



**Uncertainties in  
global aerosols and  
climate effects**

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# Uncertainties in global aerosols and climate effects due to biofuel emissions

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## Abstract

Aerosol emissions from biofuel combustion impact both health and climate; however, while reducing emissions through improvements to combustion technologies will improve health, the net effect on climate is largely unconstrained. In this study, we examine sensitivities in global aerosol concentration, direct radiative climate effect, and cloud-albedo aerosol indirect climate effect to uncertainties in biofuel emission factors, optical mixing-state, and model nucleation and background SOA. We use the Goddard Earth Observing System global chemical-transport model (GEOS-Chem) with Two Moment Aerosol Sectional (TOMAS) microphysics. The emission factors include: amount, composition, size and hygroscopicity, as well as optical mixing-state properties. We also evaluate emissions from domestic coal use, which is not biofuel but is also frequently emitted from homes. We estimate the direct radiative effect assuming different mixing states (internal, core-shell, and external) with and without absorptive organic aerosol (brown carbon). We find the global-mean direct radiative effect of biofuel emissions ranges from  $-0.02$  to  $+0.06 \text{ W m}^{-2}$  across all simulation/mixing state combinations with regional effects in source regions ranging from  $-0.2$  to  $+1.2 \text{ W m}^{-2}$ . The global-mean cloud-albedo aerosol indirect effect ranges from  $+0.01$  to  $-0.02 \text{ W m}^{-2}$  with regional effects in source regions ranging from  $-1.0$  to  $-0.05 \text{ W m}^{-2}$ . The direct radiative effect is strongly dependent on uncertainties in emissions mass, composition, emissions aerosol size distributions and assumed optical mixing state, while the indirect effect is dependent on the emissions mass, emissions aerosol size distribution and the choice of model nucleation and secondary organic aerosol schemes. The sign and magnitude of these effects have a strong regional dependence. We conclude that the climate effects of biofuel aerosols are largely unconstrained, and the overall sign of the aerosol effects is unclear due to uncertainties in model inputs. This uncertainty limits our ability to introduce mitigation strategies aimed at reducing biofuel black carbon emissions in order to counter warming effects from greenhouse-gases. To better understand the climate impact of particle emissions from biofuel combustion, we rec-

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commend field/laboratory measurements to narrow constraints on: (1) emissions mass, (2) emission size distribution, (3) mixing state, and (4) ratio of black carbon to organic aerosol.

## 1 Introduction

5 Close to half of the world's population relies on combustion of domestic solid fuel use as a source of energy (Bruce et al., 2000), creating concerns for both air quality (Bruce et al., 2006) and climate (Bond et al., 2004a; Venkataraman et al., 2005). Domestic solid fuel combustion in many parts of the world is dominated by biofuel, which includes wood, agricultural waste, animal dung, or charcoal as fuel for domestic energy  
10 needs (Bond et al., 2007; Fernandes et al., 2007). Biofuel combustion is especially prevalent in developing countries where a significant portion of the population lacks access to electricity or clean combustion technology (Bruce et al., 2000). Gaseous and particulate matter emitted from biofuel combustion degrades air quality and may lead to detrimental health risks (Akbar et al., 2011). The recent Global Burden of Disease  
15 Study ranks household air pollution from solid fuels and ambient air pollution from particulate matter (all sources) as the third and ninth largest contributors, respectively, to the global burden of disease (Lim et al., 2012). Improved combustion devices that reduce human exposure to pollutants should reduce the burden of disease from household air pollution; however, the net climate effect resulting from changing emissions  
20 remains uncertain.

Combustion of biofuel emits greenhouse gases (such as carbon dioxide and methane) (Johnson et al., 2008; Yevich and Logan, 2003) as well as carbonaceous aerosol particles, such as black carbon (BC) and organic aerosol (OA) (Bond et al., 2007). In the atmosphere, carbon dioxide and methane are generally well-mixed due  
25 to long lifetimes, and their impacts on climate are better understood than those from aerosols (Boucher et al., 2013). Conversely, BC and OA have short lifetimes with more complex climate effects necessitating the use of aerosol microphysical models

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to understand the net impacts (e.g. Pierce et al., 2013; Spracklen et al., 2011). Carbonaceous aerosols can affect climate through scattering/absorbing solar radiation (direct radiative effect), changing the radiative properties of clouds (the cloud-albedo and cloud-lifetime indirect aerosol effects), changing the absorption of snow (snow albedo effect), and changing the temperature profile of the atmosphere (semi-direct effect) (Boucher et al., 2013). The cloud-lifetime indirect effect, snow albedo effect, and semi-direct effect are more challenging to simulate than the direct effect and cloud-albedo indirect effect; such effects involve changes to meteorology or land surfaces and require fully interactive climate models to elucidate (e.g. Bauer et al., 2010; Bauer and Menon, 2012; Jacobson, 2010). In this study, we will be limited to the direct radiative effect and the cloud-albedo aerosol indirect effect but acknowledge that this is not the total aerosol climate forcing.

The direct radiative effect (DRE) refers to direct scattering and absorption of incoming solar radiation (Charlson et al., 1992). BC has a strong absorbing component while OA is usually considered to be entirely scattering; however, research has shown that under certain combustion conditions OA may have an absorbing component (Kirchstetter et al., 2004; Lack et al., 2012; Saleh et al., 2013, 2014). Absorbing OA, commonly termed brown carbon, has a strong wavelength dependence (Andrea and Gelencser, 2006), which varies with the BC to OA ratio from combustion (Saleh et al., 2014). Additionally, the efficiency with which a particle absorbs or scatters solar radiation is dependent on its size (Seinfeld and Pandis, 2006). Small particles lack sufficient size to interact efficiently with solar radiation, while large particles will only interact with radiation on the surface of the particle. There then exists a size window (generally, diameters between 100 nm and 1  $\mu$ m) that maximizes particle mass and surface area to interact with solar radiation most efficiently.

The magnitude of the DRE is strongly dependent on the mixing state of the particles, i.e. do different particle species exist in the same particle or separate particles (Jacobson, 2001; Klingmüller et al., 2014). Aerosol-climate models generally make assumptions about the mixing state of absorbing BC with scattering species rather than

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tracking the species mixing explicitly (with some exceptions e.g. Bauer et al., 2010). In these models, BC is assumed to be mixed with other particle-phase species in several different ways: homogeneously with scattering species (internal), as a core surrounded by a homogeneously mixed shell (core-shell), or as separate from other aerosol species (external) (Jacobson, 2000). In reality, the mixing state lies somewhere in between these assumptions with the mixing state dependent on microphysical processes (e.g. condensation, coagulation), time and location. For a fixed amount of BC and scattering mass, assuming an internal mixture yields the most absorption (Jacobson, 2000; Klingmüller et al., 2014). There are several methods to estimate the optical properties of a homogenous internal mixture; the most common is to volume weight the refractive indices (Bond et al., 2006; Chung and Seinfeld, 2005). Volume weighting the refractive indices decreases the real (scattering) component of the refractive index, while distributing the mass of BC around the entirety of every particle, thus increasing the surface to mass ratio of BC and assigning all particles an absorbing component (Seinfeld and Pandis, 2006). This homogenous mixture is thought to be unphysical, but it does give an upper bound on absorption efficiency (Bond et al., 2006; Jacobson, 2000). The external mixture, where absorbing BC is assumed to be in entirely different particles from scattering species, predicts the least amount of absorption (Jacobson, 2000). In an external mixture, only particles with an absorbing component contribute a positive radiative effect (Seinfeld and Pandis, 2006). However, this mixing state is also unrealistic in the atmosphere, since processes such as condensation and coagulation contribute to mixing (Pierce and Adams, 2007; Reimer et al., 2009; Roden et al., 2006). The core-shell morphology, in which a scattering shell surrounds an absorbing BC core, yields less total absorption than the internal mixture assumption but more than the external mixture assumption. The shell acts as a lens that scatters more photons into the absorbing core (Ackerman and Toon, 1981). This lensing effect amplifies absorption over an external mixture, with the magnitude of the amplification dependent on the core and shell thicknesses (Bond et al., 2006). In the case of absorptive OA, the shell can either absorb or scatter radiation (Lack and Cappa, 2010). Bond et al. (2006)



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sites for water vapor, called cloud condensation nuclei (CCN) (Pierce et al., 2007; Spracklen et al., 2011). Increasing OA and BC concentrations may lead to an increase in CDNC, which will increase cloud albedo and thus yield a negative forcing. The ability for OA and BC particles to act as CCN is a function of particle size and hygroscopicity as well as the maximum supersaturation of water vapor in the cloud (Petters and Kreidenweis, 2007). Larger particles can activate into cloud drops more easily than smaller particles (due to higher saturation vapor pressures over curved surfaces); however, larger particles may deplete water vapor concentrations, lower the maximum supersaturation, and limit activation of smaller sized particles.

Emission factors from biofuel combustion are dependent on combustion conditions, which can vary with the type and size of fuel (Li et al., 2009; L'Orange et al., 2012), the combustion device (Bond et al., 2004b; Jetter et al., 2012), and the operator (Roden et al., 2009). In general, flaming conditions tend to emit relatively more BC mass and larger sized particles (Janhäll et al., 2009) compared to smoldering. Grieshop et al. (2011) finds that the PM emission mass can vary by a factor of four based on different stove and fuel combinations. Wood and agricultural waste emit mostly carbonaceous particles, while coal (used in domestic fuel use but is not biofuel) has a higher sulfur content and so emits more SO<sub>2</sub> gas, which reacts to form condensable H<sub>2</sub>SO<sub>4</sub> vapor in the atmosphere that contributes to particle formation and growth. PM mass and composition can vary significantly between different types of technologies used mainly for cooking, heating, or lighting. Additionally, PM emission mass may be dominated during relatively short times of re-fueling and ignition (Tryner et al., 2014). Variability in emissions factors (including number of users, location of users, stove technology, cooking practices, etc.) can lead to uncertainties in global inventories.

