



Supplement of

Temperature–RH dependent viscosity of organic aerosols from 273 to 303 K: implications for global N₂O₅ uptake

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12 **S1. Flow-cell calibration and particle conditioning time determination**

13 In both bead-mobility and poke-and-flow experiments, the flow-cell was calibrated for relative humidity (RH)
14 using the deliquescence points of three salts, including K_2CO_3 , NaCl, and $(NH_4)_2SO_4$ (Winston and Bates, 1960).
15 The experimentally determined deliquescence RH values at each measurement temperature of 273, 283, 293, and
16 303 K were compared with the literature. The RH measurement uncertainty was quantified as the temperature-
17 dependent calibration root mean square error of the RH sensor. The RH accuracy was within $\pm 2.0\%$ at 293 K, and
18 the temperature was controlled within ± 1 K of the target.

19 Prior to each experiment, particles deposited on a hydrophobic substrate were conditioned at the target RH of
20 the carrier gas to allow sufficient equilibration with the surrounding gases. Conditioning times were adjusted
21 depending on the experimental temperature and technique. To evaluate whether the samples reached near-
22 equilibrium with the target RH, we applied a method previously reported in the literature (Evoy et al., 2021;
23 Maclean et al., 2021; Smith et al., 2021; Kiland et al., 2023). This method involves comparing experimental
24 conditioning time to the characteristic mixing time of water within organic aerosol (OA), τ_{mix,H_2O} , calculated:

$$\tau_{mix,H_2O} = \frac{d_p^2}{4\pi^2 D_{H_2O}(T, RH)} \quad (S1.1)$$

25 Here, d_p represents the particle diameter. $D_{H_2O}(T, RH)$ is the RH- and temperature-dependent diffusion coefficient
26 of water in OA. The value of D_{H_2O} was calculated using the fractional Stokes–Einstein equation, which accounts
27 for the link between viscosity and diffusion in cases where the diffusing species are comparable in size to, or smaller
28 the surrounding matrix molecules (Evoy, 2020):

$$D_{H_2O}(T, RH) = D_{H_2O}^\circ(T) \times \left(\frac{\eta_{H_2O}^\circ(T)}{\eta(T, RH)} \right)^\xi \quad (S1.2)$$

29 In this equation, $D_{H_2O}^\circ(T)$ is the self-diffusion of water calculated with the Stokes-Einstein equation at 293 K (2.15
30 $\times 10^{-9}$ $m^2 s^{-1}$). $\eta_{H_2O}^\circ(T)$ is the viscosity of pure water obtained from literature with values of 1.78×10^{-3} Pa·s at
31 273 K, 1.31×10^{-3} Pa·s at 283 K, 1×10^{-3} Pa·s at 293 K, and 8.14×10^{-4} Pa·s at 303 K (Weight, 2019). $\eta(T, RH)$
32 is the measured viscosity of sucrose-H₂O droplets at the corresponding RH and temperature. ξ is the fractional
33 exponent, using Eq. (S1.3), which accounts for the relative size of the diffusing molecule and the matrix:

$$\xi = 1 - \left[A \times \exp \left(-B \frac{r_{diff}}{r_{matrix}} \right) \right] \quad (S1.3)$$

34 where coefficient values of $A = 0.73$ and $B = 1.79$ (Evoy et al., 2020), r_{diff} hydrodynamic radius of water of 0.1 nm
35 (Price et al., 2016). The matrix radius, r_{matrix} , of the sucrose molecule was calculated under the assumption of
36 spherical geometry, using $342.30 \text{ g mol}^{-1}$ and a density of 1.67 g cm^{-3} (Haynes, 2016), resulting in $r_{matrix} = 0.44 \text{ nm}$
37 and $\xi = 0.51$. Conditioning durations for the other temperatures (273, 283, 293, and 303 K) are summarized in
38 Table S1.

39 In addition to water, we also estimated the mixing times of organic molecules ($\tau_{mix,org}$) to assess potential
40 diffusion limitations within the particle matrix itself. This is particularly relevant under conditions where semi-
41 solid or glassy states may impede equilibration. The diffusion coefficient of organics (D_{org}) was calculated using
42 the classical Stokes–Einstein equation (Eq. S1.4) (Evoy et al., 2019; Evoy et al., 2020):

$$D_{org}(T, RH) = \frac{kT}{6\pi\eta_{org}(RH, T)R_{diff}} \quad (\text{S1.4})$$

43 where k is the Boltzmann constant, η_{org} is the temperature- and RH-dependent viscosity of OA, and R_{diff} is the
44 radius of the diffusing molecules. We used 0.4 nm for R_{diff} , consistent with typical organic molecules of similar
45 molecular weight and density. Using the estimated diffusion coefficients, the mixing time of organics was
46 calculated as:

$$\tau_{mix,org} = \frac{d_p^2}{4\pi^2 D_{org}(T, RH)} \quad (\text{S1.5})$$

47
48 A particle diameter of 200 nm was used in all calculations, consistent with the typical size of secondary organic
49 aerosol (SOA) particles observed in the atmosphere (Pöschl et al., 2010; Riipinen et al., 2011). The $\tau_{mix,org}$ value
50 represents the time required for the concentration of the diffusing species at the particle’s center to attain $1/e$ of the
51 equilibrium concentration.

52 The diffusivity coefficient of N_2O_5 in OA, which is needed in calculations of N_2O_5 uptake coefficient, is also
53 determined using the fractional Stokes-Einstein equation (S1.2 and S1.3). In Eq. (S1.3), the hydrodynamic radius
54 of N_2O_5 is set to be 0.25 nm (Grzanic et al., 2015).

55

56 **S2. Viscosity parameterization of sucrose-H₂O droplets**

57 We parameterized the viscosity of sucrose-H₂O droplets as a function of RH and temperature, following the
58 approach proposed by Kiland et al. (2023). Experimental data at 293 ± 1 K was fitted using a mole-fraction-based
59 Arrhenius mixing rule (Eq. S2.1)

$$\log(\eta(RH, 293K)) = x_{org} \log_{10}(\eta_{org,dry}) + (1 - x_{org}) \log(\eta^{\circ}_{H_2O}) \quad (S2.1)$$

60 where $\eta(RH, 293K)$ is the viscosity of the mixture at room temperature, $\eta_{org,dry}$ is the viscosity of the organic
61 component at 0% RH (1×10^{12} Pa·s), and $\eta^{\circ}_{H_2O}$ is the viscosity of pure water (1×10^{-3} Pa·s).

62 The mole fraction of organic components, x_{org} , was calculated from the mass fraction using:

$$x_{org} = \frac{\frac{w_{org}}{M_{org}}}{\frac{w_{org}}{M_{org}} + \frac{1 - w_{org}}{M_{H_2O}}} \quad (S2.2)$$

63 where w_{org} is the mass fraction of the organic component, and M_{org} and M_{H_2O} are the molecular weights of sucrose
64 and water, respectively. w_{org} was determined from the water activity ($a_w = RH/100$) using a mass-based
65 hygroscopicity parameter κ , as shown in Eq. (S2.3).

$$w_{org} = \left(1 + \kappa \left(\frac{a_w}{1 - a_w} \right) \right)^{-1} \quad (S2.3)$$

66
67 Fitting the experimental data at 293 K yielded the value of $\kappa = 0.061 \pm 0.0023$. As shown in Fig. S8, the model
68 provides a good fit to the experimental data, supporting the robustness of the parameterization. Petters and
69 Kreidenweis (2007) reported κ values between 0.01 and 0.5 for diverse organic compounds, highlighting variability
70 in hygroscopicity with molecular structure.

