

1 **Motion-correlated flow distortion and wave-induced biases**
2 **in air-sea flux measurements from ships**

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Abstract

Direct measurements of the turbulent air-sea fluxes of momentum, heat, moisture and gases are often made using sensors mounted on ships. Ship-based turbulent wind measurements are corrected for platform motion using well established techniques, but biases at scales associated with wave and platform motion are often still apparent in the flux measurements. It has been uncertain whether this signal is due to time-varying distortion of the air flow over the platform, or to wind-wave interactions impacting the turbulence. Methods for removing such motion-scale biases from scalar measurements have previously been published but their application to momentum flux measurements remains controversial. Here we show that the measured motion-scale bias has a dependence on the horizontal ship velocity, and that a correction for it reduces the dependence of the measured momentum flux on the orientation of the ship to the wind. We conclude that the bias is due to experimental error, and that time-varying motion-dependent flow distortion is the likely source.

1. Introduction

Obtaining direct eddy covariance estimates of turbulent air-sea fluxes from ship-mounted sensors is extremely challenging. Measurements of the turbulent wind components must be corrected for the effects of platform motion and changing orientation (Edson et al., 1998; Schulze et al., 2005; Brooks, 2008; Miller et al., 2008). The ship also acts as an obstacle to the air flow forcing it to lift and change speed; this results in both the measured mean wind being biased (accelerated/decelerated) relative to the upstream flow and the effective measurement height being lower than the instrument height. This can significantly bias estimates of the 10 m neutral wind speed (U_{10n}) and the surface exchange coefficients (Yelland et al., 1998). Computational fluid dynamics (CFD) modelling studies of the flow distortion have been used to determine corrections for these mean flow distortion effects for a number of different research vessels (Yelland et al., 1998, 2002; Dupuis et al., 2003; Popinet et al., 2004; Moat et al., 2005; O’Sullivan et al., 2013, 2015) and also generic corrections for commercial vessels that report meteorological measurements (Moat et al., 2006a,b).

The modelled corrections show a strong dependence on the relative wind direction (Yelland et al., 2002; Dupuis et al., 2003) and a much weaker dependence on wind speed, but in general have been determined only for ships with zero pitch and roll angles. Weill et al. (2003) and Brut et al. (2005) reported on experiments with a 1/60 scale physical model of the *RV La Thalassa* to investigate the effect of pitch and roll angles on the mean flow distortion.

47 They found the tilt of the mean streamline to vary by more than 1° and the mean wind speed
48 by up to 12% for pitch angles between $\pm 10^\circ$; these effects were asymmetric about zero pitch.
49 Roll angle had only a small impact on the measured wind speed, about 1% for roll of up to
50 10° , but this was examined for bow-on flows only and a larger impact might be expected for
51 flows with a significant beam-on component. Comparison of in situ measurements from sonic
52 anemometers, the physical model tests, and CFD modelling also revealed that the foremast
53 itself, along with the instruments and electronics enclosures mounted on it, had a significant
54 impact on flow distortion at the location of the sonic anemometer.

55 The studies of flow distortion cited above addressed only the mean flow for a fixed
56 orientation of the ship with respect to the mean streamline; to the best of our knowledge no
57 studies have investigated the effect of time-varying flow distortion as ship attitude changes.
58 That the time-varying flow distortion has an impact can, however, be inferred from reported
59 biases of ship-based eddy covariance measurements. Edson et al. (1998) compared eddy
60 covariance estimates of the kinematic wind stress from two ships with those from a small
61 catamaran and from the stable research platform FLIP. They found that the ship-based
62 estimates were on average 15% higher than those from FLIP and the catamaran. They argued
63 that the difference resulted from flow distortion over the ship rather than from inadequate
64 motion correction because the catamaran experienced more severe platform motion. Pedreros
65 et al. (2003) similarly found momentum flux estimates from a ship to be 18% higher than
66 estimates from a nearby air-sea interaction spar buoy. Evidence of such biases, ascribed to
67 flow distortion, led to the exclusion of ship-based direct flux measurements from the most
68 recent update of the COARE bulk air-sea flux algorithm (v3.5; Edson et al., 2013).

69 Features in cospectra that manifest as significant deviations from the expected
70 spectral form (e.g. Kaimal et al., 1972) at frequencies associated with waves and platform
71 motion have been reported in observations of momentum fluxes measured from FLIP (Miller
72 et al., 2008) and from fixed platforms (Deleonibus, 1971) and towers (Drennan et al., 1999).
73 A decrease in the magnitude of the feature with height led Miller et al. (2008) to ascribe its
74 source to interactions between the waves and atmospheric turbulence. The authors also note
75 that the anemometers used were not co-mounted with inertial motion units and their tilt from
76 horizontal was determined using the planar fit method; errors in the determined tilt or in
77 estimation of anemometer and inertial motion unit alignment could also contribute to the
78 observed features via incomplete correction for platform motion (Brooks, 2008; Landwehr et
79 al., 2015). Edson et al. (2013) analysed wind profile measurements from three field
80 campaigns and found little evidence of wave influence on winds at heights above 4 m in sea

81 conditions with $c_p/U_{10n} < 2.5$, where c_p is the wave phase speed. In general, reported motion-
82 scale signals in the turbulence have been observed in measurements made either at heights
83 below 10 m (Deleonibus, 1971; Miller et al., 2008) or in conditions of fast, high swell where
84 $c_p/U_{10n} \approx 2$ and $H_{swell} \gg H_{wind}$, and H_{swell} and H_{wind} are the significant wave heights of the
85 swell and wind-wave components of the wave field respectively (Drennan et al., 1999).
86 Recent results from Large Eddy Simulations over moving wave fields also suggest that, in
87 developing sea conditions waves are not expected to significantly influence turbulent winds
88 at heights of more than about 10 m (Sullivan et al., 2014). In summary, the wave field is only
89 expected to influence the turbulent winds near the surface or in conditions where swell
90 dominates the wave field.

91 High frequency gas concentration measurements for studies of air-sea exchange have
92 been shown to suffer significant motion-correlated biases resulting from the hydrostatic
93 pressure change with vertical displacement (Miller et al., 2010), and potentially from
94 mechanical sensitivities of the sensors themselves (McGillis et al., 2001; Yelland et al., 2009;
95 Miller et al., 2010). These biases cause distortions of the cospectra between the vertical wind
96 component and gas concentration (Edson et al., 2011) apparent in the cospectra at frequencies
97 associated with the platform motion, and several recent studies have applied motion
98 decorrelation algorithms to remove this signal (Miller et al., 2010; Edson et al., 2011;
99 Blomquist et al., 2014).

100 Such an approach can also correct the apparent motion-scale bias in the momentum
101 flux, but is controversial since, as discussed above, there are circumstances in which a real
102 wave-correlated signal may be expected in the turbulence measurements. Here we present
103 measurements which demonstrate a significant motion-scale feature in momentum flux
104 measurements from a research ship. We show the impact of applying a simple regression
105 procedure to remove the bias, and provide evidence that suggests the source of the bias is
106 time-varying flow distortion correlated with ship motion and attitude.

107

108 **2. Data**

109 The measurements were made on the *RRS James Clark Ross* as part of the Waves, Aerosol
110 and Gas Exchange Study (WAGES), a programme of near-continuous measurements using
111 the autonomous AutoFlux system (Yelland et al., 2009). Turbulent wind components were
112 measured by a Gill R3 sonic anemometer installed above the forward, starboard corner of the
113 ship's foremast platform (Fig. 1). The measurement volume was approximately 16.5 m above
114 sea level. Platform motion was measured with a Systron Donner MotionPak Mk II, mounted

115 rigidly at the base of the anemometer and sampled synchronously with it. Wave field
 116 measurements were made using a WAVEX X-band radar installed above the bridge top. The
 117 WAVEX system obtains directional wave spectra and mean wave parameters every five
 118 minutes.

