1 The impact of volcanic aerosol on the Northern Hemisphere stratospheric polar

2 vortex: mechanisms and sensitivity to forcing structure

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- 13 Abstract
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15 Observations and simple theoretical arguments suggest that the Northern Hemisphere (NH) 16 stratospheric polar vortex is stronger in winters following major volcanic eruptions. However, recent 17 studies show that climate models forced by prescribed volcanic aerosol fields fail to reproduce this 18 effect. We investigate the impact of volcanic aerosol forcing on stratospheric dynamics, including the 19 strength of the NH polar vortex, in ensemble simulations with the Max Planck Institute Earth System 20 Model. The model is forced by four different prescribed forcing sets representing the radiative 21 properties of stratospheric aerosol following the 1991 eruption of Mt. Pinatubo: two forcing sets are 22 based on observations, and are commonly used in climate model simulations, and two forcing sets are 23 constructed based on coupled aerosol-climate model simulations. For all forcings, we find that 24 temperature and zonal wind anomalies in the NH high latitudes are not directly impacted by anomalous 25 volcanic aerosol heating. Instead, high latitude effects result from enhancements in stratospheric 26 residual circulation, which in turn result, at least in part, from enhanced stratospheric wave activity. 27 High latitude effects are therefore much less robust than would be expected if they were the direct 28 result of aerosol heating. Both observation-based forcing sets result in insignificant changes in vortex 29 strength. For the model-based forcing sets, the vortex response is found to be sensitive to the structure 30 of the forcing, with one forcing set leading to significant strengthening of the polar vortex in rough 31 agreement with observation-based expectations. Differences in the dynamical response to the forcing 32 sets imply that reproducing the polar vortex responses to past eruptions, or predicting the response to 33 future eruptions, depends on accurate representation of the space-time structure of the volcanic 34 aerosol forcing.

36 1. Introduction

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38 The Northern Hemisphere (NH) stratospheric winter polar vortex, which shows considerable interannual 39 and intra-seasonal variability, has been observed to be stronger than normal in winters after major 40 volcanic eruptions (Kodera, 1995; Labitzke and van Loon, 1989). While this observation is based on a 41 relatively small sample (limited to the winters after the 1963 Agung, 1982 El Chichón and 1991 Pinatubo 42 eruptions), the theoretical argument to explain such a strengthening appears clear: namely, that heating 43 of the lower stratosphere through the absorption of radiation by volcanic sulfate aerosols enhances the 44 equator-to-pole temperature gradient in the lower stratosphere, which, through the thermal wind 45 equation, leads to stronger westerly winds (Robock, 2000 and references therein). Satellite observations clearly show a warming of the tropical lower stratosphere after volcanic eruptions (Labitzke and 46 47 McCormick, 1992), so changes in meridional temperature gradients and zonal winds are logical 48 consequences. The degree to which secondary feedback mechanisms -- such as changes in ozone or 49 upward propagating planetary waves (e.g., Graf et al., 2007; Stenchikov et al., 2006)-- also affect the 50 vortex strength is at present unclear.

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52 Post-volcanic strengthening of the NH polar vortex is an important step in a proposed mechanism which 53 explains observed changes in surface climate in post-eruption winters (Kodera, 1994; Perlwitz and Graf, 54 1995). Specifically, it is thought that through stratosphere-troposphere coupling (Baldwin and Dunkerton, 2001; Gerber et al., 2012) the volcanically-induced strong stratospheric polar vortex leads to 55 56 the observed positive anomalies in surface dynamical indexes such as the Northern Annular Mode 57 (NAM) or North Atlantic Oscillation (NAO) (Christiansen, 2008). Such dynamical changes lead to the 58 "winter warming" pattern of post-volcanic temperature anomalies (Robock and Mao, 1992), which is characterized by warmer temperatures over large regions of the NH continents during winter which 59 60 oppose the overall cooling impact of the volcanic aerosols on the surface. On the other hand, in a model 61 study, Stenchikov (2002) found that a positive phase of the NAM was also produced in an experiment in 62 which only the tropospheric impact of volcanic aerosols was included, implying that aerosol heating in 63 the lower tropical stratosphere is not necessary to force a positive NAM response. Whatever the 64 mechanisms, observations show that 11 out of 13 eruptions since 1870 were followed by positive 65 wintertime NAO values (Christiansen, 2008): the apparent robustness of this post-eruption dynamical 66 response should allow for enhanced skill in seasonal prediction for winters which follow volcanic 67 eruptions (e.g., Marshall et al., 2009).

69 While a number of early model simulations reported qualified success in simulating the atmospheric 70 dynamical response to volcanic eruptions (e.g., Graf et al., 1993; Kirchner et al., 1999; Rozanov, 2002; 71 Shindell et al., 2001), assessments of the multi-model ensembles of the Coupled Model Intercomparison 72 Projects (CMIP) 3 and 5 showed no significant winter warming response to prescribed volcanic forcing, 73 nor did they show significant anomalies in post-eruption dynamical quantities in the stratosphere or at 74 the surface (Driscoll et al., 2012; Stenchikov et al., 2006). It has been suggested that in order for a model 75 to successfully respond to volcanic forcing, it should include a reasonably well-resolved stratosphere 76 (Shindell, 2004; Stenchikov et al., 2006). However, analysis of the CMIP5 ensemble revealed no 77 appreciable systematic difference in post-eruption geopotential height anomalies to volcanic aerosol forcing between models with or without well-resolved stratospheres (Charlton-Perez et al., 2013). 78 79 80 Most model simulations which incorporate the impact of volcanic eruptions, such as the CMIP3 and 5

81 historical simulations, do so using prescribed volcanic aerosol fields, which have associated 82 uncertainties. In this study we investigate how the response of the atmosphere to volcanic aerosol 83 forcing depends on the prescribed aerosol forcing used in the simulation. We use one CMIP5 model, the 84 MPI-ESM, and focus on the first winter after Mt. Pinatubo, the period of strongest volcanic forcing 85 within the era of satellite observations. We run ensemble simulations using four different forcing data 86 sets: two observation-based forcing sets, based primarily on different versions of SAGE II aerosol 87 extinction measurements, and two model-constructed aerosol forcing sets. By assessing the model 88 response to these forcing sets, we investigate (1) the mechanisms linking volcanic aerosol heating in the 89 lower tropical stratosphere and NH winter vortex strength, (2) what stratospheric circulation responses 90 to volcanic aerosol forcing are robustly produced by the model and forcings, and (3) the sensitivity of 91 the vortex response to changes in the space-time structure of the volcanic forcing.

