

1 **Phase Transition Observations and Discrimination of Small**  
2 **Cloud Particles by Light Polarisation in Expansion**  
3 **Chamber Experiments**

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### 13 **Abstract**

14 Cloud microphysical processes involving the ice phase in tropospheric clouds are among the  
15 major uncertainties in cloud formation, weather and General Circulation Models (GCMs).  
16 The detection of aerosol particles, liquid droplets, and ice crystals, especially in the small  
17 cloud-particle size range below 50  $\mu\text{m}$ , remains challenging in mixed phase, often unstable  
18 environments. The Cloud Aerosol Spectrometer with Polarisation (CASPOL) is an airborne  
19 instrument that has the ability to detect such small cloud particles and measure the variability  
20 in polarisation state of their backscattered light. Here we operate the versatile Cosmics-  
21 Leaving-OUtdoor-Droplets (CLOUD) chamber facility at the European Organisation for  
22 Nuclear Research (CERN) to produce controlled mixed phase and other clouds by adiabatic  
23 expansions in an ultraclean environment, and use the CASPOL to discriminate between  
24 different aerosols, water and ice particles. In this paper, optical property measurements of  
25 mixed phase clouds and viscous Secondary Organic Aerosol (SOA) are presented. We report  
26 observations of significant liquid - viscous SOA particle polarisation transitions under dry  
27 conditions using CASPOL. Cluster analysis techniques were subsequently used to classify  
28 different types of particles according to their polarisation ratios during phase transition. A  
29 classification map is presented for water droplets, organic aerosol (e.g., SOA and oxalic  
30 acid), crystalline substances such as ammonium sulphate, and volcanic ash. Finally, we  
31 discuss the benefits and limitations of this classification approach for atmospherically

1 relevant concentration and mixtures with respect to the CLOUD 8 - 9 campaigns and its  
2 potential contribution to Tropical Troposphere Layer (TTL) analysis.

### 3 **1 Introduction**

4 Scattering and absorption due to atmospheric particles can vary widely, leading to net  
5 radiative effect that either cool or warm the surface of the Earth. Ice crystals pose a potential  
6 challenge since their non-sphericity complicates the theoretical description of their single  
7 scattering properties (Macke et al., 1996). Several attempts have been made to model and  
8 simulate light interactions with different ice crystal habits, mixtures of crystal types,  
9 aggregates, and aerosols (Baran, 2013), but no single method can easily combine all size  
10 ranges and types of particles, making accurate, unified modelling nearly impossible.  
11 Scattering analysis is complicated further in small ice crystals and Secondary Organic  
12 Aerosol (SOA).

13 Ice crystals can have different major internal defects (e.g., stacking faults, chemical defects,  
14 molecular vacancies, interstitial molecules, ionized states, and orientation defects), surface  
15 roughness, and branching with various symmetries. These could be even more influential in  
16 small ice measurements. The optical effects of these defects depend strongly on the spatial  
17 orientation of the particle. They can lead to systematic biases since particles with a high  
18 width to height aspect ratio can have a preferred orientation in chamber measurements  
19 (Abdelmonem et al., 2011). However, single particle-by-particle analysis of the backscatter  
20 polarisation state is useful for particle discrimination as we shall show.

21 Aerosol particles found in the lower confines of the atmosphere are typically internal or  
22 external mixtures of inorganic salts, refractory components such as mineral dusts and clays,  
23 and organic species; they also contain varying quantities of water. The hygroscopicity of  
24 organic particle is derived from their composition (Cziczo et al., 2004; Jimenez et al., 2009;  
25 Duplissy et al., 2011). Pure sulphate and internally mixed organic/sulphate aerosols will have  
26 different water uptakes, and consequently different refractive indices. This may lead to mis-  
27 sizing by optical instruments if the composition is not taken into account (see Sect. 2.3). In  
28 addition to the familiar liquid and crystalline states, atmospheric aerosol may also exist in  
29 semi-solid and solid amorphous states (i.e., lacking an ordered, repeating structure) such as  
30 soft polymers, gels, or glasses (Mikhailov et al., 2009). The viscous SOA is expected to  
31 appear either in low relative humidity (RH), low temperature environments or both. A subset  
32 of these viscous particles, sometimes referred to as “glassy”, (e.g. Koop et al., 2011) or  
33 amorphous, are thought to be important components in the atmosphere because of their low

1 volatility, long lifetimes, and their potential impact on several competing processes which  
2 occur during updraft of an air parcel. These include: heterogeneous ice nucleation in the  
3 deposition mode onto the glassy solid aerosol surface; diffusion of water into the particle,  
4 inducing a gradual phase transition towards the liquid state; and immersion freezing during  
5 the transition between the states, (Berkemeier et al., 2014). Some terpenoids can affect these  
6 processes by formation of particles in the glassy state.

7 In this study we examine alpha-pinene, the most widely encountered terpenoid in nature  
8 (Noma and Asakawa, 2010); its derivatives in SOA are optically active materials (Wilberg et  
9 al., 2004; Cataldo et al., 2010) that induce a change in the polarisation state of the scattered  
10 radiation. The resulting change to the polarisation state of the back-scattered light from these  
11 aerosol particles can, therefore, be used to probe these effects. Small molecules such as water  
12 can soften the structural matrix (as water acts as a plasticizer) of SOA, thus reducing  
13 viscosity. As water molecules are removed by drying, the SOA viscosity increases. These  
14 highly viscous particles (Renbaum-Wolff et al., 2013) are, therefore, likely to be optically-  
15 anisotropic (having aspherical shape, branches, roughness, or variations in internal structure)  
16 that accentuate the polarisation shift of the incident beam in Cloud Aerosol Spectrometer  
17 with Polarisation (CASPOL). We probe this viscous state in this paper.

18 **1.1 The CLOUD Chamber and Cloudy Experiments**

19 The Cosmics Leaving OUtdoor Droplets (CLOUD) chamber was designed to simulate  
20 different atmospheric conditions to reduce the uncertainties for cloud, weather and general  
21 circulation models (Chapter 7 of IPCC 2013; Boucher et al., 2013) and provide new data for  
22 the parametrisation and modelling of atmospheric processes. The first series of CLOUD  
23 experiments at the European Organisation for Nuclear Research (CERN) began in 2006  
24 (Duplissy et al., 2010). For several years these experiments were mainly dedicated to aerosol  
25 nucleation and growth. Recently, additional experiments focusing on cloud formation have  
26 been performed at the CERN chamber. In this paper we focus particularly on results from this  
27 “Cloudy” series of experiments. This addition was driven by the importance of ice particles  
28 to the earth’s radiation budget and feedback mechanisms. The CLOUD chamber utilises the  
29 adiabatic expansion principle to generate super-cooled water and ice clouds, similar to other  
30 atmospheric cloud chambers (Möhler et al., 2006; Schnaiter, 2009; Tajiri et al., 2013). In  
31 order to investigate the microphysics of homogeneous ice nucleation in situ, two sets of  
32 Cloudy experiments were conducted during two campaigns at CERN in 2013 and 2014,  
33 hereafter referred to as CLOUD 8 and 9, respectively. In this paper we highlight results from

1 some of the mixed-phase cloud measurements as well as new polarisation transition  
2 measurements for SOA from the photo-oxidation and ozonolysis of alpha-pinene.

