAM-TOMAS simulations. The mixing depth was defined as the range of altitudes where the aerosol mass was greater than half of the peak aerosol mass:

$$\text{mixing depth} = \Delta_{\text{alt 50\% peak aerosol mass}}$$

In cases where the plume mixed down to the ground, the lower altitude bound was defined as 0 m. Runs with mixing depths greater than 2500 m were excluded to ensure that the plume did not reach the model top. In addition to mixing depth, D_{pm} and σ were calculated for each of the SAM-TOMAS time slices from the first and third integrated moments of the size distribution as detailed by Whitby et al. (1991).

We do not address new-particle formation in biomass-burning plumes in this work. In plumes where new-particle formation in biomass-burning plumes occurs, our parameterizations will underestimate the number of particles and overestimate the mean diameter of the plume particles.

2.3 Emulation of the SAM-TOMAS output

As running the full SAM-TOMAS model is too computationally expensive for implementation in global aerosol models, we built an offline emulator of the model for use as a parameterization in these global models. We created the emulator using the Gaussian Emulation Machine for Sensitivity Analysis (GEM-SA) developed by the Centre for Terrestrial Dynamics (http://www.ctcd.group.shef.ac.uk/gem.html). The GEM-SA software uses a Gaussian process to design a SAM-TOMAS simulator (the emulator) based on the behavior of the known SAM-TOMAS inputs and outputs (the training data). A complete description of GEM-SA statistics and assumptions can be found in Kennedy and O'Hagan (2001) and Kennedy et al. (2008). A description of its application as an estimator in atmospheric-aerosol modelling can be found in Lee et al. (2011). This software was previously used in sensitivity studies in atmospheric-aerosol (Lee et al., 2011, 2012) and vegetation models (Kennedy et al., 2008).

We used 400 data points from the set of 100 SAM-TOMAS simulations to train the emulator. GEM-SA assumes that the outputs are a continuous and differentiable function of the inputs to statistically emulate the model and estimate the SAM-TOMAS output (D_{pm} and σ). We used a new set of completed SAM-TOMAS simulations (624 non-training data points) to test our GEM-SA parameterization for accuracy relative to SAM-TOMAS (see Sect. 3.2-3.3).

The GEM-SA parameterization requires seven input parameters: D_{pm0}, σ_0, mass flux, fire area, wind speed, mixing depth and time, and generates predicted aged D_{pm} and σ as outputs. These estimated D_{pm} and σ describe an aged lognormal aerosol mode incorporating the sub-grid scale coagulation taking place inside concentrated biomass-burning plumes and can be used in global/regional models. We have made the GEM-SA parameterization (emulator Fortran subroutine and input files) available as Supplementary Material.

3. Results

3.1 SAM-TOMAS simulation output

Figure 4 shows the D_{pm} (panels a and c) and σ (panels b and d) as a function of distance for each of the 100 SAM-TOMAS simulations used to train the emulator (Sect. 3.2). The influence of several factors

(the distance from the source, emissions mass flux, and fire area) on the final aerosol size distributions is apparent in the output of SAM-TOMAS simulations. Panels a and b are colored by the emissions mass flux, whereas panels c and d are colored by dM/dxdz (kg m^{-2}, the amount of aerosol mass in an infinitesimally thin slice of air perpendicular to the direction of the wind, i.e. mass flux · fire area / wind speed/mixing depth). All simulations showed D_{pm} increasing with distance as coagulation progressed in each plume. The coloring in panel a shows that D_{pm} generally increases more rapidly and to higher values with higher emission fluxes. However, panel c shows that dM/dxdz appears to be a better predictor for the increase of D_{pm} with distance than the emissions flux, and the distance and dM/dxdz capture much of the variability in D_{pm}. ~~D_{pm0} appears to have little influence on D_{pm} (note, however, that the first points on these plots already include some processing and are not the initial diameters).~~

Panels b and d show that σ tends to converge with distance as simulations with large initial σ generally decrease with distance more rapidly than simulations with smaller initial σ. This convergence happens slowly relative to the times simulated, so the initial σ have a strong influence even at 200 km. The colors and panels b and d show that σ in high emissions-flux and dM/dxdz cases converge more rapidly than low-emissions cases. However, as opposed to the 1.32 modal-width asymptote in the limit of infinite coagulation found by Lee (1983), the SAM-TOMAS simulations converge to a limit of 1.2-1.25. This is likely due to the size-distribution bin-spacing in the SAM-TOMAS model, where modal widths <1.32 are smaller than a single TOMAS size bin width, which results in less accurate fits of σ for smaller σ values.

Figure 5 is a scatterplot of ~~D_{pm} vs~~ σ vs D_{pm} for each point seen in Fig. 4, excepting those at distances less than 25 km (points close to the emissions source have been removed). The points are colored by dM/dxdz. Thus, Fig. 5 shows the results of Fig. 4 panels c and d together but removes the distance information. At these distances over 25 km, D_{pm} is relatively well constrained by dM/dxdz alone, showing that the mean growth by coagulation is strongly influenced by the mass of particles in the slice of air. On the other hand, σ is unconstrained at low values of dM/dxdz but more constrained towards 1.2-1.4 at high values of dM/dxdz. At high dM/dxdz values, the convergence towards the steady-state σ proceeds much more rapidly than at low dM/dxdz as also shown in Fig. 4d.

These SAM-TOMAS results show that dM/dxdz is a powerful determinant of aged biomass-burning size. In these tests, we also explored the suitability of dM/dx (mass flux · fire area/ wind speed) and dM/dV (initial mass concentration). Large mixing depths dilute particle concentrations and reduce

coagulation, so we expected that dM/dxdz may be a better predictor of biomass-burning size-distribution aging than dM/dx. However, Fig. 4 and Fig. 5 did not look qualitatively different when using dM/dx or dM/dV. A comparison of dM/dx vs dM/dxdz vs dM/dV in predicting final size-distribution attributes is further discussed in Section 3.4. We quantitatively evaluate the fidelity of dM/dx and dM/dxdz as proxies for biomass-burning size-distribution aging in Sect. 3.4. In the following two subsections, we use the emulator to determine the contribution of the individual inputs to the changes in simulated D_{pm} and σ.

3.2 Model parameterization evaluation

We tested the GEM-SA-derived emulator parameterization against additional SAM-TOMAS model runs that were not used in the fitting of the parameterization, and we show the results in Fig. 6. We use 624 additional SAM-TOMAS-simulated data points that were not used for GEM-SA training in this evaluation. The emulator parameterization-predicted outputs corresponding to these data points for D_{pm} and σ are plotted against the SAM-TOMAS D_{pm} and σ. Predicted D_{pm} has an R^2 value of 0.83 with a slope of 0.92. Larger absolute errors in D_{pm} are found at the larger diameter sizes, but 86% are found within 10% of the SAM-TOMAS D_{pm} (76% of predicted D_{pm} are within 5% of SAM-TOMAS D_{pm}). The small mean normalized bias (MNB) of -0.06 indicates a slight negative bias in the parameterization. This bias is generally seen towards the higher final D_{pm} values in the simulations (>250 nm), which are reached only by the most aged plumes with the heaviest aerosol loads. The σ plot (Fig. 6b) shows a similar correlation coefficient (R^2=0.91) and has a slope of 0.93. The MNB is 0.01 and 77% of the predicted σ points are within 5% of the σ calculated from SAM-TOMAS. The cluster of points near σ =1.2-1.3 is indicative of the modal width steady-state limit. This limit is not captured by the σ parameterization, which assumes a smooth function towards even lower σ values.

3.3 Sensitivity of aged size distribution to input parameters

Figures 7 and 8 show the sensitivities of the parameterization outputs (D_{pm} and σ, respectively) to the input parameters (D_{pm0}, σ_0, mass flux, fire area, wind speed, time, and mixing depth) as determined by the GEM-SA emulation of the SAM-TOMAS output. (Note that distance was used as the dependent variable in Fig. 4, while we use time in the emulator. Time can be converted to distance by multiplying by the wind speed). In every panel, each line shows the change in D_{pm} (Fig. 7) or σ (Fig. 8) as an input parameter (e.g. D_{pm0} in panel a) is varied systematically from its minimum to maximum tested value

with a randomly chosen set of the other six input parameters. Each panel contains 100 lines, which means that 100 sets of the six other input parameters were randomly chosen to make these lines. We normalize each line by the value of D_{pm} or σ at the midpoint of the x-axis (i.e. where the input parameter is at the midpoint of its tested range). For time since emission (panel f) we normalize by the values at t=0 min instead of at the midpoint of the range. These plots therefore show the percent change in D_{pm} or σ, $\Delta\%_{output}$, as each input is changed from its midpoint value (or t=0 min for time), in order to emphasize the parameterization's output response to each isolated input variable.

The D_{pm} sensitivity plots (Fig. 7) show a number of well-defined responses of D_{pm} to the inputs. D_{pm} increases monotonically with increases in mass flux and fire area (Fig. 5b,d), and decreases nearly monotonically with wind speed. These trends are due to the interrelationships of these inputs with starting number concentration. These results are consistent with Fig. 4 and Fig. 5, where D_{pm} increased with increasing dM/dx in the SAM-TOMAS simulations. Additionally, the D_{pm} also decreases monotonically with mixing depth (albeit more weakly than mass flux, fire area, and wind speed), so dM/dxdz may also be a good proxy for biomass-burning size-distribution aging (evaluated in Sect. 3.4). Higher dM/dx and dM/dxdz values lead to higher initial number concentration in these plumes, which drive higher rates of coagulation due the squared dependence of coagulation rate on number concentrations.

