

1 **The interdependence of continental warm cloud properties**  
2 **derived from unexploited solar background signals in**  
3 **ground-based lidar measurements**

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12

1 **Abstract**

2 We have extensively analysed the interdependence between cloud optical depth, droplet  
3 effective radius, liquid water path (LWP) and geometric thickness for stratiform warm clouds  
4 using ground-based observations. In particular, this analysis uses cloud optical depths  
5 retrieved from untapped solar background signals that are previously unwanted and need to be  
6 removed in most lidar applications. Combining these new optical depth retrievals with radar  
7 and microwave observations at the Atmospheric Radiation Measurement (ARM) Climate  
8 Research Facility in Oklahoma during 2005–2007, we have found that LWP and geometric  
9 thickness increase and follow a power-law relationship with cloud optical depth regardless of  
10 the presence of drizzle; LWP and geometric thickness in drizzling clouds can be generally  
11 20–40% and at least 10% higher than those in non-drizzling clouds, respectively. In contrast,  
12 droplet effective radius shows a negative correlation with optical depth in drizzling clouds  
13 and a positive correlation in non-drizzling clouds, where, for large optical depths, it  
14 asymptotes to 10  $\mu\text{m}$ . This asymptotic behaviour in non-drizzling clouds is found in both  
15 droplet effective radius and optical depth, making it possible to use simple thresholds of  
16 optical depth, droplet size, or a combination of these two variables for drizzle delineation.  
17 This paper demonstrates a new way to enhance ground-based cloud observations and drizzle  
18 delineations using existing lidar networks.

## 1 **1 Introduction**

2 The response of global mean surface temperature to emissions of greenhouse gases from  
3 human activities remains highly uncertain (e.g. Hawkins and Sutton 2009). One of the  
4 primary sources of the uncertainty is how low-topped boundary-layer clouds will respond to  
5 the temperature perturbation and subsequently amplify or dampen climate change (e.g. Bony  
6 and Dufresne 2005; Bony et al. 2006). To improve representations of cloud properties and  
7 their interactions with radiation and water budget in models, sustained efforts have been made  
8 to observe and study marine low-topped clouds (e.g., Martin et al., 1994; Kubar et al., 2009;  
9 Bretherton et al., 2010; Wood, 2012). However, similar efforts have not been made for mid-  
10 latitude continental stratus and stratocumulus clouds, despite their strong links to local  
11 weather and climate (Del Genio and Wolf, 2000; Kollias et al., 2007), and their high  
12 occurrences compared to other cloud types over land (Sassen and Wang, 2008).

13 Ground-based observations for mid-latitude continental clouds are primarily provided by  
14 ARM Climate Research Facility (Stokes and Schwartz, 1994), the NASA Aerosol Robotic  
15 Network (AERONET; Holben et al., 1998), the European project Cloudnet (Illingworth et al.,  
16 2007) and its descendant ACTRIS (Aerosols, Clouds, and Trace gases Research  
17 InfraStructure Network). At the ARM Oklahoma site, low stratiform clouds have been  
18 investigated in a variety of studies, from short-period field campaigns along with airborne  
19 and/or spaceborne measurements (Sassen et al., 1999; Dong et al., 2002; Dong and Mace,  
20 2003) to long-period climatologies (Lazarus et al., 2000; Sengupta et al., 2004; Dong et al.,  
21 2006; Xi et al., 2010). These studies concentrated on variations of liquid water path (LWP),  
22 cloud base height, cloud fraction, and cloud radiative forcing. Surprisingly, little attention is  
23 given to the interdependence between cloud macrophysical, microphysical and optical  
24 properties.

25 The relationship between cloud optical depth and droplet size is of particular interest, because  
26 their correlation patterns are highly related to the stages of warm cloud developments (Suzuki  
27 et al., 2010) and have been used for drizzle delineation (Nauss and Kokhanovsky, 2006;  
28 Suzuki et al., 2011). Using satellite and airborne observations, positive correlations have been  
29 observed in non-drizzling clouds and negative correlations in drizzling clouds (Nakajima et  
30 al., 1991; Nakajima and Nakajima, 1995; Kobayashi and Masuda, 2008), though negative  
31 correlations are not always significant (Harshvardhan et al., 2002).

1 Compared to ARM fixed sites, AERONET and ACTRIS have the advantage of widespread  
2 site locations in mid-latitude continents, but these two networks are not necessarily as fully  
3 equipped as ARM sites. AERONET cloud-mode observations provide information on cloud  
4 optical depth and effective radius (Chiu et al., 2010; 2012), and therefore can be used to  
5 investigate the relationship between cloud microphysical and optical properties. ACTRIS  
6 provides sophisticated information on cloud boundary, water content and drizzle from active  
7 lidars and radars, which can be greatly enhanced by additional cloud optical depth retrievals  
8 to initiate the studies in the interdependence of cloud properties.

9 With enhancing observations of cloud optical depth in mind, this paper introduces a novel  
10 retrieval method for all-sky clouds, using the previously untapped solar background light  
11 measured by ground-based lidars. Because the active laser pulse is rapidly attenuated in thick  
12 liquid clouds, lidar applications have been limited to optically thin clouds and not used to  
13 study stratiform clouds that frequently have optical depth greater than 3. To alleviate this  
14 limitation, Chiu et al. (2007) retrieved optical depth of thick clouds using solar background  
15 light, received along with the active laser pulse but currently treated as the major source of  
16 noise in lidar applications (e.g., Campbell et al., 2002; Welton and Campbell, 2002; Dupont et  
17 al. 2011). However, since the relationship between solar background light and cloud optical  
18 depth is not monotonic (as explained in Sect. 2), Chiu et al. (2007) relied on prior knowledge  
19 of the cloud type and a manual discrimination process to provide retrievals for broken cloud  
20 scenes, an approach which is not ideal for long-term operations.

21 To address this issue, the aims of this paper are 1) to develop and evaluate an objective  
22 discrimination method that works in all-sky conditions; 2) to apply the new retrieval method  
23 to lidar measurements collected at the ARM Oklahoma site where ancillary datasets are  
24 available for intercomparisons; and 3) more importantly, to investigate the interdependence of  
25 cloud macrophysical, microphysical and optical properties. Note that there is an obvious  
26 advantage to using an instrument with a narrow field of view (FOV), typically less than 1  
27 mrad. Compared to conventional cloud optical depth retrieved from hemispheric-viewing  
28 radiometers, lidar provides properties of overhead clouds that potentially correlate better to  
29 liquid water path retrieved from microwave radiometers that have a  $6^\circ$  FOV. Additionally, the  
30 comparable  $0.5^\circ$  FOV of cloud radar, whose measurement is a good indicator of drizzle  
31 presence, significantly mitigates the issue of FOV mismatch when examining the  
32 interdependence of cloud properties for non-drizzling and drizzling clouds.

1 In section 2, we review the retrieval principle and introduce the new discrimination method.  
2 In section 3, we evaluate the performance of our new cloud optical depth against others  
3 retrieved from radiance and irradiance measurements. In section 4, we characterise properties  
4 of stratiform clouds over the ARM Oklahoma site during 2005–2007, and examine the  
5 interdependence of cloud properties for non-drizzling and drizzling clouds. Finally, key  
6 findings and implications of this work are summarised in section 5.

7

## 8 **2 Retrieval Methodology**

9 Prior to July 2006, the micropulse lidar (MPL) at the ARM Oklahoma site was operated at a  
10 wavelength of 523 nm and provided unpolarized measurements at 30 sec intervals. Since July  
11 2006, the lidar operated at 532 nm with polarized measurements at 3–10 sec temporal  
12 resolution. The FOV is 50  $\mu$ rad. Solar background light is estimated from the averaged signal  
13 at lidar range gates between 45 and 55 km where the molecular backscatter is negligible, and  
14 is calibrated against principal plane measurements from AERONET to account for lidar filter  
15 degradation and window cleanliness.

16 Note that for sites where collocated AERONET measurements are unavailable, one can  
17 calibrate solar background light by capitalising on the optical depth of thin clouds retrieved  
18 from active lidar signals. Specifically, radiance can be calculated through radiative transfer  
19 using thin cloud properties as input, and then be further used to calibrate the corresponding  
20 measured solar background light. Details of this alternative calibration approach can be found  
21 in Yang et al. (2008).

### 22 **2.1 Retrieving cloud optical depth from calibrated solar background light**

23 Solar background light received by a lidar is a function of cloud optical depth, cloud effective  
24 radius, cloud fraction, surface albedo, and solar zenith angle. Figure 1a shows that calibrated  
25 solar background light increases with cloud optical depth for optically thin clouds due to  
26 increasing scattering of solar radiation into the FOV, and decreases for optically thick clouds  
27 due to increasing attenuation, resulting in a non-monotonic relationship. For a given optical  
28 depth at lidar wavelengths, a larger effective radius and brighter surface will result in more  
29 observed solar background light. Since the FOV of lidars is small, the cloud cover for each  
30 profile is assumed to be either 0 for clear-sky or 1 for cloudy situations. This assumption is  
31 generally valid, although it becomes problematic near cloud edges when integrating signals

1 from both clear and cloudy sky, which was particularly prevalent in early observations when  
2 the lidar integration time was 30 sec.

3 Cloud optical depth is retrieved by comparing the observed calibrated solar background light  
4 with lookup tables, computed from the discrete-ordinate-method radiative transfer model  
5 (DISORT; Stamnes et al., 1988) with an assumed cloud effective radius and surface albedo  
6 over a range of solar zenith angle up to  $70^\circ$ . We assume that cloud effective radius follows a  
7 Normal distribution with a climatological mean (e.g.,  $8 \mu\text{m}$  for the ARM Oklahoma site) and  
8 a standard deviation of 25% based on the uncertainty found in effective radius retrievals (c.f.,  
9 Table 3 and 5 in Chiu et al., 2012). Surface albedo is estimated using collection 5 products  
10 from MODIS Terra/Aqua combined data at 500 m resolution with an uncertainty of 10%  
11 (Schaaf et al., 2002). We also include a 5% uncertainty in the calibrated solar background  
12 light, regarded as typical for radiance measurements (Holben et al., 1998). With the  
13 uncertainties for all input parameters defined, we perturb these parameters 40 times with  
14 values randomly drawn from Normal distributions and retrieve cloud optical depth; the final  
15 cloud optical depth is reported as the mean and standard deviation of these 40 retrievals. The  
16 choice of 40 repetitions is arbitrary, but it affects retrievals insignificantly by 2% compared to  
17 results from 1000 repetitions (Chiu et al., 2012). The overall retrieval uncertainty in cloud  
18 optical depth is  $\sim 10\%$ . Note that with an uncertainty of 10% rather than 5% in calibrated solar  
19 background light, the overall retrieval uncertainty in cloud optical depth will increase to 17–  
20 25%.

