

Environmental Controls on Pyrocumulus and Pyrocumulonimbus Initiation and Development

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Abstract:

In this paper we present the first direct observational evidence that the condensation level in pyrocumulus and pyrocumulonimbus clouds can be significantly higher than the ambient lifted condensation level. In addition, we show that the environmental thermodynamic profile, day-to-day variations in humidity, and ambient wind shear all exert significant influence over the onset and development of pyroconvective clouds. These findings are established using a scanning Doppler lidar and mobile radiosonde system during two large wildfires in Northern California, the Bald and Rocky Fires. The lidar is used to distinguish liquid water from smoke backscatter during the plume rise, and thus provides a direct detection of plume condensations levels. Plume tops are subsequently determined from both the lidar and nearby radar observations. The radiosonde data, obtained adjacent to the fires, contextualizes the lidar and radar observations, and enables estimates of the plume ascent, convective available potential energy, and equilibrium level. A note worthy finding is that in these cases the Convective Condensation Level, not the Lifted Condensation Level, provides the best estimate of the pyrocumulus initiation height.

26 **1 Introduction**

27 Pyrocumulus (pyroCu) form when wildfire convective plumes rise to their condensation
28 level and subsequently develop cumuliform cloud tops (American Meteorological
29 Society 2015). The extent of pyroCu development depends on the relationships amongst
30 atmospheric stratification, ambient moisture, and fire fluxes of heat and moisture (Potter
31 2005; Luderer et al. 2006; 2009; Frietas et al. 2007). Some pyroCu release significant
32 moist instability aloft and thereby trigger deep convective clouds that sometimes grow
33 into pyrocumulonimbus (pyroCb). Compared to their lesser counterparts, pyroCb po 
34 glaciated cloud tops and can thus generate precipitation, downdrafts, and lightning
35 (American Meteorological Society 2015). In exceptional cases, pyroCbs have been linked
36 with extreme fire growth (Peterson et al. 2015), devastating firestorms (Fromm et al.
37 2006), and even fire-induced tornados (Cunningham and Reeder 2009; McRae et al.
38 2013).

39 In addition to their impact on fire behavior, pyroCu/Cb have garnered significant
40 research attention due to their affect on vertical smoke transport, atmospheric chemistry,
41 and cloud microphysics. For example, pyroCu can cause significantly deeper smoke
42 injection than in dry convective cases (Frietas et al. 2007) and pyroCb are now
43 recognized as the source of previously unexplained aerosol layers lofted in the lower
44 stratosphere (Fromm and Servranckx 2003; Fromm et al. 2006; 2010). In addition,
45 satellite and dual polarimetric radar observations of pyroCb show that the extreme
46 aerosol loading results in high concentrations of small ice particles (Rosenfeld et al.
47 2007), especially as compared to nearby clouds forming in smoke free air. The
48 abundance of ice particles changes the radiative properties of the clouds and also favors
49 atypical positive polarity lightning strokes (Rosenfeld et al. 2007; Lang et al. 2006;
50 2012).

51 Despite the significant research on pyroCu/Cb microphysics, surprisingly little is
52 known about the environmental controls on pyroCu development. To date only a handful
53 of studies explicitly examines the thermodynamic and kinematic structure of these cloud
54 topped convective columns (Potter 2005, Trentman et al. 2006; Luderer et al. 2006; 2009;
55 Frietas et al. 2007) and no studies include direct observations of pyroCu/Cb initiation. As
56 a result, there is an open scientific debate regarding the plume condensation level, which

57 is an important parameter for modeling smoke injection height and plume evolution
58 (Frietas et al. 2007). Specifically, there are contrasting views in the literature about
59 whether the plume condensation level is expected to be higher than or lower than the
60 ambient lifted condensation level (LCL).

61 Potter (2005), for example, proposes that pyroCu/Cb should exhibit cloud bases
62 lower than the ambient LCL due to the moisture released during combustion of woody
63 fuels and from the evaporation of fuel moisture. Drawing on historical cases of
64 pyroCu/Cb, radiosonde data, and theoretical considerations, he hypothesizes that the
65 latent heat release may be the dominant factor in many moist-pyroconvective events. A
66 limitation of this study is the anecdotal treatment of condensation levels, which are
67 estimated, and the use of radiosonde observations that may not reflect the near fire
68 environment.

69 In contrast to Potter (2005), Luderer et al. (2006; 2009) use high-resolution
70 simulations and theoretical sensitivity calculations to conclude that “the combined effect
71 of released moisture and heat from the fire almost always results in a higher cloud base
72 compared to ambient conditions.” They also find that moisture released in combustion
73 constitutes less than 10% of the pyroCu/Cb water budget with the remainder of the plume
74 water resulting from entrained environmental air. While these modeled results are rather
75 convincing, they lack clear observational support.

76 To that end, the only field observations that address plume moisture are from
77 small scale grass fire experiments, where significant increases in water vapor mixing
78 ratio are documented near the surface, but then decrease rapidly with height (Clements et
79 al. 2006, 2007, Kiefer et al. 2012). While these observations are consistent with the
80 dominant role of entrainment, such small-scale plumes may not be representative of deep
81 convective plumes that extend into the upper troposphere or even lower stratosphere.

82 In this paper we present the first direct observations of condensation levels in two
83 wildfire pyroCu/Cb cases. The fires, the Bald Fire and the Rocky Fire, were located in
84 northern California, and observations were conducted on 2 August 2014 and 30 July
85 2015, respectively (Fig. 1). The pyroCu cloud bases and plume rise dynamics were
86 measured using a mobile atmospheric profiling system (Clements and Oliphant 2014)
87 that included a scanning Doppler lidar and an upper-air radiosonde system which

88 provided thermodynamic profiles immediately upstream of the fire perimeters. From
89 these data, our results clearly show that observed plume condensation levels are
90 substantially higher than the ambient LCL. Additional aspects of the plume rise,
91 including limiting factors on convective growth and the role of environmental moisture
92 are also examined.

93 **2 Data and Methods**

94 **2.1 Lidar Data**

95 In this study, data from a Halo Photonics scanning Doppler lidar are examined (Pearson
96 et al. 2009). The lidar emits a 1.5 μm laser beam and records two range resolved
97 quantities: (1) the attenuated backscatter coefficient ($\text{m}^{-1} \text{sr}^{-1}$), which is a range corrected
98 measure of backscattered energy, and (2) the Doppler velocity (m s^{-1}). The lidar also
99 reports the signal-to-noise ratio (SNR), which is useful for discriminating between
100 meteorological targets and instrument noise. The lidar range is 9.6 km and the range-gate
101 resolution is 18 m. Azimuth and elevation motors allow for scans of the full upper
102 hemisphere and the lidar is level-mounted to the bed of a pickup truck, enabling rapid
103 deployments to wildfires (Clements and Oliphant 2014).

104 **2.1.1 Lidar Sensitivity**

105 Near-IR lidars are sensitive to aerosol, cloud droplets, and forest fire smoke. Due
106 to these sensitivities numerous previous studies have used lidars to examine smoke layers
107 and smoke plumes (Banta et al. 1992; Kovalev et al. 2005; Pahlow et al. 2005 Charland
108 and Clements 2013; Lareau and Clements 2015).

109 Forest fire smoke typically exhibits a log-normal particle number distribution
110 with a peak near .13 μm and a long tail extending towards coarser particles (Radke et al.
111 1990, 1991; Banta et al. 1992; Reid and Hobbs 1998; Reid et al. 2005). The
112 corresponding mass distributions are bimodal with peaks near .1 and 10 μm and a local
113 minimum between 1 and 3 μm (Radke et al. 1990, 1991; Reid et al. 2005). During intense
114 forest fires, such as those in this study, additional “super-giant” aerosol with sizes
115 sometimes exceeding 1 mm may also be prevalent (Radke et al. 1990, 1991; Reid et al.

116 2005). These aerosols are typically composed of large ash and soil particles, which may
117 be scoured from the surface by strong fire-induced winds (Radke et al. 1991; Reid et al.
118 2005; Kavouras et al. 2012).

119 The 1.5 μm lidar beam interacts with the submicron smoke via Rayleigh
120 scattering, the micron sized smoke via Mie scattering, and with the coarsest aerosol (e.g.,
121 large ash, debris, etc.) via geometric optics. Using a radiative transfer model, Banta et al.
122 (1992) showed that the attenuated backscatter coefficient due to the numerous small
123 smoke particles was roughly comparable to the backscatter from the sparse large particles
124 in a given volume. Similar behavior is expected with the lidar used in this study. In
125 addition, based on our own experience, we expect significant attenuation for interactions
126 with very coarse debris, especially near the base of smoke plumes.

127 Near-IR lidars also record high backscatter and rapid attenuation due to cloud
128 droplets, making them an ideal tool for cloud base and cloud top detections (Hogan et al.
129 2003; Winkler et al. 2009). In this study we leverage this attribute of the lidar to
130 determine pyroCu cloud bases and edges in convective column. Similarly, Banta et al.
131 (1992) used an IR lidar to identify pyroclouds in a wildfire smoke column.

132 **2.1.2 Lidar Scan Strategy**

133 The lidar was programmed to conduct “range-height indicator” (RHI) scans
134 centered on the Bald Fire and Rocky Fire pyroconvective plumes. The scan azimuth
135 angles were determined visually. During the Bald Fire the RHI elevation step was 0.7° ,
136 whereas an elevation step of 1° was used during the Rocky Fire. Scans were conducted
137 between the horizon and $\sim 85^\circ$ in elevation, with a full RHI sweep taking ~ 1 min during
138 the Bald Fire and ~ 45 sec during the Rocky Fire. Additional scan details, including the
139 azimuth angles, are provided in the following case studies.

140 The lidar was also used to examine the velocity field near the fires and within the
141 convective plumes. For example, the Doppler radial velocity data collected during the
142 RHI scans are used to inspect the plume structure. These data have a resolution of 3-
143 4 cm s^{-1} over a range of $\pm 19 \text{ m s}^{-1}$ (Pearson et al. 2009). In addition, conical scans were
144 interspersed with RHI scans to generate vertical profiles of the horizontal wind using the
145 “velocity-azimuth display” (VAD) technique (Browning and Wexler 1968). The VADs

146 use an elevation angle of 50° and span 360° in azimuth, taking about 1 minute to
147 complete. The post-processed wind speed and direction from the VADs reflect the
148 ambient winds above the lidar.

149 **2.1.3 Plume Edge Detection**

150 The lidar data are post-processed to determine plume boundaries and beam
151 attenuation depth. The edge detection algorithm uses a combination of the lidar signal-to-
152 noise ratio (SNR) and attenuated backscatter coefficient to isolate the plume. Similar
153 approaches are presented in previous studies (Kovalev et al. 2005; Charland and
154 Clements 2013). In our analysis, we first apply a 5th order Butterworth filter with a 5-
155 point window to the SNR data along each lidar beam to eliminate some of the instrument
156 noise. Next we record the radial location of maximum attenuated backscatter coefficient.
157 Starting from that location we search inward along the beam for the first range gate
158 where the SNR+1 drops below 1.01. This point is considered to be the leading plume
159 edge. The same technique is performed searching outward along the beam to find the
160 trailing plume edge. The trailing edge is considered to be the attenuation point provided
161 the SNR+1 does not again exceed the threshold at some further distance. The SNR+1
162 threshold of 1.01 was found to best discriminate between aerosol returns and background
163 noise in our data sets, though other values (e.g., SNR+1=1.02) provide similar results.

164 To demonstrate the algorithm, Fig. 2 presents lidar data extracted from two
165 elevation angles (10.2° , 46.7°) within a full RHI scan of Bald Fire convective column.
166 The lower elevations beam (Fig. 2a) intersects the base of the smoke plume while the
167 upper beam (Fig. 2b) hits the pyroCu. Of note, the SNR+1 associated with cloud is
168 somewhat higher than in the smoke (1.105 vs. 1.089) and beam's attenuation is much
169 more rapid, penetrating only 198 m into the cloud compared to 648 m into the smoke. In
170 the following case studies we show that the sudden reduction in attenuation depth and
171 increase in attenuated backscatter coefficient aloft are robust signatures of pyroCu
172 formation.

173 2.2 Radar Data

174 Data from four National Weather Service (NWS) radars are used to examine plume
175 structure. These 10-cm radars are sensitive to large ash and hydrometeors in the
176 convective column but insensitive to cloud droplets and sub-micron smoke. Radars have
177 been used in numerous wildfires studies (Banta et al. 1992; Hufford et al. 1998; Fromm
178 et al. 2006; Rosenfeld et al. 2007; Jones and Christopher 2010a,b). Recently, dual-
179 polarization radars have been used to examine the microphysics of wildfire plumes and
180 clouds (Melnikov et al. 2008, 2009; Lang et al. 2014).