Further complicating biofuel aerosol simulations are that these particles will age in plumes on spatial scales smaller than those resolved by global models. For example, primary organic aerosol (POA) may evaporate and secondary organic aerosol (SOA) may form in woodsmoke plumes (Robinson et al., 2007; Grieshop et al., 2009; Hennigan et al., 2011). Additionally, particle number concentration is decreased by coagula-

tion, which simultaneously increases the mean diameter of the particles (Capes et al., 2008; Sakamoto et al., 2013). Since the sub-grid processes are not explicitly resolved, models must account for this processing at the emissions stage, which adds additional uncertainty to the number, size and composition of the particles beyond the uncertainties of traditional emissions inventories.

Reducing human exposure to biofuel combustion emissions will likely benefit human health. However, the climate impacts of reducing (or modifying) biofuel combustion are relatively poorly constrained due to the uncertainties described above: emissions amount, size, composition and optical properties as well as uncertainties in other model processes that affect biofuel particles. These uncertainties limit studies aimed at evaluating potential black carbon mitigation strategies from specific sources (Bond and Sun, 2005). In this paper, we quantify the contribution of various uncertainties in biofuel aerosol emissions (emissions rate, composition and size) and model processes (optical mixing state, secondary organic aerosol and nucleation) to the DRE and cloud-albedo AIE. We determine which factors pose the greatest uncertainty to our understanding of how changes to biofuel combustion will affect climate.

In Sect. 2 we discuss our methods for estimating uncertainties in the climate effects from biofuel. We present modeling results in Sect. 3. Conclusions and discussions for the results are presented in Sect. 4.

## 2 Methods

### 2.1 GEOS-Chem-TOMAS overview

We use Goddard Earth Observing System global chemical-transport model (GEOS-Chem) coupled with the TwO Moment Aerosol Sectional (TOMAS) microphysics scheme (Adams and Seinfeld, 2002) to calculate aerosol number, mass, and size distributions. This version of TOMAS uses 15 size sections ranging from 3 nm to 10  $\mu\text{m}$  with tracers for sulfate, sea-salt, hydrophobic and hydrophilic OA, externally (pure)

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and internally mixed (aged) BC, and dust (Lee and Adams, 2012; Lee et al., 2013). Conversion from externally mixed OA and BC into internally mixed occurs on a fixed ageing timescale of 1.15 days. We use GEOS-Chem version 9.02 with  $4^\circ \times 5^\circ$  horizontal resolution and 47 vertical layers with assimilated meteorology from GEOS5 (http://gmao.gsfc.nasa.gov) to simulate the year 2005.

We use black and organic carbon (OC) emissions from biofuel and other combustion-related sources for the year 2000 from Bond et al. (2007). Anthropogenic fossil fuel emissions are from the Emissions Database for Global Atmospheric Research (EDGAR) inventory (Olivier et al., 1995). The EDGAR inventory is overwritten in the United States by the Environmental Protection Agency 2005 National Emissions Inventory (NEI05; http://www.epa.gov/ttn/chief/net/2005inventory.html), in Canada by the Criteria Air Contaminants (CAC; http://www.ec.gc.ca/inrp-npri/), in Mexico and the southwestern US by the Big Bend Regional Aerosol and Visibility Study (BRAVO; Kuhns et al., 2003), in Asia by the Streets inventory (Streets et al., 2003), and in Europe by the Cooperative Programme for Monitoring and Evaluation of the Long-Range Transmission of Air Pollutants in Europe (EMEP; Auvray and Bey, 2005). Residential coal emissions in the Streets inventory are considered separately from biofuel. Open biomass burning (e.g. wildfire) emissions are from the Global Fire Emissions Database version 3 (GFEDv3; Van Der Werf et al., 2010).

Figure 1 contains the global annual biofuel BC and OC emissions from Bond et al. (2007) and sulfur-dioxide ( $\text{SO}_2$ ) emissions from EDGAR along with Asia-regional  $\text{SO}_2$  emissions from residential coal from the Street's inventory. In other parts of the world, emissions from residential coal use are combined with other sources, and thus we can only isolate this fuel use over Asia. Annual biofuel combustion emissions are  $1.6 \text{ TgC year}^{-1}$  of BC,  $6.3 \text{ TgC year}^{-1}$  of OC, and  $0.27 \text{ TgS year}^{-1}$   $\text{SO}_2$ . The emissions of  $\text{SO}_2$  from residential coal use in Asia are  $1.9 \text{ TgS year}^{-1}$  in the Streets inventory. Our lack of isolated global residential coal is a limitation of this study. Northern India and eastern China have the largest aerosol emissions, with substantial contribution from

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et al. (2014), which calculates the magnitude and wavelength dependence of the imaginary index of refraction of OA based on the BC to OA ratio.

Our values of the imaginary index of refraction at 550 nm range from 0.05 (based on Saleh et al., 2014) to 0.006 (the GADs value for non-absorbing OA). Here we use the BC to OA ratio of the model grid box based on all emissions, whereas Saleh et al. (2014) use the BC to OA ratio near the source of emissions only for biomass burning and biofuel emissions. We expect this to introduce some error, however, this method should be sufficient to show the sensitivity to OA absorption. The DRE is calculated at each grid cell for 5 wavelengths bands (380, 580, 780, 980, 3000 nm) and weighted by the solar spectrum to calculate the broadband DRE. Albedo and cloud fraction are taken as monthly averages from GEOS5. We assume no aerosol effects in columns with clouds, and our all-sky DRE is the clear-sky DRE multiplied by the cloud-free fraction.

We determine the cloud-albedo AIE due to biofuel emissions offline using the radiative transfer model of Edwards and Slingo (1996) with simulated changes to cloud droplet number concentration (CDNC). Our method is described in Scott et al. (2014) and is based on a simplification discussed in detail by Spracklen et al. (2011). The change in the number of activated particles is calculated using monthly mean aerosol distributions from GEOS-Chem-TOMAS with an activation parameterization, assuming a globally uniform updraft velocity  $0.2 \text{ m s}^{-2}$ . We calculate cloud droplet number concentrations using the mechanistic parameterization of Nenes and Seinfeld (2003), as updated by Fountoukis and Nenes (2005), which is based on modified Kohler theory. In these calculations, sea-salt and sulfate are assumed to be water soluble and assigned van't Hoff factors of 2 and 2.5 (following Wang et al., 2010), respectively, when calculating the solute term in the Kohler equation; other components present in each size bin are able to activate when they are internally mixed, which excludes the pure externally mixed BC. The change in CDNC is used to calculate a perturbation to the effective radii of cloud droplets in low- and mid-level (below 600 hPa) water clouds, which in turn leads to a change in net top of the atmosphere radiative flux. We use monthly

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averaged cloud climatology from the International Satellite Cloud Climatology Project (ISCCP-D2; Rossow and Schiffer, 1999) for the year 2000.

## 2.3 Description of simulations

We test the sensitivity of changes to global aerosol concentration and associated radiative effects due to biofuel emissions and various emission and model assumptions. The 18 simulations used in this study are outlined in Table 2. In the model, we must assume effective emissions size distributions that include the effects of sub-grid coagulation that increase the size of the particles (and reduces the number), as we do not explicitly represent coagulation within the plumes on sub-grid scales. In the BASE simulation (our “default” simulation), this assumed size distribution is a single lognormal distribution with a geometric number-mean diameter (GMD) of 100 nm and a SD of 2. Primary OC is emitted in the model as OA with a fixed OA to OC ratio of 1.8. In each grid box, 80 % of OA is emitted as hydrophilic and 50 % of BC is emitted as internally mixed. Hydrophobic OA and externally mixed BC can age through condensation and coagulation, represented in the model as a fixed timescale of 1.15 days. Nucleation rates are parameterized with binary nucleation in the free troposphere (Vehkamaki et al., 2002) along with a ternary parameterization (Napari et al., 2002) scaled globally by a  $10^{-5}$  tuning factor (Jung et al., 2010; Westervelt et al., 2013). Secondary organic aerosol (SOA) includes both a biogenic contribution ( $19 \text{ Tgyr}^{-1}$  in GEOS-Chem-TOMAS) and an anthropogenically-enhanced contribution of  $100 \text{ Tgyr}^{-1}$  correlated with anthropogenic CO emissions (D’Andrea et al., 2013), following the approach of Spracklen et al. (2011). In the NOBIOF simulation, BC, OA and  $\text{SO}_2$  emissions from biofuel are turned off, while all other emissions remain unchanged. We perform two sensitivity tests regarding emission mass. In MASSX2 and MASSX0.1 the emission mass of OA and BC from biofuel in each grid box from Bond et al. (2007) is doubled and reduced by 90 %, respectively. The purpose of increasing the upper bound by a factor of 2 is to explore general uncertainty in the emissions amount, while the lower bound represents a potentially large reduction in emissions due to

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25%. Over oceans, BC mass increases by 10–25%, except on the subtropical west coasts where frequent boundary-layer precipitation occurs. Thus, biofuel emissions are a significant source of BC in both source regions and in remote regions. As OA has additional sources beyond those of BC (e.g. secondary organic aerosol), the fractional increases in OA are smaller than those of BC. Globally-averaged SO<sub>2</sub> mass increases by 0.5% leading to a 0.02% increase in sulfate aerosol (not shown).