21 Since the relationship between zenith radiance and cloud optical depth is not monotonic, the  
22 aforementioned retrieval process results in two possible solutions at a given radiance; one  
23 corresponds to optically thin clouds, the other corresponds to optically thick. To remove this  
24 retrieval ambiguity, Chiu et al. (2007) applied a manual screening. Here we have developed  
25 an objective discrimination method using lidar backscatter measurements. We calibrated lidar  
26 backscatter signals in clear-air periods using the known molecular scattering at the lidar  
27 wavelength. Since the lidar energy was monitored and the lidar optics were assumed to not  
28 vary significantly, calibration coefficients from a suitable clear-air period were then  
29 extrapolated into cloudy periods. Figure 1b shows an example of the vertical profiles of  
30 calibrated attenuated backscatter signals for optically thin and thick clouds. For thick clouds,  
31 the attenuated backscatter signal drops dramatically above the apparent cloud top; the mean  
32 logarithm (base 10) of the lidar signal from the cloud top to the layer 1 km above is around –

1 7.5. In contrast, for optically thin clouds the mean logarithm value above cloud tops is around  
2  $-6.0$ . The difference between these two mean values is significant, suggesting that this  
3 parameter can be used to discriminate between optically thin and thick clouds; however, a  
4 proper threshold needs to be determined objectively, as described next. For convenience, the  
5 mean of the lidar attenuated backscatter signal from the apparent, or detectable, cloud top to  
6 the level 1 km above is denoted as  $\beta_{\text{ct},1\text{km}}$  hereafter.

7 The threshold of  $\beta_{\text{ct},1\text{km}}$  for discriminating cloud optical depth was determined through cases  
8 selected objectively using retrievals from shortwave narrowband irradiance measurements  
9 (Min and Harrison, 1996), available in the ARM Archive. These cases represent clear or  
10 optically thin clouds, selected when the irradiance-based cloud optical depths were smaller  
11 than 5 for at least 60 consecutive minutes. The threshold of optical depth 5 was chosen  
12 because the zenith radiance typically peaks at this optical depth, and because the lidar signal  
13 tends to be completely attenuated beyond this value. For ARM unpolarized lidar  
14 measurements, Fig. 2 shows that  $\beta_{\text{ct},1\text{km}}$  values range between  $-8.2$  and  $-5.6$ , and 94% of cases  
15 have values of  $\beta_{\text{ct},1\text{km}}$  greater than  $-7.0$ . For ARM polarized measurements, the threshold  
16  $\beta_{\text{ct},1\text{km}}$  of  $-6.8$  leads to a similar fraction 95% of clear-sky cases. Since this threshold does not  
17 vary much over time, we then used  $\beta_{\text{ct},1\text{km}}$  thresholds of  $-7.0$  and  $-6.8$  for unpolarized and  
18 polarized measurements, respectively, throughout the entire analysis.

19 Finally, since our lookup tables were based on liquid water clouds, ice clouds were excluded  
20 using the lidar depolarization ratio and cloud base height. Based on 5-year ground-based lidar  
21 and radiosonde measurements, Naud et al. (2010) suggested a depolarization ratio threshold  
22 of 11% for differentiating ice from liquid. We found that this threshold generally worked  
23 well, but occasionally missed ice clouds when cloud bases were high or clouds were not  
24 sufficiently thick. To mitigate these issues, a second criterion involving cloud base height was  
25 applied. Based on airborne lidar measurements, Hogan et al. (2004) conducted a global  
26 investigation of stratiform supercooled liquid water clouds and showed that less than 10% of  
27 supercooled liquid water clouds occurred at temperatures colder than  $-20^{\circ}\text{C}$ . This temperature  
28 threshold approximately corresponds to an altitude of 7 km at the ARM Oklahoma site during  
29 summer seasons; any clouds located higher than 7 km were excluded and not retrieved in this  
30 study. When lidar depolarization ratio was not available, we used merged sounding data and  
31 excluded cases with apparent cloud tops (identified by lidar) above the freezing level. Note  
32 that these exclusion criteria are simple yet imperfect, particularly when clouds are thick and

1 lidar cannot detect the true cloud top. Therefore, we further excluded time periods when 1-  
2 min ice water path (IWP) were greater than zero, based on retrievals from the Cloudnet  
3 algorithm that uses empirical relationships between ice water content, radar reflectivity and  
4 temperature (Hogan et al., 2006).

5

## 6 **2.2 Calculating cloud effective radius and discriminating drizzling clouds**

7 Once cloud optical depth is retrieved, cloud effective radius can be estimated by combining  
8 liquid water path (LWP) with two commonly used approaches. The first assumes a constant  
9 effective radius in the vertical (Stephens, 1978) and the second assumes a constant cloud  
10 droplet number concentration and a linear increase of liquid water content in the vertical  
11 (Wood and Hartmann, 2006). Using simultaneous retrievals of cloud optical depth and  
12 effective radius at the ARM Oklahoma site, Chiu et al. (2012) found that the second  
13 assumption led to a better agreement with LWP measured by microwave radiometers (MWR)  
14 in all sky conditions. Thus, we estimated cloud effective radius  $r_{\text{eff}}$  by:

$$15 \quad r_{\text{eff}} = \frac{9}{5} \cdot \frac{LWP}{\rho_w \tau}, \quad (1)$$

16 where  $\rho_w$  is the density of water, and  $\tau$  is cloud optical depth. LWP retrievals are available in  
17 the ARM Archive MWRRET product with an uncertainty of 20–30 g m<sup>-2</sup> and a 20-s time  
18 resolution, based on Turner et al. (2007) using 2-channel microwave radiometers.

19 To investigate how the interdependence of cloud macrophysical and microphysical properties  
20 on  $\tau$  differs between non-drizzling and drizzling clouds, we used the ARM Active Remotely  
21 Sensed Clouds Locations product (ARSCL; Clothiaux et al., 2000) for estimating cloud  
22 geometric thickness and for diagnosing drizzling clouds. Combining measurements of cloud  
23 radar, micropulse lidar, and ceilometer, ARSCL provides cloud boundary heights and  
24 reflectivity at 10-s resolution and 45-m vertical resolution. Cloud geometric thickness was  
25 derived from the lowest cloud base (typically detected by lidar) and the cloud top height  
26 (detected by radar). We restrict our analysis to single-layer warm clouds by selecting cases  
27 with geometrical thicknesses less than 1.5 km, minimising cases of multi-layer precipitating  
28 clouds that are hard to separate by radar reflectivity and could be erroneously identified as  
29 single-layer. When clouds were sufficiently thick and no significant radar returns were



1 detected, no valid geometric thickness could be obtained and thus such clouds were omitted in  
2 our analysis.

3 Additionally, drizzle discrimination was based on radar reflectivity ( $Z$ ) at the lowest cloud  
4 base. Similar to Suzuki et al. (2011), we identify clouds as ‘non-drizzling’ if  $Z$  is less than  $-15$   
5 dBZ, and ‘drizzling’ if  $Z$  is greater than  $-15$  dBZ. According to the relationship

6  $R = 0.0788 \cdot Z^{0.75}$  (rain rate  $R$  in  $\text{mm h}^{-1}$  and  $Z$  in  $\text{mm}^6 \text{m}^{-3}$ ) derived from data in Mann et al.  
7 (2013), this threshold of  $-15$  dBZ corresponds to  $\sim 0.006 \text{ mm hr}^{-1}$ .

8

### 9 **3 Evaluation of optical depth retrievals**

10 We evaluate our retrievals against a number of benchmarks. The first benchmark is retrievals  
11 using zenith radiances from AERONET cloud-mode observations that provide unambiguous  
12 cloud optical depth by capitalising on the surface reflectance contrast between 440- and 870-  
13 nm wavelengths (Chiu et al., 2010). This benchmark works for all-sky conditions, but  
14 retrievals are available only when clouds block the Sun so AERONET sunphotometers  
15 operate in cloud-mode rather than normal aerosol-mode. Cloud-mode retrievals (level 1.5) are  
16 available on the AERONET web page.

17 The second benchmark is retrievals from irradiance measurements at 20-s temporal  
18 resolution, available in the ARM Archive. This method uses direct and diffuse transmittance  
19 at 415 nm to estimate cloud optical depth with a default effective radius of  $8 \mu\text{m}$  (Min and  
20 Harrison, 1996a). With additional LWP retrievals from MWR, the estimated optical depth and  
21 initial effective radius are updated iteratively by minimizing least-squares errors in radiance  
22 along with an adjoint radiative transfer method (Min and Harrison, 1996b; Min et al. 2003).  
23 Because irradiances are measured from a hemispheric FOV, this method works best for  
24 relatively overcast homogenous clouds.

25 The third benchmark is retrievals using LWP in the ARM Archive MWRRET product (see  
26 Sec. 2.2) and Eq. (1) with an assumed effective radius of  $8 \mu\text{m}$ , a typical value for the  
27 Oklahoma site (Kim et al., 2003). However, since the true cloud effective radius is not  
28 necessarily  $8 \mu\text{m}$ , we further estimate the potential range of cloud optical depth by varying  
29 effective radius from  $6 \mu\text{m}$  to  $14 \mu\text{m}$ . Clearly, retrieval comparison to the third benchmark is  
30 intended to qualitatively evaluate cloud optical depth variations, rather than a quantitative  
31 measure.

