On the characteristics of aerosol indirect effect based on dynamic regimes in global climate models

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28 Abstract

Aerosol-cloud interactions continue to constitute a major source of uncertainty for 29 the estimate of climate radiative forcing. The variation of aerosol indirect effects (AIE) 30 in climate models is investigated across different dynamical regimes, determined by 31 monthly mean 500 hPa vertical pressure velocity (ω_{500}), lower-tropospheric stability 32 (LTS) and large-scale surface precipitation rate derived from several global climate 33 models (GCMs), with a focus on liquid water path (LWP) response to cloud 34 condensation nuclei (CCN) concentrations. The LWP sensitivity to aerosol 35 perturbation within dynamic regimes is found to exhibit a large spread among these 36 GCMs. It is in regimes of strong large-scale ascent ($\omega_{500} < -25$ hPa/d) and low clouds 37 (stratocumulus and trade wind cumulus) where the models differ most. Shortwave 38 aerosol indirect forcing is also found to differ significantly among different regimes. 39 Shortwave aerosol indirect forcing in ascending regimes is close to that in subsidence 40 regimes, which indicates that regimes with strong large-scale ascent are as important 41

as stratocumulus regimes in studying AIE. It is further shown that shortwave aerosol
indirect forcing over regions with high monthly large-scale surface precipitation rate (>
0.1 mm/d) contributes the most to the total aerosol indirect forcing (from 64% to
nearly 100%). Results show that the uncertainty in AIE is even larger within specific
dynamical regimes compared to the uncertainty in its global mean values, pointing to
the need to reduce the uncertainty in AIE in different dynamical regimes.

- 48 Key words: aerosol indirect effects, dynamic regimes, aerosol-cloud interactions
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50 **1 Introduction**

By scattering and absorbing sunlight, aerosol particles can modify the solar 51 52 radiation reaching the earth system, which is termed the direct effect. The direct radiative effect of anthropogenic aerosols combined with subsequent rapid 53 adjustments of the surface energy budget, atmospheric state variables, and cloudiness 54 to aerosol radiative effects is referred as Effective Radiative Forcing from 55 aerosol-radiation interactions (ERFari) (Boucher et. al, 2013). Apart from ERFari, 56 aerosols can also alter the Earth's radiation balance via interactions with clouds, such 57 as effects on cloud albedo and subsequent changes to the cloud lifetime and 58 thermodynamics as rapid adjustments, known as the aerosol indirect effect(s) (AIE). 59 These radiative effects are called Effective Radiative Forcing from aerosol-cloud 60 61 interactions (ERFaci) (Boucher et. al, 2013).

For liquid clouds, there are two principal ways through which aerosols interact with them in AIE. First, an increase in cloud condensation nuclei (CCN)

concentration from anthropogenic aerosols leads to smaller cloud droplet sizes 64 assuming constant liquid water content. The increased number but decreased droplet 65 sizes in turn increase cloud albedo due to more efficient backscattering. This is called 66 the cloud albedo effect or the first AIE, also known as the Twomey effect (Twomey, 67 1977). Moreover, the smaller cloud droplet sizes are hypothesized to lead to decreases 68 69 in precipitation efficiency, which may further alter cloud liquid water path (LWP) and cloud lifetime (Albrecht, 1989). These adjustments are also referred to as the cloud 70 lifetime effect or the second AIE. It is worth noting that delaying the onset of 71 precipitation may further modify latent heating profiles, which could lead to the 72 invigoration of convective clouds (e.g., Andreae et al., 2004; Rosenfeld et al., 2008). 73 There are also adjustments on mixed-phase and ice clouds (e.g., Storelymo et al., 74 2008; Lohmann and Hoose, 2009; Liu et al. 2012; Storelvmo et al., 2008; Wang et al., 75 2014). The focus of this study is on liquid cloud response to aerosol perturbation, 76 77 primarily from large-scale clouds.

AIE could be large enough to offset much of the global warming induced by anthropogenic greenhouse gases, yet its magnitude is still very uncertain (IPCC, 2013). The uncertainty in the cloud lifetime effect of aerosols is particularly large.

The complexity of microphysical-dynamical-radiative feedbacks involved in the cloud lifetime effect has been noted in previous studies. Conventional theory regarding the cloud lifetime effect suggests that higher CCN concentration slows down precipitation formation and hence leads to more LWP (Albrecht, 1989). However, this theory is inconsistent with some observations (Coakley and Walsh,

2002; Kaufman et al., 2005; Matsui et al., 2006; Chen et al., 2014) and large eddy
simulations (LESs) (e.g., Ackerman et al., 2004; Lu and Seinfeld, 2005; Wang and
Feingold, 2009b) that found either increase or decrease in LWP in responses to
increases in CCN concentration.

Further modeling studies (e.g., Ackerman et al., 2004; Stevens and Feingold, 90 2009; Guo et al., 2011) suggest that cloud top entrainment plays a critical role as a 91 dynamic feedback, to balance LWP and modify the lifetime of boundary layer clouds. 92 Ackerman et al. (2004) found that an increase in droplet number concentration (N_d) 93 reduces cloud water sedimentation while accelerating the cloud-top entrainment rate, 94 which makes the humidity of air overlying the boundary layer, wet or dry, critically 95 important in determining the response of LWP. When surface precipitation is weak 96 $(<0.1 \text{ mm day}^{-1})$ and the overlying air is dry, LWP decreases in response to increasing 97 aerosol. They showed that the entrainment rate was reduced by decreasing available 98 boundary-layer turbulence kinetic energy (TKE). However, Bretherton et al. (2007) 99 found that TKE remained unchanged and changes in entrainment rate are mainly 100 caused by reduced evaporative cooling from removing out liquid water. LES studies 101 (e.g., Wang and Feingold, 2009a) with a large model domain that is able to resolve 102 mesoscale circulations (on the order of ten kilometers) in marine stratocumulus 103 showed that aerosols can shift cloud regimes through their impact on precipitation and 104 associated dynamical feedbacks. This can represent a more significant impact on 105 cloud radiative forcing than the conventional AIE. 106

Many state-of-the-art global climate models (GCMs) appear to overestimate AIE 107 when compared with satellite observations (e.g., Quaas et al., 2009; Wang et al., 108 2012), despite of some uncertainties in satellite derived estimates (e.g., Penner et al., 109 2011; Gryspeerdt et al., 2014a; Gryspeerdt et al., 2014b). The multi-scale interactions 110 between clouds, aerosols and large-scale dynamics (Stevens and Feingold, 2009; 111 Wang et al., 2011; Ma et al., 2015) and complex microphysical processes (e.g., 112 Bretherton et al., 2007; Gettelman et al., 2013) cause uncertainties in estimating AIE 113 by GCMs. One possible source of overestimation of AIE is their inability to reproduce 114 negative LWP responses to aerosol perturbations, which are found in some 115 observations and LES studies, partly because they do not explicitly simulate the 116 droplet size effect on the entrainment process and on sub-grid cloud organizations 117 associated with changes in precipitation. Guo et al. (2011) found that this effect could 118 be captured through applying a parameterization based on multi-variate probability 119 density functions with dynamics (MVD PDFs) in single-column simulations. They 120 found decreased LWP in response to increasing aerosols concentration and suggested 121 that the implementation of MVD PDFs in GCMs may help lower the magnitude of the 122 simulated AIE. A negative correlation between LWP and aerosol loading was further 123 found for clouds with weak precipitation and dry air above the PBL in a subsequent 124 global model study (Guo et al., 2015). 125

Another likely source for the overestimation of cloud lifetime effects in GCMs is
the treatment of cloud microphysics (Penner et al., 2006; Posselt and Lohmann, 2009;
Wang et al., 2012). In warm clouds, cloud microphysical processes are dominated by

autoconversion and accretion in bulk microphysics schemes (Gettelman et al., 2013). 129 Since autoconversion acts as a sink of LWP, it is crucial in the formation of 130 precipitation, thus plays an important role in determining the cloud lifetime effect. 131 The autoconversion rate is directly dependent on droplet number concentration (N_d) 132 while the accretion rate is only weakly dependent on N_d (Khairoutdinov and Kogan, 133 2000; Gettelman et al., 2013). Furthermore, the ratio of the autoconversion rate to the 134 large-scale surface precipitation rate is found to be strongly correlated with the LWP 135 136 response to anthropogenic aerosol perturbations (e.g., Wang et al., 2012). Posselt and Lohmann (2009) suggested this ratio is related to the rain scheme adopted in GCMs. 137 They showed that the adoption of different rain schemes (prognostic vs. diagnostic) in 138 a GCM leads to a different LWP response to aerosol perturbations. A prognostic rain 139 scheme can shift the importance of (warm) rain production from autoconversion 140 process to the accretion process and therefore reduces the AIE (Posselt and Lohmann, 141 2009; Gettelman et al., 2015). However, Hill et al. (2015) shows that adding 142 prognostic rain scheme alone still cannot reduce the spread of susceptibility of 143 precipitation among different cloud microphysics parameterizations and further shows 144 that increasing the complexity of the rain representation to double-moment 145 significantly reduces the spread of precipitation sensitivity and improves overall 146 consistency between bulk and bin schemes. 147