71
72

73 **S3. Determination of D_f using global fitting of viscosity data**

74 To determine a single value of D_f , we performed a global fit of all viscosity measurements to Eq. (1). First, we
75 compiled all viscosity data points measured at different combinations of RH and Temperature. Second, for each
76 data point, we calculated $T_o(RH)$ using the RH-dependent parameterization described in Sect. S2. Third, we fitted
77 (Eq. 1) simultaneously to the entire dataset using a nonlinear least-squares routine, minimizing the sum of squared
78 differences between measured and modelled $\ln \eta(RH, T)$. In this global fit, D_f was treated as the single free
79 parameter (while η_∞ was fixed), so that one best-fit value of D_f characterizes the overall temperature sensitivity of
80 viscosity across all RH conditions.

81 The fragility parameter we derived is indicative of the unique molecular interactions and hydration properties
82 found in sucrose-rich particles. These particles are characterized by sharp changes in viscosity with temperature
83 near their glass transition point. This accounts for our slightly higher fragility value compared to the lower limit
84 reported by Derieux et al. (2018), who studied a wider variety of organic aerosol surrogates with more complex
85 functional groups, resulting in a broader range of thermodynamic behaviors. The difference in fragility values does
86 not represent a contradiction but rather highlights the influence of composition. Systems dominated by sucrose
87 typically show greater sensitivity to temperature variations near T_g than mixtures containing many functional
88 groups, leading to higher measured fragility.

89 To assess sensitivity to uncertainty in the fragility parameter, calculations using $D_f=13$ and 10 were compared
90 (Fig. S12). The predicted latitude–altitude patterns of viscosity and $\tau_{mix,org}$ of sucrose-H₂O droplets were broadly
91 similar, indicating that the large-scale phase-state structure was relatively robust to the assumed D_f . Nevertheless,
92 these results should be interpreted in light of the study’s limitations, as the viscosity parameterization is extrapolated
93 beyond 273–303 K, where sensitivities to RH and meteorological averaging increase.

94 **S4. Collection of global ambient zonal RH and temperature**

95 To evaluate the tropospheric distribution of predicted sucrose-H₂O droplet phase state and mixing times, we utilized
96 monthly mean reanalysis data for temperature and RH from the Copernicus Climate Data Store
97 (<https://cds.climate.copernicus.eu/>) (Hersbach et al., 2023). Monthly mean temperature and RH on pressure levels
98 were obtained from the ERA5 reanalysis (Copernicus Climate Data Store) for 2020–2024. Latitude–altitude
99 profiles were constructed by averaging across all longitudes at each latitude and pressure level. Pressure levels were
100 converted to geometric altitude using the barometric (hypsometric) relationship (Eq. S4.1), and the resulting zonal-
101 mean profiles are shown in Fig. S9. The dataset provides global coverage from 90°N to 90°S and 180°W to 180°E,
102 with temporal resolution spanning January 2020 to December 2024. To convert pressure-level data to geometric
103 altitude, Eq. (S4.1) represents the hypsometric relationship assuming a constant temperature lapse rate. This
104 formulation is widely used to visualize the pressure-level meteorological data on an altitude coordinate and does
105 not alter the underlying temperature or RH values (Maclean et al., 2021; Kiland et al., 2023). In this framework,
106 the height (h) associated with each pressure level was then calculated using the given Eq. (S4.1), the hypsometric
107 form of the barometric equation, which relates the vertical spacing between the pressure surfaces to the mean
108 temperature of the air column between them.

$$h = \frac{1 - \left(\frac{P}{P_o}\right)^{\frac{R\lambda}{M_{air}}}}{\lambda/T_o} \quad (\text{S4.1})$$

109 where P is the pressure at the given level (1000 hPa – 100 hPa), p_o is the pressure at sea level (101325 Pa), R is the
110 gas constant (8.314 J mol⁻¹ K⁻¹), λ is the temperature lapse rate (6.5 K km⁻¹), M_{air} is the molecular mass of the air
111 (28.97 g mol⁻¹), and T_o is the mean temperature at the surface (288.15 K). To construct latitude–altitude profiles,
112 the monthly RH and temperature values were averaged across all longitudes at each latitude and pressure level.

113 **S5. Detailed formulation of the N₂O₅ uptake calculation**

114 To calculate the N₂O₅ uptake coefficient on sucrose-H₂O droplets ($\gamma_{N_2O_5}$), we used a resistor-model framework
 115 following Ammann et al. (2013) and Gržinić et al. (2015), identical to the formulation employed in our previous
 116 work (Song et al., 2025). The overall uptake coefficient is given by Eq. (S5.1).

$$\frac{1}{\gamma_{N_2O_5}} = \frac{1}{\alpha_s} + \frac{1}{\Gamma_s + \left(\frac{1}{\Gamma_{sb}} + \frac{1}{\Gamma_b}\right)^{-1}} \quad (\text{S5.1})$$

117 where α_s is the accommodation coefficient, Γ_s represents the limiting uptake coefficient for the surface reaction,
 118 Γ_{sb} accounts for surface to bulk transport, and Γ_b describes limitations associated with bulk diffusion and chemical
 119 reaction within the particle. For organic aerosols exhibiting high viscosity, previous studies have shown that Γ_b
 120 typically dominates the overall resistance, whereas surface-related terms contribute comparatively weakly to the
 121 humidity dependence of $\gamma_{N_2O_5}$.

122 The bulk resistance term, Γ_b , is calculated using a reacto-diffusive formulation (Eq. S5.2):

$$\frac{1}{\Gamma_b} = \frac{\omega}{4HRT\sqrt{D_{N_2O_5}k^l}} \left(\coth q - \frac{1}{q}\right)^{-1} \quad (\text{S5.2})$$

123 where ω is the mean thermal velocity of N₂O₅, H is the Henry's law constant, R is the ideal gas constant, T is the
 124 temperature, $D_{N_2O_5}$ is the diffusion coefficient of N₂O₅ in the particle phase, and k^l is the first-order hydrolysis rate
 125 constant. The reacto-diffusive parameter (q) is defined as r/l , where r is the particle radius, and l is the reacto-
 126 diffusive length given by $\sqrt{D_{N_2O_5}/k^l}$. This formulation explicitly links aerosol viscosity to N₂O₅ uptake through
 127 its control on $D_{N_2O_5}$ and thus on the penetration depth of N₂O₅ into the particle prior to reaction.