1 In Sect. 3.1, we present intercomparison results from case studies, including broken cloud and  
2 overcast cloud scenes. Additionally, since irradiance-based retrievals work best for  
3 homogenous scenes, we focus on overcast stratiform clouds during 2005–2007 in Sect. 3.2.

### 4 **3.1 Case study**

5 Figure 3 shows time series of lidar backscatter signals and cloud optical depths on 19 April  
6 2005 at the ARM Oklahoma site. The penetrated signal at 17 UTC and the completely  
7 attenuated signal at 20.5 UTC indicate the presence of clear-sky and thick clouds,  
8 respectively. These indications of cloud presence by active lidar signals in Fig. 3a correspond  
9 well to optical depth retrievals in Fig. 3b. Figure 3b also shows that retrievals from calibrated  
10 lidar solar background light agree with those from AERONET cloud mode and from  
11 microwave observations for intermittent and broken cloud situations during 17–18 UTC.

12 Examining two more cases on 10 April and 2 May 2007 when both non-drizzling and  
13 drizzling periods are apparent, Figs. 4 and 5 show consistent agreements between our  
14 retrievals and the benchmark retrievals. Note that ~20% of clouds during 14–18 UTC in Fig.  
15 5 are multi-layered. Since drizzle classification is based on cloud-base reflectivity, these  
16 multilayer clouds are excluded in the following analyses to ensure that LWP, geometric  
17 thickness, optical depth and drizzling characteristics refer to the same lowest cloud layer. In  
18 short, the overall agreement between independent retrievals suggests that the calibration of  
19 solar background light and the newly developed method for distinguishing thin and thick  
20 clouds work well for all-sky conditions.

### 21 **3.2 Stratiform clouds during 2005–2007**

22 This section reports results of intercomparison between retrievals from lidar solar  
23 background, AERONET cloud-mode observations, and from narrowband irradiance  
24 measurements for relatively homogenous and overcast cloud cases. To objectively select  
25 appropriate low-level stratiform water clouds, combined measurements from cloud radar,  
26 micropulse lidar, and ceilometer in the ARSCL product were used to identify 1-hr time  
27 periods with cloud fraction greater than 0.95 and cloud top heights lower than 5 km. Since our  
28 analysis includes several datasets at various temporal resolutions, we average data points over  
29 a 1-min time period. We took a simple linear average for LWP retrievals and radar  
30 reflectivity, but used a logarithm averaging technique for lidar-based cloud optical depth  
31 because transmittance is a concave function of cloud optical depth. In other words, we  
32 averaged the natural logarithm of cloud optical depth, and then transformed the average back

1 to obtain the 1-min mean. Additionally, to use the same dataset for investigating  
2 interdependence of cloud macrophysical, microphysical and optical properties in Sect. 4, we  
3 further excluded time periods if the effective radius was outside the range between 3 and 100  
4  $\mu\text{m}$ . This exclusion process lead to a final sample size of 5,200-min of data points during  
5 2005–2007 that represents  $\sim 35\%$  of daytime stratiform cases.

6 Figure 6 shows histograms of 1-min averaged cloud optical depth and a scatterplot of  
7 retrievals from lidar solar background noise against those from flux measurements. Both  
8 datasets reveal an occurrence peak at optical depth of 15–20, but an evident discrepancy  
9 occurs in the optical depth bin of 0–5. The reason for the lack of small optical depth in lidar-  
10 based retrievals is partly because their corresponding LWP values have been always zero or  
11 negative and therefore are excluded, implying that it remains challenging for 2-channel  
12 microwave radiometers to detect very optically thin clouds.

13 The mean cloud optical depth from lidar measurements is 35, larger than that retrieved from  
14 irradiances by 2 optical depths. A high correlation coefficient of 0.94 is obtained, as shown by  
15 the majority of data points in Fig. 6b lying close to the 1:1 line. In addition, the root-mean-  
16 squared difference between the two is 8 (24% relative to the mean of irradiance-based  
17 retrievals), partly attributed to cases that have much larger lidar-based retrievals than those  
18 from irradiances. Particularly for cases where irradiance-based retrievals are less than 5, we  
19 have found that these points are associated with intermittent cloudy conditions having LWP  
20 between  $-10$  and  $80 \text{ g m}^{-2}$ . Therefore, the discrepancy in cloud optical depth for these data  
21 points is likely because lidar has a narrow FOV to capture larger variations that tend to be  
22 smeared out in irradiance-based retrievals due to a hemispheric FOV of shadowband  
23 radiometers.

24 Similarly, Fig. 7 shows a scatterplot for evaluating retrievals against the AERONET cloud-  
25 mode product. The mean cloud optical depth from lidar measurements is 30, smaller than  
26 cloud-mode retrievals by 3 optical depths. The correlation coefficient is 0.95, while the root-  
27 mean-squared difference between the two is 8 (24% relative to the mean of cloud-mode  
28 retrievals).

29

## 1 **4 Interdependence of stratiform cloud properties**

### 2 **4.1 Macrophysical properties versus optical depth**

3 Using the same stratiform cloud cases shown above, we investigate how cloud macrophysical  
4 and microphysical properties vary with optical depth in non-drizzling and drizzling stratiform  
5 clouds, categorised by a reflectivity threshold of  $-15$  dBZ as described in Sec. 2.2. Fig. 8a  
6 shows that non-drizzling clouds occur more frequently at optical depths of 10–20, while  
7 drizzling clouds have a relatively uniform frequency distribution throughout the entire optical  
8 depth range. Using an adiabatic cloud model for non-drizzling clouds, Boers and Mitchell  
9 (1994) showed that LWP, cloud geometric thickness  $H$  and optical depth  $\tau$  follow  $\text{LWP} \propto H^2$ ,  
10  $\tau \propto H^{5/3}$ , and thus  $\text{LWP} \propto \tau^{6/5}$ . Not surprisingly, Fig. 8b shows that LWP indeed increases  
11 approximately linearly with  $\tau$  for both cloud categories. LWP in non-drizzling clouds is  
12 proportional to  $\tau^{1.09 \pm 0.01}$  with 95% confidence intervals; the exponent is slightly smaller than  
13 the predicted value of 1.2 under an adiabatic assumption. LWP in drizzling clouds is  
14 generally 20–40% larger than those in non-drizzling clouds.

15 Similar to LWP, Fig. 8c shows that  $H$  also increases with  $\tau$ . Using 1-min averaged ARM data  
16 from these stratiform cloud cases, the relationship between  $H$  (in m) and  $\tau$  can be  
17 approximated by:

$$18 \quad H = (308 \pm 15) \cdot \tau^{0.25 \pm 0.01} \text{ for non-drizzling clouds; and} \quad (2)$$

$$19 \quad H = (513 \pm 51) \cdot \tau^{0.16 \pm 0.03} \text{ for drizzling clouds,} \quad (3)$$

20 corresponding to correlation coefficients of 0.95 and 0.79, respectively. These relationships  
21 indicate that the geometric thickness in drizzling clouds is at least 10% larger than that in  
22 non-drizzling clouds at a given  $\tau$ . We have also found that these relationships vary little when  
23 taking hourly means rather than 1-min averages. Using the adiabatic approximation as  
24 explained above, the exponents in non-drizzling and drizzling clouds from ARM data are both  
25 much smaller than the predicted value of 0.6.

26 Cloud geometric thickness derived from Eqs. (2) and (3) is compared to the results for marine  
27 stratocumulus off the coast of California during the First ISCCP Regional Experiment. Based  
28 on Minnis et al. (1992), their relationship between  $H$  and  $\tau$  can be re-written as:

$$29 \quad H = 58 \cdot \tau^{0.56}, \quad (4)$$

1 where  $H$  was retrieved from hourly-averaged surface ceilometer and acoustic sounder  
2 measurements;  $\tau$  was estimated from Geostationary Operational Environmental Satellite  
3 visible and infrared radiances. The relationships obtained here suggest that the geometric  
4 thicknesses in continental stratiform clouds can be thicker than marine stratocumulus by at  
5 least 35% for cloud optical depths less than 80.

## 6 **4.2 Cloud effective radius versus optical depth**

7 Unlike LWP and  $H$ , Fig. 8d shows that cloud effective radius has a different dependence on  
8 optical depth between non-drizzling and drizzling clouds. The strong positive correlation of  
9 0.8 between cloud effective radius and optical depth in non-drizzling clouds is consistent with  
10 many studies using airborne and spaceborne remote sensing measurements (e.g., Han et al.,  
11 1994; Nakajima and Nakajima, 1995; Harshvardhan et al., 2002), but the asymptotic radius  
12 from the ARM data is  $\sim 10 \mu\text{m}$ , smaller than the so-called critical radius ( $\sim 15 \mu\text{m}$ ) reported in  
13 literature for marine low clouds (Nakajima and Nakajima, 1995; Kobayashi and Masuda,  
14 2008; Painemal and Zuidema, 2011). Additionally, these non-drizzling clouds show  $r_{\text{eff}}$   
15 proportional to  $\tau^{0.11 \pm 0.01}$ . The exponent of 0.11 is smaller than the value of 0.2 derived from  
16 satellite and aircraft measurements for the eastern Pacific stratocumulus (Szczodrak et al.  
17 2001), and the theoretical value of 0.2 derived under the assumption of adiabatic and constant  
18 cloud droplet number concentration  $N_d$  (Lohmann et al., 2000), meaning that the condition at  
19 the ARM Oklahoma site may be slightly sub-adiabatic and/or  $N_d$  variation with height is not  
20 negligible.

21 For drizzling clouds, Nakajima and Nakajima (1995) showed that cloud effective radius  
22 decreased from  $20 \mu\text{m}$  to  $10 \mu\text{m}$  with an increase in  $\tau$  from 5 to 20. Similarly, our result shows  
23 a negative correlation ( $-0.75$ ) with a 99% confidence level for drizzling clouds. The negative  
24 correlations between cloud effective radius and optical depth in drizzling clouds can be  
25 explained by precipitation influence, which possibly reduces cloud optical depth through the  
26 removal of droplets (Boers and Rotstajn, 2001). Further analyses reveal that a number of  
27 drizzling clouds with small optical depths indeed have large effective radii greater than  $50$   
28  $\mu\text{m}$ , often found at the end of a precipitation system passing over. These cases, however,  
29 occurred less frequently compared to those with small effective radii, resulting in the mean  
30 cloud effective radii fluctuating between  $12\text{--}18 \mu\text{m}$  with large standard errors at small optical  
31 depths.

