

1 **Comprehensive tool for calculation of radiative fluxes:**
2 **Illustration of shortwave aerosol radiative effect**
3 **sensitivities to the details in aerosol and underlying**
4 **surface characteristics**

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12

13 **Abstract**

14 The evaluation of aerosol radiative effect on broadband hemispherical solar flux is often
15 performed using simplified spectral and directional scattering characteristics of atmospheric
16 aerosol and underlying surface reflectance. In this study we present a rigorous yet fast
17 computational tool that accurately accounts for detailed variability of both spectral and
18 angular scattering properties of aerosol and surface reflectance in calculation of direct aerosol
19 radiative effect. The tool is developed as part of the GRASP (Generalized Retrieval of
20 Aerosol and Surface Properties) project. We use the tool to evaluate instantaneous and daily
21 average radiative efficiencies (radiative effect per unit aerosol optical thickness) of several
22 key atmospheric aerosol models over different surface types. We then examine the differences
23 due to neglect of surface reflectance anisotropy, non-sphericity of aerosol particle shape and
24 accounting only for aerosol angular scattering asymmetry instead of using full phase function.
25 For example, it is shown that neglecting aerosol particle nonsphericity causes mainly
26 overestimation of the aerosol cooling effect and that magnitude of this overestimate changes
27 significantly as a function of solar zenith angle (SZA) if only asymmetry parameter is used
28 instead of detailed phase function. It was also found that the nonspherical–spherical
29 differences in the calculated aerosol radiative effect are not modified significantly if detailed

1 BRDF (Bidirectional Reflectance Distribution Function) is used instead of Lambertian
2 approximation of surface reflectance. Additionally, calculations show that usage of only
3 angular scattering asymmetry, even for case of spherical aerosols, modifies dependence of
4 instantaneous aerosol radiative effect on SZA. This effect can be canceled for daily average
5 values, but only if sun reaches the zenith, otherwise a systematic bias remains. Since the daily
6 average radiative effect is obtained by integration over a range of SZAs, the errors vary with
7 latitude and season. In summary, the present analysis showed that use of simplified
8 assumptions causes systematic biases, rather than random uncertainties, in calculation of both
9 instantaneous and daily average aerosol radiative effect. Finally, we illustrate application of
10 the rigorous aerosol radiative effect calculations performed as part of GRASP aerosol
11 retrieval from real POLDER/PARASOL satellite observations.

12

13 **1 Introduction**

14 Direct atmospheric aerosol radiative forcing remains one of the most uncertain components in
15 evaluation of Earth's climate change (Andreae et al., 2005; Hansen et al., 2011). Although
16 aerosols are generally recognized as having a negative radiative effect (cooling) on the
17 surface-atmosphere system, in some conditions aerosol can also have a positive radiative
18 effect (warming). The aerosol cooling effect is produced by reflecting solar radiation back to
19 space, i.e. scattering in the upward direction. Depending on their composition, aerosol can
20 also heat due to absorption of the incoming solar radiation. However, not only properties of
21 aerosol, but also of the underlying surface are decisive for the sign of the aerosol radiative
22 effect. For example, the same particles can decrease (warming effect) or increase (cooling
23 effect) the planetary albedo depending on whether the underlying surface is a bright desert or
24 dark ocean. Regardless of warming or cooling from the point of view of top of atmosphere
25 albedo, aerosols always warm the atmospheric layer if their absorption is not zero. In
26 addition, the aerosols generate heating effect in thermal infrared spectrum, primary caused by
27 large mineral dust particles that strongly absorb outgoing terrestrial radiation, e.g. (Legrand et
28 al., 2001). The TIR effect is similar to influence of greenhouse gasses and thus counteract the
29 scattering effect in the solar spectrum. For clarity of the analysis performed in this study it is
30 important to recall that the term aerosol direct radiative forcing, which is defined as
31 perturbation of radiative fluxes due to human-induced component only, is therefore different
32 from the term radiative effect. Aerosol radiative effect implies the difference between

1 radiative fluxes in aerosol-free and aerosol-laden atmospheric conditions, e.g. (Kaufman et
2 al., 2005; Remer and Kaufman, 2006). Using measurements, one can assess the aerosol
3 radiative effect by referring to aerosol-free conditions. In climate models, however, it is
4 feasible to evaluate forcing by referring to background or pre-industrial aerosol. Therefore,
5 because of possibility to control numerous aerosol emission and transport processes,
6 evaluation of radiative forcing of climate relies mostly on chemical transport and general
7 circulation models. In order to reduce dependence on assumptions that take place in the
8 models, important steps towards evaluation of aerosol direct radiative effect are also done
9 using global aerosol and broadband flux observations from satellite and ground-based remote
10 sensing (Bellouin et al., 2005; Boucher and Tanré, 2000; Remer and Kaufman, 2006; Su et
11 al., 2013; Yu et al., 2006; Yu et al., 2004; Zhou et al., 2005). The observation-based
12 evaluations of aerosol radiative effect open opportunities for inter-comparison with models
13 and leads to improvement in assessment of aerosol radiative effect on climate. Therefore,
14 there is an interest in continuation of the measurement-based evaluation of the aerosol
15 radiative effect and examination of possible sources of uncertainty. For example, description
16 of angular and spectral features of scattering properties of aerosol and underlying surface is
17 often simplified. The reasons for using these simplifications are usually the lack of
18 information regarding the details of these properties and a need in substantial reduction of
19 computation time required for rigorous flux computations. For instance, accurate modeling of
20 scattering by non-spherical particles and directional reflectance of surface is challenging and
21 therefore often neglected. Recent advancements in retrievals of aerosol optical characteristics
22 from ground and space remote sensing and from combination of sensors show capabilities to
23 provide more detailed properties. For example, aerosol size distribution, complex refractive
24 index, single scattering albedo and non-spherical fraction become available not only from
25 ground-based photometric observations (Dubovik et al., 2002b; Dubovik et al., 2006), but
26 also from space sensors (Dubovik et al., 2011; Dubovik et al., 2014) providing advantage of
27 large spatial coverage. The retrievals from space provide also information about the surface
28 spectral albedo or BRDF parameters. In addition, the aerosol layer height can be retrieved
29 using even passive polarimetric sensors (Dubovik et al., 2011; Tanre et al., 2011), while a
30 combination of passive and active sensors shows sensitivity to vertical profiles of extinction
31 by aerosol in fine and coarse mode fractions (Lopatin et al., 2013). These coming up
32 enhanced remote sensing retrievals imply possibility of more accurate aerosol radiative effect
33 computation that largely rely on the measurements and reduced level of assumptions. For

1 example, a close agreement is found in an inter-comparison of measured downward solar flux
2 at the surface with fluxes computed as part of the AERONET product. The studies conducted
3 in the framework of a field campaign (Derimian et al., 2008), on a global scale (Garcia et al.,
4 2008) and in specific case studies (Derimian et al., 2012) show that the computed broadband
5 solar flux generally agrees with the measured flux to within 5 to 10%; note that accuracy of
6 solar flux measurements themselves is on the order of 5 %. The agreement between simulated
7 and measured flux is remarkable yet to be expected if the computational approach, employed
8 here is understood. The main advantage of the approach is that the retrieved aerosol and
9 surface properties should fit the measured radiances at given wavelengths within few percent,
10 as it requires the inversion algorithm. Obviously, an interpolation or extrapolation outside of
11 the nominal wavelengths is needed and the errors may accumulate during spectral radiances
12 calculations and after radiances integration into broadband flux. Essentially, it also implies
13 that the retrieved aerosol models that satisfy fit of simulated to measured radiances in
14 inversion algorithms should also accurately reproduce the spectral variability of aerosol
15 properties in the simulation of broadband flux. Accurate and high spectral resolution
16 computations of radiances by accounting for spectral variability of gaseous absorption and
17 detailed aerosol characteristics, such as detailed phase function, that strongly depend on
18 particle sizes, shapes and index of refraction, should increase the accuracy of the simulated
19 flux. For example, the importance of accounting for particle nonsphericity in calculation of
20 desert dust radiative forcing is addressed in several discussions (Bellouin et al., 2004;
21 Derimian et al., 2008; Kahnert and Kylling, 2004; Kahnert et al., 2005; Mishchenko et al.,
22 1995; Yi et al., 2011). Indeed, nonsphericity of the particles shape is often neglected in
23 aerosol radiative effect computations, mainly due to necessity to reduce computational time.
24 Hence, an assumption is made that the differences in angular scattering by spherical and
25 nonspherical particles are canceled when all contributions of scattered light are summed up
26 into the total hemispherical flux. Also, the computation approach generally implies usage of
27 the asymmetry parameter, which is an integrated value and therefore differences in the aerosol
28 phase function of spheres and spheroids are expected be averaged out. However, Kahnert and
29 Kylling (2004) and Kahnert et al. (2005) conducted a detailed analysis of asymmetry
30 parameter sensitivity to particle shape and concluded that the use of spherical particles model
31 might be among the major error sources in broadband flux simulations. In the work by
32 Derimian et al. (2008) the effect of particles nonsphericity on forcing was evaluated using
33 detailed phase function in the flux calculations. The nonsphericity effect was evaluated for

1 cases of dust and mixed aerosol type during biomass burning season in western Africa. The
2 computations revealed that neglecting of particles nonsphericity leads to a systematic
3 overestimation of the aerosol cooling effect by up to 10 %; the bias was pronounced in
4 instantaneous and daily average values. It was also noted that the magnitude of the
5 overestimation depends on the magnitude of aerosol absorption and Aerosol Optical
6 Thickness (AOT or τ). Later general sensitivity tests by Yi et al (2011) evaluated the errors in
7 radiances and flux due to spherical particles approximation resulted in conclusions consistent
8 with effects observed by Derimian et al. [2008] in the specific case study. We would like to
9 emphasize here that features of aerosol directional scattering are also important for accurate
10 modeling of diurnal dependence of forcing, i.e. dependence of aerosol instantaneous forcing
11 on the SZA. This SZA dependence of aerosol radiative effect at top of atmosphere appeared
12 in an earlier simple expression developed for calculations of Earth-atmosphere albedo
13 perturbation (Lenoble et al., 1982). Later it was confirmed by exact radiative transfer
14 computations, e.g. (Bellouin et al., 2004), taken into account in space instrument forcing
15 studies using POLDER (Boucher and Tanré, 2000) and MODIS (Remer and Kaufman, 2006),
16 and using AERONET retrievals, e.g. (Derimian et al., 2012; Derimian et al., 2008; Garcia et
17 al., 2012). It is also to mention that the diurnal dependence of forcing is influenced by
18 directional properties of the underlying surface. The effect was discussed by Yu et al. [2004]
19 for land and by Bellouin et al. [2004] for ocean using the BRDF.

20 In the current study we introduce a rigorous computational tool for broadband flux
21 simulations and demonstrate the importance of detailed representation of aerosol and surface.
22 We apply our simulation for (i) evaluating radiative effect of several key aerosol models, then
23 (ii) we stress importance of diurnal dependence (dependence on SZA) of the aerosol radiative
24 effect and (iii) examine the effects of assumptions and using of simplified representations of
25 aerosol phase function, particle shape and directional properties of surface reflectance. It is
26 often expected, that the details of aerosol and surface optical properties are not really
27 important because the flux is an integral product of spectral and angular properties of
28 atmospheric radiation. Therefore we intend to clarify if any cancelations of uncertainties
29 appear in the integrated broadband hemispherical flux due to coexisting assumptions on
30 aerosol and surface directional scattering.

31 Thus, the below paper is organized as the following. Section 2 provides description of the flux
32 computational tool. Section 3 contains the description of aerosol models used in the

1 sensitivity tests. In Sect. 4 and 5 we analyze importance of diurnal dependence of
2 instantaneous aerosol radiative effect, which also varies as a function of aerosol
3 characteristics and surface albedo model. Section 6 provides the discussion about complexity
4 of evaluation of the nonspherical – spherical difference in aerosol radiative effect due to a
5 concurrent change in directional redistribution of scattering and spectral extinction cross
6 sections of volume-equivalent spheres and spheroids. Section 7 discusses the errors appearing
7 in radiative effect calculations due to use of simplified representation of aerosol directional
8 scattering by asymmetry parameter. Finally, Sect. 8 includes an example of aerosol radiative
9 effect computation for a part of Africa using GRASP (Generalized Retrieval of Aerosol and
10 Surface Properties) algorithm (Dubovik et al., 2014) applied for POLDER/PARASOL
11 observations.

12

13 **2 Computational code description**

14 The initial version of this broadband solar flux computational tool was originally built in the
15 AERONET operational code (Dubovik and King, 2000), the performances were studied and
16 inter-comparisons with the ground-based flux measurements conducted on global scale
17 (Garcia et al., 2008) and in specific case studies (Derimian et al., 2008). As described below,
18 the tool is significantly revised and integrated into the GRASP unified algorithm for
19 characterizing atmosphere and surface. Thus, at present, the calculations can be performed as
20 part of measurements processing and the radiative effect estimations can be provided in the
21 framework of GRASP retrieval product. It is also possible to use the computational tool in
22 various types of independent research calculations.