D_{pm} also increases nearly monotonically with time (the regions of slight decreases with time show that the parameterization is not necessarily always physically representative due to the statistical nature of the fit over the parameter space). The rapid rise in D_{pm} for time <2 hrs is due to the high number concentrations (N) and coagulation rates near the source. As dilution and coagulation progress, N decreases and coagulation slows, resulting in a slowing of D_{pm} increase. Mass flux has the largest range of output D_{pm} associated with the input ranges specified here (~ -50% to +100%).

The relationship between D_{pm} and the initial size parameters (D_{pm0} and σ_0) is more complicated. Neither D_{pm0} nor σ_0 show monotonic increases or decreases in D_{pm} due to changes in either of these isolated inputs. In general, there is an increasing trend in output D_{pm} with increasing D_{pm0}, but for some cases it decreases. These decreases in D_{pm} are likely due to (1) decreasing particle number concentrations with increasing D_{pm0}, which leads to reduced coagulation rates and (2) imperfections in the statistical fit of the parameter space. The larger σ_0 indicate broader emission size distributions, with more large particles and small particles. Since coagulation progresses fastest between large and small particles (as opposed to particles of approximately the same size), this favors higher D_{pm} at higher σ.

However, the initial particle number decreases with increasing σ, which lowers the coagulation rate and leads to lower D_{pm}.

The emulator-derived σ sensitivities are shown in Fig. 8. Since we expect σ to converge towards an asymptotic limit with coagulational processing (Fig. 4b,d), we see with those input parameters associated with higher plume number density (mass flux, fire area, wind speed^{-1}, mixing depth^{-1}), which gave monotonic increases for D_{pm}, show mixed results for σ due to variability in the initial σ_0. The time sensitivity plot (Fig. 8f) shows decreasing σ with time similar to Fig. 4b,d.

Emission σ_0 shows the most pronounced and largest magnitude effect on output σ (~ -30% to +30%). Thus, the timescales for σ evolving towards 1.2 is longer than the timescales tested here for even the densest plumes. These sensitivity plots show that there is less variability in σ than in D_{pm} over the tested input space.

3.4 Simplified fits to the aged size distributions

In addition to the GEM-SA emulator fits, we determined simplified fits for both D_{pm} and σ based on the behavior in Fig. 4 and Fig. 5. These fits are easier to implement in regional and global aerosol models than the full GEM-derived parameterization. These equations are meant to produce approximate estimates of D_{pm} and σ throughout plume size-distribution aging. The equations require: the initial value of the size-parameter of interest (D_{pm0} or σ_0), a value proportional to the plume aerosol loading (dM/dxdz: mass flux · fire area / wind speed / mixing depth or dM/dx: mass flux · fire area / wind speed), and time since emission from the source fire (time). (Distance may also be used in these equations rather than time, and distance/wind-speed should be used in place of time.) The functional forms fitted for D_{pm} and σ are found below.

$$D_{pm} = D_{pm0} + A\,[dM/dx]^b\,(time)^c \quad (1)$$

$$D_{pm} = D_{pm0} + A\,[dM/dxdz]^b\,(time)^c \quad (2)$$

$$\sigma = \sigma_0 + A\,[dM/dx]^b\,(time)^c\,(1.2 - \sigma_0) \quad (3)$$

$$\sigma = \sigma_0 + A\,[dM/dxdz]^b\,(time)^c\,(1.2 - \sigma_0) \quad (4)$$

where A, b and c are determined by fitting each equation to the SAM-TOMAS data. For these empirical equations, the units of dM/dx are kg m^{-1}, dM/dxdz are kg m^{-2}, D_{pm} is nm and time since emission is min. It should be noted that the equations for D_{pm} and σ are designed to be independent of each other (i.e. D_{pm} is not dependent on σ_0), which differs from the GEM-SA emulator. The aerosol loading parameter dM/dx was chosen based on the stratification seen in Fig. 4c and Fig. 5. dM/dxdz was tested as well, as it incorporates the variance associated with mixing depth into the fit. The fit to dM/dx rather than dM/dxdz may be advantageous because we expect mixing depth of the plume to be one of the more uncertain parameters in an atmospheric model, and the D_{pm} sensitivities to mixing depth tend to be smaller than those to mass flux, fire area and wind speed in the GEM-SA emulator (Fig. 7). The σ fits introduce a fourth factor, $(1.2-\sigma_0)$, which represents the difference between the SAM-TOMAS infinite-coagulation limit (Fig. 4b and d) and the initial modal width.

The scalar A, b and c variables were fit to the ensemble of SAM-TOMAS data. Their values are summarized in Table 3. The fits were tested against independent SAM-TOMAS data in Fig. 9 (D_{pm}) and Fig. 10 (σ). The simplified D_{pm} parameterizations, as expected, are not as good a fit of the SAM-TOMAS data as the GEM-SA emulator (Fig. 6). The fit statistics for the simple parameterizations are as follows: D_{pm}(dM/dx): slope = 0.82, R^2 = 0.67, MNB= 0.003, D_{pm} (dM/dxdz): slope = 0.98, R^2 = 0.77, MNB= 0.008. The fit using dM/dxdz generally performs better than that with dM/dx. The simple σ fit also did not perform as well as the GEM-SA emulator with fit statistics of: σ(dM/dx): slope = 0.64, R^2= 0.78, MNB= 0.02 and, σ(dM/dxdz): slope = 0.65, R^2 = 0.79, MNB= 0.01). Thus, dM/dxdz fits do yield better results than dM/dx (in particular for D_{pm}); however, a user may choose to use the dM/dx fit if the mixing depth is unknown. We note that these fits are only valid within the parameter ranges shown in Table 1. dM/dV was also tested as a parameter within these simplified parameterization, but did not yield better agreements for either D_{pm} or σ than dM/dxdz despite incorporating an additional plume parameter (initial plume y-extent). This is because dM/dxdz is the product of dM/dV and the initial plume width; since wider plumes are less susceptible to dilution than narrower plumes, dM/dxdz captures this plume-width effect while dM/dV does not.

3.5 OA production/loss

One of the limitations of the coagulation-only parameterizations derived in this paper is that they do not include the effects of potential condensation/evaporation of organic aerosol on the aged biomass-burning size distribution. Both condensational growth and evaporative loss of OA has been observed

406 previously in chamber studies and the field due to OA production or evaporation from
407 dilution/chemistry (Cubison et al., 2011; Hecobian et al., 2011; Hennigan et al., 2011; Grieshop et al.,
408 2009; Ortega et al., 2013; Jolleys et al., 2015; Vakkari et al., 2014). Konovalov et al. (2015) has
409 emphasized the importance of OA simulation in modeling long-range (>1000 km) plume evolution.
410 Thus, in order to predict biomass-burning aerosol mass, and thus the aerosol size distribution, we must
411 understand how OA evolves in biomass-burning plumes.

412 Here we present a simple correction to our coagulation-only parameterizations to account for
413 in-plume OA production/loss, assuming that this production/loss is known. This correction assumes all
414 SOA condenses onto existing particles (no new-particle formation). Each parameterization presented in
415 this paper may be corrected to include OA production/evaporation using the corrections below. We
416 assume that the OA production or loss does not affect the coagulation rates or σ, but acts to increase the
417 final D_{pm}. These assumptions are imperfect as irreversible condensation (evaporation) ~~in~~decreases
418 (in~~de~~creases) σ; however, σ is preserved during condensation or evaporation of semi-volatile material
419 (Pierce et al., 2011). Regardless, for the relatively small amounts of OA condensation/evaporation
420 considered here, the change in σ and coagulation rates should be minor. For a factor of 25% growth in
421 diameter from SOA, which may be expected from for a factor of 2 increase in OA mass with a small
422 change in sigma, we expect coagulation rates to stay within about 10% (Seinfeld and Pandis, 2006).
423 For larger changes in OA mass (more than a factor of ~2) due to production/loss, our simple correction
424 will have uncertainties due to these assumptions. Our correction to the final D_{pm} has the following
425 form:

$$D_{pm\,w/OA\,\mathrm{prod/loss}} = D_{pm\,w/o\,OA\,\mathrm{prod/loss}} \cdot \left(\frac{OAMass_{w/OA\,\mathrm{prod/loss}} + BCMass}{OAMass_{w/o\,OA\,\mathrm{prod/loss}} + BCMass} \right)^{1/3} \quad (5)$$

426

427 where $D_{pm\ w/o\ OA\ prod/loss}$ is the final D_{pm} from the coagulation-only GEM-SA emulator parameterization,
428 the biomass-burning aerosol OA mass (with and without additional production or loss) is in kg (per
429 particle or volume of air) and the BC mass is in kg (per particle or volume of air). Thus, for a doubling
430 of OA due to SOA production (one of the larger enhancements found in Hennigan et al., 2011),
431 particles that contain negligible BC will grow in diameter by 26% above the coagulation-only
432 predictions. If the particles contained 50% BC, then the diameter growth would only be 14%.

433 While these changes are expected to be on the large end for growth by SOA production, they

are significantly smaller than the ~200% variability in aged D_{pm} due to coagulation over the range of initial fire conditions (Fig. 7). For example, variations in wind speed, mass flux, and fire area alone can independently cause variability in the aged D_{pm} by a factor of 2 due to changes in coagulation rates while variability in condensational growth appears to cause much smaller uncertainties (~25%) in the aged D_{pm}. This indicates that although SOA condensational growth is certainly important in shaping particle composition and total particle mass, it is not among the most dominant factors determining the aged D_{pm} compared to those fire-condition parameters controlling coagulational growth. It should be noted, however, that the D_{pm} growth attributed to OA condensation is not accompanied by a change in particle number (additional OA mass is distributed among existing particles), whereas a similar increase in D_{pm} growth by coagulation only would have an accompanying decrease in particle number. Thus, the changes to the aerosol size distribution and climatic influence of a size change due to coagulation and condensation are different.