1 Since the correlation between  $\tau$  and  $r_{\text{eff}}$  is positive in non-drizzling clouds but negative in  
 2 drizzling clouds, the difference in  $r_{\text{eff}}$  between two types of clouds decreases with increasing  
 3 cloud optical depth, which is as a result of Fig. 8b. Across all optical depth bins, Fig. 8b  
 4 shows that LWP in drizzling clouds is consistently  $\sim 85 \text{ g m}^{-2}$  larger than that in non-drizzling  
 5 clouds. Compared to cases with small  $\tau$ , this extra LWP in drizzling clouds distributes to more  
 6 droplets in cases with large  $\tau$ , leading to a smaller increase in  $r_{\text{eff}}$  (as shown in Eq. (1) having  
 7 a denominator  $\tau$ ).

8 In short, the difference between non-drizzling and drizzling clouds at a given cloud optical  
 9 depth mainly ranges between 2–7  $\mu\text{m}$  with a mean of 5  $\mu\text{m}$  (Fig. 8d), similar to the finding  
 10 from satellite observations in marine stratocumulus (Kubar et al., 2009). This mean size  
 11 difference between two cloud categories is clear in Fig. 9a, showing that the distribution of  
 12 non-drizzling clouds peaks at 6–8  $\mu\text{m}$  with a mean of 8  $\mu\text{m}$ , and the distribution of drizzling  
 13 clouds peaks at 10–12  $\mu\text{m}$  with a mean of 13  $\mu\text{m}$ .

#### 14 **4.3 Interdependence derived from the ARM Min product**

15 To examine whether this interdependence is sensitive to the choice of cloud optical depth  
 16 product, we repeated the same analysis using the ARM Min product as shown in Fig. 10.  
 17 Firstly, similar to results derived from lidar-based retrievals, LWP in non-drizzling clouds is  
 18 proportional to  $\tau^{1.01 \pm 0.01}$  with 95% confidence intervals. LWP in drizzling clouds are also  
 19 larger than those in non-drizzling clouds, although the difference between two varies in a  
 20 larger range between 30 and 150  $\text{g m}^{-2}$ . Secondly, cloud geometric thickness can be  
 21 approximated by:

$$22 \quad H = (249 \pm 12) \cdot \tau^{0.30 \pm 0.01} \text{ for non-drizzling clouds; and} \quad (5)$$

$$23 \quad H = (447 \pm 33) \cdot \tau^{0.20 \pm 0.02} \text{ for drizzling clouds.} \quad (6)$$

24 Although the exponents and prefactors of Eq. (5) and (6) are different from Eq. (2) and (3),  
 25 the geometric thicknesses derived from the ARM Min and lidar-based retrievals are similar as  
 26 shown in Fig. 10c. Finally, while the negative correlation between  $\tau$  and  $r_{\text{eff}}$  in drizzling  
 27 clouds holds in Fig. 10d, the more robust positive correlation in non-drizzling clouds found in  
 28 Fig. 8d and satellite observations disappears due to a relatively flat  $r_{\text{eff}}$  of  $\sim 8 \mu\text{m}$  across all  
 29 optical depth bins.

1 In short, the relationships of LWP and  $H$  with  $\tau$  are not sensitive to the choice of cloud optical  
2 depth product, but this is not true for the correlation between  $\tau$  and  $r_{\text{eff}}$ . This highlights the  
3 potential importance of having comparable FOVs among various instruments for  
4 investigating  $\tau - r_{\text{eff}}$  correlation, although properly addressing this issue may require more  
5 thorough simulation experiments.

6

#### 7 **4.4 Implication on drizzle delineation**

8 Taking a different view, now we use the same dataset as shown in Fig. 9 to investigate how  
9 LWP,  $H$  and  $\tau$  vary with  $r_{\text{eff}}$ . Figures 9b–d show that properties between non-drizzling and  
10 drizzling clouds differ the most in the  $r_{\text{eff}}$  range of 7–11  $\mu\text{m}$ , although this could be a result of  
11 a relatively smaller sample size outside this  $r_{\text{eff}}$  range. Specifically, Fig. 9d shows optical  
12 depth of non-drizzling clouds increases with  $r_{\text{eff}}$  and changes little at  $r_{\text{eff}}$  beyond 7  $\mu\text{m}$ . The  
13 relatively small change in  $\tau$  is also found in the  $r_{\text{eff}}$  range of 7–15  $\mu\text{m}$  for drizzling clouds; this  
14 is similar to the finding in satellite observations (Kobayashi and Masuda, 2008), but their data  
15 showed such behaviour only when  $r_{\text{eff}}$  was larger than a critical value of  $\sim 15 \mu\text{m}$ . Since  
16 Kobayashi and Masuda (2008) used 21-day measurements from the Tropical Rainfall  
17 Measuring Mission satellite and sampled tropical marine warm clouds, the difference in the  
18 critical effective radius (7  $\mu\text{m}$  vs 15  $\mu\text{m}$ ) may be due to the regional variability of  
19 precipitating clouds. Additionally, the definition of this critical effective radius is rather loose  
20 and its value can strongly depend on how and at which altitude cloud effective radii were  
21 estimated. The difference in the resulting critical value of effective radius between  
22 airborne/spaceborne measurements and the ARM data can be partly due to a fact that  
23 retrievals from the former is mainly determined by droplets at cloud tops, while the latter is  
24 determined by the entire cloud layer (Platnick 2000; Chiu et al., 2012).

25 Results from Figs. 8d and 9d imply that it is plausible to delineate drizzling clouds using a  
26 simple threshold; for example, we can roughly classify clouds as drizzling when cloud  
27 effective radius exceeds a critical value  $r^*$  of 10  $\mu\text{m}$  (Fig. 8d) or when cloud optical depth  
28 exceeds 40 (Fig. 9d). Similarly, based on satellite retrievals and ground-based radar  
29 measurements, Nauss and Kokhanovsky (2006) proposed a more sophisticated delineation  
30 function, given as:

$$31 \quad r^* = \frac{A}{\tau}, \quad (7)$$

1 where coefficient  $A$  is  $920 \mu\text{m}$  and the critical value  $r^*$  varies with cloud optical depth  $\tau$ . To  
 2 evaluate how well these methods discriminate between non-drizzling and drizzling clouds  
 3 (i.e., a binary classification), we computed the Heidke skill score (HSS) from a contingency  
 4 table (Table 1), defined as:

$$5 \quad HSS = \frac{2(A \cdot D - B \cdot C)}{(A + C)(C + D) + (A + B)(B + D)}. \quad (8)$$

6 HSS not only measures the proportion of correct classifications (including both correct hits  
 7 and negatives), but more importantly, also takes into account the expected skill obtained by  
 8 chance in the absence of any skill (Barnston, 1992). In general, a HSS of 0 indicates no skill,  
 9 while 1 represents perfect skill.

10 Figure 11 summarises HSS using three different methods. Firstly, using a simple fixed cloud  
 11 effective radius as the delineation threshold (red lines), the optimal threshold that maximises  
 12 HSS is  $10 \mu\text{m}$ , agreeing with results in Fig. 8d. Secondly, applying a fixed threshold of cloud  
 13 optical depth (blue lines), the optimal threshold is  $\sim 42$  and HSS is similar in the optical depth  
 14 range between 40 and 45. Note that the maximum of HSS using the optimal optical depth  
 15 threshold is not as good as that from an effective radius threshold of  $10 \mu\text{m}$ . Thirdly, a  
 16 dynamic threshold of cloud effective radius derived by Eq. (7) apparently yields a higher HSS  
 17 ( $\sim 0.52$ ), compared to the previous two simple methods; the optimal coefficient  $A$  is  $380 \mu\text{m}$ ,  
 18 rather than  $920 \mu\text{m}$  found in satellite observations (Nauss and Kokhanovsky, 2006) for  
 19 convective systems over Central Europe taken during the extreme summer floods in 2002. It  
 20 is expected that the coefficient  $A$  varies with cloud type, site location, and more importantly,  
 21 the threshold of rain rate used to define drizzle ( $\sim 0.006 \text{ mm h}^{-1}$  in our cases).

22 Since HSS is dependent on the frequency of occurrence of an event, we further test our  
 23 delineation thresholds using Symmetric Extremal Dependence Index (SEDI) that is  
 24 independent of occurrence frequency and thus works for both common and rare events (Ferro  
 25 and Stephenson 2011). SEDI is defined as:

$$26 \quad SEDI = \frac{\ln F - \ln H + \ln(1 - H) - \ln(1 - F)}{\ln F + \ln H + \ln(1 - H) + \ln(1 - F)}, \quad (9)$$

27 where

$$28 \quad H = \frac{A}{A + C} \quad \text{and} \quad F = \frac{B}{B + D}.$$



1 Similar to HSS, a SEDI value of 0 indicates no skill, while 1 represents perfect skill. As Fig.  
2 12 shows, the optimal cloud effective radius and cloud optical depth thresholds are  $\sim 10 \mu\text{m}$   
3 and  $\sim 40$ , respectively. When considering a dynamic threshold, the optimal coefficient  $A$  of  
4  $340 \mu\text{m}$  is found. Overall, the optimal thresholds from SEDI are similar to those derived from  
5 HSS.

6 In short, depending on the availability of measurements, one can use a cloud optical depth of  
7 40 as the simplest way for drizzle delineation in the absence of LWP and radar measurements,  
8 although this threshold may depend on ambient aerosol loading. If co-incident LWP  
9 measurements are available, the dynamic threshold of cloud effective radius given in Eq. (7)  
10 with a coefficient  $A$  of  $340\text{--}380 \mu\text{m}$  is a better approach to delineating drizzle for mid-latitude  
11 continental stratiform clouds.