119 The fast-response instrumentation operated at 20 Hz, and flux estimates were
 120 calculated over 30-minute periods. The raw wind and motion measurements were first
 121 despiked and the wind components corrected for platform motion using the complementary
 122 filtering approach of Edson et al. (1998). The motion correction algorithm set out in Edson et
 123 al. (1998) and as usually applied corrects the measured horizontal winds for low frequency
 124 horizontal motions (ship's underway velocity) in the earth frame. This neglects the aliasing of
 125 the ship's horizontal speed into the vertical imposed by the non-horizontal mean streamline at
 126 the point of measurement due to flow distortion over the ship. The true vertical wind speed,
 127 w_{true} is determined from the measured, motion-corrected vertical wind, w_{rel} , and the
 128 horizontal true and relative winds (U_{true} and U_{rel}) as

$$129 \quad 130 \quad w_{true} = w_{rel} - \left(\overline{w_{rel}} \times \left[1 - \frac{\overline{U_{true}}}{\overline{U_{rel}}} \right] \right) \quad (1)$$

131
 132 where overbars indicate a time average (Tupman, 2013). The derivation of Eq. (1) and the
 133 impact of applying this correction are described in Appendix A. This correction addresses the
 134 same source of measurement error as that recently described by Landwehr et al. (2015) who
 135 address it by applying corrections for the ship's low-frequency horizontal velocity after
 136 rotating the ship-relative winds (corrected for high frequency motions) into the reference
 137 frame of the mean streamline for each flux averaging period.

138 After motion correction, each 30-minute record is rotated into a reference frame
 139 aligned with the mean streamline, wind components were linearly detrended, and eddy
 140 covariance momentum fluxes calculated. CFD modelling of the air flow over the James Clark
 141 Ross was initially undertaken by Yelland et al. (2002), but only for flow on to the bow; we
 142 have extended the CFD study for a much wider range of relative wind directions and the
 143 results were used to determine direction-dependent corrections to the mean (30-minute
 144 averaged) relative wind speed and measurement height. The new CFD study is documented
 145 in Moat and Yelland (2015) and the primary results reproduced here in Appendix B. The
 146 modelled wind speed bias at the sensor location varied between -0.9% and 8.4% for wind
 147 directions between 20° to port of the bow and 120° to starboard, and the height by which the

148 flow was raised varied between 1.3 m and 3.2 m. Wind directions beyond 20° to port of the
149 bow were affected by small-scale obstructions on the foremast platform that are not included
150 in the CFD model; these wind directions are thus excluded from the following analysis. After
151 applying the corrections, the measured winds were corrected to 10 m height and neutral
152 stability using the Businger-Dyer relationships (Businger, 1988) and the 10 m neutral drag
153 coefficient, CD_{10n} , was calculated from U_{10n} and the momentum flux estimates.

154 The measurements used here were obtained between 09 January and 16 August, 2013
155 in locations throughout the North and South Atlantic, the Southern Ocean and the Arctic
156 Ocean, at latitudes ranging from 62°S to 75°N. After excluding measurement periods when
157 the ship was within sea ice, there were 2920 individual flux estimates available for analysis.
158 Flux estimates were then rejected from the analysis where there was excessive ship
159 manoeuvring, where flux quality control criteria were failed (Foken and Wichura, 1996;
160 Vickers and Mahrt, 1997), and when the air temperature was less than 2°C when ice build up
161 may affect the sensors. Of the remaining 1054 flux estimates, 80 were removed as outliers
162 ($CD_{10n} > 5 \times 10^{-3}$). Unless otherwise indicated, mean relative wind direction limits of 20° to
163 port, and 50° to starboard of the bow were applied, a condition met by 499 flux estimates. Of
164 the removed outliers, 38 lay within acceptable relative wind direction limits; of these, 6 were
165 at winds speeds of 6 m.s⁻¹ or greater.

166

167 **3. Removal of the ship motion-scale signal**

168 Momentum flux cospectra and ogives for U_{10n} between 10 and 14 m s⁻¹, normalised (by $f/u.^2$
169 and $1/u.^2$ respectively, where f is frequency and $u.$ is the friction velocity) and averaged, are
170 shown in Fig. 2. The cospectra and ogives differ from the typical forms obtained from
171 experiments over land (e.g. Kaimal et al., 1972) at frequencies between approximately 0.06
172 and 0.25 Hz (0.09 and 0.37 in the non-dimensionalised frequency shown in Fig. 2), where a
173 significant anomalous signal is present. These are frequencies associated with surface waves
174 and with the platform motion that results, hence we term the cospectral signal at these
175 frequencies the motion-scale signal.

176 At wind speeds above 7 m s⁻¹, the CD_{10n} measurements are biased high compared
177 with previous results (Fig. 3). The bias relative to the eddy covariance-based parameterisation
178 of Smith, (1980) increases with wind speed from approximately 20% at 8 m s⁻¹ to 60% at 20
179 m s⁻¹. Note that the Smith (1980) parameterisation was derived from eddy covariance
180 measurements made from a slim floating tower moored so as to minimise platform motion

181 and inducing minimal flow distortion. The bias is smaller when compared to the COARE 3.0
182 (Fairall et al., 2003) or COARE 3.5 (Edson et al., 2013) bulk algorithms.

183 The motion-scale signal can be removed from the vertical wind component, and a
184 corrected vertical wind, w_{MSC} , obtained via a simple regression method:

$$185 \quad \quad \quad 186 \quad \quad \quad w'_{MSC} = w'_{true} - \alpha_1 acc'_z - \alpha_2 vel'_z \quad (2)$$

187
188 where acc_z and vel_z are the platform's vertical acceleration and velocity, measured at the base
189 of the sonic anemometer, and primes denote fluctuations determined from Reynolds
190 decomposition. The coefficients α_1 and α_2 are determined here by regression for each 30-
191 minute flux measurement period. This algorithm, which we term the motion-scale correction
192 (MSC), is based on the regression corrections of Yelland et al., (2009) and Miller et al.,
193 (2010). It is also similar to the motion decorrelation algorithm given in a spectral formulation
194 by Edson et al. (2011), originally utilised to remove motion biases from CO₂ flux cospectra,
195 and here termed the MSC_f. The MSC_f algorithm coefficients are defined as the ratio of
196 covariances of vertical wind and platform motion to variances of platform motion. The MSC
197 and MSC_f methods give almost identical results (Fig. 2).

198 Applying the MSC algorithm removes the motion-scale signal (Fig. 2) and results in a
199 20% to 30% decrease in CD_{10n} for wind speeds above 7 m s⁻¹ and absolute values similar to
200 those of COARE 3.0 or 3.5 (Fig. 3). The signal removed is similar in size and of the same
201 sign as the biases in ship-based momentum flux measurements reported by Edson et al.
202 (1998) and Dupuis et al. (2003).

203 Applying the MSC to the along-wind component as well as the vertical component
204 makes an insignificant (<<1%) additional difference to the measured flux (shown as MSC_{uv}
205 in Figs. 2 and 3). Interpolating the measured cospectra across the motion-scale frequencies
206 gives similar results to the MSC algorithm under most conditions (shown as "interpolated" in
207 Figs. 2 and 3: Prytherch, 2011; Tupman, 2013). However, interpolation requires selection of
208 appropriate frequencies to interpolate between, in this case 0.04 and 0.4 Hz (0.06 and 0.59 in
209 the non-dimensionalised frequency shown in Fig. 2), and is not dependent on a physical
210 variable related to the presumed source of the error (platform motion-dependent flow
211 distortion). For these reasons, correction using the MSC algorithm is preferable.

