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A better understanding of cloud optical thickness derived from the passive sensors MODIS/AQUA and POLDER/PARASOL in the A-train constellation

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Abstract

Cloud optical thickness (COT) is one of the most important parameter for the characterization of cloud in the Earth radiative budget. Its retrieval strongly depends on instrument characteristics and on many cloud and environment factors. Using coincident observations from POLDER/PARASOL and MODIS/AQUA in the A-train constellation, geographical distributions and seasonal changes of COT are presented, in good agreement with general cloud climatology characteristics. Retrieval uncertainties mainly associated to sensor spatial resolution, cloud inhomogeneity and microphysical assumptions are also discussed.

Comparisons of COT derived from POLDER and MODIS illustrate that as the primary factor, the sensor spatial resolution impacts COT retrievals and statistics through both cloud detection and sub-pixel cloud inhomogeneity sensitivity.

The uncertainties associated to cloud microphysics assumptions, namely cloud phase, particle size and shape, also impact significantly COT retrievals. For clouds with unambiguous cloud phase, strong correlations exist between the two COTs, with MODIS values comparable to POLDER ones for liquid clouds and MODIS values larger than POLDER ones for ice clouds. The large differences observed in ice phase cases are due to the use of different microphysical models in the two retrieval schemes. In cases when the two sensors disagree on cloud phase decision, COT retrieved assuming liquid phase are systematically larger.

The angular biases related to specific observation geometries are also quantified and discussed in particular based on POLDER observations. Those exhibit a clear increase of COT with decreasing sun elevation and a decrease of COT in forward scattering directions due to sub-pixel inhomogeneities and shadowing effects, this especially for lower sun. It also demonstrates unrealistic COT variations in the rainbow and backward directions due to inappropriate cloud optical properties representation and an important increase of COT in the sun-glint directions in case of broken cloud.

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1 Introduction

Cloud optical thickness (COT) is a key parameter to characterize cloud optical properties, which play an important role in the determination of cloud radiative forcing (Jensen et al., 1994; Kristiansen and Kristjansson, 1999). Some studies with GCMs simulations showed that the changes in COT result in negative radiative feedbacks as climate warms (Roeckner et al., 1987; Mitchell et al., 1989) while others showed positive feedbacks (Tselioudis and Rossow, 1994; Chang and Coakley, 2006). Long records of space-borne measurements at global scale are actually needed in order to correctly map COT and better understand the radiative effects of clouds on climate changes, especially in response to anthropogenic activities and natural variations (Rossow and Schiffer, 1991). COT is derived from the so-called solar reflective approach using a non water-absorbing band at visible wavelength under the assumption that the reflectance is a one-to-one non-linear function of COT. A number of studies showed that this widely used method is practical and effective (King, 1987; Nakajima and King, 1990; Nakajima and Nakajima, 1994). However, because the method accuracy depends on different atmospheric and surfaces conditions, associated uncertainties need careful investigations. Errors can arise from spectral radiation calibration, radiative impact of the upper molecules, gas and aerosols, surface conditions, inappropriate cloud microphysics (especially cloud phase), horizontal and vertical cloud inhomogeneities, and so on. Some of these uncertainties have already been well qualified while others are not yet fully understood. Concerning unsuitable cloud microphysics, departure in observed cloud phase function can lead to an uncertainty in COT of approximately 2 (Malkova, 1973). An error of a factor 2 in assumed water droplet radius can induce an error in COT of about 10 % (Han et al., 1994) and for cirrus using a wrong particle shape can result in an over-estimation of COT by a factor that can exceed 3 (Mishchenko et al., 1996; Zhang et al., 2009). Due to the asymptotic shape of the relation between reflectance and COT, small uncertainties in reflectance can induce large errors in COT for thick clouds while uncertainties from the surface dominates the errors for thinner clouds

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(Platnick and Valero, 1995). Due to horizontal cloud inhomogeneities, the retrieved COT can be over or under estimated depending on cloud types, sensor resolution and observation geometries. For low spatial resolution, due to the so-called plan-parallel biases (Cahalan, 1994; Szczap et al., 2000), the COT is generally underestimated.

5 Observations of liquid clouds from AVHRR or POLDER show an important decrease in retrieved COT in forward directions, especially at oblique suns (Loeb and Davies, 1996; Buriez et al., 2001), which is darkened by cloud-side shadowing in forward scattering viewing geometry. On contrary cloud-side illumination can lead to a positive bias in backscatter directions (Várnai and Marshak, 2002; Iwabuchi and Hayasaka, 2002).

10 A better understanding and assessment of the possible uncertainties are needed to obtain a better cloud climatology and also for the improvement of retrieval algorithms. Many satellites such as Advanced Very High Resolution Radiometer (Heidinger et al., 2005), Spinnig Enhanced Visible and Infrared Imager (SEVIRI) (Roebeling et al., 2006), MODIS (Platnick et al., 2003), and POLDER (Buriez et al., 1997) have developed their
15 own algorithms. All these instruments have different spectral and spatial characteristics and different assumptions (microphysics, surface reflectance, etc.) are used in the retrieval associated algorithms. They have thus different strengths and limitations to retrieve COT and an assessment of these differences is required. In a previous study, we conducted the comparisons of the cloud fractions and cloud thermodynamic phases
20 derived from POLDER/PARASOL and MODIS/AQUA and analyze their main differences (Zeng et al., 2011a,b). In this paper, we focus on the study of the COT derived from these two sensors, both in the A-train constellation.

In Sect. 2, we present quickly the POLDER and MODIS instruments, the algorithms used to retrieve COT and the products considered for our analysis. In Sect. 3, we show
25 global comparisons of the two COT products. A discussion of potential uncertainties is provided in Sect. 4 according to spatial resolution, cloud microphysics and observation geometry. At last, conclusions and perspectives are given.

2 Instruments characteristics, data and algorithms

POLDER is a component of a series of sensors (Deschamps et al., 1994), developed by CNES (French Space Agency), flying on board PARASOL since 2004. It is a multispectral imaging radiopolarimeter designed to provide global and repetitive observations of the solar radiation and polarized radiance reflected by the Earth-Atmosphere system. The instrument design consists of a wide field of view (1800 km) telecentric optics, a rotating wheel carrying spectral filters and polarizers, and a CDD (Charged Coupled Device) array of detectors that induces a moderate spatial resolution of about 6 km at ground independent of the viewing angle. When it passes over a scene, POLDER acquires up to 16 successive multiangle measurements of both the total and polarized solar radiance in eight narrow bands from 443 to 1020 nm for daytime observations only.