Figure 3c and d contains zonally averaged BC and OA mass percent changes with pressure level. When biofuel emissions are included, BC and OA mass increases throughout much of the troposphere. Black carbon increases by > 25% in the Northern Hemisphere tropics at all pressure levels. Organic aerosol increases are limited to 5–10% at higher altitudes in the Northern Hemisphere. Tropical convection lofts BC to high altitudes, which may have implications for the semi-direct and ice cloud effects (not addressed here).

The percent changes in N10 (a), N40 (b), N80 (c), and N150 (d) in the boundary layer, due to the inclusion of biofuel emissions, are shown in Fig. 4. Changes in N10, N40, and N80 vary by sign and magnitude across different regions resulting in an annual global mean change of 0.29, 0.93, and 1.59% (Table 3). Conversely, N150 increases over all land masses with percent changes of over 20% in heavy source regions and an annual global mean increase of 2.70%. The regional decreases in N10, N40, and N80 are caused through a feedback in aerosol microphysics. Biofuel BC and OA emissions (with a median diameter of 100 nm) increase total particle number and thus increase the total aerosol surface area available for condensation. This increased condensation sink leads to (1) lower concentrations of condensable vapors (sulfuric acid and secondary organics), (2) reduced nucleation rates due to reduced sulfuric acid concentrations, (3) slower growth of particles due to reduced condensable vapor concentrations, and (4) increased scavenging of small particles by coagulation due to increases in total aerosol surface area. This feedback is partly mitigated by oxidation of biofuel SO<sub>2</sub> emissions into sulfuric acid, which contributes to nucleation and growth. However, the increased condensation sink from primary BC and OA particles outweighs the contribution of

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biofuel SO<sub>2</sub> emissions, resulting in a net decrease in sulfuric acid and organic vapors. These factors combine to lower the concentration of small particles in some polluted regions where new-particle formation and growth is a significant contributor to particle concentrations. This decrease in N10 reduces the amount of particles able to grow to N40 and N80 sizes. Conversely, sub-Saharan Africa and South America have an increase in all particle sizes. Low initial sulfuric acid concentrations in these areas prevent this microphysical feedback, and therefore addition of biofuel aerosol simply increases the number of particles for all size classes. Finally, biofuel emissions do not significantly change the contribution of particles growing to N150 sizes through condensation, and so suppression of nucleation and condensational growth does not lead to any decreases at this size.

The corresponding zonally averaged percent changes in particle number concentration are plotted in Fig. 5. In all size classes, particle number concentration tends to increase near the equator and subtropics close to the surface; however, at higher altitudes and away from source regions N40 and N80 decrease by 0.2–1 %. The reason for this is a similar feedback as described above. N40 and N80 are more efficiently scavenged through wet deposition. Near the surface these particles are replaced by primary emissions; however, at higher altitudes condensational growth of nucleated particles is a significant source. With reduced nucleation and condensational growth, fewer particles are able to grow to N40 and N80 sizes. The net result is a decrease in N40 and N80 at higher altitudes. Biofuel emissions do not significantly alter the source of N150 sized particles from condensation growth, and so primary emissions lead to increases in N150 in all locations.

The DRE due to biofuel emissions is shown in Fig. 6 for the 6 different mixing state assumptions. The global-mean DRE ranges from +0.021 to  $-0.008 \text{ W m}^{-2}$  (Table 4) with strong regional variations across mixing state assumptions. Similar to past global modeling studies the external mixture gives the least absorption and the internal mixture gives the most absorption when absorptive OA is not included (Chung and Seinfeld, 2005; Jacobson, 2000; Klingmüller et al., 2014). Purely internal or external mix-

tures globally are not expected to be realistic, but they do give upper and lower bounds on our optics assumptions. The core-shell calculation lies in the middle of our calculated range with an annually averaged global DRE of  $+0.007 \text{ W m}^{-2}$ . The ext\*1.5 assumption predicts less absorption than the core-shell assumption, in this case leading to a global negative DRE of  $-0.002 \text{ W m}^{-2}$ . When absorbing OA is included, the DRE becomes more positive. The core-shell morphology with absorptive OA increases the magnitude of the DRE from  $+0.007$  to  $+0.021 \text{ W m}^{-2}$ , which results in the most positive DRE among the cases we consider here. The corresponding ext\*1.5 DRE increases from  $-0.002$  without absorptive OA to  $+0.015 \text{ W m}^{-2}$  when absorptive OA is included. Optical assumptions are one of the key uncertainties driving the variability in the DRE. In this study, we estimate the DRE assuming a single mixing state for all grid boxes (with size and composition determined by the concentrations at each location). The mixing state and optical properties of OA likely vary by region and emission source (Table 5); however, this is not explicitly explored here.

Different mixing state assumptions also lead to strong variations regionally as well as in the global mean (Fig. 6 and Table 4), In some regions, such as over China, the DRE can range from a strong positive (over  $+0.4 \text{ W m}^{-2}$ ) to negative (less than  $-0.2 \text{ W m}^{-2}$ ) in our different sensitivity tests. Some of the regional variability is explained by surface albedo variability. Over bright surfaces, such as the Arctic and Sahara, the DRE is positive in every mixing-state assumption tested. At these locations, the aerosol mixture is darker than the underlying surface across all mixing state assumptions and, therefore, planetary albedo is reduced. Over dark surfaces (oceans), a reduction in aerosol absorption efficiency (by assuming a different mixing state) makes the aerosol mixture brighter than the underlying surface and, thus, the planetary albedo increases. The negative DRE in eastern China, Southern India, and Europe in the external and ext\*1.5 mixing state is a result of the aerosols increasing the reflectivity over the relatively darker surface, but there is a positive DRE in these locations for the internal mixing state and when absorptive OA is included.

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Competing regions of positive and negative DRE limit the magnitude of the globally averaged DRE. Figure 2b contains the RMS for the different mixing states and simulations. The RMS represents the absolute model sensitivity of the climate effects to different inputs, accounting for competing regions of positive and negative effects that are not represented in a global-mean. Biofuel combustion contributes changes in the DRE on the order of  $\pm 0.1 \text{ W m}^{-2}$  around the globe. The RMS values for each mixing state are greater than the arithmetic averages; however, the relative order of the magnitude of the mixing states is slightly different. The core-shell with absorptive OA still has the largest value, but now the ext\*1.5 with absorptive OA has a noticeably stronger effect than the internal mixture. The ext\*1.5 and external mixture have the same strength of forcing, with differing amounts of positive and negative regions.

The cloud-albedo AIE due to biofuel emissions is plotted in Fig. 7. Biofuel emissions lead to a slight negative in the global mean of the indirect effect of  $-0.006 \text{ W m}^{-2}$ . The magnitude of this global mean is balanced by regional variations. In general the sign and magnitude of the AIE is a competition between increases in CDNC from the biofuel primary emissions and decreases in CDNC from an increased condensation sink of sulfuric acid, organics (suppressing nucleation and growth rates), and water vapor (suppressing supersaturation and activation into CDNC). Biofuel emissions result in a strong negative cloud-albedo AIE in the tropics, specifically in Africa and South America. In this region, the contribution of nucleation and condensational growth to N40 and N80 is less sensitive to the addition of primary biofuel aerosol, and so primary biofuel emissions lead to increases in N40 and N80 from the surface to around 600 hPa (Fig. 5). The increases in N40 and N80 aloft lead to increases in CDNC and thus cloud albedo. Conversely, over southern mid-latitude oceans, the reduced nucleation and condensational growth leading to reduced N40 wins out over transported primary emissions leading to a net positive cloud-albedo AIE. This leads to an overall reduction in column CDNC and a positive AIE signal. India and China have both significant primary emissions as well as a strong nucleation suppression feedback, which limits increases in particle number. Additionally, in this region there is strong competi-

tion for water vapor and large background aerosol concentrations suppress maximum supersaturation achieved in updrafts. Increases in N150 in this area will further limit the maximum supersaturation, as water vapor will preferentially condense on larger sized particles leaving less available for N40 and N80 sizes. As with the DRE, competing regions of positive and negative values limit the magnitude of the global mean AIE. In Fig. 2b, the RMS value ( $0.04 \text{ W m}^{-2}$ ) for the AIE is much larger in magnitude than the arithmetic mean ( $-0.006 \text{ W m}^{-2}$ ). Suppression of condensational growth and of maximum supersaturation in polluted regions explains why the magnitude of the AIE over Beijing ( $-0.004 \text{ W m}^{-2}$ ) is much smaller than over the relatively cleaner Addis Ababa ( $-0.22 \text{ W m}^{-2}$ ) (Fig. 2c and d).

