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On the relationship of polar mesospheric cloud ice water content, particle radius and mesospheric temperature and its use in multi-dimensional models

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ACPD 9, 14571-14600, 2009 **Relationship of PMC** ice water content, radius and **Temperature** A. W. Merkel et al. **Title Page** Introduction Abstract Conclusions References **Tables Figures** 14 Close Back Full Screen / Esc **Printer-friendly Version** Interactive Discussion

Abstract

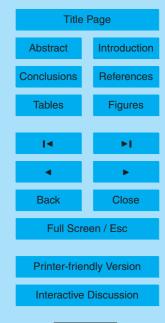
The distribution of ice layers in the polar summer mesosphere (called polar mesospheric clouds or PMCs) is sensitive to background atmospheric conditions and therefore affected by global-scale dynamics. To investigate this coupling it is necessary to simulate the global distribution of PMCs within a 3-dimensional (3-D) model that couples large-scale dynamics with cloud microphysics. However, modeling PMC microphysics within 3-D global chemistry climate models (GCCM) is a challenge due to the high computational cost associated with particle following (Lagrangian) or sectional microphysical calculations. By characterizing the relationship between the PMC radius, ice water content (*iwc*), and local temperature (*T*) from an ensemble of simulations

- ¹⁰ ice water content (*iwc*), and local temperature (*T*) from an ensemble of simulations from the sectional microphysical model, the Community Aerosol and Radiation Model for Atmospheres (CARMA), we determined that these variables have a fundamental stable relationship that is independent of the cloud time history. For our purposes we use this relationship to provide a bulk parameterization of PMC microphysics to be in-
- ¹⁵ cluded into multidimensional models. However, this relationship has the potential to be used in other applications to characterize PMCs. We use a parameterization of the relationship to predict the particle effective radius using only the local temperature and ice water content in a 3-D GCCM for decadal scale PMC simulations. Such a parameterization allows for theoretical ice cloud microphysics to be applied in the GCCM to
- simulate growth, sublimation and sedimentation of ice particles without keeping track of the time history of each ice particle size or particle size bin. This approach produces realistic PMC simulations including estimates of the optical properties of PMCs. We validate the relationship with PMC data from the Solar Occultation for Ice Experiment (SOFIE).

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1 Introduction

The Aeronomy of Ice in the Mesosphere (AIM) satellite was launched in April 2007 with its primary focus to study the formation of Polar Mesospheric Clouds (PMC) (Russell et al., 2009). Several investigations based on the first season of PMC observations have

- shown complex dynamical features, evidence of small wavelength gravity wave activity and planetary wave activity (Rusch et al., 2009; Chandran et al., 2009; Merkel et al., 2009). These studies indicate that seasonal variability of PMCs is largely affected by the dynamical conditions of the summer atmosphere. In addition, there is evidence of a long-term increase and solar-cycle variation in PMC brightness and frequency
 (Deland et al., 2007; Shettle et al., 2009). Long-term PMC variability is thought to be an
- 10 (Deland et al., 2007; Shettle et al., 2009). Long-term PMC variability is thought to be an indicator of a changing climate (Thomas et al., 2003), however satellite observations of PMCs are limited to about three solar cycles (Deland et al., 2007). Therefore modeling efforts are crucial in understanding new AIM PMC observations in the context of longer-term changes resulting from solar and anthropogenic forcing.
- Modeling efforts that use a Lagrangian ice transport scheme in a 3-D middle atmosphere model produce detailed information about cloud processes and properties on seasonal time-scales (Berger and Lübken, 2006; Lübken and Berger, 2007). However, Lagrangian models are inefficient for long-term integrations (such as several decades) due to inherent computational cost and data storage limitations. Although not as computational exhaustive, models that use sectional codes to include ice microphysics
- are also computationally expensive (Bardeen et al., 2008, 2009). For instance, the Bardeen et al. (2009) PMC ice model includes 28 dust bins and 28 ice particle bins with each advected individually at each time step. To overcome this problem, it is desirable to have a bulk parameterization for PMC microphysics to include in multi-dimensional ²⁵ atmospheric global models to facilitate simulations on decadal time scales.

A requirement for an accurate bulk PMC parameterization is a reasonable estimate of the mean ice mass, which depends on particle growth rates. The growth of ice particles not only depends on the saturation conditions and temperature but on the size

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of the particle. Therefore, given a parameterization of mean radius, the microphysical equations for cloud particle growth, sublimation and sedimentation (Gadsden et al., 1998) can be used to model PMCs within a 3-D general circulation model (GCM). Siskind et al. (2006) introduced a bulk PMC parameterization into the CHEM2D model

⁵ to study the radiative feedback from PMCs. Their simplistic parameterization assumes a constant particle size of 70 nm. This type of assumption will cause errors in the amount of mean ice mass created, destroyed and sedimented at each time step, since the microphysics is calculated for only one particle size.

By analyzing the relationship between the PMC radius, ice water content (*iwc*), and local temperature (T), we determined that these variables have a fundamental stable relationship. For our purposes we use this relationship to parameterize PMCs in a GCM. However, this relationship has the potential to be used in other applications to characterize PMCs. This paper will illustrate this relationship and its potential use in multi-dimensional global climate models.