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93 2 Experiment: Observational basis, model and methods

In this section, the hypothesis that volcanic eruptions produce a strengthened NH winter polar vortex is
briefly summarized by examining ERA-interim reanalysis data. Then, the materials and methods of the
present modelling study are presented including the model and volcanic forcing sets used, and the
design of the ensemble simulations and analysis.

- 98 2.1 NH polar vortex response to volcanic eruptions in ERA-interim
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The NH polar vortex is highly variable as a result of unforced internal variability, and the impact of external forcings such as volcanic eruptions, the 11-year solar cycle, the El Niño Southern Oscillation (ENSO), and the QBO. Isolating the vortex response to any individual forcing term can be difficult, especially in the case of major volcanic eruptions for which so few actual events have occurred within the era of satellite measurements.

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106 A common and simple method to isolate the pure volcanic impact on the state of the NH winter 107 stratosphere is to simply average post-eruptive anomalies over winters after recent major volcanic 108 eruptions. This technique is here applied to ERA-Interim reanalysis (Dee et al., 2011) temperature and 109 zonal wind fields after the eruptions of El Chichón (1982) and Pinatubo (1991). Post-eruption winter 110 anomalies are constructed for the two winters after each eruption by differencing post-eruptive 111 December-to-February (DJF) mean fields with fields averaged for 3 (El Chichón) and 5 (Pinatubo) years 112 before the eruption (the shorter reference period for El Chichón resulting from the fact that the ERA-113 Interim data set begins in 1979). Differences in post-eruption anomalies for Pinatubo are not strongly 114 dependent on the choice of a 3, 4 or 5 year reference period.

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116 Figure 1 (left) shows ERA-Interim temperature and zonal wind anomalies composited for 4 winters, the 2 117 winters each after Pinatubo and El Chichón. Mean temperature anomalies in the post-volcanic 118 composite show positive anomalies in the tropical lower stratosphere, as would be expected due to 119 aerosol heating. Temperature anomalies also show cooling of the tropical upper stratosphere, cooling of 120 the NH polar lower stratosphere, and warming of the NH polar upper stratosphere. Mean post-volcanic 121 zonal winds show a strengthening of the NH winter polar vortex by ~ 6 m/s. Such a simple average with a 122 small sample size does not completely remove the influences of variability resulting from other forcing 123 terms—notably the solar activity was at a maximum at the times of both the El Chichón and Pinatubo 124 eruptions, and El Niño events were observed in the first winters after both eruptions—but the 125 composite temperature and zonal wind anomaly structure is certainly a better approximation of the 126 direct volcanic impact than anomalies in any single post-volcanic year. For comparison, single winter 127 anomalies are shown for the first winter after the Pinatubo eruption. Temperature anomalies for DJF 128 1991/92 roughly follow the structure of the 4-year composite, albeit with tropical positive anomalies 129 located at higher altitudes, and weaker polar lower stratosphere cooling. Post-Pinatubo zonal wind 130 anomalies in the tropics reflect the state of the QBO at the time, with negative (easterly) wind 131 anomalies in the middle stratosphere (~50-15 hPa) and positive (westerly) anomalies in the upper

132 stratosphere (15-2 hPa). The polar vortex in the first post-Pinatubo winter was actually not as clearly 133 enhanced as the 4-year composite, with positive zonal wind anomalies only in the mid to lower polar 134 stratosphere centered at ~70°N. It is likely that in addition to the volcanic forcing, the vortex in DJF 135 1991/92 was weakened somewhat due to the influences of the concurrent forcing of the El Niño and 136 QBO easterly phase. Based on these arguments, it can be hypothesized that the pure response of the 137 stratosphere to Pinatubo aerosol forcing would have the approximate structure of the composite 138 response shown in Fig. 1 (left), albeit with greater amplitude, since the aerosol optical depth and hence 139 aerosol radiative heating during the first post-Pinatubo winter is the strongest of the years used in the 140 volcanic composite. This "expected" response in the first post-Pinatubo winter is based on a small 141 sample size of observations, and an assumption of linear response to the magnitude of volcanic forcing, 142 however, in light of limited evidence, it represents a best first-order, observation-based expectation, 143 consistent with that assumed explicitly or implicitly in prior studies.

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145 **2.2 MPI-ESM**

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The MPI-ESM is a full Earth System model, with atmosphere, ocean, carbon cycle, and vegetation components. Major characteristics of the full ESM and its performance in the CMIP5 experiments are described by Giorgetta et al. (2013). The "low resolution" model configuration (MPI-ESM-LR) is used here, with horizontal resolution of the atmospheric component given by a triangular truncation at 63 wave numbers (T63) and 47 vertical layers extending to 0.01 hPa. Unlike model configurations with higher vertical resolutions, the LR version has no internally generated Quasi Biennial Oscillation (QBO).

154 CMIP5 historical simulations have previously been performed with the MPI-ESM model over the time 155 period 1850-2005. Prescribed external forcings for the historical simulations, including volcanic aerosols 156 as well as greenhouse gases and ozone follow CMIP5 recommendations, and the responses to these 157 forcings are described by Schmidt et al. (2013).

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The MPI-ESM is configured to take volcanic aerosol forcing data in two formats, both of which are used in this study. One format consists of monthly zonal mean values of aerosol extinction, single scattering albedo, and asymmetry factor as a function of time, pressure, and the 30 short wave and long wave spectral bands used by the model. This format is consistent with the observation-based forcings sets introduced in Sec 2.3. A second format consists of monthly zonal mean aerosol extinction at 0.55 μm, 164 and zonal mean effective radius, both as a function of latitude, height, and time. Pre-calculated look-up 165 tables are then used to scale the 0.55 μ m extinction to the wavelengths of the model's radiation code 166 based on the effective radius. This methodology has been used to perform MPI-ESM simulations using 167 the forcing timeseries of Crowley et al. (2008, e.g., Timmreck et al., 2009) or output from prior runs of 168 the MAECHAM5-HAM coupled aerosol-climate model, (e.g., Timmreck et al., 2010) as done in this study 169 and described in Sec 2.4.