3.6 Estimating aged size distributions observed at the Mt. Bachelor Observatory

The simplified fits presented in Section 3.4 (equations 1-4) were tested against size distributions measurements of biomass-burning plumes observed at the Mt. Bachelor Observatory (MBO) in Central Oregon (43.98°N, 121.69°W, 2,764 m a.s.l.). MBO is a mountaintop site that has been in operation since 2004 (Jaffe et al., 2005). An intensive campaign was performed during the summer of 2015 to measure aerosol physical and optical properties of wildfire emissions (Laing et al., in prep). During this campaign aerosol size distributions from 14.1 to 637.8 nm were measured with a Scanning Mobility Particle Sizer (SMPS). Additional details about MBO and the sampling campaign can be found in Laing et al. (in prep).

We identified eleven biomass-burning plumes during August (Table 4). Criteria for plume selection was aerosol scattering > 20 Mm^{-1} and CO > 150 ppbv for at least an hour, a strong correlation (R^2 > 0.80) between aerosol scattering and CO, and consistent backward trajectories indicating transport over known fire locations. We calculated back-trajectories to determine fire locations using the National Oceanic and Atmospheric Administration Hybrid Single-Particle Lagrangian Integrated Trajectory (HYSPLIT) model, version 4 (Draxler, 1999; Draxler and Hess, 1997, 1998; Stein et al., 2015) with Global Data Assimilation System (GDAS, 1º x 1º) data. The Mt. Bachelor summit is located at ~1500 m amgl (above model ground level), so the back-trajectory starting heights of 1300, 1500, and 1700 m amgl were chosen (Ambrose et al., 2011). Fire locations were identified using Moderate

Resolution Imaging Spectroradiometer (MODIS) satellite-derived active fire counts (http://activefiremaps.fs.fed.us/; Justice et al., 2002).

For the plume aerosol loading parameterization inputs in equations 1-4, we used Fire INventory from NCAR (FINN) daily-averaged fire area and fire-emissions estimates (Wiedinmyer et al. 2011). Multiple FINN data points in the same vicinity were combined based on the location of large-wildfire incidents tracked by the National Interagency Fire Center (NIFC) (http://activefiremaps.fs.fed.us/). We calculated the mass flux for the aerosol-loading estimates (dM/dxdz and dM/dx) using these FINN OC+BC emissions (kg/day) and FINN fire area data (km^2). Mixing depth was defined as the mixing depth at the source location of the fire in the Global Data Assimilation System (GDAS, 1º x 1º) data. Wind speed was also extracted from GDAS data and was calculated as the average wind speed from the ground to the defined mixing height. If no data were available, the mixing height and wind speed were set to 660 m and 8.5 m/s based on the median value of the rest of the plumes. We assumed the emission diameter (D_{pm0}) to be 100 nm, and we calculated σ using initial modal widths (σ_0) of 1.6, 1.9 and 2.4, to be discussed later. We estimated the transport time from plume back-trajectories, and these values ranged from 4.5 to 35 hours.

The measured and calculated size distribution diameter and modal widths for each plume at MBO are summarized in Table 4. We calculated D_{pm} and σ as the geometric mean diameter and geometric standard deviation of the plume averaged size distribution as measured by the SMPS, respectively. The plume-averaged size distributions may be influenced by non-biomass--burning particles included along the trajectory from the wildfire. Plumes 1, 2, and 4 have bimodal distributions. The second mode (Aitken mode) of these distributions are an example of influence from a non-biomass burning source. These three bimodal distributions have inflated σ values, which will be addressed later. Due to the large number of fires in Northern California and Oregon during the summer of 2015, some of the plumes observed at MBO were influenced by more than one fire (e.g. Figure 9). For these plumes, we calculated aged D_{pm} and σ values for each fire area (black squares in Figure 9) and a weighted average based on aerosol loading (dM/dx or dM/dxdz) was taken. Column 3 in Table 4 indicates how many fire areas were averaged for each plume.

Figure 10 shows the predicted aged D_{pm} plotted against the observed values for both the dM/dx and dM/dxdz forms of the simple parameterization. An initial D_{pm0} of 100 nm was assumed. ~~used~~Equation 2 (using aerosol mass loading dM/dxdz) estimates D_{pm} somewhat more accurately ($y = 0.93x + 17.1$, $R^2 = 0.551$) than Eqn. 1, which uses aerosol mass loading dM/dx ($y = 0.62x + 53.1$, $R^2 =$

0.532). Over half of the variability in the observed D_{pm} was captured by the simplified fits. Thus, the simple parameterizations show skill at predicting the aged D_{pm} values relative to choosing a constant value of aged D_{pm} as is typically done in regional and global models.

Figure 11 shows the predicted aged σ plotted against the observed values for both parameterization forms. Both parameterizations do not predict modal width as well as D_{pm} (Figure 11). The calculated modal width changed significantly when using different emission modal-width values (σ_0). Janhäll et al. (2010) found the σ of fresh biomass burning emissions to range from ~1.6 to 1.9. When using a σ_0 of 1.6, we underestimated all of the σ values. Using a σ_0 of 1.9, we improved the estimation of aged σ ranging from 1.4-1.6 (Figure 11a). The three higher measured σ values are from the bimodal plumes mentioned previously, which have larger σ values than would be due strictly to the biomass--burning plume. We found that using a σ_0 of 2.4 provided the best fit for all of the measured plumes (Figure 11b), 2.4 being the max σ_0 value from Table 2. The σ simplified fits using $\sigma_0 = 2.4$ have statistics of: σ(dM/dx): $y = 0.50 + 1.00$, $R^2 = 0.513$, and σ(dM/dxdz): $y = 0.57 + 0.77$, $R^2 = 0.468$. Thus, both parameterizations do not predict modal width as well as D_{pm}; however, these parameterizations do show skill relative to assuming a constant value of σ.

The results from the regional fires demonstrate that the parameterizations in Eqs 1-4 can be successfully used to estimate aged biomass-burning size distributions in regional biomass-burning plumes with transport times up to 35 hours with significantly better skill than assuming fixed values for size-distribution parameters. More investigations of individual aged biomass- burning plumes, specifically with one clear source, should be completed to fully characterize this parameterization.

4. Conclusions

We used the SAM-TOMAS large-eddy simulation model and an emulation technique to explore the evolution of biomass-burning aerosol size distributions due to coagulation and build coagulation-only parameterizations of this size-distribution evolution. We have also provided a simple correction to the parameterization for cases with net OA production or loss. We used the SAM-TOMAS model to simulate plume dispersion and aerosol coagulation. The SAM-TOMAS results show that the aged D_{pm} can be largely described by dM/dx and the distance from the source (or time since emission). These

results also show that the aged σ moves from σ_0 towards a value of 1.2 at a rate that depends on dM/dx.

The GEM-SA program was used to derive a D_{pm} and σ emulator parameterization based on the SAM-TOMAS results. The parameterization requires seven input parameters: emission D_{pm0}, emission σ_0, mass flux, boundary layer wind speed, fire area, plume mixing depth, and time since emission. The predicted D_{pm} and σ can then be used as effective unimodal biomass-burning size-distribution parameters in regional and global aerosol models.

The D_{pm} parameterization showed the strongest sensitivities to those input parameters associated with the extent of aerosol loading within the plume (mass flux, fire area, wind speed). Across the fire area and wind speed ranges tested here, final D_{pm} varied by ± 50%. Mass flux had the largest associated D_{pm} sensitivity across the tested values (-50% to +100%). These sensitivities were larger than those associated with mixing depth (~ -20% to 20%) or the initial size-distribution parameters (D_{pm0}: ~ -25% to 25%, σ_0: ~ 15% to -15%). The σ parameterization showed a uniform decrease in σ with time and strong sensitivities to the emission σ_0 (-30% to 30%). This strong sensitivity to σ_0 can be attributed to the inertia in σ evolution in simulations with large modal widths and relatively small mass loading, where σ will not converge quickly to the coagulational limit (1.2).

The GEM-SA-derived parameterization performed relatively well against the SAM-TOMAS model with a correlation of R^2=0.83, slope of m=0.92 and a low mean normalized bias of MNB=-0.06 for D_{pm}. The σ parameterization has fit statistics of R^2= 0.93, slope= 0.91 and MNB= 0.01. The σ parameterization was unable to capture the coagulational limit of 1.2 seen in the SAM-TOMAS results and instead extrapolated to lower values. This 1.2 limit differs from the 1.32 σ limit proposed by Lee (1983) due to the bin-spacing in SAM-TOMAS being coarser than lognormal modes with these small modal widths.

We also provided simplified polynomial fits for D_{pm} and σ (Eqns 1-4, Table 3) for calculating aged D_{pm} and σ as independent functions of: the fresh emission parameter (D_{pm0} or σ_0), the mass loading of the aerosol (dM/dx or dM/dxdz) and the time since emission from the source fire. The σ fits also require a convergence term to account for the coagulational limit (1.2 in the SAM-TOMAS model). Tested against independent SAM-TOMAS data, the D_{pm} simplified fits performed as: D_{pm}(dM/dx): slope = 0.82, R^2 = 0.67, MNB= 0.003 and D_{pm} (dM/dxdz): slope = 0.98, R^2 = 0.77, MNB= 0.008. The σ simplified fits have statistics of σ(dM/dx): slope = 0.64, R^2= 0.78, MNB= 0.02 and σ(dM/dxdz): slope = 0.65, R^2 = 0.79, MNB= 0.01. The equations requiring (dM/dxdz) performed better than their (dM/dx)

553 counterparts as they also account for the aerosol layer depth.

554 We provided a correction for OA production/loss, and showed that significant production of
555 SOA within the plume (~ 100% OA mass enhancement) would cause a relatively small shift in the size-
556 distribution D_{pm} (14-26% increase) compared to other factors that control the coagulation rate (e.g.
557 dM/dx). We note, however, that OA production increases D_{pm} without loss of particle number while
558 coagulation increases D_{pm} with a decrease in number, thus the climatic impact of condensation and
559 coagulation are different. The simplified OA-production/loss correction assumes no change in σ with
560 condensational growth. Further testing should be done with explicit OA production and loss to better
561 quantify the effects of condensation of the size-distribution evolution.

