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**Calculation of
indirect global
warming potentials**

D. Youn et al.

Explicit calculation of indirect global warming potentials for halons using atmospheric models

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Received: 23 April 2009 – Accepted: 24 June 2009 – Published: 21 July 2009

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Published by Copernicus Publications on behalf of the European Geosciences Union.

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Abstract

The concept of Global Warming Potentials (GWPs) has been extensively used in policy consideration as a relative index for comparing the climate impact of an emitted greenhouse gas (GHG), relative to carbon dioxide with equal mass emissions. Ozone depletion due to emission of chlorinated or brominated halocarbons leads to cooling of the climate system in the opposite direction to the direct warming contribution by halocarbons as GHGs. This cooling is a key indirect effect of the halocarbons on climatic radiative forcing, which is accounted for by indirect GWPs. With respect to climate, it is critical to understand net influences considering direct warming and indirect cooling effects for Halons. Until now, the indirect GWPs have been calculated using a parameterized approach based on the concept of Equivalent Effective Stratospheric Chlorine (EESC) and the observed ozone depletion over the last few decades. As a step towards obtaining indirect GWPs through a more robust approach, we use atmospheric models to explicitly calculate the indirect GWPs of Halon-1211 and Halon-1301 for a 100-year time horizon. State-of-the-art global chemistry-transport models (CTMs) were used as the computational tools to derive more realistic ozone depletion changes caused by an added pulse emission of the two major Halons at the surface. The radiative forcings on climate from the ozone changes have been calculated for indirect GWPs using an atmospheric radiative transfer model (RTM). The simulated temporal variations of global average total column Halons after a pulse perturbation follow an exponential decay with an e-folding time which is consistent with the expected chemical lifetimes of the Halons. Our calculated indirect GWPs for the two Halons are smaller than past studies although direct GWPs agree with the published values. Nonetheless, our model-based assessment of the Halon indirect GWPs confirms the significant importance of indirect effects on climate.

ACPD

9, 15511–15540, 2009

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1 Introduction

As a family of atmospheric gases, chlorinated and brominated halocarbons, including the human-produced chlorofluorocarbons (CFCs) and bromofluorocarbons (often referred to as Halons), are all greenhouse gases (GHGs) in that they trap outgoing long-wave radiation in the troposphere that would otherwise escape to space. This trapped radiation warms the atmosphere, creating a positive forcing on the climate system, which in turn warms the Earth's surface. The concept of radiative forcing provides an estimate of the potential effect on climate from GHGs. For the given concentration of a gas, the radiative forcing depends primarily on the infrared absorption capacity of the gas.

Since it was developed for the first scientific assessment of the Intergovernmental Panel on Climate Change (IPCC) in 1990, the concept of Global Warming Potentials (GWPs) has been used as a relative index for comparing the potential climate impact of different GHGs to a reference gas such as CO₂ (IPCC, 1990, 2001, 2007; WMO, 2003, 2007). Although the concept of GWPs has well recognized uncertainties and limitations of its own (e.g., see Wuebbles, 1995), GWPs are a better measure of the relative climate impacts than radiative forcing alone as they also account for the atmospheric lifetime of the gases and thus the temporal change in concentration for a given emission. Numerical models of atmospheric chemistry and physics have been used to determine the atmospheric lifetimes and radiative forcings for the various compounds, and thus the GWPs have been evaluated and reported for a number of replacement compounds in the international assessments (IPCC, 2001, 2007; IPCC/TEAP, 2005; WMO, 1999, 2003, 2007). GWPs are also being used in policymaking considerations associated with concerns about climate change associated with GHGs, generally using the values for a 100-year integration time.

Of all GHGs, halocarbons tend to have very high GWPs, as much as ten thousand times the warming effect of carbon dioxide (see tables of Forster et al., 2007; Daniel and Velders, 2007). However, unlike carbon dioxide and some of the other GHGs,

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these gases play a significant role in the chemistry of stratospheric ozone when the halocarbon molecules are broken down to release chlorine or bromine atoms in the stratosphere. The distribution of ozone also has important implications for the Earth's climate system, not only because ozone absorbs solar radiation but also because it acts as a greenhouse gas that helps warm the Earth by absorbing some of the outgoing infrared radiation emitted from the Earth's surface. Changes in both the distribution and overall abundance of ozone can affect climate. Therefore, stratospheric ozone losses by increased halogen atoms can decrease the temperature of the lower stratosphere, resulting in a decreased infrared flux to the troposphere, reducing some or all of the positive radiative forcing effects caused by the increased infrared trapping of the halocarbons. This can even lead to a net cooling of the climate system in the opposite direction to the direct warming contribution of the halocarbons as GHGs (Lacis et al., 1990; Ramaswamy et al., 1992). This cooling effect through ozone destruction is a key indirect effect of halocarbons on radiative forcing and is particularly significant for bromine-containing gases. The indirect climate impact, accounted for by indirect GWPs, decreases the GWPs of these halocarbons.

Daniel et al. (1995) introduced a method to estimate the indirect GWPs for various halocarbons. The indirect GWP values, included in all international environmental assessments to date (IPCC/TEAP, 2005; WMO 1995, 1999, 2003, 2007), have been estimated from the traditional approach developed by Daniel et al. (1995). However, the traditional approach is based on the assumption of an approximate linear relationship between the changes in effective equivalent stratospheric chlorine (EESC) and in radiative cooling due to the lower stratospheric ozone loss. The concept of EESC has limitations because of uncertainties in the transport time, the bromine efficiency for ozone destruction versus chlorine (α), the efficiency of stratospheric halogen release of the source gas, the starting date for ozone depleting substance (ODS) release, the emissions scenarios assumed, possible temporal changes of transport times, and fractional bromine/chlorine releases (Newman et al., 2006, 2007). The reported indirect GWPs have primarily changed over time as a response to updates in the EESC and

updates to the carbon dioxide mixing ratio and response function. It is apparent that the indirect GWP calculated by the approach in Daniel et al. (1995) is very sensitive to uncertainties in EESC and associated radiative forcing. Combined “net” GWPs were reported only in IPCC/TEAP (2005) because the direct and indirect climate responses cannot be easily added to a climate response due to possibly different spatial dependencies (Joshi et al., 2003).