128 The surface accommodation coefficient (α_s) was set to 1, such that surface accommodation does not limit uptake
 129 and Γ_s was set to 2.5×10^{-4} based on previous simulations and observations at low RH conditions (Grzinić et al.,
 130 2015). This value shows a typical, experimentally supported lower bound for surface-limited uptake and contributes
 131 substantially only when bulk diffusion is suppressed, that is, under conditions of high viscosity. At moderate to
 132 high RH, where viscosity is low and diffusion is rapid, the surface term becomes negligible relative to bulk
 133 resistance. The bulk accommodation coefficient was set to $\alpha_b = 0.035$ following the recommendation of Ammann
 134 et al. (2013), representing the effective probability of N₂O₅ entering the particle phase prior to hydrolysis.

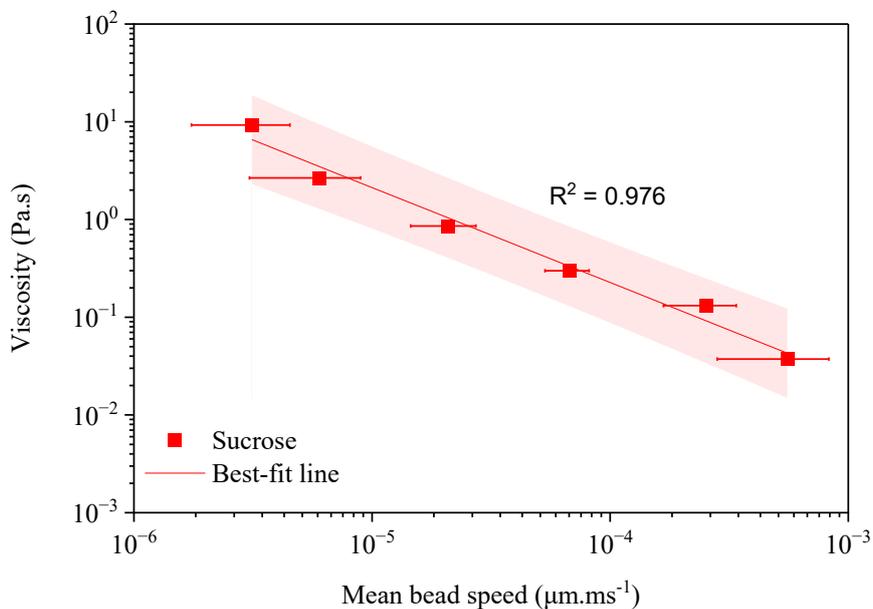
135 The diffusion coefficient of N₂O₅ in the particle phase was determined from viscosity using a fractional Stokes–
 136 Einstein relationship. A fractional exponent of 0.74 was adopted, calculated from literature-derived parameters
 137 using the equations described in Sect. S1. In this calculation, the molecular radius of N₂O₅ was taken as 0.25 nm

138 (Grzanic et al., 2015), and the matrix radius was set to 0.44 nm, corresponding to the effective molecular size of
139 sucrose.

140 The intrinsic first-order hydrolysis rate constant (k') was set to 4.1 ns, based on recent molecular dynamics
141 simulations of N_2O_5 hydrolysis in bulk water (Galib and Limmer, 2021). Although this value was derived under
142 aqueous conditions at room temperature, it provides a physically reasonable estimate of intrinsic chemical reactivity
143 once N_2O_5 is solvated. Using the viscosity-derived diffusion coefficients, the reacto-diffusive length and the
144 corresponding reacto-diffusive parameter were calculated and subsequently applied in Eqs. (S5.1) and (S5.2) to
145 determine $\gamma_{N_2O_5}$ across the tropospheric temperature and RH.

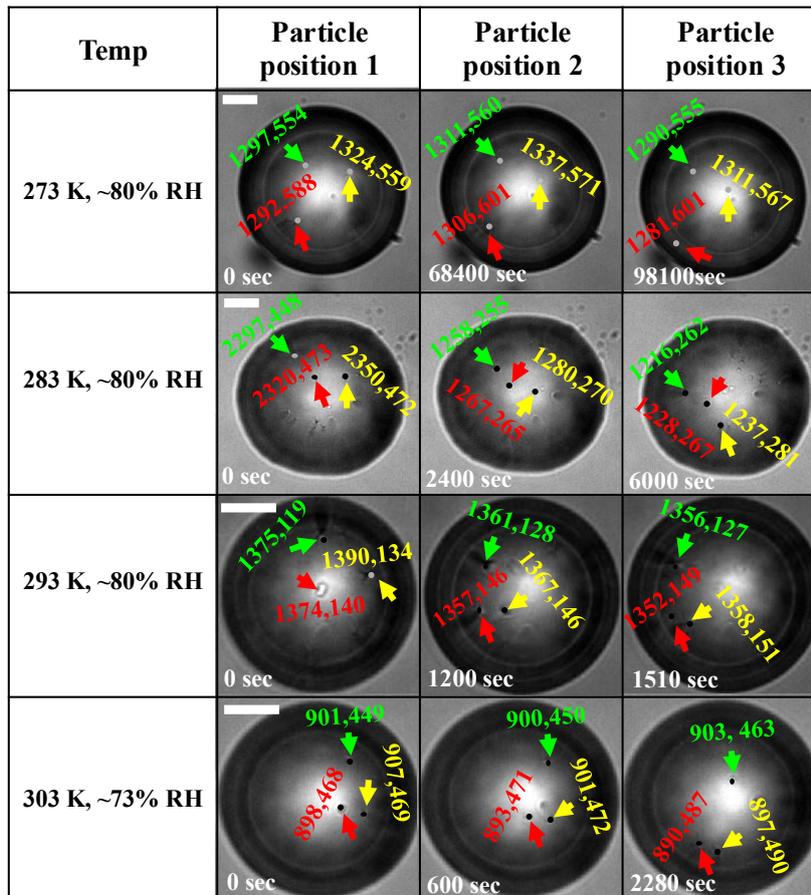
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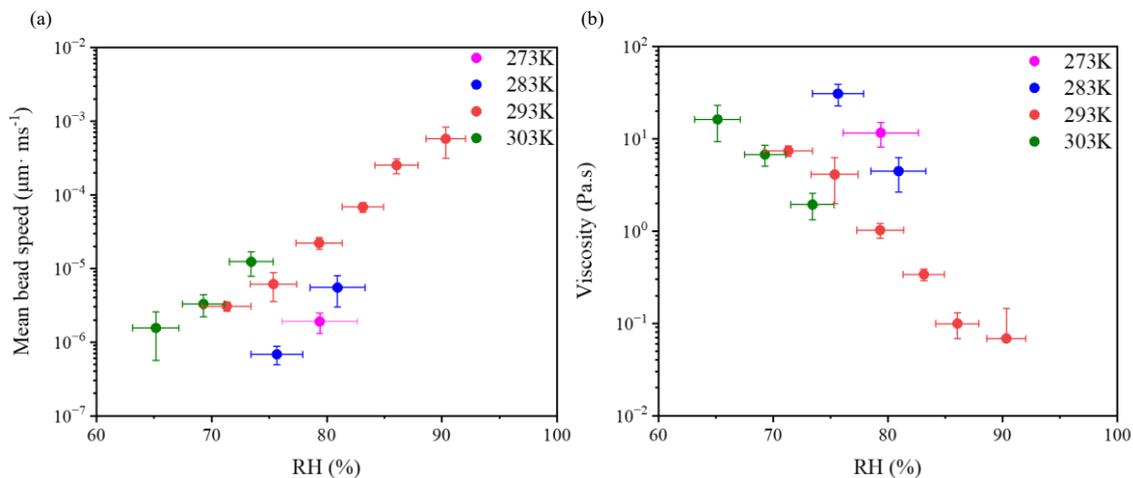
149 **Figure S1. Calibration line illustrating the relationship between mean bead speeds and viscosities of sucrose-H₂O**
150 **droplets at varying relative humidity (RH) levels. A linear regression, shown by the red solid line, fits the data with the**
151 **equation: $viscosity = 0.00003 \times (mean\ bead\ speed)^{-0.971}$. The pink area denotes 95% prediction bands of fitting**
152 **to the data in this study. The uncertainty in mean bead speed along the *x*-axis is calculated from standardization of 2–5**
153 **beads within 3–5 particles for each RH value.**