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171 2.3 Observation-based aerosol forcing sets

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173 Volcanic sulfate aerosol forcing for the MPI-ESM CMIP5 historical simulations is based on an extended 174 version of the aerosol data set developed by Stenchikov et al. (1998, hereafter S98) on the basis of SAGE 175 II measurements of aerosol extinction at 1.02 μ m and estimates of effective radii derived from 176 instruments on the Upper Atmosphere Research Satellite. The data are given at 40 pressure levels and 177 interpolated to the actual hybrid model layers during the simulations. The S98 data set is based 178 primarily on retrievals of aerosol extinction at 1.02 µm from SAGE II, with gaps filled with data from 179 ground-based lidar systems. Together, the S98 forcing set and that from Sato et al., (1993, with updates 180 http://data.giss.nasa.gov/modelforce/strataer/), both primarily based on SAGE II data, have been used 181 in roughly half of the models that performed CMIP5 historical simulations (Driscoll et al., 2012). 182 183 Subsequent updates to the SAGE II retrievals have led to significant changes in the space-time 184 morphology of the estimated aerosol extinction after Pinatubo (Arfeuille et al., 2013; Thomason and 185 Peter, 2006). A new volcanic forcing set (SAGE_4 λ , Arfeuille et al., 2013) has been compiled and made 186 available to modeling centers (http://www.pa.op.dlr.de/CCMI/CCMI_SimulationsForcings.html), 187 specifically for use in Chemistry climate simulations within the Chemistry Climate Model 188 Intercomparison (CCMI) initiative (Eyring and Lamarque, 2013). Timeseries of zonal mean aerosol optical 189 depth (AOD) at 1 µm – the wavelength closest to the original SAGE II measurements and so less impacted 190 by uncertainties in derived aerosol size distribution—are shown in Fig. 2 over the Pinatubo period for 191 the S98 and SAGE_4 λ reconstructions. 192

193 It should be noted that even for Pinatubo—the best-observed eruption in history—the observation-

194 based volcanic aerosol forcing sets suffer from significant but mostly unquantified uncertainties. Most

195 notably, gaps in the satellite record result from sparse sampling of the satellite instruments (Stenchikov

- et al., 1998) and the fact that large optical depths in the initial months after the Pinatubo eruptionreduced atmospheric transmission below detectability (Russell et al., 1996).
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199 2.4 Model-based aerosol forcing sets

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201 Two "synthetic" volcanic aerosol forcing sets were constructed based on a 12-member ensemble of 202 simulations of a Pinatubo-like eruption using the MAECHAM5-HAM coupled aerosol-climate model with 203 SO₂ injections of 17 Tg and prescribed climatological sea surface temperatures (Toohey et al., 2011). 204 Figure 3 shows lower stratospheric zonal mean temperature anomalies (at 100 hPa) and zonal wind 205 anomalies (at 50 hPa) for these simulations, in comparison with ERA-Interim post-volcanic anomalies 206 described in Sec. 2.1. Most of the MAECHAM5-HAM ensemble members (roughly 9/12) show 207 characteristics of a strengthened polar vortex in the lower stratosphere, as quantified as negative 208 temperature anomalies at polar latitudes and positive zonal wind anomalies between 60° and 80°N in 209 Fig. 3. However, the ensemble variability of the simulations is pronounced, with 3 members showing a 210 weakened polar vortex with positive temperature anomalies over the polar cap and negative wind 211 anomalies between 60° and 80°N. From the full 12-member ensemble, two subsets were defined based 212 on the zonal wind anomalies, with strong and weak vortex composites (hereafter SVC and WVC) 213 selected respectively as the 3 members with the most positive and most negative zonal wind anomalies 214 at 50 hPa and 70°N. Aerosol properties (aerosol extinction at 0.55 µm and effective radius) for these two 215 composites were collected for use in MPI-ESM simulations. SVC and WVC zonal mean AOD timeseries, 216 scaled to 1 μ m so as to be consistent with the observation-based AOD timeseries, are shown in Fig. 1. 217

218 2.5 Experiments

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The forcing sets described above were used to force four sixteen-member ensemble simulations of the Pinatubo eruption time period. The number of MPI-ESM realizations used here is therefore notably greater than the three MPI-ESM realizations used in prior single- (Schmidt et al., 2013) or multi-model (Charlton-Perez et al., 2013; Driscoll et al., 2012) investigations of the CMIP5 historical simulations. Ten of the sixteen unique initial condition states (at June 1991) were taken from ten independent, preexisting CMIP5 historical simulations. In order to increase the ensemble size, six of the historical simulations were restarted in 1980 with a small atmospheric perturbation applied, and integrated until 1991. All simulations were therefore forced with S98 volcanic forcing up until June 1991, at which point
forcing either continued as S98, or switched to one of the other forcings.

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A control ensemble (CTL) is based on the five year period 1986-1990 for the original ten historical
 simulations used to produce the initial conditions, comprising in total 50 years, during which the other
 external forcings are negligibly different from 1991-1992 conditions. Anomalies for the volcanic
 ensembles are computed by differencing ensemble mean results with the CTL ensemble mean.

235 Observations imply that the NH winter vortex response to volcanic forcing may last for two years. While 236 considering the results for the first two winters together certainly increases the ensemble size, it is quite 237 possible that the mechanisms leading to vortex response are different in the first and second winters 238 after an eruption. The magnitude of aerosol forcing will be different for the two winters for any forcing 239 set, with typically stronger aerosol forcing occurring in the first winter. Such temporal changes are likely 240 larger than differences between the different forcing sets used in this study. With these potential 241 complications in mind, in order to simplify the analysis and interpretation of the model results in this 242 study, we choose to focus our analysis only on the first winter after the Pinatubo eruption. 243 244 Unless stated, results shown are ensemble means. Confidence intervals (95% level) are calculated for

245 differences between forced and CTL ensemble means and the differences between the forced ensemble 246 means using a bootstrapping technique, with 1000 resamples of the original ensembles. For example, 247 for each latitude and pressure level, and for each forced ensemble, 1000 bootstrapped sample means 248 are produced by sampling the 16 ensemble member values with replacement, resulting in an 249 approximation of the uncertainty distribution for the mean value. The same process is applied to the CTL 250 ensemble, with n=50. The difference between the two bootstrapped distributions defines the 251 uncertainty in the mean difference, and is used to define the 95% confidence interval of the ensemble 252 mean difference. A parallel procedure is used to define the confidence intervals for the differences of 253 two volcanically forced ensembles.