15 2 Relationship between PMC *iwc*, *T* and radius

Our work is based on the development of a similar relationship found between temperature, ice water content and particle radius for tropospheric clouds (Boudala et al., 2002). Boudala et al. developed a parameterization of this relationship to predict effective radius of tropospheric ice particles using local T and *iwc* in a global model. We

- ²⁰ characterized the relationship between these PMC variables, using an ensemble of simulations from the one-dimensional version of the Community Aerosol and Radiation Model for Atmospheres (CARMA) microphysical model Rapp and Thomas, 2006. The CARMA model is used to establish this relationship because it provides full microphysical processes including nucleation, condensational growth, particle sedimentation and
- transport of ice particles on small time scales. The CARMA model handles three interactive constituents: meteoric smoke particles (distribution by Hunten et al., 1980), ice particles and mesospheric water vapor. The model is comprised of 120 altitude

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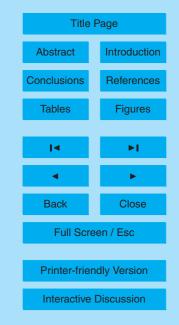
levels from 72 to 102 km and consists of 40 logarithmically spaced radius size bins (r_i) for ice particles ranging between 2 and 900 nm. The model timestep is 100 s. Ice and water vapor within CARMA is transported by the mean vertical wind, eddy diffusion and particle sedimentation.

- For this work, we use the same assumptions and settings as described for the CARMA reference case in Rapp and Thomas (2006) except our method differs in the specification of the background temperature profile. The reference case as described in Rapp and Thomas, forces the temperature altitude profile to be constant throughout a 48-h simulation. Their simulation did not take into account the changing background
 environment that occurs on timescales shorter than 48-h. To simulate the dynamical variability of the background temperature and the effects of changing temperatures on PMC ice particles, we used temperature time histories from simulations from the 3-D
- coupled chemistry-climate Whole-Atmosphere Community Climate Model (WACCM) (Garcia et al., 2007) as input to the CARMA model. In effect, our CARMA simulations were updated with a new temperature altitude profile at each timestep. The tem-
- perature histories were obtained by recording a high latitude mesosphere air parcel's movement over 4-days from a WACCM simulation. A temperature altitude profile (associated with the location of the air parcel) at each WACCM time-step was recorded. We compiled fifteen 4-days temperature histories from WACCM and used them as in-
- ²⁰ put to CARMA. A WACCM simulation over 4-days was used to capture several different dynamical timescales and to encompass particle development under hydrated and dehydrated conditions. It is noted that subgrid-scale gravity waves are not included in the simulations. The water vapor profile to initialize each CARMA run was obtained from the first time step of each WACCM time history. We note that the history curtains
- were obtained from a version of WACCM that does not contain PMC microphysics or dehydration in the mesosphere region. Therefore all the initial water vapor profiles are hydrated. Figure 1 shows the geographic trajectories that make up the WACCM temperature time histories used in this study. Each time history has a different geographic start point indicated by the diamonds in the figure. By using temperature histories as

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input to CARMA, we have indirectly simulated the changing background temperature due to atmospheric dynamics.

An example of a temperature history for one WACCM trajectory used in a CARMA simulation is shown in Fig. 2a. The temperature altitude profile varies at each CARMA

- timestep. Figure 2b shows the resulting total *iwc* calculated in CARMA for the temperature input shown in Fig. 2a. Figure 2c and d show the evolution of water vapor and effective radius. As mentioned, the first time step is forced with a hydrated water vapor profile from WACCM. In each subsequent time step, water vapor is a free variable that is modified by the microphysics, as shown in the Fig. 2c. CARMA provides a discrete pumber distribution p (r) for the 40 logarithmically appendent radius aize hims (r) ranging
- ¹⁰ number distribution $n_i(r_i)$ for the 40 logarithmically spaced radius size bins (r_i) ranging between 2 and 900 nm. We use this information to calculate an area-weighted effective radius (r_{eff}) (Wyser, 1998) at each time step and altitude level using Eq. (1). The result is plotted in Fig. 2d.

$$r_{\rm eff} = \frac{\sum r_i^3 n_i}{\sum r_i^2 n_i}$$

¹⁵ This expression for *r*_{eff} is independent of a specific observing technique and is different than the effective optical radius used by Karlsson and Rapp (2006). The latter is a measure of the optically dominant particle size, which depends on the wavelength and scattering angle of observation. We use an area-weighted radius without an observing bias to better represent the growth, sedimentation and sublimation of the particle, since these processes depend on the particle area.

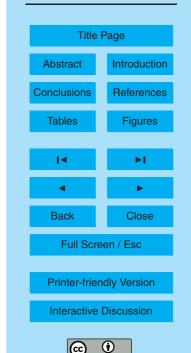
The four panels in Fig. 2 show how the distribution of ice water content, water vapor and particle size varies with the background temperature. The ice grows, sublimates and reappears as the temperature varies. This is a more realistic representation of the ice water distribution over altitude and time than the reference CARMA simula-

tion presented in Rapp and Thomas (2006) where the model was initialized with very supersaturated conditions and let to run for 48 h without changing the temperature profile. An additional advantage to including a changing temperature environment is 9, 14571–14600, 2009

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(1)



better coverage of the parameter space of T, *iwc* and the distribution of particle sizes for defining a relationship between these variables.