Previous studies are mostly confined to global averages (e.g Quaas et al., 2009;
Wang et al., 2012) or a specific dynamic environment (e.g., Bretherton et al., 2007;
Guo et al., 2011). However, aerosols, clouds, precipitation distributions and

dynamical feedbacks are all related to the prevailing meteorological environment (Stevens and Feingold, 2009). Clouds are sensitive to changes in dynamical regimes, which can be defined by large-scale circulations, thermodynamic structure and meteorological backgrounds (Bony et al., 2004). Gryspeerdt and Stier (2012) and Gryspeerdt et al. (2014c) used satellite data and found that the characteristics of aerosol cloud-albedo effect (droplet number sensitivity) vary with cloud regimes and pointed out the importance of regime-based studies of aerosol-cloud interactions.

In this study, we investigate how AIE in several GCMs varies under different 158 dynamical regimes over global oceans (60°S-60°N), with a focus on cloud lifetime 159 effects of aerosols (2nd AIE). We note that the term "cloud lifetime effects" can be 160 somehow misleading, since aerosol effects on cloud liquid water may have little to do 161 with cloud lifetime per se (e.g., Small et al., 2009). Nevertheless, this term is still used 162 in some occasions in this paper for convenience. The paper is organized as follows. 163 Methods and models are described in Section 2, and results and discussions are 164 presented in Section 3. The paper concludes with the summary in Section 4. 165

- 166 2 Methodology and models
- 167 The response of LWP to aerosol perturbations is defined as
- 168 $\lambda = d \ln LWP / d \ln CCN.$

As simulated LWP and CCN can be quite different among GCMs, the logarithmic form of LWP and CCN is adopted in the λ formula. λ is a metric to quantitatively measure cloud lifetime effect of aerosols in models. It is directly calculated as the

relative change of monthly mean LWP from pre-industrial (PI) to present day (PD) 172 divided by the relative change of CCN. Here dlnLWP=(LWPPD-LWPPI)/LWPPI and 173 dlnCCN=(CCN_{PD}-CCN_{PI})/CCN_{PI}, where LWP_{PD} and LWP_{PI} are LWP in PD and PI, 174 respectively, while CCN_{PD} and CCN_{PI} are CCN in PD and PI, respectively. This 175 parameter was used by Wang et al. (2012) to constrain the cloud lifetime effects of 176 aerosols over global oceans using precipitation frequency susceptibility (S_{pop}) derived 177 from A-Train satellite observations. Lebo and Feingold (2014) examined the 178 relationship between λ and S_{pop} to aerosol perturbations for stratocumulus and 179 trade-wind cumulus simulated by LES and found that λ may increase in marine 180 stratocumulus while decrease in the case of trade-wind cumulus in response to 181 increasing S_{pop} , suggesting a cloud regime dependence of this relationship. Note that λ 182 allows some feedbacks, for example cloud effects on CCN. 183

Dynamical regimes can be defined by environment characteristics such as 184 large-scale vertical pressure velocity (e.g., Bony and Dufresne, 2005) and 185 lower-tropospheric stability (LTS, defined as the difference in potential temperature 186 between 700hPa and the surface, θ_{700hPa} - $\theta_{surface}$) (e.g., Medeiros and Stevens, 2011). 187 Medeiros and Stevens (2011) noted that low clouds and deep convective clouds could 188 be separated by ω_{500} while different low cloud types under large-scale subsidence can 189 only be depicted by using LTS. In this study the monthly-averaged vertical pressure 190 velocity (ω) in the mid-troposphere (defined as at 500hPa) is used as a proxy for 191 large-scale motions (Bony and Dufresne, 2005). Note that ω_{500} with positive (negative) 192 193 value means descending (ascending) motions. We decompose global (60°S~60°N)

large-scale circulations over ocean as a group of dynamical regimes (equally sampled) 194 by ω_{500} (and LTS). Ascending regimes and descending regimes are defined by ω_{500} 195 and descending regimes are further divided into stratocumulus, transitional clouds and 196 trade wind cumulus regimes by LTS. This method is straight-forward to apply to 197 GCM results and gives us a direct view of the relationship between clouds and their 198 199 favorable large-scale environmental characteristics. Note however that the use of monthly means may obscure some details in the microphysical relationships, 200 especially where the variability of cloud properties is high. 201

Since vertical pressure velocity is used as a major criterion here, dynamic regimes 202 generally follow the features of vertical pressure velocity distributions. Descending 203 regimes are mostly located at subtropical regions and western coasts of continents, 204 while ascending regimes locates around ITCZ and northern Pacific where storm tracks 205 prevail. The seasonal evolution of dynamic regimes follows seasonal changes in the 206 major meteorological systems. For example, ascending regimes move north/south as 207 ITCZ move north/south and descending regimes move accompanying with subtropical 208 high move. The characteristics of dynamic and thermodynamic regimes were 209 discussed in detail in Bony et al. (2004). 210

As the perturbations in cloud radiative forcing from anthropogenic aerosols (indirect effect) are typically on the order of 1 W m⁻², which is small compared to the cloud radiative forcing (shortwave radiative effect of ~-47 W m⁻² and longwave radiative effect of ~27 W m⁻²) (Boucher et al., 2013), long integrations are required to produce statistically significant results. The Newtonian relaxation method (nudging)

provides a way to estimate AIE within a relatively short integration time, while giving 216 statistically significant results (Lohmann and Hoose, 2009; Kooperman et al., 2012). 217 Nudging here refers to the method of adding a forcing to the prognostic model 218 equations, determined by the difference between a model-computed value and a 219 prescribed value at the same time and model grid-cell, to constrain the model results 220 221 with prescribed atmospheric conditions. Kooperman et al. (2012) implemented nudging to constrain PD and PI simulations toward identical meteorological fields and 222 found that the use of nudging provided a more stable estimate of AIE in shorter 223 simulations and increased the statistical significance of the anthropogenic aerosol 224 perturbation signal. All simulations used in this study were nudged toward reanalysis 225 winds (year 2006 to 2010) provided by operational forecast centers. Some simulations 226 were further nudged toward reanalysis temperature, but this was discouraged because 227 it might affect the moist convection activities simulated in the model (Zhang et al., 228 2014). All models were driven by the same IPCC aerosol emissions for years 1850 229 and 2000 (Lamarque et al., 2010) and 5-year simulations were performed in each case 230 (PI and PD). Sea surface temperature, sea-ice extent and greenhouse gas 231 concentrations are prescribed to climatological values in all simulations. Monthly data 232 were then obtained by averaging over the 5-year integration period. 233

Only ω_{500} in PD runs is used to derive dynamical regimes and then these dynamical regimes are applied to PI simulations as well, with the assumption that ω_{500} does not change much from PI to PD. This assumption is reasonable as both PD and

237 PI runs were nudged toward the reanalysis data here, which ensures ω_{500} is very 238 similar between PD and PI.

239 A total of ten aerosol-climate models participated in this study. This includes five versions of Community Atmosphere Model (CAM) 5.3, and two versions of 240 SPRINTARS. These models show large differences in their aerosol and cloud 241 treatments. For example, while most models (CAM5, CAM5-PNNL, CAM5-MG2, 242 CAM5-CLUBB, CAM5-CLUBB-MG2, ECHAM6-HAM2, and SPRINTARS-KK) 243 use the autoconversion scheme from Khairoutdinov and Kogan (2000, hereafter KK), 244 autoconversion rate in ModelE-TOMAS is independent of cloud droplet number 245 concentration and the Berry scheme (Berry, 1967) is used for SPRINTARS. Most 246 models use diagnostic rain schemes, while an updated Morrison and Gettelman 247 microphysics scheme with a prognostic rain scheme (MG2) (Gettelman et al., 2015) is 248 adopted in CAM5-MG2 and CAM5-MG2-CLUBB. HadGEM3-UKCA also adopts a 249 prognostic rain scheme (Abel and Boutle, 2012). While most models only account for 250 aerosol effects large-scale stratiform clouds. CAM5-CLUBB 251 on and CAM5-CLUBB-MG2 use a higher-order turbulence closure (CLUBB) to unify the 252 treatment of boundary layer turbulence, stratiform clouds and shallow convection, and 253 therefore include aerosol effects on shallow convection (Bogenschutz et al., 2013). A 254 brief description of each model is provided in the Appendix. 255

256 **3 Results**

257 **3.1 Annual mean**

258 We first examine the annual climatology in different simulations to get an overall

picture of the general differences/similarities among these models (details within dynamic regimes are examined in section 3.2). All of the simulations reproduce the general pattern of large-scale circulations (ω_{500}): strong ascending motions within the inter-tropical convergence zone (ITCZ) and subsidence dominating subtropical eastern ocean regions (not shown). The similar patterns of ω_{500} (due to nudging) in these simulations ensure that dynamic regimes defined by ω_{500} do not vary much between models.

