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Trajectory model simulations of ozone (O₃) and carbon monoxide (CO) transport in the lower stratosphere (LS)

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12 Abstract

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14 A domain-filling, forward trajectory model originally developed for simulating 15 stratospheric water vapor is used to simulate ozone (O_3) and carbon monoxide (CO) in 16 the lower stratosphere (LS). Trajectories are initialized in the upper troposphere, and the 17 circulation is based on reanalysis wind fields. In addition, chemical production and loss 18 rates along trajectories are included using calculations from the Whole Atmosphere 19 Community Climate Model (WACCM). The trajectory model results show good overall 20 agreement with satellite observations from the Aura Microwave Limb Sounder (MLS) 21 and the Atmospheric Chemistry Experiment Fourier Transform Spectrometer (ACE-FTS) 22 in terms of spatial structure and seasonal variability. The trajectory model results also 23 agree well with the Eulerian WACCM simulations. Analysis of the simulated tracers 24 shows that seasonal variations in tropical upwelling exerts strong influence on O₃ and CO 25 in the tropical lower stratosphere, and the coupled seasonal cycles provide a useful test of 26 the transport simulations. Interannual variations in the tracers are also closely coupled to 27 changes in upwelling, and the trajectory model can accurately capture and explain 28 observed changes during 2005-2011. This demonstrates the importance of variability in 29 tropical upwelling in forcing chemical changes in the tropical lower stratosphere.

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31 1. Introduction

32 The influx of water vapor (H_2O) to the stratosphere is largely determined by the large-33 scale troposphere-to-stratosphere transport in the tropics, during which air is dehydrated 34 across the cold tropical tropopause (e.g., Fueglistaler et al., 2009 and references therein). 35 Observations such as the entry mixing ratios (Dessler, 1998; Dessler et al., 2013), the 36 coherent relations between water vapor and temperature (Mote et al., 1996), and the 37 extensive cirrus clouds near the tropopause (e.g., Winker and Trepte, 1998; Wang and 38 Dessler, 2012) all support this understanding. Back trajectory models have provided 39 detailed simulations of stratospheric H₂O (e.g. Fueglistaler et al., 2005). More recently, a 40 newly designed domain-filling forward trajectory model driven by reanalysis wind and 41 temperature has demonstrated success at simulating the transport of H₂O in the 42 stratosphere (Schoeberl and Dessler, 2011; Schoeberl et al., 2012; Schoeberl et al., 2013). 43 In this trajectory model, winds determine the pathways of parcels and temperature 44 determines the H₂O content through an idealized saturation calculation. This simple 45 advection-condensation strategy yields reasonable results for H₂O in the stratosphere, 46 although the detailed results depend on the wind and temperature fields utilized and 47 assumptions regarding supersaturation (Liu et al., 2010; Schoeberl et al., 2012).

48 Besides H₂O, ozone (O₃) and carbon monoxide (CO) are also important trace 49 gases in the stratosphere linked to circulation, transport and climate forcing (for O_3). 50 Ozone abundances in the stratosphere cover a wide dynamic range (~10's of ppbv to a 51 few ppmv), and are influenced by a variety of chemical and dynamical processes, 52 including photochemical production and loss and large- and small-scale transport (for 53 example, deep convective lofting of boundary layer low O_3 air to the upper troposphere, 54 e.g. Folkins et al., 2002). CO behaves as a tropospheric source gas, originating from 55 natural and anthropogenic emissions, including combustion processes near the surface 56 and oxidation of methane and other hydrocarbons within the troposphere. The main sink 57 of CO is oxidation by the hydroxyl radical (OH) (Logan et al., 1981). The CO 58 photochemical lifetime of 2-3 months makes it a useful tracer for transport studies in the 59 troposphere and lower stratosphere (e.g., Bowman, 2006; Schoeberl et al., 2006, 2008).

The purpose of this work is to simulate O_3 and CO in the lower stratosphere (LS) using the domain-filling, forward trajectory model used previously for stratospheric H₂O. Trajectory modeling of O_3 and CO can provide useful tests for simplified understanding of transport and chemical processes in the stratosphere, and provide complementary information to the H₂O simulations (which are primarily constrained by tropopause temperatures). In addition to testing circulation and transport within the trajectory model, the O_3 and CO simulations can provide understanding of mechanisms leading to observed 67 chemical behavior, including transport history and pathways, seasonal and interannual 68 variations, and relation to the stratospheric age-of-air (Waugh and Hall, 2002). Ozone 69 and CO are complementary tracers representing primarily stratospheric and tropospheric 70 sources; both tracers exhibit strong gradients across the tropopause, and they are often 71 combined using tracer correlations to understand transport and chemical behavior in the 72 upper troposphere and lower stratosphere (UTLS) (e.g., Fischer et al., 2000). Furthermore, 73 O_3 and CO exhibit relatively large out-of-phase seasonal cycles in the tropical lower 74 stratosphere (Randel et al., 2007), and these coupled variations provide a sensitive test of 75 the trajectory model simulations in this region.