181 In this study we leverage three aspects of the NWS radars. First we examine the
182 radar echo tops to estimate the maximum smoke injection height. The echo tops are the
183 highest level at which the radar reflectivity exceeds 18 dBZ (Lakshmanan et al. 2013).
184 Second we combine radar reflectivity from multiple radars to generate volume renderings
185 of the pyroconvective plumes. These volumes are constructed by creating a gridded
186 interpolant from all the available contemporaneous radar data. Data from the Medford,
187 Reno, Beale, and Sacramento radars are combined for the Bald Fire, and from the Beale
188 and Sacramento radars for the Rocky Fire. The radar locations relative to the fires are
189 shown in Fig. 1.

190 Finally, we inspect the differential reflectivity (Z_{dr}) data from the Medford, OR
191 radar during the Bald Fire. Z_{dr} is the logarithmic ratio of the reflectivity from the
192 horizontally and vertically polarized radar beams (Markowski and Richardson 2011).
193 When Z_{dr} is large and positive it indicates the presence of large horizontal targets,
194 including needle-like ash particles (Melnikov 2008, 2009). When Z_{dr} is near zero the
195 targets are more spherical (e.g. hydrometeors), and when negative the targets are
196 vertically oriented (e.g., graupel).

197 2.3 Satellite Data

198 Visible satellite observations from GOES-15 are used to characterize the presence of
199 pyroCu above each fire. These data have a spatial resolution of 1 km and a nominal
200 temporal resolution of 15 minutes, depending on the scan schedule. Data from the
201 Moderate Resolution Imaging Spectroradiometer (MODIS) Terra and Aqua satellites are
202 also used. These data include both true color visible images and fire-radiative power

203 (FRP). The nominal resolution is 500 m. FRP is derived by differencing fire pixels from
204 adjacent non-fire pixels using infrared radiance (Wooster 2002) and has been shown to
205 provide high fidelity representation of fire activity during California wildfires (Koltunov
206 et al. 2012; Peterson et al. 2015). FRP data from GOES are also examined.

207 **2.4 Radiosonde Observation and Analysis**

208 Thermodynamic profiles were collected adjacent to both fires using a GRAW™
209 GS-E radiosonde system. These sondes measure temperature, humidity, and wind from
210 the surface to the tropopause, ascending at a rate of $\sim 3 \text{ m s}^{-1}$. The balloons were launched
211 after sunset to avoid interfering with daytime fire-suppression aircraft operations, and as
212 a result the temperature profiles include surface-based stable layers that are not
213 representative of daytime conditions. To address this shortcoming, the afternoon
214 temperature from the truck weather station is used to infer the convective boundary layer
215 (CBL) depth using the “parcel method” (Holtzworth 1964).

216 The sonde data are used to examine the ambient condensation level by
217 considering three lifted parcels (1) the most unstable (MU) parcel, (2) the mixed-layer
218 (ML) parcel, and (3) the convective (CONV) parcel. The MU parcel is the parcel with the
219 highest convective available potential energy (CAPE), whereas the ML parcel is based on
220 the mean temperature and mixing ratio in the lowest 150 hPa. The CONV parcel reflects
221 the surface temperature required for free convection based on the surface mixing ratio.
222 The condensation level for each of these parcels is compared in the analyses below.

223

224 **3 The Bald Fire**

225 The Bald Fire (40.9 N, 121.3 W) was started by lightning late on 31 July 2014. It was one
226 of several lightning ignited fires in northern California and southern Oregon, including
227 the adjacent Eiler Fire. The fire growth patterns on 1 and 2 August, determined from
228 nightly U.S. Forest Service airborne infrared sensing (<http://nirops.fs.fed.us/>), are shown
229 in Fig. 3. Based on these data, the fire consumed 7275 ha of mixed conifer forest during
230 its first day, and by the end of the subsequent day had burned an additional 6821 ha. The

231 weather conditions on both days featured afternoon high temperatures near 30° C,
232 relative humidity of ~15 % and west winds gusting up to 6 m s⁻¹.

233 During its rapid expansion on 2 Aug., the Bald fire developed a towering pyroCu
234 that subsequently matured into pyroCb. Visible satellite data show the pyroCu initiation
235 occurred at 1330 PDT, with continued cloud growth through mid-afternoon (Fig. 4a-d).
236 At 1426 PDT the MODIS-Aqua satellite recorded a detailed image of the growing
237 pyroCu, showing crisp cumuliform cloud features above the fire perimeter with more
238 diffuse cloud elements extending to the northeast (Fig. 4e). The spreading cloud features
239 were detrained from the primary updraft and then advected in southwesterly flow aloft.

240 The total FRP from the Bald Fire at the time of the MODIS-Aqua image was
241 19700 MW summed over 30 fire pixels. The pixel maximum was 2258 MW, though the
242 pyroCu obscures a substantial portion of the fire. For comparison, the earlier MODIS-
243 Terra overpass at 1245 PDT yielded a maximum FRP of just 829 MW and a total FRP of
244 3836 MW summed over 13 fire pixels. Clearly the fire experienced a rapid change in size
245 and intensity during the early afternoon, coincident with the development of the pyroCu.

246

247 **3.1 Lidar Observations**

248 The truck-mounted Doppler lidar was situated ~7 km southwest of the fire where
249 it conducted RHI scans of the windward edge of the developing pyroCu from 1350 to
250 1532 PDT on 2 August. During this time, 95 RHI scans were completed, providing a time
251 and space resolved measure of the plume evolution, including the height of the
252 condensation level. The lidar location relative to the fire perimeter is indicated in Figs. 3
253 and 4e. Figure 4f provides photograph from the lidar vantage point showing the
254 windward edge of the cloud capped smoke column.

255 Figure 5 shows a sequence of lidar scans spanning the 5-minute period prior to the
256 MODIS-Aqua overpass. These data are expressed as the logarithmic attenuated
257 backscatter coefficient (hereafter backscatter) in units of m⁻¹ sr⁻¹. Red and blue dots
258 represent the leading plume edge and attenuation point, respectively, along each
259 individual beam. The backscatter is due to smoke and debris in the lower portion of the
260 plume (below 5500 m) and due to cloud droplets in the pyroCu aloft. The laser beam

261 attenuates in both the smoke and the cloud water, but the cloud edge is marked by a
262 distinct increase in the backscatter and more rapid attenuation. For example the distance
263 between the leading plume edge (red dots) and the attenuation point (blue dots) tends to
264 be much larger in the lower portion of the plume, whereas above 5500 m the attenuation
265 occurs over just a few range gates. These aspects of the data give the pyroCu cloud
266 returns a “crisper” edge.

267 While pyroCu were already present at the beginning of the scan sequence, the
268 data show the development of a new cloud element. Figures 5a, b for example, show only
269 a few points of rapid attenuation and high backscatter aloft, whereas starting at Fig. 5c a
270 new, upright cloud edge is detected. This nascent pyroCu element then rapidly expands
271 during the subsequent RHI scans, reaching a height of at least 8500 m before moving out
272 of the lidar field of view (Fig. 5d-f). As we show in the radar analysis below, the actual
273 plume top was as high as 12 km in the 10 minutes following these scans.

274 The scans, which were roughly parallel to the mean wind direction, also reveal
275 that the plume experienced significant variations in tilt with time, alternating between
276 windward (Fig. 5a) and rearward sloping geometries (Fig. 5f). In fact, the windward
277 protrusion of the plume was as much as 2 km away from its base. Large coherent vortices
278 are also apparent along the plume edge, especially in Fig. 5a,b as the “stair step” pattern
279 in the plume edge detections. Based on the radial velocity data (not shown) the inward
280 clefts in the plume edge correspond to enhanced flow into the plume and outward lobes
281 reflecting flow towards the lidar. Vortices of this sort are a well-known feature of rising
282 thermals and plumes and play a leading role in entrainment (Scorer 1957; Woodward
283 1959).

284 Following the initial plume rise, sustained pyroCu were observed with the lidar
285 until 1532 PDT, at which point the truck was relocated for safety reasons. To determine
286 the plume condensation level, we aggregate data from all of the lidar scans during this
287 period. From this larger data sample, Fig. 6a presents the time-maximum backscatter as a
288 function of height and distance, and in Fig. 6b as a function of height only. In addition,
289 Fig. 6c shows the computed percentiles (5, 50, and 95th percentiles) of the attenuation
290 depth binned into 100 m intervals. Collectively, these data reinforce many of the aspects
291 of the initial plume rise sequence discussed above. For example, there is a persistent

292 transition in backscatter near 5500 m (blue line in Fig. 6a,b). Below this level, the
293 backscatter approximately linearly decreases with height, consistent with the entrainment
294 of clear air into an aerosol-laden plume. In contrast, at 5500 m the backscatter sharply
295 increases (as does the SNR, not shown), corresponding to the condensation level and
296 development of the pyroCu. The backscatter intensity remains high there and above due
297 to the continued presence of liquid water.

298 The attenuation depth also shows a sharp transition at 5500 m (Fig. 6c). Below that
299 level the median attenuation depth increases with height, which is again consistent with
300 the dilution of the smoke plume via entrainment. At 5500 m the attenuation depth (across
301 all percentiles) sharply decreases, converging towards a median value of ~200 m. The
302 rapid attenuation aloft is consistent with the presence of liquid water drops and supports
303 our interpretation that the change in backscatter intensity is due to condensation in the
304 plume. From these data, we therefore conclude that the observed condensation level
305 occurs very near 5500 m and was nearly constant throughout the 1.5 h observation period
306 despite many changes in fire intensity.

307 **3.2 Radar Analysis**

308 Since the pyroCu cloud tops exceeded the lidar range the maximum smoke injection
309 depth is obtained from the radar echo tops product from the NWS radar in Medford, OR
310 (KMAX), which is ~200 km to the northwest. Figure 7a shows the time-maximum of the
311 echo tops above the Bald Fire on 2 August. These data indicate that smoke reached
312 altitudes in excess of 12 km, and thus the convective column rose an additional 3.5 km
313 above the maximum height resolved in the lidar scans. The radar returns also show that
314 the highest echo tops occur in a localized region above the fire perimeter where pyroCu
315 were most prevalent. In contrast, the smoke layers without pyroCu correspond to plume
316 heights closer to 6 km.

317 An additional interesting aspect of the radar data is the presence of deep echo tops
318 southwest (e.g., upwind) of the infrared fire perimeter (solid contours, Fig. 7a). This
319 observation is consistent with the periodic forward tilt of the plume as observed in the
320 lidar backscatter (Figs. 5, 6). We hypothesize that the forward tilt relates to large-scale
321 vortices that form as the plume penetrates through a stable layer at the top of the

322 boundary layer (Saunders 1961), and due to the deflection of the ambient flow around the
323 plume.

324 Figure 7b shows the corresponding time series of the maximum radar echo tops.
325 The pyroCu initiation at 1330 PDT, as shown in Fig. 4b, corresponds to a rapid rise in
326 echo tops from 6500 m to 10000 m. Following the initial plume growth, the plume tops
327 slowly rise until 1420 PDT at which point a second period of rapid plume growth occurs,
328 pushing the convective column to heights above 12 km. The onset of this deep plume
329 penetration closely corresponds to the lidar plume rise sequence shown in Fig. 6, as well
330 as the MODIS-aqua image. The plume heights subsequently subside, remaining near 10
331 km for the balance of the afternoon before diminishing more substantively at night.

332 A 3-D volume rendering of radar reflectivity from the Bald Fire at the time of
333 maximum injection height (1429 PDT) reveals additional aspects of the plume structure
334 (Fig 7c). The isosurfaces for 30, 28, 26, 24, and 18 dbZ are shown, along with the fire
335 perimeters (red shading), lidar scan plane (black dots), lidar plume edge detections
336 (yellow dots), and the lidar derived condensation level (green contour). These volume
337 data show an expansive region of high reflectivity immediately above the fire perimeter.
338 The reflectivity and plume height diminish towards the northeast, consistent with the fall
339 out of the larger soot and ash particles in the downwind direction (e.g., southwest flow
340 aloft). We note that since the radar is not sensitive to cloud droplets or micron sized
341 smoke, it is possible that the cloud edges and some smoke reside outside of the radar
342 volume rendering. It is also clear from these data that the lidar sees only the leading edge
343 of the plume before attenuating in dense smoke and cloud water, consistent with the
344 analyses presented above.

345 The shape of particles within the plume can be inferred by considering the
346 differential reflectivity (Z_{dr}) from the Medford, OR radar at different heights (Fig. 8).
347 Three elevation angles are inspected (0.5, 1.5, and 2.4 deg.), intersecting the updraft core
348 at heights of 4115, 7742, and 11009 m, respectively. The lowest scan shows very high
349 Z_{dr} , indicative of large, horizontally oriented particles, which is consistent with ash
350 (Melnikov 2008, 2009; Lange et al. 2014) (Fig. 8a). In contrast, the mid elevation scan
351 intersects the plume above the condensation level and shows a significant reduction in
352 Z_{dr} , with values between 0 and 2.5 in the updraft core (Fig. 8b). These values correspond

353 to more spherical particles and small ice, suggesting the presence of large hydrometeors.
354 Finally the upper-most portion of the plume, at ~11 km, exhibits negative Z_{dr} , posing the
355 possibility of vertically oriented graupel particles (Fig. 8c).

