## Response to all referee comments

We would like to thank both reviewers for their help in improving the manuscript. This response is structured such that we first list some *general changes made to the manuscript* (in red), followed by the referee comments with the corresponding *responses by the authors* (in blue) and where needed *the changes made in the manuscript* (again in red). The comments of referee #1 are given first followed by the ones by referee #2.

# General changes

- a) We realized that an error was made in the post-processing of the instantaneous surface precipitation now shown in Fig. 12, formerly Fig. 7. This resulted in too high rates. Now the distributions of the frequency of occurrence of the two simulations are more alike, although the prognostic scheme has a broader distribution, the reduction in the peak of low precipitation rates is still visible. The results have been adjusted accordingly.
- b) To address the changes in the hydrological cycle, we have added a Figure (Fig. 13) to the manuscript, illustrating all relevant conversion rates.
- c) Section 3.2.3 has been slightly rearranged for the analysis of the conversion rates. Now we start with the individual conversion processes and then present the global picture of the hydrological cycle. We hope that this makes the text more reader-friendly.

## Comments by anonymous referee #1

However, I have a few general concerns, and think this should be published in ACP with some important clarifications and revisions. First of all, there is not enough discussion of sensitivity of the scheme: what is the drizzle mode doing? That is really the new part here, and there is no sensitivity test that just focuses on the drizzle mode: does it matter in going from a 2 class to 3 class prognostic precipitation scheme. In general, as noted below, a few more sensitivity tests focusing on the main science questions would be helpful.

Agreed, but we would like to note that such a sensitivity was done in the publication of Sant et al. (2013). Nevertheless, with respect to the model version at hand, we see that the issue has not been discussed enough and we add a paragraph focusing on going from 2 to 3 classes.

We have added a few more sentences in the introduction with respect to the work done in Sant et al. (2013) and have added some text in the conclusions, putting the results in the context of using 2 or 3 classes.

Also, the aerosol indirect effects (AIE, or Aerosol-Cloud Interactions, ACI) are sort of a null result: less sensitivity to aerosols is shown diagnostically, but this does not seem to be the case in the simulations performed. The uncertainty levels (25% or ACI) are higher than any signal (10%) and it is of the wrong sign. I think breaking out the ACI by regime may be illuminating: why is the new scheme less sensitive, but has larger ACI (FNet).

We suppose that the reduced sensitivity to the aerosols you mention refers to the change in cloud liquid water path (CLWP) to changes in CCN. Note that this is based on process level and does not necessarily allow direct conclusions on the ACI. First, it does not include all clouds and one cannot predict how this influences the climate system as a whole. Nevertheless, the lower sensitivity did result in lower CLWPs allowing more shortwave radiation to be absorbed at TOA.

Unfortunately, we do not understand how you estimated an uncertainty of 25 %, but we see that the uncertainty in FNet is rather large. To this end, the simulations were extended.

We agree that breaking out the ACI by regime, i.e. contributions from different cloud regimes, geographical areas, etc., would be the next step in identifying reasons for the larger ACI. This will have to be part of future work.

We have extended the simulations from 5 to 12 years, which has slightly reduced the uncertainty.

With some more sensitivity tests and a bit more analysis of the details of the results, I think this will make a nice publication.

We hope that we have fulfilled the reviewers expectations and thank her/him for the comment. We have restructured section 3.2.3 and have added a paragraph on the hydrological cycle, investigating in detail the changes in source and sink rates. For this we have added an illustrative figure.

## P7784, L15: more realistic how? What metrics?

We believe that going from a diagnostic to a prognostic treatment of precipitation at time steps of around 10 min allows for a more 'realistic' representation, in what the word means, i.e. a more accurate representation of the truth. In this study (and in others) microphysical processes, such as the collection process, are better represented - e.g. the ratio of autoconversion to accretion. With a diagnostic scheme, this cannot be achieved.

With respect to the wording, we have tried to improve the English and hope that the explanations are more precise.

P7786, L10: has led to a negative... *Done.* 

P7786:L15: GCMs have been shown to overestimate...these aerosol-cloud- precipitation *Done*.

P7789, L9: the tuning is unclear to me: you imply nothing was tuned. Does that mean that both models have the same radiative balance? Or did you retune the model? If they are not in the same balance, please state that, or if tuning was done state that.

The CTRL version of the model was tuned to be in radiative balance. With the introduction of prognostic precipitation and the three liquid water classes we did not alter the tuning parameters - as not to impede the effects of the prognostic scheme. However, we wanted to note, that by implementing a new liquid water class and consequently a new warm rain scheme (for the collection rates), we inevitably changed the parameter space for tuning the model (since the autoconversion rate includes a tuning parameter). This has (obviously) led to some confusion, so we simplified the paragraph as follows:

With the introduction of the new collection scheme for the three liquid water classes and the prognostic precipitation scheme, the climate model is not necessarily in radiative equilibrium anymore as the parameter space for the tuning is changed (Lohmann and Ferrachat, 2010). To preserve comparability of the two model versions the tuning parameters were not changed. Yet, if the climate impacts for both liquid and ice phase microphysics between CTRL and PROG are comparable in terms of heating rates (Mauritsen et al., 2012), a similar climate can be expected.

P7794, L21: some further questions: does the time splitting vary or is it fixed in both the calculation and sedimentation loops? I'm guessing fixed from below. Is the run time total

model run time increase, or micro physical parameterizations only? Also: it would be nice to show some plots of autoconversion and accretion rates if you do not do it later.

The time splitting is fixed and is set at the beginning of the simulation. However, the sedimentation itself can be made subject to an extra time splitting, which would also be fixed. This could be deduced from the schematic flowchart shown in Fig. 1, but we have added this point to the model description.

The factors given of increased computational costs are given for the whole model, i.e. the atmosphere. This has been added.

We do show some plots of autoconversion and accretion rates later.

P7794,L6: so the high frequency data are from month 4 of the simulation? Yes.

That seems perhaps minimal for spin up of the land surface. Does it matter?

As we are not using an interactive land surface model (in our case JSBACH), this should not matter. Furthermore, we are only looking at instantaneous variables on the process level, the statistics will not change. If one has a few more or less clouds this will not influence the way the processes behave. Nevertheless, we will check this.

To make sure we reran the simulation for the whole year 2000, and no matter which month we chose, it always led to the same conclusion and has been stated. Nevertheless, to circumvent possible concerns, we chose the month of July instead of January.

If you only run 5 years and have large enough variability that you cannot distinguish a 15 % change in ACI, perhaps you should run longer. Another option is to remove interannual variability in SST and sea ice and fix one year or a climatology.

We do use climatological SSTs.

As mentioned, the simulations have been prolonged from 5 to 12 years.

P7796, L10: is the difference in cloud cover between CNTL and PROG significant? The same changes are still seen in the new prolonged integration period. However, we point out that the differences are small.

P7797, L2: I'm not sure your impression of the hydrologic cycle lifetime is correct. PROG has less LWP, so one would expect less cloud, and you have not constrained the in cloud LWP. I do not think you can argue that fewer clouds = faster hydrologic cycle.

Our further analysis proved us wrong and showed that the speed of the hydrological cycle has not changed. We thank the reviewer again for pointing this out to us and for (indirectly) encouraging us to pursue this in more detail. As mentioned, we have added a paragraph and an illustration of the hydrological cycle which offers a very detailed insight into the changes induced by the prognostic scheme.

P7797,L24: this clarifies my earlier point about tuning. *Ok.* 

P7798,L23: I do not think the re analysis LWC and IWC should be used for comparisons: they are not that well constrained and are the product of model physics in a different model.

However, it is difficult to get reliable observations of such height resolved variables. We agree that these are also based on model physics of a different model, but they do include observations and therefore act as a sort of guideline.

We have observations of IWC and have include these. Concerning the re-analysis we have emphasised that it is not to be taken as the truth, but that it shows a suitable and plausible

#### alternative.

P7799,L10: could you test this theory about redistribution through precipitation by a sensitivity test to change the fall speed of precipitation?

The idea is good and such a test could be done. However, the redistribution of the CIWC and the CLWC is not only coupled to the fall speed of the precipitation, but in particular to the shift in the collection rates. Especially, the collection rates for the ice phase will be revisited in future work (as we emphasise at the end of the manuscript) to better quantify the response in ACI.

Can you estimate the mass of precipation in the control model to compare precipitation?

Well, the amount of precipitation is nearly the same in both model versions, consequently we unfortunately do not understand the suggestion made.

We hope that for both points, the new paragraph on the hydrological cycle helps to show the importance of the different source and sink rates on the precipitation fields and vice versa.

P7800,L10: can you estimate the fraction of precipitaton formed that remains in the atmosphere? It would probably be the ratio of the PRECIP/tendency of formation. This would be a good way to estimate the importance of prognostic precipitation, and the difference with CNTL.

We believe that this is now nicely shown in the new illustration of the hydrological cycle.

P7800, L20: for figure 7, what region is this?

It is global and we now mention this.

It would be nice to divide this into different regimes: stratus regions, stratocumulus, storm tracks, continental mid-latitudes. A single global bar does not do this justice.

We have checked, but the overall conclusions remain the same. Naturally some small features change, but not the overall picture.

This point has been added to the text.

Also, what about comparison to observations (e.g. TRMM)?

This cannot be compared to observations as we would need instantaneous values, not daily means or even monthly means.

Also: have you shown the maximum? That is not really clear.

The point concerning the maximum is not quite clear. What maximum do you refer to?

No changes made.

P7801,128: but isn't this also a function of the width of the drop size distribution? Is that constant here? Having a few large particle may matter quite a bit?

Yes, the width of the drop size distribution, i.e. the presence of large drops, is very relevant for the autoconversion and accretion rates. However, in a bulk microphysics scheme, the width is often assumed constant or not part of the parameterization. The latter is the case for CTRL. In PROG the drop size distribution is better captured. Although the width is assumed constant for the cloud droplets, the scheme takes the tail into account, which leads to the threshold type behaviour mentioned. For effective radii below 10  $\mu$ m, the tail of the cloud droplet size distribution is small, such that the relatively high autoconversion rates for low effective radii exhibited by CTRL are not physical. We hope that answers the question.

No changes made.

P7802, L5: are these aggregation and accretion rates for ice or liquid or both: it is not clear from the text or notation ( It looks like accretion due to snow). What is ACi?

It is the accretion due to snow. The notation is introduced in the line above and relates to the processes introduced in section 2.2. Table 3 does also relate ACi again to the processes

summarized in Table 1.

No specific changes made. We hope that some of the improvements in the English, the structure and the added text help.

P7803,L28: what is the uncertainty on these regression lines? Also, does the regression vary by region or cloud regime?

We will include the correlation coefficient and the number of data points to give an idea of the statistical variability. Most of the CLWP is from clouds in the lower levels, consequently looking at different regimes would be somewhat artificial. We wanted to compare our simulations to the work by Wang et al. (2011), hence the global view. We will check the regional dependence. The figure has been modified accordingly and no regional dependence was found.

P7894, L16: can you show a map of FNet and the differences?

We could, but it is rather noisy and hence does not provide more insights.

No changes made.

How much of this is direct effect: I.e. Only the cloud part is shown in Figure 13.

The direct effect is hard to isolate (see Lohmann et al. (2010)). Thus, we prefer to only estimate  $ERF_{\rm aci+ari}$ .

This has been included in the text.

P7804, L21: what is 'this' referring to?

It refers to the difference in FNet between the two models CTRL and PROG.

We agree that the sentence is not clear enough and have rephrased it.

The explanations here are not that convincing. If the sensitivity goes down, and the accretion goes up relative to the total, why are the aerosol-cloud effects on radiation increasing? You might need to do some sensitivity tests related to precipitaton.

We have tried to improve the text with respect to the changes in radiation and in particular Fnet. We hope that this has improved the understanding for the reader.

P7804, L24: what happens to precipitation in the radiation code? Is it treated or does it not exist for purposes of radiation?

Precipitation is not seen by radiation. The sensitivity to this will also be addressed in future work.

This has been emphasized in the text.

P7895,115: a bit more on why this is valuable to have a drizzle class would be warranted: what really justifies it? Could you construct a sensitivity test without drizzle?

We have added some explanations and relate to the use of 2 prognostic liquid water classes.

P7805,L23: if you tried this, then I would probably show it. You speculate it is important, but then show it is not?

We have added some more motivation for having a drizzle class and for the type of applications it is profitable. The reference to preliminary results specifically using the drizzle class has been taken out, in particular as it is beyond the scope of the paper and to avoid confusion.

P7806,L17: versatile is not the right word here. Variable? *Yes, done.* 

## Comments by anonymous referee #2

## **Major Comments:**

Page 7784, line 16. Run the model long enough or, better yet nudge winds to a common analysis, to improve statistical significance.

The model runs have been extended from 5 to 12 years. This led to a slight decrease in variability.

Since the aerosol optical optic depth is changed significantly, this could affect the aerosol radiative forcing ERFari. Please add a diagnostic radiation calculation without aerosol to partition the ERF into ERFari and ERFaci (e.g., Ghan, 2013).

This would be a possibility, but differentiating between ERFari and ERFaci is difficult as both effects act at the same time (cf. Lohmann et al, 2010, ACP).

This point has been addressed in the text.

The manuscript has far too many problems with the English for a reviewer to identify and suggest all corrections. I've identified twelve corrections and improvements for the only the first two pages. If the corresponding author cannot improve the English in the rest of the manuscript, the coauthors should get more involved.

We note that most of the corrections are related to the style of writing. We are sorry to hear that the reviewer finds the style inappropriate and we will improve the manuscript as a whole. We thank him/her for the suggested corrections.

We have worked over the manuscript as a whole and hope that it has led to an improvement.