From the ensemble of CARMA simulations, approximately 75 000 predicted *iwc*, *T* and $r_{\rm eff}$ combinations were extracted using data at all relevant PMC altitudes. To il-⁵ lustrate the fundamental relationship between $r_{\rm eff}$, *iwc* and *T* compiled from the simulations, a surface plot of $r_{\rm eff}$ versus *iwc* and *T* is shown in Fig. 3a. Figure 3b shows the standard deviation (STD) of $r_{\rm eff}$ in each *T* and *iwc* bin from Fig. 3a. Figure 3a demonstrates that for a given *iwc*, effective radius increases with increasing temperature. This is inherent from the relationship of the background temperature with the ice particle distribution at any particular altitude. The *x*-axis (temperature) in Fig. 3a and b

- can be thought of as a proxy for altitude. Temperature is coldest near the mesopause corresponding to an altitude where the ice particles are in infancy. This altitude region relates to the lower left part of Fig. 3a and b, where the particles are generally small and have a small total *iwc*. The figures indicate that at very cold temperatures (like
- ¹⁵ near the mesopause) the PMC particles are small with very little variation. As the particles mature they fall to lower altitudes where they encounter warmer temperatures. This altitude region relates to the right part of Fig. 3a and b. The particles sizes in this area are well mixed and include larger radii with a stable standard deviation near 15–20 nm. At any specific temperature (altitude) there is a mix of particles sizes. By
- ²⁰ knowing both the *T* and *iwc* the particle size can be estimated. For example, if a line is drawn vertically at 150 K in Fig. 3a, several different groups of radii are intersected. Each radii group can be described by the *iwc* at a specific temperature. The functional form in Eq. (2) describes the fundamental relationship between r_{eff} , *T* and *iwc*. Using the data shown in Fig. 3a, a least-squares fit of r_{eff} (nm) to the following functional form was calculated:

 $r_{\rm effp} = \rho_1 i w c^{\rho_2} e^{(\rho_3 T)}$

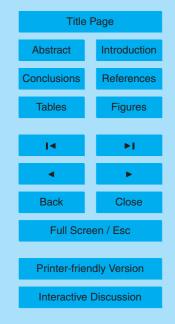
where $p_1 = 760.34$, $p_2 = 0.235$ and $p_3 = 0.034$, *iwc* is in units of g/cm³ and T in Kelvin (K). Then, using *iwc* and T pairs from the CARMA simulations and Eq. (2),

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(2)

a parameterized effective radius (r_{effp}) is calculated. Figure 4 illustrates the resulting surface plot of r_{effp} versus *iwc* and *T*. It is clear that the functional form in Eq. (2) reproduces the relationship shown in Fig. 3a. Sensitivity tests were performed to characterize the stability of the relationship. We ran simulations with the nucleation rate in

5 CARMA lowered by several orders of magnitude, the mesopause altitude displaced by 2 km, and CARMA initialized with a dehydrated water vapor profile. In all cases the relationship between these variables did not change significantly.

Figure 5 illustrates the distribution of effective radius from the CARMA simulations. The black line represents the distribution of all effective radii calculated from the ensem-

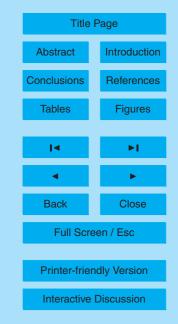
¹⁰ ble of CARMA simulations using WACCM trajectories and Eq. (1). The blue line represents the parameterized effective radius calculated using *iwc* and *T* pairs in Eq. (2). The parameterization reproduces the overall distribution of effective radii, which peaks near 20 nm.

3 Parameterization validations

15 3.1 CARMA comparisons

In this section we compare several properties of PMCs calculated as part of the ensemble of CARMA simulations to those calculated using the parameterized effective radius. Figures 6, 7, and 8 show PMC effective radius, number density and PMC optical properties as output from three of the CARMA simulations and those derived using the parameterization described in Sect. 2. Values in Fig. 6 correspond to the temperature history and *iwc* shown in Fig. 2. Figure 6a shows the evolution of effective radius as shown in Fig. 2d. In comparison, the top right panel (Fig. 6b) illustrates the parameterized effective radius calculated by Eq. (2) using the local WACCM *T* and *iwc* at each grid point in altitude and time from the panels in Fig. 2. While the error of the parameterized effective radius in 6b at any specific location in time and altitude could be rather large (in comparison to 6a), there is a striking similarity of the altitude and 9, 14571–14600, 2009

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time distribution of the effective radius, indicating that a parameterization that takes local temperature and ice water content into account, adequately models the vertical distribution of the area-weighted radius at each time step. It is this important product of the parameterization, to adequately estimate the vertical distribution of ice particle size

- at each horizontal grid point, that allows for the effective application of microphysical processes (growth, sublimation and sedimentation) to parameterized ice particles in a global model. Figures 7 (a–d) and 8 (a–d) represent the temperature history of *iwc*, r_{eff} and r_{effp} for two additional CARMA run examples, each with a dimmer PMC than the previous example. Again, the altitude and time distribution of the parameterized
- effective radius is in good agreement with the simulated distribution. The characterized relationship is able to describe the fundamental characteristics of PMC particle formation and evolution in terms of *T* and *iwc*, without needing to know the history of the ice particle. For example, in Fig. 8, a cloud is formed between the simulated hours 2 and 3. By using the local temperature structure and *iwc* at each timestep in this time range, the parameterization recreates the cloud formation around 88 km and the development
- of the cloud as the particles grow larger and move to lower altitudes.