## 3 **2 Methodology**

### 4 **2.1 CLOUD chamber and instrumentation**

5 The CLOUD chamber was designed in order to achieve excellent temperature stability and  
6 very low background aerosol and trace gas concentration levels in order to identify small  
7 changes in nucleation rates due to the influence of cosmic rays (Duplissy et al., 2010; Kirkby  
8 et al., 2011). An overview of the chamber and more detailed information is presented in the  
9 Supplement (Fig. S1). The chamber was equipped with a range of instruments that can  
10 measure atmospheric constituents. Aerosol concentrations were measured by a combination  
11 of Scanning Mobility Particle Sizer (SMPS with TSI-type, custom-built DMA), and an Ultra  
12 High Sensitivity Aerosol Spectrometer (UHSAS, DMT) to determine potential Cloud  
13 Condensation Nuclei (CCN) concentrations. During CLOUD 8 and 9, instruments for the  
14 measurements of cloud droplets and ice particles were added. Cloud particle formation and  
15 evolution was measured using several optical spectrometers including a WELAS optical  
16 particle counter (WELAS Promo 2000, Palas GmbH) (Benz et al., 2005), a Particle Phase  
17 Discriminator (PPD-2K), (Kaye et al., 2008), a 3-View Cloud Particle Imager (3V-CPI,  
18 SPEC Inc.), (Lawson et al., 2003) and a Cloud Aerosol Spectrometer – with Polarisation  
19 (CASPOL, DMT) (Baumgardner et al., 2001, 2011; Glen and Brooks, 2013). The latter will  
20 be described in more detail in Sect. 2.3. An averaged path system, the Scattering Intensity  
21 Measurements for the Optical detection of ice – SIMONE-Junior (Schnaiter et al., 2012), was  
22 installed for bulk depolarisation measurements.

23 The procedure for operation of the CLOUD facility as an expansion cloud chamber for ice  
24 nucleation studies along with full schematics are described in detail by Guida et al. (2012,  
25 2013) and will only briefly be reviewed here. Controlled supersaturated conditions are  
26 created in the chamber by allowing air to expand and cool at prescribed rates. The basic  
27 operating procedure adopted for all the cloud microphysics experiments was as follows: the  
28 chamber was slowly pressurised to +220 mb relative to ambient pressure; CCN were then  
29 vaporised and injected through the gas lines; after the CCN had mixed throughout the  
30 chamber volume, a valve was opened allowing the air to expand with the pressure reached +5  
31 mb. The pressure, temperature, and humidity traces for a typical expansion are shown in Fig.  
32 1. Super-saturation occurs due to the pressure reduction and resultant temperature decrease.

1 Before the beginning of the expansion, RH with respect to liquid water of approximately 92–  
2 96 % was achieved. The total humidity in the chamber was measured by dew point mirror  
3 instruments (model MBW973 during CLOUD 9 and MBW373LX during CLOUD 8, both  
4 from “MBW calibration Ltd.”) attached to a heated sampling line. Together with the in situ  
5 measured gas temperature (6 calibrated thermocouples, type K) these instruments provide the  
6 RH in the chamber and might overestimate it in the presence of clouds (assuming additional  
7 evaporation of cloud droplets in the heated sampling line). During CLOUD 9 a tuneable  
8 diode laser (TDL) hygrometer, comparable to the APicT instrument as described by Fahey et  
9 al. (2014), was used to measure the water vapour content with 1 Hz time resolution using a  
10 single optical path of 314 cm once across the middle plane of the CLOUD chamber. Thus,  
11 this instrument provides the RH also in the presence of clouds. Subtracting the water vapour  
12 content from the total water content results in the condensed (ice or liquid) water content.  
13 Sulphuric acid, ammonium sulphate, and oxalic acid particles were used to seed the chamber  
14 with CCN concentrations ranging from 0.5 to several thousand particles  $\text{cm}^{-3}$ . The CCN  
15 number concentrations determined the cloud droplet size, with higher CCN concentration  
16 producing higher concentrations of smaller ice particles in CLOUD.

17 Although the expansion is, ideally, adiabatic, heat is continuously transferred to the cooled air  
18 from the chamber walls, as the temperature control system is maintained at the pre-expansion  
19 temperature, resulting in eventual evaporation of the cloud. The cloud lifetime in the CLOUD  
20 experiments could be controlled (e.g., by fan speed or by number of steps in the expansion  
21 profile) from several minutes to greater than forty minutes when required (e.g., for ice  
22 evolution experiments).

## 23 **2.2 Cloud Experiment Overview**

24 A series of experiments was conducted to generate liquid clouds (Hoyle et al., 2015), mixed  
25 phase clouds, and pure ice clouds. Controlling stepwise the rate of expansion and the  
26 humidity flow into the chamber in the mixed phase experiments, it was possible to obtain  
27 water super saturation followed by ice super saturation, allowing CCN activation to form a  
28 cloud for a short period of several minutes. The adiabatic expansion experiments on which  
29 this paper focuses are summarised in Table 1, but results based on a much broader data base  
30 of several hundred CLOUD expansions will also be considered for the discussions.

31 Several additional experiments were conducted to examine any aerosol polarisation state  
32 changes arising from possible viscosity changes in response to RH variations, using

1 CASPOL. A more detailed description of these experiments can be found in the  
2 accompanying paper by Järvinen et al. (2015).

### 3 **2.3 The CASPOL Instrument**

4 The Cloud and Aerosol Spectrometer with Polarisation detection (CASPOL) is part of the  
5 Cloud Aerosol and Precipitation Spectrometer (CAPS). The first variant of the instrument  
6 was introduced in 1999 and was designed for airborne in situ cloud measurements  
7 (Baumgardner et al., 2001; Heymsfield, 2007), although it has subsequently been used for  
8 cloud chamber measurements (Krämer 2009). The version of CASPOL employed here has a  
9 linearly polarised laser to provide a collimated incident beam of light at a wavelength of 680  
10 nm.

11 The first two detectors of the instrument detect the light scattered in the forward direction.  
12 The collection of the light cone is subtended by the angles 4 to 12°. The near forward angles  
13 are used for sizing because light is preferentially scattered in the forward direction from  
14 particles whose diameters are larger than the incident wavelength. The first detector in the  
15 forward direction is used as a qualifier; it has a rectangular optical mask that restricts  
16 scattered light from particles that are outside the centre of focus of the laser beam. Only  
17 particles within the optimal view volume are counted and characterized. All data are collected  
18 on a single particle basis, thus provide a measure of particle-by-particle variability and single  
19 particle optical properties. The particle's water equivalent optical diameter in the range 0.51 –  
20 50 µm is determined from the forward scattering signal in the second detector using the  
21 standard Mie scattering assumptions, i.e., spherical geometry and isotropic refractive index.

22 The next pair of detectors measures the backscattered light with collection angles of 168 to  
23 176°. The first backscattering detector is used for qualitative particle shape discrimination.  
24 The second detector has a polarised filter (90° to the polarisation of the incident light) to  
25 measure the change in polarisation of scattered light caused by asphericity (Baumgardner et  
26 al., 2011; Glen and Brooks, 2013) or birefringence. In this configuration, spherical particles  
27 produce little response in the perpendicular polarisation backscatter detector. Conversely,  
28 frozen water droplets and aspherical ice crystals will show much more distinct signals.

29 In order to eliminate aerosol particle interference in our cloud measurements, only  
30 contributions from a subset of larger particles above 3 µm were included. This threshold is  
31 based on work by Baumgardner et al. (2001) and Lance (2012) who selected a similar size

1 range for cloud particle measurements. In the special case of SOA measurements, a subset of  
2 small particles ( $< 3 \mu\text{m}$ ) detected in the lower gain stage, was considered.

### 3 **Calibration**

4 The CASPOL was calibrated using Polystyrene Latex Spheres (PSL), as described elsewhere,  
5 e.g., DMT Manual (2011), Meyer (2011), Rosenberg et al. (2012). Size calibration relates the  
6 amplitude of the instrument's response to particle scattering cross-sections. Using the Mie–  
7 Lorenz curve, the nominal size bin limits can then be defined (Table S2 in the Supplement) in  
8 terms of the diameter of water droplets having the same scattering cross-section, giving a  
9 reasonable estimate of particle size for liquid droplets and small spherical ice particles.  
10 Aspherical particles will be mis-sized with respect to spherical particles, subject to their  
11 cross-section as shown by Borrmann et al. (2000). In our instrument this error would  
12 normally be in the order of the size bin width. The uncertainty in the derived polarisation  
13 ratio is approximately 20% as described by Baumgardner et al., (2005).