212

213 **4. Discussion**

214 Following application of the MSC the cospectral shape matches the Kaimal form expected.
215 This suggests that the motion-scale bias is being effectively removed.

216 The MSC also results in drag coefficients that lie within the range of previous
217 parameterisations. At the highest wind speeds (over 15 m s^{-1}) the parameterisations begin to
218 diverge significantly and the WAGES CD_{10m} are larger than those given by Smith (1980) and
219 lie between those of COARE 3.0 and 3.5. It should be noted that COARE 3.0 and 3.5 are
220 both defined using wind speeds in the frame of reference of the surface currents (see
221 Appendix in Edson et al., 2013), rather than in the earth frame of reference as used by Smith
222 (1980). Surface current measurements were not available for the WAGES data. For surface
223 currents aligned with the prevailing wind direction, adopting a surface current frame of
224 reference would lead to a small apparent increase in the drag coefficients presented here.

225 While several previous studies have ascribed a high bias in drag coefficient estimates
226 from ships to flow distortion (Edson et al. 1998, 2013; Pedreros et al. 2003), they have not
227 examined the effect in detail. Inaccurate tilt estimation, a related source of error, may also
228 contribute to this bias, particularly at low wind speeds (Landwehr et al., 2015). Few other
229 studies have discussed such biases at all, and it seems likely that the severity of any motion-
230 correlated bias is highly dependent on individual platforms and instrument installations in the
231 same manner as the mean flow distortion. The bias is potentially worse here than in many
232 other studies; the sonic anemometer is mounted lower on the foremast than would be ideal
233 because the long-term measurement programme made it necessary to be able to service the
234 instruments easily and without access to a crane. There are also a greater number of small-
235 scale obstructions such as searchlights near to the measurement point than would be the case
236 on lattice style masts often deployed on dedicated flux measurement campaigns. Because the
237 measurements are continuous and autonomous, a large fraction of our data is also obtained
238 with the ship underway. In contrast, dedicated eddy covariance studies of air-sea exchange
239 would usually focus almost exclusively on measurements made on station when ship motion
240 is substantially less than when underway. Finally it is possible that such biases are present in
241 some fraction of the measurements of many studies, but are excluded from final analysis by
242 quality control procedures without a close examination of the bias being made. Many studies
243 with modest data volumes have quality controlled the individual flux estimates via a visual
244 inspection of the ogive curves, rejecting those that do not closely match the expected form
245 (e.g. Fairall et al. 1997; Norris et al. 2012).

246 As discussed in Sect. 1 above, there is evidence from previous studies that the
247 influence of the wave field on the turbulent winds should be small at heights above some

248 limit which is assumed to be related to the wave properties: values between 4 and 10 m have
249 been cited (Miller et al. 2008, Sullivan et al. 2014). The Sullivan et al. (2014) results
250 correspond to a height of order 1.5 times the significant wave height. Real wave effects are
251 thus expected to be negligible for typical measurement heights of ship-based sensors (15-20
252 m) under most conditions. Below we provide more direct evidence that the wave-scale signal
253 seen in the WAGES data is due, in large part at least, to the effects of flow distortion over a
254 moving platform.

255

256 **4.1 Motion dependence of the streamline**

257 The angle to the horizontal of the airflow measured at the sonic anemometer site was found to
258 be dependent on the vertical motion of the ship (Fig. 4). Perturbations in the tilt of the
259 streamline are approximately in phase with acc_z , out of phase with the vertical displacement
260 and pitch, and lead vel_z by about 90° . There are multiple processes that may affect the
261 streamline orientation as the ship moves over the waves:

- 262 • Vertical displacement of the ship changes the vertical extent of the obstacle that the
263 ship presents to the flow and the relative height of the measurement volume with
264 respect to that of the bow above the water line.
- 265 • The ship's pitch similarly changes both the effective size of the obstacle presented to
266 the flow and the relative location of the sonic anemometer within the distorted flow
267 above the bow.
- 268 • Vertical motion of the ship will force the overlying air to move.

269

270 In the example here for 15 m s^{-1} , bow-on winds, the airflow tilt varies by about $\pm 3^\circ$ around a
271 mean of approximately 10° . The various parameters shown in Fig. 4a are all inter-dependent,
272 but streamline tilt showed slightly more consistent trends with the velocity and acceleration
273 parameters than with displacement or pitch suggesting that "pumping" of the air above the
274 moving deck may be the dominant effect.

275

276 **4.2 Characteristic frequencies of spectral features**

277 For a platform moving through a wave field aligned with the direction of travel, it would be
278 expected that the frequency of ship motion forced by the waves would differ from that for a
279 ship on station with no mean horizontal velocity. The change could be of either sign
280 depending on the ratio of wavelength to the length of the ship, with an increase in frequency

281 for wavelengths much longer than the ship. The measured frequency of atmospheric turbulent
 282 structures would also be shifted to higher frequencies relative to those measured when on
 283 station. The nature of the frequency shift should differ for turbulent air motions, which advect
 284 with the wind and have a ship-relative velocity equal to the sum of wind and ship speeds, and
 285 wave-correlated features in the turbulence field which are phase-locked to the surface waves
 286 (Sullivan et al. 2000; 2008; 2014), and will have a ship-relative velocity of the sum of wave-
 287 phase and ship speeds. A signal due to real wind-wave interaction should thus appear at a
 288 different frequency to that from a ship motion-induced measurement bias.

289 Figure 5 shows a comparison of the power spectral density of platform vertical
 290 velocity (S_{velz} , Fig. 5b) and frequency weighted cospectral densities for the streamwise
 291 momentum flux (normalised by u .) both for periods during which the ship was on station
 292 ($V_{ship} < 1 \text{ m s}^{-1}$, where V_{ship} is the speed of the ship) and when underway ($V_{ship} > 5 \text{ m s}^{-1}$). The
 293 cospectra are shown after applying the standard motion correction to the measured turbulent
 294 velocity components, but without applying the MSC correction. Also shown are the spectral
 295 densities of the surface wave field (Fig. 5a). The wave radar provides wave spectra in the
 296 earth frame, corrected for ship speed; in order to compare these directly with the measured
 297 turbulence and ship-motion spectra when underway we need to transform them into a
 298 reference frame moving with the ship. This is achieved by plotting against a modified
 299 frequency, $f_m = f_0(c_p + V_{ship})/c_p$, where f_m is the frequency that would be measured in the ship
 300 reference frame and f_0 is the true frequency in the earth frame. The periods chosen all have
 301 bow-on winds, wind speeds of between 10 and 12 m s^{-1} and similar sea states: the (true) mean
 302 peaks of the mean WAVEX-derived non-directional wave spectra (S_{zwave}) are 0.120 and 0.110
 303 Hz, and mean significant wave heights are 4.73 m and 3.51 m for the stationary and
 304 underway periods respectively.

305 For the on station measurements, the peak in the momentum flux cospectra (no MSC,
 306 Fig. 5c) is at 0.113 Hz, which matches that of the peak in ship vertical velocity (Fig. 5b) and
 307 is at slightly lower frequency than the peak in the ship-frame surface wave spectra (0.120 Hz,
 308 Fig. 5a). For the underway cases the peak in the ship-frame wave spectra is shifted to higher
 309 frequency (0.163 Hz) compared to the true spectra. The peak in the ship motion spectrum
 310 (0.148 Hz) is again lower than that of the wave spectrum and by a larger margin than for the
 311 on station case. The peak in the momentum flux cospectrum at 0.153 Hz is much closer to
 312 that of the ship motion than that of the wave spectrum.