The number density is a useful diagnostic variable that allows estimation of the observed PMC radiance to compare models to satellite observables. Given a parameterized effective radius and ice water content, it is a simple task to calculate the ice ²⁰ particle number density (N_{para}) assuming the particles are spherical:

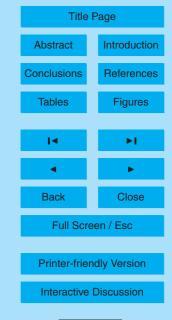
$$N_{\text{para}} = \frac{iWC}{\frac{4}{3}\pi r_{\text{effp}}^3 \rho_{\text{ice}}}$$

25

where ρ_{ice} is the density of ice (0.93 g cm^{-3}) . The derived number density can be compared to that calculated directly from the CARMA simulations (Figs. 6c, 7e, and 8e). While N_{para} , shown in 6d, 7f, and 8f are not exactly the same as the number density that comes directly out of CARMA, the vertical distribution and time histories are very similar. Since the parameterization provides a vertical distribution of particle size, the optical properties of PMCs such as backscatter coefficient (β) and albedo (A) 9, 14571-14600, 2009

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(3)

can be easily calculated using Eqs. (4) and (5),

$$\beta = \int_{0}^{\infty} \frac{d\sigma}{d\Omega} (r, \lambda, \vartheta) \cdot \frac{dN}{dr} \cdot dr$$

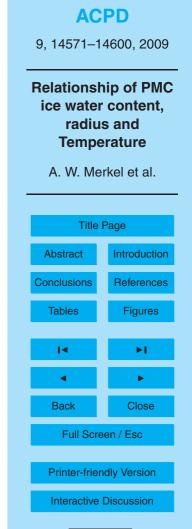
and

$$A = \int_{0}^{\infty} \int_{0}^{\infty} \frac{d\sigma}{d\Omega} (r, \lambda, \vartheta) \cdot \frac{dN}{dr} \cdot dr \cdot dz$$
(5)

- where dσ/dΩ is the Mie scattering cross-section derived applying standard Mie-scattering algorithms (Bohren and Huffman, 1983) and λ and θ are the wavelength and scattering angle under consideration. The PMC vertical backscatter coefficient at 532 nm and nadir viewing PMC albedo at 265 nm are shown in Figs. 6(e–f), 7(g–h) and 8(g–h). The "true" optical signals can be calculated directly from the discrete particle
 size distribution that comes directly from the CARMA simulations. The "true" optical
- signals are illustrated by the filled contour levels in the backscatter plots and by the solid black line in the albedo plots. Overplotted are β and A calculated using the parameterized effective radius and number density (contour lines in the backscatter plot and dashed blue line in the albedo plot). One minor draw back to estimating the optical
- ¹⁵ properties (β and A) using the parameterized r_{effp} and N_{para} , is that the particle size distribution (dN/dr) must be assumed. From an analysis of the particle size distributions that come out of the detailed CARMA simulations and using the width as a tuning parameter we determined that a Gaussian distribution (Eq. 6) with a 12 nm width (Δr) recreates the "true" optical properties.

²⁰
$$\frac{dN}{dr} = \frac{N_{\text{para}}}{\sqrt{2\pi}\Delta r} \exp\left(-\frac{\left(r - \langle r_{\text{mean}} \rangle\right)}{2\Delta r^2}\right)$$

An error is introduced by using r_{effp} (area-weighted parameterized radius) instead of a number weighted mean radius to calculate backscatter and albedo. To minimize



(4)

(6)

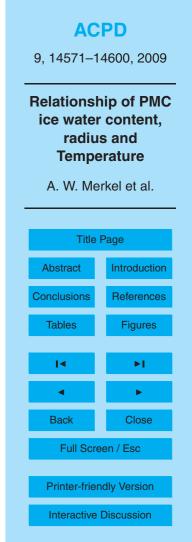


this error, an estimated mean radius is calculated as a function of $r_{\rm eff}$ from the ensemble of CARMA simulations. Using linear regression we calculated $r_{\rm mean}$ = 0.9037 $r_{\rm eff}$ -3.18 nm, and use it in the calculation in Eq. (6).

Although the error in the estimated β and A might be large at any specific timestep, the three trajectory examples demonstrate that the time history can be adequately recreated using the parameterized effective radius and number density. This has the potential to be very useful in comparing a bulk PMC parameterized GCM to satellite PMC radiance measurements such as those taken on the AIM satellite.

3.2 Comparison with AIM/SOFIE

- The Solar Occultation for Ice Experiment (SOFIE) on the AIM satellite uses the technique of satellite solar occultation to measure vertical profiles of limb path atmospheric transmission (Hervig et al., 2009). Their measurement technique allows for the first reported independent concurrent altitude profiles of PMC effective radius, *iwc* and temperature from a satellite. To validate the parameterization we use all the combinations
- ¹⁵ of r_e , *iwc*, and *T* above 79 km for SOFIE's 2007 Northern Hemisphere season (Data Version 1.022), similar to that described in Sect. 2, Fig. 3a for the ensemble of CARMA simulations. Figure 9a illustrates the surface plot of the relationship of these parameters as measured from SOFIE. While the absolute values are slightly different (SOFIE temperatures are slightly warmer and *iwc* slightly less than the CARMA values shown
- in Fig. 3) the relationship between the variables is very similar to that determined by the CARMA simulations. Figure 9a, like Fig. 3a, illustrates that for a given *iwc*, effective radius increases with increasing temperature. Figure 9b shows the standard deviation of SOFIE effective radius. It has the same general shape as Fig. 3b, where the particles are very small with not much variation at the cold extreme temperatures. The particles
- are bigger and mixed at the warmer temperatures with a standard deviation around 15–20 nm. Since SOFIE data is obtained from concurrent altitude profiles of temperature, *iwc* and effective radius, the inherent relationship of the background temperature to the effective radius distribution is upheld. Figure 9 demonstrates that the compact





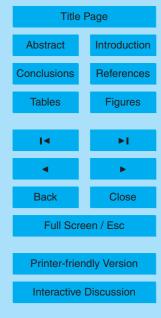
relationship between these variables is observed in SOFIE PMC measurements and validates the parameterization. Although not shown here, all four available seasons of SOFIE data (NH and SH 2007; NH and SH 2008) were analyzed. All seasons demonstrated the same stable relationship as that illustrated in Fig. 9.