A recent reformulation of the method for estimating the EESC was provided that better accounts for the age-of-air spectrum and the age-of-air dependent fractional release values (Newman et al., 2007). Although the new EESC formulation and the sensitivity to above-mentioned uncertain parameters have been explored (Newman et al., 2007), the use of the EESC is still limited by the underlying assumptions and the simplification of the complex processes occurring in the atmosphere. Of special concern is that the EESC only provides a rough estimate of the time scale of global ozone change in an unchanging atmosphere without consideration of actual chemical, physical and transport processes in the troposphere and stratosphere. Therefore, their approach should not be regarded as a perfect gauge in evaluating ozone change related to ozone depleting substances (ODS) releases and future ozone levels. Along with the uncertainties in EESC concept, the assumption of a linear relationship between changes in the EESC and changes in radiative forcing could itself lead to potentially significant errors in estimating the indirect GWPs.

As a step toward obtaining indirect GWPs through a more robust approach, the indirect cooling effects are explicitly evaluated here for several halocarbons using state-of-the-art atmospheric models. In this study, the model-based indirect GWPs of the two major Halons, Halon-1211 (CF_2ClBr) and Halon-1301 (CF_3Br), for a 100-year time horizon have been calculated. To derive ozone changes with time, following (caused by) a pulse perturbation emission for each Halon at the surface, we use the National Center for Atmospheric Research (NCAR) Model for OZone And Related Traces version 3 (MOZART-3) three-dimensional (3-D) CTM (Kinnison et al., 2007) as well as the University of Illinois at Urbana-Champaign (UIUC) two-dimensional (2-D) CTM (Li et

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al., 2006). The UIUC radiative transfer model (RTM) was then used to calculate the temporal changes in radiative forcing through the ozone loss over the CTM integration time.

The use of state-of-the-art numerical models of the chemistry and physics of the troposphere and stratosphere is indispensable to improve and expand upon the understanding of the effects of human-related emissions of atmospheric gases on the atmospheric ozone layer and its relationship to climate change. This paper is organized as follows. A brief description of the models used for our calculations is presented in Sect. 2. The methodology to obtain indirect Global Warming Potentials (GWPs) for Halons is described in Sect. 3. Time-dependent changes in each perturbed Halon and resulting ozone destruction are presented in Sect. 4. The direct GWPs from the changes in perturbed Halons and the indirect GWPs from changes in ozone radiative forcing are estimated in Sect. 5. The discussions and conclusions are found in Sect. 6.

2 Models

2.1 Global atmospheric chemistry-transport models

The University of Illinois at Urbana-Champaign (UIUC) 2-D CTM is a zonally-averaged model for the chemistry, physics, and transport of the global atmosphere. It has often been used to study human-related and natural forcings on the troposphere and stratosphere (e.g., Wuebbles et al., 2001; Li et al., 2006) and has been used in many past WMO ozone assessments. It has the advantage of being much more computationally efficient than 3-D CTMs. It captures the most important physical and chemical processes of the global atmosphere relatively well, especially for the long-lived chemical compounds. Owing to its computational efficiency, we could perform 100-year model simulations that allowed us to do initial analyses and understand effects of perturbed Halons on ozone. It is always noted that the 2-D model has limitations in adequately representing the zonally asymmetric features such as tropospheric transport processes

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and the edge and evolution of the ozone hole. However, the limitations of the 2-D model should not greatly affect this study since the Halons are long-lived species (lifetimes larger than 10 years) enough to be relatively well mixed in the troposphere. As an initial attempt at evaluating the indirect GWPs using atmospheric models, the 2-D CTM is thus an efficient tool for evaluating the time-dependent ozone changes caused by emissions of Halons instantaneously added into the atmosphere for direct and indirect GWPs over a 100-year time horizon.

The version of the UIUC 2-D model used in this study simulates the atmospheric distributions of 78 chemically active atmospheric trace constituents with 56 photolytic reactions and 161 thermal reactions (for detailed descriptions on model chemistry and physics, see Kotamarthi et al., 1999; Wuebbles et al., 2001; Li et al., 2006). Heterogeneous chemistry processes involving polar stratospheric clouds (PSCs) are included using time-dependent parameterization of PSCs (Considine et al., 2000). Reaction-rate constants and photochemical data in the model are primarily based on the recommendations compiled by the NASA Panel for Data Evaluation (DeMore et al., 1997; Sander et al., 2003). The model domain extends from pole to pole and from the ground to 84 km in altitude. A grid box covers 5 degrees of latitude and 1.5 km in log-pressure altitude. The zonally averaged temperature and wind fields are specified based on the 10-year climatology of the United Kingdom Meteorological Office (UKMO) reanalysis data (Coy and Swinbank, 1997; Youn et al., 2006).

A state-of-the-art 3-D CTM, Model for OZone And Related Tracers (MOZART) version 3 (see Horowitz et al., 2003; Kinnison et al., 2007), was also used to explicitly calculate ozone changes caused by the same pulse emission of Halons. The computationally expensive 3-D model better represents some of the physical processes affecting ozone, especially in the troposphere. The MOZART-3 is an extension of the previous MOZART version (Horowitz et al., 2003) developed for studies of both troposphere and upper atmosphere that has been documented by Kinnison et al. (2007). MOZART is a well-established three-dimensional CTM developed by a collaboration among researchers at the National Center for Atmospheric Research and other institu-

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tions including the University of Illinois at Urbana-Champaign that has been applied to numerous studies of the atmosphere (e.g. Kinnison et al., 2007; Pan et al., 2007) and to human effects on atmospheric composition (e.g. Wuebbles and Patten, 2009). The MOZART-3 coordinates follow those of the input meteorology from the NCAR Whole-
5 Atmosphere Community Climate Model version 1b (WACCM 1b) for this study. The vertical domain ranges from the surface to 5.1×10^{-6} hPa (approximately 140 km) in a hybrid sigma-pressure coordinate. The model as used in this study uses 128 longitude points and 64 latitude points for a resolution of 2.8° in both latitude and longitude and 66 vertical levels. Meteorological fields calculated in WACCM 1b (Sassi et al., 2002)
10 serve as input to the MOZART advective transport and internal calculation of convection, clouds, and wet deposition. The WACCM 1b model year of output was selected as an average meteorology at a time resolution of three hours from a run that used boundary conditions representative of 1999 (Garcia et al., 2007).