## 14 **2.4 Data processing**

### 15 **2.4.1 Particle-by-Particle analysis**

16 The polarisation ratio measured with the CASPOL instrument and reported in this paper is  
17 defined as the ratio of perpendicularly polarised backscatter intensity to total backscatter  
18 intensity and provides a measure of the combined phase, composition, and surface features of  
19 the particle. This ratio differs from the depolarisation ratio that is measured using remote  
20 sensing techniques (Groß et al., 2013). The two ratios cannot be directly compared, requiring  
21 additional calibration for this purpose (Meyer, 2011). The ratio of perpendicularly polarised  
22 backscatter to forward scatter ( $D_{pol}/Fwd$ ) indicates the contribution of particle size to the  
23 scattering. Particle by particle (PBP) measurements reveal the fraction of aspherical particles  
24 population (Fig. 2c) and its evolution. Here we employ cluster analysis on PBP data (Sect.  
25 2.4.2) for phase discrimination and for data quality assurance. This method can also be used  
26 to classify highly polarising particles. Corrections to the forward, backward and the  $D_{pol}$   
27 channels have been applied and summarized in Table S3 in the Supplement.

### 28 **2.4.2 Cluster analysis**

29 Clustering or grouping of data by the similarity in one variable or a matrix of variables  
30 reveals the size of the population with similar properties and the number of the unique groups  
31 in the dataset, as well as the spread in each group. Clustering analysis is used here to  
32 discriminate and assign unique particle properties corresponding to different phases during



1 the experiment (e.g., water, ice), primarily based on variations in the polarisation state of the  
 2 scattered light (Fig. 3). Clustering approaches have been previously used for aerosol property  
 3 classification, e.g., Omar et al. (2005), Robinson et al. (2013), Crawford et al. (2015). Here  
 4 we use the  $k$ -means cluster function (Seber, 1984; Spath, 1985) from the MATLAB<sup>®</sup> statistics  
 5 toolbox. The algorithm then calculates the minimum total intra-cluster variance (Eq. 1)

$$\sum_{i=1}^K \sum_{x_j \in S_i} d(x_j, \mu_i) \quad (1)$$

6 where  $S_i$  is the  $i$ th cluster ( $i = 1, \dots, K$ ),  $\mu_i$  is the  $i$ th centroid of all the points  $x_j$  in cluster  $S_i$ ,  
 7 and  $d$  is the distance function (e.g., squared Euclidean). In this case the function is applied to  
 8 a matrix of parameter vectors including polarisation, size, asphericity, concentration, inter-  
 9 arrival-time, time, etc. This approach should, by itself, be sufficient for discriminating a  
 10 simple mixture consisting of two discrete and well-separated phases as may be found in the  
 11 water-ice particle population. In our aerosol-cloud nucleation experiments, an a-priori  
 12 assumption of cluster number is challenging due to the variability of particles. Initial  
 13 estimates of cluster numbers (1–7) were tested in sequential iterations. A silhouette index,  
 14  $s(i)$ , was then used to quantitatively assess the quality of clustering. This is a composite  
 15 index reflecting the compactness and separation of clusters; a larger average silhouette index  
 16 indicates a better overall quality of the clustering result (Chen et al., 2002). The silhouette  
 17 value of a point is a measure of the similarity of points within a given cluster compared to  
 18 these in other clusters; it is defined as

$$s(i) = \frac{b(i) - a(i)}{\max(a(i), b(i))} \quad (2)$$

19 where  $a(i)$  is the average distance of the point  $i$  to the other points in its own cluster  $A$ .  $b(i)$   
 20 is the minimal average distance of the point  $i$  to the points in the other cluster, over all  
 21 clusters other than  $A$  (Eq. 2). For the best possible fit, the silhouette index is,  $s(i) = 1$ . This  
 22 validation is sufficient for our analysis to indicate the ability of the algorithm to group similar  
 23 data sets using the prescribed values. Following cluster analysis, asphericity thresholds are  
 24 selected based on cluster boundaries identified by the colour transition in Fig. 3 and  
 25 silhouette values greater than 0.9.

## 1 **3 Results**

### 2 **3.1 CASPOL Water-Ice measurements**

3 As the temperature in the chamber decreases in the multistep expansions, liquid cloud starts  
4 to form when the RH exceeds water saturation (Fig. 1). Figure 2a shows the formation of a  
5 mixed phase cloud as a function of time. Droplets formed at sub-zero temperatures are super-  
6 cooled and some of them freeze. During the stabilisation period, when pressure remains  
7 constant, some of the super-cooled droplets evaporate as the walls reheat the chamber.  
8 During the second step of the expansion, the ice grows further. The rapid growth of ice  
9 particles depletes the available water vapour, causing the remaining liquid droplets to  
10 evaporate by the Bergeron–Findeisen mechanism. The aspherical fraction (Fig. 2b), and the  
11 concentrations of water and ice (Fig. 2c) were calculated from the PBP cluster analysis for  
12 each of these conditions during the run. Images of some typical ice particles (diameter < 150  
13  $\mu\text{m}$ ) from the Cloudy experiments were captured by the 3VCPI. These diverse experiments  
14 produced ice habits that included needles, hexagonal plates, columns, bullets and dendrites;  
15 ice aggregates and spheroids were also detected (Fig. 4). These habits scatter the light  
16 differently. However, CASPOL data were in good agreement with ice measurements by the  
17 PPD, small water droplets measured with WELAS (Figs. S4 and S5 in the Supplement).

### 18 **3.2 ACPIM modelling**

19 Validation of ice formation was done by modelling. A modelling tool used in this analysis is  
20 the Aerosol–Cloud–Precipitation Interaction Model (ACPIM), which has been developed at  
21 the University of Manchester in collaboration with the Karlsruhe Institute of Technology  
22 (Connolly et al., 2009). Temperature time series were plotted using the initial experimental  
23 conditions (e.g., chamber temperature, pressure, RH, and CCN concentration) in the model.  
24 Subsequent fitting of the simulated temperature drop to chamber data enabled us to find the  
25 rate at which the chamber reheats after expansion ( $0.007 \text{ s}^{-1}$ ) for the runs specified in Table 1.  
26 This heat exchange coefficient is in a good agreement with the results found by Dias et al.  
27 (2016). It quantifies how effectively heat is transferred from the chamber walls and mixed  
28 throughout the gas in this chamber.

29 ACPIM was able to replicate the observed particle phase transitions in the mixed phase runs,  
30 thereby validating the phase concentration plot (Fig. 2c). Phase concentration deviations at  
31 the beginning of the expansion were probably caused by inhomogeneity in the chamber due  
32 to incomplete mixing, or by variations in the expansion rate. Ambiguous polarisation states of  
33 water, e.g., in super-cooled or frozen droplets, might be resolved by comparing ACPIM to

1 CASPOL data and examining the mismatch. This simulation of the experiment makes it  
2 possible to predict phase concentrations and sizes, supporting the planning of future  
3 experiments and validation of the theories behind the model.