MODIS is a 36-bands scanning spectroradiometer on board Aqua, launched in May 2002 as a part of NASA's Earth Observing System (EOS) (King et al., 1992). It provides 29 spectral bands at 1 km resolution, 5 bands at 500 m resolution and 2 bands at 250 m resolution. Its spectral coverage ranges from visible (VIS) to thermal infrared (IR) (0.415 to 14.235 μm). Cloud mask is based on a variety of tests using as many as 20 of these 36 spectral bands to maximize reliability of cloud detection

To determine COT, both the POLDER and MODIS operational algorithms employ the so-called plane parallel and independent pixel cloud assumptions (IPA clouds) with the solar reflective method (Nakajima and King, 1990). The official product for POLDER is derived from the 0.67 μm channel over land and 0.865 μm channel over ocean (Buriez et al., 1997) while MODIS makes use of the 0.645 μm band over land, the 0.858 μm band over ocean and the 1.24 μm band over sea ice or snow covered surfaces (Platnick et al., 2003). Concerning liquid clouds, MODIS uses different particle sizes obtained from its water absorption channel in near-infrared range while POLDER, which has no particle size information, assumes clouds composed of water droplets with a constant effective radius (9 μm over land and 11 μm over ocean) (Buriez et al., 1997; Parol et al.,

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1999). For ice clouds, MODIS Collection 5 products are derived using the Baum05 Model (Baum et al., 2005) which consists in a mixture of different ice crystal habits. It uses 12 size distributions of crystals composed of mixed habits (droxtals, aggregates, bullet rosettes, hollow columns, solid columns and hexagonal plates) with the fraction of each habit depending on particle size. These habits definitions are based on in situ observations from the FIRE-II experiment. POLDER uses a single and fixed Inhomogeneous Hexagonal Model (IHM) (C.-Labonnote et al., 2000, 2001), which assumes that light is scattered by randomly oriented hexagonal ice crystals containing air bubbles aimed at reproducing real crystals imperfections. The IHM model has been found to be in good agreement with both total and polarized reflectance measurements of POLDER and in situ observations. For POLDER, the surface albedo and bidirectional reflectance over land are obtained from surface parameters previously retrieved from POLDER observations under clear-sky conditions by the POLDER “Land surfaces” processing line (Leroy et al., 1997) while over ocean it is calculated using the Cox and Munk (1956) model depending on the surface wind velocity derived from ECMWF analysis. For MODIS, the surface albedo is from 16-day 1-km composites of clear-sky observations in MOD43 products (Strahler, 1999).

In the same way as for cloud fraction, the POLDER COT is first retrieved at the initial resolution of POLDER ($6 \times 7 \text{ km}^2$) for the 16 view directions. It is then averaged at the super-pixel resolution ($20 \times 20 \text{ km}^2$) to obtain a “directional” optical thickness (used in Sect. 4.3), which after angular averaging provides a mean COT and its standard deviation. Note that in this paper, in order to make easy comparisons with the MODIS product, we work only with the optical thickness, although using its multi-angular capability, a mean angular weighted spherical albedo is also derived for POLDER (Buriez et al., 2005).

The main differences between the POLDER and the MODIS processing lines that can affect the retrieved COT in case of thick clouds are thus the initial resolution (POLDER: $6 \times 7 \text{ km}^2$; MODIS: $1 \times 1 \text{ km}^2$) and the differences in cloud microphysics assumptions. MODIS has a higher resolution and provide particle size information. On

the other hand, POLDER takes advantage of up to 16 viewing directions to provide an average optical thickness and to assess its angular consistency. Indeed, a large angular dispersion of retrieved COT can be attributed to the departure of POLDER observations from the microphysical assumptions or from single one layer PPA model (Parol et al., 2000). The optical thickness comparisons presented in the following are made at the POLDER super-pixel scale of about $20 \times 20 \text{ km}^2$ using the PM Dataset described in Zeng et al. (2011a) which consists of official POLDER and MODIS level 2 products remapped to a common integrated sinusoidal equal area grid.

3 Results

3.1 Geographical distributions

We first compare POLDER and MODIS COT through their geographical distributions and their differences (POLDER minus MODIS) for overcast (Cloud Fraction, $CF=1$ for both sensors) clouds (Fig. 1). We present distributions for all overcast clouds (first line) and for overcast clouds separated by cloud phase classes. At first glance, one can notice that POLDER and MODIS show similar COT distributions regardless of cloud phase class with generally higher COT values for POLDER.

Also and not surprisingly, the thickest ice clouds ($COT > 20$) are found mostly in the Inter Tropical Convergence Zone (ITCZ) and in the Storm Tracks (ST) zones over ocean and over Northern South America, Southern South Africa, Southern Asia and Eastern North America and principally over land. The thickest liquid clouds ($COT > 20$) are also found in these regions except in the ITCZ where most of the convective clouds are very high and primarily composed of ice.

Secondly, regardless of cloud phase, thicker clouds are found over continent for both POLDER and MODIS. This agrees with ISCCP C product (Rossow and Schiffer, 1999) but not with ISCCP D product where the land-ocean contrast of COT has been removed primarily because a significant increase in the amount of detected thin cirrus

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has been found over land with a lower IR threshold. As we will see latter, error in phase identification for thin cirrus over liquid cloud can indeed lead to an artificial increase of COT. Moreover, low-level clouds over land can also extent to larger heights yielding larger optical thicknesses than low-level clouds over ocean (Warren et al., 1986, 1988).

For cases where POLDER and MODIS disagree on phase assumption, the POLDER-liquid/MODIS-ice class clouds have the largest COT and the POLDER-ice/MODIS-liquid class clouds tend to have the smallest COT. This suggests that the combined POLDER-liquid/MODIS-ice phase is associated with thick clouds in multi-layer systems. In these cases, angular polarized signal from POLDER is sensible to lower water level however the IR and NIR signals from MODIS give more correct information at the cloud top. The POLDER-ice/MODIS-liquid phase is suggested to be associated to thin clouds or aerosols over low clouds (Waquet et al., 2009). The angular polarized signal of the ground or non-spherical aerosols is much closer to the ice clouds.

For unambiguous ice clouds, the COT differences are negative over the whole globe. The largest differences appear over land, in the ITCZ and the STs. For confident liquid clouds the differences are also negative in the STs and over land but appear negligible over ocean in tropics and middle latitudes with slightly positive differences in the ITCZ and around the continents where they may be associated with polluted air and smaller droplets. For POLDER-ice/MODIS-liquid phase clouds, COT differences are slightly negative almost all over the globe with quasi-zero values found over some tropical oceans. For POLDER-liquid/MODIS-ice phase clouds, the differences are mostly positive over the entire globe, especially in the STs. These will be explained in Sect. 4.