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77 2. Model Description and Data Used

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79 2.1 Trajectory Model

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81 The trajectory model used here follows the details described in Schoeberl and Dessler 82 (2011) and Schoeberl et al. (2012), with trajectories calculated using the Bowman 83 trajectory code (Bowman, 1993; Bowman and Carrie, 2002; Bowman et al., 2013). 84 Because of the overly dispersive behavior of kinematic trajectories (e.g., Schoeberl et al., 85 2003; Liu et al., 2010; Ploeger et al., 2010; Schoeberl and Dessler, 2011), we perform 86 diabatic trajectories using isentropic coordinates, in which the vertical velocity is the 87 potential temperature tendency converted from the diabatic heating rates via the 88 thermodynamic equation (Andrews et al., 1987).

89 The methodology for the trajectory simulations of O_3 and CO follows Schoeberl 90 and Dessler (2011) and Schoeberl et al. (2012), wherein the parcels are initialized at the 91 370-K isentrope, below the tropical tropopause, using climatological O_3 and CO from the 92 MLS (monthly means averaged over 2005-2011) to provide approximate entry level 93 values in the upper troposphere. The 370 K level is chosen as the initialization level 94 because it is above the level of zero net heating rates and parcels there tend to ascend to 95 the stratosphere (see Discussion). To account for chemical changes along the trajectories, 96 we use chemical production and loss rates output from WACCM. Specifically, the O_3 and 97 CO concentration carried by each parcel is modified from the previous time step using98 the production and loss frequencies calculated from WACCM.

99 For each day 1350 parcels are initiated on an equal area grid covering 40° N-S 100 latitude and advected forward in time by reanalysis winds. At the end of each day, any 101 parcels that have descended below the 250-hPa level are removed since in most cases 102 they have re-entered the troposphere. The upper boundary is chosen to be the ~2200-K 103 isentrope (~1 hPa or ~50 km) to cover the entire stratosphere. Parcels are initialized and 104 added to the ensemble consecutively each day and the combined set of parcels is then run 105 forward. This process is repeated over the entire integration period. As more and more 106 parcels are injected into the model, the stratospheric domain is filled up with parcels – 107 this is the concept of domain-filling used in our model. The trajectory model simulations 108 are started in January 1990 and integrated to the end of 2011; after 3-4 years the system 109 reaches steady state with approximately 1 million parcels in the domain. We focus on 110 analyzing the model results during 2005-2011, to overlap the MLS and ACE-FTS 111 observations.

112 Upwelling across the tropical tropopause in the trajectory model is determined by 113 the reanalysis total diabatic heating rates (Q_{tot}, hereinafter Q), which include heating due 114 to long-wave and short-wave radiation, moist physics, friction, gravity wave drags, etc. 115 As shown below, our simulations of O₃ and CO are sensitive to the imposed upwelling. 116 There is substantial uncertainty in the detailed magnitude and spatial structure of Q, as 117 seen in the differences among separate reanalysis results (Schoeberl et al. 2012; Randel 118 and Jensen, 2013; Wright and Fueglistaler, 2013). Figure 1 illustrates the differences in Q in the tropics (15° N-S) based on several reanalysis data sets, highlighting large 119 120 differences in the UTLS. Given this uncertainty, we tested the sensitivity of our 121 calculations to variations in the heating rates by comparing results based on the NASA 122 Modern Era Retrospective-Analysis for Research and Applications (MERRA, Rienecker 123 et al., 2011) and the ECMWF ERA Interim reanalysis (ERAi, Dee et al., 2011). Below 124 we highlight the sensitivity of the resulting O₃ and CO simulations to these imposed 125 circulations.

126 Changes in chemical concentrations in the trajectory model are calculated using 127 the chemical continuity equation (e.g., *Dessler*, 2000):

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$$[\chi]_{current} = [\chi]_{previous} + (P_{\chi} - LF_{\chi} \cdot [\chi]_{previous}) \cdot \Delta t$$
(1)

Here, χ represents either O₃ or CO. Current chemical concentrations $[\chi]_{current}$ in volume 129 mixing ratio (VMR) are determined by concentrations in previous time step $[\chi]_{previous}$ 130 131 and the net change, derived from the production minus loss occurring in each time step. The production rate P_{χ} in VMR per unit time is obtained from WACCM. The loss rate 132 $(LF_{\chi} \cdot [\chi]_{previous})$ in VMR per unit time is calculated as a product of loss frequency LF_{χ} 133 134 (per unit time) times the chemical concentration (VMR), representing a linear chemical 135 loss. The loss frequencies are estimated from WACCM by dividing the model loss rate L_{χ} by the chemical concentration $[\chi]$, i.e., $LF_{\chi} = L_{\chi}/[\chi]$. In our simulation P_{χ} and 136 LF_{χ} are calculated from WACCM as a function of latitude, altitude, and climatological 137 138 months. The time step for calculating trajectories is 45-min while the time step for 139 calculating chemical changes is 1-day.