- 254
- 255 3 Results

- 257 3.1 Radiative forcing
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259 Latitude-pressure plots of zonal mean DJF 1 µm extinction (EXT) are shown in Fig. 4 for the S98 and 260 SAGE_4 λ forcing sets. A major difference between the two forcing sets is the vertical distribution of 261 extinction in the tropics, with SAGE 4 λ extinction being more constrained to the lower stratosphere, 262 compared to S98 which has considerable extinction extending down into the upper troposphere. This 263 difference in tropical extinction is the result of improvements in the SAGE $_4\lambda$ retrieval: the S98 vertical 264 distribution in the tropics is very likely unrealistic (Arfeuille et al., 2013). The two forcing sets also differ 265 in terms of the magnitude of tropical extinction, with SAGE 4λ having stronger extinction in the tropical lower stratosphere for all wavelengths greater than 0.55 µm. Differences in extinction magnitude are 266 267 also apparent in the high latitude lower stratosphere of both hemispheres, with SAGE 4 λ extinction 268 much smaller than that of \$98.

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Direct aerosol radiative heating rates (Q^{aer}) are computed by performing radiative transfer calculations
at each model timestep twice, once with and once without the volcanic aerosols [as in *Stenchikov et al.*,
1998]. We have calculated Q^{aer} for single realizations of each forced ensemble.

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Net (longwave + shortwave) Q^{aer} for DJF for the two observation-based forcing sets is shown in Figure 4. 274 275 Q^{aer} values are positive over most of the stratosphere for both S98 and SAGE 4 λ forcings, with highest 276 magnitude in the tropical lower stratosphere at approximately 30 hPa, just north of the equator. Like 277 the extinction values (Fig. 4, upper row), S98 heating rates are spread over a larger vertical extent than 278 SAGE $_4\lambda$ heating rates: at the equator, S98 heating rates > 0.1 K/day extend from ~100 hPa upwards whereas for the SAGE_4 λ forcing set, heating rates >0.1K/day begin ~60 hPa. Like the extinction at 1 μ m 279 (and longer wavelengths) SAGE 4 λ forcing leads to stronger Q^{aer}, with maximum values of 0.5 K/day 280 compared to max values of 0.3 K/day for S98. Although minor compared to the differences in the 281 tropical stratosphere, there are differences in Q^{aer} in the NH polar latitudes, with S98 leading to slightly 282 larger Q^{aer} values in the NH polar lower stratosphere than for SAGE 4 λ . 283

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Figure 5 shows DJF 1 µm aerosol extinction for the SVC and WVC forcing sets. Compared to the
observation-based forcing sets, both model-based forcing sets have greater magnitudes of aerosol
extinction, especially in the high latitudes. Such differences likely arise from a combination of (1)
underestimates in the observation-based forcing sets due to saturation effects, especially in the weeks
directly following the eruption, and (2) potential errors in the model simulations, either in terms of
general model deficiencies, or errors in the model formulation in regards to the actual Pinatubo

291 eruption (e.g., SO₂ injection amount or height). The model-based extinctions also have stronger 292 gradients (both vertical and horizontal) across the tropopause compared to the observation-based forcing sets. These differences, especially in the vertical, are likely due to the limited vertical resolution 293 294 of the observations. The primary difference between the SVC and WVC forcings is the hemispheric 295 partitioning of the aerosol extinction, with WVC having larger extinctions in the NH than SVC. We 296 interpret this difference as a result of wave driven circulation in the NH: in cases of strong NH wave 297 forcing in the original MAECHAM5-HAM runs, a stronger residual circulation transports more aerosol 298 from the tropical region towards the NH, while also disturbing the NH polar vortex. Therefore, by 299 selecting cases of weak polar vortex, we also select cases of strong northward aerosol transport. 300 Another major difference is the magnitude of extinction in the NH high latitudes, and therefore the 301 gradient in extinction around 60°N. The stronger aerosol extinction gradient in the SVC forcing set is 302 obviously a result of the strong vortex in the MAECHAM5-HAM simulations, which inhibits mixing of 303 aerosol into the polar cap. In the tropics, SVC and WVC have very small differences, and their vertical 304 distribution is almost identical.

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306 Q^{aer} values from the model-based forcing sets, also shown in Fig. 5, are more similar to the observationbased Q^{aer} values than the extinctions: the differences is high latitude extinction have relatively little 307 impact on the Q^{aer} differences due to the much weaker LW radiation field here than in the tropics. 308 Differences in Q^{aer} between the two model-based forcing ensembles are relatively small (compared to 309 310 differences between the two observation-based forcing ensembles), and are characterized primarily by the north-south shift between the two forcing sets, and the gradients in Q^{aer} at 60°S and 60°N. At high 311 latitudes, Q^{aer} values are apparently not strongly dependent on the aerosol extinction, e.g., SVC has 312 smaller aerosol extinction values in polar latitudes than WVC, but shows larger Q^{aer}. This is likely due to 313 314 the fact that the net (absorption-emission) long wave heating rate is strongly temperature dependent 315 because of the temperature dependence of the emission. As shown below, the SVC ensemble is characterized by a colder polar vortex than for WVC, which should lead to less emission and thus larger 316 317 net heating.

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319 **3.2 Temperature and zonal wind response**

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321 In order to determine which thermal and dynamical responses to the volcanic forcings are robust across

322 the different forcings, results are first examined by concatenating the 64 volcanic simulations into one

"grand" ensemble (VOLC), with ensemble mean anomalies of temperature and zonal wind shown in Fig.6.

325

326 Significant and robust temperature responses in the volcanic simulations include positive anomalies in 327 the tropical lower stratosphere and negative temperature anomalies in the troposphere, extending to 328 the surface between approximately 70°S-45°N. In the NH high latitude stratosphere, the temperature 329 anomaly structure for the VOLC ensemble is roughly consistent with that shown by Schmidt et al. (2013) 330 for three MPI-ESM realizations (using S98 forcing), with positive temperature anomalies in the upper 331 stratosphere and mesosphere. However, while temperature anomalies for the VOLC ensemble are 332 significant throughout most of the tropopause and lower-to-middle stratosphere, anomalies in the NH 333 polar region are generally not significant.