- ⁵ We compare the particle distribution calculated from the CARMA simulations to the particle distribution measured from SOFIE to further validate the modeled effective radius distribution. In addition we compare to the effective radius distribution measured from a lidar at ALOMAR. Figure 10 illustrates the comparison. The black and blue curves are identical to those shown in Fig. 5 for CARMA. The red curve shows the distribution of SOFIE effective radius for the whole 2007 northern hemisphere season.
- ¹⁰ distribution of SOFIE effective radius for the whole 2007 northern hemisphere season. All effective radii above 79 km are used, as advised in Hervig et al. (2009). The green curve shows the effective radii measured from the ALOMAR lidar (69° N) during the 2007 summer season (Baumgarten et al., 2008; Hervig et al., 2009). At first glance, the distributions might appear dissimilar, however there is excellent agreement of the
- ¹⁵ distributions above 40 nm. In this particle size region, observational techniques have the least error in measuring PMC particle size. Below 40 nm, the distributions slightly diverge. At the smallest particle range CARMA predicts that most of the particles lie in the 10–30 nm range, while the SOFIE and ALOMAR distributions drop off (Hervig et al., 2009). This is not surprising since the CARMA distribution is not limited by an
- ²⁰ observing technique and it is in this particle size range that observations are the least sensitive to the true particle size. Although, all the distributions have slightly different peaks, the ensemble from CARMA seems to characterize the full range of ice particle size and agrees very well with the observations from the middle to largest particles sizes. It is important that the larger particles are estimated accurately since these
- ²⁵ dominate the optical signal and contain most of the ice mass. It is clear from Fig. 10 that the distribution of larger particles is accurately represented using the effective radius parameterization.

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4 Implementation in the WACCM model

While the characterized relationship between *iwc*, T and r_{eff} can be used for a number of possible applications (i.e. estimating T from satellite measurements of r_{eff} and *iwc*), our purpose is to use it to incorporate a bulk PMC parameterization into WACCM to run

decadal scale PMC simulations. While a detailed description of how we incorporated the parameterization into WACCM (Marsh et al., 2007) is beyond the scope of this paper, we briefly describe the process and show a comparison to PMC data from the AIM satellite, to illustrate its potential use. The version of WACCM used in this study has an enhanced vertical resolution in the mesosphere-lower thermosphere region (grid spacing ~0.3 km) and 2° × 2.5° horizontal resolution. We refer to this version as WACCM-PMC.

There are many ways one could implement a nucleation process within a global model, which combined with the parameterization described in this paper, could be the basis of a PMC bulk parameterization. For example, one can take a simple approach

¹⁵ whereby a very small percentage of the available water is converted to ice when the air becomes extremely super saturated. Another approach could be to calculate the dust distribution and use that as a limit on the PMC formation rate. Yet another example could be to assume large sulphate particles as a source (Mills et al., 2005) and track their density to limit the initial ice formation. For our purposes we implemented the first nucleation approach into WACCM as the basis of our PMC bulk parameterization.

The process is started by the nucleation scheme chosen. At each subsequent timestep, using the local *iwc* and background T, we calculate the effective ice particle radius (r_{effp}) using the parameterization described in Eq. (2). The effective radius is then used to calculate the sedimentation rate and the growth rate of the particle. Growth of a particle is defined by,

 $Q_{\text{growth}} = A_{\text{ice}} \rho_{\text{ice}} \rho(\partial r / \partial t)$

25

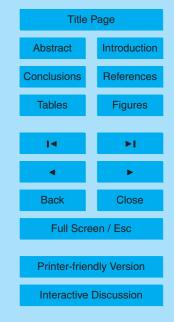
where, A_{ice} is the total surface area of the particle, ρ_{ice} is the density of ice (0.93 g cm⁻³), ρ is the density of air, and $\partial r/\partial t$ is the growth rate defined by 14583

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(7)

Gadsden (1998). The growth rate depends on the saturation ratio, background T and P_{sat} (the pressure of water vapor over ice). The total surface area (A_{ice}) of the ice particle is:

$$A_{\rm ice} = N_{\rm para} 4\pi r_{\rm effp}^2$$

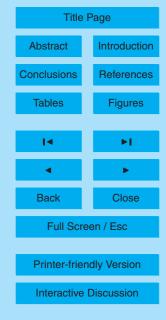
- ⁵ where the total number density (N_{para}) is calculated using Eq. (3). Using the parameterized radius, Q_{growth} modifies the *iwc* carried to the next time step. The process repeats itself, a r_{effp} is determined from the *iwc* and *T* at the next time step and processed through the microphysics. This is a very effective process since particle size bins are not stored at each time step. WACCM-PMC model runs provide PMC products such as ice mass, particle radius, and ice number density on a global scale at any specified time step. From these products a cloud albedo and frequency can be calculated to compare to satellite observations such as the Cloud Imaging and Particle Size (CIPS) experiment on the AIM mission. An example of this type of comparison is shown in Fig. 11.
- Figure 11a shows a global map of PMC albedo calculated from a daily snapshot (00:00 UT) of a WACCM run with the bulk PMC parameterization included. The cloud albedo is calculated using Eqs. (4), (5), and (6) at a scattering angle of 90° at 265 nm. In comparison, Fig. 11b shows the detected PMC albedo, the ratio of atmospheric radiance to solar irradiance in units of 1×10^{-6} sr⁻¹ at 90° scattering angle, from one day
- of CIPS data at 12:00 local time (Bailey et al., 2009; McClintock et al., 2009; Rusch et al., 2009). Since WACCM is a free running model the daily snapshot shown will indeed be different; however the global distribution of PMC occurrence and albedo match considerably well, giving confidence in the parameterization. In addition, this figure shows that a parameterized effective radius can be used to successfully estimate the optical parameterized of a built BMC.
- ²⁵ properties of PMCs, and demonstrates the potential of a bulk PMC parameterization in a multi-dimension global climate model.