In Fig. 2a–d, we present latitudinal variations of COT for overcast clouds for the four different cloud phase classes. We have also plotted latitudinal variations of the product of COT by the factor $(1 - g)$, which allows to account partially for the different assumed microphysics in POLDER and MODIS retrievals (see Sect. 4.2). The latitudinal distributions of COT for ice or liquid clouds have similar trends with large values in the Storm tracks zones (STs) and small values in the subtropics. As already mentioned,

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we see that ice and liquid COT have a different behavior in the ITCZ with larger values for ice clouds, which does not appear for liquid clouds. Confident single layer liquid clouds (i.e., without overlaying ice clouds) are rarely occurring in the ITCZ and result statistically in small mean liquid COT in this zone. COT of confident liquid clouds derived from POLDER is very similar to MODIS with maximum relative differences smaller than 10 % (Fig. 2a). For confident ice clouds, the POLDER COT is on average smaller than the MODIS one with an almost constant difference of 5 (Fig. 2d). For clouds with inconsistent phases determined, we see that the sensor using the liquid cloud model always provides larger COT than the other using ice cloud model (Fig. 2b,c). We note also that, in case of POLDER-liquid/MODIS-ice phase class clouds (Fig. 2b), the differences between the two tend to increase with latitude with the POLDER COT increasing significantly polewards.

Figure 2e–h presents latitudinal variations of COT for broken clouds. Broken clouds correspond to the cases where both sensors indicated broken clouds ($0 < CF < 1$). For these clouds, COT values are logically reduced compared to overcast clouds as clear sky part should often be present in the sub-pixel scale. For any phase class, we note that the MODIS COT is larger than the POLDER one. Differences between the two COT are a lot reduced when multiplied by CF (blue curves). We will discuss this point in Sect. 4.1. In addition, we note that POLDER COT presents smaller latitudinal variations than the MODIS one, which increase from mid-latitudes to high latitudes.

3.2 Seasonal cycles

The seasonal cycles of COT are plotted in Fig. 3 for MODIS and POLDER, separately for land and ocean. We subtract the annual mean COT from the monthly average in order to remove the systematic bias that exists between the two COTs and focus on seasonal variability analysis. The seasonal cycle is calculated for overcast clouds for both POLDER and MODIS sensors according to cloud occurrence regions (i.e. subtropics and middle latitudes over ocean and land in each hemisphere) and to their combined thermodynamic phases. From the figures, we first see that the COTs from

the two sensors depend logically on the seasons and regions. For a same region and with a consistent thermodynamic phase, MODIS and POLDER COTs have quite similar temporal variations, especially for clouds with consistent liquid phase (Fig. 3e–h). For liquid clouds (Fig. 3e–h), COT show almost asymmetrical characteristics in South and North hemispheres over both land and ocean, with thicker clouds in winter in each hemisphere. This agrees well with ISCCP climatology (Rossow and Schiffer, 1999; Rossow et al., 1989). The only noticeable exception is over land in the mid-latitude of South hemisphere certainly because of smaller sample. For ice clouds (Fig. 3a–d) in both hemispheres, seasonal variations are more pronounced with differences between ocean and land. Over ocean, thicker ice clouds appear in winter. Over land, thicker ice clouds tend to appear preferentially in summer as a result of stronger convection.

3.3 Pixel-to-pixel comparisons

We now show pixel-to-pixel comparisons between the COTs of the two sensors. Comparisons are analyzed in term of slope and correlation coefficient of the linear regression assuming linear relationship between the two dataset. These two parameters give a measure of the pertinence of the linear relationship and should be equal to one for perfect relation.

Among all conditions, overcast scenes over ocean can be considered as ideal situations for retrieving COT. Selecting these clouds, we plot the one-year statistical relationship between POLDER and MODIS COTs for four different combined thermodynamic phase classes in Fig. 4a–d. Overcast condition is determined as previously from combination of POLDER and MODIS. In these figures, we note that for the four different phase classes, the correlation coefficients are quite important with values greater to 0.8 and even above 0.92 when the two sensors agree on phase index (confident ice and liquid clouds). These results show that there is a strong linear relationship between POLDER and MODIS COTs. The slope is very good with value close to the unity for liquid water clouds but is only about 0.74 for ice clouds. For clouds with inconsistent phases between the two sensors (Fig. 4b,c), both slopes and correlation coefficients

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are logically worse compared to consistent phase cases (Fig. 4a,d). We also notice that the slope of the POLDER-ice/MODIS-liquid clouds (Fig. 4b) is below one, which means that the MODIS COT is larger than the POLDER one. On contrary, the slope of the POLDER-liquid/MODIS-ice clouds (Fig. 4c) is greater than one, which means that the POLDER COT is larger than the MODIS one. These results illustrate again, if needed, that a correct identification of phase and microphysics is critical for COT retrieval. We will discuss more precisely these differences and their associated impacts in Sect. 4.2.

Concerning broken clouds, comparisons between the two sensors are shown in Fig. 4e–h. The mean optical thickness at super-pixel scale is computed by averaging the optical thickness of the cloudy part. Regardless of thermodynamic phases, we find out that the correlation coefficients are typically smaller than 0.7 for broken clouds. The slopes are entirely below 1 with POLDER detecting much smaller COT. Compared to overcast conditions, at first appearance, the relationships between COTs from the two sensors seem thus to be much worse. Cloud detection is indeed more difficult for broken cloud scenes and some disagreements on cloud fraction between the two sensors due to the difference in spatial resolutions can explain part of the differences (Zeng et al., 2011a). This impact of sensor spatial resolution and its consequence results on COT retrieval will be discussed in Sect. 4.1.

4 COT retrieval uncertainties

4.1 Impact of the sensor spatial resolution

As already described, COT is only retrieved for the pixels detected as cloudy and the final COT is averaged at the super-pixel level considering only the cloudy pixels. Internally, the correctness of cloud identification can therefore strongly impact COT retrieval. MODIS with a higher spatial resolution identify much smaller clear scenes among clouds than POLDER. Consequently, for region of fractional cloud cover, the

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POLDER cloud cover is larger than the MODIS one (Zeng et al., 2011a). However, in such cases, except over bright surfaces, the reflectance used for the POLDER COT retrieval is reduced by the sub-pixel holes among clouds and the derived COT is thus smaller compared to the MODIS one. This bias directly associated to cloud identification can be attributed to the sensor resolution difference. It is clearly identifiable in Fig. 4e–h where the POLDER COT is much smaller than the MODIS one. In order to account for it, we present in Fig. 5 the pixel-to-pixel comparisons of the product COT \times CF for broken clouds. Compared to Fig. 4e–h, which show pixel-to-pixel comparisons of COT for the same cases, it is clear that both the correlations and the slopes of the relationship have been greatly improved for all phase classes. Linear regressions with higher confidence are much closer to the first bisecting line. This is confirmed by the curves of the product COT \times CF in Fig. 2e–h.