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2.2 WACCM Chemical Production and Loss Rates

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143 The chemical production and loss rates from the Whole Atmosphere Community Climate 144 Model (WACCM4) (Garcia et al., 2007) are applied in our model to represent chemical 145 processes. The vertical domain of WACCM4 extends from the surface to the lower thermosphere, with horizontal resolution of $2.5^{\circ} \times 1.9^{\circ}$ in longitude and latitude and 88 146 147 levels up to ~150 km. In the UTLS the vertical resolution is 1.1–1.4 km. The WACCM4 148 simulation used here is run with specified dynamics (SD) fields (Lamarque et al., 2012), 149 in which the model is nudged by the MERRA meteorological fields (Rienecker et al., 2011) from January 2004 to December 2011 (hereafter we only use WACCM to 150 151 represent SD-WACCM4).

152 Figure 2 shows the annual zonal mean ratio of net tendency (production rate 153 minus loss rate) divided by the loss rate for O_x and CO. These show the mean differences 154 of production to loss, with positive numbers indicating net chemical increase and 155 negative numbers indicating net chemical decrease; values close to zero imply balanced 156 production vs. loss. On the background are the respective photochemical lifetimes evaluated from the loss rates $(\tau_L = [\chi]/L_{\chi})$. Both O_x and CO exhibit relatively long 157

158 lifetimes (> 3 months) in the lower-middle stratosphere, and hence transport will have a 159 dominant influence on their distributions. Above ~30 km the lifetimes are shorter, and 160 the photochemical behavior will determine chemical structure.

Here O_x is the odd oxygen – the sum of O_3 and O. We use the production and loss rates of O_x because (1) the time step of trajectory model is 1 day, which is much longer than the lifetime of O_3 in the middle and upper stratosphere, and (2) the abundances of O_3 only change on time scales comparable to or longer than the lifetime of O_x (see e.g. Dessler, 2000, chapter 3). Therefore, it is reasonable to use production and loss rate of O_x instead of O_3 for our purpose of simulation. Because O_3 abundance is much greater than O, i.e., $O_x \approx O_3$, throughout the rest of the paper we use O_3 and O_x interchangeably.

 O_x production in the stratosphere is almost entirely due to the photolysis of O_2 . 168 P_{O_x} increases with altitude over most of the stratosphere because the photolysis rate 169 170 (proportional to solar radiation) increases faster with altitude than $[O_2]$ decreases, so the 171 net effect is increasing. Figure 2a shows that from 150 to ~ 10 hPa O_x production exceeds loss in the tropics, whereas a net chemical decrease of O₃ occurs in mid- to high latitudes 172 173 (transport of O_3 from the tropics to higher latitudes closes the budget). Between about 10 174 and 2 hPa the daily production and loss of O_x become comparable and the system is in 175 diurnal steady state. The system gradually reaches photochemical steady states above ~ 2 176 hPa, where the lifetime of O_x is less than a day and the instantaneous production and loss can be treated equal throughout the day. Here we simply set the $[O_x]$ to be (P_{Ox}/LF_{Ox}) in 177 178 our trajectory calculations.

In the UTLS region, the main source of CO is from the troposphere (from direct emissions and photochemical production), and this source is accounted for by initializing trajectories with observed values of CO. There is a net chemical decrease (Fig. 2b) in the UTLS, where CO is predominantly removed by oxidation with OH. CO experiences a net increase in the middle stratosphere (above \sim 24 km) due to oxidation of methane (CH₄), although this has little influence on the results shown here.

185 In our simulation P_{χ} and LF_{χ} are calculated from WACCM as a function of 186 latitude, altitude, and climatological months averaged over 2005-2011, so that a constant 187 annual cycle is applied throughout the model integration. As discussed in Sect. 5, this calculation will not accurately handle the situation where chemical losses are linked tometerological behavior, such as for ozone losses in the polar winter stratosphere.

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191 **2.3 MLS Observations of O₃ and CO**

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193 We will compare the model to observations of O₃ and CO from the Aura Microwave 194 Limb Sounder (MLS) (Waters et al., 2006). MLS climatology (constant annual cycle) of 195 O₃ and CO averaged from August 2004 to December 2012 is used to set the initial 196 abundances when parcels are initialized at the 370 K level, and the observations at higher 197 levels are used to evaluate the trajectory model results. We use the MLS version 3.3 198 (v3.3) Level 2 products, described in the data quality and description document 199 (http://mls.jpl.nasa.gov/data/v3-3 data quality document.pdf). O₃ profiles are available 200 at 12 levels per decade from 261 to 0.02 hPa and CO profiles are available between 215 201 and 0.0046 hPa at 6 levels per decade. The vertical resolution of O₃ in the UTLS is 202 approximately 2.5-3 km while for CO it is \sim 4.5-5 km. The detailed validation for these data sets can be found in Froidevaux et al., (2008), Pumphrey et al., (2007), and Livesey 203 204 et al., (2008).