313 The correspondence of the peak in momentum flux cospectra with that of the ship

314 motion rather than that of the wave field suggest that the residual signal after motion
315 correction is an artefact of motion-correlated flow distortion rather than a result of a real
316 wave-correlated signal in the turbulence.

317

318 **4.3 Directional dependence of drag coefficient bias**

319 Mean flow distortion is strongly dependent on relative wind direction (Yelland et al., 1998),
320 even for a motionless ship with zero pitch and roll angles. The dependence of the calculated
321 drag coefficients on relative wind direction before and after applying the MSC algorithm is
322 shown in Fig. 6. First, a linear fit was made between the drag coefficient and wind speed data
323 obtained for wind directions between -20 and +50 degrees of the bow. Then the drag
324 coefficient anomalies (individual minus fit) were calculated and averaged into 10° relative
325 wind direction bins, and the results were plotted against relative wind direction. It can be seen
326 that prior to applying the MSC algorithm the drag coefficient anomalies have a significant
327 dependence on relative wind direction, and that application of the algorithm significantly
328 reduces this dependence. For completeness the results are also shown without first applying
329 the direction-dependent CFD-derived correction to the mean, 30-minute averaged, wind
330 speed; this also reduces the dependence of the drag coefficient on relative wind direction.

331 Application of the MSC and the mean CFD correction does not completely remove all
332 dependence of the drag coefficient on relative wind direction. This suggests that one or both
333 corrections may need refinement. In the case of the MSC algorithm, the effect of the roll of
334 the ship is likely to become significant when the wind direction is beam-on rather than bow-
335 on. In the case of the CFD correction to the mean wind speed, the model of the ship geometry
336 may have to be refined to take into account local flow distortion caused by small objects
337 mounted on the foremast, close to the anemometer. These are areas for future investigation.

338

339 **5. Conclusions**

340 Methods for removal of motion-correlated signals from fast-response gas measurements
341 made onboard moving platforms have become more commonly applied in recent years;
342 however, these techniques remain controversial when applied to fast-response winds for the
343 purpose of momentum flux calculation. The results here demonstrate these methods and their
344 impact on ship-based momentum flux measurements where a significant motion-correlated
345 bias is present in the motion-corrected cospectra. The motion-correlated signals are shown to
346 be dependent on platform velocity relative to the wave field. In addition, the dependence of
347 the flux on wind direction relative to the ship is reduced after applying the correction

348 methods. These results suggest that the motion-correlated signal is due to the effects of time-
349 varying flow distortion. Further investigation is required to resolve the details of the physical
350 processes involved.

351 The recent revision of the COARE bulk flux algorithm (COARE 3.5, Edson et al.
352 2013) is determined only from data from platforms other than ships (buoys, towers, FLIP).
353 These data all require motion correction, and Bigorre et al. (2013) report biases of a few
354 percent in mean wind speed due to flow distortion around one of the buoys used to collect
355 data at high wind speed, but these platforms generally do not suffer such significant flow
356 distortion problems as ships.

357 For many applications ship-based measurements are the only option; for example,
358 direct eddy covariance measurements of gas transfer require instrumentation that can only
359 realistically be operated on a ship. A means of effectively dealing with biases induced by
360 flow distortion around a moving platform is thus essential. The methods demonstrated above
361 provide a successful correction; after its application the shape of the cospectra matches the
362 Kaimal form expected and our drag coefficient results lie within the range of recent leading
363 parameterisations.

364

365 **Appendix A. Underway vertical wind speed**

366 The motion correction algorithm of Edson et al. (1998) calculates a total platform velocity in
367 the earth frame as the sum of highpass filtered wave-induced motions, obtained from the
368 integration of accelerometers, and lowpass filtered velocities (the platform's underway
369 motion). The latter are applied only in the horizontal since the mean vertical velocity is zero
370 by definition. The corrected winds in the earth frame are obtained as the vector sum of
371 measured and platform velocities. This neglects the impact of flow distortion on the measured
372 winds (Fig. A1). At the point of measurement on the foremast of a ship, the mean flow is
373 forced to lift resulting in a streamline tilted upwards from the horizontal. The measured
374 along-streamline wind depends upon ship velocity as well as earth-relative wind. Since the
375 streamline is tilted, a fraction of the ship velocity affects the measured vertical as well as the
376 horizontal winds in the earth frame and must be corrected.

377 When conditions are stationary (an implicit assumption for direct flux measurement)
378 the measured, motion-corrected vertical wind, w_{rel} , can be corrected for the horizontal
379 platform mean velocity to obtain the true vertical wind speed w_{true} . The ratio of the mean true
380 to mean relative vertical winds is equal to the ratio of the mean true to mean relative
381 horizontal winds, i.e.

382

$$\frac{\overline{U_{true}}}{\overline{U_{rel}}} = \frac{\overline{w_{true}}}{\overline{w_{rel}}} \quad (\text{A1})$$

384

385 (Fig. A1). Then, as

386

$$w_{true} = w_{rel} - (\overline{w_{rel}} - \overline{w_{true}}) \quad (\text{A2})$$

388

389 w_{true} can be determined via Eq. (1)

390

$$w_{true} = w_{rel} - (\overline{w_{rel}} \times [1 - \overline{U_{true}} / \overline{U_{rel}}]).$$

392

393 Note that this affects the mean vertical wind only, not the high frequency perturbations;
394 however failure to account for the impact of flow distortion on the vertical wind
395 measurements would result in the streamline orientation being incorrectly calculated, and
396 both u' and w' values being biased after rotation into the streamline-oriented reference frame
397 in which the fluxes are calculated. We also note that at low wind speeds ($\sim < 5 \text{ m s}^{-1}$), the
398 determination of the reference frame for a particular measurement interval may be biased by
399 offsets in the vertical wind speed, leading to errors in the tilt calculation (Wilczak et al.,
400 2001; Landwehr et al., 2015).

401 The effectiveness of this correction is demonstrated through comparison of drag
402 coefficients from periods when the ship was stationary ($V_{ship} < 1 \text{ m s}^{-1}$) and underway ($V_{ship} >$
403 5 m s^{-1}). Prior to correction, measurements from the underway ship are biased high relative to
404 the stationary measurements (Fig. A2). Following correction, the stationary and underway
405 measurements are in very good agreement for all but the very lowest wind speeds.
406 Furthermore, for stationary periods (where the effect is small), the corrected and uncorrected
407 results are also in good agreement.

408

409 **Appendix B. CFD corrections for flow distortion**

410 The relative wind direction dependent CFD corrections for the mean flow distortion over the
411 ship are given in Table B1. These are strictly valid only for the location of our sonic
412 anemometer (1.24 m to starboard, 16.5 m above the waterline, and 5.0 m aft of the bow), but

413 should be broadly representative for nearby locations, and indicative of the directionally
414 dependent flow distortion that might be expected on any similar installation on other ships.