Table 1 lists the types of clouds included in LWP and rain analyzed in this study 266 and the different rain scheme (prognostic or diagnostic) in these 10 GCM simulations. 267 Table 2 lists global annual means of aerosol, precipitation and cloud parameters in PD 268 simulations and λ for each model. Note that all versions of CAM5 calculate LWP only 269 for large-scale clouds while SPRINTARS, SPRINTARS-KK and HadGEM3-UKCA 270 also count LWP from convective clouds. As for ModelE2-TOMAS, LWP includes 271 stratiform anvil clouds that formed from convective detrainment of water vapor and 272 ice. ECHAM6-HAM2 also includes the contribution of convective detrainment of 273 liquid water and ice to stratiform clouds Also note that CAM5 models with CLUBB 274 include LWP in the shallow convective regimes, which partly explains why these 275 models produce more LWP than their corresponding CAM5 models without CLUBB 276 (Table 2). 277

There are large differences among global LWP annual means. CAM5-MG2 has the lowest LWP among these simulations (30.0 g m⁻²). The LWP means over oceans are 31.1 g m⁻², 39.4 g m⁻² and 35.2 g m⁻² in CAM5, CAM5-PNNL and

CAM5-CLUBB, respectively. HadGEM3-UKCA simulates higher LWP (57.1 g m⁻²) 281 than all versions of CAM5. LWPs in ModelE2-TOMAS (80.4 g m⁻²) and 282 ECHAM6-HAM2 (84.6 g m⁻²) are greater than the aforementioned GCMs, but less 283 than in SPRINTARS and SPRINTARS-KK (139.1 g m⁻² and 98.9 g m⁻² respectively) 284 which include LWP from convective clouds. Even though CAM5-CLUBB simulates a 285 higher LWP in storm track regions and ECHAM6-HAM2 produces much more LWP 286 associated with deep convection in the ITCZ, all models here display reasonable 287 patterns of global LWP distributions (not shown). 288

The differences in CCN (at 0.1% supersaturation) among these simulations are 289 not as large as the differences in LWP (Table 2). The global annual mean CCN in 290 CAM5-PNNL, which has a different treatment of wet scavenging processes (Wang et 291 al., 2013), is slightly larger than the one in other versions of CAM5. CCN 292 concentrations simulated by CAM5-PNNL, ECHAM6-HAM2 and ModelE2-TOMAS 293 are largest among these simulations and are more than twice those simulated by 294 SPRINTARS, SPRINTARS-KK and HadGEM3-UKCA, which are the lowest. Since 295 these models are using same emissions, differences of CCN between the models are 296 mainly due to different aerosol lifetime between models. 297

The LWP response to aerosol perturbations, λ , in ECHAM6-HAM2 (0.19) is close to those derived from three CAM5 configurations (0.20 in CAM5, 0.19 in CAM5-PNNL and 0.25 in CAM5-CLUBB). Notice that λ in CAM5-MG2 and CAM5-CLUBB-MG2 is larger than that in CAM5 and CAM5-CLUBB, respectively, which indicates that the changes of LWP in the models, using the MG2 scheme, are

more sensitive to the aerosol perturbations. LWP is much less sensitive to the changes 303 of CCN in SPRINTARS and SPRINTARS-KK with λ of 0.01 and 0.04 respectively. λ 304 is also small in HadGEM3-UKCA (0.03) due to the large relative increase of CCN 305 while small relative increase of LWP. Since the aerosol effect on precipitation 306 formation is turned off in ModelE2-TOMAS (its autoconversion parameterization is 307 308 not a function of N_d), LWP barely responds to the increase of CCN (λ is -0.001). The variation in λ closely follows that of the relative enhancement of LWP (dlnLWP), as 309 the variation of the relative enhancement of CCN (dlnCCN) among the simulations is 310 generally much smaller than that of dlnLWP. 311

We should note that large differences in CCN shown in Table 2 do not 312 necessarily correspond to equally large differences in droplet concentration (N_d), since 313 N_d is primarily dependent on cloud base updraft that is an extremely uncertain 314 parameter and may vary significantly between the GCMs. It therefore seems 315 reasonable to define λ as the change in LWP vs. the change in cloud droplet number 316 concentration (N_d) , which would provide a direct insight into how clouds response to 317 N_d change since LWP directly depends on N_d , not necessarily on CCN. However, this 318 alternative definition of λ as dlnLWP/dlnN_d would be difficult to compare with 319 observations, and this also does not directly measure cloud response to anthropogenic 320 aerosols. The interactions between clouds and anthropogenic aerosols arise through a 321 chain of processes, from effects of the CCN on N_d to effects of N_d on cloud water, 322 which can be expressed as dlnLWP/dlnCCN=(dlnLWP/dlnN_d)*(dlnN_d/dlnCCN). This 323 324 chain of processes has now been examined in Ghan et al., (2016) based on the same set of model simulations documented in this study.

326 **3.2 Regime dependence**

327 a. LWP, CCN and λ

Figure 1 shows LWP and CCN as a function of vertical pressure velocity at 500 hPa (ω_{500}) derived from PD simulations. To derive Figure 1, the 12-month monthly global grid values are first sorted into 20 dynamical regimes according to their ω 500 values, keeping the number of samples in each bin equal. LWP, CCN and values of other fields for each bin are then calculated from averaging the values of all samples in that particular bin.

In general, SPRINTARS (default and KK) simulates much higher LWP in all 334 dynamic regimes and ECHAM6-HAM2/ModelE2-TOMAS in most regimes than 335 different versions of CAM5 runs (default, PNNL, CLUBB and MG2) (Figure 1a), 336 which is consistent with global means in Table 2. A peak of LWP is found around 337 $\omega_{500} = 0$ hPa/d in CAM5, ModelE2-TOMAS and ECHAM6-HAM2. For SPRINTARS, 338 LWP decreases from 190 g m $^{-2}$ to 100 g m $^{-2}$ as ω_{500} increases from -60 hPa/d to 40 339 hPa/d. In all simulations LWP is low in regimes where ω_{500} is larger than 10 hPa/d, 340 i.e., regimes dominated by low clouds. HadGEM3-UKCA simulates larger LWP than 341 CAM5 especially in ascending regimes. The model spread of LWP response is larger 342 in the ascending regimes than in the subsiding regimes. This may be partly related to 343 the fact that the types of clouds included in LWP are not the same in different models 344 (Table 1). Figure 1b shows that CCN concentrations peak at around 25 hPa/d among 345

346	all the models. This peak is partly caused by little precipitation (and therefore low wet
347	scavenging rate) in subsidence regimes as well as by the fact that these dynamic
348	regimes are located near continents where the sources of anthropogenic aerosols are
349	strong. Furthermore, CCN concentrations are low at around 0 hPa/d, which could be
350	explained by the fact that most regimes around 0 hPa/d are located over the oceans far
351	away from continents (i.e. remote marine aerosols) and anthropogenic aerosol source
352	regions (figures not shown). Generally, CCN in two versions of SPRINTARS and
353	HadGEM3-UKCA is less than other models in most regimes, consistent with Table 2.
354	All the simulations show positive λ within all dynamical regimes (Figure 2a),
355	which is consistent with the theory proposed by Albrecht (1989) that an increase in
356	aerosols leads to more liquid cloud water. However, λ can vary significantly between
357	regimes in CAM5 and ECHAM6-HAM2 (Figure 2a), which indicates that changes in
358	LWP in response to aerosol perturbations are regime-dependent in these GCMs. For
359	example, λ in CAM5-PNNL ranges from 0.35 in strong ascending regions to 0.11 in
360	strong subsidence regions, which means that LWP in strong ascending regimes is
361	more sensitive to aerosol perturbations than in strong subsidence regimes. Exceptions
362	are ModelE2-TOMAS, SPRINTARS (default and SPRINTARS-KK) and
363	HadGEM3-UKCA, in which λ is low in magnitude.(i.e., LWP changes little in
364	response to the changes of CCN, consistent with the global annual means shown in
365	Table 2).