### 3.3 Sensitivity of radiative effects to emission mass uncertainties

We test the sensitivity of the direct and indirect effects to primary biofuel particle emissions (BC and OA) to account for uncertainty in measurements, sub-grid ageing, and combustion device improvement scenarios designed to limit particle emissions. Van Donkelaar et al. (2015) find increasing particle emissions in developing countries (China, India, and the Middle East) since 1998, and due to changing emissions, emissions inventories likely carry large uncertainties. Biofuel is a significant emission source in these regions. In MASSX2, we double the BC and OA biofuel emissions mass and compare the results to the NOBIOF simulation. The DRE has a strong dependence on mass and number. Doubling the emitted BC and OA approximately doubles the increase in atmospheric concentrations of BC and OA, as well as N80 and N150 relative to the BASE simulation (Table 3). This leads to approximately doubling the magnitude of the biofuel DRE for all mixing-state assumptions compared to the BASE-NOBIOF comparison (Table 4 and Fig. 2a). The change in magnitude is in the same direction as the original sign of the DRE, so the external mixture has a larger negative DRE in MASSX2. In MASSX0.1, we emit one-tenth of the BC and OA. The percent change in atmospheric BC and OA mass is roughly one-tenth of the BASE comparison, yet the percent change in number is actually slightly greater than one-tenth of BASE for

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N80 and N150 (Table 3). MASSX0.1 leads to a larger increase in N10 (1.11 %) than BASE-NOBIOF (0.29 %). MASSX0.1 still increases the number of primary BC and OA particles and thus the condensational sink for sulfuric acid and organics over NOBIOF; however, there is less suppression of nucleation and growth compared to the BASE simulation. Therefore the relative increase in nucleation and growth relative the BASE comparison offsets some of the reduction in primary emissions (Table 3). Ultrafine particles from nucleation have little influence on the mass distribution, and so the large N10 increases have little effect on the DRE. Spatially, these changes are similar to Fig. 6 and thus are not shown. The globally and annually averaged DRE roughly doubles for MASSX2 and is reduced by one-tenth for MASSX0.1 (Fig. 2a). This highlights that the total emission mass is a key factor in determining the magnitude of the DRE (Table 5). Mixing state assumptions lead to substantial variability in the sign and magnitude in the DRE for MASSX2 (+0.039 to  $-0.016 \text{ W m}^{-2}$ ) and change the sign in MASSX0.1 (+0.003 to  $-0.001 \text{ W m}^{-2}$ ). In agreement with previous studies, our calculated DRE is roughly linearly dependent on the source emission strength (Rap et al., 2013; Scott et al., 2014).

Conversely, altering emission particle mass has non-linear effects on the AIE of biofuel aerosol. The non-linear effects complicate the response of the AIE (Table 5), such that increases in primary emission particle number do not always lead to increases in CCN and cloud reflectivity on a global scale. The AIE for MASSX2-NOBIOF and MASSX0.1-NOBIOF is shown in Fig. 8. Doubling the biofuel emission mass leads to a globally annually averaged positive cloud-albedo AIE of  $+0.002 \text{ W m}^{-2}$  (compared to  $-0.006 \text{ W m}^{-2}$  for BASE-NOBIOF). The small positive value is a result of regions experiencing a stronger negative cloud-albedo AIE due to added CDNC from primary emissions, which are more than offset by regions experiencing a stronger positive cloud-albedo AIE due to the suppressed nucleation and particle growth. The increased N40, N80, and N150 due to doubled primary emissions (Table 3) leads to increases in CDNC near source regions; however, these particle number increases also increase the condensation sink of sulfuric acid and organics, further suppressing nucleation and parti-



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size remained constant, this has little effect on particle number (Table 3). This BC:OA change leads to substantial changes in the DRE, but no change in the AIE (Table 4 and Fig. 2). The HIGHBC simulation increases the percent change in atmospheric BC from 30 % in the BASE comparison to 52 %. This large increase in BC increases the DRE for all mixing state assumptions. In this comparison all mixing state assumptions give a positive DRE ranging from  $+0.004 \text{ W m}^{-2}$  for the external mixture to  $+0.056 \text{ W m}^{-2}$  for the core-shell with absorptive OA mixture. The HIGHOA simulation increases the concentration of OA, which increases the scattering component of the aerosol mixture. This leads to a larger negative DRE relative to the BASE simulation. Relative to NO-BIOF, the core-shell,  $\text{ext} \times 1.5$  with and without absorptive OA, and the external mixture assumptions now give a negative global mean DRE. There is still enough absorptive OA in the core-shell with absorptive OA to have a small positive global mean DRE.

The cloud albedo AIE is relatively unchanged when increasing or decreasing the emissions BC to OA ratio. Both OA and internally mixed BC can contribute to the number of particles that may activate. In HIGHBC and HIGHOA the number of particles that may activate is similar to BASE. In addition, the hygroscopicity parameter ( $\kappa$ ) changes by less than 0.01 for the HIGHBC and HIGHOA simulations compared to BASE, due in part by our assumption that all species (including non-biofuel species) are internally mixed within each size bin. Thus, the composition change between BC and OA does not greatly change the activation diameters, and so the AIE is unchanged. The patterns in the globally averaged DRE and AIE (Fig. 2a) are repeated in Fig. 2b–d, showing there are no strong regional variations in this sensitivity test. Increasing the relative mass of BC results in the largest positive DRE both in the global average and regionally, where in Beijing values range from  $+0.1$  to  $+1.2 \text{ W m}^{-2}$  (Fig. 2c).

### 3.5 Sensitivity of radiative effects due to emissions size distributions

We test the sensitivity of the DRE and AIE to the emissions size distribution to account for uncertainties in fresh and aged plumes. Changes to the emissions size distributions for BC and OA lead to significant changes in both the DRE and AIE. In simulations

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SIZE200, SIZE30, SIZEWIDE, and SIZENARR, we change the emission size distribution while keeping emission mass and composition constant (see Table 2). However, shifting the emission size distribution while keeping mass constant does necessitate a change in emitted particle number and surface area. Increasing the number of primary emitted particles may increase the number of CCN near sources, while potentially decreasing the number of CCN downwind and aloft due to suppression of nucleation and growth. The sign of the AIE will depend on the relative effects from primary particles (which increase AIE) vs. suppression of nucleation/growth (which decreases AIE). Figure 9 contains the change in globally averaged differences in the (a) modeled number distribution, (b) Fuchs surface area distribution, and (c) volume distribution for the BASE-NOBIOF (black line), SIZE30-NOBIOF (blue line with squares), SIZE200-NOBIOF (red line with triangles), SIZENARR-NOBIOF (green line with diamonds), and SIZEWIDE-NOBIOF (magenta line with circles) comparisons. We will use Fig. 9 below to help understand the climate effects of changing the emissions size distribution.

The total BC and OA emissions mass in these simulations does not change relative to the BASE-NOBIOF comparison (Table 3). Altering the emission size distribution does shift the modeled volume/mass distribution relative to the NOBIOF simulation (Fig. 9c). Increasing the GMD (SIZE200) or increasing the SD of the size distribution (SIZEWIDE) predicts a greater positive DRE relative to the BASE case for all mixing states (Table 4 and Fig. 2a). In these two simulations, the mass distribution is shifted to larger size bins (Fig. 9c, red and magenta lines), which increases scattering and absorption; however, the fractional increase in the absorption is larger than that for scattering, which lowers the single scattering albedo and leads to a more positive DRE relative to BASE. The opposite is true for SIZE30 and SIZENARR. In these simulations, the mass distribution is shifted to smaller sizes (Fig. 9c, blue and green lines), causing absorption and scattering to decrease. The fractional decrease in absorption is greater than the fractional decrease in scattering, resulting in a larger single scattering albedo and lower DRE relative to BASE. The DRE ranges from positive to negative across



mixing states for all size sensitivity simulations except SIZEWIDE, which has a low value of  $+0.001 \text{ W m}^{-2}$ .

Including primary biofuel emissions (BASE) increases the Fuchs surface area (i.e. the condensation sink as a function of size) over NOBIOF (Fig. 9b, black line), which increases the condensation sink and suppresses nucleation. There is a slight negative change in the number of nucleation mode ( $< 10 \text{ nm}$ ) particles for the BASE case relative to NOBIOF (Fig. 9a). The suppressed nucleation from the increased Fuchs surface area is partly balanced by small increases in sulfur dioxide from biofuel combustion, which leads to more nucleation and growth via gas-phase sulfuric acid formation. In SIZE30 and SIZENARR, the increased number of primary emitted particles leads to larger integrated increases in Fuchs surface area compared to BASE-NOBIOF (Fig. 9b, blue and green lines), leading to a much stronger condensation sink and suppression of nucleation. The net effect is an increase in accumulation-mode particles due to primary emissions and a decrease in nucleation mode particles due to suppression of nucleation compared to the BASE-NOBIOF. Conversely, the decreased number of primary emitted particles in SIZE200 and SIZEWIDE decreases the Fuchs surface area relative the BASE-NOBIOF comparison (Fig. 9b, red and magenta lines). The unchanged sulfur dioxide emissions combined with reduced Fuchs surface area increase the rate of nucleation and condensational growth. The reduced suppression of nucleation and condensational growth leads to increases in particle number relative the BASE-NOBIOF comparison up to the 100 nm size bin.

The net result is an increased negative AIE for all four size sensitivity simulations relative BASE-NOBIOF. This is caused by either increased primary emitted particle number in source regions (SIZE30 and SIZENARR) or reduced suppression of nucleation and growth (SIZE200 and SIZEWIDE). Thus, the emissions size distribution in the BASE simulation leads to a lower magnitude AIE than if the size distribution was made larger, smaller, narrower or wider in our model. As with the sensitivity tests due to mass, this shows a non-linear relationship to primary biofuel emissions and the globally averaged cloud-albedo AIE.

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Conversely, increasing the number of primary emitted particles relative to BASE (SIZE30 and SIZENARR) does lead to a larger RMS response in the AIE, while reducing the number of primary emitted particles weakens the AIE response (Fig. 2b). This is because the RMS is dominated by the large cloud-albedo AIE in primary emissions regions. Related to this point, in Addis Ababa (Fig. 2d) where the microphysical feedback is weaker, increases in primary emitted particle number (SIZE30 and SIZENARR) greatly increase the magnitude of the negative cloud-albedo AIE relative to BASE, and the AIE is reduced relative to BASE when primary emitted number is reduced (SIZE200 and SIZEWIDE). This emphasizes that the global mean does not always capture the sign and magnitude of regional aerosol-climate effects.