415

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423

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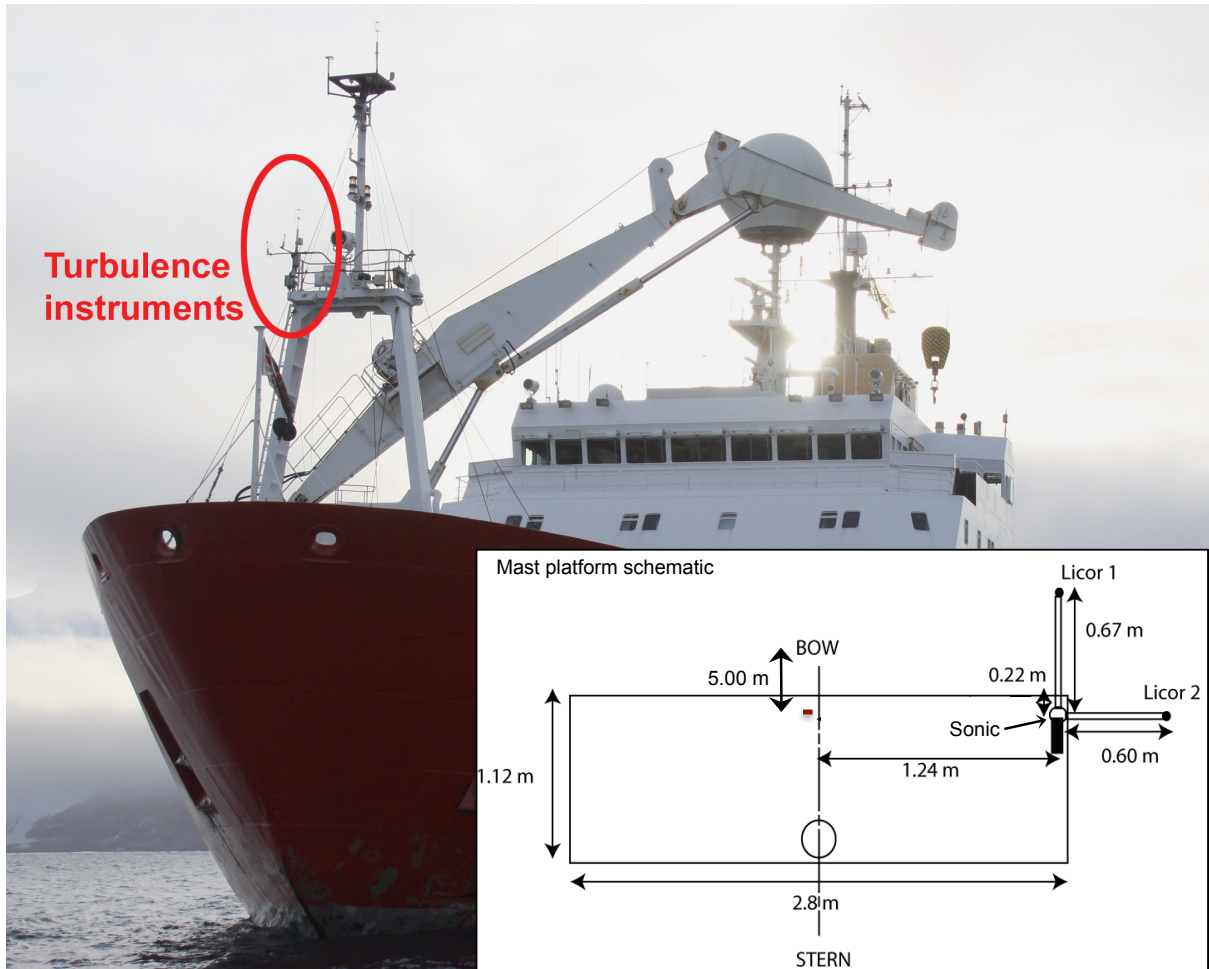
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550

Relative wind direction (°)	wind speed bias at $z - \Delta z$ (%)	Δz (m)
-20	2.98	1.44
-10	0.41	1.35
0	-0.39	1.32
10	-0.86	1.41
20	0.7	1.54
30	2.92	1.76
50	5.11	2.27
70	4.86	2.73
90	8.35	2.96
110	6.97	3.15

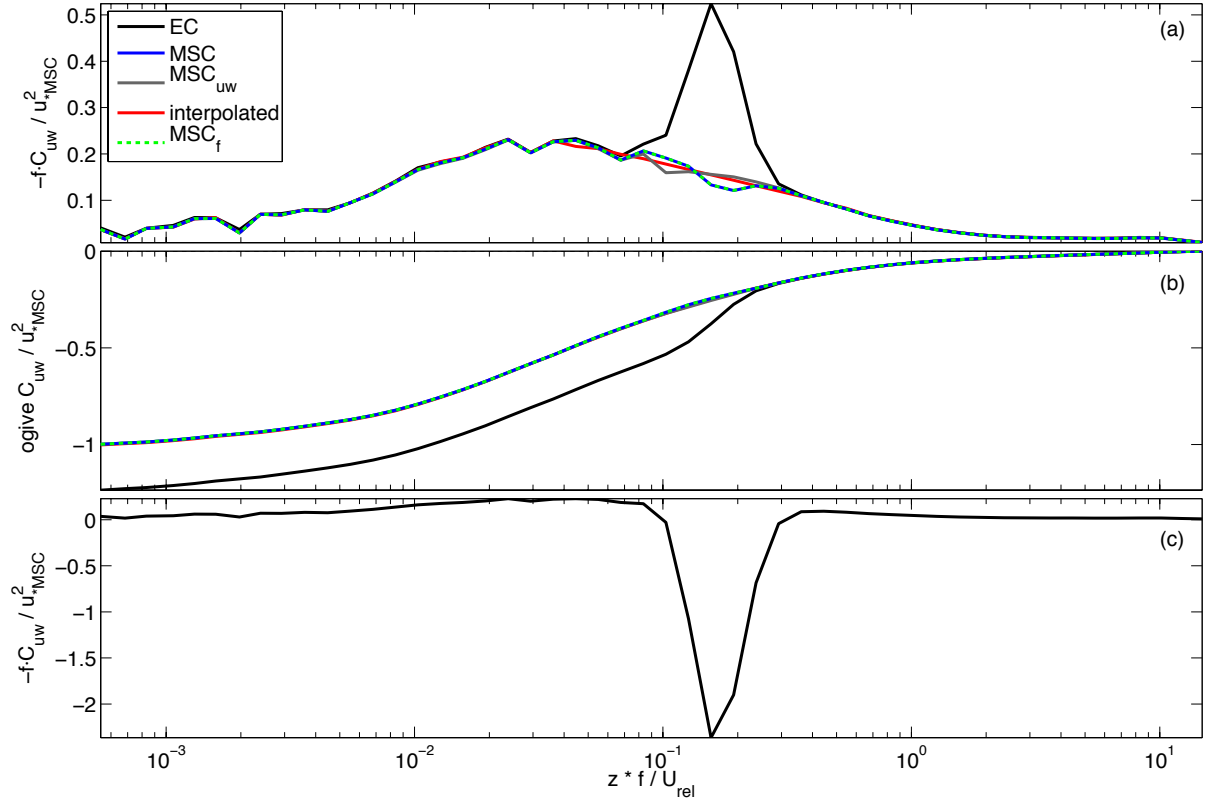
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552 Table B1. Variation of wind speed bias and vertical flow displacement with relative wind
553 direction, determined at the location of the AutoFlux anemometer (height above sea level, z ,
554 16.5 m). The wind speed bias and Δz are relative to a free stream location 2 seconds upstream
555 of the anemometer site (after Yelland et al., 2002). A negative relative wind direction
556 indicates a flow over the port side. Further details are given in Moat and Yelland (2015).
557