### 4 **3.3 Viscous SOA measurements**

5 The validated discrimination method used in water-ice phase transition analysis was  
6 subsequently applied to investigate SOA phase transition. The viscous SOA growth  
7 experiments reported here were achieved using a controlled, constant flow of precursor gases  
8 and ozone into the chamber at constant, near-ambient pressure, dry conditions, and constant  
9 temperatures, as shown in Table 2 (for details see Järvinen et al., 2015). We observe a growth  
10 in particle diameter from tens of nanometres to more than 1  $\mu\text{m}$  size particles. During these  
11 growth periods (Fig. 5), an increase in the CASPOL backscatter polarisation ratio was  
12 observed, while the  $D_{\text{pol}}/F_{\text{wd}}$  ratio did not change significantly, suggesting the change in  
13 size had less effect on the measurements than did the polarisation. A large part of the  
14 experiment produced extreme particle concentrations above the recommended CASPOL  
15 concentration limit of  $1300\text{ cm}^{-3}$ , where significant coincidence errors would be likely to  
16 occur (D. Baumgardner, personal communication, 2015). Therefore, we limit our discussion  
17 to conditions in which growth to sizes larger than  $0.56\ \mu\text{m}$  in diameter, and concentrations  
18 below  $1300\text{ cm}^{-3}$  occur (for details see Sect. 4). After the growth, RH was increased up to  
19 80% in each experiment in order to observe the phase transitions using optical depolarisation  
20 measurements made with the SIMONE instrument (Järvinen et al., 2015). Several repetitions  
21 of these growth experiments followed by humidification and phase transition were conducted.  
22 The subsequent glass transition formed liquid particles at the end of each experiment. A  
23 significantly lower particle polarisation (more optically spherical) state was detected by the  
24 CASPOL at this stage. As a consequence, we observed the presence of two distinct  
25 polarisation clusters during the growth where highly viscous SOA is expected and after the  
26 phase transition where we expect to see liquid particles. The two clusters are overlaid for  
27 several experiments as shown in Fig. 6.

28 While cooling the chamber and reducing the RH (Run #1515.16) (Fig. 7), the larger optically  
29 semi-spherical particles started to dry. Oxidized  $\alpha$ -pinene SOA compounds generally have  
30 added functional groups (oxygen containing substituents), high polarity, and, thus, lower  
31 vapour pressure (Pandis et al., 1992) than water. As a result of this drying process and the  
32 dynamics of partitioning, CASPOL measures an increase in polarisation. The detailed

1 dynamics of partitioning in SOA from alpha-pinene ozonolysis is described in Donahue et al.  
2 (2014).

3 This increase could be explained as transition to an amorphous aerosol phase with high  
4 viscosity at  $RH \sim 10\%$ ,  $T = -30$  to  $-38^\circ\text{C}$ ,  $P = 102$  kPa as suggested by the hysteresis plot of  
5 Koop et al. (2011). Our results cannot, however, be unambiguously ascribed to the viscosity  
6 transition based solely on the measurements here. We simply note the ability of the CASPOL  
7 to identify very significant polarisation shifts in the aerosol scattering properties that are  
8 likely associated with changes in their physico-chemical properties.

9 Additional support for this hypothesis comes from SMPS measurements. No particles were  
10 detected in the SMPS size range in the transition period; the upper cut-off of the  
11 measurement was about 400 nm. Small decay of the averaged diameter is observed in  
12 CASPOL (Fig. 8). These data indicate a wet to dry transformation of essentially large  
13 particles. This reversed transition of the viscosity is then followed by much slower  
14 partitioning or dissociation within these particles, and a decrease in their concentration and  
15 sizes due to constantly decreasing RH.

### 16 **3.4 Particle classification maps**

17 In order to map the whole range of atmospheric processes under future emissions scenarios, it  
18 will be necessary to identify the particles. A new strategy to categorize dust groupings was  
19 developed by Glen and Brooks (2013, 2014) whereby optical scattering signatures from  
20 CASPOL measurements were used to develop a set of threshold rules based on polarisation  
21 ratios. These rules can be used to classify types of dust sampled in the laboratory and during  
22 field campaigns. A plot of the total backscatter intensity as a function of the polarisation ratio  
23 for various types of dust clearly shows the difference in their signatures. Similar techniques  
24 for classifying aerosols are already in use by the Light Detection And Ranging (LIDAR)  
25 community (Burton et al., 2012; Petzold et al., 2010). To explore the feasibility of using the  
26 signature method in CLOUD, we have collated polarisation ratio ranges of many particles  
27 measured in the CLOUD 8 and 9 campaigns. Here we present the polarisation map (Fig. 9)  
28 combining the CLOUD campaign measurements with those obtained from aircraft flights  
29 over the North sea (Johnson et al., 2012) using the same CASPOL instrument. This map  
30 makes it possible to predict the coordinates of other potential organic compounds in the upper  
31 area. Salts, ash, and ice are in the mid-range of the  $D_{pol}/B_{ck}$  ratio; spherical liquids are at  
32 the bottom. Further separation by size might be possible on the  $x$  axis. In comparison

1 between SIMONE and CASPOL for SOA data points from CLOUD, we can see on the map  
2 that SOA – CLOUD 8 (+10°C) data points have lower polarisation ratio compared to other  
3 organic aerosols. This measurement implies lower viscosity and could explain the non-  
4 existent phase transition in SIMONE depolarisation measurements for this experiment. More  
5 experimental data is needed to fill the space for other particles, temperatures and RH.

6 Classification of small ice and water by size characteristics has limited accuracy (Heymsfield  
7 et al., 2006). As explained earlier CASPOL can differentiate between the asphericities of the  
8 particles. The ice presented on this map is aspherical. Slight changes in the polarisation state  
9 of droplets can also be observed as the droplets cool and a crystalline pattern emerges. This  
10 discrimination technique could be used in chamber measurements with mixtures of CCN and  
11 Ice Nuclei (IN) and with some limitations could be applied in explicit atmospheric  
12 measurements albeit with higher uncertainty due to potentially significant overlap in  
13 polarisation responses, particularly in real environment with high diversity of particles.

#### 14 **4 Discussion**

15 The results presented in this paper (Figs. 2, 5 and S4, S5 in the Supplement) illustrate the  
16 ability of the CASPOL instrument to provide reliable Particle Size Distribution (PSD) in  
17 expansion chamber campaigns, and to classify atmospheric particles of different phases,  
18 viscosities, shapes, and sizes. The polarisation ratio was combined with the PBP clustering  
19 technique to highlight the time resolved aspherical fraction evolution.

20 Despite the known limitations and uncertainties in these measurements, e.g., particle  
21 sedimentation (Chapter 6 in Kulkarni, 2011), electronic “ringing”, and leakage currents  
22 (Kramer, 2002), these did not affect the filtered results (Figs. 3b,3c) shown here. Another  
23 uncertainty is contributed by the extremely high aerosol concentrations  $\sim 40\,000\text{ cm}^{-3}$  (with  
24 unresolvable interarrival- times between successive particle). These concentrations may not  
25 be atmospherically relevant; their role here was solely to grow the larger SOA particles (>500  
26 nm). This was required to allow the optical detection of particles during growth and  
27 liquefaction.

28 In addition to concentration issues, a derivation of equivalent diameters from dry viscous  
29 aerosol particles may be challenging since it has been argued that spherical aerosols can be  
30 considered as purely a “figment of the imagination” (Baran et al., 2013). However, particle  
31 sizes measured by CASPOL and UHSAS during SOA growth corresponded well. The  
32 predicted SOA behaviour (Koop et al., 2011) and the measured slow increase of polarisation

1 may suggest a change in the viscosity of these particles. The polarisation transitions observed  
2 were both clear and repeatable which gives confidence in our ability to identify the  
3 hypothesised transitions and to place these observations on the general polarisation map for  
4 classification in a comparative particle analysis.

5 The general classification map presented here demonstrates a good agreement between  
6 chamber and airborne measurements (Fig. 9). Although super-cooled droplets, ice and other  
7 particle polarisation footprints seem to be quite distinct, it is clear that further spatial growth  
8 and branching of ice could lead to a significant increase in polarisation and possibly  
9 significant overlapping of different species. One of the aims of future studies would be to test  
10 aggregation and branching impacts on CASPOL signals. Slightly higher polarisation of the  
11 airborne super-cooled droplets and ice might be the result of aerosol ageing. Processes such  
12 as aerosol ageing will influence subsequent phase separation processes within the droplet but  
13 are difficult to reproduce in a chamber.