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5 Summary and Conclusions

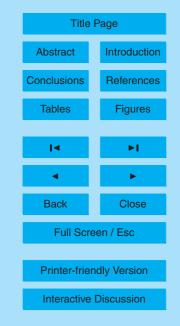
The means of including a representation of PMC microphysics in a multidimensional global climate model without using a Lagrangian or sectional microphysics code framework has been developed. This bulk PMC parameterization allows for theoretical ice cloud microphysics to be applied to ice particles to simulate growth, sublimation and 5 sedimentation without keeping track of the time history of each ice particle or particle size bin. This is accomplished by defining and parameterizing the relationship of the PMC effective radius to the local T and iwc. The relationship was characterized using an ensemble of 1-D CARMA simulations forced by temperature histories produced from the WACCM model. The characterized relationship is stable and is able 10 to describe the fundamental characteristics of PMC particle formation and evolution in terms of T and *iwc*, without knowing the history of the ice particle. Because the parameterization uses altitude profiles of both local *iwc* and T, it is able to adequately estimate an altitude profile of particle size at each grid point. It also allows for an estimate of the observed optical properties of PMCs to be compared to satellite and ground-based measurements. The proposed parameterization has the potential to be

used in any multi-dimensional climate model to investigate the drivers of PMC global variability without extensive computational costs. The relationship has the potential to be a very powerful tool in helping to characterize satellite and ground based observations of PMCs.

PMC modeling studies that include a Lagrangian work frame such as the COMMA model (Berger and Lübken, 2006; Lübken and Berger, 2007) show that PMC particles have a complicated history of variation. They nucleate, grow, sublimate, advect, and are transported all over the pole on relatively short time scales. The parameterization described here does a good job of estimating the instantaneous particle size without modeling the particle history because of the stable relationship between *iwc*, *T* and effective radius that PMCs seem to encompass as shown in Figs. 3 and 9. The history of the ice particle seems to be contained in the instantaneous *T* and *iwc* at any given time

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step. For example, if the local temperature is relatively warm, but the local ice mass is relatively high, it is likely the air parcel trajectory passed through a region of cooler temperatures that led to significant particle growth and a large mean radius. Comparisons of the parameter space from CARMA simulations used to calculate the effective

⁵ radius parameterization with that of SOFIE, shows that the relationship between these variables, appears in PMC measurements. The fact that the SOFIE measurements exhibit the same relationship between these variables gives confidence in the CARMA model simulations and the parameterization. The compact relationship between these variables could possibly be used to help interpret observations of PMC ice mass, radius and mesospheric temperature in past and future experiments.

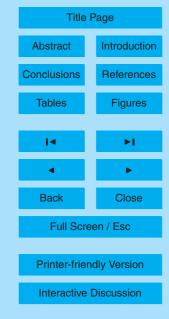
The parameterization has been included into WACCM and produces reasonable ice mass, particle radius, and ice number density distributions on a global scale. The details of the implementation in WACCM will be described in a future publication. Initial comparison of output from the WACCM model (with the PMC parameterization in-

- ¹⁵ cluded) shows that it is comparable to the CIPS measurements from AIM. Future studies will utilize this parameterization to support long-term runs to study climate change and its effect on cloud presence. Modeling efforts are crucial to understanding how AIM PMC observations relate to longer-term changes resulting from solar and anthropogenic forcing.
- Acknowledgements. This research was supported by NASA grant #05-SRT05-70 conducted at the National Center for Atmospheric Research. The National Center for Atmospheric Research is sponsored by the National Science Foundation. Any opinions, findings and conclusions or recommendations expressed in the publication are those of the author(s) and do not necessarily reflect the views of the National Science Foundation. Thanks to the AIM team for providing CIPS and SOFIE data. Thanks to Thomas G. Busch D. Bardeen C. Hervin M. and Baum-
- ²⁵ CIPS and SOFIE data. Thanks to Thomas, G., Rusch, D., Bardeen, C., Hervig, M., and Baumgarten, G., for collaboration and insight.