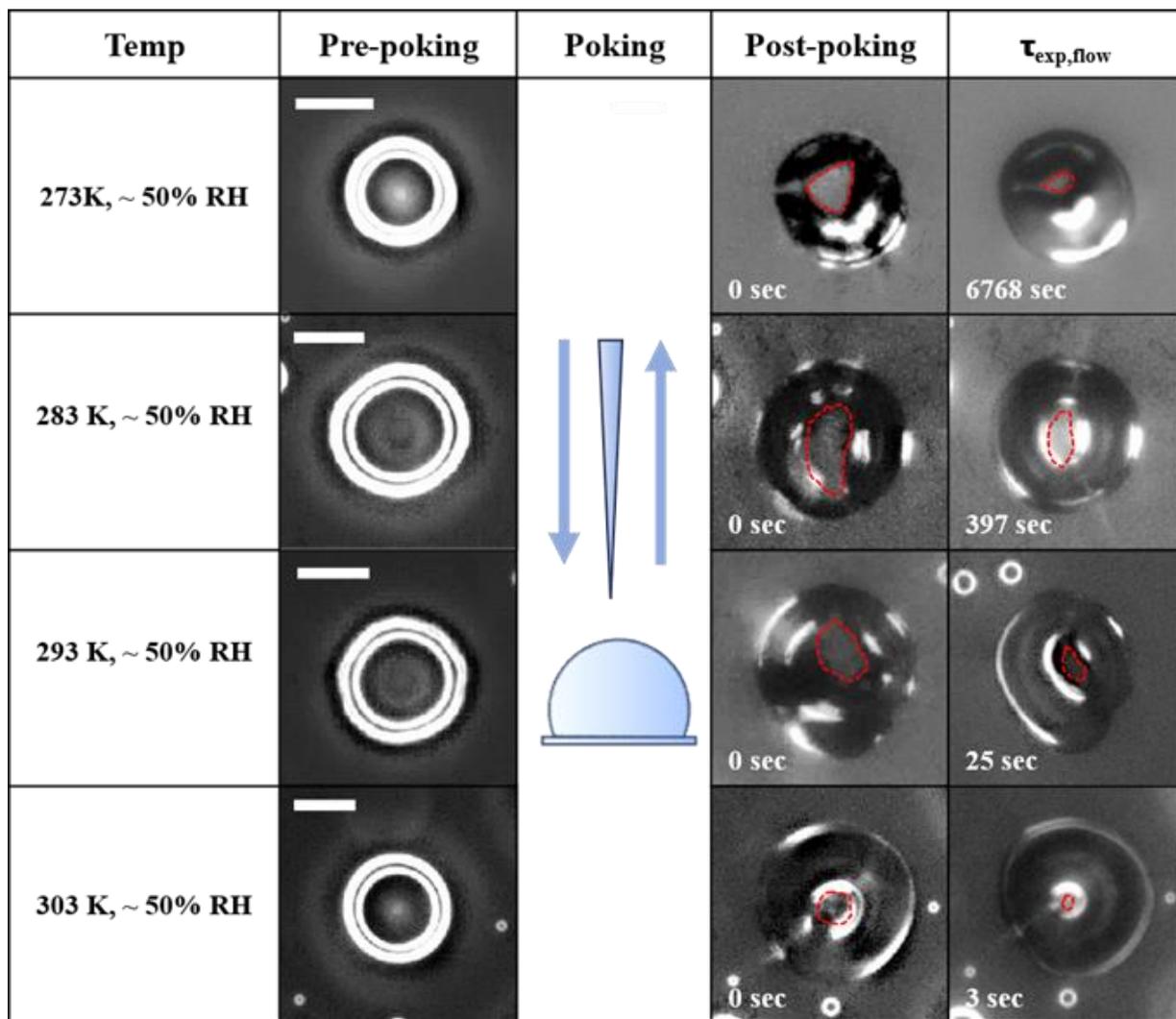
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155 **Figure S2. Optical images of sucrose-H₂O droplets during a typical bead-mobility experiment at different temperatures.**
 156 **Three labeled beads with tracked x and y coordinates were used to determine average bead speeds using ImageJ**
 157 **software. The size of the scale bar is 20 μm.**

158



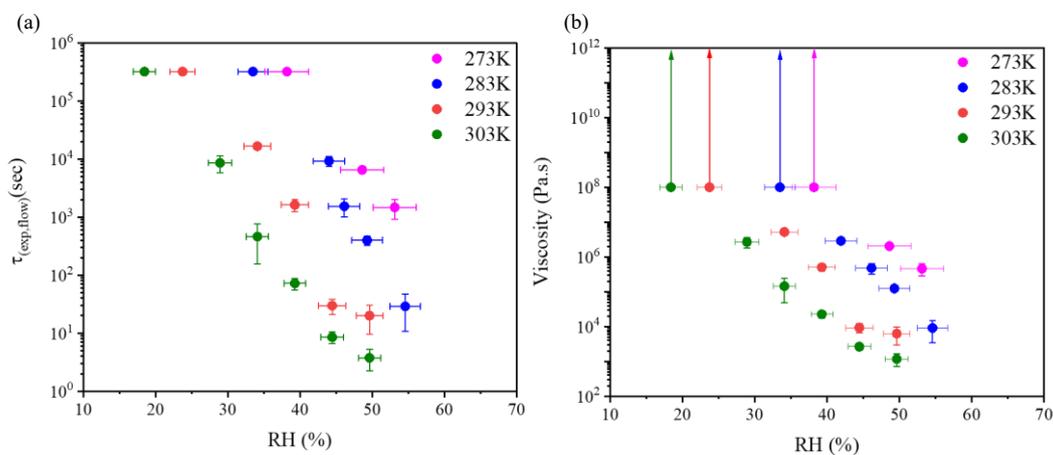
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 161 **Figure S3. (a) Mean bead speed and (b) resulting viscosities from bead-mobility experiments for sucrose-H₂O droplets**
 162 **as a function of temperature and relative humidity (RH). The x-axis error bars represent the RH range in a given**
 163 **experiment and the uncertainty in RH measurements. The y-axis error bars indicate the standard deviation of the**
 164 **measured bead speeds and viscosity calculated from 3–5 beads across 2–5 particles at each RH level.**



166

167 Figure S4. Optical images for experimental flow times ($\tau_{exp,flow}$) during poke-and-flow experiments at different
 168 temperatures. The white scale bar indicates 20 μm . The red dotted lines indicate the size of the cavity measured at the
 169 corresponding time after poking.