For the comparisons shown here, we applied the MLS O_3 and CO averaging kernels to both our trajectory simulations and WACCM model outputs when comparing with the MLS observations. We have followed the detailed instructions for applying averaging kernels as described in <u>http://mls.jpl.nasa.gov/data/v3-</u> 3 data quality document.pdf.

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211 **2.4 Other Verifying Datasets**

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Besides using chemical production and loss from the WACCM, we also compare O_3 and CO modeled by WACCM to the trajectory model simulations, which serves as a sanity check of applying the imposed WACCM chemistry, and also as a simple comparison of Lagrangian vs. Eulerian model results. We also compare the trajectory modeling to the CO measurements from the Atmospheric Chemistry Experiment Fourier Transform Spectrometer (ACE-FTS) (Bernath et al., 2005), which shows some systematic 219 differences from MLS retrievals in the stratosphere (Clerbaux et al., 2008; Park et al.,

220 2013). A detailed description of the ACE CO observations can be found in Park et al.,221 (2013).

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223 3. Trajectory Modeling Results

3.1 O₃ Results

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226 Figure 3 shows the zonal mean cross section of O₃ during December-February (DJF) 227 from the trajectory model driven by ERAi reanalysis (denoted as "TRAJ ERAi"), 228 compared to both the MLS observations and the WACCM results. Because O₃ above 10 229 hPa is in photochemical steady-state (Sect. 2) we focus on O₃ below 10 hPa. Overall the 230 trajectory simulations agree with results from both MLS and WACCM. In the lower 231 stratosphere transport plays an important role in re-distributing chemical species, 232 resulting in contours of O_3 approximately following the isentropes. The enhanced O_3 233 production due to photolysis at 30 km (~10 hPa) shifts from south during DJF towards north during JJA (not shown), following the seasonal variations of photolysis rates. 234

235 Vertical profiles of O₃ averaged over the deep tropics (18° N-S) from 2005 to 236 2011 are shown in Fig. 4a (note that the x-axis is in log scale). The trajectory model 237 driven by MERRA (denoted as "TRAJ MER") shows reasonable agreement with MLS 238 data (and WACCM) in the lower stratosphere, while the results based on ERAi (denoted 239 as "TRAJ ERAi") show relatively smaller O₃ values. Above 24 km where photochemical 240 processes dominate, different trajectory runs yield similar results and they both agree 241 with MLS and WACCM data. Note that the MERRA and ERAi simulations use identical 242 O_3 initial values at 370 K, so that the differences in Fig. 4a are primarily a result of 243 differences in upward circulation (Fig. 1). The mean differences in ozone in the lower 244 stratosphere can be explained as a result of the different heating rates (vertical 245 circulations) imposed. The ERAi heating rates are higher than MERRA up to 20 km (Fig. 246 1). Due to positive vertical gradient in O_3 the stronger circulation moves air with lower 247 O₃ upward, creating a lower relative concentration compared to MERRA.

248 Monthly time series of O_3 at 100 hPa and 68 hPa averaged over the deep tropics 249 (18° N-S) are shown in Fig. 4b. There is a strong annual cycle in ozone at these levels related to the seasonal variations in tropical upwelling (Randel et al., 2007; Abalos et al., 2012, 2013a), and this behavior is reproduced by the trajectory model, showing reasonable agreement in amplitude with the MLS observations and WACCM results (in spite of the differences in mean values). There are somewhat larger differences in annual cycle amplitude at 68 hPa, with the ERAi trajectory results showing better agreement with MLS.

256 The simulated and observed latitudinal structure of zonal mean O_3 in the lower 257 stratosphere (68 hPa) throughout the seasonal cycle is shown in Fig. 5. The overall 258 variations are reasonably well simulated by the trajectory model, although low biases are 259 found compared to MLS over middle-to-high latitudes in both hemispheres (and 260 WACCM is systematically higher over the globe). The development of the Antarctic 261 ozone hole is evident in the very low ozone values polewards of 60° S in October, which 262 is simulated in the trajectory model based on the strong chemical ozone losses in this 263 region derived from WACCM.

Figure 6 compares the horizontal structure of boreal summer (JJA) O_3 at 83 hPa from MLS data and the trajectory results driven by MERRA. The trajectory results show a reasonable simulation of the spatial patterns compared to MLS (and WACCM; not shown), with a clear minimum inside the Asian monsoon anticyclone linked to upward transport of ozone-poor air from lower levels (Park et al., 2009). There is also relatively low O_3 centered near 15° S linked to the slow ascending air from the troposphere in this region.

