

An optimal estimation based aerosol retrieval algorithm using OMI near-UV observations

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Abstract

An optimal estimation (OE) based aerosol retrieval algorithm using the OMI (Ozone Monitoring Instrument) near-ultraviolet observation was developed in this study. The OE-based algorithm has the merit of providing useful estimates of errors simultaneously with the inversion products. Furthermore, instead of using the traditional look-up tables for inversion, it performs online radiative transfer calculations with the Vector Linearized Discrete Ordinate Radiative Transfer (VLIDORT) model to eliminate interpolation errors and improve stability. The measurements and inversion products of the Distributed Regional Aerosol Gridded Observation Network campaign in Northeast Asia (DRAGON NE-Asia 2012) were used to validate the retrieved AOT and SSA. The retrieved AOT and SSA at 388 nm have a correlation with the Aerosol Robotic Network (AERONET) products that is comparable to or better than the correlation with the operational product during the campaign. The OE-based estimated error better represented the variance of actual biases of AOT at 388 nm between the retrieval and AERONET measurements than the operational error estimates. The forward model parameter errors were analyzed separately for both AOT and SSA retrievals. The surface reflectance at 388 nm, the imaginary part of the refractive index at 354 nm, and the number fine mode fraction (FMF) were found to be the most important parameters affecting the retrieval accuracy of AOT, while FMF was the most important parameter for the SSA retrieval. The additional information provided with the retrievals, including the estimated error and degrees of freedom, is expected to be valuable for relevant studies. Detailed advantages of using the OE method were described and discussed in this paper.

Key words

Aerosol; Error Analysis; Ozone Monitoring Instrument; Optimal Estimation Method; DRAGON campaign

1. Introduction

Anthropogenic aerosols have affected both the radiative and meteorological balance in the atmosphere and thus the radiative forcing of the atmosphere directly and indirectly (Ramanathan et al., 2001; Russell et al., 1999; Breon et al., 2002). To understand the role of aerosol in the atmosphere from a global perspective, reliable aerosol data from satellites are essential (Al-Saadi et al., 2005; Kinne et al., 2006). The several satellite-based aerosol retrieval methods based on multi-wavelength (Levy et al., 2007; Kim et al., 2007), multi-angle (Fisher et al., 2014), active light (Young et al., 2013), and polarization (Deuze et al., 2001) measurements have their own advantages and limitations. The inversion products from such measurements provide various parameters of aerosols at different channels. Thus, appropriate sources of aerosol information need to be employed for relevant studies.

An important advantage of using the ultraviolet (UV) channel to retrieve aerosol optical properties is that the results are less affected by uncertainties in surface reflectance (Torres et al., 1998). The retrieved aerosol products have relatively uniform quality over both land and ocean except over ice–snow surfaces (Torres et al., 2007; Herman et al., 1997). The near-UV technique for aerosol remote sensing has the additional merit of a long term data record including aerosol absorption properties of over 30 years starting from the launch of the Total Ozone Mapping Spectrometer (TOMS) on Nimbus-7 in 1978 (Torres et al., 1998; Torres et al., 2002a; Torres et al., 2005). Thus, the retrieved products using the near-UV technique from TOMS and Ozone Monitoring Instrument (OMI) measurements are appropriate for climatological research (Torres et al., 2002b; Torres et al., 2007). Information on aerosol extinction and absorption properties in the UV region is also important for estimating the air mass factor (AMF) for trace gas retrievals (Palmer et al., 2001; Lin et al., 2014). However, deriving information on aerosol using available hyperspectral measurements such as OMI and Global Ozone Monitoring Experiment (GOME) is quite challenging due to the relatively low spatial resolution compared to typical imagers. Thus, the

error estimates of retrievals using such sensors are particularly important to understand the reliability of the information, so that it can be used appropriately. The main objective of this study is to improve the applicability of the aerosol inversion products of OMI by providing the reliable error estimates of the retrievals.

Accuracy assessments of the retrieved aerosol optical properties using UV radiances have been performed by comparison with results from reference methods including ground, airborne, and satellite based remote sensing techniques (Torres et al., 2005; Jethva et al., 2014; Torres et al., 2002a; Ahn et al., 2014; Livingston et al., 2009; Ahn et al., 2008; Curier et al., 2008). The aerosol information content of selected OMI spectral radiances using a multi-wavelength algorithm has been estimated using principal component analysis for simulated radiances (Veihelmann et al., 2007). Uncertainty estimates of UV aerosol retrievals have also been calculated by perturbation analysis (Torres et al., 1998; Torres et al., 2002b). Inversion algorithms based on optimal estimation (OE) theory provide not only a constrained solution with respect to the *a priori* information but also detailed error analysis from well-categorized error sources (Rodgers, 2000). In addition, iterative inversion methods such as OE provide additional retrieval masking parameters (*e.g.*, cost function and convergence criteria). Recently developed OE-based retrieval methods have provided both improved inversion products and error estimates from the aerosol and surface error sources (Wagner et al., 2010; Govaerts et al., 2010; Wurl et al., 2010).

A large amount of aerosol is emitted from both natural and anthropogenic sources in East Asia (Lee et al., 2012). The spatial and temporal variations in aerosol optical properties are significant because of the diverse emission sources and trans-boundary transport (Jeong et al., 2011). Thus, the assumed aerosol inversion parameters may cause substantial uncertainties in the retrieval. However, there are insufficient ground-based measurements of aerosol optical properties with suitable spatial and temporal coverage in East Asia, despite their importance for global air quality and climate change. The Distributed Regional Aerosol Gridded Observation Network Northeast

51 Asia (DRAGON-NE Asia) 2012 campaign ([http://aeronet.gsfc.nasa.gov/new_web/DRAGON-](http://aeronet.gsfc.nasa.gov/new_web/DRAGON-Asia_2012_Japan_South_Korea.html)
52 [Asia_2012_Japan_South_Korea.html](http://aeronet.gsfc.nasa.gov/new_web/DRAGON-Asia_2012_Japan_South_Korea.html)) provides valuable datasets including both urban and
53 regional-scale observations at more than 40 sites in Northeast Asia. In the present study, an OE-
54 based near-UV aerosol retrieval and error analysis algorithm is developed to provide both improved
55 aerosol inversion products and estimates of their uncertainties. The retrieved aerosol products and
56 estimated uncertainties are validated against the DRAGON-NE Asia 2012 campaign measurements.

57

58 **2. Data**

59 OMI is a nadir-viewing hyperspectral spectrometer aboard the EOS (Earth Observing
60 System)-Aura spacecraft that measures upwelling radiances from the top of the atmosphere in the
61 ultraviolet and visible (270–500 nm) regions with approximate spectral resolution of 0.5 nm (Levelt
62 et al., 2006). The advantage of using OMI for aerosol retrieval is its higher spatial resolution than
63 other UV hyperspectral spectrometers (from $13 \times 24 \text{ km}^2$ at nadir to $28 \times 150 \text{ km}^2$ at the swath
64 extremes with median pixel size $15 \times 32 \text{ km}^2$) together with its 2600-km-wide swath. The
65 radiometric calibration procedure and the estimated accuracy of OMI are described in Dobber et al.
66 (2006). To determine aerosol type and vertical distribution, the current OMI near-UV aerosol
67 algorithm (OMAERUV) employs the Cloud–Aerosol Lidar with Orthogonal Polarization (CALIOP)
68 monthly climatology of aerosol layer height and real-time Atmospheric Infrared Sounder (AIRS)
69 carbon monoxide (CO) observations (Torres et al., 2013). Surface reflectances at 354 nm and 388
70 nm were assumed to be Lambertian and were taken from the TOMS climatology database. Aerosol
71 vertical distribution and surface reflectance information identical to that used in the operational
72 algorithm were used for the OE-based algorithm here.

73 In this study, the spatial and temporal domains for analysis were confined to the DRAGON-
74 NE Asia 2012 campaign as shown in Figure 1 and Table 1. The gridded observation networks had
75 high spatial resolution over the representative megacities in Northeast Asia: Seoul in South Korea

76 and Osaka in Japan. To validate and compare the retrieved aerosol products from OMI, level 2
77 campaign products were used from the aerosol robotic network (AERONET); 380 nm aerosol
78 optical thickness (AOT) from direct sun measurements and spectral single scattering albedo (SSA)
79 from almucantar inversion products (Holben et al., 1998;Dubovik and King, 2000;Dubovik et al.,
80 2000;Dubovik et al., 2006). Retrieved 388 nm AOT from OMI was validated against AERONET
81 380 nm AOT. The OMI AOT retrievals within a radius of 0.5° of the AERONET site and within
82 ± 30 minutes of the OMI overpass time (about 13:40 local time) were averaged. The resulting OMI
83 AOT average values were then compared with the time-averaged Sun photometer measurements.

84 Aerosol absorption properties are retrieved at different wavelengths by AERONET and
85 OMI. The AERONET inversion products of the SSA are available at 440, 670, 860, and 1020 nm,
86 while the OMAERUV algorithm retrieves the SSA at 354 nm and 388 nm. Earlier field studies
87 found that aerosol absorption is a continuous function of wavelength in the ultraviolet to short
88 infrared region (Kirchstetter et al., 2004;Russell et al., 2010). To compare the SSA values from OMI
89 and AERONET at the same wavelength, the AERONET SSA at 388 nm was obtained by
90 extrapolating the SSAs at 440–1020 nm using a spline function. Then the converted AERONET
91 SSA at 388 nm was compared with the retrieved OMI SSA values even though uncertainties might
92 exist in the transformation. Unlike the direct sun measurements including AOT, the inversion
93 products of AERONET from almucantar measurements are retrieved less frequently and require
94 appropriate atmospheric conditions for AOT (440 nm AOT > 0.4) and solar zenith angle (solar
95 zenith angle > 45°) (Dubovik and King, 2000;Jethva et al., 2014). Such favorable atmospheric
96 conditions for the inversion using almucantar measurements rarely overlap closely with the OMI
97 overpass time. Furthermore, too narrow a time window around the satellite overpass time reduces
98 the number of comparison samples. In this study, to secure enough data points the SSA of a region
99 at OMI overpass time was assumed to adequately represent the daily values. For the comparison,

100 the converted 388 nm SSA from AERONET was averaged over a day and the OMI retrievals of 388
101 nm SSA were spatially averaged over a grid area of $0.5^{\circ} \times 0.5^{\circ}$ centered on the AERONET site.