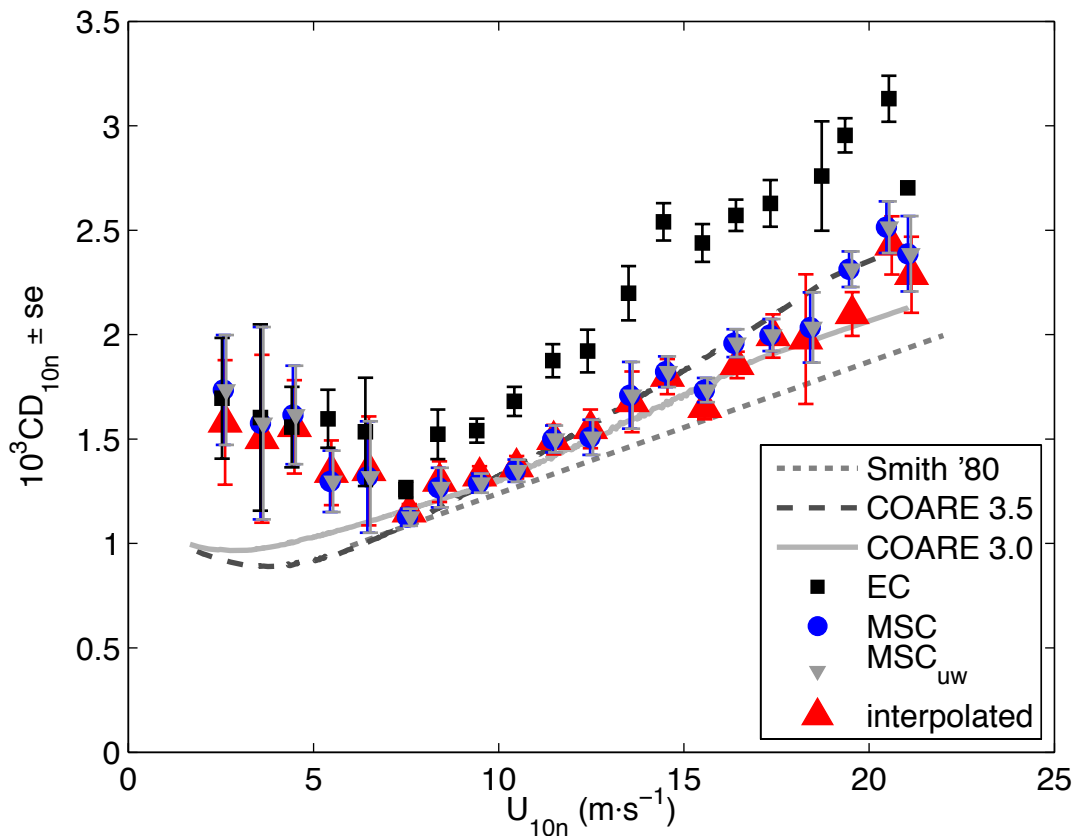


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Figure 1. Locations of the flux instrumentation on the *RRS James Clark Ross*. The sonic anemometer is 2.0 m above the starboard forward corner of the platform. Note that the forecastle crane is generally stowed close to the deck while the ship is underway or on station.



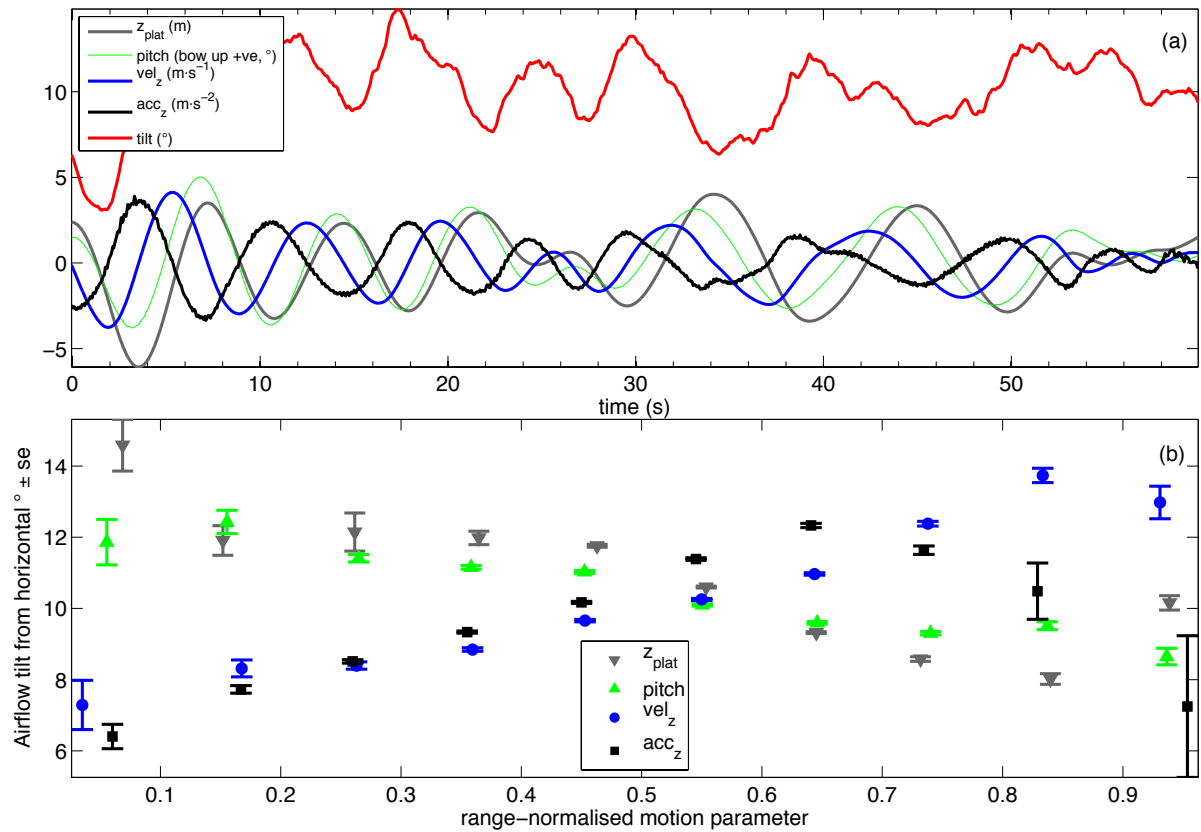
564
565 Figure 2. Frequency-weighted inverted and normalised momentum flux cospectra (a) and
566 normalised ogives (b), shown relative to frequency non-dimensionalised using measurement
567 height z and mean relative wind speed U_{rel} . Also shown are frequency-weighted, inverted and
568 normalised cospectra calculated prior to motion correcting the turbulent velocity components,
569 which results in a large upwards flux signal at the motion scale (c). Results shown are an
570 average of 131 30-minute duration measurements at mean wind speeds $10 \text{ m s}^{-1} < U_{10m} < 14$
571 m s^{-1} . EC indicates the cospectra after removing platform motion following Edson et al.
572 (1998) and shows the residual signal at scales typical of the wave field. The interpolation
573 across the wave scales has been applied between frequencies of 0.04 and 0.4 Hz. The motion-
574 scale correction (MSC) can either be applied as per Eq. (2) (MSC), with Eq. (2) applied to
575 both the along and vertical wind components (MSC_{uw}), or as described by Edson et al.,
576 (2011) (MSC_f). Normalisation of the five different sets of results is by u_* with MSC applied
577 as per Eq. (2). Note that the MSC_f line overlies the MSC line at all frequencies, and the
578 interpolated, MSC_{uw} and EC lines for frequencies away from the motion-scale.