To go further, in Fig. 6a, we also plot for liquid oceanic clouds the two constants (slopes and correlation coefficients) of the relationships between COT and the product COT \times CF as a function of MODIS cloud cover. It is clear that both the slope and the correlation coefficient for the COT relationships increase with cloud fraction almost linearly. Smaller cloud cover leads to larger the dispersion, corresponding to weaker relationships between the two sensors products. Looking at the product COT \times CF, we notice that the slope is a lot improved (closer to unity) and less dependent on cloud cover. However, the correlation coefficient is just slightly improved and still increases with cloud cover showing that COT retrieval is anyway less reliable in case of broken cloud fields.

These results stress out again that the sensor resolution affects the COT retrieval via the cloud cover difference and that COT values should not be used without knowledge of cloud fraction. This well known mechanism should absolutely be accounted for when building statistics on cloud cover from multi-sensor dataset.

In addition, compared to MODIS, POLDER with a lower resolution ignores not only the sub-pixel cloudiness but is also less sensitive to some of the sub-pixel cloud inhomogeneities. When POLDER considers a cloudy pixel at a 6 km \times 7 km resolution,

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MODIS with a pixel resolution of $1\text{ km} \times 1\text{ km}$ accounts for a part of the sub-pixel variability. The convex nature of the reflectance-COT relationship implies that the mean COT of two cloudy pixels is larger than the effective COT derived from the mean reflectance of the two pixels (Cahalan, 1994; Szczap et al., 2000). Therefore, COT retrieved from the average reflectance of the pixels is smaller than the linear average value. As POLDER cannot account for the sub-pixel heterogeneities, the POLDER COT is expected to be smaller than the MODIS one. To observe this, in Fig. 6b, we present correlation coefficients and slopes between POLDER and MODIS COTs for different inhomogeneity parameters, which are determined as the ratio of the standard deviation to the mean COT computed from MODIS product. We see for homogeneous clouds (inhomogeneity parameter close to 0), that both the correlation coefficients and the slopes values are high and close to one. When clouds become more inhomogeneous, these two constants decrease steadily. It means that the more inhomogeneous the clouds are, the smaller the POLDER COT compared to the MODIS one and the weaker the coherence between the two are. Keep also in mind that when sub-pixel cloud inhomogeneities are important for the two sensors, the COTs retrieved from the two are underestimated.

4.2 Impact of cloud microphysics

Another important factor that impacts the retrieved COT concerns the choice of the microphysical model, in particular the associated scattering phase function. As shown in Zeng (2011b), POLDER and MODIS may have inconsistent phase decisions but even in case of consistent phases, they use different particle microphysical model to retrieve COT, especially for ice clouds (see Sect. 2). To account for these differences in a first order, we introduce the scaled optical thickness (King, 1987). “Scaled” means that the optical thickness is multiplied by the factor $(1 - g)$ where g is the asymmetry coefficient (Eq. 1). Scaled optical thickness depends thus less on cloud microphysics assumption (cloud phase, cloud particles radius and shapes) and almost only on surface albedo and cloud reflection. The potential optical thickness bias associated to the uncertainties of scattering model is assessed by comparing the scaled optical thickness between the

two sensors in Fig. 7, which can be compared to Fig. 4a–d.

$$\tau^* = \tau \times (1 - g) \quad (1)$$

Looking at Figs. 4 and 7 corresponding to different phase decisions (b and c), we see a great improvement between the two scaled COT which means that a significant part of the differences between POLDER and MODIS are directly linked to an important factor concerning inconsistent phase determination. Indeed, backward scattering for ice crystals is much stronger than for water droplets ($g_{\text{ice}} < g_{\text{water}}$). Thus, to reflect the same quantity of radiation back to satellites, COT of ice clouds need to be smaller compared to water clouds. This being considered, it is fairly straightforward to understand that the slope for POLDER-ice/MODIS-liquid class clouds (Fig. 4b) is below one and the slope for POLDER-liquid/MODIS-ice class clouds (Fig. 4c) is larger than one. As the use of scaled optical thicknesses (using corresponding asymmetry factor) partially remove dependence on microphysics assumption, closer relations between the two sensors are found for the inconsistent phase cases in Fig. 7b,c compared to Fig. 4b,c and the slopes of scaled optical thickness are much closer to unity.

For overcast clouds with consistent phases over ocean (Fig. 4a,d), where the smallest bias are expected from cloud detection, phase identification and surface impact, COT is retrieved with a higher accuracy and confidence from the two sensors. COT relationships between them should appear better with both slope and correlation coefficient closer to unity. That is confirmed for confident liquid cloud cases but not for confident ice cloud cases where the correlation coefficient is good but the slope value is biased away from one.

For liquid clouds, both sensors employ a Mie scattering model. POLDER has no real-time effective radius retrievals but uses a fixed value of $9 \mu\text{m}$ over land and $11 \mu\text{m}$ over ocean whereas MODIS uses the effective radius retrieved from the near-infrared band. Consequently, the asymmetry parameter (g) is constant for POLDER but increase with particle size for MODIS. However, as the sensitivity of visible reflectance of liquid clouds to particle size is small, the impact of effective radius bias on the POLDER

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COT deviation can be most of the time ignored except for particular scattering angles, such as in rainbow and backward directions (Parol et al., 2000; Buriez et al., 2001). The good statistical relationship between the liquid COTs of the two sensors in Fig. 4d is thus not really improved for the scaled COT in Fig. 7d.

For ice clouds, a recent study from Zhang et al. (2009) has discussed ice COT differences between POLDER and MODIS. Among many potential reasons, the microphysical assumption differences appeared to be the main reason of ice COT differences between the two sensors. The slope obtained in Fig. 4a for a whole year at the global scale (0.74) is close to the value found by Zhang et al. (2009) for only one granule (0.68). The better slope (0.74 compared to 0.68) and correlation coefficient obtained may be partially due to the improvement in the matching of POLDER and MODIS coincident pixels in PM data, but also to the more extensive statistics used for the present studies. Again for scaled optical thicknesses, closer relations between the two sensors are found for ice clouds in Fig. 7 with slope value much closer to one. These statistical results confirm that the bias in ice cloud optical thickness between POLDER and MODIS comes mainly from cloud microphysical assumptions as reported by Zhang et al. (2009). These results further stress out the critical need for a better knowledge of ice clouds asymmetry parameter if we aim at a consistent description of ice cloud properties on global scale for climatological purposes.