102

103 **3. Method**

104 **3.1. Operational OMI near UV aerosol algorithm**

105 The OMAERUV uses two channel radiances at 354 nm and 388 nm to estimate aerosol
106 amount and absorption properties (Torres et al., 1998; Torres et al., 2007; Torres et al., 2013). AOT
107 and SSA at 388 nm are retrieved from pre-calculated reflectance look-up tables (LUT) for pre-
108 determined nodal points of observational geometry and aerosol optical properties, total optical
109 depth, and aerosol layer height. Three major aerosol types are considered and listed in Table 2:
110 desert dust, carbonaceous aerosols associated with biomass burning, and weakly absorbing sulfate-
111 based aerosols (hereafter dust, smoke, and sulfate, respectively). Each aerosol type has an assumed
112 particle size distribution (PSD) derived from the long-term statistics of AERONET inversion
113 products. The UV real refractive index (n_r) is obtained from the Optical Properties Of Aerosols And
114 Clouds (OPAC) database (Hess et al., 1998). In the operational algorithm, the imaginary refractive
115 indices (n_i) at 354 nm are assumed to be 1.0, 1.2, and 1.4 times the retrieved n_i at 388 nm for sulfate,
116 smoke, and dust aerosol, respectively (Torres et al., 2007; Jethva and Torres, 2011). The overall
117 concept and design of the improved OMAERUV algorithm is well described by Torres et al. (2013).

118 There have been further improvements at updated OMAERUV (version 1.5.3), which was
119 used for reprocessing the data of AOT and SSA in this study. The OMAERUV algorithm was
120 refined by adjusting thresholds of UV aerosol index (UVAI) and Atmospheric Infrared Sounder
121 (AIRS) CO data in determining aerosol types and retrieval approaches. A cloud screening scheme in
122 assigning algorithm quality flags was also modified for retaining more good retrievals of
123 carbonaceous and sulfate type aerosols when the CO level is high enough (higher than 3.2×10^{18}
124 molecules $\cdot \text{cm}^{-2}$) with various reflectivity thresholds. The UVAI threshold was changed from 0.8 to

125 0.5 over the oceans. This modification eliminates the land-ocean discontinuity in UVAI threshold. It
126 is now identical (0.5) for both conditions. The current characterization of ocean reflective properties
127 in the OMAERUV algorithm does not explicitly account for ocean color effects and, therefore, the
128 quality of the retrieved aerosol properties over the oceans for low aerosol amounts would be highly
129 uncertain. For that reason, retrievals over the oceans are only carried out for high concentrations of
130 either desert dust or carbonaceous aerosols as indicated by UVAI values larger than or equal to 0.5.

131 Depending on the magnitude of the UVAI and CO parameters as well as the aerosol type,
132 two retrieval approaches are currently used. They are referred to as two-channel and single-channel
133 retrievals. In the two-channel approach, observations at 354 and 388 nm are used to simultaneously
134 derive AOD and SSA. Over scenes when the aerosol absorption signal is low, the single-channel
135 retrieval is applied. AOD is retrieved from the 388 nm observation assuming a value of 1.0 for SSA.
136 Different CO threshold values are used for the northern and southern hemispheres to remove upper
137 tropospheric CO which may not be necessarily associated with carbonaceous aerosols. A smoothing
138 function in CO is used to transition from SH to NH threshold values. Specific criteria for retrieval
139 approaches are summarized in Table 3. More detailed information of the latest update in
140 OMAERUV is available from the Readme file at the web site
141 (http://disc.sci.gsfc.nasa.gov/Aura/data-holdings/OMI/omaerv_v003.shtml).

142

143 **3.2. OE-based OMI near UV aerosol algorithm**

144 The traditional LUT-based inversion method potentially includes errors due to interpolation
145 between the nodal points and the local minimum, despite its high numerical efficiency. Such
146 interpolation error typically depends on the interpolation method, number of the nodal points, and
147 analytic characteristics of the parameters in LUT. In order to reduce the interpolation error, higher
148 resolution of LUT nodal points is necessary which requires larger amount of numerical computation.
149 Furthermore, in order to modify the retrieval algorithm, whole LUT should be re-calculated even

150 for a few number of target retrievals. The errors from the interpolation are also hard to evaluate as
151 the LUT becomes more complicated.

152 On the contrary, online retrieval methods can reduce such errors from the interpolation and
153 are numerically efficient particularly for the smaller number of target retrievals. Thus, online
154 retrieval method is appropriate for the research purposes since retrieval sensitivity study typically
155 use smaller number of sample compared to the operational purposes and prefer rapid and accurate
156 results. In our experience, the online retrieval method was numerically more efficient compared to
157 the LUT-based retrieval method by order of 1 or 2 for less than few thousands of retrievals.
158 Furthermore, the online retrieval methods are optimized to avoid local minima by employing
159 additional constraints to find more reliable and stable solutions (Kalman, 1960;Phillips,
160 1962;Tikhonov, 1963;Twomey, 1963;Chahine, 1968). However, employing online calculation as
161 operational retrieval method requires large computation cost. Thus, using the online calculation as a
162 benchmark results for the LUT-based algorithm is recommended to develop the optimized LUT for
163 the operational purposes. Recent efforts to minimize the numerical cost of radiative transfer model
164 and to increased calculation speed are expected to make the online calculation more practical even
165 for the operational purposes.

166 Optimization for measurement error, the inclusion of *a priori* and ancillary data, and
167 employing physical constraints (*e.g.*, non-negativity of measurements and retrievals) for an
168 inversion method are important since each method has its own advantages and disadvantages
169 (Dubovik, 2004). In this study, we used OE as the inversion method (Rodgers, 2000) since it has
170 several advantages over other methods for OMI-like measurements, as discussed in Section 4. The
171 atmospheric inverse problem often suffers from both insufficient information content of the
172 measurements and imperfect measurement accuracy. Bayesian statistics provides mapping methods
173 from the measurement probability density function (pdf) into state space with prior knowledge.
174 Based on Bayes' theorem, the OE technique employs additional constraints from external sources (*a*

175 *priori*) to complement the insufficient information content of the measurements. For the nonlinear
 176 inversion case, by considering the maximum *a posteriori* approach, the general form of the
 177 Bayesian solution can be expressed as Equation (1) where measurement and *a priori* errors are
 178 assumed to be Gaussian (Rodgers, 2000):

179

$$180 \quad -2\ln P(\mathbf{x}|\mathbf{y}) = [\mathbf{y} - \mathbf{K}\mathbf{x}]^T \mathbf{S}_\epsilon^{-1} [\mathbf{y} - \mathbf{K}\mathbf{x}] + [\mathbf{x} - \mathbf{x}_a]^T \mathbf{S}_a^{-1} [\mathbf{x} - \mathbf{x}_a] + \mathbf{c} \quad (1)$$

181

182 where \mathbf{x} is the state vector and \mathbf{y} the measurement vector, \mathbf{K} is the weighting function matrix,
 183 \mathbf{S}_ϵ is the measurement error covariance matrix, \mathbf{x}_a is the *a priori* mean state, and \mathbf{S}_a is the *a*
 184 *priori* covariance matrix. The formulation finds the optimized solution that minimizes the cost
 185 function (χ):

186

$$\chi^2 = \left\| \mathbf{S}_\epsilon^{-\frac{1}{2}} (\mathbf{y} - \mathbf{K}\mathbf{x}) \right\|_2^2 + \left\| \mathbf{S}_a^{-\frac{1}{2}} (\mathbf{x} - \mathbf{x}_a) \right\|_2^2 \quad (2)$$

187

188 Detailed derivations and implications are described in previous studies (Rodgers, 2000; Wurl et al.,
 189 2010; Govaerts et al., 2010). As described above, the OMI near-UV algorithm uses radiance (I_{388})
 190 and spectral contrast (I_{354}/I_{388}) for the measurement vector, where I_{354} and I_{388} are the
 191 normalized radiances at 354 nm and 388 nm, respectively. The state vector in this study is the AOT
 192 at 388 nm (τ_{388}) and the imaginary refractive index at 388 nm ($n_{i,388}$). Then, the weighting
 193 function matrix can be expressed as follows:

194

$$\mathbf{K} = \begin{bmatrix} \frac{\partial I_{388}}{\partial \tau_{388}} & \frac{\partial I_{388}}{\partial n_{i,388}} \\ \frac{\partial}{\partial \tau_{388}} \left(\frac{I_{354}}{I_{388}} \right) & \frac{\partial}{\partial n_{i,388}} \left(\frac{I_{354}}{I_{388}} \right) \end{bmatrix} \quad (3)$$

195

196 where the weighting function of the spectral contrast can be obtained from the following derivative:

197

$$\frac{\partial}{\partial x} \left(\frac{I_{354}}{I_{388}} \right) = \frac{\frac{\partial I_{354}}{\partial x} I_{388} - I_{354} \frac{\partial I_{388}}{\partial x}}{(I_{388})^2} \quad (4)$$

198

199 In typical inversion methods, including OE, estimation of the reliable measurement error
 200 covariance matrix is important to determine the likelihood of the solution (Govaerts et al., 2010).
 201 The measurement error includes radiometric noise error and calibration accuracy. The absolute
 202 bidirectional scattering distribution function (BSDF) radiometric accuracy of the OMI instrument is
 203 reported to be about 4% for 2σ and the random noise error is provided in the level 1b product
 204 (Dobber et al., 2006). However, the reported BSDF uncertainty includes the errors in the calibration
 205 method and it represents whole wavelength domain. Thus, actual BSDF uncertainty at 354 and 388
 206 nm would be less than 2% (Jaross, 2015). In our experience, 2% of BSDF uncertainty leads to the
 207 overestimates of the error and it is still challenging to evaluate. According to multiple retrieval tests,
 208 the BSDF uncertainty was assumed to be 1% in this study. The radiometric error covariance at each
 209 wavelength was calculated from the square root of the sum of squared radiometric uncertainty and
 210 calibration accuracy. The error covariance matrix can be written as:

211

$$\mathbf{s}_\epsilon = \begin{bmatrix} \sigma(\epsilon_{388})^2 & \sigma(\epsilon_{388}, \epsilon_{354/388})^2 \\ \sigma(\epsilon_{388}, \epsilon_{354/388})^2 & \sigma(\epsilon_{354/388})^2 \end{bmatrix} \quad (5)$$

212

213 where ϵ_λ is the total error of the measured radiance at wavelength λ , $\epsilon_{354/388}$ is the error of
 214 I_{354}/I_{388} , which is described later in this section, and $\sigma(\epsilon_{388}, \epsilon_{354/388})^2$ is the covariance
 215 between the total measurement errors of I_{388} and I_{354}/I_{388} .