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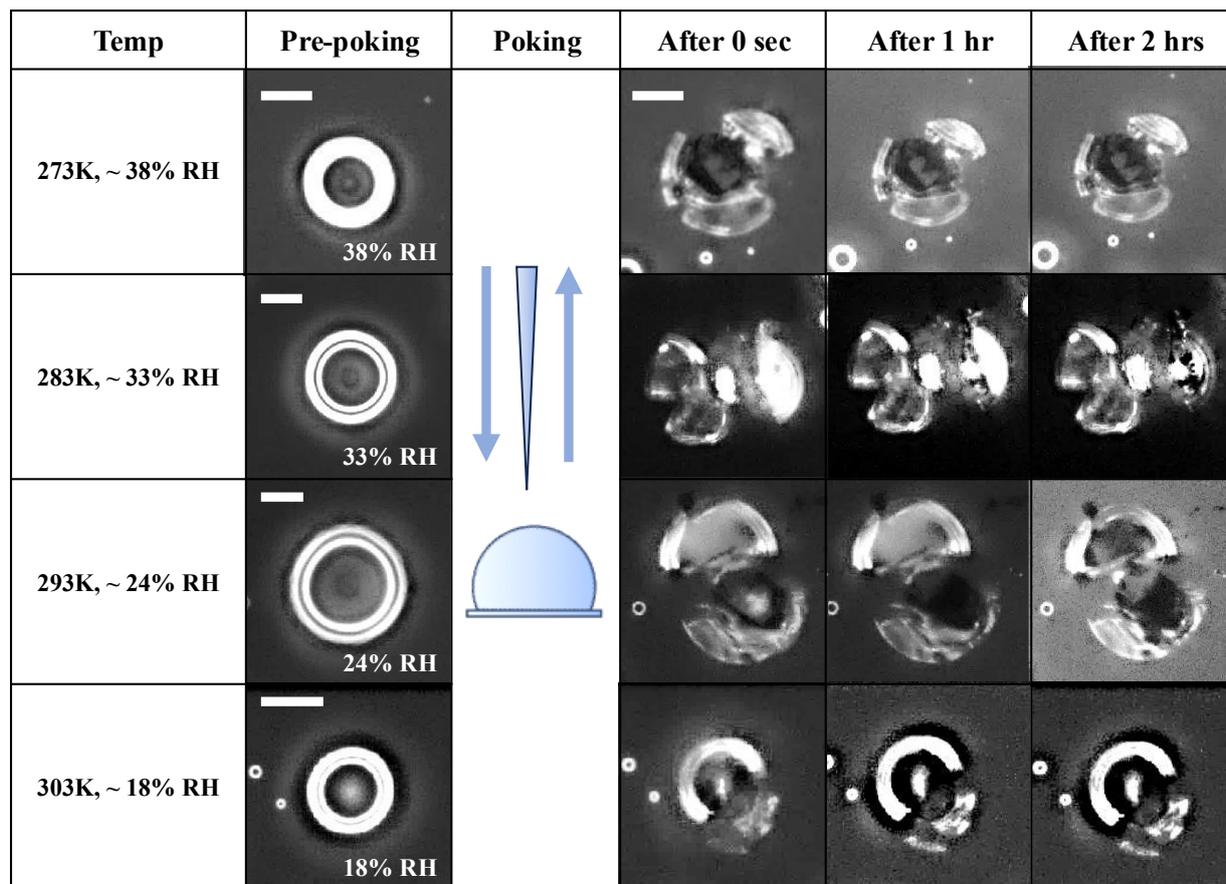


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172 **Figure S5. (a) Mean experimental flow time ($\tau_{exp, flow}$) from poke-and-flow experiments for sucrose-H₂O droplets as a**
173 **function of temperature and relative humidity (RH). The x-axis error bars represent the RH range in a given experiment**
174 **and the uncertainty in RH measurements. The y-error bars indicate the standard deviation of the measured $\tau_{exp, flow}$**
175 **calculated from measurements of 3–4 particles at each RH level. (b) Resulted viscosities from the $\tau_{exp, flow}$ using the**
176 **equation proposed by Sellier et al. (2015). The y-error bars indicate the standard deviation of the measured viscosity.**

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178



179

180 **Figure S6. Optical images of sucrose-H₂O droplets during poke-and-flow experiments at different temperatures.**

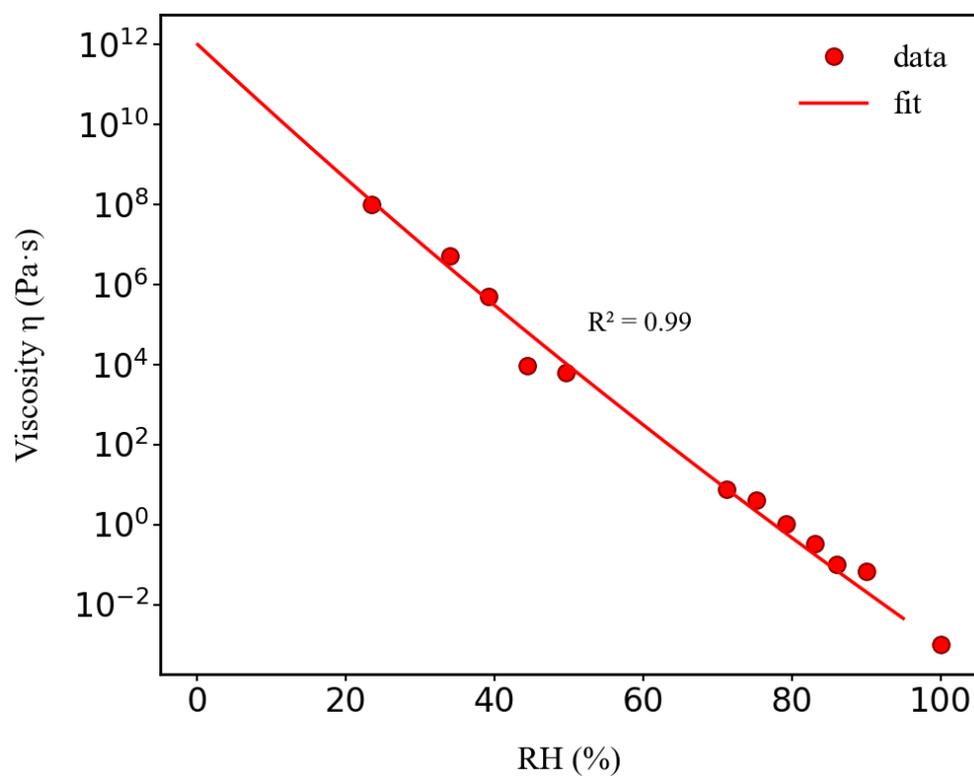
181 **Observations were made when particles cracked at a certain relative humidity (RH), and they were then monitored for**