356 **3.3 Thermodynamic Analysis**

357 The lidar observed condensation level and radar estimated plume tops provide valuable
358 constraints on the plume structure when contextualized with atmospheric profiles
359 collected adjacent to the fire. Figure 9a, for example, shows data from a radiosonde
360 launched at 2100 PDT from ~15 km to the southwest of the fire (location shown in
361 Fig. 3).

362 The sounding shows that the afternoon CBL extends from the surface (1364 m) to
363 ~4000 m and is capped by a pronounced stable layer. Within the CBL, the water vapor
364 mixing ratio is roughly constant at ~5 g kg⁻¹, whereas above the CBL a layer of very dry
365 air is observed with a mixing ratio of only ~0.5 g kg⁻¹. Further aloft, near 400 hPa, a layer
366 of higher humidity air, reflecting monsoonal moisture, is found. The height of the
367 tropopause is ~13 km.

368 Relative to the observed profile, the “in cloud” profile is estimated by pseudo-
369 adiabatically lifting a parcel from the lidar observed condensation level at 5500 m. The
370 resulting parcel possesses 910 J kg⁻¹ of CAPE, which is an upper bound on the energy
371 available for buoyant ascent. The equilibrium level (EL) of the pyroCu parcel is 11,742
372 m, which is in close agreement with the radar estimated echo tops, but does not account
373 for the inertial overshoot of the parcel, which is likely reflected in the localized region of
374 radar plume heights exceeding 12 km (Fig. 7a).

375 Also of note, the homogeneous freezing level (-38° C) in the plume profile occurs
376 at 10,158 m and the temperature at the EL is -52° C, indicating that the upper portion of
377 the cloud must be glaciated. As such, this particular pyroconvective cloud should be
378 classified as a pyroCb. In fact, pyroCb from other nearby fires on that day were known to
379 produce lightning as well as a significant and destructive fire-whirl (Muller and Herbster
380 2014)

381 **3.3.1 Lifted Parcels**

382 One of the main goals of this paper is to compare the observed plume properties with
383 conventional estimates of condensational level and convective potential. To that end, in
384 this subsection we consider each of the three lifted parcels described in Section 2.4 as
385 representations of the observed plume. The parcel ascents are shown in Fig. 9b.

386 In this case, the MU parcel (red line, Fig 9b) originates in the CBL and produces an
387 LCL of 4367 m, which is more than 1 km lower than the lidar observed condensation
388 level. In addition, compared to the observed plume structure, the MU parcel possesses
389 minimal CAPE and must overcome appreciable convective inhibition (CIN) before
390 reaching its level of free convection. Similarly, the ML parcel encounters its LCL at 4641
391 m, possesses almost no CAPE, and also must overcome appreciable CIN (cyan line, Fig.
392 9b). The LCL for the ML parcel is higher than that of the MU parcel because the layer
393 averaged mixing ratio is less than the maximum mixing ratio in the CBL.

394 Interestingly, the CONV parcel provides the best representation of the observed
395 plume (dark blue line, Fig. 9b). In this case the surface mixing ratio is 5.2 g kg^{-1} and the
396 corresponding convective condensation level (CCL) is found at 5549 m, which is very
397 close to the lidar derived 5500 m. Commensurately, the EL and CAPE for the CONV
398 parcel are also close to the observed values. The convective temperature, which is the
399 surface temperature that must be reached to support convection, is 36.4° C . The high
400 temperature for the day was 29° C , making surface based convection extremely unlikely
401 outside of the fire modified environment.

402 From these analyses it is clear that the plume condensation level is substantively
403 higher than the ambient LCL, supporting the results of Luderer et al. (2006; 2009).
404 Further, using the CCL, not the LCL, and assuming that the fire readily exceeds the
405 convective temperature, provides the best representation of the plume condensation level
406 in this case. This is a potentially useful diagnostic for forecasters and fire managers. It
407 should be noted, however, that the CONV parcel, and its associated dry-adiabat up to the
408 CCL (dark blue line, Fig. 9b), does necessarily reflect the actual properties of the lower
409 plume. Rather, the plume must be superadiabatic near its base, cooling largely due to
410 entrainment as it decays towards adiabatic ascent further aloft (Emanuel 1994;
411 Trentmann et al. 2006; Frietas et al. 2007).

412

413 **4 The Rocky Fire**

414 The Rocky Fire (38.9° N, 122.5° W) started late on 29 July 2015 (cause unknown) in the
415 coastal range of northern California and burned in complex terrain through fuels
416 consisting of grass, brush, and conifers (Figs. 1, 10). The U.S. Forest Service NIROPS
417 fire perimeters show that the fire burned 3356 ha during the first day, and then consumed
418 another 3153 ha on 30 July, the day of our observations. The fire growth on the 30th was
419 complex, expanding along multiple flanks (Fig. 10). Notably the first day's fire growth,
420 while rapid, did not generate pyroCu, whereas the second day did (Fig. 11). In addition,
421 compared to the long-lived Bald Fire pyroCu/Cb, the Rocky Fire plumes were transient,
422 repeatedly forming and dissipating in rapid succession. In this section we examine the
423 structure of these transient pyroCu along with the environmental conditions affecting
424 their evolution.

425 **4.1 Lidar Observations**

426 Lidar RHI scans were conducted between 1545 and 2008 PDT from an already
427 burned area within the Rocky Fire perimeter (Fig. 10). This location allowed for scans of
428 four separate pyroCu plumes rising from the complex fire perimeter. A total of 267 RHI
429 scans were performed.

430 PyroCu were first observed with the lidar starting at ~1600 PDT rising from the
431 northwest flank of the fire (319° azimuth, scan path #1 in Fig. 10). Figure 12 shows a
432 sequence of photographs (top panels) and contemporaneous lidar scans (bottom panels)
433 detailing the onset and expansion of this cloud topped plume. The plume was initially
434 observed as it penetrated through a stable layer at the top of the CBL, evident as a lateral
435 smoke layer at 2600 m in the backscatter data and as a diffuse haze in the photographs.
436 During this time a thin pileus cloud accompanied the developing pyroCu and the lidar
437 cloud returns were limited to a few points near the plume top (Fig. 12f).

438 By 16:03 PDT, however, a distinct cumuliform cloud had developed (Fig. 12b)
439 and the lidar backscatter showed a commensurate increase in intensity and attenuation
440 along the pyroCu edge (Fig. 12g). Based on these data the cloud base was at ~4200 m.

441 The subsequent scans show the rapid pyroCu development, and by 1609 PDT cloud
442 edges were detected as high as 7500 m. Interestingly, soon thereafter the pyroCu
443 detrained from the convective column and dissipated (not shown).

444 Another pyroCu event at 1800 PDT is detailed in Fig. 13 corresponding to a lidar
445 azimuth of 86° (scan path #3 in Fig. 11). This plume initiated ~ 2.5 km east of the lidar
446 location. As in the previous case, the rapidly growing plume was first recorded as it rose
447 through the boundary layer top, now at ~ 2300 m, and expanded into the free troposphere
448 (Fig. 13a,e). Subsequent scans revealed the onset of pyroCu with a condensation level of
449 4200 m, which is unchanged from the earlier Rocky Fire pyroCu event detailed in Fig.
450 12. In this case, however, the cloud top was not as well documented because attenuation
451 at the cloud base shielded the lidar view of the upper plume.

452 The photographs detailing the plume rise show changes in smoke colour near the
453 base of the convective column (Fig. 13a-d). For example, at 1805 PDT the smoke is a
454 dark gray (Fig. 13a), whereas later the smoke is increasingly white (Fig. 13d). We believe
455 the change in smoke coloration is associated with changes in the completeness of
456 combustion: flaming combustion produces smoke dominated by black carbon aerosols,
457 whereas smoldering combustion generates more organic carbon aerosol, which more
458 effectively backscatter sunlight and appear whiter (Bellouin 2014; Saleh et al. 2014).

459 An additional aspect of the observed plume rise is the relationship between the
460 updraft strength and the ambient wind. This relationship is examined in Fig. 14, which
461 displays VAD wind profiles (Fig. 14 a,b) and RHI radial velocities detailing the plume
462 structure (Fig. 14c-e). The wind profiles show significant shear over the lowest 2 km of
463 the atmosphere. Strong ($5\text{-}7\text{ m s}^{-1}$) northwesterly winds near the surface transition to
464 weak flow at the boundary layer top ($0\text{-}1\text{ m s}^{-1}$ near 2300 m), then reverse to easterly flow
465 aloft (Fig. 14a,b). The observed near-surface wind speed maximum is atypical in the
466 atmospheric boundary layer and “adverse” wind profiles of this character have previously
467 been linked to blow-up fires (Byram 1954).

468 Compared to the ambient wind, the flow within the plume is characterized by
469 much stronger velocities (Fig. 14c,d). For example, outbound speeds in excess of 15 m s^{-1}
470 are recorded at numerous locations within the plume at 1809 PDT. The actual updraft
471 speed likely exceeds these values since the radial velocity data only reflect the projection

472 of the updraft onto the oblique lidar beam. Significant downdrafts are also observed in
473 the upper portion of the plume, especially at 1813 PDT (blue shading, Fig. 14c,d).

474 The RHI velocity data also show that the strength of the updraft diminished with
475 time. For example, comparative histograms demonstrate that strong outbound velocities
476 were both higher and more common at 1809 than 1813 PDT (Fig. 14e). This observation
477 is consistent with the change in smoke coloration described above: flaming combustion is
478 likely to produce stronger updrafts due to more rapid heat release.

479 Changes in plume geometry also accompany the reduction in updraft strength.
480 The plume is at first mostly upright (Fig. 14c) and later becomes more sheared (Fig. 14d).
481 Specifically, the leading plume edge becomes tilted downwind within the boundary layer
482 while the upper portion of the plume bends back towards the observing location. Based
483 on these data, we hypothesize that as the fire's updraft weakens it becomes more
484 susceptible to the ambient shear. The role of wind shear as a limiting factor in plume
485 development is further discussed below.

486 A more robust examination of the plume condensation level during the Rocky
487 Fire's multiple pyroCu events is presented in Fig. 15. These analyses leverage the
488 aggregated data from all of the RHI scans on 30 July. Unsurprisingly, the time-maximum
489 backscatter exhibit a sharp transition near 4200 m (Fig. 15a, b), as was indicated in the
490 earlier plume rise sequences (Figs. 12, 13). Below 4200 m the backscatter decays roughly
491 linearly with height, and above that level the backscatter converges to a value of near -4
492 $\text{m}^{-1} \text{sr}^{-1}$ (Fig. 15b). Likewise, the attenuation depth linearly *increases* from the surface up
493 to 4200 m, then abruptly decreases to a median depth of ~ 200 m. This pattern is
494 consistent with the dilution of the smoke plume by entrainment and the onset of
495 condensation aloft. Importantly, these analyses are remarkably similar to those during the
496 Bald Fire suggesting a clear lidar signature of pyroCu onset. Moreover, the condensation
497 level is once again found to be constant throughout the observing period, indicating that
498 ambient atmospheric conditions rather than variations in water released during
499 combustion likely control its height.

500 **4.2 Radar Analysis**

501 Radar data are again used to estimate the maximum smoke injection depth. The Rocky
502 Fire was within ~100 km of both the Sacramento (KDAX) and Beale Air Force Base
503 (KBBX) NWS radars (Fig. 1), and data from both sites are in agreement.

504 The maximum echo tops (from KDAX) occur between 7000 m and 7500 m,
505 consistent with the lidar cloud detections (Fig. 16a). The spatial pattern of echo tops
506 indicate that plumes of similar height developed on all of the expanding flanks of the fire.
507 Interestingly, the corresponding time series demonstrates the plume transience, showing
508 rapid variations in plume height throughout the late afternoon (Fig. 16b). Each spike
509 corresponds to a short-lived pyroCu with durations ranging from 10-30 minutes. Satellite
510 data confirm the episodic nature of these plumes (not shown).

511 The variability in echo tops is also due to the presence of multiple updrafts. For
512 example, a volume rendering of the reflectivity data at 1609 PDT shows the two distinct
513 updrafts associated with the complex fire perimeter (Fig. 16c). The narrow updraft rising
514 from the northwestern flank of the fire is the same plume shown in Fig. 12, and the lidar
515 plume detections agree well with the radar data (yellow dots, Fig 16c). A second broader
516 plume rises from the north and northeastern flanks of the fire at the same time. Above
517 5000 m the upper portions of both plumes are tilted to the north-northwest due to
518 southeasterly flow in that layer. Later in the fire's evolution the plume growth shifted
519 towards the east and southeast (not shown).