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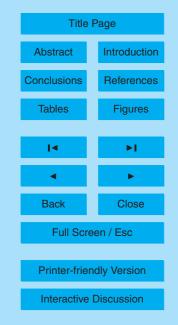
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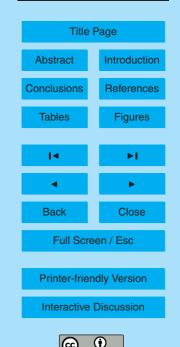
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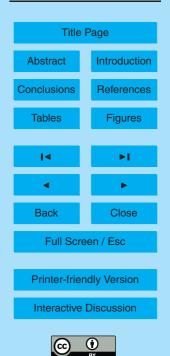
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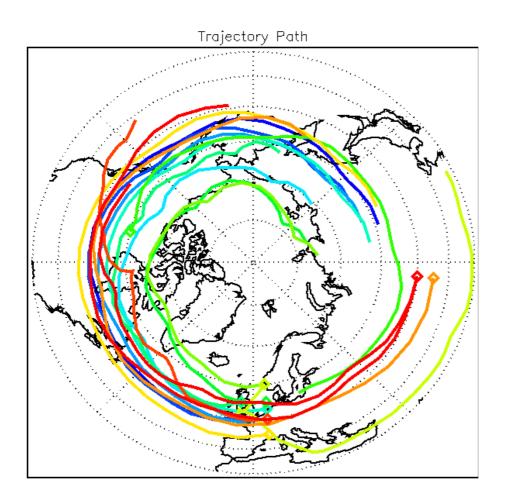
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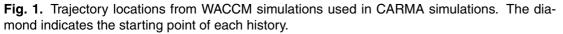
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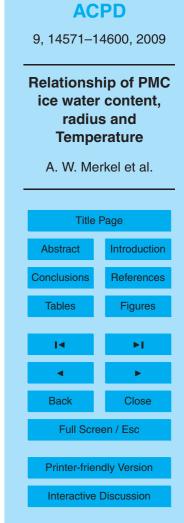
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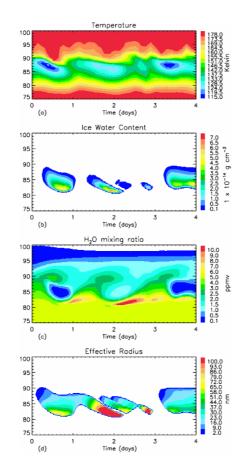


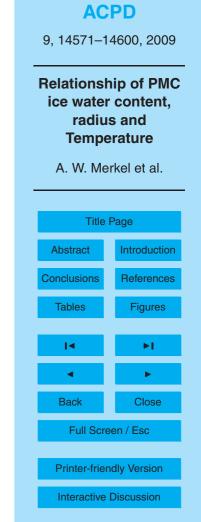












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Fig. 2. (a) Temperature time history from a WACCM trajectory simulation. (b) ice water content from the CARMA simulation associated with the temperature input in the first panel. (c) Evolution of water vapor, the first time step is forced with a water vapor profile from WACCM. (d) Evolution of effective radius.

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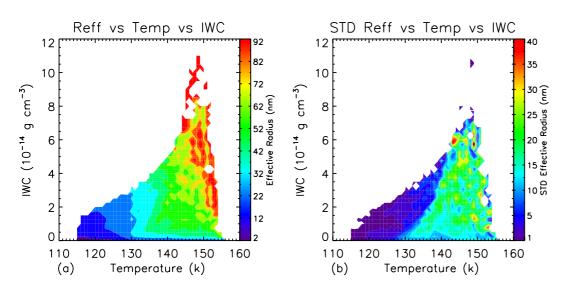
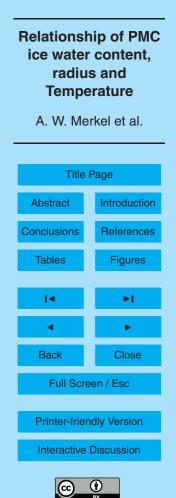
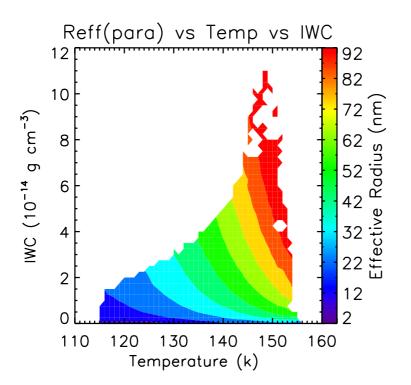
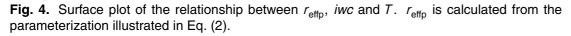


Fig. 3. Surface plots of the relationship between r_{eff} , *iwc* and *T*. (a) r_{eff} calculated directly from the particle distribution at each point in the CARMA history files. (b) standard deviation of r_{eff} in each *T* and *iwc* bin.



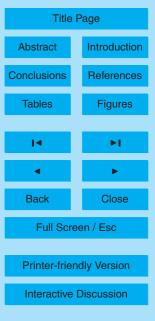




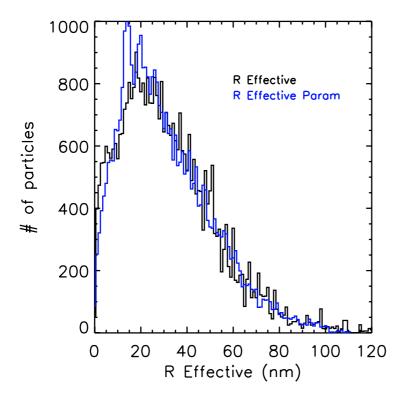
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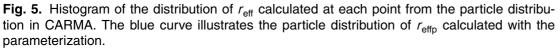
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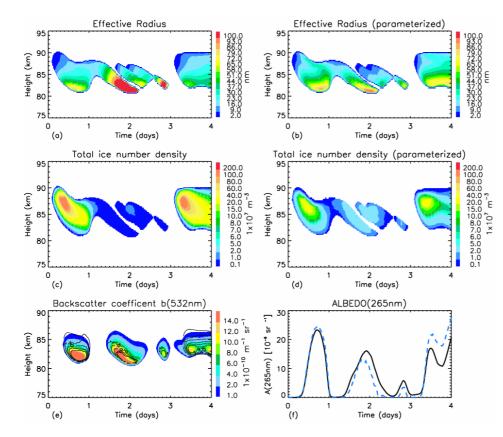
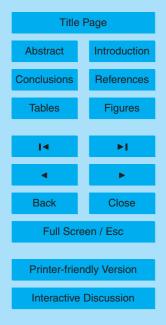


Fig. 6. Panel plots illustrating the effectiveness of the parameterization. The panels correspond to the simulation shown in Fig. 2. The left column illustrates cloud properties calculated from the output of the CARMA simulation with full microphysics and particle distribution. The right column illustrates properties calculated with the parameterized effective radius. The bottom two panels illustrate PMC optical signals. The signal calculated directly from the CARMA simulations are illustrated by the filled contour levels in the backscatter plot and by the solid black line in the albedo plot. Overplotted are β and A calculated using the parameterized effective radius and number density (contour lines in the backscatter plot and dashed blue line in the albedo plot).