182 **longer than 2 hours, with no evidence of flow restoration was detected. RH was consistently regulated throughout the**

183 **pre-poking, poking, and post-poking stages. White scale bars represent 20 μm.**

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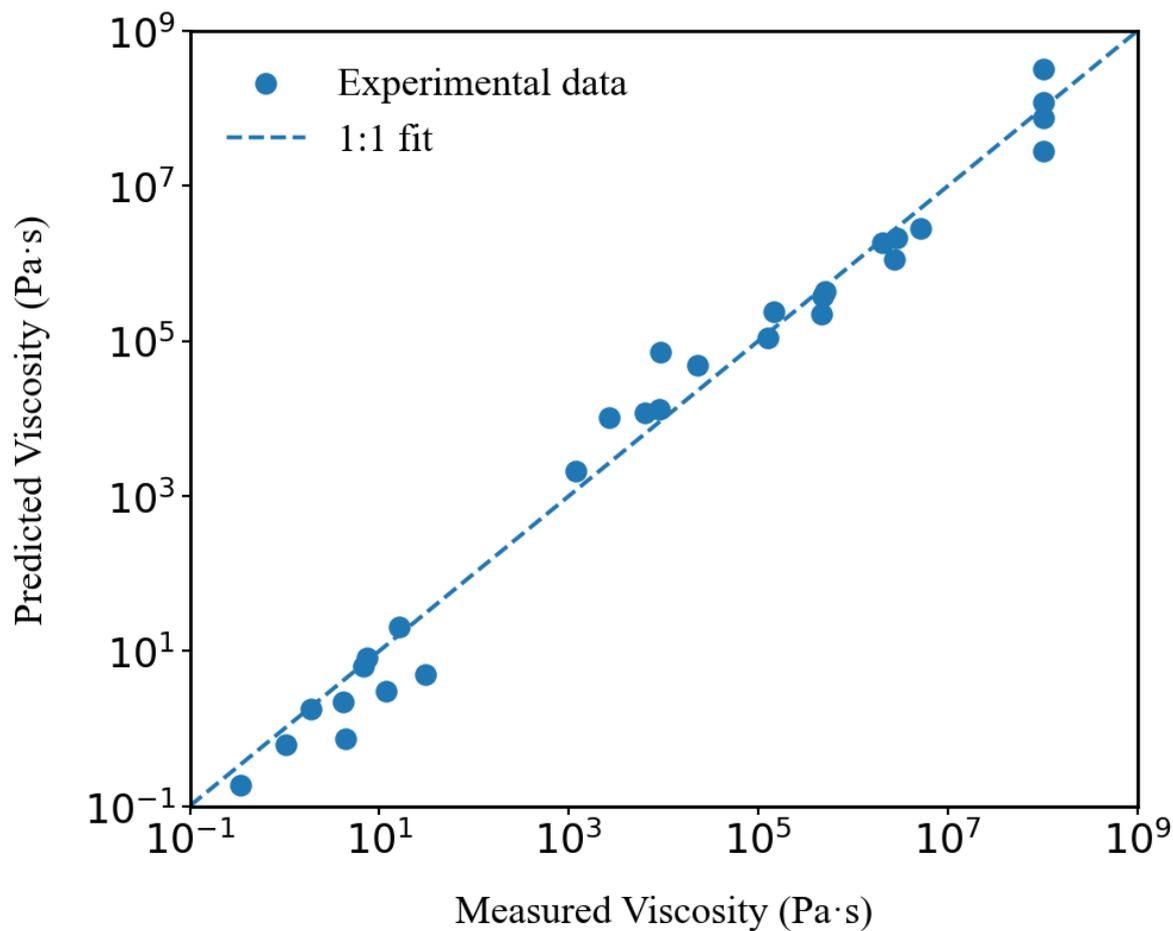
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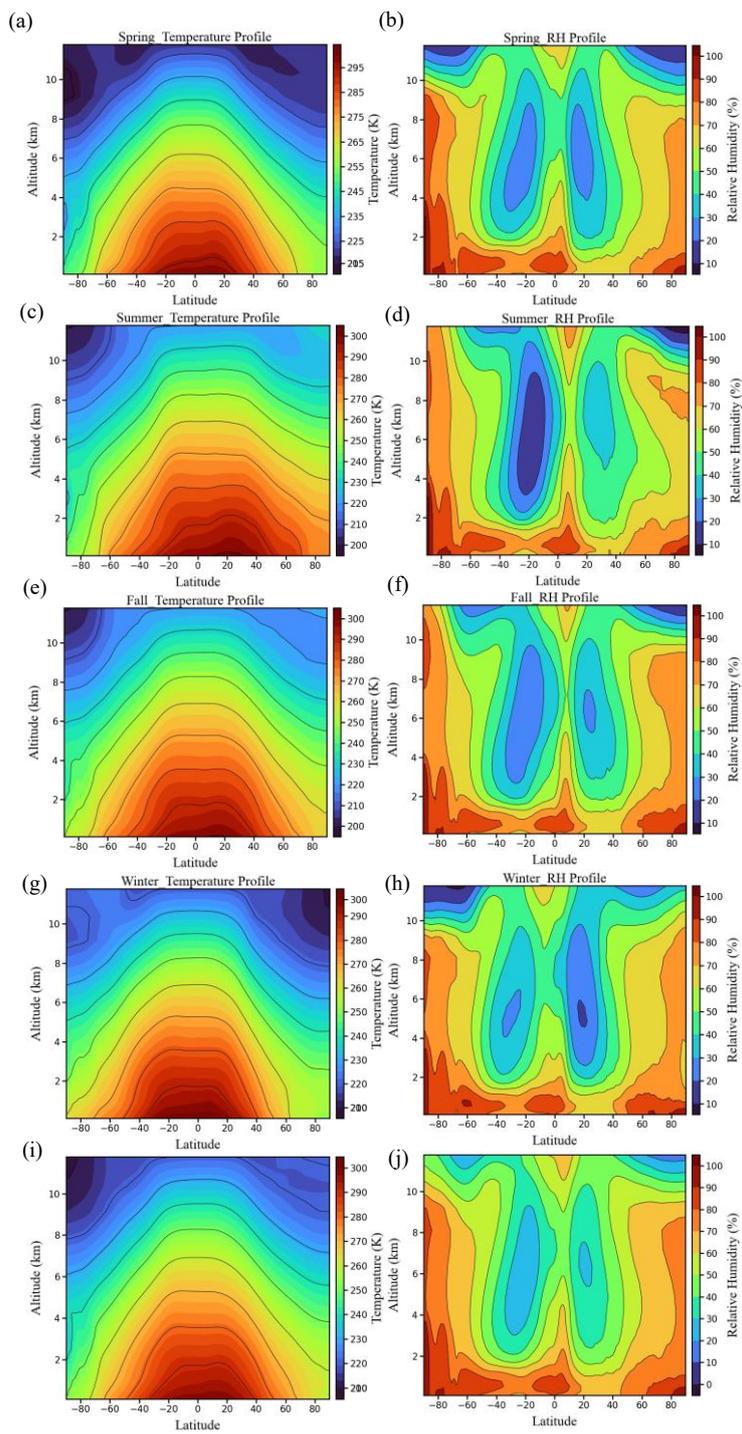
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187 **Figure S7. The viscosity of sucrose-H₂O droplets as a function of RH at room temperature. The solid line is a mole-**
188 **fraction-based Arrhenius mixing rule fit to the viscosity data (Eq S2.1), which yields a hygroscopicity parameter, $\kappa =$**
189 **0.061 ± 0.0023 . The point at 100% RH represents the viscosity of water at room temperature.**