23 Computations of broadband solar flux in spectral interval from 0.2 to 4.0 μm and of aerosol
24 radiative effect are based on forward modeling of atmospheric radiances and flux simulations
25 employed in the GRASP algorithm which inherits aerosol representation from AERONET
26 retrieval code (Dubovik and King, 2000; Dubovik et al., 2006; Sinyuk et al., 2007). Figure 1
27 shows a general structure of the aerosol radiative effect simulation logistic. The input includes
28 ozone and water vapor concentrations and a set of “retrieved parameters” (see *Dubovik et al.*,
29 2011; *Dubovik et al.*, 2014]) that includes aerosol volume size distribution; spectral real and
30 imaginary part of aerosol complex refractive index; fractions of spherical particles,
31 parameters of aerosol vertical distribution and parameters of BRDF surface reflectance. It also
32 includes information about maximal sun elevation and daylight duration that is required for

1 evaluation of 24h average radiative effect. It should be noted that in the presented studies the
2 vertical distribution of aerosol extinction was fixed and assumed as a Gaussian distribution
3 with maxima at an altitude of 1 km and standard deviation of 0.7. However, if a realistic
4 aerosol vertical profile is available, it can be included as part of the input and used in the
5 calculations. For example, GRASP retrievals provide aerosol medium height from PARASOL
6 observations (Dubovik et al., 2011) and GaRRliC/GRASP retrieval provide detailed vertical
7 profiles from joint inversion of ground-based photometer and lidar data (Lopatin et al., 2013).
8 The gaseous content in the atmospheric column is assessed from combination of retrievals,
9 climatology values and standard atmospheric models. In the presented computations, for
10 instance, instantaneous water vapor content is retrieved by AERONET using the absorption
11 differential method at the 0.94 μm channel (Smirnov et al., 2000), the total ozone content is
12 obtained from the monthly climatology values of NASA Total Ozone Mapping Spectrometer
13 (TOMS) and US standard 1976 atmosphere model is used for other gases and atmospheric
14 gaseous profiles. The aerosol optical characteristics calculated at 208 spectral intervals,
15 gaseous absorption k-distribution, and surface reflectance (Lambertian or BRDF based) are
16 then supplied into atmospheric radiative transfer calculations. The aerosol optical thickness,
17 Single Scattering Albedo (ω_0), and phase function ($P(\Theta)$) (or phase matrix) are calculated for
18 each of 208 spectral intervals using the size distribution, complex refractive index and
19 fraction of spherical particles. The missing spectral values of the complex refractive index are
20 linearly interpolated or extrapolated from the values provided in the input since spectral
21 behaviors of aerosol complex refractive index in the solar spectrum is sufficiently smooth.
22 The details of the aerosol phase function are taken into account using a 12-moment expansion
23 of the Legendre polynomial, however, usage of asymmetry parameter only (first moment
24 expansion of the Legendre polynomial) is also possible. The aerosol single scattering
25 properties are modeled using pre-computed kernel look-up tables produced for a set of size
26 parameters, complex refractive indices and fraction of spherical particles. The fixed aspect
27 ratio distribution of prolate/oblate spheroids, used for the nonspherical aerosol component, is
28 derived (Dubovik et al., 2006) as a best fit of detailed phase matrices measured in the
29 laboratory by Volten et al. [2001]. This approach enables to conduct the flux simulations in a
30 reasonable computational time even when a nonspherical aerosol model and detailed
31 representation of spectral aerosol phase function are taken into account. The effects of
32 multiple scattering in broadband integration are accounted using accurate radiative transfer
33 calculations by vector successive order of scattering code (Lenoble et al., 2007) modified by

1 adding several flexibilities desirable for aerosol retrievals (see (Dubovik et al., 2011)). It
2 should be noted that initial version of flux calculations used in the AERONET code employs
3 discrete ordinates radiative transfer code (DISORT) (Nakajima and Tanaka, 1988; Stamnes et
4 al., 1988). The gaseous absorption (H₂O, CO₂, and O₃) is accounted using the same approach
5 as the one adapted in a module of the radiative transfer model GAME (Global Atmospheric
6 Model) (Dubuisson et al., 1996; Dubuisson et al., 2006; Roger et al., 2006). Specifically,
7 gaseous absorption is calculated by utilizing the correlated k-distribution (Lacis and Oinas,
8 1991) that allows broadband flux simulations with acceptably short computational time. The
9 coefficients of the correlated k-distribution have been estimated from reference calculations
10 using a line-by-line code (Dubuisson et al., 2004). Modeling of the surface reflectance is done
11 either by BRDF model (using various models as described by Dubovik et al. [2011]) or using
12 Lambertian approximation. In current sensitivity tests we used the Li-Ross BRDF model for
13 calculation of the land surface reflectance (Rahman et al., 1993; Roujean et al., 1992; Wanner
14 et al., 1995). The surface spectral reflectance was modeled using climatological values
15 provided by MODIS, the missing spectral values are linearly interpolated or extrapolated, in a
16 similar manner as for the complex refractive index. Thus, spectral variability of aerosol
17 optical characteristics, gaseous absorption, molecular scattering and surface albedo are
18 carefully taken into account in the computation of spectral radiances that afterwards are
19 integrated into the broadband solar flux.

20 As mentioned above, several important revisions of the radiative effect computation tool were
21 done as part of GRASP project advancement (Dubovik et al., 2011). The significant reduction
22 of computational time of spectral radiances was one of these advancements. Another
23 advantage, compare to the original tool, is that the radiative transfer code implemented in the
24 GRASP also accounts for polarization and can account for both aerosol phase matrix and
25 surface BPDF (Bidirectional Polarization Distribution Function). Note, that in the presented
26 sensitivity calculation the polarization effects were not considered, but they are accounted for
27 in application for POLDER/PARASOL observations. Finally, the most important
28 advancement is that all the aerosol and surface properties, that are necessary for the
29 broadband solar flux calculation, can be derived simultaneously by GRASP as retrieval
30 products, e.g. using POLDER/PARASOL observations. In addition, there is an interest to
31 interpret new aerosol retrievals produced by GRASP on the level of direct aerosol radiative
32 effect. The radiative effect calculation strategy described above is therefore driven by this
33 motivation and is tied to the retrieved characteristics provided by GRASP. Spectral dependent

1 properties, such as aerosol complex refractive index, BRDF and BPDF parameters derived
 2 only at the fixed instrumental channels, are used after interpolation or extrapolation in the
 3 same manner as was done in initial version of the computational tool. The gas absorptions
 4 calculations using the correlated k-distribution method are adapted for convenience of
 5 satellite measurement processing. For example, in order to circumvent the time-consuming
 6 convolutions among all the gas species, only the water vapor and ozone contents are set to the
 7 real-time values obtained from satellite retrievals (e.g., POLDER, TOMS). The other gases,
 8 whose concentration ratios to the carbon dioxide vary little among different atmospheric
 9 profiles, is considered as one mixed gas based on their concentration ratios in the US standard
 10 1976 atmosphere model (the CO₂ concentration was actualized to a more recent).

11 The results of calculations include: instantaneous upward and downward fluxes (with and
 12 without aerosol); instantaneous net radiative effect at Bottom and Top of Atmosphere (BOA
 13 and TOA), and in the atmospheric layer; 24-hour average net radiative effects (BOA, TOA
 14 and atmospheric layer); and vertical profiles of aerosol radiative effect for a given aerosol
 15 profile. The aerosol net radiative effect is defined as difference between downwelling and
 16 upwelling fluxes at a given atmospheric layer in aerosol-free and aerosol-laden conditions,
 17 that is, at the BOA, the net radiative effect is defined as:

$$18 \quad \Delta F_{BOA}^{Net} = (F_{\downarrow BOA}^a - F_{\uparrow BOA}^a) - (F_{\downarrow BOA}^0 - F_{\uparrow BOA}^0), \quad (1)$$

19 where $F_{\downarrow BOA}^a$ and $F_{\uparrow BOA}^a$ are downwelling and upwelling fluxes in aerosol-laden conditions and
 20 $F_{\downarrow BOA}^0$ and $F_{\uparrow BOA}^0$ are downwelling and upwelling fluxes in aerosol-free conditions. The
 21 aerosol radiative effect at the TOA is defined similarly and can be written as follows:

$$22 \quad \Delta F_{TOA}^{Net} = (F_{\downarrow TOA}^a - F_{\uparrow TOA}^a) - (F_{\downarrow TOA}^0 - F_{\uparrow TOA}^0) = F_{\uparrow TOA}^0 - F_{\uparrow TOA}^a, \quad (2)$$

23 because at the TOA the downwelling (extraterrestrial) flux is the same either for aerosol-free
 24 or aerosol-laden conditions. Difference between the net TOA and net BOA radiative effects is
 25 the atmospheric radiative effect (ATM) that represents the energy trapped in the atmosphere
 26 due to the aerosol presence:

$$27 \quad \Delta F_{ATM}^{Net} = \Delta F_{TOA}^{Net} - \Delta F_{BOA}^{Net}. \quad (3)$$

28 The 24-hour average aerosol radiative effect is computed by integration of instantaneous
 29 values up to minimal SZA of a given day of the year and at given latitude. These

1 instantaneous values are calculated with a half-degree SZA resolution or Gaussian quadrature
2 in the GRASP version. Knowing the daylight duration and minimal SZA of that day, the SZA
3 interval is converted into a corresponding time interval by which the instantaneous values are
4 integrated over the time of the daylight duration. The obtained integral represents energy
5 perturbed by aerosols during the daylight time. This value is then divided by 24 hours to get
6 the perturbation per day – the daily average radiative effect.

7

8 **3 Aerosol models**

9 Several key aerosol models are selected in order to evaluate their radiative effect under
10 different assumptions. The average aerosol models are derived from all available years of
11 AERONET observations and include: dust and mixture of dust with biomass burning aerosol
12 in the Dakar site (also known as Mbour); biomass burning aerosol in the Mongu site;
13 urban/industrial pollution in the Paris site; and mixture of dust with urban/industrial aerosol in
14 the Kanpur site. Except for Dakar, the AERONET sites and aerosol models are selected
15 pursuing the works of Dubovik et al., (2002a) and Giles et al., (2012). The Dakar site was
16 studied in the framework of the AMMA campaign (Haywood et al., 2008) and is
17 characterized by mixture of dust with biomass burning aerosol during the dry season in
18 January and February and by desert dust only starting from March, e.g. (Derimian et al., 2008;
19 Léon et al., 2009). The aerosol characteristics are derived using version 2, level 2 almucantar
20 AERONET product and applying criteria recommended in (Dubovik et al., 2002a).
21 Additionally, a seasonal criterion is applied for the Mongu site in southern Africa, where the
22 biomass burning aerosol model is derived during the summer period that is known as a peak
23 of the biomass burning season. It has to be mentioned that at this site the aerosol absorption
24 was found as varying within the biomass burning season (Eck et al., 2013), thus variability in
25 the biomass burning radiative efficiency is also expected. For the purpose of our study we
26 take, however, only an averaged characteristic and select August and September as the
27 months with highest aerosol optical thickness and maximal number of observations. An
28 additional criterion that was used to distinguish the aerosol type is the value of Ångström
29 exponent (\AA) between wavelengths of 870 nm and 440 nm. The Ångström exponent below 0.6
30 is attributed to dust, between 0.8 and 1.2 to a mixed aerosol type in Dakar and Kanpur sites,
31 above 1.6 for urban/industrial pollution in Paris, and above 1.6 for the biomass burning in the
32 Mongu site. Average fractions of spherical particles obtained for these aerosol types were also

1 examined. The values logically correspond to the defined aerosol models, that is: 3% for dust
2 in Dakar; 5% for mixture of dust and biomass burning in Dakar; 21% for mixture of dust and
3 urban/industrial in Kanpur; 98% for urban/industrial in Paris; and 99% for biomass burning in
4 Mongu. These values were also employed in calculations of aerosol radiative effect presented
5 in this study. Other details of the selected aerosol models are presented in Table 1 and Fig. 2.
6 In order to facilitate a straightforward inter-comparison of relative importance of fine and
7 coarse modes of different aerosol models, the volume size distributions in Fig. 2a) are
8 normalized by total volume concentrations, i.e. their integration over radii is equal to unity.
9 Spectral dependences of aerosol optical thicknesses are normalized by their maximal values
10 and are intercomparable in Fig. 2b; the related values of $\hat{a}(870\text{nm}/440\text{ nm})$ are also presented
11 in the figure. Based on the derived size distributions and complex refractive index, the
12 spectral ω_0 and asymmetry parameter (g) are calculated over entire range of the solar
13 spectrum, to that end the complex refractive index is linearly interpolated between the
14 nominal wavelengths and is fixed to the last value beyond them (see Table 1).

15 Note that the computed g and ω_0 have quite strong spectral variability (Fig. 2c, d)), which
16 illustrates strong dependence of g and also of ω_0 on the ratio of particles size to wavelength.
17 For example, in the cases of biomass burning and urban aerosol models, the $\omega_0(\lambda)$ is changing
18 even if imaginary part of refractive index is spectrally constant (see Table 1 and Fig. 2c)).
19 After having a maximum at short wavelengths, the $\omega_0(\lambda)$ increases again at longer
20 wavelengths for all aerosol models where the bimodal size distribution is strongly pronounced
21 (i.e. except for dust). It is due to increasing scattering effectiveness of fine and coarse modes
22 at short and long wavelengths, respectively. The scattering effectiveness in case of dust
23 aerosol model is increasing only at long wavelengths. The spectral dependence of $g(\lambda)$ is also
24 noteworthy. For aerosol models with significant fine mode, it could be expected that with
25 decrease of the particle size relative to wavelength, the asymmetry parameter will
26 monotonously decrease. However, g starts to increase (increase of scattering in forward
27 hemisphere) at long wavelengths for all aerosol models, apparently due to bimodality of the
28 size distributions and increasing contribution of the coarse mode.