12

## 13 **5 Summary**

14 To better represent clouds in weather and climate models, long-term global measurements can  
15 provide direct constraints and improve our knowledge of cloud and precipitation formation,  
16 and their interactions with radiation and aerosol. In particular, low warm clouds strongly  
17 influence global climate through their impacts on Earth's radiation and water energy cycle.  
18 While marine low clouds have been extensively studied, continental warm clouds received  
19 relatively little attention partly due to the fact that the majority of satellite retrievals work best  
20 over oceans.

21 Using ground-based measurements at the ARM Oklahoma site during 2005–2007, we  
22 conducted an extensive analysis for mid-latitude continental low-level clouds. To retrieve  
23 cloud optical depth, we developed a novel method that capitalised on unexploited solar  
24 background light that is currently treated as noise and has largely inhibited lidar applications  
25 in all-sky conditions and during daytime. This new technique works well; when compared to  
26 other benchmarks, the mean bias of cloud optical depth is around 2 and the root-mean-  
27 squared errors is 8 (23% relative to the mean optical depth). Since lidars have a field-of-view  
28 much closer to those of microwave radiometers than conventional hemispheric-viewing  
29 radiometers, it is more appealing to use lidar-based cloud retrievals to understand the linkage  
30 between cloud macrophysical, microphysical and optical properties.

1 A number of key features are found in the relationships between LWP, geometric thickness  
2  $H$ , droplet effective radius  $r_{\text{eff}}$  and cloud optical depth  $\tau$ . Firstly, LWP and  $H$  follow a power-  
3 law relationship with positive exponents with  $\tau$ ; LWP and  $H$  in drizzling clouds are generally  
4 20–40% and at least 10% higher than those in non-drizzling clouds, respectively. Similar to  
5 LWP,  $r_{\text{eff}}$  also increases with  $\tau$  following a power-law for non-drizzling clouds, but this does  
6 not hold for drizzling clouds. In the presence of drizzle, a negative correlation is found  
7 between  $r_{\text{eff}}$  and  $\tau$ ;  $r_{\text{eff}}$  also tends to be 5  $\mu\text{m}$  larger than droplet sizes in non-drizzling clouds.  
8 While several aircraft and satellite observations have suggested  $r_{\text{eff}}$  on the order of 15  $\mu\text{m}$  may  
9 be a good indicator to distinguish between non-drizzling and drizzling marine clouds, we  
10 found that a threshold of  $\sim 10$   $\mu\text{m}$  works better for ground-based observations. The difference  
11 in threshold between various observational platforms is likely attributed to the fact that  
12 satellite retrievals are mainly determined by properties at cloud tops, and on the contrary,  
13 ground-based retrievals utilise the full cloud profile. If co-incident LWP measurements are  
14 available, a dynamic threshold of cloud effective radius given in Eq. (7) with a coefficient  $A$   
15 of 340–380  $\mu\text{m}$  is a better approach to delineating drizzle for mid-latitude continental  
16 stratiform clouds.  
17 We have demonstrated a novel retrieval method using untapped solar background signals in  
18 lidar measurements, which greatly extends lidar applications from cirrus to all types of  
19 clouds, and provides a new approach to distinguishing between non-drizzling and drizzling  
20 clouds. With collocated radar and LWP measurements, the new retrieval can also be used to  
21 compare and contrast drizzle and drizzle-free cloud properties. This new method can be easily  
22 adapted to the existing lidar networks if collocated sunphotometer measurements are available,  
23 including the high-density ceilometer networks in the United Kingdom, France and Germany  
24 that have been established for monitoring volcanic plumes (Heese et al., 2010). Combined  
25 with the ability of lidars to resolve vertical distributions of aerosol properties below cloud  
26 layers, collocated and simultaneous measurements of aerosol and cloud are also possible,  
27 which can help advance our understanding of aerosol-cloud interactions.

28

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4

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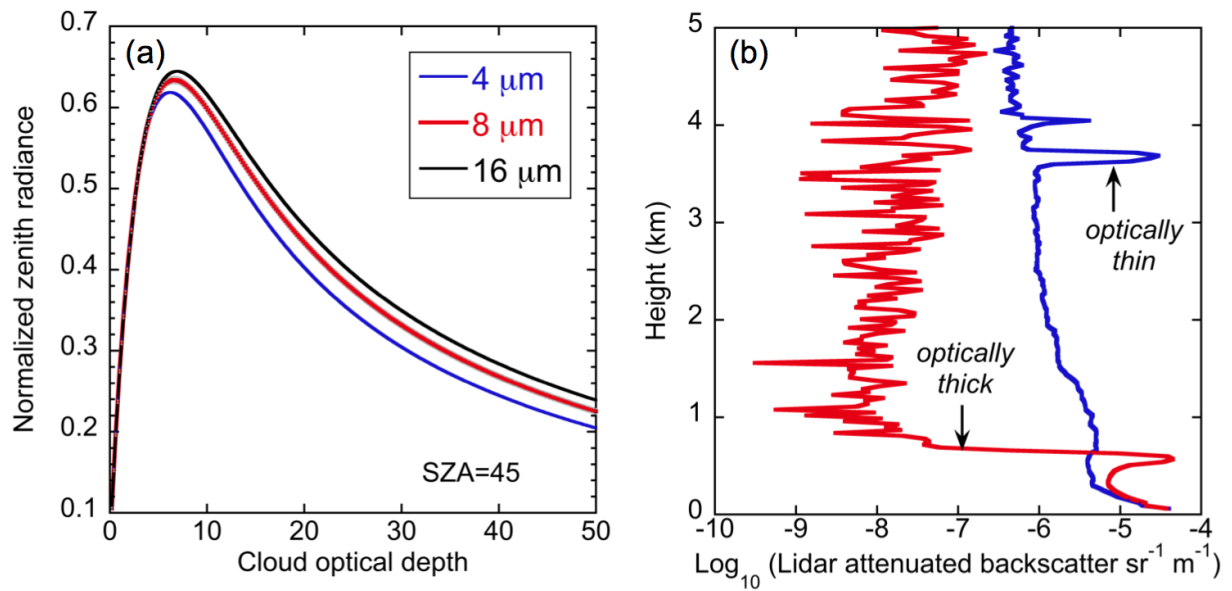
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1 Table 1. Contingency table used to evaluate drizzle delineation methods. A–D represent the  
2 number of hits, false alarms, misses and correct negatives, respectively.

New method	Reference observations	
	Yes	No
Yes	A	B
No	C	D

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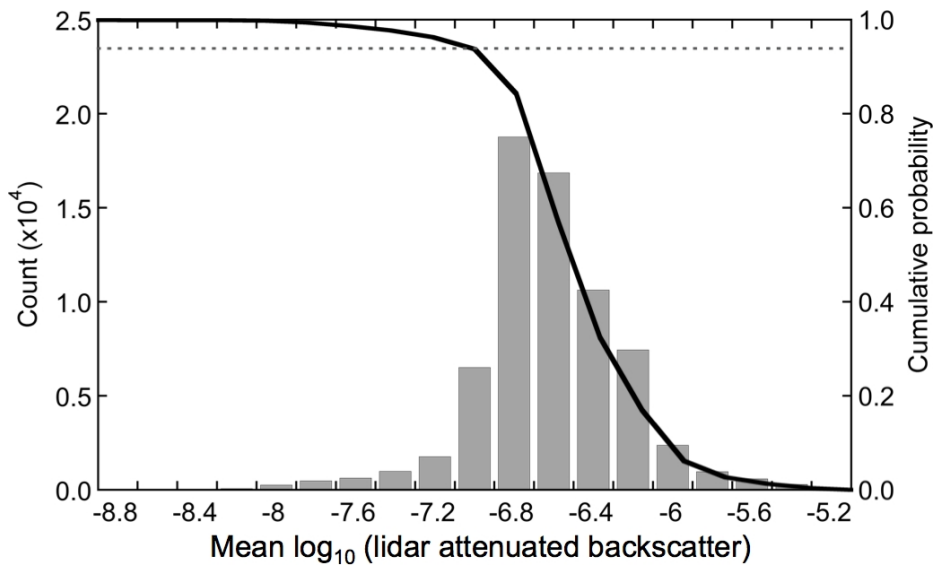
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2 Figure 1. (a) Plot of calibrated solar background light in lidar measurements versus cloud  
 3 optical depth at 523-nm wavelength and solar zenith angle (SZA) of 45° for cloud effective  
 4 radius of 4, 8 and 16 μm. (b) Vertical profiles of logarithm (with base 10) lidar attenuated  
 5 backscatter signals measured on 15th June 2007 at the ARM Oklahoma site at 19 UTC for  
 6 optically thick clouds, and at 23.5 UTC for optically thin clouds.