### 3.6 Sensitivity to changes in hydrophilicity

Altering the fraction of emitted mass that is hydrophilic caused negligible change in aerosol mass and number (Table 3) and in the DRE or AIE (Table 4 and Fig. 2). In this model, conversion from hydrophobic to hydrophilic is represented as a fixed e-folding timescale of 1.15 days. This rapid conversion prevents large changes in number concentration from enhanced wet deposition or cloud droplet activation due to changing hygroscopicity. It is plausible that with online ageing there may be a greater effect, for example if the model included aging timescales which are spatially variable due to the availability of hydrophilic material.

### 3.7 Coal as household fuel

Coal is a common household fuel in some regions of the world and is used for both heating and cooking. Household coal use is especially prevalent in China (Legros et al., 2009). Although residential coal combustion is not included in the biofuel inventory used here (Bond et al., 2007), we include an additional simulation to compare to other studies focusing on the residential sector. In this section, we compare the BASE simulation to a simulation with no biofuel emissions over the globe and no residential coal

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emissions over Asia (noSTREET) (as mentioned earlier, residential coal emissions outside of Asia are included with other sources in GEOS-Chem and cannot be isolated). Coal generally has a higher sulfur content than the biofuels (Grieshop et al., 2011), and so emits  $\text{SO}_2$  along with BC and OA. In GEOS-Chem we are further limited by only being able to isolate residential  $\text{SO}_2$  emissions and not BC and OA from coal combustion. The increased  $\text{SO}_2$  emissions lead to a stronger scattering component and thus reduced positive DRE across all mixing states. The DRE for the explicit core-shell mixture for BASE-noSTREET is shown in Fig. 10 (top). The added emissions push the DRE in the negative direction for all mixing states (Fig. 2a and Table 4). Emissions of  $\text{SO}_2$  over Asia increase the magnitude of the negative DRE over eastern China and the Indian Ocean. The magnitude of the positive DRE is generally decreased over India and Tibet. Transport of emissions leads to an increased negative DRE throughout the Northern Hemisphere mid-latitude oceans compared to the BASE-NOBIOF comparison. In Fig. 2b, the RMS value for the DRE is largely similar to the BASE comparison for the absorptive OA and internal mixing states. This is due to regions of reduced positive DRE being compensated by increased regions of negative DRE. The external, ext\*1.5, and core-shell mixtures have a larger RMS value due to an increased negative DRE over China and an increased negative DRE over oceans. The short atmospheric lifetime of aerosol limits the change in DRE. The added coal emissions lead to substantial reduction in the positive DRE over Beijing, but no change over Addis Ababa due to no changes in emissions in Africa (Fig. 2c and d).

The annually averaged percent change in N40 and N80 in the boundary layer (a and b) is positive throughout all of Asia, with heavy source regions increasing by 10–20% (Fig. 11). Increases in the Asian region are significantly greater than in the BASE-NOBIOF comparison (see Fig. 4b and c). Additionally, transported particles lead to increases in N40 and N80 over the Pacific Ocean. Figure 11c and d contain the corresponding zonally averaged N40 and N80 percent changes with pressure level. In contrast to the BASE-NOBIOF comparisons (Fig. 5c and d), addition of household coal use leads to higher increases in N40 and N80 in the Northern Hemisphere trop-

ics and mid-latitudes from the surface to around 200 hPa. The cloud-albedo AIE for BASE-noSTREET is plotted in Fig. 10 (bottom). Residential emissions lead to negative cloud-albedo AIE values of  $-0.2$  to  $-0.4 \text{ W m}^{-2}$  locally over eastern China, with transport of N40 and N80 leading to negative effects of  $-0.1$  to  $-0.3 \text{ W m}^{-2}$  over the Pacific Ocean. India also experiences a negative cloud-albedo AIE of at least  $-0.01 \text{ W m}^{-2}$  due to residential emissions. Increased  $\text{SO}_2$  mass leads to increases in sulfuric acid concentrations, which can offset the condensational sink caused by primary BC and OA particles. This leads to an increased negative cloud-albedo AIE relative to the BASE-NOBIOF comparison, both in the global arithmetic mean ( $-0.006$  to  $-0.019 \text{ W m}^{-2}$ ) and the RMS ( $0.035$  to  $0.058 \text{ W m}^{-2}$ ). The BASE-noSTREET comparison predicts the largest (negative) cloud-albedo AIE for all simulations over Beijing, but similar to the DRE, the cloud-albedo AIE over Addis Ababa is unchanged relative to BASE-NOBIOF (Fig. 2c and d) due to the lack of emissions changes in Africa.

### 3.8 Changing nucleation and background SOA

To explore the sensitivity of the cloud-albedo AIE to other common assumptions in aerosol microphysics models, we run two simulations that lead to variations in the strength of nucleation/growth feedbacks. In BASE\_ACT and NOBIOF\_ACT, we use the activation-nucleation scheme, which predicts more nucleation over oceans than the ternary scheme (used in BASE-NOBIOF) because of low  $\text{NH}_3$  concentrations over the ocean. Stronger nucleation rates mean a larger source of N40 and N80 from nucleation followed by growth, and modulations to nucleation and growth via changing the condensation sink have larger effects on N40 and N80. Addition of biofuel emissions thus reduces N40 and N80 over oceans in these activation-nucleation simulations more strongly than in simulations with ternary nucleation. The simulations with the activation scheme (BASE\_ACT-NOBIOF\_ACT) result in decreases in N10 ( $-0.52\%$ ) and smaller increases in N40 ( $0.30\%$ ) and N80 ( $1.10\%$ ) than the simulations with the ternary scheme (BASE-NOBIOF) (Table 3). The increased strength of the nucleation/growth feedbacks leads to decreases in CDNC and a positive globally aver-

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lower background concentrations lead to a larger percent increase in N80 and N150. Stronger increases in N150 for the bSOA simulations limit maximum supersaturation over polluted areas of China and Europe leading to fewer activated particles and a positive cloud-albedo AIE (+0.021 W m<sup>-2</sup> over Beijing, Fig. 2c). The stronger microphysical feedback also leads to more areas of positive AIE in southern oceans. The net result is a slight globally averaged positive cloud-albedo AIE of +0.002 W m<sup>-2</sup>.

## 4 Conclusions

In this paper, we calculate changes to simulated aerosol concentrations in a global model due to the inclusion of biofuel emissions and evaluate the associated direct and indirect radiative effects. We test the sensitivity of these changes to our assumptions about biofuel emissions mass, composition, size and optical properties, as well as model nucleation and background SOA. We find substantial variability in both the sign and magnitude of the globally and annually averaged direct radiative effect (DRE) of biofuel aerosol due to assumptions regarding mixing state across different model simulations. We find the global-mean DRE due to biofuel emissions ranges from +0.06 to -0.02 W m<sup>-2</sup> considering all simulation/mixing state combinations. The cloud-albedo aerosol indirect effect (AIE) also varies between positive and negative in the global average (-0.02 to +0.01 W m<sup>-2</sup>). Regionally, the DRE and AIE due to biofuel emissions can also vary substantially (Fig. 2c and d). In regions of heavy biofuel combustion where background pollution is also high (e.g. Beijing, Fig. 2c), the DRE can dominate over the AIE. The reduced (and slightly positive) AIE in polluted source regions compared to relatively cleaner regions is a result of an increased condensation sink of sulfuric acid/organics (suppressing nucleation and condensational growth) as well as water vapor (suppressing supersaturation and cloud drop activation). Conversely, in a relatively cleaner source region (Addis Ababa, Fig. 2d), changes to primary emissions dominate the sensitivity of the AIE. Competing regions of positive and negative cloud-albedo AIE limit the magnitude of the global average value. Root-mean-squared val-

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- BC to OA ratio,
- emissions size distribution (including the effects of sub-grid aging/coagulation),
- mixing state for optical calculations.

Without better constraints, even the sign of the net global aerosol effects is uncertain.

5 Previous work has suggested that reducing BC emissions from biofuel sources may be used as a means of countering greenhouse-gas warming effects (Shindell et al., 2012); however, if these suggested aerosol controls include removing both the OA and BC emissions from biofuel sources, it is unclear if a net global cooling will be achievable based on the range of our results.

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**Table 1.** Description of mixing state assumptions.

Mixing State	Morphology	Refractive Indices	Optical Calculation	Absorptive OA
Core-Shell with absorptive OA	Sphere composed of a homogeneous shell surrounding BC core	Shell components are volume weighted while the core is the refractive index of pure BC	BH Mie code for concentric spheres (BH-COAT)	OA absorption calculated using Saleh et al. (2014) and modeled BC to OA ratio
Internal	Homogenous sphere	Volume weighted average of individual indices	Bohren and Huffman (1983) (BH) Mie code for homogenous sphere (BHMIE)	None
ext*1.5 with absorptive OA	Aerosol components are mixed internally except BC, which is a separate particle	Volume weighted and pure BC	BHMIE with the absorption efficiency multiplied by a factor of 1.5	OA absorption calculated using Saleh et al. (2014) and modeled BC to OA ratio
Core-Shell	Same as Core-Shell with absorptive OA	Same as Core-Shell with absorptive OA	Same as Core-Shell with absorptive OA	None
ext*1.5	Same as ext*1.5 with absorptive OA	Same as ext*1.5 with absorptive OA	Same as ext*1.5 with absorptive OA	None
External	Same as ext*1.5 with absorptive OA	Same as ext*1.5 with absorptive OA	BHMIE	None



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**Table 3.** Global, annual-mean percent change in the boundary layer in N10, N40, N80, N150, mass of BC, and mass of OA for the comparisons listed.