216 The ϵ_λ typically includes both random and systematic components and its covariance can
 217 be expressed as follows:

218

$$\sigma(\epsilon_\lambda)^2 = \sigma(\epsilon_{r,\lambda})^2 + \sigma(\epsilon_{s,\lambda})^2 \quad (6)$$

219

220 where $\epsilon_{r,\lambda}$ and $\epsilon_{s,\lambda}$ are the random and systematic components of radiometric error at λ , and
 221 $\sigma(\epsilon_{r,\lambda})^2$ and $\sigma(\epsilon_{s,\lambda})^2$ are their covariance values, respectively. The $\epsilon_{354/388}$ can be
 222 approximated as follows:

223

$$\epsilon_{354/388} \cong \frac{I_{354}}{I_{388}} - \frac{I_{354} + \epsilon_{354}}{I_{388} + \epsilon_{388}} \cong \frac{I_{354}}{I_{388}} \left(\frac{\epsilon_{354}}{I_{354}} - \frac{\epsilon_{388}}{I_{388}} \right) \cong \frac{I_{354}}{I_{388}} \left(\frac{\epsilon_{s,354}}{I_{354}} - \frac{\epsilon_{s,388}}{I_{388}} + \frac{\epsilon_{r,354}}{I_{354}} - \frac{\epsilon_{r,388}}{I_{388}} \right) \quad (7)$$

224

225 When the systematic components of the measurement errors of radiances at 354 nm and 388 nm are
 226 positively correlated and their values are similar, part of the systematic uncertainties can be reduced
 227 by the $\frac{\epsilon_{s,354}}{I_{354}} - \frac{\epsilon_{s,388}}{I_{388}}$ term. However, assessment of the systematic error of OMI measurements at
 228 each pixel is still challenging despite this partial reduction of systematic errors by using I_{354}/I_{388} .
 229 In this study, the BSDF calibration uncertainties of I_{354} and I_{388} at a pixel are assumed to be
 230 systematic and similar, while the radiometric noise values of I_{354} and I_{388} are assumed to be
 231 random and independent. Then, the systematic measurement error of $\epsilon_{354/388}$ can be regarded as
 232 negligible and the covariance of $\epsilon_{354/388}$ in Equation (7) can be approximated as follows:

233

$$\sigma(\epsilon_{354/388}) \cong \frac{I_{354}}{I_{388}} \sqrt{\left(\frac{\epsilon_{r,354}}{I_{354}}\right)^2 + \left(\frac{\epsilon_{r,388}}{I_{388}}\right)^2} \quad (8)$$

The $\sigma(\epsilon_{388}, \epsilon_{354/388})^2$ can be obtained as follows:

$$\sigma(\epsilon_{388}, \epsilon_{354/388})^2 = \frac{1}{n-1} \sum_{i=1}^n \epsilon_{388}^i \epsilon_{354/388}^i \quad (9)$$

where ϵ_{388}^i and $\epsilon_{354/388}^i$ are the uncertainties in the i^{th} measurement of I_{388} and I_{354}/I_{388} for a sample of size n , respectively. Under the same assumptions used in Equations (7) and (8), $\epsilon_{354/388}^i$ has only random and independent components, and so $\sigma(\epsilon_{388}, \epsilon_{354/388})$ can be regarded as negligible. The diagonal and off-diagonal elements of the measurement error covariance matrices using two different measurement matrices are compared in Table 4.

3.3. Error characterization

Retrieved products with estimated and/or characterized error are valuable for any application. Various error sources can be categorized as shown in the following equation by linearizing the forward model with respect to associated parameters (Eyre, 1987; Rodgers, 1990, 2000):

$$\hat{\mathbf{x}} - \mathbf{x} = (\mathbf{A} - \mathbf{I}_n)(\mathbf{x} - \mathbf{x}_a) + \mathbf{G}_y \mathbf{K}_b (\mathbf{b} - \hat{\mathbf{b}}) + \mathbf{G}_y \Delta \mathbf{f}(\mathbf{x}, \mathbf{b}, \mathbf{b}') + \mathbf{G}_y \boldsymbol{\epsilon} \quad (10)$$

where $\hat{\mathbf{x}}$ and \mathbf{x} are the retrieval and true states, respectively; \mathbf{A} is the averaging kernel matrix; \mathbf{I}_n is the identity matrix; \mathbf{x}_a is the *a priori* state vector; \mathbf{G}_y is the contribution function matrix; \mathbf{K}_b

253 is the weighting function matrix of forward model parameters (\mathbf{b}); $\hat{\mathbf{b}}$ is the guessed forward model
 254 parameter; $\Delta \mathbf{f}$ is the error in the forward model relative to the real physics; and ϵ is the
 255 measurement error. The first and last term on the right-hand-side (RHS) of Equation (10) is the
 256 smoothing error and retrieval noise, respectively. Their covariance matrices can be calculated from:

$$\mathbf{S}_s = (\mathbf{A} - \mathbf{I}_n) \mathbf{S}_E (\mathbf{A} - \mathbf{I}_n)^T \quad (11)$$

$$\mathbf{S}_n = \mathbf{G}_y \mathbf{S}_\epsilon \mathbf{G}_y^T \quad (12)$$

258
 259 where \mathbf{S}_s is the smoothing error covariance matrix, \mathbf{S}_E is the covariance of the ensemble of states
 260 about the mean state, and \mathbf{S}_n is the covariance matrix of the retrieval noise. We have assumed that
 261 the climatological value is a good representation of the real ensemble of the state about the mean
 262 state, and so the covariance matrix of the *a priori* state was employed as the \mathbf{S}_E in this study. Each
 263 \mathbf{S}_s and \mathbf{S}_n has two diagonal elements that represent the variances of the smoothing error and
 264 retrieval noise for two retrievals, AOT and SSA. The smoothing error (ϵ_{sm}) and retrieval noise (ϵ_n)
 265 of AOT and SSA were defined as the square root of the corresponding diagonal elements of \mathbf{S}_s and
 266 \mathbf{S}_n , respectively. The square root of the sum of squared ϵ_{sm} and ϵ_n of AOT and SSA are defined
 267 in this study as the solution errors (ϵ_{sol}) of AOT and SSA, respectively.

268 The second term on the RHS of Equation (10) is the forward model parameter error, and its
 269 covariance matrix can be calculated as follows:

$$\mathbf{S}_f = \mathbf{G}_y \mathbf{K}_b \mathbf{S}_b \mathbf{K}_b^T \mathbf{G}_y^T \quad (13)$$

271
 272 where \mathbf{S}_f is the covariance matrix of the forward model parameter error, and \mathbf{S}_b is the error
 273 covariance matrix of forward model parameter, \mathbf{b} . The forward model parameters of the near-UV
 274 method include the aerosol microphysical model parameters, aerosol vertical distribution,

275 meteorological profile (pressure and temperature), and surface properties. These forward model
276 parameters contain both random and systematic components with different scales of spatial and
277 temporal variation. Furthermore, each forward model parameter has a different uncertainty that is
278 difficult to evaluate. In this study, \mathbf{S}_f was analyzed separately with respect to each forward model
279 parameter as suggested by Rodgers (2000). The forward model parameter error (ε_f) of AOT and
280 SSA were obtained from the square root of the corresponding diagonal elements of each \mathbf{S}_f .

281 The third term on the RHS in Equation (10) is the forward model error that is caused by
282 discrepancies between known and real physics. To simulate the earth-reflected radiance, VLIDORT
283 (linearized pseudo-spherical vector discrete ordinate radiative transfer code, version 2.6) was used.
284 This code is based on one of the most accurate radiative transfer solutions for a one-dimensional
285 atmosphere (Spurr, 2006). Linearization of state vectors and forward model parameters are
286 described in prior papers (Spurr et al., 2012; Spurr and Christi, 2014). Although the simulated
287 radiances are expected to be accurate, the forward model error depends on factors including the
288 number of streams, layers, and Legendre coefficients for the aerosol phase functions. To reduce the
289 numerical errors that can arise from an insufficient number of coefficients, 3 Stokes parameters, 75
290 layers, 16 streams, and up to 500 Legendre coefficients for the aerosol phase matrix were used for
291 the radiance simulations. However, it is still a challenge to evaluate other possible sources of
292 forward model error such as Raman scattering and the three-dimensional effect of the atmosphere
293 for retrieval. Such issues are beyond the scope of this study, and thus only smoothing error, retrieval
294 noise, and forward model parameter error are evaluated here.

295

296 **3.4. *A priori* characterization**

297 Using reliable *a priori* information is important in the OE method since the final solutions
298 are determined between the *a priori* state and the inversion space of a measurement. There are
299 several sources of *a priori* information including climatological data, reliable measurements from

more accurate instruments, and calculations from models based on theoretical or empirical statistics (Govaerts et al., 2010; Wurl et al., 2010; Rodgers, 2000). Appropriate sources of *a priori* depend on the characteristics of the state vector and the accuracy of the *a priori* database. When the *a priori* state has a systematic bias away from the true state, this bias propagates to the retrieval products. In this study, 10 years (from 2005 to 2014) of OMAERUV 388 nm AOT and SSA in spring (from March to May) were used for the *a priori* data. Figure 2 shows the collected climatological data of the 388 nm AOT and SSA in East Asia used in this study. To avoid biases due to cloud contamination in the OMI products, averaged values and standard deviations (σ) with more than 70 data points were used.

4. Results

The dust event on 28th April 2012 has been selected to compare the aerosol optical properties from the operational product with the OE-based retrievals in this study. Figure 3 shows the true color image on that day from the Moderate Resolution Imaging Spectroradiometer (MODIS) Aqua and the UV aerosol index from OMI. To see the difference between the OE-based and operational algorithm in the same area, both methods were applied to measurements with an operational algorithm flag of 0, to avoid cloud contamination and radiometric calibration uncertainties. Some of the points retrieved by the operational algorithm were rejected in the OE-based retrievals when the cost function cut-off was applied, as described in this section.

The AOT and SSA at 388 nm from the operational and OE-based products are compared in Figure 4. The ε_{sol} of the retrieved AOT and SSA at 388 nm, degree of freedom, and χ are shown in Figure 5. The OE-based and operational AOT at 388 nm are similar, as shown in Figure 4 (a) and (b). Both products seem to be affected by snow and cloud contaminated pixels around Seoul in South Korea (37°N, 126°E), Tianjin in China (39°N, 116°E) and Nagano in Japan (36°N, 138°E). In Figure 5 (a), the ε_{sol} of retrieved AOT at 388 nm has relatively high values compared with the

325 AOT level at large viewing zenith angle. Noticeable discrepancies of SSA at 388 nm from the
 326 operational and OE-based products are evident in some of the areas as shown in Figure 4 (c) and (d).
 327 The operational algorithm performed single-channel retrieval around East Mongolia (47°N, 115°E),
 328 while the OE-based algorithm performed two-channel retrieval for all cases. Since this area has a
 329 low level of AOT, the information content of aerosol absorption property is insufficient, resulting in
 330 the low degrees of freedom shown in Figure 4 (b) and Figure 5 (c). Thus the OE-based SSA in this
 331 area (Figure 4 (d)) seems noisy and the ε_{sol} of SSA appears high. Similar results for the behavior
 332 of SSA are apparent around central Japan (38°N, 138°E). Thus, SSA values with high AOT and low
 333 ε_{sol} are recommended for the analysis, similar to the retrieval conditions for AERONET inversion
 334 products (Dubovik and King, 2000). Additional information provided on retrievals at each pixel in
 335 Figure 5 is expected to be valuable for relevant studies including trace gas retrieval and data
 336 assimilation.