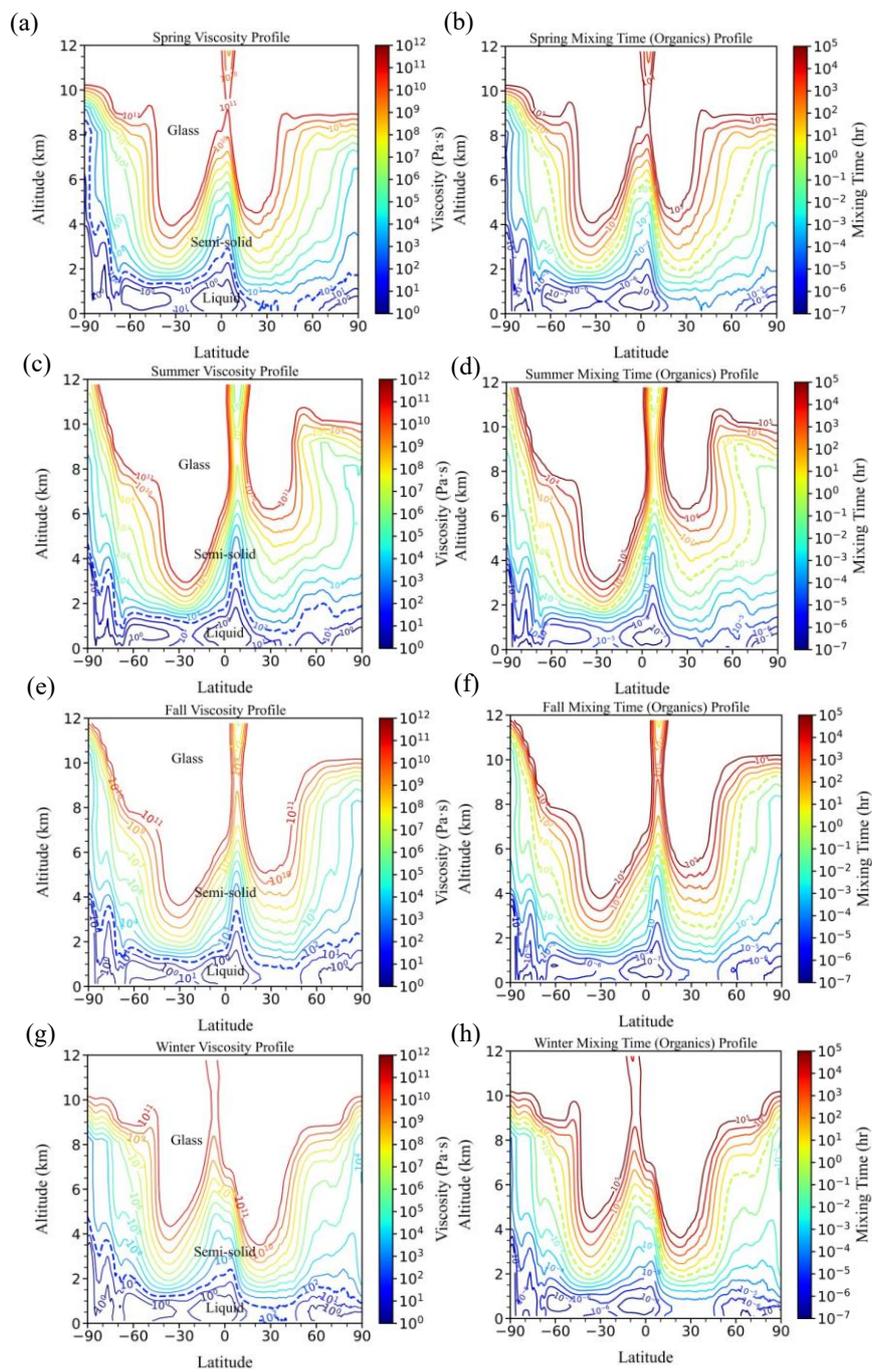
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 192 **Figure S8. Comparison of experimental and predicted viscosity values for sucrose-H₂O droplets. The experimental**
 193 **values (blue circles) were measured between 273 and 303 K and ~20 to ~90% RH, whereas predicted viscosities (dashed**
 194 **line) were calculated using the VFT equation. The fitting procedure yields the fragility parameter (D_f) value of 13 ± 1 .**
 195

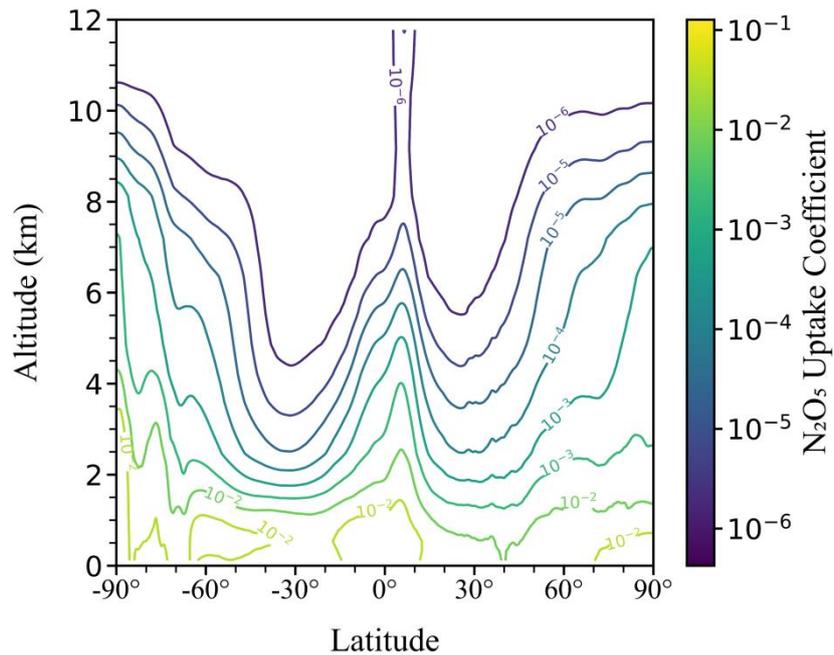


197 **Figure S9. Altitude–latitude profiles of zonal-mean temperature and RH obtained from Copernicus Climate Data Store**
198 **(<https://cds.climate.copernicus.eu/>), averaged seasonally and annually over the period 2020-01-01 to 2024-12-01. Panels**
199 **present seasonal profiles for spring (a and b), summer (c and d), fall (e and f), and winter (g and h), showing temperature**
200 **and RH, respectively. Panels (I and j) display the annual average zonal-mean profiles of temperature and RH.**