Simulation	N10 [%]	N40 [%]	N80 [%]	N150 [%]	BC mass [%]	OA mass [%]
BASE-NOBIOF	0.29	0.93	1.59	2.70	29.5	7.70
MASSX2-NOBIOF	-0.47	1.36	2.95	5.31	59.1	15.4
MASSX0.1-NOBIOF	1.11	0.71	0.55	0.46	2.95	0.75
HIGHBC-NOBIOF	0.35	0.99	1.63	2.74	59.1	6.66
HIGHOA-NOBIOF	0.26	0.90	1.57	2.68	14.8	8.22
SIZE30-NOBIOF	3.60	8.64	10.0	9.61	28.6	7.28
SIZE200-NOBIOF	0.73	0.37	0.28	0.47	29.3	7.60
SIZENARR- NOBIOF	0.16	5.53	10.1	12.2	30.2	7.71
SIZEWIDE-NOBIOF	0.86	0.56	0.48	0.56	27.4	7.07
ALLPHILIC-NOBIOF	0.31	0.95	1.61	2.72	29.5	7.70
ALLPHOBIC-NOBIOF	0.29	0.92	1.58	2.69	29.5	7.70
BASE-noSTREET	1.10	1.78	2.47	3.50	29.5	7.67
BASE_bSOA-NOBIOF_bSOA	-1.19	0.56	2.60	5.11	29.9	23.1
BASE_ACT- NOBIOF_ACT	-0.52	0.30	1.10	2.60	29.53	7.67

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**Table 4.** The global annual-mean all-sky direct radiative effect (DRE) and cloud-albedo aerosol indirect effect (AIE) due to biofuel for the various comparisons. The direct radiative effect was calculated assuming an internal, core-shell with absorptive OA, core-shell, ext\*1.5 with absorptive OA, ext\*1.5, and external mixing state (see Table 1).

Simulation	All-Sky DRE [ $\text{W m}^{-2}$ ]						AIE [ $\text{W m}^{-2}$ ]
	Internal	Core-Shell absorptive OA	Core-Shell	ext*1.5 absorptive OA	ext*1.5	External	
BASE-NOBIOF	0.015	0.021	0.007	0.015	-0.002	-0.008	-0.006
MASSX2-NOBIOF	0.029	0.039	0.013	0.027	-0.004	-0.016	0.002
MASSX0.1-NOBIOF	0.002	0.003	0.001	0.002	0.000	-0.001	-0.014
HIGHBC-NOBIOF	0.047	0.056	0.031	0.045	0.016	0.004	-0.006
HIGHOA-NOBIOF	-0.001	0.002	-0.006	-0.002	-0.011	-0.014	-0.006
SIZE30-NOBIOF	0.008	0.019	0.004	0.015	-0.001	-0.008	-0.021
SIZE200-NOBIOF	0.025	0.023	0.011	0.017	0.002	-0.002	-0.009
SIZENARR-NOBIOF	0.004	0.017	0.002	0.014	-0.002	-0.010	-0.015
SIZEWIDE-NOBIOF	0.025	0.022	0.011	0.017	0.004	0.001	-0.011
ALLPHILIC-NOBIOF	0.015	0.021	0.007	0.015	-0.002	-0.008	-0.007
ALLPHOBIC-NOBIOF	0.015	0.021	0.007	0.015	-0.002	-0.008	-0.006
BASE-noSTREET	0.010	0.016	0.002	0.010	-0.006	-0.012	-0.019
BASE_bSOA-NOBIOF_bSOA	0.013	0.013	0.005	0.008	-0.002	-0.008	0.002
BASE_act-NOBIOF_act	0.015	0.021	0.006	0.015	-0.002	-0.008	0.010

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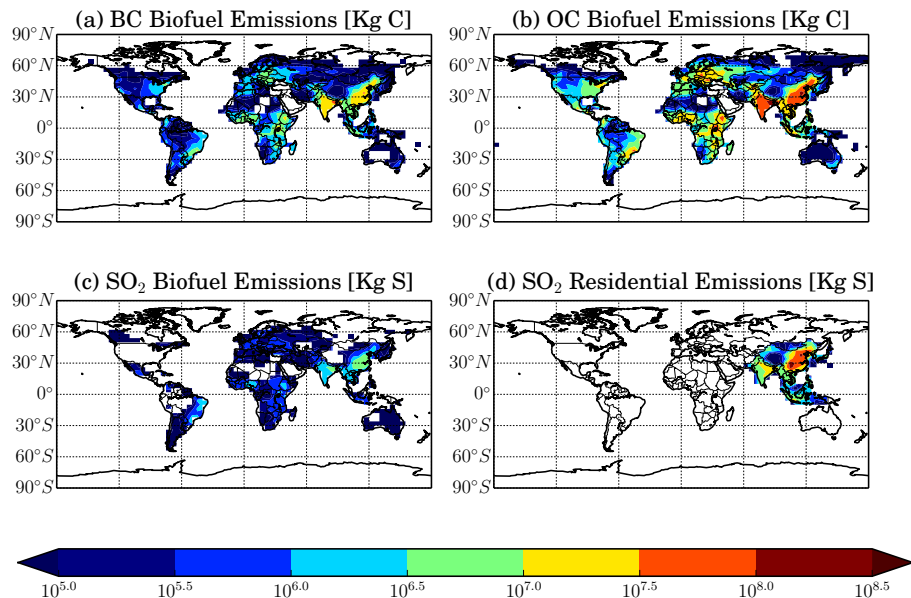
**Table 5.** Overview of key uncertainties and complicating factors that drive the variability in the direct radiative effect and cloud-albedo aerosol indirect effect.

Climate Effect	Key Uncertainties	Complicating Factors
Direct Radiative Effect	<ul style="list-style-type: none"> <li>– BC to OA emission ratio</li> <li>– Emission mass</li> <li>– Optical assumptions (mixing-state and brown carbon)</li> <li>– Emissions size</li> </ul>	Mixing-state and brown carbon vary by region, source category, burn conditions, and atmospheric processing.
Cloud-Albedo Indirect Effect	<ul style="list-style-type: none"> <li>– Emission size distribution</li> <li>– Emission mass</li> </ul>	Feedbacks on nucleation/growth create non-linear effects on CCN and indirect effect downwind and aloft of source regions.



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**Figure 1.** Annual emissions of (a) biofuel black carbon (BC) and (b) organic carbon (OC) emissions [kg C] from Bond et al. (2007), (c) biofuel sulfur-dioxide (SO<sub>2</sub>) emissions [kg S] from EDGAR, and (d) residential SO<sub>2</sub> over Asia (in kg S) from Street's inventory.

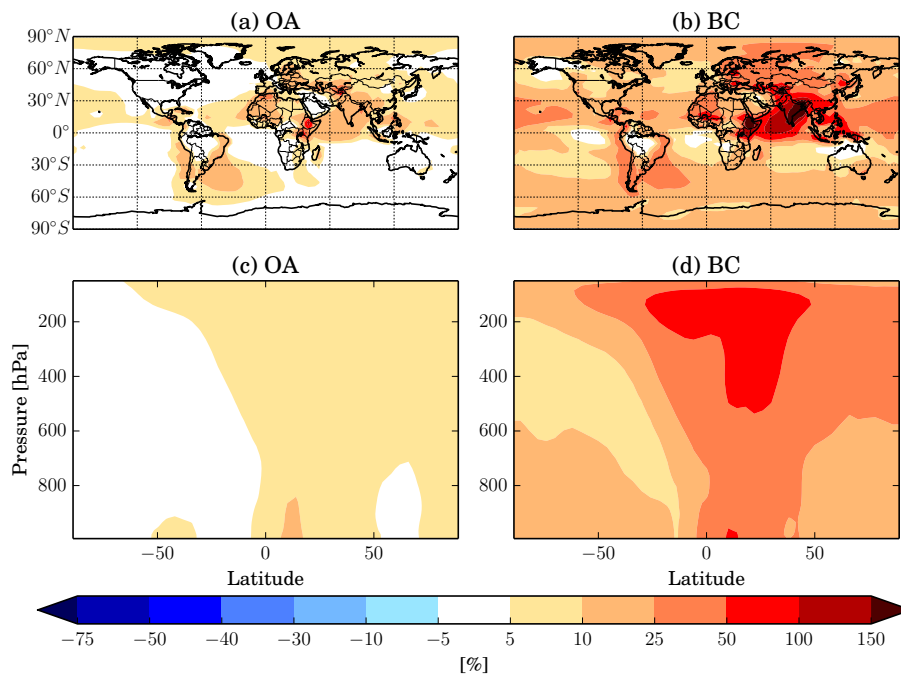
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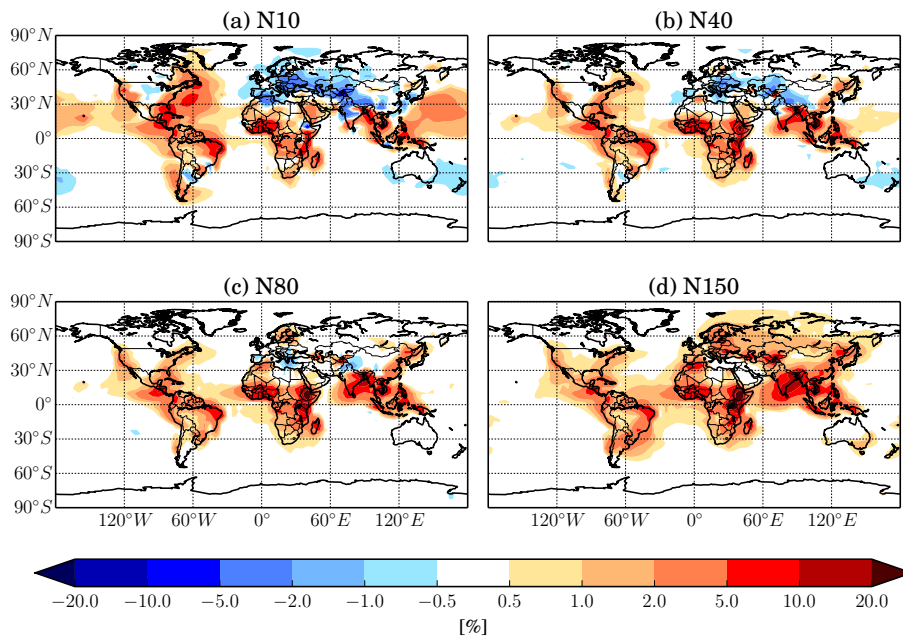


**Figure 3.** The percent change in boundary-layer (a) OA and (b) BC mass and (c and d) zonally-averaged with pressure for the BASE-NOBIOF comparison.