29 A pronounced spectral dependence in the directional scattering can also be seen in Fig. 3 that
30 shows $P11(\theta) \cdot AOT_{scat}$, where $P11(\theta)$ is the phase function that fulfill the following
31 normalization condition:

$$\frac{1}{2} \int_0^\pi P_{11}(\theta) \cdot \sin \theta d\theta = 1. \quad (4)$$

Therefore, the presented examples of significant spectral variability of ω_0 , g and directional scattering emphasize importance of accurate accounting for the aerosol spectral characteristics in the broadband flux simulations. On the other hand, it is fair to say that the solar constant is rapidly decreasing outside of the visible interval, which partially diminish inaccuracy in aerosol spectral characteristics. Another curious observation can be made regarding the single scattering albedo of the mixed aerosol type. In both cases of mixture (dust with biomass burning and dust with urban/industrial pollution) the single scattering albedo is lower than the one estimated using a simple additive combination of each component. A simple additive combination of single scattering albedos is valid for aerosol external mixture case, though apparently can hardly explain the low single scattering albedo values observed for the mixed aerosol type other than by presence of excessive absorption of pollution in Kanpur and of smoke in Dakar. The existence of internal mixture of different chemical elements (e.g. presence of absorbing material on the surface of coarse mode particles) is another explanation of that decrease the scattering effectiveness.

16

17 **4 Diurnal dependence of instantaneous forcing**

Strong dependence of instantaneous aerosol radiative effect on SZA implies importance for (i) proper inter-comparison of instantaneous values assessed in different time and location and (ii) evaluation of the daily average radiative effect, which is obtained by integration over corresponding range of SZAs in a given day and location. In order to examine dependence on SZA, diurnal radiative efficiencies are calculated for the presented above aerosol models. The radiative efficiencies are calculated with respect to AOT at 550 nm and over Lambertian ocean surface albedo. The aerosol radiative efficiency is used in order to examine influence of different aerosol type and not of concentration, which is supposed to be ruled out because efficiency is defined as radiative effect per unit AOT. One should remember, however, that the aerosol radiative effect is not a linear function of AOT, e.g. discussed by Markowicz et al. (2008). Thus, for a consistent inter-comparison of radiative efficiencies calculated for different aerosol models, we choose to set all corresponding AOTs at 550 nm to unit.

First observation that can be drawn from the Fig. 4 is that not only magnitude, but also the shape of the curves of radiative efficiency vs. $\cos(\text{SZA})$ depends on the aerosol type. Note

1 that the $\cos(\text{SZA})$ is used hereafter since this variable appears in the radiative transfer
 2 equation. This shape is essentially linked to the differences in aerosol phase functions.
 3 Significant dependence of the instantaneous radiative effect on SZA also implies that its
 4 accurate computation is important for the daily average radiative effect. Hence, a proper
 5 analysis and inter-comparison of not only instantaneous, but also of the daily average aerosol
 6 forcings should respect the range of SZAs. Consistency in the daylight time duration should
 7 also be taking into account if one intends to attribute the differences in the daily average
 8 radiative effect to differences in aerosol type or concentration. Strictly speaking, the same
 9 aerosol type and concentration over same surface and in same location, but in different time
 10 of the year, or on the same day but in different latitudes, will give different value of daily
 11 average forcing. Otherwise, for a consistent inter-comparison, a standard can be assumed, for
 12 example, the sun reaches the zenith ($\text{SZA} = 0^\circ$) and the daylight fraction is 0.5 (daylight
 13 duration is 12 hours). Coming back to the Fig. 4, a difference can also be noted in angular
 14 dependence of aerosol radiative effect at TOA and BOA. At TOA the negative radiative effect
 15 starts to decrease for higher sun elevation, but at BOA continues to increase or stays more or
 16 less constant (depending on the aerosol model). Remembering that the difference between
 17 TOA and BOA forcings is the atmospheric forcing, it means that efficiency of atmospheric
 18 layer heating due to the aerosol presence is increasing for increasing sun elevation.

19

20 **5 Directional properties of surface reflectance**

21 It is known that aerosol radiative impact on the Earth's albedo depends not only on the
 22 aerosol properties but also on reflectance of the underlying surface. In general, to describe
 23 surface reflectance accurately, the Bidirectional Reflectance Distribution Function (BRDF) is
 24 required. The BRDF depends on illumination and scattering geometries, e.g. (Litvinov et al.,
 25 2011; Litvinov et al., 2012). Therefore, diurnal dependence of aerosol radiative effect is also
 26 expected to vary with respect to SZA and directional properties of the surface reflectance. As
 27 a first approximation of surface reflectance description such characteristic as "black-sky"
 28 albedo (also known as Directional Hemispherical Reflection (DHR)) is often used. It can be
 29 defined through the integrals of BRDF [*Schaepman-Strub et al.*, 2006]:

$$30 \quad DHR(\lambda, \vartheta_0) = \int_0^{2\pi} \int_0^{\pi/2} BRDF(\lambda, \vartheta_0, \vartheta_v, \varphi) \cos \vartheta_v \sin \vartheta_v d\vartheta_v d\varphi, \quad (5)$$

1 where ϑ_v and ϑ_0 are reflected and solar zenith angles; φ is difference of azimuth angles of
2 reflected and solar directions; λ is the wavelength of incident radiation.

3 Figure 5a) shows an example of surface “black-sky” albedo dependence on SZA at three
4 AERONET sites employed in this study. These surface albedos are obtained for Ross-Li
5 based BRDF model, where the BRDF parameters are derived from MODIS climatology. As
6 can be seen, the BRDF based surface albedos significantly deviate from an isotropic
7 Lambertian surface albedo that has no dependence on SZA. Stronger directional dependence
8 for the desert sites than for a site in South Africa can be also noted, which is consistent with a
9 known general feature of soil vs. vegetation surfaces, e.g. (Litvinov et al., 2011; Litvinov et
10 al., 2012; Maignan et al., 2004). In Fig. 5b) we show dependence on SZA of Lambertian to
11 BRDF based albedo ratio for three wavelengths over the solar spectrum. The ratio is equal to
12 unity when the Lambertian albedo is equal to the BRDF based albedo, thus it shows that
13 underestimation (ratio below unity) or overestimation (ratio above unity) of the surface
14 reflectance due to simplified Lambertian model is a function of SZA and wavelength. It
15 therefore emphasizes the importance of the assumption on the surface albedo model for the
16 diurnal dependence and absolute values of the aerosol radiative effect. However, if to
17 consider the whole range of SZAs, the effect on the daily average aerosol effect can be
18 partially canceled because the values below and above unity can be quasi-symmetric. For
19 instance, for the monthly average TOA aerosol direct radiative effect over global land derived
20 from MODIS, Yu et al., (2004) found an uncertainty due to neglecting of the angular
21 dependence of the albedo of only about 5 %. However, influence of the directional properties
22 of the surface albedo is expected to vary depending on the range of SZAs over which the
23 integration is done in order to obtain the daily average forcing, we therefore draw attention to
24 the fact that the magnitude of the uncertainty will be a function of latitude and day of the year.
25 Asymmetry of the ratio around unity in Fig. 5b) is also a function of the wavelength, thus the
26 uncertainty due to Lambertian assumption is depending on spectral extinction of an aerosol
27 model.

28 Figure 6 shows calculations of diurnal aerosol radiative efficiency at top and bottom of
29 atmosphere for Lambertian and BRDF surface reflectance for different type of aerosol and
30 surface. Several observations can be done from this figure. First, diurnal radiative efficiencies
31 can be inter-compared for key aerosol types over different surfaces. It can be observed, for
32 example, that over bright desert surface, biomass burning and mixed aerosol type produce

1 mostly positive instantaneous radiative effect at TOA (Fig. 6c, g, i). Mixture of dust and
2 biomass burning over Sahel type of surface (Fig. 6g) produces positive instantaneous
3 radiative effect when SZA is less than 53° ($\cos(\text{SZA}) > 0.6$). Note that during the biomass
4 burning season in the Sahel region (January - February) the minimal SZA is in range of about
5 16° to 37° . It is also remarkable that relatively weak absorbing dust may still produce positive
6 instantaneous radiative effect at TOA over bright desert (Fig. 6a) if the SZA is less than 45°
7 or 37° ($\cos(\text{SZA}) > 0.7$ or 0.8), while absorbing biomass burning aerosol over southern Africa
8 surface always produce a negative radiative effect (Fig. 6c). With respect to the surface model
9 assumption, Fig. 6 shows an important influence of Lambertian vs. BRDF based albedo on
10 instantaneous radiative effect, which can even change the sign from negative to positive. The
11 results of calculations therefore make evident that the daily average radiative effect for a
12 given location, which is obtained by integration over relevant range of SZAs, can also be
13 significantly affected by assumption on the surface reflectance model.

14 Figure 7 shows the daily average values of aerosol radiative efficiency for the same scenarios
15 as in Fig. 6. The daily average values are calculated here for the daylight fraction of 0.5 and
16 for the minimal SZA of 0° . Similarly to Fig. 6, the daily average aerosol radiative efficiency
17 is presented for the aerosol models as a function of surface brightness. In addition, it evaluates
18 influence of the Lambertian vs. BRDF surface reflectance. For instance, Fig. 7a) shows that
19 the daily average radiative efficiency of biomass burning and both mixed aerosol models
20 switches sign at TOA when surface albedo is brighter than about 0.15 or 0.2 at 550 nm; the
21 daily values of dust and urban aerosol stay negative for the presented range of surface
22 albedos. The ratio of aerosol radiative efficiencies over Lambertian to BRDF based albedo as
23 a function of surface brightness (Fig. 7c) shows percentage of the uncertainty due to the
24 Lambertian surface assumption. When the radiative effect is negative, the ratio below unity
25 means that the daily average cooling effect is underestimated, while when the radiative effect
26 is positive, the ratio above unity means overestimation of the warming effect. At the TOA, the
27 calculated uncertainty ranges up to 30%, depending on aerosol model and surface brightness.
28 It is also evident that the magnitude of the positive radiative effect contribution is depending
29 on minimal SZA. Therefore, as follows from Fig. 6, for low sun elevation (high latitudes or
30 winter season) the Lambertian surface assumption can also cause a systematic overestimation
31 of aerosol cooling in instantaneous and daily radiative effect values. However, if to consider
32 possibly small differences between Lambertian and BRDF based albedos for vegetation
33 surfaces, which are frequent at high latitudes, the effect in this case can be diminished. At the

1 BOA, influence of the surface model is less important, however, is still distinct for the
2 instantaneous values (Fig. 6).

3 **6 Particles sphericity assumption in radiative effect calculation**

4 **6.1 Evaluation of uncertainty**

5 Phase function of spheres is known to differ from one of randomly oriented spheroidal
6 particles that are used for modeling of optical properties of nonspherical aerosol. Since
7 spheres are generally scatter stronger than spheroids at backward scattering angles, it could be
8 expected that the upward hemispherical solar flux is also stronger for spheres. However, this
9 difference is not evident without conducting a rigorous calculation. First of all, not at every
10 scattering angle the directional scattering of spheres is superior of spheroids. For example, for
11 the dust aerosol model, scattering by spheroids is dominating between $\sim 90^\circ$ and $\sim 140^\circ$ (Fig.
12 8a)). Therefore, for low sun elevations, scattering at these angles will contribute stronger to
13 the total upward flux. This also implies that the effect of nonspherical-spherical differences on
14 upward flux depends on SZA. Second, it is known that the phase function is changing
15 spectrally, thus it is possible that the nonspherical-spherical difference is also spectrally
16 dependent and contributes differently over the solar spectrum. Now, supposing that the AOT
17 is known, we would like to evaluate uncertainty in the aerosol radiative effect due to
18 differences in angular redistribution of scattering by volume equivalent spheres and spheroids.
19 The volume equivalence is often used because atmospheric aerosol particles are mainly
20 smaller than the wavelength and in this regime their scattering and absorption properties are
21 primary depend on the volume. However, while using volume equivalent spherical and
22 spheroidal particles, one has to be aware that extinction cross-section is also expected to
23 change. It is because the randomly oriented spheroid has larger geometrical cross section than
24 volume equivalent sphere. In fact, the theorem of Cauchy establishes that the average shadow
25 area of a convex body equals one-quarter the surface area of the body, while the geometry
26 prescribes that the surface area of spheroid is always larger than of volume-equivalent sphere.
27 Thus, the shadow area or the geometrical cross section of spheroids is always larger, which
28 may signify increase of the extinction cross-section as well. In fact, the nonspherical-spherical
29 extinction ratio in Fig. 8b) (black solid line) is generally above the unity. Nevertheless, in a
30 recent work by Kostinski and Mongkolsittisilp [2013] (see section 3, Fig. 4) it is discussed
31 that due to resonances in some size parameter regimes, extinction of spheroids can be smaller
32 than that of volume equivalent spheres. Of course, having realistic particles size distribution

1 instead of a single particle can smooth the effect of resonances, but computations show that
2 the phenomena exist for a realistic size distribution of dust that is employed in this study, i.e.
3 the ratio of extinctions gets below unity for long wavelengths (see black solid line in Fig.
4 8b)). Additionally, even when above the unity, the extinctions ratio is waving spectrally,
5 reflecting different contribution of the resonances as a function of size parameter. More on
6 this subject will be elaborated in further studies (Kostinski et al., personal communications).
7 However, considering that only the phase function assumption is questioned in our work, the
8 effect of different cross sections should be excluded and the AOT kept identical, which
9 appears as not evident when volume and not surface area equivalency is employed. To
10 achieve equality of the AOT in our calculations we attempt to scale the aerosol number
11 concentration in a way that it will give quasi-similar AOT values. Although the identical
12 AOTs can be achieved only at some wavelengths, fitting the AOTs at wavelength of
13 maximum intensity of the solar radiation or at peak of the extinction ratio can minimize the
14 effect of varying cross-section. Dashed black line in Fig. 8b) shows the extinction ratio after
15 the scaling, done in a way that it is equal to unity at the peak of the ratio. In this case the
16 extinction of spheres is only $\sim 1 - 2\%$ larger than of spheroids in the part of the solar spectrum
17 containing most of the energy. Despite of that, the difference becomes large in the spectrum
18 beyond $\sim 2 \mu\text{m}$ and below $\sim 0.3 \mu\text{m}$. At the same time, the gaseous absorption in this spectral
19 region becomes important – the fact minimizes influence of the difference in the AOTs.
20 Increase of the averaged projected area of volume equivalent spheroids also results in a
21 stronger forward peak of the directional scattering (see inset in Fig. 8a)). This indeed
22 contributes to an increase in the asymmetry parameter of the nonspherical relative to spherical
23 particles model (see the asymmetry parameter ratio of nonspherical to spherical model in Fig.
24 8b)). Also, the ratio of the asymmetry parameters is waving spectrally, indicating spectral
25 dependence in nonspherical-spherical difference of the directional scattering; however, it is
26 persistently superior of unity. Lower asymmetry of forward to backward scattering of spheres
27 corresponds to stronger contribution of the backward scattering fraction that hints to stronger
28 cooling effect (backward to space scattering). As for the single scattering albedo (red dashed
29 line in Fig. 8b)), although a small variation appears at short wavelengths of the solar
30 spectrum, it is indeed negligible whatever nonspherical or spherical model of particles shape
31 is used. This result is also in line with previous studies (Dubovik et al., 2006; Mishchenko et
32 al., 1997). It is worthwhile to note, however, that a recent study by Legrand et al. [2014]

1 shows that in the thermal infrared, where absorption constitutes the dominant part of the
2 extinction, the shape of particles has notable effect on the absorption.