562 We tested the simplified fits for D_{pm} and σ (Eqns 1-4, Table 3) against 11 aged biomass-burning
563 plumes observed at the Mt. Bachelor Observatory in August of 2015. D_{pm} was reasonably calculated
564 using both measures of aerosol loading, dM/dx and dM/dxdz (R^2 values above 0.7 without an outlier).
565 The fit of calculated σ and measured σ depended heavily on the assumed initial modal width, with an
566 assumed σ_0 of 2.4 working best in our case (R^2 values around 0.75 without an outlier). Despite the
567 changes in calculated D_{pm} and σ due to the estimated emission size distribution, the parameterizations
568 captured the differences from plume to plume in regional biomass-burning plumes, which is based on
569 estimated aerosol loading and transport times.

570 Our analysis does not include any cloud processing of the plume particles, i.e. the production of
571 aqueous SOA within activated plume particles is not accounted for in our simple OA mass correction.
572 The production of SOA within droplets could result in additional SOA mass being only added to the
573 larger, activated particles during activation/evaporation cycling. This extra SOA mass would favor
574 increases in the diameters of the larger particles of the size-distribution only, which could create a
575 bimodal size distribution and increase the overall coagulational rates in the plume (more, larger
576 particles coagulate more rapidly with the small-diameter particles).

577 Future work includes (1) more testing of the parameterizations ~~developed in this work~~ against
578 real world observations of size distribution aging, and (2) incorporating the parameterizations into
579 regional and global aerosol models for further evaluation against regional/global measurements.

580

5. Author Contribution

K.M. Sakamoto, R.G. Stevens, and J.R. Pierce designed the study. K. M. Sakamoto performed the SAM-TOMAS simulations, and created and evaluated the parameterizations. J.R. Laing tested the parameterizations size distributions of aged biomass burning plumes observed at the Mt. Bachelor Observatory, and D.A. Jaffe oversaw the Mt. Bachelor measurements. K. M. Sakamoto prepared the manuscript with assistance from all co-authors.

6. Acknowledgements

NCEP Reanalysis data provided by the NOAA/OAR/ESRL PSD, Boulder, Colorado, USA, from their Web site at http://www.esrl.noaa.gov/psd/. K.M. Sakamoto was funded by a Natural Sciences and Engineering Research Council of Canada (NSERC) PGS-M Fellowship. The authors gratefully acknowledge the NOAA Air Resources Laboratory (ARL) for the provision of the HYSPLIT transport model used in this publication

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Table 1. Parameter ranges for each of the seven input parameters investigated in this study.

Parameter	Description	Units	Min. Value	Max. Value
D_{pm0}	Emission median dry diameter	nm	20	100
σ_0	Emission modal width	-	1.2	2.4
Mass Flux	Emission mass flux from fire	kg m^{-2} s^{-1}	$2x10^{-8}$	$5x10^{-6}$
Fire area	Square fire emissions area	km^2	1	49
Wind speed	Mean boundary-layer wind speed	m s^{-1}	2	20
Mixing depth	Mixing depth of aerosol layer	m	**150**	2500
Time	Time since emission	min	0	300

Table 2. Parameter ranges for inputs to the SAM-TOMAS model.

Parameter	Description	Units	Min. value	Max. value
Date	Req. for Met. field selection	8-hour	July 1, 2010	July 31, 2010
Latitude		deg N	30	70
Longitude		deg W	70	135
D_{pm0}	Emission median dry diameter	nm	20	100
σ_0	Emission modal width	-	1.2	2.4
Mass Flux	Emission mass flux from fire	kg m^{-2} s^{-1}	$2x10^{-8}$	$5x10^{-6}$
Fire area	Square fire emissions area	km^2	1	49
Injection height	Lower plume injection bound	m	50	150
Injection depth	Depth of plume at emission	m	500	2000

Table 3. Best-fit parameters for the simplified D_{pm} and σ SAM-TOMAS parameterizations (Eqns. 1 to 4)

Fit	Eqn. #	Parameter		
		A	b	c
$D_{pm}(dM/dx)$	(1)	4.268	0.3854	0.4915
$D_{pm}(dM/dxdz)$	(2)	84.58	0.4191	0.4870
σ(dM/dx)	(3)	0.05940	0.1915	0.3569
σ(dM/dxdz)	(4)	0.2390	0.1889	0.3540

Table 4: Measured and calculated D_{pm} and σ of biomass-burning plumes observed at MBO during August 2015. For the calculated D_{pm} and σ of, the initial size parameters used were D_{pm0} = 100 nm and σ_0 = 1.9.

Plume	Plume date and time (UTC)	# fire areas	Measured (SMPS)		Calculated using dM/dx		Calculated using dM/dxdz	
			D_{pm} (nm)	σ	D_{pm} (nm)	σ	D_{pm} (nm)	σ
1	8/9/2015 3:00-4:00	3	136.1	1.95	140.7	1.64	151.1	1.59
2	8/9/2015 5:00-7:00	3	144.0	1.77	140.8	1.64	152.0	1.58
3	8/10/2015 3:00-5:00	3	190.1	1.50	140.9	1.63	149.7	1.58
4	8/23/2015 3:55-7:00	1	162.5	1.89	145.5	1.63	162.4	1.57
5	8/24/2015 4:00-7:25	1	201.1	1.59	167.5	1.55	184.7	1.49
6	8/24/2015 7:30-11:20	1	217.5	1.52	190.1	1.50	230.1	1.40
7	8/24/2015 13:00-18:00	1	212.5	1.49	193.9	1.48	237.8	1.37
8	8/25/2015 3:50-6:50	1	192.2	1.54	161.4	1.57	172.6	1.52
9	8/27/2015 9:00-13:00	3	192.9	1.50	194.2	1.49	220.6	1.43
10	8/28/2015 8:00-11:15	3	183.4	1.54	182.1	1.50	203.2	1.43
11	8/28/2015 17:40-19:40	3	176.7	1.60	181.4	1.50	202.0	1.43

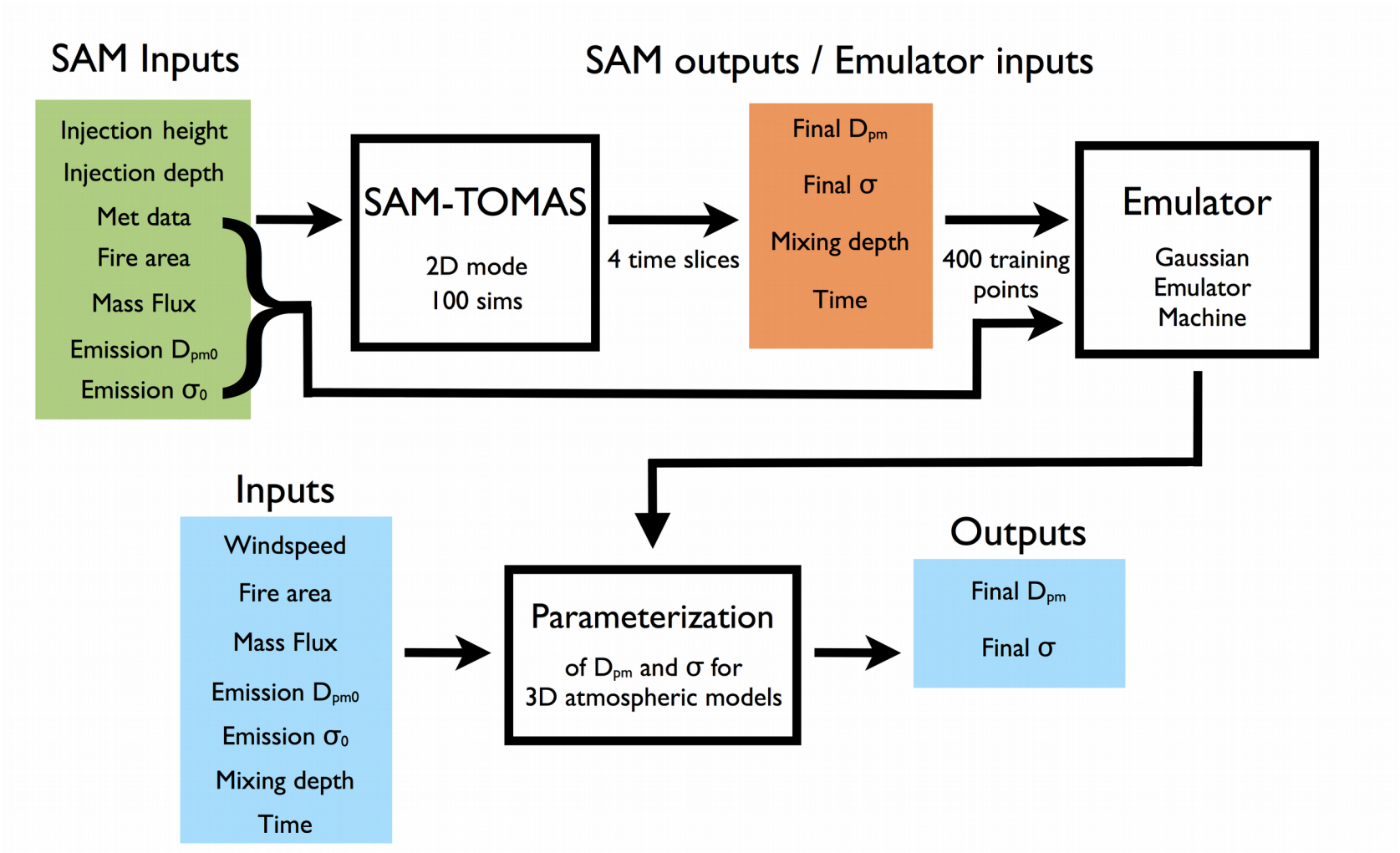

Figure 1. Schematic of the methods in this paper.