520 **4.3 Thermodynamic Analysis**

521 The Rocky Fire pyroCu development is interesting in that the thermodynamic
522 environment theoretically supports much deeper convection than was observed. Using
523 radiosonde data from ~15 km southwest of the fire at 2105 PDT, Fig. 17a shows that
524 moist adiabatic ascent from the observed 4200 m cloud base would generate 2035 J kg^{-1}
525 of CAPE and that the plume equilibrium level would be ~13 km, impinging on the
526 tropopause. The radar and lidar data indicate, however, that the plumes ascended to no
527 higher than ~7.5 km, corresponding to plume top temperature of -20° C . As such, these
528 clouds are best classified as pyroCu, and never developed as deep pyroCb.

529 What then limits the growth? There appear to be two related limiting factors in the
530 plume rise: (1) wind shear, and (2) dry air entrainment. The lidar wind profiles, presented
531 above in Fig. 14, indicate significant wind shear between the CBL and free troposphere.
532 This wind shear is also apparent in the radiosonde wind profile, which shows a 180-
533 degree wind shift at 2300 m (wind barbs and hodograph Fig 17a). The flow below this
534 level is from the west-northwest, whereas the flow above, and extending up to ~7 km, is
535 from the east-southeast. The layer of southeasterly flow is associated with a surge of
536 monsoonal moisture.

537 A second layer of significant wind shear at 7000 m separates the monsoon flow
538 from southwesterly flow in the upper troposphere. This shear also coincides with a rapid
539 decrease in dew point temperature, and thus relative humidity. It is notable then that the
540 maximum echo tops occur only about 500 m above the upper shear layer. Visual
541 observations throughout the afternoon and early evening suggest this shear zone affected
542 the pyroCu development, tending to sweep the upper portion of the cloud away from the
543 updraft core. The detraining upper portions of the cloud subsequently developed ragged
544 and wispy edges indicative of dry air entrainment as opposed to the crisp crenellations of
545 growing cumulus congestus.

546 The effect of the wind shear on a buoyant parcel is easily visualized by examining
547 the ascent track of the radiosonde, which rose at a mean rate of 2.7 m s^{-1} (Fig. 17b). The
548 ambient shear causes a pronounced zigzag pattern that is clearly detrimental to sustained
549 upright convection despite the substantive CAPE. This result is not surprising in that
550 CAPE is known to overestimate convective development and updraft strength
551 (Markowski and Richardson 2011).

552

553 **4.3.1 Lifted Parcels**

554 Despite their limited vertical development the Rocky Fire pyroCu provide
555 additional support for the hypothesis that the plume condensation level occurs above the
556 ambient LCL. Following the same procedures described for the Bald Fire we examine
557 three convective parcels, the ascents of which are shown in Fig. 18. The LCLs for the
558 MU and ML parcels are 3503 m and 3768 m, respectively (red and cyan lines, Fig. 18).

559 Both of these lifted parcels must overcome modest CIN to reach their level of free
560 convection. In contrast, the computed CCL of 4250 m is much closer to the lidar
561 observed condensation level at ~4200 m. The corresponding convective temperature is
562 ~43° C, which is higher than the observed daytime temperature of 39° C. These results,
563 like those from the Bald Fire, again suggest that the CCL is a useful parameter for
564 estimating pyroCu/Cb convective initiation heights.

565 **4.4 Fire Radiative Power and Environmental Moisture**

566 Figure 19a shows the GOES-15 and MODIS FRP in for the Rocky Fire on 29-30
567 July 2015. From these data it is clear that the diurnal cycle of fire intensity is similar
568 during the first two days of fire growth, with peak FRP values near 1500 MW in the late
569 afternoon and fire activity extending into the late evening. Interestingly, despite
570 comparable fire intensity, pyroCu were not observed on the 29th but were widespread on
571 the 30th.

572 To better understand this disparity Fig. 19b-e compares the ambient
573 meteorological conditions between days. These data are from a weather station just
574 outside the fire perimeter, the location of which is shown in Fig. 10 (data obtained from
575 MesoWest, Horel et al. 2002), and the 1700 PDT upper air soundings at Oakland
576 International Airport (KOAK, location shown in Fig. 1). The high temperature on both
577 days was ~39° C and afternoon winds were from the west with peak gusts near 6 m s⁻¹.
578 The relative humidity was higher on the 30th than the 29th.

579 Figures 19d,e show the time series of the water vapor mixing ratio and the
580 differences in the relative humidity from the KOAK soundings for two afternoons. From
581 these data it is apparent that the onset of pyroCu on 30 July corresponds to the arrival of
582 much higher humidity air, both at the surface and aloft. For example, the mixing ratio
583 increases from 4.5 g kg⁻¹ to 8 g kg⁻¹ while the relative humidity at 5500 m jumps from
584 7% to 66%. The corresponding change in the CCL is substantial, dropping from 5848 m
585 on 29th to 4267 m on the 30th. Since the fire intensity was similar on both afternoons it is
586 likely that reduction in the height of CCL due to the influx of monsoon moisture was the
587 driving factor in pyroCu formation. These observations support the conclusions of

588 Luderer et al. (2006; 2009) that environmental moisture, not water released in
589 combustion, is the primary control on pyroCu development.

590

591 **5 Summary and Conclusions**

592 The observations presented in this paper demonstrate that plume condensation levels
593 can exceed the height of the ambient LCL, sometimes substantially. For example, during
594 the Bald Fire the plume condensation level was more than 1 km higher than the
595 environmental LCL. As such, we conclude that the LCL should not be used, as it has
596 been, as a parameter for assessing pyroCu/Cb potential outside of the limiting case where
597 the CCL and LCL coincide, which is to say that widespread convective clouds are
598 possible. While our observational results span a limited portion of the parameter space,
599 they nonetheless provide strong support for the modeling results of Luderer et al. (2006;
600 2009) and Trentman et al. (2006), and seemingly contradict the results of Potter (2005).

601 While the CCL and the corresponding moist adiabatic ascent provide a useful
602 approximation for plume properties, other factors must also be considered. Specifically,
603 CAPE alone cannot determine the convective outcome. Our results from the Rocky Fire
604 show, for example, that ambient wind shear and dry air entrainment can significantly
605 curtail the convective development even in an environment that might otherwise support
606 deep pyroCb. In addition, our results show that the change in environmental humidity,
607 often in the form of a monsoonal surge, exerts a significant influence over the onset of
608 pyroCu/Cb by raising or lowering the height of CCL. These results suggest that the
609 moisture release during combustion is of secondary importance, at least in these observed
610 cases.

611 While our results mark an advance in understanding pyroCu/Cb development
612 there is a clear need for new measurement and modeling investigations of pyroconvective
613 clouds. Future field campaigns should include observations of the ambient environment
614 (e.g. radiosondes, CBL properties), the lower plume structure (temperature, moisture, and
615 momentum fluxes), and cloud properties (e.g. liquid and ice water path, particle size
616 distributions, etc.). These data should subsequently inform physical fluid dynamical
617 models in order to investigate aspects of plume dynamics that may not be observable.
618 Some potential avenues for obtaining these observations include dropsondes from

619 aircraft, surface and aircraft based dual-polarization radars, unmanned aerial vehicles,
620 and dual-Doppler lidar deployed during large-scale prescribed burn experiments where
621 the fuel loading and extent of combustion is known or can be determined after the fact.

622

623 **Author Contributions**

624 C.C. conceived of the field program, N.L. and C.C. conducted the field measurements,
625 and N.L. led the data analysis and writing.

626

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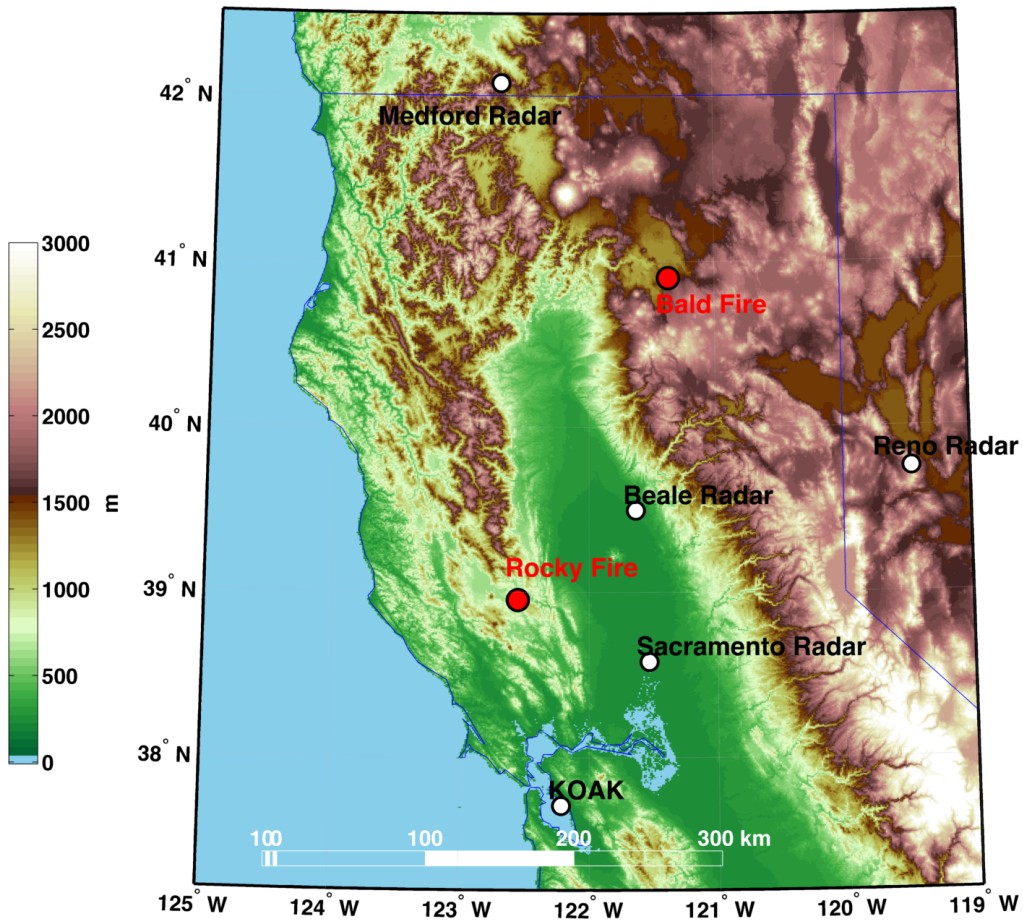
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788 Fig. 1. Overview map showing the regional topography (terrain shading), locations of the
 789 Rocky and Bald Fires (red circles), the locations and names of the NWS radars used in
 790 the plume analysis (white circles), and the KOAK sounding site (white circle).