334

Significant zonal wind responses in the VOLC ensemble include weakening of the subtropical jets at ~30°
 and ~100hPa of both hemispheres by 2-4 m/s, and a weakening of the SH stratospheric easterlies by 4-6
 m/s. Grand ensemble mean zonal wind anomalies in the NH high latitude stratosphere reach a
 maximum of ~2 m/s, and are not significant.

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340 Temperature and zonal wind anomalies are shown separately for the observation- and model-based 341 forcing ensembles in Figures 7 and 8, respectively. For the observation based forcings S98 and SAGE_ 4λ 342 (Fig. 7), the general features are consistent with the VOLC grand ensemble. In agreement with Q^{aer} of 343 Fig. 4, SAGE_ 4λ tropical temperature anomalies are greater in magnitude, with peak values of 4.8 K 344 compared to 3.6 K for S98. Temperature anomalies are also shifted in height between the two 345 ensembles, with peak temperature anomalies located at 30 hPa for SAGE_4 λ , compared to 50 hPa for 346 S98. Differences in tropical temperature anomalies between the two ensembles (right-most column of 347 Fig. 7) are significant at the 95% confidence level between approximately 200 and 20 hPa. The SAGE 4λ ensemble shows slightly larger warming in the polar upper stratosphere, although the difference 348 349 between S98 and SAGE_4 λ is not significant.

350

In the NH high latitudes, the S98 ensemble shows a weak (2-3 m/s) increase in westerly wind, while the

352 SAGE $_{4\lambda}$ ensemble shows a weak (3-4 m/s) negative wind anomaly in the upper

353 stratosphere/mesosphere. Differences between the two ensembles are not significant in the mid-to-

354 high latitudes.

356 For the model-based forcing ensembles SVC and WVC, temperature anomalies in the tropics and 357 midlatitudes are quite similar in structure between the two ensembles and similar to the grand 358 ensemble mean. In the NH high latitudes, however, the temperature responses are quite different in 359 structure between the SVC and WVC ensembles. The SVC ensemble produces a NH high latitude 360 temperature anomaly pattern with significant warming in the upper stratosphere and lower 361 mesosphere. WVC, on the other hand, gives a temperature anomaly pattern with insignificant positive 362 temperature anomalies in the polar lower and middle stratosphere. Differences between the two 363 forcing sets are significant in the polar lower and upper stratosphere.

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Zonal wind anomalies for the model-based forcing ensembles follow from the temperature anomalies.
The SVC ensemble produces a significant strengthening of the NH polar vortex, with peak zonal wind
anomalies 6-8 m/s. The WVC ensemble produces no significant change in NH vortex winds. The
difference between the vortex response in the SVC and WVC ensembles is significant in the polar lower
stratosphere, and in the mid-to-upper stratosphere at the highest latitudes (70-90°N).

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371 DJF zonal mean zonal wind at 60°N and 10 hPa for each realization of each ensemble is shown in Figure 372 9. Ensemble means with 95% confidence intervals are also shown. The mean of the control ensemble is 373 marked by a horizontal dashed line. Consistent with results discussed above, only the SVC ensemble 374 shows a significant zonal wind response to volcanic forcing, with an ensemble mean 95% interval which 375 excludes the control mean. In terms of individual ensemble members, 13/16 members of the SVC 376 ensemble have zonal mean wind greater than the control mean, although clearly all of the members lie 377 within the natural variability of the control ensemble. Three SVC members have zonal winds weaker 378 than the control; as such, the response to SVC forcing is not totally robust and the increase in ensemble 379 mean vortex strength appears to represent a change in the probability of weak vs. strong vortex states. 380 The WVC and S98 ensembles show ensemble means and spreads very similar to the control ensemble, 381 with S98 showing a weak and insignificant positive anomaly in wind strength. The SAGE_4 λ ensemble 382 shows an ensemble mean equal to that of the control ensemble, but interestingly shows some evidence 383 of a decrease in ensemble variability.

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385 **3.3 Wave-driven circulation response**

For all ensembles, we have computed transformed Eulerian mean (TEM) diagnostics (Andrews et al., 1987), including Eliassen-Palm (EP) fluxes, the meridional residual mass circulation stream function (ψ^*), the residual vertical velocity (\bar{w}^*), and temperature tendencies due to vertical residual advection. Ensemble means of these quantities have been compared to values from the control ensemble in order to compute post-volcanic anomalies in the first NH winter. As for the temperature and zonal wind anomalies, selected TEM quantities are examined first in terms of the grand VOLC ensemble in Figure 10.

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395 It is well known that the variability of polar vortex strength is largely controlled by planetary wave drag, 396 and therefore on the upward wave flux from the troposphere into the stratosphere (Newman et al., 397 2001; Polvani and Waugh, 2004). The vertical component of EP-flux (F_z) is a commonly used proxy for 398 the amount of wave activity entering and propagating through the stratosphere. Figure 10 shows DJF F_{γ} 399 for the control ensemble and anomalies for VOLC volcanic ensemble. Around the tropospause (~100-200 400 hPa), F_z anomalies in the VOLC ensemble are negative in the midlatitude (30-45°) regions of both 401 hemispheres, and generally positive in the mid-to-high latitudes (45-90°N). Positive F₂ anomalies are 402 significant throughout the SH stratosphere, while in the NH, significant positive F, anomalies occur 403 around 60°N and 100 hPa, and extend upwards slanting equatorward with height. 404

405 Convergence of EP-flux (or negative values of EP-flux divergence, EPFD) leads to wave drag, a slowing of 406 the wintertime westerly (eastward) zonal wind and a poleward residual circulation. Enhanced wave drag 407 in volcanic simulations is found throughout the SH stratosphere and in the midlatitude NH middle 408 stratosphere (around 30°N, 10 hPa, Fig. 10d). Wave drag in the latter location is especially important for 409 forcing the residual circulation (Shepherd and McLandress, 2011). Figure 10 (e,f) shows the CTL and 410 VOLC anomalies of the meridional residual circulation stream function. The poleward NH residual 411 circulation stream function is found to be enhanced in the VOLC ensemble.

412

The volcanically-induced residual circulation anomalies drive adiabatic heating anomalies where vertical motions are induced. Temperature tendency anomalies due to residual vertical velocity (hereafter dT_{w^*} , Fig. 10g,h) show clearly the tropical cooling associated with anomalous vertical upwelling, and heating at the mid and high latitudes due to anomalous downwelling.