We note that although the global means of λ in all CAM5 configurations and ECHAM6-HAM2 are close, from 0.19 in ECHAM6-HAM2 to 0.25 in

CAM5-CLUBB, λ in the different dynamical regimes can differ significantly among 368 these simulations (Figure 2). For example, LWP in CAM5-PNNL is much more 369 sensitive to CCN perturbations than in ECHAM6-HAM2 in strong ascending regimes; 370 and in strong subsidence regimes, LWP in CAM5-CLUBB and ECHAM6-HAM2 is 371 more sensitive than in CAM5-PNNL and CAM5. Models that use the MG2 with 372 prognostic rain scheme (i.e. CAM5-MG2 and CAM5-CLUBB-MG2) simulate larger 373 λ than the models that use the default MG scheme in most regimes, only except for 374 strong subsidence regimes. However, generally the shapes of the λ distribution are 375 very similar. λ in CAM5-CLUBB-MG2 is large in both ascending and subsidence 376 regimes, which explains the largest global λ in CAM5-CLUBB-MG2 among all 377 configurations (Table 2). Except for the models producing very low values of λ 378 (SPRINTARS, SPRINTARS-KK, ModelE2-TOMAS and HadGEM3-UKCA), λ from 379 the other models converges around 0 hPa/d and then diverges greatly in strong 380 ascending regimes (from 0.10 to 0.46) and, to a less extent, in strong subsidence 381 regimes. This indicates that it is in regimes with weak vertical velocity where models 382 agree most, while it is in strong ascending and descending regimes where models 383 differ most. The diversity of λ within dynamical regimes in different GCMs highlights 384 the need to distinguish different dynamical regimes in studying AIE. 385

When analyzing the numerator and denominator of λ separately, we found that this large spread in λ is mainly contributed by the numerator, dlnLWP. dlnLWP ranges from about 0 to 0.22 among the models (Figure 2a) while the denominator dlnCCN, is more stable than dlnLWP within dynamical regimes and fluctuates around 390 0.45, except for larger dlnCCN in HadGEM3-UKCA (Figure 2b). In summary, the 391 ratio of dlnLWP to dlnCCN (λ) therefore changes more consistently with dlnLWP 392 within dynamical regimes.

The decreasing trends of λ with increasing ω in CAM5, CAM5-MG2, and 393 CAM5-PNNL are similar, which is opposite to the increasing trends derived from 394 ECHAM6-HAM2, CAM5-CLUBB and CAM5-CLUBB-MG2. It is interesting that 395 the regime-dependence of λ simulated by CAM5-CLUBB and CAM5-CLUBB-MG2 396 is quite different from that simulated by CAM5, CAM5-MG2, and CAM5-PNNL 397 even though all these 5 model versions are originally from CAM5 and share many 398 CAM5, CAM5-MG2 and CAM5-PNNL, 399 similarities. In three separate parameterization schemes are used to treat planetary boundary layer (PBL) turbulence, 400 stratiform cloud macrophysics and shallow convection. In CAM5-CLUBB and 401 CAM5-CLUBB-MG2, instead, a higher-order turbulence closure, Cloud Layers 402 Unified by Binormals (CLUBB), is adopted to replace these three separate schemes to 403 provide a unified treatment of these processes (Bogenschutz et al., 2013). A major 404 improvement of CAM-CLUBB is the better simulation of the transition of 405 stratocumulus to trade wind cumulus over subtropical oceans (Bogenschutz et al., 406 2013). Fig. 2a shows that λ in CAM5-CLUBB and CAM5-CLUBB-MG2 is quite 407 different from that in CAM5 simulations without CLUBB (i.e., CAM5, CAM5-MG2 408 and CAM5-PNNL) in regimes where ω_{500} is larger than 10 hPa/d. Under such 409 suppressed conditions, low clouds such as trade wind cumulus and stratocumulus are 410 typically formed. This higher λ might be expected because CAM5-CLUBB 411

formulations apply the MG microphysics (and effects of aerosols on cloud
microphysics) to shallow convective regimes. The better representation of low clouds
in CAM5-CLUBB, and the representation of double-moment microphysics and AIE
in shallow convective regimes from the unified parameterization may help to explain
the different behaviors between CAM5 runs with CLUBB (CAM5-CLUBB and
CAM5-CLUBB-MG2) and CAM5 runs without CLUBB (CAM5, CAM5-MG2 and
CAM5-PNNL) in subsidence regimes.

In order to find out the crucial geographic locations of dynamic regimes where 419 dlnLWP differs most in Fig. 2b, we plot the global distribution of annual averaged 420 dlnLWP in different simulations, shown in Fig. 3. The ascending regimes where 421 ECHAM6-HAM2 differs significantly from the two CAM5 configurations (CAM5, 422 CAM5-PNNL) are located over the North Pacific Ocean (from 30°N to 60°N), for 423 weak ascending motions and the Southern coast of Asia for strong ascending motions. 424 The spatial patterns in ECHAM6-HAM2, CAM5-CLUBB, CAM5-CLUBB-MG2 and 425 HadGEM3-UKCA share some similarities over Northern Pacific Ocean, but the 426 magnitude in CAM5-CLUBB and CAM5-CLUBB-MG2 is larger than in 427 ECHAM6-HAM2 and HadGEM3-UKCA. Moreover, not only the spatial pattern but 428 also the magnitude of dlnLWP in ECHAM6-HAM2 differ significantly from those in 429 CAM5, CAM5-MG2 and CAM5-PNNL. For the Southern coast of Asia where strong 430 ascending motions dominate, all simulations show a relative increase of LWP. 431 However, dlnLWP in ECHAM6-HAM2 in this region is much smaller than in all 432 CAM5 simulations. This makes dlnLWP, and thus λ , in ECHAM6-HAM2 much less 433

434 than in	the five CAM5 models	(CAM5, CAM5-MG2	, CAM5-PNNL	, CAM5-CLUBB
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and CAM5-CLUBB-MG2) in ascending regimes, as shown in Fig. 2b and Fig. 2a.

436 Despite the fact that SPRINTARS (default and KK), ModelE2-TOMAS and HadGEM3-UKCA all show almost no relative change of LWP in response to aerosol 437 perturbations, the spatial patterns of dlnLWP in these four simulations shown in Fig. 3 438 are indeed different from each other. HadGEM3-UKCA simulates larger dlnLWP in 439 middle northern subtropical oceans, which is similar to CAM5-CLUBB and 440 ECHAM6-HAM2 but with smaller magnitude. However, the pattern in SPRINTARS 441 is unlike any models discussed above. SPRINTARS simulates larger dlnLWP over the 442 North Pacific Ocean, the North Atlantic Ocean and the western coasts of continents 443 than other parts of the global ocean. SPRINTARS-KK simulates the same pattern as 444 SPRINTARS only with larger values. Meanwhile, dlnLWP in ModelE2-TOMAS 445 shows no special global pattern and the values are all near zero, which indicates LWP 446 in ModelE2-TOMAS has indeed little response to aerosol perturbations as 447 autoconversion rate in ModelE2-TOMAS is not influenced by cloud droplet number 448 concentrations. 449

Figure 3 shows that the differences in subsidence regimes in Fig. 2b are mainly contributed by middle northern subtropical oceans and western coasts of continents. In middle northern subtropical oceans, the relative changes of LWP in ECHAM6-HAM2, HadGEM3-UKCA and the two CAM5 models with CLUBB (CAM5-CLUBB and CAM5-CLUBB-MG2) are much more sensitive to the aerosol perturbations than in the three CAM5 models without CLUBB (CAM5,