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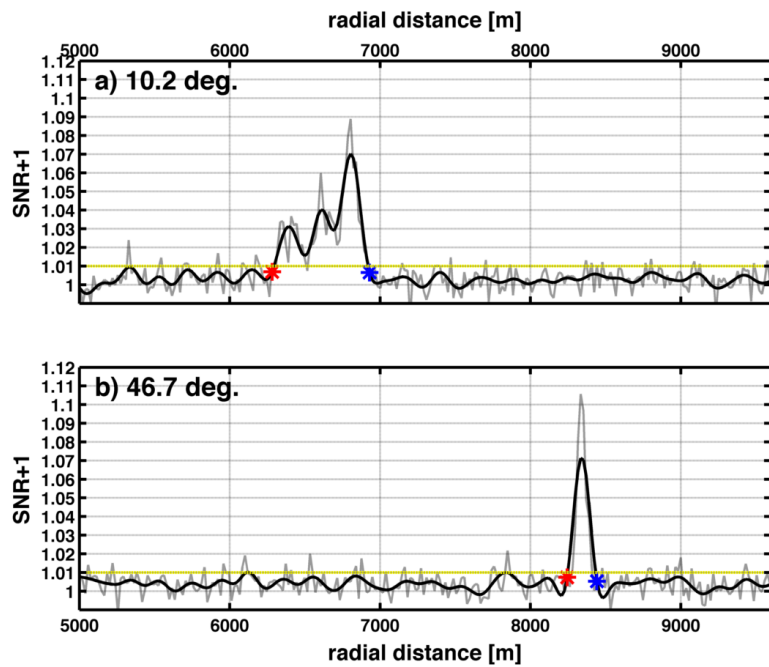
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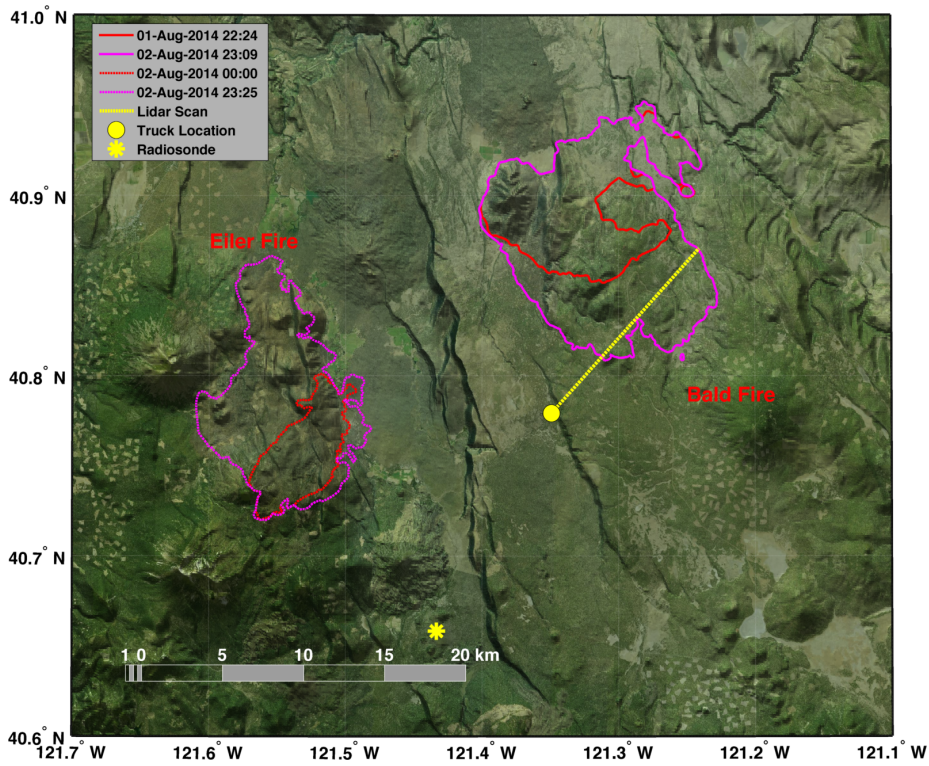
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798 Fig. 2. Examples of the plume detection and attenuation algorithm based on the filtered
 799 (solid black line) and unfiltered (gray line) lidar signal-to-noise ratio (SNR+1). (a) Low
 800 elevation angle (10.2 deg) lidar beam intersecting the base of the Bald Fire convective
 801 column. (b) High elevation angle (46.7 deg) beam intersecting the pyroCu in the upper
 802 plume. The red stars indicate the leading plume edge and the blue stars the attenuation
 803 point. The SNR+1 threshold of 1.01 is indicated with a dashed yellow line.
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806 Fig. 3. Bald and Eiler Fire progression map for 1 and 2 August. The fire perimeters are
 807 from the US National Forest Service National Infrared Operations (NIROPS) flights. The
 808 background is a satellite image draped over the terrain, which is highlighted with hill
 809 shading. Also shown are the truck location (yellow dot), lidar scan path (yellow line), and
 810 radiosonde location (yellow star).

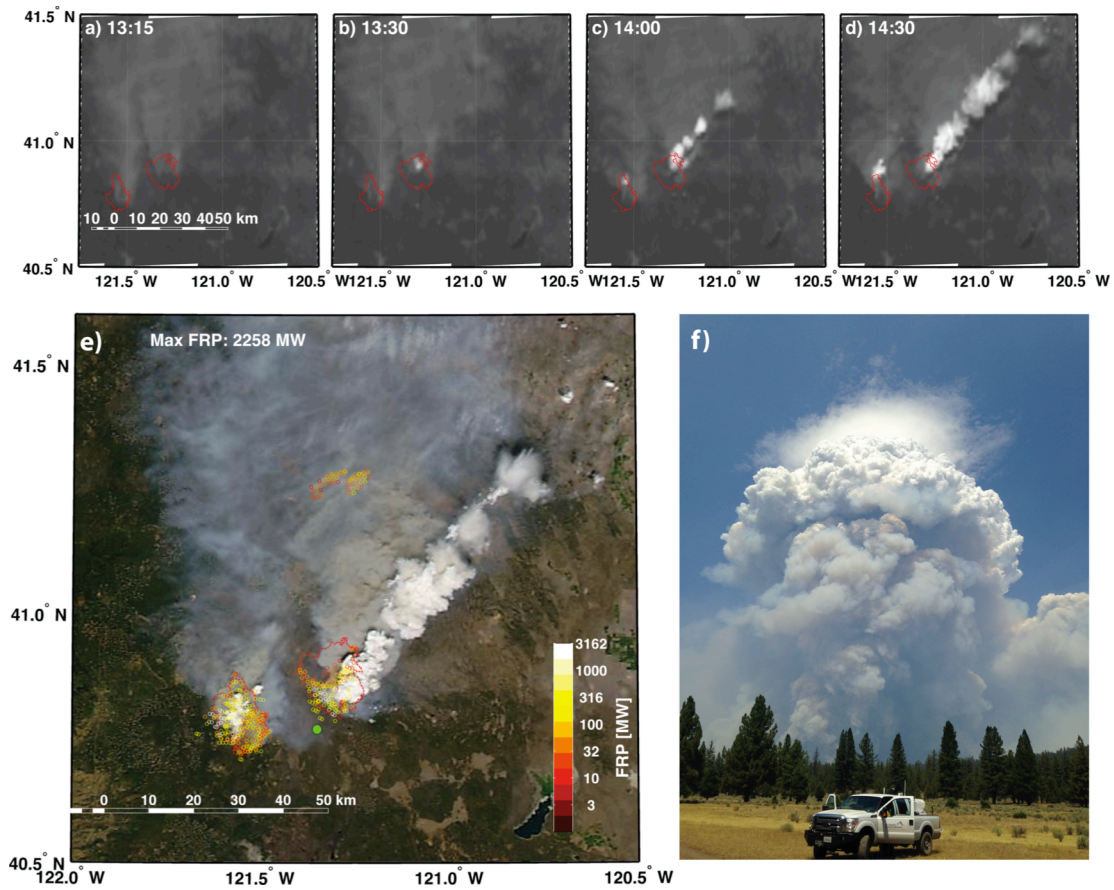


Fig. 4. Overview of the pyrocumulus initiation and growth on 2 August 2014. (a-d) GOES-15 visible imagery showing the pyroCu initiation. (e) MODIS-Aqua visible image at 1426 PDT along with fire-radiative power (FRP, colored circles). (f) Photograph of the lidar vantage point and the windward edge of convective column and pyroCu at 1401 PDT. The truck location is indicated in as a green dot in panel (e). The fire perimeters are as in Fig. 3.

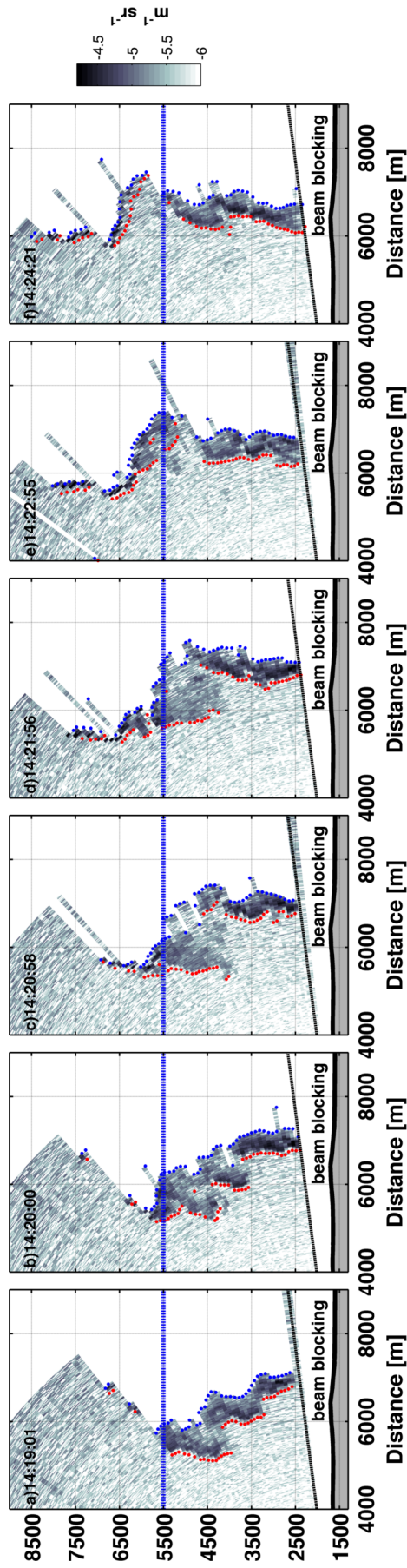


Fig. 5. Plume rise sequence recorded by successive lidar RHI scans from 1419 to 1424 PDT. The displayed data are the logarithmic attenuated backscatter coefficient ($\text{m}^{-1} \text{sr}^{-1}$). The plume edge and attenuation points are shown in red and blue dots, respectively. Data beyond the attenuation point is shadowed (blank points). The dashed blue line indicates the condensation level.

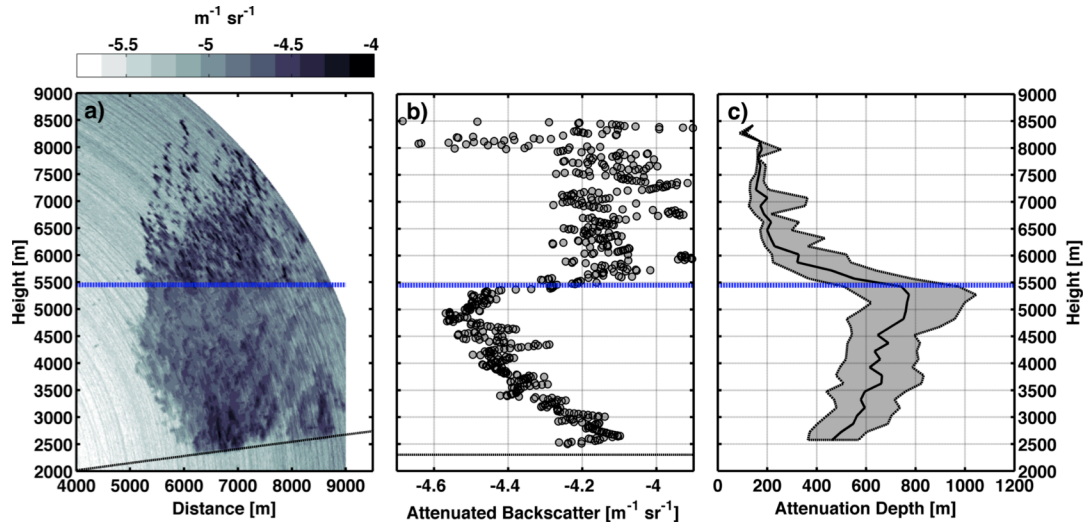


Fig. 6. Statistical analysis of lidar data between 1350 and 1502 PDT showing the plume condensation level. (a) Maximum backscatter as a function of height and distance. (b) Maximum backscatter as a function of height only. (c) 5, 50 and 95th percentiles of the attenuation depth as a function of height. The dashed blue line indicates the inferred condensation level.