580

581 Figure 3. Drag coefficients bin-averaged by wind speed, relative to U_{10n} ($n = 499$).

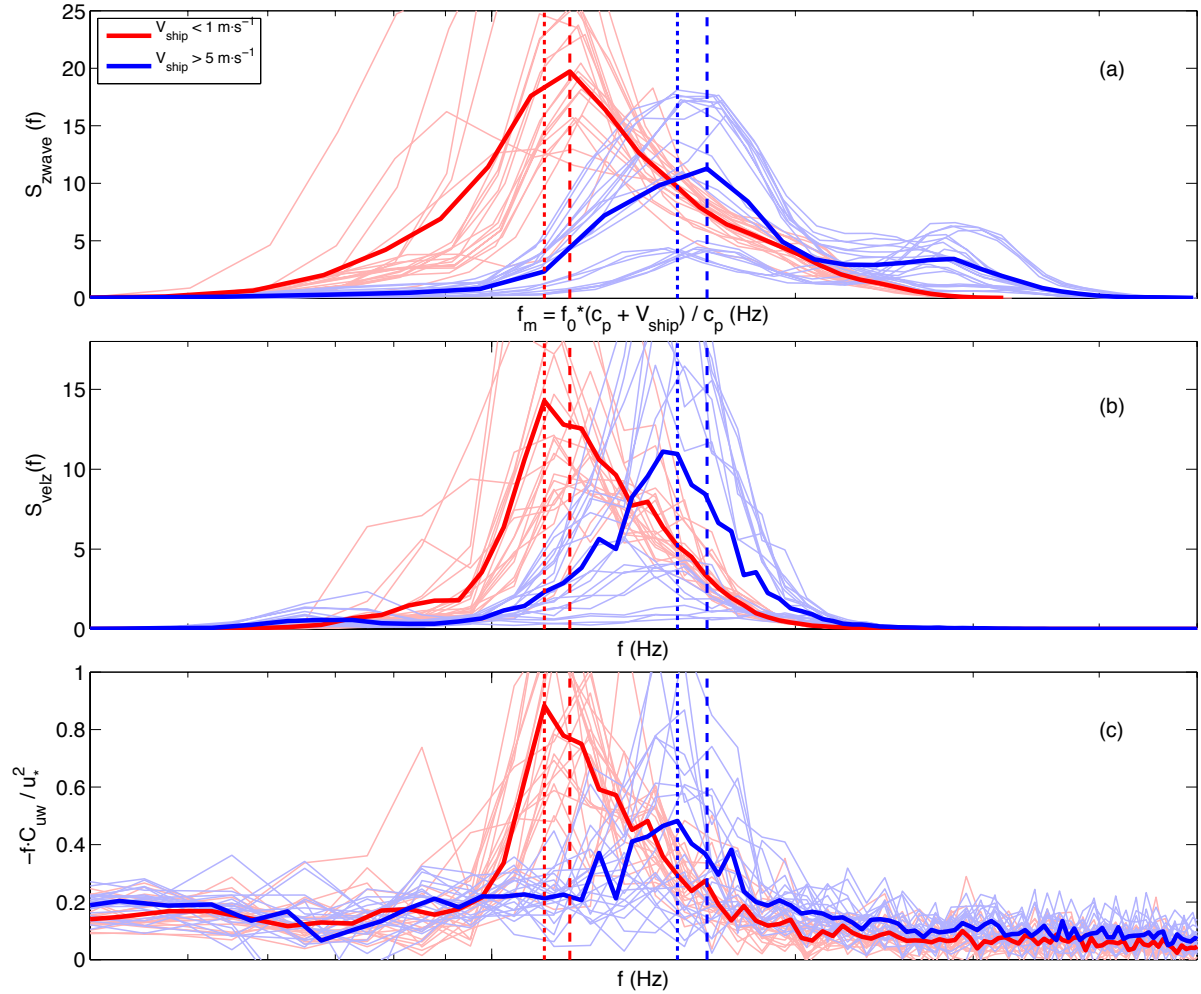
582 Measurements are shown either without correction for wave-scale bias (EC), or with
 583 correction applied to the vertical velocity only (MSC), both vertical and horizontal velocity
 584 components (MSC_{uw}), or via a simple interpolation across the wave-scale portion of the
 585 cospectra (interpolated). The bulk COARE 3.0 and 3.5 results are calculated without
 586 dependence on wave field or radiation.



588

589 Figure 4. (a): Time series (60 s) of vertical platform displacement, velocity and acceleration,
 590 platform pitch, and tilt from horizontal of the streamwise airflow measured by the AutoFlux
 591 anemometer. The tilt has been smoothed with a 40-sample moving average. The
 592 measurements are sampled from a period (23 April 2013, 21:00-21:30 UTC) with near bow-
 593 on winds and mean U_{10m} of 15.2 m s^{-1} . (b): Variation of the tilt of streamwise airflow from
 594 horizontal, relative to the vertical platform displacement, velocity, acceleration and platform
 595 pitch each normalised by their measured range. Tilt averages were made over the 30-minute
 596 period that the measurements in (a) were sampled from.

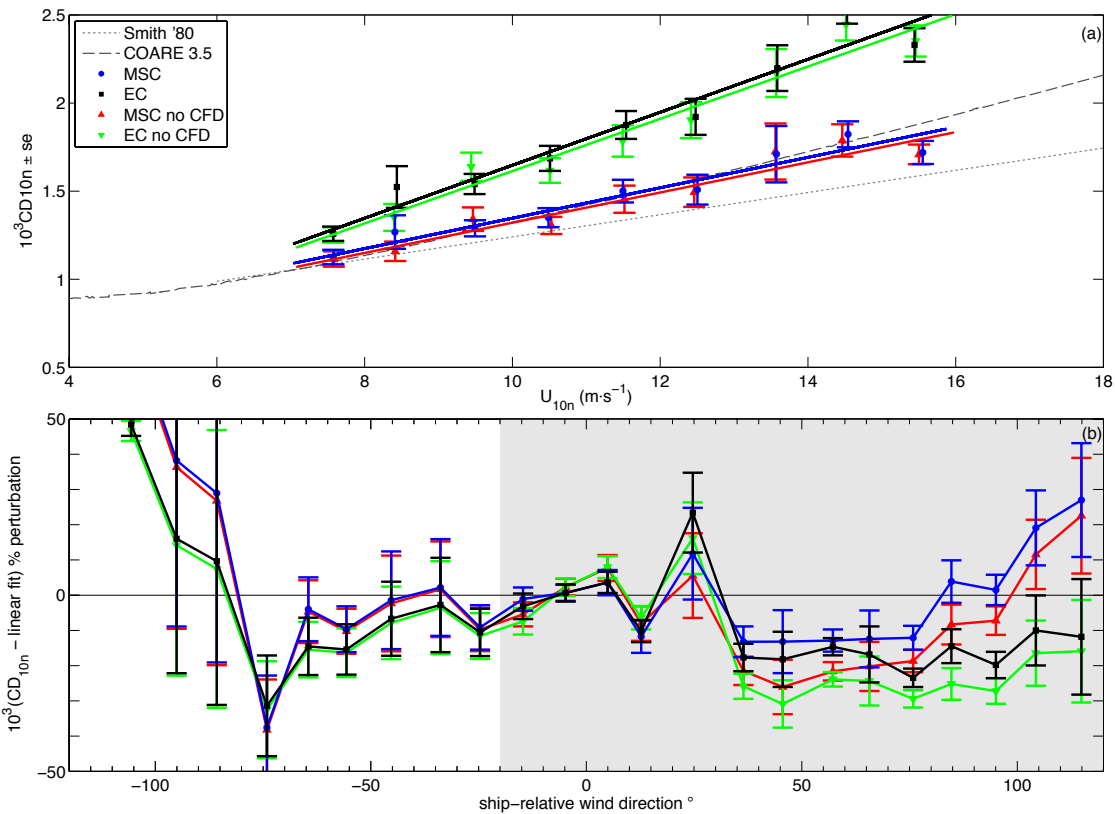
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599 Figure 5. Comparison of averaged spectra. In all panels two sets of averaged data are
600 compared, periods when the ship was stationary ($V_{ship} < 1 \text{ m s}^{-1}$, 21 periods) and periods
601 when the ship was steaming ($V_{ship} > 5 \text{ m s}^{-1}$, 20 periods); the individual spectra are shown as
602 pale lines for reference. For all measurements, U_{10m} was between 10 and 12 m s^{-1} . (a) Spectral
603 density of non-directional wave heights from WAVEX with frequency shifted to the
604 reference frame of the moving ship; (b) spectral density of platform vertical velocity as
605 measured on the foremast; (c) frequency-weighted inverted cospectral density for the
606 momentum flux (positive upwards) – turbulent velocity components are motion corrected, but
607 the MSC correction is not applied. The dashed vertical lines indicate the peak frequency of
608 the wave spectrum; dotted vertical lines indicate the peak frequency of the momentum flux
609 cospectra in (c). Note that the axis limits are set very close to the scale of the ship motion to
610 allow details to be seen clearly.