The standard MOZART-3 chemistry includes 106 gas-phase species and is largely
15 based on the 2002 JPL panel photochemical data evaluation (Sander et al., 2003). The surface boundary conditions for both mixing ratios of long-lived source gases (methane CH_4 , nitrous oxide N_2O , hydrogen H_2 , methyl chloride CH_3Cl , methyl bromide CH_3Br , and chloro- and bromofluorocarbons) and fluxes of short-lived source gases (NO_x , carbon monoxide CO , hydrocarbons including ethane and ethylene, and oxygenated
20 hydrocarbons such as acetone $\text{CH}_3\text{C}(\text{O})\text{CH}_3$) in this study are selected to represent the year 1999. Surface mixing ratios for long-lived species are based on Table 8-5 in Daniel et al. (2007), and surface fluxes of short-lived species with their geographic and seasonal distributions are based on the POET (Precursors of Ozone and their Effects in the Troposphere) emissions data base (Pfister et al., 2008).

2.2 Radiative transfer model

The UIUC radiative transfer model (RTM) version 2.3.6 calculates radiative fluxes across the tropopause for atmospheric inputs derived from the MOZART and UIUC 2-D CTMs. Earlier versions of the UIUC RTM have been used in studies of climate

effects and in the derivation of GWPs (Jain et al., 2000; Naik et al., 2000). The current version incorporates several improvements over those versions, notably that solar radiation can now be optionally handled with the DISORT multi-stream radiative transfer kernel (Stamnes et al., 1988). The UIUC RTM has been developed to evaluate radiative properties of the outputs from global atmospheric chemical-transport models, such as the UIUC 2-D, the MOZART 3-D, and the NASA Global Modeling Initiative (GMI) 3-D models, which contain air temperature and mixing ratios of the radiatively active constituents carbon dioxide (CO₂), water vapor (H₂O), ozone (O₃), methane (CH₄), nitrous oxide (N₂O), and various halocarbons.

For shortwave radiation calculation, similar to the solar and terrestrial radiative transfer models from which it was originally created (Briegleb, 1992a, b), the solar spectrum is divided into 18 bins from 0.2–5.0 micrometers, with solar flux fractions specified for each bin. Included in the calculation are molecular, cloud, and surface scattering, along with H₂O, O₃, CO₂, O₂, clouds, and surface absorption. The longwave radiation heating rates and boundary fluxes calculation utilizes a narrow band model (NBM), which computes narrow band absorptivity and emissivity over 10 cm⁻¹ intervals from 0 to 3000 cm⁻¹ for water vapor and 5 cm⁻¹ from 0 to 3000 cm⁻¹ for other gases and then sums them up to obtain the broad band radiative forcings. Orography and surface albedo are based on observations, while clouds are based on the International Satellite Cloud Climatology Project (ISCCP) climatology and tropopause height on the National Centers for Environmental Prediction (NCEP) reanalysis. The HITRAN (high resolution transmission) 2004 database (Rothman et al., 2005) provides spectral data for the radiatively active species, which are averaged to the spectral intervals within the UIUC RTM.

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3 Methodology

3.1 Indirect Global Warming Potential (GWP)

The indirect cooling effect of the two major Halons on climate arising from ozone depletion can significantly cancel the direct global warming effect of these GHGs. The indirect GWPs are of special interest because of their potential use in future policy decisions related to ozone and climate. However, the current approach to estimate the indirect GWPs has potentially serious limitations due to the parameterization approach and underlying assumptions, as explained earlier. Therefore, in this study the indirect GWPs are evaluated through the explicit calculations using state-of-the-art atmospheric models.

Our model-based analyses of indirect GWPs are derived based on the classical definition of GWPs. Since the first scientific assessment (IPCC, 1990), the GWP has been used as a relative index for comparing the potential climate impact of one greenhouse gas to another, based on globally averaged radiative forcing of the climate system over a specified time scale. GWPs are expressed as a ratio of the time-integrated radiative forcing from the instantaneous release of a kilogram of a gas relative to that of a kilogram of the reference gas, generally carbon dioxide, over a fixed time horizon (IPCC/TEAP, 2005; IPCC, 2007; WMO, 2007). Thus, the direct Global Warming Potential of a gas χ (IPCC, 1990, 2001) can be defined as:

$$\text{GWP}_{\chi}(\text{TH}) = \frac{\int_0^{\text{TH}} F_{\chi} \cdot [\chi(t)] dt}{\int_0^{\text{TH}} F_{\text{CO}_2} \cdot [\text{CO}_2(t)] dt} \quad (1)$$

where TH is the integration time (the time horizon) over which the calculation is performed, F_{χ} is the radiative forcing efficiency due to a unit mass change in the atmospheric abundance of the substance χ , and $[\chi(t)]$ describes the time-dependent decay of the abundance from an instantaneous release of gas χ . F_{χ} and F_{CO_2} are generally given in units of $\text{W m}^{-2} \text{kg}^{-1}$.