More clarity is needed in several places, including whether all microphysical modes contribute to cloud optical properties, and how those optical properties are represented.

We have added that the cloud liquid and ice water are the only contributors to the cloud optical properties and hope that this clarifies the issue.

#### **Minor Comments:**

Page 7784, first sentence. I recommend reordering this sentence so the dependent phrase follows the main part of the sentence.

Done.

Page 7784, line 5. Insert comma after and.

Done.

Page 7784, line 8. Replace "phase" with "phases".

No, this would change the meaning. We refer to the physical definition of phase, i.e. water being either in the liquid or the ice phase.

No changes made.

Replace "the prognostic" with "a prognostic". *Done.* 

Page 7784, line 10. Remove "does".

#### Done.

Page 7784, line 24. Remove "the" after "towards". *Done.* 

Page 7785, line 6. Replace "may" with "whether". *Done.* 

Page 7785, line 18. Replace "in cloud droplets" with "in the number and reduction in size of cloud droplets".

Done.

Page 7785, line 19. Replace "cloud" with "in warm clouds". Here 'cloud' is used in conjunction with droplets, i.e. 'cloud droplets'. No changes made.

Page 7785, line 20. Remove "upon".

Page 7785, line 21. Insert "produces" before "a significant".

Page 7785, line 22. Remove "is produced".

Done.

Page 7785, line 28. Start new sentence at "thus". Start new paragraph that focuses on mixed phase clouds.

Done.

Page 7788, line 15. Citing Ghan and Easter, Mon. Wea. Rev. (1992) would be useful here. *Thanks*.

Done.

Page 7790, line 20 - Page 7791, line 12. While the distinction between cloud water, drizzle and rain has a basis in physical processes, the distinction is much less clear for frozen water (Morrison and Grabowski, JAS, 2008). This issue should be discussed here. *Done.* 

Page 7794, lines 19-20. Insert "the computational cost" after "reduces" and make clear how much of the model this refers to (collection, the whole atmosphere, the coupled system). *Done.* 

Page 7796, line 11. What does "large" mean? Significant (statistically)? If so, demonstrate it. Or just noticeable? Do these paths include drizzling and precipitating particles as well as cloud particles?

'large' refers to the discussion which follows the 7 % change in cloud liquid water path. In the next paragraph we use the word 'significant' (p. 7797, line 3, not used on 7796) in the sense of large or noticeable - as suggested. We will clarify these issues.

Concerning the LWP and IWP, they only include the cloud water, i.e. the precipitating water is not included. This will be clarified by using CLWP and CIWP for cloud liquid water path and cloud ice water path, respectively.

The wording used in the manuscript has been worked on and we hope that it is more precise now. The terms LWP and IWP have been clarified by using CLWP and CIWP throughout the

## text and in all figures.

Page 7797, line 1. Is the hydrologic cycle faster if the precipitation rate and total water path don't change? Since you're discussing condensed water, not total water, I would say the cloud lifetime is lower.

As we answered to reviewer #1, our further analysis proved us wrong and showed that the speed of the hydrological cycle has not changed. We thank the reviewer here again for pointing this out to us and for (indirectly) encouraging us to pursue this in more detail. We added a paragraph and an illustration of the hydrological cycle which offers a very detailed insight into the changes induced by the prognostic scheme. We hope this clarifies the issue.

Page 7797, line 3. How do you know the differences are significant? We mean 'noticeable'. This will be edited (see answer for p. 7796, line 11). This has been changed.

Page 7797, line 10. How does cloud lifetime affect SWCRE? Do you mean cloud fraction? We have tried to clarify the points related to the changes in radiation and hope this has improved the text.

Page 7797, lines 17-19. Does your model treat the contributions of drizzling and precipitating particles to cloud optical depth? It was found to be important in CAM5. You should describe how the optical properties are determined. Are the same shapes assumed for all frozen modes? Only cloud liquid and ice water contribute to the optical properties of the clouds as has been emphasized in the text.

Section 3.2.2. Do the ice and liquid water contents include contributions from the drizzling and precipitating modes?

No.

Clarified in manuscript.

Page 7802, lines 3-21. Further discussion of Morrison and Grabowski (2008) would be useful here.

Done.

Figure 13. Please show standard error about the mean values.

Is Figure 14 meant? It would make the figure even more crowded and the information gained would not be very relevant as we already show the spread of the data. Nevertheless, we will add the correlation coefficient of the regression line and the number of data points, which should give some idea of the statistical variability.

Figure 14 has been improved.

Page 7806, lines 16-24. This might be a place to mention Morrison and Grabowski again. *Done.* 

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# Prognostic precipitation with three liquid water classes in the ECHAM5-HAM GCM

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#### **Abstract**

In order to improve the global representation of rain formation in marine stratiform clouds a A new parameterization with three prognostic liquid water classes was implemented into the general circulation model ECHAM5 with the aerosol module HAM in order to improve the global representation of rain formation in marine stratiform clouds. The additionally introduced drizzle class improves the physical representation of the droplet spectrum and, more importantly, improves the microphysical processes relevant for precipitation formation compared to the standard parameterization. In order to avoid a mismatch of the liquid and ice phase, the a prognostic treatment of snow has been introduced too. This has a significant effect on the amount and altitude of ice clouds, which in turn does not only affect not only affects in- and outgoing radiation, but also the parameterized collection rates. With the introduction of a prognostic precipitation scheme a more realistic representation of both liquid and ice phase large-scale precipitation is achieved compared to a diagnostic treatment. An encouraging finding is that the sensitivity with the prognostic treatment the increase of the liquid water path to the anthropogenic aerosol forcing with the prognostic treatment in response to anthropogenic aerosols is reduced by about 25 %. Although the total net radiative forcing is increased from  $1.4 \pm 0.4$  to  $1.6 \pm 0.4$  decreased from  $-1.3 \pm 0.3$  to  $-1.6 \pm 0.3 \,\mathrm{W}\,\mathrm{m}^{-2}$  from the control to the prognostic model version, the difference is within the interannual variability. Altogether the results suggest that the treatment of precipitation in global circulation models has a significant influence on the phase and lifetime of clouds of clouds and their conversion rates, but also hints towards the uncertainties related to a prognostic precipitation scheme.

#### 1 Introduction

The challenge of projecting future climate, its sensitivity to anthropogenic forcing, the change in global precipitation patterns and the mitigation of climate change remains difficult and is closely linked to correctly representing the radiative forcing. The greatest uncertainty

lies in the representation of clouds (e.g. Stephens, 2005), whether it is their fundamental coupling to the large-scale circulation (Stevens and Bony, 2013) or their interaction with anthropogenic aerosols on the very fine scales (Boucher et al., 2013). As resolution in general circulation models (GCMs) remains coarse  $(O(100\,\mathrm{km}))$ , the improvements made on subgrid scale processes are a priori limited, may whether this be for the parameterization of convection, turbulence or microphysics, as all parameterizations are governed by the resolved larger scales. Nevertheless, understanding the dependence of the subgrid scale parameterizations and improving crude representations of subgrid scale processes is crucial for climate models to hopefully reduce the error in climate projections.

One of the largest uncertainties is the response of clouds and precipitation to changes in aerosol concentrations (Stevens and Feingold, 2009). Aerosols are necessary for the formation of clouds by acting as cloud condensation nuclei (CCN) or ice nuclei (IN) and changes in the CCN and IN number concentrations affect the number of cloud droplets or ice crystals. Thin, low level marine stratus and stratocumulus are known to be very susceptible to changes in aerosol concentration. Being very widespread, they constitute a major source of uncertainty in climate models (e.g. Bony and Dufresne, 2005). Traditionally an increase in aerosols is thought to ultimately lead to brighter clouds (Twomey, 1974) and maybe longer-lived clouds (Albrecht, 1989), reasoning that the increase in CCN leads to an increase in cloud dropletsand consequently reduces more but smaller cloud droplets, consequently reducing the efficiency to form rain. Since cloud droplets have to first collide and coalesce to form small rain or drizzle drops (autoconversion), before upon further collection of cloud droplets by drizzle drops (accretion) a produce a significant amount of rainis produced, an increase in cloud droplets can influence the precipitation formation. Where While there is ample evidence for higher CCN concentrations leading to more cloud droplets, the subsequent effect on the cloud amount and precipitation is highly debated as the. The suppression of precipitation is found to be in competition with the evaporation by entrainment of dry air (Ackerman et al., 2004; Jiang et al., 2006), such that the suppression may not sustain (Boucher et al., 2013). (Ackerman et al., 2004; Jiang et al., 2006; Boucher et al., 2013) .

Many GCMs do not have the ability to include such compensating effects on small scales, thus. Thus, it is not surprising that for mixed or ice phase clouds, the influence of aerosols on the radiative budget is still unclear (Boucher et al., 2013). For mixed phase clouds an increase in aerosols can lead to two competing effects namely the glaciation effect and the de-activation effect (Lohmann and Hoose, 2009). With more IN the cloud would glaciate and form precipitation more readily reducing its radiative effect. On the other hand, IN may also be coated with soluble material, i.e. sulfate, reducing its nucleation ability and having the contrary effect.

Whether it is for liquid, mixed or ice phase clouds, describing the influence of aerosols on clouds and precipitation formation requires a whole chain of processes that need to be considered. Here the focus will be on the collection processes and the sensitivity of the model to the anthropogenic aerosol forcing.

The increase in aerosol concentrations since the preindustrial era, has led to a negative radiative forcing corresponding to a cooling that partly offset the greenhouse gas warming (Boucher et al., 2013). The radiative forcing is comprised of (i) a direct influence of the aerosols on the in- and outgoing radiation and (ii) an indirect influence of aerosols by changing cloud properties. GCMs have shown to overestimate the cooling related to this these aerosol-cloud-precipitation interactions compared to observations (Quaas et al., 2009). Posselt and Lohmann (2008a) showed that the diagnostic treatment of rain in GCMs results in an overemphasis of the autoconversion process, which leads to a high sensitivity of clouds and thus formation of precipitation to changes in aerosol concentrations. By introducing a prognostic treatment of rain, a more realistic partitioning between the autoconversion and accretion process can be achieved (Posselt and Lohmann, 2008a; ?) (Posselt and Lohmann, 2008a; Gettelman and Mores Further studies following suit (Reitter et al., 2011; Walters et al., 2014; ?) (Reitter et al., 2011; Walters et al., 2014; ?) using a multi-scale modelling framework (Wang et al., 2011) confirm these findings, not only reducing the sensitivity of clouds and precipitation to the anthropogenic aerosol forcing, but also improving the representation of cloud ice and snow (Reitter et al., 2011) and reducing the frequency of unrealistic light rain (Walters et al., 2014).

As a follow-on to the work by Posselt and Lohmann (2008a) and Posselt and Lohmann (2009), where the importance of the drop size distribution (DSD) in low level clouds was recognized, this study focuses on the representation of both liquid and solid phase precipitation in the ECHAM5-HAM GCM and its response to the anthropogenic aerosol forcing.

For the liquid phase, an intermediate drizzle drop class is introduced a more physical representation of (i) the drizzling conditions often found in marine stratocumulus (van Zanten et al., 2005) and following the triclass parameterization by Sant et al. (2013) for the collection processes, i.e. autoconversion, accretion and self-collection, of the three liquid water classes. Although the idea of a three class system to capture the broadening of the effect it can have DSD is older (Lüpkes et al., 1989), the concept was picked up to enable applications that benefit from an explicit drizzle class. In particular, to improve the influence of different CCN and giant CCN on the formation of drizzle drops (e.g. Saleeby and Cotton, 2004; Posselt and Lohmann, 2008b) or their effect on the cloud top effective radius (Wood, 2000), and (ii) the shape of the DSD. By introducing the scheme by Sant et al. (2013) into the ECHAM5-HAM, which. However, first and foremost the scheme allows for a more physical representation of the DSD, since it is based on truncated momentsand is believed to be a more realistic representation of the DSD, the collection processes, i. e.autoconversion, accretion and self-collection, are adapted for the three liquid water classes of cloud, drizzle and rain., and consequently the collection rates. This enables to better capture the drizzling conditions often found in marine stratocumulus (van Zanten et al., 2005). Furthermore, the triclass scheme helps represent the transition from low to larger precipitation rates and improves the sensitivity of surface precipitation to changes in CCN compared to the standard scheme used in the ECHAM5-HAM GCM (Sant et al., 2013). Therefore, here the focus will be on the triclass scheme and the prognostic precipitation as a whole, leaving the application of giant CCN and the estimation of the cloud top effective radius for the future.

A weakness of the previous studies was the biased improvement of only the liquid phase, i.e. the treatment of snow remained diagnostic. Consequently, similar to the work done

by ? and ? Gettelman and Morrison (2015) and Gettelman et al. (2015) in the CAM5 model a full prognostic treatment of precipitation is introduced for drizzle, rain and snow.

The outline of the article is as follows. Section 2 will give a general description of the ECHAM5-HAM GCM and the changes made by introducing the prognostic treatment of precipitation. Results are presented in Sect. 3 starting with global fields before looking at the cloud properties and the precipitation microphysics. The section ends with the comparison of present day and preindustrial simulations. Finally, the conclusions are presented in Sect. 4.

## 2 Model description

## 2.1 ECHAM5-HAM GCM

For this study, the ECHAM5-HAM general circulation model (GCM) (Roeckner et al., 2003) coupled to the two-moment modal aerosol scheme HAM (Stier et al., 2005) as in Lohmann and Hoose (2009) is used. The model solves the prognostic equations for temperature, surface pressure, divergence and vorticity on a spectral grid with a triangular truncation. Within HAM the aerosol size distribution is represented by a superposition of 7 lognormal distributions describing the different sizes and solubilities for the compounds sulfate, black carbon, organic carbon, sea salt and mineral dust.