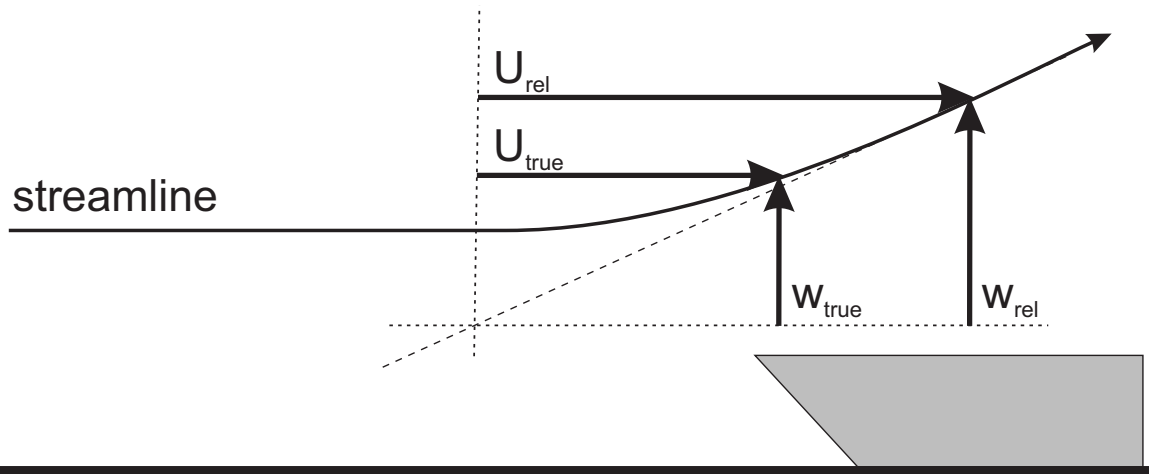
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612

613 Figure 6. (a): Measurements either without correction for wave-scale bias (EC), or with
 614 correction applied to the vertical velocity only (MSC) for wind speeds $7 m s^{-1} < U_{10m} < 16 m$
 615 s^{-1} ($n = 335$), and relative wind directions between -20° and $+50^\circ$ (where a wind on the bow
 616 is at 0°). Lines are linear fits to the measurements. (b): variation of the difference between
 617 measured drag coefficients and the linear fits against relative wind direction for the same
 618 wind speed criteria ($n = 663$). Both panels also show measurements (with and without MSC)
 619 which have not had CFD-derived corrections to mean wind speed and height applied. Note
 620 that CFD corrections were only applied for the shaded range.

621

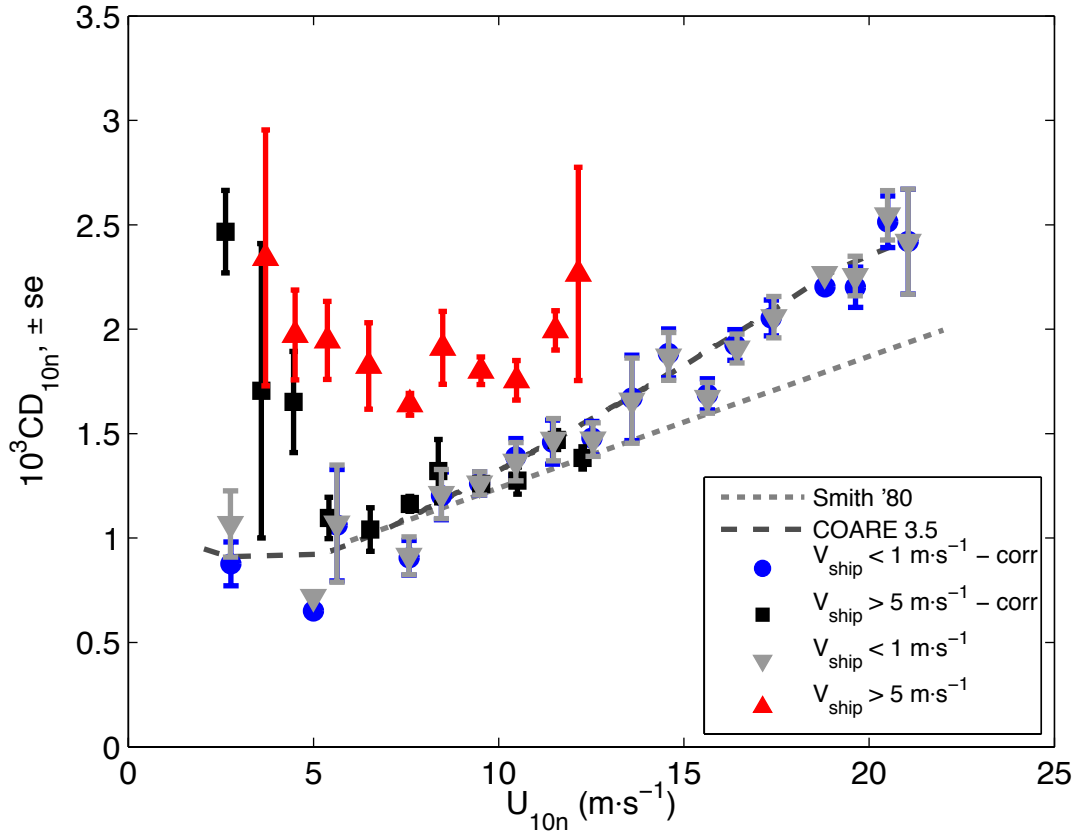


622

623 Figure A1. Schematic of the impact of ship horizontal velocity on non-horizontal airflow.

624 The measured horizontal (U_{rel}) and vertical (w_{rel}) wind components must both be corrected
625 for ship velocity to obtain the true wind components. Not correcting the measured vertical
626 wind will result in an incorrect determination of the tilt angle of the flow from horizontal.

627



628
 629 Figure A2. Wind speed-averaged drag coefficients, relative to U_{10n} . Two sets of
 630 measurements are compared: where the ship was deemed stationary ($V_{ship} < 1 m \cdot s^{-1}$, $n = 233$);
 631 and where the ship was underway ($V_{ship} > 5 m \cdot s^{-1}$, $n = 182$). The measurements are shown
 632 with ('corr') and without the vertical wind speed corrected as per Eq. (1).
 633