337 Figure 6 shows the results of validation of operational and OE-based AOT retrievals at 388
 338 nm. As shown in Figure 6 (a) and (b), the OE-based inversion method showed higher correlation
 339 coefficient ($r = 0.82$) and slightly improved slope (0.83) and offset (0.16) values than the
 340 operational algorithm ($r = 0.71$, slope = 0.71, and offset = 0.2). The Fisher's z -value between the
 341 correlation coefficients (Fisher, 1921) was 3.04 corresponding to two-tailed p -value of 0.0024. The
 342 Student's t -value for the difference between the two slopes is 2.10 with 512 degrees of freedom
 343 with the two-tailed p -value of 0.04. The statistical values show that difference between the two
 344 correlation coefficients and slopes are significant (p -value < 0.05). The Q values (percentage of
 345 AOT retrievals falling within an uncertainty envelope of $\pm 30\%$ or 0.1) of the OE-based retrievals
 346 and operational algorithm were similar (63.0%). When a measured radiance is affected by
 347 parameters that the theoretical radiative transfer model does not consider (*e.g.*, sub-pixel cloud
 348 contamination), the χ of the retrieval typically has a high value. In this study, retrievals with χ larger
 349 than a certain value (*i.e.*, 2.0 in this study) have been rejected. This limitation on retrievals imposed

350 by the χ reduced the number of retrievals with abnormally high biases, which might be associated
351 with sub-pixel cloud contamination, in the operational algorithm in Figure 6 (a).

352 The SSA values at 388 nm from OMI operational products and OE-based inversion
353 products were compared with those at 388 nm and 440 nm from AERONET inversion products as
354 shown in Figure 7. The retrieved SSA at 388 nm from the operational algorithm showed comparable
355 or higher values of $Q_{0.03}$ (59.2%) and $Q_{0.05}$ (85.1%) with those from the OE-based algorithm ($Q_{0.03}$
356 = 53.5%, $Q_{0.05}$ = 86.0%) when compared with the AERONET SSA at 440 nm (The $Q_{0.03}$ and $Q_{0.05}$
357 represent the percentage of SSA retrievals falling within an uncertainty envelope of ± 0.03 and ± 0.05 ,
358 respectively). The retrieved 388 nm SSA from both the operational and OE-based algorithms
359 showed similar correlation with the AERONET ($r = 0.27$ and 0.26 for operational and OE-based
360 algorithms, respectively. Fisher's z -value is 0.1 with two-tailed p -value of 0.92). The retrieved SSA
361 at 388 nm from the operational and OE-based algorithms showed slightly higher correlation with
362 the converted 388 nm SSA from AERONET ($r = 0.34$ and 0.33 for the operational and OE-based
363 algorithm, respectively) than with the 440 nm SSA from AERONET. However, the significances of
364 the differences in r between converted and unconverted SSA comparisons were low (Fisher's z -
365 values were 0.71 and 0.67 with two-tailed p -values of 0.48 and 0.50 for operational algorithm and
366 OE-based algorithm, respectively). The retrieved SSA at 388 nm from the operational algorithm
367 also showed comparable or higher values of $Q_{0.03}$ (59.2%) and $Q_{0.05}$ (83.9%) than those of the OE-
368 based algorithm ($Q_{0.03} = 53.5\%$, $Q_{0.05} = 82.8\%$) when compared with converted SSA at 388 nm
369 from AERONET.

370 The estimated retrieval uncertainties of the AOT at 388 nm from the operational algorithm
371 (ε_{omi} , $\pm 30\%$ or 0.1) and estimated ε_{sol} were plotted against the biases relative to AERONET
372 measurements as shown in Figure 8. The percentages of AOT retrieval biases from AERONET
373 falling within the estimated retrieval errors of operational (Q_{omi}) and OE-based method (Q_{sol}) were
374 64.8% and 65.9%, respectively. The Q_{sol} was higher than Q_{omi} despite of the lower mean value of

ε_{sol} (0.20) than that of ε_{omi} (0.21). The error bars and black squares in Figure 8 represent the moving σ and average value of the retrieval biases from AERONET as a function of estimated error, respectively. As shown in Figure 8 (b), ε_{sol} better explained the moving σ of the actual biases ($r=0.93$) than ε_{omi} in Figure 8 (a) ($r=0.52$). Fisher's z -value between the correlation coefficients was 2.33 with two-tailed p -value of 0.02. The systematic biases of ε_{sol} and ε_{omi} (represented by the moving average of each error estimates) are typically related to other error sources, including forward model parameters and sub-pixel cloud contaminations. Since the ε_{sol} of retrieved AOT considers the theoretical sensitivity of the retrieval biases to associated parameters, it explained the retrieval uncertainties better than the ε_{omi} , which only considers the retrieved AOT values.

Table 5 shows the suggested error sources and their magnitudes from the OMI ATBD (Algorithm Theoretical Basis Documents) (Torres et al., 2002b) and the values employed in this study. Although the current OMI cloud masking method is based on long-term TOMS heritage, there may still be ground pixels contaminated by sub-pixel clouds. As the TOA reflectance is greatly increased by even a small amount of cloud, cloud contamination can cause large positive biases in the AOT retrieval. Previous studies estimated the AOT retrieval errors due to 5% cloud contamination to be of the order of 0.1 to 0.2 (Torres et al., 1998; Torres et al., 2002b). They also reported an even higher error in the single scattering albedo (0 to 0.15) especially for strongly absorbing aerosols. However, estimation of sub-pixel cloud contamination is difficult because of the large spatio-temporal variability of clouds and the relatively large ground pixel size of OMI. Thus the further error analysis of cloud contamination error budget from Torres et al. (2002b) was not performed in this study. Typical uncertainties of the 354 nm and 388 nm surface reflectances were assumed to be 0.01 for both land and ocean. The BSDF accuracy was assumed to be 1% (Dobber et al., 2006; Jaross, 2015), and radiometric precision from OMI Level 1b data. To analyze the uncertainty associated with the aerosol size information and refractive index, σ values of the size parameter and n_r at 440 nm were taken from AERONET inversion products during the campaign

400 period. To analyze the assumed n_i at 354 nm, the \mathbf{S}_b was also obtained from AERONET statistics
 401 during the campaign period. Aerosol vertical distribution is important as it affects aerosol retrieval
 402 using near-UV and blue channels, particularly for absorbing aerosols (Torres et al., 1998; de Graaf et
 403 al., 2005; Torres et al., 2013). However, accuracy assessments of the aerosol height information used
 404 are still challenging. Typical uncertainties of the assumed aerosol layer peak height and half width
 405 were assumed to be 2 km and 1 km, respectively, in this study. In the OE-based near-UV aerosol
 406 retrieval algorithm, all aerosols are assumed to be spherical and the optical properties are calculated
 407 from aerosol microphysical properties using the Mie solution. However, non-sphericity may cause
 408 significant uncertainties, especially for large particles (Mishchenko and Travis, 1994; Mishchenko et
 409 al., 1995; Mishchenko et al., 1997; Mishchenko et al., 2003; Dubovik et al., 2006), and aerosol
 410 morphology is quite complicated and requires further analysis for the near-UV region. This is out of
 411 the scope of this study and thus needs to be investigated in a future study. Therefore the
 412 uncertainties due to aerosol non-sphericity were not analyzed.

413 Figure 9 shows the average and σ values of ε_f of the retrieved AOT and SSA that were
 414 sampled for the validation in Figure 6 (b) and Figure 7 (b), (d). High values of ε_f for AOT
 415 appeared in n_i at 354 nm (0.34 ± 0.25), surface reflectance at 388 nm (0.19 ± 0.07), and the number
 416 fine mode fraction (FMF) (0.16 ± 0.09). These values are higher or comparable with the mean ε_{sol}
 417 of retrieved AOT at 388 nm (0.20). Thus, the accuracy of AOT retrievals depends on not only the
 418 radiometric accuracy and information content but also the aerosol models and ancillary data of the
 419 surface reflectance, of which the effect is already well known. The FMFs of the sulfate (0.999596)
 420 and smoke type aerosols (0.999795) are similar while that for dust type aerosols is quite different
 421 (0.995650). Considering that the estimated σ value for FMF uncertainty in this study (0.0015, see
 422 Table 5) is much lower than the difference between the FMFs of dust type and other aerosols
 423 (~ 0.004 , see Table 2), the errors resulting from selection of the wrong aerosol type can be more
 424 significant. The estimated ε_f of the surface reflectance at 388 nm was higher than the previously

425 suggested value (0.07–0.09 for AOT and <0.01 for SSA) in the OMI ATBD (Torres et al., 2002b).
 426 The ε_f s of AOT with respect to the surface reflectance at 354 nm (0.12 ± 0.04), peak height of the
 427 aerosol vertical distribution (0.11 ± 0.10), fine mode n_r (0.09 ± 0.10), half width of the fine mode PSD
 428 (0.07 ± 0.06), mean radius of the coarse mode PSD (0.06 ± 0.03), and half width of the aerosol
 429 vertical distribution (0.06 ± 0.05) showed similar moderate sensitivity. Those for the mean radius of
 430 the fine mode PSD (0.02 ± 0.02) and half width of the coarse mode PSD (0.02 ± 0.01) were smaller
 431 and that of the coarse mode n_r (0.003 ± 0.004) was found to be negligible.

432 Among ε_f s of the SSA retrieval, the FMF error ($1.4 \times 10^{-2} \pm 6.4 \times 10^{-3}$) was the most
 433 important of the ε_f s of the SSA retrieval. Errors in n_i at 354 nm ($7.4 \times 10^{-3} \pm 2.8 \times 10^{-3}$) and the peak
 434 height of the aerosol vertical distribution ($5.2 \times 10^{-3} \pm 2.8 \times 10^{-3}$) were found to be the second most
 435 important. The ε_f of SSA with respect to the fine mode n_r ($3.9 \times 10^{-3} \pm 1.7 \times 10^{-3}$), width of the fine
 436 mode PSD ($3.8 \times 10^{-3} \pm 1.7 \times 10^{-3}$), the surface reflectance at 354 nm ($3.4 \times 10^{-3} \pm 1.8 \times 10^{-3}$) and 388
 437 nm ($2.8 \times 10^{-3} \pm 2.1 \times 10^{-1}$), mean radius of the coarse mode PSD ($3.1 \times 10^{-3} \pm 1.1 \times 10^{-3}$), half width of
 438 the aerosol vertical distribution ($2.3 \times 10^{-3} \pm 1.5 \times 10^{-3}$), mean radius of the fine mode PSD (1.1×10^{-3}
 439 $\pm 5.7 \times 10^{-4}$) and width of the coarse mode PSD ($8.7 \times 10^{-4} \pm 3.3 \times 10^{-4}$) were smaller and that of the
 440 coarse mode n_r ($5.2 \times 10^{-5} \pm 1.3 \times 10^{-4}$) appeared to be negligible. The estimated ε_f s of SSA were
 441 found to be about a factor magnitude lower than the ε_{sol} of SSA. The mean values of ε_{sm} and ε_n
 442 of SSA were 0.023 and 0.029, respectively. Thus, the estimated ε_{sol} of SSA at 388 nm is expected
 443 to be more reliable and represent the total uncertainties of SSA, since the uncertainty in SSA is
 444 predominantly affected by ε_{sol} , while uncertainty in AOT is affected by both ε_{sol} and ε_f . Note
 445 that the relative significances of the ε_f s of retrievals depend on their condition. It is additional
 446 merit of the error analysis using OE method that it provides specific error estimates of individual
 447 target event retrieval (e.g., dust or biomass burning event). While analysis studies using satellite
 448 inversion products have often suffered from the statistic reliabilities, more reliable error estimates in
 449 this study are expected to contribute to the assessment of significances of the analysis.