This ECHAM5 control version (CTRL) includes a two-moment stratiform cloud microphysics scheme for cloud droplets and ice crystals (Lohmann et al., 2007) coupled to the aerosol scheme HAM, and uses an empirical cloud cover scheme (Sundqvist et al., 1989). The microphysics scheme includes all phase changes between the water components (vapour, liquid and ice), the collection processes (autoconversion, accretion and aggregation), evaporation of rain, melting of snow and sedimentation of cloud ice (Lohmann et al., 2007). The activation of cloud drops follows the empirical activation scheme by Lin and Leaitch (1997), which depends on the aerosol number concentration and the sub grid-scale updraft velocity. Cirrus clouds form by homogeneous freezing of supercooled solution

drops (Lohmann et al., 2008), which has been found to be the dominant freezing mechanism for these type of clouds (Kärcher and Ström, 2003). Although recently, this view has been questioned (Cziczo et al., 2013).

CTRL version of model The the treats precipitation diagnostically (Ghan and Easter, 1992), i.e. it reaches the surface within one model time step or evaporates/sublimates in the sub-saturated air below the cloud. Following Khairoutdinov and Kogan (2000) for the collection processes, cloud and rain water are separated at a radius of 25 µm. However, the diagnostic treatment of rain is only realistic for drops larger than 100 µm, which lead Posselt and Lohmann (2008a) to the introduction of a prognostic rain scheme to account for the drizzling conditions often found in marine stratocumulus. Yet, one is still left with the rain class spanning all drops larger than 25 µm in radius. which tends to overestimate the number concentration of large drops and thus influences both the collection processes and the sedimentation. Consequently, we introduce drizzle  $(25 < r < 100 \, \mu m)$  as a third liquid water class and introduce-implement the triclass parameterization by Sant et al. (2013) to describe the collection processes based on truncated moments. As mentioned before, In this model version, which is referred to as PROG, prognostic equations for both liquid and ice phase precipitation are introduced for both mass and number for consistency, extending the CTRL version and referred to as PROG.. Note that, for both model versions only the cloud liquid and ice water contribute to the cloud optical properties, i.e. no changes have been made to the radiation scheme. Adding the prognostic precipitation hydrometeors (drizzle, rain and snow) into the radiation scheme will be part of future work.

Note that with With the introduction of the new collection scheme, due to for the three liquid water classes, and the prognostic precipitation scheme, the climate model is not necessarily in radiative equilibrium anymore as the parameter space for the tuning of is changed (Lohmann and Ferrachat, 2010). To preserve comparability of the two model versions the tuning parameters were not changed. Yet, if the climate model (?), i.e.to achieve radiative equilibrium, is changed. If the climate impacts of the schemes for both liquid and ice phase microphysics between the two model version CTRL and PROG are

comparable, in terms of heating rates (Mauritsen et al., 2012), a -similar climate can be expected. No other tuning parameters outside the cloud microphysics routine were changed as not to lose the comparability to the changes induced by the introduction of prognostic precipitation.

## 2.2 Prognostic precipitation

For the prognostic precipitation scheme the microphysical processes had to be restructured as illustrated by the flowchart in Fig. 1. In CTRL the microphysical processes are treated sequentially; first, the incoming precipitation fluxes (rain and snow) from the level above are modified by melting, sublimation and evaporation. This is followed by microphysical processes on cloud water and ice, i.e. condensation or evaporation, deposition or sublimation and freezing of cloud water. These processes do not interact with the precipitating water and for simplicity, will be referred to as the "phase changes". Then the conversion of cloud water (ice) to rain (snow) via autoconversion (aggregation) and the interaction of the two via accretion is calculated. The fluxes and tendencies of the given grid box are finally updated and the whole procedure is repeated for the next level down until reaching the surface.

For PROG the vertical loop is split into three loops; the first one computes the "phase changes", the second the formation of precipitation and its interaction with the cloud water and ice, and the third the sedimentation. The latter two are then subject to a sub-stepping over smaller time steps, i.e. time-splitting. Primarily, this is done to assure numerical stability of the sedimentation, as the Courant-Friedrichs-Levy (CFL) criterion needs to be met, and to improve the representation of the microphysical processes. Note that the sedimentation itself can be made subject to an additional time-splitting if necessary. Since the whole model could not be made subject to smaller time steps, the time-splitting is only applied to processes relevant to the formation and sedimentation of precipitation. Furthermore, the splitting of the vertical loop is necessary as some cloud processes use a moisture adjustment scheme, which is not time dependent and would lead to inconsistencies if included in the time-splitting. The number of iterations needed for the time-splitting, i.e. the length of the sub-time step, depends on the fall speeds and the new triclass parameterization. Tests in

single column mode showed that a sub-time step of 30 s leads to a good convergence in surface precipitation and conversion rates, keeping the computational costs at a reasonable amount (Sant, 2012).

The rates for drizzle (d), rain (r) and snow (s) mass mixing ratio,  $\partial q_i/\partial t$ ,  $i \in \{d, r, s\}$ , and number concentration,  $\partial N_i/\partial t$ , can be summarized as follows

$$\frac{\partial q_{\rm d}}{\partial t} = b_{\rm c}(Q_{\rm aut_c} - Q_{\rm aut_d} + Q_{\rm ac_{cd+}} - Q_{\rm ac_{cd-}} - Q_{\rm ac_{dr}}) - (1 - b_{\rm c})Q_{\rm evp_d} + b_{\rm r}Q_{\rm sed_d}$$

$$\tag{1}$$

$$\frac{\partial N_{\rm d}}{\partial t} = b_{\rm c} (P_{\rm aut_c} - P_{\rm aut_d} - P_{\rm ac_{cd}} - P_{\rm sc_d}) - (1 - b_{\rm c}) P_{\rm evp_d} + b_{\rm r} P_{\rm sed_d} \tag{2}$$

$$\frac{\partial q_{\rm r}}{\partial t} = b_{\rm c}(Q_{\rm aut_d} + Q_{\rm ac_{cr}} + Q_{\rm ac_{cd}} + Q_{\rm ac_{dr}}) - (1 - b_{\rm c})Q_{\rm evp_r} + b_{\rm r}(Q_{\rm mlt} + Q_{\rm sed_r})$$
(3)

$$\frac{\partial N_{\rm r}}{\partial t} = b_{\rm c} (P_{\rm aut_d} + P_{\rm ac_{cd.}} - P_{\rm sc_r}) - (1 - b_{\rm c}) P_{\rm evp_r} + b_{\rm r} (P_{\rm mlt} + P_{\rm sed_r}) \tag{4}$$

$$\frac{\partial q_{\rm s}}{\partial t} = b_{\rm c}(Q_{\rm agg} + Q_{\rm ac_s}) - (1 - b_{\rm c})Q_{\rm sub} + b_{\rm r}(-Q_{\rm mlt} - Q_{\rm sec} + Q_{\rm sed_s}) \tag{5}$$

$$\frac{\partial N_{\rm s}}{\partial t} = b_{\rm c} P_{\rm agg} + b_{\rm r} (-P_{\rm mlt} - P_{\rm sec} + P_{\rm sed_{\rm s}}), \tag{6}$$

where  $Q_{(.)}$  and  $P_{(.)}$  account for changes in mass and number, respectively, and the cloud and precipitation fraction are denoted by  $b_{\rm c}$  and  $b_{\rm r}$ , respectively. The source and sink terms related to the triclass parameterization for the collection processes are given in Sant et al. (2013) with some adjustments discussed later. For clarity all acronyms are summarized in Table 1.

For the ice microphysical processes the same parameterizations used in CTRL are also used in PROG. However, where the Due to the diagnostic treatment of precipitation in CTRL the number concentrations of the precipitating hydrometeors did not need to be determined in CTRL, are not needed. In PROG they are indispensable in PROG, whereby certain assumptions are necessaryand underlie certain assumptions. During the evaporation of drizzle and rain drops a constant mean mass is assumed, hence the change in mass is followed by an associated equal fractional change in number density. As pointed

out by Seifert (2008), this assumption holds rather well for drizzling stratocumulus (e.g. Khairoutdinov and Kogan, 2000). Following the formulation by Levkov et al. (1992) for the production of snow as described in Lohmann and Roeckner (1996), the created snow mass by aggregation was assumed to have a minimum diameter of 200  $\mu$ m. This assumption together with the mass-length relation by Spichtinger and Gierens (2009) was used to determine the snow flake number concentration. When snow crystals melt they are assumed to be a source term for the rain drops, where one snow crystal is assumed to produce one rain drop. Note that the gain of snow mass due to accretion is assumed not to change the number of snow crystals.

It should be remarked, that while the separation of the liquid water into distinct classes (cloud, drizzle and rain water) is mainly based on the accretion process leading to a minimum in the mass density distribution between cloud and drizzle drops, the separation of the ice phase into cloud ice and snow is not as clear (Morrison and Grabowski, 2008). More specifically, this is due to the importance of diffusional growth of ice crystals leading to a gradual growth and making the distinction less clear. Morrison and Grabowski (2008) presented a framework in which the number of ice water classes is reduced to a single one, retaining the history of rimed mass and allowing for a more physical representation of the growth of the ice particles. Recent progress has been made improving the scheme (Morrison and Milbrandt, 2015) and implementing it into a higher resolved regional model (Morrison et al., 2015), but it remains to be tested in a GCM.

#### 2.2.1 Sedimentation

Following Posselt and Lohmann (2008a, 2009) and Müller (2007), the sedimentation of all three precipitating hydrometeor types,  $i \in \{d, r, s\}$ ), is treated as a vertical 1-D-advection with the mass and number weighted fall velocities  $v_{q_i}$  and  $v_{N_i}$  for the mass and number density,  $q_i$  and  $N_i$ , respectively, allowing for gravitational sorting. Numerically an explicit Euler scheme was used, which by definition conserves mass. As described in Posselt and Lohmann (2008a) the sedimentation velocities are derived by using the flux density ap-

proach by Srivastava (1978, Eqs. 48 and 49) for both mass and number, respectively,

$$\mathfrak{F}_{q_i} = q_i v_{q_i} = \frac{\pi}{6} \rho_{\mathsf{w}} \int_0^\infty D^3 f_i(D) v_{\mathsf{t},i}(D) \, \mathrm{d}D \tag{7}$$

$$\mathfrak{F}_{N_i} = N_i v_{N_i} = \int_0^\infty f_i(D) v_{t,i}(D) \, \mathrm{d}D, \tag{8}$$

where  $\rho_{\rm w}$  is the density of water and  $v_{{\rm t},i}$  is the terminal velocity of the corresponding hydrometeor with diameter D. Note that with this approach, changes in mass density with height, due to changes in air density, are taken into account. Following Posselt and Lohmann (2009) the number density function  $f_i$  is assumed to follow a gamma distribution of the form

$$f_i(D) = \frac{N_i}{\Gamma(\mu_i)D} \left(\frac{D}{D_{0,i}}\right)^{\mu_i} \exp\left(-\frac{D}{D_{0,i}}\right),\tag{9}$$

where  $\Gamma(.)$  is the Gamma function,  $\mu_i$  is a free shape parameter and  $D_{0,i}$  is related to the mean diameter  $\overline{D}_i$  by

$$D_{0,i} = \frac{\overline{D}_i}{\sqrt[3]{\mu_i(\mu_i + 1)(\mu_i + 2)}}$$
 (10)

with

$$\overline{D}_i = \sqrt[3]{\frac{6}{\pi \rho_{\mathsf{w}}} \frac{\rho_a \, q_i}{N_i}}.\tag{11}$$

By determining the terminal velocity  $v_{t,i}(D)$  of a single drop or crystal, Eqs. (7) and (8) can be solved for the bulk fall velocities  $v_{q_i}$  and  $v_{N_i}$ .

For the liquid phase, i.e. drizzle and rain drops, the terminal velocity of a single drop was determined using an approximated expression based on Rogers et al. (1993). Integrating equations (7) and (8) yields the bulk fall velocities (cf. Posselt and Lohmann, 2009)

$$v_{q_i} = \begin{cases} (\mu_i + 3)b_v D_{0,i}, & \text{for } D_{0,i} \le D_v/(\mu_i + 3) \\ b_1, & \text{for } D_{0,i} > D_v/(\mu_i + 3) \end{cases}$$
(12)

$$v_{N_i} = \begin{cases} \mu_i b_v D_{0,i}, & \text{for } D_{0,i} \le D_v / \mu_i \\ b_1, & \text{for } D_{0,i} > D_v / \mu_i \end{cases}, \tag{13}$$

where the constant  $b_v=b_3(b_2-5(b_2-b_1))=3918\,\mathrm{s}^{-1}$  and the critical distribution parameter  $D_v$  is given by  $D_v=b_1/b_v=2463\,\mathrm{\mu m}$ . Equations (12) and (13) are used for both drizzle and rain drops, where the shape parameters  $\mu_\mathrm{d}$  and  $\mu_\mathrm{r}$  were chosen assuming a constant value of  $\mu_\mathrm{d}=\mu_\mathrm{r}=5$ . This choice follows the findings by Posselt and Lohmann (2009) and results in a distribution width, which is not too wide as to not overestimate the number of large drops, but wide enough to enable gravitational sorting. A direct link to the assumed distributions for the collection processes cannot be made, as sedimentation is based on grid mean values whereas the collection rates are based on in-cloud values. This incoherence should be avoided in future climate models. Note that for large  $D_{0,i}$  or  $\mu_i$  gravitational sorting does not occur anymore as  $v_{q_i}$  and  $v_{N_i}$  converge. In this case  $\overline{D}_i$  would have to be approximately 2 mm or larger, a regime not found in our simulations.