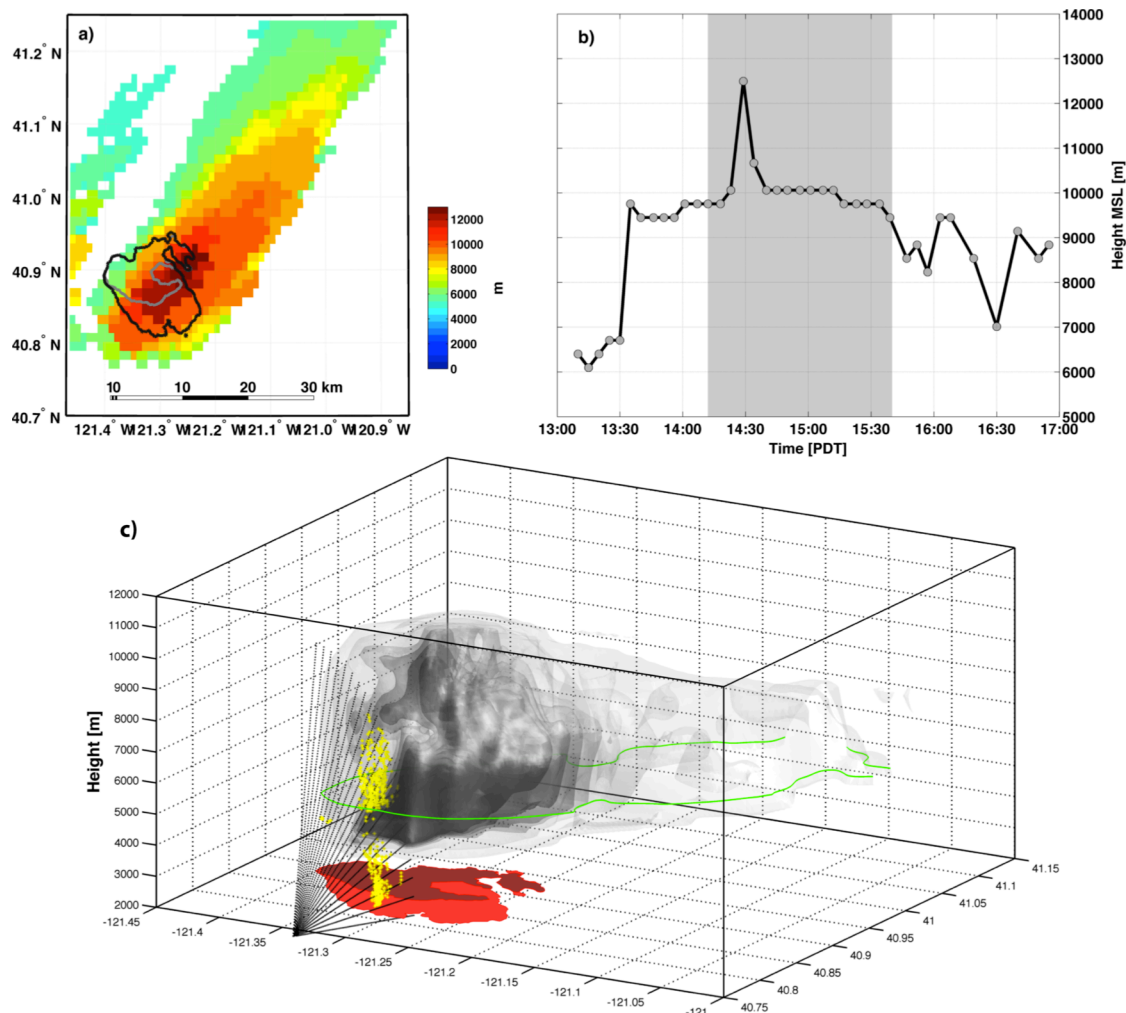
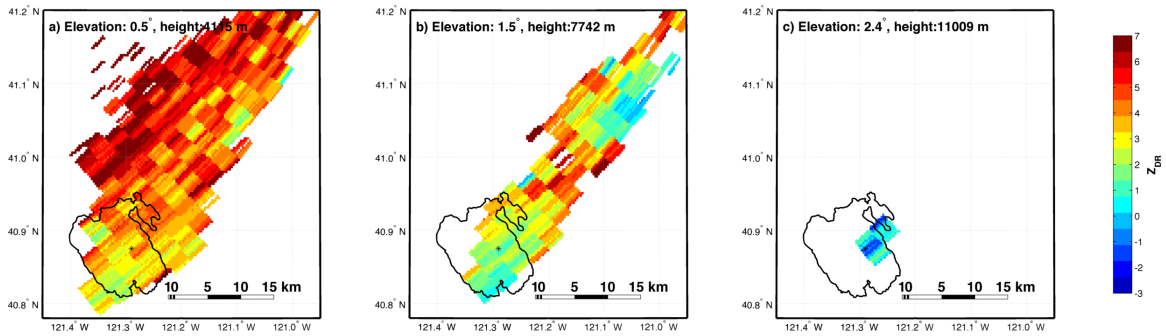


Fig. 7. Radar analysis of the Bald Fire convective column. (a) Maximum echo top heights (color shading) during the Bald Fire along with the NIROPS fire perimeters (gray and black contours). (b) Time series of the maximum echo tops heights. Gray shading shows the period of lidar observations. (c) Volume rendering of the Bald Fire plume at 1429 PDT. Reflectivity isosurfaces are displayed at 30, 28, 26, 24, 22 and 18 dbZ. The lidar scan path and plume detections are shown in black and yellow dots, respectively. Fire perimeters are shown in red shading. The lidar derived condensation level is indicated by the green contour.



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817 Fig. 8. Differential reflectivity (Z_{dr}) analysis of the Bald Fire plume at 1429 PDT. (a-c)
818 Z_{dr} at the 0.5, 1.5, and 2.4 deg. elevation sweeps. The black contour shows the fire
819 perimeter on 2 August.

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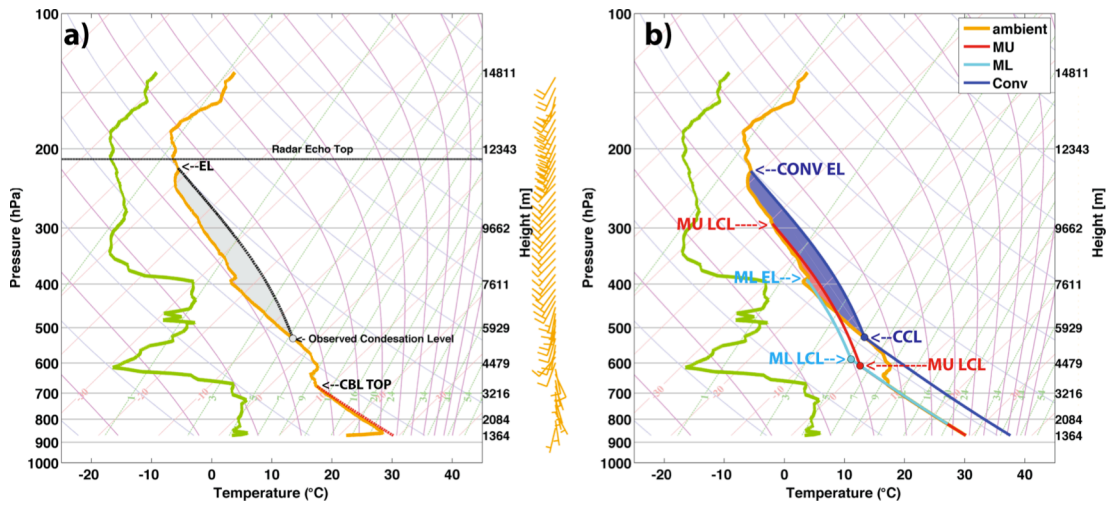
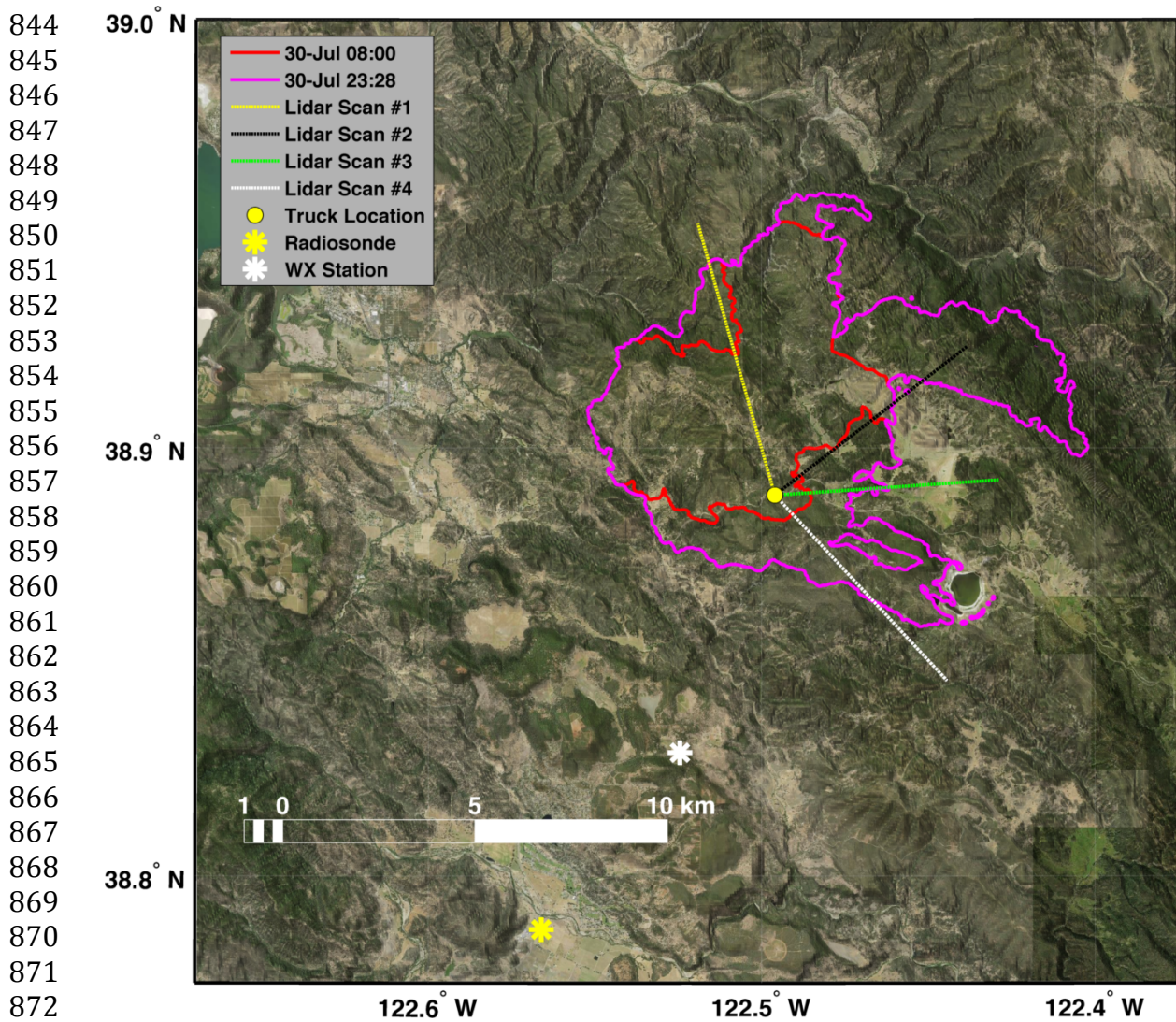


Fig. 9. Thermodynamic analysis of the ambient environment and plume parcels. (a) Observed sounding from 2 August 2014, 2100 PDT showing the adjusted boundary layer profile (dashed red line), the lidar derived condensation level (gray circle), the moist-adiabatic ascent from the condensation level, equilibrium level, and the radar derived echo tops. (b) Analysis of lifted parcels, showing the most unstable (MU), mixed-layer (ML), and convective (CONV) parcel trajectories. The condensation and equilibrium levels for each parcel are shown, and their CAPE is shaded.



873 Fig. 10. Rocky Fire progression map for 30 July 2015. The fire perimeters are from the
 874 US National Forest Service National Infrared Operations (NIROPS) flights. The
 875 background is a satellite image draped over the terrain, which is indicated with hill
 876 shading. Also shown are the truck location (yellow dot), lidar scan paths (colored lines),
 877 radiosonde location (yellow star), and weather station location (white star).
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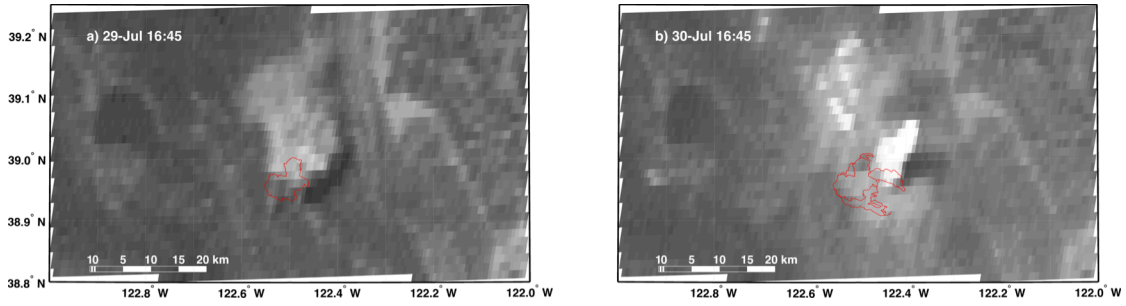


Fig. 11. Visible satellite (GOES-15) images showing the difference in the Rocky Fire plume between 1645 PDT on 29 and 30 July. The data show a pyroCu tower on 30 July that is absent on 29 July.