CAM5-PNNL and CAM5-MG2), even though dlnLWP in ECHAM6-HAM2 and 456 HadGEM3-UKCA is not as large as that in CAM5-CLUBB and 457 CAM5-CLUBB-MG2. Another difference among these models is in regions 458 dominated by more intensive subsidence, over Western coasts of North America. 459 South America and Africa. In these regions dlnLWP in ECHAM6-HAM2 and the two 460 461 CAM5 models with CLUBB is large while it is small in the three CAM5 models without CLUBB. 462

To examine the cloud lifetime effect in different cloud regimes more specifically, 463 another criterion, lower-tropospheric stability (LTS= θ_{700hPa} - $\theta_{surface}$), is added to 464 distinguish stratocumulus from trade wind cumulus regimes, following Medeiros and 465 Stevens (2011). Table 3 lists the criteria of different low cloud types conditionally 466 sampled by ω_{so} and LTS. The annual mean cloud fractions of each low cloud type in 467 CAM5-CLUBB are shown in Fig. 4; the distributions in other simulations are 468 generally similar to CAM5-CLUBB (figures not shown). The cloud type distribution 469 is consistent with satellite observations that stratocumuli occurs over subtropical 470 oceans near western continents while trade wind cumuli dominate over oceans further 471 away from continents (Medeiros and Stevens, 2011). Fig. 4 shows that some 472 differences in dlnLWP between models shown in Fig. 3 are located at regions 473 dominated by low clouds (i.e., stratocumulus and trade wind cumulus). 474

The joint distributions of LTS and ω_{500} over global oceans between 60°S and 60°N derived from the models are shown in Fig. 5. Note that the bins here are not equally sampled as in previous figures but divided into equal LTS and ω intervals. 478 LTS ranges from 8K to 24K while ω ranges from -100 hPa/d to 60 hPa/d. Instances 479 with slight downward vertical motions and moderate LTS are most frequent.

480 Fig. 5 shows that, though ω_{500} plays the primary role in determining the dlnLWP/dlnCCN distribution, LTS can reveal further details of the differences among 481 various low cloud types in subsidence regimes. The large λ in strong subsidence 482 regimes in ECHAM6-HAM2 and CAM5-CLUBB is mainly caused by stratocumulus 483 and trade wind cumulus. As for regions of ascending motions, LTS is confined 484 between 12K and 14K. λ in CAM5, CAM5-PNNL and CAM5-CLUBB in ascending 485 regimes is larger than in regimes with weak large-scale vertical velocity (ω_{500} around 486 487 0 hPa/d) and larger than in ECHAM6-HAM2 in ascending regimes. In ascending regimes, LWP is more sensitive to the change of CCN in the two CAM5 models with 488 the MG2 scheme (CAM5-MG2 and CAM5-CLUBB-MG2) than in the two 489 corresponding CAM5 models without the MG2 scheme (CAM5 and CAM5-CLUBB), 490 which is consistent with Fig. 2a. In CAM5-CLUBB-MG2, λ is larger in transitional 491 cloud regimes than in stratocumulus cloud regimes and trade wind cloud regimes, 492 which is evidently different from the low cloud regimes in CAM5-CLUBB. 493 HadGEM3-UKCA simulates higher LWP response in transitional clouds and 494 stratocumulus regimes than trade wind cloud regime. It is also interesting to note that 495 λ in SPRINTARS and SPRINTARS-KK shows stronger dependence on LTS than on 496 497 ω₅₀₀.

b. Microphysics process rates and precipitation

The balance between autoconversion and accretion is found to be critical in 23

determining cloud lifetime effect in climate models (Posselt and Lohmann, 2009; 500 Wang et al., 2012). Autoconversion rate is sensitive to cloud droplet concentration 501 while accretion has little dependence of droplet number. If the role of accretion 502 dominates over autoconversion (with all other effects equal), the effect of aerosols on 503 clouds is expected to be weakened in GCMs (Posselt and Lohmann, 2009; Gettelman 504 et al 2013). Wang et al. (2012) found that the cloud lifetime effect is highly correlated 505 with the ratio of autoconversion rate to large-scale surface precipitation rate 506 (AUTO/PRECL, where PRECL also includes ice and snow) over global oceans in 507 climate models. AUTO/PRECL for different dynamical regimes is shown in Fig. 6a. 508 Here PD monthly-averaged autoconversion rate and surface precipitation rate are used 509 in calculating AUTO/PRECL. Generally the curves of AUTO/PRECL are smoother 510 than λ (Fig. 6a and Fig. 2a). The ratio from different simulations shows large diversity 511 in ascending regimes and subsidence regimes. In all versions of CAM5 and 512 SPRINTARS the ratio decreases with increasing ω_{500} in ascending regimes and then 513 increases in descending regimes. The ratio is especially large in CAM5-CLUBB-MG2 514 HadGEM3-UKCA in descending regimes. and However, the ratio 515 in ECHAM6-HAM2 remains unchanged in ascending regimes and then increases under 516 subsidence. As discussed above, λ was shown to be highly correlated with this ratio 517 from global average results (Wang et al., 2012). According to our results, the 518 correlation also applies well for individual dynamical regimes in ECHAM6-HAM2, 519 HadGEM3-UKCA and CAM-CLUBB, in which the correlation coefficients between 520 λ and AUTO/PRECL are 0.98, 0.92 and 0.86 respectively. However, these high 521

522 correlation coefficients are not found in other simulations, in which the correlation 523 coefficients are lower than 0.7, which indicates that the relationship of AUTO/PRECL 524 and λ in these models is changing from regime to regime (i.e., this relationship is 525 regime-dependent).

Wang et al. (2012) and Gettelman et al. (2013) found that the diagnostic rain 526 scheme used in the CAM configurations might overestimate the role of 527 autoconversion over accretion. Using instantaneous microphysical process rates, 528 Gettelman et al. (2015) found that adding the new microphysics with prognostic 529 precipitation to cloud scheme (MG2) decreases the ratio of autoconversion to 530 accretion. It is in moderate regimes (-20 hPa/d $\leq \omega_{500} \leq 10$ hPa/d) where the result is 531 consistent with Gettelman et al. (2015), which shows larger AUTO/PRECL in CAM5 532 than CAM5-MG2. However, in other regimes of CAM5 and all regimes of 533 CAM5-CLUBB, adding the prognostic precipitation (MG2) increases the ratio of 534 AUTO/PRECL. The result of larger AUTO/PRECL in some regimes from models 535 with MG2 seems different from the results of Gettelman and Morrison (2015) in 536 idealized tests of MG2 and of Gettelman et al. (2015) in CAM simulations with MG2. 537 We have verified using the same model output from Gettelman et al. (2015) that the 538 difference is not due to the simulations performed. The difference is likely due to: (a) 539 the use of instantaneous output in Gettelman et al. (2015) for process rate 540 comparisons while monthly data is used here; (b) Microphysics variables and 541 precipitation are sorted by ω_{500} here while Gettelman et al. (2015) sorted them by 542 543 LWP that the microphysics sees, which includes contributions from deep convection;

(c) Vertical integrals of autoconversion rate are used here while vertical mean valuesare used in Gettelman et al. (2015).

546 As discussed in Section 1, precipitation is a key process in interactions between aerosols and clouds. A decrease in surface precipitation increases cloud water while a 547 decrease in cloud-top sedimentation increases the entrainment rate and thus dries out 548 LWP when the free troposphere air is dry (Ackerman et al., 2004). Here we 549 investigate the LWP response to aerosol perturbations under low precipitation 550 (monthly-averaged surface precipitation rate less than $0.1 \text{ mm } d^{-1}$) and high 551 precipitation (monthly-averaged surface precipitation rate larger than 0.1 mm d⁻¹). 552 Table 4 lists the occurrence frequency of each situation in different simulations. It 553 shows that instances with low PRECL occurs much less often (from 2.2% in 554 CAM5-CLUBB to 38.8% in CAM5-MG2) than those with high PRECL. The 555 occurrence frequency of low precipitation situations is increased with the MG2 556 scheme (CAM5-MG2 and CAM5-CLUBB-MG2), compared with simulations without 557 MG2. This increase is especially evident in CAM5-CLUBB (from 0.02 in 558 CAM5-CLUBB to 0.16 in CAM5-CLUBB-MG2). This is consistent with Gettelman 559 et al. (2015), who showed surface precipitation decreases slightly in GCMs with 560 MG2. 561