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Time-dependent changes in radiative forcing, resulting from ozone depletion caused by an emission of chlorine and bromine-containing species χ , are a key indirect effect of the gas on radiative forcing. If the numerator in the Eq. (1) is substituted with the time-integration of the time-dependent changes in radiative forcing due to ozone depletion following an instantaneous release of a kilogram of halocarbon χ , then the resulting GWP becomes the indirect GWP of χ through ozone depletion.

In this study, the evaluation of indirect GWPs for Halon-1211 and Halon-1301 over the 100-year time horizon has been accomplished in three steps: 1) time-dependent simulation of an instantaneously perturbed Halon and ozone response to the additional Halon pulse release using atmospheric CTMs; 2) calculation of radiative forcing using an atmospheric RTM; and 3) evaluation of a 100-year direct and indirect GWP. The MOZART-3 3-D and the UIUC 2-D CTMs are used as the computational tools for deriving ozone changes after an instantaneous surface release (one-year pulse perturbation emission) of the given Halon. Radiative forcing changes induced by the time-dependent ozone losses during the model simulations over the time-horizon are calculated with the UIUC RTM. To obtain complete 100-year changes in atmospheric concentrations and radiative forcing for the 3-D simulations, exponential curve fits have been applied as described in following sections for the two purposes of validating model simulations and filling in the tail part of remaining years.

3.2 Model setup and simulations

From initial analyses of the indirect GWPs for the Halons, we found that a correct CTM simulation of ozone responses to Halon perturbations requires two prerequisite conditions: 1) a pulse perturbation injection of Halon emissions at the surface rather than a traditional treatment of lower boundary condition (LBC), specifying mixing ratios in the lowest model layer, in CTMs for long-lived species (explained below) and 2) a starting CTM reference atmosphere that is in steady-state. A one-year surface perturbation pulse at the first year of model simulation needs to be put into the model atmosphere in a steady-state.

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The traditional CTM setup to represent the perturbation of a long-lived source gas, using a change of the specified concentration in the lowest model layer, cannot be applied to this study of evaluating time-dependent effect of a Halon perturbation on ozone. This is because the model setup mixes the chemical compound unrealistically too fast by constraining the lower boundary with a fixed concentration and so atmospheric temporal variations of the compound are confined by the lower boundary concentration of the compound. Specifying boundary concentration of a compound is also like continuously accumulating additional perturbations at the lower boundary, resulting in faster response of confining whole atmospheric distribution in shorter time. In other words, the model setup does not represent the actual atmospheric behavior since the specified concentrations at the lower boundary override the mixing processes or interaction of the compound at the surface and higher levels. This makes the model atmosphere respond to the given Halon perturbation too fast because the specified concentrations also act as another perturbation added and even accumulated up to the equilibrium of the model atmosphere. Model simulations of Halon perturbation were therefore revised to obtain reasonable time-dependent atmospheric response to a single-year pulse perturbation. In this study, the approach with surface flux LBCs for Halons was newly implemented into the CTMs as means of atmospheric loading of the Halons. However, it requires a lot more years of simulations (and thus much more computational time) for CTMs with the new Halon LBC. This is because atmospheric Halon and stratospheric ozone respond to the surface perturbation more slowly as the change in model LBCs represents more realistic mixing and transport processes from the surface to high altitudes.

The change in the LBC of the CTMs also requires additional model runs to get the perturbation simulations start from a steady-state model atmosphere with an unperturbed normal amount of surface Halon emission (comparable to the approach of specifying the mixing ratios at the lower boundary). For the 2-D surface perturbation simulations of Halon-1211 and Halon-1301, it requires about 30 and 60 years, respectively, to reach a steady-state atmosphere even before perturbations are actually injected. If the

perturbation run does not start from a steady-state, the resultant changes in radiative forcing during model simulation period, especially before peak ozone loss, will not be correctly obtained since the starting atmosphere was already in a perturbed state and the impacts of pre-existing Halons on the stratospheric ozone remain.

5 The extensive computational time in getting to a steady-state in the 3-D model led us to detour this issue by a new simulation method of simultaneously conducting two parallel simulations for the “standard” and “perturbed” atmospheres, respectively. The starting-point atmosphere of the standard simulation, same with that of the perturbation simulation, is in a perturbed state as a result of the new implementation of surface
10 flux emissions. While the standard model atmosphere comes closer to a steady-state during the model integration, the perturbation simulation responds more to the pulse injection imposed at the surface upon standard flux emissions. Ozone changes over the model integration time in the perturbation simulation, only resulting from the instantaneous pulse emission injection into the starting-point atmosphere, can be derived
15 by subtracting the standard simulation from the perturbation simulation. This subtraction minimizes the errors related to the drifting initial model atmosphere coming from the change in the model LBC setup, which are shown with our simulations in the next section.

20 In addition, it is not feasible to conduct the MOZART-3 3-D simulations for 100 years as done with the 2-D CTM, due to the high computational demand and limited computing resources. Fortunately, from the initial analyses with the UIUC 2-D CTM, we find that stratospheric ozone destructions due to Halon-1211 and Halon-1301 pulse perturbations recover exponentially after the time of peak loss. From the 2-D analyses, we also find that the 100-year cumulative decrease of ozone can be projected
25 with less than $20+\alpha$ ($30+\alpha$) years integration of Halon-1211 (Halon-1301) perturbation. Thus, we expect that $20+\alpha$ ($30+\alpha$) years of 3-D simulations enable us to estimate approximate 100 years of total radiative forcing for Halon-1211 (Halon-1301) through the application of exponential curve fitting to the tail part of the radiative forcing changes corresponding to the ozone recovery. The final radiative forcing changes over the full

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integration time should be properly obtained by calculating the differences in temporal variations of ozone radiative forcing between the perturbation and the standard simulations.

With this approach, the 3-D indirect GWP of the Halon-1211 (Halon-1301) for a 100-year time horizon has been estimated based on 23 (32) years of perturbation simulations along with 8 (14) years of standard ones.