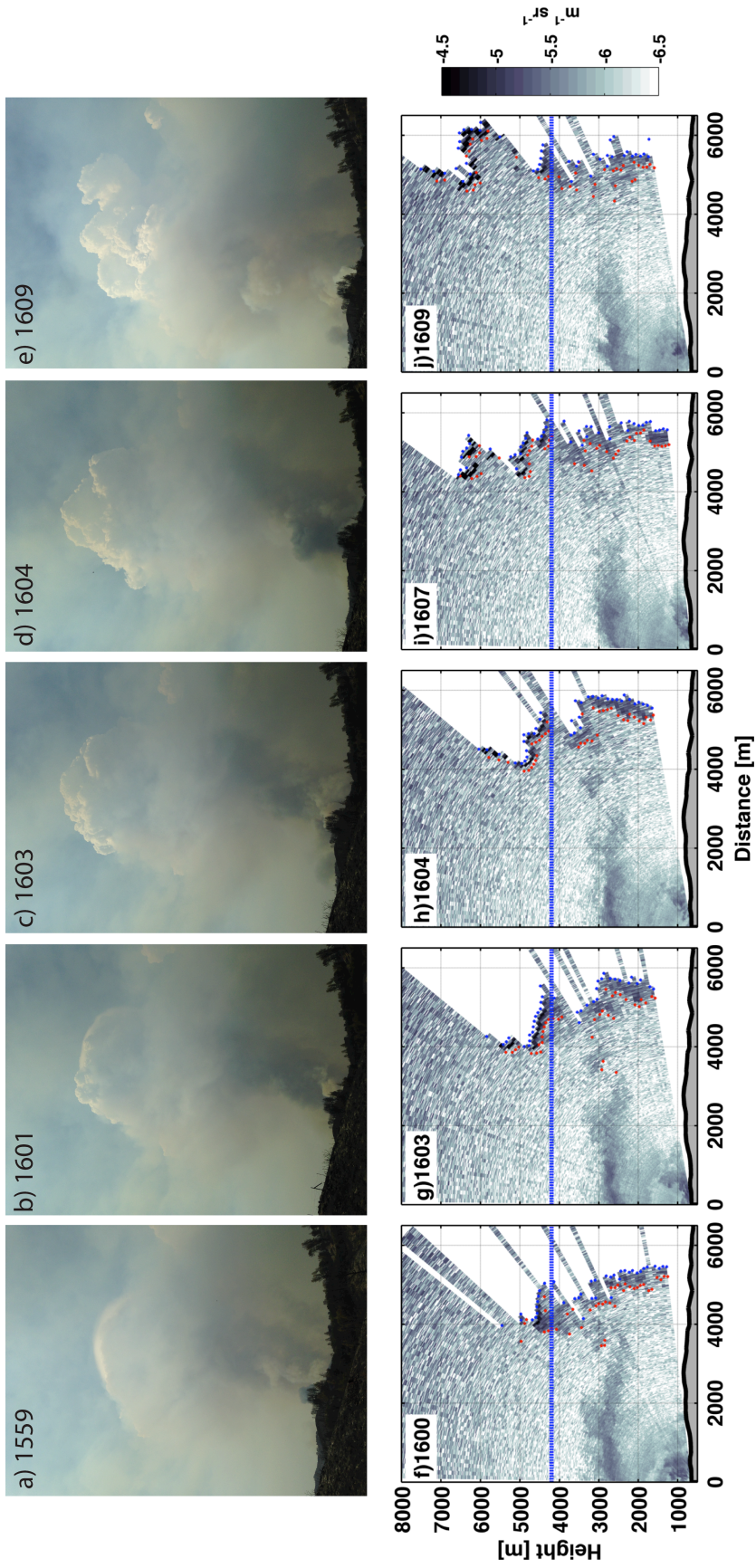


Fig 12. Pyrocumulus development from 1559-1609 PDT on 30 July 2015. (a-e) Photographs of the plume rise and pyroCu development. (f-j) Lidar backscatter showing the onset of condensation and subsequent cloud growth. The dashed blue lines shows the lidar derived condensation level.

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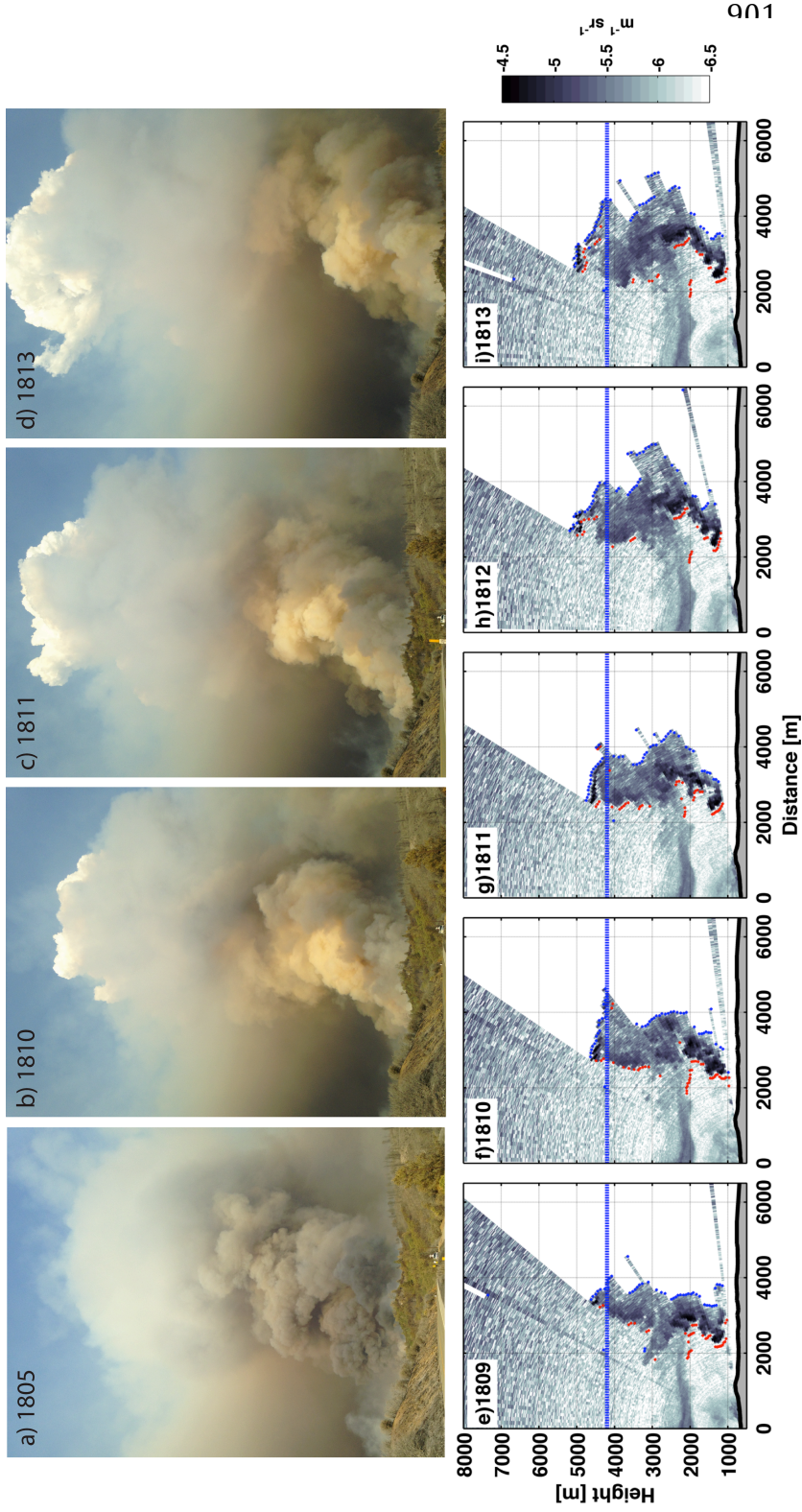
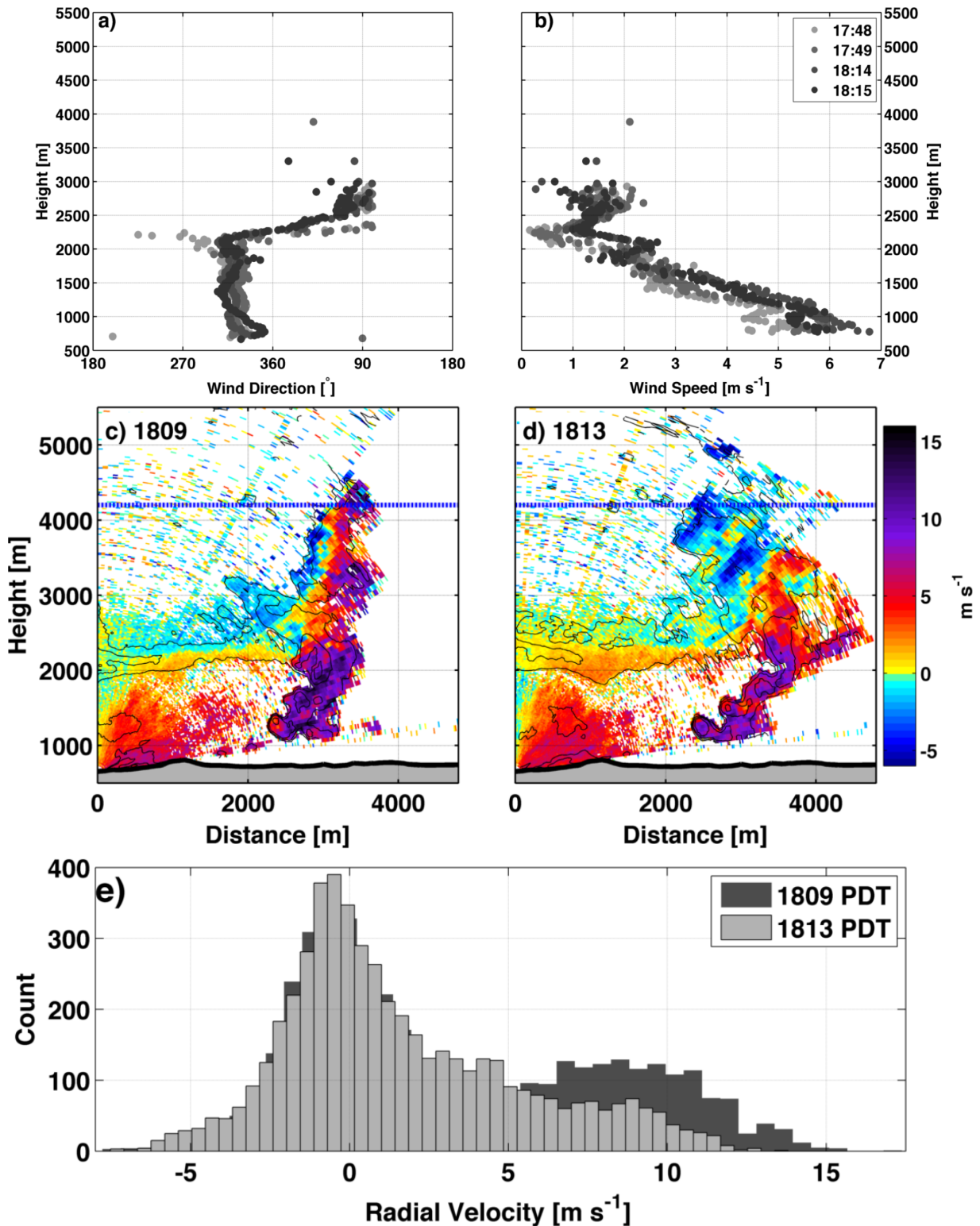
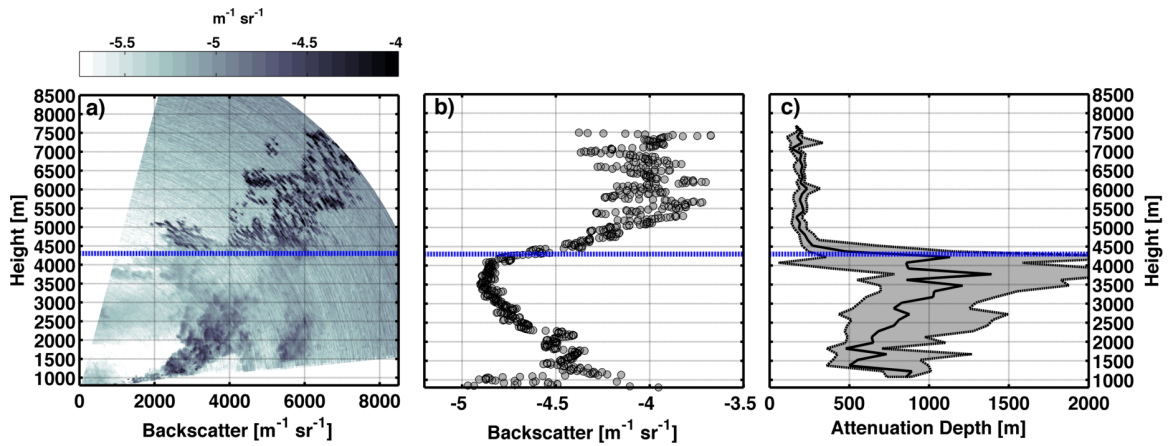


Fig 13. Pyrocumulus development from 1805-1813 PDT on 30 July 2015. (a-d) Photographs of the plume rise and pyroCu development. (e-i) Lidar backscatter showing the onset of condensation and subsequent cloud growth. The dashed blue lines show the lidar derived condensation level.



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Fig. 14. Analysis of the interaction of the plume with the ambient wind. (a-b) VAD derived profiles of the ambient wind speed and direction. (c-d). Radial velocities during the plume rise and smoke backscatter (black contours). (e) Comparative histogram of radial velocities at 1809 and 1813 PDT.



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Fig. 15. Statistical analysis of lidar data between 1545 and 2008 PDT on 30 July 2015. (1) Maximum backscatter as a function of height and distance. (b) Maximum backscatter as a function of height only. (C) Attenuation depth as a function of height. The dashed blue line indicates the inferred condensation level.