Note that low precipitation situations are only found in subsidence regimes ($\omega_{500} > 0$ hPa/d). Thus, the sensitivity of the LWP response to aerosol change under low and high precipitation is compared only in subsidence regimes. Table 4 also shows λ and the fractional occurrences of each precipitation situation in descending regimes. The

fractional occurrence of low precipitation increases evidently in subsidence regimes, 566 compared with that over global ocean. We find that the averages of λ under low 567 precipitation are larger than those under high precipitation in most models (CAM5, 568 CAM5-PNNL, CAM5-CLUBB, ECHAM6-HAM2, SPRINTARS, SPRINTARS-KK 569 and HadGEM3-UKCA) (Table4). This result is different from some LES and single 570 571 column model (SCM) results showing that smaller λ values are found for low surface precipitation rather than high precipitation due to a decrease of LWP in response to 572 increasing CCN (Ackerman et al., 2004; Guo et al., 2011). The decrease in LWP in 573 these previous studies is found to come from the entrainment drying due to increased 574 entrainment from increasing aerosol loading (e.g., Bretherton et al., 2007) and this 575 effect has not been explicitly included in most GCMs. Exceptions are CAM5 runs 576 with the prognostic precipitation scheme MG2 (CAM5-MG2, CAM5-CLUBB-MG2). 577 It can be seen from Table 4 that λ under low surface precipitation is smaller than 578 under high precipitation only when MG2 scheme is used. It is still unclear what might 579 cause this difference. It is interesting to note that λ under low surface precipitation is 580 still higher for HadGEM3-UKCA though a prognostic precipitation scheme is applied 581 in HadGEM3-UKCA. 582

583 c. Shortwave cloud radiative effect

The shortwave cloud radiative effect (SCRE) is defined as the difference between all-sky and clear sky shortwave radiative fluxes at the top of atmosphere. Here SCRE is adjusted to the "clean-sky" SCRE, which is estimated as a diagnostic with aerosol optical depth set to zero (Ghan, 2013). Recent studies on aerosol indirect effects mostly focus on stratocumulus clouds due to their significant cooling effect (e.g., Lu and Seinfeld, 2005; Bretherton et al., 2007). However, by sorting the change of SCRE (dSCRE) from PI to PD into dynamical regimes, our results suggest that the regimes of ascending motions are as important as the subsidence regimes and in some simulations dSCRE in ascending regimes is even larger than under subsidence regimes (e.g., CAM5-PNNL) (Fig. 7). This suggests that ascending regimes are crucial regimes in studying aerosol climate effect.

We also examined dSCRE contributed by low and high precipitation situations 595 (note that the total dSCRE is the sum of dSCRE under low and high precipitation 596 situation). It is found that high precipitation situations constitute most of dSCRE 597 (from 64% in CAM5-MG2 to nearly 100% in CAM5-CLUBB, Fig. 7) and the 598 contributions from clouds with low precipitation rates are generally small, ranging 599 from 0% to 36%, due to their low occurrence frequency. dSCRE is reduced by 33% 600 for high precipitation situations from CAM5 to CAM5-MG2, and 15% from 601 CAM5-CLUBB to CAM5-CLUBB-MG2 (Fig. 7), consistent with the argument that 602 prognostic precipitation schemes reduce aerosol indirect forcing (Posselt and 603 Lohmann, 2009; Wang et al., 2012; Gettelman and Morrison, 2015). However, 604 adopting a prognostic precipitation scheme is found to increase dSCRE under low 605 precipitation situations. This is partly from the increase in the occurrence frequency of 606 low precipitation instances when MG2 is adopted (Table 4). 607

608 Our sensitive tests indicate that results in Table 4 and Figure 7 can be potentially 609 sensitive to the precipitation threshold applied to separate high precipitation and low

precipitation situations (not shown). The occurrence frequency of low precipitation 610 situations increases with increasing threshold and the magnitude of increase can be 611 different for different models. For example, when the precipitation threshold increases 612 from 0.01 mm d^{-1} to 0.20 mm d^{-1} , the occurrence frequency of low precipitation 613 situations increases from 2% to 37% in CAM5-PNNL while it increases from near 0% 614 to 5% in CAM5-CLUBB. Increasing the precipitation threshold also increases the 615 contribution of low precipitation situations to the total aerosol indirect forcing as the 616 occurrence frequency of low precipitation situations increases. However, our results 617 indicate that the LWP response to aerosol perturbations under low and high 618 precipitation does not change much as the precipitation threshold changes and that 619 high precipitation situations generally contribute more to the total aerosol indirect 620 forcing for precipitation threshold in the range of 0.01 to 0.20 mm d⁻¹. More work is 621 needed to explore this further such as how results may be different when 622 instantaneous precipitation data (e.g., 3-houly data) is used. 623

624 **4 Summary**

We have examined the regime-dependence of aerosol indirect effects (AIE) over 625 global oceans (from 60° S to 60° N) in several GCMs (CAM5, CAM5-MG2, 626 CAM5-CLUBB, CAM5-CLUBB-MG2, CAM5-PNNL, ECHAM6-HAM2, 627 SPRINTARS, SPRINTARS-KK, ModelE2-TOMAS and HadGEM3-UKCA). Model 628 results are sorted into different dynamical regimes, characterized by the 629 monthly-mean mid-tropospheric 500hPa vertical pressure velocity $(\omega_{500}),$ 630 631 lower-tropospheric stability (LTS, θ_{700hPa} - $\theta_{surface}$) and surface precipitation rate.

The response of liquid water path (LWP) to aerosol perturbations, 632 λ =dlnLWP/dlnCCN, a metric to quantify cloud lifetime effect of aerosols (Wang et al., 633 2012), shows a large spread within dynamical regimes among GCMs, although the 634 global means are close. This diversity indicates that the aerosol cloud lifetime effect is 635 regime-dependent. It is in strong ascending regimes and subsidence regimes that λ 636 differs most between GCMs (Fig. 2a). Stratocumulus regimes have traditionally been 637 the focus for studying aerosol indirect effects because of their significant cooling 638 effect in climate system (e.g., Ackerman et al., 2004; Bretherton et al., 2007; 639 Gettelman et al., 2013). However, our results highlight that regimes with strong 640 large-scale ascent should be another important regime to focus on in the future. Our 641 results indicate that aerosol indirect forcing in regimes of vertical ascent is close to, or 642 even larger than that in low cloud regimes (Fig. 7). Note however that these GCMs do 643 not treat aerosol effects in their representations of deep convection that dominates 644 645 clouds and LWP in regimes with strong ascent, while new versions of CAM exist where a version of the MG microphysics has been embedded in the deep convective 646 parameterization (Song and Zhang, 2011). 647

By adding LTS as another criterion, we further separated different low cloud types under large-scale subsidence and revealed some further differences in cloud lifetime effect of aerosols on different types of low clouds. For example, the large λ in subsidence regimes in CAM5-CLUBB and ECHAM6-HAM2 comes from both stratocumulus and trade wind cumulus, while in CAM5-CLUBB-MG2 it mostly comes from trade wind cumulus (Fig. 5). It is also interesting to note that the distribution of λ in SPRINTARS and SPRINTATSKK is more likely to depend on
LTS rather than vertical pressure velocity.

Precipitation is another important factor in understanding simulated aerosol 656 indirect forcing and its spread across models. LWP is more sensitive to CCN change 657 under low precipitation situations (monthly-mean surface precipitation rate less than 658 0.1 mm d^{-1}) than under high precipitation situations (monthly-mean surface 659 precipitation rate larger than 0.1 mm d⁻¹) in all models except for CAM5 simulations 660 with prognostic rain scheme (MG2) (Table 4). Results derived from large eddy 661 simulation (LES) and single column model (SCM) (e.g., Ackerman et al., 2004; Guo 662 et al., 2011) have shown that λ could be negative under low precipitation situations, 663 which indicates that λ is expected to be smaller under low precipitation situations. 664 Further efforts are needed to understand the differences among different models and 665 the difference between global model results and results from process-level studies. 666

Our results indicate that grids with high precipitation contribute most to aerosol 667 indirect forcing (from 64% in CAM5-MG2 to nearly 100% in CAM5-CLUBB, Fig. 7) 668 and the contributions from model grids with low precipitation are relatively small, 669 ranging from 0% to 36%. Adding prognostic precipitation scheme (MG2) reduces the 670 shortwave cloud radiative effect (SCRE) for high precipitation situations. As low 671 precipitation situations are much less prevalent than high precipitation situations, total 672 673 SCRE decreases in models with prognostic rain scheme compared to those with a diagnostic rain scheme. 674