450

451 **5. Summary and Discussion**

452 An OE-based aerosol retrieval and error characterization algorithm using the OMI near-UV
453 radiances was developed in this study. The climatological values of OMAERUV products were
454 employed as *a priori* data for the inversion method. The OE-based inversion method developed
455 here provides not only the retrieved values of AOT and SSA but also estimates of their uncertainties.
456 The retrieved AOT and SSA at 388 nm were compared with the AERONET products during the
457 DRAGON-NE Asia 2012 campaign. The retrieved AOT using the OE method showed better results
458 than the operational product. The OE-based SSA at 388 nm showed consistency with AERONET
459 inversion products comparable to that of the operational SSA. The estimated retrieval noise and
460 smoothing error of OE-based AOT represented well the variances of actual biases between the
461 retrieved AOT and AERONET AOT. The forward model parameter errors were analyzed separately
462 for both AOT and SSA inversion products. Uncertainties of surface reflectance at 388 nm,
463 imaginary refractive index at 354 nm and number fine mode fraction were found to be the most
464 important parameters affecting the retrieval accuracy of AOT, while uncertainties in the coarse
465 mode real part of the refractive index had negligible effect. For SSA retrieval accuracy, number fine
466 mode fraction was found to be the most important parameter while the other parameters appeared to
467 have relatively small effects. As the FMF depends on the aerosol type, it is expected that more
468 accurate aerosol type classification might improve the retrieval accuracy of AOT and SSA. For
469 AOT retrieval, the estimated ε_f was comparable with the ε_{sol} , while the ε_f of SSA was
470 negligible compared to the ε_{sol} of the retrieved SSA. It is also found that a sufficient amount of
471 aerosol loading is necessary for reliable SSA retrieval.

472 However, there are still error sources which need to be analyzed, including the *a priori* error
473 from climatology, aerosol morphology, cloud contamination, and three dimensional effects of
474 radiative transfer. The assumed conditions in the inversion procedure also differ from the real state.

475 Validation studies for a longer period at more types of site are also necessary. Securing a more
476 reliable *a priori* database is expected to improve the OE-based aerosol retrieval algorithm.

477

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680 **Tables and Figures**

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694 **Table 1.** Positions and mean AOT and SSA at 388 nm of the AERONET sites during the
695 DRAGON-NE Asia 2012 campaign.

Site name	Latitude (°) (North)	Longitude (°) (East)	Mean 380 nm AOT	Mean 440 nm SSA
Baengnyeong	37.97	124.63	0.56	0.96
Yonsei Univ., Seoul	37.56	126.94	0.63	0.93
Anmyeon	36.54	126.33	0.57	0.94
Bokjeong, Seoul	37.46	127.13	0.75	0.90
Gangneung Wonju National Univ.	37.77	128.87	0.53	0.92
Guwol, Seoul	37.45	126.72	0.69	0.94
GIST, Gwangju	35.23	126.84	0.54	0.93
HUFS, Yongin	37.34	127.27	0.63	0.90
Kongju National Univ., Kongju	36.47	127.14	0.62	0.96
Konkuk Univ., Seoul	37.54	127.08	0.67	0.92
Korea Univ., Seoul	37.59	127.03	0.73	0.92
Kunsan National Univ., Kunsan	35.94	126.68	0.61	0.92
Kyungil Univ., Kyungsan	36.07	128.82	0.57	0.93
Mokpo National Univ., Mokpo	34.91	126.44	0.58	0.93
NIER, Incheon	37.57	126.64	0.62	0.93
Pusan National Univ., Pusan	35.24	129.08	0.62	0.93
Sanggye, Seoul	37.66	127.07	0.73	0.92
Sinjeong, Seoul	37.52	126.86	0.64	0.91
Soha, Seoul	37.45	126.89	0.69	0.91
Gosan, Jeju	33.29	126.16	0.62	0.96
Seoul National Univ., Seoul	37.46	126.95	0.65	0.93
Fukuoka	33.52	130.48	0.50	0.90
Kohriyama	37.36	140.38	0.34	0.95
Kyoto	35.03	135.78	0.47	0.94
Matsue	35.48	133.01	0.56	0.93
Mt. Ikoma	34.68	135.68	0.39	0.96
Mt. Rokko	34.76	135.23	0.41	0.95
Nara	34.69	135.83	0.48	0.94
Nishiharima	35.03	134.34	0.42	0.95
North Osaka	34.77	135.51	0.52	0.94
South Osaka	34.54	135.50	0.55	0.94

Tsukuba	36.05	140.12	0.38	0.94
Noto	37.33	137.14	0.41	0.94
Shirahama	33.69	135.36	0.41	0.96
Chiba University	35.63	140.10	0.31	0.92
Fukue	32.75	128.68	0.78	0.92

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700 **Table 2.** Aerosol number-size distribution parameters^{*} and real refractive index (n_r) for each aerosol
701 type in the OMI near-UV algorithm.

Aerosol Model	r_g m1 [μm]	r_g m2 [μm]	σ m1 [μm]	σ m2 [μm]	FMF	n_r	$n_{i,354/388}$
Sulfate	0.088	0.509	1.499	2.160	0.999596	1.40	1.0
Smoke	0.080	0.705	1.492	2.075	0.999795	1.50	1.2
Dust	0.052	0.670	1.697	1.806	0.995650	1.55	1.4

702 ^{*}Number-weighted particle size distribution parameters: fine and coarse mode radii (r_g m1 and r_g
703 m2) and variance (σ m1 and σ m2), number fine mode fraction (FMF).
704

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706

707 **Table 3.** Retrieval approach criteria of OMI near-UV algorithm version 1.5.3.

Surface Category	UVAI	CO (10^{18} molecules- cm^{-2})	Surface Type	Aerosol Type	Retrieval Approach
Ocean	≥ 0.5	> 2.2 NH (1.8 SH)	N/A [*]	Smoke	Two-channel
Ocean	≥ 0.5	≤ 2.2 NH (1.8 SH)	N/A	Dust	Two-channel
Ocean	< 0.5	-	-	-	No retrieval
Land	≥ 0.5	> 2.2 NH (1.8 SH)	All	Smoke	Two-channel
Land	≥ 0.5	≤ 2.2 NH (1.8 SH)	All	Dust	Two-channel
Land	< 0.5	> 2.2 NH (1.8 SH)	All	Sulfate	Two-channel
Land	< 0.5	≤ 2.2 NH (1.8 SH)	All but arid	Sulfate	Single channel
Land	< 0.5	≤ 2.2 NH (1.8 SH)	arid	Dust	Single Channel

708 ^{*}Not available.
709

710 **Table 4.** Diagonal and off-diagonal elements of the measurement error covariance matrices using
711 two different measurement matrices.

Measurement matrix	$[I_{354} \ I_{388}]^T$	$\begin{bmatrix} I_{354} & I_{354} \\ I_{388} & I_{388} \end{bmatrix}^T$
First diagonal term	$\sigma(\epsilon_{354})^2$	$\sigma(\epsilon_{388})^2$
Second diagonal term	$\sigma(\epsilon_{388})^2$	$\left(\frac{I_{354}}{I_{388}}\right)^2 \left\{ \left(\frac{\epsilon_{r,354}}{I_{354}}\right)^2 + \left(\frac{\epsilon_{r,388}}{I_{388}}\right)^2 \right\}$
Off-diagonal term	$\sigma(\epsilon_{354}, \epsilon_{388})^2$	0

712

713 **Table 5.** Error sources and their assumed magnitudes in the OMI ATBD (Algorithm Theoretical
714 Basis Documents) and this study.

Error source	Error perturbation (OMI ATBD)	Assumed value of σ for each error source in this study
Cloud Contamination	5% cloud contamination	NA ⁺
Surface Reflectivity	0.01 error in surface reflectivity	0.01 for both wavelengths
Radiometric uncertainty	SNR less than 1% Radiometric offset additive error of 1%	1% of BSDF calibration uncertainty (Dobber et al., 2006; Jaross, 2015)
	Radiometric scale factor multiplicative error of 1%	Radiometric precision provided by Level 1b data
Size distribution (mode radius)	5% increase of mode radius	0.019 for fine mode [*]
		0.510 for coarse mode [*]
Size distribution (width)	5% increase of width	0.265 for fine mode [*]
		0.307 for coarse mode [*]
Fine mode fraction	NA	0.0015 [*]
Refractive index	Increase with 0.05 for n_r	0.053 for n_r (for all wavelengths and size modes) [*]
	Increase with 0.01 for n_i	0.0047 for 354 nm n_i [*]
Aerosol Vertical Profile	Change of 1 km peak height	Change of 2 km peak height
		Change of 1 km half width
Particle shape	NA	NA

715 ^{*}Standard deviation of each parameter during the DRAGON-NE Asia 2012 campaign. The
716 parameters for n_r and n_i were obtained from 440 nm AERONET inversion products.

717 ⁺Not analyzed.

Figure captions

Figure 1. Mean 380 nm aerosol optical thickness (AOT) and 440 nm single scattering albedo and their probability density functions during the DRAGON-NE Asia 2012 campaign.

Figure 2. (a) Mean and (b) standard deviation of 388 nm AOT from the OMAERUV product in spring (March to May) 2005–2014. Panels (c) and (d) show the average and standard deviation of SSA, respectively, during the same period.

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Figure 6. Validation of 388 nm AOT against AERONET data from (a) operational products and (b) the OE-based algorithm during the DRAGON-NE Asia 2012 campaign.