4.3 Impact of the observation geometry

In this section, we go a little further to analyze POLDER and MODIS COT and look at the bias associated to the observation geometry already highlighted by many other works (e.g., Loeb and Coakley, 1998; Buriez et al., 2001; Iwabuchi et al., 2002). During the retrieval process, some assumptions have to be made either for simplicity and computational speed or because not enough information content is available from the data. If the cloud model was perfect and geographical samples were sufficient and unbiased, COT values should be almost constant with viewing angles. However, we see in Figs. 8 and 9 where polar graphs of COT are represented, that there are important

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angular variations. Polar graph allows to identify some deviations due to the inappropriate cloud model used for the retrieval. In addition, it can help to establish a reference to determine a suitable weighted average method or to select the direction in order to obtain the best final COT from directional COT. Hence, it is possible to avoid some systematical errors coming from 3-D effects such as, for example, shadow effect. Note that a weighted angular method is already applied to derive the cloud spherical albedo and the optical thickness from the POLDER data (Buriez et al., 2005). In Fig. 8, we present polar graphs of both POLDER-derived COT for overcast liquid clouds and different classes of sun incidences and MODIS-derived COT for all sun incidence angles. In Fig. 9, we present the same figures, for broken liquid clouds. The zero relative azimuth direction corresponds to backward scattering direction. As MODIS observes only in one direction, it has limitations in sampling different relative azimuth angles and figures by class of solar incidence angles are not very informative and thus not presented here. Figures built from POLDER cover more observing geometries thanks to the 16 observing directions.

For overcast clouds, we note that COTs increase with solar zenith angle. This is on the one hand due to geographical sampling of an increase of thicker clouds with latitude (Fig. 2) and on the other hand due to 3-D effects as observed by Loeb and Davies (1996) and reproduced with Monte-Carlo simulation by Loeb et al. (1997). In addition, we see clearly that COT values decrease in forward directions ($\phi = 180^\circ$) for both sensors and more markedly for oblique sun. In case of solar angles between 70° – 80° , derived COT in forward directions is only about 75% of the angular mean value, even lower than 50% of the angular mean value. On the contrary, larger COTs tend to be found in backscattering directions ($\phi = 0^\circ$). This comes, partially from sub-pixel heterogeneity (Loeb and Coakley, 1998; Buriez et al., 2001) and partially from cloud shadowing and illumination that respectively decreases or increases the retrieved COT (Várnai and Marshak, 2002; Iwabuchi and Haysaka, 2002). We also distinguish smaller COTs in the rainbow direction at lower sun, which are associated to both particle size assumption and cloud inhomogeneities (Buriez et al., 2001).

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For broken clouds, we see that MODIS COT decreases with viewing angle (Fig. 9) both in forward and backward directions. This decrease may come from the cloud detection as the MODIS CF increases with the viewing angle (Zeng et al., 2011a) because the pixel size increases towards the edges of the scan corresponding to large viewing angles and also because more cloud sides are being seen. As explained in Sect. 4.1, this may result in a decrease of COT with the viewing angle (70–80% of the angular mean values). For POLDER, we clearly see high values of COT around sun-glint directions where the bias can reach up to 100% of the angular mean values. It points out that the limitation to prevent COT retrieval in sun-glint regions is not conservative enough especially in case of broken clouds (Fig. 9). We also observe in the figure consistent underestimations of COT in backward and rainbow directions, which are about 75% of the angular mean values. These underestimations are mainly due to microphysical issues as the POLDER effective radius is set to a constant value.

5 Conclusions

In this paper, we studied and compared cloud optical thickness (COT) retrieved from two passive sensors in the A-train constellation, POLDER/PARASOL and MODIS/AQUA. The accurate retrieval of this parameter strongly depends on many cloud and environment factors such as cloud cover and thermodynamic phase, cloud particle information. After having reminded the main underlying principles of COT retrieval in the POLDER and the MODIS algorithms, we analyzed global geographical distributions, latitudinal variations and seasonal evolutions of COT. This helped to reveal specific uncertainties in COT retrieval. We also studied pixel-to-pixel relationships of COT between the two sensors, separating clouds by class of phase and cloud fraction determined from a combination of POLDER and MODIS. Our results show that, for overcast clouds with consistent cloud phases, it exists a strong relation between the optical thicknesses of the two sensors. We notice that for liquid clouds, MODIS and POLDER COT have similar values while for ice clouds, the MODIS COT is larger

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than the POLDER one. For overcast clouds with inconsistent phases, the sensor that employs liquid phase model retrieves logically larger COT. In case of broken cloud, we concluded not surprisingly that the COT relationships between the two sensors are much weaker. These results between the two optical thicknesses help to indicate in which conditions the retrieved COT can be considered as suitable.

As the second part of this study, we discussed many uncertainties and impacts of various factors in COT retrieval processes. Among all those factors, the sensor spatial resolution is an important one and is first discussed. Through its impact on cloud detection, the spatial resolution differences can bias cloud cover (Zeng et al., 2011a) and thus COT. Indeed, satellites with lower spatial resolution tend to detect larger cloud cover and yield as a consequence smaller retrieved COT. In addition, the spatial resolution differences lead to different plane-parallel bias due to inhomogeneous clouds. By neglecting the sub-pixel cloud inhomogeneities, satellites with lower spatial resolution will detect smaller optical thickness. Other uncertainties presented are associated to cloud microphysical models used in algorithm. A significant source of uncertainties for COT retrieval lies in the choice of thermodynamic phase, cloud particles size and shape assumptions. Using liquid model instead of ice model leads to an overestimation of COT. The influence of particles size for liquid clouds can be most of the time ignored, however for ice clouds, the differences of cloud particle sizes and shapes conduct to the main differences observed between MODIS and POLDER derived COT. In addition, to study the impact of observation geometry on COT retrieval, we presented angular variations of COT. We clearly see the underestimation of cloud optical thickness in forward directions due to shadow effect especially for the lower sun and also the underestimation of cloud optical thickness from POLDER in rainbow and backward directions due to the cloud microphysics assumption. However, our study outlines the strong interest to measure the angular variation of reflected solar radiation.

This paper provides a good overview on how COT from passive sensors can be biased. It also provides ground for selection of more confident situations when evaluating cloud representation in models or multi-sensor inter-comparison studies aimed

at building climatological records of cloud properties. Several issues concerning ice cloud microphysics and 3-D radiative effects still need to be explored to obtain better COT values, in particular a better knowledge of ice crystal asymmetry factor is critically needed. In the future, comparisons of cloud top reflectance and clear sky reflectance between POLDER, MODIS and CERES will be done in order to assess the impact of COT biases between different instruments in radiative budget computation.

Acknowledgement. This work was supported by the French Space Agency (CNES) and the Région Nord-Pas de Calais. The POLDER level2 and level3 data were processed and distributed by the French ICARE Data Management and Processing Center (CGTD). The MODIS data were obtained from the NASA/GSFC/LADS. The authors are grateful to CNES and NASA for making the POLDER and MODIS data available.