418 The terms $dT_{\tilde{w}^*}$ and Q^{aer} are found to dominate the temperature tendency budget compared to other

419 terms in the TEM diagnostics. Therefore, the temperature anomalies found in the volcanic simulations

420 can be understood to be primarily the result of the direct (diabatic) aerosol heating Q^{aer}, and the

421 indirect, dynamical (adiabatic) heating $dT_{\tilde{w}^*}$. These terms, along with the corresponding temperature

anomalies, are shown for each of the volcanic ensembles as lower stratosphere (100-20hPa) averages in

423 Figure 11.

424

425 Ensemble mean temperature anomalies show roughly similar behavior in the tropics and midlatitudes 426 (0-55°N) for all ensembles. In the NH polar latitudes, weak positive temperature anomalies are 427 simulated for the S98, SAGE 4 λ and WVC ensembles, and negative anomalies for the SVC ensemble. Q^{aer} 428 peaks in the tropics and decays to zero between 30 and 60°N. $dT_{\bar{w}^*}$ is negative in the tropics, opposing 429 the Q^{aer} heating, and becomes generally positive in the extratropics where downward advection occurs. 430 In the high latitudes, dT_{w^*} is positive for S98, SAGE_4l and WVC, but negative for SVC, consistent with the 431 temperature anomalies (Fig 8). It is thus clear that the differences in temperature anomalies at high 432 latitudes, and therefore the temperature gradients around 60°N, result from differences in dT_{w^*} 433 between the ensembles. The effect of differences in temperature tendencies on temperature anomalies 434 is likely amplified at higher latitudes since radiative damping timescales increase with latitude in the 435 lower stratosphere (Newman and Rosenfield, 1997). These results underscore the point that 436 temperature gradients at the high latitudes are controlled by the structure of the volcanically induced 437 residual circulation anomalies rather than the direct aerosol heating.

438

Differences in the dynamical responses to the SVC and WVC volcanic forcings are further explored by
 examining TEM diagnostics for these ensembles. As in the VOLC grand ensemble, both model-based
 forced ensembles show positive F_z anomalies throughout the SH stratosphere and negative anomalies in
 the subtropical tropopause region (Fig. 12a,b). Positive F_z anomalies are found at the NH high latitude
 lower stratosphere, extending upwards and slanting equatorward with height. However, these positive
 F_z anomalies are much stronger (and significant only) in the WVC ensemble.

445

As in the VOLC ensemble, negative EP-flux divergence (wave drag) is significantly enhanced in the SH
stratosphere and in the NH midlatitude middle stratosphere in both the SVC and WVC ensembles (Fig.
10d,e). NH wave EP-flux divergence anomalies are notably stronger in the WVC ensemble between 10
and 100 hPa and 30-60°N; differences between the two ensembles are significant around 100hPa and

450 60°N. Residual circulation anomalies (Fig. 12g,h) show significant poleward circulation cells in both 451 hemispheres. Notable differences in the form of the induced residual circulation cells between the two 452 ensembles exist in the lower and middle NH stratosphere, with circulation cells confined to 0-45°N in the 453 SVC ensemble, and extending 0-90°N in the WVC ensemble. These differences are significant in around 454 100 hPa and 60°N, and as such appear related to the differences in EP-flux divergence discussed above. 455 Temperature tendency anomalies due to residual vertical velocity $(dT_{w^*}, Fig. 10j,k)$ show that the broad 456 residual circulation anomaly cell in the WVC ensemble leads to dynamical heating of the lower 457 stratosphere extending from ~45N-90°N, while the narrower poleward cell in SVC leads to significant 458 heating only in the midlatitude (45-60°N) lower stratosphere. Differences in dynamical warming are 459 marginally significant in the polar lower stratosphere, consistent with the differences in EP-flux 460 divergence and residual circulation.

461

462 The TEM fields of Fig. 12 thus show that the significant difference in polar vortex zonal wind response to 463 the SVC and WVC volcanic forcings comes about due to differences in the wave activity propagating 464 upwards into the stratosphere from the troposphere. This mechanism is further explored in Fig 13 for all 465 volcanic ensembles. Fig. 13 (left) shows the relationship between polar vortex wind (u at 10 hPa, 60°N, 466 hereafter u_{vortex}) and lower stratosphere polar temperature (T at 50 hPa, 60-90°N average, hereafter 467 T_{vortex}). Individual ensemble members are shown by circular markers, while ensemble mean values are 468 shown as triangles on the bottom and right-hand axes. The compact linear relationship between uvortex 469 and T_{vortex} apparent in the control run is shifted in the volcanic simulations: for a given T_{vortex}, the 470 associated u_{vortex} is typically ~5 m/s stronger in the volcanic runs than in CTL. The SVC ensemble mean 471 T_{vortex} is similar to the CTL ensemble mean, and correspondingly, the SVC u_{vortex} is larger than the CTL 472 ensemble mean by 5-10 m/s. In contrast, the other volcanic ensembles all show increases in Tvortex, and 473 correspondingly, uvortex for these ensembles is less than that of SVC, and similar to the CTL ensemble 474 mean value.

475

Polar lower stratosphere temperatures are related to upward EP-flux in the midlatitudes (Newman et
al., 2001; Polvani and Waugh, 2004). In Fig 13b, u_{vortex} is plotted vs. F_z at 100 hPa averaged over 45-75°N
(hereafter F_{z,100}, a common scalar metric for wave activity entering the extratropical stratosphere). Polar
vortex strength, u_{vortex}, is seen to be stronger for any particular F_{z,100} value in the volcanic runs compared
to the CTL ensemble. For the SVC ensemble, F_{z,100} is comparable to the CTL value, and the u_{vortex} anomaly
is 5-10 m/s. The other ensembles show positive anomalies in F_{z,100}; the zonal wind of these volcanic

ensembles are thus reduced due to the increased wave activity and hence wave drag. Post-volcanic
 anomalies in F_{z,100} thus exert a negative feedback on the wind anomalies driven by changes in the lower
 stratosphere temperature gradient.