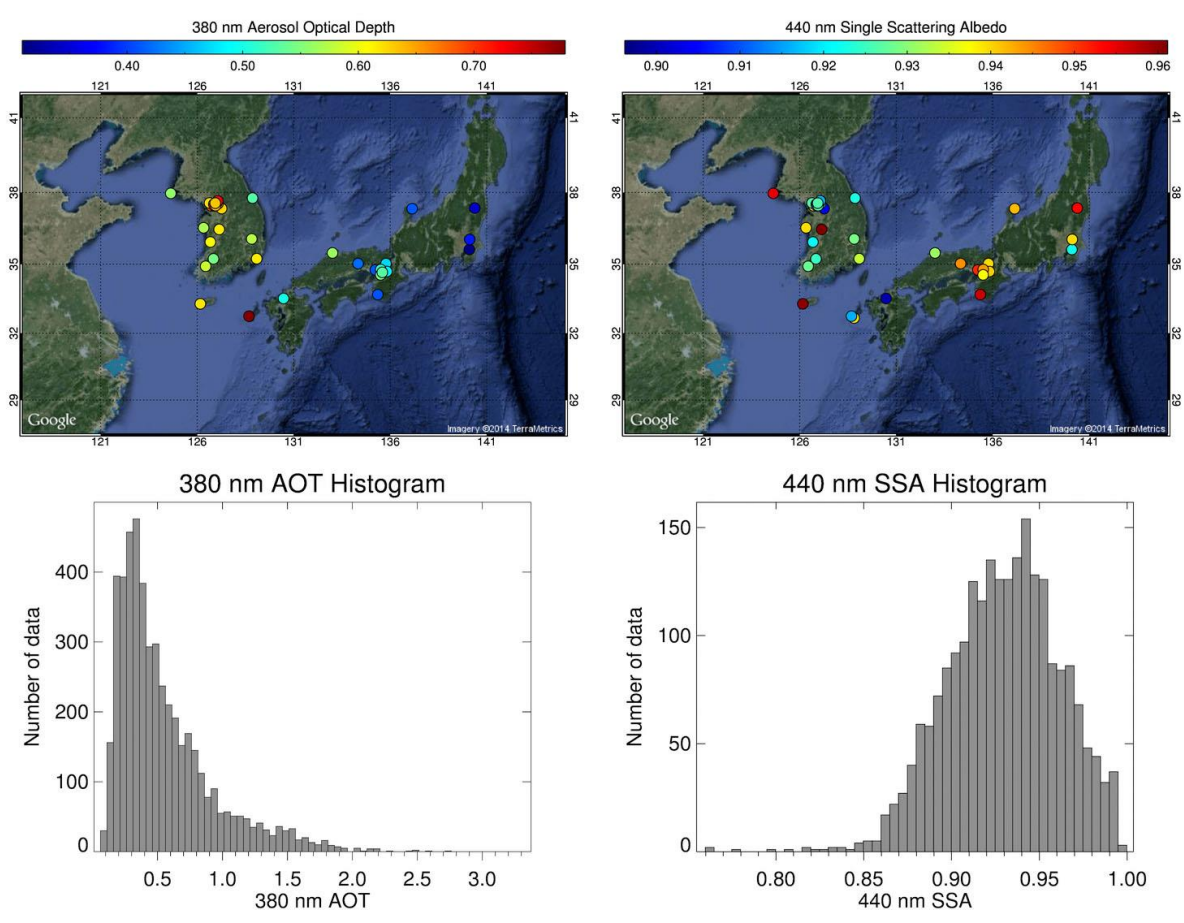
Figure 7. Comparison of the 440 nm SSA from AERONET and 388 nm SSA from (a) the operational products and (b) the OE-based algorithm, during the DRAGON-NE Asia 2012 campaign. Panels (c) and (d) compare converted 388 nm SSA from AERONET with that from (c) the operational products and (d) the OE-based algorithm.

Figure 8. Comparison between estimated uncertainties of the 388 nm AOT (x -axis) and biases of retrieved AOT from AERONET measurements (y -axis). The panels (a) and (b) are based on the operational and OE-based retrieval/error-estimation algorithm, respectively.

Figure 9. Average (gray bars) and standard deviation (black lines) of the forward model parameter errors of 388 nm (a) AOT and (b) SSA.

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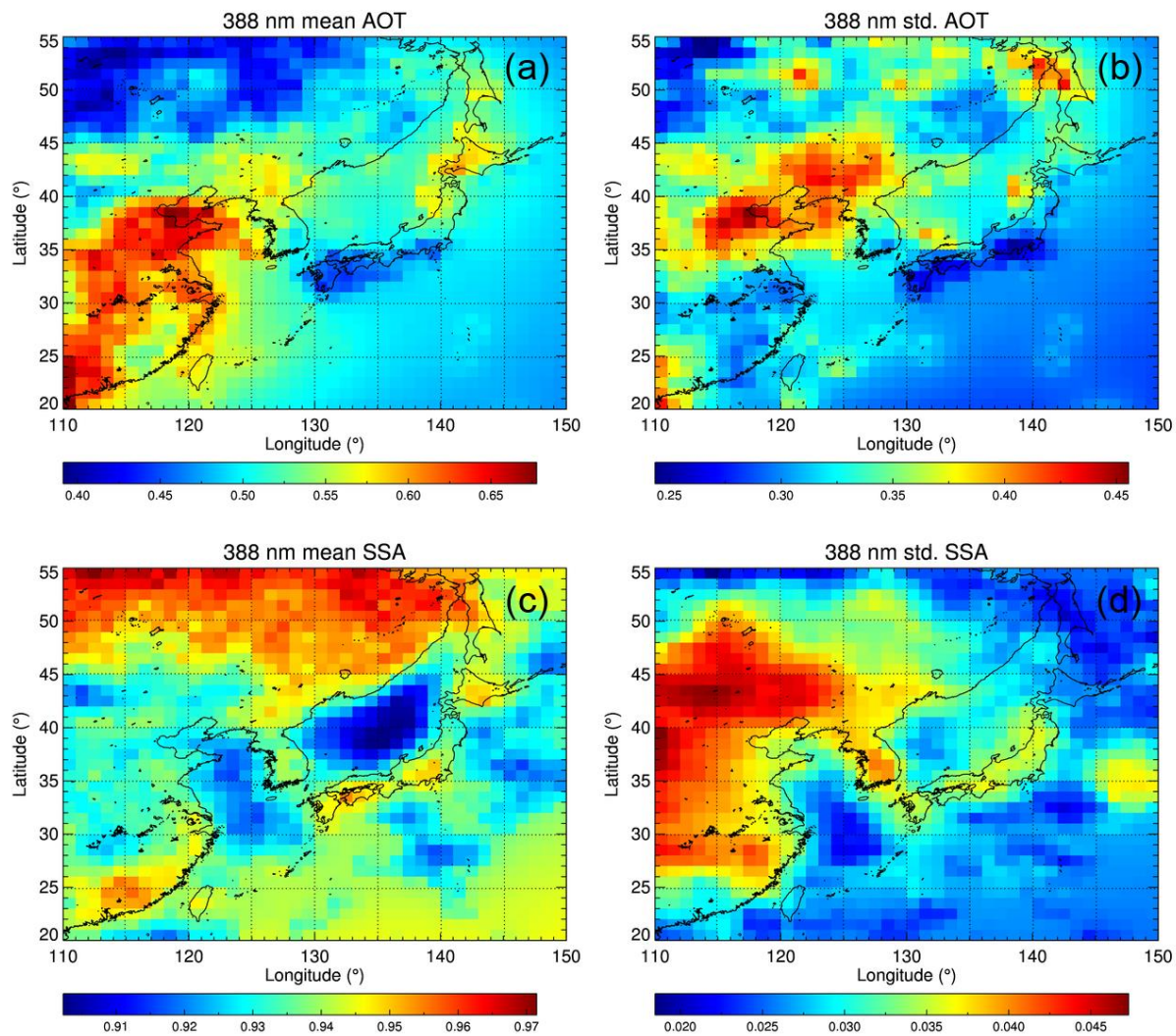


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747 their probability density functions during the DRAGON-NE Asia 2012 campaign.

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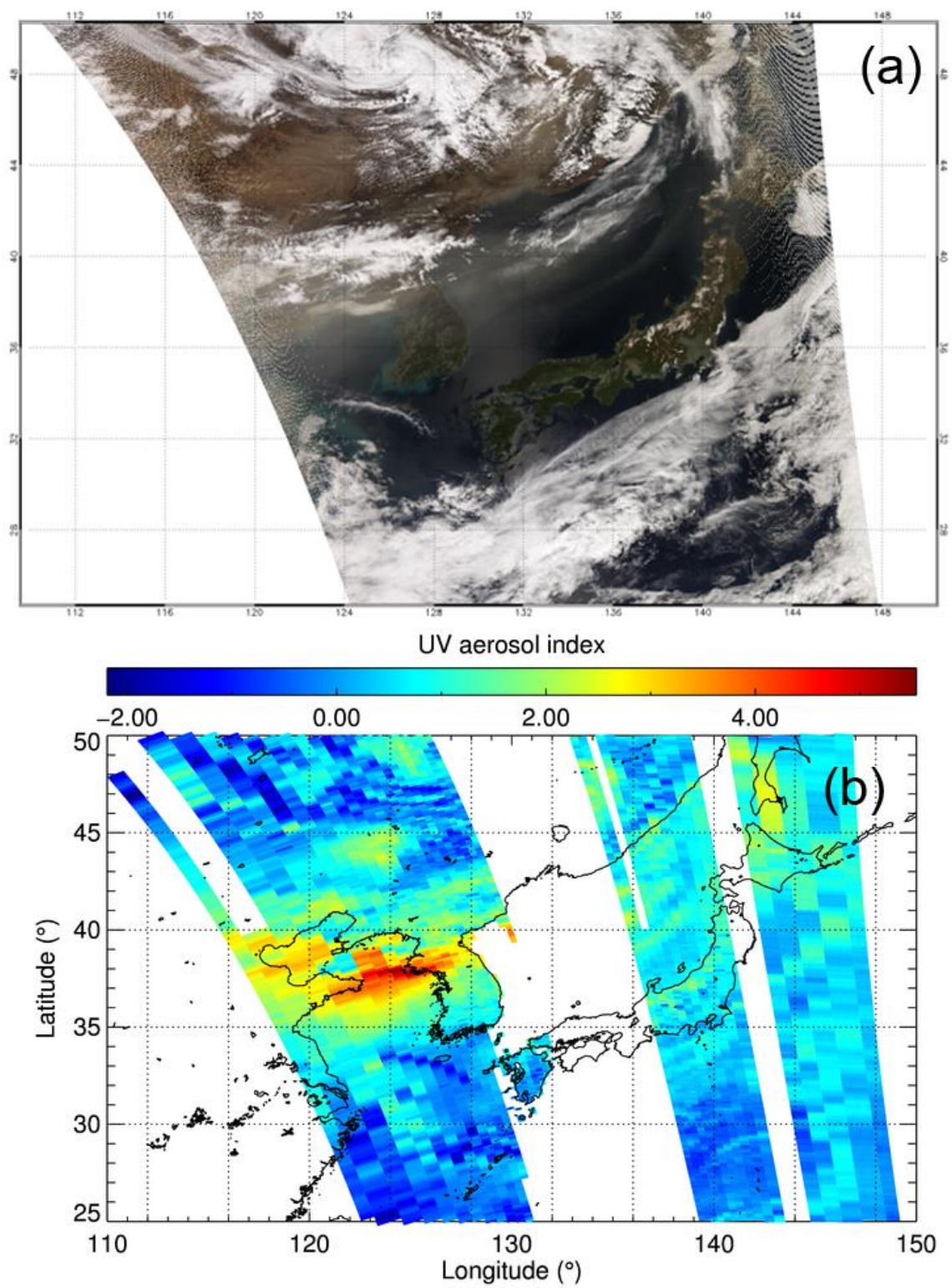
750 [Figure 2]



751
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753 spring (March to May) 2005–2014. Panels (c) and (d) show the average and standard deviation of
754 SSA, respectively, during the same period.