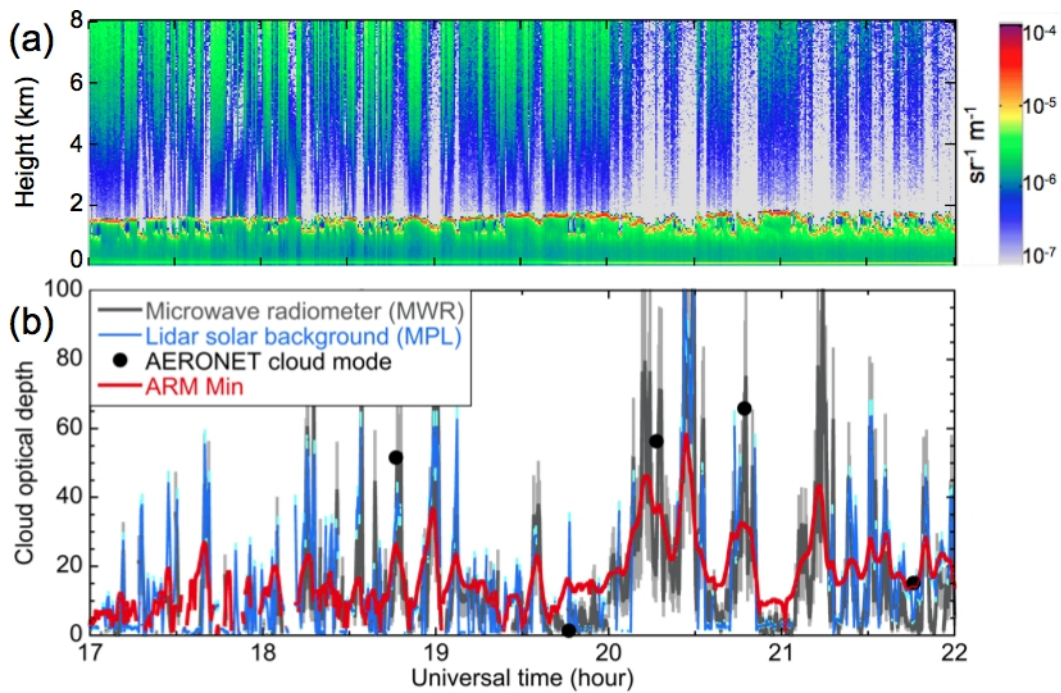
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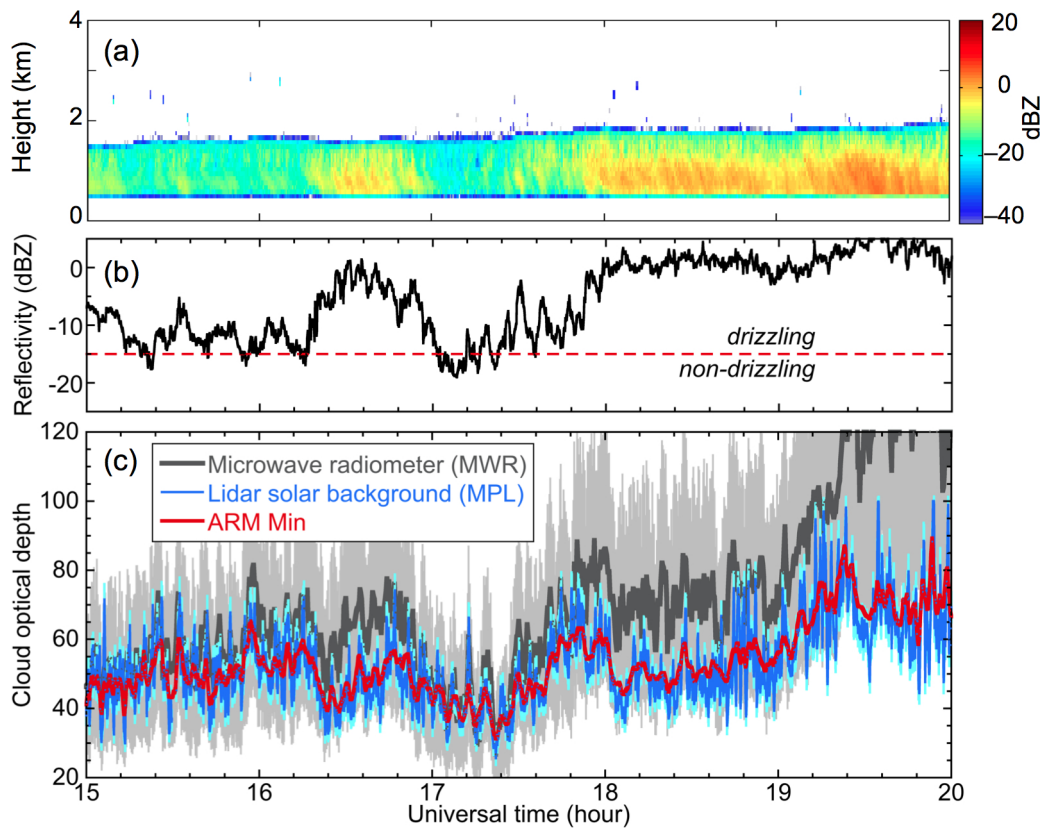
2 Figure 2. Histogram of  $\beta_{ct,1km}$  (the mean logarithm (with base 10) lidar backscatter from the  
 3 cloud top to 1 km above) and the corresponding cumulative probability (solid line) accounted  
 4 from the larger end of  $\beta_{ct,1km}$  for clear sky at the ARM Oklahoma site in 2005. The dashed  
 5 line represents the 94% cumulative probability.

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 2 Figure 3. (a) Attenuated backscatter signals from micropulse lidar on 19 April 2005. (b) Time  
 3 series of cloud optical depth retrieved from microwave radiometer (MWR), lidar solar  
 4 background signals from micropulse lidar (MPL), AERONET cloud-mode observations, and  
 5 ARM Archive Min retrievals (using narrowband irradiance measurements). MWR-based  
 6 retrievals (grey lines) are based on an assumed cloud effective radius of  $8 \mu\text{m}$ ; grey error bars  
 7 denote lower and upper limits for MWR values, respectively corresponding to a change in  
 8 droplet effective radius from  $6$  to  $14 \mu\text{m}$ .

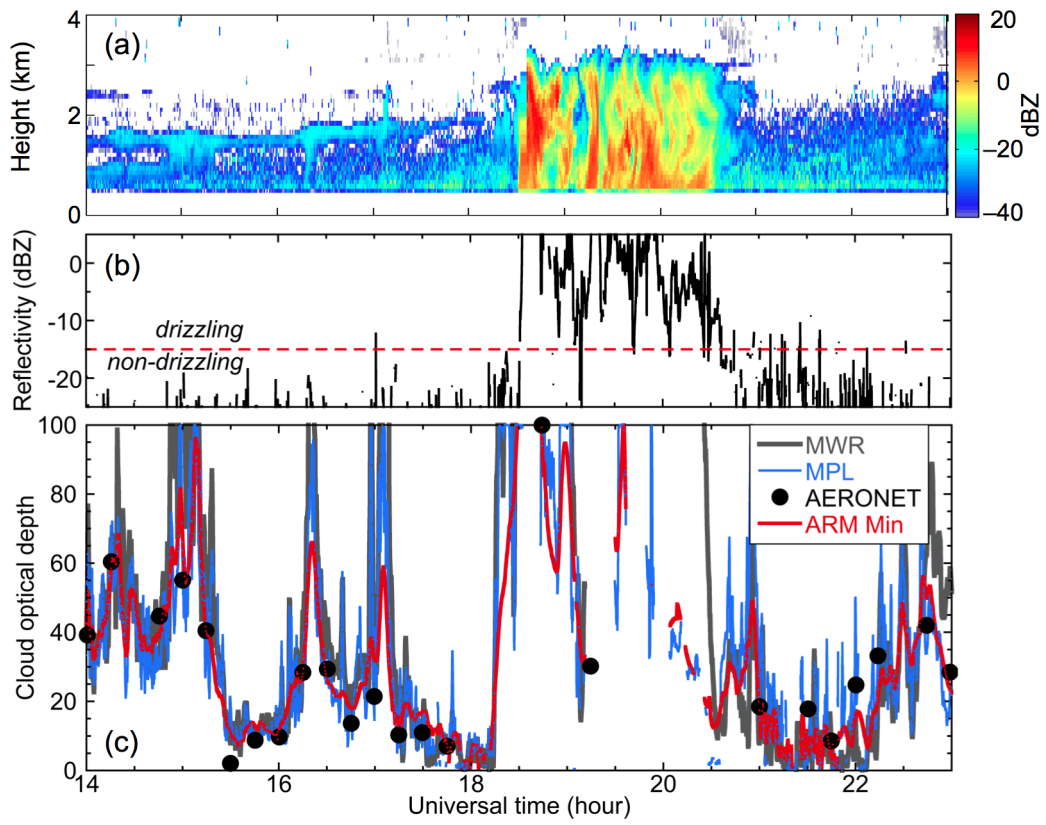
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2 Figure 4. (a) Time series of radar reflectivity profiles on 10 April 2007. (b) Cloud-base  
 3 reflectivity indicating the presence of drizzle. (c) Time series of cloud optical depth retrieved  
 4 from microwave radiometer (MWR), lidar solar background signals from micropulse lidar  
 5 (MPL) and ARM Archive Min retrievals. Grey error bars denote lower and upper limits for  
 6 MWR values, respectively corresponding to a change in droplet effective radius from 6 to 14  
 7  $\mu\text{m}$ . Note that AERONET cloud-mode observations were unavailable for this day

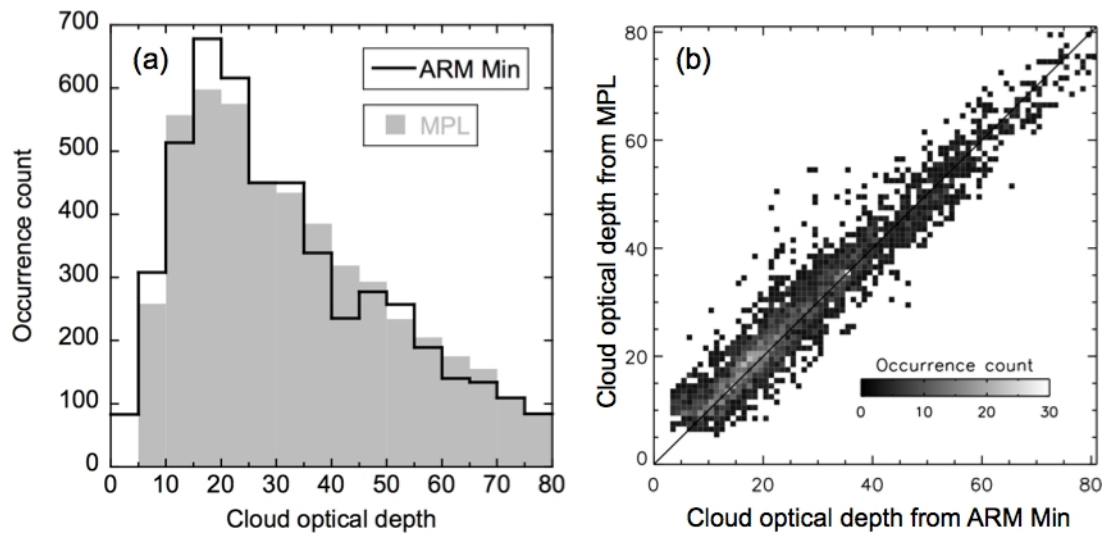
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2 Figure 5. Same as Fig. 4, but for 2 May 2007. Heavy precipitation occurs at 18.5–20.5 UTC.