3 In order to evaluate uncertainties in aerosol radiative effect due to assumption on spherical
4 particles we calculate instantaneous radiative effect for nonspherical and spherical dust
5 aerosol models. The calculations are conducted using detailed phase function or asymmetry
6 parameter and over different types of the underlying surface. The results show that, while
7 employing the detailed phase function (Fig. 9a, b)), the spherical aerosol model leads to
8 overestimation of cooling at TOA and BOA over dark surfaces; the relative differences in the
9 instantaneous values are ranging between $\sim 1\%$ to 9.5% and depend on the SZA (Fig. 9c, d)).
10 The exact calculations therefore confirm the discussed above hypothesis of overestimation of
11 the cooling effect. At the same time, neglecting of nonsphericity can also provoke
12 overestimation of the warming effect at TOA. This may happen over bright surfaces for high
13 sun elevation when surface reflectance overcomes a critical value with respect to ω_0 (Fraser
14 and Kaufman, 1985) and aerosol radiative effect becomes positive. The calculations show that
15 instantaneous radiative efficiencies at maximal sun elevation can reach overestimation of
16 warming by up to 12% . In the daily average radiative efficiencies, computed assuming
17 maximal sun elevation (SZA = 0°) and daylight fraction of 0.5, overestimation of cooling is
18 however still dominating; the differences are ranging between $2.5 - 6\%$ at TOA and $\sim 6 - 7$
19 $\%$ at BOA (Fig. 10a, b)). Based on the analysis of the differences in instantaneous values, it is
20 evident that differences in the daily average values also depend on the surface brightness; it
21 can be seen that the differences decrease as the surface brightness increase. In addition, the
22 errors are expected be influenced by multiple scattering effects that may smooth the
23 nonspherical-spherical differences in the directional scattering. To evaluate the order of the
24 multiple scattering influence, the differences were calculated for AOT(550 nm) of 0.5 and 2.0
25 (see Fig. 10). It shows that for four times increase in AOT, the error in daily average values
26 decrease by about 15% to 20% at BOA and about 30% to 40% at TOA; the decrease is
27 roughly doubled for outgoing TOA radiation that first was transmitted and then reflected by
28 the atmosphere.

29 It should be mentioned that using thoughtfully the Mie calculation for the non-spherical
30 aerosol retrievals and flux simulations, it is possible to achieve some reduction of the errors
31 due to nonspherical-spherical difference in aerosol scattering, as often expected when

1 spherical aerosol model is used in remote sensing retrievals. Nonetheless, these differences
2 cannot be fully eliminated and remain considerable, as shown in (Derimian et al., 2008).

3 **6.2 Nonspherical-spherical difference over Lambertian versus BRDF surface** 4 **model**

5 Another rather obvious aspect for the analysis is the effect of surface reflectance anisotropy
6 on the manifestation of particle non-sphericity in aerosol radiative effect. The interesting
7 practical question is: how using isotropic Lambertian surface reflectance model affects the
8 accuracy of radiative effect estimation of non-spherical aerosol? In order to answer this
9 question we re-calculated the nonspherical-spherical errors using BRDF surface models. The
10 results show that depending on the SZA the calculated errors are partially reduced or
11 increased. The errors variability also depends on the surface type. However, overall, the
12 differences stay within similar range as if Lambertian surface model is used. The conclusion
13 is valid for the instantaneous (Fig. 11) and, as a consequence, for the daily average values (not
14 shown here).

15

16 **7 Employment of detailed phase function versus asymmetry parameter**

17 A comparison was conducted between calculations of radiative effect using simplified
18 representation of aerosol directional scattering, i.e. accounting only for asymmetry parameter,
19 and using accurate calculations with detailed phase function. In this analysis two main
20 questions were posed. How large the error in calculated radiative effect is if only asymmetry
21 of phase function was accounted for? Also, what kind of uncertainty can be expected for the
22 nonspherical aerosol, if this simplification is used in the calculation of radiative effects? To
23 seek for the answers we compared the calculation using only asymmetry parameter with
24 accurate calculations where the phase function features were accounted using twelve moments
25 expansion of the Legendre polynomial. Figure 12 presents the calculated diurnal radiative
26 efficiencies of dust aerosol model over Lambertian surface using only the asymmetry
27 parameter. From comparison with Fig. 9a, b) showing the same using the detailed phase
28 function, we can notice a significant change in the shape of diurnal dependence of aerosol
29 radiative efficiency at TOA as well as at BOA. That is, the radiative efficiency varies much
30 stronger with SZA in case when the details of the directional scattering are neglected. At the
31 SZA of $\sim 60^\circ$ ($\cos(\text{SZA})$ of 0.4 – 0.5) the cooling effect appears to be systematically

1 overestimated, however, at small SZAs ($\cos(\text{SZA}) \approx 1$) the cooling is underestimated at top
2 and bottom of atmosphere. When the values are positive at top of atmosphere, the warming is
3 overestimated. Figure 12 presents the results for the nonspherical dust aerosol model, but
4 substitution by the asymmetry parameter yields similar effect for all other aerosol models
5 considered in this study. When only the asymmetry parameter is used, it is often expected the
6 most of errors in radiative effect calculations are nearly canceled for daily-integrated values.
7 However, this cancelation happens only if sun is reaching small SZAs. Evidently this is not
8 the case for high latitudes or winter season. Therefore it can be concluded that in daily
9 average values usage of the asymmetry parameter may rather produce an overestimation of
10 the aerosol cooling effect, while the magnitude of this overestimation depends on latitude and
11 season. With the respect to the errors in radiative effect of the nonspherical aerosol, the usage
12 of only the asymmetry parameter yields a significant change in dependence of the error on
13 SZA. Both, at TOA and BOA, the error increase exponentially, reaching a maximum at SZA
14 of 0° (see Fig. 12c, d)). In the daily average values, however, the errors are somewhat lower
15 than in the case of detailed phase function because of compensation of high errors at small
16 SZAs by very low errors at $\text{SZA} > 60^\circ$.

17

18 **8 Illustration of radiative effect calculations over Africa**

19 In this section we illustrate feasibility of rigorous direct aerosol radiative effect calculations
20 on large-scale using satellite observations. It is done as part of the GRASP algorithm
21 application for POLDER/PARASOL observations. The product is of particular interest
22 because it provides detailed aerosol characteristics, including absorption, also over bright
23 surfaces where information about aerosol properties is rarely available. With a goal to test the
24 computational tool and assess an observation-based aerosol radiative effect and its spatial
25 variability, the calculations were conducted for POLDER/PARASOL observations during
26 summer 2008 (June, July, August) over part of Africa known as one of the major sources of
27 the desert dust. It has to be noted, however, that the GRASP algorithm is still in completion
28 phase and that the quality of the aerosol properties retrievals is in a validation process. In this
29 works we therefore present an inter-comparison between AOT and ω_0 retrieved by GRASP
30 from POLDER/PARASOL and those of the conventional AERONET product. The inter-
31 comparison is conducted using four AERONET sites with good statistic of observations and
32 located in the area of interest (Banizoumbou, Agoufou, IER Cinzana and DMN Maine Soroa

1 sites). In order to increase the statistics of joint PARASOL and AERONET observations and
2 to cover various aerosol types and surface reflectance, one year (2008) of data was analyzed.
3 Of course, the inter-comparison at the selected sites is not fully representative for the entire
4 area. Uncertainties can appear for cases of very low AOT, in regions with complex landscape
5 (mountains, mixed land/water pixels) and failures of the cloud mask. Nevertheless, the
6 conducted inter-comparison shows very encouraging correlation coefficients and small
7 uncertainties (RMSE (Root Mean Square Error and Standard Deviation from AERONET))
8 both for AOT and ω_0 (see Fig. 13). The results are obtained for ± 15 minutes temporal
9 matching criteria between PARASOL and AERONET observations and for PARASOL pixels
10 (with about 6x6 km spatial resolution) collocated to each of the selected AERONET station.
11 In addition to comparison with AERONET, analysis of the residuals of the fit for the
12 ensemble of the retrievals employed in this work did not indicated any major problem.

13 Figure 14 presents the means for three months of: i) daily average top and bottom of
14 atmosphere net aerosol radiative effects; ii) radiative efficiencies calculated with respect to
15 AOT at 550 nm (interpolated from nominal wavelength of POLDER); iii) AOT at 565 nm; iv)
16 underlying surface albedo at 565 nm; and v) spectral ω_0 (presented by means of two
17 wavelengths, 443 nm and 1020 nm). The domain averages and standard deviations of the
18 characteristics presented in Fig. 14 are also indicated in the panels. The domain averages and
19 standard deviations are calculated for all observations during three months of summer 2008.
20 As shown in Fig. 14, fine spatial feature of aerosol radiative effect (at top of atmosphere in
21 particular) can be revealed by high spatial resolution of POLDER/PARASOL. A significant
22 amount of pixels, mostly in the northern part of Africa (e.g. central Egypt and northern part of
23 Western Sahara), shows quite strong (up to about 10 to 20 Wm^{-2}) positive radiative effect
24 with the corresponding radiative efficiency over 40 $\text{Wm}^{-2}\tau^{-1}$ (Fig. 14c), d)), despite that the
25 climatological aerosol and surface models in Fig. 7 show positive radiative efficiencies of
26 only up to 20 $\text{Wm}^{-2}\tau^{-1}$. The relatively large positive radiative effect is due to two main factors.
27 First, it happens when the surface reflectance is higher (around 0.4 at 565 nm) and the spectral
28 ω_0 is lower (around 0.8) compared to the limits assumed in calculations presented in Fig. 7.
29 Evidently, the climatological aerosol and surface models represent only an average but cannot
30 be all-inclusive of all possible variations of the properties. Second, what is more important is
31 the non-linearity of the aerosol radiative effect as function of AOT. In fact, the AOT varies
32 significantly in the real data (Fig. 14e)) and strong radiative efficiencies (Fig. 14c)) appear
33 when the AOT is low, while the AOT at 550 nm was set to one in calculations of radiative

1 efficiency presented in Fig. 7. In an attempt to illustrate and evaluate the aforementioned
2 reasons, the aerosol models presented in Sect. 3 have been slightly modified and some
3 supplementary calculations have been conducted. For example, the mixture of dust and
4 biomass burning aerosol model has been assumed to be slightly more absorbing, by changing
5 the spectral imaginary part of refractive indices k at 440/670/870/1020 nm from
6 0.021/0.016/0.013/0.013 to 0.025/0.016/0.016/0.016. This modification produces aerosol
7 properties close to those retrieved for central Egypt with the spectral ω_0 (440/670/870/1020
8 nm) of 0.80/0.81/0.81/0.81. Radiative effect and efficiency calculated for this aerosol model
9 and for corresponding to the central Egypt surface albedo of ~ 0.4 at 550 nm are presented in
10 Fig. 15 (labeled as “Absorbing mixture”). Modification of the climatological dust aerosol
11 model by increasing k (440/670/870/1020 nm) from 0.004/0.002/0.002/0.002 to
12 0.008/0.006/0.006/0.006 produces aerosol properties similar to those retrieved for northern
13 part of Western Sahara with spectral ω_0 of 0.85/0.89/0.91/0.92, for example. Results of
14 calculations for this aerosol model and for corresponding surface albedo of ~ 0.35 at 550 nm
15 are labeled in Fig. 15 as “Absorbing dust”. The radiative effect calculations presented in Fig.
16 15 show first of all that strongly absorbing aerosols over very bright surface produce
17 significant positive radiative effect at top of atmosphere and reproduce range of the radiative
18 effect values obtained over central Egypt and Western Sahara. Second, Fig. 15 illustrates that
19 because of non-linearity of the radiative effect as function of AOT, the values of the radiative
20 efficiency are strongly dependent on AOT with which were calculated. The presented
21 example shows variability in radiative efficiency up to 40% at top and 25% at bottom of
22 atmosphere due to AOT ranging from 0.2 to 1. The fact implies that one should interpret the
23 maps of radiative efficiency in Fig. 14c), d) with caution due to the spatial variation of aerosol
24 concentration.