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The regime categorization used in this study is derived from monthly mean 31

data. Giving the high variability of precipitation and microphysics processes on short 676 time scales, we acknowledge that instantaneous data (e.g. 3 hourly) might provide 677 more reliable information. For example, instantaneous data may help to reconcile 678 some of discrepancies between our studies and that of Gettelman et al. (2015) 679 regarding the prognostic rain scheme noted in Section 3.2b. However, it is challenging 680 to calculate λ and aerosol indirect forcing using instantaneous data. Here λ and aerosol 681 indirect forcing are derived from the difference between present day (PD) and 682 pre-industrial (PI) simulations. Using instantaneous data will not guarantee that the 683 sorted bins of dynamical regimes include the same instances from PI to PD, giving the 684 high variability of instantaneous data. Since the main goal in this manuscript is to 685 demonstrate the importance of examining aerosol indirect effects in different cloud 686 and dynamical regimes, the use of monthly-mean data serves this goal well. It is our 687 future plan to carry in-depth analysis to further understand some of the findings 688 documented here, such as the large spread in λ in regimes of vertical ascent in 689 different models. For example, LWP response to aerosol perturbation documented in 690 this study may include contributions from mixed-phase and ice clouds. In- depth 691 analysis of cloud macrophysics and microphysics processes will help to improve the 692 understanding of the model uncertainty. 693

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Appendix. Global aerosol-climate models

695 CAM5: This is the default version of CAM5.3. The moist turbulence scheme is 696 based on Bretherton and Park (2009), which explicitly simulates stratus-697 radiation-turbulence interactions. The shallow convection scheme is from Park and

Bretherton (2009) and the deep convection parameterization is retained from CAM4.0 698 (Neale et al., 2008). The two-moment cloud microphysics scheme from Morrison and 699 Gettelman (2008) (MG) is used to predict both the mass and number mixing ratios for 700 cloud water and cloud ice with a diagnostic formula for rain and snow. The cloud ice 701 microphysics was further modified to allow ice supersaturation and aerosol effects on 702 ice clouds (Gettelman et al., 2010). The activation of aerosol particles into cloud 703 droplets is parameterized by Abdul-Razzak and Ghan (2000, hereafter ARG) and the 704 autoconversion scheme is based on Khairoutdinov and Kogan (2000) (KK). A modal 705 approach is used to treat aerosols in CAM5 (Liu et al., 2012; Ghan et al., 2012). 706 Aerosol size distribution can be represented by using either 3 modes or 7 modes, and 707 the default 3-mode treatment is used in this study. Simulations were performed at 1.9° 708 $\times 2.5^{\circ}$ horizontal resolution with finite volume dynamical core, using 30 vertical 709 710 levels.

CAM5-PNNL: This is the same as CAM5, but a new unified treatment of vertical 711 transport and in-cloud wet removal processes in convective clouds developed by 712 Wang et al. (2013) is applied. It has a more detailed treatment of aerosol activation in 713 convective updrafts and a mechanism is added for laterally entrained aerosols to be 714 activated and then removed. In addition, a few other changes have been introduced to 715 stratiform cloud wet scavenging processes in CAM5-PNNL to improve the fidelity of 716 717 the aerosol simulation, including the vertical distribution of aerosols and their transport to remote regions (Wang et al., 2013). 718

CAM5-MG2: This is the same as CAM5, but the original two-moment MG scheme with diagnostic treatment for rain and snow in CAM5 is replaced by the updated MG scheme (MG2) with prognostic scheme for rain and snow (Gettelman et al., 2015).

CAM5-CLUBB: This is the same as CAM5, but the separate treatments of boundary layer turbulence, large-scale cloud macrophysics and shallow convection in CAM5 is replaced by CLUBB, a higher-order turbulence closure that unifies these different treatments (Bogenschutz et al., 2013). This therefore includes aerosol effects on shallow convection.

CAM5-CLUBB-MG2: This is the same as CAM5-CLUBB, but the MG2 scheme
with prognostic rain and snow treatment replaces the original MG scheme with
diagnostic rain and snow treatment (Gettelman et al., 2015). This also includes aerosol
effects on shallow convection.

ECHAM6-HAM2: ECHAM-HAMMOZ (echam6.1-ham2.2-moz0.9) is a global 732 aerosol-chemistry climate model. In this study only the global aerosol-climate model 733 part of ECHAM-HAMMOZ is used and for the sake of brevity referred to as 734 735 ECHAM6-HAM2 (Neubauer et al., 2014). It consists of the general circulation model ECHAM6 (Stevens et al., 2013) coupled to the latest version of the aerosol module 736 HAM2 (Stier et al., 2005; Zhang et al., 2012) and uses a two-moment cloud 737 microphysics scheme that includes prognostic equations for the cloud droplet and ice 738 crystal number concentrations as well as cloud water and cloud ice (Lohmann et al., 739 2007; Lohmann and Hoose, 2009). The activation of aerosol articles into cloud 740 34

droplets is parameterized by Lin and Leaitch (1997) and the autoconversion scheme is based on the KK scheme. Cumulus convection is represented by the parameterization of Tiedtke (1989) with modifications by Nordeng (1994) for deep convection. Aerosol effects on convective clouds are not included, but there is a dependence of cloud droplets detrained from convective clouds on aerosol. Simulations were performed at T63 $(1.9^{\circ} \times 1.9^{\circ})$ spectral resolution using 31 vertical levels (L31).

SPRINTARS: SPRINTARS (Takemura et al. 2005) is a global aerosol 747 transport-climate model based on a general circulation model, MIROC (Watanabe et 748 al. 2010). In this study, the horizontal and vertical resolutions are T106 (1.125° x 749 approx. 1.125°) and 56 layers, respectively. SPRINTARS is coupled with the radiation 750 and cloud microphysics schemes in MIROC to calculate the aerosol-radiation and 751 aerosol-cloud interactions. A prognostic scheme for determining the cloud droplet and 752 ice crystal number concentrations is introduced (Takemura et al. 2009). The default 753 atuoconversion scheme in MIROC-SPRINTARS is based on Berry (1967), and the 754 activation of aerosol particles into cloud droplet is based on the ARG scheme. 755

756 SPRINTARS-KK: This is the same as SPRINTARS, but the default 757 autoconversion scheme in SPRINTARS is replaced with the KK auconversion 758 scheme.

ModelE2-TOMAS: ModelE2-TOMAS is a global-scale atmospheric chemistry-climate model, which consists of the state-of-the-art NASA GISS ModelE2 general circulation model (Schmidt, 2014) coupled to the TwO-Moment Aerosol Sectional (TOMAS) microphysics model (Lee and Adams, 2012; Lee et al., 2015).

ModelE2-TOMAS has 2° latitude by 2.5° longitude resolution, with 40 vertical hybrid 763 sigma layers from the surface to 0.1 hPa (80 km). In the model, clouds are 764 distinguished into convective and large-scale stratiform clouds. The clouds 765 parameterizations are similar to Del Genio (1993) and Del Genio (1996) but have 766 been improved in several respects (see details in Schmidt, 2014; Schmidt, 2006). 767 Using a prognostic treatment of cloud droplet number concentration (CDNC) from 768 Morrison and Gettleman (2008), ModelE2-TOMAS represents the first aerosol 769 indirect effects only on large-scale stratiform clouds (Menon et al., 2010). In 770 ModelE2-TOMAS, CDNC and a critical supersaturation are computed using a 771 physical-based activation parameterization from Nenes and Seinfeld (2002) with a 772 model updraft velocity that is computed based on a large-scale vertical velocity and 773 sub-grid velocity. 774