4 Model results

4.1 Halon perturbation and direct GWP

A one-year pulse perturbation of each Halon is injected into the surface air during the first year of the 2-D and 3-D model simulations. The starting-point model atmosphere is at steady-state for the 2-D model but not for the 3-D model. The two parallel 3-D simulations of standard and perturbation cases are thus conducted as described in the previous section.

Changes in global total mass of Halon-1211 and Halon-1301 over time after the first year pulse injections into 2-D and 3-D model simulations are shown in Fig. 1. The time-dependent 2-D (3-D) simulation of Halon-1211 is denoted in red (blue) color. Global total masses of both Halons in the model atmosphere reach maxima at the second model year, while the photolytic break-down of the added Halons transported into the stratosphere leads to Halon reductions in the following years. They show exponential decrease after the second year peak. Exponential curve fitting was applied and the curve shows a lifetime (e-folding time) of 14.4 (10.9) years for 2-D (3-D) model-determined Halon-1211 variation, which is consistent with 14.8 (11.2) years of model derived chemical lifetime, total atmospheric loading over chemical loss rates (see Table 1). The modeled e-folding lifetimes of Halon-1301 are also in good agreement with the chemical lifetimes. Curve fits applied to the 3-D simulated Halons serve to fill in the remainder of a 100-year time horizon.

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Since photochemical data used in 2-D and 3-D CTMs for photolytic reactions of Halons are consistent, the shorter lifetimes of 3-D simulated Halons can be explained by the faster transport and mixing into the stratosphere in that the MOZART-3 3-D model has shorter age of air (Kinnison et al., 2007) than the UIUC 2-D model (Guillas et al., 2004). The 3-D model has about half to one year shorter age of air in the stratosphere, which means faster transport and mixing of the Halons and affected ozone in the stratosphere. Although the 3-D model lifetimes of about 11 years for Halon-1211 are particularly small compared with the 2-D lifetimes of about 14.5 years, the lifetimes are within the range of reported values 11–16 years in WMO assessments in Table 1. 2-D and 3-D model lifetimes for Halon-1301 are in reasonable agreement.

Direct GWPs for a 100-year time horizon for the simulated Halons are calculated for the comparison with previous assessments. The traditional approach to estimate absolute GWPs (AGWP) uses radiative efficiency and atmospheric mixing ratios with their lifetimes (IPCC/TEAP 2005; WMO, 2007). In this study, AGWPs for the Halons are evaluated using the actual time-dependant variations of 2-D and 3-D simulated mixing ratios. The variations of the tropospheric Halons are shown in Fig. 2. We use the CO₂ AGWP value of 0.676 W m⁻² ppmv⁻¹ yr for the time horizon of 100 years in the latest ozone assessment (WMO, 2007). A radiative efficiency of 0.30 (0.32) W m⁻¹ ppbv⁻¹ for Halon-1211 (Halon-1301) yields 2-D model-based direct GWP of 1796 (7122) and 3-D direct GWP of 1699 (6903) (see Table 1). The direct GWP values are smaller than but compare reasonably with the reported GWP of 1890 (7140) in WMO (2007).

4.2 O₃ response to the Halon perturbation

The increased bromine and chlorine in the stratosphere due to the pulse perturbation injection of Halon-1211 (only bromine for Halon-1301) lead to the stratospheric ozone depletion. Temporal changes in the stratospheric ozone destruction and recovery induced by the pulse perturbation are shown in Figs. 3 and 4. The 2-D and 3-D results show similar temporal variations of global average ozone.

Ozone loss following the Halon increase accumulates at the first phase up to peak

loss. Maximum ozone loss lags the peak Halon increase by 7–8 years for Halon-1211 and 10–13 for Halon-1301 (Fig. 3). Ozone responses to the added Halon-1301 are much slower due to much longer lifetime of Halon-1301, 70–72 years for our model simulations which are larger than but comparable to 65 years in WMO (2003, 2007).

5 Compared to the full 100-year 2-D simulations, the exponential curves for 3-D simulations are good fits for both Halons. The ozone, depleted due to an instantaneous Halon-1211 perturbation, fully recovers and comes close to the initial state within the 100 year time frame. However, this is not the case for Halon-1301, which could imply that impacts of the ozone-depleting species with longer lifetimes on the climate need
10 to be considered over a longer time horizon.

The time-dependent impacts of the pulse injected Halon-1211 on atmospheric ozone at different height levels of 2.25 (near surface), 18.75 (tropopause), 24.75 (lower stratosphere), and 45.25 km (upper stratosphere) indicate delayed ozone loss depending on the heights (Fig. 4). The globally averaged ozone concentrations were normalized
15 relative to the steady-state concentration on each level. The maximum ozone depletion appears at the level near the tropopause with peak percentage depletion of about 1.2–1.4%. The ozone loss near the tropopause (18.75 km) peaks similarly at 7–8 years after the injected pulse emission for both 2-D and 3-D runs. Decreases in ozone losses at different levels also show exponential behaviors after the peak. Ozone changes near
20 the surface lag changes at the other levels due to the transport time from the place of ozone depletion, the stratosphere, to the troposphere and to the surface.

Global total ozone masses from the MOZART-3 3-D CTM are up to 0.3% smaller than those from the UIUC 2-D throughout time. Ozone recovery after peak destruction shown in time-series of simulated ozone from the 2-D model exponentially fits almost
25 perfectly (good matches between simulated ozone and the exponential curve fits). In the 2-D CTM, the global total ozone mass at the first year, when the pulse perturbation is just injected into the surface air and has little influence on ozone, matches that at the 100th year very well. The agreement between ozone masses at starting and ending years suggests the 2-D simulations really started from a steady-state atmosphere. A

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curve fit of exponential ozone recovery for 3-D Halon simulations represent reasonably well the temporal variation of ozone recovery after the peak loss. The difference between 2-D and 3-D ozone recovery could be attributed to differences in model characteristics such as model physics and transport time (the age of air).