25 Noteworthy is also the obtained spectral ω_0 (Fig. 14g), h)). Although it is generally consistent
26 with ω_0 of mineral dust (stronger absorption at 443 nm than at 1020 nm), in some cases the ω_0
27 appears quite low (about 0.8) at 443 and 1020 nm, which indicates presence of probably
28 carbonaceous particles or mixed aerosol (e.g. over central Egypt). For the daily average BOA
29 radiative effect (Fig. 14 b)) the values show quite important spatial variability and areas with
30 strong cooling (about -60 Wm^{-2}) that generally correspond to high AOT. Overall, it can be
31 concluded that the values obtained from POLDER/PARASOL observations are in the range
32 of what could be expected from the theoretical climatological calculations presented in this

1 study. The preliminary results and spatial patterns of the aerosol radiative effect thus
2 demonstrate potential of this high advanced product of new GRASP algorithm.

3

4 **9 Conclusions**

5 A rigorous yet fast computational tool for calculations of broadband solar flux and aerosol
6 direct radiative effect was presented. The initial version of the tool developed for using
7 AERONET results and employed in the AERONET operational code was significantly
8 revised and integrated into the GRASP (Generalized Retrieval of Aerosol and Surface
9 Properties) algorithm. Therefore, the GRASP retrieval product can include the estimations of
10 radiative effect for interested users. The tool can also be used in research mode for various
11 types of sensitivity analyses.

12 Using this tool we analyzed sensitivities of the diurnal and daily average shortwave aerosol
13 radiative effects to the details in aerosol and underlying surface characteristics. Overall, the
14 obtained results showed importance of accurate accounting for details in variability of
15 atmospheric aerosol characteristics, such as AOT, ω_0 and g (or phase function) over the solar
16 spectrum in simulations of broadband solar flux and aerosol radiative effect on climate.
17 Diurnal aerosol radiative effect was found as particularly influenced by directional properties
18 of aerosol scattering and by anisotropy of underlying surface reflectance. In fact, not only
19 magnitude, but also dependence on the SZA of instantaneous radiative effect is changing for
20 different aerosol models due to differences in aerosol directional scattering. For example, the
21 changes in the directional scattering due to nonsphericity of particles are notably manifested
22 in the dependence of dust aerosol instantaneous radiative effect on SZA. Neglecting
23 nonsphericity of desert dust in the calculation of radiative effect leads to systematic errors.
24 The computations reveal that simplification of details in directional properties of aerosol
25 scattering and reflectance of underlying surface also cause systematic biases, rather than
26 uncertainties, in evaluation of aerosol radiative effect on climate. Namely, the considered here
27 simplifications are: i) accounting for the asymmetry parameter only instead of detailed phase
28 function; ii) neglecting of phase function features for nonspherical aerosol particles; and iii)
29 directional isotropy of surface reflectance with respect to SZA. We found that using only
30 asymmetry between forward and backward aerosol scattering affects quite significantly the
31 dependence of instantaneous aerosol radiative effect on SZA, relative to usage of the detailed
32 phase function. It tends to overestimate the cooling effect at SZAs around 60° , but

1 underestimate for sun near the zenith. The errors in the daily average values, therefore,
2 depend on latitude and season and minimized for low latitudes and during the summer. If only
3 asymmetry of phase function used, the change in diurnal dependence of instantaneous
4 radiative effect was observed for dust and other aerosol types. Utilization of only the
5 asymmetry parameter also significantly affects evaluation of error in radiative forcing due to
6 neglecting of aerosol nonsphericity; the errors in instantaneous values can vary from few
7 percent to up to ~ 100 %. It should be noted, though, that errors in daily average values are
8 much lower. However, once a detailed phase function is used, the observed error due to
9 neglecting particle non-sphericity is only up to ~ 10 % in instantaneous and daily average
10 aerosol radiative effect. Because of the dependence of this error on the SZA, the biases are
11 expected to vary as a function of latitude and season, having tendency of stronger
12 overestimation of cooling for higher latitudes and wintertime.

13 We emphasize also that a proper inter-comparison of radiative effects of volume equivalent
14 spherical and spheroidal aerosol particles models should account for alteration of geometrical
15 cross section together with directional redistribution of scattering. In our study we apply a
16 partial compensation of the geometrical and as a result of extinction cross-section
17 modification by scaling of concentration. The observed in this study differences between
18 nonspherical and spherical models should be considered rather as a worst-case scenario, but
19 their importance should not be underestimated because they create a notable systematic bias.
20 We also found that using Lambertian approximation instead of BRDF of surface reflectance
21 does not influence significantly the nonspherical – spherical differences, although the diurnal
22 dependence of the error is somewhat modified. The study showed that the nonspherical –
23 spherical difference at top of atmosphere is also pronouncedly depends on the magnitude of
24 surface brightness, while at bottom of atmosphere this dependence practically does not exist.
25 The differences also tend to be reduced with increase in AOT because the multiple scattering
26 effects smooth out differences in the phase functions. It is also important to mention that
27 strong variability of diurnal aerosol radiative effect signify that the minimal SZA and daylight
28 duration can overcome effects of aerosol type and concentration and thus should be taken into
29 account in inter-comparison of daily average aerosol radiative forcing in different time and
30 locations.

31 Finally, application of rigorous aerosol radiative effect calculations was illustrated as feasible
32 on a large-scale using GRASP algorithm for POLDER/PARASOL observations over Africa.

1 Results of the observation-based calculations present quite pronounced range of values and
2 spatial variability of the aerosol radiative effect. The obtained values are generally in line with
3 results of calculations for considered here climatological calculations. The effort presents one
4 more step in the measurement-based estimate of the aerosol direct radiative effect on climate.

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12

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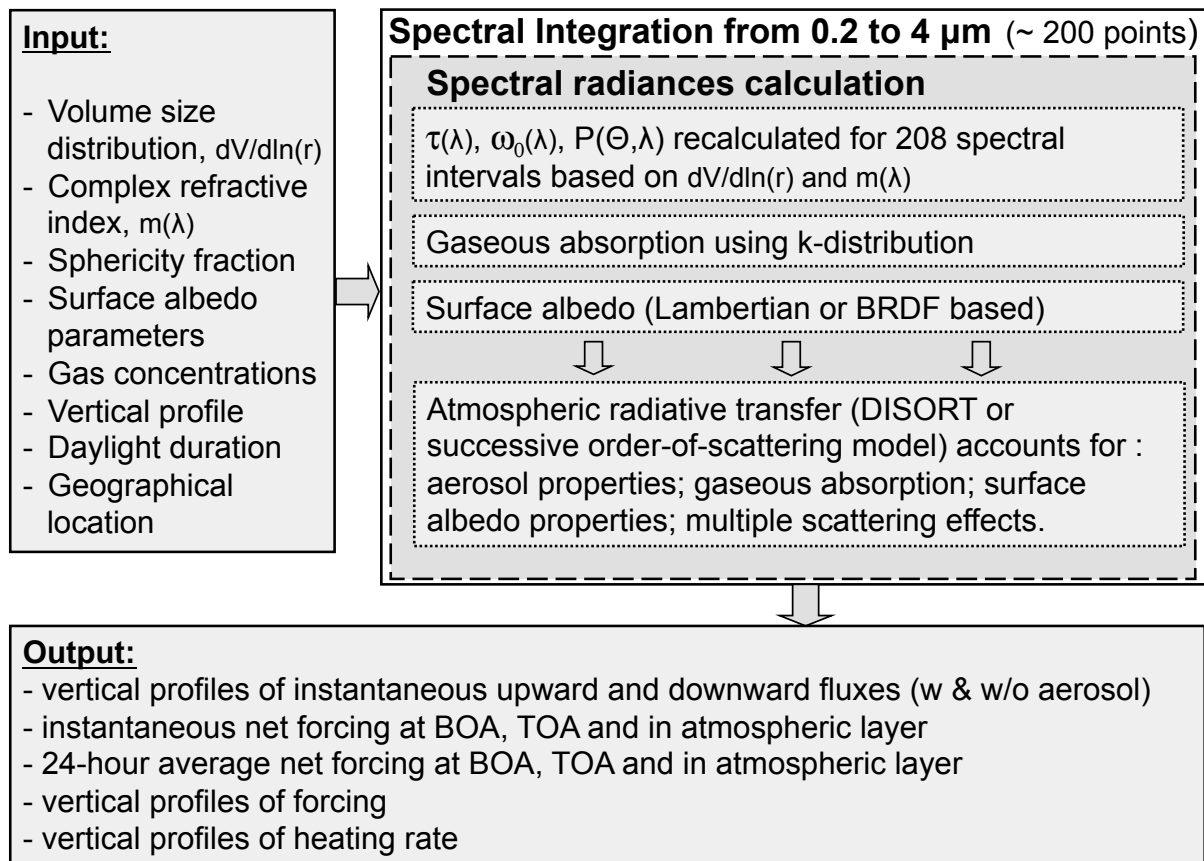
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11

1 Table 1. Complex refractive index for the employed aerosol models.

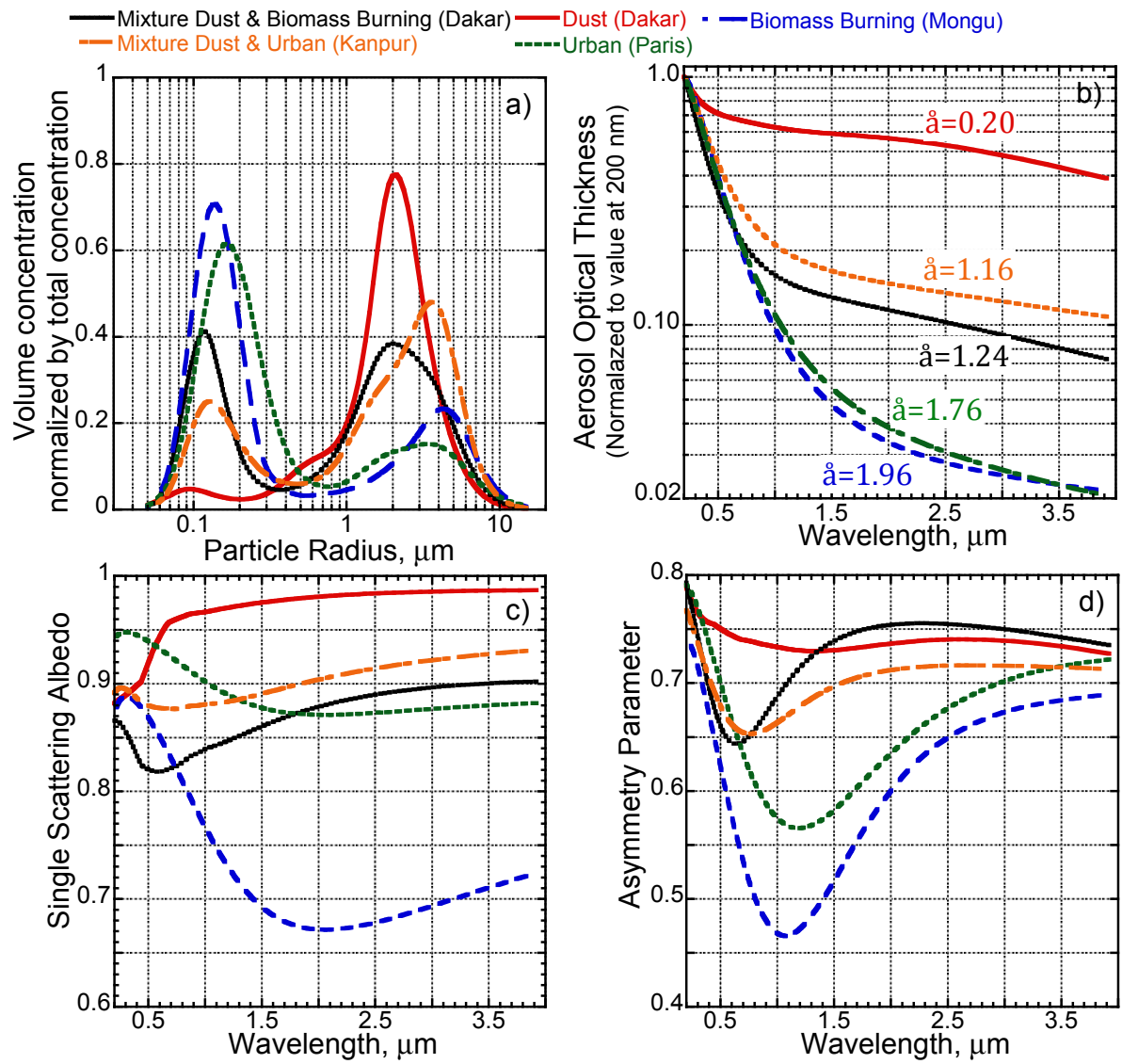
Aerosol model	Complex Refractive Index	
	n	k(440/670/870/1020)
Dust (<i>Dakar, W. Africa</i>)	1.47	0.004/0.002/0.002/0.002
Biomass Burning (<i>Mongu, S. Africa</i>)	1.51	0.023
Urban (<i>Paris</i>)	1.39	0.007
Mixture of Dust & BB (<i>Dakar, W. Africa</i>)	1.45	0.021/0.016/0.013/0.013
Mixture of Dust & Urban (<i>Kanpur, India</i>)	1.50	0.013/0.010/0.009/0.009

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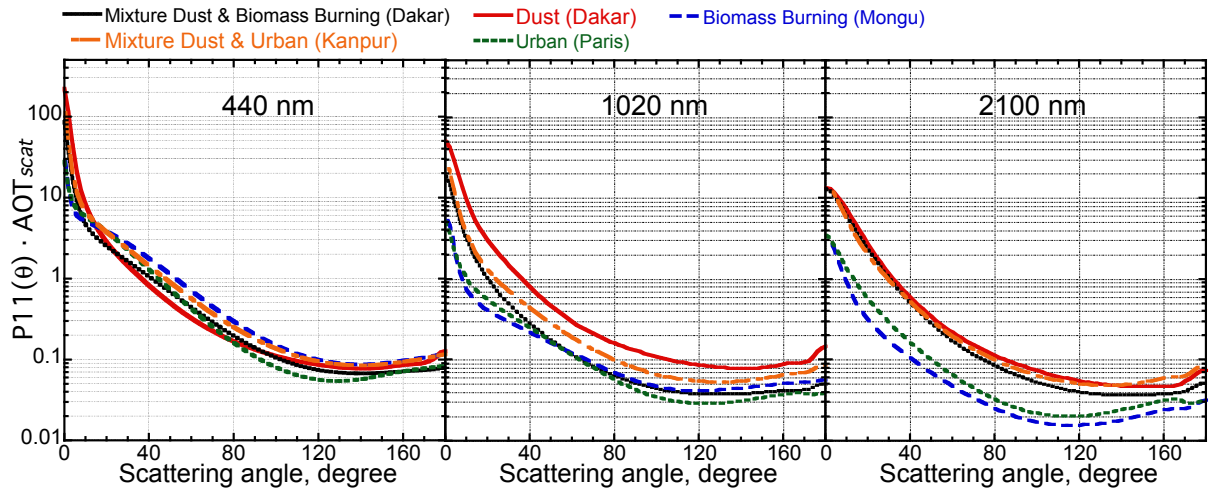
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Figure 1. General organization structure of computational code for broadband solar flux and aerosol radiative effect computations.