For snow flakes the terminal fall velocity of a single snow crystal is determined following Barthazy and Schefold (2006)

$$v_{\mathsf{t},i}(D) = \beta_1 D^{\beta_2} \tag{14}$$

assuming irregular crystals with  $\beta_1=1.23$  and  $\beta_2=0.22$ , where the diameter was again determined using the mass-length relation by Spichtinger and Gierens (2009). Repeating the procedure as for the liquid phase assuming the same number distribution  $f_i(D)$  yields the following bulk fall velocities for snow

$$v_{q_{s}} = \beta_{1} D_{0,i}^{3+\beta_{2}} \frac{\Gamma(3+\mu_{s}+\beta_{2})}{\Gamma(3+\mu_{s})}$$
(15)

$$v_{N_{\rm s}} = \beta_1 D_{0,i}^{\beta_2} \frac{\Gamma(\mu_{\rm s} + \beta_2)}{\Gamma(\mu_{\rm s})}.$$
 (16)

For consistency with Levkov et al. (1992), it is assumed that the size distribution of snow follows an exponential distribution, i.e.  $\mu_s = 1$ .

#### 2.2.2 Collection processes

In CTRL the collection processes responsible for the formation of precipitation follow the formulation by Khairoutdinov and Kogan (2000). With the introduction of drizzle as a third liquid water class in PROG, the collection processes of autoconversion, accretion and self-collection are reformulated using the scheme by Sant et al. (2013). By solving the stochastic collection equation for truncated moments, two-moment rate equations are derived for cloud, drizzle and rain mass and number densities. However, as the scheme is computationally rather expensive, especially for the numerical solution of contained integrals, an approximation is presented here. The integrals in question (cf. Sant et al., 2013, Appendix B) are normalized by the number concentration and then approximated by a polynomial in log-space with a non-linear least square fit using a Levenberg–Marquardt algorithm (Press et al., 1992). The details of the approximation are presented in the Appendix.

This approximation introduces a small error, mainly due to the integral used for the accretion between cloud and drizzle drops, as the dependence on the shape parameters and the continuity condition are larger than for the other approximated integrals (cf. Sant et al., 2013). Tests have shown that the error is within a few percent of the produced precipitation. This is very reasonable, especially in view of reduced computational cost. The approximation reduces the model calculation computational cost of the model calculation (whole atmosphere) by about a factor of 2.5, leaving the calculation of PROG to be a factor of 2.5 – instead of 6 – higher than CTRL. These numbers are based on a time step of 30 s for the time-splitting.

#### 3 Results

## 3.1 Model setup

All global simulations presented here are conducted at a horizontal resolution of T63 ( $\sim 1.875^{\circ} \times 1.875^{\circ}$ ) with 31 vertical levels, where the uppermost level interface is at 10 hPa. For PROG the model time step of 12 min is divided into 24 sub-time steps for all precipitation processes as described above, such that a time step of 30 s is reached for the time-splitting. Using climatological sea surface temperatures and sea-ice extent the simulations are integrated for 512 years (2000–20042000–2011) after a spin-up of 3 months. For a closer look at in-cloud processes further statistics were made based on 6 hourly data of the first month of the year 2000. This proved to be sufficient to get a good sampling of the parameter space.

To estimate the total anthropogenic aerosol effect, which is the change of the net radiation at the top of the atmosphere (TOA) due to anthropogenic aerosols, simulations with present-day (PD) and pre-industrial (PI) aerosol emissions integrated over the same 512 years are compared. The PI simulations use natural aerosol emissions representative of the year 1750 (Dentener et al., 2006).

## 3.2 Present day results

## 3.2.1 Global and zonal fields

The global and zonal annual mean cloud properties and TOA energy budget are presented in Table 2 and Fig. 2, respectively, for the simulations with PROG and CTRL, together with observations for comparison. The total precipitation amount is very similar in both PROG and CTRL, both overestimating the observed amount from the Global Precipitation Climatology Project (GPCP) dataset of the years 1979–2001 (Adler et al., 2003). The overestimation is mainly located in the tropics and subtropics as can be seen in Fig. 2a. Separating the total precipitation into stratiform, i.e. large-scale, and convective precipitation a shift to more stratiform precipitation in PROG can be deduced, at the same time reducing the amount of convective precipitation. Although the differences in the zonal mean are not very

pronounced (cf. Fig. 2b and d), the largest difference between PROG and CTRL is seen in stratiform precipitation in the tropics. In addition, the fraction of stratiform precipitation falling as snow is shown in Fig. 2c, where a slight shift towards the tropics is seen in PROG. In general, the prognostic treatment of both warm and cold precipitation (drizzle, rain and snow) has little effect on the precipitation amount reaching the surface, which is to be expected as we are using climatological sea surface temperatures (SSTs). Since the SSTs determine the evaporation rates over the ocean and precipitation is balanced by the amount of evaporation, the differences in precipitation amount need to be small. Nevertheless, the hydrological cycle can could be altered by increasing or decreasing the cloud lifetimes through the microphysical collection processes, which will be addressed later.

The simulated total cloud cover (TCC) is within the range of observations in the zonal mean for both PROG and CTRL, differing by about one percent in the global mean between the two simulations. The differences are located in the northern mid-latitudes as can be seen in Fig. 2e. Although small, they remain persistent throughout the whole simulation period.

The prognostic treatment of precipitation has a large effect on both liquid and ice water paths globally and on the zonal mean (cf. Fig. 2f–h). In a global mean the cloud liquid water path (LWPCLWP) is 7% smaller in PROG (69.2–69.1 compared to 74.2 g m² in CTRL), where both are within the range of observations. Over the oceans, PROG exhibits a smaller LWP-CLWP in the Southern Hemisphere, but a higher peak in the tropics compared to CTRL (cf. Fig. 2f and g). Over the Northern Hemisphere ocean the differences between the two simulations are negligible, however taking the land into account consideration (cf. Fig. 2h) reveals that the reduction in LWP-CLWP is located in the mid-latitudes of both hemispheres, where stratiform clouds are dominant. On the other hand, the cloud ice water path (HWP) is about CIWP) is nearly 60% larger in PROG and although still strongly underestimating the observations, it can be seen as an improvement. Summed up (LWPCLWP + HWPCIWP, not shown) the LWP-CLWP dominates the zonal mean, where PROG shows a reduced water path in the mid-latitudes but a larger peak in the tropics. In general, the peaks in both LWP and IWP-CLWP and CIWP in the tropics are in better agreement with observations. On

a further note, the frequency of occurrence of both liquid and ice clouds have decreased from CTRL to PROG. For the liquid phase clouds the lifetime decreased by about 30(10of the time there is a cloud in the lowest 8 model levels in PROG compared to about 14of the time in CTRL), whereas for ice clouds the decrease was about 5, respectively (same values are valid for the occurrences of the corresponding PI simulations). This shows, in particular for the liquid clouds, that the hydrological cycle is faster in PROG than in CTRL.

The prognostic scheme also has a significant impact on the vertically integrated cloud droplet number concentration (CDNC burden) and the ice crystal number concentration (ICNC burden), reducing them by. Even though the global mean CDNC burden has hardly changed (3 % and nearly 25, respectively, reduction) from CTRL to PROG. The effects are most prominent, the effect is more prominent on a zonal mean, where differences are seen in the mid-latitudes and especially large for (cf. Fig. 2i). For the ICNC burden the prognostic scheme has lead to a 25 % reduction, again located in the mid-latitudes but especially large in the northern hemisphere (cf. Fig. 2i and j). The reason for these effects and the relation to the water paths will be discussed in more detail below.

The link between the cloud properties and the TOA radiative effects is a little more intricate. The reduced LWP\_CLWP and cloud lifetime in PROG compared to CTRL has increased the shortwave cloud radiative effect (SWCRE, notation used as in Boucher et al., 2013) by nearly 3 W m<sup>-2</sup> in the global mean, which means that less shortwave radiation is reflected back to space. In the zonal-mean shown in Fig. 2k, this increase can be found again in mid-latitudes corresponding to the location of the largest changes in LWPCLWP. The longwave cloud radiative effect (LWCRE) tends to be more sensitive to the TCC. As the reduction in TCC is not very large it cannot account for the reduction of 3 W m<sup>-2</sup> in PROG compared to CTRL LWCRE from CTRL to PROG. As the increase in IWP CIWP is accompanied by a decrease in ICNC burden in PROG, the ice crystals are larger, causing a reduction in the optical depth of ice clouds. The reduction of TCC and in ice clouds optical depth are the reason for the reduction in LWCRE (Fig. 2l).

As the changes in SWCRE and LWCRE <u>nearly</u> compensate each other, the net total radiative effect at TOA  $(F_{net})$  hardly changed from CTRL to PROGonly differs by

0.2–0.3 W m<sup>-2</sup> between CTRL and PROG, which corresponds to the interannual variability (not shown). Note that CTRL is tuned to be in equilibrium and it is rather remarkable that without any further tuning the new prognostic scheme results in a similar climate. Furthermore, the TOA shortwave (SW) and longwave (LW) radiation of PROG are in better agreement with the newest estimates presented by Wild et al. (2013).

The aerosol optical depth (AOD) has been reduced in PROG by around 1715% with respect to CTRL and is certainly related to the prognostic treatment of precipitation. FirstPrimarily, the precipitation takes longer to reach the surface and can therefore be a more effective sink through wet-deposition. Second, the increased hydrological cycle mentioned above ensures that the Further changes in precipitation amount, through altered collection rates, or the hydrological cycle can influence the frequency with which aerosol particles are washed outmore frequently than in CTRL. This also shows that isolating the direct radiative effect of the aerosols (ERFari; Boucher et al., 2013) from the indirect effects on clouds (ERFaci) is very difficult, since both effects act at the same time (Lohmann et al., 2010). Consequently, we prefer to only estimate the total radiative effect (ERFari + aci).

# 3.2.2 Cloud properties

The introduction of a prognostic precipitation scheme has a particularly strong influence on the IWP-CIWP and the ICNC burden. The zonal mean Fig. 3 shows the zonal mean cloud ice water content (IWC) shown in Fig. 3 shows that the IWC-CIWC) which has not only increased in PROG compared to CTRL, but illustrates that ice clouds are more frequent in lower altitudes, in particular close to the annual mean melting line. Compared to annual mean ERA-interim reanalysis (Dee et al., 2011) for the years 2000–20042000–2011 and ensemble mean satellite observations of CloudSat and CALIPSO (Li et al., 2012), the known underestimation of IWC-CIWC by the ECHAM5-HAM model is clearly visible. ThusHowever, the higher IWC-CIWC values in PROG and the shift to lower altitudes is in better agreement with the both ERA-interim reanalysis and satellite observations. Consequently, the reduced LWCRE using prognostic precipitation is additionally also due to the

lower lying ice clouds. As the ICNC burden is reduced with higher IWCCIWC, the average ice crystal size is larger and therefore ice crystals sediment faster, resulting in clouds at lower altitudes. The reason for the reduced ICNC burden is caused by a change in collection rates and will be addressed later. Note that the ERA-interim reanalysis is based on model physics, but by using data assimilation, the comparison to the satellite observations is better. On the other hand, as discussed in more detail in Li et al. (2012), caution is also warranted for the satellite observations of mixed phase clouds in mid- and high-latitudes, where the contributions of ice and liquid water are difficult to distinguish. Consequently, the reanalysis cannot be seen as the truth, but it illustrates a suitable and plausible description of the current climate.

For the cloud liquid water content (LWCCLWC) shown in Fig. 4 the prognostic treatment has lead to an overall reduction. Comparing zonal mean LWC CLWC in PROG and CTRL (cf. Fig. 4a and b), CTRL has a higher LWCCLWC, especially in midlatitudes and above the mean melting line. Note that the annual mean melting lines are only given as a guidance for the comparison to the IWCCIWC. With respect to the ERA-interim reanalysis clear differences are seen; in the midlatitudes both PROG and CTRL overestimate the LWCCLWC, whereas towards the tropics the low altitude clouds are missing and LWC-CLWC at higher altitudes is overestimated. Although PROG shows a tendency to have a few more liquid water clouds in at lower altitudes in the subtropics, there are too many liquid water clouds close to the surface, a bias already exhibited by CTRL that remains to be addressed.

The LWC-CLWC difference between the two model versions (cf. Fig. 4c) shows that below the melting line LWC-CLWC is enhanced in PROG, but reduced above it. The hypothesis is that due to the higher IWC-CIWC in general and particularly above the melting layer, supercooled liquid water is more rapidly depleted trough the Wegener-Bergeron-Findeisen (WBF) process in PROG compared to CTRL. The difference in relative humidity depicted in Fig. 5a helps to underline this hypothesis as large regions above the melting layer are dryer by up to 6%. On one hand, this correlates with the higher IWC-CIWC values and on the other, it makes the formation of new liquid water clouds less probable, as a relative humidity based cloud cover scheme is used (Sundqvist et al., 1989). Vice versa, the enhanced LWC

CLWC in the tropics can be explained by the enhanced relative humidity in the regions below the melting layer of up to 43%, the reasons for which have not been clarified yet. Figure 5b also shows the difference in cloud cover between PROG and CTRL, which although small (below  $\pm 3\%$ ), is dominated by the reduction of ice clouds in the midlatitudes and enhanced liquid clouds in the tropics.

## 3.2.3 Precipitation microphysics

With the introduction of prognostic precipitation (PROG) the global mean precipitation rate did not change much compared to the reference simulation (CTRL). However, as precipitation is balanced by the amount of evaporation and we are using climatological sea surface temperatures SSTs, which determine the evaporation rates over the ocean, this is not surprising. Nevertheless, the microphysical processes related to precipitation may well be altered; (i) by the introduction of a third liquid water class, i.e. drizzle, and the altered collection rates, and (ii) by the time splitting of the precipitation related microphysics.