271 The trajectory model is also able to capture the spatial behavior of polar ozone. 272 Figure 7 shows a comparison of high latitude O₃ in the Northern Hemisphere (NH) during 273 winter (DJF) and in the Southern Hemisphere (SH) during spring (September) between 274 MLS and trajectory modeling driven by both MERRA and ERAi winds. During NH 275 winter, O₃ rich air (> 2.2 ppmv) occurs within the polar vortex (denoted with the 24 PVU 276 isopleth in Fig. 7a-c), and the trajectory model captures the observed isolation from 277 middle latitudes, although differences in magnitude exist. During SH springtime, the 278 Antarctic ozone hole (denoted with the 195 K isotherm in Fig. 7d-f) is reasonably well 279 reproduced in the trajectory model based on imposed WACCM chemistry. The trajectory model also captures the spatial structure of the zonal wave ozone maximum near 50° S 280

(the so-called "ozone croissant"), linked to the descending branch of the BD circulation.
The weaker extra-vortex high in ozone in the trajectory model may be related to the
weaker overall circulation in MERRA compared to observations (Schoeberl et al., 2012,
2013). Noted that due to the climatological loss frequency adopted from WACCM, this
model cannot account for interactions of circulation and chemistry, which affects O₃
mostly in winter polar regions (see Discussion).

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288 **3.2 CO Results**

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290 Figure 8 shows that zonal mean cross sections of CO from ACE-FTS, WACCM, and the 291 trajectory model agree well in the lower stratosphere. CO has a maximum in the tropical 292 upper troposphere, and decreases with altitude due to OH oxidation to a minimum near 293 22 km. CO increases above this altitude due to production from methane (CH₄) oxidation. 294 The ACE-FTS observations show high CO mixing ratios in the polar middle and upper 295 stratosphere regions, resulting from the downward transport of CO from the mesosphere 296 (from photodisassociation of CO_2); this behavior is also seen (to a weaker degree) in 297 WACCM, but is not simulated in the trajectory model, which does not include 298 mesospheric processes.

299 Figure 9a shows the CO vertical profiles and time series averaged in the deep 300 tropics (15° N-S). Due to the differences of MLS and ACE-FTS retrievals in the 301 stratosphere, here we also included ACE CO for comparison. The vertical profiles in Fig. 302 9a show broad-scale agreements, although there are differences among the trajectory 303 models (with ERAi driven results larger than those driven by MERRA) and also between 304 the ACE-FTS and MLS observations (all of the models agree better with the ACE 305 observations above 22 km). Time series of CO at 100 hPa (Fig. 9b) show a semi-annual 306 cycle linked to initialized variations in the upper troposphere (Liu et al., 2007), with 307 approximate agreement among the models and observations (with slightly larger values 308 in the MLS data). The variability changes to an annual cycle at 68 hPa, as a response to 309 variations in tropical upwelling. At 68 hPa the annual cycle is captured is a reasonable 310 manner in the models, with the ERAi results showing better agreement with MLS and 311 ACE-FTS data.

312 A further diagnostic to evaluate the model simulations is made by plotting monthly tropical (15° N-S) averages of O₃ vs. CO in the lower stratosphere (68 hPa), as 313 314 shown in Fig. 10. This includes the observations from MLS, together with trajectory 315 model simulations driven by both MERRA and ERAi, which shows the sensitivity to 316 different Q (see also Fig. 1). There is an overall anti-correlation between O_3 and CO in 317 Fig. 10, mainly representing the out-of-phase annual cycles seen in Figs. 4b and 9b. The 318 comparisons in Fig. 10 show that stronger upwelling in the ERAi simulation produces 319 slightly lower values of O₃ and higher values in CO (> 30 ppbv), in better agreement with 320 MLS data. Moreover, the slope of the CO-O₃ variations in the ERAi simulation 321 approximately matches the MLS result.

322 The DJF and JJA seasonal distributions of CO at 68 hPa from the ERAi trajectory 323 model are compared to MLS data in Fig. 11. In both seasons the trajectory model shows 324 spatial patterns consistent with MLS data. During DJF the patterns show a center of high 325 CO over central America and enhancements over South East Asia, extending to the 326 tropical western Pacific (largely attributable to fossil fuel emissions, Jiang et al., 2007). 327 The trajectory model also captures the well-known CO maximum linked to the Asian 328 monsoon anticyclone during JJA, which is substantially stronger at the 100-hPa level (e.g. 329 Randel and Park, 2006; Park et al., 2009; Randel et al., 2010).

Overall, the large-scale seasonal behavior of CO simulated by the trajectory model is in agreement with both observations (MLS and ACE-FTS) and Eulerian chemical model (WACCM), although the results are sensitive to the tropical upwelling speed (see Discussion).