Based on the simulated ozone response, temporal changes in the calculated radiative forcing resulting from the ozone destruction are expected to have an exponential variation very similar to the ozone recovery. The remaining years without 3-D simulations can thus be assumed to follow the exponential curve and so the 100 years of cooling effects caused by ozone depletion due to the Halon perturbation can be estimated. The corresponding instantaneous radiative forcing calculations are presented in the following section.

4.3 Radiative forcing and indirect GWPs

We used the UIUC RTM to calculate changes in the radiative forcings (RF) due to the ozone depletion induced by the Halon perturbation. Time-series of changes in calculated radiative forcing (down radiative flux minus up flux), net radiative imbalance by ozone loss, for a 100 year period are presented in units of $\text{W m}^{-2} \text{kg}^{-1}$ in Fig. 5. The change indicates the global annual mean radiative cooling at the tropopause arises from ozone depletions. Changes in the radiative cooling per added mass of each Halon at the tropopause clearly show an exponential decrease of the cooling during ozone recovery. Exponential curve fits were again applied to the changes in the radiative forcings so that the remaining cooling could be estimated based on the most proper fitting function, a curve fit.

The cooling at the 100-th year atmosphere for 2-D Halon-1211 simulation is negligible (almost zero) as the atmosphere is approximately in steady-state at that time. Therefore, the fitting function for 3-D Halon-1211 simulations needs to be tuned to approach zero at the infinite time, which yields the best curve fit. The magnitudes of all the 3-D radiative coolings after the peak cooling were adjusted, so that the value of applied curve fit at the infinite time is zero. However, the radiative cooling variations

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before the peak cooling were obtained just by subtracting the two RFs from the perturbation and standard runs without the adjustment. After this subtraction, the first year cooling gets to be very close to zero as shown in Fig. 5, which means the validity of two parallel simulations used to overcome the issue that the initial 3-D model atmosphere is not at the steady-state. This builds robustness of our approach doing two parallel simulations with and without the first year perturbation.

For Halon-1211 (Halon-1301), the largest 2-D and 3-D radiative cooling effects appear at the 8th (14th) year, similar to the time of the peak ozone depletion. The 3-D peak values are larger than the 2-D value. Radiative forcing changes caused by the instantaneous releases of Halons were time-integrated on per mass base over the 100-year time horizon. As described in the previous section, the ratio of the 100-year cumulative radiative forcing to the 100-year CO₂ forcing is the indirect GWP. The indirect GWPs for Halon-1211 (Halon-1301) are calculated to be -16 294 (-36 247) for 2-D model outputs and -17 050 (-37 252) for 3-D model, using the same reported CO₂ AGWP (WMO, 2007). Compared to the previously published indirect GWPs in Table 1, our calculated values are smaller but those for Halon-1301 are close to those in IPCC/TEAP (2005) and WMO (2007).

5 Discussion and conclusions

The 2-D CTM simulations and accompanying radiative forcing calculations are conducted for a complete 100 year time horizon to derive indirect GWPs of two major Halons. The 3-D CTM is computationally too expensive to be run for so many years to reach a (near) steady-state atmosphere especially when modeled atmospheric variations of Halons are driven by surface emissions instead of the specified concentrations at the lowest level of the model. The two parallel 3-D CTM simulations of standard versus perturbation cases for each Halon are shown to reasonably account for this difficulty. The huge computational demand to complete a 100-year 3-D simulation for each of Halons has been overcome by applying an exponential curve fit to 23 years of

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simulation for Halon-1211 and 32 years for Halon-1301. More years of simulations are needed for Halon-1301 with longer lifetime, about 70 to 72 years estimated in this study. The direct GWPs and lifetimes are estimated with simulated temporal variations of the Halons. The resulting direct and indirect GWPs as well as lifetimes are summarized in Table 1. Published values in major assessments are also included for comparison.

The 2-D and 3-D CTM simulations show reasonable temporal variations of the Halons and resulting ozone changes in the atmosphere, yielding consistent e-folding lifetimes and direct GWPs. The validity of the 2-D and 3-D indirect GWPs is partially guaranteed by the agreement between the e-folding lifetimes estimated from a curve fit for modeled exponential decrease of global total mass and model chemical lifetimes derived from atmospheric total burden (loading) and chemical loss. Moreover, the credibility of our current simulations seems to be achieved because direct GWP values as well as lifetimes are generally consistent relative to the previous assessments.

Our approach of explicitly calculating the indirect GWPs using atmospheric models is an improvement beyond previous evaluations in that our indirect GWP calculations are not based on a parameterization approach using the EESC index. The new assessments of indirect GWPs for Halons do not assume the relationship among stratospheric ozone depletion, inorganic chlorine and bromine loading (no need to use fractional release factors, alpha values, assumed transport time, etc.), and radiative forcing. It thus does not depend on the uncertainties unavoidable in the previous traditional approach and more closely matches the desired concept of the (indirect) GWP. It is mainly because of the explicit calculation using atmospheric models. Our model-based explicit calculation can thus be used as a more robust advancement for obtaining indirect GWPs of Halons, relative to the approach by Daniel et al. (1995). However, it is still affected by reported CO₂ AGWP and the fact that the models adopted for this study are driven by climatological meteorological fields. Also, little temperature feedback between ozone and climate is considered. As noted in previous studies (e.g., Hansen et al., 1997), the effect of radiative forcing on climate is dependent on the distribution of the perturbation of ozone both vertically and horizontally.

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The indirect GWPs for Halon-1211 derived here are much smaller than recently published results (IPCC/TEAP, 2005; WMO, 2007). However, it is still much larger by about a factor of ten compared to its direct effect. The indirect effect of Halon-1301 is more than double that of the Halon-1211 and much larger than its own direct warming effect.

5 These analyses suggest that the traditional approach for deriving indirect GWPs is not sufficiently accurate and that global CTMs can be effective tools for these derivations.