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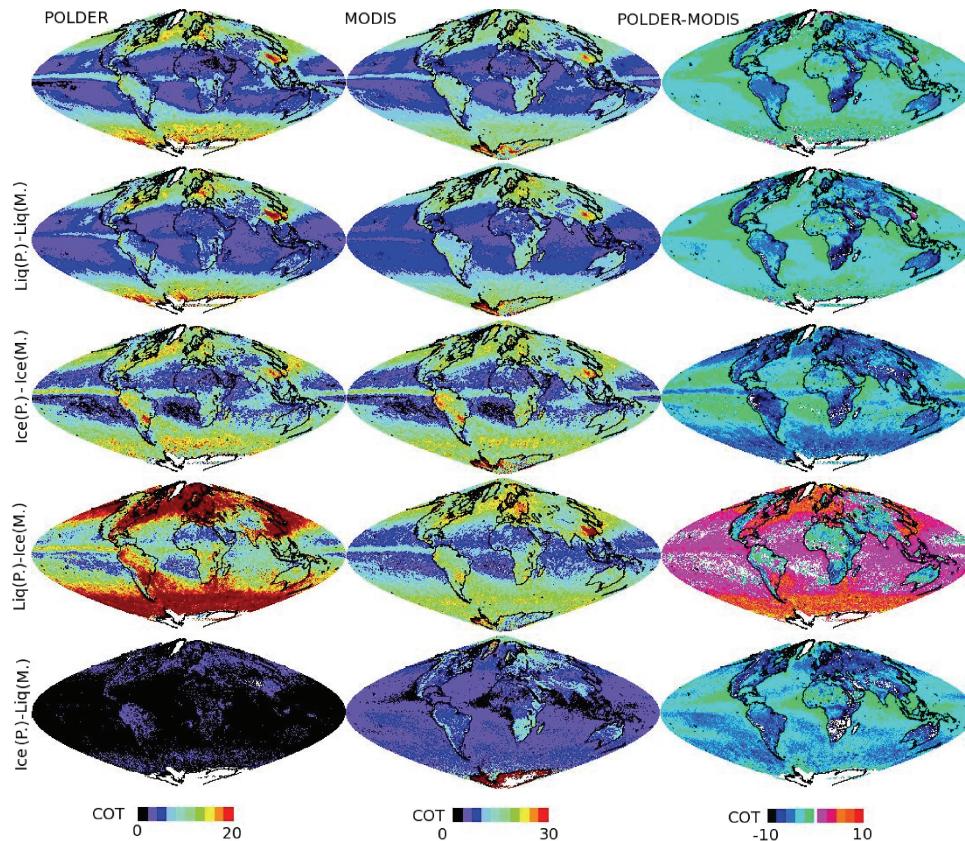


Fig. 1. Geographical distributions of POLDER COT (first column), MODIS COT (second column) and the corresponding COT differences (POLDER-MODIS) (third column) for all the overcast clouds (first line) and clouds separated in 4 different phase classes: POLDER-liquid MODIS-liquid (second line), POLDER-ice MODIS-ice (third line), POLDER-liquid MODIS-ice (fourth line) and POLDER-ice MODIS-liquid (fifth line).

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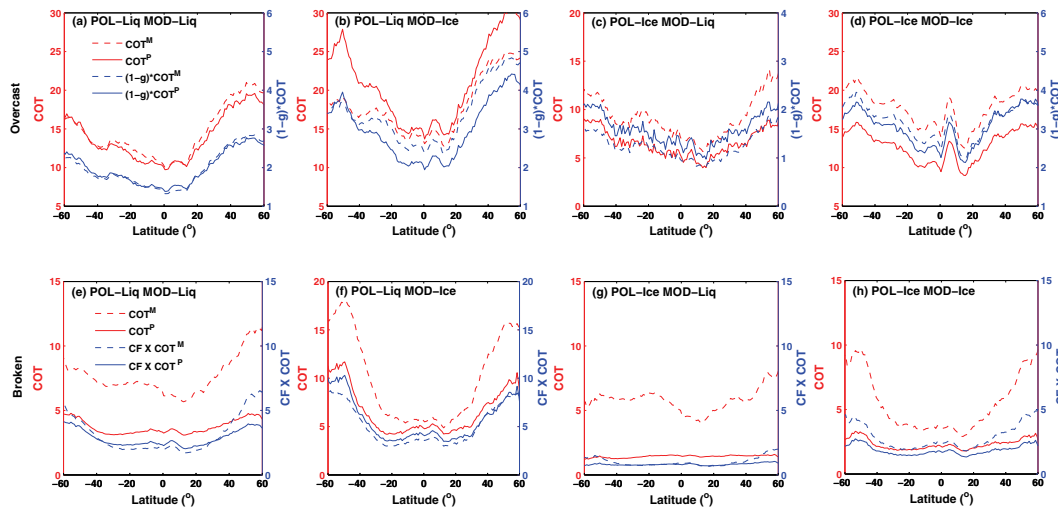


Fig. 2. First line: Latitudinal variations of MODIS and POLDER COT (left axis) and the scaled COT calculated by $(1 - g) \times COT$ (right axis) for 4 different overcast clouds classes determined by cloud combined phase: POLDER-liquid and MODIS-liquid (a), POLDER-liquid and MODIS-ice (b), POLDER-ice and MODIS-liquid (c), POLDER-ice and MODIS-ice (d). Second Line: Latitudinal variations of MODIS and POLDER COT (left axis) and the product $CF \times COT$ (right axis) for 4 different broken clouds classes determined by cloud combined phases: POLDER-liquid and MODIS-liquid (e), POLDER-liquid and MODIS-ice (f), POLDER-ice and MODIS-liquid (g), POLDER-ice and MODIS-ice (h).