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4. Interannual Variability of Tracers in the Tropical Lower Stratosphere (LS)

The coherent seasonal variations in O_3 and CO in the tropical LS demonstrate that transport processes have a large impact on the chemical concentrations in this region. The Eulerian-mean calculations of *Abalos et al.* (2012, 2013a) show that tropical upwelling is the main driver of the annual cycles in O_3 and CO above the tropical tropopause. Our Lagrangian trajectory model results (Figs. 4b and 9b) also show that the annual cycles of O_3 and CO above the tropopause (especially around 70 hPa) are strongly influenced by the tropical upwelling (Brewer-Dobson) circulation.

343 We further explore interannual variations in the chemical tracers and links to 344 changes in the upwelling circulation. Figure 12 shows the interannual anomalies (by 345 removing the annual cycle) in O₃ and CO concentrations in the tropical lower 346 stratosphere from MLS observations and from trajectory calculations, and in addition 347 anomalies in diabatic heating rates (upwelling) from reanalysis at 68 hPa. While there are 348 significant differences in time-mean diabatic heating rates between MERRA and ERAi 349 (Fig. 1), interannual changes in Q (Fig. 12c) show good agreement. Figure 12 shows that 350 interannual anomalies in O₃ and CO are strongly anti-correlated (due to oppositely signed 351 vertical gradient) and closely linked to interannual changes in diabatic heating. 352 Furthermore, Fig. 12 shows that trajectory calculations driven by both MERRA and 353 ERAi are able to simulate the observed interannual anomalies in O₃ and CO, in spite of significant differences for the background seasonal cycle (Fig. 4b and 9b). 354

355 Taking results from the MERRA run as an example, the close relationship 356 between anomalies in diabatic heating and O₃ is quantified in Fig. 13a, which shows strong anti-correlation with explained variance $r^2=0.73$ and slope $\Delta O_3/\Delta Q = -1.05\pm0.14$ 357 (ppmv)/(K/day). Similar strong correlation is found for CO and Q anomalies (Fig. 13b), 358 with explained variance $r^2=0.85$ and slope of 31.11 ± 2.86 (ppbv)/(K day⁻¹). The sign of 359 360 the slopes in Figs. 13a-b are opposite because of the different background vertical 361 gradients for O₃ and CO. These strong relationships between diabatic heating (Q) and 362 tracer anomalies highlight the dominant role of tropical upwelling in controlling species 363 with strong vertical gradients near the tropical tropopause.

Figure 12 also highlights strong anti-correlations between O₃ and CO anomalies, which is further demonstrated in Fig. 13c. Abalos et al. (2012) have shown that for the idealized case where upwelling dominates tracer transports, the ratio of tendencies for two tracers (χ_1, χ_2) is closely related to the ratio of the respective background gradients:

368 $\frac{\partial \overline{\chi}_1}{\partial t} / \frac{\partial \overline{\chi}_2}{\partial t} = \frac{\partial \overline{\chi}_1}{\partial z} / \frac{\partial \overline{\chi}_2}{\partial z} \sim constant$ (2)

369 Integrating this equation in time for monthly O₃ and CO anomalies gives the relationship:

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$$\Delta \overline{O_3} / \Delta \overline{CO} \approx \frac{\partial \overline{O_3}}{\partial z} / \frac{\partial \overline{CO}}{\partial z}$$
(3)

371 i.e. the ratio of monthly anomalies approximately follows the ratio of background vertical 372 gradients for the idealized situation where vertical transport is dominant. For the case of 373 CO and O_3 in the tropics near 68 hPa, the MERRA trajectory results yields a background 374 gradient ratio of ~ -18.34 (ppbv km⁻¹)/(ppmv km⁻¹) for (dCO/dz) / (dO₃/dz). A linear fit 375 of the observed CO/O₃ anomalies (Fig. 13c) gives a ratio of -22.45±3.40 (ppbv / ppmv), 376 which is close to the idealized result (slightly outside of the two-sigma uncertainty). This 377 approximate agreement with highly idealized theory provides further evidence for the 378 control of tropical lower stratosphere O₃ and CO by variations in upwelling.

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380 5. Summary and Discussion

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382 The results presented here demonstrate that the domain-filling, forward trajectory 383 model is useful for studying the transport of trace gases in the LS. The O_3 and CO 384 simulations are complementary to modeling H₂O (mainly controlled by tropopause 385 temperature) in that O₃ and CO rely on both initial conditions and chemical production 386 and loss rates, and are sensitive to transport. Initial conditions based on observations 387 provide entry values of chemical species into the lower stratosphere; after that the 388 chemical production and loss control the net changes of concentrations along the 389 trajectories. Because the MERRA negative heating rates at 150-130 hPa (Fig. 1) prevent 390 air ascending to the stratosphere, we chose a relatively high initialization level in the 391 upper troposphere (370 K) in this study. However, we could have initiated parcels at 392 lower altitude, such as 355 K when using ERAi circulation, to extend the results to the 393 upper troposphere (UT).