Acknowledgements. This research was supported in part by the US Environmental Protection Agency using computer resources at the University of Illinois' National Center for Supercomputing Applications (NCSA). This research also used computer resources of the Department of Energy's National Energy Research Supercomputing Center (NERSC). Daeok Youn was supported in part by the Korea Meteorological Administration Research and Development Program under grant CATER 2007-3205 and Jintai Lin by US NSF under ATM-0635548.

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Table 1. Direct and indirect Global Warming Potentials (per mass basis) of Halons relative to carbon dioxide.

Previous and present studies	Lifetimes (years)		Global Warming Potentials				for 100-years time horizon	
	Halon-1211	Halon-1301	Direct GWP		Indirect GWP		Net GWP	
			Halon-1211	Halon-1301	Halon-1211	Halon-1301	Halon-1211	Halon-1301
Daniel (1995)	20	65	N/A	4460	N/A	-92 520~-15 820	-42 020~-7100	-88 060~-11 360
WMO (1999)	11	65	1300	6900	N/A	-103 930~-16 410	N/A	-197 030~-9510
WMO (2003)	16	65	1860	7030	-35 220~-6910	-66 310~-13 000	-33 360~-5050	-59 280~-5970
IPCC /TEAP (2005)	16	65	1860±650	7030±2460	-28 200±19 600	-32 900±27 100	-26 340±20 250	-25 870±29 560
WMO (2007)	16	65	1890±660	7140±2500	-40 280±27 120	-49,090±34 280	-38,390±27 780	-41 950±36 780
2-D model	14.4 ^a (14.85) ^b	72.4 ^a (75.2) ^b	1796	7122	-16 294	-36 247	-14 498	-29 127
3-D model	10.9 ^a (11.16) ^b	70.1 (70.04) ^b	1699	6903	-17 050	-37 252	-15 351	-30 349

^a The lifetimes are e-folding times of the exponential curve fitted to the model simulated time-series of the Halons.

^b Derived lifetimes in parenthesis are global chemical lifetimes calculated using atmospheric total burden (loading) and chemical loss-rates in the models.

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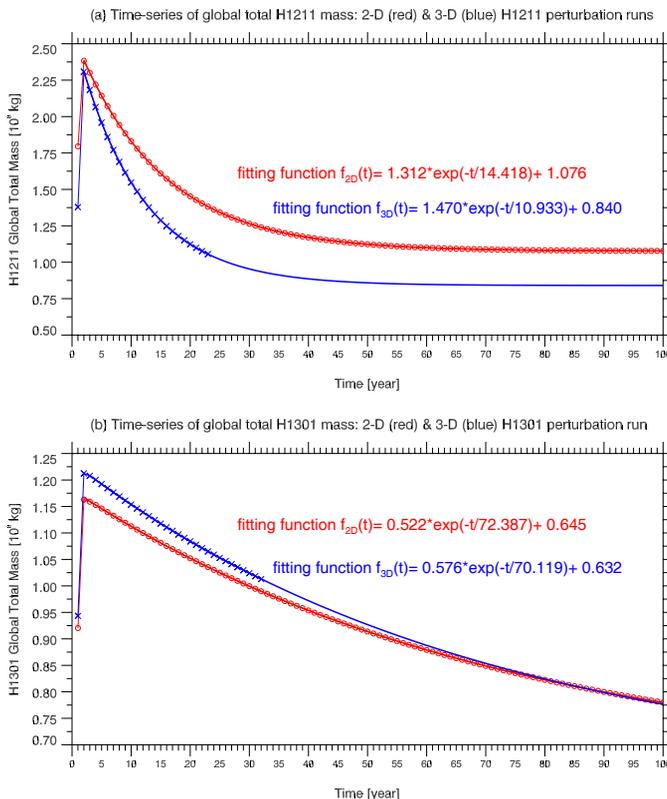


Fig. 1. Changes in global total mass of **(a)** Halon-1211 and **(b)** Halon-1301 over time after a single-year pulse injection of 10^8 kg. Red lines with circle marks denote outputs from the UIUC 2-D CTM and blue with cross from MOZART-3 3-D CTM. Exponential fitting functions are estimated for the decreasing mass.

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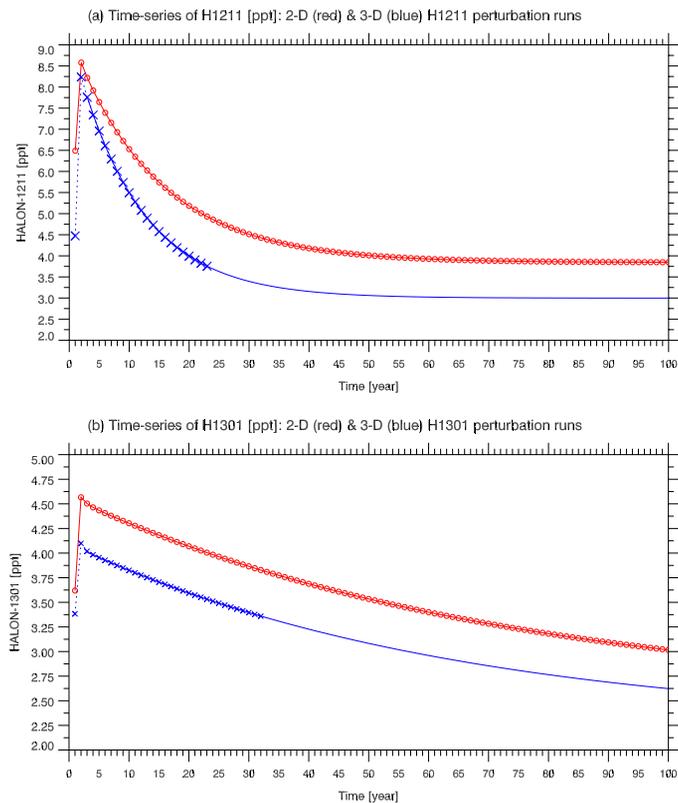


Fig. 2. Same as Fig. 1, but for tropospheric mixing ratios in ppt.