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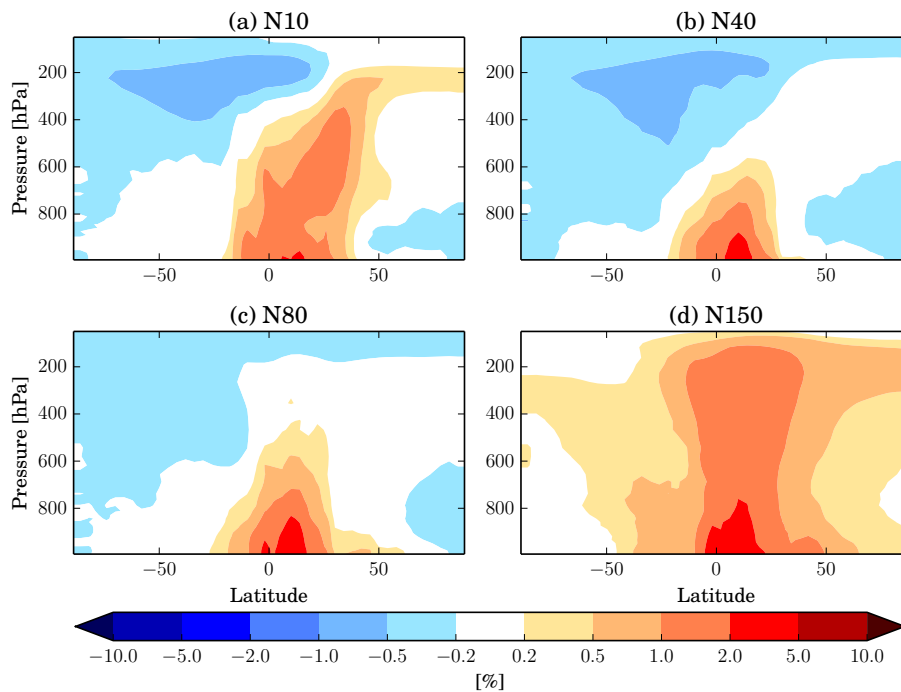


**Figure 4.** The annual-mean percent change in boundary layer (a) N10, (b) N40, (c) N80, and (d) N150 for the BASE-NOBIOF comparison.

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**Figure 5.** The annually and zonally averaged percent change throughout the troposphere for (a) N10, (b) N40, (c) N80, and (d) N150 for the BASE-NOBIOF comparison.

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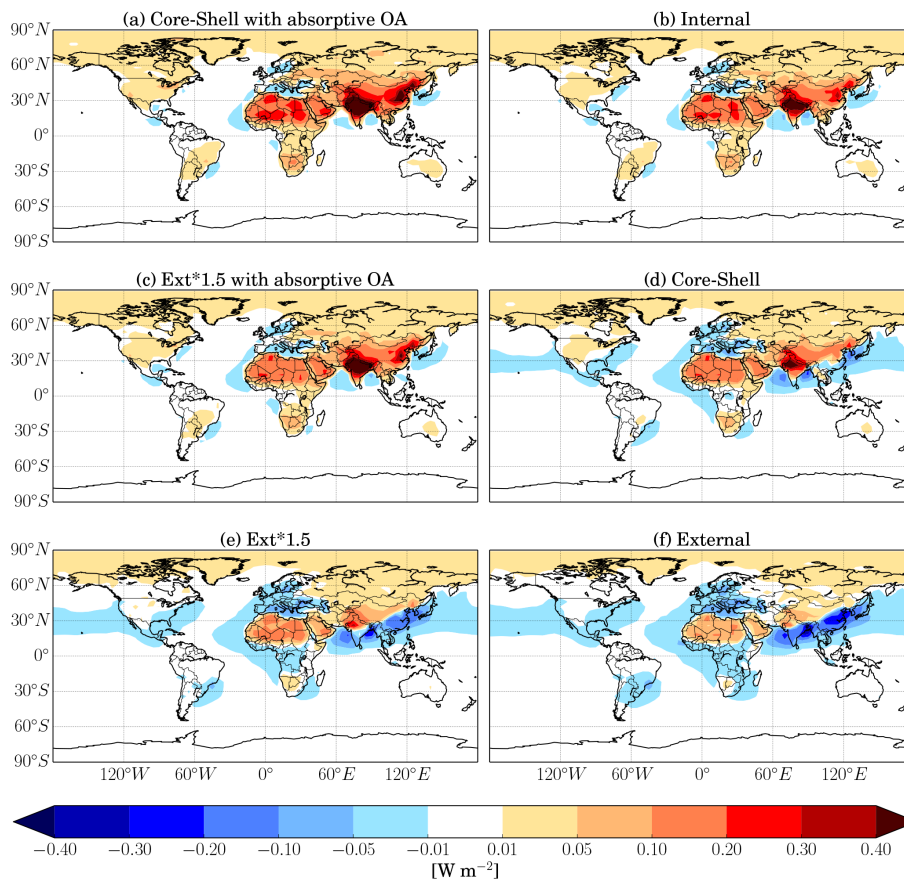
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**Figure 6.** The DRE for the BASE-NOBIOF comparison assuming a core-shell with absorptive (a) OA (mean:  $0.021 \text{ W m}^{-2}$ ), (b) internal (mean:  $0.015 \text{ W m}^{-2}$ ), (c) ext\*1.5 with absorptive OA (mean:  $0.015 \text{ W m}^{-2}$ ), (d) core-shell (mean:  $0.007 \text{ W m}^{-2}$ ), (e) ext\*1.5 (mean:  $-0.002 \text{ W m}^{-2}$ ), and (f) external (mean:  $-0.008 \text{ W m}^{-2}$ ) mixing state.

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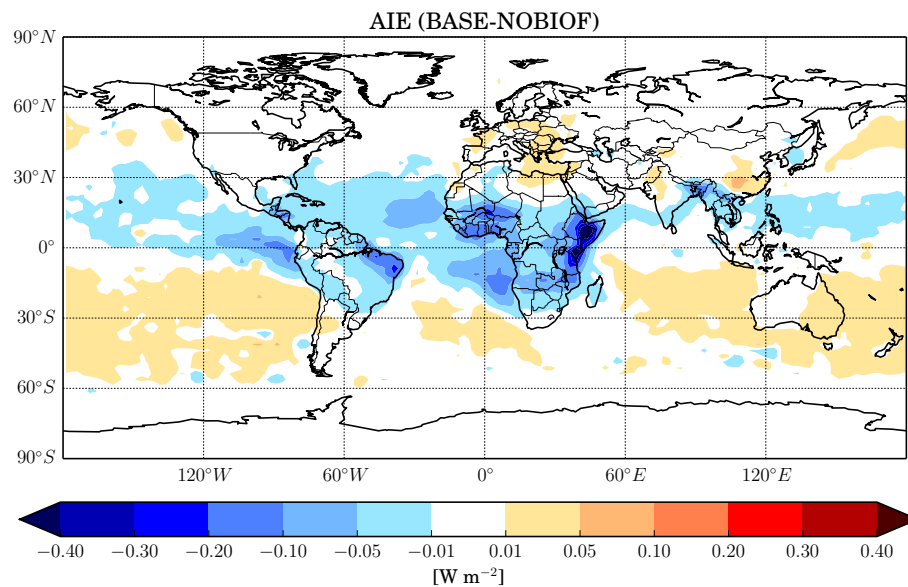
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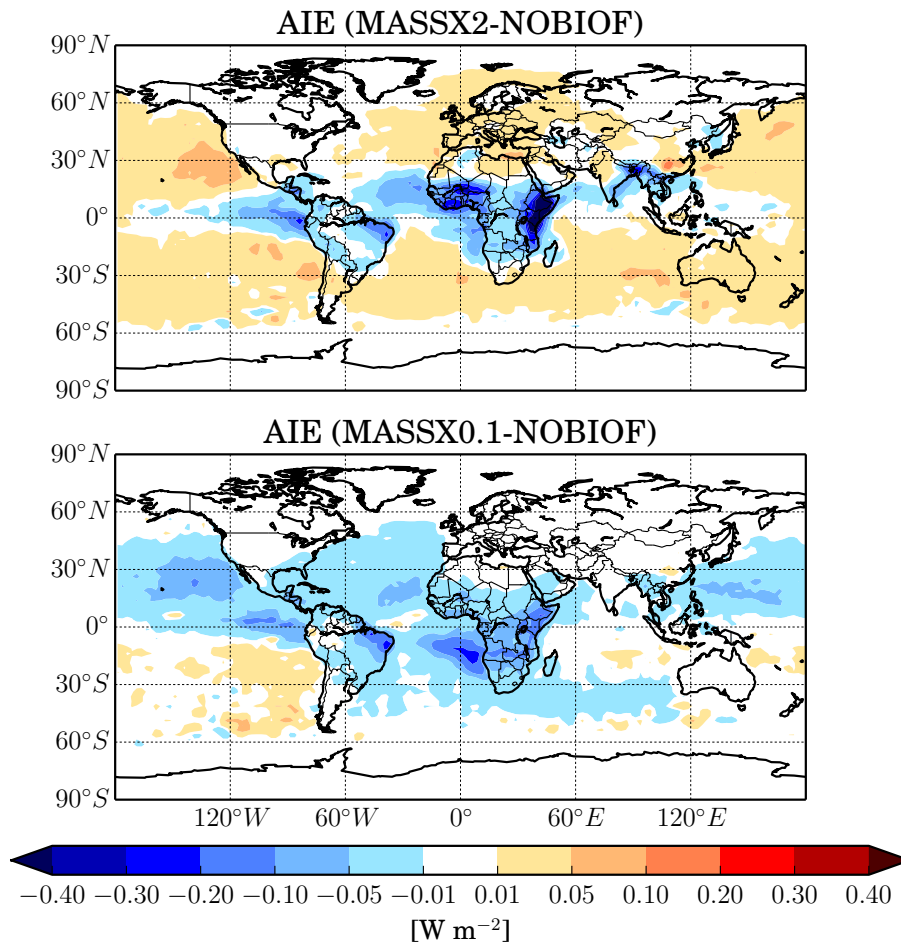
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**Figure 7.** The cloud-albedo AIE for the BASE-NOBIOF comparison (mean:  $-0.006 \text{ W m}^{-2}$ ).