14 In the real atmosphere, the particles are more complex; contain additional polarising  
15 constituents and have more branching. Froyd et al. (2010) report the coexistence of mixtures  
16 of partially or fully neutralised sulphate with organic material, nucleated ice crystals, dry  
17 ammonium sulphate, and glassy particles in the Tropical Troposphere Layer (TTL). Ice  
18 residuals were also similar in size to unfrozen aerosol. Lawson et al. (2008) suggests a  
19 thorough investigation of nucleation and growth mechanisms of ice particles in TTL at low  
20 temperatures is needed, particularly in the presence of sulphates mixed with organics and  
21 very high relative humidity. This might be difficult due to increasing anthropogenic SO<sub>2</sub>  
22 emissions which may increase the formation of sulphuric acid aerosols and thus small ice  
23 crystals in the TTL (Notholt et al., 2005). The increase in small ice concentration in presence  
24 of aerosols may complicate ice content measurements even further. The classification map  
25 presented here represents one approach to facilitate future CASPOL-PBP data analysis of the  
26 TTL and deep convective outflow regions. It could also be useful for particles like  
27 ammonium sulphate that often reach high altitudes through the seasonal biomass burning  
28 processes and initiate ice nucleation. Using a method such as the classification map presented  
29 here to discriminate between different kinds of atmospheric particles (e.g., ice crystals,  
30 ammonium sulphate, volcanic ash, SOA) will allow better insight for atmospheric transport  
31 and chemical processes.

## 1 **5 Conclusions**

2 The CLOUD 8–9 campaigns at the CERN facility, introduced a new capability of this facility  
3 for cloud particle measurements (Cloudy). In this paper the first CASPOL Cloudy  
4 measurements of mixed phase and ice clouds are presented. We discuss the advantages of  
5 particle by particle analysis of the polarisation. Single-particle polarisation was used here to  
6 discriminate water, ice, SOA, and other atmospheric particles. The capability to detect  
7 viscous oxidized alpha-pinene with the CASPOL is reported for the first time.

8 We present observation of reversed transition from liquid to viscous based on CASPOL,  
9 SMPS measurements, and SOA modelling. In our experiments, the SOA viscous to liquid  
10 transition is shown to be a reversible process. This result contributes to our understanding of  
11 viscous SOA appearance in the atmosphere, ageing and potentially to the solar radiation  
12 budget calculations.

13 Classification using the clustering technique produced a classification map that can contribute  
14 to future chamber and, possibly, atmospheric measurements of small particles with CASPOL  
15 in a heterogeneous environment. Small ice particles formed during different stages of the  
16 cloud still pose a great challenge for the optical instruments. Future efforts will focus on  
17 classification of additional cloud particles using CASPOL.

18 *Acknowledgements.* We would like to thank CERN for supporting CLOUD with important  
19 technical and financial resources, and for providing a particle beam from the CERN Proton  
20 Synchrotron. We express great appreciation for the CLOUD collaboration and the volunteers  
21 for the night shifts. We would also like to thank Darrel Baumgardner for CASPOL data  
22 filtering advice and review of the manuscript. T. B. Kristensen gratefully acknowledges  
23 funding from the German Federal Ministry of Education and Research (BMBF) through the  
24 CLOUD12 project. This research has received funding from the EC Seventh Framework  
25 Programme (Marie Curie Initial Training Network “CLOUD-TRAIN” no. 316662) and Swiss  
26 National Science Foundation (SNSF) grant no. 200 021\_140 663. The CAPS instrument used  
27 in this work was supplied by the National Centre for Atmospheric Science. The UHSAS was  
28 funded by NERC grant NE/B504873/1.

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- 1 Table 1. Experimental parameters of the expansion runs presented in this paper. Excess  
 2 pressure profile  $x$  axis is of the order of several minutes.

Run#	Seed type	Seed concentration [cm <sup>-3</sup> ]	Excess Pressure profile [mb]	T <sub>initial</sub> [°C]	$RH_{ice}^{max}$ [%]
1248.13	Ammonium Sulphate	3000		+10C	107
1291.16	Sulphuric Acid	75		-30C	168, 135
1298.20	Sulphuric Acid	700		-50C	148
1311.03	Sulphuric Acid	3260		-10C	123
1471.34	Oxalic Acid	100		-20C	165

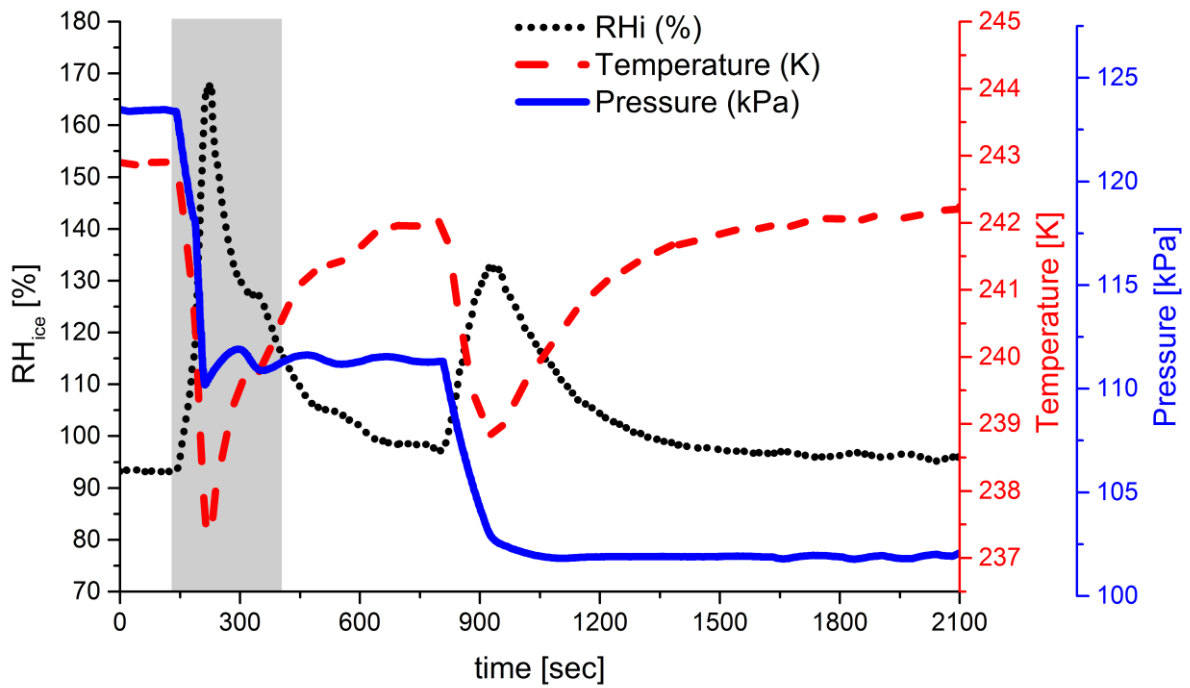
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1 Table 2. SOA growth experimental conditions of the presented runs.