485

486 4. Discussion

487

488 A major result of the preceding sections is that the temperature structure of the lower stratosphere in 489 post-volcanic winters in the high latitudes is controlled primarily by induced residual circulation 490 anomalies. Post volcanic-eruption enhancement of the stratospheric residual circulation (or BDC) have 491 been suggested based on previous model studies (Pitari and Mancini, 2002; Pitari, 1993). Graf et al., 492 (2007) assessed the observational record and found evidence of increased winter stratospheric wave 493 activity after three eruptions (Agung, El Chichón and Pinatubo). Poberaj et al., (2011) showed that 494 anomalously strong EP-fluxes occurred in the SH after Pinatubo. Such increases in winter stratospheric 495 wave activity are likely a result of changes in the wave propagation conditions of the stratosphere 496 following aerosol radiative heating, allowing more planetary waves to propagate upwards. Similar 497 arguments explain climate models' predicted increase in future stratospheric residual circulation due to 498 changes in the atmospheric temperature structure due to climate change (Garcia and Randel, 2008; 499 McLandress and Shepherd, 2009; Shepherd and McLandress, 2011), which also enhances the meridional 500 temperature gradient in the lower stratosphere, albeit at lower altitudes.

501

502 The residual circulation anomalies induced by the volcanic forcing in the ensembles strengthen the 503 climatological residual circulation, although the structure of the anomalies is different than the 504 climatology. For example, the induced upwelling is centered on the equator whereas the maximum 505 climatological upwelling is centered in the summer hemisphere, and the induced upwelling is strongest 506 above the level of maximum aerosol radiative heating, and even negative below. This latter point may 507 explain why post-Pinatubo anomalies in tropical upwelling are not apparent in observational records, 508 which are usually displayed as timeseries of upwelling in the lower stratosphere (Seviour et al., 2012). 509 We have shown that increased wave activity and wave drag in both hemispheres is a robust response to 510 volcanic aerosol forcing, and therefore a component of the residual circulation anomalies in the 511 volcanically-forced simulations results from wave drag anomalies. However, given the fact that 512 maximum anomalous tropical upwelling occurs at and above the location of maximum aerosol radiative 513 heating in the tropics, it seems possible that there is also a diabatic component to the anomalous

residual circulation, forced directly by the aerosol radiative heating, as suggested by previous studies
(e.g., Aquila et al., 2013; Pitari and Mancini, 2002).

516

517 The lack of NH polar vortex response to the Pinatubo forcings used here does not necessarily rule out 518 the possibility that other forcings could produce significant vortex responses. Most obviously, the 519 magnitude of the volcanic forcing may be essentially important. Toohey et al. (2011) found a significant 520 vortex response to a near-super eruption volcanic forcing, which likely produced a much stronger direct 521 aerosol radiative heating gradient in the mid to high latitudes. Similarly, Bittner et al., (Sensitivity of the 522 Northern Hemisphere winter stratosphere variability to the strength of volcanic eruptions, submitted 523 manuscript) find a significant response of the MPI-ESM NH polar vortex to volcanic forcing representing 524 that of the 1815 eruption of Tambora. It is also important that our simulations do not include possible 525 chemical depletion of stratospheric ozone brought about by the presence of volcanic aerosols, or ozone 526 anomalies due to changes in the residual circulation. Induced changes in the meridional structure of 527 stratospheric ozone can also affect temperature gradients in the lower stratosphere (Muthers et al., 528 2014), and may be a significant feedback in the response of the polar vortex to volcanic forcing. 529

530 While neither of the two observation-based volcanic forcing sets used here produced significant vortex 531 responses, the difference in the responses was found to be significant at the 95% level. While the 532 SAGE_4 λ is almost certainly a more accurate reconstruction of the true aerosol evolution than S98 in 533 many aspects, e.g. the vertical distribution of aerosol extinction in the tropics, it does not produce the 534 expected increase in NH polar vortex strength in simulations with the MPI-ESM. It actually leads to a 535 significantly weaker vortex than the S98 forcing, which is especially surprising given it produces stronger 536 radiative heating in the tropical lower stratosphere than S98.

537

538 While the difference between responses to SVC and WVC forcing was not significant, the SVC was the 539 only forcing to produce a significant vortex strengthening. Such differences in simulated vortex response 540 between the different ensembles imply sensitivity in vortex response to the exact structure of the volcanic forcing. These differences between the ensembles, and especially the significant response of 541 542 SVC ensemble should perhaps be taken with a grain of salt: the large variability of wave drag and vortex 543 dynamics necessitates the use of large ensembles in order to negate the impact of variability on the 544 ensemble means. While our ensemble size of 16 is larger than most prior single-model volcanic studies, 545 it may still be insufficient to unambiguously identify significant high latitude responses from volcanic

forcing. Furthermore, given the anomalous heating of the lower stratosphere in volcanic simulations, it could be the case that insignificant anomalies in residual circulation lead to significant anomalies in ensemble mean temperature gradients and therefore zonal wind. Nevertheless, planetary wave propagation through the tropopause region has been shown to be quite sensitive to local vertical and meridional gradients in zonal wind and temperature (e.g., Chen and Robinson, 1992). It is plausible that small differences in the structure of the prescribed volcanic aerosol forcing, and resulting radiative heating, temperature and wind, could have relatively large impacts on wave propagation.

553

554 **6. Conclusions**

555

In simulations of the post-Pinatubo eruption period with the MPI-ESM with four different volcanic
aerosol forcings, an enhanced polar vortex—which is expected based on limited observations and simple
theoretical arguments—was not a robust response. The responses that were significant and robust
across all four forcings in the NH winter stratospheric include: (1) positive temperature anomalies in the
lower tropical stratosphere, (2) enhanced F_z in the NH midlatitudes (40-60°N) and wave drag in the
midlatitude middle stratosphere, (3) enhanced meridional residual circulation, (4) dynamical cooling of
the tropical lower stratosphere and heating of the midlatitude lower stratosphere.

563

564 The lack of robust polar vortex response to volcanic aerosol forcing in the MPI-ESM simulations of this 565 work is consistent with the multi-model results of the CMIP5 historical experiments (Charlton-Perez et 566 al., 2013; Driscoll et al., 2012). We have shown that the meridional temperature gradient in the 567 extratropical lower stratosphere induced by volcanic aerosol forcing, and therefore the strength of the 568 induced stratospheric winter polar vortex wind anomalies, is controlled primarily by dynamical heating 569 associated with the induced residual circulation rather than the direct aerosol radiative heating. The 570 vortex response in the model is therefore much less robust than would be expected if it were directly 571 due to the aerosol radiative heating, as it is instead subject to complexity and variability of wave-mean 572 flow interaction in the winter stratosphere.