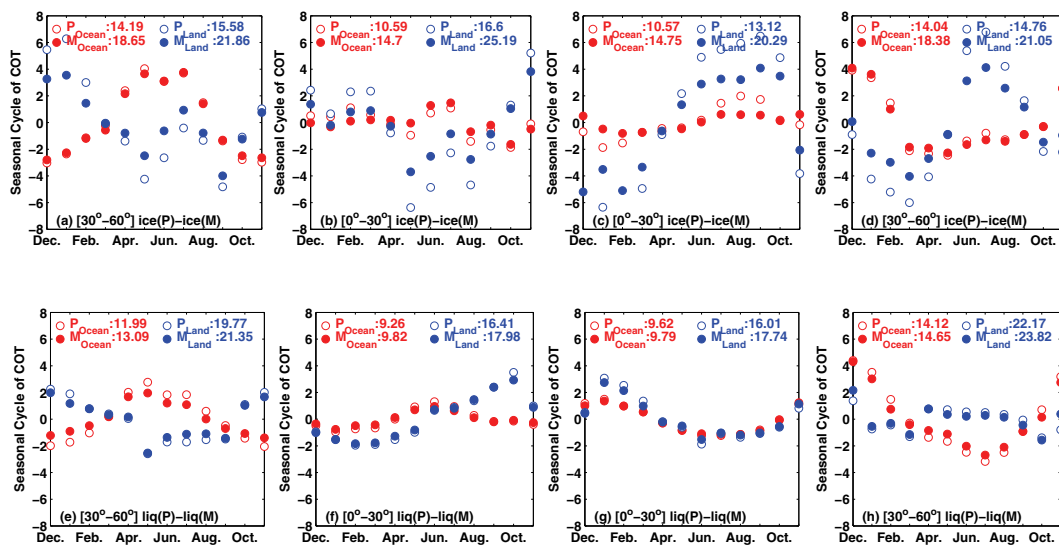


Fig. 3. Seasonal cycles of COT for overcast clouds over South Hemisphere mid-latitude regions (30° – 60° S) (first column) and tropical regions (0° – 30° S) (second column), over North Hemisphere tropical regions (0° – 30° N) (third column) and mid-latitudinal regions (30° – 60° N) (fourth line) with different combined cloud phases: POLDER-ice and MODIS-ice (first line), and POLDER-liquid and MODIS-liquid (second line). Ocean data are in red and land data in blue. The solid circles present data from MODIS and the hollow circles from POLDER. The values in the figures correspond annual mean COT in the regions.

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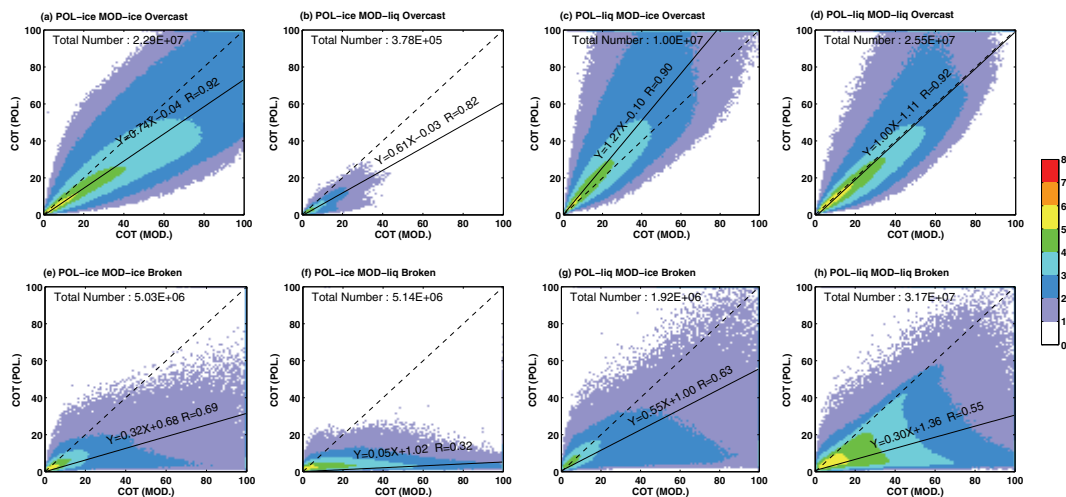


Fig. 4. Pixel-to-pixel comparisons of POLDER and MODIS COT for 4 different overcast and broken clouds separated by cloud combined phase.

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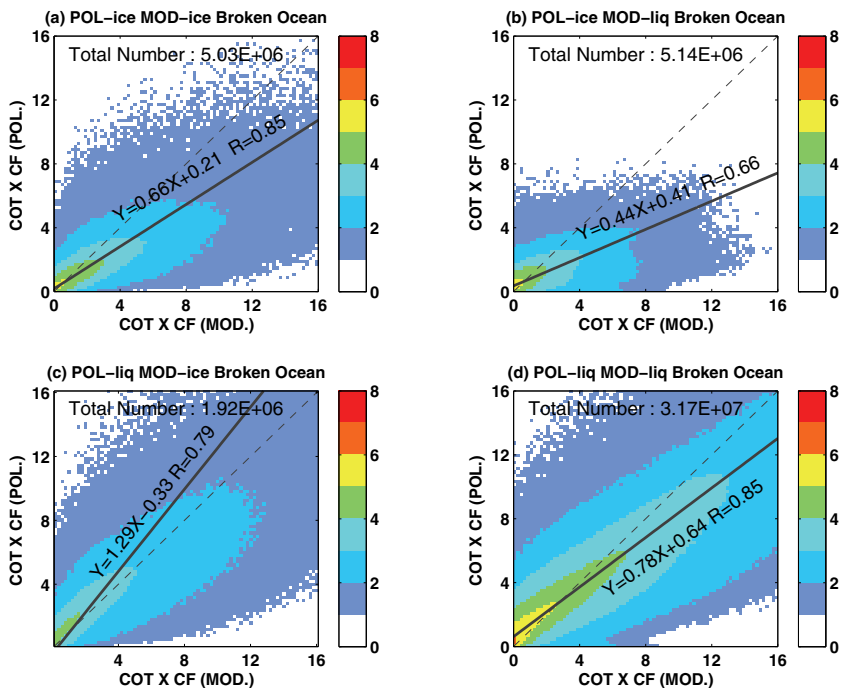


Fig. 5. Pixel-to-pixel comparisons of the product $COT \times CF$ of 4 different broken clouds separated by cloud combined phase.

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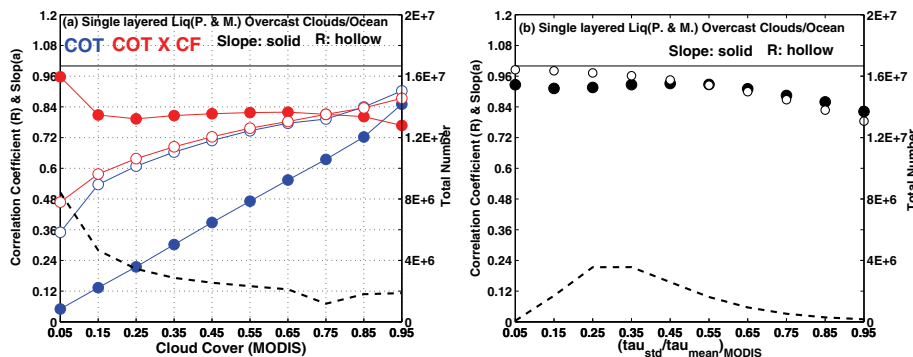


Fig. 6. Slopes and correlation coefficients, R , of the relationship between POLDER and MODIS COT in blue and between POLDER and MODIS COT \times CF in red as a function of MODIS cloud cover **(a)**. Slopes and correlation coefficients of the relationship between POLDER and MODIS COT as a function of cloud inhomogeneity from MODIS **(b)**. Only the liquid clouds over ocean are considered. Dashed curves represent the number of samples in each case.