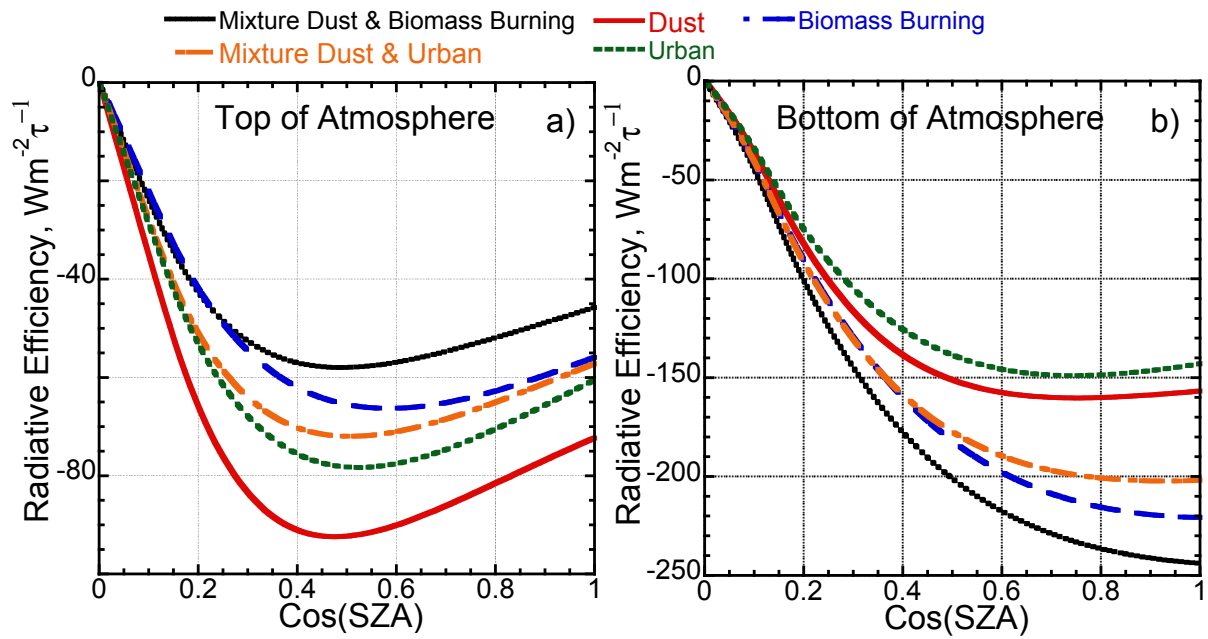
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Figure 2. Characteristics of the employed aerosol models: a) volume size distributions are normalized by total volume concentration; b) spectral aerosol optical thickness normalized by maxima at 200 nm; c) spectral single scattering albedo; d) spectral asymmetry parameter.



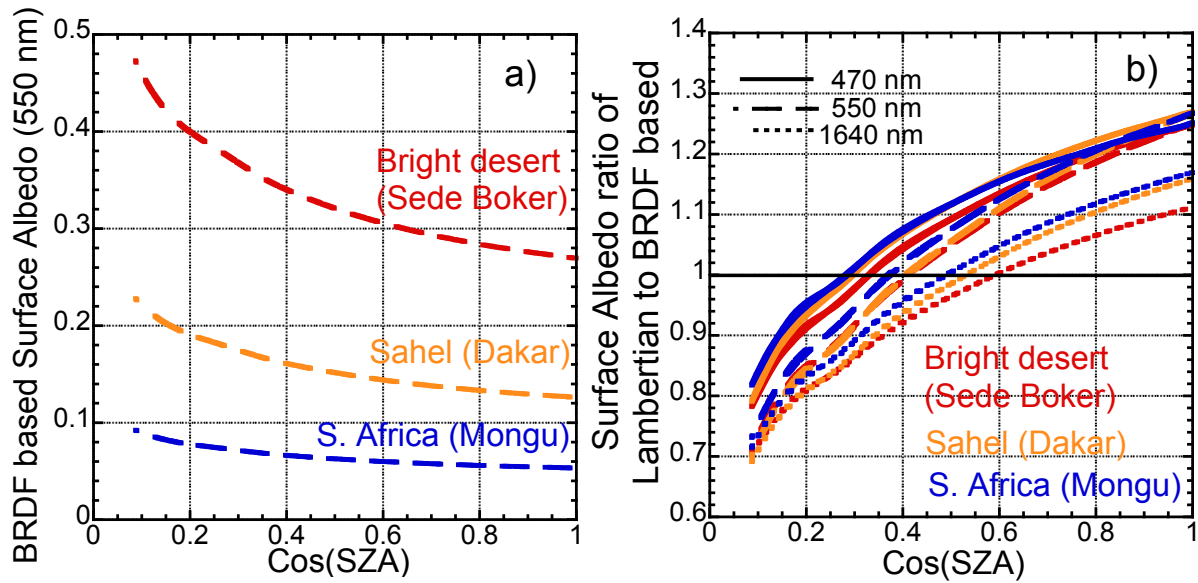
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Figure 3. The calculated directional scattering of the employed aerosol models at 440 nm, 1020 nm and 2100 nm.



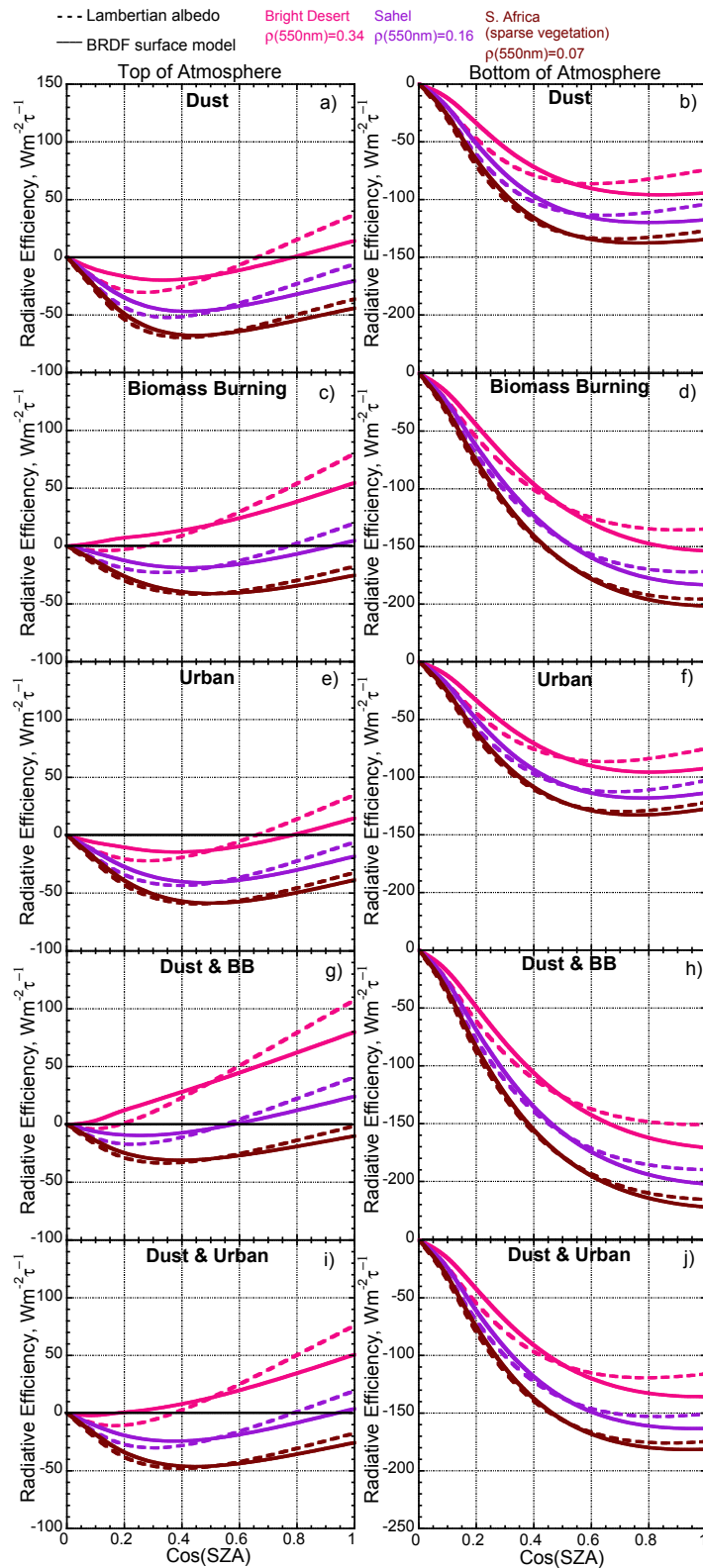
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Figure 4. Instantaneous aerosol radiative efficiencies with respect to 550 nm at a) top of atmosphere and b) bottom of atmosphere calculated over ocean Lambertian surface reflectance.



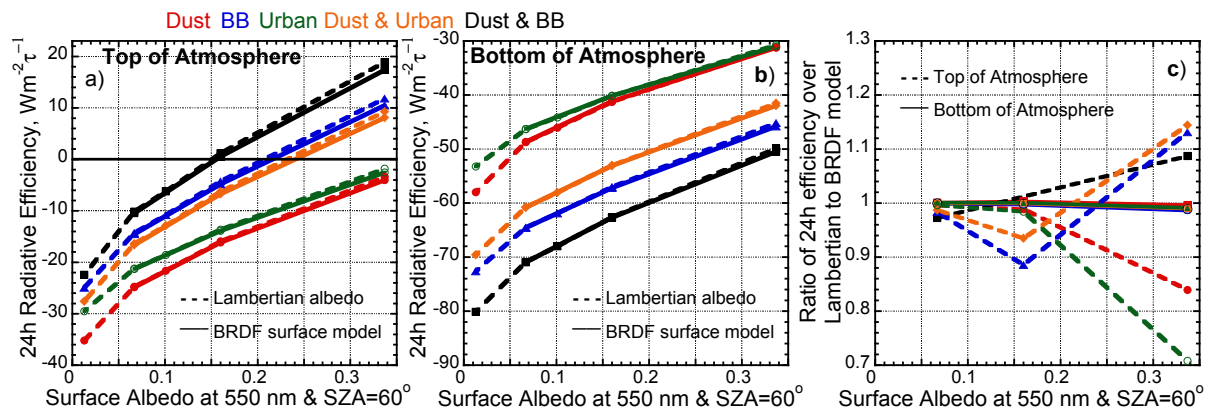
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Figure 5. a) Dependence of the BRDF based surface “black sky albedo” (here presented at 550 nm) on solar zenith angle for three different surface types. b) Ratio of Lambertian surface model to BRDF based surface model “black sky albedos” at three different wavelengths and for three surface types.