HadGEM3-UKCA: HadGEM3-UKCA is a global composition climate model 775 (http://www.ukca.ac.uk). It consists of the third generation of the Hadley Centre 776 Global Environmental Model (Hewitt et al, 2011) developed at the UK Met Office. 777 778 This general circulation model is non-hydrostatic and uses a semi-Lagrangian transport scheme. We are using the atmospheric configuration: General Atmosphere 779 (GA) 4.0 as documented in Walters et al., (2014), except for the addition of the 780 UKCA aerosol and chemistry scheme which is fully coupled with the radiation 781 scheme of HadGEM3 (Bellouin et al., 2013). UKCA is a two-moment pseudo-modal 782 scheme which carries both aerosol number concentration and component mass as 783 prognostic tracers. It calculates the evolution of five aerosol species, sulfate, 784

particulate organic matter, black carbon, sea salt and dust, in both internally and 785 externally mixed particles. The aerosol scheme in UKCA is based on the Global 786 Model of Aerosol Processes (GLOMAP-mode, Mann et al., 2010). The main 787 exception is that dust is calculated separately using 6 size bins. UKCA hence only 788 considers 5 modes. The tropospheric chemistry part of UKCA is described in 789 790 O'Connor et al. (2014). HadGEM3 uses a prognostic treatment of rain formulation (Abel and Boutle, 2012) and employs a prognostic cloud fraction and condensation 791 cloud scheme (PC2) (Wilson et al., 2008), in which the cloud droplet number 792 concentration is diagnosed from the expected number of aerosols that are available to 793 activate at each timestep (West et al., 2014). Cumulus convection is represented by a 794 mass flux convection scheme based on Gregory and Rowntree (1990) with various 795 extensions (Walters et al., 2014). Simulations were performed at N96L85 resolution, 796 a regular 1.25° latitude \times 1.875° longitude grid in the horizontal, with 85 797 hybrid-height vertical levels. 798

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1151 Tabel 1. The types of clouds included in liquid water path (LWP) and surface rain

Model	LWP	Rain	Rain scheme
CAM5	S*	S	d&
CAM5-MG2	S	S	$\mathbf{p}^{@}$
CAM5-PNNL	S	S	d
CAM5-CLUBB	S+shallow convective clouds	S+shallow convective clouds	d
CAM5-CLUBB-MG2	S+shallow convective clouds	S+shallow convective clouds	р
ECHAM6-HAM2	S+convective detrainment	S	d
SPRINTARS	$S+C^{\#}$	S+C	d
SPRINTATRS-KK	S+C	S+C	d
ModelE2-TOMAS	S+anvil clouds	S+anvil clouds	d
HadGEM3-UKCA	S+C	S	р

rate and different rain schemes in 10 participating models

1153 * S in LWP and Rain stands for stratiform clouds.

1154 # C in LWP and Rain stands for convective clouds.

1155 & d in Rain schemes represents diagnostic rain scheme.

1156 *(a)* p in Rain schemes represents prognostic rain scheme.

1157

Table 2.Global ocean (60°S-60°N) averages of LWP, column-integrated cloud condensation nuclei (CCN, at 0.1% supersaturation) concentration, precipitation rate (PRECL), shortwave cloud radiative effect (SCRE) derived from the present day (PD) cases and the relative change from pre-industrial (PI) to PD of LWP and CCN (dlnLWP and dlnCCN) and the sensitivity of LWP to CCN concentration change (λ , dlnLWP/dlnCCN) of the 10 GCM simulations.

Model	λ	LWP (g	$\begin{array}{c} \text{CCN} \\ (10^{11} \end{array}$	dlnLWP	dlnCCN	PRECL (mm	SCRE
1110.001		m^{-2})	m^{-2})	unit () i	unicert	d^{-1}	$(W m^{-2})$
CAM5	0.20	31.1	1.86	0.07	0.36	0.90	-61.9
CAM5-MG2	0.23	30.0	1.73	0.07	0.32	0.76	-67.9
CAM5-PNNL	0.19	39.4	2.51	0.08	0.42	0.91	-64.6
CAM5-CLUBB	0.25	35.2	1.88	0.11	0.45	1.26	-57.7
CAM5-CLUBB-MG2	0.27	47.1	1.66	0.11	0.42	1.08	-70.6
ECHAM6-HAM2	0.19	84.6	2.39	0.07	0.41	1.35	-54.5
SPRINTARS	0.01	139.1	1.07	0.00	0.43	1.42	-62.6
SPRINTATRS-KK	0.04	98.9	1.04	0.02	0.45	1.59	-57.0
ModelE2-TOMAS	0.00	80.4	2.66	0.00	0.43	2.17	-68.1
HadGEM3-UKCA	0.03	57.1	1.01	0.01	0.67	0.87	-58.9

1168	Table 3.	Criteria	used to	conditional	sampling	stratocumulus,	transitional	clouds	and

1169	trade wind cu	mulus regimes	(adopted from	Medeiros and	Stevens	(2011)))
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		Stratocumulus	Transitional	Trade wind
			clouds	cumulus
	LTS (K)	LTS≥18.5	18.5>LTS≥15.4	15.4>LTS≥11.3
	ω_{500hPa} (hPa d ⁻¹)	ω _{500hPa} >10	ω _{500hPa} >10	ω _{500hPa} >10
171				

1174 Table 4. The fractional occurrences of low and high surface precipitation in PD cases 1175 over downdraft regimes ($\omega_{500} > 0$ hPa/d) and global oceans and λ under these low and 1176 high surface precipitation situations only over downdraft regimes. Low precipitation 1177 situations refer to monthly surface precipitation rate (PRECL) less than 0.1 mm d⁻¹ 1178 while high precipitation situations refer to PRECL larger than 0.1mm d⁻¹.

	λ^{a}	λ^{b}	f^c	f^d	f ^e	f^{f}
Model	low,	high,	low,	high,	low,	high,
	down	down	down	down	glb	glb
CAM5	0.21	0.19	0.47	0.54	0.27	0.73
CAM5-MG2	0.19	0.24	0.57	0.43	0.39	0.61
CAM5-PNNL	0.17	0.17	0.48	0.52	0.28	0.72
CAM5-CLUBB	0.33	0.30	0.04	0.96	0.02	0.98
CAM5-CLUBB-MG2	0.26	0.33	0.22	0.78	0.16	0.84
ECHAM6-HAM2	0.25	0.23	0.31	0.69	0.18	0.82
SPARINTARS	0.06	0.01	0.06	0.94	0.03	0.97
SPARINTARS-KK	0.24	0.04	0.05	0.95	0.03	0.97
ModelE2-TOMAS	-0.011	0.001	0.002	0.998	0.001	0.999
HadGEM3-UKCA	0.04	0.03	0.11	0.89	0.06	0.94

- 1181 ^a λ under low PRECL for downdraft regimes
- 1182 ^b λ under high PRECL for downdraft regimes
- ^c Fractional occurence of low PRECL for downdraft regimes
- ^d Fractional occurrence of high PRECL for downdraft regimes
- ^e Fractional occurrence of low PRECL over all dynamical regimes
- ^f Fractional occurrence of high PRECL over all dynamical regimes



Fig 1. (a) LWP and (b) column-integrated CCN (at 0.1% supersaturation) as a function of 500 hPa vertical pressure velocity (ω_{500}) derived from different models: CAM5 (blue solid line), CAM5-MG2 (blue dashed line), CAM5-PNNL (blue dotted line), CAM5-CLUBB (cyan solid line), CAM5-CLUBB-MG2 (cyan dashed line), ECHAM6-HAM2 (red solid line), SPRINTARS (green solid line), SPRINTARS-KK (green dashed line), ModelE2-TOMAS (purple solid line) and HadGEM3-UKCA (orange solid line).

1196



1200 Fig 2. Same as Fig. 1a), but for (a) the sensitivity of LWP to the change of CCN (λ),

1201 (b) relative enhancement of liquid water path (dlnLWP) and (c) relative enhancement

1202 of cloud condensation nuclei (dlnCCN) from pre-industrial (PI) to present day (PD).





1209 Fig 3. Relative change of annual averaged LWP from PI to PD (dlnLWP)

simulations derived from the 10 GCM simulations.



1214

Fig 4. The annual mean cloud fraction (averaged on the months when the regime occurs) of stratocumulus regime (top left), transitional clouds regime (top right) and trade wind cumulus regime (bottom left) derived from PD monthly simulation in CAM5-CLUBB. The definitions of different cloud types are listed in Table 3.





Fig 5. dlnLWP/dlnCCN conditioned on vertical motion and LTS derived from the 10
GCM simulations. Solid lines are contours of grid number distribution and each line
interval is 20% of the total counted data.



Fig 6. Same as Fig. 1, but for (b) column-integrated autoconversion rate (AUTO), (c) the large-scale surface precipitation rate (PRECL) and (a) their ratio AUTO/PRECL from the 9 GCM simulations. The number marked in each simulation is the corresponding correlation coefficient between AUTO/PRECL and λ and number with mark '*' indicates the correlation is significant (at 95% confidence).

1228



Fig 7. Change in shortwave cloud radiative effect (dSCRE, shown in blue line) from
PI to PD as a function of dynamic regimes. Red patches are dSCRE contributed by
low precipitation situations while blue patches are by high precipitation situations.