Opposed In contrast to the diagnostic treatment in CTRL, PROG prognostically sediments and keeps precipitation in the atmosphere from one time step to the next. Figures 6a–c show the drizzle, rain and snow water contents (DWC, RWC and SWC), respectively, with the contour lines of the LWC and the IWC CLWC and the CIWC overlaid. The highest DWCs are found in the midlatitudes, which correlate well with the largest values in LWCCLWC, as expected. The values are larger at lower altitudes due to the effect of gravitational sorting. For the RWC the values are one order of magnitude larger, not only because of the larger drop sizes, but in particular because of the additional source from the melting of snow. Especially in the tropics this is the case, where the RWC correlates better with the IWCCIWC and SWC than with the LWCCLWC. As expected, the SWC is well correlated with the IWCCIWC, since it is the only primary source for snow. The largest values are just above the melting line, again an effect of the gravitational sorting but also due to the continuous accretion of ice water, especially as and liquid water, since the accretion rate is temperature dependent being higher and larger at warmer temperatures. Below the melting line the SWC decreases rapidly, where it melts and forms rain. Compared to other stud-

ies (Reitter et al., 2011; Wang et al., 2011; ?) (Reitter et al., 2011; Wang et al., 2011; Gettelmare the amount of snow in the atmosphere is an order of magnitude smaller. The reason for which could be the treatment of the collection processes, melting or sedimentation. Future work will address this issue.

To assess the effects of changes in the microphysics of, for both the liquid and ice phase, and the introduction of prognostic precipitation, 6 hourly data of January the year 2000 for both PROG and CTRL were used for further analysis. As we are only interested in how the different processes change for a given state of the cloud, the amount of data is sufficient For mean values the whole year was used, whereas for instantaneous values one month of data revealed to be a sufficiently long time period for good statistics.

In Fig. 12a c the frequency of occurrence of instantaneous surface precipitation, liquid water path and ice water path are shown, respectively, for both CTRL and PROG. The instantaneous surface precipitation shows a clear shift to much larger values in PROG. having its peak around 65, but all together has fewer precipitation events. In contrast, CTRL overall has more events, but the peak instantaneous surface precipitation is around 2, which includes a lot of light rain events. As shown for the global MetUM model in Walters et al. (2014), the prognostic treatment of precipitation has largely reduced these light rain events and is producing more realistic surface precipitation amounts. The distribution of instantaneous LWP has not changed significantly, except that there has been a slight shift to larger values in PROG. However, for the distribution of instantaneous IWP both simulations differ quite significantly. CTRL exhibits a somewhat bimodal distribution, with a small peak around 10and a very distinct peak around 63. In PROG the IWP is clearly unimodal with its maximum around 18, but also has a much larger tail towards larger values, all leading to the increased global annual mean values discussed before For the latter case the dependence on the season was found to be negligible. We arbitrarily chose the month of July to make sure of no remnant influence from spin-up.

For the collection processes in the liquid phase the vertically integrated monthly mean mean vertically integrated cloud autoconversion (AU = aut<sub>c</sub>) and accretion rates (AC = ac<sub>cd</sub> + ac<sub>cr</sub>) are shown in Figs. 7 and 8 for CTRL and PROG, respectively. Figure 7 shows that  $\underline{\mathsf{AU}}$  and  $\underline{\mathsf{AC}}$  in the CTRL simulation  $\underline{\mathsf{AU}}$  and  $\underline{\mathsf{AC}}$  are comparable such that the fraction of  $\underline{\mathsf{AU}}$  to the total conversion rate ( $\underline{\mathsf{AU}} + \underline{\mathsf{AC}}$ ) is well above 30 % in most regions. It increases with increasing stratiform precipitation and therefore increases towards high latitudes (cf. Fig. 2). This overestimation of  $\underline{\mathsf{AU}}$  is well known in global models with diagnostic treatment of rain (e.g. Posselt and Lohmann, 2008a; Gettelman et al., 2013). By introducing the prognostic treatment of drizzle and rain in PROG the fraction of  $\underline{\mathsf{AU}}$  to the total conversion rate is strongly reduced to values well below 6 % as can be seen in Fig. 8. Although  $\underline{\mathsf{AU}}$  is of similar magnitude,  $\underline{\mathsf{AC}}$  has increased drastically. The global mean burdens are given in Table 3 and illustrate these conclusions again. Effectively,  $\underline{\mathsf{AC}}$  has increased by nearly two orders of magnitude from CTRL to PROG, the reasons for which is the longer lifetime of drizzle and rain in the cloud.

To understand the differences in AU from CTRL to PROG, instantaneous values of AU as a function of the mean effective cloud droplet radius for both CTRL and PROG are shown in Fig. 9. At large mean radii, the autoconversion rates of PROG are an order of magnitude higher than in CTRL. However, the triclass scheme in PROG exhibits a threshold type behaviour, such that below an effective radius of  $10\,\mu m$  the amount of mass transferred to drizzle is negligible. While such a threshold does not exist in nature, model studies and observations suggest that cloud droplets need to grow beyond  $12-14\,\mu m$  for a significant amount of rain to be produced (Rosenfeld et al., 2012), i.e. for growth by collision-coalescence to dominate over diffusional growth. CTRL has a more gradual decrease and has autoconversion rates which are several orders of magnitude larger than PROG at effective radii below  $10\,\mu m$ . Although, one may argue that this is needed to produce a realistic amount of rain in a diagnostic scheme, it is not physical. Thus, the new prognostic scheme allows for a more physical representation not only of the sedimentation of hydrometeors, but consequently also for the microphysical processes.

Similarly for For the ice phase , the mean vertically integrated aggregation (AG = agg) and accretion (ACi =  $ac_s$ ) rates are shown in Figs. 10 and 11 for CTRL and PROG, respectively. In contrast to the liquid phase, ACi is lower than AG in both simulations, such that AG makes up around 6874 % and even 8586 % of the total collection rate in the global mean for

CTRL and PROG, respectively (cf. Table 3). The values for AG and ACi have both increased significantly by one or nearly two orders of magnitude in PROG, which is mainly due to the large increase in IWCCIWC, but also because of the occurrence of cloud ice at lower altitudes, since both AG and ACi increase at warmer temperatures. This increase stems from the increase in the quasi liquid layer with increased temperature, in turn causing the sticking efficiency to increase. Since the ratio of AG/ACi is not known on these scales and # constitutes a combination of complex processes, which also depend on the representation of the ice phase (Morrison and Grabowski, 2008), we cannot conclude if whether the prognostic treatment of snow has lead to a more realistic treatment of the ice phase processes. Nevertheless, these results suggest that within the framework of a prognostic precipitation scheme AG plays the dominant role in the total collection rate. However, note that as the actual collection processes for the ice phase have not been changed from CTRL to PROG. an overestimation of AG or ACi cannot be excluded due to the applied time splitting. The suitability of the collection processes needs to be further tested, e.g. whether the assumption of a smallest snow crystal size in the aggregation process needs to be adjusted under a prognostic scheme (Eidhammer et al., 2014).

To quantify the effect of prognostic precipitation on observable variables Fig. 12a–c shows the frequency of occurrence of instantaneous surface precipitation (from stratiform clouds), CLWP and CIWP, respectively, for both CTRL and PROG. The distribution of instantaneous surface precipitation in PROG has become broader and the peak values around 0.8 mm day<sup>-1</sup> are strongly reduced compared to CTRL. Although the occurrence of rates below 0.1 mm day<sup>-1</sup> has increased in PROG, the contribution to the overall surface precipitation amount is negligible. However, a larger contribution is due to the slight increase of precipitation rates above 0.2 mm day<sup>-1</sup>, which together with the increased amount of precipitation events (about 15%) could explain the increased stratiform precipitation amount in PROG. Overall, the prognostic treatment of precipitation has reduced the amount of light rain events around 1 mm day<sup>-1</sup>, similar to what was achieved for the global MetUM model by Walters et al. (2014). The distribution of instantaneous CLWP (cf. Fig. 12b) has slightly shifted to larger values in PROG. However, as there are a smaller number of values in

PROG this shift is not significant. In Fig. 12c the distribution of instantaneous CIWP values differs between both simulations. CTRL exhibits a very distinct peak around  $60\,\mathrm{g\,m^{-2}}$  and a shoulder all the way down to around  $10\,\mathrm{g\,m^{-2}}$ . In PROG the CIWP has a broader, slight bimodal distribution with a maximum at  $30\,\mathrm{g\,m^{-2}}$  and a smaller peak at  $110\,\mathrm{g\,m^{-2}}$ . For CIWP below  $10\,\mathrm{g\,m^{-2}}$  the distributions are very similar. We believe that the cut-off behavior towards larger values in CTRL is related to the higher lying ice clouds as seen in Fig. 3b. With more ice clouds in the mixed phase region, i.e. between the melting and homogeneous freezing level, the ice crystals can grow larger and hence lead to larger CIWP values. Thus, the tail of larger values in PROG corresponds to the increased global annual mean CIWC values discussed before (cf. Fig. 3a). Note that these overall conclusions of Fig. 12 can also be drawn for a different month or region, i.e. tropics, mid-latitudes or ocean and land (not shown).

To further assess the effect of prognostic precipitation on the hydrological cycle, Fig. 13 presents all the relevant sources and sinks for the large scale liquid and ice clouds of CTRL and PROG for the year 2000. The global mean values of surface evaporation, precipitation and the water paths are very similar to the multi-annual global means presented in Table 2, which suggests that the rates would not change much for an analysis over a longer time period. Due to the use of climatological SSTs the amount of surface evaporation is very similar in both simulations and is balanced by the amount of convective and large scale precipitation. The reduced CLWP in PROG is mainly due to the strongly reduced condensation rate from 18.2 mg m<sup>-2</sup> s<sup>-1</sup> in CTRL to 8.6 mg m<sup>-2</sup> s<sup>-1</sup> in PROG. This is connected to smaller amounts of CLWC in higher altitudes (cf. Fig. 4a) and is underlined by the enhanced WBF process, which has increased by 30 %. Consequently, PROG produces less warm phase precipitation than CTRL on a global mean. In contrast, the increase in CIWP is largely due to the increased deposition rate in PROG, which has nearly doubled with respect to CTRL. Again, this is correlated to the larger vertical extent of the CIWC mentioned before, but is probably also due to the longer residence time of the CIWC allowing for larger gain by deposition and the WBF process. Accordingly, the amount of freezing and sedimentation is reduced. With the larger CIWP the rate of sublimation. melting and snow production has increased substantially from CTRL to PROG. Interestingly, in PROG the amount of cloud liquid water accreted by snow has reduced drastically to  $0.6 \,\mathrm{mg}\,\mathrm{m}^{-2}\,\mathrm{s}^{-1}$  compared to  $7.5 \,\mathrm{mg}\,\mathrm{m}^{-2}\,\mathrm{s}^{-1}$  in CTRL. We believe this is simply due to the reduced amount of cloud water above the melting layer, which again is coupled to the source rates for cloud ice. It is difficult to say which process is mainly responsible for the decrease in CLWP or the increase in CIWP in PROG, as all act together. Finally, the enhanced large scale precipitation rate in PROG is ultimately due to the larger primary production of snow (around 50%) and the smaller sublimation rates. For the warm phase precipitation, the amount of loss due to evaporation is just above 65 % for both CTRL and PROG. On the other hand, the amount of sublimation is reduced from 34% for CTRL to 4% for PROG. The reduced relative humidity above the melting level (cf. Fig. 5a), which is correlated to the larger vertical extent of CIWC, is certainly one reason, but others processes related to the prognostic treatment, such as sedimentation, will need to be addressed in future work. Dividing CLWP or CIWP by the sum of all sources or sinks given in Fig. 13 yields the cloud lifetimes for liquid  $\tau_{lig}$  or ice  $\tau_{ice}$  clouds, respectively.  $\tau_{lig}$  slightly increased from CTRL to PROG from 0.6 h to 0.7 h, whereas  $\tau_{ice}$  remains the same. The conclusion is that with the introduction of prognostic precipitation the speed of the hydrological cycle is not changed, but that internal conversion rates are altered and lead to a redistribution of cloud liquid and ice. Gettelman et al. (2015) report similar findings with respect to the conversion rates, but in contrast find that the CIWP is strongly reduced.

# 3.3 Present day vs. preindustrial climate

To evaluate the total anthropogenic aerosol effect, pre-industrial (PI) simulations where performed for both CTRL and PROG. Global annual mean values are also summarized in Table 2 and the differences (PD - PI) in zonal annual mean averages means of selected quantities are shown in Fig. 14.

The differences from PD to PI in global annual mean from PI to PD in precipitation amount and total cloud cover TCC are small, i.e. on the order of a few percent , in both CTRL and PROG, which and is mainly due to the use of climatological sea surface temperatures SSTs

as mentioned before. Consequently, the slight increase in precipitation and slight reduction in total cloud cover TCC cannot be attributed to a specific process.

The LWP-CLWP and CDNC burden have increased significantly in the Northern Hemisphere due to the increase in aerosol concentration from PI to PD for both CTRL and PROG. The PROG simulation reduces the increases in both LWP and CDNC in the Northern Hemisphere, but Although PROG reduces these increases compared to CTRL, it is not as pronounced as the a factor two found in a previous study (Posselt and Lohmann, 2009). The increase in CDNC and LWP leads increases in CLWP and CDNC burden lead to an increase in SWCRE as shown in Fig. 14a, c and e. In addition to the total water path, the SWCRE is also influenced by TCC (cf. Fig. 14b), which could explain the double peak in the Northern Hemisphere of the SWCRE change in PROG. However, as these changes in TCC are below 1 %, the interpretation is to be taken with caution. Furthermore, note that the change in TCC at the equator - although small - has opposite signals in PROG and CTRL, leading to a negative SWCRE in PROG, but a positive SWCRE in 1 W m<sup>-2</sup> difference in SWCRE between PROG and CTRL.