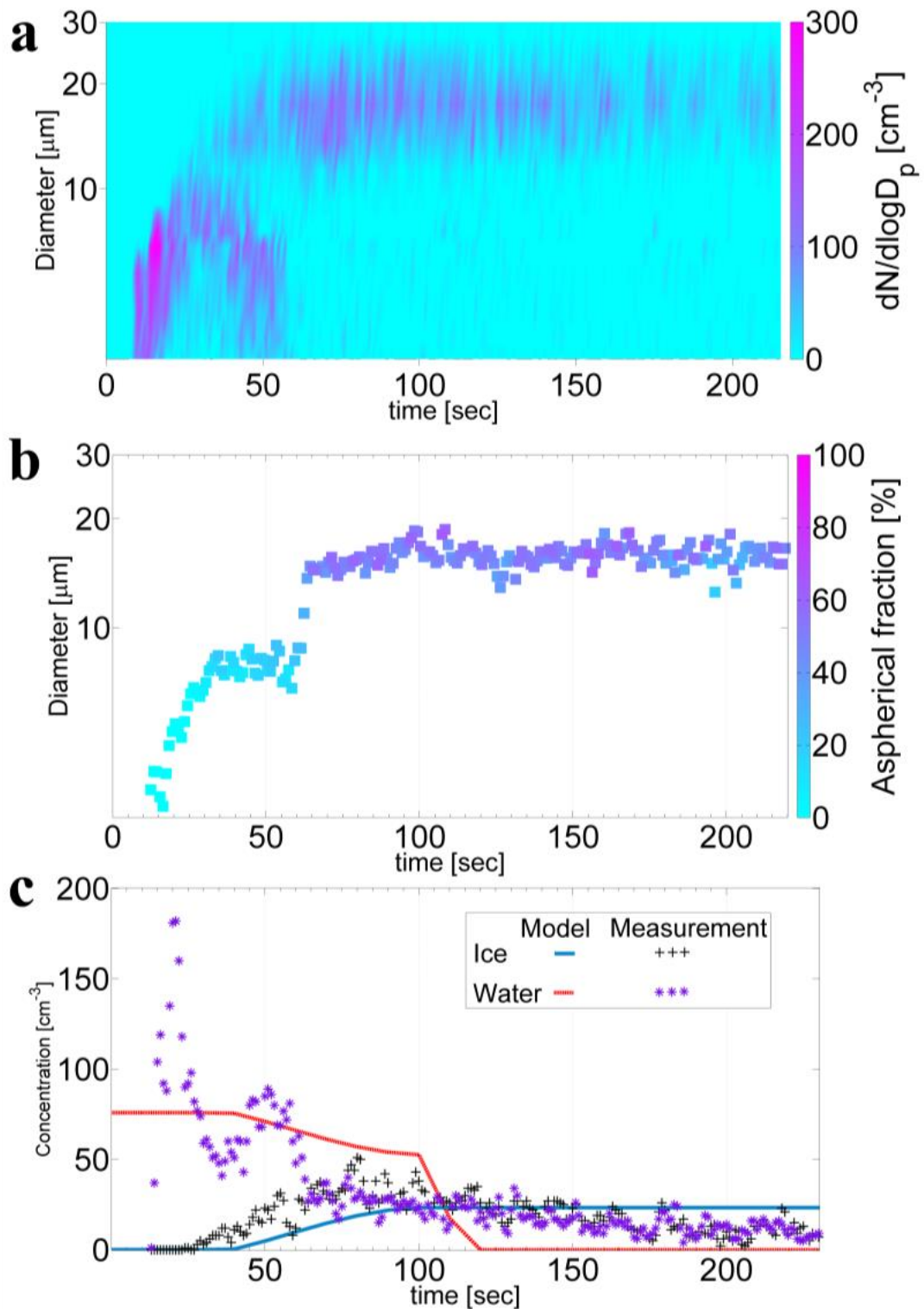
Run	T [°C]	Initial RH [%]	Max. concentration [x1000 cm <sup>-3</sup> ] (Diameter>10nm)
1313	+10	12	30
1513	-20	60	45
1514	-20	4	40
1515	-30	2	30
1516	-38	5	45

2



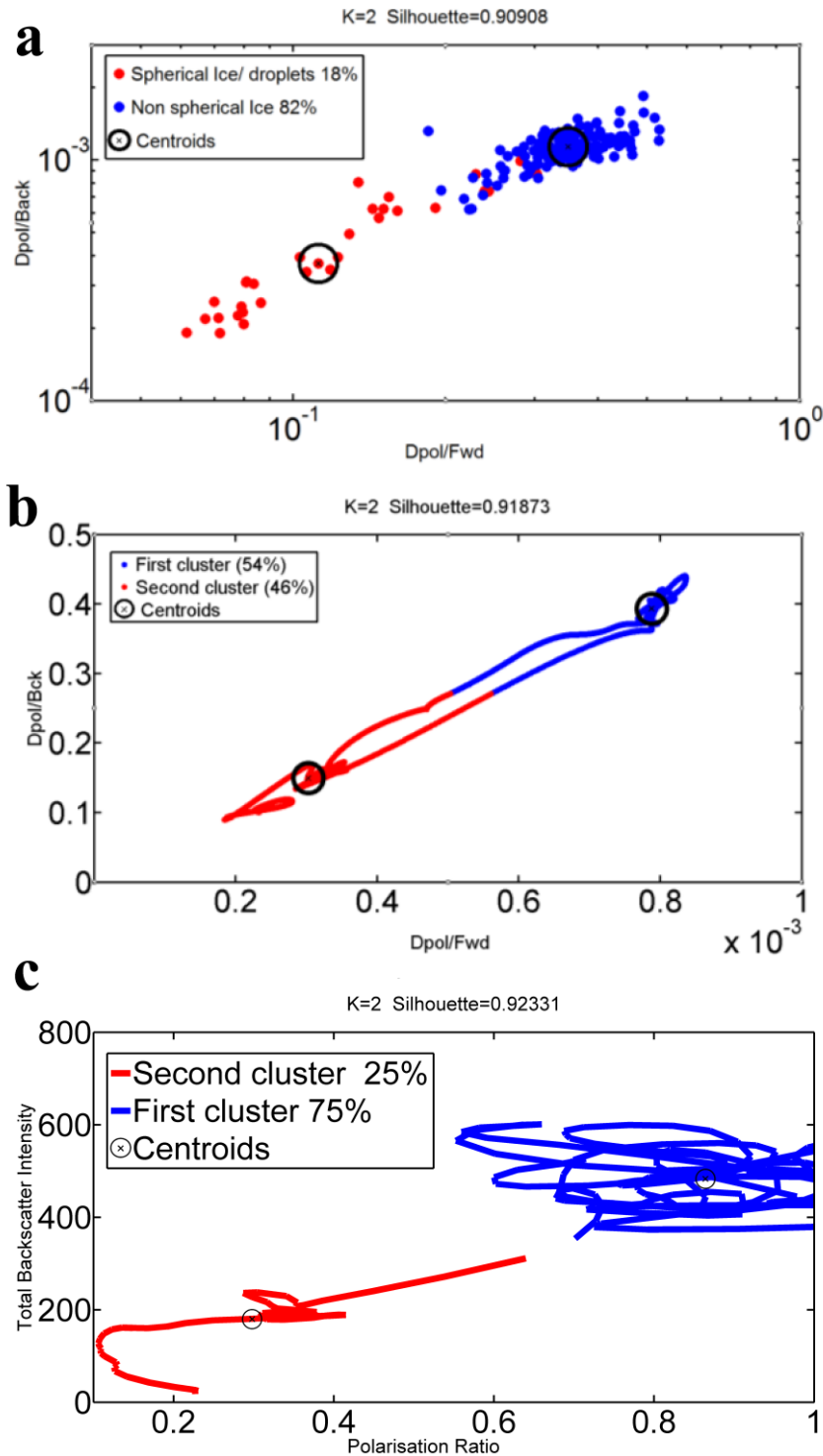
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2 Figure 1. Example of programmable multistep expansion to form a mixed phase cloud (Run  
 3 #1291.16). Relative humidity with respect to ice (RH<sub>ice</sub>) calculated from MBW and  
 4 Thermocouples. Second step grows the present ice particles in the cloud period (25 min).  
 5 Shaded time period is analysed in Fig. 3.