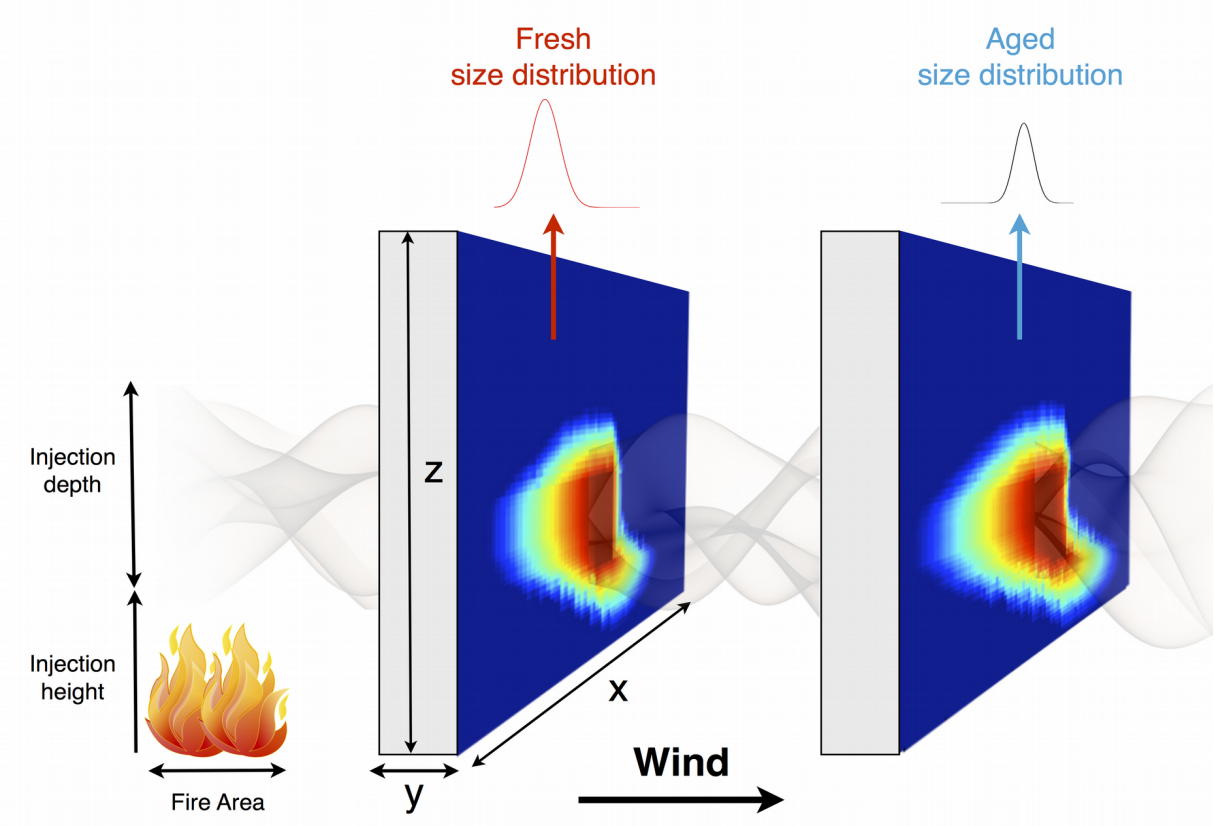

Figure 2. Schematic of a 2D SAM-TOMAS plume simulation.

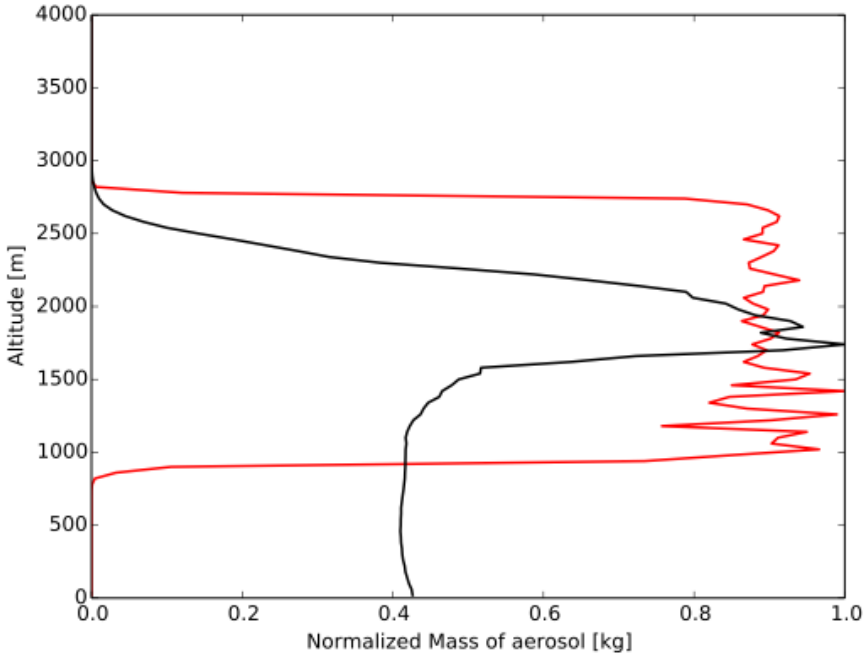

Figure 3. Final vertical profiles for two representative SAM-TOMAS simulations after four hours, normalized to individual aerosol load and averaged horizontally across the domain. The black profile shows a simulation where the aerosol mixed through the boundary layer to the ground with some aersol still trapped in a stable emission layer**, while the red profile shows a simulation where the aerosol plume is still stable at the emission injection layer.**

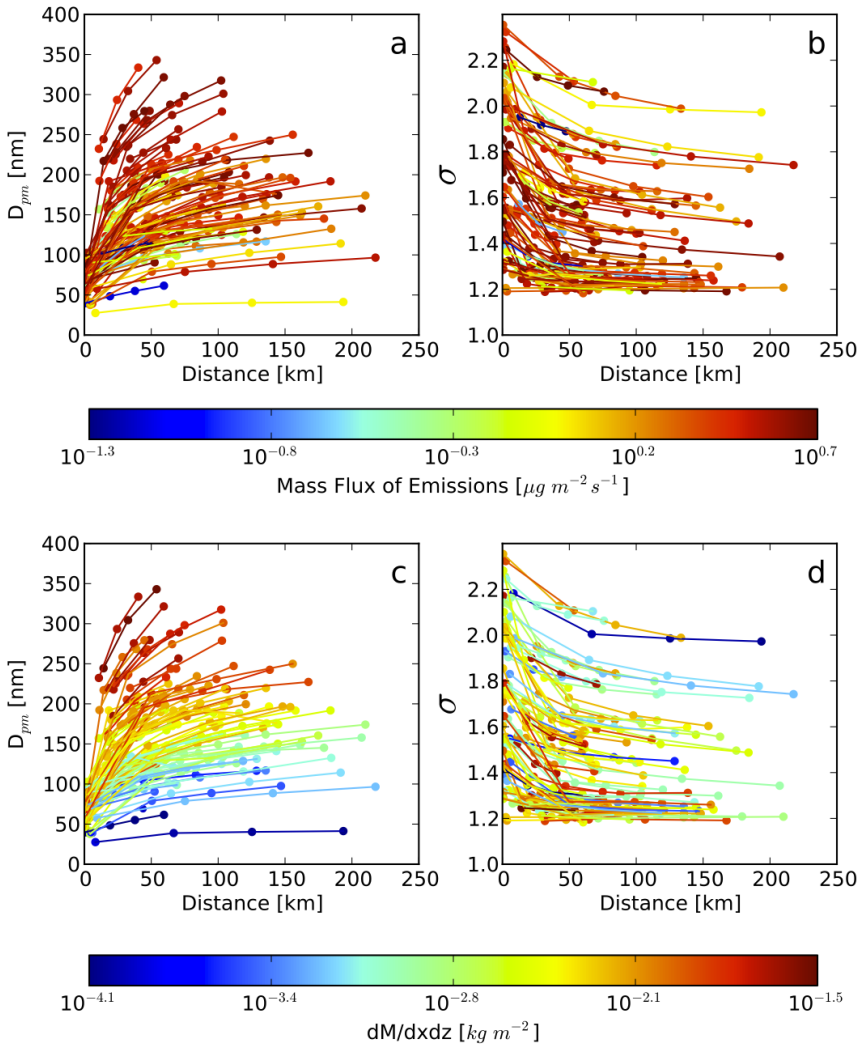

Figure 4. Wire plots showing size-distribution changes across individual SAM-TOMAS simulations colored by emission mass flux (panels a and b) and dM/dx**dz** (panels c and d) for D_{pm} (panels a and c) and σ (panels b and d).