The change in LWCRE (cf. Fig. 14f) is rather noisy and not significant in a global mean, but seems to correlate with the change in IWP\_CIWP, which has not changed significantly either (not shown). Interesting is that the rather large change in ICNC burden in CTRL (approximately 14 % with respect to PD in a global mean) has little effect on the radiation balance (cf. Fig. 14d and f). On the other hand, the changes in ICNC burden in PROG are negligible.

The total net radiative effect of the anthropogenic aerosol forcing (ERFari + ari) in terms of TOA energy balance ( $\Delta F_{\rm net}$ ) is a negative effect of around  $-1.3\pm0.3\,\rm W\,m^{-2}$  for CTRL and  $-1.6\pm0.3\,\rm W\,m^{-2}$  for PROG, respectively, where the error estimation corresponds to the interannual standard deviation. Both model versions overestimate the inverse estimate of  $-1.1\pm0.4\,\rm W\,m^{-2}$  by Murphy et al. (2009) and fall towards the negative end of the ERFaci + ari estimate in Boucher et al. (2013) . The difference of about 0.3 W m $^{-2}$  between CTRL and PROG can mainly be attributed to the differences in the tropical and subtropical regions, where the change in  $F_{\rm net}$  is more negative for PROG (cf. Fig. 14g). Here, the reason

for the enhanced CLWP and TCC in the tropics of PROG has not been clarified and will need further attention, but is likely related to the coupling of the large scale clouds to the convective cloud scheme. In addition, considering that the tuning parameters in PROG are the same as in CTRL, except for the autoconversion and accretion rates altered by the new collection scheme, the difference in  $F_{\rm net}$  could be a measure of uncertainty in the current model setup, in particular for the tropics.

As the anthropogenic aerosol forcing mainly influences liquid water clouds, the change in shortwave radiation is dominant and leads to a net negative radiative forcing from PI to PD for both CTRL and PROG. Following Wang et al. (2011, cf. their Fig. 16) the sensitivity of the LWP CLWP to changes in CCN concentrations is shown in Fig. 15a and b for CTRL and PROG, respectively. The figures show boxplots and linear regression lines of the relative changes [(PD - PI)/PI] in annual mean LWP vs. CLWP versus the relative changes in annual mean activated CCN concentrations averaged over the lowest 8 model levels (surface to about 800 hPa). From the regression lines it can be deduced that with the prognostic treatment of precipitation a 3025 % decrease in slope, and hence, lower sensitivity of the LWP CLWP to the aerosol forcing can be achieved. Furthermore, the resulting slope in PROG is very similar to the one slope of 0.11 found by Wang et al. (2011) using the MMF model. The spread in change of LWP at low cLWP at low relative changes in CCN is slightly larger in PROG than for CTRL, but at larger values the spread in CTRL is significantly notably larger. In fact, for CTRL the larger the relative change in CCN the larger and broader the change in LWP relative change in CLWP becomes, whereas for PROG the spread tends to become narrower. In PROG the spread of the data tends to level off, but not necessarily the mean values (illustrated by the stars in Fig. 15a and b).

The total net radiative effect of the anthropogenic aerosol forcing (ERFaci+ari; Boucher et al., 2013) in terms of TOA energy balance ( $F_{\text{net}}$ ) is a cooling of around  $1.4\pm0.4$ for CTRL and  $1.6\pm0.4$ for PROG, respectively, where the error estimation corresponds to the interannual SD. Both model versions overestimate the inverse estimate of  $1.1\pm0.4$ by Murphy et al. (2009) and fall towards the negative end of the ERFaci+ari estimate in Boucher et al. (2013). The difference of about 0.2between

CTRL and PROG can mainly be attributed to the differences in the Northern Hemisphere and in particular around the equator, where the two simulations have opposing  $F_{\rm net}$ . The reason for the enhanced LWP and TCC in the tropics of PROG has not been clarified and will need further attention, but is likely related to the coupling of the large scale clouds to the convective cloud scheme. In addition, considering that the tuning parameters in PROG are the same as in CTRL, except for the autoconversion and accretion rates of cloud liquid water altered by the new collection scheme, this could be a measure of uncertainty in the current model setup. Further tests and sensitivity studies are necessary to deduce what is mainly due to compensating errors and where model errors dominate The Pearson correlation coefficients are low for both simulations, due to the data being localized between -0.25 and 1.0 in relative change of CCN. However, particularly for PROG where the data tends to level off, the lower correlation coefficient could be an indication that a linear regression might not be the most appropriate. Overall, the result compares well to the study performed by Wang et al. (2011) and is encouraging regarding the use of prognostic precipitation.

## 4 Conclusions

We have introduced a prognostic precipitation scheme for both the liquid and ice phase for which the precipitation microphysics, i.e. melting and evaporation of precipitation, collection rates, and sedimentation, are subject to a time splitting. For the liquid phase, drizzle was introduced in addition to the existing cloud and rain water classes. Based on the triclass parameterization by Sant et al. (2013) the collection rates for autoconversion, accretion and self-collection were changed to account for the three classes. The ice phase microphysics scheme remained unchanged. Without any further tuning of the model the simulations compare well to observations and display clear improvements in modeled physics.

Improvements in the liquid phase are equivalent to other studies (Posselt and Lohmann, 2008a; ?) (Posselt and Lohmann, 2008a; Gettelman et al., 2015), where the prognostic treatment has lead to a better partitioning of the collection processes

and consequently autoconversion and accretion rates, and subsequently to an improved sensitivity of the cloud liquid water path to the anthropogenic aerosol forcing. Furthermore, we were able to underline as inshow similar to Walters et al. (2014), that with the prognostic treatment of rain the problem frequency of light rain events can be significantly reduced. Although the reduced. Compared to the previous studies with two prognostic liquid water classes (Posselt and Lohmann, 2008a, 2009) the triclass scheme including drizzle does not significantly influence the resulting climate and under an anthropogenic aerosol forcing leads to similar reductions in CLWP and CDNC burden changes in the Northern Hemisphere as compared to a prognostic scheme with two classes, i. e.cloud and rain waterscheme with diagnostic precipitation. However, a prognostic scheme with three liquid classes still has its benefit depending on the requirements is beneficial for certain applications. Following previous research on the effect of giant CCN on warm phase clouds (Posselt and Lohmann, 2008b) (e.g. Saleeby and Cotton, 2004; Posselt and Lohmann, giant CCN can act as a direct source for drizzle drops and would be well represented by such a scheme as it circumvents the problem of not being able to differentiate between the activation of smaller and larger drops. Furthermore, a simple coupling of drizzle to the radiative transfer calculations by taking into account both cloud and drizzle water to determine the clout top effective radius as suggested by Wood (2000), may have an influence on the SWCRE of marine stratocumulus. Although first tests with giant CCN did not lead to any new findings and the simple coupling of drizzle to the radiation only showed little effect (both not shown), this might change with improvements in the representation of low level clouds. In particular for pristine clouds, taking into account drizzle-sized drops (> 20 µm Wood, 2000) was shown to be very important in determining the optical properties of the cloud. Under the current model setup first tests of these effects have been undertaken, but have been surpassed by the effects due to the introduction of prognostic snow.

The influence of the implementation of a prognostic precipitation scheme on the icephase is largest on has lead to a redistribution of cloud liquid and ice, reducing the amount of ice water and ICNC, such that larger crystals are present leading to more ice cloudsat

liquid clouds and increasing the amount of ice clouds, especially in lower altitudes. On the other hand, larger amounts of This is caused by a more effective growth of ice crystals by deposition and the Wegener-Bergeron-Findeisen process. As a result, more ice water at lower altitudes produce produces more snow through enhanced aggregation and accretion rates. This resulted in a slightly faster hydrological cycle and is most likely the cause for the reduction in LWC above the melting levelThe speed of the hydrological cycle is not changed significantly, but changes in source and sink rates have altered the distribution of liquid and ice water in the atmosphere. As the conversion processes in description of the microphysical conversion processes for the ice phase were not changed with the introduction of prognostic snow, we do not exclude that adjustments may be necessary, especially in view of the rather small snow water mass compared to other studies (Reitter et al., 2011; Wang et al., 2011; ?) . However, the results already show that with the prognostic treatment of snow feedbacks related to both ice and liquid phase arise. (Reitter et al., 2011; Wang et al., 2011; Gettelman et al., 2015) .

The effects of the new prognostic scheme on the radiative budget of the liquid and ice clouds at the top of the atmosphere compensate each other, such that the radiation fluxes are rather well balanced. Especially the global mean short- and longwave radiation is in better agreement with observations. The prognostic treatment of liquid and solid phase precipitation is certainly more physical, as models tend to go to smaller and smaller resolutions in both time and space. Moreover, the sensitivity of the liquid water path to changes in CCN concentrations is reduced, favoring a prognostic over a diagnostic treatment of precipitation when studying aerosol-cloud-precipitation interactions. Although the total net radiative effect of the anthropogenic aerosol forcing is slightly higher in PROG than in CTRL, the values are within the interannual variabilityand further investigations related to the tuning of the model setup might lead to an improvement in the uncertainty. Furthermore, investigating the coupling of the large scale to the convection scheme, since essential differences are found in the tropical and subtropical regions, might improve the uncertainty in the anthropogenic aerosol forcing.

It has been shown that with the prognostic scheme the treatment of precipitation is more physical, but the feedbacks are very versatilevariable. Especially the representation of the ice phase within the prognostic scheme needs to be addressed in future work, as dependencies on the treatment of collection processes and sedimentation cannot be excluded. Furthermore, the influence of recent improvements in the formation of cirrus clouds (Kuebbeler et al., 2014) or the sensitivity to ice properties (Eidhammer et al., 2014) and the number of ice water classes (Morrison and Grabowski, 2008) to the prognostic treatment of precipitation will need to be studied. To conclude, with the introduction of prognostic precipitation future steps should will imply a coupling of precipitation to the radiative transfer calculations in order to ensure consistency there. In this regard, including the precipitating hydrometeors into the radiation calculation, especially for snow, should improve model fidelity and reduce compensating errors as Li et al. (2014) show.

## **Appendix: Approximated collection processes**

Following the formulation by Sant et al. (2013) the collection processes are derived from the stochastic collection equation for truncated moments

$$\frac{\partial M_i^n}{\partial t} = \frac{1}{2} \left[ \int_l^u \int_0^{u-y} (x+y)^n f(x) f(y) K(x,y) \, \mathrm{d}x \, \mathrm{d}y \right] 
+ \int_0^l \int_{l-y}^{u-y} (x+y)^n f(x) f(y) K(x,y) \, \mathrm{d}x \, \mathrm{d}y \right] 
- \int_0^\infty \int_l^u x^n f(x) f(y) K(x,y) \, \mathrm{d}x \, \mathrm{d}y, \tag{A1}$$

where x and y are the drop masses collecting each other, l and u are the lower and upper integration limits of class i ( $i \in \{c, d, r\}$ ), respectively, K(x, y) is the collection kernel (Sant et al., 2013, Eq. 9) and f is assumed to be a four parameter modified gamma distribution

$$f_i(x) = A_i x^{\tilde{\nu}_i} \exp(-\lambda_i x^{\tilde{\mu}_i}). \tag{A2}$$

To avoid solving certain integrals numerically and save computational costs the following generalized integrals

$$I_{\text{cc}}^{n} = \frac{k_{\text{cc}}}{2} \int_{0}^{x_{1}} \int_{x_{1}-y}^{x_{1}} f_{\text{c}}(x) f_{\text{c}}(y) (x^{2} + y^{2}) (x + y)^{n} dx dy$$
(A3)

$$I_{\mathsf{dd}}^{n} = \frac{k_{\mathsf{dd}}}{2} \int_{x_{1}}^{x_{2}-x_{1}} \int_{x_{2}-y_{1}}^{x_{2}} f_{\mathsf{d}}(x) f_{\mathsf{d}}(y) (x+y)^{n+1} \, \mathsf{d}x \mathsf{d}y \tag{A4}$$

$$I_{\mathsf{cd}}^{n} = k_{\mathsf{cd}} \int_{x_{2}-x_{1}}^{x_{2}} \int_{0}^{x_{2}-y} f_{\mathsf{c}}(x) f_{\mathsf{d}}(y) (x+y)^{n+1} \, \mathsf{d}x \mathsf{d}y \tag{A5}$$

are approximated for the moments of number and mass density, i.e.  $n \in [0,1]$ , using a 4 parameter polynomial of the form

$$P_4(r) = a_1(a_0 + r)^2 + a_2(a_0 + r)^3 + a_3(a_0 + r)^4,$$
(A6)

where r is the radius in  $\mu$ m and  $a_m$ ,  $m \in (0,3]$ , are the fitting parameters. To this end, from simulations of Sant et al. (2013) with the 1-D kinematic cloud model, the normalized integral values are rescaled and approximated by  $P_4(r)$  in log-space with a non-linear least square fit using a Levenberg-Marquardt algorithm (Press et al., 1992), such that

$$P_4(r)|_{jk}^n = \log\left(\frac{C_{jk}^n I_{jk}^n}{N_j N_k}\right) = \log(\mathfrak{I}_{jk}^n) \tag{A7}$$

for  $n \in [0,1]$  and  $j,k \in \{c,d\}$ . The solutions are plotted in Fig. 16 and the parameters are summarized in Tables 4 and 5. Note that  $\mathfrak{I}^n_{dd}$  and  $\mathfrak{I}^n_{cd}$  are piecewise approximated for  $r_d < 40\,\mu\text{m}$  and  $r_d > 40\,\mu\text{m}$ , because of the functional dependency of the shape parameter  $\tilde{\nu}_d$  (cf. Sant et al., 2013, Eq. 12).