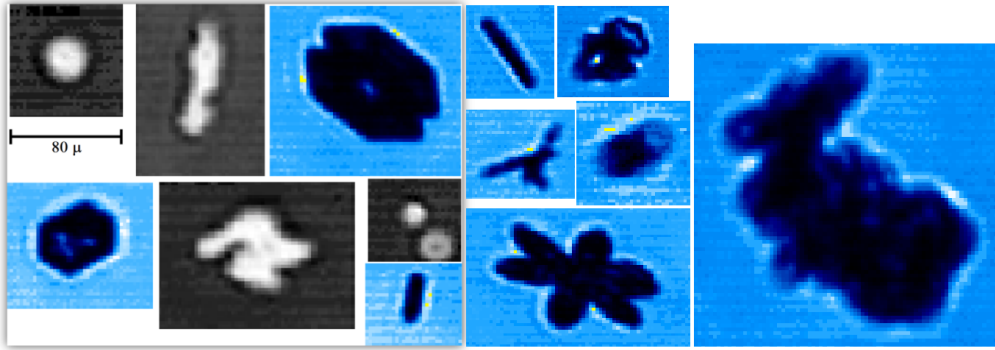
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2 Figure 2. Mixed phase cloud, phase transition period (Run #1291.16). The uncertainty in  
 3 sizing is in the order of the size bin width (Table S2). The error of the polarisation ratio and  
 4 asphericity is approximately 20 %. (a) CASPOL particle size distribution, (b) CASPOL PBP  
 5 aspherical fraction, (c) CASPOL measured water and ice concentrations derived from  
 6 asphericity compared to ACPIM.



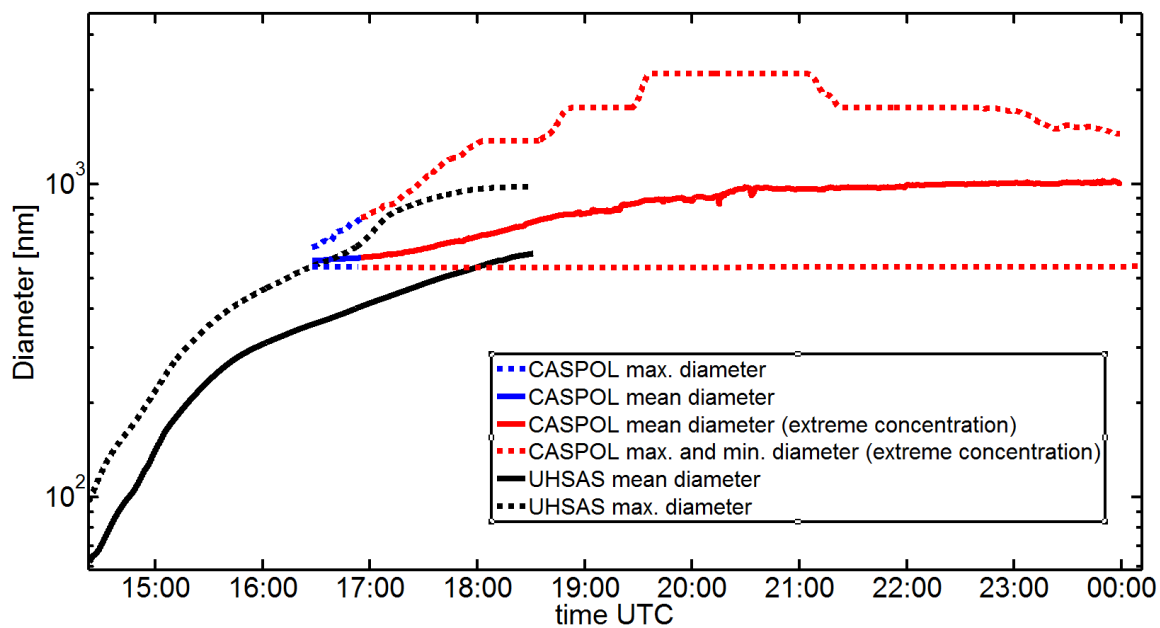
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2 Figure 3. Cluster analysis (Run #1291.16). K in the title indicates the number of clusters  
 3 found with best silhouette value. Each cluster appears with a percentage of particles in it. The  
 4 centres of clusters are marked by centroids ⊗. (a) 1 s averaged data, whole size range and all  
 5 concentration, (b) particle by particle data clustering for selected size range and concentration  
 6 thresholds, (c) particle by particle data clustering plotted in a space comparable to Glen and  
 7 Brooks (2013).



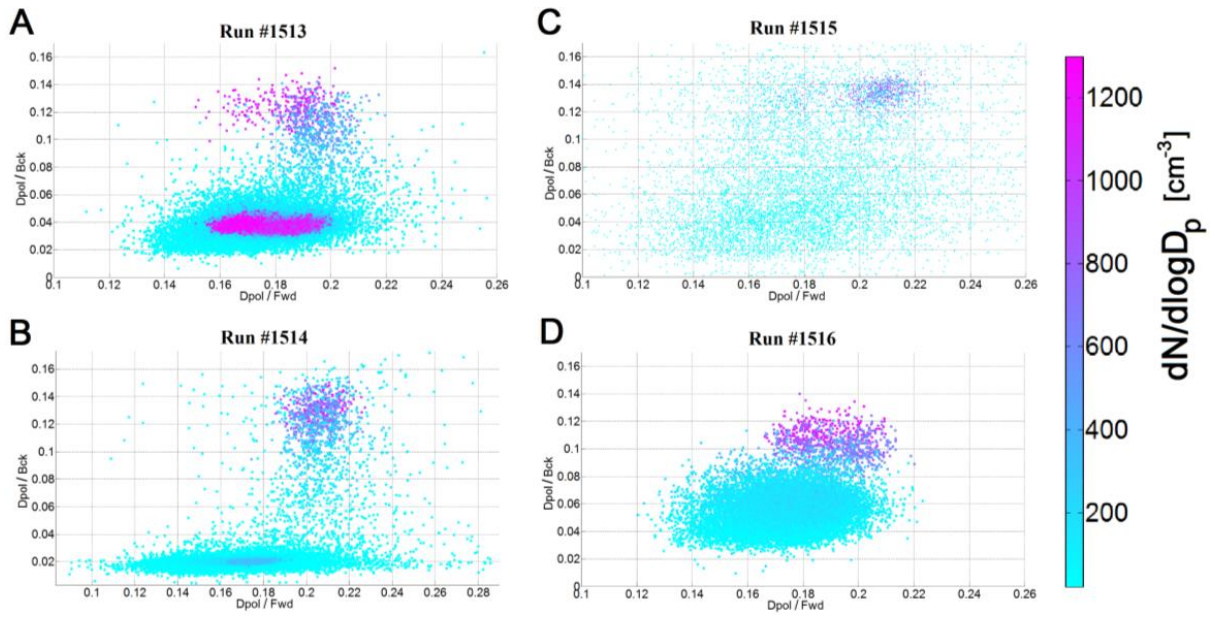
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2 Figure 4. Images of ice particles in CLOUD captured by 3VCPI with 2  $\mu\text{m}$  resolution. Most  
3 of the particles are smaller than 100  $\mu\text{m}$  (scale on the left).



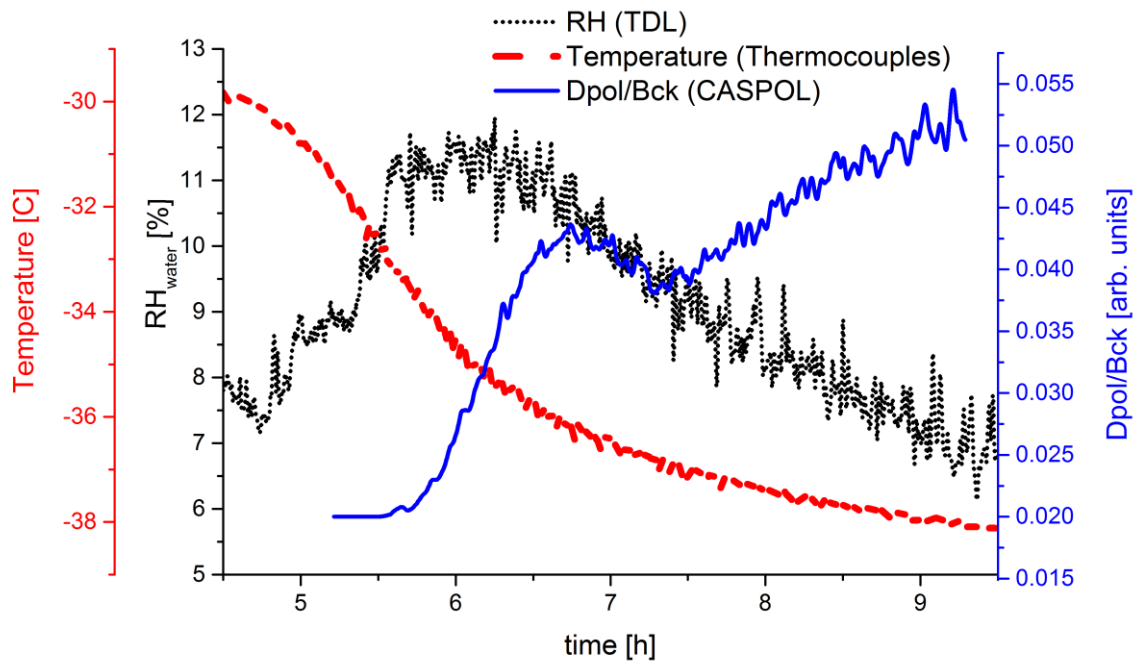
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2 Figure 5. SOA growth over a 10 h period, 1 Hz sampling rate (Run #1516). CASPOL and  
 3 UHSAS overlapped size measurements. Black lines – particles measured with UHSAS,  
 4 instrument's cut-off is at 1000 nm. Blue lines – particles measured with CASPOL. Red lines  
 5 indicate that CASPOL has passed the saturation threshold and the measurements may be  
 6 subject to coincidence errors.