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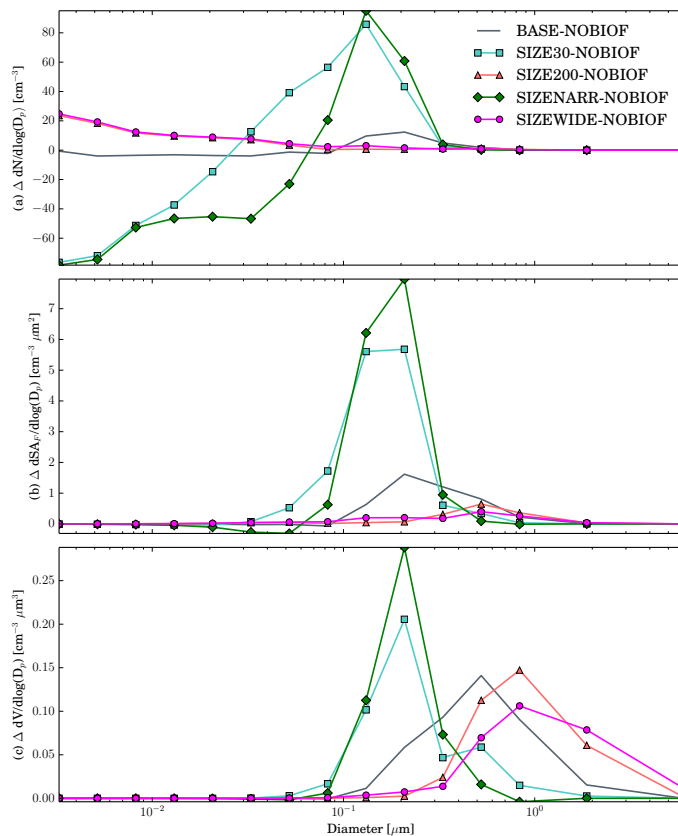


**Figure 8.** The AIE for the MASSX2-NOBIOF comparison (top, mean:  $+0.002 \text{ W m}^{-2}$ ) and MASSX0.1-NOBIOF comparison (bottom, mean:  $-0.014 \text{ W m}^{-2}$ ).

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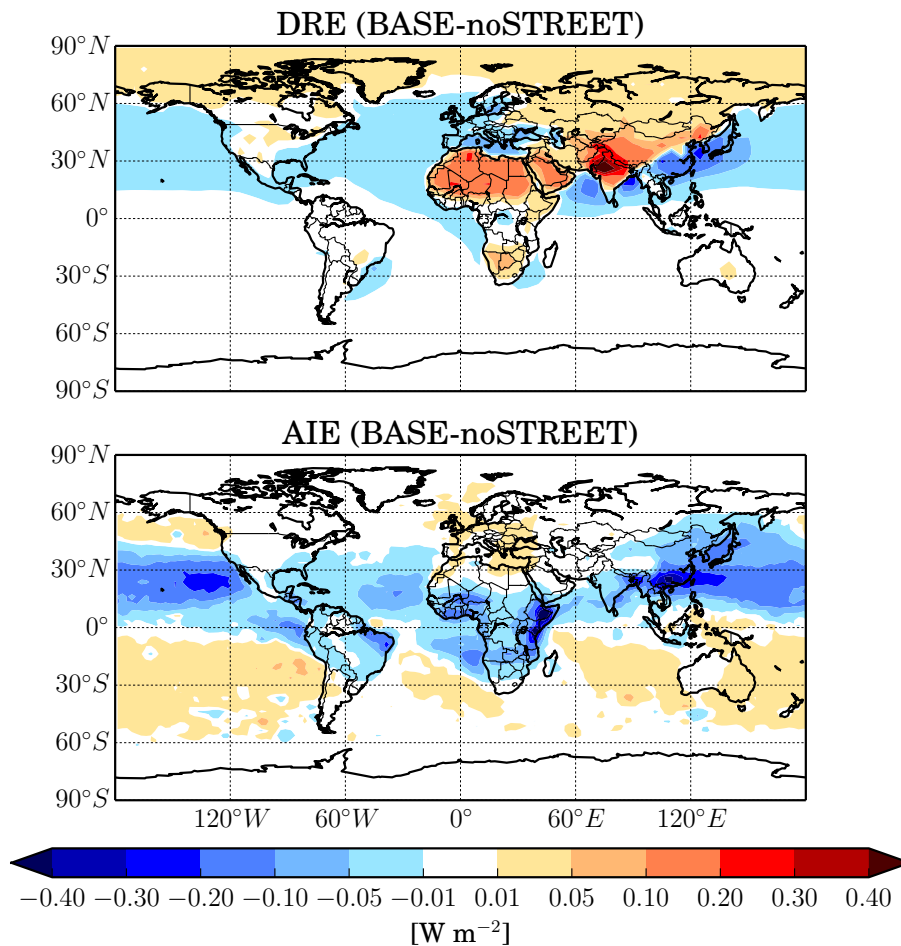


**Figure 9.** Globally averaged **(a)** change in number distribution, **(b)** change in Fuchs surface area distribution, and **(c)** change volume distribution for BASE-NOBIOF, SIZE30-NOBIOF, SIZE200-NOBIOF, SIZENARR-NOBIOF, SIZEWIDE-NOBIOF comparisons. The subtractions isolate the contributions of biofuel emissions to each distribution.



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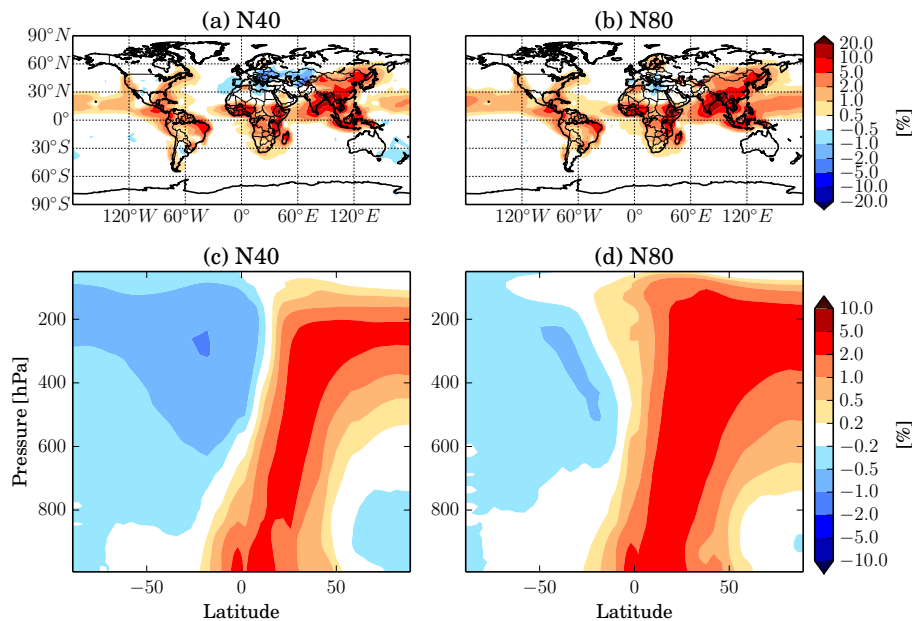
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**Figure 10.** The DRE assuming a core-shell morphology (top, mean:  $0.002 \text{ W m}^{-2}$ ) and the AIE (bottom, mean:  $-0.019 \text{ W m}^{-2}$ ) for the BASE-noSTREET comparison.

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**Figure 11.** The percent change for the BASE-noSTREET comparison in boundary layer (a) N40 and (b) N80 and (c and d) zonally-averaged with height.

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