573

574 Our results further imply that the NH polar vortex response is quite sensitive to the specific structure of 575 the volcanic forcing used. Generally, volcanic forcing leads to an overall increase in resolved wave 576 activity entering the mid- to high latitude stratosphere, and the impact of this wave activity tends to 577 weaken the polar vortex, counteracting the impact of low-latitude heating. However, one volcanic 578 forcing set used here, the model-based SVC forcing, did not significantly affect high-latitude wave 579 activity, and thereby produced a strengthened polar vortex, approximately in line with the expected 580 response. We speculate that such sensitivity is due to the role that minor differences in volcanically 581 induced temperature and wind anomalies play in wave-mean flow interactions in the stratosphere. An 582 alternate theory is that differences in volcanic forcing lead to differences in wave generation in the 583 troposphere. In either case, the results imply that—at least for a Pinatubo magnitude eruption—very accurate reconstructions of volcanic aerosol forcing would be required to reproduce any impact of the 584 585 aerosols on polar vortex strength in climate model simulations.

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- 587
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- 596
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Figure 1: ERA-Interim DJF (top) temperature and (bottom) zonal wind anomalies for (left) post-volcanic
 composite (n=4) of two winters after eruptions of El Chichón and Pinatubo, and (right) 1991/92, the first

741 winter after the Pinatubo eruption.

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Figure 2: Zonal mean aerosol optical depth at 1 µm timeseries for the Pinatubo eruption, as 745

reconstructed from observations producing the S98 and SAGE_4 λ forcing datasets (left) and based on 746

747 MAECHAM5-HAM model simulations (right) and composited according to strong (SVC) and weak (WVC)

vortex states. Gray vertical lines demark the DJF period of interest. 748





Figure 3: First post-eruption NH winter (DJF) anomalies of (left) temperature at 100 hPa and (right) zonal wind at 50 hPa for the 12 member MAECHAM5-HAM Pinatubo ensemble of (Toohey et al., 2011)(gray lines). Ensemble members comprising the strong and weak vortex composites (SVC and WVC) as defined in the main text are marked by circle and cross markers, respectively. For comparison, ERA-Interim anomalies based on a composite of the 2 winters after the eruptions of El Chichón and Pinatubo (blue), and the single winter after the Pinatubo eruption (red) are also shown.





volcanic aerosol forcing sets in first DJF after Pinatubo (winter 1991/92). (bottom) aerosol radiative

heating (Q^{aer}) as computed within the MPI-ESM model for each forcing set. Right-hand column shows

764 S98-SAGE_4 λ differences for both 1 μ m extinction and Q^{aer}.



Figure 5: (top) latitude-pressure distributions of 1 µm aerosol extinction in the SVC and WVC volcanic 768

aerosol forcing sets in first DJF after Pinatubo (winter 1991/92). (bottom) aerosol radiative heating (Q^{aer}) 769

as computed within the MPI-ESM model for each forcing set. Right-hand column shows SVC-WVC 770

771 differences for both 1 μ m extinction and Q^{aer}.



Figure 6: DJF temperature (top) and zonal wind (bottom) from the CTL ensemble (left) and anomalies for the VOLC grand ensemble (right). Hatching highlights anomalies which are significant at the 95% confidence level based on a bootstrapping algorithm.



Figure 7: Temperature and zonal wind response of the observations-based volcanic forced ensembles.
Shown are DJF anomalies for the (left) S98 and (middle) SAGE_4λ ensembles, and (right) the difference
between the two ensembles for (top) temperature and (bottom) zonal wind. Hatching highlights

anomalies which are significant at the 95% confidence level based on a bootstrapping algorithm.



789

790 Figure 8: Temperature and zonal wind response of the model-based volcanic forced ensembles. Shown

are DJF anomalies for the (left) SVC and (middle) WVC ensembles, and (right) the difference between 791

792 the two ensembles for (top) temperature and (bottom) zonal wind. Hatching highlights anomalies which

793 are significant at the 95% confidence level based on a bootstrapping algorithm.



Figure 9: DJF zonal wind at 60N and 10 hPa for the CTL ensemble and the volcanically forced ensembles

797 S98, SAGE_4 λ , SVC and WVC. Individual ensemble members shown as gray symbols. Ensemble means

shown in colored symbols, with vertical whiskers representing the 95% confidence interval of the

ensemble mean. Dashed horizontal line shows the ensemble mean of the CTL ensemble.

800



- 806 Figure 10: Selected transformed Eularian mean diagnostics for the CTL ensemble (left) and anomalies for
- the VOLC grand ensemble (right). Shown are (a,b) vertical component of EP-flux, (c,d) EP-flux
- 808 divergence, (e,f) residual circulation and (g,h) heating due to residual vertical advection. Contours in (a)
- are in units of 1×10^4 kg s⁻². Stream function contours in (e) are log-spaced, ranging from 10 to 1×10^4
- 810 kg/m/s, while those in (f) are log-spaced ranging from 1 to 100 kg/m/s.

811



814 Figure 11: Lower stratospheric (100-20 hPa), DJF average thermodynamic quantities from the S98 and

SAGE_4 λ forced ensembles: (a) temperature anomalies, (b) temperature tendencies due to aerosol

816 radiative heating (solid) and due to vertical advection (dashed).

817



Figure 12: Selected transformed Eularian mean diagnostics for the SVC and WVC forced ensembles.
Shown are DJF anomalies for the (left) SVC and (middle) WVC ensembles, and (right) the difference
between the two ensembles for (a,b,c) vertical component of EP-flux, (d,e,f) EP-flux divergence, (g,h,i)

- 825 residual circulation and (j,k,l) heating due to residual vertical advection. Stream function contours are
- 826 log-spaced ranging from 1 to 100 kg/m/s.





829 Figure 13: Scatter plots of vortex wind, temperature and wave forcing for the CTL and four volcanically

830 forced ensembles. Shown are (left) zonal wind at 10 hPa, 60°N plotted against temperature at 50 hPa,

831 60-90°N averaged, (right) zonal wind at 10 hPa, 60°N plotted against the vertical component of EP-flux

averaged over 40-75°N. Individual ensemble members are shown by colored circles or crosses, and

833 ensemble mean values are shown by colored triangles along the bottom and right-hand axes of each

834 panel.