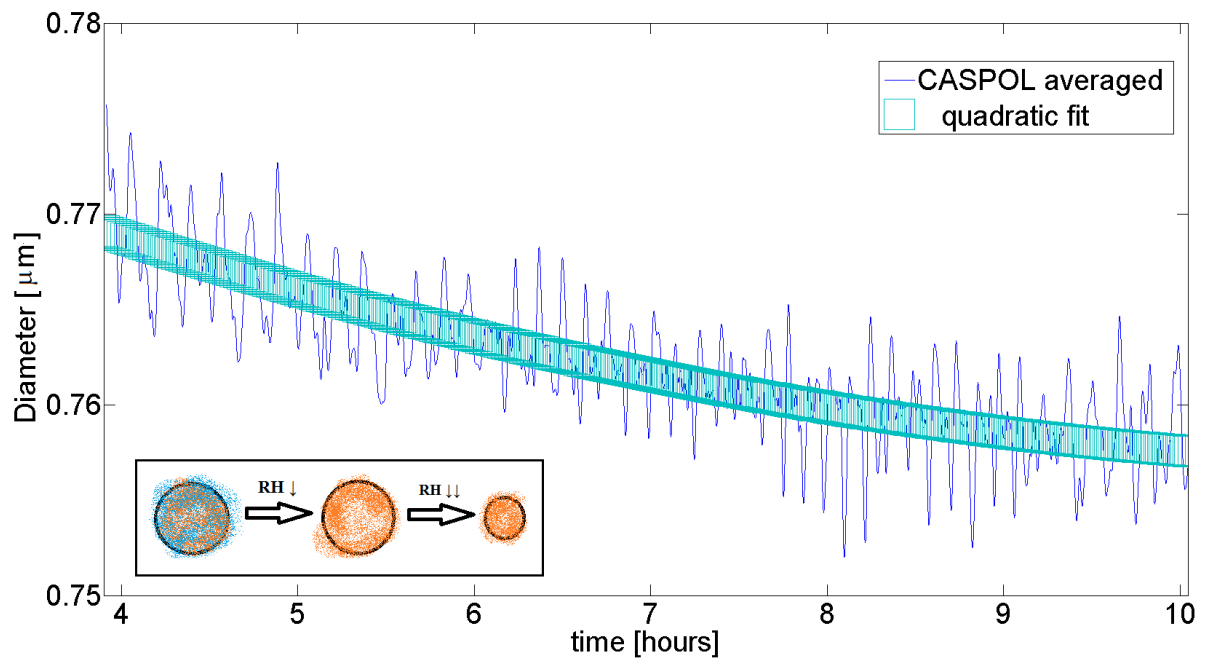
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2 Figure 6. Polarisation scatter-plots of SOA growth and liquefaction measured by CASPOL in  
 3 four experiments. Ratio of perpendicularly polarised backscatter intensity to total backscatter  
 4 intensity ( $D_{pol}/B_{ck}$ ) vs. ratio of perpendicularly polarised backscatter to forward scatter  
 5 intensity ( $D_{pol}/F_{wd}$ ), 1 s averaged run periods where the concentration was below  $1300 \text{ cm}^{-3}$   
 6  $^3$ , colour is concentration  $dN/d\log D_p [\text{cm}^{-3}]$ , (a) Run #1513, (b) Run #1514, (c) Run #1515,  
 7 (d) Run #1516.

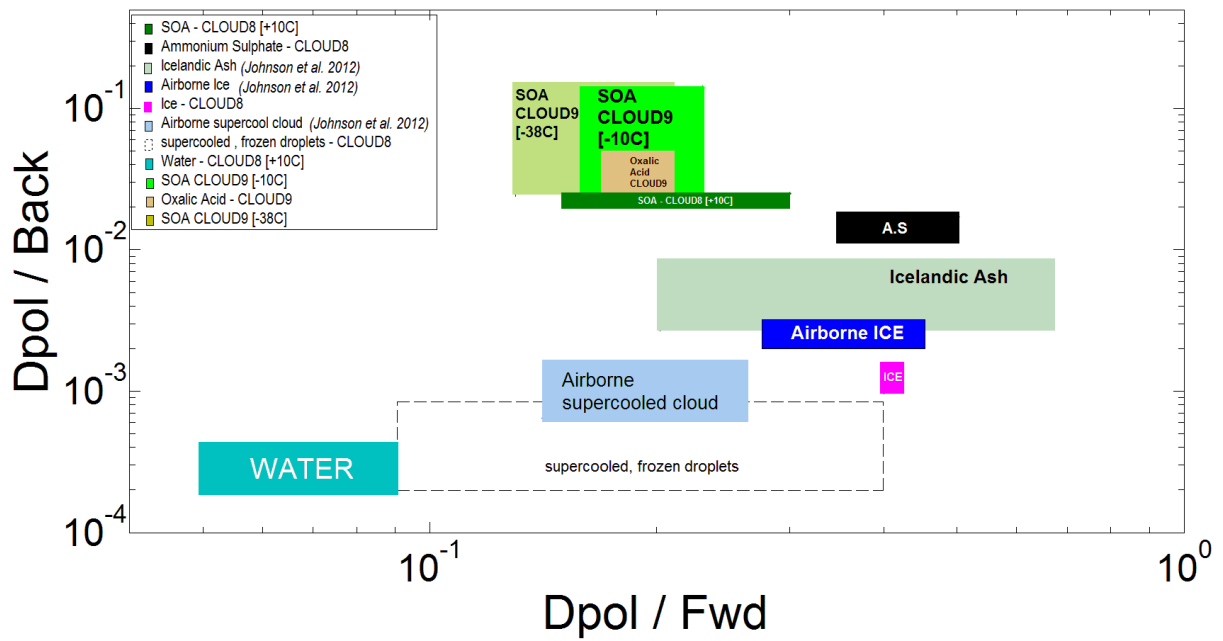


- 1
- 2 Figure 7. CASPOL polarisation ratio (blue line) increases as RH (black dotted line) decreases
- 3 during the cooling period after a SOA experiment (Run #1515.16).





1  
 2 Figure 8. Large dry particles decrease in size. Smaller frame: illustration of the hypothesised  
 3 transition sequence from CASPOL and SMPS measurements (liquid to viscous and dried  
 4 further).



1

2 Figure 9. Atmospheric particle classification map for CLOUD data. The dimensions of the  
 3 coloured rectangular boxes represent the space of measurements error and data points'  
 4 distribution. Additional CASPOL data points from aircraft measurements are presented for  
 5 comparison (Johnson et al., 2012).