394 Trajectory modeled O_3 and CO in the tropical lower stratosphere largely depend 395 on the strength of upwelling (and to a lesser degree on the amount of mixing with 396 extratropics, Abalos et al., 2013a). Stronger upwelling is linked to faster transport, which 397 results in less time for chemical production (for O₃) or loss (for CO), leading to overall 398 lower values of O₃ (Fig. 4a) and higher values of CO (Fig. 9a). The comparisons of 399 MERRA and ERAi simulations that have different tropical upwelling rates (Q), e.g. Fig. 400 1, clearly demonstrate this sensitivity. The detailed differences in diabatic heating rates 401 among reanalyses have been discussed by Wright and Fueglistaler (2013), who highlight 402 differences in the corresponding long-wave radiative heating rates in the lower 403 stratosphere, which are influenced by both temperature and ozone. We also conducted a 404 sensitivity study of increasing MERRA diabatic heating rates (Q) by constant factors, and the best overall fit to the observations is 1.5 times the MERRA Q values, consistent withthe ERAi-based results.

407 Although better agreement with observations of CO in the tropical lower 408 stratosphere are found using ERAi data, there is reason to suspect that the ERAi diabatic 409 heating in this region may be too high. For example, Ploeger et al. (2012) performed a 410 radiative calculation showing that ERAi heating rates are ~40% too high, consistent with 411 Schoeberl et al. (2012), who show that trajectory modeled water vapor tape recorder 412 signal based on ERAi heating rates is \sim 30% too fast compared with MLS observations. In spite of the detailed differences, the trajectory modeled O₃ shows reasonable simulation 413 414 of the large-scale seasonal structure compared to both MLS and WACCM, including both 415 the tropics and the polar regions. The trajectory modeled CO in the tropical stratosphere 416 is more sensitive to the MERRA vs. ERAi differences, likely because of the shorter 417 photochemical lifetime of CO in the lower stratosphere compared to O_3 .

418 The annual cycles in O_3 and CO in the tropical LS are reproduced in the trajectory 419 model simulations, and the magnitude of variations provides a useful test of the imposed 420 circulation. The variability of O₃ and CO shows significant correlations with fluctuations 421 in diabatic heating, for both seasonal and interannual time scales. These close 422 relationships support the concept that tropical upwelling plays a key role in regulating 423 variability for chemical species with strong vertical gradients in the lower stratosphere 424 (and explains the observed compact relationships among interannual anomalies in 425 diabatic heating, O_3 and CO seen in Figs. 12-13). For the idealized situation where 426 upwelling dominates tracer transports, the tracer ratios can be expressed as ratios of the 427 background vertical gradients, and the observed O₃ and CO changes are in approximate 428 agreement with this expectation (Fig. 13c).

The discussion above linking seasonal or interannual changes in O_3 and CO with chemical changes along slower or faster upward trajectories is a Lagrangian perspective on transport (appropriate for our Lagrangian trajectory model). In contrast, the discussion linking O_3 and CO variations at particular pressure levels to varying circulations acting on background vertical gradients (Sect. 4) is an Eulerian perspective. These two perspectives are complementary and do not contradict each other; Abalos et al. (2013b) have recently shown the equivalence of Lagrangian and Eulerian transport calculations inthe tropical lower stratosphere, highlighting that each can provide useful information.

One limitation of this model is that it cannot account for interactions of circulation and chemistry, due to the climatological chemical production rates and loss frequencies adopted from WACCM (Section 2.2). This is most likely important in winter polar regions, where O_3 loss chemistry due to chlorine activation on polar stratospheric clouds is highly temperature dependent. Although we cannot simulate the detailed behavior in any particular (cold) year with climatological chemical ozone loss rates, our model does a reasonable job at simulating the time-average behavior of polar regions.

444 Our simulations with O₃ and CO have demonstrated the viability of the domainfilling forward trajectory model for simulating species with relatively simple chemistry. 445 446 and extension to other species would be straightforward. There are several potential 447 applications for such a trajectory model, including describing parcel histories that 448 characterize different transport pathways, and evaluating the importance of tropical-449 extratropical exchanges. For example, trajectories can allow tracing the sources of CO-450 rich air in the summertime Asian monsoon region, and quantifying the fate of the parcels 451 after breakup of the anticyclone. The model can also allow detailed comparisons of 452 transport based on different and new reanalysis data sets, or idealized studies of the 453 chemical responses to UTLS circulation in a changing climate.

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Figures 660



Fig. 1. Comparison of total diabatic heating rates averaged over the deep tropics (18° N-S) in 2000-2010 from four different reanalysis data sets: MER (MERRA), ERAi (ECMWF ERA interim), CFSR (NCEP CFSR) and JRA25.