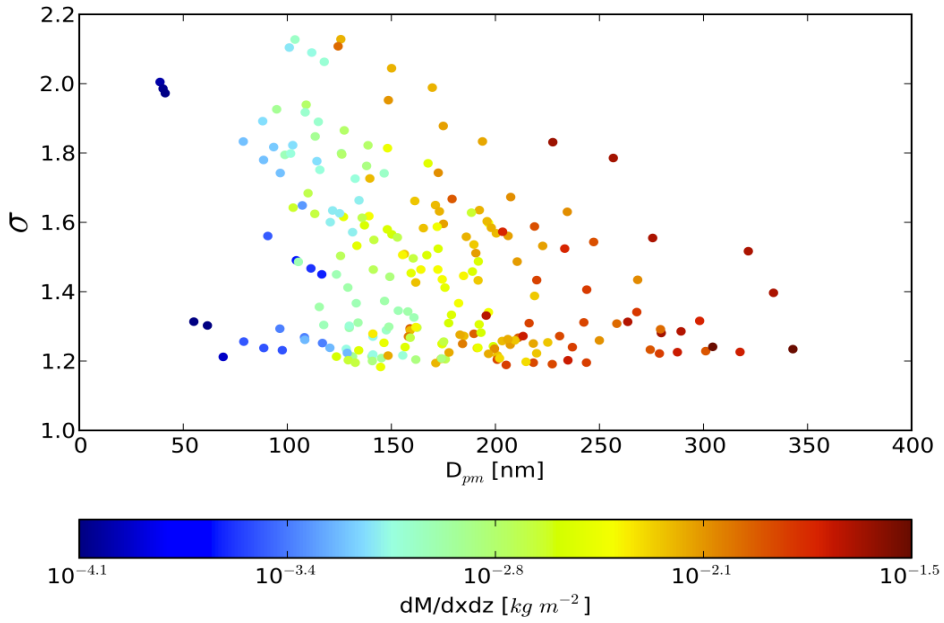

Figure 5. Scatter plot showing the relationships between final modal width (σ), final D_{pm}, and dM/dx**dz** for each of the SAM-TOMAS simulation slices at distances greater than 25 km from the fire.

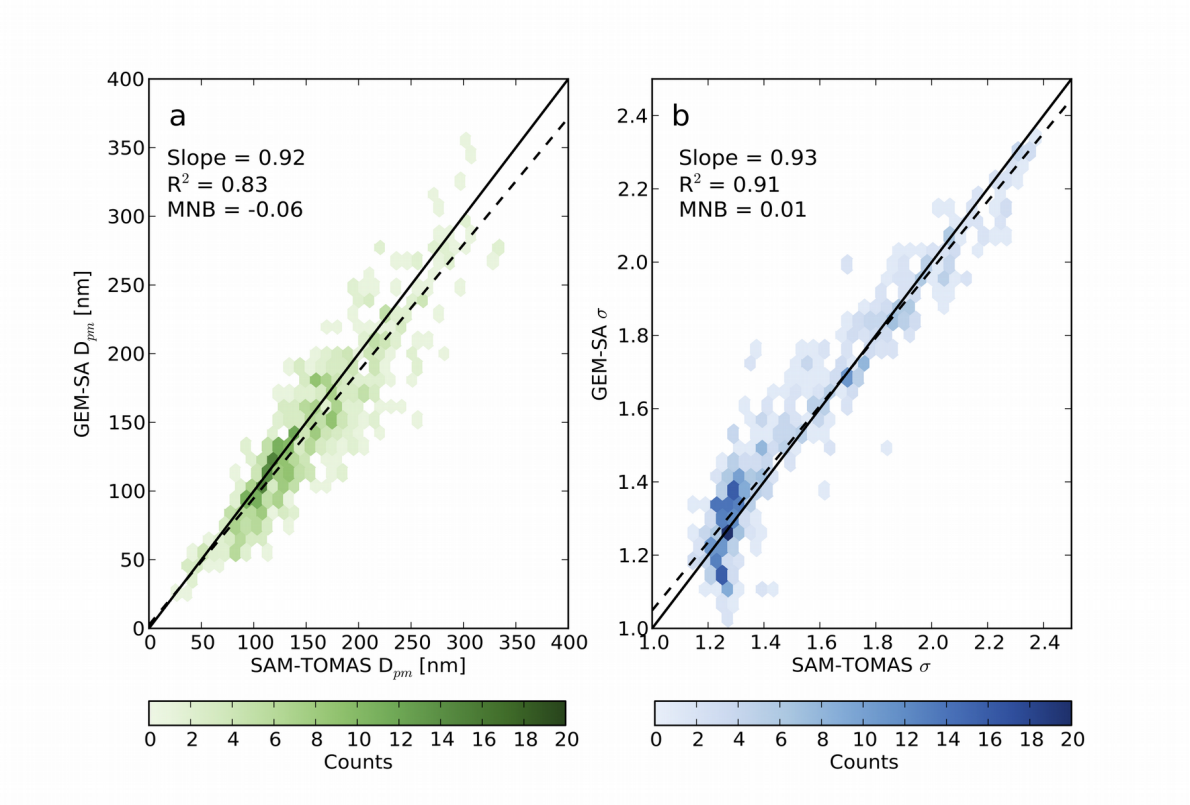

Figure 6. One-to-one plots showing GEM-SA emulator vs. SAM-TOMAS for 624 non-training simulation slices for a) final D_{pm}, and b) final modal width, σ. The black line is the one-to-one line. The dashed black line is the line of best fit.

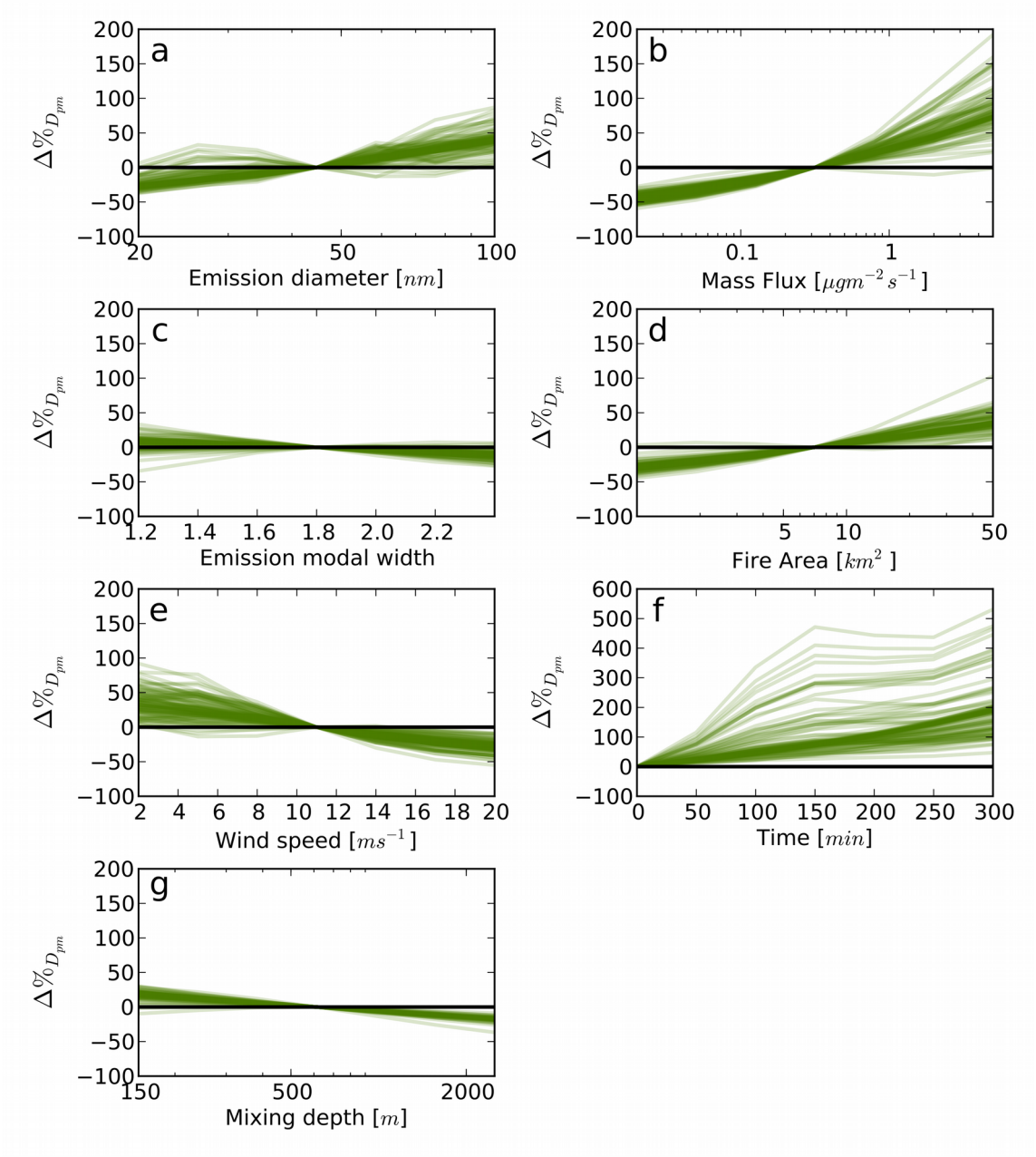

Figure 7. Sensitivity plots for the seven input parameters to the GEM-SA D_{pm} parameterization. For each panel, a single input parameter is varied systematically from its minimum to maximum value for 100 randomly chosen sets of the other six parameters (100 lines in each panel). The sensitivities are shown as percent change in final D_{pm}, individually normalized to the value at the center of the x-axis (to zero in Time).