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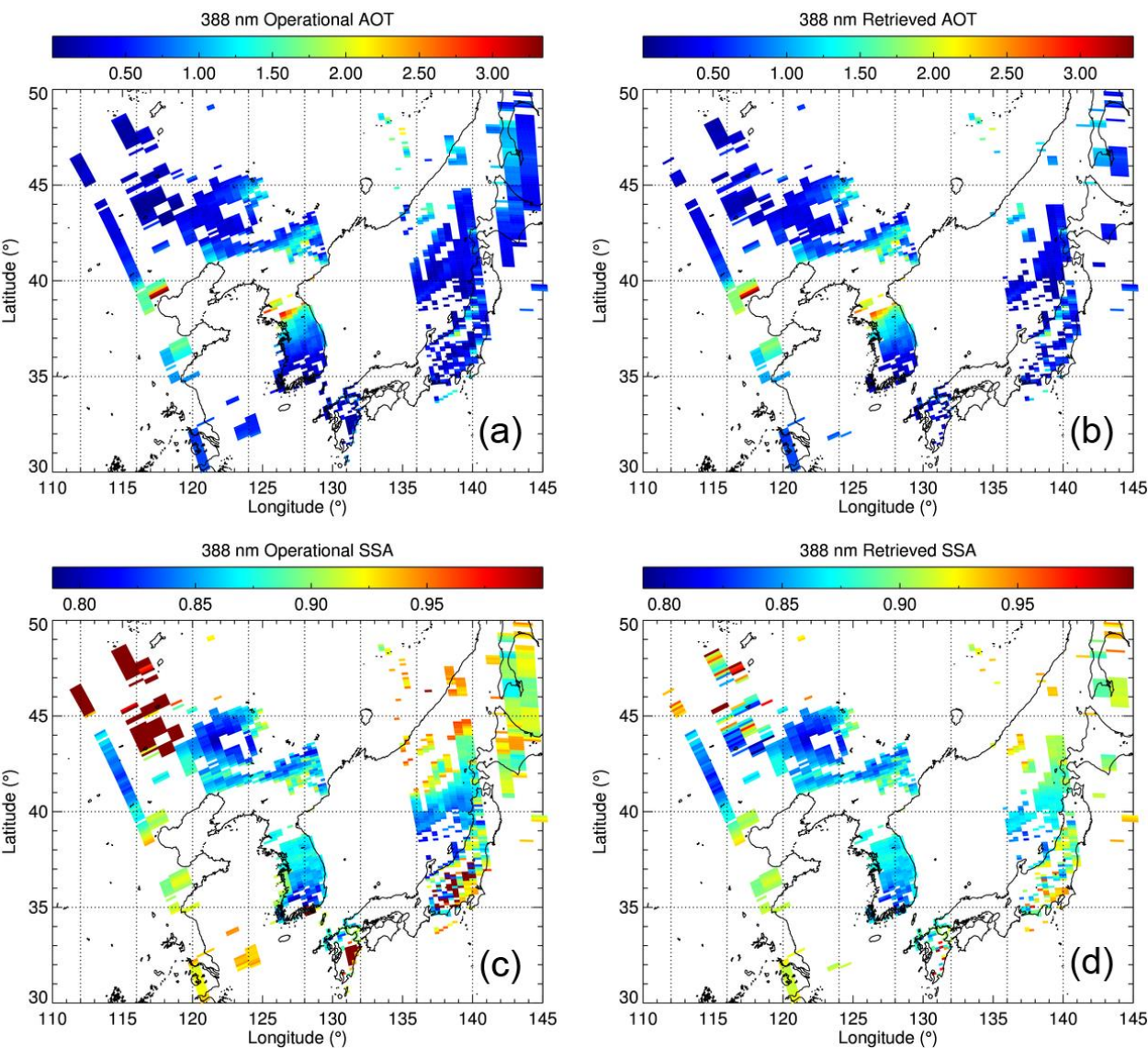


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759 Northeast Asia on 28th April 2012.

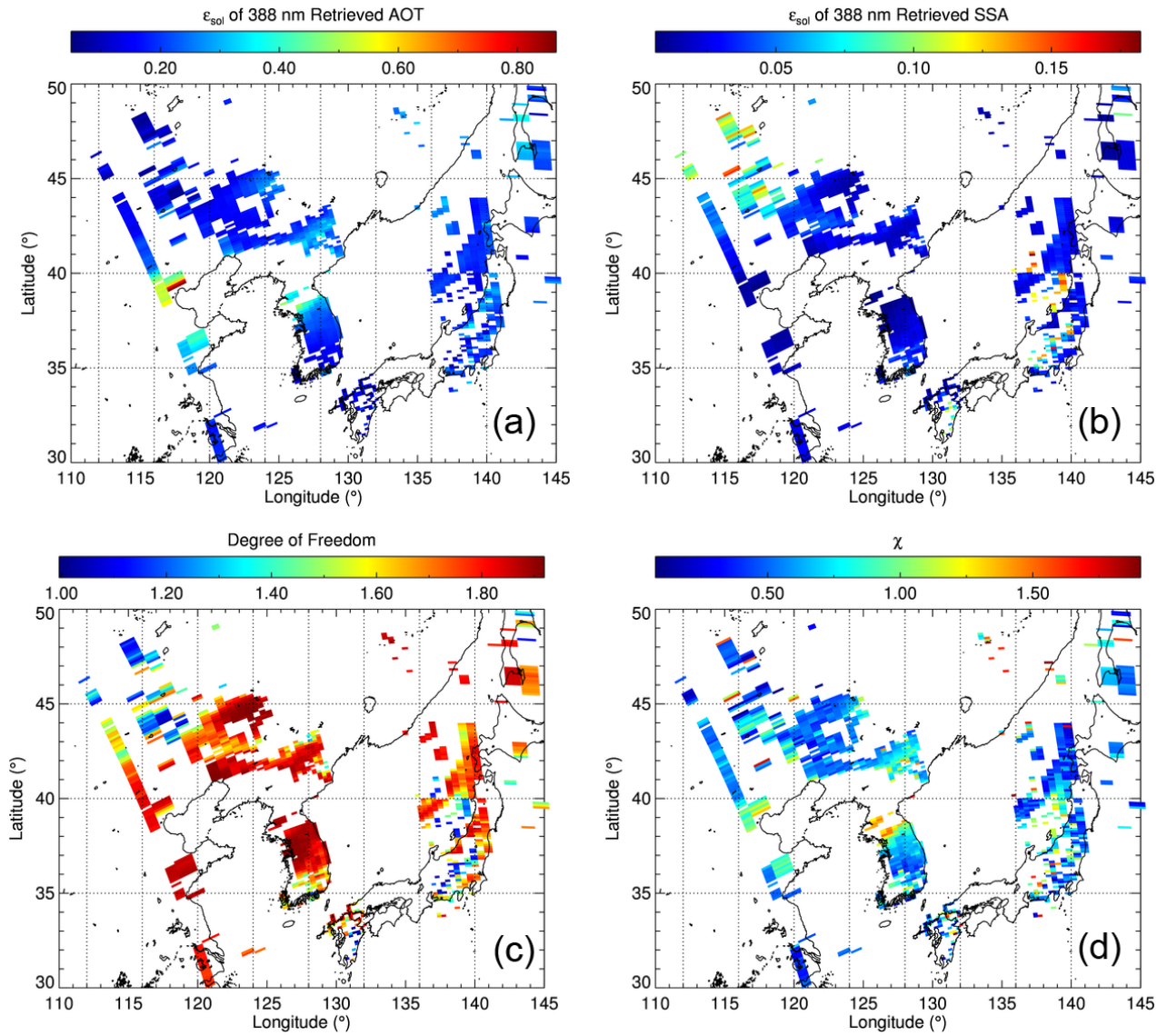
760 [Figure 4]



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762 **Figure 4.** (a) OMI operational AOT, (b) OE-based AOT, (c) operational SSA, and (d) OE-based
763 SSA on 28th April 2012 at 388 nm.

764 [Figure 5]

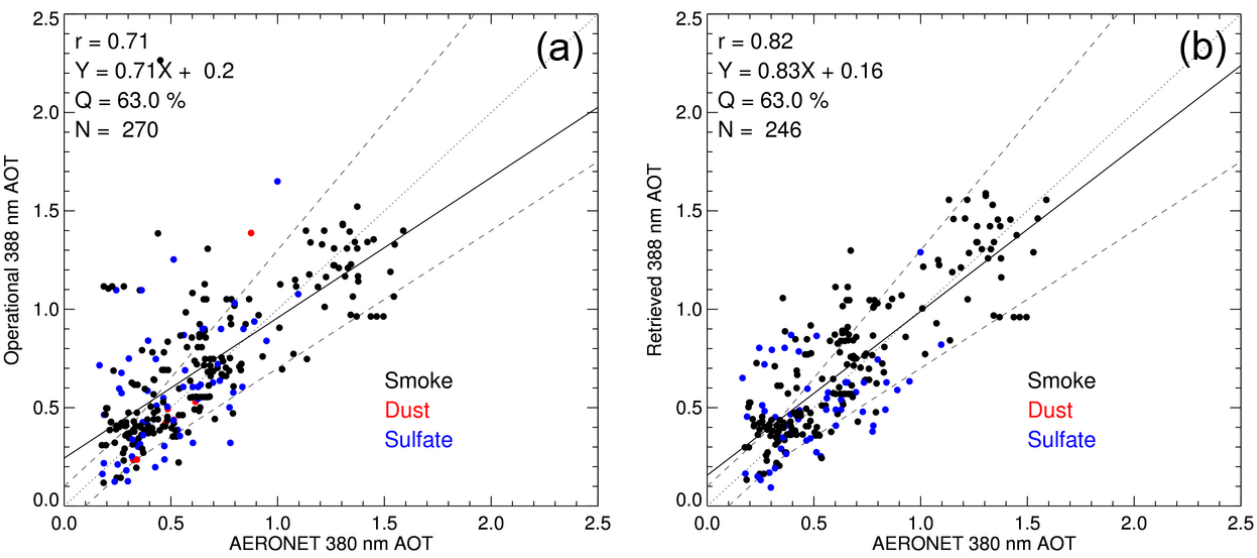


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766 **Figure 5.** Estimated solution error of (a) OE-based 388 nm AOT and (b) SSA. Panels (c) and (d)
 767 show the degrees of freedom and cost function of the retrieval, respectively.