The integrals  $I_{\rm cc}^n$  and  $I_{\rm dd}^n$  (cf. Fig. 16a, b, d and e) collapse nicely onto a line due to the normalization by the number density and therefore, allow for a very good fit. As pointed out before,  $I_{0,cd}$  and  $I_{1,cd}$  still exhibit a slight dependency on the shape parameters and the continuity condition applied to the classes (cf. Sant et al., 2013), which results in the broader range of values between 40 and 60  $\mu$ m in Fig. 16c and f depicted by the gray area. Although small, the approximations will not capture this variability. However, since the effect on produced surface precipitation is negligible (on the order of a few percent) they are a reliable alternative to the CPU intensive numerical solution of the integrals.

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**Table 1.** Description of the different acronyms used in Eqs. (1)–(6) to describe the sinks and sources of the different hydrometeors.

| Acronym                             | Description  |
|-------------------------------------|--|
| aut <sub>c</sub> , aut <sub>d</sub> | Autoconversion of cloud droplets and drizzle drops, respectively.                                  |
| $ac_{cd+},ac_{cd-}$                 | Source and sink of the accretion of cloud droplets by drizzle drops, respectively.                 |
| ac <sub>cr</sub> , ac <sub>dr</sub> | Accretion of cloud droplets and drizzle drops by rain drops, respectively.                         |
| $SC_d$ , $SC_r$                     | Selfcollection for drizzle and rain drops, respectively.   |
| agg, acs                            | Aggregation of ice crystals to form snow and accretion of ice crystals and cloud droplets by snow. |
| sec                                 | Secondary production of ice crystals by splintering of snow flakes.                                |
| evp <sub>d</sub> , evp <sub>r</sub> | Evaporation for drizzle and rain drops, respectively.  |
| sub, mlt                            | Sublimation and melting of snow, respectively.   |
| $sed_d,sed_r,sed_s$                 | Sedimentation of drizzle, rain and snow, respectively.   |

**Table 2.** Global annual mean cloud properties and TOA energy budget for the simulations performed with ECHAM5-CTRL and -PROG (2000–2011) as well as observations.

|                      |                               | PROG (PD)                       | PROG (PI)                    | CTRL (PD)                | CTRL (PI)                     | Obs                              |
|----------------------|-------------------------------|---------------------------------|------------------------------|--------------------------|-------------------------------|----------------------------------|
| $P_{tot}$            | $(mmd^{-1})$                  | <del>2.94</del> <u>2.95</u>     | 2.96                         | 2.94                     | 2.95                          | 2.61a                            |
| $P_{strat}$          | $(mmd^{-1})$                  | 1.42                            | 1.42                         | 1.35                     | 1.36                          | _                                |
| $P_{conv}$           | $(mmd^{-1})$                  | <del>1.52</del> 1.53            | 1.54                         | 1.59                     | 1.60                          | _                                |
| TCC                  | (%)                           | <del>63.7 63.5</del>            | <del>63.3</del> <u>63.2</u>  | 64.3                     | <del>64.0 64.1</del>          | 65–75 <sup>b</sup>               |
| <del>LWP</del> -CLWP | $(g m^{-2})$                  | <del>69.2</del> <u>69.1</u>     | 62.5                         | 74.2                     | 67.5                          | 50-81°                           |
| DWP                  | $(g m^{-2})$                  | 2.3                             | <del>2.8</del> <u>2.7</u>    | _                        | _                             | _                                |
| RWP                  | $(g m^{-2})$                  | 30.7                            | <del>31.7</del> -30.9        | _                        | _                             | _                                |
| IWP CIWP             | $(g m^{-2})$                  | 16.2                            | 16.1                         | 10.6                     | 10.5                          | 10-65 <sup>d</sup>               |
| SWP                  | $(g m^{-2})$                  | 2.0                             | 2.0                          | _                        | _                             | _                                |
| WVM                  | $({\rm kg}~{\rm m}^{-2})$     | 26.3                            | 26.3                         | 26.3                     | 26.4                          | 25.4e                            |
| CDNC burden          | $(10^{10}{\rm m}^{-2})$       | 3.0                             | 2.3                          | 3.1                      | 2.4                           | _                                |
| ICNC burden          | $(10^{12}{\rm m}^{-2})$       | 7.7                             | 7.4                          | 10.1                     | 8.7                           | _                                |
| AOD                  | (-)                           | 0.12                            | 0.09                         | 0.14                     | 0.11                          | 0.18 <sup>f</sup>                |
| SW                   | $({\rm W}  {\rm m}^{-2})$     | 235.0                           | 236.8                        | 231.8                    | <del>233.4</del> <u>233.3</u> | 240-244g                         |
| SWCRE                | $({\rm W}  {\rm m}^{-2})$     | - <del>51.0</del> - <u>50.9</u> | -49.7                        | -53.7                    | -52.6                         | $-46 \text{ to } -53^{\text{h}}$ |
| LW                   | $({\rm W}  {\rm m}^{-2})$     | - <del>235.5</del> 235.6        | -235.7                       | - <del>232.2</del> 232.1 | -232.3                        | $-(236-242)^g$                   |
| LWCRE                | $({\rm W}  {\rm m}^{-2})$     | 25.2                            | 25.2                         | 28.4                     | 28.3                          | 27–31 <sup>h</sup>               |
| $F_{net}$            | $(\mathrm{W}\mathrm{m}^{-2})$ | -0.50.57                        | <del>1.1</del> - <u>1.06</u> | -0.30.31                 | <del>1.0</del> - <u>1.03</u>  | _                                |
|                      |                               |                                 |                              |                          |                               |                                  |

<sup>&</sup>lt;sup>a</sup> Precipitation rate taken from the Global Precipitation Climatology Project (GPCP) for the years 1981–2010, provided by the NOAA/OAR/ESRL PSD, Boulder, Colorado, USA (Adler et al., 2003) (http://www.esrl.noaa.gov/psd/).

<sup>&</sup>lt;sup>b</sup> Total cloud cover taken from ISCCP (Rossow and Schiffer, 1999), surface observations (Hahn et al., 1994), MODIS-COLL5 (http://modis-atmos.gsfc.nasa.gov/products.html), and HIRS (Wylie et al., 2005).

<sup>&</sup>lt;sup>c</sup> Liquid water path taken from SSM/I (Greenwald et al., 1993; Weng and Grody, 1994; Wentz, 1997, only for oceans) and ISCCP (Han et al., 1998).

<sup>&</sup>lt;sup>d</sup> Ice water path taken from Fig. 18 of Waliser et al. (2009).

e Precipitable water taken from the NASA Water Vapor Project-MEaSUREs (NVAP-M) dataset as given in Vonder Haar et al. (2012).

f AOD taken from Table 2 of Bellouin et al. (2013).

g SW and LW taken from Fig. 1 of Wild et al. (2013).

h SWCRE and LWCRE taken from Table 4 of Loeb et al. (2009).

**Table 3.** Mean global Global annual mean conversion rate burdens for both cloud liquid and ice for CTRL and PROG for the month of January year 2000.

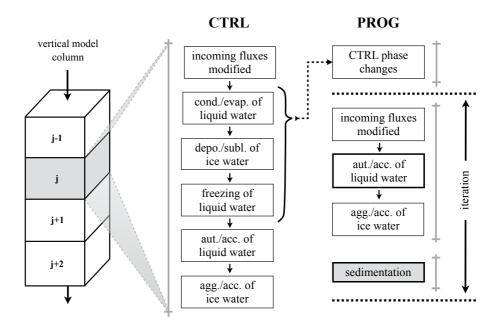
|                        |                                      | PROG                                      | CTRL                                      |
|------------------------|--------------------------------------|---|---|
| $AU(=au_c)$            | $(kg m^{-2} s^{-1})$                 | $5.0 \times 10^{-6}$ $3.9 \times 10^{-6}$ | $6.1 \times 10^{-6}$ $5.2 \times 10^{-6}$ |
| $AC(=ac_{cd}+ac_{cr})$ | $(kg m^{-2} s^{-1})$                 | $1.6 \times 10^{-4}$                      | $8.0 \times 10^{-6}$ $5.2 \times 10^{-6}$ |
| AU/(AU + AC)           | %                                    | <del>3.0_2.3</del>                        | <del>43.5</del> 50.0                      |
| AG(=agg)               | $({\rm kg}{\rm m}^{-2}{\rm s}^{-1})$ | $2.8 \times 10^{-4}$ $2.9 \times 10^{-4}$ | $7.4 \times 10^{-6} 6.8 \times 10^{-6}$   |
| $ACi(=ac_s)$           | $({\rm kg}{\rm m}^{-2}{\rm s}^{-1})$ | $4.8 \times 10^{-5} 4.6 \times 10^{-5}$   | $3.5 \times 10^{-6} 2.5 \times 10^{-6}$   |
| AG/(AG+ACi)            | %                                    | 85.4.86.4                                 | <del>68.0 73.6</del>                      |

**Table 4.** The coefficients and scaling used in the four parameter polynomials (cf. Eq. A6) for approximating the 0th moment, i.e. n = 0, of the integrals in Eqs. (A3)–(A5).

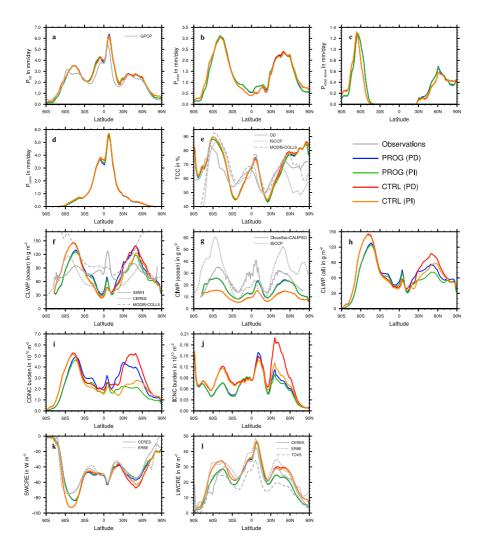
| Coeff.       | $P_4(r) _{cc}^0$         | $P_4(r) _{dd}^0$               |                                   | $P_4(r) _{cd}^0$               |                                   |
|--------------|--------------------------|--------------------------------|-----------------------------------|--------------------------------|-----------------------------------|
|              |                          | $r_{\sf d} <$ 40 $\mu {\sf m}$ | $r_{\sf d} \geq$ 40 $\mu {\sf m}$ | $r_{\sf d} <$ 40 $\mu {\sf m}$ | $r_{\sf d} \geq$ 40 $\mu {\sf m}$ |
| $a_0$        | -22.337                  | -50.158                        | -93.164                           | -49.556                        | -99.931                           |
| $a_1$        |                          |                                |                                   | $-1.8262 \times 10^{-1}$       |                                   |
| $a_2$        |                          |                                |                                   | $-1.6436 \times 10^{-2}$       |                                   |
| $a_3$        | $-2.2943 \times 10^{-3}$ | $-5.1756 \times 10^{-4}$       | $-1.4120 \times 10^{-6}$          | $-5.0922 \times 10^{-4}$       | $-1.1954 \times 10^{-6}$          |
| $C^{0}_{jk}$ | $10^{-5}$                | $10^{-2}$                      |                                   | $10^{-3}$                      |                                   |

**Table 5.** The coefficients and scaling used in the four parameter polynomials (cf. Eq. A6) for approximating the 1st moment, i.e. n = 1, of the integrals in Eqs. (A3)–(A5).

| Coeff.     | $P_4(r) _{cc}^1$         | $P_4(r) _{dd}^1$               |                                   | $P_4(r) _{cd}^1$               |                                   |
|------------|--------------------------|--------------------------------|-----------------------------------|--------------------------------|-----------------------------------|
|            |                          | $r_{\sf d} <$ 40 $\mu {\sf m}$ | $r_{\sf d} \geq$ 40 $\mu {\sf m}$ | $r_{\sf d} <$ 40 $\mu {\sf m}$ | $r_{\sf d} \geq$ 40 $\mu {\sf m}$ |
| $a_0$      | -22.557                  | -50.376                        | -101.199                          | -50.121                        | -109.893                          |
| $a_1$      | $-1.7426 \times 10^{-1}$ | $-2.0805 \times 10^{-1}$       | $-4.3356 \times 10^{-3}$          | $-1.9190 \times 10^{-1}$       | $-4.8496 \times 10^{-3}$          |
| $a_2$      | $-2.6650 \times 10^{-2}$ | $-1.8605 \times 10^{-2}$       | $-8.9089 \times 10^{-5}$          | $-1.6820 \times 10^{-2}$       | $-9.9346 \times 10^{-5}$          |
| $a_3$      | $-2.240 \times 10^{-3}$  | $-5.5578 \times 10^{-4}$       | $-8.7550 \times 10^{-7}$          | $-4.9774 \times 10^{-4}$       | $-7.4402 \times 10^{-7}$          |
| $C^1_{jk}$ | $10^{-12}$               | $10^{-7}$                      |                                   | $10^{-8}$                      |                                   |



**Figure 1.** Schematic flowchart of the CTRL and PROG microphysics schemes. The gray vertical bars denote vertical loops (illustrated by the vertical model column on the left), the boxes in bold are new or changed schemes, and the sedimentation can be subject to a further time splitting if necessary.