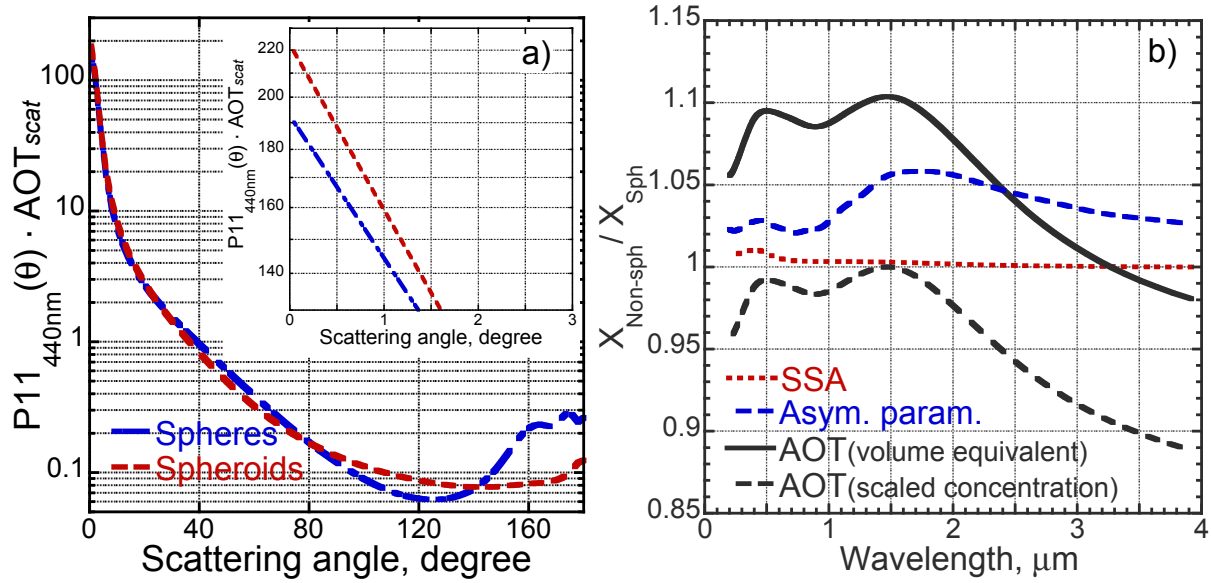
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Figure 6. Instantaneous radiative efficiencies calculated using Lambertian and BRDF surface reflectance calculated for five employed aerosol models and three surface types.



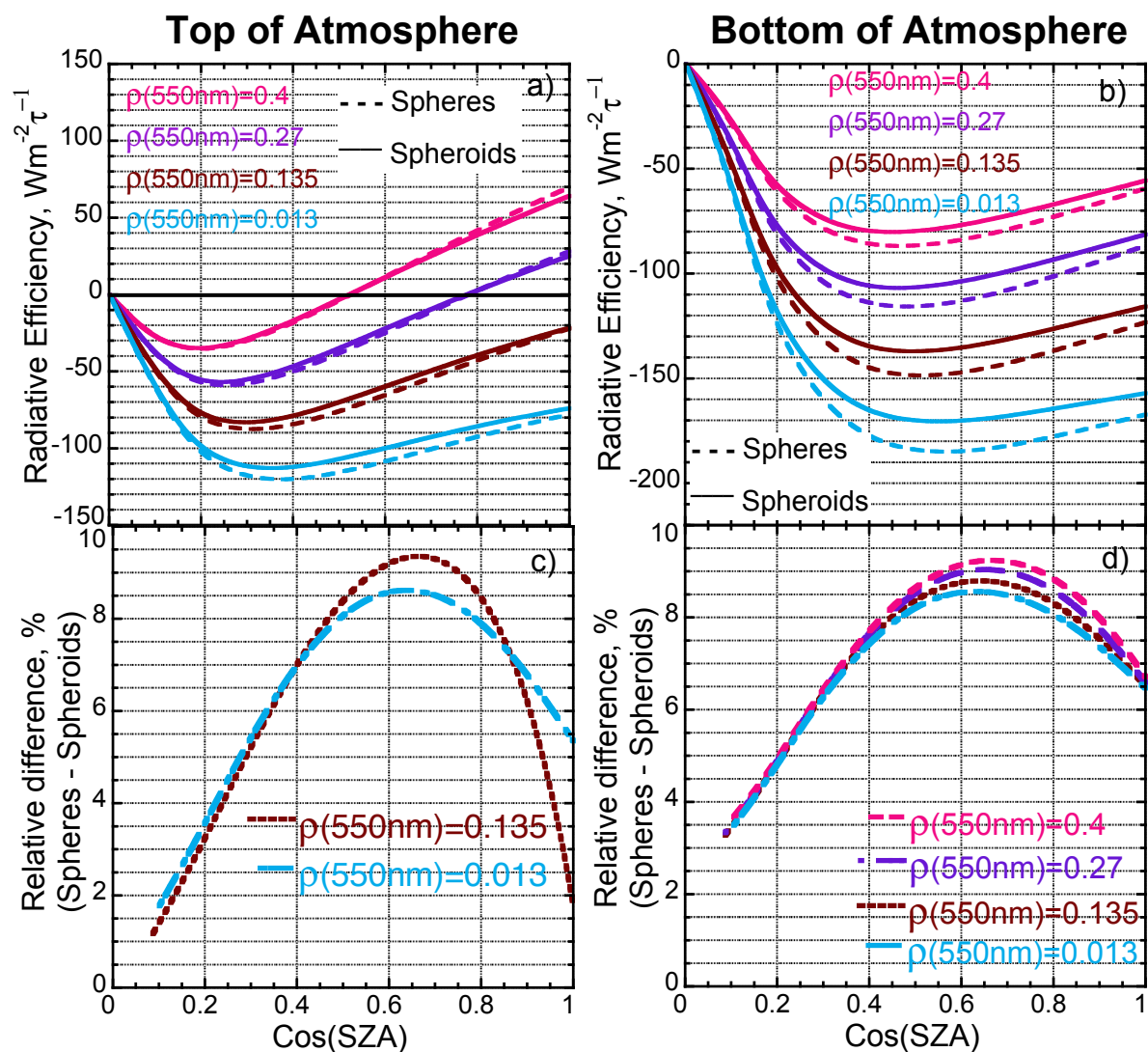
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Figure 7. Daily average aerosol radiative efficiencies at a) top and b) bottom of atmosphere calculated using Lambertian and BRDF surface reflectance. The values are presented as a function of surface albedo at 550 nm and solar zenith angle of 60°. Panel c) presents ratio of daily radiative effects calculated with Lambertian and BRDF surface models. Notes: i) mixture of dust and biomass burning is presented only by two points because small variability of values around zero gives large relative difference for the intermediate surface albedo; ii) for the ocean surface albedo (0.013 at 550 nm) calculations are done for the Lambertian model only.



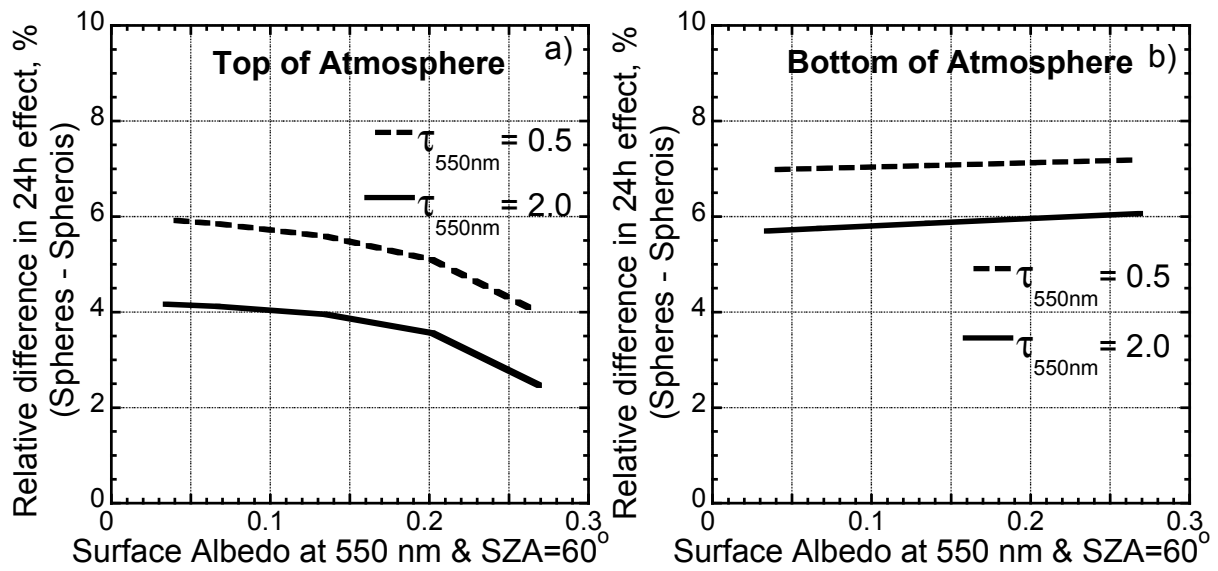
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Figure 8. a) Phase function at 440 nm of dust aerosol model calculated using ensemble of randomly oriented volume-equivalent spheroidal and spherical particles. b) Ratios of Aerosol Optical Thickness, Single Scattering Albedo and Asymmetry parameter calculated using volume-equivalent nonspherical ($X_{\text{non-sph}}$) and spherical (X_{sph}) particles (dashed line AOT - is ratio using spheres with scaled number concentration in a way it gives the same maximal AOT as the spheroid, solid line AOT - ratio without scaling).



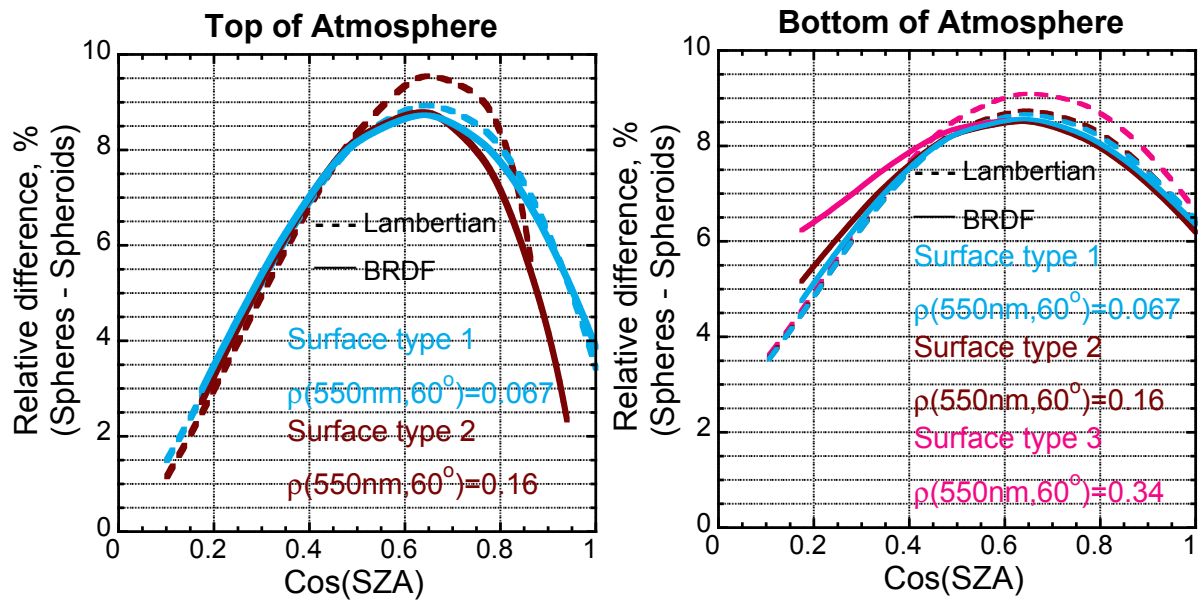
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Figure 9. Nonspherical–spherical differences in radiative efficiencies at top and bottom of atmosphere using detailed phase function of dust aerosol model. Calculations are done for different surface reflectance using Lambertian model. Panels a) and b) present instantaneous radiative efficiencies for nonspherical and spherical cases; c) and d) present relative differences over dark surfaces. The relative difference curves for high surface albedo may have very large values because small uncertainties for near zero radiative efficiencies result into relative differences of ~80-90% (not shown).



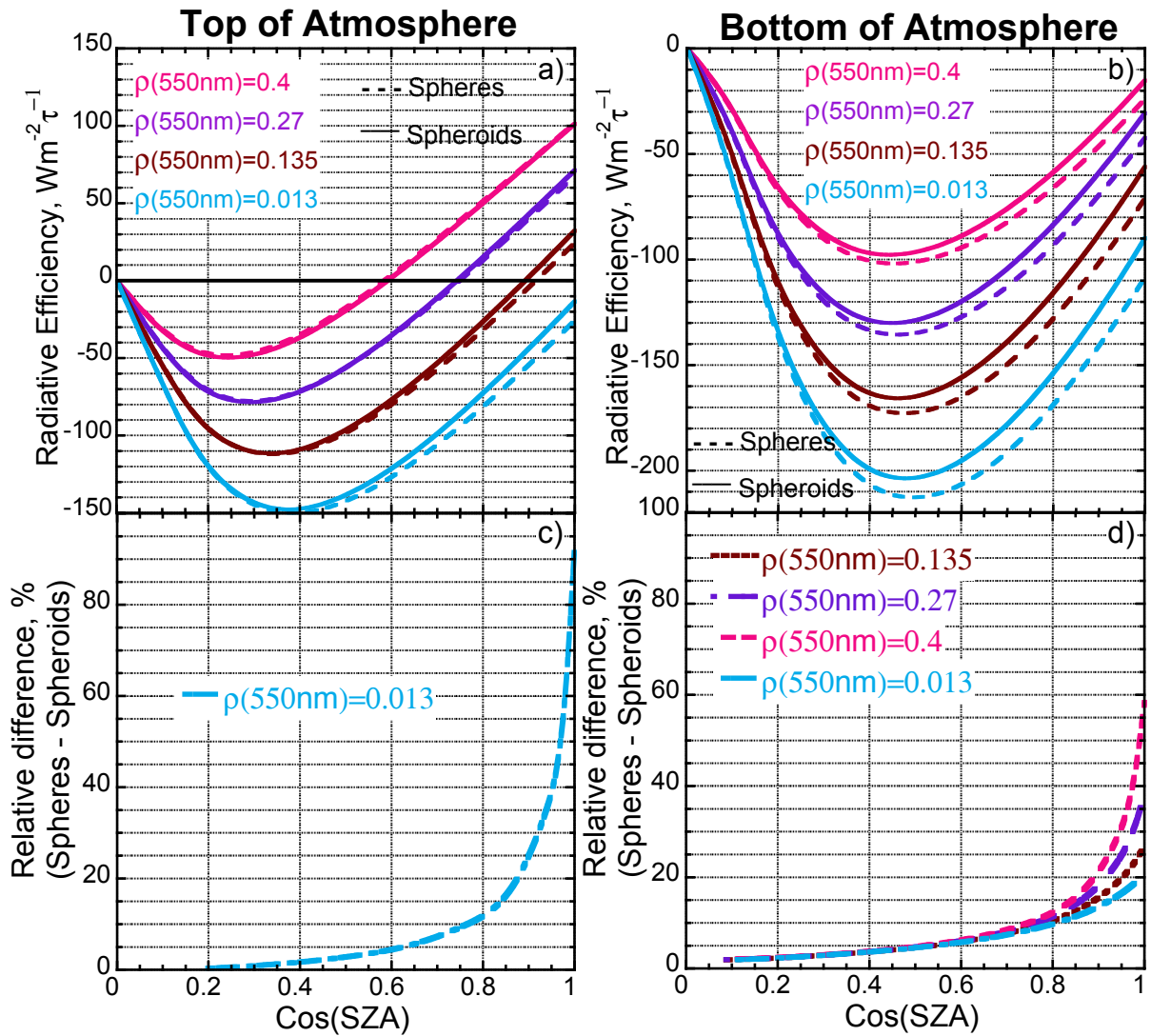
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Figure 10. Relative differences in daily average aerosol radiative effect at a) - top and b) - bottom of atmosphere due to neglecting nonsphericity as a function of surface albedo at 550 nm and solar zenith angle of 60°. The dashed and solid lines correspond to calculations with aerosol optical thickness at 550 nm of 0.5 and 2.0, respectively.