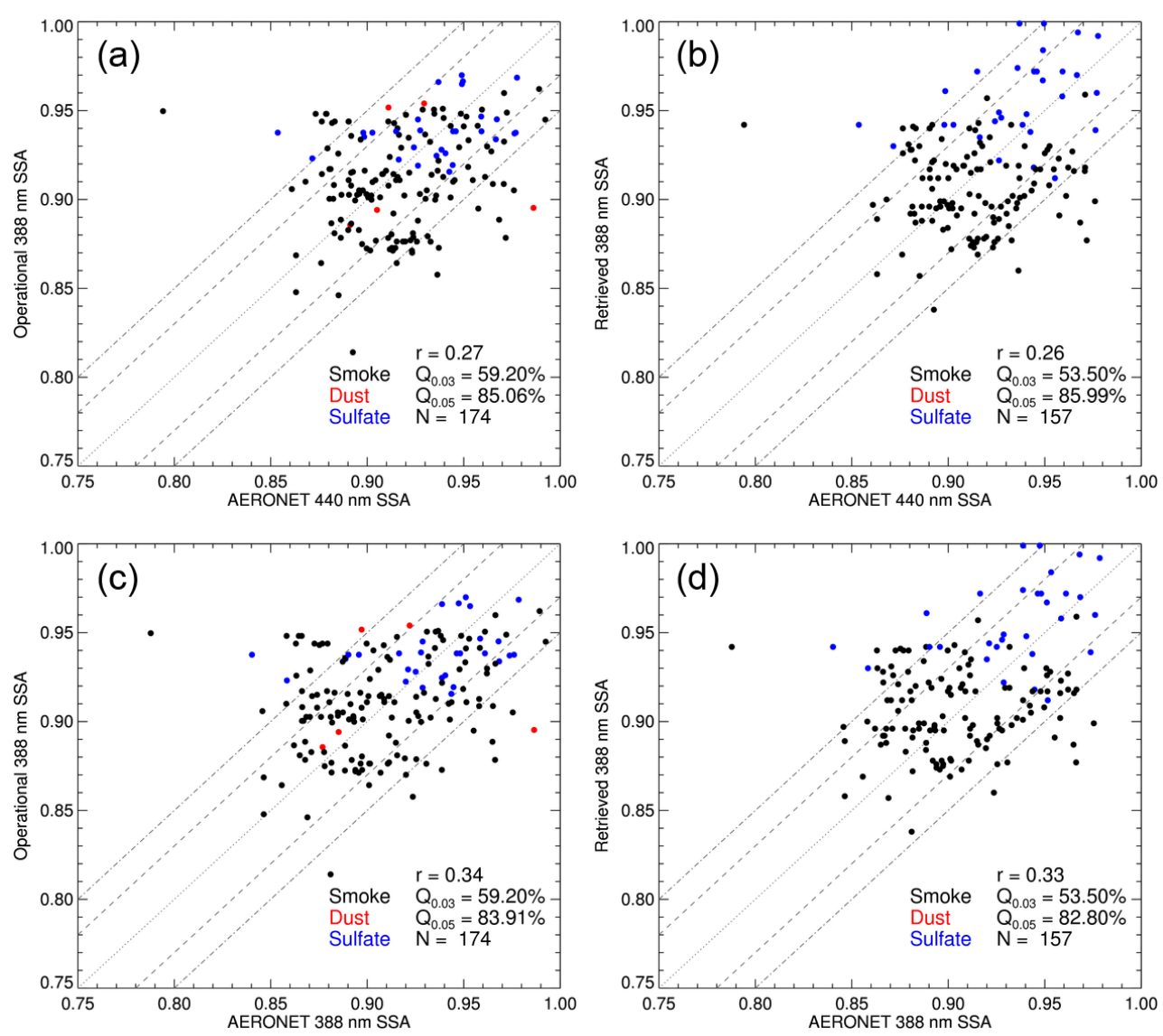
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769 [Figure 6]



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774 [Figure 7]



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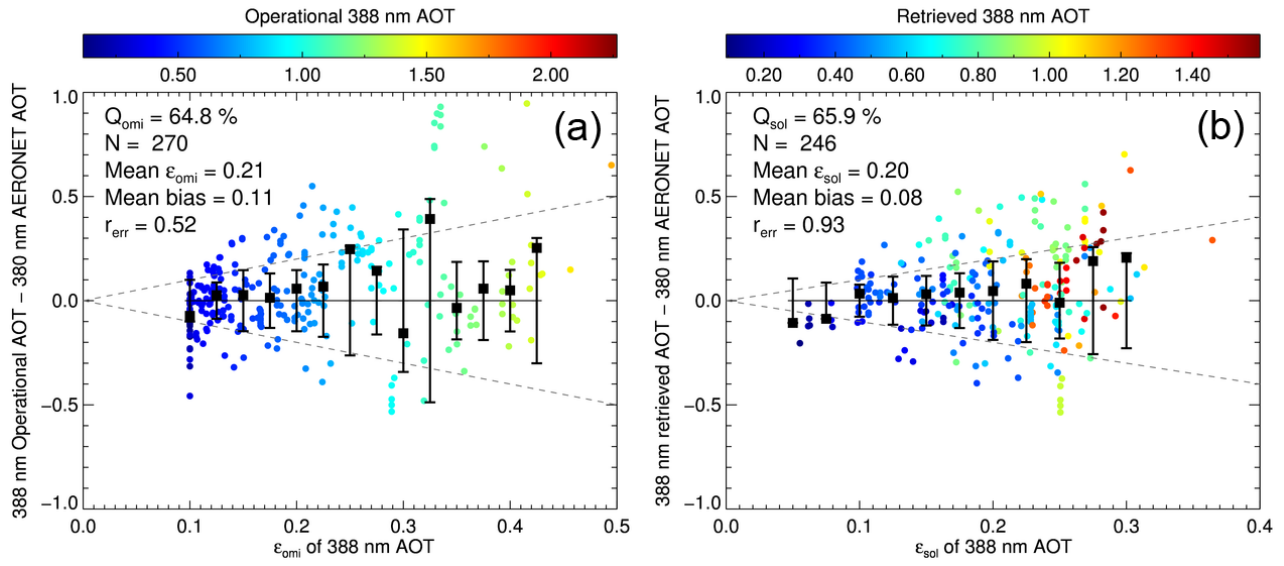
777 operational products and (b) the OE-based algorithm, during the DRAGON-NE Asia 2012

778 campaign. Panels (c) and (d) compare converted 388 nm SSA from AERONET with that from (c)

779 the operational products and (d) the OE-based algorithm.

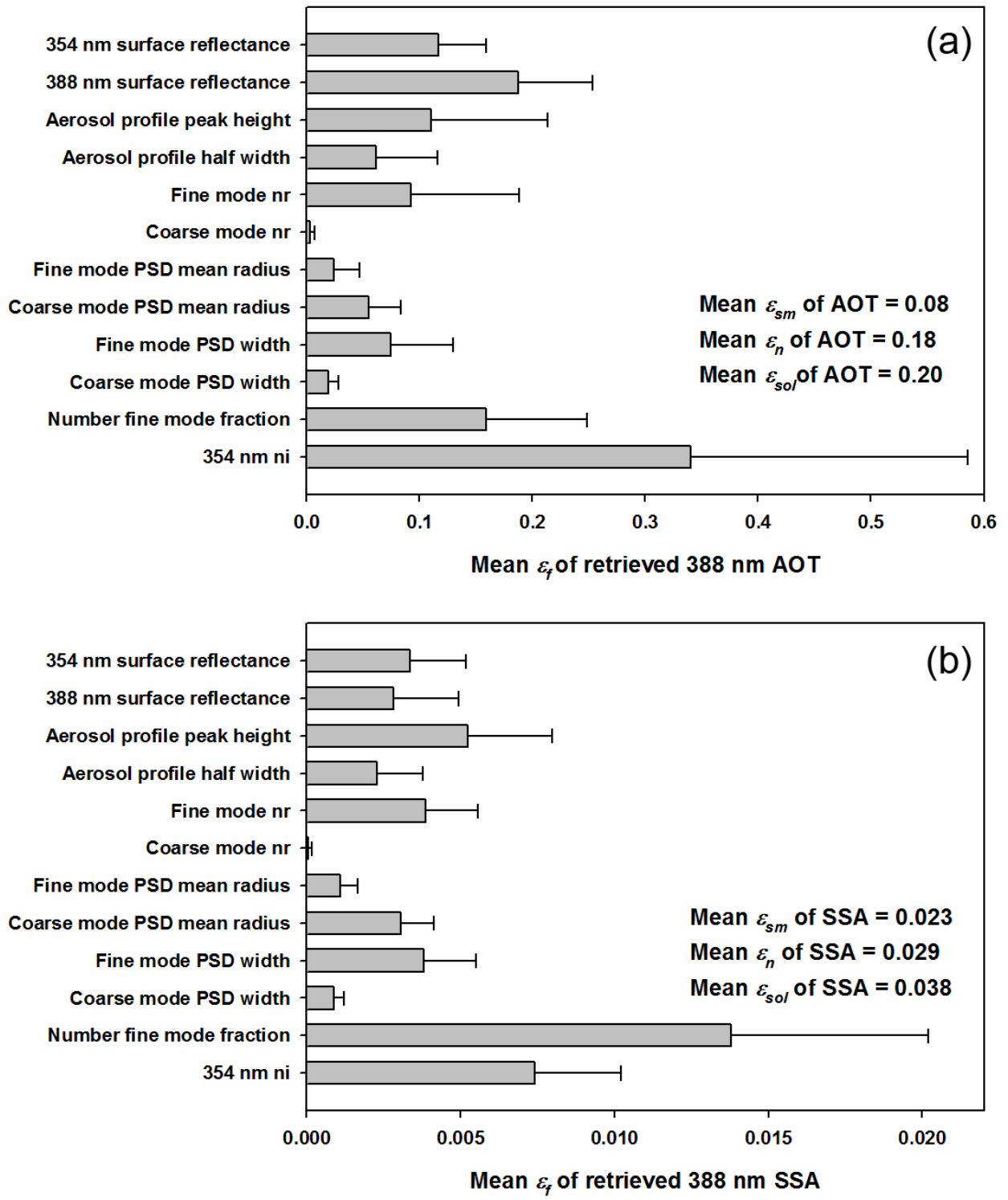
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784 retrieved AOT from AERONET measurements (y -axis). The panels (a) and (b) are based on the
785 operational and OE-based retrieval/error-estimation algorithm, respectively.
786

787 [Figure 9]



788
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