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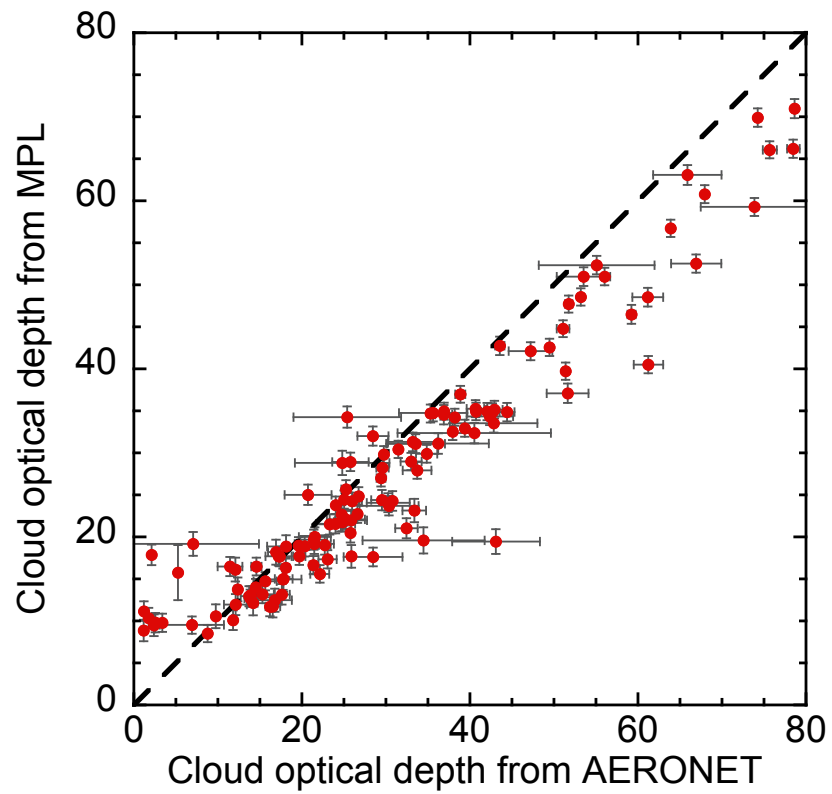


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2 Figure 6. (a) Histograms of occurrence count and (b) a scatter plot for intercomparison of  
 3 cloud optical depths retrieved from solar background signals received by micropulse lidar  
 4 (MPL) and those from the ARM Min product. Colours represent the number of occurrence  
 5 count, and the black solid line represent the 1:1 line.

6



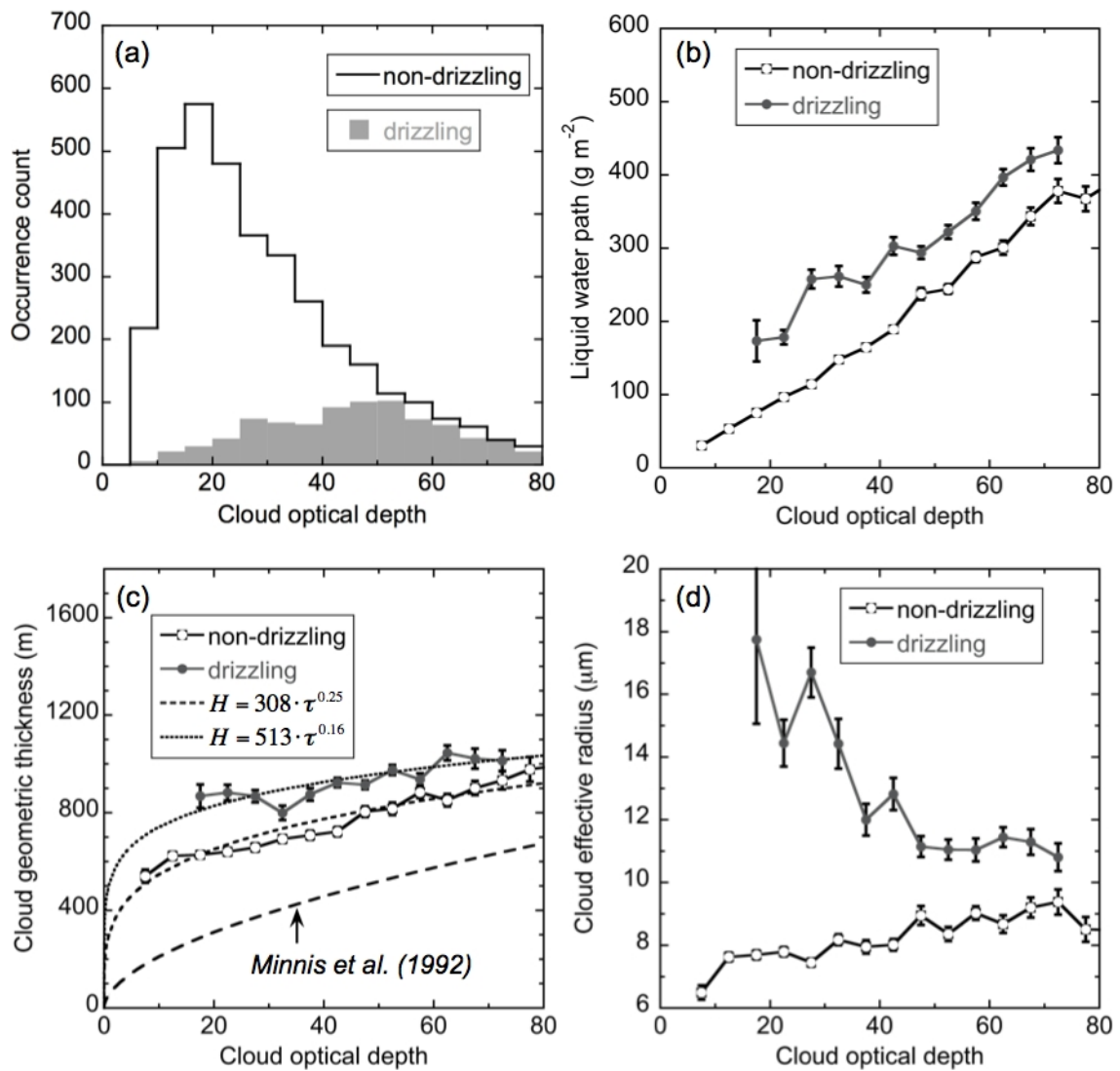


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2 Figure 7. Scatterplot for intercomparison of cloud optical depths retrieved from solar  
3 background signals received by micropulse lidar (MPL) and those from the AERONET  
4 cloud-mode product. The error bars represent one standard deviation of retrievals, while the  
5 black dashed line represents the 1:1 line.