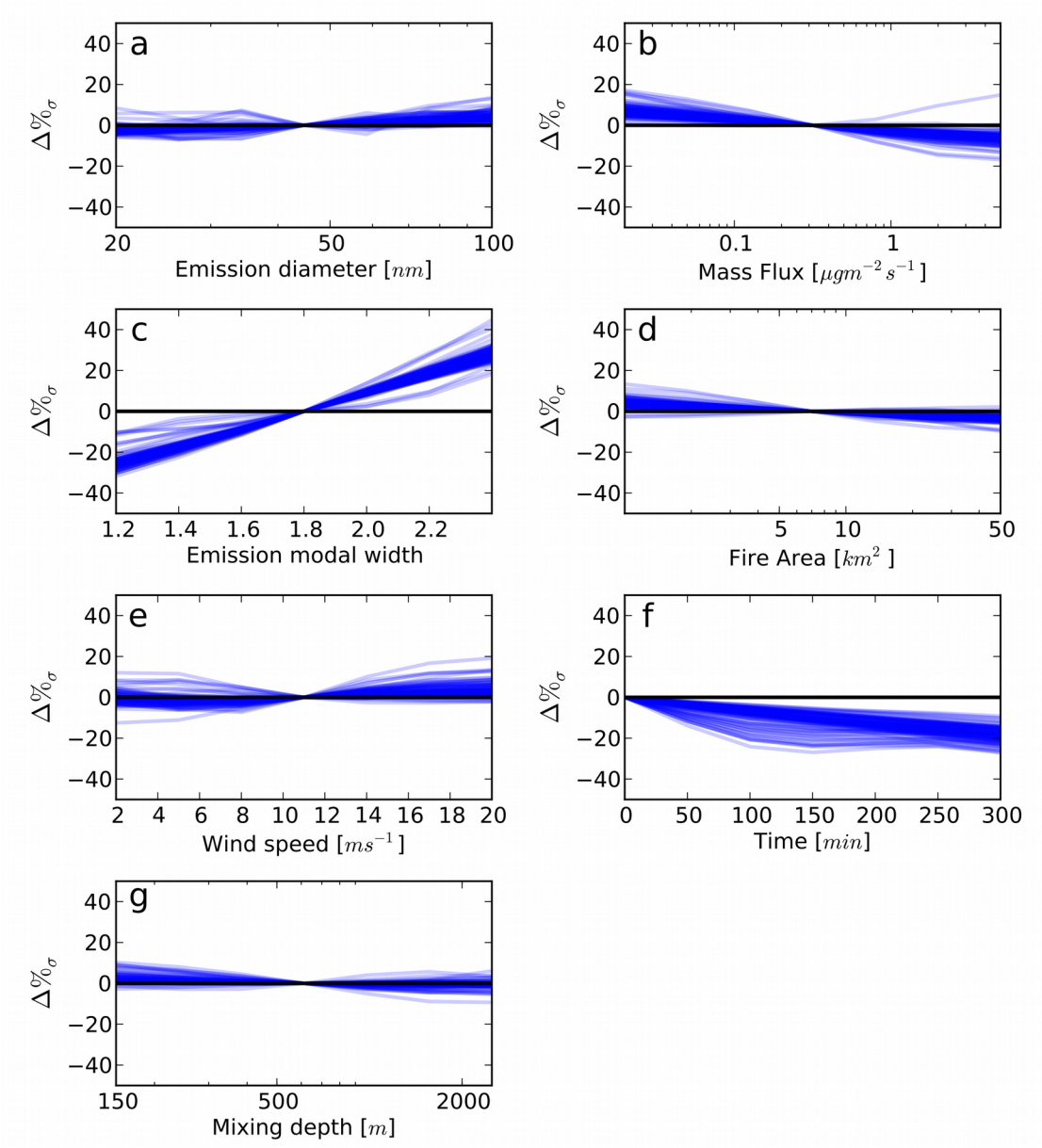

Figure 8. Sensitivity plots for the seven input parameters to the GEM-SA σ emulator parameterization. For each panel, a single input parameter is varied systematically from its minimum to maximum value for 100 randomly chosen sets of the other six parameters (100 lines in each panel). The sensitivities are shown as percent change in final σ, individually normalized to the center value of the x-axis (to zero in Time).

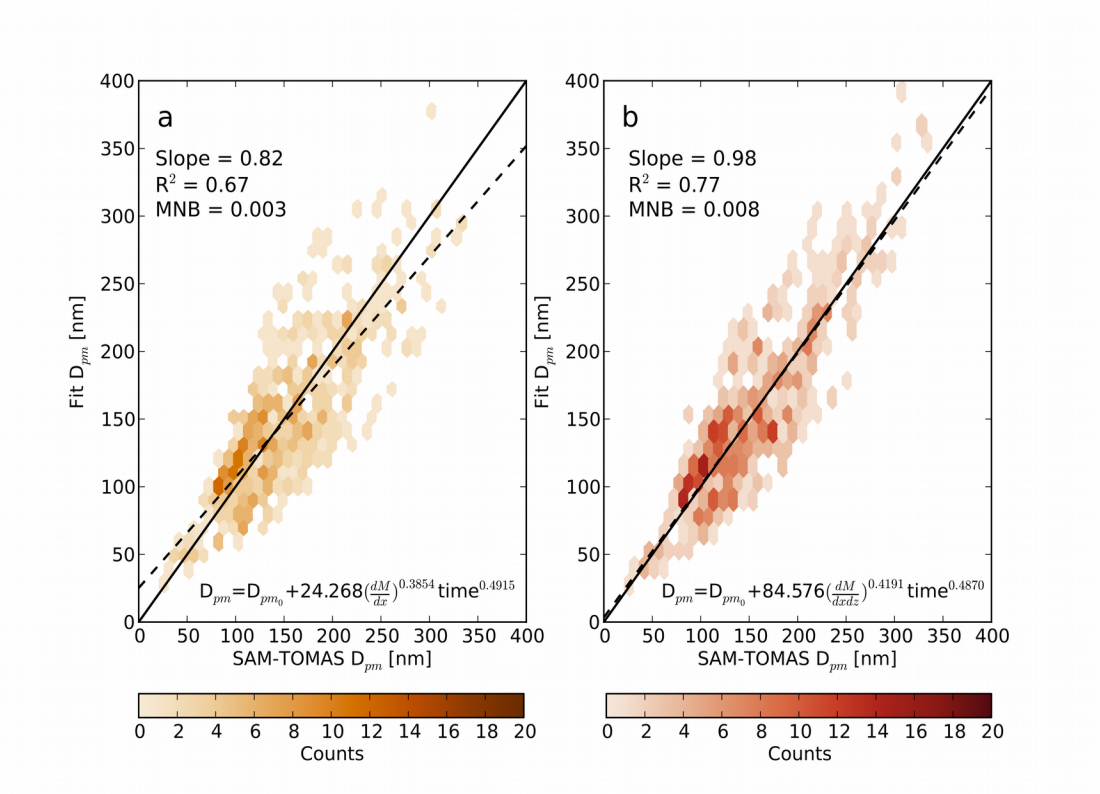

Figure 9. One-to-one plot showing simplified D_{pm} fits vs SAM-TOMAS for a) dM/dx, and b) dM/dxdz. The black line is the one-to-one line. The dashed black line is the line of best fit. N = 624.

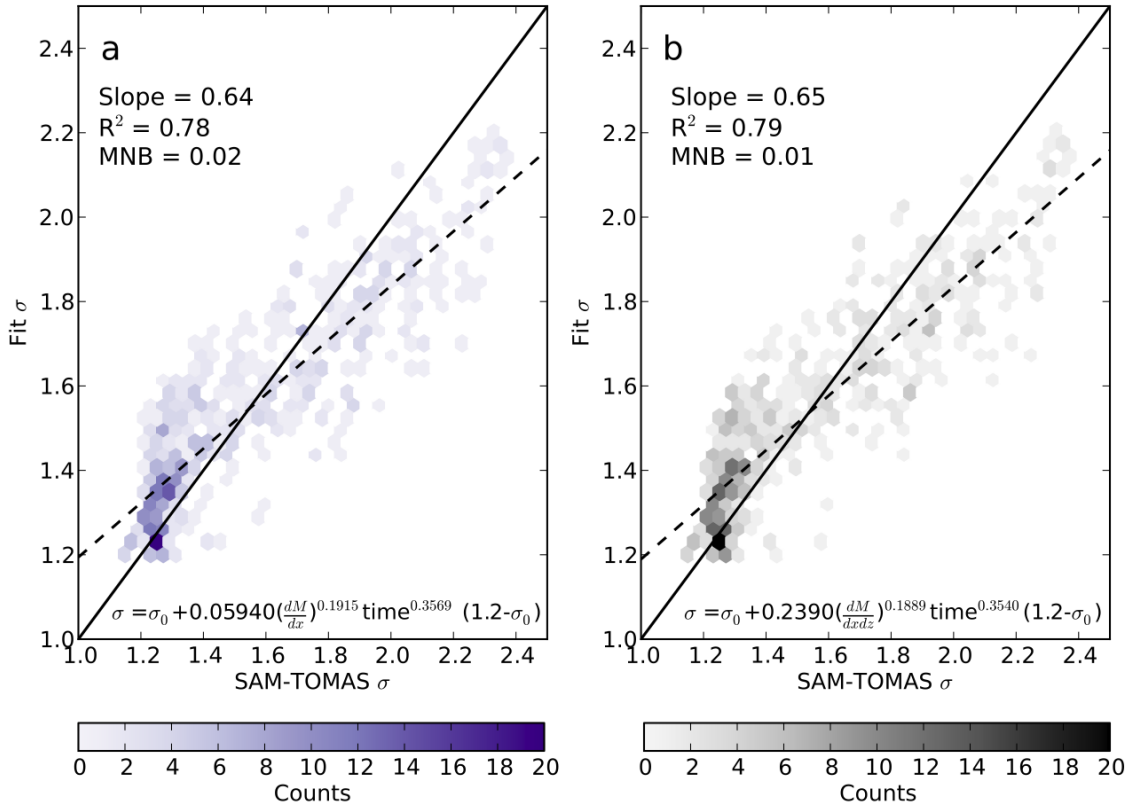

Figure 10. One-to-one plot showing simplified σ fits vs SAM-TOMAS for a) dM/dx, and b) dM/dxdz. The solid black line is the one-to-one line. The dashed black line is the line of best linear fit. N = 624.