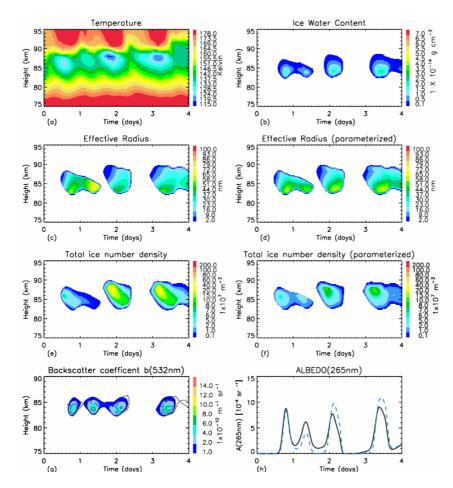
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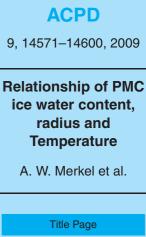
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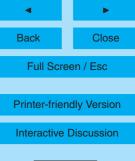




Fig. 7. Same as Fig. 6 but an alternate trajectory. The top two panels illustrate the temperature and *iwc* history.

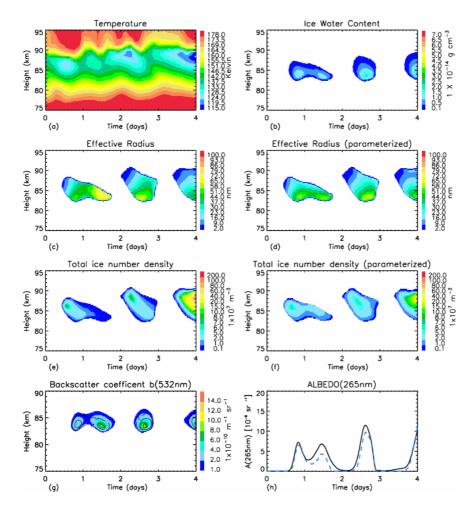
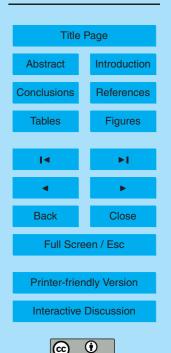


Fig. 8. Same as Fig. 7 but an alternate trajectory.

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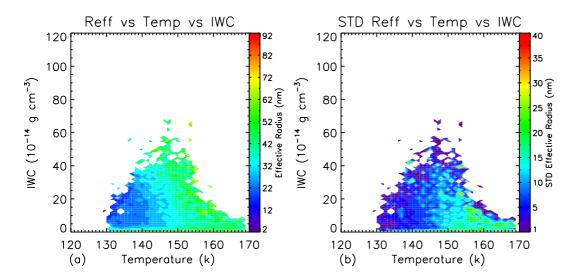
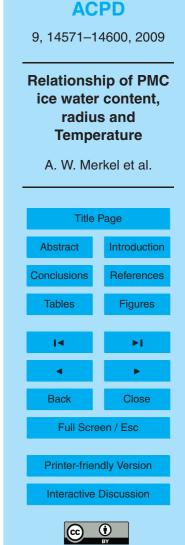


Fig. 9. Surface plot of SOFIE data showing the relationship between *T*, *iwc* and effective radius.



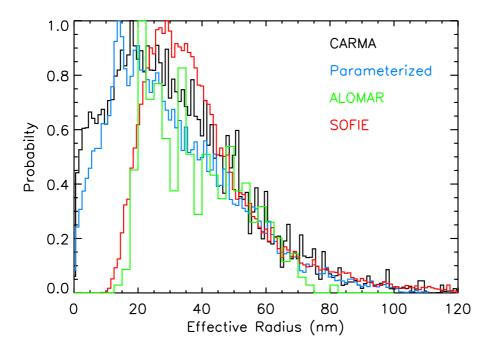
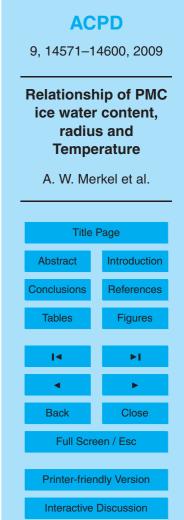


Fig. 10. Histogram of effective radius from recent satellite and ground-based measurements compared to particle distribution used to define parameterization. SOFIE NH 2007 measurements are depicted by the red line and ALOMAR lidar observations are depicted by the green line.





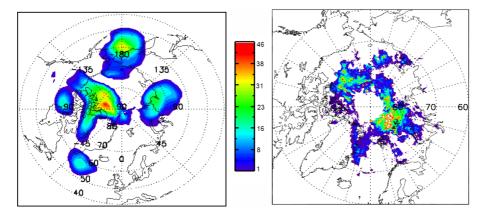


Fig. 11. PMC albedo at 265 nm wavelength in 1×10^{-6} str⁻¹ albedo units. The left globe illustrates the WACCM PMC albedo for 3 August 2007 at 00:00 UT. The right globe shows the PMC albedo from the CIPS instrument at 12:00 local time 3 August 2007.

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