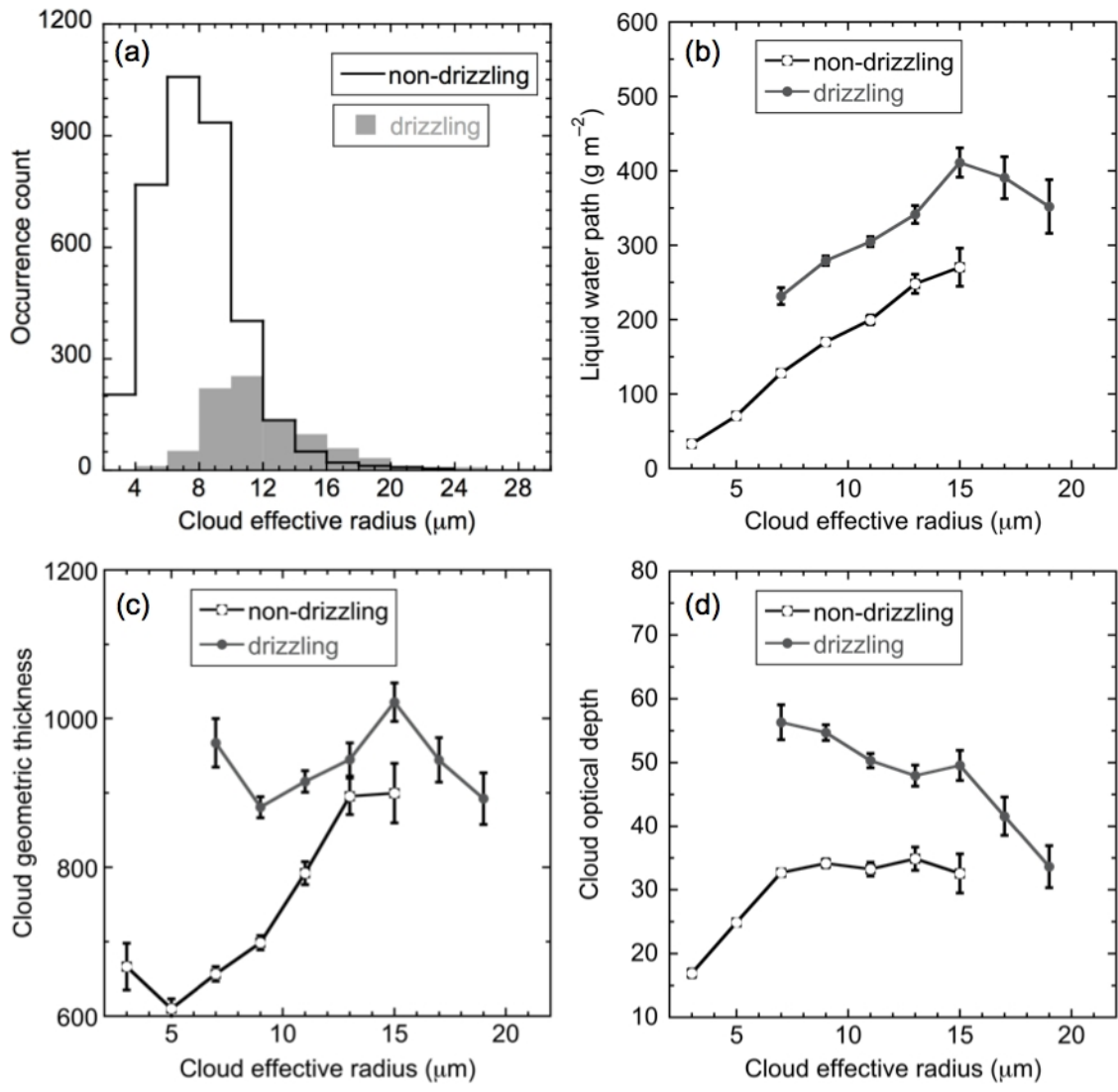
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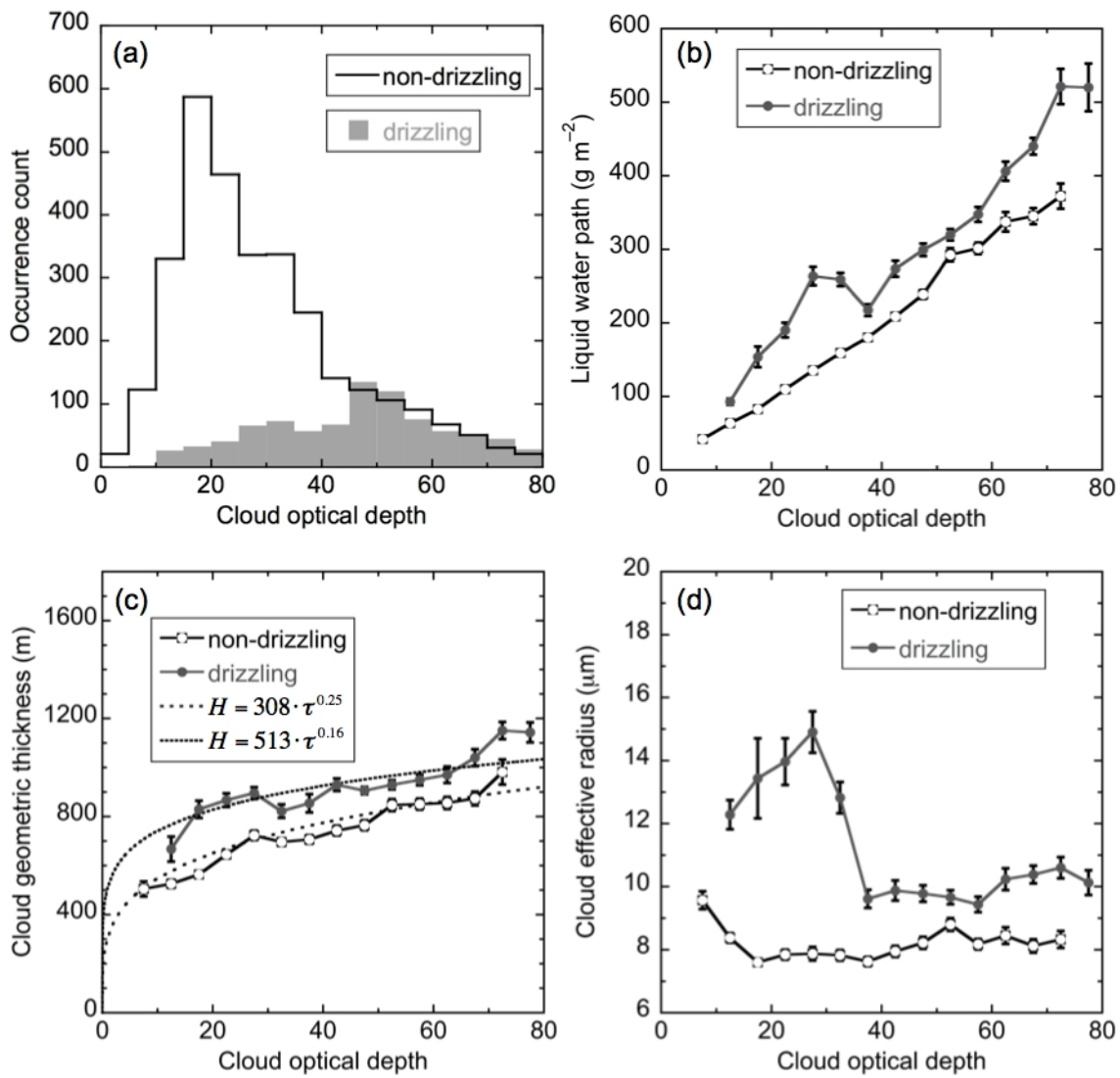
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Figure 8. (a) Occurrence histogram of cloud optical depth ( $\tau$ ); plots of (b) liquid water path, (c) geometric thickness ( $H$  in m) of cloud layer and (d) cloud effective radius versus optical depth for low-level stratiform clouds, using 1-min averaged retrievals at the ARM Oklahoma site during 2005–2007. A cloud-base radar reflectivity threshold of  $-15$  dBZ is used for drizzle classification: a cloud is drizzling if its cloud-base reflectivity exceeds the threshold, otherwise, non-drizzling. Error bars represent one standard error. Three power-law relationships are co-plotted in (c); dotted lines are based on ARM data, while the dashed line is adapted from the satellite-based finding reported in Minnis et al. (1992). (b)–(d) omit bins of cloud optical depth with a sample size smaller than 25.



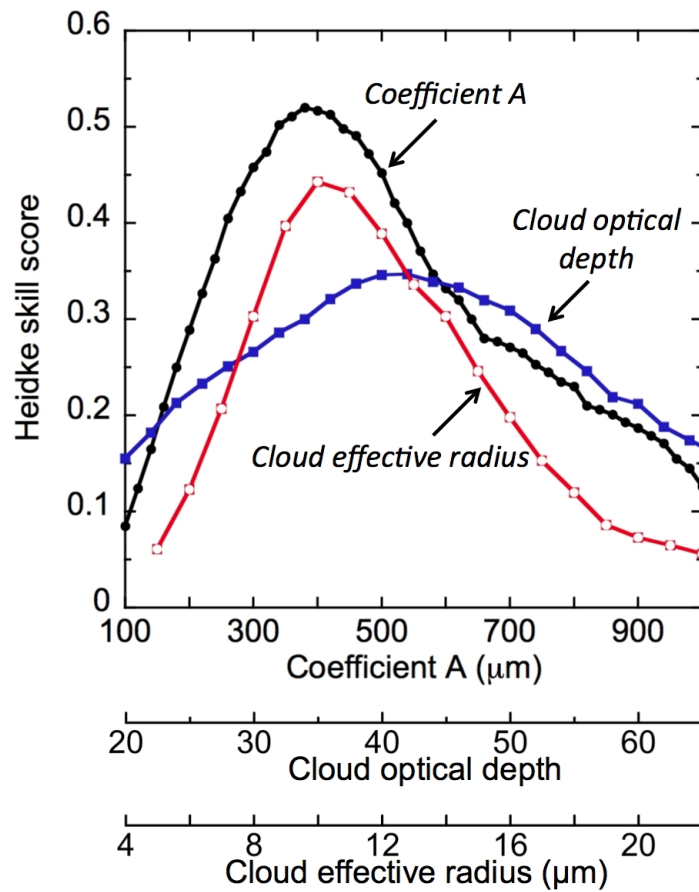
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Figure 9. Same as Fig. 8, but with plots of cloud properties versus cloud effective radius. (b)–(d) omit bins of cloud effective radius with a sample size smaller than 25.



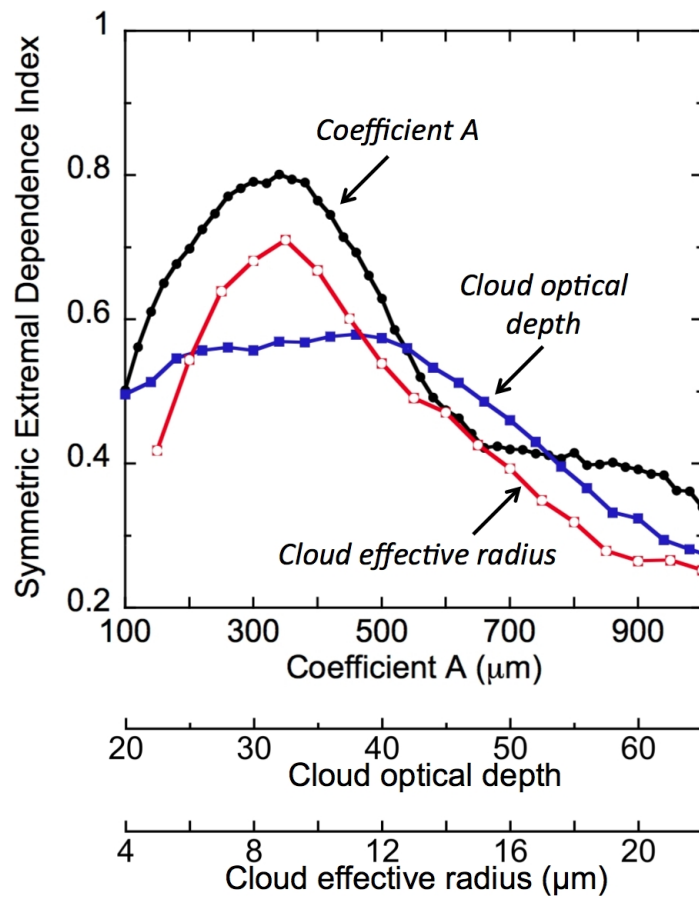
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2 Figure 10. Same as Fig. 8 but using cloud optical depths from the ARM Archive Min Product.  
 3 Note that the power-law relationships between cloud geometric thickness and optical depth by  
 4 dashed and dotted lines in (c) are derived from data in Fig. 8 for visual comparisons.



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Figure 11. Heidke skill scores for three drizzle delineation methods. The first (red) uses cloud effective radius as delineation threshold, while the second (blue) uses cloud optical depth instead. The third (black) uses a dynamic threshold as a function of both cloud optical depth and effective radius with a coefficient  $A$  (see Eq. (7) in text for details).



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2 Figure 12. Same as Fig. 11, but using Symmetric Extremal Dependence Index to optimise  
 3 thresholds for drizzle delineation.

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