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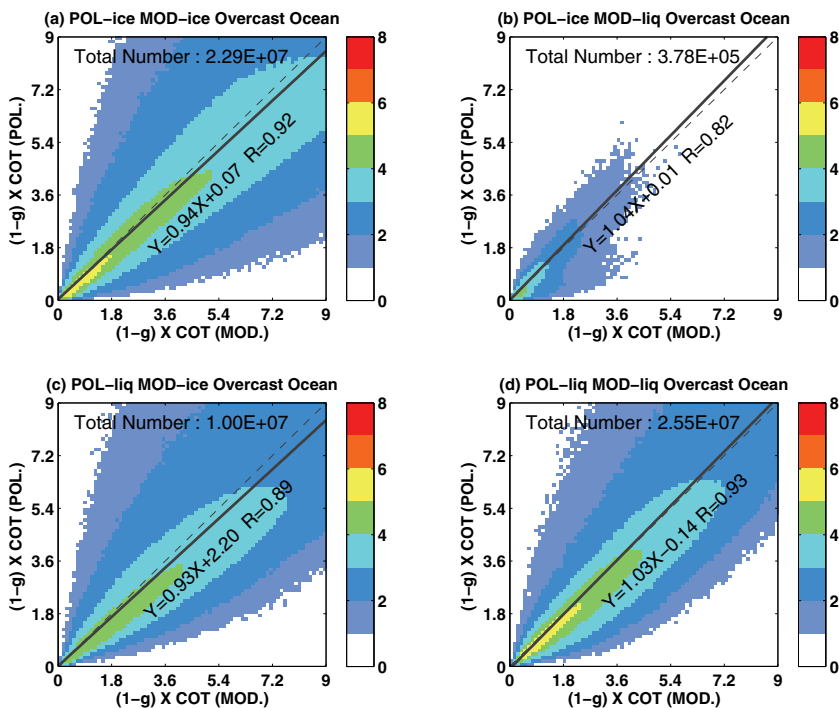


Fig. 7. Pixel-to-pixel comparisons of the product $COT \times (1 - g)$ for the 4 different overcast clouds separated by cloud phase.

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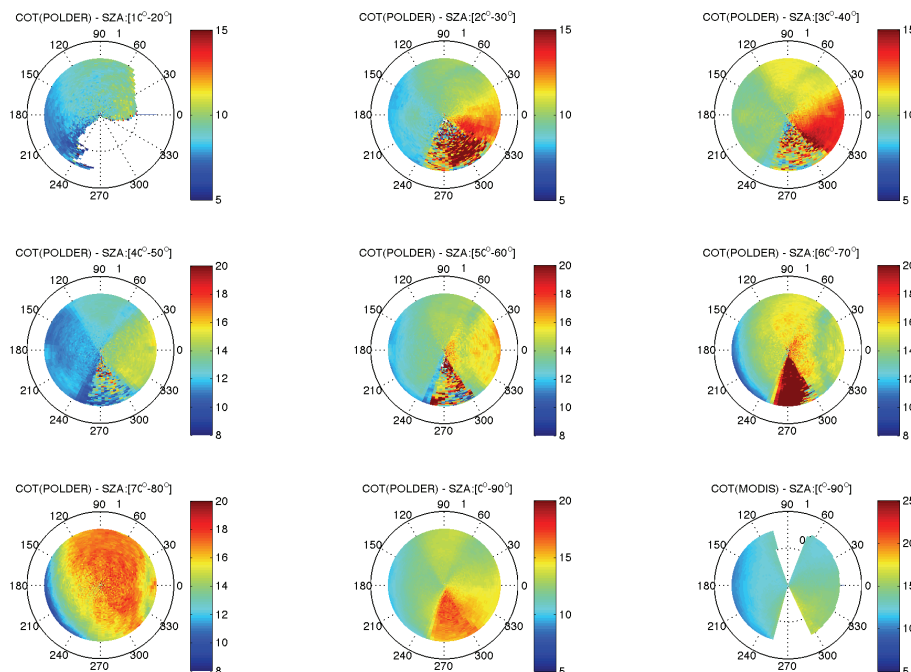


Fig. 8. Polar graphs of POLDER COT for overcast oceanic clouds for different solar incidences and of MODIS COT for all sun incidence angles (lower right corner). 0° corresponds to backscattering direction.

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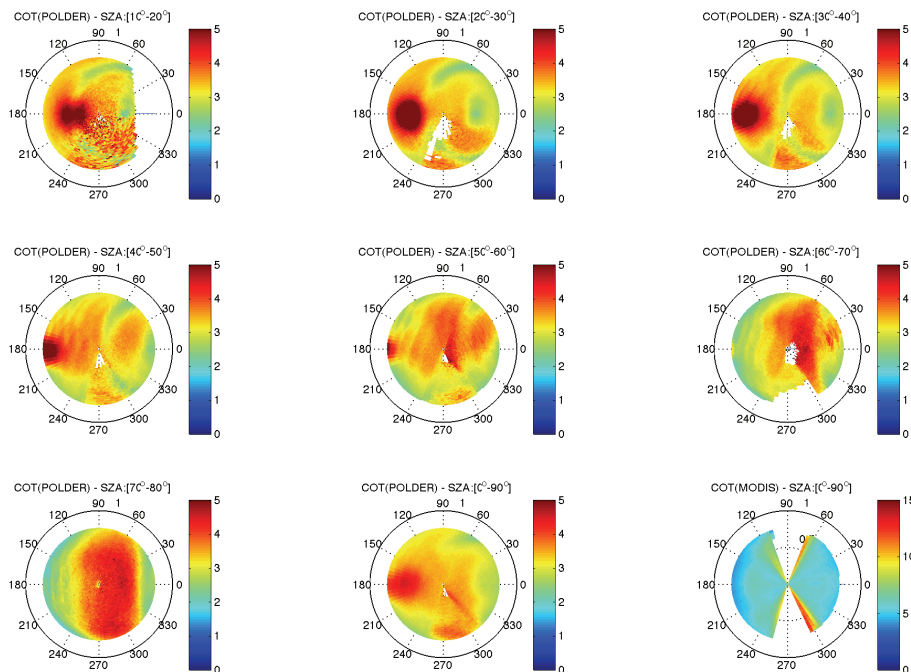


Fig. 9. Polar graphs of POLDER COT for broken oceanic clouds for different solar incidences and of MODIS COT for all sun angles (lower right corner). 0° corresponds to backscattering direction.

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