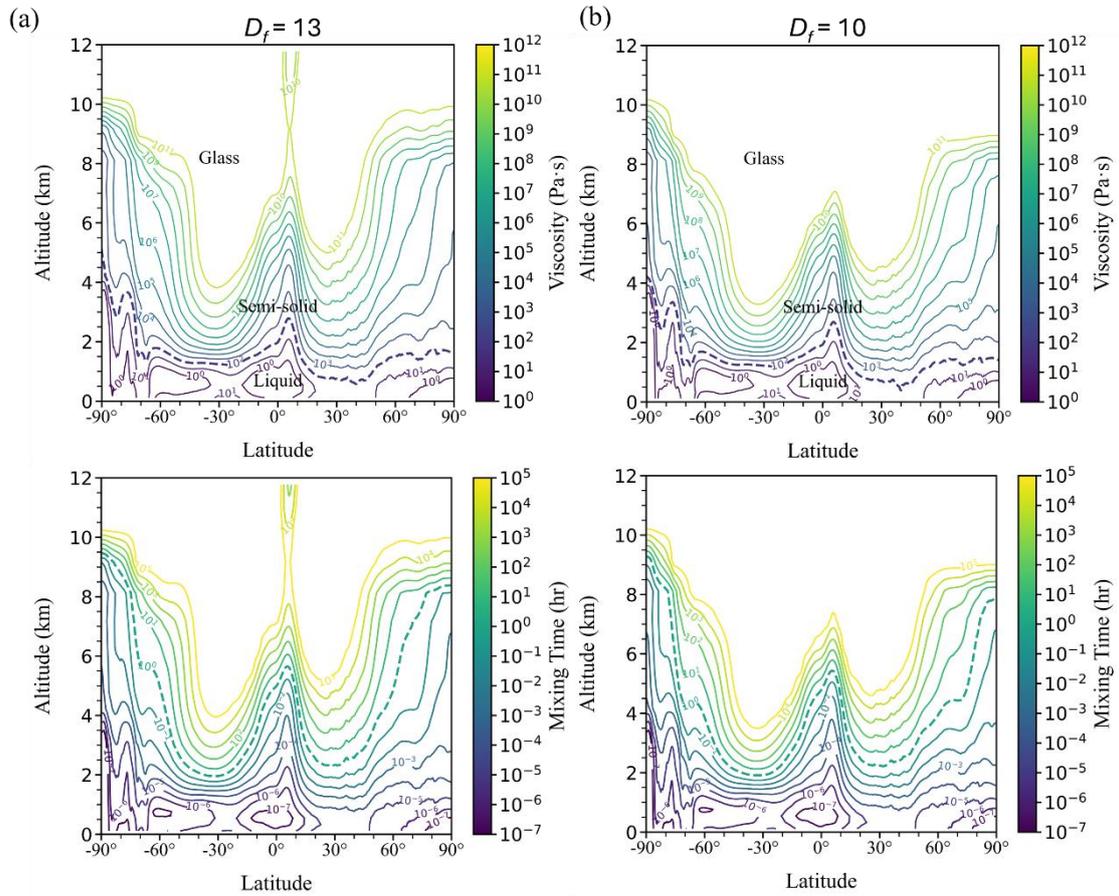
202 Figure S10. Seasonal zonal-mean profiles of viscosity and mixing time of sucrose-H₂O droplets as a function of altitude
203 and latitude, derived from temperature and RH data obtained from Copernicus Climate Data Store
204 (<https://cds.climate.copernicus.eu/>) and averaged over the period 2020-01-01 to 2024-12-01. Panels (a and b) present
205 viscosity and mixing time for spring, panels (c and d) for summer, panels (e and f) for fall, and panels (g and h) for
206 winter. The blue dashed line shows the transition from liquid to semi-solid state. The light green dashed shows the mixing
207 time of 1 hr.

208



209
 210 **Figure S11. N_2O_5 uptake coefficient in 200 nm sucrose- H_2O droplets as a function of altitude and latitude based on**
 211 **annual average zonal-mean RH and temperature fields for the years 2020 to 2024, obtained from Copernicus Climate**
 212 **Data Store (<https://cds.climate.copernicus.eu/>).** Surface hydrolysis is not considered in calculations of the N_2O_5 uptake
 213 **coefficient (Γ_s is set to be 0 in Eq. S5.1).**

214



215

216 **Figure S12. Influence of the fragility parameter D_f on the viscosity of sucrose-H₂O and related mixing-times in the**
 217 **troposphere. Panel (a) presents the viscosity of sucrose-H₂O and the mixing times calculated using $D_f=13$, whereas panel**
 218 **(b) shows the corresponding profiles obtained with $D_f=10$. All fields are shown as functions of altitude and latitude,**
 219 **derived from annual zonal-mean relative humidity and temperature for the years 2020–2024, obtained from the**
 220 **Copernicus Climate Data Store (<https://cds.climate.copernicus.eu/>).**

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225 **Table S1.** Summary of the conditioning time for sucrose-H₂O droplets measured at different temperatures and RH conditions.
 226 τ_{mix,H_2O} denotes the calculated characteristic mixing time of water, calculated from the diffusion coefficient of water from the
 227 viscosity inferred at the corresponding temperature and RH. $\tau_{conditioning}$ values represent the experimental conditioning time.
 228 Calculations were performed using the upper-limit droplet diameters employed in the experiments, namely 100 μm for bead-
 229 mobility experiments and 40 μm for poke-and-flow experiments.
 230

RH	Particle dia. (μm)	$\tau_{conditioning}$ (h)	τ_{mix,H_2O} (h)	$\tau_{conditioning} / \tau_{mix,H_2O}$
Temp = 273 K				
79	100	3.30E-01	5.69E-03	58.0
53	40	1.50E+00	2.39E-01	6.3
49	40	1.50E+00	4.58E-01	3.3
Temp = 283 K				
81	100	1.60E-01	2.89E-03	55.3
76	100	1.60E-01	7.81E-03	20.5
55	40	1.00E+00	2.19E-02	45.7
49	40	1.00E+00	8.86E-02	11.3
46	40	1.00E+00	1.94E-01	5.1
42	40	1.00E+00	4.64E-01	2.2
Temp = 293 K				
90	100	1.60E-01	2.88E-04	555.7
86	100	1.60E-01	3.48E-04	459.8

83	100	1.60E-01	6.53E-04	245.1
79	100	1.60E-01	1.15E-03	139.0
75	100	1.60E-01	2.35E-03	68.1
71	100	1.60E-01	3.18E-03	50.4
50	40	5.00E-01	1.82E-02	27.5
44	40	5.00E-01	2.14E-02	23.3
39	40	5.00E-01	1.48E-01	3.4
34	40	5.00E-01	5.08E-01	1.0

Temp = 303 K

73	100	1.60E-01	1.39E-03	115.5
69	100	1.60E-01	2.63E-03	60.7
65	100	1.60E-01	4.13E-03	38.7
50	40	5.00E-01	6.54E-03	76.5
44	40	5.00E-01	8.78E-03	56.9
39	40	5.00E-01	2.69E-02	18.6
34	40	5.00E-01	5.95E-02	8.4
29	40	5.00E-01	3.25E-01	1.5

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233

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