Fig. 2. The ratio of chemical net tendency (production rate minus loss rate) to loss rate from WACCM for (a) O_x and (b) CO. Negative numbers are dashed to highlight the net chemical decrease and positive numbers indicate net chemical increase, while zero lines indicate comparable amount of production and loss. For reference, the respective O_x and CO lifetimes are contoured in color.



Fig. 3. Trajectory modeled O_3 driven by ERAi reanalysis wind (**c**, TRAJ_ERAi) in boreal winter (DJF), compared to both MLS observations (**a**) and WACCM output (**b**). Both WACCM and trajectory model results are weighted by the MLS averaging kernels. Dashed lines are potential temperature contours, demonstrating that in the lower LS where transport dominates, chemical distributions roughly follow the isentropes.



Fig. 4. Tropical (a) vertical profile and (b) time series (100 hPa, bottom panel; 68 hPa, upper panel) of MLS, WACCM, and trajectory modeled O_3 driven by MERRA wind and ERAi wind, averaged over the deep tropics (18° N-S) from 2005 to 2011. Both WACCM and trajectory model results are weighted by the MLS averaging kernels. In (a) the x-axis is in log scale to highlight the differences; the purple square indicates the initial O_3 values carried by parcels when they were first injected into the domain at 370 K; the vertical bars in red indicate the MLS vertical resolutions at each of the MLS retrieval pressure levels.



Fig. 5. Zonal mean of O_3 at 68-hPa during January (JAN), April (APR), July (JUL), and October (OCT) averaged in 2005-2011 from MLS, WACCM, and trajectory results driven by MERRA wind and ERAi wind. Both WACCM and trajectory model results are weighted by the MLS averaging kernels. In October (Antarctic spring time), the South Pole undergoes exceptional depletion of O_3 .



Fig. 6. Northern Hemisphere Summertime (JJA) tropical O_3 distributions at 83 hPa averaged from 2005 to 2011 between MLS and the MERRA driven trajectory simulations (weighted by the MLS averaging kernels). Horizontal wind vectors from the MERRA reanalysis are overlaid to emphasis the Asian summer monsoon anticyclone.





Fig. 7. Polar O_3 distributions shown in MLS (left column), trajectory results driven by MERRA (middle column), and trajectory results driven by ERAi (right column) during Northern Hemisphere winter (DJF, **a-c**) and Southern Hemisphere spring (September, SEP, **d-f**) at 68 hPa. Trajectory results are weighted by the MLS averaging kernels. The 24-PVU potential vortices (**a-c**) and the 195-K temperature (**d-f**) are overlaid in black dashed lines for both seasons, respectively.



Fig. 8. Zonal mean cross sections of CO from **(a)** ACE-FTS, **(b)** WACCM, and **(c)** trajectory model driven by MERRA reanalysis.





Fig. 9. Tropical **(a)** vertical profile and **(b)** time series (100 hPa, bottom panel; 68 hPa, upper panel) of MLS, WACCM, ACE, and trajectory modeled CO driven by MERRA wind and ERAi wind, averaged over 15° N-S from 2005 to 2011. In **(b)** the WACCM and trajectory model results are weighted by the MLS averaging kernels. In **(a)** the purple square indicates the initial CO values carried by parcels when they were first injected into the domain at 370 K, and vertical bars in red indicate the MLS vertical resolutions at each of the MLS retrieval pressure levels.



Fig. 10. Monthly variations of O_3 vs. CO in the tropical lower stratosphere (15° N-S, 68 hPa) from MLS and trajectory modeling driven by MERRA wind and ERAi wind. Trajectory results are weighted by the MLS averaging kernels.



Fig. 11. Comparison of CO at 68 hPa (~ 20 km) during DJF (top row, **a** and **b**) and JJA (bottom row, **c** and **d**) between MLS (left) and trajectory modeling driven by ERAi wind (right, weighted by the MLS averaging kernels). Horizontal wind vectors from ERAi are overlaid as reference.



Fig. 12. Interannual anomalies of (a) O_3 and (b) CO from MLS (red) and trajectory simulations driven by MERRA and ERAi in the tropical (15° N-S) lower stratosphere (68 hPa). (c) shows interannual anomalies of total diabatic heating rates from MERRA and ERAi, which serves in our model as the vertical velocity. The trajectory results are weighted by the MLS averaging kernels.



Figure 13. Scatter plots of the anomalies of (a) O_3 vs. Q and (b) CO vs. Q, and (c) CO vs. O_3 at 68 hPa. The dots are monthly variations of tracers from trajectory modeling driven by MERRA winds and diabatic heating rates; the black lines show the linear fit. The red line in (c) is the theoretically estimated tracer ratio estimated from the respective background gradients, using simplified relation in Eq. (3).