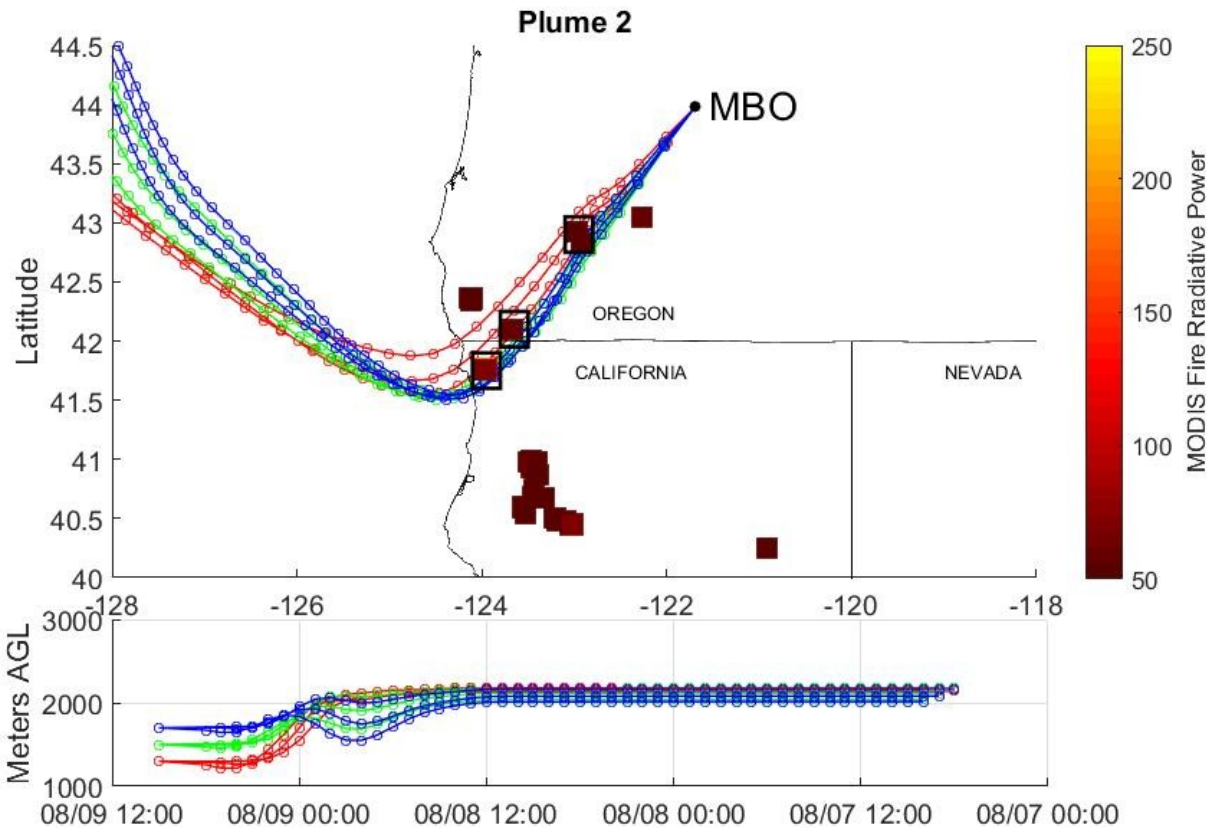

Figure 11. Back-trajectories from plume 2 observed at MBO. The colored squares represent fires during the time of the back-trajectory and are colored by Fire Radiative Power (FRP). The black squares indicate the fire areas used in the parameterization to estimate D_{pm} and σ.

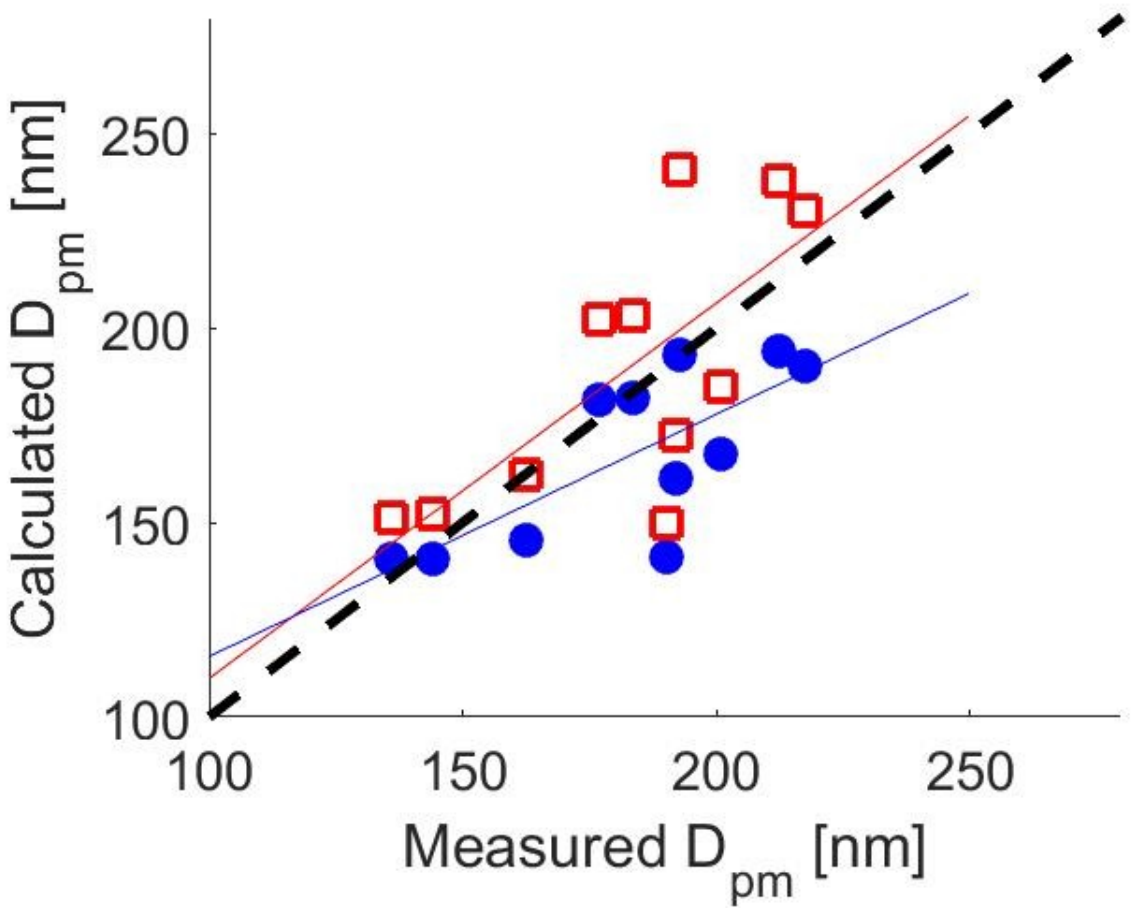

Figure 12. Scatter plot showing calculated and measured D_{pm} for biomass-burning plumes observed at MBO during August of 2015. The blue circles represent D_{pm} calculated using Eqn. 1 (dM/dx), and the red circles represent D_{pm} calculated using Eqn. 2 (dM/dxdz).

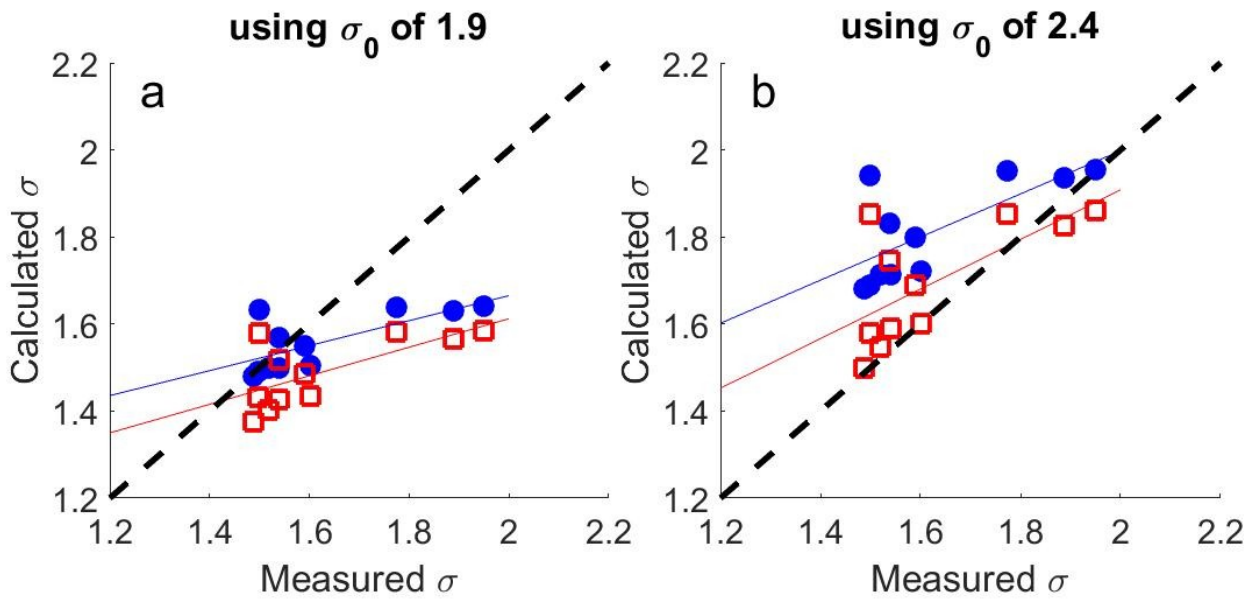

Figure 13. Scatter plots showing calculated and measured modal width (σ) for biomass-burning plumes observed at MBO during August of 2015. The blue circles represent σ calculated using Eqn. 3 (dM/dx), and the red circles represent σ calculated using Eqn 4. (dM/dxdz). Different emission modal width values (σ_0) were used to calculate σ, (a) used a σ_0 of 1.9 and (b) used a σ_0 of 2.4.