**Figure 2.** Zonal annual mean average of **(a)** total precipitation ( $P_{tot}$ ), **(b)** stratiform precipitation ( $P_{strat}$ ), **(c)** stratiform snow precipitation ( $P_{strat}$ , stratiform), **(d)** convective precipitation ( $P_{conv}$ ), **(e)** total cloud cover (TCC), **(f)** cloud liquid water (<u>LWPCLWP</u>) and **(g)** ice water path (<u>LWPCLWP</u>) over the ocean, **(h)** the <u>LWPCLWP</u> over whole globe, **(i)** the CDNC and **(j)** ICNC burden, **(k)** shortwave (SWCRE) and **(l)** longwave cloud radiative forcing (LWCRE) for PROG and CTRL (PD and PI) from 2000–2011. The following observations are used: GPCP, the Global Precipitation Climatology Project (Adler et al., 2003) in **(a)**; DD, surface observations (Hahn et al., 1994) in **(e)**; ISCCP, International Satellite Cloud Climatology Project (Han et al., 1998; Rossow and Schiffer, 1999) in **(e, g)**; MODIS-COLL5, Moderate Resolution Imaging Spectroradiometer, Collection 5 (http://modis-atmos. gsfc.nasa.gov/products.html) in **(e, f)**; SSM/I, Special Sensor Microwave/Imager (Wentz, 1997) in **(f)**; CERES, Clouds and Earth's Radiant Energy System (Wielicki et al., 1996) in **(f)**, **(k)** and **(I)**; Cloud-Sat + CALIPSO, ensemble mean satellite observations of CloudSat and the Cloud-Aerosol Lidar and Infrared Pathfinder Satellite Observationssatellite (Li et al., 2012) in **(g)**; ERBE, Earth Radiation Budget Experiment (http://science.larc.nasa.gov/erbe); and TOVS, TIROS Operational Vertical Sounder (Scott et al., 1999).

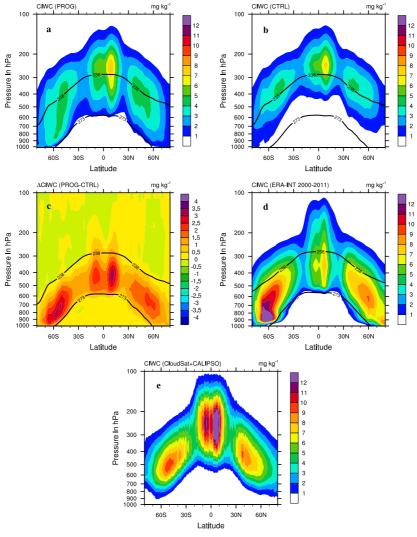


Figure 3. lee Annual mean ice water content of (a) PROG, (b) CTRL, (c) the difference PROG-CTRL and (d) ERA-interim reanalysis, where and (e) CloudSat + CALIPSO, ensemble mean satellite observations of CloudSat and the Cloud-Aerosol Lidar and Infrared Pathfinder Satellite Observations (Li et al., 2012). The solid black lines in (a)–(d) represent the melting and the homogeneous freezing levels at T=273 and T=238 K, respectively.

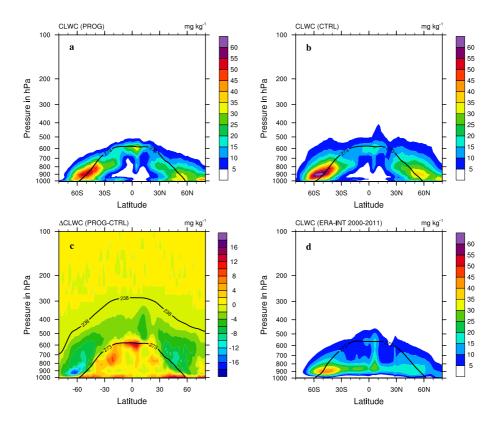
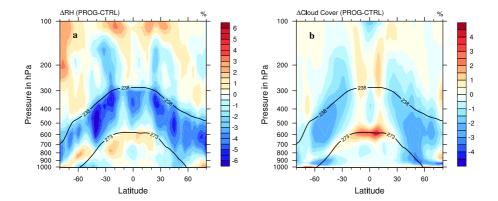


Figure 4. Liquid Annual mean liquid water content of (a) PROG, (b) CTRL, (c) the difference PROG-CTRL, and (d) ERA-interim reanalysis, where the solid black lines represent the melting and the homogeneous freezing (only in c) levels at T = 273 and T = 238 K, respectively.



**Figure 5.** Difference Annual mean difference in (a) RH and (b) TCC of PROG-CTRL, respectively. The solid black lines in (a) represent the melting and the homogeneous freezing levels at T=273 and  $T=238\,\mathrm{K}$ , respectively.

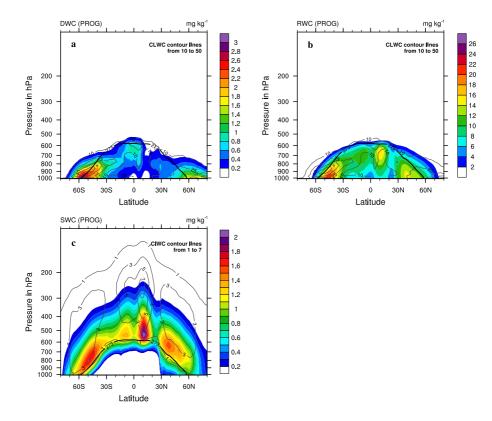
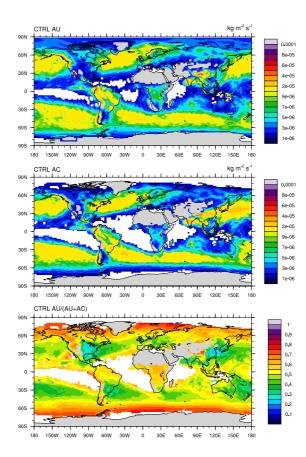
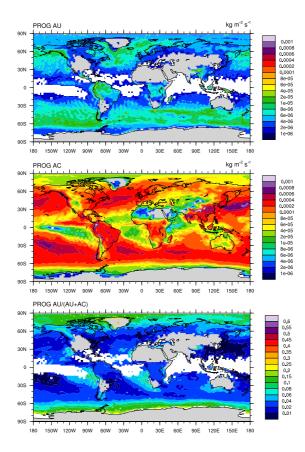


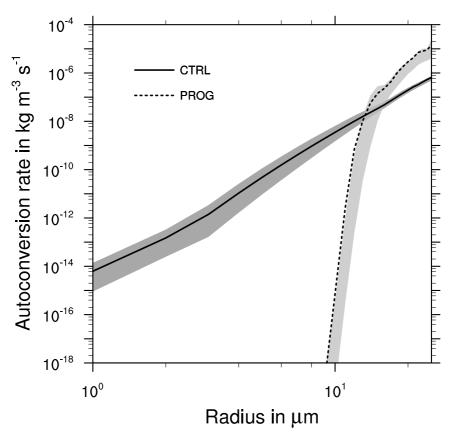
Figure 6. Precipitation Annual mean precipitation water contents from PROG of (a) drizzle and (b) rain with the contour lines of the LWC, such as (c) snow with the contour lines of the IWC and the melting line (solid black line).



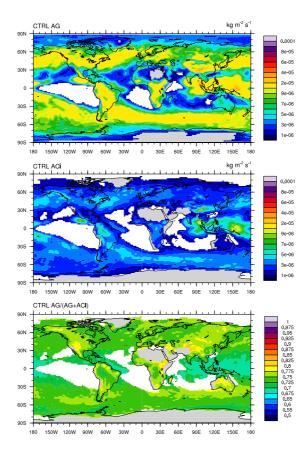
**Figure 7.** Frequency of occurrence of instantaneous **(a)** large-scale surface precipitation rate (only liquid phase), **(b)** liquid water path and **(c)** ice water path for the month of January 2000 for both CTRL (black line) and PROG (gray line) simulations. Note that the x axes are in logarithmic scale. Mean vertically integrated autoconversion (AU) and accretion rates (AC) of CTRL for the month of January year 2000, such as the fraction of AU to the total collection rate (AU + AC).



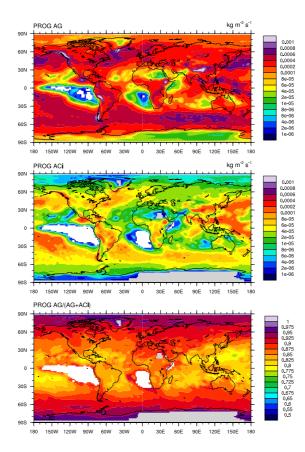
**Figure 8.** Mean vertically integrated autoconversion (AU) and accretion rates (AC) of PROG for the month of January year 2000, such as the fraction of AU to the total collection rate (AU + AC).



**Figure 9.** Dependence of autoconversion rate on effective cloud droplet radius —for the month of July 2000. The lines denote the mean and the gray areas depict the range between the 5th-10th and 95th-90th percentile, respectively.



**Figure 10.** Mean vertically integrated aggregation (AG) and (ice) accretion rates (ACi) of CTRL for the  $\frac{1}{2}$  month of  $\frac{1}{2}$  month of  $\frac{1}{2}$  month of  $\frac{1}{2}$  accretion rate (AG + ACi).



**Figure 11.** Mean vertically integrated aggregation (AG) and (ice) accretion rates (ACi) of PROG for the  $\frac{1}{2}$  month of  $\frac{1}{2}$  month of  $\frac{1}{2}$  month of  $\frac{1}{2}$  and  $\frac{1}{2}$  month of  $\frac{1}{2$ 

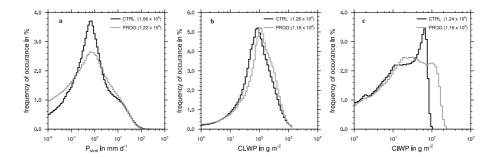
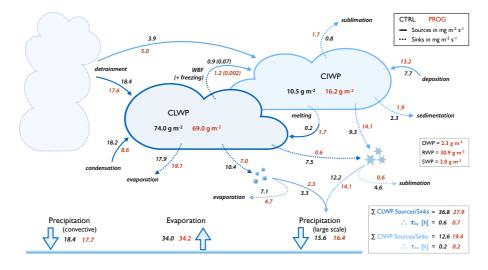


Figure 12. Zonal annual mean average difference between PD and PI Frequency of occurrence of instantaneous (a) large-scale surface precipitation rate, (b) liquid water path, (b) total cloud cover, and (c) ice water path for the CDNC month of July 2000 for both CTRL (black line) and (d) ICNC burden, (e) shortwave, (f) longwave PROG (gray line) simulations. Note that the x axes are in logarithmic scale and (g) net TOA the numbers in the brackets correspond to the total number of events plotted.



**Figure 13.** Illustration of the hydrological cycle with all relevant global annual mean (vertically integrated) conversion rates and the cloud radiative forcing for lifetimes of CTRL (black) and PROG red numbers) for the year 2000.

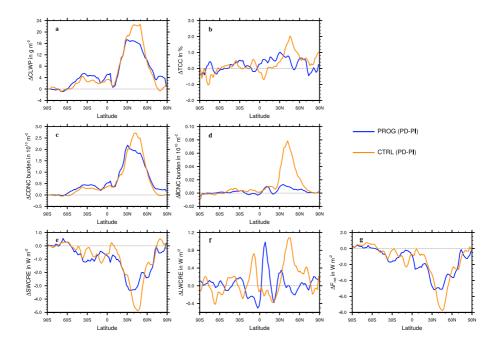
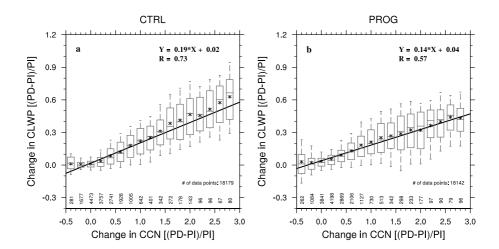
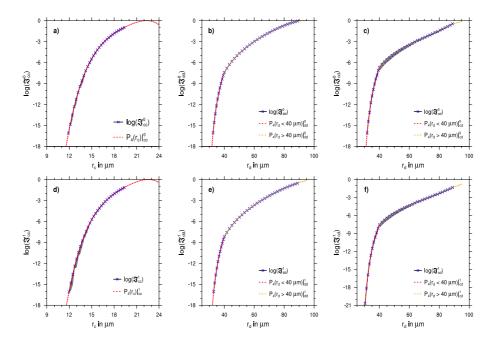


Figure 14. Zonal annual mean difference between PD and PI of (a) liquid water path. (b) total cloud cover, (c) the CDNC and (d) ICNC burden, (e) shortwave, (f) longwave and (g) net TOA cloud radiative forcing for CTRL and PROG for the years 2000–2011.



**Figure 15.** Boxplots and linear regression lines of the relative changes [(PD – PI)/PI] in annual mean LWP\_CLWP vs. the relative changes in annual mean activated CCN concentrations for (a) the (a) CTRL and (b) the PROG simulations, following the analysis by Wang et al. (2011). CCN concentrations are averaged over the lowest 8 model levels (surface to about 800 hPa). The boxes represent the median (middle line) and the 25th and 75th percentile (lower and upper line), respectively, of the data within that bin. The whiskers show the 5th 10th and 95th 90th percentile (lower and upper whisker), respectively, and the stars represent the mean. The numbers above the x-axis represent the number of data points in a given bin. The linear regression lines and the Pearson correlation coefficients (R) are given in the right upper corner, respectively.



**Figure 16.** Approximation of Eqs. (A3)–(A5) using a four parameter polynomial  $P_4(r)$  (cf. Eq. A7) from an ensemble of simulations with the 1-D kinematic cloud model (Sant et al., 2013). The gray area in the figures represents the range of the minimum and maximum values. Note that the gray area in **(a, b)** and **(e)** are very small and hardly recognizable.