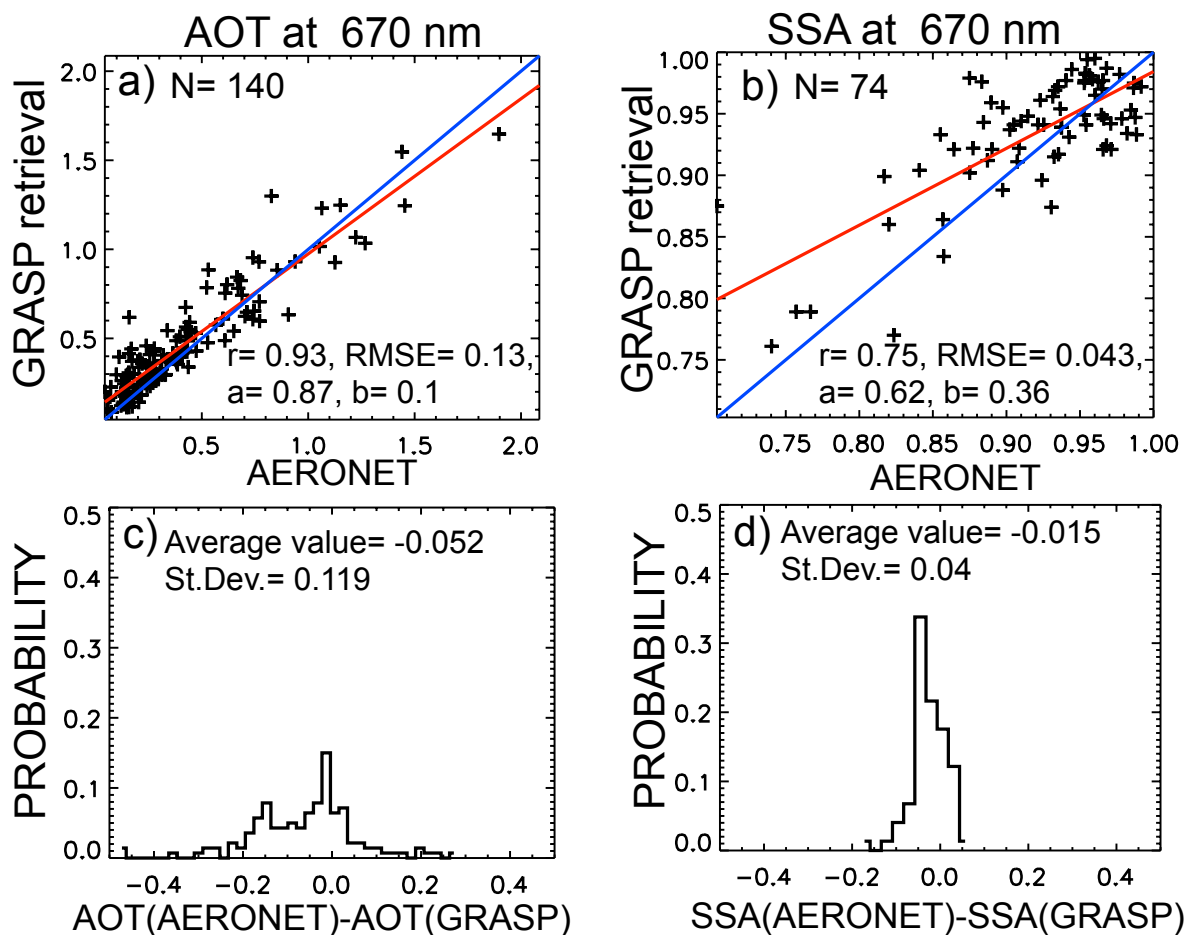
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Figure 11. Relative differences in instantaneous radiative efficiencies due to aerosol sphericity assumption at a) - top and b) - bottom of atmosphere calculated for Lambertian and BRDF surface reflectance models and for different surface types.

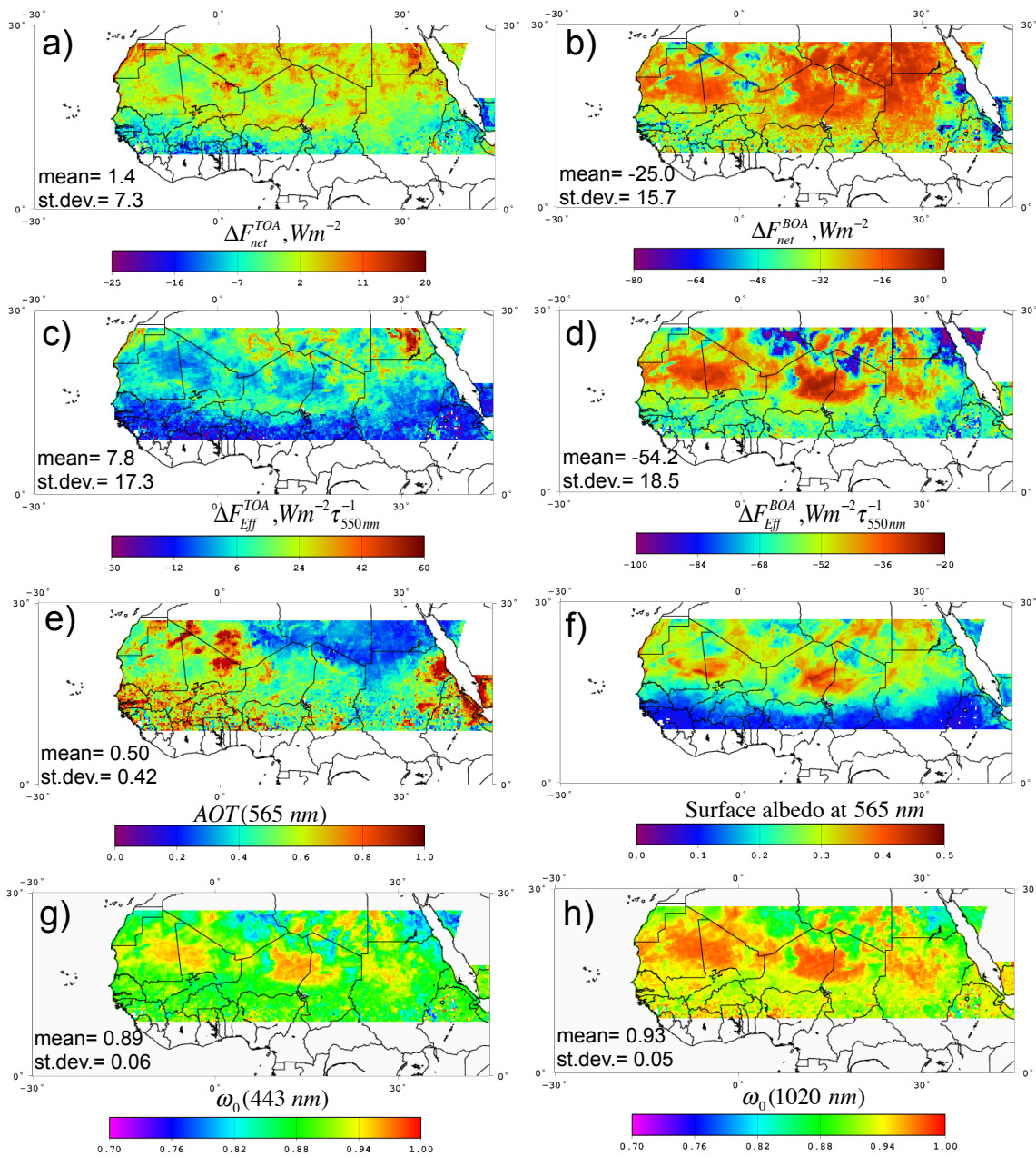


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Figure 12. Same as in Fig. 9, but using calculations of only the asymmetry parameter of the phase function. Note that the relative differences in instantaneous radiative efficiencies at top of atmosphere (panel c) are presented only for the dark surface case. For the high surface albedo the differences appear to have an opposite sign and be large because small uncertainties in the values of radiative efficiencies around zero produce large relative errors (up to ~200%).



1
 2 Figure 13. Inter-comparison between GRASP retrievals applied for POLDER/PARASOL
 3 observations and operational AERONET product during 2008 for ensemble of observations at
 4 four sites (Banizoumbou, Agoufou, IER Cinzana and DMN Maine Soroa). Panels a) and b)
 5 present correlations between AOT and ω_0 at 670 nm, respectively; c) and d) probability
 6 distributions of absolute differences for AOT and ω_0 . The temporal threshold is 15 minutes
 7 between PARASOL and AERONET observation; the products from the ground-based
 8 measurements are compared to those from the space-borne measurements of about 6x6 km
 9 pixel that includes the site.

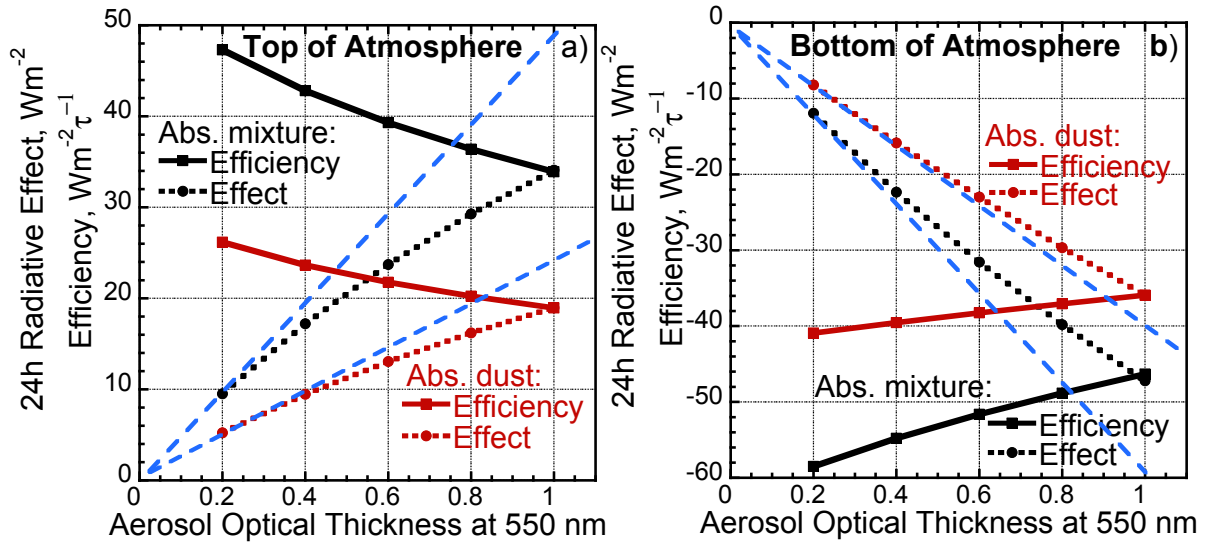


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3 Figure 14. Three months (JJA 2008) means of a) the 24h average Top and b) 24h average
 4 Bottom Of Atmosphere (TOA and BOA) net aerosol radiative effect, c) and d) the
 5 corresponding radiative efficiencies (see Sect. 8 for the interpretation), e) AOT at 565 nm, f)
 6 underlying surface albedo at 565 nm, and g) ω_0 at 443 nm and h) at 1020 nm as retrieved and
 7 calculated by GRASP algorithm applied for POLDER/PARASOL observations. The panels
 8 also include the domain averages and corresponding standard deviations.

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4 Figure 15. (dashed lines) Dependence between calculated 24h average aerosol radiative effect
5 and AOT at 550 nm; (solid lines) 24h average aerosol radiative efficiency calculated using
6 presented on the abscissa AOT. Black and red lines correspond respectively to “Absorbing
7 mixture” and “Absorbing dust” aerosol models described in Sect. 8; surface albedo at 550 nm
8 is set to 0.43 for “Absorbing mixture” and 0.34 for “Absorbing dust” scenarios; blue lines
9 represent linear dependence between 24h average aerosol radiative effect and AOT. Panel a)
10 is for top and b) for bottom of atmosphere.