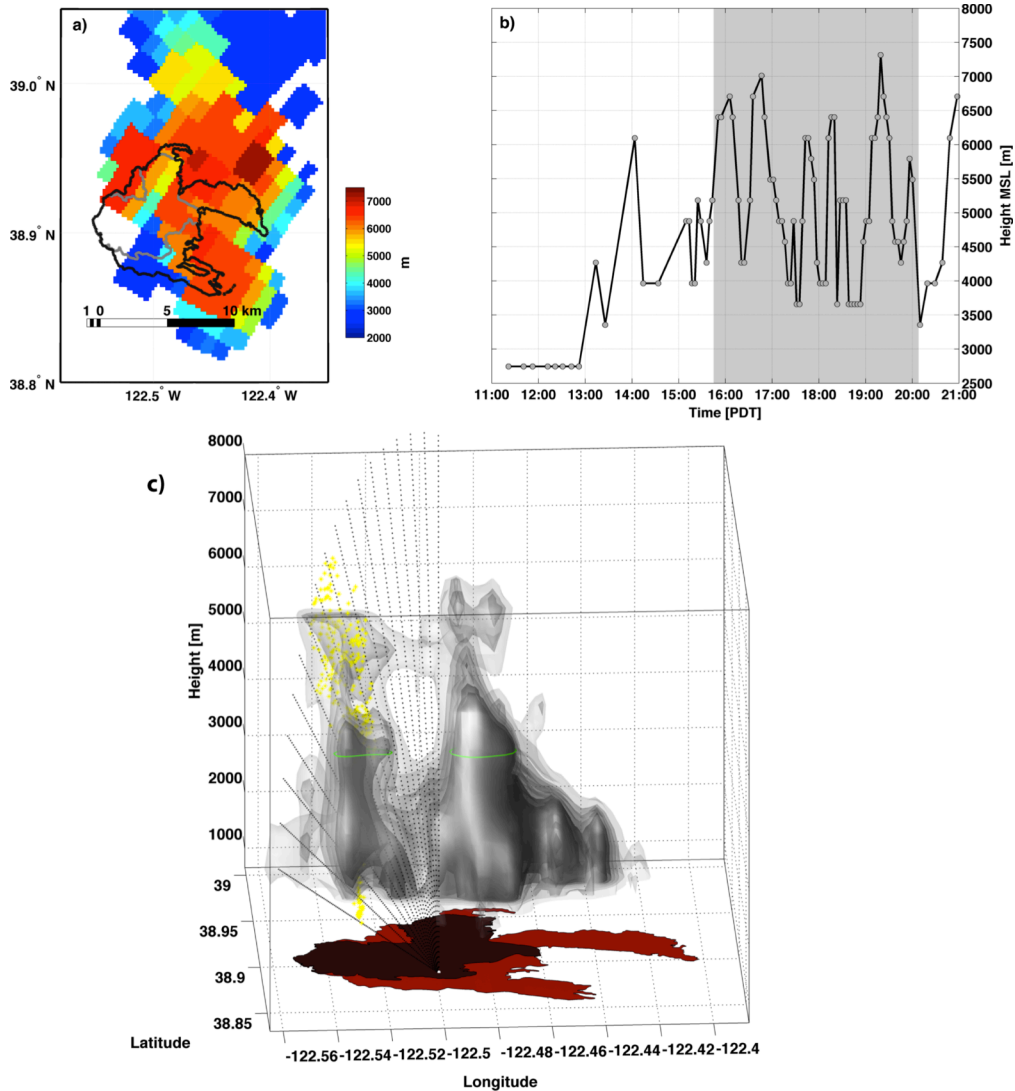


Fig. 16. Radar analysis of the Rocky Fire convective column. (a) Maximum echo top heights (color shading) along with the NIROPS fire perimeters (gray and black contours). (b) Time series of the maximum echo tops heights. Gray shading shows the period of lidar observations. (c) Volume rendering of the Rocky Fire plume at 1609 PDT. Reflectivity isosurfaces are displayed at 30, 28, 26, 24, 22 and 18 dbZ. The lidar scan path and plume detections are shown in black and yellow dots, respectively. Fire perimeters are shown in red shading.

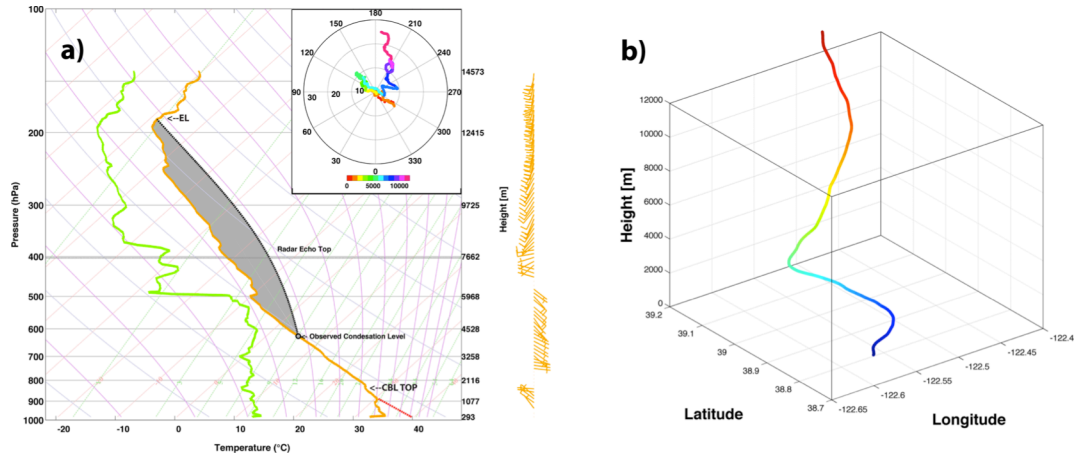
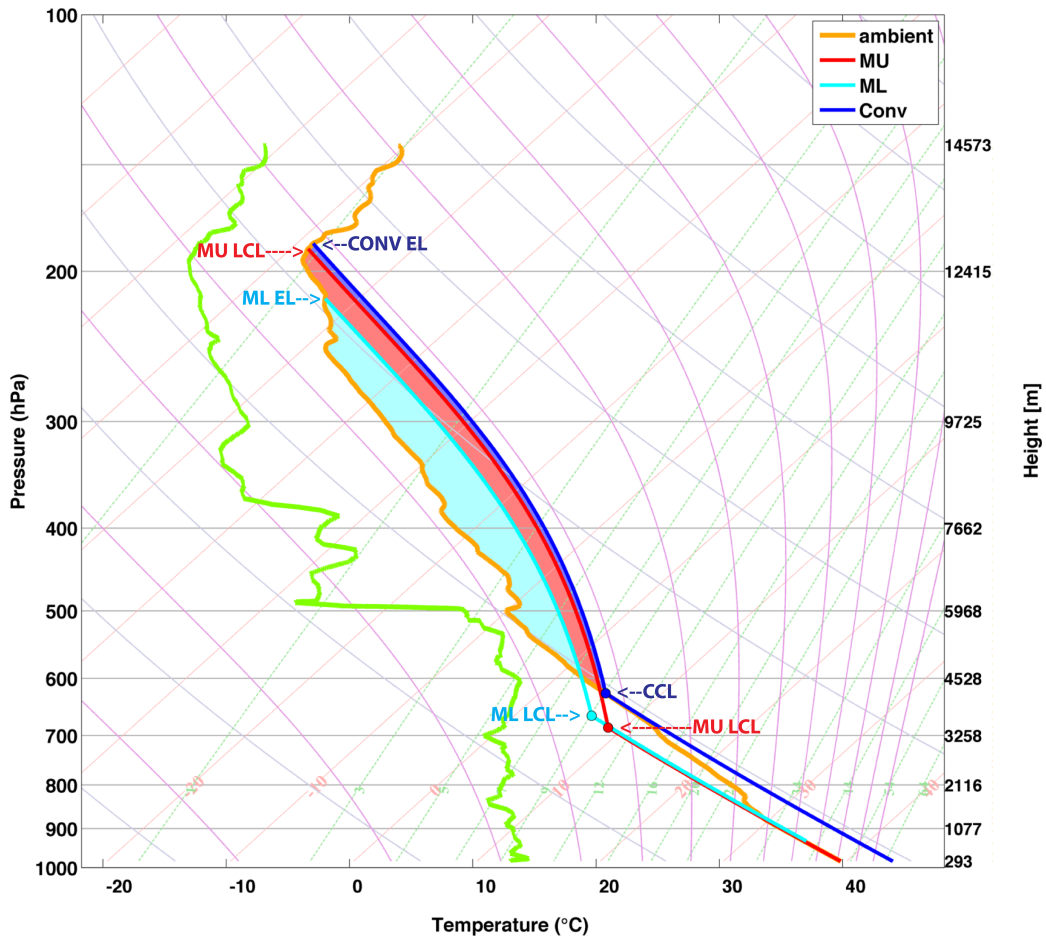


Fig. 17. Thermodynamic analysis of the ambient environment during the Rocky Fire. (a) Observed sounding from 2105 PDT showing the adjusted boundary layer profile (dashed red line), the lidar derived condensation level (gray circle), the moist-adiabatic ascent from the condensation level, and the radar derived echo tops. The inset is a hodograph and the wind barbs on the right indicate how the wind speed and direction change with height. (b) Balloon ascent path showing the affect of wind shear on a buoyant parcel.



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Fig. 18 Analysis of lifted parcels, showing the most unstable (MU), mixed-layer (ML), and convective (Conv) parcel trajectories. The condensation levels and CAPE for each parcel is described in the text.

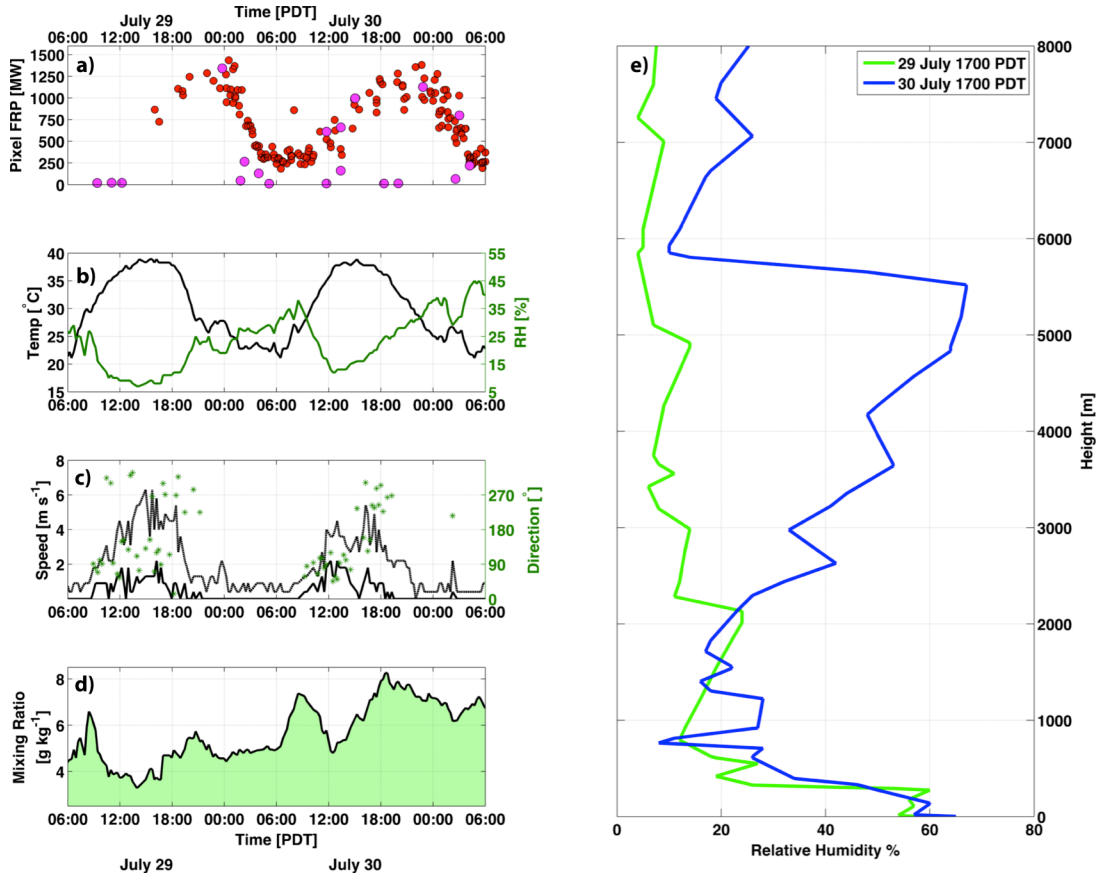


Fig 19. Analysis of the environmental conditions on 29 and 30 July. (a) Fire radiative power from GOES-15 (red dots) and MODIS (purple dots). (b) Temperature (black line) and relative humidity (green line) time series. (c) wind speed (black line), gust (dashed black line) and direction (green stars). (d) Surface mixing ratio. (e) Comparison of the vertical profile of relative humidity from the KOAK sounding at 1700 PDT on 29 and 30 July. The location of the weather station is shown in Fig. 10.