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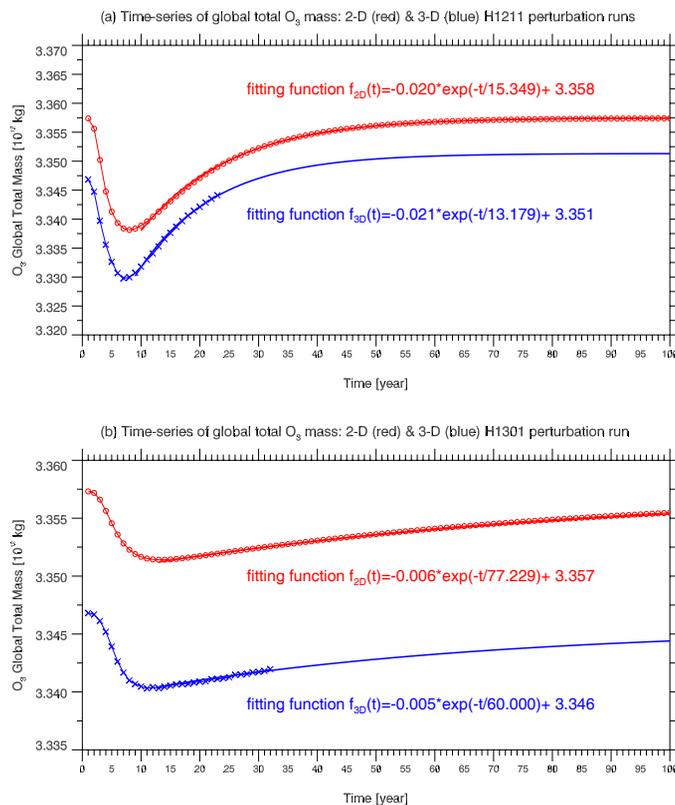


Fig. 3. Same as Fig. 1, but for ozone in 10^{12} kg. Exponential fitting functions are estimated for the years of ozone recovery.

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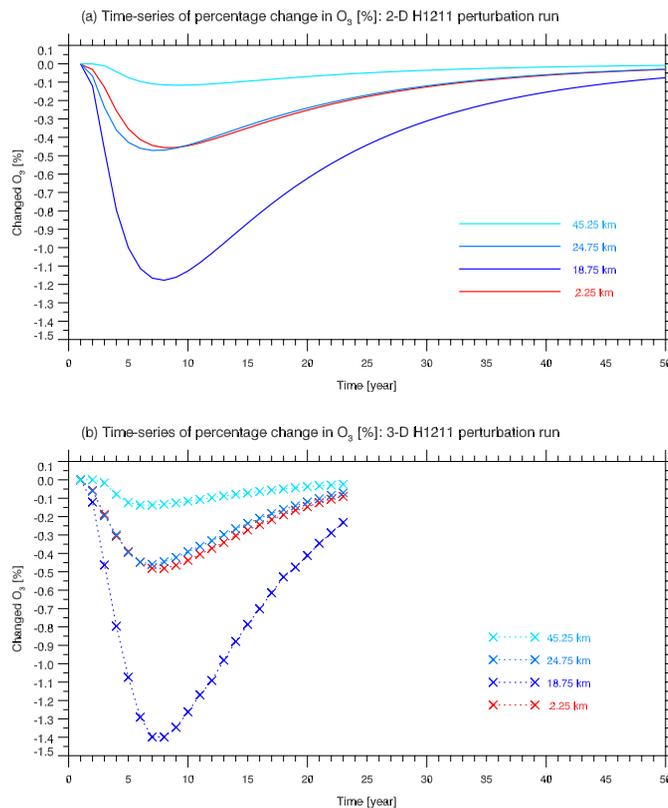


Fig. 4. Percentage change of globally averaged ozone depletion over time on selected height levels for **(a)** 2-D and **(b)** 3-D perturbation runs.

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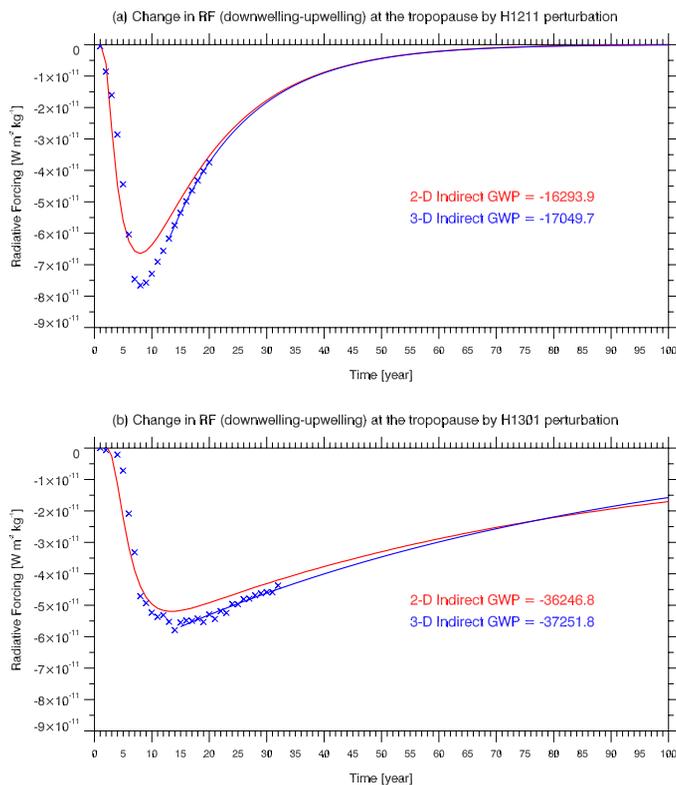


Fig. 5. Time-series of radiative cooling change due to (a) Halon-1211 and (b) Halon-1301 perturbation using the UIUC 2-D and the MOZART-3 3-D CTMs. The best fits are estimated for the years of ozone recovery.

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